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Guy Barles, Erwin Topp

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LIPSCHITZ REGULARITY FOR CENSORED SUBDIFFUSIVE INTEGRO-DIFFERENTIAL EQUATIONS WITH SUPERFRACTIONAL GRADIENT TERMS.

GUY BARLES AND ERWIN TOPP

ABSTRACT. In this paper we are interested in integro-differential elliptic and parabolic equations involving nonlocal operators with order less than one, and a gradient term whose coercivity growth makes it the leading term in the equation. We obtain Lipschitz regularity results for the associated stationary Dirichlet problem in the case when the nonlocality of the operator is confined to the domain, feature which is known in the literature as censored nonlocality. As an application of this result, we obtain strong comparison principles which allow us to prove the well-posedness of both the stationary and evolution problems, and steady/ergodic large time behavior for the associated evolution problem.

1. INTRODUCTION.

In [8], the authors, in collaboration with O. Ley and S. Koike, investigate regularity properties for subsolutions of integro-differential stationary equations with a super-fractional gradient term, providing analogous results for nonlocal equations to the ones of Capuzzo-Dolcetta, Porretta and Leoni [14] for superquadratic degenerate elliptic second-order pdes (see also Barles [1]). In the present work, our aim is to obtain analogous results but for integro-differential operators of order $\sigma < 1$, still with a super-fractional gradient term, but in the (intriguing) case of *censored operators* set in bounded domains.

In order to be more specific, we consider the model problem

$$(1.1) \quad \lambda u(x) + (-\Delta)_c^{\sigma/2} u(x) + b(x)|Du(x)|^m = f(x) \quad \text{in } \Omega,$$

where $\Omega \subset \mathbb{R}^N$ is a bounded domain, $\lambda \geq 0$ and $b, f : \bar{\Omega} \rightarrow \mathbb{R}$ are continuous functions with $b(x) > 0$ on $\bar{\Omega}$. For $\sigma \in (0, 2)$, the integro-differential operator $(-\Delta)_c^{\sigma/2}$ is known in the literature as the *censored fractional Laplacian of order σ* , and is defined through the expression

$$(1.2) \quad (-\Delta)_c^{\sigma/2} u(x) = C_{N,\sigma} \text{ P.V. } \int_{x+z \in \Omega} [u(x+z) - u(x)] |z|^{-(N+\sigma)} dz,$$

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where P.V. stands for the Cauchy principal value and $C_{N,\sigma} > 0$ is a well-known normalizing constant, see [13, 20]. The “censored” appellative is referred to the fact that the integration on the set $\{x + z \in \Omega\}$ makes the jumps outside Ω being indeed censored.

In the scope of this paper, besides the mentioned censored nonlocality feature of the problem, the main assumptions are $0 < \sigma < 1$ (subdiffusive operator of order less than 1), and $m > \sigma$ (superfractional coercivity condition). The study of this case is motivated by two principal reasons: first, under such conditions we are able to obtain regularity properties that are (maybe surprisingly) more sophisticated than in [8]; indeed, we are not only able to obtain global Hölder continuity but also global Lipschitz regularity for bounded subsolutions of (1.1). Secondly, we obtain comparison principle and well-posedness of the Dirichlet problem both for stationary and evolution equations, namely

$$(1.3) \quad u_t + \lambda u + (-\Delta)_c^{\sigma/2} u + b(x) |Du|^m = f(x) \quad \text{in } \Omega \times (0, \infty).$$

Concerning Dirichlet problems for nonlocal equations, we remark that in general the *Dirichlet boundary condition* has to be imposed on the complementary of Ω . Typically, if we replace $(-\Delta)_c^{\sigma/2}$ in (1.1) or (1.3) by the fractional Laplacian $(-\Delta)^{\sigma/2}$ defined as

$$(-\Delta)^{\sigma/2} u(x) = C_{N,\sigma} \text{ P.V. } \int_{\mathbb{R}^N} [u(x+z) - u(x)] |z|^{-(N+\sigma)} dz,$$

then this requires the value of the function in the whole space to be evaluated, and the Dirichlet problem reads

$$(1.4) \quad \begin{cases} \lambda u(x) + (-\Delta)^{\sigma/2} u(x) + b(x) |Du|^m = f(x) & \text{in } \Omega, \\ u = \varphi & \text{in } \Omega^c. \end{cases}$$

On the contrary, in the censored case, since we use only the values of u in $\bar{\Omega}$, we can complement (1.1) with a *classical* boundary condition

$$(1.5) \quad u = \varphi \quad \text{on } \partial\Omega,$$

for a boundary data $\varphi \in C(\partial\Omega)$.

A third interesting question that can be handled with the development of regularity and comparison properties on the current setting is the study of the large time behavior for Cauchy-Dirichlet problem associated to (1.3). We both study the cases when the censored parabolic problem has a steady state asymptotic behavior, and the case of the *ergodic large time behavior*, situation in which we have to solve a stationary problem with state-constraint boundary condition (the *ergodic problem*).

We want to mention immediately a very important point related to our regularity and well-posedness issues (and, as a consequence, the large time behavior issues) and which justifies our choice of the parameters $0 < \sigma < 1$ and $m > \sigma$: under these conditions, we are able to solve the censored Dirichlet problem in its full generality, but this is because we use in a key way the regularity result for (1.1). On the contrary, in the case $\sigma \geq 1$

and $m > \sigma$, we are unable to prove that the censored Dirichlet problem is well-posed, even with the regularity results of [8].

Now we detail the different points discussed above, starting with regularity. At this respect, we remark the results given in [14, 1] are concerned with superquadratic second-order degenerate elliptic problems like

$$(1.6) \quad \lambda u - \text{Tr}(A(x)D^2u) + b(x)|Du|^m = f(x) \quad \text{in } \Omega,$$

i.e. when $m > 2$. In [14, 1], the authors prove that if $u : \Omega \rightarrow \mathbb{R}$ is a bounded *viscosity subsolution* of (1.6), then u is locally Hölder continuous with exponent $\alpha := (m-2)(m-1)^{-1}$ and the local Hölder seminorm depends only on the data (L^∞ bounds on A, b and upper bound on f) and $\|\lambda u^-\|_\infty$, where $u^- = \min(u, 0)$.

In many interesting situations ($\lambda = 0$ or cases where we have a bound on λu which depends only on the data), this Hölder seminorm does not really depend on any L^∞ bound nor oscillation of u and actually provides an estimate on the L^∞ norm of u . The Hölder exponent $(m-2)(m-1)^{-1}$ just comes from a simple balance of powers in (1.6) and this Hölder regularity can be extended up to the boundary of the domain if it is regular enough.

This strategy was recently applied in [8] to get regularity results to integro-differential problems for which (1.1) is a particular case. However, the presence of the nonlocal term has an effect on the results in two main directions: first, the global Hölder exponent found by this method does not follow anymore the “natural” balance of powers in (1.1). Related to this, we must take into account the fact that (super)solutions to superfractional nonlocal problems (censored or not) may not satisfy the boundary condition in the classical sense (we come back later to this fact in a more detailed way). This creates a difficulty which is more evident in the non-censored setting because in this framework we must consider the evaluation of the nonlocal term on functions developing boundary jump discontinuities, and this gives way to unbounded terms near the boundary. A second consequence of the non locality of the operator is that the Hölder regularity results of [8] do not provide, in general, an estimate on the L^∞ norm of u but on the contrary rely on this L^∞ norm. This is where censored problems come into play: in these cases, and if $m > 1$, the Hölder regularity results of [8] provide an estimate on the L^∞ norm of u (or more precisely on its oscillation).

In this article, we obtain the global Hölder regularity results of [8] in the case $\sigma < 1$ and $m > \sigma$ through a simpler proof, and we extend these results to global Lipschitz continuity for bounded subsolutions to (1.1). As in [8], the censorship of the operator leads us to a control of the oscillation for subsolutions to (1.1) when $m > 1$.

The next step is to consider the well-posedness of the (stationary and evolution) nonlocal Dirichlet problem, and as we already mentioned above, we must deal with the presence of *loss of the boundary condition*. We shall mention that this phenomena also arises in local pdes like (1.6) because of the leading effect of the gradient term. Indeed, the *classical* Dirichlet

problem (with boundary data being really satisfied by the solution) cannot be solved in general and one has to use the *generalized* Dirichlet problem in the sense of viscosity solutions. We refer to the Users' guide [17] and references therein for an introduction of this concept and to [6, 24] for the applications to (1.6)-(1.5) where it is shown that the generalized Dirichlet problem is well-posed in $C(\bar{\Omega})$ or $C(\bar{\Omega} \times [0, T])$, and with examples where the solution is different from φ at points on $\partial\Omega$.

For nonlocal equations, the possible loss of the boundary conditions was first studied systematically in Barles, Chasseigne and Imbert [4] where an explicit example of such loss of boundary condition is provided for a non-censored operator. The first study of Dirichlet problems for nonlocal equations with loss of boundary conditions was done in [27], where well-posedness for non-censored nonlocal Hamilton-Jacobi problems which are not necessarily coercive in the gradient are obtained, and then in [12] in the non-censored but coercive case. As it can be seen in [12, 27], the non-censorship of the operator allows to incorporate the exterior data φ in (1.4) into the equation. Roughly speaking, this procedure creates “new” terms coming from the nonlocal operator, including an extra discount factor that turns out the problem to be strictly proper, from which the solvability of the Dirichlet problem can be obtained even if the λ term is negative.

The current censored setting shares some aspects of the pde framework since the problem is really set on $\bar{\Omega}$ and no extra information outside $\bar{\Omega}$ comes into play. However, the censorship of the operator must be regarded as an x -dependence, and it is known that such dependence creates special difficulties not only in the study of well-posedness [7], but also in regularity [3]. It is because of this difficulty that the case $\sigma \geq 1$ is not treated in the full generality neither for Dirichlet nor for Neumann problems (see [5]). Of course, in the Dirichlet case, the difficulty comes from the loss of boundary conditions, otherwise a more standard comparison result is enough to prove the well-posedness : this is for example the case for $\sigma \geq 1, m \leq \sigma$.

In this article the Lipschitz continuity of subsolutions allows us to turn around these difficulties and to obtain a strong comparison principle for (1.1)-(1.5) and its evolution counterpart. In this last case, we have to use a (classical) regularization in time to reduce to a situation where we have Lipschitz continuity in x . The key point is that this regularity permits to control the estimates of the nonlocal term when it is of order $\sigma < 1$.

We remark that the global Hölder regularity results of [8] would not be enough to treat the case $\sigma \geq 1$ and to treat the case $\sigma < 1$ would have been more involved. Moreover, we recall that the generalized notion of solution involves an evaluation of the equation on the boundary, and such an evaluation has no clear meaning in superdiffusive problems (that is when $\sigma \geq 1$ in (1.2)) mainly because of the asymmetry of the domain of integration. In fact, some special requirements on the test functions must be considered in this context, see [5, 20].

We finish with the application of regularity and comparison results in the study of large time behavior of the solutions of the evolution problem. In the case when $\lambda > 0$, this problem has the expected convergence to the steady-state solution, while if $\lambda = 0$ and $m > 1$, we prove that the asymptotic behavior of this problem resembles the second-order case studied by T. Tchamba in [26], i.e. a suitable ergodic asymptotic behavior

$$u(x, t) = -ct + u_\infty(x) + o(1) ,$$

where the function u_∞ is a solution to the ergodic problem.

The steady-state asymptotic behavior is closely related with the uniqueness of the associated stationary problem and the strict positiveness of λ is a nondegeneracy condition for this problem. It is interesting to compare this steady-state behavior with the corresponding asymptotic behavior in non-censored problems. As we mentioned above, the non-censorship of the operator implies the existence of an extra discount factor leading to the mentioned nondegeneracy condition even if λ is negative, provided the nonlocal operator recovers sufficient information from Ω^c .

A qualitatively different asymptotic behavior is observed in the case $\lambda = 0$ and $m > 1$, case in which we obtain robust regularity results leading to the solvability of the ergodic problem similarly to the periodic setting, where no boundary is considered, see [2, 8]. However, here we face an ergodic problem which is a *state-constraint* problem, as in [26]. The study of the ergodic problem is closely connected to the study of the (stationary) Dirichlet problem, while, again as in [26], the convergence of $u(x, t) + ct$ to u_∞ relies on a Strong Maximum Principle type argument. The Strong Maximum Principle we use is inspired by Coville [16] (see also Ciomaga [15]), and it is more related with a topological property of the support of the measure defining the nonlocal operator than with its ellipticity.

The article is organized along the same lines as this introduction : in Section 2, we provide our regularity results, proving first local Lipschitz continuity, then global Hölder regularity results and finally a global Lipschitz continuity result. Section 3 is devoted to state and prove the main Strong Comparison Result and to deduce the well-posedness of the initial boundary value problem. In Section 4, we study the ergodic problem, obtaining the existence and uniqueness (up to a constant) of the solutions invoking the Strong Maximum Principle. Finally we describe the different large time behavior of the solutions of the initial boundary value problem in Section 5.

Basic Notation and notion of solution. For $x \in \mathbb{R}^N$ and $r > 0$, we denote $B_r(x)$ the open ball centered at x with radius r , B_r if $x = 0$, and B if additionally $r = 1$.

