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L^2 -stability independent of diffusion for a finite element – finite volume discretization of a linear convection-diffusion equation.

Paul Deuring*, Robert Eymard[†] and Marcus Mildner*

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Abstract

We consider a time-dependent and a steady linear convection-diffusion equation. These equations are approximately solved by a combined finite element – finite volume method: the diffusion term is discretized by Crouzeix-Raviart piecewise linear finite elements on a triangular grid, and the convection term by upwind barycentric finite volumes. In the unsteady case, the implicit Euler method is used as time discretization. This scheme is shown to be unconditionally L^2 -stable, uniformly with respect to the diffusion coefficient.

AMS subject classifications. 65M30, 65M60, 76M10, 76M12.

Key words. convection-diffusion equation, combined finite element - finite volume method, Crouzeix-Raviart finite elements, barycentric finite volumes, upwind method, stability.

1. Introduction

We consider the convection-diffusion equation

$$\partial_t u - \nu \Delta u + \beta \cdot \nabla u = g \quad \text{in } \Omega \times (0, T), \quad (1.1)$$

supplemented by the initial and boundary conditions

$$u(x, 0) = u^{(0)}(x) \quad \text{for } x \in \Omega, \quad u|_{\partial\Omega \times (0, T)} = 0, \quad (1.2)$$

respectively. Here $\Omega \subset \mathbb{R}^2$ is a bounded open polygon with Lipschitz boundary, ν and T are positive reals, and $\beta \in H^1(\Omega)^2$, $u^{(0)} : \Omega \mapsto \mathbb{R}$ as well as $g : \Omega \times (0, T) \mapsto \mathbb{R}$ are given functions. We suppose that

$$\nabla \cdot \beta = 0, \quad -\beta \cdot \nabla \varphi \geq \underline{\beta} \quad \text{in } \Omega \quad (1.3)$$

for some function $\varphi \in C^1(\overline{\Omega})$ and some constant $\underline{\beta} > 0$. In the case where $\beta(x) = \beta_0$ for all $x \in \overline{\Omega}$ and for some $\beta_0 \in \mathbb{R}^2 \setminus \{0\}$, a suitable function φ is given by $\varphi(x) := -\beta_0 \cdot x$.

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We also consider the steady variant of problem (1.1), (1.2), that is,

$$-\nu \Delta U + \beta \cdot \nabla U = G \quad \text{in } \Omega, \quad (1.4)$$

$$U|_{\partial\Omega} = 0, \quad (1.5)$$

where $G : \Omega \mapsto \mathbb{R}$ is another given function. These problems are of particular interest in the convection-dominated regime, that is, if $\nu \ll |\beta|$, an interest that seems to be motivated by the belief that the preceding problems in the convection-dominated case show some affinity (although distant) with the Navier-Stokes system in the same regime. In this spirit, numerical schemes working well for that latter system are sometimes reduced to problem (1.1), (1.2) or (1.4), (1.5) so that they may be accessible to theoretical studies regarding stability or accuracy.

In the work at hand, we consider a discretization of (1.1), (1.2) and (1.4), (1.5), respectively, that is motivated in this way. This scheme may be described as follows: the diffusion term in (1.1) and (1.4) is discretized by piecewise linear Crouzeix-Raviart finite elements, and the convective term by an upwind finite volume method based on barycentric finite volumes on a triangular grid. Choosing an explicit time discretization, Feistauer e.a. [11, Section 7], [19, Chapter 4.4] tested this FE-FV method in the case of high-speed compressible Navier-Stokes flows in complex geometries and obtained very satisfactory results.

In [9], we applied this FE-FV method to problem (1.1), (1.2), using the implicit Euler method as time discretization. Under the assumption that β is constant, we showed that the approximate solution provided by this approach may be estimated in the L^2 -norm against the data, with the constant in this estimate being independent of the diffusion parameter ν . An analogous result was established with respect to problem (1.4), (1.5). However, we had to impose a condition on the grid ([9, (3.9)]) that goes beyond standard assumptions in a finite element or finite volume context, and which is not satisfied by certain structured regular meshes. Additional details in this respect may be found in [30].

It is the aim of the work at hand to establish the stability estimates from [9] without this inconvenient condition. In fact, we will show that even if the grid is only required to be shape-regular (minimum angle condition), an upper bound independent of ν and only involving the data may still be constructed for the L^2 -norm of the approximate solution of (1.1), (1.2) and (1.4), (1.5), respectively, obtained by the FE-FV approach described above, with the implicit Euler method being used as time discretization in the unsteady case. As an additional generalization with respect to [9], the function β need no longer be constant. Instead it will only be supposed to satisfy (1.3). For a detailed statement of our results, we refer to Theorem 2.1 and 2.2 below.

We remark that our approach also carries through if the advective velocity β in (1.1) and the associated function φ in (1.3) depend on time, provided that the spatial H^1 -norm of β is bounded uniformly with respect to time, and (1.3) holds at any time with the same constant $\underline{\beta}$. Moreover the diffusion term $\nu \Delta u$ may be replaced by an elliptic operator in divergence form $\overline{\nabla} \cdot (A \cdot \nabla u)$, where $A = A(x)$ or $A = A(x, t)$ is a symmetric matrix in $\mathbb{R}^{2 \times 2}$ which is positive definite uniformly with respect to x and (in the unsteady case) t . However, since our presentation would become more complicated while our proofs would remain without essential modification, we will not make these generalizations explicit.

Our theory should be expected to hold in the 3D case as well. Poincaré's inequality (2.5) is available in that case as well ([27], [3], [41]), so there should be no obstacle in this respect.

There is only one point that might need some effort, that is, to prove an analogue or find a replacement of equation (2.6) pertaining to the discrete L^2 -scalar product.

Condition (1.3)₁ on β , that is, $\nabla \cdot \beta = 0$, cannot be expected to be sufficient for the validity of our stability results, even if $\beta \in C^1(\overline{\Omega})^2$. This will be explained in Section 5. So inevitably an additional assumption has to be imposed on β . The one we have chosen – inequality (1.3)₂ – may look ad hoc. However, in Section 5 we will present a strong heuristic argument implying that (1.3)₂ is necessary for Theorem 2.1 and 2.2 to hold. This observation means in particular that our result is not valid in the purely diffusive case ($\beta = 0$), that is, for Poisson’s equation $\nu \Delta U = G$ in the steady and the heat equation in the evolutionary case. This is borne out by the counterexamples we present in Section 5.

Inequality (1.3)₂ yields that $\int_{\partial\Omega} (\beta \cdot n^{(\Omega)}) \varphi \, do_x \geq \underline{\beta} |\Omega| > 0$, where $n^{(\Omega)}$ denotes the outward unit normal to Ω . Thus the flux $\beta \cdot n^{(\Omega)}$ cannot be zero everywhere on $\partial\Omega$. Since $\int_{\partial\Omega} \beta \cdot n^{(\Omega)} \, do_x = 0$ by (1.3)₁, this means the advective flow β must enter the domain Ω at some places and leave it at others. In other words: condition (1.3)₂ indicates that the advective flow cannot be a purely rotational movement inside Ω .

Mathematically it is not surprising – at least in hindsight – that condition (1.3)₂ comes up. This is because we base our argument on some sort of coercivity property of the discrete bilinear form associated with the convection term (Corollary 3.1). The proof of this property is rather technical, but its main idea can easily be seen by looking at the continuous case. In fact, the assumptions in (1.3) yield $\int_{\Omega} (\beta \cdot v) \varphi v \, dx = - \int_{\Omega} (\beta \cdot \nabla \varphi) v^2 / 2 \, dx \geq \underline{\beta} \|v\|_2^2 / 2$ for $v \in H_0^1(\Omega)$.

The key argument allowing us to improve the theory from [9] consists of a weak BV estimate providing an upper bound for a certain weighted variation of our approximate solution. This upper bound involves the L^2 -norm of the solution and of the data. For technical reasons, the estimate in question will not appear explicitly, but will be split into two parts (Lemma 3.3 and (4.1) in the unsteady case, Lemma 3.3 and (4.11) in the steady one). Usually this type of inequality is exploited to pass to the limit in numerical schemes for the approximation of scalar hyperbolic equations; see [38] or [6] for example. Here, however, it is a critical part of the sequence of inequalities leading to our stability estimate.

