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Escape Rates and Singular Limiting Distributions for Intermittent Maps with Holes

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Abstract

We study the escape dynamics in the presence of a hole of a standard family of intermittent maps of the unit interval with neutral fixed point at the origin (and finite absolutely continuous invariant measure). Provided that the hole (is a cylinder that) does not contain any neighborhood of the origin, the surviving volume is shown to decay at polynomial speed with time. The associated polynomial escape rate depends on the density of the initial distribution, more precisely, on its behavior in the vicinity of the origin. Moreover, the associated normalized push forward measures are proved to converge to the point mass supported at the origin, in sharp contrast to systems with exponential escape rate. Finally, a similar result is obtained for more general systems with subexponential escape rates; namely that the Cesaro limit of normalized push forward measures is generally singular, invariant and supported on the asymptotic survivor set.

November 6, 2013.

1 Introduction and setting

The study of systems with holes finds its origin in the study of Markov chains with absorbing states [25, 34, 35, 36] and was introduced in deterministic dynamical systems by Pianigiani and Yorke [32]. It has focused on the establishment of escape rates and on the existence of conditionally invariant measures which describe the asymptotic distribution of mass conditioned on non-escape.

Since conditionally invariant measures are badly non-unique [17], physically relevant measures are usually characterized as the limit of normalized push forward iterates of a reference measure (usually Lebesgue). Such limiting distributions are typically eigenmeasures with maximal eigenvalue of the corresponding transfer operator defined on an appropriate function space. The maximal eigenvalue itself gives the exponential rate of escape of mass from the system. These limiting conditionally invariant measures have properties analogous to Sinai-Ruelle-Bowen (SRB) measures for

the corresponding closed system. Under reasonable assumptions, they converge to the SRB measure as the size of the hole tends to zero, and this establishes stability under perturbations in the form of holes.

Examples begin with open systems admitting a finite Markov partition: expanding maps in \mathbb{R}^n [14, 32], Smale horseshoes [6], Anosov diffeomorphisms [7, 8], and some unimodal maps [27]. Subsequent attempts to substitute the Markov assumption by requiring that the holes be small, have extended this analysis to Anosov diffeomorphisms with non-Markov holes [9, 10], to expanding maps of the interval [11, 15, 29], to multimodal maps satisfying a Collet-Eckmann condition [2, 16], to piecewise hyperbolic maps [18] and recently, to various classes of dispersing billiards [19, 21, 22].

The characteristic common to all these systems is that the rate of escape is exponential (the systems enjoy exponential decay of correlations before the introduction of the hole) so that the concept of conditionally invariant measure is well-defined.

Polynomial rates of escape have been studied numerically in some systems [23, 24] and via formal expansions to obtain leading order terms for the decay rate [3, 4]. However, to our knowledge, there are no analytical results regarding limiting distributions for systems with polynomial rates of escape.

The purpose of the present paper is to initiate the analysis of open systems with subexponential rates of escape. For simplicity, we consider a family of intermittent maps of the unit interval, with neutral fixed point at the origin [30]. For the hole, we take any element of a refined Markov partition for the map, not adjacent to the origin. (Of note, [26] has also considered interval maps with neutral fixed point and very specific holes which are either a neighborhood of the neutral fixed point or its complement.)

In this context, we first prove that the rates of escape must be polynomial for a large class of initial distributions, and this rate depends on the behavior of the initial distribution in the neighborhood of the origin. In particular, the polynomial rate of escape with respect to the SRB measure (before the introduction of the hole) differs from that with respect to Lebesgue measure.

In this setting, conditionally invariant measures are not physically meaningful (although plenty still exist with any desired eigenvalue [17]). We show that the limit of $\hat{T}_*^n \mu / |\hat{T}_*^n \mu|$ converges to the point mass at the neutral fixed point, again for a large class of initial distributions (including both Lebesgue and the SRB measures).

These results hold independently of the size of the hole. Thus from the point of view of the physical limit $\hat{T}_*^n \mu / |\hat{T}_*^n \mu|$, a hole of any size is always a large perturbation in the context of subexponentially mixing systems. In other words, the attracting property of the SRB measure under the action of T_*^n is unstable with respect to small leaks in the system.

Finally, we consider more general systems with subexponential rates of escape. The analysis of intermittent maps of the interval might suggest that the results are specific to this setting. Our final result Theorem 2.3 shows that this is not the case: in contrast to situations with exponential escape, in systems with slow escape, the (Cesaro) limit of $\hat{T}_*^n m / |\hat{T}_*^n m|$ will always be singular and will be supported on the survivor set of points that never escape.

1.1 Setting

We study the dynamics of the map of the unit interval $T : I \rightarrow I$ where $I = [0, 1)$ and T is defined by (see [30] and Figure 1)

$$T(x) = \begin{cases} x + 2^\gamma x^{\gamma+1} & \text{if } x \in [0, \frac{1}{2}) \\ 2x - 1 & \text{if } x \in [\frac{1}{2}, 1) \end{cases}$$

with $0 < \gamma < 1$, after the introduction of a hole H into I . In this parameter range, T preserves a

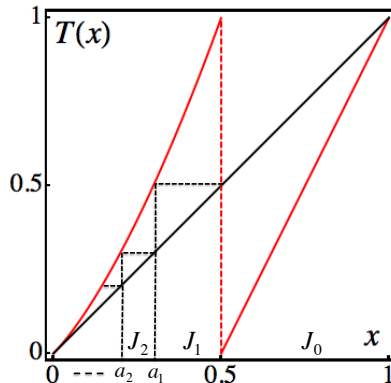


Figure 1: Graph of the map T for $\gamma = \frac{3}{4}$ (solid red branches), together with some the intervals $J_n = [a_n, a_{n-1})$.

finite invariant measure, μ_{SRB} , absolutely continuous with respect to Lebesgue.

In order to define the hole, we need to introduce the (standard) finite and countable Markov partitions of I . The finite partition is defined by $\mathcal{P} := \{J_L, J_R\}$ where $J_L = [0, \frac{1}{2})$ and $J_R = [\frac{1}{2}, 1)$. The countable partition is defined by $\mathcal{J} := \{J_n\}_{n \geq 0}$ where

$$J_n = \begin{cases} J_R & \text{if } n = 0 \\ [a_n, a_{n-1}) & \text{if } n \geq 1 \end{cases}$$

where $a_n = T_L^{-n}(\frac{1}{2})$ and T_L denotes the left branch of T . Now, given $t \geq 0$, let $\mathcal{J}^{(t)}$ be the refined partition defined as follows

$$\mathcal{J}^{(t)} := \mathcal{J} \vee \bigvee_{i=0}^t T^{-i}(\mathcal{P}).$$

The hole H is defined to be any element of $\mathcal{J}^{(\ell_H)}$ where $\ell_H \geq 0$ is arbitrary. We shall denote by $J_h \supseteq H$ with $h \geq 0$, the element of \mathcal{J} that contains H .

Of note, this assumption on the hole gives immediate access to the countable Markov partition of the open system. This is convenient for the conditioning arguments in the proof of Lemma 3.2 and for the invariance of a certain function space used in the proof of Theorem 2.2 that implies control of the structure of the singular limit. However, we believe that this cylinder assumption is purely technical and we expect our results to hold even when relaxed, although significant technical modifications will have to be made.

Define $\mathring{I} = I \setminus H$ and given $t \geq 0$, let $\mathring{I}^t = \bigcap_{i=0}^t T^{-i}(I \setminus H)$ represent those points which have not escaped by time t . We refer to $\mathring{T} := T|_{\mathring{I}^0}$ as the map with a hole and its iterates $\mathring{T}^t = T^t|_{\mathring{I}^{t-1}}$ ($t \geq 1$) describe the dynamics of the open system before escape. Notice that $\mathcal{J}^{(\ell_H)}$ is also a countable Markov partition for \mathring{T} .

One of the quantities we will be interested in studying is the rate of escape of mass from the open system. Given a measure μ on \mathring{I} , we define the *polynomial rate of escape* with respect to μ by

$$\epsilon_{\text{poly}}(\mu) = - \lim_{t \rightarrow \infty} \frac{\log \mu(\mathring{I}^t)}{\log t},$$

whenever the limit exists.

We will also study the asymptotic evolution of absolutely continuous measures that are transported under the action of \mathring{T} . Given a measure μ on \mathring{I} and $t \geq 1$, let $\mathring{T}_*^t \mu$ be the push forward measure under the action \mathring{T}^t .

