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# Time resolved terahertz spectroscopy of low frequency electronic resonances and optical pump-induced terahertz photoconductivity in reduced graphene oxide membrane

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## ABSTRACT

Towards ultrafast optoelectronic applications of single and a few layer reduced graphene oxide (RGO), we study time domain terahertz spectroscopy and optical pump induced changes in terahertz conductivity of self supported RGO membrane in the spectral window of 0.5–3.5 THz. The real and imaginary parts of conductivity spectra clearly reveal low frequency resonances, attributed to the energy gaps due to the van Hove singularities in the density of states flanking the Dirac points arising due to the relative rotation of the graphene layers. Further, optical pump induced terahertz conductivity is positive, pointing to the dominance of intraband scattering processes. The relaxation dynamics of the photo excited carriers consists of three cooling pathways: the faster ( $\sim 450$  fs) one due to optical phonon emission followed by disorder mediated large momentum and large energy acoustic phonon emission with a time constant of a few ps (called the super collision mechanism) and a very large time ( $\sim 100$  ps) arising from the deep trap states. The frequency dependence of the dynamic conductivity at different delay times is analyzed in term of Drude Smith model.

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## 1. Introduction

Graphene is a  $sp^2$  hybridized 2D honeycomb lattice with electronic linear band dispersion having zero band gap at K and K' points of the Brillouin Zone. The carriers act like massless Dirac Fermions characterized by Fermi velocity  $v_F \approx 1.1 \times 10^6$  m/s. These special attributes give rise to unusual electronic properties such as high carrier mobility,

Klein paradox and anomalous integer quantum Hall effect at room temperature [1–7]. It has been recognized that graphene is a very attractive platform for potential applications in electronics and optoelectronics devices such as terahertz modulators [8], solar photovoltaic cells [9,10], mode locked lasers [11], ultrafast photodetectors [12] etc. Therefore the production of low cost large amount of graphene becomes an important issue. The highest quality graphene devices

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are still being made using mechanical exfoliation technique. The quality of graphene synthesized using chemical vapor deposition and epitaxially grown graphene on SiC is getting better with time. In recent years the ultrasonic assisted exfoliation of graphite to get a few layer graphene [13] and chemically produced reduced graphene oxide (RGO) by reduction of exfoliated graphene oxide (GO) [14,15] have been shown to be very promising routes for low cost large scale production. To realize RGO based high performance optoelectronic devices, discerning the effects of disorder and layer misorientation on the conductivity of RGO is an important issue. Terahertz (THz) spectral regime is of particular interest in case of graphene since it directly probes the carrier dynamics close to the Dirac point.

THz conductivity using time domain spectroscopy of chemical vapor deposited and epitaxial grown monolayer graphene has been previously reported<sup>1</sup> [16–19]. Similar measurements on reduced graphene oxide network films deposited on quartz substrates have also been carried out, which showed an increase of THz conductivity by reducing graphene oxide at higher temperatures [20]. Broadband optical absorption measurements from terahertz to visible frequency range have been used to get information about the number of layers and carrier density [21–24]. The optical response of back gated graphene has also been studied from infra red to THz region [25,26]. All these experimental measurements have shown Drude like frequency dependence of the conductivity from THz to mid infrared range. In comparison to Drude conductivity, plasmon related resonant features have been recently reported in CVD grown monolayer graphene [27]. Zou et al. [28] have shown a band at  $\sim 2.7$  THz in the conductivity spectra of twisted bilayer graphene, attributed to the interband transition across the energy gap arising from relative layer rotation induced van Hove singularity in the density of states. The tunneling differential conductance measurements do show the presence of rotation induced van Hove singularities flanking at different energies, depending on the angle of rotation in twisted graphene layers [29]. RGO is an interesting system with graphene like electronic band structure affected by the presence of large disorder and misorientation in layer stacking. Thus the conductivity and carrier relaxation in RGO are expected to be strongly influenced by the defects and the modified electronic structure. Very recently, the effect of trap density on carrier relaxation pathways in graphene oxide have been studied in the THz regime showing that the Auger recombination and trap assisted dynamics play a crucial role in cooling of the photoexcited carriers [30]. In contrast, transient absorption experiments [31] have shown that optical phonon emission and trap assisted cooling in RGO may dominate. Thus, the issue of non equilibrium carrier dynamics is still not understood well in RGO.

In this work, we report time resolved terahertz spectroscopy of self supported RGO membrane in the spectral range of 0.5–3.5 THz, clearly revealing three resonances in the real and imaginary parts of terahertz conductivity. These have been assigned to the enhanced density of states due to the layer rotation induced van Hove singularities at the saddle

