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A Remark on the Stability of Peakons for the Degasperis-Procesi Equation

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Abstract

In this paper, we present a new argument (see Lemma 3.4) that allows us to simplify the proof of stability of peakons established in [7] (Theorem 1.1).

1 Introduction

In this paper, we consider the Degasperis-Procesi equation (DP)

$$u_t - u_{txx} + 4uu_x = 3u_x u_{xx} + uu_{xxx}, \quad (t, x) \in \mathbb{R}_+ \times \mathbb{R}, \tag{1.1}$$

with $u(0) = u_0 \in L^2(\mathbb{R})$ and $(1 - \partial_x^2)u_0 \in \mathcal{M}^+(\mathbb{R})$.

The DP equation is completely integrable (see [3]) and has been proved to be physically relevant for water waves (see [1]). It possesses, among others, the following conservation laws

$$E(u) = \int_{\mathbb{R}} yv = \int_{\mathbb{R}} \left(4v^2 + 5v_x^2 + v_{xx}^2 \right), \quad F(u) = \int_{\mathbb{R}} u^3 = \int_{\mathbb{R}} \left(-v_{xx}^3 + 12vv_{xx}^2 - 48v^2v_{xx} + 64v^3 \right), \quad (1.2)$$

where $y = (1 - \partial_x^2)u$ and $v = (4 - \partial_x^2)^{-1}u$. One can notice that the conservation law $E(\cdot)$ is equivalent to $\|\cdot\|_{L^2(\mathbb{R})}^2$. Indeed, using integration by parts (we assume that $u(\pm \infty) = v(\pm \infty) = v_x(\pm \infty) = 0$), it holds

$$||u||_{L^{2}(\mathbb{R})}^{2} = \int_{\mathbb{R}} u^{2} = \int_{\mathbb{R}} (4v - v_{xx})^{2} = \int_{\mathbb{R}} \left(16v^{2} + 8v_{x}^{2} + v_{xx}^{2} \right) \sim E(u).$$
(1.3)

In the sequel we will denote

$$\|u\|_{\mathcal{H}} = \sqrt{E(u)}.\tag{1.4}$$

Applying $(1 - \partial_x^2)^{-1}$ to (1.1), we obtain

$$u_t + \frac{1}{2}\partial_x u^2 + \frac{3}{2}(1 - \partial_x^2)^{-1}\partial_x u^2 = 0, \quad (t, x) \in \mathbb{R}_+ \times \mathbb{R}.$$
 (1.5)

In this form, the DP equation admits explicit solitary waves called *peakons* (see [3]) that are defined by

$$u(t,x) = \varphi_c(x-ct) = c\varphi(x-ct) = ce^{-|x-ct|}, \quad (t,x) \in \mathbb{R}_+ \times \mathbb{R}.$$
(1.6)

Our goal is to simplify the proof given in [7] of the stability of a single peakon for the DP equation. Recall that the proof of the stability for the Camassa-Holm equation (CH) in [2] follows from two integral relations between two conservation laws of CH, $\max_{\mathbb{R}} u$ and functions related to u. In [7] the proof is more complicated, since all the local maxima and minima of $v = (4 - \partial_x^2)^{-1}u$ are involved in the relations. In this paper we present a simplication of this proof, where only the maximum of v is involved in the relations. Our proof is thus closer to the proof for CH in [2]. The main idea is the following: since u is L^2 -close to the peakon $\varphi_c(\cdot - \xi)$, for some $\xi \in \mathbb{R}$, and $(1 - \partial_x^2)u \in \mathcal{M}^+(\mathbb{R})$, it is easy to check that u is actually C^0 -close to the peakon, and thus v is C^2 -close to the *smooth-peakon*:

$$\rho_c(x-\xi) = (4-\partial_x^2)^{-1}\varphi_c(x-\xi) = \frac{c}{3}e^{-|x-\xi|} - \frac{c}{6}e^{-2|x-\xi|}, \quad x \in \mathbb{R}.$$
(1.7)

First, since ρ_c , ρ'_c and ρ''_c are very small with respect to the amplitude c outside of the interval $\Theta_0 = [-6.7, 6.7]$, we can restrict ourself to study v on $\Theta_{\xi} = [\xi - 6.7, \xi + 6.7]$. Now we observe that ρ''_c has strictly negative values in the interval $\mathcal{V}_0 = [-\ln\sqrt{2}, \ln\sqrt{2}]$, with ρ'_c strictly positive on $[-6.7, -\ln\sqrt{2}]$ and ρ'_c strictly negative on $[\ln\sqrt{2}, 6.7]$. This forces v_x to change sign only one time on Θ_{ξ} , and thus v has only one local extremum (which is a maximum) on Θ_{ξ} . This fact will considerably simplify the proof of the stability.

