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A Solid State Physics Perspective on Hybrid Perovskite Semiconductors

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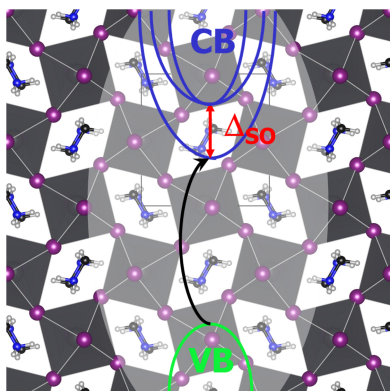
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TOC GRAPHICS



KEYWORDS Photovoltaics, optoelectronics, exciton, Bloch function, band folding, Auger effect

Biography:
Jacky Even



Jacky Even is full Professor at INSA Rennes engineering school since 1999. He received a PhD in Physics from the University of Paris VI in 1992. He was assistant professor at the Physics Department of Rennes University (1992-1999), where he gained broad experience in phase transitions and solid-state chemical reactions in molecular materials, using theoretical and experimental approaches: neutron and X-ray scattering, Raman/FTIR spectroscopy, calorimetry, among others. In 1999, he created FOTON laboratory's simulation team, to address fundamental questions on semiconductors at the atomistic level as well as to perform optoelectronic device simulation. Besides hybrid perovskite materials and colloidal nanoplatelets, his theoretical activity is now dedicated to semiconductor nanostructures, photovoltaic and light emitting (LED, laser, ...) devices for silicon photonics and optical telecommunication.

Laurent Pedesseau



Born March 12, 1975, Laurent Pedesseau obtained his MSc in condensed matter from the University of Montpellier in 2001. In 2004, he received his PhD in physics from the University of Toulouse for atomistic empirical simulations applied to Civil Engineering materials. In 2013, he was appointed as assistant-professor at FOTON laboratory (INSA Rennes) to work on III-V semiconductor nanostructures for silicon photonics, hybrid-perovskites for photovoltaics and optoelectronic device simulations for optical telecommunication.

Claudine Katan



Dr. Claudine Katan, born Hoerner, received a PhD in Physics from the University of Strasbourg in 1992, working with Pr. Albert Villaeys and Dr. Jean-Pascal Lavoine on dephasing effects in four-wave mixing experiments studied using density matrix formalism. She subsequently served

as lecturer in physics at the University of Rennes before being appointed as a CNRS Research Investigator in the Physics department in 1993. Until 2003, her research interests mainly focused on the properties of molecular charge-transfer crystals and topology of electron densities. She then joined the Chemistry department, interspersed by a two years stay at FOTON-OHM. Her current research at ISCR is devoted to theoretical approaches from modeling to state-of-the-art electronic structure calculations for structural, electronic and linear/non-linear optical properties (e.g. two-photon absorption) of molecular chromophores, but also of hybrid perovskites for light emission and photovoltaic technologies. Her theoretical research is usually conducted in close collaboration with experimentalists.

Mikaël Kepenekian



Dr. Mikaël Kepenekian graduated from the ENS de Lyon in Chemistry with a thesis work on molecular magnetism. After a PhD on bistable molecular materials conducted at ENS Lyon and CEA Grenoble under the direction of Dr. Pascale Maldivi and Prof. Vincent Robert, he joined the Centro de Investigación en Nanociencia y Nanotecnología (CIN2) in Bellaterra, Spain working with Prof. Nicolás Lorente on density functional theory applied to electron transport calculations using the non-equilibrium Green's functions formalism. In 2013, he was appointed as a CNRS Researcher at ISCR. His research is devoted to the simulation of semiconductors and interfaces such as molecules supported on surfaces.

Jean-Sébastien Lauret

After a PhD on the optical properties of carbon nanotubes and a post doc on boron nitride, he was appointed as assistant professor at ENS Cachan. There, he developed activities on the optical properties of 2D hybrid organic perovskites. He has also an activity on carbon nanostructures. J.S. Lauret is now professor at Ecole Normale Supérieure de Cachan.

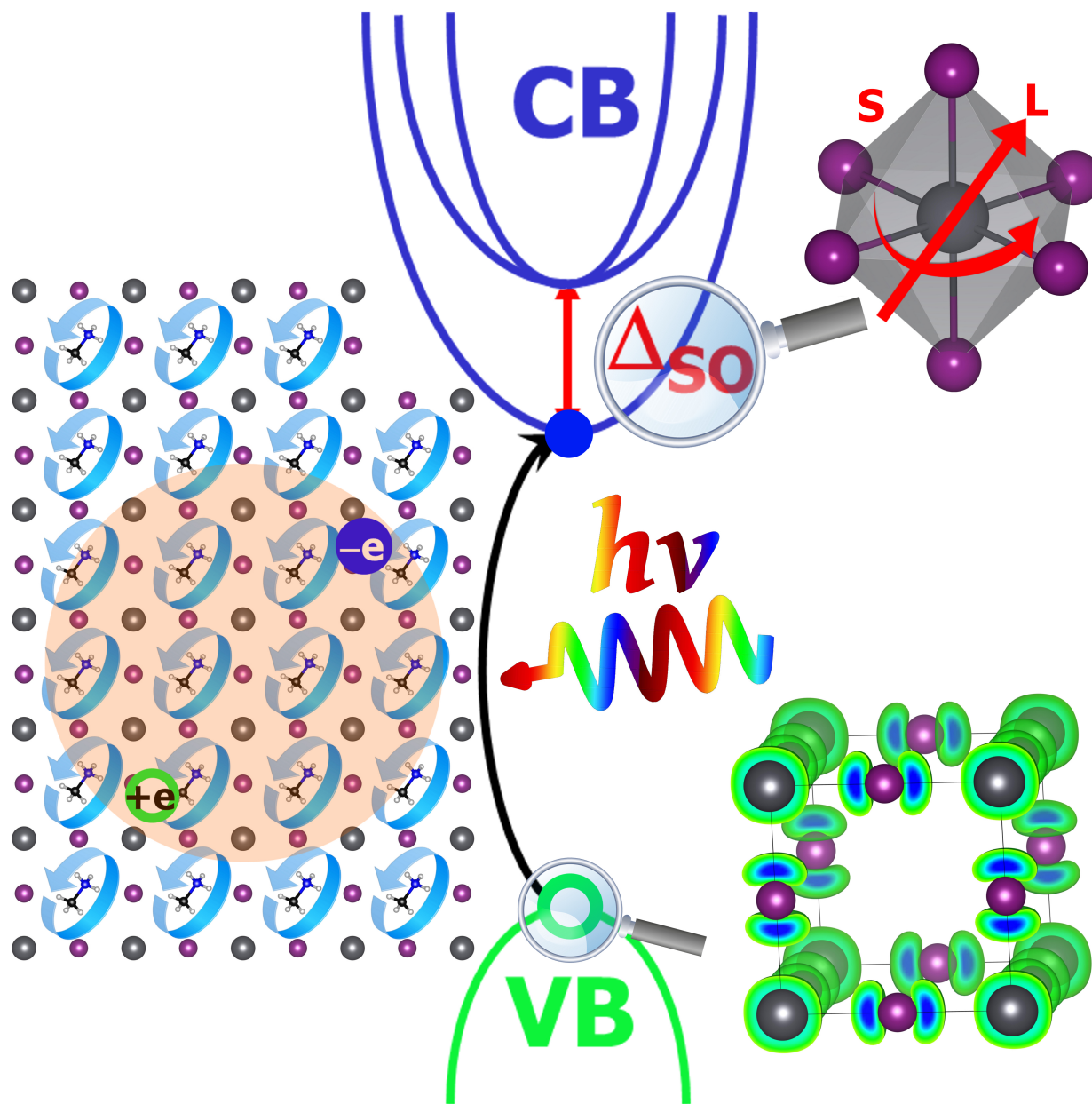
Daniel Sapori

Born on November 28th, 1991, Daniel Sapori graduated from the engineering school CPE Lyon in chemistry and chemical engineering and also obtained an MSc in nanotechnology. Now, he is a PhD student in the simulation group of the FOTON laboratory (INSA Rennes) working on electronic and optical properties of hybrid perovskites under the supervision of Jacky Even.

Emmanuelle Deleporte

Emmanuelle Deleporte is Professor at Ecole Normale Supérieure de Cachan, head of the Physics Department since 2006. Her research activities are conducted in Laboratoire Aimé Cotton. She is a specialist in optical properties of semiconductors (SCs): semimagnetic SCs, II-VI and III-V SCs, heterostructures such as quantum wells and quantum dots, hybrid perovskites. Her main interests are the study of the excitonic effects, the relaxation dynamics of the carriers, the exciton-phonon interaction and the light-matter interaction, in the framework of opto-electronics: lasers, electroluminescent diodes and photovoltaic systems.

COVER ART



ABSTRACT In this review we examine recent theoretical investigations on 2D and 3D hybrid perovskites (HOP) that combine classical solid-state physics concepts and density functional theory (DFT) simulations as a tool for studying their optoelectronic properties. Such an approach allows one to define a new class of semiconductors, where the pseudocubic high temperature perovskite structure plays a central role. Bloch states and **k.p** Hamiltonians yield new insight into the influence of lattice distortions, including loss of inversion symmetry, as well as spin-orbit coupling. Electronic band folding and degeneracy, effective masses and optical absorption are analyzed. Concepts of Bloch and envelope functions, as well as confinement potential are discussed in the context of layered HOP and 3D HOP heterostructures. Screening and dielectric confinements are important for room temperature optical properties of 3D and layered HOP, respectively. Non-radiative Auger effects are analyzed for the first time close to the electronic band gap of 3D hybrid perovskites.

1. INTRODUCTION

Most metal-halide Hybrid Organic Perovskites (HOP) have the general formulae $(R-NH_3)_nMX_m$, where R is an organic group, X a halogen atom (X=I, Br, Cl) and M a metal atom (M=Pb, Sn, Ge). Control of the stoichiometry (n, m) affords crystal packing of various dimensionalities, ranging from three-dimensional (3D, n=1 and m=3) corner-shared perovskite lattices to 0D structures with isolated inorganic octahedra. In less than 30 years, the low-cost technologies and the easy synthesis of the layered HOP (m>3) raised up the interest in the field of optoelectronics and microelectronics.¹⁻¹⁵ A scientific breakthrough for solar cells was achieved in the last three years with 3D HOP.¹⁶⁻⁵⁴ In both cases, lead-halide based HOP are currently the most investigated, but substitution of the critical metals (Pb) has to be addressed because of environmental issues.

