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# **A computationally effective dynamic hysteresis model taking into account skin effect in magnetic laminations**

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2 **Abstract**

3 We propose a simplified dynamic hysteresis model for the prediction of magnetization behavior of  
4 electrical steel up to high frequencies, taking into account the skin effect. This model has the advantage  
5 of predicting the hysteresis loop and loss behavior versus frequency with the same accuracy provided  
6 by the Dynamic Preisach Model with a largely reduced computational burden. It is here compared to  
7 experimental results obtained in Fe-Si laminations under sinusoidal flux up to 2 kHz.

8

## 9 I. INTRODUCTION

10 The problem of high frequency behavior of magnetic laminations is of primary importance in  
11 modern electrical engineering problems [1], but difficult problems arise in the prediction and  
12 assessment of magnetization process and losses, because the skin effect compounds with the nonlinear  
13 hysteretic response of the material. In order to cope with the inhomogeneous profile of the flux density  
14 over the lamination thickness, the magnetic loss is generally calculated by numerically solving the  
15 diffusion equation over the sample cross-section, and using a dynamic hysteresis model for the material  
16 constitutive law [2]. The Dynamic Preisach Model (DPM) is the model of choice, because it is accurate  
17 and solidly established from the physical viewpoint [3][4]. Its application is, however, particularly time  
18 consuming, because calculations must be done for each finite element of the spatial mesh until  
19 convergence is reached. A faster approach, preserving the special virtues of DPM, would therefore be  
20 appropriate. We apply in this paper the DPM to the broadband behavior of nonoriented Fe-Si  
21 laminations through a simplified method, drastically reducing computing time and complexity of the  
22 full method, requiring the computation of the dynamics of each elementary hysteron distributed in the  
23 Preisach plane [5]. We start our discussion from the differential relation found by Bertotti [6] (formula  
24 (9), page 4609, of reference [6]) for the excess magnetic field due to the dynamic behavior of the  
25 magnetization process

$$H_{\text{exc}}(t) \equiv H(t) - H_{\text{stat}}(t) = \text{sign}(\dot{H}_{\text{stat}}) \frac{4}{3} \sqrt{\frac{|\dot{H}|}{k_d}}, \quad (1)$$

26 where  $H(t) = H_a(t) - H_{\text{cl}}(t)$  is the difference between the applied field  $H_a(t)$  and the counterfield  $H_{\text{cl}}(t)$   
27 generated by the macroscopic eddy currents (classical field).  $H_{\text{stat}}(t)$  is the field that would provide  
28 under static conditions the same irreversible polarization  $J_{\text{irr}}$  and  $k_d$  [ $\text{A}^{-1}\text{s}^{-1}\text{m}$ ] is the DPM constant. Eq.  
29 (1) is derived under the simplifying assumptions of triangular  $H_a(t)$  and uniform Preisach distribution  
30 function. The extent to which such a restriction can be circumvented and the full DPM approach for

31 generic exciting conditions and shape of the Preisach density function can be approximated will be  
 32 discussed in the following. We find first that a numerical analysis based on (1) (Model 1) does not lead  
 33 to highly accurate dynamic loop shapes, especially at high frequencies. A slight modification of (1) is  
 34 therefore proposed (Model 2) and implemented in a non-linear magneto-dynamical model for the  
 35 computation of the field distribution inside Fe-Si laminations, taking into account skin effect. The  
 36 numerical results are compared with the experiments performed under sinusoidal flux up to 2 kHz.

