# Approximate controllability of Lagrangian trajectories of the 3D Navier-Stokes system by a finite-dimensional force 

Vahagn Nersesyan

## - To cite this version:

Vahagn Nersesyan. Approximate controllability of Lagrangian trajectories of the 3D Navier-Stokes system by a finite-dimensional force. 2015. hal-00962254v2

## HAL Id: hal-00962254 <br> https://hal.science/hal-00962254v2

Preprint submitted on 15 Jan 2015

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Approximate controllability of Lagrangian trajectories of the 3D Navier-Stokes system by a finite-dimensional force 

Vahagn Nersesyan*


#### Abstract

In the Eulerian approach, the motion of an incompressible fluid is usually described by the velocity field which is given by the Navier-Stokes system. The velocity field generates a flow in the space of volume-preserving diffeomorphisms. The latter plays a central role in the Lagrangian description of a fluid, since it allows to identify the trajectories of the individual particles. In this paper, we show that the velocity field of the fluid and the corresponding flow of diffeomorphisms can be simultaneously approximately controlled using a finite-dimensional external force. The proof is based on some methods from the geometric control theory introduced by Agrachev and Sarychev.


AMS subject classifications: 35Q30, 93B05.
Keywords: Navier-Stokes system, Agrachev-Sarychev method, approximate controllability

## Contents

0 Introduction ..... 2
1 Preliminaries ..... 7
1.1 Particle trajectories ..... 7
1.2 Existence of strong solutions ..... 11
2 Approximate controllability of the NS system ..... 14
3 Proof Theorem 2.5 ..... 16
4 Examples of saturating spaces ..... 20
4.0.1 Saturating spaces associated with the generators of $\mathbb{Z}^{3}$. ..... 20
4.0.2 Controls with two vanishing components ..... 26
4.0.3 6-dimensional example ..... 27

[^0]
## 0 Introduction

The motion of an incompressible fluid is described by the following NavierStokes (NS) system

$$
\begin{align*}
\partial_{t} u-\nu \Delta u+\langle u, \nabla\rangle u+\nabla p & =f(t, x), \quad \operatorname{div} u=0  \tag{0.1}\\
u(0) & =u_{0}, \tag{0.2}
\end{align*}
$$

where $\nu>0$ is the kinematic viscosity, $u=\left(u_{1}(t, x), u_{2}(t, x), u_{3}(t, x)\right)$ is the velocity field of the fluid, $p=p(t, x)$ is the pressure, and $f$ is an external force. Throughout this paper, we shall assume that the space variable $x=\left(x_{1}, x_{2}, x_{3}\right)$ belongs to the torus $\mathbb{T}^{3}=\mathbb{R}^{3} / 2 \pi \mathbb{Z}^{3}$.

The well-posedness of the 3D NS system (0.1) is a famous open problem. Given smooth data $\left(u_{0}, f\right)$, the existence and uniqueness of a smooth solution is known to hold only locally in time. The global existence is established in the case of small data. For large data the global existence holds in the case of a weak solution, but in that case the uniqueness is open.

The flow generated by a sufficiently smooth velocity field $u$ gives the Lagrangian trajectories of the fluid:

$$
\begin{equation*}
\dot{x}=u(t, x), \quad x(0)=x_{0} \in \mathbb{T}^{3} \tag{0.3}
\end{equation*}
$$

Since the fluid is assumed to be incompressible, for any $t \geq 0$, the mapping $\phi_{t}^{u}$ : $x_{0} \mapsto x(t)$ belongs to the group SDiff( $\left.\mathbb{T}^{3}\right)$ of orientation and volume preserving diffeomorphisms on $\mathbb{T}^{3}$ isotopic to the identity. This group is often referred as configuration space of the fluid (cf. [AK98, KW09]). Thus for sufficiently smooth data, we have a path $\left(u(t), \phi_{t}^{u}\right)$, which is defined locally in time. The main issue addressed in this paper is the approximate controllability of the couple $\left(u(T), \phi_{T}^{u}\right)$ for any $T>0$. We shall assume that the external force is of the following form

$$
f(t, x)=h(t, x)+\eta(t, x)
$$

where $h$ is the fixed part of the force (given function) and $\eta$ is a control force. To state the main result of this paper, we need to introduce some notation. Let us define the space

$$
\begin{equation*}
H:=\left\{u \in L^{2}\left(\mathbb{T}^{3}, \mathbb{R}^{3}\right): \operatorname{div} u=0, \quad \int_{\mathbb{T}^{3}} u(x) \mathrm{d} x=0\right\} \tag{0.4}
\end{equation*}
$$

and denote by $\Pi$ the orthogonal projection onto $H$ in $L^{2}\left(\mathbb{T}^{3}, \mathbb{R}^{3}\right)$. Consider the projection of system (0.1) onto $H$ :

$$
\begin{equation*}
\dot{u}+L u+B(u)=h(t, x)+\eta(t, x), \tag{0.5}
\end{equation*}
$$

where $L=-\Delta$ is the Stokes operator and $B(u):=\Pi(\langle u, \nabla\rangle u)$. Let us set $H_{\sigma}^{k}:=H^{k}\left(\mathbb{T}^{3}, \mathbb{R}^{3}\right) \cap H$, where $H^{k}\left(\mathbb{T}^{3}, \mathbb{R}^{3}\right)$ is the space of vector functions $v=$
$\left(v_{1}, v_{2}, v_{3}\right)$ with components in the usual Sobolev space of order $k$ on $\mathbb{T}^{3}$. Let $E$ a be subset of $H$. We shall say that system (0.5) is approximately controllable by an $E$-valued control, if for any $\nu>0, k \geq 3, \varepsilon>0, T>0, u_{0}, u_{1} \in H_{\sigma}^{k}$, $h \in L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right)$, and $\psi \in \operatorname{SDiff}\left(\mathbb{T}^{3}\right)$, there is a control $\eta \in L^{2}([0, T], E)$ and a solution $u$ of (0.5), (0.2) defined for any $t \in[0, T]$ and satisfying

$$
\left\|u(T)-u_{1}\right\|_{H^{k}\left(\mathbb{T}^{3}\right)}+\left\|\phi_{T}^{u}-\psi\right\|_{C^{1}\left(\mathbb{T}^{3}\right)}<\varepsilon
$$

The following theorem is a simplified version of our main result (see Section 2, Corollary 2.3).

Main Theorem. There is a finite-dimensional subspace $E \subset H$ such that (0.5) is approximately controllable by an E-valued control.

Roughly speaking, this shows that, using a finite-dimensional external force, one can drive the fluid flow (which starts at the identity) arbitrarily close to any configuration $\psi \in \operatorname{SDiff}\left(\mathbb{T}^{3}\right)$. Moreover, near the final position $\psi(x)$, the particle starting from $x$ will have approximately the prescribed velocity $v_{1}(x):=$ $u_{1}(\psi(x))$. Note that $\phi_{T}^{u}$ depends not only on $u(T)$, but on the whole path $u(t), t \in$ $[0, T]$. Thus one needs a path-controllability property for the velocity field in order to prove controllability for $\phi_{T}^{u}$. This path-controllability is one of the novelties of this paper, it is established in Theorem 2.2.

We give some explicit examples of finite-dimensional subspaces $E$ which ensure the above approximate controllability property. For instance, for any $\ell \in$ $\mathbb{Z}^{3}$, let $\{l(\ell), l(-\ell)\}$ be an arbitrary orthonormal basis in $\left\{x \in \mathbb{R}^{3}:\langle x, \ell\rangle=0\right\}$. We show that our problem is controllable by $\eta$ taking values in a space of the form

$$
\begin{equation*}
E=E(\mathcal{K}):=\operatorname{span}\{l( \pm \ell) \cos \langle\ell, x\rangle, l( \pm \ell) \sin \langle\ell, x\rangle: \ell \in \mathcal{K}\}, \mathcal{K} \subset \mathbb{Z}^{3} \tag{0.6}
\end{equation*}
$$

if and only if $\mathcal{K}$ is a generator of $\mathbb{Z}^{3}$ (i.e., any $a \in \mathbb{Z}^{3}$ is a finite linear combination of the elements of $\mathcal{K}$ with integer coefficients). The simplest example of a generator of $\mathbb{Z}^{3}$ is

$$
\mathcal{K}=\{(1,0,0),(0,1,0),(0,0,1)\}
$$

in which case $\operatorname{dim} E(\mathcal{K})=12$. We also establish approximate controllability of the system in question by controls having two vanishing components. More precisely, the space $E$ can be chosen of the form

$$
\begin{equation*}
E=\Pi\{(0,0,1) \zeta: \zeta \in \mathcal{H}\} \tag{0.7}
\end{equation*}
$$

where

$$
\mathcal{H}:=\operatorname{span}\{\sin \langle m, x\rangle, \cos \langle m, x\rangle: m \in \mathcal{K}\}
$$

and $\mathcal{K}:=\{(1,0,0),(0,1,0),(1,0,1),(0,1,1)\}$ (i.e., $\operatorname{dim} E=8)$. In (4.22) an example of a 6 -dimensional subspace is given which guarantees the controllability of the 3D NS system.

The strategy of the proof of Main Theorem is based on some methods introduced by Agrachev and Sarychev in [AS05, AS06] (see also the survey [AS08]).

In that papers they prove approximate controllability for the 2 D NS and Euler systems by a finite-dimensional force. This method is then developed and generalised by several authors for various PDE's. Rodrigues proves in [Rod06] controllability for the 2D NS system on a rectangle with the Lions boundary conditions, and in [Rod07, Rod08] he extends the results to the case of more general Navier boundary conditions and the Hemisphere under the Lions boundary conditions. The controllability for the 3D NS system on the torus is studied in [Shi06, Shi07a] by Shirikyan. He also considers the case of the Burgers equation on the real line in [Shi13] and on an interval with the Dirichlet boundary conditions in [Shi07b, Shi10]. Incompressible and compressible 3D Euler equations are considered by Nersisyan in [Ner10, Ner11], and the controllability for the 2D defocusing cubic Schrödinger equation is established by Sarychev in [Sar12]. In [Shi08a] Shirikyan proves that the Euler equations are not exactly controllable by a finite-dimensional external force.

All the above papers are concerned with the problem of controllability of the velocity field. The controllability of the Lagrangian trajectories of 2D and 3D Euler equations is studied by Glass and Horsin [GH10, GH12], in the case of boundary controls. For given two smooth contractible sets $\gamma_{1}$ and $\gamma_{2}$ of fluid particles which surround the same volume, they construct a control such that the corresponding flow drives $\gamma_{1}$ arbitrarily close to $\gamma_{2}$. In the context of our paper, a similar property can be derived from our main result. Indeed, Krygin shows in [Kry71] that there is a diffeomorphism $\psi \in \operatorname{SDiff}\left(\mathbb{T}^{3}\right)$ such that $\psi\left(\gamma_{1}\right)=\gamma_{2}$. Thus we can find an $E$-valued control $\eta$ such that $\phi_{T}^{u}\left(\gamma_{1}\right)$ is arbitrarily close to $\gamma_{2}$, and, moreover, at time $T$ the particles will have approximately the desired velocity.

When $E$ is of the form (0.7), our Main Theorem is related to the recent paper [CL12] by Coron and Lissy. In that paper, the authors establish local null controllability of the velocity for the 3D NS system controlled by a distributed force having two vanishing components (i.e., the controls are valued in a space of the form ( 0.7 ), where $\mathcal{H}$ is the space of space-time $L^{2}$-functions supported in a given open subset). The reader is referred to the book [Cor07] for an introduction to the control theory of the NS system by distributed controls and for references on that topic.

Let us give a brief (and not completely accurate) description of how the Agrachev-Sarychev method is adapted to establish approximate controllability in the above-defined sense. We assume that $E$ is given by ( 0.6 ) for some generator $\mathcal{K}$ of $\mathbb{Z}^{3}$. Let $\psi \in \operatorname{SDiff}\left(\mathbb{T}^{3}\right)$ and let $I(t, x)$ be a smooth isotopy connecting it to the identity: $I(0, x)=x$ and $I(T, x)=\psi(x)$. Then $\hat{u}(t, x):=\partial_{t} I\left(t, I^{-1}(t, x)\right)$ is a divergence-free vector field such that $\phi_{t}^{\hat{u}}(x)=I(t, x)$ for all $t \in[0, T]$. In particular, $\phi_{T}^{\hat{u}}=\psi$. The mapping $u \mapsto \phi_{T}^{u}$ is continuous from $L^{1}\left([0, T], H_{\sigma}^{k}\right)$ to $C^{1}\left(\mathbb{T}^{3}\right)$, where $L^{1}\left([0, T], H_{\sigma}^{k}\right)$ is endowed with the relaxation norm

$$
\|u\|_{T, k}:=\sup _{t \in[0, T]}\left\|\int_{0}^{t} u(s) \mathrm{d} s\right\|_{H^{k}\left(\mathbb{T}^{3}\right)}
$$

Hence we can choose a smooth vector field $u$ sufficiently close to $\hat{u}$ with respect
to this norm, so that

$$
u(0)=u_{0}, \quad u(T)=u_{1}, \quad\left\|\phi_{T}^{u}-\psi\right\|_{C^{1}\left(\mathbb{T}^{3}\right)}<\varepsilon .
$$

Then $u$ is a solution of our system corresponding to a control $\eta_{0}$, which can be explicitly expressed in terms of $u$ and $h$ from equation (0.5). In general, this control $\eta_{0}$ is not $E$-valued, so we need to approximate $u$ appropriately with solutions corresponding to $E$-valued controls. To this end, we define the sets

$$
\mathcal{K}_{0}:=\mathcal{K}, \quad \mathcal{K}_{j}=\mathcal{K}_{j-1} \cup\left\{m \pm n: m, n \in \mathcal{K}_{j-1}\right\}, \quad j \geq 1
$$

As $\mathcal{K}$ is a generator of $\mathbb{Z}^{3}$, one easily gets that $\cup_{j \geq 1} \mathcal{K}_{j}=\mathbb{Z}^{3}$, hence $\cup_{j \geq 1} E\left(\mathcal{K}_{j}\right)$ is dense in $H_{\sigma}^{k}$. Let $P_{N}$ be the orthogonal projection onto $E\left(\mathcal{K}_{N}\right)$ in $\bar{H}$. Then a perturbative result implies that, for a sufficiently large $N \geq 1$, system (0.5), (0.2) with control $P_{N} \eta_{0}$ has a strong solution $u_{N}$ verifying (see Theorem 1.3 and Lemma 1.1)

$$
\left\|u_{N}(T)-u_{1}\right\|_{H^{k}\left(\mathbb{T}^{3}\right)}+\left\|\phi_{T}^{u_{N}}-\psi\right\|_{C^{1}\left(\mathbb{T}^{3}\right)}<\varepsilon .
$$

On the other hand, if we consider the following auxiliary system

$$
\begin{equation*}
\dot{u}+\nu L(u+\zeta)+B(u+\zeta)=h+\eta \tag{0.8}
\end{equation*}
$$

with two controls $\zeta$ and $\eta$, then the below two properties hold true
Convexification principle. For any $\varepsilon>0$ and any solution $u_{j}$ of (0.5), (0.2) with an $E\left(\mathcal{K}_{j}\right)$-valued control $\eta_{1}$, there are $E\left(\mathcal{K}_{j-1}\right)$-valued controls $\zeta$ and $\eta$ and a solution $\tilde{u}_{j-1}$ of (0.8), (0.2) such that

$$
\left\|u_{j}(T)-\tilde{u}_{j-1}(T)\right\|_{H^{k}\left(\mathbb{T}^{3}\right)}+\left\|u_{j}-\tilde{u}_{j-1}\right\|_{T, k}<\varepsilon
$$

Extension principle. For any $\varepsilon>0$ and any solution $\tilde{u}_{j}$ of (0.8), (0.2) with $E\left(\mathcal{K}_{j}\right)$-valued controls $\zeta$ and $\eta$, there is an $E\left(\mathcal{K}_{j}\right)$-valued control $\eta_{2}$ and a solution $u_{j}$ of (0.5), (0.2) such that

$$
\left\|u_{j}(T)-\tilde{u}_{j}(T)\right\|_{H^{k}\left(\mathbb{T}^{3}\right)}+\left\|u_{j}-\tilde{u}_{j}\right\|_{T, k}<\varepsilon
$$

These two principles and the above-mentioned continuity property of $\phi_{T}^{u}$ with respect to the relaxation norm imply that, for any solution $u_{j}$ of (0.5), (0.2) with an $E\left(\mathcal{K}_{j}\right)$-valued control $\eta_{1}$, there is an $E\left(\mathcal{K}_{j-1}\right)$-valued control $\eta_{2}$ and a solution $u_{j-1}$ of (0.5), (0.2) such that

$$
\left\|u_{j}(T)-\tilde{u}_{j-1}(T)\right\|_{H^{k}\left(\mathbb{T}^{3}\right)}+\left\|\phi_{T}^{u_{j}}-\phi_{T}^{u_{j-1}}\right\|_{C^{1}\left(\mathbb{T}^{3}\right)}<\varepsilon
$$

Combining this with the above-constructed solution $u_{N}$, we get the approximate controllability of (0.5) by a control valued in $E(\mathcal{K})=E$. The proofs of convexification and extension principles are strongly inspired by [Shi06].