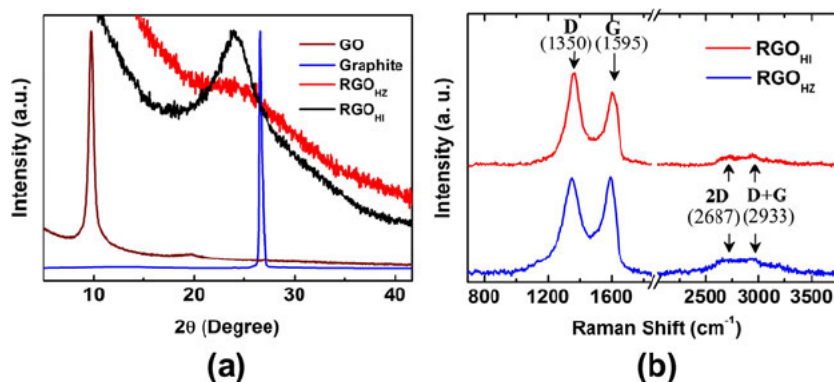
points in electronic dispersion of twisted two graphene layers in RGO [28]. The electronic part of the complex conductivity spectra is explained in terms of Drude Smith model incorporating the localization of carriers due to defects. Optical pump terahertz probe (OPTP) experiments show positive dynamic conductivity due to the dominance of intraband scattering processes. After pumping, the hot carriers thermalize generally within a few tenths of femtoseconds to a quasi Fermi Dirac distribution with electron temperature  $T_e$ . Afterwards, the hot carriers cool down via different scattering mechanisms like Auger scattering, plasmon emission, optical phonon emission, acoustic phonon emission etc. We observe three relaxation times: fast component with time constant of  $\sim 450$  fs due to emission of optical phonons, a slower component ( $\sim$  a few ps decay time) due to disorder mediated carrier acoustic phonon scattering and a very weak and slow component ( $\sim 100$  ps). The frequency dependence of the dynamic conductivity is also measured and explained in terms of Drude Smith model. The fitted parameters lead to the conclusion that after pump excitation the carriers gain enough energy to cross the nearby barriers resulting in higher average momentum relaxation time and provide an important view of the cooling dynamics of photoexcited carriers in RGO.

## 2. Experimental details

### 2.1. Sample preparation and characterization

Graphene oxide was prepared using simplified Hummer's method [15]. Briefly, graphite flakes (2 g) were added to a mixture of  $\text{H}_2\text{SO}_4$  (110 ml) and  $\text{H}_3\text{PO}_4$  (12 ml). After stirring the mixture for some time  $\text{KMnO}_4$  (11.1 g) was added slowly, followed by stirring at room temperature for 3 days and then adding  $\text{H}_2\text{O}_2$  (20 ml). Graphite oxide formed was washed repeatedly with 1 M HCl and distilled water, until the pH of the solution was 3–4. The resulting solution was centrifuged at 10,000 rpm to separate graphene oxide (GO) which was reduced chemically by two methods to get reduced graphene oxide: one using hydrazine hydrate (denoted by  $\text{RGO}_{\text{HZ}}$ ) and the other using hydroiodic (HI) acid (denoted by  $\text{RGO}_{\text{HI}}$ ). For hydrazine method [14], GO (50 mg) was dispersed in water (100 ml) and centrifuged at 3000 rpm for 30 min to remove any unexfoliated GO. Hydrazine hydrate (30  $\mu\text{L}$ , 64 wt% in water) and ammonia solution (0.8 ml, 25%) were added to the above dispersion. The resulting mixture was stirred at 100  $^\circ\text{C}$  for 5 h. The solution was centrifuged at 10,000 rpm to obtain  $\text{RGO}_{\text{HZ}}$ , following by filtering the dispersion of  $\text{RGO}_{\text{HZ}}$  in DMF (Dimethylformamide) through filter paper assisted by a vacuum pump. For the HI reduction, the procedure is as follows [32]: GO (100 mg) was dispersed in acetic acid (38 ml) followed by addition of the HI (2 ml) and stirring for 40 h at 40  $^\circ\text{C}$ . This was followed by centrifuging and washing the product with sodium bicarbonate, deionized water and acetone. The average thickness of  $\text{RGO}_{\text{HZ}}$  is  $\sim 370$  nm and  $\text{RGO}_{\text{HI}}$  is  $\sim 15$   $\mu\text{m}$ .

<sup>1</sup> Kar S., Mohapatra D. R., Freysz E., Sood A. K. Tuning photoinduced terahertz conductivity in monolayer graphene: Optical pump terahertz probe spectroscopy 2014;arXiv:1407.2752



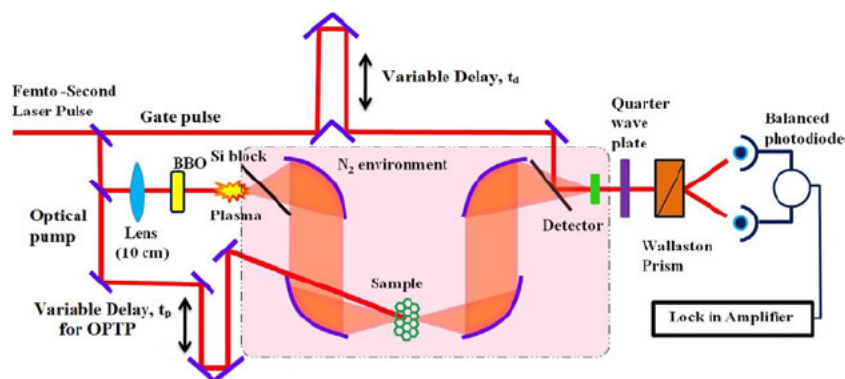
**Fig. 1 – (a)** X-ray diffraction patterns of Graphite (blue solid line), graphene oxide (GO) (wine solid line), RGO<sub>HZ</sub> (red solid line) and RGO<sub>HI</sub> (black solid line). **(b)** Raman spectra of RGO<sub>HZ</sub> (blue solid line) and RGO<sub>HI</sub> (red solid line). (A color version of this figure can be viewed online.)

