

Regularized 4D-CT reconstruction from a single dataset with a spatio-temporal prior

Fabien Momey, Eric Thiébaut, Catherine Burnier-Mennessier, Loïc Denis,

Jean-Marie Becker, Laurent Desbat

▶ To cite this version:

Fabien Momey, Eric Thiébaut, Catherine Burnier-Mennessier, Loïc Denis, Jean-Marie Becker, et al.. Regularized 4D-CT reconstruction from a single dataset with a spatio-temporal prior. 2014. hal-00998291v2

HAL Id: hal-00998291 https://hal.science/hal-00998291v2

Preprint submitted on 8 Dec 2014

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Regularized 4D-CT reconstruction from a single dataset with a spatio-temporal prior

Fabien Momey*, Éric Thiébaut, Catherine Burnier, Loïc Denis, Jean-Marie Becker, and Laurent Desbat

Abstract—X-ray Computerized Tomography (CT) reconstructions can be severely impaired by the patient's respiratory motion and cardiac beating. Motion must thus be recovered in addition to the 3D reconstruction problem. The approach generally followed to reconstruct dynamic volumes consists of largely increasing the number of projections so that independent reconstructions be possible using only subsets of projections from the same phase of the cyclic movement. Apart from this major trend, motion compensation (MC) aims at recovering the object of interest and its motion by accurately modeling its deformation over time, allowing to use the whole dataset for 4D reconstruction in a coherent way.

We consider a different approach for dynamic reconstruction based on inverse problems, without any additional measurements, nor explicit knowledge of the motion. The dynamic sequence is reconstructed out of a single data set, only assuming the motion's continuity and periodicity. This inverse problem is solved by the minimization of the sum of a data-fidelity term, consistent with the dynamic nature of the data, and a regularization term which implements an efficient spatio-temporal version of the total variation (TV). We demonstrate the potential of this approach and its practical feasibility on 2D and 3D+t reconstructions of a mechanical phantom and patient data.

Index Terms—Dynamic tomography, Reconstruction, Inverse Problems, Regularization, Signal processing.

I. INTRODUCTION

I N X-ray CT imaging, the motion induced by the breathing and the heart beating of the patient implies that the acquired projections, *i.e.* the data, are not related to the same "static

This work was supported by the MiTiV project (Méthodes Inverses pour le Traitement en Imagerie du Vivant), funded by the French ANR (N $^{\circ}$ ANR-09-EMER-008).

Fabien Momey* was with the Université de Lyon, F-42023, Saint-Etienne, France; CNRS, UMR5516, Laboratoire Hubert Curien, F-42000, Saint-Etienne, France; Université de Saint-Etienne, Jean Monnet, F-42000, Saint-Etienne, France: He was also with the Université de Lyon, Lyon, F-69003, France; Université Lyon 1, Observatoire de Lyon, 9 avenue Charles André, Saint-Genis Laval, F-69230, France; CNRS, UMR 5574, Centre de Recherche Astrophysique de Lyon; École Normale Supérieure de Lyon, Lyon, F-69007, France; CEA, LETI, MINATEC Campus, F-38054 Grenoble, France (e-mail: fabien.momey@cea.fr).

Éric Thiébaut is with the Université de Lyon, Lyon, F-69003, France; Université Lyon 1, Observatoire de Lyon, 9 avenue Charles André, Saint-Genis Laval, F-69230, France; CNRS, UMR 5574, Centre de Recherche Astrophysique de Lyon; École Normale Supérieure de Lyon, Lyon, F-69007, France (e-mail: eric.thiebaut@univ-lyon1.fr).

Loïc Denis, Catherine Burnier and Jean-Marie Becker are with the Université de Lyon, F-42023, Saint-Etienne, France; CNRS, UMR5516, Laboratoire Hubert Curien, F-42000, Saint-Etienne, France; Université de Saint-Etienne, Jean Monnet, F-42000, Saint-Etienne, France (e-mail: {loic.denis,catherine.burnier.jean-marie.becker}@univ-st-etienne.fr).

Laurent Desbat is with the Université Grenoble Alpes, TIMC-IMAG, F-38000 Grenoble, France; CNRS, TIMC-IMAG, F-38000 Grenoble, France; CHU de Grenoble,TIMC-IMAG, F-38000 Grenoble, France (e-mail: laurent.desbat@imag.fr). object". Reconstructing the patient's anatomy ignoring motion causes noticeable artifacts [2], [39], [10]. In other words, the inverse problem of reconstruction has to be addressed in 4D; more specifically in 3D+t, where the motion and deformation of anatomical structures as a function of time are recovered. In image guided radiotherapy (IGRT), this is a critical matter because a X-ray treatment protocol requires a precise localization of the lung tumor, in order to prevent at best healthy tissues from irradiation. As a result, the "object" to be reconstructed is modelized as a spatio-temporal signal f(x, t), defined for spatial coordinates $x \in \mathbb{R}^3$ and time $t \in \mathbb{R}$.

The so-called 4D-CT reconstruction problem has benefited from active research for the past twenty years. Some investigations were made on the scanning protocol itself [17], [31], [46], [47]. Ritchie et al. [31] established early that artifacts were still present even with ultrafast scanning. Such an observation still holds twenty years later. A standard solution is based on a correlation between the 3D-CT data and a 1D temporal record of the patient's pseudo-periodical movements. The acquired projections are then sorted, according to the phase of the cycle to which they are related. Data subsets are extracted, and independent static reconstructions are then performed using each of these subsets. Such a method, called gated 4D-CT or phase-correlated 4D-CT [37], [24], [16], [43], [22], [3], [40], [28], requires a sufficient number of projections for each reconstructed phase to avoid motion artifacts; as a consequence it increases the amount of X-ray dose delivered to the patient for the reconstruction of the whole sequence.

Another widely investigated class of methods aims at recovering the motion itself, *i.e.* models the deformation of the object as a function of time, and uses it to reproduce the motion. The motion is generally estimated in the form of a deformation vector field $\Gamma_t(\mathbf{x})$ that maps the volume at any time t back to its state at a reference time t_0 :

$$f(\boldsymbol{x},t) = f(\Gamma_t(\boldsymbol{x}),t_0) .$$
⁽¹⁾

Such a motion model can be incorporated into the tomographic projector in order to use all the available projections to reconstruct the patient's anatomy at this reference state. This class of methods is called *motion compensation* (MC). Research in this area has been active in the last decade, yielding very efficient algorithms. Such approaches rely on the quality of the estimation of the motion, a challenging inverse problem, especially because of the complex modeling of deformations such as breathing or cardiac beating. They generally require a separate high resolution 3D+t dataset for motion estimation.

In the IGRT context, 4D Cone-Beam CT (4D-CBCT) constitutes a particular modality because phase correlation is made difficult due to the poor number of remaining projections for the reconstruction of a single phase. One must turn to more robust approaches such as motion compensation. In 4D-CBCT, motion estimation is often done on an additional 4D-CT reconstruction [6], [5], [23], [29], [30]. IGRT makes this prior information naturally available thanks to the necessary 4D planning CT prior to treatment. This so-called *planning CT* dataset is reconstructed by a phase-correlated method using a sufficient, *i.e.* enhanced, number of projections. Hence the total amount of data required for MC methods in this context is very high, meaning higher X-ray doses delivered to the patient. However, recent works have established that a very efficient MC reconstruction can be obtained from a single dataset, by a direct estimation of the motion model directly on the current measurements. In [7], [8], the authors estimate the motion from a phase-correlated 4D-CT reconstruction on CBCT data, correcting the motion artifacts to recover accurately the deformation. Then they use the motion model to get the 4D MC reconstruction in a second step. In [45], [49], the two inverse problems (motion estimation and reconstruction) are solved in a single algorithm using alternating schemes.

