



HAL
open science

Dynamics of periodic Toda chains with a large number of particles

Dario Bambusi, Thomas Kappeler, Thierry Paul

► **To cite this version:**

Dario Bambusi, Thomas Kappeler, Thierry Paul. Dynamics of periodic Toda chains with a large number of particles. *Journal of Differential Equations*, 2015, 258 (12), pp.4103-4490. hal-00864330

HAL Id: hal-00864330

<https://hal.science/hal-00864330v1>

Submitted on 20 Sep 2013

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Dynamics of periodic Toda chains with a large number of particles

D. Bambusi*, T. Kappeler†, T. Paul‡

September 21, 2013

Abstract

For periodic Toda chains with a large number N of particles we consider states which are N^{-2} -close to the equilibrium and constructed by discretizing any given C^2 -functions with mesh size N^{-1} . For such states we derive asymptotic expansions of the Toda frequencies $(\omega_n^N)_{0 < n < N}$ and the actions $(I_n^N)_{0 < n < N}$, both listed in the standard way, in powers of N^{-1} as $N \rightarrow \infty$. At the two edges $n \sim 1$ and $N - n \sim 1$, the expansions of the frequencies are computed up to order N^{-3} with an error term of higher order. Specifically, the coefficients of the expansions of ω_n^N and ω_{N-n}^N at order N^{-3} are given by a constant multiple of the n 'th KdV frequencies ω_n^- and ω_n^+ of two periodic potentials, q_- respectively q_+ , constructed in terms of the states considered. The frequencies ω_n^N for n away from the edges are shown to be asymptotically close to the frequencies of the equilibrium. For the actions $(I_n^N)_{0 < n < N}$, asymptotics of a similar nature are derived.

1 Introduction

In this paper we study the asymptotics of the dynamics of periodic Toda chains with a large number of particles of equal mass for initial data close

*Dipartimento di Matematica, Università degli Studi di Milano, Via Saldini 50, I-20133 Milano

†Institut für Mathematik, Universität Zürich, Winterthurerstrasse 190, CH-8057 Zürich. Supported in part by the Swiss National Science Foundation.

‡CNRS and CMLS, École Polytechnique, F-91128 Palaiseau

to the equilibrium. Toda chains were introduced by Toda [20] as a class of special Fermi Pasta Ulam (FPU) chains with the main feature that they are *integrable* Hamiltonian systems. The Hamiltonian of the Toda chain with N particles is given by

$$\mathcal{H} = \frac{1}{2} \sum_{n=1}^N p_n^2 + \sum_{n=1}^N e^{q_n - q_{n+1}}$$

where q_n denotes the displacement of the n 'th particle from its equilibrium position and p_n its momentum. It is convenient to define (q_n, p_n) for any $n \in \mathbb{Z}$ by requiring that $(q_{n+N}, p_{n+N}) = (q_n, p_n) \forall n \in \mathbb{Z}$. Note that the total momentum $\sum_{i=1}^N p_i$ is a conserved quantity of the Toda flow and hence the Hamiltonian equations of motion imply that the center of mass $\frac{1}{N} \sum_{n=1}^N q_n$ moves at constant speed. Hence when considered relative to the motion of its center of mass, the Toda chain has $N - 1$ degrees of freedom.

Flaschka [9] introduced the variables $b = (b_n)_{1 \leq n \leq N}$, $a = (a_n)_{1 \leq n \leq N}$ defined by $b_n = -p_n \in \mathbb{R}$, $a_n = e^{\frac{1}{2}(q_n - q_{n+1})} > 0$ and showed that when evolved along the Toda flow the corresponding equations of motion of b, a can be described by a Lax pair (L, B) ,

$$\dot{L} = [B, L] \tag{1.1}$$

where the $N \times N$ matrices $L = L(b, a)$ and $B = B(a)$ are given by

$$\begin{pmatrix} b_1 & a_1 & 0 & \dots & 0 & a_N \\ a_1 & b_2 & a_2 & \dots & 0 & 0 \\ 0 & a_2 & b_3 & \dots & 0 & 0 \\ \vdots & & & & & \vdots \\ 0 & \dots & \dots & b_{N-1} & a_{N-1} & \\ a_N & 0 & \vdots & a_{N-1} & b_N & \end{pmatrix} \text{ and } \begin{pmatrix} 0 & a_1 & 0 & \dots & 0 & -a_N \\ -a_1 & 0 & a_2 & \dots & 0 & 0 \\ 0 & -a_2 & 0 & \dots & 0 & 0 \\ \vdots & & & & & \vdots \\ 0 & \dots & \dots & 0 & a_{N-1} & \\ a_N & 0 & \vdots & -a_{N-1} & 0 & \end{pmatrix}$$

respectively. Note that by the periodicity of the q_n 's, $\prod_{n=1}^N a_n = 1$ and hence the position of the center of mass cannot be recovered from the a_n 's. However the latter is easily reconstructed from its initial condition and the total momentum $-\sum_{n=1}^N b_n$. In the sequel, we exclusively concentrate on the system described by (1.1), viewing it as equivalent to the Toda chain and referring to it by the same name. Actually, it is convenient to choose a slightly larger space for the variable a by requiring only that $a = (a_n)_{1 \leq n \leq N} \in \mathbb{R}_{>0}^N$. First we note that b, a are not canonical variables. In fact, when expressing the Hamiltonian equations of motion of the Toda chain in terms of the

variables b, a , the corresponding Poisson bracket is degenerate. The level sets

$$\mathcal{L}_{c_1, c_2} := \{(b, a) \in \mathbb{R}^N \times \mathbb{R}_{>0}^N : \sum_{n=1}^N b_n = c_1; \prod_{n=1}^N a_n = c_2\}$$

are the corresponding symplectic leaves. As already mentioned, when restricted to such a leaf, the system has $N - 1$ degrees of freedom. Using (1.1), it can be shown to be an integrable system of $N - 1$ coupled oscillators, meaning that it admits globally defined Birkhoff coordinates – see [13] for details. In particular, its invariant manifolds are smooth tori. As a consequence, the dynamics of such chains is quasi-periodic in time. Furthermore, in terms of the spectrum of $L(b, a)$, which by (1.1) is preserved by the Toda flow, there is a canonical choice of $N - 1$ globally defined actions $I = (I_n)_{0 < n < N} \in \mathbb{R}_{\geq 0}^{N-1}$ – see [13] for details as well as Section 3. They parametrize the invariant tori \mathcal{T}_I and allow to read off their dimension, $\dim \mathcal{T}_I = |\{0 < n < N | I_n > 0\}|$. Furthermore, the Toda Hamiltonian \mathcal{H} can be expressed as a function of I which we again denote by \mathcal{H} . In fact, \mathcal{H} is a real analytic function of I and the total momentum. The frequencies corresponding to this canonical choice of actions are denoted by $\omega_n(I) = \partial_{I_n} \mathcal{H}$. We refer to $\omega(I) = (\omega_n(I))_{0 < n < N}$ as the frequency vector corresponding to I . As an aside we mention that actions are uniquely determined only up to unimodular transformations. Nevertheless, the dynamics on such a torus can be described in a coordinate free way in terms of the frequency map. More precisely, let $x = (b, a)$ be an arbitrary point on a leaf \mathcal{L}_{c_1, c_2} and denote by $\mathcal{T}(x)$ the invariant torus containing x . Let $\mathcal{M} \equiv \mathcal{M}(x)$ be the frequency module of $\mathcal{T}(x)$, consisting of all integer combinations of $\omega_1, \dots, \omega_{N-1}$,

$$\mathcal{M} = \{k \cdot \omega : k \in \mathbb{Z}^{N-1}\} \subset \mathbb{R}.$$

Note that \mathcal{M} doesn't depend on the choice of the actions and thus is invariantly defined. It plays an important role when analyzing data of the evolution of systems such as Toda chains – see e.g. [18]. For any given complex valued function f defined on \mathcal{L}_{c_1, c_2} and any $x \in \mathcal{L}_{c_1, c_2}$ let $f_x(t) := f(\Phi^t(x))$ where $x \mapsto \Phi^t(x)$ denotes the Toda flow. The support of the Fourier transform of $f_x(t)$ with respect to the real variable t is then contained in the frequency module \mathcal{M} and if f is sufficiently regular, $f_x(t)$ takes the form

$$f_x(t) = \sum_{k \in \mathbb{Z}^{N-1}} a_x(k) \exp(i\varphi_x(k)t)$$

where $\varphi_x : \mathbb{Z}^{N-1} \rightarrow \mathcal{M}$ is a linear map and $a_x : \mathbb{Z}^{N-1} \rightarrow \mathbb{C}$ has appropriate decay conditions. In the case at hand, φ_x can be chosen to be $\varphi_x(k) = k \cdot \omega(I)$ where I is the action of $\mathcal{T}(x)$. For more details see e.g. [16], Section 3.

The aim of this paper is to compute the asymptotics as $N \rightarrow \infty$ of the frequencies and actions of states which are N^{-2} -close to the equilibrium $0_N = (0, \dots, 0)$, $1_N = (1, \dots, 1)$. As KdV frequencies will play an important role for describing these asymptotics we first need to review some features of the KdV equation in the periodic set-up. It turns out that for our purposes, the period of the space variable is equal to $\frac{1}{2}$. Recall that the (generalized) KdV equation with parameters $d_1, d_2 \in \mathbb{R}$

$$\partial_t q = d_1(-\partial_x^3 q + 6q\partial_x q) + d_2\partial_x q, \quad t \in \mathbb{R}, \quad x \in \mathbb{R}/(\mathbb{Z}/2)$$

is an integrable Hamiltonian PDE with Hamiltonian

$$\mathcal{H}_{d_1, d_2} = d_1 \int_0^{\frac{1}{2}} \left(\frac{1}{2}(\partial_x q)^2 + q^3 \right) dx + d_2 \frac{1}{2} \int_0^{\frac{1}{2}} q^2 dx$$

Note that given any solution, the average is a conserved quantity. For our purposes it suffices to consider solutions in Sobolev spaces $H_0^s, s \geq 2$, of elements in H^s with mean zero. Such solutions exist globally in time and are almost periodic, evolving on invariant sets which are tori, generically of infinite dimension. These invariant tori are parametrized by globally defined actions which have the property that on H_0^1 , the KdV Hamiltonian can be expressed as a real analytic function of them alone. Furthermore, the KdV frequencies are given by the partial derivatives of the KdV Hamiltonian with respect to these actions. See [16] for details. Note that \mathcal{H}_{d_1, d_2} is a linear combination of the first two Hamiltonians \mathcal{H}_1 and \mathcal{H}_2 of the so called KdV hierarchy, where

$$\mathcal{H}_1(q) = \frac{1}{2} \int_0^{\frac{1}{2}} q^2 dx \quad \text{and} \quad \mathcal{H}_2(q) = \int_0^{\frac{1}{2}} \left(\frac{1}{2}(\partial_x q)^2 + q^3 \right) dx.$$

To state our first result, assume that β, α are in the space $C_0^2(\mathbb{T}) \equiv C_0^2(\mathbb{T}, \mathbb{R})$ of one periodic, real valued functions of class C^2 and mean 0. The space $C_0^2(\mathbb{T})$ is endowed with the standard supremum norm $\|f\|_{C^2}$.

For any $N \geq 3$, we then introduce the periodic Toda chain with N particles, defined in terms of Flaschka coordinates by

$$b_n^N = \frac{1}{4N^2} \beta\left(\frac{n}{N}\right) \quad \text{and} \quad a_n^N = 1 + \frac{1}{4N^2} \alpha\left(\frac{n}{N}\right). \quad (1.2)$$

Alternatively, one can consider

$$p_n^N = -\frac{1}{4N^2}\beta\left(\frac{n}{N}\right) \quad \text{and} \quad q_n^N = -\frac{2}{4N}\xi\left(\frac{n}{N}\right)$$

where ξ is the element in $C_0^3(\mathbb{T})$, satisfying $\xi' = \alpha$. Using that

$$\exp\left(\frac{q_n^N - q_{n+1}^N}{2}\right) = 1 + \frac{q_n^N - q_{n+1}^N}{2} + O(N^{-4}) = a_n^N + O(N^{-3})$$

one can show that our results stated below hold for either of the two discretizations. See the remark at the end of Section 2 for details. In this paper, we concentrate on the case where the data is given in the form (1.2) Denote by ω_n^N and I_n^N , $0 < n < N$, the frequencies and actions, corresponding to (b^N, a^N) . To describe their asymptotics as $N \rightarrow \infty$ introduce the two potentials

$$q_{\pm}(x) = -2\alpha(2x) \mp \beta(2x).$$

Note that they have period $\frac{1}{2}$ and are of class C^2 . Let I_j^{\pm} and $\omega_j^{\pm} = \partial_{I_j^{\pm}} \mathcal{H}_2$, $j \geq 1$, be the corresponding KdV actions and frequencies and define the following sequence of KdV Hamiltonians ($N \geq 2$)

$$\mathcal{H}_{KdV}^N := \frac{1}{2N} \mathcal{H}_1 - \frac{1}{24} \frac{1}{(2N)^3} \mathcal{H}_2$$

whose frequencies can be computed to be

$$\partial_{I_n^{\pm}} \mathcal{H}_{KdV}^N(q_{\pm}) = \frac{2\pi n}{N} - \frac{1}{24} \frac{1}{(2N)^3} \omega_n^{\pm}.$$

