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Hypergroup properties for the deltoid model

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Abstract

We investigate the hypergroup property for the family of orthogonal polynomials associated with the deltoid curve in the plane, related to the A_2 root system. This provides also the same property for another family of polynomials related to the G_2 root system.

1 Introduction

The hypergroup property is a property shared by some orthonormal bases in probability spaces which allows for the complete description of all Markov sequences and for multiplication formulas (see Section 2). It has been established for some families of orthogonal polynomials in dimension 1, and relates in general to some special structure of the underlying space. This is a quite powerful tool in many areas, ranging from pure analysis and Lie group to statistics and computer algorithms. Gasper's theorem establishes this property for dissymmetric Jacobi polynomials in dimension 1 (see [10, 11, 12, 20]), and is an extension of an earlier result due to Bochner concerning the symmetric case [6, 7]. Although many authors revisited this result, a very elegant proof of Gasper's theorem was recently proposed by Carlen, Gerónimo and Loss in [8], arising from the study of Kac's model in statistical mechanics. This proof relies on the construction of an auxiliary model, and may be extended in many other situations. The aim of this paper is to show how this technique applies in the particular case of one of the 2 dimensional extension of Jacobi polynomials, namely orthogonal polynomials on the deltoid domain (defined below in Section 4) and associated to the A_2 root system.

Jacobi polynomials is the unique family of orthogonal polynomials in dimension 1 (together with their scaled limits Hermite and Laguerre polynomials) which are at the same time eigenvectors of a second order differential operator, and more precisely of a diffusion operator [21]. Although there is no reason a priori for these two properties to be related, it is worth to study the hypergroup property for other families

of such orthogonal polynomials in higher dimension. Among these orthogonal polynomials associated with diffusion operators, some of them arise from root systems, the Heckman-Opdam Jacobi polynomials [13, 14]. It is an open and challenging problem to analyse this question for these families in particular. Partial results have been obtained in this direction in [22, 23] in the BC_n case.

We investigate in this paper the family of orthogonal polynomials in dimension 2 related to the deltoid curve and the A_2 and G_2 root systems, introduced by [16, 17, 18], as one as the special families of orthogonal polynomials in dimension 2 which are at the same time eigenvectors of symmetric diffusion operators (see [4] and Section 4). The scheme of Carlen-Geronimo-Loss relies in a crucial way on the fact that the polynomials are eigenvectors of some operator, and that the corresponding eigenspaces have dimension one. In the context that we investigate, this last property fails to be true, and this introduces some extra complexity in the formulation of the main hypergroup property result in Section 6. This difficulty arises from a symmetry invariance in the deltoid model. To get rid of this difficulty, we may look at simpler forms of the deltoid model, that is consider only functions which are invariant under this symmetry. This leads to investigate a new model, related to the G_2 root system, for which the hypergroup property takes the usual form.

The paper is organized follows. In Section 2, we present the hypergroup property and the elegant approach initiated by Carlen-Geronimo-Loss [8] to obtain it. In Section 3, we introduce the language of symmetric diffusion operators, which be used through the rest of the paper. The deltoid model is described in Section 4, where we give some details about the structure of the eigenspaces, and show the relation with the A_2 root system, through some geometric interpretations of the associated operators for two distinct values of the parameter. The Carlen-Geronimo-Loss methods relies on the construction of some other model space (in general in higher dimension), which projects on the model under study and have some extra symmetry. This model (in our case 6-dimensional) is presented in Section 5. Then we give in Section 6 the hypergroup property for the deltoid model itself, and finally in Section 7, we introduce the projected model related to the G_2 root system on which the hypergroup property presents a simpler form.

2 Hypergroup property and Markov sequences

The world hypergroup had been first introduced in 1959 by H.S Wall [24], to generalize the notion of a group, when the product of two elements is the sum of a finite numbers of elements. Different generalizations had been put forward by C. Dunkl,[9] Jewett [15], see also the exposition book [5]. In our context, we shall mostly be concerned by the hypergroup property, as defined in [2], and which concerns some properties of orthonormal bases in $\mathcal{L}^2(\mu)$ spaces, where μ is a probability measure, and that we describe below.

Let (X, \mathcal{A}, μ) a probability space, for which some orthonormal basis of $\mathcal{L}^2(\mu)$ is given, which we suppose countable, and therefore ranked as $(f_0, f_1, \dots, f_n, \dots)$. We suppose moreover that $f_0 = 1$. A central question which arises quite often is to determine all sequences (λ_n) such that, if one defines a linear operator $K :$

$\mathcal{L}^2(\mu) \mapsto \mathcal{L}^2(\mu)$ through $K(f_n) = \lambda_n f_n$, then this operator is Markov. This means that $K(\mathbf{1}) = \mathbf{1}$ and $K(f) \geq 0$ whenever $f \geq 0$. Of course, the first condition reduces immediately to $\lambda_0 = 1$ and the difficulty is to check the positivity preserving property. We call such sequences Markov sequences, and the problem is known as the Markov Sequence Problem (MSP in short).

When $\sum_n \lambda_n^2 < \infty$, the operator K may be represented through the $\mathcal{L}^2(\mu \otimes \mu)$ kernel $k(x, y) = \sum_n \lambda_n f_n(x) f_n(y)$ and, since all the functions f_n (except f_0) satisfy $\int f_n d\mu = 0$, the previous series is oscillating and positivity may not be checked directly from this representation.

The real parameters λ_n are the eigenvalues of the symmetric operator K , and it is well known (and quite immediate) that, if K is Markov, they must satisfy $|\lambda_n| \leq 1$. It is obvious that if $\lambda = (\lambda_n)$ and $\mu = (\mu_n)$ are Markov sequences, the $(\lambda_n \mu_n)$ is a Markov sequence. Moreover, for any $\theta \in [0, 1]$, $\theta \lambda + (1 - \theta) \mu = (\theta \lambda_n + (1 - \theta) \mu_n)$ is again a Markov sequence, and the simple limit of Markov sequences is again a Markov sequence. Therefore the question boils down to the determination the extremal points in the set of Markov sequences. The hypergroup property allows for such description.

The hypergroup property holds when there exists some point $x_0 \in X$ such that, for any $x \in X$, the sequence $\frac{f_n(x)}{f_n(x_0)}$ is a Markov sequence. Of course, for this to make sense, one needs for example some topology on X and require the functions (f_n) to be continuous for this topology.

When the hypergroup property holds, then these sequences $\frac{f_n(x)}{f_n(x_0)}$ are automatically the extremal points for the Markov sequence problem (see [2]) and then, for any Markov sequence (λ_n) , there exists a probability measure ν on Ω such that

$$(2.1) \quad \lambda_n = \int_X \frac{f_n(x)}{f_n(x_0)} \nu(dx).$$

To understand this representation, one should first extend the operator K by duality as an operator acting on probability measures. Since formally, δ_{x_0} has a density which may be written as $\sum_n f_n(x) f_n(x_0)$, then $K(\delta_{x_0})$ may be written as $(\sum_n \lambda_n f_n(x) f_n(x_0)) d\mu(x)$ and, with $\nu(dx) = K(\delta_{x_0})$, one has

$$\int \frac{f_n(x)}{f_n(x_0)} \nu(dx) = \lambda_n.$$

This of course is not really meaningful beyond the case of finite sets, and one should replace δ_{x_0} by a smooth approximation of it. In what follows, we shall have a symmetric diffusion operator L with eigenvectors f_n , and the associated heat kernel $P_t = e^{tL}$ will be such that $P_t(\delta_x)$ has a bounded density with respect to μ for any $t > 0$ and any x . One may then replace δ_{x_0} by $P_t(\delta_{x_0})$ in the previous argument, and let then $t \rightarrow 0$ to get the representation

When the space X is a finite set, is a consequence of the definition that this point x_0 must be the point with minimal measure (see [2]). Moreover, provided we chose the signs of (f_n) such that $f_n(x_0) \geq 0$, it is immediate (and does not require the finiteness of the space X) that the functions f_n reach their maximum (and the maximum value of their modulus) at x_0 .

This hypergroup property is strongly related with product formulae (see [19]). Indeed, assume that for any $x \in \Omega$, the Markov kernel K_x with eigenvalues $\frac{f_n(x)}{f_n(x_0)}$ has a \mathcal{L}^2 integrable density

$$K_x(f)(y) = \int f(z)k(x, y, z)\mu(dz).$$

Then, it is readily seen that

$$K(x, y, z) = \sum_n \frac{f_n(x)f_n(y)f_n(z)}{f_n(x_0)}$$

is a symmetric function of (x, y, z) and we get a product formula

$$(2.2) \quad f_n(x)f_n(y) = f_n(x_0) \int f_n(z)K(x, y, z)\mu(dz).$$

In [8], the authors provide a new and elegant proof of Gasper's result, through a method which proves to be efficient in many other situations, and that we shall describe now. For this, we first require some topology on X and the functions (f_n) to be continuous for this topology.

We assume that there exists a self-adjoint operator (in general unbounded and defined on a dense subset $\mathcal{D} \subset \mathcal{L}^2(\mu)$) $L : \mathcal{D} \mapsto \mathcal{L}^2(\mu)$ for which the functions f_n are eigenvectors, that is $Lf_n = \mu_n f_n$. Moreover, we assume that the eigenspaces of L for the eigenvalues μ_n are simple. In our context, one may chose \mathcal{D} as the set of the finite linear combinations of the functions f_i .

We also assume the existence of an auxiliary probability space (Y, \mathcal{B}, ν) with a self adjoint operator \hat{L} , again defined on a subset $\hat{\mathcal{D}} \subset \mathcal{L}^2(\nu)$, $\hat{L} : \hat{\mathcal{D}} \mapsto \mathcal{L}^2(\nu)$, together with some application $\pi : Y \mapsto X$. For a function $f : X \mapsto \mathbb{R}$, we denote by $\pi(f) : Y \mapsto \mathbb{R}$, the function $\pi(f)(y) = f(\pi(y))$. We suppose that $\pi(\mathcal{D}) \subset \hat{\mathcal{D}}$ and that $\hat{L}\pi = \pi L$. This property is often described as " L is the image of \hat{L} under π ", or that " \hat{L} projects onto L under π ", or even that " π intertwines L and \hat{L} ".