In parallel with the development of MC methods, phasecorrelated reconstruction approaches have also benefited from very interesting improvements. Following the emergence of the theory of compressive sensing, robust reconstruction algorithms were proposed, based on sparsity constraints, e.g. piecewise continuity, to allow efficient motion artifacts-free 3D reconstructions of each phase from a limited number of projections [11]. However such methods still reconstruct each phase independently from each other. Only the most recent approaches propose to take into account the temporal correlation between the cyclic phases of the object under motion, extending the spatial prior knowledge to the temporal dimension. Their basic asumption is the continuity - or piecewise continuity — in both spatial and temporal dimensions, of the 4D image to be reconstructed. In particular, the approach proposed in [42] regularizes a 4D-CT reconstruction by temporal correlation of the different phases thanks to the minimization of nonlocal means between patches taken over the temporal dimension. Other methods [48], [32] use the Rudin et al. [33] total variation prior in 4D, applying it on numerical 2D XCAT phantom CBCT-like data [48] and in vivo cardiac micro-CT mouse data [32].

Our approach follows the same principles as these spatiotemporally regularized 4D phase-correlated reconstruction methods: we use only the current dataset for the reconstruction of the 3D+t object, suppressing in this way any need for additional projections. No explicit knowledge of the motion is required and the whole 3D+t sequence is directly reconstructed, globally addressing the inverse problem of reconstruction from dynamic projections as the joint minimization of a data-fidelity term and a regularization term, as detailed in (2). The method is therefore very general and can readily be applied to many types of 4D-CT modalities. One of the novel contribution in our work is that we developed a global 4D tomographic projection model, based on classical static projection approach coupled with a temporal interpolation of the 4D object, allowing to calculate a projection at any given date t. Thus we overcome the classical phase correlation approach which only sorts the projections phase by phase with a given temporal gap. In addition, the specific continuity of the spatio-temporal object is enhanced by a 3D+t edge-preserving smoothness regularization module based on relaxed Rudin et al. [33] total variation (TV).

In this introduction, we have presented the state-of-the-art of 4D-CT reconstruction, from the phase-correlated approaches to motion compensation, and presented the method we describe in this paper (*cf.* Section III). As explained in Section II, the demonstration of our method is made in the context of 4D-CBCT. In Section IV, we demonstrate its efficiency on 2D+*t* reconstructions from numerically simulated 2D+*t* data. Then we apply our method to real data, performing 3D+t reconstructions of a mechanical phantom and patient data, for which we confront our results to a standard MC reconstruction presented in [29]. Applying our spatio-temporally regularized 4D phase-correlated reconstruction method to real human data acquired in the context of IGRT scanning protocol, is the other major contribution of our work, proposing to explore the feasibility of such an approach, in contrast to MC methods.

II. CONE-BEAM CT: DEMONSTRATION FRAMEWORK

A. Motion is critical in 4D-CBCT

4D-CBCT is an especially interesting modality for the study of 4D respiratory motion CT. Indeed, CBCT scanners have been implanted on linear accelerators for radiotherapy in the early 2000's [18], [19], [20]. The CT scan acquired during the treatment is for example used to control a posteriori the delivered therapeutic X-ray dose.

This type of system has 2 specificities: the slow period of rotation of its gantry (from 1 to 2 minutes), and the fact that all the projections (about 600-700) are acquired on a single rotation of its flat detector. On the one hand, this slow rotation speed allows several respiratory cycles to be completed during the acquisition, giving a satisfactory angular coverage for each reconstructed phase. But, on the other hand, the lack of projections makes 4D-CT reconstructions difficult.

Many MC reconstruction methods have been proposed specifically for this type of data [23], [50], [29], [30]. Li et al. [23] have incorporated the deformation field in the analytic reconstruction algorithm FDK [15]. Rit et al. [29], [30] have compared the Li et al.'s approach with an iterative reconstruction method based on the SART algorithm [1], inserting the deformation field into the tomographic projection model. These approaches assume that motion is unchanged between the acquisition on which motion is estimated and the CBCT acquisition for which motion is compensated. Zeng *et al.* [50] have proposed to avoid a preliminary 4D-CT reconstruction step by a direct estimation of the deformation model on the CBCT projections, using the complementary information given by an a priori 3D static reconstruction of the patient. The 3D anatomical model still has to be obtained from a previous reconstruction step, generally from breath-hold acquired data, a constraining step that generates additional irradiation. However, as already mentioned in Section I, very efficient MC reconstruction methods have been recently developed that

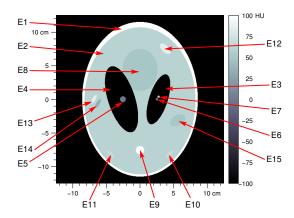


Fig. 1. Temporal phase at t = 0 of our Shepp-Logan-like 2D dynamic phantom. It simulates a 2D object with 256×256 spatial voxels of size $0.1 \times 0.1 \text{ cm}^2$, *i.e.* an global virtual size of $25.6 \times 25.6 \text{ cm}^2$. It is composed of 15 ellipses Ei identified on the image. The values are shown in Hounsfield units, with $\mu_{\text{eau}} = 0.1928 \text{ cm}^{-1}$.

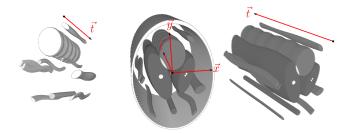


Fig. 2. Segmented 3D views of one motion cycle of our dynamic 2D Shepp-Logan phantom. The depth corresponds to the temporal dimension, along which the deformations of some of its ellipses over time are clearly visualized by oscillations.

perform the motion estimation and the 4D reconstruction by using only the CBCT dataset [7], [8], [45], [49]. Thus they overcome the approximations and possible errors induced by the hypothesis of reproducibility of the motion between two scans. Moreover they reduce the effective X-ray dose required for the reconstruction, a common objective with our approach.

B. "2D CBCT" dynamic data simulation

For testing purposes, we developed a simulated dynamic 2D phantom based on the well known Shepp-Logan model [38], [21]. The components of this phantom are ellipses whose position and shape may vary following a temporal sinusoidal signal with a given period. More precisely, these ellipses undergo changes (translations, rotations, dilations) that simulate the anatomical variations induced by the respiratory and cardiac motions at any instant t. Fig. 1 shows the phantom and the positions of its ellipses at the phase t = 0 of the motion cycle. The details of the motion features of each ellipse are summarized in Tab. I. The column Mot indicates the motion type : T for a translation, R for a rotation and D for a dilation/contraction. The shape parameters are : the center of the ellipse (x_0, y_0) (in cm), the half-axis lengths a and b (in cm), the tilt from the horizontal axis of the major half-axis α (in °), and the attenuation coefficient μ given as a factor of the attenuation coefficient of water (in

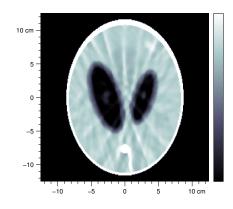


Fig. 3. Static reconstruction from the set of dynamic projections generated by the simulated dynamic Shepp-Logan phantom.

cm⁻¹). The amplitude \mathcal{A} of the motion type is attached to the corresponding shape parameter of the moving ellipse. Note that some ellipses are kept static (motion type "–"). The effective motion is given by the variation of the shape parameter following a temporal periodic sine or cosine signal (column *Sig*). The respiratory motion is simulated by giving to some of the ellipses a movement period of $t_a = 5$ seconds. The cardiac cycle is simulated by giving ellipse E8 a 1 second period. Hence the phantom can be calculated at any date *t* of a cycle. The attenuation value at a given spatial point is computed as the sum of all the ellipses levels of gray passing on it. Fig. 2 illustrates, on 3D views, the anatomical evolution of one motion cycle of this phantom. The observed waves in the temporal direction show the anatomical variations of the ellipses.

