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PERSISTENCE OF INTEGRATED STABLE PROCESSES

CHRISTOPHE PROFETA AND THOMAS SIMON

ABSTRACT. We compute the persistence exponent of the integral of a stable Lévy process in terms of its self-similarity and positivity parameters. This solves a problem raised by Z. Shi (2003). Along the way, we investigate the law of the stable process L evaluated at the first time its integral X hits zero, when the bivariate process (X, L) starts from a coordinate axis. This extends classical formulæ by McKean (1963) and Gor'kov (1975) for integrated Brownian motion.

1. INTRODUCTION AND STATEMENT OF THE RESULTS

Let $X = \{X_t, t \geq 0\}$ be a real process starting at zero and $T_x = \inf\{t > 0, X_t > x\}$ be its first-passage time above a positive level x . Studying the law of T_x is a classical problem in probability theory. In general, it is difficult to obtain an explicit expression of this law. However, it has been observed that in many interesting cases the survival function has a polynomial decay:

$$\mathbb{P}[T_x > t] = t^{-\theta+o(1)}, \quad t \rightarrow +\infty, \quad (1.1)$$

where θ is a positive constant which is called the persistence exponent, and usually does not depend on x . The computation of persistence exponents has many connections to various problems in probability and mathematical physics, and we refer to the recent surveys [2, 3] for more information on this topic. In this paper we consider this problem for the process

$$X_t = \int_0^t L_s ds,$$

where $L = \{L_t, t \geq 0\}$ is a strictly α -stable Lévy process starting from zero, with law \mathbb{P} . Our process L is normalized to have characteristic exponent

$$\Psi(\lambda) = \log(\mathbb{E}[e^{i\lambda L_1}]) = -(i\lambda)^\alpha e^{-i\pi\alpha\rho \operatorname{sgn}(\lambda)}, \quad \lambda \in \mathbb{R}, \quad (1.2)$$

where $\alpha \in (0, 2]$ is the self-similarity parameter and $\rho = \mathbb{P}[L_1 \geq 0]$ is the positivity parameter. We refer to [14, 21] for classic accounts on stable laws and processes. The strict stability implies the $(1/\alpha)$ -self-similarity of L and the $(1 + 1/\alpha)$ -self-similarity of X , in other words that

$$\{L_{kt}, t \geq 0\} \stackrel{d}{=} \{k^{1/\alpha} L_t, t \geq 0\} \quad \text{and} \quad \{X_{kt}, t \geq 0\} \stackrel{d}{=} \{k^{1+1/\alpha} X_t, t \geq 0\}$$

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for all $k > 0$. When $\alpha = 2$, one has $\rho = 1/2$ and $\Psi(\lambda) = -\lambda^2$, so that $L = \sqrt{2}B$ is a rescaled Brownian motion. When $\alpha = 1$, one has $\rho \in (0, 1)$ and L is a Cauchy process with a linear drift. When $\alpha \in (0, 1) \cup (1, 2)$ the characteristic exponent takes the more familiar form

$$\Psi(\lambda) = -\kappa_{\alpha,\rho}|\lambda|^\alpha(1 - i\beta \tan(\pi\alpha/2) \operatorname{sgn}(\lambda)),$$

where $\beta \in [-1, 1]$ is an asymmetry parameter, whose connection with the positivity parameter is given by Zolotarev's formula:

$$\rho = \frac{1}{2} + \frac{1}{\pi\alpha} \arctan(\beta \tan(\pi\alpha/2)),$$

and $\kappa_{\alpha,\rho} = \cos(\pi\alpha(\rho - 1/2)) > 0$ is a scaling constant. The latter could have taken any positive value, changing the normalization (1.2) accordingly, without incidence on our purposes below. One has $\rho \in [0, 1]$ if $\alpha < 1$ and $\rho \in [1 - 1/\alpha, 1/\alpha]$ if $\alpha > 1$. When $\alpha > 1$ and $\rho = 1/\alpha$ the process L has no positive jumps, whereas it has no negative jumps when $\alpha > 1$ and $\rho = 1 - 1/\alpha$. When $\alpha < 1$ and $\rho = 0$ or $\rho = 1$, the process $|L|$ is a stable subordinator and has increasing sample paths, a situation which will be implicitly excluded throughout this paper. In this case, the process X is indeed also monotonous and the survival function in (1.1) either is one or decays towards zero at an exponential speed - see [2] p.4 for details.

When $\alpha = 2$, the bivariate process (X, L) is Gaussian with explicit covariance function and transition density, providing also the basic example of a degenerate diffusion process - see [10] for details and references. When $\alpha < 2$, the process (X, L) is Non-Gaussian α -stable in the broad sense of [14]. The process (X, L) is a strong Markov process, which is sometimes called the Kolmogorov process in the literature. In the following we will set $\mathbb{P}_{(x,y)}$ for the law of (X, L) starting at $(x, y) \in \mathbb{R}^2$. Our main concern in this paper is the hitting time of zero for X :

$$T_0 = \inf\{t > 0, X_t = 0\}.$$

Since $|L|$ is not a subordinator, a simple argument using self-similarity and the zero-one law for Markov processes - see below Lemma 6 for details - shows that $\mathbb{P}_{(0,0)}[T_0 = 0] = 1$, in other words that the origin is regular for the vertical axis. If $x < 0$ or $x = 0$ and $y < 0$, the continuity of the sample paths of X show that a.s. $T_0 = \inf\{t > 0, X_t > 0\}$, and it will be checked in Lemma 6 below that T_0 is also a.s. finite. If $x > 0$ or $x = 0$ and $y > 0$, the law of T_0 is obviously deduced from that of the latter situation in considering the dual Lévy process $-L$.

When $(x, y) \neq (0, 0)$, the difficulty to obtain concrete informations on the law of T_0 under $\mathbb{P}_{(x,y)}$ comes from the fact that X itself is not a Markov process. In the Brownian case for example, the density function of T_0 is expressed through quite intricate integral formulæ - see [2] pp.15-16 and the references therein. On the other hand, some universal estimates can be obtained for the behaviour of the distribution function $\mathbb{P}_{(x,y)}[T_0 \leq t]$ as $t \rightarrow 0$, using self-similarity and Gaussian or stable upper tails for the supremum process - see e.g. Section 10.4 in [14]. But it is well-known that the study of $\mathbb{P}_{(x,y)}[T_0 > t]$ as $t \rightarrow +\infty$ is a harder problem, where a more exotic behaviour is expected.

Throughout the paper, for any real functions f and g we will use the standard notation $f(t) \asymp g(t)$ as $t \rightarrow +\infty$ to express the fact that there exist κ_1, κ_2 positive finite constants such that $\kappa_1 f(t) \leq g(t) \leq \kappa_2 f(t)$ as $t \rightarrow +\infty$. Our main result is the following.

Theorem A. *Assume $x < 0$ or $x = 0, y < 0$. One has*

$$\mathbb{P}_{(x,y)}[T_0 > t] \asymp t^{-\theta}, \quad t \rightarrow +\infty,$$

with $\theta = \rho/(1 + \alpha(1 - \rho))$.

In the Brownian case $\alpha = 2$, one has $\theta = 1/4 = \rho/2$ and this estimate has been known since the works of M. Goldman - see Proposition 2 in [6], with a more precise formulation on the density function of T_0 , following the seminal article of McKean [12]. The universality of the persistence exponent $1/4$ for integrals of real Lévy processes having exponential moments on both sides has been shown in [1], with the help of strong approximation arguments. Recently, it was proved in [4] that all integrated real random walks with finite variance have also $1/4$ as persistence exponent, extending [20] for the particular case of the integrated simple random walk. Let us also mention that the survival function of the n^{th} hitting time of zero for the integrated Brownian motion exhibits the same power decay up to a logarithmic term in $ct^{-1/4}(\ln(t))^{n-1}$ with an explicit constant c , as shown by the first author in [13].

In the case $1 < \alpha < 2$ and with no negative jumps, that is $\rho = 1 - 1/\alpha$, one obtains $\theta = (\alpha - 1)/2\alpha = \rho/2$, an estimate which had been proved by the second author in [18] with different techniques and a less precise formulation than Theorem A for the lower bound, involving a logarithmic correction term. It is worth mentioning that the same persistence exponent $(\alpha - 1)/2\alpha$ appears for the integrals of random walks attracted towards this spectrally positive Lévy process - see Remark 1.2 in [4].

It has been conjectured in [2] - see Conjecture 4 therein - that the persistence exponent should be $\rho/2$ in general. This expected value should be compared with a classical result of Bingham stating that the persistence exponent is ρ for the stable process L - see (2.16) in [2] and the references of Section 2.2 therein. The admissible set of (α, ρ) and Theorem A entail that $\theta > \rho/2$ as soon as L has negative jumps, hence providing a negative answer to this conjecture. The fact that θ is an increasing function of the positivity parameter ρ matches the intuition, however it is harder to explain heuristically why it is also a decreasing function of α .

