



HAL
open science

A many-body RAGE theorem

Jonas Lampart, Mathieu Lewin

► **To cite this version:**

| Jonas Lampart, Mathieu Lewin. A many-body RAGE theorem. 2015. hal-01121656v1

HAL Id: hal-01121656

<https://hal.science/hal-01121656v1>

Preprint submitted on 2 Mar 2015 (v1), last revised 20 Jul 2015 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

A MANY-BODY RAGE THEOREM

JONAS LAMPART AND MATHIEU LEWIN

ABSTRACT. We prove a generalized version of the RAGE theorem for N -body quantum systems. The result states that only bound states of systems with $0 \leq n \leq N$ particles persist in the long time average. The limit is formulated by means of an appropriate weak topology for many-body systems, which was introduced by the second author in a previous work, and is based on reduced density matrices. This topology is connected to the weak- $*$ topology of states on the algebras of canonical commutation or anti-commutation relations, and we give a formulation of our main result in this setting.

1. INTRODUCTION AND MAIN RESULT

The RAGE theorem, due to Ruelle [17], Amrein-Georgescu [2] and Enss [6] is a famous result relating the long time behavior of solutions to the Schrödinger equation and the spectral properties of the corresponding self-adjoint Hamiltonian. In particular, it states that for any fixed x in the ambient Hilbert space \mathfrak{H} ,

$$\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \|K P_c e^{-itH} x\|^2 dt = 0, \quad (1)$$

where P_c is the spectral projector to the continuous spectral subspace of H , K is any compact operator and $x(t) = e^{-itH} x$ is the unique (weak) solution to Schrödinger's equation

$$\begin{cases} i \dot{x}(t) = Hx(t), \\ x(0) = x. \end{cases}$$

An equivalent way of formulating the same result is as follows: for every positive self-adjoint operator γ in the trace class $\mathfrak{S}^1(\mathfrak{H})$, consider the ergodic mean

$$M(T) := \frac{1}{T} \int_0^T e^{-itH} \gamma e^{itH} dt$$

which is uniformly bounded in $\mathfrak{S}^1(\mathfrak{H})$. Then, the projection $P_c M(T) P_c$ converges weakly- $*$ to 0 as $T \rightarrow \infty$. Here, the operator $\gamma(t) = e^{-itH} \gamma e^{itH}$ is the unique (weak) solution to von Neumann's formulation of Schrödinger's equation

$$\begin{cases} i \dot{\gamma}(t) = [H, \gamma(t)], \\ \gamma(0) = \gamma. \end{cases} \quad (2)$$

The previous formulation (1) corresponds to $\gamma = |x\rangle\langle x|$. Put differently, any weakly- $*$ convergent subsequence of $M(T)$ has a limit M_∞ which is

supported on the space spanned by the eigenvectors of H . It can also be proved that M_∞ commutes with H , that is,

$$M_\infty = \sum_j \alpha_j |\varphi_j\rangle \langle \varphi_j|$$

where φ_j is an orthonormal system of eigenvectors of H .

The RAGE theorem is a very important result in quantum mechanics. For an infinite-dimensional Hamiltonian system such as Schrödinger's equation, strong convergence to stationary states (that is, eigenvectors of H) cannot hold in general, due to the conservation laws. The RAGE theorem states that, on the contrary, *weak* convergence towards stationary states holds. Hence, in this sense, only bound states persist in the long time average. By virtue of its generality, self-adjointness of H being the only hypothesis, the RAGE theorem is a fundamental tool in the spectral theory of self-adjoint operators. For instance, it is the first step towards a more precise scattering theory of quantum systems [16] and it is also often used in the study of Anderson localization [7, 12].

However, the information it provides for an interacting many-body system is often not very precise. Consider for instance three electrons in the field of a proton, described by the three-body Hamiltonian

$$H_3 = \sum_{j=1}^3 \left(-\Delta_{x_j} - \frac{1}{|x_j|} \right) + \sum_{1 \leq j < \ell \leq 3} \frac{1}{|x_j - x_\ell|}. \quad (3)$$

It is known that H_3 has no bound state [11, 15]. The RAGE theorem therefore tells us that the ergodic mean $M(T)$ tends weakly-* to zero for every initial condition γ (a trace-class operator on $L^2(\mathbb{R}^9)$). On the other hand, the corresponding Hamiltonian H_2 for two electrons has finitely many bound states [9, 10, 8] and, of course, the hydrogen atom H_1 for one electron has infinitely many. The physical picture is that some of the three particles escape, whereas the rest remain in a bound state of H_2 or H_1 , a phenomenon that is not captured at all by the RAGE theorem. The precise description of this process through scattering theory has been the object of several works [4, 18, 13]. It was proved that, asymptotically, the time-evolved wavefunction can be approximated by a sum of products of the form $e^{-it\lambda} \psi_t \otimes \varphi$ where φ is a λ -eigenfunction of H_1 or H_2 and $\psi_t \rightharpoonup 0$. A tensor product of this form tends weakly to zero in $L^2(\mathbb{R}^9)$, which is why the weak-* limit of $M(T)$ is always zero.

In this paper, we would like to prove a new general version of the RAGE theorem that retains some information on the particles that do not escape, without addressing all the details of the scattering process. This is done by using another weak topology, for which the tensor product $\psi_t \otimes \varphi$ converges to φ , and which was introduced by the second author in [14]. In this topology, the ergodic mean will converge to states with possibly less particles, which are supported in the point spectrum of the Hamiltonians H_n for $n \leq N$. This captures the principal physical ideas, even in situations where scattering is not known or not believed to hold, for instance for potentials with an arbitrarily slow decay at infinity.

Geometric convergence. In order to state our main result, let us quickly describe the notion of convergence used in [14], where all the details may be found. Let \mathfrak{H} be any separable Hilbert space and denote by \mathfrak{H}^N the symmetric (or antisymmetric) N -fold tensor product, denoted by \otimes_s , of \mathfrak{H} . Let Γ_k be a sequence of N -particle states, that is, $\Gamma_k \geq 0$ and $\text{Tr}(\Gamma_k) = 1$. For instance, for a pure state $\Gamma_k = |\Psi_k\rangle\langle\Psi_k|$ for a normalized $\Psi_k \in \mathfrak{H}^N$. The n -particle density matrix of Γ_k is obtained by taking the partial trace with respect to $N - n$ variables and multiplying by an appropriate normalization constant:

$$\Gamma_k^{(n)} = \binom{N}{n} \text{Tr}_{n+1, \dots, N}(\Gamma_k).$$

The sequence Γ_k is said to *converge geometrically* to a state Γ if the reduced density matrices of Γ_k all converge weakly- $*$ to those of Γ . Except if convergence holds in trace-norm, the state Γ can never be an N -particle state. It is necessary to work with states on the truncated Fock space

$$\mathcal{F}^{\leq N}(\mathfrak{H}) := \mathbb{C} \oplus \mathfrak{H} \oplus \dots \oplus \mathfrak{H}^N$$

and this corresponds to the picture that some particles can be lost. For simplicity, all the states we consider in this paper are assumed to commute with the particle number, the theory for the general case is essentially the same. Such states can be written in block form as $\Gamma = G_0 \oplus \dots \oplus G_N$, and if we start with a sequence of N -body states and investigate its geometric limits, these are the only states that can be obtained.