Fan beam projections are computed in close form under the hypothesis of an instantaneous measurement. This yields CBCT-like dynamic data in 2D. We simulated 600 such projections for a 360° angle coverage corresponding to a total acquisition time of 120 seconds. As a result, 24 complete motion cycles are done during the acquisition, with 25 projections per cycle.

Fig. 3 displays a static reconstruction using the generated set of dynamic projections. The strong motion artifacts due to the inappropriate "static hypothesis" highlight the need for a 2D+t sequence reconstruction.

III. METHOD

A. Reconstruction criterion

We solve the inverse problem [41] of 3D+t reconstruction by minimizing the joint criterion:

$$f^+ = \operatorname*{arg\,min}_{f} \left\{ \mathcal{J}_{\mathsf{data}}(f) + \mathcal{J}_{\mathsf{prior}}(f) \right\} \,,$$
 (2)

where \mathcal{J}_{data} and \mathcal{J}_{prior} are respectively the data-fidelity and regularization terms (detailed in the following paragraphs) and f denotes the parameters of the 3D+t sequence to be reconstructed, *i.e.* a 4D image of "spatio-temporal" voxels, corresponding to one period of motion.

\mathcal{E}_i	Mot	Sig	Position (cm)				half-axis (cm)				Tilt (°)		Attenuation
			x_0	\mathcal{A}_{x_0}	y_0	\mathcal{A}_{y_0}	a	\mathcal{A}_{a}	b	\mathcal{A}_b	α	\mathcal{A}_{lpha}	$(\mu_{\rm water}{\rm cm}^{-1})$
E1	-	-	0	-	0	-	11.73	-	8.80	-	90	-	2.0
E2	-	-	0	-	-0.23	-	11.14	-	8.44	-	90	-	-0.94
E3	D	\sin	2.80	-	0	-	3.95	0.51	1.40	0.51	72	-	-0.16
E4	D	\sin	-2.80	-	0	-	5.23	0.51	2.04	0.51	108	-	-0.16
E5	Т	$x : \cos y : \sin y$	-2.80	0.20	0	0.76	0.44	-	0.44	-	72	-	0.1
E6	Т	$x : \cos y : \sin y$	2.68	-0.10	0	0.38	0.19	-	0.19	-	72	-	0.1
E7	Т	$x : \cos y : \sin y$	2.93	-0.10	0.38	0.38	0.19	-	0.19	-	72	-	0.15
E8	D	sin	0	-	4.46	-	3.19	0.1275	2.68	0.1275	90	-	-0.01
E9	-	-	0	-	-7.71	-	0.64	-	0.64	-	0	-	0.94
E10	T&R	\sin	4.46	0.51	-8.29	0.51	0.76	-	0.25	-	45	15	0.01
E11	T&R	sin	-4.46	0.51	-8.29	0.51	0.76	-	0.25	-	135	-15	0.01
E12	D	sin	3.82	-	7.65	-	1.02	0.25	0.51	0.13	-45	-	0.025
E13	Т	sin	-7.01	0.25	0	0.51	0.76	-	0.25	-	60	-	0.025
E14	Т	sin	-6.37	-0.25	-0.64	-0.51	0.76	-	0.25	-	60	-	-0.025
E15	D	sin	5.74	-	-3.19	- T	1.27 ABLE I	-0.25	8.80	-0.13	30	-	-0.01

TABLE OF ATTRIBUTES OF THE 15 ELLIPSES OF THE SHEPP-LOGAN-LIKE 2D DYNAMIC PHANTOM OF FIG. 1.

B. Data fidelity and direct model

Assuming independent projections and Gaussian noise, the data fidelity can be written:

$$\mathcal{J}_{\mathsf{data}}(\boldsymbol{f}) = \sum_{\theta \in \Theta} \left\| \boldsymbol{y}^{\theta} - \mathbf{R}^{\theta} \cdot \mathbf{S}^{t_{\theta}} \cdot \boldsymbol{f} \right\|_{\mathbf{W}_{\theta}}^{2}, \qquad (3)$$

where y^{θ} denotes the measured projections at orientation θ of the detector and time t_{θ} , Θ is the set of projection angles. The term $\mathbf{R}^{\theta} \cdot \mathbf{S}^{t_{\theta}} \cdot \mathbf{f}$ is the model of the data, operator $\mathbf{S}^{t_{\theta}}$ interpolates its arguments at time t_{θ} while operator \mathbf{R}^{θ} performs the projection for orientation θ (as further detailed below). The discrepancy between the data and its model is measured by the squared Mahalanobis distance:

$$\|oldsymbol{u}\|_{\mathbf{W}_{ heta}}^2 = oldsymbol{u}^ op\cdot\mathbf{W}_{ heta}\cdotoldsymbol{u}$$

with \mathbf{W}_{θ} a weighting matrix equal to the inverse of the covariance of the noise.

As already mentioned, we assume that the acquisition of a projection is instantaneous. This hypothesis motivates the decomposition of the dynamic projection of the 3D+t object f in two successive operations:

1) The operator $S^{t_{\theta}}$ performs a temporal interpolation of the 3D+t voxels in order to extract the 3D object $f(x, t_{\theta})$ at

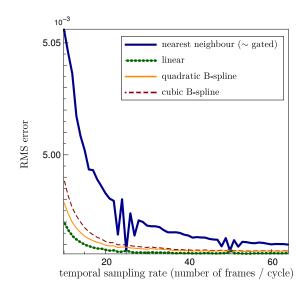


Fig. 4. RMS errors between simulated data of our dynamic Shepp-Logan phantom and the output of our dynamic projection model, $\mathbf{R}^{\theta} \cdot \mathbf{S}^{t_{\theta}}$, as a function of the temporal sampling rate. Each curve is representative of a given temporal interpolation method implemented by the operator $\mathbf{S}^{t_{\theta}}$.

the projection time t_{θ} . Our choice is justified by the fact that the temporal variation of the object is a continuous phenomenon. Hence it is possible to approximate the anatomical state of the object at any date by interpolating the frames sampled by the 4D voxels. Of course, the quality of the interpolation depends on the temporal sampling rate, as well as on the type of interpolation used. In this way we go further from classical 4D phase-correlated methods that simply divide the period of a cycle into equally spaced intervals and sort the projections.

