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# A TWO-DIMENSIONAL OBLIQUE EXTENSION OF BESSEL PROCESSES 

Dominique Lépingle ${ }^{1}$


#### Abstract

We consider a Brownian motion forced to stay in the quadrant by an electrostatic oblique repulsion from the sides. We tackle the question of hitting the corner or an edge, and find product-form stationary measures under a certain condition, which is reminiscent of the skew-symmetry condition for a reflected Brownian motion.


## 1. Introduction

In the present paper we study existence and properties of a new process with values in the nonnegative quadrant $S=\mathbb{R}_{+} \times \mathbb{R}_{+}$where $\mathbb{R}_{+}:=[0, \infty)$. It may be seen as a two-dimensional extension of a usual Bessel process. It is a two-dimensional Brownian motion forced to stay in the quadrant by electrostatic repulsive forces, in the same way as in the one-dimensional case where a Brownian motion which is prevented from becoming negative by an electrostatic drift becomes a Bessel process. Note here and now that the corner $\mathbf{0}=(0,0)$ will play a crucial role and in some cases it will be necessary to restrict the state space to the punctured nonnegative quadrant $S^{\mathbf{0}}=S \backslash\{\mathbf{0}\}$.

Let $\left(\Omega, \mathcal{F},\left(\mathcal{F}_{t}\right)_{t \geq 0}, \mathbb{P}\right)$ be a complete probability space endowed with a filtration $\left(\mathcal{F}_{t}\right)_{t \geq 0}$ satisfying the usual conditions. Let $\left(B_{t}, C_{t}\right)$ be an adapted driftless Brownian motion in the plane starting from $\mathbf{0}$, with covariance matrix

$$
\left(\begin{array}{ll}
1 & \rho \\
\rho & 1
\end{array}\right)
$$

and $\rho \in[-1,+1]$.
Definition 1. Let $\alpha, \beta, \gamma, \delta$ be four real constants with $\alpha>0, \delta>0$. We say that an $\left(\mathcal{F}_{t}\right)$ adapted continuous process $(X, Y)$ with values in $S$ is an Oblique Two-dimensional Bessel Process (O2BP) if for any $t \geq 0$

$$
\begin{align*}
& X_{t}=X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}}+\beta \int_{0}^{t} \frac{d s}{Y_{s}} \geq 0 \\
& Y_{t}=Y_{0}+C_{t}+\gamma \int_{0}^{t} \frac{d s}{X_{s}}+\delta \int_{0}^{t} \frac{d s}{Y_{s}} \geq 0 \tag{1}
\end{align*}
$$

where $X_{0}$ and $Y_{0}$ are nonnegative $\mathcal{F}_{0}$-measurable random variables, and

$$
\begin{array}{ll}
\int_{0}^{t} \mathbf{1}_{\left\{X_{s}=0\right\}} d s=0 & \int_{0}^{t} \mathbf{1}_{\left\{Y_{s}=0\right\}} d s=0 \\
\int_{0}^{t} \mathbf{1}_{\left\{X_{s}>0\right\}} \frac{d s}{X_{s}}<\infty & \int_{0}^{t} \mathbf{1}_{\left\{Y_{s}>0\right\}} \frac{d s}{Y_{s}}<\infty
\end{array}
$$

This stochastic differential system is very singular at the edges of the quadrant and the question of existence and uniqueness of a solution is not simple. The particular case when $\beta=\gamma=0$

$$
\begin{align*}
U_{t} & =X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{U_{s}} \geq 0 \\
V_{t} & =Y_{0}+C_{t}+\delta \int_{0}^{t} \frac{d s}{V_{s}} \geq 0 \tag{2}
\end{align*}
$$

[^0]is already known: the processes $U$ and $V$ are Bessel processes. Actually, $U$ is a Bessel process of dimension $2 \alpha+1$, and the point 0 is instantaneously reflecting for $U$ if $\alpha<\frac{1}{2}$ and polar if $\alpha \geq \frac{1}{2}$. If $\rho=0, U^{2}+V^{2}$ is the square of a Bessel process of dimension $2 \alpha+2 \delta+2$ [32], and so the corner $\mathbf{0}$ is polar for $(U, V)$ in this case. Comparison between $X$ and $U, Y$ and $V$ will play a key role in the construction of the solution $(X, Y)$ and the study of its behavior close to the edges of the quadrant. The process $(U, V)$ is an example of Brownian motion perturbed by a drift deriving from a convex potential. More generally, stochastic differential systems including such a singular drift have been studied in $[2,26,34,27]$, where strong existence and uniqueness were obtained. They are examples of so-called multivalued stochastic differential equations, also called stochastic variational inequalities in convex analysis.

We will use these results to study the solutions to (1) in the oblique case where $\beta$ or $\gamma$ do not vanish. We obtain strong existence and uniqueness for a large set of parameters and initial conditions, but not for all possible values. In the proofs we naturally fall into the crucial question of hitting the corner, that is the non-smooth part of the boundary. Using McKean's argument on the asymptotic behavior of continuous local martingales obtained by time change from the real driftless Brownian motion, we are able to state several sufficient conditions to prevent the processes from hitting the corner. Our methods are not powerful enough to allow for necessary conditions. However we shall not restrict to processes avoiding the corner and, depending on the parameters, we will get existence and uniqueness (in a strong sense) sometimes in the whole quadrant, sometimes in the punctured quadrant, that is, the quadrant without a corner.

We will also obtain some partial results about the attainibility of the edges of the quadrant. It is interesting to see whether the boundary behavior of one component may be modified by the interaction with the other component.

In the one-dimensional case, the so-called scale functions transform the solution of a stochastic differential equation into local martingales and one may infer some information on the boundary behavior of the process. This technique was very successful in the study of Bessel processes. Here we still obtain functions of an O2BP which are local martingales or supermartingales for some values of the parameters and we derive some information on its asymptotic behavior.

The laws of Bessel processe with different parameters [28] are mutually abolutely continuous when their paths do not reach the origin. Here we obtain two partial results about the absolute continuity of the laws of O2BPs for some values of the parameters.

Finally we follow the way explored in [23], where a drifted Brownian motion is endowed with another drift term that is continuous and depends obliquely, via a regular potential function, on the position of the process relative to an orthant. Under an additional condition, which is called a skew-symmetry condition, an invariant density was given in an explicit product form. We obtain the same result for an O2BP (where the Brownian motion is now drifted) under an anologous condition. This time, the invariant density is the product of two gamma densities, which is consistent with the one-dimensional case where the invariant measure of a drifted Bessel process is a gamma distribution.
1.1. Comparison with obliquely reflected Brownian motion. Studying O2BPs makes appear a strong connection with the properties of a semimartingale reflecting Brownian motion in the quadrant. We briefly recall the definition. Let $\left(B_{t}, C_{t}\right)_{t \geq 0}$ be the Brownian motion in Definition 1.

Definition 2. An semimartingale reflecting Brownian motion (SRBM) in the quadrant is a continuous adapted process $\left(X_{t}, Y_{t}\right)$ which is a solution to the system

$$
\begin{align*}
& X_{t}=X_{0}+B_{t}+L_{t}^{1}+r_{1} L_{t}^{2} \geq 0 \\
& Y_{t}=Y_{0}+C_{t}+r_{2} L_{t}^{1}+L_{t}^{2} \geq 0 \tag{3}
\end{align*}
$$

where $L^{1}$ and $L^{2}$ are two continuous adapted nondecreasing processes with $L_{0}^{1}=L_{0}^{2}=0$ such that for any $t \geq 0$

$$
\int_{0}^{t} \mathbf{1}_{\left\{X_{s}>0\right\}} d L_{s}^{1}=\int_{0}^{t} \mathbf{1}_{\left\{Y_{s}>0\right\}} d L_{s}^{2}=0
$$

Here $X_{0}$ and $Y_{0}$ are nonnegative $\mathcal{F}_{0}$-measurable random variables, $r_{1}$ and $r_{2}$ are real numbers.
There was an extensive literary output on that topic in the eighties, with a more general domain : a wedge, an orthant or a convex polyhedron. We mention the works of Harrison, Reiman, Varadhan, Williams, Dai $[7,9,36,39,4]$ to cite a few of them. For a more complete bibliography we refer to [40].

We note a first analogy between their results and ours: we prove a necessary condition of existence of an O2BP is the existence of a convex combination of the directions of interaction

$$
\begin{equation*}
\mathbf{r}_{\mathbf{x}}=\binom{\alpha}{\gamma} \text { and } \mathbf{r}_{\mathbf{y}}=\binom{\beta}{\delta} \tag{4}
\end{equation*}
$$

that points into the quadrant from the corner. In the reflection setting it was proved in [37] that a necessary an sufficient condition for the existence of a SRBM is the existence of a convex combination of the directions of reflection

$$
\binom{1}{r_{2}} \text { and }\binom{r_{1}}{1}
$$

with the same property.
Another analogy is met in the delicate and important question of attainibility of the corner. The authors in [36] have found an explicit harmonic function that provides a full answer to the question of hitting the corner for a reflected Brownian motion in a wedge of angle $\xi \in(0,2 \pi)$ with the identity matrix as covariance matrix. We do not have any such convenient function. However, our second condition

$$
\begin{equation*}
2 \rho \leq \frac{\beta}{\delta}+\frac{\gamma}{\alpha} \tag{5}
\end{equation*}
$$

in Corollary 10 is reminiscent of the necessary and sufficient condition in [29] which writes

$$
\begin{equation*}
2 \rho \leq r_{1}+r_{2} \tag{6}
\end{equation*}
$$

in the setting of Definition 2. Restricting to $\xi \in(0, \pi)$, we may consider a linear transformation in the plane that changes the wedge in [36] into the nonnegative quadrant and the initial Brownian motion into a Brownian motion with covariance coefficient $\rho=-\cos \xi$. Then the non-attainibility condition in [36] is transformed into the above condition (6).

For theoretical as well as practical reasons, a great deal of interest was taken in the question of recurrence of the Brownian motion with a constant drift vector and oblique reflection, and in the computation of the invariant measure $[38,41,8,31]$. Under the assumption that the directions of reflection satisfy a skew-symmetry condition, it was proved that the invariant measure has exponential product form density $[10,11,39,40]$. This result has a practical interest because of the interpretation of the SRBM as approximation of the queue length processes for networks of queues in heavy traffic [7, 10]. There is also a financial reason for studying SRBMs. Motivated by the so-called Atlas model of equity market presented in [6], some authors $[13,17,14,15,16,29]$ have recently studied competing Brownian particles on the line with rank dependent local characteristics. In fact the gaps between adjacent particles are proportional to the components of a SRBM in an orthant. There is an invariant probability density with an explicit exponential product form when the volatility coefficients are constant [24]. More generally, this is still true if their squares depend on rank linearly since one may infer from the statements in Section 2 of [29] that the skew-symmetry condition is still satisfied. Following the way in [11], [5] and [8], a more general question is the recurrence
or transience of an O2BP. Answering this question does not seem to be an easy task. The method in [12], which provides a full answer for the obliquely reflected Brownian motion in the quadrant, appears to break down here. Mimicking the computation in [23], we just calculate an invariant measure in product form under a skew-symmetry condition which is the equality condition in the inequality (5), whereas the skew-symmetry condition in the reflecting case is the equality condition in (6). Now the terms of the product are gamma distributions with explicit parameters.

Another topic of interest in [23] lies in its Remark 4.12: when a scale parameter goes to zero, the exponentially reflected Brownian motion should converge to the obliquely reflected Brownian motion. It is typically a penalty method. This kind of approximation by a sequence of diffusions with regular drifts living on the whole Euclidean space has been used for instance in $[20,30]$ for oblique reflection in domains with smooth boundary and in $[22,33]$ for normal reflection in convex domains. With regard to our framework, it can be shown that a sequence of real Bessel processes with dimension decreasing to 1 pathwise decreases to a reflected real Brownian motion. It could be interesting to consider a sequence of oblique two-dimensional Bessel processes with interaction matrices

$$
n^{-1}\left(\begin{array}{cc}
1 & r_{1} \\
r_{2} & 1
\end{array}\right)
$$

and to let $n$ go to infinity. It should converge to the SRBM in Definition 2. This time, this would be an interior approximation, as it was done in [41] and [25] in the particular setting of stationary Markov processes associated with Dirichlet forms. We leave aside this point for further investigation.
1.2. Organization of the paper. The rest of the paper is organized as follows. In Section 2 we recall some trajectorial properties of usual Bessel processes. In Section 3 we state and prove three main lemmas of repeated use in the sequel. We also recall an existence and uniqueness result for a special case of multivalued stochastic differential equation that will be useful in our construction of an OB2P. Section 4 is devoted to the proof of sufficient conditions to avoid the corner of the quadrant. The main theorems of existence and uniqueness of an O2BP are given in Section 5. In Section 6 we discuss the question of hitting the edges of the quadrant. In Section 7 we consider two particular cases where there exist simple functions of O2BPs that are local martingales and obtain some information on the asymptotic behavior of the paths. We leave the trajectorial point of view in Section 8 to tackle questions of absolute continuity of the law of an O2BP with respect to the product of laws of real Bessel processes. The final Section 9 introduces a skew-symmetry condition that allows us to obtain existence of a stationary probability in form of the product of two gamma distributions.

