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UNIQUENESS AND RIGIDITY IN NONLINEAR ELLIPTIC EQUATIONS, INTERPOLATION INEQUALITIES, AND SPECTRAL ESTIMATES

by Jean DOLBEAULT and Michał KOWALCZYK (*)

ABSTRACT. This paper is devoted to the Lin-Ni conjecture for a semi-linear elliptic equation with a super-linear, sub-critical nonlinearity and homogeneous Neumann boundary conditions. We establish a new rigidity result, that is, we prove that the unique positive solution is a constant if the parameter of the problem is below an explicit bound that we relate with an optimal constant for a Gagliardo-Nirenberg-Sobolev interpolation inequality and also with an optimal Keller-Lieb-Thirring inequality. Our results are valid in a sub-linear regime as well. The rigidity bound is obtained by nonlinear flow methods inspired by recent results on compact manifolds, which unify nonlinear elliptic techniques and the *carré du champ* method in semi-group theory. Our method requires the convexity of the domain. It relies on integral quantities, takes into account spectral estimates and provides improved functional inequalities.

Unicité et rigidité pour des équations elliptiques non linéaires, inégalités d'interpolation et estimations spectrales

RÉSUMÉ. Cet article est consacré à la conjecture de Lin-Ni pour une équation semi-linéaire elliptique avec non-linéarité super-linéaire, sous-critique et des conditions de Neumann homogènes. Nous établissons un résultat de rigidité, c'est-à-dire nous prouvons que la seule solution positive est constante si le paramètre du problème est en-dessous d'une borne explicite, reliée à la constante optimale d'une inégalité d'interpolation de Gagliardo-Nirenberg-Sobolev et aussi à une inégalité de Keller-Lieb-Thirring optimale. Nos résultats sont également valides dans un régime sous-linéaire. La borne de rigidité est obtenue par des méthodes de flots non-linéaires inspirées de résultats récents sur les variétés compactes, qui unifient des techniques d'équations elliptiques non-linéaires et la méthode du carré du champ en théorie des semi-groupes. Notre méthode requiert la convexité du domaine. Elle repose sur des quantités intégrales, prend en compte des estimations spectrales et fournit des inégalités améliorées.

Keywords: semilinear elliptic equations; Lin-Ni conjecture; Sobolev inequality; interpolation; Gagliardo-Nirenberg inequalities; Keller-Lieb-Thirring inequality; optimal constants; rigidity results; uniqueness; *carré du champ* method; $CD(\rho, N)$ condition; bifurcation; multiplicity; generalized entropy methods; heat flow; nonlinear diffusion; spectral gap inequality; Poincaré inequality; improved inequalities; non-Lipschitz nonlinearity; compact support principle.

Math. classification (2010): 35J60; 26D10; 46E35.

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1. Introduction and main results

Let us assume that Ω is bounded domain in \mathbb{R}^d with smooth boundary. To avoid normalization issues, we shall assume throughout this paper that

$$|\Omega| = 1.$$

The unit outgoing normal vector at the boundary is denoted by n and $\partial_n u = \nabla u \cdot n$. We shall denote by $2^* = \frac{2d}{d-2}$ the critical exponent if $d \geq 3$ and let $2^* = \infty$ if $d = 1$ or 2 . Assume first that p is in the range $1 < p < 2^* - 1 = (d+2)/(d-2)$ if $d \geq 3$, $1 < p < \infty$ if $d = 1$ or 2 , and let us consider the three following problems.

(P1) For which values of $\lambda > 0$ does the equation

$$-\Delta u + \lambda u = u^p \quad \text{in } \Omega, \quad \partial_n u = 0 \quad \text{on } \partial\Omega \quad (1.1)$$

has a unique positive solution ?

(P2) For any $\lambda > 0$, let us define

$$\mu(\lambda) := \inf_{u \in H^1(\Omega) \setminus \{0\}} \frac{\|\nabla u\|_{L^2(\Omega)}^2 + \lambda \|u\|_{L^2(\Omega)}^2}{\|u\|_{L^{p+1}(\Omega)}^2}.$$

For which values of $\lambda > 0$ do we have $\mu(\lambda) = \lambda$?

(P3) Assume that ϕ is nonnegative function in $L^q(\Omega)$ with $q = \frac{p+1}{p-1}$ and denote by $\lambda_1(\Omega, -\phi)$ the lowest eigenvalue of the Schrödinger operator $-\Delta - \phi$. Let us consider the optimal inequality

$$\lambda_1(\Omega, -\phi) \geq -\nu(\|\phi\|_{L^q(\Omega)}) \quad \forall \phi \in L^q_+(\Omega)$$

For which values of μ do we know that $\nu(\mu) = \mu$?

The three problems are related. Uniqueness in (P1) means that $u = \lambda^{1/(p-1)}$ while equality cases $\mu(\lambda) = \lambda$ in (P2) and $\nu(\mu) = \mu$ and (P3) are achieved by constant functions and constant potentials respectively. We define a threshold value μ_i with $i = 1, 2, 3$ such that the answer to (Pi) is yes if $\mu < \mu_i$ and no if $\mu > \mu_i$.

Our method is not limited to the case $p > 1$. If p is in the range $0 < p < 1$, the three problems can be reformulated as follows.

(P1) For which values of $\lambda > 0$ does the equation

$$-\Delta u + u^p = \lambda u \quad \text{in } \Omega, \quad \partial_n u = 0 \quad \text{on } \partial\Omega \quad (1.2)$$

has a unique nonnegative solution ?

(P2) For any $\mu > 0$, let us define

$$\lambda(\mu) := \inf_{u \in H^1(\Omega) \setminus \{0\}} \frac{\|\nabla u\|_{L^2(\Omega)}^2 + \mu \|u\|_{L^{p+1}(\Omega)}^2}{\|u\|_{L^2(\Omega)}^2}.$$

For which values of $\mu > 0$ do we have $\lambda(\mu) = \mu$?

(P3) Assume that ϕ is nonnegative function in $L^q(\Omega)$ with $q = \frac{1+p}{1-p}$ and still denote by $\lambda_1(\Omega, \phi)$ the lowest eigenvalue of the Schrödinger operator $-\Delta + \phi$. Let us consider the optimal inequality

$$\lambda_1(\Omega, \phi) \geq \nu(\|\phi^{-1}\|_{L^q(\Omega)}) \quad \forall \phi \in L_+^q(\Omega)$$

For which values of μ do we know that $\nu(\mu) = \mu$?

The problems of the range $0 < p < 1$ and $1 < p < 2^*$ can be unified. Let us define

$$\varepsilon(p) = \frac{p-1}{|p-1|}$$

and observe that $\|u\|_{L^{p+1}(\Omega)} \leq \|u\|_{L^2(\Omega)}$ if $p < 1$, $\|u\|_{L^2(\Omega)} \leq \|u\|_{L^{p+1}(\Omega)}$ if $p > 1$, so that $\varepsilon(p) (\|u\|_{L^{p+1}(\Omega)} - \|u\|_{L^2(\Omega)})$ is nonnegative. Our three problems can be reformulated as follows.

(P1) Let us consider the equation

$$-\varepsilon(p) \Delta u + \lambda u - u^p = 0 \quad \text{in } \Omega, \quad \partial_n u = 0 \quad \text{on } \partial\Omega \quad (1.3)$$

and define

$$\mu_1 := \inf\{\lambda > 0 : (1.3) \text{ has a non constant solution}\}.$$

We shall say that *rigidity* holds in (1.3) if $u = \lambda^{1/(p-1)}$ is its unique positive solution.

(P2) For any $\mu > 0$, take $\lambda(\mu)$ as the best (i.e. the smallest if $\varepsilon(p) > 0$ and the largest if $\varepsilon(p) < 0$) constant in the inequality

$$\|\nabla u\|_{L^2(\Omega)}^2 \geq \varepsilon(p) \left[\mu \|u\|_{L^{p+1}(\Omega)} - \lambda(\mu) \|u\|_{L^2(\Omega)} \right] \quad \forall u \in H^1(\Omega). \quad (1.4)$$

Here we denote by $\mu \mapsto \lambda(\mu)$ the inverse function of $\lambda \mapsto \mu(\lambda)$. Let

$$\mu_2 := \inf\{\lambda > 0 : \mu(\lambda) \neq \lambda \text{ in (1.4)}\}.$$

(P3) Let us consider the optimal *Keller-Lieb-Thirring inequality*

$$\nu(\mu) = -\varepsilon(p) \inf_{\phi \in \mathcal{A}_\mu} \lambda_1(\Omega, \varepsilon(p) \phi) \quad (1.5)$$

where the admissible set for the potential ϕ is defined by

$$\mathcal{A}_\mu := \left\{ \phi \in L_+^q(\Omega) : \|\phi^{\varepsilon(p)}\|_{L^q(\Omega)} = \mu \right\}$$

and $q = (p + 1)/|p - 1|$. Let

$$\mu_3 := \inf\{\mu > 0 : \nu(\mu) \neq \mu \text{ in (1.5)}\}.$$

Finally let us define Λ_\star as the *best constant in the interpolation inequality*

$$\|\nabla u\|_{L^2(\Omega)}^2 \geq \frac{\Lambda_\star}{p-1} \left[\|u\|_{L^{p+1}(\Omega)} - \|u\|_{L^2(\Omega)} \right] \quad \forall u \in H^1(\Omega). \quad (1.6)$$

THEOREM 1. — *Assume that $d \geq 2$, $p \in (0, 1) \cup (1, 2^* - 1)$ and Ω is bounded domain in \mathbb{R}^d with smooth boundary such that $|\Omega| = 1$. With the above notations, we have*

$$0 < \mu_1 \leq \mu_2 = \mu_3 = \frac{\Lambda_\star}{|p-1|}$$

and the following properties hold:

- (P1) *Rigidity holds in (1.3) for any $\lambda \in (0, \mu_1)$.*
- (P2) *The function $\mu \mapsto \lambda(\mu)$ is monotone increasing, concave if $p \in (0, 1)$, convex if $p \in (1, 2^* - 1)$ and $\lambda(\mu) = \mu$ if and only if $\mu \leq \mu_2$.*
- (P3) *For any $\mu > 0$, $\nu(\mu) = \lambda(\mu)$.*

This result is inspired from a series of recent papers on interpolation inequalities, rigidity results and Keller-Lieb-Thirring estimates on compact manifolds that will be discussed later. However, the absence of curvature and the presence of a boundary induce a number of changes compared to these papers, that we shall study next. Beyond the properties of Theorem 1 which are not very difficult to prove, our main goal is to get explicit estimates of μ_i and Λ_\star .

