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Determining the waveguide conductivity in a hyperbolic equation from a single measurement on the lateral boundary

¹MICHEL CRISTOFOL, ²SHUMIN LI, ³ERIC SOCCORSI

Abstract

We consider the multidimensional inverse problem of determining the conductivity coefficient of a hyperbolic equation in an infinite cylindrical domain, from a single boundary observation of the solution. We prove Hölder stability with the aid of a Carleman estimate specifically designed for hyperbolic waveguides.

AMS 2010 Mathematics Subject Classification: 35R30.

Keywords: Inverse problem, hyperbolic equation, conductivity, Carleman estimate, infinite cylindrical domain.

1 Statement of the problem and results

1.1 Introduction

The present paper deals with the inverse problem of determining the time-independent isotropic conductivity coefficient $c : \Omega \rightarrow \mathbb{R}$ appearing in the hyperbolic partial differential equation $\partial_t^2 - \nabla \cdot c \nabla = 0$, where $\Omega := \omega \times \mathbb{R}$ is an infinite cylindrical domain whose cross section ω is a bounded open subset of \mathbb{R}^{n-1} , $n \geq 2$. Namely, $\ell > 0$ being arbitrarily fixed, we seek Hölder stability in the identification of c in $\Omega_\ell := \omega \times (-\ell, \ell)$ from the observation of u on the lateral boundary $\Gamma_L := \partial\omega \times (-L, L)$ over the course of time $(0, T)$, for $L > \ell$ and $T > 0$ sufficiently large.

Several stability results in the inverse problem of determining one or several unknown coefficients of a hyperbolic equation from a finite number of measurements of the solution are available in the mathematics literature [1, 5, 6, 7, 17, 18, 26, 29]. Their derivation relies on Bukhgeim-Klibanov's method [9], which is by means of a Carleman inequality specifically designed for hyperbolic systems. More precisely, [17, 29] study the determination of the zero-th order term $p : \Omega \rightarrow \mathbb{R}$ appearing in $\partial_t - \Delta + p = 0$, while [1, 7] deal with the identification of the speed $c : \Omega \rightarrow \mathbb{R}$ in the hyperbolic equation $\partial_t - cA = 0$ where $A = A(x, D_x)$ is a second order differential operator. The case of a principal matrix term in the divergence form, arising from anisotropic media, was treated by Bellassoued, Jellali and Yamamoto in [5], using the full data (i.e. the measurements are performed on the whole boundary). Using the FBI transform Bellassoued and Yamamoto claimed logarithmic stability in [6] from arbitrarily small

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boundary observations. Imanuvilov and Yamamoto derived stability results in [18] by means of H^{-1} Carleman inequality, from data observation on subdomains fulfilling specific geometric assumptions. In [26] Klibanov and Yamamoto employed a different approach inspired by [25].

Similarly, numerous authors have used the Dirichlet-to-Neumann operator to claim stability in the determination of unknown coefficients of a hyperbolic equation. We refer to [4, 19, 28] for a non-exhaustive list of such references.

In all the above mentioned papers, the domain was bounded. Recently, the Bukhgeim-Klibanov method was adapted to the framework of infinite quantum cylindrical domains in [10, 23, 24]. Inverse boundary value problems stated in unbounded waveguides were also studied in [11, 12, 13] with the help of the Dirichlet-to-Neumann operator. In all the six previous articles, the observation is taken on the infinitely extended lateral boundary of the waveguide. The approach developed in this paper is rather different in the sense that we aim to retrieve any arbitrary bounded subpart of the non-compactly supported conductivity c from one data taken on a compact subset of the lateral boundary. This requires that suitable smoothness properties of the solution to (1.1) be preliminarily established in the context of the unbounded domain Ω .

The paper is organized as follows. Section 2 is devoted to the analysis of the direct problem associated with the hyperbolic system under study. In Section 3 we prove a global Carleman estimate specifically designed for hyperbolic systems in the cylindrical domain Ω . Finally Section 4 contains the analysis of the inverse problem and the proof of the main result.

1.2 Settings

1.2.1 Notations

Throughout this text we write $x = (x', x_n) \in \Omega$ for every $x' := (x_1, \dots, x_{n-1}) \in \omega$ and $x_n \in \mathbb{R}$. Further, we denote by $|y| := \left(\sum_{i=1}^m y_i^2\right)^{1/2}$ the Euclidian norm of $y = (y_1, \dots, y_m) \in \mathbb{R}^m$, $m \in \mathbb{N}^*$, and we put $\mathbb{S}^{n-1} := \{x' = (x_1, \dots, x_{n-1}) \in \mathbb{R}^{n-1}, |x'| = 1\}$.

For the sake of shortness we write ∂_j for $\partial/\partial x_j$, $j \in \mathbb{N}_{n+1}^* := \{m \in \mathbb{N}^*, m \leq n+1\}$. For convenience the time variable t is sometimes denoted by x_{n+1} so that $\partial_t = \partial/\partial t = \partial_{n+1}$. We set $\nabla := (\partial_1, \dots, \partial_n)^T$, $\nabla_{x'} := (\partial_1, \dots, \partial_{n-1})^T$ and $\nabla_{x,t} = (\partial_1, \dots, \partial_n, \partial_t)^T$.

For any open subset D of \mathbb{R}^m , $m \in \mathbb{N}^*$, we denote by $H^p(D)$ the p -th order Sobolev space on D for every $p \in \mathbb{N}$, where $H^0(D)$ stands for $L^2(D)$. We write $\langle \cdot, \cdot \rangle_{p,D}$ (resp., $\|\cdot\|_{p,D}$) for the usual scalar product (resp., norm) in $H^p(D)$ and we denote by $H_0^1(D)$ the closure of $C_0^\infty(D)$ in the topology of $H^1(D)$.

Finally, for $d > 0$ we put $\Omega_d := \omega \times (-d, d)$, $Q_d := \Omega_d \times (0, T)$, $\Gamma_d := \partial\omega \times (-d, d)$ and $\Sigma_d := \partial\omega \times (-d, d) \times (0, T)$.

1.2.2 Statement of the problem

We examine the following initial boundary value problem (IBVP in short)

$$\begin{cases} \partial_t^2 u - \nabla \cdot c \nabla u = 0 & \text{in } Q := \Omega \times (0, T), \\ u(\cdot, 0) = \theta_0, \partial_t u(\cdot, 0) = \theta_1 & \text{in } \Omega, \\ u = 0 & \text{on } \Sigma := \Gamma \times (0, T), \end{cases} \quad (1.1)$$

with initial conditions (θ_0, θ_1) , where c is the unknown conductivity coefficient we aim to retrieve. This is by means of the Bukhgeim-Klibanov method imposing that the solution u to (1.1) be sufficiently smooth and appropriately bounded.

Throughout the entire text we shall suppose that c fulfills the ellipticity condition

$$c \geq c_m \text{ in } \Omega, \quad (1.2)$$

for some positive constant c_m . Notice that we may assume, and this will be systematically the case in the sequel, without restricting the generality of the foregoing, that $c_m \in (0, 1)$.

Let us now say a few words on the solution to (1.1). In order to exhibit sufficient conditions on the initial conditions (θ_0, θ_1) (together with the cross section ω and the conductivity c) ensuring that the solution to (1.1) is within an appropriate functional class we shall make precise further, we need to introduce the self-adjoint operator $A = A_c$, associated with c , generated in $L^2(\Omega)$ by the closed sesquilinear form

$$q_A[u] := \|c^{1/2}u\|_{0,\Omega}^2 = \int_{\Omega} c(x)|u(x)|^2 dx, \quad u \in D(q_A) := H_0^1(\Omega).$$

Evidently, A acts on its domain as $-\nabla \cdot c \nabla$. Since A is positive in $L^2(\Omega)$, by (1.2), the operator $A^{1/2}$ is well defined from the spectral theorem, and $D(A^{1/2}) = D(q_A) = H_0^1(\Omega)$. For the sake of definiteness, we set $A^0 := I$ and $D(A^0) := L^2(\Omega)$, where I denotes the identity operator in $L^2(\Omega)$, and we put

$$A^{m/2}v := A^{(m-1)/2}(A^{1/2}v), \quad v \in D(A^{m/2}) := \{v \in D(A^{(m-1)/2}), A^{1/2}v \in D(A^{(m-1)/2})\},$$

for each $m \in \mathbb{N}^*$. It turns out that the linear space $D(A^{m/2})$ endowed with the scalar product

$$\langle v, w \rangle_{D(A^{m/2})} := \sum_{j=0}^m \langle A^{j/2}v, A^{j/2}w \rangle_{0,\Omega},$$

is Hilbertian, and it is established in Proposition 2.6 that

$$D(A^m) = \{v \in H^{2m}(\Omega); v, Av, \dots, A^{p-1}v \in H_0^1(\Omega)\}, \quad m \in \{p-1/2, p\}, \quad p \in \mathbb{N}^*, \quad (1.3)$$

provided $\partial\omega$ is C^{2m} and $c \in W^{2m-1,\infty}(\Omega)$. As a matter of fact we know from Corollary 2.7 for any natural number m , that the system (1.1) admits a unique solution

$$u \in \bigcap_{k=0}^{m+1} C^k([0, T]; H^{m+1-k}(\Omega)), \quad (1.4)$$

provided the boundary $\partial\omega$ is C^{m+1} , the conductivity $c \in W^{m,\infty}(\Omega; \mathbb{R})$ fulfills (1.2) and $(\theta_0, \theta_1) \in D(A^{(m+1)/2}) \times D(A^{m/2})$. Moreover, if $\|c\|_{W^{m,\infty}(\Omega)} \leq c_M$ for some constant $c_M > 0$, then the solution u to (1.1) satisfies the estimate

$$\sum_{k=0}^{m+1} \|u\|_{C^k([0,T]; H^{m+1-k}(\Omega))} \leq C (\|\theta_0\|_{m+1,\Omega} + \|\theta_1\|_{m,\Omega}), \quad (1.5)$$

where $C > 0$ depends only on T , ω and c_M .