## 37 II. THE SIMPLIFIED MODEL

### 38 A. The simplified DPM-Model 2

39 The Preisach distribution function, including the reversible contribution, was determined in a 0.35  
 40 mm thick Fe-(3.2wt%)Si-(0.5wt%)Al lamination, and the dynamic constant  $k_d=350 \text{ A}^{-1}\text{s}^{-1}\text{m}$  was  
 41 identified. The full and the simplified (Model 1) DPM are compared in terms of the excess field. In this  
 42 model benchmark, a sinusoidal dynamic field  $H(t) = H_a(t) - H_{cl}(t)$  (peak value  $H_p = 100\text{A}\cdot\text{m}^{-1}$ ,  
 43 frequency  $f = 200 \text{ Hz}$ ) has been applied. For the full DPM case, the excess field  $H_{\text{exc}}(t) = H(t) - H_{\text{stat}}(t)$   
 44 is derived applying the DPM to  $H(t)$  in order to get the irreversible polarization  $J_{\text{irr}}(t)$ , and then using  
 45 the inverse static Preisach model to compute the static field  $H_{\text{stat}}(t)$  from  $J_{\text{irr}}(t)$ . For the model 1, a  
 46 numerical solution of (1) directly provides  $H_{\text{stat}}(t)$  from the sinusoidal  $H(t)$ . A comparison between the  
 47 results of two approaches is illustrated in Fig. 1 (the  $J_{\text{irr}}$  waveform obtained from the DPM has also  
 48 been represented). Discrepancies are found between the two predicted  $H_{\text{exc}}(t)$  waveforms, particularly  
 49 around the reversal points of the irreversible magnetization  $J_{\text{irr}}$ . With the simplified Model 1 the zero of  
 50  $H_{\text{exc}}$  is, according to (1), coincident with the maximum of  $H(t)$ , whereas from the same picture, it  
 51 appears that it occurs when  $J_{\text{irr}}$  is maximum. Consequently, (1) is formulated as

$$H_{\text{exc}}(t) \equiv H(t) - H_{\text{stat}}(t) = \text{sign}(\dot{H}_{\text{stat}}) \frac{4}{3} \sqrt{\frac{|\dot{H}_{\text{stat}}|}{k_d}}, \quad (2)$$

52 on account of the fact that the sign of  $\dot{H}_{\text{stat}}$  is the same as that of  $\dot{J}_{\text{irr}}$ . The simplified DPM based on

53 (2) (DPM-Model 2) appears now to provide (Fig. 1) an  $H_{\text{exc}}(t)$  waveform in good agreement with the  
 54 one provided by the full DPM.

55 Once the static field is known,  $J_{\text{irr}}(H_{\text{stat}})$  is computed by means of the Static Preisach Model, while  
 56 the reversible component  $J_{\text{rev}}(H)$  is calculated ignoring any dynamic effect linked to the reversible  
 57 contribution [1]. This procedure, which permits to obtain the constitutive law of the material  $J(H)$ , is  
 58 summed up in Fig. 2. An example of hysteresis loop prediction (in this case with nested minor loop) is  
 59 illustrated in Fig. 3, confirming the good agreement between the results provided by the full DPM and  
 60 the simplified DPM-Model 2.

#### 61 A. Numerical implementation of DPM-Model 2

62 This model requires solving the non-linear differential equation (2) for  $H_{\text{stat}}$  knowing the field  $H$ .  
 63 This can be done putting (2) under an equivalent canonical form and adding the periodicity conditions  
 64 (period  $T=1/f$ ):

$$\begin{cases} \dot{H}_{\text{stat}} = \text{sign}(H(t) - H_{\text{stat}}(t)) \cdot \frac{9k_d}{16} \cdot (H(t) - H_{\text{stat}}(t))^2 \\ H_{\text{stat}}(0) = H_{\text{stat}}(T) \end{cases} \quad (3)$$

65 and numerically solving it using a Runge-Kutta method. The application of this newly defined  
 66 dynamical model to the computation of flux distribution inside the steel lamination can dramatically  
 67 reduce the computational burden, as demonstrated in the next section.

### 68 III. MAGNETO DYNAMICAL MODELING AND EXPERIMENTAL

#### 69 A. Magneto-dynamical model of the lamination

70 The solution of the diffusion equation over the lamination thickness with a dynamic hysteretic  
 71 constitutive law requires a special numerical treatment of the non-linearity using the Fixed Point  
 72 method [3][7]. The method proposed in [3] has been here implemented. The diffusion problem on the  
 73 lamination thickness is one-dimensional, with the spatial coordinate  $x$  ranging over the lamination

