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# Estimates $L^{r}-L^{s}$ for solutions of the $\bar{\partial}$ equation in strictly pseudo convex domains in $\mathbb{C}^{n}$. 

Eric Amar


#### Abstract

We prove some new estimates for solutions of the $\bar{\partial} u=\omega$ equation in a strictly pseudo convex domain $\Omega$ in $\mathbb{C}^{n}$. For instance if the $(p, q)$ current $\omega$ has its coefficients in $L^{r}(\Omega)$ with $1 \leq r<2(n+1)$ then there is a solution $u$ in $L^{s}(\Omega)$ with $\frac{1}{s}=\frac{1}{r}-\frac{1}{2(n+1)}$. These results are based on Carleson measures of order $\alpha$ introduced and studied in [4] and on the subordination lemma [5].


## 1 Introduction.

Let $\Omega$ be a bounded strictly pseudo convex domain with smooth $\mathcal{C}^{\infty}$ boundary. We shall denote these domains as s.p.c. domains in the sequel.

Ovrelid [10] proved that if we have a $(p, q)$ current $\omega, \bar{\partial}$ closed in $\Omega$ and such that its coefficients are in $L^{r}(\Omega)$ then there is a $(p, q-1)$ current $u$ solution of the equation $\bar{\partial} u=\omega$ and with coefficients still in $L^{r}(\Omega)$. Let us define a norm on these currents :

$$
\omega \in L_{(p, q)}^{r}(\Omega), \omega=\sum_{|I|=p,|J|=q} \omega_{I, J} d z^{I} \wedge d \bar{z}^{J} \Rightarrow\|\omega\|_{r}^{r}:=\sum_{|I|=p,|J|=q}\left\|\omega_{I, J}\right\|_{r}^{r} .
$$

Then Ovrelid proved that $\|u\|_{r}<C\|\omega\|_{r}$, where the constant $C$ does not depend on $\omega$.
In the case of $r=\infty$, this was done before by Lieb [9] and Romanov and Henkin [11] proved that still for $r=\infty$, there is a solution $u$ in the space Lipschitz $1 / 2$. In the book of Henkin and Leiterer [8] we can find precise references for these topics.

The aim of this work is to improve the $L^{r}$ estimates of Ovrelid in the direction of Romanov and Henkin but with a different approach.

We already got this kind of results in [2] by use of Skoda's kernels [12] but we where dealing with boundary values instead of inside ones. Nevertheless using Skoda results we shall prove the following theorem, where $A \lesssim B$ means that there is a constant $C>0$ independent of $A$ and $B$ such that $A \leq C B$.

Theorem 1.1 Let $\Omega$ be a s.p.c. domain in $\mathbb{C}^{n}$ then for $1<r<2 n+2$ we have $\forall \omega \in L_{(p, q)}^{r}(\Omega), \bar{\partial} \omega=0, \exists u \in L_{(p, q-1)}^{s}(\Omega):: \bar{\partial} u=\omega,\|u\|_{L^{s}(\Omega)} \lesssim\|\omega\|_{L^{r}(\Omega)}$, for any $s$ such that $\frac{1}{s}>\frac{1}{r}-\frac{1}{2(n+1)}$.

We shall also show :
Theorem 1.2 Let $\Omega$ be a s.p.c. domain in $\mathbb{C}^{n}$ then for $1<r<2 n+2$ we have

- $\forall \omega \in L_{(p, q)}^{r}(\Omega), \bar{\partial} \omega=0, \exists u \in L_{(p, q-1)}^{s}(\Omega):: \bar{\partial} u=\omega,\|u\|_{L^{s}(\Omega)} \lesssim\|\omega\|_{L^{r}(\Omega)}$,
with $\frac{1}{s}=\frac{1}{r}-\frac{1}{2(n+1)}$.
- For $r=2 n+2$ we have

$$
\exists u \in \operatorname{BMO}_{(p, q)}(\Omega):: \bar{\partial} u=\omega,\|u\|_{B M O(\Omega)} \lesssim\|\omega\|_{L^{2 n+2}(\Omega)}
$$

If $\omega$ is a $(p, 1)$ form we have also :

- for $r=1$,

$$
\exists u \in L_{(r, 0)}^{s, \infty}(\Omega):: \bar{\partial} u=\omega,\|u\|_{L^{s, \infty}(\Omega)} \lesssim\|\omega\|_{L^{1}(\Omega)}
$$

with $\frac{1}{s}=1-\frac{1}{2(n+1)}$.