Fig. 4 shows the evolution of the root mean square (RMS) error between the data simulated in Section II-B and the output of our dynamic projection model as a function of the temporal sampling rate, for different interpolation methods. As expected, the error decreases when the number of frames increases. In phase correlated 4D-CT, the selection of projections can be compared to using a nearest neighbor interpolation, which generates much more modeling errors than linear or B-spline-based interpolations as can be seen on Fig. 4. Oscillations in the RMS error curves (notably for the nearest neighbor interpolation) correspond to the sampling rates where some frames are exactly synchronized with the actual projections. In these particular cases, $S^{t_{\theta}}$ is just the identity operator and the remaining errors are due to the spatial sampling of f and to the approximations made by the projector \mathbf{R}^{θ} . In our implementation, we use linear temporal interpolation as it gives the least approximation errors (see Fig. 4).

The coefficients of $S^{t_{\theta}}$ are determined from a 1D temporal signal giving the periodic evolution of the motion, similar to the signals used for 4D-CT reconstructions. It is easy from this signal to extract the period of a cycle, and to identify the phase of each temporal frame. Then the projection dates t_{θ} give the locations of the interpolation points on the temporal axis, from which the coefficients of $S^{t_{\theta}}$ are deduced. Since the period of the respiratory or cardiac motion may slightly vary from one cycle to another, the sequence f has to be normalized to a mean period. Hence the cycles in the 1D signal are dilated or contracted in order to coincide with the normalized cycle. This permits to achieve a finer and more realistic calculation of the interpolation coefficients. A similar temporal registration method is used by Blondel et al. [4], [5] for cardiac SPECT. In our tests, we made the same simplification as Blondel et al. [4], [5], i.e. we have neglected the variation of the spatial amplitude of the motion from one cycle to another. Our results on empirical data discussed in sections IV-B and IV-C support the use of this simplifying hypothesis.

2) The operator \mathbf{R}^{θ} is the tomographic projector which performs a static projection of the 3D object interpolated at t_{θ} . \mathbf{R}^{θ} can be any numerical model implemented for iterative CT reconstruction. As already mentioned, projections must be separated according to the motion phase in order to reconstruct the temporal sequence. This implies a drastic reduction of the number of projections available per reconstructed temporal frame. In this context of reconstruction from a few projections, we have shown [25], [26] that the accuracy of the projector is a critical issue to make

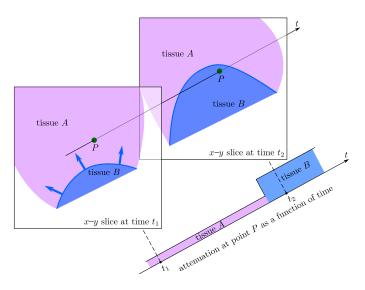


Fig. 5. A spatio-temporal representation of a situation where a same voxel P is "seen" at different times in 2 different neighbouring tissues, due to the motion of the interface between tissues A and B, point P is in tissue A at time t_1 , and in tissue B at time t_2 . The attenuation at point P then quickly changes between t_1 and t_2 when the boundary reaches P.

the best of available data. In order to cope with such cases, we have developed an accurate and fast model. It exploits B-spline basis functions to represent the 3D object and to approximate their projections. This model will be designated henceforth under the name *spline driven* [26]. When reconstructing objects from a limited number of projections, we have shown *spline driven* to be more accurate than classical projectors such as *distance driven* [13], with a satisfactory computional time. In order to use dynamic data in a better way, we therefore used the *spline driven* projector to implement our dynamic reconstruction approach.

C. Spatio-temporal regularization

The term $\mathcal{J}_{\text{prior}}$ in equation (2) is essential to regularize the otherwise ill-conditioned inverse problem of dynamic reconstruction. In contrast to phase-correlated CT reconstructions that are performed independently, the regularization \mathcal{J}_{prior} can enforce a temporal coherence in the reconstructions. The spatial regularization must be chosen so as to favor smooth areas while preserving discontinuities at the boundary of anatomical structures. An edge-preserving smoothness prior such as TV [33] is typically chosen in this context. Having a look at the 3D representation of our 2D+t simulated phantom displayed on Fig. 2, we observe that the variations of shapes are exclusively noticeable at the interfaces between the different parts of the object. For such an object, the high amplitude spatial gradients give the tissues' boundaries. They are therefore intrinsically correlated to the temporal gradients from which moving boundaries can be traced back. Fig. 5 illustrates the case where a given spatial position "sees" 2 different tissues at time t_1 and time t_2 . The temporal profile of attenuation sketched in Fig. 5 corresponds to a large motion where the change of the intensity value at position P is sharp. On the contrary, for sub-voxel motions, the sharp boundary will be more progressive and thus will give a smoother transition. Cardiac and respiratory movements also cause smooth spatial and temporal changes because they induce fluctuations of the absorption, particularly at the interfaces between different tissues, *e.g.* the lung walls or the myocardium's tissues. Consequently, we advocate the use of an $\ell_2 - \ell_1$ edge preserving regularization [9], [14], [12], [44] which is known to favor smoothness for small amplitude changes while preserving larger discontinuities. Being concerned with spatio-temporal (dis)continuity, we propose to use the following generalized form of relaxed TV:

$$\begin{aligned} \mathcal{J}_{\text{prior}}(\boldsymbol{f}) &= \sum_{\boldsymbol{k},\ell} \left(\epsilon^2 + \mu_{\text{space}}^2 \left\| \nabla_{\boldsymbol{k},\ell}^{\text{space}} \cdot \boldsymbol{f} \right\|^2 \\ &+ \mu_{\text{time}}^2 \left(\partial_{\boldsymbol{k},\ell}^{\text{time}} \cdot \boldsymbol{f} \right)^2 \right)^{1/2}. \end{aligned}$$
(4)

Here operators $\nabla_{k,\ell}^{\text{space}}$ et $\partial_{k,\ell}^{\text{time}}$ are respectively the 3D spatial gradient and the temporal derivative of the 3D+*t* object at the voxel (k, ℓ) with k and ℓ the spatial index and temporal index (i.e., the frame) of the voxel. In order to tune the strength of the regularization (relatively to the data fidelity term) and to account for the heterogeneity of the physical dimensions of the spatial gradient and of the temporal derivative, we introduce two different hyperparameters $\mu_{\text{space}} \geq 0$ and $\mu_{\text{time}} \geq 0$ in Eq. (4). Note that the ratio of these two hyperparameters is homogeneous to a velocity. In Eq. (4), the parameter $\epsilon > 0$ is not to be mistaken with a relaxation parameter introduced to avoid a singularity: it has to be appropriately tuned to set the trade-off between sharp and smooth changes. Note that our 4D regularizer is similar to those used in [48], [32].

Having a convex differentiable criterion (2), we perform the numerical minimization with the VMLM algorithm [27] which is a limited memory quasi-Newton optimization method with BFGS updates.