## Acknowledgments

The author would like to thank the referee for his careful reading and valuable comments. This research was supported by the ANR grants EMAQS (No ANR 2011 BS01 017 01) and STOSYMAP (No ANR 2011 BS01 015 01).

## Notation

We denote by $\mathbb{T}^{d}$ the standard $d$-dimensional torus $\mathbb{R}^{d} / 2 \pi \mathbb{Z}^{d}$. It is endowed with the metric and the measure induced by the usual Euclidean metric and the Lebesgue measure on $\mathbb{R}^{d}$. More precisely, if $\Pi: \mathbb{R}^{d} \rightarrow \mathbb{T}^{d}$ denotes the canonical projection, we have

$$
\begin{aligned}
d(x, y) & =\inf \left\{|\tilde{x}-\tilde{y}|: \Pi \tilde{x}=x, \Pi \tilde{y}=y, \tilde{x}, \tilde{y} \in \mathbb{R}^{d}\right\} \quad \text { for any } x, y \in \mathbb{T}^{d} \\
\mathrm{~d}(A) & =(2 \pi)^{-d} \mathrm{~d}_{\mathbb{R}^{d}}\left(\Pi^{-1}(A) \cap[0,2 \pi]^{d}\right) \quad \text { for any Borel subset } A \subset \mathbb{T}^{d}
\end{aligned}
$$

where $|x|=\left|x_{1}\right|+\cdots+\left|x_{d}\right|, x=\left(x_{1}, \ldots, x_{d}\right) \in \mathbb{R}^{d}$ and $\mathrm{d}_{\mathbb{R}^{d}}$ is the Lebesgue measure on $\mathbb{R}^{d}$. We denote by $\langle\cdot, \cdot\rangle$ the scalar product in $\mathbb{R}^{d}$. $L^{p}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)$ and $H^{s}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)$ stand for spaces of vector functions $u=\left(u_{1}, \ldots, u_{d}\right)$ with components in the usual Lebesgue and Sobolev spaces on $\mathbb{T}^{d}$.
$C^{k, \lambda}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right), k \geq 0, \lambda \in(0,1]$ is the space of vector functions $u=\left(u_{1}, \ldots, u_{d}\right)$ with components that are continuous on $\mathbb{T}^{d}$ together with their derivatives up to order $k$, and whose derivatives of order $k$ are Hölder-continuous of exponent $\lambda$, equipped with the norm

$$
\|u\|_{C^{k, \lambda}}:=\sum_{|\alpha| \leq k} \sup _{x \in \mathbb{T}^{d}}\left|D^{\alpha} u(x)\right|+\sum_{|\alpha|=k} \sup _{x, y \in \mathbb{T}^{d}, x \neq y} \frac{\left|D^{\alpha} u(x)-D^{\alpha} u(y)\right|}{d(x, y)^{\lambda}}
$$

$H_{\sigma}^{k}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right):=H^{k}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right) \cap H$ and $C_{\sigma}^{k, \lambda}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right):=C^{k, \lambda}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right) \cap H$, where $H$ is given by (0.4) (with $d$ instead of 3 ). In what follows, when the space dimension $d$ is 3 , we shall write $L^{p}, H^{k}, \ldots$ instead of $L^{p}\left(\mathbb{T}^{3}, \mathbb{R}^{3}\right), H^{k}\left(\mathbb{T}^{3}, \mathbb{R}^{3}\right), \ldots$
$C^{1}\left(\mathbb{T}^{d}\right)$ is the space of continuously differentiable maps from $\mathbb{T}^{d}$ to $\mathbb{T}^{d}$ endowed with the usual distance $\left\|\psi_{1}-\psi_{2}\right\|_{C^{1}\left(\mathbb{T}^{d}\right)}, \psi_{1}, \psi_{2} \in C^{1}\left(\mathbb{T}^{d}\right)$.

Let $X$ be a Banach space endowed with a norm $\|\cdot\|_{X}$ and $J_{T}:=[0, T]$. For $1 \leq p<\infty$, let $L^{p}\left(J_{T}, X\right)$ be the space of measurable functions $u: J_{T} \rightarrow X$ such that

$$
\|u\|_{L^{p}\left(J_{T}, X\right)}:=\left(\int_{0}^{T}\|u(s)\|_{X}^{p} \mathrm{~d} s\right)^{\frac{1}{p}}<\infty
$$

The spaces $C\left(J_{T}, X\right)$ and $W^{k, p}\left(J_{T}, X\right)$ are defined in a similar way. We define the relaxation norm on $L^{1}\left(J_{T}, X\right)$ by

$$
\begin{equation*}
\|u\|_{T, X}:=\sup _{t \in J_{T}}\left\|\int_{0}^{t} u(s) \mathrm{d} s\right\|_{X} . \tag{0.9}
\end{equation*}
$$

A mapping $\psi: \mathbb{T}^{d} \rightarrow \mathbb{T}^{d}$ is volume-preserving if $\mathrm{d}\left(\psi^{-1}(A)\right)=\mathrm{d}(A)$ for any Borel subset $A \subset \mathbb{T}^{d}$. We shall say that $\psi \in C^{1}\left(\mathbb{T}^{d}\right)$ is orientation-preserving if the differential $D_{x} \psi$ is an orientation-preserving linear map for all $x \in \mathbb{T}^{d}$. We denote by $\operatorname{SDiff}\left(\mathbb{T}^{d}\right)$ be the group of all diffeomorphisms on $\mathbb{T}^{d}$ preserving the orientation and volume and isotopic to the identity, i.e., $\operatorname{SDiff}\left(\mathbb{T}^{d}\right)$ is the set of all functions $\psi: \mathbb{T}^{d} \rightarrow \mathbb{T}^{d}$ such that there is a path $I \in W^{1, \infty}\left(J_{1}, C^{1}\left(\mathbb{T}^{d}\right)\right)$ with $I(0, x)=x, I(1, x)=\psi(x)$ for all $x \in \mathbb{T}^{d}$, and $I(t, \cdot)$ is a diffeomorphism on $\mathbb{T}^{d}$ preserving the orientation and volume for all $t \in J_{1}$.

## 1 Preliminaries

### 1.1 Particle trajectories

In this section, we study some existence and stability properties for the Lagrangian trajectories. Let us fix a time $T>0$ and an integer $d \geq 1$. For any vector field $u \in L^{1}\left(J_{T}, C^{1}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right)$, we consider the following ordinary differential equation in $\mathbb{T}^{d}$

$$
\begin{equation*}
\dot{x}=u(t, x) . \tag{1.1}
\end{equation*}
$$

By standard methods, one can show that for any $y \in \mathbb{T}^{d}$ this equation admits a unique solution $x \in W^{1,1}\left(J_{T}, \mathbb{T}^{d}\right)$ such that $x(0)=y$ (e.g., see Chapter 1 in [Hal80] and Section 2.1 in [Shi08b]). Moreover, if $\phi_{t}^{u}: \mathbb{T}^{d} \rightarrow \mathbb{T}^{d}, t \in J_{T}$ is the corresponding flow sending $y$ to $x(t)$, then $\phi_{t}^{u}$ is a $C^{1}$-diffeomorphism on $\mathbb{T}^{n}$ and

$$
\begin{equation*}
\phi: L^{1}\left(J_{T}, C^{1}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right) \rightarrow C\left(J_{T}, C^{1}\left(\mathbb{T}^{d}\right)\right), \quad u \mapsto \phi^{u} \quad \text { is continuous. } \tag{1.2}
\end{equation*}
$$

We shall also use the following stability property with respect to a weaker norm (cf. Chapter 4 in [Gam78]).

Lemma 1.1. For any $\lambda \in(0,1]$ and $R>0$, there is $C:=C(R, \lambda, T)>0$ such that

$$
\begin{equation*}
\left\|\phi^{u}-\phi^{\hat{u}}\right\|_{L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{T}^{d}\right)\right)} \leq C\|u-\hat{u}\|_{T, C^{1}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)}^{\lambda / 2} \tag{1.3}
\end{equation*}
$$

for any $u, \hat{u} \in L^{\infty}\left(J_{T}, C^{1, \lambda}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right)$ verifying

$$
\begin{gather*}
\|u-\hat{u}\|_{T, C^{1}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)}<1  \tag{1.4}\\
\|u\|_{L^{\infty}\left(J_{T}, C^{1, \lambda}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right)}+\|\hat{u}\|_{L^{\infty}\left(J_{T}, C^{1, \lambda}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right)} \leq R \tag{1.5}
\end{gather*}
$$

Proof. We shall regard $u$ and $\hat{u}$ as functions on $\mathbb{R}^{d}$ which are $2 \pi$-periodic in each variable. Clearly, it suffices to prove this lemma in the case when $\mathbb{T}^{d}$ is replaced by $\mathbb{R}^{d}$ and $\phi_{t}^{u}, \phi_{t}^{\hat{u}}: \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}, t \in J_{T}$ are the flows corresponding to $u$ and $\hat{u}$.

Step 1. Let us show that there is a constant $C:=C(R, T)>0$ such that

$$
\begin{equation*}
\left\|\phi^{u}-\phi^{\hat{u}}\right\|_{L^{\infty}\left(J_{T} \times \mathbb{R}^{d}\right)} \leq C\|u-\hat{u}\|_{T, L^{\infty}\left(\mathbb{R}^{d}\right)}^{1 / 2} . \tag{1.6}
\end{equation*}
$$

Indeed, we have

$$
\begin{align*}
\left\|\phi_{t}^{u}-\phi_{t}^{\hat{u}}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}= & \left\|\int_{0}^{t}\left(u\left(s, \phi_{s}^{u}\right)-\hat{u}\left(s, \phi_{s}^{\hat{u}}\right)\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
\leq & \left\|\int_{0}^{t}\left(u\left(s, \phi_{s}^{u}\right)-u\left(s, \phi_{s}^{\hat{u}}\right)\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
& +\left\|\int_{0}^{t}\left(u\left(s, \phi_{s}^{\hat{u}}\right)-\hat{u}\left(s, \phi_{s}^{\hat{u}}\right)\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}=: G_{1}+G_{2} . \tag{1.7}
\end{align*}
$$

Then

$$
\begin{equation*}
G_{1} \leq\|u\|_{L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{R}^{d}\right)\right)} \int_{0}^{t}\left\|\phi_{s}^{u}-\phi_{s}^{\hat{u}}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \mathrm{d} s \tag{1.8}
\end{equation*}
$$

To estimate $G_{2}$, let us first note that

$$
\left|\phi_{t_{1}}^{\hat{u}}(y)-\phi_{t_{2}}^{\hat{u}}(y)\right|=\left|\int_{t_{1}}^{t_{2}} \hat{u}\left(s, \phi_{s}^{\hat{u}}(y)\right) \mathrm{d} s\right| \leq\left|t_{2}-t_{1}\right|\|\hat{u}\|_{L^{\infty}\left(J_{T} \times \mathbb{R}^{d}\right)}
$$

for any $y \in \mathbb{R}^{d}$ and $t_{1}, t_{2} \in J_{T}$. Hence for any $\eta>0$,

$$
\begin{equation*}
\sup _{t_{1}, t_{2} \in J_{T},\left|t_{1}-t_{2}\right| \leq \eta}\left\|\phi_{t_{1}}^{\hat{u}}-\phi_{t_{2}}^{\hat{u}}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \leq \eta\|\hat{u}\|_{L^{\infty}\left(J_{T} \times \mathbb{R}^{d}\right)} \tag{1.9}
\end{equation*}
$$

Taking a partition $\tau_{i}=i t / n, i=0, \ldots, n$, we write

$$
\begin{align*}
G_{2} \leq & \sum_{i=1}^{n}\left\|\int_{\tau_{i-1}}^{\tau_{i}}\left(u\left(s, \phi_{s}^{\hat{u}}\right)-u\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
& +\sum_{i=1}^{n}\left\|\int_{\tau_{i-1}}^{\tau_{i}}\left(\hat{u}\left(s, \phi_{s}^{\hat{u}}\right)-\hat{u}\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
& +\sum_{i=1}^{n}\left\|\int_{\tau_{i-1}}^{\tau_{i}}\left(u\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)-\hat{u}\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
= & G_{2,1}+G_{2,2}+G_{2,3} . \tag{1.10}
\end{align*}
$$

To estimate $G_{2,1}+G_{2,2}$, we use (1.9):

$$
\begin{equation*}
G_{2,1}+G_{2,2} \leq \frac{T^{2}}{n}\|\hat{u}\|_{L^{\infty}\left(J_{T} \times \mathbb{R}^{d}\right)}\left(\|u\|_{L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{R}^{d}\right)\right)}+\|\hat{u}\|_{L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{R}^{d}\right)\right)}\right) \tag{1.11}
\end{equation*}
$$

We use the relaxation norm defined by $(0.9)$ to bound $G_{2,3}$ :

$$
\begin{aligned}
\left|\int_{\tau_{i-1}}^{\tau_{i}}\left(u\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)-\hat{u}\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)\right) \mathrm{d} s\right| \leq & \left|\int_{0}^{\tau_{i}}\left(u\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)-\hat{u}\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)\right) \mathrm{d} s\right| \\
& +\left|\int_{0}^{\tau_{i-1}}\left(u\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)-\hat{u}\left(s, \phi_{\tau_{i-1}}^{\hat{u}}\right)\right) \mathrm{d} s\right| \\
\leq & 2\|u-\hat{u}\|_{T, L^{\infty}\left(\mathbb{R}^{d}\right)},
\end{aligned}
$$

hence

$$
G_{2,3} \leq 2 n\|u-\hat{u}\|_{T, L^{\infty}\left(\mathbb{R}^{d}\right)} .
$$

Combining this with (1.10) and (1.11), we get

$$
\begin{align*}
G_{2} \leq & \frac{T^{2}}{n}\|\hat{u}\|_{L^{\infty}\left(J_{T} \times \mathbb{R}^{d}\right)}\left(\|u\|_{L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{R}^{d}\right)\right)}+\|\hat{u}\|_{L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{R}^{d}\right)\right)}\right) \\
& +2 n\|u-\hat{u}\|_{T, L^{\infty}\left(\mathbb{R}^{d}\right)} . \tag{1.12}
\end{align*}
$$

If $\|u-\hat{u}\|_{T, L^{\infty}\left(\mathbb{R}^{d}\right)}=0$, then (1.6) is trivial. Assume that $\|u-\hat{u}\|_{T, L^{\infty}\left(\mathbb{R}^{d}\right)}>0$. Choosing ${ }^{1} n:=\left[\|u-\hat{u}\|_{T, L^{\infty}\left(\mathbb{R}^{d}\right)}^{-1 / 2}\right]$, we derive from (1.12) and (1.4) that

$$
G_{2} \leq C\|u-\hat{u}\|_{T, L^{\infty}\left(\mathbb{R}^{d}\right)}^{1 / 2}
$$

Combining this with (1.7) and (1.8) and applying the Gronwall inequality, we obtain (1.6).