To describe the asymptotics of the frequencies ω_n^N for $n \sim 1$ and $n \sim N - 1$ consider functions $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$ satisfying

$$(F) \quad \lim_{N \rightarrow \infty} F(N) = \infty; \quad F \text{ increasing}; \quad F(N) \leq N^{\eta} \quad \text{with} \quad \eta > 0.$$

Theorem 1.1 *Let $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$ satisfy (F) with $\eta \leq 1/3$ and set $M = [F(N)]$. Then the asymptotics of the frequencies $(\omega_n^N)_{0 < n < N}$ of (b^N, a^N) defined by (1.2) are as follows:*

at the left and right edges: for $1 \leq n \leq F(M)$

$$\omega_n^N = \partial_{I_n^-} \mathcal{H}_{KdV}^N + O\left(\frac{1}{N^3} \left(\frac{n^2 F(M)}{M^{1/2}} + \frac{1}{F(M)^{5/2}}\right)\right) \quad (1.3)$$

$$\omega_{N-n}^N = \partial_{I_n^+} \mathcal{H}_{KdV}^N + O\left(\frac{1}{N^3}\left(\frac{n^2 F(M)}{M^{1/2}} + \frac{1}{F(M)^{5/2}}\right)\right) \quad (1.4)$$

in the bulk: $M < n < N - M$

$$\omega_n^N = 2 \sin \frac{\pi n}{N} \left(1 + O\left(\frac{\log M}{M^2}\right)\right). \quad (1.5)$$

Finally, for $F(M) < n \leq M$

$$0 < \omega_n^N = \frac{2\pi n}{N} + O\left(\frac{n^3}{N^3}\right), \quad 0 < \omega_{N-n}^N = \frac{2\pi n}{N} + O\left(\frac{n^3}{N^3}\right). \quad (1.6)$$

These estimates hold uniformly in $0 < n < N$ and uniformly on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

We remark that uniformly for $1 \leq n \leq F(M)$, the error terms in the asymptotics (1.3) and (1.4) have stronger decay as $N \rightarrow \infty$ than the principal terms $\partial_{I_n^\pm} \mathcal{H}_{KdV}^N = \frac{2\pi n}{N} - \frac{1}{24} \frac{1}{(2N)^3} \omega_n^\pm$ in view of the asymptotics $\omega_n^\pm = (4n\pi)^3 + O(1)$ (cf Section 4).

To state the asymptotics of the actions I_n^N , $0 < n < N$, we first need to introduce some additional notation. Recall that the Hill operators $H_\pm := -\partial_x^2 + q_\pm$, associated to the potentials q_\pm , come up in the formulation of the KdV equation in terms of a Lax pair. In particular, the spectrum of H_\pm , when considered with periodic boundary conditions on the interval $[0, 1]$, is preserved by the KdV flow. It is pure point and consists of real eigenvalues which when listed in increasing order and with their multiplicities satisfy

$$\lambda_0^\pm < \lambda_1^\pm \leq \lambda_2^\pm < \dots$$

For any $n \geq 1$, the difference $\gamma_n^\pm := \lambda_{2n}^\pm - \lambda_{2n-1}^\pm$ is referred to as n 'th gap length. It is well known that the decay properties of γ_n^\pm as $n \rightarrow \infty$ are related to the smoothness of q_\pm . In particular, as q_\pm are of class C^2 , one has $\sum_{n=1}^\infty n^4 (\gamma_n^\pm)^2 < \infty$, hence in particular $\gamma_n^\pm = O(n^{-2})$.

Theorem 1.2 *Let $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$ satisfy (F) with $\eta < 1/2$ and set $M = [F(N)]$. Then the asymptotics of the actions $I^N = (I_n^N)_{0 < n < N}$ of the states (b^N, a^N) defined by (1.2) are as follows:*

at the left and right edges: for $1 \leq n \leq F(M)$

$$8N^2 I_n^N = I_n^- + O\left(\frac{M^2 F(M)}{N M^{1/2}} + \frac{M^3}{N^{3/2}} + \gamma_n^- \left(\frac{F(M)}{M^{1/2}} + \frac{M}{N^{1/2}}\right)\right) \quad (1.7)$$

$$8N^2 I_{N-n}^N = I_n^+ + O\left(\frac{M^2 F(M)}{N M^{1/2}} + \frac{M^3}{N^{3/2}} + \gamma_n^+\left(\frac{F(M)}{M^{1/2}} + \frac{M}{N^{1/2}}\right)\right) \quad (1.8)$$

in the bulk: $M < n \leq N/2$

$$I_n^N, I_{N-n}^N = O\left(\frac{1}{nM^2} \frac{1}{N^2}\right)$$

whereas for $F(M) < n \leq M$,

$$I_n^N = O\left(\frac{1}{n}((\gamma_n^-)^2 + \frac{M^4}{N^2}) \frac{1}{N^2}\right), \quad I_{N-n}^N = O\left(\frac{1}{n}((\gamma_n^+)^2 + \frac{M^4}{N^2}) \frac{1}{N^2}\right).$$

These estimates hold uniformly in $0 < n < N$ and uniformly on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

The following remark comments on the implications of Theorem 1.1 and Theorem 1.2 on the approximation of the Toda chain with initial data given by (b^N, a^N) by KdV type solutions. For simplicity we assume that $F(N) = N^\eta$ with $0 < \eta \leq 1/3$. Then $M = [N^\eta]$ and $F(M) \sim N^{\eta^2}$.

Remark 1.3 *In Birkhoff coordinates, the n 'th component $(x_n^N(t), y_n^N(t))$ of the solution, $0 < n < N$, is of the form*

$$(x_n^N(t), y_n^N(t)) = \sqrt{2I_n^N}(\cos(\theta_n^N + t\omega_n^N), \sin(\theta_n^N + t\omega_n^N)) \quad \text{if } I_n^N \neq 0$$

and zero otherwise where θ_n^N is the n 'th angle coordinate, determined by the initial data. For any given N , it evolves on the torus

$$\mathcal{T}^N = \{(x_n^N, y_n^N)_{0 < n < N} \mid (x_n^N)^2 + (y_n^N)^2 = 2I_n^N \quad \forall 0 < n < N\}.$$

First note that according to Theorem 1.2 the size of the components of the solution in the bulk is small in the sense that

$$\sum_{F(M) < n < N-1-F(M)} I_n^N = O\left(\frac{1}{N^2} \left(\frac{1}{N^{5\eta^2}} + \frac{\log N}{N^{2\eta}}\right)\right)$$

Here we used the fact mentioned above that $\sum_{n=1}^{\infty} n^4 (\gamma_n^\pm)^2 < \infty$. Hence for N sufficiently large, the solution of the Toda chain with initial data (b^N, a^N) can be viewed as a small perturbation of a long wave obtained from the solution by setting the components $(x_n^N(t), y_n^N(t))$ with $F(M) < n < N - 1 - F(M)$ to zero.

Secondly we note that it follows from Theorem 1.1 that on a time interval of size larger than N^3 , these long waves are approximated by two KdV type solutions. More precisely, in the case where $0 < \eta < \frac{1}{11}$, one has for any $1 \leq n \leq F(M)$ by (1.3) and (1.4),

$$\omega_n^N - \partial_{I_n^-} \mathcal{H}_{\text{KdV}}^N, \quad \omega_{N-n}^N - \partial_{I_n^+} \mathcal{H}_{\text{KdV}}^N = O\left(\frac{1}{N^3 N^{5\eta^2/2}}\right).$$

Hence the approximation of solutions of Toda chains considered above is valid on a longer time interval than the one obtained in [3] (cf also [21]).

The asymptotics of the Toda frequencies up to order 5 – which in principle is possible – will incorporate in (1.3)-(1.4) further terms, conjecturally involving the third Hamiltonian of the KdV hierarchy, and thus will provide an approximation beyond the one by KdV type solutions.

Outline of the proofs of Theorem 1.1 and Theorem 1.2: A first key ingredient are the asymptotic expansions of the eigenvalues of Jacobi matrices $L(b^N, a^N)$ obtained in our paper [2], having the novel feature that they involve the periodic eigenvalues of two independent Hill operators. Secondly, we rely on the rich structure of periodic Toda chains and the KdV equation related to their integrability. In particular, we use that their spectral curves are of a very similar nature. The proof of Theorem 1.2 involves formulas, representing the Toda and KdV actions as periods of differentials on the corresponding spectral curves where the cycles involved are closely related to the spectrum of the underlying Jacobi matrices respectively Hill operators. The spectral curve as well as the differentials are expressed in terms of the discriminants, associated to the Toda chains respectively KdV equation and the asymptotics of the Toda actions are obtained from the asymptotic expansion of the discriminant $\Delta_N(\mu)$, associated to the sequence of Toda chains (1.2), derived in [2] – see Section 2 for a precise statement of this result. The proof of Theorem 1.1 is more involved. Again the Toda and the KdV frequencies can both be represented as periods of certain differentials on the corresponding spectral curves. Unlike in the case of the actions there are however significant structural differences between the differentials involved in the two cases. To continue we first need to introduce additional notation. It turns out to be more convenient to double the size of the Jacobi matrix $L(b^N, a^N)$ and to consider

$$Q_N^{\alpha, \beta} = L((b^N, b^N), (a^N, a^N)).$$

The eigenvalues of the symmetric $2N \times 2N$ matrix $Q_N^{\alpha,\beta}$ when listed in increasing order and with their multiplicities then satisfy

$$\lambda_0^N < \lambda_1^N \leq \lambda_2^N < \cdots < \lambda_{2N-3}^N \leq \lambda_{2N-2}^N < \lambda_{2N-1}^N.$$

Precise asymptotics as $N \rightarrow \infty$ of the spectrum of $Q_N^{\alpha,\beta}$ were derived in [2] and are recalled in Section 2. Furthermore, let $\chi_N(\mu)$ be the characteristic polynomial of $Q_N^{\alpha,\beta}$

$$\chi_N(\mu) = \prod_{0 \leq k \leq 2N-1} (\mu - \lambda_k^N).$$

By Floquet theory, it equals $\Delta_N^2(\mu) - 4$ up to a μ -independent factor where $\Delta_N(\mu)$ is the discriminant of the Toda chain (b^N, a^N) (see Section 2). This factor has been computed in [2],

$$\Delta_N^2(\mu) - 4 = \mathfrak{q}_N^{-2} \chi_N(\mu) \quad \text{with} \quad \mathfrak{q}_N = \prod_{n=1}^N a_n^N. \quad (1.9)$$

Finally denote by $\sqrt[c]{\Delta_N^2(\mu) - 4}$ the canonical root of $\Delta_N^2(\mu) - 4$, defined for $\mu \in \mathbb{C} \setminus ((-\infty, \lambda_0^N] \cup [\lambda_1^N, \lambda_2^N] \cup \cdots \cup [\lambda_{2N-1}^N, \infty))$ by the sign condition

$$\sqrt[c]{\Delta_N(\mu + i0)^2 - 4} > 0 \quad \forall \mu > \lambda_{2N-1}^N \quad (1.10)$$

and set $\sqrt[c]{\chi_N(\mu)} = \mathfrak{q}_N \sqrt[c]{\Delta_N(\mu)^2 - 4}$. The representation of the Toda frequencies used for deriving their asymptotics is the following one

$$iN\omega_n^N = \int_{\lambda_0^N}^{\lambda_{2n-1}^N} \frac{(\mu - \mathfrak{p}_N/N) \dot{\Delta}_N(\mu) d\mu}{\sqrt[c]{\Delta_N^2(\mu - i0) - 4}} - \sum_{\substack{0 < k < N \\ I_k^N \neq 0}} I_k^N \omega_k^N \int_{\lambda_0^N}^{\lambda_{2n-1}^N} \frac{\varphi_k^N(\mu) d\mu}{\sqrt[c]{\chi_N(\mu - i0)}}. \quad (1.11)$$

Here \mathfrak{p}_N is the trace of $L(b^N, a^N)$, $\mathfrak{p}_N = \text{tr} L(b^N, a^N)$, $\dot{\Delta}_N(\mu) = \frac{d}{d\mu} \Delta_N(\mu)$ and for any $0 < n < N$, $\varphi_n^N(\mu)$ is a polynomial of degree $N - 2$,

$$\varphi_n^N(\mu) = \prod_{\substack{0 < k < N \\ k \neq n}} (\mu - \sigma_k^{N,n}),$$

whose $N - 2$ zeroes $(\sigma_k^{N,n})_{0 < k \neq n < N}$ are uniquely determined by the $N - 2$ normalization conditions

$$\frac{1}{2\pi} \int_{\Gamma_k^N} \frac{\varphi_n^N(\mu)}{\sqrt[c]{\chi_N(\mu)}} d\mu = 0 \quad \forall 0 < k \neq n < N. \quad (1.12)$$

Here $(\Gamma_k^N)_{0 < k < N}$, denote pairwise disjoint counterclockwise oriented contours in \mathbb{C} , chosen in such a way that $\lambda_{2k-1}^N, \lambda_{2k}^N$ are inside of Γ_k^N whereas all other eigenvalues of $Q_N^{\alpha, \beta}$ are outside. (Actually, $(\frac{\varphi_k^N(\mu)}{\sqrt{\chi_N(\mu)}} d\mu)_{0 < k < N}$ are differential forms on the spectral curve defined by $\chi_N(\mu)$ – see Section 4 for more details.) Unlike for the corresponding formula for the KdV frequencies, (1.11) is a *system of equations* for the frequencies $(\omega_n^N)_{0 < n < N}$ rather than a formula as the Toda frequencies also appear on the right hand side of the identity (1.11). Fortunately, the two terms on the right hand side of (1.11) are not of the same order as $N \rightarrow \infty$. It turns out that the first one dominates the second one, suggesting to use a two step approach to prove the claimed asymptotics. In a first step (Section 6), we compute the leading order of the asymptotics of ω_n^N , using the following alternative representation of the Toda frequencies.