2 Preliminaries

In this section, we briefly recall the global well-posedness results for the Cauchy problem of the DP equation (see [5] and [8]), and its consequences. For I a finite or infinite time interval of \mathbb{R}_+ , we denote by $\mathcal{X}(I)$ the function space ¹

$$\mathcal{X}(I) = \left\{ u \in C\left(I; H^1(\mathbb{R})\right) \cap L^{\infty}\left(I; W^{1,1}(\mathbb{R})\right), \ u_x \in L^{\infty}\left(I; BV(\mathbb{R})\right) \right\}.$$
(2.1)

Theorem 2.1 (Global weak solution; see [5] and [8]). Assume that $u_0 \in L^2(\mathbb{R})$ with $y_0 = (1 - \partial_x^2)u_0 \in \mathcal{M}^+(\mathbb{R})$. Then the DP equation has a unique global weak solution $u \in \mathcal{X}(\mathbb{R}_+)$ such that

$$y(t,\cdot) = (1 - \partial_x^2)u(t,\cdot) \in \mathcal{M}^+(\mathbb{R}), \quad \forall t \in \mathbb{R}_+$$

$$(2.2)$$

and

$$|u_x(t,x)| \le u(t,x), \quad \forall (t,x) \in \mathbb{R}_+ \times \mathbb{R}.$$
(2.3)

Moreover $E(\cdot)$ and $F(\cdot)$ are conserved by the flow.

Remark 2.1 (Control of L^{∞} norm by L^2 norm). (2.3) and the well-known Sobolev embedding of $H^1(\mathbb{R})$ into $L^{\infty}(\mathbb{R})$ lead to

$$\|u\|_{L^{\infty}(\mathbb{R})} \leq \frac{1}{\sqrt{2}} \|u\|_{H^{1}(\mathbb{R})} \leq \|u\|_{L^{2}(\mathbb{R})}.$$
(2.4)

3 Stability of peakons

In this section, we present our simplification of the proof of stability of peakons for the DP equation.

Theorem 3.1 (Stability of peakons). Let $u \in \mathcal{X}([0,T[))$, with $0 < T \leq +\infty$, be a solution of the DP equation and φ_c be the peakon defined in (1.6), traveling to the right at the speed c > 0. There exist C > 0 and $\varepsilon_0 > 0$ only depending on the speed c, such that if

$$y_0 = (1 - \partial_x^2) u_0 \in \mathcal{M}^+(\mathbb{R})$$
(3.1)

 $^{{}^{1}}W^{1,1}(\mathbb{R})$ is the space of $L^{1}(\mathbb{R})$ functions with derivatives in $L^{1}(\mathbb{R})$ and $BV(\mathbb{R})$ is the space of function with bounded variation.

and

$$\|u_0 - \varphi_c\|_{\mathcal{H}} \le \varepsilon^2, \quad with \quad 0 < \varepsilon < \varepsilon_0, \tag{3.2}$$

then

$$\|u(t,\cdot) - \varphi_c(\cdot - \xi(t))\|_{\mathcal{H}} \le C\sqrt{\varepsilon}, \quad \forall t \in [0,T[,$$
(3.3)

where $\xi(t) \in \mathbb{R}$ is the only point where the function $v(t, \cdot) = (4 - \partial_x^2)^{-1}u(t, \cdot)$ attains its maximum.

We first recall that $E(u) \sim E(\varphi_c)$ and $F(u) \sim F(\varphi_c)$, if $u \sim \varphi_c$ in $L^2(\mathbb{R})$, with $y \in \mathcal{M}^+(\mathbb{R})$ (see for instance [7] or [6]).