Some additional remarks about existing literature are in order. A very large body of work deals with convection-diffusion-reaction equations, that is, with equations of the form

$$-\nu \Delta U + (\beta \cdot \nabla) U + \mu U = G \quad \text{in } \Omega \tag{1.6}$$

(often more shortly called convection-diffusion equations, too) and their time-dependent counterparts, under the assumption that

$$\mu - \nabla \cdot \beta / 2 \geq \gamma_0 \quad \text{for some } \gamma_0 > 0. \tag{1.7}$$

As is well known, in this situation various discretizations allow error and stability estimates with constants independent of ν . We refer to the monographs [33, Chapter 8 and 12], [12, Section 5.2.3, 5.4.4] and [35] for more details.

In the case of (1.1) and (1.4), however, the condition $\mu - \nabla \cdot \beta / 2 \geq \gamma_0$ reduces to the inequality $-\nabla \cdot \beta \geq \gamma_0$, which, in view of (1.3), is not satisfied. A possible strategy in this situation consists in using anisotropic meshes (“Shishkin meshes”; see [35]) that take account of boundary layers,

which, in turn, depend on ν . However, construction of such meshes requires information about the advective velocity β and becomes difficult if Ω does not exhibit a simple geometry. Moreover, for this type of mesh, the constant σ_0 in condition (2.1) relative to the grid shape is influenced by the diffusion coefficient. In view of these features, we think Shishkin's approach works best in situations that are not the main interest of the work at hand. So, when we compare our theory with what is available in existing literature, we in fact look for stability estimates for discretizations of scalar convection-diffusion equations with a structure similar to that of (1.1) or (1.4), with the discretizations in question based on conventional grids.

Let us give a brief survey of such results, starting with those that pertain to discretizations similar to the one considered here. Ohmori and Ushijima [32] treated the steady problem (1.4), (1.5) under assumptions on β that are weaker than ours. They discretized (1.4) by Crouzeix-Raviart finite elements, using an upwind version of this element in order to approximate the convective term. For constant β , this approach reduces to the FE-FV discretization of (1.4), (1.5) we study in the work at hand. The authors provided an upper bound for their discrete solutions, under the assumption that the right-hand side G in (1.4) is negative, and if the underlying triangulation is of acute type. Also an estimate of the H^1 -error is derived. The constant in this estimate depends on ν , which is inevitable for this type of error. Feistauer e.a. [1], [11] considered a scalar time-dependent nonlinear conservation law with a diffusion term. Discretizing this equation by the combined FE-FV scheme described above, with a rather general numerical flux adapted to the nonlinearity in question, and with a semi-implicit Euler method as time discretization, they established L^∞ -stability independent of ν . However, they required their triangulations to be of weakly acute type, and they supposed that the ratio of the time step to the grid size is small with respect to a quantity depending on the lifetime T of the solution, on the L^∞ -norm of the initial data and the L^∞ -norm of the right-hand side in the differential equation. In [1] the stability result in question is a key element in the proof of a convergence result, whereas in [11], it is applied for deriving $L^2(L^2)$ - and $L^2(H^1)$ -error estimates, which involve constants with exponential dependence on $1/\nu$. The articles [17], [18] and [21] establish results analogous to those in [1] and [11], but for a combined finite element - finite volume method involving piecewise linear conforming finite elements and dual finite volumes. Similar $L^2(L^2)$ - and $L^2(H^1)$ -error estimates as in [11], [18] and [21] are shown in [20], but pertaining to various discontinuous Galerkin discretizations.

Reference [14] deals with a degenerate time-dependent convection-diffusion equation which is much more general than (1.1) in some respects. A discretization is performed with a similar finite element - finite volume method as in the present article. The theory in [14] yields ν -independent L^∞ -bounds for the discrete solution, but covers (1.1) only in the case $g = 0$. This means in particular the approach in [14] does not provide an access to the steady problem (1.4). In [15], a convergence proof for a mixed FE-FV method is carried through for a similar problem as in [14]. The method in question involves conforming P1 finite elements and cell-centered finite volumes.

The stability estimates in [32], [1], [11], [17], [18], [21], [14] require rather severe assumptions because they are deduced from the maximum principle. In fact, even in the case of the continuous steady equation (1.4), this principle introduces a factor of the form $e^{diam\Omega\|\beta\|_\infty/\nu}$ if it is applied in order to yield L^∞ -bounds; see [24, Theorem 3.7]. The same factor arises for the parabolic problem (1.1) if $g \neq 0$ ([22, (3.3.10)]). So the assumptions in the references just mentioned are

needed in order to circumvent this phenomenon.

In [8] we derived L^2 - and H^1 -error estimates for the same discretization of (1.1) - (1.4) and (1.4), (1.5) as considered here. The constants in our estimates are analyzed with respect to their dependence on ν . This analysis, which includes the norms of the exact solution appearing in our error bounds, reveals a polynomial dependence on $\nu^{-1/2}$.

L^2 -stability and error bounds independent of ν are constructed in [34] for finite volume approximate solutions to (1.1). Article [34] is the most recent one in a series of papers dealing with ν -independent stability and error estimates, with previous publications treating Euler-Lagrangian and finite element methods. The corresponding references may be found in [34]. However, although constants independent of ν constitute the best possible case one might hope for, the type of theory developed in [34] and its predecessor papers also has drawbacks. It requires periodic boundary conditions instead of Dirichlet ones, it assumes Ω to be a rectangle and the grid to be structured according to this geometry, and it relies on Gronwall's inequality in a way which leads to exponential dependence of the stability and error bounds on the $W^{1,\infty}$ -norm of the convective velocity. (In the work at hand, our stability bounds depend linearly on the H^1 -norm of β .) Moreover, using Gronwall's inequality leaves open the question of how to deal with the steady case.

In [2], a mixed finite element method is applied to a unsteady nonlinear degenerate convection-diffusion equation. An argument based on Gronwall's inequality leads to a stability constant with exponential dependence on the reciprocal value of the smallest eigenvalue of the diffusion matrix. Li [29] studied the steady problem (1.6) with $\mu \geq 0$, but without assumption (1.7). The author supposed the domain Ω to be a rectangle and admitted mixed boundary conditions. He transformed equation (1.6) into a system of first-order partial differential equations and discretized them by an upwind finite element method based on Raviart-Thomas finite elements on a structured grid. H^1 -error estimates are derived, with constants depending on ν in an explicit way and via certain norms of the exact solution. Various least-squares finite element methods applied to (1.4) - again written as first-order system - were studied by Hsieh and Young [26], who established error bounds whose dependence on ν is explicit and also indirect through some norms of the exact solution. Buffa, Hughes, Sangalli [4] proved stability and error estimates for a stabilized discontinuous Galerkin method applied to (1.4). Their stability bounds are ν -uniform, and their error bounds depend on ν only insofar as certain norms of the exact solutions are concerned. However their estimates pertain to norms weighted by ν and β . Using norms weighted in a similar way, Verfürth [39] derived a posteriori error estimates for finite element discretizations of (1.6) without requiring (1.7). The unsteady case is treated in [40]. A lowest order discontinuous Petrov-Galerkin formulation with flux-upwind stabilization was considered by Causin, Sacco, Bottasso [5]. Since their discrete bilinear form associated with the convection term is positive, the authors obtained H^1 -stability. The related constant inevitably depends on ν . In addition, H^1 -error estimates are shown. A streamline-diffusion finite element method with a stability parameter incorporating the flow direction is proposed for (1.4) in [31]. An H^1 -error estimate is proved. A unsteady variant of (1.6) without condition (1.7) is discretized by a least-squares finite element method in [28]. Stability holds with a constant which, due to Gronwall's lemma, depends exponentially on the reciprocal value of the smallest eigenvalue of the diffusion matrix.

We mention that numerical tests of various FE-FV discretizations of compressible flows are reported in [7], [19, Chapter 4.4] and [25]. Of course, combined FE-FV methods are not only applied to scalar convection-diffusion equations, but also to other problems. As examples, we indicate that in [13], [23], combined FE-FV methods based on the Crouzeix-Raviart finite element yield approximate solutions to a nonlinear version of the steady compressible Stokes system. These articles present convergence results. Reference [37] uses the scheme from [32] described above, applying it to the 2D steady incompressible Navier-Stokes system. Error estimates are presented for the case of Reynolds number 1.