Let us define

$$\lambda_2 := \lambda_2(\Omega, 0)$$

which is the second (and first positive) eigenvalue of $-\Delta$ on Ω , with homogeneous Neumann boundary conditions. Recall that the lowest eigenvalue of $-\Delta$ is $\lambda_1 = 0$ and that the corresponding eigenspace is spanned by the constants. For this reason λ_2 is often called the *spectral gap* and the Poincaré inequality sometimes appears in the literature as the *spectral gap inequality*. Finally let us introduce the number

$$\theta_\star(p, d) = \frac{(d-1)^2 p}{d(d+2) + p}, \quad (1.7)$$

THEOREM 2. — *Assume that $d \geq 1$. With the above notations, we get the following estimates:*

$$\frac{1 - \theta_\star(p, d)}{|p-1|} \lambda_2 \leq \mu_1 \leq \mu_2 = \mu_3 = \frac{\Lambda_\star}{|p-1|} \leq \frac{1}{|p-1|} \lambda_2$$

for any $p \in (0, 1) \cup (1, 2^* - 1)$. The lower estimate holds only under the additional assumptions that Ω is convex and $d \geq 2$.

Before giving a brief overview of the literature related to our results, let us emphasize two points. We first notice that $\lim_{p \rightarrow (d+2)/(d-2)} \theta_*(p, d) = 1$ if $d \geq 3$ so that the lower estimate goes to 0 as the exponent p approaches the critical exponent. This is consistent with the previously known results on rigidity, that are based on Morrey's scheme and deteriorate as p approaches $(d+2)/(d-2)$. In the critical case, multiplicity may hold for any value of λ , so that one cannot expect that rigidity could hold without an additional assumption. The second remark is the fact that the convexity of Ω is essential for known results in the critical case and one should not be surprised to see this condition also in the sub-critical range. This assumption is however not required in the result of C.-S. Lin, W.-M. Ni, and I. Takagi in [35]. Compared to their paper, what we gain here is a fully explicit estimate which relies on a simple computation when $p > 1$. The case $p < 1$ has apparently not been studied yet.

It is remarkable that the case $p = 1$ is the endpoint of the two admissible intervals in p . We may notice that

$$\theta_*(1, d) = \frac{(d-1)^2}{(d+1)^2}$$

is in the interval $(0, 1)$ for any $d \geq 2$. The case $p = 1$ is a limit case, which corresponds to the logarithmic Sobolev inequality

$$\|\nabla u\|_{L^2(\Omega)}^2 - \frac{\Lambda_\star}{2} \int_{\Omega} |u|^2 \log \left(\frac{|u|^2}{\|u\|_{L^2(\Omega)}^2} \right) dx \quad \forall u \in H^1(\Omega) \quad (1.8)$$

where Λ_\star denotes the optimal constant, and by passing to the limit as $p \rightarrow 1$ in (1.6), we have the following result.

COROLLARY 3. — *If $d \geq 2$ and Ω is convex, then*

$$\frac{4d}{(d+1)^2} \lambda_2 \leq \Lambda_\star \leq \lambda_2.$$

It is also possible to define a family of logarithmic Sobolev inequalities depending on λ , or to get a parametrized Keller-Lieb-Thirring inequality and find that $\lambda = \Lambda_\star$ corresponds to a threshold value between a linear dependence of the optimal constant in λ and a regime in which this dependence is given by a strictly convex function of λ . The interested reader is invited to refer to [23, Corollaries 13 and 14] for similar results on the sphere.

The existence of a non-trivial solution to (1.1) bifurcating from the constant ones for $\lambda = \lambda_2/(p-1)$ has been established for instance in [36] when $p > 1$. This paper also contains the conjecture, known in the literature as the *Lin-Ni conjecture* and formulated in [36], that there are no nontrivial solutions for $\lambda > 0$ small enough and that there are non-trivial solutions for λ large enough, even in the super-critical case $p > 2^* - 1$. More details can be found in [35]. Partial results were obtained before in [44], when the exponent is in the range $1 < p < d/(d-2)$. These papers were originally motivated by the connection with the model of Keller and Segel in chemotaxis and the Gierer-Meinhardt system in pattern formation: see [42] for more explanations.

For completeness, let us briefly review what is known in the critical case $p = 2^* - 1$. When $d = 3$, it was proved by M. Zhu in [60] that rigidity holds true for $\lambda > 0$ small enough when one considers the positive solutions to the nonlinear elliptic equation

$$\Delta u - \lambda u + f(u) = 0$$

on a smooth bounded domain of \mathbb{R}^3 with homogeneous Neumann boundary conditions, if $f(u)$ is equal to u^5 up to a perturbation of lower order. Another proof was given by J. Wei and X. Xu in [58] and slightly extended later in [31]. The Lin-Ni conjecture is wrong for $p = 2^* - 1$ in higher dimensions: see [57], and also earlier references therein. Some of the results have been extended to the d -Laplacian in dimension d in [59].

Bifurcations, qualitative aspects of the branches of solutions and multiplicity have been widely studied in the case of homogeneous Dirichlet boundary conditions or in the whole space. For qualitative features, we may refer primarily to [37] by P.-L. Lions, to [53] by K. Schmitt and to a series of papers, [48, 47, 46, 45], by T. Ouyang and J. Shi. Further related references can be found for instance in [10, 38, 14, 39, 50, 43, 56, 7, 8, 49, 55, 11]. Much less is known in the case of homogeneous Neumann boundary conditions and we may for instance refer to [55, 52, 40, 54, 41] for some results in this direction, but only in rather simple cases (balls, intervals or rectangles).

Concerning the Lin-Ni conjecture, it is known from [35, Theorem 3, (ii)] that $u \equiv \lambda^{1/(p-1)}$ if λ is small enough (also see [44] for an earlier partial result), and that there is a non-trivial solution if λ is large enough. As already said above, the method is based on the Moser iteration technique, in order to get a uniform estimate on the solution, and then on a direct estimate based on the Poincaré inequality. In the proof of Theorem 2 we shall adopt a completely different strategy, which is inspired by the *rigidity* results for nonlinear elliptic PDEs as in [28, 9] on the one hand, and by the

carré du champ method of D. Bakry and M. Emery on the other hand, that can be traced back to [3]. More precisely, we shall rely on improved versions of these methods as in [2, 34], which involve the eigenvalues of the Laplacian, results on interpolation inequalities on compact manifolds obtained by J. Demange in [16], and a recent improvement with a computation based on traceless Hessians in [26].

From a larger perspective, our approach in Theorem 2 is based on results for compact Riemannian manifolds that can be found in various papers: the most important ones are the rigidity results of L. Véron et al. in [9, 33], the computations inspired by the *carré du champ* method of [2, 16], and the nonlinear flow approach of [26] (also see [22, 21, 25]). Using these estimates in the range $1 < p < 2^* - 1$ and the Bakry-Emery method as in [17] in the case $p \in (0, 1)$, our goal is to prove that rigidity holds in a certain range of λ without relying on uniform estimates (and the Moser iteration technique) and discuss the estimates of the threshold values. The spectral estimates of Theorem 1 are directly inspired by [23, 24].

This paper is organized as follows. Preliminary results have been collected in Section 2. The proof of Theorem 1 is given in Section 3. In Section 4, we use the heat flow to establish a first lower bound like the one of Theorem 2. A better bound that can be interpreted in terms of a nonlinear flow is given in Section 5. The last section is devoted to some results on the branches of radial solutions, various considerations on flows and, in particular, improvements based on the nonlinear flow method.

Notations

If $A = (A_{ij})_{i,j=1}^d$ and $B = (B_{ij})_{i,j=1}^d$ are two matrices, let $A : B = \sum_{i,j=1}^d A_{ij} B_{ij}$ and $|A|^2 = A : A$. If a and b take values in \mathbb{R}^d , we adopt the definitions:

$$a \cdot b = \sum_{i=1}^d a_i b_i, \quad \nabla \cdot a = \sum_{i=1}^d \frac{\partial a_i}{\partial x_i},$$

$$a \otimes b = (a_i b_j)_{i,j=1}^d, \quad \nabla \otimes a = \left(\frac{\partial a_j}{\partial x_i} \right)_{i,j=1}^d.$$

2. Preliminary results

Let us recall that $\Omega \subset \mathbb{R}^d$, $d \geq 2$ is a bounded domain with smooth boundary (or an open interval if $d = 1$) and let λ_2 be the first non-zero

eigenvalue of the Laplace operator on Ω , supplemented with homogeneous Neumann boundary conditions. We shall denote by n a unit outgoing normal vector of $\partial\Omega$ and consider a non-trivial eigenfunction $u_2 \in H^1(\Omega)$ associated with the lowest positive eigenvalue λ_2 , so that

$$-\Delta u_2 = \lambda_2 u_2 \quad \text{in } \Omega, \quad \partial_n u_2 = 0 \quad \text{on } \partial\Omega. \quad (2.1)$$

As a trivial observation we may observe that u_2 is in $H^2(\Omega)$.