Specifying $x = -1$ and $y = 0$ in Theorem A entails by self-similarity the following lower tail probability estimate

$$\mathbb{P}[X_1^* \leq \varepsilon] \asymp \varepsilon^{\frac{\theta\alpha}{\alpha+1}}, \quad \varepsilon \rightarrow 0,$$

with the notation $X_1^* = \sup\{X_t, t \leq 1\}$. Some heuristics on the subordination of X by the inverse local time of L when $\alpha > 1$ had led to the conjecture, formulated in Part 2 of [16], that in the symmetric case $\rho = 1/2$ one should have $\mathbb{P}[X_1^* \leq \varepsilon] = \varepsilon^{(\alpha-1)^+/2(\alpha+1)+o(1)}$ as $\varepsilon \rightarrow 0$. The invalidity of this conjecture as soon as α is close enough to 1 had been observed in [19]. Theorem A shows that Shi's exponent is the right one only for integrated Brownian motion: in the symmetric case one has $\theta\alpha/(\alpha+1) = \alpha/(\alpha+1)(\alpha+2) \geq (\alpha-1)^+/2(\alpha+1)$, with an equality only if $\alpha = 2$. Let us mention in passing that lower tail probabilities offer some

challenging problems for Gaussian processes - see [11, 15].

Our method to prove Theorem A hinges upon the random variable L_{T_0} , the so-called hitting place of (X, L) on the vertical axis, which has been extensively studied in the Brownian case - see [12, 6, 9, 10]. Notice that this random variable is positive under $\mathbb{P}_{(x,y)}$ if $x < 0$ or $x = 0$ and $y < 0$. The reason why it is connected to the persistence exponent comes from the following heuristical equivalence for fractional moments

$$\mathbb{E}_{(x,y)}[T_0^s] < +\infty \quad \Leftrightarrow \quad \mathbb{E}_{(x,y)}[L_{T_0}^{\alpha s}] < +\infty$$

for all $s > 0$, which had been conjectured in [18] p.176, and turns out to be true as a consequence of Theorem A and Lemma 11 below. The precise relationship between the upper tails of T_0 and that of L_{T_0} follows from a series of probabilistic estimates which are the matter of Section 4.

In this paper we also provide a rather complete description of the law of the random variable L_{T_0} when (X, L) starts from a coordinate axis. To express our second main result, we need some further notation. For every $\mu \in (0, 1)$, introduce the μ -Cauchy random variable \mathbf{C}_μ , with density

$$\frac{\sin(\pi\mu)}{\pi\mu(x^2 + 2\cos(\pi\mu)x + 1)}1_{\{x \geq 0\}}.$$

Our above denomination comes from the case $\mu = 1/2$, where $\mathbf{C}_{1/2}$ is the half-Cauchy distribution. If X is a positive random variable and $\nu \in \mathbb{R}$ is such that $\mathbb{E}[X^\nu] < \infty$, the positive random variable $X^{(\nu)}$ defined by

$$\mathbb{E}[f(X^{(\nu)})] = \frac{\mathbb{E}[X^\nu f(X)]}{\mathbb{E}[X^\nu]}$$

for all $f : \mathbb{R}^+ \rightarrow \mathbb{R}$ bounded continuous, is known as the size bias of order ν of X . Observe that when X is absolutely continuous, the density of $X^{(\nu)}$ is obtained in multiplying that of X by x^ν and renormalizing. Introduce finally the parameters

$$\gamma = \frac{\rho\alpha}{1+\alpha} \quad \text{and} \quad \chi = \frac{\rho\alpha}{1+\alpha(1-\rho)} = \alpha\theta.$$

Notice that from the admissible set for (α, ρ) , we have $\gamma \in (0, 1/2)$ and $\chi \in (0, 1)$.

Theorem B. (i) *For every $y < 0$, under $\mathbb{P}_{(0,y)}$ one has*

$$L_{T_0} \stackrel{d}{=} |y|(\mathbf{C}_\chi^{1-\gamma})^{(1)}.$$

(ii) *For every $x < 0$, under $\mathbb{P}_{(x,0)}$ the positive random variable L_{T_0} has Mellin transform*

$$\mathbb{E}_{(x,0)}[L_{T_0}^{s-1}] = \frac{(1+\alpha)^{\frac{1-s}{1+\alpha}} \Gamma(\frac{\alpha+2}{\alpha+1}) \Gamma(\frac{1-s}{\alpha+1}) \sin(\pi\gamma)}{\Gamma(\frac{s}{\alpha+1}) \Gamma(1-s) \sin(\pi s(1-\gamma))} |x|^{\frac{s-1}{\alpha+1}}, \quad |s| < 1/(1-\gamma).$$

The proof of this result is given in Section 3, following some preliminary computations involving oscillating integrals and the Fourier transform of X_t , performed in Section 2. Observe

that the density in (i) above is explicit and reads for example

$$\frac{3}{2\pi} \frac{|y|^{1/2} z^{3/2}}{|y|^3 + z^3} 1_{\{z \geq 0\}}$$

in the Brownian case, a formula originally proved by McKean in [12] - see also formulæ (1) and (2) in [9]. As is well-known, the Cauchy random variable appears in exit or winding problems for two-dimensional Brownian motion. The fact that it is also connected with similar problems for general integrated stable processes is perhaps more surprising.

An interesting consequence of (ii) is that the Mellin transform can be inverted in the Cauchy case $\alpha = 1$ and exhibits the same type of law as in (i): one obtains

$$L_{T_0} \stackrel{d}{=} \sqrt{2|x|} (\mathbf{C}_\delta^{1-\gamma})^{(1)}$$

with the notation $\delta = (1 + \chi)/2$. The Mellin transform of (ii) can also be simply inverted in the Brownian case in terms of Beta and Gamma random variables, shedding some new light on a formula by Gor'kov [7] which was of the analytical type, and in the case $\alpha < 1$ in terms of positive stable random variables. The Mellin inversion is however more complicated when $\alpha \in (1, 2)$, and involves no classical random variables in general - see Section 3.3 below for details.

2. PRELIMINARY COMPUTATIONS

The following lemma, which we could not locate in the literature, will be useful in the sequel.

Lemma 1. *Let $\nu \in (0, 1)$ and X be a real random variable such that $\mathbb{E}[|X|^{-\nu}] < \infty$. One has*

$$\int_0^\infty \lambda^{\nu-1} \mathbb{E}[\cos(\lambda X)] d\lambda = \Gamma(\nu) \cos(\pi\nu/2) \mathbb{E}[|X|^{-\nu}]$$

and

$$\int_0^\infty \lambda^{\nu-1} \mathbb{E}[\sin(\lambda X)] d\lambda = \Gamma(\nu) \sin(\pi\nu/2) \mathbb{E}[|X|^{-\nu} \text{sgn}(X)].$$

Proof. The generalized Fresnel integral which is computed e.g. in formula (37) p.13 of [5] shows that for all $u \neq 0, \nu \in (0, 1)$, one has

$$\int_0^\infty \lambda^{\nu-1} \cos(\lambda u) d\lambda = \Gamma(\nu) \cos(\pi\nu/2) |u|^{-\nu}. \quad (2.1)$$

The first statement of the lemma is hence simply a switching of the expectation and the integral. However, we cannot apply Fubini's theorem directly. Set μ for the probability distribution of X . From (2.1) and an integration by parts, we get

$$\begin{aligned} \Gamma(\nu) \cos(\pi\nu/2) \mathbb{E}[|X|^{-\nu}] &= \int_{\mathbb{R}} \mu(du) \left(\int_0^\infty \lambda^{\nu-1} \cos(\lambda u) d\lambda \right) \\ &= (1 - \nu) \int_{\mathbb{R}} \mu(du) \left(\int_0^\infty \frac{\sin(\lambda u)}{u} \lambda^{\nu-2} d\lambda \right). \end{aligned}$$

Since

$$\begin{aligned} \int_{\mathbb{R}} \mu(du) \left(\int_0^\infty \left| \frac{\sin(\lambda u)}{u} \right| \lambda^{\nu-2} d\lambda \right) &\leq \int_{\mathbb{R}} \mu(du) \left(\int_0^\infty \left(\lambda \wedge \frac{1}{|u|} \right) \lambda^{\nu-2} d\lambda \right) \\ &\leq \frac{\mathbb{E}[|X|^{-\nu}]}{\nu(1-\nu)} < +\infty, \end{aligned}$$

we may now apply Fubini's theorem and obtain

$$\Gamma(\nu) \cos(\pi\nu/2) \mathbb{E}[|X|^{-\nu}] = (1-\nu) \int_0^\infty \lambda^{\nu-2} d\lambda \left(\int_{\mathbb{R}} \frac{\sin(\lambda u)}{u} \mu(du) \right).$$

