



HAL
open science

A KINETIC MODEL FOR THE TRANSPORT OF ELECTRONS IN A GRAPHEN LAYER

Clotilde Fermanian Kammerer, Florian Méhats

► **To cite this version:**

Clotilde Fermanian Kammerer, Florian Méhats. A KINETIC MODEL FOR THE TRANSPORT OF ELECTRONS IN A GRAPHEN LAYER. 2015. hal-01160791v1

HAL Id: hal-01160791

<https://hal.science/hal-01160791v1>

Preprint submitted on 8 Jun 2015 (v1), last revised 27 Sep 2017 (v3)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

A KINETIC MODEL FOR THE TRANSPORT OF ELECTRONS IN A GRAPHEN LAYER

CLOTILDE FERMANIAN KAMMERER AND FLORIAN MÉHATS

ABSTRACT. In this article, a kinetic model for the transport of electrons in graphene is derived with the tools of semiclassical analysis. The underlying quantum model is a massless Dirac equation, whose eigenvalues display a conical singularity responsible for non adiabatic transitions between the two modes. Our kinetic model takes the form of two Boltzmann equations coupled by a collision operator modeling these transitions. This collision term includes a Landau-Zener transfer term and a jump operator whose presence is essential in order to ensure a good energy conservation during the transitions. We propose an algorithmic realization of the semi-group solving the kinetic model, by a particle method. In the last section, a series of numerical experiments are given in order to study the influences of the various sources of errors between the quantum and the kinetic models.

1. INTRODUCTION

1.1. Graphene structures. Recently, graphene based structures have been the object of intensive research in nanoelectronics, see for instance the reviews [5, 10] and references therein. Graphene is a single 2D sheet of carbon atoms in a honeycomb lattice and, differently from conventional semiconductors, the most important aspect of graphene's energy dispersion is its linear energy-momentum relationship. Electrons behave as massless relativistic particles, the conduction and valence bands intersecting at the zero energy point, with no energy gap. These features enable to observe at low energy some physical phenomena of quantum electrodynamics, such as Klein tunneling that is, the fact that Dirac fermions can be transmitted through a classically forbidden region.

We are here interested in the description of the transport of electrons in a graphene device via a kinetic model. Kinetic models are usually easier to implement numerically and have a cheaper numerical cost, compared to out-of-equilibrium full quantum models. Moreover, the treatment of boundary conditions is simpler in this framework, which also enables to enrich the description by adding collisional effects via Boltzmann-like terms. However, due to the absence of gap between the conduction and valence bands, it is not correct to describe separately electrons and holes, which remain coupled even at the semiclassical limit. The objective of this paper is to introduce a kinetic model for ballistic transport, which treats the possible transitions between bands. This kinetic model is derived rigorously in a linear setting and leads to algorithmic realizations which is tested numerically. Similar

The authors would like to express their gratitude to Caroline Lasser for her help. This work was supported by the ANR-FWF Project Lodiquas ANR-11-IS01-0003 and by the ANR project Moonrise ANR-14-CE23-0007-01.

strategy has been developed at the same moment where we were writing this paper by A. Faraj and S. Jin in [12].

1.2. The quantum model. The kinetic model that will be introduced below consists in a system of approximate equations based on the Wigner counterpart of an underlying quantum transport model. At the quantum level, the ensemble of particles is described by its density matrix $\varrho(T)$, solving the von Neumann equation

$$i\hbar\partial_T\varrho = [\mathcal{H}, \varrho].$$

The Hamiltonian reads

$$\mathcal{H} = -i\hbar v_F \sigma \cdot \nabla_X + eU = v_F \hbar \begin{pmatrix} 0 & -i\partial_{X_1} - \partial_{X_2} \\ -i\partial_{X_1} + \partial_{X_2} & 0 \end{pmatrix} + eU,$$

where $X \in \mathbf{R}^2$, v_F is the Fermi velocity, $\sigma = (\sigma_{X_1}, \sigma_{X_2})$ denotes the Pauli matrices vector and $U = U(X)$ is a potential, see [10] for physical references.

Let us first put this equation in dimensionless form. We introduce a characteristic space length L , a characteristic energy \bar{E} and a characteristic density \bar{n} , then define the associated characteristic time by $\bar{t} = \frac{\hbar}{\bar{E}}$ and denote

$$x = \frac{X}{L}, \quad t = \frac{T}{\bar{t}}, \quad V = \frac{eU}{\bar{E}}, \quad \varrho^\varepsilon = \frac{\varrho}{\bar{n}L^2}.$$

The system in dimensionless form reads

$$(1.1) \quad i\varepsilon\partial_t\varrho^\varepsilon = [(A(\varepsilon D) + V), \varrho^\varepsilon],$$

where the semiclassical dimensionless parameter is

$$\varepsilon = \frac{\hbar v_F}{\bar{E}L} \ll 1,$$

where $D = -i\nabla_x$ and A is the matrix

$$A(\xi) = \begin{pmatrix} 0 & \xi_1 - i\xi_2 \\ \xi_1 + i\xi_2 & 0 \end{pmatrix}.$$

The matrix $A(\xi)$ has two eigenvalues $|\xi|$ and $-|\xi|$ with associated eigenprojectors Π^+ and Π^- ,

$$\Pi^\pm = \frac{1}{2}\text{Id} \pm \frac{1}{2|\xi|}A(\xi)$$

where Id is the identity matrix. The singularity of the eigenvalues at the point $\xi = 0$ is called conical singularity. The applied potential $V(x)$ is supposed to be smooth.

Finally, we assume that for any $\varepsilon > 0$, the initial data $\varrho^\varepsilon(0)$ is simultaneously a nonnegative trace-class operator and a Hilbert-Schmidt one. We shall denote by $\mathcal{L}^1(L^2(\mathbf{R}^2))$ the set of trace-class operators on $L^2(\mathbf{R}^2)$ and by $\mathcal{L}^2(L^2(\mathbf{R}^2))$ the space of the Hilbert-Schmidt operators. We shall assume that the family of operator $(\varrho^\varepsilon(0))_{\varepsilon>0}$ is a bounded family of $\mathcal{L}^1(L^2(\mathbf{R}^d))$, that is

$$(1.2) \quad \exists C > 0, \quad \|\varrho^\varepsilon(0)\|_{\mathcal{L}^1(L^2(\mathbf{R}^2))} \leq C.$$

We shall not assume any bound uniform in ε on the Hilbert-Schmidt norm of $\varrho^\varepsilon(0)$. Note that under these assumptions, we obtain

$$\forall t \in \mathbf{R}, \quad \forall \varepsilon > 0, \quad \varrho^\varepsilon(t) \in \mathcal{L}^1(L^2(\mathbf{R}^2)) \cap \mathcal{L}^2(L^2(\mathbf{R}^2)),$$

with

$$\forall t \in \mathbf{R}, \quad \forall \varepsilon > 0, \quad \|\varrho^\varepsilon(t)\|_{\mathcal{L}^1(L^2(\mathbf{R}^2))} \leq C.$$

1.3. Wigner functions. Denoting now by $\rho^\varepsilon(t, x, y)$ the integral kernel of ϱ^ε , the Wigner function is defined by

$$w^\varepsilon(t, x, \xi) = \frac{1}{(2\pi)^2} \int e^{i\xi \cdot \eta} \rho^\varepsilon \left(t, x - \varepsilon \frac{\eta}{2}, x + \varepsilon \frac{\eta}{2} \right) d\eta.$$

Since $\rho^\varepsilon(t)$ is Hilbert-Schmidt, its kernel is a function of $L^2(\mathbf{R}_x^2 \times \mathbf{R}_y^2)$ and similarly for $w^\varepsilon(t)$. Note however that this fact holds for any $\varepsilon > 0$ without any uniform bound. The fact that the family $(\rho^\varepsilon(t))_{\varepsilon > 0}$ is bounded in $\mathcal{L}^1(L^2(\mathbf{R}^2))$ implies that the family of distributions $(w^\varepsilon(t))_{\varepsilon > 0}$ is bounded in the set of distributions (see Remark 2.9)

We call diagonal part of the Wigner transform the scalar distributions

$$(1.3) \quad w_\pm^\varepsilon(t, x, \xi) = \text{tr} \left(\Pi^\pm(\xi) w^\varepsilon(t, x, \xi) \Pi^\pm(\xi) \right)$$

and we have

$$w^\varepsilon(t, x, \xi) = w_+^\varepsilon(t, x, \xi) \Pi^+(\xi) + w_-^\varepsilon(t, x, \xi) \Pi^-(\xi) + \Pi^+ w^\varepsilon(t, x, \xi) \Pi^- + \Pi^- w^\varepsilon(t, x, \xi) \Pi^+.$$

When ε is small, the off-diagonal contribution to the Wigner transform is known to be highly oscillating in time so that

$$\Pi^\pm w^\varepsilon(t, x, \xi) \Pi^\mp \xrightarrow[\varepsilon \rightarrow 0]{} 0 \quad \text{in } \mathcal{D}'(\mathbf{R}_t \times (\mathbf{R}^{2d} \setminus \{\xi = 0\})),$$

(see [22]). For this reason, we focus on the quasi-distribution functions $w_\pm^\varepsilon(t, x, \xi)$. Far from the crossing set $\{\xi = 0\}$, w_+^ε and w_-^ε satisfy approximated transport equations

$$(1.4) \quad \begin{cases} \partial_t w_+^\varepsilon + \frac{\xi}{|\xi|} \cdot \nabla_x w_+^\varepsilon - \nabla V(x) \cdot \nabla_\xi w_+^\varepsilon = \mathcal{O}(\varepsilon), \\ \partial_t w_-^\varepsilon - \frac{\xi}{|\xi|} \cdot \nabla_x w_-^\varepsilon - \nabla V(x) \cdot \nabla_\xi w_-^\varepsilon = \mathcal{O}(\varepsilon), \end{cases}$$

in $\mathcal{D}'(\mathbf{R}_t \times (\mathbf{R}^{2d} \setminus \{\xi = 0\}))$. Besides, the equations (1.4) imply that, outside $\{\xi = 0\}$, the functions w_\pm^ε are constant along the integral curves $(\Phi_\pm^t)_{t \in \mathbf{R}}$ of the vector fields

$$(1.5) \quad H_\pm(x, \xi) = \pm \frac{\xi}{|\xi|} \cdot \nabla_x - \nabla_x V \cdot \nabla_\xi.$$

Such curves – also called Hamiltonian curves of $\pm |\xi| + V(x)$ – are well-defined and smooth as long as they do not reach $\{\xi = 0\}$. They satisfy

$$\dot{\Phi}_\pm^t(x, \xi) = H_\pm(\Phi_\pm^t(x, \xi)), \quad \Phi_\pm^0 = (x, \xi).$$

It is proved in [14] (see Proposition 3 therein) that such curves may reach $\{\xi = 0\}$ in finite time, and that, if $\nabla V \neq 0$ at the impact point, the curve can be prolongedated in a unique way away from $\{\xi = 0\}$ generating a continuous trajectory (which is not \mathcal{C}^1). These facts are recalled in details in Section 2.1.1 below.

It is also important to notice that the evolution of w_+^ε and w_-^ε are decoupled at leading order outside $\{\xi = 0\}$: this regime is said to be adiabatic. The singularity of the eigenvalues of $A(\xi)$ when $\xi = 0$ is known to produce non adiabatic transitions between the modes. Our aim here is to propose a kinetic model which is also valid

close to $\xi = 0$. We are going to add a collision kernel to these equations which will couple the evolutions and generate transitions between the modes.

1.4. Conical singularities. Systems presenting conical singularities have been the subject of extensive works since the early thirties with the works of Landau and Zener [27, 33]. Such singularities arise in particular when studying molecular dynamics in the frame of Born-Oppenheimer approximation (see [31, 28] for example). Pioneer works have been performed in this context by G. Hagedorn and his collaborators, with a wave-packet approach [23, 24]. Several ideas used here are due to these contributions. Ten years ago, classification of crossings for rather general systems was performed independently by [6, 7] and [15]. In the latter reference and in [14], the analysis of the crossing is made from the point of view of Wigner transform and can be adapted to our setting. This kind of analysis has led to numerical realizations for molecular propagation ([29], [17] and [18]) and we have been inspired by these results. Of course, the Dirac equation arising in the graphen context presents major difference, when compared to the Schrödinger equation which models molecular propagation. However, the transitions due to the conical intersections can be treated similarly. The collision kernel which treats the transitions arising from the conical intersections, is derived from the analysis of conical intersections performed in [14] and from the particle description derived in [29, 17, 18, 19] for molecular dynamics. Precise statements are given below.

1.5. The approximate kinetic model. We consider the set Σ defined by

$$\Sigma = \{(x, \xi) \in \mathbf{R}^4, \quad \xi \cdot \nabla V(x) = 0\}$$

which is an hypersurface of \mathbf{R}^4 under the assumption

$$(1.6) \quad \nabla V(x) \neq 0 \quad \forall x \in \mathbf{R}^2.$$

This set is the place where the gap between the two modes (i.e. the function $2|\xi|$) is minimal along the trajectories (see Remark 1.5). We notice that, assuming (1.6), the vector fields $H_{\pm}(x, \xi)$ defined in (1.5) are transverse to Σ in a neighborhood of $\{\xi = 0\}$. This comes from the observation that

$$(1.7) \quad \nabla V(x) \cdot (H_{\pm})_{\xi} + d^2V(x)\xi \cdot (H_{\pm})_x = -|\nabla V(x)|^2 \pm d^2V(x)\xi \cdot \frac{\xi}{|\xi|} < 0$$

if ξ is small enough. As a consequence, when the trajectories reach their minimal distance to the gap, they pass through Σ , arriving from the region $\{\xi \cdot \nabla V(x) > 0\}$ and going to the region $\{\xi \cdot \nabla V(x) < 0\}$.

We define $(f_+^{\varepsilon}(t), f_-^{\varepsilon}(t))$ as a pair of solutions to the following system:

$$(1.8) \quad \begin{cases} \partial_t f_+^{\varepsilon} + \frac{\xi}{|\xi|} \cdot \nabla_x f_+^{\varepsilon} - \nabla V(x) \cdot \nabla_{\xi} f_+^{\varepsilon} = K_+(f_+^{\varepsilon}, f_-^{\varepsilon}) \\ \partial_t f_-^{\varepsilon} - \frac{\xi}{|\xi|} \cdot \nabla_x f_-^{\varepsilon} - \nabla V(x) \cdot \nabla_{\xi} f_-^{\varepsilon} = K_-(f_+^{\varepsilon}, f_-^{\varepsilon}) \end{cases}$$

with initial conditions $f_+^{\varepsilon}(0) = w_+^{\varepsilon}(0)$ and $f_-^{\varepsilon}(0) = w_-^{\varepsilon}(0)$ and where K_{\pm} are two collision kernels, defined below in (1.11) and (1.12).

The collision process is involved above Σ ; as a consequence, outside Σ , the functions f_{\pm}^{ε} are constant along the curves $(\Phi_{\pm}^t)_{t \in \mathbf{R}}$ introduced previously. Starting from a initial data localized far from Σ , the solution f_{\pm}^{ε} of system (1.8) is obtained

by propagating the data by the flow $(\Phi_{\pm}^t)_{t \in \mathbf{R}}$ so that the plus and the minus modes have decoupled evolutions. Whenever trajectories reach Σ , the transition kernel will generate transfers between the modes.

Even for smooth initial data, the result of this process will not be smooth functions and they will present discontinuities on Σ . But we have to localize on Σ functions that present discontinuities through it, and thus have different traces. In order to distinguish two sides of Σ , we will take advantage from the fact that, as noticed above, the flows H_{\pm} are transverse to Σ in suitably chosen neighborhoods Ω of points (x_0, ξ_0) such that $\nabla V(x_0) \neq 0$ and $|\xi_0|$ is small enough. For a function g which is defined in such Ω , we set for $(x, \xi) \in \Omega \cap \Sigma$

$$(1.9) \quad g_{\Sigma, in}(x, \xi) = \lim_{s \rightarrow 0^-} g(\Phi_{\pm}^s(x, \xi)) \quad \text{and} \quad g_{\Sigma, out}(x, \xi) = \lim_{s \rightarrow 0^+} g(\Phi_{\pm}^s(x, \xi)).$$

Let us now describe these collision kernels K_{\pm} . They depend on a transfer coefficient

$$(1.10) \quad T_{\varepsilon}(x, \xi) = \exp\left(-\frac{\pi}{\varepsilon} \frac{|\xi|^2}{|\nabla V(x)|}\right),$$

and on two jump operators

$$J_{\pm}(x, \xi) = \left(x \pm 2|\xi| \frac{\nabla V(x)}{|\nabla V(x)|^2}, \xi\right), \quad (x, \xi) \in \Sigma.$$

Then, for functions $f \in \mathcal{C}^1(\mathbf{R}^4 \setminus \Sigma)$, which have traces on Σ , we distinguish between *entering* and *outgoing* traces by denoting by $f_{\Sigma, in}$ the restriction to Σ of the function $f \mathbf{1}_{\{\xi \cdot \nabla V(x) \geq 0\}}$ and by $f_{\Sigma, out}$ the restriction to Σ of $f \mathbf{1}_{\{\xi \cdot \nabla V(x) \leq 0\}}$. This is adapted to functions which are discontinuous through Σ . Then, the collision kernels K_{\pm} are defined by

$$(1.11) \quad K_+(f, g) = \Lambda_+(x, \xi) \delta_{\Sigma}(x, \xi) (T_{\varepsilon} f_{\Sigma, in} - (T_{\varepsilon} g_{\Sigma, in}) \circ J_+),$$

$$(1.12) \quad K_-(f, g) = \Lambda_-(x, \xi) \delta_{\Sigma}(x, \xi) (T_{\varepsilon} g_{\Sigma, in} - (T_{\varepsilon} f_{\Sigma, in}) \circ J_-),$$

where the Jacobians Λ_{\pm} are defined by

$$(1.13) \quad \forall (x, \xi) \in \Sigma, \quad \Lambda_{\pm}(x, \xi) = \frac{-|\nabla V(x)|^2 \pm |\xi|^{-1} d^2 V(x) \xi \cdot \xi}{\sqrt{|\nabla V(x)|^2 + |d^2 V(x) \xi|^2}}.$$

Remark 1.1. Note that the transfer coefficient $T_{\varepsilon}(x, \xi)$ is exponentially small as soon as $|\xi| > R\sqrt{\varepsilon}$ for some $R > 0$. Moreover, if $|\xi| \leq R\sqrt{\varepsilon}$, we have

$$J_{\pm} = \text{Id} + \mathcal{O}(R\sqrt{\varepsilon}) \quad \text{and} \quad \Lambda_{\pm}(x, \xi) = -|\nabla V(x)| + \mathcal{O}(R\sqrt{\varepsilon}),$$

under the assumption (1.6).

Since the initial data that we want to treat are not necessarily continuous outside Σ , something has to be said in order to extend the collision kernel to locally integrable functions. For this and for stating the main result, we shall need a certain number of assumptions.

- Assumption 1.2.** (1) *The initial data $(\varrho_0^{\varepsilon})_{\varepsilon > 0}$ satisfies (1.2) and its Wigner transform $(w^{\varepsilon}(0))_{\varepsilon > 0}$ is localized away from Σ .*
(2) *The potential is non degenerated: $\nabla V(x) \neq 0$.*

- (3) We have $w_-^\varepsilon(0) = \mathcal{O}(\varepsilon^{1/8})$ in $L^1(\mathbf{R}^{2d})$ and the symbol a and the time T are such that within the time interval $[0, T]$, each of the trajectories (Φ_+^t) arriving at the support of a at time T has passed through Σ at most once.

Our main result is the following theorem, which states that the functions (f_\pm^ε) provides an approximation of the Wigner transforms (w_\pm^ε) .

Theorem 1.3. *If (1) and (2) of Assumption 1.2 is satisfied on the time interval $[0, T]$, then (1.8) admits a unique weak solution.*

Moreover if $a \in \mathcal{C}_0^\infty(\mathbf{R}^2)$ is compactly supported outside $\{\xi = 0\}$ and $\chi \in \mathcal{C}_c^\infty([0, T], \mathbf{R})$ satisfy (3) of Assumption 1.2, then there exist positive constants $C, \varepsilon_0 > 0$ such that, for all $0 < \varepsilon < \varepsilon_0$,

$$(1.14) \quad \left| \operatorname{tr} \int_{\mathbf{R}^5} \chi(t) (f_\pm^\varepsilon - w_\pm^\varepsilon)(t, x, \xi) a(x, \xi) dx d\xi dt \right| \leq C\varepsilon^{1/8}.$$

Remark 1.4. (1) The existence of solutions to these kinetic equations comes from the fact that, under (1) and (2) of Assumption 1.2, these equations have a particle description relying on a Markov semi-group that is explained in the next section and is crucial for the proof of the theorem.

