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Effectivized Hölder-logarithmic stability estimates for the Gel'fand inverse problem

M.I. Isaev and R.G. Novikov

Abstract

We give effectivized Hölder-logarithmic energy and regularity dependent stability estimates for the Gel'fand inverse boundary value problem in dimension $d = 3$. This effectivization includes explicit dependence of the estimates on coefficient norms and related parameters. Our new estimates are given in L^2 and L^∞ norms for the coefficient difference and related stability efficiently increases with increasing energy and/or coefficient difference regularity. Comparisons with preceding results are given.

1 Introduction and main results

We consider the equation

$$-\Delta\psi + v(x)\psi = E\psi, \quad x \in D \subset \mathbb{R}^3, \quad (1.1)$$

where

$$D \text{ is an open bounded domain in } \mathbb{R}^3, \quad \partial D \in C^2, \quad (1.2)$$

$$v \in L^\infty(D). \quad (1.3)$$

Equation (1.1) can be regarded as the stationary Schrödinger equation of quantum mechanics at fixed energy E . Equation (1.1) at fixed E arises also in acoustics and electrodynamics.

As in Section 5 of Gel'fand's work [9] we consider an operator establishing a relationship between ψ and $\partial\psi/\partial\nu$ on ∂D for all sufficiently regular solutions ψ of equation (1.1) in $\bar{D} = D \cup \partial D$ at fixed E , where ν is the outward normal to ∂D . As in [26], [16] (for example) we represent such an operator as the Dirichlet-to-Neumann map $\hat{\Phi}(E)$ defined by the relation

$$\hat{\Phi}(E)(\psi|_{\partial D}) = \frac{\partial\psi}{\partial\nu}|_{\partial D}, \quad (1.4)$$

where we assume also that

$$E \text{ is not a Dirichlet eigenvalue for operator } -\Delta + v \text{ in } D. \quad (1.5)$$

The map $\hat{\Phi} = \hat{\Phi}(E)$ can be regarded as all possible boundary measurements for the physical model described by equation (1.1) at fixed energy E under assumption (1.5).

We consider the following inverse boundary value problem for equation (1.1):

Problem 1.1. Given $\hat{\Phi}$ for some fixed E , find v .

This problem is known as the Gel'fand inverse boundary value problem for the Schrödinger equation at fixed energy E in three dimensions (see [9], [26]). For $E = 0$ this problem can be regarded also as a generalization of the Calderón problem of the electrical impedance tomography in three dimensions (see [5], [26]). Problem 1.1 can be also considered as an example of ill-posed problem; see [4], [23] for an introduction to this theory.

Let, for real $m \geq 0$,

$$\begin{aligned} H^m(\mathbb{R}^3) &= \{w \in L^2(\mathbb{R}^3) : \mathcal{F}^{-1}(1 + |\xi|^2)^{\frac{m}{2}} \mathcal{F}w \in L^2(\mathbb{R}^3)\}, \\ \|w\|_{H^m(\mathbb{R}^3)} &= \|\mathcal{F}^{-1}(1 + |\xi|^2)^{\frac{m}{2}} \mathcal{F}w\|_{L^2(\mathbb{R}^3)}, \end{aligned} \quad (1.6)$$

where \mathcal{F} denote the Fourier transform

$$\mathcal{F}w(\xi) = \frac{1}{(2\pi)^3} \int_{\mathbb{R}^3} e^{i\xi x} w(x) dx, \quad \xi \in \mathbb{R}^3.$$

In addition, for real $m \geq 0$, we consider the spaces $W^m(\mathbb{R}^3)$ defined by

$$\begin{aligned} W^m(\mathbb{R}^3) &= \{w \in L^1(\mathbb{R}^3) : (1 + |\xi|^2)^{\frac{m}{2}} \mathcal{F}w \in L^\infty(\mathbb{R}^3)\}, \\ \|w\|_{W^m(\mathbb{R}^3)} &= \|(1 + |\xi|^2)^{\frac{m}{2}} \mathcal{F}w\|_{L^\infty(\mathbb{R}^3)}. \end{aligned} \quad (1.7)$$

We note that for integer m the space $W^m(\mathbb{R}^3)$ contains the standard Sobolev space $W^{m,1}(\mathbb{R}^3)$ of m -times smooth functions in L^1 on \mathbb{R}^3 .

In the present work we obtain, in particular, the following theorems:

Theorem 1.1. *Suppose that D satisfies (1.2) and v_1, v_2 satisfy (1.3), (1.5) for some real E . Suppose also that: $\|v_j\|_{L^\infty(D)} \leq N$ for some $N > 0$, $j = 1, 2$; $\text{supp}(v_2 - v_1) \subset D$, $v_2 - v_1 \in H^m(\mathbb{R}^3)$, $\|v_2 - v_1\|_{H^m(\mathbb{R}^3)} \leq N_{H^m}$ for some $m > 0$ and $N_{H^m} > 0$. Let*

$$\delta = \|\hat{\Phi}_2(E) - \hat{\Phi}_1(E)\|_{L^\infty(\partial D) \rightarrow L^\infty(\partial D)}, \quad (1.8)$$

where $\hat{\Phi}_1(E)$, $\hat{\Phi}_2(E)$ denote the Dirichlet-to-Neumann maps for v_1 , v_2 , respectively. Then, there exist some positive constants A, B, α, β depending on D only such that

$$\begin{aligned} \|v_2 - v_1\|_{L^2(D)} &\leq A \left(\alpha E + \beta(1 - \tau)^2 (\ln(3 + \delta^{-1}))^2 \right)^{\frac{1}{2}} \delta^\tau + \\ &+ B(1 + N)^{\frac{4m}{3}} N_{H^m} \left(\alpha E + \beta(1 - \tau)^2 (\ln(3 + \delta^{-1}))^2 \right)^{-\frac{m}{3}} \end{aligned} \quad (1.9)$$

for any $\tau \in (0, 1]$ and $E \geq 0$.

Besides, estimate (1.9) is also fulfilled for any $\tau \in (0, 1)$ and $E < 0$ under the following additional condition:

$$\alpha E + \beta(1 - \tau)^2 (\ln(3 + \delta^{-1}))^2 > 0. \quad (1.10)$$

Theorem 1.2. *Suppose that D satisfies (1.2) and v_1, v_2 satisfy (1.3), (1.5) for some real E . Suppose also that: $\|v_j\|_{L^\infty(D)} \leq N$ for some $N > 0$, $j = 1, 2$; $\text{supp}(v_2 - v_1) \subset D$, $v_2 - v_1 \in W^m(\mathbb{R}^3)$, $\|v_2 - v_1\|_{W^m(\mathbb{R}^3)} \leq N_{W^m}$ for some*

$m > 3$ and $N_{W^m} > 0$. Let δ be defined by (1.8). Then, there exist some positive constants $\tilde{A}, \tilde{B}, \tilde{\alpha}, \tilde{\beta}$ depending on D only such that

$$\begin{aligned} \|v_2 - v_1\|_{L^\infty(D)} &\leq \tilde{A} \left(\tilde{\alpha}E + \tilde{\beta}(1 - \tau)^2 (\ln(3 + \delta^{-1}))^2 \right)^{\frac{1}{2}} \delta^\tau + \\ &+ \tilde{B} \frac{(1 + N)^{\frac{2(m-3)}{3}} N_{W^m}}{m-3} \left(\tilde{\alpha}E + \tilde{\beta}(1 - \tau)^2 (\ln(3 + \delta^{-1}))^2 \right)^{-\frac{m-3}{6}} \end{aligned} \quad (1.11)$$

for any $\tau \in (0, 1]$ and $E \geq 0$.