Theorem 1.4 *The frequencies ω_n^N , $0 < n < N$, of the states (b^N, a^N) defined by (1.2) are strictly positive, $\omega_n^N > 0$, and given by the expression*

$$\omega_n^N = \left(\frac{1}{2\pi} \int_{\Gamma_n^N} \frac{\varphi_n^N(\mu)}{\sqrt{\chi_N(\mu)}} d\mu \right)^{-1}. \quad (1.13)$$

We use Theorem 1.4 to prove the asymptotics (1.5) of the frequencies in the bulk in Section 6, Proposition 6.2. In a second step we establish the asymptotics (1.3), (1.4), and (1.6) of Theorem 1.1. The proof of these asymptotics is based on formula (1.11), uses the first order asymptotics mentioned above and relies on a comparison with corresponding formulas for the KdV frequencies (Section 4) as well as the asymptotics of the zeroes of the functions φ_n^N (Section 5) and the Toda actions I_n^N (Section 3). The formulas (1.11) of the Toda frequencies are reviewed in Section 4. There we also prove Theorem 1.4. The higher order asymptotics (1.6) are shown in Proposition 7.1. Using a symmetry of the Toda chain, discussed in Appendix E, the proof of (1.4) can be reduced to the one of (1.3). The latter asymptotics are proved in Section 7 using auxiliary estimates derived in Appendix D.

Related work: In the fifties, Fermi, Pasta, and Ulam introduced and studied a model consisting of a chain of particles interacting with their nearest neighbors through nonlinear strings, referred to as FPU chain. Expecting that the thermalization of small energy solution of such chains is valid they wanted to compute in numerical experiments the partition of energy among the normal modes of the linearized chain to see at what rate equipartition of energy is achieved. Much to their surprise, instead of equipartition of energy,

they observed recurrent features of these solutions. Their report [8] had a far reaching impact. To our days, the questions raised with their experiments have been intensely studied. However, the question if thermalization occurs has not been settled so far. It requires to compute delicate asymptotics for large number of particles. For status reports on the investigations of the FPU chains at the occasion of the 50th anniversary of the FPU experiments see the articles/books [4], [6], [11].

Formally, the dynamics of FPU chains of particles corresponding to long wave initial data can be approximated by solutions of the KdV equation at least for some interval of time – see e.g. [22], [20], [7]. But only recently this has been rigorously proved for certain classes of initial data ([21], [3]). More precisely, in [21] it was proved that for long waves on the real line with initial data of size ϵ^2 and sufficiently strong decay, the corresponding solutions of the FPU chain can be approximated up to a translation by the superposition of a right-going and a left-going wave, both evolving according to KdV, over a time interval of order ϵ^{-3} . Despite the fact that in the periodic set-up, waves interact much stronger, an analogous result for special type of analytic initial data could be proved in this set-up as well by [3].

In the case of the Toda chain, Toda himself formally addressed the problem of approximating solutions of the Toda chain by solutions of KdV with the aim of explaining the recurrent features of FPU chains at least in this integrable case – see [20]. Typically, in investigations of this kind, a limit of Toda chains of the type describe above is obtained by computing the formal limit of the Jacobi matrices L for a sequence of initial data, obtained by discretizing two given smooth functions, one for the diagonal and the other for the upper (and hence also lower) diagonal of L . In such a way one gets formally that L approaches a Hill operator which via the Lax pair formalism then leads to a solution of the KdV equation. The fundamental question then arose how a second Hill operator can be found providing the second solution of the KdV equation needed to describe the above mentioned approximation of long waves of periodic Toda chains.

To answer this question, we used as a key ingredient the discovery made in [5] that the Jacobi matrix describing the Toda chain with N particles can be viewed as the (geometric) quantization with Planck constant \hbar of a Hamiltonian system on the torus \mathbb{T}^2 , and the limit of such periodic Toda chains as the semiclassical limit of this quantization where the number N of particles of a chain coincides with the inverse of the mesh size of the discretization as well as with $\frac{1}{\hbar}$. This observation was used in a crucial way in paper [2] (cf

also [1]) where among other results we proved an asymptotic expansion in N^{-1} of the eigenvalues of the Jacobi matrices in terms of the eigenvalues of two Hill operators with potentials q_+ respectively q_- , introduced above.

Acknowledgments: The authors would like to thank the CNRS, the École polytechnique, the Swiss National Science Foundation, the University of Milan and the University of Zürich for financial support during the elaboration of this work.

2 Spectral asymptotics

For the convenience of the reader, we recall in this section results established in [2] which will be used throughout in the sequel. The first result concerns the asymptotics of the eigenvalues

$$\lambda_0^N < \lambda_1^N \leq \lambda_2^N < \cdots < \lambda_{2N-3}^N \leq \lambda_{2N-2}^N < \lambda_{2N-1}^N$$

of the Jacobi matrices

$$Q_N^{\alpha,\beta} \equiv Q(b^N, a^N) = L((b^N, b^N), (a^N, a^N)) \quad (2.1)$$

where (b^N, a^N) are the states introduced in (1.2),

$$b_n^N = \frac{1}{4N^2} \beta\left(\frac{n}{N}\right), \quad a_n^N = 1 + \frac{1}{4N^2} \alpha\left(\frac{n}{N}\right) \quad \text{with} \quad \beta, \alpha \in C_0^2(\mathbb{T}).$$

Explicitly, the symmetric $2N \times 2N$ matrix $Q_N^{\alpha,\beta}$ is given by

$$Q_N^{\alpha,\beta} = \begin{pmatrix} b_1^N & a_1^N & 0 & \cdots & \cdots & \cdots & \cdots & \cdots & 0 & a_N^N \\ a_1^N & b_2^N & a_2^N & 0 & \cdots & \cdots & \cdots & \cdots & \cdots & 0 \\ 0 & a_2^N & b_3^N & a_3^N & 0 & \cdots & \cdots & \cdots & \cdots & 0 \\ \vdots & \vdots \\ 0 & \cdots & 0 & a_{N-1}^N & b_N^N & a_N^N & 0 & \cdots & \cdots & 0 \\ 0 & \cdots & \cdots & 0 & a_N^N & b_1^N & a_1^N & 0 & \cdots & 0 \\ \vdots & \vdots \\ 0 & \cdots & \cdots & \cdots & \cdots & \cdots & 0 & a_{N-2}^N & b_{N-1}^N & a_{N-1}^N \\ a_N^N & 0 & \cdots & \cdots & \cdots & \cdots & \cdots & 0 & a_{N-1}^N & b_N^N \end{pmatrix}$$

Using Floquet theory (cf. e.g. [12]) one sees that the eigenvalues of $L(b^N, a^N)$ can be identified among the ones of $Q(b^N, a^N)$ as follows: if N is even they are

$$\lambda_0^N < \lambda_3^N \leq \lambda_4^N < \cdots < \lambda_{2N-5}^N \leq \lambda_{2N-4}^N < \lambda_{2N-1}^N \quad (2.2)$$

and hence

$$\Delta_N(\mu) - 2 = \mathbf{q}_N^{-1}(\mu - \lambda_0^N)(\mu - \lambda_{2N-1}^N) \prod_{0 < k < N/2} (\mu - \lambda_{4k}^N)(\mu - \lambda_{4k-1}^N)$$

whereas if N is odd

$$\lambda_1^N \leq \lambda_2^N < \lambda_5^N \leq \lambda_6^N < \cdots < \lambda_{2N-5}^N \leq \lambda_{2N-4}^N < \lambda_{2N-1}^N \quad (2.3)$$

leading to

$$\Delta_N(\mu) - 2 = \mathbf{q}_N^{-1}(\mu - \lambda_{2N-1}^N) \prod_{0 < k < N/2} (\mu - \lambda_{4k-2}^N)(\mu - \lambda_{4k-3}^N).$$

To describe the asymptotics of λ_n^N at the edges, $n \sim 1$ or $n \sim 2N - 1$, we introduced in [2] two Hill operators $H_\pm := -\partial_x^2 + q_\pm$ with potentials

$$q_\pm(x) = -2\alpha(2x) \mp \beta(2x). \quad (2.4)$$

The periodic eigenvalues $(\lambda_n^\pm)_{n \geq 0}$ of H_\pm on $[0, 1]$, when listed in increasing order and with multiplicities are known to satisfy

$$\lambda_0^\pm < \lambda_1^\pm \leq \lambda_2^\pm < \lambda_3^\pm \leq \lambda_4^\pm < \cdots .$$

As q_\pm have average zero, one has in addition that $\lambda_0^\pm < 0$.

Theorem 2.1 *Let $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$ satisfy (F) with $\eta < 1/2$ and set $M = [F(N)]$. Then for any α, β in $C_0^2(\mathbb{T})$, the asymptotics of the eigenvalues $(\lambda_n^N)_{0 \leq n \leq 2N-1}$ of $Q_N^{\alpha, \beta}$ are as follows:*

at the left and right edges: for $0 \leq n \leq 2M$

$$\lambda_n^N = -2 + \frac{1}{4N^2} \lambda_n^- + O\left(\frac{M^2}{N^3}\right), \quad \lambda_{2N-1-n}^N = 2 - \frac{1}{4N^2} \lambda_n^+ + O\left(\frac{M^2}{N^3}\right)$$

in the bulk: for $n = 2\ell, 2\ell - 1$ with $M < \ell < N - M$,

$$\lambda_n^N = -2 \cos \frac{\ell\pi}{N} + O\left(\frac{1}{MN^2}\right).$$

These estimates hold uniformly in $0 < n < N$ and uniformly on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

The second result which we recall from [2] concerns the asymptotics of the discriminant associated to the difference equation

$$a_{k-1}^N y(k-1) + b_k^N y(k) + a_k^N y(k+1) = \mu y(k) \quad \forall k \in \mathbb{Z}. \quad (2.5)$$

The discriminant is the trace of the Floquet operator of (2.5) and given by

$$\Delta_N(\mu) = y_1(N, \mu) + y_2(N+1, \mu)$$

where y_1^N and y_2^N are the fundamental solutions of (2.5) determined by

$$y_1(0, \mu) = 1, \quad y_1(1, \mu) = 0 \quad \text{and} \quad y_2(0, \mu) = 0, \quad y_2(1, \mu) = 1.$$

Analogously, the discriminant $\Delta_{\pm}(\lambda) \equiv \Delta(\lambda, q_{\pm})$ of

$$-y''(x, \lambda) + q_{\pm}(x)y(x, \lambda) = \lambda y(x, \lambda), \quad (2.6)$$

is defined as the trace of the Floquet operator associated to (2.6),

$$\Delta_{\pm}(\lambda) = y_1^{\pm}(1/2, \lambda) + (y_2^{\pm})'(1/2, \lambda)$$

where $y_1^{\pm}(x, \lambda)$ and $(y_2^{\pm})'(x, \lambda)$ are the fundamental solutions of (2.6) determined by

$$y_1^{\pm}(0, \lambda) = 1, \quad (y_1^{\pm}(0, \lambda))' = 0 \quad \text{and} \quad y_2^{\pm}(0, \lambda) = 0, \quad (y_2^{\pm}(0, \lambda))' = 1.$$

For $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$ as in Theorem 2.1 and $M = [F(N)]$ let

$$\Lambda_2^{\pm, M} := [\lambda_0^{\pm} - 2, \lambda_{2[F(M)]}^{\pm} + 2] + i[-2, 2]$$

and choose $N_0 \in \mathbb{Z}_{\geq 1}$ so that

$$\lambda_{2k+1}^{\pm} - \lambda_{2k}^{\pm} \geq 6 \quad \forall k \geq F(F(N_0)) \quad (2.7)$$

By the Counting Lemma for eigenvalues of Hill operators, N_0 can be chosen uniformly for subsets of function α, β in $C_0^2(\mathbb{T})$.