Let m denote Lebesgue measure on I and given $f \in L^1(m)$, let μ_f be the absolutely continuous measure with density f . As in the standard case (without hole), we have $\mathring{T}_*^t \mu = \mu_{\mathring{\mathcal{L}}^t f}$ where $\mathring{\mathcal{L}}^t f := 1_{j_0} \mathcal{L}^t(f 1_{j_{t-1}})$ and \mathcal{L} is the transfer operator associated with T , defined by the expression

$$\mathcal{L}f(x) = \sum_{y \in T^{-1}(x)} \frac{f(y)}{DT(y)},$$

where $DT > 0$ is the (first) derivative of T . In addition, the change of variable formula implies in this case the following relation

$$|\mathring{\mathcal{L}}^t f|_1 = \mu_f(\mathring{I}^t), \quad \forall t \geq 0,$$

where $|\cdot|_1$ denotes the L^1 -norm with respect to Lebesgue measure m .

2 Statement of Results

2.1 Results for the open system \mathring{T}

Throughout this section, the hole H is fixed as in the previous section (and so are ℓ_H and h). Our first result describes a common set of escape rates for initial distributions depending on their behavior near 0. Following [37], the notation $u_t \approx v_t$ (resp. $u_t \lesssim v_t$) means there exists $C > 0$ such that $C^{-1}v_t \leq u_t \leq Cv_t$ (resp. $u_t \leq Cv_t$) for all t .

Theorem 2.1. *For any non-negative $f \in L^1(m)$ for which there exist $x_0 \in (0, 1)$ and $\alpha \in [0, 1)$ such that*

$$0 < \inf_{x \in (0, x_0)} x^\alpha f(x) \leq \sup_{x \in (0, 1)} x^\alpha f(x) < +\infty,$$

we have

$$\mu_f(\mathring{I}^t) \approx t^{-\frac{1-\alpha}{\gamma}}.$$

Consequently, the associated measure μ_f has polynomial escape rate, $\epsilon_{\text{poly}}(\mu_f) = \frac{1-\alpha}{\gamma}$.

The proof is given in Section 3.2. Of note, to obtain the lower bound on $\mu_f(\mathring{I}^t)$ is rather immediate (see relation (9)). Moreover, ergodicity of the map T implies $\mu_f(\mathring{I}^t) \xrightarrow{t \rightarrow +\infty} 0$. Thus, most of the proof consists in proving the upper bound. This part is inspired by the proof in [37] of the speed of convergence to the equilibrium measure.

Theorem 2.1 implies in particular that the polynomial escape rate associated with Lebesgue measure is given by $\epsilon_{\text{poly}}(m) = \frac{1}{\gamma}$. Interestingly, since $\frac{d\mu_{\text{SRB}}}{dm}(x) \approx x^{-\gamma}$ for x near 0 [37], this rate differs from the one associated with the SRB measure, $\epsilon_{\text{poly}}(\mu_{\text{SRB}}) = \frac{1-\gamma}{\gamma}$.

That the escape rate is polynomial depends on the assumptions on both H and the initial density f . Indeed, if the hole included a neighborhood of the neutral fixed point 0, then the corresponding open system \hat{T} would be uniformly expanding and the escape rate would be exponential, for any initial density $f \in L^1(m)$, see [26, 32] for the Markov case and any of [11, 15, 29] for the non-Markov case. (Obviously, such holes do not belong to $\mathcal{J}^{(t)}$ for any $t \geq 0$.)

Alternatively, $H = [\frac{1}{2}, \frac{1}{2} + d_{\ell_H}) \in \mathcal{J}^{(\ell_H)}$ for some $d_{\ell_H} > 0$ can also create exponential behavior for some initial densities f . Indeed, the map $T|_{[2d_{\ell_H}, 1)}$ is uniformly expanding and no point in $[2d_{\ell_H}, 1)$ can enter the interval $[0, 2d_{\ell_H})$ without first falling into H . Hence, the measure associated with any smooth density f satisfying $f|_{[0, 2d_{\ell_H})} \equiv 0$ must experience an exponential rate of escape in this case. (Such densities do not satisfy the assumption of Theorem 2.1.)

However, for any hole not blocking repeated passes through a neighborhood of 0, $\hat{\mathcal{L}}^t f$ will eventually be positive in a neighborhood of 0 (and bounded) for any ‘typical’ smooth density f ; hence Theorem 2.1 implies that the associated measure will experience a polynomial escape rate $\frac{1}{\gamma}$.

Our next result describes the limiting behavior of the sequence $\{\frac{\hat{T}_*^t \mu_f}{\mu_f(\hat{I}^t)}\}_{t \in \mathbb{N}}$ of push forward probability measures, for initial densities f that are log-Hölder continuous on elements of the partition $\mathcal{J}^{(\ell_H)}$. To be precise, let $C^0(\mathcal{J}^{(\ell_H)})$ denote the set of functions defined in the interior of I and continuous on each element of $\mathcal{J}^{(\ell_H)}$. Given $f \in C^0(\mathcal{J}^{(\ell_H)})$, $f \geq 0$, $p \in \mathbb{R}^+$ and $J \in \mathcal{J}^{(\ell_H)}$, define the quantity $H_J^p(f)$ as follows

$$H_J^p(f) = \begin{cases} 0 & \text{if } f \equiv 0 \text{ on } J \\ +\infty & \text{if } f(x) = 0 < f(y) \text{ for some } x, y \in J \\ \sup_{x \neq y \in J} \frac{\log f(x) - \log f(y)}{|x - y|^p} & \text{if } f > 0 \text{ on } J \end{cases}$$

and let $\|f\|_p := \sup_{J \in \mathcal{J}^{(\ell_H)}} H_J^p(f)$. Consider also the set of functions,

$$\mathcal{F}_p = \{f \in C^0(\mathcal{J}^{(\ell_H)}) : f \geq 0, |f|_1 = 1 \text{ and } \|f\|_p < +\infty\}.$$

and its subset

$$\mathcal{F}_p^0 = \{f \in \mathcal{F}_p : \exists C > 0 \text{ and } n_0 \in \mathbb{N} \text{ such that } f|_{\bigcup_{n \geq n_0} J_n} \geq C\}.$$

of functions which are bounded away from zero in a neighborhood of 0 (so that their escape rate captures the behavior of the neutral fixed point).

Theorem 2.2. *Let $f \in \mathcal{F}_p^0$ for some $p \in (0, \frac{\gamma}{\gamma+1}]$. Then the sequence $\{\frac{\hat{T}_*^t \mu_f}{\mu_f(\hat{I}^t)}\}_{t \in \mathbb{N}}$ of absolutely continuous measures behaves asymptotically as follows*

$$\lim_{t \rightarrow +\infty} \frac{\hat{T}_*^t \mu_f}{\mu_f(\hat{I}^t)} = \delta_0,$$

where δ_0 denotes the point mass at 0 and the convergence is in the weak sense. Moreover, we have

$$\lim_{t \rightarrow +\infty} \frac{\mu_f(\hat{I}^{t+1})}{\mu_f(\hat{I}^t)} = 1.$$

Of note, this last expression of the theorem can be alternatively formulated as

$$\lim_{t \rightarrow +\infty} \frac{\mu_f(\mathring{I}^t \setminus \mathring{I}^{t+1})}{\mu_f(\mathring{I}^t)} = 0.$$

Theorem 2.2 applies in particular to Lebesgue measure, since 1 belongs to \mathcal{F}_p^0 for every $p \in \mathbb{R}^+$. (More generally, one easily checks that any positive Hölder continuous density on I with exponent p belongs to \mathcal{F}_p^0). Theorem 2.2 also applies to μ_{SRB} since the density $f_{\text{SRB}} = \frac{d\mu_{\text{SRB}}}{dm}$ belongs to \mathcal{F}_p^0 for every $p \in (0, \frac{\gamma}{\gamma+1}]$ as does any normalized density $f/|f|_1$, where $f(x) = x^{-\alpha}$ for some $\alpha \in (0, 1)$ (see Lemma 3.7 in Section 3.3).

As mentioned in the introduction, this theorem implies that arbitrarily small holes in systems with polynomial rates of escape can act as large perturbations from the point of view of the physical limit $\frac{\mathring{T}_*^t \mu_f}{\mu_f(\mathring{I}^t)}$.

As part of the broader analysis of the dependence of the escape rate on the size and location of the hole [1, 5, 20, 28], one can also think of stability of the open system in terms of the continuity of limiting invariant and conditionally invariant objects as we ‘slide’ a hole continuously in I . If we take $H_\varepsilon = (\varepsilon, \varepsilon + \delta)$ for some small $\delta > 0$, then for each $\varepsilon > 0$, the sequence $\frac{\mathring{T}_*^t m}{m(\mathring{I}^t)}$ tends to δ_0 for large n . But when $\varepsilon = 0$, $H_0 = (0, \delta)$, the escape rate is exponential and the sequence $\frac{\mathring{T}_*^t m}{m(\mathring{I}^t)}$ tends to a conditionally invariant measure that is absolutely continuous with respect to Lebesgue. Thus once again, the hole is seen as a discontinuous perturbation.