1.2.3 Admissible conductivity coefficients and initial data

In order to solve the inverse problem associated with (1.1) we seek solutions belonging to $\cap_{k=3}^4 C^k([0, T]; H^{5-k}(\Omega))$. Hence we chose $m = 4$ in (1.4) and impose on c to be in $W^{4,\infty}(\Omega; \mathbb{R})$ and satisfy (1.2). In what follows we note c_M a positive constant fulfilling

$$\|c\|_{W^{4,\infty}(\Omega)} \leq c_M. \quad (1.6)$$

Since our strategy is based on a Carleman estimate for the hyperbolic system (1.1), it is also required that the condition

$$a' \cdot \nabla_{x'} c \geq \mathbf{a}_0 \text{ in } \Omega, \quad (1.7)$$

hold for some $a' = (a_1, \dots, a_{n-1}) \in \mathbb{S}^{n-1}$ and $\mathbf{a}_0 > 0$. Hence, given \mathcal{O}_* , a neighborhood of Γ in \mathbb{R}^{n-1} , and $c_* \in W^{4,\infty}(\mathcal{O}_* \cap \Omega; \mathbb{R})$ satisfying

$$c_* \geq c_m \text{ and } a' \cdot \nabla_{x'} c_* \geq \mathbf{a}_0 \text{ in } \mathcal{O}_* \cap \Omega, \quad (1.8)$$

we introduce the set $\Lambda_{\mathcal{O}_*} = \Lambda_{\mathcal{O}_*}(a', \mathbf{a}_0, c_*, c_m, c_M)$ of admissible conductivity coefficients as

$$\Lambda_{\mathcal{O}_*} := \{c \in W^{4,\infty}(\Omega; \mathbb{R}) \text{ obeying (1.2) and (1.6) – (1.7); } c = c_* \text{ in } \mathcal{O}_* \cap \Omega\}. \quad (1.9)$$

Notice that the above choice of $m = 4$ dictates that (θ_0, θ_1) be taken in $D(A^{5/2}) \times D(A^2)$, which is embedded in $H^5(\Omega) \times H^4(\Omega)$ according to (1.3). Furthermore, it is required by the analysis of the inverse problem carried out in this article that θ_0 be in $W^{3,\infty}(\Omega)$ and satisfy

$$-a' \cdot \nabla_{x'} \theta_0 \geq \eta_0 \text{ in } \omega_*, \quad (1.10)$$

for some $\eta_0 > 0$ and some open subset ω_* in \mathbb{R}^{n-1} , with C^2 boundary, satisfying

$$\overline{\omega \setminus \mathcal{O}_*} \subset \omega_* \subset \omega. \quad (1.11)$$

Thus, for $M_0 > 0$ such that

$$\|\theta_0\|_{W^{3,\infty}(\Omega)} + \sum_{j=0}^1 \|\theta_j\|_{5-j,\Omega} \leq M_0, \quad (1.12)$$

we define the set $\Theta_{\omega_*} = \Theta_{\omega_*}(a', M_0, \eta_0)$ of admissible initial conditions (θ_0, θ_1) as

$$\Theta_{\omega_*} := \left\{ (\theta_0, \theta_1) \in \left(D(A^{5/2}) \cap W^{3,\infty}(\Omega) \right) \times D(A^2), \text{ fulfilling (1.10) and (1.12)} \right\}. \quad (1.13)$$

Having introduced all these notations we may now state the main result of this paper.

1.2.4 Main result

The following result claims Hölder stability in the inverse problem of determining c in Ω_ℓ , where $\ell > 0$ is arbitrary, from the knowledge of one boundary measurement of the solution to (1.1), performed on Σ_L for $L > \ell$ sufficiently large. The corresponding observation is viewed as a vector of the Hilbert space

$$\mathcal{H}(\Sigma_L) := H^4(0, T; L^2(\Gamma_L)) \cap H^3(0, T; H^1(\Gamma_L)),$$

endowed with the norm,

$$\|u\|_{\mathcal{H}(\Sigma_L)}^2 := \|u\|_{H^4(0,T;L^2(\Gamma_L))}^2 + \|u\|_{H^3(0,T;H^1(\Gamma_L))}^2, \quad u \in \mathcal{H}(\Sigma_L).$$

Theorem 1.1. *Assume that $\partial\omega$ is C^5 and let \mathcal{O}_* be a neighborhood of Γ in \mathbb{R}^{n-1} . For $a' = (a_1, \dots, a_{n-1}) \in \mathbb{S}^{n-1}$, $\mathfrak{a}_0 > 0$, $c_m \in (0, 1)$, $c_M > c_m$ and $c_* \in W^{4,\infty}(\mathcal{O}_* \cap \Omega; \mathbb{R})$ fulfilling (1.8), pick c_j , $j = 1, 2$, in $\Lambda_{\mathcal{O}_*}(a', \mathfrak{a}_0, c_*, c_m, c_M)$, defined by (1.9). Further, given $M_0 > 0$, $\eta_0 > 0$ and an open subset $\omega_* \subset \mathbb{R}^{n-1}$ obeying (1.11), chose $\theta_0 \in \Theta_{\omega_*}(a', M_0, \eta_0)$, defined in (1.13), and $\theta_1 = 0$.*

Then for any $\ell > 0$ we may find $L > \ell$ and $T > 0$, such that the $\bigcap_{k=0}^5 C^k([0, T], H^{5-k}(\Omega))$ -solution u_j , $j = 1, 2$, to (1.1) associated with (θ_0, θ_1) , where c_j is substituted for c , satisfies

$$\|c_1 - c_2\|_{H^1(\Omega_\ell)} \leq C \|u_1 - u_2\|_{\mathcal{H}(\Sigma_L)}^\kappa.$$

Here $C > 0$ and $\kappa \in (0, 1)$ are two constants depending only on ω , ℓ , M_0 , η_0 , a' , \mathfrak{a}_0 , c_* , c_m and c_M .

We stress out that the measurement of the observation data is performed on Γ_L and not on the whole boundary $\partial\Omega_L$.

2 Analysis of the direct problem

In this section we establish existence and uniqueness results as well as regularity properties, for the solution to hyperbolic (1.1)-like IBVP systems. The corresponding results are similar to the ones obtained for hyperbolic equations in bounded domains (see e. g. [15, Sect. 7.2, Theorem 6]) but since Ω is unbounded here, they cannot be derived from them.

2.1 Existence and uniqueness result

With reference to (1.1) we consider the boundary value problem

$$\begin{cases} \partial_t^2 v + Av = f & \text{in } Q \\ v(\cdot, 0) = g, \partial_t v(\cdot, 0) = h & \text{in } \Omega, \end{cases} \quad (2.1)$$

where f , g and h are suitable data, and we recall from ⁴[27, Sect. 3, Theorem 8.2] the following existence and uniqueness result.

Proposition 2.1. *Assume that $f \in C^0([0, T]; D(A^0))$, $g \in D(A^{1/2})$ and $h \in D(A^0)$. Then there exists a unique solution v to (2.1) such that*

$$\partial_t^k v \in C^0([0, T]; D(A^{(1-k)/2})), \quad k = 0, 1. \quad (2.2)$$

Moreover we have the estimate

$$\sum_{k=0}^1 \sup_{t \in [0, T]} \|\partial_t^k v(\cdot, t)\|_{D(A^{(1-k)/2})} \leq C \left(\|f\|_{C^0([0, T]; D(A^0))} + \|g\|_{D(A^{1/2})} + \|h\|_{D(A^0)} \right). \quad (2.3)$$

⁴Upon taking $H := D(A^0) = L^2(\Omega)$, $V := D(A^{1/2}) = H_0^1(\Omega)$ and $a(t; u, v) := \langle A^{1/2}u, A^{1/2}v \rangle_{0, \Omega} = \int_{\Omega} c(x) \nabla u(x) \cdot \overline{\nabla v(x)} dx$ for $u, v \in V$ and all $t \in [0, T]$.

2.2 Improved regularity

Proposition 2.2. *Assume that $f \in C^1([0, T]; D(A^0))$, $g \in D(A)$ and $h \in D(A^{1/2})$. Then the solution v to (2.1) satisfies*

$$\partial_t^k v \in C^0([0, T]; D(A^{(2-k)/2})), \quad k = 0, 1, 2, \quad (2.4)$$

and we have the estimate

$$\sum_{k=0}^2 \sup_{t \in [0, T]} \|\partial_t^k v(\cdot, t)\|_{D(A^{(2-k)/2})} \leq C \left(\|f\|_{C^1([0, T]; D(A^0))} + \|g\|_{D(A)} + \|h\|_{D(A^{1/2})} \right). \quad (2.5)$$

Proof. By differentiating (2.1) with respect to t , we check that $w := \partial_t v$ obeys

$$\begin{cases} \partial_t^2 w + Aw = \partial_t f & \text{in } Q \\ w(\cdot, 0) = h, \quad \partial_t w(\cdot, 0) = f(\cdot, 0) - Ag & \text{in } \Omega. \end{cases} \quad (2.6)$$

Since $f(\cdot, 0) - Ag$ is lying in $D(A^0)$ then we have $\partial_t^{k+1} v = \partial_t^k w \in C^0([0, T]; D(A^{(1-k)/2}))$ for $k = 0, 1$, with the estimate

$$\begin{aligned} \sum_{k=0}^1 \sup_{t \in [0, T]} \|\partial_t^{k+1} v(\cdot, t)\|_{D(A^{(1-k)/2})} &\leq C \left(\|\partial_t f\|_{C^0([0, T]; D(A^0))} + \|h\|_{D(A^{1/2})} + \|f(\cdot, 0) - Ag\|_{D(A^0)} \right) \\ &\leq C \left(\|f\|_{C^1([0, T]; D(A^0))} + \|g\|_{D(A)} + \|h\|_{D(A^{1/2})} \right) \end{aligned} \quad (2.7)$$

by Proposition 2.1. Further, as $Av = f - \partial_t^2 v$ from the first line in (2.1), we get that $v \in C^0([0, T]; D(A))$, and that $\|v(\cdot, t)\|_{D(A)}$ is majorized by the right hand side of (2.7), uniformly in $t \in [0, T]$. This and (2.7) yield the desired result. \square