The dominated convergence theorem entails that the function

$$\psi(\lambda) = \int_{\mathbb{R}} \frac{\sin(\lambda u)}{u} \mu(du)$$

is differentiable, with derivative

$$\psi'(\lambda) = \int_{\mathbb{R}} \cos(\lambda u) \mu(du) = \mathbb{E}[\cos(\lambda X)].$$

Thus, another integration by parts yields

$$\Gamma(\nu) \cos(\pi\nu/2) \mathbb{E}[|X|^{-\nu}] = \int_0^\infty \lambda^{\nu-1} \mathbb{E}[\cos(\lambda X)] d\lambda - \left[\lambda^{\nu-1} \psi(\lambda) \right]_0^{+\infty}$$

and it remains to prove that the bracket is zero. On the one hand, one has

$$\lambda^{\nu-1} |\psi(\lambda)| \leq \lambda^\nu \rightarrow 0 \quad \text{as } \lambda \rightarrow 0.$$

On the other hand, using

$$\lambda^{\nu-1} \left| \frac{\sin(\lambda u)}{u} \right| \leq \lambda^{\nu-1} \frac{|\sin(\lambda u)|^{1-\nu}}{|u|} \leq |u|^{-\nu},$$

and the dominated convergence theorem, we see that $\lambda^{\nu-1} |\psi(\lambda)| \rightarrow 0$ as $\lambda \rightarrow +\infty$. This completes the proof of the first statement of the lemma. The second statement may be handled similarly with the help of the formula

$$\int_0^\infty \lambda^{\nu-1} \sin(\lambda x) d\lambda = \Gamma(\nu) \sin(\pi\nu/2) \operatorname{sgn}(x) |x|^{-\nu}, \quad |\nu| < 1, \quad (2.2)$$

which is given e.g. in (38) p.13 in [5].

□

Lemma 2. *For all $x, y \in \mathbb{R}$ and $t \geq 0$ one has*

$$\log(\mathbb{E}_{(x,y)}[e^{i\lambda X_t}]) = i\lambda(x + yt) - \frac{t^{\alpha+1}}{\alpha+1} (i\lambda)^\alpha e^{-i\pi\alpha\rho \operatorname{sgn}(\lambda)}, \quad \lambda \in \mathbb{R}.$$

Proof. It is clearly enough to consider the case $x = y = 0$. Integrating by parts yields the following representation of X_t as a stable integral:

$$X_t = \int_0^\infty (t-s)^+ dL_s = \int_0^\infty (t-x)^+ M(dx),$$

where M is an α -stable random measure on \mathbb{R}^+ with Lebesgue control measure and constant skewness intensity $\beta(x) = \beta$ - see Example 3.3.3 in [14]. In the case $\alpha \neq 1$, the statement of the lemma is a direct consequence of Proposition 3.4.1 (i) in [14], reformulated with the (α, ρ) parametrization. In the case $\alpha = 1, \rho = 1/2$ we use Proposition 3.4.1 (ii) in [14] (with $\beta = 0$). The case $\alpha = 1, \rho \neq 1/2$ follows from the symmetric case in adding a drift coefficient μt for some $\mu \neq 0$, which integrates in $\mu t^2/2$. \square

We now set

$$s_{\alpha, \rho} = \frac{\sin(\pi\alpha(\rho - 1/2))}{\alpha + 1} \in (-1, 1) \quad \text{and} \quad c_{\alpha, \rho} = \frac{\cos(\pi\alpha(\rho - 1/2))}{\alpha + 1} \in (0, 1).$$

The proposition gives a representation for the Mellin transform of $X_t^+ = X_t \mathbf{1}_{\{X_t > 0\}}$.

Proposition 3. *For all $x, y \in \mathbb{R}, t > 0$ and $\nu \in (0, 1)$ one has*

$$\mathbb{E}_{(x, y)}[(X_t^+)^{-\nu}] = \frac{\Gamma(1 - \nu)}{\pi} \int_0^\infty \lambda^{\nu-1} e^{-c_{\alpha, \rho} \lambda^{\alpha+1}} \sin(\lambda(x + yt) + s_{\alpha, \rho} \lambda^{\alpha+1} + \pi\nu/2) d\lambda.$$

Proof. Since X_t is a stable random variable, it has a bounded density and $\mathbb{E}_{(x, y)}[(X_t^+)^{-\nu}]$ is hence finite for all $\nu \in (0, 1)$. By Lemma 2 we have

$$\log(\mathbb{E}_{(x, y)}[e^{i\lambda X_t}]) = i\lambda(x + yt) - \lambda^{\alpha+1} (c_{\alpha, \rho} - i s_{\alpha, \rho}), \quad \lambda \geq 0.$$

Taking the real part and integrating with respect to $\lambda^{\nu-1}$ on $]0, +\infty[$, we deduce

$$\begin{aligned} \int_0^\infty \lambda^{\nu-1} e^{-c_{\alpha, \rho} \lambda^{\alpha+1}} \cos(\lambda(x + yt) + s_{\alpha, \rho} \lambda^{\alpha+1}) d\lambda &= \int_0^\infty \lambda^{\nu-1} \mathbb{E}_{(x, y)}[\cos(\lambda X_t)] d\lambda \\ &= \Gamma(\nu) \cos\left(\frac{\pi\nu}{2}\right) \mathbb{E}_{(x, y)}[|X_t|^{-\nu}], \end{aligned}$$

where the second equality comes from Lemma 1. Similarly, taking the imaginary part entails

$$\int_0^\infty \lambda^{\nu-1} e^{-c_{\alpha, \rho} \lambda^{\alpha+1}} \sin(\lambda(x + yt) + s_{\alpha, \rho} \lambda^{\alpha+1}) d\lambda = \Gamma(\nu) \sin\left(\frac{\pi\nu}{2}\right) \mathbb{E}_{(x, y)}[|X_t|^{-\nu} \text{sgn}(X_t)].$$

Multiplying the first relation by $\sin(\pi\nu/2)$, the second by $\cos(\pi\nu/2)$, and summing, we finally obtain

$$\Gamma(\nu) \sin(\pi\nu) \mathbb{E}_{(x, y)}[(X_t^+)^{-\nu}] = \int_0^\infty \lambda^{\nu-1} e^{-c_{\alpha, \rho} \lambda^{\alpha+1}} \sin(\lambda(x + yt) + s_{\alpha, \rho} \lambda^{\alpha+1} + \pi\nu/2) d\lambda,$$

which yields the required expression by the complement formula for the Gamma function. \square

Our last proposition provides some crucial computations for the proof of Theorem B.

Proposition 4. *Set $\nu \in (\alpha/(\alpha + 1), 1)$ and $s = (1 - \nu)(\alpha + 1) \in (0, 1)$.*

(i) *For every $y > 0$, one has*

$$\int_0^\infty \mathbb{E}_{(0, y)}[(X_t^+)^{-\nu}] dt = (\alpha + 1)^{1-\nu} \Gamma(1 - s) \sin(\pi s(1 - \gamma)) \frac{\Gamma(1 - \nu)^2}{\pi} y^{s-1}.$$

(ii) For every $y < 0$, one has

$$\int_0^\infty \mathbb{E}_{(0,y)}[(X_t^+)^{-\nu}] dt = (\alpha + 1)^{1-\nu} \Gamma(1-s) \sin(\pi\gamma s) \frac{\Gamma(1-\nu)^2}{\pi} |y|^{s-1}.$$

(iii) For every $x < 0$, one has

$$\int_0^\infty \mathbb{E}_{(x,0)}[(X_t^+)^{-\nu}] dt = (\alpha + 1)^{-\frac{\alpha}{\alpha+1}} \Gamma\left(\frac{1-s}{\alpha+1}\right) \sin(\pi\gamma) \Gamma\left(\frac{1}{\alpha+1}\right) \frac{\Gamma(1-\nu)}{\pi} |x|^{\frac{s-1}{\alpha+1}}.$$

Proof. Suppose first $x = 0$ and $y \in \mathbb{R}$. Integrating the expression on the right-hand side of Proposition 3 yields a double integral of the form

$$\begin{aligned} I_\nu &= \int_0^\infty \left(\int_0^\infty \lambda^{\nu-1} e^{-c_{\alpha,\rho} \lambda^\alpha t^{1+\alpha}} \sin(\lambda y t + s_{\alpha,\rho} \lambda^\alpha t^{1+\alpha} + \nu\pi/2) d\lambda \right) dt \\ &= \int_0^\infty \left(\int_0^\infty r^{\nu-1} e^{-c_{\alpha,\rho} r^\alpha} \sin(ry t^{-1/\alpha} + s_{\alpha,\rho} r^\alpha + \nu\pi/2) dr \right) t^{-\nu(1+1/\alpha)} dt \\ &= \int_0^\infty \left(\int_0^\infty r^{\nu-1} e^{-c_{\alpha,\rho} r^\alpha} \sin(ry u + s_{\alpha,\rho} r^\alpha + \nu\pi/2) dr \right) \alpha u^{-s} du \\ &= \int_0^\infty \left(\int_0^\infty \alpha u^{-s} \sin(ry u + s_{\alpha,\rho} r^\alpha + \nu\pi/2) du \right) r^{\nu-1} e^{-c_{\alpha,\rho} r^\alpha} dr, \end{aligned}$$

where the first, resp. second, equality comes from the change of variable $\lambda t^{1+1/\alpha} = r$, resp. $u = t^{-1/\alpha}$, and the switching of the integrals in the third equality is made exactly as in Lemma 1, using the fact that $s \in (0, 1)$ and $s + \nu > 1$.