2.2 General open systems: Consequence of a subexponential escape rate

Convergence of $\frac{\mathring{T}_*^t \mu_f}{\mu_f(\mathring{I}^t)}$ to a singular limit as in Theorem 2.2 is not limited to the map \mathring{T} above. Indeed, as we show now, this phenomenon occurs very generally when the rate of escape is subexponential.

To see this, let X be a compact, separable metric space and let $T : X \rightarrow X$ be now an arbitrary Borel measurable map. Assume there exists a probability measure μ with respect to which T is nonsingular (but not necessarily invariant). This will be our reference measure.

Let an open set $H \subset X$ be the hole and let $\mathring{X}^t = \bigcap_{i=0}^t T^{-i}(X \setminus H)$ denote the survivor set up until time $t \in \mathbb{N} \cup \{+\infty\}$. Our main assumption on the open system is that μ -almost every point escapes and that the escape rate is subexponential, i.e. we assume

$$\mu(\mathring{X}^\infty) = 0, \quad \mu(\mathring{X}^t) > 0 \text{ for all } t \geq 0 \quad \text{and} \quad \limsup_{t \rightarrow +\infty} \frac{\log \mu(\mathring{X}^t)}{t} = 0. \quad (1)$$

In particular, this includes both polynomial and stretched exponential rates of escape (and does not assume that the limit above converges).

By assumption, all push forward (probability) measures $\frac{\mathring{T}_*^t \mu}{\mu(\mathring{X}^t)}$ are nonsingular with respect to μ . Hence, the same is true for

$$\mu_t = \frac{1}{t} \sum_{k=0}^{t-1} \frac{\mathring{T}_*^k \mu}{\mu(\mathring{X}^k)}, \quad (2)$$

for all $t \geq 1$. As the next result shows, any limit point however must be singular. Convergence here is also understood in the weak sense.

Theorem 2.3. *Any limit point μ_∞ of the sequence $\{\mu_t\}_{t \in \mathbb{N}}$ is singular with respect to μ and is supported on \dot{X}^∞ . If, in addition, μ_∞ gives zero measure to the discontinuity set of \dot{T} , then μ_∞ is T -invariant.*

Interestingly, the averaging method presented here does not work so easily in the case of exponential escape (unless a priori one knows that the limit of $\frac{\dot{T}_*^t \mu}{\mu(\dot{X}^t)}$ itself exists). Indeed, in this case, the ratio of consecutive normalizations $\frac{\mu(\dot{X}^{t+1})}{\mu(\dot{X}^t)}$ does not converge to 1 and the terms appearing in the sum must be weighed to compensate for this. For an example of an averaging method in the exponential case under stronger assumptions, see [12, 13]. Theorem 2.3 is proved in Section 3.5.

3 Proofs

3.1 Preliminary estimates

In proving the theorems, we shall repeatedly use the following bounds [37]

$$a_n \approx n^{-\frac{1}{\gamma}} \quad \text{and} \quad |J_n| \approx n^{-\frac{\gamma+1}{\gamma}}. \quad (3)$$

We shall also rely on the following lemma.

Lemma 3.1. *Given $n \geq 0$ and $t \geq 1$, let $x, y \in \dot{I}^t$ lie in the same element of $\mathcal{J}^{(t+1)}$ such that $T^t(x), T^t(y) \in J_n$. Then we have*

$$(a) \quad \frac{1}{DT^t(x)} \lesssim \left(\frac{n}{n+t}\right)^{\frac{\gamma+1}{\gamma}};$$

$$(b) \quad \text{for any } p \in (0, \frac{\gamma}{\gamma+1}], \text{ one has } \left| \log \frac{DT^t(x)}{DT^t(y)} \right| \lesssim |T^t(x) - T^t(y)|^p.$$

Proof. (a) Given $i \in \{0, \dots, t-1\}$, let J_{n_i} denote the element of \mathcal{J} containing $T^i(x)$ and $T^i(y)$. Let also $B_{n_i} = 2^\gamma \gamma (\gamma+1) a_{n_i}^{\gamma-1}$ be the maximum value of $|D^2 T|$ and $M_{n_i}^{(j)}$ be the minimum value of $|DT^j|$ on J_{n_i} , respectively. We have

$$\begin{aligned} \left| \log \frac{DT^t(x)}{DT^t(y)} \right| &\leq \sum_{i=0}^{t-1} |\log DT \circ T^i(x) - \log DT \circ T^i(y)| \\ &\leq \sum_{i=0}^{t-1} \frac{B_{n_i}}{M_{n_i}^{(1)}} |T^i(x) - T^i(y)| \leq \sum_{i=0}^{t-1} B_{n_i} |T^i(x) - T^i(y)|. \end{aligned} \quad (4)$$

Following [37], we write $|T^t(x) - T^t(y)| = |DT^{t-i}(z)| |T^i(x) - T^i(y)|$ for some $z \in J_{n_i}$ and use that the expansion $DT^{t-i}(z)$ decreases as n_i increases to conclude that the last sum here is maximised for $n_i = n + t - i$, i.e. $T^i(x), T^i(y) \in J_{n+t-i}$. Using equation (3) we obtain

$$\left| \log \frac{DT^t(x)}{DT^t(y)} \right| \lesssim \sum_{i=0}^{t-1} (n+t-i)^{\frac{1-\gamma}{\gamma}} (n+t-i)^{-\frac{\gamma+1}{\gamma}} = \sum_{i=0}^{t-1} (n+t-i)^{-2} \leq n^{-1},$$

where the last inequality follows from $\sum_{i=n+1}^{\infty} i^{-2} \leq \int_n^{+\infty} x^{-2} dx$.

Now, $T^t(x)$ has no preimage in $\bigcup_{i>n+t} J_i$, so the weakest expansion occurs when $x \in J_{n+t}$. The previous distortion estimate implies the existence of $C' > 0$ such that

$$DT^t(x) \geq C' \frac{|J_n|}{|J_{n+t}|}.$$

Using equation (3) again, statement (a) easily follows.

(b) Adopting the same notation as in (a) and starting from (4), we fix $p \in (0, \frac{\gamma}{\gamma+1}]$ and write

$$\begin{aligned} \left| \log \frac{DT^t(x)}{DT^t(y)} \right| &\leq \sum_{i=0}^{t-1} B_{n_i} |T^i(x) - T^i(y)|^{1-p} \frac{|T^i(x) - T^i(y)|^p}{|T^t(x) - T^t(y)|^p} |T^t(x) - T^t(y)|^p \\ &\leq \sum_{i=0}^{t-1} \frac{B_{n_i} |J_{n_i}|^{1-p}}{(M_{n_i}^{(t-i)})^p} |T^t(x) - T^t(y)|^p. \end{aligned}$$

Using statement (a) and, as in the previous proof, that the worst case scenario in the upper bounds of equation (3) occurs for $n_i = n + t - i$, we obtain

$$\left| \log \frac{DT^t(x)}{DT^t(y)} \right| \lesssim n^{p \frac{\gamma+1}{\gamma}} \sum_{i=0}^{t-1} (n+t-i)^{-2} |T^t(x) - T^t(y)|^p \leq n^{p \frac{\gamma+1}{\gamma} - 1} |T^t(x) - T^t(y)|^p$$

and statement (b) follows from the assumption $p \leq \frac{\gamma}{\gamma+1}$. \square

3.2 Estimating escape rates - proof of Theorem 2.1

Recall that H is a cylinder in $\mathcal{J}^{(\ell_H)}$ and $H \subseteq J_h$ for some $h \geq 0$. The main estimate of this section is the following lemma.

Lemma 3.2. $m(\hat{I}^t) \lesssim t^{-\frac{1}{\gamma}}$.

The proof of this lemma is based on the fact that an induced map related to T has exponential escape rate. To formulate this property, choose $n_S \geq h$, let $I_S = [a_{n_S}, 1) \supset H$ and consider the induced map $S = T^R : I_S \rightarrow I_S$, where R is the first return time to I_S .

Let $\hat{I}_S^t = \bigcap_{i=0}^t S^{-i}(I_S \setminus H)$ denote the set of points in I_S which do not enter H before time t under the action of S . The induced open system $S|_{\hat{I}_S^t}$ is uniformly expanding with countably many branches and admits a countable Markov partition which is formed by joining $\mathcal{J}^{(\ell_H)}$ with the partition into sets on which R is constant. The action of S on this partition satisfies the large images condition [15]; hence the following property holds.¹

Lemma 3.3. [15] *There exists $\sigma < 1$ such that $m(\hat{I}_S^t) \lesssim \sigma^t$.*

¹The full results of [15] also require a ‘‘smallness’’ condition on the size of the hole. This condition is not needed here since we are not invoking any results regarding a spectral gap for the transfer operator associated with S , but just an exponential rate of escape, which does not require the hole to be small.