LEMMA 4. — *With the above notations, for any $u \in H^2(\Omega)$ such that $\partial_n u = 0$ on $\partial\Omega$, we have*

$$\left(\int_{\Omega} |\nabla u|^2 dx \right)^2 \leq \int_{\Omega} |\Delta u|^2 dx \int_{\Omega} |u|^2 dx,$$

and equality holds if $u = u_2$. As a consequence, we also have

$$\lambda_2 \int_{\Omega} |\nabla u|^2 dx \leq \int_{\Omega} |\Delta u|^2 dx. \quad (2.2)$$

Proof. — By expanding the square and integrating by parts the cross term, we notice that

$$\begin{aligned} 0 &\leq \int_{\Omega} \left| \frac{1}{\sqrt{\mu}} \Delta u + \lambda_2 \sqrt{\mu} u \right|^2 dx \\ &= \frac{1}{\mu} \|\Delta u\|_{L^2(\Omega)}^2 - 2\lambda_2 \|\nabla u\|_{L^2(\Omega)}^2 + \lambda_2^2 \mu \|u\|_{L^2(\Omega)}^2, \end{aligned}$$

where μ is an arbitrary positive real parameter. After optimizing on $\mu > 0$, we arrive at

$$0 \leq 2\lambda_2 \left(\|\Delta u\|_{L^2(\Omega)} \|u\|_{L^2(\Omega)} - \|\nabla u\|_{L^2(\Omega)}^2 \right).$$

To check the equality case with $u = u_2$, it is enough to multiply (2.1) by u_2 and by $-\Delta u_2$, and then integrate by parts. By definition of λ_2 , we know that

$$\int_{\Omega} |\nabla u|^2 dx \geq \lambda_2 \int_{\Omega} |u|^2 dx$$

with equality again if $u = u_2$, which concludes the proof. \square

A key result for this paper is based on the computation of $(\Delta u)^2$ in terms of the Hessian matrix of u , which involves integrations by parts and boundary terms. The following result can be found in [27, Lemma 5.1] or [29].

LEMMA 5. — *If Ω is a smooth convex domain in \mathbb{R}^d and if $u \in C^3(\overline{\Omega})$ such that $\partial_n u = 0$ on $\partial\Omega$, then*

$$- \sum_{i,j=1}^d \int_{\partial\Omega} \partial_{i_j}^2 u \partial_i u n_j d\mathcal{H}^{d-1} \geq 0.$$

As a consequence, if $u \in W^{2,2}(\Omega)$ is such that $\partial_n u = 0$, then we have that

$$\int_{\Omega} |\Delta u|^2 dx \geq \int_{\Omega} |\text{Hess } u|^2 dx.$$

Consider on $H^1(\Omega)$ the functional

$$\mathcal{J}_{\Lambda}[u] := \|\nabla u\|_{L^2(\Omega)}^2 - \frac{\Lambda}{p-1} \left[\|u\|_{L^{p+1}(\Omega)}^2 - \|u\|_{L^2(\Omega)}^2 \right]$$

if $p \neq 1$, and

$$\mathcal{J}_{\Lambda}[u] := \|\nabla u\|_{L^2(\Omega)}^2 - \frac{\Lambda}{2} \int_{\Omega} |u|^2 \log \left(\frac{|u|^2}{\|u\|_{L^2(\Omega)}^2} \right) dx$$

if $p = 1$.

LEMMA 6. — *There exists a function $u \in H^1(\Omega)$ such that $\partial_n u = 0$ and $\mathcal{J}_{\Lambda}[u] < 0$ if $\Lambda > \lambda_2$.*

Proof. — A simple computations shows that

$$\mathcal{J}_{\Lambda}[1 + \epsilon w] \sim \epsilon^2 \left[\|\nabla w\|_{L^2(\Omega)}^2 - \Lambda \|w\|_{L^2(\Omega)}^2 \right]$$

as $\epsilon \rightarrow 0$. By choosing w to be an eigenfunction associated with λ_2 , we get that

$$\mathcal{J}_{\Lambda}[1 + \epsilon w] \sim \epsilon^2 (\lambda_2 - \Lambda) \|w\|_{L^2(\Omega)}^2$$

is negative for $\epsilon > 0$ small enough if $\Lambda > \lambda_2$. \square

Lemma 6 provides the upper bound in Theorem 2. It proves indeed that

$$\Lambda_{\star} \leq \lambda_2.$$

This method has been widely exploited and a similar argument can be found for instance in [44].

3. Proof of Theorem 1

Assume that $p > 1$ and let us recall that

$$\mu(\lambda) := \inf_{u \in H^1(\Omega) \setminus \{0\}} \mathcal{Q}_{\lambda}[u] \quad \text{with} \quad \mathcal{Q}_{\lambda}[u] := \frac{\|\nabla u\|_{L^2(\Omega)}^2 + \lambda \|u\|_{L^2(\Omega)}^2}{\|u\|_{L^{p+1}(\Omega)}^2}.$$

We denote by $\kappa_{p,d}$ the optimal constant in the following Gagliardo-Nirenberg inequality on \mathbb{R}^d :

$$\|\nabla v\|_{L^2(\mathbb{R}^d)}^2 + \|v\|_{L^2(\mathbb{R}^d)}^2 \geq \kappa_{p,d} \|v\|_{L^{p+1}(\mathbb{R}^d)}^2 \quad \forall v \in H^1(\mathbb{R}^d).$$

LEMMA 7. — *If $p \in (1, 2^* - 1)$, the function $\lambda \mapsto \mu(\lambda)$ is monotone increasing, concave, such that $\mu(\lambda) \leq \lambda$ for any $\lambda > 0$, and $\mu(\lambda) = \lambda$ if and only if $0 < \lambda < \mu_2 = \Lambda_*/(p-1)$. Moreover, we have*

$$\mu(\lambda) \sim 2^{\frac{1-p}{1+p}} \kappa_{p,d} \lambda^{1-\frac{d}{2}} \frac{p-1}{p+1} \quad \text{as } \lambda \rightarrow +\infty.$$

Proof. — For any given $u \in H^1(\Omega) \setminus \{0\}$, $\lambda \mapsto \mathcal{Q}_\lambda[u]$ is affine, increasing. By taking an infimum, we know that μ is concave, non-decreasing. Using $u \equiv 1$ as a test function, we know that $\mu(\lambda) \leq \lambda$ for any $\lambda > 0$. By standard variational methods, we know that there is an optimal function $u \in H^1(\Omega) \setminus \{0\}$, so that

$$\|\nabla u\|_{L^2(\Omega)}^2 + \lambda \|u\|_{L^2(\Omega)}^2 = \mu(\lambda) \|u\|_{L^{p+1}(\Omega)}^2.$$

On the other hand, we know from (1.6) that

$$\|\nabla u\|_{L^2(\Omega)}^2 + \frac{\Lambda_*}{p-1} \|u\|_{L^2(\Omega)}^2 \geq \frac{\Lambda_*}{p-1} \|u\|_{L^{p+1}(\Omega)}^2.$$

Hence we have the inequality

$$\left(1 - \frac{\lambda}{\Lambda_*} (p-1)\right) \|\nabla u\|_{L^2(\Omega)}^2 \leq (\mu(\lambda) - \lambda) \|u\|_{L^{p+1}(\Omega)}^2.$$

If $\lambda < \Lambda_*/(p-1)$, the l.h.s. is nonnegative while the r.h.s. is nonpositive because $\mu(\lambda) \leq \lambda$, so that we conclude at once that $\mu(\lambda) = \lambda$ and u is constant. As a consequence, $\mu_2 \geq \Lambda_*/(p-1)$. On the other hand, by definition of Λ_* , we know that $\mu_2 \leq \Lambda_*/(p-1)$.

The regime as $\lambda \rightarrow \infty$ is easily studied by a rescaling. If u_λ denotes an optimal function such that $\mathcal{Q}_\lambda[u_\lambda] = \mu(\lambda)$, then $v_\lambda(x) := u_\lambda(x/\sqrt{\lambda})$ is an optimal function for

$$\|\nabla v\|_{L^2(\Omega_\lambda)}^2 + \|v\|_{L^2(\Omega_\lambda)}^2 \geq \frac{\mu(\lambda)}{\lambda^{1-\frac{d}{2}} \frac{p-1}{p+1}} \|v\|_{L^{p+1}(\Omega_\lambda)}^2 \quad \forall v \in H^1(\Omega_\lambda)$$

where $\Omega_\lambda := \{x \in \mathbb{R}^d : \lambda^{-1}x \in \Omega\}$. Using truncations of the optimal functions for the Gagliardo-Nirenberg inequality on $\mathbb{R}_+^d = \{(x_1, x_2, \dots, x_d) \in \mathbb{R}^d : x_d > 0\}$ and an analysis of the convergence of an extension of v_λ in $H^1(\mathbb{R}^d)$ as $\lambda \rightarrow \infty$ based on standard concentration-compactness methods, up to the extraction of subsequences and translations, we get that the limit function v is optimal for the inequality

$$\|\nabla v\|_{L^2(\mathbb{R}_+^d)}^2 + \|v\|_{L^2(\mathbb{R}_+^d)}^2 \geq 2^{\frac{1-p}{1+p}} \kappa_{p,d} \|v\|_{L^{p+1}(\mathbb{R}_+^d)}^2 \quad \forall v \in H^1(\mathbb{R}_+^d).$$

See [23, Lemma 5] for more details in a similar case.