Moreover, one requires some measurable transformation $\Phi : Y \mapsto Y$ which commutes with the action of \hat{L} . That is, once again denoting $\Phi(f)(y) = f(\Phi(y))$, we assume that $\Phi : \hat{\mathcal{D}} \mapsto \hat{\mathcal{D}}$ and $\hat{L}\Phi = \Phi\hat{L}$.

In order for the next proposition to make sense, we shall impose some topology on Y , and require π and Φ to be continuous.

Let ξ be a random variable with values in Y , distributed according to ν . From the hypotheses, it is quite clear that the laws of $\pi(\xi)$ and of $\pi(\Phi(\xi))$ are μ . We look at the conditional law $k(x, dy)$ of $\pi(\Phi(\xi))$ given $\pi(\xi)$, where $x \mapsto k(x, dy)$ is continuous for the weak topology.

Proposition 2.1. *Assume that there exist some point $x_0 \in X$ such that the conditional law $k(x_0, dy)$ is the Dirac mass at some point $x_1 \in X$. Then, the sequence $\lambda_n = \frac{f_n(x_1)}{f_n(x_0)}$ is a Markov sequence.*

Proof. —

We consider the correlation operator, defined as

$$K(f)(x) = \int f(y)k(x, dy).$$

It is a Markov operator by construction. The assumption tells us that, for f continuous, $K(f)(x_0) = f(x_1)$. We shall see that this operator K corresponds to the choice of the Markov sequence $(\mu_n) = \frac{f_n(x_1)}{f_n(x_0)}$. By definition of the conditional expectation, for any pair of $\mathcal{L}^2(\mu)$ functions

$$\int_X fK(g)d\mu = \int_Y f(\pi(y))g(\pi(\Phi(y)))\nu(dy) = \int_Y \pi(f)\Phi\pi(g)d\nu.$$

We want to show first that for any n , f_n is an eigenvector for K . For this, chose any p and consider

$$\begin{aligned} \int_X Lf_pK(f_n) &= \int_Y \pi L(f_p)\Phi\pi(f_n)d\nu = \int_Y \hat{L}\pi(f_p)\Phi\pi(f_n)d\nu \\ &= \int_Y \pi(f_p)\hat{L}\Phi\pi(f_n)d\nu = \int_Y \pi(f_p)\Phi\pi L(f_n)d\nu \\ &= \lambda_n \int_Y \pi(f_p)\Phi\pi(f_n)d\nu = \lambda_n \int_X f_pK(f_n)d\mu. \end{aligned}$$

Therefore, this being valid for any p , $K(f_n)$ is an eigenvector for L with eigenvalue λ_n . Since the eigenspaces are one dimensional, we deduce that $K(f_n) = \mu_n f_n$. Applying this at the point x_0 , we see that $\mu_n = \frac{f_n(x_1)}{f_n(x_0)}$. ■

If we have at disposal a full family of such transformations Φ such that the associated points x_1 (depending on Φ) cover X , we conclude to the semigroup property. The main challenge for proving the hypergroup property, when we have a basis given as eigenvectors for some operator L , is then to construct the space Y , the operator \hat{L} , the projection π and the transformations $\Phi : Y \mapsto Y$, which satisfy the required properties. This is what we are going to do for the deltoid model in Section 5.

3 Symmetric diffusion operators

We briefly recall in this Section the context of symmetric diffusion operators, following [1], a specific context adapted to our setting.

For a given probability space (X, \mathcal{X}, μ) , we suppose given an algebra \mathcal{A} of functions such that $\mathcal{A} \subset \cap_{1 \leq p < \infty} \mathcal{L}^p(\mu)$, where \mathcal{A} is dense in $\mathcal{L}^2(\mu)$. A bilinear application $\Gamma : \mathcal{A} \times \mathcal{A} \mapsto \mathcal{A}$ is given such that, $\forall f \in \mathcal{A}$, $\Gamma(f, f) \geq 0$. If $\Phi : \mathbb{R}^n \mapsto \mathbb{R}$ is a smooth function such that for any $(f_1, \dots, f_n) \in \mathcal{A}^n$, $\Phi(f_1, \dots, f_n) \in \mathcal{A}$, then

$$\Gamma(\Phi(f_1, \dots, f_k), g) = \sum_i \partial_i \Phi \Gamma(f_i, g).$$

A linear operator L is defined through

$$(3.3) \quad \int_X fL(g)d\mu = - \int_X \Gamma(f, g)d\mu$$

and we assume that L maps \mathcal{A} into \mathcal{A} . We extend L into a self adjoint operator and we suppose and that \mathcal{A} is dense in the domain of L .

Then, for $f = (f_1, \dots, f_k)$, and again if Φ is smooth with $\Phi(f_1, \dots, f_n)$ in \mathcal{A} whenever $f_i \in \mathcal{A}$, then

$$(3.4) \quad L(\Phi(f)) = \sum_1^k \partial_i \Phi(f) L(f_i) + \sum_{i,j=1}^k \partial_{ij}^2 \Phi(f) \Gamma(f_i, f_j).$$

We have from (3.4)

$$\Gamma(f, g) = \frac{1}{2} \left(L(fg) - fL(g) - gL(f) \right).$$

If X is an open domain in \mathbb{R}^d , of some smooth manifold, then, in a local system of coordinates (x^1, \dots, x^d) , L may in general be written as

$$(3.5) \quad L(f) = \sum_{ij} g^{ij}(x) \partial_{ij}^2 f + \sum_i b^i(x) \partial_i f,$$

while

$$(3.6) \quad \Gamma(f, g) = \sum_{ij} g^{ij}(x) \partial_i f \partial_j g,$$

where $g^{ij} = \Gamma(x^i, x^j)$ and $b^i = L(x^i)$. The non negativity of Γ translates into the fact that the symmetric matrix $(g^{ij})(x)$ is everywhere non negative. Moreover, $g^{ij}(x) = \Gamma(x^i, x^j)$ and $b^i(x) = L(x^i)$. Observe then that in order to describe L , we just have to describe $L(x^i)$ and $\Gamma(x^i, x^j)$.

When the measure μ has a positive density ρ with respect to the Lebesgue measure, then the symmetry property of L translates into

$$(3.7) \quad L = \frac{1}{\rho} \sum_{ij} \partial_i \left(g^{ij} \rho \partial_j \right),$$

which shows that, in formula (3.5),

$$(3.8) \quad b_i = \sum_j \partial_j g^{ij} + g^{ij} \partial_j \log \rho,$$

and this formula may be applied in many circumstance to identify the measure density ρ up to some normalizing constant. Since we shall mainly use this setting for finite measure, we may always assume with no loss of generality that μ is a probability measure.

A central question is to determine on which set of functions one apply the operator L , particularly when Ω is bounded. Indeed, this requires to look at some self adjoint extension of L , and one needs in general to describe an algebra of functions which is dense in the domain of L and on which the integration by parts formula (3.3) holds true. This is done in general by prescribing boundary conditions on the functions f , such as Neumann or Dirichlet conditions.

In our setting however, we will always work on bounded open sets $\Omega \subset \mathbb{R}^d$ with piecewise smooth boundary $\partial\Omega$. Our functions g^{ij} and b^i will be smooth in Ω . Moreover, $g^{ij}(x)$ will be defined and smooth in some neighborhood of Ω . Then, suppose that in a neighborhood \mathcal{V} of any regular point x of the boundary, the boundary may be described through $\{F = 0\}$, where F is a smooth function, defined in \mathcal{V} and with real values. Then, our fundamental assumption is that

$$(3.9) \quad \Gamma(F, x_i) = 0 \text{ on } \partial\Omega \cap \mathcal{V}.$$

When this happens, we may chose for \mathcal{A} the algebra of smooth compactly supported functions defined in a neighborhood of Ω , referred in what follows as "smooth functions", and the integration by parts formula (3.3) holds true for those functions. In other words, for such operators, there is no need to consider boundary conditions of the functions in \mathcal{A} (see [4]). In the context of orthogonal polynomials on bounded domains that we shall consider, this property is always satisfied (see [4]).

We shall make a strong use of the notion of image operator, to fit with the setting described in Section 2. Suppose that we have a set of functions $X = (X_1, \dots, X_k)$ for which

$$L(X^i) = B^i(X), \Gamma(X^i, X^j) = G^{ij}(X),$$

then, of any smooth function $\Phi : \mathbb{R}^k \mapsto \mathbb{R}$, and thanks to equation (3.4), one has

$$L(\Phi(X)) = L_1(\Phi)(X),$$

where

$$L_1 = \sum_{ij} G^{ij} \partial_{ij}^2 + \sum_i B^i \partial_i.$$

This is again a symmetric diffusion operator, defined on the image $\Omega_1 = X(\Omega)$, and its reversible measure is the image of μ through X . This new diffusion operator L_1 is therefore the image of L under Φ , in the sense described in Section 2 : $\Phi L_1 = L \Phi$ by construction.

In the next sections, we shall always work on polynomials, moreover in even dimensions $2k$. We shall suppose that the coordinates are paired as $(x_p, y_p), p = 1, \dots, k$. In this context, it is often quite simpler to use complex coordinates (that is identify $\mathbb{R}^{2k} \simeq \mathbb{C}^k$), setting $z_p = x_p + iy_p, \bar{z}_p = x_p - iy_p$ and one has to describe then, using linearity and bilinearity $\Gamma(z_p, z_q), \Gamma(\bar{z}_p, \bar{z}_q), \Gamma(z_p, \bar{z}_q)$, together with $L(z_p)$ and $L(\bar{z}_p)$.

For example, $\Gamma(z_p, z_p) = \Gamma(x_p, x_p) - \Gamma(y_p, y_p) + 2i\Gamma(x_p, y_p)$.

The positivity of the metric here may be checked according to the parity of k . For example, a careful inspection shows that indeed, the determinant of the metric in the variables (z_p, \bar{z}_p) is $(-1)^k 4^k \det(g)$, where the determinant computed in real variables.