The samples were characterized by UV visible spectroscopy, Raman spectroscopy, atomic force microscopy and X ray diffraction. The suspension has mostly single layer graphene with average lateral size of 2  $\mu\text{m}$ . Fig. 1a shows X ray diffraction pattern revealing that the RGO<sub>HZ</sub> has a broad peak at  $2\theta \sim 25^\circ$  (d spacing  $\sim 3.56 \text{ \AA}$ ) which correspond to the partial restacking of exfoliated graphene layers with two dimensional sheets along [002] direction. A more prominent peak at  $2\theta \sim 25^\circ$  is observed for RGO<sub>HI</sub> (Fig. 1a) indicating better quality graphene. In comparison, graphite exhibits a sharp peak at  $2\theta \sim 26.5^\circ$ , corresponding to the interlayer distance  $\sim 3.34 \text{ \AA}$  and GO shows a peak at  $2\theta \sim 9.7^\circ$  (d spacing  $\sim 9.11 \text{ \AA}$ ) due to oxygen containing functional group and intercalated water molecules. Fig. 1b shows the Raman spectra of the RGO<sub>HZ</sub> and RGO<sub>HI</sub> membranes recorded at room temperature in backscattering geometry using  $\lambda = 514 \text{ nm}$  laser light and Witec Raman spectrometer. Raman bands are observed at  $1595 \text{ cm}^{-1}$  (G band),  $1350 \text{ cm}^{-1}$  (D band) and  $2687 \text{ cm}^{-1}$  (2D band). The size of defect free domain,  $L_D$ , can be estimated by using the intensity ratios of the D and G bands [33]:  $L_D^2 (\text{nm}^2) = (1.8 \pm 0.5) \times 10^{-9} \lambda^4 (I_D/I_G)^{-1}$  giving  $L_D \sim 11 \pm 2 \text{ nm}$  for RGO<sub>HZ</sub> and  $\sim 8 \pm 2 \text{ nm}$  for RGO<sub>HI</sub>.

## 2.2. THz experimental set up

THz radiation is generated via two color laser filament using Ti: Sapphire laser delivering  $\sim 70 \text{ fs}$  pulse of maximum energy

$\sim 1.6 \text{ mJ}$  per pulse at a central wavelength of  $785 \text{ nm}$  with repetition rate  $1 \text{ kHz}$ . The beam is divided into three parts: (1) to generate terahertz, (2) to probe terahertz and (3) for pumping the sample. The pump excitation energy is varied by using neutral density filters to give a fluence varying from  $3.4$  to  $340 \mu\text{J}/\text{cm}^2$ . The fundamental laser beam and its second harmonic laser beams are focused in air to produce plasma. Terahertz radiation generated from the plasma [34] is focused at the sample position by using a pair of off axis parabolic mirrors of  $10 \text{ cm}$  focal length. The unwanted light generated from the plasma is blocked by  $1 \text{ mm}$  thick high resistive silicon wafer. After propagating through the sample the terahertz beam is collected and focused at  $1 \text{ mm}$  thick ZnTe detector by a pair of parabolic mirrors. The whole terahertz beam path was purged with dry nitrogen to avoid unwanted absorptions by water vapors. Terahertz time domain measurements were performed in transmission geometry. Photo induced changes in terahertz transmission due to optical pump at  $785 \text{ nm}$  was measured at the peak position of the terahertz field. Fig. 2 shows a schematic of the experimental set up to perform optical pump terahertz probe transmission measurements. While measuring the terahertz field passing through unexcited sample, the chopper was placed in the terahertz generation path and the gate delay line  $t_d$  was used. In case of the photoinduced transmission measurements, the chopper was placed in the path of the pump beam exciting the RGO membrane and the pump delay line  $t_p$  was used.



**Fig. 2 – Schematic of the THz experimental set up for time resolved terahertz spectroscopy and optical pump-terahertz probe spectroscopy.** (A color version of this figure can be viewed online.)

The spot size of the terahertz beam on the sample position was 2 mm while the size of pump beam (785 nm) to excite the sample was 5 mm. The estimated terahertz peak field was  $\sim 26$  kV/cm.

### 3. Results and discussion

#### 3.1. THz conductivity measurements

Fig. 3a shows the terahertz electric field through air  $E_{\text{ref}}(t)$  and through the RGO membrane  $E_{\text{sam}}(t)$ . After taking the Fourier Transform, the spectral function is evaluated as  $T(\omega) = \frac{E_{\text{sam}}(\omega)}{E_{\text{ref}}(\omega)} A(\omega) \exp(i\theta(\omega))$ , where  $A(\omega)$  is the amplitude and  $\theta(\omega)$  is the phase of the transmission function given by [35,36]