Theorem 2.2 *Let $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$ satisfy (F) with $\eta < 1/2$ and set $M = [F(N)]$. Then for any $\alpha, \beta \in C_0^2(\mathbb{T})$ and $N \geq N_0$, one has uniformly for λ in $\Lambda_2^{-, M}$,*

$$\Delta_N\left(-2 + \frac{1}{4N^2}\lambda\right) = (-1)^N \Delta_-(\lambda) + O\left(\frac{F(M)^2}{M}\right). \quad (2.8)$$

Similarly, uniformly for λ in $\Lambda_2^{+, M}$

$$\Delta_N\left(2 - \frac{1}{4N^2}\lambda\right) = \Delta_+(\lambda) + O\left(\frac{F(M)^2}{M}\right). \quad (2.9)$$

These estimates hold uniformly on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

The latter theorem has been applied in [2] to obtain asymptotics for the derivatives of the discriminant. In this paper we will need such asymptotic estimates for the first derivative, $\dot{\Delta}_N(\mu) := \frac{d}{d\mu}\Delta_N(\mu)$. Let

$$\Lambda_1^{\pm, M} := [\lambda_0^{\pm} - 1, \lambda_{2[F(M)]}^{\pm} + 1] + i[-1, 1].$$

Then, under the same assumptions as in Theorem 2.2 the following holds:

Corollary 2.3 *Uniformly for λ in $\Lambda_1^{-, M}$*

$$\frac{1}{4N^2}\dot{\Delta}_N\left(-2 + \frac{1}{4N^2}\lambda\right) = (-1)^N\dot{\Delta}_-(\lambda) + O\left(\frac{F(M)^2}{M}\right)$$

and similarly, for λ in $\Lambda_1^{+, M}$

$$-\frac{1}{4N^2}\dot{\Delta}_N\left(2 - \frac{1}{4N^2}\lambda\right) = \dot{\Delta}_+(\lambda) + O\left(\frac{F(M)^2}{M}\right).$$

These estimates hold uniformly on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

For later reference we record that in view of (1.9) and (2.2) - (2.3), $\dot{\Delta}_N(\mu)$ admits the product representation

$$\dot{\Delta}_N(\mu) = Nq_N^{-1} \cdot \prod_{0 < k < N} (\mu - \dot{\lambda}_k^N) \quad (2.10)$$

where the zeros $(\dot{\lambda}_n^N)_{0 < n < N}$ of $\dot{\Delta}_N(\mu)$ are listed in increasing order and satisfy

$$\lambda_{2n-1}^N \leq \dot{\lambda}_n^N \leq \lambda_{2n}^N \quad \forall 0 < n < N.$$

Remark 2.4 *As mentioned in the introduction, instead of considering the states (b^N, a^N) , alternatively, one could consider (d^N, c^N) , obtained by expressing (p^N, q^N) in Flaschka variables. Recall that the components of (p^N, q^N) were defined in Section 1 by*

$$p_n^N = -\frac{1}{4N^2}\beta\left(\frac{n}{N}\right) \quad \text{and} \quad q_n^N = -\frac{2}{4N}\xi\left(\frac{n}{N}\right)$$

where ξ is the element in $C_0^3(\mathbb{T})$, satisfying $\xi' = \alpha$. Note that $d^N = b^N$ whereas $c_n^N = a_n^N + O(N^{-3})$ for any $1 \leq n \leq N$. Indeed, as

$$q_n^N - q_{n+1}^N = \frac{1}{4N} \int_0^{N^{-1}} \alpha\left(\frac{n}{N} + t\right) dt \quad \text{and} \quad \alpha\left(\frac{n}{N} + t\right) = \alpha\left(\frac{n}{N}\right) + \int_0^t \alpha'\left(\frac{n}{N} + s\right) ds$$

one concludes that

$$c_n^N = \exp\left(\frac{q_n^N - q_{n+1}^N}{2}\right) = 1 + \frac{q_n^N - q_{n+1}^N}{2} + O(N^{-4}) = a_n^N + O(N^{-3})$$

where the error term is uniform for α in a bounded subset of the Banach space $C_0^2(\mathbb{T})$. Let $Q(b^N, a^N) = L((b^N, b^N), (a^N, a^N))$ and similarly, define $Q(d^N, c^N)$. The operator norm of the difference $Q(b^N, a^N) - Q(d^N, c^N)$ can thus be estimated by

$$\|Q(b^N, a^N) - Q(d^N, c^N)\| = \|Q(0_N, a^N - c^N)\| = O(N^{-3})$$

Therefore the spectrum of $Q(c^N, d^N)$ is N^{-3} -close to the one of $Q(b^N, a^N)$. Theorem 2.1 thus remains true if (b^N, a^N) is replaced by (d^N, c^N) . In view of Proposition 8.1 in [2], it also follows that Theorem 2.2 and Corollary 2.3 remain true for (d^N, c^N) . We leave it to the reader to go through the arguments used in the proofs of the results of the paper to verify that they remain valid if (b^N, a^N) is replaced by (d^N, c^N) .

3 Asymptotics of the actions

The aim of this section is to prove the asymptotics of the action variables for the states (b^N, a^N) stated in Theorem 1.2. Without further reference, we will use the notation introduced in the previous sections and assume that the assumptions of Theorem 1.2 hold. To prove it we first derive asymptotic estimates for I_n^N not involving KdV actions and valid for any $0 < n < N$. In the case where $\lambda_{2n-1}^N < \lambda_{2n}^N$, the action variable I_n^N is given by (cf [10], [12, Section 3])

$$I_n^N = \frac{1}{\pi} \int_{\lambda_{2n-1}^N}^{\lambda_{2n}^N} \frac{(\mu - \dot{\lambda}_n^N) \dot{\Delta}_N(\mu)}{c \sqrt{\Delta_N^2(\mu - i0) - 4}} d\mu \quad (3.1)$$

and is zero otherwise. Recall that $\dot{\lambda}_n^N$ denotes the n 'th zero of $\dot{\Delta}_N(\mu)$ (Section 2) and that the canonical root $\sqrt{c \Delta_N^2(\mu) - 4}$ was introduced in Section 1. The above formula for the actions leads to the following estimates.

Proposition 3.1 (i) *For any $1 \leq n \leq M$*

$$I_n^N = O\left(\frac{1}{n}((\gamma_n^-)^2 + \frac{M^4}{N^2})\frac{1}{N^2}\right), \quad I_{N-n}^N = O\left(\frac{1}{n}((\gamma_n^+)^2 + \frac{M^4}{N^2})\frac{1}{N^2}\right).$$

(ii) For any $M < n \leq \frac{N}{2}$

$$I_n^N, I_{N-n}^N = O\left(\frac{1}{nM^2} \frac{1}{N^2}\right).$$

The estimates hold uniformly in n and uniformly on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

Proof of Proposition 3.1: Clearly, we only need to consider $0 < n < N$ with $\lambda_{2n-1}^N < \lambda_{2n}^N$, as otherwise, $I_n^N = 0$. As by (2.10) and (1.9),

$$\dot{\Delta}_N(\mu) = N\mathfrak{q}_N^{-1} \cdot \prod_{0 < k < N} (\mu - \dot{\lambda}_k^N), \quad \Delta_N^2(\mu) - 4 = \mathfrak{q}_N^{-2} \prod_{0 \leq k \leq 2N-1} (\mu - \lambda_k^N)$$

one then gets by the above formula for the actions

$$I_n^N = \frac{N}{\pi} \int_{\lambda_{2n-1}^N}^{\lambda_{2n}^N} \frac{(\mu - \dot{\lambda}_n^N)^2}{\sqrt[+]{(\lambda_{2N-1}^N - \mu)(\mu - \lambda_0^N)}} \prod_{k \neq n} \frac{\mu - \dot{\lambda}_k^N}{w_k^N(\mu)} \cdot \frac{1}{\sqrt[+]{(\lambda_{2n}^N - \mu)(\mu - \lambda_{2n-1}^N)}} d\mu$$

where $w_k^N(\mu)$ denotes the standard root, $w_k^N(\mu) = \sqrt[+]{(\mu - \lambda_{2k}^N)(\mu - \lambda_{2k-1}^N)}$, defined on $\mathbb{C} \setminus [\lambda_{2k-1}^N, \lambda_{2k}^N]$ – see Section 5 for the determination of the sign. Using the identity $\int_{\lambda_{2n-1}^N}^{\lambda_{2n}^N} \frac{1}{\sqrt[+]{(\lambda_{2n}^N - \mu)(\mu - \lambda_{2n-1}^N)}} d\mu = \pi$ and the mean value theorem (cf Lemma C.6) one sees that there exists $\lambda_{2n-1}^N \leq \mu_* \leq \lambda_{2n}^N$ so that

$$I_n^N \leq \frac{N(\gamma_n^N)^2}{(\lambda_{2N-1}^N - \lambda_{2n}^N)^{1/2} (\lambda_{2n-1}^N - \lambda_0^N)^{1/2}} \prod_{k \neq n} \frac{\mu_* - \dot{\lambda}_k^N}{w_k^N(\mu_*)}.$$

Using that $1 + x \leq e^x$ for $x \geq 0$ one has, with $\mu \equiv \mu_*$,

$$\prod_{k < n} \frac{(\mu - \dot{\lambda}_k^N)^2}{w_k^N(\mu)^2} \leq \prod_{k < n} \frac{\mu - \lambda_{2k-1}^N}{\mu - \lambda_{2k}^N} = \prod_{k < n} \left(1 + \frac{\gamma_k^N}{\mu - \lambda_{2k}^N}\right) \leq \exp\left(\sum_{k < n} \frac{\gamma_k^N}{\lambda_{2n-1}^N - \lambda_{2k}^N}\right),$$

and similarly $\prod_{k > n} \frac{(\mu - \dot{\lambda}_k^N)^2}{w_k^N(\mu)^2} \leq \exp\left(\sum_{k > n} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N}\right)$. By Proposition A.1 it then follows that $\prod_{k \neq n} \frac{\mu_* - \dot{\lambda}_k^N}{w_k^N(\mu_*)} = O(1)$, yielding

$$I_n^N = O\left(\frac{N(\gamma_n^N)^2}{(\lambda_{2N-1}^N - \lambda_{2n}^N)^{1/2} (\lambda_{2n-1}^N - \lambda_0^N)^{1/2}}\right). \quad (3.2)$$

We now discuss the asymptotics of I_n^N , $N \rightarrow \infty$, in the four cases listed in the statement. To get the claimed estimate for I_n^N with $1 \leq n \leq M$ note that by Theorem 2.1, $(\gamma_n^N)^2 = O(\frac{1}{N^4}((\gamma_n^-)^2 + \frac{M^4}{N^2}))$ and

$$(\lambda_{2N-1}^N - \lambda_{2n}^N)^{-1/2} = O(1), \quad (\lambda_{2n-1}^N - \lambda_0^N)^{-1/2} = O\left(\frac{N}{n}\right) \quad (3.3)$$

yielding the claimed estimate $I_n^N = O(\frac{1}{N^4} \frac{1}{n} ((\gamma_n^-)^2 + \frac{M^4}{N^2}))$. The one for I_{N-n}^N is shown similarly. In the case $M < n \leq \frac{N}{2}$, Theorem 2.1 implies $(\gamma_n^N)^2 = O(\frac{1}{N^4 M^2})$ as well as $(\lambda_{2n-1}^N - \lambda_0^N)^{-1/2} = O(\frac{N}{n})$. Indeed, as by (C.2), $2(1 - \cos(\frac{n\pi}{N})) \geq \frac{\pi n^2}{N^2}$ one gets

$$\lambda_{2n-1}^N - \lambda_0^N = \frac{\pi n^2}{N^2} - \frac{\lambda_0^-}{8N^2} + O\left(\left(\frac{M^2}{N} + \frac{1}{M}\right) \frac{1}{N^2}\right)$$

yielding in view of the inequality $\lambda_0^- \leq 0$ (cf Section 2) the stated estimate. Combining the above estimates leads to $I_n^N = O(\frac{1}{nM^2} \frac{1}{N^2})$. In the case $\frac{N}{2} \leq n < N - M$ one argues in a similar fashion. Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

To prove the asymptotic estimates (1.7) and (1.8) of Theorem 1.2, we will use the following version of the above formulas of the actions (cf [10], [13])

$$I_n^N = \frac{1}{\pi} \int_{\lambda_{2n-1}^N}^{\lambda_{2n}^N} \operatorname{arcosh}\left((-1)^{N-n} \frac{\Delta_N(\mu)}{2}\right) d\mu \quad \forall 1 \leq n \leq N-1. \quad (3.4)$$

The action variables of KdV can be expressed in a similar way (cf [10])

$$I_n^\pm = \frac{2}{\pi} \int_{\lambda_{2n-1}^\pm}^{\lambda_{2n}^\pm} \operatorname{arcosh}\left((-1)^n \frac{\Delta_\pm(\lambda)}{2}\right) d\lambda \quad \forall n \geq 1. \quad (3.5)$$

In [16] one finds a detailed derivation of formula (3.5) for 1-periodic potentials. As indicated in [10], the formula remains valid for potentials of any given period.

