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Simon Conan, Etienne Thoret, Mitsuko Aramaki, Olivier Derrien, Charles
Gondre, Richard Kronland-Martinet, Sølvi Ystad

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NAVIGATING IN A SPACE OF SYNTHESIZED INTERACTION-SOUNDS: RUBBING, SCRATCHING AND ROLLING SOUNDS

S. Conan, E. Thoret, M. Aramaki, O. Derrien, C. Gondre, R. Kronland-Martinet, S. Ystad[*]

Laboratoire de Mécanique et d'Acoustique

Aix-Marseille Univ.

Centrale Marseille

Marseille, France

`conan@lma.cnrs-mrs.fr`

ABSTRACT

In this paper, we investigate a control strategy of synthesized interaction-sounds. The framework of our research is based on the action/object paradigm that considers that sounds result from an action on an object. This paradigm presumes that there exists some sound invariants, i.e. perceptually relevant signal morphologies that carry information about the action or the object. Some of these auditory cues are considered for rubbing, scratching and rolling interactions. A generic sound synthesis model, allowing the production of these three types of interaction together with a control strategy of this model are detailed. The proposed control strategy allows the users to navigate continuously in an "action space", and to morph between interactions, e.g. from rubbing to rolling.

1. INTRODUCTION

Synthesis of everyday sounds is still a challenge and especially the control of sound synthesis processes. Indeed, it is of interest to control intuitively the sounds obtained with a synthesis model, i.e. to be able to create sounds that carry a specific evocation, information. To achieve this, we need to offer the users the possibility to create sounds from semantic descriptors of the sound events or from gestures. Previous studies by Hoffman and Cook [1] have approached this problem by proposing intuitive control of sounds through acoustic descriptors or features (so-called feature synthesis) and an intuitive synthesizer of impact sounds has been developed by Aramaki et al. [2] to control the perceived material, size and shape of impacted objects.

This paper is devoted to the synthesis of continuous-interaction sounds. By continuous interaction we mean any kind of friction [3] or rolling. In this paper, we look at a subset of continuous interaction sounds: rubbing, scratching and rolling. Synthesis models for such sounds have already been proposed in previous studies, some based on physical modeling or physically informed considerations [4, 5, 6, 7, 8], others on analysis-synthesis schemes [9, 10]. The synthesizer proposed in [2] allows the user to control the perceived impacted material sounds and to morph continuously from one material to another (e.g. from glass to metal trough a continuum of ambiguous materials). We would like to control the synthesis of interaction sounds in the same way, (e.g. being able to synthesize a rubbing sound and slowly transform it into a rolling one). Such

a tool could be useful for sound design and at a more fundamental level be used to study sound perception as done by Aramaki et al. [11] and Micoulaud et al. [12].

We adopted the {action/object} paradigm in order to propose such a control strategy of interaction-sounds synthesis. This paradigm is based on the assumption that a sound results from an action on an object, e.g. "plucking a metal string" or "hitting a wood plate", as proposed in [4]. The underlying hypothesis is the existence of sound invariants, i.e. sound signal morphologies that carry specific meanings about the action or the object. For instance, a string produces a particular spectrum that allows the listener to recognize it, even if it is bowed (violin), plucked (guitar) or hit (piano). Likewise, it is possible to recognize that a cylinder bounces even if it is made of glass, wood or metal [13]. This is inspired from the ecological theory of visual perception as proposed by Gibson [14] (for a more accessible introduction to Gibson's theory, refer to [15]), and formalized as an explanation of sound sources recognition by McAdams [16].

Some studies have already identified sound invariants. For instance, to a certain extent, it has been shown that impact sounds contain sufficient information to discriminate the materials of impacted objects [17]. Other authors showed that the perceived material is mainly related to the damping of spectral components [18, 19, 20] and to sound roughness [11]. A study by Warren and Verbrugge revealed that from the rhythm of a series of impact sounds, it is possible to predict if a glass will break or bounce [21].

In this paper, we will focus on sound invariants related to the actions of rubbing, scratching and rolling to propose an intuitive control of these actions. For this purpose we will first identify some invariants related to the auditory perception of these interactions. Then, we will describe how the concerned interactions can be reproduced by synthesis. Further, a control strategy of the proposed synthesis model is presented. Then in the last section we give some general conclusion and propose future research.

2. INVARIANTS

In this section, we examine sound signal morphologies that are related to the perception of rubbing, scratching and rolling. The results on rubbing and scratching sounds are based on preliminary studies by Conan et al. [22], which are briefly recalled below[1].