4 The deltoid model

The deltoid curve is a degree 4 algebraic plane curve which may be parametrized as

$$x(t) = \frac{1}{3}(2 \cos t + \cos 2t), \quad y(t) = \frac{1}{3}(2 \sin t - \sin 2t)$$

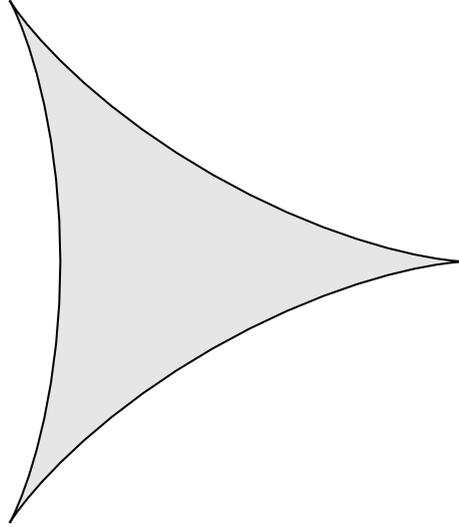


Figure 1: The deltoid domain.

The factor $1/3$ in the previous formulae are just here to simplify future computations, but play no fundamental rôle. The connected component Ω of the complementary of the curve which contains 0 is a bounded open set, that we refer to as the deltoid domain. It turns out that there exist on this domain a one parameter family $L^{(\lambda)}$ of symmetric diffusion operator which may be diagonalized in a basis of orthogonal polynomials. It was introduced in [16, 17], and further studied in [25]. This is one of the 11 families of sets carrying such diffusion operators, as described in [4].

In order to describe the operator, and thanks to the diffusion property (3.4), it is enough to describe $\Gamma(x, x)$, $\Gamma(x, y)$, $\Gamma(y, y)$, $L^{(\lambda)}(x)$ and $L^{(\lambda)}(y)$ (the operator Γ does not depend on λ here).

The symmetric matrix $\begin{pmatrix} \Gamma(x, x) & \Gamma(x, y) \\ \Gamma(y, x) & \Gamma(y, y) \end{pmatrix}$ is referred to in what follows as the metric associated with the operator, although properly speaking it is in fact a co-metric. We may also use the complex structure of $\mathbb{R}^2 \simeq \mathbb{C}$, and the complex variables $Z = x + iy$, $\bar{Z} = x - iy$, and it turns out that the formulas are much simpler under this description.

The operator $L^{(\lambda)}$ is then described as

$$(4.10) \quad \begin{cases} \Gamma(Z, Z) = \bar{Z} - Z^2, \\ \Gamma(\bar{Z}, \bar{Z}) = Z - \bar{Z}^2, \\ \Gamma(\bar{Z}, Z) = 1/2(1 - Z\bar{Z}), \\ L(Z) = -\lambda Z, L(\bar{Z}) = -\lambda \bar{Z}, \end{cases}$$

where $\lambda > 0$ is a real parameter.

The boundary of the domain Ω turns out to be the curve with equation

$$P(Z, \bar{Z}) := \Gamma(Z, \bar{Z})^2 - \Gamma(Z, Z)\Gamma(\bar{Z}, \bar{Z}) = 0,$$

and inside the domain Ω , the associated metric is positive definite, so that it corresponds to some elliptic operator in Ω . Moreover, for this function P , the boundary condition (3.9) is satisfied, with

$$\Gamma(P, Z) = -3ZP, \quad \Gamma(P, \bar{Z}) = -3\bar{Z}P.$$

The reversible measure associated with it, easily identified through equation (3.8), has density $C_\alpha P(Z, \bar{Z})^\alpha$ with respect to the Lebesgue measure, where $\lambda = \frac{1}{2}(6\alpha + 5)$, and is a probability measure exactly when $\lambda > 0$ (see [25] for more details). We shall refer to this probability measure on Ω as $\mu^{(\lambda)}$.

It is quite immediate that the operator $L^{(\lambda)}$ commutes with the transformation $Z \mapsto \bar{Z}$. Indeed, there is another invariance : the transformation $Z \mapsto e^{i\theta}Z$ commutes with $L^{(\lambda)}$ provided $e^{3i\theta} = 1$. Therefore, everything is invariant under $Z \mapsto jZ$ and $Z \mapsto \bar{j}Z$, where j and \bar{j} are the third roots of unity in the complex plane.

From the form of the operator, we easily see that $L^{(\lambda)}$ maps the set \mathcal{P}_n of polynomials with total degree n (in the variables (Z, \bar{Z}) or equivalently in the variables (x, y)) into itself, and being symmetric, may be diagonalized in a basis of orthogonal polynomials. We refer to [25] for a complete description of these polynomials. In what follows, and in the rest of the paper, we shall forget the dependence in λ of the polynomials, in order to have lighter notations.

The eigenspaces of $L^{(\lambda)}$ are described in [25]. For any (k, n) , there is a unique polynomial $R_{n,k}$ which is an eigenvector with eigenvalue $-\lambda_{n,k} = -((\lambda - 1)(n + k) + n^2 + k^2 + nk)$, which is a polynomial in the variables (Z, \bar{Z}) with real coefficients, and has a unique highest degree term $Z^n \bar{Z}^k$. When $n \neq k$, the eigenspaces have dimension 2, and we want to distinguish between the symmetric and antisymmetric part, under the symmetry $Z \mapsto \bar{Z}$. We therefore chose when $n \neq k$ the basis $\frac{1}{2}(R_{n,k} + R_{k,n})$ and $\frac{1}{2i}(R_{n,k} - R_{k,n})$. The following proposition summarizes a few properties of this basis of eigenvectors.

Proposition 4.1. *For any $\lambda > 0$, for any $n, k \in \mathbb{N}^2$, with $k \neq n$, there are exactly, up to a sign, two real valued polynomials $P_{n,k}(Z, \bar{Z})$ and $Q_{n,k}(Z, \bar{Z})$, with degree $n+k$, $P_{n,k}$ being symmetric and $Q_{n,k}$ antisymmetric under the symmetry $(Z, \bar{Z}) \mapsto (\bar{Z}, Z)$, with norm 1 in $\mathcal{L}^2(\mu_\lambda)$, such that*

$$L^{(\lambda)}P_{n,k} = -\lambda_{n,k}P_{n,k}, \quad L^{(\lambda)}Q_{n,k} = -\lambda_{n,k}Q_{n,k},$$

where $\lambda_{n,k} = (\lambda - 1)(n + k) + n^2 + k^2 + nk$.

When $n = k$, there is exactly one such eigenvector $P_{n,n}$, with eigenvalue $-\lambda_{n,n}$, and it is symmetric in (Z, \bar{Z}) .

Moreover, $P_{n,k}$ has real coefficients, $Q_{n,k}$ purely imaginary ones, and they satisfy

$$(4.11) \quad P_{n,k}(jZ, \bar{j}\bar{Z}) + iQ_{n,k}(jZ, \bar{j}\bar{Z}) = \bar{j}^{n-k}(P_{n,k}(Z, \bar{Z}) + iQ_{n,k}(Z, \bar{Z})).$$

Observe that $P_{n,k}$ and $Q_{n,k}$ are real valued. We shall investigate the hypergroup property in terms of this basis. In order to have lighter notations, we shall often write $P_{n,k}(Z)$ and $Q_{n,k}(Z)$ instead of $P_{n,k}(Z, \bar{Z})$ and $Q_{n,k}(Z, \bar{Z})$, although they are really polynomials of both variables Z and \bar{Z} (in particular, they are not harmonic in \mathbb{C}). Moreover, recall that we shall by convention set $Q_{n,n} = 0$.

Proof. — Let $R_{n,k}(Z, \bar{Z})$ the unique eigenvector with unique highest degree term $Z^n \bar{Z}^k$. From the invariance under conjugacy, then when $n \neq k$, the conjugate of $R_{n,k}$, that is $\bar{R}_{n,k}(Z, \bar{Z})$ is an eigenvector with same eigenvalue, and looking at the highest degree term, it is therefore $R_{k,n}(Z, \bar{Z})$. Due to the conjugacy invariance, $R_{n,k}$ and $R_{k,n}$ have the same $\mathcal{L}^2(\mu_\lambda)$ norm.

For $n \neq k$, let $\hat{P}_{n,k}$ and $\hat{Q}_{n,k}$ be the symmetric and antisymmetric eigenvectors of $L^{(\lambda)}$ with dominant terms $\frac{1}{2}(Z^n \bar{Z}^k + Z^k \bar{Z}^n)$ and $\frac{-i}{2}(Z^n \bar{Z}^k - Z^k \bar{Z}^n)$ respectively, so that $\hat{P}_{n,k}$ and $\hat{Q}_{n,k}$ take real values, with $\hat{Q}_{n,k}$ vanishing on the real axis. (By convention, set $\hat{Q}_{n,n} = 0$). The fact that $\int (P_{n,k}^{(\lambda)})^2 d\mu_\lambda = \int (P_{k,n}^{(\lambda)})^2 d\mu_\lambda$ (from the conjugacy invariance) shows that $\hat{P}_{n,k}$ and $\hat{Q}_{n,k}$ are orthogonal. Moreover, due to the invariance under $Z \mapsto jZ$, we know that $\hat{P}_{n,k}(jZ)$ and $\hat{Q}_{n,k}(jZ)$ are again eigenvectors for $L^{(\lambda)}$ with the same eigenvalue, and therefore are linear combinations of $\hat{P}_{n,k}$ and $\hat{Q}_{n,k}$. Looking at the highest degree terms, one sees that

$$(4.12) \quad \begin{cases} \hat{P}_{n,k}(jZ) = \left(\frac{j^{n-k} + \bar{j}^{n-k}}{2}\right) \hat{P}_{n,k}(Z) + i \left(\frac{j^{n-k} - \bar{j}^{n-k}}{2}\right) \hat{Q}_{n,k}(Z), \\ \hat{Q}_{n,k}(jZ) = -i \left(\frac{j^{n-k} - \bar{j}^{n-k}}{2}\right) \hat{P}_{n,k}(Z) + \left(\frac{j^{n-k} + \bar{j}^{n-k}}{2}\right) \hat{Q}_{n,k}(Z). \end{cases}$$

In other words, $(\hat{P}_{n,k} + i\hat{Q}_{n,k})(jZ) = \bar{j}^{n-k}(\hat{P}_{n,k} + i\hat{Q}_{n,k})(Z)$. As a consequence, since $\int \hat{P}_{n,k}(jZ)^2 d\mu_\lambda = \int \hat{P}_{n,k}(Z)^2 d\mu_\lambda$, one sees that $\|\hat{P}_{n,k}\|^2 = \|\hat{Q}_{n,k}\|^2$ when $n - k \not\equiv 0 \pmod{3}$. We do not know if this is true for $n = k \pmod{3}$. Of course, the problem does not exist when $n = k$.