IV. RESULTS

A. Numerical data

To test our approach, we reconstructed our dynamic Shepp-Logan phantom presented in section II-B, from a simulated set of 600 projections. We chose to sample the 2D+t sequence with 25 frames. In this configuration, the frames are exactly synchronized with the projections which, as explained below, lets us mimic a phase-correlated reconstruction. Fig. 6 displays some frames of two different reconstructions from this dataset:

- 1) The sequence in the second column was reconstructed while imposing only the spatial regularization for each frame. This was achieved by setting $\mu_{time} = 0$ in Eq. (4). Thanks to the perfect synchronisation of the frames with the projections, the result corresponds to an ideal spatially TV regularized *gated*-like reconstruction, very close to the PICCS approach [11]. The spatial hyperparameter $\mu_{space} > 0$ was tuned to achieve the best visual quality for the result.
- 2) The sequence in the third column was reconstructed using the spatio-temporal regularization in Eq. (4) with $\mu_{\text{time}} > 0$ and $\mu_{\text{space}} > 0$ both tuned to improve the visual quality of the result.

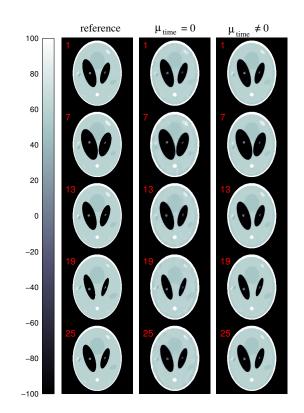


Fig. 6. 2D+*t* reconstructions (25 frames) of the dynamic Shepp-Logan phantom (values are given in Hounsfield units). 24 projections are available per frame of the reconstructed cycle, with which they are synchronized (perfect spatially TV regularized *gated*-like case). 5 frames have been selected on a period-6 base (numbers 1, 7, 13, 19, 25). The first column shows the reference images. The second column shows the reconstructed images without temporal regularization ($\mu_{time} = 0$), *i.e.* the frames are independently reconstructed with a spatial regularization. The third column shows the reconstruction with a 3D (2D+*t*) spatio-temporal regularization ($\mu_{time} \neq 0$).

To accurately compare the quality of the 2 reconstructions, Fig. 7 zooms into 2 regions of interest (ROI) centered at small and fine moving structures, which are very close from each other.

In the first case, each frame has been reconstructed independently from the others because there was no temporal correlation imposed between frames in this configuration. A set of only 24 projections is available per frame, which explains the rather poor quality of the reconstruction. This is particularly noticeable on the fine moving structures. Indeed we can observe on Fig. 7 a strong blurring effect due to the motion which makes the discrimination between the structures very difficult all along the cycle.

The second case, $\mu_{\text{time}} \neq 0$, shows the noticeable improvement brought by the temporal regularization of the 2D+*t* sequence in the quality of the reconstruction. The imposed continuity between frames helps to reduce artifacts due to the lack of available projections per frame. As a result, structures are better recovered, as well as their motion, in a non-ambiguous way. The gain is particularly clear on the ROIs in Fig. 7 where the fine structures are well resolved and can be discriminated without ambiguity at each frame. As a result their respective motion is very well recovered.

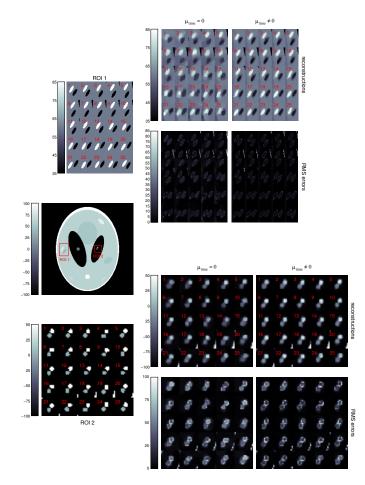


Fig. 7. Zoom into 2 regions of interest (ROI) of the 2D+*t* reconstructions of Fig. 6, showing small and fine moving structures of the phantom. Each of the 25 frames are shown for each ROI and compared to the truth with the root mean square error (RMS) map.

B. Mechanical phantom data

We have reconstructed a mechanical phantom from a dataset acquired at the *Centre Léon Bérard*, on a scanner *Elekta Synergy Cone-Beam CT*. The flat panel is composed of a grid of 512×512 detector pixels of size 0.08×0.08 cm². The acquisition process is similar to that described for CBCT scanners in section II-A.

The phantom has the features of a human thorax. Several zones of various densities mimic the lungs, the muscles and the spine. A small sphere of 2 cm in diameter is inserted in the right lung. It can be mechanically animated with a periodic circular motion in the transversal plane, and a periodic translation in the cranio-caudal direction. The two motions can be combined, reproducing the trajectory of a tumor during the respiratory cycle.

1) 2D+t reconstructions: A set of 630 projections has been acquired regularly in time and angle on 360° during 116 seconds. The spherical insert was animated with a periodical motion in the transversal plane. The period of this "respiratory" cycle was 4 seconds. Thus, about 29 cycles are done during the acquisition, corresponding to about 22 projections per cycle.

The phantom was positioned so that the isocenter of the

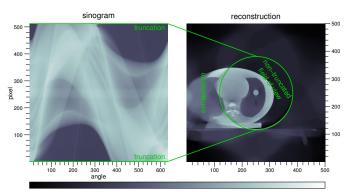


Fig. 8. 2D+*t* reconstruction of a mechanical phantom. *Left:* A sinogram extracted from the CB projections of the mechanical phantom. *Right:* The corresponding reconstruction. The disk's interior corresponds to the part of the object that is projected, hence "visible", on all the frames. A remarkable feature of our inverse approach is its ability to produce a rather satisfactory reconstruction outside the "visible" part where there are some missing data.

scanner is contained in the same plane as the trajectory of the spherical insert. Extracting the median detector lines corresponding to this plane from the CB projections, we obtained a fan beam sinogram from which we reconstructed a 2D+t sequence of the corresponding phantom's slice. A single temporal frame of this reconstruction is shown in Fig. 8. The projections are clearly truncated by the detector. With iterative reconstruction methods, this problem can be partially addressed by reconstructing a larger volume of interest in order to encompass the whole object, the regularization helping to extend the reconstruction to regions with missing data. The static reconstruction displayed by Fig. 8 shows that, despite the truncation of the field of view, the object support is correctly restored even though there are some artifacts and a loss of details outside the fully seen region. However note that these defects have no impact on the quality of the reconstruction inside the fully seen region.

Our dynamic 2D+t reconstructions are presented on Fig. 9. As for the dynamic Shepp-Logan phantom, we have performed two reconstructions: a plain spatially regularized one (Fig. 9(d)), and a fully spatial and time regularized one (Fig. 9(e)). We chose to reconstruct a 2D+t sequence of 22 frames. The operator $S^{t_{\theta}}$ of our direct model (see Section III-A) is implemented by a linear temporal interpolation. Hence, as the direct model links each object frame to two projections, about 29×2 projections per frame are available for the reconstruction. For comparison, Fig. 9 also displays static reconstructions from a varying number of projections with the mechanical insert at still (a-c). Thanks to the spatio-temporal regularization, the reconstructed frames of the dynamic reconstruction in Fig. 9(e) have almost the same visual quality as the reconstruction from 630 projections in Fig. 9(a). This visual quality is in neat improvement over the static reconstruction from 64 projections shown in Fig. 9(b) or from the gatedlike reconstruction in Fig. 9(d). The motion of the insert is correctly recovered and the anatomical structures are better reconstructed, in particular the spine.

