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# Uniform Hausdorff measure of the level sets of the Brownian tree. 

Xan Duhalde*

July 21, 2014


#### Abstract

Let $(\mathcal{T}, d)$ be the random real tree with root $\rho$ coded by a Brownian excursion. So $(\mathcal{T}, d)$ is (up to normalisation) Aldous CRT [1] (see Le Gall [10]). The $a$-level set of $\mathcal{T}$ is the set $\mathcal{T}(a)$ of all points in $\mathcal{T}$ that are at distance $a$ from the root. We know from Duquesne and Le Gall [7] that for any fixed $a \in(0, \infty)$, the measure $\ell^{a}$ that is induced on $\mathcal{T}(a)$ by the local time at $a$ of the Brownian excursion, is equal, up to a multiplicative constant, to the Hausdorff measure in $\mathcal{T}$ with gauge function $g(r)=r \log \log 1 / r$, restricted to $\mathcal{T}(a)$. As suggested by a result due to Perkins $[14,15]$ for super-Brownian motion, we prove in this paper a more precise statement that holds almost surely uniformly in $a$, and we specify the multiplicative constant. Namely, we prove that almost surely for any $a \in(0, \infty), \ell^{a}(\cdot)=\frac{1}{2} \mathscr{H}_{g}(\cdot \cap \mathcal{T}(a))$, where $\mathscr{H}_{g}$ stands for the $g$-Hausdorff measure.


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## 1 Introduction.

The Continuum Random Tree was introduced by Aldous [1] as a random compact metric space $\left(\mathcal{T}_{1}, d, \mathbf{m}_{1}\right)$, endowed with a mass measure $\mathbf{m}_{1}$ such that almost surely $\mathbf{m}_{1}\left(\mathcal{T}_{1}\right)=1$. It appears as the scaling limit of a large class of discrete models of random trees, and can be alternatively encoded by a normalised Brownian excursion (see Le Gall [10]). This encoding procedure will be the viewpoint of the present paper, but for the sake of simplicity, we will not ask the total mass to be equal to one. Instead, we work on the tree encoded by a Brownian excursion ( $e_{t}, t \geqslant 0$ ), under its excursion measure $\mathbf{N}$. Let us mention that our result remains true for the CRT.

The Brownian tree has a distinguished vertex $\rho$ called the root, so it makes sense to define, for all $a \in(0, \infty)$ the $a$-level set $\mathcal{T}(a)=\{\sigma \in \mathcal{T}: d(\rho, \sigma)=a\}$. Moreover, one can define the collection of measures ( $\ell(\mathrm{d} \sigma), \sigma \in \mathcal{T}, a \in(0, \infty)$ ), as the image of the local times on the levels of the excursion. Those measures are called local time measures. Indeed, $\mathbf{N}$-a.e. for all $a \in(0, \infty)$, the topological support of $\ell^{a}$ is included in $\mathcal{T}(a)$. Duquesne and Le Gall [7] showed that for a fixed level $a$, one has

$$
\begin{equation*}
\mathbf{N} \text {-a.e. } \quad \ell^{a}(\cdot)=c \mathscr{H}_{g}(\cdot \cap \mathcal{T}(a)), \tag{1}
\end{equation*}
$$

[^0]where $\mathscr{H}_{g}$ stands for the Hausdorff measure associated with the gauge function $g(r)=r \log \log 1 / r$ and $c \in(0, \infty)$ is a multiplicative constant. In this paper, we prove that $c=\frac{1}{2}$ and that the result holds $\mathbf{N}$-a.e. simultaneously for all levels $a$. Let us mention that the value $\frac{1}{2}$ depends on the normalisation chosen for the excursion measure $\mathbf{N}$. A similar result has been obtained by Perkins [14, 15] for Super Brownian Motion. Briefly, let ( $\left.Z_{a}, a \geqslant 0\right)$ a version of this measure-valued process on $\mathbb{R}^{d}$, defined on $(\Omega, \mathcal{F}, \mathbf{P})$. Perkins proves that if the dimension $d$ of the space is such that $d \geqslant 3$ (which corresponds to the supercritical dimension case), there exists two constants $c_{d}, C_{d}$ in $(0, \infty)$, only depending on $d$ such that the following holds
\[

$$
\begin{equation*}
\text { P-a.s. } \quad \forall a \in(0, \infty) \quad c_{d} \mathscr{H}_{g}\left(\cdot \cap \operatorname{Supp}\left(Z_{a}\right)\right) \leqslant Z_{a}(\cdot) \leqslant C_{d} \mathscr{H}_{g}\left(\cdot \cap \operatorname{Supp}\left(Z_{a}\right)\right), \tag{2}
\end{equation*}
$$

\]

where $\operatorname{Supp}\left(Z_{a}\right)$ is the topological support of the measure $Z_{a}$ and $\mathscr{H}_{g}$ is the Hausdorff measure associated to the gauge function $g(r)=r^{2} \log \log 1 / r$. In this paper, we use the ideas and techniques of $[14,15]$ to get a result similar to (2), an equality being accessible in the setting of trees.

Before stating formally our result, let us recall precisely basic facts. A metric space $(T, d)$ is a real tree if and only if the following two properties hold for any $\sigma_{1}, \sigma_{2}$ in $T$ :
(i) There is a unique isometric map $f_{\sigma_{1}, \sigma_{2}}$ from $\left[0, d\left(\sigma_{1}, \sigma_{2}\right)\right]$ into $T$ such that $f_{\sigma_{1}, \sigma_{2}}(0)=\sigma_{1}$ and $f_{\sigma_{1}, \sigma_{2}}\left(d\left(\sigma_{1}, \sigma_{2}\right)\right)=\sigma_{2}$. We set $\llbracket \sigma_{1}, \sigma_{2} \rrbracket=f_{\sigma_{1}, \sigma_{2}}\left(\left[0, d\left(\sigma_{1}, \sigma_{2}\right)\right]\right)$ that is the geodesic path joining $\sigma_{1}$ and $\sigma_{2}$.
(ii) If $q$ is a continuous injective map from $[0,1]$ into $T$, such that $q(0)=\sigma_{1}$ and $q(1)=\sigma_{2}$, we have

$$
q([0,1])=f_{\sigma_{1}, \sigma_{2}}\left(\left[0, d\left(\sigma_{1}, \sigma_{2}\right)\right]\right)
$$

If $\sigma_{1} \in \llbracket \rho, \sigma_{2} \rrbracket$, we will say that $\sigma_{1}$ is an ancestor of $\sigma_{2}\left(\sigma_{2}\right.$ is a descendant of $\left.\sigma_{1}\right)$.
Real trees can be derived from continuous functions that represent their contour functions. Namely, let us consider a (deterministic) excursion $e$, that is to say a continuous function for which there exists $\zeta \in(0, \infty)$ such that : $\forall t \geqslant \zeta, e(0)=e(t)=0$, and $\forall t \in(0, \zeta), e(t)>0$. A real tree $T$ can be associated with $e$ in the following way. For $s, t \in[0, \zeta]$, we set

$$
d(s, t)=e(s)+e(t)-2 \inf _{r \in[s \wedge t, s \vee t]} e(r)
$$

It is easy to see that $d$ is a pseudo-distance on $[0, \zeta]$. Defining the equivalence relation $s \sim t$ iff $d(s, t)=0$, one can set

$$
\begin{equation*}
T=[0, \zeta] / \sim \tag{3}
\end{equation*}
$$

The function $d$ induces a distance on the quotient set $T$. For a fixed excursion $e$, let

$$
\begin{equation*}
p:[0, \zeta] \longrightarrow(T, d) \tag{4}
\end{equation*}
$$

be the canonical projection. Clearly $p$ is continuous, which implies that $(T, d)$ is a compact metric space. Moreover, it can be shown (see [6] for a proof) that $(T, d)$ is a real tree
We take $\rho=p(0)$ as the root of $T$. For all $a \in(0, \infty)$, the $a$-level set $T(a)=\{\sigma \in T: d(\rho, \sigma)=a\}$ is the image by $p$ of the set $\{t \in[0, \zeta]: e(t)=a\}$. The total height of the tree is defined by

$$
\begin{equation*}
h(T)=\sup \{d(\rho, \sigma) ; \sigma \in T\} \tag{5}
\end{equation*}
$$

We define the Brownian tree as the metric space $(\mathcal{T}, d)$ coded by the Brownian excursion. More precisely, let $(\Omega, \mathcal{F}, \mathbf{P})$ a probability space, large enough to carry all the random variables we need. We consider on that space a process $\left(X_{t}, t \in[0, \infty)\right)$ such that $\left(\frac{1}{\sqrt{2}} X_{t}, t \in[0, \infty)\right)$ is a standard real-valued Brownian motion (the choice of the normalizing constant $\sqrt{2}$ is explained below). Let us set $\underline{X}_{t}=\inf _{s \in[0, t]} X_{s}$. Then, the reflected process $X-\underline{X}$ is a strong Markov process, and the state 0 is instantaneous in $(0, \infty)$ and recurrent (see [2], chapter VI). We denote by $\mathbf{N}$ the excursion measure associated with the local time $-\underline{X} ; \mathbf{N}$ is a sigma-finite measure on the space of continuous functions on $[0, \infty)$, denoted $\mathbf{C}^{0}$ in this work. More precisely, let $\bigcup_{j \in \mathcal{J}}\left(l_{j}, r_{j}\right)=\left\{t>0: X_{t}-\underline{X}_{t}>0\right\}$ be the excursion intervals of the reflected process, and for all $j \in \mathcal{J}$, we set $e_{j}(s)=X_{\left(l_{j}+s\right) \wedge d_{j}}-\underline{X}_{l_{j}}, s \in[0, \infty)$. Then,

$$
\mathcal{M}(\mathrm{d} t, \mathrm{~d} e)=\sum_{j \in \mathcal{J}} \delta_{\left(-\underline{X}_{l_{j}}, e_{j}\right)}
$$

is a Poisson point measure on $[0, \infty) \times \mathbf{C}^{0}$ of intensity $\mathrm{d} t \mathbf{N}(\mathrm{~d} e)$. Let us recall that the two processes $\left(\left|X_{t}\right|, 2 L_{t}\right)_{t \geqslant 0}$ and $\left(X_{t}-\underline{X}_{t},-\underline{X}_{t}\right)_{t \geq 0}$ have the same law under $\mathbf{P}$ by a celebrated result of Lévy (see Blumenthal [4], Th. II 2.2) where the process $\left(L_{t}, t \geqslant 0\right)$ is defined by the approximation $L_{t}=\lim _{\varepsilon \rightarrow 0}(2 \varepsilon)^{-1} \int_{0}^{t} \mathbf{1}_{\left\{\left|X_{s}\right| \leqslant \varepsilon\right\}} \mathrm{d} s$ that holds uniformly in $t$ on compact subsets of $[0, \infty)$.

We shall denote by $\left(e_{t}, t \geqslant 0\right)$ the canonical process on $\mathbf{C}^{0}$. Under $\mathbf{N}$, it is a strong Markov process, with transition kernel of the original process $X$ killed when it hits 0 (see [4] III 3(f)). The following properties hold for the process $\mathbf{N}$-a.e. : there exists a unique real $\zeta \in(0, \infty)$ such that $\forall t \in(0, \zeta), e(t)>0$, and $\forall t \in[\zeta, \infty), e(t)=e(0)=0$. Moreover, with our normalization, one has (see [4] IV 1.1)

$$
\begin{equation*}
\forall \lambda \in[0, \infty), \mathbf{N}\left(1-e^{-\lambda \zeta}\right)=\sqrt{\lambda} \quad \text { and } \quad \mathbf{N}(\zeta \in \mathrm{d} r)=\frac{r^{-3 / 2}}{2 \sqrt{\pi}} \mathrm{~d} r \tag{6}
\end{equation*}
$$

One can show that $\mathbf{N}(\cdot \mid \zeta \in[1-\varepsilon, 1+\varepsilon])$ converges when $\varepsilon$ goes to 0 , towards a probability measure that is denoted by $\mathbf{N}(\cdot \mid \zeta=1)$. It can be seen as the law of the excursion of $X-\underline{X}$ conditioned to have length one. The tree encoded by $e$ under $\mathbf{N}(\cdot \mid \zeta=1)$ is the CRT defined in [1]. The choice of the normalising constant $\sqrt{2}$ is explained by the following. Let $\tau_{n}$ be uniformly distributed as the set of rooted planar trees with $n$ vertices. We view it as a real tree, the edges of $\tau_{n}$ being intervals of length one, and we denote by $\left(\tau_{n}, d_{n}\right)$ the resulting metric space. Denote by $\left(C_{t}^{(n)}, t \in[0,2(n-1)]\right)$ its contour function that is (informally) defined as follows. We let a particle explore the planar tree at speed one, from the left to the right, beginning at the root. We set $C_{t}^{(n)}$ as the distance from the root of the particle at time $t$. It can be shown (see [12] Th. 1.17) that $\left(C_{t}^{(n)}, t \in[0,2(n-1)]\right)$ has the law of a simple random walk conditioned to be positive on $[1,2(n-1)-1]$ and null at $2(n-1)$. Using Donsker invariance principle, the rescaled contour function $\left(n^{-1 / 2} C_{2(n-1) t}^{(n)}, t \in[0,1]\right)$ converges in law towards the law of $\left(e_{t}, t \in[0,1]\right)$ under $\mathbf{N}(\cdot \mid \zeta=1)$. In terms of trees, $\left(\tau_{n}, n^{-1 / 2} d_{n}\right)$ converges towards the CRT, that is the tree $\left(\mathcal{T}_{1}, d\right)$ coded by $e$ under $\mathbf{N}(\cdot \mid \zeta=1)$. The latter convergence can be stated using the distance of Gromov-Hausdorff (see Evans, Pitman, Winter [9]).

Recalling definition (5), we get from [4] IV 1.1 that with our normalization,

$$
\begin{equation*}
\forall a \in(0, \infty) \quad \mathbf{N}\left(\sup _{t \in[0, \zeta]} e_{t}>a\right)=\mathbf{N}(h(\mathcal{T})>a)=\frac{1}{a} \tag{7}
\end{equation*}
$$

In the paper, for $a \in(0, \infty)$ we shall use the probability measure,

$$
\begin{equation*}
\mathbf{N}_{a}=\mathbf{N}(\cdot \mid h(\mathcal{T})>a)=a \mathbf{N}\left(\cdot \mathbf{1}_{\{h(\mathcal{T})>a\}}\right) \tag{8}
\end{equation*}
$$

Recall that the $a$-level set of the Brownian tree is defined by

$$
\begin{equation*}
\mathcal{T}(a)=\{\sigma \in \mathcal{T}: d(\rho, \sigma)=a\} \tag{9}
\end{equation*}
$$

As a consequence of Trotter's theorem on the regularity of Brownian local time ([4] sec VI.3) there exists a $[0, \infty)$-valued process $\left(L_{t}^{a}\right)_{a, t \in[0, \infty)}$ such that $\mathbf{N}$-a.e. the following holds true:

- $(a, t) \mapsto L_{t}^{a}$ is continuous,
- for all $a \in[0, \infty), t \mapsto L_{t}^{a}$ is non-decreasing,
- for all $a \in[0, \infty)$, for all $t \in[0, \infty)$ and for all $b \in(0, \infty)$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \mathbf{N}\left(1_{\{\sup e>b\}} \sup _{0 \leqslant s \leqslant t \wedge \zeta}\left|\frac{1}{\varepsilon} \int_{0}^{s} 1_{\{a-\varepsilon<e(u) \leqslant a\}} \mathrm{d} u-L_{s}^{a}\right|\right)=0 \tag{10}
\end{equation*}
$$

We refer to [5], Proposition 1.3.3. for details in a more general setting.

The image by the projection $p:[0, \zeta] \rightarrow \mathcal{T}$ of those local times defines the collection of local time measures on the tree, $\left(\ell^{a}(\mathrm{~d} \sigma), \sigma \in \mathcal{T}, a \in(0, \infty)\right)$. More precisely,

$$
\begin{equation*}
\mathbf{N} \text {-a.e. } \quad \text { for all } f: \mathcal{T} \xrightarrow{\text { meas. }}[0, \infty) \quad \forall a \in(0, \infty) \quad \int_{\mathcal{T}} f(\sigma) \ell^{a}(\mathrm{~d} \sigma)=\int_{0}^{\zeta} f(p(t)) \mathrm{d} L_{t}^{a} \tag{11}
\end{equation*}
$$

See [6], Th. 4.2 for an intrinsic definition of the measure $\ell^{a}$ (for fixed $a$ ). Let $\mathcal{G}_{a}$ the $\sigma$-field generated by the excursion below level $a$ (formal definitions and details on what follows are given in section 3.1). The approximation (10) entails that for fixed $a, \ell^{a}(\mathcal{T})=L_{\zeta}^{a}$ is $\mathcal{G}_{a}$ measurable. Moreover, the Ray-Knight theorem ([4] VI 2.10) entails that under $\mathbf{N}_{a}(\cdot)$ conditionally on $\mathcal{G}_{a}$, the process $\left(\ell^{a+a^{\prime}}(\mathcal{T}), a^{\prime} \geqslant 0\right)$ is a Feller diffusion started at $\ell^{a}(\mathcal{T})$. In particular, one has

$$
\begin{equation*}
\forall a, \lambda \in(0, \infty) \quad \mathbf{N}\left[1-e^{-\lambda \ell^{a}(\mathcal{T})}\right]=\frac{\lambda}{1+a \lambda} \tag{12}
\end{equation*}
$$

which implies that under $\mathbf{N}_{a}, \ell^{a}(\mathcal{T})$ is exponentially distributed with mean $a$. The regularity of $a \mapsto \ell^{a}(\mathcal{T})$ is extended by Duquesne and Le Gall [6] : they prove that $\mathbf{N}$-a.e. the process $a \mapsto \ell^{a}$ is continuous for the weak topology of measures. In the same work, the topological support of the level set measures is precised as follows. A vertex $\sigma \in \mathcal{T}$ is called an extinction point if there exists $\varepsilon \in(0, \infty)$ such that $d(\rho, \sigma)=\sup \{d(\rho, \tau), \tau \in B(\sigma, \varepsilon)\}$, where $B(\sigma, \varepsilon)$ is the open ball in $\mathcal{T}$ with center $\sigma$ and radius $\varepsilon$. For $s \in[0, \zeta]$, the vertex $p(s) \in \mathcal{T}$ is an extinction time iff $s \in[0, \zeta]$ is a local maximum of $e$. As a consequence, the set of all extinction points, denoted $\mathscr{E}$, is countable. Let us denote $\operatorname{Supp}(\mu)$ for the topological support of the measure $\mu$. The result states that

$$
\begin{equation*}
\mathbf{N} \text {-a.e. } \quad \forall a \in(0, \infty) \backslash \mathscr{E}, \operatorname{Supp}\left(\ell^{a}\right)=\mathcal{T}(a), \quad \text { and } \quad \forall a \in \mathscr{E}, \operatorname{Supp}\left(\ell^{a}\right)=\mathcal{T}(a) \backslash\left\{\sigma_{a}\right\} \tag{13}
\end{equation*}
$$

where $\sigma_{a}$ is the (unique) extinction point at level $a$ (see Perkins [16] for previous results on SuperBrownian motion).

Let us briefly introduce the construction of the Hausdorff measure. We set the gauge function $g$ as

$$
\begin{equation*}
g(r)=r \log \log 1 / r, \quad r \in\left(0, e^{-1}\right) . \tag{14}
\end{equation*}
$$

In all the paper it will be assumed implicitly that $g(r)$ is considered only for $r \in\left(0, e^{-1}\right)$. On that interval, $g$ is an increasing continuous function. For any subset $A$ of $\mathcal{T}$, one can define

$$
\begin{equation*}
\mathscr{H}_{g}(A)=\lim _{\varepsilon \rightarrow 0} \inf \left\{\sum_{i \in \mathbb{N}} g\left(\operatorname{diam}\left(E_{i}\right)\right) ; A \subset \bigcup_{i \in \mathbb{N}} E_{i}, \operatorname{diam}\left(E_{i}\right)<\varepsilon\right\} . \tag{15}
\end{equation*}
$$

Standard results on Hausdorff measures (see e.g. [17]) ensure that $\mathscr{H}_{g}$ defines a Borel-regular outer measure on $\mathcal{T}$ called the $g$-Hausdorff measure on $\mathcal{T}$. The main result of the paper is the following.