Step 2. We turn to the proof of (1.3). It is easy to verify that there is a constant $C_{1}:=C_{1}(R, T)>0$ such that

$$
\begin{equation*}
\left\|\phi^{\hat{u}}\right\|_{L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{R}^{d}\right)\right)}+\left\|\partial_{t} \phi^{\hat{u}}\right\|_{L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{R}^{d}\right)\right)} \leq C_{1} . \tag{1.13}
\end{equation*}
$$

For $j=1, \ldots, d$, we have

$$
\begin{align*}
\left\|\partial_{j} \phi_{t}^{u}-\partial_{j} \phi_{t}^{\hat{u}}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}= & \left\|\int_{0}^{t}\left(\left\langle\nabla u\left(s, \phi_{s}^{u}\right), \partial_{j} \phi_{s}^{u}\right\rangle-\left\langle\nabla \hat{u}\left(s, \phi_{s}^{\hat{u}}\right), \partial_{j} \phi_{s}^{\hat{u}}\right\rangle\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
\leq & \left\|\int_{0}^{t}\left(\left\langle\nabla u\left(s, \phi_{s}^{u}\right), \partial_{j} \phi_{s}^{u}-\partial_{j} \phi_{s}^{\hat{u}}\right\rangle\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
& +\left\|\int_{0}^{t}\left(\left\langle\nabla u\left(s, \phi_{s}^{u}\right)-\nabla u\left(s, \phi_{s}^{\hat{u}}\right), \partial_{j} \phi_{s}^{\hat{u}}\right\rangle\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
& +\left\|\int_{0}^{t}\left(\left\langle\nabla u\left(s, \phi_{s}^{\hat{u}}\right)-\nabla \hat{u}\left(s, \phi_{s}^{\hat{u}}\right), \partial_{j} \phi_{s}^{\hat{u}}\right\rangle\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
= & I_{1}+I_{2}+I_{3} . \tag{1.14}
\end{align*}
$$

From (1.5) it follows that

$$
I_{1} \leq R \int_{0}^{t}\left\|\partial_{j} \phi_{s}^{u}-\partial_{j} \phi_{s}^{\hat{u}}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \mathrm{d} s .
$$

Using (1.5), (1.13), and (1.6), we get

$$
I_{2} \leq C_{1} R \int_{0}^{t}\left\|\phi_{s}^{u}-\phi_{s}^{\hat{u}}\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}^{\lambda} \mathrm{d} s \leq C_{2}\|u-\hat{u}\|_{T, L^{\infty}\left(\mathbb{R}^{d}\right)}^{\lambda / 2}
$$

[^1]To estimate $I_{3}$, we integrate by parts and use (1.13)

$$
\begin{aligned}
I_{3} \leq & \left\|\left\langle\int_{0}^{t}\left(\nabla u\left(s, \phi_{s}^{\hat{u}}\right)-\nabla \hat{u}\left(s, \phi_{s}^{\hat{u}}\right)\right) \mathrm{d} s, \partial_{j} \phi_{t}^{\hat{u}}\right\rangle\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
& +\left\|\int_{0}^{t}\left(\left\langle\int_{0}^{s}\left(\nabla u\left(\theta, \phi_{\theta}^{\hat{u}}\right)-\nabla \hat{u}\left(\theta, \phi_{\theta}^{\hat{u}}\right)\right) \mathrm{d} \theta, \partial_{j}\left(\partial_{t} \phi_{s}^{\hat{u}}\right)\right\rangle\right) \mathrm{d} s\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} \\
\leq & C_{3} \sup _{s \in[0, t]}\left\|\int_{0}^{s}\left(\nabla u\left(\theta, \phi_{\theta}^{\hat{u}}\right)-\nabla \hat{u}\left(\theta, \phi_{\theta}^{\hat{u}}\right)\right) \mathrm{d} \theta\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} .
\end{aligned}
$$

Repeating the arguments of the proof of (1.12) and using the fact that $\nabla u$ and $\nabla \hat{u}$ are Hölder continuous with exponent $\lambda$, we obtain that

$$
\begin{aligned}
\sup _{s \in[0, t]}\left\|\int_{0}^{s}\left(\nabla u\left(\theta, \phi_{\theta}^{\hat{u}}\right)-\nabla \hat{u}\left(\theta, \phi_{\theta}^{\hat{u}}\right)\right) \mathrm{d} \theta\right\|_{L^{\infty}\left(\mathbb{R}^{d}\right)} & \leq \frac{C_{4}}{n^{\lambda}}+2 n\|u-\hat{u}\|_{T, C^{1}\left(\mathbb{R}^{d}\right)} \\
& \leq C_{5}\|u-\hat{u}\|_{T, C^{1}\left(\mathbb{R}^{d}\right)}^{\lambda / 2}
\end{aligned}
$$

for $n:=\left[\|u-\hat{u}\|_{T, C^{1}\left(\mathbb{R}^{d}\right)}^{-1 / 2}\right]$. Combining this with the estimates for $I_{1}, I_{2}$ and (1.14), and applying the Gronwall inequality, we arrive at the required result.

By the Liouville theorem (see Corollary 1 in [Arn78, p.198]), if we assume additionally that $u$ is divergence-free, then the flow $\phi_{t}^{u}$ preserves the orientation and the volume. Thus if $u \in L^{\infty}\left(J_{T}, C_{\sigma}^{1}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right.$ ), then $\phi_{t}^{u} \in \operatorname{SDiff}\left(\mathbb{T}^{d}\right)$ for any $t \in J_{T}$. The following proposition shows that, using a suitable divergencefree field $u$, the flow $\phi_{t}^{u}$ can be driven approximately to any position $\psi \in$ $\operatorname{SDiff}\left(\mathbb{T}^{d}\right)$ at time $T$.
Proposition 1.2. For any $\varepsilon>0, k>1+d / 2, u_{0}, u_{1} \in H_{\sigma}^{k}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)$, and $\psi \in \operatorname{SDiff}\left(\mathbb{T}^{d}\right)$, there is a vector field $u \in C^{\infty}\left(J_{T}, H_{\sigma}^{k}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right)$ such that $u(0)=$ $u_{0}, u(T)=u_{1}$, and

$$
\left\|\phi_{T}^{u}-\psi\right\|_{C^{1}\left(\mathbb{T}^{d}\right)}<\varepsilon
$$

Proof. Step 1. We first forget about the endpoint conditions $u(0)=u_{0}, u(T)=$ $u_{1}$ and show that there is a divergence-free vector field $\hat{u} \in C^{\infty}\left(\mathbb{R} \times \mathbb{T}^{d}, \mathbb{R}^{d}\right)$ such that

$$
\left\|\phi_{T}^{\hat{u}}-\psi\right\|_{C^{1}\left(\mathbb{T}^{d}\right)}<\varepsilon / 2 .
$$

Since $\psi \in \operatorname{SDiff}\left(\mathbb{T}^{d}\right)$, there is a path $I \in W^{1, \infty}\left(J_{T}, C^{1}\left(\mathbb{T}^{d}\right)\right)$ such that $I(0, x)=$ $x, I(T, x)=\psi(x)$ for all $x \in \mathbb{T}^{d}$, and $I(t, \cdot)$ is a $C^{1}$-diffeomorphism on $\mathbb{T}^{d}$ preserving the orientation and the volume for all $t \in J_{T}$. Let us define the vector field $\hat{u}(t, x)=\partial_{t} I\left(t, I^{-1}(t, x)\right)$. Then we have $\hat{u} \in L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right)$ and $I(t, x)=\phi_{t}^{\hat{u}}(x), t \in J_{T}$. As $\phi_{t}^{\hat{u}}$ preserves the orientation and the volume, for any $g \in C^{1}\left(\mathbb{T}^{d}, \mathbb{R}\right)$

$$
\begin{aligned}
0 & =\frac{\mathrm{d}}{\mathrm{~d} t} \int_{\mathbb{T}^{d}} g\left(\phi_{t}^{\hat{u}}(y)\right) \mathrm{d} y=\int_{\mathbb{T}^{d}}\left\langle\nabla g\left(\phi_{t}^{\hat{u}}(y)\right), \partial_{t} \phi_{t}^{\hat{u}}(y)\right\rangle \mathrm{d} y \\
& =\int_{\mathbb{T}^{d}}\left\langle\nabla g\left(\phi_{t}^{\hat{u}}(y)\right), \hat{u}\left(t, \phi_{t}^{\hat{u}}(y)\right)\right\rangle \mathrm{d} y=\int_{\mathbb{T}^{d}} g(y) \operatorname{div} \hat{u}(t, y) \mathrm{d} y .
\end{aligned}
$$

This shows that $\hat{u}$ is divergence-free. Taking a sequence of mollifying kernels $\rho_{n} \in C^{\infty}\left(\mathbb{R} \times \mathbb{T}^{d}, \mathbb{R}\right), n \geq 1$, we consider $\hat{u}_{n}:=\rho_{n} * \hat{u}=\left(\rho_{n} * \hat{u}_{1}, \ldots, \rho_{n} * \hat{u}_{d}\right) \in$ $C^{\infty}\left(\mathbb{R} \times \mathbb{T}^{d}, \mathbb{R}^{d}\right)$. Then $\hat{u}_{n}$ is also divergence-free, since $\operatorname{div} \hat{u}_{n}=\rho_{n} * \operatorname{div} \hat{u}=0$, and $\left\|\hat{u}_{n}-\hat{u}\right\|_{L^{\infty}\left(J_{T}, C^{1}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right)} \rightarrow 0$ as $n \rightarrow \infty$. By (1.2), this implies that $\left\|\phi_{T}^{\hat{u}_{n}}-\phi_{T}^{\hat{u}}\right\|_{C^{1}\left(\mathbb{T}^{d}\right)} \rightarrow 0$ as $n \rightarrow \infty$. Since $\phi_{T}^{\hat{u}}=\psi$, we get the required result with $\hat{u}=\hat{u}_{n}$ for sufficiently large $n \geq 1$.

Step 2. By the Sobolev embedding, $H^{k} \subset C^{1}\left(\mathbb{T}^{d}\right)$ for $k>1+d / 2$ (e.g., see [Ada75]). For any $\delta>0$, we take an arbitrary $u \in C^{\infty}\left(J_{T}, H_{\sigma}^{k}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right)$ satisfying

$$
\begin{aligned}
& u(0)=u_{0}, \quad u(T)=u_{1} \\
& \|u-\hat{u}\|_{L^{1}\left(J_{T}, C^{1}\left(\mathbb{T}^{d}, \mathbb{R}^{d}\right)\right)}<\delta
\end{aligned}
$$

Then by Step 1 and (1.2), we have

$$
\begin{aligned}
\left\|\phi_{T}^{u}-\psi\right\|_{C^{1}\left(\mathbb{T}^{d}\right)} & \leq\left\|\phi_{T}^{u}-\phi_{T}^{\hat{u}}\right\|_{C^{1}\left(\mathbb{T}^{d}\right)}+\left\|\phi_{T}^{\hat{u}}-\psi\right\|_{C^{1}\left(\mathbb{T}^{d}\right)} \\
& <\varepsilon / 2+\varepsilon / 2=\varepsilon
\end{aligned}
$$

for sufficiently small $\delta>0$.

### 1.2 Existence of strong solutions

In what follows, we shall assume that $d=3, k \geq 3$, and $\nu=1$. In this section, we prove a perturbative result on existence of strong solutions for the evolution equation

$$
\begin{equation*}
\dot{u}+L u+B(u)=g \tag{1.15}
\end{equation*}
$$

where $B(a, b):=\Pi\{\langle a, \nabla\rangle b\}$ and $B(a):=B(a, a)$. Along with (1.15), we consider the following more general equation

$$
\begin{equation*}
\dot{u}+L(u+\zeta)+B(u+\zeta)=g \tag{1.16}
\end{equation*}
$$

Let us fix any $T>0$ and introduce the space $\mathcal{X}_{T, k}:=C\left(J_{T}, H_{\sigma}^{k}\right) \cap L^{2}\left(J_{T}, H_{\sigma}^{k+1}\right)$ endowed with the norm

$$
\|u\|_{\mathcal{X}_{T, k}}:=\|u\|_{L^{\infty}\left(J_{T}, H^{k}\right)}+\|u\|_{L^{2}\left(J_{T}, H^{k+1}\right)} .
$$

The following result is a version of Theorem 1.8 and Remark 1.9 in [Shi06] and Theorem 2.1 in [Ner10] in the case of the 3D NS system in the spaces $H^{k}, k \geq 3$. For the sake of completeness, we give all the details of the proof, even though it is very close to the proofs of the previous results.
Theorem 1.3. Suppose that for some functions $\hat{u}_{0} \in H_{\sigma}^{k}, \hat{\zeta} \in L^{4}\left(J_{T}, H_{\sigma}^{k+1}\right)$, and $\hat{g} \in L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right)$ problem (1.16), (0.2) with $u_{0}=\hat{u}_{0}, \zeta=\hat{\zeta}$, and $g=\hat{g}$ has a solution $\hat{u} \in \mathcal{X}_{T, k}$. Then there are positive constants $\delta$ and $C$ depending only on

$$
\|\hat{\zeta}\|_{L^{4}\left(J_{T}, H^{k+1}\right)}+\|\hat{g}\|_{L^{2}\left(J_{T}, H^{k-1}\right)}+\|\hat{u}\|_{\mathcal{X}_{T, k}}
$$

such that the following statements hold.
(i) If $u_{0} \in H_{\sigma}^{k}, \zeta \in L^{4}\left(J_{T}, H_{\sigma}^{k+1}\right)$, and $g \in L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right)$ satisfy the inequality

$$
\begin{equation*}
\left\|u_{0}-\hat{u}_{0}\right\|_{k}+\|\zeta-\hat{\zeta}\|_{L^{4}\left(J_{T}, H^{k+1}\right)}+\|g-\hat{g}\|_{L^{2}\left(J_{T}, H^{k-1}\right)}<\delta \tag{1.17}
\end{equation*}
$$

then problem (1.16), (0.2) has a unique solution $u \in \mathcal{X}_{T, k}$.
(ii) Let

$$
\mathcal{R}: H_{\sigma}^{k} \times L^{4}\left(J_{T}, H_{\sigma}^{k+1}\right) \times L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right) \rightarrow \mathcal{X}_{T, k}
$$

be the operator that takes each triple $\left(u_{0}, \zeta, g\right)$ satisfying (1.17) to the solution $u$ of (1.16), (0.2). Then

$$
\begin{aligned}
\| \mathcal{R}\left(u_{0}, \zeta, g\right) & -\mathcal{R}\left(\hat{u}_{0}, \hat{\zeta}, \hat{g}\right) \|_{\mathcal{X}_{T, k}} \leq C\left(\left\|u_{0}-\hat{u}_{0}\right\|_{k}\right. \\
& \left.+\|\zeta-\hat{\zeta}\|_{L^{4}\left(J_{T}, H^{k+1}\right)}+\|g-\hat{g}\|_{L^{2}\left(J_{T}, H^{k-1}\right)}\right)
\end{aligned}
$$

Proof. We use the following standard estimates for the bilinear form $B$

$$
\begin{align*}
\|B(a, b)\|_{k} \leq C\|a\|_{k}\|b\|_{k+1} & \text { for } k \geq 2  \tag{1.18}\\
\left|\left(B(a, b), L^{k} b\right)\right| \leq C\|a\|_{k}\|b\|_{k}^{2} & \text { for } k \geq 3 \tag{1.19}
\end{align*}
$$

for any $a \in H_{\sigma}^{k}$ and $b \in H_{\sigma}^{k+1}$ (see Chapter 6 in [CF88]). We are looking for a solution of (1.16), (0.2) of the form $u=\hat{u}+w$. We have the following equation for $w$ :

$$
\begin{align*}
& \dot{w}+L(w+\eta)+B(w+\eta, \hat{u}+\hat{\zeta})+B(\hat{u}+\hat{\zeta}, w+\eta)+B(w+\eta)=q \\
& w(0, x)=w_{0}(x) \tag{1.20}
\end{align*}
$$

where $w_{0}:=u_{0}-\hat{u}_{0}, \eta:=\zeta-\hat{\zeta}$, and $q:=g-\hat{g}$. Setting $\tilde{B}(u, v):=B(u, v)+$ $B(v, u)$, we get that
$\dot{w}+L w+B(w)+\tilde{B}(w, \eta)+\tilde{B}(w, \hat{u})+\tilde{B}(w, \hat{\zeta})=q-(L \eta+B(\eta)+\tilde{B}(\hat{u}, \eta)+\tilde{B}(\hat{\zeta}, \eta))$,
Using (1.18), we see that for any $\varepsilon>0$, we can choose $\delta \in(0,1)$ in (1.17) such that

$$
\left\|w_{0}\right\|_{k}+\|q-(L \eta+B(\eta)+\tilde{B}(\hat{u}, \eta)+\tilde{B}(\hat{\zeta}, \eta))\|_{L^{2}\left(J_{T}, H^{k-1}\right)}<\varepsilon
$$

Then, using some standard methods, one gets that system (1.21), (1.20) has a solution $w \in \mathcal{X}_{T, k}$ for sufficiently small $\varepsilon>0$ (see Section 4 of Chapter 17 in [Tay96]).

