

# Characterization of Few Mode Fibers by OLCI Technique

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**Abstract :** All LP modes of a few mode fiber are simultaneously characterized using phase-sensitive optical low-coherence interferometry. The differential modal group delay and absolute chromatic dispersion values of each mode are retrieved from a single measurement without spatial mode transformers.

## Introduction

In recent years, advanced modulation formats with coherent detection have resulted in a remarkable capacity increase of optical communication links, enabling 100 Gbit/s and 400 Gbit/s per channel transmissions. Despite these significant advances, the capacity of optical fibers is expected to reach its limit unless other multiplexing strategies, such as mode division multiplexing (MDM) with multiple-input multiple-output (MIMO) digital signal processing, are considered [1]. Performance optimization of MDM transmissions requires management of channel parameters [3] for which precise knowledge of the LP mode characteristics, including their group delay (GD), dispersion (GVD) and mode-dependent loss, is particularly important.

Different methods have been proposed to characterize LP modes of few mode fibers (FMF). The well-known phase shift method is well suited for characterizing each mode individually. Other methods that can characterize FMF, without mode converters, include interferometric techniques with either RF or spectral measurements [5]. The main drawback with most of these approaches is that the measurement only extracts the dispersion of each mode relative to the fundamental mode. Moreover, these methods can fail if the group delays of the modes are too close.

In this paper, we show that the OLCI technique, in combination with a numerical method called "time-wavelength mapping" [7], can retrieve the mode properties of FMF without mode converters, irrespective of the relative power level of each mode and even if their group indices are close. We will compare these results with the well-known S<sup>2</sup> method [8].

## Experimental setup

The OLCI is a Mach-Zehnder interferometer, illuminated by a broadband source, with a

scanning reference mirror in one arm and the FMF under test in the other (Fig. 1). An Er<sup>3+</sup> superfluorescent source with a flat-top spectrum allows characterizing device-under-test (DUT) over the entire C+L band (i.e., 1525-1605nm). The interferogram is obtained by scanning the optical path difference between the two arms at a constant velocity (~1.0mm/s).

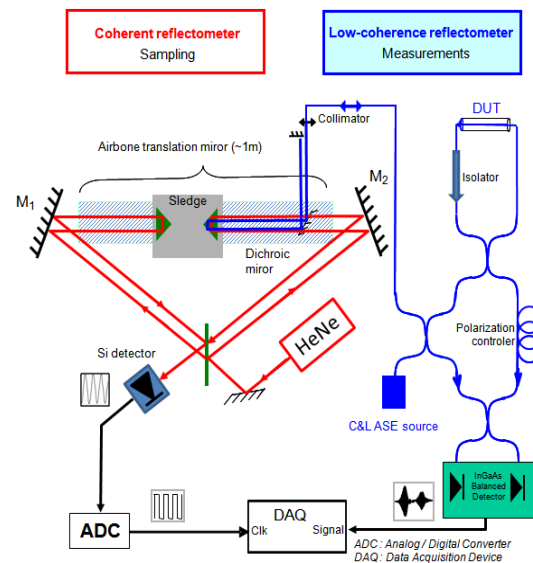
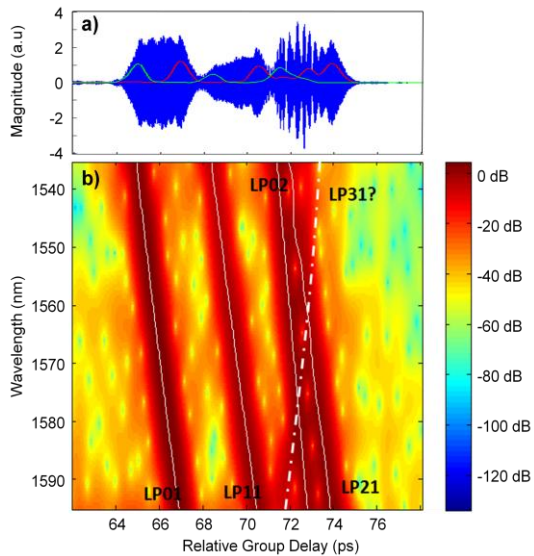