For $x \in \mathcal{O}$, we denote $d_{\mathcal{O}}(x) = \text{dist}(x, \partial\mathcal{O})$. In the case $\mathcal{O} = \Omega$ we simply write $d(x)$. For $\delta > 0$, we denote $\mathcal{O}_\delta = \{x \in \mathcal{O} : d_{\mathcal{O}}(x) < \delta\}$. By the smoothness of the boundary of the domain, we can find $\delta_0 > 0$ such that the distance function is of class C^2 in Ω_{δ_0} , see [21].

We use the notion of viscosity solution with generalized boundary condition, see [17] for an introduction of this notion in the second-order setting and [4] in the nonlocal setting. For stationary problems set in a domain $\Omega \subset \mathbb{R}^N$, for a subsolution we mean an upper semicontinuous (usc for short) function on $\bar{\Omega}$ which satisfies the viscosity inequality in Ω with generalized boundary condition on $\partial\Omega$. We analogously define a supersolution as a lower semicontinuous (lsc for short) function on $\bar{\Omega}$ satisfying the corresponding viscosity inequality in Ω with generalized boundary condition on $\partial\Omega$. For a viscosity solution we mean a function which is both a viscosity sub and supersolution.

We consider the corresponding notion of viscosity solutions for the parabolic setting.

2. REGULARITY FOR THE STATIONARY PROBLEM.

In this section, we are going to provide general regularity results for equations in which the nonlocal operator has the general form

$$(2.1) \quad \mathcal{I}(u, x) = \int_{\mathbb{R}^N} [u(x+z) - u(x)] \nu_x(dz),$$

for $x \in \bar{\Omega}$, ν_x a nonnegative regular measure of \mathbb{R}^N , and for any bounded function $u : \mathbb{R}^N \rightarrow \mathbb{R}$ which has a sufficient regularity at x depending on the singularity of the measure at $z = 0$. Because of the assumptions we are going to introduce below, $\mathcal{I}(u, x)$ is well-defined if u is C^1 in $B_r(x)$ for some $r > 0$.

To be more precise, we describe the assumptions on the nonlocal term.

(M0) If $\Omega_\nu = \bigcup_{x \in \bar{\Omega}} \{x + \text{supp}\{\nu_x\}\}$, the family $\{\nu_x\}_x$ satisfies

$$\Omega_\nu \subseteq \bar{\Omega}.$$

Assumption (M0) means that, for each $x \in \bar{\Omega}$, $x+z \in \bar{\Omega}$ if $z \in \text{supp}\{\nu_x\}$ and then the domain of integration in (2.1) can be replaced just by $\bar{\Omega} - x$. For this reason we say the nonlocal operator \mathcal{I} has a *censored nature* since the jumps outside $\bar{\Omega}$ are actually censored.

Next we impose restrictions on the possible singularities of the measures ν_x .

(M1) There exists a constant $C_1 > 0$ and $\sigma \in (0, 1)$ such that, for all $\beta \in [0, 2]$ and all $\delta > 0$, we have

$$\sup_{x \in \bar{\Omega}} \int_{B_\delta^c} \min\{1, |z|^\beta\} \nu_x(dz) \leq C_1 h_{\beta, \sigma}(\delta),$$

where $h_{\beta, \sigma}(\delta)$ is defined for $\delta > 0$ as

$$h_{\beta, \sigma}(\delta) = \begin{cases} \delta^{\beta-\sigma} & \text{if } \beta < \sigma \\ |\ln(\delta)| + 1 & \text{if } \beta = \sigma \\ 1 & \text{if } \beta > \sigma, \end{cases}$$

and where we use the convention $|z|^\beta = 1, z \in \mathbb{R}^N$ when $\beta = 0$.

(M2) *There exists $C_2 > 0$ and $\sigma \in (0, 1)$ such that, for all $\beta \in (\sigma, 2]$ and all $\delta \in (0, 1)$ we have*

$$\sup_{x \in \Omega} \int_{B_\delta} |z|^\beta \nu_x(dz) \leq C_2 \delta^{\beta-\sigma}.$$

Roughly speaking, assumptions (M1) and (M2) say \mathcal{I} has at most order $\sigma \in (0, 1)$. Below we always assume that (M1) and (M2) are satisfied with the same σ .

Examples of censored nonlocal operator satisfying (M0)-(M2) are the censored fractional Laplacian of order σ (see (1.2)) and regional operators depending on the distance to the boundary of the domain (see [23]), typically

$$(2.2) \quad \mathcal{I}(u, x) = \int_{B_{d(x)}} [u(x+z) - u(x)] |z|^{-(N+\sigma)} dz.$$

For viscosity evaluation purposes, for a set $D \subset \mathbb{R}^N$, $x \in \Omega$ and $\phi \in C(\Omega)$, we define

$$\mathcal{I}[D](\phi, x) = \int_D [\phi(x+z) - \phi(x)] \nu_x(dz),$$

each time the integral makes sense.

In what follows, we are going to argue on the simpler equation

$$(2.3) \quad -\mathcal{I}(u, x) + b_0 |Du|^m = A_0, \quad \text{in } \Omega,$$

where $b_0 > 0$ and $A_0 \geq 0$.

We start with the following preliminary result concerning the interior regularity for bounded subsolutions of (2.3).

Lemma 2.1. *Let \mathcal{I} be a nonlocal operator as in (2.1) satisfying (M0)-(M2) with the same $\sigma \in (0, 1)$, and $m > \sigma$. If u is an usc, bounded viscosity subsolution to Equation (2.3), then, there exists $K, r_0 > 0$ such that, for all $0 < r < r_0$, we have*

$$|u(x) - u(y)| \leq K r^{-1} |x - y|, \quad \text{for all } x, y \in \Omega \setminus \Omega_r.$$

The constant K only depends on the data and $\text{osc}_{\bar{\Omega}}(u)$.

Proof: For $x \in \Omega$, we consider the function

$$\Phi_x : y \mapsto u(y) - \phi_x(y), \quad y \in \bar{\Omega}$$

with $\phi_x(y) := u(x) + K d^{-1}(x) |x - y|$ and $K > 0$ to be fixed later. Our aim is to prove that $\Phi_x(y) \leq 0$ for any $y \in \bar{\Omega}$ for some K large enough which does not depend on x . Indeed, if this is true for any $x \in \Omega \setminus \Omega_r$, the result is proved.

Note that the function Φ_x is usc on $\bar{\Omega}$ and therefore it attains its (non-negative) maximum at a point $\bar{y} \in \bar{\Omega}$. We argue by contradiction assuming that $\Phi_x(\bar{y}) > 0$. Since $\Phi_x(x) = 0$, we clearly have $\bar{y} \neq x$.

Furthermore, if we take $K \geq 4\text{osc}_{\bar{\Omega}}(u)$, from the inequality $\Phi_x(x) < \Phi_x(\bar{y})$ we see that $\bar{y} \in \Omega$. Indeed, it implies

$$Kd(x)^{-1}|x - \bar{y}| \leq 2\text{osc}_{\bar{\Omega}}(u),$$

and therefore $|x - \bar{y}| \leq d(x)/2$, which leads to $d(\bar{y}) \geq d(x)/2$.

We can write the viscosity subsolution inequality for u at \bar{y} since ϕ_x is smooth enough at \bar{y} and, for each $0 < \delta \leq d(x)/2$, we have

$$(2.4) \quad -\mathcal{I}[B_\delta](\phi_x, \bar{y}) - \mathcal{I}[B_\delta^c](u, \bar{y}) + b_0 K^m d^{-m}(x) \leq A_0.$$

By the Lipschitz continuity ϕ_x and (M2) we readily get

$$\mathcal{I}[B_\delta](\phi_x, \bar{y}) \leq C_2 K d^{-1}(x) \delta^{1-\sigma},$$

where the constant C_2 appears in (M2). On the other hand, by (M1) we have

$$\mathcal{I}[B_\delta^c](u, \bar{y}) \leq 2C_1 \text{osc}_{\bar{\Omega}}(u) \delta^{-\sigma},$$

and then, replacing these estimates into (2.4) we obtain

$$(2.5) \quad b_0 K^m d(x)^{-m} \leq C_2 K d(x)^{-1} \delta^{1-\sigma} + C_1 \text{osc}_{\bar{\Omega}}(u) \delta^{-\sigma} + A_0.$$

and we end up with

$$K^m \leq C \left(K d(x)^{m-1} \delta^{1-\sigma} + \text{osc}_{\bar{\Omega}}(u) d(x)^m \delta^{-\sigma} + d(x)^m \right),$$

for some constant $C > 0$ depending only on the data.

Now we choose $\delta = (1 + \text{osc}_{\bar{\Omega}}(u))d(x)/K$ (assuming that $K > 2(1 + \text{osc}_{\bar{\Omega}}(u))$) in the above expression which becomes

$$K^m \leq C \left(K^\sigma (1 + \text{osc}_{\bar{\Omega}}(u))^{1-\sigma} d(x)^{m-\sigma} + 1 \right),$$

From this last inequality and since $m > \sigma$, we conclude the result by choosing K large enough to get a contradiction. We point out that the above analysis drives us to an estimate of the size of K like

$$C \max\{(1 + \text{osc}_{\bar{\Omega}}(u))^{(1-\sigma)/(m-\sigma)}, 1 + \text{osc}_{\bar{\Omega}}(u)\},$$

for some $C > 0$ not depending on r nor $\text{osc}_{\bar{\Omega}}(u)$. The proof is complete. \square

The next lemma improves the Lipschitz seminorm.

Lemma 2.2. *Assume that Ω is a bounded, C^1 -domain, let \mathcal{I} be a nonlocal operator as in (2.1) satisfying (M0)-(M2) with the same $\sigma \in (0, 1)$, and $m > \sigma$. If u is an usc, bounded viscosity subsolution to Equation (2.3), then, there exists $C, r_0 > 0$ such that, for all $0 < r < r_0$, we have*

$$|u(x) - u(y)| \leq \bar{C} r^{-\sigma/m} |x - y| \quad \text{for } x, y \in \Omega \setminus \Omega_r,$$

where the constant \bar{C} depends on the data and the constant K of Lemma 2.1.

Proof: From Lemma 2.1 and Rademacher's Theorem, we know that any bounded subsolution of (2.3) is differentiable a.e. in Ω . Hence, we can evaluate the equation a.e. and therefore we can write

$$b_0|Du(x)|^m \leq A_0 + \mathcal{I}(u, x) \quad \text{a.e. in } \Omega.$$

Writing

$$\mathcal{I}(u, x) = \mathcal{I}[B_{d(x)/2}^c](u, x) + \mathcal{I}[B_{d(x)/2}](u, x),$$

by (M1) and (M2) we arrive at

$$b_0|Du(x)|^m \leq A_0 + C \text{osc}_{\bar{\Omega}}(u) d(x)^{-\sigma} + C d(x)^{1-\sigma} \operatorname{esssup}_{y \in B_{d(x)/2}(x)} \{|Du(y)|\},$$

for some $C > 0$ not depending on x, K or $\text{osc}_{\bar{\Omega}}(u)$. Applying Lemma 2.1 over $|Du(y)|$ we conclude

$$(2.6) \quad |Du(x)| \leq \bar{C} d(x)^{-\sigma/m},$$

where

$$(2.7) \quad \bar{C} = C(1 + K)^{1/m},$$

with $C > 0$ depending on the data, and K is given by Lemma 2.1.

The result is a direct consequence of this inequality. In fact, consider $0 < r < \min\{\delta_0, r_0\}$, where r_0 is given in Lemma 2.1 and δ_0 is such that the distance function is smooth in Ω_{δ_0} . We recall that, for any $x \in \Omega_{\delta_0}$, there exists a unique “projection” $\hat{x} \in \partial\Omega$ such that $d(x) = |x - \hat{x}|$ and \hat{x} is equal to $x - d(x)Dd(x)$. The “ \sim ” notation always denotes below such a projection.

Let $x, y \in \Omega \setminus \Omega_r$. If $|x - y| \leq r/2$, then we have

$$[x, y] := \{tx + (1 - t)y : t \in (0, 1)\} \subset \Omega.$$

Moreover, for each $z \in [x, y]$ we have $d(z) \geq r/2$. In fact

$$\begin{aligned} d(z) &= d(tx + (1 - t)y) = d(x) + (d(tx + (1 - t)y) - d(x)) \\ &\geq d(x) - |x - (tx + (1 - t)y)| \\ &\geq d(x) - (1 - t)|x - y| \\ &\geq r - r/2(1 - t), \end{aligned}$$

concluding $d(z) \geq r/2$ since $t \in [0, 1]$. We use this in the following formal computation (arguing as if u were C^1 , but the justification is more than classical) together with (2.6) to write down

$$\begin{aligned} u(x) - u(y) &= \int_0^1 \frac{d}{dt} u(tx + (1 - t)y) dt \\ &\leq C \int_0^1 d(tx + (1 - t)y)^{-\sigma/m} |x - y| dt \\ &\leq C \min\{d(x), d(y)\}^{-\sigma/m} |x - y|, \end{aligned}$$

implying that

$$(2.8) \quad |u(x) - u(y)| \leq Cr^{-\sigma/m}|x - y|.$$

On the other hand, if $|x - y| > r/2$, the only difficulty is close to the boundary and we may assume $d(x), d(y) < \delta_0$. Then, we denote $x_0, y_0 \in \Omega$ as the points such that $\hat{x}_0 = \hat{x}, \hat{y}_0 = \hat{y}$ and $d(x_0) = d(y_0) = \delta_0$. Then, we have

$$(2.9) \quad u(x) - u(y) = (u(x) - u(x_0)) + (u(x_0) - u(y_0)) + (u(y_0) - u(y)).$$

For the first and third term on the right hand side of (2.9) we have clearly $[x, x_0], [y, y_0] \subset \Omega$ and then we can apply the same argument leading to (2.8) to conclude

$$|u(x) - u(x_0)|, |u(y) - u(y_0)| \leq Cr^{-\sigma/m}|x - y|.$$

On the other hand, for the second term of the right-hand side of (2.9) we apply Lemma 2.1 and the regularity of the boundary to conclude

$$|u(x_0) - u(y_0)| \leq C|x - y|$$

with a constant C independent of r . This concludes the proof. \square

With this last lemma we are in position to prove the Hölder regularity up to the boundary.

