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# Positive nonlinear CVFE scheme for degenerate anisotropic Keller-Segel system 

Clément Cancès, Moustafa Ibrahim, Mazen Saad


#### Abstract

In this paper, a nonlinear control volume finite element (CVFE) scheme for a degenerate KellerSegel model with anisotropic and heterogeneous diffusion tensors is proposed and analyzed. In this scheme, degrees of freedom are assigned to vertices of a primal triangular mesh, as in finite element methods. The diffusion term which involves an anisotropic and heterogeneous tensor is discretized on a dual mesh (Donald mesh) using the diffusion fluxes provided by the conforming finite element reconstruction on the primal mesh. The other terms are discretized using a nonclassical upwind finite volume scheme on the dual mesh. The scheme ensures the validity of the discrete maximum principle without any restriction on the transmissibility coefficients. The convergence of the scheme is proved under very general assumptions. Finally, some numerical experiments are carried out to prove the ability of the scheme to tackle degenerate anisotropic and heterogeneous diffusion problems over general meshes.


## 1 Introduction and model

In this paper, we are interested in degenerate nonlinear parabolic reaction-convection-diffusion systems modeling the chemotaxis process over general mesh, with anisotropic and heterogeneous diffusion tensors. From the numerical point of view, the convergence analysis of the finite volume scheme for this type of systems is carried out in [3] for the isotropic case (i.e. the diffusion tensor is considered to be proportional to the identity matrix) and under the "admissibility" assumption on the mesh used for the space discretization in the sense of satisfying the orthogonality condition (see e.g. [14]). Although its ability to ensure stability, the classical upwind finite volume method does not permit to handle anisotropic diffusion even if the mesh verifies the orthogonality condition. Various "multi-point" schemes, where the approximation of the flux through an edge involves several scalar unknowns, have been proposed for anisotropic diffusion problems, see for example $[16,15,12,10,2,11]$ for a detailed review of modern finite volume methods for diffusion equations. However, nonlinear corrections have been proposed in [5] in order to enforce the monotony, but no complete convergence proof have been provided for such methods yet.

Let us introduce the chemotaxis model. For that, let $\Omega$ be a connected open bounded polygonal domain of $\mathbb{R}^{2}$, and $t_{\mathrm{f}}>0$ be a fixed time. The modified Keller-Segel system (e.g., see [22, 23]) modeling the chemotaxis process is given by the following set of equations

[^0]\[

$$
\begin{cases}\partial_{t} u-\operatorname{div}(\Lambda(\mathbf{x}) a(u) \nabla u-\Lambda(\mathbf{x}) \chi(u) \nabla v)=f(u) & \text { in } Q_{t_{\mathrm{f}}}=\Omega \times\left(0, t_{\mathrm{f}}\right)  \tag{1}\\ \partial_{t} v-\operatorname{div}(D(\mathbf{x}) \nabla v)=g(u, v) & \text { in } Q_{t_{\mathrm{f}}}=\Omega \times\left(0, t_{\mathrm{f}}\right)\end{cases}
$$
\]

The system is complemented with zeros-flux boundary conditions on $\Sigma_{t_{\mathrm{f}}}:=\partial \Omega \times\left(0, t_{\mathrm{f}}\right)$ given by

$$
\begin{equation*}
(\Lambda(\mathbf{x}) a(u) \nabla u-\Lambda(\mathbf{x}) \chi(u) \nabla v) \cdot \mathbf{n}=0, \quad D(\mathbf{x}) \nabla v \cdot \mathbf{n}=0 \tag{2}
\end{equation*}
$$

and the initial conditions on $\Omega$ :

$$
\begin{equation*}
u(\mathbf{x}, 0)=u_{0}(\mathbf{x}), \quad v(\mathbf{x}, 0)=v_{0}(\mathbf{x}) \tag{3}
\end{equation*}
$$

In the above model, the density of the cell-population and the chemoattractant concentration are represented by $u=u(\mathbf{x}, t)$ and $v=v(\mathbf{x}, t)$ respectively. Next, $a(u)$ is a density-dependent diffusion coefficient, and $\Lambda(\mathbf{x})$ is the diffusion tensor in a heterogeneous medium. Furthermore, the function $\chi$ is the chemoattractant sensitivity, and $D(\mathbf{x})$ is the diffusion tensor for $v$. The function $f$ describes the cell density proliferation and the cell density death. The function $g$ describes the production and the degradation of the chemoattractant concentration; for simplicity, we assume that it is a linear function given by

$$
\begin{equation*}
g(u, v)=\alpha u-\beta v, \quad \alpha, \beta \geq 0 \tag{4}
\end{equation*}
$$

$\alpha$ and $\beta$ represent respectively the production and the degradation rate of the chemical concentration.
Let us state the main assumptions made about system (1)-(3):
(A1) The cell-density diffusion coefficient $a:[0,1] \longrightarrow \mathbb{R}^{+}$is a continuous function such that, $a(0)=a(1)=$ 0 , and $a(u)>0$ for $0<u<1$.
(A2) The chemosensitivity $\chi:[0,1] \longrightarrow \mathbb{R}^{+}$is a continuous function such that, $\chi(0)=\chi(1)=0$. Furthermore, we assume that there exists a function $\mu \in \mathscr{C}\left([0,1] ; \mathbb{R}^{+}\right)$, such that $\mu(u)=\frac{\chi(u)}{a(u)}$ for all $u \in(0,1)$ and $\mu(0)=\mu(1)=0$.
(A3) The diffusion tensors $\Lambda$ and $D$ are two bounded, uniformly positive symmetric tensors on $\Omega$, that is: $\forall \mathbf{w} \neq 0,0<T_{-}|\mathbf{w}|^{2} \leq\langle T(\mathbf{x}) \mathbf{w}, \mathbf{w}\rangle \leq T_{+}|\mathbf{w}|^{2}<\infty, \quad T=\Lambda$ or $D$.
(A4) The cell density proliferation $f$ is a continuous function such that $f(0) \geq 0$ and $f(1) \leq 0$.
(A5) The initial function $u_{0}$ and $v_{0}$ are two functions in $L^{2}(\Omega)$ such that, $0 \leq u_{0} \leq 1$ and $v_{0} \geq 0$.
In the sequel, we use the Lipschitz continuous nondecreasing function $\xi: \mathbb{R} \longrightarrow \mathbb{R}$ defined by

$$
\begin{equation*}
\xi(u):=\int_{0}^{u} \sqrt{a(s)} \mathrm{d} s, \quad \forall u \in \mathbb{R} \tag{5}
\end{equation*}
$$

We recall the definition of a weak solution of system (1)-(3).
Definition 1.1 (weak solution) Under the assumptions (A1)-(A5), we say that the couple of measurable functions ( $u, v$ ) is a weak solution of system (1)-(3) if

$$
\begin{aligned}
& 0 \leq u(\mathbf{x}, t) \leq 1, \quad 0 \leq v(\mathbf{x}, t) \quad \text { for a.e. in } Q_{t_{\mathrm{f}}}, \\
& \xi(u) \in L^{2}\left(0, t_{\mathrm{f}} ; H^{1}(\Omega)\right), \\
& v \in L^{\infty}\left(Q_{t_{\mathrm{f}}}\right) \cap L^{2}\left(0, t_{\mathrm{f}} ; H^{1}(\Omega)\right),
\end{aligned}
$$

and for all $\varphi, \psi \in \mathscr{D}\left(\bar{\Omega} \times\left[0, t_{\mathrm{f}}\right)\right)$, one has

$$
\begin{align*}
& -\int_{\Omega} u_{0}(\mathbf{x}) \varphi(\mathbf{x}, 0) \mathrm{d} \mathbf{x}-\iint_{Q_{t_{\mathrm{f}}}} u \partial_{t} \varphi \mathrm{~d} \mathbf{x} \mathrm{~d} t+\iint_{Q_{t_{\mathrm{f}}}} \sqrt{a(u)} \Lambda(\mathbf{x}) \nabla \xi(u) \cdot \nabla \varphi \mathrm{d} \mathbf{x} \mathrm{~d} t \\
& \quad-\iint_{Q_{t_{\mathrm{f}}}} \Lambda(\mathbf{x}) \chi(u) \nabla v \cdot \nabla \varphi \mathrm{~d} \mathbf{x} \mathrm{~d} t=\iint_{Q_{t_{\mathrm{f}}}} f(u) \varphi(\mathbf{x}, t) \mathrm{d} \mathbf{x} \mathrm{~d} t
\end{aligned} \quad \begin{aligned}
& -\int_{\Omega} v_{0}(\mathbf{x}) \psi(\mathbf{x}, 0) \mathrm{d} \mathbf{x}-\iint_{Q_{t_{\mathrm{f}}}} v \partial_{t} \psi \mathrm{~d} \mathbf{x} \mathrm{~d} t+\iint_{Q_{t_{\mathrm{f}}}} D(\mathbf{x}) \nabla v \cdot \nabla \psi \mathrm{~d} \mathbf{x} \mathrm{~d} t=\iint_{Q_{t_{\mathrm{f}}}} g(u, v) \psi \mathrm{d} \mathbf{x} \mathrm{~d} t . \tag{6}
\end{align*}
$$

A standard weak formulation uses the Kirchhoff transform $\kappa(u)$ as a primitive of the function $a(u)$. According to [7, 6], we approximate the degenerate diffusion term in its original form in (1). Next, we use the specific form of the chemoattractant function to propose a new scheme preserving the positivity of solutions and convergent.

Schemes with mixed conforming piecewise linear finite elements on triangles for the diffusion term and finite volume on dual elements were proposed and analyzed in [4, 9, 1] for fluid mechanics equations, and in [21] for a degenerate nonlinear chemotaxis model. The convergence analysis for these schemes is carried out for the case of anisotropic and heterogeneous diffusion problems under an essential assumption that all the transmissibility coefficients are nonnegative. However, there is no sufficient conditions for nonnegativity of transmissibility coefficients and therefore the schemes do not permit to tackle general anisotropic diffusion problems. Nevertheless, in [8] the authors propose a combined nonconforming finite elements finite volumes scheme for which they add a monotone regularization permitting positiveness of discrete solution; the convergence of the scheme, introduced in [5], is ensured under a numerical condition depending on the mesh size and on the discrete solutions.

Recently, Cancès and Guichard proposed and analyzed in [6] a nonlinear Control Volume Finite Element (CVFE) scheme for solving degenerate anisotropic parabolic diffusion equations modeling flows in porous media. The convergence analysis is carried out without any restriction on the transmissibility coefficients, and the efficiency of the scheme is tested using anisotropic diffusion tensors over an unstructured mesh.
Our aim is to elaborate a general approach, inspired from [6] and [21], to approximate a nonlinear degenerate parabolic system modeling the chemotaxis process over general mesh, with anisotropic and heterogeneous diffusion tensors. Especially, the diffusion terms are discretized by means of a conforming piecewise linear finite element method on a primal triangular mesh and using the Godunov scheme to approximate the diffusion fluxes provided by the conforming finite element reconstruction. The others terms are discretized by means of a nonclassical upwind finite volume method on a dual mesh (Donald mesh or Median dual mesh).

The rest of this paper is organized as follows. In section 2, we define a primal triangular mesh and its corresponding Donald dual mesh, next, we define standard $\mathbb{P}_{1}$ finite element and finite volume reconstructions. Then, we introduce the nonlinear CVFE scheme and specify the discretization of the degenerate diffusion and convection terms. In Section 3, we prove the existence of a discrete solution to the CVFE scheme based on the establishment of a priori estimates on the discrete solution as well as the discrete maximum principle. In Section 4, we give estimates on differences of time and space translates for the approximate solutions. In Section 5, using the Kolmogorov relative compactness criterion, we prove the convergence of a subsequent of discrete solutions to the weak solution (Definition 1.1). Finally, some numerical simulations are carried out, in Section 6, to show the effectiveness of the scheme to tackle degenerate anisotropic and heterogeneous diffusion problems over general unstructured mesh.

## 2 The numerical scheme and main result

In this section, we describe the space and time discretizations of $Q_{t_{\mathrm{f}}}$, define the approximate spaces, introduce useful properties on discrete $H^{1}$-norms stemming from finite elements discretizations as well as the nonlinear CVFE scheme, and state the main result.

### 2.1 Space-time discretization and notations

### 2.1.1 Space discretizations of $\Omega$.

In order to discretize problem (1)-(3), we perform a finite element triangulation $\mathscr{T}$ of the polygonal domain $\Omega$, consisting of open bounded triangles such that $\bar{\Omega}=\bigcup_{T \in \mathscr{T}} \bar{T}$ and such that for all $T, T^{\prime} \in \mathscr{T}, \bar{T} \cap \overline{T^{\prime}}$ is either an empty set or a common vertex or edge of $T$ and $T^{\prime}$. We denote by $\mathscr{V}$ the set of vertices of the
discretization $\mathscr{T}$, located at positions $\left(\mathbf{x}_{K}\right)_{K \in \mathscr{V}}$, and by $\mathscr{E}$ the set of edges of $\mathscr{T}$ joining two vertices of $\mathscr{V}$. The edge joining two vertices $K$ and $L$ is denoted by $\sigma_{K L}$.

For a given triangle $T \in \mathscr{T}$, we denote by $\mathbf{x}_{T}$ the centre of gravity of $T$, by $\mathscr{E}_{T}$ the set of the edges of $T$, by $h_{T}$ the diameter of $T$, and by $\rho_{T}$ the diameter of the largest ball inscribed in the triangle $T$. We denote by $h$ the size of the triangulation $\mathscr{T}$ defined by $h:=\max _{T \in \mathscr{T}} h_{T}$ and and by $\theta_{\mathscr{T}}$ the shape regularity of the triangulation $\mathscr{T}$, defined by $\theta_{\mathscr{T}}:=\max _{T \in \mathscr{T}} \frac{h_{T}}{\rho_{T}}$.

For $K \in \mathscr{V}$, we denote by $\mathscr{E}_{K}$ the set of the edges having $K$ as an extremity, and by $\mathscr{T}_{K}$ the subset of $\mathscr{T}$ including the triangles having $K$ as a vertex. We also define a barycentric dual mesh $\mathscr{M}$ (known as


Fig. 1 Triangular mesh $\mathscr{T}$ and Donald dual mesh $\mathscr{M}$ : dual volumes, vertices, interfaces.

Donald dual or Median dual mesh) generated by the triangulation mesh $\mathscr{T}$. There is one dual element $\omega_{K}$ associated with each vertex $K \in \mathscr{V}$. We construct it around the vertex $K$ by connecting the barycenter $\mathbf{x}_{T}$ of each surrounding triangle $T \in \mathscr{T}_{K}$ with the barycenters $\mathbf{x}_{\sigma}$ of the edges $\sigma \in \mathscr{E}_{K}$. We refer to Fig. 1 for an illustration of the primal and the barycentric dual mesh in a two-dimensional space. Note that $\bar{\Omega}=\bigcup_{K \in \mathscr{V}} \bar{\omega}_{K}$. The 2-dimensional Lebesgue measure of $\omega_{K}$ is denoted by $m_{K}$.