To prove (ii), we multiply (1.21) by $L^{k} w$ and use estimates (1.18) and (1.19)

$$
\begin{aligned}
\frac{1}{2} \frac{d}{d t}\|w\|_{k}^{2}+ & \|w\|_{k+1}^{2} \leq C\left(\|w\|_{k}^{3}+\|w\|_{k+1}\|w\|_{k}\left(\|\eta\|_{k}+\|\hat{u}\|_{k}+\|\hat{\zeta}\|_{k}\right)\right. \\
& \left.+\|w\|_{k+1}\left(\|q\|_{k-1}+\|\eta\|_{k+1}+\|\eta\|_{k}\left(\|\eta\|_{k-1}+\|\hat{u}\|_{k}+\|\hat{\zeta}\|_{k}\right)\right)\right)
\end{aligned}
$$

This implies that

$$
\begin{gathered}
\frac{1}{2} \frac{d}{d t}\|w\|_{k}^{2}+\frac{1}{2}\|w\|_{k+1}^{2} \leq C_{1}\left(\|w\|_{k}^{3}+\|w\|_{k}^{2}\left(\|\eta\|_{k}^{2}+\|\hat{u}\|_{k}^{2}+\|\hat{\zeta}\|_{k}^{2}\right)\right. \\
\left.+\left[\|q\|_{k-1}^{2}+\|\eta\|_{k+1}^{4}+\|\eta\|_{k}^{2}\left(\|\hat{u}\|_{k}^{2}+\|\hat{\zeta}\|_{k}^{2}\right)\right]\right)
\end{gathered}
$$

Integrating this inequality and setting

$$
A:=\left\|w_{0}\right\|_{k}^{2}+\int_{0}^{T}\left[\|q\|_{k-1}^{2}+\|\eta\|_{k+1}^{4}+\|\eta\|_{k}^{2}\left(\|\hat{u}\|_{k}^{2}+\|\hat{\zeta}\|_{k}^{2}\right)\right] \mathrm{d} t
$$

we obtain

$$
\begin{equation*}
\|w\|_{k}^{2}+\int_{0}^{t}\|w\|_{k+1}^{2} \leq 2 A+2 C_{1} \int_{0}^{t}\left(\|w\|_{k}^{3}+\|w\|_{k}^{2}\left(\|\eta\|_{k}^{2}+\|\hat{u}\|_{k}^{2}+\|\hat{\zeta}\|_{k}^{2}\right)\right) \mathrm{d} t \tag{1.22}
\end{equation*}
$$

By (1.17), we have that $\|\eta\|_{L^{4}\left(J_{T}, H^{k+1}\right)} \leq \delta<1$. So the Gronwall inequality gives that

$$
\begin{equation*}
\|w\|_{k}^{2} \leq C_{2} A+C_{2} \int_{0}^{t}\|w(s, \cdot)\|_{k}^{3} \mathrm{~d} s, \quad t \in J_{T} \tag{1.23}
\end{equation*}
$$

where $C_{2}>0$ depends only on $\|\hat{u}\|_{L^{2}\left(J_{T}, H^{k}\right)}+\|\hat{\zeta}\|_{L^{2}\left(J_{T}, H^{k}\right)}$. Let us denote

$$
\Phi(t):=A+\int_{0}^{t}\|w(s, \cdot)\|_{k}^{3} \mathrm{~d} s, \quad t \in J_{T}
$$

Since the case $A=0$ is trivial, we can assume that $A>0$ and $\Phi(t)>0$ for all $t \in J_{T}$. Thus (1.23) can be written as

$$
(\dot{\Phi}(t))^{2 / 3} \leq C_{2} \Phi(t)
$$

which is equivalent to

$$
\frac{\dot{\Phi}(t)}{(\Phi(t))^{3 / 2}} \leq C_{3}, \quad C_{3}:=C_{2}^{3 / 2}
$$

Integrating this inequality, we get

$$
\begin{equation*}
\Phi(t) \leq \frac{A}{\left(1-t C_{3} \sqrt{A} / 2\right)^{2}} \leq 4 A \text { for any } t \leq \frac{1}{C_{3} \sqrt{A}} \tag{1.24}
\end{equation*}
$$

Choosing $\delta>0$ so small that $\frac{1}{C_{3} \sqrt{A}} \geq T$ and using (1.23) and (1.24), we obtain

$$
\|w(t)\|_{k}^{2} \leq C_{2} \Phi(t) \leq 4 C_{2} A \text { for any } t \in J_{T}
$$

Combining this with (1.22), we get for any $t \in J_{T}$

$$
\|w\|_{k}^{2}+\int_{0}^{t}\|w\|_{k+1}^{2} \leq C_{4} A \leq C_{5}\left(\left\|w_{0}\right\|_{k}^{2}+\|\eta\|_{L^{4}\left(J_{T}, H^{k+1}\right)}^{2}+\|q\|_{L^{2}\left(J_{T}, H^{k-1}\right)}^{2}\right)
$$

This completes the proof of the theorem.

## 2 Approximate controllability of the NS system

In this section, we state the main results of this paper. Let us fix any $T>0$ and $k \geq 3$, and consider the NS system

$$
\begin{align*}
\dot{u}+L u+B(u) & =h(t)+\eta(t)  \tag{2.1}\\
u(0, x) & =u_{0}(x) \tag{2.2}
\end{align*}
$$

where $h \in L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right)$ and $u_{0} \in H_{\sigma}^{k}$ are given functions and $\eta$ is a control taking values in a finite-dimensional space $E \subset H_{\sigma}^{k+1}$. We denote by $\Theta\left(h, u_{0}\right)$ the set of functions $\eta \in L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right)$ for which (2.1), (2.2) has a solution $u$ in $\mathcal{X}_{T, k}$. By Theorem $1.3, \Theta\left(h, u_{0}\right)$ is an open subset of $L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right)$. Recall that $\mathcal{R}(\cdot, \cdot, \cdot)$ is the operator defined in Theorem 1.3. To simplify notation, we write $\mathcal{R}\left(u_{0}, h+\eta\right)$ instead of $\mathcal{R}\left(u_{0}, 0, h+\eta\right)$ for any $\eta \in \Theta\left(h, u_{0}\right)$. The embedding $H^{3} \subset C^{1,1 / 2}$ implies that the flow $\phi_{t}^{\mathcal{R}\left(u_{0}, h+\eta\right)}$ is well defined for any $t \in J_{T}$. We set

$$
Y_{T, k}:=\mathcal{X}_{T, k} \cap W^{1,2}\left(J_{T}, H_{\sigma}^{k-1}\right)
$$

We shall use the following notion of controllability.
Definition 2.1. Equation (2.1) is said to be approximately controllable at time $T$ by an $E$-valued control if for any $\varepsilon>0$ and any $\varphi \in Y_{T, k}$ there is a control $\eta \in \Theta\left(h, u_{0}\right) \cap L^{2}\left(J_{T}, E\right)$ such that
$\left\|\mathcal{R}_{T}\left(u_{0}, h+\eta\right)-\varphi(T)\right\|_{k}+\left\|\mathcal{R}\left(u_{0}, h+\eta\right)-\varphi\right\|_{T, k}+\left\|\phi^{\mathcal{R}\left(u_{0}, h+\eta\right)}-\phi^{\varphi}\right\|_{L^{\infty}\left(J_{T}, C^{1}\right)}<\varepsilon$,
where $u_{0}=\varphi(0)$ and $\||\cdot|\|_{T, k}:=\|\mid \cdot\|_{T, H^{k}}$.
Let us recall some notation introduced in [AS05, AS06], and [Shi06]. For any finite-dimensional subspace $E \subset H_{\sigma}^{k+1}$, we denote by $\mathcal{F}(E)$ the largest vector space $F \subset H_{\sigma}^{k+1}$ such that for any $\eta_{1} \in F$ there are vectors ${ }^{2} \eta, \zeta^{1}, \ldots, \zeta^{p} \in E$ satisfying the relation

$$
\begin{equation*}
\eta_{1}=\eta-\sum_{i=1}^{p} B\left(\zeta^{i}\right) \tag{2.4}
\end{equation*}
$$

As $E$ is a finite-dimensional subspace and $B$ is a bilinear operator, the set of all vectors $\eta_{1} \in H_{\sigma}^{k+1}$ of the form (2.4) is contained in a finite-dimensional space. It is easy to see that if subspaces $G_{1}, G_{2} \subset H_{\sigma}^{k+1}$ are composed of elements $\eta_{1}$ of the form (2.4), then so does $G_{1}+G_{2}$. Thus the space $\mathcal{F}(E)$ is well defined. We define $E_{j}$ by the rule

$$
\begin{equation*}
E_{0}=E, \quad E_{j}=\mathcal{F}\left(E_{j-1}\right) \quad \text { for } \quad j \geq 1, \quad E_{\infty}=\bigcup_{j=1}^{\infty} E_{j} \tag{2.5}
\end{equation*}
$$

Clearly, $E_{j}$ is a non-decreasing sequence of subspaces. We say that $E$ is saturating in $H_{\sigma}^{k-1}$ if $E_{\infty}$ is dense in $H_{\sigma}^{k-1}$. The following theorem is the main result of this paper.

[^2]Theorem 2.2. Assume that $E$ is a finite-dimensional subspace of $H_{\sigma}^{k+1}$ and $h \in L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right)$. If $E$ is saturating in $H_{\sigma}^{k-1}$, then (2.1) is approximately controllable at time $T$ by controls $\eta \in C^{\infty}\left(J_{T}, E\right)$ in the sense of Definition 2.1.

We have the following two corollaries of this result.
Corollary 2.3. Under the conditions of Theorem 2.2, if $E$ is saturating in $H_{\sigma}^{k-1}$, then for any $\varepsilon>0, u_{0}, u_{1} \in H_{\sigma}^{k}$, and $\psi \in \operatorname{SDiff}\left(\mathbb{T}^{3}\right)$ there is a control $\eta \in$ $\Theta\left(h, u_{0}\right) \cap C^{\infty}\left(J_{T}, E\right)$ such that

$$
\left\|\mathcal{R}_{T}\left(u_{0}, h+\eta\right)-u_{1}\right\|_{k}+\left\|\phi_{T}^{\mathcal{R}\left(u_{0}, h+\eta\right)}-\psi\right\|_{C^{1}}<\varepsilon
$$

Let us denote by $\operatorname{VPM}\left(\mathbb{T}^{3}\right)$ the set of all volume-preserving mappings from $\mathbb{T}^{3}$ to $\mathbb{T}^{3}$. According ${ }^{3}$ to Corollary 1.1 in [BG03], we have that $\operatorname{VPM}\left(\mathbb{T}^{3}\right)$ is the closure of $\operatorname{SDiff}\left(\mathbb{T}^{3}\right)$ in $L^{p}\left(\mathbb{T}^{3}\right)$ for any $p \in[1,+\infty)$. Thus we get the following result.
Corollary 2.4. Under the conditions of Theorem 2.2, if $E$ is saturating in $H_{\sigma}^{k-1}$, then for any $\varepsilon>0, p \in[1,+\infty)$, $u_{0}, u_{1} \in H_{\sigma}^{k}$, and $\psi \in \operatorname{VPM}\left(\mathbb{T}^{3}\right)$ there is a control $\eta \in \Theta\left(h, u_{0}\right) \cap C^{\infty}\left(J_{T}, E\right)$ such that

$$
\left\|\mathcal{R}_{T}\left(u_{0}, h+\eta\right)-u_{1}\right\|_{k}+\left\|\phi_{T}^{\mathcal{R}\left(u_{0}, h+\eta\right)}-\psi\right\|_{L^{p}}<\varepsilon .
$$

The rest of this subsection is devoted to the proofs of Theorem 2.2 and Corollary 2.3. They are based on the following result which is proved in Section 3.
Theorem 2.5. Assume that $E$ is an arbitrary finite-dimensional subspace of $H_{\sigma}^{k+1}$ and $h \in L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right)$. Then for any $\varepsilon>0, u_{0} \in H_{\sigma}^{k}$, and $\eta_{1} \in \Theta\left(h, u_{0}\right) \cap$ $L^{2}\left(J_{T}, E_{1}\right)$ there is $\eta \in \Theta\left(h, u_{0}\right) \cap C^{\infty}\left(J_{T}, E\right)$ such that

$$
\begin{aligned}
\| \mathcal{R}_{T}\left(u_{0}, h+\eta_{1}\right) & -\mathcal{R}_{T}\left(u_{0}, h+\eta\right)\left\|_{k}+\right\| \mathcal{R}\left(u_{0}, h+\eta_{1}\right)-\mathcal{R}\left(u_{0}, h+\eta\right) \|_{T, k} \\
& +\left\|\phi^{\mathcal{R}\left(u_{0}, h+\eta_{1}\right)}-\phi^{\mathcal{R}\left(u_{0}, h+\eta\right)}\right\|_{L^{\infty}\left(J_{T}, C^{1}\right)}<\varepsilon
\end{aligned}
$$

Proof of Theorem 2.2. Let us take any $\varepsilon>0, \delta>0$, and $\varphi \in Y_{T, k}$. Then

$$
\eta_{0}:=\dot{\varphi}+L \varphi+B(\varphi)-h
$$

belongs to $\Theta\left(u_{0}, h\right)$ and $\varphi(t)=\mathcal{R}_{t}\left(u_{0}, h+\eta_{0}\right)$ for any $t \in J_{T}$, where $u_{0}=\varphi(0)$. Since $E_{\infty}$ is dense in $H_{\sigma}^{k-1}$, we have that

$$
\left\|P_{E_{N}} \eta_{0}-\eta_{0}\right\|_{L^{2}\left(J_{T}, H^{k-1}\right)} \rightarrow 0 \text { as } N \rightarrow \infty
$$

where $P_{E_{N}}$ is the orthogonal projection onto $E_{N}$ in $H$. By Theorem 1.3, for sufficiently large $N$, we have $P_{E_{N}} \eta_{0} \in \Theta\left(h, u_{0}\right)$ and

$$
\left\|\mathcal{R}\left(u_{0}, h+P_{E_{N}} \eta_{0}\right)-\varphi\right\|_{X_{T, k}}<\delta
$$

By (1.2), we can choose $\delta>0$ so small that

$$
\left\|\phi^{\mathcal{R}\left(u_{0}, h+P_{E_{N}} \eta_{0}\right)}-\phi^{\varphi}\right\|_{L^{\infty}\left(J_{T}, C^{1}\right)}<\varepsilon .
$$

Applying $N$ times Theorem 2.5, we complete the proof of Theorem 2.2.