Proof of Lemma 3.2. We first assume there exists $d > 0$ such that $[\frac{1}{2}, \frac{1}{2} + d) \cap H = \emptyset$. The complementary case is much simpler and will be addressed at the end of the proof.

Without loss of generality, we can choose the index n_S that defines I_S sufficiently large so that $\frac{a_{n_S}}{2} < d$, viz. the open system makes full returns to the interval $J_S = [\frac{1}{2}, \frac{1}{2} + \frac{a_{n_S}}{2})$ before entering $I \setminus I_S$. In order to obtain the estimate on $m(\dot{I}^t)$, we consider separately the sets $\dot{I}^t \cap I_S$ and $\dot{I}^t \cap [0, a_{n_S})$.

Case I. Estimate for points in $\dot{I}^t \cap I_S$. Consider the decomposition of $\dot{I}^t \cap I_S$ into subsets E_k^t of points having made k passes through $I \setminus I_S$ before time t . After each pass through $I \setminus I_S$, an orbit must spend at least $n_S + 1$ iterates within I_S before making its next pass. It results that the index k here is at most $\lfloor \frac{t-2}{n_S+2} \rfloor$.

In order to estimate the measure of the sets E_k^t , we consider separately the cases $k > b \log t$ and $k \leq b \log t$, where $b = \frac{1}{(n_S+1)\gamma \log \sigma^{-1}}$ and σ is from Lemma 3.3.

For $k > b \log t$, we observe that every point in E_k^t must spend at least $(k-1)(n_S+1)$ iterates in I_S before hitting the hole. Hence we have $E_k^t \subset \dot{I}_S^{(k-1)(n_S+1)}$ and using Lemma 3.3 and the definition of b , we get

$$\sum_{k > b \log t} m(E_k^t) \lesssim \sum_{k > b \log t} \sigma^{(k-1)(n_S+1)} \lesssim t^{-\frac{1}{\gamma}},$$

as desired.

For $k \leq b \log t$, we first note that the case $k = 0$ is easily estimated using $E_0^t \subset \dot{I}_S^t$ and Lemma 3.3. From now on, we assume $k \in \{1, \dots, b \log t\}$ and observe as before that the subset $E_k^{t,+} \subset E_k^t$ of points whose orbit spends at least $b(n_S+1) \log t$ iterates in I_S up to time t is included in $\dot{I}_S^{b(n_S+1) \log t}$. This inclusion implies $m\left(\bigcup_{k=1}^{b \log t} E_k^{t,+}\right) \lesssim t^{-\frac{1}{\gamma}}$.

It remains to consider the complementary subset $E_k^{t,-} = E_k^t \setminus E_k^{t,+}$ of points whose orbits spend more than $t - b(n_S+1) \log t$ iterates in $[0, n_S)$. Given $x \in E_k^{t,-}$ and $i \in \{1, \dots, k\}$, let $n_i > n_S$ be such that J_{n_i} is the element of \mathcal{J} where $T^j(x)$ begins its i th pass through $[0, a_{n_S})$. We have

$$m(E_k^{t,-}) \leq \sum_{i=1}^k m\left(x \in E_k^t : n_i - n_S > \frac{t - b(n_S+1) \log t}{k}\right). \quad (5)$$

The sets in this sum can be decomposed using symbolic dynamics. Given two integers $t_1 < t_2$, let

$J_{\theta^{t_1}^{t_2}} = \bigcap_{\ell=t_1}^{t_2} T^{-\ell}(J_{\theta_\ell})$. We have

$$\{x \in E_k^t : n_i = n\} = \bigcup_{j, \{\theta_\ell\}_{\ell=1}^{j-2}, \{\theta_\ell\}_{\ell=j+1}^t} J_{\theta_0^{j-1}} \cap T^{-j}(J_n) \cap J_{\theta_j^{t_1}}, \quad (6)$$

where, by an abuse of notation²

$$J_{\theta_0^{j-1}} := I_S \cap J_{\theta_1^{j-2}} \cap T^{-(j-1)}(J_S),$$

²For this expression to be meaningful, we should decompose the sets I_S and $T^{-(j-1)}(J_S)$ into (standard) cylinder sets prior to define the intermediate cylinder $J_{\theta_1^{j-2}}$. This abuse of notation has no impact on the reasoning here. (The same comment applies to the set $T^{-j}(J_n)$ and to the decomposition in equation (6).)

and where the union on $\{\theta_\ell\}_{\ell=0}^{j-2}$ (resp. $\{\theta_\ell\}_{\ell=j+1}^t$) is taken over all admissible words compatible with $i-1$ (resp. $k-i$) passes through $[0, n_S)$ and avoiding H until at least time t . The sets in (6) are pairwise disjoint; hence it suffices to estimate each quantity $m\left(J_{\theta_0^{j-1}} \cap T^{-j}(J_n) \cap J_{\theta_{j+1}^t}\right)$.

To proceed, notice first that the property $T^n(J_n) = J_0$ implies

$$J_{\theta_0^{j-1}} \cap T^{-j}(J_n) = J_{\theta_0^{j-1}} \cap T^{-j}(I \setminus I_S) \cap T^{-(j+n)}(J_0).$$

Moreover, the map T^{j+n} is one-to-one on each element of $T^{-(j+n)}(J_0)$. Applying the bounded distortion estimate in the proof of Lemma 3.1, we obtain

$$\frac{m\left(J_{\theta_0^{j-1}} \cap T^{-j}(J_n) \cap J_{\theta_{j+1}^t}\right)}{m\left(J_{\theta_0^{j-1}} \cap T^{-j}(I \setminus I_S) \cap T^{-(j+n)}(J_0)\right)} \approx \frac{m\left(T^{j+n}\left(J_{\theta_0^{j-1}} \cap T^{-j}(J_n) \cap J_{\theta_{j+1}^t}\right)\right)}{m\left(T^{j+n}\left(J_{\theta_0^{j-1}} \cap T^{-j}(I \setminus I_S) \cap T^{-(j+n)}(J_0)\right)\right)}.$$

The second ratio here is equal to

$$\frac{m\left(J_0 \cap T^{j+n}(J_{\theta_{j+n+1}^t})\right)}{m(J_0)},$$

from where our first estimate follows

$$m\left(J_{\theta_0^{j-1}} \cap T^{-j}(J_n) \cap J_{\theta_{j+1}^t}\right) \approx m\left(J_0 \cap T^{j+n}(J_{\theta_{j+n+1}^t})\right) m\left(J_{\theta_0^{j-1}} \cap T^{-j}(J_n)\right).$$

Proceeding similarly for the second factor above and using $T(J_S) = I \setminus I_S$, we get

$$\frac{m\left(J_{\theta_0^{j-1}} \cap T^{-j}(J_n)\right)}{m\left(J_{\theta_0^{j-1}}\right)} \approx \frac{m\left(T^j\left(J_{\theta_0^{j-1}} \cap T^{-j}(J_n)\right)\right)}{m\left(T^j\left(J_{\theta_0^{j-1}}\right)\right)} = \frac{m(J_n)}{a_{n_S}},$$

from which equation (3) implies

$$m\left(J_{\theta_0^{j-1}} \cap T^{-j}(J_n) \cap J_{\theta_{j+1}^t}\right) \approx n^{-\frac{\gamma+1}{\gamma}} m\left(J_0 \cap T^{j+n}(J_{\theta_{j+n+1}^t})\right) m\left(J_{\theta_0^{j-1}}\right). \quad (7)$$

Now use that imposing $i-1$ passes through $I \setminus I_S$ before time j implies at least $(i-1)(n_S+1)$ iterates in I_S before j to obtain the following relation

$$m\left(\bigcup_{j, \{\theta_\ell\}_{\ell=1}^{j-2}} J_{\theta_0^{j-1}}\right) \subset m\left(\dot{I}_S^{(i-1)(n_S+1)}\right) \lesssim \sigma^{(i-1)(n_S+1)}.$$

Similarly, let q be the number of iterates that the orbits of points in $J_0 \cap T^{j+n}(J_{\theta_{j+n+1}^t})$ spend in I_S . Each pass in $I \setminus I_S$ from $i+1$ through $k-1$ must be followed by at least n_S+1 iterates in I_S ; hence $q \geq (k-i-1)(n_S+1)$ (also $q \leq t - (j+n) + 1 - (k-i-1)$ where the maximum is obtained when each pass in $I \setminus I_S$ consists of a single iterate) and then by Lemma 3.3,

$$m\left(\bigcup_{j, \{\theta_\ell\}_{\ell=j+1}^t} J_0 \cap T^{j+n}(J_{\theta_{j+n+1}^t})\right) \leq \sum_{q \geq (k-i-1)(n_S+1)} m(\dot{I}_S^q) \lesssim \sigma^{(k-i)(n_S+1)}.$$