Proposition 2.3. (GLOBAL HÖLDER ESTIMATES) *Assume that Ω is a bounded, C^1 -domain, let \mathcal{I} be a nonlocal operator as in (2.1) satisfying (M0)-(M2) with the same $\sigma \in (0, 1)$, and $m > \sigma$. If u is an usc, bounded viscosity subsolution to Equation (2.3), then, there exists a constant $C_0 > 0$ such that*

$$|u(x) - u(y)| \leq C_0|x - y|^{\frac{m-\sigma}{m}}, \quad \text{for all } x, y \in \Omega,$$

where C_0 depends on the data and $\text{osc}_{\bar{\Omega}}(u)$.

In particular, $u : \Omega \rightarrow \mathbb{R}$ can be extended as a continuous function on $\bar{\Omega}$.

Proof: If $|x - y| \leq \min\{d(x), d(y)\}$, then we use directly Lemma 2.2 to conclude

$$|u(x) - u(y)| \leq \bar{C} \min\{d(x), d(y)\}^{-\sigma/m}|x - y| \leq C\bar{C}|x - y|^{1-\sigma/m},$$

where \bar{C} is given by (2.7) and $C > 0$ depends only on the data.

Now, assuming $\min\{d(x), d(y)\} < |x - y|$ (which implies that both x and y are close to the boundary since $|x - y|$ can be chosen less than (say) $\delta_0/2$), we proceed in a similar way as Lemma 2.2, considering a parameter δ with $|x - y| \leq \delta \leq \delta_0$ and define $x_\delta, y_\delta \in \Omega$ such that $d(x_\delta) = d(y_\delta) = \delta$ and $\hat{x}_\delta = \hat{x}, \hat{y}_\delta = \hat{y}$. Then, following (2.9) we write

$$u(x) - u(y) = u(x) - u(x_\delta) + u(x_\delta) - u(y_\delta) + u(y_\delta) - u(y),$$

and using Lemma 2.2 we get $|u(x_\delta) - u(y_\delta)| \leq \bar{C}\delta^{-\sigma/m}|x - y|$, meanwhile, using again Lemma 2.2 we can write (again formally but this is easy to

justify)

$$\begin{aligned}
u(x) - u(x_\delta) &\leq \int_0^1 |Du(tx + (1-t)x_\delta)| |x - x_\delta| dt \\
&\leq \bar{C} \int_0^1 (td(x) + (1-t)\delta)^{-\sigma/m} (\delta - d(x)) dt \\
&\leq \bar{C} (1 - \sigma/m)^{-1} (\delta^{1-\sigma/m} - d(x)^{1-\sigma/m}),
\end{aligned}$$

and in the same way for $u(y_\delta) - u(y)$. This estimates imply the existence of a constant $C > 0$ such that

$$|u(x) - u(y)| \leq C\bar{C}(\delta^{1-\sigma/m} + \delta^{-\sigma/m}|x - y|) \quad \text{for all } \delta \geq |x - y|.$$

Taking infimum over these δ we arrive to the inequality

$$|u(x) - u(y)| \leq C\bar{C}|x - y|^{\frac{m-\sigma}{m}}$$

where $C > 0$ depends only on the data. We finish recalling that by the definition of $\bar{C} = C(1 + K)^{1/m}$, where the constraints on K are given at the end of the proof of Lemma 2.1. \square

The first consequence of the last proposition is the following

Corollary 2.4. (OSCILLATION BOUND) *Assume that Ω is a bounded, C^1 -domain, let \mathcal{I} be a nonlocal operator as in (2.1) satisfying (M0)-(M2) with the same $\sigma \in (0, 1)$, and $m > 1$. Let u be an usc, bounded viscosity subsolution to Equation (2.3) and consider the function*

$$\tilde{u}(x) = \begin{cases} u(x) & \text{if } x \in \Omega \\ \limsup_{y \rightarrow x, y \in \Omega} u(y) & \text{if } x \in \partial\Omega \end{cases}$$

Then, there exists $C > 0$ depending only on the data such that

$$\text{osc}_{\bar{\Omega}}(\tilde{u}) \leq C.$$

Proof: Note $\tilde{u} = u$ in Ω . In view of Proposition 2.3, we see that \tilde{u} is well-defined in $\bar{\Omega}$ and it is Hölder continuous on $\bar{\Omega}$. Moreover, we can write

$$|\tilde{u}(x) - \tilde{u}(y)| \leq C(1 + K)^{1/m} |x - y|^{(m-\sigma)/m},$$

for each $x, y \in \bar{\Omega}$. Thus, we can take $x_{\min}, x_{\max} \in \bar{\Omega}$ making

$$\tilde{u}(x_{\max}) - \tilde{u}(x_{\min}) = \text{osc}_{\bar{\Omega}}(\tilde{u}),$$

from which we obtain

$$\text{osc}_{\bar{\Omega}}(\tilde{u}) \leq C(1 + K)^{1/m} \text{diam}(\Omega)^{(m-\sigma)/m}.$$

But the end of the proof of Lemma 2.1 provides the constraint

$$K = C \max\{(1 + \text{osc}_{\bar{\Omega}}(u))^{(1-\sigma)/(m-\sigma)}, 1 + \text{osc}_{\bar{\Omega}}(u)\}$$

for some constant $C > 0$ depending only on the data. Since we assume $m > 1$, the exponent $(1 - \sigma)/(m - \sigma)$ is less than 1 and therefore we can

choose K as $C'(1 + \text{osc}_{\bar{\Omega}}(u))$ for some C' large enough depending only on the data. Hence, changing perhaps the constant, the above inequality reads

$$\text{osc}_{\bar{\Omega}}(\tilde{u}) \leq C(1 + \text{osc}_{\bar{\Omega}}(\tilde{u}))^{1/m} \text{diam}(\Omega)^{(m-\sigma)/m},$$

providing an estimate on the oscillation of u . \square

The following is the main result of this paper

Theorem 2.5. (GLOBAL LIPSCHITZ REGULARITY) *Assume that Ω is a bounded, C^1 -domain, let \mathcal{I} be a nonlocal operator as in (2.1) satisfying (M0)-(M2) with the same $\sigma \in (0, 1)$, and $m > \sigma$. If u is an usc, bounded viscosity subsolution to Equation (2.3), then, there exists a constant $L > 0$ such that*

$$|u(x) - u(y)| \leq L|x - y|, \quad \text{for all } x, y \in \Omega,$$

where the constant L depends on the data and $\text{osc}_{\bar{\Omega}}(u)$.

We point out that, in this result, m can be less or equal to 1. Notice that, while for $m > 1$ Corollary 2.4 provides a bound on the oscillation of u , we do not have such bound in the $m = 1$ -case. Therefore, in Theorem 2.5, the oscillation can be seen as an external and additional data.

In proving this theorem we require the following

Lemma 2.6. *Let \mathcal{I} be a nonlocal operator as in (2.1) satisfying (M0)-(M2) with the same $\sigma \in (0, 1)$. Then, there exists $\bar{\delta} \in (0, \delta_0)$ and $C > 0$ such that, for each $\beta \in (0, \sigma)$*

$$\mathcal{I}(d^\beta, x) \leq C d^{\beta-\sigma}(x), \quad \text{for all } x \in \Omega_{\bar{\delta}}.$$

Proof: Note that $d^\beta : \bar{\Omega} \rightarrow \mathbb{R}$ is bounded and smooth in Ω_{δ_0} . We start considering $\bar{\delta} < \delta_0/4$ and for $x \in \Omega_{\bar{\delta}}$ we write

(2.10)

$$\mathcal{I}(d^\beta, x) = \mathcal{I}[B_{d(x)/2}](d^\beta, x) + \mathcal{I}[B_{\delta_0/4} \setminus B_{d(x)/2}](d^\beta, x) + \mathcal{I}[B_{\delta_0/4}^c](d^\beta, x),$$

and we estimate each term separately. For the first integral term in the right-hand side of (2.10), we can perform a Taylor expansion to write

$$\mathcal{I}[B_{d(x)/2}](d^\beta, x) \leq \beta(d(x)/2)^{\beta-1} \int_{B_{d(x)/2}} |z| \nu_x(dz),$$

from which, applying (M2) we arrive at

$$\mathcal{I}[B_{d(x)/2}](d^\beta, x) \leq C d^{\beta-\sigma}(x).$$

For the second integral term in (2.10), we use that $z \mapsto d^\beta(z)$ is β -Hölder continuous in Ω_{δ_0} and therefore we can write

$$\mathcal{I}[B_{\delta_0/4} \setminus B_{d(x)/2}](d^\beta, x) \leq C \int_{B_{\delta_0/4} \setminus B_{d(x)/2}} |z|^\beta \nu_x(dz),$$

with $C > 0$ not depending on x or β . Applying (M1), we conclude

$$\mathcal{I}[B_{\delta_0/4} \setminus B_{d(x)/2}](d^\beta, x) \leq C d^{\beta-\sigma}(x).$$

Finally, using the boundedness of d^β in Ω , (M0) and (M1), for the third integral term in the right-hand side of (2.10), we conclude

$$\mathcal{I}[B_{\delta_0/4}^c](d^\beta, x) \leq C\delta_0^{-\sigma},$$

where $C > 0$ depends only on $\text{diam}(\Omega)$. Joining the above estimates in the right-hand side of (2.10), we conclude the result taking $\bar{\delta}$ smaller if it is necessary. \square

Proof of Theorem 2.5: Denote $\gamma = 1 - \sigma/m$ the Hölder exponent of Lemma 2.3. Let $x \in \Omega$ and $\beta \in (0, \min\{\gamma, \sigma\})$ fixed. For $L > 1$ and $0 < \eta \ll 1$, we consider the function

$$y \mapsto \Phi_x(y) := u(y) - \phi_x(y),$$

where the function ϕ_x has the form

$$\phi_x(y) = L|y - x| - \eta d^\beta(y).$$

Our aim is to prove that, for L large enough (with a size which does not depend on x), we have $\Phi_x(y) \leq 0$ on $\bar{\Omega}$, for all η small enough (possibly depending on L and/or x). If this is true, we deduce the Lipschitz continuity from this property by letting η tend to 0. Of course, the difficulty is to get such property close to the boundary and therefore we may assume that $d(x) \leq \bar{\delta}/2$ with $\bar{\delta}$ as in Lemma 2.6.

We argue by contradiction, assuming that

$$(2.11) \quad \max_{\bar{\Omega}} \Phi_x \geq \epsilon_L > 0,$$

for each L large and η small.

By the upper semicontinuity of u , there exists $\bar{y} \in \bar{\Omega}$ attaining the maximum in (2.11), and this maximum cannot be equal to x since $\epsilon_L > 0$. Moreover, $\bar{y} \in \Omega$ since otherwise, for $a > 0$ small enough, the inequality $\Phi(\bar{y} + aDd(\bar{x})) \leq \Phi_x(\bar{y})$ leads to

$$\eta a^\beta \leq u(\bar{y}) - u(\bar{y} + aDd(\bar{y})) + L(|\bar{y} + aDd(\bar{y}) - x| - |\bar{y} - x|),$$

and by the Hölder regularity given by Proposition 2.3, we arrive at

$$\eta a^\beta \leq C_0 a^\gamma + La,$$

which is a contradiction when $a > 0$ is small since $\beta < \gamma$.

