



HAL
open science

Expansion of the almost sure spectrum in the weak disorder regime

Denis Borisov, Francisco Hoecker-Escuti, Ivan Veselić

► **To cite this version:**

Denis Borisov, Francisco Hoecker-Escuti, Ivan Veselić. Expansion of the almost sure spectrum in the weak disorder regime. 2014. hal-01053573

HAL Id: hal-01053573

<https://hal.science/hal-01053573>

Preprint submitted on 31 Jul 2014

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

EXPANSION OF THE ALMOST SURE SPECTRUM IN THE WEAK DISORDER REGIME

DENIS BORISOV, FRANCISCO HOECKER-ESCIUTI, AND IVAN VESELIĆ

ABSTRACT. The spectrum of random ergodic Schrödinger-type operators is almost surely a deterministic subset of the real line. The random operator can be considered as a perturbation of a periodic one. As soon as the disorder is switched on via a global coupling constant, the spectrum expands. We estimate how much the spectrum expands at its bottom for operators on $\ell^2(\mathbb{Z}^d)$.

CONTENTS

1. Introduction	1
2. General model	6
3. Main results	8
4. Periodic comparison operators	9
5. Perturbation calculation	12
5.1. Sign-changing random variables	13
5.2. A simple non-degeneracy condition	14
5.3. Proof of Theorem 5.1	15
5.4. Positive random variables	19
5.5. Application to the discrete alloy type model	23
6. Appendix	24
6.1. An interesting example: Proof of Theorem 1.2	24
6.2. Proof of Lemma 4.1	25
Acknowledgment	27
References	27

1. INTRODUCTION

Due to the self-averaging property of ergodic Schrödinger operators the resulting spectrum is almost surely a fixed subset of the real line. If a random operator is a perturbation of a periodic operator, it is of interest to know how the spectrum expands once we switch on the disorder via a global coupling constant. Apart from the genuine interest to identify the location of the spectrum, this is also of central importance when identifying energy regions corresponding to localized wavepackets.

Otherwise it may happen that one proves a Wegner estimate, a Lifschitz tail bound or a similar statement related to localization, and then later discovers that the considered energy regime belongs to the resolvent set.

In this paper we consider an ϵ -small random perturbation of a discrete translation-invariant operator and we study how the bottom of its spectrum behaves. By symmetry, similar estimates apply to the location of the maximum of the spectrum, in a weak disorder regime. To fix the ideas, let us introduce a prototypical example. Let $\mathcal{H} = \ell^2(\mathbb{Z}^d)$ and $\Delta_{\mathbb{Z}^d} : \mathcal{H} \rightarrow \mathcal{H}$ the (negative definite) discrete Laplacian on \mathbb{Z}^d , i.e.

$$(\Delta_{\mathbb{Z}^d} u)(n) := \sum_{|n-m|_\infty=1} (u(m) - u(n)).$$

We define the operator $H_0 : \mathcal{H} \rightarrow \mathcal{H}$ by

$$H_0 := -\Delta_{\mathbb{Z}^d} + W,$$

where W is the multiplication operator by a real-valued function, which we also denote by W and which we assume periodic with respect to the subgroup $\gamma := N\mathbb{Z}^d$.

Let $\square := [0, N-1]^d \subset \mathbb{Z}^d$ and $V^\square \in \ell^\infty(\mathbb{Z}^d)$ be a non-trivial, compactly supported single-site potential satisfying

$$\text{supp}(V^\square) \subset \square.$$

Let $(\omega_k)_{k \in \gamma}$ be a sequence of non-trivial, bounded, independent, identically distributed random variables. For the sake of the introduction, assume that $\{-1, 1\} \in \text{supp } \omega_0 \subset [-1, 1]$. From now on we denote by $V_\omega : \ell^2(\mathbb{Z}^d) \rightarrow \ell^2(\mathbb{Z}^d)$ the diagonal operator defined, for $f \in \ell^2(\mathbb{Z}^d)$, as

$$(1) \quad (V_\omega f)(x) = \sum_{k \in N\mathbb{Z}^d} \omega_k V^\square(x-k) f(x).$$

To motivate our results, let us consider the following discrete alloy-type random Schrödinger operator defined by

$$(2) \quad H_{\omega, \epsilon} := H_0 + \epsilon V_\omega.$$

Under the stated assumptions, this operator is ergodic, and thus there exists a set $\Sigma_\epsilon \subset \mathbb{R}$ such that

$$\sigma(H_{\omega, \epsilon}) = \Sigma_\epsilon$$

with probability 1 (see e.g. [19]). From now on we refer to Σ_ϵ as the *almost-sure spectrum* of $H_{\omega, \epsilon}$. The best known example of this kind of operators is the celebrated *Anderson model*, where H_0 is the discrete Laplacian on \mathbb{Z}^d (i.e. $W \equiv 0$), $V^\square = \delta_0$ and $N = 1$. In this case, it is not hard to see ([19]) that the bottom of the spectrum of the perturbed operator $E_\epsilon := \inf(\Sigma_\epsilon)$ moves away from the bottom of the spectrum of the free operator $E_0 := \inf(\Sigma_0)$ as

$$E_\epsilon = E_0 - \epsilon.$$

If one considers instead, for example, the dipole model, i.e. $V^\square = \delta_0 - \delta_{e_1}$ with $e_1 = (1, 0, \dots, 0)$, it is proven in [4] that

$$E_\epsilon \leq E_0 - C\epsilon^2.$$

In this note we study this question for a very general, wide class of operators (see assumptions in Section 2). More precisely, we prove some upper bounds of the quantity $E_\epsilon - E_0$, which in turns gives us information on the location of the spectrum of the perturbed operator. We also discuss some partial results on the lower bound.

In order to state the result in this setting, we need to consider the operator H_0 with $N\mathbb{Z}^d$ -periodic boundary conditions. Because of the translation invariance, the subspace of $N\mathbb{Z}^d$ -periodic functions in $\ell^\infty(\mathbb{Z}^d)$ is invariant under the action of H_0 . This subspace is N^d -dimensional, so that the action of the operator corresponds to a matrix we denote by

$$(3) \quad H_0^\square : \ell^2(\square) \rightarrow \ell^2(\square).$$

We now state the result.

Theorem 1.1. *Let $H_{\omega,\epsilon}$ be the alloy-type random Schrödinger operator defined by (2) and E_ϵ the bottom of its corresponding almost-sure spectrum. To the $N\mathbb{Z}^d$ -periodic operator H_0 we associate a Hermitian matrix $H_0^\square \in \mathbb{C}^{N^d \times N^d}$, defined as in (3), and we let $\psi_1 \in \ell^2(\square)$ be the (unique normalized) positive ground state of H_0^\square . Define*

$$A_1 := -\langle \psi_1, V^\square \psi_1 \rangle_{\ell^2(\square)}.$$

There exists $A_2 \leq 0$ such that for $\epsilon > 0$ small enough

$$E_\epsilon \leq E_0 + \epsilon A_1 + \epsilon^2 A_2.$$

Furthermore, if $A_1 = 0$ then $|A_2|$ is non-zero and larger than the spectral gap of H_0^\square , i.e. the difference between its two smallest eigenvalues.

We provide an explicit formula for the constant A_2 only in the next section as it requires the introduction of additional notation. We have an analogous estimate for (fibers of) periodic operators, see Theorems 5.1 and 5.8. In fact, the estimate for periodic operators is one step in the proof of Theorem 1.1. In the context of periodic operators we have a related, complementary lower bound, see Lemmas 5.7 and 5.9.

We would like to make some remarks on the relevance of this result. First, the location of the bottom of the spectrum with respect to the coupling constant has been the subject of several papers: with periodic potentials in dimension one [22] and in arbitrary dimension [13], [6], as well as with random positive potentials [14] and under some generic assumptions on W [16]. Understanding the spectrum provides valuable information on the solutions of partial differential equations. In particular, if one considers the Schrödinger equation for the Hamiltonian $H_{\omega,\epsilon}$, the spectral type of the Anderson model characterizes the

transport properties of the underlying disordered medium. For this model, the spectrum is expected to exhibit a transition from localized states at the bottom of the spectrum (pure point spectrum with exponentially localized eigenfunctions) to extended states (absolutely continuous spectrum) in the bulk of the spectrum. This *Anderson transition* is still a conjecture in the setting of this article. The existence of localized states at the bottom of the spectrum has been studied in many papers. We invite the reader to consult the monographs [5], [19], [21], [10] and their extensive bibliography. The perturbative regime $\epsilon \ll 1$ has attracted much attention [1], [23], [15], [16], [7], [4], [11], [9], [2], [3]. In this regime one can prove very precise estimates of the interval of localization, namely that states with energies in

$$I_\eta(\epsilon) := (-\infty, -C_0\epsilon^\eta] \cap \Sigma_\epsilon = (-\infty, -C_0\epsilon^\eta] \cap [E_\epsilon, +\infty)$$

are localized. In [7] it was proved that in dimension $d = 3$ one may take η to be as large as 2 and in [11] that for $d \geq 2$ this holds with $\eta < 2$. These results are meaningful, as for the Anderson model $E_\epsilon = -C_1\epsilon$. If we now consider different potentials, we may have a quadratic expansion of the bottom of the spectrum E_ϵ , and understanding where the bottom of the spectrum lies appears to be crucial, so that the interval of localization is non-trivial. Some of the issues addressed in this note were already explored in [16] where it is assumed that the single-site potential has a non-zero mean and the Floquet eigenvalues of the underlying periodic discrete Schrödinger operator $H_0 = -\Delta + W$ are assumed to be non-degenerated, as well as in [4] for the dipole potential. These are special cases of our models. The general operator we study corresponds roughly to tridiagonal block matrices of the form

$$(4) \quad \begin{pmatrix} \ddots & \ddots & \ddots & & \\ B^* & A & B & 0 & \ddots \\ 0 & B^* & A & B & 0 \\ \ddots & 0 & B^* & A & B \\ & & \ddots & \ddots & \ddots \end{pmatrix} + \begin{pmatrix} \ddots & \ddots & \ddots & & \\ 0 & \omega_{n-1}V^\square & 0 & \ddots & \\ \ddots & 0 & \omega_n V^\square & 0 & \ddots \\ & \ddots & 0 & \omega_{n+1}V^\square & 0 \\ & & \ddots & \ddots & \ddots \end{pmatrix},$$

where A and V^\square are Hermitian matrices and $\{\omega_n\}$ i.i.d. random variables. We introduce in Section 2 the general framework in which our results are obtained.

To complete the description of the obtained results, let us briefly address the question of the optimality of the lower bound (or at least its exponent). As far as the authors know, there is no result in the literature in this direction in a discrete setting (apart from the Anderson model, where the bottom of the spectrum is known explicitly). One may naively expect, from perturbation theory, that the behavior should be linear or quadratic. The question turns out to be more subtle as the

behavior may depend on the speed at which the Floquet eigenvalues associated to the bottom of the spectrum approach their minimum, as the following example shows.