Suppose first $y > 0$. We start by computing the integral in u with the help of formulæ (2.1) and (2.2) and some trigonometry:

$$\alpha \int_0^\infty u^{-s} \sin(ry u + s_{\alpha,\rho} r^\alpha + \nu\pi/2) du = \alpha \Gamma(1-s) \cos((s-\nu)\pi/2 - s_{\alpha,\rho} r^\alpha) (yr)^{s-1}.$$

We then compute the integral in r with the change of variable $z = r^\alpha$, using the notation $Z = e^{i\pi\alpha(\rho-1/2)}$:

$$\begin{aligned} I_\nu &= \alpha \Gamma(1-s) y^{s-1} \int_0^\infty r^{\alpha(1-\nu)-1} e^{-c_{\alpha,\rho} r^\alpha} \cos((s-\nu)\pi/2 - s_{\alpha,\rho} r^\alpha) dr \\ &= \Gamma(1-s) y^{s-1} \int_0^\infty z^{-\nu} e^{-c_{\alpha,\rho} z} \cos((s-\nu)\pi/2 - s_{\alpha,\rho} z) dz \\ &= (1+\alpha)^{1-\nu} \Gamma(1-s) \Gamma(1-\nu) \Re(e^{i\pi(s-\nu)/2} Z^{\nu-1}) y^{s-1} \\ &= (1+\alpha)^{1-\nu} \Gamma(1-s) \Gamma(1-\nu) \sin(\pi s(1-\gamma)) y^{s-1}, \end{aligned}$$

where the third line follows after some algebraic simplifications. By Proposition 3, this completes the proof of (i).

Suppose now $y < 0$. An analogous computation to the above shows that

$$I_\nu = \alpha \int_0^\infty u^{-s} \sin(ry u + s_{\alpha,\rho} r^\alpha + \nu\pi/2) du = \alpha \Gamma(1-s) \sin((s+\nu-1)\pi/2 + s_{\alpha,\rho} r^\alpha) |yr|^{s-1}.$$

The integral in r is then computed in the same way and yields the formula

$$\begin{aligned} I_\nu &= (1 + \alpha)^{1-\nu} \Gamma(1-s) \Gamma(1-\nu) \Im(e^{i\pi(s+\nu-1)/2} \bar{Z}^{\nu-1}) |y|^{s-1} \\ &= (1 + \alpha)^{1-\nu} \Gamma(1-s) \Gamma(1-\nu) \sin(\pi\gamma s) |y|^{s-1}, \end{aligned}$$

which completes the proof of (ii) by Proposition 3.

We last suppose $x < 0$ and $y = 0$. We again integrate the expression on the right-hand side of Proposition 3, making the changes of variable $\lambda t^{1+1/\alpha} = r$ and $u = t^{-(1+1/\alpha)}$. This yields a double integral of the form

$$\frac{\alpha}{\alpha+1} \int_0^\infty \left(\int_0^\infty r^{\nu-1} e^{-c_{\alpha,\rho} r^\alpha} \sin(rxu + s_{\alpha,\rho} r^\alpha + \nu\pi/2) dr \right) u^{-(s+\alpha)/(1+\alpha)} du,$$

where we can switch the orders of integration as in Lemma 1 because $(s+\alpha)/(1+\alpha) \in (0, 1)$ and $(s+\alpha)/(1+\alpha) + \nu > 1$. We then compute the integral in u similarly as above and find

$$\frac{\alpha}{\alpha+1} \Gamma\left(\frac{1-s}{\alpha+1}\right) \sin(\pi\alpha/2(\alpha+1) + s_{\alpha,\rho} r^\alpha) |xr|^{\frac{s-1}{\alpha+1}}.$$

We finally compute the integral in r with the change of variable $r = z^{1/\alpha}$, and get after some algebraic manipulations

$$I_\nu = (\alpha+1)^{-\frac{\alpha}{\alpha+1}} \sin(\pi\gamma) \Gamma\left(\frac{1}{\alpha+1}\right) \Gamma\left(\frac{1-s}{\alpha+1}\right) |x|^{\frac{s-1}{\alpha+1}},$$

which completes the proof of (iii) by Proposition 3. □

Remark 5. It seems hard to find an explicit formula in general for

$$\int_0^\infty \mathbb{E}_{(x,y)}[(X_t^+)^{-\nu}] dt$$

when (x, y) is not on a coordinate axis. In the symmetric Cauchy case, some further computations show that the integral equals

$$\frac{1}{\sin(\pi\nu)} \Im\left(\int_0^\infty (-(x + yt + it^2/2))^{-\nu} dt\right).$$

This can be rewritten with the hypergeometric function, apparently not in a tractable manner when $xy \neq 0$.

3. PROOF OF THEOREM B

The following lemma shows the aforementioned and intuitively obvious fact that T_0 is a proper random variable for any starting point.

Lemma 6. *For all $x, y \in \mathbb{R}$ one has $\mathbb{P}_{(x,y)}[T_0 < +\infty] = 1$.*

Proof. Suppose first $x = -1$ and $y = 0$. Then

$$\mathbb{P}_{(-1,0)}[T_0 = +\infty] = \mathbb{P}_{(0,0)}[X_\infty^* < 1] = \mathbb{P}_{(0,0)}[X_\infty^* = 0] \leq \mathbb{P}_{(0,0)}[X_1 \leq 0] < 1,$$

where the second equality comes from the self-similarity of X and the strict inequality from the fact that X_1 is a two-sided stable random variable - see Lemma 2. On the other hand, setting $T = \inf\{t > 0, X_t > 0\}$, it is clear by self-similarity that under the probability measure $\mathbb{P}_{(0,0)}$ one has

$$T \stackrel{d}{=} kT$$

for all $k > 0$. In particular, $\mathbb{P}_{(0,0)}[T \in \{0, +\infty\}] = 1$. Moreover, the zero-one law for the Markov process (X, L) entails that $\mathbb{P}_{(0,0)}[T = 0]$ is 0 or 1. Since $\mathbb{P}_{(0,0)}[T = +\infty] = \mathbb{P}_{(0,0)}[X_\infty^* = 0] < 1$, we get $\mathbb{P}_{(0,0)}[T = +\infty] = 0$ whence $\mathbb{P}_{(-1,0)}[T_0 = +\infty] = 0$ as desired. Notice that it also entails $\mathbb{P}_{(0,0)}[T = 0] = 1$, as mentioned in the introduction.

Using again self-similarity, this entails $\mathbb{P}_{(x,0)}[T_0 < +\infty] = 1$ for all $x \leq 0$, and also for all $x \geq 0$ in considering the dual process $-L$. The fact that $\mathbb{P}_{(x,y)}[T_0 < +\infty] = 1$ for all x, y such that $xy < 0$ follows then by a comparison of the sample paths.

Suppose now that $x \leq 0, y < 0$. Introduce the stopping time $S = \inf\{t > 0, L_t > 0\}$, which is finite a.s. under $\mathbb{P}_{(x,y)}$ because $|L|$ is not a subordinator. It is clear that $L_S \geq 0$ and $X_S < 0$ a.s. Applying the strong Markov property, we see from the above cases that

$$\mathbb{P}_{(x,y)}[T_0 = +\infty] \leq \mathbb{P}_{(x,y)}[\mathbb{P}_{(X_S, L_S)}[T_0 = +\infty]] = 0.$$

The same argument holds for $x \geq 0, y > 0$. □

Assume now $x < 0$ or $x = 0$ and $y < 0$. It is clear that at T_0 the process X has a non-negative speed, which entails by right-continuity that $L_{T_0} \geq 0$ a.s. Applying the Markov property at T_0 entails

$$\mathbb{P}_{(x,y)}[X_t \in du] = \int_0^\infty \int_0^t \mathbb{P}_{(0,z)}[X_{t-s} \in du] \mathbb{P}_{(x,y)}[T_0 \in ds, L_{T_0} \in dz] \quad (3.1)$$

for all $t, u > 0$. Integrating in time yields then after a change of variable and Fubini's theorem

$$\int_0^\infty \mathbb{P}_{(x,y)}[X_t \in du] dt = \int_0^\infty \left(\int_0^\infty \mathbb{P}_{(0,z)}[X_t \in du] dt \right) \mathbb{P}_{(x,y)}[L_{T_0} \in dz]$$

for all $u > 0$. Integrating in space along $u^{-\nu}$ and applying again Fubini's theorem shows finally the general formula

$$\int_0^\infty \mathbb{E}_{(x,y)}[(X_t^+)^{-\nu}] dt = \int_0^\infty \mathbb{P}_{(x,y)}[L_{T_0} \in dz] \left(\int_0^\infty \mathbb{E}_{(0,z)}[(X_t^+)^{-\nu}] dt \right) \quad (3.2)$$

which is valid for all $\nu \in \mathbb{R}$, with possibly infinite values on both sides.