## 2. Some properties of Bessel processes

A Bessel process of dimension $d>1$, starting at $r \geq 0$, is the unique solution to the stochastic differential equation

$$
\begin{equation*}
R_{t}=r+W_{t}+\frac{d-1}{2} \int_{0}^{t} \frac{1}{R_{s}} d s \tag{7}
\end{equation*}
$$

where $W$ is a standard driftless real-valued Brownian motion starting at 0 .
We know that:

- $d \geq 2:$ the point 0 is a polar set ([28], Proposition V.2.7);
- $d=2: \lim \sup _{t \rightarrow \infty} R_{t}=+\infty, \liminf _{t \rightarrow \infty} R_{t}=0$ ([28], Theorem V.2.8);
- $1<d<2$ : the point 0 is instantaneously reflecting ([28], Proposition XI.1.5).

More precisely ([18], p. 337 and p.339),

- Dimension $d>2$ :

$$
\begin{aligned}
& \mathbb{P}\left(R_{t}>0, \forall t>0\right)=1, \\
& \mathbb{P}\left(R_{t} \rightarrow \infty, t \rightarrow \infty\right)=1, \\
& R^{2-d} \text { is a local martingale. }
\end{aligned}
$$

- Dimension $d=2$ :

$$
\begin{aligned}
& \mathbb{P}\left(R_{t}>0, \forall t>0\right)=1, \\
& \mathbb{P}\left(\sup _{t>a} R_{t}=\infty, \inf _{t>a} R_{t}=0\right)=1 \text { for any } a>0, \\
& \ln R \text { is a local martingale. }
\end{aligned}
$$

- Dimension $1<d<2$ :

$$
P\left(R_{t}>0, \forall t>a\right)=0 \text { for any } a>0 .
$$

Corollary 3. When $d \geq 2$, for any $a>0$,

$$
\int_{a}^{\infty} \frac{d s}{R_{s}^{2}}=\infty \text { a.s. }
$$

Proof. From Itô's formula

$$
\ln R_{t}=\ln R_{a}+\int_{a}^{t} \frac{d W_{s}}{R_{s}}+\frac{d-2}{2} \int_{a}^{t} \frac{d s}{R_{s}^{2}} .
$$

From ([28], Proposition IV.1.26) we know that on the set $\left\{\int_{a}^{\infty} \frac{d s}{R_{s}^{2}}<\infty\right\}$, the continuous local martingale $\int_{a}^{t} \frac{d W_{s}}{R_{s}}$ converges as $t \rightarrow \infty$. But we have seen that with unit probability $\ln R_{t}$ does not converge.

We shall also use an absolute continuity result ([28], Exercise XI.1.22 or [18], Proposition 6.1.5.1). On the canonical space $\Omega=C([0, \infty),[0, \infty))$, we denote by $R$ the canonical map $R_{t}(\omega)=\omega(t)$, by $\mathcal{R}_{t}=\sigma\left(R_{s}, s \leq t\right)$ the canonical filtration and by $\mathbb{P}_{r}^{d}$ the law of the Bessel process of dimension $d \geq 2$ starting at $r>0$. Then,

$$
\begin{equation*}
\left.\mathbb{P}_{r}^{d}\right|_{\mathcal{R}_{t}}=\left.\left(\frac{R_{t}}{r}\right)^{\frac{d-2}{2}} \exp \left(-\frac{(d-2)^{2}}{8} \int_{0}^{t} \frac{d s}{R_{s}^{2}}\right) \cdot \mathbb{P}_{r}^{2}\right|_{\mathcal{R}_{t}} \tag{8}
\end{equation*}
$$

## 3. Four useful tools

The following simple lemma is a comparison lemma. It will play an ubiquitous role in our proofs.

Lemma 4. For $T>0, \alpha>0$, let $x_{1}$ and $x_{2}$ be nonnegative continuous solutions on $[0, T]$ to the equations

$$
\begin{aligned}
& x_{1}(t)=v_{1}(t)+\alpha \int_{0}^{t} \frac{d s}{x_{1}(s)} \\
& x_{2}(t)=v_{2}(t)+\alpha \int_{0}^{t} \frac{d s}{x_{2}(s)}
\end{aligned}
$$

where $v_{1}, v_{2}$ are continuous functions such that $0 \leq v_{1}(0) \leq v_{2}(0)$, and $v_{2}-v_{1}$ is nondecreasing. Then $x_{1}(t) \leq x_{2}(t)$ on $[0, T]$.

Proof. Assume there exists $t \in(0, T]$ such that $x_{2}(t)<x_{1}(t)$. Set

$$
\tau:=\max \left\{s \leq t: x_{1}(s) \leq x_{2}(s)\right\} .
$$

Then,

$$
\begin{aligned}
x_{2}(t)-x_{1}(t) & =x_{2}(\tau)-x_{1}(\tau)+\left(v_{2}(t)-v_{1}(t)\right)-\left(v_{2}(\tau)-v_{1}(\tau)\right)+\alpha \int_{\tau}^{t}\left(\frac{1}{x_{2}(s)}-\frac{1}{x_{1}(s)}\right) d s \\
& \geq 0,
\end{aligned}
$$

a contradiction.

The following elementary lemma will also be repeatedly used.
Lemma 5. Let $Q(x, y)=a x^{2}+b x y+c y^{2}$ be a second degree homogeneous polynomial. Then $Q$ is nonnegative on the whole $S$ if and only if $a \geq 0, c \geq 0$ and $b \geq-2 \sqrt{a c}$.

## Proof.

Taking $x=0$ (resp. $y=0$ ) we see that $a$ (resp. c) must be nonnegative. Then we rewrite

$$
Q(x, y)=(\sqrt{a} x-\sqrt{c} y)^{2}+(b+2 \sqrt{a c}) x y .
$$

So $Q$ is nonnegative for all $x \geq 0$ and $y \geq 0$ iff $b \geq-2 \sqrt{a c}$.
Another main tool will be the following convergence result, whose argument goes back to McKean ([21], p. 31 and p.47). The statement and the proof below are borrowed from ([28], Theorem V.1.7 and Proposition V.1.8). The only change lies in the introduction of a stopping time $\tau$ up to which the local martingale $M$ is now defined.
Lemma 6. Let $M$ be a continuous local martingale defined on $[0, \tau)$ where $\tau$ is a stopping time. Let $\left(\langle M\rangle_{t}\right)_{0 \leq t<\tau}$ be its quadratic variation and let $\langle M\rangle_{\tau}:=\lim _{t \rightarrow \tau}\langle M\rangle_{t} \leq \infty$.
(1) $O n\left\{\langle M\rangle_{\tau}<\infty\right\}, \quad \lim _{t \rightarrow \tau} M_{t}$ exists a.s. in $\mathbb{R}$.
(2) On $\left\{\langle M\rangle_{\tau}=\infty\right\}$, $\quad \limsup \sup _{t \rightarrow \tau} M_{t}=-\liminf _{t \rightarrow \tau} M_{t}=+\infty$ a.s.

Proof. (1)For any $p \geq 1$ let

$$
\sigma_{p}=\inf \left\{t \geq 0:\langle M\rangle_{t} \geq p\right\} .
$$

In order every term in the following to be well defined we introduce a nondecreasing sequence of stopping times $\left(\tau_{n}\right)_{n \geq 1}$ with limit $\tau$ such that each stopped process $\left(M_{\tau_{n} \wedge t}\right)_{t \geq 0}$ is a uniformly integrable martingale for any $n \geq 1$. The stopped process $\left(M_{t \wedge \tau_{n} \wedge \sigma_{p}}\right)_{t \geq 0}$ is a $L^{2}$-bounded martingale and as $m, n \rightarrow \infty$

$$
\mathbb{E}\left[\left(M_{\tau_{n} \wedge \sigma_{p}}-M_{\tau_{m} \wedge \sigma_{p}}\right)^{2}\right]=\mathbb{E}\left[\left|\langle M\rangle_{\tau_{n} \wedge \sigma_{p}}-\langle M\rangle_{\tau_{m} \wedge \sigma_{p}}\right|\right] \rightarrow 0 .
$$

We set

$$
M^{(p)}:=\lim _{n \rightarrow \infty} M_{\tau_{n} \wedge \sigma_{p}}=\lim _{t \rightarrow \tau} M_{t \wedge \sigma_{p}} .
$$

On $\left\{\langle M\rangle_{\tau}<\infty\right\}$, the stopping times $\sigma_{p}$ are a.s. infinite from some $p$ on and we can set $M_{\tau}:=\lim _{p \rightarrow \infty} M^{(p)}$. Thus on this set

$$
M_{t} \rightarrow M_{\tau} \quad \text { as } \quad t \rightarrow \tau .
$$

(2) Let for any $t \geq 0$

$$
\left.T_{t}=\inf _{\tilde{\sim}} 0 \leq s \leq \tau:\langle M\rangle_{s}>t\right\} .
$$

There exist an enlargement $\left(\tilde{\Omega},\left(\tilde{\mathcal{F}}_{t}\right), \tilde{\mathbb{P}}\right)$ of $\left(\Omega,\left(\mathcal{F}_{T_{t}}\right), \mathbb{P}\right)$ and a Brownian motion $\tilde{\beta}$ on $\tilde{\Omega}$ independent of $M$ such that the process

$$
B_{t}=\left\{\begin{array}{lll}
M_{T_{t}} & \text { if } \quad t<\langle M\rangle_{\tau} \\
M_{\tau}+\tilde{\beta}_{t-\langle M\rangle_{\tau}} & \text { if } \quad t \geq\langle M\rangle_{\tau}
\end{array}\right.
$$

is a standard linear Brownian motion. As

$$
\limsup _{t \rightarrow \infty} B_{t}=-\liminf _{t \rightarrow \infty} B_{t}=+\infty \quad \text { a.s. }
$$

we obtain on $\left\{\langle M\rangle_{\tau}=\infty\right\}$

$$
\limsup _{t \rightarrow \infty} M_{T_{t}}=-\liminf _{t \rightarrow \infty} M_{T_{t}}=+\infty
$$

and therefore

$$
\limsup _{t \rightarrow \tau} M_{t}=-\liminf _{t \rightarrow \tau} M_{t}=+\infty
$$

We will also need the following consequence of the results in [2] on multivalued stochastic differential systems, completed with the method used in [3] and developed in [19] to check the lack of additional boundary process.

Proposition 7. Let $\alpha>0, \delta \geq 0, \sigma=\left(\sigma_{j}^{i} ; i, j=1,2\right)$ a $2 \times 2$-matrix, $(B, C)$ a Brownian motion in the plane, $b_{1}$ and $b_{2}$ two Lipschitz functions on $\mathbb{R}^{2}, Z_{0}^{1}$ and $Z_{0}^{2}$ two $\mathcal{F}_{0}$-measurable nonnegative random variables. There exists a unique strong solution $\left(Z^{1}, Z^{2}\right)$ to the system

$$
\begin{align*}
& Z_{t}^{1}=Z_{0}^{1}+\sigma_{1}^{1} B_{t}+\sigma_{2}^{1} C_{t}+\alpha \int_{0}^{t} \frac{d s}{Z_{1}^{1}}+\int_{0}^{t} b_{1}\left(Z_{s}^{1}, Z_{s}^{2}\right) d s \\
& Z_{t}^{2}=Z_{0}^{2}+\sigma_{1}^{2} B_{t}+\sigma_{2}^{2} C_{t}+\delta \int_{0}^{t} \frac{d_{s}^{2}}{Z_{s}^{2}}+\int_{0}^{t} b_{2}\left(Z_{s}^{1}, Z_{s}^{2}\right) d s \tag{9}
\end{align*}
$$

with the conditions $Z_{t}^{1} \geq 0$ if $\delta=0$ and $Z_{t}^{1} \geq 0, Z_{t}^{2} \geq 0$ if $\delta>0$.

It is worth noticing that the solutions to (1) enjoy the Brownian scaling property. It means that if $(X, Y)$ is a solution to (1) starting from $\left(X_{0}, Y_{0}\right)$ with driving Brownian motion $\left(B_{t}, C_{t}\right)$, then for any $c>0$ the process ( $X_{t}^{\prime}:=c^{-1} X_{c^{2} t}, Y_{t}^{\prime}:=c^{-1} Y_{c^{2}} ; t \geq 0$ ) is a solution to (1) starting from ( $c^{-1} X_{0}, c^{-1} Y_{0}$ ) with driving Brownian motion $\left(c^{-1} B_{c^{2} t}, c^{-1} C_{c^{2} t}\right)$.

## 4. Avoiding the corner

We shall see in Section 5 that existence and uniqueness of the solution to (1) are easily obtained as soon as the solution process keeps away from the corner. Thus the question of attaining the corner in finite time is of great interest. For some class of reflection matrices, a necessary and sufficient condition is given in [29]. Unfortunately, we are not able to provide such a complete answer and we have to be content with a collection of sufficient conditions ensuring nonattainability. We will just see at the end of Section 5 a degenerated case where the corner is reached in finite time with full probability.

Our sufficient conditions are stated in the following theorem. In some cases (conditions $C_{2 a}$ and $C_{2 b}$ below), the comparison with Bessel processes suffices to conclude. In other cases (conditions $C_{1}$ and $C_{3}$ below), we are first looking for a $C^{2}$-function $f$ on the punctured quadrant $S^{\mathbf{0}}$ with limit $-\infty$ at the corner. Then we use Lemma 6 to show that $f\left(X_{t}, Y_{t}\right)$ cannot converge to $-\infty$ in finite time.