2) 3D+t reconstructions: Fig. 10 displays a 3D+t reconstruction of the phantom from another set of CB projections,

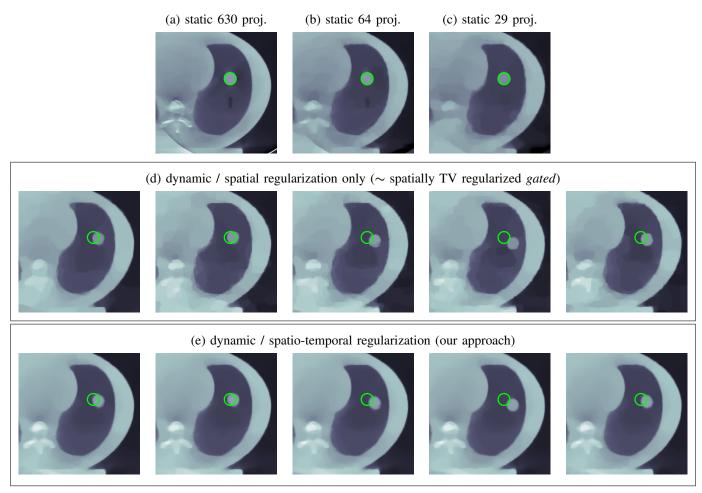


Fig. 9. 2D+t reconstructions of the median transversal slice of the mechanical phantom. Each restored sequence consists in 22 temporal frames with 500×500 pixels of size $1 \times 1 \text{ mm}^2$ (the region of interest is zoomed in). (a,b,c): Static reconstructions from respectively 630, 64 and 29 projections of the static phantom (*i.e.*, the insert was motionless). (d): Only spatially regularized reconstruction. (e): Spatio-temporally regularized reconstruction. The reference position of the spherical insert is identified on each frame by a small circle.

acquired in the same conditions as previously. This time, the spherical insert was periodically animated with two kinds of motion: circular in the transversal plane and translational in the cranio-caudal direction. Hence the insert had a realistic 3D movement. We reconstructed a sequence of 22 frames made of $225 \times 225 \times 75$ voxels of size $2 \times 2 \times 2 \text{ mm}^3$. Outside the field of view, the reconstruction appears corrupted by motion and truncation artifacts. However, inside the field of view, the quality of the reconstruction is comparable to that of our 2D+*t* reconstructions, thanks to the spatio-temporal regularization. To focus on the quality of the restored images, Fig. 10 only shows the non-truncated field of view. Transition from 2D+*t* to 3D+*t* involves no additional difficulties, but a bigger computional burden (many more variables, computations of 4D instead of 3D gradients in the regularization process).

C. Patient data

We have processed the same CBCT dataset of a patient's thorax as the one used in the article of Rit *et al.* [29]. The CB projections were acquired at the Centre Léon Bérard, (Lyon, France), on a scanner *Elekta Synergy Cone-Beam CT*. Thus the acquisition process is the same as for the mechanical phantom

in section IV-B. In [29], the reconstructions were performed with a MC method developed by the authors and with a *gated* method. These reconstructions have been made available to us, in order to make a comparison with our method. We also obtained the recorded 1D temporal signal giving the periodic linear phase of the respiratory cycles, from which we have calibrated the temporal interpolator $S^{t_{\theta}}$. The mean period of a cycle was evaluated to 2.4 s. We only kept the full cycles, reducing the number of required CB projections to 625. The duration of acquisition was 116 seconds, involving about 48 cycles, and 13 projections per cycle. We chose to reconstruct a 3D+t sequence of 13 frames. Using again linear interpolation for $S^{t_{\theta}}$, about 48×2 projections were available per reconstructed frame.

Figure 11 displays the full 3D+t reconstruction. Each frame consists in $275 \times 200 \times 135$ voxels of $2 \times 2 \times 2 \text{ mm}^3$. Inside the field of view, the thorax is correctly reconstructed, and one can identify the different anatomical structures: the right lung containing a tumor, its bronchi, and also the heart and the ribs. The respiratory motion is well recovered, particularly that of the diaphragm and of the tumor, without ambiguities nor noticeable artifacts.

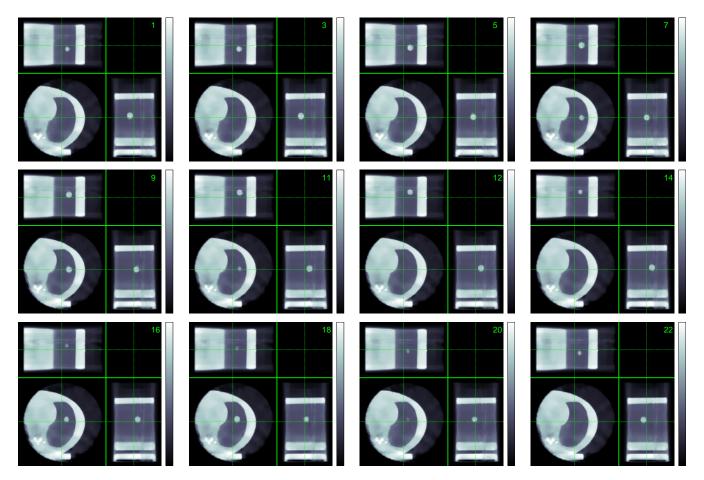


Fig. 10. 3D+t reconstruction of the mechanical phantom from the CB projections. A sliced view of each frame in the axial (center), sagittal (right), and coronal (top) directions is displayed. Only the non-truncated field of view is shown.

For comparison, Fig. 12 shows side by side the MC and *gated* reconstructions by Rit *et al.* in [29] and the result with our method. The end-exhale phase reconstructed by the MC method has a size of $262 \times 261 \times 132$, *i.e.* a voxel's size of $0.98 \times 0.98 \times 2 \text{ mm}^3$. The *gated* 3D+*t* sequence was reconstructed with the same parameters. We compare the reconstructed frame corresponding to the end-exhale phase. We have re-interpolated our reconstruction on the grid of finer voxels used by Rit *et al.*, using a cubic B-spline interpolation kernel.

Regarding the *gated* reconstruction, our method has eliminated the motion artifacts and better recovered the anatomical structures. Indeed *gated* methods have to be applied on extended datasets (*4D-CT*) to be efficient, which is not the case for CBCT and explains the poor quality of the *gated* reconstruction. Note that this *gated* is basic, and that it could be improved using spatial regularization of each phase. However it evidences the gain of our approach which achieves a good quality of reconstruction using only the CBCT dataset.

The MC reconstruction reveals a finer resolution of recovered structures than ours, which looks smoother. Note that this MC method requires an accurate motion model computed from a high resolution 4D-CT data set. Our reconstruction has only needed the current set of CB projections to reconstruct the patient's thorax and the respiratory motion, recovering the motion as well as the static structures without ambiguity or motion artifact. As a result, a lower effective X-ray dose is necessary with our method than with this particular MC method. Indeed the latter method does not show the best result that could be obtained with the MC approach, since efficient MC reconstructions can be performed on the single CBCT dataset [7], [8], [45], [49]. However these results also demonstrate the potential of the spatio-temporally regularized 4D phase-correlated reconstruction approach, which is able to recover a motion artifacts-free 3D+t sequence with very simple and general hypotheses.

V. CONCLUSION

This paper described an approach for phase-correlated 4D-CT reconstruction with a spatio-temporal regularization. We address the 4D reconstruction problem with very simple and general hypothesis. Our motivation was to perform 3D+t reconstructions from a limited number of projections, such as provided by a CBCT system, without resorting to additional 4D-CT data. Our results on CBCT data show that satisfying reconstructions can be obtained from limited data sets using a simple prior model based on spatial and temporal coherence.