Theorem 1.2. *For $H_0 := (-\Delta_{\mathbb{Z}})^2$ defined on $\ell^2(\mathbb{Z})$ and V^\square the multiplication operator given by the following single-site potential:*

$$V^\square := -\frac{1}{2}\delta_{-1} + \delta_0 - \frac{1}{2}\delta_1.$$

Set as before $H_{\omega,\epsilon} := H_0 + \epsilon V_\omega$, cf. (1) and (2). Pick some $\xi > 1/4$. Then for $\epsilon > 0$ small enough we have

$$(5) \quad E_\epsilon := \inf \sigma(H_0 + \epsilon V_\omega) \leq -\frac{1}{6}\epsilon^{1+2\xi}.$$

For this example, which is of the form (4), the coefficient A_1 corresponding to the linear term vanishes. The bound in Theorem 1.2 is nevertheless better than quadratic thanks to the quartic behavior of the Floquet eigenvalues in a neighbourhood of their minimum. Unfortunately, apart from the trivial linear bound, we have no corresponding lower bound, although some results in this article provide a first step in this direction.

This work can be extended in several directions. It would be very interesting to find the corresponding lower bounds, or at least conditions under which the infimum of the spectrum does not expand linearly. A related question concerns the expansion of the spectrum near a band edge, where one can also prove Anderson localisation. Note that if one studies the expansion of the spectrum from a band edge instead of the bottom of the spectrum, the Floquet eigenvalues may vanish faster than quadratically when approaching the edge, even for the operator defined by (2). Rather than a pathological example, Theorem 1.2 provides a model for this situation. Another question of interest is the study of overlapping single-site potentials. Under some non-degeneracy condition (see Remark 4.5) the results stated here can be extended to this situation, but a full understanding needs to consider periodic approximations of every order, something we also believe necessary to prove the lower bounds complementing Theorem 1.1.

In a forthcoming project we consider the same questions for operators of Schrödinger type in the continuum setting, i.e. for operators acting on (dense subspaces) of $L^2(\mathbb{R}^d)$. Most of our findings are similar. In the continuum, it is more natural to define the operators via quadratic forms, and then formulate appropriate regularity conditions. Also, certain additional compactness arguments are necessary, due to the infinite dimensionality of the Hilbert space over the periodicity cell. On the other hand, in specific situations, better results are possible in the continuum setting, due to unique continuation principles for solutions of partial differential equations.

2. GENERAL MODEL

Let $d \geq 1$ be the space dimension, $D = \mathbb{Z}^d$ be the physical space and $\gamma = N\mathbb{Z}^d$ a sub-lattice of D . We denote by \square its periodicity cell, i.e. $\square := [0, N-1]^d \cap \mathbb{Z}^d$. Note that $D = \bigcup_{k \in \gamma} \{x \in D : x - k \in \square\}$. We

also denote the reciprocal periodicity cell as $\square^* := [0, \frac{2\pi}{N}]^d$. From now on we assume the following hypotheses to hold.

(HA) Let $H_0 : \ell^2(D) \rightarrow \ell^2(D)$ be a bounded, non-negative Hermitian operator defined by the matrix

$$H_0 := (H_0(k, k'))_{k, k' \in D},$$

satisfying the following properties:

- for all $k, k' \in D$, we have $H_0(k, k') = \overline{H_0(k', k)}$;
- there exists $k_0 \neq 0$ such that $H_0(0, k_0) \neq 0$;
- the associated operator is γ -invariant, i.e. for every $k \in \gamma$

$$\langle \tau_k u, H_0 \tau_k v \rangle = \langle u, H_0 v \rangle,$$

where $u, v \in \ell^2(D)$ and τ_k is the translation by $k \in \gamma$ operator; and

- the associated operator is of *finite hopping range* with hopping range N , i.e. if $|k - k'| \geq N$ then

$$H_0(k, k') = 0.$$

- Through a global energy shift we may assume, with no loss of generality, that $E_0 := \inf \sigma(H_0) = 0$.

Note that if an operator is of finite hopping range with hopping range R , for some $R > 0$, then it also is of finite hopping range with hopping range R' for any $R' > R$. On the other hand, any γ -invariant operator is also $n\gamma$ -invariant, $n \in \mathbb{N}$. This means that we can always assume that $R = N$, without loss of generality.

(HB) Let $V^\square : \square \rightarrow \mathbb{R}$ be a non-trivial Hermitian matrix (we call it the *single-cell potential*, even when V^\square is not diagonal). For any bounded sequence $(\omega_k)_{k \in \gamma}$ of real numbers, we define the block diagonal operator

$$\begin{aligned} V_\omega &: \ell^2(D) \rightarrow \ell^2(D) \\ V_\omega &:= \sum_{k \in \gamma} \omega_k \tau_{-k} V^\square \tau_k. \end{aligned}$$

For any real number $q \in \mathbb{R}$, we denote also by q the constant sequence indexed by γ , equal to q on every site in \mathbb{Z}^d . We thus have, for example, that

$$(6) \quad V_q := q \sum_{k \in \gamma} \tau_{-k} V^\square \tau_k$$

and V_q is γ -invariant.

From now on, the values of ω will be drawn from a sequence of bounded, non-trivial, independent and identically distributed random variables with distribution measure μ . We will write $S_\mu := \text{supp } \mu$ and we assume that

$$\{s_-, s_+\} \in S_\mu \subset [s_-, s_+],$$

where s_- and s_+ satisfy one of the following alternatives:

(HC) The random variables change sign, i.e. $s_- < 0 < s_+$.

(HC') The random variables are positive, i.e. $0 \leq s_- < s_+$.

The methods in this paper may also be adapted to negative random variables.

Remark 2.1. It looks tempting, in order to achieve $s_- = 0$, to renormalize the random variables by adding and subtracting some periodic potential, but in this case the underlying non-random operator depends on ϵ . On the other hand, it is indeed allowed to rescale the random variables by absorbing the scaling factor in the single site potential V^\square .

Let us now define our object of study. For each $\epsilon > 0$, we let

$$H_{\omega,\epsilon} := H_0 + \epsilon V_\omega$$

which is a self-adjoint, ergodic operator. We denote its almost-sure spectrum by Σ_ϵ and by

$$(7) \quad E_\epsilon := \inf \Sigma_\epsilon$$

the bottom of the spectrum. We also write $H_{q,\epsilon} := H_0 + \epsilon V_q$ the corresponding operator with V_ω replaced by the periodic potential V_q (defined as in (6)) and $E_{q,\epsilon} := \inf \sigma(H_{q,\epsilon})$. In the following we will study the bottom of the spectrum E_ϵ of the random operator for small ϵ .

We define a finite dimensional matrix associated to the above objects. Define the $(|\square| \times |\square|)$ -matrix $H_0^\square(\theta)$ by its coefficients

$$(8) \quad \begin{aligned} (H_0^\square(\theta))(k, k') &:= \sum_{m \in \gamma} e^{i\theta \cdot m} H_0(k, k' - m) \\ &= \sum_{\substack{|m| \leq N \\ m \in N\mathbb{Z}^d}} e^{i\theta \cdot m} H_0(k, k' - m), \end{aligned}$$

where $k, k' \in \square$. Note that the second line is a consequence of the finite hopping range and the sum in (8) is thus finite. Now define the matrix $H_{q,\epsilon}^\square(\theta)$ by

$$H_{q,\epsilon}^\square(\theta) := H_0^\square(\theta) + \epsilon q V^\square.$$

Remark 2.2. The matrix $H_{q,\epsilon}^\square$ represents the action of $H_{q,\epsilon}$ on the fiber of θ -quasiperiodic functions in the Floquet–Bloch direct integral decomposition. More precisely, let (abusing notation) $\varphi \in \ell^2(\square) \subset \ell^2(\mathbb{Z}^d)$. Then, regarding $H_{q,\epsilon}$ as an operator $\ell^\infty(\mathbb{Z}^d) \rightarrow \ell^\infty(\mathbb{Z}^d)$,

$$(9) \quad H_{\epsilon,q}^\square(\theta)\varphi = \chi_\square H_{\epsilon,q} \left(\sum_{m \in \gamma} e^{i\theta \cdot m} \tau_m \varphi \right) \in \ell^2(\square), \quad (\text{abusing notation})$$

where χ_\square is the indicator function of $\square \subset \mathbb{Z}^d$.

3. MAIN RESULTS

Recall that, by the continuity of the Floquet–Bloch eigenvalues ([17], [20]), there exists some θ such that

$$\inf \sigma(H_0^\square(\theta)) = E_0 := \inf \sigma(H_0) = 0.$$

We denote by $\Theta \subset \square^*$ the compact set of θ for which the last equality holds. From now on we fix some $\theta \in \Theta$, so the quantities below will depend on θ . Let \mathcal{V}_0 be the eigenspace of $H_0(\theta)$ associated to the eigenvalue $E_0 = 0$ and p its multiplicity. Choose an orthonormal basis ψ_j , $j = 1, \dots, p$ spanning \mathcal{V}_0 and diagonalizing the Hermitian matrix $A \in \mathbb{C}^{p \times p}$, given by the coefficients

$$A_{ij} := \langle \psi_i, V^\square \psi_j \rangle.$$

We take the eigenvalues of the matrix A in ascending order counting multiplicities so that $P_1 := A_{11} = \langle \psi_1, V^\square \psi_1 \rangle$ is the minimal eigenvalue and $P_p := A_{pp} = \langle \psi_p, V^\square \psi_p \rangle$ is the maximal eigenvalue of A .

Our result for sign-changing random variables reads as follows.

Theorem 3.1. *Assume (HA), (HB) and (HC). Fix $\theta \in \Theta$ and define*

$$(10) \quad A_1 := \inf_{q \in S_\mu} \inf_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)}=1}} q \langle \psi, V^\square \psi \rangle = \min(s_+ P_1, s_- P_p) \leq 0,$$

and

$$(11) \quad A_2 := -\max(s_-^2, s_+^2) \sup_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)}=1}} \sup_{\substack{\varphi \in \mathcal{V}_0^\perp \\ \|\varphi\|_{\ell^2(\square)}=1}} \frac{|\langle \psi, V^\square \varphi \rangle|^2}{\langle H_0^\square(\theta) \varphi, \varphi \rangle} \leq 0.$$

For any $\epsilon > 0$ small enough the following holds: if $A_1 \neq 0$,

$$E_\epsilon \leq \epsilon A_1,$$

whereas if $A_1 = 0$, then

$$E_\epsilon \leq \epsilon^2 A_2 + O(\epsilon^3).$$

Finally, if $A_1 = A_2 = 0$, then

$$E_\epsilon \leq 0.$$

Our result for positive random variables reads as follows.