(Notice that this estimate holds even in the case $i = k$.) Putting these estimates together with (6) and (7), we have obtained

$$m(x \in E_k^t : n_i = n) \lesssim \sigma^{(k-1)(n_S+1)} n^{-\frac{\gamma+1}{\gamma}}.$$

Using the inequality

$$\sum_{n > n_S + \frac{t - b(n_S+1) \log t}{k}} n^{-\frac{\gamma+1}{\gamma}} \leq \left(n_S + \frac{t - b(n_S+1) \log t}{k} \right)^{-\frac{1}{\gamma}}$$

it follows from (5) that

$$m(E_k^{t,-}) \lesssim \sigma^{(k-1)(n_S+1)} k^{\frac{\gamma+1}{\gamma}} (kn_S + t - b(n_S+1) \log t)^{-\frac{1}{\gamma}}.$$

It remains to sum over k . We finally have

$$m\left(\bigcup_{k=1}^{b \log t} E_k^{t,-}\right) \lesssim (n_S + t - b(n_S+1) \log t)^{-\frac{1}{\gamma}} \sum_{k=1}^{b \log t} \sigma^{(k-1)(n_S+1)} k^{\frac{\gamma+1}{\gamma}} \lesssim t^{-\frac{1}{\gamma}},$$

as desired, where we used $\sum_{k=1}^{b \log t} \sigma^{(k-1)(n_S+1)} k^{\frac{\gamma+1}{\gamma}} < +\infty$ and $(n_S + t - b(n_S+1) \log t)^{-1} \lesssim t^{-1}$.

Case II. Estimate for points in $\dot{I}^t \cap [0, a_{n_S})$. Recall that $[0, a_{n_S}) = \bigcup_{n > n_S} J_n$ and by definition of the J_n , we have $\dot{I}^t \supset \bigcup_{n > t+h} J_n$ so that using equation (3) yields $m\left(\dot{I}^t \cap \bigcup_{n > t+h} J_n\right) \lesssim (t+h)^{-\frac{1}{\gamma}} \leq t^{-\frac{1}{\gamma}}$.

It remains to estimate $m\left(\dot{I}^t \cap \bigcup_{n=n_S+1}^{t+h} J_n\right)$. For every $n > n_S$, we have $T^{n-n_S}(\dot{I}^t \cap J_n) = \dot{I}^{t-n+n_S} \cap J_{n_S}$. Using bounded distortion again, we get

$$\frac{m(\dot{I}^t \cap J_n)}{m(J_n)} \approx \frac{m(\dot{I}^{t-n+n_S} \cap J_{n_S})}{m(J_{n_S})},$$

which, together with the inclusion $J_{n_S} \subset I_S$ and the conclusion in Case I, implies

$$m\left(\dot{I}^t \cap \bigcup_{n=n_S+1}^{t+h} J_n\right) \lesssim \sum_{n=n_S+1}^{t+h} n^{-\frac{\gamma+1}{\gamma}} (t-n+n_S)^{-\frac{1}{\gamma}}. \quad (8)$$

Using the inequality $\sum_{n=a}^b f(n) \leq \int_{a-1}^{b+1} f(x) dx$ which holds for every $f \geq 0$, we estimate the sum up to $t+h-1$ via the following integral (for the last term, we have $(t+h)^{-\frac{\gamma+1}{\gamma}} (n_S-h)^{-\frac{1}{\gamma}} \lesssim t^{-\frac{1}{\gamma}}$),

$$\int_{n_S}^{t+h} x^{-\frac{\gamma+1}{\gamma}} (t+n_S-x)^{-\frac{1}{\gamma}} dx = \int_{n_S}^{t/2} x^{-\frac{\gamma+1}{\gamma}} (t+n_S-x)^{-\frac{1}{\gamma}} dx + \int_{t/2}^{t+h} x^{-\frac{\gamma+1}{\gamma}} (t+n_S-x)^{-\frac{1}{\gamma}} dx.$$

In the first integral, the second factor in the integrand is at most $(n_S + \frac{t}{2})^{-\frac{1}{\gamma}}$ while the first factor integrates to something less than $\gamma n_S^{-\frac{1}{\gamma}}$. In the second integral, the first factor is at most

$(\frac{t}{2})^{-\frac{\gamma+1}{\gamma}}$ while the second factor integrates to something less than $\frac{\gamma}{1-\gamma}(n_S - h)^{-\frac{1}{\gamma}+1}$. It results that $m\left(\dot{I}^t \cap \bigcup_{n=n_S+1}^{t+h} J_n\right) \lesssim t^{-\frac{1}{\gamma}}$ and this concludes the proof of Lemma 3.2 in the case in which H is disjoint from $[\frac{1}{2}, \frac{1}{2} + d)$ for some $d > 0$.

The case in which H contains $[\frac{1}{2}, \frac{1}{2} + d)$ for some $d > 0$ is much simpler since points starting in $[2d, 1)$ never enter $[0, 2d)$ before escaping. Thus the estimates of Case I with $k = 0$ together with Case II imply that the upper bound on $m(\dot{I}^t)$ in this case is the same. \square

With Lemma 3.2 established, we are ready to prove Theorem 2.1.

Proof of Theorem 2.1. Given a density f as in the theorem, let $C_0 = \inf_{x \in (0, x_0)} x^\alpha f(x) > 0$ and $n_0 = \min\{n : a_n \leq x_0\}$. We have $\dot{I}^t \supset \bigcup_{n>t+h} J_n$ for all $t \geq 0$; hence for t sufficiently large so that $t + h \geq n_0$, the equation (3) implies

$$\mu_f(\dot{I}^t) = \int_{\dot{I}^t} f dm \geq C_0 \int_0^{a_{t+h}} x^{-\alpha} dx \geq \frac{C_0}{1-\alpha} (h+t)^{-\frac{1-\alpha}{\gamma}}, \quad (9)$$

from which the lower bound immediately follows.

For the upper bound, we split $[0, 1)$ into 3 intervals: $[0, a_{h+t}) \cup [a_{h+t}, a_{n_S}) \cup I_S$ and estimate the intersection of \dot{I}^t with each of these separately.³

On $[0, a_{h+t})$, we estimate,

$$\int_{\dot{I}^t \cap [0, a_{h+t})} f dm \lesssim \int_0^{a_{h+t}} x^{-\alpha} dx \lesssim (h+t)^{-\frac{1-\alpha}{\gamma}},$$

while on I_S , we have, using that f is bounded on I_S and Lemma 3.2,

$$\int_{\dot{I}^t \cap I_S} f dm \lesssim m(\dot{I}^t) \lesssim t^{-\frac{1}{\gamma}}.$$

On $[a_{h+t}, a_{n_S}) = \bigcup_{n=n_S+1}^{t+h} J_n$, we proceed as in Case II of the previous proof,

$$\int_{\dot{I}^t \cap [a_{h+t}, a_{n_S})} f dm = \sum_{n=n_S+1}^{t+h} \int_{\dot{I}^t \cap J_n} f dm \lesssim \sum_{n=n_S+1}^{t+h} a_n^{-\alpha} m(\dot{I}^t \cap J_n) \lesssim \sum_{n=n_S+1}^{t+h} n^{-\frac{\gamma+1-\alpha}{\gamma}} (t-n+n_S)^{-\frac{1}{\gamma}},$$

where we have used (8). As before, the last sum (except its last term) is estimated after splitting into two integrals, to give

$$\sum_{n=n_S+1}^{t+h} n^{-\frac{\gamma+1-\alpha}{\gamma}} (t-n+n_S)^{-\frac{1}{\gamma}} \lesssim t^{-\frac{1-\alpha}{\gamma}}.$$

\square

³For $\alpha = 0$, the upper bound also directly follows from the fact that f is uniformly bounded on I together with Lemma 3.2,

$$\mu_f(\dot{I}^t) \lesssim m(\dot{I}^t) \lesssim t^{-\frac{1}{\gamma}}.$$

3.3 Properties of the function spaces \mathcal{F}_p

The definition of the quantity H_J^p before Theorem 2.1 implies the following simple facts about the set \mathcal{F}_p ($p \in \mathbb{R}^+$), whose proof we leave to the reader.

Lemma 3.4. (1) $\|\cdot\|_p$ is scale invariant, i.e. $\|Cf\|_p = \|f\|_p$ for any $C > 0$.