Indeed, tin-based layered HOP exhibit hole transport and natural doping, and were shown to be suitable for the demonstration of field effect transistor operation.¹⁰⁻¹² First results on tin-based solar cells were recently obtained with 3D HOP.^{29,41,42}

The quantum and dielectric confinement expected in the 2D layered HOPs (n=2, m=4) prompted active research to develop optoelectronic devices with enhanced performances.¹¹ Since 2009, 3D HOPs (n=1, m=3), with relatively small organic cations, have been suggested as a novel class of low-cost solution-processable materials for high efficiency hybrid photovoltaic cells. Tremendous progress has been made in a short period of time since 2012,²¹⁻³⁵ leading to both meso-superstructured solar cells and solid thin-film planar heterojunctions with record solar-to-electrical power conversion efficiencies. In 2013, NREL recognized the so-called ‘Perovskite Cells’ as a new and specific technology different from dye-sensitized solar cell (DSSC).³⁶⁻⁴⁶ Recent theoretical works show that better understanding could benefit from tools and concepts developed both in the fields of organic optoelectronics and conventional semiconductors.⁵⁵⁻¹²¹ Indeed, since the initial use of 3D HOP as the sensitizer in conventional DSSC, technological developments have simultaneously led to a gradual shift of fundamental issues from materials chemistry to solid-state physics. This can be envisioned as a change of paradigm.⁷⁸ Further improvements should make the best profits from both fields. As quoted by H. Snaith in a 2013 review: “*Solar cells based on perovskite absorbers promise to break the prevailing paradigm by combining both ultimately low cost and high efficiency*”.³⁴ The efficiency of hybrid perovskites solar cells rose rapidly from 3.8% in 2009 and 6.5% in 2011 with DSSC-like technologies to about 10% in 2012 using new concepts. These record values steeply increased to about 15% in 2013 thanks to new deposition techniques.^{22,23,27,28} At the end of 2014, optimized conversion efficiency amounts to 19.3%³⁹ and then 20% in the last (2014) NREL efficiency chart.¹²² A quick

insight into the possible applications of hybrid perovskites for solar cells can be gained from a couple of recent review papers.³³⁻³⁸ Besides, the feasibility of light emitting diodes or lasers with 3D HOP were recently investigated by several research groups.⁴⁷⁻⁵⁰

Clearly, understanding physical properties of operational HOPs and design of novel devices with improved performances may greatly benefit from theoretical work. Prior to 2013, most theoretical studies either focused on layered HOP,^{55,57,61} purely inorganic 3D perovskites,^{56,60,62-65} and 3D HOP,^{55,56,58,59} but, in the latter case, missing inherent physical properties such as effect of spin-orbit coupling. In 2013,⁶⁶⁻⁷² a few papers have appeared on 3D HOP in relation with the breakthrough in the field of photovoltaics, followed by a large number of papers in 2014: a near-exhaustive list is given in the reference section.⁷³⁻¹²¹ Currently, most of the theoretical work relies on the use of Density Functional Theory (DFT) based simulations. They have proved their effectiveness to analyze qualitatively band gaps, effective masses or optical absorption of 2D and 3D materials. It is however difficult to make theoretical predictions or relevant comparisons with physical parameters derived from experiments, especially for structures suited from a device perspective such as HOP heterostructures. In the case of conventional semiconductors materials and heterostructures (*e.g.* silicon and III-V semiconductors), knowledge often profit most from semi-empirical approaches based on basic solid-state physics concepts and symmetry considerations.^{61,62,66,74,75,123} In such a case, DFT including many-body corrections affords a tool to assess the validity of the semi-empirical approach. However, comparison is limited to particular situations (*e.g.* small heterostructures) that remain computationally feasible.

This paper aims to review our recent work on HOP that is essentially based on concepts and tools already well deployed in the fields of optoelectronics and conventional semiconductors. But, from the introduction we would like to emphasis major contributions from several other

research groups, who were very active from the start in 2013,⁶⁶⁻⁷⁰ and in early 2014.⁹⁴⁻¹⁰⁰ Close attention will be paid on solid-state physics concepts in the context of HOP. The validity and limitations of mainstream approaches developed for conventional semiconductors will be discussed. Properties that make HOP a new and distinct class of semiconductors will also be emphasized.

The paper is organized as follows. First and foremost, to ensure easier understanding of underlying concepts discussed in the first two sections, effect of spin orbit coupling (SOC) will be temporary left aside. But readers must keep in mind that SOC is essential to capture the main physics in this class of hybrid materials. In Section 2 we introduce the pseudocubic phase as a reference for the investigation of 3D HOP, introduce the corresponding **k.p** Hamiltonian and list the major differences between conventional and HOP semiconductors. Section 3 is devoted to clarify and distinguish the impact of Brillouin Zone (BZ) folding and lattice strain. Next, importance of SOC will be considered in Section 4 and discussed both for 3D and 2D HOP. In these first four sections, particular attention will be paid to symmetry and how these effects can be accounted for within the **k.p** method. Section 5 summarizes available information on possible loss of inversion symmetry in crystal structures, and addresses its interplay with SOC that suggests potential applications in spintronics. Sections 6 and 7 deal with quantum confinement and band alignment in layered and 3D HOP, respectively. The nature of photoexcited species (excitons versus free carriers) is analyzed in Section 8 with a special focus on dielectric screening and confinement. Section 9 presents a discussion on the importance of Auger effects, including comparison with III-V semiconductors. Finally, main conclusions are drawn in Section 10.

2. THE PSEUDOCUBIC PEROVSKITE STRUCTURE AS A REFERENCE FOR 3D HOP

Bulk structures of well-known conventional semiconductors adopt either diamond cubic (Fd3m, space group no 227) and zinc-blende (F43m, no 216) (Si, Ge, C, GaAs, InP, CdSe ...) unit cells, or hexagonal graphite (P6₃/mmc, no 194) and würtzite (P6₃mc, no 186) (C, GaN, ZnO, CdSe, ZnS...) lattices. For zinc-blende and würtzite structures, many of the structural properties are similar to that of diamond and graphite, but loss of inversion symmetry has a subtle impact on the optoelectronic and piezoelectric properties of the crystal.

Let's start with GaAs, which is a reference among conventional cubic semiconductors for optoelectronic applications that exhibits both good optical and transport properties. As stated in the introduction we remind the reader that, for simplicity, in this and the next sections SOC will be disregarded. The density of states (DOS) of cubic GaAs computed within the GGA (Generalized Gradient Approximation) with the SIESTA code is reported in Figure 1a. It shows that both valence (VB) and conduction bands (CB) are associated to hybridized atomic *s*- and *p*-orbitals of Ga and As atoms. In order to take advantage of translational symmetry, it is useful to describe the energy eigenfunctions in a bulk crystal as Bloch waves:

$$\psi_{n,\vec{k}}(\vec{r}) = e^{i\vec{k}\cdot\vec{r}} u_{n,\vec{k}}(\vec{r}) \quad , \quad (1)$$

where *n* is the band index, *r* the position vector, *k* the wavevector inside the BZ and *u_{n,k}(r)* the periodic part of the eigenfunction. The corresponding energy dispersion curves are represented within the BZ (Figure 1b). A plane wave representation of the Bloch waves would be possible, but Wannier functions based on local atomic orbitals are a good alternative to provide an intuitive avenue of chemical bonding in solids starting from DFT or tight binding results. Triply degenerated electronic states form the valence band maximum (VBM) at the center of the BZ, Γ (Figure 1b). The non-degenerate conduction band minimum (CBM) of GaAs also shows up at Γ

and leads to a direct electronic band gap that, at room temperature, is almost the same as that of $\text{CH}_3\text{NH}_3\text{PbI}_3$ (1.5 eV),³³ the prototype 3D HOP for solar cells.

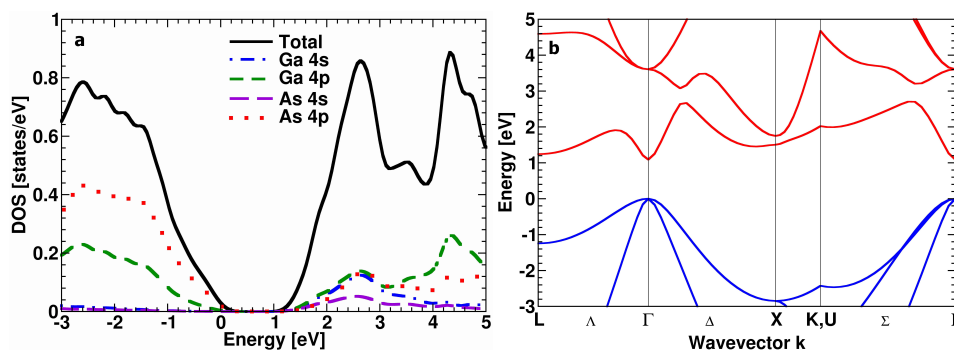


Figure 1. (a) Density of states (DOS) and (b) electronic band dispersion diagram of GaAs in the zinc-blende structure. The energy at the top of the valence band (VB) at Γ point is set to 0 eV. The DOS close to the band gap is mostly associated to hybridized s and p states of the Ga and As atoms. The conduction bands (CB) and VB are represented respectively as red and blue lines in the band dispersion diagram (b). The computation is performed within the GGA with the SIESTA code. Spin-orbit coupling (SOC) is not considered.

However the situation is more complicated in 3D HOP than in semiconductors like GaAs, due to orientational disorder of the organic cation. Electronic properties of the cubic phase of the prototype HOP $\text{CH}_3\text{NH}_3\text{PbI}_3$ are best captured taking its all-inorganic analog, CsPbI_3 . In fact, many effects, such as strain, phase transitions, relativistic effects, local distortions due to the organic cations, loss of inversion symmetry, etc, can be considered as perturbations to a certain extent (vide infra). Figure 2a represents a real space 3D view of the $\text{Pm}\bar{3}\text{m}$ reference cubic crystal structure of metal-halide perovskites of general formulae AMX_3 where $\text{A}=\text{Cs}$, CH_3NH_3^+ , $\text{M}=\text{Pb}$, Sn and $\text{X}=\text{I}$, Br , Cl . The reciprocal space 3D view (Figure 2b) shows the first BZ of the $\text{Pm}\bar{3}\text{m}$ space group, with some points of high symmetry. The electronic band structure of the cubic ($\text{Pm}\bar{3}\text{m}$, no 221) phase of $\text{CH}_3\text{NH}_3\text{PbI}_3$ is shown on Figure 2c. An upward energy shift of 0.3eV has been applied to match the experimental bandgap value at R. The CH_3NH_3^+ cation is located at the center of the cube with an averaged position. It has been replaced by a Cs^+ cation

for the computation of the band structure. The electronic structure of the reference cubic phase of $\text{CH}_3\text{NH}_3\text{PbI}_3$ reveals two main differences when compared with GaAs (Figure 1):⁶⁶

- Direct band gap located at the high symmetry point R instead of Γ ;
- Inverted band structure: the triply degenerate electronic state is located at CBM instead of VBM.