- (2) Note that the role of the indexes *plus* and *minus* can be inverted in (3) of Assumptions 1.2 and result of Theorem 1.3 still holds.
- (3) The limitation induced by (3) to the range of validity of Theorem 1.3 comes from the fact that the kinetic kernels K_\pm are not adapted in some situations where the modes interfere too much. It appears nevertheless that these kernels' description encounter a larger range of situation than those satisfying (3), as it appears in the numerical realizations of Section 3. Some example of situation where (3) is not satisfied and where the description by the kernels K_\pm fails is given in [20] in the context of conical intersections for molecular dynamics.
- (4) We refer to Remark 3.1 below where additional comments are performed in view of the semi-group description of $w_\pm^\varepsilon(t)$.

1.6. A Markov semi-group. We consider the Hamiltonian flows

$$\Phi_+^t = (x_+^t(x, \xi), \xi_+^t(x, \xi)) \quad \text{and} \quad \Phi_-^t = (x_-^t(x, \xi), \xi_-^t(x, \xi))$$

with $(x_\pm^0(x, \xi), \xi_\pm^0(x, \xi)) = (x, \xi)$ and

$$(1.15) \quad \begin{cases} \frac{d}{dt} x_+^t = \frac{\xi_+^t}{|\xi_+^t|} \quad \text{and} \quad \frac{d}{dt} \xi_+^t = -\nabla V(x_+^t), \\ \frac{d}{dt} x_-^t = -\frac{\xi_-^t}{|\xi_-^t|} \quad \text{and} \quad \frac{d}{dt} \xi_-^t = -\nabla V(x_-^t). \end{cases}$$

As long as $\xi \neq 0$, the smoothness of the Hamiltonians $|\xi| \pm V(x)$ yields local existence and uniqueness of the trajectory passing through (x, ξ) for any $x \in \mathbf{R}^2$. However, it may happens that $\xi_\pm^t \xrightarrow[t \rightarrow t^*]{} 0$ for some $t^* \in \mathbf{R}$ and some index $+$ or $-$. If at the point $x_\pm^{t^*}$, the assumption (1.6) is satisfied (that is if $\nabla V(x_\pm^{t^*}) \neq 0$), then one can prove that there exists a unique continuation to the map $t \mapsto \Phi_\pm^t$ when $t > t^*$ (see Proposition 1 in [15] and Proposition 2.1 below where a precise statement and a proof are given for the convenience of the reader). As a consequence, the assumption (1.6) guarantees the existence and uniqueness of the solutions to (1.15). However, these trajectories are no longer smooth when passing through $\xi = 0$; more

precisely, the vector $\dot{\Phi}^t$ has a discontinuity at $t = t^*$ whenever $\xi_{\pm}^{t^*} = 0$. It is also interesting to notice that if the latter assumption (1.6) fails at $x_{\pm}^{t^*}$, then uniqueness is no longer guaranteed.

Remark 1.5. For any trajectory Φ_{\pm}^t , the quantity $|\xi_{\pm}^t|$ reaches its minimum on Σ . Indeed, we can check that

$$\frac{d}{dt} (|\xi_{\pm}^t|^2) = -\xi_{\pm}^t \cdot \nabla V(x_{\pm}^t),$$

which is equal to 0 whenever $\Phi_{\pm}^t \in \Sigma$. This motivates the introduction of the set Σ .

We are now going to introduce a branching process between both types of trajectories. We attach the labels $+$ and $-$ to the phase space and for points

$$(x, \xi, j) \in \mathbf{R}_{\pm}^4 := \mathbf{R}^4 \times \{+1, -1\},$$

we consider trajectories

$$\mathcal{T}_{\varepsilon}^{(x, \xi, j)} : [0, +\infty) \rightarrow \mathbf{R}_{\pm}^4,$$

which combine deterministic classical transport and random jumps between the levels at the manifold Σ . More precisely, we set

$$\mathcal{T}_{\varepsilon}^{(x, \xi, j)}(t) = (\Phi_j^t(x, \xi), j)$$

as long as $\Phi_j^t(x, \xi) \notin \Sigma$. Whenever the deterministic flow $\Phi_j^t(x, \xi)$ hits the manifold Σ at a point (x^*, ξ^*) , a random jump from

$$(x^*, \xi^*, j) \text{ to } (J_j(x^*, \xi^*), -j)$$

occurs with probability $T^{\varepsilon}(x^*, \xi^*)$.

The jump aims at preserving at order $\mathcal{O}(\varepsilon)$ the energy of the trajectory for points ξ where the transfer coefficient is relevant, that is points where $|\xi| \leq R\sqrt{\varepsilon}$ according to Remark 1.1. This is an important ingredient of the proof (see Remark 2.13). Indeed, set

$$E_{\pm}(x, \xi) := \pm|\xi| + V(x),$$

then, if $|\xi| \leq R\sqrt{\varepsilon}$, at a jump from $+$ to $-$, we have

$$E_{-}(J_{+}(x, \xi)) = E_{+}(x, \xi) + \mathcal{O}(|\xi|^2) = E_{+}(x, \xi) + \mathcal{O}(R^2\varepsilon).$$

Note that the importance of the jump has been illustrated numerically in [19] in the context of molecular propagation ; this jump was already performed in [24] for the construction of gaussian wave packets which are approximated solutions of a Schrödinger equation with matrix valued potential presenting a conical intersection.

Since $\nabla V(x) \neq 0$, the trajectories which reach the manifold Σ arrive there transversally to Σ . As a consequence, in each bounded time interval $[0, T]$, each path

$$(x, \xi, j) \rightarrow \mathcal{T}_{\varepsilon}^{(x, \xi, j)}(t)$$

has a finite number of jumps and remains in a bounded region of the phase space \mathbf{R}_{\pm}^4 . Besides, away from the jump manifold $\Sigma \times \{-1, +1\}$ each path is smooth.

Following [9, 30], we define the function $P_{\varepsilon}(x, \xi, j; t, \cdot)$ as the function which associate to a measurable set $\Gamma \subset \mathbf{R}_{\pm}^4$ the probability $P_{\varepsilon}(x, \xi, j; t, \Gamma)$ of being at

time t in Γ having started in (x, ξ, j) . And we define a time-dependent Markov process $(\mathcal{L}_\varepsilon^t)_{t \geq 0}$ acting on bounded measurable scalar functions $f : \mathbf{R}_\pm^4 \rightarrow \mathbf{C}$ by

$$(\mathcal{L}_\varepsilon^t f)(x, \xi, j) := \int_{\mathbf{R}^4 \times \{-1, +1\}} f(q, p, k) P_\varepsilon(x, \xi, j; t, d(q, p, k)).$$

An explicit expression of $\mathcal{L}_\varepsilon^t$ is written on short interval times close to jump points in Section 2.2.2 (see equations (2.7) and (2.8) below).

In order to define its action on Wigner functions, we need to identify pairs of functions (a_+, a_-) with some function a on \mathbf{R}_\pm^4 , which is done by the identification

$$(1.16) \quad a(x, \xi, \pm 1) = a_\pm(x, \xi), \quad \forall (x, \xi) \in \mathbf{R}^4.$$

Through this identification, the action of $(\mathcal{L}_\varepsilon^t)_{t \geq 0}$ on $\mathcal{C}_0^\infty(\mathbf{R}^4 \setminus \{\xi = 0\})$ is given by

$$(\mathcal{L}_\varepsilon^t a)(x, \xi) := ((\mathcal{L}_\varepsilon^t a)(x, \xi, +1), (\mathcal{L}_\varepsilon^t a)(x, \xi, -1)),$$

We extend this action to $\mathcal{D}'(\mathbf{R}^4 \setminus \{\xi = 0\})$ by duality by setting

$$\forall f \in \mathcal{D}'(\mathbf{R}^4), \quad \forall a \in \mathcal{C}_0^\infty(\mathbf{R}^4 \setminus \{\xi = 0\}), \quad \langle (\mathcal{L}_\varepsilon^t f)(\pm 1), a \rangle = \langle f, (\mathcal{L}_\varepsilon^t a)(\pm 1) \rangle.$$

Proposition 1.6 (Resolution of the kinetic model). *Set $w^\varepsilon(0) = (w_+^\varepsilon(0), w_-^\varepsilon(0))$ and assume that (1) and (2) of Assumptions 1.2 are satisfied. Then the function $f^\varepsilon(t, x, \xi) = (\mathcal{L}_\varepsilon^t w^\varepsilon(0))(x, \xi)$ is the unique solution to system (1.8) in $\mathcal{D}'(\mathbf{R}^4 \setminus \{\xi = 0\})$.*

Then, Theorem 1.3 is a corollary of the following proposition.

Proposition 1.7 (Approximation by the semi-group). *Set $w^\varepsilon(0) = (w_+^\varepsilon(0), w_-^\varepsilon(0))$ and assume for $\chi \in \mathcal{C}_c^\infty([0, T], \mathbf{R})$ and $a \in \mathcal{C}_0^\infty(\mathbf{R}^4 \setminus \{\xi = 0\})$, Assumptions 1.2 are satisfied, then, there exist positive constants $C, \varepsilon_0 > 0$ such that for all $0 < \varepsilon < \varepsilon_0$,*

$$\left| \operatorname{tr} \int_{\mathbf{R}^{2d+1}} \chi(t) (w_\pm^\varepsilon(t, x, \xi) - (\mathcal{L}_\varepsilon^t w^\varepsilon(0))(x, \xi, \pm 1)) a(x, \xi) dx d\xi dt \right| \leq C \varepsilon^{1/8}.$$

Remark 1.8. Note that the hypothesis (3) of Assumptions 1.2 imply that on the interval $[0, T]$, the trajectories which reach the support of a has performed at most one jump.

This Markov semi-group suits to an algorithmic resolution which is developed in Section 3.1, where we provide some numerical experiments which testify to the performance of the algorithm.

2. PROOFS

In this section, we prove the main result, Theorem 1.3. After a preliminary section devoted to the study of the classical flow, we proceed in two steps. We first prove Proposition 1.6 which relates the approximated system (1.8) and the Markov semi-group described in Section 1.6. Then we give a sketch of the proof of Proposition 1.7, which explains why this semi-group describes the evolution of the Wigner transform of the solutions of (1.1); for this, we follow the proof of Theorem 2.2 in [17] which has been performed for a system of Schrödinger equations coupled by a matrix-valued potential.

2.1. Preliminaries. In this section, we analyze precisely the geometry close to a point of Σ in order to precise the setting in which the proofs will be performed.

2.1.1. *The generalized flow.* In this section, we gather some properties of the flows Φ_{\pm}^t that will be useful in the next sections. We first focus on the existence and uniqueness of the generalized trajectories of the Hamiltonian vector fields H_{\pm} and recall the arguments of the proof given in [14].

Proposition 2.1. *For any $x_0 \in \mathbf{R}^2$ such that $\nabla V(x_0) \neq 0$, there exists $\tau_0 > 0$ and a unique Lipschitz continuous map*

$$t \mapsto (x_{\pm}^t(x_0, 0), \xi_{\pm}^t(x_0, 0)), \quad t \in [-\tau_0, \tau_0]$$

satisfying (1.15) for $t \neq 0$ and such that $x_{\pm}^0(x_0, 0) = x_0$, $\xi_{\pm}^0(x_0, 0) = 0$ and

$$\dot{x}_{\pm}^t(x_0, 0) \xrightarrow[t \rightarrow 0]{} \mp \frac{\nabla V(x_0)}{|\nabla V(x_0)|}, \quad \dot{\xi}_{\pm}^t(x_0, \xi_0) \xrightarrow[t \rightarrow 0]{} -\nabla V_0(x_0).$$

Corollary 2.2. *With the notations of Proposition 2.1, we have*

$$\begin{aligned} \lim_{t \rightarrow 0^-} H_+(x_{\pm}^t, \xi_{\pm}^t) &= \lim_{t \rightarrow 0^+} H_-(x_{\pm}^t, \xi_{\pm}^t) = -\nabla V(x_0) \cdot \nabla_{\xi} + \frac{\nabla V(x_0)}{|\nabla V(x_0)|} \cdot \nabla_x, \\ \lim_{t \rightarrow 0^+} H_-(x_{\pm}^t, \xi_{\pm}^t) &= \lim_{t \rightarrow 0^-} H_-(x_{\pm}^t, \xi_{\pm}^t) = -\nabla V(x_0) \cdot \nabla_{\xi} - \frac{\nabla V(x_0)}{|\nabla V(x_0)|} \cdot \nabla_x. \end{aligned}$$

In the following and with the notations of Proposition 2.1, we shall set

$$(2.1) \quad H(x_0) := \lim_{t \rightarrow 0^-} H_+(x_{\pm}^t, \xi_{\pm}^t) = \lim_{t \rightarrow 0^+} H_-(x_{\pm}^t, \xi_{\pm}^t),$$

$$(2.2) \quad H'(x_0) := \lim_{t \rightarrow 0^+} H_-(x_{\pm}^t, \xi_{\pm}^t) = \lim_{t \rightarrow 0^-} H_-(x_{\pm}^t, \xi_{\pm}^t).$$

We observe that if $\omega = d\xi \wedge dx$ is the canonical skew-symmetric 2-form of the cotangent space of \mathbf{R}^2 , we have

$$(2.3) \quad \omega(H, H') = 2|\nabla V(x_0)| > 0.$$

Proof of Proposition 2.1. Following Proposition 3 in [14], we introduce two flows $\Psi_1^t = (x_1^t, \xi_1^t)$ and $\Psi_2^t = (x_2^t, \xi_2^t)$ which are defined by

$$\begin{cases} \dot{x}_j^t = (-1)^{j+1} \operatorname{sgn}(t) \frac{\xi_j^t}{|\xi_j^t|}, & x_j^0 = x_0, \\ \dot{\xi}_j^t = -\nabla V(x_j^t), \end{cases}$$

or equivalently

$$\xi_j^t = -t \int_0^1 \nabla V(x_j^{ts}) ds, \quad x_j^t = x_0 + (-1)^{j+1} \int_0^t \frac{\int_0^1 \nabla V(x_j^{s\sigma}) d\sigma}{\left| \int_0^1 \nabla V(x_j^{s\sigma}) d\sigma \right|} ds.$$

The last system can be solved on short time by a fixed point argument in an open subset of $\{\nabla V(x) \neq 0\}$ and the resulting map $x \mapsto \Psi_j^t(x, 0)$ is smooth. As a consequence, there exists a neighborhood Ω of $(x_0, 0)$ such that $\Omega \subset \{\nabla V(x) \neq 0\}$ and $\tau_0 > 0$ such that for all $(x, 0) \in \Omega$, the maps

$$\begin{aligned} (\Phi_+^t(x, 0))_{t \in [0, \tau_0]} &= (\Psi_1^t(x, 0))_{t \in [0, \tau_0]} \quad \text{and} \quad (\Phi_+^t(x, 0))_{t \in [-\tau_0, 0]} = (\Psi_2^t(x, 0))_{t \in [-\tau_0, 0]}, \\ (\Phi_-^t(x, 0))_{t \in [0, \tau_0]} &= (\Psi_2^t(x, 0))_{t \in [0, \tau_0]} \quad \text{and} \quad (\Phi_-^t(x, 0))_{t \in [-\tau_0, 0]} = (\Psi_1^t(x, 0))_{t \in [-\tau_0, 0]}, \end{aligned}$$

solve our problem. The flows Φ_{\pm}^t (which are well defined for $\xi \neq 0$) extend to Lipschitz continuous maps

$$t \mapsto \Phi_{\pm}^t(x, \xi), \quad t \in [-\tau_0, \tau_0], \quad (x, \xi) \in \Omega.$$

□

Remark 2.3. Following the arguments of Section 6.2 in [16], one can prove that for $|t| < \tau_0$ and $\alpha \in \mathbf{N}^d$, the maps $(x, \xi) \mapsto \partial_x^\alpha \Phi_\pm^t(x, \xi)$ are continuous maps on Ω with bounded locally integrable time derivatives $\partial_t \partial_x^\alpha \Phi^t(x, \xi)$.

2.1.2. *Local analysis.* In what follows, we shall associate with points (x, ξ) , which are close enough to the set Σ , a number $\tau_\pm(x, \xi)$ which is the time that separates (x, ξ) from the point of the trajectory $\Phi_\pm^t(x, \xi)$ which belongs to Σ .

Proposition 2.4. *Let $x_0 \in \mathbf{R}^2$ such that $\nabla V(x_0) \neq 0$, there exists an open set $\Omega \subset \{\nabla V(x) \neq 0\} \subset \mathbf{R}_{(x, \xi)}^4$ containing $(x_0, 0)$ and such that the relations $\Phi^{\tau_\pm(x, \xi)}(x, \xi) \in \Sigma$ define two continuous functions on Ω*

$$(x, \xi) \mapsto \tau_\pm(x, \xi).$$

Remark 2.5. By definition of $\tau^\pm(x, \xi)$, we have $\nabla V(x_{\tau_\pm(x, \xi)}^\pm) \cdot \xi_{\tau_\pm(x, \xi)}^\pm = 0$ for any $(x, \xi) \in \Omega$.

Proof. Let us study the plus mode. We observe that

$$(2.4) \quad \frac{d}{dt} |\xi_+^t|^2 = -\nabla V(x_+^t) \cdot \xi_+^t \quad \text{and} \quad \frac{d^2}{dt^2} |\xi_+^t|^2 = |\nabla V(x_+^t)|^2 - d^2 V(x_+^t) \xi_+^t \cdot \frac{\xi_+^t}{|\xi_+^t|}$$

Since $\nabla V(x_0) \neq 0$, we can find a neighborhood U of $(x_0, 0)$, $c_0 > 0$ and $\tau_1 > 0$ such that

$$(2.5) \quad \forall (x, \xi) \in \Omega_1, \quad \forall |t| \leq \tau_1, \quad \frac{d^2}{dt^2} |\xi_+^t(x, \xi)|^2 \geq c_0 > 0.$$

Because of (2.5), the map $t \mapsto |\xi_+^t|$ reaches its minimum at most once in U . With any $(x, \xi) \in U$, we associate an interval $[t_i(x, \xi), t_f(x, \xi)]$ of maximal size such that $\Phi^t(x, \xi) \in U$ for all $t \in]t_i(x, \xi), t_f(x, \xi)[$. Because of the first relation of (2.4), we are interested to the times where the curves $\Phi^t(x, \xi)$ crosses the hypersurface Σ , which happens at most once in U . We set

$$\Omega^+ = \{(x, \xi) \in U, \exists (y, \eta) \in \Sigma \cap U, \exists s \in]t_i(y, \eta), t_f(y, \eta)[, (x, \xi) = \Phi_+^s(y, \eta)\}.$$

Then Ω^+ is a neighborhood of $(x_0, 0)$ included in U and such that the relation $\Phi_+^{\tau_+(x, \xi)}(x, \xi) \in \Sigma$ defines a map τ^+ from Ω^+ into \mathbf{R} . We define similarly Ω^- and the map $\tau^-(x, \xi)$ and we choose $\Omega = \Omega^+ \cap \Omega^-$ (see Fig. 1).

It is classical to prove that τ^+ is a continuous map on Ω . We set

$$z = (x, \xi) \quad \text{and} \quad \phi(t, z) = \frac{d}{dt} |\xi_+^t|^2 = -\nabla V(x_+^t(z)) \cdot \xi_+^t(z)$$

and we argue by contradiction. We assume that there exist α_0, z_0 and a sequence $(z_n)_{n \in \mathbf{N}^*}$ going to z_0 as n goes to $+\infty$ and such that

$$|\tau_+(z_n) - \tau_+(z_0)| > \alpha_0.$$

Then, by the continuity of ϕ , we get $\phi(\tau_+(z_0), z_n) \xrightarrow{n \rightarrow +\infty} 0$ and since $\partial_t \phi \geq c_0$ (by (2.5)), we have

$$c_0 |\tau_+(z_0) - \tau_+(z_n)| \leq |\phi(\tau_+(z_0), z_n) - \phi(\tau_+(z_n), z_n)| = |\phi(\tau_+(z_0), z_n)|.$$

Taking n large enough, we get $|\tau_+(z_n) - \tau_+(z_0)| < \alpha_0/2$, whence a contradiction and the continuity of the map $z \mapsto \tau_+(z)$ is proved.

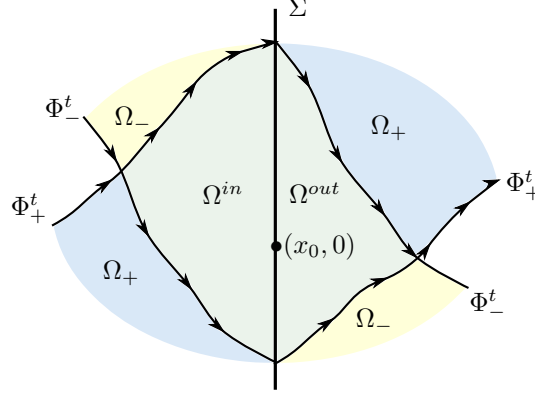


Figure 1: Domains Ω_+ (blue+green), Ω_- (yellow+green) and Ω (green) = $\Omega_+ \cap \Omega_- = \Omega^{in} \cup \Omega^{out}$

One argues similarly for the mode minus. \square

Remark 2.6. The hypersurface Σ parts Ω into two distinct connected regions (see Fig. 1):

$$(2.6) \quad \begin{aligned} \Omega^{out} &:= \{\tau_+ < 0, \tau_- < 0\} \cap \Omega = \{\nabla V(x) \cdot \xi < 0\} \cap \Omega \\ \Omega^{in} &:= \{\tau_+ > 0, \tau_- > 0\} \cap \Omega = \{\nabla V(x) \cdot \xi > 0\} \cap \Omega. \end{aligned}$$

Indeed, if $\nabla V(x) \cdot \xi > 0$, we have simultaneously, $\frac{d}{dt} |\xi_+^t|^2 < 0$ and $\frac{d}{dt} |\xi_-^t|^2 < 0$, thus we have $\tau_+(x, \xi) > 0$ and $\tau_-(x, \xi) > 0$. Besides, for $(x, \xi) \in \Omega^{in}$, we have $\Phi^t(x, \xi) \in \Omega^{in}$ for $t \in [0, \tau_+(x, \xi)]$.