2. Notation. FE-FV discretization of (1.1), (1.2) and (1.4), (1.5), respectively. Statement of main results.

If $\epsilon > 0$ and $x \in \mathbb{R}^2$, we write $B_\epsilon(x)$ for the open disk in \mathbb{R}^2 centered in x and with radius ϵ . If $U \subset \mathbb{R}^2$, we put $\text{diam}(U) := \sup\{|x - y| : x, y \in U\}$. We write $|U|$ for the measure of a measurable set $U \subset \mathbb{R}^2$. For $u, v \in L^2(\Omega)$, we use the abbreviation $(u, v) := \int_\Omega uv \, dx$. We write $\|\cdot\|_2$ for the standard norm of $L^2(\Omega)$. The Sobolev space $H^1(\Omega)$ and its subspace $H_0^1(\Omega)$ are defined in the usual way. The standard norm of $H^1(\Omega)$ is denoted by $\|\cdot\|_{1,2}$.

As already indicated in Section 1, we assume that $\Omega \subset \mathbb{R}^2$ is a bounded open polygon with Lipschitz boundary. The functions β, φ, g, u_0 and G and the constants $\nu, T, \underline{\beta} \in (0, \infty)$ were also introduced in Section 1, with $\beta \in H^1(\Omega)^2$, $\varphi \in C^1(\overline{\Omega})$ being such that (1.3) is satisfied. We assume that $g \in C^0([0, T], H^1(\Omega))$ and $u^{(0)}, G \in H^1(\Omega)$. The functions $g(\cdot, t)$ and G are required to belong to $H^1(\Omega)$ instead of only to $L^2(\Omega)$ so that they admit traces on edges, in view of an interpolation operator we will introduce below.

By adding a constant to φ , we may suppose without loss of generality that $\varphi(x) \geq \varphi_0$ ($x \in \overline{\Omega}$) for some $\varphi_0 > 0$. For example, in the case $\beta = \beta_0$ for some $\beta_0 \in \mathbb{R}^2 \setminus \{0\}$, we may put $\varphi(x) := 2|\beta_0| \text{diam}(\Omega) - \beta_0 \cdot (x - x_0)$, where x_0 is an arbitrary but fixed point in Ω . Obviously there is a constant $\varphi_1 > 0$ with $\varphi(x) \leq \varphi_1$ and $|\nabla\varphi(x)| \leq \varphi_1$ for $x \in \overline{\Omega}$. We further introduce a parameter $\sigma_0 \in (0, 1)$, which will appear in condition (2.1) below. The set Ω , the functions β, φ, g, u_0 and G as well as the numbers $\nu, T, \varphi_0, \varphi_1, \underline{\beta}$ and σ_0 will be kept fixed throughout.

By the symbol \mathfrak{C} , we denote constants that may depend on $\sigma_0, \text{diam} \Omega, \underline{\beta}, \varphi_0$ and φ_1 , with polynomial dependence on $\underline{\beta}^{-1}, \varphi_0$ and φ_1 . This last feature is important because we want to control how our estimates are influenced by β . This influence not only manifests itself by the factor $1 + \|\beta\|_{1,2}$ appearing in Theorem 2.1 and 2.2 and their proof, but also via the three quantities $\underline{\beta}, \varphi_0$ and φ_1 .

We consider triangulations \mathfrak{T} of Ω with the following three properties: Firstly, \mathfrak{T} is a finite set of open triangles $K \subset \mathbb{R}^2$ with $\overline{\Omega} = \cup\{\overline{K} : K \in \mathfrak{T}\}$. Secondly, if $K_1, K_2 \in \mathfrak{T}$ with $\overline{K_1} \cap \overline{K_2} \neq \emptyset$ and $K_1 \neq K_2$, then $\overline{K_1} \cap \overline{K_2}$ is a common vertex or a common side of K_1 and K_2 . And thirdly, for any $K \in \mathfrak{T}$, the relation

$$B_{\sigma_0 \text{diam} K}(x) \subset K \tag{2.1}$$

is valid for some $x \in K$. All estimates appearing in the following involve constants that do not depend on the grid except via the parameter σ_0 in (2.1). Therefore we may simplify our

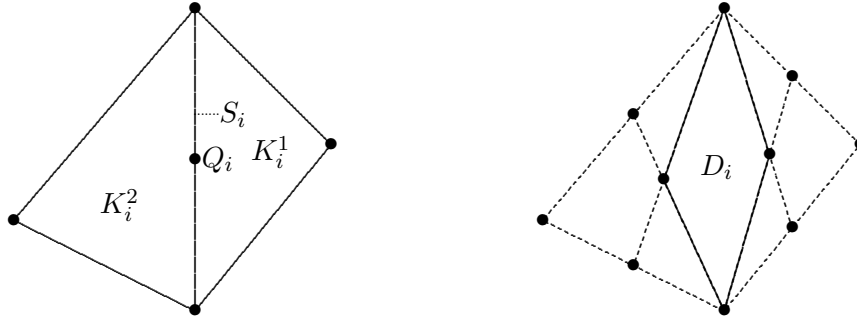


Figure 1: A quadrangle $\overline{K_i^1} \cup \overline{K_i^2}$ (left) and the quadrangle D_i inside $\overline{K_i^1} \cup \overline{K_i^2}$ (right).

presentation by fixing \mathfrak{T} already at this point. So let \mathfrak{T} be given such that the preceding assumptions, in particular (2.1), are verified. Below, when we introduce some spaces, norms and functions related to \mathfrak{T} , we will indicate this relation by an index h . This parameter h may be considered as a quantity related to the size of the triangles $K \in \mathfrak{T}$, for example $h := \max\{\text{diam } K : K \in \mathfrak{T}\}$, but it will not play any role in our theory. As a consequence of (2.1), we have

$$(\text{diam } K)^2 \leq \mathfrak{C}|K| \quad \text{for } K \in \mathfrak{T}. \quad (2.2)$$

Let \mathfrak{S} be the set of the sides of the triangles $K \in \mathfrak{T}$. Put $J := \{1, \dots, \#\mathfrak{S}\}$, where $\#\mathfrak{S}$ denotes the number of elements of \mathfrak{S} . Let $(S_i)_{i \in J}$ be a numbering of \mathfrak{S} , and denote the midpoint of S_i by Q_i ($i \in J$). Set $J^\circ := \{i \in J : Q_i \in \Omega\}$, so that $J \setminus J^\circ = \{i \in J : Q_i \in \partial\Omega\}$. Note that for $i \in J \setminus J^\circ$, we have $S_i \subset \partial\Omega$.

We further introduce a barycentric mesh $(D_i)_{i \in J}$ on the triangular grid \mathfrak{T} : If $i \in J^\circ$, there are two triangles in \mathfrak{T} , denoted by K_i^1, K_i^2 , such that $\overline{K_i^1} \cap \overline{K_i^2} = S_i$. We join the barycenter of each of these triangles with the endpoints of S_i . In this way we obtain a closed quadrilateral containing S_i (Fig. 1). This quadrilateral is denoted by D_i . If $i \in J \setminus J^\circ$ (hence $Q_i \in \partial\Omega$), let D_i be the closed triangle whose sides are the segment S_i and the segments joining the endpoints of S_i with the barycenter of the (unique) triangle $K \in \mathfrak{T}$ with $S_i \subset \overline{K}$. If $i, j \in J$ with $i \neq j$ are such that the set $D_i \cap D_j$ contains more than one point, then this set is a common side of D_i and D_j . In this case, the quadrilaterals D_i and D_j are called ‘‘adjacent’’, and their common side is denoted Γ_{ij} . For $i \in J$, we set

$$s(i) := \{j \in J \setminus \{i\} : D_i \text{ and } D_j \text{ are adjacent}\}.$$

If $i \in J$ and $j \in s(i)$, let n_{ij} denote the outward unit normal to D_i on Γ_{ij} . This means that n_{ij} points from D_i into D_j . We will use the abbreviation

$$\Theta_{ij}^+ := \int_{\Gamma_{ij}} \max\{\beta(x) \cdot n_{ij}, 0\} \, do_x \quad (i \in J, j \in s(i)).$$

Since $\nabla \cdot \beta = 0$ (see (1.3)) and $n_{ij} = -n_{ji}$ ($i \in J, j \in s(i)$), we get

$$\sum_{j \in s(i)} (\Theta_{ij}^+ - \Theta_{ji}^+) = \sum_{j \in s(i)} \int_{\Gamma_{ij}} \beta(x) \cdot n_{ij} \, do_x = 0 \quad (i \in J^\circ). \quad (2.3)$$