Theorem 8. Let $(X, Y)$ be a solution to (1). We set

$$
\tau^{\mathbf{0}}:=\inf \left\{t>0:\left(X_{t}, Y_{t}\right)=\mathbf{0}\right\}
$$

with the usual convention $\inf \emptyset=\infty$. Then $\mathbb{P}\left(\tau^{0}<\infty\right)=0$ if one of the following conditions is satisfied:
(1) $C_{1}: \beta \geq 0, \gamma \geq 0$ and $-1 \leq \rho \leq \alpha+\delta$.
(2) $C_{2 a}: \alpha \geq \frac{1}{2}$ and $\beta \geq 0$
(3) $C_{2 b}: \delta \geq \frac{1}{2}$ and $\gamma \geq 0$
(4) $C_{3}$ : There exist $\lambda>0$ and $\mu>0$ such that

- $\lambda \alpha+\mu \gamma \geq 0$
- $\lambda \beta+\mu \delta \geq 0$
- $(\sqrt{\lambda(\lambda \alpha+\mu \gamma)}+\sqrt{\mu(\lambda \beta+\mu \delta)})^{2} \geq \frac{1}{2}\left(\lambda^{2}+\mu^{2}+2 \rho \lambda \mu\right)$.


## Proof.

Condition $C_{1}$. For $\epsilon>0$ let

$$
\begin{aligned}
& \sigma^{\epsilon}=\mathbf{1}_{\left\{\left(X_{0}, Y_{0}\right)=\mathbf{0}\right\}} \inf \left\{t>0: X_{t}+Y_{t} \geq \epsilon\right\} \\
& \boldsymbol{\tau}^{\mathbf{0}, \epsilon}=\inf \left\{t>\sigma^{\epsilon}:\left(X_{t}, Y_{t}\right)=\mathbf{0}\right\} .
\end{aligned}
$$

As $\epsilon \downarrow 0, \sigma^{\epsilon} \downarrow 0$ and $\tau^{\mathbf{0}, \epsilon} \downarrow \tau^{\mathbf{0}}$. We set $R_{t}=X_{t}^{2}+Y_{t}^{2}$. From Itô's formula we get for $t \in\left[\sigma^{\epsilon}, \tau^{\mathbf{0}, \epsilon}\right)$

$$
\begin{aligned}
& \ln R_{t} \\
= & \ln R_{\sigma^{\epsilon}}+2 \int_{\sigma^{\epsilon}}^{t} \frac{X_{s} d B_{s}+Y_{s} d C_{s}}{R_{s}}+2 \int_{\sigma^{\epsilon}}^{t} \frac{(\alpha+\delta) d s}{R_{s}}+2 \int_{\sigma^{\epsilon}}^{t}\left(\beta \frac{X_{s}}{Y_{s}}+\gamma \frac{Y_{s}}{X_{s}}\right) \frac{d s}{R_{s}}-4 \rho \int_{\sigma^{\epsilon}}^{t} \frac{X_{s} Y_{s}}{R_{s}^{2}} d s \\
\geq & \ln R_{\sigma^{\epsilon}}+M_{t}+2 \int_{\sigma^{\epsilon}}^{t} \frac{P_{1}\left(X_{s}, Y_{s}\right)}{R_{s}^{2}} d s
\end{aligned}
$$

where $M$ is a continuous local martingale on $\left[\sigma^{\epsilon}, \tau^{\mathbf{0}, \epsilon}\right)$ and $P_{1}(x, y)$ is the second degree homogeneous polynomial

$$
P_{1}(x, y)=(\alpha+\delta)\left(x^{2}+y^{2}\right)-2 \rho x y
$$

Using Lemma 5 we check that this polynomial is nonnegative on $S$ if $\rho \leq \alpha+\delta$. Therefore

$$
0 \leq \int_{\sigma^{\epsilon}}^{\tau^{0, \epsilon}} \frac{P_{1}\left(X_{s}, Y_{s}\right)}{R_{s}^{2}} d s \leq \infty
$$

From Lemma 6 we know that depending on whether $\langle M\rangle_{\tau^{0, \epsilon}}$ is finite or not, the local martingale $M_{t}$ either converges in $\mathbb{R}$ as $t \rightarrow \tau^{\mathbf{0}, \epsilon}$ or oscillates between $+\infty$ and $-\infty$. It cannot converge to $-\infty$. Thus $R_{\tau^{0, \epsilon}}>0$ on $\left\{\tau^{\mathbf{0}, \epsilon}<\infty\right\}$, which contradicts the definition of the moment $\tau^{\mathbf{0}, \epsilon}$. This proves that $\tau^{\mathbf{0}, \epsilon}=\infty$ for every $\epsilon>0$, and therefore, $\tau^{\mathbf{0}}=\infty$ a.s.
Condition $C_{2 a}$ (resp. $C_{2 b}$ ). Recall the Bessel processes $U$ and $V$ in (2). From Lemma 1 we get $X_{t} \geq U_{t}$ (resp. $Y_{t} \geq V_{t}$ ) and in this case 0 is polar for $U$ (resp. V), so $X_{t}>0$ (resp. $\left.Y_{t}>0\right)$ for $t>0$.
Condition $C_{3}$. We use again the notation for $\epsilon>0, \sigma^{\epsilon}$ and $\tau^{0, \epsilon}$. We set $S_{t}=\lambda X_{t}+\mu Y_{t}$ for $t \geq 0, \lambda>0$ and $\mu>0$. From Itô's formula we get for $t \in\left[\sigma^{\epsilon}, \tau^{\mathbf{0}, \epsilon}\right)$

$$
\begin{aligned}
& \ln S_{t} \\
= & \ln S_{\sigma^{\epsilon}}+\int_{\sigma^{\epsilon}}^{t} \frac{\lambda d B_{s}+\mu d C_{s}}{S_{s}}+(\lambda \alpha+\mu \gamma) \int_{\sigma^{\epsilon}}^{t} \frac{d s}{X_{s} S_{s}}+(\lambda \beta+\mu \delta) \int_{\sigma^{\epsilon}}^{t} \frac{d s}{Y_{s} S_{s}} \\
& \quad-\frac{1}{2}\left(\lambda^{2}+\mu^{2}+2 \rho \lambda \mu\right) \int_{\sigma^{\epsilon}}^{t} \frac{d s}{S_{s}^{2}} \\
= & \ln S_{\sigma^{\epsilon}}+M_{t}+\int_{\sigma^{\epsilon}}^{t} \frac{P_{2}\left(X_{s}, Y_{s}\right)}{X_{s} Y_{s} S_{s}^{2}} d s
\end{aligned}
$$

where $M$ is a continuous local martingale on $\left[\sigma^{\epsilon}, \tau^{\mathbf{0}, \epsilon}\right)$ and $P_{2}(x, y)$ is the second degree homogeneous polynomial

$$
\begin{aligned}
& P_{2}(x, y) \\
= & \lambda(\lambda \beta+\mu \delta) x^{2}+\mu(\lambda \alpha+\mu \gamma) y^{2}+\left[\lambda(\lambda \alpha+\mu \gamma)+\mu(\lambda \beta+\mu \delta)-\frac{1}{2}\left(\lambda^{2}+\mu^{2}+2 \rho \lambda \mu\right)\right] x y
\end{aligned}
$$

Using again Lemma 5 , we see that $P_{2}$ is nonnegative on $S$ if

$$
\frac{1}{2}\left(\lambda^{2}+\mu^{2}+2 \rho \lambda \mu\right)-[\lambda(\lambda \alpha+\mu \gamma)+\mu(\lambda \beta+\mu \delta)] \leq 2 \sqrt{\lambda(\lambda \beta+\mu \delta) \mu(\lambda \alpha+\mu \gamma)}
$$

This is exactly the condition $C_{3}$. Therefore

$$
0 \leq \int_{\sigma^{\epsilon}}^{t} \frac{P_{2}\left(X_{s}, Y_{s}\right)}{X_{s} Y_{s} S_{s}^{2}} d s<\infty
$$

and so

$$
0 \leq \int_{\sigma^{\epsilon}}^{\tau^{0, \epsilon}} \frac{P_{2}\left(X_{s}, Y_{s}\right)}{X_{s} Y_{s} S_{s}^{2}} d s \leq \infty
$$

Similarly, using Lemma 6 again, we see that the continuous local martingale $M$ either converges to a finite limit or oscillates between $+\infty$ and $-\infty$ when $t \rightarrow \tau^{\mathbf{0}, \epsilon}$. It cannot converge to $-\infty$ and thus $S_{\tau^{\mathbf{0}, \epsilon}}>0$ on $\left\{\tau^{\mathbf{0}, \epsilon}<\infty\right\}$, proving $\mathbb{P}\left(\tau^{\mathbf{0}, \epsilon}<\infty\right)=0$. Letting finally $\epsilon \rightarrow 0$ we obtain $\mathbb{P}\left(\tau^{0}<\infty\right)=0$.

Condition $C_{3}$ is not explicit. We give two concrete examples when this condition holds true.

Corollary 9. Assume $\rho=0, \alpha=\delta$ and $|\beta|=|\gamma|$. Then the condition $C_{3}$ is satisfied if

$$
\begin{array}{ll}
-\quad \beta^{2} \leq \alpha-\frac{1}{4} & \text { when } \quad \beta=-\gamma  \tag{10}\\
\text { - }-\beta \leq \alpha-\frac{1}{4} & \text { when } \quad \beta=\gamma<0 .
\end{array}
$$

Proof. In both cases we take $\lambda=\mu>0$.
When $\beta=-\gamma \geq 0$, the condition $C_{3}$ writes

- $\alpha-\beta \geq 0$
- $1 \leq(\sqrt{\alpha+\beta}+\sqrt{\alpha-\beta})^{2}=2 \alpha+2 \sqrt{\alpha^{2}-\beta^{2}}$.

If $\beta^{2} \leq \alpha-\frac{1}{4}$, then

$$
\alpha^{2}-\beta^{2} \geq \alpha^{2}-\alpha+\frac{1}{4} \geq 0
$$

and

$$
(1-2 \alpha)^{2}=1-4 \alpha+4 \alpha^{2} \leq 4\left(\alpha^{2}-\beta^{2}\right)
$$

When $\beta=\gamma<0$, the condition $C_{3}$ writes

- $\alpha+\beta \geq 0$
- $1 \leq(2 \sqrt{\alpha+\beta})^{2}=4(\alpha+\beta)$
and this is $-\beta \leq \alpha-\frac{1}{4}$.

Corollary 10. Assume

- $\max \{\alpha, \delta\} \geq \frac{1}{2}$
- $2 \rho \leq \frac{\beta}{\delta}+\frac{\gamma}{\alpha}$.

Then $\mathbb{P}\left(\tau^{\mathbf{0}}<\infty\right)=0$.
Proof. We may assume $\alpha \geq \frac{1}{2}$. If $\beta \geq 0$, then condition $C_{2 a}$ holds true and the conclusion follows. If $\beta<0$, we will use condition $C_{3}$. We take $\lambda=\delta$ and $\mu=-\beta$. Then $\lambda \beta+\mu \delta=0$,

$$
\begin{aligned}
\lambda \alpha+\mu \gamma & =\delta \alpha-\beta \gamma \\
& \geq \delta \alpha-\beta \alpha\left(2 \rho-\frac{\beta}{\delta}\right) \\
& =\frac{\alpha}{\delta}\left(\delta^{2}+\beta^{2}-2 \rho \beta \delta\right) \\
& \geq \frac{\alpha}{\delta}(\delta-\rho \beta)^{2} \\
& \geq 0
\end{aligned}
$$

and

$$
\begin{aligned}
& (\sqrt{\lambda(\lambda \alpha+\mu \gamma)}+\sqrt{\mu(\lambda \beta+\mu \delta)})^{2}-\frac{1}{2}\left(\lambda^{2}+\mu^{2}+2 \rho \lambda \mu\right) \\
= & \lambda(\lambda \alpha+\mu \gamma)-\frac{1}{2}\left(\lambda^{2}+\mu^{2}+2 \rho \lambda \mu\right) \\
\geq & \alpha\left(\delta^{2}+\beta^{2}-2 \rho \beta \delta\right)-\frac{1}{2}\left(\delta^{2}+\beta^{2}-2 \rho \beta \delta\right) \\
\geq & \left(\alpha-\frac{1}{2}\right)(\delta-\rho \beta)^{2} \\
\geq & 0
\end{aligned}
$$

## 5. EXISTENCE AND UNIQUENESS

We now proceed to the question of existence and uniqueness of a global solution to (1). We consider separately the three cases: first $\beta \geq 0$ and $\gamma \geq 0$, second $\beta \gamma<0$, third $\beta \leq 0$ and $\gamma \leq 0$. In the first and second cases, we construct the solution by switching from one edge to the other and patching the paths together. Thus it is essential to avoid the corner, as it was supposed in [30] in order to weakly approximate an obliquely reflected Brownian motion. The third case uses a different proof and does not requires avoiding the corner. All three proofs heavily use the comparison method of Lemma 4.

### 5.1. Case $\beta \geq 0$ and $\gamma \geq 0$.

Theorem 11. Assume $\beta \geq 0, \gamma \geq 0$ and one of the conditions $C_{1} . C_{2 a}, C_{2 b}, C_{3}$ is satisfied.
(1) There is a unique solution to (1) in $S^{\mathbf{0}}$.
(2) There is a solution to (1) in $S$ starting from $\mathbf{0}$.
(3) If $\alpha \delta \geq \beta \gamma$, there is a unique solution to (1) in $S$.

