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► **To cite this version:**

Florent Berthelin, Stéphane Junca. Averaging lemmas with a force term in the transport equation. *Journal de Mathématiques Pures et Appliquées*, 2010, 93 (2), pp.113-131. 10.1016/j.matpur.2009.10.009 . hal-00874401

**HAL Id: hal-00874401**

**<https://hal.science/hal-00874401v1>**

Submitted on 3 Sep 2019

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# Averaging lemmas with a force term in the transport equation

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## Abstract

We obtain several averaging lemmas for transport operator with a force term. These lemmas improve the regularity yet known by not considering the force term as part of an arbitrary right-hand side. Two methods are used: local variable changes or stationary phase. These new results are subjected to two non-degeneracy assumptions. We characterize the optimal conditions of these assumptions to compare the obtained regularities according to the space and velocity variables. Our results are mainly in  $L^2$ , and for constant force, in  $L^p$  for  $1 < p \leq 2$ .

## Résumé

Nous obtenons plusieurs lemmes de moyenne pour des équations de transport avec un terme de force. Ces résultats améliorent la régularité connue en ne considérant pas le terme de force comme terme source arbitraire. Deux techniques sont utilisées : des changements de variables locaux ou des phases stationnaires. Ces résultats sont quantifiées par deux hypothèses de non dégénérescence. Nous caractérisons les conditions optimales de ces hypothèses pour comparer les régularités obtenues, par rapport aux variables d'espace et de vitesse. Les résultats sont principalement dans  $L^2$ , et pour le cas constant, dans  $L^p$  pour  $1 < p \leq 2$ .

*MSC:* 35B65; 42B20; 82C40

*Keywords:* Averaging lemma; Force term; Kinetic equation; Stationary phase; Non-degeneracy conditions; Fourier series; Hardy space

## 1. Introduction

Averaging lemma is a major tool to get compactness from a kinetic equation ([7], etc.). Such results have been used in a lot of papers during these last years. Among this literature, an important result using an averaging lemma as a key argument is the proof of the hydrodynamic limits of the Boltzmann or BGK equations to the incompressible Euler or Navier–Stokes equations [16]. Another major application consists in obtaining the compactness for nonlinear scalar conservation laws (in [25]) which allows, for instance, to study the propagation of high frequency waves [6].

Basically, averaging lemma is a result which says that the macroscopic quantities  $\int f(t, x, v)\psi(v) dv$  have a better regularity with respect to  $(t, x)$  than the microscopic quantity  $f(t, x, v)$  where  $f$  is solution of a kinetic equation.

For example, in [9] and [2], the following result is established.

**Theorem** (DiPerna, Lions, Meyer–Bézard). *Let  $f, g_k \in L^p(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M)$  with  $1 < p \leq 2$  such that*

$$\partial_t f + \operatorname{div}_x [a(v)f] = \sum_{|k| \leq m} \partial_v^k g_k, \quad (1.1)$$

*with  $a \in W^{m,\infty}(\mathbb{R}^M, \mathbb{R}^N)$  for  $m \in \mathbb{N}$ . Let  $\psi \in W^{m,\infty}(\mathbb{R}^M)$  with compact support. Let  $A > 0$  such that the support of  $\psi$  is included in  $[-A, A]^M$ . We assume the following non-degeneracy for  $a(\cdot)$ : there exists  $0 < \alpha \leq 1$  and  $C > 0$  such that for any  $(u, \sigma) \in S^N$  and  $\varepsilon > 0$ ,*

$$\operatorname{meas}(\{v \in [-A, A]^M; u - \varepsilon < a(v) \cdot \sigma < u + \varepsilon\}) \leq C\varepsilon^\alpha.$$

*Then*

$$\rho_\psi(t, x) = \int_{\mathbb{R}^M} f(t, x, v) \psi(v) dv$$

*is in  $W^{s,p}(\mathbb{R}_t \times \mathbb{R}_x^N)$  where  $s = \frac{\alpha}{(m+1)p'}$ ,  $p'$  being the conjugated exponent for  $p$ .*

Regarding Eq. (1.1), the obtained regularity is proved to be optimal, see [23] and [24]. In [11], the gain of a half-derivative in  $L^2$  context was proved as optimal. A study in the case of a full derivative with respect to  $x$  in the second member is done in [21]. We also refer to [10] and [4] for other results about averaging lemmas. Regularity of  $f$  itself is also challenging, for example by assuming some regularity with respect to  $v$ , see [3,18] and [1] for such results.

Theorem here above says for example with  $m = 1$  that for the equation,

$$\partial_t f + a(v) \cdot \nabla_x f = g - F(t, x, v) \cdot \nabla_v \tilde{g}, \quad (1.2)$$

the obtained regularity is  $W^{s,p}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with  $s = \frac{\alpha}{2p'}$ . When we consider equation,

$$\partial_t f + a(v) \cdot \nabla_x f + F(t, x, v) \cdot \nabla_v f = g, \quad (1.3)$$

that is to say that  $\tilde{g} = f$ , it is classical to consider the term  $F(t, x, v) \cdot \nabla_v f$  being part of the right-hand side and to obtain the regularity  $W^{s,p}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with  $s = \frac{\alpha}{2p'}$ . But for (1.3), the derivation with respect to  $v$  is only on  $f$  through the transport equation and not on an arbitrary term  $\tilde{g}$ . That is to say, the conventional method is losing information because this term is part of characteristics and the right-hand side terms are in  $L^2$ , i.e. for  $m = 0$ , and the obtained regularity should be  $W^{s,p}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with  $s = \frac{\alpha}{p'}$ .

This is the first motivation of this paper and one of the result we get.

Few other papers deal with averaging lemma avoiding to consider the acceleration term as a source, namely [12,14]. But they are based on a transversality assumption on  $a(\cdot)$  restricting the generality to the case  $\alpha = 1$ .

Notations for (1.3) are  $f(t, x, v) \in \mathbb{R}$  with  $t \in \mathbb{R}$ ,  $x \in \mathbb{R}^N$ ,  $v \in \mathbb{R}^M$ ,  $a : \mathbb{R}^M \rightarrow \mathbb{R}^N$ ,  $F : \mathbb{R} \times \mathbb{R}^N \times \mathbb{R}^M \rightarrow \mathbb{R}^M$ , and

$$a(v) \cdot \nabla_x f = \sum_{i=1}^N a_i(v) \partial_{x_i} f, \quad F(t, x, v) \cdot \nabla_v f = \sum_{i=1}^M F_i(t, x, v) \partial_{v_i} f.$$

In this paper, we will prove the following averaging lemmas on Eq. (1.3).

**Theorem 1** ( $L^2$  result). *Let  $a \in C^{N+3}(\mathbb{R}_v^M, \mathbb{R}_x^N)$ ,  $F \in C^{N+3}(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M, \mathbb{R}^M)$ ,  $f, g \in L^2(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M)$ , satisfying (1.3). Let  $A > 0$  and  $\psi \in C_c^{N+2}(\mathbb{R}_v^M)$  be such that the support of  $\psi$  is included in  $[-A, A]^M$ . We assume that there exists  $0 < \alpha \leq 1$  and  $C > 0$  such that for any  $(u, \sigma) \in S^N$  and  $\varepsilon > 0$ ,*

$$\operatorname{meas}(\{v \in [-A, A]^M; u - \varepsilon < a(v) \cdot \sigma < u + \varepsilon\}) \leq C\varepsilon^\alpha. \quad (1.4)$$

*Then the averaging,*

$$\rho_\psi(t, x) = \int_{\mathbb{R}^M} f(t, x, v) \psi(v) dv,$$

*is in  $H_{loc}^{\alpha/2}(\mathbb{R}_t \times \mathbb{R}_x^N)$ .*

**Remark 1.1.** We notice that we obtain  $\alpha/2$  instead of the well known  $\alpha/4$  when the acceleration term  $F \cdot \nabla_x f$  is considered as a right-hand side with no particular relation to  $f$ .

**Remark 1.2.** For Vlasov equation, the classical application of averaging lemma is the DiPerna, Lions, Meyer Theorem which gives the compactness for  $\rho_\psi$  with an operator of the kind (1.3) applying the result with  $g_1 = -F \cdot f$  when  $F \in L_{loc}^\infty$ . More precisely, if  $f^n, g_0^n$  and  $g_1^n = -F_n \cdot f^n$  are solutions of (1.1) with some bounds in  $L^p$ , then  $\rho_\psi^n$  is bounded in  $W^{s,p}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with  $s = \frac{\alpha}{2p}$ , and thus is compact in  $W^{s',p}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with  $s' < s$ . For  $p = 2$ , it is compact in  $H^{s'}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with  $s' < \frac{\alpha}{4}$ . By this way, paper [8] proves the existence of weak solutions for Vlasov–Maxwell. With Theorem 1, the obtained compactness is in  $H_{loc}^{s'}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with  $s' < \frac{\alpha}{2}$ .

When the force is constant, we obtain a global regularity result with a less smooth test function.

**Theorem 2** ( *$L^2$  result with  $F$  constant*). Let  $a \in C^\gamma(\mathbb{R}_v^M, \mathbb{R}_x^N)$ ,  $F(t, x, v) = F \in \mathbb{R}^M$ ,  $F \neq 0$ ,  $f, g \in L^2(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M)$  satisfying (1.3) where we assume that function  $a(\cdot)$  satisfies the following condition with  $\gamma$ , which is a positive integer, such that  $\forall (v, \sigma) \in \mathbb{R}^M \times S^N$ ,  $\sigma = (\sigma_0, \sigma_1, \dots, \sigma_N)$ ,  $\tilde{\sigma} = (\sigma_1, \dots, \sigma_N)$ ,

$$\left| \sigma_0 + a(v) \cdot \tilde{\sigma} \right| + \sum_{k=1}^{\gamma-1} \left| (F \cdot \nabla_v)^k a(v) \cdot \tilde{\sigma} \right| > 0. \quad (\gamma ND) \quad (1.5)$$

Let  $\psi \in C_c^1(\mathbb{R}_v^M)$ , then the averaging,

$$\rho_\psi(t, x) = \int_{\mathbb{R}^M} f(t, x, v) \psi(v) dv,$$

is in  $H^{1/\gamma}(\mathbb{R}_t \times \mathbb{R}_x^N)$ .

**Remark 1.3.** The proof of Theorem 2 is not valid when  $F = 0$ . So this theorem does not give an averaging lemma for the kinetic equation  $\partial_t f + a(v) \cdot \nabla_x f = g$ .

**Remark 1.4.** The case of a nonzero constant force field is not without interest, as it appears for instance when considering gravity effects in the kinetic theory of neutral gases.