2.3 Higher regularity

Theorem 2.3. *Let m be a nonnegative integer. We assume that $g \in D(A^{(m+1)/2})$, $h \in D(A^{m/2})$, and*

$$\partial_t^k f \in C^0([0, T]; D(A^{(m-k)/2})), \quad k = 0, \dots, m.$$

Then there exists a unique solution v to (2.1), such that

$$\partial_t^k v \in C^0([0, T]; D(A^{(m+1-k)/2})), \quad k = 0, 1, \dots, m+1. \quad (2.8)$$

Moreover we have the estimate

$$\begin{aligned} &\sum_{k=0}^{m+1} \sup_{t \in [0, T]} \|\partial_t^k v(\cdot, t)\|_{D(A^{(m+1-k)/2})} \\ &\leq C \left(\sum_{k=0}^m \|\partial_t^k f\|_{C^0([0, T]; D(A^{(m-k)/2})} + \|g\|_{D(A^{(m+1)/2})} + \|h\|_{D(A^{m/2})} \right). \end{aligned} \quad (2.9)$$

Proof. a) The proof is by an induction on m , the case $m = 0$ following from Proposition 2.1.
b) We assume that the theorem is valid for some $m \in \mathbb{N}$ and suppose that

$$\begin{cases} g \in D(A^{(m+2)/2}), h \in D(A^{(m+1)/2}), \\ \partial_t^k f \in C^0([0, T]; D(A^{(m+1-k)/2})), k = 0, \dots, m+1. \end{cases} \quad (2.10)$$

We use the same strategy as in the proof of Proposition 2.2. That is we differentiate (2.1) with respect to t and get that $w := \partial_t v$ is solution to (2.6). Next, using that $h \in D(A^{(m+1)/2})$, $f(0) - Ag \in D(A^{m/2})$ and

$$\partial_t^k(\partial_t f) = \partial_t^{k+1} f \in C^0([0, T]; D(A^{(m-k)/2})), k = 0, \dots, m,$$

from (2.10), we get that $\partial_t^{k+1} v = \partial_t^k w \in C^0([0, T]; D(A^{(m+1-k)/2}))$ for $k = 0, 1, \dots, m+1$, and the estimate:

$$\begin{aligned} & \sum_{k=0}^{m+1} \sup_{t \in [0, T]} \|\partial_t^{k+1} v(\cdot, t)\|_{D(A^{(m+1-k)/2})} \\ & \leq C \left(\sum_{k=0}^m \|\partial_t^{k+1} f\|_{C^0([0, T]; D(A^{(m-k)/2})} + \|h\|_{D(A^{(m+1)/2})} + \|f(0) - Ag\|_{D(A^{m/2})} \right). \end{aligned}$$

This entails $\partial_t^k v \in C^0([0, T]; D(A^{(m+2-k)/2}))$ for $k = 1, \dots, m+2$, and

$$\begin{aligned} & \sum_{k=1}^{m+2} \sup_{t \in [0, T]} \|\partial_t^k v(\cdot, t)\|_{D(A^{(m+2-k)/2})} \\ & \leq C \left(\sum_{k=0}^{m+1} \|\partial_t^k f\|_{C^0([0, T]; D(A^{(m+1-k)/2})} + \|h\|_{D(A^{(m+1)/2})} + \|g\|_{D(A^{(m+2)/2})} \right). \end{aligned} \quad (2.11)$$

Further, as $Av = f - \partial_t^2 v$ from the first line in (2.1), we find out that

$$\begin{aligned} \|v(\cdot, t)\|_{D(A^{(m+2)/2})} & \leq C \left(\|Av(\cdot, t)\|_{D(A^{m/2})} + \|v\|_{D(A^0)} \right) \\ & \leq C \left(\|f(\cdot, t)\|_{D(A^{m/2})} + \|\partial_t^2 v(\cdot, t)\|_{D(A^{m/2})} + \|v\|_{D(A^0)} \right). \end{aligned} \quad (2.12)$$

Here we used the identity $\|v(\cdot, t)\|_{D(A^{(m+2)/2})}^2 = \|Av(\cdot, t)\|_{D(A^{m/2})}^2 + \sum_{k=0}^1 \|A^{k/2} v(\cdot, t)\|_{D(A^0)}^2$ and the estimate $\|A^{1/2} v(\cdot, t)\|_{D(A^0)} \leq \sum_{k=0}^1 \|A^k v(\cdot, t)\|_{D(A^0)}$. Since $\|v\|_{D(A^0)}$ and $\|\partial_t^2 v(\cdot, t)\|_{D(A^{m/2})}$ are majorized by the right hand side of (2.11), uniformly in $t \in [0, T]$, (2.11)-(2.12) yield the assertion of the theorem for $m+1$. \square

Remark 2.4. The result of Theorem 2.3 is similar to the one of [15, Sect. 7.2, Theorem 6], which holds for a bounded domain. This can be seen from the characterization of the $D(A^{k/2})$ for $k = 0, \dots, m$, displayed in Subsection 2.4. Namely, it is worth noticing that the m^{th} -order compatibility conditions [15, Sect. 7.2, Eq. (62)] imposed on f , g and h , are actually hidden in the operatorial formulation of Theorem 2.3.

2.4 Characterizing the domain of $A^{m/2}$ for $m \in \mathbb{N}^*$

2.4.1 Elliptic boundary regularity

In this subsection we extend the classical elliptic boundary regularity result for the operator $\nabla \cdot c \nabla$, which is well known in any sufficiently smooth bounded subdomain of \mathbb{R}^n (see e.g. [15, Sect. 6.3, Theorem 5]), to the case of the infinite waveguide Ω under study. The proof of this result boils down to [24, Lemma 2.4] which claims elliptic boundary regularity for the Dirichlet Laplacian in Ω .

Lemma 2.5. *Let r be a nonnegative integer. We assume that $\partial\omega$ is C^{r+2} and that $c \in W^{r+1,\infty}(\Omega)$ obeys (1.2). Then, for any $\varphi \in H^r(\Omega)$, there exists a unique solution $v \in H^{r+2}(\Omega)$ to the boundary problem*

$$\begin{cases} -\nabla \cdot c \nabla v = \varphi & \text{in } \Omega \\ v = 0 & \text{on } \partial\Omega. \end{cases} \quad (2.13)$$

Moreover we have the estimate

$$\|v\|_{r+2,\Omega} \leq C_r \|\varphi\|_{r,\Omega}, \quad (2.14)$$

where C_r is a positive constant depending only on r , ω , the constant c_m appearing in (1.2) and $\|c\|_{W^{r+1,\infty}(\Omega)}$.

Proof. The proof is by induction on r .

a) We first consider the case $r = 0$. Due to (1.2) there is a unique solution $v \in H_0^1(\Omega)$ to (2.13) by the Lax-Milgram theorem. Moreover v satisfies the energy estimate

$$\|v\|_{1,\Omega} \leq C \|\varphi\|_{0,\Omega}, \quad (2.15)$$

where the constant $C > 0$ depends only on ω and c_m . Here we used (1.2) and the Poincaré inequality, which holds true in Ω since ω is bounded. Furthermore, v is solution to the boundary value problem

$$\begin{cases} -\Delta v = f & \text{in } \Omega \\ v = 0 & \text{on } \partial\Omega, \end{cases} \quad (2.16)$$

where

$$f := c^{-1} (\varphi + \nabla c \cdot \nabla v). \quad (2.17)$$

Since $f \in L^2(\Omega)$ then $v \in H^2(\Omega)$ by [24, Lemma 2.4], and $\|v\|_{2,\Omega}$ is upper bounded, up to some multiplicative constant depending only on ω , by $\|f\|_{0,\Omega}$. As a consequence we have

$$\|v\|_{2,\Omega} \leq C' (\|\varphi\|_{0,\Omega} + \|v\|_{1,\Omega}),$$

from (2.17), the constant $C' > 0$ depending only on ω , c_m and $\|c\|_{W^{1,\infty}(\Omega)}$. This and (2.15) yield (2.14) with $r = 0$.

b) Suppose that the statement of the lemma is true for $r \in \mathbb{N}$ fixed, and assume that $\partial\omega$ is C^{r+3} , $c \in W^{r+2,\infty}(\Omega)$ and $\varphi \in H^{r+1}(\Omega)$. Hence the solution v to (2.16) belongs to $H^{r+2}(\Omega)$ and satisfies the estimate (2.14), by induction assumption, and we have $f \in H^{r+1}(\Omega)$ in virtue of (2.17). Further, v being solution to (2.16) where the boundary $\partial\omega$ is C^{r+3} then $v \in H^{r+3}(\Omega)$

by [24, Lemma 2.4]. Moreover $\|v\|_{r+3,\Omega}$ is upper bounded (up to some multiplicative constant depending only on r and ω) by $\|f\|_{r+1,\Omega}$. From this and (2.17) then follows that

$$\|v\|_{r+3,\Omega} \leq C'' (\|\varphi\|_{r+1,\Omega} + \|v\|_{r+2,\Omega}),$$

where the constant $C'' > 0$ depends only on r , ω , c_m and $\|c\|_{W^{r+2,\infty}(\Omega)}$. Putting this together with (2.14), we obtain (2.14) where r is replaced by $r + 1$, proving that the statement of the lemma remains valid upon substituting $r + 1$ for r . \square

2.4.2 The domain of $A^{m/2}$ for $m \in \mathbb{N}^*$

In this subsection we characterize the domain of $A^{m/2}$ for $m \in \mathbb{N}^*$.

