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Extreme Value Theory for Piecewise Contracting Maps with Randomly Applied Stochastic Perturbations

Davide Faranda * Jorge Milhazes Freitas † Pierre Guiraud ‡ Sandro Vaienti § January 14, 2015

Abstract

We consider globally invertible and piecewise contracting maps in higher dimensions and we perturb them with a particular kind of noise introduced by Lasota and Mackey. We got random transformations which are given by a stationary process: in this framework we develop an extreme value theory for a few classes of observables and we show how to get the (usual) limiting distributions together with an extremal index depending on the strength of the noise.

1 Introduction

In a previous note [7] we announced a few results on the application of Extreme Value Theory (EVT) to onedimensional piecewise contracting maps (PCM) whenever the latter are perturbed with a particular kind of noise. This noise was introduced in the book by Lasota and Mackey [16] and it was called randomly applied stochastic perturbation (RASP). It turns out that it is particularly useful to study PCM especially because it allows very efficient representations of the transfer operator (Perron-Frobenius) and of the stationary measures. We now generalize to higher dimensional PCM our previous results by proving the existence of extreme value laws (EVL) for some kinds of observables associated to rare visits in the neighborhood of a given point. As a matter of fact, we first consider globally invertible maps defined on compact subsets of \mathbb{R}^n and which are diffeomorphisms but on some singular (Borel) sets of zero Lebesgue measure, so they are not necessarily contractions. The RASP belongs to the class of random transformations (see next section), which naturally define a Markov chain with associated a unique absolutely continuous stationary measure (see Th. 10.4.2 in [16]). By taking the direct product of this stationary measure with the probability distribution of the concatenated maps, we get the (annealed) probability \mathbb{P} which will govern the statistical properties of our system. In particular the stochastic process given by the observation along a random orbit will be stationary according to P. We will first prove a general result about decay of correlations which will not require strong assumptions on the functional spaces of the observables (Proposition 2.1). We will therefore use this result to recover some sort of asymptotic independence in order to prove convergence of the distribution of the maxima toward the Gumbel's law. This will be accomplished in two other steps: first we will get an explicit expression of the affine coefficients defining the sequence of levels u_n needed to rescale the distribution of the maximum to avoid degenerate limits (Propositions 4.1,4.2,4.3); then we will control the short returns in Proposition 4.5. Whenever we will consider rare visits in the neighborhood of a periodic cycle, we will prove the existence of an extremal index less than 1, and in particular equal to the magnitude of the noise.

We would like to stress that we already investigated the EVT for random transformations obtained by perturbing with additive noise; we applied it to uniformly expanding maps [1] and even to rotations [8].

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In both cases we got the existence of an EVL around any point and with extremal index equal to 1. It seems that RASP is much more suitable for piecewise contractions and does not smoothen out completely the periodicity features of the unperturbed map, as this is reflected in the persistency of an extremal index less than one. A few numerical simulations seem to suggest however the existence of an EVL even when we perturb PCM with additive noise [7].

We now briefly recall the main statements of the EVT. Let us suppose that $(Y_n)_{n\in\mathbb{N}}$ is a sequence of real-valued random variables defined on the probability space (Ψ, \mathbb{P}) . We will be interested in the distribution of the maximum $M_n := \max\{Y_0, Y_1, \dots, Y_{n-1}\}$ when $n \to \infty$. It is well known that the limiting distribution is degenerate unless one proceed to a suitable re-scaling of the levels of exceedances. The precise formulation is the following: we have an Extreme Value Law for $(M_n)_{n\in\mathbb{N}}$ if there is a non-degenerate distribution function $H: \mathbb{R} \to [0,1]$ with H(0) = 0 and, for every $\tau > 0$, there exists a sequence of levels $(u_n(\tau))_{n\in\mathbb{N}}$ such that

$$\lim_{n \to \infty} n \mathbb{P}(Y_0 > u_n) \to \tau, \tag{1}$$

and for which the following holds:

$$\lim_{n\to\infty} \mathbb{P}(M_n \le u_n) \to 1 - H(\tau).$$

The motivation for using a normalizing sequence $(u_n)_{n\in\mathbb{N}}$ satisfying (1) comes from the case when $(Y_n)_{n\in\mathbb{N}}$ are independent and identically distributed. In this i.i.d. setting, it is clear that $\mathbb{P}(M_n \leq u) = (F(u))^n$, being F(u) the cumulative distribution function for the variable u. Hence, condition (1) implies that

$$\mathbb{P}(M_n \leq u_n) = (1 - \mathbb{P}(Y_0 > u_n))^n \sim \left(1 - \frac{\tau}{n}\right)^n \to e^{-\tau},$$

as $n \to \infty$. Note that in this case $H(\tau) = 1 - e^{-\tau}$ is the standard exponential distribution function. Let us now choose the sequence $u_n = u_n(y)$ as the one parameter family $u_n = y/a_n + b_n$, where $y \in \mathbb{R}$ and $a_n > 0$, for all $n \in \mathbb{N}$. Whenever the variables Y_i are i.i.d. and for some constants $a_n > 0$, $b_n \in \mathbb{R}$, we have $\mathbb{P}(a_n(M_n - b_n) \leq y) \to G(y)$, where the convergence occurs at continuity points of G, and G is non-degenerate, then G_n will converge to one of the three EVLs: Gumbel, Fréchet or Weibull. The law obtained depends on the the common distribution of the random variables.

When Y_0, Y_1, Y_2, \ldots are not independent, the standard exponential law still applies under some conditions on the dependence structure. These conditions will be stated in detail later and they are usually designated with D' and D_2 ; when they hold for $(Y_n)_{n\in\mathbb{N}}$ then there exists an extreme value law for M_n and $H(\tau) = 1 - e^{-\tau}$, see Theorem 1 in [11]. We want to stress that these two conditions alone do not imply the existence of an extreme value law; they require, even to be checked, that the limit (1) holds. It turns out that for the kind of observables we are going to introduce, and which are related to the local properties of the invariant measure, the limit (1) is difficult to prove when the invariant measure is not absolutely continuous, since one needs the exact asymptotic behavior of that measure on small balls. Instead it turns out that whenever the systems is randomly perturbed and the invariant measure is absolutely continuous, the existence of the limit (1) can be insured with general arguments and in the majority of the cases (see below for a precise meaning), it can be expressed by a closed formula in terms of the strength of the noise, Proposition 4.2 and 4.3 below. Moreover that formula could be used in a reversed way: since the sequence $(u_n)_{n\in\mathbb{N}}$ is now uniquely determined for any n, a numerical sampling for $(u_n)_{n\in\mathbb{N}}$ which provides convergence to the extreme value law, will bring information on the local geometrical properties of the stationary measure: this approach was successfully used, for instance, in [9, 10].

2 Random dynamical systems: RASP perturbation

Let us consider a sequence of i.i.d. random variables $(W_k)_{k\in\mathbb{N}}$ with values $(\omega_k)_{k\in\mathbb{N}}$ in a space Ω_{ε} and with common probability distribution θ_{ε} . Let $X\subset\mathbb{R}^D$ be a compact set equipped with the Lebesgue measure m defined on the Borel σ -algebra, and $(f_{\omega})_{\omega\in\Omega_{\varepsilon}}$ a family of measurable transformations such that $f_{\omega}:X\to X$

for all $\omega \in \Omega_{\varepsilon}^{-1}$. Given a point $x \in X$ and a realization $\underline{\omega} = (\omega_1, \omega_2, \dots) \in \Omega_{\varepsilon}^{\mathbb{N}}$ of the stochastic process $(W_k)_{k \in \mathbb{N}}$, we define the random orbit of x as the sequence $(f_{\omega}^n(x))_{n \in \mathbb{N}}$, where

$$f_{\underline{\omega}}^{0}(x) = x$$
 and $f_{\underline{\omega}}^{n}(x) = f_{\omega_{n}} \circ f_{\omega_{n-1}} \circ \cdots \circ f_{\omega_{1}}(x)$ $\forall n \ge 1$.

The transformations f_{ω} will be considered as stochastic perturbations of a deterministic map f, in the sense that they will be taken in a suitable neighborhood of f whose *size* will be determined by the value of ε , see below. We could therefore define a Markov process on X with transition function

$$L_{\varepsilon}(x,A) = \int_{\Omega_{\varepsilon}} \mathbf{1}_{A}(f_{\omega}(x)) d\theta_{\varepsilon}(\omega), \qquad (2)$$

where $A \in X$ is a measurable set, $x \in X$ and $\mathbf{1}_A$ is the indicator function of the set A. A probability measures μ_{ε} is called *stationary* if for any measurable set A we have:

$$\mu_{\varepsilon}(A) = \int_{X} L_{\varepsilon}(x, A) d\mu_{\varepsilon}(x).$$

We call it an absolutely continuous stationary measure, if it has a density with respect to the Lebesgue measure.

We are now ready to introduce the randomly applied stochastic perturbations. We first have a Borel measurable map f acting on the compact set X; then we operate an aleatory reset of the state of the system at each failure of a Bernoulli random variable: if $(x_n)_{n\in\mathbb{N}}$ denotes the successive states of such a random dynamical systems, then at each time $n\in\mathbb{N}$ we have $x_{n+1}=f(x_n)$ with probability $(1-\epsilon)$ and $x_{n+1}=\xi_n$ with probability ϵ , where ξ_n is the realization of a random variable with value in X. This kind of perturbation corresponds to the family $(f_{\omega})_{\omega\in\Omega_{\varepsilon}}$ of random transformations defined by

$$f_{\omega}(x) = \eta f(x) + (1 - \eta)\xi \qquad \forall x \in X, \tag{3}$$

where $\omega = (\eta, \xi)$ is a random vector with value in $\Omega_{\varepsilon} = \{0, 1\} \times X$. The two components η and ξ of ω are independent and η is a Bernoulli variable with the probability of being 0 equal to ε , while ξ is a random variable that we will suppose Lebesgue-uniformly distributed on X^2 . The joint distribution θ_{ε} of these two components is the product of the Bernoulli measure with weights $(\varepsilon, 1 - \varepsilon)$ and the uniform measure on X.

