



HAL
open science

Global Continuation beyond Singularities on the Boundary for a Degenerate Diffusive Hamilton-Jacobi Equation

Amal Attouchi, Guy Barles

► **To cite this version:**

Amal Attouchi, Guy Barles. Global Continuation beyond Singularities on the Boundary for a Degenerate Diffusive Hamilton-Jacobi Equation. *Journal de Mathématiques Pures et Appliquées*, 2015, 104 (2), pp.383-402. hal-00904365

HAL Id: hal-00904365

<https://hal.science/hal-00904365v1>

Submitted on 14 Nov 2013

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Global Continuation beyond Singularities on the Boundary for a Degenerate Diffusive Hamilton-Jacobi Equation

A. Attouchi* G. Barles†

Abstract

In this article, we are interested in the Dirichlet problem for parabolic viscous Hamilton-Jacobi Equations. It is well-known that the gradient of the solution may blow up in finite time on the boundary of the domain, preventing a classical extension of the solution past this singularity. This behavior comes from the fact that one cannot prescribe the Dirichlet boundary condition for all time and, in order to define a solution globally in time, one has to use “generalized boundary conditions” in the sense of viscosity solution. In this work, we treat the case when the diffusion operator is the p-Laplacian where the gradient dependence in the diffusion creates specific difficulties. In this framework, we obtain the existence and uniqueness of a continuous, global in time, viscosity solution. For this purpose, we prove a Strong Comparison Result between semi-continuous viscosity sub and super-solutions. Moreover, the asymptotic behavior of $\frac{u(x, t)}{t}$ is analyzed through the study of the associated ergodic problem.

Key-words : Viscous Hamilton-Jacobi Equations, generalized Dirichlet problem, maximum principle, viscosity solutions, p-Laplacian.

AMS subject classifications : 35K10, 35K20, 49L25, 53C44 , 35B50, 35B05.

*Université Paris 13, Sorbonne Paris Cité, Laboratoire Analyse, Géométrie et Applications (UMR CNRS 7539), 99, avenue Jean-Baptiste Clément 93430 - Villetaneuse, France.

†Laboratoire de Mathématiques et Physique Théorique (UMR CNRS 7350), Fédération Denis Poisson (FR CNRS 2964), Université François Rabelais, Parc de Grandmont, 37200 Tours, France.

1 Introduction and Main Results

In this article we are interested in the following generalized Dirichlet problem for second-order degenerate parabolic partial differential equations

$$u_t - \operatorname{div}(|Du|^{p-2}Du) + |Du|^q = f(x, t) \quad \text{in } \Omega \times (0, +\infty) \quad (1.1)$$

$$u(x, 0) = u_0(x) \quad \text{on } \overline{\Omega}, \quad (1.2)$$

$$u(x, t) = g(x, t) \quad \text{on } \partial\Omega \times (0, +\infty), \quad (1.3)$$

where $q > p \geq 2$, u_0 and g are continuous functions satisfying the compatibility condition

$$u_0(x) = g(x, 0) \quad \text{on } \partial\Omega \quad (1.4)$$

Most of works devoted to this degenerate diffusive Hamilton-Jacobi equation concerned the case where $\Omega = \mathbb{R}^N$, providing results on well-posedness, gradient estimates and asymptotic behavior of either classical or weak solutions in the sense of distributions (see [10, 1, 24, 22] and the references therein).

Some other works are concerned with the solvability of the Cauchy-Dirichlet problem. They proved that, under suitable assumptions on u_0 and g , there exists a weak solution on some time interval $[0, T_{max}(u_0))$, with the property that its gradient blows up on the boundary $\partial\Omega$ while the solution itself remains bounded. We refer the reader to [2] [19] and [23] for the degenerate parabolic case and to [25] for the uniformly parabolic case. This singularity is a difficulty to extend the solution past $T_{max}(u_0)$. A natural question is then: Can we extend the weak solution past $t = T_{max}(u_0)$ and in which sense ?

Let us mention here that a result in this direction where the continuation beyond gradient blow-up does not satisfy the original boundary conditions was obtained in [16, 17].

Recently, for the linear diffusion case ($p = 2$), Barles and Da Lio [6] showed that such gradient blow-up is related to a loss of boundary condition and address the problem through a viscosity solutions approach. They proved a "Strong Comparison Result" (that is a comparison result between *discontinuous* viscosity sub and supersolutions) which allowed them to obtain the existence of a unique continuous, *global in time viscosity solution* of (1.1)–(1.3), the Dirichlet boundary condition being understood in the *generalized sense of viscosity solution theory*. They also provided an explicit expression of the solution of (1.1)–(1.3) in terms of a value function of some exit time control problem, which allows a simple explanation of the losses of boundary condition when it arises.

We recall that the formulation of the generalized Dirichlet boundary condition for (1.1)–(1.3) in the viscosity sense reads

$$\min(u_t - \operatorname{div}(|Du|^{p-2}Du) + |Du|^q - f(x, t), u - g) \leq 0 \quad \text{on } \partial\Omega \times (0, +\infty), \quad (1.5)$$

and

$$\max(u_t - \operatorname{div}(|Du|^{p-2}Du) + |Du|^q - f(x, t), u - g) \geq 0 \quad \text{on } \partial\Omega \times (0, +\infty). \quad (1.6)$$

Our first result mainly extends the investigation of [6] to the degenerate diffusion case $p > 2$.

Theorem 1.1. *Assume that $q > p \geq 2$ and that Ω is a bounded domain with a C^2 -boundary. For any $u_0 \in C(\overline{\Omega})$, $f \in C(\overline{\Omega} \times [0, T])$ and $g \in C(\partial\Omega \times [0, T])$ satisfying (1.4), there exists a unique continuous solution u of (1.1)–(1.3) which is defined globally in time.*

As it is classical in viscosity solutions theory, the proof of Theorem 1.1 relies on a Strong Comparison Result (SCR in short), the existence of the global solution u being an almost immediate consequence of the Perron's method introduced in the context of viscosity solutions by Ishii [18] (see also [13]).

The most important difficulties in the proof of Strong Comparison Results come from the formulation of the boundary condition in the viscosity sense, the discontinuity of the sub and the supersolution to be compared and the strong nonlinearity of the Hamiltonian term $|Du|^q$. A key argument in the proof of the SCR in [6] is the "cone condition" which is useful in the treatment of boundary points. Roughly speaking the "cone condition" holds if at any point (\tilde{x}, \tilde{t}) of the boundary $\partial\Omega \times (0, T)$, an usc subsolution u satisfies $u(\tilde{x}, \tilde{t}) = \lim_{k \rightarrow \infty} u(x_k, t_k)$ where $\{(x_k, t_k)\}_k$ is a sequence of points of $\Omega \times (0, T)$ with the following properties

$$(x_k, t_k) \rightarrow (\tilde{x}, \tilde{t}) \quad \text{and} \quad d_{\partial\Omega}(x_k, t_k) \geq b(|x_k - \tilde{x}| + |t_k - \tilde{t}|),$$

where b is a positive constant.

Our approach is slightly different: instead of directly proving the "cone condition" for any viscosity subsolution of (1.1)–(1.3) as it was done in [6], we use a combination of a $C^{0,\beta}$ regularity result for subsolutions of stationary problems, strongly inspired by the result of Capuzzo Dolcetta, Leoni and Porretta [11], together with a regularization by a sup-convolution in time. These arguments provide an approximation of the (a priori only usc) subsolution by a continuous subsolution, which automatically satisfies the "cone condition", allowing to borrow the methods of [9] to conclude.

The generalisation of the $C^{0,\beta}$ regularity result of [11] is the following.

Theorem 1.2. *If u is a locally bounded, usc viscosity subsolution of*

$$- \operatorname{div}(|Du|^{p-2}Du) + |Du|^q \leq C \quad \text{in } \Omega, \tag{1.7}$$

where Ω is an open subset of \mathbb{R}^N and C is a positive constant, and if $q > p \geq 2$, then $u \in C_{loc}^{0,\beta}(\Omega)$ with $\beta = \frac{q-p}{q-p+1}$.

Moreover, if Ω is a bounded domain with a C^2 -boundary, then u is bounded on $\overline{\Omega}$ and it can be extended as a $C^{0,\beta}$ -function on $\overline{\Omega}$ and

$$|u(x) - u(y)| \leq M|x - y|^\beta \quad \text{for all } x, y \in \overline{\Omega}, \tag{1.8}$$

for some positive constant M depending only on p, q, C and $\partial\Omega$.

The regularity result of [11] was revisited in [4], where an interpretation was given in terms of state-constraint problems together with several possible applications. Our proof will rely on the arguments of [4].

A second motivation where such regularity results are useful, is the asymptotic behavior as $t \rightarrow +\infty$ of solutions of the evolution equation. For this purpose, one has first to study the ergodic (or additive eigenvalue) problem

$$-\operatorname{div}(|Du_\infty|^{p-2}Du_\infty) + |Du_\infty|^q - \tilde{f}(x) = c \quad \text{in } \Omega, \quad (1.9)$$

associated to a state-constraint boundary condition on $\partial\Omega$

$$-\operatorname{div}(|Du_\infty|^{p-2}Du_\infty) + |Du_\infty|^q - \tilde{f}(x) \geq c \quad \text{on } \partial\Omega. \quad (1.10)$$

We recall that, in this type of problems, both the solution u_∞ and the constant c (the ergodic constant) are unknown. First we have the following result.

Theorem 1.3. *Assume that Ω is a bounded domain with a C^2 -boundary, $\tilde{f} \in C(\overline{\Omega})$ and $q > p \geq 2$, then there exists a unique constant c such that the state-constraints problem (1.9)–(1.10) has a continuous viscosity solution u_∞ .*

A typical result that connects the study of the ergodic problem to the large time behavior of the solution u of (1.1)–(1.3) is the following.