Theorem 3.2. *Assume (HA), (HB) and (HC'). Fix $\theta \in \Theta$. Let us define the subspace $\mathcal{V}_{01} \subset \mathcal{V}_0$ as*

$$\mathcal{V}_{01} := \text{span} \{ \psi_i : i \in \mathbb{N}, P_i = P_1 \},$$

i.e. the eigenspace of A associated to its minimal eigenvalue P_1 . Define

$$(12) \quad A'_1 := \inf_{q \in S_\mu} \inf_{\substack{\psi \in \mathcal{V}_{01} \\ \|\psi\|_{\ell^2(\square)}=1}} q \langle \psi, V^\square \psi \rangle = \min(s_+ P_1, s_- P_1) \in \mathbb{R},$$

and

$$(13) \quad A'_2 := -s_+^2 \sup_{\substack{\psi \in \mathcal{V}_{01} \\ \|\psi\|_{\ell^2(\square)}=1}} \sup_{\substack{\varphi \in \mathcal{V}_0^\perp \\ \|\varphi\|_{\ell^2(\square)}=1}} \frac{|\langle \psi, V^\square \varphi \rangle|^2}{\langle H_0^\square(\theta) \varphi, \varphi \rangle} \leq 0.$$

For any $\epsilon > 0$ small enough the following holds: if $P_1 \neq 0$,

$$E_\epsilon \leq \epsilon A'_1,$$

whereas if $P_1 = 0$, then

$$E_\epsilon \leq \epsilon^2 A'_2 + O(\epsilon^3).$$

Finally, if $P_1 = A'_2 = 0$, then

$$E_\epsilon \leq 0.$$

Note that A_1 is always non-positive but A'_1 may be positive. Both Theorems 3.1 and 3.2 are a consequence of similar upper bounds for perturbations of periodic operators restricted to a fiber, see Theorems 5.1 and 5.8. For these upper bounds, we present complementary lower bounds in Lemmas 5.7 and 5.9.

4. PERIODIC COMPARISON OPERATORS

In the present section we reduce the problem of studying E_ϵ to that of understanding certain adapted operators which are periodic with respect to a sublattice. Define

$$\square_n := \bigcup_{\substack{m \in \gamma \\ |m| \leq nN}} \square + m$$

and $\chi_n := \chi_{\square_n}$, i.e.

$$\chi_n(x) = \begin{cases} 1 & \text{for } x \in \square_n \\ 0 & \text{otherwise.} \end{cases}$$

Note that $\square_0 = \square$ and that \square_n is just the collection of $(2n+1)^d$ disjoint translates of \square . Let us start by stating the following lemma.

Lemma 4.1. *Let u be a θ -quasi- γ -periodic function, i.e. such that for all $n \in \mathbb{Z}^d$ and $k \in \gamma$ we have*

$$u(n+k) = e^{-i\theta \cdot k} u(n).$$

Define,

$$u_n := \chi_n u, \quad n = 0, 1, 2, \dots$$

Then

$$\lim_{n \rightarrow \infty} \frac{\langle u_n, H_{q,\epsilon} u_n \rangle_{\ell^2(D)}}{\|u_n\|_{\ell^2(D)}^2} = \frac{\langle u_0, H_{q,\epsilon}^\square(\theta) u_0 \rangle_{\ell^2(\square)}}{\|u_0\|_{\ell^2(\square)}^2}.$$

The proof of this lemma is found in the appendix.

For the definition of periodic comparison operators we introduce

$$(14) \quad \Omega_{\text{per}}^n := \{\omega \in \Omega : \omega \text{ is periodic w.r.t. } n\gamma\}.$$

We now state the first comparison theorem.

Theorem 4.2. *Assume (HA), (HB) and either (HC) or (HC'). Let $\epsilon \geq 0$, $n \in \mathbb{N}$ and let $\omega \in \Omega_{\text{per}}^n$ be a $n\gamma$ -periodic sequence of real numbers satisfying $\omega \in (S_\mu)^\gamma$, i.e. $\omega_k \in S_\mu$ for all $k \in \gamma$. Then, we have*

$$\sigma(H_{\omega,\epsilon}) \subset \Sigma_\epsilon.$$

We immediately deduce the following upper bound on the minimum of the spectrum.

Corollary 4.3. *Assume (HA), (HB) and either (HC) or (HC'). Let $\epsilon \geq 0$, then*

$$E_\epsilon \leq \inf_{q \in S_\mu} E_{q,\epsilon}.$$

Proof of Theorem 4.2. For the calculation below, we need a Weyl sequence of compactly supported functions. This can indeed be done, since we only deal with bounded operators. Fix $\omega \in \Omega_{\text{per}}^n$ and $E \in \sigma(H_{\omega,\epsilon})$. By Floquet–Bloch theory, there exists some θ and some normalized state $f \in \ell^2(\square)$ for which

$$E = \langle H_{q,\epsilon}^\square(\theta) f, f \rangle.$$

We extend f as a θ -quasi- γ -periodic function, i.e. for any $x \in \mathbb{Z}^d$ let $k \in \gamma$ such that $x - k \in \square$ and let

$$f(x) := e^{i\theta \cdot k} f(x - k).$$

Using Lemma 4.1, extract a sub-sequence $\{f_n\}$ from the sequence of functions $\left\{ \frac{\chi_n f}{\|\chi_n f\|_2} \right\}$, such that

$$|\langle (H_{\omega,\epsilon} - E) f_n, f_n \rangle| \leq 1/n$$

and satisfying, for a sequence $l_n \in \mathbb{N}$,

$$\text{supp } f_n \subset \Lambda_{l_n},$$

where Λ_{l_n} is a cube centered at zero and sidelength l_n . For $x \in \gamma$ we define

$$\Omega(x, n) := \{\omega' \in \Omega : \forall k \in (x + \Lambda_{l_n}) \cap \gamma : |\epsilon(\omega'_k - \omega_k)| \leq 1/n\}.$$

Now, since $\omega \in (S_\mu)^\gamma$,

$$\mathbb{P}[\Omega(x, n)] > 0,$$

and for $x, y \in \gamma$ satisfying $|x - y| > l_n$, the events $\Omega(x, n)$ and $\Omega(y, n)$ are independent (and identically distributed). Using Borel–Cantelli lemma, we see that the event

$$\Omega' := \bigcap_{n \in \mathbb{N}} \bigcup_{x \in \gamma} \Omega(x, n)$$

has probability one.

From the definition of $\Omega(x, n)$, we have that given $\omega' \in \Omega'$ and $n \in \mathbb{N}$, there exists a $x(n, \omega')$ such that $\omega' \in \Omega(x, n)$. We write from now on $\tau_{x(n, \omega')} f_n$ for the translated function $f_n(\cdot - x(n, \omega'))$. Let $\omega' \in \Omega'$ and $n \in \mathbb{N}$, and calculate

$$\begin{aligned} & \langle (H_{\omega', \epsilon} - E) \tau_{x(n, \omega')} f_n, \tau_{x(n, \omega')} f_n \rangle \\ &= \langle (H_0 - E) \tau_{x(n, \omega')} f_n, \tau_{x(n, \omega')} f_n \rangle + \epsilon \langle V_{\omega'} \tau_{x(n, \omega')} f_n, \tau_{x(n, \omega')} f_n \rangle \\ &= \langle (H_0 - E) f_n, f_n \rangle + \epsilon \langle V_\omega \tau_{x(n, \omega')} f_n, \tau_{x(n, \omega')} f_n \rangle \\ & \quad + \epsilon \langle V_{\omega' - \omega} \tau_{x(n, \omega')} f_n, \tau_{x(n, \omega')} f_n \rangle \\ &= \langle (H_{\omega, \epsilon} - E) f_n, f_n \rangle + \epsilon \langle V_{\omega' - \omega} \tau_{x(n, \omega')} f_n, \tau_{x(n, \omega')} f_n \rangle \end{aligned}$$

Note that $|\epsilon V_{\omega' - \omega}(x)| \leq \|V^\square\|_\infty / n$ if $x \in \text{supp } f_n(\cdot - x(n, \omega'))$, so that

$$\left| \langle (H_{\omega', \epsilon} - E) \tau_{x(n, \omega')} f_n, \tau_{x(n, \omega')} f_n \rangle - \langle (H_{\omega, \epsilon} - E) f_n, f_n \rangle \right| \leq \frac{1}{n} \|V^\square\|_\infty.$$

In particular, we see that $\tau_{x(n, \omega')} f_n$ is a Weyl sequence. \square

Remark 4.4. This is an adaptation of a well known argument of Kirsch and Martinelli [12] in the continuous setting, with S_μ connected and V^\square a multiplication operator.

Remark 4.5. When the random potential is diagonal (as in the introduction), the proof above can be adapted to *overlapping*, but compactly supported single-site potentials $V^\square \in \ell^\infty(D)$ as long as

$$\sum_{n \in \gamma} V^\square(\cdot - n) \not\equiv 0.$$

Note that if this condition does not hold then $H_q = H_0$ for all q . One way around this problem would be to consider periodic (non-constant) sequences of coupling constants ω_n such that the resulting periodic potential is not zero.

To prove the following converse to Theorem 4.2 we define

$$\Omega_{\text{per}} := \{\omega \in \Omega : \exists n \in \mathbb{N} \text{ such that } \omega \text{ is periodic w.r.t. } n\gamma\} = \bigcup_{n \in \mathbb{N}} \Omega_{\text{per}}^n.$$

Lemma 4.6. *Denote by Σ_ϵ the almost sure spectrum of $H_{\omega,\epsilon}$. Then:*

$$\Sigma_\epsilon \subset \overline{\bigcup_{\omega \in \Omega_{per}} \sigma(H_{\omega,\epsilon})}.$$

Proof. Let $n \in \mathbb{N}$ and set

$$\begin{aligned} \omega_k^{(n)} &= \omega_k & \text{for } f \in \square_N \\ \omega_k^{(n)} &= \omega_j & \text{if } j - k \in N\gamma. \end{aligned}$$

Let $C_0(\mathbb{Z}^d)$ be the set of compactly supported functions in $\ell^2(\mathbb{Z}^d)$. Choose any $\varphi \in C_0(\mathbb{Z}^d)$. Then

$$\lim_{n \rightarrow \infty} \|H_{\omega,\epsilon}\varphi - H_{\omega^{(n)},\epsilon}\varphi\| = 0,$$

i.e. we have strong convergence $H_{\omega^{(n)}} \rightarrow H_\omega$. Since the operators H_ω are bounded, the set C_0 is an operator core for H_ω . This implies that we have strong convergence on the whole $\ell^2(\mathbb{Z}^d)$.

By the resolvent equation, for any $E \in \mathbb{R} \setminus \Sigma$,

$$\begin{aligned} & (H_{\omega,\epsilon} - E)^{-1} - (H_{\omega^{(n)},\epsilon} - E)^{-1} \\ &= (H_{\omega,\epsilon} - E)^{-1}(V_\omega - V_{\omega^{(n)}})(H_{\omega^{(n)},\epsilon} - E)^{-1} \\ &= (H_{\omega^{(n)},\epsilon} - E)^{-1}(V_\omega - V_{\omega^{(n)}})(H_{\omega,\epsilon} - E)^{-1}, \end{aligned}$$

which converges strongly to 0. We know that if $E \in \mathbb{R} \setminus \Sigma$, then $(H_{\omega,\epsilon} - E)^{-1}\varphi \in \ell^2(D)$ for any $\varphi \in \ell^2(D)$ and that, using Theorem 4.2, the inclusion $\sigma(H_{\omega,\epsilon}(n)) \subset \Sigma$ holds for any ω in the support of the product measure $\bigotimes_D \mu$. To conclude, we apply Theorem VIII.24 in [20]

which tells us that

$$\sigma(H_{\omega,\epsilon}) \subset \overline{\bigcup_{n \in \mathbb{N}} \sigma(H_{\omega^{(n)}})}.$$

This finishes the proof. \square

In particular we obtain the following corollary.