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[1]The related paper and all the stimuli used in the experiments and the listening test interfaces are available on `http://www.lma.cnrs-mrs.fr/~kronland/RubbingScratching`
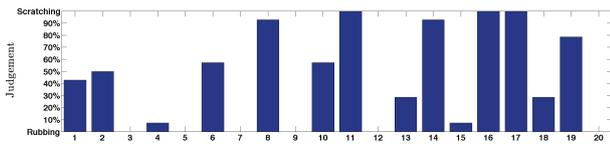
Figure 1: *Results of the experiment with recorded sounds. On the X-axis, the number of the sound. The Y-axis represents the percentage of association to scratching for each sound.*
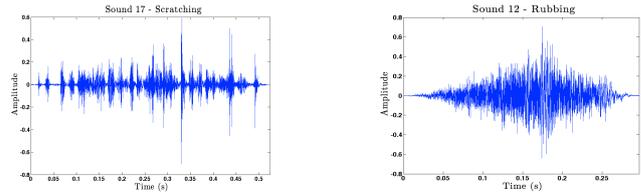

Figure 2: *Left: A recorded sound 100% associated to scratching. Right: A recorded sound 100% associated to rubbing.*

Figure 3: *Left: A nail which is scratching a surface. Right: A finger which is rubbing a surface (asperities which are not "seen" by the finger are circled). For clarity, dimensions of the asperities are exaggerated on the scheme.*

2.1. Rubbing and Scratching Interactions

To our knowledge, the auditory ability to distinguish rubbing sounds from scratching sounds has not been investigated yet. A listening test with recorded sounds was therefore set up to identify signal properties that enable to distinguish between rubbing and scratching from the auditory point of view.

- *Participants.* 14 partipants took part in the experiment (4 women, 10 men, mean aged=30.64 years ; SD=12.05).

- *Stimuli.* Twenty monophonic recordings of friction sounds on ten different surfaces (2 different sounds on each surface, one by acting with the fingertips, the other with the nails) were recorded with a Roland R05 recorder at 44.1 kHz sampling rate.

- *Apparatus.* The listening test interface was designed using MAX/MSP[2] and sounds were presented through Sennheiser HD-650 headphones.

- *Procedure.* Subjects were placed at a desk in front of a computer screen in a quiet room. They were informed that they were to classify twenty sounds in two categories, rubbing and scratching. Before the session started, the twenty stimuli were played once. Then, the subjects had to evaluate the evoked action for each sound by classifying each sound in one of the two categories "rub" or "scratch" using a drag and drop graphical interface. They could listen to each stimulus as often as they desired. No time constraints were imposed and sounds were placed in a random position on the graphical interface across subjects.

Results are presented in figure 1. Three sounds were 100% associated to scratching (number 11, 16, 17) and six sounds were 100% associated to rubbing (number 3, 4, 7, 9, 12, 20).

Qualitative signal analysis led us to suppose that, if considering that friction sounds result from successive impacts of a plectrum on an irregular surface [4, 6], rubbing sounds result from a higher temporal density of impacts than scratching sounds (*i.e.* impacts occur more frequently in rubbing sounds than in scratching sounds). A sound signal that is 100% associated to scratching and of a sound that is 100% associated to rubbing are plotted in figure 2 (left and right parts respectively). The differences can be explained as follows. On one hand, scratching a surface results from a deeper scan of the surface, which imply encountering each surface's irregularity one after another and more intensely. On the other hand, rubbing a surface results from a soft contact between the finger and several surface irregularities, implying a more noisy and stationnary sound because of the higher density of impacts (see figure 3).

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[2]http://cycling74.com/

To validate the previous hypothesis made from the qualitative observations, an experiment with controlled synthetic stimuli was set up. The stimuli were synthetic friction sounds generated with different impact densities to test the hypothesis. The synthesis model was based on the pioneering work of Gaver [4], improved by Van den Doel et al. [6] and consisted in simulating the sound as a result of successive micro-impacts of a plectrum on the asperities of a surface. The successive impacts were modeled by a low pass filtered noise with a cutoff frequency related to the gesture velocity while the roughness of the surface was controlled by the nature of the noise. By generating low to high density impact series, continuous transitions between scratching and rubbing were then generated. The synthesized stimuli were presented randomly to the subjects who were to associate each stimulus to one of the two interaction categories (rubbing or scratching). The results confirmed the empirical hypothesis: small impact densities are associated to scratching and high impact densities to rubbing (more details in [22]).

In summary, rubbing and scratching interactions can be simulated by series of impacts on a surface. One possible invariant (among others), that contributes to the discrimination between these two interactions is the temporal density of the impacts, i.e. the more (respectively the less) impacts occur in the sound signal, the more the sound is related to rubbing (respectively to scratching).

2.2. Rolling Interaction

The aim of the previous section was to highlight the perceptual information which is relevant to perceptually differentiate rubbing interactions from scratching, i.e. the acoustical morphology which characterized these interactions. In this section we address the same question concerning the auditory perception of rolling: what is the signal information which allows us to recognize a rolling ball? To answer this question, we investigate a physical model of rolling.

