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Error estimates in periodic homogenization with a non-homogeneous Dirichlet condition.

G. Griso

Laboratoire J.-L. Lions-CNRS, Boîte courrier 187, Université Pierre et Marie Curie, 4 place Jussieu, 75005 Paris, France, Email: griso@ann.jussieu.fr

Abstract

In this paper we investigate the homogenization problem with a non-homogeneous Dirichlet condition. Our aim is to give error estimates with boundary data in $H^{1/2}(\partial\Omega)$. The tools used are those of the unfolding method in periodic homogenization.

1 Introduction

We consider the following homogenization problem:

$$\phi^{\varepsilon} \in H^1(\Omega), \quad -\text{div}(A_{\varepsilon} \nabla \phi^{\varepsilon}) = f \quad \text{in} \quad \Omega, \quad \phi^{\varepsilon} = g \quad \text{on} \quad \partial \Omega$$

where A_{ε} is a periodic matrix satisfying the usual condition of uniform ellipticity and where $f \in L^2(\Omega)$ and $g \in H^{1/2}(\partial\Omega)^1$. We know (see e.g. [9]) that the function ϕ^{ε} weakly converges in $H^1(\Omega)$ towards the solution Φ of the homogenized problem

$$\Phi \in H^1(\Omega), \quad -\text{div}(\mathcal{A}\nabla\Phi) = f \quad \text{in} \quad \Omega, \quad \Phi = g \quad \text{on} \quad \partial\Omega$$

where \mathcal{A} is the homogenized matrix (see (4.4) and (4.5)). Using the results of [9] we can give an approximation of ϕ^{ε} belonging to $H^{1}(\Omega)$ and we easily obtain

$$\phi^{\varepsilon} - \Phi - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{i}} \right) \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \longrightarrow 0 \quad \text{strongly in} \quad H^{1}(\Omega)$$

where Q_{ε} is the *scale-splitting operator* (see also Subsection 2.4) and where the χ_i are the correctors (see (4.2)).

The aim of this paper is to give error estimates. Obviously, if we have $g \in H^{3/2}(\partial\Omega)$ and the appropriate assumptions on the boundary of the domain then we can apply the

¹The homogenization problem with a L^p boundary data is investigated in [3].

results obtained in [4], [12], [13], [14], [15] and [20] to deduce error estimates. All of them require that the function Φ belongs at least to $H^2(\Omega)$. Here, the solution Φ of the homogenized problem is only in $H^1(\Omega) \cap H^2_{loc}(\Omega)$. In this paper we must work with this lack of regularity; this is the main difficulty.

Our method to obtain error estimates (see [13], [14] and [15]) is mainly based on projection theorems. This is why we prove here two new projection theorems. In the second one, for a function $\phi \in H_0^1(\Omega)$ satisfying $1/\rho \nabla \phi \in L^2(\Omega; \mathbb{R}^n)$ -where $\rho(x)$ is the distance between x and the boundary of Ω - we obtain an upper bound for the distance between $\mathcal{T}_{\varepsilon}(\nabla \phi)$ and the space $\nabla H^1(\Omega) \oplus \nabla_y L^2(\Omega; H^1_{per}(Y))$ for an appropriate norm (see Section 3). Then, due to the result recalled in Subsection 7.1 of the Appendix we introduce a lifting in $H^1(\Omega)$ for the function $g \in H^{1/2}(\partial \Omega)$ and we show estimates in $H^1(\Omega)$ and in $H^1_{loc}(\Omega)$ using $||g||_{H^{1/2}(\partial \Omega)}$ and $||g||_{H^{-1/2}(\partial \Omega)}$. Afterwards, all the tools of the unfolding method (see [9]) and the results obtained in the first sections allow to derive the main results of the paper (Theorems 5.1 and 5.3). In both theorems we give L^2 error estimates and H^1_{loc} error estimates. It is worth noting that the error estimates are only of order $\varepsilon^{1/2}$ while in [15] the obtained error was of order ε . We end the paper by investigating the case where the boundary data are strongly oscillating. A forthcoming paper we will be devoted to the homogenization problem with other strongly oscillating boundary data.

In Section 2 we introduce a few general notations, we also give some recalls on lemmas, definitions and results about the unfolding method in periodic homogenization; this complements the paper which presents the unfolding method (see [9]). Section 3 is devoted to the proofs of two new projection theorems which supplement the ones given in [14] and [15]. In Section 4, we recall the main results on the classical homogenization problem. In Section 5 we derive the error estimate results (Theorems 5.1 and 5.3) with a non-homogenous Dirichlet condition and in the last section we investigate a first case of a strongly oscillating boundary data (Theorem 6.1). In the Appendix we introduce an operator from $H^{-1/2}(\partial\Omega)$ into $L^2(\Omega)$.

As general references on the homogenization theory we refer to [1], [4] and [12]. The reader is referred to [9], [11] and [12] for an introduction of the unfolding method in periodic homogenization. The following papers [5], [6], [7], [10] give various applications of the unfolding method in periodic homogenization. As far as the error estimates are concerned, we refer to [2], [4], [13], [14], [15], [18], [20] and [21].

Keywords: periodic homogenization, error estimate, non-homogeneous Dirichlet condition, unfolding method.

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2 Preliminaries

2.1 Notations

• We denote by Ω a bounded domain in \mathbb{R}^n with a Lipschitz boundary.² Let $\rho(x)$ be the distance between $x \in \mathbb{R}^n$ and the boundary of Ω , we set

$$\widetilde{\Omega}_{\gamma} = \left\{ x \in \Omega \mid \rho(x) < \gamma \right\} \qquad \widetilde{\widetilde{\Omega}}_{\gamma} = \left\{ x \in \mathbb{R}^n \mid \rho(x) < \gamma \right\} \qquad \gamma \in \mathbb{R}^{*+}.$$

• There exist constants a, A and γ_0 strictly positive and $M \geq 1$, a finite number N of local euclidian coordinate systems $(O_r; \mathbf{e}_{1r}, \dots, \mathbf{e}_{nr})$ and mappings $f_r : [-a, a]^{n-1} \longrightarrow \mathbb{R}$, Lipschitz continuous with ratio M, $1 \leq r \leq N$, such that (see e.g. [16] or [17])

$$\partial\Omega = \bigcup_{r=1}^{N} \left\{ x = x'_r + x_{nr} \mathbf{e}_{nr} \in \mathbb{R}^n \mid x'_r \in \Delta_a \text{ and } x_{nr} = f_r(x'_r) \right\},$$
where $x'_r = x_{1r} \mathbf{e}_{1r} + \ldots + x_{n-1r} \mathbf{e}_{n-1r}, \quad \Delta_a = \left\{ x'_r \mid x_{ir} \in] - a, a[, i \in \{1, \ldots, n-1\} \right\}$

$$\widetilde{\Omega}_{\gamma_0} \subset \bigcup_{r=1}^{N} \Omega_r \subset \Omega, \quad \Omega_r = \left\{ x \in \mathbb{R}^n \mid x'_r \in \Delta_a \text{ and } f_r(x'_r) < x_{nr} < f_r(x'_r) + A \right\}$$

$$\widetilde{\widetilde{\Omega}}_{\gamma_0} \subset \bigcup_{r=1}^{N} \left\{ x \in \mathbb{R}^n \mid x'_r \in \Delta_a \text{ and } f_r(x'_r) - A < x_{nr} < f_r(x'_r) + A \right\}$$

$$\forall r \in \{1, \ldots, N\}, \quad \forall x \in \Omega_r \text{ we have } \frac{1}{2M} (x_{nr} - f_r(x'_r)) \leq \rho(x) \leq x_{nr} - f_r(x'_r).$$

$$(2.1)$$

• We set

$$Y =]0,1[^n, \qquad \Xi_{\varepsilon} = \{ \xi \in \mathbb{Z}^n \mid \varepsilon(\xi + Y) \subset \Omega \},$$

$$\widehat{\Omega}_{\varepsilon} = \operatorname{interior} \left(\bigcup_{\xi \in \Xi_{\varepsilon}} \varepsilon(\xi + \overline{Y}) \right), \qquad \Lambda_{\varepsilon} = \Omega \setminus \widehat{\Omega}_{\varepsilon},$$

where ε is a strictly positive real.

• We define

$$\begin{split} \star \ & H^1_\rho(\Omega) = \Big\{\phi \in L^2(\Omega) \mid \rho \nabla \phi \in L^2(\Omega;\mathbb{R}^n) \Big\}, \\ \star \ & L^2_{1/\rho}(\Omega) = \Big\{\phi \in L^2(\Omega) \mid \phi/\rho \in L^2(\Omega) \Big\}, \\ \star \ & H^1_{1/\rho}(\Omega) = \Big\{\phi \in H^1_0(\Omega) \mid \nabla \phi/\rho \in L^2(\Omega;\mathbb{R}^n) \Big\}. \end{split}$$

We endow $H^1_{\rho}(\Omega)$ (resp. $H^1_{1/\rho}(\Omega)$) with the norm

$$\forall \phi \in H^1_{\rho}(\Omega), \quad ||\phi||_{\rho} = ||\phi||_{L^2(\Omega)} + ||\rho \nabla \phi||_{L^2(\Omega;\mathbb{R}^n)}$$
 (resp.
$$\forall \phi \in H^1_{1/\rho}(\Omega), \quad ||\phi||_{1/\rho} = ||\nabla \phi/\rho||_{L^2(\Omega;\mathbb{R}^n)}).$$

²In Section 7.1 and those which follow, we will assume that Ω is a bounded domain of class $\mathcal{C}^{1,1}$ or an open bounded convex polygon (n=2) or polyhedral (n=3).

Remark 2.1. If ϕ belongs to $H^1_{\rho}(\Omega)$ then the function $\psi = \rho \phi$ is in $H^1_0(\Omega)$ and vice versa if the function ψ belongs to $H^1_0(\Omega)$ then $\phi = \psi/\rho$ is in $H^1_{\rho}(\Omega)$ since we have (see [8] or [19])

$$\forall \psi \in H_0^1(\Omega), \qquad \left\| \psi/\rho \right\|_{L^2(\Omega)} \le C||\nabla \psi||_{L^2(\Omega;\mathbb{R}^n)}. \tag{2.2}$$

The space \mathbb{R}^k $(k \geq 1)$ is endowed with the standard basis $(\mathbf{e}_1, \dots, \mathbf{e}_k)$; the euclidian norm is denoted $|\cdot|$.

2.2 A characterization of the functions belonging to $H^1_{1/\rho}(\Omega)$

Observe first that if a function ϕ satisfies $\phi/\rho \in H_0^1(\Omega)$ then ϕ belongs to $H_{1/\rho}^1(\Omega)$. The reverse is true.

Lemma 2.2. Let Ω be a bounded open set with a Lipschitz boundary, we have

$$\phi \in H^1_{1/\rho}(\Omega) \iff \phi/\rho \in H^1_0(\Omega).$$

Furthermore there exists a constant which depends only on $\partial\Omega$ such that

$$\forall \phi \in H^{1}_{1/\rho}(\Omega) \qquad \|\phi/\rho^{2}\|_{L^{2}(\Omega)} + \|\phi/\rho\|_{H^{1}(\Omega)} \le C||\phi||_{1/\rho}. \tag{2.3}$$

Proof. Step 1. Let ϕ be in $H^1(]-a,a[^{n-1}\times]0,A[)$ (a,A>0) satisfying $\frac{1}{x_n}\nabla\phi(x)\in L^2(]-a,a[^{n-1}\times]0,A[)$ and $\phi(x)=0$ for a.e. x in $]-a,a[^{n-1}\times\{0\}\cup]-a,a[^{n-1}\times\{A\}.$ We have

$$\int_{|-a,a|^{n-1}\times|0,A|} \frac{|\phi(x)|^2}{x_n^4} dx \le \frac{1}{2} \int_{|-a,a|^{n-1}\times|0,A|} \frac{|\nabla\phi(x)|^2}{x_n^2} dx. \tag{2.4}$$

To prove (2.4), we choose $\eta > 0$ and we integrate by parts $\int_{]-a,a]^{n-1}\times]0,A[} \frac{|\phi(x)|^2}{(\eta+x_n)^4} dx$, then thanks to the identity relation $2bc \leq b^2 + c^2$ we obtain

$$\int_{]-a,a]^{n-1}\times]0,A[} \frac{|\phi(x)|^2}{(\eta+x_n)^4} dx \le \frac{1}{2} \int_{]-a,a]^{n-1}\times]0,A[} \frac{1}{(\eta+x_n)^2} \left| \frac{\partial \phi}{\partial x_n}(x) \right|^2 dx
\le \frac{1}{2} \int_{]-a,a]^{n-1}\times]0,A[} \frac{|\nabla \phi(x)|^2}{x_n^2} dx.$$

Passing to the limit $(\eta \to 0)$ it leads to (2.4).

Step 2. Let h be in $W^{1,\infty}(\Omega)$ such that

$$h(x) \in [0, 1]$$
 if $x \in \Omega$,
 $h(x) = 1$ if $\rho(x) \ge \gamma_0$,
 $h(x) = 0$ if $\rho(x) \le \gamma_0/2$.