Theorem 1.4. *Assume that Ω is a bounded domain with a C^2 -boundary, $u_0 \in C(\overline{\Omega})$, $g \in C(\partial\Omega)$ satisfying (1.4) and assume that $f(x, t) = \tilde{f}(x)$ with $\tilde{f} \in C(\overline{\Omega})$ and $q > p \geq 2$. If (c, u_∞) is the solution of (1.9)–(1.10) and if u is the unique viscosity solution of (1.1)–(1.3), then $u + c^+t$ is bounded, where $c^+ = \max(c, 0)$. In particular*

$$\lim_{t \rightarrow \infty} \frac{u(x, t)}{t} = -c^+$$

uniformly on $\overline{\Omega}$.

The next step in the study of the asymptotic behavior would be to show that $u(x, t) + ct \rightarrow u_\infty(x)$ as $t \rightarrow \infty$ where u_∞ solves (1.9)–(1.10). The main difficulty to prove such more precise asymptotic behavior comes from the fact that (1.9)–(1.10) does not admit a unique solution ((1.9)–(1.10) is invariant by addition of constants). Such results were obtained recently in [26] for the uniformly elliptic case $p = 2$ through the use of the Strong Comparison Principle (i.e. a result which allows to apply the Strong Maximum Principle to the difference of solutions) and the Lipschitz regularity of u_∞ . But, for $p > 2$, such Strong Comparison Principle is not available since the equation is quasilinear and not semilinear. We recall that a Strong Maximum Principle is available for $p > 2$, see [3]. Another difficulty comes from the proof of a strong comparison result for the steady problem in case of an operator that does not fulfill a monotonicity property, even if there exists a strict subsolution. Let us mention the works of [23, 7] for more results on the asymptotic behavior of global solutions.

Finally we point out that it was shown in [8] that the expected asymptotic behavior, namely $u(x, t) + ct \rightarrow u_\infty(x)$, is not always true in the $p = 2$ -case when the nonlinearity is sub quadratic in Du .

This article is organized as follows: in Section 2, we present the needed results on viscosity solutions for the stationary and evolution problems we consider; in particular, we analyze the losses of boundary conditions for subsolutions. In Section 3 we prove the Hölder regularity result of Theorem 1.2. In Section 4 we study the ergodic problem. Section 5 is devoted to the proof of Theorem 1.1 and the asymptotic behavior of solutions of the evolution equation.

2 Preliminaries and Analysis of Boundary Conditions

In this section we collect some preliminary properties of viscosity subsolutions (the boundary conditions being always understood in the viscosity sense) and we also formulate SCR under different forms, some of them being only useful as a step in the proof of the complete regularity result. These results are concerned with either problem (1.1)–(1.3) or the following two nonlinear elliptic problem

$$-(p-1)|Du|^{p-2} \sum_{\lambda_i(D^2u) > 0} \lambda_i(D^2u) + |Du|^q = C \quad \text{in } \Omega, \quad u = \tilde{g} \quad \text{in } \partial\Omega. \quad (2.1)$$

and

$$- \operatorname{div} (|Du|^{p-2} Du) + |Du|^q + \lambda u - \tilde{f} = 0 \quad \text{in } \Omega, \quad u = \tilde{g} \quad \text{in } \partial\Omega. \quad (2.2)$$

where $q > p \geq 2$, $C, \lambda \geq 0$, $\tilde{f} \in C(\overline{\Omega})$ and $\tilde{g} \in C(\partial\Omega)$.

From now on, we assume that Ω is a smooth domain with a C^2 -boundary. We define the distance from $x \in \overline{\Omega}$ to $\partial\Omega$ by $d_{\partial\Omega}(x) := \operatorname{dist}(x, \partial\Omega)$. For $\delta > 0$, we denote by

$$\Omega^\delta := \{x \in \Omega \mid d_{\partial\Omega}(x) < \delta\}, \quad (2.3)$$

$$\Omega_\delta := \{x \in \Omega \mid d_{\partial\Omega}(x) > \delta\}. \quad (2.4)$$

As a consequence of the regularity of $\partial\Omega$, $d_{\partial\Omega}$ is a C^2 -function in a neighborhood Ω^δ of the boundary for all $0 < \delta \leq \delta_0$. We denote by d a C^2 -function agreeing with $d_{\partial\Omega}$ in Ω^δ such that $|Dd(x)| \leq 1$ in Ω_δ . We also denote by $n(x)$ the C^1 -function defined by $n(x) = -Dd(x)$ in Ω^δ ; if $x \in \partial\Omega$, then $n(x)$ is just the unit outward normal vector to $\partial\Omega$ at x .

Our first result says that there is no loss of boundary conditions for the subsolutions, namely that the subsolutions satisfy the boundary condition in the classical sense.

Proposition 2.1. *Assume that $q > 0$ and $p \geq 2$. We have the following*

i) If u is a bounded, usc subsolution of (1.1)–(1.3) on a time interval $(0, T)$, then

$$u \leq g \quad \text{on } \partial\Omega \times (0, T). \quad (2.5)$$

ii) If u is a bounded, usc subsolution of (2.1) or (2.2), then

$$u \leq \tilde{g} \quad \text{on } \partial\Omega. \quad (2.6)$$

Proof. We only give the proof for the time dependent problem, the proof for the stationary problems being similar. We use a result of Da Lio [14, Corollary 6.2]. We denote by \mathcal{S}^N the space of real symmetric $N \times N$ matrices. For $x \in \Omega$, $t \in (0, T)$, $\xi \in \mathbb{R}^N$ and $M \in \mathcal{S}^N$, we define the function F by

$$F(x, t, \xi, M) = -|\xi|^{p-2} \text{Tr}(M) - (p-2)|\xi|^{p-4} \langle M\xi, \xi \rangle + |\xi|^q - f(x, t),$$

so that the equation can be written as $u_t + F(x, t, Du, D^2u) = 0$. From [14], we know that, if $u(x_0, t_0) > g(x_0, t_0)$ at some point $(x_0, t_0) \in \partial\Omega \times (0, T)$, then the following conditions hold

$$\begin{aligned} \liminf_{\substack{(y,t) \rightarrow (x_0,t_0) \\ \alpha \downarrow 0}} \left\{ \left[\frac{o(1)}{\alpha} + F\left(y, t, \frac{Dd(y) + o(1)}{\alpha}, -\frac{Dd(y) \otimes Dd(y) + o(1)}{\alpha^2}\right) \right] \right\} &\leq 0 \\ \liminf_{\substack{(y,t) \rightarrow (x_0,t_0) \\ \alpha \downarrow 0}} \left\{ \left[\frac{o(1)}{\alpha} + F\left(y, t, \frac{Dd(y) + o(1)}{\alpha}, \frac{D^2d(y) + o(1)}{\alpha}\right) \right] \right\} &\leq 0. \end{aligned} \quad (2.7)$$

But the first condition cannot hold since

$$F\left(y, t, \frac{Dd(y) + o(1)}{\alpha}, -\frac{Dd(y) \otimes Dd(y) + o(1)}{\alpha^2}\right) \geq \frac{(p-1)}{\alpha^p} (1 + o(1)) + \frac{1 - o(1)}{\alpha^q} - f(y, t),$$

and the right hand side is going to $+\infty$ as $\alpha \rightarrow 0$ since $p \geq 2$, $q > 0$ and all terms converge to $+\infty$. \square

Let us point out that the above computation shows that there is no competition between the nonlinear Hamiltonian term and the slow diffusion operator since they both produce positive contribution which prevent any loss of boundary conditions for the subsolution.

Next, we remark that there cannot be loss of initial condition.

Lemma 2.1. *Assume that $q > p \geq 2$, $f \in C(\overline{\Omega} \times [0, T])$ and $u_0 \in C(\overline{\Omega})$, $g \in C(\partial\Omega \times [0, T])$ satisfy (1.4). Let u and v be respectively a bounded usc viscosity subsolution and a bounded lsc super-solution of (1.1)–(1.3) then*

$$u(x, 0) \leq u_0(x) \leq v(x, 0) \quad \text{on } \overline{\Omega}. \quad (2.8)$$

Proof Fix $x_0 \in \overline{\Omega}$ and define for $\varepsilon > 0$ and $C_\varepsilon > 0$ the function $\phi_\varepsilon(x, t)$ by

$$\phi_\varepsilon(x, t) = u(x, t) - \frac{|x - x_0|}{\varepsilon^2} - C_\varepsilon t.$$

This function attains a global maximum on $\overline{\Omega} \times [0, T)$ at $x_\varepsilon, t_\varepsilon$. Using the boundedness of u , it is easy to see that, for any $C_\varepsilon > 0$, $(x_\varepsilon, t_\varepsilon) \rightarrow (x_0, 0)$ as $\varepsilon \rightarrow 0$. Arguing as in [9], choosing C_ε sufficiently large depending on ε , we are left with $(x_\varepsilon, t_\varepsilon) \in (\partial\Omega \times (0, T)) \cup (\overline{\Omega} \times \{0\})$ and the two following possibilities

$$\begin{aligned} &\text{either } t_\varepsilon = 0 && \text{and } u(x_\varepsilon, 0) \leq u_0(x_\varepsilon), \\ \text{or } &t_\varepsilon > 0, x_\varepsilon \in \partial\Omega && \text{and } u(x_\varepsilon, t_\varepsilon) \leq g(x_\varepsilon, t_\varepsilon). \end{aligned}$$

In either case, since $u(x_0, 0) \leq \phi_\varepsilon(x_\varepsilon, t_\varepsilon) \leq u(x_\varepsilon, t_\varepsilon)$, we get the desired result for u letting $\varepsilon \rightarrow 0$ and using the continuity of u_0 and g . The argument for v is similar.

Now we claim that under some assumptions (set out below), a SCR holds for semicontinuous viscosity sub- and supersolutions of (1.1)–(1.3) or (2.1) or (2.2). The proof being somehow technical we refer the reader to the appendice for a detailed proof of the following two propositions.