Since $\bar{y} \in \Omega$ and $\bar{y} \neq x$, we can write the viscosity subsolution inequality for u with test function ϕ_x at \bar{y} . Thus, for all $\delta > 0$ we can write

$$(2.12) \quad -\mathcal{I}[B_\delta](\phi_x, \bar{y}) - \mathcal{I}[B_\delta^c](u, \bar{y}) + b_0 |D\phi_x(\bar{y})|^m \leq A_0,$$

and the idea is to estimate each term in the right-hand side separately. Note that

$$\mathcal{I}[B_\delta](\phi_x, \bar{y}) = LJ_1 - \eta J_2,$$

with

$$J_1 = \int_{B_\delta} (|\bar{y} + z - x| - |\bar{y} - x|) \nu_{\bar{y}}(dz),$$

$$J_2 = \int_{B_\delta} [d^\beta(\bar{y} + z) - d^\beta(\bar{y})] \nu_{\bar{y}}(dz).$$

Since $\sigma < 1$ and applying (M1), we easily conclude that

$$J_1 \leq C_1 \delta^{1-\sigma}.$$

Now, using that $\Phi_x(x) \leq \Phi_x(\bar{y})$ and considering η small in terms of $\text{diam}(\Omega)$, we see that

$$(2.13) \quad |\bar{y} - x| \leq L^{-1}(\text{osc}_{\bar{\Omega}}(u) + 1).$$

Since we are considering x such that $d(x) \leq \bar{\delta}/2$ with $\bar{\delta}$ as in Lemma 2.6, by the last inequality we can take L large enough (depending on $\text{osc}_{\bar{\Omega}}(u)$ and $\bar{\delta}$, and, a fortiori, on δ_0) to get $d(\bar{y}) \leq \bar{\delta}$. Then, by Lemma 2.6, we can write

$$J_2 \leq C d^{\beta-\sigma}(\bar{y}),$$

where C not depends on \bar{x}, y, L or η . Then, by the estimates for J_1 and J_2 we obtain

$$\mathcal{I}[B_\delta](\phi_x, \bar{y}) \leq CL\delta^{1-\sigma} + C\eta d^{\beta-\sigma}(\bar{y}).$$

Concerning the integral term outside B_δ , the Hölder continuity of u on $\bar{\Omega}$ given by Lemma 2.3 allows us to write

$$\mathcal{I}[B_\delta^c](u, \bar{y}) \leq C \int_{B_\delta^c} |z|^\gamma \nu_{\bar{y}}(dz) \leq C\delta^{\gamma-\sigma},$$

where $C > 0$ depends only on the data and the constant C_0 of Proposition 2.3. In summary, concerning the integral terms we have the estimate

$$(2.14) \quad \mathcal{I}[B_\delta](\phi_x, \bar{y}) + \mathcal{I}[B_\delta^c](u, \bar{y}) \leq C \left(L\delta^{1-\sigma} + \eta d^{\beta-\sigma}(\bar{x}) + \delta^{\gamma-\sigma} \right).$$

Now we deal with the first-order term. A straightforward computation gives us

$$D\phi_x(\bar{y}) = L(\widehat{\bar{y} - x}) - \eta\beta d^{\beta-1}(\bar{y})Dd(\bar{y}),$$

with

$$\widehat{\bar{y} - x} := \frac{(\bar{y} - x)}{|\bar{y} - x|}.$$

Recalling \bar{y} depends on η , at this point we have two different situations : either there exists a subsequence such that $d(\bar{y}) \geq d(x)$ as η tends to zero, or $d(\bar{y}) \leq d(x)$ as η tends to zero. In the former case, $\eta d(\bar{y})^{\beta-1} \leq \eta d(x)^{\beta-1}$ and therefore $\eta d(\bar{y})^{\beta-1} \rightarrow 0$ as $\eta \rightarrow 0$. This yields $|D\phi_x(\bar{y})| \geq L + o_\eta(1)$. In the second case, computing $|D\phi_x(\bar{y})|^2$, we see that

$$|D\phi_x(\bar{y})|^2 = L^2 + \eta^2 \beta^2 d^{2(\beta-1)}(\bar{y}) - 2\eta\beta L d^{\beta-1}(\bar{y}) \langle Dd(\bar{y}), \widehat{\bar{y} - x} \rangle.$$

But $x, \bar{y} \in \Omega_{\delta_0}$ where d is C^1 , and performing a Taylor expansion we get

$$d(x) - d(\bar{y}) = \langle Dd(\bar{y}), (x - \bar{y}) \rangle + o(|\bar{y} - x|),$$

which replaced into the above expression for $|D\phi_x(\bar{x})|^2$ drives us to

$$(2.15) \quad |D\phi_x(\bar{y})|^2 = L^2 + \eta^2 \beta^2 d^{2(\beta-1)}(\bar{y}) - 2\eta\beta L d^{\beta-1}(\bar{y})((d(\bar{y}) - d(x))/|\bar{y} - x| - o(1)),$$

where $o(1) \rightarrow 0$ when $|\bar{x} - y| \rightarrow 0$. But taking into account that $d(\bar{y}) \leq d(x)$, we are led to

$$|D\phi_x(\bar{y})|^2 \geq L^2 + \eta^2 \beta^2 d^{2(\beta-1)}(\bar{y}) - 2\eta\beta L d^{\beta-1}(\bar{y})o(1),$$

and by using Cauchy-Schwarz inequality, we arrive at

$$(2.16) \quad |D\phi_x(\bar{y})|^2 \geq L^2/2 + \eta^2 \beta^2 d^{2(\beta-1)}(\bar{y})/2,$$

for all L large and η small enough.

Then, depending on the case we are looking at, we replace (2.14) and (2.16) or $|D\phi_x(\bar{y})| \geq L + o_\eta(1)$ into (2.12) to obtain, for all $x \in \Omega, \delta > 0, L$ large (depending only on the data and the oscillation of u) and η small depending on L and x , the inequality

$$(2.17) \quad L^m + (\eta d^{\beta-1}(\bar{y}))^m \leq C \left(A_0 + L\delta^{1-\sigma} + \delta^{\gamma-\sigma} + \eta d^{\beta-\sigma}(\bar{y}) \right),$$

for some constant $C > 0$ depending only on the data.

From this inequality, we claim that there exists an universal constant $\Lambda_0 > 0$ such that

$$(2.18) \quad C\eta d^{\beta-\sigma}(\bar{y}) - (\eta d^{\beta-1}(\bar{y}))^m \leq \Lambda_0,$$

for all η small. We postpone the justification of this claim until the end of this proof.

Using this claim into (2.17), we observe that

$$L^m \leq C \left(A_0 + L\delta^{1-\sigma} + \delta^{\gamma-\sigma} + \Lambda_0 \right).$$

Hence, in the case $m \geq 1$, since $\sigma < 1$ we can fix $\delta > 0$ small in terms of C to conclude the result by taking L large in terms of the data and the oscillation of u . In the case $m < 1$, we take $\delta = L^{-1} > 0$ and since $m > \sigma$, we conclude taking L large in terms of the data and the oscillation of u . This concludes the proof of the theorem.

Now we address the claim leading to (2.18). We write

$$\theta = \eta d^{\beta-\sigma}(\bar{y}),$$

from which, we obtain that

$$C\eta d^{\beta-\sigma}(\bar{y}) - (\eta d^{\beta-1}(\bar{y}))^m = C\theta - \theta^\tau \eta^{-\alpha},$$

where $\tau := m(1-\beta)/(\sigma-\beta) > 0$ and $\alpha := m(1-\sigma)/(\sigma-\beta) > 0$. Note that since $m > \sigma$, we can fix $\beta > 0$ small enough to get $\tau > 1$. With this choice we can consider $\eta \leq C^{-1/\alpha}$ and with this we arrive at

$$C\eta d^{\beta-\sigma}(\bar{y}) - (\eta d^{\beta-1}(\bar{y}))^m \leq C(\theta - \theta^\tau) \leq C\tau^{-\tau/(\tau-1)}(\tau-1) \leq C(\tau-1) =: \Lambda_0,$$

which is a constant depending only on the data, but not on η, L or x . This concludes the claim. \square

3. COMPARISON PRINCIPLE AND WELL-POSEDNESS.

In this section we use the notation $Q = \Omega \times (0, +\infty)$ and $\partial Q = \partial\Omega \times (0, +\infty)$. Here we address the well-posedness of the following Cauchy-Dirichlet evolution problem

$$(CP) \quad \begin{cases} u_t - \mathcal{I}(u(\cdot, t), x) + H(x, u, Du) &= f, & \text{in } Q \\ u &= \varphi, & \text{in } \partial Q \\ u &= u_0, & \text{in } \bar{\Omega} \times \{0\}. \end{cases}$$

where \mathcal{I} is a nonlocal operator with the form (2.1) satisfying conditions (M0)-(M2), $H \in C(\bar{\Omega} \times \mathbb{R} \times \mathbb{R}^N)$, $f \in C(\bar{Q})$, $\varphi \in C(\partial Q)$ and $u_0 \in C(\bar{\Omega})$. We also request the following *compatibility condition* between φ and u_0

$$(3.1) \quad u_0 \equiv \varphi \quad \text{on } \partial\Omega.$$

On H we additionally assume the properness condition

(H1) For all $R > 0$, there exists $\lambda_R \geq 0$ nonnegative such that, for all $x \in \bar{\Omega}$, $u \geq v$ with $|u|, |v| \leq R$ and $p \in \mathbb{R}^N$, we have

$$H(x, u, p) - H(x, v, p) \geq \lambda_R(u - v).$$

In order to apply the regularity results of the previous section, we consider the coercivity condition

(H2) There exists $m > 0$ and $\underline{c} > 0$ such that, for each $R > 0$ there exists $C_R \geq 0$ satisfying

$$H(x, r, p) \geq \underline{c}|p|^m - C_R,$$

for all $x \in \bar{\Omega}$, $p \in \mathbb{R}^N$ and $|r| \leq R$.

Finally, in addition to (M0)-(M2), we will require the following continuity assumption over the measures $\{\nu_x\}$ defining \mathcal{I}

(M3) Let $\sigma \in (0, 1)$ given in (M1), (M2) and h as in (M1). Then, for each $\beta > 0$, there exists a modulus of continuity $\omega_{\beta, \sigma}$ such that for all $x, y \in \bar{\Omega}$ and $\delta > 0$, we have

$$\int_{B_\delta^c} |x|^\beta (\nu_x(dz) - \nu_y(dz)) \leq (1 + h_{\beta, \sigma}(\delta)) \omega_{\beta, \sigma}(|x - y|)$$

Note that (M3) holds in the censored fractional Laplacian case (see (1.2)) and regional operators depending on the distance to the boundary (see (2.2)).

The main result of this section reads as follows

Proposition 3.1. Assume that Ω is a bounded, C^1 -domain and let \mathcal{I} be a nonlocal operator as in (2.1) satisfying (M0)-(M3) with the same $\sigma \in (0, 1)$. Assume also that $H \in C(\bar{\Omega} \times \mathbb{R} \times \mathbb{R}^N)$ satisfies (H1)-(H2) with $m > \sigma$, $f \in C(\bar{Q})$, $\varphi \in C(\partial Q)$, $u_0 \in C(\bar{\Omega})$ and that (3.1) holds. If u, v are respectively

viscosity subsolution and supersolution to (CP) which are bounded in each compact subset of \bar{Q} , then

$$u \leq v \quad \text{in } Q \cup \bar{\Omega} \times \{0\}.$$

Moreover, if \tilde{u} is defined by

$$(3.2) \quad \tilde{u}(x, t) = \begin{cases} u(x, t) & \text{if } (x, t) \in Q \cup \bar{\Omega} \times \{0\} \\ \limsup_{(y, s) \rightarrow (x, t), y \in \Omega} u(y, s) & \text{if } (x, t) \in \partial Q, \end{cases}$$

then $\tilde{u} \leq v$ in \bar{Q} .

Before giving the proof of this result, we start with two lemmas concerning the (parabolic) boundary condition of problem (CP). Next result states classical initial condition holds mainly by the compatibility condition (3.1)

Lemma 3.2. *Under the assumptions of Proposition 3.1, for all $x \in \bar{\Omega}$ we have $u(x, 0) \leq u_0(x) \leq v(x, 0)$.*

We refer to [18] for a proof of Lemma 3.2 in the second-order case which can be readily adapted to the current framework. We continue with the classical boundary condition for subsolutions

Lemma 3.3. *Under the assumptions of Proposition 3.1, $u(x, t) \leq \varphi(x, t)$ for each $(x, t) \in \partial Q$.*

We may refer to [4, 12, 27] for a proof of this result in very similar settings.

Remark 3.4. *A classical strategy to deal with the well-posedness of (CP) in $C(\bar{Q})$ is to argue over the redefined function \tilde{u} in (3.2) instead of the original subsolution u , see [4, 12, 27] and the second-order references therein. Note that $\tilde{u} = u$ in Q , $\tilde{u} \leq u$ on $\partial\Omega \cup \bar{\Omega} \times \{0\}$ and therefore Lemmas 3.2 and 3.3 hold for \tilde{u} . Moreover, \tilde{u} is a (generalized) viscosity subsolution to the problem if u is. Thus, for simplicity, we are going to assume that $u = \tilde{u}$ on \bar{Q} in order to avoid the superscript “ \sim ”.*

The arguments to come are carried out on the finite time horizon problem

$$(CP_T) \quad \begin{cases} \partial_t u - \mathcal{I}(u(\cdot, t), x) + H(x, u, Du) = f, & \text{in } Q_T \\ u = \varphi, & \text{in } \partial Q_T \\ u = u_0, & \text{in } \bar{\Omega} \times \{0\}, \end{cases}$$

where, for $T > 0$, we denote $Q_T = \Omega \times (0, T]$ and $\partial Q_T = \partial\Omega \times (0, T]$. The infinite horizon setting for Proposition 3.1 is readily obtained by taking a sequence of problems (CP_T) with $T \rightarrow +\infty$.