Lemma 3.1 (Control of distances between energies; see [6]). Let $u \in L^2(\mathbb{R})$ with $y = (1 - \partial_x^2)u \in \mathcal{M}^+(\mathbb{R})$. If $||u - \varphi_c||_{\mathcal{H}} \leq \varepsilon^2$, then

$$|E(u) - E(\varphi_c)| \le O(\varepsilon^2) \tag{3.4}$$

and

$$|F(u) - F(\varphi_c)| \le O(\varepsilon^2), \tag{3.5}$$

where $O(\cdot)$ only depends on the speed c.

To prove Theorem 3.1, by the conservation of $E(\cdot)$, $F(\cdot)$ and the continuity of the map $t \mapsto u(t)$ from [0, T[to \mathcal{H} (since $\mathcal{H} \simeq L^2$), it suffices to prove that for any function $u \in L^2(\mathbb{R})$ satisfying $y = (1 - \partial_x^2)u \in \mathcal{M}^+(\mathbb{R})$, (3.4) and (3.5), if

$$\inf_{z \in \mathbb{R}} \|u - \varphi_c(\cdot - z)\|_{\mathcal{H}} \le \varepsilon^{1/4},\tag{3.6}$$

then

$$\|u - \varphi_c(\cdot - \xi)\|_{\mathcal{H}} \le C\sqrt{\varepsilon},\tag{3.7}$$

where $\xi \in \mathbb{R}$ is the only point of maximum of v.

Let us present some important properties of smooth-peakons, defined in (1.7), which will play a crutial role in the proof of Theorem 3.1. The smooth-peakon ρ_c belong to $H^3(\mathbb{R}) \hookrightarrow C^2(\mathbb{R})$ (by the Sobolev embedding) since φ_c belong to $H^1(\mathbb{R})$ (defined in (1.6)). It is a positive even function, which admits a single maximum c/6 at point 0, and decays at infinity to 0 (see Figure 1a). Its derivative ρ'_c belong to $H^2(\mathbb{R}) \hookrightarrow C^1(\mathbb{R})$, it is an odd function, which vanishes only at the origin and has negative values on $[0, +\infty[$. It admits a a single minimum -c/12 at point ln2 and tends at infinity to 0 (see Figure 1b). Its second derivative ρ''_c belongs to $H^1(\mathbb{R}) \hookrightarrow C^0(\mathbb{R})$, it is an even function, which vanishes at $\pm \ln 2$, takes positive values on $] -\infty, -\ln 2[\cup]\ln 2, +\infty[$ and negative values on $[-\ln 2, \ln 2]$. It admits a single minimum -c/3 at point 0 and two maxima c/24 at points $\pm \ln 4$, and decays at infinity to 0 (see Figure 1c).

Next, we will need the following estimates.

Lemma 3.2 (C^0 , C^1 and C^2 approximations). Let $u \in L^2(\mathbb{R})$ with $y = (1 - \partial_x^2)u \in \mathcal{M}^+(\mathbb{R})$. If $||u - \varphi_c||_{\mathcal{H}} \leq \varepsilon^{1/4}$, then

$$||u - \varphi_c||_{C^0(\mathbb{R})} + ||v - \rho_c||_{C^2(\mathbb{R})} \le O(\varepsilon^{1/8})$$
(3.8)

and

$$\|v - \rho_c\|_{C^1(\mathbb{R})} \le O(\varepsilon^{1/4}).$$
 (3.9)

Proof. Let us begin with the second estimate. From the definition of $E(\cdot)$ and \mathcal{H} (see respectively (1.2) and (1.4)), one can see that $||u||_{\mathcal{H}}$ is equivalent to $||v||_{H^2(\mathbb{R})}$, since $||v||_{H^2(\mathbb{R})} \leq ||u||_{\mathcal{H}} \leq 5||v||_{H^2(\mathbb{R})}$. Then, assumption u is \mathcal{H} -close to φ_c implies that v is H^2 -close to ρ_c . Now, using the Sobolev embedding of $H^2(\mathbb{R})$ into $C^1(\mathbb{R})$, we deduce (3.9).