Fig. 1: OLCI Experimental setup

The interferometric signal is detected using a low-noise balanced photodetector. Interference only occurs when the optical path difference is less than the coherence length of the OLCI broadband source. In order to keep track of the absolute position of the reference mirror, the zero-crossing of a coherent interferometric signal at 632.8nm is used to trigger sampling of the optical path difference with a ~80nm period. The interferogram  $I(\tau)$  can be expressed as function of the optical path difference expressed in time domain  $\tau$  as:

$$I(\tau) \propto \frac{1}{4\pi} \text{Re} \left\{ \int_{-\infty}^{+\infty} S(\omega) \tilde{t}(\omega) \exp(-j\omega\tau) d\omega \right\} \quad (1)$$

where  $S(\omega)$  and  $\tilde{t}(\omega) = |\tilde{t}(\omega)|\exp[j\varphi(\omega)]$  are the power spectrum of the OLCI source and the complex transmission coefficient of the DUT, respectively. It is therefore possible to deduce the GD by extraction and derivation of the wavelength dependent phase function [9]. However, major difficulties can occur whenever the group indices of the different LP modes are sufficiently close to induce overlapping signals in the interferogram thus producing beating in the measured GD evolution, e.g. in Fig.2.a (blue trace) between 70-74ps relative GD. This beating makes identifying each LP mode on the interferogram difficult and sometimes impossible. To overcome modal overlapping issues, a numerical method called "time-wavelength mapping" has been proposed to extract GD properties of each mode independently [7]. Having access to the complete interferogram signal, we numerically apply a sliding-wavelength filter in the spectral domain. The Gaussian filter bandwidth (~few nm) is chosen as a trade-off between source coherence and expected DUT dispersion to optimize the spatial resolution. Finally, after inverse Fourier transform of the filtered signal, we obtain the wavelength dependent modulus of the interferogram. The result is the 3D graph shown in Fig. 2b. This method is more versatile than the one presented in [6] which need mode converter to treat each mode separately.



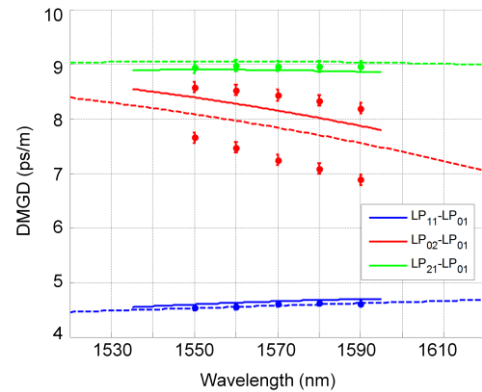
**Fig. 2:** a) Recorded interferogram (blue), interferogram after 10nm BW numerical filtering at 1535/1595nm (green/red). b) Corresponding time-wavelength mapping. White solid lines, GD fit of  $LP_{01}$ ,  $LP_{11}$ ,  $LP_{21}$  and  $LP_{02}$ ; dot-dash white line, possible  $LP_{31}$ .

## Results and comments

Fig. 2.a. presents two "cuts" of the interferogram of a 778mm of the characterized fiber reconstructed after numerical filtering with a 10nm 3dB-bandwidth (BW) filter centered at

1535nm (in green) and at 1595nm (in red). Note that a ~10nm filter allows maximizing resolution in terms of GD evolution.

Unlike previously published work [8], this time, the fiber is analyzed in a transmission configuration (OLCI). This was found to be necessary in order to remove partial mode coupling resulting from reflection at the output fiber end. This mode coupling multiplies the number of possible optical path per round trip and creates confusion in the data interpretation. The "time-wavelength mapping" of the recorded interferogram is shown in Fig. 2.b. It then becomes easy to numerically extract the position of each mode with respect to wavelength (white solid lines) and to deduce GD evolution. In the present case, we can clearly identify four LP modes:  $LP_{01}$ ,  $LP_{11}$ ,  $LP_{21}$  and  $LP_{02}$ . On this short length of fiber, we also observe a  $LP_{31}$  mode that is very close to cut-off in the simulations and is not analyzed further (dot-dash line in Fig. 2).



**Fig. 3:** Differential modal group delay of each mode with respect to  $LP_{01}$  mode. Solid lines: experiment. Dashed lines: simulations. Points with error-bars :  $S^2$  method.