We therefore may use the basis $P_{n,k} = a_{n,k} \hat{P}_{n,k}$ and $Q_{n,k} = b_{n,k} \hat{Q}_{n,k}$ as an orthonormal basis for the eigenspace associated with the eigenvalue $\lambda_{n,k}$, with $a_{n,k} = b_{n,k}$ when $n - k \not\equiv 0 \pmod{3}$, and then equation (4.12) translates into (4.11). If $n - k \equiv 0 \pmod{3}$, then (4.12) is trivial since in this case $P_{n,k}(jZ) = P_{n,k}(Z)$, $Q_{n,k}(jZ) = Q_{n,k}(Z)$. ■

There are two particular cases which are worth understanding, namely $\lambda = 1$ and $\lambda = 4$, corresponding to the parameters $\alpha = \pm 1/2$. These two models show the relation with the A_2 root system, and indeed our polynomials are the examples of Heckman-Opddam polynomials associated to this root system. We briefly present those two models, referring to [25] for more details.

In the first case $\lambda = 1$, one sees that this operator is nothing else that the image of the Euclidean Laplace operator on \mathbb{R}^2 acting on the functions which are invariant under the symmetries around the lines of a regular triangular lattice.

Indeed, consider the three unit roots of identity in \mathbb{C} , say $(e_1, e_2, e_3) = (1, j, \bar{j})$. Then, consider the functions $z_k : \mathbb{C} \mapsto \mathbb{C}$ which are defined as

$$(4.13) \quad z_k(z) = e^{i\Re(z\bar{e}_k)}.$$

They satisfy $|z_k| = 1$ and $z_1 z_2 z_3 = 1$.

For the 2 dimensional Laplace operator, the functions z_i satisfy

$$\Delta(z_i) = -z_i, \Delta(\bar{z}_i) = -\bar{z}_i,$$

and, if we denote for $i \neq j \in \{1, 2, 3\}$ by $c(i, j)$ the index in $\{1, 2, 3\}$ which differs from i and j , that is $\{c(i, j), i, j\} = \{1, 2, 3\}$.

$$(4.14) \quad \begin{cases} \Gamma(z_i, z_i) = -z_i^2, \Gamma(\bar{z}_i, \bar{z}_i) = -\bar{z}_i^2, \Gamma(z_i, \bar{z}_i) = 1 \\ \Gamma(z_i, z_j) = \frac{1}{2}\bar{z}_{c(i,j)}, \Gamma(\bar{z}_i, \bar{z}_j) = \frac{1}{2}z_{c(i,j)}, \Gamma(z_i, \bar{z}_j) = \frac{1}{2}z_i\bar{z}_j, i \neq j \end{cases}$$

Let now $Z = \frac{1}{3}(z_1 + z_2 + z_3)$. It is easily seen that, for the Euclidean Laplace operator on \mathbb{R}^2 , Z and \bar{Z} satisfy the relations (4.10) with $\lambda = 1$. Moreover, the function $Z : \mathbb{C} \mapsto \mathbb{C}$ is a diffeomorphism between the interior Ω_1 of the triangle T and Ω , where T is one of the equilateral triangles which contains the two edges 0 and $4\pi/3$. The functions which are invariant under the symmetries of the triangular lattice generated by this triangle T are exactly functions of Z . In particular, the application $(z_1, z_2) \mapsto Z = \frac{1}{3}(z_1 + z_2 + \bar{z}_1\bar{z}_2)$ maps $\mathcal{S}^1 \times \mathcal{S}^1$ onto the closure of the deltoid domain Ω . The boundary of Ω is the image under Z of the boundary of the triangle, and also of the set where 2 of the three variables (z_1, z_2, z_3) coincide. The cusps of the deltoid model are the images of the points $z_1 = z_2 = z_3$, that is $z_i = 1, j, \bar{j}$.

In particular, the set of all the points $Z = \frac{1}{3}(z_1 + z_2 + z_3)$ such that $|z_i| = 1$ and $z_1 z_2 z_3 = 1$ is the closure of Ω . Indeed, given $Z \in \Omega$, there exist, up to permutation, 3 unique and distinct complex numbers z_i satisfying $|z_i| = 1$ and $z_1 z_2 z_3 = 1$ such that $Z = \frac{1}{3}(z_1 + z_2 + z_3)$. They are the three distinct roots of the equation $P(X) = X^3 - 3ZX^2 + 3\bar{Z}X - 1 = 0$. It is worth for this to observe that the discriminant of P is, up to a numerical constant, equal to $P(Z, \bar{Z})$, and therefore does not vanish in Ω .

The second description comes from the Casimir operator on $SU(3)$. For $SU(d)$, one may describe this operator through its action on the various entries z_{ij} of the matrices $g \in SU(d)$. Up to some scaling factor, one has

$$\begin{cases} \Gamma^{SU(d)}(z_{ij}, z_{kl}) = z_{ij}z_{kl} - dz_{il}z_{kj}, \\ \Gamma^{SU(d)}(\bar{z}_{ij}, \bar{z}_{kl}) = \bar{z}_{ij}\bar{z}_{kl} - d\bar{z}_{il}\bar{z}_{kj}, \\ \Gamma^{SU(d)}(z_{ij}, \bar{z}_{kl}) = d\delta_{ik}\delta_{jl} - z_{ij}\bar{z}_{kl} \\ \mathbb{L}^{SU(d)}(z_{ij}) = -(d^2 - 1)z_{ij}, \mathbb{L}^{SU(d)}(\bar{z}_{ij}) = -(d^2 - 1)\bar{z}_{ij} \end{cases}$$

Then, choosing $Z = \frac{1}{3}(z_{11} + z_{22} + z_{33})$, one obtains the relations (4.10) with $\lambda = 4$ when applying $\frac{3}{4}\mathbb{L}^{SU(3)}$ to functions of Z and \bar{Z} . For this particular case, we may make full use of the group structure of $SU(3)$ to obtain the hypergroup property for the deltoid model, at any of the points $Z = 1, Z = e^{2i\pi/3}, Z = e^{4i\pi/3}$. It is enough to follow the scheme described in Section 2, choosing $Y = SU(3)$, π being the application which to some matrix $g \in SU(3)$ associates $\frac{1}{3}\text{trace}(g)$, and Φ being $g \mapsto g_0 g$, for any $g_0 \in SU(3)$. . If $\pi(g) = 1$ for example, then $g = \text{Id}$, and the

conditional of $\pi(\Phi(g))$ knowing that $\pi(g) = 1$ is the Dirac mass at $\pi(g_0)$. However, the eigenspaces being two dimensional, this introduces some extra complexity that we shall examine in Section 6.

Similarly, using the representation for $\lambda = 1$, we may also prove the hypergroup property in this case, using the translations in \mathbb{R}^2 .

Our aim in what follows is to propose another 6-dimensional model which projects onto the deltoid model in the general case (that is for $\lambda \neq 1, 4$), and on which we have enough symmetry to use the machinery described in Section 2.

5 A 6-dimensional model for the deltoid

In this section, we construct a symmetric diffusion operator \hat{L}_λ in dimension 6 (or more precisely on an open bounded set $\Omega_1 \subset \mathbb{C}^3$), such that the image of the operator \hat{L}_λ under the projection $\pi : \mathbb{C}^3 \mapsto \mathbb{C}$ which is $\pi(z_1, z_2, z_3) = \frac{1}{3}(z_1 + z_2 + z_3)$ is exactly our deltoid model with parameter λ . Moreover, this operator may be diagonalized in a complete system of orthogonal polynomials for its reversible measure.

We consider a diffusion operator in \mathbb{C}^3 , acting on 3 complex variables z_1, z_2, z_3 , defined as follows. We set

$$(5.15) \quad \begin{cases} \hat{\Gamma}(z_i, z_j) = \frac{3}{2}\bar{z}_{c(i,j)} - z_i z_j & \text{for } i \neq j \\ \hat{\Gamma}(\bar{z}_i, \bar{z}_j) = \frac{3}{2}z_{c(i,j)} - \bar{z}_i \bar{z}_j & \text{for } i \neq j \\ \hat{\Gamma}(z_i, z_i) = -z_i^2 \\ \hat{\Gamma}(\bar{z}_i, \bar{z}_i) = -\bar{z}_i^2 \\ \hat{\Gamma}(z_i, \bar{z}_j) = \frac{3}{2}\delta_{ij} - \frac{1}{2}z_i \bar{z}_j \\ \hat{L}(z_i) = -\lambda z_i, \hat{L}(\bar{z}_i) = -\lambda \bar{z}_i, \end{cases}$$

where as before $c(i, j)$ the index in $\{1, 2, 3\}$ which differs from i and j . It is worth to observe that those relations fit with (4.14) whenever $|z_i| = 1$ and $z_1 z_2 z_3 = 1$, so that those equations may be seen as an extension of (4.14) in the interior of some domain bounded by those equations. We shall see that it is indeed the case, when we shall describe the domain.

The first task is to observe that \hat{L} projects onto the deltoid model in the sense of Section 2, with the same parameter λ .