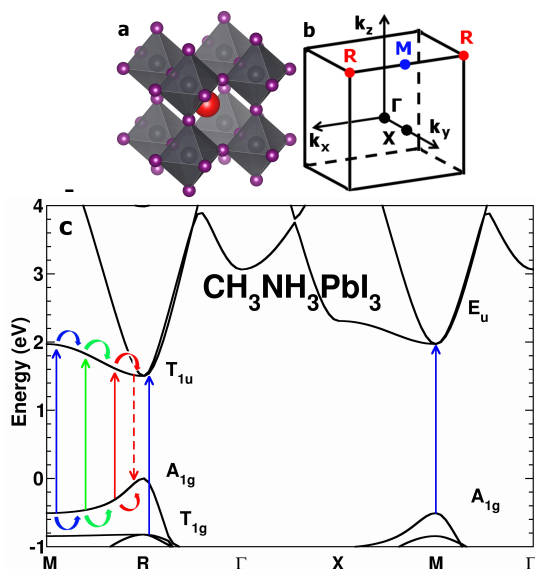


Figure 2. (a) Real space 3D view of the $\text{Pm}\bar{3}\text{m}$ reference cubic crystal structure of metal-halide hybrid perovskites of general formulae AMX_3 where $\text{A}=\text{Cs}$, CH_3NH_3^+ , $\text{M}=\text{Pb}$, Sn and $\text{X}=\text{I}$, Br , Cl . In hybrid perovskites, the CH_3NH_3^+ cation is located at the center of the cube with an averaged position sketched by the red ball. (b) Reciprocal space 3D view showing the first BZ of the $\text{Pm}\bar{3}\text{m}$ space group. Points of high symmetry in the cubic BZ are indicated by conventional letters: Γ denotes the origin of the BZ; X is the center of a square face at the BZ boundary, M is a center of a cube edge; and R are vertices of the cube. (c) Electronic band structure for the high temperature cubic $\text{Pm}\bar{3}\text{m}$ phase of $\text{CH}_3\text{NH}_3\text{PbI}_3$ without SOC at the LDA (Local Density Approximation) level of theory. An upward energy shift of 0.3eV has been applied to match the experimental bandgap value at R. Irreducible representations obtained from a $\text{Pm}\bar{3}\text{m}$ simple group analysis, are given at R and M points for the electronic states close to the band gap. Vertical arrows show various possible optical transitions close to the band gap energy. Optical transitions along the line between the M and R points generate carriers that easily relax toward the R point. Reprinted with permission from ref 75. Copyright 2014 American Chemical Society.

The dispersion curves close to the critical point R of the BZ can be analysed within a multiband effective mass model ($\mathbf{k}\cdot\mathbf{p}$ method)⁷⁴ as a function of the \mathbf{k} -vector, which represents the

difference with the wavevector at R point \mathbf{k}_R . The corresponding \mathbf{k} -dependent Hamiltonian for the triply degenerate CBM deduced from a symmetry analysis reads:^{66,74,75}

$$H_{\bar{k}} = E_{CBM}^0 + \begin{pmatrix} Lk_x^2 + M(k_y^2 + k_z^2) & Nk_x k_y & Nk_x k_z \\ Nk_x k_y & Lk_y^2 + M(k_x^2 + k_z^2) & Nk_y k_z \\ Nk_x k_z & Nk_y k_z & Lk_z^2 + M(k_x^2 + k_y^2) \end{pmatrix} \quad (1)$$

for the simple group T_{1u} vectorial irreducible representation (IR) of the CBM at R. L , M and N are parameters that can be extracted either from experiment or DFT calculations. The basis Bloch functions at R for this IR are associated to p_x, p_y, p_z atomic orbitals of Pb and can thus be labeled $|X\rangle, |Y\rangle, |Z\rangle$ by analogy to the Bloch functions at VBM in diamond and zinc-blende lattices. In the Pm3m cubic structure at R, VBM corresponds to the totally symmetric A_{1g} IR of the simple group. The basis Bloch function for this IR is associated to p_x, p_y, p_z orbitals of iodine and the s orbital of lead and can thus be labeled $|S\rangle$ by analogy to the Bloch functions at the CBM in zinc-blende lattices with direct band gap.¹²³

The symmetry analysis of the electronic states at R and M points allows identifying the optically allowed transitions. The strength of the optical transitions can be evaluated from energy parameters for the optical matrix elements, obtained for each pair of [CB,VB] states from the expression: $\frac{2}{m_e} |\langle \psi_{CB} | -i\hbar\nabla | \psi_{VB} \rangle|^2$. In the case of the fundamental optical transition they can be related to the Kane energy, which is commonly used in semiconductor physics.¹²³ For $\text{CH}_3\text{NH}_3\text{PbI}_3$, they amount to 40 and 120 eV for the fundamental transitions at R and M, respectively,⁷⁵ but next sections will show that strength of optical absorption crucially depends on other effects, in particular SOC that modifies both band degeneracies and effective masses.

3. BRILLOUIN ZONE FOLDING FROM THE PSEUDO-CUBIC PHASE IN 3D HOP EVIDENCING EFFECT OF LATTICE STRAIN

At high temperature, experimental crystal structures of HOP do not reveal strictly ordered and symmetric phases as a consequence of molecular symmetry of the organic cations that do not fit the site symmetry of the lattice. For 3D HOP, the reference high temperature phase is the cubic phase of CsPbI_3 ($Z=1$). In fact, for $\text{CH}_3\text{NH}_3\text{PbX}_3$ a similar pseudo-cubic $\text{Pm}3\text{m}$ phase is observed by diffraction measurements at high temperature, but the positions of all the atoms in the organic molecules are dynamically averaged and occupy the same special Wyckoff position at the center of the cubic cell. A first ad-hoc attempt is to completely locate the organic cation within the cubic cell. However, this procedure leads to an artificial reduction of the space group symmetry and corresponds to strong long range dipole-dipole interactions in the lattice, which are not expected to play a role in the static limit. Indeed, the high symmetry of the $\text{Pm}3\text{m}$ phase is related to a dynamical disorder. A more practical way to make a rigorous theoretical description including all the atoms in an ordered lattice, is to consider the low temperature orthorhombic phases (Pnma , no 62) of $\text{CH}_3\text{NH}_3\text{PbX}_3$ 3D HOP compounds ($Z=4$ for $X=\text{I}$ or Br and $Z=8$ for $X=\text{Cl}$). The electronic band gap remains direct but shows up at Γ and the number of bands is multiplied by a factor of 4 or 8 with respect to the cubic case. At a first sight, the electronic band structure of the Pnma phase of $\text{CH}_3\text{NH}_3\text{PbI}_3$ (Figure 3a) appears very different from that of the reference cubic phase (Figure 3c). In order to properly understand the connection between the two diagrams, it is necessary to treat both phases on an equal footing. When replacing the organic cations by cesium atoms located at nitrogen positions, the electronic band diagram close to the band gap of HOP remains unchanged.^{61,66,73} This remarkable property is due to the fact that the CBM and VBM electronic states are mostly associated to atomic orbitals of the inorganic octahedra and that the Cs^+ cations correctly mimic the ionic interactions. This computational trick cannot be used to predict the structural properties of $\text{CH}_3\text{NH}_3\text{PbX}_3$, which depend on other

interactions not accounted for by this substitution. Nevertheless, the electronic band diagram of the pseudo-cubic Pm3m phase of CH₃NH₃PbX₃ can be conveniently studied by putting a cesium atom at the center of the cell (average position of the N atom in the Pm3m lattice) and using the experimental

lattice

constants.

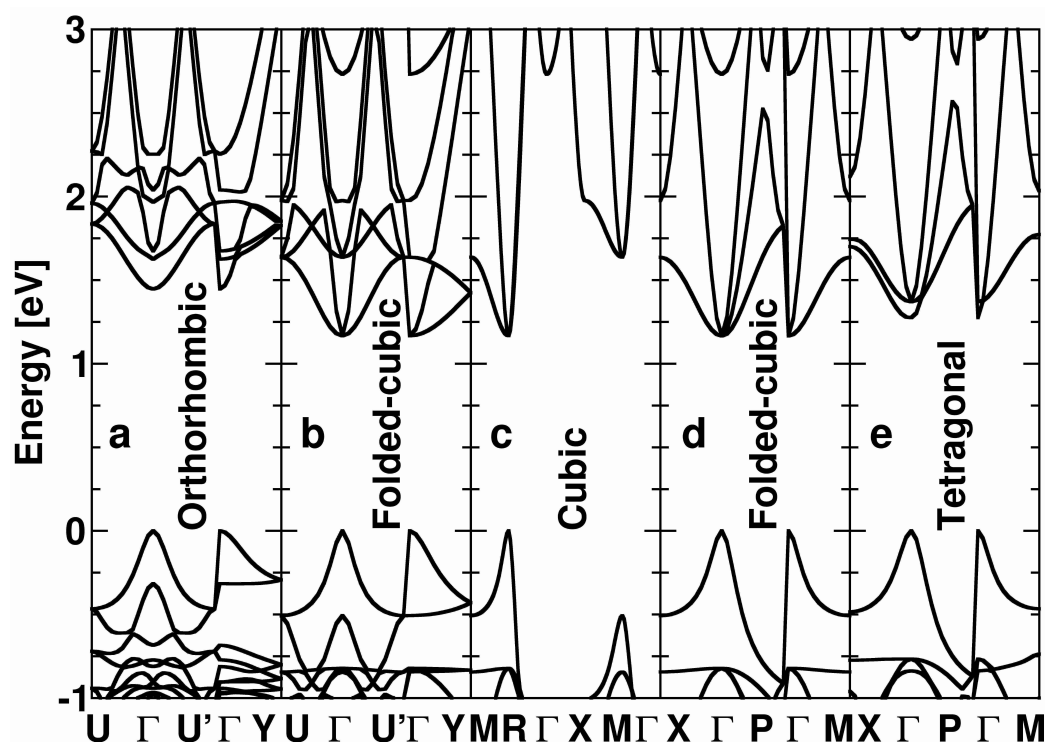


Figure 3. Band folding from the pseudo-cubic phase to the low temperature phases of CH₃NH₃PbI₃. (a-c) Electronic band dispersion diagram for the Pm3m cubic phase of CH₃NH₃PbI₃: (c) in its BZ, (b) folded back to the BZ of the Pnma orthorhombic phase, and (a) for the real Pnma orthorhombic phase of CH₃NH₃PbI₃. (d) Electronic band dispersion diagram for the Pm3m cubic phase folded back to the BZ of the I4/mcm tetragonal phase and (e) for the real I4/mcm tetragonal phase. The computation is performed at the LDA level of theory and SOC is not considered.

Next, the effect of the high to low temperature (Pm3m/Pnma) cell transformation in CH₃NH₃PbI₃ can be explored by considering a $(\vec{a} + \vec{b}, 2\vec{c}, \vec{a} - \vec{b})$ supercell of the cubic lattice while keeping the atoms at the special positions of the Pm3m phase (Figure 3b, folded-cubic band structure). As the cell size is increased by a factor of 4 ($Z=4$), it leads to a reduction of the BZ volume by a

factor of 4. The electronic states close to the electronic band gap are folded from R of the cubic phase BZ to Γ of the orthorhombic supercell BZ. This band folding effect clarifies the seeming complexity of the band diagram in the low temperature Pnma phase (Figure 3a). In addition, electronic states from M of the Pm3m BZ are also folded back to Γ and mixed with R-states at the Γ point of the low temperature Pnma BZ. This procedure reveals that, besides electronic band folding effect, the most important transformation of the diagram is due to atomic displacements.⁶⁶ The increase of the electronic band gap is thus clearly induced by lattice distortions that are mainly related to tilts of the inorganic octahedra and space group symmetry reduction.⁸³ Lattice transformations simultaneously lift the CBM degeneracy (Figure 3a). This is consistent with the group-subgroup relationship between the cubic Pm3m phase and the orthorhombic Pnma phase. Similarly, effects of band folding can be analyzed for the phase transition between the cubic Pm3m phase and, the tetragonal I4/mcm (Figure 3c-e). Here again, lattice distortions induced by octahedral tilt leads only to a partial lifting of the CBM degeneracy and are consistent with the space group symmetry reduction.⁷⁴

More generally, the influence of strain (ϵ_{ij} tensor) and degeneracy lifting can be understood and treated empirically within the $|X\rangle, |Y\rangle, |Z\rangle$ basis of the simple group T_{1u} vectorial IR of the cubic CBM at R,⁷⁴ by adding the following strain perturbation Hamiltonian to the multiband effective mass ($\mathbf{k}\cdot\mathbf{p}$) Hamiltonian (equation 1):

$$H_{\epsilon} = \begin{pmatrix} l\epsilon_{xx} + m(\epsilon_{yy} + \epsilon_{zz}) & n\epsilon_{xy} & n\epsilon_{xz} \\ n\epsilon_{xy} & l\epsilon_{yy} + m(\epsilon_{xx} + \epsilon_{zz}) & n\epsilon_{yz} \\ n\epsilon_{xz} & n\epsilon_{yz} & l\epsilon_{zz} + m(\epsilon_{xx} + \epsilon_{yy}) \end{pmatrix} \quad (2)$$

where ϵ_{ij} components of the lattice strain tensors are used and l , m and n parameters are deformation potentials that can be derived from experimental data or DFT results.