- for $r>2 n+2$,

$$
\exists u \in \Gamma_{(p, 0)}^{\beta}(\Omega):: \bar{\partial} u=\omega,\|u\|_{\Gamma^{\beta}(\Omega)} \lesssim\|\omega\|_{L^{r}(\Omega)}
$$

where $\beta=1-\frac{2 n+2}{r}$ and $\Gamma^{\beta}$ is an anisotropic Lipschitz class of functions.
Moreover the solution $u$ is linear on the data $\omega$.
The classes $B M O(\Omega)$ and $\Gamma^{\beta}(\Omega)$ will be defined later. The space $L_{(p, 0)}^{s, \infty}(\Omega)$ is the Lorentz space [7].
This theorem is stronger than theorem 1.1 because here, in the case $1 \leq r<2(n+1)$ we get the result for the end point $s$ such that $\frac{1}{s}=\frac{1}{r}-\frac{1}{2(n+1)}$.
Of course if $u \in L_{(p, q-1)}^{s}(\Omega)$ for $s>r$ then $u \in L_{(p, q-1)}^{r}(\Omega)$ hence we also have an strong improvement to Ovrelid's theorem.

Because the class Lipschitz $1 / 2$ is contained in $\Gamma^{1}(\Omega)$ we see that we recover the RomanovHenkin result when $r=\infty$ in the case of $(p, 1)$ forms.

Even if they do not appear in the statement, the Carleson measures of order $\alpha$, A. Bonami and I introduced in [4], are at the heart of the proof.

## 2 Proof of the first theorem.

Let $\Omega$ be a s.p.c. in $\mathbb{C}^{n}$, defined by the function $\rho \in \mathcal{C}^{\infty}\left(\mathbb{C}^{n}\right)$, i.e. $\Omega:=\left\{z \in \mathbb{C}^{n}:: \rho(z)<0\right\}$ and $\forall z \in \partial \Omega, \partial \rho(z) \neq 0$.

Let $\Omega^{\prime}:=\left\{(z, w) \in \mathbb{C}^{n} \times \mathbb{C}:: \rho^{\prime}(z, w):=\rho(z)+|w|^{2}<0\right\}$ and lift a current $\omega$ to $\Omega^{\prime}$ this way: $\omega^{\prime}(z, w):=\omega(z)$.

Lemma 2.1 Let $\Omega$ be a s.p.c. domain in $\mathbb{C}^{n}$, with the above notations we have $\omega \in L_{(p, q)}^{r}(\Omega) \Rightarrow \omega^{\prime}(z, w) \in L_{(p, q)}^{r}\left(\partial \Omega^{\prime}\right)$.

Proof.
This is an instance of the subordination principle [1], [5]. Let $f(z) \in L^{r}(\Omega)$ and set $f^{\prime}(z, w):=f(z)$ in $\Omega^{\prime}$, then, by the main lemma in [5], p. 6,

$$
\left\|f^{\prime}\right\|_{L^{r}\left(\partial \Omega^{\prime}\right)}^{r}:=\int_{\partial \Omega^{\prime}}\left|f^{\prime}(z, w)\right|^{r} d \sigma(z, w)=\int_{\Omega}|f(z)|^{r} \sqrt{-\rho(z)+\frac{|\operatorname{grad} \rho(z)|^{2}}{4}}\left\{\int_{|w|^{2}=-\rho(z)} d|w|\right\} d m(z)
$$

where $d|w|$ is the normalized Lebesgue measure [5] on the circle $|w|^{2}=-\rho(z)$. Because $\bar{\Omega}$ is compact, we have $\forall z \in \bar{\Omega}, \sqrt{-\rho(z)+\frac{|\operatorname{grad} \rho(z)|^{2}}{4}} \leq C(\rho)<\infty$ hence we have

$$
\left\|f^{\prime}\right\|_{L^{r}\left(\partial \Omega^{\prime}\right)}^{r} \leq C(\rho) \int_{\Omega}|f(z)|^{r} d m(z)=C(\rho)\|f\|_{L^{r}(\Omega)}
$$

It remains to apply this taking for $f$ any coefficient of $\omega$.