As a consequence of the behavior at infinity and of the concavity property, μ is a monotone increasing function of λ . \square

Assume that $p < 1$ and let us recall that

$$\lambda(\mu) := \inf_{u \in H^1(\Omega) \setminus \{0\}} \mathcal{Q}^\mu[u] \quad \text{with} \quad \mathcal{Q}^\mu[u] := \frac{\|\nabla u\|_{L^2(\Omega)}^2 + \mu \|u\|_{L^{p+1}(\Omega)}^2}{\|u\|_{L^2(\Omega)}^2}.$$

We denote by $\kappa_{p,d}^+$ the optimal constant in the following Gagliardo-Nirenberg inequality on \mathbb{R}_+^d :

$$\|\nabla v\|_{L^2(\mathbb{R}_+^d)}^2 + \|v\|_{L^{p+1}(\mathbb{R}_+^d)}^2 \geq \kappa_{p,d}^+ \|v\|_{L^2(\mathbb{R}_+^d)}^2 \quad \forall v \in H^1(\mathbb{R}_+^d).$$

LEMMA 8. — *If $p \in (0, 1)$, the function $\mu \mapsto \lambda(\mu)$ is monotone increasing, concave, such that $\lambda(\mu) \leq \mu$ for any $\mu > 0$, and $\lambda(\mu) = \mu$ if and only if $0 < \mu < \mu_2 = \Lambda_*/(1-p)$. Moreover, we have*

$$\lambda(\mu) \sim \kappa_{p,d}^+ \mu^{\frac{1}{1+\frac{d}{2}\frac{1-p}{p+1}}} \quad \text{as } \lambda \rightarrow +\infty.$$

Proof. — The proof follows the same strategy as in the proof of Lemma 7. See [23, Lemma 11] for more details in a similar case. \square

Recall that we denote by $\mu \mapsto \lambda(\mu)$ the inverse function of $\lambda \mapsto \mu(\lambda)$ and get in both cases, $p > 1$ and $p < 1$, the fact that

$$\mu(\lambda) = O\left(\lambda^{1-\frac{d}{2}\frac{p-1}{p+1}}\right) \quad \text{as } \lambda \rightarrow +\infty.$$

LEMMA 9. — *With the notations of Theorem 2, we have $\nu(\mu) = \lambda(\mu)$ for any $\mu > 0$ and, as a consequence, $\mu_2 = \mu_3$.*

Proof. — Assume first that $p > 1$. The proof is based on two ways of estimating the quantity

$$\mathcal{A} = \|\nabla u\|_{L^2(\Omega)}^2 + \lambda \|u\|_{L^2(\Omega)}^2 - \int_{\Omega} \phi |u|^2 dx.$$

On the one hand we may use Hölder's inequality to estimate

$$\int_{\Omega} \phi |u|^2 dx \leq \|\phi\|_{L^q(\Omega)} \|u\|_{L^{p+1}(\Omega)}^2$$

with $q = \frac{p+1}{p-1}$ and get

$$\mathcal{A} \geq \|\nabla u\|_{L^2(\Omega)}^2 + \lambda \|u\|_{L^2(\Omega)}^2 - \mu \|u\|_{L^{p+1}(\Omega)}^2$$

with $\mu = \|\phi\|_{L^q(\Omega)}$. Using $u \equiv 1$ as a test function, we observe that the lowest eigenvalue $\lambda_1(\Omega, -\phi)$ of the Schrödinger operator $-\Delta - \phi$ is nonpositive. With $u = u_1$ an eigenfunction associated with $\lambda_1(\Omega, -\phi)$, we know

that

$$\mathcal{A} = \lambda - |\lambda_1(\Omega, -\phi)| \geq \|\nabla u\|_{L^2(\Omega)}^2 + \lambda \|u\|_{L^2(\Omega)}^2 - \mu \|u\|_{L^{p+1}(\Omega)}^2$$

is nonnegative if $\lambda = \lambda(\mu)$, thus proving that $\lambda(\mu) - |\lambda_1(\Omega, -\phi)| \geq 0$ and hence

$$\lambda(\mu) \geq \nu(\mu).$$

On the other hand, with $\phi = \mu u^{p-1} / \|u\|_{L^{p+1}(\Omega)}^{p-1}$, we observe that

$$\begin{aligned} \mathcal{A} &= \|\nabla u\|_{L^2(\Omega)}^2 + \lambda \|u\|_{L^2(\Omega)}^2 - \mu \|u\|_{L^{p+1}(\Omega)}^2 \\ &\geq (\lambda - |\lambda_1(\Omega, -\phi)|) \|u\|_{L^2(\Omega)}^2 \geq (\lambda - \nu(\mu)) \|u\|_{L^2(\Omega)}^2 \end{aligned}$$

and \mathcal{A} is nonpositive if we take $\mu = \mu(\lambda)$ and u the corresponding optimal function. This proves that

$$\lambda(\mu) \leq \nu(\mu),$$

which concludes the proof when $p > 1$.

A similar computation can be done if $p < 1$, based on the Hölder inequality

$$\int_{\Omega} u^{p+1} dx \leq \int_{\Omega} u^{p+1} \phi^{\frac{p+1}{2}} \phi^{-\frac{p+1}{2}} dx \leq \left(\int_{\Omega} |u|^2 \phi dx \right)^{\frac{p+1}{2}} \|\phi^{-1}\|_{L^q(\Omega)}^{\frac{p+1}{2}}$$

with $q = \frac{1+p}{1-p}$, that is

$$\int_{\Omega} |u|^2 \phi dx \geq \mu \|u\|_{L^{p+1}(\Omega)}^2$$

with $\mu^{-1} = \|\phi^{-1}\|_{L^q(\Omega)}$. With

$$\mathcal{A} = \|\nabla u\|_{L^2(\Omega)}^2 - \lambda \|u\|_{L^2(\Omega)}^2 + \int_{\Omega} \phi |u|^2 dx,$$

the computation is parallel to the one of the case $p > 1$. Also see [23] for similar estimates. \square

LEMMA 10. — *With the notations of Theorem 2, we have $\mu_1 \leq \mu_2$.*

Proof. — Let u be an optimal function for (1.4). It can be taken nonnegative without restriction and solves the Euler-Lagrange equation

$$-\varepsilon(p) \Delta u + \lambda u - \mu \frac{u^p}{\|u\|_{L^{p+1}(\Omega)}^{p-1}} = 0 \quad \text{in } \Omega, \quad \partial_n u = 0 \quad \text{on } \partial\Omega, \quad (3.1)$$

where $\lambda = \lambda(\mu)$ or equivalently $\mu = \mu(\lambda)$. By homogeneity, we can fix $\|u\|_{L^{p+1}(\Omega)}$ as we wish and may choose $\|u\|_{L^{p+1}(\Omega)}^{p-1} = \mu$, hence concluding that u is constant if $\mu \leq \mu_1$ and, as a consequence, $\lambda(\mu) = \mu$, thus proving that $\mu \leq \mu_2$. The conclusion follows. \square

4. Estimates based on the heat equation

We use the Bakry-Emery method to prove some results that are slightly weaker than the assertion of Theorem 2 but the method is of its own independent interest. Except for the precise value of the constant, this following result can be found in [17] (also see earlier references therein).

LEMMA 11. — *Assume that Ω is convex and such that $|\Omega| = 1$. For any $p \in (0, 1)$, for any $u \in H^1(\Omega)$ such that $\partial_n u = 0$ on $\partial\Omega$, we have*

$$\|\nabla u\|_{L^2(\Omega)}^2 \geq \lambda_2 \left[\|u\|_{L^2(\Omega)}^2 - \|u\|_{L^{p+1}(\Omega)}^2 \right].$$

Proof. — We give a proof based on the entropy – entropy production method. It is enough to prove the result for nonnegative functions u since the inequality for $|u|$ implies the inequality for u . By density, we may assume that u is smooth. According to [17], if v is a nonnegative solution of the heat equation

$$\frac{\partial v}{\partial t} = \Delta v$$

on Ω with homogeneous Neumann boundary conditions, then $v = u^{p+1}$ is such that

$$\frac{d}{dt} \int_{\Omega} \frac{v^r - M^r}{r-1} dx = -\frac{4}{r} \int_{\Omega} |\nabla u|^2 dx$$

with $M := \int_{\Omega} v dx$ and $r = 2/(p+1)$. With this change of variables, u solves

$$\frac{\partial u}{\partial t} = \Delta u + \frac{2-r}{r} \frac{|\nabla u|^2}{u}$$

and we find that

$$\begin{aligned} & -\frac{1}{2} \frac{d}{dt} \int_{\Omega} |\nabla u|^2 dx \\ &= \int_{\Omega} |\Delta u|^2 dx + p \int_{\Omega} \frac{|\nabla u|^4}{u^2} dx - 2p \int_{\Omega} \text{Hess } u : \frac{\nabla u \otimes \nabla u}{u} dx, \end{aligned} \quad (4.1)$$

that is,

$$\begin{aligned} & -\frac{1}{2} \frac{d}{dt} \int_{\Omega} |\nabla u|^2 dx = 2 \frac{r-1}{r} \int_{\Omega} |\Delta u|^2 dx + p \int_{\Omega} (|\Delta u|^2 - |\text{Hess } u|^2) dx \\ & \quad + p \int_{\Omega} \left| \text{Hess } u - \frac{1}{u} \nabla u \otimes \nabla u \right|^2 dx, \end{aligned}$$