Proposition 2.6. *Let $m \in \mathbb{N}^*$ and let k be either $m - 1/2$ or m . Assume that $\partial\omega$ is C^{2k} and that $c \in W^{2k-1,\infty}(\Omega)$ fullfills (1.2). Then we have*

$$D(A^k) = \{u \in H^{2k}(\Omega), u, Au, \dots, A^{m-1}u \in H_0^1(\Omega)\}.$$

Moreover, the norm associated with $D(A^k)$ is equivalent to the usual one in $H^{2k}(\Omega)$: we may find a constant $c(k) > 1$, depending only on k , ω , the constant c_m defined in (1.2) and $\|c\|_{W^{2k-1,\infty}(\Omega)}$, such that we have

$$c(k)^{-1}\|u\|_{D(A^k)} \leq \|u\|_{2k,\Omega} \leq c(k)\|u\|_{D(A^k)}, \quad u \in D(A^k).$$

Proof. It suffices to show that

$$D(A^k) \subset \{u \in H^{2k}(\Omega), u, Au, \dots, A^{m-1}u \in H_0^1(\Omega)\}, \quad (2.18)$$

and

$$\|u\|_{2k,\Omega} \leq c(k)\|u\|_{D(A^k)}, \quad u \in D(A^k). \quad (2.19)$$

The proof is by induction on m .

a) We start with $m = 1$ and notice from the very definition of $A^{1/2}$ that $D(A^{1/2}) = D(q_A) = H_0^1(\Omega)$. Moreover we have

$$\|A^{1/2}u\|_{0,\Omega}^2 = q_A[u] \geq c_m \|\nabla u\|_{0,\Omega}^2, \quad u \in D(A^{1/2}),$$

in virtue of (1.2). Bearing in mind that $c_m \in (0, 1)$, we obtain that $\|u\|_{1,\Omega} \leq c_m^{-1/2}\|u\|_{D(A^{1/2})}$ for every $u \in D(A^{1/2})$. This establishes (2.19) for $k = 1/2$.

Similarly, bearing in mind that $D(A) = \{u \in H_0^1(\Omega), Au \in L^2(\Omega)\}$, we apply Lemma 2.5 with $r = 0$ and $\varphi = Au$, where $u \in D(A)$ is arbitrary. We find that $u \in H^2(\Omega)$ satisfies $\|u\|_{2,\Omega} \leq C_0\|Au\|_{0,\Omega}$, which entails (2.18)-(2.19) for $k = 1$.

b) Let us now suppose that the statement of the lemma is true for some $m \in \mathbb{N}^*$ fixed. Pick $k \in \{m - 1/2, m\}$ and assume that $\partial\omega$ is $C^{2(k+1)}$ and that $c \in W^{2k+1,\infty}(\Omega)$ satisfies (1.2). As $D(A^{k+1}) = \{u \in D(A^k), Au \in D(A^k)\}$, we deduce from the induction assumption that we have

$$D(A^{k+1}) = \{u \in H^{2k}(\Omega), Au \in H^{2k}(\Omega) \text{ and } u, \dots, A^m u \in H_0^1(\Omega)\},$$

with $\|A^j u\|_{2k, \Omega} \leq c(k) \|A^j u\|_{D(A^k)}$ for $j = 0, 1$. Thus, applying Lemma 2.5, with $r = 2k$ and $\varphi = Au$, for $u \in D(A^{k+1})$, we get that $u \in H^{2(k+1)}(\Omega)$, proving (2.18) where $(k+1, m+1)$ is substituted for (k, m) . Moreover, it holds true that

$$\|u\|_{2(k+1), \Omega} \leq C_{2k} \|Au\|_{2k, \Omega},$$

and since $Au \in D(A^k)$, the induction assumption yields $\|Au\|_{2k, \Omega} \leq c(k) \|Au\|_{D(A^k)}$. Therefore $\|u\|_{2(k+1), \Omega}$ is majorized, up to a multiplicative constant depending only on ω , c and m , by $\|u\|_{D(A^{k+1})}$, which is (2.19) where $k+1$ is substituted for k . \square

In view of Theorem 2.3 and Proposition 2.6 we obtain the following result.

Corollary 2.7. *Let m be a natural number. Assume that $\partial\omega$ is C^{m+1} , that $c \in W^{m, \infty}(\Omega)$ fulfills (1.2) and that $(\theta_0, \theta_1) \in D(A^{(m+1)/2}) \times D(A^{m/2})$. Then the initial boundary value problem (1.1) admits a unique solution*

$$u \in \bigcap_{k=0}^{m+1} C^k([0, T]; H^{m+1-k}(\Omega)).$$

Moreover, we have

$$\sum_{k=0}^{m+1} \|u\|_{C^k([0, T]; H^{m+1-k}(\Omega))} \leq C (\|\theta_0\|_{m+1, \Omega} + \|\theta_1\|_{m, \Omega}), \quad (2.20)$$

for some constant $C > 0$ depending only on T , ω and $\|c\|_{W^{m, \infty}(\Omega)}$.

3 Global Carleman estimate for hyperbolic equations in cylindrical domains

In this section we establish a global Carleman estimate for the system (1.1). To this purpose we start by time-symmetrizing the solution u of (1.1). Namely, we put

$$u(x, t) := u(x, -t), \quad x \in \Omega, \quad t \in (-T, 0). \quad (3.1)$$

Under the conditions of Theorem 1.1, and since $\theta_1 = 0$, it is not hard to check that

$$u \in \bigcap_{k=3}^4 C^k([-T, T]; H^{5-k}(\Omega)).$$

With a slight abuse of notations we put $Q := \Omega \times (-T, T)$, $\Sigma := \Gamma \times (-T, T)$ and $Q_L := \Omega_L \times (-T, T)$, $\Sigma_L := \partial\omega \times (-L, L) \times (-T, T)$ for any $L > 0$, in the remaining part of this text.

3.1 The case of second order hyperbolic operators

In view of establishing a Carleman estimate for the operator

$$A := A(x, t, \partial) = \partial_t^2 - \nabla \cdot c \nabla + R, \quad (3.2)$$

where R is a first-order partial differential operator with $L^\infty(Q)$ coefficients, we define for every $\delta > 0$ and $\gamma > 0$ the following weight functions:

$$\psi(x, t) = \psi_\delta(x, t) := |x' - \delta a'|^2 - x_n^2 - t^2 \text{ and } \varphi(x, t) = \varphi_{\delta, \gamma}(x, t) := e^{\gamma \psi(x, t)}, \quad (x, t) \in Q. \quad (3.3)$$

Further for $L > 0$ and $T > 0$ we introduce the space

$$\mathcal{X}_{L, T} := \left\{ u \in H^2(Q_L); \partial_{x_n}^j u(\cdot, \pm L, \cdot) = 0 \text{ in } \omega \times (-T, T), \partial_t^j u(\cdot, \pm T) = 0 \text{ in } \Omega_L, j = 0, 1 \right\}, \quad (3.4)$$

and state the:

Proposition 3.1. *Let A be defined by (3.2), where c obeys (1.2)–(1.7), and let ℓ be positive. Then there exist $\delta_0 > 0$ and $\gamma_0 > 0$ such that for all $\delta \geq \delta_0$ and $\gamma \geq \gamma_0$, we may find $L > \ell$, $T > 0$ and $s_0 > 0$ for which the estimate*

$$s \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j v\|_{0, Q_L}^2 \leq C \left(\|e^{s\varphi} A v\|_{0, Q_L}^2 + s \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j v\|_{0, \partial Q_L}^2 \right), \quad (3.5)$$

holds for any $s \geq s_0$ and $v \in \mathcal{X}_{L, T}$. Here C is a positive constant depending only on ω , a' , \mathbf{a}_0 , δ_0 , γ_0 , s_0 , c_m and c_M .

Moreover there exists a constant $d_\ell > 0$, depending only on ω , ℓ , δ_0 and γ_0 , such that the weight function φ defined by (3.3) satisfies

$$\varphi(x', x_n, 0) \geq d_\ell, \quad (x', x_n) \in \bar{\omega} \times [-\ell, \ell], \quad (3.6)$$

and we may find $\epsilon \in (0, (L - \ell)/2)$ and $\nu > 0$ so small that we have:

$$\max_{x \in \bar{\omega} \times [-L, L]} \varphi(x', x_n, t) \leq \tilde{d}_\ell := d_\ell e^{-\gamma \nu^2}, \quad |t| \in [T - 2\epsilon, T], \quad (3.7)$$

$$\max_{(x', t) \in \bar{\omega} \times [-T, T]} \varphi(x', x_n, t) \leq \tilde{d}_\ell, \quad |x_n| \in [L - 2\epsilon, L]. \quad (3.8)$$

Proof. The proof is divided in three parts and essentially boils down to [20, theorem 3.2.1’].

First part: Definition of δ_0 , L and T . Bearing in mind that $|x' - \delta a'|^2 - |y' - \delta a'|^2 = |x'| - |y'|^2 - 2\delta a' \cdot (x' - y')$ for all $x', y' \in \omega$, we see that $\sup_{x' \in \omega} |x' - \delta a'|^2 - \inf_{x' \in \omega} |x' - \delta a'|^2 \leq |\omega|(|\omega| + 4\delta|a'|)$ for every $\delta > 0$, where $|\omega| := \sup_{x' \in \bar{\omega}} |x'|$. Hence the function

$$g_\ell(\delta) = \left(\sup_{x' \in \omega} |x' - \delta a'|^2 - \inf_{x' \in \omega} |x' - \delta a'|^2 + \ell^2 \right)^{1/2} \quad (3.9)$$

scales at most like $\delta^{1/2}$, proving that there exists $\delta_0 > 0$ so large that

$$\delta \mathbf{a}_0 > \left(\left(1 + \frac{2}{c_m^{1/2}} \right) g_\ell(\delta) + |\omega| + 2 \right) c_M + 2, \quad \delta \geq \delta_0. \quad (3.10)$$

Further, since ω is bounded and $a' \neq 0_{\mathbb{R}^{n-1}}$ by (1.7), we may as well assume upon possibly enlarging δ_0 , that we have in addition $c_m^{1/2} \inf_{x' \in \overline{\omega}} |x' - \delta a'| > g_\ell(\delta)$ for all $\delta \geq \delta_0$. This and (3.10) yield that there exists $\vartheta > 0$ so small that the two following inequalities

$$\delta \mathbf{a}_0 - \left(L + |\omega| + 2 \left(1 + \frac{T}{c_m^{1/2}} \right) \right) c_M - 2 > 0, \quad (3.11)$$

and

$$c_m^{1/2} \inf_{x' \in \overline{\omega}} |x' - \delta a'| > T, \quad (3.12)$$

hold simultaneously for every L and T in $(g_\ell(\delta), g_\ell(\delta) + \vartheta)$, uniformly in $\delta \geq \delta_0$.