Proposition 2.2 (Parabolic SCR). *Assume that $q > p \geq 2$, $f \in C(\overline{\Omega} \times [0, T])$ and $u_0 \in C(\overline{\Omega})$, $g \in C(\partial\Omega \times [0, T])$ satisfy (1.4). Let u and v be respectively a bounded usc viscosity subsolution and a bounded lsc super-solution of (1.1)–(1.3), then $u \leq v$ in $\Omega \times [0, T]$. Moreover, if we define \tilde{u} on $\overline{\Omega} \times [0, T]$ by setting*

$$\tilde{u}(x, t) := \begin{cases} \limsup_{\substack{(y,s) \rightarrow (x,t) \\ (y,s) \in \Omega \times (0,T)}} u(y, s) & \text{for all } (x, t) \in \partial\Omega \times (0, T] \\ u(x, t) & \text{otherwise,} \end{cases} \quad (2.9)$$

then \tilde{u} remains an usc subsolution of (1.1)–(1.3) and

$$\tilde{u} \leq v \quad \text{on } \overline{\Omega} \times [0, T]. \quad (2.10)$$

The stationary version of the SCR is used either in the proof of the $C^{0,\beta}$ -regularity or for solving the ergodic problem.

Proposition 2.3 (Elliptic SCR). *Assume that $q > p \geq 2$, $\tilde{f} \in C(\overline{\Omega})$ and $\tilde{g} \in C(\partial\Omega)$.*

(i) *Let u and v be respectively a bounded usc viscosity subsolution and a bounded lsc super-solution of (2.1). If v is continuous on $\overline{\Omega}$ and is a strict supersolution of (2.1), then*

$$u \leq v \quad \text{on } \overline{\Omega}. \quad (2.11)$$

(ii) *Let u and v be respectively a bounded usc viscosity subsolution and a bounded lsc super-solution of (2.2). Assume that either $\lambda > 0$ or $\lambda = 0$ and v is a strict supersolution. We define \tilde{u} on $\overline{\Omega}$ by setting*

$$\tilde{u}(x) := \begin{cases} \limsup_{\substack{y \rightarrow x \\ y \in \Omega}} u(y) & \text{for all } x \in \partial\Omega \\ u(x) & \text{otherwise,} \end{cases} \quad (2.12)$$

then \tilde{u} remains an usc subsolution of (1.1)–(1.3) and

$$\tilde{u} \leq v \quad \text{on } \overline{\Omega}. \quad (2.13)$$

3 Hölder Regularity of Viscosity Subsolutions for the Degenerate Elliptic Problem

In this section we are going to prove that equation of type (1.7) enters into the general framework described in [4] which allows us to state that, if u is a locally bounded, usc viscosity subsolution of (1.7), then u is Hölder continuous with exponent $\beta = \frac{q-p}{q-p+1}$. The key point is that the strong growth of the first order term balances the degeneracy of the second order term, providing a control on $|Du|$.

Proof of Theorem 1.2. If u is a subsolution of (1.7), then it is a subsolution in $B_r(x) = \{y \in \mathbb{R}^N; |y-x| < r\}$ of the simpler equation

$$-(p-1)|Du|^{p-2} \sum_{\lambda_i(D^2u) > 0} \lambda_i(D^2u) + |Du|^q \leq C.$$

Now we are going to check the required hypotheses in [4].

H1. For $0 < r < 1$, $s \in \mathbb{R}^N$ and $M \in \mathcal{S}^N$, \mathcal{S}^N denoting the space of $N \times N$ real valued symmetric matrices, define the function $G_r(s, M)$ by

$$G_r(s, M) := -(p-1)|s|^{p-2} \sum_{\lambda_i(M) > 0} \lambda_i(M) + |s|^q - C.$$

Then, for any $x \in \Omega$ with $d_{\partial\Omega}(x) \geq r$, $G_r(Du, D^2u) \leq 0$ in $B_r(x)$.

H2. There exists a super-solution up to the boundary $w_r \in C(\overline{B_r(0)})$ such that $w_r(0) = 0$, $w_r(x) \geq 0$ in $B_r(x)$ and

$$G_r(Dw_r, D^2w_r) \geq \eta_r > 0 \quad \text{on } \overline{B_r(0)} \setminus \{0\}, \quad (3.1)$$

for some $\eta_r > 0$.

Despite the construction of the functions w_r is a rather easy adaptation of [4], we reproduce it for the sake of completeness and for the reader's convenience. In order to build w_r , we first build w_1 and then use the scale invariance of the equation. To do so, we borrow arguments from [4]. For $C_1, C_2 > 0$ to be chosen later on and for $\beta = \frac{q-p}{q-p+1}$, we consider the function

$$w_1(x) := \frac{C_1}{\beta}|x|^\beta + \frac{C_2}{\beta}(d^\beta(0) - d^\beta(x)),$$

where $d(x) = 1 - |x|$ on $B_1(0) \setminus B_{1/2}(0)$ and we regularize it in $B_{1/2}(0)$ by changing it into $h(1 - |x|)$ where h is a smooth, non-decreasing and concave function such that $h(s)$ is constant for $s \geq 3/4$ and $h(s) = s$ for $s \leq 1/2$. Obviously we have $w_1(0) = 0$, $w_1 \geq 0$ in $\overline{B_1(0)}$ and w_1 is smooth in $\overline{B_1(0)} \setminus \{0\}$.

We first remark that $-(p-1)|Dw_1(x)|^{p-2}\lambda_i(D^2w_1(x)) + |Dw_1(x)|^q$ can be written as

$$|Dw_1(x)|^{p-2} [-(p-1)\lambda_i(D^2w_1(x)) + |Dw_1(x)|^{q-p+2}] .$$

Therefore, in order to prove the claim, we are going to show that, for $C_1, C_2 > 0$ large enough, the bracket is positive and bounded away from 0 and that $|Dw_1(x)|^{p-2}$ remains large.

Computing the derivatives of w_1 in $B_1(0) \setminus \{0\}$, we have

$$\begin{aligned} Dw_1(x) &= C_1|x|^{\beta-2}x - C_2d^{\beta-1}(x)Dd(x), \\ D^2w_1(x) &= C_1|x|^{\beta-2}Id + (\beta-2)C_1|x|^{\beta-4}x \otimes x \\ &\quad - C_2d^{\beta-1}(x)D^2d(x) - (\beta-1)C_2d^{\beta-2}(x)Dd(x) \otimes Dd(x). \end{aligned}$$

Using that $-Dd(x) = \mu(x)x$ for some $\mu(x) \geq 0$ and that $q > p > 2$, we have

$$\begin{aligned} |Dw_1(x)|^{q-p+2} &= (|C_1|x|^{\beta-2}x| + |C_2d^{\beta-1}(x)Dd(x)|)^{q-p+2} \\ &\geq |C_1|x|^{\beta-2}x|^{q-p+2} + |C_2d^{\beta-1}(x)Dd(x)|^{q-p+2} \\ &= C_1^{q-p+2}|x|^{(\beta-1)(q-p+2)} + C_2^{q-p+2}d^{(\beta-1)(q-p+2)}(x)|Dd(x)|^{q-p+2}, \end{aligned}$$

and

$$|Dw_1(x)|^{p-2} \geq C_1^{p-2}|x|^{(\beta-1)(p-2)} + C_2^{p-2}d^{(\beta-1)(p-2)}(x)|Dd(x)|^{p-2}.$$

Using that $d(x) = h(1 - |x|)$, with h being C^2 , non-decreasing and concave, $0 < \beta < 1$, we have

$$D^2w_1(x) \leq C_1|x|^{\beta-2}Id + C_2d^{\beta-1}(x) \left(\frac{h'}{|x|}Id - h''\frac{x}{|x|} \otimes \frac{x}{|x|} \right) + (1-\beta)C_2d^{\beta-2}(x)Dd(x) \otimes Dd(x),$$

and

$$\lambda_i(D^2w_1(x)) \leq C_1|x|^{\beta-2} + C_2d^{\beta-1}(x) \left(\frac{h'}{|x|} - h'' \right) + (1-\beta)C_2d^{\beta-2}(x)|Dd(x)|^2.$$

At this point, it is worth noticing that because of the properties of h , the term $\left(\frac{h'}{|x|} - h''\right)$ is bounded.

These properties imply that, we can (almost) consider the two terms (in $|x|$ and in $d(x)$) separately. Since $(\beta-1)(q-p+2) = (\beta-2)$, the $\frac{C_1}{\beta}|x|^\beta$ term yields

$$-(p-1)C_1|x|^{\beta-2} + |C_1|x|^{\beta-2}x|^{q-p+2} = |x|^{\beta-2} (-(p-1)C_1 + C_1^{q-p+2}).$$

By choosing C_1 large enough, we can have for any $K_1 > 0$

$$|x|^{\beta-2} (-(p-1)C_1 + C_1^{q-p+2}) \geq K_1|x|^{\beta-2} \quad \text{in } B_1(0) \setminus \{0\}.$$

On the other hand the $\frac{C_2}{\beta}(d^\beta(0) - d^\beta(x))$ term yields

$$-(p-1)C_2d^{\beta-1}(x)\left(\frac{h'}{|x|} - h''\right) + (\beta-1)(p-1)C_2d^{\beta-2}(x)|Dd(x)|^2 + C_2^{q-p+2}|d^{\beta-1}Dd(x)|^{q-p+2}. \quad (3.2)$$

We have to consider two cases: either $|x| \geq \frac{1}{2}$ and then $h' = 1$, $h'' = 0$ and $Dd(x) = -\frac{x}{|x|}$; hence the above quantity is given by

$$-(p-1)C_2d^{\beta-1}(x)\left(\frac{1}{|x|}\right) - (p-1)(1-\beta)C_2d^{\beta-2}(x) + C_2^{q-p+2}d^{(\beta-1)(q-p+2)}.$$

Recalling that $(\beta-1)(q-p+2) = (\beta-2)$, then for C_2 large enough we have for any $K_2 > 0$

$$-(p-1)C_2d^{\beta-1}(x)\left(\frac{1}{|x|}\right) - (p-1)(1-\beta)C_2d^{\beta-2}(x) + C_2^{q-p+2}d^{(\beta-1)(q-p+2)} \geq K_2d^{\beta-2}(x).$$

Now for $|x| \leq \frac{1}{2}$, the quantity (3.2) coming from the $d(x)$ -term is bounded and can be controlled by the $|x|$ -term. Hence, for any constant $C > 0$, choosing first C_2 large enough and then C_1 large enough, we have in $\overline{B_1(0)} \setminus \{0\}$

$$\begin{aligned} -(p-1)|Dw_1(x)|^{p-2}\lambda_i(D^2w_1(x)) + |Dw_1(x)|^q &\geq |Dw_1(x)|^{p-2}(K_1|x|^{\beta-2} + K_2d^{\beta-2}(x)) \\ &\geq (K_1|x|^{(\beta-1)(p-1)-1}) \geq C. \end{aligned}$$

Next we set

$$w_r(x) := r^\beta w_1\left(\frac{x}{r}\right).$$

It is easy to check that for $0 < r \leq 1$, $G(Dw_r, D^2w_r) \geq r^{(\beta-1)(p-1)-1}C - C \geq 0$ on $\overline{B_r(0)} \setminus \{0\}$.