In the literature, most authors consider that the physics of a rolling ball is close to the physics of a bouncing ball. The models
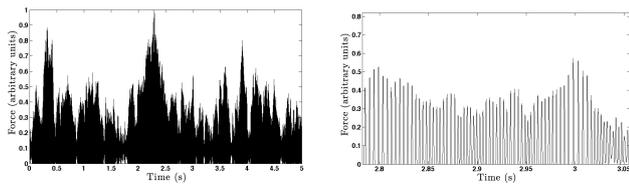
Figure 4: *Simulated rolling force (a zoom on the right figure). The force parameters are $\alpha = \gamma = 3/2$, $k = 10^7 N/m^{3/2}$, $\lambda = 10^7 Ns/m^{5/2}$. The ball has a mass of $5g$ and a velocity of $0.5m/s$. The surface is assumed to be fractal with $\beta = 1.2$ and a maximum amplitude of $10^{-9}m$.*

Figure 5: *Left: Notations for the impacts' amplitudes and time interval series. Right: Fourier spectrum of Δ_T series (Fourier spectrum of A series is similar).*

Figure 6: *Cross-correlation between A^n and Δ_T^n.*

generally take into account a nonlinear sphere-plane interaction that relates the force f applied to the sphere to the penetration x and the penetration velocity $\dot{x}$ of the sphere into the surface:

$$f(x, \dot{x}) = \begin{cases} kx^\alpha + \lambda x^\gamma \dot{x} & , x > 0 \\ 0 & , x \le 0 \end{cases} \quad (1)$$

where k is the stiffness and λ damping weight of the force. By taking into account the effect of the gravity on the ball, this model can simulate the behavior of a bouncing ball (see for instance [23], and [24] for an application to bouncing sound synthesis).

Studies that addressed the simulation of a rolling ball [7, 8] or a rolling wheel [25] consider that the rolling object moves along an irregular surface. The height of this irregular surface is added as a perturbation to the penetration term x in equation 1. It can be seen as a bouncing ball whose height changes randomly.

We simulated the rolling of a ball over an irregular surface (we do not consider here the vibration of the surface as in [7]) thanks to a Fourth-Order Runge-Kutta scheme (as described in [26], where numerical issues of the model (1) are studied), instead of the K-method [27]. An example of a simulated nonlinear interaction force f is plotted in figure 4, where the used parameters are listed[3]. In the numerical model, the surface is modeled by a noise with a specific spectrum adjusted according to tribological observations [30] and phenomenological considerations [6, 31]. In practice, the spectrum is $S(\omega) \propto 1/\omega^\beta$. Also called fractal noise, such a spectrum models accurately most surfaces. β controlled the perceived roughness, and the amplitude was normalized to provide a maximum of $10^{-9}m$.

From informal listening tests, it has been concluded that the force f evokes a rolling object such as a small hard marble. One can note that this force can be considered as a series of impacts, see figure 4. This assumption has already been exploited by Hermes [5] for sound synthesis purposes and by Lagrange et al. [9] in an analysis-synthesis context. The first observation is that the temporal structure of the impact series seems to follow a specific pattern (see figure 4 left). Let us consider (A, Δ_T) as time series, with A^n and Δ_T^n respectively the amplitude of the nth impact and the time interval between the nth and the $(n + 1)^{th}$ impact (see figure 5, left). The Fourier spectrum of Δ_T time series is plotted in figure 5, right. One can note that this time series has a strong auto-correlation, i.e. that successive impacts have strong mutual dependencies. The amplitude series A exhibits the same behavior. A and Δ_T are also strongly cross-correlated as can be seen in figure 6. These observations are coherent with the physics of

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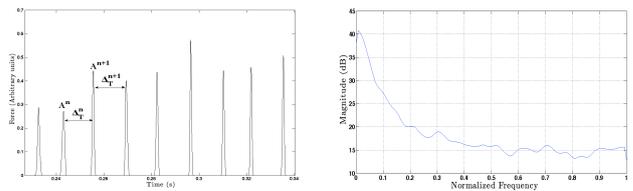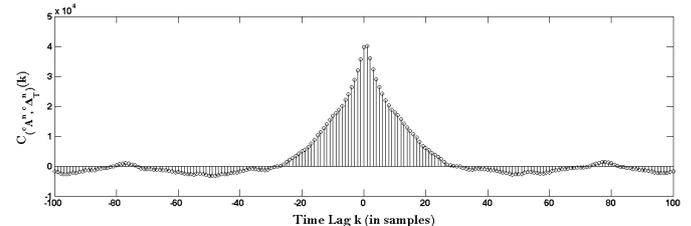[3]When $\alpha = \gamma$, the model is the Hunt and Crossley one [28, 29]

a bouncing ball (recalling that the rolling model is derived from a bouncing model) which relates the impacts to each other.

Another important aspect is the fact that the contact time of the impact depends on the impact velocity, related to the amplitude of the impact in f. This dependency has been studied by several authors [32, 24] and seems to be an important auditory cue that is responsible for the rolling interaction evocation. Indeed, to informally test this assumption, we first detected all the impact time locations and amplitudes in the simulated force in figure 4. We then created two modified versions of this force signal by replacing the impact windows by a raised cosine window [33] that fits well the original simulated impacts, one with a duration depending on the impact amplitude, the other with a fixed duration. For different fixed impact durations, it was always found that the version with varying impact durations depending on the impact amplitudes clearly produced the most realistic evocations of the rolling interaction, which confirms previous findings.

2.3. Invariants Related to Continuous Interactions

As described in this section, rubbing, scratching and rolling interactions can all be represented as impacts series. Thanks to listening tests, it was shown that the temporal distribution of successive impacts convey the information of rubbing or scratching sounds. For the rolling interaction, the temporal structure of the impacts (their correlations and statistics) may also convey the rolling information. As opposed to rubbing and scratching interactions, a specific relation between the impact durations and amplitudes seems to be an important signal morphology that conveys information about the rolling interaction. These assumptions need to be formally verified but this is out of the scope of this paper. In the next section, we will propose a model to simulate these impact series.