We require the following lemmas in order to apply the regularity results of the previous section

Lemma 3.5. *Assume $\varphi \in C(\partial Q_T)$, $f \in C(\bar{Q}_T)$ and $H \in C(\bar{\Omega} \times \mathbb{R} \times \mathbb{R}^N)$ satisfies (H2). Let u be bounded usc viscosity subsolution to problem*

$$(3.3) \quad \begin{cases} \partial_t u - \mathcal{I}(u, x) + H(x, u, Du) = f & \text{in } Q_T \\ u = \varphi & \text{in } \partial Q_T, \end{cases}$$

For $\gamma > 0$ and $(x, t) \in \bar{Q}_T$, consider u^γ , the sup-convolution in time of u , which is given by the expression

$$(3.4) \quad u^\gamma(x, t) = \sup_{s \in [0, T]} \{u(x, s) - \gamma^{-1}(s - t)^2\}.$$

Then, there exists a constant $a_\gamma > 0$ with $a_\gamma \rightarrow 0$ as $\gamma \rightarrow 0$ such that u^γ is a viscosity subsolution to problem

$$\begin{cases} u_t^\gamma - \mathcal{I}(u^\gamma(\cdot, t), x) + H(x, u^\gamma, Du^\gamma) = f + o_\gamma(1), & \text{in } \Omega \times (a_\gamma, T] \\ u^\gamma = \varphi^\gamma, & \text{in } \partial\Omega \times [a_\gamma, T], \end{cases}$$

where $o_\gamma(1) \rightarrow 0$ as $\gamma \rightarrow 0$ and depends exclusively on the modulus of continuity of f .

The proof of the above lemma follows closely the arguments of [9, 12]. Note that by Lemma 3.3 the boundary condition in the above equation is satisfied in the classical sense.

It is a well-known fact that for each $\gamma > 0$ and $x \in \bar{\Omega}$, $t \mapsto u^\gamma(x, t)$ is Lipschitz continuous in $[0, T]$, with Lipschitz constant $C_\gamma := 4T\gamma^{-1}$. Therefore, for any $t \in (a_\gamma, T]$, the function $x \mapsto u^\gamma(x, t)$ is a subsolution to

$$-\mathcal{I}(u, x) + H(x, u, Du) \leq f + o_\gamma(1) + C_\gamma \text{ in } \Omega.$$

Applying Theorem 2.5, we conclude the following

Lemma 3.6. *Let u be a bounded viscosity subsolution to problem (3.3). Let $\gamma > 0$, u^γ defined as in (3.4) and a_γ the constant given in Lemma 3.5. Then, $u^\gamma \in \text{Lip}(\bar{\Omega} \times [a_\gamma, T])$.*

Now we are in position to provide the

Proof of Proposition 3.1: By contradiction, we assume that $u - v$ is positive at some point on \bar{Q}_T . Then, for all $\eta > 0$ small in terms of T and the supremum of $u - v$ on \bar{Q}_T , we see that

$$M := \sup_{\bar{Q}_T} \{u - v - \eta t\} > 0.$$

By semicontinuity, this supremum is attained at some point $(x_0, t_0) \in \bar{Q}_T$. Taking η smaller if it is necessary, by Lemma 3.2 we have $t_0 > 0$. From this point, we fix such an $\eta > 0$.

For $\alpha > 0$, notice the function

$$(x, t) \mapsto u(x, t) - v(x, t) - \eta t - \alpha(|x - x_0|^2 + (t - t_0)^2)$$

attains its unique maximum on \bar{Q}_T at (x_0, t_0) , and in fact this maximum equals M . Recalling the definition of u^γ in (3.4), we define

$$(3.5) \quad M_\gamma := \sup_{(x, t) \in \bar{Q}_T} \{u^\gamma(x, t) - v(x, t) - \eta t - \alpha(|x - x_0|^2 + (t - t_0)^2)\}.$$

By definition of u^γ we see that $M_\gamma \geq M$ and this supremum is attained at some point $(x_\gamma, t_\gamma) \in \bar{Q}_T$. Since (x_0, t_0) is a strict maximum point, properties of the sup-convolution imply that $(x_\gamma, t_\gamma) \rightarrow (x_0, t_0)$ as $\gamma \rightarrow 0$.

We are going to consider the case $x_\gamma \in \partial\Omega$ for all γ . The case $x_\gamma \in \Omega$ with $d(x_\gamma)$ uniformly positive with respect to γ is obtained by easier arguments and computations, see [12, 27].

We follow Soner's penalization procedure introduced in [25] (see also [17, 12, 27]). We double variables, consider a parameter $\epsilon > 0$ and define the function

$$\Phi(x, y, s, t) := u^\gamma(x, s) - v(y, t) - \phi(x, y, s, t),$$

where, denoting $\bar{\xi} = (Dd(x_\gamma), 0) \in \mathbb{R}^{N+1}$, we define

$$\phi(x, y, s, t) := |\epsilon^{-1}((x, s) - (y, t)) - \bar{\xi}|^2 + \eta s + \alpha(|y - x_0|^2 + (t - t_0)^2).$$

Since Φ is upper semicontinuous in $\bar{Q}_T \times \bar{Q}_T$, there exists a point $(\bar{x}, \bar{y}, \bar{s}, \bar{t}) \in \bar{Q}_T \times \bar{Q}_T$ where the function Φ attains its maximum. By the inequality

$$\Phi(\bar{x}, \bar{y}, \bar{s}, \bar{t}) \geq \Phi(x_\gamma + \epsilon Dd(x_\gamma), x_\gamma, t_\gamma, t_\gamma)$$

and the continuity of u^γ , classical viscosity arguments lead us to

(3.6)

$$(\bar{x}, \bar{s}), (\bar{y}, \bar{t}) \rightarrow (x_\gamma, t_\gamma), \quad |\epsilon^{-1}(\bar{x} - \bar{y}) - Dd(x_\gamma)| \rightarrow 0, \quad \epsilon^{-1}|\bar{s} - \bar{t}| \rightarrow 0,$$

$$\text{and } u^\gamma(\bar{x}, \bar{s}) \rightarrow u^\gamma(x_\gamma, t_\gamma), \quad v(\bar{y}, \bar{t}) \rightarrow v(x_\gamma, t_\gamma),$$

as $\epsilon \rightarrow 0$. Moreover, for all ϵ small, the same inequality implies that

$$\Phi(\bar{x}, \bar{y}, \bar{s}, \bar{t}) \geq M/2.$$

This inequality implies that $u^\gamma(\bar{x}, \bar{s}) > v(\bar{y}, \bar{t})$ and therefore, using Lemma 3.3 and the continuity of φ , we deduce that, if $\bar{y} \in \partial\Omega$,

$$(3.7) \quad \varphi(\bar{y}, \bar{t}) > v(\bar{y}, \bar{t}).$$

Thus, the supersolution inequality holds at (\bar{y}, \bar{t}) for v even if $\bar{y} \in \partial\Omega$. On the other hand, the second convergence in (3.6) leads us to

$$(3.8) \quad \bar{x} = \bar{y} + \epsilon Dd(x_\gamma) + o(\epsilon),$$

from which we conclude that $\bar{x} \in \Omega$ for all ϵ small. Indeed, by a simple Taylor's expansion, d being C^1

$$d(\bar{x}) = d(\bar{y}) + \epsilon Dd(x_\gamma) \cdot Dd(\bar{y}) + o(\epsilon),$$

and since $\bar{y} \rightarrow x_\gamma$ as $\epsilon \rightarrow 0$, we have

$$d(\bar{x}) \geq \epsilon |Dd(x_\gamma)|^2 + o(\epsilon),$$

and $d(\bar{x}) > 0$ for ϵ small enough since $|Dd(x_\gamma)| = 1$. In conclusion, $\bar{x} \in \Omega$ and we can write down a subsolution viscosity inequality for u^γ at (\bar{x}, \bar{s}) .

Now we subtract the corresponding viscosity inequalities for u^γ at (\bar{x}, \bar{s}) (see Lemma 3.5) and v at (\bar{y}, \bar{t}) and then, for all $\delta > 0$ we can write

$$(3.9) \quad \mathcal{A} - \mathcal{I}^\delta - \mathcal{I}_\delta \leq o_\gamma(1),$$

where

$$\mathcal{I}_\delta = \mathcal{I}[B_\delta](\phi(\cdot, \bar{y}, \bar{s}, \bar{t}), \bar{x}) - \mathcal{I}[B_\delta](-\phi(\bar{x}, \cdot, \bar{s}, \bar{t}), \bar{y}),$$

$$\mathcal{I}^\delta = \mathcal{I}[B_\delta^c](u^\gamma(\cdot, \bar{s}), \bar{x}) - \mathcal{I}[B_\delta^c](v(\cdot, \bar{t}), \bar{y}),$$

and

$$\begin{aligned}\mathcal{A} &= (\partial_s \phi + \partial_t \phi)(\bar{x}, \bar{y}, \bar{s}, \bar{t}) - f(\bar{x}, \bar{s}) + f(\bar{y}, \bar{t}) \\ &\quad + H(\bar{x}, u^\gamma(\bar{x}, \bar{s}), \bar{p}) - H(\bar{y}, v(\bar{y}, \bar{t}), \bar{q}),\end{aligned}$$

with

$$\begin{aligned}\bar{p} &= D_x \phi(\bar{x}, \bar{y}, \bar{s}, \bar{t}) = \epsilon^{-1}(\epsilon^{-1}((\bar{x}, \bar{s}) - (\bar{y}, \bar{t})) - \bar{\xi}), \\ \bar{q} &= -D_y \phi(\bar{x}, \bar{y}, \bar{s}, \bar{t}) = \bar{p} - 2\alpha(\bar{y} - x_0).\end{aligned}$$

The core of the remaining of the proof is to estimate $\mathcal{I}_\delta, \mathcal{I}^\delta$ and \mathcal{A} . We start with the later. Recalling that $u^\gamma \in \text{Lip}(\bar{\Omega} \times [a_\gamma, T])$ and that $(\bar{x}, \bar{y}, \bar{s}, \bar{t})$ is a maximum point for Φ , we have $|\bar{p}| \leq L_\gamma$ for all ϵ small, where $L_\gamma > 0$ is the Lipschitz constant for u^γ given in Lemma 3.6. Denoting $R = \|u^\gamma\|_\infty + \|v\|_\infty$, applying (H1) together with the (uniform) continuity of f in \bar{Q}_T and (3.6), we see that

$$\mathcal{A} \geq \eta - o_\epsilon(1) + \lambda_R(u^\gamma(\bar{x}, \bar{s}) - v(\bar{y}, \bar{t})) + H(\bar{x}, v(\bar{y}, \bar{t}), \bar{p}) - H(\bar{y}, v(\bar{y}, \bar{t}), \bar{q}),$$

with $o_\epsilon(1) \rightarrow 0$ as $\epsilon \rightarrow 0$ depending only on the modulus of continuity of f . Hence, by the uniform continuity of H on compact sets of $\bar{\Omega} \times \mathbb{R} \times \mathbb{R}^N$ we arrive at

$$\mathcal{A} \geq \eta - o_\alpha(1) - o_\epsilon(1) + \lambda_R(u^\gamma(\bar{x}, \bar{s}) - v(\bar{y}, \bar{t})),$$

and by the definition of M_γ in (3.5) together with (3.6) and the continuity of λ_R , we conclude

$$(3.10) \quad \mathcal{A} \geq \eta - o_\alpha(1) - o_\epsilon(1) + \lambda_R M.$$

Now we address the nonlocal terms. By the smoothness of ϕ , we have

$$(3.11) \quad \mathcal{I}_\delta \leq \epsilon^{-1} o_\delta(1),$$

where $o_\delta(1) \rightarrow 0$ as $\delta \rightarrow 0$ and does not depend on ϵ .

On the other hand, for \mathcal{I}^δ we can write

$$\begin{aligned}\mathcal{I}^\delta &= \int_{B_\delta^c} [u^\gamma(\bar{x} + z, \bar{s}) - u^\gamma(\bar{x}, \bar{s})](\nu_{\bar{x}}(dz) - \nu_{\bar{y}}(dz)) \\ &\quad + \int_{B_\delta^c} [u^\gamma(\bar{x} + z, \bar{s}) - v(\bar{y} + z, \bar{t}) - (u^\gamma(\bar{x}, \bar{s}) - v(\bar{y}, \bar{t}))]\nu_{\bar{y}}(dz) \\ &=: \mathcal{I}_1^\delta + \mathcal{I}_2^\delta.\end{aligned}$$

For \mathcal{I}_1^δ , by Lemma 3.6 and (M3) we have

$$\mathcal{I}_1^\delta \leq C L_\gamma \omega_{1,\sigma}(|\bar{x} - \bar{y}|),$$

for some $C > 0$ not depending on ϵ or δ . Thus, by (3.6) we can write

$$\mathcal{I}_1^\delta \leq L_\gamma o_\epsilon(1).$$

Now, for \mathcal{I}_2^δ , we divide the region of integration as

$$B_\delta^c = (B_\delta^c \cap (\Omega - \bar{x})) \cup (B_\delta^c \setminus (\Omega - \bar{x})) =: \Theta_1 \cup \Theta_2.$$

At this point we notice that (M0) says that $\text{supp}\{\nu_{\bar{y}}\} \subset \bar{\Omega} - \bar{y}$ and therefore, by using that $(\bar{x}, \bar{y}, \bar{s}, \bar{t})$ is a maximum for Φ , we can write

$$\mathcal{I}_2^\delta \leq C\alpha \int_{\Theta_1} |z| \nu_{\bar{y}}(dz) + 2R \int_{\Theta_2} \nu_{\bar{y}}(dz).$$

By (M1) we can control the first integral term in the right-hand side of the last inequality by $C\alpha$ for some universal constant $C > 0$. Using again that $\text{supp}\{\nu_{\bar{y}}\} \subset \bar{\Omega} - \bar{y}$, applying (3.8) and the arguments given in [12, 27] we conclude that Θ_2 is uniformly away the origin and vanishes as $\epsilon \rightarrow 0$. This last fact and (M2) lead us to the following estimate

$$\int_{\Theta_2} \nu_{\bar{y}}(dz) = o_\epsilon(1),$$

and therefore we arrive at

$$\mathcal{I}_2^\delta \leq C\alpha + Ro_\epsilon(1).$$

Using this, (3.11) and (3.10) into (3.9) we can write

$$\eta + \lambda_R M - o_\alpha(1) - o_\epsilon(1) - \epsilon^{-1} o_\delta(1) \leq o_\gamma(1).$$

Recalling that λ_R is nonnegative, we arrive to a contradiction with $\eta > 0$ by taking $\delta \rightarrow 0, \epsilon \rightarrow 0, \gamma \rightarrow 0$ and $\alpha \rightarrow 0$. \square

As it is usual in the viscosity theory, comparison principle lead us to the following well-posedness result for (CP).