Besides, estimate (1.11) is also fulfilled for any $\tau \in (0, 1)$ and $E < 0$ under the following additional condition

$$\tilde{\alpha}E + \tilde{\beta}(1 - \tau)^2 (\ln(3 + \delta^{-1}))^2 > 0. \quad (1.12)$$

Theorems 1.1 and 1.2 are proved in Sections 3 and 4, respectively. These proofs are based on Lemmas 2.1, 2.2 and 2.3 given in Section 2. Then these proofs are based on the intermediate estimates (3.7), (4.8) which may be of independent interest.

Remark 1.1. The estimates of Theorem 1.2 can be regarded as a significant effectivization of the following estimates of [16] for the three-dimensional case:

$$\|v_2 - v_1\|_{L^\infty(D)} \leq C_1(N_m, D, m, E) (\ln(3 + \delta^{-1}))^{-s_1} \quad (1.13)$$

for $E \in \mathbb{R}$;

$$\begin{aligned} \|v_2 - v_1\|_{L^\infty(D)} &\leq C_2(N_m, D, m, \tau)(1 + \sqrt{E})\delta^\tau + \\ &+ C_3(N_m, D, m, \tau)(1 + \sqrt{E})^{s-s_1} (\ln(3 + \delta^{-1}))^{-s} \end{aligned} \quad (1.14)$$

for $E \geq 0$, $\tau \in (0, 1)$ and any $s \in [0, s_1]$. Here δ is defined by (1.8) and $s_1 = (m - 3)/3$.

In addition, estimates (1.13) and (1.14) were obtained in [16] under the assumptions that: D satisfies (1.2), v_j satisfies (1.3), (1.5), $\text{supp } v_j \subset D$, $v_j \in W^{m,1}(\mathbb{R}^3)$, $\|v_j\|_{W^{m,1}(\mathbb{R}^3)} \leq N_m$, $j = 1, 2$, for some integer $m > 3$ and $N_m > 0$.

Actually, Theorem 1.2 was obtained in the framework of finding the dependence of C_1, C_2, C_3 of (1.13), (1.14) on N_m, m and τ . One can see that the estimates of Theorem 1.2 depend explicitly on coefficient norms N, N_{W^m} and parameters m, τ and imply (1.13), (1.14) with some C_1, C_2, C_3 explicitly dependent on N_m, m, τ as a corollary. Besides, in Theorem 1.2 we do not assume that each of potentials v_1, v_2 is m -times differentiable and is supported in D (in a similar way with Theorem 2.1 of [34]).

By the way we would like to note also that even for $E = 0$ the reduction of Hölder-logarithmic stability estimates like (1.9), (1.11) to pure logarithmic estimates like (1.13) is not optimal for large m because of the following asymptotic formula:

$$\sup_{\delta \in (0, 1]} \left| \frac{\delta^\tau}{(\ln(3 + \delta^{-1}))^{-\mu}} \right| = O\left(\left(\frac{\mu}{\tau}\right)^\mu e^{-\frac{\mu}{\tau}}\right) \quad \text{as } \mu \rightarrow +\infty.$$

In particular, even for $E = 0$ the Hölder-logarithmic estimates (1.9), (1.11) are much more informative than their possible pure logarithmic reductions.

Remark 1.2. Theorem 1.1 was obtained as an extension of Theorem 1.2 to the L^2 -norm case. In addition, it is important to note that the second (“logarithmic”) term of the right-hand side of (1.9) is considerably better than the analogous term of (1.11). In particular,

$$\begin{aligned} R &= O\left(\left(\ln(3 + \delta^{-1})\right)^{-\frac{2m}{3}}\right) \quad \text{for } \delta \rightarrow 0, \\ R &= O\left(E^{-\frac{m}{3}}\right) \quad \text{for } E \rightarrow +\infty, \end{aligned}$$

whereas

$$\begin{aligned} \tilde{R} &= O\left(\left(\ln(3 + \delta^{-1})\right)^{-\frac{m-3}{3}}\right) \quad \text{for } \delta \rightarrow 0, \\ \tilde{R} &= O\left(E^{-\frac{m-3}{6}}\right) \quad \text{for } E \rightarrow +\infty, \end{aligned}$$

where R and \tilde{R} denote the second (“logarithmic”) terms of the right-hand sides of (1.9) and (1.11), respectively.

Remark 1.3. The estimates of Theorem 1.1 should be compared also with the following estimate of [21] for the three-dimensional case:

$$\|v_2 - v_1\|_{H^{-m}(\mathbb{R}^3)} \leq C \left(E^2 \delta + \left(\sqrt{E} + \ln \delta^{-1} \right)^{-(2m-3)} \right), \quad (1.15)$$

where $C = C(N_m, D, \text{supp}(v_2 - v_1), m) > 0$, $\|v_j\|_{H^m(D)} \leq N_m$ ($j = 1, 2$), $\text{supp}(v_2 - v_1) \subset D$, $m > 3/2$, δ is the distance between the boundary measurements (Cauchy data) for v_1, v_2 and is, roughly speaking, similar to δ of (1.8) and where $\delta \leq 1/e$.

A principal advantage of (1.9) in comparison with (1.15) consists in estimation $v_2 - v_1$ in the L^2 -norm instead of the H^{-m} -norm. Besides, estimate (1.9) depends explicitly on coefficient norms N, N_{W^m} and parameters m, τ in contrast with (1.15). In addition, in (1.9) we do not assume that each of v_1, v_2 belongs to H^m .

Remark 1.4. In the literature on Problem 1.1 estimates of the form (1.13) are known as global logarithmic stability estimates. The history of these estimates goes back to [1] for the case when $s_1 \leq 1$ and to [33] for the case when $s_1 > 1$.

In addition, estimates of the form (1.9), (1.11), (1.14), (1.15) are known in the literature as Hölder-logarithmic energy and regularity dependent stability estimates. For the case when $\tau = 1$ in (1.9), (1.11) or when $s = 0$ in (1.14) the history of such estimates in dimension $d = 3$ goes back to [29], [31], where such energy and regularity dependent rapidly convergent approximate stability estimates were given for the inverse scattering problem.

Then for Problem 1.1 energy dependent stability estimates changing from logarithmic type to Hölder type for high energies were given in [20]. However, this high energy stability increasing of [20] is slow. The studies of [29], [31], [33], [20] were continued, in particular, in [25], [16], [21] and in the present work.

Remark 1.5. In Theorems 1.1, 1.2 we consider the three-dimensional case for simplicity only. Similar results hold in dimension $d > 3$.

As regards to logarithmic and Hölder-logarithmic stability estimates for Problem 1.1 in dimension $d = 2$, we refer to [35], [37], [38]. In addition, for problems like Problem 1.1 the history of energy and regularity dependent rapidly convergent approximate stability estimates in dimension $d = 2$ goes back to [28].

Remark 1.6. In a similar way with results of [17], [18] and subsequent studies of [36], estimates (1.9), (1.11) can be extended to the case when we do not assume that condition (1.5) is fulfilled and consider an appropriate impedance boundary map (Robin-to-Robin map) instead of the Dirichlet-to-Neumann map.

Remark 1.7. Apparently, estimates analogous to estimates of Theorems 1.1 and 1.2 hold if we replace the difference of DtN maps by the difference of corresponding near field scattering data in a similar way with results of [10], [14], [19].