Proposition 5.1. *For $z = (z_1, z_2, z_3) \in \mathbb{C}^3$, let $\pi(z) := Z = \frac{1}{3}(z_1 + z_2 + z_3)$. Then, one has*

$$\begin{cases} \hat{\Gamma}(Z, Z) = \bar{Z} - Z^2 \\ \hat{\Gamma}(\bar{Z}, \bar{Z}) = Z - \bar{Z}^2 \\ \hat{\Gamma}(Z, \bar{Z} = \frac{1}{2}(1 - Z\bar{Z})) \\ \hat{L}(Z) = -\lambda Z, \hat{L}(\bar{Z}) = -\lambda \bar{Z} \end{cases}$$

This comes from an immediate computation. We see that these computations fit with formulae (4.10).

The next task is to show that this corresponds to some elliptic symmetric diffusion generator in some bounded domain $\Omega_1 \subset \mathbb{C}^3$.

For this, let us introduce the following notations. We write $z_j = r_j e^{i\theta_j}$, with $r_j = |z_j|$, and

$$\begin{cases} \theta = \theta_1 + \theta_2 + \theta_3 \\ S_1 = r_1^2 + r_2^2 + r_3^2, \\ S_2 = r_1^4 + r_2^4 + r_3^4 \\ \sigma = r_1 r_2 r_3 \cos(\theta) \end{cases}$$

and let

$$(5.16) \quad \begin{cases} P_1(z_1, z_2, z_3) = 2 - (S_1 + 1)^2 + 2S_2 + 8\sigma. \\ P_2(z_1, z_2, z_3) = 2(S_2 - 1) - (S_1 - 1)^2. \end{cases}$$

Let us describe the domain first. Let D be the determinant of the matrix $\hat{\Gamma}$ (in both coordinates (z_i, \bar{z}_i)). Then, define Ω_1 to be the connected component of $\mathbb{C}^3 \setminus \{D = 0\}$ which contains $(0, 0, 0)$.

Proposition 5.2. *We have*

1. $D = \frac{3^5}{2^6} P_1 P_2$;
2. The operator \hat{L} is elliptic in Ω_1 ;
3. Ω_1 is bounded;
4. $P_1 > 0$ on Ω_1 and $P_2 < 0$ on Ω_1 .
5. The boundary of Ω_1 is included in $\{P_1 = 0\}$.

Proof. — The first item results from a direct computation. Since it is straightforward although quite technical, we recommend that the reader uses a computer program to check it.

For the second item, let us observe that the metric at $z = (0, 0, 0)$ is the classical Euclidean metric of \mathbb{C}^3 , up to some scaling factor. Therefore, it remains elliptic as long as no eigenvalue of the metric vanishes, that is in the connected component of the set $\{D \neq 0\}$ which contains $\{0, 0, 0\}$. It is therefore elliptic in Ω_1 .

For the third assertion, observe that as long as $\hat{\Gamma}$ is positive definite, then one has $\Gamma(z_i, \bar{z}_i)^2 > \Gamma(z_i, z_i)\Gamma(\bar{z}_i, \bar{z}_i)$. This translates into $(3 + r_i^2)(1 - r_i^2) > 0$ so that $r_i \leq 1$ in Ω_1 .

We have $P_2(0, 0, 0) < 0$ and $P_1(0, 0, 0) > 0$, so that, in Ω_1 , $P_1 > 0$ and $P_2 < 0$.

It remains to prove the last assertion. Any point $z \in \partial\Omega_1$ satisfies $P_1(z)P_2(z) = 0$. Let $z = (z_1, z_2, z_3) \in \partial\Omega_1$ such that $P_2(z) = 0$. We shall prove that we also have $P_1(z) = 0$ at this point.

We know that $S_1 \geq 0$, and it is immediate that $\frac{1}{3}S_1^2 \leq S_2 \leq S_1^2$. From the upper bound, we get $P_2 \leq (S_1 - 1)(S_1 + 3)$. Therefore, at a point where $P_2 = 0$, one has $S_1 \geq 1$. At such a point, we also have $P_1 = 4(1 - S_1 + 2\sigma)$.

We now want to estimate the maximum possible value of $r_1 r_2 r_3$ on the set where $P_2 = 0$ and $0 \leq r_i \leq 1$. We shall see that this maximum value is $(S_1 - 1)/2$. Setting $x_i = r_i^2$, it amounts to look for the maximum value of $x_1 x_2 x_3$ on the set where

$S_2 := \sum_i x_i^2 = 1 + \frac{1}{2}(S - 1)^2$, where $S = \sum_i x_i$ and $0 \leq x_i \leq 1$. Looking at the Lagrange multiplier, one sees that, if the maximum is not attained at the boundary of the set, then one must have

$$\begin{cases} x_2 x_3 = \lambda(1 + x_1 - x_2 - x_3) \\ x_1 x_3 = \lambda(1 + x_2 - x_1 - x_3) \\ x_1 x_2 = \lambda(1 + x_3 - x_1 - x_2) \end{cases}$$

from which using $S_2 = 1 + \frac{1}{2}(S - 1)^2$, one deduces that $3\lambda(3 - S) = (S - 1)(S + 3)$ and therefore $\lambda > 0$. Multiplying the first equation by x_1 we find then that

$$\pi := x_1 x_2 x_3 = \lambda(x_1^2 + x_1 + 2\lambda(x_1 - 1)).$$

The same relation being true for x_2 and x_3 , the values for x_i must lie among the two solutions of the equation

$$\pi = \lambda(x^2 + x + 2\lambda(x - 1)).$$

But $\pi > 0$ and $\lambda > 0$. Unless they are all equal, they may not be all positive. Therefore, either the maximum is attained at a point (x, x, x) , in which case this common value is 1 and $(S_1 - 1)/2$, either one of the values for x_i is 1, and the fact that $P_2 = 0$ implies that the two other values are equal, in which case again the value is $(S - 1)/2$.

Therefore, when $P_2 = 0$, $P_1 \leq 1 - S_1 + \frac{1}{2}(S_1 - 1) = \frac{1}{2}(1 - S_1) \leq 0$. Since $P_1 > 0$ in Ω_1 , then any point in $\partial\Omega_1$ which satisfies $P_2(z) = 0$ also satisfies $P_1(z) = 0$. ■

Remark 5.3. *It may be worth to observe that P_1 may take a somewhat simpler form. Setting*

$$\begin{cases} \sigma_0 = r_1 + r_2 + r_3 \\ \sigma_1 = -r_1 + r_2 + r_3 \\ \sigma_2 = r_1 - r_2 + r_3 \\ \sigma_3 = r_1 + r_2 - r_3 \end{cases}$$

and

$$(5.17) \quad \begin{cases} S = (1 + \sigma_0)(1 - \sigma_1)(1 - \sigma_2)(1 - \sigma_3) \\ D = (1 - \sigma_0)(1 + \sigma_1)(1 + \sigma_2)(1 + \sigma_3) \end{cases}$$

one may write

$$P_1 = S \cos^2(\theta/2) + D \sin^2(\theta/2).$$

Our next result shows that indeed the operator $\hat{\Gamma}$ defined on Ω_1 satisfies the boundary condition (3.9).

Proposition 5.4.

1. *With P_1 defined in (5.16), one has*

$$(5.18) \quad \hat{\Gamma}(\log(P_1), z_i) = -3z_i, \hat{\Gamma}(\log(P_1), \bar{z}_i) = -3\bar{z}_i.$$

2. For $i = 1, \dots, 3$, one has

$$(5.19) \quad \begin{cases} \sum_{j=1}^3 \partial_{z_j} \hat{\Gamma}(z_j, z_i) + \partial_{\bar{z}_j} \hat{\Gamma}(\bar{z}_j, z_i) = -\frac{11}{2} z_i \\ \sum_{j=1}^3 \partial_{z_j} \hat{\Gamma}(z_j, \bar{z}_i) + \partial_{\bar{z}_j} \hat{\Gamma}(\bar{z}_j, \bar{z}_i) = -\frac{11}{2} \bar{z}_i \end{cases}$$

Proof. — These formulas may be checked with a direct and tedious computation. However, we have no simple interpretation, beyond mere calculus, why this is true. The first one is the condition required for the operator \hat{L} on Ω_1 to be diagonalizable in a system of orthogonal polynomials. The second will be used to identify the reversible measure. ■

Remark 5.5. Together with the fact that $\Gamma(z_i, z_j)$, $\Gamma(z_i, \bar{z}_j)$ and $\Gamma(\bar{z}_i, \bar{z}_j)$ are polynomials in the variables (z_i, \bar{z}_i) , and following [4], equation (5.18) is the key identity which insures that the operator \hat{L} may be diagonalized in a basis of orthogonal polynomials.

We now prove the following lemma

Lemma 5.6. For any $\beta > -1$, $\int_{\Omega_1} P_1^\beta dx < \infty$, where dx is the Lebesgue measure on $\Omega_1 \subset \mathbb{C}^3$.

Before proving Lemma 5.6, we deduce the following

Corollary 5.7. For any $\lambda > 5/2$, the operator $\hat{L}^{(\lambda)}$ is reversible with respect to the probability measure $C_\beta P_1^\beta dx$, where $\beta = \frac{1}{6}(2\lambda - 11)$, dx is the Lebesgue measure in Ω_1 , and C_β is the normalizing constant.

Proof. — (Of Corollary 5.7) This is a direct consequence of the general formula (3.8) and Proposition 5.4. ■

Remark 5.8. It is worth to observe that $\beta = -1$ corresponds to $\lambda = 5/2$, while

$$\hat{L}(P_1) = -4(\lambda + 2)P_1 + (2\lambda - 5)(P_2 - 3P_1)/6.$$

This suggest that, for this precise value $\lambda = 5/2$, the measure is concentrated on the surface $\{P_1 = 0\}$, and the associated process lives indeed on the boundary of the domain Ω_1 .

Moreover, this limit $\beta = 5/2$ corresponds for the projected model on the deltoid domain to the case $\alpha = 0$, that is when the reversible measure of the image operator on Ω is the Lebesgue measure.