Proposition 3.7. *Assume hypotheses of Proposition 3.1 hold. Then, there exists a unique $u \in C(\bar{Q})$, viscosity solution to problem (CP). For each $T > 0$, this solution satisfies*

$$|u(x, t)| \leq (\|H(\cdot, 0, 0)\|_{L^\infty(\bar{\Omega})} + \|f\|_{L^\infty(\bar{Q}_T)}) t + \|\varphi\|_{L^\infty(\partial Q_T)} + \|u_0\|_{L^\infty(\bar{\Omega})},$$

for all $(x, t) \in \bar{Q}_T$. Moreover, if there exists $\lambda_0 > 0$ such that, for all $R > 0$, $\lambda_R \geq \lambda_0$, and f, φ are uniformly bounded, then u is uniformly bounded in \bar{Q} , with

$$|u(x, t)| \leq \lambda_0^{-1} (\|H(\cdot, 0, 0)\|_{L^\infty(\bar{\Omega})} + \|f\|_{L^\infty(\bar{Q})}) + \|\varphi\|_{L^\infty(\partial Q)} + \|u_0\|_{L^\infty(\bar{\Omega})},$$

for all $(x, t) \in \bar{Q}$.

The existence of a solution in the case u_0, φ are smooth follows from the application of Perron's method for discontinuous solutions with generalized boundary conditions and comparison principle, see [17, 22, 18]. The existence and uniqueness for general continuous initial and boundary data is obtained by approximation through smooth data and viscosity stability. The L^∞ -estimates for the solution are easily obtained by taking sub and supersolutions for the problem with the form $(x, t) \mapsto \pm(C_1 t + C_2)$, with $C_1 \geq 0$ and $C_2 > 0$ suitable constants depending on the data.

Following the same lines of the proof of Proposition 3.1 we can get the corresponding results for the state-constraint evolution problem and for the stationary equation. Both results are presented next.

Proposition 3.8. *Under the assumptions of Proposition 3.1, comparison principle holds for bounded viscosity sub and supersolutions of the state-constraint problem*

$$(3.12) \quad \begin{cases} \partial_t u - \mathcal{I}(u, x) + H(x, u, Du) = f & \text{in } Q_T, \\ \partial_t u - \mathcal{I}(u, x) + H(x, u, Du) \geq f & \text{in } \partial Q_T, \\ u = u_0 & \text{on } \bar{\Omega} \times \{0\}. \end{cases}$$

Roughly speaking, we understand a (viscosity) solution to (3.12) as a function which is a subsolution in Q_T and a viscosity supersolution in \bar{Q}_T , see [19].

Proposition 3.9. *Let \mathcal{I} be as in (2.1) defined through a family of measures $\{\nu_x\}$ satisfying (M0)-(M3) with the same $\sigma < 1$. Let $\varphi \in C(\partial\Omega)$, $H \in C(\bar{\Omega} \times \mathbb{R} \times \mathbb{R}^n)$ satisfying (H1) and (H2) with $m > \sigma$. We consider the problem*

$$(3.13) \quad \begin{cases} -\mathcal{I}(u) + H(x, u, Du) = 0 & \text{in } \Omega, \\ u = \varphi & \text{in } \partial\Omega, \end{cases}$$

Further, assume one of the following additional hypotheses on H holds

- (i) There exists $\lambda_0 > 0$ such that, for all $R > 0$, we have $\lambda_R \geq \lambda_0$.
- (ii) For each $x \in \bar{\Omega}$, the function $(r, p) \mapsto H(x, r, p)$ is convex and problem (3.13) has a bounded strict subsolution.

If u, v are bounded, respective sub and supersolution to the problem (3.13), then $u \leq v$ in Ω . Moreover, if we define

$$(3.14) \quad \tilde{u}(x, t) = \begin{cases} u(x) & \text{if } x \in \Omega, \\ \limsup_{y \rightarrow x, y \in \Omega} u(y) & \text{if } x \in \partial\Omega, \end{cases}$$

then $\tilde{u} \leq v$ in $\bar{\Omega}$. In the above setting, there exists a unique solution $u \in C(\bar{\Omega})$ for (3.13).

We finish this section with the following well-known consequence of comparison principle. See [26] for a proof.

Lemma 3.10. *Assume the hypotheses of Proposition 3.1 hold. Assume further that $\varphi \in C(\partial Q)$, $f \in C(\bar{Q})$ are bounded and $u_0 \in C^1(\bar{\Omega})$ with $\|u_0\|_{C^1(\bar{\Omega})} \leq \Lambda$ for some $\Lambda > 0$. Then, the unique viscosity solution $u \in C(\bar{Q})$ to problem (CP) satisfies*

$$|u(x, t) - u(x, s)| \leq C\Lambda|t - s|, \quad \text{for all } s, t > 0; x \in \bar{\Omega},$$

where $C > 0$ depends only on the data.

The importance of this lemma is that once we assume the initial data is smooth, we have the time derivative of the solution to (CP) is bounded. Thus, by Theorem 2.5 we see that the solution to (CP) is uniformly Lipschitz (in space and time) on \bar{Q} .

4. THE ERGODIC PROBLEM.

In this section we are interested in the existence and uniqueness of a constant $c \in \mathbb{R}$ for which the state-constraint problem

$$(4.1) \quad \begin{cases} -\mathcal{I}(u) + H_0(x, Du) = c & \text{in } \Omega, \\ -\mathcal{I}(u) + H_0(x, Du) \geq c & \text{on } \partial\Omega, \end{cases}$$

has a bounded viscosity solution. Here and below H_0 is a continuous function which satisfies (H2) with $m > 1$. Such a problem is known in the literature as the *ergodic problem*. The key questions are the uniqueness of c as well as the structure of the solutions u to this problem. For definition and further concepts associated to state-constraint problems we refer to [19].

4.1. Solvability. Existence of c comes as a consequence of comparison and regularity results given in the previous sections.

Proposition 4.1. *Let \mathcal{I} as in (2.1) satisfying (M0)-(M3) with the same $\sigma < 1$ and $H_0 \in C(\bar{\Omega} \times \mathbb{R}^N)$ satisfying (H2) with $m > 1$. Then, there exists a unique constant $c \in \mathbb{R}$ for which the problem (4.1) has a viscosity solution in $C(\bar{\Omega})$, which is Lipschitz continuous on $\bar{\Omega}$.*

Following Tchamba [26], the above proposition can be obtained through the next lemmas.

Lemma 4.2. *Let \mathcal{I} as in (2.1) satisfying (M0)-(M3) with the same $\sigma < 1$, and $H_0 \in C(\bar{\Omega} \times \mathbb{R}^N)$ satisfying (H2). For each $\alpha \in (0, 1)$, there exists a unique viscosity solution $u_\alpha \in C(\bar{\Omega})$ to the problem*

$$(4.2) \quad \begin{cases} \alpha u - \mathcal{I}(u) + H_0(x, Du) = 0 & \text{in } \Omega, \\ \alpha u - \mathcal{I}(u) + H_0(x, Du) \geq 0 & \text{on } \partial\Omega, \end{cases}$$

and such a solution satisfies the inequality

$$(4.3) \quad \|\alpha u_\alpha\|_{L^\infty(\bar{\Omega})} \leq \tilde{C},$$

for $\tilde{C} > 0$ independent of α .

Proof: For $R > 0$, we are going to introduce the intermediate problem

$$(4.4) \quad \begin{cases} \alpha u - \mathcal{I}(u) + H_0(x, Du) = 0 & \text{in } \Omega, \\ u = R & \text{on } \partial\Omega, \end{cases}$$

Since $\alpha > 0$, by Proposition 3.9 there exists a unique solution $u_{\alpha,R} \in C(\bar{\Omega})$ to (4.4). Since the constant function equals to $-\alpha^{-1}\|H_0(\cdot, 0)\|_{L^\infty(\bar{\Omega})}$ is a subsolution for this problem, we have $u_{\alpha,R} \geq -\alpha^{-1}\|H_0(\cdot, 0)\|_{L^\infty(\bar{\Omega})}$. The idea is to provide an upper bound which is independent of R . For this, we recall $\delta_0 > 0$ is such that the function $x \mapsto \text{dist}(x, \partial\Omega)$ is smooth in the set Ω_{δ_0} . We are going to denote $d \in C^2(\Omega) \cap C(\bar{\Omega})$ a nonnegative function defined as $d(x) = \text{dist}(x, \partial\Omega)$ if $x \in \bar{\Omega}_{\delta_0}$ and which is strictly positive (depending only on δ_0) in $\Omega \setminus \Omega_{\delta_0}$. Then, for $\beta \in (0, 1)$ and by Lemma 2.6 we see that

$$\mathcal{I}(d^\beta, x) \leq C d^{\beta-\sigma}(x), \quad \text{for all } x \in \Omega.$$

For $x \in \Omega$ and by the coercivity of H_0 , evaluating the function $-d^\beta$ into (4.4) we see that

$$\begin{aligned} & \alpha(-d^\beta(x)) - \mathcal{I}(-d^\beta, x) + H_0(x, -\beta d^{\beta-1}(x)Dd(x)) \\ & \geq -\text{diam}(\Omega)^\beta - Cd^{\beta-\sigma}(x) + c\beta^m d^{m(\beta-1)}(x) - C, \end{aligned}$$

where $C, c > 0$ are universal constants. Recalling the coercivity degree of H_0 given by $m > \sigma$ in (H2), fixing $0 < \beta < (m - \sigma)/(m - 1)$ if $m > 1$ and any $\beta > 0$ if $m \leq 1$, and taking $\bar{\delta} < \delta_0/4$ very small in terms of the data, we conclude $-d^\beta$ is a supersolution to (4.4) in the set $\Omega_{\bar{\delta}}$. Moreover, this function is a supersolution up to $\partial\Omega$ in the generalized sense because there is no smooth function touching $-d^\beta$ from below at the boundary. Additionally, there exists a constant $C_1 \geq 0$ (depending only on the data and $\bar{\delta}$) such that, for all $x \in \Omega \setminus \Omega_{\bar{\delta}/2}$ we have

$$\lambda(-d^\beta(x)) - \mathcal{I}(-d^\beta, x) + H_0(x, -\beta d^{\beta-1}(x)Dd(x)) \geq -C_1.$$

Then, considering the function

$$\psi(x) = 2\alpha^{-1} \left(\|H_0(\cdot, 0)\|_{L^\infty(\bar{\Omega})} + C_1 \right) - d^\beta(x),$$

we see by the above discussion that ψ is a viscosity supersolution (in the generalised boundary sense) to (4.4) and therefore, by strong comparison principle, for all $R > 0$ we get the estimate

$$-\alpha^{-1} \|H_0(\cdot, 0)\|_{L^\infty(\bar{\Omega})} \leq u_{\alpha, R} \leq 2\alpha^{-1} \left(\|H_0(\cdot, 0)\|_{L^\infty(\bar{\Omega})} + C_1 \right).$$

Hence, defining $\tilde{C} = 2(\|H_0(\cdot, 0)\|_{L^\infty(\bar{\Omega})} + C_1)$ and taking $R > \tilde{C}/\alpha$, the solution $u_{\alpha, R}$ to (4.4) cannot satisfy the boundary condition for supersolutions in the classical sense and then it satisfies the state-constraint problem (4.2). By uniqueness, we conclude the result labeling $u_\alpha = u_{\alpha, R}$ for all $R > \tilde{C}/\alpha$. The estimate (4.3) for u_α is inherited from the corresponding inequality for $u_{\alpha, R}$. \square

Lemma 4.3. *Let \mathcal{I} as in (2.1) satisfying (M0)-(M3) and $H_0 \in C(\bar{\Omega} \times \mathbb{R}^N)$ satisfying (H2). Then, for each $\alpha \in (0, 1)$, the solution $u_\alpha \in C(\bar{\Omega})$ of the problem (4.2) is Lipschitz continuous in $\bar{\Omega}$ with Lipschitz constant depending on the data and $\text{osc}_{\bar{\Omega}}(u_\alpha)$.*

In particular, if $m > 1$ in (H2), the Lipschitz constant depends only on the data and not on α nor $\|u_\alpha\|_{L^\infty(\bar{\Omega})}$.

The proof comes as a combination of (4.3) and Theorem 2.5. The additional property for superlinear Hamiltonians comes from Corollary 2.4.