In order to obtain the stationary measure μ_{ε} , let us introduce the random Koopman operator U_{ε} : $L^{\infty} \to L^{\infty}$ defined for all $\phi \in L^{\infty 3}$ by

$$U_{\varepsilon}\phi(x) := \int \phi(f_{\omega}(x))d\theta_{\varepsilon}.$$

Now, if we take two observables $\phi \in L^{\infty}$ and $\psi \in L^{1}$, it easy to check that

$$\int U_{\varepsilon}\phi(x)\psi(x)dx = \int \int \phi(f_{\omega}(x))\psi(x)d\theta_{\varepsilon}dx =$$

$$\varepsilon \int \int \phi(x)\psi(y)dxdy + (1-\varepsilon)\int \phi(f(x))\psi(x)dx = \int \phi(x)P_{\varepsilon}\psi(x)dx,$$

where P_{ε} is the adjoint operator of U_{ε} , that is the random transfer operator. If we denote P the transfer operator associated to f and $\overline{\psi} = \int \psi(y)dy$, then we have

$$P_{\varepsilon}\psi(x) = (1 - \varepsilon)P\psi(x) + \varepsilon\overline{\psi}.$$
 (4)

¹In the following when we will refer to a dynamical system (X, f, μ) we will mean that f is defined on X and preserves the Borel probability measure μ ; if we will write (X, f), this will simply correspond to the action of f on X.

²In the book of Lasota Mackay [16] the variable ξ is more generally distributed with a density with respect to the Lebesgue measure; this will not change the results of our paper.

³From now on L^1 and L^{∞} will be referred to the Lebesgue measure m and the integral with respect to the latter will be denote as $\int (\cdot) dx$.

The stationary measure μ_{ε} verifies $\int \phi(x)d\mu_{\varepsilon} = \int U_{\varepsilon}\phi(x)d\mu_{\varepsilon}$ and in our case is given by $\mu_{\varepsilon} = h_{\varepsilon}m$ where $h_{\varepsilon} \in L^1$ is a density such that $h_{\varepsilon} = P_{\varepsilon}h_{\varepsilon}$. Such a density exists and is given by [16]:

$$h_{\varepsilon} = \varepsilon \sum_{k=0}^{\infty} (1 - \varepsilon)^k P^k \mathbf{1}. \tag{5}$$

In the random setting we are developing the decay of correlations will be of annealed type and it will be formulated in terms of the iterates of the random Koopman operator, in particular we have:

Proposition 2.1. If $\phi \in L^{\infty}$ and $\psi \in L^1 \cap L^{\infty}$, then

$$\left| \int \int \phi(f_{\underline{\omega}}^{n}(x))\psi(x)d\mu_{\varepsilon}d\theta_{\varepsilon}^{\mathbb{N}} - \int \phi(x)d\mu_{\varepsilon} \int \psi(x)d\mu_{\varepsilon} \right| \leq 2(1-\varepsilon)^{n}||\phi||_{L^{\infty}}||\psi h_{\varepsilon}||_{L^{1}}.$$
 (6)

Proof. To use the duality of the Koopman operator with the transfer operator P, let us introduce the following quantity:

$$Cor_{m}(\phi, \psi, n) := \left| \int U_{\varepsilon}^{n}(\phi(x))\psi(x)dx - \int \phi(x)d\mu_{\varepsilon} \int \psi(x)dx \right| = \left| \int \phi(x)P_{\varepsilon}^{n}\psi(x)dx - \int \phi(x)d\mu_{\varepsilon} \int \psi(x)dx \right|.$$

$$(7)$$

By recurrence, it is easy to obtain the following formula for the iterates of P_{ε} :

$$P_{\varepsilon}^{n}\psi = (1 - \varepsilon)^{n}P^{n}\psi + \varepsilon\overline{\psi}\sum_{k=0}^{n-1}(1 - \varepsilon)^{k}P^{k}\mathbf{1}.$$
 (8)

By replacing this expression into (7) and writing there $d\mu_{\varepsilon} = h_{\varepsilon} dx$, with h_{ε} given by (5), we get

$$Cor_{m}(\phi, \psi, n) = \left| (1 - \varepsilon)^{n} \int \phi(x) P^{n} \psi(x) dx - \varepsilon \overline{\psi} \int \phi(x) \sum_{k=n}^{\infty} (1 - \varepsilon)^{k} P^{k} \mathbf{1}(x) dx \right|$$

$$\leqslant (1 - \varepsilon)^{n} ||\phi||_{L^{\infty}} ||P^{n} \psi||_{L^{1}} + \varepsilon |\overline{\psi}|||\phi||_{L^{\infty}} ||\sum_{k=n}^{\infty} (1 - \varepsilon)^{k} P^{k} \mathbf{1}||_{L^{1}}$$

$$\leqslant 2(1 - \varepsilon)^{n} ||\phi||_{L^{\infty}} ||\psi||_{L^{1}},$$

by Hölder inequality and the fact that the infinite sum is of positive terms. Therefore, if $\psi \in L^1 \cap L^\infty$, then the correlations (6) are equal to

$$Cor_m(\phi, \psi h_{\varepsilon}, n) \leq 2(1 - \varepsilon)^n ||\phi||_{L^{\infty}} ||\psi h_{\varepsilon}||_{L^1}.$$

Remark 2.2. We will argue later on that in the EVT the preceding decay of correlations result plays a fundamental role whenever ψ is a specific characteristic function: in this case it will be sufficient to have the density h_{ε} in L^1 . For observables ψ which are solely in L^1 we will need a density in L^{∞} . Note that the exponential decay of the correlations does not depend on the map f, which is a remarkable property of RASP.

We are now ready to introduce the extreme value statistics; at this regard we define a stochastic process $(Y_n)_{n\in\mathbb{N}},\ n\in\mathbb{N}$, by composing a given observable $\phi:X\to\mathbb{R}$ as: $Y_n=\phi\circ f_\omega^n$, with

$$f_{\omega_n}(x) = \eta_n f(x) + (1 - \eta_n) \xi_n, \quad \forall x \in X,$$

where $\omega_n = (\eta_n, \xi_n) \in \{0, 1\} \times X$ is random variable which distribution is the product of the Bernoulli measure with weights $(\varepsilon, 1 - \varepsilon)$ and the uniform measure on X.

In this case $(Y_n)_{n\in\mathbb{N}}$ will be a stationary process if we consider the probability $\mathbb{P} = \mu_{\varepsilon} \times \theta_{\varepsilon}^{\mathbb{N}}$, which corresponds to the *annealed* situation where we average over the initial condition and over the realization of the noise.

We will in particular choose the observable $\phi(x) = -\log d(x,z)$, where $d(\cdot,\cdot)$ denotes (some) distance on X and z a point of X. As we anticipated in the Introduction, such an observable is related to recurrence in small sets, since the distribution of the maximum of the random observable $\phi \circ f_{\underline{\omega}}^k, k = 0, \ldots, n-1$ up to the level u_n , coincides with the distribution of the first entrance of the random orbit into the ball (in the metric d), $B(z, e^{-u_n})$ centered at z and of radius e^{-u_n}

We now state the two conditions which insure weak dependence of the process and which allow to get the limiting distribution of the maxima.

Condition $[D_2(u_n)]$ We say that $D_2(u_n)$ holds for the sequence Y_0, Y_1, \ldots if for all ℓ, t and n,

$$|\mathbb{P}(Y_0 > u_n \cap \max\{Y_t, \dots, Y_{t+\ell-1} \leqslant u_n\}) - \mathbb{P}(Y_0 > u_n)\mathbb{P}(M_\ell \leqslant u_n)| \leqslant \gamma(n, t),$$

where $\gamma(n,t)$ is decreasing in t for each n, and $n\gamma(n,t_n) \to 0$ when $n \to \infty$ for some sequence $t_n = o(n)$.

Now, let $(k_n)_{n\in\mathbb{N}}$ be a sequence of integers such that

$$k_n \to \infty$$
 and $k_n t_n = o(n)$. (9)

Condition $[D'(u_n)]$ We say that $D'(u_n)$ holds for the sequence $Y_0, Y_1, Y_2, ...$ if there exists a sequence $(k_n)_{n\in\mathbb{N}}$ satisfying (9) and such that

$$\lim_{n \to \infty} n \sum_{j=1}^{\lfloor n/k_n \rfloor} \mathbb{P}(Y_0 > u_n, Y_j > u_n) = 0.$$
 (10)

As we said in the Introduction when these two conditions hold for $(Y_n)_{n\in\mathbb{N}}$ then there exists an extreme value law for M_n and $H(\tau) = 1 - e^{-\tau}$, but provided that the limit (1) is true.

3 Stationary measure for injective piecewise diffeomorphisms with RASP

We now introduce the dynamical system upon which we will perform the randomly applied stochastic perturbation.

Definition 3.1. The map. Let us X be a compact subset of \mathbb{R}^D which is the closure of its interior X = int(X), equipped with the metric d and with the normalized Lebesgue measure m; suppose $\{X_i\}_{i=1}^N$ is a collection of N disjoint open subsets of X such that $X = \bigcup_{i=1}^N \overline{X_i}$ and $m(\Delta) = 0$, where $\Delta := X \setminus \bigcup_{i=1}^N X_i$ is the singular set. We will consider non singular maps $f: X \to X$ which are injective and such that $f|_{X\setminus \Delta}$ is a C^1 -diffeomorphism. We say that f is continuous in a point w of the boundary of X if there is an open ball of radius κ , $B(w,\kappa)$ such that f is continuous on $B(w,\kappa) \cap int(X)$ and it can be extended continuously on $B(w,\kappa)$.

Remark 3.2. In the examples we will treat, the set Δ will be composed by points where f is discontinuous and by the points of the boundary of X. We would like to stress that a point of that boundary is not necessarily a discontinuity point of f.

The following proposition give the expression of the stationary density.