and finally, using Lemma 5,

$$\frac{d}{dt} \int_{\Omega} |\nabla u|^2 dx \leq -4 \frac{r-1}{r} \int_{\Omega} |\Delta u|^2 dx \leq -4 \frac{r-1}{r} \lambda_2 \int_{\Omega} |\nabla u|^2 dx$$

where the last inequality follows from Lemma 4, thus proving the result for any $p = (2 - r)/r \in (0, 1)$. Indeed, with previous notations, we have shown that $\int_{\Omega} |\nabla u|^2 dx$ is exponentially decaying. Hence

$$\int_{\Omega} \frac{v^r - M^r}{r - 1} dx = \frac{1}{r - 1} \left(\|u\|_{L^2(\Omega)}^2 - \|u\|_{L^{p+1}(\Omega)}^2 \right)$$

also converges to 0 as $t \rightarrow \infty$ and

$$\frac{d}{dt} \left[\|\nabla u\|_{L^2(\Omega)}^2 - \mu \int_{\Omega} \frac{v^r - M^r}{r - 1} dx \right] \leq \left(-4 \frac{r-1}{r} \lambda_2 + \frac{4}{r} \mu \right) \int_{\Omega} |\nabla u|^2 dx$$

is nonpositive if $\mu \leq (r - 1) \lambda_2$. Altogether, we have shown that

$$\|\nabla u\|_{L^2(\Omega)}^2 - \lambda_2 \left(\|u\|_{L^2(\Omega)}^2 - \|u\|_{L^{p+1}(\Omega)}^2 \right)$$

is nonincreasing with limit 0, which concludes the proof. \square

A better result can be obtained by considering the traceless quantities as in [26]. Let us introduce

$$\begin{aligned} M u &:= \frac{\nabla u \otimes \nabla u}{u} - \frac{1}{d} \frac{|\nabla u|^2}{u} \text{Id}, \\ L u &:= \text{Hess } u - \frac{1}{d} \Delta u \text{Id}, \end{aligned}$$

and define $p^\sharp := \frac{d(d+2)}{(d-1)^2}$ so that

$$\vartheta(p, d) := p \frac{(d-1)^2}{d(d+2)}$$

satisfies $\vartheta(p, d) < 1$ for any $p \in (0, p^\sharp)$. Notice that $p^\sharp + 1 = \frac{2d^2+1}{(d-1)^2}$ is the threshold value that has been found in [4] (also see [21, 22]).

LEMMA 12. — *Assume that Ω is convex and such that $|\Omega| = 1$. For any $p \in (0, p^\sharp)$, for any $u \in H^1(\Omega)$ such that $\partial_n u = 0$ on $\partial\Omega$, we have*

$$\|\nabla u\|_{L^2(\Omega)}^2 \geq \frac{1}{2} (1 - \vartheta(p, d)) \lambda_2 \frac{\|u\|_{L^{p+1}(\Omega)}^2 - \|u\|_{L^2(\Omega)}^2}{p - 1}$$

if $p \neq 1$ and, in the limit case $p = 1$,

$$\|\nabla u\|_{L^2(\Omega)}^2 \geq \frac{1}{4} (1 - \vartheta(1, d)) \lambda_2 \int_{\Omega} |u|^2 \log \left(\frac{|u|^2}{\|u\|_{L^2(\Omega)}^2} \right) dx.$$

Proof. — We use the same conventions as in the proof of Lemma 11. Let us first observe that

$$\begin{aligned} |\mathbf{M}u|^2 &= \left(1 - \frac{1}{d}\right) \frac{|\nabla u|^4}{u^2}, \\ |\mathbf{L}u|^2 &= |\text{Hess } u|^2 - \frac{1}{d} (\Delta u)^2. \end{aligned}$$

Since $\text{Hess } u = \mathbf{L}u + \frac{1}{d} \Delta u \text{Id}$, we have that

$$\text{Hess } u : \frac{\nabla u \otimes \nabla u}{u} = \mathbf{L}u : \frac{\nabla u \otimes \nabla u}{u} + \frac{1}{d} \Delta u \frac{|\nabla u|^2}{u} = \mathbf{L}u : \mathbf{M}u + \frac{1}{d} \Delta u \frac{|\nabla u|^2}{u}$$

because $\mathbf{L}u$ is traceless. An integration by parts shows that

$$\begin{aligned} \int_{\Omega} \Delta u \frac{|\nabla u|^2}{u} dx &= \frac{|\nabla u|^4}{u^2} - 2 \text{Hess } u : \frac{\nabla u \otimes \nabla u}{u} \\ &= \frac{d}{d-1} |\mathbf{M}u|^2 - 2 \text{Hess } u : \frac{\nabla u \otimes \nabla u}{u} \end{aligned}$$

so that we get

$$\frac{d+2}{d} \text{Hess } u : \frac{\nabla u \otimes \nabla u}{u} = \mathbf{L}u : \mathbf{M}u + \frac{1}{d-1} |\mathbf{M}u|^2,$$

hence

$$\text{Hess } u : \frac{\nabla u \otimes \nabla u}{u} = \frac{d}{d+2} \mathbf{L}u : \mathbf{M}u + \frac{d}{(d-1)(d+2)} |\mathbf{M}u|^2.$$

Now let us come back to the proof of Lemma 11. From (4.1), we read that

$$\begin{aligned} & -\frac{1}{2} \frac{d}{dt} \int_{\Omega} |\nabla u|^2 dx \\ &= \int_{\Omega} |\Delta u|^2 dx + p \int_{\Omega} \frac{|\nabla u|^4}{u^2} dx - 2p \int_{\Omega} \text{Hess } u : \frac{\nabla u \otimes \nabla u}{u} dx \\ &= \int_{\Omega} |\Delta u|^2 dx + \frac{pd}{d-1} \int_{\Omega} |\mathbf{M}u|^2 dx \\ &\quad - 2p \left(\frac{d}{d+2} \int_{\Omega} \mathbf{L}u : \mathbf{M}u dx + \frac{d}{(d-1)(d+2)} \int_{\Omega} |\mathbf{M}u|^2 dx \right) \\ &= \int_{\Omega} |\Delta u|^2 dx - p \frac{d-1}{d+2} \int_{\Omega} |\mathbf{L}u|^2 dx \\ &\quad + \frac{pd^2}{(d-1)(d+2)} \int_{\Omega} \left| \mathbf{M}u - \frac{d-1}{d} \mathbf{L}u \right|^2 dx \end{aligned}$$

We know from Lemma 5 that

$$\int_{\Omega} (\Delta u)^2 dx \geq \int_{\Omega} |\text{Hess } u|^2 dx = \int_{\Omega} |\mathbf{L}u|^2 dx + \frac{1}{d} \int_{\Omega} (\Delta u)^2 dx,$$

i.e.,

$$\int_{\Omega} (\Delta u)^2 dx \geq \frac{d}{d-1} \int_{\Omega} |\mathbb{L}u|^2 dx.$$

Altogether, this proves that, for any $\theta \in (0, 1)$,

$$-\frac{1}{2} \frac{d}{dt} \int_{\Omega} |\nabla u|^2 dx \geq (1-\theta) \int_{\Omega} |\Delta u|^2 dx + \left(\frac{\theta d}{d-1} - p \frac{d-1}{d+2} \right) \int_{\Omega} |\mathbb{L}u|^2 dx$$

and finally, with $\theta = \vartheta(p, d)$ and using (2.2),

$$\frac{d}{dt} \left[\|\nabla u\|_{L^2(\Omega)}^2 - \mu \int_{\Omega} \frac{v^r - M^r}{r-1} dx \right] \leq \left(-(1-\theta) \lambda_2 + \frac{4}{r} \mu \right) \int_{\Omega} |\nabla u|^2 dx$$

is nonpositive if

$$\mu \leq \frac{r}{4} (1 - \vartheta(p, d)) \lambda_2 = \frac{1 - \vartheta(p, d)}{2(p+1)} \lambda_2.$$

Since $r-1 = (1-p)/(1+p)$, this concludes the proof if $p \neq 1$. The case $p = 1$ is obtained by passing to the limit as $p \rightarrow 1$. \square

5. Estimates based on nonlinear diffusion equations

LEMMA 13. — *With the notations of Theorem 2, we have*

$$\frac{1 - \theta_*(p, d)}{|p-1|} \lambda_2 \leq \mu_1.$$

Proof. — This bound is inspired from [2, 34, 33, 26]. Let us give the main steps of the proof. Here we do it at the level of the nonlinear elliptic PDE. Flows will be introduced afterwards, with the intent of providing improvements.