## Proof of theorem 1.1.

Since $\Omega$ is a s.p.c. domain so is $\Omega^{\prime}$ by the subordination lemma [5]. By use of lemma 2.1 we have that $\omega^{\prime} \in L_{(p, q)}^{r}\left(\partial \Omega^{\prime}\right)$ and still $\bar{\partial} \omega^{\prime}=0$, hence we can apply Skoda's theorem 2 in [12] to get that there is a solution $u^{\prime}$ of $\bar{\partial}_{b} u^{\prime}=\omega^{\prime}$ such that

$$
u^{\prime} \in L_{(p, q-1)}^{s}\left(\partial \Omega^{\prime}\right) \text { with } \frac{1}{s}>\frac{1}{r}-\frac{1}{2(n+1)}
$$

We have

$$
u^{\prime}(z, w)=\sum_{I, J} a_{I, J}^{\prime}(z, w) d z^{I} \wedge d \bar{z}^{J}
$$

Because $\omega^{\prime}$ does not depend on $w$ we have that the coefficients of $u^{\prime}$ are holomorphic in $w$, hence we can set (recall that $u^{\prime}$ is defined on $\partial \Omega^{\prime}$ )

$$
\forall z \in \Omega, a_{I, J}(z):=\int_{|w|^{2}=-\rho(z)} a_{I, J}^{\prime}(z, w) d|w|
$$

and

$$
u(z):=\sum_{I, J} a_{I, J}(z) d z^{I} \wedge d \bar{z}^{J}
$$

then exactly as in [3] we still have

$$
\bar{\partial} u=\omega \text { in } \Omega
$$

Moreover the subordination lemma [5] gives again $u \in L_{(p, q-1)}^{s}(\Omega)$, because $u^{\prime} \in L_{(p, q-1)}^{s}\left(\partial \Omega^{\prime}\right)$.

## 3 Carleson measures of order $\alpha$.

For $\Omega$ a s.p.c. domain in $\mathbb{C}^{n}$, let $V^{0}(\Omega)$ be the space of bounded measures in $\Omega$, and $V^{1}(\Omega)$ the space of Carleson measures in $\Omega$ as defined for instance in [4]. We know that these spaces form a interpolating scale for the real method [4], and we set

$$
V^{\alpha}(\Omega):=\left(V^{0}, V^{1}\right)_{(\alpha, \infty)} ; W^{\alpha}(\Omega):=\left(V^{0}, V^{1}\right)_{(\alpha, p)} \text { with } p=\frac{1}{1-\alpha}
$$

Recall that a $(p, q)$ form $\omega$ is in $W_{(p, q)}^{\alpha}(\Omega)\left(\operatorname{resp} . V_{(p, q)}^{\alpha}(\Omega)\right)$ if its coefficients and the coefficients of $\frac{\omega \wedge \bar{\partial} \rho}{\sqrt{-\rho}}$ are measures in $W^{\alpha}(\Omega)$ (resp. $\left.V^{\alpha}(\Omega)\right)$ see [4] and [6]. A $(p, q)$ form is in $L_{(p, q)}^{r}(\Omega)$ if just its coefficients are in $L^{r}(\Omega)$.

Let $\Omega^{\prime}:=\left\{(z, w) \in \mathbb{C}^{n} \times \mathbb{C}:: \rho^{\prime}(z, w):=\rho(z)+|w|^{2}<0\right\}$ and lift a current $\omega$ to $\Omega^{\prime}$ as before : $\omega^{\prime}(z, w):=\omega(z)$.
Our first result links $L^{r}$ estimates to Carleson $\alpha$ ones.
Theorem 3.1 Let $\Omega$ be a s.p.c. domain in $\mathbb{C}^{n}$ then we have