Remark 1.8. The optimality (in different senses) of estimates like (1.13), (1.14) was proved in [24], [12], [13]. See also [6], [15] and references therein for the case of inverse scattering problems.

Remark 1.9. Estimates (1.9), (1.11) for $\tau = 1$ are roughly speaking coherent with stability properties of the approximate monochromatic inverse scattering reconstruction of [29], [31], implemented numerically in [2]. Estimates (1.9), (1.11) for $E = 0$ are roughly speaking coherent with stability properties of the reconstruction of [32].

In addition, estimates (1.9), (1.11) can be used for the convergence rate analysis for iterative regularized reconstructions for Problem 1.1 in the framework of an effectivization of the approach of [10] for monochromatic inverse scattering problems.

2 Lemmas

Let \hat{v} denote the Fourier transform of v :

$$\hat{v}(\xi) = \mathcal{F}v(\xi) = \frac{1}{(2\pi)^3} \int_{\mathbb{R}^3} e^{i\xi x} v(x) dx, \quad \xi \in \mathbb{R}^3. \quad (2.1)$$

Lemma 2.1. *Suppose that D satisfies (1.2) and v_1, v_2 satisfy (1.3), (1.5) for some real E . Suppose also that $\|v_j\|_{L^\infty(D)} \leq N$, $j = 1, 2$, for some $N > 0$. Let δ be defined by (1.8). Then*

$$|\hat{v}_2(\xi) - \hat{v}_1(\xi)| \leq c_1(1 + N)^2 \left(e^{2\rho L} \delta + \frac{\|v_1 - v_2\|_{L^2(D)}}{\sqrt{E + \rho^2}} \right) \quad (2.2)$$

for any $\rho > 0$ such that

$$|\xi| \leq 2\sqrt{E + \rho^2}, \quad E + \rho^2 \geq (1 + N)^2 r_1^2,$$

where $L = \max_{x \in \partial D} |x|$ and constants $c_1, r_1 > 0$ depend on D only.

Some version of estimate (2.2) was given in [16] (see formula (4.13) of [16]). Lemma 2.1 is proved in Section 6. This proof is based on results presented in Section 5.

Lemma 2.2. *Let $w \in H^m(\mathbb{R}^3)$, $\|w\|_{H^m(\mathbb{R}^3)} \leq N_{H^m}$ for some real $m > 0$ and $N_{H^m} > 0$, where the space $H^m(\mathbb{R}^3)$ is defined in (1.6). Then, for any $r > 0$,*

$$\left(\int_{|\xi| \geq r} |\mathcal{F}w(\xi)|^2 d\xi \right)^{1/2} \leq c_2 N_{H^m} r^{-m}, \quad (2.3)$$

where $\mathcal{F}w$ is defined according to (2.1) and $c_2 = (2\pi)^{-3/2}$.

Proof of Lemma 2.2. Note that

$$\int_{|\xi| \geq r} |\mathcal{F}w(\xi)|^2 d\xi \leq \left\| \frac{(1 + |\xi|^2)^{\frac{m}{2}} \mathcal{F}w}{r^m} \right\|_{L^2(\mathbb{R}^3)}^2. \quad (2.4)$$

Using (1.6), (2.4) and the Parseval theorem

$$\|\mathcal{F}\tilde{w}\|_{L^2(\mathbb{R}^3)} = (2\pi)^{-3/2} \|\tilde{w}\|_{L^2(\mathbb{R}^3)} \quad (2.5)$$

for $\tilde{w} \equiv \mathcal{F}^{-1}(1 + |\xi|^2)^{\frac{m}{2}} \mathcal{F}w$, we get estimate (2.3). ■

Lemma 2.3. *Let $w \in W^m(\mathbb{R}^3)$, $\|w\|_{W^m(\mathbb{R}^3)} \leq N_{W^m}$ for some real $m > 3$ and $N_{W^m} > 0$, where the space $W^m(\mathbb{R}^3)$ is defined in (1.7). Then, for any $r > 0$,*

$$\int_{|\xi| \geq r} |\mathcal{F}w(\xi)| d\xi \leq \tilde{c}_2 \frac{N_{W^m}}{m-3} r^{3-m}, \quad (2.6)$$

where $\mathcal{F}w$ is defined according to (2.1) and $\tilde{c}_2 = 4\pi$.

Proof of Lemma 2.3. Note that

$$r^m |\mathcal{F}w(\xi)| \leq (1 + |\xi|^2)^{m/2} |\mathcal{F}w(\xi)| \leq N_{W^m} \quad \text{for } |\xi| \geq r. \quad (2.7)$$

Using (2.7), we obtain that

$$\int_{|\xi| \geq r} |\mathcal{F}w(\xi)| d\xi \leq \int_r^{+\infty} \frac{N_{W^m}}{t^m} 4\pi t^2 dt \leq \frac{4\pi N_{W^m}}{m-3} r^{3-m}. \quad (2.8)$$

■

3 Proof of Theorem 1.1

Using the Parseval formula (2.5), we get that

$$\|v_2 - v_1\|_{L^2(D)} = (2\pi)^{3/2} \|\hat{v}_2 - \hat{v}_1\|_{L^2(\mathbb{R}^3)} \leq (2\pi)^{3/2} (I_1(r) + I_2(r)), \quad (3.1)$$

for $r > 0$, where \hat{v}_j is defined according to (2.1) with $v_j \equiv 0$ on $\mathbb{R}^3 \setminus D$, $j = 1, 2$,

$$I_1(r) = \left(\int_{|\xi| \leq r} |\hat{v}_2(\xi) - \hat{v}_1(\xi)|^2 d\xi \right)^{1/2},$$

$$I_2(r) = \left(\int_{|\xi| \geq r} |\hat{v}_2(\xi) - \hat{v}_1(\xi)|^2 d\xi \right)^{1/2}.$$

Let

$$r = q(1+N)^{-4/3}(E+\rho^2)^{1/3}, \quad q = \frac{1}{2\pi} \left(\frac{16\pi c_1^2}{3} \right)^{-1/3}, \quad (3.2)$$

where c_1 is the constant of Lemma 2.1.

Then, using Lemma 2.1 for $|\xi| \leq r$, we get that

$$\begin{aligned} I_1(r) &\leq \left(\frac{4\pi r^3}{3} c_1^2 (1+N)^4 \left(e^{2\rho L} \delta + \frac{\|v_1 - v_2\|_{L^2(D)}}{\sqrt{E+\rho^2}} \right)^2 \right)^{1/2} \leq \\ &\leq (2\pi)^{-3/2} \left(\frac{\sqrt{E+\rho^2} e^{2\rho L} \delta}{2} + \frac{\|v_1 - v_2\|_{L^2(D)}}{2} \right) \end{aligned} \quad (3.3)$$

for $q(1+N)^{-4/3}(E+\rho^2)^{1/3} \leq 2\sqrt{E+\rho^2}$ and $E+\rho^2 \geq (1+N)^2 r_1^2$.