Proposition 3.3. Let $\{\Lambda_k\}_{k\in\mathbb{N}}$ be the sequence of sets defined by $\Lambda_0 := X$ and $\Lambda_{k+1} := f(\Lambda_k \setminus \Delta)$ for all $k \ge 1$. Denote $J_k(x) := \prod_{l=1}^k |\det(f'(f^{-l}(x))|^{-1})$ for all $k \ge 1$ and $J_0(x) := \mathbf{1}(x)$. Then,

$$h_{\varepsilon}(x) = \varepsilon \sum_{k=0}^{\infty} (1 - \varepsilon)^k J_k(x) \mathbf{1}_{\Lambda_k}(x) \quad \forall x \in X.$$
 (11)

Note that $J_k(x)$ is well defined for every $x \in X$ such that $x \notin f^l(\Delta)$ for all $l \in \{1, ..., k\}$. Among other basics properties of the sets Λ_k which will be helpful all along this paper, the following lemma ensures that $J_k(x)$ (and therefore $h_{\varepsilon}(x)$) is well defined for all $x \in \Lambda_k$.

Lemma 3.4. The sequence of sets $\{\Lambda_k\}_{k\in\mathbb{N}}$ has the following properties:

- 1. Λ_k is open for all $k \ge 1$
- 2. $\Lambda_{k+1} \subset \Lambda_k$ for all $k \in \mathbb{N}$,
- 3. $\Lambda_k \cap f^p(\Delta) = \emptyset$, for all $p \in \{1, \dots, k\}$.

Proof. 1) The set Λ_1 is open since $X \setminus \Delta$ is open and f is a diffeomorphism in $X \setminus \Delta$. Using the same arguments, we obtain by induction that Λ_k is open for any $k \geq 1$. 2) We have $\Lambda_2 = f(\Lambda_1 \setminus \Delta) \subset f(X \setminus \Delta) = \Lambda_1$ and if we suppose $\Lambda_k \subset \Lambda_{k-1}$, then $\Lambda_{k+1} = f(\Lambda_k \setminus \Delta) \subset f(\Lambda_{k-1} \setminus \Delta) = \Lambda_k$. 3) By injectivity we have $f(X \setminus \Delta) \cap f(\Delta) = \emptyset$, i.e. $\Lambda_1 \cap f(\Delta) = \emptyset$. Suppose now that $\Lambda_k \cap f^k(\Delta) = \emptyset$ and that there exists $x \in \Lambda_{k+1} \cap f^{k+1}(\Delta)$. Then there exist $y \in \Lambda_k \setminus \Delta$ and $y' \in f^k(\Delta)$ such that f(y) = f(y') = x. As f is injective, we must have y = y', which contradicts the recurrence hypothesis. It follows that $\Lambda_k \cap f^k(\Delta) = \emptyset$ for all $k \geq 1$ and applying 1) we obtain the desired result.

Proof. (Proposition 3.3). We compute the transfer operator P of f in order to use the formula (5). Since f is a diffeomorphism on the full Lebesgue measure open set $X \setminus \Delta$, its associated transfer operator acts on a function $\psi \in L^1(X)$ as

$$P\psi(x) = \mathbf{1}_{f(X \setminus \Delta)}(x)\psi(f^{-1}(x))|\det(f'(f^{-1}(x)))|^{-1} = \mathbf{1}_{\Lambda_1}(x)\psi(f^{-1}(x))J_1(x) \quad \forall x \in X.$$

Since the set $f(\Delta)$ has zero Lebesgue measure, we can suppose without loss of generality that $P\psi(x) = 0$ if $x \in f(\Delta)$. Then, $P\psi(x)$ is well defined for all $x \in X$, since $f(X \setminus \Delta) = \Lambda_1$ does not contain point of $f(\Delta)$. Now, we show by induction that for the uniform density $\mathbf{1} = \mathbf{1}_X$ we have

$$P^k \mathbf{1}(x) = \mathbf{1}_{\Lambda_k}(x) J_k(x) \qquad \forall k \geqslant 1.$$

It is easy to check that $P\mathbf{1}(x) = \mathbf{1}_{\Lambda_1}(x)J_1(x)$ and that

$$P^{k+1}\mathbf{1}(x) = \mathbf{1}_{\Lambda_1}(x)\mathbf{1}_{f(\Lambda_k)}(x)J_{k+1}(x),$$

if we suppose the property true at rank k. Now, we have

$$\Lambda_1 \cap f(\Lambda_k) = (\Lambda_1 \cap \Lambda_{k+1}) \cup (\Lambda_1 \cap f(\Lambda_k \cap \Delta)) = \Lambda_{k+1} \cup (\Lambda_1 \cap f(\Lambda_k \cap \Delta)).$$

Since $\Lambda_1 \cap f(\Lambda_k \cap \Delta) \subset \Lambda_1 \cap f(\Delta) = \emptyset$, we have $\Lambda_1 \cap f(\Lambda_k) = \Lambda_{k+1}$. It follows that $P^{k+1}\mathbf{1}(x) = \mathbf{1}_{\Lambda_{k+1}}(x)J_{k+1}(x)$, which ends the induction. Using (5) we obtain that

$$h_{\varepsilon}(x) = \varepsilon \sum_{k=0}^{\infty} (1 - \varepsilon)^k J_k(x) \mathbf{1}_{\Lambda_k}(x) \quad \forall x \in X,$$

which ends the proof of the proposition.

Now we decompose the space X as the disjoint union $X = \Lambda_1 \cup f(\Delta) \cup (X \setminus f(X))$ and we use (11) to study the density in each of these sets. Since the last two sets do not belong to any Λ_k with $k \neq 0$, we obtain that $h_{\varepsilon} = \varepsilon$ on these sets. For $x \in \Lambda_1$, by property 2) of Lemma 3.4, two situations can occur: either there exists $p \geqslant 1$ such that $x \in \Lambda_p \setminus \Lambda_{p+1}$ or $x \in \Lambda := \bigcap_{k=1}^{\infty} \Lambda_k$. So in Λ_1 we have on one hand

$$h_{\varepsilon}(x) = \varepsilon \sum_{k=0}^{p} (1 - \varepsilon)^{k} J_{k}(x) \qquad \forall x \in \Lambda_{p} \setminus \Lambda_{p+1}, \tag{12}$$

and the by continuity of the J_k , the density is continuous. On the other hand we have

$$h_{\varepsilon}(x) = \varepsilon \sum_{k=0}^{\infty} (1 - \varepsilon)^k J_k(x) \qquad \forall x \in \Lambda.$$
 (13)

In this last case the density may not be bounded, depending on the behavior of $J_k(x)$. However, if we suppose that there exists $0 < \lambda < (1 - \varepsilon)^{-1}$ such that $|\det(f'(x)|^{-1} \leq \lambda)$ for all $x \in X \setminus \Delta$, then the series (13) converges uniformly in Λ and h_{ε} is also continuous in Λ and thus on the whole set Λ_1 (see Proposition 4.3).

In the next sections we decompose our study of the existence of EVL in two parts. The first part concerns the point of $\Lambda_1 \setminus \tilde{\Lambda}$, where

$$\tilde{\Lambda} := \bigcap_{k \in \mathbb{N}} \overline{\Lambda}_k. \tag{14}$$

The set $\tilde{\Lambda}$ is called the *global attractor* of f and contains all the point of Λ . In $\Lambda_1 \setminus \tilde{\Lambda}$ we can verify the conditions $D_2(u_n)$ and $D'(u_n)$ to show the existence of EVL, but in $\tilde{\Lambda}$ condition $D'(u_n)$ is not verified. For this reason in the second part, we introduce other conditions which allow us to show the existence of EVL but with an extremal index different from 1; this is achieved by giving additional property, namely by choosing f as a piecewise contracting map. It has been shown in [4] that for any piecewise contracting map defined in a compact space, if the global attractor $\tilde{\Lambda}$ does not intersect the set of the discontinuities, then it is composed of a finite number of periodic orbits. Moreover, this condition, for injective maps as our, is generic in C^0 topology [3]. All these properties will help to get the existence of EVL in the set $\tilde{\Lambda}$. Let us give a few example to illustrate that

Example 3.1. As an easy example of an uncountable hyperbolic attractor, let us consider the baker's map f defined on the unit square $[0,1] \times [0,1]$ iteratively for $n \ge 1$, by

$$x_{n+1} = \begin{cases} \gamma_a x_n, & \text{for } y_n < \alpha \\ \frac{1}{2} + \gamma_b x_n, & \text{for } y_n > \alpha \end{cases}$$
$$y_{n+1} = \begin{cases} \frac{1}{\alpha} y_n, & \text{for } y_n < \alpha \\ \frac{1}{1-\alpha} (y_n - \alpha) + \gamma_b x_n, & \text{for } y_n > \alpha \end{cases}$$

with $0 \le x_n, y_n \le 1$, $0 < \gamma_a < \gamma_b < \frac{1}{2}$, $\alpha \le \frac{1}{2}$. In this case Δ is given by the union of the boundary of the square and the segment $y = \alpha$. Moreover $\overline{\Lambda}_1$ is given by the two rectangles $[0, \gamma_a] \times [0, 1]$, $[\frac{1}{2}, \frac{1}{2} + \gamma_b] \times [0, 1]$. Notice that the derivative f' is a diagonal matrix with diagonal elements γ_a (or γ_b), and $\frac{1}{\alpha}$ (or $\frac{1}{1-\alpha}$). We can choose these parameters in order to get $\det(f') < 1$.

Typical examples of PCM for which the attractor is generically a finite set of periodic orbits are piecewise affine [2, 17]. That is, if for each $i \in \{1, ..., N\}$ the restriction $f_i := f|_{X_i}$ of the map $f: X \to X$ of Definition 3.1 to a piece X_i has the following form:

$$f_i(x) = A_i x + c_i \quad \forall x \in X_i,$$

where $A_i : \mathbb{R}^D \to \mathbb{R}^D$ is a linear contraction and $c_i \in X$.