3. IMPLEMENTATION OF A GENERIC MODEL OF CONTINUOUS INTERACTIONS

As described in the previous section, the force applied to a surface carries the information about the type of interaction. For rubbing, scratching and rolling, the forces are series of impacts, with specific relations between each other. The aim of this section is to propose a synthesis process which is generic enough to simulate different continuous interactions and navigate continuously between them. In order to cohere with the {action/object} paradigm proposed in the section 1, the force previously described is be modeled as the source part of a source-filter model. In fact, by convolving the force with an impulse response of a resonant surface, the resulting sound conveys the information on both the action and the object. The object properties such as material and shape are included into the filter part of the source-filter model (as proposed by Gaver [4]). The implemented filters are described in [34]. The control strategy is similar to the one proposed by Aramaki et al. [2], which ables to navigate continuously between the perceived materials. Aspects linked to the perceived object are not presented here.

3.1. Modeling the Source Signal

As presented in the section 2, the force that conveys the information about the interaction type (rolling, scratching or rubbing) can be considered as impacts series. The specific behavior of the time series A^n and Δ_T^n, respectively representing the amplitude of the nth impact and the time interval between the nth and the $(n+1)^{th}$ impact, seems to be an important cue to discriminate these interactions. We will experimentally characterize the series ${}^cA$ or ${}^c\Delta_T$, centered version of A and Δ_T, in order to propose a synthesis scheme.

As pointed out in 2.2, A and Δ_T can be strongly autocorrelated (so are ${}^cA$ or ${}^c\Delta_T$). We consider these two time series as autoregressive moving average (ARMA) processes. To characterize their behavior, these series are whitened (the method that we use is described in [35]), and we experimentally noted that no more than 2 poles - 2 zeros whitening filters are needed. So let X be one of the two processes ${}^cA$ or ${}^c\Delta_T$, we can write in the z domain:

$$X(z) \approx H_X(z)\tilde{X}(z) \quad , \quad H_X(z) = \frac{B(z)}{A(z)} \qquad (2)$$

where $B(z) = 1 + \sum_{i=1}^p b_i z^{-i}$, $A(z) = \sum_{i=0}^q a_i z^{-i}$ and $\tilde{X}$ is the whitened version of X. Then, as $\tilde{X}$ is white, we can model its probability density function, and transform $\tilde{X}$ into W which follows a uniform law, thanks to the inverse transform sampling method (*i.e.* $W = F_X(\tilde{X})$, where $F_X(.)$ is the cumulative distribution function of $\tilde{X}$). The analysis scheme is presented in figure 7.

Results of the analysis on Δ_T time series (from the force f computed in the section 2.2) are presented on top and middle part of figure 8 (results on A series are similar and therefore not displayed). The top figure displays the autocorrelation of ${}^c\tilde{\Delta}_T$, whitened version of ${}^c\Delta_T$, showing that a 2 poles - 2 zeros whitening filter is well suited to whiten the time series. The probability density of ${}^c\tilde{\Delta}_T$ is plotted in the middle figure, showing that for roll force, ${}^c\tilde{\Delta}_T$ and ${}^c\tilde{A}$ can be modeled as gaussian processes. in the bottom figure, $C_{(W_A^n, W_{\Delta_T}^n)}(k)$ the cross correlation between W_A and W_{Δ_T} is displayed.
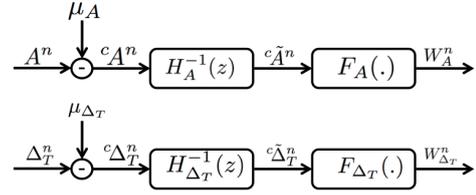

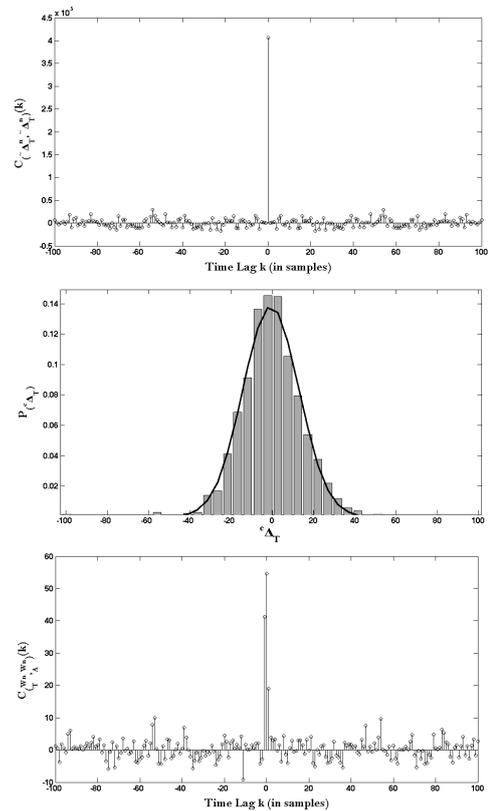Figure 7: *Analysis scheme of A^n and Δ_T^n series.*