74 cross-section  $[-e/2; e/2]$  (where  $e = 0.35$  mm is the lamination thickness). To solve the diffusion  
 75 equation, the non-linearity of the material constitutive law  $J(H)$  is contained in a spatio-temporal  
 76 function  $R(x,t)$ , called the residual [3] and the relationship between  $H$  and  $J$  is written as

$$H(x,t) = \nu_{\text{FP}} \cdot B(x,t) + R(x,t), \quad (4)$$

77 where  $\nu_{\text{FP}}$  is a properly chosen constant, ensuring convergence [1]. Using (4), the diffusion equation  
 78 can be written, with imposed mean flux density  $B_{\text{MEAN}}(t)$ , in terms of vector potential  $A$  on a half  
 79 lamination  $[0; e/2]$ :

$$\begin{cases} \nu_{\text{FP}} \cdot \frac{\partial^2 A}{\partial x^2} - \sigma \frac{\partial A}{\partial t} = \frac{\partial R}{\partial x} \\ A(0,t) = 0 \text{ and } A(e/2,t) = -B_{\text{MEAN}}(t) \cdot e/2 \\ A(x,t) \text{ is } T\text{-periodic} \end{cases} \quad (5)$$

80 The time derivative and the time periodic condition are dealt with using temporal Fourier series [3].  
 81 More precisely, the residual function  $R(x,t)$  is decomposed into complex Fourier series for what  
 82 concerns the time dependence. Equation (5) is then solved for each time harmonic, the time derivative  
 83 becoming an algebraic multiplication by the corresponding harmonic pulsation. An inverse Fourier  
 84 transformation permits to retrieve the function  $A(x,t)$ . For each time harmonic, the spatial second order  
 85 derivative and the boundary conditions are dealt with using a numerical finite difference scheme (the  
 86 half lamination is subdivided into 50 intervals). At the beginning of the iterative procedure, the residual  
 87  $R$  is initialized to zero and at each iteration step (index number  $i$ ), the following process is performed:

88 1. Knowing the residual at the previous stage  $R^{(i-1)}(x,t)$ , the differential equation (5) is solved to  
 89 compute the vector potential  $A^{(i)}(x,t)$ . The magnetic field is obtained as

$$90 \quad H^{(i)}(x,t) = -\nu_{\text{FP}} \frac{\partial A^{(i)}}{\partial x} + R^{(i-1)}(x,t).$$

91 2. A dynamic hysteresis model, providing the dynamic constitutive law  $J(H)$ , is used to evaluate  $J^{(i)}$

92 from  $H^{(i)}$  and the new induction value  $B^{(i)}(x,t)=J^{(i)}(x,t)+\mu_0\cdot H^{(i)}(x,t)$  is calculated. In [3], the dynamic  
 93 hysteresis model is the DPM, and is replaced in this paper by the simplified model proposed in Fig.  
 94 2.

95 3. The new residual is then computed by  $R^{(i)}(x,t)=H^{(i)}(x,t)-v_{FP}\cdot B^{(i)}(x,t)$ . The process is repeated  
 96 until convergence of the residuals is obtained.

## 97 B. Results

98 The dynamic hysteretic constitutive law  $J(H)$  of the material, which was in [3] given by the full  
 99 DPM, has been here replaced by the model proposed in Fig. 2 where the Model 2 of the local excess  
 100 field is applied. The Preisach distribution function has been identified following the method proposed  
 101 in [1], in the previous 0.35 mm thick Fe-(3.2wt%)Si-(0.5wt%)Al lamination (conductivity  $\sigma =$   
 102  $1.773\cdot 10^6 \text{ S}\cdot\text{m}^{-1}$ ,  $k_d=350 \text{ A}^{-1}\text{s}^{-1}\text{m}$ ). The magnetic loss and hysteresis loops have been measured under  
 103 sinusoidal flux at different frequencies and peak inductions up to 2 kHz. The magnetic flux distribution  
 104 and loss have been computed by applying the dynamical model, using both the full and the simplified  
 105 DPM (Model 2). Fig. 4 provides an example of computed and experimental loops ( $J_p = 0.5 \text{ T}$ ,  
 106 frequency  $f = 1 \text{ kHz}$ ). In Fig. 5 the comparison is made for energy loss versus frequency behavior. It  
 107 shows the good predicting capability of the simplified dynamic Model 2, in spite of a computing time  
 108 reduced by a factor around 60 with respect to the full DPM. Fig. 5 also shows the predicted loss  
 109 behavior if the skin effect is ignored. In this case the induction is assumed uniform across the sample  
 110 thickness and equal to the mean flux density  $B_{\text{MEAN}}(t)$ . The classical loss is thus calculated according to

$$W_{\text{class}} = \frac{\sigma e^2}{12} \int_0^T \left( \frac{dB_{\text{MEAN}}}{dt} \right)^2 dt \quad (6)$$

111 It is found that ignoring the skin effect brings about a large overestimation of the loss above about  $f =$   
 112 400 Hz, for the considered  $B_{\text{MEAN}}(t)$  peak value of 0.5 T considered in Fig. 5. An example of calculated  
 113 induction profile  $B(x, t_0)$  across the lamination thickness at a given instant of time  $t_0$  is shown in Fig. 6.