Proof of Theorem 1.2: In view of Proposition 3.1 it remains to prove (1.7) and (1.8). As both estimates can be proven in a similar way, we concentrate on (1.7) only. So let $1 \leq n \leq M$. By the change of variable $\mu = -2 + \frac{1}{4N^2} \lambda$ in the integral in the above formula for I_n^N one gets

$$4N^2 I_n^N = \frac{1}{\pi} \int_{\nu_{2n-1}^N}^{\nu_{2n}^N} \operatorname{arcosh}\left(\frac{(-1)^{N-n}}{2} \Delta_N\left(-2 + \frac{\lambda}{4N^2}\right)\right) d\lambda \quad (3.6)$$

where for any $0 \leq j \leq 2N - 1$ we set $\nu_j^N = 4N^2(\lambda_j^N + 2)$. By Theorem 2.1 one has for any $0 \leq j \leq 2M$,

$$\nu_j^N = \lambda_j^- + O\left(\frac{M^2}{N}\right). \quad (3.7)$$

Let us consider first the case where $\gamma_n^- = \lambda_{2n}^- - \lambda_{2n-1}^- > 0$. Then with the change of variable

$$\lambda(x) = \nu_{2n-1}^N + \frac{4N^2\gamma_n^N}{\gamma_n^-} x, \quad 0 \leq x \leq \gamma_n^-$$

one gets in view of the identity (3.4)

$$8N^2 I_n^N = \frac{8N^2\gamma_n^N}{\gamma_n^-} I_n^- + \frac{8N^2\gamma_n^N}{\gamma_n^-} \frac{1}{\pi} \int_0^{\gamma_n^-} f(x) dx \quad (3.8)$$

where

$$f(x) := \operatorname{arcosh}\left(\frac{(-1)^{N-n}}{2} \Delta_N\left(-2 + \frac{\lambda(x)}{4N^2}\right)\right) - \operatorname{arcosh}\left(\frac{(-1)^n}{2} \Delta_-(\lambda_{2n-1}^- + x)\right).$$

Note that by Theorem 2.2 one has

$$\frac{(-1)^{N-n}}{2} \Delta_N\left(-2 + \frac{\lambda(x)}{4N^2}\right) = \frac{(-1)^n}{2} \Delta_-(\lambda(x)) + O\left(\frac{F(M)^2}{M}\right).$$

By (3.7), the first term on the right hand side of (3.8) can be estimated as

$$\frac{8N^2\gamma_n^N}{\gamma_n^-} I_n^- = I_n^- + O\left(\frac{I_n^-}{\gamma_n^-} \frac{M^2}{N}\right).$$

As $I_n^- = O\left(\frac{1}{n}(\gamma_n^-)^2\right)$ (cf [16, Theorem 7.3]) one concludes that

$$\frac{8N^2\gamma_n^N}{\gamma_n^-} I_n^- = I_n^- + O\left(\gamma_n^- \frac{M^2}{nN}\right). \quad (3.9)$$

The second term on the right hand side of (3.8) is estimated by using that arcosh is $\frac{1}{2}$ -Hölder continuous (cf Lemma C.1). Indeed, together with the fact that $\Delta_-(\lambda)$ is Lipschitz continuous and therefore

$$\Delta_-(\lambda(x)) - \Delta_-(\lambda_{2n-1}^- + x) = O(\lambda(x) - \lambda_{2n-1}^- - x) = O\left(\frac{M^2}{N}\right)$$

one gets

$$f(x) = O\left(\frac{F(M)}{M^{1/2}} + \frac{M}{N^{1/2}}\right).$$

Hence

$$\frac{8N^2\gamma_n^N}{\gamma_n^-} \frac{1}{\pi} \int_0^{\gamma_n^-} f(x)dx = O\left(\left(\gamma_n^- + \frac{M^2}{N}\right)\left(\frac{F(M)}{M^{1/2}} + \frac{M}{N^{1/2}}\right)\right). \quad (3.10)$$

Substituting (3.9) - (3.10) into (3.8) leads to

$$8N^2I_n^N = I_n^- + O\left(\left(\frac{M^2}{N} + \gamma_n^-\right)\left(\frac{F(M)}{M^{1/2}} + \frac{M}{N^{1/2}}\right)\right)$$

as claimed. Since the estimates are uniform with respect to the size of the gap length γ_n^- one can use an approximation argument to extend the result to the case where $\gamma_n^- = 0$. Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

To prove the asymptotics of the Toda frequencies stated in Theorem 1.1 we will need in addition the asymptotics of the quantities $J_n^N, 0 < n < N$, given by

$$J_n^N := \begin{cases} I_n^N/\gamma_n^N & \text{if } \gamma_n^N > 0 \\ 0 & \text{if } \gamma_n^N = 0 \end{cases}$$

in terms of $J_n^\pm, n \geq 1$, given by

$$J_n^\pm := \begin{cases} I_n^\pm/(2\gamma_n^\pm) & \text{if } \gamma_n^\pm > 0 \\ 0 & \text{if } \gamma_n^\pm = 0 \end{cases}.$$

Theorem 3.2 *Uniformly in $1 \leq n \leq F(M)$ and uniformly on bounded subsets of α, β in $C_0^2(\mathbb{T})$,*

$$J_n^N - J_n^-, \quad J_{N-n}^N - J_n^+ = O\left(\frac{M}{\sqrt{N}} + \frac{F(M)}{\sqrt{M}}\right).$$

Remark 3.3 *The proof of Theorem 3.2 will show that both $(J_n^N)_{1 \leq n \leq M}$ and $(J_n^-)_{1 \leq n \leq M}$ respectively $(J_{N-n}^N)_{1 \leq n \leq M}, (J_n^+)_{1 \leq n \leq M}$ are uniformly bounded with respect to n and with respect to bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.*

Remark 3.4 In [12] and [16] it is established that for any fixed N , $I_n^N = O((\gamma_n^N)^2)$ and $I_n^\pm = O((\gamma_n^\pm)^2)$ uniformly on closed bounded subsets of vectors (b, a) in $\mathbb{R}^N \times \mathbb{R}_{>0}^N$ respectively of potentials in $L^2(\mathbb{T})$. Theorem 3.2 states that the estimates for $(J_n^N)_{1 \leq n \leq M}, (J_{N-n}^N)_{1 \leq n \leq M}$ are uniform in N .

Proof: The estimates for $J_n^N - J_n^-$ and $J_{N-n}^N - J_n^+$ are proven in a similar way and so we concentrate on the first one only. We argue similarly as in the proof of Theorem 3.2. Let $1 \leq n \leq F(M)$. First consider the case where $\gamma_n^- > 0$. With the change of variable $\lambda = \lambda_{2n-1}^- + t\gamma_n^-$, one gets

$$\begin{aligned} I_n^- &= \frac{2}{\pi} \int_{\lambda_{2n-1}^-}^{\lambda_{2n}^-} \operatorname{arcosh}\left((-1)^n \frac{\Delta_-(\lambda)}{2}\right) d\lambda \\ &= 2\gamma_n^- \frac{1}{\pi} \int_0^1 \operatorname{arcosh}\left(\frac{(-1)^n}{2} \Delta_-(\lambda_{2n-1}^- + t\gamma_n^-)\right) dt. \end{aligned}$$

Note that the assumption $\gamma_n^- > 0$ together with the asymptotics $\gamma_n^N = \frac{\gamma_n^-}{4N^2} + O\left(\frac{M^2}{N^3}\right)$ implies that for N sufficiently large one has $\gamma_n^N > 0$. Hence with the change of variable $\mu = \lambda_{2n-1}^N + t\gamma_n^N$ one gets

$$I_n^N = \gamma_n^N \frac{1}{\pi} \int_0^1 \operatorname{arcosh}\left(\frac{(-1)^{N-n}}{2} \Delta_N(\lambda_{2n-1}^N + t\gamma_n^N)\right) dt.$$

We then conclude that

$$\begin{aligned} \frac{I_n^N}{\gamma_n^N} - \frac{I_n^-}{2\gamma_n^-} &= \frac{1}{\pi} \int_0^1 \left(\operatorname{arcosh}\left(\frac{(-1)^{N+n}}{2} \Delta_N(\lambda_{2n-1}^N + t\gamma_n^N)\right) \right. \\ &\quad \left. - \operatorname{arcosh}\left(\frac{(-1)^n}{2} \Delta_-(\lambda_{2n-1}^- + t\gamma_n^-)\right) \right) dt \end{aligned}$$

and as arcosh is $\frac{1}{2}$ -Hölder continuous (cf Lemma C.1) one then has

$$\frac{I_n^N}{\gamma_n^N} - \frac{I_n^-}{2\gamma_n^-} = O\left(\sup_{0 \leq t \leq 1} \left| (-1)^N \Delta_N(\lambda_{2n-1}^N + t\gamma_n^N) - \Delta_-(\lambda_{2n-1}^- + t\gamma_n^-) \right|^{1/2}\right).$$

As $\lambda_{2n-1}^N + t\gamma_n^N = -2 + (\lambda_{2n-1}^- + t\gamma_n^- + O(\frac{M^2}{N})) \frac{1}{4N^2}$ Theorem 2.2 implies that

$$\begin{aligned} &\left| (-1)^N \Delta_N(\lambda_{2n-1}^N + t\gamma_n^N) - \Delta_-(\lambda_{2n-1}^- + t\gamma_n^-) \right| \\ &= \left| \Delta_-(\lambda_{2n-1}^- + t\gamma_n^- + O(\frac{M^2}{N})) - \Delta_-(\lambda_{2n-1}^- + t\gamma_n^-) \right| + O\left(\frac{F(M)^2}{M} + \frac{M^2}{N}\right) \end{aligned}$$

and by the Lipschitz continuity of Δ_- it then follows that

$$\left| \frac{I_n^N}{\gamma_n^N} - \frac{I_n^-}{2\gamma_n^-} \right| = O\left(\frac{F(M)}{\sqrt{M}} + \frac{M}{\sqrt{N}}\right).$$

Since the estimates are uniform with respect to the size of the gap length γ_n^- one can use an approximation argument to extend the result to the case where $\gamma_n^- = 0$. Going through the arguments of the proof one verifies that the claimed uniformity holds. \square

4 Formulas for the frequencies

In this section we prove Theorem 1.4 and review formulas expressing the frequencies of the periodic Toda chain and the KdV equation as periods of certain differentials of cycles on the corresponding spectral curves.

Spectral curves. The Toda curve $\mathcal{C}_N \equiv \mathcal{C}_N^{(\alpha,\beta)}$ corresponding to the state (b^N, a^N) is defined as the affine curve

$$\mathcal{C}_N = \{(\mu, z) \in \mathbb{C}^2 : z^2 = \Delta_N^2(\mu) - 4\}.$$

In case the spectrum of $Q_N^{\alpha,\beta}$ is simple, it is a two sheeted open Riemann surface whose ramification points are the eigenvalues $(\lambda_n^N)_{0 \leq n \leq 2N-1}$ of $Q_N^{\alpha,\beta}$. If λ_{2n}^N is a double eigenvalue of $Q_N^{\alpha,\beta}$, then $(\lambda_{2n}^N, 0)$ is a singular point of \mathcal{C}_N . On \mathcal{C}_N we define a set of one forms $\frac{\psi_n^N(\mu)}{\sqrt{\Delta_N^2(\mu) - 4}} d\mu$ where $\psi_n^N(\mu)$ are polynomials of degree $N - 2$, uniquely determined by the normalisation conditions

$$\frac{1}{2\pi} \int_{\Gamma_k^N} \frac{\psi_n^N(\mu)}{\sqrt{\Delta_N^2(\mu) - 4}} d\mu = \delta_{nk} \quad \forall 1 \leq k \leq N - 1.$$

Here $\sqrt{\Delta_N^2(\mu) - 4}$ denotes the canonical root of

$$\Delta_N^2(\mu) - 4 = \mathfrak{q}_N^{-2} \chi_N(\mu), \quad \mathfrak{q}_N := \prod_{i=1}^N a_i^N, \quad (4.1)$$

and $(\Gamma_k^N)_{1 \leq k \leq N-1}$ the counterclockwise oriented contours introduced at the end of Section 1. In view of (4.1) it is convenient to write ψ_n^N in the form

$$\psi_n^N(\mu) = M_n^N \mathfrak{q}_N^{-1} \varphi_n^N(\mu), \quad \varphi_n^N(\mu) = \prod_{\substack{0 < k < N \\ k \neq n}} (\mu - \sigma_k^{N,n}) = (-1)^N \prod_{\substack{0 < k < N \\ k \neq n}} (\sigma_k^{N,n} - \mu).$$

Note that if the spectrum of $Q_N^{\alpha,\beta}$ is simple $(\frac{\psi_n^N(\mu)}{\sqrt{\Delta_N^2(\mu)-4}}d\mu)_{0 < n < N}$ are linearly independent holomorphic differentials on the Riemann surface \mathcal{C}_N . If λ_{2n}^N is a double eigenvalue, $\frac{\psi_n^N(\mu)}{\sqrt{\Delta_N^2(\mu)-4}}d\mu$ has a pole of order 1 at $(\lambda_{2n}^N, 0)$.