### 2.2 Discrete finite elements space $\mathscr{H}_{\mathscr{T}}$, control volumes space $\mathscr{X}_{\mathscr{M}}$.

We define two discrete functional spaces associated with each mesh of the above meshes. The first one, denoted by $\mathscr{H}_{\mathscr{T}}$, is the usual $\mathbb{P}_{1}$-finite element space corresponding to the triangular mesh $\mathscr{T}$, consisting of piecewise affine finite elements.

$$
\mathscr{H}_{\mathscr{T}}:=\left\{\varphi \in \mathscr{C}^{0}(\bar{\Omega}) ; \varphi_{T} \in \mathbb{P}_{1}(\mathbb{R}), \forall T \in \mathscr{T}\right\} \subset H^{1}(\Omega)
$$

The canonical basis of $\mathscr{H}_{\mathscr{T}}$ is spanned by the shape functions $\left(\varphi_{K}\right)_{K \in \mathscr{V}}$, such that

$$
\varphi_{K}\left(\mathbf{x}_{K}\right)=1, \quad \varphi_{K}\left(\mathbf{x}_{L}\right)=0 \text { if } L \neq K, \quad \forall K \in \mathscr{V}
$$

On the other hand, we denote by $\mathscr{X}_{\mathscr{M}}$ the discrete control volumes space consisting of piecewise constant functions on the dual mesh $\mathscr{M}$.

$$
\mathscr{X}_{\mathscr{M}}=\left\{\varphi: \Omega \longrightarrow \overline{\mathbb{R}} \text { measurable; } \varphi_{\left.\right|_{\omega_{K}}} \in \overline{\mathbb{R}} \text { is constant, } \forall K \in \mathscr{V}\right\} .
$$

Given a vector $\left(u_{K}\right)_{K \in \mathscr{V}} \in \mathbb{R}^{\# \mathscr{V}}\left(\right.$ resp. $\left.\left(v_{K}\right)_{K \in \mathscr{V}} \in \mathbb{R}^{\# \mathscr{V}}\right)$, there exists a unique function $u_{\mathscr{T}} \in \mathscr{H}_{\mathscr{T}}$ (resp. $v_{\mathscr{T}} \in \mathscr{H}_{\mathscr{T}}$ ) and a unique $u_{\mathscr{M}} \in \mathscr{X}_{\mathscr{M}}\left(\right.$ resp. $\left.v_{\mathscr{M}} \in \mathscr{X}_{\mathscr{M}}\right)$ such that

$$
\begin{array}{ll}
u_{\mathscr{T}}\left(\mathbf{x}_{K}\right)=u_{\mathscr{M}}\left(\mathbf{x}_{K}\right)=u_{K}, & \forall K \in \mathscr{V}, \\
v_{\mathscr{T}}\left(\mathbf{x}_{K}\right)=v_{\mathscr{M}}\left(\mathbf{x}_{K}\right)=v_{K}, & \forall K \in \mathscr{V} . \tag{8}
\end{array}
$$

For all $(K, L) \in \mathscr{V}^{2}$, we define the transmissibility coefficient $T_{K L}$ by

$$
\begin{equation*}
T_{K L}=-\int_{\Omega} T(\mathbf{x}) \nabla \varphi_{K}(\mathbf{x}) \cdot \nabla \varphi_{L}(\mathbf{x}) \mathrm{d} \mathbf{x}=T_{L K}, \quad T(\mathbf{x})=\Lambda(\mathbf{x}) \text { or } D(\mathbf{x}) \tag{9}
\end{equation*}
$$

We have $T_{K K}=-\sum_{L \neq K} T_{K L}$, since $\sum_{K \in \mathscr{V}} \nabla \varphi_{K}=0$. As a consequence, one has

$$
\int_{\Omega} T(\mathbf{x}) \nabla u_{\mathscr{T}} \cdot \nabla v_{\mathscr{T}} \mathrm{d} \mathbf{x}=\sum_{\sigma_{K L} \in \mathscr{E}} T_{K L}\left(u_{K}-u_{L}\right)\left(v_{K}-v_{L}\right), \quad T(\mathbf{x})=\Lambda(\mathbf{x}) \text { or } D(\mathbf{x})
$$

### 2.3 Time discretization of $\left(0, t_{\mathrm{f}}\right)$.

For the time discretization of the interval $\left(0, t_{\mathrm{f}}\right)$, we consider a uniform time discretization, and we do not impose any restriction on the time step. In addition, we assume that the spatial meshes do not change with the time step. We note that all the results presented in this paper can be extended to the case of general time discretization.
Let $N$ be a nonnegative integer, we define the uniform time step $\Delta t=t_{\mathrm{f}} /(N+1)$, and $t^{n}=n \Delta t$ for all $n \in\{0, \ldots, N+1\}$, so that $t^{0}=0$, and $t^{N+1}=t_{\mathrm{f}}$.

### 2.4 Space-time discretization of $Q_{t_{\mathrm{f}}}$.

Here, we define the space and time discrete spaces $\mathscr{H}_{\mathscr{T}, \Delta t}$ and $\mathscr{X}_{\mathscr{M}, \Delta t}$ as the set of piecewise constant functions in time with values in $\mathscr{H}_{\mathscr{T}}$ and $\mathscr{X}_{\mathscr{T}}$ respectively.

$$
\begin{aligned}
& \mathscr{H}_{\mathscr{T}, \Delta t}=\left\{\varphi \in L^{2}\left(0, t_{\mathrm{f}} ; H^{1}(\Omega)\right), \varphi(\mathbf{x}, t)=\varphi\left(\mathbf{x}, t^{n+1}\right) \in \mathscr{H}_{\mathscr{T}}, \quad \forall t \in\left(t^{n}, t^{n+1}\right]\right\} \\
& \mathscr{X}_{\mathscr{M}, \Delta t}=\left\{\varphi: Q_{t_{\mathrm{f}}} \longrightarrow \overline{\mathbb{R}} \text { measurable, } \varphi(\mathbf{x}, t)=\varphi\left(\mathbf{x}, t^{n+1}\right) \in \mathscr{X}_{\mathscr{M}}, \quad \forall t \in\left(t^{n}, t^{n+1}\right]\right\} .
\end{aligned}
$$

For a given $\left(u_{K}^{n}\right)_{n \in\{0, \cdots, N+1\}, K \in \mathscr{V}} \in \mathbb{R}^{(N+2) \#^{\mathscr{V}}}$ (resp. $\left.\left(v_{K}^{n}\right)_{n \in\{0, \cdots, N+1\}, K \in \mathscr{V}}\right)$, there exists a unique function $u_{\mathscr{T}, \Delta t} \in \mathscr{H}_{\mathscr{T}, \Delta t}\left(\right.$ resp. $\left.v_{\mathscr{T}, \Delta t} \in \mathscr{H}_{\mathscr{T}, \Delta t}\right)$ and a unique $u_{\mathscr{M}, \Delta t} \in \mathscr{X}_{\mathscr{M}, \Delta t}$ (resp. $\left.v_{\mathscr{M}, \Delta t} \in \mathscr{X}_{\mathscr{M}, \Delta t}\right)$ such that

$$
\begin{array}{ll}
u_{\mathscr{T}, \Delta t}\left(\mathbf{x}_{K}, t\right)=u_{\mathscr{M}, \Delta t}\left(\mathbf{x}_{K}, t\right)=u_{K}^{n+1}, & \forall K \in \mathscr{V}, \forall t \in\left(t^{n}, t^{n+1}\right], \\
v_{\mathscr{T}, \Delta t}\left(\mathbf{x}_{K}, t\right)=v_{\mathscr{M}, \Delta t}\left(\mathbf{x}_{K}, t\right)=v_{K}^{n+1}, & \forall K \in \mathscr{V}, \forall t \in\left(t^{n}, t^{n+1}\right] . \tag{10}
\end{array}
$$

### 2.5 The nonlinear CVFE scheme

The discretizations of the initial data $u_{K}^{0}$ and $v_{K}^{0}, K \in \mathscr{V}$ are defined by

$$
\begin{equation*}
u_{\mathscr{M}}^{0}(\mathbf{x})=u_{K}^{0}=\frac{1}{m_{K}} \int_{\omega_{K}} u_{0}(\mathbf{y}) \mathrm{d} \mathbf{y}, \quad \forall \mathbf{x} \in \omega_{K} \tag{11}
\end{equation*}
$$

$$
\begin{equation*}
v_{\mathscr{M}}^{0}(\mathbf{x})=v_{K}^{0}=\frac{1}{m_{K}} \int_{\omega_{K}} v_{0}(\mathbf{y}) \mathrm{d} \mathbf{y}, \quad \forall \mathbf{x} \in \omega_{K} \tag{12}
\end{equation*}
$$

### 2.5.1 Discretization of the first equation of system (1)

For all $K \in \mathscr{V}$, and $n \in\{0, \ldots, N\}$, we define the discretization of the diffusion term by

$$
\sum_{\sigma_{K L} \in \mathscr{E}_{K}} a_{K L}^{n+1} \Lambda_{K L}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)
$$

where,

$$
a_{K L}^{n+1}= \begin{cases}\max _{u \in I_{K L}^{n+1}} a(u) & \text { if } \Lambda_{K L} \geq 0  \tag{13}\\ \min _{u \in I_{K L}^{n+1}} a(u) & \text { if } \Lambda_{K L} \leq 0\end{cases}
$$

and $I_{K L}^{n+1}$ denotes the interval defined by

$$
I_{K L}^{n+1}= \begin{cases}{\left[u_{K}^{n+1}, u_{L}^{n+1}\right]} & \text { if } u_{K}^{n+1} \leq u_{L}^{n+1} \\ {\left[u_{L}^{n+1}, u_{K}^{n+1}\right]} & \text { otherwise }\end{cases}
$$

Let us focus on the discretization of the convection term, and recall that the function $\chi(u)$ is defined to be the product of the continuous functions $\mu(u)$ and $a(u)$. To handle the discretization of the convection term in order to obtain a robust and stable scheme, we perform a nonclassical upwind finite volume scheme which consists of considering an upwind scheme for the function $\mu(u)$ according to the discrete gradient of $v$, and an upwind finite volume scheme for the function $a(u)$ with respect to $u$. These choices of discretization are crucial to ensure the discrete maximum principle as well as the energy estimates on the approximate solutions.

We are now in a position to introduce what we call nonlinear control volume finite element (CVFE) scheme. For all $K \in \mathscr{V}$, and all $n \in\{0, \ldots, N\}$,

$$
\begin{equation*}
\frac{u_{K}^{n+1}-u_{K}^{n}}{\Delta t} m_{K}+\sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)-\sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)=f\left(u_{K}^{n+1}\right) m_{K} \tag{14}
\end{equation*}
$$

where, the transmissibility coefficients $\Lambda_{K L}$ and $D_{K L}$ are given by equality (9), and $\mu_{K L}^{n+1}$ denotes an approximation of $\mu(u)$ on the interfaces of $\omega_{K}$ with respect to the discrete gradient of $v$. We give here two examples on the construction of $\mu_{K L}^{n+1}$. The first example consists of taking the Engquist-Osher scheme and the second example consists of taking the Godunov scheme (see e.g. [20, 24]).

## Engquist-Osher scheme

- $\mu_{K L}^{n+1}= \begin{cases}\mu_{\downarrow}\left(u_{K}^{n+1}\right)+\mu_{\uparrow}\left(u_{L}^{n+1}\right), & \text { if } \Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right) \geq 0, \\ \mu_{\uparrow}\left(u_{K}^{n+1}\right)+\mu_{\downarrow}\left(u_{L}^{n+1}\right), & \text { if } \Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)<0 .\end{cases}$

The functions $\mu_{\uparrow}$ and $\mu_{\downarrow}$ are given by

$$
\mu_{\uparrow}(z):=\int_{0}^{z}\left(\mu^{\prime}(s)\right)^{+} \mathrm{d} s, \quad \mu_{\downarrow}(z):=-\int_{0}^{z}\left(\mu^{\prime}(s)\right)^{-} \mathrm{d} s .
$$

Herein, $s^{+}=\max (s, 0)$ and $s^{-}=\max (-s, 0)$.

## Godunov scheme

$$
\bullet \mu_{K L}^{n+1}= \begin{cases}\max _{\left[u_{K}^{n+1}, u_{L}^{n+1}\right]} \mu(u), & \text { if } \Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right) \geq 0 \\ \min _{\left[u_{L}^{n+1}, u_{K}^{n+1}\right]} \mu(u), & \text { if } \Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right) \geq 0 \\ \max _{\left[u_{L}^{n+1}, u_{K}^{n+1}\right]} \mu(u), & \text { if } \Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)<0 \\ \min _{\left[u_{K}^{n+1}, u_{L}^{n+1}\right]} \mu(u), & \text { if } \Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)<0\end{cases}
$$

Remark 1 The approximate function $\mu_{K L}^{n+1}=z\left(u_{K}^{n+1}, u_{L}^{n+1}\right)$ must be nonincreasing (resp. nondecreasing) with respect to the first variable $u_{K}^{n+1}$ and nondecreasing (resp. nonincreasing) with respect to the other variable $u_{L}^{n+1}$ when $\Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right) \geq 0\left(\right.$ resp. $\left.\Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right) \leq 0\right)$. We also have $z\left(u_{K}^{n+1}, u_{K}^{n+1}\right)=$ $\mu\left(u_{K}^{n+1}\right)$.

### 2.5.2 Discretization of the second equation of system (1)-(3)

Here, we focus on the discretization of the second equation of system (1)-(3). We note that a classical discretization of this equation is given by the following form

$$
\begin{equation*}
m_{K} \frac{v_{K}^{n+1}-v_{K}^{n}}{\Delta t}+\sum_{\sigma_{K L} \in \mathscr{E}_{K}} D_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)=m_{K}\left(\alpha u_{K}^{n}-\beta v_{K}^{n+1}\right) \tag{15}
\end{equation*}
$$

However, this discretization does not guaranty the positivity of the discrete solutions without any restriction on the transmissibility coefficients, for instance, one can get the discrete maximum principle by assuming that all the transmissibility coefficients $D_{K L}$ are nonnegative (see [21]).
Here, we propose a numerical discretization in order to ensure the discrete maximum principle without any restriction on the transmissibility coefficients. To do this, we introduce the following set of functions: $\eta(v)$, $p(v), \Gamma(v)$ and $\phi(v)$ defined by

$$
\begin{align*}
& \eta(v)=\max (0, \min (v, 1)),  \tag{16}\\
& p(v)=\int_{1}^{v} \frac{1}{\eta(s)} \mathrm{d} s= \begin{cases}\ln (v) & \text { if } v \in(0,1), \\
v-1 & \text { if } v \geq 1,\end{cases}  \tag{17}\\
& \Gamma(v)=\int_{1}^{v} p(s) \mathrm{d} s= \begin{cases}v \ln (v)-v+1 & \text { if } v \in[0,1), \\
\frac{(v-1)^{2}}{2} & \text { if } v \geq 1,\end{cases}  \tag{18}\\
& \phi(v)=\int_{0}^{v} \frac{1}{\sqrt{\eta(s)}} \mathrm{d} s= \begin{cases}2 \sqrt{v} & \text { if } v \in[0,1), \\
v+1 & \text { if } v \geq 1\end{cases} \tag{19}
\end{align*}
$$

In the sequel, we adopt the convention

$$
\begin{equation*}
\eta(v) p(v)=0 \quad \text { if } v \leq 0 \tag{20}
\end{equation*}
$$

We give the discretization of the second equation of system (1)-(3); specifically, we have

$$
\begin{equation*}
m_{K} \frac{v_{K}^{n+1}-v_{K}^{n}}{\Delta t}+\sum_{\sigma_{K L} \in \mathscr{E}_{K}} D_{K L} \eta_{K L}^{n+1}\left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)=m_{K}\left(\alpha u_{K}^{n}-\beta v_{K}^{n+1}\right) \tag{21}
\end{equation*}
$$

where, denoting by $J_{K L}^{n+1}=\left[\min \left(v_{K}^{n+1}, v_{L}^{n+1}\right), \max \left(v_{K}^{n+1}, v_{L}^{n+1}\right)\right]$, we have set

$$
\eta_{K L}^{n+1}= \begin{cases}\max _{s \in J_{K L}^{n+1}} \eta(s) & \text { if } D_{K L} \geq 0 \\ \min _{s \in J_{K L}^{n+1}}^{n+} \eta(s) & \text { if } D_{K L}<0\end{cases}
$$

We can show that the scheme (14)-(21), whose construction is based on finite elements for the diffusion term and a nonclassical upwind finite volume for the convection term, can be interpreted as a finite volume scheme. Indeed, denoting by

$$
\begin{aligned}
& F_{K L}^{n+1}=\Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)-\Lambda_{K L} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right), \\
& \Phi_{K L}^{n+1}=D_{K L} \eta_{K L}^{n+1}\left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right) .
\end{aligned}
$$