Second part: Proof of (3.5). We first introduce the following notations, we shall use in the remaining part of the proof. For notational simplicity we put $\mathbf{x} := (x, t)$ for $(x, t) \in Q_L$ and $\nabla_{\mathbf{x}} = (\partial_1, \dots, \partial_n, \partial_{n+1})^T$. We also write $\xi' = (\xi_1, \dots, \xi_{n-1})^T \in \mathbb{R}^{n-1}$, $\xi = (\xi_1, \dots, \xi_n)^T \in \mathbb{R}^n$ and $\tilde{\xi} = (\xi_1, \dots, \xi_n, \xi_{n+1})^T \in \mathbb{R}^{n+1}$. We call A_2 the principal part of the operator A , that is $A_2 = A_2(\mathbf{x}, \partial) = \partial_t^2 - c(x)\Delta$, and denote its symbol by $A_2(x, \tilde{\xi}) = c(x)|\xi|^2 - \xi_{n+1}^2$, where $|\xi| = \left(\sum_{j=1}^n \xi_j^2 \right)^{1/2}$. Since $A_2(x, \nabla_{\mathbf{x}}\psi(\mathbf{x})) = 4(c(x)(|x' - \delta a'|^2 + x_n^2) - x_{n+1}^2)$ for every $\mathbf{x} \in \overline{Q}_L$, we have

$$A_2(x, \nabla_{\mathbf{x}}\psi(\mathbf{x})) > 0, \quad \mathbf{x} \in \overline{Q}_L, \quad (3.13)$$

by (1.7) and (3.12). For all $\mathbf{x} \in \overline{Q}_L$ and $\tilde{\xi} \in \mathbb{R}^{n+1}$, put

$$J(\mathbf{x}, \tilde{\xi}) = J = \sum_{j,k=1}^{n+1} \frac{\partial A_2}{\partial \xi_j} \frac{\partial A_2}{\partial \xi_k} \partial_j \partial_k \psi + \sum_{j,k=1}^{n+1} \left(\left(\partial_k \frac{\partial A_2}{\partial \xi_j} \right) \frac{\partial A_2}{\partial \xi_k} - (\partial_k A_2) \frac{\partial^2 A_2}{\partial \xi_j \partial \xi_k} \right) \partial_j \psi, \quad (3.14)$$

where, for the sake of shortness, we write ∂_j , $j \in \mathbb{N}_{n+1}^*$, instead of $\partial/\partial x_j$, and x_{n+1} stands for t . Assuming that

$$A_2(x, \tilde{\xi}) = c(x)|\xi|^2 - \xi_{n+1}^2 = 0, \quad x \in \overline{\Omega}, \quad \tilde{\xi} \in \mathbb{R}^{n+1} \setminus \{0\}, \quad (3.15)$$

and

$$\nabla_{\tilde{\xi}} A_2(x, \tilde{\xi}) \cdot \nabla_{\mathbf{x}}\psi(\mathbf{x}) = 4 [c(x)(\xi' \cdot (x' - \delta a') - \xi_n x_n) + \xi_{n+1} x_{n+1}] = 0, \quad \mathbf{x} \in \overline{Q}_L, \quad \tilde{\xi} \in \mathbb{R}^{n+1} \setminus \{0\}, \quad (3.16)$$

we shall prove that $J(\mathbf{x}, \tilde{\xi}) > 0$ for any $(\mathbf{x}, \tilde{\xi}) \in \overline{Q}_L \times \mathbb{R}^{n+1}$. To this end we notice that the first sum in the rhs of (3.14) reads $\langle \text{Hess}(\psi) \nabla_{\tilde{\xi}} A_2, \nabla_{\tilde{\xi}} A_2 \rangle = 8(c^2(|\xi'|^2 - \xi_n^2) - \xi_{n+1}^2)$, and that

$$\begin{aligned} & \sum_{j,k=1}^{n+1} \left(\left(\partial_k \frac{\partial A_2}{\partial \xi_j} \right) \frac{\partial A_2}{\partial \xi_k} - (\partial_k A_2) \frac{\partial^2 A_2}{\partial \xi_j \partial \xi_k} \right) \partial_j \psi = \sum_{j,k=1}^n \left(\left(\partial_k \frac{\partial A_2}{\partial \xi_j} \right) \frac{\partial A_2}{\partial \xi_k} - (\partial_k A_2) \frac{\partial^2 A_2}{\partial \xi_j \partial \xi_k} \right) \partial_j \psi \\ & = 2c \sum_{j,k=1}^n \left(\left(2\xi_j \xi_k - |\xi|^2 \frac{\partial \xi_j}{\partial \xi_k} \right) \partial_k c \right) \partial_j \psi = 2(2(\nabla c \cdot \xi)(\nabla \psi \cdot \xi) - (\nabla c \cdot \nabla \psi)|\xi|^2), \end{aligned}$$

since $(\partial_k \frac{\partial A_2}{\partial \xi_j}) \frac{\partial A_2}{\partial \xi_k} - (\partial_k A_2) \frac{\partial^2 A_2}{\partial \xi_j \partial \xi_k} = 0$ if either j or k is equal to $n+1$. Therefore we have

$$\begin{aligned} J &= 2 \left[4(c^2(|\xi'|^2 - \xi_n^2) - \xi_{n+1}^2) + 2c(\nabla c \cdot \xi)(\nabla \psi \cdot \xi) - c(\nabla c \cdot \nabla \psi)|\xi|^2 \right] \\ &= 4 \left[2c^2(|\xi'|^2 - \xi_n^2) - (2 + (x' - \delta a') \cdot \nabla_{x'} c - x_n \partial_n c) \xi_{n+1}^2 - 2x_{n+1} \xi_{n+1} \nabla c \cdot \xi \right], \end{aligned}$$

from (3.15)-(3.16). Further, in view of (3.15) we have $c^2(|\xi'|^2 - \xi_n^2) \geq -c^2|\xi|^2 \geq -c\xi_{n+1}^2$ and $|\nabla c \cdot \xi| \leq |\nabla c||\xi| \leq (|\nabla c|/c^{1/2})|\xi_{n+1}|$, whence

$$J \geq 4 \left[\delta a' \cdot \nabla_{x'} c - \left(x \cdot \nabla c + 2c + 2T \frac{|\nabla c|}{c^{1/2}} + 2 \right) \right] \xi_{n+1}^2. \quad (3.17)$$

Here we used that fact that $x_{n+1} = t \in [0, T]$. Due to (1.6)-(1.7), the rhs of (3.17) is lower bounded, up to the multiplicative constant $4\xi_{n+1}^2$, by the lhs of (3.11). Since ξ_{n+1} is non zero by (1.6) and (3.15), then we obtain $J(\mathbf{x}, \tilde{\xi}) > 0$ for all $(\mathbf{x}, \tilde{\xi}) \in \overline{Q}_L \times \mathbb{R}^{n+1}$. With reference to (3.13), we may apply [20, Theorem 3.2.1'], getting two constants $s_0 = s_0(\gamma) > 0$ and $C > 0$ such that (3.5) holds for any $s \geq s_0$ and $v \in H^2(Q_L)$.

Third part: Proof of (3.6)-(3.8). First, (3.6) follows readily from (3.3), with $d_\ell := e^{\gamma\beta_\ell}$ and $\beta_\ell := \inf_{x' \in \omega} |x' - \delta a'|^2 - \ell^2$. Next, for $\nu \in (0, \vartheta)$ arbitrarily fixed, we put

$$L = T = g_\ell(\delta) + \nu. \quad (3.18)$$

Notice for further reference from (3.9), (3.12) and (3.18), that we have

$$\beta_\ell \geq \frac{g_\ell(\delta)^2}{c_m} - \ell^2 \geq \left(\frac{1 - c_m}{c_m} \right) \ell^2 > 0, \quad (3.19)$$

since $c_m \in (0, 1)$, by assumption. Similarly, as $T^2 \geq g_\ell(\delta)^2 + \nu^2 = \sup_{x' \in \omega} |x' - \delta a'|^2 - (\beta_\ell^2 - \nu^2)$, we deduce from (3.3) that

$$\varphi(x', x_n, \pm T) \leq e^{\gamma(\sup_{x' \in \omega} |x' - \delta a'|^2 - x_n^2 - T^2)} \leq e^{\gamma(\beta_\ell - \nu^2)} e^{-\gamma x_n^2}, \quad (x', x_n) \in \overline{\omega} \times [-L, L].$$

With reference to (3.19) we may thus chose $\epsilon \in (0, (L - \ell)/2)$ so small that

$$\varphi(x', x_n, t) \leq d_\ell e^{-\gamma\nu^2} e^{-\gamma x_n^2}, \quad (x', x_n) \in \overline{\omega} \times [-L, L], \quad |t| \in [T - 2\epsilon, T],$$

which entails (3.7). Finally, since t and x_n play symmetric roles in (3.3), and since $T = L$, we obtain (3.8) by substituting (T, t) for (L, x_n) in (3.7). \square

3.2 A Carleman estimate for the system (1.1)

In this subsection we derive from Proposition 3.1 a global Carleman estimate for the solution to the boundary value problem

$$\begin{cases} \partial_t^2 u - \nabla \cdot c \nabla u = f & \text{in } Q \\ u = 0 & \text{on } \Sigma, \end{cases} \quad (3.20)$$

where $f \in L^2(Q)$. To this purpose we introduce a cut-off function $\chi \in C^2(\mathbb{R}; [0, 1])$, such that

$$\chi(x_n) := \begin{cases} 1 & \text{if } |x_n| < L - 2\epsilon, \\ 0 & \text{if } |x_n| \geq L - \epsilon, \end{cases} \quad (3.21)$$

where ϵ is the same as in Proposition 3.1, and we set

$$u_\chi(x, t) := \chi(x)u(x, t) \text{ and } f_\chi(x, t) := \chi(x)f(x, t), \quad (x, t) \in Q.$$

Corollary 3.2. *Let $f \in L^2(Q)$. Then, under the conditions of Proposition 3.1, there exist two constants $s_* > 0$ and $C > 0$, depending only on ω , ℓ , M_0 , η_0 , a' , \mathbf{a}_0 , c_m and c_M , such that the estimate*

$$s \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j u\|_{0,Q_L}^2 \leq C \left(\|e^{s\varphi} f\|_{0,Q_L}^2 + s^3 e^{2s\tilde{d}_\ell} \|u\|_{1,Q_L}^2 + s \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j u_\chi\|_{0,\Sigma_L}^2 \right),$$

holds for any solution u to (3.20), uniformly in $s \geq s_*$.