**Remark 1.5** ( $M = 1$ , one-dimensional velocity).

1. The Sobolev estimate for  $\rho_\psi$  comes from optimal bounds in stationary phase lemma. Then, with only  $f, g \in L^2$  and  $M = 1$ , we expect Theorem 2 to give the best Sobolev's exponent.
2. Since  $\gamma \geq N + 1$  (see Proposition 6 for this inequality), with only  $f, g \in L^2$ , we expect  $\rho_\psi$  to belong at most to  $H^{1/(N+1)}(\mathbb{R}_x^{N+1})$  when  $M = 1$ .
3. With scalar velocity, the condition  $(\gamma ND)$  is similar to a non-degeneracy condition given in [13] about averaging for operators with real principal symbols. More precisely it is the condition (5) of Theorem 4 with  $t = v$  and  $\xi_0 = F$  in [13]. But our result yields a better smoothing effect, the gain of regularity for the average is  $1/\gamma$  instead of  $1/(2(\gamma - 1))$  in [13].

Next theorem is a comparison between the two previous results. It shows that Theorem 1 does not give the best Sobolev exponent when  $M = 1$  and that Theorem 2 is not optimal for  $M > 1$ .

**Theorem 3.** For  $N \geq 2$  and  $M = 1$ , Theorem 2 gives a stronger smoothing effect than Theorem 1 for the best  $\gamma = \gamma_{opt}$  compared with the best  $\alpha = \alpha_{opt}$  since

$$\frac{1}{\gamma_{opt}} = \frac{1}{N+1} > \frac{\alpha_{opt}}{2} = \frac{1}{2N}.$$

Conversely, for  $N = M$ , Theorem 1 can give one half derivative with the best  $\alpha = 1$ .

**Remark 1.6.** For scalar velocity ( $v \in \mathbb{R}$ ,  $M = 1$ ), we characterize in Theorem 3 the best parameter  $\alpha$  for the classical non-degeneracy condition, namely condition (1.4). This characterization is mentioned in few works, see [25,17], but the proof of optimality is a new result. This kind of characterization also gives new results for scalar conservation laws, see [19].

Finally, we find out two results in  $L^p$  framework.

**Theorem 4** (First  $L^p$  result with  $F$  constant). Let  $a \in C^{N+3}(\mathbb{R}_v^M, \mathbb{R}_x^N)$ ,  $F(t, x, v) = F \in \mathbb{R}_v^M$ ,  $f, g \in L^p(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M)$ , satisfying (1.3). Let  $A > 0$  and  $\psi \in C_c^{N+2}(\mathbb{R}_v^M)$  be such that the support of  $\psi$  is included in  $[-A, A]^M$ . We assume that there exists  $0 < \alpha \leq 1$  and  $C > 0$  such that for any  $(u, \sigma) \in S^N$  and  $\varepsilon > 0$ ,

$$\text{meas}(\{v \in [-A, A]^M; u - \varepsilon < a(v) \cdot \sigma < u + \varepsilon\}) \leq C\varepsilon^\alpha. \quad (1.6)$$

Then the averaging,

$$\rho_\psi(t, x) = \int_{\mathbb{R}^M} f(t, x, v) \psi(v) dv,$$

is in  $W_{loc}^{s,p}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with  $s = \frac{\alpha}{p'}$ .

**Theorem 5** (Second  $L^p$  result with  $F$  constant). Let  $a \in C^\gamma(\mathbb{R}_v^M, \mathbb{R}_x^N)$ ,  $F(t, x, v) = F \in \mathbb{R}_v^M$ ,  $F \neq 0$ ,  $f, g \in L^p(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M)$  with  $1 < p \leq 2$ , satisfying (1.3), where we assume that  $a(\cdot)$  satisfies the following condition with  $\gamma$ , which is a positive integer, such that

$$\forall (v, \sigma) \in \mathbb{R}^M \times S^N, \quad \sigma = (\sigma_0, \sigma_1, \dots, \sigma_N), \quad \tilde{\sigma} = (\sigma_1, \dots, \sigma_N),$$

$$|\sigma_0 + a(v) \cdot \tilde{\sigma}| + \sum_{k=1}^{\gamma-1} |(F \cdot \nabla_v)^k a(v) \cdot \tilde{\sigma}| > 0. \quad (\gamma ND)$$

Let  $\psi \in C_c^1(\mathbb{R}_v^M)$ , then the averaging,

$$\rho_\psi(t, x) = \int_{\mathbb{R}} f(t, x, v) \psi(v) dv,$$

is in  $W^{s,p}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with  $s = \frac{2}{\gamma p'}$ .

**Remark 1.7.** These results are presented with time dependence because it is more useful in applications.

In the proof of next sections, we take the following notations. We set  $X = (t, x)$  and  $b(v) = (1, a(v))$ . Then (1.3) can be rewritten as follows:

$$b(v) \cdot \nabla_X f + F(X, v) \cdot \nabla_v f = g, \quad (1.7)$$

where  $X \in \mathbb{R}^{N+1}$ ,  $v \in \mathbb{R}^M$ .

Here is how the paper is structured.

In Section 2, we prove Theorem 1 for a smooth force field. In Section 3, we prove Theorem 2 for a constant and nonzero force field. In Section 4, we compare both results (Theorem 3) and finally in Section 5, we prove the extension to  $L^p$  spaces for constant force (Theorems 4 and 5).

## 2. First theorem in the $L^2$ framework

We first recall the following classical averaging lemma (see [15,5]).

**Proposition 1** (Golse, Lions, Perthame, Sentis). Let  $a \in L_{loc}^\infty(\mathbb{R}^M, \mathbb{R}^N)$ ,  $f, g \in L^2(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M)$ , such that

$$\partial_t f + a(v) \cdot \nabla_x f = g. \quad (2.1)$$

Let  $\psi \in L^\infty(\mathbb{R}_v^M)$ , with compact support in some  $[-A, A]^M$ , such that there exists  $0 < \alpha \leq 1$  and  $C > 0$  such that

$$\text{meas}(\{v \in [-A, A]^M; u - \varepsilon < a(v) \cdot \sigma < u + \varepsilon\}) \leq C\varepsilon^\alpha, \quad (2.2)$$

for any  $(u, \sigma) \in S^N$  and  $\varepsilon > 0$ . Then the averaging,

$$\rho_\psi(t, x) = \int_{\mathbb{R}^M} f(t, x, v) \psi(v) dv,$$

is in  $H^{\alpha/2}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with the estimate:

$$\|\rho_\psi\|_{H^{\alpha/2}} \leq \tilde{C}(N) (\|\psi\|_{L^2} + \sqrt{K} \|\psi\|_{L^\infty}) (\|f\|_{L^2} + \|g\|_{L^2}).$$

We use this averaging lemma to prove an other result, which deals with test function depending on  $(t, x, v)$ .

**Proposition 2** (Averaging lemma with test function in  $(X, v)$ ). Let  $a \in L_{loc}^\infty(\mathbb{R}_v^M, \mathbb{R}_x^N)$ ,  $f, g \in L^2(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M)$ , such that

$$\partial_t f + a(v) \cdot \nabla_x f = g. \quad (2.3)$$

Let  $\psi \in L_c^\infty(\mathbb{R}_v^M, W^{N+2, \infty}(\mathbb{R}_{tx}^{N+1}))$  with compact support with respect to  $v$  in some  $[-A, A]^M$ . We assume that there exists  $0 < \alpha \leq 1$  and  $C > 0$  such that

$$\text{meas}(\{v \in [-A, A]^M; u - \varepsilon < a(v) \cdot \sigma < u + \varepsilon\}) \leq C\varepsilon^\alpha, \quad (2.4)$$

for any  $(u, \sigma) \in S^N$  and  $\varepsilon > 0$ .

Then, for any compact  $K$ , there exists a constant  $C(N, K)$  such that the averaging,

$$\rho_\psi(t, x) = \int_{\mathbb{R}} f(t, x, v) \psi(t, x, v) dv,$$

is in  $H_{loc}^{\alpha/2}(\mathbb{R}_t \times \mathbb{R}_x^N)$  with the bound

$$\|\rho_\psi\|_{H_K^{\alpha/2}} \leq C(N, K) (\|f\|_{L^2} + \|g\|_{L^2}) \|\psi\|_{(L^2 \cap L^\infty)_v(W_{tx}^{N+2, \infty})}.$$

**Proof.** We fix a compact  $K$  on  $X$ . We take  $\tilde{K} = [-S, S]^{N+1}$  such that  $K \subset \tilde{K}$  and  $\chi$  a  $C^\infty$  function such that  $\chi = 1$  on  $K$  and 0 outside  $\tilde{K}$ . Finally, we set  $\tilde{\psi} = \psi \chi$ .

Since  $\tilde{\psi}$  has a compact support with respect to  $X$ , we can extend it by periodicity in these variables. Then the Fourier expansion with respect to  $X$  gives:

$$\tilde{\psi}(X, v) = \sum_{\beta \in \mathbb{Z}^{N+1}} c_\beta(v) e^{iS\beta \cdot X}.$$

We write this formula through,

$$\tilde{\psi}(X, v) = \sum_{\beta \in \mathbb{Z}^{N+1}} ((1 + |\beta|^r) c_\beta(v)) \cdot \frac{e^{iS\beta \cdot X}}{1 + |\beta|^r},$$

with  $r = N/2 + 1$ . We set:

$$\phi_\beta(X) = \frac{e^{iS\beta \cdot X}}{1 + |\beta|^r}, \quad \text{and} \quad \psi_\beta(v) = (1 + |\beta|^r) c_\beta(v).$$

We use the decreasing of Fourier coefficients for  $W^{N+2, \infty}(\mathbb{R}_x^{N+1})$  function, that is to say,

$$|c_\beta(v)| \leq \frac{C_1}{(S|\beta|)^{N+2}} \|\tilde{\psi}(\cdot, v)\|_{W_X^{N+2, \infty}}.$$

Thus we have:

$$\begin{aligned}
\int_{\mathbb{R}^M} \sum_{\beta \in (\mathbb{Z}^{N+1})^*} |\psi_\beta(v)|^2 dv &\leq \int_{\mathbb{R}^M} \sum_{\beta \in (\mathbb{Z}^{N+1})^*} (1 + |\beta|^r)^2 |c_\beta(v)|^2 dv \\
&\leq \frac{C_2}{S^{2N+4}} \int_{\mathbb{R}^M} \sum_{\beta \in (\mathbb{Z}^{N+1})^*} \frac{4|\beta|^{2r}}{|\beta|^{2N+4}} \|\tilde{\psi}(\cdot, v)\|_{W_X^{N+2, \infty}}^2 dv \\
&\leq \frac{4C_2}{S^{2N+4}} \sum_{\beta \in (\mathbb{Z}^{N+1})^*} \frac{1}{|\beta|^{N+2}} \|\psi\|_{L_v^2(W_X^{N+2, \infty})}^2 < +\infty.
\end{aligned} \tag{2.5}$$