3.1. Proof of (i). Assume $x = 0$ and $y < 0$. Setting $\nu \in (\alpha/(\alpha + 1), 1)$, a straightforward application of Proposition 4 (i) and (ii) is that both sides of (3.2) are finite, which leads to

$$\mathbb{E}_{(0,y)}[L_{T_0}^{s-1}] = |y|^{s-1} \left(\frac{\sin(\pi\gamma s)}{\sin(\pi(1-\gamma)s)} \right)$$

for all $s \in (0, 1)$. The formula extends then to $\{|s| < 1/(1-\gamma)\}$ by analytic continuation. On the other hand, for all $\mu \in (0, 1)$ and $s \in (-1, 1)$, the formula

$$\int_0^\infty \frac{\sin(\pi\mu)x^s}{\pi\mu(x^2 + 2\cos(\pi\mu)x + 1)} dx = \frac{\sin(\pi\mu s)}{\mu \sin(\pi s)}$$

is a simple and well-known consequence of the residue theorem. Recalling that

$$\chi = \frac{\gamma}{1-\gamma} \in (0, 1)$$

and the definition of \mathbf{C}_μ , we deduce

$$\mathbb{E}_{(0,y)}[L_{T_0}^{s-1}] = |y|^{s-1} \mathbb{E}[\mathbf{C}_\chi^{(1-\gamma)s}]$$

for all $|s| < 1/(1-\gamma)$, which concludes the proof of (i) by Mellin inversion. \square

3.2. Proof of (ii). Assume $x < 0$ and $y = 0$. Another application of (3.2) combined with Proposition 4 (i) and (iii) shows that

$$\mathbb{E}_{(x,0)}[L_{T_0}^{s-1}] = \frac{(1+\alpha)^{\frac{1-s}{1+\alpha}} \Gamma(\frac{\alpha+2}{\alpha+1}) \Gamma(\frac{1-s}{\alpha+1}) \sin(\pi\gamma)}{\Gamma(\frac{s}{\alpha+1}) \Gamma(1-s) \sin(\pi s(1-\gamma))} |x|^{\frac{s-1}{\alpha+1}} \quad (3.3)$$

for all $s \in (0, 1)$. A simple analysis on the Gamma factors shows that the above expression remains finite for all $|s| < 1/(1-\gamma)$. \square

3.3. Some further Mellin inversions. In this paragraph we would like to invert (3.3) for certain values of the parametrization (α, ρ) . Without loss of generality we set $x = -1, y = 0$. Applying the complement formula for the Gamma function, we first deduce from (3.3)

$$\mathbb{E}_{(-1,0)}[L_{T_0}^{s-1}] = (1+\alpha)^{\frac{1-s}{1+\alpha}} \frac{\Gamma(\frac{\alpha+2}{\alpha+1}) \Gamma(\frac{1-s}{\alpha+1}) \Gamma(s(1-\gamma)) \Gamma(1-s(1-\gamma))}{\Gamma(\frac{s}{\alpha+1}) \Gamma(1-s) \Gamma(\gamma) \Gamma(1-\gamma)} \quad (3.4)$$

for $|s| < 1/(1-\gamma)$.

3.3.1. The Cauchy case. We have $\alpha = 1$ and $\rho \in (0, 1)$, whence $\gamma = \rho/2 \in (0, 1/2)$. As mentioned in the introduction, set

$$\delta = \frac{1}{2(1-\gamma)} = \frac{1+\chi}{2} \in (1/2, 1).$$

Applying the Legendre-Gauss multiplication formula transforms (3.4) into

$$\mathbb{E}_{(-1,0)}[L_{T_0}^{s-1}] = 2^{\frac{s-1}{2}} \times \frac{\sin(\pi\gamma) \sin(\pi s/2)}{\sin(\pi s(1-\gamma))}.$$

As above, this entails that under $\mathbb{P}_{(x,0)}$ one has

$$L_{T_0} \stackrel{d}{=} \sqrt{2|x|} (\mathbf{C}_\delta^{1-\gamma})^{(1)},$$

which provides an striking similarity with the law of L_{T_0} under $\mathbb{P}_{(0,y)}$ for $y < 0$. Notice that these two laws are however never the same, because $\delta \neq \chi$.

3.3.2. The Brownian case. We have $\alpha = 2, \rho = \chi = 1/2$ and $\gamma = 1/3$. Applying three times the Legendre-Gauss multiplication formula and simplifying the quotients shows

$$\mathbb{E}_{(-1,0)}[L_{T_0}^{s-1}] = 9^{\frac{s-1}{3}} \times \frac{\Gamma(1/2 + s/3)}{\Gamma(5/6)} \times \frac{\Gamma(1/2 - s/3)\Gamma(1/3)}{\Gamma(2/3 - s/3)\Gamma(1/6)}$$

for all $s \in (0, 1)$. Inverting the Mellin transform, this entails that under $\mathbb{P}_{(x,0)}$ one has

$$L_{T_0} \stackrel{d}{=} |9x|^{1/3} \left(\frac{\mathbf{\Gamma}_{5/6}}{\mathbf{B}_{1/6,1/6}} \right)^{1/3},$$

where $\mathbf{\Gamma}_c$ resp. $\mathbf{B}_{a,b}$ stands for the standard Gamma resp. Beta random variable, and the quotient is assumed independent. Gor'kov [7] provides an expression of the density of L_{T_0} under $\mathbb{P}_{(x,y)}$ in terms of the confluent hypergeometric function - see also formula (3) in [9]. It seems however that the above simple identity in law has passed unnoticed in the literature on integrated Brownian motion.

Remark 7. It is well-known that $\log(\mathbf{\Gamma}_c)$ and $\log(\mathbf{B}_{a,b})$ are infinitely divisible random variables, and this property is hence also shared by $\log(L_{T_0})$ under $\mathbb{P}_{(x,0)}$. The question whether L_{T_0} itself is infinitely divisible is an interesting open problem for Brownian motion.

3.3.3. The case $\alpha < 1$. We have $\rho \in (0, 1), \gamma \in (0, 1/2)$ and $\chi \in (0, 1)$. Set

$$\eta = \frac{1}{(\alpha+1)(1-\gamma)} = \frac{1}{1+\alpha(1-\rho)} \in (1/2, 1) \quad \text{and} \quad \sigma = \frac{\alpha+1}{2} \in (1/2, 1).$$

To express our result, we need some further notation. For every $\mu \in (0, 1)$ set \mathbf{Z}_μ for the standard positive μ -stable random variable [21], which is characterized through its Mellin transformation by

$$\mathbb{E}[\mathbf{Z}_\mu^s] = \frac{\Gamma(1 - \frac{s}{\mu})}{\Gamma(1 - s)}, \quad s < \mu.$$

Applying again the Legendre-Gauss formula entails

$$\begin{aligned} \mathbb{E}_{(-1,0)}[L_{T_0}^{s-1}] &= \kappa \left(\frac{2}{(1+\alpha)^{\frac{1}{1+\alpha}}} \right)^{s-1} \times \frac{\Gamma(1 + s(1-\gamma))}{\Gamma(1 + \frac{s}{\alpha+1})} \times \frac{\Gamma(1 - s(1-\gamma))}{\Gamma(1 - \frac{s}{2})} \times \frac{\Gamma(1 + \frac{1-s}{\alpha+1})}{\Gamma(1 + \frac{1-s}{2})} \\ &= \kappa \left(\frac{2}{(1+\alpha)^{\frac{1}{1+\alpha}}} \right)^{s-1} \mathbb{E}[\mathbf{Z}_\eta^{-\frac{s}{\alpha+1}}] \times \mathbb{E}[\mathbf{Z}_\delta^{\frac{s}{2}}] \times \mathbb{E}[\mathbf{Z}_\sigma^{\frac{s-1}{2}}] \end{aligned}$$

for all $s \in (0, 1)$, where κ is the normalizing constant. Identifying, this shows that under $\mathbb{P}_{(x,0)}$ one has

$$L_{T_0} \stackrel{d}{=} 2 \left| \frac{x}{\alpha + 1} \right|^{\frac{1}{\alpha+1}} \mathbf{Z}_\sigma^{\frac{1}{2}} \times \left(\frac{\mathbf{Z}_\delta^{\frac{1}{2}}}{\mathbf{Z}_\eta^{\frac{1}{\alpha+1}}} \right)^{(1)},$$

where the product and the quotient are assumed independent.