Proof. Since u is solution to (3.20) we have

$$\partial_t^2 u_\chi - \nabla \cdot c \nabla u_\chi = f_\chi + R_1 u \text{ in } Q,$$

where

$$R_1 = R_1(x, \partial) := [\chi, \nabla \cdot c \nabla] = -(c \Delta \chi + \nabla c \cdot \nabla \chi + 2c \nabla \chi \cdot \nabla), \quad (3.22)$$

is a first-order differential operator. Therefore, the function $v(x, t) := \eta(t)u_\chi(x, t)$, where $\eta \in C^2(\mathbb{R}; [0, 1])$ is such that

$$\eta(t) := \begin{cases} 1 & \text{if } |t| < T - 2\epsilon, \\ 0 & \text{if } |t| \geq T - \epsilon, \end{cases}$$

satisfies

$$\partial_t^2 v - \nabla \cdot c \nabla v = g := \eta f_\chi + \eta R_1 u + \eta'' u_\chi + 2\eta' \partial_t u_\chi \text{ in } Q.$$

Moreover, as $v(\cdot, \pm L, \cdot) = \partial_{x_n} v(\cdot, \pm L, \cdot) = 0$ in $\omega \times (-T, T)$ and $v(\cdot, \pm T) = \partial_t v(\cdot, \pm T) = 0$ in Ω_L , we may apply Proposition 3.1, getting

$$s \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j v\|_{0,Q_L}^2 \leq C \left(\|e^{s\varphi} g\|_{0,Q_L}^2 + s \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j v\|_{0,\partial Q_L}^2 \right). \quad (3.23)$$

Further, bearing in mind that $\partial Q_L = \Sigma_L \cup (\omega \times \{\pm L\} \times (-T, T)) \cup (\Omega_L \times \{\pm T\})$, we deduce from the vanishing of $v(\cdot, \pm L, \cdot)$ and $\nabla_{x,t} v(\cdot, \pm L, \cdot)$ in $\omega \times (-T, T)$, and the one of $v(\cdot, \pm T)$ and $\nabla_{x,t} v(\cdot, \pm T)$ in Ω_L , that

$$\|e^{s\varphi} \nabla_{x,t}^j v\|_{0,\partial Q_L} = \|e^{s\varphi} \nabla_{x,t}^j v\|_{0,\Sigma_L}, \quad j = 0, 1. \quad (3.24)$$

Next we know from (3.8) and (3.22) that

$$\|e^{s\varphi} \eta R_1 u\|_{0,Q_L} \leq C e^{s\tilde{d}_\ell} \|u\|_{L^2(-T,T; H^1(\Omega_L))}, \quad (3.25)$$

and from (3.7) that

$$\|e^{s\varphi}(\eta''u_\chi + 2\eta'\partial_t u_\chi)\|_{0,Q_L} \leq Ce^{s\tilde{d}_\ell}\|u_\chi\|_{H^1(-T,T;L^2(\Omega_L))}. \quad (3.26)$$

Hence, putting (3.23)–(3.26) together, we find out that

$$s \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j v\|_{0,Q_L}^2 \leq C \left(\|e^{s\varphi} f\|_{0,Q_L}^2 + s \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j u_\chi\|_{0,\Sigma_L}^2 + e^{2s\tilde{d}_\ell} \|u\|_{H^1(Q_L)}^2 \right). \quad (3.27)$$

The next step of the proof involves noticing from (3.7) that $\|e^{s\varphi}(1-\eta)\nabla_x^j u_\chi\|_{0,Q_L} \leq e^{s\tilde{d}_\ell}\|\nabla_x^j u_\chi\|_{0,Q_L}$ for $j = 0, 1$, hence

$$\begin{aligned} \|e^{s\varphi} \nabla_x^j u_\chi\|_{0,Q_L} &\leq \|e^{s\varphi}(1-\eta)\nabla_x^j u_\chi\|_{0,Q_L} + \|e^{s\varphi} \nabla_x^j v\|_{0,Q_L} \\ &\leq e^{s\tilde{d}_\ell} \|\nabla_x^j u_\chi\|_{0,Q_L} + \|e^{s\varphi} \nabla_x^j v\|_{0,Q_L}, \quad j = 0, 1. \end{aligned} \quad (3.28)$$

Furthermore, by combining the identity $\partial_t u_\chi = (1-\eta)\partial_t u_\chi + \eta\partial_t u_\chi = (1-\eta)\partial_t u_\chi - \eta'\partial_t u_\chi + \partial_t v$ with (3.7), we get that

$$\begin{aligned} \|e^{s\varphi} \partial_t u_\chi\|_{0,Q_L} &\leq \|e^{s\varphi}(1-\eta)\partial_t u_\chi\|_{0,Q_L} + \|e^{s\varphi} \eta' u_\chi\|_{0,Q_L} + \|e^{s\varphi} \partial_t v\|_{0,Q_L} \\ &\leq e^{s\tilde{d}_\ell} (\|\partial_t u_\chi\|_{0,Q_L} + \|\eta'\|_{L^\infty(-T,T)} \|u_\chi\|_{0,Q_L}) + \|e^{s\varphi} \partial_t v\|_{0,Q_L}, \end{aligned}$$

which, together with (3.28), yields

$$\sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j u_\chi\|_{0,Q_L}^2 \leq C \sum_{j=0,1} s^{2(1-j)} \left(e^{2s\tilde{d}_\ell} \|\nabla_{x,t}^j u_\chi\|_{0,Q_L}^2 + \|e^{s\varphi} \nabla_{x,t}^j v\|_{0,Q_L}^2 \right). \quad (3.29)$$

Similarly, using (3.8), we derive from the identity $\partial_t^j u = \partial_t^j u_\chi + (1-\chi)\partial_t^j u$ for $j = 0, 1$, that

$$\|e^{s\varphi} \partial_t^j u\|_{0,Q_L} \leq \|e^{s\varphi} \partial_t^j u_\chi\|_{0,Q_L} + e^{s\tilde{d}_\ell} \|\partial_t^j u\|_{0,Q_L}, \quad j = 0, 1,$$

and from $\nabla_x u = \nabla_x u_\chi + (1-\chi)\nabla_x u - \chi'(0, \dots, 0, u)^T$, that

$$\|e^{s\varphi} \nabla_x u\|_{0,Q_L} \leq \|e^{s\varphi} \nabla_x u_\chi\|_{0,Q_L} + e^{s\tilde{d}_\ell} (\|\nabla_x u\|_{0,Q_L} + \|\chi'\|_{L^\infty(-L,L)} \|u\|_{0,Q_L}).$$

As a consequence we have

$$\sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j u\|_{0,Q_L}^2 \leq C \sum_{j=0,1} s^{2(1-j)} \left(e^{2s\tilde{d}_\ell} \|\nabla_{x,t}^j u\|_{0,Q_L}^2 + \|e^{s\varphi} \nabla_{x,t}^j u_\chi\|_{0,Q_L}^2 \right). \quad (3.30)$$

Finally we obtain the desired result by gathering (3.27) and (3.29)–(3.30). \square

4 Inverse problem

4.1 Linearized inverse problem and preliminary estimate

In this subsection we introduce the linearized inverse problem associated with (1.1) and relate the first Sobolev norm of the conductivity to some suitable initial condition of this boundary problem.

Namely, given $c_i \in \Lambda_\Gamma$ for $i = 1, 2$, we note u_i the solution to (1.1) where c_i is substituted for c , suitably extended to $(-T, 0)$ in accordance with (3.1). Thus, putting

$$c := c_1 - c_2 \text{ and } f_c := \nabla \cdot (c \nabla u_2), \quad (4.1)$$

it is clear from (1.1) that the function $u := u_1 - u_2$ is solution to the linearized system

$$\begin{cases} \partial_t^2 u - \nabla \cdot (c_1 \nabla u) = f_c & \text{in } Q \\ u = 0 & \text{on } \Sigma \\ u(\cdot, 0) = \partial_t u(\cdot, 0) = 0 & \text{in } \Omega. \end{cases} \quad (4.2)$$

By differentiating k -times (4.2) with respect to t , for $k \in \mathbb{N}^*$ fixed, we see that $u^{(k)} := \partial_t^k u$ is solution to

$$\begin{cases} \partial_t^2 u^{(k)} - \nabla \cdot (c_1 \nabla u^{(k)}) = f_c^{(k)} & \text{in } Q \\ u^{(k)} = 0 & \text{on } \Sigma, \end{cases} \quad (4.3)$$

with $f_c^{(k)} := \partial_t^k f_c = \nabla \cdot (c \nabla u_2^{(k)})$, where $u_2^{(k)}$ stands for $\partial_t^k u_2$.

We stick with the notations of Corollary 3.2. In particular, for any function v , we denote χv by v_χ , where χ is defined in (3.21). Upon multiplying both sides of the identity (4.3) by χ , we obtain that

$$\begin{cases} \partial_t^2 u_\chi^{(k)} - \nabla \cdot (c_1 \nabla u_\chi^{(k)}) = f_{c_\chi}^{(k)} - g_k & \text{in } Q \\ u_\chi^{(k)} = 0 & \text{on } \Sigma, \end{cases} \quad (4.4)$$

with

$$f_{c_\chi} := \nabla \cdot (c_\chi \nabla u_2) \text{ and } g_k := \nabla \cdot (c_1 (\nabla \chi) u^{(k)}) + c_1 \nabla \chi \cdot \nabla u^{(k)} + c \nabla \chi \cdot \nabla u_2^{(k)}. \quad (4.5)$$

Notice that g_k is supported in $\tilde{\Omega}_\epsilon := \{x = (x', x_n), x' \in \omega \text{ and } |x_n| \in (L - 2\epsilon, L - \epsilon)\}$.

Having said that we may now upper bound, up to suitable additive and multiplicative constants, the $e^{s\varphi(\cdot, 0)}$ -weighted first Sobolev norm of the conductivity c_χ in Ω_L , by the corresponding norm of the initial condition $u_\chi^{(2)}(\cdot, 0)$.