For the first estimate, note that the assumption $y = (1 - \partial_x^2)u \ge 0$ implies that $u = (1 - \partial_x^2)^{-1}y \ge 0$ and satisfies $|u_x| \le u$ on \mathbb{R} (see (2.3)). Then, applying triangular inegality, and using that $|\varphi'_c| = \varphi_c$ on \mathbb{R} and (2.4), we have

$$\begin{aligned} \|u - \varphi_c\|_{H^1(\mathbb{R})} &\leq \|u\|_{H^1(\mathbb{R})} + \|\varphi_c\|_{H^1(\mathbb{R})} \\ &\leq 2\|u\|_{L^2(\mathbb{R})} + 2\|\varphi_c\|_{L^2(\mathbb{R})} \\ &\leq 2\|u - \varphi_c\|_{L^2(\mathbb{R})} + 4\|\varphi_c\|_{L^2(\mathbb{R})} \\ &\leq O(\varepsilon^{1/4}) + O(1), \end{aligned}$$

where $\|\varphi_c\|_{L^2(\mathbb{R})} = c$. Therefore, applying the Gagliardo-Nirenberg inegality and using that $\|u - \varphi_c\|_{\mathcal{H}} \leq \varepsilon^{1/4}$ (with $\mathcal{H} \simeq L^2$), we obtain

$$\begin{aligned} \|u - \varphi_c\|_{C^0(\mathbb{R})} &\leq \|u - \varphi_c\|_{L^2(\mathbb{R})}^{1/2} \|u - \varphi_c\|_{H^1(\mathbb{R})}^{1/2} \\ &\leq O(\varepsilon^{1/8}) \left(O(\varepsilon^{1/8}) + O(1) \right) \\ &\leq O(\varepsilon^{1/8}) . \end{aligned}$$

Finally to estimate the second term of the left-hand side of (3.8), we first notice that the continuity of $(4 - \partial_x^2)^{-1}(\cdot)$ from $H^s(\mathbb{R})$ to $H^{s+1}(\mathbb{R})$ and the above estimates ensures that $||v - \rho_c||_{H^3} = O(1)$ and $||v - \rho_c||_{H^2} = O(\varepsilon^{1/4})$. These last estimates combined with the Gagliardo-Nirenberg inequality yields the result as above.

The following lemma specifies the distance to minimize for stability.

Lemma 3.3 (Quadratic identity; see [7]). For any $u \in L^2(\mathbb{R})$ and $\xi \in \mathbb{R}$, it holds

$$E(u) - E(\varphi_c) = \|u - \varphi_c(\cdot - \xi)\|_{\mathcal{H}}^2 + 4c\left(v(\xi) - \frac{c}{6}\right),$$
(3.10)

where $v = (4 - \partial_x^2)^{-1}u$ and $\rho_c(0) = c/6$.

Proof. We follow the idea of Constantin and Strauss with the CH equation (see [2], Lemma 1). We compute

$$E(u - \varphi_c(\cdot - \xi)) = E(u) + E(\varphi_c) - 2\left\langle (1 - \partial_x^2)\varphi_c(\cdot - \xi), (4 - \partial_x^2)^{-1}u \right\rangle_{H^{-1}, H^1} = E(u) + E(\varphi_c) - 2\left\langle (1 - \partial_x^2)\varphi_c(\cdot - \xi), v \right\rangle_{H^{-1}, H^1},$$
(3.11)

where $\langle \cdot, \cdot \rangle_{H^{-1}, H^1}$ denotes the duality bracket $H^{-1}(\mathbb{R})$, $H^1(\mathbb{R})$. Now, using the definition of $\varphi'_c(\cdot - \xi)$ and integration by parts, we have

$$\left\langle (1 - \partial_x^2)\varphi_c(\cdot - \xi), v \right\rangle_{H^{-1}, H^1} = \int_{\mathbb{R}} v\varphi_c(\cdot - \xi) + \int_{\mathbb{R}} v_x \varphi'_c(\cdot - \xi)$$
$$= \int_{\mathbb{R}} v\varphi_c(\cdot - \xi) + \int_{-\infty}^{\xi} v_x \varphi_c(\cdot - \xi) - \int_{\xi}^{+\infty} v_x \varphi_c(\cdot - \xi)$$
$$= 2cv(\xi).$$
(3.12)