Due to relation (2.1) and because K_i^1 and K_i^2 have a common side, we obtain

$$\text{diam } K_i^\nu \leq \mathfrak{C} \text{diam } K_i^\mu \quad \text{for } i \in J^o, \nu, \mu \in \{1, 2\}. \quad (2.4)$$

Next we introduce two finite element spaces by setting

$$\begin{aligned} X_h &:= \{v \in L^2(\Omega) : v|_K \in P_1(K) \text{ for } K \in \mathfrak{T}_h, v \text{ continuous at } Q_i \text{ for } i \in J\}, \\ V_h &:= \{v_h \in X_h : v_h(Q_i) = 0 \text{ for } i \in J \setminus J^o\}, \end{aligned}$$

where $P_1(A)$, for $A \subset \mathbb{R}^2$, denotes the set of all polynomials of degree at most 1 over A . The spaces X_h and V_h are nonconforming finite element spaces based on the piecewise linear Crouzeix-Raviart finite element. For $i \in J$, let w_i be the function from X_h that is uniquely determined by the requirement that $w_i(Q_j) = \delta_{ij}$ for $j \in J$. The family $(w_i)_{i \in J}$ is a basis of X_h . For $v_h, w_h \in X_h$, we put

$$((v_h, w_h))_{X_h} := \sum_{K \in \mathfrak{T}} \int_K \nabla(v_h|_K) \cdot \nabla(w_h|_K) dx, \quad \|v_h\|_{X_h} := \left(((v_h, v_h))_{X_h} \right)^{1/2}.$$

Obviously

$$((v_h, v_h))_{X_h} = \|v_h\|_{X_h}^2, \quad |((v_h, w_h))_{X_h}| \leq \|v_h\|_{X_h} \|w_h\|_{X_h} \quad \text{for } v_h, w_h \in X_h.$$

The discrete Poincaré inequality

$$\|v_h\|_2 \leq \mathfrak{C} \|v_h\|_{X_h} \quad \text{for } v_h \in V_h \quad (2.5)$$

was proved in [10], [27], [3] and [41], with a detailed analysis of the constant in [41].

From [1, (3.29), (3.31), (3.33)], we take two formulas for the L^2 -scalar product of $v_h, w_h \in X_h$, that is,

$$(v_h, w_h) = (1/3) \sum_{K \in \mathfrak{T}} |K| \sum_{r=1}^3 v_h(Q_K^r) w_h(Q_K^r), \quad (v_h, w_h) = \sum_{i \in J} v_h(Q_i) w_h(Q_i) |D_i|, \quad (2.6)$$

where $Q_K^r \in \mathbb{R}^2$ for $r \in \{1, 2, 3\}$ are the three vertices of $K \in \mathfrak{T}$. Put $H^1(\Omega) \oplus X_h := \{v + w_h : v \in H^1(\Omega), w_h \in X_h\}$, and let $I_h : H^1(\Omega) \oplus X_h \mapsto X_h$ be the interpolation operator introduced in [16, 8.9.79]; it is defined by

$$I_h(v) := \sum_{i \in J} l_i^{-1} \int_{S_i} v(x) do_x w_i \quad \text{for } v \in H^1(\Omega) \oplus X_h,$$

where l_i denotes the length of S_i ($i \in J$). Note that a function $v \in H^1(\Omega)$ admits a trace on S_i for $i \in J$, and a function $v_h \in X_h$ verifies the equation

$$\int_{S_i} E(v_h|_{K_i^1}) do_x = l_i v_h(Q_i) = \int_{S_i} E(v_h|_{K_i^2}) do_x \quad (i \in J^o),$$

where $E(v_h|K_i^s)$ denotes the continuous extension of $v_h|K_i^s$ to S_i ($s \in \{1, 2\}$). Thus the operator I_h is well defined. By [16, Lemma 8.9.81], it satisfies the estimate

$$\|I_h(v)\|_2 \leq \mathfrak{C}\|v\|_{1,2} \quad \text{for } v \in H^1(\Omega).$$

It will be useful to introduce another interpolation operator besides I_h . In fact, for $v \in L^2(\Omega)$ with $v|K \in C^0(K)$ for $K \in \mathfrak{T}_h$, v continuous at Q_i for $i \in J$, we set $\varrho_h(v) := \sum_{i \in J} v(Q_i)w_i$. Next we define a discrete convection term b_h , which is to approximate the variational form

$$b(v, w) := \int_{\Omega} (\beta \cdot \nabla) v w \, dx \quad (v, w \in H^1(\Omega))$$

associated with the convection term $\beta \cdot \nabla u$ in (1.1) and (1.4). We put

$$b_h(v_h, w_h) := \sum_{i \in J} w_h(Q_i) \sum_{j \in s(i)} \left(\Theta_{ij}^+ v_h(Q_i) - \Theta_{ji}^+ v_h(Q_j) \right) \quad \text{for } v_h, w_h \in X_h.$$

This definition means that we discretize b by an upwind finite volume method on the barycentric grid $(D_i)_{i \in J}$. A motivation for this discretization may be found in [11, p. 308], where a general numerical flux H is considered instead of the upwind method we use here. If β is constant, our upwind method corresponds to a choice of H given by

$$H(a, b, n) = a \text{ if } n \cdot \beta > 0, \quad H(a, b, n) = b \text{ if } n \cdot \beta \leq 0, \quad \text{for } a, b \in \mathbb{R}, n \in \mathbb{R}^2 \text{ with } |n| = 1.$$

The approach here and the one in [11], as any upwind method, is based on the idea of taking into account the direction from which the information arrives. In view of discretizing the time variable, we fix $N \in \mathbb{N}$ and choose $t_1, \dots, t_N \in (0, T)$ with $t_1 < \dots < t_N$. Put $t_0 := 0$, $t_{N+1} := T$, $\tau_k := t_k - t_{k-1}$ for $1 \leq k \leq N+1$. For brevity, we put $G_h := I_h(G)$ and introduce functions $g_h^{(k)} : \bar{\Omega} \mapsto \mathbb{R}$ by setting

$$g_h^{(k)}(x) := I_h(g(\cdot, t_k))(x) \quad \text{for } k \in \{0, \dots, N+1\}, x \in \bar{\Omega}.$$

Now we are in a position to introduce the finite element – finite volume discretization of problem (1.1), (1.2) and (1.4), (1.5), respectively, that we want to study in the work at hand. Concerning (1.1), (1.2), we consider functions $u_h^{(0)}, \dots, u_h^{(N+1)} \in V_h$ with

$$\begin{aligned} \tau_k^{-1}(u_h^{(k+1)} - u_h^{(k)}, v_h) + \nu((u_h^{(k+1)}, v_h))_{X_h} + b_h(u_h^{(k+1)}, v_h) &= (g_h^{(k+1)}, v_h) \\ \text{for } v_h \in V_h, k \in \{0, \dots, N\}, \quad u_h^{(0)} &= I_h(u^{(0)}). \end{aligned} \quad (2.7)$$

This scheme is implicit because both the diffusion and the convection term are discretized implicitly. For the steady problem (1.4), (1.5), we consider an approximate solution $U_h \in V_h$ satisfying

$$\nu((U_h, v_h))_{X_h} + b_h(U_h, v_h) = (G_h, v_h) \quad \text{for } v_h \in V_h. \quad (2.8)$$

In view of (2.5), and because $b_h(v_h, v_h) \geq 0$ for $v_h \in V_h$ (Lemma 3.3) and $\dim V_h < \infty$, both problems admit a unique solution. Our main results may now be stated as follows.

Theorem 2.1 Let $u_h^{(0)}, \dots, u_h^{(N+1)} \in V_h$ be a system of functions satisfying (2.7). Then

$$\begin{aligned} & \left(\sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_2^2 \right)^{1/2} + \max_{1 \leq l \leq N+1} \|u_h^{(l)}\|_2 + \nu^{1/2} \left(\sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_{X_h}^2 \right)^{1/2} \\ & \leq \mathfrak{C}(1 + \|\beta\|_{1,2}) \left[\left(\sum_{l=1}^{N+1} \tau_l \|g_h^{(l)}\|_2^2 \right)^{1/2} + \|u_h^{(0)}\|_2 \right]. \end{aligned} \quad (2.9)$$

Theorem 2.2 Let $U_h \in V_h$ be the solution of (2.8). Then

$$\|U_h\|_2 + \nu^{1/2} \|U_h\|_{X_h} \leq \mathfrak{C}(1 + \|\beta\|_{1,2}) \|G_h\|_2,$$

3. Auxiliary results.

We begin with a simple observation.