114 It refers to  $f = 1$  kHz and a time  $t_0$  corresponding to the average flux density  $B_{\text{MEAN}}(t_0) = 0$ . Again, the  
115 full and simplified DPMs predict very close results.

#### 116 **IV. CONCLUSIONS**

117 In this work we have discussed a novel dynamic model for energy losses and hysteresis loops in soft  
118 magnetic laminations based on simplified treatment of the Dynamic Preisach Model (DPM). It is  
119 applied on experiments performed up to the kHz range in Fe-Si sheets, in the presence of substantial  
120 skin effect, showing good agreement both with the experimental results and the prediction made using  
121 the full machinery of the DPM. A remarkable advantage in computing time, which is reduced by a  
122 factor around 60, is obtained substituting the full DPM with the present model.

123

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- 138

139 **Figure captions**

140 Fig. 1- Excess fields provided by the full DPM, Model 1, and Model 2, for a sinusoidal dynamic  
141 field  $H(t) = H_a(t) - H_{cl}(t)$  (frequency  $f = 200$  Hz, field peak value  $H_p = 100\text{A}\cdot\text{m}^{-1}$ ) and corresponding  
142 irreversible polarization  $J_{irr}$  provided by the DPM in a 0.35 mm thick Fe-(3.2wt%)Si-(0.5wt%)Al  
143 lamination (dynamic constant  $k_d=350\text{ A}^{-1}\text{s}^{-1}\text{m}$ ).

144 Fig. 2 - Block diagram of the simplified constitutive law  $J(H)$  of the material, using the Model 2 as a  
145 link between the field  $H$  and the static field  $H_{stat}$  (computation of the local excess field).

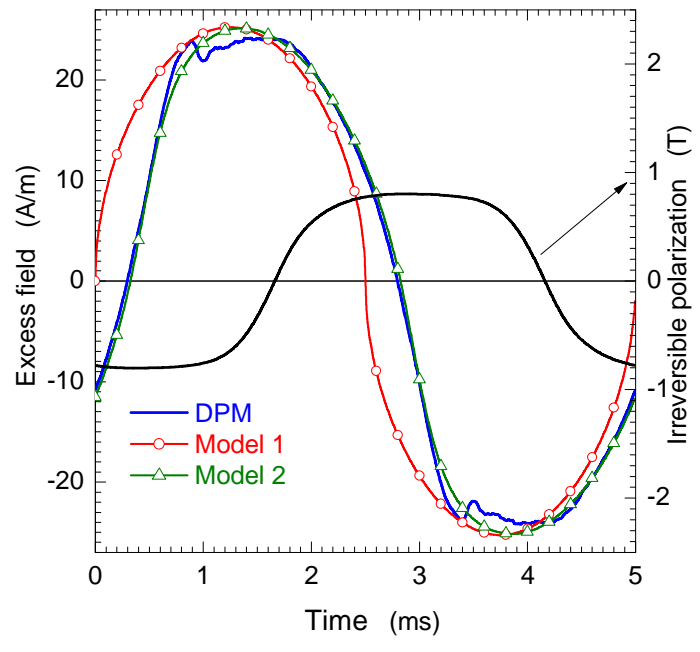
146 Fig. 3 - Example of the reconstructed  $J(H)$  hysteresis loop with nested minor loops ( $f = 200$  Hz,  $H_p =$   
147  $150\text{A}\cdot\text{m}^{-1}$ )

148 Fig. 4 – Experimental and theoretical hysteresis loops at  $f = 200$  Hz in the 0.35 mm thick Fe-  
149 (3.2wt%)Si-(0.5wt%)Al lamination for sinusoidal induction ( $B_p=0.5$  T). The measurements have been  
150 performed on Epstein strips. Closed results are obtained by the full and the simplified DPM.