Let ϕ be in $H^1_{1/\rho}(\Omega)$. The function $\phi h/\rho^4$ belongs to $H^1_0(\Omega)$, therefore as a consequence of the Poincaré's inequality we obtain

$$\int_{\Omega} \frac{|\phi(x)h(x)|^2}{\rho(x)^4} dx \le C \int_{\Omega} \left| \nabla \left(\frac{\phi(x)h(x)}{\rho(x)^4} \right) \right|^2 dx \le C \int_{\Omega} \left(|\nabla \phi(x)|^2 + |\phi(x)|^2 \right) dx
\le C \int_{\Omega} |\nabla \phi(x)|^2 dx \le C \int_{\Omega} \frac{|\nabla \phi(x)|^2}{\rho(x)^2} dx.$$
(2.5)

Then, due to the covering (2.1) of $\widetilde{\Omega}_{\gamma_0}$, the inequality (2.4) and thanks to a simple change of variables we get

$$\int_{\Omega_r} \frac{|\phi(x)(1 - h(x))|^2}{\rho(x)^4} dx \le C \int_{\Omega_r} \frac{|\nabla(\phi(x)(1 - h(x)))|^2}{\rho(x)^2} dx \le C \int_{\Omega_r} \frac{|\nabla\phi(x)|^2 + |\phi(x)|^2}{\rho(x)^2} dx.$$

Since $\phi \in H_0^1(\Omega)$ the function ϕ/ρ belongs to $L^2(\Omega)$ and we have (2.2). Hence, adding these inequalities (r = 1, ..., N) we obtain

$$\int_{\Omega} \frac{|\phi(x)(1-h(x))|^2}{\rho(x)^4} dx \le C \int_{\Omega} \frac{|\nabla \phi(x)|^2}{\rho(x)^2} dx. \tag{2.6}$$

Finally $\phi/\rho^2 \in L^2(\Omega)$ and (2.5)-(2.6) lead to $\|\phi/\rho^2\|_{L^2(\Omega)} \leq C||\phi||_{1/\rho}$ and then (2.3). \square

2.3 Three lemmas

In the below lemma we give sharp estimates of a function on the boundary and in a neighborhood of the boundary.

Lemma 2.3. Let Ω be a bounded open set with a Lipschitz boundary, there exists $\gamma_0 > 0$ (see Subsection 2.2) such that for any $\gamma \in]0, \gamma_0]$ and for any $\phi \in H^1(\Omega)$ we have

$$||\phi||_{L^{2}(\partial\Omega)} \leq \frac{C}{\gamma^{1/2}} (||\phi||_{L^{2}(\widetilde{\Omega}_{\gamma})} + \gamma||\nabla\phi||_{L^{2}(\widetilde{\Omega}_{\gamma};\mathbb{R}^{n})}),$$

$$||\phi||_{L^{2}(\widetilde{\Omega}_{\gamma})} \leq C (\gamma^{1/2}||\phi||_{L^{2}(\partial\Omega)} + \gamma||\nabla\phi||_{L^{2}(\widetilde{\Omega}_{\gamma};\mathbb{R}^{n})}).$$

$$(2.7)$$

The constants do not depend on γ .

Proof. Let ψ be in $H^1(]-a,a[^{n-1}\times]0,A[)$. For any $\eta\in]0,A[$ we have

$$\begin{aligned} ||\psi||_{L^2(]-a,a[^{n-1}\times\{0\})}^2 &\leq \frac{C}{\eta} ||\psi||_{L^2(]-a,a[^{n-1}\times]0,\eta[)}^2 + C\eta ||\nabla\psi||_{L^2(]-a,a[^{n-1}\times]0,\eta[;\mathbb{R}^n)}^2, \\ ||\psi||_{L^2(]-a,a[^{n-1}\times]0,\eta[)}^2 &\leq C\eta ||\psi||_{L^2(]-a,a[^{n-1}\times\{0\})}^2 + C\eta^2 ||\nabla\psi||_{L^2(]-a,a[^{n-1}\times]0,\eta[;\mathbb{R}^n)}^2. \end{aligned}$$

The constants do not depend on η . Now, let ϕ be in $H^1(\Omega)$. We use the above estimates, the covering of $\widetilde{\Omega}_{\gamma_0}$ given by (2.1) and a simple change of variables to obtain (2.7). \square

Below we recall a classical extension lemma which is proved for example in [14] or which can be proved using the covering (2.1).

Lemma 2.4. Let Ω be a bounded open set with a Lipschitz boundary, there exist $c_0 \geq 1$ (which depends only on the boundary of Ω) and a linear and continuous extension operator \mathcal{P} from $L^2(\Omega)$ into $L^2(\mathbb{R}^n)$ which also maps $H^1(\Omega)$ into $H^1(\mathbb{R}^n)$ such that

$$\forall \phi \in L^{2}(\Omega), \quad \mathcal{P}(\phi)|_{\Omega} = \phi, \qquad ||\mathcal{P}(\phi)||_{L^{2}(\mathbb{R}^{n})} \leq C||\phi||_{L^{2}(\Omega)}, \\ ||\mathcal{P}(\phi)||_{L^{2}(\widetilde{\Omega}_{\gamma})} \leq C||\phi||_{L^{2}(\widetilde{\Omega}_{c_{0}\gamma})}$$

$$(2.8)$$

and moreover we have

$$\forall \phi \in H^1(\Omega), \qquad ||\nabla \mathcal{P}(\phi)||_{L^2(\mathbb{R}^n;\mathbb{R}^n)} \le C||\nabla \phi||_{L^2(\Omega;\mathbb{R}^n)}.$$

From now on, if need be, a function ϕ belonging to $L^2(\Omega)$ (resp. $H^1(\Omega)$) will be extended to a function belonging to $L^2(\mathbb{R}^n)$ (resp. $H^1(\mathbb{R}^n)$) using the above lemma. The extension will be still denoted ϕ .

In the third lemma we show that a function in $H_0^1(\Omega)$ can be approached by functions vanishing outside of $\widetilde{\Omega}_{6\sqrt{n}\varepsilon}$.

Lemma 2.5. Let ϕ be in $H_0^1(\Omega)$, there exists $\phi_{\varepsilon} \in H^1(\mathbb{R}^n)$ satisfying

$$\phi_{\varepsilon}(x) = 0 \quad \text{for a.e. } x \notin \widetilde{\Omega}_{6\sqrt{n}\varepsilon},$$

$$||\phi - \phi_{\varepsilon}||_{L^{2}(\Omega)} \le C\varepsilon ||\nabla \phi||_{L^{2}(\Omega;\mathbb{R}^{n})}, \qquad ||\phi_{\varepsilon}||_{H^{1}(\Omega)} \le C||\phi||_{H^{1}(\Omega)}.$$

$$(2.9)$$

Moreover, if $\phi \in H^1_{1/\rho}(\Omega)$ then we have

$$\left\| \left(\phi - \phi_{\varepsilon} \right) / \rho \right\|_{L^{2}(\Omega)} \le C \varepsilon ||\nabla \phi||_{1/\rho}, \qquad ||\phi_{\varepsilon}||_{1/\rho} \le C ||\phi||_{1/\rho}. \tag{2.10}$$

The constant C is independent of ε .

Proof. Let ϕ be in $H_0^1(\Omega)$. We define ϕ_{ε} by

$$\phi_{\varepsilon}(x) = \begin{cases} \frac{(\rho(x) - 6\sqrt{n}\varepsilon)^{+}}{\rho(x)} \phi(x) & \text{for a. e. } x \in \Omega, \\ 0 & \text{for a. e. } x \in \mathbb{R}^{n} \setminus \overline{\Omega}. \end{cases}$$

where $\delta^+ = \max\{0, \delta\}$. The above function ϕ_{ε} belongs to $H^1(\mathbb{R}^n)$ and satisfies $\phi_{\varepsilon} = 0$ outside of $\widetilde{\Omega}_{6\sqrt{n}\varepsilon}$. Then, due to the fact that ϕ/ρ belongs to $L^2(\Omega)$ and verifies $\|\phi/\rho\|_{L^2(\Omega)} \leq C\|\nabla\phi\|_{L^2(\Omega;\mathbb{R}^n)}$ we obtain the estimates in (2.9). If $\phi \in H^1_{1/\rho}(\Omega)$ we use the estimate (2.3) to obtain (2.10).

2.4 Recalls and complements on the unfolding operators

In the sequel, we will make use of some definitions and results from [9] concerning the periodic unfolding method. For almost every $x \in \mathbb{R}^n$, there exists an unique element in \mathbb{Z}^n denoted [x] such that

$$x = [x] + \{x\}, \qquad \{x\} \in Y.$$

• The unfolding operator $\mathcal{T}_{\varepsilon}$.

For any $\phi \in L^1(\Omega)$, the function $\mathcal{T}_{\varepsilon}(\phi) \in L^1(\Omega \times Y)$ is given by

$$\mathcal{T}_{\varepsilon}(\phi)(x,y) = \begin{cases} \phi\left(\varepsilon\left[\frac{x}{\varepsilon}\right] + \varepsilon y\right) & \text{for a.e. } (x,y) \in \widehat{\Omega}_{\varepsilon} \times Y, \\ 0 & \text{for a.e. } (x,y) \in \Lambda_{\varepsilon} \times Y. \end{cases}$$
(2.11)

Since $\Lambda_{\varepsilon} \subset \widetilde{\Omega}_{\sqrt{n}\varepsilon}$, using Proposition 2.5 in [9] we get

$$\left| \int_{\Omega} \phi(x) dx - \int_{\Omega \times Y} \mathcal{T}_{\varepsilon}(\phi)(x, y) dx dy \right| \le \int_{\Lambda_{\varepsilon}} |\phi(x)| dx \le ||\phi||_{L^{1}(\widetilde{\Omega}_{\sqrt{n}\varepsilon})}$$
 (2.12)

For $\phi \in L^2(\Omega)$ we have

$$||\mathcal{T}_{\varepsilon}(\phi)||_{L^{2}(\Omega)} \le ||\phi||_{L^{2}(\Omega)}. \tag{2.13}$$

We also have (see Proposition 2.5 in [9]) for $\phi \in H^1(\Omega)$ (resp. $\psi \in H^1(\Omega)$)

$$||\mathcal{T}_{\varepsilon}(\phi) - \phi||_{L^{2}(\widehat{\Omega}_{\varepsilon} \times Y)} \leq C\varepsilon ||\nabla \phi||_{L^{2}(\Omega; \mathbb{R}^{n})}$$
(resp.
$$||\mathcal{T}_{\varepsilon}(\psi) - \psi||_{L^{2}(\Omega \times Y)} \leq C\varepsilon ||\nabla \psi||_{L^{2}(\Omega; \mathbb{R}^{n})}$$
). (2.14)

• The local average operator $\mathcal{M}_{\varepsilon}$

For $\phi \in L^1(\mathbb{R}^n)$, the function $\mathcal{M}_{\varepsilon}(\phi) \in L^{\infty}(\mathbb{R}^n)$ is defined by

$$\mathcal{M}_{\varepsilon}(\phi)(x) = \int_{Y} \phi\left(\varepsilon\left[\frac{x}{\varepsilon}\right] + \varepsilon y\right) dy$$
 for a.e. $x \in \mathbb{R}^{n}$. (2.15)

The value of $\mathcal{M}_{\varepsilon}(\phi)$ in the cell $\varepsilon(\xi + Y)$ $(\xi \in \mathbb{Z}^n)$ will be denoted $\mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi)$. In [9] we proved the following results:

For $\phi \in L^2(\Omega)$ we have

$$||\mathcal{M}_{\varepsilon}(\phi)||_{L^{2}(\Omega)} \leq C||\phi||_{L^{2}(\Omega)}, \qquad ||\mathcal{M}_{\varepsilon}(\phi) - \phi||_{H^{-1}(\Omega)} \leq C\varepsilon||\phi||_{L^{2}(\Omega)}$$
(2.16)

and for $\psi \in H_0^1(\Omega)$ (resp. $\phi \in H^1(\Omega)$) we have

$$||\mathcal{M}_{\varepsilon}(\psi) - \psi||_{L^{2}(\Omega)} \leq C\varepsilon ||\nabla \psi||_{L^{2}(\Omega;\mathbb{R}^{n})}$$
(resp.
$$||\mathcal{M}_{\varepsilon}(\phi) - \phi||_{L^{2}(\widehat{\Omega}_{\varepsilon})} \leq C\varepsilon ||\nabla \phi||_{L^{2}(\Omega;\mathbb{R}^{n})}.$$
(2.17)

Lemma 2.6. For $\phi \in H^1_{\rho}(\Omega)$ we have

$$||\rho(\mathcal{M}_{\varepsilon}(\phi) - \phi)||_{L^{2}(\Omega)} \leq C\varepsilon||\phi||_{\rho},$$

$$\forall i \in \{1, \dots, n\}, \qquad ||\rho(\phi(\cdot + \varepsilon \mathbf{e}_{i}) - \phi)||_{L^{2}(\Omega)} \leq C\varepsilon||\phi||_{\rho},$$

$$||\rho(\mathcal{M}_{\varepsilon}(\phi)(\cdot + \varepsilon \mathbf{e}_{i}) - \mathcal{M}_{\varepsilon}(\phi))||_{L^{2}(\Omega)} \leq C\varepsilon||\phi||_{\rho}.$$

$$(2.18)$$

For $\phi \in L^2_{1/\rho}(\Omega)$ we have

$$||\mathcal{M}_{\varepsilon}(\phi) - \phi||_{(H_{\sigma}^{1}(\Omega))'} \le C\varepsilon||\phi/\rho||_{L^{2}(\Omega)}.$$
(2.19)

The constants do not depend on ε .

Proof. Step 1. We prove $(2.18)_1$. Let ϕ be in $H^1_{\rho}(\Omega)$ and let $\varepsilon(\xi + Y)$ be a cell included in Ω .