In addition, using (2.3), we have that

$$I_2(r) \leq c_2 N_{H^m} r^{-m}. \quad (3.4)$$

Let $r_2 = r_2(D) \geq r_1$ be such that

$$E + \rho^2 \geq r_2^2 \implies q(E + \rho^2)^{1/3} \leq 2\sqrt{E + \rho^2}. \quad (3.5)$$

Using (3.1), (3.3)–(3.5) with r defined in (3.2), we obtain that

$$\begin{aligned} \|v_2 - v_1\|_{L^2(D)} &\leq \frac{\sqrt{E+\rho^2} e^{2\rho L} \delta}{2} + \frac{\|v_1 - v_2\|_{L^2(D)}}{2} + \\ &+ (2\pi)^{3/2} c_2 \frac{(1+N)^{\frac{4m}{3}}}{q^m} N_{H^m} (E + \rho^2)^{-\frac{m}{3}}, \end{aligned} \quad (3.6)$$

$$\begin{aligned} \frac{1}{2} \|v_2 - v_1\|_{L^2(D)} &\leq \frac{\sqrt{E+\rho^2} e^{2\rho L} \delta}{2} + \\ &+ \frac{(1+N)^{\frac{4m}{3}}}{q^m} N_{H^m} (E + \rho^2)^{-\frac{m}{3}}, \end{aligned} \quad (3.7)$$

for $E + \rho^2 \geq (1+N)^2 r_2^2$, where L, c_2 are the constants of Lemmas 2.1, 2.2 and q, r_2 are the constants of formulas (3.2), (3.5).

Let $\tau \in (0, 1)$ and

$$\gamma = \frac{1-\tau}{2L}, \quad \rho = \gamma \ln(3 + \delta^{-1}). \quad (3.8)$$

Due to (3.7), for δ such that

$$E + (\gamma \ln(3 + \delta^{-1}))^2 \geq (1+N)^2 r_2^2, \quad (3.9)$$

the following estimate holds:

$$\begin{aligned} \frac{1}{2} \|v_1 - v_2\|_{L^2(D)} &\leq \\ &\leq \frac{1}{2} \left(E + (\gamma \ln(3 + \delta^{-1}))^2 \right)^{1/2} (3 + \delta^{-1})^{2\gamma L} \delta + \\ &+ \frac{(1+N)^{\frac{4m}{3}}}{q^m} N_{H^m} \left(E + (\gamma \ln(3 + \delta^{-1}))^2 \right)^{-\frac{m}{3}}, \end{aligned} \quad (3.10)$$

where γ is defined in (3.8). Note that

$$(3 + \delta^{-1})^{2\gamma L} \delta = (1 + 3\delta)^{1-\tau} \delta^\tau \leq 4\delta^\tau \quad \text{for } \delta \leq 1. \quad (3.11)$$

Combining (3.10), (3.11), we get that

$$\begin{aligned} \|v_2 - v_1\|_{L^2(D)} &\leq A_1 \left(\lambda \left(E + \gamma^2 (\ln(3 + \delta^{-1}))^2 \right) \right)^{\frac{1}{2}} \delta^\tau + \\ &+ B_1 (1 + N)^{\frac{4m}{3}} N_{H^m} \left(\lambda \left(E + \gamma^2 (\ln(3 + \delta^{-1}))^2 \right) \right)^{-\frac{m}{3}} \end{aligned} \quad (3.12)$$

for $\delta \leq 1$ satisfying (3.9) and some positive constants A_1, B_1, λ depending on D only.

In view of definition (1.6), we have that

$$\|v_2 - v_1\|_{L^2(D)} \leq \|v_2 - v_1\|_{H^m(\mathbb{R}^3)} \leq N_{H^m}.$$

Hence, we get that, for $0 < E + (\gamma \ln(3 + \delta^{-1}))^2 \leq (1 + N)^2 r_2^2$,

$$\|v_2 - v_1\|_{L^2(D)} \leq (1 + N)^{\frac{4m}{3}} N_{H^m} \left(\frac{E + \gamma^2 (\ln(3 + \delta^{-1}))^2}{r_2^2} \right)^{-\frac{m}{3}}. \quad (3.13)$$

On other hand, in the case when $E + (\gamma \ln(3 + \delta^{-1}))^2 \geq (1 + N)^2 r_2^2$ and $\delta > 1$ we have that

$$\begin{aligned} \|v_2 - v_1\|_{L^2(D)} &\leq c_3 \|v_2 - v_1\|_{L^\infty(D)} \leq c_3 2N \leq \\ &\leq 2c_3 \left(\frac{E + \gamma^2 (\ln(3 + \delta^{-1}))^2}{r_2^2} \right)^{\frac{1}{2}} \delta^\tau, \end{aligned} \quad (3.14)$$

where

$$c_3 = \left(\int_D 1 \, dx \right)^{1/2}. \quad (3.15)$$

Combining (3.8), (3.12)–(3.14), we obtain estimate (1.9). This completes the proof of Theorem 1.1.

4 Proof of Theorem 1.2

Due to the inverse Fourier transform formula

$$v(x) = \int_{\mathbb{R}^3} e^{-i\xi x} \hat{v}(\xi) d\xi, \quad x \in \mathbb{R}^3, \quad (4.1)$$

we have that

$$\|v_1 - v_2\|_{L^\infty(D)} \leq \sup_{x \in D} \left| \int_{\mathbb{R}^3} e^{-i\xi x} (\hat{v}_2(\xi) - \hat{v}_1(\xi)) d\xi \right| \leq \tilde{I}_1(r) + \tilde{I}_2(r) \quad (4.2)$$

for $r > 0$, where

$$\begin{aligned}\tilde{I}_1(r) &= \int_{|\xi| \leq r} |\hat{v}_2(\xi) - \hat{v}_1(\xi)| d\xi, \\ \tilde{I}_2(r) &= \int_{|\xi| \geq r} |\hat{v}_2(\xi) - \hat{v}_1(\xi)| d\xi.\end{aligned}$$

Let

$$r = \tilde{q}(1+N)^{-2/3}(E+\rho^2)^{1/6}, \quad \tilde{q} = \left(\frac{8\pi c_1 c_3}{3}\right)^{-1/3}, \quad (4.3)$$

where c_1 is the constant of Lemma 2.1 and c_3 is defined by (3.15).

Then, combining the definition of \tilde{I}_1 , Lemma 2.1 for $|\xi| \leq r$ and the inequality

$$\|v_2 - v_1\|_{L^2(D)} \leq c_3 \|v_2 - v_1\|_{L^\infty(D)},$$

we get that

$$\begin{aligned}\tilde{I}_1(r) &\leq \frac{4\pi r^3}{3} c_1 (1+N)^2 \left(e^{2\rho L} \delta + \frac{c_3 \|v_1 - v_2\|_{L^\infty(D)}}{\sqrt{E+\rho^2}} \right) \leq \\ &\leq \frac{1}{2c_3} \sqrt{E+\rho^2} e^{2\rho L} \delta + \frac{\|v_1 - v_2\|_{L^\infty(D)}}{2}\end{aligned} \quad (4.4)$$

for $\tilde{q}(1+N)^{-2/3}(E+\rho^2)^{1/6} \leq 2\sqrt{E+\rho^2}$ and $E+\rho^2 \geq (1+N)^2 r_1^2$.