Lemma 4.1. *Let u be the solution to the linearized problem (4.2) and let χ be defined by (3.21). Then there exist two constants $s_* > 0$ and $C > 0$, depending only on ω , ε and the constant M_0 defined by (1.12), such that the estimate*

$$\sum_{j=0,1} \|e^{s\varphi(\cdot, 0)} \nabla^j c_\chi\|_{0, \Omega_L}^2 \leq C s^{-1} \left(\sum_{j=0,1} \|e^{s\varphi(\cdot, 0)} \nabla^j u_\chi^{(2)}(\cdot, 0)\|_{0, \Omega_L}^2 + e^{2s\bar{d}_\ell} \right),$$

holds for all $s \geq s_*$.

Proof. Let Ω_* be an open subset of \mathbb{R}^n with C^2 boundary, such that

$$\omega_* \times (-(L - \epsilon), L - \epsilon) \subset \Omega_* \subset \omega_* \times (-L, L), \quad (4.6)$$

where ϵ is defined by Proposition 3.1. We notice from (3.21) and (1.11) that $\partial_i^j c_\chi \in H_0^1(\Omega_*)$ for all $i \in \mathbb{N}_n^*$ and $j = 0, 1$.

Further, with reference to (1.10) we may assume upon possibly enlarging $\delta \in [\delta_0, +\infty)$, where δ_0 is the same as in Proposition 3.1, that we have

$$|\nabla\theta_0 \cdot (x_1 - \delta a_1, \dots, x_{n-1} - \delta a_{n-1}, -x_n)| \geq \mu_0 > 0, \quad x \in \Omega_*.$$

Thus applying [18, Proposition 2.2]⁵ with $D = \Omega_*$, $P(x, \partial)v = \nabla \cdot ((\nabla\theta_0)v)$ and $v = \partial_i^j c_\chi \in H_0^1(\Omega_*)$ since $\chi(x_n) = 0$ for $x_n \geq L - \epsilon$, for $i \in \mathbb{N}_n^*$ and $j = 0, 1$, we obtain that

$$s \|e^{s\varphi(\cdot, 0)} \partial_i^j c_\chi\|_{0, \Omega_*}^2 \leq C \|e^{s\varphi(\cdot, 0)} \nabla \cdot ((\partial_i^j c_\chi) \nabla\theta_0)\|_{0, \Omega_*}^2, \quad i \in \mathbb{N}_n^*, \quad j = 0, 1. \quad (4.7)$$

Since $c_\chi(x', x_n) = 0$ a.e. in $\omega_* \times ((-L, -(L - \epsilon)) \cup (L - \epsilon, L))$ by (3.21), we have $\|e^{s\varphi(\cdot, 0)} \partial_i^j c_\chi\|_{0, \Omega_*} = \|e^{s\varphi(\cdot, 0)} \partial_i^j c_\chi\|_{0, \Omega_L}$ for each $i \in \mathbb{N}_n^*$ and $j = 0, 1$, from (4.6). We derive from this and (4.7) that

$$s \|e^{s\varphi(\cdot, 0)} \partial_i^j c_\chi\|_{0, \Omega_L}^2 \leq C \|e^{s\varphi(\cdot, 0)} \nabla \cdot ((\partial_i^j c_\chi) \nabla\theta_0)\|_{0, \Omega_L}^2, \quad i \in \mathbb{N}_n^*, \quad j = 0, 1. \quad (4.8)$$

Further, taking $t = 0$ in the first line of (4.4) with $j = 0$, we get that

$$\nabla \cdot (c_\chi \nabla\theta_0) = u_\chi^{(2)}(\cdot, 0) + c \nabla\chi \cdot \nabla\theta_0. \quad (4.9)$$

From this, (4.8) and (3.8) then follows that

$$\begin{aligned} s \|e^{s\varphi(\cdot, 0)} c_\chi\|_{0, \Omega_L}^2 &\leq C \left(\|e^{s\varphi(\cdot, 0)} u_\chi^{(2)}(\cdot, 0)\|_{0, \Omega_L}^2 + \|e^{s\varphi(\cdot, 0)} c \nabla\chi \cdot \nabla\theta_0\|_{0, \Omega_L}^2 \right) \\ &\leq C \left(\|e^{s\varphi(\cdot, 0)} u_\chi^{(2)}(\cdot, 0)\|_{0, \Omega_L}^2 + e^{2s\bar{d}_\ell} \right). \end{aligned} \quad (4.10)$$

Similarly, since $\nabla \cdot ((\partial_i c_\chi) \nabla\theta_0) = \partial_i \nabla \cdot (c_\chi \theta_0) - \nabla \cdot (c_\chi \nabla \partial_i \theta_0)$ for every $i \in \mathbb{N}_n^*$, we derive from (4.9) that

$$\nabla \cdot ((\partial_i c_\chi) \nabla\theta_0) = \partial_i u_\chi^{(2)}(\cdot, 0) + \partial_i (c \nabla\theta_0 \cdot \nabla\chi) - \nabla c_\chi \cdot \nabla \partial_i \theta_0 - c_\chi \Delta \partial_i \theta_0.$$

As a consequence we have,

$$\begin{aligned} &\|e^{s\varphi(\cdot, 0)} \nabla \cdot ((\partial_i c_\chi) \nabla\theta_0)\|_{0, \Omega_L}^2 \\ &\leq C \left(\|e^{s\varphi(\cdot, 0)} \partial_i u_\chi^{(2)}(\cdot, 0)\|_{0, \Omega_L}^2 + \sum_{j=0,1} \|e^{s\varphi(\cdot, 0)} \nabla^j c_\chi\|_{0, \Omega_L}^2 + e^{2s\bar{d}_\ell} \right), \quad i \in \mathbb{N}_n^*, \end{aligned}$$

according to (3.8). Summing up the above estimate over i in \mathbb{N}_n^* , it follows from (4.8) that

$$s \|e^{s\varphi(\cdot, 0)} \nabla c_\chi\|_{0, \Omega_L}^2 \leq C \left(\|e^{s\varphi(\cdot, 0)} \nabla u_\chi^{(2)}(\cdot, 0)\|_{0, \Omega_L}^2 + \sum_{j=0,1} \|e^{s\varphi(\cdot, 0)} \nabla^j c_\chi\|_{0, \Omega_L}^2 + e^{2s\bar{d}_\ell} \right).$$

This and (4.10) yield the desired result. \square

⁵Let D be a bounded open subset of \mathbb{R}^n , $n \geq 1$, with C^2 boundary, and consider the first-order operator $P(x, \partial) := \sum_{i=1}^n p_i(x) \partial_i + p_0(x)$, where $p_0 \in C^0(\bar{D})$ and $p := (p_1, \dots, p_n) \in C^1(\bar{D})^n$. Assume that

$$|p(x) \cdot (x_1 - \delta a_1, \dots, x_{n-1} - \delta a_{n-1}, -x_n)| \geq \mathfrak{p}_m, \quad x \in \bar{D},$$

for some $\mathfrak{p}_m > 0$. Then for any $\mathfrak{p}_M \geq \max\{\|p_0\|_{C^0(\bar{D})}, \|p_i\|_{C^1(\bar{D})}, i \in \mathbb{N}_n^*\}$, there exist $s_* > 0$ and $C > 0$, depending only on \mathfrak{p}_M , such that the estimate

$$\|e^{s\varphi(\cdot, 0)} v\|_{0, D}^2 \leq C s^{-1} \|e^{s\varphi(\cdot, 0)} P v\|_{0, D}^2,$$

holds for all $s \geq s_*$ and $v \in H_0^1(D)$.

4.2 Completion of the proof

The proof is divided into three steps.

Step 1. The first step of the proof is to upper bound $u_\chi^{(2)}(\cdot, 0)$ in the $e^{s\varphi(\cdot, 0)}$ -weighted $H^1(\Omega_L)$ -norm topology, by the corresponding norms of $u_\chi^{(2)}$ and $u_\chi^{(3)}$ in Q_L , with the aid of the following technical result, borrowed from [3][Lemma 3.2].

Lemma 4.2. *There exists a constant $s_* > 0$ depending only on T such that we have*

$$\|z(\cdot, 0)\|_{0, \Omega_L}^2 \leq 2 \left(s \|z\|_{0, Q_L}^2 + s^{-1} \|\partial_t z\|_{0, Q_L}^2 \right),$$

for all $s \geq s_*$ and $z \in H^1(-T, T; L^2(\Omega_L))$.

Namely, we apply Lemma 4.2 with $z = e^{s\varphi} \partial_i^j u_\chi^{(2)}$ for $i \in \mathbb{N}_n^*$ and $j = 0, 1$, getting

$$\|e^{s\varphi(\cdot, 0)} \partial_i^j u_\chi^{(2)}(\cdot, 0)\|_{0, \Omega_L}^2 \leq 2 \left(s \|e^{s\varphi} \partial_i^j u_\chi^{(2)}\|_{0, Q_L}^2 + s^{-1} \|e^{s\varphi} \partial_i^j u_\chi^{(3)}\|_{0, Q_L}^2 \right), \quad s \geq s_*.$$

Summing up the above estimate over $i \in \mathbb{N}_n^*$ and $j = 0, 1$, we obtain for all $s \geq s_*$ that

$$\sum_{j=0,1} \|e^{s\varphi(\cdot, 0)} \nabla^j u_\chi^{(2)}(\cdot, 0)\|_{0, \Omega_L}^2 \leq 2 \sum_{j=0,1} \left(s \|e^{s\varphi} \nabla^j u_\chi^{(2)}\|_{0, Q_L}^2 + s^{-1} \|e^{s\varphi} \nabla^j u_\chi^{(3)}\|_{0, Q_L}^2 \right). \quad (4.11)$$

Step 2. The next step involves majorizing the right hand side of (4.11) with

$$\mathfrak{h}_k(s) := \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j u_\chi^{(k)}\|_{0, \Sigma_L}^2, \quad k = 2, 3. \quad (4.12)$$