Theorem 1 Let $\mathcal{T}$ be the Brownian tree, that is the tree encoded by the excursion e under $\mathbf{N}$. Let ( $\left.\ell^{a}(\mathrm{~d} \sigma), \sigma \in \mathcal{T}, a \in(0, \infty)\right)$ the collection of local time measures and $\mathscr{H}_{g}$ the $g$-Hausdorff measure on $\mathcal{T}$, where $g(r)=r \log \log 1 / r$. Then, the following holds:

$$
\begin{equation*}
\mathbf{N} \text {-a.e. } \quad \forall a \in(0, \infty) \quad \ell^{a}(\cdot)=\frac{1}{2} \mathscr{H}_{g}(\cdot \cap \mathcal{T}(a)) . \tag{16}
\end{equation*}
$$

Comment 1.1 Thanks to the scaling properties of the Brownian excursion, one can derive from Theorem 1 a similar statement for the tree coded by $e$ under $\mathbf{N}(\cdot \mid \zeta=1)$, that is Aldous CRT.

Comment 1.2 Our result seems close to a theorem of Perkins [13] on linear Brownian motion. Let $\left(L_{t}^{a}, t \geqslant 0, a \in \mathbb{R}\right)$ be the bi-continuous version of the local times for the process ( $X_{t}, t \geqslant 0$ ) defined above. Those local times are given by an approximation of the type of (10). Perkins proves that almost surely, uniformly in $a$, one has $L_{t}^{a}=\mathscr{H}_{g}\left(\left\{s \in[0, t]: X_{s}=x\right\}\right)$, where $\mathscr{H}_{g}$ stands for the Hausdorff measure on the line associated with the gauge $g(r)=\sqrt{r \log \log 1 / r}$ (the result for fixed $a$ had been obtain by Taylor and Wendel in [19]). The Brownian tree being coded by the Brownian excursion, everything happens as if the projection mapping $p:[0, \zeta] \rightarrow \mathcal{T}$ is $1 / 2$-HÃPlder and induces a strong "doubling", such that the entire gauge function is squared. Nevertheless, we don't see how to derive our result from [13].

The paper is organised as follows. In section 2, we state some deterministic facts on the geometry of the level sets for a real tree. In particular, we provide two comparison lemmas with respect to Hausdorff measure on real trees. The second one, that is specific to our setting, seems new to us. In section 3, we recall basic facts on the Brownian tree and we establish some technical estimates. Section 4 is devoted to the proof of Theorem 1. As a first step, we prove Theorem 2 , which gives an upper bound for the local time measures. To that end, we need to control the total mass of the balls that are "too large". The second step is the proof of Theorem 3, which requires a control of the number of balls that are "too small". Let us mention again that our strategy and many ideas in this work were borrowed from [14, 15].

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## 2 Geometric properties of the level sets of real trees.

### 2.1 The balls of the level sets of real trees.

Let ( $T, d, \rho$ ) be a compact rooted real tree as defined in the introduction. Recall that for any $\sigma, \sigma^{\prime} \in T, \llbracket \sigma, \sigma^{\prime} \rrbracket$ stands for the unique geodesic path joining $\sigma$ to $\sigma^{\prime}$. We shall view $T$ as a family tree whose ancestor is the root $\rho$ and we then denote by $\sigma \wedge \sigma^{\prime}$ the most recent common ancestor of $\sigma$ and $\sigma^{\prime}$ that is formally defined by

$$
\llbracket \rho, \sigma \wedge \sigma^{\prime} \rrbracket=\llbracket \rho, \sigma \rrbracket \cap \llbracket \rho, \sigma^{\prime} \rrbracket .
$$

Observe that

$$
\begin{equation*}
\forall \sigma, \sigma^{\prime} \in T, \quad d\left(\sigma, \sigma^{\prime}\right)=d(\rho, \sigma)+d\left(\rho, \sigma^{\prime}\right)-2 d\left(\rho, \sigma \wedge \sigma^{\prime}\right) . \tag{17}
\end{equation*}
$$

Let $a \in[0, \infty)$. Recall that the $a$-level set of $T$ is given by

$$
T(a)=\{\sigma \in T: d(\rho, \sigma)=a\} .
$$

Subtrees above level $b$. Let $b \in[0, \infty)$ and denote by $\left(T_{j}^{o, b}\right)_{j \in \mathcal{J}_{b}}$ the connected components of the open set $\{\sigma \in T: d(\rho, \sigma)>b\}$ :

$$
\bigcup_{j \in \mathcal{J}_{b}} T_{j}^{o, b}=\{\sigma \in T: d(\rho, \sigma)>b\} .
$$

Then for any $j \in \mathcal{J}_{b}$, there exists a unique point $\sigma_{j} \in T(b)$ such that $T_{j}^{b}:=T_{j}^{o, b} \cup\left\{\sigma_{j}\right\}$ is the closure of $T_{j}^{o, b}$ in $T$. Note that $\left(T_{j}^{b}, d, \sigma_{j}\right)$ is a compact rooted real tree and that

$$
\forall j \in \mathcal{J}_{b}, \forall \sigma \in T_{j}^{b}, \quad \sigma_{j} \in \llbracket \rho, \sigma \rrbracket .
$$

Open balls in $T(a)$. Recall that $B(\sigma, r)$ stands for the open ball in $T$ with center $\sigma$ and radius $r$. We shall also denote by $\Gamma(\sigma, r)$ the open ball with center $\sigma$ and radius $r$ in the level set of $\sigma$, namely

$$
\begin{equation*}
\Gamma(\sigma, r)=B(\sigma, r) \cap T(a), \quad \text { where } a=d(\rho, \sigma) . \tag{18}
\end{equation*}
$$

If $\sigma \in T(a)$, then we call $\Gamma(\sigma, r)$ a $T(a)$-ball with radius $r$; we denote by $\mathscr{B}_{a, r}$ the set of all the $T(a)$-balls with radius $r$ :

$$
\begin{equation*}
\mathscr{B}_{a, r}=\{\Gamma(\sigma, r) ; \sigma \in T(a)\} . \tag{19}
\end{equation*}
$$

The following proposition provides the geometric properties of $T(a)$-balls that we shall use.
Proposition 1 Let $(T, d, \rho)$ be a compact rooted real tree. Let $a, r \in(0, \infty)$ be such that $a \geqslant r / 2$. Then, the number of $T(a)$-balls with radius $r$ is finite. We set

$$
\begin{equation*}
Z_{a, r}=\# \mathscr{B}_{a, r} \quad \text { and } \quad\left\{\Gamma_{i}, 1 \leqslant i \leqslant Z_{a, r}\right\}=\mathscr{B}_{a, r} . \tag{20}
\end{equation*}
$$

Then, the following holds true.
(i) Set $b=a-\frac{1}{2} r$. Then, there are $Z_{a, r}$ distinct subtrees above $b$ denoted by $\left(T_{j_{i}}^{b}, d, \sigma_{j_{i}}\right), j_{i} \in \mathcal{J}_{b}$, $1 \leqslant i \leqslant Z_{a, r}$ such that

$$
\Gamma_{i}=T(a) \cap T_{j_{i}}^{b}=\left\{\sigma^{\prime} \in T_{j_{i}}^{b}: d\left(\sigma_{j_{i}}, \sigma^{\prime}\right)=r / 2\right\} .
$$

Thus, the $T(a)$-balls with radius $r$ are pairwise disjoint.
(ii) For all $\sigma \in T(a)$, one has $\operatorname{diam}(\Gamma(\sigma, r)) \leqslant r$. If furthermore $r \in(0,2 a)$, then $\operatorname{diam}(\Gamma(\sigma, r))<r$ and

$$
\begin{equation*}
\forall r^{\prime} \in(\operatorname{diam}(\Gamma(\sigma, r)), r) \quad \Gamma\left(\sigma, r^{\prime}\right)=\Gamma(\sigma, r) \tag{21}
\end{equation*}
$$

Therefore, the set of all $T(a)$-balls is countable.
(iii) Two $T(a)$-balls are either contained one in the other or disjoint. Namely, for all $r^{\prime}<r$ and all $\sigma, \sigma^{\prime} \in T(a)$, either $\Gamma\left(\sigma^{\prime}, r^{\prime}\right) \subset \Gamma(\sigma, r)$ or $\Gamma\left(\sigma^{\prime}, r^{\prime}\right) \cap \Gamma(\sigma, r)=\emptyset$.

Proof. Let us prove $(i)$. Let $\sigma, \sigma^{\prime} \in T(a)$ and set $b=a-\frac{1}{2} r$. By (17), $d\left(\sigma, \sigma^{\prime}\right)=2 a-2 d\left(\rho, \sigma \wedge \sigma^{\prime}\right)$. Thus, $d\left(\sigma, \sigma^{\prime}\right)<r$ iff $d\left(\rho, \sigma \wedge \sigma^{\prime}\right)>b$. Let $j \in \mathcal{J}_{b}$ be such that $\sigma \in T_{j}^{b}$; namely, $T_{j}^{b}$ is the unique subtree above $b$ containing $\sigma$ and $\sigma_{j}$ is the unique point $\gamma \in \llbracket \rho, \sigma \rrbracket$ such that $d(\rho, \gamma)=b$. Now observe that for all $\sigma^{\prime} \in T(a)$,

$$
d\left(\rho, \sigma \wedge \sigma^{\prime}\right)>b \Longleftrightarrow \sigma \wedge \sigma^{\prime} \in \rrbracket \sigma_{j}, \sigma \rrbracket \Longleftrightarrow \sigma^{\prime} \in T_{j}^{b} .
$$

This proves that

$$
\begin{equation*}
\Gamma(\sigma, r)=T(a) \cap T_{j}^{b} \tag{22}
\end{equation*}
$$

Conversely, let $j \in \mathcal{J}_{b}$ be such that $h\left(T_{j}^{b}\right):=\max \left\{d\left(\sigma_{j}, \gamma\right) ; \gamma \in T_{j}^{b}\right\} \geqslant r / 2$. Let $\sigma \in T(a) \cap T_{j}^{b}$; then the previous arguments imply (22). Since $T$ is compact, the set $\left\{j \in \mathcal{J}_{b}: h\left(T_{j}^{b}\right) \geqslant r / 2\right\}$ is finite, which completes the proof of $(i)$.

Let us prove (ii): let $\sigma \in T(a)$, let $r \in(0,2 a)$ and set $\delta=\operatorname{diam}(\Gamma(\sigma, r))$. Then (22) implies that $\Gamma(\sigma, r)$ is compact and there are $\sigma_{1}, \sigma_{2} \in \Gamma(\sigma, r)$ such that $d\left(\sigma_{1}, \sigma_{2}\right)=\delta$. Observe that it implies

$$
\Gamma(\sigma, r)=\left\{\sigma^{\prime} \in T(a): \sigma_{1} \wedge \sigma_{2} \in \llbracket \rho, \sigma^{\prime} \rrbracket\right\}
$$

Thus, $\Gamma(\sigma, r)=\bar{\Gamma}(\sigma, \delta)$, that is the closure of $\Gamma(\sigma, \delta)$, and it implies (21). The set of all $T(a)$-balls is therefore $\bigcup_{q \in \mathbb{Q} \cap[0, \infty)} \mathscr{B}_{a, q}$, which is a countable set.

Let us prove (iii): $r^{\prime}<r$ and $\sigma, \sigma^{\prime} \in T(a)$ and suppose that $\Gamma\left(\sigma^{\prime}, r^{\prime}\right) \cap \Gamma(\sigma, r) \neq \emptyset$. Then (i) and (ii) implies that $\Gamma(\sigma, r)=\Gamma\left(\sigma^{\prime}, r\right)$, which implies that $\Gamma\left(\sigma^{\prime}, r^{\prime}\right) \subset \Gamma(\sigma, r)$.

### 2.2 Comparison lemmas for Hausdorff measures on real trees.

Let $(T, d, \rho)$ be a compact real tree. We briefly recall the definition of Hausdorff measures on $T$ and we state two comparison lemmas that are used in the proofs. Let $r_{0} \in(0, \infty)$ and let $g:\left[0, r_{0}\right) \rightarrow[0, \infty)$ be a function that is assumed to be increasing, continuous and such that $g(0)=0$. For all $\varepsilon \in\left(0, r_{0}\right)$ and all $A \subset T$, we set

$$
\mathscr{H}_{g}^{(\varepsilon)}(A)=\inf \left\{\sum_{n \in \mathbb{N}} g\left(\operatorname{diam}\left(E_{n}\right)\right) ; A \subset \bigcup_{n \in \mathbb{N}} E_{n}, \operatorname{diam}\left(E_{n}\right)<\varepsilon\right\}
$$

and

$$
\mathscr{H}_{g}(A)=\lim _{\varepsilon \downarrow 0} \uparrow \mathscr{H}_{g}^{(\varepsilon)}(A)
$$

Under our assumptions, $\mathscr{H}_{g}$ is a Borel-regular outer measure : this is the $g$-Hausdorff measure on $T$ (see Rogers [17]). The following comparison lemma was first stated for Euclidean spaces by Rogers and Taylor [18]. The proof can be easily adapted to general metric spaces (see Edgar [8]). We include a brief proof of it in order to make the paper self-contained.

Lemma 2 Let ( $T, d, \rho$ ) be a compact rooted real tree. Let $\mu$ be a Borel measure on $T$. Let $A$ be a Borel subset of $T$ and let $c \in(0, \infty)$. Assume that

$$
\forall \sigma \in A \quad \underset{r \rightarrow 0}{\limsup } \frac{\mu(B(\sigma, r))}{g(r)}<c .
$$

Then, $\mu(A) \leqslant c \mathscr{H}_{g}(A)$.
Proof. For any $\varepsilon \in\left(0, r_{0}\right)$, set

$$
A_{\varepsilon}=\left\{\sigma \in A: \sup _{r \in(0, \varepsilon)} \frac{\mu(B(\sigma, r))}{g(r)}<c\right\} .
$$

Observe that for all $\varepsilon^{\prime}<\varepsilon, A_{\varepsilon} \subset A_{\varepsilon^{\prime}} \subset A$ and $A=\bigcup_{\varepsilon \in\left(0, r_{0}\right)} A_{\varepsilon}$. Let $\left(E_{n}\right)_{n \in \mathbb{N}}$ be a $\varepsilon$-covering of $A_{\varepsilon}$ : namely $A_{\varepsilon} \subset \bigcup_{n \in \mathbb{N}} E_{n}$ and $\operatorname{diam}\left(E_{n}\right)<\varepsilon$, for all $n \in \mathbb{N}$. Set $I=\left\{n \in \mathbb{N}: E_{n} \cap A_{\varepsilon} \neq \emptyset\right\}$ and for all $n \in I$, fix $\sigma_{n} \in E_{n} \cap A_{\varepsilon}$. Since $g$ is continuous, for all $n \in I$ there exists $r_{n} \in\left(\operatorname{diam}\left(E_{n}\right), \varepsilon\right)$ such that

$$
E_{n} \subset B\left(\sigma_{n}, r_{n}\right) \quad \text { and } \quad g\left(r_{n}\right) \leqslant 2^{-n-1} \varepsilon+g\left(\operatorname{diam}\left(E_{n}\right)\right) .
$$

Observe that $\mu\left(B\left(\sigma_{n}, r_{n}\right)\right)<c g\left(r_{n}\right)$ and that $A_{\varepsilon} \subset \bigcup_{n \in I} B\left(\sigma_{n}, r_{n}\right)$. Thus,

$$
\begin{aligned}
\mu\left(A_{\varepsilon}\right) & \leqslant \mu\left(\bigcup_{n \in I} B\left(\sigma_{n}, r_{n}\right)\right) \leqslant \sum_{n \in I} \mu\left(B\left(\sigma_{n}, r_{n}\right)\right) \\
& \leqslant \sum_{n \in I} c g\left(r_{n}\right) \leqslant c \varepsilon+\sum_{n \in \mathbb{N}} c g\left(\operatorname{diam}\left(E_{n}\right)\right) .
\end{aligned}
$$

Taking the infimum over all the possible $\varepsilon$-coverings of $A_{\varepsilon}$ yields

$$
\mu\left(A_{\varepsilon}\right) \leqslant c \varepsilon+c \mathscr{H}_{g}^{(\varepsilon)}\left(A_{\varepsilon}\right) \leqslant c \varepsilon+c \mathscr{H}_{g}\left(A_{\varepsilon}\right) \leqslant c \varepsilon+\mathscr{H}_{g}(A),
$$

which implies the desired result since $\mu(A)=\lim _{\varepsilon \downarrow 0} \uparrow \mu\left(A_{\varepsilon}\right)$.
In the next comparison lemma, that seems new to us, we restrict our attention to the level sets of real trees. A more general variant of this result involves a multiplicative constant depending on the gauge function. It has been first stated in Euclidian spaces by Rogers and Taylor [18] (see also Perkins [14]) and in general metric spaces (see Edgar [8]).

Lemma 3 Let $(T, d, \rho)$ be a compact rooted real tree. Let $a \in(0, \infty)$ be such that the a-level set $T(a)$ is not empty. Let $\mu$ be a finite Borel measure on $T$ such that $\mu(T \backslash T(a))=0$. Let $A \subset T(a)$ be a Borel subset and let $c \in(0, \infty)$. Assume that

$$
\forall \sigma \in A \quad \underset{r \rightarrow 0}{\limsup } \frac{\mu(B(\sigma, r))}{g(r)}>c
$$

Then, $\mu(A) \geqslant c \mathscr{H}_{g}(A)$.
Proof. Let $\varepsilon \in\left(0,(2 a) \wedge r_{0}\right)$. Let $U$ be an open set of $T$ such that $A \subset U$. For all $\sigma \in A$, there exists $r_{\sigma} \in(0, \varepsilon)$ such that

$$
\mu\left(\Gamma\left(\sigma, r_{\sigma}\right)\right)=\mu\left(B\left(\sigma, r_{\sigma}\right)\right)>c g\left(r_{\sigma}\right) \quad \text { and } \quad \Gamma\left(\sigma, r_{\sigma}\right) \subset U .
$$

Thus, $A \subset \bigcup_{\sigma \in A} \Gamma\left(\sigma, r_{\sigma}\right) \subset U$. Then, Proposition 1 (ii) asserts that the set of all $T(a)$-balls is countable and Proposition 1 (iii) asserts that two $T(a)$-balls are either contained one in the other or disjoint. Therefore, there exists $I \subset \mathbb{N}$ and $\sigma_{n} \in A, n \in I$, such that the $\Gamma\left(\sigma_{n}, r_{\sigma_{n}}\right)$, $n \in I$, are pairwise disjoint and $A \subset \bigcup_{n \in I} \Gamma\left(\sigma_{n}, r_{\sigma_{n}}\right) \subset U$. Moreover, by Proposition 1 (ii), $\operatorname{diam}\left(\Gamma\left(\sigma_{n}, r_{\sigma_{n}}\right)\right) \leqslant r_{\sigma_{n}}$. Thus, we get

$$
\begin{aligned}
c \mathscr{H}_{g}^{(\varepsilon)}(A) & \leqslant \sum_{n \in I} c g\left(\operatorname{diam}\left(\Gamma\left(\sigma_{n}, r_{\sigma_{n}}\right)\right)\right) \leqslant \sum_{n \in I} c g\left(r_{\sigma_{n}}\right) \\
& \leqslant \sum_{n \in I} \mu\left(\Gamma\left(\sigma_{n}, r_{\sigma_{n}}\right)\right)=\mu\left(\bigcup_{n \in I} \Gamma\left(\sigma_{n}, r_{\sigma_{n}}\right)\right) \leqslant \mu(U) .
\end{aligned}
$$

As $\varepsilon \rightarrow 0$, it entails $c \mathscr{H}_{g}(A) \leqslant \mu(U)$, for all open set $U$ containing $A$. Since $\mu$ is a finite Borel measure, it is outer-regular for the open subsets, which implies the desired result.

## 3 Preliminary results on the Brownian tree.

### 3.1 Basic facts on the Brownian excursion.

We work under the excursion measure $\mathbf{N}$ defined in the introduction and $e$ denote the canonical excursion whose duration is denoted by $\zeta$ (see (6)). We shall denote by ( $\mathcal{T}, d, \rho$ ) the compact rooted real tree coded by $e$.