4. IMPORTANCE OF SPIN ORBIT COUPLING IN 3D AND LAYERED HOP

A giant SOC operates on the CB, which dominates the band gap of lead-based HOP, both for the high (Figure 4) and low temperature phases.^{66,74} About three times smaller for $\text{CH}_3\text{NH}_3\text{SnX}_3$ than for $\text{CH}_3\text{NH}_3\text{PbX}_3$, it remains sizeable in tin-based perovskites and the following conclusions apply both to lead and tin HOP.⁷⁴ Recent analysis of tri-chlorides, $\text{CH}_3\text{NH}_3\text{MCl}_3$ (M= Pb, Sn, Ge), show that SOC splittings are consistent with metal atomic energy level tables with an increase down group-14 of the periodic table.⁷⁷ Moreover, SOC interactions stemming from halogens are shown to be sizable, especially for iodine-based perovskites, consistently with atomic data and computed atomic charges from the Atom In Molecules (AIM) quantum theory.⁷⁷

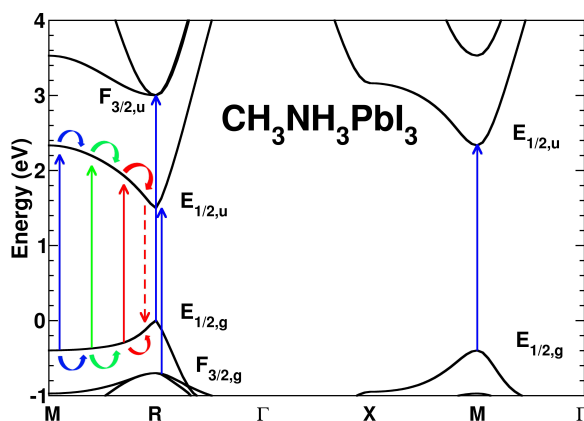


Figure 4. Electronic band structure for the high temperature cubic Pm3m phase of CH₃NH₃PbI₃ with SOC at the LDA level of theory. An upward energy shift of 1.4eV has been applied to match the experimental bandgap value at R. Irreducible representations obtained from a Pm3m double group analysis, are given at R and M points for the electronic states close to the band gap. Vertical arrows show various possible optical transitions close to the band gap energy. Optical transitions along the line between the M and R points generate carriers that easily relax toward the R point. Reprinted with permission from ref 75. Copyright 2014 American Chemical Society.

In order to account for spin effect and lattice symmetry, a double space group description is needed.⁷⁵ The starting point is the $|X\rangle, |Y\rangle, |Z\rangle$ basis of the simple group T_{1u} vectorial IR of the cubic CBM at R and corresponding energy E_{CBM}^0 . At R, the spinor state basis at CBM converts into $|X \uparrow\rangle, |Y \uparrow\rangle, |Z \uparrow\rangle, |X \downarrow\rangle, |Y \downarrow\rangle, |Z \downarrow\rangle$ and SOC leads to an energy splitting Δ_{SO} . CBM becomes a two-fold degenerate spin-orbit split-off (SO) state whose energy reads $E_{(SO)} = E_{CBM}^0 - 2\Delta_{SO}/3$. The remaining four-fold degenerate CB states undergo an upward energy shift: $E = E_{CBM}^0 + \Delta_{SO}/3$.

In other words, the simple group T_{1u} vectorial representation of the CBM at R is split in a doubly degenerated E_{1/2u} (SO states) and a fourfold degenerated F_{3/2u} state as illustrated in Figure 4. The ground state isotropic optical transition is thus predicted at R between doubly degenerated E_{1/2g} VBM and E_{1/2u} CBM states. At R, a series of other transitions are optically allowed, the lowest in energy corresponding to two secondary transitions F_{3/2g} → E_{1/2u} and E_{1/2g} → F_{3/2u} (Figure 4). The transition between the doubly degenerated E_{1/2g} VB and E_{1/2u} CB states at M is also optically allowed and polarized with transverse isotropy. A total isotropic activity is found for the transition at the M point in the pseudo-cubic phase by considering that the star of the **k**_M point contains three arms.⁷⁵ Inclusion of SOC leads to a dramatic reduction of oscillator strengths that decreases to 17 eV for the fundamental transition at R.⁷⁵ In terms of Kane energy, it drops to ca 6 eV which is about four times smaller than that of GaAs, thus evidencing an additional difference between HOP and conventional semiconductors.

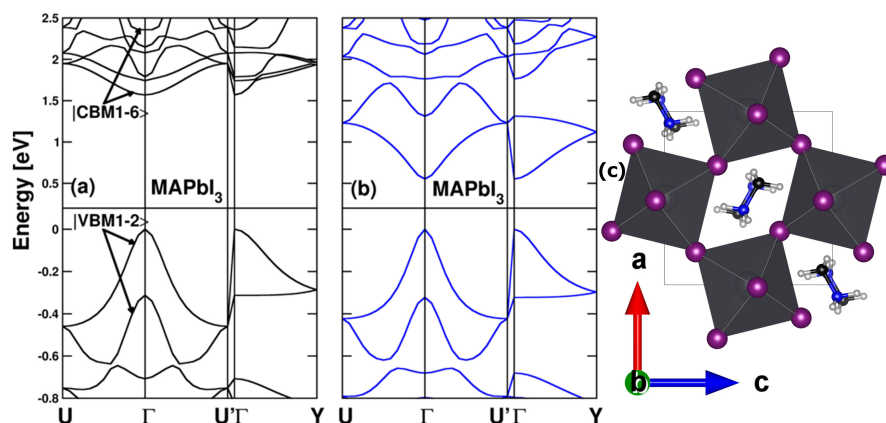


Figure 5. Electronic band structure of $\text{CH}_3\text{NH}_3\text{PbI}_3$ (MAPbI₃) in the low temperature orthorhombic Pnma phase (a) without and (b) with SOC at the GGA level of theory with the ABINIT code. The energy levels are referenced to the valence band maximum. (c) Overview of the crystal structure. Reprinted with permission from ref 66. Copyright 2013 American Chemical Society.

Impact of SOC can be further analyzed by comparing the electronic structures of pseudocubic Pm3m, orthorhombic (Figure 5) and tetragonal phases of $\text{CH}_3\text{NH}_3\text{PbI}_3$.⁷⁴ These dispersion diagrams can be empirically rationalized by considering the spinor state basis and adding a perturbation Hamiltonian accounting for SOC to the $\mathbf{k}\cdot\mathbf{p}$ multiband Hamiltonian of the CBM close to the critical point R of the cubic BZ (equation 1) and the perturbative strain Hamiltonian (equation 2). Noteworthy, the lattice distortion at the cubic-tetragonal phase transition combined with SOC leads to a complete splitting of the vectorial representation, whereas only a partial splitting is predicted when the SOC is neglected (Figure 3e).⁷⁴ Indeed for such a tetragonal lattice distortion, the three eigenvalues of the total Hamiltonian including all effects are doubly degenerated and can be analytically calculated at the critical point of the BZ (R-point in the cubic phase):

$$\begin{aligned}
 E_0 &= E_{CB}^0 + (l + m)\epsilon_{xx} + m\epsilon_{zz} + \Delta_{SO}/3 \\
 E_{\pm} &= E_0 - \frac{\delta E + \Delta_{SO}}{2} \pm \sqrt{\left(\frac{\delta E - \Delta_{SO}/3}{2}\right)^2 + 2\Delta_{SO}^2/9}
 \end{aligned}
 \tag{3}$$

with $\delta E = (l - m)(\epsilon_{xx} - \epsilon_{zz})$. The effect of strain is systematically smaller than that induced by SOC in lead halide compounds. If SOC is neglected ($\Delta_{SO}=0$ in equation 3), one of the CBM eigenvalues remains doubly degenerated (Figure 3e).

The most important effect of SOC is the reduction of the DOS in the conduction band by a factor of 3 as compared to DOS calculated without SOC.^{74,75} This DOS reduction has a deep impact on both transport and optical properties, which leads to a net reduction of the optical absorption by the same amount. However, since both valence and conduction bands are now associated with two dispersion parabola with almost the same effective masses (*i.e.* band curvatures),^{89,92} well-balanced carrier transport is obtained for both electrons and holes. Moreover, the optical absorption is related to the joint density of states close to the critical R-point, and thus to reduced mass μ . This quantity is smaller in GaAs than in 3D HOP, because the mass of the hole and of the electron, are respectively larger and smaller. This important property directly connected to SOC adequately compensates the smaller Kane energy (oscillator strength) of 3D HOP.⁷⁵

When including SOC, the calculated band-gaps of 3D HOP appear substantially under-estimated. This is a well-known limitation of DFT that can be corrected by including many-body effects using GW self-energy corrections for the band gap⁶⁶ and the Bethe–Salpeter equation for the exciton.⁷⁴ Accurate GW self-energy corrections have been reported by several research groups.^{80,92,108} These affords more accurate effective masses of the carrier both for conventional semiconductors and HOP.

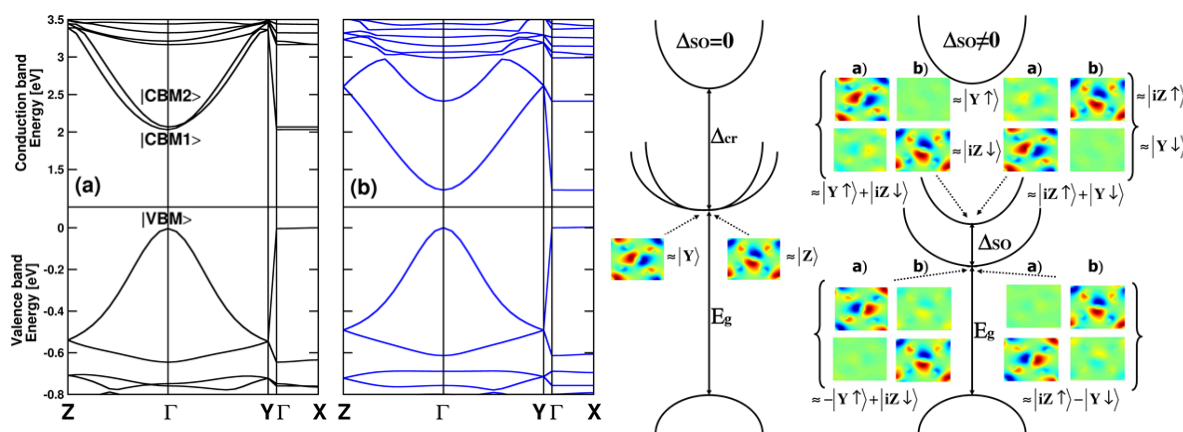


Figure 6. (left) Electronic band structure of the layered HOP ([pFC₆H₅C₂H₄NH₃]₂PbI₄), (a) without and (b) with SOC at the GGA level of theory with the ABINIT code.⁶¹ Energy levels are referenced to the valence band maximum. (right) General schematic representations of a layered HOP electronic band diagram without and with SOC (Δ_{SO}). Δ_{cr} represents the anisotropy of the crystal field. This panel includes the a) real and b) imaginary parts of the complex spinorial components of the first and second CBM states with spin up component on top of the spin down component. Reprinted with permission from ref 61. Copyright 2012 American Physical Society.