Reconstruction of a 3D+t sequence from limited data is possible by considering jointly all projections in a global inverse problem. We carefully designed the model of pro-

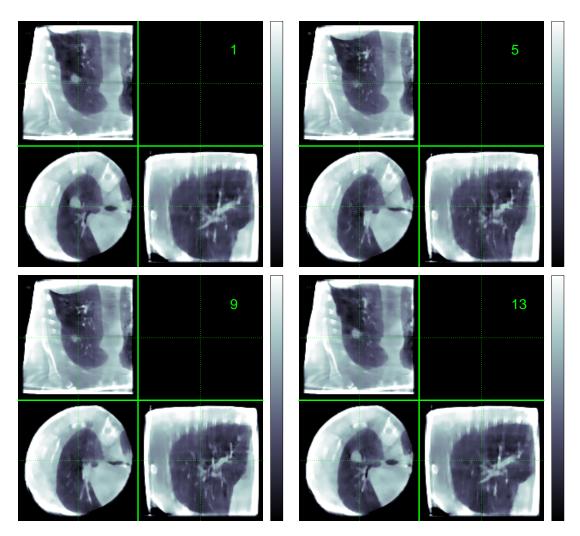
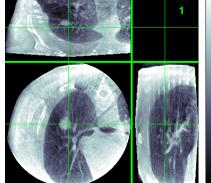


Fig. 11. Our 3D+t reconstruction of the thorax of a patient (*Elekta Synergy Cone-Beam CT* data from [29], used by courtesy of Centre Léon Bérard). A sliced view of each frame in the axial (center), sagittal (right), and coronal (top) directions is shown.

(a) Rit et al. gated (from a single scan)

(b) Rit et al. MC (from two scans)



(c) Our method (from a **single** scan)

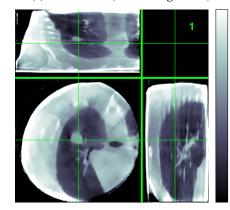


Fig. 12. Comparison of the reconstructions from the patient's dataset performed by a *gated method* (a) and the MC method of Rit *et al.* (b), with the reconstruction by our spatio-temporally regularized method (c). The end-exhale phase is shown. For better readability, only the fully seen field of view is displayed.

jections, based on spatial and temporal interpolation. Temporal frames are linearly interpolated to match phases of the cyclic movement under reconstruction. A joint spatio-temporal regularization exploits spatial and temporal coherence of the dynamic volume to provide images with smooth regions and motion. In order to preserve sharp boundaries and prevent from blurring motion, we proposed a relaxed 3D+t total-variation regularization.

Such spatio-temporally regularized phase-correlated 4D-CT reconstruction method has already been explored in the context of numerical 2D XCAT phantom CBCT-like data [48] and in vivo cardiac micro-CT mouse data [32]. We applied our method to 4D-CBCT reconstruction problems from a mechanical phantom and real patient data. Our results both confirmed that our method can produce 3D+t reconstructions using only CBCT data, with a significantly improved quality compared to pure and even spatially regularized gated reconstructions. The motion of a lung tumor can be clearly identified in our reconstruction, which is confirmed by comparison with a reconstruction obtained using a standard MC method based on an accurate estimation of motion from high-resolution 4D-CT images. The reconstruction quality does not match that achieved using accurate motion models but may be a simple and general alternative when high-resolution 4D-CT datasets are not available.

By avoiding the use of additional 4D-CT data in dynamic CBCT reconstructions, the approach we followed should reduce the effective X-ray dose required for a given 4D reconstruction.

ACKNOWLEDGMENTS

The empirical data were kindly provided by Simon Rit from the *Centre Léon Bérard* (Lyon, France). Our algorithms were developped in Yorick [http://yorick.sourceforge. net/index.php], a freely available data processing language provided by David Munro. Computations for the 3D+t reconstructions were carried out on the Horizon cluster [http: //www.projet-horizon.fr].

REFERENCES

- A. Andersen and A. Kak. Simultaneous Algebraic Reconstruction Technique (SART): a superior implementation of the ART algorithm. Ultrasonic Imaging, 1984. 2
- [2] J. Balter, R. Ten Haken, T. Lawrence, K. Lam and J. Robertson. Uncertainties in CT-based radiation therapy treatment planning associated with patient breathing. International Journal of Radiation Oncology Biology Physics, 1996. 1
- [3] R. Berbeco, H. Mostafavi, G. Sharp and S. Jiang. *Towards fluoroscopic respiratory gating for lung tumours without radiopaque markers*. Physics in Medicine and Biology, 2005. 1
- [4] C. Blondel, R. Vaillant, G. Malandain and N. Ayache. 4D deformation field of coronary arteries from monoplane rotational X-ray angiography. International Congress Series, 2003. 5
- [5] C. Blondel, R. Vaillant, G. Malandain and N. Ayache. 3D tomographic reconstruction of coronary arteries using a precomputed 4D motion field. Physics in Medicine and Biology, 2004. 2, 5
- [6] J. Brankov, Y. Yang, M. Narayanan and M. Wermck. *Motion-compensated* 4D processing of gated SPECT perfusion studies. IEEE Nuclear Science Symposium Conference Record, 2002. 2
- [7] M. Brehm, P. Paysan, M. Oelhafen, P. Kunz and M. Kachelriess. Selfadapting cyclic registration for motion-compensated cone-beam CT in image-guided radiation therapy. Medical Physics, 2012. 2, 3, 9