Indeed, since $u_\chi^{(k)}$, for $k = 2, 3$, is solution to (3.20) with $c = c_1$ and $f = f_{c_\chi} - g_k$, according to (4.4), then Corollary 3.2 yields

$$s \sum_{j=0,1} s^{2(1-j)} \|e^{s\varphi} \nabla_{x,t}^j u_\chi^{(k)}\|_{0, Q_L}^2 \leq C \left(\|e^{s\varphi} f_{c_\chi}^{(k)}\|_{0, Q_L}^2 + \|e^{s\varphi} g_k\|_{0, \tilde{Q}_\epsilon}^2 + s^3 e^{2s\tilde{d}_\ell} \|u_\chi^{(k)}\|_{1, Q_L}^2 + s \mathfrak{h}_k(s) \right),$$

for s large enough. In light of (4.11) this entails that

$$\begin{aligned} & \sum_{j=0,1} \|e^{s\varphi(\cdot, 0)} \nabla^j u_\chi^{(2)}(\cdot, 0)\|_{0, \Omega_L}^2 \\ & \leq C \sum_{k=2,3} \left(\|e^{s\varphi} f_{c_\chi}^{(k)}\|_{0, Q_L}^2 + \|e^{s\varphi} g_k\|_{0, \tilde{Q}_\epsilon}^2 + s^3 e^{2s\tilde{d}_\ell} \|u_\chi^{(k)}\|_{1, Q_L}^2 + s \mathfrak{h}_k(s) \right). \end{aligned} \quad (4.13)$$

Further, recalling (4.5), we see from (1.5) and (1.12) (resp., from (1.5), (1.6), (1.12) and (3.8)) that the first (resp., second) term of the sum in the right hand side of (4.13) is upper bounded up to some multiplicative constant, by $\sum_{j=0,1} \|e^{s\varphi} \nabla^j c_\chi\|_{0, Q_L}^2$ (resp., $e^{2s\tilde{d}_\ell} (\|u^{(k)}\|_{1, Q_L}^2 + 1)$). From this and Lemma 4.1 then follows for s sufficiently large that

$$\begin{aligned} & s \sum_{j=0,1} \|e^{s\varphi(\cdot, 0)} \nabla^j c_\chi\|_{0, \Omega_L}^2 \\ & \leq C \left(\sum_{j=0,1} \|e^{s\varphi} \nabla^j c_\chi\|_{0, Q_L}^2 + e^{2s\tilde{d}_\ell} + \sum_{k=2,3} \left(s^3 e^{2s\tilde{d}_\ell} \|u^{(k)}\|_{1, Q_L}^2 + s \mathfrak{h}_k(s) \right) \right). \end{aligned} \quad (4.14)$$

Step 3. Finally, we notice from (3.3) that

$$\|e^{s\varphi}\nabla^j c_\chi\|_{0,Q_L} = \|\rho_s^{1/2}e^{s\varphi(\cdot,0)}\nabla^j c_\chi\|_{0,\Omega_L}, \quad j = 0, 1, \quad (4.15)$$

where $\rho_s(x) := \int_{-T}^T e^{2s(\varphi(x,t)-\varphi(x,0))} dt = \int_{-T}^T e^{-2s\varphi(x,0)(1-\exp(-\gamma t^2))} dt$ for all $x \in \Omega_L$. Bearing in mind that $\varphi(x,0) \geq \kappa := e^{\gamma(\inf_{x' \in \omega} |x' - \delta a'|^2 - L^2)} > 0$ for all $x \in \Omega_L$, we get that

$$\|\rho_s\|_{L^\infty(\Omega_L)} \leq \int_{-T}^T e^{-2s\kappa(1-\exp(-\gamma t^2))} dt, \quad s > 0.$$

Therefore we have $\lim_{s \rightarrow +\infty} \rho_s = 0$, uniformly in Ω_L , by the dominated convergence theorem, so we derive from (4.14)-(4.15) that

$$s \sum_{j=0,1} \|e^{s\varphi(\cdot,0)}\nabla^j c_\chi\|_{0,\Omega_L}^2 \leq C \left(e^{2s\tilde{d}_\ell} + \sum_{k=2,3} \left(s^3 e^{2s\tilde{d}_\ell} \|u^{(k)}\|_{1,Q_L}^2 + s\mathfrak{h}_k(s) \right) \right), \quad (4.16)$$

upon taking s sufficiently large. With reference to (3.6)-(3.8), this entails that

$$\sum_{j=0,1} \|\nabla^j c_\chi\|_{0,\Omega_L}^2 \leq C \sum_{k=2,3} \left(s^2 e^{-2s(d_\ell - \tilde{d}_\ell)} + \mathfrak{h}_k(s) \right). \quad (4.17)$$

Here we used (1.5)-(1.12) and the embedding $\Omega_\ell \subseteq \Omega_L$ in order to substitute Ω_ℓ for Ω_L in the left hand side of (4.16). Now, taking into account that $\tilde{d}_\ell < d_\ell$, we end up getting the desired result from (4.17).

References

- [1] M. Bellassoued, *Global logarithmic stability in inverse hyperbolic problem by arbitrary boundary observation*, Inverse Problems, **20** (2004), 1033-1052.
- [2] M. Bellassoued, *Uniqueness and stability in determining the speed of propagation of second-order hyperbolic equation with variable coefficients*, Applicable Analysis, **83** (2004), 983-1014.
- [3] M. Bellassoued, M. Cristofol, and E. Soccorsi, *Inverse boundary value problem for the dynamical heterogeneous Maxwell's system*, Inverse Problems **28** (2012), 095009.
- [4] M. Bellassoued, M. Choulli and M. Yamamoto, *Stability estimate for an inverse wave equation and a multidimensional Borg-Levinson theorem*, J. Diff. Equat., **247**, 2 (2009), 465-494.
- [5] M. Bellassoued, D. Jellali and M. Yamamoto, *Lipschitz stability in an inverse problem for a hyperbolic equation with a finite set of boundary data*, Applicable Analysis **87** (2008), 1105-1119.
- [6] M. Bellassoued and M. Yamamoto, *Logarithmic stability in determination of a coefficient in an acoustic equation by arbitrary boundary observation*, J. Math. Pures Appl. **85** (2006), 193-224.

- [7] M. Bellassoued and M. Yamamoto, *Determination of a coefficient in the wave equation with a single measurement*, *Applicable Analysis* **87** (2008), 901-920.
- [8] M. Bellassoued and M. Yamamoto, *Carleman estimates with second large parameter for second order hyperbolic operators in a Riemannian manifolds*, *Applicable Analysis* (2011).
- [9] A. L. Bukhgeim and M. V. Klibanov, *Global uniqueness of class of multidimensional inverse problems*, *Soviet Math. Dokl.* **24** (1981), 244-247.
- [10] L. Cardoulis, M. Cristofol and L. Cardoulis, *Inverse problem for the Schrödinger operator in an unbounded strip*, *J. Inv. Ill-Posed Problems* **15** (2007), 1-20.
- [11] M. Choulli, Y. Kian and E. Soccorsi, *Determining the scalar potential in a periodic quantum waveguide from the DN map*, *New Prospects in Direct, Inverse and Control Problems for Evolution Equations*, *INdAM Series* **10**, chap. 5, Springer, 2014.
- [12] M. Choulli, Y. Kian and E. Soccorsi, *Stable determination of time-dependent scalar potential from boundary measurements in a periodic quantum waveguide*, arXiv:1306.6601.
- [13] M. Choulli and E. Soccorsi, *Recovering the twisting function in a twisted waveguide from the DN map*, to appear in *Jour. of Spectral Theory*.
- [14] R. Dautray and J.-L. Lions, *Mathematical analysis and Numerical Methods for Science and Technology*, Vol. 5, *Evolution problems I*, Springer Verlag, Berlin Heidelberg, 1992.
- [15] L. C. Evans, *Partial Differential Equations*, Amer. Math. Soc., Graduate Studies in Mathematics, vol. 19.
- [16] O. Imanuvilov, *On Carleman estimates for hyperbolic equations*, *Asymptotic Analysis* **32** (2002), 185–220.
- [17] O. Imanuvilov and M. Yamamoto, *Global Lipschitz stability in an inverse hyperbolic problem by interior observations*, *Inverse Problems* **17** (2001), 717-728.
- [18] O. Imanuvilov and M. Yamamoto, *Determination of a coefficient in an acoustic equation with single measurement*, *Inverse Problems* **19** (2003), 157-171.
- [19] V. Isakov and Z. Sun, *Stability estimates for hyperbolic inverse problems with local boundary data*, *Inverse Problems* **8** (1992), 193-206.
- [20] V. Isakov, *Inverse Problems for Partial Differential Equations*, Springer-Verlag, Berlin, 2006.
- [21] M. A. Kazemi and M. V. Klibanov, *Stability estimates for ill-posed Cauchy problems involving hyperbolic equations and inequality*, *Applicable Analysis* **50** (1993), 93-102.
- [22] Y. Kian, *Stability of the determination of a coefficient for the wave equation in an infinite waveguide*, arXiv:1305.4888.
- [23] Y. Kian, Q. S. Phan and E. Soccorsi, *Carleman estimate for infinite cylindrical quantum domains and application to inverse problems*, *Inverse Problems* **30**, 5 (2014), 055016.

- [24] Y. Kian, Q. S. Phan and E. Soccorsi, *Hölder stable determination of a quantum scalar potential in unbounded cylindrical domains*, arXiv:1311.5323.
- [25] M. V. Klibanov and J. Malinsky, *Newton-Kantorovich method for 3-dimensional potential inverse scattering problem and stability of the hyperbolic Cauchy problem with time dependent data*, Inverse Problems **7** (1991), 577–595.
- [26] M. V. Klibanov and M. Yamamoto, *Lipschitz stability of an inverse problem for an acoustic equation*, Applicable Analysis **85** (2006), 515-538.
- [27] J.-L. Lions, E. Magenes, *Problèmes aux limites non homogènes et applications*, vol. 1, Dunod (1968).
- [28] P. Stefanov and G. Uhlmann, *Stability estimates for the hyperbolic Dirichlet to Neumann map in anisotropic media*, J. Funct. Anal. **154** (1998), 330-358.
- [29] M. Yamamoto, *Uniqueness and stability in multidimensional hyperbolic inverse problems*, J. Math. Pures Appl. **78** (1999), 65-98.