Proof. 1. Let $a>0, \epsilon>0$ and define for $(x, z) \in \mathbb{R}_{+} \times \mathbb{R}$

$$
\psi_{\epsilon}(x, z):=\frac{1}{\max (\gamma x+z, \alpha \epsilon)}
$$

This is a Lipschitz function. From Proposition 7 we know that the system

$$
\begin{align*}
X_{t}^{\epsilon} & =X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}^{\epsilon}}+\alpha \beta \int_{0}^{t} \psi_{\epsilon}\left(X_{s}^{\epsilon}, Z_{s}^{\epsilon}\right) d s \geq 0 \\
Z_{t}^{\epsilon} & =-\gamma X_{0}+\alpha\left(Y_{0}+\mathbf{1}_{\left\{Y_{0}=0\right\}} a\right)-\gamma B_{t}+\alpha C_{t}+\alpha(\alpha \delta-\beta \gamma) \int_{0}^{t} \psi_{\epsilon}\left(X_{s}^{\epsilon}, Z_{s}^{\epsilon}\right) d s \tag{11}
\end{align*}
$$

has a unique solution. Let

$$
\tau_{Y}^{\epsilon}:=\inf \left\{t>0: \gamma X_{t}^{\epsilon}+Z_{t}^{\epsilon}<\alpha \epsilon\right\}
$$

If $0<\eta<\epsilon<a$ we deduce from the uniqueness that $\left(X^{\epsilon}, Z^{\epsilon}\right)$ and $\left(X^{\eta}, Z^{\eta}\right)$ are identical on $\left[0, \tau_{Y}^{\epsilon}\right]$. Patching together we can set

$$
\begin{aligned}
X_{t} & :=\lim _{\epsilon \rightarrow 0} X_{t}^{\epsilon} \\
Y_{t} & :=\lim _{\epsilon \rightarrow 0} \frac{1}{\alpha}\left(\gamma X_{t}^{\epsilon}+Z_{t}^{\epsilon}\right)
\end{aligned}
$$

on $\left\{(\omega, t) \in \Omega \times[0, \infty): Y_{0}(\omega)>0\right.$ and $\left.0 \leq \mathrm{t}<\tau_{\mathrm{Y}}^{0}(\omega)\right\}$, where

$$
\tau_{Y}^{0}:=\lim _{\epsilon \rightarrow 0} \tau_{Y}^{\epsilon}
$$

On this set, $(X, Y)$ is the unique solution to (1). As we already noted, we have $X_{t} \geq U_{t}$ and $Y_{t} \geq V_{t}$. Therefore, on $\left\{Y_{0}>0\right\} \cap\left\{\tau_{Y}^{0}<\infty\right\}$,

$$
\int_{0}^{\tau_{Y}^{0}} \frac{d s}{X_{s}} \leq \int_{0}^{\tau_{Y}^{0}} \frac{d s}{U_{s}}<\infty \quad \text { and } \quad \int_{0}^{\tau_{Y}^{0}} \frac{d s}{Y_{s}} \leq \int_{0}^{\tau_{Y}^{0}} \frac{d s}{V_{s}}<\infty
$$

and we can define

$$
\begin{align*}
& X_{\tau_{Y}^{0}}:=\lim _{t \rightarrow \tau_{Y}^{0}} X_{t}=X_{0}+B_{\tau_{Y}^{0}}+\alpha \int_{0}^{\tau_{Y}^{0}} \frac{d s}{X_{s}}+\beta \int_{0}^{\tau_{Y}^{0}} \frac{d s}{Y_{s}} \\
& Y_{\tau_{Y}^{0}}:=\lim _{t \rightarrow \tau_{Y}^{0}} Y_{t}=Y_{0}+{\tau_{Y}^{0}}^{\tau_{Y}^{0}}+\gamma \int_{0}^{\tau_{s}}+\delta \int_{0}^{\tau_{Y}^{0}} \frac{d s}{Y_{s}} \tag{12}
\end{align*}
$$

We have $Y_{\tau_{Y}^{0}}=0$. From Theorem 8 we know that $X_{\tau_{Y}^{0}}>0$. In exactly the same way we can construct a solution on $\left\{Y_{0}>0\right\}$ in the interval $\left[T_{1}, T_{2}\right]$, where $T_{1}=\tau_{Y}^{0}, T_{2}=\inf \left\{t>T_{1}\right.$ : $\left.X_{t}=0\right\}$. Iterating, we get a solution on $\left\{Y_{0}>0\right\} \times\left[0, \lim _{n \rightarrow \infty} T_{n}\right)$ where

$$
\begin{array}{ll}
T_{2 p} & :=\inf \left\{t>T_{2 p-1}: X_{t}=0\right\} \\
T_{2 p+1} & :=\inf \left\{t>T_{2 p}: Y_{t}=0\right\}
\end{array}
$$

On $\left\{Y_{0}>0\right\} \cap\left\{\lim _{n \rightarrow \infty} T_{n}<\infty\right\}$ we set $X_{\lim _{n \rightarrow \infty} T_{n}}:=\lim _{p \rightarrow \infty} X_{T_{2 p}}=0$ and $Y_{\lim _{n \rightarrow \infty} T_{n}}:=$ $\lim _{p \rightarrow \infty} Y_{T_{2 p+1}}=0$. The polarity of $\mathbf{0}$ entails this is not possible in finite time and thus $\lim _{n \rightarrow \infty} T_{n}=\infty$. So we have obtained a unique global solution on $\left\{Y_{0}>0\right\}$. In the same way we obtain a unique global solution on $\left\{X_{0}>0\right\}$ and as $\mathbb{P}\left(\left(X_{0}, Y_{0}\right)=\mathbf{0}\right)=0$ the proof is complete.
2. Assume now $X_{0}=Y_{0}=0$. Let $\left(y_{n}\right)_{n \geq 1}$ be a sequence of real numbers (strictly) decreasing to 0 . From the above paragraph it follows there exists for any $n \geq 1$ a unique solution ( $X^{n}, Y^{n}$ ) with values in $S^{0}$ to the system

$$
\begin{aligned}
X_{t}^{n} & =B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}^{n}}+\beta \int_{0}^{t} \frac{d s}{Y_{s}^{n}} \\
Y_{t}^{n} & =y_{n}+C_{t}+\gamma \int_{0}^{t} \frac{d s}{X_{s}^{n}}+\delta \int_{0}^{t} \frac{d s}{Y_{s}^{n}} .
\end{aligned}
$$

Let

$$
\tau:=\inf \left\{t>0: X_{t}^{n+1}<X_{t}^{n}\right\}
$$

Using Lemma 4 we obtain $Y_{t}^{n+1} \leq Y_{t}^{n}$ on $[0, \tau]$. We note that $\left(X_{\tau}^{n}, Y_{\tau}^{n}\right) \in S^{\mathbf{0}}$ on $\{\tau<\infty\}$. On $\left\{Y_{\tau}^{n+1}=Y_{\tau}^{n}\right\} \cap\{\tau<\infty\}$, since $X_{\tau}^{n+1}=X_{\tau}^{n}$ and the solution starting at time $\tau$ is unique, it follows that $X_{t}^{n+1}=X_{t}^{n}$ and $Y_{t}^{n+1}=Y_{t}^{n}$ on $[\tau, \infty)$. On $\left\{Y_{\tau}^{n+1}<Y_{\tau}^{n}\right\} \cap\{\tau<\infty\}$, the continuity of solutions at time $\tau$ entails there exists $\rho>0$ such that $Y_{t}^{n+1} \leq Y_{t}^{n}$ on $[\tau, \tau+\rho]$. A second application of Lemma 4 proves that $X_{t}^{n+1} \geq X_{t}^{n}$ on $[\tau, \tau+\rho]$, a contradiction to the definition of $\tau$. Therefore $\mathbb{P}(\tau=\infty)=1$. It follows that $X_{t}^{n+1} \geq X_{t}^{n}$ and $Y_{t}^{n+1} \leq Y_{t}^{n}$ for any $t \in[0, \infty)$, and we may define

$$
X_{t}:=\lim _{n \rightarrow \infty} \uparrow X_{t}^{n} \quad Y_{t}:=\lim _{n \rightarrow \infty} \downarrow Y_{t}^{n}
$$

As $Y_{t}^{n} \geq V_{t}$ where $(U, V)$ is the solution to (2) with $X_{0}=Y_{0}=0$, we have

$$
\begin{aligned}
X_{t} & =B_{t}+\alpha \lim _{n \rightarrow \infty} \int_{0}^{t} \frac{d s}{X_{s}^{n}}+\beta \lim _{n \rightarrow \infty} \int_{0}^{t} \frac{d s}{Y_{s}^{n}} \\
& =B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}}+\beta \int_{0}^{t} \frac{d s}{Y_{s}} \\
& <\infty
\end{aligned}
$$

and also

$$
\begin{aligned}
Y_{t} & =\lim _{n \rightarrow \infty} y_{n}+C_{t}+\gamma \lim _{n \rightarrow \infty} \int_{0}^{t} \frac{d s}{X_{s}^{n}}+\delta \lim _{n \rightarrow \infty} \int_{0}^{t} \frac{d s}{Y_{s}^{n}} \\
& =C_{t}+\gamma \int_{0}^{t} \frac{d s}{X_{s}}+\delta \int_{0}^{t} \frac{d s}{Y_{s}} \\
& <\infty
\end{aligned}
$$

3. Assume finally $\alpha \delta-\beta \gamma \geq 0$. As the conclusion holds true if $\beta=\gamma=0$, we may also assume $\beta>0$. Let $(X, Y)$ be the solution to (1) with $X_{0}=Y_{0}=0$ obtained in the previous paragraph and let $\left(X^{\prime}, Y^{\prime}\right)$ be another solution. Considering ( $X^{n}, Y^{n}$ ) again and replacing $\left(X^{n+1}, Y^{n+1}\right)$ with $\left(X^{\prime}, Y^{\prime}\right)$, the previous proof works and we finally obtain $X_{t}^{\prime} \geq X_{t}$ and $Y_{t}^{\prime} \leq Y_{t}$. Then,

$$
\begin{equation*}
0 \leq \delta\left(X_{t}^{\prime}-X_{t}\right)-\beta\left(Y_{t}^{\prime}-Y_{t}\right)=\int_{0}^{t}(\alpha \delta-\beta \gamma)\left(\frac{1}{X_{s}^{\prime}}-\frac{1}{X_{s}}\right) d s \leq 0 \tag{13}
\end{equation*}
$$

Thus $X_{t}^{\prime}=X_{t}$ and $Y_{t}^{\prime}=Y_{t}$, proving uniqueness. Replacing $(\delta, \beta)$ with $(\gamma, \alpha)$ in equation (13) we obtain the same conclusion if $\gamma>0$.

Remark 12. The statement in Theorem 11 is not complete since the problem of uniqueness when starting at the corner and $\alpha \delta<\beta \gamma$ is not solved. When considering the solution $(X, Y)$ in the above proof of existence, we have noted that $X^{\prime} \geq X$ and $Y^{\prime} \leq Y$ for any other solution $\left(X^{\prime}, Y^{\prime}\right)$. Thus uniqueness in law would be sufficient to obtain path uniqueness. A possible way to prove weak uniqueness could be the method in [1]. This would be far from our trajectorial methods and we don't go further in that direction.

### 5.2. Case $\beta \gamma<0$.

Theorem 13. Assume $\beta \gamma<0$ and one of the conditions $C_{2 a}$ or $C_{3}$ is satisfied. Then, there exists a unique solution to (1) in $S^{\mathbf{0}}$.

Proof. Assume first $\beta>0, \gamma<0$. The proof is similar to the proof of 1 in Theorem 11. The only change is that now $Y_{t} \leq V_{t}$. Therefore, on $\left\{Y_{0}>0\right\} \cap\left\{\tau_{Y}^{0}<\infty\right\}$,

$$
\delta \int_{0}^{\tau_{Y}^{0}} \frac{d s}{Y_{s}} \leq V_{\tau_{Y}^{0}}-Y_{0}-C_{\tau_{Y}^{0}}-\gamma \int_{0}^{\tau_{Y}^{0}} \frac{d s}{U_{s}}<\infty
$$

and we can define $X_{\tau_{Y}^{0}}$ and $Y_{\tau_{Y}^{0}}$ as previously done. The application of Theorem 8 to the process on the time interval $\left[0, \tau_{Y}^{0}\right]$ shows that $X_{\tau_{Y}^{0}}>0$ and we can iterate the construction as in Theorem 11. The proof if $\beta<0, \gamma>0$ is analogous.
5.3. Case $\beta \leq 0$ and $\gamma \leq 0$. In this case we can give a full answer to the question of existence and uniqueness. When $|\rho|<1$, our condition of existence is analogous to the condition found in [37] for the reflected Brownian in a wedge being a semimartingale, i.e. there is a convex combination of the directions of reflection that points into the wedge from the corner. It amounts to saying that the interaction matrix

$$
\left(\begin{array}{ll}
\alpha & \beta \\
\gamma & \delta
\end{array}\right)
$$

is completely- $\mathcal{S}$ in the terminology of $[35,5,40,4,29,30,31]$.
Theorem 14. Assume $\beta \leq 0$ and $\gamma \leq 0$.
(1) If $\alpha \delta>\beta \gamma$, there exists a unique solution to (1) in $S$.
(2) If $\alpha \delta \leq \beta \gamma$ and $1+\rho+|\alpha+\gamma|+|\beta+\delta|>0$, there is no solution.
(3) If $1+\rho=\alpha+\gamma=\beta+\delta=0$ and $\left(X_{0}, Y_{0}\right) \neq \mathbf{0}$ there exists a unique solution.
(4) If $1+\rho=\alpha+\gamma=\beta+\delta=0$ and $\left(X_{0}, Y_{0}\right)=\mathbf{0}$ there is no solution.