[^3]Proof of Corollary 2.3. Let us take any $\varepsilon>0, \psi \in \operatorname{SDiff}\left(\mathbb{T}^{3}\right)$, and $u_{0}, u_{1} \in H_{\sigma}^{k}$. By Proposition 1.2, there is a vector field $u \in C^{\infty}\left(J_{T}, H_{\sigma}^{k}\right)$ such that $u(0)=$ $u_{0}, u(T)=u_{1}$, and

$$
\begin{equation*}
\left\|\phi_{T}^{u}-\psi\right\|_{C^{1}}<\varepsilon \tag{2.6}
\end{equation*}
$$

Applying Theorem 2.2, we find a control $\eta \in \Theta\left(h, u_{0}\right) \cap C^{\infty}\left(J_{T}, E\right)$ such that (2.3) holds with $\varphi=u$. In particular,

$$
\left\|\mathcal{R}_{T}\left(u_{0}, h+\eta\right)-u(T)\right\|_{k}+\left\|\phi_{T}^{\mathcal{R}\left(u_{0}, h+\eta\right)}-\phi_{T}^{u}\right\|_{C^{1}}<\varepsilon
$$

Combining this with (2.6), we get the required result.

## 3 Proof Theorem 2.5

The proof follows the arguments of [AS05, AS06], and [Shi06]. We consider the following system

$$
\begin{equation*}
\dot{u}+L(u+\zeta)+B(u+\zeta)=h+\eta \tag{3.1}
\end{equation*}
$$

with two $E$-valued controls $\eta, \zeta$. We denote by $\hat{\Theta}\left(u_{0}, h\right)$ the set of $(\eta, \zeta) \in$ $L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right) \times L^{4}\left(J_{T}, H_{\sigma}^{k+1}\right)$ for which problem (3.1), (0.2) has a solution in $\mathcal{X}_{T, k}$. Theorem 2.5 is deduced from the following proposition which is proved at the end of this section (cf. Proposition 3.2 in [Shi06]).

Proposition 3.1. For any $\eta_{1} \in \Theta\left(u_{0}, h\right) \cap L^{2}\left(J_{T}, E_{1}\right)$, there is a sequence $\left(\eta_{n}, \zeta_{n}\right) \in \hat{\Theta}\left(u_{0}, h\right) \cap C^{\infty}\left(J_{T}, E \times E\right)$ such that

$$
\begin{align*}
& \left\|\mathcal{R}\left(u_{0}, 0, h+\eta_{1}\right)-\mathcal{R}\left(u_{0}, \zeta_{n}, h+\eta_{n}\right)\right\|_{L^{\infty}\left(J_{T}, H^{k}\right)}+\left\|\zeta_{n}\right\|_{T, k} \rightarrow 0 \text { as } n \rightarrow \infty  \tag{3.2}\\
& \sup _{n \geq 1}\left(\left\|\mathcal{R}\left(u_{0}, \zeta_{n}, h+\eta_{n}\right)\right\|_{\mathcal{X}_{T, k}}+\left\|\zeta_{n}\right\|_{L^{\infty}\left(J_{T}, H^{k+1}\right)}+\left\|\eta_{n}\right\|_{L^{2}\left(J_{T}, H^{k-1}\right)}\right)<\infty \tag{3.3}
\end{align*}
$$

Proof of Theorem 2.5. Let us take any $u_{0} \in H_{\sigma}^{k}$ and $\eta_{1} \in \Theta\left(h, u_{0}\right) \cap L^{2}\left(J_{T}, E_{1}\right)$, and let $\left(\eta_{n}, \zeta_{n}\right) \in \hat{\Theta}\left(u_{0}, h\right) \cap C^{\infty}\left(J_{T}, E \times E\right)$ be any sequence satisfying (3.2) and (3.3). Let $\hat{\zeta}_{n} \in C^{\infty}\left(J_{T}, E\right)$ be such that $\hat{\zeta}_{n}(0)=\hat{\zeta}_{n}(T)=0$ and

$$
\begin{align*}
& \left\|\zeta_{n}-\hat{\zeta}_{n}\right\|_{L^{4}\left(J_{T}, H^{k+1}\right)} \rightarrow 0 \text { as } n \rightarrow \infty  \tag{3.4}\\
& \sup _{n \geq 1}\left\|\hat{\zeta}_{n}\right\|_{L^{\infty}\left(J_{T}, H^{k+1}\right)}<+\infty \tag{3.5}
\end{align*}
$$

Then (3.2) and (3.4) imply that

$$
\begin{align*}
\left\|\hat{\zeta}_{n}\right\|_{T, k} & \leq\left\|\hat{\zeta}_{n}-\zeta_{n}\right\|_{T, k}+\left\|\zeta_{n}\right\|_{T, k} \\
& \leq \int_{0}^{T}\left\|\hat{\zeta}_{n}(s)-\zeta_{n}(s)\right\|_{k} \mathrm{~d} s+\left\|\zeta_{n}\right\|_{T, k} \rightarrow 0 \quad \text { as } n \rightarrow \infty \tag{3.6}
\end{align*}
$$

By Theorem 1.3 and (3.3), for sufficiently large $n \geq 1$, we have $\left(\eta_{n}, \hat{\zeta}_{n}\right) \in$ $\hat{\Theta}\left(u_{0}, h\right)$ and

$$
\begin{equation*}
\left\|\mathcal{R}\left(u_{0}, \zeta_{n}, h+\eta_{n}\right)-\mathcal{R}\left(u_{0}, \hat{\zeta}_{n}, h+\eta_{n}\right)\right\|_{\mathcal{X}_{T, k}} \rightarrow 0 \quad \text { as } n \rightarrow \infty \tag{3.7}
\end{equation*}
$$

Note that

$$
\begin{align*}
\mathcal{R}_{t}\left(u_{0}, \hat{\zeta}_{n}, h+\eta_{n}\right) & =\mathcal{R}_{t}\left(u_{0}, 0, h+\hat{\eta}_{n}\right)-\hat{\zeta}_{n}(t) \quad \text { for } t \in J_{T}  \tag{3.8}\\
\mathcal{R}_{T}\left(u_{0}, \hat{\zeta}_{n}, h+\eta_{n}\right) & =\mathcal{R}_{T}\left(u_{0}, 0, h+\hat{\eta}_{n}\right) \tag{3.9}
\end{align*}
$$

where $\hat{\eta}_{n}:=\eta+\partial_{t} \hat{\zeta}_{n}$. From (3.2), (3.7), and (3.9) it follows that

$$
\left\|\mathcal{R}_{T}\left(u_{0}, 0, h+\eta_{1}\right)-\mathcal{R}_{T}\left(u_{0}, 0, h+\hat{\eta}_{n}\right)\right\|_{k} \rightarrow 0 \quad \text { as } n \rightarrow \infty
$$

Using (3.2), (3.6)-(3.8), we obtain

$$
\begin{aligned}
\| \mathcal{R}\left(u_{0}, 0, h+\eta_{1}\right)-\mathcal{R}\left(u_{0}, 0,\right. & \left.h+\hat{\eta}_{n}\right) \|_{T, k} \\
\leq & T\left\|\mathcal{R}\left(u_{0}, 0, h+\eta_{1}\right)-\mathcal{R}\left(u_{0}, \zeta_{n}, h+\eta_{n}\right)\right\|_{L^{\infty}\left(J_{T}, H^{k}\right)} \\
& +T\left\|\mathcal{R}\left(u_{0}, \zeta_{n}, h+\eta_{n}\right)-\mathcal{R}\left(u_{0}, \hat{\zeta}_{n}, h+\eta_{n}\right)\right\|_{\mathcal{X}_{T, k}} \\
& +\left\|\mathcal{R}\left(u_{0}, \hat{\zeta}_{n}, h+\eta_{n}\right)-\mathcal{R}\left(u_{0}, 0, h+\hat{\eta}_{n}\right)\right\|_{T, k} \\
\rightarrow & 0 \quad \text { as } n \rightarrow \infty
\end{aligned}
$$

Combining this with the embedding $H^{3} \subset C^{1,1 / 2}$, (3.3), (3.8), and applying Lemma 1.1 with $\lambda=1 / 2$, we get that

$$
\left\|\phi^{\mathcal{R}\left(u_{0}, 0, h+\eta_{1}\right)}-\phi^{\mathcal{R}\left(u_{0}, 0, h+\hat{\eta}_{n}\right)}\right\|_{L^{\infty}\left(J_{T}, C^{1}\right)} \rightarrow 0 \quad \text { as } n \rightarrow \infty .
$$

This completes the proof of Theorem 2.5.
Proof of Proposition 3.1. Step 1. Without loss of generality, we can assume that $\eta_{1} \in \Theta\left(u_{0}, h\right) \cap E_{1}$ is constant. Indeed, the general case is then obtained by approximation in $L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right)$ by piecewise constant controls with finite number of intervals of constancy and successive applications of the result on the intervals of constancy.

By the definition of $\mathcal{F}(E)$, for any $\eta_{1} \in E_{1}$, there are vectors $\xi^{1}, \ldots, \xi^{p}, \eta \in E$ such that

$$
\eta_{1}=\eta-\sum_{i=1}^{p} B\left(\xi^{i}\right)
$$

Choosing $m=2 p$ and

$$
\zeta^{i}:=-\zeta^{i+p}:=\frac{1}{\sqrt{2}} \xi^{i}, \quad i=1, \ldots, p
$$

it is easy to see that

$$
\begin{equation*}
B(u)-\eta_{1}=\frac{1}{m} \sum_{j=1}^{m}\left(B\left(u+\zeta^{j}\right)+L \zeta^{j}\right)-\eta \quad \text { for any } u \in H_{\sigma}^{1} \tag{3.10}
\end{equation*}
$$

Then $u_{1}:=\mathcal{R}\left(u_{0}, 0, h+\eta_{1}\right) \in \mathcal{X}_{T, k}$ satisfies the following equation

$$
\begin{equation*}
\dot{u}_{1}+L u_{1}+\frac{1}{m} \sum_{j=1}^{m}\left(B\left(u+\zeta^{j}\right)+L \zeta^{j}\right)=h(t)+\eta . \tag{3.11}
\end{equation*}
$$

Let us define $\zeta_{n}(t)=\zeta\left(\frac{n t}{T}\right)$, where $\zeta(t)$ is a 1-periodic function such that

$$
\zeta(s)=\zeta^{j} \text { for } s \in[(j-1) / m, j / m), j=1, \ldots, m
$$

Equation (3.11) is equivalent to

$$
\dot{u}_{1}+L\left(u_{1}+\zeta_{n}\right)+B\left(u_{1}+\zeta_{n}\right)=h(t)+\eta+f_{n}(t),
$$

where

$$
\begin{equation*}
f_{n}(t):=L \zeta_{n}+B\left(u_{1}+\zeta_{n}\right)-\frac{1}{m} \sum_{j=1}^{m}\left(B\left(u_{1}+\zeta^{j}\right)+L \zeta^{j}\right) \tag{3.12}
\end{equation*}
$$

For any $f \in L^{2}\left(J_{T}, H\right)$, let us set

$$
K f(t)=\int_{0}^{t} e^{-(t-s) L} f(s) \mathrm{d} s
$$

It is easy to check that
$K$ is continuous from $L^{2}\left(J_{T}, H_{\sigma}^{p-1}\right)$ to $\mathcal{X}_{T, p}$ for any $p \geq 1$,
and $v_{n}=u_{1}-K f_{n}$ is a solution of the problem

$$
\begin{align*}
\dot{v}_{n}+L\left(v_{n}+\zeta_{n}\right)+B\left(v_{n}+\zeta_{n}+K f_{n}\right) & =h(t)+\eta,  \tag{3.14}\\
v_{n} & =u_{0} .
\end{align*}
$$

Step 2. Let us show that

$$
\begin{equation*}
\left\|K f_{n}\right\|_{L^{\infty}\left(J_{T}, H^{k}\right)} \rightarrow 0 \quad \text { as } n \rightarrow \infty \tag{3.15}
\end{equation*}
$$

Indeed, the definition of $\zeta_{n}$ gives that

$$
\begin{equation*}
\sup _{n \geq 1}\left\|\zeta_{n}\right\|_{L^{\infty}\left(J_{T}, H^{k+1}\right)}<\infty \tag{3.16}
\end{equation*}
$$

Combining this with (3.12), (1.18), and the fact that $u_{1} \in \mathcal{X}_{T, k}$, we get

$$
\begin{equation*}
\sup _{n \geq 1}\left\|f_{n}\right\|_{L^{\infty}\left(J_{T}, H^{k-1}\right)}<\infty \tag{3.17}
\end{equation*}
$$

This implies that

$$
\begin{align*}
\sup _{n \geq 1}\left\|K f_{n}\right\|_{L^{\infty}\left(J_{T}, H_{\sigma}^{k+1 / 2}\right)} & \leq C \sup _{n \geq 1, t \in[0, T]} \int_{0}^{t}\left\|L^{3 / 4} e^{-(t-s) L}\right\|_{\mathcal{L}(H)}\left\|f_{n}(s)\right\|_{k-1} \mathrm{~d} s \\
& \leq C_{1} \sup _{n \geq 1, t \in[0, T]} \int_{0}^{t}(t-s)^{-3 / 4}\left\|f_{n}(s)\right\|_{k-1} \mathrm{~d} s \\
& \leq C_{2} \sup _{n \geq 1}\left\|f_{n}\right\|_{L^{\infty}\left(J_{T}, H^{k-1}\right)}<\infty \tag{3.18}
\end{align*}
$$

where we used the inequality ${ }^{4}$

$$
\left\|L^{r} e^{-t L}\right\|_{\mathcal{L}(H)} \leq C_{r} t^{-r} \quad \text { for any } r \geq 0, t>0
$$

In Step 4 of the proof of Proposition 3.2 in [Shi06], it is established that

$$
\left\|K f_{n}\right\|_{L^{\infty}\left(J_{T}, H^{1}\right)} \rightarrow 0
$$

Using this with (3.18) and an interpolation inequality, we get (3.15). Combining (3.13) with (3.17), we obtain also that

$$
\begin{equation*}
\sup _{n \geq 1}\left\|K f_{n}\right\|_{\mathcal{X}_{T, k}}<\infty \tag{3.19}
\end{equation*}
$$

Step 3. Equation (3.14) can be rewritten as

$$
\begin{equation*}
\dot{v}_{n}+L\left(v_{n}+\zeta_{n}\right)+B\left(v_{n}+\zeta_{n}\right)=h(t)+\eta+g_{n}(t) \tag{3.20}
\end{equation*}
$$

where

$$
g_{n}(t):=-\left(B\left(v_{n}+\zeta_{n}, K f_{n}\right)+B\left(K f_{n}, v_{n}+\zeta_{n}\right)+B\left(K f_{n}\right)\right)
$$

From (3.15), (3.3), and (1.18) it is easy to deduce that $\left\|g_{n}\right\|_{L^{2}\left(J_{T}, H^{k-1}\right)} \rightarrow 0$ as $n \rightarrow \infty$. From (3.19) it follows that

$$
\sup _{n \geq 1}\left\|v_{n}\right\|_{\mathcal{X}_{T, k}}<\infty
$$

Therefore, by Theorem 1.3 and (3.16), we have $\left(\eta, \zeta_{n}\right) \in \hat{\Theta}\left(u_{0}, h\right)$ for sufficiently large $n \geq 1$ and

$$
\left\|\mathcal{R}\left(u_{0}, \zeta_{n}, \eta\right)-v_{n}\right\|_{\mathcal{X}_{T, k}} \rightarrow 0 \text { as } n \rightarrow \infty
$$

On the other hand, by (3.15),

$$
\left\|v_{n}-u_{1}\right\|_{L^{\infty}\left(J_{T}, H^{k}\right)} \rightarrow 0 \text { as } n \rightarrow \infty
$$

whence

$$
\begin{aligned}
& \left\|\mathcal{R}\left(u_{0}, \zeta_{n}, \eta\right)-u_{1}\right\|_{L^{\infty}\left(J_{T}, H^{k}\right)} \rightarrow 0 \text { as } n \rightarrow \infty \\
& \sup _{n \geq 1}\left\|\mathcal{R}\left(u_{0}, \zeta_{n}, \eta\right)\right\|_{\mathcal{X}_{T, k}}<+\infty
\end{aligned}
$$

Step 4. Let us show that

$$
\begin{equation*}
\left\|\zeta_{n}\right\|_{T, k} \rightarrow 0 \quad \text { as } n \rightarrow \infty \tag{3.21}
\end{equation*}
$$

We set $\mathcal{L} \zeta_{n}(t):=\int_{0}^{t} \zeta_{n}(s) \mathrm{d} s$. It suffices to check that

[^4](i) the sequence $\mathcal{L} \zeta_{n}$ is relatively compact in $C\left(J_{T}, H_{\sigma}^{k}\right)$.
(ii) for any $t \in J_{T}, \mathcal{L} \zeta_{n}(t) \rightarrow 0$ in $H_{\sigma}^{k}$ as $n \rightarrow \infty$.