The branching property. Fix $b \in(0, \infty)$. We discuss here a decomposition of $e$ in terms of its excursions above level $b$; this yields a decomposition of the Brownian tree called the branching property. To that end we first introduce the following time change: for all $t \in[0, \infty)$, we set

$$
\begin{equation*}
\tau_{b}(t)=\inf \left\{s \in[0, \infty): \int_{0}^{s} \mathbf{1}_{\left\{e_{u} \leqslant b\right\}} \mathrm{d} u>t\right\} \quad \text { and } \quad \widetilde{e}_{b}(t)=e\left(\tau_{b}(t)\right) . \tag{23}
\end{equation*}
$$

Note that $\left(\widetilde{e}_{b}(t)\right)_{t \in[0, \infty)}$ codes the tree below $b$ namely $\{\sigma \in \mathcal{T}: d(\rho, \sigma) \leqslant b\}$ that is the closed ball with center $\rho$ and radius $b$. We denote by $\mathcal{G}_{b}$, the sigma-field generated by $\left(\widetilde{e}_{b}(t)\right)_{t \in[0, \infty)}$ and completed with the $\mathbf{N}$-negligible sets. The approximation (10) implies that $L_{\zeta}^{b}$ is $\mathcal{G}_{b}$-measurable. Then denote by $\left(\alpha_{j}, \beta_{j}\right), j \in \mathcal{J}_{b}$, the connected components of the time-set $\{s \in[0, \infty): e(s)>b\}$. Namely,

$$
\bigcup_{j \in \mathcal{J}_{b}}\left(\alpha_{j}, \beta_{j}\right)=\{s \in[0, \infty): e(s)>b\}
$$

and we call $\left(\alpha_{j}, \beta_{j}\right)$ the excursion intervals of $e$ above level $b$. For all $j \in \mathcal{J}_{b}$, we next set

$$
l_{j}^{b}=L_{\tau_{b}\left(\alpha_{j}\right)}^{b} \quad \text { and } \quad \forall s \in[0, \infty), \quad e_{j}^{b}(s)=e_{\left(\alpha_{j}+s\right) \wedge \beta_{j}}-b
$$

Then, the $\left(e_{j}^{b}\right)_{j \in \mathcal{J}_{b}}$ are the excursions of $e$ above level $b$. Recall from (7) and (8) the notation $\mathbf{N}_{b}=\mathbf{N}(\cdot \mid \sup e>b)$, that is a probability measure. The branching property asserts the following: under $\mathbf{N}_{b}$ and conditionally on $\mathcal{G}_{b}$, the measure

$$
\begin{equation*}
\mathcal{M}_{b}(\mathrm{~d} l, \mathrm{~d} e)=\sum_{j \in \mathcal{J}_{b}} \delta_{\left(l_{j}^{b}, e_{j}^{b}\right)} \tag{24}
\end{equation*}
$$

is a Poisson point measure on $\left[0, L_{\zeta}^{b}\right] \times \mathbf{C}^{0}$ with intensity $\mathbf{1}_{\left[0, L_{\zeta}^{b}\right.}(l) \mathrm{d} l \mathbf{N}(\mathrm{~d} e)$.

The following decomposition of $e$ is interpreted in terms of the Brownian tree $\mathcal{T}$ as follows. Recall that $p:[0, \zeta] \rightarrow \mathcal{T}$ stands for the canonical projection. Then for all $j \in \mathcal{J}_{b}$, we set

$$
\sigma_{j}=p\left(\alpha_{j}\right)=p\left(\beta_{j}\right), \quad \mathcal{T}_{j}^{o, b}=p\left(\left(\alpha_{j}, \beta_{j}\right)\right) \quad \text { and } \quad \mathcal{T}_{j}^{b}=p\left(\left[\alpha_{j}, \beta_{j}\right]\right) .
$$

Then, we easily check that the $\mathcal{T}_{j}^{o, b}, j \in \mathcal{J}_{b}$, are the connected components of the open subset $\{\sigma \in \mathcal{T}: d(\rho, \sigma)>b\}$ and that $\mathcal{T}_{j}^{o, b}=\mathcal{T}_{j}^{b} \backslash\left\{\sigma_{j}\right\}$. Namely, the $\left(\mathcal{T}_{j}, d, \sigma_{j}\right), j \in \mathcal{J}_{b}$ are the subtrees above level $b$ of $\mathcal{T}$ as introduced in Section 2.1. Moreover note that for all $j \in \mathcal{J}_{b}$, the rooted compact real tree $\left(\mathcal{T}_{j}, d, \sigma_{j}\right)$ is isometric to the tree coded by the excursion $e_{j}^{b}$. We next use this and Proposition 1 to discuss the balls in a fixed level of $\mathcal{T}$.

To that end, we fix $a, r \in(0, \infty)$ such that $a>r / 2$ and we conveniently set $b=a-r / 2$. Recall that $\mathcal{T}(a)=\{\sigma \in \mathcal{T}: d(\rho, \sigma)=a\}$ and that for all $\sigma \in \mathcal{T}(a)$, we have set $\Gamma(\sigma, r)=\mathcal{T}(a) \cap B(\sigma, r)$ that is the ball in $\mathcal{T}(a)$ with center $\sigma$ and radius $r$. We also recall that $\mathscr{B}_{a, r}=\{\Gamma(\sigma, r) ; \sigma \in \mathcal{T}(a)\}$ stands for the set of all $\mathcal{T}(a)$-balls with radius $r$. By Proposition $1, \mathscr{B}_{a, r}$ is a finite set and that

$$
\mathscr{B}_{a, r}=\left\{\mathcal{T}(a) \cap \mathcal{T}_{j}^{b} ; j \in \mathcal{J}_{b}: h\left(\mathcal{T}_{j}^{b}\right) \geqslant r / 2\right\}
$$

where the trees $\left(\mathcal{T}_{j}^{b}, d, \sigma_{j}\right), j \in \mathcal{J}_{b}$, are the subtrees of $\mathcal{T}$ above level $b$ as previously defined; here $h\left(\mathcal{T}_{j}^{b}\right)=\sup _{\sigma \in \mathcal{T}_{j}^{b}} d\left(\sigma_{j}, \sigma\right)$ stands for the total height of $\mathcal{T}_{j}^{b}$. Note that $h\left(\mathcal{T}_{j}^{b}\right)=\sup e_{j}^{b}$ that is maximum of the excursion corresponding to $\mathcal{T}_{j}^{b}$, as explained above.

Then, we set $Z_{a, r}=\# \mathscr{B}_{a, r}$, that is the number of $\mathcal{T}(a)$-ball with radius $r$. Assume that $Z_{a, r} \geqslant 1$. We then define the indices $j_{1}, \ldots, j_{Z_{a, r}} \in \mathcal{J}_{b}$ by

$$
\left\{j_{1}, \ldots, j_{Z_{a, r}}\right\}=\left\{j \in \mathcal{J}_{b}: h\left(\mathcal{T}_{j}^{b}\right) \geqslant r / 2\right\} \quad \text { and } \quad \alpha_{j_{1}}<\ldots<\alpha_{j_{Z_{a, r}}}
$$

and we set

$$
\begin{equation*}
\forall i \in\left\{1, \ldots, Z_{a, r}\right\}, \quad \Gamma_{i}:=\mathcal{T}(a) \cap \mathcal{T}_{j_{i}}^{b} . \tag{25}
\end{equation*}
$$

Namely $\mathscr{B}_{a, r}=\left\{\Gamma_{i} ; 1 \leqslant i \leqslant Z_{a, r}\right\}$ is the set of the $\mathcal{T}(a)$-balls with radius $r$ listed in their order of visit by the excursion $e$ coding $\mathcal{T}$.

Lemma 4 Let $a, r \in(0, \infty)$ such that $a>r / 2$. Let $\left\{\Gamma_{i} ; 1 \leqslant i \leqslant Z_{a, r}\right\}$ is the set of the $\mathcal{T}(a)$-balls with radius $r$ listed in their order of visit as explained above. Then the following holds true.
(i) Under $\mathbf{N}_{a}=\mathbf{N}(\cdot \mid \sup e>a), Z_{a, r}$ has a geometric law with parameter $2 a / r$. Namely,

$$
\forall k \geqslant 1, \quad \mathbf{N}_{a}\left[Z_{a, r}=k\right]=\left(1-\frac{r}{2 a}\right)^{k-1} \frac{r}{2 a} .
$$

(ii) For all $k \geqslant 1$, under $\mathbf{N}_{a}\left(\cdot \mid Z_{a, r}=k\right)$, the r.v. $\left(\ell^{a}\left(\Gamma_{i}\right)\right)_{1 \leqslant i \leqslant k}$ are independent and exponentially distributed with mean $r / 2$.

Proof. Let $a \in(0, \infty)$ and denote $b=a-r / 2$. Let $k \geqslant 1$ and $F_{1}, \ldots F_{k}: \mathbf{C}^{0} \rightarrow[0, \infty)$ be measurable functionals. Recall from (7) that $\mathbf{N}(\sup e \geqslant r / 2)=2 / r$. Then, the definition of the $j_{i}$ combined with the branching property and basic results on Poisson point measures entail

$$
\begin{equation*}
\mathbf{N}_{b}\left[\mathbf{1}_{\left\{Z_{a, r}=k\right\}} \prod_{1 \leqslant i \leqslant k} F_{i}\left(e_{j_{i}}^{b}\right) \mid \mathcal{G}_{b}\right]=\frac{\left(\frac{2}{r} L_{\zeta}^{b}\right)^{k}}{k!} e^{-\frac{2}{r} L_{\zeta}^{b}} \prod_{1 \leqslant i \leqslant k} \mathbf{N}_{r / 2}\left[F_{i}(e)\right] . \tag{26}
\end{equation*}
$$

Then recall (12) that implies that $L_{\zeta}^{b}$ under $\mathbf{N}_{b}$ is exponentially distributed with mean $b$. Thus,

$$
\frac{1}{k!} \mathbf{N}_{b}\left[\left(\frac{2}{r} L_{\zeta}^{b}\right)^{k} e^{-\frac{2}{r} L_{\zeta}^{b}}\right]=\frac{\left(\frac{2}{r} b\right)^{k}}{\left(1+\frac{2}{r} b\right)^{k+1}}=\frac{r}{2 a}\left(1-\frac{r}{2 a}\right)^{k}
$$

because $b=a-r / 2$ and $\left(1+\frac{2}{r} b\right)^{-1}=r /(2 a)$. It implies

$$
\mathbf{N}_{b}\left[\mathbf{1}_{\left\{Z_{a, r}=k\right\}} \prod_{1 \leqslant i \leqslant k} F_{i}\left(e_{j_{i}}^{b}\right)\right]=\frac{r}{2 a}\left(1-\frac{r}{2 a}\right)^{k} \prod_{1 \leqslant i \leqslant k} \mathbf{N}_{r / 2}\left[F_{i}(e)\right]
$$

Next observe that $\mathbf{N}_{b-\text {-a.s. }} \mathbf{1}_{\{\sup e>a\}}=\mathbf{1}_{\left\{Z_{a, r} \geqslant 1\right\}}$. Thus, we get

$$
\begin{align*}
\mathbf{N}_{a}\left[\mathbf{1}_{\left\{Z_{a, r}=k\right\}} \prod_{1 \leqslant i \leqslant k} F_{i}\left(e_{j_{i}}^{b}\right)\right] & =\frac{a}{b} \mathbf{N}_{b}\left[\mathbf{1}_{\left\{Z_{a, r}=k\right\}} \prod_{1 \leqslant i \leqslant k} F_{i}\left(e_{j_{i}}^{b}\right)\right] \\
& =\frac{r}{2 a}\left(1-\frac{r}{2 a}\right)^{k-1} \prod_{1 \leqslant i \leqslant k} \mathbf{N}_{r / 2}\left(F_{i}(e)\right) \tag{27}
\end{align*}
$$

because $a / b=\left(1-\frac{r}{2 a}\right)^{-1}$. Recall that (12) implies that under $\mathbf{N}_{r / 2}, \ell^{r / 2}(\mathcal{T})=L_{\zeta}^{r / 2}$ is exponentially distributed with mean $r / 2$. By taking $F_{i}(e)=f_{i}\left(L_{\zeta}^{r / 2}\right)$ in (27) we then get

$$
\mathbf{N}_{a}\left[\mathbf{1}_{\left\{Z_{a, r}=k\right\}} \prod_{1 \leqslant i \leqslant k} f_{i}\left(\ell^{a}\left(\Gamma_{i}\right)\right)\right]=\frac{r}{2 a}\left(1-\frac{r}{2 a}\right)^{k-1} \prod_{1 \leqslant i \leqslant k} \int_{0}^{\infty} f_{i}(s) \frac{2}{r} e^{-\frac{2}{r} s} \mathrm{~d} s
$$

with entails the desired result.

Ray-Knight theorem under N. We first recall the definition of Feller diffusion, namely a Continuous States space Branching Process (CSBP) with branching mechanism $\psi(\lambda)=\lambda^{2}$. Let $x \in[0, \infty)$ and let $\left(Y_{a}^{x}\right)_{a \in[0, \infty)}$ be a $[0, \infty)$-valued continuous process defined on the probability space $(\Omega, \mathcal{F}, \mathbf{P})$. It is a Feller diffusion with branching mechanism $\psi(\lambda)=\lambda^{2}$ and initial value $Y_{0}^{x}=x$ if it is a Markov process such that

$$
\mathbf{E}\left[\exp \left(-\lambda Y_{a+a^{\prime}}^{x}\right) \mid Y_{a}^{x}\right]=\exp \left(-\frac{Y_{a}^{x} \lambda}{1+a^{\prime} \lambda}\right), \quad a, a^{\prime}, \lambda \in[0, \infty)
$$

Recall notation $\mathbf{N}_{a}=\mathbf{N}(\cdot \mid \sup e>a)$ and $\mathcal{G}_{a}$ for the sigma-field generated by the excursion $\widetilde{e}_{a}$ defined in (23). Recall that $\ell^{a}(\mathcal{T})=L_{\zeta}^{a}$, the total mass of the local-time measure at level $a$, is $\mathcal{G}_{a}$-measurable.
We shall use the following statement of Ray-Knight theorem. Let $a \in(0, \infty)$.
(i) $\mathbf{N}_{a}\left[\exp \left(-\lambda \ell^{a}(\mathcal{T})\right)\right]=\frac{1}{1+a \lambda}$.
(ii) Under $\mathbf{N}_{a}$ and conditionally given $\mathcal{G}_{a}$, the process $\left(\ell^{a+a^{\prime}}(\mathcal{T})\right)_{a^{\prime} \in[0, \infty)}$ is a Feller diffusion with branching mechanism $\psi(\lambda)=\lambda^{2}$ and initial value $\ell^{a}(\mathcal{T})$.

This is an immediate consequence of the Ray-Knight theorem for standard Brownian motion and of the Markov property under $\mathbf{N}$ : see [4] III 3 and VI 2.10.

Combined with the branching property, the above Ray-Knight theorem, has the following consequence. Let us recall that we enumerate the $\mathcal{T}(a)$-balls of $\mathscr{B}_{a, r}$ as $\left\{\Gamma_{i}, 1 \leqslant i \leqslant Z_{a, r}\right\}$ (see (25)). Let $\Gamma$ such a $\mathcal{T}(a)$-ball. For $a^{\prime} \geqslant 0$, we define

$$
\begin{equation*}
\Gamma^{a+a^{\prime}}=\left\{\sigma \in \mathcal{T}\left(a+a^{\prime}\right) \exists \sigma^{\prime} \in \Gamma: \sigma^{\prime} \in \llbracket \rho, \sigma \rrbracket\right\}, \tag{28}
\end{equation*}
$$

the set of vertices at level $a+a^{\prime}$ that have an ancestor in $\Gamma$ (notice that $\Gamma^{a}=\Gamma$ ). The following lemma is a straightforward consequence of Ray-Knight theorem.

Lemma 5 Let $a \in(0, \infty), r \in[0,2 a]$. Let $\left\{\Gamma_{i}, 1 \leqslant i \leqslant Z_{a, r}\right\}$ the set of $\mathcal{T}(a)$-balls of radius $r$. Under $\mathbf{N}_{a}$ conditionally on $\mathcal{G}_{a}$, the processes $\left(\ell^{a+a^{\prime}}\left(\Gamma_{i}^{a+a^{\prime}}\right), a^{\prime} \geqslant 0\right), 1 \leqslant i \leqslant Z_{a, r}$, are independent Feller diffusions started at $\left(\ell^{a}\left(\Gamma_{i}\right)\right), 1 \leqslant i \leqslant Z_{a, r}$.
Proof. Recalling for $b=a-r / 2$ the decomposition (25), we see that

$$
\begin{equation*}
\forall i \in\left\{1, \ldots, Z_{a, r}\right\}, \quad \Gamma_{i}^{a+a^{\prime}}:=\mathcal{T}\left(a+a^{\prime}\right) \cap \mathcal{T}_{j_{i}}^{b} . \tag{29}
\end{equation*}
$$

Hence, one can use (26), and the Ray-Knight theorem (see (ii) above) to get the desired result.
Spinal decomposition. We recall another decomposition of the Brownian tree called spinal decomposition. This is a consequence of Bismut's decomposition of the Brownian excursion that we recall here.

Let $X$ be a real valued process defined on $(\Omega, \mathcal{F}, \mathbf{P})$ such that $\left(\frac{1}{\sqrt{2}} X_{t}\right)_{t \in[0, \infty)}$ is distributed as a standard Brownian motion with initial value 0 . Let $X^{\prime}$ be an independent copy of $X$ on $(\Omega, \mathcal{F}, \mathbf{P})$. We fix $a \in(0, \infty)$ and we set

$$
T_{a}=\inf \left\{t \in[0, \infty): X_{t}=-a\right\} \quad \text { and } \quad T_{a}^{\prime}=\inf \left\{t \in[0, \infty): X_{t}^{\prime}=-a\right\}
$$

We next set for any $s \in[0, \infty)$,

$$
\check{e}_{s}^{t}=e_{(t-s)_{+}} \quad \text { and } \quad \hat{e}_{s}^{t}=e_{t+s}
$$

Then the Bismut's identity (see [3] or [11]) states that for any non-negative measurable functional $F$ on $\left(\mathbf{C}^{0}\right)^{2}$,

$$
\begin{equation*}
\mathbf{N}\left[\int_{0}^{\zeta} \mathrm{d} L_{t}^{a} F\left(\tilde{e}^{t} ; \hat{e}^{t}\right)\right]=\mathbf{E}\left[F\left(a+X_{\cdot \wedge T_{a}} ; a+X_{\cdot \wedge T_{a}^{\prime}}^{\prime}\right)\right] \tag{30}
\end{equation*}
$$

We derive from (30) an identity involving the excursions above the infimum of $\hat{e}^{t}$ and $\check{e}^{t}$. To that end, we introduce the following. Let $h:[0, \infty) \rightarrow[0, \infty)$ with compact support. We define a point measure point $\mathcal{N}(h)$ as follows: set $\underline{h}(t)=\inf _{[0, t]} h$ and denote by $\left(g_{i}, d_{i}\right), i \in \mathcal{I}(h)$ the excursion intervals of $h-\underline{h}$ away from 0 that are the connected component of the open set $\{t \geqslant 0: h(t)-\underline{h}(t)>0\}$. For any $i \in \mathcal{I}(h)$, set $h^{i}(s)=\left((h-\underline{h})\left(\left(g_{i}+s\right) \wedge d_{i}\right), s \geqslant 0\right)$. We then define $\mathcal{N}(h)$ as the point measure on $[0, \infty) \times \mathbf{C}^{0}$ given by

$$
\mathcal{N}(h)=\sum_{i \in \mathcal{I}(h)} \delta_{\left(h\left(g_{i}\right), h^{i}\right)} .
$$

Then, for any $t, a \in(0, \infty)$,

$$
\begin{equation*}
\mathcal{N}_{t}:=\mathcal{N}\left(e^{t}\right)+\mathcal{N}\left(\hat{e}^{t}\right)=: \sum_{j \in \mathcal{J}_{t}} \delta_{\left(h_{j}^{t}, e^{t, j}\right)} \tag{31}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathcal{N}_{a}^{*}:=\mathcal{N}\left(a+X_{\cdot \wedge T_{a}}\right)+\mathcal{N}\left(a+X_{\cdot \wedge T_{a}^{\prime}}^{\prime}\right)=: \sum_{j \in \mathcal{J}_{a}^{*}} \delta_{\left(h_{j}^{*}, e^{* j}\right)} . \tag{32}
\end{equation*}
$$

We deduce from (30) that for all any $a$ and for all nonnegative measurable function $F$ on the set of positive measures on $[0, \infty) \times \mathcal{C}^{0}$, one has

$$
\begin{equation*}
\mathbf{N}\left[\int_{0}^{\zeta} \mathrm{d} L_{t}^{a} F\left(\mathcal{N}_{t}\right)\right]=\mathbf{E}\left[F\left(\mathcal{N}_{a}^{*}\right)\right] \tag{33}
\end{equation*}
$$

and as consequence of Itô's decomposition of Brownian motion above its infimum, $\mathcal{N}_{a}^{*}$ is a Poisson point measure on $[0, \infty) \times \mathbf{C}^{0}$ with intensity $2 \mathbf{1}_{[0, a]}(h) \mathrm{d} h \mathbf{N}(\mathrm{~d} e)$.