On  $K$ , we notice that

$$\begin{aligned}
\rho_\psi(X) &= \int_{\mathbb{R}} f(X, v) \psi(X, v) dv \quad \chi(X) = \int_{\mathbb{R}} f(X, v) \tilde{\psi}(X, v) dv \\
&= \int_{\mathbb{R}^M} f(X, v) \sum_{\beta \in \mathbb{Z}^{N+1}} \phi_\beta(X) \psi_\beta(v) dv.
\end{aligned}$$

To apply Fubini's Theorem, we need that, for *a.e.*  $X$ ,

$$\int_{\mathbb{R}^M} \sum_{\beta \in (\mathbb{Z}^{N+1})^*} |f(X, v) \phi_\beta(X) \psi_\beta(v)| dv < +\infty.$$

It comes from,

$$\begin{aligned}
&\int_{\mathbb{R}^M} \sum_{\beta \in (\mathbb{Z}^{N+1})^*} |f(X, v) \phi_\beta(X) \psi_\beta(v)| dv \\
&\leq \int_{\mathbb{R}^M} |f(X, v)| \sum_{\beta \in (\mathbb{Z}^{N+1})^*} |\phi_\beta(X) \psi_\beta(v)| dv \\
&\leq \sqrt{\int_{\mathbb{R}^M} |f(X, v)|^2 dv} \sqrt{\int_{\mathbb{R}^M} \left( \sum_{\beta \in (\mathbb{Z}^{N+1})^*} |\phi_\beta(X) \psi_\beta(v)| \right)^2 dv} \\
&\leq \|f(X, \cdot)\|_{L_v^2} \sqrt{\sum_{\beta \in (\mathbb{Z}^{N+1})^*} |\phi_\beta(X)|^2 \int_{\mathbb{R}^M} \sum_{\beta \in (\mathbb{Z}^{N+1})^*} |\psi_\beta(v)|^2 dv} \\
&\leq \|f(X, \cdot)\|_{L_v^2} \sqrt{\sum_{\beta \in (\mathbb{Z}^{N+1})^*} \frac{1}{(1 + |\beta|^r)^2} \int_{\mathbb{R}^M} \sum_{\beta \in (\mathbb{Z}^{N+1})^*} |\psi_\beta(v)|^2 dv} \\
&< +\infty,
\end{aligned}$$

since  $2r > N + 1$  and from (2.5). Thus we can write, on  $K$ ,

$$\begin{aligned}
\rho_\psi(X) &= \sum_{\beta \in \mathbb{Z}^{N+1}} \phi_\beta(X) \rho_{\psi_\beta}(X), \\
\rho_{\psi_\beta}(X) &= \int_{\mathbb{R}} f(X, v) \psi_\beta(v) dv.
\end{aligned}$$

The classical averaging lemma (Proposition 1) gives that

$$\|\rho_{\psi_\beta}\|_{H_K^{\alpha/2}} \leq \tilde{C}(N) (\|\psi_\beta\|_{L^2} + \sqrt{C} \|\psi_\beta\|_{L^\infty}) (\|f\|_{L^2} + \|g\|_{L^2}).$$

We now use the following property: For  $u_1 \in C^s(\Omega)$ ,  $u_2 \in H^s(\Omega)$ , with  $s \in ]0, 1[$ , with  $\Omega$  a bounded open set of  $\mathbb{R}^{N+1}$ , we have  $u_1 u_2 \in H^s(\Omega)$  with

$$\|u_1 u_2\|_{H^s} \leq C_3 \|u_1\|_{C^s} \|u_2\|_{H^s}.$$

This result gives, for  $s = \alpha/2$ ,

$$\begin{aligned} \|\rho_\psi\|_{H_K^{\alpha/2}} &\leq C_3 \sum_{\beta \in \mathbb{Z}^{N+1}} \|\phi_\beta\|_{C_K^{\alpha/2}} \|\rho_\psi\|_{H_K^{\alpha/2}} \\ &\leq C_4 \sum_{\beta \in (\mathbb{Z}^{N+1})^*} \|\phi_\beta\|_{C_K^{\alpha/2}} (\|\psi_\beta\|_{L^2} + \|\psi_\beta\|_{L^\infty}) (\|f\|_{L^2} + \|g\|_{L^2}) + C_3 \|\rho_\psi\|_{H_K^{\alpha/2}} \\ &\leq C_5 \left( \sum_{\beta \in (\mathbb{Z}^{N+1})^*} \frac{1}{|\beta|^{r-\alpha/2}} \frac{\|\psi\|_{(L^2 \cap L^\infty)_v(W_X^{N+2, \infty})}}{|\beta|^{N+2-r}} (\|f\|_{L^2} + \|g\|_{L^2}) + \|\tilde{\psi}\|_{C_K^1} \right) \\ &\leq C_5 \left( \sum_{\beta \in (\mathbb{Z}^{N+1})^*} \frac{1}{|\beta|^{N+2-\alpha/2}} (\|f\|_{L^2} + \|g\|_{L^2}) \|\psi\|_{(L^2 \cap L^\infty)_v(W_X^{N+2, \infty})} + \|\psi\|_{C_c^{N+2}} \right). \end{aligned}$$

Since  $N + 2 - \alpha/2 > N + 1$ , the proof is completed.  $\square$

With this proposition now stated, we can go into the proof of our first theorem.

**Proof of Theorem 1.** Let  $K$  be a compact in  $\mathbb{R}_X^{N+1}$ . We set  $\mathcal{K} = K \times [-A, A]^M$ . We perform locally a change in variables in order to rewrite Eq. (1.7) without the term  $\nabla_v f$  and to apply previous result. For any  $(X_0, v_0) \in \mathcal{K}$ , using the characteristics since  $b(v) = (1, a(v)) \neq 0$ , there exists  $\mathcal{B}_0 \subset \mathcal{K}$  a neighborhood of  $(X_0, v_0)$  and a  $C^{N+3}$  diffeomorphism:

$$\begin{aligned} \Phi_0 : \mathcal{B}_0 &\rightarrow \mathcal{B}^0, \\ (X, w) &\mapsto \Phi_0(X, w) = (X, V_0(X, w)), \end{aligned}$$

such that on  $\mathcal{B}_0$  we have

$$b(V_0(X, w)) \cdot \nabla_X V_0(X, w) = F(X, V_0(X, w)). \quad (2.6)$$

Let us explain more precisely how to define the diffeomorphism  $\Phi_0$  from Eq. (2.6). Since  $b(v) = (1, a(v))$ ,  $X = (t, x)$  and  $X_0 = (t_0, x_0)$ , Eq. (2.6) can be reformulated as a nonlinear hyperbolic system (where  $w$  is a parameter):

$$\partial_t V_0(t, x; w) + a(V_0(t, x; w)) \cdot \nabla_x V_0(t, x; w) = F(t, x, V_0(t, x; w)), \quad (2.7)$$

completed by the initial data,

$$V_0(t_0, x; w) = w. \quad (2.8)$$

By the classical method of characteristics, for each  $w$ , there exists a neighborhood of  $(t_0, x_0)$  where  $V_0$  is well defined and smooth. The characteristics are smooth with respect to the parameter  $w$ , thus  $V_0(t, x, w)$  is well defined on a neighborhood of  $(t_0, x_0; v_0)$ . Notice that  $\partial_w V_0(t_0, x; w) = id_{\mathbb{R}^M}$ , with  $id_{\mathbb{R}^M}$  the identity operator on  $\mathbb{R}_v^M$ , and  $\det(D\Phi_0) = \det(\partial_w V_0)$ , so reducing if necessary the previous neighborhood,  $\Phi_0$  is a diffeomorphism on  $\mathcal{B}_0$ .

Denoting by  $\tilde{f}_0(X, w) = f(X, V_0(X, w))$ ,  $\tilde{g}_0(X, w) = g(X, V_0(X, w))$ ,  $\tilde{b}_0(w) = b(V_0(X, w))$ , Eq. (1.7) rewrites,

$$\tilde{b}_0(w) \cdot \nabla_X \tilde{f}_0 = \tilde{g}_0. \quad (2.9)$$

Now, there exists a finite number of  $\mathcal{B}^l$  to recover this compact, i.e. there exists  $\{(X_l, v_l)\}_{l=1, \dots, L}$ , with the associated diffeomorphism  $\Phi_l : \mathcal{B}_l \rightarrow \mathcal{B}^l$ ,  $\Phi_l(X, w) = (X, V_l(X, w))$ , such that  $\mathcal{K} \subset \bigcup_{l=1, \dots, L} \mathcal{B}^l$ . For this recovering, we use a partition of unity, we have:

$$f(X, v) = f(X, v) \mathbb{1}_{\mathcal{K}}(X, v) = \sum_{l=1}^L f(X, v) \chi_l(X, v),$$

where function  $\chi_l$  are  $C^\infty$  and have a compact support in  $\mathcal{B}^l$ .



Denoting again by  $\tilde{f}_l(X, w) = f(X, V_l(X, w))$ ,  $\tilde{g}_l(X, w) = g(X, V_l(X, w))$ ,  $\tilde{b}_l(w) = b(V_l(X, w))$  on  $\mathcal{B}_l$ , and

$$\begin{aligned} I^l[X] &= \{v \in \mathbb{R}^M \text{ such that } (X, v) \in \mathcal{B}^l\}, \\ I_l[X] &= \{w \in \mathbb{R}^M \text{ such that } (X, w) \in \mathcal{B}_l\}, \end{aligned}$$

we have the following decomposition. It is,

$$\begin{aligned} \rho_\psi(X) &= \sum_{l=1}^L \int_{\mathbb{R}^M} f_l(X, v) \chi_l(X, v) \psi(v) dv \\ &= \sum_{l=1}^L \int_{I^l[X]} f(X, v) \chi_l(X, v) \psi(v) dv \\ &= \sum_{l=1}^L \int_{I_l[X]} \tilde{f}(X, w) \chi_l(X, V_l(X, w)) \psi(V_l(X, w)) J_l(X, w) dw, \end{aligned}$$

where we can perform the variable change  $v \mapsto w = V(X, v)$  on every neighborhood  $\mathcal{B}^l$  corresponding to  $l$  and denoting by  $J_l(X, w)$  the associated Jacobian, i.e.  $J_l = |\det D\Phi_l| = |\det \partial_w V_l|$ .