This remark is fundamental to prove that the kinetic system (1.8) has a meaning in Lebesgue space $L^1(\mathbf{R}^4)$ and then to solve it. Indeed, the collision kernel requires to take the trace of the functions on Σ , which is not necessarily doable for L^1 -functions.

We consider the open set Ω of Proposition 2.4 where H^\pm are transverse to Σ . We introduce the sets Ω^{in} and Ω^{out} defined by (2.6). We can find an open set \mathcal{V} , $\mathcal{V} \subset \Omega$ such that there exists $t_i < t_f$ (i for initial and f for final) with

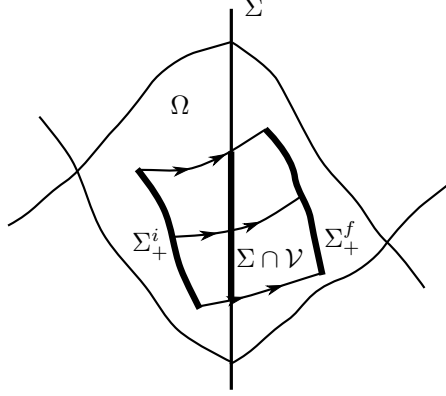
$$\begin{aligned} \forall (y, \eta) \in \Sigma \cap \mathcal{V}, \quad \Phi_\pm^{t_i}(y, \eta) &\in \Omega^{in}, \\ \forall (y, \eta) \in \Sigma \cap \mathcal{V}, \quad \Phi_\pm^{t_f}(y, \eta) &\in \Omega^{out}, \end{aligned}$$

and we define

$$\begin{aligned} \Sigma_\pm^i &= \{\Phi_\pm^{t_i}(y, \eta), (y, \eta) \in \Sigma \cap \mathcal{V}\} \subset \Omega^{in}, \\ \Sigma_\pm^f &= \{\Phi_\pm^{t_f}(y, \eta), (y, \eta) \in \Sigma \cap \mathcal{V}\} \subset \Omega^{out}, \end{aligned}$$

(see Fig. 2).

Lemma 2.7. *With the above notations, let f_\pm^ε be two functions which are invariant by Φ_\pm^ε in $\Omega^{in} \cap \mathcal{V} = \{\nabla V(x) \cdot \xi > 0\} \cap \mathcal{V}$. Then f_\pm^ε have traces on $\Sigma \cap \mathcal{V}$ that we*

Figure 2: Surfaces Σ_+^i , Σ_+^f and $\Sigma \cap \mathcal{V}$

denote by $(f_\pm^\varepsilon)_{\Sigma, in}$.

Similarly, if f_\pm^ε are invariant by Φ_\pm^s in $\Omega^{out} \cap \mathcal{V} = \{\nabla V(x) \cdot \xi < 0\} \cap \mathcal{V}$. Then f_\pm^ε have traces on $\Sigma \cap \mathcal{V}$ that we denote by $(f_\pm^\varepsilon)_{\Sigma, out}$.

Proof. We first use the functions τ^\pm for defining the four following continuous maps

$$\begin{aligned} \kappa_\pm^i &: \Sigma_\pm^i \rightarrow \Sigma, \\ (x, \xi) &\mapsto \Phi_\pm^{\tau_\pm^i(x, \xi)}(x, \xi), \end{aligned}$$

and

$$\begin{aligned} \kappa_\pm^f &: \Sigma_\pm^f \rightarrow \Sigma, \\ (x, \xi) &\mapsto \Phi_\pm^{\tau_\pm^f(x, \xi)}(x, \xi). \end{aligned}$$

These four maps are homeomorphisms. Set

$$(f_\pm^\varepsilon)_{\Sigma, in} = (\kappa_\pm^i)^* f_\pm^\varepsilon(t_i).$$

These two functions are L^1 functions of $\Sigma \cap \Omega$ which are the traces of f_\pm^ε on Σ .

We argue similarly in Ω^{out} . \square

2.2. Analysis of the kinetic equations (1.8). In this section we prove the existence and uniqueness of solutions to the kinetic equations (1.8) and we analyze the link between the Markov semi-group and this system of equations.

2.2.1. Uniqueness of solutions to (1.8). As a corollary of the analysis of the previous subsection, we obtain that the solutions of (1.8) are unique if they do exist. Indeed, by (1) of Assumptions 1.2, the data has been chosen supported outside Σ so that for short time the kinetic system (1.8) reduces to classical transport by the two flows. We cut the data in a sum of compactly supported pieces so that each of these pieces satisfy the hypothesis of Lemma 2.7 after a certain amount of time. We only need to consider one of these pieces and we take the notations of Lemma 2.7. Because of the previous decomposition of the data, we may assume that $f_\pm^\varepsilon(t_i)$ is compactly

supported in $\mathcal{V} \cap \Omega_{in}$. By Lemma 2.7, if the solution does exist on the interval of time $[t_i, t_f]$, we must have

$$f_+^\varepsilon(t, x, \xi) = \mathbf{1}_{\tau_+ > 0} f_{\Sigma, in}^\varepsilon \left(\Phi^{\tau_+(x, \xi)}(x, \xi) \right) + \mathbf{1}_{\tau_+ < 0} f_{\Sigma, out}^\varepsilon \left(\Phi^{\tau_+(x, \xi)}(x, \xi) \right).$$

As a consequence, the solution for the plus mode will be unique (if it exists) if and only if the outgoing trace $(f_+^\varepsilon)_{\Sigma, out}$ is uniquely determined by the entering traces $(f_+^\varepsilon)_{\Sigma, -}$ and $(f_-^\varepsilon)_{\Sigma, +}$. We shall use the following Lemma, the proof of which is postponed to the end of the section.

We gather in this Lemma properties that we will use in the following and we postpone their proof at the end of the section.

Lemma 2.8. *In the set of distributions, we have*

$$\begin{aligned} (\partial_t + H_\pm) (\tau_\pm(x, \xi) + t) &= 0, & (\partial_t + H_\pm) \Phi_\pm^{\tau_\pm(x, \xi)}(x, \xi) &= 0, \\ H_+ (\mathbf{1}_{\tau_+(x, \xi) < 0}) &= -\Lambda_+(x, \xi) \delta_\Sigma(x, \xi), \end{aligned}$$

where Λ_+ is given by (1.13).

As a consequence of this Lemma, we also have in $\mathcal{D}'(\Omega_{out})$, $H_+ \tau_+ = -1$, which implies

$$H_+ \left(\Phi^{\tau_+(x, \xi)}(x, \xi) \right) = H_+ \tau_+ \dot{\Phi}^{\tau_+(x, \xi)}(x, \xi) + (H_+ \Phi)^{\tau_+(x, \xi)}(x, \xi) = 0.$$

Therefore, the fact that

$$(\partial_t + H_+) f_+^\varepsilon = K_+((f_+^\varepsilon)_{\Sigma, in}, (f_-^\varepsilon)_{\Sigma, out}),$$

implies

$$\Lambda_+ \left((f_+^\varepsilon)_{\Sigma, in} - (f_+^\varepsilon)_{\Sigma, out} \right) = \Lambda_+ \left(T_\varepsilon (f_+^\varepsilon)_{\Sigma, in} - T_\varepsilon (f_-^\varepsilon)_{\Sigma, in} \circ J_+ \right).$$

As a consequence,

$$(f_+^\varepsilon)_{\Sigma, out} = (1 - T_\varepsilon)(f_+^\varepsilon)_{\Sigma, in} - T_\varepsilon(f_-^\varepsilon)_{\Sigma, in} \circ J_+,$$

and the outgoing and ingoing traces are linked. A similar argument holds for the minus mode and as a consequence, the solution of (1.8) is unique if it exists. \square

2.2.2. Link between the particle description and the kinetic model. Let us prove Proposition 1.6. We aim at proving that the semigroup $\mathcal{L}_\varepsilon^t$ provides the unique solution of (1.8). In fact, since the uniqueness property has been proved at the end of Subsection 2.2.1, we only have to prove that if

$$g^\varepsilon(t) = (\mathcal{L}_\varepsilon^t w^\varepsilon(0)), \quad t \in \mathbf{R}^+,$$

and $g_\pm^\varepsilon(t, x, \xi) = g^\varepsilon(t, x, \xi, \pm 1)$, then $(g_+^\varepsilon, g_-^\varepsilon)$ satisfies system (1.8) in $\mathcal{D}'(\mathbf{R}^4 \setminus \{\xi = 0\})$. By density of compactly supported continuous functions in $L^1(\mathbf{R}^4)$, it is enough to prove it for continuous initial data.

By definition, the solutions of system (1.8) include classical transport and jumps on Σ . Let us first consider a point (x_0, ξ_0) which is far from Σ , i.e. such that $\nabla V(x_0) \cdot \xi_0 \neq 0$. Then, there exists a neighborhood Ω_2 of (x_0, ξ_0) and $\varepsilon_2, \tau_2 > 0$ such that for $\varepsilon \in (0, \varepsilon_2)$,

$$\forall (x, \xi) \in \Omega_2, \quad \forall t \in \mathbf{R}, \quad \forall s \in]t - \tau_2, t + \tau_2[, \quad g_\pm^\varepsilon(t, x, \xi) = g_\pm^\varepsilon(s, \Phi_\pm^{-t+s}(x, \xi)).$$

As a consequence,

$$\partial_t g_\pm^\varepsilon + H_\pm g_\pm^\varepsilon = 0 \quad \text{in } \Omega_2,$$

and more generally, $\partial_t g_{\pm}^{\varepsilon} + H_{\pm} g_{\pm}^{\varepsilon} = 0$ for $\nabla V(x) \cdot \xi \neq 0$.

Consider now a point (x_0, ξ_0) such that $\xi_0 \cdot \nabla V(x_0) = 0$. Then, there exists a neighborhood Ω_2 of (x_0, ξ_0) and $\varepsilon_2, \tau_2 > 0$ such that for $\varepsilon \in (0, \varepsilon_2)$, $(x, \xi) \in \Omega_2$, $t \in \mathbf{R}$ and $s \in (t - \tau_2, t + \tau_2)$ (the time τ_2 corresponds to the length of an interval of time during which trajectories issued from points of Ω_2 have at most one jump),

$$\begin{aligned} g_+^{\varepsilon}(t, x, \xi) &= g_+^{\varepsilon}(s, \Phi_+^{-t+s}(x, \xi)) \\ &- \mathbf{1}_{-t+s \leq \tau_+(x, \xi) < 0} T_{\varepsilon}(\Phi_+^{\tau_+(x, \xi)}(x, \xi)) g_+^{\varepsilon}(\tau_+(x, \xi) + t, \Phi_+^{\tau_+(x, \xi)}(x, \xi)) \\ &+ \mathbf{1}_{-t+s \leq \tau_+(x, \xi) < 0} T_{\varepsilon} \circ J_+(\Phi_+^{\tau_+(x, \xi)}(x, \xi)) \\ &\quad \times g_-^{\varepsilon}(\tau_+(x, \xi) + t, J_+(\Phi_+^{\tau_+(x, \xi)}(x, \xi))). \end{aligned}$$

We obtain

$$(2.7) \quad g_+^{\varepsilon}(t, x, \xi) = g_+^{\varepsilon}(s, \Phi_+^{-t+s}(x, \xi)) + \mathbf{1}_{-t+s \leq \tau_+(x, \xi) < 0} \left[(T_{\varepsilon} g_-^{\varepsilon})(\tau_+(x, \xi) + t, J_+(\Phi_+^{\tau_+(x, \xi)}(x, \xi))) - (T_{\varepsilon} g_+^{\varepsilon})(\tau_+(x, \xi) + t, \Phi_+^{\tau_+(x, \xi)}(x, \xi)) \right].$$

Similarly, we have

$$(2.8) \quad g_-^{\varepsilon}(t, x, \xi) = g_-^{\varepsilon}(s, \Phi_-^{-t+s}(x, \xi)) + \mathbf{1}_{-t+s \leq \tau_-(x, \xi) < 0} \left[(T_{\varepsilon} g_+^{\varepsilon})(\tau_-(x, \xi) + t, J_-(\Phi_-^{\tau_-(x, \xi)}(x, \xi))) - (T_{\varepsilon} g_-^{\varepsilon})(\tau_-(x, \xi) + t, \Phi_-^{\tau_-(x, \xi)}(x, \xi)) \right].$$

Note that equations (2.7) and 2.8 are equivalent to the relation $g^{\varepsilon}(t) = \mathcal{L}_{\varepsilon}^{t-s}(g^{\varepsilon}(s))$ and give an explicit expression for the semigroup on a small time $\tau := t - s$ during which the trajectories jump at most once.

The result is then straightforward by Lemma 2.8. For the $+$ mode (the proof for the $-$ mode is similar), we have

$$(\partial_t + H_+)g_+^{\varepsilon} = \mathbf{1}_{\tau_+(x, \xi) + t \geq 0} H_+ (\mathbf{1}_{\tau_+(x, \xi) \leq 0}) \left[(T_{\varepsilon} g_-^{\varepsilon})(\tau_+(x, \xi) + t, J_+(\Phi_+^{\tau_+(x, \xi)}(x, \xi))) - (T_{\varepsilon} g_+^{\varepsilon})(\tau_+(x, \xi) + t, \Phi_+^{\tau_+(x, \xi)}(x, \xi)) \right].$$

Observing that $\Phi_+^{\tau_+(x, \xi)}(x, \xi) = (x, \xi)$ for $(x, \xi) \in \Sigma$, we obtain for $t \geq 0$,

$$(\partial_t + H_+)g_+^{\varepsilon} = \Lambda_+(x, \xi) \delta_{\Sigma}(x, \xi) [T_{\varepsilon}(g_+^{\varepsilon})_{\Sigma, in} - (T_{\varepsilon}(g_-^{\varepsilon})_{\Sigma, in}) \circ J_+]$$

where we have used that for $\tau_+(x, \xi) = 0$ and $t \geq 0$, we have $\mathbf{1}_{\tau_+(x, \xi) + t \geq 0} = 1$. This implies that g_+^{ε} satisfies (1.8). \square

2.2.3. *Proof of Lemma 2.8.* Let us begin with the first line. Recall that we have

$$(\partial_t + H_{\pm})\Phi_{\pm}^{-t} = 0.$$

Similarly, in view of $\tau_{\pm}(\Phi_{\pm}^{-t}(x, \xi)) = \tau_{\pm}(x, \xi) + t$ and writing

$$\Phi_{\pm}^{\tau_{\pm}(x, \xi)}(x, \xi) = \Phi_{\pm}^{\tau_{\pm}(x, \xi) + t}(\Phi_{\pm}^{-t}(x, \xi)),$$

we get $(\partial_t + H_{\pm})(\Phi_{\pm}^{\tau_{\pm}}(x, \xi)) = 0$ and

$$H_{\pm}(\Phi_{\pm}^{\tau_{\pm}}(x, \xi)) = 0.$$

Let us now prove the second line. Note that $\tau_+(x, \xi) = 0$ is an equation of the hypersurface Σ . For calculating $H_+(\mathbf{1}_{\tau_+(x, \xi) < 0})$, we write

$$z = (x, \xi), \quad H_+ a = F(z) \cdot \nabla_z a,$$

where a is a test-function, and we observe that $\nabla_z \cdot F(z) = 0$. Therefore, using that a is compactly supported, Green's formula reads

$$\begin{aligned} \int_{\{\tau_+(z) < 0\}} F(z) \cdot \nabla_z a(z) dz &= \int_{\Sigma} F(z) \cdot n(z) a(z) d\sigma - \int_{\{\tau_+(z) < 0\}} \nabla_z \cdot F(z) a(z) dz \\ &= \int_{\Sigma} F(z) \cdot n(z) a(z) d\sigma \end{aligned}$$

where $n(z)$ is the unitary exterior normal vector to Σ :

$$n(x, \xi) = (|d^2V(x)\xi|^2 + |\nabla V(x)|^2)^{-1/2} \begin{pmatrix} d^2V(x)\xi \\ \nabla V(x) \end{pmatrix}.$$

Note that H_+ is transverse to Σ and points towards the region $\tau_+ < 0$. Besides, one can check that $H_+(x, \xi) \cdot n(x, \xi) = \Lambda_+(x, \xi)$, where Λ_+ is defined by (1.13). At this stage of the proof, we have obtained for any smooth compactly supported function a ,

$$\begin{aligned} \int (H_+ \mathbf{1}_{\tau_+(x, \xi) < 0}) a dx d\xi &= - \int \mathbf{1}_{\tau_+(x, \xi) < 0} (H_+ a) dx d\xi \\ &= - \int_{\Sigma} a (H_+ \cdot n) d\sigma \\ &= - \int_{\Sigma} \Lambda_+ a d\sigma, \end{aligned}$$

which gives the result. \square

2.3. Mathematical justification of the approximation by the semi-group.

2.3.1. *Strategy.* In this section, we prove Proposition 1.7. The proof relies on a characterization of $w^{\varepsilon}(t)$ via pseudodifferential operators. Recall that if $a \in \mathcal{C}_0^{\infty}(\mathbf{R}^2)$, the semiclassical pseudodifferential operator of symbol a is defined by the Weyl quantization rule

$$\text{op}_{\varepsilon}(a)f(x) = (2\pi)^{-2} \int e^{i\xi \cdot (x-y)} a\left(\frac{x+y}{2}, \varepsilon\xi\right) f(y) dy d\xi, \quad f \in \mathcal{S}(\mathbf{R}^2).$$

This operator extends to functions $f \in L^2(\mathbf{R}^2)$ and one can prove that $(\text{op}_{\varepsilon}(a))$ is a uniformly bounded family of operators of $\mathcal{L}(L^2(\mathbf{R}^2))$ since there exists a constant $C > 0$ and an integer $N \in \mathbf{N}^*$ such that

$$(2.9) \quad \forall a \in \mathcal{C}_0^{\infty}(\mathbf{R}^4), \quad \|\text{op}_{\varepsilon}(a)\|_{\mathcal{L}(L^2(\mathbf{R}^2))} \leq C \sup_{|\beta| \leq 3} \sup_{\xi \in \mathbf{R}^2} \int_{\mathbf{R}^2} |\partial_x^{\beta} a(x, \xi)| dx.$$

We refer to the books [1, 8, 34] for a complete study of pseudodifferential operators. The estimate (2.9) is not the standard Calderon-Vaillancourt estimate (see [4]) that is usually used. It has the advantage not to differentiate in the variable ξ and is inspired from [21] (see also the survey [2]). A short proof is given in the Appendix

for the convenience of the reader, we also recall the single symbolic calculus result that we shall use.

Denote by $\mathbf{C}^{2,2}$ the set of 2×2 complex matrices and consider symbols a that are matrix-valued : $a \in \mathcal{C}_0^\infty(\mathbf{R}^2, \mathbf{C}^{2,2})$. Then, the operator $\text{op}_\varepsilon(a)$ is a matrix-valued operator acting on functions of $L^2(\mathbf{R}^2, \mathbf{C}^2)$. If w^ε is the Wigner transform of the matrix density ϱ^ε , we have the relation

$$(2.10) \quad \langle a, w^\varepsilon \rangle = \text{tr}(\text{op}_\varepsilon(a)\varrho^\varepsilon),$$

where the bracket between the two matrices $a = (a_{i,j})$ and $w^\varepsilon = (w_{i,j}^\varepsilon)$ is defined by

$$(2.11) \quad \langle a, w^\varepsilon \rangle = \sum_{i,j} \int_{\mathbf{R}^4} a_{i,j}(x, \xi) w_{j,i}^\varepsilon(x, \xi) dx d\xi.$$

We shall use this description of $\langle a, w^\varepsilon \rangle$ in order to prove Proposition 1.7.

Remark 2.9. The relations (2.10) and (2.9) imply that, under Assumption 1.2, the family $(w^\varepsilon(t))_{\varepsilon>0}$ is a bounded family in the set of distributions.