We rephrase the previous discussion in the following:

Definition 1 (Geometric convergence). *A sequence of states $\{\Gamma_k\}_{k=1}^\infty$ on \mathfrak{H}^N converges geometrically to a state $\Gamma = G_0 \oplus \dots \oplus G_N$ on $\mathcal{F}^{\leq N}(\mathfrak{H})$, if*

$$\Gamma_k^{(n)} \xrightarrow{*} \Gamma^{(n)} = G_n + \sum_{m=n+1}^N \binom{m}{n} \text{Tr}_{n+1, \dots, m}(G_m)$$

weakly- $$ in the trace-class $\mathfrak{S}^1(\mathfrak{H}^n)$ for all $n = 0, \dots, N$. That is, $\text{Tr}(K\Gamma_k^{(n)}) \rightarrow \text{Tr}(K\Gamma^{(n)})$ for every compact operator K on \mathfrak{H}^k . This notion of convergence is denoted as $\Gamma_k \xrightarrow{g} \Gamma$ and extended by linearity to sequences of states on the truncated Fock space $\mathcal{F}^{\leq N}(\mathfrak{H})$.*

Note that the geometric limit Γ is always a state, that is, satisfies $\Gamma \geq 0$ and

$$\text{Tr}_{\mathcal{F}^{\leq N}}(\Gamma) = G_0 + \sum_{n=1}^N \text{Tr}_{\mathfrak{H}^n}(G_n) = 1.$$

If convergence does not hold in $\mathfrak{S}_1(\mathfrak{H}^N)$, then the final state has to live over spaces with less particles, but its trace is always equal to one. If all the particles are lost, then Γ is the vacuum state $\Gamma = 1 \oplus 0 \oplus \dots \oplus 0$. It is proved in [14, Lemma 3] that every sequence of states Γ_k has a geometrically convergent subsequence.

Main result. We are now able to state our main result. We consider an abstract many-body Hamiltonian of the form

$$H_n = \sum_{j=1}^n h_j + \sum_{1 \leq j < \ell \leq n} w_{j\ell} \quad (4)$$

acting on the n -particle space \mathfrak{H}^n (for us symmetric or antisymmetric). Here h is a given self-adjoint operator acting on the one-particle space \mathfrak{H} and h_j acts on the j th factor. On the other hand, w is a self-adjoint operator on the two particle space \mathfrak{H}^2 and $w_{j\ell}$ acts on the j th and ℓ th factors. We make rather general assumptions on w in order to give a proper meaning to H_n for all n . As will be clear from the rest of the paper, these can be weakened in specific examples or if one is only interested in a particular $n = N$. We assume that

$$h \text{ is bounded from below (without loss of generality } h > 1) \quad (5)$$

and that

$$|w| \text{ is infinitesimally form bounded with respect to } h_1 + h_2 \quad (6)$$

which means that

$$\varepsilon(h_1 + h_2) - C_\varepsilon \leq w \leq \varepsilon(h_1 + h_2) + C_\varepsilon \quad (7)$$

for all $\varepsilon > 0$. Under these assumptions, H_n is bounded from below and may be realized as a self-adjoint operator for all n , by the method of Friedrichs.

Theorem 1 (Many-body RAGE). *In addition to (5) and (6), assume that for every compact operator K on \mathfrak{H}*

$$K \otimes_s 1(h_1 + h_2)^{-1/2} w (h_1 + h_2)^{-1/2} \quad (8)$$

is compact on \mathfrak{H}^2 . Let $\Gamma \geq 0$ with $\text{Tr}(\Gamma) = 1$ be a state on the (symmetric or antisymmetric) N -particle space \mathfrak{H}^N . Then, the ergodic mean

$$\frac{1}{T} \int_0^T e^{-itH_N} \Gamma e^{itH_N} dt$$

has geometrically convergent subsequences as $T \rightarrow \infty$, and, for every such sequence, the limit is a convex combination of projections to eigenspaces of the n -body Hamiltonians H_n for $0 \leq n \leq N$.

If we consider a sequence of times $T_k \rightarrow \infty$ for which

$$\frac{1}{T_k} \int_0^{T_k} e^{-itH_N} \Gamma e^{itH_N} dt \xrightarrow{\text{g}} M_\infty = G_0 \oplus \cdots \oplus G_N,$$

the result states that

$$G_n = \sum_j \alpha_{n,j} |\varphi_{n,j}\rangle \langle \varphi_{n,j}| \quad (9)$$

where $\{\varphi_{n,j}\}_{j \geq 1}$ is an orthonormal system of eigenvectors of the n -body operator H_n . By definition, geometric convergence means that the density matrices converge

$$\frac{1}{T_k} \int_0^{T_k} \left(e^{-itH_N} \Gamma e^{itH_N} \right)^{(n)} dt \xrightarrow{*} M_\infty^{(n)} = \sum_{m=n}^N \binom{m}{n} \text{Tr}_{n+1, \dots, m} G_m$$

weakly-* in the trace-class for all $n = 0, \dots, N$. Let us emphasize that, because of the partial traces, the $M_\infty^{(n)}$ are in general not supported on the point spectrum of H_n . The density matrix $M_\infty^{(n)}$ should not be confused with the restriction G_n of the state M_∞ to the n -particle space.

The condition (8) on w ensures that the interaction between a particle that stays and a particle that escapes vanishes in a weak sense. A condition of this type is clearly necessary to be sure that the remaining n particles are described by the Hamiltonian H_n when the other $N - n$ escape. The operator $B := (h_1 + h_2)^{-1/2} w (h_1 + h_2)^{-1/2}$ is always bounded by Assumption (6). However, in the applications, it is usually not compact. The picture is that w only decays in one direction and it is necessary to multiply by one compact operator K in another variable as in (8) to make $K \otimes 1B$ and $1 \otimes KB$ compact.

Our theorem is stated for any initial datum Γ , possibly with an infinite energy. However, our proof does use the conservation of energy for smooth initial data. It is an interesting open problem to derive a similar result when h is not bounded from below. We will make more comments on this below.

We also remark that the exact same theorem holds if the ergodic mean is replaced by

$$\frac{1}{T} \int_{\mathbb{R}} \chi(t/T) e^{-iH_N t} \Gamma e^{iH_N t} dt$$

where χ is any nonnegative function such that $\int_{\mathbb{R}} \chi(t) dt = 1$.

Reformulation in terms of the CAR and CCR algebras. Theorem 1 has a natural extension to the Fock space and the associated algebra of canonical (anti-) commutation relations (see [5] for a detailed introduction). This extension is based on the observation (see [14, Remark 2.1]) that for the rank-one operator

$$K = |f_1 \otimes_s \dots \otimes_s f_n\rangle \langle g_1 \otimes_s \dots \otimes_s g_n|, \quad (10)$$

geometric convergence of $\Gamma_k \in \mathfrak{S}_1(\mathcal{F}^{\leq N})$ just means that

$$\begin{aligned} \text{Tr} \left(K \Gamma_k^{(n)} \right) &= \left\langle g_1 \otimes_s \dots \otimes_s g_n, \Gamma_k^{(n)} f_1 \otimes_s \dots \otimes_s f_n \right\rangle \\ &= \text{Tr}_{\mathcal{F}} \left(a^*(f_1) \dots a^*(f_n) a(g_1) \dots a(g_n) \Gamma_k \right) \end{aligned}$$

converges. Here, $a^*(f), a(f)$ denote, respectively, bosonic or fermionic creation and annihilation operators, depending on whether \otimes_s is the symmetric or antisymmetric tensor product. We thus have:

Corollary 2. *Let $\mathcal{A} \in \{\text{CCR}(\mathfrak{H}), \text{CAR}(\mathfrak{H})\}$ be the C^* -subalgebra of the bounded operators of $\mathcal{F}(\mathfrak{H})$ satisfying canonical commutation relations, if $\mathcal{F}(\mathfrak{H})$ is the symmetric Fock-space, or anti-commutation relations, if $\mathcal{F}(\mathfrak{H})$ is the anti-symmetric Fock-space. Denote by Π_N the projection of $\mathcal{F}(\mathfrak{H})$ to \mathfrak{H}^N and let $\rho \in \mathcal{A}'$ be a normal state on \mathcal{A} , that is $\rho(A) = \text{Tr}(\Gamma_\rho A)$ for some $\Gamma_\rho \in \mathfrak{S}_1(\mathcal{F})$, satisfying $[\Gamma_\rho, \Pi_N] = 0$ for all $N \in \mathbb{N}$. Then the ergodic mean*

$$\mu_T(A) := \frac{1}{T} \int_0^T \rho \left(e^{i\mathbb{H}t} A e^{-i\mathbb{H}t} \right) dt.$$

has weakly- convergent subsequences as $T \rightarrow \infty$ and, for every such sequence, the limit state is normal. For the density operator Γ_∞ of this*

state, $\Pi_N \Gamma_\infty \Pi_N$ is a convex combination of projections to eigenspaces of H_N for all $N \geq 1$.

Proof. The existence of weakly- $*$ convergent subsequences T_k follows from the fact that \mathcal{A} is separable and the Banach-Alaoglu theorem. By a diagonal argument we may additionally assume that

$$\mu_{T_k}^N(A) := \mu_{T_k}(\Pi_N A \Pi_N) \rightarrow \mu_\infty^N(A)$$

for some (unnormalized) state $\mu_\infty^N \in \mathcal{A}'$ and every $N \in \mathbb{N}$.

The action of $\Pi_N A \Pi_N$ on \mathfrak{H}^N is completely determined by the action of operators of the form (10) with $n \leq N$: For the fermionic case this is an immediate consequence of the commutation relations and the fact that $K \Pi_N = 0$ if $n > N$; For the bosonic case one uses additionally that on \mathfrak{H}^N the Weyl-operators may be expressed as convergent power series in creation and annihilation operators (see [5, chapter 9]). Thus $\Pi_N A \Pi_N$ is generated by operators of the form $K \otimes_s 1$ with compact operators K on \mathfrak{H}^n , $n \leq N$. This shows that weak- $*$ convergence of $\mu_{T_k}^N$ is equivalent to geometric convergence of the sequence $\Pi_N M(T_k) \Pi_N$ of ergodic means with $M(0) = \Gamma_\rho$. Thus, the geometric limit satisfies $\mu_\infty^N(A) = \text{Tr}(AM_\infty^N)$ for some $M_\infty^N \in \mathfrak{S}_1(\mathcal{F}^{\leq N})$. Using that $\mu_T([\Pi_N, A]) = 0$ by hypothesis, and that $\|\mu_{T_k}^N\|_{\mathcal{A}'} = \|\Pi_N \Gamma_\rho \Pi_N\|_{\mathfrak{S}_1}$ is independent of T_k , we obtain

$$\mu_\infty(A) := \lim_{k \rightarrow \infty} \mu_{T_k}(A) = \lim_{k \rightarrow \infty} \sum_{N=0}^{\infty} \mu_{T_k}^N(A) = \sum_{N=0}^{\infty} \text{Tr}_{\mathcal{F}}(AM_\infty^N).$$

Consequently, $\Gamma_\infty = \sum_{N=0}^{\infty} M_\infty^N$ and, by Theorem 1, $\Pi_n M_\infty^N \Pi_n$ is a convex combination of projections to eigenspaces of H_n . \square

Application to Schrödinger operators. A typical example to which our result applies is that of a non-relativistic system of N fermions or bosons, for which the interaction w is a function of the relative position of the particles, that tends to zero at infinity in a weak sense.

Corollary 3 (Non-relativistic Schrödinger operators). *Let $\mathfrak{H} = L^2(\mathbb{R}^d)$, $h = -\Delta + V(x) + e$ and w be the multiplication operator by an even function $w(x_1 - x_2)$. We assume that $V = f_1 + f_2$ and that $w = f_3 + f_4$ where $f_i \in L^{p_i}(\mathbb{R}^d)$ for some $\max(1, d/2) < p_i < \infty$, or $f_i \in L^\infty(\mathbb{R}^d)$ and $f_i \rightarrow 0$ at infinity. Then Theorem 1 holds for the many-body Hamiltonian*

$$H_N = \sum_{j=1}^N -\Delta_{x_j} + V(x_j) + \sum_{1 \leq j < \ell \leq N} w(x_j - x_\ell).$$

Proof. Under the assumptions of the statement, $(1 - \Delta)^{-1/2} V(x) (1 - \Delta)^{-1/2}$ and $(1 - \Delta)^{-1/2} w(x) (1 - \Delta)^{-1/2}$ are compact on $L^2(\mathbb{R}^d)$ [3, Chap 8], hence the hypothesis (5), (6) are satisfied and $h > 1$ for an appropriate choice of e .

The verification of (8) is more involved, the intuition however is rather simple. It relies on the fact that $(h_1 + h_2)^{-1/2} w (h_1 + h_2)^{-1/2}$ is compact in the relative coordinate $x_1 - x_2$ parametrizing the subspace $V = \{(x_1, x_2) \in \mathbb{R}^{2d} : x_1 + x_2 = 0\}$, i.e. its action on the space $L^2(V)$ is compact. This, together with the fact that $K \otimes 1$ is compact in the direction $\mathbb{R}^d \times \{0\}$

parametrized by x_1 and $(\mathbb{R}^d \times \{0\}) \oplus V = \mathbb{R}^{2d}$, implies that the product is compact. More precisely, the property we use is the content of the following lemma, which we prove in Appendix A.

Lemma 4. *There exist compact operators K_j on $L^2(V)$ and bounded operators B_j on $L^2(V^\perp)$ such that*

$$(h_1 + h_2)^{-1/2} w (h_1 + h_2)^{-1/2} = \sum_{j=1}^{\infty} K_j \otimes_V B_j \quad (11)$$

and the sum converges in the operator-norm.

The notation \otimes_V emphasizes that the tensor product of operators is induced by $L^2(\mathbb{R}^{2d}) = L^2(V) \otimes L^2(V^\perp)$. The property (11) is known as V -semicompactness and implies hypothesis (8) by [1, Proposition 9.2.2].