Let us consider the solution u to (1.3) and define a function v such that $v^\beta = u$ for some exponent β to be chosen later. Then v solves

$$-\varepsilon(p) \left(\Delta v + (\beta-1) \frac{|\nabla v|^2}{v} \right) + \lambda v - v^\kappa = 0 \quad \text{in } \Omega, \quad \partial_n u = 0 \quad \text{on } \partial\Omega$$

where

$$\kappa = \beta(p-1) + 1. \tag{5.1}$$

If we multiply the equation by $(\Delta v + \kappa |\nabla v|^2/v)$ and integrate by parts, then the nonlinear term disappears and we are left with the identity

$$\begin{aligned} \int_{\Omega} (\Delta v)^2 dx + (\kappa + \beta - 1) \int_{\Omega} \Delta v \frac{|\nabla v|^2}{v} dx + \kappa(\beta - 1) \int_{\Omega} \frac{|\nabla v|^4}{v^2} dx \\ - \lambda |p-1| \int_{\Omega} |\nabla v|^2 dx = 0. \end{aligned}$$

Let us define

$$Qv := \text{Hess } v - \frac{1}{d} \Delta v \text{Id} - \frac{(d-1)(p-1)}{\theta(d+3-p)} \left[\frac{\nabla v \otimes \nabla v}{v} - \frac{1}{d} \frac{|\nabla v|^2}{v} \text{Id} \right].$$

Exactly the same computation as in [26], based on the identities that can be found in the proof of Lemma 13, shows that

$$\begin{aligned} \theta \left(\int_{\Omega} (\Delta v)^2 dx - \int_{\Omega} |\text{Hess } v|^2 dx \right) + \frac{\theta d}{d-1} \int_{\Omega} |Qv|^2 dx \\ + (1-\theta) \int_{\Omega} (\Delta v)^2 dx - \lambda |p-1| \int_{\Omega} |\nabla v|^2 dx = 0 \end{aligned}$$

if $\theta = \theta_*(p, d) = \frac{(d-1)^2 p}{d(d+2)+p}$ and $\beta = \frac{d+2}{d+2-p}$. In the previous identity, the first term is nonnegative by Lemma 5, the second term is the integral of a square and is therefore nonnegative, and the last one is positive according to Lemma 4 if $(1-\theta) \int_{\Omega} (\Delta v)^2 dx - \lambda |p-1| \int_{\Omega} |\nabla v|^2 dx > 0$, unless $\nabla v = 0$ a.e.

We may notice that $p = d+2$ has to be excluded in order to define β , and this may occur if $d = 2$. However, by working directly on u , it is possible to cover this case as well. \square

Since the exponent p is in the sub-critical range, it is classical that the functional \mathcal{J}_{Λ} has a minimizer, with $\Lambda = \lambda |p-1|$, which therefore satisfies the Euler-Lagrange equations. On the other hand, any optimal function with $|p-1|\lambda < \Lambda_*$ is a constant by Lemma 13, and we are therefore in the case $\lambda = \mu$ of Theorem 1 if $\lambda \leq \frac{1-\theta_*(p,d)}{|p-1|} \lambda_2$. This completes the proof of Theorem 2.

For later purpose, let us consider the proof based on the flow. With $\lambda = \Lambda/|p-1|$, we may consider the functional

$$\mathcal{J}_{\Lambda}[u] = \|\nabla u\|_{L^2(\Omega)}^2 - \frac{\Lambda}{p-1} \left[\|u\|_{L^{p+1}(\Omega)}^2 - \|u\|_{L^2(\Omega)}^2 \right]$$

with $u = v^{\beta}$ and evolve it according to

$$\frac{\partial v}{\partial t} = v^{2-2\beta} \left(\Delta v + \kappa \frac{|\nabla v|^2}{v} \right). \quad (5.2)$$

This flow has the nice property that

$$\frac{d}{dt} \int_{\Omega} u^p dx = \frac{d}{dt} \int_{\Omega} v^{\beta p} dx = 0$$

if κ is given by (5.1), and a simple computation shows that

$$\begin{aligned} -\frac{1}{\beta^2} \frac{d}{dt} \mathcal{J}_\Lambda[v^\beta] &= \int_\Omega (\Delta v)^2 dx + (\kappa + \beta - 1) \int_\Omega \Delta v \frac{|\nabla v|^2}{v} dx \\ &\quad + \kappa(\beta - 1) \int_\Omega \frac{|\nabla v|^4}{v^2} dx - \lambda|p - 1| \int_\Omega |\nabla v|^2 dx = 0. \end{aligned}$$

The same choices of β and θ as in the proof of Lemma 13 allow us to conclude, but it is interesting to discuss the possible values of β and θ which guarantee that $\frac{d}{dt} \mathcal{J}_\Lambda[v^\beta] \leq 0$ unless v is a constant. As in [26], elementary computations show that

$$\begin{aligned} -\frac{1}{\beta^2} \frac{d}{dt} \mathcal{J}_\Lambda[v^\beta] &= \theta \left(\int_\Omega (\Delta v)^2 dx - \int_\Omega |\text{Hess } v|^2 dx \right) + \frac{\theta d}{d-1} \int_\Omega |\mathbb{Q}v|^2 dx \\ &\quad + \mathfrak{R} \int_\Omega \frac{|\nabla v|^4}{v^2} dx + (1-\theta) \int_\Omega (\Delta v)^2 dx - \lambda|p-1| \int_\Omega |\nabla v|^2 dx \end{aligned}$$

where $\mathbb{Q}u$ is now defined by

$$\mathbb{Q}u := \text{L}u - \frac{1}{\theta} \frac{d-1}{d+2} (\kappa + \beta - 1) \left[\frac{\nabla u \otimes \nabla u}{u} - \frac{1}{d} \frac{|\nabla u|^2}{u} \text{Id} \right]$$

and

$$\mathfrak{R} := -\frac{1}{\theta} \left(\frac{d-1}{d+2} \right)^2 (\kappa + \beta - 1)^2 + \kappa(\beta - 1) + (\kappa + \beta - 1) \frac{d}{d+2}.$$

After replacing κ by its value according to (5.1), we obtain that

$$\mathfrak{R} = \left[\left(\frac{d-1}{d+2} \right)^2 \frac{p^2}{\theta} - p + 1 \right] \beta^2 - 2 \left(1 - \frac{p}{d+1} \right) \beta + 1$$

has two roots $\beta_\pm(\theta, p, d)$ if $\theta \in (\theta_\star(p, d), 1)$ and $\mathfrak{R} > 0$ if $\beta \in (\beta_-, \beta_+)$.

PROPOSITION 14. — *With the above notations, for any $\theta \in (\theta_\star(p, d), 1)$ and any $\beta \in (\beta_-(\theta, p, d), \beta_+(\theta, p, d))$, we have*

$$\frac{d}{dt} \mathcal{J}_\Lambda[v^\beta] \leq -\mathfrak{R} \beta^2 \int_\Omega \frac{|\nabla v|^4}{v^2} dx$$

if v is a solution to (5.2).

The computations in the proof of Lemma 13 can now be reinterpreted in the framework of the flow defined by (5.2). When $\theta = \theta_\star(p, d)$, the reader is invited to check that $\beta_- = \beta_+ = \beta$. A solution to (1.3) is stationary with respect to (5.2) and then all computations amount to write that $\mathcal{J}_\Lambda[v^\beta] = 0$ is possible only if v is a constant.

6. Further considerations

Let us conclude this paper by a series of remarks.

6.1. The threshold case

PROPOSITION 15. — *With $\theta_*(p, d)$ defined by (1.7) and Λ_* given as the best constant in (1.6), we have that*

$$[1 - \theta_*(p, d)] \lambda_2 < \Lambda_* \leq \lambda_2.$$

Proof. — The proof goes along the same lines as [26, Theorem 4]. We argue by contradiction and assume first that

$$[1 - \theta_*(p, d)] \lambda_2 = \Lambda_*$$

and that there is a nontrivial solution to (1.3) for $\lambda |p - 1| = \Lambda = \Lambda_*$. Then $\mathcal{J}_\Lambda[v^\beta]$ is constant with respect to t if $\lambda = \Lambda_* |p - 1|$ and v is a solution of (5.2) with initial datum v_0 such that $u = v_0^\beta$ is optimal for the functional inequality (1.6). Since the limit of $v(t, \cdot)$ as $t \rightarrow \infty$ is a positive constant that can be approximated by the average of $v(t, \cdot)$, then

$$0 = \mathcal{J}_\Lambda[v(t, \cdot)] \sim \|\nabla w\|_{L^2(\Omega)}^2 - \lambda \|w\|_{L^2(\Omega)}^2 \geq \theta_*(p, d) \lambda_2 \|w\|_{L^2(\Omega)}^2$$

with $w = v - \int_\Omega v(t, \cdot) dx \neq 0$, a contradiction.

Next we can consider non-trivial optimal functions u_λ with $\lambda > \Lambda_*$. As $\lambda \rightarrow \Lambda_*$, u_λ has to converge to a constant and we again reach a contradiction. \square

6.2. An interpolation between Poincaré and logarithmic Sobolev inequalities

It is well known that inequalities (1.6) with $p \in (0, 1)$ can be seen as a family of inequalities which interpolate between Poincaré and logarithmic Sobolev inequalities. See for instance [32]. Next, using the method introduced by W. Beckner in [6] in case of a Gaussian measure and later used for instance in [1, 21], we are in position to get an estimate of the best constant in (1.6) for any $p \in (0, 1)$ in terms of the best constant at the endpoints $p = 0$ and $p = 1$. To emphasize the dependence if the optimal constant in p , we shall denote it by $\Lambda_*(p)$ and consistently use $\Lambda_*(1)$ as

the optimal constant in the logarithmic Sobolev inequality (1.8). We recall that optimal constant in (1.6) are such that

$$\Lambda_*(p) \leq \Lambda_*(0) = \lambda_2$$

for any $p \in (0, 2^* - 1)$, including in the case $p = 1$ of the logarithmic Sobolev inequality. This can be checked easily as in the proof of Lemma 6 by using $u = 1 + \epsilon w$ as a test function, where w is an eigenfunction associated with λ_2 , and by taking the limit as $\epsilon \rightarrow 0$.

PROPOSITION 6.1. — *Assume that $p \in (0, 1)$ and $d \geq 1$. Then we have the estimate*

$$\Lambda_*(p) \geq \frac{1-p}{1-p^\alpha} \lambda_2 \quad \text{with} \quad \alpha = \frac{\lambda_2}{\Lambda_*(1)}.$$

Proof. — Let us briefly sketch the proof which is based on two main steps.