To prove the first assertion, we use the Arzelà-Ascoli theorem. The functions $\zeta_{n}$ are piecewise constant and the set $\zeta_{n}(t), t \in J_{T}$ is contained in a finite subset of $H_{\sigma}^{k+1}$ not depending on $n$. This implies that there is a compact set $F \subset H_{\sigma}^{k+1}$ such that

$$
\mathcal{L} \zeta_{n}(t) \in F \text { for all } t \in J_{T}, n \geq 1
$$

From (3.3) it follows that the sequence $\mathcal{L} \zeta_{n}$ is uniformly equicontinuous on $J_{T}$. Thus, by the Arzelà-Ascoli theorem, $\mathcal{L} \zeta_{n}$ is relatively compact in $C\left(J_{T}, H_{\sigma}^{k}\right)$.

Let us prove (ii). Let $t=t_{l}+\tau$, where $t_{l}=\frac{l T}{n}, l \in \mathbb{N}$ and $\tau \in\left[0, \frac{T}{n}\right.$ ). In view of the construction of $\zeta_{n}$, we have that $\mathcal{L} \zeta_{n}(l T / n)=0$. Combining this with (3.3), we get

$$
\mathcal{L} \zeta_{n}(t)=\int_{\frac{I T}{n}}^{t} \zeta_{n}(s) \mathrm{d} s \rightarrow 0
$$

which completes the proof of (3.21).
Finally, taking an arbitrary sequence $\hat{\zeta}_{n} \in C^{\infty}\left(J_{T}, E\right)$ such that

$$
\left\|\zeta_{n}-\hat{\zeta}_{n}\right\|_{L^{\infty}\left(J_{T}, E\right)} \rightarrow 0 \text { as } n \rightarrow \infty
$$

and using Theorem 1.3, we see that the conclusions of Proposition 3.1 hold for the sequence $\left(\eta, \hat{\zeta}_{n}\right) \in C^{\infty}\left(J_{T}, E \times E\right)$.

## 4 Examples of saturating spaces

In this section, we provide three types of examples of saturating spaces which ensure the controllability of the 3D NS system in the sense of Definition 2.1.

### 4.0.1 Saturating spaces associated with the generators of $\mathbb{Z}^{3}$

Let us first introduce some notation. Denote by $\mathbb{Z}_{*}^{3}$ the set of non-zero integer vectors $\ell=\left(l_{1}, l_{2}, l_{3}\right) \in \mathbb{Z}^{3}$. For any $\ell \in \mathbb{Z}_{*}^{3}$, let us define the functions

$$
\begin{equation*}
c_{\ell}(x)=l(\ell) \cos \langle\ell, x\rangle, s_{\ell}(x)=l(\ell) \sin \langle\ell, x\rangle \tag{4.1}
\end{equation*}
$$

where $\{l(\ell), l(-\ell)\}$ is an arbitrary orthonormal basis in

$$
\ell^{\perp}:=\left\{x \in \mathbb{R}^{3}:\langle x, \ell\rangle=0\right\}
$$

Then $c_{\ell}$ and $s_{\ell}$ are eigenfunctions of $L$ and the family $\left\{c_{\ell}, s_{\ell}\right\}_{\ell \in \mathbb{Z}_{*}^{3}}$ is an orthonormal basis in $H$. Let $c_{0}=s_{0}=0$. For any subset $\mathcal{K} \subset \mathbb{Z}^{3}$, we denote

$$
\begin{equation*}
E(\mathcal{K}):=\operatorname{span}\left\{c_{\ell}, c_{-\ell}, s_{\ell}, s_{-\ell}: \ell \in \mathcal{K}\right\} \tag{4.2}
\end{equation*}
$$

When $\mathcal{K}$ is finite, the spaces $E_{j}(\mathcal{K})$ and $E_{\infty}(\mathcal{K})$ are defined by $(2.5)$ with $E=$ $E(\mathcal{K})$. We denote by $\mathbb{Z}_{\mathcal{K}}^{3}$ the set of all vectors $a \in \mathbb{Z}^{3}$ which can be represented as finite linear combination of elements of $\mathcal{K}$ with integer coefficients. We shall say that $\mathcal{K} \subset \mathbb{Z}^{3}$ is a generator if $\mathbb{Z}_{\mathcal{K}}^{3}=\mathbb{Z}^{3}$. The following theorem provides a characterisation of saturating spaces of the form (4.2).

Theorem 4.1. For any finite set $\mathcal{K} \subset \mathbb{Z}^{3}$, we have the equality

$$
\begin{equation*}
E\left(\mathbb{Z}_{\mathcal{K}}^{3}\right)=E_{\infty}(\mathcal{K}) . \tag{4.3}
\end{equation*}
$$

Moreover, $E(\mathcal{K})$ is saturating in $H$ if and only if $\mathcal{K}$ is a generator of $\mathbb{Z}^{3}$. If $E(\mathcal{K})$ is saturating in $H$, then it is saturating in $H_{\sigma}^{k}$ for any $k \geq 0$.

In [Rom04] a similar result is conjectured in the case of finite-dimensional approximations of the 3D NS system and a proof is given for the saturating property of $E(\mathcal{K})$ when $\mathcal{K}=\{(1,0,0),(0,1,0),(0,0,1)\}^{5}$. A 2 D version of Theorem 4.1 is established in [EM01] and [HM06]. In that case, the set $\mathcal{K}$ is a generator of $\mathbb{Z}^{2}$ containing at least two vectors with different Euclidian norms (the reader is referred to the original papers for the exact statement). The proof in the 3D case, as well as the statement of the result, differ essentially from the 2 D case.

In view of Theorem 4.1, the following simple criterion is useful for constructing saturating spaces (see Section 3.7 in [Jac85]).

Theorem 4.2. $A$ set $\mathcal{K} \subset \mathbb{Z}^{3}$ is a generator if and only if the greatest common divisor of the set $\{\operatorname{det}(a, b, c): a, b, c \in \mathcal{K}\}$ is 1 , where $\operatorname{det}(a, b, c)$ is the determinant of the matrix with rows $a, b$ and $c$.

The proof of Theorem 4.1 is deduced from the following auxiliary result.
Proposition 4.3. Assume that $\mathcal{W} \subset \mathbb{Z}^{3}$ is a finite set containing a linearly independent family $\{p, q, r\} \subset \mathbb{Z}^{3}$. Then for any non-parallel vectors $m, n \in \mathcal{W}$ we have $\mathcal{A}_{m \pm n}, \mathcal{B}_{m \pm n} \subset E_{3}(\mathcal{W})$, where

$$
\mathcal{A}_{\ell}:=\operatorname{span}\left\{c_{\ell}, c_{-\ell}\right\}, \quad \mathcal{B}_{\ell}:=\operatorname{span}\left\{s_{\ell}, s_{-\ell}\right\}, \quad \ell \in \mathbb{Z}_{*}^{3}
$$

Proof of Proposition 4.3. We shall confine ourselves to the proof of the inclusion

$$
\begin{equation*}
\mathcal{A}_{m+n} \subset E_{3}(\mathcal{W}) \tag{4.4}
\end{equation*}
$$

The other conclusions in the proposition are checked in the same way.
Step 1. We shall write $m \nVdash n$ when the vectors $m, n \in \mathbb{R}^{3}$ are non-parallel. For any $m, n \in \mathcal{W}$ such that $m \nVdash n$, let us denote by $\delta:=\delta(m, n)$ one of the two unit vectors belonging to $m^{\perp} \cap n^{\perp}$. In this step we show that

$$
\begin{equation*}
\delta \cos \langle m \pm n, x\rangle, \delta \sin \langle m \pm n, x\rangle \in E_{1}(\mathcal{W}) \tag{4.5}
\end{equation*}
$$

[^5]Indeed, for any $a \in \mathbb{R}_{*}^{3}$, let us denote by $P_{a}$ the orthogonal projection in $\mathbb{R}^{3}$ onto $a^{\perp}$. Then we have

$$
\Pi(a \cos \langle l, x\rangle)=\left(P_{l} a\right) \cos \langle l, x\rangle, \quad \Pi(a \sin \langle l, x\rangle)=\left(P_{l} a\right) \sin \langle l, x\rangle
$$

for any $l \in \mathbb{Z}_{*}^{3}$. Combining this with some trigonometric identities and the definition of $B$, one gets that

$$
\begin{align*}
2 B(a \cos \langle m, x\rangle+b \sin \langle n, x\rangle)= & \cos \langle m-n, x\rangle P_{m-n}(\langle a, n\rangle b-\langle b, m\rangle a) \\
& +\cos \langle m+n, x\rangle P_{m+n}(\langle a, n\rangle b+\langle b, m\rangle a) \tag{4.6}
\end{align*}
$$

for any $a \in m^{\perp}$ and $b \in n^{\perp}$ (see Step 1 of the proof of Proposition 2.8 in [Shi06]). This implies that

$$
\begin{align*}
2 B(b \cos \langle n, x\rangle+a \sin \langle m, x\rangle)= & -\cos \langle m-n, x\rangle P_{m-n}(\langle a, n\rangle b-\langle b, m\rangle a) \\
& +\cos \langle m+n, x\rangle P_{m+n}(\langle a, n\rangle b+\langle b, m\rangle a) \tag{4.7}
\end{align*}
$$

Taking the sum of (4.6) and (4.7), we obtain that

$$
\begin{align*}
\cos \langle m+n, x\rangle P_{m+n}(\langle a, n\rangle b+\langle b, m\rangle a)= & B(a \cos \langle m, x\rangle+b \sin \langle n, x\rangle) \\
& +B(b \cos \langle n, x\rangle+a \sin \langle m, x\rangle) \tag{4.8}
\end{align*}
$$

Let us fix any $\lambda \in \mathbb{R}$ and choose in this equality $a=\delta$ and $\langle b, m\rangle=\lambda$. This choice is possible since $m \nVdash n$. Then we have

$$
\lambda \delta \cos \langle m+n, x\rangle=B(\delta \cos \langle m, x\rangle+b \sin \langle n, x\rangle)+B(b \cos \langle n, x\rangle+\delta \sin \langle m, x\rangle)
$$

Since $\lambda \in \mathbb{R}$ is arbitrary, from the definition of $E_{1}(\mathcal{W})$ we get that $\delta \cos \langle m+$ $n, x\rangle \in E_{1}(\mathcal{W})$. To prove that $\delta \cos \langle m-n, x\rangle \in E_{1}(\mathcal{W})$, it suffices to replace $b$ by $-b$ in (4.7), take the sum of the resulting equality with (4.6):

$$
\begin{align*}
\cos \langle m-n, x\rangle P_{m-n}(\langle a, n\rangle b-\langle b, m\rangle a)= & B(a \cos \langle m, x\rangle+b \sin \langle n, x\rangle) \\
& +B(-b \cos \langle n, x\rangle+a \sin \langle m, x\rangle) \tag{4.9}
\end{align*}
$$

and choose $a=\delta$ and $\langle b, m\rangle=-\lambda$

$$
\lambda \delta \cos \langle m-n, x\rangle=B(\delta \cos \langle m, x\rangle+b \sin \langle n, x\rangle)+B(-b \cos \langle n, x\rangle+\delta \sin \langle n, x\rangle) .
$$

The fact that $\delta \sin \langle m \pm n, x\rangle \in E_{1}(\mathcal{W})$ is proved in a similar way using the following identities

$$
\begin{aligned}
2 B(a \cos \langle m, x\rangle+b \cos \langle n, x\rangle)= & \sin \langle m-n, x\rangle P_{m-n}(\langle a, n\rangle b-\langle b, m\rangle a) \\
& -\sin \langle m+n, x\rangle P_{m+n}(\langle a, n\rangle b+\langle b, m\rangle a) \\
2 B(a \sin \langle m, x\rangle+b \sin \langle n, x\rangle)= & \sin \langle m-n, x\rangle P_{m-n}(\langle a, n\rangle b-\langle b, m\rangle a) \\
& +\sin \langle m+n, x\rangle P_{m+n}(\langle a, n\rangle b+\langle b, m\rangle a) .
\end{aligned}
$$

Step 2. To prove (4.4), let us take any vector $r \in \mathcal{W}$ such that $\mathcal{E}:=\{m, n, r\}$ is a linearly independent family ${ }^{6}$ in $\mathbb{R}^{3}$. This choice is possible, by the conditions of the proposition. For any $\alpha, \beta, \gamma \in \mathbb{R}$, we shall write $(\alpha, \beta, \gamma)_{\mathcal{E}}$ instead of $\alpha m+\beta n+\gamma r$. Then we have also that the family $\left\{(1,1,-1)_{\mathcal{E}},(1,-1,1)_{\mathcal{E}}\right.$, $\left.(-1,1,1)_{\mathcal{E}}\right\}$ is independent, hence

$$
\{0\}=(1,1,-1)_{\mathcal{\mathcal { E }}}^{\perp} \cap(1,-1,1)_{\mathcal{\mathcal { E }}}^{\perp} \cap(-1,1,1)_{\mathcal{E}}^{\perp} .
$$

We are going to prove (4.4) under the assumption

$$
\begin{equation*}
(1,1,1)_{\mathcal{E}} \notin(1,1,-1)_{\mathcal{E}}^{\perp} \tag{4.10}
\end{equation*}
$$

The other two cases $(1,1,1)_{\mathcal{E}} \notin(1,-1,1)_{\mathcal{E}}^{\perp}$ and $(1,1,1)_{\mathcal{E}} \notin(-1,1,1)_{\mathcal{E}}^{\perp}$ are similar. As

$$
(1,1,0)_{\mathcal{E}}=(1,0,0)_{\mathcal{E}}+(0,1,0)_{\mathcal{E}}=m+n
$$

by (4.5), we have

$$
\begin{equation*}
\delta(m, n) \cos \left\langle(1,1,0)_{\mathcal{E}}, x\right\rangle \in E_{1}(\mathcal{W}) \tag{4.11}
\end{equation*}
$$