Proof. 1. Assume first $\alpha \delta>\beta \gamma$.
a) Existence. Let $\left(h_{n}, n \geq 1\right)$ be a (strictly) increasing sequence of bounded positive nonincreasing Lipschitz functions converging to $1 / x$ on $(0, \infty)$ and to $+\infty$ on $(-\infty, 0]$. For instance we can take

$$
\begin{aligned}
h_{n}(x) & =\left(1-\frac{1}{n}\right) \frac{1}{x} & & \text { on } \quad\left[\frac{1}{n}, \infty\right) \\
& =n-1 & & \text { on } \quad\left(-\infty, \frac{1}{n}\right] .
\end{aligned}
$$

We consider for each $n \geq 1$ the system

$$
\begin{align*}
X_{t}^{n} & =X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}^{n}}+\beta \int_{0}^{t} h_{n}\left(Y_{s}^{n}\right) d s \\
Y_{t}^{n} & =Y_{0}+C_{t}+\gamma \int_{0}^{t} h_{n}\left(X_{s}^{n}\right) d s+\delta \int_{0}^{t} \frac{d s}{Y_{s}^{n}} \tag{14}
\end{align*}
$$

From Proposition 7 it follows there exists a unique solution to this system. We set

$$
\tau:=\inf \left\{s>0: X_{s}^{n+1}>X_{s}^{n}\right\}
$$

We have $h_{n+1}\left(X_{t}^{n+1}\right) \geq h_{n}\left(X_{t}^{n}\right)$ on $[0, \tau]$. A first application of Lemma 4 shows that $Y_{t}^{n+1} \leq$ $Y_{t}^{n}$ on $[0, \tau]$. Since $h_{n+1}\left(Y_{\tau}^{n+1}\right)>h_{n}\left(Y_{\tau}^{n}\right)$ on $\{\tau<\infty\}$, we deduce from the continuity of solutions that there exists $\sigma>0$ such that $h_{n+1}\left(Y_{t}^{n+1}\right) \geq h_{n}\left(Y_{t}^{n}\right)$ on $[\tau, \tau+\sigma]$. A second application of Lemma 4 shows that $X_{t}^{n+1} \leq X_{t}^{n}$ on $[\tau, \tau+\sigma]$, a contradiction to the definition of $\tau$. Thus $\mathbb{P}(\tau=\infty)=1$ proving that on the whole $[0, \infty)$ we have $X_{t}^{n+1} \leq X_{t}^{n}$ and $Y_{t}^{n+1} \leq Y_{t}^{n}$. Then we can set for any $t \in[0, \infty)$

$$
X_{t}:=\lim _{n \rightarrow \infty} X_{t}^{n} \quad \text { and } \quad Y_{t}:=\lim _{n \rightarrow \infty} Y_{t}^{n} .
$$

If $\alpha \delta>\beta \gamma$, there is a convex combination of the directions of repulsion pointing into the positive quadrant, i.e. there exist $\lambda>0$ and $\mu>0$ such that $\lambda \alpha+\mu \gamma>0$ and $\mu \delta+\lambda \beta>0$. For $n \geq 1$ and $t \geq 0$,

$$
\begin{align*}
\lambda U_{t}+\mu V_{t} & \geq \lambda X_{t}^{n}+\mu Y_{t}^{n} \\
& \geq \lambda X_{0}+\mu Y_{0}+\lambda B_{t}+\mu C_{t}+(\lambda \alpha+\mu \gamma) \int_{0}^{t} \frac{d s}{X_{s}^{n}}+(\mu \delta+\lambda \beta) \int_{0}^{t} \frac{d s}{Y_{s}^{n}} \tag{15}
\end{align*}
$$

Letting $n \rightarrow \infty$ in (15) we obtain

$$
\int_{0}^{t} \frac{d s}{X_{s}}<\infty \quad \text { and } \quad \int_{0}^{t} \frac{d s}{Y_{s}}<\infty
$$

Then we may let $n$ go to $\infty$ in (14) proving that $(X, Y)$ is a solution to (1).
b) Uniqueness. Let $\left(X^{\prime}, Y^{\prime}\right)$ be another solution to (1). Replacing $\left(X^{n+1}, Y^{n+1}\right)$ with $\left(X^{\prime}, Y^{\prime}\right)$ we follow the above proof to obtain for $t \in[0, \infty)$ and $n \geq 1$

$$
X_{t}^{\prime} \leq X_{t}^{n} \quad \text { and } \quad Y_{t}^{\prime} \leq Y_{t}^{n}
$$

Letting $n \rightarrow \infty$ we conclude

$$
X_{t}^{\prime} \leq X_{t} \quad \text { and } \quad Y_{t}^{\prime} \leq Y_{t}
$$

With the same $\lambda>0$ and $\mu>0$ as above,

$$
0 \leq \lambda\left(X_{t}-X_{t}^{\prime}\right)+\mu\left(Y_{t}-Y_{t}^{\prime}\right)=\int_{0}^{t}\left[(\lambda \alpha+\mu \gamma)\left(\frac{1}{X_{s}}-\frac{1}{X_{s}^{\prime}}\right)+(\mu \delta+\lambda \beta)\left(\frac{1}{Y_{s}}-\frac{1}{Y_{s}^{\prime}}\right)\right] d s \leq 0
$$

and therefore $X_{t}^{\prime}=X_{t}, Y_{t}^{\prime}=Y_{t}$.
2. If $\alpha \delta \leq \beta \gamma$ there exist $\lambda>0$ and $\mu>0$ such that $\lambda \alpha+\mu \gamma \leq 0$ and $\mu \delta+\lambda \beta \leq 0$. For that, just take

$$
\begin{equation*}
\frac{\alpha}{-\gamma} \leq \frac{\mu}{\lambda} \leq \frac{-\beta}{\delta} \tag{16}
\end{equation*}
$$

Let us consider the nonnegative quadratic $\left(\lambda^{2}+\mu^{2}+2 \rho \lambda \mu\right)$. It is positive if $\rho \geq 0$. It is larger than $\left(1-\rho^{2}\right) \lambda^{2}>0$ if $-1<\rho<1$. It may be positive if $\rho=-1$ and $|\alpha+\gamma|+|\beta+\delta|>0$ because from (16) we may take $\lambda \neq \mu$. Thus, if $(X, Y)$ is a solution to (1), for any $t \geq 0$,

$$
0 \leq \lambda X_{t}+\mu Y_{t} \leq \lambda X_{0}+\mu Y_{0}+\lambda B_{t}+\mu C_{t}
$$

This is not possible since the paths of the Brownian martingale $\left(\lambda^{2}+\mu^{2}+2 \rho \lambda \mu\right)^{-1 / 2}\left(\lambda B_{t}+\mu C_{t}\right)$ are not bounded below. So there is no global solution.
3. Assume now $1+\rho=|\alpha+\gamma|+|\beta+\delta|=0$ and $X_{0}+Y_{0}>0$. The system becomes

$$
\begin{align*}
& X_{t}=X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}}-\delta \int_{0}^{t} \frac{d s}{Y_{s}} \geq 0  \tag{17}\\
& Y_{t}=Y_{0}-B_{t}-\alpha \int_{0}^{t} \frac{d s}{X_{s}}+\delta \int_{0}^{t} \frac{d s}{Y_{s}} \geq 0
\end{align*}
$$

This entails for any $t \geq 0$

$$
X_{t}+Y_{t}=X_{0}+Y_{0}
$$

and the first equation in (17) reduces to

$$
\begin{equation*}
0 \leq X_{t}=X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}}-\delta \int_{0}^{t} \frac{d s}{X_{0}+Y_{0}-X_{s}} \leq X_{0}+Y_{0} \tag{18}
\end{equation*}
$$

Clearly pathwise uniqueness holds for equation (18) since if there are two solutions $X$ and $X^{\prime}$, for any $t \geq 0$

$$
\begin{aligned}
\left(X_{t}-X_{t}^{\prime}\right)^{2} & =2 \alpha \int_{0}^{t}\left(X_{s}-X_{s}^{\prime}\right)\left(\frac{1}{X_{s}}-\frac{1}{X_{s}^{\prime}}\right) d s-2 \delta \int_{0}^{t}\left(X_{s}-X_{s}^{\prime}\right)\left(\frac{1}{X_{0}+Y_{0}-X_{s}}-\frac{1}{X_{0}+Y_{0}-X_{s}^{\prime}}\right) d s \\
& \leq 0 .
\end{aligned}
$$

Consider now for any $c>0$ and $0 \leq c_{0} \leq c$ the equation

$$
\begin{equation*}
0 \leq Z_{t}=c_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{Z_{s}}-\delta \int_{0}^{t} \frac{d s}{c-Z_{s}} \leq c \tag{19}
\end{equation*}
$$

The solution $Z$ is a Brownian motion perturbed by a drift deriving from the concave potential

$$
\alpha \ln z+\delta \ln (c-z) \quad \text { for } \quad 0<z<c .
$$

From [2] we know that equation (19) has a unique solution living on the interval $[0, c]$. Thus there is a weak and then a unique strong solution to (18). Setting

$$
Y_{t}=X_{0}+Y_{0}-X_{t}
$$

we obtain a unique strong solution to (17).
4. We must have for any $t \geq 0$

$$
X_{t}+Y_{t}=X_{0}+Y_{0}=0
$$

and thus $X_{t}=Y_{t}=0$. But this is not possible for a solution to (1).
We end this section by considering a degenerate case where the solution hits the corner with probability one.

Proposition 15. Assume $\rho=1, \alpha \delta>\beta \gamma$, $\max \{\alpha, \delta\}<\frac{1}{2}$ and $\max \{\beta, \gamma\} \leq 0$. Then, $\mathbb{P}\left(\tau^{0}<\infty\right)=1$.

Proof. We set

$$
\begin{aligned}
\tau_{X}^{0} & :=\inf \left\{t>0: X_{t}=0\right\} \\
\tau_{Y}^{0} & :=\inf \left\{t>0: Y_{t}=0\right\}
\end{aligned}
$$

The dimension of each Bessel process $U$ and $V$ in (2) is less than 2 , and $X \leq U, Y \leq V$. Then $\mathbb{P}\left(\tau_{X}^{0}<\infty=1\right)$ and $\mathbb{P}\left(\tau_{Y}^{0}<\infty=1\right)$. Assume first $\alpha+\beta \leq \gamma+\delta$. Using Lemma 4 we obtain $X_{t} \leq Y_{t} \leq V_{t}$ on the time interval $\left[\tau_{X}^{0}, \infty\right)$. Therefore $\mathbb{P}\left(\tau^{\mathbf{0}}<\infty\right)=1$. Same result when $\alpha+\beta \geq \gamma+\delta$.

In the following pictures, we display the directions of interaction $\mathbf{r}_{\mathbf{x}}$ and $\mathbf{r}_{\mathbf{y}}$ defined in (4) in three illustrative instances.

$\beta>0, \gamma>0, \alpha \delta>\beta \gamma$

$\beta<0, \gamma>0, \alpha \beta+\gamma \delta=0$

$\beta<0, \gamma<0, \alpha \delta>\beta \gamma$

## 6. Avoiding The edges

We now consider the question of hitting an edge of the quadrant. Remember the definitions

$$
\begin{align*}
\tau_{X}^{0} & :=\inf \left\{t>0: X_{t}=0\right\}  \tag{20}\\
\tau_{Y}^{0} & :=\inf \left\{t>0: Y_{t}=0\right\}
\end{align*}
$$

We already know that $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=0$ if $\alpha \geq \frac{1}{2}$ and $\beta \geq 0$. Conversely, a comparison with the Bessel process $U$ shows that $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=1$ if $\alpha<\frac{1}{2}$ and $\beta \leq 0$. If we know that the corner is not hit and $\alpha \geq \frac{1}{2}$, we can get rid of the nonnegativity assumption on $\beta$. Since we are only interested in one coordinate, we are looking for a function that is $C^{2}$ on $(0, \infty)$ and goes to $-\infty$ when approaching 0 . A natural candidate is the logarithmic function.
Proposition 16. Assume $\mathbb{P}\left(\tau^{\mathbf{0}}<\infty\right)=0$. If $\alpha \geq \frac{1}{2}$, then $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=0$.
Proof. For $\eta>0$ let

$$
\begin{aligned}
\theta_{X}^{\eta} & =\mathbf{1}_{\left\{X_{0}=0\right\}} \inf \left\{t>0: X_{t} \geq \eta\right\} \\
\tau_{X}^{0, \eta} & =\inf \left\{t>\theta_{X}^{\eta}: X_{t}=0\right\}
\end{aligned}
$$

As $\eta \downarrow 0, \theta_{X}^{\eta} \downarrow 0$ and $\tau_{X}^{0, \eta} \downarrow \tau_{X}^{0}$. For $t \in\left[\theta_{X}^{\eta}, \tau_{X}^{0, \eta}\right)$, from Itô's formula

$$
\begin{equation*}
\ln X_{t}=\ln X_{\theta_{X}^{\eta}}+\int_{\theta_{X}^{\eta}}^{t} \frac{d B_{s}}{X_{s}}+\left(\alpha-\frac{1}{2}\right) \int_{\theta_{X}^{\eta}}^{t} \frac{d s}{X_{s}^{2}}+\beta \int_{\theta_{X}^{\eta}}^{t} \frac{d s}{X_{s} Y_{s}} \tag{21}
\end{equation*}
$$

Since $\mathbb{P}\left(\tau^{0}<\infty\right)=0$, on the set $\left\{\tau_{X}^{0, \eta}<\infty\right\}$ we have $Y_{\tau_{X}^{0, \eta}}>0$ because $X_{\tau_{X}^{0, \eta}}=0$. On this set, from the definition of an O2BP,

$$
\int_{\theta_{X}^{\eta}}^{\tau_{X}^{0, \eta}} \frac{d s}{X_{s}}<\infty, \quad \int_{\theta_{X}^{\eta}}^{\tau_{X}^{0, \eta}} \frac{d s}{Y_{s}}<\infty
$$

As $Y_{s}>0$ on some interval $\left[\chi, \tau_{X}^{0, \eta}\right]$ with positive measure and $X_{s}>0$ on $\left[\theta_{X}^{\eta}, \chi\right]$, we see that

$$
|\beta| \int_{\theta_{X}^{\eta}}^{\tau_{X}^{0, \eta}} \frac{d s}{X_{s} Y_{s}}<\infty
$$

As $t \rightarrow \tau_{X}^{0, \eta}$, the local martingale in the r.h.s. of (21) cannot converge to $-\infty$. This entails that $\mathbb{P}\left(\tau_{X}^{0, \eta}<\infty\right)=0$ and therefore $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=0$.