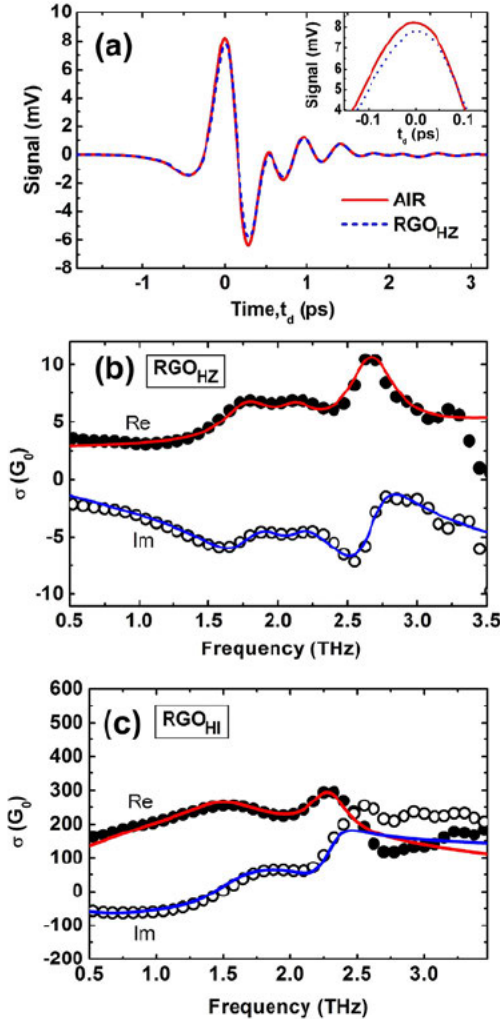


Fig. 3 – (a) Terahertz time domain signal through the RGO<sub>HZ</sub> membrane (blue dotted line) and air (continuous line). Inset: enlarged part of the THz pulse near the maximum. (b,c) Real and imaginary part of the conductivity spectra. Filled circles represent real part of conductivity and open circles represent imaginary part. The solid lines represent the fitted curves of the Drude-Smith model along with the Lorentz model. (A color version of this figure can be viewed online.)

$$T(\omega) = \frac{\frac{4n^*(\omega)}{[n^*(\omega)+1]^2} \exp[i\omega(n^*(\omega)-1)d/c]}{1 - \frac{[n^*(\omega)-1]^2}{[n^*(\omega)+1]^2} \exp[2i\omega n^*(\omega)d/c]} \quad (1)$$

where  $n^*(\omega)$  is the complex refractive index of the membrane,  $d$  is the thickness of the membrane and  $c$  is velocity of light. In the limit  $n^*\omega d/c \ll 1$ , one can approximate Eq. (1) for the membrane as  $T(\omega) \approx \frac{2}{Z_0\sigma(\omega)}$ , where  $Z_0 = 377 \Omega$  is impedance of free space and  $\sigma(\omega)$  is complex conductivity of the membrane. Now we will present the results of RGO<sub>HZ</sub>. Fig. 3b shows the real (filled circles) and imaginary part (open circles) of the complex conductivity of the RGO<sub>HZ</sub> membrane in the spectral range of 0.5 to 3.5 THz in units of  $G_0$  ( $2e^2/h = 77.3 \mu S$ ). It can be noted that the imaginary part of conductivity is negative throughout the whole spectral range and three resonances are clearly discernible at 1.8, 2.2 and 2.7 THz in the real and imaginary parts of the terahertz conductivity. We find that Drude model used for single layer graphene [16,17,20] cannot account for the contribution of the carriers to  $\sigma(\omega)$ . One should note that the graphene basal planes of GO are decorated non uniformly by oxygen containing functional groups (mainly epoxy and hydroxyl). The presence of such functional groups results in small ordered  $sp^2$  clusters isolated within  $sp^3$  C O matrix [37]. The thermal reduction of GO creates new  $sp^2$  domains, which provide percolation pathways between  $sp^2$  clusters [37]. These new  $sp^2$  clusters can provide hopping pathways for the carrier transport. Such localized finite size molecular  $sp^2$  clusters embedded within  $sp^3$  matrix lead to confinement of  $\pi$  electrons in RGO. Drude Smith model takes into account such carrier localization [38]. Accordingly, the complex conductivity spectra shown in Fig. 3b has been fitted using

$$\sigma = \frac{D_0\tau}{1 - i\omega\tau} \left( 1 + \frac{C_1}{1 - i\omega\tau} \right) + \sum_{j=1}^3 \frac{iF_j\omega}{(\omega^2 - \omega_{0j}^2) + i\omega\gamma_j} \quad (2)$$

where  $D_0$  is the Drude weight,  $\tau$  is momentum relaxation time,  $F_j$  is the oscillator strength of the  $j$ 'th oscillator with angular frequency  $\omega_{0j}$  and lifetime broadening parameter  $\gamma_j$ . The first term is the Drude Smith part for the contribution of the carriers to the conductivity where  $C_1$  ( $-1 \leq C_1 \leq 0$ ) is related to the degree of localization of the carriers;  $C_1 = 1$  corresponds to complete localization and  $C_1 = 0$  is the Drude limit giving free carrier response. The fitted curves shown by the red and blue solid lines in Fig. 3b, following Eq. (2) reproduce the conductivity spectra very well. The best fit is obtained by using the parameter  $C_1 = 0.947 \pm 0.004$ , indicating a significant degree of carrier localization. As mentioned above, the  $sp^2$  hybridized carbon regions are separated by the  $sp^3$  matrix [37] resulting in a potential barrier between the adjacent  $sp^2$  regions. The carriers can either cross the barrier to reach the adjacent  $sp^2$  domains, or they can backscatter resulting in carrier localization. The  $C_1$  parameter takes into account the relative probabilities of these two extremes [38]. In case of  $sp^2$  domains fully isolated by large  $sp^3$  regions, the probability of reflection is high giving  $|C_1| \sim 1$ . If there are other  $sp^2$  clusters in the immediate vicinity, the carriers with sufficient energy may cross the  $sp^2$  domain boundaries by either hopping or tunneling processes, resulting in  $C_1$  between 0 and  $-1$ . We further note that the fitted value of the momentum relaxation time  $\tau$  is as small as  $6.7 \pm 0.5$  fs.