Case 1: $\rho(\varepsilon\xi) \geq 2\sqrt{n\varepsilon}$. In this case, observing that

$$1 \le \frac{\max_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}}{\min_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}} \le 3$$

and thanks to the Poincaré-Wirtinger's inequality we obtain

$$\int_{\varepsilon(\xi+Y)} [\rho(x)]^{2} |\mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi) - \phi(x)|^{2} dx \leq \left[\max_{z \in \varepsilon(\xi+Y)} {\{\rho(z)\}}\right]^{2} \int_{\varepsilon(\xi+Y)} |\mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi) - \phi(x)|^{2} dx
\leq \left[\max_{z \in \varepsilon(\xi+Y)} {\{\rho(z)\}}\right]^{2} C\varepsilon^{2} \int_{\varepsilon(\xi+Y)} |\nabla\phi(x)|^{2} dx
\leq C\varepsilon^{2} \int_{\varepsilon(\xi+Y)} [\rho(x)]^{2} |\nabla\phi(x)|^{2} dx.$$

Case 2: $\rho(\varepsilon \xi) \leq 2\sqrt{n}\varepsilon$. In this case we have

$$\int_{\varepsilon(\xi+Y)} [\rho(x)]^2 |\mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi) - \phi(x)|^2 dx \le C\varepsilon^2 \int_{\varepsilon(\xi+Y)} |\phi(x)|^2 dx.$$

The cases 1 and 2 lead to

$$\int_{\widehat{\Omega}_{\varepsilon}} [\rho(x)]^2 |\mathcal{M}_{\varepsilon}(\phi)(x) - \phi(x)|^2 dx \le C\varepsilon^2 \int_{\widehat{\Omega}_{\varepsilon}} ([\rho(x)]^2 |\nabla \phi(x)|^2 + |\phi(x)|^2) dx. \tag{2.20}$$

Then, since $\Lambda_{\varepsilon} \subset \widetilde{\Omega}_{\sqrt{n}\varepsilon}$ and thanks to (2.8) we get

$$\int_{\Lambda_{\varepsilon}} [\rho(x)]^2 |\mathcal{M}_{\varepsilon}(\phi)(x) - \phi(x)|^2 dx \le C\varepsilon^2 \int_{\widetilde{\Omega}_{co\sqrt{\eta_{\varepsilon}}}} |\phi(x)|^2 dx$$

which due to (2.20) gives $(2.18)_1$. Proceeding in the same way we obtain $(2.18)_2$ and $(2.18)_3$.

Step 2. We prove (2.19). Let ϕ be in $L^2_{1/\rho}(\Omega)$ and $\psi \in H^1_{\rho}(\Omega)$. We have

$$\int_{\widehat{\Omega}_{\varepsilon}} (\mathcal{M}_{\varepsilon}(\phi) - \phi) \psi = \int_{\widehat{\Omega}_{\varepsilon}} (\mathcal{M}_{\varepsilon}(\psi) - \psi) \phi.$$

Consequently we obtain

$$\left| \int_{\Omega} \left(\mathcal{M}_{\varepsilon}(\phi) - \phi \right) \psi - \int_{\Omega} \left(\mathcal{M}_{\varepsilon}(\psi) - \psi \right) \phi \right| \leq \int_{\Lambda_{\varepsilon}} \left| \left(\mathcal{M}_{\varepsilon}(\phi) - \phi \right) \psi \right| + \int_{\Lambda_{\varepsilon}} \left| \left(\mathcal{M}_{\varepsilon}(\psi) - \psi \right) \phi \right| \\ \leq C \left(||\phi||_{L^{2}(\Lambda_{\varepsilon})} + ||\mathcal{M}_{\varepsilon}(\phi)||_{L^{2}(\Lambda_{\varepsilon})} \right) ||\psi||_{L^{2}(\Omega)}.$$

The inclusion $\Lambda_{\varepsilon} \subset \widetilde{\Omega}_{\sqrt{n}\varepsilon}$, the fact that $\phi \in L^2_{1/\rho}(\Omega)$ and $(2.8)_1$ - $(2.18)_1$ lead to

$$\int_{\Omega} (\mathcal{M}_{\varepsilon}(\phi) - \phi) \psi \leq C \varepsilon ||\phi/\rho||_{L^{2}(\Omega)} ||\psi||_{\rho}.$$

Hence we get (2.19).

- The scale-splitting operator Q_{ε} .
- * For $\phi \in L^1(\mathbb{R}^n)$, the function $\mathcal{Q}_{\varepsilon}(\phi) \in W^{1,\infty}(\mathbb{R}^n)$ is given by

$$Q_{\varepsilon}(\phi)(x) = \sum_{\xi \in \mathbb{Z}^n} \mathcal{M}_{\varepsilon}(\phi)(\varepsilon \xi) H_{\varepsilon,\xi}(x)$$
 for a.e. $x \in \mathbb{R}^n$,

where

$$H_{\varepsilon,\xi}(x) = H\left(\frac{x - \varepsilon\xi}{\varepsilon}\right) \quad \text{with}$$

$$H(z) = \begin{cases} (1 - |z_1|)(1 - |z_2|) \dots (1 - |z_n|) & \text{if } z \in [-1, 1]^n, \\ 0 & \text{if } z \in \mathbb{R}^n \setminus [-1, 1]^n. \end{cases}$$

Below we recall some results about Q_{ε} proved in [9] and [15].

 \star For $\phi \in L^2(\mathbb{R}^n)$ we have

$$||\mathcal{Q}_{\varepsilon}(\phi)||_{L^{2}(\mathbb{R}^{n})} \leq C||\phi||_{L^{2}(\mathbb{R}^{n})}, \qquad ||\nabla \mathcal{Q}_{\varepsilon}(\phi)||_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} \leq \frac{C}{\varepsilon}||\phi||_{L^{2}(\mathbb{R}^{n})}$$
(2.21)

and

$$Q_{\varepsilon}(\phi) \longrightarrow \phi$$
 strongly in $L^2(\mathbb{R}^n)$.

 \star For $\phi \in H^1(\mathbb{R}^n)$ we have

$$||\nabla \mathcal{Q}_{\varepsilon}(\phi)||_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} \leq C||\nabla \phi||_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})},$$

$$||\phi - \mathcal{Q}_{\varepsilon}(\phi)||_{L^{2}(\mathbb{R}^{n})} \leq C\varepsilon||\nabla \phi||_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})}$$
(2.22)

and

$$Q_{\varepsilon}(\phi) \longrightarrow \phi$$
 strongly in $H^1(\mathbb{R}^n)$. (2.23)

* For $\phi \in L^2(\mathbb{R}^n)$ and $\chi \in L^2(Y)$ we have $\mathcal{Q}_{\varepsilon}(\phi)\chi\left(\left\{\frac{\cdot}{\varepsilon}\right\}\right) \in L^2(\mathbb{R}^n)$, $\nabla \mathcal{Q}_{\varepsilon}(\phi)\chi\left(\left\{\frac{\cdot}{\varepsilon}\right\}\right) \in L^2(\mathbb{R}^n)$ and

$$\|\mathcal{Q}_{\varepsilon}(\phi)\chi\left(\left\{\frac{\cdot}{\varepsilon}\right\}\right)\|_{L^{2}(\mathbb{R}^{n})} \leq C\|\phi\|_{L^{2}(\mathbb{R}^{n})}\|\chi\|_{L^{2}(Y)},$$

$$\|\mathcal{Q}_{\varepsilon}(\phi)\chi\left(\left\{\frac{\cdot}{\varepsilon}\right\}\right)\|_{L^{2}(\widetilde{\Omega}\sqrt{n}\varepsilon)} \leq C\|\phi\|_{L^{2}(\widetilde{\widetilde{\Omega}}_{3\sqrt{n}\varepsilon})}\|\chi\|_{L^{2}(Y)}.$$

$$(2.24)$$

Moreover, if $\phi \in H^1(\mathbb{R}^n)$ then we have

$$\begin{split} & \left\| \left(\mathcal{Q}_{\varepsilon}(\phi) - \mathcal{M}_{\varepsilon}(\phi) \right) \chi \left(\left\{ \frac{\cdot}{\varepsilon} \right\} \right) \right\|_{L^{2}(\mathbb{R}^{n})} \le C \varepsilon \| \nabla \phi \|_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} \| \chi \|_{L^{2}(Y)}, \\ & \left\| \nabla \mathcal{Q}_{\varepsilon}(\phi) \chi \left(\left\{ \frac{\cdot}{\varepsilon} \right\} \right) \right\|_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} \le C \| \nabla \phi \|_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} \| \chi \|_{L^{2}(Y)}, \\ & \left\| \nabla \mathcal{Q}_{\varepsilon}(\phi) \chi \left(\left\{ \frac{\cdot}{\varepsilon} \right\} \right) \right\|_{L^{2}(\widetilde{\Omega}_{\sqrt{n}\varepsilon};\mathbb{R}^{n})} \le C \| \nabla \phi \|_{L^{2}(\widetilde{\widetilde{\Omega}}_{3\sqrt{n}\varepsilon};\mathbb{R}^{n})} \| \chi \|_{L^{2}(Y)}, \end{split}$$
 (2.25)

Lemma 2.7. For $\phi \in H^1_{\rho}(\Omega)$ we have

$$||\rho(\mathcal{Q}_{\varepsilon}(\phi) - \phi)||_{L^{2}(\Omega)} \le C\varepsilon||\phi||_{\rho} \tag{2.26}$$

For $\phi \in H^1_{1/\rho}(\Omega)$ and ϕ_{ε} given by Lemma 2.5 we have

$$\|\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})\|_{1/\rho} \leq C\|\phi\|_{1/\rho}, \qquad \|(\phi - \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon}))/\rho\|_{L^{2}(\Omega)} \leq C\varepsilon\|\phi\|_{1/\rho},$$

$$\forall \mathbf{i} = i_{1}\mathbf{e}_{1} + \ldots + i_{n}\mathbf{e}_{n}, \qquad (i_{1}, \ldots, i_{n}) \in \{0, 1\}^{n}$$

$$\|(\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\cdot + \varepsilon \mathbf{i}) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}))/\rho\|_{L^{2}(\Omega)} \leq C\varepsilon\|\phi\|_{1/\rho}.$$

$$(2.27)$$

For $\phi \in L^2(\mathbb{R}^n)$ and $\chi \in L^2(Y)$

$$\left\| \left(\mathcal{M}_{\varepsilon}(\rho\phi) - \rho \mathcal{M}_{\varepsilon}(\phi) \right) \chi \left(\left\{ \frac{\cdot}{\varepsilon} \right\} \right) \right\|_{L^{2}(\mathbb{R}^{n})} \le C \varepsilon \|\phi\|_{L^{2}(\mathbb{R}^{n})} \|\chi\|_{L^{2}(Y)}. \tag{2.28}$$

For $\phi \in H^1_\rho(\Omega)$ and $\chi \in L^2(Y)$

$$\|\rho(\mathcal{Q}_{\varepsilon}(\phi) - \mathcal{M}_{\varepsilon}(\phi))\chi(\left\{\frac{\cdot}{\varepsilon}\right\})\|_{L^{2}(\Omega)} \leq C\varepsilon \|\phi\|_{\rho} \|\chi\|_{L^{2}(Y)},$$

$$\|\rho\nabla\mathcal{Q}_{\varepsilon}(\phi)\chi(\left\{\frac{\cdot}{\varepsilon}\right\})\|_{L^{2}(\Omega)} \leq C\|\phi\|_{\rho} \|\chi\|_{L^{2}(Y)}.$$
(2.29)

The constants do not depend on ε .

Proof. Step 1. Let ϕ be in $H^1_{\rho}(\Omega)$. We first prove

$$||\rho(\mathcal{Q}_{\varepsilon}(\phi) - \mathcal{M}_{\varepsilon}(\phi))||_{L^{2}(\Omega)} \le C\varepsilon||\phi||_{\rho}.$$
 (2.30)

To do that, we proceed as in the proof of $(2.18)_1$. Let $\varepsilon(\xi + Y)$ be a cell included in Ω . Case 1: $\rho(\varepsilon\xi) \geq 3\sqrt{n}\varepsilon$. In this case we have

$$1 \le \frac{\max_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}}{\min_{z \in \varepsilon(\xi+2Y)} \{\rho(z)\}} \le 4 \quad \text{and} \quad 1 \le \frac{\max_{z \in \varepsilon(\xi+2Y)} \{\rho(z)\}}{\min_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}} \le \frac{5}{2}.$$

By definition of $Q_{\varepsilon}(\phi)$ we deduce that

$$\int_{\varepsilon(\xi+Y)} [\rho(x)]^{2} |\mathcal{Q}_{\varepsilon}(\phi)(x) - \mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi)|^{2} dx \leq \left[\max_{z \in \varepsilon(\xi+Y)} {\{\rho(z)\}}\right]^{2} \int_{\varepsilon(\xi+Y)} |\mathcal{Q}_{\varepsilon}(\phi)(x) - \mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi)|^{2} dx
\leq \left[\max_{z \in \varepsilon(\xi+Y)} {\{\rho(z)\}}\right]^{2} C\varepsilon^{2} \int_{\varepsilon(\xi+2Y)} |\nabla\phi(x)|^{2} dx
\leq C\varepsilon^{2} \int_{\varepsilon(\xi+2Y)} [\rho(x)]^{2} |\nabla\phi(x)|^{2} dx.$$

Case 2: $\rho(\varepsilon\xi) \leq 3\sqrt{n\varepsilon}$. Then again by definition of $\mathcal{Q}_{\varepsilon}(\phi)$ we get

$$\int_{\varepsilon(\xi+Y)} [\rho(x)]^2 |\mathcal{Q}_{\varepsilon}(\phi)(x) - \mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi)|^2 dx \le C\varepsilon^2 \int_{\varepsilon(\xi+2Y)} |\phi(x)|^2 dx.$$

As a consequence of the cases 1 and 2 we get

$$\int_{\widehat{\Omega}_{\varepsilon}} [\rho(x)]^2 |\mathcal{Q}_{\varepsilon}(\phi)(x) - \mathcal{M}_{\varepsilon}(\phi)(x)|^2 dx \le C\varepsilon^2 \int_{\Omega} ([\rho(x)]^2 |\nabla \phi(x)|^2 + |\phi(x)|^2) dx. \quad (2.31)$$

Furthermore we have

$$\int_{\Lambda_{\varepsilon}} [\rho(x)]^2 |\mathcal{Q}_{\varepsilon}(\phi)(x)|^2 dx \le C\varepsilon^2 \int_{\Lambda_{\varepsilon}} |\mathcal{Q}_{\varepsilon}(\phi)(x)|^2 dx \le C\varepsilon^2 \int_{\Omega} |\phi(x)|^2 dx$$

which with (2.31) lead to (2.30). Then as a consequence of $(2.18)_1$ and (2.30) we get (2.26).