In addition, using (2.6), we get that

$$\tilde{I}_2(r) \leq \tilde{c}_2 \frac{N W^m}{m-3} r^{3-m}. \quad (4.5)$$

Let $\tilde{r}_2 = \tilde{r}_2(D) \geq r_1$ be such that

$$E + \rho^2 \geq \tilde{r}_2^2 \implies \tilde{q}(E + \rho^2)^{1/6} \leq 2\sqrt{E + \rho^2}. \quad (4.6)$$

Using (4.2), (4.4)–(4.6) with r defined in (4.3), we obtain that

$$\begin{aligned}\|v_2 - v_1\|_{L^\infty(D)} &\leq \frac{1}{2c_3} \sqrt{E + \rho^2} e^{2\rho L} \delta + \frac{\|v_1 - v_2\|_{L^\infty(D)}}{2} + \\ &+ \tilde{c}_2 \frac{(1+N)^{\frac{2(m-3)}{3}}}{(m-3)\tilde{q}^{m-3}} N W^m (E + \rho^2)^{-\frac{m-3}{6}},\end{aligned} \quad (4.7)$$

$$\begin{aligned}\frac{1}{2} \|v_2 - v_1\|_{L^\infty(D)} &\leq \frac{1}{2c_3} \sqrt{E + \rho^2} e^{2\rho L} \delta + \\ &+ 4\pi \frac{(1+N)^{\frac{2(m-3)}{3}}}{(m-3)\tilde{q}^{m-3}} N W^m (E + \rho^2)^{-\frac{m-3}{6}}\end{aligned} \quad (4.8)$$

for $E + \rho^2 \geq (1+N)^2 \tilde{r}_2^2$, where L, \tilde{c}_2 are the constants of Lemmas 2.1, 2.3 and $c_3, \tilde{q}, \tilde{r}_2$ are the constants of formulas (3.15), (4.3), (4.6).

Let $\tau \in (0, 1)$ and

$$\gamma = \frac{1-\tau}{2L}, \quad \rho = \gamma \ln(3 + \delta^{-1}). \quad (4.9)$$

Due to (4.8), for δ such that

$$E + (\gamma \ln(3 + \delta^{-1}))^2 \geq (1 + N)^2 \tilde{r}_2^2, \quad (4.10)$$

the following estimate holds:

$$\begin{aligned} \frac{1}{2} \|v_1 - v_2\|_{L^\infty(D)} &\leq \\ &\leq \frac{1}{2c_3} \left(E + (\gamma \ln(3 + \delta^{-1}))^2 \right)^{1/2} (3 + \delta^{-1})^{2\gamma L} \delta + \\ &\quad + 4\pi \frac{(1 + N)^{\frac{2(m-3)}{3}}}{(m-3)\tilde{q}^m} N_{W^m} \left(E + (\gamma \ln(3 + \delta^{-1}))^2 \right)^{-\frac{m-3}{6}}, \end{aligned} \quad (4.11)$$

where γ is defined in (4.9). Note that

$$(3 + \delta^{-1})^{2\gamma L} \delta = (1 + 3\delta)^{1-\tau} \delta^\tau \leq 4\delta^\tau \quad \text{for } \delta \leq 1. \quad (4.12)$$

Combining (4.11), (4.12), we get that

$$\begin{aligned} \|v_2 - v_1\|_{L^\infty(D)} &\leq \tilde{A}_1 \left(\tilde{\lambda} \left(E + \gamma^2 (\ln(3 + \delta^{-1}))^2 \right) \right)^{\frac{1}{2}} \delta^\tau + \\ &\quad + \tilde{B}_1 \frac{(1 + N)^{\frac{2(m-3)}{3}}}{m-3} N_{W^m} \left(\tilde{\lambda} \left(E + \gamma^2 (\ln(3 + \delta^{-1}))^2 \right) \right)^{-\frac{m-3}{6}} \end{aligned} \quad (4.13)$$

for $\delta \leq 1$ satisfying (4.10) and some positive constants $\tilde{A}_1, \tilde{B}_1, \tilde{\lambda}$ depending on D only.

Using (1.7) and (4.2), we get that

$$\begin{aligned} \|v_2 - v_1\|_{L^\infty(D)} &\leq \int_{\mathbb{R}^3} \left((1 + |\xi|^2)^{-m/2} \|v_2 - v_1\|_{W^m(\mathbb{R}^3)} \right) d\xi \leq \\ &\leq N_{W^m} \int_0^{+\infty} \frac{4\pi t^2}{(1 + t^2)^{m/2}} dt \leq c_4 \frac{e^{m-3}}{m-3} N_{W^m} \end{aligned} \quad (4.14)$$

for some $c_4 > 0$. Here we used also that

$$\int_0^{+\infty} \frac{4\pi t^2}{(1 + t^2)^{m/2}} dt \leq \int_0^1 4\pi t^2 dt + \int_1^{+\infty} \frac{4\pi t^2}{t^m} dt \leq c_4 \left(1 + \frac{1}{m-3} \right) \leq c_4 \frac{e^{m-3}}{m-3}.$$

Using (4.14), we get that, for $0 < E + (\gamma \ln(3 + \delta^{-1}))^2 \leq (1 + N)^2 \tilde{r}_2^2$,

$$\begin{aligned} \|v_2 - v_1\|_{L^\infty(D)} &\leq \\ &\leq c_4 \frac{(1 + N)^{\frac{2(m-3)}{3}}}{m-3} N_{W^m} \left(\frac{E + \gamma^2 (\ln(3 + \delta^{-1}))^2}{e^6 \tilde{r}_2^2} \right)^{-\frac{m-3}{6}}. \end{aligned} \quad (4.15)$$

On other hand, in the case when $E + (\gamma \ln \delta^{-1})^2 \geq (1 + N)^2 \tilde{r}_2^2$ and $\delta > 1$ we have that

$$\|v_2 - v_1\|_{L^\infty(D)} \leq 2N \leq 2 \left(\frac{E + \gamma^2 (\ln(3 + \delta^{-1}))^2}{\tilde{r}_2^2} \right)^{\frac{1}{2}} \delta^\tau. \quad (4.16)$$

Combining (4.9), (4.13), (4.15) and (4.16), we obtain estimate (1.11). This completes the proof of Theorem 1.2.

5 Faddeev functions

Suppose that

$$v \in L^\infty(D), \quad v \equiv 0 \text{ on } \mathbb{R}^3 \setminus D, \quad (5.1)$$

where D satisfies (1.2). More generally, one can assume that

$$\begin{aligned} v \text{ is a sufficiently regular function on } \mathbb{R}^3 \\ \text{with sufficient decay at infinity.} \end{aligned} \quad (5.2)$$

Under assumptions (5.2), we consider the functions ψ , μ , h :

$$\psi(x, k) = e^{ikx} \mu(x, k), \quad (5.3)$$

$$\mu(x, k) = 1 + \int_{\mathbb{R}^3} g(x-y, k) v(y) \mu(y, k) dy, \quad (5.4)$$

$$g(x, k) = -(2\pi)^{-3} \int_{\mathbb{R}^3} \frac{e^{i\xi x} d\xi}{\xi^2 + 2k\xi},$$

where $x \in \mathbb{R}^3$, $k \in \mathbb{C}^3$, $\text{Im } k \neq 0$,

$$h(k, l) = (2\pi)^{-3} \int_{\mathbb{R}^3} e^{i(k-l)x} v(x) \mu(x, k) dx, \quad (5.5)$$

where $k, l \in \mathbb{C}^3$, $k^2 = l^2$, $\text{Im } k = \text{Im } l \neq 0$. Here, (5.4) at fixed k is considered as a linear integral equation for μ , where μ is sought in $L^\infty(\mathbb{R}^3)$.

The functions ψ , h and $G = e^{ikx} g$ are known as the Faddeev functions, see [7], [8], [11], [26]. These functions were introduced for the first time in [7], [8].