Example 3.2. 1-D case The simplest example of a piecewise contracting map of this class is given in dimension 1 in the interval [0,1] by f(x) = ax + c with $a, c \in (0,1)$ (here $X_1 = (0,(1-c)/a)$, $X_2 = ((1-c)/a,1)$ and $\Delta = \{0,1,(1-c)/a\}$). For almost all the values of the parameters a and c, the global attractor $\tilde{\Lambda}$ is composed of a unique periodic orbit (which period depends on the specific values of the parameters), but for the remaining set of parameters the global attractor is a Cantor set supporting a minimal dynamics. We refer to [5] for a detailed description of the asymptotic dynamics of this map. We remind that in our paper [7] we investigated in detail the case with c = 0 and 0 the (unique) fixed point; we briefly quote here the simple structure of the density

$$h_{\varepsilon}(x) = \varepsilon \sum_{k=0}^{p-1} \frac{(1-\varepsilon)^k}{a^k} \quad \forall x \in (a^p, a^{p-1}], \ p \ge 1$$

and the formulae for the coefficients of the linear scaling. At this regard let $\tau > 0$, $y = -\ln(\tau)$ and $u_n = \frac{y}{a_n} + b_n$ with

$$a_n = 1$$
 and $b_n = \log \left(2n\varepsilon \sum_{k=0}^{p-1} \frac{(1-\varepsilon)^k}{\alpha^k} \right)$ $\forall n \in \mathbb{N},$

If.

(i) $z \neq 0$ on the interval but not in the countably many discontinuity points of h_{ε} , namely $z \notin \bigcup_{j \in \mathbb{N}} \{a^j\}$; (ii) $p \geqslant 1$ such that $z \in (a^p, \alpha^{p-1})$ and n is large enough such that the ball $B(z, e^{-u_n}) \subset (a^p, a^{p-1})$, then:

$$n\mathbb{P}(X_0 > u_n) = \tau$$

Example 3.3. 2-D case In the unit square $X = [0,1]^2$, another interesting example is given by a map with four pieces $X_1 := (0,T_1) \times (0,T_2)$, $X_2 := (T_1,1) \times (0,T_2)$, $X_3 := (0,T_1) \times (T_2,1)$ and $X_4 := (T_1,1) \times (T_2,1)$, singular set Δ given by the two segments $\{x=T_1\}$, $\{y=T_2\}$ and the boundary of the unit square, and restricted maps $f_1(x,y) := a(x,y) + (1-a)(1,0)$, $f_2(x,y) := a(x,y) + (1-a)(1,1)$, $f_3(x,y) := a(x,y)$ and $f_4(x,y) := a(x,y) + (1-a)(0,1)$. Once again, for almost all values of the parameter $a \in [0,1)$, T_1 and T_2 , the global attractor is composed of a finite number of periodic orbits. However, contrarily to the 1 dimensional case, where the number of periodic orbits is bounded by the number of pieces N [18], the number of periodic orbits increases with the parameter a and tends to infinity when a goes to 1. This example belongs to a larger class of higher dimension piecewise affine contracting maps which are models for genetic regulatory networks and are studied in [6].

Question An interesting question is the check if the stationary measure is stochastically stable, in the sense that the measure μ_{ε} will converge weakly to the SRB measure supported on the global attractor. Let us remind that it is *not* the case for the simple contraction f(x) = ax, $x \in [0,1]$, a < 1, as it was proved in Lasota Mackay [16].

4 Convergence to the extreme value law

We prove in this section the convergence towards an extreme value distribution in the complement of the attractor defined in (14). We will work with the process $Y_n(x,\underline{\omega}) := -\log d(f_{\underline{\omega}}^n(x),z)$. Our first task will be to prove the existence of the limit (1); this will be accomplished in two manners: we will first give a general result showing the existence of that limit for any point of X, but without an explicit expression for the scaling coefficients a_n and b_n . Such coefficients could instead be computed on the set $\Lambda_1 \setminus \tilde{\Lambda}$ and even on the attractor itself provided the density of the stationary measure will be essentially bounded. The first general result uses Theorem 1.7.13 in the book by Leadbetter, Lindgren and Rootzen [15]: according to it, a sufficient condition to guarantee the existence of the limit (1) for $0 < \tau < \infty$, is that $\frac{1-F(x)}{1-F(x-)} \to 1$, as $x \to x_F$, where F is the distribution function of Y_0 , the term F(x-) in the denominator denotes the left limit of F at x and finally $x_F := \sup\{x; F(x) < 1\}$. For our particular observable Y_0 , $x_F = \infty$ and the probability $\mathbb P$ reduces to the stationary measure μ_{ε} which is absolutely continuous and therefore non atomic. Therefore F is continuous and the above ratio goes to 1. We have thus proved the following proposition:

Proposition 4.1. Suppose $f: X \to X$ satisfies Definition 3.1. Then, for any $z \in X$ defining the observable $Y_0(\cdot) = -\log d(\cdot, z)$ and any $0 < \tau < \infty$, there exists a sequence u_n such that

$$\lim_{n \to \infty} n \mathbb{P}(X_0 > u_n) = \tau.$$

As we anticipated above the previous result could be strengthened whenever z lies in the complement of the global attractor under the general assumption that h_{ε} is summable. Actually, we will suppose that $z \in (\Lambda_0 \setminus \tilde{\Lambda}) \setminus \bigcup_{k \geqslant 2}^{\infty} \partial \Lambda_k$, where $\partial \Lambda_k := \overline{\Lambda}_k \setminus \Lambda_k$; this in turn implies that $z \in \Lambda_p \setminus \overline{\Lambda}_{p+1} \subset \Lambda_p \setminus \Lambda_{p+1}$ for some $p \geqslant 1$.⁴ We can take advantage of the continuity of the density in such points to obtain a possible scaling sequence for which the limit (1) holds.

Proposition 4.2. Suppose $f: X \to X$ satisfies Definition 3.1. Let $p \geqslant 1$ such that $z \in \Lambda_p \setminus \overline{\Lambda}_{p+1}$. Let $\tau > 0$, $y = -\ln(\tau)$ and $u_n = \frac{y}{a_n} + b_n$ with $a_n = D$ and b_n such that

$$m(B(z, e^{-b_n})) = \left(n\varepsilon \sum_{k=0}^p (1-\varepsilon)^k J_k(z)\right)^{-1} = \frac{1}{nh_\varepsilon(z)} \quad \forall n \in \mathbb{N},$$

⁴We notice that in the case of the baker's transformation of Example 3.1 the boundaries of the Λ_k belong to the attractor and therefore we can take simply $z \in (\Lambda_0/\tilde{\Lambda})$.

where B(z,r) denotes a ball of center z and radius r. Then,

$$\lim_{n \to \infty} n \mathbb{P}(X_0 > u_n) = \tau. \tag{15}$$

Proof. For all $n \in \mathbb{N}$, we have

$$n\mathbb{P}(X_0 > u_n) = n \int \mathbf{1}_{B(z,e^{-u_n})}(x) d\mu_{\epsilon} = n \int \mathbf{1}_{B(z,e^{-u_n})}(x) h_{\varepsilon}(x) dx.$$

Since $(u_n)_{n\in\mathbb{N}}$ is an increasing sequence for any $\tau>0$ and z is in the open set $\Lambda_p\setminus\overline{\Lambda}_{p+1}$, if n is large enough $B(z,e^{-u_n})\subset\Lambda_p\setminus\Lambda_{p+1}$. Using formula (12) and $m(B(z,e^{-u_n}))=\tau m(B(z,e^{-b_n}))$ we obtain:

$$n\tau m(B(z, e^{-b_n}))\varepsilon \sum_{k=0}^{p} (1-\varepsilon)^k \inf_{x \in B(z, e^{-u_n})} J_k(x) \leqslant n\mathbb{P}(X_0 > u_n) \leqslant n\tau m(B(z, e^{-b_n}))\varepsilon \sum_{k=0}^{p} (1-\varepsilon)^k \sup_{x \in B(z, e^{-u_n})} J_k(x).$$
(16)

Since J_k is continuous in $B(z, e^{-u_n})$ for all $k \in \{0, \dots, p\}$, we have

$$\lim_{n \to \infty} \inf_{x \in B(z, e^{-u_n})} J_k(x) = J_k(z) = \lim_{n \to \infty} \sup_{x \in B(z, e^{-u_n})} J_k(x).$$

Using the expression of the sequence $(u_n)_{n\in\mathbb{N}}$, we obtain

$$\lim_{n \to \infty} n \mathbb{P}(X_0 > u_n) = \tau.$$

We note that an inspection of the proof of Proposition 4.2 easily shows that similar results can be obtained for the points of Λ , but supposing a continuous bounded density. The following proposition shows that the density has such properties in Λ under some condition on the derivative of f.

Proposition 4.3. Let $0 < \lambda < (1 - \varepsilon)^{-1}$ and suppose that $|\det(f'(x))|^{-1} \le \lambda$, for all $x \in X/\Delta$, then h_{ε} is continuous in Λ and the same conclusions of Proposition 4.2 hold.

Proof. For any $k \geq 0$ and $x \in X$ let $g_k(x) := \varepsilon(1-\varepsilon)^k J_k(x) \mathbf{1}_{\Lambda_k}(x)$. Then, $g_k(x) = 0$ if $x \neq \Lambda_k$ and $g_k(x) \leq \varepsilon((1-\varepsilon)\lambda)^k$ if $x \in \Lambda_k$ and therefore $\sum_{k=0}^p g_k(x)$ converges uniformly on X when $p \to \infty$, since $\sum_{k=0}^{\infty} ((1-\varepsilon)\lambda)^k$ converges. Now, let $x_0 \in \Lambda$, then

$$\lim_{x \to x_0} h_{\varepsilon}(x) = \lim_{x \to x_0} \sum_{k=0}^{\infty} g_k(x) = \sum_{k=0}^{\infty} \lim_{x \to x_0} g_k(x) = \varepsilon \sum_{k=0}^{\infty} (1 - \varepsilon)^k \lim_{x \to x_0} J_k(x) = \varepsilon \sum_{k=0}^{\infty} (1 - \varepsilon)^k J_k(x_0) = h_{\varepsilon}(x_0)$$

We have used that $x_0 \in \Lambda_k$ for all $k \geq 0$, which implies on one hand by openness of the Λ_k that $\lim_{x \to x_0} \mathbf{1}_{\Lambda_k}(x) = 1$ and in the other hand that J_k is continuous in x_0 , since for any $x \in \Lambda_k$ we have $f^l(x) \notin f(\Delta)$ for all $l \leq k$ and f is C^1 in $X \setminus \Delta$.