Since the initial density matrix $\varrho^\varepsilon(0)$ is supposed to be a Hilbert-Schmidt operator, there exists a sequence $(\lambda_j)_{j \in \mathbf{N}}$ of $\ell^2(\mathbf{N})$ and a sequence of bounded normalized families $(\psi_{0,j}^\varepsilon)_{j \in \mathbf{N}}$ of $L^2(\mathbf{R}^2)$ such that

$$\varrho^\varepsilon(0) = \sum_{j \in \mathbf{N}} \lambda_j |\psi_{0,j}^\varepsilon\rangle \langle \psi_{0,j}^\varepsilon|.$$

As a consequence, for $t \in \mathbf{R}$,

$$\varrho^\varepsilon(t) = \sum_{j \in \mathbf{N}} \lambda_j |\psi_j^\varepsilon(t)\rangle \langle \psi_j^\varepsilon(t)|,$$

where for any $j \in \mathbf{N}$, the family $(\psi_j^\varepsilon(t))_{\varepsilon>0}$ is a family of solutions to the Dirac equation

$$(2.12) \quad i\varepsilon \partial_t \psi_j^\varepsilon = (A(\varepsilon D) + V(x)) \psi_j^\varepsilon,$$

with initial data $\psi_j^\varepsilon(t) = \psi_{0,j}^\varepsilon$. Besides, the relation (2.10) yields

$$\langle a, w^\varepsilon \rangle = \text{tr}(\text{op}_\varepsilon(a)\varrho^\varepsilon) = \sum_{j \in \mathbf{N}} \lambda_j (\text{op}_\varepsilon(a)\psi_j^\varepsilon(t), \psi_j^\varepsilon(t))_{L^2(\mathbf{R}_x^2)}.$$

We denote by $w_j^\varepsilon(t)$ the Wigner transform of the family $(\psi_j^\varepsilon(t))_{\varepsilon>0}$ which is defined by the relation

$$\forall a \in \mathcal{C}_0^\infty(\mathbf{R}^4), \quad \langle a, w_j^\varepsilon(t) \rangle = (\text{op}_\varepsilon(a)\psi_j^\varepsilon(t), \psi_j^\varepsilon(t))_{L^2(\mathbf{R}_x^2)}.$$

The Wigner function $w_j^\varepsilon(t)$ is a 2 by 2 matrix and the bracket involved in the preceding relation is also the one defined in (2.11). In the following, we will characterize $w_j^\varepsilon(t)$ in terms of the flow $\mathcal{L}_{\varepsilon,R}^t$. More precisely, we are going to prove that for any $j \in \mathbf{N}$, $w_j^\varepsilon(t)$ satisfies Proposition 1.7, which gives the result for $w^\varepsilon(t)$.

For this purpose, we use the results of [29, 17, 18, 19] which are stated for a Schrödinger equation with matrix-valued potential. This comes from the following observation : whenever $V(x) = |x|^2$, the operator $A(\varepsilon D) + V(x)$ becomes a Schrödinger operator with a matrix-valued potential by taking the Fourier transform. As a consequence, the methods developed in [29, 17] for Schrödinger equation

with matrix-valued potential can be adapted to our setting. Furthermore, conical intersections have been classified in [6] and [15] and the Dirac-type equation (1.1), like the Schrödinger equations of [17, 18, 19], enters in the same class of crossings. Thus, it is not surprising that similar methods do apply. Note however that the jumps were omitted in [29] and [17]; as mentioned in [19], these jumps are required for the correctness of the proof of [17].

Then, the main steps of the proof will consist in :

- (1) The transport outside Σ .
- (2) Localization in energy and use of space-time variables.
- (3) A normal form which reduces to a simple model called the Landau-Zener system.
- (4) The computation of the transitions on Σ that is performed via the normal form and the Landau-Zener system.

In the following, we will take advantage of Remark 1.1 for taking into account only the jumps which occur inside the set

$$\mathcal{U}_{\varepsilon,R} = \{(x, \xi) \in \mathbf{R}^4, \quad |\xi| \leq R\sqrt{\varepsilon}\}.$$

Let us introduce the semi-group $\mathcal{L}_{\varepsilon,R}$ which restricts the jumps to those occurring inside $\mathcal{U}_{\varepsilon,R}$, the semi-group $\mathcal{L}_{\varepsilon,R}$ differs from $\mathcal{L}_{\varepsilon}$ by exponentially small terms. It is this semi-group that we shall consider now. The real number R will be chosen as $R = \varepsilon^{-1/8}$ (see Notation 1).

Let us now detail these steps. For simplicity, we omit the index “ j ” and simply consider a family $(\psi^\varepsilon(t))_{\varepsilon>0}$, uniformly bounded in $L^2(\mathbf{R}^2)$, of solutions to the Dirac equation (2.12) with initial data $(\psi_0^\varepsilon)_{\varepsilon>0}$ and we denote by $w^\varepsilon(t)$ its Wigner transform at time t . We also denote by $w_\pm^\varepsilon(t)$ the scalar quantities $w_\pm^\varepsilon(t) = \text{tr}(\Pi^\pm w^\varepsilon(t))$.

2.3.2. *The transport outside the transition region.* The analogue of Proposition 2.3 in [19] is the following

Proposition 2.10. *Let $c \in \mathcal{C}_c^\infty(\mathbf{R}^4, \mathbf{C})$, and let $b \in \mathcal{C}^\infty(\mathbf{R}^2, \mathbf{C})$ with ∇b compactly supported. If there exist $C > 0$ and $s_0 > 0$ such that*

$$\forall r \in [-s_0, s_0] : \quad \Phi_\pm^r(\text{supp}(c)) \cap \mathcal{U}_{\varepsilon,R} = \emptyset$$

then for all $\chi \in \mathcal{C}_c^\infty(\mathbf{R}, \mathbf{R})$ and for all $s \in [t - s_0, t + s_0]$

$$\begin{aligned} \text{tr} \int_{\mathbf{R}^{2d+1}} \chi(t) c(x, \xi) b\left(\frac{\xi}{R\sqrt{\varepsilon}}\right) w_\pm^\varepsilon(t, x, \xi) dx d\xi dt = \\ \text{tr} \int_{\mathbf{R}^{2d+1}} \chi(t) c(x, \xi) b\left(\frac{\xi}{R\sqrt{\varepsilon}}\right) (w_\pm^\varepsilon(s, \Phi_\pm^{-t+s})(x, \xi)) dx d\xi dt \\ + \mathcal{O}(1/(\sqrt{\varepsilon}R^5)) + \mathcal{O}(1/R^2) + \mathcal{O}(\varepsilon). \end{aligned}$$

This proposition is a refined version of the resolution of the kinetic system (1.4). Indeed, solving (1.4) in a subset of the domain $\{|\xi| > \delta_0\}$ for some δ_0 on an interval of time $[-s_0, s_0]$ such that the trajectories $(\Phi_\pm^s)_{s \in [-s_0, s_0]}$ do not reach $\{\xi = 0\}$ gives

$$w_\pm^\varepsilon(t, x, \xi) = w_\pm^\varepsilon(t - s, \Phi_\pm^{-s}(x, \xi)) + \mathcal{O}(\varepsilon) \text{ in } \mathcal{D}'(\mathbf{R}^4).$$

Proposition 2.10 authorizes to be at a distance of order $\mathcal{O}(R\sqrt{\varepsilon})$ of $\{\xi = 0\}$

Since for scalar symbols a , we have

$$\frac{d}{dt} \langle a\Pi^\pm, w^\varepsilon(t) \rangle = \left(\frac{1}{i\varepsilon} [\text{op}_\varepsilon(a\Pi^\pm), A(\varepsilon D) + V] \psi^\varepsilon(t), \psi^\varepsilon(t) \right)_{L^2(\mathbf{R}_x^2)},$$

the proof of this proposition relies on a good understanding of the operator

$$L_\varepsilon = \frac{1}{i\varepsilon} [\text{op}_\varepsilon(a\Pi^\pm), A(\varepsilon D) + V].$$

The main ingredients are the two following observations :

- For $a(x, \xi) = c(x, \xi) b\left(\frac{\xi}{R\sqrt{\varepsilon}}\right)$, the symbol $a(x, \xi)\Pi^+(\xi)$ is smooth and we have

$$(2.13) \quad \forall \alpha, \beta \in \mathbf{N}^2, \exists C_{\alpha, \beta} > 0, \left| \partial_x^\alpha \partial_\xi^\beta (a(x, \xi)\Pi^+(\xi)) \right| \leq C (R\sqrt{\varepsilon})^{-|\beta|},$$

so that we can use the symbolic calculus theorems of the Appendix, paying attention to the rest terms.

- If B is an off-diagonal symbol, that is a symbol which satisfies

$$B = \Pi^+ B \Pi^- + \Pi^- B \Pi^+,$$

the quantities

$$\int \chi(t) (\text{op}_\varepsilon(B) \psi^\varepsilon(t), \psi^\varepsilon(t)) dt,$$

which seems to be of order $\mathcal{O}(1)$, can be proved to be of smaller order than expected by re-using the equation satisfied by $\psi^\varepsilon(t)$.

Proof of Proposition 2.10. Let us now focus on the proof itself. Using Proposition 4.1 and observing that $A(\xi) = |\xi|(\Pi^+ - \Pi^-)$ with $\Pi^+ + \Pi^- = 1$, we obtain

$$L_\varepsilon = -\text{op}_\varepsilon \left((\nabla V(x) \cdot \nabla_\xi a - \frac{\xi}{|\xi|} \cdot \nabla_x a) \Pi^+ \right) + \text{op}_\varepsilon(B) + \mathcal{O}(R^{-2}) + \mathcal{O}(\varepsilon),$$

with

$$\begin{aligned} B &= -a \nabla V \cdot \nabla \Pi^+ - \frac{1}{2} (|\xi|(\{a\Pi^+, \Pi^+\} - \{\Pi^+, a\Pi^+\})) \\ &\quad - \frac{1}{2} (|\xi|(\{a\Pi^+, \Pi^-\} - \{\Pi^-, a\Pi^+\})) \\ &= -a \nabla V \cdot \nabla \Pi^+ - |\xi| (\{a\Pi^+, \Pi^+\} - \{\Pi^+, a\Pi^+\}) \\ &= -a \nabla V \cdot \nabla \Pi^+ + |\xi| (\Pi^+ \nabla_x a \cdot \nabla \Pi^+ + \nabla_x a \cdot \nabla \Pi^+ \Pi^+) \\ &= -a \nabla V \cdot \nabla \Pi^+ + |\xi| \nabla_x a \cdot \nabla \Pi^+. \end{aligned}$$

Here we have used $\nabla \Pi^+ = \Pi^+ \nabla \Pi^+ + \nabla \Pi^+ \Pi^+$. As a consequence, B is an off-diagonal symbol. We write $B = B_0 + B_1$ with $B_0 = -a \nabla V \cdot \nabla \Pi^+$, we have

$$(2.14) \quad \forall \alpha, \beta \in \mathbf{N}^d, \exists C_{\alpha, \beta} > 0, \forall x, \xi \in \mathbf{R}^d, \left| \partial_\xi^\beta \partial_x^\alpha B_j \right| \leq C_{\alpha, \beta} (R\sqrt{\varepsilon})^{-|\beta|-1+j}.$$

The result comes from the next lemma which concludes the proof. \square

Lemma 2.11. *For any $\chi \in \mathcal{C}_0^\infty(\mathbf{R})$, we have*

$$\begin{aligned} \int \chi(t) (\text{op}_\varepsilon(B_1) \psi^\varepsilon(t), \psi^\varepsilon(t)) &= \mathcal{O}(R^{-2}) + \mathcal{O}(\varepsilon), \\ \int \chi(t) (\text{op}_\varepsilon(B_0) \psi^\varepsilon(t), \psi^\varepsilon(t)) &= \mathcal{O}(R^{-5} \varepsilon^{-1/2}) + \mathcal{O}(\sqrt{\varepsilon}). \end{aligned}$$

Proof. We begin with B_1 . Since B_1 is off-diagonal, we can write

$$\begin{aligned} B_1 &= [(\Pi^- B_1 \Pi^+ - \Pi^+ B_1 \Pi^-)(2|\xi|)^{-1}, A(\xi)] \\ &= [(\Pi^- B_1 \Pi^+ - \Pi^+ B_1 \Pi^-)(2|\xi|)^{-1}, \tau + V(x) + A(\xi)]. \end{aligned}$$

After quantization, we get

$$\begin{aligned} \text{op}_\varepsilon(B_1) &= \left[\text{op}_\varepsilon((\Pi^- B_1 \Pi^+ - \Pi^+ B_1 \Pi^-)(2|\xi|)^{-1}), \frac{\varepsilon}{i} \partial_t + V(x) + A(\varepsilon D) \right] \\ &\quad + \mathcal{O}(R^{-2}) + \mathcal{O}(\sqrt{\varepsilon}/R). \end{aligned}$$

Once applied to ψ^ε which satisfies the Dirac equation (2.12), we obtain the announced relation.

Note that we have obtained more generally that if B_j is off-diagonal and satisfies the relation (2.14), then

$$\begin{aligned} (2.15) \quad \int \chi(t) (\text{op}_\varepsilon(B_j) \psi^\varepsilon(t), \psi^\varepsilon(t)) &= \mathcal{O}((R\sqrt{\varepsilon})^{-1+j}) (R^{-2} + \sqrt{\varepsilon}/R) \\ &= \mathcal{O}((R\sqrt{\varepsilon})^{-1+j}) (R^{-2} + \varepsilon) \end{aligned}$$

In particular, for B_0 , we obtain

$$\int \chi(t) (\text{op}_\varepsilon(B_0) \psi^\varepsilon(t), \psi^\varepsilon(t)) = \mathcal{O}(R^{-3} \varepsilon^{-1/2}) + \mathcal{O}(\sqrt{\varepsilon}/R)$$

that we want to improve. Therefore, we go one step further in the symbolic calculus and we write

$$\begin{aligned} \text{op}_\varepsilon(B_0) &= \left[\text{op}_\varepsilon((\Pi^- B_0 \Pi^+ - \Pi^+ B_0 \Pi^-)(2|\xi|)^{-1}), \frac{\varepsilon}{i} \partial_t + V(x) + A(\varepsilon D) \right] \\ &\quad + \frac{\varepsilon}{i} \text{op}_\varepsilon(\{(\Pi^- B_0 \Pi^+ - \Pi^+ B_0 \Pi^-)(2|\xi|)^{-1}, \tau + V(x)\}) \\ &\quad + \frac{\varepsilon}{2i} \text{op}_\varepsilon(\{(\Pi^- B_0 \Pi^+ - \Pi^+ B_0 \Pi^-)(2|\xi|)^{-1}, A(\xi)\}) \\ &\quad - \{A(\xi), (\Pi^- B_0 \Pi^+ - \Pi^+ B_0 \Pi^-)(2|\xi|)^{-1}\} \\ &\quad + \mathcal{O}(R^{-4}) + \mathcal{O}(\sqrt{\varepsilon}/R^3). \end{aligned}$$

Paying attention to all these terms, we observe that

$$\begin{aligned} \text{op}_\varepsilon(B_0) &= \left[\text{op}_\varepsilon((\Pi^- B_0 \Pi^+ - \Pi^+ B_0 \Pi^-)(2|\xi|)^{-1}), \frac{\varepsilon}{i} \partial_t + V(x) + A(\varepsilon D) \right] \\ &\quad - i\varepsilon \text{op}_\varepsilon(\nabla V \cdot \nabla_\xi ((\Pi^- B_0 \Pi^+ - \Pi^+ B_0 \Pi^-)(2|\xi|)^{-1})) + \mathcal{O}(R^{-2}) + \mathcal{O}(\sqrt{\varepsilon}). \end{aligned}$$

The matrix

$$B_{-2} = \nabla V \cdot \nabla_\xi ((\Pi^- B_0 \Pi^+ - \Pi^+ B_0 \Pi^-)(2|\xi|)^{-1})$$

satisfies (2.14) with $j = -2$ and we claim that B_{-2} is also off-diagonal. As a consequence, equation (2.15) gives

$$\int \chi(t) (\text{op}_\varepsilon(B_{-2}) \psi^\varepsilon(t), \psi^\varepsilon(t)) dt = \mathcal{O}(R^{-5} \varepsilon^{-3/2}) + \mathcal{O}(R^{-3} \varepsilon^{-1/2}),$$

which concludes the proof of Lemma 2.11.

It remains to prove the claim, a simple calculus shows that

$$B_{-2} = -\frac{1}{2} |\xi|^{-3} (\nabla V \cdot \xi) (\Pi^- B_0 \Pi^+ - \Pi^+ B_0 \Pi^-) + (2|\xi|)^{-1} \nabla V \cdot \nabla_\xi (\Pi^- B_0 \Pi^+ - \Pi^+ B_0 \Pi^-).$$

Therefore,

$$\begin{aligned}
\Pi^\pm B_{-2} \Pi^\pm &= (2|\xi|)^{-1} \Pi^\pm (\nabla V \cdot \nabla \Pi^- B_0 \Pi^+ + \Pi^- B_0 \nabla V \cdot \nabla \Pi^+) \Pi^\pm \\
&\quad - (2|\xi|)^{-1} \Pi^\pm (\nabla V \cdot \nabla \Pi^+ B_0 \Pi^- + \Pi^+ B_0 \nabla V \cdot \nabla \Pi^-) \Pi^\pm \\
&= |\xi|^{-1} \Pi^\pm (-\nabla V \cdot \nabla \Pi^+ B_0 \Pi^+ + \Pi^+ B_0 \nabla V \cdot \nabla \Pi^+) \Pi^\pm \\
&= |\xi|^{-1} \Pi^\pm [B_0, \nabla V \cdot \nabla \Pi^+] \Pi^\pm \\
&= 0
\end{aligned}$$

since $B_0 = a \nabla V \cdot \nabla \Pi^+$, which proves that B_{-2} is off-diagonal. \square

Notation 1. In the following, it will be convenient to denote by η_ε any rest term smaller than $\mathcal{O}(1/(R^5\sqrt{\varepsilon})) + \mathcal{O}(1/R^2) + \mathcal{O}(\sqrt{\varepsilon}) + \mathcal{O}(R^3\sqrt{\varepsilon})$. The term in $R^3\sqrt{\varepsilon}$ will be useful in the following. Note that when $R = \varepsilon^{-1/8}$, we have $\eta_\varepsilon = \mathcal{O}(\varepsilon^{1/8})$.

2.3.3. Localization in energy. The memory of the mode by use of a matrix-valued symbol of the form $a\Pi^+$ or $a\Pi^-$, with a scalar, can be replaced by a localization in energy. This requires to work in space time variables and has the advantage that we are reduced to use scalar symbols. Using scalar symbols will be convenient in the next section when we will perform a normal form and use a Fourier Integral Operator. The energy surfaces of the space-time phase space $\mathbf{R}_{t,x,\tau,\xi}^6$ are the sets

$$(2.16) \quad E^\pm = \{(t, x, \tau, \xi) \in \mathbf{R}^6, \tau = \mp|\xi| - V(x)\}.$$

Recall that the dual variable of the time t is interpreted as an energy τ .

In the following, we shall use semi-classical pseudo differential operators with symbols depending on the variable $(t, x, \tau, \xi) \in \mathbf{R}^6$ with the choice of the Weyl quantization in the time variables, as it was already the case for the space variables.

The localization in energy is done by use of a cut-off function $\theta \in \mathcal{C}_0^\infty(\mathbf{R})$ such that $0 \leq \theta \leq 1$, $\theta(\tau) = 0$ for $|\tau| > 1$ and $\theta(\tau) = 1$ for $|\tau| < 1/2$. This function θ is fixed from now on.

Lemma 2.12. *Let $a \in \mathcal{C}_0^\infty(\mathbf{R}^4)$ and set $c_{\varepsilon,R}(x, \xi) = a(x, \xi)(1 - \theta)(2\xi/(R\sqrt{\varepsilon}))$, then for all $\chi \in \mathcal{C}_0^\infty(\mathbf{R})$,*

$$\begin{aligned}
\int_{\mathbf{R}} \chi(t) (\text{op}_\varepsilon(c_{\varepsilon,R}\Pi^\pm)\psi^\varepsilon(t), \psi^\varepsilon(t))_{L^2(\mathbf{R}^2)} dt &= \mathcal{O}(1/R^2) + \mathcal{O}(\sqrt{\varepsilon}) \\
&\quad + \left(\text{op}_\varepsilon \left(\chi(t)c_{\varepsilon,R}(x, \xi) \theta \left(\frac{\tau \pm |\xi| + V(x)}{R\sqrt{\varepsilon}} \right) \right) \psi^\varepsilon, \psi^\varepsilon \right)_{L^2(\mathbf{R}_{t,x}^3)}.
\end{aligned}$$

Remark 2.13. It is because the localization in energy is made in balls of size $\sqrt{\varepsilon}$ that we need to perform the jumps: they guarantee that the energy of the created trajectory do not differ at order $\mathcal{O}(\sqrt{\varepsilon})$ but at least at order $\mathcal{O}(\varepsilon)$.