H3. Comparison result. Let v be any bounded usc viscosity subsolution of $G_r(Dv, D^2v) \leq 0$ in $B_r(0) \setminus \{0\}$ then

$$v(y) \leq v(x) + r^\beta w_1\left(\frac{y-x}{r}\right). \quad (3.3)$$

We use the fact that $v(0) + w_r(x)$ is a **strict** super-solution up to the boundary and that it is a **continuous** function. It follows that the comparison is a direct consequence of Proposition 2.3.

Since the hypotheses are satisfied, we can apply Proposition 2.1 of [4] to obtain the $C^{0,\beta}$ regularity of subsolutions, both locally and globally with further assumptions on Ω .

Remark 3.1. As far as the exponent β is concerned, the value is the best one can expect in the assumption of the above theorem (see [11]).

It is well-known that the degeneracy of the p -Laplacian is an obstruction to the solvability

of the Dirichlet problem in the classical sense. The presence of the strongly non-linear term with $q > p$ is another source of obstruction, even in the uniformly elliptic case since examples of boundary layers can occur [6, 20]. By the previous result, we know that every continuous solution to (1.7) is Hölder continuous up to the boundary. Hence, a necessary condition in order that the solution can attain continuously the boundary data g is the existence of some $C \geq 0$ such that

$$|g(x) - g(y)| \leq C|x - y|^\beta \quad \text{for all } x, y \in \partial\Omega, \quad \beta = \frac{q - p}{q - p + 1}.$$

For the uniformly elliptic case $p = 2$, a more detailed study including several gradient bounds and applications can be found in [20].

As an application of the previous regularity result, we consider the generalized Dirichlet problem consisting in solving (2.2).

Theorem 3.1. *Let $\Omega \subset \mathbb{R}^N$ be a bounded domain with a C^2 -boundary. Assume that $q > p \geq 2$, $\tilde{f} \in C(\overline{\Omega})$, $\tilde{g} \in C(\partial\Omega)$ and $\lambda > 0$. Let u and v be respectively a bounded usc subsolution and a bounded lsc super-solution of (2.2) with u satisfying for $x \in \partial\Omega$*

$$u(x) = \limsup_{\substack{y \rightarrow x \\ y \in \Omega}} u(y).$$

Then, $u \leq v$ on $\overline{\Omega}$. Moreover Problem (2.2) has a unique viscosity solution which belongs to $C^{0,\beta}(\overline{\Omega})$.

Proof. For the comparison part, Theorem 1.2 implies that u is Hölder continuous, hence the comparison $u \leq v$ is a direct consequence of Proposition 2.3. Once noticed that $-\left(\lambda^{-1}\|\tilde{f}\|_{L^\infty} + \|\tilde{g}\|_{L^\infty}\right)$ and $+\left(\lambda^{-1}\|\tilde{f}\|_{L^\infty} + \|\tilde{g}\|_{L^\infty}\right)$ are respectively sub and super-solution, we can apply the Perron's method with the version up to the boundary (see [13]). Since a solution is also a subsolution, the Hölder regularity is a direct consequence of Theorem 1.2.

4 The Ergodic Problem

4.1 Existence of the pair (c, u_∞)

In this part we study the existence of a pair $(c, u_\infty) \in \mathbb{R} \times C(\overline{\Omega})$ for which u_∞ is a viscosity solution of the state-constraints problem (1.9)-(1.10), to gather with the uniqueness of the ergodic constant c . For this purpose, we introduce a λu -term in the equation, as it is classical, with the aim to let λ tend toward 0. This key step is described by the following Lemma.

Lemma 4.1. *Let $\tilde{f} \in C(\overline{\Omega})$ and $\beta = \frac{q-p}{q-p+1}$. For $0 < \lambda < 1$ and $q > p$, there exists a unique viscosity solution $u_\lambda \in C^{0,\beta}(\overline{\Omega})$ of the state constraint problem*

$$- \operatorname{div}(|Du_\lambda|^{p-2}Du_\lambda) + |Du_\lambda|^q + \lambda u_\lambda = \tilde{f}(x) \quad \text{in } \Omega, \quad (4.1)$$

$$- \operatorname{div}(|Du_\lambda|^{p-2}Du_\lambda) + |Du_\lambda|^q + \lambda u_\lambda \geq \tilde{f}(x) \quad \text{on } \partial\Omega. \quad (4.2)$$

Moreover there exists a constant $\tilde{C} > 0$ such that, for all $0 < \lambda < 1$,

$$|\lambda u_\lambda| \leq \tilde{C} \quad \text{in } \overline{\Omega}. \quad (4.3)$$

Proof. For $R > 0$, we consider the following generalized Dirichlet problem

$$\begin{cases} -\operatorname{div}(|Du_{R,\lambda}|^{p-2}Du_{R,\lambda}) + |Du_{R,\lambda}|^q + \lambda u_{R,\lambda} = f(x) & \text{in } \Omega, \\ u_{R,\lambda} = R & \text{in } \partial\Omega. \end{cases} \quad (4.4)$$

By Theorem 3.1, this problem admits a unique viscosity solution $u_{R,\lambda}$.

Moreover, $u_{R,\lambda}$ satisfies

$$-\lambda^{-1} \|f\|_{L^\infty} \leq u_{R,\lambda} \leq -\frac{M_1}{\beta} d^\beta(x) + \frac{M_2}{\lambda} \quad \text{in } \Omega. \quad (4.5)$$

Indeed, on the one hand, it is easy to see that $-\lambda^{-1} \|f\|_{L^\infty}$ is a subsolution. On the other hand, borrowing arguments from [26], we claim that for some $M_1, M_2 > 0$ chosen large enough, $\bar{u}(x) = -\frac{M_1}{\beta} d^\beta(x) + \frac{M_2}{\lambda}$ is a supersolution of (4.1)-(4.2). Indeed, using that $q(\beta-1) = (p-2)(\beta-1) + (\beta-2)$, we have

$$\begin{aligned} -\operatorname{div}(|D\bar{u}|^{p-2}D\bar{u}) + |D\bar{u}|^q + \lambda\bar{u} - \tilde{f}(x) &= M_1^{p-1} |Dd|^{p-2} d^{(p-2)(\beta-1)} \left[(p-1)(\beta-1) d^{\beta-2} |Dd|^2 \right. \\ &\quad \left. + d^{\beta-1} \Delta d + (p-2) d^{\beta-1} \langle D^2 d \hat{D}d, \hat{D}d \rangle \right] \\ &\quad + M_1^q d^{q(\beta-1)} |Dd|^q - \lambda \frac{M_1}{\beta} d^\beta + M_2 - \tilde{f} \\ &= M_1^{p-1} |Dd|^{p-2} d^{q(\beta-1)} \left[(p-1)(\beta-1) |Dd|^2 + d \Delta d \right. \\ &\quad \left. + (p-2) d \langle D^2 d \hat{D}d, \hat{D}d \rangle + M_1^{q-p+1} |Dd|^{q-p+2} \right] \\ &\quad - \lambda \frac{M_1}{\beta} d^\beta + M_2 - \tilde{f}. \end{aligned}$$

In Ω^δ where $|Dd| = 1$ and $0 \leq d \leq \delta$, we have

$$\begin{aligned} -\operatorname{div}(|D\bar{u}|^{p-2}D\bar{u}) + |D\bar{u}|^q + \lambda\bar{u} - \tilde{f}(x) &= M_1^{p-1} d^{q(\beta-1)} \left[(p-1)(\beta-1) + d \Delta d \right. \\ &\quad \left. + (p-2) d \langle D^2 d Dd, Dd \rangle + M_1^{q-p+1} \right. \\ &\quad \left. - \lambda \frac{M_1^{2-p}}{\beta} d^{\beta(2-p)+p} \right] + M_2 - \tilde{f}. \end{aligned}$$

Taking $M_1 > 1$ and $M_2 > 0$ such that

$$M_1^{q-p+1} \geq (p-1)(1-\beta) + (p-2 + \sqrt{N})\delta \|D^2d\|_{L^\infty} + \frac{\delta^{\beta(2-p)+p}}{\beta} \quad (4.6)$$

$$\text{and } M_2 \geq 2 \|\tilde{f}\|_{L^\infty}, \quad (4.7)$$

then we have $-\operatorname{div}(|D\bar{u}|^{p-2}D\bar{u}) + |D\bar{u}|^q + \lambda\bar{u} - \tilde{f}(x) \geq 0$ in Ω^δ .