We set  $\overline{\psi}_l(X, w) = \chi_l(X, V(X, w)) \psi(V(X, w)) J_l(X, w)$ . Since  $a$  and  $F$  have  $C^{N+3}$  regularity,  $J_l$  has  $C^{N+2}$  one. Furthermore  $\psi \in C_c^{N+2}$ , thus  $\overline{\psi}_l \in (L^2 \cap L^\infty)_c(\mathbb{R}_v^M, W^{N+2, \infty}(\mathbb{R}_X^{N+1}))$ . We apply previous result, namely Proposition 2, on the averaging,

$$\rho_{\overline{\psi}_l}(X) = \int_{\mathbb{R}^M} \tilde{f}(X, w) \overline{\psi}_l(X, w) dw \quad \text{which is in } H_{loc}^{\alpha/2}(\mathbb{R}_X^{N+1}).$$

Finally the inequality  $\|\rho_\psi\|_{H_K^{\alpha/2}} \leq \sum_{l=1}^L \|\rho_{\overline{\psi}_l}\|_{H_K^{\alpha/2}}$  concludes the proof.  $\square$

### 3. Case of a constant force field

When  $F$  is a nonzero constant vector, we can obtain a different result. The way to get it is quite different and we have to be restricted to the case of a constant force field. A key tool here is a generalized uniform version of the classical method of the stationary phase. We work on Eq. (1.7) with  $F$  constant,  $F \in \mathbb{R}^M$ ,  $F \neq 0$ . Let us denote a directional  $v$ -derivative along vector  $F$  by:

$$D = F \cdot \nabla_v. \quad (3.1)$$

The smoothing effect depends on  $(\gamma ND)$  assumption of Theorem 2. Indeed, it is exactly the following non-degeneracy condition about  $D$ -derivatives of  $b(\cdot)$ :

$$\forall (v, \sigma) \in \mathbb{R}^M \times S^N, \quad \sum_{k=0}^{\gamma-1} |D^k b(v) \cdot \sigma| > 0. \quad (\gamma ND)$$

Before proving Theorem 2 we give some useful results about oscillatory integrals following Stein's book [26].

**Proposition 3.** (See [26].) Suppose  $\phi \in C^{k+1}(\mathbb{R}, \mathbb{R})$  so that, for some  $k \geq 1$ ,

$$\frac{d^k \phi}{dv^k}(v) \geq 1, \quad \forall v \in ]\alpha, \beta[. \quad (3.2)$$

Then

$$\left| \int_{\alpha}^{\beta} e^{i\lambda\phi(v)} dv \right| \leq c_k \cdot \frac{1}{|\lambda|^{1/k}}$$

holds when,

1.  $k \geq 2$ , or
2.  $k = 1$  and  $\phi'$  is monotonous.

Furthermore, the bound  $c_k$  is independent of  $\lambda$  and  $\phi$ .

This proposition can be found in [26, p. 332]. Elias M. Stein obtains  $c_k \leq 5 \cdot 2^{k-1} - 2$  in his proof. Notice that  $c_k$  is independent of the length of the interval  $]\alpha, \beta[$ . For  $|\lambda| < 1$ , the bound for the oscillatory integral blows up. Indeed, for  $k = 1$ , we can relax the monotonous assumption on  $\phi$  by the following bounds:

$$|\phi'(v)| \geq \delta > 0, \quad \forall v \in ]\alpha, \beta[, \quad \tilde{c}_1 = 2 + \delta^{-1} \int_{\alpha}^{\beta} |\phi''(v)| dv.$$

Indeed, integrating by parts and using the inequality:

$$\min(a, \beta b) \leq \min(1, \beta) \max(a, b) \quad \text{for all non-negative } a, b, \beta,$$

we get,

$$\left| \int_{\alpha}^{\beta} e^{i\lambda\phi(v)} dv \right| \leq \max(|\beta - \alpha|, \tilde{c}_1) \cdot \max\left(1, \frac{1}{\delta}\right) \cdot \min\left(1, \frac{1}{|\lambda|}\right).$$

Furthermore, the bound given in Proposition 3 blows up for small  $\lambda$ , so we replace it by the length of the interval and get the following corollary.

**Corollary 1.** Let  $\delta > 0$ . Suppose  $\phi \in C^{k+1}(\mathbb{R}, \mathbb{R})$  so that, for some  $k \geq 1$ ,

$$\left| \frac{d^k \phi}{dv^k}(v) \right| \geq \delta, \quad \forall v \in ]\alpha, \beta[. \quad (3.3)$$

Then  $|\int_{\alpha}^{\beta} e^{i\lambda\phi(v)} dv| \leq \max(|\beta - \alpha|, \tilde{c}_k) \cdot \max(1, \frac{1}{\delta^{1/k}}) \min(1, \frac{1}{|\lambda|^{1/k}})$ , where  $\tilde{c}_k$  is independent of  $\lambda$ ,  $\phi$  and  $]\alpha, \beta[$  for  $k \geq 2$  and  $\tilde{c}_1 = 2 + \delta^{-1} \int_{\alpha}^{\beta} |\phi''(v)| dv$ .

Notice that, for  $k \geq 2$ ,  $\tilde{c}_k = c_k$  is given in Proposition 3.

Following Stein's book (Corollary, p. 334), we obtain the following proposition.

**Proposition 4.** (See [26].) Let  $\psi \in W^{1,1}(\alpha, \beta)$ ,  $\phi \in C^{k+1}(\mathbb{R}, \mathbb{R})$  such that, for some  $\delta > 0$  and  $k \geq 1$ ,

$$\left| \frac{d^k \phi}{dv^k}(v) \right| \geq \delta, \quad \forall v \in ]\alpha, \beta[.$$

Then

$$\left| \int_{\alpha}^{\beta} \psi(v) e^{i\lambda\phi(v)} dv \right| \leq \frac{\max(|\beta - \alpha|, \tilde{c}_k)}{\min(1, \delta^{1/k}) \max(1, |\lambda|^{1/k})} (\|\psi\|_{L^{\infty}(\alpha, \beta)} + \|\psi'\|_{L^1(\alpha, \beta)}),$$

where  $\tilde{c}_k$  is independent of  $\lambda$ ,  $\phi$ ,  $\psi$  and  $]\alpha, \beta[$  for  $k \geq 2$ , and  $\tilde{c}_1 = 2 + \delta^{-1} \int_{\alpha}^{\beta} |\phi''(v)| dv$ .

**Proof.** This is classically proved in writing the integral  $\int_{\alpha}^{\beta} \psi(v) e^{i\lambda\phi(v)} dv$  as  $\int_{\alpha}^{\beta} \psi(v) I'(v) dv$ , with  $I(v) = \int_{\alpha}^v e^{i\lambda\phi(u)} du$ , integrating by parts and using the uniform estimate for  $|I(v)|$  from previous corollary.  $\square$

Now we generalize Proposition 4 in the case with parameters and a  $(\gamma ND)$  like assumption.

**Proposition 5.** Suppose  $P$  is a compact set of parameter  $p$ ,  $A > 0$ ,  $\psi(u; p)$  belongs to  $L_p^\infty(P, W_u^{1,1}(\mathbb{J}-A, A])$  and  $\phi(u; p) \in C^{\gamma+1}(\mathbb{R}_u \times P_p, \mathbb{R})$ , such that for all  $(u, p)$  in  $K = [-A, A] \times P$ ,

$$\sum_{k=1}^{\gamma} \left| \frac{\partial^k \phi}{\partial u^k} \right| (u; p) > 0. \quad (3.4)$$

Then, for any  $]\alpha, \beta[ \subset ]-A, A[$ ,

$$\left| \int_{\alpha}^{\beta} \psi(u; p) e^{i\lambda\phi(u; p)} du \right| \leq d_{\gamma} \cdot \min\left(1, \frac{1}{|\lambda|^{1/\gamma}}\right) \cdot \left( \|\psi\|_{L^\infty(K)} + \left\| \frac{\partial\psi}{\partial u} \right\|_{L^\infty(P, L^1(\mathbb{J}-A, A])} \right),$$

where constant  $d_{\gamma}$  is independent of  $\lambda$  and only depends on  $A$ ,  $\sup_K \left| \frac{\partial^2 \phi}{\partial u^2} \right|$ ,  $\inf_K \frac{1}{\gamma} \sum_{k=1}^{\gamma} \left| \frac{\partial^k \phi}{\partial u^k} \right|$ .

**Proof.** Since  $K$  is a compact set, we can choose  $0 < \delta \leq 1$  such that everywhere on  $K$ :

$$0 < \delta < \frac{1}{\gamma} \sum_{k=1}^{\gamma} \left| \frac{\partial^k \phi}{\partial u^k} \right| (u; p).$$

Let us define the open set  $Z_k = \{(u; p), |\partial_u^k \phi(u; p)| > \delta\}$ , for  $k = 1, \dots, \gamma$ . Necessarily  $K \subset \bigcup_{k=1}^{\gamma} Z_k$ , and then there exists a partition of unity such that  $\sum_{k=1}^{\gamma} \rho_k \equiv 1$  on  $K$  and such that the support of  $\rho_k$  is included in  $Z_k$ . Let us define  $\psi_k = \rho_k \psi$  and  $I = I_1 + \dots + I_{\gamma}$  where  $I_k(p) = \int_a^b \psi_k(u; p) e^{i\lambda\phi(u; p)} du$ . We apply Proposition 4 on each  $I_k$  where the exponent “ $\cdot$ ” denotes  $\partial_u$ :

$$|I_k| \leq \frac{\max(2A, \tilde{c}_k)}{\delta^{1/k} \max(1, |\lambda|^{1/k})} \sup_P (\|\psi_k(\cdot, p)\|_{L^\infty(\mathbb{J}-A, A]) + \|\psi_k'(\cdot, p)\|_{L^1(\mathbb{J}-A, A]}).$$

Since for any fixed  $p$  and  $J = ]-A, A[$ , we have:

$$(\|\psi_k(\cdot, p)\|_{L^\infty(J)} + \|\psi_k'(\cdot, p)\|_{L^1(J)}) \leq (\|\rho_k\|_{L^\infty(J)} + \|\rho_k'\|_{L^1(J)}) (\|\psi(\cdot, p)\|_{L^\infty(J)} + \|\psi'(\cdot, p)\|_{L^1(J)}),$$

it is enough to take,

$$d_{\gamma} = \sum_k \frac{\max(2A, \tilde{c}_k)}{\delta^{1/k}} (\|\rho_k\|_{L^\infty(K)} + \|\partial_u \rho_k\|_{L^\infty(P, L^1(J))}),$$

to conclude the proof.  $\square$

We are now able to prove the second theorem.

**Proof of Theorem 2.** The proof is splitted in three steps. First, we choose a suitable variable associated to  $D$ . Secondly, we use Fourier transform with respect to  $X$  and solve a linear ordinary differential equation with respect to  $v_1$ . Third, we obtain Sobolev estimates for  $\rho_{\psi}$  with Proposition 5.