Remark 8. The above argument shows that the function

$$\mathcal{M}_\alpha(s) = \frac{\Gamma(\frac{1}{\alpha+1} + \frac{s-1}{\alpha+1})\Gamma(1 + \frac{s-1}{\alpha+1})}{\Gamma(\frac{1}{\alpha+1})\Gamma(1 + (s-1))}$$

is the Mellin transform of a positive random variable for all $\alpha \leq 1$. This is equivalent to the fact that the independent product $\mathbf{X}_u = \mathbf{\Gamma}_u^u \times \mathbf{\Gamma}_1^u$ is an exponential mixture for all $u = 1/(\alpha + 1) \in [1/2, 1)$. It is easy to see that \mathbf{X}_u is also an exponential mixture for all $u \geq 1$. However, this property is not true in general for $u < 1/2$. Taking $u = 1/3$ viz. $\alpha = 2$, we see indeed from the Legendre-Gauss multiplication formula that

$$\mathcal{M}_2(s) = \left(\frac{4}{27} \right)^{\frac{s-1}{3}} \frac{\Gamma(\frac{2}{3})}{\Gamma(\frac{2}{3} + \frac{s-1}{3})}$$

is log-concave and hence not the Mellin transform of a positive measure.

3.3.4. *The case $1 < \alpha < 2$.* We first separate (3.4) into

$$\begin{aligned} \mathbb{E}_{(-1,0)}[L_{T_0}^{s-1}] &= (1 + \alpha)^{\frac{1-s}{1+\alpha}} \times \frac{\Gamma(\frac{1}{\alpha+1})\Gamma(s(1-\gamma))}{\Gamma(1-\gamma)\Gamma(\frac{s}{\alpha+1})} \times \frac{(\frac{1}{\alpha+1})\Gamma(\frac{1-s}{\alpha+1})\Gamma(1-s(1-\gamma))}{\Gamma(\gamma)\Gamma(1-s)} \\ &= (1 + \alpha)^{\frac{1-s}{1+\alpha}} \times \mathbb{E} \left[\mathbf{Z}_{1/(1+\alpha(1-\rho))}^{-s/(\alpha+1)} \right] \times \frac{(\frac{1}{\alpha+1})\Gamma(\frac{1-s}{\alpha+1})\Gamma(1-s(1-\gamma))}{\Gamma(\gamma)\Gamma(1-s)}. \end{aligned}$$

We next make the following assumption

$$\gamma \leq 1/3.$$

Notice that this assumption is fulfilled in the spectrally negative case, where $\rho = 1 - 1/\alpha$ viz. $\gamma = (\alpha - 1)/(\alpha + 1) < 1/3$. Setting $\mathcal{M}_{\alpha,\gamma}(s)$ for the second multiplicand on the right-hand side, we use again the Legendre-Gauss multiplication formula to get the transformation

$$\begin{aligned} \mathcal{M}_{\alpha,\gamma}(s) &= \kappa 3^s \frac{\Gamma(\frac{1-s}{\alpha+1})}{\Gamma(\frac{1-s}{3})} \times \frac{\Gamma(\frac{1}{2} - \frac{s}{3})}{\Gamma(\frac{2}{3} - \frac{s}{3})} \times \frac{\Gamma(1-s(1-\gamma))}{\Gamma(\frac{1}{2} - \frac{s}{3})\Gamma(1 - \frac{s}{3})} \\ &= \tilde{\kappa} (3 \cdot 2^{-2/3})^s \mathbb{E} \left[\mathbf{Z}_{\frac{3}{\alpha+1}}^{\frac{s-1}{3}} \right] \times \mathbb{E} \left[\mathbf{B}_{1/6,1/6}^{\frac{1-s}{3}} \right] \times \frac{\Gamma(1-s(1-\gamma))}{\Gamma(1 - \frac{2s}{3})} \\ &= \tilde{\kappa} (3 \cdot 2^{-2/3})^s \mathbb{E} \left[\mathbf{Z}_{\frac{3}{\alpha+1}}^{\frac{s-1}{3}} \right] \times \mathbb{E} \left[\mathbf{B}_{1/6,1/6}^{\frac{1-s}{3}} \right] \times \mathbb{E} \left[\mathbf{Z}_{2/3(1-\gamma)}^{2s/3} \right] \end{aligned}$$

where $\kappa, \tilde{\kappa}$ are normalizing constants. This leads to the identity in law

$$L_{T_0} \stackrel{d}{=} 3 \cdot 2^{-2/3} \left| \frac{x}{\alpha + 1} \right|^{\frac{1}{\alpha+1}} \times \left(\frac{\mathbf{Z}_{\frac{\alpha+1}{3}}}{\mathbf{B}_{1/6, 1/6}} \right)^{\frac{1}{3}} \times \left(\frac{\mathbf{Z}_{2/3(1-\gamma)}^{2/3}}{\mathbf{Z}_{1/(1+\alpha(1-\rho))}^{1/(\alpha+1)}} \right)^{(1)},$$

an extension of the Brownian case because when $\alpha = 2$ the first multiplicand is $\mathbf{B}_{1/6, 1/6}^{-1/3}$, whereas the second one reads

$$\left(\mathbf{Z}_{1/2}^{-1/3} \right)^{(1)} \stackrel{d}{=} 2^{2/3} \left(\mathbf{\Gamma}_{1/2}^{1/3} \right)^{(1)} \stackrel{d}{=} 2^{2/3} \mathbf{\Gamma}_{5/6}^{1/3}.$$

Remark 9. We do not know whether $\mathcal{M}_{\alpha, \gamma}(s)$ is still the Mellin transform of a positive random variable in the remaining case $1/3 < \gamma \leq 1/(\alpha + 1)$. This would be a consequence of the exponential mixture property of the independent product $\mathbf{\Gamma}_u^u \times \mathbf{\Gamma}_1^{1-u}$ for all $u \in (1/3, 1/2)$. It is easily shown that the latter property holds for all $u = 1/n, n \geq 2$, and we think that it does for all $u \in [0, 1/2]$.

4. PROOF OF THEOREM A

We first reduce the problem to the situation where the bivariate process (X, L) starts from a coordinate axis.

Lemma 10. *Assume $x < 0$. For all $y \in \mathbb{R}$ one has*

$$\mathbb{P}_{(x, y)}[T_0 > t] \asymp \mathbb{P}_{(x, 0)}[T_0 > t], \quad t \rightarrow +\infty.$$

Proof. Fix $t > 1$ and suppose first $y > 0$. One has $\mathbb{P}_{(x, y)}[T_0 > t] \leq \mathbb{P}_{(x, 0)}[T_0 > t]$ by a direct comparison of the sample paths. On the other hand,

$$\begin{aligned} \mathbb{P}_{(x, y)}[T_0 > t] &\geq \mathbb{P}_{(x, y)}[X_1 < x, L_1 < 0, T_0 > t] \\ &= \mathbb{P}_{(x, y)}[X_1 < x, X_1^* < 0, L_1 < 0, \mathbb{P}_{(X_1, L_1)}[T_0 > t - 1]] \\ &\geq \mathbb{P}_{(x, y)}[X_1 < x, X_1^* < 0, L_1 < 0] \times \mathbb{P}_{(x, 0)}[T_0 > t - 1] \geq c \mathbb{P}_{(x, 0)}[T_0 > t] \end{aligned}$$

for some $c > 0$, where the equality follows from the Markov property, the second inequality from a comparison of the sample paths, and the third inequality from a support theorem in uniform norm for the Lévy stable process L . More precisely, the latter process can be approximated by any continuous function in uniform norm, because the support of its Lévy measure is \mathbb{R} - see Corollary 1 in [17].