Proof. — (Of Lemma 5.6)

When looking at the behavior of P_1 near a regular point of the boundary $\{P_1 = 0\}$, it is clear that the condition $\beta > -1$ is necessary for P_1^β to be locally integrable

near such a point, in the domain $\{P_1 > 0\}$. But we have also to consider the behavior near singular points. Set as before

$$\begin{cases} S_1 = r_1^2 + r_2^2 + r_3^2, \\ S_2 = r_1^4 + r_2^4 + r_3^4, \\ \tau = r_1 r_2 r_3. \end{cases}$$

and, writing $z_j = r_j e^{i\theta_j}$, set $t = \theta_1 + \theta_2 + \theta_3$. Then we have

$$P_1 = 2 - (S_1 + 1)^2 + 2S_2 + 8\tau \cos(t),$$

and we want to show that for $\beta > -1$, P_1^β is locally integrable with respect to the measure $r_1 r_2 r_3 dr_1 dr_2 dr_3 dt$. on the domain $0 < r_i < 1$, $t \in (0, 2\pi)$, and $P_1 > 0$. We set

$$A = 2 - (S_1 + 1)^2 + 2S_2, \quad B = 8\tau,$$

The first thing is to compute

$$(5.20) \quad I(A, B) = \int_{A+B \cos(t) > 0, t \in (0, 2\pi)} (A + B \cos(t))^\beta dt$$

$$(5.21) \quad = 2 \int_{-1}^1 \mathbb{1}_{A+Bu > 0} (A + Bu)^\beta \frac{du}{\sqrt{1-u^2}},$$

that we want to estimate up to some constants depending only on β . For this we write $I \simeq J$ when $c_\beta \leq \frac{I}{J} \leq C_\beta$, for two positive constants c_β and C_β .

We cut the integral in (5.21) into \int_0^1 and \int_{-1}^0 . Then we change variables to write both integrals as \int_0^1 , and again change variables $v \mapsto 1 - v$. Finally, we get

$$I(A, B) \simeq \int_0^1 \mathbb{1}_{A+B-Bv > 0} (A + B - Bv)^\beta \frac{dv}{v^{1/2}} + \int_0^1 \mathbb{1}_{A-B+Bv > 0} (A - B + Bv)^\beta \frac{dv}{v^{1/2}}.$$

Write this $I(A, B) = I_1 + I_2$.

Then

$$\begin{cases} I_1(A, B) = 0 \text{ if } A + B < 0, \\ I_1(A, B) \simeq (A + B)^{\beta+1/2} B^{-1/2} \text{ if } A + B > 0 \end{cases}$$

and

$$\begin{cases} I_2 = 0 \text{ if } A < 0, \\ I_2 \simeq A^{\beta+1} B^{-1} \text{ if } 0 < A < B \\ I_2 \simeq A^\beta \text{ if } 0 < B < A \end{cases}$$

We now have to consider the integral of I_1 and I_2 on the image of the domain Ω_1 through the projection $(z_1, z_2, z_3) \mapsto (r_1, r_2, r_3)$, that is the integral of I_1 and I_2 with respect to the measure $r_1 r_2 r_3 dr_1 dr_2 dr_3$, on the domain $\{0 \leq r_i \leq 1\} \cap \{A + B > 0\} \cap \{A - B > 0\}$. An easy inspection show that $A + B = S$ and $A - B = D$, where S and D are given in formulae (5.17). We see that $I_1 \geq C(l_1 l_2 l_3)^{\beta+1/2} B^{-1/2}$, where l_i are 3 independent affine forms, restricted to a domain where they are positive,

and C is a constant. Such a function is integrable with respect to the measure $B^{1/2} \mathbf{1}_{0 < r_1 \leq 1} \mathbf{1}_{0 < r_2 \leq 1} \mathbf{1}_{0 < r_3 \leq 1} dr_1 dr_2 dr_3$. This integral is finite when $\beta > -1$.

The only problem comes from $\int_{0 < B < A, 0 < r_i \leq 1} A^\beta r_1 r_2 r_3 dr_1 dr_2 dr_3$. We see that one may reduce to a neighborhood of a point where both A and B vanish, which means that (up to a permutation of the indices) $r_1 = 0, r_2 + r_3 = 1$. Once again, integrating in a area where $0 \ll r_2 \ll 1$ causes no problem, and it remains to consider the integral on an area where $r_2 = 0, r_3 = 1$ (or symmetrically). Then, in a neighborhood of this point, we may bound $A \geq A - B = D$, and consider the integral of D^β on $D > 0$. In this neighborhood, with the notations of Remark 5.3, D is equivalent to $(1 - \sigma_0)(1 + \sigma_3)$, which is the product of two independent linear forms (in the variables r_1, r_2, r_3), and is locally integrable as soon as $\beta > -1$. ■

6 Hypergroup property for the deltoid model

For $\lambda > 5/2$, the operator $\hat{L}^{(\lambda)}$ provides a symmetric diffusion operator on the set Ω_1 which projects onto the deltoid model $L^{(\lambda)}$ with the same parameter through the map $\pi(z_1, z_2, z_3) = \frac{1}{3}(z_1 + z_2 + z_3)$.

We now observe that, for any $\theta = (\theta_1, \theta_2) \in \mathbb{R}^2$, the operator \hat{L} is invariant under the transformations $\Phi_\theta(z_1, z_2, z_3) = (e^{i\theta_1} z_1, e^{i\theta_2} z_2, e^{-i(\theta_1 + \theta_2)} z_3)$. It is worth to notice that $\Phi_{\theta_1} \Phi_{\theta_2} = \Phi_{\theta_1 + \theta_2}$, and that the measure $\mu_\alpha = P_1^\alpha dz$ is invariant under Φ_θ . Therefore, if we define $\Phi_\theta(f)(z) = f(\Phi_\theta(z))$, one see that the adjoint of Φ_θ is $\Phi_{-\theta}$.

To fit with the situation described in Section 2, it remains to identify the point x_0 of Proposition 2.1. Let then $x_0 = \pi(1, 1, 1) = 1$: it belongs to the boundary of Ω . According to the analysis of the flat model ($\lambda = 1$), $\theta \mapsto \Phi_\theta(1, 1, 1)$ is onto the deltoid domain Ω , and, when θ varies in \mathbb{R}^2 . Indeed, we have seen that any point in Ω may be written as $\frac{1}{3}(z_1 + z_2 + z_3)$, where $|z_i| = 1$ and $z_1 z_2 z_3 = 1$, which corresponds to points $\Phi_\theta(1, 1, 1)$.

Since $|z_i|^2 \leq 1$ for any point in Ω , then, $\forall z \in \Omega, |\pi(z)| \leq 1$, and if $\pi(z) = 1$, then $z_1 = z_2 = z_3 = 1$. Therefore, $\pi(\Phi_\theta(z_1, z_2, z_3)) = \Phi_\theta(1, 1, 1)$. This shows that the conditional law of $\pi(\Phi_\theta(z))$ when $\pi(z) = 1$ is a Dirac mass at $\Phi_\theta(1, 1, 1)$.

Indeed, we may as well chose for x_0 the image of the points (j, j, j) or $(\bar{j}, \bar{j}, \bar{j})$, and those three points correspond to the cusps of the deltoid curve. Define $Z(\theta) = \pi(\Phi_\theta(1, 1, 1))$. Then, $\pi(\Phi_\theta(j, j, j)) = jZ(\theta)$, and $\pi(\Phi_\theta(\bar{j}, \bar{j}, \bar{j})) = \bar{j}Z(\theta)$. Moreover, the conditional law of $z = \frac{1}{3}(z_1 + z_2 + z_3)$ knowing that $\pi(z) = 1, j, \bar{j}$ is a Dirac mass at $Z(\theta), jZ(\theta)$ and $\bar{j}Z(\theta)$ respectively. The point $Z(\theta)$ is in the interior of the deltoid domain as soon as $\theta_1 \neq \theta_2 \pmod{2\pi}$ and $2\theta_1 \neq -\theta_2 \pmod{2\pi}$. Observe also that $Z(-\theta) = \bar{Z}(\theta)$.

To apply the method described in Section 2, the only problem is that the eigenspaces associated with $L^{(\lambda)}$ have dimension 2, when $n \neq k$. They are invariant under the symmetry $Z \mapsto \bar{Z}$, and the model $\hat{L}^{(\lambda)}$ also shares the symmetry $\mathcal{S} : (z_1, z_2, z_3) \mapsto (\bar{z}_1, \bar{z}_2, \bar{z}_3)$. So instead of looking at eigenvectors of $L^{(\lambda)}$ alone, one may look at eigenvectors of $L^{(\lambda)}$ which are symmetric, or antisymmetric, through the

transformation $Z \mapsto \bar{Z}$, which leads us to consider the basis $(P_{n,k}, Q_{n,k})$ described in section 4.

The transformation $z \mapsto \Phi_\theta(z)$ is not invariant under the symmetry \mathcal{S} . We have $\overline{\Phi_\theta(z)} = \Phi_{-\theta}(\bar{z})$. Observe also that the operators $\Phi_\theta(f)(z) = f(\Phi_\theta(z))$ satisfy $\langle \Phi_\theta f, g \rangle = \langle f, \Phi_{-\theta} g \rangle$, which comes from the fact that the measure is invariant under the symmetry \mathcal{S} .

From the general scheme described in Section 2, one sees that, writing the Markov operator $K_\theta(f) = \mathbb{E}(f(\pi(\Phi_\theta z)) / \pi(z) = Z)$, that both $K_\theta(P_{n,k})$ and $K_\theta(Q_{n,k})$ belong to the eigenspace associated to the eigenvalue $\lambda_{n,k}$, defined in Proposition 4.1. We have

Proposition 6.1. *With $K_\theta(f)(Z) = \mathbb{E}(f(\pi(\Phi_\theta z)) / \pi(z) = Z)$, one has, for any $\lambda \geq 5/2$,*

$$\begin{pmatrix} K_\theta(P_{n,k}) \\ K_\theta(Q_{n,k}) \end{pmatrix} = \begin{pmatrix} \alpha_{n,k}(\theta) & \beta_{n,k}(\theta) \\ \gamma_{n,k}(\theta) & \delta_{n,k}(\theta) \end{pmatrix} \begin{pmatrix} P_{n,k} \\ Q_{n,k} \end{pmatrix},$$

with

$$\begin{cases} \alpha_{n,k}(\theta) = \frac{P_{n,k}(Z(\theta))}{P_{n,k}(1)}; \\ \gamma_{n,k}(\theta) = -\beta_{n,k}(\theta) = \frac{Q_{n,k}(Z(\theta))}{P_{n,k}(1)}. \end{cases}$$

Moreover, one has $\alpha_{n,k}(\theta) = \alpha_{n,k}(-\theta)$, $\delta_{n,k}(-\theta) = \delta_{n,k}(\theta)$, and $\beta_{n,k}(\theta) = -\beta_{n,k}(-\theta)$.