Proof of Proposition 4.1: Fix $x^* \in \Omega$ and denoting u_α the unique solution to (4.2) given in Lemma 4.2, we define

$$v_\alpha(x) = u_\alpha(x) - u_\alpha(x^*), \quad x \in \bar{\Omega},$$

which is a viscosity solution to

$$\begin{cases} \alpha v_\alpha - \mathcal{I}(v_\alpha) + H_0(x, Dv_\alpha) = -\alpha u_\alpha(x^*) & \text{in } \Omega, \\ \alpha v_\alpha - \mathcal{I}(v_\alpha) + H_0(x, Dv_\alpha) \geq -\alpha u_\alpha(x^*) & \text{on } \partial\Omega, \end{cases}$$

By (4.3) we see that $\alpha u_\alpha(x^*)$ is uniformly bounded as $\alpha \rightarrow 0$. Thus, by Corollary 2.4 and Lemma 4.3, the family $\{v_\alpha\}_{\alpha \in (0,1)}$ is uniformly bounded and equi-Lipschitz. Then, letting $\alpha \rightarrow 0$, classical stability results provide us the existence of a pair $(v, c) \in \text{Lip}(\bar{\Omega}) \times \mathbb{R}$, viscosity solution to (4.1).

Concerning the uniqueness of c , we consider $(v_1, c_1), (v_2, c_2) \in \text{Lip}(\bar{\Omega}) \times \mathbb{R}$ solving the ergodic problem. It is direct to see that for each $i = 1, 2$, the function $w_i(x, t) := v_i(x) - c_i t$ solves the parabolic state-constraint problem

$$\begin{cases} \partial_t w_i - \mathcal{I}(w_i, x) + H_0(x, Dw_i) = 0 & \text{in } Q_T, \\ \partial_t w_i - \mathcal{I}(w_i, x) + H_0(x, Dw_i) \geq 0 & \text{in } \partial Q_T, \\ w_i = v_i & \text{on } \bar{\Omega} \times \{0\}. \end{cases}$$

Applying comparison principle given in Proposition 3.8, we see that

$$w_1(x, t) \leq w_2(x, t) + \|v_1 - v_2\|_{L^\infty(\bar{\Omega})} \quad \text{for all } (x, t) \in \bar{Q}.$$

Thus, $c_2 - c_1 \leq 2\|v_1 - v_2\|_{L^\infty(\bar{\Omega})}/t$ and we arrive at $c_2 \leq c_1$ making $t \rightarrow \infty$. Exchanging the roles of c_1 and c_2 we conclude the uniqueness of c . \square

4.2. Strong Maximum Principle. Here we provide a version of the Strong Maximum Principle which is going to play a key role in the arguments to come. This Strong Maximum Principle is inspired by arguments from Coville [16] (see also Ciomaga [15]).

For this, we introduce the evolutive counterpart of (4.1)

$$(4.5) \quad \begin{cases} \partial_t u - \mathcal{I}(u) + H_0(x, Du) = c & \text{in } Q, \\ \partial_t u - \mathcal{I}(u) + H_0(x, Du) \geq c & \text{on } \partial Q. \end{cases}$$

Recalling $H_0 \in C(\bar{\Omega} \times \mathbb{R}^N)$ trivially satisfies (H1), Proposition 3.8 provides us comparison principle for (4.5) and therefore we can get

Lemma 4.4. *Assume (M0)-(M3) holds with the same $\sigma < 1$ and $H_0 \in C(\bar{\Omega} \times \mathbb{R}^N)$ satisfies (H2). Let u usc in \bar{Q} , v lsc in \bar{Q} with $u, v \in L^\infty(\bar{Q}_T)$ for all $T > 0$ be respective sub and supersolutions to problem (4.5). For $t \in [0, +\infty)$, define*

$$(4.6) \quad \kappa(t) = \sup_{x \in \bar{\Omega}} \{u(x, t) - v(x, t)\}.$$

Then, for all $0 \leq s \leq t$, we have $\kappa(t) \leq \kappa(s)$.

Next we introduce some notation: let $\{\nu_x\}_x$ in the definition of \mathcal{I} and for $x \in \mathbb{R}^N$ we define inductively

$$X_0(x) = \{x\}, \quad X_{n+1}(x) = \bigcup_{\xi \in X_n(x)} \{\xi + \text{supp}\{\nu_\xi\}\}, \quad \text{for } n \in \mathbb{N},$$

and

$$(4.7) \quad \mathcal{X}(x) = \overline{\bigcup_{n \in \mathbb{N}} X_n}.$$

The Strong Maximum Principle presented here relies in the nonlocality of the operator under the “iterative covering property” which is close to the ideas of Coville [16] but it has to be combined here with different arguments. This property is established through the condition

$$(4.8) \quad \mathcal{X}(x) = \bar{\Omega}, \quad \text{for all } x \in \bar{\Omega}.$$

Notice this condition is satisfied by our two main examples (1.2) and (2.2).

Proposition 4.5. (STRONG MAXIMUM PRINCIPLE) *Let \mathcal{I} as in (2.1), where $\{\nu_x\}$ satisfies (M0)-(M3) and (4.8), and $H_0 \in C(\bar{\Omega} \times \mathbb{R}^N)$ satisfying (H2).*

1.- *Strong Maximum Principle - Parabolic Version: Let u, v be respective bounded, viscosity sub and supersolution to (4.5), with $u = \tilde{u}$ as in (3.2). Let κ as in (4.6) and assume there exists $t_0 > 0$ satisfying*

$$\kappa(t_0) = \sup_{t \geq 0} \{\kappa(t)\}.$$

Then, the function $u - v$ is constant in $\bar{\Omega} \times [0, t_0]$. Moreover, we have

$$(u - v)(x, t) = \kappa(0), \quad \text{for all } (x, t) \in \bar{\Omega} \times [0, t_0].$$

2.- *Strong Maximum Principle - Stationary Version: Let u, v be respective bounded viscosity sub and supersolution to (4.1), with $u = \tilde{u}$ as in (3.14). Then, $u - v$ is constant in $\bar{\Omega}$.*

Proof: We focus in the parabolic version, and we start defining $T = t_0 + 1$. As we did it in the proof of comparison principle Proposition 3.1, we may assume u is Lipschitz continuous (in space and time) in \bar{Q}_T by replacing u by its sup-convolution u^γ . We avoid the direct use of u^γ for simplicity.

Notice that by Lemma 4.4, we see that $\kappa(t) = \kappa(0)$ for all $t \in [0, t_0]$. Fix $\tau \in (0, t_0)$, denote

$$\mathcal{M}_\tau = \{x \in \bar{\Omega} : (u - v)(x, \tau) = \kappa(\tau)\},$$

and choose a point $x_\tau \in \mathcal{M}_\tau$. For simplicity, we also assume that $x_\tau \in \Omega$, otherwise we apply Soner type penalization procedure used in the proof of the comparison principle instead of the function Φ below. In that case, the control of the integral terms can be made in the same way as in the proof of the comparison principle.

Consider $\epsilon, \alpha > 0$ and define the function

$$\Phi(x, y, s, t) = u(x, s) - v(y, t) - \phi(x, y, s, t),$$

with $\phi(x, y, s, t) := \epsilon^{-2}(|x - y|^2 + (s - t)^2) + \alpha((s - \tau)^2 + |x - x_\tau|^2)$. This function attains its maximum $(\bar{x}, \bar{y}, \bar{s}, \bar{t})$ in $\bar{Q}_T \times \bar{Q}_T$, and using the inequality

$$\Phi(\bar{x}, \bar{y}, \bar{s}, \bar{t}) \geq \Phi(x_0, x_0, \tau, \tau),$$

we see that, keeping $\alpha > 0$,

$$(4.9) \quad \bar{s}, \bar{t} \rightarrow \tau, \quad \bar{x}, \bar{y} \rightarrow x_\tau, \quad \text{and} \quad u(\bar{x}, \bar{s}) \rightarrow u(x_\tau, \tau), \quad v(\bar{y}, \bar{t}) \rightarrow v(x_\tau, \tau),$$

as $\epsilon \rightarrow 0$. Moreover, by the Lipschitz continuity of u , $\bar{p} := 2\epsilon^{-2}(\bar{x} - \bar{y})$ is uniformly bounded as $\epsilon \rightarrow 0$. Thus, properly using ϕ as a test function for u at (\bar{x}, \bar{s}) and for v at (\bar{y}, \bar{t}) and subtracting the corresponding viscosity inequalities, for each $\delta > 0$ we can write

$$(4.10) \quad 2\alpha(\bar{s} - \tau) - (\epsilon^{-1} + \alpha)o_\delta(1) - \mathcal{I}^\delta + \mathcal{A} \leq 0,$$

where

$$\begin{aligned} \mathcal{I}^\delta &:= \mathcal{I}[B_\delta^c](u(\cdot, \bar{s}), \bar{x}) - \mathcal{I}[B_\delta^c](v(\cdot, \bar{t}), \bar{y}) \\ \mathcal{A} &:= H_0(\bar{x}, \bar{p} + \alpha(\bar{x} - x_\tau)) - H_0(\bar{y}, \bar{p}). \end{aligned}$$

Note that by the continuity of H_0 , using (4.9) and the boundedness of \bar{p} , we readily conclude that

$$\mathcal{A} = o_\epsilon(1) + o_\alpha(1).$$

Now we deal with the nonlocal terms. Using (M3), the Lipschitz continuity of u and (4.9), we can write

$$\begin{aligned} \mathcal{I}^\delta &\leq L o_\epsilon(1) + \int_{B_\delta^c \cap (\Omega - \bar{x})} (u(\bar{x} + z, \bar{s}) - v(\bar{y} + z, \bar{t}) - (u(\bar{x}, \bar{s}) - v(\bar{y}, \bar{t}))) \nu_{\bar{y}}(dz) \\ &\quad - \int_{B_\delta^c \setminus (\Omega - \bar{x})} (v(\bar{y} + z, \bar{t}) - v(\bar{y}, \bar{t})) \nu_{\bar{y}}(dz) \\ &=: L o_\epsilon(1) + \mathcal{I}_1^\delta - \mathcal{I}_2^\delta, \end{aligned}$$

where $L > 0$ is the Lipschitz constant for u . But using that $x_\tau \in \Omega$, by the second convergence in (4.9) we see that $B_{\delta'} \subset (\Omega - \bar{x})$ for each $0 < \delta' \leq d(x_\tau)/2$. By (M3) we have $\text{supp}\{\nu_{\bar{y}}\} \subset (\Omega - \bar{y})$ and using again the second convergence in (4.9) we obtain $|\text{supp}\{\nu_{\bar{y}}\} \setminus (\Omega - \bar{x})| \rightarrow 0$ as $\epsilon \rightarrow 0$. This together with the boundedness of v and (M1) allows us to write, for all $0 < \delta < \delta'$

$$\mathcal{I}_2^\delta = o_\epsilon(1),$$

meanwhile, using that $(\bar{x}, \bar{y}, \bar{s}, \bar{t})$ is maximum for Φ and (M2) lead us to

$$\mathcal{I}_1^\delta \leq \alpha o_{\delta'}(1) + \int_{B_{\delta'}^c \cap (\Omega - \bar{x})} (u(\bar{x} + z, \bar{s}) - v(\bar{y} + z, \bar{t}) - (u(\bar{x}, \bar{s}) - v(\bar{y}, \bar{t}))) \nu_{\bar{y}}(dz),$$

Replacing the above estimates for \mathcal{A} and \mathcal{I}^δ in (4.10), we make $\delta \rightarrow 0$ to get rid of the term $\epsilon^{-1} o_\delta(1)$. Recalling (4.9), by the boundedness assumptions for u and v and the continuity of the measure given by (M3),

keeping $\delta' > 0$ and making $\epsilon \rightarrow 0$ and finally $\alpha \rightarrow 0$, Dominated Convergence Theorem leads us to

$$- \int_{B_{\delta'}^c \cap (\Omega - x_\tau)} (u(x_\tau + z, \tau) - v(x_\tau + z, \tau) - (u(x_\tau, \tau) - v(x_\tau, \tau))) \nu_{x_\tau}(dz) \leq 0,$$

but since $x_\tau \in \mathcal{M}_\tau$ and by the upper semicontinuity of $u - v$, we arrive at

$$\int_{B_{\delta'}^c \cap (\Omega - x_\tau)} (u(x_\tau + z, \tau) - v(x_\tau + z, \tau) - (u(x_\tau, \tau) - v(x_\tau, \tau))) \nu_{x_\tau}(dz) = 0.$$

Thus, since $\delta' > 0$ is arbitrary, we conclude that $u - v \equiv \kappa(\tau)$ in $X_1(x_\tau)$ and proceeding inductively as above, we obtain

$$(u - v)(x, \tau) = \kappa(\tau) \quad \text{for each } x \in \bigcup_{n \in \mathbb{N}} X_n(x_\tau),$$

from which we obtain $(u - v)(\cdot, \tau) = \kappa(\tau)$ in $\bar{\Omega}$ by applying (4.8). The result for $\tau = 0$ and $\tau = t_0$ can be easily obtained by Lemma 4.4 and the upper semicontinuity of $u - v$. \square

As a consequence of the Strong Maximum Principle, we obtain the following

Proposition 4.6. *Let \mathcal{I} as in (2.1) satisfying (M0)-(M3) and the iterative covering property (4.8), let $H_0 \in C(\bar{\Omega} \times \mathbb{R}^N)$ satisfying (H2) with $m > 1$ and c the unique ergodic constant given in Proposition 4.1. Then, the solution to (4.1) is unique up to an additive constant.*

Proof: Consider $(v_i, c) \in \text{Lip}(\bar{\Omega}) \times \mathbb{R}$, $i = 1, 2$ be two solutions to (4.1). Then, we may cast v_1 as a subsolution (in Ω) to the problem for which v_2 is a solution. Thus, we conclude $v_2 = v_1 + C$ for some $C \in \mathbb{R}$ by the Strong Maximum Principle. \square

5. LARGE TIME BEHAVIOR.

For simplicity, in this section we concentrate on the parabolic problem

$$(CP') \quad \begin{cases} \partial_t u + \lambda u - \mathcal{I}(u) + H_0(x, Du) &= 0 & \text{in } Q, \\ u &= \varphi & \text{on } \partial Q, \\ u &= u_0 & \text{on } \bar{\Omega} \times \{0\}, \end{cases}$$

where $\lambda \geq 0$. In the rest of this section, we always assume \mathcal{I} is as in (2.1) and satisfies (M0), (M1)-(M2) with the same $\sigma \in (0, 1)$, and (M3). We also assume $H_0 \in C(\bar{\Omega} \times \mathbb{R}^N)$ satisfies (H2) with $m > \sigma$, and $u_0 \in C(\bar{\Omega})$, $\varphi \in C(\partial\Omega)$ satisfy the compatibility condition (3.1). In this setting, problem (CP') can be uniquely solved in $C(\bar{Q})$.