Step 2. We prove $(2.27)_1$. Let ϕ be in $H^1_{1/\rho}(\Omega)$ and ϕ_{ε} given by Lemma 2.5. Due to the fact that $\phi_{\varepsilon}(x) = 0$ for a.e. $x \in \mathbb{R}^n \setminus \overline{\Omega}_{6\sqrt{n}\varepsilon}$, hence $\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) = 0$ for all $x \in \Omega$ such that $\rho(x) \leq 4\sqrt{n}\varepsilon$. Again we take a cell $\varepsilon(\xi + Y)$ included in Ω such that $\rho(\varepsilon\xi) \geq 3\sqrt{n}\varepsilon$. The values taken by $\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})$ in the cell $\varepsilon(\xi + Y)$ depend only on the values of ϕ_{ε} in $\varepsilon(\xi + 2Y)$. Then we have

$$\begin{split} &\int_{\varepsilon(\xi+Y)} \frac{1}{[\rho(x)]^2} |\nabla \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x)|^2 dx \leq \frac{C}{[\min_{x \in \varepsilon(\xi+Y)} \{\rho(x)\}]^2} \int_{\varepsilon(\xi+2Y)} |\nabla \phi_{\varepsilon}(x)|^2 dx \\ &\leq C \frac{[\max_{x \in \varepsilon(\xi+2Y)} \{\rho(x)\}]^2}{[\min_{x \in \varepsilon(\xi+Y)} \{\rho(x)\}]^2} \int_{\varepsilon(\xi+2Y)} \frac{1}{[\rho(x)]^2} |\nabla \phi_{\varepsilon}(x)|^2 dx \leq C \int_{\varepsilon(\xi+2Y)} \frac{1}{[\rho(x)]^2} |\nabla \phi_{\varepsilon}(x)|^2 dx. \end{split}$$

Adding all these inequalities gives

$$\int_{\widetilde{\Omega}_{4\sqrt{\rho}\varepsilon}} \frac{1}{[\rho(x)]^2} |\nabla \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x)|^2 dx \le C \int_{\Omega} \frac{1}{[\rho(x)]^2} |\nabla \phi_{\varepsilon}(x)|^2 dx$$

Since $Q_{\varepsilon}(\phi_{\varepsilon})(x) = 0$ for every $x \in \Omega$ such that $\rho(x) \leq 4\sqrt{n\varepsilon}$, we get $||Q_{\varepsilon}(\phi_{\varepsilon})||_{1/\rho} \leq C||\phi_{\varepsilon}||_{1/\rho}$. We conclude using $(2.10)_2$.

Step 3. Now we prove $(2.27)_2$. Again we consider a cell $\varepsilon(\xi + Y)$ included in Ω such that $\rho(\varepsilon\xi) \geq 3\sqrt{n}\varepsilon$. We have

$$\int_{\varepsilon(\xi+Y)} \frac{1}{[\rho(x)]^2} |\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) - \phi_{\varepsilon}(x)|^2 dx \leq \frac{C}{[\min_{x \in \varepsilon(\xi+Y)} {\{\rho(x)\}}]^2} \int_{\varepsilon(\xi+Y)} |\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) - \phi_{\varepsilon}(x)|^2 dx
\leq \frac{C}{[\min_{x \in \varepsilon(\xi+Y)} {\{\rho(x)\}}]^2} \sum_{\mathbf{i} \in {\{0,1\}}^n} \int_{\varepsilon(\xi+\mathbf{i}+Y)} |\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon\xi + \varepsilon\mathbf{i}) - \phi_{\varepsilon}(x)|^2 dx
\leq C\varepsilon^2 \frac{[\max_{z \in \varepsilon(\xi+2Y)} {\{\rho(z)\}}]^2}{[\min_{z \in \varepsilon(\xi+Y)} {\{\rho(z)\}}]^2} \int_{\varepsilon(\xi+2Y)} \frac{1}{[\rho(x)]^2} |\nabla \phi_{\varepsilon}(x)|^2 dx \leq C\varepsilon^2 \int_{\varepsilon(\xi+2Y)} \frac{1}{[\rho(x)]^2} |\nabla \phi_{\varepsilon}(x)|^2 dx.$$

Hence we get

$$\int_{\widetilde{\Omega}_{4,\sqrt{n}\varepsilon}} \frac{1}{[\rho(x)]^2} |\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) - \phi_{\varepsilon}(x)|^2 dx \le C\varepsilon^2 \int_{\Omega} \frac{1}{[\rho(x)]^2} |\nabla \phi_{\varepsilon}(x)|^2 dx$$

The above estimate and the fact that $Q_{\varepsilon}(\phi_{\varepsilon})(x) - \phi_{\varepsilon}(x) = 0$ for a.e. $x \in \Omega$ such that $\rho(x) \leq 4\sqrt{n\varepsilon}$ yield $||(\phi_{\varepsilon} - Q_{\varepsilon}(\phi_{\varepsilon}))/\rho||_{L^{2}(\Omega)} \leq C\varepsilon||\phi_{\varepsilon}||_{1/\rho}$. We conclude using both estimates in (2.10).

Proceeding as in the Steps 2 and 3 we obtain $(2.27)_3$, (2.28) and (2.29).

3 Two new projection theorems

Theorem 3.1. Let ϕ be in $H^1_{1/\rho}(\Omega)$. There exists $\widehat{\psi}_{\varepsilon} \in H^1_{per}(Y; L^2(\Omega))$ such that

$$\begin{cases}
||\widehat{\psi}_{\varepsilon}||_{H^{1}(Y;L^{2}(\Omega))} \leq C\{||\phi||_{L^{2}(\Omega)} + \varepsilon||\nabla\phi||_{[L^{2}(\Omega)]^{n}}\} \\
||\mathcal{T}_{\varepsilon}(\phi) - \widehat{\psi}_{\varepsilon}||_{H^{1}(Y;(H^{1}_{\rho}(\Omega))')} \leq C\varepsilon(||\phi/\rho||_{L^{2}(\Omega)} + \varepsilon||\phi||_{1/\rho}).
\end{cases}$$
(3.1)

The constants depend only on n and $\partial\Omega$.

Proof. Here, we proceed as in the proof of the Proposition 3.3 of [5]. We first reintroduce the open sets $\widehat{\Omega}_{\varepsilon,i}$ and the unfolding operators $\mathcal{T}_{\varepsilon,i}$. We set

$$\widehat{\Omega}_{\varepsilon,i} = \widehat{\Omega}_{\varepsilon} \cap (\widehat{\Omega}_{\varepsilon} - \varepsilon \mathbf{e}_i), \qquad K_i = \operatorname{interior}(\overline{Y} \cup (\mathbf{e}_i + \overline{Y})), \quad i \in \{1, \dots, n\}.$$

The unfolding operator $\mathcal{T}_{\varepsilon,i}$ from $L^2(\Omega)$ into $L^2(\Omega \times K_i)$ is defined by

$$\forall \psi \in L^{2}(\Omega), \qquad \mathcal{T}_{\varepsilon,i}(\psi)(x,y) = \begin{cases} \psi\left(\varepsilon\left[\frac{x}{\varepsilon}\right]_{Y} + \varepsilon y\right) & \text{for } x \in \widehat{\Omega}_{\varepsilon,i} \text{ and for a.e. } y \in K_{i}, \\ 0 & \text{for } x \in \Omega \setminus \overline{\widehat{\Omega}}_{\varepsilon,i} \text{ and for a.e. } y \in K_{i}. \end{cases}$$

The restriction of $\mathcal{T}_{\varepsilon,i}(\psi)$ to $\widehat{\Omega}_{\varepsilon,i} \times Y$ is equal to $\mathcal{T}_{\varepsilon}(\psi)$.

Step 1. Let us first take $\phi \in L^2_{1/\rho}(\Omega)$. We set $\psi = \frac{1}{\rho}\phi$ and we evaluate the difference $\mathcal{T}_{\varepsilon,i}(\phi)(.,..+\mathbf{e}_i) - \mathcal{T}_{\varepsilon,i}(\phi)$ in $L^2(Y;(H^1_{\rho}(\Omega))')$. For any $\Psi \in H^1_{\rho}(\Omega)$ a change of variables gives for a. e. $y \in Y$

$$\begin{split} \int_{\Omega} \mathcal{T}_{\varepsilon,i}(\phi)(x,y+\mathbf{e}_i)\Psi(x)dx &= \int_{\widehat{\Omega}_{\varepsilon,i}} \mathcal{T}_{\varepsilon}(\phi)(x+\varepsilon\mathbf{e}_i,y)\Psi(x)dx \\ &= \int_{\widehat{\Omega}_{\varepsilon,i}+\varepsilon\mathbf{e}_i} \mathcal{T}_{\varepsilon}(\phi)(x,y)\Psi(x-\varepsilon\mathbf{e}_i)dx. \end{split}$$

Then we obtain for a. e. $y \in Y$

$$\left| \int_{\Omega} \left\{ \mathcal{T}_{\varepsilon,i}(\phi)(.,y + \mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\phi)(.,y) \right\} \Psi - \int_{\widehat{\Omega}_{\varepsilon,i}} \mathcal{T}_{\varepsilon}(\psi)(.,y) \rho \left\{ \Psi(. - \varepsilon \mathbf{e}_{i}) - \Psi \right\} \right| \\
\leq \left| \int_{\widehat{\Omega}_{\varepsilon,i}} \mathcal{T}_{\varepsilon}(\psi)(.,y) \left(\mathcal{T}_{\varepsilon}(\rho) - \rho \right) \left\{ \Psi(. - \varepsilon \mathbf{e}_{i}) - \Psi \right\} \right| + C ||\mathcal{T}_{\varepsilon}(\phi)(.,y)||_{L^{2}(\widetilde{\Omega}_{2\sqrt{n}\varepsilon})} ||\Psi||_{L^{2}(\widetilde{\Omega}_{2\sqrt{n}\varepsilon})}.$$

From $(2.18)_2$ we obtain

$$||\rho(\Psi(.-\varepsilon\mathbf{e}_i)-\Psi)||_{L^2(\widehat{\Omega}_{\varepsilon,i})} \le C\varepsilon||\Psi||_{\rho} \quad \forall i \in \{1,\ldots,n\}.$$

We have

$$||\mathcal{T}_{\varepsilon}(\rho) - \rho||_{L^{\infty}(\Omega)} \le C\varepsilon.$$
 (3.2)

The above inequalities lead to

$$\langle \mathcal{T}_{\varepsilon,i}(\phi)(.,y+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\phi)(.,y), \Psi \rangle_{(H_{\rho}^{1}(\Omega))',H_{\rho}^{1}(\Omega)}$$

$$= \int_{\Omega} \{ \mathcal{T}_{\varepsilon,i}(\phi)(.,y+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\phi)(.,y) \} \Psi$$

$$\leq C\varepsilon ||\Psi||_{\rho} ||\mathcal{T}_{\varepsilon}(\psi)(.,y)||_{L^{2}(\Omega)} + C\varepsilon ||\Psi||_{L^{2}(\Omega)} ||\mathcal{T}_{\varepsilon}(\psi)(.,y)||_{L^{2}(\Omega)}$$

$$+ C||\mathcal{T}_{\varepsilon}(\phi)(.,y)||_{L^{2}(\widetilde{\Omega}_{2},\sqrt{\eta_{\varepsilon}})} ||\Psi||_{L^{2}(\widetilde{\Omega}_{2},\sqrt{\eta_{\varepsilon}})}.$$

Therefore, for a.e. $y \in Y$ we have

$$||\mathcal{T}_{\varepsilon,i}(\phi)(.,y+\mathbf{e}_i)-\mathcal{T}_{\varepsilon,i}(\phi)(.,y)||_{(H^1_\rho(\Omega))'}\leq C\varepsilon||\mathcal{T}_{\varepsilon}(\psi)(.,y)||_{L^2(\Omega)}+C||\mathcal{T}_{\varepsilon}(\phi)(.,y)||_{L^2(\widetilde{\Omega}_{2\sqrt{\eta_{\varepsilon}}})}$$

which leads to the following estimate of the difference between $\mathcal{T}_{\varepsilon,i}(\phi)_{|_{\Omega\times Y}}$ and one of its translated:

$$||\mathcal{T}_{\varepsilon,i}(\phi)(.,..+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\phi)||_{L^{2}(Y;(H^{1}_{\rho}(\Omega))')} \leq C\varepsilon||\phi/\rho||_{L^{2}(\Omega)} + C||\phi||_{L^{2}(\widetilde{\Omega}_{2\sqrt{n}\varepsilon})}$$

$$\leq C\varepsilon||\phi/\rho||_{L^{2}(\Omega)}.$$
(3.3)

The constant depends only on the boundary of Ω .