Let us interpret this decomposition in terms of the Brownian tree. Choose $t \in(0, \zeta)$ such that $e_{t}=a$ and set $\sigma=p(t) \in \mathcal{T}$ (namely $\sigma \in \mathcal{T}(a)$ ). Then, the geodesic $\llbracket \rho, \sigma \rrbracket$ is interpreted as the ancestral line of $\sigma$. Let us denote by $\mathcal{T}_{j}^{o}, j \in \mathcal{J}$, the connected components of the open set $\mathcal{T} \backslash \llbracket \rho, \sigma \rrbracket$ and denote by $\mathcal{T}_{j}$ the closure of $\mathcal{T}_{j}^{o}$. Then, there exists a point $\sigma_{j} \in \llbracket \rho, \sigma \rrbracket$ such that $\mathcal{T}_{j}=\left\{\sigma_{j}\right\} \cup \mathcal{T}_{j}^{o}$. Recall notation $\left(h_{j}^{t}, e^{t, j}\right), j \in \mathcal{J}_{t}$ from (31). The specific coding of $\mathcal{T}$ by $e$ entails that for any $j \in \mathcal{J}$ there exists a unique $j^{\prime} \in \mathcal{J}_{t}$ such that $d\left(\rho, \sigma_{j}\right)=h_{j^{\prime}}^{t}$ and such that the rooted compact real tree $\left(\mathcal{T}_{j}, d, \sigma_{j}\right)$ is isometric to the tree coded by $e^{t, j^{\prime}}$.

Recall that $p(t)=\sigma$. We fix $r, r^{\prime} \in[0,2 a)$ such that $r^{\prime} \leqslant r$. We now compute the mass of the ring $B(\sigma, r) \backslash B\left(\sigma, r^{\prime}\right)$ in terms of $\mathcal{N}_{t}$. First, observe that for any $s \in[0, \zeta]$ such that $e_{s}=a$, we have

$$
r^{\prime} \leqslant d(s, t)<r \Longleftrightarrow a-\left(r^{\prime} / 2\right) \geqslant \inf _{u \in[s \wedge t, s \vee t]} e_{u}>a-(r / 2)
$$

We then get

$$
\begin{equation*}
\ell^{a}\left(B(\sigma, r) \backslash B\left(\sigma, r^{\prime}\right)\right)=\sum_{j \in \mathcal{J}_{t}} \mathbf{1}_{\left(a-\frac{r}{2}, a-\frac{r^{\prime}}{2}\right]}\left(h_{j}^{t}\right) L_{\zeta_{j}^{t}}^{a-h_{j}^{t}}(t, j) \tag{34}
\end{equation*}
$$

where $L_{\zeta_{j}^{t}}^{a-h_{j}^{t}}(t, j)$ stands for the local time at level $a-h_{j}^{t}$ of the excursion $e^{t, j}$.
Then, for any $a \in(0, \infty)$ and any $r, r^{\prime} \in(0,2 a)$ such that $r^{\prime} \leqslant r$, we also set

$$
\begin{equation*}
\Lambda_{r^{\prime}, r}^{a}=\sum_{j \in \mathcal{J}_{a}^{*}} \mathbf{1}_{\left(a-\frac{r}{2}, a-\frac{r^{\prime}}{2}\right]}\left(h_{j}^{*}\right) L_{\zeta_{j}^{*}}^{a-h_{j}^{*}} \tag{35}
\end{equation*}
$$

where, $L_{\zeta_{j}^{*}}^{a-h_{j}^{*}}$ stands for the local time at level $a-h_{j}^{*}$ of the excursion $e^{* j}$ defined in (32). Then, (33) implies that for any $a \in(0, \infty)$ and for all non-negative measurable $F$ function

$$
\begin{equation*}
\mathbf{N}\left[\int_{\mathcal{T}} \ell^{a}(\mathrm{~d} \sigma) F\left(\ell^{a}\left(B(\sigma, r) \backslash B\left(\sigma, r^{\prime}\right)\right) ; 0 \leqslant r^{\prime} \leqslant r \leqslant 2 a\right)\right]=\mathbf{E}\left[F\left(\Lambda_{r^{\prime}, r}^{a} ; 0 \leqslant r^{\prime} \leqslant r \leqslant 2 a\right)\right] \tag{36}
\end{equation*}
$$

On the right-hand-side, the dependency with respect to the level $a$ is a bit artificial. Indeed, for $a \in(0, \infty)$, the Poisson point measure $\mathcal{N}_{a}^{*}(\mathrm{~d} h \mathrm{~d} e)$ has its law invariant under the transformation $(h, e) \mapsto(a-h, e)$. Thus, let us consider on $(\Omega, \mathcal{F}, \mathbf{P})$ a new Poisson point measure $\mathcal{M}^{*}=$ $\sum_{j \in \mathcal{I}^{*}} \delta_{\left(h_{j}^{*}, e_{j}^{*}\right)}$ with intensity $2 \mathrm{~d} h \mathbf{N}(\mathrm{~d} e)$ (we abuse notations and keep the notation $\left(h_{j}^{*}, e_{j}^{*}\right)$ for the atoms). We set

$$
\begin{equation*}
\Lambda_{r^{\prime}, r}^{*}=\sum_{j \in \mathcal{I}^{*}} \mathbf{1}_{\left[\frac{r^{\prime}}{2}, \frac{r}{2}\right)}\left(h_{j}^{*}\right) L_{\zeta_{j}^{*}}^{h_{j}^{*}} \tag{37}
\end{equation*}
$$

where $L_{\zeta_{j}^{*}}^{h_{*}^{*}}$ stands for the local time at height $h_{j}^{*}$ for the excursion $e_{j}^{*}$. One can now rewrite (36) as

$$
\begin{align*}
& \mathbf{N}\left[\int_{\mathcal{T}} \ell^{a}(\mathrm{~d} \sigma) F\left(\ell^{a}\left(B(\sigma, r) \backslash B\left(\sigma, r^{\prime}\right)\right) ; 0 \leqslant r^{\prime} \leqslant r \leqslant 2 a\right)\right] \\
&=\mathbf{E}\left[F\left(\Lambda_{r^{\prime}, r}^{*} ; 0 \leqslant r^{\prime} \leqslant r \leqslant 2 a\right)\right] \tag{38}
\end{align*}
$$

The law of the $\Lambda_{r^{\prime}, r}^{*}$ is quite explicit as shown by the following lemma.
Lemma 6 Let $0 \leqslant r_{n} \leqslant r_{n-1} \leqslant \ldots \leqslant r_{1} \leqslant 2 a$. Then,

$$
\Lambda_{r_{n}, r_{n-1}}^{*}, \Lambda_{r_{n-1}, r_{n-2}}^{*}, \ldots, \Lambda_{r_{2}, r_{1}}^{*}
$$

are independent. Moreover, for any $0 \leqslant r^{\prime} \leqslant r \leqslant 2 a$,

$$
\forall y \in(0, \infty) \quad \mathbf{P}\left(\Lambda_{r^{\prime}, r}^{*}>y\right)=\left(1-\frac{r^{\prime}}{r}\right)^{2} \frac{2 y}{r} e^{-2 y / r}+\left(1-\left(\frac{r^{\prime}}{r}\right)^{2}\right) e^{-2 y / r}
$$

and $\mathbf{P}\left(\Lambda_{r^{\prime}, r}^{*}=0\right)=\left(r^{\prime} / r\right)^{2}$.
Proof. The intervals $\left[r_{k+1} / 2, r_{k} / 2\right)$ being pairwise disjoint, the independence of the increments is a straightforward consequence of the properties the Poisson point measure $\mathcal{M}^{*}$. Using Campbell formula and (12), we compute, for all $\lambda \geqslant 0$,

$$
\begin{aligned}
\mathbf{E}\left[e^{-\lambda \Lambda_{2 r^{\prime}, 2 r}^{*}}\right] & =\exp \left(-\int_{r^{\prime}}^{r} 2 \mathrm{~d} h \mathbf{N}\left[1-e^{-\lambda \ell^{h}(\mathcal{T})}\right]\right) \\
& =\exp \left(-\int_{r^{\prime}}^{r} 2 \mathrm{~d} h \frac{\lambda}{1+h \lambda}\right)=\left(\frac{1+r^{\prime} \lambda}{1+r \lambda}\right)^{2}
\end{aligned}
$$

Thus, $\Lambda_{2 r^{\prime}, 2 r}^{*} \stackrel{(\text { law })}{=} X_{1}+X_{2}$, where $X_{1}$ and $X_{2}$ are i.i.d random variables where

$$
\mathbf{E}\left[e^{-\lambda X_{1}}\right]=\frac{r^{\prime}}{r}+\left(1-\frac{r^{\prime}}{r}\right) \frac{1}{1+r \lambda}
$$

Thus, $X_{1}=0$ with probability $r^{\prime} / r$ and conditionally on being non-zero, it is exponentially distributed with mean $r$. Thus, for $y>0$,

$$
\begin{aligned}
\mathbf{P}\left(\Lambda_{2 r^{\prime}, 2 r}^{*}>y\right) & =2 \mathbf{P}\left(X_{1}=0 ; X_{2}>y\right)+\mathbf{P}\left(X_{1}>0 ; X_{2}>0 ; X_{1}+X_{2}>y\right) \\
& =2 \frac{r^{\prime}}{r} \mathbf{P}\left(X_{1}>y\right)+\left(1-\frac{r^{\prime}}{r}\right)^{2} \mathbf{P}(Z>y)
\end{aligned}
$$

where $Z$ has law $\operatorname{Gamma}(2,1 / r)$. The result proceeds now from elementary computations.

### 3.2 Estimates.

The following elementary computation is needed twice in our proofs.

Lemma 7 Let $\left(X_{n}\right)_{n \geqslant 1}$ a sequence of i.i.d real valued random variables on $(\Omega, \mathcal{F}, \mathbf{P})$, with mean 0 and a moment of order 4 . Let $Z$ be a random variable taking its values in $\mathbb{N}$, independent of the sequence $\left(X_{n}\right)$. Then

$$
\mathbf{E}\left[\left(X_{1}+X_{2}+\cdots+X_{Z}\right)^{4}\right] \leqslant 3 \mathbf{E}\left[X_{1}^{4}\right] \mathbf{E}\left[Z^{2}\right] .
$$

Moreover, the following holds : $\mathbf{E}\left[\left(X_{1}-\mathbf{E}\left[X_{1}\right]\right)^{4}\right] \leqslant 2 \mathbf{E}\left[X_{1}^{4}\right]$.
Proof. One has

$$
\mathbf{E}\left[\left(X_{1}+X_{2}+\cdots+X_{Z}\right)^{4} \mid Z\right]=\sum_{1 \leqslant i_{1}, i_{2}, i_{3}, i_{4} \leqslant Z} \mathbf{E}\left[X_{i_{1}} X_{i_{2}} X_{i_{3}} X_{i_{4}}\right]
$$

When $\left(i_{1}, i_{2}, i_{3}, i_{4}\right)$ contains an index that is distinct of the three others, then the contribution of the corresponding term will be null. Thus the latter mean equals $Z \mathbf{E}\left[X_{1}^{4}\right]+3 Z(Z-1) \mathbf{E}\left[X_{1}^{2}\right]^{2} \leqslant$ $3 Z^{2} \mathbf{E}\left[X_{1}^{4}\right]$ (using Jensen's inequality). The second statement follows from

$$
\begin{aligned}
\mathbf{E}\left[\left(X_{1}-\mathbf{E}\left[X_{1}\right]\right)^{4}\right] & =\mathbf{E}\left[\left(X_{1}-\mathbf{E}\left[X_{1}\right]\right)^{4} \mathbf{1}_{\left\{X_{1} \geqslant \mathbf{E}\left[X_{1}\right]\right\}}\right]+\mathbf{E}\left[\left(X_{1}-\mathbf{E}\left[X_{1}\right]\right)^{4} \mathbf{1}_{\left\{X_{1}<\mathbf{E}\left[X_{1}\right]\right\}}\right] \\
& \leqslant \mathbf{E}\left[X_{1}^{4}\right]+\mathbf{E}\left[X_{1}\right]^{4},
\end{aligned}
$$

and using Jensen's inequality.
We explained in Section 3.1 the link between the process $\left(\ell^{a}(\mathcal{T}), a \in(0, \infty)\right)$ and the Feller diffusion, for which we provide here some basic estimates.

Lemma 8 Let $\left(Y_{a}^{x}\right)_{a \geqslant 0}$ be a Feller diffusion starting at $x \geqslant 0$, defined on $(\Omega, \mathcal{F}, \mathbf{P})$. For all $x, y \in[0, \infty)$, for all $a \in(0, \infty)$, the following inequalities hold :
(i) If $y \leqslant x$, then $\mathbf{P}\left(\inf _{b \in[0, a]} Y_{b}^{x} \leqslant y\right) \leqslant \exp \left(-\frac{1}{a}(\sqrt{x}-\sqrt{y})^{2}\right)$.
(ii) If $y \geqslant x$, then $\mathbf{P}\left(\sup _{b \in[0, a]} Y_{b}^{x} \geqslant y\right) \leqslant \exp \left(-\frac{1}{a}(\sqrt{y}-\sqrt{x})^{2}\right)$.

Proof. Let us prove (i). Recall that for all $x, b, \lambda \in[0, \infty), \mathbf{E}\left[e^{-\lambda Y_{b}^{x}}\right]=\exp \left(-\frac{\lambda x}{1+b \lambda}\right)$. Thus, for fixed $a \in(0, \infty)$, and for $\lambda \in\left[0, \frac{1}{a}\right)$, we set

$$
\begin{equation*}
\forall b \in[0, a], \quad M_{b}^{(\lambda, x)}:=\exp \left(-\frac{\lambda Y_{b}^{x}}{1-b \lambda}\right) . \tag{39}
\end{equation*}
$$

We stress that for $b \in[0, a]$, one has $1-b \lambda \geqslant 1-a \lambda>0$, and one can compute

$$
\mathbf{E}\left[M_{b}^{(\lambda, x)}\right]=\exp \left(-\frac{\lambda}{1-b \lambda} x /\left(1+\frac{b \lambda}{1-b \lambda}\right)\right)=e^{-\lambda x} .
$$

Combined with the Markov property, this entails that $\left(M_{b}^{(\lambda, x)}, b \in[0, a]\right)$ is a martingale. Moreover, on $\left\{\inf _{b \in[0, a]} Y_{b}^{x} \leqslant y\right\}$, one has $\inf _{b \in[0, a]} \frac{\lambda Y_{b}^{x}}{1-b \lambda} \leqslant \inf _{b \in[0, a]} \frac{\lambda Y_{b}^{x}}{1-a \lambda} \leqslant \frac{\lambda y}{1-a \lambda}$. Hence, the maximal inequality
for sub-martingales entails

$$
\begin{aligned}
\mathbf{P}\left(\inf _{b \in[0, a]} Y_{b}^{x} \leqslant y\right) & \leqslant \mathbf{P}\left(\sup _{b \in[0, a]} M_{b}^{(\lambda, x)} \geqslant e^{-\frac{\lambda y}{1-a \lambda}}\right) \\
& \leqslant e^{\frac{\lambda y}{1-a \lambda}} \mathbf{E}\left[M_{a}^{(\lambda, x)}\right]=\exp \left(\frac{\lambda y}{1-a \lambda}-\lambda x\right)
\end{aligned}
$$

The reader can check using elementary computations that the function $\lambda \mapsto \frac{\lambda y}{1-a \lambda}-\lambda x$ has a negative minimum on $(0,1 / a)$ at the value $\lambda=\frac{1}{a}\left(1-\sqrt{\frac{y}{x}}\right)$, and this minimum is $-\frac{1}{a}(\sqrt{x}-\sqrt{y})^{2}$, which completes the proof.

In order to prove $(i i)$, one could extend the definition of $\left(M_{b}^{(\lambda, x)}, b \in[0, a]\right)$ for $\lambda \in(-1 / a, 0)$. In what follows, we use a simpler argument. Let us begin with the following remark: let $b \in(0, \infty)$, let $\mathcal{E}$ be a r.v. on $(\Omega, \mathcal{F}, \mathbf{P})$ that is exponentially distributed with mean $b$, then for all $\lambda \geqslant 0$, $\mathbf{E}\left[e^{-\lambda \mathcal{E}}\right]=\frac{1}{1+b \lambda}$, and this Laplace transform remains finite for $\lambda \in(-1 / b, 0)$. Moreover, one can plainly check that for $x, b \in(0, \infty), Y_{b}^{x}$ has the same law as $\sum_{i=1}^{N} \mathcal{E}_{i}$, where the $\mathcal{E}_{i}$ are independent copies of $\mathcal{E}$ and $N$ is an independent Poisson r.v. with mean $x / b$. Thus, one has

$$
\begin{equation*}
\forall \mu \in(0,1 / b), \quad \mathbf{E}\left[e^{\mu Y_{b}^{x}}\right]=\exp \left(\frac{\mu x}{1-\mu b}\right) \tag{40}
\end{equation*}
$$

The Feller diffusion $\left(Y_{b}^{x}, b \geqslant 0\right)$ is a martingale, so by convexity ( $e^{\mu Y_{b}^{x}}, b \geqslant 0$ ) is a submartingale. Thus, for all $\mu \in(0,1 / a)$, and $y \geqslant x \geqslant 0$, one has

$$
\begin{aligned}
\mathbf{P}\left(\sup _{b \in[0, a]} Y_{b}^{x} \geqslant y\right) & \leqslant \mathbf{P}\left(\sup _{b \in[0, a]} e^{\mu Y_{b}^{x}} \geqslant e^{\mu y}\right) \\
& \leqslant e^{-\mu y} \mathbf{E}\left[e^{\mu Y_{a}^{x}}\right]=\exp \left(\frac{\mu x}{1-a \mu}-\mu y\right)
\end{aligned}
$$

and the result follows by optimizing the same function as before.
The next result is a corollary of Lemma 8 (ii).
Lemma 9 Let $m \in(0,1 / 2)$. For all $y \in(0, \infty)$,

$$
\mathbf{N}\left(\sup _{b \in\left[m, m^{-1}\right]} \ell^{b}(\mathcal{T})>y\right) \leqslant(2 / m) \exp (-m y / 2)
$$

Proof. Let $m \in(0,1 / 2)$ and recall from (23) the definition of $\mathcal{G}_{m}$. As recalled in Section 3.1, under $\mathbf{N}_{m}$, conditionally on $\mathcal{G}_{m}$, the process $\left(\ell^{b}(\mathcal{T}), b \geqslant m\right)$ is a Feller diffusion started at $\ell^{m}(\mathcal{T})$. Hence, conditioning with respect to $\mathcal{G}_{m}$ and using Lemma 8 (ii), we get

$$
\mathbf{N}_{m}\left(\sup _{b \in\left[m, m^{-1}\right]} \ell^{b}(\mathcal{T})>y\right) \leqslant \mathbf{N}_{m}\left[\exp \left(-m\left(\sqrt{y}-\sqrt{\ell^{m}(\mathcal{T})}\right)^{2}\right)\right]
$$

Expanding $(\sqrt{u / 2}-\sqrt{2 v})^{2}$, one shows that for all $u, v \geqslant 0,(\sqrt{u}-\sqrt{v})^{2} \geqslant u / 2-v$. Thus, $\mathbf{N}_{m}\left(\sup _{b \in\left[m, m^{-1}\right]} \ell^{m}(\mathcal{T})>y\right) \leqslant \exp \left(-m \frac{y}{2}\right) \mathbf{N}_{m}\left[e^{m \ell^{m}(\mathcal{T})}\right]$. Recalling from (12) that under $\mathbf{N}_{m}$, $\ell^{m}(\mathcal{T})$ is exponentially distributed with mean $m$, we get $\mathbf{N}_{m}\left[e^{m \ell^{m}(\mathcal{T})}\right]=\left(1-m^{2}\right)^{-1} \leqslant 2$, because $m<1 / 2$. This entails the desired result, recalling that $\mathbf{N}_{m}(\cdot)=m \mathbf{N}\left(\cdot \mathbf{1}_{\{h(\mathcal{T})>m\}}\right)$ and that the events $\{h(\mathcal{T})>m\}$ and $\left\{\ell^{m}(\mathcal{T})>0\right\}$ are equal, up to a $\mathbf{N}$ negligible set.