151 Fig. 5 – Same as Fig. 4 for the energy loss versus magnetizing frequency (DC – 2 kHz).

152 Fig. 6: Instantaneous induction profile  $B(x,t_o)$  across the lamination thickness as obtained by the full  
153 and the simplified (Model 2) DPM ( $f = 1$  kHz,  $-0.175\text{ mm} \leq x \leq 0.175\text{ mm}$ ). The time  $t_o$  considered in  
154 this figure is the one for which the mean flux density  $B_{MEAN}(t_o) = 0$ .

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Fig. 1

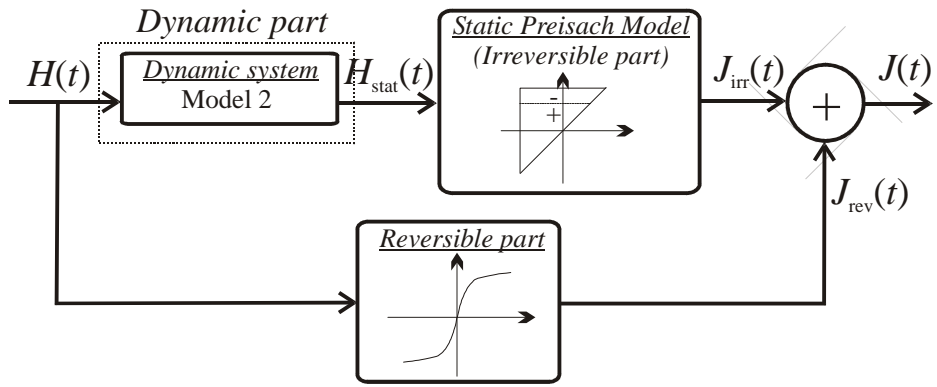


Fig. 2

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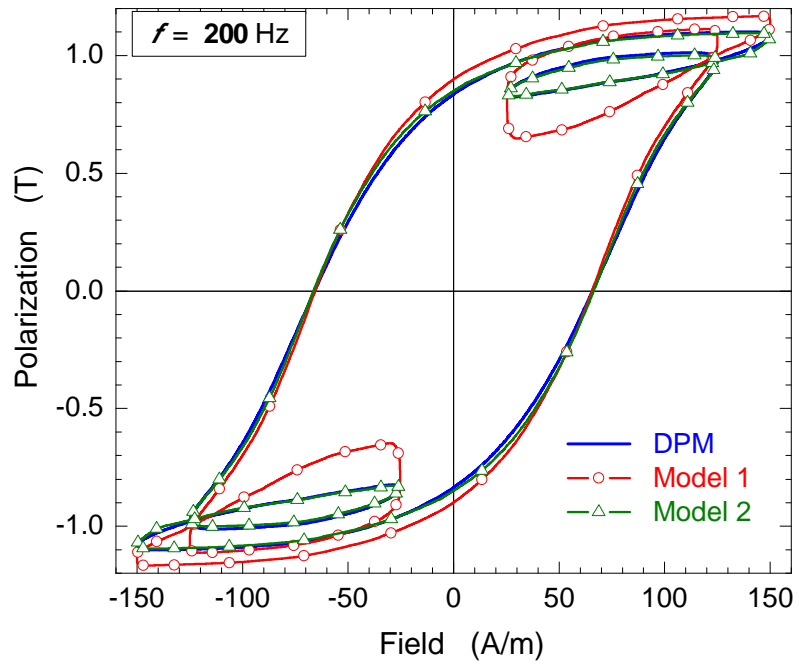