We are now ready to check conditions $D_2(u_n)$ and $D'(u_n)$ which will guarantee the convergence of our process toward an extreme value law, the Gumbel's one in our case. As we said above condition $D_2(u_n)$ is insured by the decay of correlations result established in Proposition 2.1, see the proof of Th. D in Sect. 4.1 in [1] for the very general approach. We are thus left with the proof of $D'(u_n)$. We first need:

Lemma 4.4. Let $n \in \mathbb{N}$ and $U_n = B(z, e^{-u_n})$ be a ball centered at a point $z \in X$ with radius e^{-u_n} . Let $p \ge 1$ be such that $z \in \Lambda_p \setminus \overline{\Lambda}_{p+1}$ then,

$$\mathbb{P}((x,\omega_1,\ldots,\omega_j):x\in U_n,f_{\underline{\omega}}^j(x)\in U_n)=\epsilon\mu_{\epsilon}(U_n)\sum_{k=0}^{\min\{j-1,p\}}(1-\epsilon)^k\int J_k(x)\mathbf{1}_{U_n}(x)dx\qquad\forall j\geqslant 1,\quad(17)$$

for n large enough. In particular, under the same hypothesis we have,

$$\mathbb{P}((x,\omega_1,\ldots,\omega_j):x\in U_n,f^j_{\underline{\omega}}(x)\in U_n)\leqslant \mu_{\epsilon}(U_n)^2 \qquad \forall j\geqslant 1.$$
 (18)

Proof. Let us denote

$$\Pr_j := \mathbb{P}((x, \omega_1, \dots, \omega_j) : x \in U_n, f_\omega^j(x) \in U_n) \quad \forall j \geqslant 1.$$

First of all, for all $j \ge 1$ we can write

$$\Pr_{j} = \int \int \mathbf{1}_{U_{n}}(x)\mathbf{1}_{U_{n}}(f_{\underline{\omega}}^{j}(x))d\theta_{\epsilon}^{\mathbb{N}}d\mu_{\epsilon} = \int \mathbf{1}_{U_{n}}(x)U_{\epsilon}^{j}(\mathbf{1}_{U_{n}})(x)d\mu_{\epsilon} = \int P_{\epsilon}^{j}(\mathbf{1}_{U_{n}}h_{\epsilon})(x)\mathbf{1}_{U_{n}}(x)dx.$$

Using (8), we obtain that for all $j \ge 1$

$$\operatorname{Pr}_{j} = \int P_{\epsilon}^{j}(\mathbf{1}_{U_{n}}h_{\epsilon})(x)\mathbf{1}_{U_{n}}(x)dx
= (1-\epsilon)^{j} \int P^{j}(\mathbf{1}_{U_{n}}h_{\epsilon})(x)\mathbf{1}_{U_{n}}(x)dx + \epsilon\mu_{\epsilon}(U_{n}) \sum_{k=0}^{j-1} (1-\epsilon)^{k} \int P^{k}\mathbf{1}(x)\mathbf{1}_{U_{n}}(x)dx
= (1-\epsilon)^{j} \int \mathbf{1}_{U_{n}}(x)\mathbf{1}_{U_{n}}(f^{j}(x))h_{\epsilon}(x)dx + \epsilon\mu_{\epsilon}(U_{n}) \sum_{k=0}^{j-1} (1-\epsilon)^{k} \int J_{k}(x)\mathbf{1}_{\Lambda_{k}}(x)\mathbf{1}_{U_{n}}(x)dx
= (1-\epsilon)^{j} \mu_{\epsilon}(U_{n} \cap f^{-j}(U_{n})) + \epsilon\mu_{\epsilon}(U_{n}) \sum_{k=0}^{j-1} (1-\epsilon)^{k} \int J_{k}(x)\mathbf{1}_{\Lambda_{k} \cap U_{n}}(x)dx.$$

Now suppose n large enough for U_n being a subset of $\Lambda_p \setminus \Lambda_{p+1}$. Then one can show that $f^j(U_n) \subset \bigcup_{k=1}^j f^k(\Delta) \cup \Lambda_{p+j}$ (using recurrence on j). Now, in the union $\bigcup_{k=1}^j f^k(\Delta)$ we have that $\Lambda_p \cap f^k(\Delta) = \emptyset$ for any $k \leq p$. We also have $f^l(\Lambda_p) \cap f^{p+l}(\Delta) = \emptyset$ for any $l \in \mathbb{N}$, since $\Lambda_p \cap f^p(\Delta) = \emptyset$ and f is injective. Together with $\Lambda_{p+j} \subset \Lambda_p$, this implies that $U_n \cap f^j(U_n) = \emptyset$, and therefore $U_n \cap f^{-j}(U_n) = \emptyset$ for all $j \geq 1$. On the other hand, $U_n \cap \Lambda_k = \emptyset$ for all k > p. It follows that

$$\Pr_{j} = \epsilon \mu_{\epsilon}(U_{n}) \sum_{k=0}^{\min\{j-1,p\}} (1-\epsilon)^{k} \int J_{k}(x) \mathbf{1}_{U_{n}}(x) dx \qquad \forall j \geqslant 1,$$

which prove (17). Now, it easy to show using (12) that $\Pr_j \leqslant \Pr_{p+1} = \mu_{\epsilon}(U_n)^2$ for any $j \geqslant 1$.

Proposition 4.5. Let $p \ge 1$ and $z \in \Lambda_p \setminus \overline{\Lambda}_{p+1}$, then the maxima of the process $X_n(x,\underline{\omega}) := -\ln d(f_{\underline{\omega}}^n(x),z)$ verifies condition $D'(u_n)$.

Proof. Suppose $z \in \Lambda_p \setminus \overline{\Lambda}_{p+1}$ for some $p \ge 1$. Consider the sequence $(u_n(y))_{n \in \mathbb{N}}$ of Proposition 4.2. Let $(k_n)_{n \in \mathbb{N}}$ be such that $k_n \to \infty$ when $n \to \infty$ and $k_n = o(n)$. Supposing n large enough and using (18), we deduce that

$$n\sum_{j=1}^{\left\lfloor \frac{n}{k_n}\right\rfloor} \mathbb{P}((x,\omega_1,\ldots,\omega_j): x\in U_n, f_{\underline{\omega}}^j(x)\in U_n)\leqslant n\left\lfloor \frac{n}{k_n}\right\rfloor \mu_{\epsilon}(U_n)^2\leqslant \frac{(n\mu_{\epsilon}(U_n))^2}{k_n}.$$

Using $\lim n\mu_{\varepsilon}(U_n) = \tau$, we obtain that

$$\lim_{n \to \infty} n \sum_{j=1}^{\left[\frac{n}{k_n}\right]} \mathbb{P}((x, \omega_1, \dots, \omega_j) : x \in U_n, f_{\omega_j} \circ \dots \circ f_{\omega_1}(x) \in U_n) = 0,$$

since $n\mu_{\epsilon}(U_n) \to \tau$ and $k_n \to \infty$ when $n \to \infty$. In such a case, $(X_n(x,\underline{\omega}))_{n \in \mathbb{N}}$ verifies conditions $D_2(u_n)$ and $D'(u_n)$ for our sequence $(u_n(y))_{n \in \mathbb{N}}$, wich implies that $\mathbb{P}(M_n \leq u_n(y))$ converges towards the Gumbel law $e^{-e^{-y}}$.

Collecting all the results of this section we finally proved the following theorem

Theorem 4.6. Suppose $f: X \to X$ satisfies Definition 3.1. Then, for any $z \in (\Lambda_1 \setminus \hat{\Lambda}) \setminus \bigcup_{k \geq 2}^{\infty} \partial \Lambda_k$, the sequence $\{M_n\}_{n \in \mathbb{N}}$ of the maxima of the process defined for every $n \in \mathbb{N}$ by $Y_n(x, \underline{\omega}) := -\log d(f_{\underline{\omega}}^n(x), z)$ admits the Gumbel's law as extreme values distribution.

5 Clustering in the attractor of piecewise contracting maps

In this section, we complete the results of the results of Theorem 5.1 by studying the case when the maximum of the observable is achieved in the set $\tilde{\Lambda}$ of a piecewise contracting map. This creates clustering of exceedances which are responsible for the appearance of an Extremal Index less than 1. To prove the existence of such an EVL, we follow the techniques developed in [13] that were recently upgraded in [14]. In particular, we will use the notation used in the latter. Actually in [13] the authors introduced two conditions on the dependence structure of the stochastic processes in order to obtain convergence to an EVL under the presence of clustering. These conditions were called $D^p(u_n)$ and $D'_p(u_n)$ and essentially were adaptations of conditions $D_2(u_n)$ and $D'(u_n)$, specifically designed to cope with clustering. Later, in [14], the authors introduced some general conditions called $\mathcal{I}(u_n)$ and $\mathcal{I}'_q(u_n)$, which polished and joined all the previous conditions, allowing to address the presence $(q \ge 1)$ and absence (q = 0) of clustering at once. We recall this conditions here.