Now in Ω_δ , we have $|Dd| \leq 1$ and $\delta \leq d(x) \leq C(\Omega)$. Using that $0 < \lambda < 1$, then we have

$$\begin{aligned} -\operatorname{div}(|D\bar{u}|^{p-2}D\bar{u}) + |D\bar{u}|^q + \lambda\bar{u} - \tilde{f}(x) &\geq M_1^{p-1} \left[(p-1)(\beta-1) \|d^{(\beta-1)(p-1)-1}\|_{L^\infty} \right. \\ &\quad \left. - (p-2 + \sqrt{N}) \|d^{(\beta-1)(p-1)}\|_{L^\infty} \|D^2d\|_{L^\infty} \right] \\ &\quad - \frac{M_1}{\beta} \|d^\beta\|_{L^\infty} + M_2 - \|\tilde{f}\|_{L^\infty}. \end{aligned}$$

Hence if we take M_1 as in (4.6) and M_2 such that

$$\begin{aligned} M_2 &\geq M_1^{p-1} \left[(p-1)(1-\beta) \|d^{(\beta-1)(p-1)-1}\|_{L^\infty} + (p-2 + \sqrt{N}) \|d^{(\beta-1)(p-1)}\|_{L^\infty} \|D^2d\|_{L^\infty} \right] \\ &\quad + \frac{M_1}{\beta} \|d^\beta\|_{L^\infty} + 3 \|\tilde{f}\|_{L^\infty}, \end{aligned} \quad (4.8)$$

then the function \bar{u} satisfies the supersolution inequality in Ω_δ . The estimate follows by applying the SCR to $-\lambda^{-1} \|f\|_{L^\infty}$, $u_{R,\lambda}$ and \bar{u} .

It is worth pointing out that, if M_2 is as in (4.8), then

$$u_{R,\lambda} < R \quad \text{on } \bar{\Omega} \quad \text{for any } R > \frac{M_2}{\lambda}.$$

It follows that $u_{R,\lambda}$ is a viscosity solution of (4.1)-(4.2) for all $R > \frac{M_2}{\lambda}$. Theorem 3.1

implies that $u_\lambda = u_{R,\lambda}$ for $R > \frac{M_2}{\lambda}$.

We have

$$-\max\left(\|\tilde{f}\|_{L^\infty}, M_2\right) \leq \lambda u_\lambda \leq \max\left(\|\tilde{f}\|_{L^\infty}, M_2\right).$$

□

Now we are in position to prove Theorem 1.3. Using that $u_\lambda \geq -\lambda^{-1} \|f\|_{L^\infty}$ in $\bar{\Omega}$, we have

$$-\operatorname{div}(|Du_\lambda|^{p-2}Du_\lambda) + |Du_\lambda|^q - \tilde{f} \leq \|\tilde{f}\|_{L^\infty} \quad \text{in } \Omega.$$

Theorem 1.2 implies uniform Hölder estimates with respect to λ for the functions u_λ . Consequently if x_0 is an arbitrary point in $\bar{\Omega}$, we get that $w_\lambda := u_\lambda(x) - u_\lambda(x_0)$ is also uniformly bounded in $C^{0,\beta}(\bar{\Omega})$ (recall that Ω is connected).

From (4.3), we also know that $\{-\lambda u_\lambda(x_0)\}_\lambda$ is bounded. It follows that, by Ascoli's Theorem, we can extract a uniformly converging subsequence from $\{w_\lambda\}_\lambda$ and we can assume that $\{-\lambda u_\lambda(x_0)\}_\lambda$ converges along the same subsequence. Denoting by u_∞ and c , the limits of $\{w_\lambda\}_\lambda$ and $\{-\lambda u_\lambda(x_0)\}_\lambda$ respectively and taking into account that w_λ solves

$$-\operatorname{div}(|Dw_\lambda|^{p-2}Dw_\lambda) + |Dw_\lambda|^q - \tilde{f}(x) + \lambda w_\lambda = -\lambda u_\lambda(x_0) \quad \text{in } \Omega,$$

we can pass into the limit $\lambda \rightarrow 0$ and conclude by the stability result for viscosity solutions that, (c, u_∞) solves the ergodic problem.

Now let (c_1, u_∞^1) and (c_2, u_∞^2) be two solutions of the ergodic problem. If $c_1 < c_2$ or $c_1 > c_2$, we could use Proposition 2.3 to obtain either $u_\infty^1 \leq u_\infty^2$ or $u_\infty^2 \leq u_\infty^1$. But such comparison cannot hold since, for all $k \in \mathbb{R}$, $u_\infty^i + k$ are solutions as well of the ergodic problem, proving the uniqueness of c .

5 Proof of Theorem 1.1 and Study of the Large Time Behavior

5.1 Proof of Theorem 1.1

Once one noticed that $u_1(x, t) = t \|f\|_{L^\infty} + \|g\|_{L^\infty} + \|u_0\|_{L^\infty}$ and $u_2(x, t) = -t \|f\|_{L^\infty} - \|g\|_{L^\infty} - \|u_0\|_{L^\infty}$ are respectively super-solution and subsolution of (1.1)–(1.3), the existence and uniqueness of a continuous global solution can be obtained by Perron's method, combining classical arguments of [12] (see also [18]), the version up to the boundary of Da Lio [13] and the Strong Comparison Result of the Proposition 2.2 on any time interval $[0, T]$.

5.2 Large Time Behavior

Let u_∞ be a bounded solution of (1.9)–(1.10). If $c \leq 0$, then u is uniformly bounded. Indeed, if $C > \|u_\infty\|_{L^\infty} + \|u_0\|_{L^\infty} + \|g\|_{L^\infty}$, then $u_\infty - C$ is a subsolution of (1.1)–(1.3). On the other hand, if \bar{x} is a point far enough from Ω , then $|x - \bar{x}|^2$ is a super-solution. To see this, it suffices to take \bar{x} such that $B(\bar{x}, R) \cap \bar{\Omega} = \emptyset$ with $R > \max(1, (\|f\|_{L^\infty} + (p-1))^{\frac{1}{q-p+2}})$.

Hence applying the Strong Comparison Result, we have

$$u_\infty(x) - C \leq u(x, t) \leq |x - \bar{x}|^2 + C \quad \text{on } \bar{\Omega} \times (0, +\infty),$$

and therefore

$$\lim_{t \rightarrow \infty} \frac{u(x, t)}{t} = 0.$$

If $c > 0$, then $u_\infty - ct + C$ is a supersolution of (1.1)–(1.3) with state constraint condition on $\partial\Omega$. On the other hand, $u_\infty - ct - C$ is a subsolution of (1.1)–(1.3) which is below u_0 at $t = 0$ and below g on $\partial\Omega$. Applying the Strong Comparison Result, we have

$$-ct + u_\infty - C \leq u(x, t) \leq u_\infty - ct + C \quad \text{on } \bar{\Omega} \times (0, +\infty).$$

The result follows by dividing by t and then letting $t \rightarrow +\infty$.

Appendice: A General Strong Comparison Result

A: Properties of the Regularization by Sup-convolution of Viscosity Subsolutions

To circumvent the lack of smoothness of the viscosity subsolution u , we consider instead the more regular time sup-convolution u^α . Such regularization was first introduced by Lasry and Lions [21], and for $0 < \alpha \leq 1$ and u a bounded usc, viscosity subsolution is defined by

$$u^\alpha(x, t) = \sup_{s \geq 0} \left\{ u(x, s) - \frac{|t - s|^2}{\alpha^2} \right\}. \quad (5.1)$$

We have the following useful properties on u^α .

Proposition 5.1. *If u is a bounded usc viscosity subsolution u of (1.1)–(1.3), the following properties are true*

i) *Set $K = \sqrt{2\|u\|_{L^\infty}}$. Then up to $o_\alpha(1)$, u^α is an usc viscosity subsolution of (1.1)–(1.3) on $\Omega \times (K\alpha, T - K\alpha)$. Moreover u^α is locally Lipschitz w.r.t to the time variable and*

$$\|u_t^\alpha\|_{L^\infty} \leq \frac{2K}{\alpha}.$$

ii) *We have $u^\alpha(x, K\alpha) \leq u_0(x) + o_\alpha(1)$ in $\overline{\Omega}$ and $u^\alpha(x, t) \leq g(x, t) + o_\alpha(1)$ on $\partial\Omega \times (K\alpha, T - K\alpha)$.*

iii) *u^α is Hölder continuous w.r.t the space variable x on $\overline{\Omega}$ uniformly w.r.t the time for $t > K\alpha$.*

Proof. Since $u(x)$ is bounded, the supremum in (5.1) is attained at some point $s^*(t)$ which belong to the interval $(t - K\alpha, t + K\alpha)$. Let $\varphi \in C^2(\overline{\Omega} \times [K\alpha, T - K\alpha])$ and assume that $u^\alpha - \varphi$ has a local maximum at $(x_0, t_0) \in \Omega \times (K\alpha, T - K\alpha)$. Denote by $s^*(t_0)$ a point such that $u^\alpha(x_0, t_0) = u(x_0, s^*(t_0)) - \frac{|t_0 - s^*(t_0)|^2}{\alpha^2}$, then the function

$$\tau \mapsto u(x, \tau) - \varphi(x, \tau - s^*(t_0) + t_0)$$

reaches a local maximum at $(x_0, s^*(t_0))$. Recalling that u is a viscosity subsolution of (1.1), we get by definition

$$\varphi_t(x_0, t_0) - \operatorname{div}(|D\varphi|^{p-2}D\varphi(x_0, t_0)) + |D\varphi(x_0, t_0)|^q \leq f(x_0, s^*(t_0)) \leq f(x_0, t_0) + o_\alpha(1),$$

by using the uniform continuity of f on $\overline{\Omega} \times [0, T]$.

Next, let $h > 0$ small enough, then

$$\begin{aligned} u^\alpha(x, t \pm h) - u^\alpha(x, t) &\geq u(x, s^*(t)) - \frac{|t \pm h - s^*(t)|^2}{\alpha^2} - u(x, s^*(t)) + \frac{|t - s^*(t)|^2}{\alpha^2} \\ &= - \left(\frac{h^2 \pm 2h(t - s^*(t))}{\alpha^2} \right) \geq - \left(\frac{h^2 + 2Kh\alpha}{\alpha^2} \right). \end{aligned}$$

A first estimate of u_t^α (from below) follows by dividing the previous inequality by h and sending $h \rightarrow 0$. Exchanging the role of $t + h$ and t provides the estimate from above.