Step 2. Let $\phi \in H^1_{1/\rho}(\Omega)$. The above estimate (3.3) applied to ϕ and its partial derivatives give

$$||\mathcal{T}_{\varepsilon,i}(\phi)(.,..+\mathbf{e}_i) - \mathcal{T}_{\varepsilon,i}(\phi)||_{L^2(Y;(H^1_{\rho}(\Omega))')} \le C\varepsilon||\phi/\rho||_{L^2(\Omega)}$$
$$||\mathcal{T}_{\varepsilon,i}(\nabla\phi)(.,..+\mathbf{e}_i) - \mathcal{T}_{\varepsilon,i}(\nabla\phi)||_{[L^2(Y;(H^1_{\sigma}(\Omega))']^n)} \le C\varepsilon||\phi||_{1/\rho}.$$

which in turn lead to (we recall that $\nabla_y (\mathcal{T}_{\varepsilon,i}(\phi)) = \varepsilon \mathcal{T}_{\varepsilon,i}(\nabla \phi)$).

$$||\mathcal{T}_{\varepsilon,i}(\phi)(.,..+\mathbf{e}_i) - \mathcal{T}_{\varepsilon,i}(\phi)||_{H^1(Y;(H^1_0(\Omega))')} \le C\varepsilon (||\phi/\rho||_{L^2(\Omega)} + \varepsilon||\phi||_{1/\rho}).$$

From these inequalities, for any $i \in \{1, ..., n\}$, we deduce the estimate of the difference of the traces of $y \longrightarrow \mathcal{T}_{\varepsilon}(\phi)(., y)$ on the faces Y_i and $\mathbf{e}_i + Y_i$

$$||\mathcal{T}_{\varepsilon}(\phi)(.,..+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon}(\phi)||_{H^{1/2}(Y_{i};(H^{1}_{\varrho}(\Omega))')} \le C\varepsilon(||\phi/\rho||_{L^{2}(\Omega)} + \varepsilon||\phi||_{1/\varrho}). \tag{3.4}$$

It measures the periodic defect of $y \longrightarrow \mathcal{T}_{\varepsilon}(\phi)(.,y)$. We decompose $\mathcal{T}_{\varepsilon}(\phi)$ into the sum of an element belonging to $H^1_{per}(Y;L^2(\Omega))$ and an element belonging to $(H^1(Y;L^2(\Omega)))^{\perp}$ (the orthogonal of $H^1_{per}(Y;L^2(\Omega))$ in $H^1(Y;L^2(\Omega))$, see [5])

$$\mathcal{T}_{\varepsilon}(\phi) = \widehat{\psi}_{\varepsilon} + \overline{\phi}_{\varepsilon}, \qquad \widehat{\psi}_{\varepsilon} \in H^{1}_{per}(Y; L^{2}(\Omega)), \qquad \overline{\phi}_{\varepsilon} \in (H^{1}(Y; L^{2}(\Omega)))^{\perp}. \tag{3.5}$$

The function $y \longrightarrow \mathcal{T}_{\varepsilon}(\phi)(.,y)$ takes its values in a finite dimensional space,

$$\overline{\phi}_{\varepsilon}(.,..) = \sum_{\xi \in \Xi_{\varepsilon}} \overline{\phi}_{\varepsilon,\xi}(..) \chi_{\xi}(.)$$

where $\chi_{\xi}(.)$ is the characteristic function of the cell $\varepsilon(\xi + Y)$ and where $\overline{\phi}_{\varepsilon,\xi}(..) \in (H^1(Y))^{\perp}$ (the orthogonal of $H^1_{per}(Y)$ in $H^1(Y)$, see [5]). The decomposing (3.5) is the same in $H^1(Y; (H^1_{\rho}(\Omega))')$ and we have

$$||\widehat{\psi}_{\varepsilon}||_{H^{1}(Y;L^{2}(\Omega))}^{2} + ||\overline{\phi}_{\varepsilon}||_{H^{1}(Y;L^{2}(\Omega))}^{2} = ||\mathcal{T}_{\varepsilon}(\phi)||_{H^{1}(Y;L^{2}(\Omega))}^{2} \le C\{||\phi||_{L^{2}(\Omega)} + \varepsilon||\nabla\phi||_{[L^{2}(\Omega)]^{n}}\}^{2}$$

which is the first inequality in (3.1) and the estimate of $\overline{\phi}_{\varepsilon}$ in $H^1(Y; L^2(\Omega))$. From Theorem 2.2 of [5] and (3.4) we obtain a finer estimate of $\overline{\phi}_{\varepsilon}$ in $H^1(Y; (H^1_{\rho}(\Omega))')$

$$||\overline{\phi}_{\varepsilon}||_{H^{1}(Y;(H^{1}_{\rho}(\Omega))')} \leq C\varepsilon (||\phi/\rho||_{L^{2}(\Omega)} + \varepsilon ||\phi||_{1/\rho}).$$

It is the second inequality in (3.1).

Theorem 3.2. For any $\phi \in H^1_{1/\rho}(\Omega)$, there exists $\widehat{\phi}_{\varepsilon} \in H^1_{per}(Y; L^2(\Omega))$ such that

$$||\widehat{\phi}_{\varepsilon}||_{H^{1}(Y;L^{2}(\Omega))} \leq C||\nabla\phi||_{[L^{2}(\Omega)]^{n}},$$

$$||\mathcal{T}_{\varepsilon}(\nabla\phi) - \nabla\phi - \nabla_{y}\widehat{\phi}_{\varepsilon}||_{[L^{2}(Y;(H^{1}_{\rho}(\Omega))')]^{n}} \leq C\varepsilon||\phi||_{1/\rho}.$$
(3.6)

The constants depend only on $\partial\Omega$.

Proof. Let ϕ be in $H^1_{1/\rho}(\Omega)$ and $\psi = \phi/\rho \in H^1_0(\Omega)$. The function ϕ is extended by 0 outside of Ω . We decompose ϕ as

$$\phi = \Phi + \varepsilon \underline{\phi}$$
, where $\Phi = \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})$ and $\underline{\phi} = \frac{1}{\varepsilon} \Big(\phi - \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon}) \Big)$

where ϕ_{ε} is given by Lemma 2.5. We have Φ and $\underline{\phi} \in H_0^1(\Omega)$ and due to (2.27) we get the following estimates:

$$||\Phi||_{1/\rho} + \varepsilon ||\phi||_{1/\rho} + ||\phi/\rho||_{L^2(\Omega)} \le C||\phi||_{1/\rho}. \tag{3.7}$$

The projection Theorem 3.1 applied to $\underline{\phi} \in H^1_{1/\rho}(\Omega)$ gives an element $\widehat{\phi}_{\varepsilon}$ in $H^1_{per}(Y; L^2(\Omega))$ such that

$$||\widehat{\phi}_{\varepsilon}||_{H^{1}(Y;L^{2}(\Omega))} \leq C||\phi||_{1/\rho},$$

$$||\mathcal{T}_{\varepsilon}(\underline{\phi}) - \widehat{\phi}_{\varepsilon}||_{H^{1}(Y;(H^{1}_{\rho}(\Omega))')} \leq C\varepsilon||\phi||_{1/\rho}.$$
(3.8)

Now we evaluate $||\mathcal{T}_{\varepsilon}(\nabla\Phi) - \nabla\Phi||_{[L^{2}(Y;(H_{\sigma}^{1}(\Omega))')]^{n}}$.

From (2.19), $(2.27)_1$ and (3.7)we get

$$\|\nabla\Phi - \mathcal{M}_{\varepsilon}(\nabla\Phi)\|_{(H_{\rho}^{1}(\Omega;\mathbb{R}^{n}))'} \le C\varepsilon\|\phi\|_{1/\rho}.$$
(3.9)

We set

$$H^{(1)}(z) = \begin{cases} (1 - |z_2|)(1 - |z_3|) \dots (1 - |z_n|) & \text{if } z = (z_1, z_2, \dots, z_n) \in [-1, 1]^n, \\ 0 & \text{if } z \in \mathbb{R}^n \setminus [-1, 1]^n. \end{cases}$$
$$\mathbf{I} = \left\{ \mathbf{i} \mid \mathbf{i} = i_2 \mathbf{e}_2 + \dots + i_n \mathbf{e}_n, \quad (i_2, \dots, i_n) \in \{0, 1\}^{n-1} \right\}$$

For $\xi \in \mathbb{Z}^n$ and for every $(x, y) \in \varepsilon(\xi + Y) \times Y$ we have

$$\mathcal{T}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{1}} \right) (x, y) = \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon} (\phi_{\varepsilon}) \left(\varepsilon (\xi + \mathbf{e}_{1} + \mathbf{i}) \right) - \mathcal{M}_{\varepsilon} (\phi_{\varepsilon}) \left(\varepsilon (\xi + \mathbf{i}) \right)}{\varepsilon} H^{(1)} (y - \mathbf{i})$$

$$\mathcal{M}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{1}} \right) (\varepsilon \xi) = \frac{1}{2^{n-1}} \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon} (\phi_{\varepsilon}) \left(\varepsilon (\xi + \mathbf{e}_{1} + \mathbf{i}) \right) - \mathcal{M}_{\varepsilon} (\phi_{\varepsilon}) \left(\varepsilon (\xi + \mathbf{i}) \right)}{\varepsilon}.$$

Now, let us take $\psi \in H^1_{\rho}(\Omega)$. We recall that $\phi_{\varepsilon}(x) = 0$ for a.e. $x \in \mathbb{R}^n \setminus \overline{\widetilde{\Omega}}_{6\sqrt{n}\varepsilon}$, hence $\Phi(x) = 0$ for all $x \in \mathbb{R}^n \setminus \overline{\widetilde{\Omega}}_{3\sqrt{n}\varepsilon}$. As a first consequence $\mathcal{M}_{\varepsilon}\left(\frac{\partial \Phi}{\partial x_1}\right)$ in Λ_{ε} . For $y \in Y$ we have

$$<\mathcal{T}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)(.,y) - \mathcal{M}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big), \psi>_{(H_{\rho}^{1}(\Omega))',H_{\rho}^{1}(\Omega)} = \int_{\Omega}\Big\{\mathcal{T}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)(.,y) - \mathcal{M}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)\Big\}\psi$$

$$= \int_{\widehat{\Omega}_{\varepsilon}}\Big\{\mathcal{T}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)(.,y) - \mathcal{M}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)\Big\}\mathcal{M}_{\varepsilon}(\psi).$$

We have

$$\int_{\widehat{\Omega}_{\varepsilon}} \mathcal{M}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{1}} \right) \mathcal{M}_{\varepsilon}(\psi) = \varepsilon^{n} \sum_{\xi \in \mathbb{Z}^{n}} \mathcal{M}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{1}} \right) (\varepsilon \xi) \mathcal{M}_{\varepsilon}(\psi) (\varepsilon \xi)
= \frac{\varepsilon^{n}}{2^{n-1}} \sum_{\xi \in \mathbb{Z}^{n}} \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \left(\varepsilon (\xi + \mathbf{e}_{1} + \mathbf{i}) \right) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \left(\varepsilon (\xi + \mathbf{i}) \right)}{\varepsilon} \mathcal{M}_{\varepsilon}(\psi) (\varepsilon \xi)
= \frac{\varepsilon^{n}}{2^{n-1}} \sum_{\xi \in \mathbb{Z}^{n}} \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon}(\psi) \left(\varepsilon (\xi - \mathbf{e}_{1}) \right) - \mathcal{M}_{\varepsilon}(\psi) \left(\varepsilon \xi \right)}{\varepsilon} \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) (\varepsilon (\xi + \mathbf{i}))$$

and

$$\int_{\widehat{\Omega}_{\varepsilon}} \mathcal{T}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{1}} \right) (., y) \mathcal{M}_{\varepsilon}(\psi)
= \varepsilon^{n} \sum_{\xi \in \mathbb{Z}^{n}} \sum_{\mathbf{i} \in \mathbf{I}} \left[\frac{\mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \left(\varepsilon (\xi + \mathbf{e}_{1} + \mathbf{i}) \right) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \left(\varepsilon (\xi + \mathbf{i}) \right)}{\varepsilon} \right] H^{(1)}(y - \mathbf{i}) \mathcal{M}_{\varepsilon}(\psi) (\varepsilon \xi)
= \varepsilon^{n} \sum_{\xi \in \mathbb{Z}^{n}} \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon}(\psi) \left(\varepsilon (\xi - \mathbf{e}_{1}) \right) - \mathcal{M}_{\varepsilon}(\psi) \left(\varepsilon \xi \right)}{\varepsilon} H^{(1)}(y - \mathbf{i}) \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) (\varepsilon (\xi + \mathbf{i}))$$

Due to the fact that $\phi_{\varepsilon}(x) = 0$ for a.e. $x \in \mathbb{R}^n \setminus \overline{\widetilde{\Omega}}_{6\sqrt{n}\varepsilon}$, in the above summations we only need to take $\xi \in \Xi_{\varepsilon}$ satisfying $\rho(\varepsilon\xi) \geq 3\sqrt{n}\varepsilon$. Hence

$$<\mathcal{T}_{\varepsilon}\left(\frac{\partial\Phi}{\partial x_{1}}\right)(.,y)-\mathcal{M}_{\varepsilon}\left(\frac{\partial\Phi}{\partial x_{1}}\right),\psi>_{(H_{\rho}^{1}(\Omega))',H_{\rho}^{1}(\Omega)}$$

$$=\varepsilon^{n}\sum_{\xi\in\mathbb{Z}^{n}}\frac{\mathcal{M}_{\varepsilon}(\psi)\left(\varepsilon(\xi-\mathbf{e}_{1})\right)-\mathcal{M}_{\varepsilon}(\psi)\left(\varepsilon\xi\right)}{\varepsilon}\sum_{\mathbf{i}\in\mathbf{I}}\left[H^{(1)}(y-\mathbf{i})-\frac{1}{2^{n-1}}\right]\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon(\xi+\mathbf{i})).$$

Since
$$\sum_{\mathbf{i}\in\mathbf{I}} \left[H^{(1)}(y-\mathbf{i}) - \frac{1}{2^{n-1}} \right] = 0$$
 we obtain that

$$\Big| \sum_{\mathbf{i} \in \mathbf{I}} \Big[H^{(1)}(y - \mathbf{i}) - \frac{1}{2^{n-1}} \Big] \mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon(\xi + \mathbf{i})) \Big| \leq \sum_{\mathbf{i} \in \mathbf{I}} \Big| \mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon(\xi + \mathbf{i})) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon\xi) \Big|.$$

Taking into account the above equality and inequality we deduce that

$$< \mathcal{T}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{1}} \right) (., y) - \mathcal{M}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{1}} \right), \psi >_{(H_{\rho}^{1}(\Omega))', H_{\rho}^{1}(\Omega)} \\
= \varepsilon^{n} \sum_{\xi \in \mathbb{Z}^{n}} \sum_{\mathbf{i} \in \mathbf{I}} \left| \frac{\mathcal{M}_{\varepsilon}(\psi) \left(\varepsilon(\xi - \mathbf{e}_{1}) \right) - \mathcal{M}_{\varepsilon}(\psi) \left(\varepsilon \xi \right)}{\varepsilon} \right| \left| \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) (\varepsilon(\xi + \mathbf{i})) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) (\varepsilon \xi) \right| \\
= \frac{1}{\varepsilon} \sum_{\mathbf{i} \in \mathbf{I}} \int_{\Omega} \left| \mathcal{M}_{\varepsilon}(\psi) (\cdot - \varepsilon \mathbf{e}_{1}) - \mathcal{M}_{\varepsilon}(\psi) \right| \left| \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) (\cdot + \varepsilon \mathbf{i}) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \right| \\
\leq \frac{C}{\varepsilon} \sum_{\mathbf{i} \in \mathbf{I}} \left\| \rho \left(\mathcal{M}_{\varepsilon}(\psi) (\cdot - \varepsilon \mathbf{e}_{1}) - \mathcal{M}_{\varepsilon}(\psi) \right) \right\|_{L^{2}(\Omega)} \left\| \frac{1}{\rho} \left(\mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) (\cdot + \varepsilon \mathbf{i}) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \right) \right\|_{L^{2}(\Omega)}.$$