- [8] M. Brehm, P. Paysan, M. Oelhafen and M. Kachelriess. Artifact-resistant motion estimation with a patient-specific artifact model for motioncompensated cone-beam CT. Medical Physics, 2013. 2, 3, 9
- [9] P. Charbonnier, L. Blanc-Féraud, G. Aubert and M. Barlaud. *Deterministic edge-preserving regularization in computed imaging*. IEEE Transactions on Image Processing, 1997. 6
- [10] G. Chen, J. Kung and K. Beaudette. Artifacts in computed tomography scanning of moving objects. Seminars in Radiation Oncology, 2004. 1
- [11] G. Chen, J. Tang and S. Leng. Prior image constrained compressed sensing (PICCS): A method to accurately reconstruct dynamic CT images from highly undersampled projection data sets. Medical Physics, 2008. 2, 6
- [12] W. Chlewicki, F. Hermansen and S. Hansen. Noise reduction and convergence of Bayesian algorithms with blobs based on the Huber function and median root prior. Physics in Medicine and Biology, 2004.
- [13] B. DeMan and S. Basu. *Distance driven projection and backprojection in three dimensions.* Physics in Medicine and Biology, 2004. 5
- [14] I. Elbakri and J. Fessler. Statistical image reconstruction for polyenergetic X-ray computed tomography. EEE Transactions on Medical Imaging, 2002. 6
- [15] L. Feldkamp, L. Davis and J. Kress. Practical cone-beam algorithm. Journal of the Optical Society of America A, 1984. 2
- [16] E. Ford, G. Mageras, E. Yorke and C. Ling. Respiration-correlated spiral CT: a method of measuring respiratory-induced anatomic motion for radiation treatment planning. Medical Physics, 2003. 1
- [17] H. Goldberg, R. Gould, I. Feuerstein, J. Sigeti and M. Lipton. Evaluation of ultrafast CT scanning of the adult abdomen. Investigative Radiology, 1989. 1
- [18] D. Jaffray, D. Drake, M. Moreau, A. Martinez and J. Wong. A radiographic and tomographic imaging system integrated into a medical linear accelerator for localization of bone and soft-tissue targets. International Journal of Radiation Oncology Biology Physics, 1999. 2
- [19] D. Jaffray and J. Siewerdsen. Cone-beam computed tomography with a flat-panel imager: initial performance characterization. Medical Physics, 2000. 2
- [20] D.A. Jaffray, J.H. Siewerdsen, J.W. Wong and A.A. Martinez. Flat-panel cone-beam computed tomography for image-guided radiation therapy. International journal of Radiation Oncology Biology Physics, 2002. 2
- [21] A. Kak and M. Slaney. Principles of Computerized Tomographic Imaging. IEEE Press, 1988. 3
- [22] P. Keall, G. Starkschall, H. Shukla, K. Forster, V. Ortiz, C. Stevens, S. Vedam, R. George, T. Guerrero and R. Mohan. *Acquiring 4D thoracic CT scans using a multislice helical method*. Physics in Medicine and Biology, 2004. 1
- [23] T. Li, E. Schreibmann, Y. Yang and L. Xing. Motion correction for improved target localization with on-board cone-beam computed tomography. Physics in Medicine and Biology, 2006. 2
- [24] D. Low, M. Nystrom, E. Kalinin, P. Parikh, J. Dempsey, J. Bradley, S. Mutic, S. Wahab, T. Islam and G. Christensen. A method for the reconstruction of four-dimensional synchronized CT scans acquired during free breathing. Medical Physics, 2003. 1
- [25] F. Momey, L. Denis, C. Burnier, É. Thiébaut, J.M. Becker and L. Desbat. A new representation and projection model for tomography, based on separable B-splines. IEEE Nuclear Science Symposium and Medical Imaging Conference, 2011 [https://hal-ujm.archives-ouvertes.fr/ujm-00670025]. 5
- [26] F. Momey, L. Denis, C. Burnier, É. Thiébaut, J.M. Becker and L. Desbat. Spline driven: high accuracy projectors for tomographic reconstruction. Manuscript under peer review for IEEE Transactions on Image Processing, 2013 [https://hal.archives-ouvertes.fr/hal-00990015]. 5
- [27] J. Nocedal. Updating quasi-Newton matrices with limited storage. Mathematics of computation, 1980. 6
- [28] E. Rietzel, T. Pan and G. Chen. Four-dimensional computed tomography: image formation and clinical protocol. Medical Physics, 2005. 1
- [29] S. Rit, D. Sarrut and L. Desbat. Comparison of analytic and algebraic methods for motion-compensated cone-beam CT reconstruction of the thorax. IEEE Transactions on Medical Imaging, 2009. 2, 8, 9, 10
- [30] S. Rit, J. Wolthaus, M. Van Herk and J. Sonke. On-the-fly motioncompensated cone-beam CT using an a priori model of the respiratory motion. Medical Physics, 2009. 2
- [31] C. Ritchie, J. Godwin, C. Crawford, W. Stanford, H. Anno, and K. Yongmin. *Minimum scan speeds for suppression of motion artifacts* in CT. Radiology, Radiological Society of North America, 1992. 1
- [32] L. Ritschl, S. Sawall, M. Knaup, A. Hess and M. Kachelriess. Iterative 4D cardiac micro-CT image reconstruction using an adaptive spatiotemporal sparsity prior. Physics in Medicine and Biology, 2012. 2, 6, 11

- [33] L.I. Rudin, S. Osher and E. Fatemi. Nonlinear total variation based noise removal algorithms. Physica D: Nonlinear Phenomena, 1992. 2, 5
- [34] W. Segars, D. Lalush and B. Tsui. A realistic spline-based dynamic heart phantom. IEEE Transactions on Nuclear Science, 1999.
- [35] W. Segars, D. Lalush and B. Tsui. Modeling respiratory mechanics in the MCAT and spline-based MCAT phantoms. IEEE Transactions on Nuclear Science, 2001.
- [36] W. Segars, D. Lalush and B. Tsui. 4D XCAT phantom for multimodality imaging research. Medical physics, 2010.
- [37] Y. Seppenwoolde, H. Shirato, K. Kitamura, S. Shimizu, M. Van Herk, J. Lebesque and K. Miyasaka. *Precise and real-time measurement of* 3D tumor motion in lung due to breathing and heartbeat, measured during radiotherapy. International Journal of Radiation Oncology Biology Physics, 2002. 1
- [38] L. Shepp and B. Logan. *The Fourier reconstruction of a head section*. IEEE Transactions on Nuclear Science, 1974. 3
- [39] S. Shimizu, H. Shirato, K. Hagei, T. Nishioka, X. Bo, H. Dosaka-Akita, S. Hashimoto, H. Aoyama, K. Tsuchiya and K. Miyasaka. *Impact of respiratory movement on the computed tomographic images of small lung tumors in three-dimensional (3D) radiotherapy*. International Journal of Radiation Oncology Biology Physics, 2000. 1
- [40] J.J. Sonke, L. Zijp, P. Remeijer and M. Van Herk. *Respiratory correlated cone beam CT*. Medical Physics, 2005. 1
- [41] A. Tarantola. Inverse Problem Theory and Methods for Model Parameter Estimation. SIAM Pub., 2005. 3
- [42] Z. Tian, X. Jia, B. Dong, Y. Lou and S.B. Jiang. Low-dose 4DCT reconstruction via temporal nonlocal means. Medical Physics, 2011. 2
- [43] S. Vedam, P. Keall, V. Kini, H. Mostafavi, H. Shukla and R. Mohan. Acquiring a four-dimensional computed tomography dataset using an external respiratory signal. Physics in Medicine and Biology, 2003. 1
- [44] J. Wang, T. Li and L. Xing. Iterative image reconstruction for CBCT using edge-preserving prior. Medical Physics, 2009. 6
- [45] J. Wang and X. Gu. Simultaneous motion estimation and image reconstruction (SMEIR) for 4D cone-beam CT. Medical Physics, 2013. 2, 3, 9
- [46] P. Willis and Y. Bresler. Optimal scan for time-varying tomography. I. Theoretical analysis and fundamental limitations. IEEE Transactions on Image Processing, 1995. 1
- [47] P. Willis and Y. Bresler. Optimal scan for time-varying tomography. II. Efficient design and experimental validation. IEEE Transactions on Image Processing, 1995. 1
- [48] H. Waibo, A. Maier, R. Fahrig and J. Hornegger. Spatial-temporal Total Variation Regularization (STTVR) for 4D-CT reconstruction. Proceedings of SPIE, 2012. 2, 6, 11
- [49] H. Yan, X. Zhen, M. Folkerts, Y. Li, T. Pan, L. Cervino, S.B. Jiang and X. Jia. A hybrid reconstruction algorithm for fast and accurate 4D cone-beam CT imaging. Medical Physics, 2014. 2, 3, 9
- [50] R. Zeng, J. Fessler and J. Balter. *Estimating 3-D respiratory motion from orbiting views by tomographic image registration*. IEEE Transactions on Medical Imaging, 2007. 2