Let us begin to introduce a few notations. For $q \in \mathbb{N}_0$ we put

$$A_n^{(q)} := \{ Y_0 > u_n, Y_1 \le u_n, \dots, Y_q \le u_n \}$$

where again $Y_k, k \in \mathbb{N}_0$ is a \mathbb{P} -stationary process, that we will take as that studied in the preceding chapters. For $s, \ell \in \mathbb{N}$ and an event B, let

$$\mathcal{W}_{s,\ell}(B) = \bigcap_{i=s}^{s+\ell-1} T^{-i}(B^c).$$
 (19)

Condition $(\mathcal{A}(u_n))$. We say that $\mathcal{A}(u_n)$ holds for the sequence Y_0, Y_1, \ldots if for every $\ell, t, n \in \mathbb{N}$ and $q \in \mathbb{N}_0$,

$$\left| \mathbb{P}\left(A_n^{(q)} \cap \mathcal{W}_{t,\ell}\left(A_n^{(q)} \right) \right) - \mathbb{P}\left(A_n^{(q)} \right) \mathbb{P}\left(\mathcal{W}_{0,\ell}\left(A_n^{(q)} \right) \right) \right| \leqslant \gamma(q,n,t), \tag{20}$$

where $\gamma(q,n,t)$ is decreasing in t for each q,n and, for every $q \in \mathbb{N}_0$, there exists a sequence $(t_n)_{n \in \mathbb{N}}$ such that $t_n = o(n)$ and $n\gamma(q,n,t_n) \to 0$ when $n \to \infty$.

For some fixed $q \in \mathbb{N}_0$, consider the sequence $(t_n)_{n \in \mathbb{N}}$, given by condition $\mathcal{A}(u_n)$ and let $(k_n)_{n \in \mathbb{N}}$ be another sequence of integers such that

$$k_n \to \infty$$
 and $k_n t_n = o(n)$. (21)

Condition $(A'_q(u_n))$. We say that $A'_q(u_n)$ holds for the sequence $Y_0, Y_1, Y_2, ...$ if there exists a sequence $(k_n)_{n\in\mathbb{N}}$ satisfying (21) and such that

$$\lim_{n \to \infty} n \sum_{j=q+1}^{\lfloor n/k_n \rfloor - 1} \mathbb{P}\left(A_n^{(q)} \cap T^{-j}\left(A_n^{(q)}\right)\right) = 0.$$
 (22)

Now let

$$\vartheta = \lim_{n \to \infty} \vartheta_n = \lim_{n \to \infty} \frac{\mathbb{P}(A_n^{(q)})}{\mathbb{P}(U_n)}.$$
 (23)

From [14, Corollary 2.4], it follows that if the stochastic process Y_0, Y_1, \ldots satisfies conditions $\mathcal{A}(u_n)$ and $\mathcal{A}'_q(u_n)$ and the limit in (23) exists then $\lim_{n\to\infty} \mathbb{P}(M_n \leqslant u_n) = \mathrm{e}^{-\vartheta \tau}$, where the sequence $(u_n)_{n\in\mathbb{N}}$ is such that $\lim_{n\to\infty} n\mathbb{P}(X_0 > u_n) = \tau$.

5.1 Analysis in periodic attractors

Using the result just mentioned we will show that

Theorem 5.1. Suppose that $f: X \to X$ satisfies Definition 3.1. Suppose moreover that f is piecewise contracting, that is, there exists $\alpha \in (0,1)$ such that for every $i \in \{1,\ldots,N\}$ we have

$$d(f(x), f(y)) \leq \alpha d(x, y) \quad \forall x, y \in X_i.$$

Suppose that the set $\tilde{\Lambda}$ does not contain any discontinuity point of f, and let $z \in \tilde{\Lambda}$. Consider the random perturbations defined in (3) and let $\mathbb{P} = \mu_{\varepsilon} \times \theta_{\varepsilon}^{\mathbb{N}}$, where μ_{ε} is a stationary probability measure. Let Y_0, Y_1, \ldots be the stochastic process given by $Y_n(x,\underline{\omega}) := -\log d(f_{\underline{\omega}}^n(x), \mathcal{O}(z))$, where $\mathcal{O}(z)$ is the orbit of z. Let $(u_n)_{n \in \mathbb{N}}$ be a sequence such that $\lim_{n \to \infty} n\mathbb{P}(X_0 > u_n) = \tau \geqslant 0$. Then we have

$$\lim_{n\to\infty} \mathbb{P}(M_n \leqslant u_n) = e^{-\varepsilon\tau}.$$

It has been shown in [4] that for any piecewise contracting map defined in a compact space, if the global attractor $\bigcap_{k=1}^{\infty} \overline{\Lambda}_k$ does not intersect the set of the discontinuities, then it is composed of finite number of periodic orbits. Therefore, under this hypothesis $\mathcal{O}(z)$ is a periodic orbit of the map. If p is the period of this orbit, then $\{Y_0 > u_n\}$ occurs if and only if $x \in U_n$, where

$$U_n = \bigcup_{i=0}^{p-1} B(f^i(z), e^{-u_n}).$$

Remark 5.2. Note that on the contrary to the usual choice for the observable function, here, we are considering that the maximum is achieved at every point of the whole periodic orbit, instead of at a single periodic point. In an ongoing work concerning EVLs for observables achieving a maximum on multiple correlated points, the second-named author with D. Azevedo, A.C.M. Freitas and F.B. Rodrigues are developing a thorough study regarding the consequences of having multiple maxima in terms of the extremal behavior of the system. The important fact to retain here is that by considering the whole orbit we are performing a sort of reduction to the case where the period is 1.

As in the previous section, Proposition 4.1 guaranties the existence of the a sequence $\{u_n\}_{n\in\mathbb{N}}$ such that the limit (1) holds. However, as a side comment, we note that we can obtain an estimation of the sequence $\{u_n\}_{n\in\mathbb{N}}$ using a similar proof as the one of Proposition 4.2, provided the density is bounded and continuous. Since $\tilde{\Lambda} \cap \Delta = \emptyset$, it is composed of finite set of periodic points that are not in Δ and using the definition of the sets Λ_k , it is easy to show that all the points of $\tilde{\Lambda}$ also belongs to Λ , i.e $\tilde{\Lambda} = \Lambda$. Thus, we can argue that, under the hypothesis of Theorem 5.1 and of Proposition 4.3, the density has the desired properties on $\tilde{\Lambda}$ to obtain an estimation of the sequence $\{u_n\}_{n\in\mathbb{N}}$.

The proof of the theorem consists in showing that Y_0, Y_1, \ldots satisfies conditions $\mathcal{A}(u_n)$ and $\mathcal{A}'_q(u_n)$ and moreover ϑ given in (23) is well defined and equals ε . Since $\mathcal{O}(z) \cap \Delta = \emptyset$, for n large enough each ball of the union defining U_n is contained in one of the open pieces X_i , and because of the contraction we have that $f(U_n) \subset U_n$. Therefore, we will consider q = 1 and use $A_n^{(1)} = \{(x, \underline{\omega}) : x \in U_n, f_{\underline{\omega}}(x) \notin U_n\}$.

Verification of $A(u_n)$:

From the proof of Proposition 2.1 we have:

$$Cor_{m}(\phi, \psi, n) := \left| \iint \psi(x) \phi(f_{\underline{\omega}}^{n}(x)) d\theta_{\varepsilon}^{\mathbb{N}}(\underline{\omega}) dx - \int \psi(x) dx \int \phi(x) d\mu_{\varepsilon}(x) \right| \leq 2(1 - \varepsilon)^{n} \|\phi\|_{\infty} \|\psi\|_{1}. \tag{24}$$

We will use this information about decay of correlations to verify condition $\mathcal{A}(u_n)$, in the case q=1. We have

$$\mathbb{P}\left(A_{n}^{(1)} \cap \mathcal{W}_{t,\ell}(A_{n}^{(1)})\right) = \iint \mathbf{1}_{U_{n}}(x)\mathbf{1}_{U_{n}^{c}}(f_{\omega_{1}})(x)\mathbf{1}_{\mathcal{W}_{0,\ell}(A_{n}^{(1)})}(f_{\underline{\omega}}^{t}(x))d\theta_{\varepsilon}^{\mathbb{N}}(\underline{\omega})d\mu_{\varepsilon}(x) \\
= \varepsilon \iiint \mathbf{1}_{U_{n}}(x)\mathbf{1}_{U_{n}^{c}}(y)\mathbf{1}_{\mathcal{W}_{0,\ell}(A_{n}^{(1)})}(f_{\underline{\sigma}\underline{\omega}}^{t-1}(y))dyd\theta_{\varepsilon}^{\mathbb{N}}(\underline{\omega})d\mu_{\varepsilon}(x) \\
+ (1 - \varepsilon) \iint \mathbf{1}_{U_{n}}(x)\mathbf{1}_{U_{n}^{c}}(f(x))\mathbf{1}_{\mathcal{W}_{0,\ell}(A_{n}^{(1)})}(f_{\underline{\sigma}\underline{\omega}}^{t-1}(f(x)))d\theta_{\varepsilon}^{\mathbb{N}}(\underline{\omega})d\mu_{\varepsilon}(x) \\
= \varepsilon \mu_{\varepsilon}(U_{n}) \iint \mathbf{1}_{U_{n}^{c}}(y)\mathbf{1}_{\mathcal{W}_{0,\ell}(A_{n}^{(1)})}(f_{\underline{\sigma}\underline{\omega}}^{t-1}(y))d\theta_{\varepsilon}^{t+\ell-1}(\omega_{2}, \dots, \omega_{t+\ell})dy \\
= \varepsilon \mu_{\varepsilon}(U_{n}) \iint \mathbf{1}_{U_{n}^{c}}(y)\phi(f_{\underline{\sigma}\underline{\omega}}^{t-1}(y))d\theta_{\varepsilon}^{t-1}(\omega_{2}, \dots, \omega_{t})dy,$$

where

$$\phi(x) = \int \mathbf{1}_{\mathcal{W}_{0,\ell}(A_n^{(1)})}(x)\theta_{\varepsilon}^{\ell}(\tilde{\omega}_1, \dots, \tilde{\omega}_{\ell}).$$