Thanks to $(2.18)_3$ and $(2.27)_3$ we finally get

$$<\mathcal{T}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)(.,y)-\mathcal{M}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big),\psi>_{(H_{\rho}^{1}(\Omega))',H_{\rho}^{1}(\Omega)}\leq C\varepsilon||\phi_{\varepsilon}||_{1/\rho}||\psi||_{\rho}$$

which leads to

$$\left\| \mathcal{T}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{1}} \right) - \mathcal{M}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{1}} \right) \right\|_{L^{\infty}(Y; (H_{\sigma}^{1}(\Omega))')} \leq C \varepsilon ||\phi_{\varepsilon}||_{1/\rho}. \tag{3.10}$$

Besides we have

$$\int_{\Omega} \frac{\partial \underline{\phi}}{\partial x_1}(x)\psi(x)dx = -\int_{\Omega} \underline{\phi}(x)\frac{\partial \psi}{\partial x_1}(x)dx \le C||\underline{\phi}/\rho||_{L^2(\Omega)}||\psi||_{\rho} \le C||\phi||_{1/\rho}||\psi||_{\rho}.$$

Hence $\left\| \varepsilon \frac{\partial \underline{\phi}}{\partial x_1} \right\|_{(H^1_{\rho}(\Omega;\mathbb{R}^n))'} \leq C\varepsilon ||\phi||_{1/\rho}$. This last estimate with $(2.10)_2$, (3.9) and (3.10) yield

$$\left\| \mathcal{T}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_1} \right) - \frac{\partial \phi}{\partial x_1} \right\|_{L^{\infty}(Y; (H_0^1(\Omega))')} \leq C \varepsilon ||\phi_{\varepsilon}||_{1/\rho}.$$

Proceeding in the same way we obtain the same estimates for the partial derivatives with respect to x_i , $i \in \{2, ..., n\}$. Hence we get $\|\mathcal{T}_{\varepsilon}(\nabla \Phi) - \nabla \phi\|_{[L^{\infty}(Y;(H^{1}_{\rho}(\Omega))')]^{n}} \leq C\varepsilon ||\phi_{\varepsilon}||_{1/\rho}$. Then, thanks to (3.8) the second estimate in (3.6) is proved.

4 Recalls about the classical periodic homogenization problem

We consider the homogenization problem

$$\phi^{\varepsilon} \in H_0^1(\Omega), \qquad \int_{\Omega} A_{\varepsilon}(x) \nabla \phi^{\varepsilon}(x) \nabla \psi(x) dx = \int_{\Omega} f(x) \psi(x) dx, \qquad \forall \psi \in H_0^1(\Omega), \quad (4.1)$$

where

- $A_{\varepsilon}(x) = A\left(\left\{\frac{x}{\varepsilon}\right\}\right)$ for a.e. $x \in \Omega$, here A is a square matrix belonging to $L^{\infty}(Y; \mathbb{R}^{n \times n})$, satisfying the condition of uniform ellipticity $c|\xi|^2 \leq A(y)\xi \cdot \xi \leq C|\xi|^2$ for a.e. $y \in Y$, with c and C strictly positive constants,
 - $f \in L^2(\Omega)$.

We showed in [9] that

$$\mathcal{T}_{\varepsilon}(\nabla \phi^{\varepsilon}) \longrightarrow \nabla \Phi + \nabla_y \widehat{\phi}$$
 strongly in $L^2(\Omega \times Y; \mathbb{R}^n)$

where $(\Phi, \widehat{\phi}) \in H_0^1(\Omega) \times L^2(\Omega; H_{per}^1(Y))$ is the solution of the problem of unfolding homogenization

$$\begin{split} \forall (\Psi, \widehat{\psi}) &\in H^1_0(\Omega) \times L^2(\Omega; H^1_{per}(Y)) \\ &\int_{\Omega} \int_Y A(y) \big\{ \nabla \Phi(x) + \nabla_y \widehat{\phi}(x,y) \big\} \, \big\{ \nabla \Psi(x) + \nabla_y \widehat{\psi}(x,y) \big\} dx dy = \int_{\Omega} f(x) \Psi(x) dx. \end{split}$$

The correctors χ_i , $i \in \{1, ..., n\}$, are the solutions of the following variational problems:

$$\chi_i \in H^1_{per}(Y), \qquad \int_Y \chi_i = 0,$$

$$\int_Y A(y) \nabla_y (\chi_i(y) + y_i) \nabla_y \psi(y) dy = 0, \qquad \forall \psi \in H^1_{per}(Y).$$
(4.2)

They allow to express $\widehat{\phi}$ in terms of the partial derivatives of Φ

$$\widehat{\phi} = \sum_{i=1}^{n} \frac{\partial \Phi}{\partial x_i} \chi_i \tag{4.3}$$

and to give the homogenized problem verified by Φ

$$\Phi \in H_0^1(\Omega), \qquad \int_{\Omega} \mathcal{A} \nabla \Phi(x) \nabla \Psi(x) dx = \int_{\Omega} f(x) \Psi(x) dx, \qquad \forall \Psi \in H_0^1(\Omega)$$
 (4.4)

where (see [9])

$$\mathcal{A}_{ij} = \sum_{k,l=1}^{n} \int_{Y} a_{kl}(y) \frac{\partial (y_j + \chi_j(y))}{\partial y_l} \frac{\partial (y_i + \chi_i(y))}{\partial y_k} dy. \tag{4.5}$$

5 Error estimates with a Dirichlet condition

Theorem 5.1. Let $(\phi^{\varepsilon})_{\varepsilon>0}$ be a sequence of functions belonging to $H^1(\Omega)$ such that

$$div(A_{\varepsilon}\nabla\phi^{\varepsilon}) = 0 \qquad in \quad \Omega. \tag{5.1}$$

Setting $g_{\varepsilon} = \phi_{|\partial\Omega}^{\varepsilon}$ and $\phi_{g_{\varepsilon}} = \mathbf{T}(g_{\varepsilon}) \in H^1(\Omega)$, there exists $\varepsilon_0 > 0$ such that for every $\varepsilon \leq \varepsilon_0$ we have

$$||\phi^{\varepsilon}||_{H^{1}(\Omega)} \leq C||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}, \qquad ||\phi^{\varepsilon} - \phi_{g_{\varepsilon}}||_{L^{2}(\Omega)} \leq C\varepsilon^{1/2}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)},$$

$$||\rho\left(\nabla\phi^{\varepsilon} - \nabla\phi_{g_{\varepsilon}} - \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon}\left(\frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}}\right)\nabla_{y}\chi_{i}\left(\frac{\cdot}{\varepsilon}\right)\right)||_{L^{2}(\Omega;\mathbb{R}^{n})} \leq C\varepsilon^{1/2}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$

$$(5.2)$$

Moreover we have

$$||\phi^{\varepsilon}||_{\rho} \le C(\varepsilon^{1/2}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} + ||g_{\varepsilon}||_{H^{-1/2}(\partial\Omega)}). \tag{5.3}$$

Proof. Step 1. We prove the first estimate in (5.2). From Section 7.1 we get

$$||\phi_{g_{\varepsilon}}||_{H^{1}(\Omega)} \le C||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} \qquad ||\phi_{g_{\varepsilon}}||_{\rho} \le C||g_{\varepsilon}||_{H^{-1/2}(\partial\Omega)}. \tag{5.4}$$

We write (5.1) in the following weak form:

$$\phi^{\varepsilon} = \overline{\phi}_{\varepsilon} + \phi_{g_{\varepsilon}}, \quad \overline{\phi}_{\varepsilon} \in H_0^1(\Omega)$$

$$\int_{\Omega} A_{\varepsilon} \nabla \overline{\phi}_{\varepsilon} \nabla v = -\int_{\Omega} A_{\varepsilon} \nabla \phi_{g_{\varepsilon}} \nabla v \qquad \forall v \in H_0^1(\Omega).$$
(5.5)

The solution $\overline{\phi}_{\varepsilon}$ of the above variational problem satisfies

$$||\overline{\phi}_{\varepsilon}||_{H^1(\Omega)} \le C||\nabla \phi_{g_{\varepsilon}}||_{L^2(\Omega;\mathbb{R}^n)}.$$

Hence, from $(5.4)_1$ and the above estimate we get the first inequality in (5.2).

Step 2. We prove the second estimate in (5.2). For every test function $v \in H_0^1(\Omega)$ we have

$$\int_{\Omega} A_{\varepsilon} \nabla \phi^{\varepsilon} \nabla v = 0. \tag{5.6}$$

Now, in order to obtain the L^2 error estimate we proceed as in the proof of the Theorem 3.2 in [15]. We first recall that for any $\phi \in H^1(\Omega)$ we have (see Lemma 2.3) for every $\varepsilon \leq \varepsilon_0 = \gamma_0/3\sqrt{n}$

$$||\phi||_{L^2(\widetilde{\Omega}_{3c_0\sqrt{n}\varepsilon})} \le C\varepsilon^{1/2}||\phi||_{H^1(\Omega)}.$$

Let U be a test function belonging to $H_0^1(\Omega) \cap H^2(\Omega)$. The above estimate yields

$$||\nabla U||_{L^2(\widetilde{\Omega}_{3c_0,\sqrt{n}\varepsilon};\mathbb{R}^n)} \le C\varepsilon^{1/2}||U||_{H^2(\Omega)}$$
(5.7)

which in turn with (2.12)-(2.13)-(2.14)₁ and (5.2)₁-(5.6) lead to

$$\left| \int_{\Omega \times Y} A(y) \mathcal{T}_{\varepsilon}(\nabla \phi^{\varepsilon})(x, y) \nabla U(x) dx dy \right| \le C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)} ||U||_{H^{2}(\Omega)}. \tag{5.8}$$

The Theorem 2.3 in [15] gives an element $\widehat{\phi}_{\varepsilon} \in L^2(\Omega; H^1_{per}(Y))$ such that

$$||\mathcal{T}(\nabla \phi^{\varepsilon}) - \nabla \phi^{\varepsilon} - \nabla_{y} \widehat{\phi}_{\varepsilon}||_{[L^{2}(Y;(H^{1}(\Omega))')]^{n}} \leq C \varepsilon^{1/2} ||\nabla \phi^{\varepsilon}||_{L^{2}(\Omega;\mathbb{R}^{n})}$$

$$\leq C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$
(5.9)

The above inequality and (5.8) yield

$$\left| \int_{\Omega \times Y} A(\nabla \phi^{\varepsilon} + \nabla_y \widehat{\phi}_{\varepsilon}) \nabla U \right| \le C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)} ||U||_{H^2(\Omega)}.$$
 (5.10)

We set

$$\forall x \in \mathbb{R}^n, \qquad \rho_{\varepsilon}(x) = \inf \left\{ 1, \frac{\rho(x)}{\varepsilon} \right\}.$$

Now, we take $\overline{\chi} \in H^1_{per}(Y)$ and we consider the test function $u_{\varepsilon} \in H^1_0(\Omega)$ defined for a.e. $x \in \Omega$ by

$$u_{\varepsilon}(x) = \varepsilon \rho_{\varepsilon}(x) \mathcal{Q}_{\varepsilon} \left(\frac{\partial U}{\partial x_i} \right) (x) \overline{\chi} \left(\frac{x}{\varepsilon} \right).$$

Due to $(2.24)_2$ and (5.7) we get

$$\left\| \mathcal{Q}_{\varepsilon} \left(\frac{\partial U}{\partial x_{i}} \right) \nabla_{y} \overline{\chi} \left(\frac{\cdot}{\varepsilon} \right) \right\|_{L^{2}(\widetilde{\Omega}_{\sqrt{n}\varepsilon}; \mathbb{R}^{n})} \leq C \varepsilon^{1/2} ||U||_{H^{2}(\Omega)} ||\overline{\chi}||_{H^{1}(Y)}$$
(5.11)

Then by a straightforward calculation and thanks to $(2.24)_2$ - $(2.25)_2$ and (5.7)-(5.11) we obtain

$$\left\| \nabla u_{\varepsilon} - \mathcal{Q}_{\varepsilon} \left(\frac{\partial U}{\partial x_{i}} \right) \nabla_{y} \overline{\chi} \left(\frac{\cdot}{\varepsilon} \right) \right\|_{L^{2}(\Omega; \mathbb{R}^{n})} \leq C \varepsilon^{1/2} ||U||_{H^{2}(\Omega)} ||\overline{\chi}||_{H^{1}(Y)}$$

which in turn with again (5.11) give

$$\|\nabla u_{\varepsilon}\|_{L^{2}(\widetilde{\Omega}, \overline{\eta_{\varepsilon}}; \mathbb{R}^{n})} \le C\varepsilon^{1/2} ||U||_{H^{2}(\Omega)} ||\overline{\chi}||_{H^{1}(Y)}$$
(5.12)

and then with $(2.25)_1$ that yields

$$\left\| \nabla u_{\varepsilon} - \mathcal{M}_{\varepsilon} \left(\frac{\partial U}{\partial x_{i}} \right) \nabla_{y} \overline{\chi} \left(\frac{\cdot}{\varepsilon} \right) \right\|_{L^{2}(\Omega; \mathbb{R}^{n})} \leq C \varepsilon^{1/2} ||U||_{H^{2}(\Omega)} ||\overline{\chi}||_{H^{1}(Y)}.$$