(2) For any $J \in \mathcal{J}^{(\ell_H)}$, E a subinterval of J and $f \in \mathcal{F}_p$, we have

$$\sup_{x \in J} f(x) \leq e^{H_J^p(f)|E|^p} \inf_{x \in E} f(x) \leq e^{H_J^p(f)|E|^p} |E|^{-1} \int_E f dm.$$

Now, we equip the set of measures with the topology of weak convergence, consider the ball $\mathcal{B}_p = \{\mu_f : f \in \mathcal{F}_p, \|f\|_p \leq 1\}$ and notice that this ball is not closed. Indeed, given $\ell \in \mathbb{N}$, let the density f_ℓ be defined by

$$f_\ell(x) = \begin{cases} a_\ell^{-1} & \text{if } x \in \bigcup_{n \geq \ell+1} J_n \\ 0 & \text{elsewhere.} \end{cases}$$

Then we have, $f_\ell \geq 0$, $\int f_\ell dm = 1$ and $\|f_\ell\|_p = 0$ so that $\mu_{f_\ell} \in \mathcal{B}_p$. However, we clearly have

$$\lim_{\ell \rightarrow \infty} \mu_{f_\ell} = \delta_0 \notin \mathcal{B}_p.$$

The Dirac measure at 0 turns out to be the only possible singular component to where sequences in \mathcal{B}_p can accumulate.

Lemma 3.5. The set $\{(1-s)\mu_f + s\delta_0 : s \in [0, 1], f \in \mathcal{B}_p\}$ is compact.

Proof. Let $\{\mu_{f_\ell}\}_{\ell \in \mathbb{N}} \subset \mathcal{B}_p$ be an arbitrary sequence. Since $\int f_\ell dm = 1$, there exists a subsequence $\{\mu_{f_{\ell_k}}\}$ which converges weakly to a probability measure μ_∞ on I . Now fix $J \in \mathcal{J}^{(\ell_H)}$. By Lemma 3.4, the sequence of densities $\{f_{\ell_k}\}$ is a bounded, equicontinuous family on J . By the Arzelà-Ascoli theorem, there exists a subsequence that converges uniformly to a function $f_J^{(\infty)}$ on J .⁴ Note that $H_J^p(f_J^{(\infty)}) \leq 1$.

Diagonalizing, we obtain a subsequence $\{f_{\ell_{k_j}}\}$ converging to $f_J^{(\infty)}$ on each $J \in \mathcal{J}^{(\ell_H)}$. Letting $f^{(\infty)} = \sum_J f_J^{(\infty)}$, we have $f^{(\infty)} \geq 0$, $\|f^{(\infty)}\|_p \leq 1$, and by Fatou's lemma, $\int f^{(\infty)} dm \leq 1$.

Let $s = 1 - \int f^{(\infty)} dm$. If $s < 1$, let $f_\infty = (1-s)^{-1} f^{(\infty)}$. By the above observations, we have $\mu_{f_\infty} \in \mathcal{B}_p$. Since $\{0\}$ is the only accumulation point of the sequence of sets $\{J\}_{J \in \mathcal{J}^{(\ell_H)}}$, we must have $\mu_\infty = (1-s)\mu_{f_\infty} + s\delta_0$, as required. \square

For the next statement, we need to introduce the (nonlinear) normalized transfer operator and its iterates,

$$\mathring{\mathcal{L}}_1^t f := \frac{\mathring{\mathcal{L}}^t f}{|\mathring{\mathcal{L}}^t f|_1}, \quad \forall t \geq 1. \quad (10)$$

⁴To be precise, note that $\|f_\ell\|_p \leq 1$ implies that f_ℓ is uniformly continuous on J so that f can be extended to a continuous function f_ℓ on the closure \bar{J} and it still holds that $\|f_\ell\|_p \leq 1$.

Proposition 3.6. *Let $p \in (0, \frac{\gamma}{\gamma+1}]$. We have $\mathring{\mathcal{L}}_1(\mathcal{F}_p^0) \subset \mathcal{F}_p^0$. In addition, there exist two constants $C_1, C_2 \geq 0$, such that for every $f \in \mathcal{F}_p^0$,*

$$\|\mathring{\mathcal{L}}_1^t f\|_p \leq C_1 \|f\|_p + C_2 \text{ for all } t \geq 1.$$

Proof. Every $J \subset \mathcal{J}^{(\ell_H)}$ has at most two pre-images under \mathring{T} and each pre-image is included in some element of $\mathcal{J}^{(\ell_H)}$. This implies that $\mathring{\mathcal{L}}(C^0(\mathcal{J}^{(\ell_H)})) \subset C^0(\mathcal{J}^{(\ell_H)})$. Also we obviously have

$$\mathring{\mathcal{L}}f \geq 0 \quad \text{and} \quad |\mathring{\mathcal{L}}f|_1 \leq |f|_1,$$

for every $f \geq 0$.

Now, assume that f is such that $f|_{\bigcup_{n \geq n_0} J_n} \geq C$ for some $n_0 \in \mathbb{N}$ and $C > 0$, and let $n' = \max\{n_0, h\} + 1$. Clearly, we have $\mathring{\mathcal{L}}f|_{\bigcup_{n \geq n'} J_n} \geq \frac{C}{1+2^\gamma(1+\gamma)a_{n'-1}^\gamma}$. In particular, $|\mathring{\mathcal{L}}f|_1 > 0$ so that $\mathring{\mathcal{L}}_1^t f$ is well defined. Anticipating the proof below that $\|\mathring{\mathcal{L}}_1^t f\|_p < +\infty$ for every f with $\|f\|_p < +\infty$, we obtain $\mathring{\mathcal{L}}_1(\mathcal{F}_p^0) \subset \mathcal{F}_p^0$.

In order to check the estimate on $\|\mathring{\mathcal{L}}_1^t f\|_p$, it suffices to prove the inequality for $\|\mathring{\mathcal{L}}^t f\|_p$ due to the scale invariance property from Lemma 3.4(1).

Let $f \in \mathcal{F}_p^0$. Fix $n \geq 0$, $J \in \mathcal{J}^{(\ell_H)}$, $J \subset J_n$ and $x, y \in J$. Let also $t \geq 1$ and $\{x_i\}$ (resp. $\{y_i\}$) be an enumeration of the pre-images $\mathring{T}^{-t}(x)$ (resp. $\mathring{T}^{-t}(y)$) such that each pair x_i, y_i lies in the same branch of \mathring{T}^{-t} . Then

$$\log \mathring{\mathcal{L}}^t f(x) - \log \mathring{\mathcal{L}}^t f(y) = \log \frac{\sum_i f(x_i)/DT^t(x_i)}{\sum_i f(y_i)/DT^t(y_i)} \leq \max_i \log \frac{f(x_i)}{f(y_i)} + \log \frac{DT^t(y_i)}{DT^t(x_i)}, \quad (11)$$

where we have used the fact that $\frac{\sum_i b_i}{\sum_i c_i} \leq \max_i \frac{b_i}{c_i}$ for two series of positive terms. The first term on the right hand side of (11) is estimated by

$$\log \frac{f(x_i)}{f(y_i)} \leq \|f\|_p |x_i - y_i|^p,$$

while the second term on the right side of (11) is estimated by

$$\log \frac{DT^t(x_i)}{DT^t(y_i)} \lesssim |T^t(x_i) - T^t(y_i)|^p = |x - y|^p,$$

according to Lemma 3.1(b). Putting these two estimates together and using Lemma 3.1(a), we obtain that there exist $C_1, C_2 > 0$ such that

$$\begin{aligned} \frac{|\log \mathring{\mathcal{L}}^t f(x) - \log \mathring{\mathcal{L}}^t f(y)|}{|x - y|^p} &\leq \|f\|_p \max_i \frac{|x_i - y_i|^p}{|x - y|^p} + C_2 \\ &\leq \|f\|_p C_1 \left(\frac{n}{n+t} \right)^{p \frac{\gamma+1}{\gamma}} + C_2 \\ &\leq C_1 \|f\|_p + C_2, \end{aligned}$$

as desired. □

The last statement of this section provides examples of unbounded densities that belong to the sets \mathcal{F}_p^0 .

Lemma 3.7. *Let $p \in (0, \frac{\gamma}{\gamma+1}]$. Every function $\frac{f}{|f|_1}$ where $f(x) = x^{-\alpha}$ for all $x \in (0, 1)$, $\alpha \in [0, 1)$, belongs to \mathcal{F}_p^0 . The same is true for the density f_{SRB} associated to the SRB measure.*

Proof. Letting $f(x) = x^{-\alpha}$ for all $x \in (0, 1)$, we clearly have $f \in C^0(\mathcal{J}^{(\ell_H)})$, $f \geq 0$ and $|f|_1 < +\infty$. Therefore, we only need to check that $\|f\|_p < +\infty$.

Fix $J \in \mathcal{J}^{(\ell_H)}$, $J \subset J_n$ for some $n \geq 0$ and let $x < y \in J$. Then

$$\frac{\log f(x) - \log f(y)}{|x - y|^p} = \alpha \frac{\log(y/x)}{|x - y|^p} \leq \frac{\alpha}{x} |x - y|^{1-p}.$$

Since $x \in J_n$, we have $\frac{1}{x} \lesssim n^{-\frac{1}{\gamma}}$ and $|x - y| \lesssim n^{-\frac{\gamma+1}{\gamma}}$ by equation (3). Thus

$$H_J^p(f) \lesssim n^{\frac{1-(\gamma+1)(1-p)}{\gamma}}$$

and the exponent of n is non-positive when $p \leq \gamma/(\gamma+1)$. Taking the supremum over n , we have $\|f\|_p < \infty$ as required.