In particular, we have that

$$(\Delta + k^2)G(x, k) = \delta(x),$$

$$(-\Delta + v(x))\psi(x, k) = k^2 \psi(x, k),$$

where $x \in \mathbb{R}^3$, $k \in \mathbb{C}^3 \setminus \mathbb{R}^3$.

We recall also that the Faddeev functions G , ψ , h are some extension to the complex domain of the functions of the classical scattering theory for the Schrödinger equation (in particular, h is an extension of the classical scattering amplitude).

Note also that G , ψ , h in their zero energy restriction, that is for $k^2 = 0$, $l^2 = 0$, were considered for the first time in [3]. The Faddeev functions G , ψ , h were, actually, rediscovered in [3].

For further considerations we will use the following notations:

$$\Sigma_E = \{k \in \mathbb{C}^3 : k^2 = k_1^2 + k_2^2 + k_3^2 = E\},$$

$$\Theta_E = \{k \in \Sigma_E, l \in \Sigma_E : \text{Im } k = \text{Im } l\},$$

$$|k| = (|\text{Re } k|^2 + |\text{Im } k|^2)^{1/2} \text{ for } k \in \mathbb{C}^3.$$

Under assumptions (5.2), we have that:

$$\mu(x, k) \rightarrow 1 \quad \text{as } |k| \rightarrow \infty, \quad (5.6)$$

where $x \in \mathbb{R}^3$, $k \in \Sigma_E$;

$$\hat{v}(\xi) = \lim_{\substack{(k,l) \in \Theta_E, k-l=\xi \\ |\operatorname{Im} k| = |\operatorname{Im} l| \rightarrow \infty}} h(k,l) \quad \text{for any } \xi \in \mathbb{R}^3, \quad (5.7)$$

where \hat{v} is defined by (2.1).

Results of the type (5.6) go back to [3]. Results of the type (5.7) go back to [11]. These results follow, for example, from equation (5.4), formula (5.5) and the following estimates:

$$\begin{aligned} g(x,k) &= O(|x|^{-1}) \quad \text{for } x \in \mathbb{R}^3, \\ &\text{uniformly in } k \in \mathbb{C}^3 \setminus \mathbb{R}^3, \end{aligned} \quad (5.8)$$

$$\begin{aligned} \|\Lambda^{-s} g(k) \Lambda^{-s}\|_{L^2(\mathbb{R}^3) \rightarrow L^2(\mathbb{R}^3)} &= O(|k|^{-1}), \quad \text{for } s > 1/2, \\ &\text{as } |k| \rightarrow \infty, \quad k \in \mathbb{C}^3 \setminus \mathbb{R}^3, \end{aligned} \quad (5.9)$$

where $g(x,k)$ is defined in (5.4), $g(k)$ denotes the integral operator with the Schwartz kernel $g(x-y,k)$ and Λ denotes the multiplication operator by the function $(1+|x|^2)^{1/2}$. Estimate (5.8) was given in [11]. Estimate (5.9) was formulated, first, in [22]. Concerning proof of (5.9), see [40]. In addition, estimate (5.9) in its zero energy restriction goes back to [39].

In the present work we use the following lemma:

Lemma 5.1. *Let D satisfy (1.2) and v satisfy (5.1). Let $\|v\|_{L^\infty(D)} \leq N$ for some $N > 0$. Then*

$$|\mu(x,k)| \leq c_5(1+N) \quad \text{for } x \in \mathbb{R}^3, |k| \geq r_3(1+N), \quad (5.10)$$

where $\mu(x,k)$ is the Faddeev function of (5.4) and constants $c_5, r_3 > 0$ depend on D only.

Lemma 5.1 is proved in Section 6. This proof is based on estimates (5.8) and (5.9).

In addition, we have that (see [27], [30]):

$$\begin{aligned} h_2(k,l) - h_1(k,l) &= (2\pi)^{-3} \int_{\mathbb{R}^3} \psi_1(x,-l)(v_2(x) - v_1(x))\psi_2(x,k)dx \\ &\quad \text{for } (k,l) \in \Theta_E, |\operatorname{Im} k| = |\operatorname{Im} l| \neq 0, \\ &\quad \text{and } v_1, v_2 \text{ satisfying (5.2),} \end{aligned} \quad (5.11)$$

$$\begin{aligned} h_2(k,l) - h_1(k,l) &= (2\pi)^{-3} \int_{\partial D} \psi_1(x,-l) \left[(\hat{\Phi}_2 - \hat{\Phi}_1) \psi_2(\cdot, k) \right] (x) dx \\ &\quad \text{for } (k,l) \in \Theta_E, |\operatorname{Im} k| = |\operatorname{Im} l| \neq 0, \\ &\quad \text{and } v_1, v_2 \text{ satisfying (1.5), (5.1),} \end{aligned} \quad (5.12)$$

where ψ_j, h_j denote ψ and h of (5.3) and (5.5) for $v = v_j$, and $\hat{\Phi}_j$ denotes the Dirichlet-to-Neumann map $\hat{\Phi}$ for $v = v_j$ in D , where $j = 1, 2$.

In the present work we also use the following lemma:

Lemma 5.2. *Let D satisfy (1.2). Let v_j satisfy (5.1), $\|v_j\|_{L^\infty(D)} \leq N$, $j = 1, 2$, for some $N > 0$. Then*

$$\begin{aligned} |\hat{v}_1(\xi) - \hat{v}_2(\xi) - h_1(k, l) + h_2(k, l)| &\leq \frac{c_6 N(1+N)\|v_1 - v_2\|_{L^2(D)}}{(E + \rho^2)^{1/2}} \\ \text{for } (k, l) \in \Theta_E, \quad \xi = k - l, \quad |\operatorname{Im} k| = |\operatorname{Im} l| = \rho, \\ E + \rho^2 &\geq r_4^2(1+N)^2, \end{aligned} \quad (5.13)$$

where $E \in \mathbb{R}$, \hat{v}_j is the Fourier transform of v_j , h_j denotes h of (5.5) for $v = v_j$, ($j = 1, 2$) and constants $c_6, r_4 > 0$ depend on D only.

Some versions of estimate (5.13) were given in [16], [27], [30] (see, for example, formula (3.18) of [16]). Lemma 5.2 is proved in Section 6.

6 Proofs of Lemmas 2.1, 5.1 and 5.2

Proof of Lemma 5.1. Using (5.1), (5.4) and (5.9), we get that

$$\begin{aligned} \|\mu(\cdot, k) - 1\|_{L^2(D)} &\leq \left\| \int_{\mathbb{R}^3} g(\cdot - y)v(y)\mu(y, k)dy \right\|_{L^2(D)} \leq \\ &\leq c_7 \frac{N}{|k|} \|\mu(\cdot, k)\|_{L^2(D)}, \end{aligned} \quad (6.1)$$

$$\|\mu(\cdot, k)\|_{L^2(D)} \leq c_3 + c_7 \frac{N}{|k|} \|\mu(\cdot, k)\|_{L^2(D)}, \quad (6.2)$$

where c_3 is defined by (3.15) and c_7 is some positive constant depending on D only. Hence, we obtain that

$$\|\mu(\cdot, k)\|_{L^2(D)} \leq 2c_3 \quad \text{for } |k| \geq 2c_7 N. \quad (6.3)$$