Figure 8: Top: *${}^c\tilde{\Delta}_T$ autocorrelation.* Middle: *probability density of ${}^c\tilde{\Delta}_T$ (grey bars are the measure and the black line is the gaussian fit).* Bottom: *cross correlation between W_A and W_{Δ_T}.*
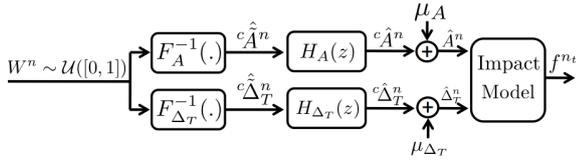
Figure 9: *Proposed synthesis scheme. The hat stands for the variables that are estimated. W is white and follows a uniform distribution on $[0, 1]$. First $\hat{A}^n$ and $\hat{\Delta}_T^n$ are estimated which allow the nonlinear force f^{n_t} construction. n_t is the time index, while n is the impact number index.*

Given the long term autocorrelation of ^cA and of $^c\Delta_T$, it seems reasonable to assume that W_A^{n-1} has a poor influence on W_A^n (and respectively that $W_{\Delta_T}^{n-1}$ has a poor influence on $W_{\Delta_T}^n$), and then that $C_{(W_A^n, W_{\Delta_T}^n)}(k)$ is proportional to $\delta(k)$, the Dirac function that equals 1 if $k = 0$ and equals 0 elsewhere. This led us us propose the synthesis scheme presented in figure 9 where, according to the previous observations on $C_{(W_A^n, W_{\Delta_T}^n)}(k)$, let us consider that we can start from the same white and uniform process W to synthesize $\hat{A}$ and $\hat{\Delta}_T$ (the hat stands for the 'estimation', unlike no hat stands for 'measures'). The generation process of $\hat{A}$ and $\hat{\Delta}_T$ is summarized by the following equations:

$$\begin{cases} W \sim \mathcal{U}([0, 1]) \text{ and } S_{WW}(\omega) = B \\ \hat{A}(n) = \mu_A + \left[F_A^{-1}(W) * h_A\right](n) \\ \hat{\Delta}_T(n) = \mu_{\Delta_T} + \left[F_{\Delta_T}^{-1}(W) * h_{\Delta_T}\right](n) \end{cases} \quad (3)$$

where $h_X(.)$ is the impulse response of the filter $H_X(.)$, $S_{WW}(\omega)$ is the power spectral density of W and B is a constant.

To finally reconstruct the whole force a raised cosine impact model is used [33]:

$$f^{n_t}(n_t) = \begin{cases} \frac{F_{\max}}{2^\xi} \left[1 - \cos\left(\frac{2\pi n_t}{N}\right)\right]^\xi & , n_t \in [\![0, N]\!] \\ 0 & , \text{otherwise} \end{cases} \quad (4)$$

An additional exponent ξ is used that controls the sharpness of the impact and then helps to better fit the measured impacts. $F_{\max}$ is the impact amplitude and N is the duration of the impact. N is a function of the impact amplitude:

$$N = K \times A^{-\nu} \quad (5)$$

where K is a constant depending on the mass of the ball and on the stiffness k (see eq. 1), A is the impact amplitude and ν is a positive exponent (equals to $1/5$ when $\gamma = 0$ and $\alpha = 3/2$ in eq. 1, as shown by Chaigne and Doutaut [32]).

This model also reproduces rubbing and scratching sounds as the synthesized one in [22]. Indeed, in the perceptual experiment conducted in this paper, we do not consider correlations between impacts, i.e. in the present model this corresponds to setting to zero the (a_i, b_i) coefficients of the filters $H_A(.)$ and $H_{\Delta_T}(.)$ for $i \geq 1$ (eq. 2). The amplitude series A follows a gaussian process and Δ_T follows an exponential distribution. The prototypes of each interactions rubbing, scratching and rolling, and a strategy of navigation between these prototypes will be detailed in the next section.

4. CONTROL STRATEGY

In this section, we will detail our control strategy for continuous navigation between rubbing, scratching and rolling. First, the prototypes of each interaction will be given. Then, we will present the navigation strategy between these prototypes.

Let us first recall the low level parameters of the model:

- **Impact Model.** Two parameters control the impact duration (eq. 5) in the chosen impact model (eq. 4): K and ν.

- **Probability Density.** The probability density is defined as a list of discrete values, which are used to derive the cumulative distribution function F_X (cumulative sum of the probability density), leading to obtain $^c\hat{A}$ and $^c\hat{\Delta}_T$ series (see figure 9). These lists, respectively for $^c\hat{A}$ series and $^c\hat{\Delta}_T$ series will be noted $\mathbb{P}_A$ and $\mathbb{P}_{\Delta_T}$.

- **Filters.** As pointed out in 3, no more than 2 poles and 2 zeros are needed for the filters. Then, each filter will be described by a set of 5 coefficients $(a_0, a_1, a_2, b_1, b_2)$, each set will be noted $\mathcal{C}_A$ and $\mathcal{C}_{\Delta_T}$, respectively for $H_A(.)$ and $H_{\Delta_T}(.)$.

- **Offset coefficients.** μ_A and μ_{Δ_T} respectively for $\hat{A}$ and $\hat{\Delta}_T$.