$$
\omega \in L_{(p, q)}^{r}(\Omega) \Rightarrow \omega^{\prime}(z, w):=\omega(z) \in W_{(p, q)}^{\alpha}\left(\Omega^{\prime}\right)
$$

with $\alpha=\frac{1}{r^{\prime}}+\frac{1}{2(n+1)}$.
Proof.
Let $U^{\prime}:=\bigcup_{j=1}^{N} Q^{\prime}\left(\zeta_{j}^{\prime}, h_{j}\right) \cap \partial \Omega^{\prime}$ be an open set in $\partial \Omega^{\prime}$ and $T\left(U^{\prime}\right)=\bigcup_{j=1}^{N} Q^{\prime}\left(\zeta_{j}^{\prime}, h_{j}\right)$ be its associated "tent" set inside [4] ; in order to see that a measure $d \mu=f d m$, with $m$ the Lebesgue measure in $\mathbb{C}^{n}$, belongs to $V^{\alpha}\left(\Omega^{\prime}\right)$ we have to show, see [4],

$$
\int_{T\left(U^{\prime}\right)}\left|f\left(z^{\prime}\right)\right| d m\left(z^{\prime}\right) \leq C\left|U^{\prime}\right|^{\alpha}
$$

where $\left|U^{\prime}\right|:=\sigma\left(U^{\prime}\right)$ is the Lebesgue measure of $U^{\prime}$ on $\partial \Omega$, and with a constant $C$ independent of $U^{\prime}$.
Because we are dealing with $(p, q)$ currents here, this means that we have to estimate

$$
A:=\int_{T\left(U^{\prime}\right)} \frac{|\omega(z)|}{\sqrt{-\rho^{\prime}(z, w)}} d m(z, w)
$$

with $\rho^{\prime}(z, w):=\rho(z)+|w|^{2}$ is equivalent to the distance of $(z, w) \in \Omega^{\prime}$ to the boundary $\partial \Omega^{\prime}$.
Back to $A$,

$$
A:=\int_{T\left(U^{\prime}\right)} \frac{|\omega(z)|}{\sqrt{-\rho^{\prime}(z, w)}} d m(z, w) \leq \sum_{j=1}^{N} \int_{Q_{j}^{\prime}} \frac{|\omega(z)|}{\sqrt{-\rho^{\prime}(z, w)}} d m(z, w)
$$

The Carleson window $Q_{j}^{\prime}$ is equivalent to the product $\left(Q_{j}^{\prime} \cap \partial \Omega^{\prime}\right) \times\left[h_{j}\right]_{\nu_{j}}$ with $\left[h_{j}\right]_{\nu_{j}}$ the real segment of length $h_{j}$ supported by the real normal $\nu_{j}$ to $\partial \Omega^{\prime}$ at $\zeta_{j}^{\prime}$. Set $h:=\max _{j=1, \ldots, N} h_{j}$, we shall replace $Q_{j}^{\prime}$ by $Q_{j}^{\prime \prime}:=\left(Q_{j}^{\prime} \cap \partial \Omega^{\prime}\right) \times[h]_{\nu_{j}}$.

So we have

$$
A \leq \sum_{j=1}^{N} \int_{Q_{j}^{\prime}} \frac{|\omega(z)|}{\sqrt{-\rho^{\prime}(z, w)}} d m(z, w) \leq \sum_{j=1}^{N} \int_{Q_{j}^{\prime \prime}} \frac{|\omega(z)|}{\sqrt{-\rho^{\prime}(z, w)}} d m(z, w),
$$

where now all the depths have the same value $h$. Hence by Fubini we have

$$
A \leq \int_{0}^{h} \frac{1}{\sqrt{t}}\left\{\int_{U_{t}^{\prime}}|\omega(z)| d \sigma(z, w)\right\} d t
$$

with $U_{t}^{\prime}:=\bigcup_{j=1}^{N} Q_{j}^{\prime \prime} \cap \partial \Omega^{\prime}{ }_{t}$ and $\partial \Omega_{t}^{\prime}:=\left\{(z, w) \in \Omega^{\prime}:: \rho(z)+|w|^{2}=-t\right\}$.
We can estimate the inner integral by Hölder