We now estimate the Fermi energy  $E_F$  by extrapolating fitted real part of conductivity to  $\omega = 0$  ( $\sigma_{DC} = 2.9G_0$ ). Since the RGO membrane consists of randomly stacked graphene layers, we can assume that the carriers are distributed equally in each layer. As a consequence, the DC conductivity can be expressed [23] as  $\sigma_{DC} = \tau \frac{v_F^2}{\pi\hbar} \sqrt{n\pi N}$ , where  $n$  is carrier density in each layer and  $N$  is the number of graphene layers. The approximate number of graphene layers in the RGO<sub>HZ</sub> membrane is  $N \sim 370/.356 \approx 1040$ . Further, the carrier density  $n$  is related to Fermi energy  $E_F$  by  $E_F = \hbar v_F \sqrt{n\pi}$ . By taking  $\tau = 6.7$  fs and  $\sigma_{DC} = 2.9G_0$ , we get Fermi energy to be  $\sim 8$  meV.

Now we focus to understand the three resonances located at  $\nu_1 \sim 1.8$  THz (7.4 meV),  $\nu_2 \sim 2.2$  THz (9.0 meV) and  $\nu_3 \sim 2.7$  THz (11.0 meV). There are recent theoretical studies on low frequency ( $< 150$  cm<sup>-1</sup>) phonon modes in layered crystals (such as a few layer Bernal stacked graphene and h BN) associated with the movements of planes perpendicular to the hexagonal axis (shear modes) and parallel to the hexagonal axis (compression modes) [39]. Some of these modes have been observed recently by femto second coherent phonon spectroscopy [40] and by X ray scattering [41]. The three resonances observed by us in the RGO membrane, though very close to the calculated values of the compression modes for a few layer graphene, may not be associated with these phonons because the graphene layers in RGO do not have perfect Bernal stacking, as assumed in the theoretical calculations [39]. Recently, terahertz conductivity of twisted bilayer graphene has shown a resonance in real part of  $\sigma(\omega)$  at  $\sim 2.7$  THz [28]. This has been assigned to an enhanced density of electronic states from the van Hove singularities (VHS) arising from the two twisted graphene layers. We make a similar assignment: The RGO membrane is expected to have graphene layers stacked in turbostratic manner [42]. The rotation between the layers gives rise to Moiré Pattern and hence a superlattice on the graphite surface [29]. For simplicity we consider a graphene bilayer. In Bernal stacking A1 (B1) atom of layer 1 sits on B2 (A2) atom of layer 2 and B1 (A2) atom is at the middle of layer 2 (1). The condition for commensurate rotation by an angle  $\theta$  about B2 (A1) to produce Moiré pattern is given by [29]  $\cos(\theta_j) = \frac{3j^2+3j+1/2}{3j^2+3j+1}$  where  $j$  being an integer.  $j = 0$ ,  $\theta = 60^\circ$  corresponds to AA stacked and  $j \rightarrow \infty$ ,  $\theta = 0^\circ$  corresponds to AB stacked bilayer. The relative rotation of the layers causes a shift between Dirac points in momentum space by  $\Delta K = 2K \sin(\theta/2)$  with  $K = 4\pi/3a_0$  where  $a_0 = 2.46$  Å. The two Dirac cones intersect near middle of the superlattice Brillouin zone and hybridize resulting in saddle point in energy dispersion and VHSs in the density of states. The hybridized states near zero energy result from the coupling of the zero energy states of one layer with states of energy  $\pm \hbar v_F \Delta K$  in other layer with interlayer hopping amplitude [43],  $t_\perp^0 \approx 0.4t_\perp$ . It is important to note that the cases where  $t_\perp^0 = 0$  but  $\theta \neq 0$ , the VHS do not appear even though Moiré patterns can form [29]. For small angle of rotation the energy difference between the VHSs is given by  $\Delta E_{VHS} = \frac{8\pi\hbar v_F}{3a_0} |\sin(\theta/2)| \cdot 2t_\perp^0$ . Thus  $\Delta E_{VHS}$  can be easily tuned by changing the twisting angle, the interlayer coupling and strain [44]. The rotation induced generation of VHS is not limited to graphene bilayer but also to multi layer graphene, provided  $t_\perp^0 \neq 0$ . The value of  $t_\perp^0$  is close to 0.1 eV [45]. Taking  $\Delta E_{VHS}$  as the three observed resonances,

a consistence choice of the parameters are:  $j = 28$ ,  $\theta = 1.16^\circ$  and  $t_\perp^0 = 0.110, 0.109$  and  $0.108$  eV, corresponding to the observed resonances at 1.8, 2.2 and 2.7 THz, respectively. A small spread in the value of  $t_\perp^0$  can arise from the small variation in the interlayer spacing.

To see the effect of film thickness and reduction condition, we have measured the terahertz conductivity of RGO<sub>HI</sub>. It can be seen from Fig. 3c that the two electronic resonances occur at 1.5 and 2.3 THz which are again due to interband transition across band gap arising from rotation induced generation of VHS in density of states. The zero frequency conductivity of RGO<sub>HI</sub> in unit of S/m is 500 S/m which is similar to 600 S/m in RGO<sub>HZ</sub>. The value of the momentum relaxation time  $\tau$  is  $98 \pm 1$  fs and the parameter  $C_1 = 0.77 \pm 0.01$  which corroborate the fact that the oxygen functional groups in RGO<sub>HI</sub> are less than in RGO<sub>HZ</sub> [32].