1st step: *Nelson's hypercontractivity result.* Based on the strategy of L. Gross in [30], we first establish an optimal hypercontractivity result using (1.8). Let us consider the heat equation of Ω

$$\frac{\partial f}{\partial t} = \Delta f$$

with initial datum $f(t=0, \cdot) = u$, Neumann homogeneous boundary conditions and let $F(t) := \|f(t, \cdot)\|_{L^{Q(t)}(\Omega)}$. The key computation goes as follows.

$$\begin{aligned} \frac{F'}{F} &= \frac{d}{dt} \log F(t) = \frac{d}{dt} \left[\frac{1}{Q(t)} \log \left(\int_{\Omega} |f(t, \cdot)|^{Q(t)} dx \right) \right] \\ &= \frac{Q'}{Q^2 F Q} \left[\int_{\Omega} v^2 \log \left(\frac{v^2}{\int_{\Omega} v^2 dx} \right) dx + 4 \frac{Q-1}{Q'} \int_{\Omega} |\nabla v|^2 dx \right] \end{aligned}$$

with $v := |f|^{Q(t)/2}$. Assuming that $4 \frac{Q-1}{Q'} = \frac{2}{\Lambda_*(1)}$, we find that

$$\log \left(\frac{Q(t) - 1}{p} \right) = 2 \Lambda_*(1) t$$

if we require that $Q(0) = p + 1$. Let $t_* > 0$ be such that $Q(t_*) = 2$. As a consequence of the above computation, we observe that F is non increasing by the logarithmic Sobolev inequality (1.8) and get that

$$\|f(t_*, \cdot)\|_{L^2(\Omega)} \leq \|u\|_{L^p(\Omega)} \quad \text{if} \quad \frac{1}{p} = e^{2 \Lambda_*(1) t_*}. \quad (6.1)$$

2nd step: *Spectral decomposition.* Let $u = \sum_{k \geq 1} u_k$ be a decomposition of the initial datum on the eigenspaces of $-\Delta$ with Neumann boundary conditions and denote by λ_k the ordered sequence of the eigenvalues: $-\Delta u_k =$

$\lambda_k u_k$. Let $a_k = \|u_k\|_{L^2(\Omega)}^2$. As a straightforward consequence of this decomposition, we know that $\|u\|_{L^2(\Omega)}^2 = \sum_{k \geq 1} a_k$, $\|\nabla u\|_{L^2(\Omega)}^2 = \sum_{k \geq 1} \lambda_k a_k$,

$$\|f(t_*, \cdot)\|_{L^2(\Omega)}^2 = \sum_{k \geq 1} a_k e^{-2\lambda_k t_*}.$$

Using (6.1), it follows that

$$\frac{\|u\|_{L^2(\Omega)}^2 - \|u\|_{L^p(\Omega)}^2}{1-p} \leq \frac{\|u\|_{L^2(\Omega)}^2 - \|f(t_*, \cdot)\|_{L^2(\Omega)}^2}{1-p}$$

where the right hand side can be rewritten as

$$\begin{aligned} \frac{1}{1-p} \sum_{k \geq 2} \lambda_k a_k \frac{1 - e^{-2\lambda_k t_*}}{\lambda_k} &\leq \frac{1 - e^{-2\lambda_2 t_*}}{(1-p)\lambda_2} \sum_{k \geq 2} \lambda_k a_k \\ &= \frac{1 - e^{-2\lambda_2 t_*}}{(1-p)\lambda_2} \|\nabla u\|_{L^2(\Omega)}^2. \end{aligned}$$

□

Notice that the estimate of Proposition 6.1 allows us to recover the optimal values λ_2 and $\Lambda_*(1)$ when passing to the limit in $\frac{1-p}{1-p^\alpha} \lambda_2$ as $p \rightarrow 0$ and $p \rightarrow 1$ respectively. Hence any improvement on the estimate of $\Lambda_*(1)$ automatically produces an improvement upon the lower estimate in Theorem 2 at least in a neighborhood of $p = 1_-$.

6.3. A family of special critical points for $0 < p < 1$

The equation

$$-\Delta u = f(u) \quad \text{in } \mathbb{R}^d$$

with $f(u) = \lambda u - u^p$ admits nonnegative solutions with compact support. According to [12, 13], all nonnegative solutions such that $\lim_{|x| \rightarrow \infty} u(x) = 0$ are compactly supported and the support is made of a union of disjoint balls on the interior of which the solutions are radial, and monotone decaying along a radius. Both u and ∇u vanish at the boundary of the support. Hopf's lemma does not apply because f is not Lipschitz at $u = 0$. On the contrary, we are in a case where the *compact support principle* holds: see for instance [51]. Assume now that u is radially symmetric and supported in the centered ball of radius $\rho > 0$, with $u(\rho) = u'(\rho) = 0$ and $u''(r) \sim u^p(r)$ as $r \rightarrow \rho^-$. A few integrations show that

$$u(r) \sim \left(\frac{1-p}{1+p} (\rho - r) \right)^{\frac{2}{1-p}} \quad \text{as } r \rightarrow \rho^-.$$

It is now relatively easy to describe the set of all radial solutions with compact support: see for instance [5]. We may first consider the case $\lambda = 1$ and get all solutions (including sign-changing ones) by a shooting method

$$u'' + \frac{d-1}{r} u' + |u|^{p-1} u - u = 0 \quad \text{on } (0, \infty),$$

$$\lim_{r \rightarrow 0^+} u'(r) = 0, \quad \lim_{r \rightarrow 0^+} u(r) = a,$$

in terms of the shooting parameter $a > 0$. Elementary estimates that can be found in [51] show that there is some $a > 0$ such that u is positive on $(0, \rho)$ and $u(\rho) = u'(\rho) = 0$, so that the solution u on $(0, \rho)$ can be extended by 0 on (ρ, ∞) . Let us denote this solution by u_* , by R_* the radius of the largest ball contained in Ω , and consider

$$u_\lambda(r) := \lambda^{\frac{1}{1-p}} u_*(\sqrt{\lambda} r) \quad \forall r > 0.$$

Up to translations, u_λ restricted to Ω is a nonnegative solution to (1.2) as long as $\lambda \geq \rho^2/R_*^2$. This provides us with a new upper estimate of Λ_* .

PROPOSITION 16. — *With the above notations, we have that*

$$\Lambda_* \leq \left(\frac{\rho}{R_*} \right)^{\frac{d+2-p(d-2)}{p+1}} \frac{\|\nabla u_*\|_{L^2(\mathbb{R}^d)}^2 + \|u_*\|_{L^{p+1}(\mathbb{R}^d)}^2}{\|u_*\|_{L^2(\mathbb{R}^d)}^2}.$$

Proof. — We use u_λ as a test function, undo the change of variables and optimize on the value of λ . Details are left to the reader. \square

6.4. The radial case

When Ω is a ball, all radial solutions can be obtained by a shooting method similar to the one in Section 6.3. We can indeed consider the solutions to

$$\varepsilon(p) \left(u'' + \frac{d-1}{r} u' \right) + |u|^{p-1} u - u = 0 \quad \text{on } (0, \infty),$$

$$\lim_{r \rightarrow 0^+} u'(r) = 0, \quad \lim_{r \rightarrow 0^+} u(r) = a,$$

in terms of the shooting parameter $a > 0$, and denote by $\rho_i(a)$ the zeros of u' . At least for a solution with negative initial *potential energy*, there are infinitely many such zeros. The energy is defined as

$$\mathcal{E}[u] = \frac{1}{2} |u'|^2 + \frac{1}{2} |u|^2 - \frac{1}{p+1} |u|^{p+1} \quad \text{if } p > 1,$$

$$\mathcal{E}[u] = \frac{1}{2} |u'|^2 - \frac{1}{2} |u|^2 + \frac{1}{p+1} |u|^{p+1} \quad \text{if } p < 1,$$

and it is straightforward to check that

$$\frac{d}{dr}\mathcal{E}[u(r)] = \frac{d-1}{r} |u'|^2 \leq 0$$

so that $\mathcal{E}[u(r)] < 0$ for any $r > 0$ if $\mathcal{E}[a] < 0$: a simple phase diagram analysis in (u, u') coordinates gives the result. A scaling similar to the one in Section 6.3 associates to any of the zeros $\rho_i(a)$ a branch with $\lambda = \lambda_i(a)$. Notice that for $p < 1$, one may have solutions whose support is not the whole ball, but an inner ball and eventually a disjoint annulus touching the boundary in dimension $d \geq 2$. In case of dimension $d = 1$, see next section.

6.5. The one-dimensional case

When $d = 1$ and $\Omega = (0, 1)$, the computations of Section 6.5 are explicit. Theorem 2 is not valid but Lemma 11 holds. We also get that $\frac{d}{dr}\mathcal{E}[u(r)] = 0$ and it is straightforward to check that branches with solutions having sign changing derivatives can be build from the branch of monotone solutions. Such solutions are obtained by gluing monotone solutions and scaling them. Increasing and decreasing solutions are symmetric and uniquely defined by their energy if $p > 1$. On the opposite, if $p < 1$, one may have solutions whose support is not the whole interval, but a union of disjoint intervals. On the interior of their supports, the solutions are anyhow uniquely defined. Much more is of course known, for instance on the size of the interval when λ is fixed: see for instance [55, 52].

6.6. Improvements based on the nonlinear flow

Let us define the exponent

$$\delta := \frac{p+1-(3-p)\beta}{2\beta(p-1)}.$$

Improvements of (1.6) can be obtained as in [15, 22], using the following interpolation lemma.