For the KdV equation, similar objects have been introduced. Denote by \mathcal{C}_\pm the curve corresponding to q_\pm ,

$$\mathcal{C}_\pm = \{(\lambda, w) \in \mathbb{C}^2 : w^2 = \Delta_\pm^2(\lambda) - 4\}.$$

If the spectrum of the Hill operator H_\pm is simple, \mathcal{C}_\pm is a two sheeted open Riemann surface of infinite genus whose ramification points are the eigenvalues $(\lambda_n^\pm)_{n \geq 0}$ of H_\pm . On \mathcal{C}_\pm we define a set of one forms $\frac{\psi_n^\pm(\lambda)}{\sqrt{\Delta_\pm^2(\lambda)-4}}d\lambda$. Note that $\Delta_\pm(\lambda)^2 - 4$ admits the product representation (cf [2, Appendix A])

$$\Delta_\pm^2(\lambda) - 4 = (\lambda_0^\pm - \lambda) \prod_{n \geq 1} \frac{(\lambda_{2n}^\pm - \lambda)(\lambda_{2n-1}^\pm - \lambda)}{(4\pi^2 n^2)^2}, \quad (4.2)$$

whereas $\psi_n^\pm(\lambda)$, $n \geq 1$, are entire functions of λ of the form

$$\psi_n^\pm(\lambda) = \frac{1}{2\pi n} \prod_{k \geq 1, k \neq n} \frac{\sigma_k^{\pm, n} - \lambda}{4\pi^2 k^2}$$

constructed in [16, Proposition D.10] for potentials of period 1 - see [2, Appendix A] for the adjustments needed in the case of potentials of period 1/2. Its zeros $\sigma_k^{\pm, n}$ are determined by the conditions

$$\int_{\Gamma_k^\pm} \frac{\psi_n^\pm(\lambda)}{\sqrt[3]{\Delta_\pm^2(\lambda) - 4}} d\lambda = 0 \quad \forall k \geq 1, k \neq n \quad (4.3)$$

and the normalisation factor $\frac{1}{2\pi n}$ by the requirement that

$$\int_{\Gamma_n^\pm} \frac{\psi_n^\pm(\lambda)}{\sqrt[3]{\Delta_\pm^2(\lambda) - 4}} d\lambda = 1. \quad (4.4)$$

Here $(\Gamma_k^\pm)_{k \geq 1}$ denote pairwise disjoint counterclockwise oriented contours chosen in such a way that $\lambda_{2k-1}^\pm, \lambda_{2k}^\pm$ are inside of Γ_k^\pm whereas all other eigenvalues are outside and $\sqrt[3]{\Delta_\pm^2(\lambda) - 4}$ is the canonical root of $\Delta_\pm(\lambda)^2 - 4$ determined by the sign condition $i\sqrt[3]{\Delta_\pm^2(\lambda) - 4} > 0$ for $\lambda_0^\pm < \lambda < \lambda_1^\pm$ or

$$\sqrt[3]{\Delta_\pm^2(\lambda + i0) - 4} > 0 \quad \forall \lambda < \lambda_0^\pm.$$

Toda frequencies. As a first step we prove Theorem 1.4.

Proof of Theorem 1.4. Let $0 < n < N$. Using a continuity argument, it suffices to prove the claimed identity in the case where $\lambda_{2n-1}^N < \lambda_{2n}^N$. By the definitions of ψ_n^N and φ_n^N one has

$$\frac{\psi_n^N(\mu)}{\sqrt{\Delta_N^2(\mu) - 4}} = M_n^N \frac{\varphi_n^N(\mu)}{\sqrt{\chi_N(\mu)}}, \quad (4.5)$$

implying that the normalisation factor M_n^N is given by

$$M_n^N \cdot \frac{1}{\pi} \int_{\lambda_{2n-1}^N}^{\lambda_{2n}^N} \frac{\varphi_n^N(\mu)}{\sqrt[{}^c]{\chi_N(\mu - i0)}} d\mu = 1. \quad (4.6)$$

To see that $M_n^N > 0$ note that $\lambda_{2k-1}^N \leq \sigma_k^{N,n} \leq \lambda_{2k}^N$, $0 < k < N$, and thus

$$(-1)^{N-1-n} \varphi_n^N(\mu) = \prod_{k < n} (\mu - \sigma_k^{N,n}) \cdot \prod_{k > n} (\sigma_k^{N,n} - \mu) > 0 \quad \forall \lambda_{2n-1}^N \leq \mu \leq \lambda_{2n}^N.$$

On the other hand, if $\lambda_{2n-1}^N < \lambda_{2n}^N$ one has by the definition of the c -root

$$\sqrt[{}^c]{\chi_N(\mu - i0)} = (-1)^{N+1-n} \sqrt[{}^c]{\chi_N(\mu)} \quad \forall \lambda_{2n-1}^N < \mu < \lambda_{2n}^N.$$

Altogether it follows that $M_n^N > 0$. Going through the arguments of [15], but without assuming that the trace of $L(b^N, a^N)$ vanishes, one concludes from [15, Lemma 4.4, formulas (28), (33), (38)] that $|\omega_n^N| = M_n^N$. In particular, it follows that $\omega_n^N \neq 0$. Taking into account that at the equilibrium $(0_N, 1_N)$, $\omega_n^N = 2 \sin(\frac{n\pi}{N}) > 0$, a deformation arguments yields the claimed identity $\omega_n^N = M_n^N$ and positivity $\omega_n^N > 0$. \square

Remark 4.1 Assume that $\lambda_{2n-1}^N < \lambda_{2n}^N$. Applying the mean value theorem (cf Lemma C.6) to (4.6) it follows that there is $\lambda_{2n-1}^N \leq \mu_* \leq \lambda_{2n}^N$ such that

$$\omega_n^N = \sqrt[{}^c]{(\lambda_{2N-1}^N - \mu_*)(\mu_* - \lambda_0^N)} \prod_{\substack{k \neq n \\ 1 \leq k < N}} \frac{\sqrt[{}^c]{(\lambda_{2k}^N - \mu_*)(\lambda_{2k-1}^N - \mu_*)}}{|\sigma_k^{N,n} - \mu_*|}. \quad (4.7)$$

By a continuity argument, (4.7) continues to hold in the case $\lambda_{2n-1}^N = \lambda_{2n}^N$ with $\mu_* = \tau_n^N$ where

$$\tau_n^N = (\lambda_{2n}^N + \lambda_{2n-1}^N)/2. \quad (4.8)$$

Remark 4.2 Actually, formula (4.7) is valid for any state $(b, a) \in \mathbb{R}^N \times \mathbb{R}_{>0}^N$. It allows to compute the frequencies corresponding to the equilibrium states

$$b = r1_N, \quad a = s1_N \quad \text{with} \quad 1_N = (1, \dots, 1), \quad r \in \mathbb{R}, \quad s > 0.$$

By [12, Lemma 2.6] one has $\lambda_0^N = -2s + r$, $\lambda_{2N-1}^N = 2s + r$ and

$$\lambda_{2n}^N = \lambda_{2n-1}^N = -2s \cos \frac{n\pi}{N} + r \quad \forall 0 < n < N.$$

Hence $\sigma_k^{N,n} = -2s \cos \frac{k\pi}{N} + r$ for any $1 \leq n, k \leq N-1$ and thus in this case, for any $0 < n < N$,

$$\omega_n^N = 2s \sqrt{\left(1 + \cos \frac{n\pi}{N}\right) \left(-\cos \frac{n\pi}{N} + 1\right)} = 2s \cdot \sin \frac{n\pi}{N}.$$

Formula (4.7) is suited for deriving asymptotic estimates of ω_n^N as $N \rightarrow \infty$ for n in the range $M < n < N-M$. However to get asymptotic estimates for the remaining frequencies we need alternative formulas, obtained with the help of Riemann bilinear relations. Such formulas were derived in [15], Section 4, for trace free Jacobi matrices. In addition it is proved there that they are independent of the trace. Hence these formulas hold in general modulo a spectral shift. It is however more convenient for us to have formulas which do not involve such a shift.

Proposition 4.3 The frequencies $(\omega_n^N)_{0 < n < N}$ of the state (b^N, a^N) satisfy

$$iN\omega_n^N = \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{(\mu - \frac{1}{N}\mathbf{p}_N) \dot{\Delta}_N(\mu) d\mu}{\sqrt[3]{\Delta_N^2(\mu) - 4}} - \sum_{k \in \mathcal{J}_N} I_k^N \omega_k^N \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\varphi_k^N(\mu) d\mu}{\sqrt[3]{\chi_N(\mu)}} \quad (4.9)$$

where $\mathcal{J}_N := \{1 \leq k \leq N-1 \mid I_k^N \neq 0\}$ and $\mathbf{p}_N = \sum_{i=1}^N b_i^N$.

Remark 4.4 In the formula (4.9), the quotients, $\frac{\dot{\Delta}_N(\mu)}{\sqrt[3]{\Delta_N^2(\mu) - 4}}$ and $\frac{\varphi_k^N(\mu)}{\sqrt[3]{\chi_N(\mu)}}$, $k \in \mathcal{J}_N$, are desingularized in the sense that factors which appear both in the numerator and in the denominator are canceled. Both denominators then consist of the square root of a polynomial with simple roots.

Proof. In the case where all the eigenvalues $(\lambda_j^N)_{0 \leq j \leq 2N-1}$ are simple such formulas have been computed in [15, Theorem 4.5]. Here we will use a slightly modified version, not involving a translation of the vector b^N . Going through the arguments [15, Section 4] – in particular Theorem 4.5 and Lemma 4.2 – one sees that if $(\lambda_j^N)_{0 \leq j \leq 2N-1}$ are simple, one has for any $1 < n < N$

$$iN\omega_n^N = \int_{\lambda_0^N}^{\lambda_{2n-1}^N} \frac{(\mu - \frac{1}{N}\mathbf{p}_N)\dot{\Delta}_N(\mu)d\mu}{\sqrt{c\Delta_N^2(\mu - i0) - 4}} - \sum_{k=1}^{N-1} I_k^N \omega_k^N \int_{\lambda_0^N}^{\lambda_{2n-1}^N} \frac{\varphi_k^N(\mu)d\mu}{\sqrt{c\chi_N(\mu - i0)}}. \quad (4.10)$$

Here we used Theorem 1.4 and formula (4.5) together with the fact that the one form

$$\frac{(\mu - \frac{1}{N}\mathbf{p}_N)\dot{\Delta}_N(\mu)}{\sqrt{\Delta_N^2(\mu) - 4}}d\mu$$

has an expansion in $\mu = \frac{1}{z}$ at ∞^\pm of the form $(\epsilon_\pm \frac{N}{z^2} + O(1))dz$ with $\epsilon_\pm \in \{+, -\}$ – see [15, Lemma 4.1 and Appendix A]. The integrals in (4.10) are split up into integrals over bands and gaps,

$$\int_{\lambda_0^N}^{\lambda_{2n-1}^N} = \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} + \sum_{j=1}^{n-1} \int_{\lambda_{2j-1}^N}^{\lambda_{2j}^N}. \quad (4.11)$$

Recall that $\int_{\lambda_{2j-1}^N}^{\lambda_{2j}^N} \frac{\omega_k^N \varphi_k^N(\mu)}{\sqrt{c\chi_N(\mu - i0)}}d\mu = \pi\delta_{jk}$ and, as $\int_{\lambda_{2j-1}^N}^{\lambda_{2j}^N} \frac{\dot{\Delta}_N(\mu)}{\sqrt{\Delta_N^2(\mu) - 4}}d\mu = 0$,

$$\int_{\lambda_{2j-1}^N}^{\lambda_{2j}^N} \frac{(\mu - \frac{1}{N}\mathbf{p}_N)\dot{\Delta}_N(\mu)}{\sqrt{c\Delta_N^2(\mu - i0) - 4}}d\mu = \pi I_j^N.$$

Hence the integrals in (4.11) over the gaps cancel and $iN\omega_n^N$ equals

$$\sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{(\mu - \frac{1}{N}\mathbf{p}_N)\dot{\Delta}_N(\mu)}{\sqrt{c\Delta_N^2(\mu) - 4}}d\mu - \sum_{k=1}^{N-1} I_k^N \omega_k^N \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\varphi_k^N(\mu)}{\sqrt{c\chi_N(\mu)}}d\mu.$$

In case some of the eigenvalues $(\lambda_j^N)_{0 \leq j \leq 2N-1}$ are double a continuity argument shows that [15, Theorem 4.5] continues to hold. Going through the arguments above and using that $I_k^N = O((\gamma_k^N)^2)$ by [12, Proposition 3.6], a limiting argument shows that the claimed formula for the frequencies holds also in this case. \square

KdV Frequencies. It turns out that there are formulas for the KdV frequencies which are of a very similar form as the ones of the Toda chains.

Lemma 4.5 *The KdV frequencies of a periodic potential of period T and mean 0 are $\omega_{T,n} = -\frac{24}{T}W_n$, $n \geq 1$, where*

$$W_n = \int_{\lambda_0}^{\lambda_{2n-1}} \frac{\lambda \dot{\Delta}(\lambda) d\lambda}{i \sqrt[3]{\Delta^2(\lambda - i0) - 4}} - \sum_{k \geq 1, I_k \neq 0} \frac{I_k}{2} \int_{\lambda_0}^{\lambda_{2n-1}} \frac{\psi_n(\lambda) d\lambda}{i \sqrt[3]{\Delta^2(\lambda - i0) - 4}} \quad (4.12)$$

where $(\lambda_n)_{n \geq 0}$ denote the periodic eigenvalues of $-\partial_x^2 + q$ on $[0, 2T]$, Δ the discriminant of $-\partial_x^2 + q$ on $[0, T]$, and I_k the action variables.