Then the scheme (14)-(21) rewrites

$$
\begin{cases}F_{K L}^{n+1}+F_{L K}^{n+1}=0=\Phi_{K L}^{n+1}+\Phi_{L K}^{n+1}, & \text { for all } \sigma_{K L} \in \mathscr{E}, \\ m_{K} \frac{u_{K}^{n+1}-u_{K}^{n}}{\Delta t}+\sum_{\sigma_{K L} \in \mathscr{E}_{K}} F_{K L}^{n+1}=f\left(u_{K}^{n+1}\right) m_{K}, & \text { for all } K \in \mathscr{V} \\ m_{K} \frac{v_{K}^{n+1}-v_{K}^{n}}{\Delta t}+\sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Phi_{K L}^{n+1}=g\left(u_{K}^{n}, v_{K}^{n+1}\right) m_{K}, & \text { for all } K \in \mathscr{V}\end{cases}
$$

### 2.6 Main result

Let $\left(\mathscr{T}_{m}\right)_{m \geq 1}$ be a sequence of triangulations of $\Omega$ such that

$$
h_{m}=\max _{T \in \mathscr{T}_{m}} \operatorname{diam}(T) \rightarrow 0 \text { as } m \rightarrow \infty .
$$

We assume that the sequence of triangulations has a bounded regularity, i.e., there exists a constant $\theta>0$ such that

$$
\theta_{\mathscr{T}_{m}} \leq \theta, \quad \forall m \geq 1
$$

As before, a sequence of dual meshes $\left(\mathscr{M}_{m}\right)_{m \geq 1}$ is given.
Let $\left(N_{m}\right)_{m}$ be an increasing sequence of integers, then we define the corresponding sequence of time steps $\left(\Delta t_{m}\right)_{m}$ such that $\Delta t_{m} \rightarrow 0$ as $m \rightarrow \infty$. The intention of this paper is to prove the following main result.
Theorem 2.1 Let $\left(u_{\mathscr{M}_{m}, \Delta t_{m}}, v_{\mathscr{M}_{m}, \Delta t_{m}}\right)_{m}$ be a sequence of solutions to the scheme (14)-(21), such that $0 \leq$ $u_{\mathscr{M}_{m}, \Delta t_{m}} \leq 1$ and $0 \leq v_{\mathscr{M}_{m}, \Delta t_{m}}$ for almost everywhere in $Q_{\mathrm{t}_{\mathrm{f}}}$, then

$$
u_{\mathscr{M}_{m}, \Delta t_{m}} \rightarrow u \text { and } v_{\mathscr{M}_{m}, \Delta t_{m}} \rightarrow v \quad \text { a.e. in } Q_{t_{\mathrm{f}}} \text { as } m \rightarrow \infty,
$$

where the couple $(u, v)$ is a weak solution to the system (1)-(3) in the sense of Definition 1.1.

## 3 Discrete properties, a priori estimates and existence of a discrete solution

In this section, we first bring up some technical lemmas presented by Cancès and Guichard [6], that we reproduce here for clarity. Then, we establish the a priori estimates necessary to prove later the existence of a solution to the discrete problem.
Lemma 3.1 Let $\left(u_{K}^{n+1}\right)_{K, n} \in \mathbb{R}^{(N+1) \#^{\mathscr{V}}}\left(\right.$ resp. $\left(v_{K}^{n+1}\right)_{K, n} \in \mathbb{R}^{(N+1) \# \mathscr{V}^{\prime}}$ ), then denoting by $\xi_{\mathscr{T}, \Delta t}\left(\right.$ resp. $\left.\phi_{\mathscr{T}, \Delta t}\right)$ the unique function of $\mathscr{H}_{\mathscr{T}, \Delta t}$ with nodal values $\left(\xi\left(u_{K}^{n+1}\right)\right) \in \mathbb{R}^{(N+1) \# \mathscr{V}}\left(\right.$ resp. $\left.\left(\phi\left(v_{K}^{n+1}\right)\right) \in \mathbb{R}^{(N+1) \# \mathscr{V}}\right)$, one has

$$
\begin{align*}
\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L} a_{K L}^{n+1} & \left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2}  \tag{22}\\
& \geq \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L}\left(\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right)^{2}=\iint_{Q_{t_{\mathrm{f}}}} \Lambda \nabla \xi_{\mathscr{T}, \Delta t} \cdot \nabla \xi_{\mathscr{T}, \Delta t} \mathrm{~d} \mathbf{x} \mathrm{~d} t
\end{align*}
$$

and

$$
\begin{align*}
\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} D_{K L} \eta_{K L}^{n+1} & \left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)^{2} \\
& \geq \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} D_{K L}\left(\phi\left(v_{K}^{n+1}\right)-\phi\left(v_{L}^{n+1}\right)\right)^{2}=\iint_{Q_{\mathrm{f}}} D \nabla \phi_{\mathscr{T}, \Delta t} \cdot \nabla \phi_{\mathscr{T}, \Delta \mathrm{t}} \mathrm{~d} \mathbf{x} \mathrm{~d} t . \tag{23}
\end{align*}
$$

Proof. We refer to [6, Lemma 3.1] for the proof of this lemma.
Let $T \in \mathscr{T}$, and let $(K, L) \in \mathscr{V}^{2}$, we denote by

$$
\begin{aligned}
& \lambda_{K L}^{T}:=-\int_{T} \Lambda \nabla \varphi_{K} \cdot \nabla \varphi_{L} \mathrm{~d} \mathbf{x}=\lambda_{L K}^{T} \\
& \delta_{K L}^{T}:=-\int_{T} D \nabla \varphi_{K} \cdot \nabla \varphi_{L} \mathrm{~d} \mathbf{x}=\delta_{L K}^{T}
\end{aligned}
$$

As a consequence, $\Lambda_{K L}=\sum_{T \in \mathscr{T}} \lambda_{K L}^{T}$ and $D_{K L}=\sum_{T \in \mathscr{T}} \delta_{K L}^{T}$ for all $\sigma_{K L} \in \mathscr{E}$.
Lemma 3.2 Let $\Psi_{\mathscr{T}}=\sum_{K \in \mathscr{V}} \psi_{K} \varphi_{K} \in \mathscr{H}_{\mathscr{T}}$, then there exists a quantity $C_{0}$ depending only on $\Lambda, D$, and $\theta_{\mathscr{T}}$ such that

$$
\begin{equation*}
\sum_{\sigma_{K L} \in \mathscr{E}} \sum_{T \in \mathscr{T}}\left|\lambda_{K L}^{T}\right|\left(\psi_{K}-\psi_{L}\right)^{2} \leq C_{0} \int_{\Omega} \Lambda(\mathbf{x}) \nabla \Psi_{\mathscr{T}} \cdot \nabla \Psi_{\mathscr{T}} \mathrm{d} \mathbf{x} \tag{24}
\end{equation*}
$$

and

$$
\begin{equation*}
\sum_{\sigma_{K L} \in \mathscr{E}} \sum_{T \in \mathscr{T}}\left|\delta_{K L}^{T}\right|\left(\psi_{K}-\psi_{L}\right)^{2} \leq C_{0} \int_{\Omega} D(\mathbf{x}) \nabla \Psi_{\mathscr{T}} \cdot \nabla \Psi_{\mathscr{T}} \mathrm{d} \mathbf{x} \tag{25}
\end{equation*}
$$

Proof. We refer to [6, Lemma 3.2] for the proof of this lemma.
Lemma 3.3 There exists a quantity $C_{1}$ depending only on $\Lambda, D$, and $\theta_{\mathscr{T}}$ such that

$$
\begin{align*}
& \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|\Lambda_{K L}\right| a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} \leq C_{1} \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} .  \tag{26}\\
& \text { and } \\
& \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|D_{K L}\right| \eta_{K L}^{n+1}\left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)^{2} \leq C_{1} \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} D_{K L} \eta_{K L}^{n+1}\left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)^{2} . \tag{27}
\end{align*}
$$

Proof. We denote by $\mathscr{E}^{-}:=\left\{\sigma_{K L} \in \mathscr{E} ; \Lambda_{K L}<0\right\}$, then since $|\mathbf{x}|=\mathbf{x}+2 \mathbf{x}^{-}, \mathbf{x}^{-}=\max (-\mathbf{x}, 0)$, one has

$$
\begin{aligned}
\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|\Lambda_{K L}\right| a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} & =\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} \\
& +2 \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}^{-}}\left|\Lambda_{K L}\right| a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2}
\end{aligned}
$$

Now, from the definition (9) of $a_{K L}^{n+1}$, there exists $c \in I_{K L}^{n+1}$ such that

$$
\left(\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right)^{2}=a(c)\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} \geq a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2}, \quad \forall \sigma_{K L} \in \mathscr{E}^{-}
$$

Therefore,

$$
\begin{align*}
\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|\Lambda_{K L}\right| a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} \leq & \leq \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2}  \tag{28}\\
& +2 \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|\Lambda_{K L}\right|\left(\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right)^{2}
\end{align*}
$$

Lemma 3.2 ensures the existence of a quantity $C_{0}>0\left(=C_{0}\left(\Lambda, \theta_{\mathscr{T}}\right)\right)$ such that

$$
\begin{array}{r}
\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|\Lambda_{K L}\right|\left(\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right)^{2} \leq \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \sum_{T \in \mathscr{T}}\left|\lambda_{K L}\right|\left(\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right)^{2} \\
\leq C_{0} \int_{Q_{t}} \Lambda \nabla \xi_{\mathscr{T}, \Delta t} \cdot \nabla \xi_{\mathscr{T}, \Delta t} \mathrm{~d} \mathbf{x} \mathrm{~d} t
\end{array}
$$

and from lemma 3.1, we deduce that

$$
\begin{equation*}
\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|\Lambda_{K L}\right|\left(\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right)^{2} \leq C_{0} \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} \tag{29}
\end{equation*}
$$

Plugging estimate (29) into estimate (28), then estimate (26) holds with $C_{1}=1+2 C_{0}$. The proof of estimate (27) is similar.

### 3.1 Discrete maximum principle

Lemma 3.4 Let $\left(u_{K}^{n+1}, v_{K}^{n+1}\right)_{K \in \mathscr{V}, n \in\{0, \ldots, N\}}$ be a solution to the CVFE scheme (14)-(21). Then, for all $K \in \mathscr{V}$, and all $n \in\{0, \ldots, N+1\}$, we have $0 \leq u_{K}^{n} \leq 1$ and $v_{K}^{n} \geq 0$.
Proof. We show this property using an induction on $n$. The property is true for $n=0$ thanks to the definitions (11) and (12) of $u_{K}^{0}$ and $v_{K}^{0}$ and to the assumptions on $u_{0}$ and $v_{0}$. Now, assume that the claim is true up to time step $n$. Consider a dual control volume $\omega_{K}$ such that $u_{K}^{n+1}=\min _{L \in \mathscr{V}}\left\{u_{L}^{n+1}\right\}$, we want to show that $u_{K}^{n+1} \geq 0$ i.e. $\left(u_{K}^{n+1}\right)^{-}=0$. Multiplying equation (14) by $-\left(u_{K}^{n+1}\right)^{-}$, one has

$$
\begin{align*}
-m_{K} \frac{u_{K}^{n+1}-u_{K}^{n}}{\Delta t} & \left(u_{K}^{n+1}\right)^{-}-\sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)\left(u_{K}^{n+1}\right)^{-}  \tag{30}\\
& +\sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)\left(u_{K}^{n+1}\right)^{-}=-m_{K} f\left(u_{K}^{n+1}\right)\left(u_{K}^{n+1}\right)^{-} \leq 0
\end{align*}
$$

to which, we have used the extension by $f(0) \geq 0$ (see assumption (A4)) of the continuous function $f$ for $u \leq 0$.
In view of the definition (13) of $a_{K L}^{n+1}$, and of the fact that $a(u)=0$ if $u \leq 0$, one has $a_{K L}^{n+1}=0$, if $\Lambda_{K L} \leq 0$. Therefore, the second term in the left hand side of equation (30) reads to

$$
-\sum_{\sigma_{K L} \in \mathscr{E}_{K}} a_{K L}^{n+1}\left(\Lambda_{K L}\right)^{+}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)\left(u_{K}^{n+1}\right)^{-} \geq 0
$$

Let us now focus on the third term of equation (30), and denote by $\mathscr{A}$ this term. Since $a_{K L}^{n+1}=0$ for $\Lambda_{K L} \leq 0$, then $\mathscr{A}$ rewrites

$$
\mathscr{A}=\sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L}^{+} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{+}\left(u_{K}^{n+1}\right)^{-}-\sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L}^{+} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{-}\left(u_{K}^{n+1}\right)^{-} .
$$

The second term of $\mathscr{A}$ is nonpositive, but in view of Remark 1 on the approximation $\mu_{K L}^{n+1}$ and in view of the extension by zero of the function $\mu$ for $u \leq 0$ since $\mu(0)=0$, one can deduce that

$$
\mu_{K L}^{n+1} \Lambda_{K L}^{+}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{-} \leq \mu\left(u_{K}^{n+1}\right) \Lambda_{K L}^{+}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{-}=0,
$$

thus, the second term of $\mathscr{A}$ is equal to zero and consequently $\mathscr{A} \geq 0$ since the first term of $\mathscr{A}$ is nonnegative. Finally, we use the identity $u_{K}^{n+1}=\left(u_{K}^{n+1}\right)^{+}-\left(u_{K}^{n+1}\right)^{-}$and the nonnegativity of $u_{K}^{n}$, one can deduce from equation (30) that $\left(u_{K}^{n+1}\right)^{-}=0$. According to the choice of the dual control volume $\omega_{K}$, then $\min _{L \in \mathscr{V}}\left\{u_{L}^{n+1}\right\}$ is non-negative. Consequently, $u_{K}^{n} \geq 0, \quad \forall K \in \mathscr{V}$, and all $n \in\{0, \ldots, N+1\}$.
In order to prove by induction that $u_{K}^{n} \leq 1, \forall K \in \mathscr{V}, \forall n \in\{0, \ldots, N+1\}$, we proceed in the same way as before, so that we consider a dual control volume $\omega_{K}$ such that $u_{K}^{n+1}=\max _{L \in \mathscr{Y}}\left\{u_{L}^{n+1}\right\}$. We get up the result using Remark 1 , the extension by zero of each of the function $a$, and $\mu$ for $u \geq 1$, and the extension by $f(1) \leq 0$ of the continuous function $f$ for $u \geq 1$.

Let us prove now that $v_{K}^{n+1} \geq 0$ for all $K \in \mathscr{V}$. Let $K_{m} \in \mathscr{V}$ be such that $v_{K_{m}}^{n+1}=\min _{K \in \mathscr{V}} v_{K}^{n+1}$, and assume that $v_{K_{m}}^{n} \geq 0$. Thanks to the convention (20), we can claim that

$$
D_{K_{m} L} \eta_{K_{m} L}^{n+1}\left(p\left(v_{K_{m}}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)=0 \quad \text { if } D_{K_{m} L} \leq 0
$$

and that

$$
D_{K_{m} L} \eta_{K_{m} L}^{n+1}\left(p\left(v_{K_{m}}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right) \leq 0 \quad \text { if } D_{K_{m} L} \geq 0
$$

Therefore, the scheme (21) together with the positivity of $u_{K_{m}}^{n}$ yields

$$
v_{K_{m}}^{n+1} \geq \frac{v_{K_{m}}^{n}}{1+\beta \Delta t} \geq 0
$$

This achieves the proof of Lemma 3.4.