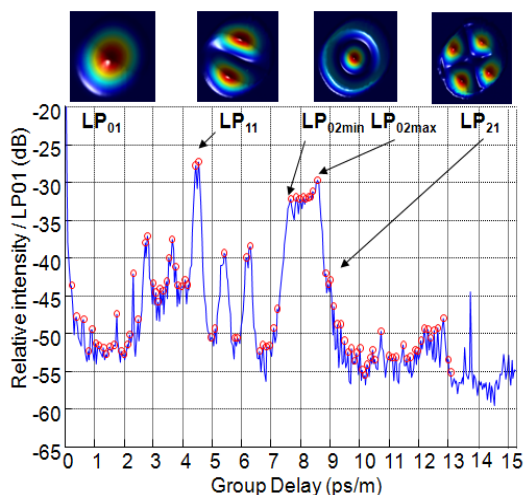
In Fig.3, we plot the corresponding differential modal group delay (DMGD) of each mode (solid lines) with respect to the  $LP_{01}$  mode (Fig. 3).

The interferogram is analyzed with reference to the modes predicted by numerical simulations performed with a commercial mode solver (Optifiber). Calculations are done using a step index profile ( $NA=0.17$ ,  $\phi_{core}=15 \mu m$ ) with parameters based on the fiber specifications and measurements done using a refracted near-field analyzer (Exfo NR-9200HR, refractive index resolution of  $1.10^{-4}$ , spatial resolution of  $0.1 \mu m$ ). The main benefit of the simulation is to provide a sound framework to identify the modes and we observe a generally good agreement between experiment (solid lines in Fig. 3) and simulations (dashed lines in Fig. 3).

To validate these measurements, we compare the extended OLCI with the well-known  $S^2$  method (points with error bars in Fig.3). The

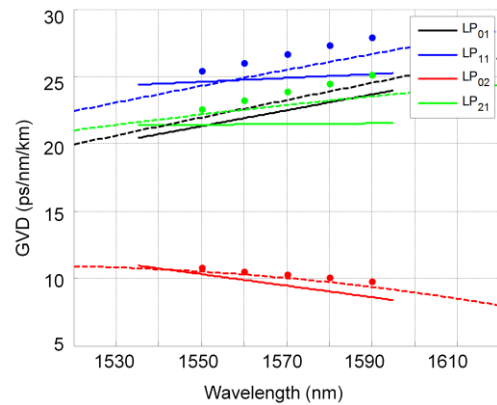
latter technique is based on a spatially resolved measurement of the spectral interference between coherent modes propagating with different group delays in the DUT [8]. The  $S^2$  measurement method gives the intensity and group delay of higher order modes with respect to the  $LP_{01}$  mode at a given central wavelength. This method requires that the power injected in the fundamental mode be dominant.

Fig.4 shows an example of the measurement centered at a wavelength of 1550nm. The resolution is 0.04ps/m. On this figure we can see that  $LP_{02}$  mode appears over a certain range of group delays referred to as  $LP_{02}$  min and  $LP_{02}$  max. DMGD of each mode are reported in Fig. 3 for five different wavelengths from 1550 to 1590nm. Note that for the  $S^2$  method, the group delay and dispersion of the fundamental mode are deduced by numerical simulations or needs to be measured by another technique. The results obtained with  $S^2$  technique confirm those obtained by OLCI.



**Fig. 4:** Intensity and group delay of each mode with respect to  $LP_{01}$  mode with associated transverse intensity profiles obtained with  $S^2$  method (wavelength: 1550 nm).

Fig. 5 presents the dispersion of each LP mode, calculated from group delay curves measured by OLCI. We see that three modes ( $LP_{01}$ ,  $LP_{11}$  and  $LP_{21}$ ) are dominated by material dispersion (20-25 ps/nm/km), which is not the case of the  $LP_{02}$  mode that presents a lower dispersion, around 10 ps/nm/km. It is important to note that the OLCI technique allows to extract absolute dispersion values of each mode whereas  $S^2$  technique gives access only to the value of relative dispersion compared to the fundamental mode, which is determined in this case by simulation [8]. These results are in agreement with each other and validate the OLCI technique. Also, wavelength-dependent losses



**Fig. 5:** Dispersion of each mode: solid lines: experiments, dashed lines: simulations. Points with error-bars :  $S^2$  method.

of each mode can be extracted but in this case, no significant variations were observed over the measured spectral range.

## Conclusion

We demonstrated that the OLCI method, in a transmission configuration, combined to numerical spectral filtering, is a powerful method to rapidly characterize few modes fibers over the whole C+L band. Absolute values of group delay and dispersion of each mode are obtained in a single scan without mode converters. These results are confirmed by  $S^2$  method and numerical simulations.

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