The global set of parameters will be called:
$\mathfrak{P} = \{K, \nu, \mathbb{P}_A, \mathbb{P}_{\Delta_T}, \mathcal{C}_A, \mathcal{C}_{\Delta_T}, \mu_A, \mu_{\Delta_T}\}$.

4.1. Interaction Sounds Prototypes

4.1.1. Rubbing Prototype

As showed in [22] and recalled in section 2.1, small intervals between impacts are associated to rubbing sounds (a maximum impact density implies a source signal which is a white noise). For the rubbing prototype, we then tune the parameters to obtain a gaussian white noise: setting the coefficients of $\mathcal{C}_A$ and $\mathcal{C}_{\Delta_T}$ to zero for $i \geq 1$, $\mu_A = \mu_{\Delta_T} = 0$, $\nu = 0$ and $K = 1$ (to get an impact duration of 1 sample), $\mathbb{P}_{\Delta_T}$ to get 1 impact at each sample and $\mathbb{P}_A$ following a gaussian distribution. This set of parameters will be called $\mathfrak{P}_{\text{rub}}$.

Following the model proposed by Van den Doel et al. [6], the generated source signal is lowpass filtered with a cutoff frequency which is directly related to the relative transversal velocity between the object that rubs (hand, plectrum...) and the rubbed surface. This enabled the simulation of the sound producing gesture and turned out to be an efficient synthesis strategy to convey the velocity of a human gesture when associated to a biological gesture law [36].

4.1.2. Scratching Prototype

The scratching prototype is associated to a low impact density, i.e. to high Δ_T^n values and, as proposed empirically in [22], the Δ_T series follows an exponential distribution. The rest of the parameters set is the same as in the rubbing prototype. Again, the velocity between the interacting objects controls the cutoff frequency of a lowpass filter. The parameter set will be called $\mathfrak{P}_{\text{scratch}}$.
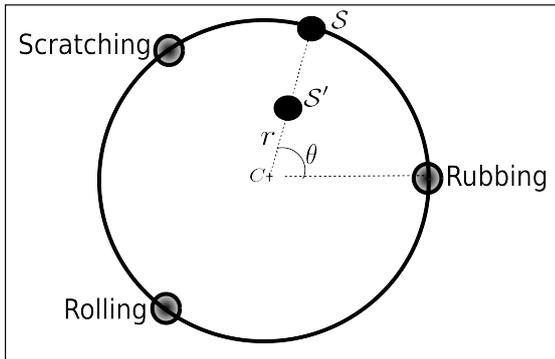
Figure 10: *Schematic control space of the interaction sound synthesizer.*

4.1.3. Rolling Prototype

In section 3, we developed our synthesis model based on an experimental analysis of a simulated rolling force. We saw that $\mathbb{P}_A$ and $\mathbb{P}_{\Delta_T}$ are gaussian distributions. The filter coefficients $\mathcal{C}_A$ and $\mathcal{C}_{\Delta_T}$ are here non zero, and we take $\nu = 1/5$. As the whole rolling model necessitates a more precise perceptual calibration that requires an individual study, we will only give qualitative remarks on the parameters' influence in this article. Presently, the complete perceptual mapping of the rolling model is not fully established and is left for further studies, so for now the synthesis parameters are controled by the following observations. As K controls the contact time, it is related to the size of the rolling ball. The mean value of $\mathbb{P}_{\Delta_T}$ is related to the perceived velocity. The filters coefficients $\mathcal{C}_A$ and $\mathcal{C}_{\Delta_T}$ are linked to surface parameters such as roughness. As in the scratching and rubbing prototypes, the resulting source signal is also lowpass filtered according to the ball's transversal velocity. The parameter set will be called $\mathfrak{P}_{\text{roll}}$.

4.2. Navigation Strategy

The control strategy that we adopt is based on the one proposed by Aramaki et al. to control the perceived material in an impact sound synthesizer [37]. In this paper, the authors proposed to allow the users to navigate continuously in a material space that comprises wood, metal and glass. The three prototypes of perceived material (wood, metal and glass), which synthesis parameters were determined thanks to behavioral and electrophysiological experiments, are placed on the border of a circle which represent the so-called material space. Hereby, the user can move a cursor in this space to get the desired impact sound, and the synthesis parameters are weighting of the three prototypes of perceived material (wood, metal and glass).

A similar control space of interaction is schematized in figure 10. On the circumference of the circle, a sound $\mathcal{S}$, characterized by its angle θ in the control space, is generated with the parameters $\mathfrak{P}_{\mathcal{S}}(\theta)$ as follows:

$$\mathfrak{P}_{\mathcal{S}}(\theta) = T(\theta)\,\mathfrak{P}_{\text{rub}} + T\left(\theta - \frac{2\pi}{3}\right)\mathfrak{P}_{\text{scratch}} + T\left(\theta - \frac{4\pi}{3}\right)\mathfrak{P}_{\text{roll}}$$
(6)

The function θ is defined as follows:

$$T(\theta) = \begin{cases} -\frac{3}{2\pi}\theta + 1 & , \quad \theta \in \left[0; \frac{2\pi}{3}\right[\\ 0 & , \quad \theta \in \left[\frac{2\pi}{3}; \frac{4\pi}{3}\right[\\ \frac{3}{2\pi}\theta - 2 & , \quad \theta \in \left[\frac{4\pi}{3}; 2\pi\right[\end{cases}$$
(7)