The argument is as follows: It is clearly sufficient to prove that $(K \otimes 1)(K_j \otimes_V 1)$ is compact for every j . Then, since Hilbert-Schmidt operators are dense in the compact operators, it suffices to show that $C := (A \otimes 1)(B \otimes_V 1)$ is compact (and actually Hilbert-Schmidt) for operators $A \in \mathfrak{S}_2(L^2(\mathbb{R}^d))$, $B \in \mathfrak{S}_2(L^2(V))$. To prove this, let $a(x, x')$, $b(v, v')$ be the integral kernels of A and B , respectively. The operator C acts on $\psi \in L^2(\mathbb{R}_{x_1}^d \times \mathbb{R}_{x_2}^d)$ as

$$(C\psi)(x_1, x_2) = \int_{\mathbb{R}^{2d}} a(x_1, x') b(x' - x_2, x' - y') \psi(x', y') dx' dy'.$$

So C is an integral operator, whose kernel is easily seen to be in $L^2(\mathbb{R}^{4d})$. \square

If w and V decay fast enough at infinity, the existence of scattering [16, 4, 18, 13] means that there exists vectors $\psi_{n,j}$ such that

$$\lim_{t \rightarrow +\infty} \left\| e^{-itH_N} \Psi - \sum_{n=0}^N \sum_j e^{-i\lambda_{n,j}t} (e^{it\Delta_{\mathbb{R}^d(N-n)}} \psi_{n,j}) \otimes \varphi_{n,j} \right\|_{L^2(\mathbb{R}^{dN})} = 0 \quad (12)$$

where $H_n \varphi_{n,j} = \lambda_{n,j} \varphi_{n,j}$. Furthermore, asymptotic completeness states that one can prepare the initial datum Ψ such as to get any chosen $\psi_{n,j}$ in the large time limit. The convergence (12) implies that

$$e^{-itH_N} |\Psi\rangle \langle \Psi| e^{itH_N} \xrightarrow{g} \bigotimes_{n=0}^N \left(\sum_j \|\psi_{n,j}\|^2 |\varphi_{n,j}\rangle \langle \varphi_{n,j}| \right) \quad (13)$$

geometrically. Hence, in this case, the coefficients in (9) are $\alpha_{n,j} = \|\psi_{n,j}\|^2$. Asymptotic completeness therefore tells us that any combination G_0, \dots, G_N of the eigenvectors of the n -particles Hamiltonians H_n may be obtained in the geometric limit. We do not know if the same property is true under the more general assumptions of Theorem 1.

Results similar to Corollary 3 hold for Schrödinger operators with magnetic fields, for pseudo-relativistic operators, etc. We do not state them here for shortness.

The rest of the paper will be devoted to the proof of Theorem 1.

2. PROOF OF THEOREM 1

Let Γ be a state on the N -particle space \mathfrak{H}^N and $T_k \rightarrow \infty$ be a sequence of times such that the ergodic means $M(T_k)$, with initial condition $M(0) = \Gamma$, converge geometrically to a limit $M_\infty = \bigoplus_{n=0}^N G_n$. Such a sequence exists for every Γ by [14, Lemma 3].

The proof of Theorem 1 proceeds by showing that M_∞ is left invariant by the Hamiltonian $\mathbb{H} := \bigoplus_{n=0}^N H_n$ (with $H_0 = 0$ and $H_1 = h$) on the truncated Fock space $\mathcal{F}^{\leq N} = \mathbb{C} \bigoplus_{n=1}^N \mathfrak{H}^n$. That is, we have $M_\infty = e^{-i\mathbb{H}t} M_\infty e^{i\mathbb{H}t}$ and thus $e^{-itH_n} G_n e^{itH_n} \equiv G_n$, for every $t \in \mathbb{R}$. Consequently, the eigenspaces of G_n are H_n -invariant. For a non-zero eigenvalue of G_n , the eigenspace has finite dimension and is thus a direct sum of eigenspaces of H_n .

Step 1. The first step to proving invariance of M_∞ is to note that, for every $s > 0$, the sequence $e^{-isH_N} M(T_k) e^{isH_N}$ also converges to M_∞ geometrically. This holds because geometric convergence is controlled by the trace-norm and

$$\begin{aligned} e^{-isH_N} M(T_k) e^{isH_N} &= \frac{1}{T_k} \int_0^{T_k} e^{-i(t+s)H_N} \Gamma e^{i(t+s)H_N} dt \\ &= \frac{1}{T_k} \int_s^{T_k+s} e^{-itH_N} \Gamma e^{itH_N} dt \\ &= M(T_k) + \underbrace{\frac{1}{T_k} \left(\int_{T_k}^{T_k+s} e^{-itH_N} \Gamma e^{itH_N} dt - \int_0^s e^{-itH_N} \Gamma e^{itH_N} dt \right)}_{\leq 2s/T_k \text{ in } \mathfrak{S}_1} \\ &\xrightarrow{\mathfrak{g}} M_\infty. \end{aligned}$$

Having established this, we would like to prove that also

$$e^{-isH_N} M(T_k) e^{isH_N} = e^{-is\mathbb{H}} M(T_k) e^{is\mathbb{H}} \xrightarrow{\mathfrak{g}} e^{-is\mathbb{H}} M_\infty e^{is\mathbb{H}}. \quad (14)$$

As the left hand side depends only on H_N and the right may depend on \mathbb{H} on all the sectors with $n \leq N$ particles, this will certainly not be true for arbitrary Hamiltonians on Fock space. In fact, the proof of (14) will depend crucially both on the properties of \mathbb{H} and the sequence $M(T_k)$.

Step 2. Using the particular form of the sequence $M(T_k)$ we can reduce the proof of (14) to sequences of bounded energy. That is, let e_N be such that $H_N + e_N > 1$ and assume the initial condition $M(0) = \Gamma$ satisfies

$$\mathrm{Tr} \left((H_N + e_N)^{1/2} \Gamma (H_N + e_N)^{1/2} \right) < E$$

for some constant $E > 0$. Then, since e^{-itH_N} commutes with H_N and preserves the trace-norm, we have

$$\mathrm{Tr} \left((H_N + e_N)^{1/2} M(T_k) (H_N + e_N)^{1/2} \right) < E, \quad (15)$$

for every $k \in \mathbb{N}$. Since any initial condition $M(0) = \Gamma$ can be approximated to arbitrary precision by states of finite energy, it is sufficient to prove (14) for sequences satisfying (15), by virtue of the following lemma.

Lemma 5. *There exists a constant c , depending only on N , such that for every sequence $\{\gamma_k\}_{k=1}^\infty$ in $\mathfrak{S}_1(\mathfrak{H}^N)$ that converges geometrically to $\gamma \in \mathfrak{S}_1(\mathcal{F}^{\leq N})$ and satisfies $\|\gamma_k\|_{\mathfrak{S}_1(\mathfrak{H}^N)} < \varepsilon$ we have $\|\gamma\|_{\mathfrak{S}_1(\mathcal{F}^{\leq N})} < c\varepsilon$.*

Proof. Since the trace-class is the dual of the compact operators, we have

$$\begin{aligned} \|\gamma^{(n)}\|_{\mathfrak{S}_1(\mathfrak{H}^n)} &= \sup_{K \text{ compact}, \|K\|=1} \lim_{k \rightarrow \infty} \left| \text{Tr}_{\mathfrak{H}^n} \left(K \gamma_k^{(n)} \right) \right| \\ &= \binom{N}{n} \sup_{\|K\|=1} \lim_{k \rightarrow \infty} |\text{Tr}_{\mathfrak{H}^N} (K \otimes 1_{N-n} \gamma_k)| < \binom{N}{n} \varepsilon. \end{aligned}$$

Now $\gamma = \bigoplus_{n=0}^N G_n$ is completely determined by its reduced density matrices, explicitly (cf. [14, Lemma 1.1]),

$$G_n = \gamma^{(n)} + \sum_{j=1}^{N-n} (-1)^j \binom{n+j}{n} \text{Tr}_{n+1, \dots, n+j} \gamma^{(n+j)},$$

which proves the claim. \square

It will often be useful to state the energy bound using a Hamiltonian without interaction. Let

$$H_n^0 := \sum_{j=1}^n h_j > 1$$

and

$$\mathfrak{S}_{1, H^0}(\mathfrak{H}^n) := \left\{ \gamma \in \mathfrak{S}_1(\mathfrak{H}^n) : \|\sqrt{H_n^0} \gamma \sqrt{H_n^0}\|_{\mathfrak{S}_1} < \infty \right\}.$$

This space has a natural norm given by $\|\gamma\| := \|\gamma\|_{\mathfrak{S}_1} + \|\sqrt{H_n^0} \gamma \sqrt{H_n^0}\|_{\mathfrak{S}_1}$. Using H_n instead of H_n^0 gives an equivalent norm, because of the inequalities (7). We also define

$$\mathfrak{S}_{1, \mathbb{H}^0}(\mathcal{F}^{\leq N}) := \left\{ \gamma \in \mathfrak{S}_1(\mathcal{F}^{\leq N}) : \|\sqrt{\mathbb{H}^0} \gamma \sqrt{\mathbb{H}^0}\|_{\mathfrak{S}_1} < \infty \right\},$$

with $\mathbb{H}^0 = \bigoplus_{n=0}^N H_n^0$, $H_0^0 = 0$. It will be important that geometric convergence preserves such energy estimates. The following lemma proves that the unit ball of $\mathfrak{S}_{1, \mathbb{H}^0}(\mathcal{F}^{\leq N})$ is closed under geometric convergence. This is the only step in the proof of Theorem 1 for which positivity of H_n^0 is essential.