$$
\begin{equation*}
\int_{U_{t}^{\prime}}|\omega(z)| d \sigma(z, w) \leq\left(\int_{U_{t}^{\prime}}|\omega(z)|^{r} d \sigma(z, w)\right)^{1 / r}\left(\int_{U_{t}^{\prime}} d \sigma(z, w)\right)^{1 / r^{\prime}} \tag{3.1}
\end{equation*}
$$

but

$$
\int_{U_{t}^{\prime}}|\omega(z)|^{r} d \sigma(z, w) \leq \int_{\partial \Omega_{t}}|\omega(z)|^{r} d \sigma(z, w) \leq C(\rho) \int_{\Omega_{t}}|\omega(z)|^{r}\left\{\int_{|w|^{2}=-\rho(z)-t} d|w|\right\} d m(z)
$$

where $d|w|$ is the normalized Lebesgue measure on the circle $|w|^{2}=-\rho(z)-t$. Hence, with $\Omega_{t}:=\{z \in \Omega:: \rho(z)<-t\}$,

$$
\int_{U_{t}^{\prime}}|\omega(z)|^{r} d \sigma(z, w) \leq C(\rho) \int_{\Omega_{t}}|\omega(z)|^{r} d m(z)=C(\rho)\|\omega\|_{L^{r}(\Omega)}^{r}
$$

For the last factor of (3.1) we have

$$
\int_{U_{t}^{\prime}} d \sigma(z, w)=\sigma\left(U_{t}^{\prime}\right) \lesssim \sigma\left(U^{\prime}\right)
$$

so

$$
A \leq \int_{0}^{h} \frac{1}{\sqrt{t}}\left\{\int_{U_{t}^{\prime}}|\omega(z)| d \sigma(z, w)\right\} d t \lesssim\|\omega\|_{L^{r}(\Omega)}\left(\sigma\left(U^{\prime}\right)\right)^{1 / r^{\prime}} \int_{0}^{h} \frac{d t}{\sqrt{t}}=\frac{1}{2}\|\omega\|_{L^{r}(\Omega)} \sqrt{h} \sigma\left(U^{\prime}\right)^{1 / r^{\prime}}
$$

Recall that $\sigma\left(Q_{j}^{\prime}\right) \simeq h_{j}^{(n+1)}$ then we have

$$
\sqrt{h}=\sqrt{\max _{j} h_{j}} \lesssim \max \sigma\left(Q_{j}^{\prime}\right)^{1 / 2(n+1)} \leq \sigma\left(\bigcup_{j=1}^{N} Q_{j}^{\prime} \cap \partial \Omega^{\prime}\right)^{1 / 2(n+1)}
$$

so finally we get

$$
A:=\int_{T\left(U^{\prime}\right)} \frac{|\omega(z)|}{\sqrt{-\rho^{\prime}(z, w)}} d m(z, w) \lesssim\|\omega\|_{L^{r}(\Omega)^{\prime}} \sigma\left(U^{\prime}\right)^{\frac{1}{r^{\prime}}+\frac{1}{2(n+1)}}
$$

This means that $\frac{|\omega(z)|}{\sqrt{-\rho^{\prime}(z, w)}}$ is a Carleson measure in $\Omega^{\prime}$ of order $\alpha$ with

$$
\alpha=\frac{1}{r^{\prime}}+\frac{1}{2(n+1)}
$$

To get a usual Carleson measure, we need $\alpha=1$ hence

$$
\frac{1}{r^{\prime}}+\frac{1}{2(n+1)}=1 \Longleftrightarrow r=2(n+1)
$$

We have by theorem 1 in [4], written in our situation, that if $\mu \in V^{\alpha}\left(\Omega^{\prime}\right)$ then $P^{0 *}(\mu) \in$ $L^{r, \infty}\left(\partial \Omega^{\prime}\right)$, where $P^{0 *}(\mu)$ is the "balayage" of $\mu$ by the Hardy Littlewood kernel $P_{t}^{0}$. Hence we have that the linear operator $P^{0 *}$ sends $V^{\alpha_{0}}\left(\Omega^{\prime}\right)$ to $L^{r_{0}, \infty}\left(\partial \Omega^{\prime}\right)$, and $V^{\alpha_{1}}\left(\Omega^{\prime}\right)$ to $L^{r_{1}, \infty}\left(\partial \Omega^{\prime}\right)$ with, as usual, $\alpha_{j}=1-\frac{1}{r_{j}}$. This means that