Corollary 4.7. *As before we set*

$$\Omega_{per} := \{\omega \in \Omega : \exists n \in \mathbb{N} \text{ such that } \omega \text{ is periodic w.r.t. } n\gamma\}$$

and denote by Σ_ϵ the almost sure spectrum of $H_{\omega,\epsilon}$. Then:

$$\inf \Sigma_\epsilon = \inf \bigcup_{\omega \in \Omega_{per}} \sigma(H_{\omega,\epsilon}).$$

5. PERTURBATION CALCULATION

For the readers convenience we recall the definition of the constants A_1 and A_2 , the notation and the statement of the theorems before the proofs. By the continuity of the Floquet–Bloch eigenvalues there exists some θ such that

$$E_0 := \inf \sigma(H_0) = \inf \sigma(H_0^\square(\theta)) = 0.$$

We denote by $\Theta \subset \square^*$ the compact set of θ for which the last equality holds. From now on we fix some $\theta \in \Theta$, so the quantities below will depend on θ . Let \mathcal{V}_0 be the eigenspace of $H_0(\theta)$ associated to the eigenvalue $E_0 = 0$, p its multiplicity and choose an orthonormal basis ψ_j , $j = 1, \dots, p$ spanning \mathcal{V}_0 and diagonalizing the Hermitian matrix $A \in \mathbb{C}^{p \times p}$, given by the coefficients

$$A_{ij} := \langle \psi_i, V^\square \psi_j \rangle.$$

We take the eigenvalues of the matrix A in the ascending order counting multiplicities so that $P_1 := A_{11} = \langle \psi_1, V^\square \psi_1 \rangle$ is the minimal eigenvalue and $P_p := A_{pp} = \langle \psi_p, V^\square \psi_p \rangle$ is the maximal eigenvalue of A .

5.1. Sign-changing random variables. In this subsection we assume (HC) to hold. We will only treat this case in detail as the calculation for positive random variables is very similar. Recall from (HC) that $s_- < 0 < s_+$. We define the following quantities :

$$(15) \quad A_1 := \inf_{q \in S_\mu} \inf_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)}=1}} q \langle \psi, V^\square \psi \rangle = \min(s_+ P_1, s_- P_p) \leq 0,$$

and

$$(16) \quad A_2 := -\max(s_-^2, s_+^2) \sup_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)}=1}} \sup_{\substack{\varphi \in \mathcal{V}_0^\perp \\ \|\varphi\|_{\ell^2(\square)}=1}} \frac{|\langle \psi, V^\square \varphi \rangle|^2}{\langle H_0^\square(\theta) \varphi, \varphi \rangle} \leq 0.$$

Note that the sign of A_1 and A_2 is fixed. We will prove the following theorem, which is only a restatement of Theorem 3.1.

Theorem 5.1. *Assume (HA), (HB) and (HC). Fix $\theta \in \Theta$. Then, if $A_1 \neq 0$,*

$$E_\epsilon \leq \epsilon A_1,$$

whereas if $A_1 = 0$, then

$$E_\epsilon \leq \epsilon^2 A_2 + O(\epsilon^3).$$

Finally, if $A_1 = A_2 = 0$, then

$$E_\epsilon \leq 0.$$

Remark 5.2.

- We remind that we have fixed θ to simplify notations, but A_1 and A_2 depend on θ . The best bound for the behavior of the bottom of the spectrum is obtained by looking at each $\theta \in \Theta$ and taking the minimum.
- We see that our bound on the bottom of the spectrum behaves linearly, quadratically or it doesn't move with ϵ . In the analogous setting in continuum space, if the unique continuation principle is not violated, then the analogous result does not allow the third case $A_1 = A_2 = 0$. This leaves only the cases of a linear or a quadratic bound.

- The definition of the quantities A_1, A_2 may seem complicated at first sight, but these choices are optimal, in the sense of Lemma 5.7 below, which is a converse of Lemma 5.6 in the regime $\epsilon \ll 1$.

Before proving the theorem, let us provide a much simpler, non-optimal upper bound for A_2 as well as a condition ensuring that $|A_1| + |A_2| \neq 0$.

5.2. A simple non-degeneracy condition. Theorem 5.1 tells us that if $A_2 \neq 0$, then the expansion of the bottom of the spectrum is at least quadratic, but if $A_1 = A_2 = 0$, we can only say that the spectrum starts at zero. When V^\square is diagonal this only happens if the support of the single-cell potential and the eigenfunctions ψ_1, \dots, ψ_p are disjoint (the ψ_i were defined at the beginning of this section).

Note, that in the continuous configuration space this can only happen if the potential violates the unique continuation principle. For a discussion on the validity of the unique continuation principle see for instance [24].

Let us discuss the condition in our general setting. First let us remark that if $A_1 = 0$, then the matrix $A \in \mathbb{C}^{p \times p}$ vanishes identically, i.e.

$$(17) \quad A_1 = 0 \implies \sup_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)}=1}} |\langle \psi, V^\square \psi \rangle| = 0.$$

and thus

$$(\forall \psi \in \mathcal{V}_0) \quad V^\square \psi \in \mathcal{V}_0^\perp.$$

The operator H_0^\square is invertible on \mathcal{V}_0^\perp and thus there exists some $\varphi \in \mathcal{V}_0^\perp$ such that

$$(18) \quad H_0^\square(\theta)\varphi = V^\square \psi^*.$$

Hence, we have that

$$\langle V^\square \varphi, \psi^* \rangle = \langle H_0^\square(\theta)\varphi, \varphi \rangle.$$

Now, assume there exists some $\psi^* \in \mathcal{V}_0$ such that

$$(19) \quad V^\square \psi^* \neq 0.$$

Then φ in (18) does not vanish and

$$A_2 \leq -\max(s_-^2, s_+^2) \frac{|\langle V^\square \varphi, \psi^* \rangle|^2}{\langle H_0^\square(\theta)\varphi, \varphi \rangle} = -\max(s_-^2, s_+^2) \langle H_0^\square(\theta)\varphi, \varphi \rangle < 0,$$

because $\varphi \notin \ker H_0^\square(\theta)$.

Remark 5.3. Formally, we have

$$A_2 \leq -\max(s_-^2, s_+^2) \langle \psi^*, V^\square H_0^\square(\theta)^{-1} V^\square \psi^* \rangle$$

when $A_1 = 0$.

In the converse direction, $A_1 = 0$ together with $A_2 = 0$ implies that

$$(\forall \psi \in \mathcal{V}_0 \text{ and } \forall \varphi \in \ell^2(\square)) \quad \langle V^\square \psi, \varphi \rangle = 0,$$

i.e. that

$$(\forall \psi \in \mathcal{V}_0) \quad V^\square \psi = 0.$$

We summarize the above discussion as follows.

Lemma 5.4. *Under the assumptions of Theorem 5.1 we have that*

$$A_1 = 0 \text{ and } A_2 = 0 \quad \text{if and only if} \quad (\forall \psi^* \in \mathcal{V}_0) \quad V^\square \psi^* = 0.$$

5.3. Proof of Theorem 5.1. We subdivide the proof of Theorem 5.1 into two lemmas. The first covers both types of sign assumptions on the random variables.

Lemma 5.5. *Assume (HA), (HB), and either (HC) or (HC'). Let $u \in \ell^2(\square)$ and E_ϵ as in (7). Then,*

$$E_\epsilon \leq \inf_{q \in S_\mu} \inf_{u \in \ell^2(\square)} \frac{\langle H_{\epsilon,q}^\square(\theta)u, u \rangle}{\|u\|_{\ell^2(\square)}} \quad \text{for any } \theta \in \square^*$$

Proof. By Corollary 4.3 it is enough to consider the periodic realizations of the potential. By the Courant–Weyl–Fischer min–max principle,

$$(20) \quad E_\epsilon \leq E_{q,\epsilon} = \min \sigma(H_{\epsilon,q}) = \inf_{\substack{a \in \ell^2(\mathbb{Z}^d) \\ \|a\|_2=1}} \langle H_{\epsilon,q} a, a \rangle.$$

Finally, by Lemma 4.1,

$$(21) \quad \inf_{\substack{a \in \ell^2(\mathbb{Z}^d) \\ \|a\|=1}} \langle H_{\epsilon,q} a, a \rangle \leq \inf_{u \in \ell^2(\square)} \frac{\langle H_{\epsilon,q}^\square(\theta)u, u \rangle}{\|u\|_{\ell^2(\square)}}.$$

This proves the lemma. \square

We state now the second lemma. It applies to the case of sign-changing random variables.

Lemma 5.6. *Let A_1, A_2 as in (10), (11), assume (HA), (HB), and (HC) and fix $\theta \in \Theta$. Then, for $\epsilon > 0$ small enough, if $A_1 \neq 0$,*

$$\inf_{q \in S_\mu} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle \leq \epsilon A_1,$$

whereas if $A_1 = 0$, then

$$\inf_{q \in S_\mu} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle \leq \epsilon^2 A_2 + O(\epsilon^3)$$

Finally, if $A_1 = A_2 = 0$, then

$$\inf_{q \in S_\mu} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle \leq 0.$$

Proof. It is enough to show that for some $q \in S_\mu$, there is some normalized state $u \in \ell^2(\square)$ satisfying

$$\langle H_{\epsilon,q}^\square(\theta)u, u \rangle \leq \epsilon A_1 \quad \text{or} \quad \epsilon^2 A_2 + O(\epsilon^3) \quad \text{or} \quad 0 \quad \text{resp.}$$

Let $\psi \in \mathcal{V}_0$ and $\varphi \in \mathcal{V}_0^\perp$, to be chosen later, and $u = \psi + \epsilon q \varphi$. We assume furthermore $\|\psi\| = 1$. We expand

$$\|u\|^2 = \|\psi\|^2 + \epsilon^2 q^2 \|\varphi\|^2$$

and thus

$$(22) \quad 1/\|u\|^2 = 1 - \epsilon^2 q^2 \|\varphi\|^2 + O(\epsilon^4 \|\varphi\|^4).$$

We calculate the kinetic energy of this state, i.e.

$$(23) \quad \langle H_0^\square(\theta)u, u \rangle = \langle H_0^\square(\theta)\psi, \psi \rangle + 2\epsilon q \operatorname{Re} \langle H_0^\square(\theta)\psi, \varphi \rangle + \epsilon^2 q^2 \langle H_0^\square(\theta)\varphi, \varphi \rangle.$$

Because $\psi \in \mathcal{V}_0$ and $E_0 = 0$, we see that (23) becomes

$$\langle H_0^\square(\theta)u, u \rangle = \epsilon^2 q^2 \langle H_0^\square(\theta)\varphi, \varphi \rangle.$$

We expand the potential energy as

$$\epsilon q \langle V^\square u, u \rangle = \epsilon q \langle V^\square \psi, \psi \rangle + 2\epsilon^2 q^2 \operatorname{Re} \langle V^\square \varphi, \psi \rangle + \epsilon^3 q^3 \langle V^\square \varphi, \varphi \rangle.$$

Thus,

$$(24) \quad \begin{aligned} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle &= \epsilon q \langle V^\square \psi, \psi \rangle + \epsilon^2 q^2 \left(\langle H_0^\square(\theta)\varphi, \varphi \rangle + 2 \operatorname{Re} \langle \psi, V^\square \varphi \rangle \right) \\ &+ \epsilon^3 q^3 \langle V^\square \varphi, \varphi \rangle. \end{aligned}$$

Case $A_1 \neq 0$. Note that in this case $P_1 P_p \neq 0$. From now on we assume that $s_+ P_1 \leq s_- P_p$. If this is not the case, we can always replace $V^\square \mapsto -V^\square$ and $\omega_n \mapsto -\omega_n$ to get an equivalent model. In this case, we take $\psi = \psi_1$, $\varphi = 0$ and $q = s_+$. Then, (24) becomes

$$\langle H_{\epsilon,q}^\square(\theta)u, u \rangle = \epsilon q \langle V^\square \psi_1, \psi_1 \rangle = \epsilon s_+ P_1,$$

which proves the result in this case, as u is normalized.