### 3.2. Optical pump - terahertz probe studies

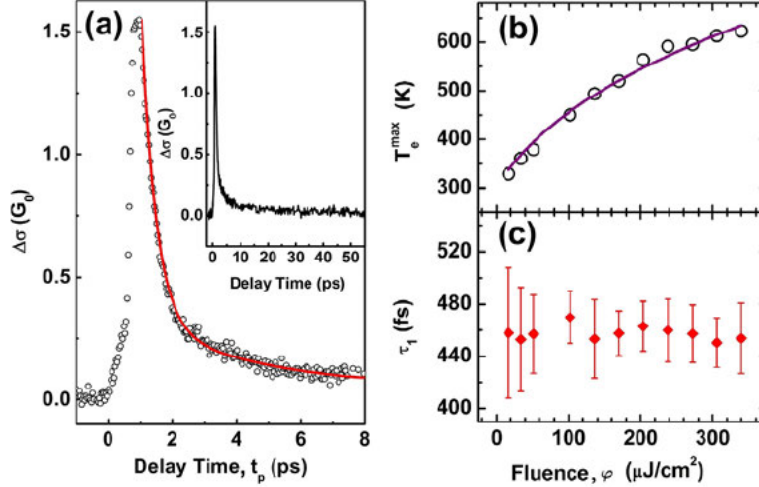
We now investigate the changes in terahertz conductivity by optical pump using time resolved optical pump terahertz probe spectroscopy. To measure the photoinduced THz electric field transmission, we follow the procedure shown in Fig. 2. The sample is photoexcited with 785 nm ( $\sim 70$  fs pulses) pump. The gate line is kept fixed at the peak of the THz electric field ( $t_d = 0$ ). The pump induced change of the peak of THz waveform is measured by changing the delay line  $t_p$  of optical pump in a way such that the optical pump reaches the sample earlier than the THz probe. The photoinduced change in terahertz transmission:  $\Delta T = T(\text{pump on}) - T(\text{pump off})$ , is related to photoinduced conductivity  $\Delta\sigma$  by the relation  $\frac{\Delta T}{T} \approx \frac{Z_0}{2} \Delta\sigma$ . Now we will present the results of RGO<sub>HZ</sub>. Fig. 4a shows the photoinduced change of conductivity measured at the peak of the terahertz electric field ( $t_d = 0$ ) as a function of the delay time between optical pump and terahertz probe pulses using the pump fluence of  $238 \mu\text{J}/\text{cm}^2$ . Here the photoconductivity  $\Delta\sigma$  is positive, similar to earlier reports on single layer graphene [22,46]. We must note that negative  $\Delta\sigma$  has also been observed in graphene [16,17,47,48] and tuning of  $\Delta\sigma$  has been addressed in recent studies [16,17,49]. The positive  $\Delta\sigma$  is due to dominance of intraband scattering of photoinduced carriers over the other cooling channel involving secondary hot carrier generation [147].  $\Delta\sigma$  can be related to corresponding electron temperature  $T_e$  by using

$$\Delta\sigma = \sigma_{\text{intra}}(T_e) - \sigma_{\text{intra}}(T_0),$$

$$\sigma_{\text{intra}}(T) = \frac{2G_0\tau k_B T}{\hbar(\omega^2\tau^2 + 1)} \ln \left[ 2 \cosh \left( \frac{E_F}{2k_B T} \right) \right] \quad (3)$$

Here  $T_0$  ( $\sim 300$  K) is the lattice temperature. The value of  $\tau$  is taken 30 fs obtained from fitting the frequency dependence of the dynamic conductivity (to be shown later). Fig. 4b shows the fluence dependence of  $T_e^{\text{max}}$  corresponding to  $\Delta\sigma^{\text{max}}$ . The solid line is a fit to the equation  $T_e^{\text{max}} = (T_0^3 + B\phi)^{1/3}$  [50,51] whose origin is in the linear dependence of electronic density of states on energy as in graphene. The best fit is obtained for  $B = 6.7 \times 10^7 \text{ J}^{-1} \text{ m}^2 \text{ K}^3$ .

We now focus on the cooling dynamics of the hot carriers. The decay of  $\Delta\sigma$  consists of three components; a fast relaxa



**Fig. 4 – (a)** Photoinduced change of Terahertz conductivity  $\Delta\sigma$  of RGO<sub>HZ</sub>. The inset shows the decay dynamics up to 55 ps. The solid red line in the main panel is a fit to super-collision model including the fast relaxation component  $\tau_1$ . **(b)**  $T_e^{\max}$  plotted as a function of pump fluence. Solid line is fit to function  $(T_0^3 + B\phi)^{1/3}$ . **(c)** The fast decay component  $\tau_1$  is displayed as function of  $\phi$ . The fitted  $\tau_1$  shows  $\phi$  independent behavior. (A color version of this figure can be viewed online.)

tion with time constant  $\tau_1$  (a few hundred fs), a slow component  $\tau_2$  (a few ps) and a small third component with long decay time  $\tau_3$ . The origin of  $\tau_1$  can be Auger recombination [30] or optical phonon emission [31,50]. For the former process,  $\tau_1$  should decrease with  $\phi$  as observed in graphene oxide. This is however not the case in the present study in Fig. 4c. Thus,  $\tau_1$  can be associated with the optical phonon emission instead of Auger recombination. This difference from Ref. [30] may arise due to different method of reducing GO (Ref. [30] has used thermal annealing as compared to chemical reduction by us).