Lemma 3.1 Let $a_{ij} \in \mathbb{R}$ for $i, j \in J$. Then $\sum_{i \in J} \sum_{j \in s(i)} a_{ij} = \sum_{i \in J} \sum_{j \in s(i)} a_{ji}$.

Proof: For $i, j \in J$, we have $j \in s(i)$ if and only if $i \in s(j)$. As a consequence, we obtain $\{(i, j) \in J^2 : j \in s(i)\} = \{(i, j) \in J^2 : i \in s(j)\}$. Thus the left-hand side of the equation in Lemma 3.1 equals $\sum_{j \in J} \sum_{i \in s(j)} a_{ij}$. Lemma 3.1 now follows by renaming indices. \square

The ensuing lemma deals with the interpolation operator ϱ_h .

Lemma 3.2 Let $v_h, w_h \in X_h$. Then

$$\begin{aligned} \|v_h\|_2^2 & \leq \mathfrak{C}(v_h, \varrho_h(v_h \varphi)), \quad (v_h, \varrho_h(v_h \varphi)) \leq \mathfrak{C}\|v_h\|_2^2, \\ (v_h, \varrho_h(w_h \varphi)) & \leq (v_h, \varrho_h(v_h \varphi))^{1/2} (w_h, \varrho_h(w_h \varphi))^{1/2}, \\ \|\varrho_h(v_h \varphi)\|_{X_h} & \leq \mathfrak{C}\|v_h\|_{X_h}. \end{aligned}$$

Proof: The first three inequalities stated in Lemma 3.2 are an immediate consequence of (2.6), the relation $0 < \varphi_0 \leq \varphi(x) \leq \varphi_1$ ($x \in \Omega$) and – as concerns the third estimate – of the Cauchy-Schwarz inequality.

In order to prove the fourth, take $K \in \mathfrak{T}$. There are three elements $i_1, i_2, i_3 \in J$ such that the points Q_{i_r} with $r \in \{1, 2, 3\}$ are the three vertices of K . Then $\varrho_h(v_h \varphi)|_K = \sum_{r=1}^3 (v_h \varphi)(Q_{i_r}) w_{i_r}|_K$ by the definition of $\varrho_h(v_h \varphi)$. As a consequence

$$\|\nabla(\varrho_h(v_h \varphi)|_K)\|_2^2 = \int_K \left| \sum_{r=1}^3 (v_h \varphi)(Q_{i_r}) \nabla(w_{i_r}|_K) \right|^2 dx.$$

An estimate as in the proof of [9, Lemma 3.1] now yields that $\|\nabla(\varrho_h(v_h \varphi)|_K)\|_2^2$ is bounded by

$$\begin{aligned} & 2 \int_K \left| \sum_{r=1}^3 v_h(Q_{i_r}) \int_0^1 \left(\nabla \varphi(Q_{i_1} + \vartheta(Q_{i_r} - Q_{i_1})) (Q_{i_r} - Q_{i_1}) \right) d\vartheta \nabla(w_{i_r}|_K) \right|^2 dx \\ & + 2\varphi(Q_{i_1})^2 \|\nabla(v_h|_K)\|_2^2. \end{aligned}$$

Therefore we obtain the upper bound

$$\mathfrak{C} \|\nabla \varphi\|_\infty^2 \sum_{r=1}^3 v_h(Q_{i_r})^2 (\text{diam } K)^2 \int_K |\nabla(w_{i_r}|K)|^2 dx + 2 \|\varphi\|_\infty^2 \|\nabla(v_h|K)\|_2^2$$

for $\|\nabla(\varrho_h(v_h \varphi)|K)\|_2^2$. But in view of (2.2) $|\nabla(w_{i_r}|K)| \leq \mathfrak{C}(\text{diam } K)^{-1}$, so by (2.6) and because $\|\nabla \varphi\|_\infty \leq \varphi_1$, $\|\varphi\|_\infty \leq \varphi_1$,

$$\|\nabla(\varrho_h(v_h \varphi)|K)\|_2^2 \leq \mathfrak{C}(\|v_h|K\|_2^2 + \|\nabla(v_h|K)\|_2^2).$$

The fourth estimate in Lemma 3.2 now follows with (2.5). \square

Next we prove a basic fact regarding the discrete convection term b_h . The proof in question, as well as the proof of Lemma 3.4 below, applies techniques similar to those used by Sacco, Saleri [36].

Lemma 3.3 *For $v_h \in V_h$, the equation $b_h(v_h, v_h) = \mathfrak{K}_h/2$ holds, where*

$$\mathfrak{K}_h := \mathfrak{K}_h(v_h) := \sum_{i \in J} \sum_{j \in s(i)} \Theta_{ji}^+ (v_h(Q_i) - v_h(Q_j))^2.$$

In particular $b_h(v_h, v_h) \geq 0$.

Proof: By (2.3) and because $v_h(Q_i) = 0$ for $i \in J \setminus J^o$, we can write

$$b_h(v_h, v_h) = \sum_{i \in J} v_h(Q_i) \sum_{j \in s(i)} (v_h(Q_i) - v_h(Q_j)) \Theta_{ji}^+.$$

Using the equation $v_h(Q_i)(v_h(Q_i) - v_h(Q_j)) = \frac{1}{2}v_h(Q_i)^2 + \frac{1}{2}(v_h(Q_i) - v_h(Q_j))^2 - \frac{1}{2}v_h(Q_j)^2$, we get

$$b_h(v_h, v_h) = \frac{1}{2} \sum_{i \in J} \sum_{j \in s(i)} (v_h(Q_i) - v_h(Q_j))^2 \Theta_{ji}^+ + \frac{1}{2} \sum_{i \in J} \sum_{j \in s(i)} (v_h(Q_i)^2 - v_h(Q_j)^2) \Theta_{ji}^+.$$

On the other hand, by (2.3), the relation $v_h(Q_i) = 0$ for $i \in J \setminus J^o$ and Lemma 3.1, we obtain

$$\sum_{i \in J} \sum_{j \in s(i)} (v_h(Q_i)^2 - v_h(Q_j)^2) \Theta_{ji}^+ = \sum_{i \in J} \sum_{j \in s(i)} (\Theta_{ij}^+ v_h(Q_i)^2 - \Theta_{ji}^+ v_h(Q_j)^2) = 0.$$

This concludes the proof. \square

In the remainder of this section, we prove some sort of coercivity inequality for b_h (Corollary 3.1). We begin by establishing a lower bound for $b_h(v_h, \varrho_h(v_h \varphi))$.

Lemma 3.4 *Let $v_h \in V_h$. Then $b_h(v_h, \varrho_h(v_h \varphi)) \geq A_h/2$, with*

$$A_h := A_h(v_h) := \sum_{j \in J} v_h(Q_j)^2 \sum_{i \in s(j)} (\Theta_{ji}^+ \varphi(Q_i) - \Theta_{ij}^+ \varphi(Q_j)).$$

Proof: By (2.3) and the equation $v_h(Q_i) = 0$ for $i \in J \setminus J^o$,

$$b_h(v_h, \varrho_h(v_h \varphi)) = \sum_{i \in J} v_h(Q_i) \varphi(Q_i) \sum_{j \in s(i)} (v_h(Q_i) - v_h(Q_j)) \Theta_{ji}^+.$$

Thus, with an argument as in the proof of Lemma 3.3,

$$\begin{aligned} b_h(v_h, \varrho_h(v_h \varphi)) &= \frac{1}{2} \sum_{i \in J} \varphi(Q_i) \sum_{j \in s(i)} (v_h(Q_i) - v_h(Q_j))^2 \Theta_{ji}^+ \\ &+ \frac{1}{2} \sum_{i \in J} \varphi(Q_i) \sum_{j \in s(i)} (v_h(Q_i)^2 - v_h(Q_j)^2) \Theta_{ji}^+. \end{aligned}$$

The first term in the above right hand side is non-negative, and the second equals $A_h/2$. In fact, by (2.3), the equation $v_h(Q_i) = 0$ for $i \in J \setminus J^o$ and Lemma 3.1,

$$\sum_{i \in J} \varphi(Q_i) \sum_{j \in s(i)} (v_h(Q_i)^2 - v_h(Q_j)^2) \Theta_{ji}^+ = \sum_{i \in J} \varphi(Q_i) \sum_{j \in s(i)} (v_h(Q_i)^2 \Theta_{ij}^+ - v_h(Q_j)^2 \Theta_{ji}^+) = A_h.$$

This concludes the proof. \square

Lemma 3.5 *Let $v_h \in V_h$, and put*

$$B_h := B_h(v_h) := - \sum_{i \in J} v_h(Q_i)^2 \sum_{j \in s(i)} \int_{\Gamma_{ij}} \varphi(x) \beta(x) \cdot n_{ij} \, d\sigma_x.$$

Then $B_h \geq \underline{\beta} \|v_h\|_2^2$.