Proof: In the case $T = 1$, formula (4.12) has been proven in [16, Lemma 2.5]. For an arbitrary period we argue similarly, using that $\omega_{T,n} = K_T W_n$, $n \geq 1$ for some constant $K_T > 0$ which can be computed by following the arguments in [16] for the case $T = 1$: according to the computations at the end of [2, Appendix A], the zero potential, viewed as potential of period T has discriminant $\Delta(\lambda) = 2 \cos(\sqrt{\lambda}T)$. By the definition of the c-root one has for $\lambda \geq 0$, $\sqrt[3]{\Delta^2(\lambda) - 4} = -2i \sin(\sqrt[3]{\lambda}T)$ and the periodic eigenvalues are

$$\lambda_0 = 0, \quad \lambda_{2n} = \lambda_{2n-1} = \frac{n^2 \pi^2}{T^2} \quad \forall n \geq 1$$

and, for $\lambda \geq 0$, $\dot{\Delta}(\lambda) = -\frac{T}{\sqrt[3]{\lambda}} \sin \sqrt[3]{\lambda}T$. As $I_k = 0$ for any $k \geq 1$ it then follows that

$$W_n = \int_{\lambda_0}^{\lambda_{2n-1}} \frac{\lambda \dot{\Delta}(\lambda) d\lambda}{i \sqrt[3]{\Delta^2(\lambda - i0) - 4}} = -\frac{T}{3} \lambda^{\frac{3}{2}} \Big|_0^{(\frac{n\pi}{T})^2} = -\frac{T}{3} \left(\frac{n\pi}{T}\right)^3.$$

Arguing as in [16], the frequencies of the zero potential on $[0, T]$ can be computed as $\omega_{T,n} = (\frac{2n\pi}{T})^3$. Hence $\omega_{T,n} = -\frac{24}{T}W_n$ as claimed. \square

Applying Lemma 4.5 in the case $T = \frac{1}{2}$ and $q = q_{\pm}$ and arguing as in the proof of Proposition 4.3 leads to the following formulas.

Proposition 4.6 *The KdV frequencies ω_n^{\pm} , $n \geq 1$, of the potentials q_{\pm} of period $1/2$ and mean 0 are*

$$-48 \sum_{j=1}^n \int_{\lambda_{2j-2}^{\pm}}^{\lambda_{2j-1}^{\pm}} \frac{\lambda \dot{\Delta}_{\pm}(\lambda) d\lambda}{i \sqrt[3]{\Delta_{\pm}^2(\lambda) - 4}} + 24 \sum_{k=1, I_k^{\pm} \neq 0}^{\infty} I_k^{\pm} \sum_{j=1}^n \int_{\lambda_{2j-2}^{\pm}}^{\lambda_{2j-1}^{\pm}} \frac{\psi_n^{\pm}(\lambda) d\lambda}{i \sqrt[3]{\Delta_{\pm}^2(\lambda) - 4}} \quad (4.13)$$

5 Asymptotics of the zeroes of φ_k^N

In this section we compute the asymptotics of the zeroes $\sigma_\ell^{N,k}$ of the polynomials φ_ℓ^N , introduced in Section 1, in terms of the zeroes $\sigma_\ell^{\pm,k}$ of the entire functions ψ_k^\pm (cf Section 4). To define the edges of the spectrum of $Q_N^{\alpha,\beta}$, we use an arbitrary function $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$, satisfying (F) with $\eta \leq 1/3$. For the remainder of the paper we set $M = [F(N)]$ and $L = [F(M)]$. Then $L^3 \leq F(M)^3 \leq M$ and $M \leq \frac{N}{M^2}$. Furthermore, as $L \geq 1$, $M^3 \leq NL^2$ or $\frac{M}{L^2} \leq \frac{N}{M^2}$. We record for reference the following inequalities

$$\frac{M^2}{N} \leq L^{-3} \leq L^{-2}, \quad \frac{M^2}{N} \leq \frac{L^2}{M}, \quad \frac{M}{N^{1/2}} \leq \frac{L}{M^{1/2}} \quad (5.1)$$

Theorem 5.1 *Let $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$ satisfy (F) with $\eta \leq 1/3$. For $1 \leq k, \ell \leq L = [F(M)]$ with $\ell \neq k$*

$$(i) \quad \sigma_\ell^{N,k} = -2 + \frac{\sigma_\ell^{-,k}}{4N^2} + O\left(\frac{1}{N^2 L^{5/2}}\right) \quad (ii) \quad \sigma_{N-\ell}^{N,k} = 2 - \frac{\sigma_\ell^{+,k}}{4N^2} + O\left(\frac{1}{N^2 L^{5/2}}\right).$$

The error estimate is uniform in k, ℓ and uniform on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

To prove Theorem 5.1 we first need to introduce additional notation and derive auxiliary results. It turns out to be convenient to define the following roots: for $0 < \ell < N$ denote by $w_\ell^N(\mu)$ the standard root

$$w_\ell^N(\mu) = \sqrt[\zeta]{(\lambda_{2\ell}^N - \mu)(\lambda_{2\ell-1}^N - \mu)}, \quad \mu \in \mathbb{C} \setminus [\lambda_{2\ell-1}^N, \lambda_{2\ell}^N], \quad (5.2)$$

determined by the sign condition $w_\ell^N(\mu) < 0$ for any $\mu > \lambda_{2\ell}^N$. The standard root is related to the canonical root, introduced in (1.10) as follows

$$\sqrt[\zeta]{\Delta_N^2(\mu - i0) - 4} = \mathfrak{q}_N^{-1} \sqrt[\dagger]{\mu - \lambda_0^N} \sqrt[\dagger]{\lambda_{2N-1}^N - \mu} (-1)^N \prod_{0 < \ell < N} w_\ell^N(\lambda) \quad \forall \mu > \lambda_{2N-1}^N,$$

leading to the formula $\frac{\psi_k^N(\mu)}{\sqrt[\zeta]{\Delta_N^2(\mu - i0) - 4}} = \omega_k^N \frac{\varphi_k^N(\mu)}{\sqrt[\zeta]{\chi_N(\mu - i0)}}$ where

$$\frac{\varphi_k^N(\mu)}{\sqrt[\zeta]{\chi_N(\mu - i0)}} = \frac{1}{\sqrt[\dagger]{\mu - \lambda_0^N} \sqrt[\dagger]{\lambda_{2N-1}^N - \mu}} \frac{1}{w_k^N(\mu)} \prod_{\ell \neq k} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)}. \quad (5.3)$$

For later reference we record that

$$(-1)^N i \sqrt[{}^c]{\chi_N(\mu)} > 0 \quad \forall \mu \in (\lambda_0^N, \lambda_1^N). \quad (5.4)$$

Similarly, introduce the standard root

$$w_\ell^\pm(\lambda) := \sqrt[{}^s]{(\lambda_{2\ell}^\pm - \lambda)(\lambda_{2\ell-1}^\pm - \lambda)}, \quad \lambda \in \mathbb{C} \setminus [\lambda_{2\ell-1}^\pm, \lambda_{2\ell}^\pm], \quad (5.5)$$

with the sign condition $w_\ell^\pm(\lambda) < 0$ for any $\lambda > \lambda_{2\ell}^N$. The canonical root $\sqrt[{}^c]{\Delta_\pm^2(\lambda) - 4}$ is related to the standard roots by

$$\sqrt[{}^c]{\Delta_\pm^2(\lambda + i0) - 4} = \sqrt[{}^+]{\lambda_0^\pm - \lambda} \prod_{\ell \geq 1} \frac{w_\ell^\pm(\lambda)}{4\pi_\ell^2} \quad \forall \lambda < \lambda_0^\pm \quad (5.6)$$

where $\pi_\ell = \ell\pi$ for $\ell \neq 0$ and $\pi_0 = 1$, yielding for $\lambda < \lambda_0^\pm$

$$\frac{\psi_k^\pm(\lambda)}{\sqrt[{}^c]{\Delta_\pm^2(\lambda) - 4}} = \frac{2k\pi}{\sqrt[{}^+]{\lambda_0^\pm - \lambda}} \frac{1}{w_k^\pm(\lambda)} \prod_{\ell \neq k} \frac{\sigma_\ell^{\pm, k} - \lambda}{w_\ell^\pm(\lambda)}. \quad (5.7)$$

Again for later reference we record that

$$i \sqrt[{}^c]{\Delta_-^2(\lambda) - 4} > 0 \quad \forall \lambda \in (\lambda_0^-, \lambda_1^-). \quad (5.8)$$

As items (i) and (ii) are proved in the same way we concentrate on (i) only. For the sequel it turns out to be convenient to choose the contours Γ_n^- in (4.3) as follows. Choose $\rho > 0$ in such a way that

$$\lambda_{2\ell}^- + 3\rho < \lambda_{2\ell+1}^- - 3\rho \quad \forall \ell \geq 0$$

and define Γ_n^- to be the rectangle with top and bottom side given by $[\lambda_{2n-1}^- - 2\rho, \lambda_{2n}^- + 2\rho] \pm i\rho$. (Recall from [16, Proposition B.11] that the periodic eigenvalues of Hill operators are compact functions of the potential on $L^2(\mathbb{T})$. It then follows from the Counting Lemma that for any $k \geq 0$, the k 'th band length $\lambda_{2k+1}^- - \lambda_{2k}^-$ can be estimated from below – and hence $\rho > 0$ chosen – uniformly on bounded subsets of functions $\alpha, \beta \in C_0^2(\mathbb{T})$.) Furthermore, let

$$\tau_\ell^- = \frac{\lambda_{2\ell-1}^- + \lambda_{2\ell}^-}{2} \quad (5.9)$$

and denote by V the set of all sequences $s = (s_\ell)_{\ell \geq 1}$ in ℓ^2 such that for any $\ell \geq 1$, the sequence $\sigma_\ell = \tau_\ell^- + s_\ell$ satisfies $\lambda_{2\ell-1}^- - \rho \leq \sigma_\ell \leq \lambda_{2\ell}^- + \rho$. In particular the sequence σ_ℓ is strictly increasing. Note that V is a closed, connected subset of ℓ^2 . Furthermore, for any $k \geq 1$, let $\mathfrak{F}^k \equiv (\mathfrak{F}_n^k)_{n \geq 1}$ denote the map from ℓ^2 to ℓ^2 whose n 'th component is given by

$$\mathfrak{F}_n^k : s = (s_\ell)_{\ell \geq 1} \mapsto \begin{cases} \int_{\Gamma_n^-} \left(\prod_{\ell \neq k} \frac{\sigma_\ell - \lambda}{w_\ell^-(\lambda)} \right) \cdot \frac{\pi_k^2 - \pi_n^2}{w_k^-(\lambda)} \cdot \frac{2n\pi}{\sqrt{\lambda - \lambda_0^-}} d\lambda & \text{if } n \neq k \\ s_k & \text{if } n = k \end{cases}$$

In [16, Appendix D] it is shown that this map is real analytic and its differential at any point $s \in V$ a linear isomorphism (cf [16, Lemma D.6]). Furthermore, $s^{-,k} = (s_\ell^{-,k})_{\ell \geq 1}$, given by

$$s_\ell^{-,k} = \begin{cases} \sigma_\ell^{-,k} - \tau_\ell^- = 0 & \ell \neq k \\ 0 & \ell = k \end{cases} \quad (5.10)$$

is the unique solution of $\mathfrak{F}^k(s) = 0$ in V . Let $V^k := \bigcup_{N \geq 3} V^{k,N}$ where $(V^{k,N})_{N \geq 3}$ is a sequence of subsets of V with $V^{k,N}$ given by

$$\{s \in V : s_\ell = \sigma_\ell^{-,k} - \tau_\ell^- \ \forall \ell > L; \ -\frac{\gamma_\ell^-}{2} - C\frac{M^2}{N} \leq s_\ell \leq \frac{\gamma_\ell^-}{2} + C\frac{M^2}{N} \ \forall \ell \leq L\}$$

where $C > 0$ denotes the constant in the error estimate of Theorem 2.1 for the asymptotics of $\lambda_\ell^N, \ell \leq M$. Next we want to show that for any $k \geq 1$, the restriction of \mathfrak{F}^k to V^k is 1-1. To this end we first verify that V^k is relatively compact in ℓ^2 . To see it note that as $F(N) \leq N^\eta$ with $\eta \leq \frac{1}{3}$ one has $\sum_{\ell \leq F(M)} (\frac{M^2}{N})^2 \leq N^{\eta^2} \cdot N^{4\eta-2} = O(N^{-5/9})$ and hence V^k is bounded in ℓ^2 . To see that the sequences in V^k are uniformly summable, choose a strictly increasing sequence $(\ell_j)_{j \geq 1}$ in \mathbb{N} so that $\sum_{\ell > \ell_j} (\gamma_\ell^-)^2 \leq \frac{1}{j}$. By the first part of assumption (F), there then exists another strictly increasing sequence $(N_j)_{j \geq 1}$ in \mathbb{N} so that $F[F(N_j)] \leq \ell_j < F[F(N_j + 1)]$. Then, for any sequence $s = (s_\ell)_{\ell \geq 1} \in V^{k,N}$ with $N \leq N_j$ one has $\sum_{\ell > \ell_j} s_\ell^2 \leq \sum_{\ell > \ell_j} (\gamma_\ell^-)^2 \leq \frac{1}{j}$ whereas for sequences in $V^{k,N}$ with $N > N_j$ one has

$$\sum_{\ell > \ell_j} s_\ell^2 \leq \sum_{\ell > \ell_j} (\gamma_\ell^-)^2 + 2C^2 \sum_{\ell \leq F(M)} \frac{M^4}{N^2} \leq \frac{1}{j} + 2C^2 N_j^{\eta^2} \cdot N_j^{4\eta-2} \leq \frac{1}{j} + 2C^2 N_j^{-5/9}.$$

Altogether, we thus have proved that V^k is a relatively compact subset of V and hence, as V is closed, $\overline{V^k} \subset V$ is compact. Further note that V^k and

hence $\overline{V^k}$ is star shaped and thus connected. The following Lemma is the key ingredient of the proof of Theorem 5.1.