The electron optical phonon scattering continues till the hot carriers energy is reduced to less than the optical phonon energy ( $\sim 200$  meV). Below 200 meV, emission of acoustic phonons is the only possible way of energy relaxation permitted by momentum conservation. At each scattering event, very small energy ( $k_B T_{BG} - 2v_s E_F / v_F$ ,  $v_s$  is sound velocity) is dissipated and thus the cooling process can last as long as 300 ps [46]. Here,  $T_{BG}$  is Bloch Gruneisen temperature defining a boundary between the low and the high temperature response. A recent proposal for the relaxation process with time constant  $\tau_2$  is the emission of high energy ( $\sim k_B T_e$ ) and high momentum ( $\sim k_B T_e / \hbar v_s$ ) acoustic phonons mediated by disorder (for wave vector conservation) and is termed as super collision cooling (SC) with relaxation time of a few ps [52]. For the SC mechanism the rate equation governing the cooling of the carriers is given by [52–54]

$$\frac{\partial T_e(t)}{\partial t} = \frac{A}{\alpha} \frac{T_e^3(t) - T_0^3}{T_e(t)} \quad (4)$$

Here  $A/\alpha$  is the SC rate coefficient which is related to deformation potential,  $D_{ac}$  by [53,54]

$$\frac{A}{\alpha} = \frac{6\zeta(3)}{\pi^2} \frac{g}{k_F \ell} \frac{k_B}{\hbar} \simeq \frac{2}{3} \frac{g}{k_F \ell} \frac{k_B}{\hbar} \quad (5)$$

where  $g$ , electron phonon coupling strength is  $2D_{ac}^2 E_F / \rho v_s^2 \pi (\hbar v_F)^2$ ,  $\zeta$  is Riemann zeta function and  $\rho \sim 7.6 \times 10^{-7} \text{ kg/m}^2$  is

density of graphene. The fitted curve shown by solid line in Fig. 4a consists of the initial decay, taken in the form  $A_1 \exp(-t/\tau_1)$  to represent electron phonon scattering process along with decay of  $T_e$  due to SC mechanism given by Eq. (4). The best fit is obtained using  $A_1 = 0.78$ ,  $\tau_1 = 450$  fs,  $\tau = 30$  fs and  $A/\alpha = 2.6 \times 10^8 \text{ K}^{-1} \text{ s}^{-1}$ . The value of  $A/\alpha$  is very close to the reported rate coefficient for the suspended graphene [53]. Using Eq. (5) the deformation potential  $D_{ac}$  is  $\sim 15$  eV, which is close to the reported experimental values [55,56].

A very slow relaxation time of  $\tau_3 \sim 100$  ps is observed in the relaxation dynamics [ $A_3 \exp(-t/\tau_3)$  with  $A_3 = 0.02$ ]. This term can be related to the trapping of the electrons in deep trap states of RGO<sub>HZ</sub> and was also observed by transient absorption measurements [31]. The trap density  $N_t$  can be estimated from  $\tau_3$  by using  $N_t^{-1} = \tau_3 B v_F$  where  $B$  is the absorption cross section [57] taken as  $\sim 10^{-14} \text{ cm}^2$ . The estimated trap density  $N_t$  is thus  $9 \times 10^{15} / \text{cm}^3$ , close to the trap densities observed by current voltage measurements [58,59]. The OPTP measurements on RGO<sub>HI</sub> show similar relaxation dynamics and fluence dependence.

We now measure the frequency dependence of  $\Delta\sigma(\omega)$  at different times of carrier relaxation. At time  $t_p$  after photoexcitation, the THz pulse was sent through the membrane. The mapping of the photoinduced change in transmission ( $\Delta T$ ) of the THz waveform in time domains was done by varying the arrival time of the gate pulse and the THz pulse using the time delay  $t_d$ . The mapped  $\Delta T$  is shown in Fig. 5. Since the gate line was used, each point of the terahertz waveform experiences a different delay after photoexcitation and thus probe the RGO<sub>HZ</sub> system in different states over a time scale of the duration of the THz waveform. Thus, the information recorded is an average of the non equilibrium dynamics of the hot carriers over a few ps. The complex  $\Delta\sigma(\omega)$ , extracted from  $\Delta T(t_d)$  at different  $t_p$  after photoexcitation are shown in Fig. 6. The real part of  $\Delta\sigma$  is positive, increasing with frequency while the imaginary part is completely negative and decreases with frequency. The  $\Delta\sigma(\omega)$  is well described

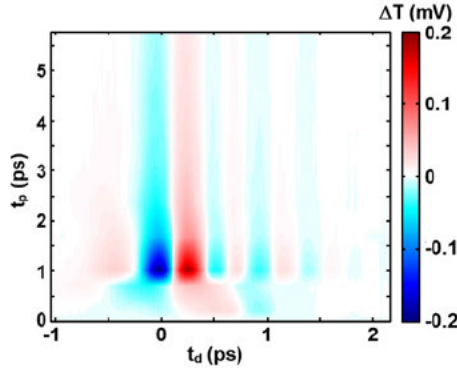


Fig. 5 – Mapping of the Photo-induced change of THz transmission ( $\Delta T$ ). (A color version of this figure can be viewed online.)