In (5.6) we replace ∇u_{ε} with $\mathcal{M}_{\varepsilon}\left(\frac{\partial U}{\partial x_{i}}\right)\nabla_{y}\overline{\chi}\left(\frac{\cdot}{\varepsilon}\right)$; we continue using (2.12)-(2.13) and (5.2)₁-(5.12) to obtain

$$\left| \int_{\Omega \times Y} A(y) \mathcal{T}_{\varepsilon}(\nabla \phi^{\varepsilon})(x,y) \mathcal{M}_{\varepsilon}\left(\frac{\partial U}{\partial x_{i}}\right)(x) \nabla_{y} \overline{\chi}(y) dx dy \right| \leq C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)} ||U||_{H^{2}(\Omega)} ||\overline{\chi}||_{H^{1}(Y)}$$

which with $(2.17)_2$ and then (5.9) give

$$\left| \int_{\Omega \times V} A(\nabla \phi^{\varepsilon} + \nabla_y \widehat{\phi}_{\varepsilon}) \frac{\partial U}{\partial x_i} \nabla_y \overline{\chi} \right| \le C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)} ||U||_{H^2(\Omega)} ||\overline{\chi}||_{H^1(Y)}. \tag{5.13}$$

As in [15] we introduce the correctors $\overline{\chi}_i \in H^1_{per}(Y)$, $i \in \{1, \dots, n\}$, defined by

$$\int_{Y} A\nabla_{y} \psi \nabla_{y}(\overline{\chi}_{i} + y_{i}) = 0 \qquad \forall \psi \in H^{1}_{per}(Y).$$
 (5.14)

From (5.13) we get

$$\left| \int_{\Omega \times Y} A(\nabla \phi^{\varepsilon} + \nabla_{y} \widehat{\phi}_{\varepsilon}) \nabla_{y} \left(\sum_{i=1}^{n} \frac{\partial U}{\partial x_{i}} \overline{\chi}_{i} \right) \right| \leq C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)} ||U||_{H^{2}(\Omega)}$$

and from the definition (4.2) of the correctors χ_i we have

$$\int_{\Omega \times Y} A \left(\nabla \phi^{\varepsilon} + \sum_{i=1}^{n} \frac{\partial \phi^{\varepsilon}}{\partial x_{i}} \nabla_{y} \chi_{i} \right) \nabla_{y} \left(\sum_{j=1}^{n} \frac{\partial U}{\partial x_{j}} \overline{\chi}_{j} \right) = 0.$$

Thus

$$\Big| \int_{\Omega \times Y} A \nabla_y \Big(\widehat{\phi}_{\varepsilon} - \sum_{i=1}^n \frac{\partial \phi^{\varepsilon}}{\partial x_i} \chi_i \Big) \nabla_y \Big(\sum_{j=1}^n \frac{\partial U}{\partial x_j} \overline{\chi}_j \Big) \Big| \le C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} ||U||_{H^2(\Omega)}$$

and thanks to (5.14) we obtain

$$\left| \int_{\Omega \times Y} A \nabla_y \left(\widehat{\phi}_{\varepsilon} - \sum_{i=1}^n \frac{\partial \phi^{\varepsilon}}{\partial x_i} \chi_i \right) \nabla U \right| \le C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)} ||U||_{H^2(\Omega)}.$$

The above estimate, (5.10) and the expression (4.5) of the matrix \mathcal{A} yield

$$\left| \int_{\Omega} \mathcal{A} \nabla \phi^{\varepsilon} \nabla U \right| \leq C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)} ||U||_{H^{2}(\Omega)}.$$

Finally, since we have $\int_{\Omega} \mathcal{A} \nabla \phi_{g_{\varepsilon}} \nabla v = 0$ for any $v \in H_0^1(\Omega)$, we deduce that

$$\forall U \in H_0^1(\Omega) \cap H^2(\Omega), \qquad \left| \int_{\Omega} \mathcal{A} \nabla (\phi^{\varepsilon} - \phi_{g_{\varepsilon}}) \nabla U \right| \leq C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} ||U||_{H^2(\Omega)}.$$

Now, let $U_{\varepsilon} \in H_0^1(\Omega)$ be the solution of the following variational problem:

$$\int_{\Omega} \mathcal{A} \nabla v \nabla U_{\varepsilon} = \int_{\Omega} v(\phi^{\varepsilon} - \phi_{g_{\varepsilon}}), \qquad \forall v \in H_0^1(\Omega).$$

Under the assumption on the boundary of Ω , we know that U_{ε} belongs to $H_0^1(\Omega) \cap H^2(\Omega)$ and satisfies $||U_{\varepsilon}||_{H^2(\Omega)} \leq C||\phi^{\varepsilon} - \phi_{g_{\varepsilon}}||_{L^2(\Omega)}$. Therefore, the second estimate in (5.2) is proved.

Step 3. We prove the third estimate in (5.2) and (5.3). The partial derivative $\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_i}$ satisfies

$$\operatorname{div}\left(\mathcal{A}\nabla\left(\frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}}\right)\right) = 0 \quad \text{in} \quad \Omega, \quad \frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}} \in L^{2}(\Omega).$$

Thus, from Remark 7.7 and estimate $(5.4)_2$ we get

$$\left\| \rho \nabla \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \right\|_{L^{2}(\Omega; \mathbb{R}^{n})} \leq C \left\| \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right\|_{L^{2}(\Omega)} \leq C ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)}. \tag{5.15}$$

Now, let U be in $H^1_0(\Omega)$, the function ρU belongs to $H^1_{1/\rho}(\Omega)$. Applying the Theorem 3.2 with the function ρU , there exists $\widehat{u}_{\varepsilon} \in L^2(\Omega; H^1_{per}(Y))$ such that

$$||\mathcal{T}_{\varepsilon}(\nabla(\rho U)) - \nabla(\rho U) - \nabla_{y}\widehat{u}_{\varepsilon}||_{L^{2}(Y;(H^{1}_{\rho}(\Omega;\mathbb{R}^{n}))')} \leq C\varepsilon||\rho U||_{H^{1}_{1/\rho}(\Omega)} \leq C\varepsilon||U||_{H^{1}(\Omega)}. \quad (5.16)$$

The above estimates (5.15) and (5.16) lead to

$$\left| \int_{\Omega \times Y} A \left(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y} \chi_{i} \right) \left(\mathcal{T}_{\varepsilon} \left(\nabla (\rho U) \right) - \nabla (\rho U) - \nabla_{y} \widehat{u}_{\varepsilon} \right) \right| \leq C \varepsilon ||U||_{H^{1}(\Omega)} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)}$$

By definition of the correctors χ_i we have

$$\int_{\Omega \times Y} A \left(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y} \chi_{i} \right) \nabla_{y} \widehat{u}_{\varepsilon} = 0.$$

Besides, from the definitions of the function $\phi_{g_{\varepsilon}}$ and the homogenized matrix \mathcal{A} we have

$$0 = \int_{\Omega} \mathcal{A} \nabla \phi_{g_{\varepsilon}} \nabla (\rho U) = \int_{\Omega \times Y} A \Big(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y} \chi_{i} \Big) \nabla (\rho U).$$

The above inequality and equalities yield

$$\left| \int_{\Omega \times Y} A \left(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y} \chi_{i} \right) \mathcal{T}_{\varepsilon} \left(\nabla (\rho U) \right) \right| \leq C \varepsilon ||\nabla U||_{L^{2}(\Omega; \mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)}.$$
 (5.17)

We have

$$\nabla(\rho U) = \rho \left(\nabla U + \nabla \rho \frac{U}{\rho}\right)$$

and (3.2), therefore, since $U/\rho \in L^2(\Omega)$ and $||U/\rho||_{L^2(\Omega)} \leq C||\nabla U||_{L^2(\Omega;\mathbb{R}^n)}$ we get

$$\left\| \mathcal{T}_{\varepsilon} \left(\nabla (\rho U) \right) - \rho \mathcal{T}_{\varepsilon} \left(\nabla U + \nabla \rho \frac{U}{\rho} \right) \right\|_{L^{2}(\Omega; \mathbb{R}^{n})} \leq C \varepsilon \left\| \nabla U + \nabla \rho \frac{U}{\rho} \right\|_{L^{2}(\Omega; \mathbb{R}^{n})} \leq C \varepsilon ||U||_{H^{1}(\Omega)}.$$

From (5.17) and the above inequalities we deduce that

$$\Big| \int_{\Omega \times Y} A \Big(\rho \nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \rho \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y} \chi_{i} \Big) \mathcal{T}_{\varepsilon} \Big(\nabla U + \nabla \rho \frac{U}{\rho} \Big) \Big| \leq C \varepsilon ||\nabla U||_{L^{2}(\Omega; \mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)}.$$

We recall that $\rho \nabla \phi_{g_{\varepsilon}} \in H_0^1(\Omega; \mathbb{R}^n)$, hence from $(2.14)_2$, $(2.17)_1$ and (5.15) we get

$$\left| \int_{\Omega \times Y} A \left(\rho \nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \rho \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y} \chi_{i} \right) \mathcal{T}_{\varepsilon} \left(\nabla U + \nabla \rho \frac{U}{\rho} \right) \right|$$

$$- \int_{\Omega \times Y} A \left(\mathcal{T}_{\varepsilon} (\rho \nabla \phi_{g_{\varepsilon}}) + \sum_{i=1}^{n} \mathcal{M}_{\varepsilon} \left(\rho \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \nabla_{y} \chi_{i} \right) \mathcal{T}_{\varepsilon} \left(\nabla U + \nabla \rho \frac{U}{\rho} \right) \right| \leq C \varepsilon ||\nabla U||_{L^{2}(\Omega; \mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)}.$$

Then, transforming by inverse unfolding we obtain

$$\left| \int_{\widehat{\Omega}_{\varepsilon}} A_{\varepsilon} \left(\rho \nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \mathcal{M}_{\varepsilon} \left(\rho \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \nabla_{y} \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \right) \left(\nabla U + \nabla \rho \frac{U}{\rho} \right) \right| \leq C \varepsilon ||\nabla U||_{L^{2}(\Omega; \mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)}.$$

Now, thanks to (2.28) and (5.15) we get

$$\left| \int_{\Omega} A_{\varepsilon} \rho \left(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \mathcal{M}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \nabla_{y} \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \right) \left(\nabla U + \nabla \rho \frac{U}{\rho} \right) \right| \leq C \varepsilon ||\nabla U||_{L^{2}(\Omega; \mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$

Then using $(2.29)_1$ it leads to

$$\left| \int_{\Omega} A_{\varepsilon} \left(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \nabla_{y} \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \right) \nabla(\rho U) \right| \leq C \varepsilon ||\nabla U||_{L^{2}(\Omega; \mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)}.$$

We recall that
$$\int_{\Omega} A_{\varepsilon} \nabla \phi^{\varepsilon} \nabla (\rho U) = 0$$
. We choose $U = \rho \Big(\phi^{\varepsilon} - \phi_{g_{\varepsilon}} - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \Big(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \Big) \chi_{i} \Big(\frac{\dot{\cdot}}{\varepsilon} \Big) \Big)$

which belongs to $H_0^1(\Omega)$. Due to the second estimate in (5.2), the third one in (5.2) follows immediately. The estimates (5.3) are the consequences of $(5.2)_2$, $(5.2)_3$ and $(5.4)_2$.

Corollary 5.2. Let $(\phi^{\varepsilon})_{{\varepsilon}>0}$ be a sequence of functions belonging to $H^1(\Omega)$ and satisfying (5.1). We set $g_{\varepsilon} = \phi^{\varepsilon}_{|\partial\Omega}$, if we have

$$g_{\varepsilon} \rightharpoonup g$$
 weakly in $H^{1/2}(\partial\Omega)$

then we obtain

$$\phi^{\varepsilon} \rightharpoonup \phi_g \quad \text{weakly in} \quad H^1(\Omega),$$

$$\phi^{\varepsilon} - \phi_g - \varepsilon \sum_{i=1}^n \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_g}{\partial x_i}\right) \chi_i \left(\frac{\cdot}{\varepsilon}\right) \longrightarrow 0 \quad \text{strongly in} \quad H^1_{\rho}(\Omega).$$

$$(5.18)$$

Moreover, if

$$g_{\varepsilon} \longrightarrow g \quad strongly \ in \quad H^{1/2}(\partial\Omega)$$
 (5.19)

then we have

$$\phi^{\varepsilon} - \phi_g - \varepsilon \sum_{i=1}^n \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_g}{\partial x_i} \right) \chi_i \left(\frac{\cdot}{\varepsilon} \right) \longrightarrow 0 \quad strongly \ in \quad H^1(\Omega).$$
 (5.20)

Proof. Thanks to $(5.2)_1$ the sequence $(\phi^{\varepsilon})_{\varepsilon>0}$ is uniformly bounded in $H^1(\Omega)$. Then, due to Lemma 7.1 and Remark 7.7 we get

$$||\phi_g - \phi_{g_{\varepsilon}}||_{\rho} \le C||g - g_{\varepsilon}||_{H^{-1/2}(\partial\Omega)}$$

which with $(5.2)_2$ (resp. $(5.2)_3$) give the convergence $(5.18)_1$ (resp. $(5.18)_2$). Under the assumption (5.19), we use (7.1) and we proceed as in the proof of Theorem 6.1 of [9] to obtain the strong convergence (5.20).