As for f_{SRB} , observe that the sequence $\{\mathcal{L}^t 1\}_{t \in \mathbb{N}}$ (transfer operator for the system without hole) converges to the density f_{SRB} . On the other hand, the proof of Proposition 3.6 can be repeated *mutatis mutandis* to conclude $\mathcal{L}(\mathcal{F}_p^0) \subset \mathcal{F}_p^0$ and $\sup_{t \in \mathbb{N}} \|\mathcal{L}^t f\|_p < +\infty$ for all $f \in \mathcal{F}_p^0$. Since $1 \in \mathcal{F}_p^0$, it follows from Lemma 3.5 that $f_{\text{SRB}} \in \mathcal{F}_p$. However, $f_{\text{SRB}}(x) \approx x^{-\gamma}$; hence we must have $f_{\text{SRB}} \in \mathcal{F}_p^0$ as desired. \square

3.4 Proof of Theorem 2.2

Throughout this proof, we choose an arbitrary number $p \in (0, \frac{\gamma}{\gamma+1}]$ and density $f \in \mathcal{F}_p^0$.

By Proposition 3.6, the sequence $\{\frac{\hat{T}_*^t \mu_f}{\mu_f(\hat{I}^t)}\}_{t \in \mathbb{N}}$ is composed of absolutely continuous probability measures with densities in \mathcal{F}_p^0 . By Lemma 3.5, any of its limit points must have the form $\mu_\infty = (1 - s_\infty)\mu_{f_\infty} + s_\infty \delta_0$ for some $f_\infty \in \mathcal{F}_p$ and $s_\infty \in [0, 1]$. We want to prove that $s_\infty = 1$ for any limit point.

To proceed, we first prove that $s_\infty > 0$ in all cases. Take any cylinder $J \subset \mathcal{J}^{(\ell_H)}$ and consider a converging subsequence $\{\hat{\mathcal{L}}_1^{t_j} f\}_{j \in \mathbb{N}}$ with limit point $(1 - s_\infty)f_\infty$. (Recall that $\hat{\mathcal{L}}_1^t$ is the normalized transfer operator, see equation (10) above.) The limit density f belongs to \mathcal{F}_p ; hence the convergence $\hat{\mathcal{L}}_1^{t_j} f|_J \rightarrow (1 - s_\infty)f_\infty|_J$ holds in the uniform topology of functions defined on this interval. In particular, its integrals against any bounded measurable function converge as well on each $J \in \mathcal{J}^{(\ell_H)}$. Thus for any $k \in \mathbb{N}$, since $\hat{I}^k \supset [0, a_{h+k})$, we have on the one hand,

$$\begin{aligned} \lim_{j \rightarrow \infty} \frac{\int \hat{\mathcal{L}}_1^{t_j} f \cdot 1_{\hat{I}^k} dm}{|\hat{\mathcal{L}}_1^{t_j} f|_1} &= \lim_{j \rightarrow \infty} \frac{\int_{[0, a_{h+k})} \hat{\mathcal{L}}_1^{t_j} f dm}{|\hat{\mathcal{L}}_1^{t_j} f|_1} + \frac{\int_{[a_{h+k}, 1)} \hat{\mathcal{L}}_1^{t_j} f \cdot 1_{\hat{I}^k} dm}{|\hat{\mathcal{L}}_1^{t_j} f|_1} \\ &= \mu_\infty([0, a_{h+k})) + \mu_\infty([a_{h+k}, 1) \cap \hat{I}^k) = \mu_\infty(\hat{I}^k). \end{aligned}$$

On the other hand, by applying a change of variable, we get

$$\int \hat{\mathcal{L}}_1^{t_j} f \cdot 1_{\hat{I}^k} dm = \int f \cdot 1_{\hat{I}^{t_j}} \cdot 1_{\hat{I}^k} \circ \hat{T}^{t_j} dm = \int f \cdot 1_{\hat{I}^{t_j}} \cdot 1_{\hat{I}^{k+t_j}} dm = \mu_f(\hat{I}^{k+t_j}).$$

Using Theorem 2.1 immediately yields,

$$\frac{\int \mathcal{L}^{t_j} f \cdot 1_{\hat{I}^k} dm}{\mu_f(\hat{I}^{t_j})} = \frac{\mu_f(\hat{I}^{k+t_j})}{\mu_f(\hat{I}^{t_j})} \approx \left(\frac{k+t_j}{t_j} \right)^{-\frac{1-\alpha}{\gamma}},$$

which implies the existence of $C > 0$ such that $\mu_\infty(\hat{I}^k) \geq C$ for all $k \geq 1$. However, ergodicity ensures that we must have

$$\lim_{k \rightarrow +\infty} \mu_{f_\infty}(\hat{I}^k) = 0, \quad (12)$$

therefore we conclude that $s_\infty \geq C > 0$ as desired.

For the second step of the proof, we assume, without loss of generality, that the subsequence $\{t_j\}_{j \in \mathbb{N}}$ along which $\frac{\hat{T}_*^{t_j} \mu_f}{\mu_f(\hat{I}^{t_j})}$ converge to μ_∞ is such that $\lim_{j \rightarrow +\infty} t_{j+1} - t_j = +\infty$. We are going to show that

$$\lim_{j \rightarrow +\infty} \frac{\hat{T}_*^{t_{j+1}-t_j} \mu_\infty}{\mu_\infty(\hat{I}^{t_{j+1}-t_j})} = \delta_0. \quad (13)$$

To that purpose, given a function $\varphi \in C^0(I)$, we write using the expression for μ_∞ and a change of variable again

$$\begin{aligned} \int \varphi d\hat{T}_*^{t_{j+1}-t_j} \mu_\infty &= (1 - s_\infty) \int \varphi \cdot \mathcal{L}^{t_{j+1}-t_j} f_\infty dm + s_\infty \varphi(0) \\ &= (1 - s_\infty) \int \varphi \circ \hat{T}^{t_{j+1}-t_j} \cdot f_\infty \cdot 1_{\hat{I}^{t_{j+1}-t_j}} dm + s_\infty \varphi(0). \end{aligned}$$

By assumption on φ , the function $|\varphi \circ \hat{T}^{t_{j+1}-t_j}|$ is bounded. Using also equation (12) and the assumption $\lim_{j \rightarrow +\infty} t_{j+1} - t_j = +\infty$, we obtain

$$\lim_{j \rightarrow +\infty} \int \varphi \circ \hat{T}^{t_{j+1}-t_j} \cdot f_\infty \cdot 1_{\hat{I}^{t_{j+1}-t_j}} dm = \lim_{j \rightarrow +\infty} \mu_{f_\infty}(\hat{I}^{t_{j+1}-t_j}) = 0,$$

so that

$$\lim_{j \rightarrow +\infty} \int \varphi d\hat{T}_*^{t_{j+1}-t_j} \mu_\infty = s_\infty \varphi(0).$$

Applying this relation with $\varphi = 1$ and using that $s_\infty > 0$ yields

$$\lim_{j \rightarrow +\infty} \frac{\int \varphi d\hat{T}_*^{t_{j+1}-t_j} \mu_\infty}{\mu_\infty(\hat{I}^{t_{j+1}-t_j})} = \varphi(0), \quad \forall \varphi \in C^0(I)$$

which is nothing but the desired behavior (13).

Now, the measure μ_∞ clearly gives zero weight to the singularity set for \hat{T} . By continuity, letting $\mu_t := \frac{\hat{T}_*^t \mu_f}{\mu_f(\hat{I}^t)}$, it follows that

$$\lim_{j \rightarrow +\infty} \frac{\hat{T}_*^{t_{k+1}-t_k} \mu_{t_j}}{\mu_{t_j}(\hat{I}^{t_{k+1}-t_k})} = \frac{\hat{T}_*^{t_{k+1}-t_k} \mu_\infty}{\mu_\infty(\hat{I}^{t_{k+1}-t_k})}, \quad \forall k \in \mathbb{N}.$$

By construction, we have

$$\frac{\hat{T}_*^{t_{k+1}-t_k} \mu_{t_k}}{\mu_{t_k}(\hat{I}^{t_{k+1}-t_k})} = \frac{\hat{T}_*^{t_{k+1}} \mu_f}{\mu_f(\hat{I}^{t_{k+1}})}.$$

Hence, using a diagonal argument together with equation (13), we conclude

$$\mu_\infty = \lim_{k \rightarrow +\infty} \frac{T_*^{\circ t_{k+1}} \mu_f}{\mu_f(\mathring{I}^{t_{k+1}})} = \delta_0,$$

as desired.