Lemma 6. *Let $\{\gamma_k\}_{k=1}^\infty$ in $\mathfrak{S}_{1, \mathbb{H}^0}(\mathcal{F}^{\leq N})$ be a bounded sequence that converges geometrically to γ . Then for every $n \leq N$*

$$\sqrt{H_n^0} \gamma_k^{(n)} \sqrt{H_n^0} \rightharpoonup_* \sqrt{H_n^0} \gamma^{(n)} \sqrt{H_n^0}$$

in $\mathfrak{S}_1(\mathfrak{H}^n)$.

Proof. By linearity, it is sufficient to prove the claim for $\gamma_k \in \mathfrak{S}_{1, H^0}(\mathfrak{H}^N)$. For any $n \leq N$ we have

$$0 \leq (H_n^0 \otimes 1_{\mathfrak{H}^{N-n}})(H_N^0)^{-1} = 1 - (1_{\mathfrak{H}^n} \otimes H_{N-n}^0)(H_N^0)^{-1} \leq 1,$$

because $H_n^0 \geq 1$. Thus, for an arbitrary $\eta \in \mathfrak{S}_{1, H^0}(\mathfrak{H}^N)$,

$$\begin{aligned} \left\| \sqrt{H_n^0} \eta^{(n)} \sqrt{H_n^0} \right\|_{\mathfrak{S}_1} &= \sup_{B \text{ bounded, } \|B\|=1} \left| \text{Tr} \left(B \sqrt{H_n^0} \eta^{(n)} \sqrt{H_n^0} \right) \right| \\ &\leq \binom{N}{n} \left\| (\sqrt{H_n^0} \otimes 1) \eta (\sqrt{H_n^0} \otimes 1) \right\|_{\mathfrak{S}_1} \\ &\leq \binom{N}{n} \left\| \sqrt{H_n^0} \eta \sqrt{H_n^0} \right\|_{\mathfrak{S}_1}. \end{aligned} \quad (16)$$

So $\sqrt{H_n^0} \gamma_k^{(n)} \sqrt{H_n^0}$ is uniformly bounded in $\mathfrak{S}_1(\mathfrak{H}^n)$ and there exists a subsequence such that

$$\sqrt{H_n^0} \gamma_k^{(n)} \sqrt{H_n^0} \rightharpoonup_* \tilde{\gamma}^{(n)}.$$

Now for every compact operator K on \mathfrak{H}^n we have

$$\begin{aligned} \text{Tr} \left((H_n^0)^{-1/2} K (H_n^0)^{-1/2} \tilde{\gamma}^{(n)} \right) &= \lim_{k \rightarrow \infty} \text{Tr} \left((H_n^0)^{-1/2} K \gamma_k^{(n)} (H_n^0)^{1/2} \right) \\ &= \text{Tr} \left(K \gamma^{(n)} \right), \end{aligned}$$

whence $\tilde{\gamma}^{(n)} = \sqrt{H_n^0} \gamma^{(n)} \sqrt{H_n^0}$. \square

Step 3. The reduction of the problem to sequences of bounded energy in step 2 will now allow us to study the sequence $e^{-isH_N} M(T_k) e^{isH_N}$ via the differential equation it satisfies. If $\gamma(t)$ is a solution to the von Neumann equation (2), its reduced density matrices (formally) satisfy the finite BBGKY hierarchy

$$i \frac{d}{dt} \gamma^{(n)}(t) = [H_n, \gamma^{(n)}(t)] + (n+1) \sum_{j=1}^n \text{Tr}_{n+1} \left([w_{j, n+1}, \gamma^{(n+1)}(t)] \right).$$

Note, however, that the equation has no clear meaning (not even in a weak sense) if w is not bounded, due to the partial trace. If w is a bounded operator $D((H_2^0)^\alpha) \rightarrow \mathfrak{H}^2$ for some $\alpha > 0$, this problem can be handled by considering only initial conditions satisfying $\text{Tr}((H_N^0)^\alpha \gamma (H_N^0)^\alpha) < E$, but to deal with potentials that are only H_2^0 -form-bounded we will have to define a modified equation, that is equivalent to the original one for bounded w .

Proposition 7 (Well-posedness of the truncated BBGKY hierarchy).

- (1) For every $\gamma \in \mathfrak{S}_{1, \mathbb{H}^0}(\mathcal{F}^{\leq N})$ the family of reduced density matrices $\gamma(t)^{(n)} = (e^{-it\mathbb{H}} \gamma e^{it\mathbb{H}})^{(n)}$ is the unique solution $\{\gamma^{(n)}(t) : n = 0, \dots, N\}$ to the the system of equations

$$\begin{aligned} \gamma^{(n)}(t) &= e^{-itH_n} \gamma^{(n)} e^{itH_n} \\ &\quad - i(n+1) \sum_{j=1}^n \int_0^t e^{-i(t-s)H_n} \text{Tr}_{n+1} \left(\mathcal{L}_{jn} \left(w, \gamma^{(n+1)}(s) \right) \right) e^{i(t-s)H_n} ds, \\ \mathcal{L}_{jn} \left(w, \gamma^{(n+1)} \right) &= \left[h_{n+1}^{-1/2} w_{j, n+1} h_{n+1}^{-1/2}, h_{n+1}^{1/2} \gamma^{(n+1)}(s) h_{n+1}^{1/2} \right], \end{aligned} \quad (17)$$

such that

$$\gamma^{(n)}(t) \in L^\infty(\mathbb{R}, \mathfrak{S}_{1, H_n^0}(\mathfrak{H}^n)).$$

- (2) For every bounded sequence $\gamma_k \in \mathfrak{S}_{1, \mathbb{H}^0}(\mathcal{F}^{\leq N})$ that converges geometrically to γ_∞ , the corresponding solutions $\gamma_k(t)$ of (17) with initial condition γ_k converge geometrically,

$$\forall t \in \mathbb{R} : \quad \gamma_k(t) \xrightarrow{g} \gamma_\infty(t).$$

Proof. Note that $\gamma^{(N+1)} = 0$ and this is a triangular system, which can be solved starting with $\gamma^{(N)}(t) = e^{-itH_N} \gamma^{(N)} e^{itH_N}$. This immediately gives us uniqueness, for assume we have two solutions with $\gamma^{(n)}(0) = \gamma^{(n)} = \theta^{(n)}(0)$, then $\gamma^{(N)}(t) = \theta^{(N)}(t)$. If $\gamma^{(N)}, \theta^{(N)} \in L^\infty(\mathbb{R}, \mathfrak{S}_{1, H_n^0})$ this then implies $\gamma^{(N-1)}(t) = \theta^{(N-1)}(t)$.