Proof: By the Divergence theorem and (1.3),

$$- \sum_{j \in s(i)} \int_{\Gamma_{ij}} \varphi(x) \beta(x) \cdot n_{ij} \, d\sigma_x = - \int_{D_i} \beta(x) \cdot \nabla \varphi(x) \, dx \geq |D_i| \underline{\beta} \quad \text{for } i \in J^o.$$

The lemma follows with (2.6) and the equation $v_h(Q_i) = 0$ for $i \in J \setminus J^o$. \square

Lemma 3.6 *Let $v_h \in V_h$. Then*

$$B_h \leq 2b_h(v_h, \varrho_h(v_h \varphi)) + \mathfrak{C} \|\beta\|_{1,2}^{1/2} \|v_h\|_2 \mathfrak{K}_h^{1/2},$$

where $B_h = B_h(v_h)$ and $\mathfrak{K}_h = \mathfrak{K}_h(v_h)$ were introduced in Lemma 3.5 and 3.3, respectively.

Proof: By Lemma 3.4, we have

$$B_h = B_h - A_h + A_h \leq |A_h - B_h| + 2b_h(v_h, \varrho_h(v_h \varphi)), \quad (3.1)$$

with A_h from that lemma. On the other hand, since $n_{ij} = -n_{ji}$ ($i \in J$, $j \in s(i)$),

$$\begin{aligned} |A_h - B_h| &= \left| \sum_{i \in J} v_h(Q_i)^2 \sum_{j \in s(i)} \left[\int_{\Gamma_{ij}} (\varphi(Q_i) - \varphi(x)) \max\{\beta(x) \cdot n_{ji}, 0\} \, d\sigma_x \right. \right. \\ &\quad \left. \left. - \int_{\Gamma_{ij}} (\varphi(Q_j) - \varphi(x)) \max\{\beta(x) \cdot n_{ij}, 0\} \, d\sigma_x \right] \right| \\ &= \left| \sum_{i \in J} \sum_{j \in s(i)} (v_h(Q_i)^2 - v_h(Q_j)^2) \int_{\Gamma_{ij}} (\varphi(Q_i) - \varphi(x)) \max\{\beta(x) \cdot n_{ji}, 0\} \, d\sigma_x \right|, \end{aligned} \quad (3.2)$$

where the last equation follows from Lemma 3.1. For $i \in J$, $j \in s(i)$ and $x \in \Gamma_{ij}$, we find with the relation $|\nabla\varphi(x)| \leq \varphi_1$ ($x \in \bar{\Omega}$) that

$$|\varphi(Q_i) - \varphi(x)| \leq \varphi_1 \max_{\nu \in \{1,2\}} \text{diam } K_i^\nu.$$

Here and in the following, if $i \in J \setminus J^o$, we use both the notation K_i^1 and K_i^2 for the unique triangle $K \in \mathfrak{T}$ with $Q_i \in K$. Now we may deduce from (3.2) that

$$|A_h - B_h| \leq \varphi_1 \sum_{i \in J} \sum_{j \in s(i)} \Theta_{ji}^+ \max_{\nu \in \{1,2\}} \text{diam } K_i^\nu |v_h(Q_j) + v_h(Q_i)| |v_h(Q_j) - v_h(Q_i)|,$$

hence by the Cauchy-Schwarz inequality,

$$|A_h - B_h| \leq \varphi_1 \left(\sum_{i \in J} \sum_{j \in s(i)} \Theta_{ji}^+ \left(\max_{\nu \in \{1,2\}} \text{diam } K_i^\nu \right)^2 (v_h(Q_j) + v_h(Q_i))^2 \right)^{1/2} \mathfrak{K}_h^{1/2}.$$

For $i \in J$, $j \in s(i)$, there is $\nu \in \{1, 2\}$ with $\Gamma_{ji} \subset K_j^\nu$, so

$$\Theta_{ji}^+ \leq \int_{\Gamma_{ji}} |\beta(x)| dx \leq (\text{diam } K_j^\nu)^{1/2} \|\beta|_{\Gamma_{ji}}\|_2 \leq \mathfrak{C} \|\beta|_{D_j \cap K_j^\nu}\|_{1,2} \leq \mathfrak{C} \cdot \|\beta\|_{1,2},$$

where the second last inequality follows by a change of the domain of integration from $D_j \cap K_j^\nu$ to the unit triangle $\{x \in [0, 1]^2 : x_1 + x_2 \leq 1\}$, an application of a standard trace inequality on that triangle, and then another change of the domain of integration back to $D_j \cap K_j^\nu$. Concerning that last transformation, note that $(\text{diam } K_j^\nu)^2 \leq \mathfrak{C} |K_j^\nu|$ by (2.2) and $|K_j^\nu|/3 = |D_j \cap K_j^\nu|$. As a consequence, we arrive at the estimate

$$|A_h - B_h| \leq \mathfrak{C} \|\beta\|_{1,2}^{1/2} \left(\sum_{i \in J} \sum_{j \in s(i)} \left(\max_{\nu \in \{1,2\}} \text{diam } K_i^\nu \right)^2 (v_h(Q_j)^2 + v_h(Q_i)^2) \right)^{1/2} \mathfrak{K}_h^{1/2}. \quad (3.3)$$

On the other hand, with (2.4),

$$\begin{aligned} & \sum_{i \in J} \sum_{j \in s(i)} \left(\max_{\nu \in \{1,2\}} \text{diam } K_i^\nu \right)^2 (v_h(Q_j)^2 + v_h(Q_i)^2) \\ & \leq \sum_{i \in J} \sum_{j \in s(i)} \left((\text{diam } K_j^1)^2 v_h(Q_j)^2 + (\text{diam } K_i^1)^2 v_h(Q_i)^2 \right). \end{aligned}$$

Therefore by Lemma 3.1,

$$\sum_{i \in J} \sum_{j \in s(i)} \left(\max_{\nu \in \{1,2\}} \text{diam } K_i^\nu \right)^2 (v_h(Q_j)^2 + v_h(Q_i)^2) \leq 2 \sum_{i \in J} \sum_{j \in s(i)} (\text{diam } K_i^1)^2 v_h(Q_i)^2.$$

Again we use that $(\text{diam } K_i^\nu)^2 \leq \mathfrak{C} |K_i^\nu|$ by (2.2) and $|K_i^\nu|/3 = |D_i \cap K_i^\nu|$ for $i \in J^o$, $\nu \in \{1, 2\}$. It follows with (2.6) that

$$\sum_{i \in J} \sum_{j \in s(i)} \left(\max_{\nu \in \{1,2\}} \text{diam } K_i^\nu \right)^2 (v_h(Q_j)^2 + v_h(Q_i)^2) \leq \mathfrak{C} \sum_{i \in J} v_h(Q_i)^2 |D_i| \leq \mathfrak{C} \|v_h\|_2^2.$$

Thus we obtain from (3.3) that $|A_h - B_h| \leq \mathfrak{C} \|\beta\|_{1,2}^{1/2} \|v_h\|_2 \mathfrak{K}_h^{1/2}$. Lemma 3.6 now follows from (3.1). \square

Corollary 3.1 *Let $v_h \in V_h$. Then*

$$\|v_h\|_2^2 \leq \mathfrak{C}(1 + \|\beta\|_{1,2}) (b_h(v_h, \varrho(v_h \varphi)) + b_h(v_h, v_h)). \quad (3.4)$$

Proof: Lemma 3.5, 3.6 and 3.3 imply

$$\|v_h\|_2^2 \leq \mathfrak{C} (b_h(v_h, \varrho(v_h \varphi)) + \|\beta\|_{1,2}^{1/2} b_h(v_h, v_h)^{1/2} \|v_h\|_2),$$

so that

$$\|v_h\|_2^2 \leq \mathfrak{C} (b_h(v_h, \varrho(v_h \varphi)) + \|\beta\|_{1,2} b_h(v_h, v_h)) + \|v_h\|_2^2/2.$$