Recalling that the energy of peakons is given by

$$E(\varphi_c) = \left\langle (1 - \partial_x^2)\varphi_c, (4 - \partial_x^2)^{-1}\varphi_c \right\rangle_{H^{-1}, H^1} = \int_{\mathbb{R}} \rho_c \varphi_c + \int_{\mathbb{R}} \rho'_c \varphi'_c$$
$$= \int_{\mathbb{R}} \rho_c \varphi_c + \int_{-\infty}^0 \rho'_c \varphi_c - \int_0^{+\infty} \rho'_c \varphi_c = 2c\rho_c(0) = \frac{c^2}{3} , \qquad (3.13)$$

we obtain the lemma.

Now we will study carefully the local extrema of $v = (4 - \partial_x^2)^{-1}u$. Let $u \in L^2(\mathbb{R})$ with $y = (1 - \partial_x^2)u \in \mathcal{M}^+(\mathbb{R})$, and assume that (3.6) holds for some $z \in \mathbb{R}$. We consider the interval in which the mass of

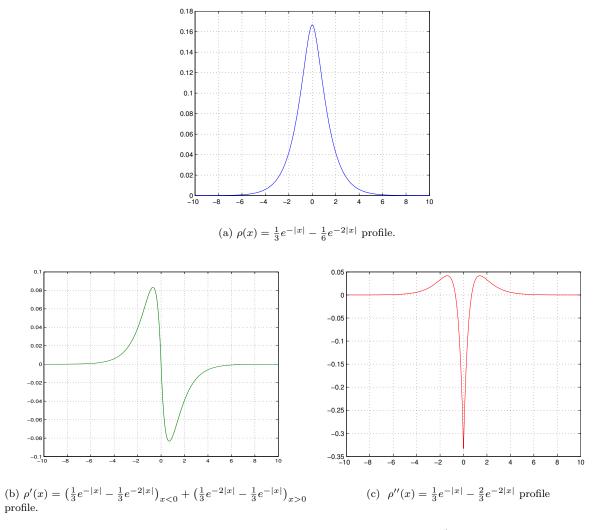


Figure 1: Variation of the smooth-peakon with the amplitude 1/6 at initial time.

smooth-peakons is concentrated, and the interval in which the mass of second derivative of smooth-peakons is strictly negative. In the sequel of this paper, the notation $\alpha \simeq \beta$ means that $0.9 \times \beta \le \alpha \le 1.1 \times \beta$. We set, for any $z \in \mathbb{R}$,

$$\Theta_z = [z - 6.7, z + 6.7], \text{ where } 6.7 \simeq \ln\left(\frac{20}{20 - \sqrt{399}}\right),$$
(3.14)

and

$$\mathcal{V}_z = \left[z - \ln\sqrt{2}, z + \ln\sqrt{2} \right]. \tag{3.15}$$

One can clearly see that \mathcal{V}_0 is a subset of Θ_0 (since $20/(20 - \sqrt{399}) > \sqrt{2}$). We chose the values ± 6.7 such that $\rho_c(\pm 6.7) \simeq c/2400 \simeq 4.1 \times 10^{-4}c$ as in [6]. Also, we have $\rho'_c(-6.7) = -\rho'_c(6.7) \simeq 4.1 \times 10^{-4}c$ and $\rho''_c(\pm 6.7) \simeq 4.1 \times 10^{-4}c$. Then ρ_c , ρ'_c and ρ''_c are very small with respect to the amplitude c on $\mathbb{R} \setminus \Theta_0$. We claim the following result.

Lemma 3.4 (Uniqueness of the local maximum). Let $u \in L^2(\mathbb{R})$, with $y = (1 - \partial_x^2)u \in \mathcal{M}^+(\mathbb{R})$, that satisfies (3.6) for some $z \in \mathbb{R}$. There exist $\varepsilon_0 > 0$ only depending on the speed c, such that if $0 < \varepsilon < \varepsilon_0$,

then the function $v = (4 - \partial_x^2)^{-1}u$ admits a unique local extremum on Θ_z . This extremum is a maximum, and it holds

$$v(x) \le \frac{c}{300}, \quad \forall x \in \mathbb{R} \setminus \Theta_z,$$
(3.16)

$$u(x) \le \frac{c}{300}, \quad \forall x \in \mathbb{R} \setminus \Theta_z.$$
 (3.17)