We now check that $\gamma(t)^{(n)} = (e^{-it\mathbb{H}} \gamma e^{it\mathbb{H}})^{(n)}$ is indeed a solution. By linearity, we may restrict to initial conditions in $\gamma \in \mathfrak{S}_{1, H^0}(\mathfrak{H}^N)$. Since both $(H_N + e_N)^{1/2} (H_N^0)^{-1/2}$ and $(H_N^0)^{1/2} (H_N + e_N)^{-1/2}$ are bounded, due to hypothesis (6), we have $\gamma(t) = \gamma(t)^{(N)} \in L^\infty(\mathbb{R}, \mathfrak{S}_{1, H^0}(\mathfrak{H}^N))$. Equation (16) then implies that $\gamma(t)^{(n)} \in L^\infty(\mathbb{R}, \mathfrak{S}_{1, H^0}(\mathfrak{H}^n))$.

Now assume for a moment that w is bounded and let

$$W_n := \sum_{1 \leq j \leq n < \ell \leq N} w_{j\ell}$$

be the interaction of the first n particles with the remaining $N - n$. Then we have

$$H_N = H_n \otimes 1_{\mathfrak{H}^{N-n}} + 1_{\mathfrak{H}^n} \otimes H_{N-n} + W_n,$$

and we can write $\gamma(t)^{(N)}$ using Duhamel's formula,

$$\begin{aligned} \gamma(t)^{(N)} &= e^{-itH_{N-n}} e^{-itH_n} \gamma e^{itH_n} e^{itH_{N-n}} \\ &\quad - i \int_0^t e^{-i(t-s)H_{N-n}} e^{-i(t-s)H_n} \left[W_n, \gamma(s)^{(N)} \right] e^{i(t-s)H_{N-n}} e^{i(t-s)H_n} ds \end{aligned}$$

(with the $1 \otimes$ omitted for shortness). This implies that

$$\begin{aligned} &\gamma(t)^{(n)} - e^{-itH_n} \gamma^{(n)} e^{itH_n} \\ &= -i \binom{N}{n} \int_0^t e^{-i(t-s)H_n} \text{Tr}_{n+1, \dots, N} \left(\left[W_n, \gamma(s)^{(N)} \right] \right) e^{i(t-s)H_n} ds \\ &= -i(n+1) \int_0^t \sum_{j=1}^n e^{-i(t-s)H_n} \text{Tr}_{n+1} \left(\left[w_{j, n+1}, \gamma(s)^{(n+1)} \right] \right) e^{i(t-s)H_n} ds. \end{aligned}$$

This gives (17) as $h_{n+1}^{1/2} \gamma^{(n+1)} h_{n+1}^{1/2}$ is trace-class. In order to account for unbounded interactions, let $w^\delta := (\delta H_2^0 + 1)^{-1/2} w (\delta H_2^0 + 1)^{-1/2}$ for $\delta > 0$, which is bounded by hypothesis (6). We clearly have

$$\lim_{\delta \rightarrow 0} (H_2^0)^{-1/2} w^\delta (H_2^0)^{-1/2} = (H_2^0)^{-1/2} w (H_2^0)^{-1/2}$$

in the strong topology of operators on \mathfrak{H}^2 . Hence, the constant C_ε of equation (7) may be chosen in such a way that the inequalities hold for all $0 < \delta \leq 1$. Then, for an appropriate choice of $e_N \geq 0$,

$$\left(H_N^0 + \sum_{1 \leq j < \ell \leq N} w_{j\ell}^\delta + e_N \right)^{-1/2} : \mathfrak{H}^N \rightarrow D((H_N^0)^{1/2})$$

is bounded, uniformly in δ . This implies that, as operators on $D((H_N^0)^{1/2})$,

$$\lim_{\delta \rightarrow 0} \left(H_N^0 + \sum_{1 \leq j < \ell \leq N} w_{j\ell}^\delta + e_N \right)^{-1} = (H_N + e_N)^{-1}$$

strongly. This in turn implies that the unitary groups U_δ generated by these operators also converge in the strong operator topology on $D((H_N^0)^{1/2})$. So, for $\gamma \in \mathfrak{S}_{1, H^0}$, we have

$$\lim_{\delta \rightarrow 0} (H_N^0)^{1/2} U_\delta(t) \gamma = (H_N^0)^{1/2} e^{-itH_N} \gamma$$

in trace-norm. We thus have, using (16), that

$$\lim_{\delta \rightarrow 0} (U_\delta(t) \gamma U_\delta(t)^*)^{(n)} = \gamma^{(n)}(t)$$

in $\mathfrak{S}_{1, H^0}(\mathfrak{H}^n)$. As w^δ is bounded, the left hand side solves (17) with this interaction. To take the limit on the right hand side of (17) observe that

$$\lim_{\delta \rightarrow 0} \text{Tr}_{n+1} \mathcal{L}_{jn} \left((w^\delta, \gamma^{(n+1)}) \right) = \text{Tr}_{n+1} \mathcal{L}_{jn} \left((w, \gamma^{(n+1)}) \right)$$

in the space of operators for which $(H_n^0)^{-1/2} A (H_n^0)^{-1/2}$ is trace-class. As \mathfrak{S}_{1, H^0} is obviously contained in this space, this shows that $\gamma(t)^{(n)}$ is indeed a solution.

We now prove the continuous dependence on the initial condition, item (2). Since we have already proved uniqueness, it is sufficient to show that the geometric limit of $\gamma_k(t) = e^{-it\mathbb{H}} \gamma_k e^{it\mathbb{H}}$ is a solution of (17). On the space \mathfrak{H}^N this is trivial, since for any compact operator K , $K(t) := e^{itH_N} K e^{-itH_N}$ is also compact and thus

$$\lim_{k \rightarrow \infty} \text{Tr} \left(K e^{-itH_N} \gamma_k e^{itH_N} \right) = \text{Tr} \left(K e^{-itH_N} \gamma_\infty e^{itH_N} \right),$$

for every $t \in \mathbb{R}$. Now let $m < N$ and assume that $\gamma_k(t)^{(n)} \rightarrow_* \gamma_\infty(t)^{(n)}$ in $\mathfrak{S}_1(\mathfrak{H}^n)$ for every $t \in \mathbb{R}$ and $n = m+1, \dots, N$. Then, we have for every compact operator K on \mathfrak{H}^m

$$\begin{aligned} & \frac{i}{m+1} \left(\text{Tr}_{\mathfrak{H}^m} \left(K \gamma_k(t)^{(m)} \right) - \text{Tr}_{\mathfrak{H}^m} \left(e^{-itH_m} K \gamma_k(0)^{(m)} e^{itH_m} \right) \right) \\ &= \int_0^t \sum_{j=1}^m \text{Tr}_{\mathfrak{H}^{m+1}} \left((K(t-s) \otimes_s 1) \mathcal{L}_{jm}(w, \gamma^{(m+1)}) \right) ds. \end{aligned}$$

It is clearly sufficient to prove convergence for K in a dense set of compact operators, so we may assume that $K(t-s) = \tilde{K}(t-s)(H_m)^{-1/2}$ with compact \tilde{K} . For such an operator K , the integrand is uniformly bounded, and we will show that it converges pointwise. Using Lemma 6, the induction hypothesis gives us

$$(H_{m+1}^0)^{1/2} \gamma_k^{(m+1)}(s) (H_{m+1}^0)^{1/2} \rightarrow_* (H_{m+1}^0)^{1/2} \gamma_\infty^{(m+1)}(s) (H_{m+1}^0)^{1/2}.$$