### 3.2 Entropy estimates on $v_{\mathscr{M}, \Delta t}$

Lemma 3.5 There exists $C>0$ depending only on $\left\|v_{0}\right\|_{L^{2}(\Omega)}, \Omega, t_{\mathrm{f}}, \alpha$ and $\beta$ such that, for all $n^{\star} \in$ $\{0, \ldots, N\}$, one has

$$
\sum_{K \in \mathscr{V}} m_{K} \Gamma\left(v_{K}^{n^{\star}+1}\right)+\sum_{n=0}^{n^{\star}} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} D_{K L} \eta_{K L}^{n+1}\left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)^{2} \leq C
$$

Proof. It follows from Jensen's inequality - recall that $\Gamma$ is convex - that

$$
\sum_{K \in \mathscr{V}} m_{K} \Gamma\left(v_{K}^{0}\right) \leq \int_{\Omega} \Gamma\left(v_{0}(\mathbf{x})\right) \mathrm{d} \mathbf{x}
$$

Since $\Gamma(v) \leq(v-1)^{2}$ for all $v \geq 0$, we obtain that

$$
\begin{equation*}
\sum_{K \in \mathscr{V}} m_{K} \Gamma\left(v_{K}^{0}\right) \leq \int_{\Omega}\left(v_{0}(\mathbf{x})-1\right)^{2} \mathrm{~d} \mathbf{x} \leq C \tag{31}
\end{equation*}
$$

Multiplying the scheme (21) by $p\left(v_{K}^{n+1}\right) \Delta t$ and summing of $K \in \mathscr{V}$ and $n=0, \ldots, n^{\star}$ provides

$$
\begin{equation*}
\mathscr{A}+\mathscr{B}=\mathscr{C}, \tag{32}
\end{equation*}
$$

where we have set

$$
\begin{aligned}
\mathscr{A} & =\sum_{n=0}^{n^{\star}} \sum_{K \in \mathscr{V}} m_{K}\left(v_{K}^{n+1}-v_{K}^{n}\right) p\left(v_{K}^{n+1}\right) \\
\mathscr{B} & =\sum_{n=0}^{n^{\star}} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} D_{K L} \eta_{K L}^{n+1}\left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)^{2} \\
\mathscr{C} & =\sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} m_{K}\left(\alpha u_{K}^{n}-\beta v_{K}^{n+1}\right) p\left(v_{K}^{n+1}\right)
\end{aligned}
$$

Since, thanks to Lemma 3.4, $u_{K}^{n}$ is non-negative for all $K \in \mathscr{V}$ and all $n \geq 0$, and since $p(v) \leq(v-1)$ for all $v \geq 0$ (with the convention $p(0)=-\infty$ ), one has

$$
\alpha u_{K}^{n} p\left(v_{K}^{n+1}\right) \leq \alpha u_{K}^{n}\left(v_{K}^{n+1}-1\right)
$$

On the other hand, there exists an absolute constant $c^{\star}$ such that $v p(v) \geq(v-1)^{2}-c^{\star}$ for all $v \geq 0$. Therefore,

$$
\beta v_{K}^{n+1} p\left(v_{K}^{n+1}\right) \geq \beta\left(v_{K}^{n+1}-1\right)^{2}-c^{\star}
$$

As a consequence, we obtain that

$$
\mathscr{C} \leq t_{\mathrm{f}}|\Omega| c^{\star}+\sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} m_{K}\left(\alpha u_{K}^{n}\left(v_{K}^{n+1}-1\right)-\beta\left(v_{K}^{n+1}-1\right)^{2}\right)
$$

Using the weighted Young's inequality $\alpha a b \leq \beta b^{2}+\frac{\alpha^{2}}{4 \beta} a^{2}$ for all $(a, b) \in \mathbb{R}^{2}$ provides

$$
\alpha u_{K}^{n}\left(v_{K}^{n+1}-1\right)-\beta\left(v_{K}^{n+1}-1\right)^{2} \leq \frac{\alpha^{2}}{4 \beta} u_{K}^{n} \leq \frac{\alpha^{2}}{4 \beta}
$$

thanks to Lemma 3.4. Hence, we obtain that

$$
\begin{equation*}
\mathscr{C} \leq t_{\mathrm{f}}|\Omega|\left(c^{\star}+\frac{\alpha^{2}}{4 \beta}\right) \tag{33}
\end{equation*}
$$

The function $p$ being increasing, an elementary convexity inequality provides that

$$
(a-b) p(a) \geq \Gamma(a)-\Gamma(b), \quad \forall(a, b) \in\left(\mathbb{R}_{+}\right)^{2}
$$

ensuring that

$$
\begin{equation*}
\mathscr{A} \geq \sum_{n=0}^{n^{\star}} \sum_{K \in \mathscr{V}} m_{K}\left(\Gamma\left(v_{K}^{n+1}\right)-\Gamma\left(v_{K}^{n}\right)\right)=\sum_{K \in \mathscr{V}} m_{K}\left(\Gamma\left(v_{K}^{n^{\star}+1}\right)-\Gamma\left(v_{K}^{0}\right)\right) . \tag{34}
\end{equation*}
$$

Using (33), (34) and (31) in (32) concludes the proof of Lemma 3.5.
Lemma 3.6 There exists $C$ depending only on $\left\|v_{0}\right\|_{L^{2}(\Omega)}, \Omega, t_{\mathrm{f}}, \alpha, \beta, D_{ \pm}, \Lambda_{+}$and $\theta_{\mathscr{T}}$, such that

$$
\begin{equation*}
\iint_{Q_{t_{\mathrm{f}}}} \Lambda(\mathbf{x}) \nabla v_{\mathscr{T}, \Delta t}(\mathbf{x}, t) \cdot \nabla v_{\mathscr{T}, \Delta t}(\mathbf{x}, t) \mathrm{d} \mathbf{x} \mathrm{~d} t=\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{2} \leq C \tag{35}
\end{equation*}
$$

Proof. Thanks to Lemmas 3.1 and 3.5, we know that

$$
\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} D_{K L}\left(\phi\left(v_{K}^{n+1}\right)-\phi\left(v_{L}^{n+1}\right)\right)^{2} \leq \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} D_{K L} \eta_{K L}^{n+1}\left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)^{2} \leq C
$$

Therefore, it follows from Lemma 3.2 that

$$
\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|D_{K L}\right|\left(\phi\left(v_{K}^{n+1}\right)-\phi\left(v_{L}^{n+1}\right)\right)^{2} \leq C
$$

since $\left(\phi\left(v_{K}^{n+1}\right)-\phi\left(v_{L}^{n+1}\right)\right)^{2} \geq\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{2}$, we obtain that

$$
\begin{aligned}
\iint_{Q_{t_{\mathrm{f}}}} D(\mathbf{x}) \nabla v_{\mathscr{T}, \Delta t}(\mathbf{x}, t) \cdot \nabla v_{\mathscr{T}, \Delta t}(\mathbf{x}, t) \mathrm{d} \mathbf{x} \mathrm{~d} t & =\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} D_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{2} \\
& \leq \sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|D_{K L}\right|\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{2} \leq C .
\end{aligned}
$$

It only remains to check that

$$
\iint_{Q_{t_{\mathrm{f}}}} \Lambda(\mathbf{x}) \nabla v_{\mathscr{T}, \Delta t}(\mathbf{x}, t) \cdot \nabla v_{\mathscr{T}, \Delta t}(\mathbf{x}, t) \mathrm{d} \mathbf{x} \mathrm{~d} t \leq \frac{\Lambda_{+}}{D_{-}} \iint_{Q_{t_{\mathrm{f}}}} D(\mathbf{x}) \nabla v_{\mathscr{T}, \Delta t}(\mathbf{x}, t) \cdot \nabla v_{\mathscr{T}, \Delta t}(\mathbf{x}, t) \mathrm{d} \mathbf{x} \mathrm{~d} t
$$

in order to conclude the proof of Lemma 3.6.
Remark 2 A careful analysis allows to prove, after a more involving proof very similar to the one of Lemma 3.5, that the constant $C$ depends neither on $D_{+}$nor on $\theta_{\mathscr{T}}$. The cornerstone of the corresponding proof is the inequality

$$
D_{K L} \eta_{K L}^{n+1}\left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)\left(v_{K}^{n+1}-v_{L}^{n+1}\right) \geq D_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{2}
$$

that relies on the definition of $\eta_{K L}^{n+1}$ and on the link between the functions $\eta$ and $p$.

### 3.3 Energy estimates on $u_{\mathscr{M}, \Delta t}$

In the following, $C$ denotes a "generic" constant, that may vary throughout the proofs.
Proposition 3.7 There exists a constant $C>0$ depending only on $\left\|v_{0}\right\|_{L^{2}(\Omega)}, \Omega, t_{\mathrm{f}}, \alpha, \beta, \Lambda, D$, and $\theta_{\mathscr{T}}$ such that, for all $n^{\star} \in\{0, \ldots, N\}$, one has

$$
\begin{equation*}
\sum_{K \in \mathscr{V}} m_{K}\left(u_{K}^{n^{\star}+1}\right)^{2}+\sum_{n=0}^{n^{\star}} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} \leq C \tag{36}
\end{equation*}
$$

Proof. We multiply equation (14) by $\Delta t u_{K}^{n+1}$ and sum over $K \in \mathscr{V}$ and $n \in\left\{0, \ldots, n^{\star}\right\}$. This yields

$$
\begin{equation*}
E_{1}+E_{2}+E_{3}=E_{4} \tag{37}
\end{equation*}
$$

where

$$
\begin{array}{ll}
E_{1}=\sum_{n=0}^{n^{\star}} \sum_{K \in \mathscr{V}} m_{K}\left(u_{K}^{n+1}-u_{K}^{n}\right) u_{K}^{n+1}, & E_{2}=\sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right) u_{K}^{n+1}, \\
E_{4}=\sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} m_{K} f\left(u_{K}^{n+1}\right) u_{K}^{n+1}, & E_{3}=-\sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right) u_{K}^{n+1} .
\end{array}
$$

For the time evolution term, we use the following inequality: $(a-b) a \geq \frac{1}{2}\left(a^{2}-b^{2}\right), \forall a, b \in \mathbb{R}$, to get

$$
\begin{equation*}
E_{1} \geq \frac{1}{2} \sum_{n=0}^{n^{\star}} \sum_{K \in \mathscr{V}} m_{K}\left(\left(u_{K}^{n+1}\right)^{2}-\left(u_{K}^{n}\right)^{2}\right)=\frac{1}{2} \sum_{K \in \mathscr{Y}} m_{K}\left(\left(u_{K}^{n^{\star}+1}\right)^{2}-\left(u_{K}^{0}\right)^{2}\right) \tag{38}
\end{equation*}
$$

Next, for the diffusion term, we reorganize the sum over the edges, we find

$$
\begin{equation*}
E_{2}=\sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right) u_{K}^{n+1}=\sum_{n=0}^{n^{\star}} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} . \tag{39}
\end{equation*}
$$

Similarly, we reorganize the sum over the edges for the convection term. We obtain, using the uniform boundedness of the function $\mu$, that

$$
\begin{aligned}
E_{3} & =-\sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right) u_{K}^{n+1} \\
& =-\sum_{n=0}^{n^{\star}} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}} \Lambda_{K L} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)\left(u_{K}^{n+1}-u_{L}^{n+1}\right) .
\end{aligned}
$$

Using the weighted Young inequality, we deduce

$$
\begin{aligned}
\left|E_{3}\right| & \leq C \sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}}\left|\Lambda_{K L}\right| a_{K L}^{n+1}\left|v_{K}^{n+1}-v_{L}^{n+1}\right|\left|u_{K}^{n+1}-u_{L}^{n+1}\right| \\
& \leq C \sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}}\left|\Lambda_{K L}\right|\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{2}+\frac{1}{2 C_{1}} \sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}}\left|\Lambda_{K L}\right| a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2},
\end{aligned}
$$

where $C_{1}$ is the same constant introduced in Lemma 3.3.
Thanks to estimates (24) and (26), one has

$$
\left|E_{3}\right| \leq C \sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{2}+\frac{1}{2} \sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2}
$$

Therefore, Lemma 3.6 provides that

$$
\begin{equation*}
\left|E_{3}\right| \leq C+\frac{1}{2} \sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}} \sum_{\sigma_{K L} \in \mathscr{E}_{K}} \Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2} \tag{40}
\end{equation*}
$$

Finally, for the reaction term, since $0 \leq u_{K}^{n+1} \leq 1$ thanks to Lemma 3.4, one has

$$
\begin{equation*}
E_{4}=\sum_{n=0}^{n^{\star}} \Delta t \sum_{K \in \mathscr{V}_{m}} m_{K} f\left(u_{K}^{n+1}\right) u_{K}^{n+1} \leq|\Omega|\|f\|_{L^{\infty}(0,1)} t_{\mathrm{f}} \tag{41}
\end{equation*}
$$

Plugging estimates (38)-(41) into equation (37), one deduces that estimate (36) holds.

### 3.4 Enhanced estimate on $v_{\mathcal{M}, \Delta t}$

The goal of this section is to prove a refined estimate on $v_{\mathscr{M}, \Delta t}$ inspired from [6, Lemma 3.10], claiming that either $v_{\mathscr{M}, \Delta t}$ is constant equal to 0 , or $v_{\mathscr{M}, \Delta t} \geq r_{h}>0$ for some $r_{h}$ depending on the discretization parameters. The first step consists of bounding from below the $L^{\infty}\left(\left(0, t_{\mathrm{f}}\right) ; L^{1}(\Omega)\right)$ norm of $v_{\mathscr{M}, \Delta t}$.
Lemma 3.8 Assume that $\int_{\Omega} u_{0}(\mathbf{x}) \mathrm{d} \mathbf{x}>0$ or $\int_{\Omega} v_{0}(\mathbf{x}) \mathrm{d} \mathbf{x}>0$, then there exists $\kappa>0$ depending on the discretization and on the data such that

$$
\int_{\Omega} v_{\mathscr{M}, \Delta t}(\mathbf{x}, t) \mathrm{d} \mathbf{x} \geq \kappa, \quad \forall t \in\left[0, t_{\mathrm{f}}\right]
$$

Proof. Summing equation (21) over $K \in V$ ensures that

$$
\begin{equation*}
\sum_{K \in \mathscr{V}} m_{K}(1+\beta \Delta t) v_{K}^{n+1}=\sum_{K \in \mathscr{V}} m_{K} v_{K}^{n}+\alpha \Delta t \sum_{K \in \mathscr{V}} m_{K} u_{K}^{n}, \quad \forall n \in\{0, \ldots, N\} . \tag{42}
\end{equation*}
$$

Assume that $v_{K_{\star}^{n}}^{n}>0$ or $u_{K_{\star}^{n}}^{n}>0$ for some $K_{\star}^{n} \in \mathscr{V}$, as this is the case for $n=0$ because of the assumption on the initial data $u_{0}$ and $v_{0}$, then we deduce from (42) and from the non-negativity of $v_{K}^{n}$ and $u_{K}^{n}$ proved in Lemma 3.4 that

$$
\sum_{K \in \mathscr{V}} m_{K}(1+\beta \Delta t) v_{K}^{n+1}>0 .
$$

In particular, there exists $K_{\star}^{n+1} \in \mathscr{V}$ such that $v_{K_{\star}^{n+1}}^{n+1}$ is (strictly) positive and

$$
\sum_{K \in \mathscr{V}} m_{K} v_{K}^{n+1}:=\kappa_{n+1}>0 .
$$

One concludes the proof by setting $\kappa=\min _{n=1, \ldots, N+1} \kappa_{n}$.
We give now the definition of $D$-transmissive path, which was introduced in [6, Definition 3.4].
Definition 3.1 A D-transmissive path w joining $K_{\mathrm{i}} \in \mathscr{V}$ to $K_{\mathrm{f}} \in \mathscr{V}$ consists in a list of vertices $\left(K_{q}\right)_{0 \leq q \leq M}$ such that $K_{\mathrm{i}}=K_{0}, K_{\mathrm{f}}=K_{M}$, with $K_{q} \neq K_{\ell}$ if $q \neq \ell$, and such that $\sigma_{K_{q} K_{q+1}} \in \mathscr{E}$ with $D_{K_{q} K_{q+1}}>0$ for all $q \in\{0, \ldots, M-1\}$. We denote by $\mathscr{W}\left(K_{\mathrm{i}}, K_{\mathrm{f}}\right)$ the set of the transmissive path joining $K_{\mathrm{i}} \in \mathscr{V}$ to $K_{\mathrm{f}} \in \mathscr{V}$.