The second assertion comes from the upper semi-continuity of u and the fact that $u(x, 0) \leq u_0(x)$. Indeed

$$u^\alpha(x, K\alpha) = u(x, s^*(K\alpha)) - \frac{|K\alpha - s^*(K\alpha)|^2}{\alpha^2} \leq u(x, s^*(K\alpha))$$

with $s^*(K\alpha) \rightarrow 0$ as $\alpha \rightarrow 0$. Taking the lim sup we get

$$\limsup_{\alpha \rightarrow 0} u^\alpha(x, K\alpha) \leq \limsup_{\alpha \rightarrow 0} u(x, s^*(K\alpha)) \leq u(x, 0) \leq u_0(x).$$

Similarly, we use the semi-continuity of u and Proposition 2.1 to prove that

$$u^\alpha(x, t) \leq g(x, t) + o_\alpha(1).$$

The last assertion is a consequence of Theorem 1.2 where $C = \|f\|_{L^\infty} + \|u_t^\alpha\|_{L^\infty}$. \square

Let us note that, using the lower semi-continuity of v and $v(x, 0) \geq u_0(x)$, we have $v(x, K\alpha) \geq u_0(x) - o_\alpha(1)$. Hence

$$u^\alpha(x, K\alpha) \leq v(x, K\alpha) + \omega(\alpha) \quad \text{for all } x \in \overline{\Omega}, \quad (5.2)$$

for some $\omega(\alpha)$ satisfying $\lim_{\alpha \rightarrow 0} \omega(\alpha) = 0$. \square

B: Proof of Proposition 2.2

In order to prove the SCR, we are going to show that $\tilde{u}^\alpha - v \leq \omega(\alpha)$ in $\overline{\Omega} \times [K\alpha, T - K\alpha]$. Inequality (2.10) follows by passing to the limit as $\alpha \rightarrow 0$. To do so, the continuity of u^α is a key point since it allows to use the arguments of [9, 4].

For the sake of simplicity of notations, we drop the $\tilde{\cdot}$ on \tilde{u}^α . The key idea is to compare $u_\mu^\alpha := \mu u^\alpha$ and v with $0 < \mu < 1$ close to 1 in order to take care of the difficulty due to the $|Du|^q$ term.

We argue by contradiction assuming that $M^\alpha = \max_{\overline{\Omega} \times [K\alpha, T - K\alpha]} (u^\alpha - v - \omega(\alpha)) > 0$.

If μ is sufficiently close to 1 and if $\eta_\alpha > 0$ is a constant small enough, then we have $M_{\mu, \eta}^\alpha = \max_{\overline{\Omega} \times [K\alpha, T - K\alpha]} (u_\mu^\alpha - v - \omega(\alpha) - \eta_\alpha(t - K\alpha)) > M^\alpha/2$.

We denote by (x_0, t_0) a point of $\overline{\Omega} \times [K\alpha, T - K\alpha]$ such that $M_{\mu, \eta}^\alpha = u_\mu^\alpha(x_0, t_0) - v(x_0, t_0) - \omega(\alpha) - \eta_\alpha(t_0 - K\alpha)$. The existence of (x_0, t_0) is guaranteed by the upper and lower semi-continuity of u^α and v respectively (we drop the dependence of (x_0, t_0) on η_α , α and μ for the sake of simplicity of notations). Since $M_{\mu, \eta}^\alpha > 0$, we necessarily have $t_0 > K\alpha$ in view of (5.2). By the Maximum Principle of the "Users guide" [12], we have $(x_0, t_0) \in \partial\Omega \times (K\alpha, T - K\alpha)$.

Next, using the regularity of the boundary, we can find a C^2 -function $\xi : \mathbb{R}^N \rightarrow \mathbb{R}^N$ which is equal to $n = -Dd$ in a neighborhood of $\partial\Omega$. Now we consider the auxiliary function $\Phi_\varepsilon : \bar{\Omega} \times \bar{\Omega} \times [K\alpha, T - K\alpha] \times [K\alpha, T - K\alpha] \rightarrow \mathbb{R}$ defined by

$$\begin{aligned} \Phi_\varepsilon(z, w, t, s) &= u_\mu^\alpha(z, t) - v(w, s) - \omega(\alpha) - \eta_\alpha(t - K\alpha) - \left| \frac{z - w}{\varepsilon} - \chi\left(\frac{z + w}{2}\right) \right|^4 \\ &\quad + \frac{|t - s|^2}{\varepsilon^2} \end{aligned}$$

Let $(\bar{z}, \bar{w}, \bar{t}, \bar{s})$ be a global maximum point of Φ_ε on $\bar{\Omega} \times \bar{\Omega} \times [K\alpha, T - K\alpha] \times [K\alpha, T - K\alpha]$. For notational simplicity we drop again the dependance of $(\bar{z}, \bar{w}, \bar{t}, \bar{s})$ on ε, μ and η . Using the inequality $\Phi_\varepsilon(\bar{z}, \bar{w}, \bar{t}, \bar{s}) \geq \Phi_\varepsilon(x_0, x_0, t_0, t_0)$ and the boundedness of u_μ^α, v and χ , we have

$$\left| \frac{\bar{z} - \bar{w}}{\varepsilon} \right| \leq C, \quad \left| \frac{\bar{t} - \bar{s}}{\varepsilon} \right| \leq C,$$

for some constant $C > 0$ depending on $\|u\|_{L^\infty}, \|v\|_{L^\infty}$ and α . By the compactness of $\bar{\Omega} \times [K\alpha, T - K\alpha]$, we can assume that $(\bar{z}, \bar{t}), (\bar{w}, \bar{s})$ converge to $(\tilde{x}, \tilde{t}) \in \bar{\Omega} \times [K\alpha, T - K\alpha]$. Moreover, using the continuity of u^α , we have

$$\Phi_\varepsilon(\bar{z}, \bar{w}, \bar{t}, \bar{s}) \geq \Phi_\varepsilon(x_0 - \varepsilon\xi(x_0), x_0, t_0, t_0) = M_{\mu, \eta}^\alpha - o_\varepsilon(1), \quad \text{as } \varepsilon \rightarrow 0,$$

and hence

$$\liminf_{\varepsilon \rightarrow 0} \Phi_\varepsilon(\bar{z}, \bar{w}, \bar{t}, \bar{s}) \geq M_{\mu, \eta}^\alpha. \quad (5.3)$$

On the other hand, we have also

$$\begin{aligned} \limsup_{\varepsilon \rightarrow 0} \Phi_\varepsilon(\bar{z}, \bar{w}, \bar{t}, \bar{s}) &\leq \limsup_{\varepsilon \rightarrow 0} (u_\mu^\alpha(\bar{z}, \bar{t}) - v(\bar{w}, \bar{s}) - \eta_\alpha(\bar{t} - K\alpha) - \omega(\alpha)) \\ &\quad - \liminf_{\varepsilon \rightarrow 0} \left| \frac{\bar{z} - \bar{w}}{\varepsilon} - \chi\left(\frac{\bar{z} + \bar{w}}{2}\right) \right|^4 \\ &\quad - \liminf_{\varepsilon \rightarrow 0} \frac{|\bar{t} - \bar{s}|^2}{\varepsilon^2} \\ &\leq M_{\mu, \eta}^\alpha. \end{aligned} \quad (5.4)$$

Therefore, combining (5.3) and (5.4) with classic arguments, we have

$$\left| \frac{\bar{z} - \bar{w}}{\varepsilon} - \chi\left(\frac{\bar{z} + \bar{w}}{2}\right) \right|^4 = o_\varepsilon(1), \quad \frac{|\bar{t} - \bar{s}|^2}{\varepsilon^2} = o_\varepsilon(1), \quad (5.5)$$

$$\begin{aligned} u_\mu^\alpha(\bar{z}, \bar{t}) - v(\bar{w}, \bar{s}) - \eta_\alpha(\bar{t} - K\alpha) - \omega(\alpha) &\rightarrow u_\mu^\alpha(\tilde{x}, \tilde{t}) - v(\tilde{x}, \tilde{t}) - \eta_\alpha(\tilde{t} - K\alpha) - \omega(\alpha) \\ &= M_{\mu, \eta}^\alpha \quad \text{as } \varepsilon \rightarrow 0. \end{aligned} \quad (5.6)$$

It follows that $u_\mu^\alpha(\bar{z}, \bar{t}) \rightarrow u_\mu^\alpha(\tilde{x}, \tilde{t})$ and $v(\bar{w}, \bar{s}) \rightarrow v(\tilde{x}, \tilde{t})$.

Now, recalling the properties of u^α and v at $t = K\alpha$, we have $\bar{t}, \bar{s} > K\alpha$ for ε small enough. Next we claim that, for ε small enough the viscosity inequalities hold for u^α and v . This is obviously the case for v if $\bar{w} \in \Omega$. If on the contrary $\bar{w} \in \partial\Omega$, then we necessarily have $v(\bar{w}, \bar{s}) < g(\bar{w}, \bar{s})$. Indeed if $\bar{w} \in \partial\Omega$ then $\tilde{x} \in \partial\Omega$. Since there is no loss of boundary conditions for subsolution s as clearly specified in Proposition 2.1, we have

$$\mu u^\alpha(\tilde{x}, \tilde{t}) \leq \mu(g(\tilde{x}, \tilde{t}) + \omega(\alpha)).$$

Using that $M_{\mu, \eta}^\alpha > 0$, we cannot have $v(\tilde{x}, \tilde{t}) \geq g(\tilde{x}, \tilde{t})$ since we would then have

$$\frac{M^\alpha}{2} \leq M_{\mu, \eta}^\alpha = \mu u^\alpha(\tilde{x}, \tilde{t}) - v(\tilde{x}, \tilde{t}) - \eta_\alpha(\tilde{t} - K\alpha) - \omega(\alpha) \leq (\mu - 1)(g(\tilde{x}, \tilde{t}) + \omega(\alpha)),$$

a contradiction by sending $\mu \rightarrow 1$.