LEMMA 17. — *Assume that $\beta > 1$, and $\beta \leq \frac{2}{3-p}$ if $p < 3$. For any $u = v^\beta \in H^1(\Omega)$ such that $\|u\|_{L^{p+1}(\Omega)} = 1$, we have*

$$\int_{\Omega} \frac{|\nabla v|^4}{v^2} dx \geq \frac{1}{\beta^2} \frac{\int_{\Omega} |\nabla u|^2 dx}{\left(\int_{\Omega} u^2 dx\right)^\delta}.$$

Proof. — With $\frac{1}{2} + \frac{\beta-1}{2\beta} + \frac{1}{2\beta} = 1$, Hölder's inequality shows that

$$\begin{aligned} \int_{\Omega} |\nabla v|^2 dx &= \int_{\Omega} \frac{|\nabla v|^2}{v} 1 v dx \\ &\leq \left(\int_{\Omega} \frac{|\nabla v|^4}{v^2} dx \right)^{\frac{1}{2}} \left(\int_{\Omega} 1 dx \right)^{\frac{\beta-1}{2\beta}} \left(\int_{\Omega} v^{2\beta} dx \right)^{\frac{1}{2\beta}}, \end{aligned}$$

from which we deduce that

$$\left(\int_{\Omega} \frac{|\nabla v|^4}{v^2} dx \right)^{\frac{1}{2}} \geq \frac{\int_{\Omega} |\nabla v|^2 dx}{\left(\int_{\Omega} u^2 dx \right)^{\frac{1}{2\beta}}}$$

because $|\Omega| = 1$. With $\frac{1}{2} + \frac{\beta-1}{\beta(p-1)} + \frac{2-(3-p)\beta}{2\beta(p-1)} = 1$, Hölder's inequality shows that

$$\begin{aligned} \frac{1}{\beta^2} \int_{\Omega} |\nabla(v^\beta)|^2 dx &= \int_{\Omega} v^{2(\beta-1)} |\nabla v|^2 dx = \int_{\Omega} \frac{|\nabla v|^2}{v} v^{\frac{(p+1)(\beta-1)}{p-1}} \cdot v^{\frac{\beta(p-3)+2}{p-1}} dx \\ &\leq \left(\int_{\Omega} \frac{|\nabla v|^4}{v^2} dx \right)^{\frac{1}{2}} \left(\int_{\Omega} v^{\beta(p+1)} dx \right)^{\frac{\beta-1}{\beta(p-1)}} \left(\int_{\Omega} v^{2\beta} dx \right)^{\frac{\beta(p-3)+2}{2\beta(p-1)}}, \end{aligned}$$

from which we deduce that

$$\left(\int_{\Omega} \frac{|\nabla v|^4}{v^2} dx \right)^{\frac{1}{2}} \geq \frac{1}{\beta^2} \frac{\int_{\Omega} |\nabla u|^2 dx}{\left(\int_{\Omega} u^2 dx \right)^{\frac{\beta(p-3)+2}{2\beta(p-1)}}}.$$

□

For any $\beta > 1$, we define

$$\varphi(s) := \int_0^s \exp \left[\kappa \left((1 - (p-1)z)^{1-\delta} - (1 - (p-1)s)^{1-\delta} \right) \right] dz$$

where $\kappa = \frac{\mathfrak{R}}{\beta(\beta-1)(p+1)}$ and \mathfrak{R} appears in Proposition 14, and let

$$\Phi(s) := (1 + (p-1)s) \varphi \left(\frac{s}{1 + (p-1)s} \right).$$

THEOREM 18. — *Assume that one of the following conditions is satisfied:*

- (i) $d = 2$ and $p \in (0, 1) \cup (1, 8 + 4\sqrt{3})$,
- (ii) $d \geq 3$ and $p \in (0, 1) \cup (1, 2^* - 1)$.

For any $u \in H^1(\Omega)$ be such that $\|u\|_{L^2(\Omega)} = 1$, we have the inequality

$$\frac{1-\theta}{p-1} \lambda_2 \Phi \left(\frac{\|u\|_{L^p(\Omega)}^2 - 1}{p-1} \right) \leq \|\nabla u\|_{L^2(\Omega)}^2,$$

for any $\theta \in (\theta_*(p, d), 1)$ and $\beta > 1$ such that $\beta_-(\theta, p, d) < \beta < \beta_+(\theta, p, d)$, with the notations of Proposition 14.

Proof. — Let us define $\Lambda = \frac{1-\theta}{p-1} \lambda_2$

$$\mathbf{e} = \frac{1}{p-1} \left[\|u\|_{L^{p+1}(\Omega)}^2 - \|u\|_{L^2(\Omega)}^2 \right], \quad i := \|\nabla u\|_{L^2(\Omega)}^2$$

so that

$$\mathcal{J}_\Lambda[u] = i - \Lambda \mathbf{e}.$$

Using Proposition 14 and Lemma 17, we obtain the differential inequality

$$i' - \Lambda \mathbf{e}' - \frac{\mathfrak{R}}{2\beta^2} \frac{i \mathbf{e}'}{(1 - (p-1)\mathbf{e})^\delta} \leq 0$$

which can be rewritten as

$$\frac{d}{dt} (i \psi'(\mathbf{e}) - \Lambda \psi(\mathbf{e})) \leq 0$$

if φ and ψ are related by

$$\varphi(\mathbf{e}) := \frac{\psi(\mathbf{e})}{\psi'(\mathbf{e})}.$$

It is then elementary to check that φ satisfies the ODE

$$\varphi' = 1 - \varphi \frac{\psi''(\mathbf{e})}{\psi'(\mathbf{e})} = 1 + \varphi \frac{\mathfrak{R}}{2\beta^2} (1 - (p-1)\mathbf{e})^{-\delta}$$

and $\varphi(0) = 0$. □

The reader interested in the precise ranges of the exponent β and the values of θ is invited to refer to [22] for more details.

6.7. Some concluding remarks and open questions

Beyond the fact that we deal with a bounded domain with Neumann boundary conditions instead of a compact manifold with positive curvature, the lower estimate in Theorem 2 differs from the existing literature in several aspects. First of all we emphasize the fact that the convexity is needed for the computation but is certainly not necessary. The whole range of exponents corresponding to $2 < p+1 < 2^*$ is covered as in [2, 9, 34, 33] and the flow interpretation gives a nice framework, which is already present

in the results of J. Demange in [16] and has been emphasized in [26, 22]. Even better, the range $1 < p+1 < 2$ is also covered, which is new in the context of bounded domains. As the problem is set on the Euclidean space, we have neither a curvature assumption nor pointwise $CD(\rho, N)$ conditions. What matters is the Poincaré constant, which was already taken into account in the papers of J.R. Licois and L. Véron in [33, 34] and D. Bakry and M. Ledoux in [2] in the case of compact manifolds. However, we deal only with integral quantities and integrations by parts, as was emphasized in [26], still in the compact manifolds case. Last but not least, the nonlinear flow approach is also based on the methods of [21, 26] for compact manifolds, but the results of Section 6.6 on the improved inequalities as the ones obtained in [22] go beyond the results that can be achieved by standard techniques of nonlinear elliptic equations.

Problems (P2) and (P3) are equivalent. An optimal function for (P2) solves (3.1), and any solution of (3.1) is optimal as can be checked by multiplying the equation by u and integrating on Ω . The threshold for *rigidity* in (3.1) is therefore $\lambda = \mu_2$. However, this problem is of different nature than the rigidity problem in (P1). Because all terms in (3.1) are 1-homogenous (which is a consequence of the invariance of $\mathcal{Q}_\lambda[u]$ if $p > 1$) the normalization of u in $L^{p+1}(\Omega)$ is free and one can of course take $\|u\|_{L^{p+1}(\Omega)}^{p-1} = \mu$ so that u solves (1.3). Rigidity in (P1) implies rigidity in (P2). The reverse implication is not true and, up to the multiplication by a constant, all solutions of (3.1) solve (1.3), while the opposite is not true. In that sense, the set of solutions to (3.1) is larger, which explains why $\mu_1 \leq \mu_2$.

As a conclusion, let us mention a few open questions. First of all a natural question would be to try to prove that $\mu_1 = \mu_2$ in the statement of Theorem 2: under which assumptions can this be done? In the framework of compact manifolds, this is true in the case of the sphere, but it turns out that the lower estimate on μ_1 given by the nonlinear flow method is then equal to $\Lambda_\star/|p-1|$, which is definitely a very peculiar case.

Are there cases for which $\mu_1 < \mu_2$? For more complex interpolation inequalities on a cylinder, it has been established in [18] that this happens and the interested reader is invited to refer to [19, 20] for more details of qualitative nature. If Ω is a ball numerical computations when $d = 2$ and $p = 2$ also show that $\mu_1 < \mu_2$ as long as the study is done within the radial setting, but the branch of solutions corresponding to $\mu(\lambda) < \lambda$ is generated by non-radial functions. If $\mu_1 = \mu_2$, then μ_1 is also a threshold value for the existence of non-constant solutions: for any $\mu > \mu_1$ such solutions

indeed exist. Is this also what happens if $\mu_1 < \mu_2$, or are there values of $\mu \in (\mu_1, \mu_2)$ such that all positive, or at least nonnegative, solutions are in fact constants ?

Branches of solutions and bifurcations have been the subject of numerous papers and we did not review the existing literature, but at least one can mention an interesting problem. We know that optimal potentials in (P3) are related with optimal functions in (P2). Is it possible to take advantage of the spectral information in Problem (P3) to get information on branches of solutions associated with (P1) ?

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