Proof. — It is enough to check this property for $\lambda > 5/2$, since everything is continuous and we would have the same property in the limit $\lambda = 5/2$.

Equation $\langle K_\theta(f), g \rangle = \langle f, K_{-\theta}(g) \rangle$ proves the parity relations, choosing f and g either $P_{n,k}$ or $Q_{n,k}$.

If we apply $K_\theta(f)(1) = f(Z(\theta))$, and using the fact that $Q_{n,k}(1) = 0$, we get $P_{n,k}(1)\alpha_{n,k}(\theta) = P_{n,k}(Z(\theta))$, and $P_{n,k}(1)\gamma_{n,k}(\theta) = Q_{n,k}(Z(\theta))$. This shows that $P_{n,k}(1)$ may not vanish (since it would imply that $P_{n,k} = 0$ everywhere), and leads to the representation formula for $\alpha_{n,k}$ and $\gamma_{n,k}$. The formula for $\beta_{n,k}$ follows by symmetry, since $Z(-\theta) = \bar{Z}(\theta)$ and $Q_{n,k}(\bar{Z}) = -Q_{n,k}(Z)$. ■

Remark 6.2. *The value of $\delta_{n,k}(\theta)$ however is more difficult to obtain. One may apply the identity $K_\theta(f)(j) = f(jZ(\theta))$ and formulae (4.12) whenever $n - k \not\equiv 0 \pmod{3}$, to get*

$$\delta_{n,k}(\theta) = \frac{\cos(2(n-k)\pi/3) P_{n,k}(Z(\theta))}{\sin(2(n-k)\pi/3) P_{n,k}(1)}.$$

This method, however, does not provide any information on $\delta_{n,k}(\theta)$ when $n - k \equiv 0 \pmod{3}$.

The point $x_0 = 1$ corresponds to one of the 3 cusps of the deltoid curve, and is the image of one of the points of $\mathcal{S}^1 \times \mathcal{S}^1 \times \mathcal{S}^1$ where $z_1 = z_2 = z_3$ and $z_1 z_2 z_3 = 1$. There are 3 such points, corresponding to the three cusps of the deltoid, and we could have similarly proved the hypergroup property for any of those points, but for another polynomial basis. The choice of the point $x_0 = 1$ corresponds to the choice

of the basis $(P_{n,k}, Q_{n,k})$ such that, under the symmetry $\mathcal{S} : Z \mapsto \bar{Z}$, $SP_{n,k} = P_{n,k}$ and $SQ_{n,k} = -Q_{n,k}$.

We may as well consider symmetries which leave the two other cusps invariant, and this provides new bases for the eigenspace in which the operator K_θ has a similar expression. A polynomial R is symmetric with respect to the symmetry through the j axis if $R(\bar{j}\bar{Z}) = R(Z)$. Then, thanks to equation (4.11), a basis $(R_{n,k}, S_{n,k})$ of the eigenspace associated with $\lambda_{n,k}$ for which the first element is symmetric and the second antisymmetric under the symmetry around the j axis would be $R_{n,k} + iS_{n,k} = j^{n-k}(P_{n,k} + iQ_{n,k})$. One also have the hypergroup property for the family $(R_{n,k}, S_{n,k})$ and similarly for the family $(U_{n,k}, V_{n,k})$ corresponding to the point \bar{j} . This leads to other representations of the Markov operator K_θ .

As usual, the operators K_θ lead to a full representation of Markov kernels.

Theorem 6.3. *Let K be a symmetric Markov operator, bounded in $\mathcal{L}^2(\mu^{(\lambda)})$, with $\lambda \geq 5/2$. Assume that K commutes with $L^{(\lambda)}$. Then, with the notations of Proposition 4.1, for any $n, k \in \mathbb{N}^2$, it satisfies*

$$(6.22) \quad \begin{pmatrix} K(P_{n,k}) \\ K(Q_{n,k}) \end{pmatrix} = \begin{pmatrix} a_{n,k} & b_{n,k} \\ b_{n,k} & c_{n,k} \end{pmatrix} \begin{pmatrix} P_{n,k} \\ Q_{n,k} \end{pmatrix}$$

and there exists a probability measure ν_1 on the deltoid domain Ω such that

$$(6.23) \quad a_{n,k} = \int \frac{P_{n,k}(z)}{P_{n,k}(1)} \nu(dz), \quad b_{n,k} = \int \frac{Q_{n,k}(z)}{P_{n,k}(1)} \nu(dz).$$

Proof. —

We follow the lines of the proof described in Section 2. Equation (6.22) is immediate from the fact that K commute with $L^{(\lambda)}$ and the description of the eigenspaces of $L^{(\lambda)}$ given in Proposition 4.1. Extending the operator K to act on probability measures, we choose $\nu_1(dz) = K(\delta_1)$. If P_t denotes the heat kernel associated with $L^{(\lambda)}$, let $\nu_t = K(P_t(\delta_1))$. Following [3], we know that $P_t(\delta_1)$ has a bounded density with respect to $\mu^{(\lambda)}$, which may be written as $\sum_{n,k} e^{-\lambda_{n,k}t} P_{n,k}(z) P_{n,k}(1)$, where this simplified form comes from the fact that $Q_{n,k}(1) = 0$. Then, the density h_t of ν_t with respect to $\mu^{(\lambda)}$ may be written as

$$h_t = \sum_{n,k} e^{-\lambda_{n,k}t} P_{n,k}(1) (a_{n,k} P_{n,k}(Z) + b_{n,k} Q_{n,k}(Z)),$$

and we see that

$$\int \frac{P_{n,k}(Z)}{P_{n,k}(1)} d\nu_t = e^{-\lambda_{n,k}t} a_{n,k}, \quad \int \frac{Q_{n,k}(Z)}{P_{n,k}(1)} d\nu_t = e^{-\lambda_{n,k}t} b_{n,k}.$$

Since ν_t converges to $\nu = K(\delta_1)$ when $t \rightarrow 0$, we get the result in the limit. \blacksquare

The previous representation relies in an essential way on the fact that $Q_{n,k}(1) = 0$. This choice comes from the symmetry properties of the operator under $Z \mapsto \bar{Z}$, that is the symmetry around the real axis.

In order to get informations about the coefficients $c_{n,k}$, one may use the invariance of the model under $Z \mapsto jZ$ and $Z \mapsto \bar{j}Z$ whenever $n - k \not\equiv 0 \pmod{3}$. One may use similarly the invariance through multiplication by j and \bar{j} and symmetries with the corresponding axis. This means that a similar presentation is valid in the two other bases $(R_{n,k}, S_{n,k})$ and $(U_{n,k}, V_{n,k})$ (using the symmetries leaving j and \bar{j} invariant respectively).

In this new basis, the matrix of the operator is unchanged when $n - k \equiv 0 \pmod{3}$, while the matrix of the operator K in this new basis becomes,

$$\frac{1}{4} \begin{pmatrix} a_{n,k} + 2\epsilon\sqrt{3}b_{n,k} + 3c_{n,k} & -\epsilon\sqrt{3}a_{n,k} - 2b_{n,k} + \epsilon\sqrt{3}c_{n,k} \\ -\epsilon\sqrt{3}a_{n,k} - 2b_{n,k} + \epsilon\sqrt{3}c_{n,k} & 3a_{n,k} - 2\epsilon\sqrt{3}b_{n,k} + c_{n,k} \end{pmatrix}$$

where $\epsilon = 1$ when $n - k \equiv 1 \pmod{3}$ and $\epsilon = -1$ when $n - k \equiv 2 \pmod{3}$.

If we observe that $(R_{n,k} + iS_{n,k})(jZ) = (P_{n,k} + iQ_{n,k})(Z)$, so that $R_{n,k}(j) = U_{n,k}(\bar{j}) = P_{n,k}(1)$, we get a new representation, with the measure $\nu_1 = K(\delta_j)$, when $n - k \not\equiv 0 \pmod{3}$

$$\frac{1}{4}(a_{n,k} + 2\epsilon\sqrt{3}b_{n,k} + 3c_{n,k}) = \int \frac{R_{n,k}(z)}{P_{n,k}(1)} \nu_1(dz) = \int \frac{U_{n,k}(z)}{P_{n,k}(1)} \nu_2(dz)$$

and

$$\frac{1}{4}(-\epsilon\sqrt{3}a_{n,k} - 2b_{n,k} + \epsilon\sqrt{3}c_{n,k}) = \int \frac{S_{n,k}(z)}{P_{n,k}(1)} \nu_1(dz) = \int \frac{V_{n,k}(z)}{P_{n,k}(1)} \nu_2(dz).$$

This may be rewritten as

$$\begin{cases} -\frac{a_{n,k}}{2} - \epsilon\frac{\sqrt{3}}{2}b_{n,k} = \int \frac{P_{n,k}(z)}{P_{n,k}(1)} d\nu_1(z) \\ -\frac{b_{n,k}}{2} - \epsilon\frac{\sqrt{3}}{2}c_{n,k} = \int \frac{Q_{n,k}(z)}{P_{n,k}(1)} d\nu_1(z). \end{cases}$$

Changing j into \bar{j} amounts to change ϵ into $-\epsilon$ in the previous formulas, with the measure $\nu_2 = K(\delta_{\bar{j}})$.