It follows that, if $\tilde{x} \in \partial\Omega$, then we have necessarily that

$$v(\tilde{x}, \tilde{t}) < g(\tilde{x}, \tilde{t}) \quad \text{and} \quad \mu u^\alpha(\tilde{x}, \tilde{t}) \leq \mu(g(\tilde{x}, \tilde{t}) + \omega(\alpha)). \quad (5.7)$$

Hence, using that $v(\bar{w}, \bar{s}) \rightarrow v(\tilde{x}, \tilde{t}) < g(\tilde{x}, \tilde{t})$, we deduce that if $\bar{w} \in \partial\Omega$, then $v(\bar{w}, \bar{s}) < g(\bar{w}, \bar{s})$ for ε small enough and the viscosity inequality holds also in this case.

On the other hand, from (5.5) we get that

$$\bar{z} = \bar{w} + \varepsilon \chi \left(\frac{\bar{z} + \bar{w}}{2} \right) + o_\varepsilon(1), \quad (5.8)$$

which implies by the smoothness of the domain and the properties of χ that \bar{z} lies in Ω for ε small enough and hence the viscosity inequality for u_μ^α holds too.

Next, we notice that u_μ^α satisfies

$$\frac{1}{\mu}(u_\mu^\alpha)_t - \frac{1}{\mu^{p-1}} \operatorname{div}(|Du_\mu^\alpha|^{p-2} Du_\mu^\alpha) + \frac{1}{\mu^q} |Du_\mu^\alpha|^q \leq f + o_\alpha(1) \quad \text{in } \Omega \times (K\alpha, T - K\alpha),$$

and we can also re-write it as

$$\begin{aligned} \mu^{p-2}(u_\mu^\alpha)_t - |Du_\mu^\alpha|^{p-2} \left\{ \Delta u_\mu^\alpha + (p-2)(D^2 u_\mu^\alpha \widehat{Du}_\mu^\alpha, \widehat{Du}_\mu^\alpha) - \mu^{p-1-q} |Du_\mu^\alpha|^{q-p+2} \right\} \\ \leq \mu^{p-1}(f + o_\alpha(1)), \end{aligned}$$

where $\widehat{\xi} = \frac{\xi}{|\xi|}$ for $\xi \neq 0$ and $\widehat{\xi} = 0$ if $\xi \equiv 0$.

The Jensen-Ishii's Lemma [12] ensures the existence of $X, Y \in \mathcal{S}^N$, $a, b \in \mathbb{R}$, $q_1, q_2 \in \mathbb{R}^N$ such that

$$(a, q_1, X) \in \overline{\mathcal{P}}^{2,1,+} u_\mu^\alpha(\bar{z}, \bar{t}), \quad (b, q_2, Y) \in \overline{\mathcal{P}}^{2,1,-} v(\bar{w}, \bar{s}), \quad (5.9)$$

$$a - b \geq \eta_\alpha > 0, \quad |q_1 - q_2| \leq C\varepsilon(|q_1| \wedge |q_2|), \quad (5.10)$$

$$-\frac{o(1)}{\varepsilon^2} I_{2N} \leq \begin{pmatrix} X & 0 \\ 0 & -Y \end{pmatrix} \leq \frac{o(1)}{\varepsilon^2} \begin{pmatrix} I_N & -I_N \\ -I_N & I_N \end{pmatrix} + o(1)I_{2N}. \quad (5.11)$$

where $|q_1| \wedge |q_2|$ denotes the minimum of $|q_1|$ and $|q_2|$. Indeed, for (5.10), we remark that

$$q_1 = \left(I_N - \frac{\varepsilon}{2} D\chi \left(\frac{\bar{z} + \bar{w}}{2} \right) \right) q \quad \text{and} \quad q_2 = \left(I_N + \frac{\varepsilon}{2} D\chi \left(\frac{\bar{z} + \bar{w}}{2} \right) \right) q,$$

with

$$q = \frac{4}{\varepsilon} \left| \frac{z-w}{\varepsilon} - \chi \left(\frac{z+w}{2} \right) \right|^2 \left(\frac{z-w}{\varepsilon} - \chi \left(\frac{z+w}{2} \right) \right),$$

and (5.10) is an easy consequence of the boundedness of $D\chi$.

Moreover the viscosity inequalities for u_μ^α and v read

$$\begin{aligned} \mu^{p-2}a - |q_1|^{p-2} \{ \text{tr}([Id + (p-2)(\hat{q}_1 \otimes \hat{q}_1)]X) - \mu^{p-1-q}|q_1|^{q-p+2} \} \\ \leq \mu^{p-1}(f(\bar{z}, \bar{t}) + o_\alpha(1)), \end{aligned} \quad (5.12)$$

$$b - |q_2|^{p-2} \{ \text{tr}([Id + (p-2)(\hat{q}_2 \otimes \hat{q}_2)]Y) - |q_2|^{q-p+2} \} \geq f(\bar{w}, \bar{s}). \quad (5.13)$$

In the sequel we fix $\eta_\alpha > 2o_\alpha(1)$ (recall that the $o_\alpha(1)$ comes from the sup-convolution procedure and is fixed, therefore we can choose in such a way η_α). Since we may have a singularity at $q_1 = 0$ or $q_2 = 0$, we have to consider separately three cases. First we assume that there exists a constant $\gamma > 0$ such that

$$|q_1|, |q_2| \geq \gamma.$$

In this case the matrix $A(\xi) = Id + (p-2)(\hat{\xi} \otimes \hat{\xi})$ is positive definite, so that its matrix square root σ exists and satisfies

$$|\sigma(\xi_1) - \sigma(\xi_2)| \leq c \frac{|\xi_1 - \xi_2|}{|\xi_1| \wedge |\xi_2|}.$$

Combining (5.10) with the fact that (5.11) implies that $X \leq Y + o_\varepsilon(1)$, we have

$$\text{tr}(A(q_1)X) - \text{tr}(A(q_2)Y) \leq \frac{o_\varepsilon(1)}{\varepsilon^2} |\sigma(q_1) - \sigma(q_2)|^2 + o_\varepsilon(1) \leq o_\varepsilon(1), \quad (5.14)$$

$$\begin{aligned} |q_2|^{q-p+2} - \mu^{p-1-q}|q_1|^{q-p+2} &= |q_2|^{q-p+2} - |q_1|^{q-p+2} + (1 - \mu^{p-1-q})|q_1|^{q-p+2} \\ &\leq (q-p+2)|q_1|^{q-p+1}|q_2 - q_1| + (1 - \mu^{p-1-q})|q_1|^{q-p+2} \\ &\leq o(\varepsilon)|q_1|^{q-p+2} + (1 - \mu^{p-1-q})|q_1|^{q-p+2}. \end{aligned} \quad (5.15)$$

Multiplying (5.12) by $\frac{|q_2|^{p-2}}{|q_1|^{p-2}}$ which is of order $1 + O(\varepsilon)$ and subtracting from it (5.13), we have

$$\begin{aligned}
(1 + O(\varepsilon))\mu^{p-2}a - b &\leq |q_2|^{p-2} \{o_\varepsilon(1) + o(\varepsilon)|q_1|^{q-p+2} + (1 - \mu^{p-1-q})|q_1|^{q-p+2}\} \\
&+ \mu^{p-1}(1 + O(\varepsilon))(f(\bar{z}, \bar{t}) + o_\alpha(1)) - f(\bar{w}, \bar{s}). \tag{5.16}
\end{aligned}$$

At this point, we recall that the Lipschitz continuity of u^α implies that $|a| \leq \frac{2\mu K}{\alpha}$. On the other hand we remark that, since $1 - \mu^{p-1-q} < 0$, for fixed μ the term $o(\varepsilon)|q_1|^{q-p+2}$ is controlled by the $(1 - \mu^{p-1-q})|q_1|^{q-p+2}$ term.

Now we are going to let $\varepsilon \rightarrow 0$: if we assume that q_1, q_2 (which depend on ε) are bounded, we may assume that they converge (we still denote their limits as q_1, q_2 respectively). For μ close enough to 1, we get as $\varepsilon \rightarrow 0$

$$0 < \eta_\alpha/2 \leq \mu^{p-2}a - b \leq (\mu^{p-1} - 1)f(\tilde{x}, \tilde{t}) + \mu^{p-1}o_\alpha(1) + |q_2|^{p-2}(1 - \mu^{p-1-q})|q_1|^{q-p+2}$$

Recalling that $\eta_\alpha > 2o_\alpha(1)$, we get a contradiction when $\mu \rightarrow 1$ since the last term of the right-hand side is negative. Of course, we get the same contradiction if (at least for some subsequence) q_1 or $q_2 \rightarrow \infty$.

If $q_1, q_2 \neq 0$ but $q_1 \rightarrow 0, q_2 \rightarrow 0$ then, noticing that $\frac{|q_2|^{p-2}}{|q_1|^{p-2}}$ is still of order $1 + O(\varepsilon)$, we can pass to the limit $\varepsilon \rightarrow 0$ in the same way and obtain

$$0 < \eta_\alpha/2 \leq \mu^{p-2}a - b \leq (\mu^{p-1} - 1)f(\tilde{x}, \tilde{t}) + \mu^{p-1}o_\alpha(1),$$

and we also get a contradiction.

If $q_1 = 0$ or $q_2 = 0$, then necessarily $q_1 = q_2 = 0$ and, by subtracting (5.13) from (5.12), we have

$$\eta_\alpha/2 \leq \mu^{p-2}a - b \leq \mu^{p-1}(f(\bar{z}, \bar{t}) + o_\alpha(1)) - f(\bar{w}, \bar{s}).$$

We get a contradiction when $\varepsilon \rightarrow 0$.

In all cases fixing $\eta_\alpha > 2o_\alpha(1)$ we get a contradiction for ε small enough and μ close to 1 and the conclusion follows.