Case $A_1 = 0$ and $A_2 \neq 0$. First let us remark that if $A_1 = 0$ then the matrix $A \in \mathbb{C}^{p \times p}$ vanishes identically, i.e.

$$(25) \quad A_1 = 0 \implies \sup_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)} = 1}} |\langle \psi, V^\square \psi \rangle| = 0.$$

In this case we have that, for any $\psi \in \mathcal{V}_0$ and $\varphi \in \mathcal{V}_0^\perp$, the expansion (24) becomes

$$(26) \quad \langle H_{\epsilon,q}^\square(\theta)u, u \rangle = \epsilon^2 q^2 \left(\langle H_0^\square(\theta)\varphi, \varphi \rangle + 2 \operatorname{Re} \langle \psi, V^\square \varphi \rangle \right) + \epsilon^3 q^3 \langle V^\square \varphi, \varphi \rangle.$$

Note that, for $\psi \in \mathcal{V}_0$ and $\varphi \in \mathcal{V}_0^\perp$ such that

$$\|\psi\|_2 = \|\varphi\|_2 = 1$$

the map

$$(\varphi, \psi) \mapsto \frac{|\langle \psi, V^\square \varphi \rangle|^2}{\langle H_0^\square(\theta) \varphi, \varphi \rangle}$$

is continuous. Given that the spaces involved are finite-dimensional and their respective unit balls thus compact, we know that there exists a couple (ψ^*, φ^*) maximizing this quantity, i.e.

$$\frac{|\langle \psi^*, V^\square \varphi^* \rangle|^2}{\langle H_0^\square(\theta) \varphi^*, \varphi^* \rangle} = \sup_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)}=1}} \sup_{\substack{\varphi \in \mathcal{V}_0^\perp \\ \|\varphi\|_{\ell^2(\square)}=1}} \frac{|\langle \psi, V^\square \varphi \rangle|^2}{\langle H_0^\square(\theta) \varphi, \varphi \rangle}.$$

Let $\psi = \psi^*$ and $\varphi = \lambda \varphi^*$ in the definition of u , where

$$\lambda = -\frac{\langle \psi^*, V^\square \varphi^* \rangle}{\langle H_0^\square(\theta) \varphi^*, \varphi^* \rangle} \in \mathbb{C}.$$

Replacing, we see that

$$\begin{aligned} \langle H_0^\square(\theta) \varphi, \varphi \rangle + 2\operatorname{Re} \langle \psi, V^\square \varphi \rangle &= |\lambda|^2 \langle H_0^\square(\theta) \varphi^*, \varphi^* \rangle + 2\operatorname{Re} \bar{\lambda} \langle \psi^*, V^\square \varphi^* \rangle \\ &= \frac{|\langle \psi^*, V^\square \varphi^* \rangle|^2}{\langle H_0^\square(\theta) \varphi^*, \varphi^* \rangle} - 2 \frac{|\langle \psi^*, V^\square \varphi^* \rangle|^2}{\langle H_0^\square(\theta) \varphi^*, \varphi^* \rangle} \\ &= -\frac{|\langle \psi^*, V^\square \varphi^* \rangle|^2}{\langle H_0^\square(\theta) \varphi^*, \varphi^* \rangle}. \end{aligned}$$

Using this in (26) and letting $q^2 = \max(s_-^2, s_+^2)$, we obtain

$$\begin{aligned} \langle H_{\epsilon, q}^\square(\theta) u, u \rangle &= -\epsilon^2 \max(s_-^2, s_+^2) \frac{|\langle \psi^*, V^\square \varphi^* \rangle|^2}{\langle H_0^\square(\theta) \varphi^*, \varphi^* \rangle} + O(\epsilon^3 q^3 \|\varphi^*\|^2) \\ &= \epsilon^2 A_2 + O(\epsilon^3 q^3 \|\varphi\|^2). \end{aligned}$$

Normalizing u by multiplying by (22) gives the result.

Case $A_1 = 0$ and $A_2 = 0$. Choose $\varphi = 0$ and any normalized $\psi \in \mathcal{V}_0$. The development using u in this case gives

$$\langle H_{\epsilon, q}^\square(\theta) u, u \rangle = \epsilon^3 q^3 \langle V^\square \varphi, \varphi \rangle = 0$$

and this yields the desired result. \square

We prove the following converse lemma.

Lemma 5.7. *Let A_1, A_2 as in (10), (11), assume (HA), (HB) and (HC), and fix $\theta \in \Theta$. Then, for $\epsilon > 0$ small enough, if $A_1 \neq 0$,*

$$\inf_{q \in S_\mu} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon, q}^\square(\theta) u, u \rangle \geq \epsilon A_1 + O(\epsilon^{3/2}),$$

whereas if $A_1 = 0$, then

$$\inf_{q \in S_\mu} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon, q}^\square(\theta) u, u \rangle \geq \epsilon^2 A_2 + O(\epsilon^3)$$

Finally, if $A_1 = A_2 = 0$, then

$$\inf_{q \in S_\mu} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle \geq 0.$$

Proof. Fix $\epsilon > 0$ and let $q \in S_\mu$ be the value which minimizes the map

$$q \mapsto \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle,$$

i.e. $q \in \{s_-, s_+\}$. We lower bound this quantity by minimizing over a larger set by writing

$$\inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle \geq \inf_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)} \leq 1}} \inf_{\substack{\varphi \in \mathcal{V}_0^\perp \\ \|\varphi\|_{\ell^2(\square)} \leq 1}} \langle H_{\epsilon,q}^\square(\theta)(\psi + \varphi), (\psi + \varphi) \rangle.$$

By continuity and compactness, there exists some pair (ψ^*, φ^*) in $\mathcal{V}_0 \times \mathcal{V}_0^\perp$ realizing the infimum on the right hand side. We see that

$$(27) \quad \langle H_0^\square(\theta)(\psi^* + \varphi^*), (\psi^* + \varphi^*) \rangle = \langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle \geq g\|\varphi^*\|_{\ell^2(\square)}^2,$$

where the constant g is the spectral gap of H_0^\square . Due to our normalization g coincides with the (positive) second eigenvalue of H_0^\square . We study the different cases.

Case $A_1 \neq 0$. From Lemma 5.6, we know already that

$$(28) \quad |A_1| + |A_2| \neq 0 \implies \langle H_{\epsilon,q}^\square(\theta)(\psi^* + \varphi^*), (\psi^* + \varphi^*) \rangle < 0.$$

Using (27) and (28) we get that

$$\|\varphi^*\|_{\ell^2(\square)}^2 \leq -g^{-1}\epsilon q \langle V^\square(\psi^* + \varphi^*), (\psi^* + \varphi^*) \rangle \leq 4g^{-1}\|V^\square\|_\infty \epsilon q.$$

We deduce then that

$$\begin{aligned} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle &\geq \epsilon q \langle V^\square \psi^*, \psi^* \rangle + 2\epsilon q \operatorname{Re} \langle V^\square \varphi^*, \psi^* \rangle + \epsilon q \langle V^\square \varphi^*, \varphi^* \rangle \\ &\geq \epsilon A_1 - 4g^{-1/2}\epsilon^{3/2}q^{3/2}\|V^\square\|_\infty^{3/2} - 4g^{-1}\epsilon^2q^2\|V^\square\|_\infty^2. \end{aligned}$$

Case $A_1 = 0$ and $A_2 \neq 0$. In this case

$$\begin{aligned} \langle H_{\epsilon,q}^\square(\theta)(\psi^* + \varphi^*), (\psi^* + \varphi^*) \rangle &= \langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle + 2\epsilon q \operatorname{Re} \langle V^\square \varphi^*, \psi^* \rangle + \epsilon q \langle V^\square \varphi^*, \varphi^* \rangle. \end{aligned}$$

Using (28) we see that $\varphi^* \neq 0$. Furthermore, (27) and (28) together imply that

$$\begin{aligned} \|\varphi^*\|_{\ell^2(\square)}^2 &\leq \epsilon q g^{-1}\|V^\square\|_\infty (2\|\varphi^*\|_{\ell^2(\square)} + \|\varphi^*\|_{\ell^2(\square)}^2) \\ &\leq 3\epsilon q g^{-1}\|V^\square\|_\infty \|\varphi^*\|_{\ell^2(\square)}. \end{aligned}$$

Note that $\|\varphi^*\|_{\ell^2(\square)}$ is on both sides of the inequality. Simplifying,

$$(29) \quad \|\varphi^*\|_{\ell^2(\square)} \leq 3\epsilon q g^{-1}\|V^\square\|_\infty.$$

Expanding as $\epsilon \rightarrow 0$, employing (29) and then simply multiplying by $1 = |\lambda|^2/|\lambda|^2 = \bar{\lambda}/\lambda$, we write

$$\begin{aligned} & \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle \\ & \geq \langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle + 2 \operatorname{Re} \epsilon q \langle V^\square \psi^*, \varphi^* \rangle + \epsilon q \langle V^\square \varphi^*, \varphi^* \rangle \\ & = |\lambda|^2 \frac{\langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle}{|\lambda|^2} + 2 \operatorname{Re} \bar{\lambda} \epsilon q \frac{\langle V^\square \psi^*, \varphi^* \rangle}{\bar{\lambda}} + O(\epsilon^3) \end{aligned}$$

We choose λ as

$$\lambda = -\frac{\langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle}{\langle V^\square \psi^*, \varphi^* \rangle}.$$

It is well defined for small ϵ . Indeed, using (27) and (28) we see that

$$\begin{aligned} -2\epsilon q \operatorname{Re} \langle V^\square \varphi^*, \psi^* \rangle & \geq \langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle + \epsilon q \langle V^\square \varphi^*, \varphi^* \rangle \\ & \geq c\|\varphi^*\|_{\ell^2(\square)}^2 - \epsilon q \|V^\square\|_\infty \|\varphi^*\|_{\ell^2(\square)}^2 \end{aligned}$$

and we know that $\varphi^* \neq 0$. We see that $\lambda \neq 0$ for ϵ small enough.