**Step 1. Change of coordinates.** With a suitable choice of orthonormal coordinates, we assume, without loss of generality that

$$D = F \cdot \nabla_v = |F| \frac{\partial}{\partial v_1},$$

where  $|F|$  is the Euclidean norm of vector  $F$  and  $v = (v_1, v_2, \dots, v_M) \equiv (v_1; w)$ . Notice that the Jacobian for an orthonormal change of variables is one, thus the estimates on  $\rho_{\psi}$  are invariant through such choice for  $v_1$ . With such notations, Eq. (1.7) becomes:

$$b(v) \cdot \nabla_X f + |F| \frac{\partial f}{\partial v_1} = g. \quad (3.5)$$

**Step 2.** *Linear o.d.e.* Denoting by  $\mathcal{F}(f)$  the Fourier transform of  $f$  with respect to  $X$ , and by  $Y$  the dual variable of  $X$ , Eq. (3.5) becomes:

$$|F| \frac{\partial}{\partial v_1} \mathcal{F}(f) + i(b(v) \cdot Y) \mathcal{F}(f) = \mathcal{F}(g). \quad (3.6)$$

For almost all fixed  $Y$ , we solve an ordinary differential equation with respect to  $v_1$ . For this purpose, we chose the initial  $v_1$ , namely  $v_1^0 \in ]0, 1[$ , such that

$$\int_{\mathbb{R}_Y^{N+1} \times \mathbb{R}_w^{M-1}} |\mathcal{F}(f)|^2(Y; v_1^0; w) dY dw \leq \int_{\mathbb{R}_{v_1}} \int_{\mathbb{R}_Y^{N+1} \times \mathbb{R}_w^{M-1}} |\mathcal{F}(f)|^2(Y; v_1; w) dY dw dv_1. \quad (3.7)$$

Existence of such  $v_1^0$  is a consequence of Fubini's Theorem.

Indeed, let  $h(v_1) = \int_{\mathbb{R}_Y^{N+1}} \int_{\mathbb{R}_w^{M-1}} |\mathcal{F}(f)|^2(Y; v_1; w) dY dw \geq 0$ .

Function  $h$  is defined almost everywhere, belongs to  $L^1(\mathbb{R}_{v_1})$  and satisfies  $\|h\|_{L^1(\mathbb{R}_{v_1})} = \|f\|_{L_{X,v}^2}^2$ . Since  $h$  function cannot be everywhere greater than its mean value on  $]0, 1[$ , there exists  $v_1^0 \in ]0, 1[$  such that  $h(v_1^0) \leq \int_0^1 h(v_1) dv_1$ , which confirms (3.7).

We finally write an explicit formula for  $\mathcal{F}(f)$  with  $B(v)$  being a primitive with respect to  $v_1$  of  $-b/|F|$ :

$$B(v) = B(v_1; w) = - \int_{v_1^0}^{v_1} \frac{b(u; w)}{|F|} du,$$

$$\mathcal{F}(f)(Y, v_1; w) = \mathcal{F}(f)(Y, v_1^0; w) e^{iB(v) \cdot Y} + \frac{1}{|F|} \int_{v_1^0}^{v_1} \mathcal{F}(g)(Y, u; w) e^{i(B(v_1; w) - B(u; w)) \cdot Y} du.$$

**Step 3.**  *$H^{1/\gamma}$  estimates with oscillatory integrals.* We decompose  $\rho_\psi(t, x) = \int_{\mathbb{R}^M} f(t, x, v) \psi(v) dv$  in two parts from the explicit expression of  $\mathcal{F}(f)$  in Step 2:  $\mathcal{F}(\rho_\psi) = \hat{\rho}_f + \hat{\rho}_g$ . The first term is

$$\hat{\rho}_f(Y) = \int_{\mathbb{R}_w^{M-1}} \mathcal{F}(f)(Y, v_1^0; w) \int_{\mathbb{R}_u} \psi(u; w) e^{iB(u; w) \cdot Y} du dw.$$

In this integral, there is an oscillatory integral which is parametrized by  $w$  and  $Y = \lambda \sigma$  with  $\lambda = |Y|$  and  $\sigma \in S^N$ ; it is

$$Osc(Y, w) = \int_{\mathbb{R}_u} \psi(u; w) e^{i\lambda B(u; w) \cdot \sigma} du. \quad (3.8)$$

To use Proposition 5, we set  $p = (\sigma, w)$  which belongs to the compact set  $P = S^N \times [-A, A]^{M-1}$  with  $A > 1 > v_1^0 > 0$  such that  $\text{supp } \psi \subset [-A, A]^M$ . Condition (3.4) of Proposition 5 for oscillatory integral (3.8) is

$$\sum_{k=1}^{\gamma} \left| \frac{\partial^k B(u; w)}{\partial u^k} \cdot \sigma \right| > 0$$

which is exactly the  $(\gamma ND)$  assumption for  $b(\cdot)$ . Thanks to the  $(\gamma ND)$  assumption and Proposition 5, there exists a constant  $L$  such that for all  $(Y, w) \in \mathbb{R}^d \times [-A, A]^{M-1}$ , and for all  $\alpha, \beta$  such that  $-A < \alpha < \beta < A$ , we have

$$\max(1, |Y|^{1/\gamma}) \left| \int_{\alpha}^{\beta} \psi(u; w) e^{i\lambda B(u; w) \cdot \sigma} du \right| \leq L. \quad (3.9)$$

Using constant  $L$  and the compact support of  $\psi$  we have:

$$\max(1, |Y|^{1/\gamma}) |\hat{\rho}_f(Y)| \leq L \int_{[-A, A]^{M-1}} |\mathcal{F}(f)(Y, v_1^0; w)| dw.$$

By Cauchy–Schwarz inequality, we get:

$$\max(1, |Y|^{2/\gamma}) |\hat{\rho}_f(Y)|^2 \leq (2A)^{M-1} L^2 \int_{[-A, A]^{M-1}} |\mathcal{F}(f)(Y, v_1^0; w)|^2 dw.$$

Finally, since  $v_1^0$  satisfies (3.7), we obtain:

$$\int_{\mathbb{R}^{N+1}} \max(1, |Y|^{2/\gamma}) |\hat{\rho}_f(Y)|^2 dY \leq (2A)^{M-1} L^2 \int_{\mathbb{R}^{N+1} \times \mathbb{R}^M} |\mathcal{F}(f)(Y, v)|^2 dv dY,$$

which gives  $\hat{\rho}_f \in H^{1/\gamma}$ .

The second term  $\hat{\rho}_g$  is bounded in the same way. More precisely, we set:

$$\hat{\rho}_g(Y) = \int_{\mathbb{R}^{M-1}} H(Y, w) dw,$$

with

$$H(Y, w) = \frac{1}{|F|} \int_{-A}^A \int_{v_1^0}^{v_1} \mathcal{F}(g)(Y, u; w) e^{i(B(v_1; w) - B(u; w)) \cdot Y} du dv_1.$$

Using Fubini's Theorem and notation,

$$\Psi(Y, u; w) = \psi(u; w) e^{iB(u; w) \cdot Y},$$

we have another expression of  $H(Y, w)$ :

$$\begin{aligned} H(Y, w) &= \frac{1}{|F|} \int_{v_1^0}^A \mathcal{F}(g)(Y, u; w) e^{-iB(u; w) \cdot Y} \left( \int_u^A \Psi(Y, v_1; w) dv_1 \right) du \\ &\quad + \frac{1}{|F|} \int_{-A}^{v_1^0} \mathcal{F}(g)(Y, u; w) e^{-iB(u; w) \cdot Y} \left( \int_{-A}^u \Psi(Y, v_1; w) dv_1 \right) du, \end{aligned}$$

where there are two oscillatory integrals  $\int_u^A \Psi(Y, v_1; w) dv_1$  and  $\int_{-A}^u \Psi(Y, v_1; w) dv_1$  which are uniformly bounded thanks to inequality (3.9). Then we have:

$$\max(1, |Y|^{1/\gamma}) |H(Y; w)| \leq \frac{L}{|F|} \int_{-A}^A |\mathcal{F}(g)(Y, u; w)| du.$$

With Cauchy–Schwarz inequality, we obtain:

$$\max(1, |Y|^{2/\gamma}) |H(Y; w)|^2 \leq \frac{2AL^2}{|F|^2} \int_{-A}^A |\mathcal{F}(g)(Y, u; w)|^2 du,$$

and finally,

$$\max(1, |Y|^{2/\gamma}) |\hat{\rho}_g(Y)|^2 \leq (2A)^M \frac{L^2}{|F|^2} \int_{\mathbb{R}^M} |\mathcal{F}(g)(Y, v)|^2 dv.$$

Then  $\rho_g \in H^{1/\gamma}$ , thus finally  $\rho_\psi$  is also in this space, which concludes the proof of the theorem.  $\square$

#### 4. About non-degeneracy conditions

Theorems 1 and 2 assume two different non-degeneracy conditions on vector field  $a(v) \in \mathbb{R}^N$ ,  $v \in \text{supp } \psi \subset \mathbb{R}^M$ . Those conditions involve two parameters, namely  $\alpha = \alpha_{a(\cdot)} \in ]0, 1]$  in (1.4) and  $\gamma = \gamma_{a(\cdot), F} \in \mathbb{N}^*$  in (1.5), directly linked to the smoothing effect for the averaging in  $H_{loc}^{\alpha/2}$  or  $H^{1/\gamma}$ . In this section, we give some optimal upper bounds for  $\alpha$  and  $1/\gamma$  to compare both results obtained by different ways. Indeed, for  $M = 1$  and  $N \geq 2$ , Theorem 2 gives a better smoothing effect than Theorem 1. Conversely, when  $N = M$ , Theorem 1 is stronger than Theorem 2. In this part, we study these various properties and in particular, we prove Theorem 3.

More precisely, let  $A$  be positive, we obtain the optimal  $\alpha$  and  $\gamma$ , namely:

$$\begin{aligned}\alpha_{opt}(N, M) &= \sup_{a(\cdot) \in C^\infty([-A, A]_v^M, \mathbb{R}^N)} \alpha, \\ \gamma_{opt}(N, M) &= \min_{a(\cdot) \in C^\infty(\mathbb{R}_v^M, \mathbb{R}^N), F \in \mathbb{R}^N \setminus \{0\}} \gamma.\end{aligned}$$

We start by obtaining the easiest estimate which is a lower bound for  $\gamma$ .

**Proposition 6.** *For all  $N, M$ , we have  $\gamma \geq \gamma_{opt}(N, M) = N + 1$ .*

**Proof.** We use notations from Section 3. Following this section, the  $(\gamma ND)$  condition can be rewritten and means that we cannot find  $\sigma \in S^N$  such that  $\sigma \perp b(v)$ ,  $\sigma \perp Db(v)$ ,  $\dots$ ,  $\sigma \perp D^{\gamma-1}b(v)$ . There are  $\gamma$  conditions to satisfy. Since  $b(v)$  belongs to  $\mathbb{R}^{N+1}$ , we necessarily have  $\gamma \geq N + 1$ . Indeed  $N + 1$  is the minimal possible value for  $\gamma$ . For instance, if  $D = \frac{\partial}{\partial v_1}$ ,  $b(v) = (1, v_1, v_1^2, \dots, v_1^N)$ , with  $v = (v_1, v_2, \dots, v_M)$ , we have  $\gamma_{opt} = N + 1$ .  $\square$

The optimal  $\alpha$  is more difficult to get and it is obtained in the following subsections, see also [19]. The evaluation of exponent  $\alpha$  also implies new asymptotic expansions involving piecewise smooth functions in [20].