The corollary follows from this estimate. \square

4. Proof of Theorem 2.1 and 2.2.

The largest part of this section is taken by the

Proof of Theorem 2.1: Putting $v_h = u_h^{(k+1)}$ in (2.7), we get

$$\begin{aligned} & \nu \|u_h^{(k+1)}\|_{X_h}^2 + \tau_{k+1}^{-1} \|u_h^{(k+1)}\|_2^2 + b_h(u_h^{(k+1)}, u_h^{(k+1)}) \\ & \leq (g_h^{(k+1)}, u_h^{(k+1)}) + (2\tau_{k+1})^{-1} \|u_h^{(k+1)}\|_2^2 + (2\tau_{k+1})^{-1} \|u_h^{(k)}\|_2^2 \quad \text{for } 0 \leq k \leq N. \end{aligned}$$

Taking the sum with respect to $k \in \{0, \dots, s\}$, for $s \in \{0, \dots, N\}$, and using a simple shoestring argument, we find with $s = N$ that

$$\sum_{l=1}^{N+1} \tau_l b_h(u_h^{(l)}, u_h^{(l)}) + \nu \sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_{X_h}^2 \leq \sum_{l=1}^{N+1} \tau_l \|g_h^{(l)}\|_2 \|u_h^{(l)}\|_2 + \|u_h^{(0)}\|_2^2/2, \quad (4.1)$$

and for $s \in \{0, \dots, N\}$,

$$\begin{aligned} & \sum_{l=1}^{s+1} \tau_l (b_h(u_h^{(l)}, u_h^{(l)}) + \nu \|u_h^{(l)}\|_{X_h}^2) + \|u_h^{(s+1)}\|_2^2/2 \\ & \leq \sum_{l=1}^{s+1} \tau_l \|g_h^{(l)}\|_2 \|u_h^{(l)}\|_2 + \|u_h^{(0)}\|_2^2/2. \end{aligned}$$

The term $\|u_h^{(N+1)}\|_2^2/2$ was dropped on the left-hand side of (4.1) because it will not be needed. Since $b_h(v_h, v_h) \geq 0$ for $v_h \in V_h$ (Lemma 3.3), we may conclude that

$$\max_{1 \leq l \leq N+1} \|u_h^{(l)}\|_2^2/2 \leq \sum_{l=1}^{N+1} \tau_l \|g_h^{(l)}\|_2 \|u_h^{(l)}\|_2 + \|u_h^{(0)}\|_2^2/2. \quad (4.2)$$

Next we use (2.7) once more, this time with $v_h = \varrho_h(u_h^{(k+1)} \varphi)$, to obtain with Lemma 3.2 that

$$\begin{aligned} & \nu \tau_{k+1} ((u_h^{(k+1)}, \tilde{u}_h^{(k+1)}))_{X_h} + \tau_{k+1} b_h(u_h^{(k+1)}, \tilde{u}_h^{(k+1)}) + (u_h^{(k+1)}, \tilde{u}_h^{(k+1)}) \\ & \leq \tau_{k+1} (g_h^{(k+1)}, \tilde{u}_h^{(k+1)}) + (u_h^{(k)}, \tilde{u}_h^{(k)})^{1/2} (u_h^{(k+1)}, \tilde{u}_h^{(k+1)})^{1/2} \\ & \leq \tau_{k+1} (g_h^{(k+1)}, \tilde{u}_h^{(k+1)}) + (u_h^{(k+1)}, \tilde{u}_h^{(k+1)})/2 + (u_h^{(k)}, \tilde{u}_h^{(k)})/2 \quad \text{for } k \in \{0, \dots, N\}. \end{aligned}$$

Here and in the following, we use the abbreviation $\tilde{u}_h^{(k)} := \varrho_h(u_h^{(k)} \varphi)$ for $0 \leq k \leq N+1$. Now we again take the sum with respect to $k \in \{0, \dots, N\}$. By a simple shoestring argument and because $(u_h^{(N+1)}, \tilde{u}_h^{(N+1)}) \geq 0$ according to Lemma 3.2, we then get

$$\sum_{l=1}^{N+1} \tau_l b_h(u_h^{(l)}, \tilde{u}_h^{(l)}) \leq \sum_{l=1}^{N+1} \tau_l |(g_h^{(l)}, \tilde{u}_h^{(l)})| + \nu \sum_{l=1}^{N+1} \tau_l |((u_h^{(l)}, \tilde{u}_h^{(l)}))_{X_h}| + (u_h^{(0)}, \tilde{u}_h^{(0)})/2. \quad (4.3)$$

On the other hand, from Corollary 3.1 with $v_h = u_h^{(l)}$,

$$\|u_h^{(l)}\|_2^2 \leq \mathfrak{C}(1 + \|\beta\|_{1,2}) (b_h(u_h^{(l)}, \tilde{u}_h^{(l)}) + b_h(u_h^{(l)}, u_h^{(l)})) \quad \text{for } 1 \leq l \leq N+1.$$

Therefore, with (4.1) and (4.3),

$$\sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_2^2 \leq \mathfrak{C}(1 + \|\beta\|_{1,2}) \mathfrak{R}, \quad (4.4)$$

where

$$\begin{aligned} \mathfrak{R} := & \sum_{l=1}^{N+1} \tau_l |(g_h^{(l)}, \tilde{u}_h^{(l)})| + \nu \sum_{l=1}^{N+1} \tau_l |((u_h^{(l)}, \tilde{u}_h^{(l)}))_{X_h}| + (u_h^{(0)}, \tilde{u}_h^{(0)}) \\ & + \sum_{l=1}^{N+1} \tau_l \|g_h^{(l)}\|_2 \|u_h^{(l)}\|_2 + \|u_h^{(0)}\|_2^2. \end{aligned}$$

By adding the left- and right-hand side of (4.1) and (4.2) to respectively the left- and right-hand side of (4.4), and taking account of the fact that $b_h(v_h, v_h) \geq 0$ for $v_h \in V_h$ (Lemma 3.3), we get

$$\sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_2^2 + \max_{1 \leq l \leq N+1} \|u_h^{(l)}\|_2^2 + \nu \sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_{X_h}^2 \leq \mathfrak{C}(1 + \|\beta\|_{1,2}) \mathfrak{R}. \quad (4.5)$$

But with Lemma 3.2,

$$|(g_h^{(l)}, \tilde{u}_h^{(l)})| \leq (g_h^{(l)}, \varrho_h(g_h^{(l)} \varphi))^{1/2} (u_h^{(l)}, \tilde{u}_h^{(l)})^{1/2} \leq \mathfrak{C} \|g_h^{(l)}\|_2 \|u_h^{(l)}\|_2 \quad \text{for } 1 \leq l \leq N+1, \quad (4.6)$$

$$(u_h^{(0)}, \tilde{u}_h^{(0)}) \leq \mathfrak{C} \|u_h^{(0)}\|_2^2, \quad (4.7)$$

$$\nu \sum_{l=1}^{N+1} \tau_l |((u_h^{(l)}, \tilde{u}_h^{(l)}))_{X_h}| \leq \nu \sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_{X_h} \|\tilde{u}_h^{(l)}\|_{X_h} \leq \mathfrak{C} \nu \sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_{X_h}^2 \quad (4.8)$$

$$\leq \mathfrak{C} \left(\sum_{l=1}^{N+1} \tau_l \|g_h^{(l)}\|_2 \|u_h^{(l)}\|_2 + \|u_h^{(0)}\|_2^2 \right),$$

where the last inequality holds according to (4.1). Thus the left-hand side of (4.5) is bounded by $\mathfrak{C}(1 + \|\beta\|_{1,2}) \left(\sum_{l=1}^{N+1} \tau_l \|g_h^{(l)}\|_2 \|u_h^{(l)}\|_2 + \|u_h^{(0)}\|_2^2 \right)$. Noting that

$$\mathfrak{C}(1 + \|\beta\|_{1,2}) \tau_l \|g_h^{(l)}\|_2 \|u_h^{(l)}\|_2 \leq \mathfrak{C}(1 + \|\beta\|_{1,2})^2 \tau_l \|g_h^{(l)}\|_2^2 + \tau_l \|u_h^{(l)}\|_2^2 / 2 \quad (4.9)$$

for $1 \leq l \leq N + 1$, we may deduce from (4.5) that the left-hand side of (4.5) is bounded by

$$\mathfrak{C}(1 + \|\beta\|_{1,2})^2 \left(\sum_{l=1}^{N+1} \tau_l \|g_h^{(l)}\|_2^2 + \|u_h^{(0)}\|_2^2 \right) + \sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_2^2 / 2. \quad (4.10)$$