$$
f \in L^{r}(\Omega) \Rightarrow \mu:=f / \sqrt{-\rho^{\prime}} d m \in V^{\alpha}\left(\Omega^{\prime}\right) \Rightarrow P^{0 *}(\mu) \in L^{s, \infty}\left(\partial \Omega^{\prime}\right)
$$

with control of the norms.
So we have a linear operator $T$ such that, with $r_{0}<r_{1}$,

$$
\begin{aligned}
& T: L^{r_{0}}(\Omega) \rightarrow L^{s_{0}, \infty}\left(\partial \Omega^{\prime}\right), \text { with } \frac{1}{s_{0}}=\frac{1}{r_{0}}-\frac{1}{2(n+1)} \\
& T: L^{r_{1}}(\Omega) \rightarrow L^{s_{1}, \infty}\left(\partial \Omega^{\prime}\right), \text { with } \frac{1}{s_{1}}=\frac{1}{r_{1}}-\frac{1}{2(n+1)}
\end{aligned}
$$

hence we can apply Marcinkiewich interpolation theorem between these two values of $r \in] 1,2(n+1)$ [ i.e.

$$
T: L^{r}(\Omega) \rightarrow L^{s}\left(\partial \Omega^{\prime}\right), \text { with } \frac{1}{s}=\frac{1}{r}-\frac{1}{2(n+1)} \text { and } r \leq s
$$

which is needed to apply Marcinkiewich theorem, with control of norms. But this implies by theorem 2 in [4], that $\mu:=f / \sqrt{-\rho^{\prime}} d m \in W^{\alpha}\left(\Omega^{\prime}\right)$.

## 4 The main result.

Let $\Omega$ be a domain in $\mathbb{C}^{n}$ defined by the function $\rho$ as above ; define $\Omega^{\prime} \subset \mathbb{C}^{n+1}$ the lifted domain : we shall define the anisotropic class $\Gamma^{\beta}\left(\partial \Omega^{\prime}\right)$ as in [4] ; we say that a vector field $X$ on $\partial \Omega^{\prime}$ is admissible if $X$ is of class $\mathcal{C}^{k}$ and at any point of $\zeta \in \partial \Omega^{\prime}, X(\zeta)$ belongs to the complex tangent space of $\partial \Omega^{\prime}$ at $\zeta$.

We say that $u \in \Gamma^{\beta}\left(\partial \Omega^{\prime}\right)$ if $u$ is bounded on $\partial \Omega^{\prime}$ and $u$ belongs to the usual Lipschitz $\Lambda^{\beta / 2}\left(\partial \Omega^{\prime}\right)$, where $\partial \Omega^{\prime}$ is viewed as a real manifold, and on any integral curve of an admissible vector field, $t \in[0,1] \rightarrow \gamma(t) \in \partial \Omega^{\prime}$, the function $u \circ \gamma$ belongs to $\Lambda^{\beta}(0,1)$.

We can now define the class $\Gamma^{\beta}(\Omega):$ take a function $u$ defined in $\Omega$ and lift it as $u^{\prime}(z, w):=u(z)$ in $\Omega^{\prime}$; then $u \in \Gamma^{\beta}(\Omega)$ if $u^{\prime} \in \Gamma^{\beta}\left(\partial \Omega^{\prime}\right)$. We have that $u \in \Gamma^{\beta}(\Omega)$ implies that $u \in L^{\infty}(\Omega)$ and $u \in \Lambda^{\beta / 2}(\Omega)$ with a Lipschitz constant uniform in $\Omega$.

The same way we define function $u \in B M O(\Omega)$ if $u^{\prime} \in B M O\left(\partial \Omega^{\prime}\right)$. We have that $u \in B M O(\Omega)$ implies that $u \in \bigcap_{r \geq 1} L^{r}(\Omega)$.