Using our choice of λ gives

$$\begin{aligned} & c\|\varphi^*\|_{\ell^2(\square)}^2 - \epsilon q \|V^\square\|_\infty \|\varphi^*\|_{\ell^2(\square)}^2 \\ & = (|\lambda|^2 - 2 \operatorname{Re} \bar{\lambda} \epsilon q) \frac{|\langle V^\square \psi^*, \varphi^* \rangle|^2}{\langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle} + O(\epsilon^3) \\ & \geq -\epsilon^2 A_2 + O(\epsilon^3), \end{aligned}$$

where we have used that $|\lambda|^2 - 2 \operatorname{Re} \bar{\lambda} \epsilon q \geq |\lambda|^2 - 2|\lambda|\epsilon q \geq -\epsilon^2 q^2$.

Case $A_1 = 0$ and $A_2 = 0$. In this case

$$(30) \quad 0 \geq \inf_{\|u\|_{\ell^2(\square)}} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle = \langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle + \epsilon q \langle V^\square \varphi^*, \varphi^* \rangle$$

$$(31) \quad \geq c\|\varphi^*\|^2 - O(\epsilon)\|\varphi^*\|^2,$$

where the first inequality comes from Lemma 5.6. It is now clear that $\psi^* = 0$ and this finishes the proof. \square

5.4. Positive random variables. We study in this subsection the case involving positive random variables. We remind the reader the definition of the constants involved, for which we use the functions ψ_i , the matrix A , its eigenvalues P_i and the linear space \mathcal{V}_0 , which can be found at the beginning of this section. We define the subspace $\mathcal{V}_{01} \subset \mathcal{V}_0$ as

$$\mathcal{V}_{01} := \operatorname{span}_{\{i:P_i=P_1\}} \langle \psi_i \rangle,$$

i.e. the eigenspace of A associated to its minimal eigenvalue P_1 .

We recall the following quantities :

$$(32) \quad A'_1 := \inf_{q \in S_\mu} \inf_{\substack{\psi \in \mathcal{V}_{01} \\ \|\psi\|_{\ell^2(\square)}=1}} q \langle \psi, V^\square \psi \rangle = \min(s_+ P_1, s_- P_1) \in \mathbb{R},$$

and

$$(33) \quad A'_2 := -s_+^2 \sup_{\substack{\psi \in \mathcal{V}_{01} \\ \|\psi\|_{\ell^2(\square)}=1}} \sup_{\substack{\varphi \in \mathcal{V}_0^\perp \\ \|\varphi\|_{\ell^2(\square)}=1}} \frac{|\langle \psi, V^\square \varphi \rangle|^2}{\langle H_0^\square(\theta) \varphi, \varphi \rangle} \leq 0.$$

Note that, unlike the other coefficients, A'_1 may take on both signs. We also restate Theorem 3.2 for the reader's convenience.

Theorem 5.8. *Assume (HA), (HC) and (HC'). Fix $\theta \in \Theta$. Then, for $\epsilon > 0$ small enough, if $P_1 \neq 0$,*

$$E_\epsilon \leq \epsilon A'_1,$$

whereas if $P_1 = 0$, then

$$E_\epsilon \leq \epsilon^2 A'_2 + O(\epsilon^3).$$

Finally, if $P_1 = A'_2 = 0$, then

$$E_\epsilon \leq 0.$$

Sketch of proof. The proof of this theorem is very similar to the proof of Theorem 5.1. Indeed, Lemma 5.5 is also valid in this setting. We proceed then as in Lemma 5.6 up to equation (24). If $A'_1 \neq 0$ we let

$$u = \psi_1, \quad \varphi = 0, \quad q = \begin{cases} s_+ & \text{if } P_1 < 0 \\ s_- & \text{if } P_1 \geq 0 \end{cases}$$

in (24). If $A'_1 = 0$ but $A'_2 \neq 0$, then we find $\psi^* \in \mathcal{V}_{01}$ and $\varphi \in \mathcal{V}_0$ realizing the supremum in the definition of A'_2 and then we proceed as in Lemma 5.6. Finally, if $A'_1 = A'_2 = 0$ we take $u = \psi_1$ in (24). \square

We prove the following converse lemma.

Lemma 5.9. *Let A'_1, A'_2 as in (12), (13), assume (HA), (HB) and (HC'), and fix $\theta \in \Theta$. Then, for $\epsilon > 0$ small enough, if $A'_1 \neq 0$,*

$$\inf_{q \in S_\mu} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta) u, u \rangle \geq \epsilon A'_1 + O(\epsilon^{3/2}),$$

whereas if $P_1 = 0$, then

$$\inf_{q \in S_\mu} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta) u, u \rangle \geq \epsilon^2 A'_2 + O(\epsilon^3)$$

Finally, if $P_1 = A'_2 = 0$, then

$$\inf_{q \in S_\mu} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta) u, u \rangle \geq 0.$$

Proof. We adapt here the proof of Lemma 5.7.

Fix $\epsilon > 0$ and let $q \in S_\mu$ be the value which minimizes the map

$$q \mapsto \inf_{\|u\|_{\ell^2(\square)}=1} \langle (H_{\epsilon,q}^\square(\theta) - \epsilon A'_1)u, u \rangle,$$

i.e. $q \in \{s_-, s_+\}$. We lower bound this quantity by minimizing over a larger set by writing

$$\begin{aligned} & \inf_{\|u\|_{\ell^2(\square)}=1} \langle (H_{\epsilon,q}^\square(\theta) - \epsilon A'_1)u, u \rangle \\ & \geq \inf_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)} \leq 1}} \inf_{\substack{\varphi \in \mathcal{V}_0^\perp \\ \|\varphi\|_{\ell^2(\square)} \leq 1}} \langle (H_{\epsilon,q}^\square(\theta) - \epsilon A'_1)(\psi + \varphi), (\psi + \varphi) \rangle. \end{aligned}$$

By continuity and compactness, there exists some pair (ψ^*, φ^*) in $\mathcal{V}_0 \times \mathcal{V}_0^\perp$ realizing the infimum on the right hand side. We see that

$$(34) \quad \langle H_0^\square(\theta)(\psi^* + \varphi^*), (\psi^* + \varphi^*) \rangle = \langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle \geq g \|\varphi^*\|_{\ell^2(\square)}^2,$$

where the constant g is the spectral gap of H_0^\square , which is also its (positive) second eigenvalue.

We study the different cases.

Case $A'_1 \neq 0$. We know already that

$$(35) \quad \langle (H_{\epsilon,q}^\square(\theta) - \epsilon A'_1)(\psi^* + \varphi^*), (\psi^* + \varphi^*) \rangle \leq 0.$$

Using (34), (35) and $|A'_1| \leq q \|V^\square\|$ we get

$$\|\varphi^*\|_{\ell^2(\square)}^2 \leq -g^{-1}\epsilon \langle (qV^\square - A'_1)(\psi^* + \varphi^*), (\psi^* + \varphi^*) \rangle \leq 4g^{-1} \|V^\square\|_\infty \epsilon q.$$

Note that

$$q \langle V^\square \psi^*, \psi^* \rangle \geq q \inf_{\substack{\psi \in \mathcal{V}_0 \\ \|\psi\|_{\ell^2(\square)} \leq 1}} \langle V^\square \psi, \psi \rangle = q \inf_{\substack{\psi \in \mathcal{V}_{00} \\ \|\psi\|_{\ell^2(\square)} \leq 1}} \langle V^\square \psi, \psi \rangle \geq E'_1.$$

It implies

$$\begin{aligned} & \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle \\ & \geq \epsilon q \langle V^\square \psi^*, \psi^* \rangle + 2\epsilon q \operatorname{Re} \langle V^\square \varphi^*, \psi^* \rangle + \epsilon q \langle V^\square \varphi^*, \varphi^* \rangle \\ & \geq \epsilon A'_1 - 4g^{-1/2} \epsilon^{3/2} q^{3/2} \|V^\square\|_\infty^{3/2} - 4g^{-1} \epsilon^2 q^2 \|V^\square\|_\infty^2. \end{aligned}$$

Case $A'_1 = 0$ and $A'_2 \neq 0$. We know that

$$(36) \quad \langle H_{\epsilon,q}^\square(\theta)(\psi^* + \varphi^*), (\psi^* + \varphi^*) \rangle < 0.$$

We will decompose further $\psi^* = \psi_{01}^* + \psi_{0\perp}^* \in \mathcal{V}_0$, with $\psi_{01}^* \in \mathcal{V}_{01}$ and $\psi_{0\perp}^* \in \mathcal{V}_{01}^\perp$. We have

$$(37) \quad \langle V^\square u, u \rangle \geq 2 \operatorname{Re} \langle V^\square \psi_{01}^*, \varphi^* \rangle + O(\|\varphi^*\|^2)$$

Indeed, using the definition of u , we see that,

$$\begin{aligned}
\langle V^\square u, u \rangle &= \langle V^\square \psi_{01}^*, \psi_{01}^* \rangle + 2 \operatorname{Re} \langle V^\square \psi_{01}^*, \psi_{0\perp}^* \rangle + \langle V^\square \psi_{0\perp}^*, \psi_{0\perp}^* \rangle \\
&\quad + 2 \operatorname{Re} \langle V^\square \psi_{01}^*, \varphi^* \rangle + 2 \operatorname{Re} \langle V^\square \psi_{0\perp}^*, \varphi^* \rangle + \langle V^\square \varphi^*, \varphi^* \rangle \\
(38) \quad &= \langle V^\square \psi_{0\perp}^*, \psi_{0\perp}^* \rangle + 2 \operatorname{Re} \langle V^\square \psi_{01}^*, \varphi^* \rangle \\
&\quad + 2 \operatorname{Re} \langle V^\square \psi_{0\perp}^*, \varphi^* \rangle + \langle V^\square \varphi^*, \varphi^* \rangle.
\end{aligned}$$

Let us also note that

$$A_1 = 0 \implies P_1 \geq 0,$$

and, in particular

$$\langle V^\square \psi_{0\perp}^*, \psi_{0\perp}^* \rangle \geq 0.$$

Hence if $\langle V^\square \psi_{0\perp}^*, \varphi^* \rangle = 0$, we immediately get (37) from (38).