Estimates for small balls. We consider here a level $a \in(0, \infty)$ and recall that $\mathcal{T}(a)$ is the $a$-level set of the Brownian tree $\mathcal{T}$. If $r \in[0,2 a]$, we recall from (18) the notation $\Gamma(\sigma, r)$ for the $\mathcal{T}(a)$-ball of radius $r$ and center $\sigma \in \mathcal{T}(a)$, the set of $\mathcal{T}(a)$-balls of radius $r$ being denoted $\mathscr{B}_{a, r}$. Let $\Gamma$ be a $\mathcal{T}(a)$-ball of radius $r^{\prime}$, where $r^{\prime} \in[0,2 a]$. From Proposition 1 (iii), we know that if $r \in\left[r^{\prime}, 2 a\right]$, there exists a unique $\mathcal{T}(a)$-ball of radius $r$ that contains $\Gamma$, and we shall denote this "enlarged" ball by

$$
\begin{equation*}
\Gamma[r]:=\Upsilon \quad \text { where } \Upsilon \in \mathscr{B}_{a, r} \text { and } \Gamma \subset \Upsilon \tag{41}
\end{equation*}
$$

We consider positive real numbers $r_{1}>r_{2}>\ldots>r_{n}>0$, and $\varepsilon_{1}>\ldots>\varepsilon_{n-1}>0$, where $n \in \mathbb{N}^{*}$. We set $\mathbf{r}=\left\{r_{1}, \ldots, r_{n}\right\}$ and $\varepsilon=\left\{\varepsilon_{1}, \ldots, \varepsilon_{n-1}\right\}$. We shall say that $\Gamma$, a $\mathcal{T}(a)$-ball of radius $r_{n}$, is $(\mathbf{r}, \boldsymbol{\varepsilon})$-small if and only if for all $1 \leqslant k \leqslant n-1$, the enlarged ball of radius $k$ has a local time smaller than $\varepsilon_{k}$, namely

$$
\begin{equation*}
\forall k \in\{1, \ldots, n-1\} \quad \ell^{a}\left(\Gamma\left[r_{k}\right]\right) \leqslant \varepsilon_{k} \tag{42}
\end{equation*}
$$

We denote by $S_{a, \mathbf{r}, \varepsilon}$ the total number of $\operatorname{such}(\mathbf{r}, \boldsymbol{\varepsilon})$-small balls at level $a$ :

$$
\begin{equation*}
S_{a, \mathbf{r}, \varepsilon}:=\sum_{\Gamma \in \mathscr{B}_{a, r_{n}}} \mathbf{1}_{\{\Gamma \text { is }(\mathbf{r}, \varepsilon)-\text { small }\}} . \tag{43}
\end{equation*}
$$

To control that number, we introduce

$$
\begin{equation*}
\mu(\mathbf{r}, \varepsilon):=\mathbf{N}\left[S_{r_{1} / 2, \mathbf{r}, \varepsilon}\right] \tag{44}
\end{equation*}
$$

Let us stress that its definition does not depend on $a$.
Lemma 10 Let $a \in(0, \infty), \boldsymbol{r}=\left\{r_{1}, \ldots, r_{n}\right\}$, and $\varepsilon=\left\{\varepsilon_{1}, \ldots, \varepsilon_{n-1}\right\}$, where $r_{1}>\ldots>r_{n}>0$, and $\varepsilon_{1}>\ldots>\varepsilon_{n-1}>0$. There exists a constant $c_{0} \in\left(0,10^{4}\right]$ such that if a/r $r_{1}>1$ and $r_{1} / r_{n}>2$,

$$
\mathbf{N}\left[\left(S_{a, \boldsymbol{r}, \boldsymbol{\varepsilon}}-\mu(\boldsymbol{r}, \boldsymbol{\varepsilon}) \ell^{a}(\mathcal{T})\right)^{4}\right] \leqslant c_{0} a \frac{r_{1}^{2}}{r_{n}^{4}}
$$

Proof. Let $a, \mathbf{r}, \boldsymbol{\varepsilon}$ as above. From Proposition 1 (iii), we know that the $\mathcal{T}(a)$-balls of radius $r_{n}$ are disjoint and that for all $\Upsilon \in \mathscr{B}_{a, r_{n}}$, there exists a unique $\mathcal{T}(a)$-ball $\Gamma \in \mathscr{B}_{a, r_{1}}$ such that $\Upsilon \subset \Gamma$. Let us enumerate $\mathscr{B}_{a, r_{1}}$ as $\left\{\Gamma_{i}, 1 \leqslant i \leqslant Z_{a, r_{1}}\right\}$, and set

$$
\begin{aligned}
& \forall i \in\left\{1 \ldots Z_{a, r_{1}}\right\}, \mathscr{B}_{a, r_{n}}^{(i)}=\left\{\Upsilon \in \mathscr{B}_{a, r_{n}}: \Upsilon \subset \Gamma_{i}\right\} \\
& \quad \text { and } \quad S_{a, \mathbf{r}, \varepsilon}^{(i)}=\#\left\{\Upsilon \in \mathscr{B}_{a, r_{n}}^{(i)}: \Upsilon \text { is }(\mathbf{r}, \varepsilon) \text {-small }\right\} .
\end{aligned}
$$

One has

$$
\begin{equation*}
S_{a, \mathbf{r}, \varepsilon}-\mu(\mathbf{r}, \varepsilon) \ell^{a}(\mathcal{T})=\sum_{i=1}^{Z_{a, r_{1}}}\left(S_{a, \mathbf{r}, \varepsilon}^{(i)}-\mu(\mathbf{r}, \varepsilon) \ell^{a}\left(\Gamma_{i}\right)\right)=: \sum_{i=1}^{Z_{a, r_{1}}} X_{i} \tag{45}
\end{equation*}
$$

Let us denote $b=a-r_{1} / 2$ and recall from (23) the definition of the sigma-field $\mathcal{G}_{b}$. Adapting the proof of Lemma 4 , it is not difficult to see that under $\mathbf{N}_{b}$, conditionally on $\mathcal{G}_{b}$, and conditionally on $\left\{Z_{a, r_{1}}=k\right\}$, the r.v. $X_{1}, \ldots X_{k}$ are independent and have the same law as $S_{r_{1} / 2, \mathbf{r}, \boldsymbol{\varepsilon}}-\mu(\mathbf{r}, \boldsymbol{\varepsilon}) \ell^{r_{1} / 2}(\mathcal{T})$ under $\mathbf{N}_{r_{1} / 2}$. Recalling from (12) that $\mathbf{N}\left[\ell^{r_{1} / 2}(\mathcal{T})\right]=\left(2 / r_{1}\right) \mathbf{N}_{r_{1} / 2}\left[\ell^{r_{1} / 2}(\mathcal{T})\right]=$ 1, we see that

$$
\mathbf{N}_{b}\left[X_{1} \mid \mathcal{G}_{b}\right]=\mathbf{N}_{r_{1} / 2}\left[S_{r_{1} / 2, \mathbf{r}, \boldsymbol{\varepsilon}}-\mu(\mathbf{r}, \boldsymbol{\varepsilon}) \ell^{r_{1} / 2}(\mathcal{T})\right]=0
$$

which explains the definition (44). We thus apply Lemma 7 to get from (45):

$$
\begin{equation*}
\mathbf{N}_{b}\left[\left(S_{a, \mathbf{r}, \boldsymbol{\varepsilon}}-\mu(\mathbf{r}, \boldsymbol{\varepsilon}) \ell^{a}(\mathcal{T})\right)^{4} \mid \mathcal{G}_{b}\right] \leqslant 3 \mathbf{N}_{r_{1} / 2}\left[X_{1}^{4}\right] \mathbf{N}_{b}\left[Z_{a, r_{1}}^{2} \mid \mathcal{G}_{b}\right] \tag{46}
\end{equation*}
$$

The second assertion in Lemma 7 entails $\mathbf{N}_{r_{1} / 2}\left[X_{1}^{4}\right] \leqslant 2 \mathbf{N}_{r_{1} / 2}\left[S_{r_{1} / 2, \mathbf{r}, \varepsilon}^{4}\right]$. Moreover, we can use that $S_{r_{1} / 2, \mathbf{r}, \varepsilon}$ is smaller than $Z_{r_{1} / 2}\left(r_{n}\right)$, the total number of $\mathcal{T}\left(r_{1} / 2\right)$-balls of radius $r_{n}$ which has under $\mathbf{N}_{r_{1} / 2}$ a geometric distribution with success probability $r_{n} / r_{1}<1 / 2$. Thus,

$$
\begin{equation*}
\left.\mathbf{N}_{r_{1} / 2}\left[X_{1}^{4}\right] \leqslant 2 \mathbf{N}_{r_{1} / 2}\left[S_{r_{1} / 2, \mathbf{r}, \varepsilon}^{4}\right] \leqslant 2 \mathbf{N}_{r_{1} / 2}\left[Z_{r_{1} / 2}\left(r_{n}\right)\right)^{4}\right] \leqslant \frac{48}{1-r_{n} / r_{1}}\left(\frac{r_{1}}{r_{n}}\right)^{4} \leqslant 96\left(\frac{r_{1}}{r_{n}}\right)^{4} \tag{47}
\end{equation*}
$$

In addition, according to the branching property, under $\mathbf{N}_{b}$, conditionally on $\mathcal{G}_{b}, Z_{a, r_{1}}$ is a Poisson variable with mean $\mathbf{N}\left(h(\mathcal{T})>r_{1} / 2\right) \ell^{b}(\mathcal{T})=\left(2 / r_{1}\right) \ell^{b}(\mathcal{T})$. Thus,

$$
\begin{equation*}
\mathbf{N}_{b}\left[Z_{a, r_{1}}^{2}\right]=\left(2 / r_{1}\right) \mathbf{N}_{b}\left[\ell^{b}(\mathcal{T})\right]+\left(2 / r_{1}\right)^{2} \mathbf{N}_{b}\left[\ell^{b}(\mathcal{T})^{2}\right] \tag{48}
\end{equation*}
$$

Recalling that $b=r_{1} / 2$, we get $\left(2 / r_{1}\right) \mathbf{N}_{b}\left[\ell^{b}(\mathcal{T})\right]=1$ and $\left(2 / r_{1}\right)^{2} \mathbf{N}_{b}\left[\ell^{b}(\mathcal{T})^{2}\right]=8 a^{2} / r_{1}^{2}$. We assumed that $a / r_{1}>1$, thus $\mathbf{N}_{b}\left[Z_{a, r_{1}}^{2}\right] \leqslant 9 a^{2} / r_{1}^{2}$. Combined with (46) and (47) it entails

$$
\mathbf{N}_{b}\left[\left(S_{a, \mathbf{r}, \varepsilon}-\mu(\mathbf{r}, \varepsilon) \ell^{a}(\mathcal{T})\right)^{4}\right] \leqslant c_{0} a^{2} \frac{r_{1}^{2}}{r_{n}^{4}}
$$

with $c_{0}$ a positive constant smaller than $(1 / 2) 10^{4}$. This implies the desired result, using that $\mathbf{N}(h(\mathcal{T})>b)=1 / b \leqslant 2 / a$.

We state now the main technical Lemma of the paper. Let us recall from (44) the definition of $\mu(\mathbf{r}, \boldsymbol{\varepsilon})$. The proof of the lemma makes use of the spinal decomposition described in Section 3.1. In particular, a geometric argument allows to rely the problem to the variables introduced in (37).

Lemma 11 Let $r=\left\{r_{1}, \ldots, r_{n}\right\}$, where $r_{1}>\ldots>r_{n}>0$, and $\varepsilon=\left\{\varepsilon_{1}, \ldots, \varepsilon_{n-1}\right\}$, where $\varepsilon_{1}>\ldots>\varepsilon_{n-1}>0$. The following inequality holds :

$$
\begin{equation*}
\mu(\boldsymbol{r}, \boldsymbol{\varepsilon}) \leqslant \frac{5}{r_{n}} \sqrt{\prod_{k=1}^{n-1} \mathbf{P}\left(\Lambda_{r_{k+1}, r_{k}}^{*} \leqslant \varepsilon_{k}\right)} \tag{49}
\end{equation*}
$$

Proof. Let $\mathbf{r}=\left\{r_{1}, \ldots, r_{n}\right\}$ and $\left\{\varepsilon_{1}, \ldots \varepsilon_{n-1}\right\}$ as above. In that proof, we denote, for convenience, $b=r_{1} / 2$; hence, a dependency with respect to $b$ is actually a dependency with respect to $\mathbf{r}$. Let
us consider $\Gamma$ a $\mathcal{T}(b)$-ball of radius $r_{n}$ and recall the notation (41). The ball $\Gamma$ is ( $\mathbf{r}, \varepsilon$ )-small iff ((42)) holds. But, for all $\sigma \in \Gamma, k \in \llbracket 1, n-1 \rrbracket$,

$$
\Gamma\left[r_{k}\right]=\Gamma\left(\sigma, r_{k}\right) \supset \Gamma\left(\sigma, r_{k}\right) \backslash \Gamma\left(\sigma, r_{k+1}\right)
$$

Thus, if $\Gamma$ is $(\mathbf{r}, \boldsymbol{\varepsilon})$-small, then all the vertices in $\Gamma$ belong to the set

$$
\begin{equation*}
\mathscr{S}(\mathbf{r}, \boldsymbol{\varepsilon}):=\left\{\sigma \in \mathcal{T}(b): \forall k \in\{1 \ldots n-1\} \quad \ell^{b}\left(\Gamma\left(\sigma, r_{k}\right) \backslash \Gamma\left(\sigma, r_{k+1}\right)\right) \leqslant \varepsilon_{k}\right\} \tag{50}
\end{equation*}
$$

The last set is easy to handle using the spinal decomposition. Indeed, according to (38) and the independence stated in Lemma 6, one has

$$
\begin{equation*}
\nu(\mathbf{r}, \varepsilon):=\mathbf{N}\left[\int \ell^{b}(\mathrm{~d} \sigma) \mathbf{1}_{\{\sigma \in \mathscr{S}(\mathbf{r}, \varepsilon)\}}\right]=\prod_{k=1}^{n-1} \mathbf{P}\left(\Lambda_{r_{k+1}, r_{k}}^{*} \leqslant \varepsilon_{k}\right) \tag{51}
\end{equation*}
$$

To rely $\mu(\mathbf{r}, \boldsymbol{\varepsilon})$ and $\nu(\mathbf{r}, \boldsymbol{\varepsilon})$, one can write

$$
\begin{equation*}
\mathbf{1}_{\{\Gamma \text { is }(\mathbf{r}, \boldsymbol{\varepsilon})-\text { small }\}} \leqslant \mathbf{1}_{\left\{\ell^{b}(\Gamma) \leqslant r_{n} \sqrt{\nu(\mathbf{r}, \boldsymbol{\varepsilon})}\right\}}+\frac{\ell^{b}(\Gamma)}{r_{n} \sqrt{\nu(\mathbf{r}, \boldsymbol{\varepsilon})}} \mathbf{1}_{\{\Gamma \text { is }(\mathbf{r}, \boldsymbol{\varepsilon})-\text { small }\}} \tag{52}
\end{equation*}
$$

Moreover, (50) entails that $\ell^{b}(\Gamma) \mathbf{1}_{\{\Gamma \text { is }(\mathbf{r}, \boldsymbol{\varepsilon})-\text { small }\}} \leqslant \int_{\Gamma} \ell^{b}(\mathrm{~d} \sigma) \mathbf{1}_{\{\sigma \in \mathscr{S}(\mathbf{r}, \boldsymbol{\varepsilon})\}}$. Recall now from Proposition $1(i)$ that the balls of the set $\mathscr{B}_{b, r_{n}}$ are pairwise disjoint. Summing in (52) over this set entails

$$
\begin{equation*}
S_{b, \mathbf{r}, \varepsilon} \leqslant \sum_{\Gamma \in \mathscr{B}_{b, r_{n}}} \mathbf{1}_{\left\{\ell^{b}(\Gamma) \leqslant r_{n} \sqrt{\nu(\mathbf{r}, \boldsymbol{\varepsilon})}\right\}}+\frac{\int \ell^{b}(\mathrm{~d} \sigma) \mathbf{1}_{\{\sigma \in \mathscr{S}(\mathbf{r}, \varepsilon)\}}}{r_{n} \sqrt{\nu(\mathbf{r}, \varepsilon)}} \tag{53}
\end{equation*}
$$

Now, recalling Lemma 4, we compute

$$
\begin{aligned}
\mathbf{N}_{b}\left[\sum_{\Gamma \in \mathscr{B}_{b, r_{n}}} \mathbf{1}_{\left\{\ell^{b}(\Gamma) \leqslant r_{n} \sqrt{\nu(\mathbf{r}, \boldsymbol{\varepsilon})}\right\}}\right] & =\mathbf{N}_{b}\left[Z_{b, r_{n}}\right]\left(1-\exp \left(-\left(2 / r_{n}\right) r_{n} \sqrt{\nu(\mathbf{r}, \boldsymbol{\varepsilon})}\right)\right) \\
& \leqslant \frac{r_{1}}{r_{n}} 2 \sqrt{\nu(\mathbf{r}, \boldsymbol{\varepsilon})}
\end{aligned}
$$

so the $\mathbf{N}$-measure of the first term in (53) is smaller than $\frac{2 b^{-1} r_{1}}{r_{n}} \sqrt{\nu(\mathbf{r}, \boldsymbol{\varepsilon})}$. Recalling that $b=r_{1} / 2$, we get that the latter equals $\frac{4}{r_{n}} \sqrt{\nu(\mathbf{r}, \varepsilon)}$. Moreover, by the mere definition (51), the $\mathbf{N}$-measure of the second term in (53) equals $\frac{1}{r_{n}} \sqrt{\nu(\mathbf{r}, \varepsilon)}$, so the first inequality is checked.

## 4 Proof of Theorem 1.

The proof of Theorem 1 will combine the following two theorems.
Theorem 2 Let $\kappa \in\left(\frac{1}{2}, \infty\right)$ and $m \in\left(0, \frac{1}{2}\right)$. Then, there exists a Borel subset $\mathbf{V}=\mathbf{V}(\kappa, m) \subset \mathbf{C}^{0}$ such that $\mathbf{N}\left(\mathbf{C}^{0} \backslash \mathbf{V}\right)=0$ and such that
on $\mathbf{V}, \quad$ for all Borel subset $\mathcal{A} \subset \mathcal{T}, \quad \forall a \in\left[m, m^{-1}\right], \quad \ell^{a}(\mathcal{A}) \leqslant \kappa \mathscr{H}_{g}(\mathcal{A} \cap \mathcal{T}(a))$.

For all $a, \alpha \in(0, \infty)$, let us set

$$
\begin{equation*}
\Delta_{a}^{\alpha}:=\left\{\sigma \in \mathcal{T}(a): \limsup _{r \rightarrow 0} \frac{\ell^{a}(B(\sigma, r))}{g(r)}<\alpha\right\} . \tag{54}
\end{equation*}
$$

Theorem 3 Let $\alpha \in\left(0, \frac{1}{2}\right)$ and $m \in(0,1 / 2)$. Then, there exists a Borel subset $\mathbf{V}^{\prime}=\mathbf{V}^{\prime}(\alpha, m) \subset$ $\mathbf{C}^{0}$ such that $\mathbf{N}\left(\mathbf{C}^{0} \backslash \mathbf{V}^{\prime}\right)=0$ and such that

$$
\text { on } \mathbf{V}^{\prime}, \quad \forall a \in\left[m, m^{-1}\right], \quad \mathscr{H}_{g}\left(\Delta_{a}^{\alpha}\right)=0 .
$$

The proofs of Theorem 2 and 3 share a common strategy, taken from Perkins [14, 15]. We need to control the mass, or the number of "bad" $\mathcal{T}(a)$-balls where "bad" means too large or too small. And we want to do it uniformly for all levels $a$. This problem will be linked with a discrete one using a finite grid, and the measure or the number of bad $\mathcal{T}(a)$-balls will be compared with a convenient multiple of $\ell^{a}(\mathcal{T})$, the total mass at level $a$.

### 4.1 Proof of Theorem 2.

### 4.1.1 Large balls.

Let us fix a level $a \in(0, \infty)$, and recall from section 3.1 the definition of the sigma-field $\mathcal{G}_{a}$, generated by the excursion below level $a$. We also recall the definition of $\mathcal{T}(a)$-balls (18). We fix a threshold $y \in(0, \infty)$ and we consider the following set of "large" points on $\mathcal{T}(a)$ :

$$
\begin{equation*}
\mathcal{L}_{a, r, y}=\left\{\sigma \in \mathcal{T}(a): \ell^{a}(\Gamma(\sigma, r))>y\right\} . \tag{55}
\end{equation*}
$$

According to Lemma 4, the "total large mass" $\ell^{a}\left(\mathcal{L}_{a, r, y}\right)=\sum_{\Gamma \in \mathscr{B}_{a, r}} \ell^{a}(\Gamma) \mathbf{1}_{\left\{\ell^{a}(\Gamma)>y\right\}}$ is $\mathcal{G}_{a^{-}}$ measurable.