We may also be interested in hitting either edge. This time we are looking for a function that is $C^{2}$ in the interior of $S$ and goes to $-\infty$ when approaching either edge.

Proposition 17. Assume $\alpha \geq \frac{1}{2}$ and $\delta \geq \frac{1}{2}$. Then $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=\mathbb{P}\left(\tau_{Y}^{0}<\infty\right)=0$ if one of the following conditions is satisfied:
(1) $\beta \geq 0$
(2) $\gamma \geq 0$
(3) $0<\beta \gamma \leq\left(\alpha-\frac{1}{2}\right)\left(\delta-\frac{1}{2}\right)$.

Proof. The proofs in the cases (1) and (2)) are direct consequences of Theorem 8 (under conditions $C_{2 a}$ or $C_{2 b}$ ) and Proposition 16. Assume now the conditions in (3) hold true. For $\epsilon>0$ let

$$
\begin{aligned}
\sigma^{\epsilon} & =1_{\left\{X_{0} Y_{0}=0\right\}} \inf \left\{t>0: X_{t} Y_{t} \geq \epsilon\right\} \\
\tau^{\epsilon} & =\inf \left\{t>\sigma^{\epsilon}: X_{t} Y_{t}=0\right\}
\end{aligned}
$$

For $\lambda>0$ and $\mu>0$ we set

$$
R_{t}=\lambda \ln X_{t}+\mu \ln Y_{t}
$$

From Itô's formula we get for $t \in\left[\sigma^{\epsilon}, \tau^{\epsilon}\right)$

$$
\begin{aligned}
R_{t} & =R_{\sigma^{\epsilon}}+\int_{\sigma^{\epsilon}}^{t}\left(\frac{\lambda}{X_{s}} d B_{s}+\frac{\mu}{Y_{s}} d C_{s}\right)+\int_{\sigma^{\epsilon}}^{t}\left[\frac{\lambda\left(\alpha-\frac{1}{2}\right)}{X_{s}^{2}}+\frac{\mu\left(\delta-\frac{1}{2}\right)}{Y_{s}^{2}}+\frac{(\lambda \beta+\mu \gamma)}{X_{s} Y_{s}}\right] d s \\
& =R_{\sigma^{\epsilon}}+N_{t}+\int_{\sigma^{\epsilon}}^{t} \frac{P_{3}\left(X_{s}, Y_{s}\right)}{X_{s}^{2} Y_{s}^{2}} d s
\end{aligned}
$$

where $N$ is a continuous local martingale and $P_{3}(x, y)$ is the second degree homogeneous polynomial

$$
P_{3}(x, y)=\mu\left(\delta-\frac{1}{2}\right) x^{2}+\lambda\left(\alpha-\frac{1}{2}\right) y^{2}+(\lambda \beta+\mu \gamma) x y
$$

If the conditions (3) are satisfied, we may take

$$
\begin{array}{ll}
\lambda=2\left(\alpha-\frac{1}{2}\right)\left(\delta-\frac{1}{2}\right)-\beta \gamma & >0 \\
\mu=\beta^{2} & >0
\end{array}
$$

Then, using again conditions (3), we check that
$P_{3}(x, y)=\beta^{2}\left(\delta-\frac{1}{2}\right) x^{2}+\left[2\left(\alpha-\frac{1}{2}\right)\left(\delta-\frac{1}{2}\right)-\beta \gamma\right]\left(\alpha-\frac{1}{2}\right) y^{2}+2\left(\alpha-\frac{1}{2}\right)\left(\delta-\frac{1}{2}\right) \beta x y$ is nonnegative on $S$. The proof terminates as previously in Theorem 8 .

## 7. Associated local martingales

We easily check that if $\max \{\beta, \gamma, \alpha \delta-\beta \gamma\} \geq 0$ and if $\rho>-1$, there exist some $\lambda \geq 0$ and $\mu \geq 0$ with $\lambda+\mu>0$ such that $\lambda X_{t}+\mu Y_{t}$ is not less than $\lambda\left(X_{0}+B_{t}\right)+\mu\left(Y_{0}+C_{t}\right)$ which is proportional to a real driftless Brownian motion. Therefore

$$
\limsup _{t \rightarrow \infty}\left(\lambda X_{t}+\mu Y_{t}\right)=+\infty
$$

An usual way in the study the asymptotic behavior of the solutions to stochastic differential equations is introducing associated martingales. This is carried out through scale functions ([28] Section VII.3). There is no equivalent functions on the plane. However, in some particular cases we can find simple functions of O2BPs that are supermartingales or local martingales.

Proposition 18. Assume the following set of conditions:

$$
\begin{align*}
& \alpha>1 / 2 \\
& \delta>1 / 2 \\
& \frac{\beta}{2 \delta-1}+\frac{\gamma}{2 \alpha-1} \geq \rho  \tag{22}\\
& \rho>-1
\end{align*}
$$

Then

$$
M_{t}:=X_{t}^{1-2 \alpha} Y_{t}^{1-2 \delta}
$$

is a positive supermartingale on $(0, \infty)$ which tends to 0 as $t \rightarrow \infty$. It is a local martingale if the third inequality in (22) is an equality.

Proof. For $\epsilon>0$ let

$$
\begin{aligned}
\sigma^{\epsilon} & =\mathbf{1}_{\left\{X_{0} Y_{0}=0\right\}} \inf \left\{t>0: X_{t} Y_{t} \geq \epsilon\right\} \\
\tau^{\epsilon} & =\inf \left\{t>\sigma^{\epsilon}: X_{t} Y_{t}=0\right\}
\end{aligned}
$$

As $\epsilon \downarrow 0, \sigma^{\epsilon} \downarrow 0$. Up to $\tau^{\epsilon}$, we get from Itô's formula on $\left\{\sigma^{\epsilon}<\infty\right\}$

$$
\begin{aligned}
M_{t}= & M_{\sigma^{\epsilon}}+\int_{\sigma^{\epsilon}}^{t} \frac{(1-2 \alpha) Y_{s} d B_{s}+(1-2 \delta) X_{s} d C_{s}}{X_{s}^{2 \alpha} Y_{s}^{2 \delta}} \\
& +[(1-2 \alpha) \beta+(1-2 \delta) \gamma+\rho(2 \alpha-1)(2 \delta-1)] \int_{\sigma^{\epsilon}}^{t} \frac{d s}{X_{s}^{2 \alpha} Y_{s}^{2 \delta}}
\end{aligned}
$$

The function $f(x, y)=x^{1-2 \alpha} y^{1-2 \delta}$ is $C^{2}$ in the interior of $S$ and goes to $+\infty$ when approaching the edges of the quadrant, whereas the finite variation part in the semimartingale decomposition of $M_{t}$ is nonincreasing. We apply again Lemma 6 and obtain that $\mathbb{P}\left(\tau^{\epsilon}<\infty\right)=0$. Then, letting $\epsilon \rightarrow 0$, we see that $M$ is a positive supermartingale on $(0, \infty)$. As such it tends to a limit $H \geq 0$ when $t \rightarrow \infty$ ([28], Corollary II.2.11 ). If $\beta \geq 0$ and $\gamma \geq 0$, then $X \geq U$ and $Y \geq V$ where $U$ and $V$ are the Bessel processes in (2). We have seen in Section 2 that $U_{t} \rightarrow \infty$ and $V_{t} \rightarrow \infty$ as $t \rightarrow \infty$; so $M_{t} \rightarrow 0$ as $t \rightarrow \infty$. If now $\min \{\beta, \gamma\}<0$, say $\gamma<0$, then we have $Y \leq V$. From Corollary 3 we deduce that for any $\epsilon>0$

$$
\begin{equation*}
\int_{\sigma^{\epsilon}}^{\infty} \frac{d s}{Y_{s}^{2}} \geq \int_{\sigma^{\epsilon}}^{\infty} \frac{d s}{V_{s}^{2}}=\infty \tag{23}
\end{equation*}
$$

We consider the quadratic variation $\langle M\rangle^{\epsilon}$ given by

$$
\begin{aligned}
\langle M\rangle_{t}^{\epsilon} & =\int_{\sigma^{\epsilon}}^{t} X_{s}^{-4 \alpha} Y_{s}^{-4 \delta}\left[(1-2 \alpha)^{2} Y_{s}^{2}+(1-2 \delta)^{2} X_{s}^{2}+2 \rho(1-2 \alpha)(1-2 \delta) X_{s} Y_{s}\right] d s \\
& =\int_{\sigma^{\epsilon}}^{t} M_{s}^{2}\left[\frac{(1-2 \alpha)^{2}}{X_{s}^{2}}+\frac{(1-2 \delta)^{2}}{Y_{s}^{2}}+2 \rho \frac{(1-2 \alpha)(1-2 \delta)}{X_{s} Y_{s}}\right] d s
\end{aligned}
$$

If $\rho \geq 0$,

$$
\langle M\rangle_{t}^{\epsilon} \geq \int_{\sigma^{\epsilon}}^{t} M_{s}^{2} \frac{(1-2 \delta)^{2}}{Y_{s}^{2}} d s
$$

and if $\rho^{2}<1$,

$$
\langle M\rangle_{t}^{\epsilon} \geq\left(1-\rho^{2}\right) \int_{\sigma^{\epsilon}}^{t} M_{s}^{2} \frac{(1-2 \delta)^{2}}{Y_{s}^{2}} d s
$$

In both cases, using (23) we should have $\langle M\rangle_{\infty}^{\epsilon}=\infty$ on the set $\{H>0\} \cap\left\{\sigma^{\epsilon}<\infty\right\}$, and then $\limsup \sup _{t} M_{t}=-\liminf _{t} M_{t}=\infty$, a contradiction with $M_{t} \rightarrow H$. Thus $H=0$ a.s.

During the proof we have seen that $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=\mathbb{P}\left(\tau_{Y}^{0}<\infty\right)=0$. In fact this is not a new result since here the conditions in Proposition 17 are in force. There is another case of interest.

Proposition 19. Assume $\alpha=\delta=1 / 2, \rho \beta \gamma<|\beta \gamma|$, and there exists a solution to (1) satisfying $\mathbb{P}\left(\tau^{\mathbf{0}}<\infty\right)=0$. Then

$$
M_{t}:=\gamma \ln X_{t}-\beta \ln Y_{t}
$$

is a continuous local martingale on $(0, \infty)$ and

$$
\limsup _{t \rightarrow \infty} M_{t}=-\liminf _{t \rightarrow \infty} M_{t}=\infty
$$

Proof. We deduce from Proposition 17 that $X_{t}>0$ and $Y_{t}>0$ for any $t>0$. Let again for $\epsilon>0$

$$
\sigma^{\epsilon}=\mathbf{1}_{\left\{X_{0} Y_{0}=0\right\}} \inf \left\{t>0: X_{t} Y_{t} \geq \epsilon\right\}
$$

Now Itô's formula gives

$$
M_{t}=M_{\sigma^{\epsilon}}+\int_{\sigma^{\epsilon}}^{t}\left(\gamma \frac{d B_{s}}{X_{s}}-\beta \frac{d C_{s}}{Y_{s}}\right)
$$

and the quadratic variation is

$$
\langle M\rangle_{t}^{\epsilon}=\int_{\sigma^{\epsilon}}^{t}\left(\frac{\gamma^{2}}{X_{s}^{2}}+\frac{\beta^{2}}{Y_{s}^{2}}-2 \rho \frac{\beta \gamma}{X_{s} Y_{s}}\right) d s
$$

Then, if $\rho \beta \gamma \leq 0$,

$$
\langle M\rangle_{t}^{\epsilon} \geq \int_{\sigma^{\epsilon}}^{t}\left(\frac{\gamma^{2}}{X_{s}^{2}}+\frac{\beta^{2}}{Y_{s}^{2}}\right) d s
$$

and if $\rho^{2}<1$,

$$
\langle M\rangle_{t}^{\epsilon} \geq\left(1-\rho^{2}\right) \max \left\{\int_{\sigma^{\epsilon}}^{t} \frac{\gamma^{2}}{X_{s}^{2}} d s, \int_{\sigma^{\epsilon}}^{t} \frac{\beta^{2}}{Y_{s}^{2}} d s\right\}
$$

If $\beta>0$, then $X \geq U$ where $U$ is the Bessel process of dimension two in (2) and

$$
\limsup _{t \rightarrow \infty}\left(\ln X_{t}\right) \geq \limsup _{t \rightarrow \infty}\left(\ln U_{t}\right)=\infty
$$

If $\beta<0$, then $X \leq U$ and

$$
\liminf _{t \rightarrow \infty}\left(\ln X_{t}\right) \leq \liminf _{t \rightarrow \infty}\left(\ln U_{t}\right)=-\infty
$$

But

$$
\ln X_{t}=\ln X_{\sigma^{\epsilon}}+\int_{\sigma^{\epsilon}}^{t} \frac{d B_{s}}{X_{s}}+\beta \int_{\sigma^{\epsilon}}^{t} \frac{d s}{X_{s} Y_{s}}
$$

and

$$
\int_{\sigma^{\epsilon}}^{t} \frac{d s}{X_{s} Y_{s}} \leq \frac{1}{2} \int_{\sigma^{\epsilon}}^{t}\left(\frac{1}{X_{s}^{2}}+\frac{1}{Y_{s}^{2}}\right) d s
$$

Thus $\ln X_{t}$ would converge a.s. as $t \rightarrow \infty$ if

$$
\int_{\sigma^{\epsilon}}^{\infty} \frac{d s}{X_{s}^{2}}<\infty \quad \text { and } \quad \int_{\sigma^{\epsilon}}^{\infty} \frac{\mathrm{ds}}{\mathrm{Y}_{\mathrm{s}}^{2}}<\infty
$$

Therefore $\langle M\rangle_{\infty}^{\epsilon}=\infty$ on $\left\{\sigma^{\epsilon}<\infty\right\}$ and the assertion is proved.