Finally, since again δ_0 is an invariant measure that sits outside the singularity set of \mathring{T} , we have

$$\delta_0 = \mathring{T}_*(\delta_0) = \mathring{T}_* \left(\lim_{t \rightarrow +\infty} \frac{\mathring{T}_*^t \mu_f}{\mu_f(\mathring{I}^t)} \right) = \lim_{t \rightarrow +\infty} \frac{\mathring{T}_*^{t+1} \mu_f}{\mu_f(\mathring{I}^t)},$$

from which the relation

$$\lim_{t \rightarrow +\infty} \frac{\mu_f(\mathring{I}^{t+1})}{\mu_f(\mathring{I}^t)} = 1,$$

immediately follows.

3.5 General open systems - proof of Theorem 2.3

The proof requires a preliminary statement. We say a set $A \subset \mathbb{N}$ has density $\rho \in [0, 1]$ if $\lim_{n \rightarrow \infty} \frac{\#(A \cap [1, n])}{n} = \rho$, where $\#$ denotes the cardinality of a set.

Define $\beta_t = \frac{\mu(\mathring{X}^{t+1})}{\mu(\mathring{X}^t)}$ for $t \geq 0$ and $\beta_{-1} = \mu(\mathring{X}^0)$. We obviously have $\beta_t \leq 1$ for all $t \geq 0$ and if it happens that $\liminf_{t \rightarrow \infty} \beta_t < 1$, then the subexponential escape rate assumption (1) imposes a strong constraint on any index set with accumulation point below 1.

Lemma 3.8. *For $\lambda \in (0, 1)$, define $A_\lambda = \{k \in \mathbb{N} : \beta_k \leq \lambda\}$. Then A_λ must have zero density.*

Proof. Fix $\lambda \in (0, 1)$ and note that if A_λ is finite, it has zero density. So assume A_λ is infinite. Then $\limsup_{\substack{k \rightarrow \infty \\ k \in A_\lambda}} \beta_k \leq \lambda$. Equation (1) implies

$$\limsup_{t \rightarrow +\infty} \mu(\mathring{X}^t)^{\frac{1}{t}} = \limsup_{t \rightarrow +\infty} \exp \left(\frac{1}{t} \log \mu(\mathring{X}^t) \right) = 1. \quad (14)$$

To derive a contradiction, assume A has density $\rho > 0$. The relation $\mu(\mathring{X}^t) = \prod_{k=0}^{t-1} \beta_k$ then implies

$$\mu(\mathring{X}^t)^{\frac{1}{t}} \leq \lambda^{\frac{\#(A_\lambda \cap [0, t-1])}{t}} \text{ for all } t > 0,$$

from which it follows that $\limsup_{t \rightarrow \infty} \mu(\mathring{X}^t)^{\frac{1}{t}} \leq \lambda^\rho < 1$, contradicting equation (14). \square

Passing now to the proof of the theorem, we consider the sequence of probability measures $\{\mu_t\}_{t \in \mathbb{N}}$ where μ_t is defined by equation (2). By assumption on X , this sequence is compact and any limit point μ_∞ is necessarily a probability measure.

Our proof has two steps: (A) μ_∞ gives zero weight to $T^{-i}(H)$ for every $i \geq 0$; (B) if μ_∞ gives zero weight to the discontinuity set of \mathring{T} , then μ_∞ is an invariant measure for T .

Since T is assumed to be nonsingular with respect to μ , the associated transfer operator \mathcal{L} acting on $L^1(\mu)$ is well-defined. As before, we also consider the transfer operator $\mathring{\mathcal{L}}$ for the system with hole defined by $\mathring{\mathcal{L}}^t f = 1_{\mathring{X}^0} \mathcal{L}(f 1_{\mathring{X}^{t-1}})$ for all $t \geq 1$.

(A) First note that $\mathring{T}_*^t \mu$ is supported on $X \setminus H$ for each $t \geq 0$ so that $\mu_\infty(H) = 0$. Now define $g_k = \frac{\mathring{\mathcal{L}}^{k-1} 1}{\mu(\mathring{X}^k)}$ for $k \geq 1$. For each $i \geq 1$, $k \geq 1$, we have

$$\beta_{k+i-1} \cdots \beta_k = \frac{\int_X \mathring{\mathcal{L}}^i(\mathring{\mathcal{L}}^{k-1} 1) d\mu}{\mu(\mathring{X}^k)} = \int_{\mathring{X}^i} g_k d\mu = 1 - \int_{\bigcup_{j=0}^i T^{-j}(H)} g_k d\mu,$$

since $\int_{X \setminus H} g_k = 1$. Choose $\lambda \in (0, 1)$ and let A_λ be as defined in Lemma 3.8. Given $i \geq 1$, let

$$A_\lambda^i = A - [0, i-1]$$

denote the translates of elements of A_λ by some integer at most $i-1$. Note that A_λ^i still has frequency zero. Then since $\liminf_{k \rightarrow +\infty} \beta_{k+i-1} \cdots \beta_k \geq \lambda^i$ as long as we avoid the exceptional set A_λ^i , we have

$$\limsup_{\substack{k \rightarrow \infty \\ k \notin A_\lambda^i}} \int_{\bigcup_{j=0}^i T^{-j}(H)} g_k d\mu \leq 1 - \lambda^i.$$

Since $\mu_{t_j} = \frac{1}{t_j} \sum_{k=0}^{t_j-1} g_k \mu$ and A_λ^i has zero density, we conclude $\mu_\infty \left(\bigcup_{j=0}^i T^{-j}(H) \right) \leq 1 - \lambda^i$. Since λ

is arbitrary, it must be that $\mu_\infty \left(\bigcup_{j=0}^i T^{-j}(H) \right) = 0$. Since i is arbitrary, μ_∞ must be supported on \mathring{X}^∞ , and therefore is singular with respect to μ by assumption (1).

(B) For any test function $\varphi \in C^0(X)$, we have

$$\mathring{\mathcal{L}} \mu_{t_j}(\varphi) = \frac{1}{t_j} \sum_{k=0}^{t_j-1} \frac{\mathring{\mathcal{L}}^{k+1} 1(\varphi)}{\mu(\mathring{X}^k)} = \frac{1}{t_j} \sum_{k=0}^{t_j-1} \frac{\mathring{\mathcal{L}}^{k+1} 1(\varphi)}{\mu(\mathring{X}^{k+1})} \beta_k = \frac{1}{t_j} \sum_{k=0}^{t_j-1} \frac{\mathring{\mathcal{L}}^k 1(\varphi)}{\mu(\mathring{X}^k)} \beta_{k-1} + \frac{1}{t_j} \left(\frac{\mathring{\mathcal{L}}^{t_j} 1(\varphi)}{\mu(\mathring{X}^{t_j-1})} - \varphi \right).$$

Again choose $\lambda \in (0, 1)$, define A_λ as in Lemma 3.8, and let $A_\lambda^c = \mathbb{N} \setminus A_\lambda$. Thus, assuming that μ_∞ gives zero measure to the discontinuity set of \mathring{T} , we can pass to the limit to obtain

$$\begin{aligned} |\mathring{\mathcal{L}} \mu_\infty(\varphi) - \mu_\infty(\varphi)| &\leq \lim_{j \rightarrow +\infty} \frac{1}{t_j} \sum_{\substack{k=0 \\ k \in A_\lambda}}^{t_j-1} \frac{|\mathring{\mathcal{L}}^{k+1} 1(\varphi)|}{|\mathring{\mathcal{L}}^{k+1} 1|_1} |\beta_{k-1} - 1| + \frac{1}{t_j} \left(\frac{|\varphi|_\infty |\mathring{\mathcal{L}}^{t_j} 1|_1}{|\mathring{\mathcal{L}}^{t_j-1} 1|_1} + |\varphi|_\infty \right) \\ &\quad + \lim_{j \rightarrow +\infty} \frac{1}{t_j} \sum_{\substack{k=0 \\ k \in A_\lambda^c}}^{t_j-1} \frac{|\mathring{\mathcal{L}}^{k+1} 1(\varphi)|}{|\mathring{\mathcal{L}}^{k+1} 1|_1} |\beta_{k-1} - 1| \\ &\leq |\varphi|_\infty (1 - \lambda) \end{aligned}$$

where we have used the fact that the set of k for which $\beta_k \leq \lambda$ has density 0, $\beta_k \leq 1$, and $|\mathring{\mathcal{L}}^{k+1} 1(\varphi)| / |\mathring{\mathcal{L}}^{k+1} 1|_1 \leq |\varphi|_\infty$ for each k . Since λ is arbitrary, we must have $\mathring{\mathcal{L}} \mu_\infty(\varphi) = \mu_\infty(\varphi)$.

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