##### 4.1. $M = 1$ , one-dimensional velocity

**Proposition 7.** *For  $M = 1$ , we have  $\alpha \leq \alpha_{opt}(N, 1) = \frac{1}{N}$ .*

To obtain this optimal  $\alpha$  for  $M = 1$ , we need some other notations and the following results. The proof of Proposition 7 is achieved at the end of Section 4.1.

Let  $\varphi \in C^\infty([a, b], \mathbb{R})$  and  $v \in [a, b]$ , the multiplicity of  $\varphi$  on  $v$  is defined by:

$$m_\varphi[v] = \inf\{k \in \mathbb{N}, \varphi^{(k)}(v) \neq 0\} \in \bar{\mathbb{N}} = \mathbb{N} \cup \{+\infty\}.$$

It means that if  $k = m_\varphi$  then  $\varphi^{(k)}(v) \neq 0$  and  $\varphi^{(j)}(v) = 0$  for  $j = 0, 1, \dots, k - 1$ . For instance  $m_\varphi[v] = 0$  means  $\varphi(v) \neq 0$ ;  $m_\varphi[v] = 1$  means  $\varphi(v) = 0$ ,  $\varphi'(v) \neq 0$  and  $m_\varphi[v] = +\infty$  means  $\varphi^{(j)}(v) = 0$  for all  $j \in \mathbb{N}$ .

Set the multiplicity of  $\varphi$  on  $[a, b]$  by:

$$m_\varphi = \sup_{v \in [a, b]} m_\varphi[v] \in \bar{\mathbb{N}}.$$

Notice that the case where  $\varphi$  only belongs to  $C^k$ ,  $m_\varphi$  is well defined only if  $m_\varphi[v] \leq k$  for all  $v \in [a, b]$ .

**Lemma 1.** *Let  $\varphi \in C^k([a, b], \mathbb{R})$  with  $a < b$ , and*

$$Z(\varphi, \varepsilon) = \{v \in [a, b], |\varphi(v)| \leq \varepsilon\}.$$

*If  $m_\varphi$  is well defined ( $m_\varphi \leq k$ ) then there exists  $C > 0$  such that, for all  $\varepsilon > 0$ ,*

$$\text{meas}(Z(\varphi, \varepsilon)) \leq C \varepsilon^\alpha \quad \text{with } \alpha = \frac{1}{m_\varphi}. \quad (4.1)$$

*Furthermore, if  $m_\varphi$  is positive, for all  $\beta > \alpha$ , we have  $\lim_{\varepsilon \rightarrow 0} \frac{\text{meas}(Z(\varphi, \varepsilon))}{\varepsilon^\beta} = +\infty$  (Optimality).*

**Proof.** The case  $m_\varphi = 0$  is clear enough since there is no zero in this situation. Quantity  $m_\varphi$  is positive simply means that the set  $Z(\varphi, 0)$  of roots of  $\varphi$  is not empty. Since any root of  $\varphi$  has a finite multiplicity, the compact set  $Z(\varphi, 0)$  is discrete and then finite:  $Z(\varphi, 0) = \{z_1, \dots, z_v\}$ . For each  $z_i$  and  $h > 0$ , let  $V_i(h)$  be  $]z_i - h, z_i + h[ \cap [a, b]$ . For any  $0 < h < |b - a|$ , we have:

$$h \leq \text{meas}(V_i(h)) \leq 2h.$$

For any root  $z_i$ , there exists  $h_i \in ]0, |b - a|[$ ,  $A_i > 0$  and  $\delta_i > 0$  such that

$$\delta_i |h|^{k_i} \leq |\varphi(z_i + h)| \leq A_i |h|^{k_i} \quad \text{for all } h \in V_i(h_i), \quad (4.2)$$

with  $k_i = m_\varphi[z_i]$ . This is a direct consequence of Taylor–Lagrange formula. Let  $V$  be  $\bigcup_i V_i(h_i)$  and  $\varepsilon_0 = \min(1, \min_{v \in [a, b] \setminus V} |\varphi(v)|)$ . By the continuity of  $\varphi$  on the compact set  $[a, b] \setminus V$ ,  $\varepsilon_0$  is positive. Then for all  $0 < \varepsilon < \varepsilon_0$ , we have  $Z(\varphi, \varepsilon) \subset V$ . If  $\varepsilon \geq |\varphi(z_i + h)|$  for  $|h| < h_i$ , then from (4.2), we have  $(\varepsilon/\delta_i)^{1/k_i} \geq |h|$ . This last inequality implies for  $0 < \varepsilon < \varepsilon_0 \leq 1$  that  $Z(\varphi, \varepsilon)$  is a subset of  $\bigcup_i V_i((\varepsilon/\delta_i)^{1/k_i})$  and then

$$\text{meas}(Z(\varphi, \varepsilon)) \leq 2 \sum_{i=1}^v (\varepsilon/\delta_i)^{1/k_i} \leq \left( 2 \sum_{i=1}^v \delta_i^{-1/k_i} \right) \varepsilon^{1/m_\varphi}.$$

It gives inequality (4.1). To obtain the optimality of  $\alpha$ , let  $z_j$  be a root of  $\varphi$  with maximal multiplicity i.e.  $m_\varphi[z_j] = m_\varphi = k$ . Again from (4.2),  $V_j((\varepsilon/A_j)^{1/k})$  is a subset of  $Z(\varphi, \varepsilon)$  for all  $\varepsilon \in ]0, \varepsilon_0[$ . Then we have  $(\varepsilon/A_j)^{1/k} \leq \text{meas}(Z(\varphi, \varepsilon))$ , which is enough to get the optimality of  $\alpha = 1/k$  and concludes the proof.  $\square$

An upper bound of  $\alpha_{opt}(N, 1)$  is a consequence of previous lemma.

**Lemma 2.** For all  $N$ , we have  $\alpha_{opt}(N, 1) \leq 1/N$ .

**Proof.** For any  $a(\cdot) \in C^\infty(\mathbb{R}_v, \mathbb{R}_x^N)$  and  $A > 0$ , we set:

$$\varphi(v; u, \sigma) = a(v) \cdot \sigma - u = b(v) \cdot (-u, \sigma),$$

defined for  $v \in [-A, A]$ , with  $u \in \mathbb{R}$ ,  $(-u, \sigma) \in S^N$ ,  $b(v) = (1, a(v)) \in \mathbb{R}^{N+1}$ , and  $m = \sup_{(-u, \sigma) \in S^N} m_{\varphi(\cdot; u, \sigma)}$ .

Let  $v$  be fixed, we choose  $(-u, \sigma)$  such that  $m_\varphi[v] \geq N$  in order to obtain a lower bound for  $m$ .

Since  $\text{rank}\{b(v), b'(v), \dots, b^{(N-1)}(v)\} \leq N$ , there exists  $(-u, \sigma)$  such that  $u^2 + |\sigma|^2 = 1$ , and  $(-u, \sigma) \perp \{b(v), b'(v), \dots, b^{(N-1)}(v)\}$ . Then with such  $u$  and  $\sigma$ ,  $m_{\varphi(\cdot; u, \sigma)}[v] \geq N$  which implies  $m \geq N$  and consequently, from the optimality obtained in Lemma 1, we get  $\alpha \leq \alpha_{opt}(N, 1) \leq \frac{1}{N}$ .  $\square$

When function  $v \rightarrow \varphi(v; p)$  depends on a parameter  $p$ , some results are obtained in the two following lemmas to bound quantity  $C$  of Lemma 1 independently of  $p$  parameter.

**Lemma 3.** Let  $k \geq 1$ ,  $I$  an interval of  $\mathbb{R}$ ,  $\phi \in C^k(I, \mathbb{R})$  and  $\delta > 0$ .

If  $|\phi^{(k)}(v)| \geq \delta > 0$  for all  $x \in I$  then there exists a constant  $\bar{c}_k$  independent of  $\phi, I, \delta$  such that

$$\text{meas}(Z(\phi, \varepsilon)) \leq \bar{c}_k (\varepsilon/\delta)^{1/k}, \quad \text{where } Z(\phi, \varepsilon) = \{v \in I, |\phi(v)| \leq \varepsilon\}.$$

**Proof.** Since the result is independent of interval  $I$  and of  $\phi^{(k-1)}(0)$  sign, let us suppose that  $I = \mathbb{R}$  with  $|\phi^{(k)}(v)| \geq \delta > 0$  on  $\mathbb{R}$ , and  $\phi^{(k-1)}(0) \leq 0$ .

We first treat the case  $k = 1$ . If  $\phi'(v)$  stays positive, we have  $\phi(0) + \delta v \leq \phi(v)$  for  $0 \leq v$  and since  $\phi(0) \leq 0$ , there exists a unique  $c \geq 0$  such that  $\phi(c) = 0$ . In the other case,  $\phi'(v)$  stays negative, and we find a unique  $c \leq 0$  such that  $\phi(c) = 0$ . Then  $|\phi(v)| \geq \delta|v - c|$  for all  $v$ , and  $|\phi(v)| \leq \varepsilon$  implies  $|v - c| \leq \varepsilon/\delta$  i.e.  $Z(\phi, \varepsilon) \subset [c - \varepsilon/\delta, c + \varepsilon/\delta]$ . So the lemma is proved for  $k = 1$  with  $\bar{c}_1 = 2$ .

We now prove the lemma by induction on  $k$ . Let us suppose that the case  $k$  is known. As for  $k = 1$ , there exists a unique  $c$  such that  $\phi^{(k)}(c) = 0$ . Thus for all  $v$  we have  $|\phi^{(k)}(v)| \geq \delta|v - c|$ . Let  $\eta > 0$  and set  $W = Z(\phi, \varepsilon) \cap [c - \eta, c + \eta]$ ,  $U = Z(\phi, \varepsilon) \cap (]-\infty, c - \eta[ \cup ]c + \eta, +\infty[)$ . We have  $\text{meas}(W) \leq 2\eta$  and by our inductive hypothesis, since  $|\phi^{(k)}(v)| \geq \delta|v - c| \geq \delta\eta$  on  $U$ ,  $\text{meas}(U) \leq \bar{c}_k (\varepsilon/(\delta\eta))^{1/k}$ . Now the relation

$Z(\phi, \varepsilon) = W \cup U$  gives  $\text{meas}(Z(\phi, \varepsilon)) \leq \inf_{\eta > 0} (2\eta + \bar{c}_k(\varepsilon/(\delta\eta))^{1/k})$  which implies by a simple computation of the minimum that  $\text{meas}(Z(\phi, \varepsilon)) \leq \bar{c}_{k+1}(\varepsilon/\delta)^{1/(k+1)}$ , where  $\bar{c}_{k+1} = 2^{1/(k+1)}(k+1)k^{1/(k+1)-1}\bar{c}_k^{1-1/(k+1)}$  which concludes the proof.  $\square$

**Lemma 4.** *Let  $P$  be a compact set of parameters,  $k$  a positive integer,  $A > 0$ ,  $V = [-A, A]$ ,  $K = V \times P$ ,  $\phi(v; p) \in C^0(P, C^k(V, \mathbb{R}))$ , such that, for all  $(v, p)$  in the compact  $K$ , we have:*

$$\sum_{j=1}^k \left| \frac{\partial^j \phi}{\partial v^j} \right| (v; p) > 0.$$

Let  $Z(\phi(\cdot; p), \varepsilon) = \{v \in V, |\phi(v; p)| \leq \varepsilon\}$ , then there exists a constant  $C$  such that

$$\sup_{p \in P} \text{meas}(Z(\phi(\cdot; p), \varepsilon)) \leq C\varepsilon^{1/k}.$$

**Proof.** Since  $K$  is a compact set, we can choose  $0 < \delta \leq 1$  such that, everywhere on  $K$ , we have  $0 < 2\delta < \frac{1}{k} \sum_{i=1}^k \left| \frac{\partial^i \phi}{\partial v^i} \right| (v; p)$ .