Inside the circle, a sound $\mathcal{S}'$, characterized by its angle θ and radius r, is generated with the parameters $\mathfrak{P}_{\mathcal{S}'}$:

$$\mathfrak{P}_{\mathcal{S}'} = (1 - r)\mathfrak{P}_C + r\mathfrak{P}_{\mathcal{S}}(\theta)$$
(8)

where

$$\mathfrak{P}_C = \frac{1}{3}\left(\mathfrak{P}_{\text{rub}} + \mathfrak{P}_{\text{scratch}} + \mathfrak{P}_{\text{roll}}\right)$$
(9)

and $\mathfrak{P}_{\mathcal{S}}(\theta)$ defined in equation 6. Even if this control space is not perceptually validated and calibrated, it yields good results (see the demonstration video on the webpage associated to this paper: `http://www.lma.cnrs-mrs.fr/~kronland/InteractionSpace`).

5. CONCLUSIONS

In this paper, we described specific signal morphologies that are related to the auditory perception of rubbing, scratching and rolling interactions. Both phenomenological considerations and physical modeling were investigated as well as qualitative signal analysis. A generic synthesis model of these three interactions that enables continuous transitions between the categories has been proposed. Although a precise perceptual calibration is still to be carried out, the control is already convincing.

Further studies will be done to expand this control space to other interactions such as nonlinear friction (squeaking, squealing...) [38]. The rubbing and scratching models are also quite simple and could be improved by an analysis method as proposed by Lagrange et al. [9]. Finally, the influence of the physical velocity profiles (rolling marble in a bowl, sliding object on an inclined plate...) is to be studied, as already done by [36] with the velocity of human gestures.

6. REFERENCES

[1] M. Hoffman and P.R. Cook, "Feature-based synthesis: Mapping acoustic and perceptual features onto synthesis parameters," in *Proceedings of the 2006 International Computer Music Conference (ICMC), New Orleans*. Citeseer, 2006, vol. 33.

[2] M. Aramaki, C. Gondre, R. Kronland-Martinet, T. Voinier, and S. Ystad, "Thinking the sounds: an intuitive control of an impact sound synthesizer," in *Proceedings of the 15th International Conference on Auditory Display, Copenhagen, Denmark May 18 - 22*, 2009.

[3] A. Akay, "Acoustics of friction," *The Journal of the Acoustical Society of America*, vol. 111, pp. 1525, 2002.

[4] W.W. Gaver, "How do we hear in the world? explorations in ecological acoustics," *Ecological psychology*, vol. 5, no. 4, pp. 285–313, 1993.

[5] D.J. Hermes, "Synthesis of the sounds produced by rolling balls," Internal IPO report no. 1226, IPO, Center for User-System Interaction, Eindhoven, The Netherlands, September 1998.

[6] K. Van Den Doel, P.G. Kry, and D.K. Pai, "Foleyautomatic: physically-based sound effects for interactive simulation and animation," in *Proceedings of the 28th annual conference on Computer graphics and interactive techniques*. ACM, 2001, pp. 537–544.

[7] M. Rath and D. Rocchesso, "Informative sonic feedback for continuous human–machine interaction—controlling a sound model of a rolling ball," *IEEE Multimedia Special on Interactive Sonification*, vol. 12, no. 2, pp. 60–69, 2004.

[8] C. Stoelinga and A. Chaigne, "Time-domain modeling and simulation of rolling objects," *Acta Acustica united with Acustica*, vol. 93, no. 2, pp. 290–304, 2007.

[9] M. Lagrange, G. Scavone, and P. Depalle, "Analysis/synthesis of sounds generated by sustained contact between rigid objects," *IEEE Transactions on Audio, Speech, and Language Processing*, vol. 18, no. 3, pp. 509–518, 2010.

[10] J.S. Lee, P. Depalle, and G. Scavone, "Analysis/synthesis of rolling sounds using a source-filter approach," in *13th Int. Conference on Digital Audio Effects (DAFx-10), Graz, Austria*, 2010.

[11] M. Aramaki, M. Besson, R. Kronland-Martinet, and S. Ystad, "Timbre perception of sounds from impacted materials: behavioral, electrophysiological and acoustic approaches," *Computer Music Modeling and Retrieval. Genesis of Meaning in Sound and Music*, pp. 1–17, 2009.

[12] J.M. Micoulaud-Franchi, M. Aramaki, A. Merer, M. Cermolacce, S. Ystad, R. Kronland-Martinet, and J. Vion-Dury, "Categorization and timbre perception of environmental sounds in schizophrenia," *Psychiatry research*, vol. 189, no. 1, pp. 149–152, 2011.

[13] G. Lemaitre and L.M. Heller, "Auditory perception of material is fragile while action is strikingly robust," *Journal of the Acoustical Society of America*, vol. 131, pp. 1337–1348, February 2012.

[14] J. J. Gibson, "The ecological approach to visual perception.," 1979.

[15] C.F. Michaels and C. Carello, *Direct perception*, Prentice-Hall Englewood Cliffs, NJ, 1981.