Lemma 12 For all $a, l, y, r, \delta \in(0, \infty)$, for all $c \in(1, \infty)$,

$$
\mathbf{N}\left(\ell^{a}\left(\mathcal{L}_{a, r, y / c}\right) \leqslant l ; \sup _{b \in[a, a+\delta]} \ell^{b}\left(\mathcal{L}_{b, r, y}\right)>4 l\right) \leqslant \frac{1}{a} \exp (-l / \delta)+\frac{2}{r} \exp \left(-\left(1-c^{-1 / 2}\right)^{2} y / \delta\right) .
$$

Proof. For all $a \in(0, \infty)$, and all $c \in(1, \infty)$, we define $A_{0}$, a Borel subset of $\mathbf{C}^{0}$, as the event

$$
\begin{equation*}
A_{0}=\left\{\ell^{a}\left(\mathcal{L}_{a, r, y / c}\right) \leqslant l ; \sup _{b \in[a, a+\delta]} \ell^{b}\left(\mathcal{L}_{b, r, y}\right)>4 l\right\} . \tag{56}
\end{equation*}
$$

We recall from Proposition 1 that for $r \in(0, \infty), \mathscr{B}_{a, r}=\left\{\Gamma_{i}, 1 \leqslant i \leqslant Z_{a, r}\right\}$ is the collection of $\mathcal{T}(a)$-balls of radius $r$ at level $a$. For $\Gamma$ a $\mathcal{T}(a)$-ball and $b \in[a, \infty)$, we defined $\Gamma^{b}=\{\sigma \in \mathcal{T}(b)$ : $\left.\exists \sigma^{\prime} \in \Gamma, \sigma^{\prime} \in \llbracket \rho, \sigma \rrbracket\right\}$ as the set of vertices at level $b$ having an ancestor in $\Gamma$ (see (29) for details). Next we define $A_{1}$ a Borel subset of $\mathbf{C}^{0}$ as the event

$$
\begin{equation*}
A_{1}:=\left\{\exists i \in\left\{1, \ldots, Z_{a, r}\right\}, \quad \ell^{a}\left(\Gamma_{i}\right) \leqslant y / c \text { and } \sup _{b \in[a, a+\delta]} \ell^{b}\left(\Gamma_{i}^{b}\right)>y\right\}, \tag{57}
\end{equation*}
$$

and set

$$
\begin{equation*}
\mathcal{L}_{a, r, y / c}^{b}:=\bigcup_{\substack{i \in\left\{1, \ldots, Z_{a, r}\right\} \\ \ell^{a}\left(\Gamma_{i}\right) \geqslant y / c}} \Gamma_{i}^{b} \subset \mathcal{T}(b), \tag{58}
\end{equation*}
$$

which is the set of all vertices at level $b$ having an "large" ancestor at level $a$. We prove the following :

$$
\begin{equation*}
\text { on } \mathbf{C}^{0} \backslash A_{1}, \quad \forall b \in[a, a+\delta] \quad \mathcal{L}_{b, r, y} \subset \mathcal{L}_{a, r, y / c}^{b} . \tag{59}
\end{equation*}
$$

Proof of (59). Let $b \geqslant a \geqslant r / 2>0$ and let $\sigma \in \mathcal{L}_{b, r, y}$. Thus, the ball $\Gamma:=\Gamma(\sigma, r) \in \mathscr{B}_{b, r}$ is such that $\ell^{b}(\Gamma) \geqslant y$. Let $\sigma_{a}$ the unique ancestor of $\sigma$ at level $a$, namely $d\left(\rho, \sigma_{a}\right)=a$ and $\sigma_{a} \in \llbracket \rho, \sigma \rrbracket$. We set $\Upsilon:=\Gamma\left(\sigma_{a}, r\right) \in \mathscr{B}_{a, r}$ and we first claim that $\Gamma \subset \Upsilon^{b}$. Indeed, let $\sigma^{\prime} \in \Gamma$ (so $d\left(\sigma, \sigma^{\prime}\right)<r$ ) and let $\sigma_{a}^{\prime}$ the unique ancestor of $\sigma^{\prime}$ at level $a$. Recalling that $\sigma \wedge \sigma^{\prime}$ stands for the most recent common ancestor of $\sigma$ and $\sigma^{\prime}$, one get $d\left(\rho, \sigma \wedge \sigma^{\prime}\right)=\frac{1}{2}\left(2 b-d\left(\sigma, \sigma^{\prime}\right)\right)$. Then two cases may occur. First, if $d\left(\sigma, \sigma^{\prime}\right) \leqslant 2(b-a)$, then $d\left(\rho, \sigma \wedge \sigma^{\prime}\right) \geqslant a$, thus $\sigma_{a}=\sigma_{a}^{\prime}$ and $\sigma^{\prime} \in \Upsilon^{b}$. If $d\left(\sigma, \sigma^{\prime}\right) \in(2(b-a), r)$. Then, one has $d\left(\rho, \sigma \wedge \sigma^{\prime}\right)<a$. We deduce from that inequality that $\sigma_{a} \neq \sigma_{a}^{\prime}$ and that $\sigma_{a} \wedge \sigma_{a}^{\prime}=\sigma \wedge \sigma^{\prime}$. Hence,

$$
\begin{aligned}
d\left(\sigma_{a}, \sigma_{a}^{\prime}\right) & =2 a-2 d\left(\rho, \sigma_{a} \wedge \sigma_{a}^{\prime}\right) \\
& =2 b-2 d\left(\rho, \sigma \wedge \sigma^{\prime}\right)+2 a-2 b \\
& =d\left(\sigma, \sigma^{\prime}\right)-2(b-a)<r
\end{aligned}
$$

Thus $\sigma_{a}^{\prime} \in \Upsilon$ and $\sigma^{\prime} \in \Upsilon^{b}$, which ends the proof of the inclusion $\Gamma \subset \Upsilon^{b}$. We then get $\ell^{b}\left(\Upsilon^{b}\right) \geqslant \ell^{b}(\Gamma)>y$. On $\mathbf{C}^{0} \backslash A_{1}$, one cannot have both $\ell^{a}(\Upsilon) \leqslant y / c$ and $\ell^{b}\left(\Upsilon^{b}\right)>y$, which entails that here, $\ell^{a}(\Upsilon)>y / c$.
To sum up, on $\mathbf{C}^{0} \backslash A_{1}$, a vertex $\sigma$, taken in $\mathcal{L}_{b, r, y}$, has an ancestor in a ball $\Upsilon$, such that $\ell^{a}(\Upsilon) \geqslant y / c$. Thus, this ancestor belongs to $\mathcal{L}_{a, r, y / c}$ and $\sigma \in \mathcal{L}_{a, r, y / c}^{b}$.
End of the proof of (59).
Let us finish the proof of the lemma. From (59), we see that

$$
\begin{equation*}
\mathbf{N}\left(A_{0}\right) \leqslant \mathbf{N}\left(A_{1}\right)+\mathbf{N}\left(A_{0} \cap\left(\mathbf{C}^{0} \backslash A_{1}\right)\right) \leqslant \mathbf{N}\left(A_{1}\right)+\mathbf{N}\left(A_{2}\right) \tag{60}
\end{equation*}
$$

where $A_{2}$ is defined by

$$
\begin{equation*}
A_{2}:=\left\{\ell^{a}\left(\mathcal{L}_{a, r, y / c}\right) \leqslant l ; \sup _{b \in[a, a+\delta]} \ell^{b}\left(\mathcal{L}_{a, r, y / c}^{b}\right) \geqslant 4 l\right\} . \tag{61}
\end{equation*}
$$

We control $\mathbf{N}\left(A_{1}\right)$ and $\mathbf{N}\left(A_{2}\right)$ thanks to Lemma 8 (ii). Indeed, Lemma 5 states that under $\mathbf{N}_{a}$, conditionally on $\mathcal{G}_{a}$, the processes $\left(\ell^{a+a^{\prime}}\left(\Gamma_{i}^{a+a^{\prime}}\right), a^{\prime} \geqslant 0\right), 1 \leqslant i \leqslant Z_{a, r}$ are independent Feller diffusions started at $\ell^{a}\left(\Gamma_{i}\right), 1 \leqslant i \leqslant Z_{a, r}$. Since $\ell^{b}\left(\mathcal{L}_{a, r, y / c}^{b}\right)=\sum_{i=1}^{Z_{a, r}} \mathbf{1}_{\left\{\ell^{a}\left(\Gamma_{i}\right)>y / c\right\}} \ell^{b}\left(\Gamma_{i}^{b}\right)$, it implies that $\left(\ell^{a+a^{\prime}}\left(\mathcal{L}_{a, r, y / c}^{a+a^{\prime}}\right), a^{\prime} \geqslant 0\right)$ is a Feller diffusion started at $\ell^{a}\left(\mathcal{L}_{a, r, y / c}\right)$. Thus, on the one hand, sub-additivity and Lemma 8 (ii) entails

$$
\begin{align*}
\mathbf{N}\left(A_{1}\right) & \leqslant \frac{1}{a} \mathbf{N}_{a}\left[\sum_{i=1}^{Z_{a, r}} \mathbf{1}_{\left\{\ell^{a}\left(\Gamma_{i}\right) \leqslant y / c\right\}} \exp \left(-\delta^{-1}\left(\sqrt{y}-\sqrt{\ell^{a}\left(\Gamma_{i}\right)}\right)^{2}\right)\right] \\
& \leqslant \frac{1}{a} \exp \left(-\left(1-c^{-1 / 2}\right)^{2} \delta^{-1} y\right) \mathbf{N}_{a}\left[Z_{a, r}\right]=\frac{2}{r} \exp \left(-\left(1-c^{-1 / 2}\right)^{2} \delta^{-1} y\right) . \tag{62}
\end{align*}
$$

On the other hand,

$$
\begin{align*}
\mathbf{N}\left(A_{2}\right) & \leqslant \frac{1}{a} \mathbf{N}_{a}\left[\mathbf{1}_{\left\{\ell^{a}\left(\mathcal{L}_{a, r, y / c}\right) \leqslant l\right\}} \exp \left(-\delta^{-1}\left(2 \sqrt{l}-\sqrt{\ell^{a}\left(\mathcal{L}_{a, r, y / c}\right)}\right)^{2}\right)\right] \\
& \leqslant \frac{1}{a} \exp (-l / \delta) . \tag{63}
\end{align*}
$$

Hence, the desired result follows from (60), (62), and (63).
Recall that $g(r)=r \log \log (1 / r)$. We fix $\kappa \in\left(\frac{1}{2}, \infty\right)$, and we shall apply the previous lemma with $y=\kappa g(r)$. The next lemma allows to control $\ell^{a}\left(\mathcal{L}_{a, r, \kappa g(r)}\right)$ uniformly for all levels $a$. Its proof involves a discrete grid : for $m<1 / 2$ and $r \in(0, \infty)$, we set

$$
\begin{equation*}
G(r, m):=\left\{m+k \delta_{r}, k \in \mathbb{N}^{*}\right\} \cap\left[m, m^{-1}\right], \tag{64}
\end{equation*}
$$

where $\delta_{r}$ is the mesh of the grid, defined by

$$
\begin{equation*}
\delta_{r}=r^{3 / 2} . \tag{65}
\end{equation*}
$$

Note that $G(r, m)$ contains less than $\left(m \delta_{r}\right)^{-1}$ points.
Lemma 13 Let $m \in(0,1 / 2)$. Let $\kappa \in\left(\frac{1}{2}, \infty\right)$ and $\beta \in(1, \infty)$ such that $2 \kappa-\beta>0$. There exists $a$ constant $r_{1} \in(0, \infty)$ only depending on $\kappa, \beta, m$, such that

$$
\begin{equation*}
\forall r \in\left(0, r_{1}\right), \quad \mathbf{N}\left(\sup _{b \in\left[m, m^{-1}\right]} \ell^{b}\left(\mathcal{L}_{b, r, \kappa g(r)}\right)>4 \log (1 / r)^{-\beta}\right) \leqslant \log (1 / r)^{-2} . \tag{66}
\end{equation*}
$$

Proof. In what follows, we denote $T_{0}$ the left-hand-side of (66). Let us consider $c \in(1, \infty)$ such that $2 \kappa / c-\beta>0$. Recall that $G(r, m)$ stands for the grid defined by (64). Then we have $T_{0} \leqslant T_{1}+T_{2}$, where :

$$
\begin{aligned}
& T_{1}=\mathbf{N}\left(\sup _{a \in G(r, m)} \ell^{a}\left(\mathcal{L}_{a, r, \kappa k g(r) / c}\right) \leqslant \log (1 / r)^{-\beta} ; \sup _{b \in\left[m, m^{-1}\right]} \ell^{b}\left(\mathcal{L}_{b, r, \kappa g(r)}\right) \geqslant 4 \log (1 / r)^{-\beta}\right) \\
& T_{2}=\mathbf{N}\left(\sup _{a \in G(r, m)} \ell^{a}\left(\mathcal{L}_{a, r, \kappa g(r) / c}\right)>\log (1 / r)^{-\beta}\right) .
\end{aligned}
$$

Using sub-additivity and Lemma 12, one get

$$
\begin{aligned}
T_{1} & =\mathbf{N}\left(\bigcup_{a \in G(r, m)}\left\{\ell^{a}\left(\mathcal{L}_{a, r, \kappa g(r) / c}\right) \leqslant \log (1 / r)^{-\beta} ; \sup _{b \in\left[a, a+\delta_{r}\right]} \ell^{b}\left(\mathcal{L}_{b, r, \kappa g(r)}\right) \geqslant 4 \log (1 / r)^{-\beta}\right\}\right) \\
& \leqslant\left(m \delta_{r}\right)^{-1} \sup _{a \in G(r, m)} \mathbf{N}\left(\ell^{a}\left(\mathcal{L}_{a, r, \kappa g(r) / c}\right) \leqslant \log (1 / r)^{-\beta} ; \sup _{b \in\left[a, a+\delta_{r}\right]} \ell^{b}\left(\mathcal{L}_{b, r, \kappa g(r)}\right) \geqslant 4 \log (1 / r)^{-\beta}\right) \\
& \leqslant\left(m \delta_{r}\right)^{-1}\left(m^{-1} \exp \left(-\delta_{r}^{-1} \log (1 / r)^{-\beta}\right)+\frac{2}{r} \exp \left(-\left(1-c^{-1 / 2}\right)^{2} \kappa \delta_{r}^{-1} g(r)\right)\right)
\end{aligned}
$$

One has $\delta_{r}^{-1} \log (1 / r)^{-\beta} \geqslant r^{-1}$ and $\delta_{r}^{-1} g(r) \geqslant r^{-1 / 2}$ for all $r$ sufficiently small. Thus, for example, $T_{1} \leqslant \exp \left(-r^{-1 / 4}\right) \leqslant(1 / 2) \log (1 / r)^{-2}$ for all $r$ sufficiently small.

Let us bound $T_{2}$. To that end, we set

$$
\begin{equation*}
\lambda(r, \kappa, c):=(2 / r) \mathbf{E}\left[\mathcal{E} \mathbf{1}_{\{\mathcal{E}>\kappa g(r) / c\}}\right] \tag{67}
\end{equation*}
$$

where $\mathcal{E}$ is a r.v. on $(\Omega, \mathcal{F}, \mathbf{P})$ exponentially distributed with mean $r / 2$. For fixed $\kappa$ and $c$, elementary computations entail

$$
\begin{align*}
\lambda(r, \kappa, c) & =(2 / r) \mathbf{P}(\mathcal{E}>\kappa g(r) / c) \mathbf{E}[\kappa g(r) / c+\mathcal{E}] \\
& =(2 / r) \exp (-2(\kappa / c) \log \log 1 / r)((\kappa / c) r \log \log 1 / r+r / 2) \\
& \underset{r \rightarrow 0}{\sim}(2 \kappa / c) \log (1 / r)^{-2 \kappa / c} \log \log 1 / r . \tag{68}
\end{align*}
$$

We set $T_{2} \leqslant T_{3}+T_{4}$, where

$$
\begin{aligned}
& T_{3}=\mathbf{N}\left(\sup _{a \in G(r, m)}\left|\ell^{a}\left(\mathcal{L}_{a, r, \kappa g(r) / c}\right)-\lambda(r, \kappa, c) \ell^{a}(\mathcal{T})\right|>\frac{1}{2} \log (1 / r)^{-\beta}\right) \\
& T_{4}=\mathbf{N}\left(\sup _{a \in G(r, m)} \lambda(r, \kappa, c) \ell^{a}(\mathcal{T})>\frac{1}{2} \log (1 / r)^{-\beta}\right)
\end{aligned}
$$

By sub-additivity and a Markov inequality involving a moment of order 4, we get

$$
\begin{align*}
T_{3} & \leqslant\left(m \delta_{r}\right)^{-1} \sup _{a \in G(r, m)} \mathbf{N}\left(\left|\ell^{a}\left(\mathcal{L}_{a, r, \kappa g(r) / c}\right)-\lambda(r, \kappa, c) \ell^{a}(\mathcal{T})\right|>\frac{1}{2} \log (1 / r)^{-\beta}\right) \\
& \leqslant\left(m \delta_{r}\right)^{-1} 2^{4} \log (1 / r)^{4 \beta} \sup _{a \in G(r, m)} a^{-1} \mathbf{N}_{a}\left[\left(\ell^{a}\left(\mathcal{L}_{a, r, \kappa g(r) / c}\right)-\lambda(r, \kappa, c) \ell^{a}(\mathcal{T})\right)^{4}\right] \tag{69}
\end{align*}
$$

Recall notation $\mathscr{B}_{a, r}=\left\{\Gamma_{i}, 1 \leqslant i \leqslant Z_{a, r}\right\}$ for the set of $\mathcal{T}(a)$-balls with radius $r$. Then, consider the decomposition

$$
\ell^{a}\left(\mathcal{L}_{a, r, \kappa g(r) / c}\right)-\lambda(r, \kappa, c) \ell^{a}(\mathcal{T})=\sum_{i=1}^{Z_{a, r}} X_{i}
$$

where $X_{i}:=\ell^{a}\left(\Gamma_{i}\right)\left(\mathbf{1}_{\left\{\ell^{a}\left(\Gamma_{i}\right) \geqslant \kappa g(r) / c\right\}}-\lambda(r, \kappa, c)\right)$. Using Lemma 4, we see that under $\mathbf{N}_{a}$, conditionally on $Z_{a, r}$, the random variables $\ell^{a}\left(\Gamma_{1}\right), \ldots \ell^{a}\left(\Gamma_{Z_{a, r}}\right)$ are independent and exponentially distributed with mean $r / 2$. Thus, the definition (67) of $\lambda(r, \kappa, c)$ entails that under $\mathbf{N}_{a}$, conditionally on $Z_{a, r}$, the r.v. $X_{1}, \ldots X_{Z_{a, r}}$ are i.i.d., with mean 0 and a moment of order 4 . Then, by Lemma 7,

$$
\begin{equation*}
\mathbf{N}_{a}\left[\left(\sum_{i=1}^{Z_{a, r}} X_{i}\right)^{4}\right] \leqslant 3 \mathbf{N}_{a}\left(X_{1}^{4}\right) \mathbf{N}_{a}\left[Z_{a, r}^{2}\right] \tag{70}
\end{equation*}
$$

From (68), we know that $\lambda(r, \kappa, c) \xrightarrow{r \rightarrow 0} 0$, so for all sufficiently small $r, \lambda(r, \kappa, c) \leqslant 1 / 2$ and $\left|X_{1}^{4}\right| \leqslant \ell^{a}\left(\Gamma_{1}\right)$, which implies $\mathbf{N}_{a}\left[X_{1}\right] \leqslant \mathbf{N}_{a}\left[\ell^{a}\left(\Gamma_{1}\right)^{4}\right]=\frac{3}{2} r^{4}$ for all sufficiently small $r$. Moreover, $Z_{a, r}$ is under $\mathbf{N}_{a}$ a geometric r.v. with "success" probability $p=r / 2 a$ (see Lemma 4), thus $\mathbf{N}_{a}\left[Z_{a, r}^{2}\right]=(2-p) / p^{2} \leqslant 8 a^{2} / r^{2}$. Combining (69) and (70), we get, for all sufficiently small $r$,

$$
\begin{equation*}
T_{3} \leqslant 3.2^{4} \cdot\left(m \delta_{r}\right)^{-1} \log (1 / r)^{4 \beta} \sup _{a \in G(r, m)} a^{-1} \frac{3 r^{4}}{2} \frac{8 a^{2}}{r^{2}} \leqslant 10^{3} m^{-2} \log (1 / r)^{4 \beta} r^{1 / 2} \tag{71}
\end{equation*}
$$

recalling that $\delta_{r}=r^{3 / 2}$. Observe now that the right hand side is smaller than $(1 / 4) \log (1 / r)^{-2}$ for all sufficiently small $r$.