In the third equality above we used the fact that if $f(U_n) \subset U_n$ and consequently $\mathbf{1}_{U_n}(x) \cdot \mathbf{1}_{U_n^c}(f(x)) = 0$ for all $x \in X$. We will use this fact again in the following computation:

$$\mathbb{P}\left(A_n^{(1)}\right) = \iint \mathbf{1}_{U_n}(x)\mathbf{1}_{U_n^c}(f_{\omega_1})(x))d\theta_{\varepsilon}^{\mathbb{N}}(\underline{\omega})d\mu_{\varepsilon}(x) = \varepsilon \iint \mathbf{1}_{U_n}(x)\mathbf{1}_{U_n^c}(y)dyd\mu_{\varepsilon}(x) + (1-\varepsilon)\int \mathbf{1}_{U_n}(x)\mathbf{1}_{U_n^c}(f(x))d\mu_{\varepsilon}(x) = \varepsilon\mu_{\varepsilon}(U_n)\int \mathbf{1}_{U_n^c}(y)dy. \tag{25}$$

Also

$$\mathbb{P}\left(W_{0,\ell}(A_n^{(1)})\right) = \iint \mathbf{1}_{\mathcal{W}_{0,\ell}(A_n^{(1)})}(x)\theta_{\varepsilon}^{\mathbb{N}}(\underline{\omega})d\mu_{\varepsilon}(x) = \int \phi(x)d\mu_{\varepsilon}(x).$$

It follows by (24)

$$\begin{split} \left| \mathbb{P}\left(A_n^{(1)} \cap \mathscr{W}_{t,\ell}(A_n^{(1)}) \right) - \mathbb{P}\left(A_n^{(1)} \right) \mathbb{P}\left(W_{0,\ell}(A_n^{(1)}) \right) \right| \\ &= \varepsilon \mu_{\varepsilon}(U_n) \left| \iint \mathbf{1}_{U_n^c}(y) \phi(f_{\sigma\underline{\omega}}^{t-1}(y)) d\theta_{\varepsilon}^{\mathbb{N}}(\underline{\omega}) dy - \int \mathbf{1}_{U_n^c}(y) dy \int \phi(x) d\mu_{\varepsilon}(x) \right| \\ &= \varepsilon \mu_{\varepsilon}(U_n) Cor_m(\phi, \mathbf{1}_{U_n^c}, t-1) \leqslant 2\varepsilon \mu_{\varepsilon}(U_n) (1-\varepsilon)^t \end{split}$$

Hence, we may take $t_n = \log n$, for example, and $\mathcal{A}(u_n)$ is easily verified with $\gamma(n,t) = 2\varepsilon \mu_{\varepsilon}(U_n)(1-\varepsilon)^t$.

Verification of $\coprod_{a}'(u_n)$:

For
$$j > 1$$
, let $\Pr_{j}^{(1)} = \mathbb{P}(\{(x,\underline{\omega}) : x \in A_{n}^{(1)}, f_{\underline{\omega}}^{j}(x) \in A_{n}^{(1)}\})$. We have
$$\Pr_{j}^{(1)} = \iint \mathbf{1}_{U_{n}}(x) \cdot \mathbf{1}_{U_{n}^{c}}(f_{\underline{\omega}}(x)) \cdot \mathbf{1}_{U_{n}}(f_{\underline{\omega}}^{j}(x)) \cdot \mathbf{1}_{U_{n}^{c}}(f_{\underline{\omega}}^{j+1}(x)) d\theta_{\varepsilon}^{\mathbb{N}} h_{\varepsilon}(x) dx$$

$$= \varepsilon \iiint \mathbf{1}_{U_{n}}(x) \cdot \mathbf{1}_{U_{n}^{c}}(y) \cdot \mathbf{1}_{U_{n}}(f_{\underline{\sigma}\underline{\omega}}^{j-1}(y)) \cdot \mathbf{1}_{U_{n}^{c}}(f_{\underline{\sigma}\underline{\omega}}^{j}(y)) dy d\theta_{\varepsilon}^{\mathbb{N}} h_{\varepsilon}(x) dx$$

$$+ (1 - \varepsilon) \iint \mathbf{1}_{U_{n}}(x) \cdot \mathbf{1}_{U_{n}^{c}}(f(x)) \cdot \mathbf{1}_{U_{n}}(f_{\underline{\sigma}\underline{\omega}}^{j-1}(f(x))) \cdot \mathbf{1}_{U_{n}^{c}}(f_{\underline{\sigma}\underline{\omega}}^{j}(f(x))) d\theta_{\varepsilon}^{\mathbb{N}} h_{\varepsilon}(x) dx$$

$$= \varepsilon \int \mathbf{1}_{U_{n}}(x) h_{\varepsilon}(x) dx \iint \mathbf{1}_{U_{n}^{c}}(y) \cdot \mathbf{1}_{U_{n}}(f_{\underline{\sigma}\underline{\omega}}^{j-1}(y)) \cdot \mathbf{1}_{U_{n}^{c}}(f_{\underline{\sigma}\underline{\omega}}^{j}(y)) dy d\theta_{\varepsilon}^{\mathbb{N}}$$

$$\leqslant \varepsilon \mu_{\varepsilon}(U_{n}) \iint \mathbf{1}_{U_{n}^{c}}(y) \cdot \mathbf{1}_{U_{n}}(f_{\underline{\sigma}\underline{\omega}}^{j-1}(y)) d\theta_{\varepsilon}^{\mathbb{N}} dy = \varepsilon \mu_{\varepsilon}(U_{n}) \iint \mathbf{1}_{U_{n}^{c}}(y) \cdot U_{\varepsilon}^{j-1}(\mathbf{1}_{U_{n}})(y) dy$$

$$= \varepsilon \mu_{\varepsilon}(U_{n}) \iint P_{\varepsilon}^{j-1}(\mathbf{1}_{U_{n}^{c}})(y) \cdot \mathbf{1}_{U_{n}}(y) dy,$$

where in the third equality we used the fact that if $x \in U_n$ then $f(x) \notin U_n^c$ and consequently $\mathbf{1}_{U_n}(x) \cdot \mathbf{1}_{U_n^c}(f(x)) = 0$ for all $x \in X$.

Now, using (8), we define:

$$\iint P_{\varepsilon}^{j-1}(\mathbf{1}_{U_{n}^{c}})(y) \cdot \mathbf{1}_{U_{n}}(y) \, dy = (1-\varepsilon)^{j-1} \int P^{j-1}(\mathbf{1}_{U_{n}^{c}})(y) \mathbf{1}_{U_{n}}(y) dy + \varepsilon m(U_{n}^{c}) \sum_{k=0}^{j-2} (1-\varepsilon)^{k} \int P^{k} \mathbf{1}(y) \mathbf{1}_{U_{n}}(y) dy$$

$$= (1-\varepsilon)^{j-1} m(U_{n}^{c} \cap f^{j-1}(U_{n})) + \varepsilon m(U_{n}^{c}) \sum_{k=0}^{j-2} (1-\varepsilon)^{k} m(f^{-k}(U_{n}))$$

$$=: I_{j} + II_{j}$$

Note that to check (22) we need to show that $n \sum_{j=2}^{\lfloor n/k_n \rfloor - 1} \Pr_j^{(1)}$ vanishes. But $n \sum_{j=2}^{\lfloor n/k_n \rfloor - 1} \Pr_j^{(1)} \leqslant n \varepsilon \mu_{\varepsilon}(U_n) \sum_{j=2}^{\lfloor n/k_n \rfloor - 1} I_j + II_j$. Since $\lim_{n \to \infty} n \mu_{\varepsilon}(U_n) = \tau > 0$, then to verify $\mathcal{I}_q'(u_n)$ we only need to check that both $\sum_{j=2}^{\lfloor n/k_n \rfloor - 1} I_j$ and $\sum_{j=2}^{\lfloor n/k_n \rfloor - 1} I_j$ vanish as $n \to \infty$.

Since $\Lambda \cap \Delta = \emptyset$, there exists $k_0 := \min\{k \in \mathbb{N} : \Lambda_k \cap \Delta = \emptyset\}$ and since the sets Λ_k are open and satisfy $\Lambda_{k+1} \subset \Lambda_k$ for all $k \in \mathbb{N}$, for a n large enough there is a $k' \geqslant k_0$ such that $U_n \subset \Lambda'_k$. Therefore, we can define $d_n := k_0 + d'_n$, where $d'_n := \max\{k \in \mathbb{N} : U_n \subset \Lambda_k\}$. Note that $\lim_{n \to \infty} d_n = \infty$, since the distance $d(\Lambda_k, \Lambda) \to 0$ when $k \to \infty$. Moreover, as for any $k \geqslant k_0$ we have $\Lambda_{k+1} = f(\Lambda_k)$, it follows that $f^{-k}(U_n) \subset f^{-k}(\Lambda_{k_0+d_n}) = f^{d_n-k}(\Lambda_{k_0})$ for all $k \in \{0, 1, \ldots, d_n\}$. Also, for any $k \in \{0, 1, \ldots, d_n\}$, we have

$$m(f^{-k}(U_n)) \leqslant m(f^{d_n-k}(\Lambda_{k_0})) \leqslant \alpha^{-k}\beta_n \text{ where } \beta_n := \alpha^{d_n}m(\Lambda_{k_0}).$$

Now we can check (22) starting with $\sum_{j=2}^{\lfloor n/k_n\rfloor-1} I_j$ and splitting the sum

$$\sum_{j=2}^{\lfloor n/k_n \rfloor - 1} I_j = \sum_{i=1}^{\lfloor n/k_n \rfloor - 2} I_{i+1} = \sum_{i=1}^{d_n} I_{i+1} + \sum_{i=d_n+1}^{\lfloor n/k_n \rfloor - 2} I_{i+1}.$$

For the second sum on the right we have:

$$\sum_{i=d_n+1}^{\lfloor n/k_n\rfloor-2} I_{i+1} \leqslant \sum_{i=d_n+1}^{\infty} (1-\varepsilon)^i \leqslant \varepsilon^{-1} (1-\varepsilon)^{d_n+1} \xrightarrow[n\to\infty]{} 0, \text{ because } \lim_{n\to\infty} d_n = \infty.$$