We now state a result which is proved in [6, Lemma 3.5].
Lemma 3.9 For all $\left(K_{\mathrm{i}}, K_{\mathrm{f}}\right) \in \mathscr{V}^{2}$ there exists a transmissive path $\mathrm{w} \in \mathscr{W}\left(K_{\mathrm{i}}, K_{\mathrm{f}}\right)$.
We have now introduced all the necessary tools for proving the main result of this section.
Lemma 3.10 Assume that $\int_{\Omega} u_{0}(\mathbf{x}) \mathrm{d} \mathbf{x}>0$ or $\int_{\Omega} v_{0}(\mathbf{x}) \mathrm{d} \mathbf{x}>0$, then there exists $r_{h}>0$ depending on the data as well as on the mesh $\mathscr{T}$ and $\Delta t$ such that

$$
\begin{equation*}
v_{K}^{n+1} \geq r_{h}, \quad \forall K \in \mathscr{V}, \forall n \in\{0, \ldots, N\} . \tag{43}
\end{equation*}
$$

Proof. Thanks to Lemma 3.8, we know that there exists $K_{\mathrm{i}}$ such that $v_{K_{\mathrm{i}}}^{n+1}>0$. Let $K_{\mathrm{f}} \in \mathscr{V}$, then there exists a D-transmissive path $\mathrm{w}=\left(K_{q}\right)_{0 \leq q \leq M} \in \mathscr{W}\left(K_{\mathrm{i}}, K_{\mathrm{f}}\right)$ thanks to Lemma 3.9, with $K_{0}=K_{\mathrm{i}}$ and $K_{M}=K_{\mathrm{f}}$.

Thanks to Lemmas 3.3 and $3 . \overline{5}$, we know that there exists $C$ such that

$$
\sum_{n=0}^{N} \Delta t \sum_{\sigma_{K L} \in \mathscr{E}}\left|D_{K L}\right| \eta_{K L}^{n+1}\left(p\left(v_{K}^{n+1}\right)-p\left(v_{L}^{n+1}\right)\right)^{2} \leq C
$$

In particular, this ensures that

$$
D_{K_{q} K_{q+1}} \eta_{K_{q} K_{q+1}}^{n+1}\left(p\left(v_{K_{q}}^{n+1}\right)-p\left(v_{q+1}^{n+1}\right)\right)^{2} \leq \frac{C}{\Delta t}, \quad \forall q \in\{0, \ldots, M-1\}
$$

Assume now that $v_{K_{q}}^{n+1}>0$, as this is the case for $q=0$, then $\eta_{K_{q} K_{q+1}}^{n+1} \geq \eta\left(v_{K_{q}}^{n+1}\right)>0$. Then one has

$$
\begin{equation*}
\left(p\left(v_{K_{q}}^{n+1}\right)-p\left(v_{K_{q+1}}^{n+1}\right)\right)^{2} \leq \frac{C}{\Delta t D_{K_{q} K_{q+1}} \eta_{K_{q} K_{q+1}}^{n+1}}<\infty \tag{44}
\end{equation*}
$$

Since $\lim _{v \rightarrow 0} p(v)=-\infty$, we deduce from (44) that $p\left(v_{K_{q+1}}^{n+1}\right)>-\infty$, hence $v_{K_{q+1}}^{n+1}>0$. A straightforward induction provides that $v_{K_{\mathrm{f}}}^{n+1}>0$, and since $K_{\mathrm{f}}$ was chosen arbitrarily, we obtain that

$$
v_{K}^{n+1}>0, \quad \forall K \in \mathscr{V}
$$

Since the set $\mathscr{V} \times\{0, \ldots, N\}$ is finite, we can conclude that there exists $r_{h}$ such that (43) holds.

### 3.5 Existence of a discrete solution

Proposition 3.11 Given $\left(u_{K}^{n}, v_{K}^{n}\right)_{K \in \mathscr{V}}$ such that $u_{\mathscr{M}, \Delta t}(\cdot, n \Delta t)$ and $v_{\mathscr{M}, \Delta t}(\cdot, n \Delta t)$ are non-negative, then there exists (at least) one solution $\left(u_{K}^{n+1}, v_{K}^{n+1}\right)_{K \in \mathscr{V}}$ of the scheme (14),(21). Moreover, $u_{\mathscr{M}, \Delta t}(\cdot, n \Delta t)$ and $v_{\mathscr{M}, \Delta t}(\cdot, n \Delta t)$ are non-negative.

Proof. The case where $\left(u_{K}^{n}, v_{K}^{n}\right)_{K \in \mathscr{V}} \equiv 0$ has to be treated apart. In this very particular case, it is easy to check that $\left(u_{K}^{n+1}, v_{K}^{n+1}\right)_{K \in \mathscr{V}} \equiv 0$ is a solution to the scheme.

Let us now focus on the case where $u_{K}^{n}$ or $v_{K}^{n}$ is strictly positive for some $K \in \mathscr{V}$. Because of the weak coupling on the numerical scheme, we can first solve (21), and afterwards (14). The existence of a solution $\left(v_{K}^{n+1}\right)_{K \in \mathscr{V}}$ can be proved by slightly adapting the proof of [6, Proposition 3.11], which relies on a topological argument. The main difficulty comes from the fact the scheme (21) is not continuous w.r.t. $\left(v_{K}^{n+1}\right)_{K \in \mathscr{V}}$ on $\left(\mathbb{R}_{+}\right)^{\# \mathscr{V}}$, but Lemma 3.10 ensures that no component $v_{K}^{n+1}$ of the discrete solution can go close to 0 . Let us detail now the proof.

Let $\gamma \in[0,1]$, we denote by $\left(v_{K, \gamma}^{n+1}\right)_{K \in \mathscr{V}}$ the solution (if it exists) to the numerical scheme

$$
\begin{align*}
& \frac{v_{K, \gamma}^{n+1}-v_{K}^{n}}{\Delta t} m_{K}+\gamma \sum_{\sigma_{K L} \in \mathscr{E}_{K}} D_{K L} \eta_{K L, \gamma}^{n+1}\left(p\left(v_{K, \gamma}^{n+1}\right)-p\left(v_{L, \gamma}^{n+1}\right)\right)  \tag{45}\\
& \quad+(1-\gamma) \sum_{\sigma_{K L} \in \mathscr{E}_{K}}\left|D_{K L}\right|\left(p\left(v_{K, \gamma}^{n+1}\right)-p\left(v_{L, \gamma}^{n+1}\right)\right)=\alpha u_{K}^{n} m_{K}-\beta v_{K, \gamma}^{n+1} m_{K}
\end{align*}
$$

In the above scheme, we have set

$$
\eta_{K L, \gamma}^{n+1}= \begin{cases}\max _{v \in J_{K L, \gamma}^{n+1}} \eta(v) & \text { if } D_{K L} \geq 0, \\ \min _{v \in J_{K L, \gamma}^{n+1}} \eta(v) & \text { if } D_{K L}<0,\end{cases}
$$

where $J_{K L, \gamma}^{n+1}=\left[\min \left(v_{K, \gamma}^{n+1}, v_{L, \gamma}^{n+1}\right), \max \left(v_{K, \gamma}^{n+1}, v_{L, \gamma}^{n+1}\right)\right]$. Reproducing carefully the analysis carried out in $\S 3.2$ and $\S 3.4$, we get that for all $\gamma \in[0,1]$,

$$
\begin{equation*}
\sum_{\sigma_{K L} \in \mathscr{E}} D_{K L}\left(\phi\left(v_{K, \gamma}^{n+1}\right)-\phi\left(v_{L, \gamma}^{n+1}\right)\right)^{2} \leq \sum_{\sigma_{K L} \in \mathscr{E}} D_{K L} \eta_{K L, \gamma}^{n+1}\left(p\left(v_{K, \gamma}^{n+1}\right)-p\left(v_{L, \gamma}^{n+1}\right)\right)^{2} \leq C \tag{46}
\end{equation*}
$$

and, that there exists $\varepsilon>0$ such that

$$
\begin{equation*}
v_{K, \gamma}^{n+1} \geq \varepsilon>0, \quad \forall K \in \mathscr{V} \tag{47}
\end{equation*}
$$

This ensures in particular that for all $\gamma \in[0,1]$, the solutions of equation (45) stay in the interior of a compact subset $\mathscr{K}$ of $\mathbb{R}^{\# \mathscr{V}}$ such that

$$
\operatorname{dist}\left(\mathscr{K},\left(\mathscr{R}_{-}\right)^{\# \mathscr{Y}}\right) \geq \frac{\varepsilon}{2} .
$$

Define the function $\Upsilon: \mathscr{K} \times[0,1] \rightarrow \mathbb{R}^{\#^{\mathscr{V}}}$ by: $\forall K \in \mathscr{V}$,

$$
\begin{aligned}
\Upsilon_{K}\left(\left(w_{K}\right)_{K}, \gamma\right)= & \frac{w_{K}-v_{K}^{n}}{\Delta t} m_{K}+\gamma \sum_{\sigma_{K L} \in \mathscr{E}_{K}} D_{K L} \eta_{K L, \gamma}^{n+1}\left(p\left(w_{K}\right)-p\left(w_{L}\right)\right) \\
& \left.+(1-\gamma) \sum_{\sigma_{K L} \in \mathscr{E}_{K}}\left|D_{K L}\right|\left(w_{K}\right)-p\left(w_{L}\right)\right)-\alpha u_{K}^{n} m_{K}+\beta w_{K} m_{K}
\end{aligned}
$$

The function $\Upsilon$ is uniformly continuous on $\mathscr{K} \times[0,1]$, and for all $\gamma \in[0,1]$ the solution $v_{K, \gamma}^{n+1}$ of the nonlinear system

$$
\begin{equation*}
\Upsilon\left(\left(v_{K, \gamma}^{n+1}\right)_{K \in \mathscr{V}}, \gamma\right)=0 \tag{48}
\end{equation*}
$$

cannot reach $\partial \mathscr{K}$. For $\gamma=0$, the system is monotone, so that the system (48) admits a unique solution, whose topological degree is equal to 1 (we refer to [13, Proposition 3.1] for a proof of this property). The topological degree being constant w.r.t. $\gamma \in[0,1]$, the system (48) admits at least one solution for $\gamma=1$, concluding the proof of the existence of $\left(v_{K}^{n+1}\right)_{K \in \mathscr{V}}$.

The existence proof for $\left(u_{K}^{n+1}\right)_{K \in \mathscr{V}}$ is similar but simpler since
i. the a priori estimate $0 \leq u_{K}^{n+1} \leq 1$ is sufficient for the claim, and no energy estimate is needed here;
ii. the scheme (14) depends in a uniformly continuous way on $\left(u_{K}^{n+1}\right)_{K \in \mathscr{V}}$ on the compact subset $[-1,2]^{\# \mathscr{V}}$ of $\mathscr{R}^{\# \mathscr{V}}$.
Therefore, we let to the reader the care of checking the proof for self-conviction.

## 4 Compactness estimates on the family of discrete solutions.

In this section, we derive estimates on differences of time and space translates of the discrete solutions necessary to prove the relative compactness property of the sequence of approximate solutions. To do this, we use the increasing functions $\xi:[0,1] \longrightarrow \mathbb{R}$ and $\phi:[0,+\infty) \longrightarrow \mathbb{R}$ defined in equation (5) and (19).

For all $K \in \mathscr{V}_{m}$ and for all $n \geq 1$, we denote by $\xi_{K}^{n}=\xi\left(u_{K}^{n}\right)$, and by $\xi_{\mathscr{T}_{m}, \Delta t_{m}}$ the corresponding piecewise affine in space and constant in time reconstruction in $\mathscr{H}_{\mathscr{T}_{m}, \Delta t_{m}}$, and by $\xi_{\mathscr{M}_{m}, \Delta t_{m}}$ the piecewise constant reconstruction in $\mathscr{X}_{\mathscr{M}_{m}, \Delta t_{m}}$.

### 4.1 Time translate estimate.

We give below the time translate estimate for the family $\left(\xi_{M_{m}, \Delta t_{m}}\right)_{m}$ expressed using the function $\xi$ defined in (5). We denote by $Q_{t_{\mathrm{f}}-\tau}=\Omega \times\left(0, t_{\mathrm{f}}-\tau\right)$, for all $\tau \in\left(0, t_{\mathrm{f}}\right)$.

Lemma 4.1 There exists two constants $C_{\xi, \mathrm{t}}$ and $C_{v, t}$ independent of $h$ and $\tau$ such that,

$$
\begin{align*}
& \iint_{Q_{t_{\mathrm{f}}-\tau}}\left|\xi_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t+\tau)-\xi_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right|^{2} \mathrm{~d} \mathbf{x} \mathrm{~d} t \leq C_{\xi, \mathrm{t}}\left(\tau+\Delta t_{m}\right),  \tag{49}\\
& \iint_{Q_{\mathrm{t}_{\mathrm{f}}-\tau}}\left|v_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t+\tau)-v_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right|^{2} \mathrm{~d} \mathbf{x} \mathrm{~d} t \leq C_{v, \mathrm{t}}\left(\tau+\Delta t_{m}\right) \tag{50}
\end{align*}
$$

for all $\tau \in\left(0, t_{\mathrm{f}}\right)$.
Proof. We consider the quantity $A_{m}(t)$ defined by

$$
A_{m}(t)=\int_{\Omega}\left|\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t+\tau)-\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t)\right|^{2} \mathrm{~d} \mathbf{x}, \quad \text { for all } t \in\left(0, t_{\mathrm{f}}-\tau\right)
$$

which implies, that

$$
\iint_{Q_{t_{\mathrm{f}}-\tau}}\left|\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t+\tau)-\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t)\right|^{2} \mathrm{~d} \mathbf{x} \mathrm{~d} t=\int_{0}^{t_{\mathrm{f}}-\tau} A_{m}(t) \mathrm{d} t
$$

For $t \in\left(0, t_{\mathrm{f}}\right]$, we denote by $v(t) \in\{0, \ldots, N\}$ the unique positive integer such that $t_{v(t)}<t \leq t_{v(t+\tau)}$, so that, we can rewrite $A_{m}(t)$ as

$$
A_{m}(t)=\sum_{K \in \mathscr{V}_{m}}\left(\xi_{K}^{v(t+\tau)+1}-\xi_{K}^{v(t)+1}\right)^{2} m_{K}, \quad \text { for all } t \in\left(0, t_{\mathrm{f}}-\tau\right) .
$$

We observe, using the definition (5) of the function $\xi$, that

$$
\begin{aligned}
\left(\xi_{K}^{v(t+\tau)+1}-\xi_{K}^{v(t)+1}\right)^{2} \leq C\left(u_{K}^{v(t+\tau)+1}-u_{K}^{v(t)+1}\right) & \times\left(\xi\left(u_{K}^{v(t+\tau)+1}\right)-\xi\left(u_{K}^{v(t)+1}\right)\right) \\
& =C \sum_{n=v(t)+1}^{v(t+\tau)}\left(u_{K}^{n+1}-u_{K}^{n}\right)\left(\xi\left(u_{K}^{v(t+\tau)+1}\right)-\xi\left(u_{K}^{v(t)+1}\right)\right)
\end{aligned}
$$