For each  $(v; p) \in K$ , there exists an integer  $i \in \{1, \dots, k\}$ , a number  $r > 0$  and an open set  $O_p$  with  $p \in O_p \subset P$  such that  $|\partial_v^i \phi| > \delta$  on  $U(v, p) = ]v - r, v + r[ \times O_p$ . Therefore, we have

$$\text{meas}(Z(\phi(\cdot; p), \varepsilon) \cap ]v - r, v + r]) \leq \bar{c}_i(\varepsilon/\delta)^{1/i} \leq \bar{c}\varepsilon^{1/k}/\delta$$

using Lemma 3, where  $\bar{c} = \max_{i=1, \dots, k} \bar{c}_i$ .

By compactness of  $K$ , there exists a finite number of such sets  $U_j = U(v_j, p_j)$  such that  $K \subset \bigcup_{j=1}^v U_j$ . Thus, for each  $p$ ,  $Z(\phi(\cdot; p), \varepsilon)$  intersects at most  $v$  intervals  $]v_j - r_j, v_j + r_j[$  where Lemma 3 is applied. This allows to write  $\text{meas}(Z(\phi(\cdot; p), \varepsilon)) \leq v\varepsilon^{1/k}/\delta$  for all  $p$  and to conclude the proof.  $\square$

**Lemma 5.** *Let  $a(v)$  be the field  $(v^1, v^2, \dots, v^N)$  then  $\alpha_{a(\cdot)} = 1/N$ .*

**Proof.** From Lemma 2, we have yet  $\alpha_{a(\cdot)} \leq 1/N$ . So, we just have to prove that  $\alpha = 1/N$  satisfies (1.4) to conclude.

For all  $v$ ,  $\text{rank}\{a'(v), \dots, a^{(N)}(v)\} = N$ , thus it is impossible to find  $\sigma \in S^{N-1}$  such that  $\sigma \perp \{a'(v), \dots, a^{(N)}(v)\}$ . Let  $\varphi(v; u, \sigma)$  be  $a(v) \cdot \sigma - u$ . Since  $\partial_v^j \varphi(v; u, \sigma) = a^{(j)}(v) \cdot \sigma$  for  $j \geq 1$ , we have everywhere  $\sum_{j=1}^N |\partial_v^j \varphi(v; u, \sigma)| > 0$ .

Furthermore, for  $|u| > 1 + a_{\max}$ , where  $a_{\max} = \sup_{|v| \leq A} |a(v)|$ , we have  $|\varphi(v; u, \sigma)| > 1$  for any  $v \in [-A, A]$  and  $\sigma \in S^{N-1}$ . Thus we can apply Lemma 4 with  $0 < \varepsilon \leq 1$  on the compact set  $[-A, A] \times [-a_{\max} - 1, a_{\max} + 1]_u \times S_{\sigma}^{N-1}$  which concludes the proof with  $\alpha_{a(\cdot)} = 1/N$ .  $\square$

**Proof of Proposition 7.** With Lemma 2, we have  $\alpha_{opt}(N, 1) \leq 1/N$ . From Lemma 5, necessarily  $\alpha_{opt}(N, 1) = 1/N$  which concludes the proof.  $\square$

## 4.2. $M = N$

The case when space dimension is equal to velocity dimension is the most physical one and then is very important. In this case, we can get the best smoothing effect with  $\alpha = 1$ .

**Proposition 8.** *For  $N = M$ , we have  $\alpha_{opt}(N, N) = 1$ .*

**Proof.** Since  $\alpha \leq 1$ , it is enough to find  $a(\cdot)$  such that  $\alpha = 1$ .

Let  $a(\cdot) : \mathbb{R}_v^N \rightarrow \mathbb{R}_x^N$  be a global diffeomorphism,  $A > 0$ ,  $(u, \sigma) \in S^N$  and  $\varphi(v) = a(v) \cdot \sigma - u$ . Let  $Z(\varphi, \varepsilon) = \{|v| \leq A, |\varphi(v)| \leq \varepsilon\}$ . Since  $Da(v) \in GL_N(\mathbb{R})$ , and  $\sigma \neq 0$ , then  $\nabla_v \varphi \neq 0$  and the set  $Z(\varphi, 0)$  is empty or a manifold of dimension  $N - 1$ .

Notice that for any  $v$ , there exists  $(u, \sigma) \in S^N$  such that  $a(v) \cdot \sigma - u = 0$ , i.e.  $Z(\varphi, 0) \neq \emptyset$ . For instance, let  $\tilde{\sigma}$  belong to  $S^{N-1}$  and set  $\tilde{u} = a(v) \cdot \tilde{\sigma}$ , then  $(u, \sigma) = \frac{1}{\sqrt{\tilde{u}^2 + 1}}(\tilde{u}, \tilde{\sigma})$  satisfies the conditions.

We thus consider that  $Z(\varphi, 0)$  is not empty.



There exists  $\delta$  such that  $0 < \delta < |\nabla_v \varphi(v)| < 1/\delta$  for all  $|v| \leq A$ ,  $u^2 + |\sigma|^2 = 1$ .

Using the mean inequality, we obtain  $\delta|v - v'| \leq |\varphi(v) - \varphi(v')| \leq \frac{|v-v'|}{\delta}$ , which implies for all  $\varepsilon < 1$ , with  $B(x, r) = \{y, |x - y| \leq r\} \subset \mathbb{R}^N$ , that

$$\bigcup_{z \in Z(\varphi, 0)} B(z, \delta\varepsilon) \subset Z(\varphi, \varepsilon) \subset \bigcup_{z \in Z(\varphi, 0)} B(z, \varepsilon/\delta),$$

and  $Z(\varphi, 0)$  is diffeomorph to a piece of a hyperplane, so  $\text{meas}(Z(\varphi, \varepsilon))$  is of order  $\varepsilon$ . More precisely, there exists a constant  $C > 0$ , only dependent on  $A$ ,  $\delta$  and  $\|Da(\cdot)\|_{B(0, A)}$  such that  $0 < C < \frac{\text{meas}(Z(\varphi, \varepsilon))}{\varepsilon} < C^{-1}$ .

Notice that if  $a(\cdot)$  is a local diffeomorphism,  $\alpha$  is still 1.  $\square$

Incidentally, we also have  $\alpha_{opt}(N, M) = 1$  for all  $M \geq N$ .

## 5. Theorem in the $L^p$ framework

Let us now deal with  $L^p$  case. It will be an interpolation result of the  $L^2$  obtained bound and an estimate in  $L^1$  using some operators in Hardy spaces. We note  $\mathcal{H}^1(\mathbb{R}^{N+1})$  the Hardy space and  $\mathcal{H}^1(\mathbb{R}^N \times \mathbb{R})$  the product Hardy space as done in [2] (see [26] for more details about such spaces).

We will use the two following propositions. The first one is an interpolation result (see [22,2] and [5]) and the second one is about multiplier [2].

**Proposition 9** (Bézard, Interpolation). *Let  $T$  be a  $\mathbb{C}$ -linear operator, bounded in,*

$$L^2(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M) \rightarrow W^{\beta, 2}(\mathbb{R}_t \times \mathbb{R}_x^N),$$

and in

$$L^1(\mathbb{R}_v^M, \mathcal{H}^1(\mathbb{R}^N \times \mathbb{R})) \rightarrow \mathcal{H}^1(\mathbb{R}_{t,x}^{N+1}),$$

for some  $\gamma \geq 0$ . Then  $T$  is bounded,

$$L^p(\mathbb{R}_t \times \mathbb{R}_x^N \times \mathbb{R}_v^M) \rightarrow W^{s, p}(\mathbb{R}_t \times \mathbb{R}_x^N),$$

for  $1 < p \leq 2$ , with  $s = 2\beta/p'$ .

**Proposition 10** (Bézard, Multiplier on  $\mathcal{H}^1$ ). *Let  $m(y, y_{N+1})$  be a function of  $(y, y_{n+1}) \in \mathbb{R}^N \times \mathbb{R}$  which is  $C^\infty$  out of  $[y = 0 \text{ or } y_{N+1} = 0]$ , and verifying for all  $\alpha, \beta$ ,*

$$|\partial_y^\alpha \partial_{y_{N+1}}^\beta m(y, y_{N+1})| \leq \frac{C_{\alpha\beta}}{|y|^\alpha |y_{N+1}|^\beta},$$

then  $m$  defines a bounded Fourier multiplier on  $\mathcal{H}^1(\mathbb{R}^N \times \mathbb{R})$ .

**Proof of Theorems 4 and 5.** For Theorem 4 (respectively Theorem 5), we use the averaging lemma of Theorem 1 (respectively Theorem 2) which gives that  $T(f, g) = \rho_\psi$  is bounded from  $L^2$  to  $H_{loc}^{\alpha/2}$  (respectively  $H^{1/\gamma}$ ).