Writing

$$
(1,1,1)_{\mathcal{E}}=(0,0,1)_{\mathcal{E}}+(1,1,0)_{\mathcal{E}}
$$

and applying (4.8) and (4.11), we obtain for any $b \in(0,0,1)_{\mathcal{E}}$ that

$$
\begin{align*}
& \cos \left\langle(1,1,1)_{\mathcal{E}}, x\right\rangle P_{(1,1,1)_{\mathcal{E}}}\left(\left\langle\delta(m, n),(0,0,1)_{\mathcal{E}}\right\rangle b+\left\langle b,(1,1,0)_{\mathcal{E}}\right\rangle \delta(m, n)\right) \\
& =B\left(\delta(m, n) \cos \left\langle(1,1,0)_{\mathcal{E}}, x\right\rangle+b \sin \left\langle(0,0,1)_{\mathcal{E}}, x\right\rangle\right) \\
& \quad+B\left(b \cos \left\langle(0,0,1)_{\mathcal{E}}, x\right\rangle+\delta(m, n) \sin \left\langle(1,1,0)_{\mathcal{E}}, x\right\rangle\right) \in E_{2}(\mathcal{W}) . \tag{4.12}
\end{align*}
$$

Let us define the set

$$
\mathcal{G}:=\left\{\left\langle\delta(m, n),(0,0,1)_{\mathcal{E}}\right\rangle b+\left\langle b,(1,1,0)_{\mathcal{E}}\right\rangle \delta(m, n): b \in(0,0,1)_{\mathcal{E}}^{\perp}\right\} .
$$

Since $m, n, r$ are linearly independent, we have $\left\langle\delta(m, n),(0,0,1)_{\mathcal{E}}\right\rangle \neq 0$. Thus $\mathcal{G}$ is a two-dimensional subspace of $\mathbb{R}^{3}$ contained in $(1,1,-1) \frac{\perp}{\mathcal{E}}$. This shows that $\mathcal{G}=(1,1,-1) \frac{\perp}{\mathcal{E}}$. Assumption (4.10) implies that the orthogonal projection $P_{(1,1,1)_{\mathcal{E}} \mathcal{G}}$ coincides with $(1,1,1)_{\mathcal{E}}^{\perp}$, so (4.12) proves that

$$
\begin{equation*}
\mathcal{A}_{(1,1,1)_{\mathcal{E}}} \subset E_{2}(\mathcal{W}) \tag{4.13}
\end{equation*}
$$

Similarly, one can show that $\mathcal{B}_{(1,1,1)_{\mathcal{E}}} \subset E_{2}(\mathcal{W})$. Finally, writing

$$
(1,1,0)_{\mathcal{E}}=(1,1,1)_{\mathcal{E}}-(0,0,1)_{\mathcal{E}}
$$

and applying (4.5) and (4.11) to the set $\mathcal{W}_{1}:=\mathcal{W} \cup\left\{(1,1,1)_{\mathcal{E}},(0,0,1)_{\mathcal{E}}\right\}$, we see that

$$
\begin{aligned}
\delta\left((1,1,1)_{\mathcal{E}},(0,0,1)_{\mathcal{E}}\right) \cos \left\langle(1,1,0)_{\mathcal{E}}, x\right\rangle & \in E_{1}\left(\mathcal{W}_{1}\right)=\mathcal{F}\left(E\left(\mathcal{W}_{1}\right)\right) \\
& \subset \mathcal{F}\left(E_{2}(\mathcal{W})\right)=E_{3}(\mathcal{W}) .
\end{aligned}
$$

Combining this with (4.11) and the fact that $\delta\left((1,1,1)_{\mathcal{E}},(0,0,1)_{\mathcal{E}}\right) \nVdash \delta(m, n)$, we get (4.4). This completes the proof of the proposition.

[^6]Proof of Theorem 4.1. Step 1. Let us show that

$$
\begin{equation*}
E\left(\mathbb{Z}_{\mathcal{K}}^{3}\right) \subset E_{\infty}(\mathcal{K}) \tag{4.14}
\end{equation*}
$$

To this end, we introduce the sets

$$
\mathcal{K}_{0}:=\mathcal{K}, \quad \mathcal{K}_{j}=\mathcal{K}_{j-1} \cup\left\{m \pm n: m, n \in \mathcal{K}_{j-1}, m \nVdash n\right\}, \quad j \geq 1
$$

From Proposition 4.3 it follows that

$$
\begin{align*}
E\left(\mathcal{K}_{j}\right) \subset E_{3}\left(\mathcal{K}_{j-1}\right) & =\mathcal{F}^{3}\left(E\left(\mathcal{K}_{j-1}\right)\right) \subset \mathcal{F}^{6}\left(E\left(\mathcal{K}_{j-2}\right)\right) \subset \ldots \subset \mathcal{F}^{3 j}(E(\mathcal{K})) \\
& =E_{3 j}(\mathcal{K}) \tag{4.15}
\end{align*}
$$

On the other hand, since $\mathcal{K}$ is a generator of $\mathbb{Z}_{\mathcal{K}}^{3}$, one easily checks that $\cup_{j=1}^{\infty} \mathcal{K}_{j}=$ $\mathbb{Z}_{\mathcal{K}}^{3}$. Combining this with (4.15), we get (4.14).

Step 2. Now let us prove that

$$
\begin{equation*}
E_{\infty}(\mathcal{K}) \subset E\left(\mathbb{Z}_{\mathcal{K}}^{3}\right) \tag{4.16}
\end{equation*}
$$

For any $\eta_{1} \in E_{1}\left(\mathcal{K}_{j-1}\right)$ and $j \geq 1$, there are vectors $\eta, \zeta^{1}, \ldots, \zeta^{p} \in E\left(\mathcal{K}_{j-1}\right)$ satisfying the relation

$$
\eta_{1}=\eta-\sum_{i=1}^{p} B\left(\zeta^{i}\right)
$$

Here we use the following simple lemma.
Lemma 4.4. For any $j \geq 1$, we have

$$
\left\{B(\zeta): \zeta \in E\left(\mathcal{K}_{j-1}\right)\right\} \subset E\left(\mathcal{K}_{j}\right)
$$

This lemma implies that

$$
E_{1}\left(\mathcal{K}_{j-1}\right) \subset E\left(\mathcal{K}_{j}\right)
$$

Iterating this, we get

$$
E_{j}(\mathcal{K}) \subset E\left(\mathcal{K}_{j}\right)
$$

hence

$$
E_{\infty}(\mathcal{K})=\bigcup_{j=1}^{\infty} E_{j}(\mathcal{K}) \subset \bigcup_{j=1}^{\infty} E\left(\mathcal{K}_{j}\right) \subset E\left(\bigcup_{j=1}^{\infty} \mathcal{K}_{j}\right)=E\left(\mathbb{Z}_{\mathcal{K}}^{3}\right)
$$

This proves (4.16) and (4.3).
Step 3. If $\mathcal{K}$ is a generator of $\mathbb{Z}^{3}$, then (4.3) implies that $E(\mathcal{K})$ is saturating in $H_{\sigma}^{k}$ for any $k \geq 0$.

Now let us assume that $\mathcal{K}$ is not a generator of $\mathbb{Z}^{3}$, i.e., there is $\ell \in \mathbb{Z}^{3}$ such that $\ell \notin \mathbb{Z}_{\mathcal{K}}^{3}$. Then it follows from (4.3) that $c_{\ell}$ is orthogonal to $E_{\infty}(\mathcal{K})$ in $H$. This shows that $E(\mathcal{K})$ is not saturating in $H$ and completes the proof of the theorem.

Proof of Lemma 4.4. For any $\zeta \in E\left(\mathcal{K}_{j-1}\right)$, we have

$$
\zeta=\sum_{\ell \in \pm \mathcal{K}_{j-1}}\left(a_{\ell} c_{\ell}+b_{\ell} s_{\ell}\right)
$$

for some $a_{\ell}, b_{\ell} \in \mathbb{R}$. It follows that

$$
\begin{aligned}
B(\zeta)= & \sum_{m, n \in \pm \mathcal{K}_{j-1}}\left(a_{m} a_{n} B\left(c_{m}, c_{n}\right)+b_{m} b_{n} B\left(s_{m}, s_{n}\right)\right. \\
& \left.+a_{m} b_{n} B\left(c_{m}, s_{n}\right)+b_{m} a_{n} B\left(s_{m}, c_{n}\right)\right)
\end{aligned}
$$

Using some trigonometric identities, it is easy to verify that

$$
B\left(c_{m}, c_{n}\right) \in \operatorname{span}\left\{s_{m+n}, s_{m-n}\right\} \subset E\left(\mathcal{K}_{j}\right)
$$

In a similar way, one gets $B\left(s_{m}, s_{n}\right), B\left(c_{m}, s_{n}\right), B\left(s_{m}, c_{n}\right) \in E\left(\mathcal{K}_{j}\right)$.
For any finite set $\mathcal{K} \subset \mathbb{Z}^{3}$ and $k \geq 3$, let us define the space

$$
H_{\sigma, \mathcal{K}}^{k}:={\overline{E_{\infty}(\mathcal{K})}}^{H^{k}}
$$

From the structure of the nonlinearity it follows that $H_{\sigma, \mathcal{K}}^{k}$ is invariant for (2.1) when $h, \eta \in L^{2}\left(J_{T}, H_{\sigma, \mathcal{K}}^{k-1}\right)$. Moreover, $H_{\sigma, \mathcal{K}}^{k}=H_{\sigma}^{k}$ if and only if $\mathcal{K}$ is a generator of $\mathbb{Z}^{3}$. As a corollary we get the following characterisation of the controllability in $H_{\sigma}^{k}$.
Theorem 4.5. Let $\mathcal{K} \subset \mathbb{Z}^{3}$ be a finite set and $h \in L^{2}\left(J_{T}, H_{\sigma, \mathcal{K}}^{k-1}\right)$. Then equation (2.1) is approximately controllable in the space $H_{\sigma}^{k}$ at time $T$ by controls $\eta \in C^{\infty}\left(J_{T}, E(\mathcal{K})\right)$ if and only if $\mathcal{K}$ is a generator of $\mathbb{Z}^{3}$.

It is also interesting to study the controllability properties of the NS system when $E(\mathcal{K})$ given by (4.2) is not saturating (i.e., $\mathcal{K}$ is not a generator of $\mathbb{Z}^{3}$ ). Let us note that the space $E(\mathcal{K})$ is saturating in $H_{\sigma, \mathcal{K}}^{k}$ for any $\mathcal{K} \subset \mathbb{Z}^{3}$ and $k \geq 0$ (in the sense that $E_{\infty}(\mathcal{K})$ is dense in $\left.H_{\sigma, \mathcal{K}}^{k}\right)$. We have the following refined version of Theorem 2.2.

Theorem 4.6. For any non-empty finite $\mathcal{K} \subset \mathbb{Z}^{3}$ and $h \in L^{2}\left(J_{T}, H_{\sigma, \mathcal{K}}^{k-1}\right)$, equation (2.1) is approximately controllable in the space $H_{\sigma, \mathcal{K}}^{k}$ at time $T$ by controls $\eta \in C^{\infty}\left(J_{T}, E(\mathcal{K})\right)$, i.e., for any $\varepsilon>0$ and any

$$
\varphi \in C\left(J_{T}, H_{\sigma, \mathcal{K}}^{k}\right) \cap L^{2}\left(J_{T}, H_{\sigma, \mathcal{K}}^{k+1}\right) \cap W^{1,2}\left(J_{T}, H_{\sigma, \mathcal{K}}^{k-1}\right)
$$

there is a control $\eta \in \Theta\left(h, u_{0}\right) \cap C^{\infty}\left(J_{T}, E(\mathcal{K})\right)$ such that
$\left\|\mathcal{R}_{T}\left(u_{0}, h+\eta\right)-\varphi(T)\right\|_{k}+\left\|\mathcal{R}\left(u_{0}, h+\eta\right)-\varphi\right\|_{T, k}+\left\|\phi^{\mathcal{R}\left(u_{0}, h+\eta\right)}-\phi^{\varphi}\right\|_{L^{\infty}\left(J_{T}, C^{1}\right)}<\varepsilon$, where $u_{0}=\varphi(0)$.

The proof of this result literally repeats the arguments of the proof of Theorem 2.2 , so we omit the details.

### 4.0.2 Controls with two vanishing components

In this section, we consider the NS system

$$
\begin{align*}
\partial_{t} u-\nu \Delta u+\langle u, \nabla\rangle u+\nabla p & =h(t, x)+(0,0,1) \eta(t, x), \quad \operatorname{div} u=0  \tag{4.17}\\
u(0) & =u_{0} \tag{4.18}
\end{align*}
$$

where $\eta$ is a control taking values in a finite-dimensional space of the form

$$
\mathcal{H}(\mathcal{K}):=\operatorname{span}\{\cos \langle m, x\rangle, \sin \langle m, x\rangle: m \in \mathcal{K}\}
$$

where $\mathcal{K}$ is a subset of $\mathbb{Z}^{3}$, and $h$ is a given smooth divergence-free function. Let us rewrite (4.17) in an equivalent form

$$
\begin{equation*}
\dot{u}-\nu \Delta u+B(u)=h(t, x)+\tilde{\eta}(t, x), \tag{4.19}
\end{equation*}
$$

where $\tilde{\eta}:=\Pi(e \eta)$ and $e:=(0,0,1)$. Then the control $\tilde{\eta}$ takes values in the space

$$
\begin{equation*}
\tilde{E}(\mathcal{K}):=\operatorname{span}\left\{\left(P_{m} e\right) \cos \langle m, x\rangle,\left(P_{m} e\right) \sin \langle m, x\rangle: m \in \mathcal{K}\right\} \tag{4.20}
\end{equation*}
$$

For an appropriate choice of $\mathcal{K}$, this space is saturating.
Proposition 4.7. Let

$$
\begin{equation*}
\mathcal{K}:=\{(1,0,0),(0,1,0),(1,0,1),(0,1,1)\} \tag{4.21}
\end{equation*}
$$

Then $\tilde{E}(\mathcal{K})$ is an 8-dimensional saturating space in $H_{\sigma}^{k}$ for any $k \geq 0$.
Combining this proposition with Theorem 2.2, we get immediately the following result.

Theorem 4.8. Let $h \in L^{2}\left(J_{T}, H_{\sigma}^{k-1}\right), k \geq 3$, and $T>0$. If $\mathcal{K}$ is defined by (4.21), then system (4.19) is approximately is controllable at time $T$ by controls $\tilde{\eta} \in C^{\infty}\left(J_{T}, \tilde{E}(\mathcal{K})\right)$.

Proof. Step 1. Let us first show that $\mathcal{A}_{(0,0,1)} \subset \mathcal{F}(\tilde{E}(\mathcal{K}))$. Using (4.9), we get for any $\lambda \in \mathbb{R}$

$$
\begin{aligned}
& \lambda(-1 / 2,0,0) \cos \langle(0,0,1), x\rangle \\
&= B\left(\lambda\left(P_{(1,0,0)} e\right) \cos \langle(1,0,0), x\rangle+\left(P_{(1,0,1)} e\right) \sin \langle(1,0,1), x\rangle\right) \\
& \quad+B\left(-\left(P_{(1,0,1)} e\right) \cos \langle(1,0,1), x\rangle+\lambda\left(P_{(1,0,0)} e\right) \sin \langle(1,0,0), x\rangle\right), \\
& \lambda(0,-1 / 2,0) \cos \langle(0,0,1), x\rangle \\
&= B\left(\lambda\left(P_{(0,1,0)} e\right) \cos \langle(0,1,0), x\rangle+\left(P_{(0,1,1)} e\right) \sin \langle(0,1,1), x\rangle\right) \\
& \quad+B\left(-\left(P_{(0,1,1)} e\right) \cos \langle(0,1,1), x\rangle+\lambda\left(P_{(0,1,0)} e\right) \sin \langle(0,1,0), x\rangle\right) .
\end{aligned}
$$

The definition of $\mathcal{F}$ implies that $\mathcal{A}_{(0,0,1)} \subset \mathcal{F}(\tilde{E}(\mathcal{K}))$. A similar computation gives that $\mathcal{B}_{(0,0,1)} \subset \mathcal{F}(\tilde{E}(\mathcal{K}))$.