Remark 2.14. Note that the presence of the eigenprojector in the symbol induces restriction on both components of the function $\psi^\varepsilon(t)$. Indeed, by the symbolic calculus of the Appendix and by equation (2.13) we have

$$\begin{aligned}
\int_{\mathbf{R}} \chi(t) (\text{op}_\varepsilon(c_{\varepsilon,R}\Pi^\pm)\psi^\varepsilon(t), \psi^\varepsilon(t))_{L^2(\mathbf{R}^2)} dt \\
= \int_{\mathbf{R}} \chi(t) (\text{op}_\varepsilon(c_{\varepsilon,R}\Pi^\pm)\psi^\varepsilon(t), \Pi^\pm\psi^\varepsilon(t))_{L^2(\mathbf{R}^2)} dt + \mathcal{O}(R\sqrt{\varepsilon}).
\end{aligned}$$

Proof of Lemma 2.12. We set

$$\theta_{\varepsilon,R}^{\pm}(x, \tau, \xi) = \theta\left(\frac{\tau \pm |\xi| + V(x)}{R\sqrt{\varepsilon}}\right).$$

Following the lines of the proof of Lemma 5.1 in [18], we observe that since $1 - \theta$ vanishes identically close to 0, one can write

$$1 - \theta_{\varepsilon,R}^+(x, \tau, \xi) = \frac{1}{R\sqrt{\varepsilon}}(\tau + |\xi| + V(x))G\left(\frac{\tau + |\xi| + V(x)}{R\sqrt{\varepsilon}}\right)$$

for some smooth function G , with

$$(\tau + |\xi| + V(x))\Pi^+(\xi) = \Pi^+(\xi)(\tau + A(\xi) + V(x)).$$

Therefore, we can use the equation satisfied by $\psi^\varepsilon(t)$, symbolic calculus and the estimate (2.13) to obtain

$$\left(\text{op}_\varepsilon(\chi_{c_{\varepsilon,R}}\Pi^+)\psi^\varepsilon, \psi^\varepsilon\right) = \left(\text{op}_\varepsilon(\chi_{c_{\varepsilon,R}}\theta_{\varepsilon,R}^+\Pi^+)\psi^\varepsilon, \psi^\varepsilon\right) + \mathcal{O}(R^{-2}) + \mathcal{O}(\varepsilon).$$

It remains to get rid of the matrix $\Pi^+(\xi)$. In view of

$$\chi_{c_{\varepsilon,R}}\theta_{\varepsilon,R}^+ = \chi_{c_{\varepsilon,R}}\theta_{\varepsilon,R}^+\Pi^+ + \chi_{c_{\varepsilon,R}}\theta_{\varepsilon,R}^+\Pi^-,$$

we only need to prove that $\left(\text{op}_\varepsilon(\chi_{c_{\varepsilon,R}}\theta_{\varepsilon,R}^+\Pi^-)\psi^\varepsilon, \psi^\varepsilon\right) = \mathcal{O}(\eta_\varepsilon)$. We observe that $\theta_{\varepsilon,R}^+c_{\varepsilon,R} = \theta_{\varepsilon,R}^+(1 - \theta_{\varepsilon,R}^-)c_{\varepsilon,R}$, and

$$\begin{aligned} (1 - \theta_{\varepsilon,R}^-)\Pi^- &= \frac{1}{R\sqrt{\varepsilon}}G\left(\frac{\tau - |\xi| + V(x)}{R\sqrt{\varepsilon}}\right)(\tau - |\xi| + V(x))\Pi^- \\ &= \frac{1}{R\sqrt{\varepsilon}}G\left(\frac{\tau - |\xi| + V(x)}{R\sqrt{\varepsilon}}\right)\Pi^-(\tau + A(\xi) + V(x)). \end{aligned}$$

By using again the equation, symbolic calculus and estimate (2.13), we can write

$$\begin{aligned} &\left(\text{op}_\varepsilon(\chi_{a^+}\theta_{\varepsilon,R}^+\Pi^-)\psi^\varepsilon, \psi^\varepsilon\right) \\ &= \left(\text{op}_\varepsilon(\chi_{a^+}\theta_{\varepsilon,R}^+(1 - \theta_{\varepsilon,R}^-)\Pi^-)\psi^\varepsilon, \psi^\varepsilon\right) + \mathcal{O}(R^{-2}) + \mathcal{O}(\varepsilon) = \mathcal{O}(\eta_\varepsilon). \end{aligned}$$

The proof for the minus mode is similar. \square

2.3.4. The normal form. For computing the transitions, we use a normal form result. For this, we need to work microlocally in space-time variables. Following [6, 14, 13], close to a point $(t_0, x_0, \xi_0 = 0, \tau_0 = -V(x_0))$, there exist a change of coordinates

$$\kappa : (s, z, \sigma, \zeta) \mapsto (t, x, \tau, \xi)$$

with $s, \sigma \in \mathbf{R}$ and $z = (z_1, z_2)$, $\zeta = (\zeta_1, \zeta_2) \in \mathbf{R}^2$, and a matrix B such that

$$(2.17) \quad (\tau + V + A(\xi)) \circ \kappa = {}^t B \left(-\sigma + \tilde{A}(s, z_1) \right) B,$$

with

$$\tilde{A}(s, z_1) = \begin{pmatrix} s & z_1 \\ z_1 & -s \end{pmatrix}.$$

Moreover, this change of coordinates preserves the symplectic structure of the phase space $\mathbf{R}_{t,x}^3 \times \mathbf{R}_{\tau,\xi}^3$: the variables σ and ζ are respectively the dual variables of s and z . Besides, in view of Section 6.2 of [14], there exists a function $\gamma > 0$ such that ${}^t B B = \gamma \text{Id}$.

The construction of the canonical transform is based on the vectors H and H' defined in 2.1 and 2.2 (see [6] and the analysis performed in [17, 13]). The variable s is chosen such that the trajectories which reach $\{\xi = 0\}$ are included in $\{s < 0\}$ and those which leave ($\xi = 0$) are included in $\{s > 0\}$. Besides, one extends the vectors H and H' as vectors of $T^*\mathbf{R}_{t,x}^3$ by adding the coordinate 1 along ∂_t and the coordinate 0 along ∂_τ and we keep calling them H and H' . The resulting vectors are the limit on $\{\xi = 0\}$ and along the flows of the Hamiltonian vector fields associated with the functions $\tau + V(x) \pm |\xi|$. They are sent by $d\kappa$ on the limit on $\{s = z_1 = 0\}$ and along the flow of the Hamiltonian vector fields associated with $\gamma^2(-\sigma \pm \sqrt{s^2 + z_1^2})$. A simple calculus shows that since the canonic symplectic form ω is preserved by canonical transform, the relation (2.3) implies that H' is sent on $\gamma^2(-\partial_s + \partial_\sigma)$, the limit as s goes to 0^- of the Hamiltonian field associated with $\gamma^2(-\sigma - \sqrt{s^2 + z_1^2})$ and H' is sent on on $\gamma^2(-\partial_s - \partial_\sigma)$, the limit as s goes to 0^- of the Hamiltonian field associated with $\gamma^2(-\sigma + \sqrt{s^2 + z_1^2})$. This observation allows to relate the modes after the change of coordinates.

As a consequence, in these new variables (s, z, σ, ζ) , the geometry of the crossing is simple and we have

$$(2.18) \quad \begin{aligned} S &:= \{\xi = 0, \tau + V(x) = 0\} = \kappa(\{s = 0, z_1 = 0, \sigma = 0\}), \\ E^\pm &= \kappa\left(\{-\sigma \mp \sqrt{s^2 + z_1^2} = 0\}\right), \end{aligned}$$

where the energy sets E_\pm are defined by (2.16).

Finally, in the construction of the canonical form κ , the function $z_1(x, \xi)$ can be related to the variables x and ξ according to

$$(2.19) \quad z_1(x, \xi) = \xi \wedge \frac{\nabla V(x)}{|\nabla V(x)|^{3/2}} + \mathcal{O}(|\xi|^2).$$

Similar formula can be written for the functions s and σ . However, in the sequel, we will only use the formula for z_1 .

Then, thanks to Theorem 3 of [6], it is possible to pass equation (2.17) at the quantum level : there exists a unitary operator K_ε and a matrix B_1 such that

$$\begin{aligned} K_\varepsilon \text{op}_\varepsilon({}^t B_\varepsilon) \text{op}_\varepsilon(-\sigma + \tilde{A}(s, z_1)) \text{op}_\varepsilon(B_\varepsilon) (K_\varepsilon)^* \\ = \text{op}_\varepsilon(\tau + V(x) + A(\xi)) + \mathcal{O}(\varepsilon^2), \end{aligned}$$

where $B_\varepsilon = B + \varepsilon B_1$. The operator K_ε is a Fourier Integral Operator associated with the canonical transform κ (see [8] or [14]). It allows to pass at the quantum level the relation (2.17) induced by the change of variables κ . An important property of these Fourier Integral Operators is that they are compatible with pseudo differential calculus in the sense that for all $a \in \mathcal{C}_0^\infty(\mathbf{R}^6)$,

$$K_\varepsilon \text{op}_\varepsilon(a \circ \kappa^{-1}) K_\varepsilon^* = \text{op}_\varepsilon(a) + \mathcal{O}(\varepsilon N_\varepsilon(a)),$$

where

$$N_\varepsilon(a) = \sup_{|\alpha| \leq K, |\beta| \leq 2, |\alpha| + |\beta| \geq 2} \sup_{(\tau, \xi) \in \mathbf{R}^3} \int_{\mathbf{R}^3} \left| \partial_{t,x}^\alpha \partial_{\tau, \xi}^\beta a(t, x, \tau, \xi) dt dx \right|$$

for some $K \in \mathbf{N}$. In particular, when one applies this relation to a symbol of the form

$$a_{\varepsilon,R}(x, \xi) = \chi(t)c_{\varepsilon,R}(x, \xi)\theta\left(\frac{\tau \pm |\xi| + V(x)}{R\sqrt{\varepsilon}}\right)\Pi^\pm(\xi),$$

one gets

$$N_\varepsilon(a_{\varepsilon,R}) = \mathcal{O}((R\sqrt{\varepsilon})^{-2}) + \mathcal{O}(1),$$

and $\varepsilon N_\varepsilon(a_{\varepsilon,R}) = \mathcal{O}(\eta_\varepsilon)$. We will use this property to translate the quantities that we want to study in the variables (t, x, τ, ξ) in these new variables (s, z, σ, ζ) . More precisely, we set

$$v^\varepsilon = \text{op}_\varepsilon(B_\varepsilon)K_\varepsilon^*\psi^\varepsilon,$$

then v^ε solves (microlocally in $L^2(\mathbf{R}_{s,z}^3)$) the system

$$(2.20) \quad \frac{\varepsilon}{i}\partial_s v^\varepsilon = \tilde{A}(s, z_1)v^\varepsilon + \mathcal{O}(\varepsilon^2)$$

and

$$(\text{op}_\varepsilon(({}^tBaB) \circ \kappa)\psi^\varepsilon, \psi^\varepsilon)_{L^2(\mathbf{R}_{t,x}^3)} = (\text{op}_\varepsilon(a)v^\varepsilon, v^\varepsilon)_{L^2(\mathbf{R}_{s,z}^3)} + \mathcal{O}(\eta_\varepsilon)$$

where η_ε denotes a rest term as defined in Notation 1. In particular, for scalar functions a , we have

$$(\text{op}_\varepsilon((\gamma a) \circ \kappa)\psi^\varepsilon, \psi^\varepsilon)_{L^2(\mathbf{R}_{t,x}^3)} = (\text{op}_\varepsilon(a)v^\varepsilon, v^\varepsilon)_{L^2(\mathbf{R}_{s,z}^3)} + \mathcal{O}(\eta_\varepsilon)$$

In what follows, we shall focus on the analysis of this family v^ε .

Let us now write the Markov process $\mathcal{L}_{\varepsilon,R}$ in the new coordinates, we shall denote by $\tilde{\mathcal{L}}_{\varepsilon,R}$ the resulting semi-group.

- The *hopping region* is a neighborhood of size $R\sqrt{\varepsilon}$ of the intersection of the energy sets E^+ and E^- , therefore it can be replaced by

$$\tilde{\mathcal{U}}_{\varepsilon,R} = \{z_1^2 + s^2 \leq C_1 R^2 \varepsilon\}$$

for some constant $C_1 > 0$. As already noticed before, turning C_1 into some $C_2 > 0$ will only induce exponentially small additional transitions, so that one can take any neighborhood of size $R\sqrt{\varepsilon}$ of the set S .

- As we have already observed, by the geometric properties of canonical transforms, the *Hamiltonian trajectories* of our system are preserved by κ and one is able to identify each branch of the trajectories: the trajectories for the plus mode are Hamiltonian trajectories of $-\sigma - \sqrt{s^2 + z_1^2}$ and the trajectories for the minus mode are those of the Hamiltonian $-\sigma + \sqrt{s^2 + z_1^2}$. We denote by $\tilde{\Phi}_\pm$ these trajectories and we observe that they write

$$(2.21) \quad \tilde{\Phi}_\pm^\mathfrak{N}(s, z, \sigma, \zeta) = \left(s - \mathfrak{N}, z, \tilde{\sigma}_\pm^\mathfrak{N}(s, z_1, \sigma), (\tilde{\zeta}_1)_\pm^\mathfrak{N}(s, z_1, \sigma), \zeta'\right),$$

where we set $\zeta = (\zeta_1, \zeta')$ and $\tilde{\sigma}_\pm^\mathfrak{N}(s, z_1, \sigma) = \sigma \mp \sqrt{(s - \mathfrak{N})^2 + z_1^2 \pm \sqrt{s^2 + z_1^2}}$ by the conservation of the energy and $(\tilde{\zeta}_1)_\pm^\mathfrak{N}(s, z_1, \sigma) = \zeta_1 + \mathcal{O}(\mathfrak{N})$.

- The *transitions* occur when the gap is minimal along the trajectories, that is when $s = 0$. Besides, when the transitions occur, one has $\xi \cdot \nabla V(x) = 0$, which implies $|\xi \wedge \nabla V(x)| = |\xi| |\nabla V(x)|$ and the relation (2.19) gives

$$T_\varepsilon(x, \xi) = \exp\left(-\frac{\pi}{\varepsilon} \frac{|\xi \wedge \nabla V(x)|^2}{|\nabla V(x)|^3}\right) = T_{LZ}\left(\frac{z_1}{\sqrt{\varepsilon}}\right) (1 + \mathcal{O}(R^3 \sqrt{\varepsilon}))$$

where $T_{LZ}(\eta) = e^{-\pi\eta^2}$. Since $T_\varepsilon(x, \xi)$ and $T_{LZ}(z_1/\sqrt{\varepsilon})$ differ of a term of order η_ε , we define the flow $\tilde{\mathcal{L}}_{\varepsilon, R}$ with the transition rate $T_{LZ}(z_1/\sqrt{\varepsilon})$.

- Finally, we observe that the *drift* is made in the direction of $H_+ - H_-$. By the description above, $d\kappa$ sends $H_+ - H_-$ on a vector collinear to ∂_σ . As a consequence, we deduce that there exists a map $\sigma \mapsto \delta\sigma$ such that

$$(2.22) \quad \tilde{J}_\pm := \kappa^{-1} \circ J_\pm \circ \kappa(0, z, \sigma, \zeta) = (0, z, \sigma + \delta\sigma, \zeta).$$

Using (2.18), we deduce $\delta\sigma = \mp 2|z_1|$.

Let us now reformulate our problem in these new variables. Recall that we work in the region $\tilde{\mathcal{U}}_{\varepsilon, R}$. Let $b^\pm(s, z, \zeta)$ be two smooth functions compactly supported in $\{s > 0\}$ and such that the trajectories reaching their support have only experienced one transition during an interval of time of length \aleph . We also suppose that the functions $b^\pm(s + \aleph, z, \zeta)$ are supported in $\{s < 0\}$. We consider the symbol $c_{\varepsilon, R}^{out} = (c_{\varepsilon, R}^{+, out}, c_{\varepsilon, R}^{-, out})$ defined by

$$c_{\varepsilon, R}^{\pm, out}(s, z, \sigma, \zeta) = b^\pm(s, z, \zeta) \theta\left(\frac{\tilde{\lambda}^\pm(s, z_1, \sigma)}{R\sqrt{\varepsilon}}\right)$$

where θ is the cut-off function of Lemma 2.12 and $\tilde{\lambda}^\pm(s, z_1, \sigma)$ is the energy

$$\tilde{\lambda}^\pm(s, z_1, \sigma) = -\sigma \mp \sqrt{z_1^2 + s^2}.$$

Note that the localization in energy and the fact that we work in the hopping region yield that $\sigma = \mathcal{O}(R\sqrt{\varepsilon})$ in the zone of interest; for this reason we do not need to assume that b^\pm depends on the variable σ .

We now want to compute $\tilde{\mathcal{L}}_{\varepsilon, R}^\aleph c_{\varepsilon, R}^{out}$ the pull back by the semi-group $\tilde{\mathcal{L}}_{\varepsilon, R}$ in the normal coordinates. The observable $c_{\varepsilon, R}^{out}$ has two parts $c_{\varepsilon, R}^{+, out}$ and $c_{\varepsilon, R}^{-, out}$ and we have to consider the random trajectories that reach the support of each of these functions. More precisely, for $c_{\varepsilon, R}^{+, out}$, we consider the plus trajectories that reach its support; however, these trajectories may have known a jump and either they result from plus trajectories, either they result for minus trajectory. Similar description holds for trajectories reaching the support of $c_{\varepsilon, R}^{-, out}$.

At that point of the analysis, we notice that by point (3) of Assumptions, in the ingoing region, one of the mode is negligible. Without loss of generality, we can assume that the contribution of trajectories which arise from the minus mode is negligible. For summarizing, the picture is the following :

- for calculating the backward image by the semigroup of $c_{\varepsilon, R}^{+, out}$, we only need to consider the plus trajectories which reach its support,
- for calculating the backward image by the semigroup of $c_{\varepsilon, R}^{-, out}$, we need to consider the minus trajectories which reach its support and these trajectories arises from plus trajectories which have had a jump.

We denote by $\tilde{\Phi}_\pm^s$ the Hamiltonian trajectories associated with $\tilde{\lambda}^\pm$ and we observe that along a trajectory, the variable ζ_1 is constant up to $\mathcal{O}(R\sqrt{\varepsilon})$ term, since $\dot{\zeta}_1 = \mathcal{O}(1)$ with observation time of order $R\sqrt{\varepsilon}$.

Let us now calculate $c_{\varepsilon, R}^{+, out}$. We observe that for $(s, z, \sigma, \zeta = (\zeta_1, \zeta')) \in \text{supp } c_{\varepsilon, R}^{+, out}$, we have

$$(s, z, \sigma, \zeta_1 + \mathcal{O}(R\sqrt{\varepsilon}), \zeta') = \tilde{\Phi}_+^\aleph(s - \aleph, z, \sigma, \zeta).$$

Since the variable z_1 is constant along the trajectory, by applying the transition rate at time 0, we have

$$c_{\varepsilon,R}^{+,in}(s - \aleph, z, \sigma, \zeta) = \left(1 - T_{LZ} \left(\frac{z_1}{R\sqrt{\varepsilon}}\right)\right) c_{\varepsilon,R}^{+,out} \circ \tilde{\Phi}_+^{\aleph}(s - \aleph, z, \sigma, \zeta).$$

We deduce

$$c_{\varepsilon,R}^{+,in}(s - \aleph, z, \sigma, \zeta) = \left(1 - T_{LZ} \left(\frac{z_1}{R\sqrt{\varepsilon}}\right)\right) b^+(s, z, \zeta_1 + \mathcal{O}(R\sqrt{\varepsilon}), \zeta') \\ \times \theta \left(\frac{\tilde{\lambda}^+ \left(\tilde{\Phi}_+^{\aleph}(s - \aleph, z, \sigma, \zeta) \right)}{R\sqrt{\varepsilon}} \right).$$

Using moreover the conservation of the energy $\tilde{\lambda}^+$ along trajectories, we obtain (2.23)

$$c_{\varepsilon,R}^{+,in}(s, z, \sigma, \zeta) = \left(1 - T_{LZ} \left(\frac{z_1}{\sqrt{\varepsilon}}\right)\right) b^+(s + \aleph, z, \zeta) \theta \left(\frac{\tilde{\lambda}^+(s, z_1, \sigma)}{R\sqrt{\varepsilon}} \right) + \mathcal{O}(R\sqrt{\varepsilon}).$$

The component $c_{\varepsilon,R}^{-,in}$ is more intricate since it incorporates classical transport through both modes, application of the transfert coefficient and of the drift. Indeed, the branches of minus trajectories which reach the support of $c_{\varepsilon,R}^{-,out}$ results from plus trajectories that have been drifted : more precisely, for $(s, z, \sigma, \zeta) \in \text{supp } c_{\varepsilon,R}^{-,out}$,

$$(s, z, \sigma, \zeta_1 + \mathcal{O}(R\sqrt{\varepsilon}), \zeta') = \tilde{\Phi}_-^s \circ \tilde{J}_- \circ \tilde{\Phi}_+^{-s+\aleph}(s - \aleph, z, \sigma, \zeta).$$

By applying the transition rate at time 0, we obtain

$$c_{\varepsilon,R}^{-,in}(s - \aleph, z, \sigma, \zeta) = T_{LZ} \left(\frac{z_1}{R\sqrt{\varepsilon}}\right) c_{\varepsilon,R}^{-,out} \circ \tilde{\Phi}_-^s \circ \tilde{J}_- \circ \tilde{\Phi}_+^{-s+\aleph}(s - \aleph, z, \sigma, \zeta).$$

We deduce

$$c_{\varepsilon,R}^{-,in}(s - \aleph, z, \sigma, \zeta) = \left(1 - T_{LZ} \left(\frac{z_1}{R\sqrt{\varepsilon}}\right)\right) b^-(s, z, \zeta_1 + \mathcal{O}(R\sqrt{\varepsilon}), \zeta') \\ \times \theta \left(\frac{\tilde{\lambda}^- \left(\tilde{\Phi}_-^s \circ \tilde{J}_- \circ \tilde{\Phi}_+^{-s+\aleph}(s - \aleph, z, \sigma, \zeta) \right)}{R\sqrt{\varepsilon}} \right).$$

The crucial point is that (2.22) implies that for all $(z, \sigma, \zeta) \in \mathbf{R}^5$,

$$\tilde{\lambda}^-(\tilde{J}_-(0, z, \sigma, \zeta)) = \tilde{\lambda}^+(0, z_1, \sigma),$$

so, by using also the conservation of the energy along trajectories, we obtain

$$\tilde{\lambda}^- \left(\tilde{\Phi}_-^s \circ \tilde{J}_- \circ \tilde{\Phi}_+^{-s+\aleph}(s - \aleph, z, \sigma, \zeta) \right) = \tilde{\lambda}^+(s - \aleph, z_1, \sigma).$$

As a consequence,

$$c_{\varepsilon,R}^{-,in}(s - \aleph, z, \sigma, \zeta) = \left(1 - T_{LZ} \left(\frac{z_1}{R\sqrt{\varepsilon}}\right)\right) b^-(s, z, \zeta_1, \zeta') \\ \times \theta \left(\frac{\tilde{\lambda}^+(s - \aleph, z, \sigma, \zeta)}{R\sqrt{\varepsilon}} \right) + \mathcal{O}(R\sqrt{\varepsilon}),$$

whence

$$(2.24) \quad c_{\varepsilon,R}^{-,in}(s, z, \sigma, \zeta) = T_{LZ} \left(\frac{z_1}{\sqrt{\varepsilon}} \right) b^-(s + \aleph, z, \zeta) \theta \left(\frac{\tilde{\lambda}^+(s, z_1, \sigma)}{R\sqrt{\varepsilon}} \right) + \mathcal{O}(R\sqrt{\varepsilon}).$$

As a conclusion, in order to prove our result, we only need to prove the following relations:

$$(2.25) \quad \left(\text{op}_{\varepsilon}(c_{\varepsilon,R}^{\pm,out})v^{\varepsilon}, v^{\varepsilon} \right)_{L^2(\mathbf{R}^3_{s,z})} = \left(\text{op}_{\varepsilon}(c_{\varepsilon,R}^{\pm,in})v^{\varepsilon}, v^{\varepsilon} \right)_{L^2(\mathbf{R}^3_{s,z})} + \mathcal{O}(\eta_{\varepsilon}).$$

2.3.5. *The transitions.* The claim (2.25) is proved by use of the following Landau-Zener type formula (see [27, 33, 14, 17]).