Now, using the scheme (14) then gathering by edges and using the weighted Young inequality, we obtain

$$
\begin{aligned}
A_{m}(t) \leq C \sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}} \Delta t_{m} & {\left[\left(\Lambda_{K L} a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)-\Lambda_{K L} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)\right)\right.} \\
& \left.\times\left(\left(\xi\left(u_{K}^{v(t)+1}\right)-\xi\left(u_{L}^{v(t)+1}\right)\right)-\left(\xi\left(u_{K}^{v(t+\tau)+1}\right)-\xi\left(u_{L}^{v(t+\tau)+1}\right)\right)\right)\right] \\
+ & C \sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m} \sum_{K \in \mathscr{V}_{m}} m_{K} f\left(u_{K}^{n+1}\right)\left(\xi\left(u_{K}^{v(t+\tau)+1}\right)-\xi\left(u_{K}^{v(t)+1}\right)\right) \\
\leq & C\left(A_{1, m}(t)+A_{2, m}(t)+A_{3, m}(t)+A_{4, m}(t)+A_{5, m}(t)\right)
\end{aligned}
$$

where, we have set

$$
\begin{aligned}
& A_{1, m}(t)=\sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}}\left|\Lambda_{K L}\right| a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2}, \\
& A_{2, m}(t)=\frac{\|a\|_{\infty}^{2}+\|\chi\|_{\infty}^{2}}{2} \sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}}\left|\Lambda_{K L}\right|\left(\xi\left(u_{K}^{v(t)+1}\right)-\xi\left(u_{L}^{v(t)+1}\right)\right)^{2}, \\
& A_{3, m}(t)=\frac{\|a\|_{\infty}^{2}+\|\chi\|_{\infty}^{2}}{2} \sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}}\left|\Lambda_{K L}\right|\left(\xi\left(u_{K}^{v(t+\tau)+1}-\xi\left(u_{L}^{v(t+\tau)+1}\right)\right)^{2},\right. \\
& A_{4, m}(t)=\sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}}\left|\Lambda_{K L}\right|\left(v_{K}^{n+1}-v_{L}^{n+1}\right)^{2}, \\
& A_{5, m}(t)=\sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m} \sum_{K \in \mathscr{Y}_{m}} m_{K}\left|f\left(u_{K}^{n+1}\right)\right|\left|\xi\left(u_{K}^{v(t+\tau)+1}\right)-\xi\left(u_{K}^{v(t)+1}\right)\right| .
\end{aligned}
$$

Now, we introduce the characteristic function $\rho(n, t, \tau)=1$ if $t<n \Delta t_{m} \leq t+\tau$ and $\rho(n, t, \tau)=0$ otherwise. Let $\left(a^{n}\right)_{n \in\{0, \ldots, N\}}$ be a family of non negative real values, we have the following properties

$$
\int_{0}^{t_{\mathrm{f}}-\tau} \rho(n, t, \tau) \mathrm{d} t=\int_{t^{n}-\tau}^{t^{n}} \mathrm{~d} t=\tau, \quad \sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m}=\sum_{n ; t<t^{n}<t+\tau} t^{n+1}-t^{n} \leq \tau+\Delta t_{m}
$$

and

$$
\int_{0}^{t_{\mathrm{f}}-\tau} \sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m} a^{n+1} \mathrm{~d} t=\int_{0}^{t_{\mathrm{f}}-\tau} \sum_{n=0}^{N} \Delta t_{m} a^{n+1} \rho(n, t, \tau) \mathrm{d} t=\tau \sum_{n=0}^{N} \Delta t_{m} a^{n+1}
$$

Using these properties, Lemma 3.3, the a priori estimates (35)-(36) on $u_{K}^{n+1}$ and $v_{K}^{n+1}$, and the $L^{\infty}$ bound of the function $\xi$, one can deduce that

$$
\int_{0}^{t_{\mathrm{f}}-\tau} A_{1, m}(t) \leq C\left(\tau+\Delta t_{m}\right), \quad \int_{0}^{t_{\mathrm{f}}-\tau} A_{4, m}(t) \leq C\left(\tau+\Delta t_{m}\right), \quad \text { and } \int_{0}^{t_{\mathrm{f}}-\tau} A_{5, m}(t) \leq C\left(\tau+\Delta t_{m}\right)
$$

for some constant $C>0$. It remains to show that $\int_{0}^{t_{\mathrm{f}}-\tau} A_{2, m}(t) \leq C\left(\tau+\Delta t_{m}\right)$, and $\int_{0}^{t_{\mathrm{f}}-\tau} A_{3, m}(t) \leq$ $C\left(\tau+\Delta t_{m}\right)$, for some constant $C>0$.
Consider the function $\zeta$ defined by $\zeta(n, t)=1$ if $v(t)=n$ and $\zeta(n, t)=0$ otherwise. We have the following result (see e.g. [18])

$$
\begin{aligned}
\int_{0}^{t_{\mathrm{f}}-\tau}\left(\sum_{n=v(t)+1}^{v(t+\tau)} \Delta t_{m}\right) a^{v(t+\tau)+1} \mathrm{~d} t & \leq\left(\tau+\Delta t_{m}\right) \int_{0}^{t_{\mathrm{f}}-\tau} \sum_{m=0}^{N} a^{m+1} \zeta(m, t+\tau) \mathrm{d} t \\
& =\left(\tau+\Delta t_{m}\right) \sum_{m=0}^{N} a^{m+1} \int_{t^{m}-\tau}^{t^{m+1}-\tau} \mathrm{d} t=\left(\tau+\Delta t_{m}\right) \sum_{m=0}^{N} a^{m+1} \Delta t_{m}
\end{aligned}
$$

One can conclude the proof of estimate (49) using this property, Lemma 3.1, and Lemma 3.3. The proof of estimate (50) is similar.

We now extend by zero the functions $\xi_{M_{m}, \Delta t_{m}}$ and $\xi_{\mathscr{T}_{m}, \Delta t_{m}}$ outside of $Q_{t_{\mathrm{f}}}$ and give the time translate estimate over $\mathbb{R}^{3}$ on $\xi_{M_{m}, \Delta t_{m}}$. Indeed, there exists a constant $C>0$ independent of $m$ and $\tau$ such that

$$
\int_{\mathbb{R}} \int_{\mathbb{R}^{2}}\left|\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t+\tau)-\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t)\right|^{2} \mathrm{~d} \mathbf{x} \mathrm{~d} t \leq C\left(\tau+\Delta t_{m}\right), \quad \text { for all } \tau \in\left(0, t_{\mathrm{f}}\right)
$$

Proof. Using the extension by zero of $\xi_{M_{m}, \Delta t_{m}}$ outside of $Q_{t_{\mathrm{f}}}$, one has

$$
\begin{aligned}
\int_{\mathbb{R}} \int_{\mathbb{R}^{2}}\left|\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t+\tau)-\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t)\right|^{2} \mathrm{~d} \mathbf{x} \mathrm{~d} t & =\iint_{Q_{t_{\mathrm{f}}-\tau}}\left|\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t+\tau)-\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t)\right|^{2} \mathrm{~d} \mathbf{x} \mathrm{~d} t \\
& +\int_{t_{\mathrm{f}}-\tau}^{t_{\mathrm{f}}}
\end{aligned} \int_{\Omega}\left|\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t)\right|^{2} \mathrm{~d} \mathbf{x} \mathrm{~d} t .
$$

One can deduce the proof using Lemma 4.1 and the $L^{\infty}$ bound of the function $\xi$.
We give now the space translate estimate on $\xi_{\mathscr{M}_{m}, \Delta t_{m}}$.

### 4.2 Space translate estimate.

Lemma 4.2 There exists two constants $C_{\xi, \mathrm{s}}$ and $C_{v, s}$ independent of $m$ and $\mathbf{y}$ such that,

$$
\begin{align*}
& \int_{0}^{t_{\mathrm{f}}} \int_{\mathbb{R}^{2}}\left|\xi_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}+\mathbf{y}, t)-\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t)\right| \mathrm{d} \mathbf{x} \mathrm{~d} t \leq C_{\xi, \mathrm{s}}\left(|\mathbf{y}|+h_{m}\right),  \tag{51}\\
& \int_{0}^{t_{\mathrm{f}}} \int_{\mathbb{R}^{2}}\left|v_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}+\mathbf{y}, t)-v_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right| \mathrm{d} \mathbf{x} \mathrm{~d} t \leq C_{\xi, \mathrm{s}}\left(|\mathbf{y}|+h_{m}\right), \tag{52}
\end{align*}
$$

for all $\mathbf{y} \in \mathbb{R}^{2}$.
Proof. We follow the same proof as in [6]. We first state that

$$
\int_{0}^{t_{\mathrm{f}}} \int_{\mathbb{R}^{2}}\left|\xi_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}+\mathbf{y}, t)-\xi_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right| \mathrm{d} \mathbf{x} \mathrm{~d} t \leq C|\mathbf{y}|, \quad \text { for all } \mathbf{y} \in \mathbb{R}^{2}
$$

then, using the extension by zero of $\xi_{M_{m}, \Delta t_{m}}$ outside of $Q_{t_{\mathrm{f}}}$ and the triangle inequality, we get

$$
\begin{aligned}
\int_{0}^{t_{\mathrm{f}}} \int_{\mathbb{R}^{2}}\left|\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}+\mathbf{y}, t)-\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t)\right| \mathrm{d} \mathbf{x} \mathrm{~d} t & \leq \int_{0}^{t_{\mathrm{f}}} \int_{\mathbb{R}^{2}}\left|\xi_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}+\mathbf{y}, t)-\xi_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right| \mathrm{d} \mathbf{x} \mathrm{~d} t \\
& +2 \iint_{Q_{t_{\mathrm{f}}}}\left|\xi_{M_{m}, \Delta t_{m}}(\mathbf{x}, t)-\xi_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right| \mathrm{d} \mathbf{x} \mathrm{~d} t
\end{aligned}
$$

One can conclude the proof of estimate (51) using the following estimate which is a direct consequence of [6, Lemma A.2]

$$
\begin{equation*}
\iint_{Q_{t_{\mathrm{f}}}}\left|\xi_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t)-\xi_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right| \mathrm{d} \mathbf{x} \mathrm{~d} t \leq C h_{m} \tag{53}
\end{equation*}
$$

The proof of estimate (52) follows the same lines. This ends the proof of the lemma.
Lemma 4.3 (Strong convergence in $L^{1}\left(Q_{t_{\mathrm{f}}}\right)$ ) There exists a subsequence of the sequence $\left(\xi(u)_{\mathscr{M}_{m}, \Delta t_{m}}\right)_{m \geq 1}$ (resp. $\left.\left(v_{M_{m}, \Delta t_{m}}\right)_{m>1}\right)$ which converges in $L^{1}\left(Q_{t_{\mathrm{f}}}\right)$ to some function $\xi(u) \in L^{2}\left(0, t_{\mathrm{f}} ; H^{1}(\Omega)\right)$ (resp. $v \in$ $\left.L^{2}\left(0, t_{\mathrm{f}} ; H^{1}(\Omega)\right)\right)$.
Proof. Let us consider the function $\xi_{M_{m}, \Delta t_{m}}$ on $Q_{t_{\mathrm{f}}}$ and its extension by zero outside of $Q_{t_{\mathrm{f}}}$. Lemma 4.1, Lemma 4.2, and the boundedness of $\xi$ due to Lemma 3.1 ensure that the sequence $\left(\xi(u)_{\mathscr{M}_{m}, \Delta t_{m}}\right)_{m \geq 1}$ verifies the assumptions of the Kolmogorov compactness criterion (see e.g. [19, 15]), where the third item of the theorem is satisfied using the triangle inequality: for any $\eta \in \mathbb{R}^{2}$ and $\tau \in \mathbb{R}$,

$$
\begin{aligned}
\left\|\xi_{M_{m}, \Delta t_{m}}(\cdot+\eta, \cdot+\tau)-\xi_{M_{m}, \Delta t_{m}}(\cdot, \cdot)\right\|_{L^{1}\left(\mathbb{R}^{2} \times \mathbb{R}\right)} \leq\left\|\xi_{M_{m}, \Delta t_{m}}(\cdot+\eta, \cdot)-\xi_{M_{m}, \Delta t_{m}}(\cdot, \cdot)\right\|_{L^{1}\left(\mathbb{R}^{2} \times \mathbb{R}\right)} \\
+\left\|\xi_{M_{m}, \Delta t_{m}}(\cdot, \cdot+\tau)-\xi_{M_{m}, \Delta t_{m}}(\cdot, \cdot)\right\|_{L^{1}\left(\mathbb{R}^{2} \times \mathbb{R}\right)}
\end{aligned}
$$

Kolmogorov's theorem ensures that the sequence $\left(\xi(u)_{\mathscr{M}_{m}, \Delta t_{m}}\right)_{m \geq 1}$ is relatively compact in $L^{1}\left(Q_{t_{\mathrm{f}}}\right)$, that implies the existence of a subsequence of $\left(\xi(u)_{\mathscr{M}_{m}, \Delta t_{m}}\right)_{m \geq 1}$ such that

$$
\begin{equation*}
\xi(u)_{\mathscr{M}_{m}, \Delta t_{m}} \longrightarrow \xi^{\star} \quad \text { strongly in } L^{1}\left(Q_{t_{\mathrm{f}}}\right) \tag{54}
\end{equation*}
$$

Furthermore, as $\xi$ is a continuous and non decreasing function on $[0,1]$, there exists a unique $u(\mathbf{x}, t)$ defined by

$$
u(\mathbf{x}, t)=\xi^{-1}\left(\xi^{\star}(\mathbf{x}, t)\right), \quad \text { for a.e. }(\mathbf{x}, t) \in Q_{t_{\mathrm{f}}-\tau}
$$

Since $\xi^{-1}$ is well defined and continuous, applying the $L^{\infty}$ bound on $u_{\mathscr{M}_{m}, \Delta t_{m}}$ and the dominated convergence theorem to $u_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t)=\xi^{-1}\left(\xi_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right), \quad \forall(\mathbf{x}, t) \in \omega_{K} \times\left(0, t_{\mathrm{f}}\right), \forall K \in \mathscr{V}_{m}$, we get

$$
u_{\mathscr{M}_{m}, \Delta t_{m}} \longrightarrow u \text { a.e. in } Q_{t_{\mathrm{f}}} \text { and strongly in } L^{p}\left(Q_{t_{\mathrm{f}}}\right) \text { for } p<+\infty
$$

It remains to show that $\xi(u) \in L^{2}\left(0, t_{\mathrm{f}} ; H^{1}(\Omega)\right)$. Indeed, plugging estimate (36) into estimate (22), one gets that $\nabla \xi(u)_{\mathscr{T}_{m}, \Delta t_{m}}$ is uniformly bounded in $\left(L^{2}\left(Q_{t_{\mathrm{f}}}\right)\right)^{2}$. It follows that $\left(\xi(u)_{\mathscr{T}_{m}, \Delta t_{m}}\right)_{m}$ is bounded in $L^{2}\left(0, t_{\mathrm{f}} ; H^{1}(\Omega)\right)$ since $\xi(u)_{\mathscr{T}_{m}, \Delta t_{m}}$ is uniformly bounded in $L^{2}\left(Q_{t_{\mathrm{f}}}\right)$. Consequently, the sequence $\left(\xi(u)_{\mathscr{T}_{m}, \Delta t_{m}}\right)_{m}$ converges weakly, up to an unlabeled subsequence, to a function $\tilde{\xi}$ in $L^{2}\left(0, t_{\mathrm{f}} ; H^{1}(\Omega)\right)$. According to estimate (53), the sequences $\left(\xi(u)_{\mathscr{T}_{m}, \Delta t_{m}}\right)_{m}$ and $\left(\xi(u)_{\mathscr{M}_{m}, \Delta t_{m}}\right)_{m}$ have the same limit, as a consequence $\xi(u)=\xi^{\star}=\tilde{\xi} \in L^{2}\left(0, t_{\mathrm{f}} ; H^{1}\left(Q_{t_{\mathrm{f}}}\right)\right)$.