Assume that $\langle V^\square \psi_{0\perp}^*, \varphi^* \rangle \neq 0$. For each $\mu \in \mathbb{C}$,

$$\begin{aligned}
(38) &= |\mu|^2 \frac{\langle V^\square \psi_{0\perp}^*, \psi_{0\perp}^* \rangle}{|\mu|^2} + 2 \operatorname{Re} \bar{\mu} \frac{\langle V^\square \psi_{0\perp}^*, \varphi^* \rangle}{\bar{\mu}} \\
&\quad + 2 \operatorname{Re} \langle V^\square \psi_{01}^*, \varphi^* \rangle + O(\|\varphi^*\|^2).
\end{aligned}$$

We choose μ as

$$\mu := -\frac{\langle V^\square \psi_{0\perp}^*, \psi_{0\perp}^* \rangle}{\langle V^\square \psi_{0\perp}^*, \varphi^* \rangle}.$$

We obtain

$$\begin{aligned}
(38) &= (|\mu|^2 - 2 \operatorname{Re} \bar{\mu}) \frac{|\langle V^\square \psi_{0\perp}^*, \varphi^* \rangle|^2}{\langle V^\square \psi_{0\perp}^*, \psi_{0\perp}^* \rangle} + 2 \operatorname{Re} \langle V^\square \psi_{01}^*, \varphi^* \rangle + O(\|\varphi^*\|^2) \\
&\geq -\frac{1}{2} \frac{|\langle V^\square \psi_{0\perp}^*, \varphi^* \rangle|^2}{\langle V^\square \psi_{0\perp}^*, \psi_{0\perp}^* \rangle} + 2 \operatorname{Re} \langle V^\square \psi_{01}^*, \varphi^* \rangle + O(\|\varphi^*\|^2) \\
&= 2 \operatorname{Re} \langle V^\square \psi_{01}^*, \varphi^* \rangle + O(\|\varphi^*\|^2),
\end{aligned}$$

which is also (37).

Using this result, we obtain

$$\begin{aligned}
&\langle H_{\epsilon, q}^\square(\theta)(\psi^* + \varphi^*), (\psi^* + \varphi^*) \rangle \\
&\geq \langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle + 2\epsilon q \operatorname{Re} \langle V^\square \psi_{01}^*, \varphi^* \rangle + O(\epsilon q \|\varphi^*\|^2)
\end{aligned}$$

Equality (38) and inequalities (34) and (36) together imply that

$$\begin{aligned}
\|\varphi^*\|_{\ell^2(\square)}^2 &\leq \epsilon q g^{-1} \|V^\square\|_\infty (2\|\varphi^*\|_{\ell^2(\square)} + \|\varphi^*\|_{\ell^2(\square)}^2) \\
&\leq 3\epsilon q g^{-1} \|V^\square\|_\infty \|\varphi^*\|_{\ell^2(\square)}.
\end{aligned}$$

Note that $\|\varphi^*\|_{\ell^2(\square)}$ is involved in both sides of the inequality. Simplifying,

$$(39) \quad \|\varphi^*\|_{\ell^2(\square)} \leq 3\epsilon q g^{-1} \|V^\square\|_\infty.$$

Using (39) and then simply multiplying by $1 = |\lambda|^2/|\lambda|^2 = \bar{\lambda}/\lambda$, we get

$$\begin{aligned} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle \\ \geq |\lambda|^2 \frac{\langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle}{|\lambda|^2} + 2 \operatorname{Re} \bar{\lambda} \epsilon q \frac{\langle V^\square \psi_{01}^*, \varphi^* \rangle}{\bar{\lambda}} + O(\epsilon^3) \end{aligned}$$

In view of this inequality and (36), we see that $\langle V^\square \psi_{01}^*, \varphi^* \rangle \neq 0$ for small ϵ . We choose λ as

$$\lambda = -\frac{\langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle}{\langle V^\square \psi_{01}^*, \varphi^* \rangle}.$$

It is well defined for small ϵ . Using our choice of λ , it gives

$$\begin{aligned} \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle &\geq (|\lambda|^2 - 2 \operatorname{Re} \bar{\lambda} \epsilon q) \frac{|\langle V^\square \psi_{01}^*, \varphi^* \rangle|^2}{\langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle} + O(\epsilon^3) \\ &\geq -\epsilon^2 A'_2 + O(\epsilon^3), \end{aligned}$$

where we have used that $|\lambda|^2 - 2 \operatorname{Re} \bar{\lambda} \epsilon q \geq |\lambda|^2 - 2|\lambda|\epsilon q \geq -\epsilon^2 q^2$.

Case $A_1 = 0$ and $A_2 = 0$. In this case, re-using (37) and employing $A'_2 = 0$,

$$\begin{aligned} 0 &\geq \inf_{\|u\|_{\ell^2(\square)}=1} \langle H_{\epsilon,q}^\square(\theta)u, u \rangle \\ &= \langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle + 2\epsilon q \operatorname{Re} \langle V^\square \psi_{01}^*, \varphi^* \rangle + \epsilon q O(\|\varphi^*\|^2) \\ &= \langle H_0^\square(\theta)\varphi^*, \varphi^* \rangle + \epsilon q O(\|\varphi^*\|^2) \\ &\geq c\|\varphi^*\|^2 - |O(\epsilon)|\|\varphi^*\|^2, \end{aligned}$$

where the first inequality comes from Lemma 5.6. It is now clear that $\psi^* = 0$ and this finishes the proof. \square

5.5. Application to the discrete alloy type model.

Proof of Theorem 1.1. It is enough to verify that the assumptions of Theorem 5.1 are satisfied. Let $C_W := \inf \sigma(-\Delta_{\mathbb{Z}^d} + W)$. It is clear that the operator $H_0 := -\Delta_{\mathbb{Z}^d} + W - C_W$ satisfies hypothesis (HA). In this case the set Θ consists of the single point $\theta = 0$ (see Theorem 5.10 below).

Let us check property (19) for the operator

$$H_0(0) := -\Delta_\square + W_\square + C_W,$$

where Δ_\square is the Laplacian on \square with periodic boundary conditions and W_\square is the restriction of W to \square . To check the property we use Perron–Frobenius theorem [18]. For $m > |\square|$, we verify that

$$\langle \delta_x, (\Delta_\square - W_\square + \|W\|_\infty + 2d + 1)^m \delta_y \rangle \geq \langle \delta_x, (\Delta_\square + 2d + 1)^m \delta_y \rangle \geq 1.$$

This implies that the largest eigenvalue of the matrix $\Delta_\square - W_\square + \|W\|_\infty + 2d + 1$ is simple and its corresponding eigenfunction ψ_1 is

positive (i.e. $(\forall n \in \square) \psi_1(n) > 0$). Because of this strict positivity, condition (19) is satisfied as soon as $V^\square \neq 0$. The subspace \mathcal{V}_0 is thus one-dimensional and contains only ψ_1 . The theorem is now proven, by simply stating the consequences of Theorem 5.1. \square

We now recall Theorem 2.4 in [13], with our notations. It implies that 0 is the unique $\theta \in \square^*$ realizing the minimum of the spectrum.

Theorem 5.10. *Let $H_0 = -\Delta_{\mathbb{Z}^d} + W$ with W a periodic potential with respect to $\gamma = N\mathbb{Z}^d$, and $E_0(\theta)$ be the smallest eigenvalue of $H_0(\theta)$. Then*

$$(a_-/a_+)^2 \left(2d - \sum_{i=1}^d \cos(\theta_i) \right) \leq E_0(\theta) - E_0(0) \leq \left(2d - \sum_{i=1}^d \cos(\theta_i) \right).$$

Here, $a_\pm = \pm \max \pm \psi_1$ and ψ_1 is the positive ground state of $H_0(0)$.

The band structure of the spectrum of discrete alloy type models has been analyzed in [8]

6. APPENDIX

6.1. An interesting example: Proof of Theorem 1.2. Let $H_0 := \Delta_{\mathbb{Z}}^2$ defined on $\ell^2(\mathbb{Z})$. This operator has hopping range $N = 3$ (see (HA)) and thus $\square = \{-1, 0, 1, \}$. We define V^\square as the multiplication operator given by the following single-site potential:

$$\begin{aligned} V^\square : \ell^2(\square) &\rightarrow \mathbb{R} \\ V^\square(n) &:= -\frac{1}{2}\delta_{-1}(n) + \delta_0(n) - \frac{1}{2}\delta_1(n). \end{aligned}$$

With these definitions, we see that, for $\theta \in [-\pi/3, \pi/3]^d$,

$$H_0^\square(\theta) = \begin{pmatrix} 6 & -4 + e^{-3i\theta} & 1 - 4e^{-3i\theta} \\ -4 + e^{3i\theta} & 6 & -4 + e^{-3i\theta} \\ 1 - 4e^{-3i\theta} & -4 + e^{3i\theta} & 6 \end{pmatrix},$$

after (8). This matrix has a simple ground state

$$\psi_0(\theta) := (e^{-i\theta}, 1, e^{i\theta})/\sqrt{3}$$

with eigenvalues $E_0(\theta) = (2 - 2\cos(\theta))^2$. Let now $\tilde{f}_n(\theta) = \chi_n \tilde{\psi}_0(\theta) \in \ell^2(\mathbb{Z}^d)$ where $\tilde{\psi}_0(\theta)$ is the θ -quasi- γ -periodic extension of $\psi_0(\theta)$. Finally, for $\xi > 1/4$, let

$$u_n := f_n(0) + \epsilon^\xi f_n(\epsilon^\xi).$$

Let us calculate the kinetic energy. We see that

$$\begin{aligned} (40) \quad \langle H_0 u_n, u_n \rangle &= \langle H_0 f_n(0), f_n(0) \rangle + 2\epsilon^\xi \operatorname{Re} \langle H_0 f_n(0), f_n(\epsilon^\xi) \rangle + \epsilon^{2\xi} \langle H_0 f_n(\epsilon^\xi), f_n(\epsilon^\xi) \rangle. \end{aligned}$$

Let $\delta > 0$ and pick n so large so that

$$\left| \frac{\langle f_n(0), H_0 f_n(0) \rangle_{\ell^2(D)}}{\|f_n(0)\|_{\ell^2(D)}^2} - \frac{\langle \psi_0(0), H_0^\square(\theta) \psi_0(0) \rangle_{\ell^2(\square)}}{\|\psi_0(0)\|_{\ell^2(\square)}^2} \right| \leq \delta,$$

$$\left| \frac{\langle f_n(\epsilon^\xi), H_0 f_n(\epsilon^\xi) \rangle_{\ell^2(D)}}{\|f_n(\epsilon^\xi)\|_{\ell^2(D)}^2} - \frac{\langle \psi_0(\epsilon^\xi), H_0^\square(\theta) \psi_0(\epsilon^\xi) \rangle_{\ell^2(\square)}}{\|\psi_0(\epsilon^\xi)\|_{\ell^2(\square)}^2} \right| \leq \delta,$$

and

$$\left| \frac{\|H_0 f_n(0)\|_{\ell^2(D)}^2}{\|f_n(0)\|_{\ell^2(D)}^2} \right| = \left| \frac{\langle f_n(0), H_0^2 f_n(0) \rangle_{\ell^2(D)}}{\|f_n(0)\|_{\ell^2(D)}^2} \right| \leq \delta.$$

Then, from (40) we see that

$$\begin{aligned} & \langle H_0 u_n, u_n \rangle \\ & \leq \delta \|f_n(0)\|^2 + 2\epsilon^\xi \delta \|f_n(\epsilon^\xi)\| + \epsilon^{2\xi} \langle \psi_0(\epsilon^\xi), H_0^\square(\theta) \psi_0(\epsilon^\xi) \rangle_{\ell^2(\square)} \|f_n(\theta)\|^2 \\ & \leq 3\delta + \epsilon^{2\xi} E_0(\epsilon^\xi) \|f_n(\theta)\|^2. \end{aligned}$$