For the term $T_{4}$, Lemma 9 entails

$$
\begin{align*}
T_{4} & \leqslant \mathbf{N}\left(\sup _{b \in\left[m, m^{-1}\right]} \ell^{b}(\mathcal{T})>\frac{1}{2} \lambda(r, \kappa, c)^{-1} \log (1 / r)^{-\beta}\right) \\
& \leqslant(2 / m) \exp \left(-(m / 4) \lambda(r, \kappa, c)^{-1} \log (1 / r)^{-\beta}\right) \tag{72}
\end{align*}
$$

By (68),

$$
\lambda(r, \kappa, c)^{-1} \log (1 / r)^{-\beta} \underset{r \rightarrow 0}{\sim} \frac{c}{2 \kappa} \log (1 / r)^{2 \kappa / c-\beta} \log \log (1 / r)^{-1}
$$

Recall that $2 \kappa / c>0$ and take $\varepsilon \in(0,2 \kappa / c-\beta)$. Thus, for all sufficiently small $r$,

$$
T_{4} \leqslant(2 / m) \exp \left(-\log (1 / r)^{\varepsilon}\right)
$$

which is smaller than $(1 / 4) \log (1 / r)^{-2}$ for all sufficiently small $r$.

### 4.1.2 Proof of Theorem 2.

Let $\kappa \in(1 / 2, \infty)$, and let $m \in(0,1 / 2)$. Let $\beta \in(1, \infty)$ such that $2 \kappa-\beta>0$. For all $a \in(0, \infty)$, $y \in(1, \infty)$ recall from (54) the definition :

$$
\begin{equation*}
\Delta_{a}^{y \kappa}=\left\{\sigma \in \mathcal{T}(a): \limsup _{r \rightarrow 0} \frac{\ell^{a}(B(\sigma, r))}{g(r)}<y \kappa\right\} \tag{73}
\end{equation*}
$$

For any $p \in \mathbb{N}$, set $r_{p}:=y^{-p}$. By Lemma 13 , for all sufficiently large $p$,

$$
\begin{equation*}
\mathbf{N}\left(\sup _{a \in\left[m, m^{-1}\right]} \ell^{a}\left(\mathcal{L}_{a, r_{p}, \kappa g\left(r_{p}\right)}\right)>4 \log \left(1 / r_{p}\right)^{-\beta}\right) \leqslant \log \left(1 / r_{p}\right)^{-2}=\log (y)^{-2} p^{-2} \tag{74}
\end{equation*}
$$

whose sum over $p$ is finite. By Borel Cantelli lemma,

$$
\begin{equation*}
\mathbf{N} \text {-a.e., for all sufficiently large } p, \sup _{a \in\left[m, m^{-1}\right]} \ell^{a}\left(\mathcal{L}_{a, r_{p}, \kappa g\left(r_{p}\right)}\right) \leqslant 4 \log \left(1 / r_{p}\right)^{-\beta} \tag{75}
\end{equation*}
$$

Moreover, $\log \left(1 / r_{p}\right)^{-\beta}=\log (y)^{-\beta} p^{-\beta}$, and recall that $\beta>1$. Thus, (75) entails that there exists a Borel subset $\mathbf{V}_{y} \subset \mathbf{C}^{0}$, such that $\mathbf{N}\left(\mathbf{C}^{0} \backslash \mathbf{V}_{y}\right)=0$, and on $\mathbf{V}_{y}$ :

$$
\forall a \in\left[m, m^{-1}\right], \quad \sum_{p=1}^{\infty} \ell^{a}\left(\mathcal{L}_{a, r_{p}, \kappa g\left(r_{p}\right)}\right)=\sum_{p=1}^{\infty} \ell^{a}\left(\left\{\sigma: \ell^{a}\left(B\left(\sigma, r_{p}\right)\right)>\kappa g\left(r_{p}\right)\right\}\right)<\infty
$$

We can apply again the Borel-Cantelli Lemma, to the finite measures $\ell^{a}$ to get that,

$$
\begin{equation*}
\text { on } \mathbf{V}_{y} \quad \forall a \in\left[m, m^{-1}\right], \quad \ell^{a}(\mathrm{~d} \sigma) \text {-a.e. } \quad \exists p_{0}(a, \sigma), \quad \forall p \geqslant p_{0}(a, \sigma), \quad \frac{\ell^{a}\left(B\left(\sigma, r_{p}\right)\right)}{g\left(r_{p}\right)} \leqslant \kappa . \tag{76}
\end{equation*}
$$

If $u \in\left(r_{p+1}, r_{p}\right]$, one has $\frac{\ell^{a}(B(\sigma, u))}{g(u)}<\frac{\ell^{a}\left(B\left(\sigma, r_{p}\right)\right)}{g\left(r_{p+1}\right)} \leqslant y \frac{\ell^{a}\left(B\left(\sigma, r_{p}\right)\right)}{g\left(r_{p}\right)}$. Combined with (76), this entails that on $\mathbf{V}_{y}$, for all $a$ in $\left[m, m^{-1}\right]$, for $\ell^{a}$-almost every $\sigma$ in $\mathcal{T}(a), \lim \sup _{r \rightarrow 0} \ell^{a}(B(\sigma, r)) / g(r)<y \kappa$. This can be rewritten in

$$
\begin{equation*}
\text { on } \mathbf{V}_{y}, \quad \forall a \in\left[m, m^{-1}\right], \quad \ell^{a}\left(\mathcal{T}(a) \backslash \Delta_{a}^{y \kappa}\right)=0 . \tag{77}
\end{equation*}
$$

Now set $\mathbf{V}=\bigcap\left\{\mathbf{V}_{y} ; y>1 ; y \in \mathbb{Q}\right\}$. Clearly, $\mathbf{N}\left(\mathbf{C}^{0} \backslash \mathbf{V}\right)=0$ and by monotonicity, for all $\kappa^{\prime} \in(\kappa, \infty), \mathcal{T}(a) \backslash \Delta_{a}^{\kappa^{\prime}} \subset \bigcup_{y>1 ; y \in \mathbb{Q}}\left\{\mathcal{T}(a) \backslash \Delta_{a}^{y \kappa}\right\}$. It follows easily from (77) that

$$
\begin{equation*}
\text { on } \mathbf{V}, \quad \forall a \in\left[m, m^{-1}\right], \quad \forall \kappa^{\prime} \in(\kappa, \infty) \quad \ell^{a}\left(\mathcal{T}(a) \backslash \Delta_{a}^{\kappa^{\prime}}\right)=0 \tag{78}
\end{equation*}
$$

Thus, using Lemma 2, we get :

$$
\begin{aligned}
& \text { on } \mathbf{V} \quad \forall \mathcal{A} \text { Borel subset of } \mathcal{T} \quad \forall a \in\left[m, m^{-1}\right] \quad \forall \kappa^{\prime} \in(\kappa, \infty) \\
& \quad \ell^{a}(\mathcal{A})=\ell^{a}\left(\mathcal{A} \cap \Delta_{a}^{\kappa^{\prime}}\right) \leqslant \kappa^{\prime} \mathscr{H}_{g}\left(\mathcal{A} \cap \Delta_{a}^{\kappa^{\prime}}\right) \leqslant \kappa^{\prime} \mathscr{H}_{g}(\mathcal{A} \cap \mathcal{T}(a)) .
\end{aligned}
$$

This ends the proof of Theorem 2 letting $\kappa^{\prime} \searrow \kappa$.

### 4.2 Proof of Theorem 3.

### 4.2.1 Small balls.

For given level $a \in(0, \infty)$ and $r \in(0, \infty)$ we recall the notation $\mathscr{B}_{a, r}$ for the set of $\mathcal{T}(a)$-balls of radius $r$. We recall from (41) that for $r \geqslant r^{\prime}>0$, a ball $\Gamma \in \mathscr{B}_{a, r^{\prime}}$ is contained in a unique ball in $\mathscr{B}_{a, r}$, denoted $\Gamma[r]$. Let $\mathbf{r}=\left\{r_{1}>\ldots>r_{n}\right\}$ and $\boldsymbol{\varepsilon}=\left\{\varepsilon_{1}>\ldots>\varepsilon_{n-1}\right\}$. Recall from (42) that $\Gamma$, a $\mathcal{T}(a)$-ball of radius $r_{n}$ is $(\mathbf{r}, \boldsymbol{\varepsilon})$-small iff

$$
\forall k \in \llbracket 1, n-1 \rrbracket \quad \ell^{a}\left(\Gamma\left[r_{k}\right]\right) \leqslant \varepsilon_{k} .
$$

The total number of $(\mathbf{r}, \boldsymbol{\varepsilon})$-small balls at level $a$ is denoted by $S_{a, \mathbf{r}, \boldsymbol{\varepsilon}}$ (see (43)). For $u \in(0, \infty)$, we write $u \mathbf{r}$ for the set $\left\{u r_{1}>\ldots>u r_{n}\right\}$. We recall from (28) the following notation : if $\Gamma$ is a $\mathcal{T}(a)$-ball, then, for all $b \geqslant a, \Gamma^{b}$ is the subset of all the vertices in $\mathcal{T}(b)$ having an ancestor in $\Gamma$. Namely, $\Gamma^{b}=\left\{\sigma \in \mathcal{T}(b), \exists \sigma^{\prime} \in \Gamma: \sigma^{\prime} \in \llbracket \rho, \sigma \rrbracket\right\}$.

Lemma 14 Let $a, \delta \in(0, \infty)$, and $n \geqslant 2$. Let $\boldsymbol{r}=\left\{r_{1}>\ldots>r_{n}\right\}$ and $\varepsilon=\left\{\varepsilon_{1}>\ldots>\varepsilon_{n-1}\right\}$ and $c \in(1, \infty)$. Let $\alpha \in(0,1 / 2)$ and $\tilde{\alpha} \in(\alpha, 1 / 2)$. If $\delta<\frac{c-1}{2 c} r_{n}$, then

$$
\mathbf{N}\left(\sup _{b \in[a, a+\delta]} S_{b, \boldsymbol{r}, \alpha \varepsilon}>S_{a, c^{-1} r, \tilde{\alpha} \varepsilon}\right) \leqslant \frac{2 n}{r_{n}} \exp \left(-(\sqrt{\tilde{\alpha}}-\sqrt{\alpha})^{2} \varepsilon_{n-1} / \delta\right) .
$$

Proof. Let us denote $B_{0}=\left\{\sup _{b \in[a, a+\delta]} S_{b, \mathbf{r}, \alpha \varepsilon}>S_{a, c^{-1} \mathbf{r}, \tilde{\alpha} \varepsilon}\right\}$. Next, we define the event $B_{1}$ by

$$
\begin{equation*}
B_{1}=\left\{\exists k \in\{1 \ldots, n-1\}, \exists \Gamma \in \mathscr{B}_{a, r_{k} / c}: \ell^{a}(\Gamma) \geqslant \tilde{\alpha} \varepsilon_{k} \text { and } \inf _{b \in[a, a+\delta]} \ell^{b}\left(\Gamma^{b}\right)<\alpha \varepsilon_{k}\right\} . \tag{79}
\end{equation*}
$$

We will prove that $B_{0} \subset B_{1}$, that is to say

$$
\begin{equation*}
\text { on } \mathbf{C}^{0} \backslash B_{1}, \sup _{b \in[a, a+\delta]} S_{b, \mathbf{r}, \alpha \varepsilon} \leqslant S_{a, c^{-1} \mathbf{r}, \tilde{\alpha} \varepsilon} . \tag{80}
\end{equation*}
$$

Proof of (80). We work deterministically on $\mathbf{C}^{0} \backslash B_{1}$. The inequality (80) follows from the following claim.

For every $b \in[a, a+\delta]$, for every $\Gamma$ a $\mathcal{T}(b)$-ball of radius $r_{n}$ which is $(\boldsymbol{r}, \alpha \boldsymbol{\varepsilon})$-small, there exists $\Upsilon$ a $\mathcal{T}(a)$-ball of radius $r_{n} / c$ such that $\Upsilon$ is $\left(c^{-1} \boldsymbol{r}, \tilde{\alpha} \varepsilon\right)$-small and $\Upsilon^{b} \subset \Gamma$.

Assume that the latter is true. Then, to any ( $\mathbf{r}, \alpha \varepsilon$ )-small ball at level $b$ corresponds a $\left(c^{-1} \mathbf{r}, \tilde{\alpha} \varepsilon\right)$ small ball at level $a$ and the correspondence is injective. Summing over all $\mathcal{T}(b)$-ball, we obtain (80).
Now let $b \in[a, a+\delta]$ and $\Gamma \in \mathscr{B}_{b, r_{n}}$ such that $\Gamma$ is $(\mathbf{r}, \alpha \boldsymbol{\varepsilon})$-small. Let $\sigma \in \Gamma$ and let $\sigma_{a}$ its unique ancestor at level $a$. Namely $\sigma_{a} \in \mathcal{T}(a)$ and $\sigma_{a} \in \llbracket \rho, \sigma \rrbracket$. We denote $\Upsilon=\Gamma\left(\sigma_{a}, r_{n} / c\right) \in \mathscr{B}_{a, r_{n} / c}$ the $\mathcal{T}(a)$-ball of radius $r_{n} / c$ that contains $\sigma_{a}$. We claim that $\Upsilon$ is $\left(c^{-1} \mathbf{r}, \tilde{\alpha} \varepsilon\right)$-small and that $\Upsilon^{b} \subset \Gamma$. To prove this, we show

$$
\begin{equation*}
\forall k \in\{1, \ldots, n\} \quad\left(\Upsilon\left[r_{k} / c\right]\right)^{b} \subset \Gamma\left[r_{k}\right] . \tag{81}
\end{equation*}
$$

Let $k \in\{1 \ldots n\}$ and let $\gamma \in\left(\Upsilon\left[r_{k} / c\right]\right)^{b}$. Its unique ancestor at level $a$, denoted $\gamma_{a}$, is such that $\gamma_{a} \in \Upsilon\left[r_{k} / c\right]$. Two cases may occur. First, if $d(\sigma, \gamma) \leqslant 2(b-a)$, then we have $2(b-a) \leqslant$ $2 \delta<\frac{c-1}{c} r_{n}<r_{n} \leqslant r_{k}$. The other case corresponds to $d(\sigma, \gamma)>2(b-a)$. Then $d(\rho, \sigma \wedge \gamma)=$ $\frac{1}{2}(2 b-d(\sigma, \gamma))<a$. Thus, $\sigma \wedge \gamma=\sigma_{a} \wedge \gamma_{a}$ and we have

$$
\begin{aligned}
d(\sigma, \gamma) & =2 b-2 d(\rho, \sigma \wedge \gamma) \\
& =2 a-2 d\left(\rho, \sigma_{a} \wedge \gamma_{a}\right)+2 b-2 a \\
& \leqslant d\left(\sigma_{a}, \gamma_{a}\right)+2 \delta \\
& <\frac{r_{k}}{c}+\frac{c-1}{c} r_{n} \leqslant r_{k},
\end{aligned}
$$

where we used in the last line that $\sigma_{a} \in \Gamma \subset \Gamma\left[r_{k} / c\right]$. In both cases, $d(\sigma, \gamma)<r_{k}$ so $\gamma \in$ $\Gamma\left(\sigma, r_{k}\right)=\Gamma\left[r_{k}\right]$, the last equality being a consequence of Proposition 1 (ii), and the definition of $\Gamma=\Gamma\left(\sigma, r_{n}\right)$. Thus, (81) is proved and it implies

$$
\forall k \in\{1 \ldots n-1\} \quad \ell^{b}\left(\left(\Upsilon\left[r_{k} / c\right]\right)^{b}\right) \leqslant \ell^{a}\left(\Gamma\left[r_{k}\right]\right) \leqslant \alpha \varepsilon_{k},
$$

which, on $\mathbf{C}^{0} \backslash B_{1}$, implies

$$
\forall k \in\{1 \ldots n-1\} \quad \ell^{a}\left(\Upsilon\left[r_{k} / c\right]\right) \leqslant \tilde{\alpha} \varepsilon_{k} .
$$

This entails that $\Upsilon$ is $\left(c^{-1} \mathbf{r}, \tilde{\alpha} \varepsilon\right)$-small. The inclusion $\Upsilon^{b} \subset \Gamma$ was proved at line (81) with $k=n$ because $\Upsilon=\Upsilon\left[r_{n} / c\right] \subset \Gamma\left[r_{n}\right]=\Gamma$.
End of the proof of (80)
As in the proof of Lemma 12 , we can use the fact that under $\mathbf{N}_{a}$, conditionally on $\mathcal{G}_{a}$, if $\Gamma$ is a $\mathcal{T}(a)$-ball, then the process $\left\{\ell^{a+a^{\prime}}\left(\Gamma^{a+a^{\prime}}\right), a^{\prime} \geqslant 0\right\}$ is a Feller diffusion started at $\ell^{a}(\Gamma)$. Using sub-additivity and Lemma $8(i)$, we get

$$
\begin{align*}
\mathbf{N}\left(B_{1}\right) & \leqslant \sum_{k=1}^{n-1} \frac{1}{a} \mathbf{N}_{a}\left[\sum_{i=1}^{Z_{a, r_{k}}} \mathbf{1}_{\left\{\ell^{a}\left(\Gamma_{i}\right) \geqslant \tilde{\alpha} \varepsilon_{k}\right\}} \exp \left(-\delta^{-1}\left(\sqrt{\ell^{a}\left(\Gamma_{i}\right)}-\sqrt{\alpha \varepsilon_{k}}\right)^{2}\right)\right]  \tag{82}\\
& \leqslant \frac{1}{a} \exp \left(-\delta^{-1}\left(\sqrt{\tilde{\alpha} \varepsilon_{k}}-\sqrt{\alpha \varepsilon_{k}}\right)^{2}\right) \sum_{k=1}^{n-1} \mathbf{E}\left[Z_{a, r_{k}}\right] \tag{83}
\end{align*}
$$

The proof is completed recalling that for all $k \in\{1 \ldots n-1\}, \varepsilon_{k} \leqslant \varepsilon_{n-1}$, and that, by Lemma 4, $\mathbf{N}_{a}\left[Z_{a, r_{k}}\right]=\frac{2 a}{r_{k}} \leqslant \frac{2 a}{r_{n}}$.