Fix again $t > 1$ and suppose now $y < 0$. Then $\mathbb{P}_{(x, y)}[T_0 > t] \geq \mathbb{P}_{(x, 0)}[T_0 > t]$, and similarly as above one has

$$\begin{aligned} \mathbb{P}_{(x, 0)}[T_0 > t] &\geq \mathbb{P}_{(x, 0)}[X_1 < x, X_1^* < 0, L_1 < y, \mathbb{P}_{(X_1, L_1)}[T_0 > t - 1]] \\ &\geq \mathbb{P}_{(x, 0)}[X_1 < x, X_1^* < 0, L_1 < y] \times \mathbb{P}_{(x, y)}[T_0 > t - 1] \geq c \mathbb{P}_{(x, y)}[T_0 > t] \end{aligned}$$

for some $c > 0$. This completes the proof. \square

In the remainder of this section, we will implicitly assume, without loss of generality, that

$$\{x = 0, y < 0\} \quad \text{or} \quad \{x < 0, y = 0\}.$$

We start by studying the asymptotics at infinity of the density function of L_{T_0} under $\mathbb{P}_{(x,y)}$, which we denote by $f_{x,y}^0$.

Lemma 11. *There exists $c > 0$ such that*

$$f_{x,y}^0(z) \sim cz^{-1/(1-\gamma)}, \quad z \rightarrow +\infty.$$

Proof. If $x = 0$, the asymptotic is a direct consequence of the explicit expression of $f_{(0,y)}^0$ which is given in Theorem B (i). If $y = 0$, Theorem B (ii) shows that the first positive pole of the Mellin transform of L_{T_0} under $\mathbb{P}_{(x,0)}$ is at $1/(1-\gamma)$, and is simple. The required asymptotic for $f_{(x,0)}^0$ is then a consequence of a converse mapping theorem for Mellin transforms - see e.g. Theorem 6.4 in [8]. □

Remark 12. (a) The converse mapping theorem for Mellin transforms yields also an explicit expression for the constant c , but we shall not need this information in the sequel.

(b) We believe that the above asymptotic remains true for $x < 0$ and all $y \neq 0$. However, the Mellin transform of L_{T_0} under $\mathbb{P}_{(x,y)}$ is then expressed with the help of a double integral which is absolutely divergent, and whose singularities are difficult to study at first sight.

(c) The lemma entails by integration that

$$\mathbb{P}_{(x,y)}[L_{T_0} > z] \sim c\chi^{-1} z^{-\chi}, \quad z \rightarrow +\infty.$$

Heuristically, it is tempting to write by scaling $L_{T_0} = T_0^{1/\alpha} |L_1|$ and since $\mathbb{P}_{(x,y)}[|L_1| > z] \sim cz^{-\alpha} \ll z^{-\chi}$ at infinity, we may infer that

$$\mathbb{P}_{(x,y)}[T_0 > t] \asymp t^{-\chi/\alpha} = t^{-\theta}, \quad t \rightarrow +\infty.$$

This explains the equivalence between finite moments stated in the introduction. We will prove in the remainder of this section that this heuristic is actually correct.

The following lemma provides our key-estimate.

Lemma 13. *For all $\nu \in (\alpha(1-\theta)/(\alpha+1), 1)$ there exists $c > 0$ such that*

$$\mathbb{E}_{(x,y)} \left[\int_0^t \mathbf{1}_{\{T_0 > t-u\}} \mathbb{E}_{(0,L_{T_0})} [(X_u^+)^{-\nu}] du \right] \sim ct^{1-(1+1/\alpha)\nu-\theta}, \quad t \rightarrow +\infty.$$

Proof. We first assume $\nu \in (\alpha/(\alpha+1), 1)$ and transform the expression on the left-hand side. From (3.1), Fubini's theorem, and the Markov property, we obtain

$$\int_0^\infty e^{-\lambda t} \mathbb{E}_{(x,y)} [(X_t^+)^{-\nu}] dt = \mathbb{E}_{(x,y)} \left[e^{-\lambda T_0} \int_0^\infty e^{-\lambda t} \mathbb{E}_{(0,L_{T_0})} [(X_t^+)^{-\nu}] dt \right], \quad \lambda \geq 0,$$

both sides being finite because of Proposition 4. Integrating by parts shows then, with the help of (3.2) and Proposition 4, that

$$\begin{aligned} \lambda \int_0^\infty e^{-\lambda t} \int_t^\infty & \left(\mathbb{E}_{(x,y)} [(X_u^+)^{-\nu}] - \mathbb{E}_{(x,y)} \left[\mathbb{E}_{(0,L_{T_0})} [(X_u^+)^{-\nu}] \right] \right) du dt \\ &= \mathbb{E}_{(x,y)} \left[(1 - e^{-\lambda T_0}) \int_0^\infty e^{-\lambda t} \mathbb{E}_{(0,L_{T_0})} [(X_t^+)^{-\nu}] dt \right] \\ &= \mathbb{E}_{(x,y)} \left[\int_0^\infty \lambda e^{-\lambda t} \left(\int_0^t \mathbf{1}_{\{T_0 > t-u\}} \mathbb{E}_{(0,L_{T_0})} [(X_t^+)^{-\nu}] du \right) dt \right]. \end{aligned}$$

Inverting the Laplace transforms shows that

$$\mathbb{E}_{(x,y)} \left[\int_0^t \mathbf{1}_{\{T_0 > t-u\}} \mathbb{E}_{(0,L_{T_0})} [(X_u^+)^{-\nu}] du \right] = H_{(x,y)}(t), \quad (4.1)$$

with the notation

$$H_{(x,y)}(t) = \int_t^{+\infty} \left(\mathbb{E}_{(x,y)} [(X_u^+)^{-\nu}] - \mathbb{E}_{(x,y)} \left[\mathbb{E}_{(0,L_{T_0})} [(X_u^+)^{-\nu}] \right] \right) du, \quad t > 0. \quad (4.2)$$

It remains therefore to compute the asymptotics of the function $H_{(x,y)}$, which only depends on the law of L_{T_0} under $\mathbb{P}_{(x,y)}$. From Proposition 3, the additive property of sine, and a change of variable, we get

$$\begin{aligned} \mathbb{E}_{(x,y)} [(X_u^+)^{-\nu}] - \mathbb{E}_{(x,y)} \left[\mathbb{E}_{(0,L_{T_0})} [(X_u^+)^{-\nu}] \right] &= \frac{2\Gamma(1-\nu)}{\pi} \int_0^\infty \lambda^{\nu-1} e^{-c_{\alpha,\rho} \lambda^\alpha u^{\alpha+1}} \Phi_u(\lambda u^{1+1/\alpha}) d\lambda \\ &= \frac{2\Gamma(1-\nu)}{\pi} \frac{1}{u^{(1+1/\alpha)\nu}} \int_0^\infty \xi^{\nu-1} e^{-c_{\alpha,\rho} \xi^\alpha} \Phi_u(\xi) d\xi \end{aligned} \quad (4.3)$$

where the function Φ_u is defined by

$$\begin{aligned} \Phi_u(\xi) &= \mathbb{E}_{(x,y)} \left[\cos \left(\frac{\xi x}{2u^{1+1/\alpha}} + s_{\alpha,\rho} \xi^\alpha + \frac{\nu\pi}{2} + \frac{\xi(y + L_{T_0})}{2u^{1/\alpha}} \right) \sin \left(\frac{\xi x}{2u^{1+1/\alpha}} + \frac{\xi(y - L_{T_0})}{2u^{1/\alpha}} \right) \right] \\ &= \int_0^\infty \cos \left(\frac{\xi x}{2u^{1+1/\alpha}} + s_{\alpha,\rho} \xi^\alpha + \frac{\nu\pi}{2} + \frac{\xi(y + z)}{2u^{1/\alpha}} \right) \sin \left(\frac{\xi x}{2u^{1+1/\alpha}} + \frac{\xi(y - z)}{2u^{1/\alpha}} \right) f_{x,y}^0(z) dz. \end{aligned}$$

Setting $F_u(\xi, z)$ for the trigonometric function inside the integral, a change of variable entails

$$u^\theta \left(\int_0^\infty F_u(\xi, z) f_{x,y}^0(z) dz \right) = \int_0^\infty F_u(\xi, ru^{1/\alpha}) u^{\frac{1}{\alpha(1-\gamma)}} f_{x,y}^0(ru^{1/\alpha}) dr. \quad (4.4)$$