Similar effects are predicted for layered HOP with the important difference that the CBM without SOC is only twice degenerated.⁶¹ SOC leads to a non-degenerate state at the CBM, and band dispersion at both CBM and VBM correspond to parabola (Figure 6) in a favorable configuration for optical activity. The study of the monocrystalline [pFC₆H₅C₂H₄NH₃]₂PbI₄ as a model for layered HOP demonstrates that the optical process is governed by three active Bloch states at the Γ point of the reduced BZ (two Bloch states for the CBM and one for the VBM) with a reverse ordering compared to anisotropically bonded semiconductors. Giant spin-orbit coupling effects (Figure 6d) and optical activities can be subsequently inferred from symmetry analysis. Transverse electric (TE) optical activity has indeed been predicted theoretically,⁶¹ in good agreement with previous experimental results.¹ TE optical activity indicates that the light is absorbed at the band gap energy only for an electrical polarization parallel to the layer.

5. LOSS OF INVERSION SYMMETRY AND INTERPLAY WITH SOC IN HOP

The room (high)-temperature crystal structure of 3D HOP remains under debate. For long time it was assumed that it is a Pm3m pseudo-cubic phase, with dynamical disorder of the organic cations CH_3NH_3^+ .¹²⁴ Recent single crystal X-ray diffraction experiments, performed on $\text{CH}_3\text{NH}_3\text{MI}_3$ close to room temperature, support a non-centrosymmetric tetragonal space group (P4mm, no 99).²⁹ These compounds exhibit a phase transition to a tetragonal non-centrosymmetric and centered I4cm phase (no 108),²⁹ instead of the previously reported symmetric I4mcm phase (no 140).¹²⁴ This phase transition is associated to a group-subgroup relationship between P4mm and I4cm, similar to the previously proposed Pm3m-I4/mcm phase sequence analyzed for $\text{CH}_3\text{NH}_3\text{PbI}_3$ and $\text{CH}_3\text{NH}_3\text{PbBr}_3$.¹²⁴ For $\text{CH}_3\text{NH}_3\text{PbI}_3$, this transition occurs above room temperature ($T_c = 333$ K).

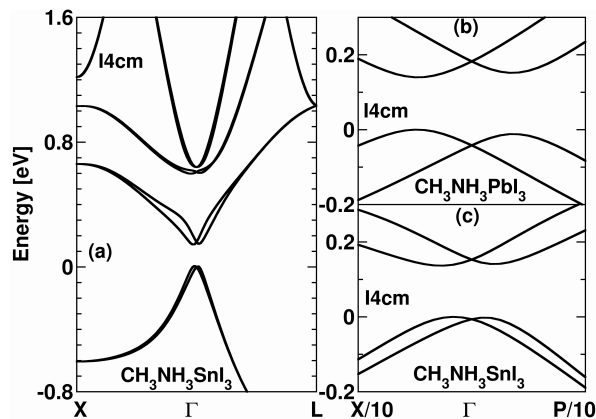


Figure 7. (a) Band structure of the I4cm phase of $\text{CH}_3\text{NH}_3\text{SnI}_3$ at the LDA level of theory with SOC and zoom close to the critical Γ point, (b) for $\text{CH}_3\text{NH}_3\text{PbI}_3$ and (c) for $\text{CH}_3\text{NH}_3\text{SnI}_3$ showing the Rashba-Dresselhaus-like spinor splitting due to the loss of inversion symmetry. Reprinted with permission from ref 74. Copyright 2014 John Wiley and Sons.

The loss of the inversion symmetry combined with SOC is known to induce a splitting of the spinor bands away from the critical point Γ of the BZ in zinc-blende and wurtzite structures.¹²⁵ It was shown, that even a small symmetry breaking might lead to a strong spinor splitting in 3D

HOP, since the SOC effect is giant by comparison to conventional semiconductors.⁷⁴ Indeed, the electronic band diagram of the P4mm phase of both CH₃NH₃MI₃ HOPs close to the critical point A (corresponding to R of the ideal cubic phase) exhibits a splitting of the spinor bands, even though the deviation from the centrosymmetric Pm3m phase is small.⁷⁴ This effect is more important for the tetragonal I4cm CH₃NH₃PbI₃ and CH₃NH₃SnI₃ 3D HOP (Figure 7).⁷⁴ The bands are folded from A to Γ . This result was confirmed and analyzed, using different DFT approaches and subsequent many-body corrections.^{83,92,114} This effect still deserves some experimental investigation, and various names were used for the interplay between SOC and loss of inversion symmetry in 3D HOP: “band splitting”,⁷⁴ Rashba/Dresselhaus,⁸³ Dresselhaus,⁸³ Rashba.¹¹⁴ Besides, we underline that most works on the low temperature phase of CH₃NH₃PbI₃ use a refined centrosymmetric orthorhombic Pnma structure. However, the structure of the low temperature-phase of CH₃NH₃MI₃ is still debated: alternative monoclinic, triclinic or non-centrosymmetric Pna21 (no 33) orthorhombic space groups have been suggested, which might also lead to spinor band splitting away from critical points of the BZ.

6. QUANTUM CONFINEMENT AND BAND ALIGNEMENT IN LAYERED HYBRID PEROVSKITES

Most available theoretical studies of layered HOP either give a general description of the electronic band structure, mostly using DFT,^{55,57,61} or focus on the excitonic coupling using effective mass parameters for the carrier dispersion and abrupt dielectric confinement schemes.^{126,127} HOP are very large systems for DFT simulations, and may not be readily simulated due to computational resources limitations. A fruitful schematic representation was proposed earlier for the layered HOP, introducing the quantum well (QW) concept.¹¹ This

qualitative picture is based on the observation that these HOP are built from semiconducting inorganic sheets alternating with organic layers having much wider band gaps. It has led a qualitative picture of a type-I QW-like heterostructure but without quantitative evaluation of the quantum confinement effect.

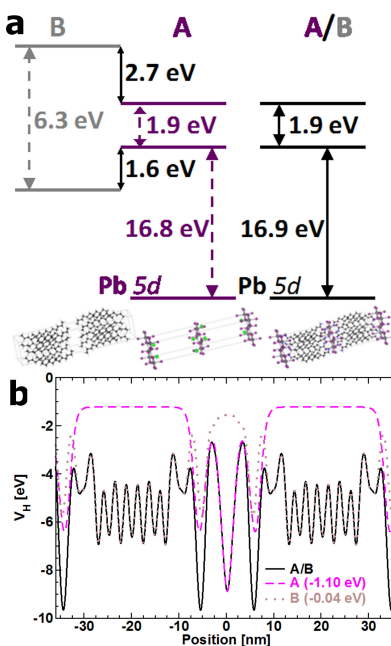


Figure 8. (a) Schematic representation of the CB and VB band alignment considering the $(C_{10}H_{21}NH_3)_2PbI_4$ layered HOP (A/B) as a composite structure made of A and B modeled by $(Na)_2PbI_4$ and $C_{10}H_{21}CH_3$, respectively. (b) Computed potential profiles for the real layered material (A/B, straight line) and A (dashed line) and B (dotted line) bulk-like materials. Band alignment and match of potential profiles of A/B are obtained thanks to downward shifts of the data computed for A and B amounting to 1.1 and 0.04 eV, respectively. Alignment of Pb 5d orbitals between A/B and A also require a 1.1 eV shift. Computations were performed at the GGA level of theory with the SIESTA code. Reprinted with permission from ref 128. Copyright 2014 John Wiley and Sons.

However, quantitative description of layered HOP faces a special problem.¹²⁸ In fact, defining the A/B heterostructure (the layered HOP) is straightforward, but the A and B bulk partners are not easily defined. A reverse situation applies in heterostructures grown from conventional semiconductors, since in this case the A and B bulk partners are usually well-defined, whereas the A/B heterostructure (a QW for example) deserves experimental or theoretical investigations

to capture accurately effects of strain or interfaces.¹²³ At a first glance, A and B of layered HOP may be roughly defined as the inorganic and organic layers, respectively. It is however not conceptually possible to identify the A and B bulk partners with the inorganic PbI_4^{2-} layer and the sheet of organic cations, respectively, because the Coulomb interactions have a large impact on the electronic structure of the hybrid material, especially for electronic states located close to the electronic band gap.

A description of the whole layered HOP structure as a composite material in a slightly different way leads to a rigorous composite approach to compute the valence band-lineup of layered HOP.¹²⁸ It is illustrated for $(\text{C}_{10}\text{H}_{21}\text{NH}_3)_2\text{PbI}_4$ in Figure 8. Structure A corresponds to an inorganic PbI_4^{2-} perovskite layer with additional Na cations. The latter are located at the position of the nitrogen atom of each organic cation, so as to mimic the strong Coulomb interactions close to the inorganic layer. This procedure is similar to the one described in section 3 for 3D HOP. Structure B is made from organic neutral molecules, namely $\text{C}_{10}\text{H}_{21}\text{CH}_3$, where nitrogen has been replaced by a carbon atom, thus ensuring electroneutrality and avoiding electron double counting. This analysis shows that layered HOP can be considered as composite materials with very weak interactions between the inorganic layers, with a reconstruction of the whole Hartree potential profile by pieces. This procedure also affords the complete conduction and valence band alignments of A and B “bulk like” materials with respect to each other and leads to a confinement potential both for holes and electrons in the inorganic layer. Interestingly, alignment of the Pb *5d* orbitals of both structures is obtained by applying the very same offsets. In fact, Pb *5d* orbitals are low lying orbitals virtually unaffected by chemical substitution in such HOP and have already been suggested to offer an electronic marker to deduce absolute VB energies.⁶⁶ It gives a first quantitative basis to the schematic picture initially proposed by Mitzi and

coworkers.¹¹ It should be noted that this quantitative approach for layered HOP can be extended to other classes of 2D functional materials.¹²⁸