C: Proof of Proposition 2.3

The proof of (i) and (ii) are very similar. Indeed we know by Theorem 1.2 that subsolutions of (2.2) are Hölder continuous, so we are always in a case where the subsolution or the supersolution are continuous. We will only give details of the proof for the p -Laplacian operator in the case where $\lambda > 0$ and $v \in C(\bar{\Omega})$. The other cases are an easy adaptation (the equation (2.1) is even easier to study). Since v is assumed to be continuous, we follow the proof of [9] with the same trick as before in order to take care of the strong growth of the gradient term. We argue by contradiction assuming that $M = \max_{\bar{\Omega}}(u - v) > 0$. If μ is sufficiently close to 1, then we have $M_\mu = \max_{\bar{\Omega}}(u_\mu - v) > M/2 > 0$. Since u is usc and

v is continuous this maximum is achieved at x_0 . We may assume that $x_0 \in \partial\Omega$. We drop the dependence of x_0 on μ .

Next, using the regularity of the boundary, we can find a C^2 -function $\xi : \mathbb{R}^N \rightarrow \mathbb{R}^N$ which is equal to n in a neighborhood of $\partial\Omega$. Now we consider the test function $\Phi_\varepsilon : \overline{\Omega} \times \overline{\Omega} \rightarrow \mathbb{R}$ defined by

$$\Phi_\varepsilon(z, w) = \mu u(z) - u(w) + \left| \frac{z - w}{\varepsilon} + \xi \left(\frac{z + w}{2} \right) \right|^4.$$

Let (x, y) be a global maximum point of Φ_ε on $\overline{\Omega} \times \overline{\Omega}$. For notational simplicity we drop the dependence of x and y on ε . and μ . Using the boundedness of u and v , it is clear that $x - y = O(\varepsilon)$ and it follows that, along a subsequence, $x, y \rightarrow \bar{x} \in \overline{\Omega}$. Since u is lsc and v is continuous, we get that $\limsup_{\varepsilon \rightarrow 0} \Phi_\varepsilon(x, y) \leq M_\mu$.

On the other hand we have, using the **continuity** of v , we have

$$\Phi_\varepsilon(x, y) \geq \Phi_\varepsilon(x_0, x_0 - \varepsilon\xi(x_0)) \geq M_\mu + O(\varepsilon).$$

It follows that $\liminf_{\varepsilon \rightarrow 0} \Phi_\varepsilon \leq M_\mu$. Hence we get that

$$\Phi_\varepsilon(x, y) \rightarrow M_\mu \quad \text{as } \varepsilon \rightarrow 0. \quad (5.17)$$

Standard arguments allow us to deduce from (5.17) that

$$\left| \frac{x - y}{\varepsilon} + \xi \left(\frac{x + y}{2} \right) \right|^4 = o_\varepsilon(1), \quad (5.18)$$

$$\mu u(x) - v(y) \rightarrow \mu u(\bar{x}) - v(\bar{x}) = M_\mu \quad \text{as } \varepsilon \rightarrow 0. \quad (5.19)$$

Next we claim that, for ε small enough the viscosity inequalities hold for u and v . This is obviously the case if $y \in \Omega$. Using Proposition 2.1 and arguing similarly as in the previous proof, we get that, if $y \in \partial\Omega$ than $v(y) \leq \tilde{g}(y)$ and the viscosity inequality holds also in this case. On the other hand, using (5.19), we get that

$$x = y - \varepsilon\xi(y) + o_\varepsilon(1) \quad (5.20)$$

which implies by the smoothness of the domain and the properties of ξ that x lies in Ω for ε small enough and hence the viscosity inequality holds for μu .

Using the same arguments as the previous proof, we get that the elements $(q_1, X) \in \overline{\mathcal{F}}^{2,+} u_\mu$ and $(q_2, Y) \in \overline{\mathcal{F}}^{2,-} v$ given by the Jensen-Ishii's Lemma satisfy

$$\begin{aligned} \lambda(1 + O(\varepsilon))\mu^{p-2}u_\mu(x) - v(y) &\leq |q_2|^{p-2} \{o_\varepsilon(1) + o(\varepsilon)|q_1|^{q-p+2} + (1 - \mu^{p-1-q})|q_1|^{q-p+2}\} \\ &\quad + \mu^{p-1}(1 + O(\varepsilon))\tilde{f}(x) - \tilde{f}(y). \end{aligned} \quad (5.21)$$

Letting $\varepsilon \rightarrow 0$ and then $\mu \rightarrow 1$, we get a contradiction. It follows that $\tilde{u} \leq v$ on $\overline{\Omega}$.

References

- [1] Laurent Amour and Matania Ben-Artzi. Global existence and decay for viscous hamilton-jacobi equations. *Nonlinear Analysis: Theory, Methods & Applications*, 31(5):621–628, 1998.
- [2] Amal Attouchi. Well-posedness and gradient blow-up estimate near the boundary for a hamilton-jacobi equation with degenerate diffusion. *Journal of Differential Equations*, 253:2474–2492, 2012.
- [3] Martino Bardi and Francesca Da Lio. On the strong maximum principle for fully nonlinear degenerate elliptic equations. *Arch. Math. (Basel)*, 73(4):276–285, 1999.
- [4] Guy Barles. A short proof of the $C^{0,\alpha}$ -regularity of viscosity subsolutions for superquadratic viscous Hamilton-Jacobi equations and applications. *Nonlinear Anal.*, 73(1):31–47, 2010.
- [5] Guy Barles and Francesca Da Lio. Remarks on the dirichlet and state-constraint problems for quasilinear parabolic equations. *Advances in Differential Equations*, 8(8):897–922, 2003.
- [6] Guy Barles and Francesca Da Lio. On the generalized Dirichlet problem for viscous Hamilton-Jacobi equations. *J. Math. Pures Appl. (9)*, 83(1):53–75, 2004.
- [7] Guy Barles, Philippe Laurençot, and Christian Stinner. Convergence to steady states for radially symmetric solutions to a quasilinear degenerate diffusive Hamilton-Jacobi equation. *Asymptot. Anal.*, 67(3-4):229–250, 2010.
- [8] Guy Barles, Alessio Porretta, and Thierry Tabet Tchamba. On the large time behavior of solutions of the dirichlet problem for subquadratic viscous hamilton-jacobi equations. *J. Math. Pures Appl. (9)*, 94(5):497–519, 2010.
- [9] Guy Barles, Elisabeth Rouy, and Panagiotis E. Souganidis. Remarks on the Dirichlet problem for quasilinear elliptic and parabolic equations. In *Stochastic analysis, control, optimization and applications*, Systems Control Found. Appl., pages 209–222. Birkhäuser Boston, Boston, MA, 1999.
- [10] Matania Ben-Artzi, Philippe Souplet, and Fred B Weissler. The local theory for viscous hamilton-jacobi equations in lebesgue spaces. *Journal de mathématiques pures et appliquées*, 81(4):343–378, 2002.
- [11] I. Capuzzo Dolcetta, F. Leoni, and A. Porretta. Hölder estimates for degenerate elliptic equations with coercive Hamiltonians. *Trans. Amer. Math. Soc.*, 362(9):4511–4536, 2010.

- [12] Michael G. Crandall, Hitoshi Ishii, and Pierre-Louis Lions. User's guide to viscosity solutions of second order partial differential equations. *Bull. Amer. Math. Soc. (N.S.)*, 27(1):1–67, 1992.
- [13] Francesca Da Lio. Comparison results for quasilinear equations in annular domains and applications. *Comm. Partial Differential Equations*, 27:283–323, 2002.
- [14] Francesca Da Lio. Comparison results for quasilinear equations in annular domains and applications. *Comm. Partial Differential Equations*, 27(1-2):283–323, 2002.
- [15] Francesca Da Lio. Remarks on the strong maximum principle for viscosity solutions to fully nonlinear parabolic equations. *Commun. Pure Appl. Anal.*, 3(3):395–415, 2004.
- [16] Marek Fila and Gary M. Lieberman. Derivative blow-up and beyond for quasilinear parabolic equations. *Differential Integral Equations*, 7(3-4):811–821, 1994.
- [17] Marek Fila and Michael Winkler. rate of convergence to a singular steady state of a supercritical parabolic equation. *J. Evol. Equ.*, 8(4):673–692, 2008.
- [18] H. Ishii. Perron's method for hamilton-jacobi equations. *Duke Mathematical Journal*, 55(2):369–384, 1987.
- [19] Bernard Kawohl and Nikolai Kutev. A study on gradient blow up for viscosity solutions of fully nonlinear, uniformly elliptic equations. *Acta Math. Sci. Ser. B Engl. Ed.*, 32(1):15–40, 2012.
- [20] J.-M. Lasry and P.-L. Lions. Nonlinear elliptic equations with singular boundary conditions and stochastic control with state constraints. I. The model problem. *Math. Ann.*, 283(4):583–630, 1989.
- [21] Jean-Michel Lasry and Pierre-Lions Lions. A remark on regularization in hilbert spaces. *Israel Journal of Mathematics*, 55(3):257–266, 1986.
- [22] Philippe Laurençot. Non-diffusive large time behavior for a degenerate viscous Hamilton-Jacobi equation. *Comm. Partial Differential Equations*, 34(1-3):281–304, 2009.
- [23] Philippe Laurençot and Christian Stinner. Convergence to separate variables solutions for a degenerate parabolic equation with gradient source. *J. Dynam. Differential Equations*, 24(1):29–49, 2012.
- [24] Pavol Quittner and Philippe Souplet. *Superlinear parabolic problems: Blow-up, global existence and steady states*. Springer, 2007.
- [25] Ph. Souplet and Q.S Zhang. Global solutions of inhomogeneous hamilton-jacobi equations. *J. Anal. Math.*, 99:355–396, 2006.

- [26] Thierry Tabet Tchamba. Large time behavior of solutions of viscous Hamilton-Jacobi equations with superquadratic Hamiltonian. *Asymptot. Anal.*, 66(3-4):161–186, 2010.