Therefore again from (4.5), by a simple shoestring argument,

$$\begin{aligned} & \sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_2^2 / 2 + \max_{1 \leq l \leq N+1} \|u_h^{(l)}\|_2^2 + \nu \sum_{l=1}^{N+1} \tau_l \|u_h^{(l)}\|_{X_h}^2 \\ & \leq \mathfrak{C}(1 + \|\beta\|_{1,2})^2 \left(\sum_{l=1}^{N+1} \tau_l \|g_h^{(l)}\|_2^2 + \|u_h^{(0)}\|_2^2 \right). \end{aligned}$$

This proves Theorem 2.1. \square

We finally turn to the

Proof of Theorem 2.2: From (2.8) with v_h replaced by U_h , we get

$$\nu \|U_h\|_{X_h}^2 + b_h(U_h, U_h) = (G_h, U_h) \leq \|G_h\|_2 \|U_h\|_2. \quad (4.11)$$

Thus Corollary 3.1 with U_h in the place of v_h yields

$$\|U_h\|_2^2 \leq \mathfrak{C}(1 + \|\beta\|_{1,2}) (b_h(U_h, W_h) + \|G_h\|_2 \|U_h\|_2), \quad (4.12)$$

where we used the abbreviation $W_h := \varrho_h(U_h \varphi)$. Moreover, from (2.8) with $v_h = W_h$,

$$b_h(U_h, W_h) \leq \nu |((U_h, W_h))_{X_h}| + |(G_h, W_h)|. \quad (4.13)$$

On the other hand, with Lemma 3.2,

$$|(G_h, W_h)| \leq (G_h, \varrho(G_h \varphi))^{1/2} (U_h, W_h)^{1/2} \leq \mathfrak{C} \|G_h\|_2 \|U_h\|_2, \quad (4.14)$$

and by Lemma 3.2, 3.3 and (4.11),

$$\nu |((U_h, W_h))_{X_h}| \leq \mathfrak{C} \nu \|U_h\|_{X_h}^2 \leq \mathfrak{C} (\nu \|U_h\|_{X_h}^2 + b_h(U_h, U_h)) \leq \mathfrak{C} \|G_h\|_2 \|U_h\|_2.$$

This estimate together with (4.12) – (4.14) yield

$$\|U_h\|_2^2 \leq \mathfrak{C}(1 + \|\beta\|_{1,2}) \|G_h\|_2 \|U_h\|_2,$$

hence with (4.11) and the relation $b_h(U_h, U_h) \geq 0$ (Lemma 3.3),

$$\|U_h\|_2^2 + \nu |((U_h, W_h))_{X_h}| \leq \mathfrak{C}(1 + \|\beta\|_{1,2}) \|G_h\|_2 \|U_h\|_2 \leq \mathfrak{C}(1 + \|\beta\|_{1,2})^2 \|G_h\|_2^2 + \|U_h\|_2 / 2.$$

Theorem 2.2 follows from this inequality. \square

5. Final comments.

It is natural to ask what are the minimal hypotheses on β such that Theorem 2.1 and 2.2 remain valid. In this respect we first remark that the assumptions $\beta \in C^1(\overline{\Omega})^2$, $\nabla \cdot \beta = 0$ are not sufficient for Theorem 2.2 to hold.

To see this, consider the case $\Omega = \{x \in \mathbb{R}^2 : |x| < 1\}$, $\beta(x) = (-x_2, x_1)$ for $x \in \overline{\Omega}$. Obviously $\beta \in C^1(\overline{\Omega})^2$ and $\nabla \cdot \beta = 0$. Note that since $\beta(x) \rightarrow 0$ for $|x| \rightarrow 0$, there is no function $\varphi \in C^1(\overline{\Omega})$ such that the second relation in (1.3) is valid. We further put $u_\nu(x) := |x|^2/\nu - 1/\nu$ for $x \in \overline{\Omega}$, $\nu \in (0, \infty)$. Then $-\nu \Delta u_\nu + \beta \cdot \nabla u_\nu = -4$ and $u_\nu|_{\partial\Omega} = 0$ for any $\nu \in (0, \infty)$. Since $\|u_\nu\|_2 = \sqrt{\pi/(3\nu)}$ for $\nu \in (0, \infty)$, it follows there cannot be a constant $C > 0$ such that $\|u\|_2 \leq C\|f\|_2$ for $\nu \in (0, \infty)$, $f \in L^2(\Omega)$ and $u \in H^2(\Omega)$ the solution of the boundary value problem

$$-\nu \Delta u + \beta \cdot \nabla u = f \text{ in } \Omega, \quad u|_{\partial\Omega} = 0. \quad (5.1)$$

With some additional effort, the preceding example may be extended to the case that Ω is an annular domain, for example $\Omega = \{y \in \mathbb{R}^2 : 1/2 < |y| < 1\}$, and to rotationally symmetric functions $f \in L^2(\Omega)$. In fact, for such a domain Ω and such a function f , the solution $u \in H^2(\Omega)$ of the boundary value problem

$$-\Delta u = f/\nu \text{ in } \Omega, \quad u|_{\partial\Omega} = 0, \quad (5.2)$$

is rotationally symmetric as well. Due to our choice of β , this means that u solves problem (5.1). On the other hand, u as a solution of (5.2) may be represented as a series involving eigenfunctions of the Laplacian on Ω . This representation shows that $\|u\|_2$ cannot be bounded uniformly in $\nu > 0$ when ν tends to zero.

It is true that in the above cases Ω is not a polygon, and that the exact solution of (5.1) and an approximate solution provided by (2.8) are not the same thing. Still the preceding observations on L^2 -estimates of solutions to (5.1) indicate it cannot be expected that an inequality as in Theorem 2.2 holds with a constant C independent of h and ν if $\beta \in H^1(\Omega)^2$ is only supposed to satisfy the relation $\nabla \cdot \beta = 0$.

Going a step further, one might conjecture that if $\nabla \cdot \beta = 0$, the second condition in (1.3) is not only sufficient (as we have shown), but even necessary for existence of a constant C as in Theorem 2.2. In fact, if such a constant exists, it should be expected there is $C > 0$ with $\|u\|_2 \leq C\|f\|_2$ for $\nu \in (0, \infty)$, $f \in L^2(\Omega)$ and $u \in H^2(\Omega)$ the solution of (5.1). This means in particular that the preceding constant C is independent of ν . In this situation, let $\underline{\beta} \in (0, \infty)$ and let (ν_n) be a sequence in $(0, \infty)$ with $\nu_n \rightarrow 0$. Let the solution of (5.1) with f, ν replaced by $\underline{\beta}, \nu_n$, respectively, be denoted by u_n , for $n \in \mathbb{N}$. Then, due to the estimate $\|u\|_2 \leq C\|f\|_2$ for ν, f, u as above, the sequence $(\|u_n\|_2)$ is bounded and thus possesses a subsequence which is weakly convergent in $L^2(\Omega)$. Let this subsequence be also denoted by (u_n) . Let $u \in L^2(\Omega)$ be the limit function of this sequence, and take $\phi \in C_0^\infty(\Omega)$. Choosing $f = -\underline{\beta}$, $u = u_n$ in (5.1), multiplying this equation by ϕ , integrating over Ω and performing some integrations by parts, we arrive at the equation

$$\int_{\Omega} u_n (-\nu_n \Delta \phi - \beta \cdot \nabla \phi + \underline{\beta} \phi) = 0 \quad \text{for } n \in \mathbb{N}.$$

Letting n tend to infinity, we see that $\int_{\Omega} u(-\beta \cdot \nabla \phi + \underline{\beta} \phi) dx = 0$. This means that u satisfies the equation $-\beta \cdot \nabla u = \underline{\beta}$ in a weak sense. Thus, if $u \in C^1(\overline{\Omega})$, the second condition in (1.3) would hold with u in the role of φ . However, the question of how to prove C^1 -regularity of u is open.

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