For the first sum on the right, observe that

$$\sum_{i=1}^{d_n} I_{i+1} = \sum_{i=1}^{d_n} (1-\varepsilon)^i m(U_n^c \cap f^{-i}(U_n)) \leqslant \sum_{i=1}^{d_n} (1-\varepsilon)^i m(f^{-i}(U_n)) \leqslant \beta_n \sum_{i=1}^{d_n} (1-\varepsilon)^i \alpha^{-i}.$$

Now to analyze the asymptotic behavior of $\beta_n \sum_{i=1}^{d_n-1} (1-\varepsilon)^i \alpha^{-i}$ we consider three different cases. We assume first that $(1-\varepsilon)\alpha^{-1} > 1$.

$$\beta_n \sum_{i=1}^{d_n} (1-\varepsilon)^i \alpha^{-i} = \beta_n \frac{[(1-\varepsilon)\alpha^{-1}]^{d_n+1} - (1-\varepsilon)\alpha^{-1}}{(1-\varepsilon)\alpha^{-1} - 1} \leqslant \alpha^{d_n} m(\Lambda_{k_0}) \frac{[(1-\varepsilon)\alpha^{-1}]^{d_{n+1}}}{(1-\varepsilon)\alpha^{-1} - 1}$$
$$\leqslant \frac{m(\Lambda_{k_0})}{(1-\varepsilon)-\alpha} (1-\varepsilon)^{d_{n+1}} \xrightarrow[n\to\infty]{} 0, \text{ since } \lim_{n\to\infty} d_n = \infty.$$

In the case $(1-\varepsilon)\alpha^{-1} < 1$, then

$$\beta_n \sum_{i=1}^{d_n} (1-\varepsilon)^i \alpha^{-i} \leqslant \beta_n \sum_{i=1}^{\infty} (1-\varepsilon)^i \alpha^{-i} \xrightarrow[n \to \infty]{} 0, \text{ since } \lim_{n \to \infty} \beta_n = 0.$$

In the case $(1 - \varepsilon)\alpha^{-1} = 1$, then

$$\beta_n \sum_{i=1}^{d_n} (1 - \varepsilon)^i \alpha^{-i} \leqslant \beta_n d_n \leqslant d_n \alpha^{d_n} \xrightarrow[n \to \infty]{} 0, \text{ since } \lim_{n \to \infty} d_n = \infty.$$

Now, we turn to $\sum_{j=2}^{\lfloor n/k_n\rfloor-1} II_j$. Let β_n and d_n be defined as above, then

$$\sum_{j=2}^{\lfloor n/k_n\rfloor - 1} H_j = \sum_{i=0}^{\lfloor n/k_n\rfloor - 3} H_{i+2} = \sum_{i=0}^{d_n} \varepsilon m(U_n^c) \sum_{k=0}^i (1 - \varepsilon)^k m(f^{-k}(U_n)) + \sum_{i=d_n+1}^{\lfloor n/k_n\rfloor - 3} \varepsilon m(U_n^c) \sum_{k=0}^i (1 - \varepsilon)^k m(f^{-k}(U_n))$$

$$\leqslant \sum_{i=0}^{d_n} \varepsilon \sum_{k=0}^i (1 - \varepsilon)^k \alpha^{-k} \beta_n + \sum_{i=1}^{\lfloor n/k_n\rfloor} \varepsilon \sum_{k=0}^{\infty} (1 - \varepsilon)^k m(f^{-k}(U_n))$$

$$\leqslant \varepsilon \beta_n \left(1 + \sum_{i=1}^{d_n} \sum_{k=0}^i (1 - \varepsilon)^k \alpha^{-k} \right) + \frac{n}{k_n} \mu_{\varepsilon}(U_n)$$

Here we have used (5) to obtain

$$\mu_{\varepsilon}(U_n) = \int \mathbf{1}_{U_n}(x)h_{\varepsilon}(x)dx = \varepsilon \sum_{k=0}^{\infty} (1-\varepsilon)^k \int P^k \mathbf{1}(x)\mathbf{1}_{U_n}(x)dx = \varepsilon \sum_{k=0}^{\infty} (1-\varepsilon)^k m(f^{-k}(U_n)).$$

The second term on the right tends to 0 when n goes to infinity, since $n\mu_{\varepsilon}(U_n) \to \tau$ and $k_n \to \infty$ when n goes to ∞ . For the first term we assume first that $(1 - \varepsilon) \neq \alpha$. Then,

$$\varepsilon \beta_n \left(1 + \sum_{i=1}^{d_n} \sum_{k=0}^i (1 - \varepsilon)^k \alpha^{-k} \right) \leqslant \varepsilon \beta_n + \frac{\varepsilon \alpha}{(1 - \varepsilon) - \alpha} \beta_n \sum_{i=1}^{d_n} [(1 - \varepsilon) \alpha^{-1}]^{i+1}$$

$$= \varepsilon \beta_n + \frac{\varepsilon (1 - \varepsilon)}{(1 - \varepsilon) - \alpha} \beta_n \sum_{i=1}^{d_n} (1 - \varepsilon)^i \alpha^{-i}$$

Since, as we have seen above, $\lim_{n\to\infty} \beta_n \sum_{i=1}^{d_n} (1-\varepsilon)^i \alpha^{-i} = 0$, then $\lim_{n\to\infty} \sum_{j=2}^{\lfloor n/k_n\rfloor-1} II_j = 0$. Now, if $(1-\varepsilon) \neq \alpha$, we have

$$\varepsilon \beta_n \left(1 + \sum_{i=1}^{d_n} \sum_{k=0}^i (1 - \varepsilon)^k \alpha^{-k} \right) = \varepsilon m(\Lambda_{k_0}) \left(\alpha^{d_n} + d_n \alpha^{d_n} + \frac{1}{2} d_n^2 \alpha^{d_n} \right) \xrightarrow[n \to \infty]{} 0.$$

Convergence of ϑ_n :

As in equation (23) we consider $\vartheta_n = \frac{\mathbb{P}(A_n^{(1)})}{\mathbb{P}(U_n)}$. By (25) we have

$$\vartheta_n = \frac{\mathbb{P}(A_n^{(1)})}{\mathbb{P}(U_n)} = \frac{\varepsilon \mu_{\varepsilon}(U_n) m(U_n^c)}{\mu_{\varepsilon}(U_n)} = \varepsilon m(U_n^c) \xrightarrow[n \to \infty]{} \varepsilon =: \vartheta.$$

This means that the Extremal Index ϑ is given by the noise level ε .

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References

- [1] H. Aytac, J. M. Freitas, and S. Vaienti. Laws of rare events for deterministic and random dynamical systems. *Trans. Amer. Math.*, Published online with DOI: 10.1090/S0002-9947-2014-06300-9, 2014.
- [2] Bruin H and Deane J 2009 Piecewise contractions are asymptotically periodic *Proceedings of the American Mathematical Society* **137** (4) 1389-1395
- [3] Catsigeras E and Budelli R 2011 Topological dynamics of generic piecewise contracting maps in n dimensions International Journal of Pure and Applied Mathematics 68 (1) 61-83
- [4] E. Catsigeras, P. Guiraud, A. Meyroneinc and E. Ugalde, On the asymptotic properties of piecewise contracting maps, arXiv:1108.1501.
- [5] R. Coutinho, Dinâmica simbólica linear. PhD Thesis, Technical University of Lisbon, 1999.

- [6] Coutinho R, Fernandez B, Lima R and Meyroneinc A 2006 Discrete time piecewise affine models of genetic regulatory networks *Journal of Mathematical Biology* **52** (4) 524-570
- [7] D. Faranda, J.M. Freitas, P. Guiraud, S. Vaienti, Sampling local properties of attractors via extreme value theory, to appear on *Chaos, Solitons and Fractals*, http://arxiv.org/pdf/1407.0412.pdf
- [8] D. Faranda, J. Freitas, V. Lucarini, G. Turchetti, S. Vaienti, Extreme Value Statistics for Dynamical Systems with Noise, *Nonlinearity*, 26, 2597-2622, (2013)
- [9] D. Faranda, S. Vaienti, Extreme Value laws for dynamical systems under observational noise, *Physica D*, 280: 86-94, (2014).
- [10] D. Faranda, S. Vaienti, A new recurrences based technique for detecting robust extrema in long temperature records, *Geoph. Resear. Lett.* **40.21**: 5782-5786 (2013).
- [11] A. C. M. Freitas and J. M. Freitas, On the link between dependence and independence in extreme value theory for dynamical systems, *Statist. Probab. Lett.* **78** (2008), no. 9, 1088-1093.
- [12] A. C. M. Freitas, J. M. Freitas, and M. Todd, Extreme value laws in dynamical systems for non-smooth observations, *J. Stat. Phys.*, **142** (2011), 108-126.
- [13] A. C. M. Freitas, J. M. Freitas, and M. Todd, The extremal index, hitting time statistics and periodicity, Adv. Math., 231 (2012), 2626-2665.
- [14] A. C. M. Freitas, J. M. Freitas, and M. Todd. Speed of convergence for laws of rare events and escape rates. *Stochastic Process. Appl.*, Published online with DOI: 10.1016/j.spa.2014.11.011, 2014.
- [15] M. R. Leadbetter, G. Lindgren, and H. Rootzen, Extremes and related properties of random sequences and processes, Springer Series in Statistics, New York: Springer-Verlag (1983).
- [16] Lasota, A., and Mackey, M. C. (1994). Chaos, fractals, and noise: stochastic aspects of dynamics (Vol. 97). Springer.
- [17] Lima R and Ugalde E 2006 Dynamical complexity of discrete-time regulatory networks *Nonlinearity* 19 (1) 237-259
- [18] A. Nogueira, B. Pires and A. Rafael, Rosales Piecewise contractions defined by iterated function systems, arXiv:1408.1663