## 8. Absolute continuity properties

In this section we suppose $\rho=0$. In some cases we easily obtain an absolute continuity property between the laws of O2BPs with various parameters. When there exists a unique solution to (1), we denote by $\mathbb{P}_{x, y}^{\alpha, \beta, \gamma, \delta}$ the law on $C\left(\mathbb{R}_{+}, S\right)=C\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right) \times C\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)$of the solution starting at $(x, y) \in S$. We denote by $(U, V)$ the canonical map $\left(U_{t}(u, v), V_{t}(u, v)\right)=$ $(u(t), v(t))$ and by $\mathcal{U}_{t}=\sigma\left(\left(U_{s}, V_{s}\right), s \leq t\right)$ the canonical filtration. Recall that we denote by $\mathbb{P}_{r}^{d}$ the law of one-dimensional Bessel process of dimension $d$ starting at $r \geq 0$.

Proposition 20. Assume $\rho=0, \delta \geq \frac{1}{2}$ and $y>0$. Then

$$
\begin{equation*}
\left.\mathbb{P}_{x, y}^{\alpha, \beta, 0, \delta}\right|_{\mathcal{U}_{t}}=\exp \left\{\beta \int_{0}^{t} \frac{d U_{s}}{V_{s}}-\alpha \beta \int_{0}^{t} \frac{d s}{U_{s} V_{s}}-\frac{\beta^{2}}{2} \int_{0}^{t} \frac{d s}{V_{s}^{2}}\right\} \cdot \mathbb{P}_{x}^{2 \alpha+1} \otimes \mathbb{P}_{y}^{2 \delta+1}{\mid \mathcal{U}_{t}} \tag{24}
\end{equation*}
$$

Proof. Under $\mathbb{P}_{x}^{2 \alpha+1} \otimes \mathbb{P}_{y}^{2 \delta+1}$ the process

$$
B_{t}:=U_{t}-x-\alpha \int_{0}^{t} \frac{d s}{U_{s}}
$$

is a one-dimensional Brownian motion. The assumptions on $\delta$ and on $y$ imply that $V_{t}>0$ for any $t \geq 0$ and thus

$$
\int_{0}^{t} \frac{d s}{v^{2}(s)}<\infty \quad \mathbb{P}_{y}^{2 \delta+1}-a . s .
$$

Therefore, for $\mathbb{P}_{y}^{2 \delta+1}$-almost every $v, \int_{0}^{t} \frac{d B_{s}}{v(s)}$ is a $\mathbb{P}_{x}^{2 \alpha+1}$-centered Gaussian variable with variance $\int_{0}^{t} \frac{d s}{v^{2}(s)}$ and

$$
\int\left(\int \exp \left\{\beta \int_{0}^{t} \frac{d B_{s}}{v(s)}-\frac{\beta^{2}}{2} \int_{0}^{t} \frac{d s}{v(s)^{2}}\right\} d \mathbb{P}_{x}^{2 \alpha+1}(u)\right) d \mathbb{P}_{y}^{2 \delta+1}(v)=1
$$

We see that

$$
Z_{t}:=\exp \left\{\beta \int_{0}^{t} \frac{d B_{s}}{V_{s}}-\frac{\beta^{2}}{2} \int_{0}^{t} \frac{d s}{V_{s}^{2}}\right\}
$$

is a $\mathbb{P}_{x}^{2 \alpha+1} \otimes \mathbb{P}_{y}^{2 \delta+1}$-positive martingale with expectation 1 . Setting for any $T>0$

$$
\mathbb{Q}_{T}:=Z_{T} \cdot \mathbb{P}_{x}^{2 \alpha+1} \otimes \mathbb{P}_{y}^{2 \delta+1} \mid \mathcal{u}_{T}
$$

we check that under $\mathbb{Q}_{T}$

$$
U_{t}-x-\alpha \int_{0}^{t} \frac{d s}{U_{s}}-\beta \int_{0}^{t} \frac{d s}{V_{s}}, \quad 0 \leq t \leq T
$$

is a real Brownian motion independent of the Brownian motion $\left\{V_{t}-y-\delta \int_{0}^{t} \frac{d s}{V_{s}}, 0 \leq t \leq T\right\}$. Thus,

$$
\mathbb{Q}_{T}=\left.\mathbb{P}_{x, y}^{\alpha, \beta, 0, \delta}\right|_{\mathcal{U}_{T}}
$$

A second set of conditions is obtained using Novikov's criterion.
Proposition 21. Assume $\rho=0$ and the following set of conditions is satisfied:

$$
\begin{array}{cc}
|\beta| \leq \delta-1 / 2 & |\gamma| \leq \alpha-1 / 2  \tag{25}\\
x>0 & y>0 .
\end{array}
$$

The process

$$
Z_{t}:=\exp \left\{\beta \int_{0}^{t} \frac{d U_{s}}{V_{s}}-\gamma \int_{0}^{t} \frac{d V_{s}}{U_{s}}-(\alpha \beta+\gamma \delta) \int_{0}^{t} \frac{d s}{U_{s} V_{s}}-\frac{\beta^{2}}{2} \int_{0}^{t} \frac{d s}{V_{s}^{2}}-\frac{\gamma^{2}}{2} \int_{0}^{t} \frac{d s}{U_{s}^{2}}\right\}
$$

is a $\mathbb{P}_{x}^{2 \alpha+1} \otimes \mathbb{P}_{y}^{2 \delta+1}$-positive martingale with expectation 1 and

$$
\begin{equation*}
\left.\mathbb{P}_{x, y}^{\alpha, \beta, \gamma, \delta}\right|_{\mathcal{U}_{t}}=\left.Z_{t} \cdot \mathbb{P}_{x}^{2 \alpha+1} \otimes \mathbb{P}_{y}^{2 \delta+1}\right|_{\mathcal{U}_{t}} \tag{26}
\end{equation*}
$$

Proof. Under $\mathbb{P}_{x}^{2 \alpha+1} \otimes \mathbb{P}_{y}^{2 \delta+1}$ the processes

$$
B_{t}:=U_{t}-x-\alpha \int_{0}^{t} \frac{d s}{U_{s}}
$$

and

$$
C_{t}:=V_{t}-y-\delta \int_{0}^{t} \frac{d s}{V_{s}}
$$

are independent Brownian motions. Then $Z_{t}$ may be written

$$
Z_{t}=\exp \left\{\beta \int_{0}^{t} \frac{d B_{s}}{V_{s}}+\gamma \int_{0}^{t} \frac{d C_{s}}{U_{s}}-\frac{\beta^{2}}{2} \int_{0}^{t} \frac{d s}{V_{s}^{2}}-\frac{\gamma^{2}}{2} \int_{0}^{t} \frac{d s}{U_{s}^{2}}\right\}
$$

which shows that $Z_{t}$ is a positive local martingale. Using (8) we compute

$$
\begin{aligned}
\iint \exp \left\{\frac{\gamma^{2}}{2} \int_{0}^{t} \frac{d s}{U_{s}^{2}}\right\} d \mathbb{P}_{x}^{2 \alpha+1} \otimes d \mathbb{P}_{y}^{2 \delta+1} & =\int \exp \left\{\frac{\gamma^{2}}{2} \int_{0}^{t} \frac{d s}{R_{s}^{2}}\right\} d \mathbb{P}_{x}^{2 \alpha+1} \\
& =\int\left(\frac{R_{t}}{x}\right)^{\alpha-1 / 2} \exp \left\{-\frac{(\alpha-1 / 2)^{2}-\gamma^{2}}{2} \int_{0}^{t} \frac{d s}{R_{s}^{2}}\right\} d \mathbb{P}_{x}^{2} \\
& <\infty
\end{aligned}
$$

since a Bessel process of dimension two has finite moments of any order. Finally,

$$
\begin{aligned}
& \iint \exp \left\{\frac{\beta^{2}}{2} \int_{0}^{t} \frac{d s}{V_{s}^{2}}\right\} \exp \left\{\frac{\gamma^{2}}{2} \int_{0}^{t} \frac{d s}{U_{s}^{2}}\right\} d \mathbb{P}_{x}^{2 \alpha+1} \otimes d \mathbb{P}_{y}^{2 \delta+1} \\
= & \left(\int \exp \left\{\frac{\beta^{2}}{2} \int_{0}^{t} \frac{d s}{R_{s}^{2}}\right\} d \mathbb{P}_{x}^{2 \delta+1}\right)\left(\int \exp \left\{\frac{\gamma^{2}}{2} \int_{0}^{t} \frac{d s}{R_{s}^{2}}\right\} d \mathbb{P}_{x}^{2 \alpha+1}\right) \\
< & \infty
\end{aligned}
$$

and Novikov's criterion ([28], Proposition VIII.1.15) proves that $Z_{t}$ has expectation 1 with respect to $\mathbb{P}_{x}^{2 \alpha+1} \otimes \mathbb{P}_{y}^{2 \delta+1}$. We easily see that $\left\{B_{t}-\beta \int_{0}^{t} \frac{d s}{V_{s}}, 0 \leq t \leq T\right\}$ and $\left\{C_{t}-\gamma \int_{0}^{t} \frac{d s}{U_{s}}, 0 \leq\right.$ $t \leq T\}$ are independent Brownian motions under the probability with density $Z_{T}$ and this proves that the new probability is $\left.\mathbb{P}_{x, y}^{\alpha, \beta, \gamma, \delta}\right|_{\mathcal{U}_{T}}$.

## 9. Product form stationary distribution

For $a>0$ and $c>0$, let $\Gamma(a, c)$ be the probability measure on $[0, \infty)$ with density

$$
\gamma(x ; a, c):=\frac{c^{a}}{\Gamma(a)} x^{a-1} e^{-c x}
$$

and characteristic function

$$
\begin{equation*}
\phi(\lambda ; a, c):=\int_{0}^{\infty} e^{i \lambda x} \gamma(x ; a, c) d x=\left(1-\frac{i \lambda}{c}\right)^{-a} . \tag{27}
\end{equation*}
$$

For $a>0, b>0$, let $B(a, b)$ be the probability measure on $[0,1]$ with density

$$
\beta(x ; a, b):=\frac{\Gamma(a+b)}{\Gamma(a) \Gamma(b)} x^{a-1}(1-x)^{b-1} .
$$

It is well-known that the function $g(x)=x^{d-1}$ is an invariant measure density on $[0, \infty)$ for the Bessel process of dimension $d$. To get a stationary probability we have to introduce a negative drift that entails positive recurrence. Then, the process

$$
X_{t}=X_{0}+B_{t}+\frac{d-1}{2} \int_{0}^{t} \frac{d s}{X_{s}}-\theta t
$$

with $\theta>0$, has $\gamma(x ; d, 2 \theta)$ as stationary density. As for the Brownian motion reflected at 0 with constant drift $-\theta$, it has the stationary exponential density $2 \theta \exp \{-2 \theta x\}$. In the bidimensional case, the drifted obliquely reflected Brownian motion has a stationary density in the form of product of two exponential densities if and only if it satisfies at once [10, 40]:

- an invertibility condition on the reflection matrix;
- a positivity condition on the exponential coefficients;
- a skew-symmetry condition.

Therefore it is natural to ask wether one can find similar conditions for drifted O2BPs ensuring existence of a stationary distribution in the form of product of two gamma distributions. The answer is positive and given in the next theorem. In fact, this may be considered as a consequence of the study in [23] which introduced a generalised reflected Brownian motion associated with a potential $U$ regular on the whole real line. The difference is that here the logarithmic potential is defined only on the positive axis. The proof below is an adaptation of the proof in [23] to this special case.

We introduce an additional constant drift $(-\theta,-\eta)$ in order to make the solution a recurrent process in the nonnegative quadrant. We consider the system

$$
\begin{align*}
& X_{t}=X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}}+\beta \int_{0}^{t} \frac{d s}{Y_{s}}-\theta t \\
& Y_{t}=Y_{0}+C_{t}+\gamma \int_{0}^{t} \frac{d s}{X_{s}}+\delta \int_{0}^{t} \frac{d s}{Y_{s}}-\eta t \tag{28}
\end{align*}
$$

with the conditions $X_{t} \geq 0, Y_{t} \geq 0$. Changing probability through a Girsanov transformation, we easily check that Theorem 8 in Section 4 is still valid for this drifted system. Moreover, the proofs in Section 5 do not bother whether the Brownian motions $B$ and $C$ are drifted or not drifted. Therefore existence and uniqueness results in Section 5 hold true for the solution to (28).