A further application of the converse mapping theorem and of Theorem B shows that $f_{x,y}^0(z) \sim cz^{1/(1-\gamma)}$ as $z \rightarrow 0+$, for some $c > 0$. This estimate and Lemma 11 yield the uniform bound

$$f_{x,y}^0(z) \leq Kz^{-1/(1-\gamma)}, \quad z > 0$$

for some $K > 0$. Hence, for u large enough, the integrated function on the right-hand side of (4.4) is dominated by $K(r \wedge 1) r^{-(\chi+1)}$ for some $K > 0$, which is an integrable function

because $\chi \in (0, 1)$. Applying the dominated convergence theorem in (4.4) and using Lemma 11 shows that there exists $c > 0$ such that for all $\xi > 0$

$$\begin{aligned} u^\theta \Phi_u(\xi) &\rightarrow -c \int_0^\infty \cos(s_{\alpha,\rho} \xi^\alpha + (\nu\pi + \xi r)/2) \sin(\xi r/2) r^{-(1+\chi)} dr \\ &= -(c/2) \Gamma(-\chi) \sin(s_{\alpha,\rho} \xi^\alpha + (\nu - \chi)\pi/2) \xi^\chi \end{aligned}$$

as $u \rightarrow \infty$, where the equality comes from (2.2), an integration by parts and some trigonometry. Plugging back this expression in (4.3), we deduce by dominated convergence that

$$\begin{aligned} \mathbb{E}_{(x,y)} [(X_u^+)^{-\nu}] - \mathbb{E}_{(x,y)} [\mathbb{E}_{(0,L_{T_0})} [(X_u^+)^{-\nu}]] \\ \underset{u \rightarrow +\infty}{\sim} c \left(\int_0^\infty \xi^{\chi+\nu-1} e^{-c_{\alpha,\rho} \xi^\alpha} \sin(s_{\alpha,\rho} \xi^\alpha + (\nu - \chi)\pi/2) d\xi \right) u^{-((1+1/\alpha)\nu+\theta)} \quad (4.5) \end{aligned}$$

for some $c > 0$. A last computation shows that the integral on the right-hand side equals

$$(\alpha + 1)^{\frac{\chi+\nu}{\alpha}} \Gamma\left(\frac{\chi + \nu}{\alpha}\right) \sin(\pi(\rho\nu + (\rho - 1)\chi))$$

and is positive because $\rho\nu + (\rho - 1)\chi \in (0, 1)$ for all $\nu \in (\alpha(1 - \theta)/(\alpha + 1), 1)$ as can be readily checked. The final result follows then by integration, and the proof is complete for $\nu \in (\alpha/(\alpha + 1), 1)$.

Suppose now $\nu \in (\alpha(1 - \theta)/(\alpha + 1), \alpha/(\alpha + 1))$. The left-hand side of (4.1) is well-defined and the estimate (4.5), which does not require the lower bound $\nu > \alpha/(\alpha + 1)$, together with the positivity of the constant entails that the integral in (4.2) is absolutely convergent, because $(1 + 1/\alpha)\nu + \theta > 1$. By analytic continuation this shows that (4.1) remains valid for $\nu \in (\alpha(1 - \theta)/(\alpha + 1), \alpha/(\alpha + 1))$, and the estimate (4.5) holds as well. This completes the proof, again by integration of (4.5). \square

4.1. Proof of the upper bound. Fix $A > 0$ and $\nu \in (\alpha/(\alpha + 1), 1)$. By continuity and positivity there exists $\varepsilon > 0$ such that for all $z \in [0, A]$,

$$\int_0^1 \mathbb{E}_{(0,z)} [(X_u^+)^{-\nu}] du \geq \varepsilon.$$

For all $t > 0$, we get from (4.1), a change of variable and the self-similarity

$$\begin{aligned} t^{(1+1/\alpha)\nu+\theta-1} H_{(x,y)}(t) &\geq t^{(1+1/\alpha)\nu+\theta-1} \mathbb{E}_{(x,y)} \left[\mathbf{1}_{\{T_0 > t\}} \int_0^t \mathbb{E}_{(0,L_{T_0})} [(X_u^+)^{-\nu}] du \right] \\ &= t^\theta \mathbb{E}_{(x,y)} \left[\mathbf{1}_{\{T_0 > t\}} \int_0^1 \mathbb{E}_{(0, \frac{1}{t^{1/\alpha}} L_{T_0})} [(X_u^+)^{-\nu}] du \right] \\ &\geq \varepsilon t^\theta \mathbb{P}_{(x,y)} [T_0 > t, L_{T_0} \leq At^{1/\alpha}] \\ &\geq \varepsilon t^\theta (\mathbb{P}_{(x,y)} [T_0 > t] - \mathbb{P}_{(x,y)} [T_0 > t, L_{T_0} \geq At^{1/\alpha}]) \\ &\geq \varepsilon t^\theta (\mathbb{P}_{(x,y)} [T_0 > t] - \mathbb{P}_{(x,y)} [T_0 > t]^{1-1/p} \mathbb{P}_{(x,y)} [L_{T_0} \geq At^{1/\alpha}]^{1/p}), \end{aligned}$$

where the last inequality follows from Hölder's inequality and is valid for all $p > 1$. We now take $t > 1$ and

$$\frac{1}{p} = 1 - \frac{1}{\ln(t)}.$$

On the one hand, Lemma 11 entails

$$\begin{aligned} \limsup_{t \rightarrow +\infty} t^\theta \mathbb{P}_{(x,y)}[T_0 > t]^{\frac{1}{\ln(t)}} \mathbb{P}_{(x,y)}[L_{T_0} \geq At^{1/\alpha}]^{1 - \frac{1}{\ln(t)}} \\ \leq \limsup_{t \rightarrow +\infty} t^\theta \mathbb{P}_{(x,y)}[L_{T_0} \geq At^{1/\alpha}]^{1 - \frac{1}{\ln(t)}} = K < +\infty. \end{aligned}$$

On the other hand, Lemma 13 shows that

$$t^{(1+1/\alpha)\nu + \theta - 1} H_{(x,y)}(t) \rightarrow c > 0 \quad \text{as } t \rightarrow +\infty.$$

Putting everything together entails

$$t^\theta \mathbb{P}_{(x,y)}[T_0 > t] \leq \tilde{K}$$

for some finite \tilde{K} as soon as t is large enough. □

4.2. Proof of the lower bound. We start with the following lemma :

Lemma 14. *One has*

$$\int_0^t \mathbb{P}_{(x,y)}[T_0 > u] du \asymp t^{1-\theta} \quad \text{as } t \rightarrow +\infty.$$

Proof. Firstly, integrating the above upper bound for $\mathbb{P}_{(x,y)}[T_0 > t]$ entails the existence of a finite κ_2 such that

$$\int_0^t \mathbb{P}_{(x,y)}[T_0 > u] du \leq \kappa_2 t^{1-\theta} \quad \text{as } t \rightarrow +\infty.$$

To prove the lower inequality, we fix $\nu \in (\alpha(1-\theta)/(1+\alpha), \alpha/(1+\alpha))$ and deduce from Proposition 3 the uniform bound

$$\mathbb{E}_{(0,y)}[(X_u^+)^{-\nu}] \leq \frac{\Gamma(1-\nu)}{\pi} \int_0^\infty \lambda^{\nu-1} e^{-c_{\alpha,\rho} \lambda^{\alpha} u^{\alpha+1}} d\lambda \leq K u^{-\nu(1+1/\alpha)}, \quad u > 0, \quad (4.6)$$

for some finite constant K . Set $\eta = \nu(1+1/\alpha) \in (0, 1)$ and fix $\varepsilon \in (0, 1)$. Using (4.1) and (4.6) we decompose

$$\begin{aligned} t^{\eta+\theta-1} H_{(x,y)}(t) &\leq K t^{\eta+\theta-1} \left(\int_0^{t(1-\varepsilon)} \frac{\mathbb{P}_{(x,y)}[T_0 > u]}{(t-u)^\eta} du + \int_{t(1-\varepsilon)}^t \frac{\mathbb{P}_{(x,y)}[T_0 > u]}{(t-u)^\eta} du \right) \\ &\leq K \varepsilon^{-\eta} t^{\theta-1} \int_0^t \mathbb{P}_{(x,y)}[T_0 > u] du + \frac{K t^\theta \varepsilon^{1-\eta}}{1-\eta} \mathbb{P}_{(x,y)}[T_0 > t(1-\varepsilon)] \\ &\leq \tilde{K} \varepsilon^{-\eta} \left(t^{\theta-1} \int_0^t \mathbb{P}_{(x,y)}[T_0 > u] du + \varepsilon \right) \end{aligned}$$

for some finite \tilde{K} , where the third inequality follows from the upper bound. Applying Lemma 13 and taking ε small enough shows finally that there exists $\kappa_1 > 0$ such that

$$\int_0^t \mathbb{P}_{(x,y)}[T_0 > u] du \geq \kappa_1 t^{1-\theta} \quad \text{as } t \rightarrow +\infty.$$

□

We can now finish the proof. Fixing $A > 0$ and applying the mean value theorem entails

$$A t^\theta \mathbb{P}_{(x,y)}[T_0 > t] \geq t^{\theta-1} \int_t^{t+tA} \mathbb{P}_{(x,y)}[T_0 > u] du \geq \kappa_1 (1+A)^{1-\theta} - \kappa_2$$

as $t \rightarrow +\infty$, for some constants $0 < \kappa_1 < \kappa_2 < \infty$ given by Lemma 14. Since $\theta < 1$, the lower bound follows in choosing A large enough.

□

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