This in turn provides a representation of $c_{n,k}$ when $n - k \not\equiv k \pmod{3}$, of the form

$$c_{n,k} = \epsilon \int \frac{Q_{n,k}(z)}{P_{n,k}(1)} \left(\frac{2}{\sqrt{3}} d\nu_1(z) - \frac{1}{\sqrt{3}} d\nu(z) \right).$$

Unfortunately, this does not carry any information about $c_{n,k}$ when $n \equiv k \pmod{3}$.

Remark 6.4. *As we already observed, for $\lambda = 5/2$, the 6 dimensional model is in fact carried by the algebraic hypersurface $\{P_1 = 0\}$, and in fact is a 5-dimensional model. However, for other values of λ (with the sole exception of $\lambda = 1$), we do not know if the property is true. It would be interesting to construct lower dimensional models for these values, but this seems quite hard.*

Remark 6.5. *It would be interesting to have an explicit expression for the kernel $K_\theta(x, dy)$, in order to have explicit representation for the product formula (2.2). Unfortunately, the law of $(Z, R_\theta(Z))$ is already apparently quite out of reach through simple formulas.*

7 Projection of the deltoid model and the G_2 -root system

As we saw in Section 6, the fact that the eigenspaces for $L^{(\lambda)}$ are two dimensional introduce extra complexity in the representation of the eigenvalues of the Markov operators which commute with $L^{(\lambda)}$. This works much better if we concentrate on functions which are symmetric in (Z, \bar{Z}) , which correspond to the symmetric polynomials $P_{n,k}(Z, \bar{Z})$. It turns out that these symmetric polynomials are again orthogonal polynomials corresponding to another bounded set $\Omega_2 \subset \mathbb{R}^2$, bounded by a cuspidal cubic and a parabola, tangent to each other at the second order. Going back to the triangle model, remember that the deltoid model in the case $\lambda = 1$ corresponds to functions which are invariant under the symmetries of a triangular lattice, corresponding to the root system A_2 . Adding this new invariance under $Z \mapsto \bar{Z}$ amounts then to add new symmetries, namely with respect to the medians of the triangles, corresponding to the root system G_2 . Let us describe this new polynomial system. Setting $s = Z + \bar{Z}$ and $p = Z\bar{Z}$, formulae (4.10) give

$$(7.24) \quad \begin{cases} \Gamma(s, s) = p - s^2 + s + 1, \\ \Gamma(s, p) = s^2 - 2p - \frac{3}{2}sp + \frac{s}{2}, \\ \Gamma(p, p) = s^3 - 3p^2 - 3sp + p \\ L^{(\lambda)}(s) = -\lambda s, \quad L^{(\lambda)}(p) = 1 - (2\lambda + 1)p \end{cases}$$

Let us call $\tilde{L}^{(\lambda)}$ this operator acting on functions (indeed polynomials) in the variables (s, p) .

From equation (7.24), it is clear that the operator $\tilde{L}^{(\lambda)}$ preserves the set of polynomials in the variables (s, p) (this just translates the invariance of $L^{(\lambda)}$ under $Z \mapsto \bar{Z}$). However, because of the term p^3 in the coefficient $\Gamma(p, p)$, it does not preserve the degree of the polynomial. But things work better if we decide that the degree of $s^r p^t$ is $s + 2t$. Then, with this new notion of degree, a polynomial $Q(s, p)$ of degree k is transformed under $\tilde{L}^{(\lambda)}$ into a polynomial of degree k . One may therefore find an orthonormal basis for $\tilde{L}^{(\lambda)}$ as polynomials in the variables (s, p) . In fact, it is nothing else than the symmetric polynomials $P_{n,k}(Z, \bar{Z})$ expressed as polynomials in (s, p) . Since now, the eigenspaces are one dimensional, one may play the same operation with the operator $\hat{L}^{(\lambda)}$, but now with the projection $\Psi : Z \mapsto \frac{1}{2}(Z + \bar{Z}) = s, Z\bar{Z} = p$. We shall obtain the true hypergroup property for this polynomial family, through the projection $\Omega_1 \mapsto R_\theta(z) + \overline{R_\theta(z)} = R_\theta(z) + R_{-\theta}(\bar{z})$.

The image operator may be diagonalized in a family of orthogonal polynomials (in the variables (s, p)) for the image measure. This model does not appear in the [4] classification, since in this case the orthogonal polynomials must be ranked according to a degree which is $2\deg(p) + \deg(s)$, whereas in [4], the polynomials are ranked according to their usual degree. The boundary of Ω_2 is indeed the set where the determinant of the metric vanishes. This determinant may be written as

$$\frac{1}{4}(s^2 - 4p)(3s^2 + 12sp + 6p - 4s^3 - 1),$$

and the boundary of Ω_2 is a degree 5 algebraic curve. The curves $s^2 - 4p = 0$ and $3p^2 + 12sp + 6p - 4s^3 - 1 = 0$ correspond respectively to the images under Ψ of the line $Z = \bar{Z}$ (the real axis), and of the boundary of Ω (the deltoid curve). It turns out that this last curve is a cuspidal cubic. Setting

$$s = x - 1, p = y - 2x + 1,$$

(which corresponds to some affine change of coordinates), they are transformed in

$$4x^3 - 3y^2 = 0, \quad x^2 + 6x - 4y - 3 = 0.$$

These two curves (cubic and parabola) cross in the point $(1/3, -2/9)$, and are tangent to the second order at the point $(3, 6)$ (in the variables (x, y)). The cuspidal point in the cubic $(0, 0)$ is the image of j and \bar{j} in the deltoid (two of the cuspidal points of the deltoid), while the point $(3, 6)$ is the image of the third cusp of the deltoid, that is the point 1, which is also on the line $Z = \bar{Z}$.

Moreover, with $Q_1 = s^2 - 4p$ and $Q_2 = 3p^2 + 12sp + 6p - 4s^3 - 1$, both Q_1 and Q_2 satisfy an equation similar to equation (5.18), and more precisely

$$\begin{cases} \Gamma(\log Q_1, s) = -2s - 2, & \Gamma(\log Q_1, p) = -3p - 2s + 1, \\ \Gamma(\log Q_2, s) = -3s, & \Gamma(\log Q_2, p) = -6p \end{cases}$$

This shows that the operator defined through equations (7.24) has reversible measure $C_\lambda Q_1^{-1/2} Q_2^{(2\lambda-5)/6} dsdp$, which is of course the image of the measure $\mu^{(\lambda)}$ through the projection Ψ . This operator therefore satisfies the usual hypergroup property, with reference point the image of 1, which is $(2, 1)$, that is the point where the cuspidal cubic and the parabola are bi-tangent to each other.

But from the general presentation of [4], there is now a two-parameter family of measures, namely $\mu_{\alpha, \beta}(ds, dp) = C_{\alpha, \beta} Q_1^{\alpha_1} Q_2^{\alpha_2} dsdp$ on this set Ω_2 , for which there exist a family of orthogonal polynomials which are eigenvectors of a diffusion operator.

The conditions under which those measures are finite are $\alpha_1 > -1$, $\alpha_2 > -5/6$, and $\alpha_1 + \alpha_2 > -4/3$. The condition $\alpha_2 > -5/6$ and $\alpha_1 + \alpha_2 > -4/3$ correspond respectively to the integrability conditions around the cusp of the cubic and the double tangent point. For the double tangent point, to check the integrability condition, one may reduce, up to an affine transformation of the plane, to check the integration condition for

$$\int_0^1 \int_0^2 \mathbf{1}_{y^2 < x < y^2 + cy^3} (x - y^2)_1^\alpha (y^2 + cy^3 - x)_2^\alpha dx dy.$$

After a few change of variables, this reduces to check the integrability condition for $\int_0^1 y^{3(\alpha_1 + \alpha_2 + 1)} dy$. The cusp is easier to deal with and quite immediate.

For these measures μ_{α_1, α_2} , we have an associated operator L_{α_1, α_2} (sharing the same Γ operator, and defined through formula (3.3)) and associated orthogonal polynomials which are eigenvectors of L_{α_1, α_2} .

We already mentioned that this model reflects in fact the symmetries of the root system G_2 . In this context, the Weyl group acting on the roots has two orbits,

corresponding to the two irreducible factors of the boundary of Ω_1 , and to the two parameters in the choice of the measure. When $\alpha_1 = -1/2$, the model is the direct image of the deltoid model through the projection Ψ , and the eigenvectors of the associated operators are just the rewriting of $P_{n,k}$ as polynomials in the variables $s = Z + \bar{Z}$ and $p = Z\bar{Z}$. The hypergroup property for this model is the direct consequence of the previous Section 6, but may also be reproved directly using the 6 dimensional model of Section 5.

The case where $\alpha_1 \neq -1/2$ remains open. Let us show however how to deal with a different value of α_1 , whenever $\alpha_2 = -1/2$. This relies on a very specific property of this G_2 polynomial model. We first observe that the domain Ω_2 is stable under the transformation

$$\Psi_1 : (s, p) \mapsto (S = 3p - 1, P = 1 + s^3 - 3ps - 6p).$$

Under this transformation, the parabolic part of the boundary is mapped onto the cubic one, and conversely. The transformation Ψ_1 comes in fact from the invariance $Z \mapsto jZ$ in the deltoid model, where instead of looking at functions of $(Z\bar{Z}, Z + \bar{Z})$, we looked at functions of $(Z\bar{Z}, Z^3 + \bar{Z}^3)$, under a slight change of coordinates such that the domain Ω_2 is invariant under the transformation.

Moreover, one may check that, when $\alpha_1 = -1/2$, the image of the operator $L_{-1/2, \alpha_2}$ is $\frac{1}{3}L_{\alpha_2, -1/2}$. It may be easily checked looking at $\Gamma(S, S), \Gamma(S, P), \Gamma(P, P), L_{-1/2, \alpha_2}(S)$ and $L_{-1/2, \alpha_2}(P)$. There is nothing similar when $\alpha_1 \neq -1/2$. The transformation Ψ_1 is not a diffeomorphism, and the operators L_{α_1, α_2} in general do not have images under Ψ_1 when $\alpha_1 \neq -1/2$. From the previous scheme, we may now conclude to the hypergroup property for the orthogonal polynomials associated with $L_{\alpha, -1/2}$ when $\alpha \geq 0$.

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