## 5 Identification as a weak solution

It remains to be shown that $(u, v)$ satisfies the weak formulation (6)-(7). To do this, we consider a test function $\psi \in \mathscr{D}\left(\bar{\Omega} \times\left[0, t_{\mathrm{f}}\right)\right)$, and denote by $\psi_{K}^{n}=\psi\left(\mathbf{x}_{K}, t^{n}\right)$, for all $K \in \mathscr{V}_{m}$ and all $n \in\left\{0, \ldots, N_{m}\right\}$. Let
us focus on the convergence of the first equation of scheme (14)-(21), i.e., we show that equation (6) is verified when $m \rightarrow \infty$. We note that the convergence of the second equation of the scheme is similar and major difficulties that we can met are discussed hereafter.

Multiplying the first equation (14) by $\Delta t_{m} \psi_{K}^{n}$ and summing over $n \in\left\{0, \ldots, N_{m}\right\}$ and $K \in \mathscr{V}_{m}$ yields, after a reorganization of the sum,

$$
\begin{equation*}
\mathscr{A}_{m}+\mathscr{B}_{m}+\mathscr{C}_{m}+\mathscr{D}_{m}=\mathscr{F}_{m}, \tag{55}
\end{equation*}
$$

where

$$
\begin{aligned}
& \mathscr{A}_{m}=\sum_{n=0}^{N_{m}} \sum_{K \in \mathscr{V}_{m}}\left(u_{K}^{n+1}-u_{K}^{n}\right) \psi_{K}^{n} m_{K}, \quad \mathscr{F}_{m}=\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{K \in \mathscr{Y}_{m}} f\left(u_{K}^{n+1}\right) \psi_{K}^{n} m_{K}, \\
& \mathscr{B}_{m}=\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}} \Lambda_{K L}\left(a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)-\sqrt{a_{K L}^{n+1}}\left(\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right)\right)\left(\psi_{K}^{n}-\psi_{L}^{n}\right), \\
& \mathscr{C}_{m}=\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}} \Lambda_{K L} \sqrt{a_{K L}^{n+1}}\left(\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right)\left(\psi_{K}^{n}-\psi_{L}^{n}\right), \\
& \mathscr{D}_{m}=-\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}} \Lambda_{K L} \mu_{K L}^{n+1} a_{K L}^{n+1}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)\left(\psi_{K}^{n}-\psi_{L}^{n}\right) .
\end{aligned}
$$

## Accumulation term

Note that $\psi_{K}^{N_{m}+1}=0$ for all $K \in \mathscr{V}_{m}$, then, performing summation by parts in time, the term $\mathscr{A}_{m}$ can be rewritten

$$
\begin{aligned}
\mathscr{A}_{m} & =\sum_{n=0}^{N_{m}} \sum_{K \in \mathscr{V}_{m}} u_{K}^{n+1} \psi_{K}^{n} m_{K}-\sum_{n=1}^{N_{m}} \sum_{K \in \mathscr{V}_{m}} u_{K}^{n} \psi_{K}^{n} m_{K}-\sum_{K \in \mathscr{V}_{m}} u_{K}^{0} \psi_{K}^{0} m_{K} \\
& =-\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{K \in \mathscr{V}_{m}} u_{K}^{n+1} \frac{\psi_{K}^{n+1}-\psi_{K}^{n}}{\Delta t_{m}} m_{K}-\sum_{K \in \mathscr{V}_{m}} u_{K}^{0} \psi_{K}^{0} m_{K} \\
& =-\iint_{Q_{t_{\mathrm{f}}}} u_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t) \partial_{t} \psi_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, t) \mathrm{d} \mathbf{x} \mathrm{~d} t-\int_{\Omega} u_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, 0) \psi_{\mathscr{M}_{m}, \Delta t_{m}}(\mathbf{x}, 0) \mathrm{d} \mathbf{x}
\end{aligned}
$$

Thanks to the regularity of $\psi$, and the convergence in $L^{1}\left(Q_{t_{\mathrm{f}}}\right)$ of the sequence $\left(u_{\mathscr{M}_{m}, \Delta t_{m}}\right)_{m}$ towards $u$, it follows that (see e.g. [17])

$$
\mathscr{A}_{m} \longrightarrow-\iint_{Q_{\mathrm{f}_{\mathrm{f}}}} u(\mathbf{x}, t) \partial_{t} \psi(\mathbf{x}, t) \mathrm{d} \mathbf{x} \mathrm{~d} t-\int_{\Omega} u(\mathbf{x}, 0) \psi(\mathbf{x}, 0) \mathrm{d} \mathbf{x}, \quad \text { as } m \rightarrow \infty
$$

## Diffusion term

Let us first prove that $\lim _{m \rightarrow \infty} \mathscr{B}_{m}=0$.
For all $\sigma_{K L} \in \mathscr{E}_{m}$ and all $n \in\left\{0, \ldots, N_{m}\right\}$, we denote by $\bar{a}_{K L}^{n+1}$ the quantity defined by

$$
\bar{a}_{K L}^{n+1}=\left\{\begin{array}{cl}
\left(\frac{\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)}{u_{K}^{n+1}-u_{L}^{n+1}}\right)^{2} & \text { if } u_{K}^{n+1} \neq u_{L}^{n+1} \\
a\left(u_{K}^{n+1}\right) & \text { if } u_{K}^{n+1}=u_{L}^{n+1}
\end{array}\right.
$$

Then, the term $\mathscr{B}_{m}$ rewrites

$$
\mathscr{B}_{m}=\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}} \Lambda_{K L} \sqrt{a_{K L}^{n+1}}\left(\sqrt{a_{K L}^{n+1}}-\sqrt{\bar{a}_{K L}^{n+1}}\right)\left(u_{K}^{n+1}-u_{L}^{n+1}\right)\left(\psi_{K}^{n}-\psi_{L}^{n}\right)
$$

Now, using the Cauchy-Schwarz inequality, we get

$$
\left|\mathscr{B}_{m}\right| \leq\left(\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}}\left|\Lambda_{K L}\right| a_{K L}^{n+1}\left(u_{K}^{n+1}-u_{L}^{n+1}\right)^{2}\right)^{\frac{1}{2}} \times \mathscr{R}_{m}^{\frac{1}{2}}
$$

where, $\mathscr{R}_{m}$ is given by

$$
\mathscr{R}_{m}=\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{\sigma_{K L} \in \mathscr{E}_{m}}\left|\Lambda_{K L}\right|\left(\sqrt{a_{K L}^{n+1}}-\sqrt{\bar{a}_{K L}^{n+1}}\right)^{2}\left(\psi_{K}^{n}-\psi_{L}^{n}\right)^{2}
$$

Using Lemma 3.3 and Proposition 3.7, one has $\left|\mathscr{B}_{m}\right| \leq C \mathscr{R}_{m}{ }^{\frac{1}{2}}$. Hence, in order to prove that $\lim _{m \rightarrow \infty} \mathscr{B}_{m}=0$, it suffices to prove that $\lim _{m \rightarrow \infty} \mathscr{R}_{m}=0$.
For all $T \in \mathscr{T}_{m}$, we denote by

$$
\bar{\xi}_{T}^{n+1}=\max _{\mathbf{x} \in T}\left(\xi(p)_{\mathscr{T}_{m}, \Delta t_{m}}\left(\mathbf{x}, t^{n+1}\right)\right), \quad \underline{\xi}_{T}^{n+1}=\min _{\mathbf{x} \in T}\left(\xi(p)_{\mathscr{T}_{m}, \Delta t_{m}}\left(\mathbf{x}, t^{n+1}\right)\right)
$$

and for all $(\mathbf{x}, t) \in T \times\left(t^{n}, t^{n+1}\right)$, by

$$
\bar{\xi}_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)=\bar{\xi}_{T}^{n+1}, \quad \underline{\xi}_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)=\underline{\xi}_{T}^{n+1}
$$

Consider the uniform continuous function $\sqrt{a \circ \xi^{-1}}$ defined on the closed bounded interval $[0, \xi(1)]$, and let $\eta$ be its modulus of continuity, then we have

$$
\begin{equation*}
\left|\sqrt{a_{K L}^{n+1}}-\sqrt{\bar{a}_{K L}^{n+1}}\right| \leq \eta\left(\bar{\xi}_{T}^{n+1}-\underline{\xi}_{T}^{n+1}\right), \quad \text { for all } \sigma_{K L} \in \mathscr{E}_{T} \tag{56}
\end{equation*}
$$

Therefore, using this inequality in the definition of $\mathscr{R}_{m}$, we get

$$
\begin{equation*}
0 \leq \mathscr{R}_{m} \leq \mathscr{Q}_{m} \tag{57}
\end{equation*}
$$

where,

$$
\begin{equation*}
\mathscr{Q}_{m}=\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{T \in \mathscr{T}_{m}}\left(\eta\left(\bar{\xi}_{T}^{n+1}-\underline{\xi}_{T}^{n+1}\right)\right)^{2} \sum_{\sigma_{K L} \in \mathscr{E}_{T}}\left|\lambda_{K L}^{T}\right|\left(\psi_{K}^{n}-\psi_{L}^{n}\right)^{2} \tag{58}
\end{equation*}
$$

and $\lambda_{K L}^{T}$ is the constant defined by (24).
Thanks to Lemma 3.2, one can deduce that the inequality (57) implies that

$$
0 \leq \mathscr{R}_{m} \leq C \iint_{Q_{t_{\mathrm{f}}}} \eta\left(\bar{\xi}_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)-\underline{\xi}_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right) \mathrm{d} \mathbf{x} \mathrm{~d} t
$$

where $C$ is independent of $h_{m}$, and $\Delta t_{m}$. Therefore, it suffices to show that $\bar{\xi}_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)-\underline{\xi}_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t) \rightarrow 0$ a.e. in $Q_{t_{\mathrm{f}}}$ to consequently prove that $\lim _{m \rightarrow \infty} \mathscr{R}_{m}=0$. By a simple generalization of [6, Lemma A.1] and by the help of Lemma 3.1 and Proposition 3.7, it follows that

$$
\iint_{Q_{t_{\mathrm{f}}}}\left|\bar{\xi}_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)-\underline{\xi}_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right| \mathrm{d} \mathbf{x} \mathrm{~d} t \leq C h\left(\iint_{Q_{\mathrm{f}_{\mathrm{f}}}}\left|\nabla \xi(u)_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right|^{2} \mathrm{~d} \mathbf{x} \mathrm{~d} t\right)^{\frac{1}{2}} \leq C h
$$

As a consequence, up to a subsequence, one has

$$
\lim _{m \rightarrow \infty} \mathscr{B}_{m}=\lim _{m \rightarrow \infty} \mathscr{R}_{m}=\lim _{m \rightarrow \infty} \mathscr{Q}_{m}=0
$$

We now focus on the term $\mathscr{C}_{m}$ and prove that

$$
\lim _{m \rightarrow \infty} \mathscr{C}_{m}=\iint_{Q_{t_{\mathrm{f}}}} \Lambda(\mathbf{x}) a(u) \nabla u \cdot \nabla \psi \mathrm{~d} \mathbf{x} \mathrm{~d} t
$$

To do this, we introduce the term $\mathscr{C}_{m}^{\prime}$ defined by

$$
\mathscr{C}_{m}^{\prime}:=\iint_{Q_{t_{\mathrm{f}}}} \Theta_{\mathscr{T}_{m}, \Delta t_{m}} \Lambda(\mathbf{x}) \nabla \xi(u)_{\mathscr{T}_{m}, \Delta t_{m}} \cdot \nabla \psi_{\mathscr{T}_{m}, \Delta t_{m}}\left(\cdot, t-\Delta t_{m}\right) \mathrm{d} \mathbf{x} \mathrm{~d} t
$$

where $\Theta_{\mathscr{T}_{m}, \Delta t_{m}}$ is a piecewise constant function (on the triangular mesh) function given by

$$
\Theta_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)=\sqrt{a \circ \xi^{-1}}\left(\Upsilon_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right), \quad \forall \mathbf{x} \in T, \forall t \in\left(t^{n}, t^{n+1}\right], \forall T \in \mathscr{T}_{m},
$$

where $\Upsilon_{\mathscr{T}_{m}, \Delta t_{m}}$ is defined by

$$
\Upsilon_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)=\xi(u)_{\mathscr{T}_{m}, \Delta t_{m}}\left(\mathbf{x}_{T}, t\right), \quad \forall \mathbf{x} \in T, \forall t \in\left(t^{n}, t^{n+1}\right], \forall T \in \mathscr{T}_{m}
$$

Using again a slight generalization of [6, Lemma A.1] as well as the boundedness of the continuous function $\sqrt{a \circ \xi^{-1}}$, we obtain

$$
\begin{align*}
r_{\mathscr{T}_{m}, \Delta t_{m}} \longrightarrow \xi(u) & \text { in } L^{2}\left(Q_{t_{\mathrm{f}}}\right) \text { as } m \rightarrow \infty  \tag{59}\\
\Theta_{\mathscr{T}_{m}, \Delta t_{m}} \longrightarrow \sqrt{a(u)} & \text { in } L^{2}\left(Q_{t_{\mathrm{f}}}\right) \text { as } m \rightarrow \infty
\end{align*}
$$

It remains to verify that $\left|\mathscr{C}_{m}-\mathscr{C}_{m}^{\prime}\right| \longrightarrow 0$, when $m$ tends to infinity.
We denote by

$$
a_{T}^{n+1}=\left(\Theta_{\mathscr{T}_{m}, \Delta t_{m}}\left(\mathbf{x}_{T}, t^{n+1}\right)\right)^{2}, \quad \forall T \in \mathscr{T}_{m}, \forall n \in\left\{0, \ldots, N_{m}\right\}
$$

The discretization of the term $\mathscr{C}_{m}^{\prime}$ is written as

$$
\mathscr{C}_{m}^{\prime}=\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{T \in \mathscr{T}_{m}} \sqrt{a_{T}^{n+1}} \sum_{\sigma_{K L} \in \mathscr{E}_{T}} \lambda_{K L}^{T}\left(\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right)\left(\psi_{K}^{n}-\psi_{L}^{n}\right)
$$

Similar arguments as for obtaining inequality (56) yield

$$
\left|\sqrt{a_{K L}^{n+1}}-\sqrt{a_{T}^{n+1}}\right| \leq \eta\left(\bar{\xi}_{T}^{n+1}-\underline{\xi}_{T}^{n+1}\right), \quad \text { for all } \sigma_{K L} \in \mathscr{E}_{T}
$$

Therefore, using the Cauchy-Schwarz inequality, Lemma 3.1, Lemma 3.2, and Proposition 3.7, we deduce that there exists a constant $C$ does not depend on $h_{m}$ such that

$$
\begin{aligned}
\left|\mathscr{C}_{m}-\mathscr{C}_{m}^{\prime}\right|^{2} & \leq\left(\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{T \in \mathscr{T}_{m}} \eta\left(\bar{\xi}_{T}^{n+1}-\underline{\xi}_{T}^{n+1}\right) \sum_{\sigma_{K L} \in \mathscr{E}_{T}}\left|\lambda_{K L}^{T}\right|\left|\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right|\left|\psi_{K}^{n}-\psi_{L}^{n}\right|\right)^{2} \\
& \leq \mathscr{Q}_{m} \times \sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{T \in \mathscr{T}_{m}} \sum_{\sigma_{K L} \in \mathscr{E}_{T}}\left|\lambda_{K L}^{T}\right|\left|\xi\left(u_{K}^{n+1}\right)-\xi\left(u_{L}^{n+1}\right)\right|^{2} \leq C \mathscr{Q}_{m} \longrightarrow 0 \quad \text { as } m \rightarrow \infty .
\end{aligned}
$$

## Convection term

For all $T \in \mathscr{T}_{m}$, we define the piecewise constant function $\kappa_{\mathscr{T}_{m}, \Delta t_{m}}$ by

$$
\kappa_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)=\chi \circ \xi^{-1}\left(\Upsilon_{\mathscr{T}_{m}, \Delta t_{m}}(\mathbf{x}, t)\right), \quad \forall \mathbf{x} \in T, \forall t \in\left(t^{n}, t^{n+1}\right]
$$