Now we are in position to prove our main result.
Theorem 4.1 Let $\Omega$ be a s.p.c. domain in $\mathbb{C}^{n}$ then for $1<r<2 n+2$ we have $\forall \omega \in L_{(p, q)}^{r}(\Omega), \bar{\partial} \omega=0, \exists u \in L_{(p, q-1)}^{s}(\Omega):: \bar{\partial} u=\omega,\|u\|_{L^{s}(\Omega)} \lesssim\|\omega\|_{L^{r}(\Omega)}$,
with $\frac{1}{s}=\frac{1}{r}-\frac{1}{2(n+1)}$.
For $r=2 n+2$ we have

$$
\exists u \in B M O_{(p, q)}(\Omega):: \bar{\partial} u=\omega,\|u\|_{B M O(\Omega)} \lesssim\|\omega\|_{L^{2 n+2}(\Omega)}
$$

If $\omega$ is a $(p, 1)$ form we have also :
for $r=1$, we have

$$
\exists u \in L_{(p, 0)}^{s, \infty}(\Omega):: \bar{\partial} u=\omega,\|u\|_{L^{s, \infty}(\Omega)} \lesssim\|\omega\|_{L^{1}(\Omega)}
$$

with $\frac{1}{s}=1-\frac{1}{2(n+1)}$.
for $r>2 n+2$ we have

$$
\exists u \in \Gamma_{(p, 0)}^{\beta}(\Omega):: \bar{\partial} u=\omega,\|u\|_{\Gamma^{\beta}(\Omega)} \lesssim\|\omega\|_{L^{r}(\Omega)},
$$

where $\beta=1-\frac{2(n+1)}{r}$ and $\Gamma^{\beta}$ is an anisotropic Lipschitz class of functions.
Moreover the solution $u$ is linear on the data $\omega$.
Proof.
By use of theorem 3.1 we have that $\omega^{\prime} \in W_{(p, q)}^{\alpha}\left(\Omega^{\prime}\right)$ with $\alpha=\frac{1}{r^{\prime}}+\frac{1}{2(n+1)}$ where $\Omega^{\prime}$ is still s.p.c. [5], hence we can apply the theorem 7 in [4] if $\omega$ is a $(p, 1)$ current or the generalisation to $(p, q)$ current done in theorem 4.1 in [6] to get that there is a solution $u^{\prime}$ of $\bar{\partial}_{b} u^{\prime}=\omega^{\prime}$ such that

$$
u^{\prime} \in L_{(p, q-1)}^{s}\left(\partial \Omega^{\prime}\right) \text { with } \frac{1}{s}=1-\alpha=\frac{1}{r}-\frac{1}{2(n+1)}
$$

Because $\omega^{\prime}$ does not depend on $w$ we have that the coefficients of $u^{\prime}$ are holomorphic in $w$ hence with

$$
u^{\prime}(z, w)=\sum_{I, J} a_{I, J}^{\prime}(z, w) d z^{I} \wedge d \bar{z}^{J}
$$

we can set (recall that $u^{\prime}$ is defined on $\partial \Omega^{\prime}$ )

$$
\forall z \in \Omega, a_{I, J}(z):=\int_{|w|^{2}=-\rho(z)} a_{I, J}^{\prime}(z, w) d|w|
$$

and we set also

$$
u(z):=\sum_{I, J} a_{I, J}(z, w) d z^{I} \wedge d \bar{z}^{J}
$$

then exactly as in [3] we still have

$$
\bar{\partial} u=\omega \text { in } \Omega .
$$

Moreover the subordination lemma [5], gives us $u \in L_{(p, q-1)}^{s}(\Omega)$.
The last two results came directly from [4], theorem 7 and theorem 8 with the fact that we apply them in $\Omega^{\prime} \subset \mathbb{C}^{n+1}$ so we have from theorem 8 that $\beta=2(n+1)(\alpha-1)$.

Remark 4.2 In the range $1<r<2 n+2$ theorem 1.2 is stronger than theorem 1.1 because we get the result with $\frac{1}{s}=\frac{1}{r}-\frac{1}{2(n+1)}$ and not only for $\frac{1}{s}>\frac{1}{r}-\frac{1}{2(n+1)}$.

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