Theorem 22. Assume there exists a unique solution to (28) in $S^{\mathbf{0}}$ or $S$. This process has an invariant distribution in the form $\Gamma(a, c) \otimes \Gamma(b, d)$ with $a>1$ and $b>1$ if and only if at once

- $\alpha \delta-\beta \gamma \neq 0 \quad$ (invertibility of the interaction matrix)
- $\alpha \eta-\gamma \theta>0$ and $\quad \delta \theta-\beta \eta>0 \quad$ (positivity of exponents)
- $2 \rho=\frac{\beta}{\delta}+\frac{\gamma}{\alpha} \quad$ (skew-symmetry)

Under these conditions, the unique solution is given by

- $a=1+2 \alpha$
- $b=1+2 \delta$
- $c=2 \alpha \frac{\delta \theta-\beta \eta}{\alpha \delta-\beta \gamma}$
- $d=2 \delta \frac{\alpha \eta-\gamma \theta}{\alpha \delta-\beta \gamma}$.

Proof. We first remark that under the above conditions $\alpha \delta-\beta \gamma$ cannot be $<0$, because in that case there is no solution if $\beta$ and $\gamma$ are $<0$ and if $\beta$ and $\gamma$ are $>0$ the skew-symmetry condition is not satisfied since $2 \rho \leq 2<\beta / \delta+\gamma / \alpha$. So we could have written $\alpha \delta-\beta \gamma>0$ as first condition. Let now

$$
q(x, y)=\gamma(x ; a, c) \gamma(y ; b, d) \quad \text { for } x \geq 0, y \geq 0
$$

The infinitesimal generator of the diffusion (28) is given by

$$
L=\frac{1}{2}\left(\frac{\partial^{2}}{\partial x^{2}}+\frac{\partial^{2}}{\partial y^{2}}+2 \rho \frac{\partial^{2}}{\partial x \partial y}\right)+\left(\frac{\alpha}{x}+\frac{\beta}{y}-\theta\right) \frac{\partial}{\partial x}+\left(\frac{\gamma}{x}+\frac{\delta}{y}-\eta\right) \frac{\partial}{\partial y} .
$$

Assume $\left(X_{0}, Y_{0}\right)$ has density $q$ and characteristic function

$$
\mathbb{E}\left[e^{i\left(\lambda X_{0}+\mu Y_{0}\right)}\right]=\phi(\lambda ; a, c) \phi(\mu ; b, d)
$$

We set $f(x, y)=e^{i(\lambda x+\mu y)}$. We want to prove that for any $t \geq 0, \lambda$, and $\mu$,

$$
\mathbb{E}\left[f\left(X_{t}, Y_{t}\right)\right]=\mathbb{E}\left[f\left(X_{0}, Y_{0}\right)\right](=\phi(\lambda ; a, c) \phi(\mu ; b, d))
$$

It is enough to prove that

$$
\int_{0}^{\infty} \int_{0}^{\infty} L f(x, y) q(x, y) d x d y=0
$$

for any $\lambda$ and $\mu$. Let

$$
R(\lambda, \mu):=\left[-\frac{1}{2}\left(\lambda^{2}+\mu^{2}+2 \rho \lambda \mu\right)+i \lambda\left(\frac{\alpha}{x}+\frac{\beta}{y}-\theta\right)+i \mu\left(\frac{\gamma}{x}+\frac{\delta}{y}-\eta\right)\right]
$$

and compute

$$
\begin{aligned}
\int_{0}^{\infty} \int_{0}^{\infty} L f(x, y) q(x, y) d x d y & =\int_{0}^{\infty} \int_{0}^{\infty} R(\lambda, \mu) e^{i(\lambda x+\mu y)} \gamma(x ; a, c) \gamma(y ; b, d) d x d y \\
& =S(\lambda, \mu) \phi(\lambda ; a, c) \phi(\mu ; b, d)
\end{aligned}
$$

where $S(\lambda, \mu)$ is a second degree polynomial. In the last computation we used formula (27) several times. Setting to zero the coefficients of $\lambda, \mu, \lambda^{2}, \mu^{2}$ and $\lambda \mu$ in polynomial $S$ we obtain the set of conditions:

$$
\begin{aligned}
& 0=-\theta+\frac{\alpha c}{a-1}+\frac{\beta d}{b-1} \\
& 0=-\eta+\frac{\gamma c}{a-1}+\frac{\delta d}{b-1} \\
& 0=-\frac{1}{2}+\frac{\alpha}{a-1} \\
& 0=-\frac{1}{2}+\frac{\delta}{b-1} \\
& 0=-\rho+\frac{\gamma}{a-1}+\frac{\beta}{b-1} .
\end{aligned}
$$

The solution is given by the specified values for $a, b, c, d$ and the skew-symmetry condition.

It was proved in [39] (see also [29]) that under a skew-symmetry condition the obliquely reflected Brownian motion does not reach the non-smooth part of the boundary. In the same way the above skew-symmetry equality is reminiscent of the second condition in Corollary 10 , which is an inequality. We therefore get a partial but handy statement.

Corollary 23. Assume

- $\max \{\alpha, \delta\} \geq \frac{1}{2}$
- $\alpha \delta-\beta \gamma>0$
- $\alpha \eta-\gamma \theta>0 \quad$ and $\quad \delta \theta-\beta \eta>0$
- $2 \rho=\frac{\beta}{\delta}+\frac{\gamma}{\alpha}$.

Then the unique solution to (28) has an invariant distribution given by

$$
\Gamma\left(1+2 \alpha, 2 \alpha \frac{\delta \theta-\beta \eta}{\alpha \delta-\beta \gamma}\right) \otimes \Gamma\left(1+2 \delta, 2 \delta \frac{\alpha \eta-\gamma \theta}{\alpha \delta-\beta \gamma}\right)
$$

Proof. This is a straightforward consequence of Corollary 10, Theorems 11, 13, 14 and 22.

Remark 24. When $\rho=0$, the skew-symmetry condition $\alpha \beta+\gamma \delta=0$ means that the directions of interaction $\mathbf{r}_{\mathbf{x}}$ and $\mathbf{r}_{\mathbf{y}}$ are orthogonal. Then the second condition in Theorem 22 means that $(\theta, \eta)$ points into the interior of the quadrant designed by $\mathbf{r}_{\mathbf{x}}$ and $\mathbf{r}_{\mathbf{y}}$. This condition ensures recurrence of the process, while the first condition is now a consequence of the skew-symmetry condition.

Remark 25. With the same proof, we may check that under the skew-symmetry condition, when $\theta=\eta=0$, the function $q(x, y)=x^{2 \alpha} y^{2 \delta}$ is a non-integrable invariant density that does not depend on the other parameters.

A simple change of variables (beta-gamma algebra) provides the following result.
Corollary 26. With the conditions and notations of Theorem 22, the two-dimensional process

$$
\begin{align*}
& W_{t}:=\frac{c X_{t}}{c X_{t}+d Y_{t}}  \tag{29}\\
& Z_{t}:=c X_{t}+d Y_{t}
\end{align*}
$$

has $B(a, b) \otimes \Gamma(a+b, 1)$ for invariant distribution on $[0,1] \times[0, \infty)$.

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## References

[1] Bass R.F., Pardoux E. Uniqueness for diffusions with piecewise constant coefficients. Probab. Th. Rel. Fields 76, 557-572, 1987.
[2] Cépa E. Equations différentielles stochastiques multivoques. Sém. Probab. XXIX, Lecture Notes in Math. 1613, 86-107, Springer 1995.
[3] Cépa E., Lépingle D. Diffusing particles with electrostatic repulsion. Probab. Th. Rel. Fields 107, 429-449, 1997.
[4] Dai J.G., Williams R.J. Existence and uniqueness of semimartingale reflecting Brownian motions in convex polyhedra. Theory Probab. Appl. 40, 1-40, 1996.
[5] Dupuis P., Williams R.J., Lyapunov functions for semimartingale reflecting Brownian motions. Ann. Probab. 22, 680-702, 1994.
[6] Fernholz E.R. Stochastic portfolio theory. Applications of Mathematics 48, Springer Verlag, New York, 2002.
[7] Harrison J.M. The diffusion approximation for tandem queues in heavy traffic. Adv. Appl. Probab. 10, 886-905, 1978.
[8] Harrison J.M., Hasenbein J.J. Reflected Brownian motion in the quadrant: tail behavior of the stationary distribution. Queueing Syst. 61, 113-138, 2009.
[9] Harrison J.M., Reiman M.I. Reflected Brownian motion in an orthant. Ann. Probab. 9, 302-308, 1981.
[10] Harrison J.M., Williams R.J. Multidimensional reflected Brownian motion having exponential stationary distributions. Ann. Probab. 15, 115-137, 1987.
[11] Harrison J.M., Williams R.J. Brownian models of open queueing networks with homogeneous customer populations. Stochastics 22, 77-115, 1987.
[12] Hobson D.G., Rogers L.C.G. Recurrence and transience of reflecting Brownian motion in the quadrant. Math. Proc. Camb. Phil. Soc. 113, 387-399, 1993.
[13] Ichiba T., Karatzas I. On collisions of Brownian particles. Ann. Appl. Probab. 20, 951-977, 2010.
[14] Ichiba T., Karatzas J., Shkolnikov M. Strong solutions to stochastic equations with rank-biased coefficients. Probab. Th. Rel. Fields 156, 229-248, 2013.
[15] Ichiba T., Karatzas I., Prokaj V. Diffusions with rank-based characteristics and values in the nonnegative quadrant. Bernoulli 19, 2455-2493, 2013.
[16] Ichiba T., Pal S., Shkolnikov M. Convergence rates for rank-based models with applications to portfolio theory. Probab. Th. Rel. Fields 156, 415-448, 2013.
[17] Ichiba T., Papathanakos V., Banner A., Karatzas I., Fernholz R. Hybrid Atlas models. Ann. Appl. Probab. 21, 609-644, 2011.
[18] Jeanblanc M., Yor M., Chesney M. Mathematical Methods for Financial Markets. Springer Verlag, Berlin, 2009.
[19] Lépingle D. Boundary behavior of a constrained Brownian motion between reflecting-repelling walls. Probab. Math. Stat. 30, 273-287, 2010.
[20] Lions P.L., Sznitman A.S. Stochastic differential equations with reflecting boundary conditions. Comm. Pure Appl. Math. 37 (1984) 511-537, 1984.
[21] McKean H.P. Stochastic Integrals. Academic Press, New York, 1969.
[22] Menaldi J.L. Stochastic variational inequality for reflected diffusion. Indiana Univ. Math. J. 32, 733-744, 1983.
[23] O'Connell N., Ortmann J. Product-form invariant measures for Brownian motion with drift satisfying a skew-symmetry type condition. ALEA, Lat. Am. J. Probab. Math. Stat. 11, 307-329, 2014.
[24] Pal S., Pitman J. One-dimensional Brownian particle systems with rank-dependent drifts. Ann. Appl. Probab. 18, 2179-2207, 2008.
[25] Pardoux E., Williams R.J. Symmetric reflected diffusions. Ann. Inst. Henri Poincaré 30, 13-62, 1994.
[26] Pettersson R. Yosida approximations for multivalued stochastic differential equations. Stoch. Stoch. Reports 52, 107-120, 1995.
[27] Răşcanu A. Deterministic and stochastic differential equations in Hilbert spaces involving multivalued maximal monotone operator. PanAmerican Math. J. 6, 83-119, 1996.
[28] Revuz D., Yor M. Continuous Martingales and Brownian Motion. Second Edition, Springer Verlag, Berlin, 1994.
[29] Sarantsev A. Triple and simultaneous collisions of competing Brownian particles. Electr. J. Probab. 20 (29), 1-28, 2015.
[30] Sarantsev A. Penalty method for obliquely reflected diffusions. Preprint. Available at arXiv:1509.01777v5.
[31] Sarantsev A. Reflected Brownian motion in a convex polygonal cone: tail estimates for the stationary distribution. Preprint. Available at arXiv: 1509.01781v4.
[32] Shiga T., Watanabe S. Bessel diffusions as a one-parameter family of diffusion processes. Z. Wahrsch. Verw. Geb. 27, 37-46, 1973.
[33] Slomiński L. Weak and strong approximations of reflected diffusions via penalization methods. Stochastic Process. Appl. 123, 752-763, 2013.
[34] Storm A. Stochastic differential equations with convex constraint. Stoch. Stoch. Rep. 53, 241-274, 1995.
[35] Taylor L.M., Williams R.J. Existence and uniqueness of semimartingale reflecting Brownian motions in an orthant. Probab. Th. Rel. Fields 96, 283-317, 1993.
[36] Varadhan S.R.S., Williams R.J. Brownian motion in a wedge with oblique reflection. Comm. Pure Appl. 38, 405-443, 1985.
[37] Williams R.J. Reflected Brownian motion in a wedge: semimartingale property. Z. Wahr. verw. Geb. 69, 161-176, 1985.
[38] Williams R.J. Recurrence classification and invariant measure for reflected Brownian motion in a wedge. Ann. Probab. 13, 758-778, 1985.
[39] Williams R.J. Reflected Brownian motion with skew symmetric data in a polyhedral domain. Probab. Th. Rel. Fields 75, 459-485, 1987.
[40] Williams R.J. Semimartingale reflecting Brownian motions in the orthant. A survey. In Stochastic Networks. Springer Verlag, New York, 1995.
[41] Williams R.J., Zheng W.A. On reflecting Brownian motion - a weak convergence approach. Ann. Inst. H.P. 26, 461-488, 1990.


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