by Drude Smith model (given by the first term in Eq. (2)) after  $t_p = 3$  ps. Below 3 ps the model does not fit well as expected due to very high rate of carrier relaxation just after achieving the quasi equilibrium condition. The fitted parameter, momentum relaxation time  $\tau$  is  $30 \pm 2$  fs (and is independent of  $t_p$ ) and  $|C_1|$  increases with  $t_p$  as shown in Fig. 7a. The increase of  $\tau$  from 6.7 to 30 fs after photoexcitation can arise since the hot carriers have enough energy to cross the nearby barriers resulting in higher average momentum relaxation time. Further, the increase of  $|C_1|$  with  $t_p$  confirms the same. The extracted Drude weight ( $D_0$ ) is plotted in Fig. 7b as a function of  $t_p$ . The red solid line is a fit to the function  $D_0 = D_0^0 + D_t \exp(-\frac{t_p - t_0}{\tau_D})$  with  $t_0 = 1$  ps (corresponds to the maximum in OPTP signal shown in Fig. 4a) and the fitted parameters are  $D_0^0 = 2.0 \times 10^9 \Omega^{-1} \text{s}^{-1}$ ,  $D_t = 1.1 \times 10^{10} \Omega^{-1} \text{s}^{-1}$  and

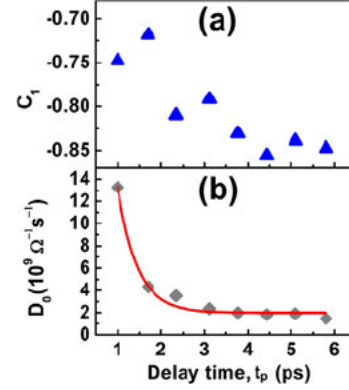


Fig. 7 – Fitted parameters (a)  $C_1$  and (b)  $D_0$  are plotted as a function of  $t_p$ . The red solid line is the fit to the exponential function with time constant  $\sim 450$  fs as described in the text. (A color version of this figure can be viewed online.)

$\tau_D = 450$  fs. The time constant  $\tau_D$  is same as  $\tau_1$  due to electron optical phonon scattering processes.

#### 4. Conclusions

Terahertz time domain measurements of two kinds of RGO membranes are carried out to measure the conductivity from 0.5 to 3.5 THz. The terahertz conductivity spectra reveal deviation from the Drude form and are explained in terms of Drude Smith model which takes into account the back scattering of the carriers. In addition, three resonances are observed in the conductivity spectra of RGO<sub>HZ</sub>. The origin of these resonances is best attributed to the enhanced density of states near the Dirac cone due to rotation of adjacent

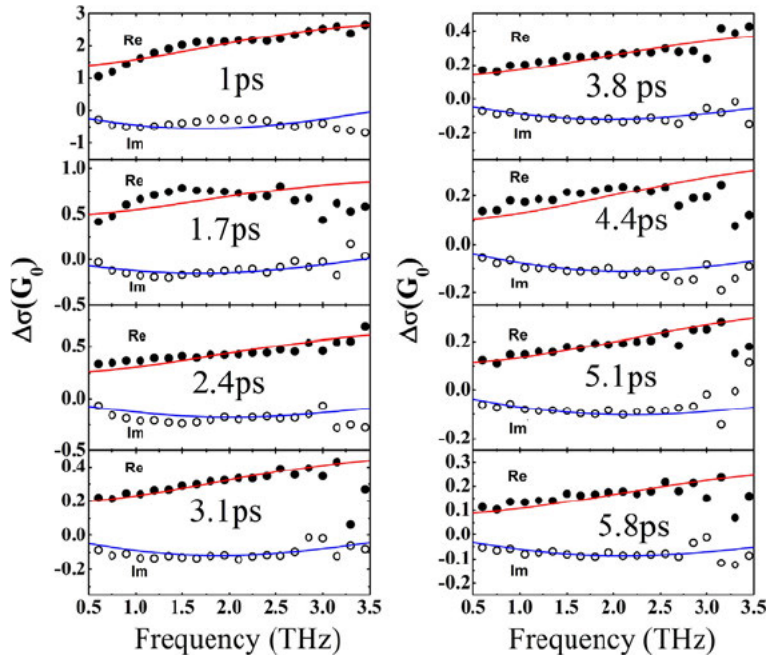


Fig. 6 – Spectral dependence of photo-induced dynamic conductivity  $\Delta\sigma(\omega)$  at different delay time  $t_p$ . The real part is shown by closed circles and imaginary part by open circles. The solid lines are fit to the Drude-Smith model. (A color version of this figure can be viewed online.)

graphene layers. A control on the relative orientation of the layers in graphene is important to tune the electronic structure of RGO. Photoexcited dynamics was studied by optical pump-terahertz probe spectroscopy where a positive value of dynamic photoconductivity is observed that can be explained by intraband scattering processes. The cooling dynamics is explained in terms of fast cooling involving optical phonons along with the super collision cooling involving acoustic phonons and disorder. The fitted parameters of super collision cooling model are used to estimate the deformation potential ( $\sim 15$  eV). We hope that the present study will help in evaluating the application potential of RGO for ultrafast optoelectronic applications. It will be interesting to study reduced graphene oxide with systematic quantitative variation of degree of oxygen functionalities.

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