Step 2. Again using (4.9), we obtain for any $b:=\left(b_{1}, b_{2}, 0\right) \in \mathbb{R}^{3}$

$$
\begin{aligned}
& \left(0, b_{2} / 2,-b_{1} / 2\right) \cos \langle(1,0,0), x\rangle=B\left(\left(P_{(1,0,1)} e\right) \cos \langle(1,0,1), x\rangle+b \sin \langle(0,0,1), x\rangle\right) \\
& \quad+B\left(-b \cos \langle(0,0,1), x\rangle+\left(P_{(1,0,1)} e\right) \sin \langle(1,0,1), x\rangle\right) \in \mathcal{F}^{2}(\tilde{E}(\mathcal{K}))
\end{aligned}
$$

This shows that $\mathcal{A}_{(1,0,0)} \subset \mathcal{F}^{2}(\tilde{E}(\mathcal{K}))$. Similarly one proves also

$$
\mathcal{B}_{(1,0,0)}, \mathcal{A}_{(0,1,0)}, \mathcal{B}_{(0,1,0)} \subset \mathcal{F}^{2}(\tilde{E}(\mathcal{K}))
$$

Thus the result follows from the fact that $\{(1,0,0),(0,1,0),(0,0,1)\}$ is a generator of $\mathbb{Z}^{3}$.

### 4.0.3 6-dimensional example

The following result, combined with Theorem 2.2, shows that that the 3D NS system can be approximately controlled with $\eta$ taking values in a 6 -dimensional space.

Proposition 4.9. Let us define the following 6-dimensional space

$$
\begin{align*}
\hat{E}:=\operatorname{span}\{ & a \cos \langle(1,0,1), x\rangle, a \sin \langle(1,0,1), x\rangle, \\
& e \cos \langle(0,1,1), x\rangle, e \sin \langle(0,1,1), x\rangle, \\
& b \cos \langle(0,0,1), x\rangle, b \sin \langle(0,0,1), x\rangle\}, \tag{4.22}
\end{align*}
$$

where $a:=(1,1,1), b:=(1,0,0), e:=(0,0,1)$. Then $\hat{E}$ is saturating in $H_{\sigma}^{k}$ for any $k \geq 0$.
Proof. Step 1. Let us first show that $\mathcal{A}_{(1,-1,0)} \subset \mathcal{F}^{2}(\hat{E})$. Using (4.9), we get for any $\lambda \in \mathbb{R}$

$$
\begin{aligned}
\lambda(0,-1,-1) & \cos \langle(1,0,0), x\rangle=B(\lambda a \cos \langle(1,0,1), x\rangle+b \sin \langle(0,0,1), x\rangle) \\
& +B(-b \cos \langle(0,0,1), x\rangle+\lambda a \sin \langle(1,0,1), x\rangle) \in \mathcal{F}(\hat{E}) \\
\lambda(1,0,0) & \cos \langle(0,1,0), x\rangle=B(\lambda e \cos \langle(0,1,1), x\rangle+b \sin \langle(0,0,1), x\rangle) \\
& +B(-b \cos \langle(0,0,1), x\rangle+\lambda e \sin \langle(0,1,1), x\rangle) \in \mathcal{F}(\hat{E})
\end{aligned}
$$

and $(0,-1,-1) \sin \langle(1,0,0), x\rangle,(1,0,0) \sin \langle(0,1,0), x\rangle \in \mathcal{F}(\hat{E})$, similarly. Writing

$$
(1,-1,0)=(1,0,0)-(0,1,0)=(1,0,1)-(0,1,1)
$$

and applying (4.9), we see that

$$
\begin{aligned}
& \lambda(0,0,1) \cos \langle(1,-1,0), x\rangle=B(\lambda(0,-1,-1) \cos \langle(1,0,0), x\rangle+b \sin \langle(0,1,0), x\rangle) \\
&+B(-b \cos \langle(0,1,0), x\rangle+\lambda(0,-1,-1) \sin \langle(1,0,0), x\rangle) \in \mathcal{F}^{2}(\hat{E}) \\
& \lambda(-1,-1,1) \cos \langle(1,-1,0), x\rangle=B(\lambda a \cos \langle(1,0,1), x\rangle+e \sin \langle(0,1,1), x\rangle) \\
&+B(-e \cos \langle(0,1,1), x\rangle+\lambda a \sin \langle(1,0,1), x\rangle) \in \mathcal{F}^{2}(\hat{E})
\end{aligned}
$$

This proves that $\mathcal{A}_{(1,-1,0)} \subset \mathcal{F}^{2}(\hat{E})$. A similar computation establishes that $\mathcal{B}_{(1,-1,0)} \subset \mathcal{F}^{2}(\hat{E})$.

Step 2. Let us show that $\mathcal{A}_{(1,0,0)}, \mathcal{B}_{(1,0,0)} \subset \mathcal{F}^{3}(\hat{E})$. Taking any vector $f:=\left(f_{1}, f_{1}, f_{2}\right) \in(1,-1,0)^{\perp}$, we apply (4.8)

$$
\begin{array}{r}
\left(0, f_{1}, f_{2}\right) \cos \langle(1,0,0), x\rangle=B(f \cos \langle(1,-1,0), x\rangle+b \sin \langle(0,1,0), x\rangle) \\
+B(b \cos \langle(0,1,0), x\rangle+f \sin \langle(1,-1,0), x\rangle) \in \mathcal{F}^{3}(\hat{E})
\end{array}
$$

This proves that $\mathcal{A}_{(1,0,0)} \subset \mathcal{F}^{3}(\hat{E})$, and $\mathcal{B}_{(1,0,0)} \subset \mathcal{F}^{3}(\hat{E})$ is similar.
Step 3. Let us show that $\mathcal{A}_{(0,0,1)}, \mathcal{B}_{(0,0,1)} \subset \mathcal{F}^{4}(\hat{E})$. Again we shall prove only the first inclusion. For any $g:=\left(0, g_{1}, g_{2}\right) \in \mathbb{R}^{3}$, we apply (4.9)

$$
\begin{aligned}
& \left(-g_{2}, g_{1}-g_{2}, 0\right) \cos \langle(0,0,1), x\rangle=B(a \cos \langle(1,0,1), x\rangle+g \sin \langle(1,0,0), x\rangle) \\
& +B(-g \cos \langle(1,0,0), x\rangle+a \sin \langle(1,0,1), x\rangle) \in \mathcal{F}^{3}(\hat{E}) .
\end{aligned}
$$

This proves that $\mathcal{A}_{(0,0,1)} \subset \mathcal{F}^{4}(\hat{E})$ and $\mathcal{B}_{(0,0,1)} \subset \mathcal{F}^{4}(\hat{E})$ is similar. By Theorem 4.2 , we have that the family $\{(1,0,0),(0,0,1),(1,-1,0)\}$ is a generator of $\mathbb{Z}^{3}$. Thus applying Theorem 4.1, we complete the proof.

It would be interesting to get a characterisation of finite-dimensional saturating spaces of the following general form

$$
E\left(\mathcal{K}_{c}, \mathcal{K}_{s}, a, b\right):=\operatorname{span}\left\{a_{m} \cos \langle m, x\rangle ; b_{n} \sin \langle n, x\rangle: m \in \mathcal{K}_{c}, n \in \mathcal{K}_{s}\right\}
$$

where $\mathcal{K}_{c}, \mathcal{K}_{s} \subset \mathbb{Z}^{3}, a:=\left\{a_{m}\right\}_{m \in \mathcal{K}_{c}} \subset \mathbb{R}_{*}^{3}$, and $b:=\left\{b_{n}\right\}_{n \in \mathcal{K}_{s}} \subset \mathbb{R}_{*}^{3}$. From the results of Subsection 4.0.1 it follows that both $\mathcal{K}_{c}$ and $\mathcal{K}_{s}$ are necessarily generators of $\mathbb{Z}^{3}$.

## References

[Ada75] R. A. Adams, Sobolev Spaces, Academic Press, New York-London, 1975.
[Arn78] V. I. Arnold, Ordinary Differential Equations, MIT Press, Cambridge, MA, 1998.
[AK98] V. I. Arnold and B. A. Khesin, Topological Methods in Hydrodynamics, Applied Mathematical Sciences, vol. 125, Springer-Verlag, New York, 1998.
[AS05] A. A. Agrachev and A. V. Sarychev, Navier-Stokes equations: controllability by means of low modes forcing, J. Math. Fluid Mech. 7 (2005), no. 1, 108-152.
[AS06]
__, Controllability of $2 D$ Euler and Navier-Stokes equations by degenerate forcing, Comm. Math. Phys. 265 (2006), no. 3, 673-697.
[AS08] , Solid controllability in fluid dynamics, Instability in models connected with fluid flows. I, 1-35, Springer, 2008.
[BG03] Y. Brenier and W. Gangbo, $L^{p}$ approximation of maps by diffeomorphisms, Calc. Var. Partial Differential Equations 16 (2003), no. 2, 147-164.
[CF88] P. Constantin and C. Foias, Navier-Stokes Equations, University of Chicago Press, Chicago, IL, 1988.
[CL12] J.-M. Coron and P. Lissy, Local null controllability of the threedimensional Navier-Stokes system with a distributed control having two vanishing components, Invent. math. 198 (2014), no. 3, 833-880.
[Cor07] J.-M. Coron, Control and Nonlinearity, Mathematical Surveys and Monographs, vol. 136, AMS, Providence, RI, 2007.
[EM01] W. E and J. C. Mattingly, Ergodicity for the Navier-Stokes equation with degenerate random forcing: finite-dimensional approximation, Comm. Pure Appl. Math. 54 (2001), no. 11, 1386-1402.
[Gam78] R. V. Gamkrelidze, Principles of Optimal Control Theory, Plenum Press, New York, 1978.
[GH10] O. Glass and T. Horsin, Approximate Lagrangian controllability for the 2-D Euler equation. Application to the control of the shape of vortex patches, J. Math. Pures Appl. (9) 93 (2010), no. 1, 61-90.
[GH12] _, Prescribing the motion of a set of particles in a threedimensional perfect fluid, SIAM J. Control Optim. 50 (2012), no. 5, 2726-2742.
[Hal80] J. K. Hale, Ordinary Differential Equations, second ed., Robert E. Krieger Publishing Co. Inc., Huntington, N.Y., 1980.
[HM06] M. Hairer and J. C. Mattingly, Ergodicity of the 2D Navier-Stokes equations with degenerate stochastic forcing, Ann. of Math. (2) 164 (2006), no. 3, 993-1032.
[Jac85] N. Jacobson, Basic Algebra. I, second ed., W. H. Freeman and Company, New York, 1985.
[Kry71] A. B. Krygin, Extension of diffeomorphisms that preserve volume, Funkcional. Anal. i Priložen. 5 (1971), no. 2, 72-76.
[KW09] B. Khesin and R. Wendt, The Geometry of Infinite-dimensional Groups, Series of Modern Surveys in Mathematics, vol. 51, SpringerVerlag, Berlin, 2009.
[Ner10] H. Nersisyan, Controllability of 3D incompressible Euler equations by a finite-dimensional external force, ESAIM Control Optim. Calc. Var. 16 (2010), no. 3, 677-694.
[Ner11] , Controllability of the 3D compressible Euler system, Comm. Partial Differential Equations 36 (2011), no. 9, 1544-1564.
[Rod06] S. S. Rodrigues, Navier-Stokes equation on the Rectangle: Controllability by means of low modes forcing, J. Dyn. Control Syst., 12 (2006), no. 4, 517-562.
[Rod07] , Controllability of nonlinear PDE's on compact Riemannian manifolds, Proceedings WMCTF'07, (2007), 462-493.
[Rod08] _ Methods of Geometric Control Theory in Problems of Mathematical Physics, PhD Thesis, Universidade de Aveiro, Portugal, 2008.
[Rom04] M. Romito, Ergodicity of the finite dimensional approximation of the 3D Navier-Stokes equations forced by a degenerate noise, J. Statist. Phys. 114 (2004), no. 1-2, 155-177.
[Sar12] A. Sarychev, Controllability of the cubic Schrödinger equation via a low-dimensional source term, Math. Control Relat. Fields 2 (2012), no. $3,247-270$.
[Shi06] A. Shirikyan, Approximate controllability of three-dimensional NavierStokes equations, Comm. Math. Phys. 266 (2006), no. 1, 123-151.
[Shi07a] , Contrôlabilité exacte en projections pour les équations de Navier-Stokes tridimensionnelles, Ann. Inst. H. Poincaré Anal. Non Linéaire 24 (2007), no. 4, 521-537.
[Shi07b] , Controllability of nonlinear PDEs: Agrachev-Sarychev approach, Journées EDP, Evian, (2007), Exposé no. IV, pages 1-11.
[Shi08a] , Euler equations are not exactly controllable by a finitedimensional external force, Physica D 237 (2008), no. 10-11, 13171323.
[Shi08b] , Navier-stokes and Euler equations: Cauchy problem and controllability, PhD course delivered at SISSA, Trieste, Italy, 2008. URL: http://shirikyan.u-cergy. fr/lecturenotes.html.
[Shi10] , Control theory for the Burgers equation: Agrachev-Sarychev approach, PhD course delivered at the University of Iasi, Romania, 2010. URL: http://shirikyan. u-cergy.fr/lecturenotes.html.
[Shi13] , Approximate controllability of the viscous Burgers equation on the real line, Geometric control theory and sub-Riemannian geometry, Springer INdAM Ser. 5 (2014), 351-370.
[Tay96] M. E. Taylor, Partial Differential Equations III, Springer-Verlag, N.Y., 1996.


[^0]:    *Laboratoire de Mathématiques, UMR CNRS 8100, Université de Versailles-Saint-Quentin-en-Yvelines, F-78035 Versailles, France

[^1]:    ${ }^{1}$ Here $[a]$ stands for the integer part of $a \in \mathbb{R}$.

[^2]:    ${ }^{2}$ The integer $p$ may depend on $\eta_{1}$.

[^3]:    ${ }^{3}$ The result of [BG03] is stated for a cube, but it remains valid also in the case of a torus.

[^4]:    ${ }^{4}$ This inequality is proved with the help of a decomposition in the eigenbasis $\left\{e_{j}\right\}$ of $L$ :

    $$
    \left\|L^{r} e^{-t L} u\right\|^{2}=\sum_{j=1}^{\infty} \alpha_{j}^{2 r} e^{-2 t \alpha_{j}} u_{j}^{2} \leq C_{r} t^{-r}\|u\|^{2}
    $$

    where $u_{j}:=\left\langle u, e_{j}\right\rangle$ and $\alpha_{j}$ is the eigenvalue corresponding to $e_{j}$.

[^5]:    ${ }^{5}$ In that case one has $\operatorname{dim} E(\mathcal{K})=12$. In Proposition 4.9, we give an example of a 6 dimensional saturating space.

[^6]:    ${ }^{6}$ Note that $\mathcal{E}$ is not necessarily a generator of $\mathbb{Z}^{3}$. For example, $m=(2,0,0), n=$ $(0,1,0), r=(0,0,1)$ is a basis in $\mathbb{R}^{3}$, but not a generator of $\mathbb{Z}^{3}$, since the greatest common divisor in Theorem 4.2 is equal to 2 .