Proposition 2.15. *Let v^{ε} be a solution of (2.20). There exist two vector-valued functions $k^{\varepsilon,\pm} \in L^2(\mathbf{R}^d, \mathbf{C}^2)$ such that $k^{\varepsilon,\pm} \mathbf{1}_{|z_1| \leq R\sqrt{\varepsilon}}$ are bounded and such that we have for $\pm s > 0$ and $|z_1| \leq R\sqrt{\varepsilon}$*

$$\begin{aligned} v_1^{\varepsilon}(z, s) &= e^{is^2/(2\varepsilon)} \left| \frac{s}{\sqrt{\varepsilon}} \right|^{i\frac{z_1^2}{2\varepsilon}} k_1^{\varepsilon,\pm}(z) + \mathcal{O} \left(\frac{R^2\sqrt{\varepsilon}}{s} \right), \\ v_2^{\varepsilon}(z, s) &= e^{-is^2/(2\varepsilon)} \left| \frac{s}{\sqrt{\varepsilon}} \right|^{-i\frac{z_1^2}{2\varepsilon}} k_2^{\varepsilon,\pm}(z) + \mathcal{O} \left(\frac{R^2\sqrt{\varepsilon}}{s} \right) \end{aligned}$$

Moreover, $k^{\varepsilon,+} = S \left(\frac{z_1}{\sqrt{\varepsilon}} \right) k^{\varepsilon,-}$ where the unitary matrix S is given by

$$\begin{aligned} S(\lambda) &= \begin{pmatrix} a(\lambda^2) & -\lambda \bar{b}(\lambda^2) \\ \lambda b(\lambda^2) & a(\lambda^2) \end{pmatrix}, \\ a(\lambda) &= e^{-\pi\lambda/2}, \quad b(\lambda) = \frac{2ie^{i\pi/4}}{\lambda\sqrt{\pi}} 2^{-i\lambda/2} e^{-\pi\lambda/4} \Gamma(1 + i\frac{\lambda}{2}) \sinh(\frac{\pi\lambda}{2}). \end{aligned}$$

We are now in position to conclude the proof of Proposition 1.7 by proving (2.25). We first take advantage of the localization near the energy surfaces to translate it as a focalization: by the analogous of Lemma 2.12 and Remark 2.14 in (s, z) variables, we obtain

$$\begin{aligned} \left(\text{op}_{\varepsilon}(c_{\varepsilon,R}^{\pm,out})v^{\varepsilon}, v^{\varepsilon} \right)_{L^2(\mathbf{R}^3_{s,z})} &= \left(\text{op}_{\varepsilon}(c_{\varepsilon,R}^{\pm,out} \tilde{\Pi}^{\pm}(s, z_1))v^{\varepsilon}, v^{\varepsilon} \right)_{L^2(\mathbf{R}^3_{s,z})} + \mathcal{O}(\eta_{\varepsilon}) \\ &= \left(\text{op}_{\varepsilon}(c_{\varepsilon,R}^{\pm,out} \tilde{\Pi}^{\pm}(s, z_1))v^{\varepsilon}, \tilde{\Pi}^{\pm}(s, z_1)v^{\varepsilon} \right)_{L^2(\mathbf{R}^3_{s,z})} + \mathcal{O}(\eta_{\varepsilon}) \end{aligned}$$

where we denote by $\tilde{\Pi}^{\pm}(s, z_1)$ the eigenprojectors of the matrix

$$\tilde{A}(s, z_1) = \begin{pmatrix} s & z_1 \\ z_1 & -s \end{pmatrix}$$

associated with the eigenvalues $\mp\sqrt{s^2 + z_1^2}$. For $|z_1| \leq R\sqrt{\varepsilon}$, we have

$$(2.26) \quad \begin{aligned} \tilde{\Pi}^+(s, z_1) &= \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} + \mathcal{O}(R\sqrt{\varepsilon}) \text{ in } \{s > 0\}, \\ \tilde{\Pi}^+(s, z_1) &= \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + \mathcal{O}(R\sqrt{\varepsilon}) \text{ in } \{s < 0\} \end{aligned}$$

(whence similar asymptotics for $\tilde{\Pi}^-$ since $\text{Id} = \tilde{\Pi}^+ + \tilde{\Pi}^-$). Therefore,

$$\begin{aligned} \left(\text{op}_\varepsilon(c_{\varepsilon,R}^{+,out})v^\varepsilon, v^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} &= \left(\text{op}_\varepsilon(b_{\varepsilon,R}^{+,out})v_2^\varepsilon, v_2^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} + \mathcal{O}(\eta_\varepsilon), \\ \left(\text{op}_\varepsilon(c_{\varepsilon,R}^{-,out})v^\varepsilon, v^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} &= \left(\text{op}_\varepsilon(b_{\varepsilon,R}^{-,out})v_1^\varepsilon, v_1^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} + \mathcal{O}(\eta_\varepsilon) \end{aligned}$$

with $b_{\varepsilon,R}^{\pm,out} = b^\pm(s, z, \zeta)$.

The symbols $c_{\varepsilon,R}^{\pm,in}$ are supported in the region $\{s < 0\}$, i.e. before the transitions. By (3) of Assumptions 1.2, we know that the mode plus only has mass of order 1. By (2.26), in $\{s < 0\}$, the mode plus corresponds to the component v_1^ε and we deduce $v_2^\varepsilon = \mathcal{O}(\eta_\varepsilon)$. Therefore, we have

$$\begin{aligned} \left(\text{op}_\varepsilon(c_{\varepsilon,R}^{+,in})v^\varepsilon, v^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} &= \left(\text{op}_\varepsilon \left(\left(1 - T_{LZ} \left(\frac{z_1}{\sqrt{\varepsilon}} \right) \right) b_{\varepsilon,R}^{+,in} \right) v_1^\varepsilon, v_1^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} + \mathcal{O}(\eta_\varepsilon), \\ \left(\text{op}_\varepsilon(c_{\varepsilon,R}^{-,in})v^\varepsilon, v^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} &= \left(\text{op}_\varepsilon \left(T_{LZ} \left(\frac{z_1}{\sqrt{\varepsilon}} \right) b_{\varepsilon,R}^{-,in} \right) v_1^\varepsilon, v_1^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} + \mathcal{O}(\eta_\varepsilon), \end{aligned}$$

with

$$b_{\varepsilon,R}^{+,in} = b^+(s + \aleph, z, \zeta), \quad b_{\varepsilon,R}^{-,in} = b^-(s + \aleph, z, \zeta).$$

We now use Proposition 2.15 in order to relate $v_{1,2}^\varepsilon$ with $k_{1,2}^{\varepsilon,\pm}$. For the term in function of $c_{\varepsilon,R}^{+,out}$, we have

$$\begin{aligned} \left(\text{op}_\varepsilon(c_{\varepsilon,R}^{+,out})v^\varepsilon, v^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} &= \mathcal{O}(R^2\sqrt{\varepsilon}) \\ &+ \left(\text{op}_\varepsilon \left(b_{\varepsilon,R}^{+,out} \right) e^{-i\frac{s}{2\varepsilon}} \left| \frac{s}{\sqrt{\varepsilon}} \right|^{-i\frac{z_1^2}{2\varepsilon}} k_2^{\varepsilon,+}, e^{-i\frac{s}{2\varepsilon}} \left| \frac{s}{\sqrt{\varepsilon}} \right|^{-i\frac{z_1^2}{2\varepsilon}} k_2^{\varepsilon,+} \right)_{L^2(\mathbf{R}_{s,z}^3)} \end{aligned}$$

We use the following Lemma (see Lemma 8 and Lemma 9 in [15]) in order to commute the pseudo differential operator and the phases.

Lemma 2.16. *In $\mathcal{L}(L^2(\mathbf{R}_{s,z}^3))$, we have for $|z_1| \leq R\sqrt{\varepsilon}$*

$$\left| \frac{s}{\sqrt{\varepsilon}} \right|^{\pm i\frac{z_1^2}{2\varepsilon}} \text{op}_\varepsilon(b(s, z, \zeta)) \left| \frac{s}{\sqrt{\varepsilon}} \right|^{\mp i\frac{z_1^2}{2\varepsilon}} = \text{op}_\varepsilon(b(s, z, \zeta)) + \mathcal{O}(\sqrt{\varepsilon}|\ln \varepsilon|).$$

As a consequence, we obtain

$$\left(\text{op}_\varepsilon(c_{\varepsilon,R}^{+,out})v^\varepsilon, v^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} = \left(\text{op}_\varepsilon \left(b_{\varepsilon,R}^{+,out} \right) k_2^{\varepsilon,+}, k_2^{\varepsilon,+} \right)_{L^2(\mathbf{R}_{s,z}^3)} + \mathcal{O}(\eta_\varepsilon).$$

Using the change of variable $s \mapsto s + \aleph$, we find

$$\left(\text{op}_\varepsilon(c_{\varepsilon,R}^{+,out})v^\varepsilon, v^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} = \left(\text{op}_\varepsilon \left(b_{\varepsilon,R}^{+,in} \right) k_2^{\varepsilon,+}, k_2^{\varepsilon,+} \right)_{L^2(\mathbf{R}_{s,z}^3)} + \mathcal{O}(\eta_\varepsilon).$$

By Proposition 2.15, we have

$$k_2^{\varepsilon,+} = \frac{z_1}{\sqrt{\varepsilon}} b \left(\frac{z_1^2}{\varepsilon} \right) k_1^{\varepsilon,-} + a \left(\frac{z_1^2}{\varepsilon} \right) k_2^{\varepsilon,-}$$

Besides, $b^\pm(s + \aleph)$ is supported in $\{s < 0\}$ and, as we said before, by (3) of Assumptions 1.2, near $\{\sigma + s = 0, s < 0\}$ (i.e. near minus trajectories entering in the hopping zone), $v_2^\varepsilon(z, s) = \mathcal{O}(\eta_\varepsilon)$. Therefore, because of the link between v_2^ε

and $k_2^{\varepsilon,-}$ in the region $\{s < 0\}$, $k_2^{\varepsilon,-}(z) = \mathcal{O}(\eta_\varepsilon)$ and, using the relation $\lambda^2|b(\lambda)|^2 = 1 - a(\lambda)^2 = 1 - e^{-\pi\lambda}$, we obtain

$$\frac{z_1}{\sqrt{\varepsilon}} \bar{b} \left(\frac{z_1^2}{\varepsilon} \right) \text{op}_\varepsilon \left(b_{\varepsilon,R}^{+,in} \right) \frac{z_1}{\sqrt{\varepsilon}} \bar{b} \left(\frac{z_1^2}{\varepsilon} \right) = \text{op}_\varepsilon \left(\left((1 - T_{LZ} \left(\frac{z_1}{\sqrt{\varepsilon}} \right)) b_{\varepsilon,R}^{+,in} \right) \right) + \mathcal{O}(\sqrt{\varepsilon})$$

in $\mathcal{L}(L^2(\mathbf{R}_{s,z}^3))$. Finally, in view of the relations satisfied by $c_{\varepsilon,R}^{\pm,in}$, and arguing as before, we can conclude that

$$\begin{aligned} \left(\text{op}_\varepsilon (c_{\varepsilon,R}^{+,out}) v^\varepsilon, v^\varepsilon \right)_{L^2(\mathbf{R}_{s,z}^3)} &= \left(\text{op}_\varepsilon \left(\left(1 - T_{LZ} \left(\frac{z_1}{\sqrt{\varepsilon}} \right) \right) b_{\varepsilon,R}^{+,in} \right) k_1^{\varepsilon,-}, k_1^{\varepsilon,-} \right)_{L^2(\mathbf{R}_{s,z}^3)} \\ &\quad + \mathcal{O}(\eta_\varepsilon) \\ &= \left(\text{op}_\varepsilon (c_{\varepsilon,R}^{+,in}) v^\varepsilon, v^\varepsilon \right)_{L^2(\mathbf{R}^{d+1})} + \mathcal{O}(\eta_\varepsilon). \end{aligned}$$

Arguing similarly for $c_{\varepsilon,R}^{-,out}$, we obtain (2.25). \square

3. NUMERICAL EXPERIMENTS

3.1. Algorithmic realization of the semi-group.

The algorithm. The mechanism describing the evolution of $f_{\pm,app}^\varepsilon(t, x, \xi)$ has a simple algorithmic description as proposed in [18] :

- (1) Far from Σ , $f_+^\varepsilon(t, x, \xi)$ propagates along the trajectories Φ_+^t and $f_-^\varepsilon(t, x, \xi)$ propagates along the trajectories Φ_-^t .
- (2) Whenever a trajectory reaches Σ at time t^* in a point (x^*, ξ^*) , one may transmit some energy to the other mode according to a random process. One takes a random number r between 0 and 1 and one compares r and the transfer coefficient $T_\varepsilon(x^*, \xi^*)$:
 - If $r > T_\varepsilon(x^*, \xi^*)$, one continues with the same trajectory and propagate the mass $f_+^\varepsilon(t^*, x^*, \xi^*)$ on the trajectory $\Phi_+^t(x, \xi)$.
 - If $r < T_\varepsilon(x^*, \xi^*)$, one initiates a trajectory Φ_-^t from the point

$$(3.1) \quad (x_-^{t^*}, \xi_-^{t^*}) := J_+(x^*, \xi^*) = \left(x^* + 2|\xi^*| \frac{\nabla V(x^*)}{|\nabla V(x^*)|^2}, \xi^* \right)$$

and propagate the mass $f_-^\varepsilon(t^*, x^*, \xi^*) = f_+^\varepsilon(t^*, x^*, \xi^*)$ on the new trajectory $\Phi_-^{t-t^*}(x_-^{t^*}, \xi_-^{t^*})$.

A similar process is performed on the other mode.

Remark 3.1. Note that the hypothesis (3) of Assumptions 1.2 imply that at a transition point, only one of the trajectory is weighted.

3.2. The simulated models. In this section, we present various numerical experiments in order to validate our kinetic model (1.8). We assume that the initial data has a Wigner transform supported in the domain $\Omega \subset \{|x_1| \leq a, |\xi| \leq M\}$. In practice, we will take $a = M = 10$. We will not discuss the behavior close to the boundary and we will consider a time schedule $[0, T]$ such that the trajectories issued from Ω do not reach any boundary. For the quantum model, we will use periodic boundary conditions. In the case of the barrier potential, we will also assume that the trajectories issued from points of Ω do not reach the top of the

barrier with velocity 0; thus, we are in the frame of Assumption 1.2 and the above description of the trajectories is valid.

Several sources of error can be identified in this model, if we compare it to the original quantum equation (1.1):

- the error made in the computation of the initial data and on its sampling by a bunch of particles,
- the error made during the plus/minus transition processes; on the computation of the transmission coefficients and the position of particles after transitions (presence or not of the jump process),
- the error made during the transport phase, when the quantum transport is replaced by the classical transport induced by the Hamiltonian (1.5) and when the coherence effects between particles are neglected,
- the error made when we neglect the transport of the antidiagonal part of the Wigner function.

In the numerical tests that we present, we concentrate our study on these modeling errors, instead of numerical errors due to time and space discretization. Hence, we have taken sufficiently small time steps and space steps, such that the error associated to these numerical parameters is negligible compared to the modeling errors.

In order to characterize these different modeling errors, let us identify the models that we simulate:

- (i) The *quantum graphene model* is the original quantum equation (1.1), computed with the Strang splitting method. In the first series of experiments (Subsections 3.3, 3.4 and 3.5), the initial data is a gaussian coherent quantum state localized at the position x_1^0 in the x_1 direction, multiplied by a plane wave in the x_2 direction, with the momentum $\xi^0 = (\xi_1^0, \xi_2^0)$, and polarized on the plus mode:

$$(3.2) \quad \rho_0^\varepsilon(x, y) = \psi_0^\varepsilon(x_1 - x_1^0, \xi_1^0) (\psi_0^\varepsilon(y_1 - x_1^0, \xi_1^0))^* \exp\left(i \frac{(x_2 - y_2)\xi_2^0}{\varepsilon}\right),$$

with

$$(3.3) \quad \psi_0^\varepsilon(x_1, \xi_1^0) = \Pi_+(\varepsilon D) \begin{pmatrix} \sqrt{2} u_0^\varepsilon(x_1, \xi_1^0) \\ 0 \end{pmatrix},$$

and

$$(3.4) \quad u_0^\varepsilon(x_1, \xi_1^0) = (\pi\varepsilon)^{-1/4} \exp\left(-\frac{x_1^2}{2\varepsilon} + i \frac{x_1 \xi_1^0}{\varepsilon}\right).$$

In the last experiments (Subsection 3.6), the initial data is a mixture of coherent quantum states

$$(3.5) \quad \rho_0^\varepsilon(x, y) = \iint_{\mathbf{R}^2} \psi_0^\varepsilon(x_1 - x_1^0, \xi_1^0) (\psi_0^\varepsilon(y_1 - x_1^0, \xi_1^0))^* \exp\left(i \frac{(x_2 - y_2)\xi_2^0}{\varepsilon}\right) \times f^0(x_1^0, \xi_1^0) dx_1^0 d\xi_1^0,$$

where f^0 is a given distribution density and where, for simplicity, all the states are taken with the same momentum ξ_2^0 in the x_2 direction.

- (ii) The *kinetic graphene model* is (1.8), discretized according to the particle algorithm described in Section 3.1, with randomly computed transitions between plus and minus modes, with or without jumps (3.1). The initial

data $f_{+,0}^\varepsilon$ and $f_{-,0}^\varepsilon$ are described in Subsection 3.4. The time integrator is the triple jump method of order 4 [25].

- (iii) The *quantum pseudo-graphene model* is the following modified quantum equation, computed with the Strang splitting method:

$$(3.6) \quad i\varepsilon\partial_t\varrho^\varepsilon = [(\tilde{A}(\varepsilon D) + V), \varrho^\varepsilon], \quad \text{with } \tilde{A}(\xi) = \begin{pmatrix} |\xi| & 0 \\ 0 & -|\xi| \end{pmatrix}$$

with the same initial data ρ_0^ε as (i). Far from the crossing set $\{\xi = 0\}$, this model displays the same dynamics as (i), but the major difference is that no transition occur between plus and minus modes with (3.6).