We now focus on estimate in  $L^1$ . We denote by  $\mathcal{F}$  the Fourier transform with respect to  $X$ . Taking this Fourier transform in  $b(v) \cdot \nabla_X f + F(X) \cdot \nabla_v f = g$ , we have:

$$\mathcal{F}(f) = \frac{\mathcal{F}(g) - \mathcal{F}(F \cdot \nabla_v f)}{i(b(v) \cdot Y)}.$$

Let  $\chi \in C_c^\infty(\mathbb{R})$ ,  $\chi(0) = 1$ ,  $\chi'(0) = 0$  and  $\chi''(0) \neq 0$  be an even, non-increasing function in  $[0, +\infty[$ . We set  $L$  such that  $\text{supp } \chi \subset [-L, L]$ . We have:

$$\begin{aligned} f(Y, v) &= \mathcal{F}^{-1}[\chi(b(v) \cdot Y)\mathcal{F}(f)(Y, v) + (1 - \chi(b(v) \cdot Y))\mathcal{F}(f)(Y, v)] \\ &= \mathcal{F}^{-1}[\chi(b(v) \cdot Y)\mathcal{F}(f)(Y, v)] + \mathcal{F}^{-1}\left[(1 - \chi(b(v) \cdot Y))\frac{\mathcal{F}(g) - \mathcal{F}(F \cdot \nabla_v f)}{i(b(v) \cdot Y)}\right], \end{aligned}$$

and then, in order to bound operator  $f \mapsto \int_{\mathbb{R}^M} f(Y, v) \psi(v) dv$ , we have to bound the three following operators:

$$Q : f \mapsto \int_{\mathbb{R}^M} \mathcal{F}^{-1}[\chi(b(v) \cdot Y) \mathcal{F}(f)(Y, v)] \psi(v) dv, \quad (5.1)$$

$$W : g \mapsto \int_{\mathbb{R}^M} \mathcal{F}^{-1} \left[ \frac{1 - \chi(b(v) \cdot Y)}{i(b(v) \cdot Y)} \mathcal{F}(g)(Y, v) \right] \psi(v) dv, \quad (5.2)$$

and

$$R : f \mapsto - \int_{\mathbb{R}^M} \mathcal{F}^{-1} \left[ \frac{1 - \chi(b(v) \cdot Y)}{i(b(v) \cdot Y)} \mathcal{F}(F \cdot \nabla_v f)(Y, v) \right] \psi(v) dv. \quad (5.3)$$

As in the classical case (by this we refer to [2,5]), we transform the operators in order for them to involve only one direction in  $X$ . Indeed, the manipulation of product structure for Hardy space which depends on a moving direction is difficult to deal with. Thus, for any  $v$ , we take  $R_v$  an orthogonal transform in  $\mathbb{R}^{N+1}$  such that

$$R_v \left( \frac{b(v)}{|b(v)|} \right) = e_{N+1},$$

where  $e_{N+1}$  is the very last vector of the canonical base, and we set:

$$f_*(X, v) = f(R_v^{-1}(X), v),$$

and

$$Q_* f_* = Qf.$$

Since  $f \mapsto f_*$  is an isometry on  $L^p_{Xv}$ , we have now to study  $Q_*$  instead of  $Q$ . We perform similar transformations for the two other operators and we get  $W_*$  and  $R_*$ .

For the two first operators, as in the classical proof, we have:

$$\|Qf\|_{\mathcal{H}^1(\mathbb{R}^{N+1})} \leq C \|f\|_{L^1(\mathbb{R}_v^M, \mathcal{H}^1(\mathbb{R}^N \times \mathbb{R}))},$$

and

$$\|Wg\|_{\mathcal{H}^1(\mathbb{R}^{N+1})} \leq C \|g\|_{L^1(\mathbb{R}_v^M, \mathcal{H}^1(\mathbb{R}^N \times \mathbb{R}))}.$$

The new term is the third one (operator  $R$ ). We use the following rewrite of  $R(f)$  in order to bound it. This is

$$\begin{aligned} (Rf)(Y) &= - \mathcal{F}^{-1} \int_{\mathbb{R}^M} \left[ \frac{1 - \chi(b(v) \cdot Y)}{i(b(v) \cdot Y)} F \cdot \nabla_v \mathcal{F}(f)(Y, v) \right] \psi(v) dv \\ &= \mathcal{F}^{-1} \left( F \cdot \int_{\mathbb{R}^M} \mathcal{F}(f)(Y, v) \nabla_v \left[ \frac{1 - \chi(b(v) \cdot Y)}{i(b(v) \cdot Y)} \psi(v) \right] dv \right) \\ &= \mathcal{F}^{-1} \left( F \cdot \int_{\mathbb{R}^M} \mathcal{F}(f)(Y, v) \frac{1 - \chi(b(v) \cdot Y)}{i(b(v) \cdot Y)} \nabla_v \psi(v) dv \right) \\ &\quad + \mathcal{F}^{-1} \left( F \cdot \int_{\mathbb{R}^M} \mathcal{F}(f)(Y, v) m_0(b(v) \cdot Y) \nabla_v (b(v) \cdot Y) \psi(v) dv \right), \end{aligned} \quad (5.4)$$

with

$$m_0(y) = \frac{-y\chi'(y) - 1 + \chi(y)}{iy^2}. \quad (5.5)$$

We denote by  $\mathcal{F}(R_1 f)$  and  $\mathcal{F}(R_2 f)$  the two terms of this decomposition. We perform as previously orthogonal transformations and we have to study the obtained  $(R_1)_*$  and  $(R_2)_*$ .

The term  $(R_1)_*$  is the same than  $W_*$  but with  $\nabla_v \psi$  instead of  $\psi$ . Thus we have the same result thanks to the regularity assumption on  $\psi$ .

Now, setting  $T = m_0 \nabla_v$ , we have:

$$\begin{aligned} (R_2)_*(f_*)(Y) &= F \cdot \int_{\mathbb{R}^M} \mathcal{F}^{-1}(\mathcal{F}(f_*)(R_v(Y), v) T(b(v) \cdot Y)) \psi(v) dv \\ &= F \cdot \int_{\mathbb{R}^M} \mathcal{F}^{-1}(\mathcal{F}(f_*)(R_v(Y), v) T(R_v(b(v)) \cdot R_v(Y))) \psi(v) dv \\ &= F \cdot \int_{\mathbb{R}^M} \mathcal{F}^{-1}(\mathcal{F}(f_*)(R_v(Y), v) T(|b(v)|e_{N+1} \cdot R_v(Y))) \psi(v) dv, \end{aligned}$$

thus, setting  $T_j = m_0 \partial_{v_j}$ , we get,

$$\begin{aligned} \|(R_2)_*(f_*)\|_{\mathcal{H}^1(\mathbb{R}^{N+1})} &\leq \sum_j |F_j| \int_{\mathbb{R}^M} \|\mathcal{F}^{-1}(\mathcal{F}(f_*)(R_v(Y), v) T_j(|b(v)|e_{N+1} \cdot R_v(Y)))\|_{\mathcal{H}^1(\mathbb{R}^{N+1})} |\psi(v)| dv \\ &\leq \sum_j |F_j| \int_{\mathbb{R}^M} \|\mathcal{F}^{-1}(\mathcal{F}(f_*)(Y, v) T_j(|b(v)|e_{N+1} \cdot Y))\|_{\mathcal{H}^1(\mathbb{R}^{N+1})} |\psi(v)| dv \\ &\leq C_1 \sum_j |F_j| \int_{\mathbb{R}^M} \|\mathcal{F}^{-1}(\mathcal{F}(f_*)(Y, v) T_j(|b(v)|e_{N+1} \cdot Y))\|_{\mathcal{H}^1(\mathbb{R}^N \times \mathbb{R})} |\psi(v)| dv, \end{aligned}$$

using the invariance under orthogonal transformation in  $\mathcal{H}^1(\mathbb{R}^{N+1})$  and thanks to the continuous injection of  $\mathcal{H}^1(\mathbb{R}^N \times \mathbb{R})$  in  $\mathcal{H}^1(\mathbb{R}^{N+1})$ .

We use now Proposition 10 with the term:

$$m_j(y, y_{N+1}) = T_j(|b(v)|e_{N+1} \cdot Y) = m_0(|b(v)|_{y_{N+1}}) \partial_{v_j}(|b(v)|)_{y_{N+1}}, \quad \text{for } j = 1, \dots, M.$$

Those terms rewrite

$$m_j(y, y_{N+1}) = m_0(|b(v)|_{y_{N+1}}) \frac{a(v) \cdot \partial_{v_j} a(v)}{|b(v)|} y_{N+1}.$$

Now  $m_0(z) \rightarrow_{z \rightarrow 0} -\frac{1}{2i} \chi''(0)$ , therefore  $m_0$  is  $C^\infty$ . The terms in (5.5) with  $\chi$  have a compact support and the other term is  $1/y^2$ , then every derivatives of  $m_0$  is bounded at infinity.

We differentiate  $m_j$  with respect to  $y_{N+1}$ , it gives:

$$\partial_{y_{N+1}}^k m_j(y, y_{N+1}) = \frac{a(v) \cdot \partial_{v_j} a(v)}{|b(v)|} (m_0^{(k)}(|b(v)|_{y_{N+1}}) |b(v)|^k y_{N+1} + k m_0^{(k-1)}(|b(v)|_{y_{N+1}}) |b(v)|^{k-1}).$$

There exists some constants  $C$  and  $C_k$  such that

$$|b(v)| \leq C, \quad |b(v)|^{k-2} |a(v) \cdot \partial_{v_j} a(v)| \leq C_k$$

for  $v$  in the compact support of  $\psi$ . Thus

$$|\partial_{y_{N+1}}^k m_j(y, y_{N+1})| |y_{N+1}|^k \leq C_k (C m_0^{(k)}(|b(v)|_{y_{N+1}}) y_{N+1} + k m_0^{(k-1)}(|b(v)|_{y_{N+1}})).$$

For  $|y_{N+1}| \geq (R+1)/C$ , we have  $m_0^{(j)}(|b(v)|_{y_{N+1}}) = 0$  for any  $j$ , and then

$$m_0^{(k)}(|b(v)|_{y_{N+1}}) y_{N+1} + k m_0^{(k-1)}(|b(v)|_{y_{N+1}}) = 0 \quad \text{for } |y_{N+1}| \geq (R+1)/C.$$

Furthermore  $|m_0^{(k)}(|b(v)|_{y_{N+1}}) y_{N+1} + k m_0^{(k-1)}(|b(v)|_{y_{N+1}})| \leq \|m_0^{(k)}\|_\infty \frac{R+1}{C} + k \|m_0^{(k-1)}\|_\infty$  for  $|y_{N+1}| < (R+1)/C$ . Finally, for any  $(y, y_{N+1})$ , we get:

$$|\partial_{y_{N+1}}^k m_j(y, y_{N+1})| |y_{N+1}|^k \leq C_k (\|m_0^{(k)}\|_\infty (R+1) + k \|m_0^{(k-1)}\|_\infty)$$

uniformly with respect to  $v$  in the support of  $\psi$ . Then, we can apply Proposition 10 to get the boundary of  $(R_2)_*$ .

The interpolation result concludes, since  $\beta = \alpha/2$  (respectively  $\beta = 1/\gamma$ ), that the obtained regularity is  $s = \alpha/p'$  (respectively  $s = 2/(\gamma p')$ ).  $\square$

## Acknowledgements

We thank Gilles Lebeau and Jeffrey Rauch for fruitful discussions on oscillatory integrals. We also thank referees for their useful comments.

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