(3.18)

Proof. The key is to study the impact of the assumption $y \in \mathcal{M}^+(\mathbb{R})$ on v. First, let us show that $|v_x| \leq 2v$ on \mathbb{R} . We recall that from the assumption $y \geq 0$, we have $u \geq 0$ and $v \geq 0$ on \mathbb{R} . According to the definition of v, we have for all $x \in \mathbb{R}$,

$$v(x) = \frac{e^{-2x}}{4} \int_{-\infty}^{x} e^{2x'} u(x') dx' + \frac{e^{2x}}{4} \int_{x}^{+\infty} e^{-2x'} u(x') dx'$$

and

$$v_x(x) = -\frac{e^{-2x}}{2} \int_{-\infty}^x e^{2x'} u(x') dx' + \frac{e^{2x}}{2} \int_x^{+\infty} e^{-2x'} u(x') dx',$$
$$|v_x(x)| \le 2v(x), \quad \forall x \in \mathbb{R}.$$
(3.18)

which yields

Second, let us show that $u \leq 6v$ on \mathbb{R} . Using the Fourier transform, one can check that

$$(1 - \partial_x^2)^{-1} (4 - \partial_x^2)^{-1} (\cdot) = \mathcal{F}^{-1} \left[\frac{1}{3(1 + \omega^2)} - \frac{1}{3(4 + \omega^2)} \right] (\cdot)$$
$$= \frac{1}{3} (1 - \partial_x^2)^{-1} (\cdot) - \frac{1}{3} (4 - \partial_x^2)^{-1} (\cdot), \tag{3.19}$$

and one can rewrite v as

$$v = (4 - \partial_x^2)^{-1} (1 - \partial_x^2)^{-1} y = \frac{1}{3} (1 - \partial_x^2)^{-1} y - \frac{1}{3} (4 - \partial_x^2)^{-1} y.$$
(3.20)

Then for all $x \in \mathbb{R}$,

$$\begin{aligned} u(x) - 6v(x) &= -(1 - \partial_x^2)^{-1} y(x) + 2(4 - \partial_x^2)^{-1} y(x) \\ &= -\frac{1}{2} \int_{\mathbb{R}} e^{-|x - x'|} y(x') dx' + \frac{1}{2} \int_{\mathbb{R}} e^{-2|x - x'|} y(x') dx' \\ &\le 0, \end{aligned}$$
(3.21)

since $e^{-2|\cdot|} \leq e^{-|\cdot|}$ on \mathbb{R} .

We are now ready to prove the uniqueness of local maxima in Θ_z . Let us first study the sign of v_{xx} on \mathcal{V}_z . One can easy check that for all $x \in \mathcal{V}_0$,

$$\rho_c''(x) \le \frac{\sqrt{2} - 2}{6}c. \tag{3.22}$$

Then, combining (3.8) and (3.22), taking $0 < \varepsilon < \varepsilon_0$ with $\varepsilon_0 \ll 1$, we have for all $x \in \mathcal{V}_z$,

$$v_{xx}(x) \le rac{\sqrt{2}-2}{6}c + O(\varepsilon^{1/4}) \le rac{\sqrt{2}-2}{600}c < 0,$$

which implies that v_x is strictly decreasing on \mathcal{V}_z . Let us study the sign of v_x on $\Theta_z \setminus \mathcal{V}_z$. One can easily check that

$$\rho_c'\left(-\ln\sqrt{2}\right) = \frac{\sqrt{2}-1}{6}c \text{ and } \rho_c'\left(\ln\sqrt{2}\right) = -\frac{\sqrt{2}-1}{6}c,$$
(3.23)

and that $\rho'_c(x) \ge 10^{-4}c$ for all $x \in [-6.7, -\ln\sqrt{2}]$. Then using (3.9) and taking $0 < \varepsilon < \varepsilon_0$ with $\varepsilon_0 \ll 1$, we have $v_x(x) \ge 4 \times 10^{-5}c > 0$ for all $x \in [z - 6.7, z - \ln\sqrt{2}]$. Proceeding in the same way, we obtain $v_x(x) \le -4 \times 10^{-5}c < 0$ for all $x \in [z + \ln\sqrt{2}, z + 6.7]$. Since v_x is strictly decreasing on \mathcal{V}_z and changes sign, v_x vanishes once on \mathcal{V}_z and thus on Θ_z . Hence, v admits a single local extremum on Θ_z , which is a maximum since $v_{xx} < 0$ on \mathcal{V}_z .