- (iv) The *kinetic pseudo-graphene model* is the same as (ii), without transition process. It is the classical counterpart of the quantum pseudo-graphene model (iv).

3.3. The Klein effect. Our aim in this subsection is to observe qualitatively the Klein effect (the tunneling of particles through a classically forbidden potential barrier) with both models: the quantum graphene model and the kinetic graphene model. For the simulations presented in this subsection, we have taken a sufficiently large number of particles (4×10^6) and assume that we are at numerical convergence. The potential is the following smooth barrier potential, depending only on x_1 :

$$V(x_1) = 4 \sin^3\left(\frac{\pi}{4}(x_1 + 1)\right) \text{ for } x_1 \in [-1, 3], \text{ and } V(x_1) = 0 \text{ otherwise.}$$

Let us describe the phenomenology that can be observed on Figures 3 and 4. Here, $\varepsilon = 0.064$. We represent on Figure 3 the contour plots of the trajectories of the plus and minus modes, computed with the quantum graphene model. On Figure 4, the plus and minus densities are represented at four instants, computed with the quantum graphene model (plotted in plain lines) and computed with our kinetic graphene model (plotted with the 'X'). It can be seen that the results given by both models are in very good agreement.

Initially, a coherent wavepacket in the plus mode is at position $x_1^0 = -2$ with the momentum $\xi^0 = (\xi_1^0, \xi_2^0) = (1.3, 0.1)$ (Figure 4(a)). The wavefunction first propagates freely, then enters inside the barrier and its momentum ξ_1 decreases until it vanishes. At the instant t_1 , the wavefunction is partially reflected and partially transferred through the barrier in the minus mode. Then, the minus wavepacket propagates inside the barrier (see Figure 4(b)) until it reaches the other side of the barrier (instant t_2), through which it is partially transferred into a plus wavepacket and partially reflected. Finally, the remaining minus wavepacket propagates from the right to the left (see Figure 4(c)) and hits again the barrier (instant t_3), where it is almost integrally transferred into a third plus wavepacket outside the barrier.

At the end of the simulation (Figure 4(d)), almost all the mass have been redistributed into three plus wavepackets. In Table 3.3, we give the numerical transfer rates (i.e. the ratio of mass in each wavepacket over the initial mass) of the initial mass into wavepackets 1, 2 and 3 for different values of ε , computed with the quantum model and with the kinetic model. Here again, we observe a good agreement between our kinetic model and the reference one.

3.4. Choice of the initial data. Let us concentrate on the computation of the initial data for the kinetic model. Recall that, in the pure-state case, the initial data for the quantum graphene model is the density matrix given by (3.2), where the wave function (3.3) is the gaussian wavefunction (3.4) projected on the plus

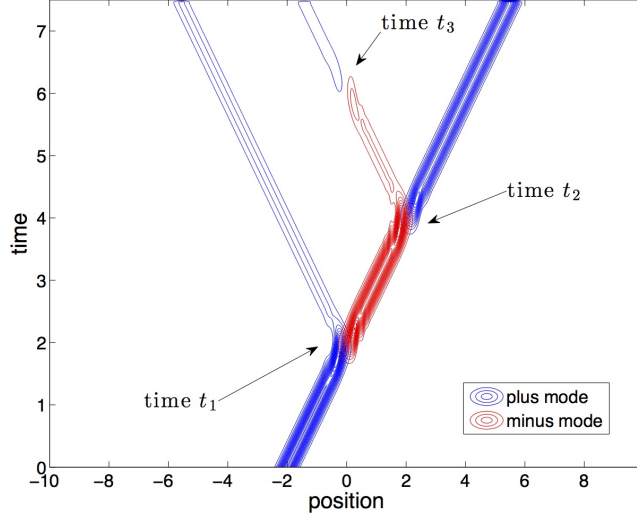


Figure 3: Contour plots of the plus and minus densities, quantum graphene model with $\varepsilon = 0.064$.

Table 1: Transfer rates

ε	0.128	0.064	0.032	0.016	0.008	0.004
wavepacket 1 (quantum)	7.42%	14.18%	26.25%	45.55%	70.31%	91.17%
wavepacket 1 (kinetic)	7.64%	14.64%	27.03%	46.73%	71.47%	91.92%
wavepacket 2 (quantum)	85.71%	73.65%	54.39%	29.66%	8.82%	0.78%
wavepacket 2 (kinetic)	85.69%	73.49%	54.17%	29.46%	8.79%	0.75%
wavepacket 3 (quantum)	6.38%	10.46%	14.28%	13.50%	6.20%	0.71%
wavepacket 3 (kinetic)	6.20%	10.2%	13.98%	13.15%	6.01%	0.68%

mode. Let us denote by $f_0^\varepsilon(x, \xi)$ the Wigner function of ϱ_0^ε and by $\widetilde{f}_0^\varepsilon(x, \xi)$ the Wigner function of the gaussian wavepacket

$$\widetilde{\rho}_0^\varepsilon(x, y) = u_0^\varepsilon(x_1 - x_1^0, \xi_1^0) (u_0^\varepsilon(y_1 - x_1^0, \xi_1^0))^* \exp\left(i \frac{(x_2 - y_2)\xi_2^0}{\varepsilon}\right),$$

which is

$$\begin{aligned} \widetilde{f}_0^\varepsilon(x, \xi) &= \frac{1}{2\pi} \int e^{i\xi\eta} u_0^\varepsilon\left(t, x_1 - \varepsilon \frac{\eta}{2}\right) \overline{u_0^\varepsilon}\left(x_1 + \varepsilon \frac{\eta}{2}\right) d\eta \delta_{\xi_2 = \xi_2^0} \\ (3.7) \quad &= \frac{1}{\pi\varepsilon} \exp\left(-\frac{(x_1 - x_1^0)^2}{\varepsilon} - \frac{(\xi_1 - \xi_1^0)^2}{\varepsilon}\right) \delta_{\xi_2 = \xi_2^0}. \end{aligned}$$

The function $\text{tr}(\Pi^+(\xi) f_0^\varepsilon(x, \xi) \Pi^+(\xi))$ is asymptotically close to $\widetilde{f}_0^\varepsilon(x, \xi)$ as $\varepsilon \rightarrow 0$. More precisely, using pseudo-differential calculus, the following expansion can be

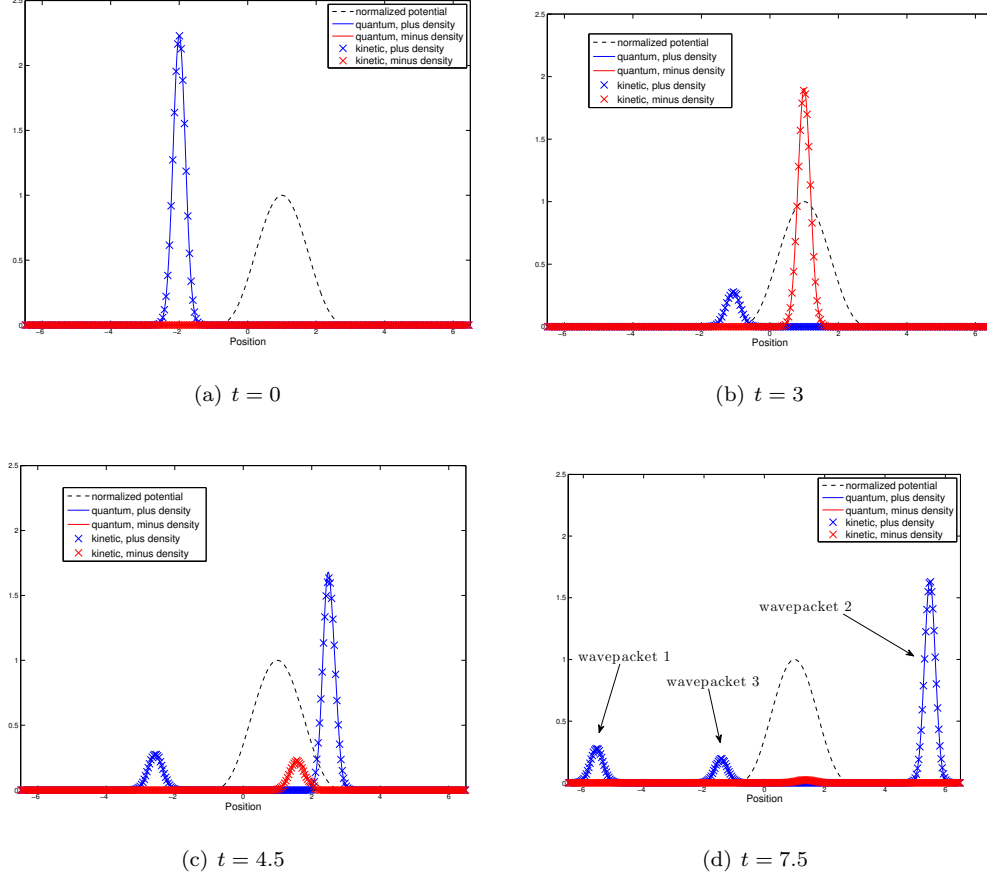


Figure 4: Propagation of a coherent wavepacket through a barrier for $\varepsilon = 0.064$.

obtained:

$$(3.8) \quad \text{tr}(\Pi^+(\xi) f_0^\varepsilon(x, \xi) \Pi^+(\xi)) = \widetilde{f}_0^\varepsilon(x, \xi) - \frac{\varepsilon}{2} \frac{\xi_2}{|\xi|^2} \partial_x \widetilde{f}_0^\varepsilon(x, \xi) + \mathcal{O}(\varepsilon)$$

$$(3.9) \quad = \widetilde{f}_0^\varepsilon \left(x_1 - \frac{\varepsilon}{2} \frac{\xi_2}{|\xi|^2}, \xi \right) + \mathcal{O}(\varepsilon)$$

in $L^1(\mathbf{R}^4)$. Notice that this expansion is in powers of $\varepsilon^{1/2}$ (indeed, the L^1 norm of $\partial_x \widetilde{f}_0^\varepsilon(x, \xi)$ is of order $\varepsilon^{-1/2}$).

As initial data for our kinetic graphene model, in the pure-state case, let us experiment these two approximations of f_0^ε : the first one is simply $\widetilde{f}_0^\varepsilon(x, \xi)$ and the second one is the shifted function

$$(3.10) \quad f_{+,0}^\varepsilon(x, \xi) = \widetilde{f}_0^\varepsilon \left(x_1 - \frac{\varepsilon}{2} \frac{\xi_2}{|\xi|^2}, \xi \right).$$

The advantage of this last choice compared to the right-hand side of (3.8) is that the distribution function is always positive. On Figure 5, we plot in logarithmic

scale the errors on the densities, i.e. the quantities

$$\text{err}_0 = \int \left| |\psi_0^\varepsilon(x)|^2 - \int \widetilde{f}_0^\varepsilon(x, \xi) d\xi \right| dx_1, \quad \text{err}_1 = \int \left| |\psi_0^\varepsilon(x)|^2 - \int f_{+,0}^\varepsilon(x, \xi) d\xi \right| dx_1,$$

for $\varepsilon = 2^N \times 2.5 \times 10^{-3}$, with $N \in \{0, 1, 2, 3, 4, 5, 6, 7, 8, 9\}$. Again, these integrals have been computed with enough discretization points such that the numerical integration errors is negligible: the (x, ξ) domain $[-10, 10]^2$ is discretized with 2^{18} grid points in the x direction and 2^{11} grid points in the ξ direction with a non uniform cartesian mesh refined near the point (x^0, ξ_1^0) . As above, we have taken $(x_1^0, \xi_1^0, \xi_2^0) = (-2, 1.3, 0.1)$. Figure 5 confirms the estimate (3.8) and (3.9): the function $\widetilde{f}_0^\varepsilon$ is an $\mathcal{O}(\varepsilon^{1/2})$ approximation of f_0^ε and the shifted function $f_{+,0}^\varepsilon(x, \xi)$ is an order $\mathcal{O}(\varepsilon)$ approximation. From now on, we choose this shifted function $f_{+,0}^\varepsilon$ defined by (3.10) (and $f_{-,0}^\varepsilon = 0$) as initial data for the kinetic graphene model.

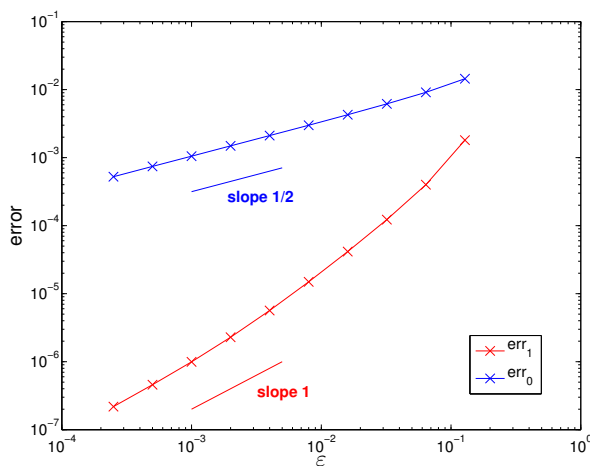


Figure 5: L^1 error on the initial density, for the two choices $\widetilde{f}_0^\varepsilon$ and $f_{+,0}^\varepsilon$.

Let us now discuss briefly the initial sampling step for the distribution function $f_{+,0}^\varepsilon$ and compare the rates of convergence of Monte Carlo and quasi-Monte Carlo sampling. Monte Carlo sampling can be achieved by a rescaled and shifted sampling of a bidimensional Gaussian distribution, taking advantage of the tensorial structure of the function. Its convergence rate is known to be of order $\mathcal{O}(1/\sqrt{N_{part}})$, where N_{part} denotes the number of particles. Quasi-Monte Carlo methods use quasi-random sequences, also known as low-discrepancy, which are deterministic approximation of the uniform distribution on $[0, 1]^2$ and can be transformed into a Gaussian distribution by the cumulative distribution function. Such methods display better convergence rate, of the form $C(\log N_{part})^k/N_{part}$. In both cases, densities are reconstructed by a fifth order spline interpolation method.

On Figure 6, we have represented with the same scales, for the quasi-Monte Carlo method (using a 2D Hammersley set) and the Monte Carlo method, the L^1 error between the reconstructed density and the reference density $|\psi_0^\varepsilon|^2$, for four different values of ε . For the quasi-Monte Carlo method, we observe a convergence

rate which is compatible with $\mathcal{O}(C(\log N_{part})^2/N_{part})$ (with a saturation due to the difference between $f_{+,0}^\varepsilon$ and f_0^ε , studied above). For the Monte Carlo method, we observe a slower convergence, of the form $\mathcal{O}(1/\sqrt{N})$. Therefore, in the sequel we systematically use the quasi-Monte Carlo method, with 4×10^6 particles.

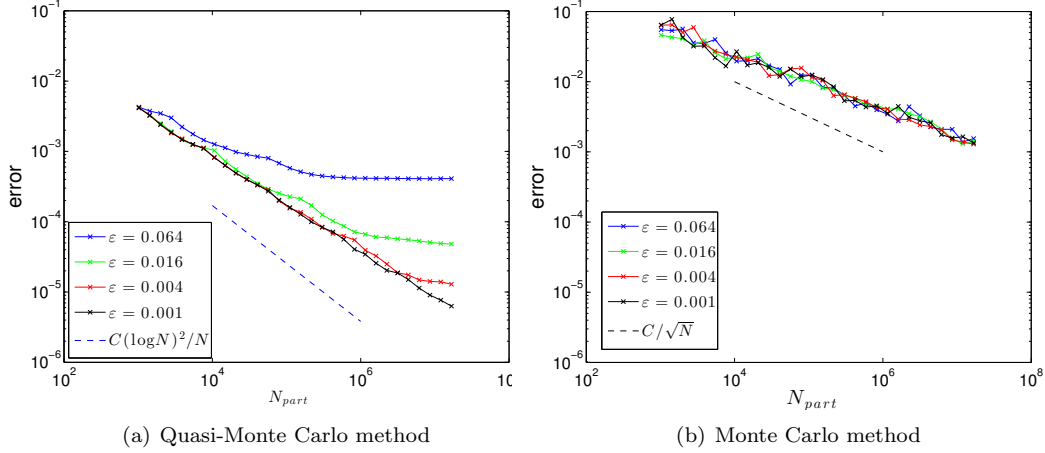


Figure 6: L^1 error on the initial density with respect to the number of particles for the quasi-Monte Carlo and the Monte Carlo methods

3.5. Transport and transition phases. In this section, we compare the dynamics computed with the quantum and the kinetic models. The initial data are the same as in the previous subsections: $\psi_0^\varepsilon(x_1 - x_1^0)$ given by (3.3) and $f_{+,0}^\varepsilon$ defined by (3.10). Here also, we take $(x_1^0, \xi_1^0, \xi_2^0) = (-2, 1.3, 0.1)$. The final time of the simulation is $t_f = 4.5$.

Transport phase. In a first step, in order to characterize the error made during the transport phase, we simulate the pseudo-graphene model, which displays the same transport properties as the graphene model but which induces no transition between plus and minus modes. We simulate the actions of two potentials:

$$(3.11) \quad V_1(x_1) = \frac{1}{20}(x_1 + 10)^2 \quad \text{and} \quad V_2(x_1) = \alpha \operatorname{atan}(2x_1 + \frac{\pi}{2}),$$

where $\alpha \approx 0.643$ has been adjusted such that, at the points where a particle is stopped by the potential barrier (resp. x_1^* for V_1 and x_2^* for V_2), we have $V_1'(x_1^*) = V_2'(x_2^*)$. On Figure 7, we plot the error

$$\begin{aligned} \text{error} = & \int \left| |\Pi^+ \psi^\varepsilon(t_f, x)|^2 - \int f_+^\varepsilon(t_f, x, \xi) d\xi \right| dx_1 \\ & + \int \left| |\Pi^- \psi^\varepsilon(t_f, x)|^2 - \int f_-^\varepsilon(t_f, x, \xi) d\xi \right| dx_1 \end{aligned}$$

as a function of ε . We observe two features. First, the error behaves as $\mathcal{O}(\varepsilon^{1/2})$ in both cases. This can be explained by the fact that the derivative in x and ξ of

the coherent wavepacket are of order $\varepsilon^{-1/2}$ in L^1 norm. Second, the error for the transport in V_2 is 5 times higher than the error for the transport in V_1 . This can be explained by the fact that V_1 is a harmonic potential, so the quantum transport operator coincides with its Wigner counterpart. Hence, for the transport by V_1 , the main source of error comes from the fact that coherence effects between particles are not taken into account. For the transport by V_2 , the main source of error comes from the replacement of the quantum transport operator by the classical one.

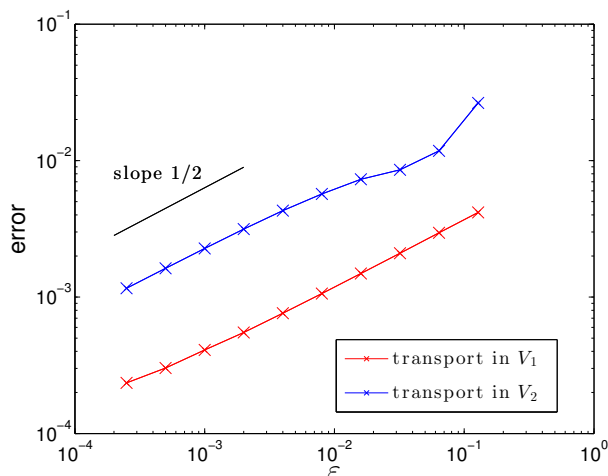


Figure 7: L^1 error on the densities for the pseudo-graphene model as a function of ε , at time $t_f = 4.5$, with the two potentials V_1 (harmonic) and V_2 (anharmonic).

Transition phase. Let us now come back to the real graphene models, with transitions. We compute the transport by the potential V_1 of a coherent wavepacket initially at $(x_1^0, \xi_1^0, \xi_2^0) = (-2, 1.3, 0.1)$. We plot on Figure 8 the plus and minus trajectories in the potential V_1 . When the plus particles hits the potential barrier at the position x_1^* and at time t^* , its momentum ξ_1 vanishes and a transition occurs: the mass is partially transferred to a minus particle, with the transfer rate T_ε . This phenomenology can be observed on Figure 8 where, for $\varepsilon = 0.064$, the contour plots of plus and minus wavepackets are represented, computed with the quantum graphene model. In order to check the formula (1.10), we give in Table 3.5 the numerical transfer rate T_ε (i.e. the ratio of the transferred mass over the initial mass) and the quantity $-\frac{\pi(\xi_2^0)^2}{\varepsilon \log T_\varepsilon}$ (theoretically equal to $V'(x_1^*)$) for 6 values of ε , computed with the quantum model and with the kinetic model. Note that x_1^* is defined by $V_1(x_1^*) + \xi_2^0 = V_1(x_1^0) + |\xi^0|$, which gives $x_1^* \approx -0.61$ and $V_1'(x_1^*) \approx 0.94$. The numerical results corroborate the predicted rates.