Let us introduce

$$
\begin{equation*}
\forall j \in \mathbb{N}, \quad r_{j}=2^{-j} \quad \text { and } \quad \varepsilon_{j}=g\left(r_{j}\right) \tag{84}
\end{equation*}
$$

and then

$$
\begin{equation*}
\forall p \in \mathbb{N}, \quad j_{p}=\left\lfloor(4 / 3)^{p}\right\rfloor, \quad \mathbf{r}^{(p)}=\left\{r_{j}, j_{p} \leqslant j \leqslant j_{p+1}-1\right\} \quad \text { and } \quad \varepsilon^{(p)}=\left\{\varepsilon_{j} ; j_{p} \leqslant j<j_{p+1}-1\right\} . \tag{85}
\end{equation*}
$$

Let $m \in(0,1 / 2)$, we also introduce the following discrete grid

$$
\begin{equation*}
G^{\prime}(p, m):=\left\{m^{-1}+k \delta_{p}, k \in \mathbb{N}^{*}\right\} \cap\left[m, m^{-1}\right] \tag{86}
\end{equation*}
$$

where $\delta_{p}$ is the mesh of the grid, given by

$$
\begin{equation*}
\delta_{p}=r\left(j_{p+1}\right)^{5 / 4} \tag{87}
\end{equation*}
$$

Note that $G^{\prime}(p, m)$ contains less than $\left(m \delta_{p}\right)^{-1}$ points.
Lemma 15 Let $\alpha \in(0,1 / 2), m \in(0,1 / 2)$. For $p \in \mathbb{N}$, denote $u_{p}:=g\left(r\left(j_{p+1}\right)\right)^{-1} p^{-2}$. Then there exists $p_{0} \in \mathbb{N}$ only depending on $\alpha, m$ such that for all $p \geqslant p_{0}$,

$$
\begin{equation*}
\mathbf{N}\left(\sup _{b \in\left[m, m^{-1}\right]} S_{b, \boldsymbol{r}^{(p)}, \alpha \boldsymbol{\varepsilon}(p)}>u_{p}\right) \leqslant p^{-2} \tag{88}
\end{equation*}
$$

Proof. Let $\tilde{\alpha} \in(\alpha, 1 / 2)$ and $c$ in $(1, \infty)$ such that $2 c \tilde{\alpha} \in(0,1)$. In what follows, we denote $T_{0}^{\prime}$ the left-hand-side of (88). Observe that $T_{0}^{\prime} \leqslant T_{1}^{\prime}+T_{2}^{\prime}$, where we have set

$$
\begin{aligned}
& T_{1}^{\prime}=\mathbf{N}\left(\sup _{a \in G^{\prime}(p, m)} S_{a, c^{-1} \mathbf{r}^{(p)}, \tilde{\alpha} \varepsilon^{(p)}} \leqslant u_{p} ; \sup _{b \in\left[m, m^{-1}\right]} S_{b, \mathbf{r}^{(p)}, \alpha \boldsymbol{\varepsilon}^{(p)}}>u_{p}\right), \\
& T_{2}^{\prime}=\mathbf{N}\left(\sup _{a \in G^{\prime}(p, m)} S_{a, c^{-1} \mathbf{r}^{(p)}, \tilde{\alpha} \boldsymbol{\varepsilon}^{(p)}}>u_{p}\right) .
\end{aligned}
$$

Using sub-additivity and Lemma 14, we get

$$
\begin{aligned}
T_{1}^{\prime} & \leqslant \mathbf{N}\left(\bigcup_{a \in G^{\prime}(p, m)}\left\{\sup _{b \in\left[a, a+\delta_{p}\right]} S_{b, \mathbf{r}^{(p)}, \alpha \boldsymbol{\varepsilon}^{(p)}}>S_{a, c^{-1} \mathbf{r}^{(p)}, \tilde{\alpha} \boldsymbol{\varepsilon}^{(p)}}\right\}\right) \\
& \leqslant\left(m \delta_{p}\right)^{-1} \sup _{a \in G^{\prime}(p, m)} \mathbf{N}\left(\sup _{b \in\left[a, a+\delta_{p}\right]} S_{b, \mathbf{r}^{(p)}, \alpha \varepsilon^{(p)}}>S_{a, c^{-1} \mathbf{r}^{(p)}, \tilde{\alpha} \varepsilon^{(p)}}\right) \\
& \leqslant\left(m \delta_{p}\right)^{-1} \frac{2\left(j_{p+1}-j_{p}\right)}{r\left(j_{p+1}\right)} \exp \left(-(\sqrt{\tilde{\alpha}}-\sqrt{\alpha})^{2} \delta_{p}^{-1} g\left(r\left(j_{p+1}-2\right)\right)\right) .
\end{aligned}
$$

One has $\delta_{p}^{-1} g\left(r\left(j_{p+1}-2\right)\right) \geqslant \delta_{p}^{-1} g\left(r\left(j_{p+1}\right)\right)=r\left(j_{p+1}\right)^{-1 / 4} \log \log 1 / r\left(j_{p+1}\right)$, which implies that $T_{1}^{\prime}$ is smaller than $(1 / 2) p^{-2}$, for all $p$ sufficiently large (it is obviously not a sharp bound).

Recalling the definitions (44), we set

$$
\begin{equation*}
\mu_{p}=\mu\left(c^{-1} \mathbf{r}^{(p)}, \tilde{\alpha} \varepsilon^{(p)}\right)=\mathbf{N}\left(S_{r\left(j_{p}\right) /(2 c), c^{-1} \mathbf{r}^{(p)}, \tilde{\alpha} \boldsymbol{\varepsilon}^{(p)}}\right) \tag{89}
\end{equation*}
$$

We will prove that $T_{2}^{\prime} \leqslant T_{3}^{\prime}+T_{4}^{\prime}$, where

$$
\begin{aligned}
& T_{3}^{\prime}=\mathbf{N}\left(\sup _{a \in G^{\prime}(p, m)}\left|S_{a, c^{-1} \mathbf{r}^{(p)}, \tilde{\alpha} \varepsilon^{(p)}}-\mu_{p} \ell^{a}(\mathcal{T})\right|>u_{p} / 2\right) \\
& T_{4}^{\prime}=\mathbf{N}\left(\sup _{a \in G^{\prime}(p, m)} \mu_{p} \ell^{a}(\mathcal{T})>u_{p} / 2\right)
\end{aligned}
$$

By sub-additivity and a Markov inequality involving a moment of order 4, we get

$$
\begin{align*}
T_{3}^{\prime} & \leqslant\left(m \delta_{p}\right)^{-1} \sup _{a \in G^{\prime}(p, m)} \mathbf{N}\left(\left|S_{a, c^{-1} \mathbf{r}^{(p)}, \tilde{\alpha} \varepsilon^{(p)}}-\mu_{p} \ell^{a}(\mathcal{T})\right|>u_{p} / 2\right) \\
& \leqslant\left(m \delta_{p}\right)^{-1} 2^{4} u_{p}^{-4} \sup _{a \in G^{\prime}(p, m)} \mathbf{N}\left[\left(S_{a, c^{-1} \mathbf{r}^{(p)}, \tilde{\alpha} \varepsilon^{(p)}}-\mu_{p} \ell^{a}(\mathcal{T})\right)^{4}\right] \tag{90}
\end{align*}
$$

We want to apply Lemma 10 with $\mathbf{r}=, c^{-1} \mathbf{r}^{(p)}$ and $\boldsymbol{\varepsilon}=\tilde{\alpha} \boldsymbol{\varepsilon}^{(p)}$. Thus, recalling (84) and (85), we check that for all sufficiently large $p, m / r\left(j_{p}\right)>1$ and $r\left(j_{p}\right) / r\left(j_{p+1}-1\right)$. Recalling that $c_{0} \in\left(0,10^{4}\right]$ is the universal constant given by Lemma 10, we get from (90)

$$
\begin{equation*}
T_{3}^{\prime} \leqslant\left(m \delta_{p}\right)^{-1} 2^{4} u_{p}^{-4} \sup _{a \in G^{\prime}(p, m)} c_{0} a \frac{r\left(j_{p}\right)^{2}}{r\left(j_{p+1}-1\right)^{4}} \leqslant 2^{4} c_{0} m^{-2} \frac{r\left(j_{p}\right)^{2}}{\delta_{p} u_{p}^{4} r\left(j_{p+1}\right)^{4}} \tag{91}
\end{equation*}
$$

Recall that $u_{p}=g\left(r\left(j_{p+1}\right)\right)^{-1} p^{-2}$, and by (84) and (85), we get $\log \log \left(1 / r\left(j_{p}\right)\right) \underset{p \rightarrow \infty}{\sim} p \log (4 / 3)$. Hence, $u_{p} \geqslant p^{-3} r\left(j_{p+1}\right)^{-1}$ and (91) implies

$$
\begin{equation*}
T_{3}^{\prime} \leqslant 2^{4} c_{0} m^{-2} p^{12} \frac{r\left(j_{p}\right)^{2} r\left(j_{p+1}\right)^{4}}{r\left(j_{p+1}\right)^{5 / 4} r\left(j_{p+1}\right)^{4}} \tag{92}
\end{equation*}
$$

Now, one can plainly check that $\frac{r\left(j_{p}\right)^{2}}{r\left(j_{p+1}\right)^{5 / 4}}$ is smaller than $r\left(j_{p}\right)^{1 / 3}$. Thus, $T_{3}^{\prime}$ is smaller than $(1 / 4) p^{-2}$ for all $p$ sufficiently large.

For the term $T_{4}^{\prime}$, we use Lemma 9 to obtain

$$
\begin{equation*}
T_{4}^{\prime} \leqslant(2 / m) \exp \left(-(m / 4) u_{p} \mu_{p}^{-1}\right) \tag{93}
\end{equation*}
$$

Recalling (89) and Lemma 11, we get that for all $p$,

$$
\begin{equation*}
\mu_{p} \leqslant \frac{5}{r\left(j_{p+1}\right)}\left(\prod_{j=j_{p}}^{j_{p+1}-2} \mathbf{P}\left(\Lambda_{r_{j+1} / c, r_{j} / c}^{*} \leqslant \tilde{\alpha} r_{j} \log \log \left(1 / r_{j}\right)\right)\right)^{1 / 2} \tag{94}
\end{equation*}
$$

We want to get an lower bound of $u_{p} \mu_{p}^{-1}$, so we compute an upper bound for $u_{p}^{-1} \mu_{p}$. Recalling that $u_{p} \geqslant p^{-3} r\left(j_{p+1}\right)^{-1}$, one has

$$
\begin{equation*}
u_{p}^{-1} \mu_{p} \leqslant 5 p^{3} \exp \left(\frac{1}{2} \sum_{j=j_{p}}^{j_{p+1}-2} \log \left(1-q_{j}\right)\right) \leqslant 5 p^{3} \exp \left(-\frac{1}{2} \sum_{j=j_{p}}^{j_{p+1}-2} q_{j}\right) \tag{95}
\end{equation*}
$$

where $q_{j}=\mathbf{P}\left(\Lambda_{r_{j+1} / c, r_{j} / c}^{*}>\tilde{\alpha} r_{j} \log \log \left(1 / r_{j}\right)\right)$. Recalling that $r_{j}=2^{-j}$, it follows from Lemma 6 that

$$
\begin{aligned}
q_{j} & =\left(1-\frac{1}{2}\right)^{2} \frac{2 \tilde{\alpha} r / \rho \log \log 1 / r_{j}}{r / / j c} \exp \left(-\frac{2 \tilde{\alpha} r / j \log \log 1 / r_{j}}{r / \rho / c}\right)+\left(1-\frac{1}{4}\right) \exp \left(-\frac{2 \tilde{\alpha} r / \rho \log \log 1 / r_{j}}{r / / / c}\right) \\
& \underset{j \rightarrow \infty}{\sim} \frac{\tilde{\alpha} c}{2} \log \log \left(1 / r_{j}\right) e^{-2 \tilde{\alpha} c \log \log \left(1 / r_{j}\right)} \\
& \underset{j \rightarrow \infty}{\sim} c^{\prime} \log (j) j^{-2 \tilde{\alpha} c},
\end{aligned}
$$

where $c^{\prime}$ is a positive constant depending on $\alpha, \tilde{\alpha}, c$. We stress that the particular choice of $c$ was made to ensure that $\chi:=1-2 \tilde{\alpha} c$ is strictly positive, so that the following is true for all large $p$ :

$$
\sum_{j_{p}}^{j_{p+1}-2} q_{j} \geqslant \sum_{j_{p}}^{j_{p+1}-2} j^{-2 \tilde{\alpha} c} \geqslant \int_{j_{p}}^{j_{p+1}-1} x^{-2 \tilde{\alpha} c} \mathrm{~d} x \underset{p \rightarrow \infty}{\sim} \chi^{-1}\left((4 / 3)^{\chi}-1\right)\left(\frac{4}{3}\right)^{p \chi}
$$

Thus, for all $p$ sufficiently large, $\sum_{j_{p}}^{j_{p+1}} q_{j} \geqslant 2 p$ which, combined with (95), entails that $u_{p}^{-1} \mu_{p} \leqslant 5 p^{3} \exp (-p)$. Thus, $u_{p} \mu_{p}^{-1} \geqslant 5^{-1} p^{-3} e^{p}$. Finally, we see from (93) that $T_{3}^{\prime}$ is smaller than $(1 / 4) p^{-2}$ for all $p$ sufficiently large, which ends the proof.

### 4.2.2 Proof of Theorem 3.

Let $\alpha \in(0,1 / 2)$. For a level $a \in(0, \infty)$, we recall the definition (54) of $\Delta_{a}^{\alpha}$. To show that the $g$-Hausdorff measure of $\Delta_{a}^{\alpha}$ is null, we need an efficient covering of this set. Let us recall the integer sequence $j_{p}=\left\lfloor(4 / 3)^{p}\right\rfloor$ and the radii $r_{j}=2^{-j}$. For $p \in \mathbb{N}$, we recall the definition of the finite subsets $\mathbf{r}^{(p)}=\left\{r_{j}, j_{p} \leqslant j \leqslant j_{p+1}-1\right\}$, and $\varepsilon^{(p)}=\left\{\varepsilon_{j}, j_{p} \leqslant j<j_{p+1}-1\right\}$ where $\varepsilon_{j}=g\left(r_{j}\right)$. Recalling the definition (42) for small balls, we set

$$
\mathscr{C}_{n}:=\bigcup_{p=n}^{\infty}\left\{\Gamma \in \mathscr{B}_{a, r\left(j_{p+1}\right)}: \quad \Gamma \text { is }\left(\mathbf{r}^{(p)}, \alpha \varepsilon^{(p)}\right)-\text { small }\right\}
$$

Observe that if $\sigma \in \Delta_{a}^{\alpha}$, then the $\mathcal{T}(a)$-ball $\Gamma\left(\sigma, r\left(j_{p+1}\right)\right)$ is $\left(\mathbf{r}^{(p)}, \alpha \boldsymbol{\varepsilon}^{(p)}\right)$-small for all large $p$, thus for all $n \in \mathbb{N}$, we have $\Delta_{a}^{\alpha} \subset \mathscr{C}_{n}$. Let us recall the definition (15) of Hausdorff measures, and the fact that the diameter of a $\mathcal{T}(a)$-ball is smaller than its radius. We get

$$
\begin{equation*}
\forall a \in\left[m, m^{-1}\right] \quad \mathscr{H}_{g}^{\left(r\left(j_{p+1}\right)\right)}\left(\Delta_{a}^{\alpha}\right) \leqslant \sum_{p=n}^{\infty} S_{a, \mathbf{r}^{(p)}, \alpha \varepsilon^{(p)}} \cdot g\left(r\left(j_{p+1}\right)\right) \tag{96}
\end{equation*}
$$

because $\Delta_{a}^{\alpha} \subset \mathscr{C}_{n}$. Thus,

$$
\begin{equation*}
\forall a \in\left[m, m^{-1}\right] \quad \mathscr{H}_{g}\left(\Delta_{a}^{\alpha}\right) \leqslant \limsup _{n \rightarrow \infty} \sum_{p=n}^{\infty} S_{a, \mathbf{r}^{(p)}, \alpha \varepsilon^{(p)}} \cdot g\left(r\left(j_{p+1}\right)\right) \tag{97}
\end{equation*}
$$

Now, let $m \in(0,1 / 2)$. Applying Lemma 15 , we easily get that

$$
\sum_{p=1}^{\infty} \mathbf{N}\left(\sup _{a \in\left[m, m^{-1}\right]} S_{a, \mathbf{r}^{(p)}, \alpha \boldsymbol{\varepsilon}^{(p)}}>u_{p}\right)<\infty
$$

where we recall the notation $u_{p}=g\left(r\left(j_{p+1}\right)\right)^{-1} p^{-2}$. By Borel-Cantelli lemma there exists a subset $\mathbf{V}^{\prime} \subset \mathbf{C}^{0}$ such that $\mathbf{N}\left(\mathbf{C}^{0} \backslash \mathbf{V}^{\prime}\right)=0$ and such that

$$
\text { on } \mathbf{V}^{\prime}, \quad g\left(r\left(j_{p+1}\right)\right) \sup _{a \in\left[m, m^{-1}\right]} S_{a, \mathbf{r}^{(p)}, \alpha \varepsilon^{(p)}} \leqslant p^{-2}, \quad \text { for all suff. large } p .
$$

Combined with (96), we deduce on $\mathbf{V}^{\prime}$, for $a \in\left[m, m^{-1}\right]$, one has

$$
\mathscr{H}_{g}\left(\Delta_{a}^{\alpha}\right) \leqslant \lim _{n \rightarrow \infty} \sum_{p=n}^{\infty} p^{-2}=0
$$

which is the desired result.

### 4.3 Proof of Theorem 1.

Let $\kappa \in\left(\frac{1}{2}, \infty\right), \alpha \in\left(0, \frac{1}{2}\right)$, and $m \in(0,1 / 2)$. Theorem 2 entails that there exists a Borel subset $\mathbf{V}=\mathbf{V}(\kappa, m) \subset \mathbf{C}^{0}$ such that $\mathbf{N}\left(\mathbf{C}^{0} \backslash \mathbf{V}\right)=0$ and
on $\mathbf{V}(\kappa, m), \quad$ for all Borel subset $\mathcal{A} \subset \mathcal{T}, \quad \forall a \in\left[m, m^{-1}\right], \quad \ell^{a}(\mathcal{A}) \leqslant \kappa \mathscr{H}_{g}(\mathcal{A} \cap \mathcal{T}(a))$.
Now, let us rewrite the definition (54)

$$
\begin{equation*}
\Delta_{a}^{\alpha}=\left\{\sigma \in \mathcal{T}(a): \limsup _{r \rightarrow 0} \frac{\ell^{a}(B(\sigma, r))}{g(r)}<\alpha\right\} \tag{99}
\end{equation*}
$$

According to Theorem 3, there exists a Borel subset $\mathbf{V}^{\prime}=\mathbf{V}^{\prime}(\alpha, m) \subset \mathbf{C}^{0}$ such that $\mathbf{N}\left(\mathbf{C}^{0} \backslash \mathbf{V}^{\prime}\right)=$ 0 and

$$
\begin{equation*}
\text { on } \mathbf{V}^{\prime}(\alpha, m) \quad \forall a \in\left[m, m^{-1}\right] \quad \mathscr{H}_{g}\left(\Delta_{a}^{\alpha}\right)=0 \tag{100}
\end{equation*}
$$

Let $\alpha^{\prime}<\alpha$ and notice that $\mathcal{T}(a) \backslash \Delta_{a}^{\alpha} \subset\left\{\sigma: \limsup _{r \rightarrow 0} \frac{\ell^{a}(B(\sigma, r))}{g(r)}>\alpha^{\prime}\right\}$. Moreover, from (13), we know that $\mathbf{N}$-a.e. for all $a \in(0, \infty), \ell^{a}(\mathcal{T} \backslash \mathcal{T}(a))=0$. Thus, on $\mathbf{V}^{\prime}$, for all Borel subset $\mathcal{A} \subset \mathcal{T}$, and for all $a \in\left[m, m^{-1}\right]$ and all $\tilde{\alpha}<\alpha$, Lemma 3 entails

$$
\begin{equation*}
\ell^{a}(\mathcal{A}) \geqslant \ell^{a}\left(\mathcal{A} \cap\left(\mathcal{T}(a) \backslash \Delta_{a}^{\alpha}\right)\right) \geqslant \alpha^{\prime} \mathscr{H}_{g}\left(\mathcal{A} \cap\left(\mathcal{T}(a) \backslash \Delta_{a}^{\alpha}\right)\right)=\alpha^{\prime} \mathscr{H}_{g}(\mathcal{A} \cap \mathcal{T}(a)) \tag{101}
\end{equation*}
$$

where we used (100) for the last equality. Letting $\alpha^{\prime} \rightarrow \alpha$, we get

$$
\begin{equation*}
\text { on } \mathbf{V}^{\prime}(\alpha, m) \quad \text { for all Borel subset } \mathcal{A} \subset \mathcal{T}, \quad \forall a \in\left[m, m^{-1}\right] \quad \ell^{a}(\mathcal{A}) \geqslant \alpha \mathscr{H}^{g}(\mathcal{A} \cap \mathcal{T}(a)) . \tag{102}
\end{equation*}
$$

Now, let us set

$$
\begin{equation*}
\tilde{\mathbf{V}}:=\left(\bigcap_{\substack{\kappa \in(1 / 2, \infty) \cap \mathbb{Q} \\ m \in(0,1 / 2) \cap \mathbb{Q}}} \mathbf{V}(\kappa, m)\right) \bigcap\left(\bigcap_{\substack{\alpha \in(0,1 / 2) \cap \mathbb{Q} \\ m \in(0,1 / 2) \cap \mathbb{Q}}} \mathbf{V}^{\prime}(\alpha, m)\right) . \tag{103}
\end{equation*}
$$

Clearly, $\tilde{\mathbf{V}}$ is a Borel subset of $\mathbf{C}^{0}$ such that $\mathbf{N}\left(\mathbf{C}^{0} \backslash \tilde{\mathbf{V}}\right)=0$. Moreover, combining (98) and (102), we get that on $\tilde{\mathbf{V}}$, for all Borel subset $\mathcal{A} \subset \mathcal{T}$, and for all level $a \in(0, \infty)$, one has $\ell^{a}(\mathcal{A})=\frac{1}{2} \mathscr{H}_{g}(\mathcal{A} \cap \mathcal{T}(a))$.

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