Now, using that ρ_c is increasing on \mathbb{R}^- , (3.9) and taking $0 < \varepsilon < \varepsilon_0$ with $\varepsilon_0 \ll 1$, it holds for all $x \in]-\infty, z-6.7[$,

$$v(x) = \rho_c(x-z) + O(\varepsilon^{1/4}) \le \frac{c}{2400} + O(\varepsilon^{1/4}) \le \frac{c}{300}.$$

Proceeding in the same way for $x \in [z + 6.7, +\infty)$, we obtain (3.16).

Combining (3.8), (3.21) and proceeding as for the estimate (3.16), we get (3.17). Note that $\varphi_c(\pm 6.7) \simeq 1.2 \times 10^{-3}c$. This completes the proof of the lemma.

Under the assumptions of Lemma 3.4, v has got a unique point of global maximum on \mathbb{R} . In the sequel of this section, we will denote by ξ this point of global maximum and we set $M = v(\xi) = \max_{x \in \mathbb{R}} v(x)$. The next two lemmas can be directly deduced from the similar lemmas established in [7] (see also [6]).

Lemma 3.5 (Connection between $E(\cdot)$ and M^2 ; see [7]). Let $u \in L^2(\mathbb{R})$ and $v = (4 - \partial_x^2)^{-1}u \in H^2(\mathbb{R})$. Define the function g by

$$g(x) = \begin{cases} 2v + v_{xx} - 3v_x, & x < \xi, \\ 2v + v_{xx} + 3v_x, & x > \xi. \end{cases}$$
(3.24)

Then it holds

$$\int_{\mathbb{R}} g^2(x) dx = E(u) - 12M^2.$$
(3.25)

Lemma 3.6 (Connection between $F(\cdot)$ and M^3 ; see [7]). Let $u \in L^2(\mathbb{R})$ and $v = (4 - \partial_x^2)^{-1} u \in H^2(\mathbb{R})$. Define the function h by

$$h(x) = \begin{cases} -v_{xx} - 6v_x + 16v, & x < \xi, \\ -v_{xx} + 6v_x + 16v, & x > \xi. \end{cases}$$
(3.26)

Then it holds

$$\int_{\mathbb{R}} h(x)g^2(x)dx = F(u) - 144M^3.$$
(3.27)

We can now connect the conservation laws.

Lemma 3.7 (Connection between $E(\cdot)$ and $F(\cdot)$). Let $u \in L^2(\mathbb{R})$, with $y = (1 - \partial_x^2)u \in \mathcal{M}^+(\mathbb{R})$, that satisfies (3.6) for some $z \in \mathbb{R}$. There exists $\varepsilon_0 > 0$ only depending on the speed c, such that if $0 < \varepsilon < \varepsilon_0$, then it holds

$$M^{3} - \frac{1}{4}E(u)M + \frac{1}{72}F(u) \le 0.$$
(3.28)

Proof. The key is to show that $h \leq 18M$ on \mathbb{R} . Note that by (3.9) we know that $18M \geq c/4$ and that Lemma 3.4 ensures that $\xi \in \theta_z$ for ε_0 small enough. Let us set $\lambda = z - 6.7$, $\mu = z + 6.7$, and rewrite the function h as

$$h(x) = \begin{cases} -v_{xx} - 6v_x + 16v, & x < \lambda, \\ u - 6v_x + 12v, & \lambda < x < \xi, \\ u + 6v_x + 12v, & \xi < x < \mu, \\ -v_{xx} + 6v_x + 16v, & x > \mu. \end{cases}$$