It is quite tempting to apply effective mass models based on an ultrathin quantum well with finite confinement barriers for carriers, in order to predict quantum confinement effects in layered HOP.¹²⁹ In fact, these empirical models predict superlattice (SL) effects that are not observed experimentally or in DFT simulations.^{61,128} A first fundamental limitation of such approaches is that the computed energies of the confined charge carriers lie in a range where strong non-parabolicity occurs.¹²⁸ Indeed, the commonly used reference bulk material, namely (CH₃NH₃)PbI₃, taken for the inorganic QW¹²⁹ is inappropriate.¹²⁸ An even more important limitation is that the effective mass model for conventional semiconductor heterostructures is linked with a definition of envelope functions to describe the electronic eigenfunctions of the whole heterostructure:

$$\psi(\mathbf{r}) = \sum_n F_n(\mathbf{r}) U_n(\mathbf{r}), \quad (4)$$

where $F_n(\mathbf{r})$ are expected to be slowly varying functions and $U_n(\mathbf{r})$ are periodic and rapidly oscillating functions. In order to further simplify the theoretical approach, the envelope function approximation (EFA) is usually implemented in these heterostructures. This approximation relies on the assumption that many conventional bulk semiconductors are similar in their chemical nature, and that the Bloch functions of the bulk materials at high symmetry points of the BZ differ only slightly. At the Γ point, it corresponds to the following approximation:

$$U_n(\mathbf{r}) \approx U_{n,\mathbf{0}}^A(\mathbf{r}) \approx U_{n,\mathbf{0}}^B(\mathbf{r}) \quad (5)$$

If this condition is verified, boundary conditions related to the A/B heterostructure, namely $\psi(\mathbf{r})$ total wavefunction matching and current conservation, are reduced to the envelope

functions only. For a simple effective mass model in a A/B 2D structure perpendicular to the z axis having its interface at $z = z_i$:

$$F_n^A(z_i) = F_n^B(z_i) \quad (6)$$

$$\frac{1}{m^A} \frac{dF_n^A(z)}{dz} \Big|_{z=z_i} = \frac{1}{m^B} \frac{dF_n^B(z)}{dz} \Big|_{z=z_i}$$

The condition on the Bloch functions of A and B partners (equation 5) is obviously not fulfilled in layered hybrid perovskites. The confinement potentials are correctly defined but cannot be combined with effective mass modeling to predict quantum confinement effects.¹²⁸ Thus, DFT remains the best tool to explore theoretically electronic properties of layered HOP.⁶¹

7. QUANTUM CONFINEMENT AND BAND ALIGNMENT IN 3D HYBRID PEROVSKITE HETEROSTRUCTURES

Heterostructures built from 3D HOP have not yet been explored experimentally, but may afford interesting possibilities for band gap engineering and tuning of the optical and transport properties. The electronic band dispersion diagrams of two short SL of $\text{CH}_3\text{NH}_3\text{PbI}_3/\text{CH}_3\text{NH}_3\text{PbBr}_3$ are compared in Figure 9. The stacking axis for the two SL correspond to the long axis (b) of pure $\text{CH}_3\text{NH}_3\text{PbI}_3$ and $\text{CH}_3\text{NH}_3\text{PbBr}_3$ in their low temperature Pnma orthorhombic phases. The electronic dispersions of the two SL along the Γ -Y direction are more flat than for pure compounds, indicative of quantum confinement. Thus, the $\text{CH}_3\text{NH}_3\text{PbI}_3$ domain of a $\text{CH}_3\text{NH}_3\text{PbI}_3/\text{CH}_3\text{NH}_3\text{PbBr}_3$ SL can be considered as a QW.

This finding is consistent with the CB and VB confinement potentials reported recently in this system (0.09 eV in the CB and 0.26 eV in the VB).⁹³ However, the electronic dispersions along the Γ -Y direction are not completely flat (Figure 9) like in a 2D layered HOP (Figure 6). This is a clear indication that the wavefunctions both in CB and VB are coupled from $\text{CH}_3\text{NH}_3\text{PbI}_3$ QW

to $\text{CH}_3\text{NH}_3\text{PbI}_3$ QW (so-called SL effect), and are delocalized across the $\text{CH}_3\text{NH}_3\text{PbBr}_3$ barrier. This is confirmed when plotting along the SL axis, the electronic density of the VBM of the short SL (SL2) containing two $\text{CH}_3\text{NH}_3\text{PbI}_3$ cells and one $\text{CH}_3\text{NH}_3\text{PbBr}_3$ cell (Figure 10). The electronic density is maximum in the $\text{CH}_3\text{NH}_3\text{PbI}_3$ QW, but non-zero in the $\text{CH}_3\text{NH}_3\text{PbBr}_3$ barrier.

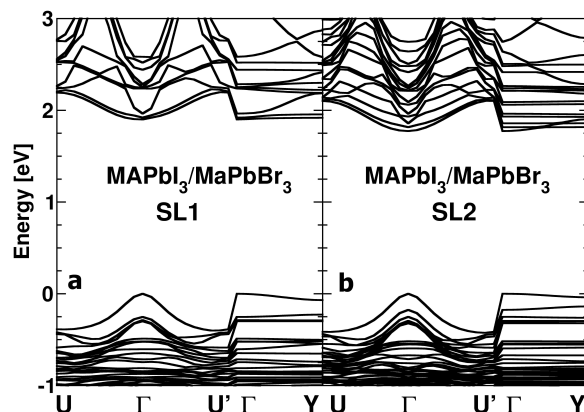


Figure 9. Electronic band dispersion diagrams of (a) a short SL (SL1) containing one $\text{CH}_3\text{NH}_3\text{PbI}_3$ cell and one $\text{CH}_3\text{NH}_3\text{PbBr}_3$ cell, and (b) a second short SL (SL2) containing two $\text{CH}_3\text{NH}_3\text{PbI}_3$ cells and one $\text{CH}_3\text{NH}_3\text{PbBr}_3$ cell. The diagrams are computed at the DFT level with a GGA exchange-correlation functional. SOC is not considered.

Moreover, one may notice that the electronic density variation is based on the same repeated pattern close to lead atoms (Figure 10a). This pattern is characteristic of the VBM Bloch functions of both $\text{CH}_3\text{NH}_3\text{PbI}_3$ and $\text{CH}_3\text{NH}_3\text{PbBr}_3$ 3D bulk HOP.⁶⁶ The $U_n(\mathbf{r}) \approx U_{n,0}^A(\mathbf{r}) \approx U_{n,0}^B(\mathbf{r})$ condition (equation 5) is thus fulfilled with $A=\text{CH}_3\text{NH}_3\text{PbI}_3$ and $B=\text{CH}_3\text{NH}_3\text{PbBr}_3$. We may thus infer that the EFA is justified for an heterostructure built from 3D HOP. We must point out that computation of the electronic properties, including quantum confinement and electronic band gap opening, are only qualitatively computed at the DFT level, since the effective masses at CBM and VBM are not well predicted at this level of theory. Figure 10c is a representation of the CBM and VBM energies computed for a single $\text{CH}_3\text{NH}_3\text{PbI}_3/\text{CH}_3\text{NH}_3\text{PbBr}_3$ QW within the EFA as a function of the QW thickness. The effective masses were taken from ref. 92 (DFT

including GW corrections) and the quantum confinement potentials from ref. 93. The quantum confinement effect is expected to show off mostly for $\text{CH}_3\text{NH}_3\text{PbI}_3$ QW thicknesses below 4nm, where the correction to the band gap is larger than 0.1eV. This is consistent with the fact that the DFT computation for SL2 (Figure 9b) is different from that of bulk $\text{CH}_3\text{NH}_3\text{PbI}_3$ (Figure 5a).

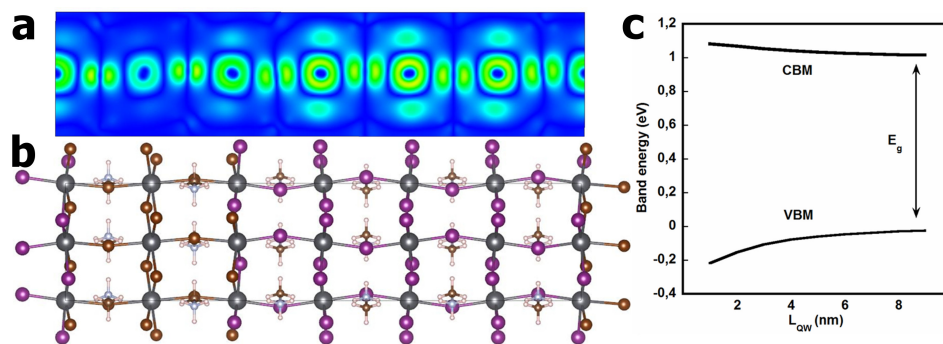


Figure 10. (a) Representation of the electronic density for the VBM of the short superlattice SL2 containing two $\text{CH}_3\text{NH}_3\text{PbI}_3$ cells and one $\text{CH}_3\text{NH}_3\text{PbBr}_3$ cell (Figure 9). (b) Overview of the atomic structure. The diagrams are computed at the DFT level with a GGA exchange-correlation functional. SOC is not considered. (c) Electron (CBM) and hole (VBM) ground state energies as a function of QW thickness computed with effective masses and confinement potentials.

8. EXCITONIC PROPERTIES OF 3D AND LAYERED HYBRID PEROVSKITES

Excitonic effects in 3D HOP can be accounted for using the Bethe-Salpeter Equation (BSE) starting from the mono-electronic states calculated at the DFT level. When taking the high temperature cubic phases of $\text{CH}_3\text{NH}_3\text{PbI}_3$ and $\text{CH}_3\text{NH}_3\text{PbCl}_3$ enhancement of absorption at the bandgap was clearly evidenced in the first BSE simulations reported in the literature (Figure 11).⁷⁴ In the perturbative BSE/DFT approach, screening of the electron-hole interaction due to atomic motion⁷⁵ is however not taken into account. Thus it is more suited to the low temperature phases where these motions are frozen. Such calculations are computationally more demanding and only resonant states may be addressed accurately.¹¹⁷ This means that at higher temperature experimental results on the exciton cannot be fully understood from BSE /DFT calculations, if atomic motions are important.

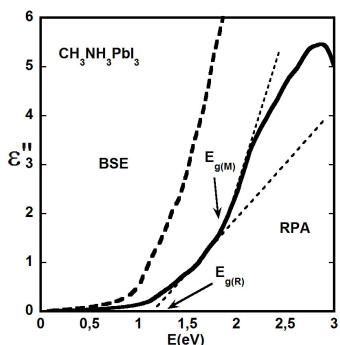


Figure 11. Comparison of the dielectric constant variation of $\text{CH}_3\text{NH}_3\text{PbI}_3$ in the cubic phase with (dotted line, BSE calculation) and without (straight line, RPA calculation) the excitonic interaction. Onsets of optical transitions at R and M points are indicated. Reprinted with permission from ref 74. Copyright 2014 John Wiley and Sons.