Fig. 3

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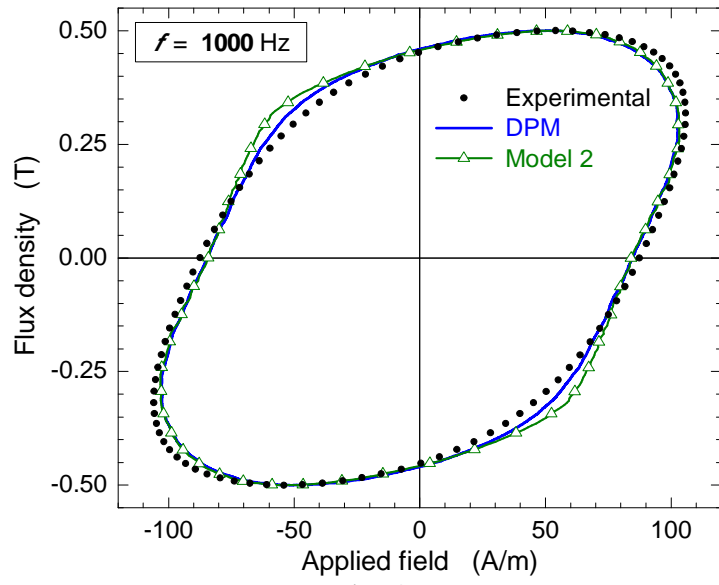


Fig. 4

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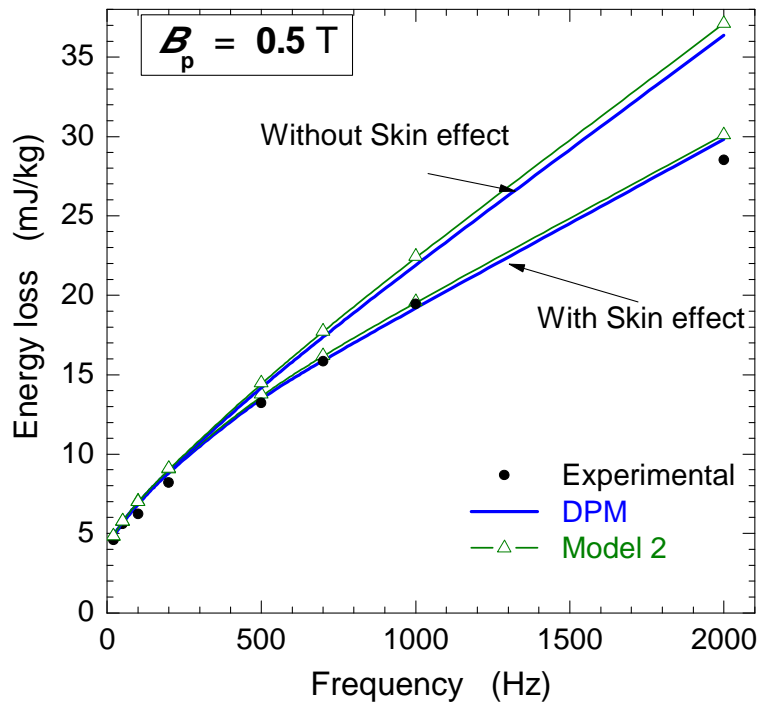
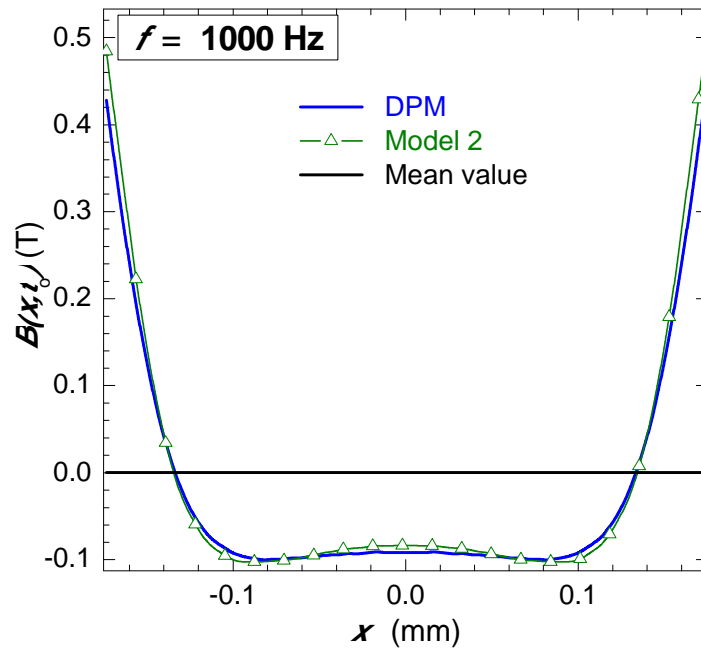


Fig. 5

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Fig. 6