If $x \in \mathbb{R} \setminus \Theta_z$, using that $v_{xx} = 4v - u$, (3.16) and (3.17), it holds

$$h \le |v_{xx}| + 6|v_x| + 16v \le u + 32v \le \frac{c}{9} \le 18M.$$

If $\lambda < x < \xi$, then $v_x \ge 0$, and using that $u \le 6v$ on \mathbb{R} , we have

$$h = u - 6v_x + 12v \le 18v.$$

If $\xi < x < \mu$, then $v_x \leq 0$, and similarly using that $u \leq 6v$ on \mathbb{R} , we get

$$h = u + 6v_x + 12v \le 18v.$$

Therefore, it holds

$$h(x) \le 18 \max_{x \in \mathbb{R}} v(x) = 18M, \quad \forall x \in \mathbb{R}.$$
(3.29)

Now, combining (3.25), (3.27) and (3.29), we get

$$F(u) - 144M^3 = \int_{\mathbb{R}} h(x)g^2(x)dx \le \|h\|_{L^{\infty}(\mathbb{R})} \int_{\mathbb{R}} g^2(x)dx \le 18M(E(u) - 12M^2),$$

and we obtain the lemma.

Proof of Theorem 3.1. We argue as El Dika and Molinet in [4]. As noticed after the statement of the theorem, it suffices to prove (3.7) assuming that $u \in L^2(\mathbb{R})$ satisfies (3.1), (3.2) and (3.6). We recall that $M = v(\xi) = \max_{x \in \mathbb{R}} v(x)$ and we set $\delta = c/6 - M$. We first remark that if $\delta \leq 0$, combining (3.4) and (3.10), it holds

$$|u - \varphi_c(\cdot - \xi)||_{\mathcal{H}} \le |E(u_0) - E(\varphi_c)|^{1/2} \le O(\varepsilon)$$

that yields the desired result. Now suppose that $\delta > 0$, that is the maximum of the function v is less than the maximum of ρ_c . Combining (3.4), (3.5) and (3.28), we get

$$M^3 - \frac{1}{4}E(\varphi_c)M + \frac{1}{72}F(\varphi_c) \le O(\varepsilon^2).$$

Using that $E(\varphi_c) = c^2/3$ and $F(\varphi_c) = 2c^3/3$, our inequality becomes

$$\left(M - \frac{c}{6}\right)^2 \left(M + \frac{c}{3}\right) \le O(\varepsilon^2)$$

Next, substituting M by $c/6 - \delta$ and using that $[M + c/3]^{-1} < 3/c$, we obtain

$$\delta^2 \le O(\varepsilon^2) \Rightarrow \delta \le O(\varepsilon). \tag{3.30}$$

Finally, combining (3.4), (3.10) and (3.30), we infer that

$$\|u - \varphi_c(\cdot - \xi)\|_{\mathcal{H}} \le C\sqrt{\varepsilon},$$

where C > 0 only depends on the speed c. This completes the proof of the stability of peakons.

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References

- Adrian Constantin and David Lannes. The hydrodynamical relevance of the Camassa-Holm and Degasperis-Procesi equations. Arch. Ration. Mech. Anal., 192(1):165–186, 2009.
- [2] Adrian Constantin and Walter A. Strauss. Stability of peakons. Comm. Pure Appl. Math., 53(5):603-610, 2000.

- [3] A. Degasperis, D. D. Kholm, and A. N. I. Khon. A new integrable equation with peakon solutions. *Teoret. Mat. Fiz.*, 133(2):170–183, 2002.
- [4] Khaled El Dika and Luc Molinet. Stability of multipeakons. Ann. Inst. H. Poincaré Anal. Non Linéaire, 26(4):1517–1532, 2009.
- [5] Joachim Escher, Yue Liu, and Zhaoyang Yin. Global weak solutions and blow-up structure for the Degasperis-Procesi equation. J. Funct. Anal., 241(2):457–485, 2006.
- [6] André Kabakouala. Stability in the energy space of the sum of N peakons for the Degasperis-Procesi equation. J. Differential Equations, 259(5):1841–1897, 2015.
- [7] Zhiwu Lin and Yue Liu. Stability of peakons for the Degasperis-Processi equation. Comm. Pure Appl. Math., 62(1):125–146, 2009.
- [8] Yue Liu and Zhaoyang Yin. Global existence and blow-up phenomena for the Degasperis-Procesi equation. *Comm. Math. Phys.*, 267(3):801–820, 2006.