[16] S.E. McAdams and E.E. Bigand, *Thinking in sound: The cognitive psychology of human audition.*, Oxford Science Publications, 1993.

[17] R. P. Wildes and W. A. Richards, "Recovering material properties from sound," *Natural computation*, pp. 356–363, 1988.

[18] S. Tucker and G.J. Brown, "Investigating the perception of the size, shape and material of damped and free vibrating plates," *University of Sheffield, Department of Computer Science Technical Report CS-02-10*, 2002.

[19] R.L. Klatzky, D.K. Pai, and E.P. Krotkov, "Perception of material from contact sounds," *Presence: Teleoperators & Virtual Environments*, vol. 9, no. 4, pp. 399–410, 2000.

[20] B.L. Giordano and S. McAdams, "Material identification of real impact sounds: Effects of size variation in steel, glass, wood, and plexiglass plates," *The Journal of the Acoustical Society of America*, vol. 119, pp. 1171–1181, 2006.

[21] W.H. Warren and R.R. Verbrugge, "Auditory perception of breaking and bouncing events: A case study in ecological

acoustics.," *Journal of Experimental Psychology: Human Perception and Performance*, vol. 10, no. 5, pp. 704–712, 1984.

[22] S. Conan, M. Aramaki, R. Kronland-Martinet, E. Thoret, and S. Ystad, "Perceptual differences between sounds produced by different continuous interactions," in *Acoustics 2012, Nantes*, 23-27 april 2012.

[23] E. Falcon, C. Laroche, S. Fauve, and C. Coste, "Behavior of one inelastic ball bouncing repeatedly off the ground," *The European Physical Journal B-Condensed Matter and Complex Systems*, vol. 3, no. 1, pp. 45–57, 1998.

[24] F. Avanzini and D. Rocchesso, "Modeling collision sounds: Non-linear contact force," in *Proceedings of the COST-G6 Conference Digital Audio Effects (DAFx-01)*. Citeseer, 2001, pp. 61–66.

[25] A. Nordborg, "Wheel/rail noise generation due to nonlinear effects and parametric excitation," *The Journal of the Acoustical Society of America*, vol. 111, pp. 1772, 2002.

[26] S. Papetti, F. Avanzini, and D. Rocchesso, "Numerical methods for a nonlinear impact model: a comparative study with closed-form corrections," *IEEE Transactions on Audio, Speech, and Language Processing*, vol. 19, no. 7, pp. 2146–2158, 2011.

[27] G. Borin, G. De Poli, and D. Rocchesso, "Elimination of delay-free loops in discrete-time models of nonlinear acoustic systems," *IEEE Transactions on Speech and Audio Processing*, vol. 8, no. 5, pp. 597–605, 2000.

[28] K.H. Hunt and F.R.E. Crossley, "Coefficient of restitution interpreted as damping in vibroimpact," *J. Appl. Mech.*, 1975.

[29] D.W. Marhefka and D.E. Orin, "A compliant contact model with nonlinear damping for simulation of robotic systems," *IEEE Transactions on Systems, Man and Cybernetics, Part A: Systems and Humans*, vol. 29, no. 6, pp. 566–572, 1999.

[30] H. Ben Abdelounis, A. Le Bot, J. Perret-Liaudet, and H. Zahouani, "An experimental study on roughness noise of dry rough flat surfaces," *Wear*, vol. 268, no. 1, pp. 335–345, 2010.

[31] D. Rocchesso and F. Fontana, *The sounding object*, Mondo estremo, 2003.

[32] A. Chaigne and V. Doutaut, "Numerical simulations of xylophones. i. time-domain modeling of the vibrating bars," *Journal of the Acoustical Society of America*, vol. 101, no. 1, pp. 539–557, 1997.

[33] S. Bilbao, *Numerical Sound Synthesis: Finite Difference Schemes and Simulation in Musical Acoustics*, John Wiley & Sons, 2009.

[34] M. Mathews and J.O. Smith, "Methods for synthesizing very high q parametrically well behaved two pole filters," in *Proceedings of the Stockholm Musical Acoustics Conference (SMAC 2003)(Stockholm), Royal Swedish Academy of Music (August 2003)*, 2003.

[35] B. Ninness, A. Wills, and S. Gibson, "The university of newcastle identification toolbox (unit)," in *Proc. IFAC World Congress*, 2005.

[36] E. Thoret, M. Aramaki, R. Kronland-Martinet, J.L. Velay, and S. Ystad, "From shape to sound: sonification of two dimensional curves by reenaction of biological movements," in *9th International Symposium on Computer Music Modeling and Retrieval, London*, 2012.

[37] M. Aramaki, M. Besson, R. Kronland-Martinet, and S. Ystad, "Controlling the perceived material in an impact sound synthesizer," *IEEE Transactions on Audio, Speech, and Language Processing*, vol. 19, no. 2, pp. 301–314, 2011.

[38] F. Avanzini, S. Serafin, and D. Rocchesso, "Interactive simulation of rigid body interaction with friction-induced sound generation," *Speech and Audio Processing, IEEE Transactions on*, vol. 13, no. 5, pp. 1073–1081, 2005.