We use also that

$$\int_D \frac{1}{|x-y|^2} dy \leq \int_D 1 dy + \int_{|x-y| \leq 1} \frac{1}{|x-y|^2} dy \leq c_8^2, \quad x \in D, \quad (6.4)$$

where $c_8 = c_8(D) > 0$. Using (5.1), (5.4), (5.8), (6.3), (6.4), we get that

$$\begin{aligned} |\mu(x, k)| &\leq 1 + \left| \int_D g(x-y)v(y)\mu(y, k)dy \right| \leq \\ &\leq 1 + \left(\int_D |g(x-y)|^2 dy \right)^{1/2} N \|\mu(\cdot, k)\|_{L^2(\mathbb{R}^3)} \leq \\ &\leq c_5(D)(1+N) \quad \text{for } x \in D, |k| \geq 2c_7 N. \end{aligned} \quad (6.5)$$

■

Proof of Lemma 5.2. Due to (5.1), (5.11), we have that

$$\begin{aligned}
h_2(k, l) - h_1(k, l) &= (2\pi)^{-3} \int_D \psi_1(x, -l)(v_2(x) - v_1(x))\psi_2(x, k)dx = \\
&= (2\pi)^{-3} \int_D e^{i(k-l)x} \mu_1(x, -l)(v_2(x) - v_1(x))\mu_2(x, k)dx = \\
&= \hat{v}_2(k - l) - \hat{v}_1(k - l) + I_\Delta
\end{aligned} \tag{6.6}$$

for $(k, l) \in \Theta_E$, $|\operatorname{Im} k| = |\operatorname{Im} l| \neq 0$, where

$$\begin{aligned}
I_\Delta &= (2\pi)^{-3} \int_D (\mu_1(x, -l) - 1)(v_2(x) - v_1(x))\mu_2(x, k)dx + \\
&\quad + (2\pi)^{-3} \int_D \mu_1(x, -l)(v_2(x) - v_1(x))(\mu_2(x, k) - 1)dx + \\
&\quad + (2\pi)^{-3} \int_D (\mu_1(x, -l) - 1)(v_2(x) - v_1(x))(\mu_2(x, k) - 1)dx.
\end{aligned} \tag{6.7}$$

Note that, for $(k, l) \in \Theta_E$, $E \in \mathbb{R}$, $|\operatorname{Im} k| = |\operatorname{Im} l| = \rho$,

$$|k| = \sqrt{|\operatorname{Re} k|^2 + |\operatorname{Im} k|^2} = \sqrt{k^2 + 2|\operatorname{Im} k|^2} = \sqrt{E + 2\rho^2} = |l|. \tag{6.8}$$

Using estimates (6.1), (6.3), (6.5) in (6.7), we get that

$$\begin{aligned}
I_\Delta &\leq (2\pi)^{-3} \left(\|\mu_1(\cdot, -l) - 1\|_{L^2(D)} \|v_2 - v_1\|_{L^2(D)} \|\mu_2(\cdot, -l)\|_{L^\infty(D)} + \right. \\
&\quad \left. + \|\mu_1(\cdot, -l)\|_{L^\infty(D)} \|v_2 - v_1\|_{L^2(D)} \|\mu_2(\cdot, -l) - 1\|_{L^2(D)} + \right. \\
&\quad \left. + \|\mu_1(\cdot, -l) - 1\|_{L^\infty(D)} \|v_2 - v_1\|_{L^2(D)} \|\mu_2(\cdot, -l) - 1\|_{L^2(D)} \right) \leq \\
&\leq \frac{2c_3c_7N \|v_1 - v_2\|_{L^2(D)} c_5(1 + N)}{(2\pi)^3 |k|} + \frac{c_5(1 + N) \|v_1 - v_2\|_{L^2(D)} 2c_3c_7N}{(2\pi)^3 |l|} + \\
&\quad + \frac{(1 + c_5(1 + N)) \|v_1 - v_2\|_{L^2(D)} 2c_3c_7N}{(2\pi)^3 |l|} \leq \\
&\leq c_8(D) \frac{N(1 + N) \|v_1 - v_2\|_{L^2(D)}}{\sqrt{E + 2\rho^2}}
\end{aligned} \tag{6.9}$$

for $(k, l) \in \Theta_E$, $|\operatorname{Im} k| = |\operatorname{Im} l| = \rho$ and $|k| = |l| = \sqrt{E + 2\rho^2} \geq 2c_7N$.

Formula (6.6) and estimate (6.9) imply (5.13). ■

Proof of Lemma 2.1. Due to (5.11), we have that

$$\begin{aligned}
|h_2(k, l) - h_1(k, l)| &\leq c_9 \|\psi_1(\cdot, -l)\|_{L^\infty(\partial D)} \delta \|\psi_2(\cdot, k)\|_{L^\infty(\partial D)}, \\
(k, l) &\in \Theta_E, \quad |\operatorname{Im} k| = |\operatorname{Im} l| \neq 0,
\end{aligned} \tag{6.10}$$

where

$$c_9 = (2\pi)^{-3} \int_{\partial D} dx.$$

Using formula (5.3) and Lemma 5.1, we find that

$$\begin{aligned} \|\psi_j(\cdot, k)\|_{L^\infty(\partial D)} &\leq c_5(1+N)e^{|\operatorname{Im} k|L}, \quad j = 1, 2, \\ &\text{for } k \in \Sigma_E, \quad |k| \geq r_3(1+N), \end{aligned} \quad (6.11)$$

where $L = \max_{x \in \partial D} |x|$. Combining (6.8), (6.10) and (6.11), we get that

$$\begin{aligned} |h_2(k, l) - h_1(k, l)| &\leq c_9 c_5^2 (1+N)^2 e^{2\rho L} \delta, \\ &\text{for } (k, l) \in \Theta_E, \quad \rho = |\operatorname{Im} k| = |\operatorname{Im} l|, \\ &E + \rho^2 \geq r_3^2 (1+N)^2. \end{aligned} \quad (6.12)$$

Note that for any $\xi \in \mathbb{R}^3$ satisfying $|\xi| \leq 2\sqrt{E + \rho^2}$ (where $\rho > 0$) there exist some pair $(k, l) \in \Theta_E$ such that $\xi = k - l$ and $|\operatorname{Im} k| = |\operatorname{Im} l| = \rho$. Therefore, estimates (5.13) and (6.12) imply (2.2). \blacksquare

References

- [1] G. Alessandrini, *Stable determination of conductivity by boundary measurements*, Appl.Anal. 27, 1988, 153-172.
- [2] N.V. Alexeenko, V.A. Burov, O.D. Rumyantseva, *Solution of three-dimensional acoustical inverse scattering problem, II: modified Novikov algorithm*, Acoust. J. 54(3), 2008, 469–482 (in Russian); English transl.: Acoust. Phys. 54(3), 2008, 407–419.
- [3] R. Beals, R. Coifman, *Multidimensional inverse scattering and nonlinear partial differential equations*, Proc. Symp. Pure Math., 43, 1985, 45-70.
- [4] L. Beilina, M.V. Klibanov, *Approximate global convergence and adaptivity for coefficient inverse problems*, Springer (New York), 2012. 407 pp.
- [5] A.P. Calderón, *On an inverse boundary problem*, Seminar on Numerical Analysis and its Applications to Continuum Physics, Soc. Brasileira de Matematica, Rio de Janeiro, 1980, 61–73.
- [6] M. Di Cristo, L. Rondi *Examples of exponential instability for inverse inclusion and scattering problems* Inverse Problems 19, 2003, 685–701.
- [7] L.D. Faddeev, *Growing solutions of the Schrödinger equation*, Dokl. Akad. Nauk SSSR, 165, N.3, 1965, 514-517 (in Russian); English Transl.: Sov. Phys. Dokl. 10, 1966, 1033-1035.
- [8] L.D. Faddeev, *The inverse problem in the quantum theory of scattering. II*, Current problems in mathematics, Vol. 3, 1974, pp. 93-180, 259. Akad. Nauk SSSR Vsesojuz. Inst. Nauch. i Tehn. Informacii, Moscow (in Russian); English Transl.: J.Sov. Math. 5, 1976, 334-396.
- [9] I.M. Gelfand, *Some problems of functional analysis and algebra*, Proceedings of the International Congress of Mathematicians, Amsterdam, 1954, pp.253-276.