Given the large carrier mobilities reported for 3D HOP at room temperature, it prompts the question of the nature of photoexcited species: are they free charges or excitons? If excitons, what role do they play in transport properties?⁷⁵ Importance and origin of exciton screening can be gauged from experimental absorption spectra of $\text{CH}_3\text{NH}_3\text{PbI}_3$ recorded two decades ago by Ishihara et al (Figure 12).^{4,75} These data are consistent with an exciton quenching between 159K and 212K, in relation with the structural transition at $T_c = 162\text{K}$. The exciton quenching includes both exciton screening and gap-switching (ill-named exciton-switching), that also appears in 2D hybrids. A convenient way to simulate exciton screening is to consider a two-particle wave function $\psi(\mathbf{r}_e, \mathbf{r}_h)$, where \mathbf{r}_e (\mathbf{r}_h) is the electron (hole) position. The effective mass approximation works well close to the bandgap in 3D HOP both for electrons and holes, thanks to the giant SOC in the CB leading to a non-degenerate band instead of a triply degenerate one obtained without SOC. For a Wannier exciton, the two-particle Hamiltonian including electrostatic interaction and effective mass approximation for electrons and holes reads:

$$H = E_g - \frac{\hbar^2}{2m_e} \Delta_{r_e} - \frac{\hbar^2}{2m_h} \Delta_{r_h} - \frac{e^2}{4\pi\epsilon_{\text{eff}}|\mathbf{r}_e - \mathbf{r}_h|}$$

E_g is the bandgap energy, ϵ_{eff} the effective dielectric constant and $m_e(m_h)$ the electron (hole) effective mass. Introducing the center-of-mass coordinate $\mathbf{R} = \frac{m_e \mathbf{r}_e + m_h \mathbf{r}_h}{m_e + m_h}$ and the difference coordinate $\mathbf{r} = \mathbf{r}_e - \mathbf{r}_h$, the general solution becomes of separable form, $\psi(\mathbf{r}_e, \mathbf{r}_h) = \frac{e^{i\mathbf{K}\cdot\mathbf{R}}}{\sqrt{V}} \varphi(\mathbf{r})$, where $\varphi(\mathbf{r})$ is the wave function for the relative motion. The corresponding Hamiltonian reads:

$$H_r = E_g - \frac{\hbar^2}{2\mu} \Delta_r - \frac{e^2}{4\pi\epsilon_{eff}r}$$

The reduced mass is defined by $1/\mu = 1/m_e + 1/m_h$. For the relative electron-hole motion, the Hamiltonian has two types of solutions: bound pair states and a continuum of pair states having energies larger than the bandgap energy. The optical absorption reads:

$$I(\tilde{E})/I_0 = 4 \sum_{n=1}^{\infty} \frac{\gamma/n^3}{(\tilde{E} + 1/n^2)^2 + \gamma^2} + \int_0^{\infty} \frac{d\tilde{E}'}{\pi} \frac{\gamma S(\tilde{E}') \sqrt{\tilde{E}'}}{(\tilde{E} - \tilde{E}')^2 + \gamma^2}$$

$\tilde{E} = \frac{\hbar\omega - E_g}{Ry}$ is a reduced energy, γ a broadening factor that depends on the temperature and

$S(\tilde{E}) = \frac{2\pi}{\sqrt{\tilde{E}}(1 - e^{-2\pi/\sqrt{\tilde{E}}})}$ the absorption enhancement factor of the continuum of pair states.

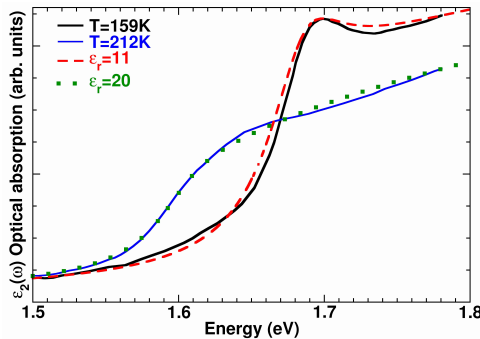


Figure 12. Optical absorption spectra of $\text{CH}_3\text{NH}_3\text{PbI}_3$ highlighting exciton screening. Experimental data taken from Ref. 4 recorded at 159 K (black line) and 212 K (blue line) and computed spectra for bound and continuum pair states, considering two-particle wave function and effective mass equations for electron and hole. The effect of dielectric screening is shown for $\epsilon_{eff} = 11$ (red dash line) and 20 (green dot line) and leads to a good fit of the experimental

spectra below (159K) and above T_c (212K), respectively. Reprinted with permission from ref 75. Copyright 2014 American Chemical Society.

In the limit of free carriers or a totally screened interaction, the first sum (bound exciton pair states) disappears in the optical absorption expression and $S(\tilde{E})$ reduces to 1. Simulated optical absorption spectra using this expression are compared to experimental data in Figure 12 considering different dielectric constants to highlight screening effects. Experimental absorption spectra above T_c are well reproduced with a larger effective dielectric constant than below T_c . It yields a reduction of the 1S exciton binding energy for all the spectra above T_c , from *ca* 15 to 5meV.⁷⁵ The value of the exciton binding energy being smaller than kT ($\sim 26\text{meV}$) at room temperature, we may infer that most electron-hole pairs are ionized yielding free carriers. This also leads to *ca* 80meV shift of the electronic bandgap at T_c . For $\text{CH}_3\text{NH}_3\text{PbX}_3$, we thus concluded that the exciton resonance below T_c is related to a Wannier like exciton already partially screened by polar modes, that becomes completely screened above T_c , yielding almost free carriers when the orientational motions of the cations are collectively activated by a structural phase transition. It shall be pointed that recent experimental results,⁵³ similar to the one reported in Figure 12 and carefully measured up to the room temperature, were discussed on the basis of a commonly accepted larger exciton binding energy ($\sim 37\text{-}50\text{ meV}$) reported earlier at helium temperature.^{4,130} In order to overcome this apparent paradox, a complete experimental study was performed on a single crystal from helium temperature up to the room temperature (Figure 13, left part).¹³¹ In the [25-75K] temperature range, multiple exciton lines were indeed observed. The helium temperature exciton line was attributed to a bound exciton with a long lifetime, possibly arising from a strong coupling with the organic cations.

The Wannier exciton screening phenomenon was connected with a renormalization of the dielectric constant at high temperature above T_c (Figure 13, right part).⁷⁵ The high frequency

dielectric constant ϵ_∞ is equal to about 6.5 for $\text{CH}_3\text{NH}_3\text{PbI}_3$, it has a smooth variation (~ 30) as a function of the temperature in the medium frequency range ($\sim 90\text{GHz}$).¹²⁴ The difference between these two values can be associated to the contributions of polar optical phonons by comparison to purely inorganic 3D perovskites.⁷⁵ The low frequency value ($\sim 1\text{KHz}$) undergoes a steep increase above the critical temperature T_c ,¹³² and amounts to about 60 at room temperature. This is attributed to the thermal activation of the collective tumblings of methyl ammonium cations.⁷⁵ More recently, a giant dielectric constant phenomenon was shown in these materials consisting on a low frequency dielectric constant in the dark of the order 1000.¹³³ With illumination under 1 sun, an additional factor of 1000 was measured. This effect was attributed to very slowly moving photoinduced carriers assisted by molecular rotations.

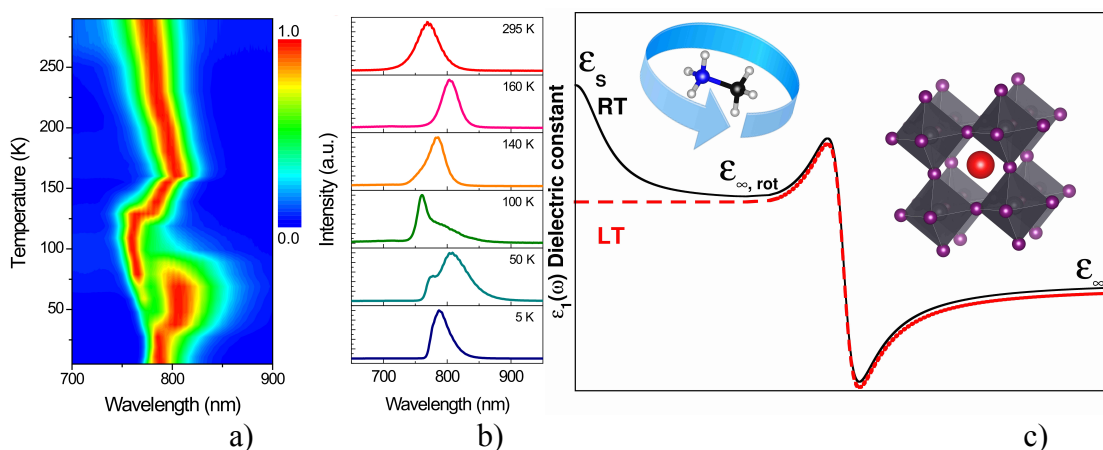


Figure 13. (Left) Temperature dependent PL of $\text{CH}_3\text{NH}_3\text{PbI}_3$ single crystal. (a) Contour plot of PL spectra of $\text{CH}_3\text{NH}_3\text{PbI}_3$ single crystal at different temperatures under weak excitation density ($0.42 \mu\text{J}/\text{cm}^2$) and (b) PL spectra at different temperatures. Reprinted with permission from ref 131. Copyright 2015 John Wiley and Sons. (Right, c) Schematic drawing of the dielectric constant as a function of frequency displaying respective contributions of CH_3NH_3^+ rotations and perovskite's vibrational phonons. Frequency dependence of the dielectric constant of $\text{CH}_3\text{NH}_3\text{PbX}_3$ as derived from available experimental data. The high frequency response is related to vibrational polar phonons stemming from modes of the perovskite lattice (right-hand side). The low frequency region (left-hand side) shows significant differences for the low (LT, red dash line) and room (RT, black line) temperature phases. The static dielectric constant (ϵ_S) increase at RT is primarily attributed to rotational motion of CH_3NH_3^+ cations. Reprinted with permission from ref 75. Copyright 2014 American Chemical Society.

For layered HOP, resonance energies of Wannier-series excitons have been measured and empirically predicted up to the 4s state.¹³⁴ Figure 14-a shows the low temperature absorption spectrum of the layered HOP $(\text{C}_6\text{H}_5\text{C}_2\text{H}_4\text{-NH}_3)_2\text{PbI}_4$ (PEPI) layer. The peak located at about 2.35-2.40eV correspond to the excitonic resonance, and the band to band absorption edge is located at about 2.60eV. The exciton binding energy is thus very large on the order of 200-250meV. For that reason, clear exciton signatures can be observed at room temperature by photoluminescence or absorption on the same compound (Figure 14-b).¹³⁵ This behavior is clearly different from the one of 3D HOP (Figures 12 and 13).

In 2D HOP, the enhanced exciton binding is related to dielectric confinement and image-charge effects, where the dielectric mismatch between both the inorganic and the organic layers is the most important parameter.^{126,127} The 2D Wannier exciton model leads to a consistent picture where the in-plane exciton Bohr radius exceeds by far the thickness of the inorganic layer where it is confined. It allows a consistent theoretical analysis of the influence of halogen alloying on the excitonic properties of layered HOP exactly like in $\text{Ga}_{1-x}\text{Al}_x\text{As}$, $\text{Cd}_{1-x}\text{Hg}_x\text{Te}$ conventional semiconductor pseudobinary alloys.⁷⁹ Figure 14-b shows the room temperature optical tunability and inhomogeneous broadening of $(\text{C}_6\text{H}_5\text{C}_2\text{H}_4\text{-NH}_3)_2\text{PbI}_{4(1-x)}\text{Br}_{4x}$ (PEPI/PEPB) alloys.