It is thus enough to prove that

$$(\tilde{K}_j(t-s) \otimes_s 1) (h_{m+1} h_j)^{-1/2} w_{j, m+1} (H_{m+1}^0)^{-1/2},$$

where $\tilde{K}_j(t-s) = \tilde{K}(t-s)(H_m^0)^{-1/2}h_j^{1/2}$, is compact. Denote by $S : \mathfrak{H}^{\otimes m+1} \rightarrow \mathfrak{H}^{m+1}$ the symmetrization operator. Hypothesis (8) guarantees that

$$S(L \otimes 1)(h_{m+1} + h_j)^{-1/2}w_{j,m+1}(H_{m+1}^0)^{-1/2}$$

is compact for any compact operator L . Hence, it is sufficient to show that

$$(\tilde{K}_j(t-s) \otimes 1)(h_{m+1}h_j)^{-1/2}(h_j + h_{m+1})^{1/2} = \sum_{l=1}^{\infty} L_l \otimes B_l,$$

as a norm-convergent sum, with bounded operators B_l on \mathfrak{H} and compact operators L_l on \mathfrak{H}^m . Since h_j and h_{m+1} commute, we have

$$((h_{m+1} + h_j)/h_j h_{m+1})^{-1/2} = \left(1 + h_j^{-1}\right)^{-1/2} \left(1 - \frac{1 - h_m^{-1}}{1 + h_j^{-1}}\right)^{-1/2},$$

and because $0 < (1 - h_m^{-1})/(1 + h_j^{-1}) < 1$ we can write this using the convergent power series of $(1-x)^{-1/2}$ on $|x| < 1$, which gives the desired form.

We have thus shown that $\gamma_k(t)^{(m)}$ converges weakly-* to the right hand side of equation (17) with $\gamma^{(m+j)}(t) = (e^{-it\mathbb{H}}\gamma_{\infty}e^{it\mathbb{H}})^{(m+j)}$, which proves the claim by the uniqueness of solutions to (17). \square

This completes the proof of Theorem 1. To summarize: We noted that the Theorem is implied by invariance of the limit state M_{∞} under H . As a first step we then showed that this holds if the limit of the time-evolved sequence is the evolved limit state (equation (14)). In the second step we reduced the problem to considering initial conditions with bounded energy, and thus to proving (14) for sequences of uniformly bounded energy, i.e. satisfying (15). The final step consisted in studying the BBGKY-hierarchy satisfied by the reduced density matrices under the condition of finite energy, in its weak form (17). We proved uniqueness of solutions to this equation in Proposition 7. Additionally, we proved that that geometric convergence for the initial conditions, together with a uniform energy bound, implies geometric convergence of the corresponding solutions at any finite time. This gives (14) for the sequence $M(T_k)$ of ergodic means and thus Theorem 1.

Remark 1. *Instead of studying the evolution equation for the density matrices as in step 3, one could also study the dual evolution on the space of operators that we test against. The condition that needs to be verified is that this space is invariant under conjugation by $e^{-it\mathbb{H}}$. Geometric convergence of γ_k is defined via convergence of $\text{Tr}(K \otimes 1\gamma_k)$, so the natural space of operators would be $K \otimes 1$, where K is compact on \mathfrak{H}^n for some $n \leq N$. However, this space will in general not be invariant under $e^{-it\mathbb{H}}$. With a bound on the energy of the sequence γ_k one can instead consider operators of the form $A \otimes 1$ such that A is $(H_n^0)^{1/2}$ -relatively compact. The space spanned by such operators is invariant by a reasoning similar to that applied in part two of Proposition 7.*

The picture is that the map $\gamma \mapsto e^{-it\mathbb{H}}\gamma e^{it\mathbb{H}}$ is not continuous with respect to the geometric topology on $\mathfrak{S}_1(\mathcal{F}^{\leq N})$, it is however continuous for a similarly defined topology on the energy space $\mathfrak{S}_{1,H^0}(\mathcal{F}^{\leq N})$. This is sufficient to

obtain Theorem 1, since the sequence $M(T_k)$ can be uniformly approximated in that space.

APPENDIX A. PROOF OF LEMMA 4

Recall the notation $V := \{(x_1, x_2) \in \mathbb{R}^d \times \mathbb{R}^d : x_1 + x_2 = 0\}$ and parametrize this subspace by the coordinate $v = x_1 - x_2$. We will begin by showing that

$$(1 - \Delta_{x_1} - \Delta_{x_2})^{-1/2} w(v) (1 - \Delta_v)^{-1/2} = \sum_{j=1}^{\infty} K_j \otimes_V B_j, \quad (18)$$

with compact K_j and bounded B_j .

Let ξ denote the conjugate Fourier-variable to v and ζ that to $v^\perp := x_2 + x_1$. Let $\eta(\xi, \zeta) = (1 + 2\xi^2 + 2\zeta^2)^{-1/2}$, so $(1 - \Delta)^{-1/2}$ is just the Fourier-multiplier by η . Since η tends to zero at infinity, there exist functions $f_j(\xi), g_j(\zeta)$ with compact support, such that $\sum_{j=1}^m f_j(\xi)g_j(\zeta) \xrightarrow{k \rightarrow \infty} \eta(\xi, \zeta)$ in L^∞ . We can additionally arrange to have $\sum_{j=1}^k f_j(\xi)g_j(\zeta) \leq \eta(\xi, \zeta)$ for every k . Now define $B_j := \mathcal{F}_\zeta^{-1} g_j(\zeta) \mathcal{F}_{v^\perp}$ and $K_j := \mathcal{F}_\xi^{-1} f_j(\xi) \mathcal{F}_v w(v) (1 - \Delta_v)^{-1/2}$, where \mathcal{F}_x is the Fourier transform in the variable x . These operators clearly have the desired properties, and we now need to check convergence of the sum (18), which is not immediately obvious if w is not bounded. First, note that $K_0 := \mathcal{F}_\xi^{-1} (1 + \xi^2)^{-1/2} \mathcal{F}_v w(v) (1 - \Delta_v)^{-1/2}$ is a compact operator on $L^2(V)$ and $f_n \sqrt{1 + \xi^2}$ is bounded. In order to exploit the fact that $K_0 \otimes_V 1$ is compact in the first factor, we identify the operators on $L^2(V) \otimes L^2(V^\perp)$ with operators from $L^2(V)$ to $\mathcal{L}(L^2(V^\perp), L^2(\mathbb{R}^{2d}))$ via $A(\psi)\varphi := A\psi \otimes \varphi$. After this identification, $K_0 \otimes_V 1$ defines a compact operator, since for any weakly convergent sequence $\psi_k \rightharpoonup \psi$ in $L^2(V)$ and $\varphi \in L^2(V^\perp)$

$$\|(K \otimes_V 1)(\psi_k)\varphi - (K \otimes 1)(\psi)\varphi\| \leq \|K\psi_k - K\psi\|_{L^2(V)} \|\varphi\|_{L^2(V^\perp)}.$$