Using the same guidelines as for the convergence results (59), one has

$$
\kappa_{\mathscr{T}_{m}, \Delta t_{m}} \longrightarrow \chi(u) \quad \text { in } L^{2}\left(Q_{t_{\mathrm{f}}}\right) \text { as } m \rightarrow \infty
$$

We introduce the term

$$
\mathscr{D}_{m}^{\prime}:=-\iint_{{Q_{\mathrm{f}}}} \kappa_{\mathscr{T}_{m}, \Delta t_{m}} \Lambda(\mathbf{x}) \nabla v_{\mathscr{T}_{m}, \Delta t_{m}} \cdot \nabla \psi_{\mathscr{T}_{m}, \Delta t_{m}}\left(\cdot, t-\Delta t_{m}\right) \mathrm{d} \mathbf{x} \mathrm{~d} t
$$

Thanks to the weakly convergence in $L^{2}\left(Q_{t_{\mathrm{f}}}\right)$ of the sequence $\nabla v_{\mathscr{T}_{m}, \Delta t_{m}}$ towards $\nabla v$, and to the uniform convergence of $\nabla \psi_{\mathscr{T}_{m}, \Delta t_{m}}$ towards $\nabla \psi$, we obtain

$$
\mathscr{D}_{m}^{\prime} \longrightarrow-\iint_{Q_{t_{\mathrm{f}}}} \chi(u) \Lambda(\mathbf{x}) \nabla v \cdot \nabla \psi \mathrm{~d} \mathbf{x} \mathrm{~d} t \quad \text { as } m \rightarrow \infty .
$$

Let us prove, using the same guidelines as before, that $\left|\mathscr{D}_{m}-\mathscr{D}_{m}^{\prime}\right| \longrightarrow 0$, when $m$ tends to infinity. We denote by

$$
\begin{array}{ll}
\chi_{T}^{n+1}=\kappa_{\mathscr{T}_{m}, \Delta t_{m}}\left(\mathbf{x}_{T}, t^{n+1}\right), & \forall T \in \mathscr{T}_{m}, \forall n \in\left\{0, \ldots, N_{m}\right\} \\
\mu_{T}^{n+1}=\mu_{\mathscr{T}_{m}, \Delta t_{m}}\left(\mathbf{x}_{T}, t^{n+1}\right), & \forall T \in \mathscr{T}_{m}, \forall n \in\left\{0, \ldots, N_{m}\right\}
\end{array}
$$

Therefore,

$$
\mathscr{D}_{m}-\mathscr{D}_{m}^{\prime}=\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{T \in \mathscr{T}_{m}} \sum_{\sigma_{K L} \in \mathscr{E}_{T}}\left(a_{T}^{n+1} \mu_{T}^{n+1}-a_{K L}^{n+1} \mu_{K L}^{n+1}\right) \lambda_{K L}^{T}\left(v_{K}^{n+1}-v_{L}^{n+1}\right)\left(\psi_{K}^{n}-\psi_{L}^{n}\right)
$$

Thanks to the the triangle inequality and to the existence of a continuity moduli $\eta$ and $\delta$ of the continuous functions $\sqrt{a \circ \xi^{-1}}$ and $\mu \circ \xi^{-1}$ respectively, one has

$$
\begin{aligned}
\left|a_{K L}^{n+1} \mu_{K L}^{n+1}-a_{T}^{n+1} \mu_{T}^{n+1}\right| & \leq \mu_{K L}^{n+1}\left|a_{K L}^{n+1}-a_{T}^{n+1}\right|+a_{T}^{n+1}\left(\mu_{K L}^{n+1}-\mu_{T}^{n+1}\right) \\
& \leq C\left(\eta\left(\bar{\xi}_{T}^{n+1}-\underline{\xi}_{T}^{n+1}\right)+\delta\left(\bar{\xi}_{T}^{n+1}-\underline{\xi}_{T}^{n+1}\right)\right),
\end{aligned}
$$

where the constant $C$ does not depend on $h_{m}$. Therefore, using the Cauchy-Schwarz inequality, Lemma 3.1, Lemma 3.2, and Proposition 3.7, we deduce that there exists a constant $C$ independent of $h_{m}$ such that

$$
\left|\mathscr{D}_{m}-\mathscr{D}_{m}^{\prime}\right|^{2} \leq C\left(\mathscr{Q}_{m}+\mathscr{W}_{m}\right) \times \sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{T \in \mathscr{T}_{m}} \sum_{\sigma_{K L} \in \mathscr{E}_{T}}\left|\lambda_{K L}^{T}\right|\left|v_{K}^{n+1}-v_{L}^{n+1}\right|^{2}
$$

where $\mathscr{Q}_{m}$ is given by equation (58), and $\mathscr{W}_{m}$ is given by

$$
\mathscr{W}_{m}=\sum_{n=0}^{N_{m}} \Delta t_{m} \sum_{T \in \mathscr{T}_{m}}\left(\delta\left(\bar{\xi}_{T}^{n+1}-\underline{\xi}_{T}^{n+1}\right)\right)^{2} \sum_{\sigma_{K L} \in \mathscr{E}_{T}}\left|\lambda_{K L}^{T}\right|\left(\psi_{K}^{n}-\psi_{L}^{n}\right)^{2}
$$

Now, using the same proof as for the diffusion term, one can deuce that $\mathscr{W}_{m} \leq C h_{m}$. Therefore

$$
\lim _{m \rightarrow \infty}\left|\mathscr{D}_{m}-\mathscr{D}_{m}^{\prime}\right|=0
$$

and consequently,

$$
\lim _{m \rightarrow \infty} \mathscr{D}_{m}=-\iint_{Q_{t_{\mathrm{f}}}} \chi(u) \Lambda(\mathbf{x}) \nabla v \cdot \nabla \psi \mathrm{~d} \mathbf{x} \mathrm{~d} t
$$

## Reaction term

We would now like to show that

$$
\mathscr{F}_{m} \longrightarrow \iint_{Q_{t_{\mathrm{f}}}} f(u(\mathbf{x}, t)) \psi(\mathbf{x}, t) \mathrm{d} \mathbf{x} \mathrm{~d} t \quad \text { as } m \rightarrow \infty
$$

For this purpose, we denote, for all $K \in \mathscr{V}_{m}$ and for all $n \geq 1$, by $f_{K}^{n}=f\left(u_{K}^{n}\right)$, and by $f_{\mathscr{M}_{m}, \Delta t_{m}}$ the piecewise constant reconstruction in $\mathscr{X}_{\mathscr{M}_{m}, \Delta t_{m}}$. Thus we have

$$
\mathscr{F}_{m}=\iint_{Q_{t_{\mathrm{f}}}} f_{\mathscr{M}_{m}, \Delta t_{m}} \psi_{\mathscr{M}_{m}, \Delta t_{m}}\left(\cdot, t-\Delta t_{m}\right) \mathrm{d} \mathbf{x} \mathrm{~d} t \longrightarrow \iint_{Q_{t_{\mathrm{f}}}} f(u(\mathbf{x}, t)) \psi(\mathbf{x}, t) \mathrm{d} \mathbf{x} \mathrm{~d} t \quad \text { as } m \rightarrow \infty
$$

since $f(u)_{\mathscr{M}_{m}, \Delta t_{m}}$ converges strongly in $L^{2}\left(Q_{t_{\mathrm{f}}}\right)$ towards $f(u)$, and as $\psi_{\mathscr{M}_{m}, \Delta t_{m}}$ converges uniformly towards $\psi$. This ends the proof of Theorem 2.1.

## 6 Numerical results

In this section, we establish various 2-D numerical results provided by the nonlinear CVFE scheme (14), (21). Newton's algorithm is carried out for the implementation of the scheme, coupled with a biconjugate gradient method to solve linear systems arising from the Newton algorithm. We provide three tests to show the effectiveness of the nonlinear CVFE scheme (14), (21). For these tests, we consider the following data: $L_{\mathbf{x}}=1, L_{\mathbf{y}}=1$ (the length and the width of the domain). We fix: $\Delta t=0.002, \alpha=0.01, \beta=0.05, a(u)=$ $d_{u} u(1-u), d_{u}=0.0005, \chi(u)=\zeta \times(u(1-u))^{2}, \zeta=0.05$. By definition, we have $\mu(u)=\frac{\zeta}{d_{u}} u(1-u)$ then, the numerical flux function $\mu_{K L}^{n+1}$ is given using the following functions:

$$
\mu_{\uparrow}(z)=\mu\left(\min \left\{z, \frac{1}{2}\right\}\right), \quad \text { and } \mu_{\downarrow}(z)=\mu\left(\max \left\{z, \frac{1}{2}\right\}\right)-\mu\left(\frac{1}{2}\right), \quad \forall z \in(0,1) \times(0,1)
$$

Unless stated otherwise and throughout the tests, we assume that $f(u)=0$, that the initial conditions are defined by regions, and we assume zero-flux boundary conditions. For instance, the cell density is initially defined by $u_{0}(\mathbf{x}, \mathbf{y})=1$ in the square region given by $(\mathbf{x}, \mathbf{y}) \in[0.45,0.55]$ and 0 otherwise. The initial chemeoattractant concentration is defined by $v_{0}(\mathbf{x}, \mathbf{y})=5$ in the space region given by $(\mathbf{x}, \mathbf{y}) \in$ $[0.2,0.3] \times[0.45,0.55] \cup[0.45,0.55] \times[0.2,0.3] \cup[0.45,0.55] \times[0.7,0.8] \cup[0.7,0.8] \times[0.45,0.55]$.

Test 1 (Weak anisotropic case). In this test, we assume that the diffusion tensors are given by

$$
\Lambda(\mathbf{x})=\left(\begin{array}{ll}
1 & 0 \\
0 & \theta
\end{array}\right), \quad D(\mathbf{x})=d\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right), \quad d=0.0001
$$

Further, we consider an admissible triangular primary mesh made of 14336 triangles, the corresponding Donald dual mesh consists of 7297 dual control volumes. In a admissible triangular mesh, all the angles of triangles are acute, then one can deduce that the maximum principle is verified for $v$ since the transmissibility coefficients are nonnegative, which it is not the case for $u$. In Tab. 1, we present minimum and maximum values obtained with each of the scheme (14)-(15), the nonlinear CVFE scheme (14),(21), and the finite volume scheme.

|  |  | scheme (14)-(15) | scheme (14),(21) | FV scheme |
| :---: | :---: | :---: | :---: | :---: |
| After 1 iteration | Min. Val. u | 0.0 | 0.0 | 0.0 |
| $\theta=1$ | Max. Val. u | 1.0 | 1.0 | 1.0 |
| After 10 iterations | Min. Val. u | 0.0 | 0.0 | 0.0 |
| $\theta=1$ | Max. Val. u | 0.971110 | 1.0 | 1.0 |
| After 1 iteration | Min. Val. u | $-1.73001 \times 10^{-3}$ | $8.68789 \times 10^{-20}$ | - |
| $\theta=5$ | Max. Val. u | 0.99722922 | 1.0 | - |
| After 10 iterations | Min. Val. u | $-1.62500 \times 10^{-2}$ | 0.00 | - |
| $\theta=5$ | Max. Val. u | 0.9715705 | 1.0 | - |
| After 1 iteration | Min. Val. u | $-4.46953 \times 10^{-3}$ | $6.30555 \times 10^{-16}$ | - |
| $\theta=10$ | Max. Val. u | 1.00018368 | 1.0 | - |
| After 10 iterations | Min. Val. u | $-3.91245 \times 10^{-2}$ | $6.30554 \times 10^{-16}$ | - |
| $\theta=10$ | Max. Val. u | 0.98342428 | 0.9999999 | - |

Table 1 Numerical results after 1 and 10 iterations.


Fig. 2 Meshes: admissible mesh for Test 1(left), initial primal mesh for Test 2 and 3 (center) and barycentric dual mesh for Test 2 and 3 (right).

Test 2 (Weak anisotropic case/obtuse angles). In this test, we consider a general unstructured mesh that contains obtuse angles, this mesh is made of 5193 triangles and 2665 dual control volumes. The discrete maximum principle is not guarantied for $v$, hence we cannot expect the maximum principle for $u$ since the computation of $u$ depends on the values of $v$, for that we consider the nonlinear discretization (21) of $v$.


Fig. 3 Initial condition for the cell density $u$ (left) with $0 \leq u \leq 1$ and for the chemeoattractant concentration $v$ (right) with $0 \leq v \leq 5$.

The diffusion tensors are defined, for all $\mathbf{x} \in(0,1) \times(0,1)$, by

$$
\Lambda(\mathbf{x})=\left(\begin{array}{cc}
7 & 2 \\
2 & 10
\end{array}\right), \quad D(\mathbf{x})=d\left(\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right), \quad d=0.0001
$$

Figure 3 represents initial distributions of the cell density $u$ and the chemeoattractant concentration $v$ over the initial triangular mesh as well as the corresponding dual mesh.

Figures 4-5 represent the evolution of the cell density at time $t=0.4, t=1.4, t=2.4$, and $t=4$. At moment $t=0.4$, it is clear that the cell density diffuses in the space without any interactions with the chemeoattractant which diffuses uniformly in the space. Then, after a while, and when the chemeoattractant diffusion reaches the cell density location, we see that the latter changes its direction to be absorbed by the chemeoattractant located vertically. This process continues and the cells accumulate into the location of the chemeoattractant and finally we obtain the cell density aggregations as shown at $t=4$.


Fig. 4 Evolution of the cell density $u$ at time $t=0.4$ with $0 \leq u \leq 0.667$ (left), and at time $t=1.4$ with $0 \leq u \leq 0.632$ (right).


Fig. 5 Evolution of the cell density $u$ at time $t=2.4$ with $0 \leq u \leq 0.972$ (left), and at time $t=4$ with $0 \leq u \leq 0.987$ (right).

Test 3 (Anisotropic case/obtuse angles). In this test, we consider an unstructured mesh consisting of 15 568 primal triangles and 7912 dual control dual volumes. Further, we assume that the diffusion tensors are anisotropic and are given by:

$$
\Lambda(\mathbf{x})=\left(\begin{array}{cc}
8 & -7 \\
-7 & 20
\end{array}\right), \quad D(\mathbf{x})=d\left(\begin{array}{ll}
1 & 0 \\
0 & 3
\end{array}\right), \quad d=0.0001
$$

|  |  | CVFE scheme: (14)-(15) | CVFE scheme: (14),(21) |
| :---: | :---: | :---: | :---: |
| After 1 iteration | Min. Val. u | 0.0 | 0.0 |
|  | Max. Val. u | 1.0 | 1.0 |
|  | Min. Val. v | $-1.141912 \mathrm{E}-002$ | $1.922764 \mathrm{E}-051$ |
|  | Max. Val. v | 5.012383 | 4.999982 |
| After 200 iterations | Min. Val. u | 0.0 | 0.0 |
|  | Max. Val. u | 0.5298226 | 0.5312562 |
|  | Min. Val. v | $-1.731068 \mathrm{E}-003$ | $1.297192 \mathrm{E}-080$ |
|  | Max. Val. v | 4.8053827 | 4.8018742 |
| After 1000 iterations | Min. Val. u | 0.0 | 0.0 |
|  | Max. Val. u | 0.9957580 | 0.9974757 |
|  | Min. Val. v | $6.265859 \mathrm{E}-023$ | $3.171769 \mathrm{E}-080$ |
|  | Max. Val. v | 2.961761 | 2.910828 |

Table 2 Numerical results after 1, 200 and 1000 iterations over an unstructured mesh with obtuse angles.

Table 2 provides a comparison between the nonlinear CVFE scheme coupled on the one hand with the discretization (15) of $v$ and with the discretization (21) of $v$ on the other hand. We see that the discretization (21) carries out a better approximation than the discretization (15) in terms of ensuring the discrete maximum principle property.

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