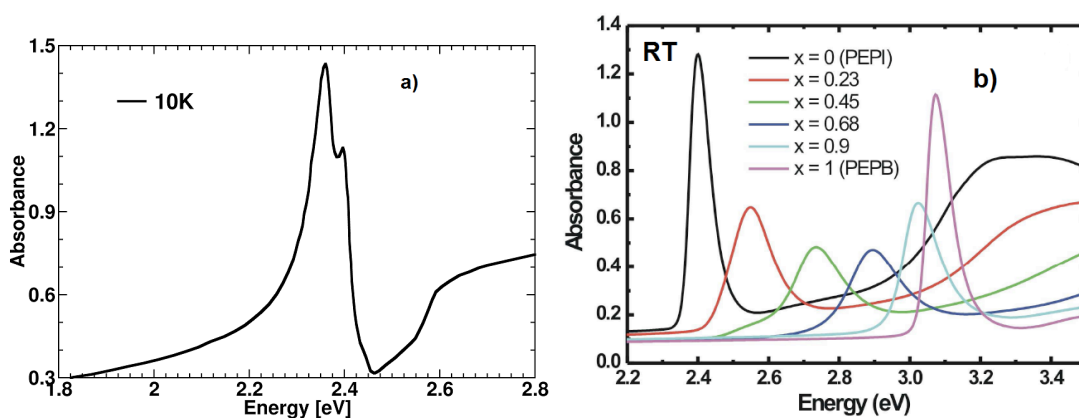


Figure 14. (a) Absorbance spectrum of a 50 nm thick $(\text{C}_6\text{H}_5\text{C}_2\text{H}_4\text{-NH}_3)_2\text{PbI}_4$ (PEPI) layer deposited by spin-coating on a glass substrate at $T=10\text{K}$. The peak located at about 2.35-2.40eV correspond to the excitonic resonance, and the band to band absorption edge is located at about

2.60eV. (b) Room temperature (RT) optical absorbance spectra of $(\text{C}_6\text{H}_5\text{C}_2\text{H}_4\text{-NH}_3)_2\text{PbI}_{4(1-x)}\text{Br}_{4x}$ (PEPI/PEPB) alloys. Reprinted with permission from ref 79. Copyright 2014 American Chemical Society.

Until recently, empirical modeling of dielectric confinement in layered HOP relied on crude models with ad-hoc abrupt interfaces. Recently, we have introduced the basic concepts needed to model dielectric confinement beyond an abrupt interface, based on the atomic structure of the interface.¹³⁶ It solves the well-known problem of self-energy divergence at abrupt dielectric interfaces and was first implemented for CdSe colloidal nanoscale platelets.¹³⁶ This method allows to bridge the gap between empirical and DFT approaches used to investigate hybrid nano-objects with 2D electronic properties.¹³⁷ It may also contribute to further unravel electronic and dielectric properties of HOP, especially of layered structures such as $(\text{R-NH}_3)_2\text{MX}_4$ and nanocrystals or slabs of $(\text{R-NH}_3)\text{MX}_3$.

9. AUGER EFFECTS IN 3D HYBRID PEROVSKITES

Next, in order to gauge the importance of third-order non-radiative Auger effects on the recombination in hybrid perovskites, let us recall a few basic properties of these effects in the case of free carriers for III-V semiconductors (Fig. 15a). An electron in the CB may recombine with a heavy hole (HH) in the VB, either by transferring energy and momenta to an electronic transition in the CB (CCCH or CHCC process) or to a hole jumping from the HH band to the spin-orbit split-off (SO) band (CHHS or CHSH process). Among the numerous other possible processes, a third process (CHLH or CHHL) not represented here is usually considered in addition. This is weaker than the CHHS process, except when the spin-orbit splitting energy is larger than the band gap value. A complete computation involves repeated summations over the particle momenta and a careful description of the electronic band structure since electronic states

away from the band gap are involved. The dependence on the electronic band gap, effective masses and carrier densities can be estimated using parabolic approximations for the electronic dispersions:¹³⁸⁻¹⁴⁰

$$R_{Auger,CCCH} \propto n^2 p e^{-(E_T - E_{gap})/kT}$$

$$R_{Auger,CHHS} \propto np^2 e^{-(E_T - E_{gap})/kT}$$

Where

$$E_{T,CCCH} = \frac{2m_C + m_{HH}}{m_C + m_{HH}} E_{gap}$$

$$E_{T,CHHS} = \frac{2m_{HH} + m_C}{2m_{HH} + m_C - m_{SO}} (E_{gap} - \Delta)$$

The Auger effect decreases exponentially as the band gap increases in III-V semiconductors. Moreover, the CHHS process prevails in most cases, when the SOC is of the same order as the band gap value. However, in n-type III-V semiconductors, the CCCH process is dominant through the $n^2 p$ term. For large band gap materials like GaAs, where Auger processes are very weak, phonon-assisted Auger processes (not described here) give the most important contribution.¹³⁹

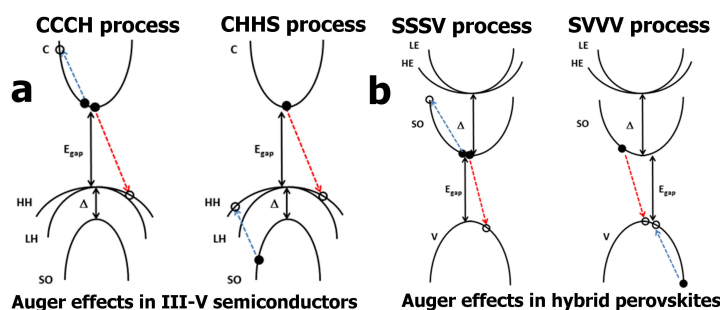


Figure 15. Dominant Auger processes in (a) III-V semiconductors and (b) hybrid perovskites. (a) HH, LH and SO correspond to heavy hole, light hole and hole spin-orbit split-off states of the VB in III-V semiconductors. (b) HE, LE and SO correspond to heavy electron, light electron and electron spin-orbit split-off states of the CB in hybrid perovskites.

The electronic band structure of hybrid perovskites (Fig. 15b) has two important differences with respect to that of III-V semiconductors: 1) the band structure has a reverse band ordering, with SOC in the CB, and 2) the SOC energy is much larger than that found in III-V semiconductors. Using the method developed for III-V semiconductors, we obtain the dependence on the electronic band gap, effective masses and carrier densities for the main two processes:

$$R_{Auger,SSSV} \propto n^2 p e^{-(E_T - E_{gap})/kT}$$

$$R_{Auger,SVVV} \propto np^2 e^{-(E_T - E_{gap})/kT}$$

where

$$E_{T,SSSV} = \frac{2m_S + m_V}{m_S + m_V} E_{gap}$$

$$E_{T,SVVV} = \frac{2m_V + m_S}{m_S + m_V} E_{gap}$$

An electron in the split-off state (S) of CB may recombine with a hole (V) in the VB, either by transferring energy and momenta to an electronic transition in the split-off state of the CB (SSSV process) or to a hole in the valence band (SVVV process). Using effective masses of 0.12 and 0.15 for the split-off CB and for the VB respectively, we deduce that the SVVV process slightly dominates over the SSSV process, especially when considering the natural p-doping of hybrid perovskites. However, in HOP the overall effect is expected to be small like in large band gap III-V semiconductors (GaAs, InP, ...). A process involving SOC, and heavy electrons (HE) or light electrons (LE) in the CB (HSSV or LSSV processes), requires a more detailed analysis, since HE and LE electronic states, lying at high energy in the CB, are hybridized with molecular states.⁶⁶

Besides, one may also expect that some of the carriers generated by the secondary optical transitions at R may be trapped in the $F_{3/2u}$ (CB) and $F_{3/2g}$ (VB) states (Figure 4). Further

relaxation $F_{3/2g} \rightarrow E_{1/2g}$ (holes) and $F_{3/2u} \rightarrow E_{1/2u}$ (electrons) may occur by optical phonon assisted relaxations, provided that the energy conservation rule is fulfilled. Such a “*phonon bottleneck*” is well-known in the case of carrier injection in semiconductor quantum dots. However, in quantum dots, carrier assisted (Auger relaxation) is the dominant effect in the high injection regime, yielding very fast carrier relaxation. A similar hot hole cooling (intraband process) from $F_{3/2g}$ to $E_{1/2g}$ is consistent with the relaxation after optical excitation reported for $\text{CH}_3\text{NH}_3\text{PbI}_3$.²⁸

10. COMPUTATIONAL METHODS

Some calculations were performed using the DFT implementation available in the ABINIT package,¹⁴¹ with the LDA or the GGA PBE gradient correction for exchange-correlation¹⁴² and relativistic, norm-conserving, separable, dual-space Gaussian type pseudopotentials of Goedecker, Teter, and Hutter for all atoms.¹⁴³ The SIESTA code was used for the other simulations at the DFT-GGA level,¹⁴⁴ calculations being performed with and without SOC.¹⁴⁵

11. CONCLUSIONS AND PERSPECTIVES

In summary, this paper reviews some of the recent theoretical investigations that have contributed to gain better physical understanding of 2D and 3D HOP materials. It turns out that solid-state physics concepts are adapted to this “*new*” class of semiconductors, and prove to provide valuable tools when combined with DFT simulations of the optoelectronic properties. The pseudocubic perovskite structure plays a central role in this approach, allowing to define reference Bloch states and $\mathbf{k}\cdot\mathbf{p}$ Hamiltonians close to the electronic band gap. Lattice distortions, including loss of inversion symmetry, as well as spin-orbit coupling have a strong impact on the electronic band structure, in particular regarding folding, degeneracy, effective masses and optical absorption. A careful examination of the concepts of Bloch function and envelope

function is particularly helpful to understand the difference between layered HOP, and heterostructures made of 3D HOP. Carrier confinement potentials can be defined in both cases. Many-body interactions are discussed starting with excitonic interactions, where screening and dielectric confinements are of outmost importance for the optical properties at room temperature in 3D and layered HOP, respectively. The Wannier exciton picture holds for HOP in the medium to high temperature range (typically [100-300K]). The Wannier exciton wavefunction is mostly built from the Bloch wave electronic states located on the inorganic lattice. By contrast to layered HOP, the Wannier exciton in 3D HOP is strongly screened by collective molecular rotations as well as vibrations of the inorganic lattice. This mechanism, which leads to almost free carriers at room temperature, is essential to explain the good transport properties of these materials for photovoltaic applications. In the low temperature range, the interaction of the electron-hole pairs with neighboring molecules is expected to yield bound excitons with long lifetime in 3D HOP as shown by recent experimental studies.¹³¹ Further investigations are needed to explain the microscopic origin of this strong coupling mechanism, which might be connected to the giant dielectric constant observed in 3D HOP at low frequency and under illumination.¹³³ Theoretical studies on many-body effects in HOP are scarce due to the complexity of these materials. However, we have shown that semi-empirical approaches previously developed for conventional semiconductors may provide with new perspectives to analyze these effects. As an example, the detrimental non-radiative Auger effect close to the electronic band gap of 3D hybrid perovskites, is expected to have a moderate influence on the performances of optoelectronic devices. Finally, we hope to have convinced the reader that theoretical methods developed for both molecular materials and inorganic semiconductors may be efficiently combined to make further progress in the understanding of HOP.

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Notes

The authors declare no competing financial interests.

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