But the rate T_ε is not the only parameter appearing in the transition phenomenon. We highlight the importance of the jump operator J_\pm , designed in order to ensure the energy conservation during the transition process. On Figure 9, we plot the L^1 error between the densities for the quantum model and the densities

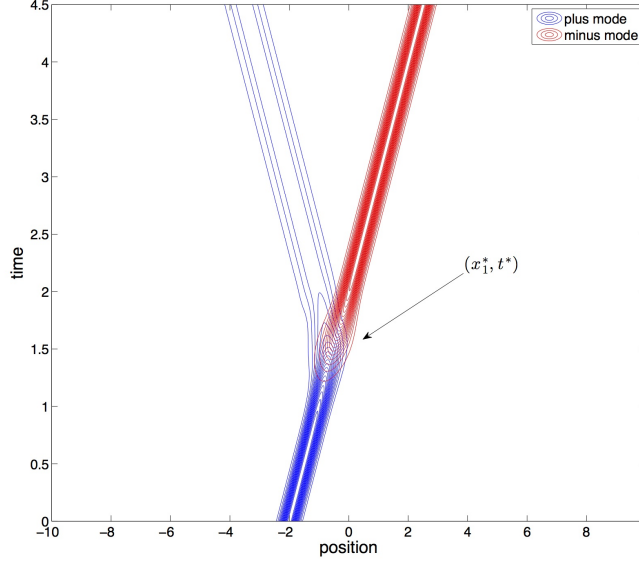


Figure 8: Contour plots of the plus and minus densities, quantum graphene model with $\varepsilon = 0.064$.

Table 2: Transfer rates

ε	0.128	0.064	0.032	0.016	0.008	0.004
T_ε (quantum)	0.772	0.596	0.355	0.126	1.60×10^{-2}	2.55×10^{-4}
T_ε (kinetic)	0.769	0.593	0.351	0.123	1.52×10^{-2}	2.38×10^{-4}
$-\frac{\pi(\xi_2^0)^2}{\varepsilon \log T_\varepsilon}$ (quantum)	0.9481	0.9486	0.9489	0.9491	0.9491	0.9492
$-\frac{\pi(\xi_2^0)^2}{\varepsilon \log T_\varepsilon}$ (kinetic)	0.9364	0.9387	0.9384	0.9383	0.9385	0.9415

for the kinetic model, with and without the jump. For clarity, we plot separately the error on the plus density (continuous red and blue lines) and the error on the minus density (dashed red and blue lines). The final time of the simulation is still $t_f = 4.5$.

We observe on the minus density curves (dashed lines) that the jump process improves significantly the precision of the computation of the post-transition density: without the jump operator and for $\varepsilon > 10^{-2}$, the error on the minus density is of order 1. Note that, for $\varepsilon < 4 \times 10^{-3}$, the transfer rate becomes negligible and the major part of the error becomes the error made on the transport process, studied above: the plus density curves (continuous curves) are comparable to the red curve of Figure 7 (the scales in these two figures are the same).

3.6. N -particles system. In this last numerical experiments, we take the initial data as a mixture of coherent quantum states, with a double bump distribution

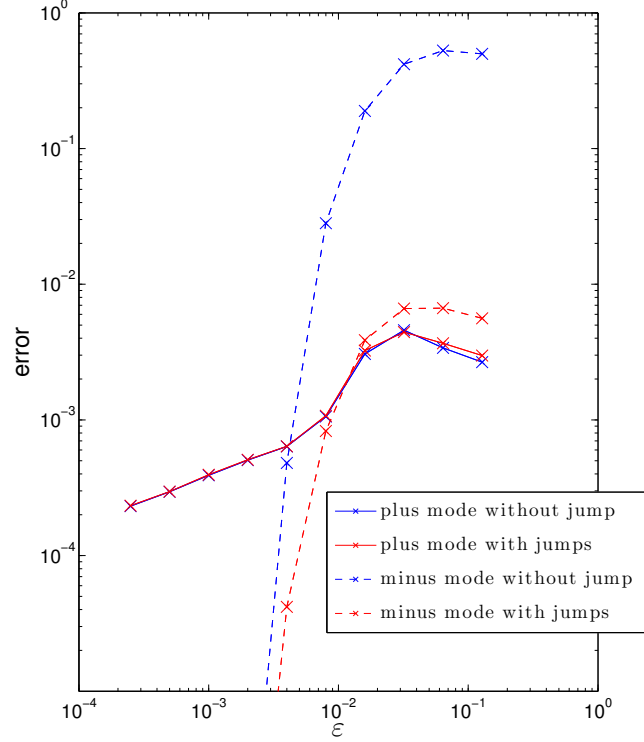


Figure 9: L^1 error on the densities for the kinetic graphene model at time $t_f = 4.5$, as a function of ε , with and without the jump process.

function, i.e. the matrix density is (3.5) with

$$f^0(x^0, \xi^0) = \max(\cos(2\pi x)\mathbb{1}_{-7/4 \leq x \leq -5/4}, \cos(2\pi x)\mathbb{1}_{-11/4 \leq x \leq -9/4}) \times (5/\pi)^{-1/2} \exp(-5(\xi_1 - 1.3)^2).$$

The potential is V_1 defined in (3.11). We plot on Figure 10 the plus and minus densities computed with the quantum and the kinetic graphene model, for $\varepsilon = 0.016$ at times $t = 0, 1.8, 2.7$ and 4.5 . The kinetic model is still used with 4×10^6 particles, and the quantum model is now used with 5000 wavefunctions.

Finally, in Table 3.6, we provide the numerical transfer rates between the plus and minus mixture of coherent states. Again, on Figure 10 and in Table 3.6, we

Table 3: Transfer rates

ε	0.128	0.064	0.032	0.016	0.008	0.004
Transfer rate (quantum)	0.771	0.595	0.355	0.126	1.62×10^{-2}	2.76×10^{-4}
Transfer rate (kinetic)	0.768	0.592	0.350	0.123	1.54×10^{-2}	2.57×10^{-4}

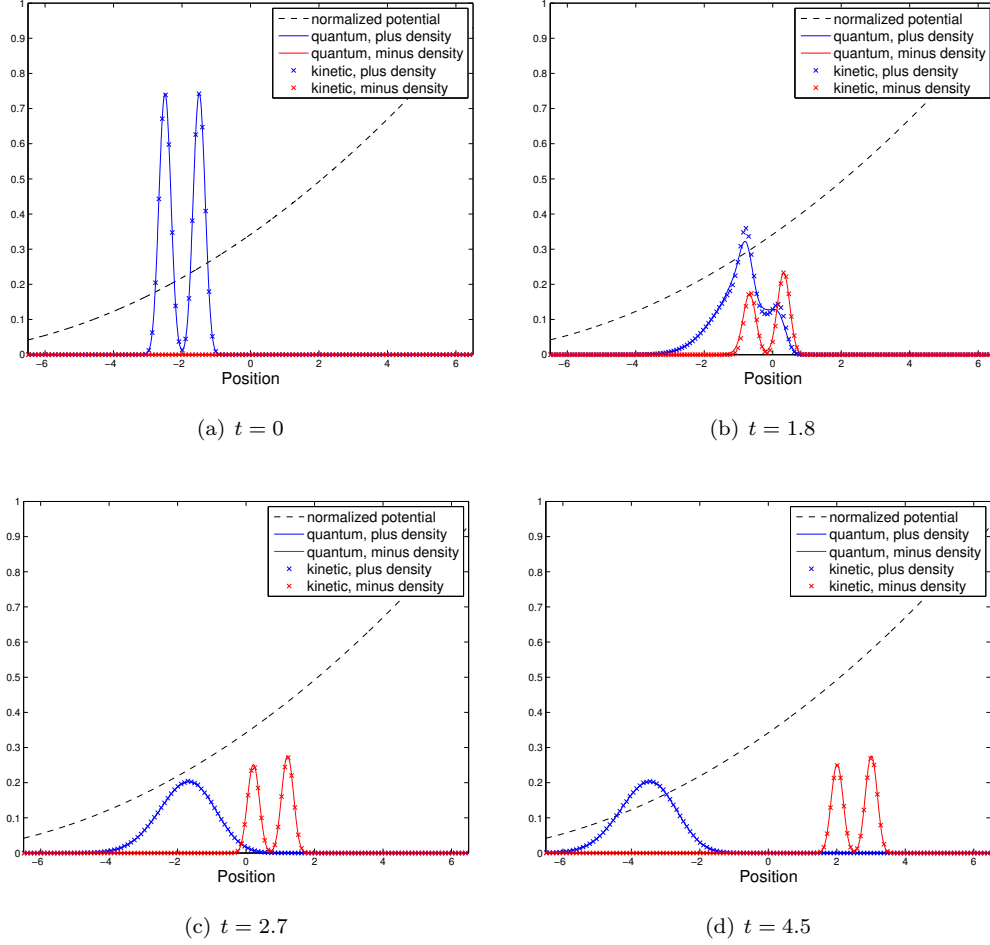


Figure 10: Propagation of a mixture of coherent wavepackets through a barrier for $\varepsilon = 0.016$.

observe a very good agreement between both models.

4. APPENDIX: PSEUDO-DIFFERENTIAL CALCULUS

In this Appendix, we recall a few results of symbolic calculus that we use in this article and we prove the estimate (2.9), which is at the core of these results.

The estimate (2.9) relies on the Schur Lemma. With $f \in L^2(\mathbf{R}^d)$, we associate

$$\mathcal{F}_\varepsilon(f)(\xi) := (2\pi\varepsilon)^{-d} \widehat{f}\left(\frac{\xi}{\varepsilon}\right) = (2\pi\varepsilon)^{-d/2} \int f(x) e^{-\frac{i}{\varepsilon} \xi \cdot x} dx,$$

and we observe that for any $a \in \mathcal{C}_0^\infty(\mathbf{R}^{2d})$,

$$\begin{aligned}
(\text{op}_\varepsilon(a)f, f) &= (2\pi\varepsilon)^{-d} \int a\left(\frac{x+y}{2}, \xi\right) e^{\frac{i}{\varepsilon}\xi \cdot (x-y)} f(y) \bar{f}(x) dx dy d\xi \\
&= (2\pi\varepsilon)^{-3d} \int a\left(\frac{x+y}{2}, \xi\right) e^{\frac{i}{\varepsilon}(\xi \cdot (x-y) + i\eta \cdot y - ix \cdot \zeta)} \\
&\quad \times \widehat{f}\left(\frac{\eta}{\varepsilon}\right) \overline{\widehat{f}}\left(\frac{\zeta}{\varepsilon}\right) dx dy d\xi d\eta d\zeta \\
&= (2\pi\varepsilon)^{-2d} \int a(X, \xi) e^{\frac{i}{\varepsilon}(\xi \cdot v + i\eta \cdot (X - \frac{v}{2}) - i\zeta \cdot (X + \frac{v}{2}))} \\
&\quad \times \mathcal{F}_\varepsilon(f)(\eta) \overline{\mathcal{F}_\varepsilon(f)}(\zeta) dX dv d\xi d\eta d\zeta \\
&= (2\pi\varepsilon)^{-d} \int a\left(-X, \frac{\eta + \zeta}{2}\right) e^{\frac{i}{\varepsilon}X \cdot (\zeta - \eta)} \mathcal{F}_\varepsilon(f)(\eta) \overline{\mathcal{F}_\varepsilon(f)}(\zeta) dX d\eta d\zeta \\
&= (K_\varepsilon \mathcal{F}_\varepsilon(f), \mathcal{F}_\varepsilon(f)),
\end{aligned}$$

where K_ε is the operator of kernel $k_\varepsilon(\xi, \xi')$,

$$k_\varepsilon(\xi, \xi') = \frac{1}{\varepsilon^d} \tilde{a}\left(\frac{\xi + \xi'}{2}, \frac{\xi - \xi'}{\varepsilon}\right),$$

with

$$\tilde{a}(\xi, v) = (2\pi)^{-d} \int a(-x, \xi) e^{ix \cdot v} dx.$$

By the Plancherel theorem, the norm of $\text{op}_\varepsilon(a)$ and of K_ε are the same. Besides, using the Schur Lemma, we obtain

$$\begin{aligned}
\|K_\varepsilon\| &\leq \text{Max} \left(\sup_{\xi \in \mathbf{R}^d} \int |k_\varepsilon(\xi, \xi')| d\xi', \sup_{\xi' \in \mathbf{R}^d} \int |k_\varepsilon(\xi, \xi')| d\xi \right) \\
&\leq \int_{\mathbf{R}^d} \sup_{\xi \in \mathbf{R}^d} |\tilde{a}(\xi, v)| dv \\
&\leq C \sup_{|\beta| \leq d+1} \sup_{\xi \in \mathbf{R}^d} \int |\partial_x^\beta a(x, \xi)| dx.
\end{aligned}$$

for some constants C independent of a and ε , which gives the result.

By use of the same techniques and of the Taylor formula, one can prove the following result about the composition of pseudo differential operators :

Proposition 4.1. *Let $a, b \in \mathcal{C}_0^\infty(\mathbf{R}^{2d}, \mathbf{C}^{N,N})$, $N \in \mathbf{N}$, then*

$$\text{op}_\varepsilon(a)\text{op}_\varepsilon(b) = \text{op}_\varepsilon(ab) + \frac{\varepsilon}{2i} \text{op}_\varepsilon(\{a, b\}) + \varepsilon^2 R_\varepsilon,$$

with $\{a, b\} = \nabla_\xi a \cdot \nabla_x b - \nabla_x a \cdot \nabla_\xi b$ and

$$\|R_\varepsilon\|_{\mathcal{L}(L^2(\mathbf{R}^d))} \leq C \sup_{c=a,b} \sup_{|\alpha|=2} \sup_{2 \leq |\beta| \leq d+3} \sup_{\xi \in \mathbf{R}^d} \int |\partial_\xi^\alpha \partial_x^\beta c(x, \xi)| dx$$

for some constant $C > 0$ independent of a, b and ε .

REFERENCES

- [1] S. Alinhac, P. Gérard, Pseudo-differential operators and the Nash-Moser Theorem, Graduate Studies in Mathematics, **82**, AMS, 2007.
- [2] N. Anantharaman, F. Faure, C. Fermanian Kammerer, Chaos Quantique, Journées mathématiques X-UPS 2014. Comité Editorial : Pascale Harinck, Alain Plagne, Claude Sabbah, *Les Éditions de l'École polytechnique* (2014) <http://www.math.polytechnique.fr/xups/volumes.html>.
- [3] D. Berdebes, T. Low, M. Lundstrom, Low Bias Transport in Graphene: An Introduction, Lecture Notes of the 2009 NCN Purdue Summer School: Electronics from the Bottom Up, available at <http://nanohub.org/topics/ElectronicsFromTheBottomUp>.
- [4] A. P. Calderón, R. Vaillancourt, On the boundedness of pseudodifferential operators, *J. Math. Soc. Japan.* **23**, 2, 374–378 (1971).
- [5] A. H. Castro Neto, F. Guinea, N. M. R. Peres, K. S. Novoselov, A. K. Geim, The electronic properties of graphene, *Rev. Mod. Phys.* **81**, 109–162 (2009).
- [6] Y. Colin de Verdière, The level crossing problem in semi-classical analysis I. The symmetric case, *Ann. Inst. Fourier*, 53, no. 4 (2003), pp. 1023–1054.
- [7] Y. Colin de Verdière, The level crossing problem in semi-classical analysis II. The hermitian case, *Ann. Inst. Fourier*, 54, no. 5 (2004), pp. 1423–1441.
- [8] M. Dimassi, J. Sjöstrand, Spectral asymptotics in the semi-classical limit, London Mathematical Society Lecture Note Series, 268, Cambridge University Press, 1999.
- [9] E. B. Dynkin, Markov processes. Vol. I. Die Grundlehren der Mathematischen Wissenschaften, Bände 121, 122 Academic Press Inc., Publishers, New York; Springer-Verlag, Berlin-Göttingen-Heidelberg 1965.
- [10] S. Das Sarma, S. Adam, E. H. Hwang, E. Rossi, Electronic transport in two-dimensional graphene, *Rev. Mod. Phys.* **83**, 407–470 (2011).
- [11] R. El Hajj, F. Méhats, Analysis of models for quantum transport of electrons in graphene layers, *Math. Models Methods Appl. Sci.* **24**, 11, 2287–2310 (2014).
- [12] A. Faraj, S. Jin, Some studies on asymptotic and numerical approximation of 2D Dirac equation for graphene, preprint.
- [13] C. Fermanian Kammerer, Normal forms for conical intersections in quantum chemistry, *Math. Phys. Elect. Jour.*, **13**, N° 4 (2007).
- [14] C. Fermanian Kammerer, P. Gérard, Mesures semi-classiques et croisements de modes, *Bull. Soc. math. France*, **130**, N°1, 123–168 (2002).
- [15] C. Fermanian Kammerer, P. Gérard, A Landau-Zener formula for non-degenerated involutive codimension 3 crossings., *Ann. Henri Poincaré*, **4**, N°1, 513–552 (2003).
- [16] C. Fermanian Kammerer, P. Gérard, C. Lasser, Wigner measure propagation and Lipschitz conical singularity for general initial data, *Arch. Ration. Mech. Anal.*, **209**, 209–236 (2013).
- [17] C. Fermanian Kammerer, C. Lasser, Propagation through generic level crossings: a surface hopping semigroup. *SIAM J. of Math. Anal.* , **140**, 1, p. 103-133 (2008).
- [18] C. Fermanian Kammerer, C. Lasser, Single switch surface hopping for molecular dynamics with transitions. *Journal of Chemical Physics*, **128**, 144102 (2008).
- [19] C. Fermanian Kammerer, C. Lasser, Single switch surface hopping for molecular dynamics. Proceedings of the *International meeting on "Mathematical Methods for Ab Initio Quantum Chemistry" – Nice, 2008* .
- [20] C. Fermanian Kammerer, C. Lasser, Egorov Theore for avoided crossings of eigenvalue surfaces (*work in progress*).
- [21] P. Gérard, E. Leichtnam, Ergodic properties of eigenfunctions for the Dirichlet problem, *Duke Math. J.* **71**, no. 2 (1993), pp. 559–607.
- [22] P. Gérard, P. A. Markowich, N. J. Mauser, F. Poupaud, Homogenization Limits and Wigner Transforms, *Comm. Pure Appl. Math.*, **50**, 4, 323–379 (1997) and Erratum: Homogenization limits and Wigner Transforms, *Comm. Pure Appl. Math.*, **53**, 280–281 (2000).
- [23] G. Hagedorn, Molecular Propagation through Electron Energy Level Crossings, *Mem. Am. Math. Soc.*, 111, no. 536, 1994.
- [24] G. A. Hagedorn, A. Joye, Landau-Zener transitions through small electronic eigenvalue gaps in the Born-Oppenheimer approximation. *Ann. Inst. Henri Poincaré*, **68**, N°1, p. 85-134 (1998).

- [25] E. Hairer, C. Lubich, G. Wanner, Geometric numerical integration. Structure-preserving algorithms for ordinary differential equations. Reprint of the second (2006) edition. Springer Series in Computational Mathematics, 31. Springer, Heidelberg, 2010.
- [26] V. Hung Nguyen, A. Bournel, C. Chassat, P. Dollfus, Quantum transport of Dirac fermions in graphene field effect transistors, International Conference on Simulation of Semiconductor Processes and Devices (SISPAD), Bologne, Italie, 6–8 September 2010, SISPAD 2010 Proceedings, 9–12 (2010).
- [27] L. Landau, Collected papers of L. Landau, Pergamon Press, 1965.
- [28] H. Spohn and S. Teufel, Adiabatic decoupling and time-dependent Born-Oppenheimer theory, *Commun. Math. Phys.*, 224, no. 1 (2001), pp. 113–132.
- [29] C. Lasser, S. Teufel, Propagation through Conical Crossings: an Asymptotic Semigroup, *Comm. Pure Appl. Math.* 58, 9, 1188–1230 (2005).
- [30] T. Liggett, Interacting particle systems. Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences], 276. Springer-Verlag, New York, 1985.
- [31] A. Martinez, V. Sordani, Twisted Pseudodifferential Calculus and Application to the Quantum Evolution of Molecules, *Memoirs of American Mathematical Society*, 200, p. 1 - 82 (2009).
- [32] O. Morandi, F. Schürer, Wigner model for quantum transport in graphene, *J. Phys. A: Math. Theor.* 44 (2011) 265301.
- [33] C. Zener, Non-adiabatic crossing of energy levels, *Proc. Roy. Soc. Lond.*, 137 (1932), pp. 696–702.
- [34] M. Zworski, *Semi-classical analysis*, Graduate Studies in Mathematics, 138, AMS, 2012.

(C. Fermanian) LAMA UMR CNRS 8050, UNIVERSITÉ PARIS EST, 61, AVENUE DU GÉNÉRAL DE GAULLE, 94010 CRÉTEIL CEDEX, FRANCE

E-mail address: Clotilde.Fermanian@u-pec.fr

(F. Méhats) IRMAR, UNIVERSITÉ RENNES 1 AND IPSO INRIA TEAM, CAMPUS DE BEAULIEU, 35042 RENNES CEDEX, FRANCE

E-mail address: florian.mehats@univ-rennes1.fr