For simplicity, we assume a smooth initial data u_0 in (CP') and then the solution to this parabolic problem will be Lipschitz continuous in \bar{Q} with Lipschitz constant depending on the data and the $\|u_0\|_{C^1(\bar{\Omega})}$, see Lemma 3.10.

We get the large time behavior for the general case $u_0 \in C(\bar{\Omega})$ by approximation through smooth initial data, see [26] for a complete exposition of these arguments.

We start with the study of parabolic problems for which we have steady state asymptotic behavior.

Theorem 5.1. *Assume one of the following conditions hold*

(i) $\lambda > 0$.

(ii) $\lambda = 0$, the ergodic constant c associated to the ergodic problem (4.1) is negative, and the function $p \mapsto H_0(x, p)$ is convex for all $x \in \bar{\Omega}$.

In both cases, there exists a unique $C(\bar{\Omega})$ -viscosity solution to the problem

$$(S_\lambda) \quad \begin{cases} \lambda u - \mathcal{I}(u) + H_0(x, Du) &= 0 & \text{in } \Omega, \\ u &= \varphi & \text{on } \partial\Omega, \end{cases}$$

and therefore, the unique solution to (CP') converges uniformly on $\bar{\Omega}$ as $t \rightarrow \infty$ to the unique viscosity solution to (S_λ) .

Proof: The solvability of (S_λ) in the case $\lambda > 0$ comes as a consequence of the case (i) in Proposition 3.9. In case (ii), we notice that a solution to (4.1) is a strict subsolution to problem (S_0) and we fall in case (ii) of Proposition 3.9, from which the solvability of (S_0) holds.

We note that in both cases (i) and (ii), the solution u to (CP') is bounded in \bar{Q} . In fact, this can be easily seen in case (i) through Proposition 3.7. For case (ii), we see that the function

$$(x, t) \mapsto u_\infty(x) - \|u_\infty - \varphi\|_{L^\infty(\partial\Omega)} - \|u_\infty - u_0\|_{L^\infty(\bar{\Omega})},$$

where u_∞ is a solution to the ergodic problem (4.1), is a viscosity subsolution to the problem satisfied by u , and by comparison principle we conclude

$$u_\infty - \|u_\infty - \varphi\|_{L^\infty(\partial\Omega)} - \|u_\infty - u_0\|_{L^\infty(\bar{\Omega})} \leq u \quad \text{on } \bar{Q}.$$

To find an upper bound for u , we fix $\bar{x} \in \Omega$ and denote $K = 2\text{diam}(\Omega)$. By (H2) there exists a constant $C_K > 0$ depending on the data and K such that the function

$$(x, t) \mapsto C_K(1 - K^{-1}|x - \bar{x}|) + \|\varphi\|_{L^\infty(\partial\Omega)} + \|u_0\|_{L^\infty(\bar{\Omega})},$$

is a viscosity supersolution to (CP') and from this we see that

$$u \leq C_K + \|\varphi\|_{L^\infty(\partial\Omega)} + \|u_0\|_{L^\infty(\bar{\Omega})} \quad \text{on } \bar{Q}.$$

From this point, we argue simultaneously for cases (i) and (ii). The above analysis implies u is uniformly bounded on \bar{Q} . Then, the functions

$$\bar{u}(x, t) = \limsup_{\epsilon \rightarrow 0, z \rightarrow x} u(z, t/\epsilon), \quad \underline{u}(x, t) = \liminf_{\epsilon \rightarrow 0, z \rightarrow x} u(z, t/\epsilon),$$

are well-defined in \bar{Q} . Naturally we have $\underline{u} \leq \bar{u}$ in $Q \cup \bar{\Omega} \times \{0\}$. Besides, for each $t > 0$, applying half-relaxed limits method [10, 11] we have the functions

$$x \mapsto \bar{u}(x, t) \quad \text{and} \quad x \mapsto \underline{u}(x, t)$$

are respective viscosity sub and supersolution to the problem (S_λ) . Thus, by comparison we have $\tilde{u} \leq \underline{u}$ and then they coincide in Ω .

We claim that $\tilde{u} = \bar{u}$ on $\partial\Omega$ and postpone its justification until the end of the convergence proof. If this happens, then the function

$$x \mapsto U(x, t) := \bar{u}(x, t) = \underline{u}(x, t) = \lim_{z \rightarrow x, \epsilon \rightarrow 0} u(z, t/\epsilon), \quad x \in \bar{\Omega}$$

is in $C(\bar{\Omega})$ and it is a viscosity solution to (S_λ) , which is the unique one by Proposition 3.9. From this, we easily conclude the uniform convergence on $\bar{\Omega}$ of the solution of the parabolic problem (CP') to the stationary problem (S_λ) as $t \rightarrow \infty$.

Now we deal with the claim. Let $x_0 \in \partial\Omega$. Note that by definition, for each $\eta > 0$ small, there exist $y_\eta, z_\eta \in \Omega$ and $\epsilon_\eta > 0$ satisfying $|x_0 - y_\eta|, |x_0 - z_\eta| \leq \eta$ and such that

$$\bar{u}(x_0, t) - \tilde{u}(x_0, t) \leq \eta + u(y_\eta, t/\epsilon_\eta) - u(z_\eta, t/\epsilon_\eta).$$

But we have u is uniformly Lipschitz in \bar{Q} (see Lemma 3.10) and then

$$\bar{u}(x_0, t) - \tilde{u}(x_0, t) \leq \eta + C\|u_0\|_{C^1(\bar{\Omega})}\eta,$$

for some $C > 0$ depending only on the data. A similar lower bound can be stated and since η is arbitrary, we conclude the equality. \square

Now we address the ergodic large time behavior.

Theorem 5.2. *Assume $\lambda = 0$, \mathcal{I} satisfies the iterative covering property (4.8) and that H_0 satisfies (H2) with $m > 1$. Let $u \in C(\bar{Q})$ be the unique solution to (CP') and $(u_\infty, c) \in \text{Lip}(\bar{\Omega}) \times \mathbb{R}$ be an ergodic pair solution to (4.1). Then*

1.- CASE I. *If $c > 0$, then (S_0) has no bounded viscosity solution. Moreover, the function $(x, t) \mapsto u(x, t) + ct$ is uniformly bounded in \bar{Q} and*

$$u + ct \rightarrow u_\infty + K \quad \text{uniformly in } \bar{\Omega}, \quad \text{as } t \rightarrow +\infty,$$

for some constant $K \in \mathbb{R}$ depending on H, u_0, φ and c .

2.- CASE II. *If $c = 0$, then any bounded solution to (S_0) has the form $u_\infty + K$ for some constant $K \in \mathbb{R}$ such that $u_\infty + K \leq \varphi$ on $\partial\Omega$. Moreover,*

$$u \rightarrow u_\infty + K \quad \text{uniformly in } \bar{\Omega}, \quad \text{as } t \rightarrow +\infty,$$

for some constant $K \in \mathbb{R}$ depending on H, u_0, φ and c , and such that $u_\infty + K \leq \varphi$ on $\partial\Omega$.

Proof: 1.- Case I: The nonexistence of a bounded viscosity solution for (S_0) can be proved by contradiction. If ϕ is a bounded viscosity solution to (S_0) , since $c > 0$ we have ϕ is an strict subsolution for the equation in Ω corresponding to (SC_c) . The function $u_\infty - 2(\|u_\infty\|_\infty + \|\phi\|_\infty) - 1$ is a solution for this last problem, and therefore, by comparison principle we arrive at $\phi \leq u_\infty - 2(\|u_\infty\|_\infty + \|\phi\|_\infty) - 1$ in $\bar{\Omega}$, which is a contradiction.

For the convergence, we start noting that the function $v(x, t) = u(x, t) + ct$ is a viscosity solution to the problem

$$\begin{cases} \partial_t v - \mathcal{I}(v) + H_0(x, Dv) &= c & \text{in } Q, \\ v &= \varphi + ct & \text{on } \partial Q, \\ v &= u_0 & \text{on } \bar{\Omega} \times \{0\}. \end{cases}$$

Now, considering the functions

$$\psi_{\pm}(x, t) = u_{\infty}(x, t) \pm (\|u_{\infty}\|_{L^{\infty}(\bar{\Omega})} + \|\varphi\|_{L^{\infty}(\partial\Omega)} + \|u_0\|_{L^{\infty}(\bar{\Omega})}),$$

we note that ψ_- is a viscosity subsolution for the above problem satisfying the boundary condition in the classical sense because $c > 0$, meanwhile ψ_+ is a viscosity supersolution for the same problem satisfying the boundary condition in the generalized sense. Thus, we obtain that v is uniformly bounded in \bar{Q} (by a constant depending on the data and $\|u_{\infty}\|_{L^{\infty}(\bar{\Omega})}$).

Now, considering v as a subsolution to the equation

$$(5.1) \quad \partial_t u - \mathcal{I}(u) + H_0(x, Du) = c \quad \text{in } Q,$$

and u_{∞} as a supersolution for the state-constraint problem associated to this equation. Defining

$$(5.2) \quad \kappa(t) = \max_{x \in \bar{\Omega}} \{v(x, t) - u_{\infty}(x)\}, \quad t \geq 0,$$

we can apply Lemma 4.4 to conclude that κ is decreasing in t . Since it is also bounded, we conclude that $\kappa(t) \rightarrow \bar{\kappa} \in \mathbb{R}$, as $t \rightarrow \infty$.

Considering $\{v_k\}_k$ the sequence of functions in $\text{Lip}(\bar{Q})$ given by $v_k(x, t) = v(x, t + t_k)$, we see that

$$|v_k(x, t) - v_k(x', t')| \leq C|x - x'| + (C + c)|t - t'|, \quad \text{for each } (x, t), (x', t') \in \bar{Q},$$

where $C > 0$ depends on the data and the smoothness of u_0 . Then, $\{v_k\}_k$ is equicontinuous and bounded, and by Arzela-Ascoli Theorem, up to subsequences, $v_k \rightarrow \bar{v} \in \text{Lip}(\bar{Q})$, uniformly in \bar{Q}_T for each $T > 0$. Stability results for viscosity solutions imply that the limit function \bar{v} satisfies (5.1). Evaluating (5.2) in $t + t_k$ and taking limit as $k \rightarrow \infty$ we arrive at

$$\bar{\kappa} = \max_{x \in \bar{\Omega}} \{\bar{v}(x, t) - u_{\infty}(x)\} \quad \text{for each } t \geq 0,$$

and since u_{∞} is a solution for the state-constraint problem associated to (5.1), by the application of the Strong Maximum Principle given by Proposition 4.5, we conclude that

$$\bar{v}(x, t) = u_{\infty}(x) + \bar{\kappa}, \quad \text{for all } (x, t) \in \bar{Q}.$$

This, together with the definition of \bar{v} and the nonincreasing property of κ allows us to conclude the result.

2.- *Case II:* Let $\phi \in C(\bar{\Omega})$ be a bounded solution to (S_0) . Since u_{∞} is a solution to (SC_0) it can be regarded as a viscosity supersolution to (S_0) with generalized boundary condition. Thus, the stationary version of the Strong Maximum Principle leads us to the existence of a constant $K \in \mathbb{R}$

such that $\phi = u_\infty + K$. But the stationary version of Lemma 3.3 implies that $\phi \leq \varphi$ on $\partial\Omega$, from which we conclude the result. The convergence proof follows the same lines of the previous case. \square

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GUY BARLES: LABORATOIRE DE MATHÉMATIQUES ET PHYSIQUE THÉORIQUE (UMR CNRS 7350), FÉDÉRATION DENIS POISSON (FR CNRS 2964), UNIVERSITÉ FRANÇOIS RABELAIS TOURS, PARC DE GRANDMONT, 37200 TOURS, FRANCE. Guy.Barles@lmpt.univ-tours.fr

ERWIN TOPP: DEPARTAMENTO DE INGENIERÍA MATEMÁTICA (UMI 2807 CNRS), UNIVERSIDAD DE CHILE, CASILLA 170, CORREO 3, SANTIAGO, CHILE. etopp@dim.uchile.cl