Theorem 5.3. Let ϕ^{ε} be the solution of the following homogenization problem:

$$-div(A_{\varepsilon}\nabla\phi^{\varepsilon}) = f \quad in \quad \Omega, \qquad \phi^{\varepsilon} = g \quad on \quad \partial\Omega$$

where $f \in L^2(\Omega)$ and $g \in H^{1/2}(\partial\Omega)$. We have

$$||\phi^{\varepsilon} - \Phi||_{L^{2}(\Omega)} \le C\{\varepsilon||f||_{L^{2}(\Omega)} + \varepsilon^{1/2}||g||_{H^{1/2}(\partial\Omega)}\},$$

$$\left\| \rho \left(\nabla \phi^{\varepsilon} - \nabla \Phi - \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{i}} \right) \nabla_{y} \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \right) \right\|_{L^{2}(\Omega; \mathbb{R}^{n})} \leq C \left\{ \varepsilon ||f||_{L^{2}(\Omega)} + \varepsilon^{1/2} ||g||_{H^{1/2}(\partial \Omega)} \right\}$$

where Φ is the solution of the homogenized problem

$$-\operatorname{div}(\mathcal{A}\nabla\Phi)=f\quad \text{in}\quad \Omega, \qquad \Phi=g\quad \text{on}\quad \partial\Omega.$$

Moreover we have

$$\phi^{\varepsilon} - \Phi - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{i}} \right) \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \longrightarrow 0 \quad strongly \ in \quad H^{1}(\Omega). \tag{5.21}$$

Proof. Let $\widetilde{\phi}^{\varepsilon}$ be the solution of the homogenization problem

$$\widetilde{\phi}^{\varepsilon} \in H_0^1(\Omega), \quad -\operatorname{div}(A_{\varepsilon} \nabla \widetilde{\phi}^{\varepsilon}) = f \quad \text{in} \quad \Omega$$

and $\widetilde{\Phi}$ the solution of the homogenized problem

$$\widetilde{\Phi} \in H_0^1(\Omega), \quad -\text{div}(A\nabla \widetilde{\Phi}) = f \quad \text{in} \quad \Omega$$

The Theorem 3.2 in [15] gives the following estimate:

$$||\widetilde{\phi}^{\varepsilon} - \widetilde{\Phi}||_{L^{2}(\Omega)} + \left\| \rho \nabla \left(\widetilde{\phi}^{\varepsilon} - \widetilde{\Phi} - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \widetilde{\Phi}}{\partial x_{i}} \right) \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \right) \right\|_{L^{2}(\Omega; \mathbb{R}^{n})} \leq C\varepsilon ||f||_{L^{2}(\Omega)} \quad (5.22)$$

while the Theorem 4.1 in [14] gives

$$\left\| \widetilde{\phi}^{\varepsilon} - \widetilde{\Phi} - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \widetilde{\Phi}}{\partial x_{i}} \right) \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \right\|_{H^{1}(\Omega)} \leq C \varepsilon^{1/2} ||f||_{L^{2}(\Omega)}.$$
 (5.23)

Then $\phi^{\varepsilon} - \widetilde{\phi}^{\varepsilon}$ satisfies

$$\operatorname{div}(A_{\varepsilon}\nabla(\phi^{\varepsilon} - \widetilde{\phi}^{\varepsilon})) = 0 \quad \text{in} \quad \Omega, \quad \phi^{\varepsilon} - \widetilde{\phi}^{\varepsilon} = g \quad \text{on} \quad \partial\Omega.$$

Thanks to the inequalities (5.2) and (5.22) we deduce the estimates of the theorem. The strong convergence (5.21) is a consequence of (5.23) and the strong convergence (5.20) after having observed that $\Phi - \widetilde{\Phi} = \phi_q$.

Remark 5.4. In Theorem 5.3, if $g \in H^{3/2}(\partial\Omega)$ then in the estimates therein, we can replace $\varepsilon^{1/2}||g||_{H^{1/2}(\partial\Omega)}$ with $\varepsilon||g||_{H^{3/2}(\partial\Omega)}$. Moreover we have the following H^1 -global error estimate:

$$\left\|\phi^{\varepsilon} - \Phi - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{i}}\right) \chi_{i} \left(\frac{\cdot}{\varepsilon}\right) \right\|_{H^{1}(\Omega)} \leq C \varepsilon^{1/2} \left\{ ||f||_{L^{2}(\Omega)} + ||g||_{H^{3/2}(\partial \Omega)} \right\}.$$

6 A first result with strongly oscillating boundary data

Now, we consider the solution ϕ^{ε} of the following homogenization problem:

$$\operatorname{div}(A_{\varepsilon}\nabla\phi^{\varepsilon}) = 0 \quad \text{in} \quad \Omega$$

$$\phi^{\varepsilon} = g_{\varepsilon} \quad \text{on} \quad \partial\Omega$$
(6.1)

where $g_{\varepsilon} \in H^{1/2}(\partial\Omega)$. As a consequence of the Theorem 5.1 we first obtain the following theorem:

Theorem 6.1. Let ϕ^{ε} be the solution of the Problem (6.1). If we have

$$g_{\varepsilon} \rightharpoonup g \quad weakly \ in \quad H^{-1/2}(\partial\Omega)$$

and

$$\varepsilon^{1/2}g_{\varepsilon} \longrightarrow 0 \quad strongly \ in \quad H^{1/2}(\partial\Omega)$$
 (6.2)

then

$$\phi^{\varepsilon} \rightharpoonup \phi_g \quad weakly \ in \quad H_{\rho}^1(\Omega).$$
 (6.3)

Furthermore, if we have

$$g_{\varepsilon} \longrightarrow g$$
 strongly in $H^{-1/2}(\partial\Omega)$

then

$$\phi^{\varepsilon} - \phi_g - \varepsilon \sum_{i=1}^n \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_i} \right) \chi_i \left(\frac{\cdot}{\varepsilon} \right) \longrightarrow 0 \quad strongly \ in \quad H^1_{\rho}(\Omega). \tag{6.4}$$

Proof. Due to (5.3) the sequence $(\phi^{\varepsilon})_{\varepsilon>0}$ is strictly bounded in $H^1_{\rho}(\Omega)$. From the estimates (5.2)₃ and (5.4)₂ we get

$$\left\|\phi^{\varepsilon} - \phi_{g_{\varepsilon}} - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}}\right) \chi_{i} \left(\frac{\cdot}{\varepsilon}\right) \right\|_{H_{\rho}^{1}(\Omega)} \leq C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$

Then, using the variational problem (7.4) and $(5.4)_2$ we obtain

$$\phi_{g_{\varepsilon}} \rightharpoonup \phi_g$$
 weakly in $H^1_{\rho}(\Omega)$.

Besides we have $\varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \rightharpoonup 0$ weakly in $H_{\rho}^{1}(\Omega)$. The weak convergence (6.3) is proved.

In the case $g_{\varepsilon} \longrightarrow g$ strongly in $H^{-1/2}(\partial\Omega)$, (7.3) and (7.7) lead to

$$||\phi_{g_{\varepsilon}} - \phi_g||_{H^{1}_{\rho}(\Omega)} \le C||g_{\varepsilon} - g||_{H^{-1/2}(\partial\Omega)}.$$

Hence with $(2.29)_2$ they yield the strong convergence (6.4).

In a forthcoming paper we will show that in both cases (weak or strong convergence of g_{ε} towards g in $H^{-1/2}(\partial\Omega)$) the assumption (6.2) is essential in order to obtain at least (6.3).

7 Appendix

7.1 An operator from $H^{-1/2}(\partial\Omega)$ into $L^2(\Omega)$

In this section Ω is an open bounded set with a $\mathcal{C}^{1,1}$ boundary or an open bounded convex polygon (n=2) or polyhedral (n=3). Let \mathcal{A} be a $n \times n$ definite positive constant matrix (e.g. the matrix \mathcal{A} given by (4.5)). For every couple $(\psi, \Psi) \in [\mathcal{C}^{\infty}(\overline{\Omega})]^2$, integrating by parts over Ω gives

$$\int_{\Omega} \mathcal{A}\nabla \psi(x)\nabla \Psi(x)dx = -\int_{\Omega} \psi(x)\operatorname{div}(\mathcal{A}^{T}\nabla \Psi)(x)dx + \int_{\partial \Omega} \psi(x)(\mathcal{A}^{T}\nabla \Psi)(x) \cdot \nu(x)d\sigma.$$

The space $\mathcal{C}^{\infty}(\overline{\Omega})$ being dense in $H^1(\Omega)$ and $H^2(\Omega)$, hence the above equality holds true for any $\psi \in H^1(\Omega)$ and any $\Psi \in H^2(\Omega)$. Now, let g be in $H^{1/2}(\partial\Omega)$, there exists one $\phi_g \in H^1(\Omega)$ such that

$$\operatorname{div}(\mathcal{A}\nabla\phi_q) = 0$$
 in Ω , $\phi_q = g$ on $\partial\Omega$

and we have

$$||\phi_g||_{H^1(\Omega)} \le C||g||_{H^{1/2}(\partial\Omega)}.$$
 (7.1)

Besides, for any function $\Psi \in H_0^1(\Omega) \cap H^2(\Omega)$ we get

$$\int_{\Omega} \phi_g(x) \operatorname{div}(\mathcal{A}^T \nabla \Psi)(x) dx = \int_{\partial \Omega} g(x) \left(\mathcal{A}^T \nabla \Psi \right)(x) \cdot \nu(x) d\sigma. \tag{7.2}$$

Under the assumption on Ω the function $\Psi(g)$ defined by

$$\Psi(g) \in H_0^1(\Omega), \quad \operatorname{div}(\mathcal{A}^T \nabla \Psi(g)) = \phi_g \quad \text{in} \quad \Omega$$

also belongs to $H^2(\Omega)$ and satisfies

$$||\Psi(g)||_{H^2(\Omega)} \le C||\phi_g||_{L^2(\Omega)}.$$

Taking $\Psi = \Psi(g)$ in the above equality (7.2) we obtain

$$\int_{\Omega} |\phi_g(x)|^2 dx = \int_{\partial \Omega} g(x) \left(\mathcal{A}^T \nabla \Psi(g)(x) \right) \cdot \nu(x) d\sigma \le ||g||_{H^{-1/2}(\partial \Omega)} ||\left(\mathcal{A}^T \nabla \Psi(g) \right) \cdot \nu||_{H^{1/2}(\partial \Omega)}$$

$$\le C||g||_{H^{-1/2}(\partial \Omega)} ||\Psi(g)||_{H^2(\Omega)}.$$

This leads to

$$||\phi_g||_{L^2(\Omega)} \le C||g||_{H^{-1/2}(\partial\Omega)}.$$
 (7.3)

We denote by **T** the operator from $H^{1/2}(\partial\Omega)$ into $H^1(\Omega)$ which associates to $g \in H^{1/2}(\partial\Omega)$ the function $\phi_g \in H^1(\Omega)$. Due to (7.3), this operator admits an extension (still denoted **T**) from $H^{-1/2}(\partial\Omega)$ into $L^2(\Omega)$ and we have

$$\forall g \in H^{-1/2}(\partial\Omega), \qquad ||\mathbf{T}(g)||_{L^2(\Omega)} \le C||g||_{H^{-1/2}(\partial\Omega)}.$$

Hence, for $g \in H^{-1/2}(\partial\Omega)$, we also denote $\phi_g = \mathbf{T}(g)$. This function is the "very weak" solution of the problem

$$\phi_g \in L^2(\Omega), \quad \operatorname{div}(\mathcal{A}\nabla\phi_g) = 0 \quad \text{in } \Omega, \quad \phi_g = g \quad \text{on } \partial\Omega$$

or the solution of the following:

$$\phi_g \in L^2(\Omega),$$

$$\int_{\Omega} \phi_g(x) \operatorname{div}(\mathcal{A}^T \nabla \psi(x)) dx = \langle g, (\mathcal{A}^T \nabla \psi) \cdot \nu \rangle_{H^{-1/2}(\partial\Omega), H^{1/2}(\partial\Omega)}, \tag{7.4}$$

$$\forall \psi \in H_0^1(\Omega) \cap H^2(\Omega).$$

Lemma 7.1. The operator **T** is a bicontinuous linear operator from $H^{-1/2}(\partial\Omega)$ onto

$$\mathbf{H} = \Big\{ \phi \in L^2(\Omega) \mid \operatorname{div}(\mathcal{A} \nabla \phi) = 0 \quad \text{in} \quad \Omega \Big\}.$$

There exists a constant $C \geq 1$ such that

$$\forall g \in H^{-1/2}(\partial\Omega), \qquad \frac{1}{C}||g||_{H^{-1/2}(\partial\Omega)} \le ||\mathbf{T}(g)||_{L^{2}(\Omega)} \le C||g||_{H^{-1/2}(\partial\Omega)}. \tag{7.5}$$

Proof. Let ϕ be in **H** we are going to prove that there exists an element $g \in H^{-1/2}(\partial\Omega)$ such that $\mathbf{T}(g) = \phi$. To do that, we consider a continuous linear lifting operator **R** from $H^{1/2}(\partial\Omega)$ into $H^1_0(\Omega) \cap H^2(\Omega)$ satisfying for any $h \in H^{1/2}(\partial\Omega)$

$$\mathbf{R}(h) \in H_0^1(\Omega) \cap H^2(\Omega),$$

$$\mathcal{A}^T \nabla \mathbf{R}(h)_{|\partial\Omega} \cdot \nu = h \quad \text{on} \quad \partial\Omega,$$

$$||\mathbf{R}(h)||_{H^2(\Omega)} \le C||h||_{H^{1/2}(\partial\Omega)}.$$

The map $h \mapsto \int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \mathbf{R}(h))$ is a continuous linear form defined over $H^{1/2}(\partial \Omega)$. Thus, there exists $g \in H^{-1/2}(\partial \Omega)$ such that

$$\int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \mathbf{R}(h)) = \langle g, h \rangle_{H^{-1/2}(\partial \Omega), H^{1/2}(\partial \Omega)}.$$
 (7.6)

Since $\phi \in \mathbf{H}$, we deduce that for any $\psi \in \mathcal{C}_0^{\infty}(\Omega)$ we have $\int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \psi) = 0$. Therefore, for any $\psi \in H_0^2(\Omega)$ we have $\int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \psi) = 0$. Taking into account (7.6) we get

$$\int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \psi) = \langle g, (\mathcal{A}^T \nabla \psi) \cdot \nu \rangle_{H^{-1/2}(\partial\Omega), H^{1/2}(\partial\Omega)}, \qquad \forall \psi \in H^1_0(\Omega) \cap H^2(\Omega).$$

It yields $\phi = \phi_g$ and then (7.5).

Remark 7.2. It is well known (see e.g. [17]) that every function $\phi \in \mathbf{H}$ also belongs to $H^1_{\rho}(\Omega)$ and verifies

$$||\phi||_{\rho} \le C||\phi||_{L^2(\Omega)}. \tag{7.7}$$

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