Furthermore, the operators

$$A_m := \sum_{j=1}^m \mathcal{F}^{-1} f_j(\xi) \sqrt{1 + \xi^2} g_j(\zeta) \mathcal{F}$$

converge to $\mathcal{F}^{-1} \eta(\xi, \zeta) \sqrt{1 + \xi^2} \mathcal{F}$ in the strong operator topology of $\mathcal{L}(L^2(V), \mathcal{L}(L^2(V^\perp), L^2(\mathbb{R}^{2d})))$, since for every $\psi \in L^2(V)$

$$\begin{aligned} & \lim_{m \rightarrow \infty} \left\| A_m(\psi) - \mathcal{F}_{(\xi, \zeta)}^{-1} \eta(\xi, \zeta) \sqrt{1 + \xi^2} \mathcal{F}_{(v, v^\perp)}(\psi) \right\|_{\mathcal{L}(L^2(V^\perp), L^2(\mathbb{R}^{2d}))} \\ &= \lim_{m \rightarrow \infty} \sup_{\zeta \in \mathbb{R}^d} \int_{\mathbb{R}^d} \left| \left(\sum_{j=1}^m f_j(\xi) g_j(\zeta) - \eta(\xi, \zeta) \right) \sqrt{1 + \xi^2} \widehat{\psi}(\xi) \right|^2 d\xi = 0, \end{aligned}$$

as the integrand is bounded by $|2\eta \sqrt{1 + \xi^2} \widehat{\psi}|^2 \in L^1$ and converges to zero, pointwise in ξ and uniformly in ζ . Thus, since $\sum_{j=1}^m K_j \otimes_V B_j = A_m(K_0 \otimes_V 1)$, the sum in (18) converges in norm. To deduce the statement of lemma 4

from (18) we use the formula

$$\begin{aligned} & (h_1 + h_2)^{-1/2} - (2e - \Delta_{x_1} - \Delta_{x_2})^{-1/2} \\ &= \frac{1}{2\pi i} \int_{\sigma} z^{-1/2} (h_1 + h_2 - z)^{-1} (V(x_1) + V(x_2)) (2e - \Delta_{x_1} - \Delta_{x_2} - z)^{-1} dz, \end{aligned} \quad (19)$$

where σ is the boundary of a sector in the right half plane $\operatorname{Re}(z) > 0$ with $\sigma \cap \{\operatorname{Im}(z) = 0\} = \{e\}$ and the integral converges in the operator norm. The argument we used to prove (18) also implies that $(1 - \Delta_{x_1} - z)^{-1/2} V(x_1) (2e - \Delta_{x_1} - \Delta_{x_2} - z)^{-1/2}$ is x_1 -semicompact, and thus (cf. Lemma 3)

$$(1 - \Delta_{x_1} - z)^{-1/2} V(x_1) (2e - \Delta_{x_1} - \Delta_{x_2} - z)^{-1} w(v) (1 - \Delta_v)^{-1/2}$$

is a compact operator for every $z \in \sigma$. Now for $f \in L^p(\mathbb{R}^d)$ with $p > \max(1, d/2)$ the Kato-Seiler-Simon inequality gives

$$\left\| (1 - \Delta_x - z)^{-1/2} f(x) (1 - \Delta_x - z)^{-1/2} \right\|_{\mathfrak{S}_p} \leq C(p, d) (1 + |z|^2)^{(d/4p - 1/2)} \|f\|_{L^p}.$$

Using the decomposition $V = f_1 + f_2$ we see that the integral in (19) still converges in norm if multiplied by $w(v) (1 - \Delta_v)^{-1/2}$ from the right. Consequently, it defines a compact operator, so the difference between equations (11) and (18) is compact and the proof is complete.

Acknowledgment. We thank Laurent Bruneau, Vladimir Georgescu and Phan Thành Nam for useful discussions. We acknowledge financial support from the European Research Council under the European Community's Seventh Framework Programme (FP7/2007-2013 Grant Agreement MNIQS 258023).

REFERENCES

- [1] W. O. AMREIN, A. BOUTET DE MONVEL, AND V. GEORGESCU, *C₀-groups, commutator methods and spectral theory of N-body Hamiltonians*, vol. 135 of Progress in Mathematics, Birkhäuser Verlag, Basel, 1996.
- [2] W. O. AMREIN AND V. GEORGESCU, *On the characterization of bound states and scattering states in quantum mechanics*, Helv. Phys. Acta, 46 (1973/74), pp. 635–658.
- [3] E. DAVIES, *Spectral theory and differential operators*, vol. 42 of Cambridge Studies in Advanced Mathematics, Cambridge University Press, Cambridge, 1995.
- [4] J. DEREZIŃSKI, *Asymptotic completeness of long-range N-body quantum systems*, Ann. of Math. (2), 138 (1993), pp. 427–476.
- [5] J. DEREZIŃSKI AND C. GÉRARD, *Mathematics of quantization and quantum fields*, Cambridge monographs on mathematical physics, Cambridge University Press, 2013.
- [6] V. ENSS, *Asymptotic completeness for quantum mechanical potential scattering. I. Short range potentials*, Commun. Math. Phys., 61 (1978), pp. 285–291.
- [7] G. M. GRAF, *Anderson localization and the space-time characteristic of continuum states*, J. Statist. Phys., 75 (1994), pp. 337–346.
- [8] H. GROSSE AND L. PITTNER, *On the number of unnatural parity bound states of the H⁻ ion*, J. Math. Phys., 24 (1983), pp. 1142–1147.
- [9] R. N. HILL, *Proof that the H⁻ ion has only one bound state*, Phys. Rev. Lett., 38 (1977), pp. 643–646.
- [10] ———, *Proof that the H⁻ ion has only one bound state. details and extension to finite nuclear mass*, J. Math. Phys., 18 (1977), pp. 2316–2330.

- [11] ———, *Proof that the H^- ion has only one bound state: A review, a new result, and some related unsolved problems*, in *Mathematical Problems in Theoretical Physics*, K. Osterwalder, ed., vol. 116 of *Lecture Notes in Physics*, Springer Berlin Heidelberg, 1980, pp. 52–56.
- [12] D. HUNDERTMARK, *On the time-dependent approach to Anderson localization*, *Math. Nachr.*, 214 (2000), pp. 25–38.
- [13] W. HUNZIKER AND I. M. SIGAL, *Time-Dependent Scattering Theory of N -Body Quantum Systems*, *Rev. Math. Phys.*, 12 (2000), pp. 1033–1084.
- [14] M. LEWIN, *Geometric methods for nonlinear many-body quantum systems*, *J. Funct. Anal.*, 260 (2011), pp. 3535–3595.
- [15] E. H. LIEB, *Bound on the maximum negative ionization of atoms and molecules*, *Phys. Rev. A*, 29 (1984), pp. 3018–3028.
- [16] M. REED AND B. SIMON, *Methods of Modern Mathematical Physics. III. Scattering theory*, Academic Press, New York, 1979.
- [17] D. RUELLE, *A remark on bound states in potential-scattering theory*, *Nuovo Cimento A* (10), 61 (1969), pp. 655–662.
- [18] I. M. SIGAL AND A. SOFFER, *Asymptotic completeness of N -particle long-range scattering*, *J. Amer. Math. Soc.*, 7 (1994), pp. 307–334.

E-mail address: lampart@ceremade.dauphine.fr

CNRS AND CEREMADE (UMR CNRS 7534), UNIVERSITY OF PARIS-DAUPHINE,
PLACE DE LATTRE DE TASSIGNY, 75775 PARIS CEDEX 16, FRANCE

E-mail address: mathieu.lewin@math.cnrs.fr

CNRS AND CEREMADE (UMR CNRS 7534), UNIVERSITY OF PARIS-DAUPHINE,
PLACE DE LATTRE DE TASSIGNY, 75775 PARIS CEDEX 16, FRANCE