- [10] P. Hähner, T. Hohage, *New stability estimates for the inverse acoustic inhomogeneous medium problem and applications*, SIAM J. Math. An., 33(3), 2001, 670–685.
- [11] G.M. Henkin and R.G. Novikov, *The $\bar{\partial}$ -equation in the multidimensional inverse scattering problem*, Uspekhi Mat. Nauk 42(3), 1987, 93–152 (in Russian); English Transl.: Russ. Math. Surv. 42(3), 1987, 109–180.
- [12] M.I. Isaev, *Exponential instability in the Gel'fand inverse problem on the energy intervals*, J. Inverse Ill-Posed Probl., Vol. 19(3), 2011, 453–473.
- [13] M.I. Isaev, *Instability in the Gel'fand inverse problem at high energies*, Applicable Analysis, Vol. 92, No. 11, 2013, 2262–2274.
- [14] M.I. Isaev, *Energy and regularity dependent stability estimates for near-field inverse scattering in multidimensions*, Journal of Mathematics, Article ID 318154, 2013, 10 p.; DOI:10.1155/2013/318154.
- [15] M.I. Isaev, *Exponential instability in the inverse scattering problem on the energy interval*, Func. Anal. Prilozh., Vol. 47(3), 2013, 28–36 (in Russian); Engl. translation: Functional Analysis and Its Applications, Vol. 47(3), 2013, 187–194.
- [16] M.I. Isaev, R.G. Novikov *Energy and regularity dependent stability estimates for the Gel'fand inverse problem in multidimensions*, J. of Inverse and Ill-posed Probl., Vol. 20(3), 2012, 313–325.
- [17] M.I. Isaev, R.G. Novikov, *Stability estimates for determination of potential from the impedance boundary map*, Algebra and Analysis, Vol. 25(1), 2013, 37–63 (in Russian); Engl. Transl.: St. Petersburg Mathematical Journal, Vol. 25, 2014, 23–41.
- [18] M.I. Isaev, R.G. Novikov, *Reconstruction of a potential from the impedance boundary map*, Eurasian Journal of Mathematical and Computer Applications, Vol. 1(1), 2013, 5–28.
- [19] M.I. Isaev, R.G. Novikov, *New global stability estimates for monochromatic inverse acoustic scattering*, SIAM Journal on Mathematical Analysis, Vol. 45(3), 2013, 1495–1504.
- [20] V. Isakov, *Increasing stability for the Schrödinger potential from the Dirichlet-to-Neumann map*, Discrete Contin. Dyn. Syst. Ser. S 4, 2011, no. 3, 631-640.
- [21] V. Isakov, S. Nagayasu, G. Uhlmann, J.-N. Wang, *Increasing stability of the inverse boundary value problem for the Schrödinger equation*, Contemporary Mathematics (to appear), e-print arXiv:1302.0940.
- [22] R.B. Lavine and A.I. Nachman, *On the inverse scattering transform of the n -dimensional Schrödinger operator* Topics in Soliton Theory and Exactly Solvable Nonlinear Equations ed M Ablowitz, B Fuchssteiner and M Kruskal (Singapore: World Scientific), 1987, pp 33-44.

- [23] M.M. Lavrentev, V.G. Romanov, S.P. Shishatskii, *Ill-posed problems of mathematical physics and analysis*, Translated from the Russian by J. R. Schulenberger. Translation edited by Lev J. Leifman. Translations of Mathematical Monographs, 64. American Mathematical Society, Providence, RI, 1986. vi+290 pp.
- [24] N. Mandache, *Exponential instability in an inverse problem for the Schrödinger equation*, Inverse Problems 17, 2001, 1435-1444.
- [25] S. Nagayasu, G. Uhlmann, J.-N. Wang, *Increasing stability in an inverse problem for the acoustic equation*, Inverse Problems 29, 2013, 025013(11pp).
- [26] R.G. Novikov, *Multidimensional inverse spectral problem for the equation $-\Delta\psi + (v(x) - Eu(x))\psi = 0$* , Funkt. Anal. Prilozhen. 22(4), 1988, 11-22 (in Russian); Engl. Transl. Funct. Anal. Appl. 22, 1988, 263-272.
- [27] R.G. Novikov, *$\bar{\partial}$ -method with nonzero background potential. Application to inverse scattering for the two-dimensional acoustic equation*, Comm. Partial Differential Equations 21, 1996, no. 3-4, 597-618.
- [28] R.G. Novikov, *Rapidly converging approximation in inverse quantum scattering in dimension 2*, Physics Letters A 238, 1998, 73-78.
- [29] R.G. Novikov, *The $\bar{\partial}$ -approach to approximate inverse scattering at fixed energy in three dimensions*. IMRP Int. Math. Res. Pap. 2005, no. 6, 287-349.
- [30] R.G. Novikov, *Formulae and equations for finding scattering data from the Dirichlet-to-Neumann map with nonzero background potential*, Inverse Problems 21, 2005, 257-270.
- [31] R.G. Novikov, *The $\bar{\partial}$ -approach to monochromatic inverse scattering in three dimensions*, J. Geom. Anal 18, 2008, 612-631.
- [32] R.G. Novikov, *An effectivization of the global reconstruction in the Gel'fand-Calderon inverse problem in three dimensions*, Contemporary Mathematics, 494, 2009, 161-184.
- [33] R.G. Novikov, *New global stability estimates for the Gel'fand-Calderon inverse problem*, Inverse Problems 27, 2011, 015001(21pp).
- [34] R.G. Novikov, *Approximate Lipschitz stability for non-overdetermined inverse scattering at fixed energy*, J.Inverse Ill-Posed Probl., Vol. 21(6), 2013, 813-823.
- [35] R. Novikov and M. Santacesaria, *A global stability estimate for the Gelfand-Calderon inverse problem in two dimensions*, J.Inverse Ill-Posed Probl., Volume 18, Issue 7, 2010, 765-785.
- [36] L. Päivärinta, M. Zubeldia, *The inverse Robin boundary value problem in a half-space*, e-print arXiv:1311.6947
- [37] M. Santacesaria, *Stability estimates for an inverse problem for the Schrödinger equation at negative energy in two dimensions*, Applicable Analysis, 2013, Vol. 92, No. 8, 1666-1681.

- [38] M. Santacesaria, *A Holder-logarithmic stability estimate for an inverse problem in two dimensions*, J. Inverse Ill-Posed Probl. (to appear), e-print arXiv:1306.0763.
- [39] J. Sylvester and G. Uhlmann, *A global uniqueness theorem for an inverse boundary value problem*, Ann. of Math. 125, 1987, 153-169.
- [40] R. Weder, *Generalized limiting absorption method and multidimensional inverse scattering theory*, Mathematical Methods in the Applied Sciences, 14, 1991, 509-524.

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