Letting $n \rightarrow \infty$ and $\delta \rightarrow 0$ we see that

$$\langle H_0 u_n, u_n \rangle \leq \epsilon^{2\xi} E_0(\epsilon^\xi) \leq C \epsilon^{6\xi}.$$

Now let us calculate the potential energy.

$$\begin{aligned} & \epsilon \langle V_q u_n, u_n \rangle \\ & = \epsilon \langle V_q f_n(0), f_n(0) \rangle + 2\epsilon^{1+\xi} \operatorname{Re} \langle V_q f_n(0), f_n(\epsilon^\xi) \rangle + \epsilon^{1+2\xi} \langle V_q f_n(\epsilon^\xi), f_n(\epsilon^\xi) \rangle \\ & = 2\epsilon^{1+\xi} \operatorname{Re} \langle V_q f_n(0), f_n(\epsilon^\xi) \rangle. \end{aligned}$$

Now we can calculate explicitly

$$\langle V_q f_n(\epsilon^\xi), f_n(\epsilon^\xi) \rangle = \frac{1}{6} (-e^{-i\epsilon^\xi} + 2 - e^{i\epsilon^\xi}) = -\frac{1}{3} \epsilon^\xi + O(\epsilon^{2\xi}).$$

This shows that, for small ϵ ,

$$\langle H_{\epsilon, q} u_n, u_n \rangle \leq C \xi^{6\xi} - \frac{1}{3} \epsilon^{1+2\xi} + O(\epsilon^{1+2\xi}) \leq -\frac{1}{6} \epsilon^{1+2\xi},$$

where we have used that $6\xi > 1 + 2\xi$.

6.2. Proof of Lemma 4.1. As the V_q is block-diagonal, it is enough to do the calculation for the free operator H_0 . Let us first calculate some norms. Because of the quasi-periodicity, we easily see that

$$(41) \quad \|u_n\|_{\ell^2(D)}^2 = (2n+1)^d \|u_0\|_{\ell^2(\square)}^2.$$

and

$$\|u_n - u_{n-1}\|_{\ell^2(D)}^2 \leq C n^{d-1} \|u_0\|_{\ell^2(\square)}^2.$$

So we have that

$$(42) \quad \langle H_0 u_n, u_n \rangle = \langle H_0 u_n, u_{n-1} \rangle + O(n^{d-1}) \|u_0\|_{\ell^2(\square)}^2.$$

For any $k \in \square_{n-1}$ and $k' \in \mathbb{Z}^d \setminus \square_n$, we have that $|k - k'| \geq N$ and thus, because of the finite hopping range (assumption (HB)),

$$\langle H_0 u_n, u_{n-1} \rangle = \langle H_0 u, u_{n-1} \rangle.$$

Now, we develop

$$\begin{aligned}
& \langle H_0 u, u_{n-1} \rangle \\
(43) \quad &= \sum_{k \in \mathbb{Z}^d} \sum_{k' \in \mathbb{Z}^d} H_0(k, k') u(k') \overline{u_{n-1}(k)} \\
&= \sum_{\substack{m \in \gamma \\ |m| \leq (n-1)N}} \sum_{m' \in \gamma} \sum_{k \in \square + m} \sum_{k' \in \square + m'} H_0(k, k') u(k') \overline{u_{n-1}(k)} \\
&= \sum_{\substack{m \in \gamma \\ |m| \leq (n-1)N}} \sum_{m' \in \gamma} \sum_{k \in \square} \sum_{k' \in \square} H_0(k + m, k' + m') u(k' + m') \overline{u_{n-1}(k + m)}
\end{aligned}$$

Using the translation invariance (assumption (HB)), the last quantity is equal to

$$\begin{aligned}
(44) \quad & \sum_{\substack{m \in \gamma \\ |m| \leq (n-1)N}} \sum_{m' \in \gamma} \sum_{k \in \square} \sum_{k' \in \square} H_0(k, k' + m' - m) u(k' + m' - m) \overline{u_{n-1}(k)} \\
&= \sum_{\substack{m \in \gamma \\ |m| \leq (n-1)N}} \sum_{m' \in \gamma} \sum_{k \in \square} \sum_{k' \in \square} e^{i\theta \cdot (m - m')} H_0(k, k' - m' + m) u_0(k') \overline{u_0(k)} \\
&= \sum_{\substack{m \in \gamma \\ |m| \leq (n-1)N}} \sum_{m'' \in \gamma} \sum_{k \in \square} \sum_{k' \in \square} e^{i\theta \cdot m''} H_0(k, k' - m'') u_0(k') \overline{u_0(k)} \\
&= (2n - 3)^d \langle H_0(\theta) u_0, u_0 \rangle.
\end{aligned}$$

We see from this calculation and (42) thus that

$$\left| \langle H_0 u_n, u_n \rangle - (2n - 3)^d \langle H_0(\theta) u_0, u_0 \rangle \right| \leq C n^{d-1} \|u_0\|_{\ell^2(\square)}^2.$$

As $(2n - 3)/(2n - 1) \rightarrow 1$, dividing by $\|u_n\|_{\ell^2(D)}^2$, using (41) and taking the limit proves the lemma. \square

ACKNOWLEDGMENT

The research of D.B. was supported by Russian Scientific Foundation (project no. 14-11-00078). F.H-E and I.V. were supported by DFG-Projects *Zufällige und periodische Quantengraphen* and *Eindeutige-Fortsetzungsprinzipien und Gleichverteilungseigenschaften von Eigenfunktionen*.

We thank M. Täufer for reading a previous version of the manuscript.

REFERENCES

- [1] M. Aizenman. Localization at weak disorder: some elementary bounds. *Rev. Math. Phys.*, 06:1163, 1994.
- [2] D. Borisov and I. Veselić. Low lying spectrum of weak-disorder quantum waveguides. *J. Stat. Phys.*, 142(1):58–77, 2011.
- [3] D. Borisov and I. Veselić. Low lying eigenvalues of randomly curved quantum waveguides. *J. Funct. Anal.*, 265(11):2877–2909, 2013.
- [4] Z. Cao and A. Elgart. The weak localization for the alloy-type Anderson model on a cubic lattice. *J. Stat. Phys.*, 148(6):1006–1039, 2012.
- [5] R. Carmona and J. Lacroix. *Spectral theory of random Schrödinger operators*. Probability and its Applications. Birkhäuser Boston, 1990.
- [6] Y. Colin de Verdière. Sur les singularités de Van Hove génériques. (On generic Van Hove singularities). *Mém. Soc. Math. Fr., Nouv. Sér.*, 46:99–109, 1991.
- [7] A. Elgart. Lifshitz tails and localization in the three-dimensional Anderson model. *Duke Math. J.*, 146(2):331–360, 2009.
- [8] A. Elgart, H. Krüger, M. Tautenhahn, and I. Veselić. Discrete Schrödinger operators with random alloy-type potential. In R. Benguria, E. Friedman, and M. Mantoiu, editors, *Spectral Analysis of Quantum Hamiltonians*, volume 224 of *Operator Theory: Advances and Applications*, pages 107–131. Springer Basel, 2012.
- [9] A. Elgart, M. Tautenhahn, and I. Veselić. Anderson localization for a class of models with a sign-indefinite single-site potential via fractional moment method. *Ann. Henri Poincaré*, 12(8):1571–1599, 2011.
- [10] P. D. Hislop. Lectures on Random Schrödinger Operators. *Fourth Summer School in Analysis and Mathematical Physics: topics in spectral theory and quantum mechanics, May 2005, Universidad Nacional Autónoma de México, Cuernavaca, México*, 476:41, 2008.
- [11] F. Hoecker-Escuti. *Le modèle d'Anderson en régime de faible désordre*. PhD thesis, Université Paris 13, June 2013.
- [12] W. Kirsch and F. Martinelli. On the spectrum of Schrödinger operators with a random potential. *Commun. Math. Phys.*, 85:329–350, 1982.
- [13] W. Kirsch and B. Simon. Comparison theorems for the gap of Schrödinger operators. *J. Funct. Anal.*, 75:396–410, 1987.
- [14] W. Kirsch, P. Stollmann, and G. Stolz. Localization for random perturbations of periodic Schrödinger operators. *Random Oper. Stoch. Equ.*, 6(3):241–268, 1998.
- [15] F. Klopp. Weak disorder localization and Lifshitz tails. *Commun. Math. Phys.*, 232(1):125–155, 2002.
- [16] F. Klopp. Weak disorder localization and Lifshitz tails: continuous Hamiltonians. *Annales Henri Poincaré*, 3(4):711–737, 2002.

- [17] H. Krueger. Periodic and limit-periodic discrete Schrödinger operators. *arXiv:1108.1584*, 2011.
- [18] C. D. Meyer. *Matrix analysis and applied linear algebra*. Society for Industrial and Applied Mathematics, 2000.
- [19] L. Pastur and A. Figotin. *Spectra of random and almost-periodic operators*, volume 297 of *Grundlehren der mathematischen Wissenschaften*. Springer, 1992.
- [20] M. Reed and B. Simon. *Analysis of Operators*, volume 1 : Functional Analysis of *Methods of Modern Mathematical Physics*. Academic Press, Inc., 1981.
- [21] P. Stollmann. *Caught by disorder. Bound states in random media.*, volume 20 of *Progress in Mathematical Physics*. Birkhäuser, 2001.
- [22] E. C. Titchmarsh. *Eigenfunction expansions associated with second-order differential equations*. Oxford University Press, 1946.
- [23] W.-M. Wang. Localization and universality of Poisson statistics for the multi-dimensional Anderson model at weak disorder. *Invent. Math.*, 146(2):365–398, 2001.
- [24] T. H. Wolff. Recent work on sharp estimates in second-order elliptic unique continuation problems. *J. Geom. Anal.*, 3(6):621–650, 1993.

(D. Borisov) INSTITUTE OF MATHEMATICS OF UFA SCIENTIFIC CENTER OF RAS, CHERNYSHEVSKOGO, 112, UFA, 450008, RUSSIA

&

DEPARTMENT OF PHYSICS AND MATHEMATICS, BASHKIR STATE PEDAGOGICAL UNIVERSITY, OCTOBER REV. ST. 3A, UFA, 450000, RUSSIA

URL: <http://borisovdi.narod.ru/>

E-mail address: borisovdi@yandex.ru

(F. Hoecker-Escuti) TECHNISCHE UNIVERSITÄT CHEMNITZ, FAKULTÄT FÜR MATHEMATIK, REICHENHAINER STR. 41, CHEMNITZ, D-09126, GERMANY

E-mail address: francisco.hoecker-escuti@mathematik.tu-chemnitz.de

(I. Veselić) TECHNISCHE UNIVERSITÄT CHEMNITZ, FAKULTÄT FÜR MATHEMATIK, REICHENHAINER STR. 41, CHEMNITZ, D-09126, GERMANY

URL: <http://www.tu-chemnitz.de/mathematik/stochastik/>