Lemma 5.2 *For any $k \geq 1$, the map $\mathfrak{F}^k|_{V^k} : V^k \rightarrow \mathfrak{F}^k(V^k)$ is one-to-one and together with its inverse uniformly Lipschitz.*

Proof. Let $\tilde{V}^k := \{s \in \overline{V^k} : \#((\mathfrak{F}^k)^{-1}(\mathfrak{F}^k(s)) \cap \overline{V^k}) = 1\}$. By the observation above, $(s_\ell^{-,k})_{\ell \geq 1}$ is in \tilde{V}^k and hence $\tilde{V}^k \neq \emptyset$. As \mathfrak{F}^k is a local diffeomorphism, \tilde{V}^k is closed. Furthermore by the compactness of $\overline{V^k}$ and the local diffeomorphism property of \mathfrak{F}^k one concludes that the complement of \tilde{V}^k in V^k , $\overline{V^k} \setminus \tilde{V}^k$, is also closed. As $\overline{V^k}$ is connected we thus have proved that $\overline{V^k} = \tilde{V}^k$, i.e. $\mathfrak{F}^k|_{\overline{V^k}}$ is 1 – 1. As $\overline{V^k}$ is compact one can argue as in the proof of [16, Proposition D.8] that $d_s \mathfrak{F}^k$ is uniformly boundedly invertible for any $s \in \overline{V^k}$ and any $k \geq 1$, implying that $\mathfrak{F}^k|_{\overline{V^k}}$ and its inverse are uniformly Lipschitz for any $k \geq 1$. \square

For any $1 \leq k \leq L$ introduce

$$\tilde{\sigma}_\ell^{N,k} = \begin{cases} 4N^2(\sigma_\ell^{N,k} + 2) & 1 \leq \ell \leq L, \ell \neq k \\ \tau_k^- & \ell = k \\ \sigma_\ell^{-,k} & \ell > L \end{cases},$$

and the corresponding shifted sequence $\tilde{s}^{N,k} = (\tilde{s}_\ell^{N,k} = \tilde{\sigma}_\ell^{N,k} - \tau_\ell^-)_{\ell \geq 1} \in \ell^2$.

Lemma 5.3 *There exists $N_0 \geq 3$ so that $\tilde{s}^{N,k} \in V^{k,N}$ for any $1 \leq k \leq L$ and $N \geq N_0$.*

Proof. By Theorem 2.1 that there exist $N_0 \geq 3$ and $C > 0$ so that for any $N \geq N_0$ and $1 \leq \ell \leq M$

$$\lambda_{2\ell-1}^- - C \frac{M^2}{N} \leq \nu_{2\ell-1}^N, \nu_{2\ell}^N \leq \lambda_{2\ell}^- + C \frac{M^2}{N} \quad \text{and} \quad C \frac{M^2}{N} < \rho$$

where $\nu_j^N = 4N^2(\lambda_j^{N,k} + 2)$. As $\lambda_{2\ell-1}^N \leq \sigma_\ell^{N,k} \leq \lambda_{2\ell}^N$ for $0 < k, \ell < N$ the definition of $\tilde{\sigma}_\ell^{N,k}$ implies that for any $1 \leq k \leq L$, $N \geq N_0$ one has for $1 \leq \ell \leq L$

$$\lambda_{2\ell-1}^- - C \frac{M^2}{N} \leq \tilde{\sigma}_\ell^{N,k} \leq \lambda_{2\ell}^- + C \frac{M^2}{N} \quad (5.11)$$

or, by subtracting τ_ℓ^-

$$-\gamma_\ell^-/2 - C\frac{M^2}{N} \leq \tilde{s}_\ell^{N,k} \leq \gamma_\ell^-/2 + C\frac{M^2}{N}$$

and for any $\ell \geq 1$, $\lambda_{2\ell-1}^- - \rho \leq \tilde{\sigma}_\ell^{N,k} \leq \lambda_{2\ell}^- + \rho$. Altogether we thus have shown that $\tilde{s}^{N,k} \in V^{k,N}$ for any $N \geq N_0$ and $1 \leq \ell \leq L$. \square

Lemma 5.4 *The first L elements of the sequence $\mathfrak{F}^k(\tilde{s}^{N,k})$ satisfy*

$$\left(\sum_{n=1}^L |\mathfrak{F}_n^k(\tilde{s}^{N,k})|^2 \right)^{1/2} = O\left(\frac{1}{L^{5/2}}\right) \quad (5.12)$$

uniformly for $1 \leq k \leq L$ and on bounded sets of potentials α, β in $C_0^2(\mathbb{T})$.

Proof. Recall that the $N-1$ zeroes $(\sigma_\ell^{N,k})_{0 < \ell < N}$ of φ_k^N are determined by the $N-1$ equations ($0 < n < N, n \neq k$)

$$0 = \int_{\Gamma_n^N} \left(\prod_{\substack{0 < \ell < N \\ \ell \neq k}} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} \right) \frac{\pi_k^2 - \pi_n^2}{4N^2 w_k^N(\mu)} \frac{2n\pi}{2N \sqrt[+]{\mu - \lambda_0^N}} \cdot \frac{2}{\sqrt[+]{\lambda_{2N-1}^N - \mu}} d\mu. \quad (5.13)$$

Introduce the products

$$Q_k^{N,L}(\mu) = \frac{2}{\sqrt[+]{\lambda_{2N-1}^N - \mu}} \cdot \prod_{L < \ell < N} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} \quad (5.14)$$

$$Q_k^{-,L}(\lambda) = \prod_{\ell > L} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} \quad (5.15)$$

$$\mathfrak{G}_k^{N,L}(\lambda) = \frac{Q_k^{-,L}(\lambda)}{Q_k^{N,L}(\mu)} \cdot \frac{2N \sqrt[+]{\mu - \lambda_0^N}}{\sqrt[+]{\lambda - \lambda_0^-}} \cdot \frac{4N^2 w_k^N(\mu)}{w_k^-(\lambda)} \quad (5.16)$$

where $\mu \equiv \mu(\lambda) = -2 + \frac{\lambda}{4N^2}$. Then

$$\begin{aligned} \mathfrak{F}_n^k(\tilde{s}^{N,k}) &= \int_{\Gamma_n^-} \left(\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\tilde{\sigma}_\ell^{N,k} - \lambda}{4N^2 w_\ell^N(\mu)} \right) \cdot Q_k^{N,L}(\mu) \frac{\pi_k^2 - \pi_n^2}{4N^2 w_k^N(\mu)} \frac{2n\pi}{2N \sqrt[+]{\mu - \lambda_0^N}} \\ &\quad \cdot \frac{2}{\sqrt[+]{\lambda_{2N-1}^N - \mu}} \mathfrak{G}_k^{N,L}(\lambda) \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{4N^2 w_\ell^N(\mu)}{w_\ell^-(\lambda)} d\lambda. \end{aligned}$$

The terms in the latter integral are estimated separately. By Theorem 2.1 (cf also Proposition A.1), one has for any $\lambda \in \Gamma_n^-$,

$$\begin{aligned} \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{4N^2 w_\ell^N(\mu)}{w_\ell^-(\lambda)} &= \frac{4N^2 w_n^N(\mu)}{w_n^-(\lambda)} \cdot \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k, n}} \frac{4N^2 w_\ell^N(\mu)}{w_\ell^-(\lambda)} \quad (5.17) \\ &= \left(1 + O\left(\frac{M^2}{N}\right)\right) \exp\left(\sum_{\substack{1 \leq \ell \leq L \\ \ell \neq k, n}} \log\left(1 + O\left(\frac{M^2}{N} \frac{1}{|\ell^2 - n^2|}\right)\right)\right) = 1 + O\left(\frac{M^2}{N}\right). \end{aligned}$$

By Lemma B.7 one then has $\mathfrak{G}_k^{N,L}(\lambda) \cdot \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{4N^2 w_\ell^N(\mu)}{w_\ell^-(\lambda)} = 1 + O\left(\frac{1}{L^3}\right)$ uniformly for $1 \leq n, k \leq L, n \neq k$ and $\lambda \in \Gamma_n^-$. By (5.13) it then follows that

$$\begin{aligned} \mathfrak{F}_N^k(\tilde{s}^{N,k}) &= \int_{\Gamma_n^-} \left(\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\tilde{\sigma}_\ell^{N,k} - \lambda}{4N^2 w_\ell^N(\mu)} \right) \cdot Q_k^{N,L}(\mu) \\ &\cdot \frac{\pi_k^2 - \pi_n^2}{4N^2 w_k^N(\mu)} \cdot \frac{2n\pi}{2N \sqrt[+]{\mu - \lambda_0^N}} \cdot \frac{2}{\sqrt[+]{\lambda_{2N-1}^N - \mu}} \cdot O\left(\frac{1}{L^3}\right) d\lambda. \end{aligned}$$

Furthermore, again by Theorem 2.1 one has uniformly for $1 \leq n, k \leq L$ and $\lambda \in \Gamma_n^-$

$$\frac{2n\pi}{2N \sqrt[+]{\mu - \lambda_0^N}}, \quad \frac{2}{\sqrt[+]{\lambda_{2N-1}^N - \mu}}, \quad \frac{\tilde{\sigma}_n^{N,k} - \lambda}{4N^2 w_n^N(\mu)} = O(1).$$

It remains to estimate the term $\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\tilde{\sigma}_\ell^{N,k} - \lambda}{4N^2 w_\ell^N(\mu)}$. By Theorem 2.1 (cf also Proposition A.1), $\tilde{\sigma}_\ell^{N,k} - \nu_j^N = O\left(\gamma_\ell^- + \frac{M^2}{N}\right)$ for $j \in \{2\ell, 2\ell - 1\}$ and $1 \leq \ell \leq L$ where $\nu_j^N = 4N^2(\lambda_j^{N,k} + 2)$. Hence, for $\lambda \in \Gamma_n^-$ and $1 \leq \ell \leq L$ with $\ell \neq k, n$

$$\frac{\tilde{\sigma}_\ell^{N,k} - \lambda}{4N^2 w_\ell^N(\mu)} = \left(\frac{(\tilde{\sigma}_\ell^{N,k} - \lambda)(\tilde{\sigma}_\ell^{N,k} - \lambda)}{(\nu_{2\ell}^N - \lambda)(\nu_{2\ell-1}^N - \lambda)} \right)^{1/2} = 1 + O\left(\left(\gamma_\ell^- + \frac{M^2}{N}\right) \frac{1}{|\ell^2 - n^2|}\right)$$

and thus

$$\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k, n}} \frac{\lambda - \tilde{\sigma}_\ell^{N,k}}{4N^2 w_\ell^N(\mu)} = \exp\left(\sum_{\substack{1 \leq \ell \leq L \\ \ell \neq k, n}} \log\left(1 + O\left(\left(\gamma_\ell^- + \frac{M^2}{N}\right) \frac{1}{|\ell^2 - n^2|}\right)\right)\right) = O(1).$$

In view of (5.1) we thus proved that for any $1 \leq n \leq L, n \neq k$, $\mathfrak{F}_n^k(\tilde{s}^{N,k}) = O\left(\frac{1}{L^3}\right)$. Going through the arguments of the proof one verifies the claimed uniformity statement. \square

Lemma 5.5 *The tail $(\mathfrak{F}_n^k(\tilde{s}^{N,k}))_{n>L}$ of $\mathfrak{F}^k(\tilde{s}^{N,k})$ satisfies the estimate*

$$\left(\sum_{n>L} |\mathfrak{F}_n^k(\tilde{s}^{N,k})|^2\right)^{1/2} = O\left(\frac{1}{L^4}\right) \quad (5.18)$$

uniformly in $1 \leq k \leq L$ and on bounded sets of potentials α, β in $C_0^2(\mathbb{T})$.

Proof. By the definition of $\tilde{s}^{N,k}$ and \mathfrak{F}_n^k one has

$$\mathfrak{F}_n^k(\tilde{s}^{N,k}) = \int_{\Gamma_n^-} \left(\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\tilde{\sigma}_\ell^{N,k} - \lambda}{\sigma_\ell^{-,k} - \lambda} \right) \zeta_n^{-,k}(\lambda) \frac{\sigma_n^{-,k} - \lambda}{w_n^-(\lambda)} d\lambda$$

where

$$\zeta_n^{-,k}(\lambda) := \left(\prod_{\substack{\ell \geq 1 \\ \ell \neq k, n}} \frac{\sigma_\ell^{-,k} - \lambda}{w_n^-(\lambda)} \right) \frac{\pi_k^2 - \pi_n^2}{w_k^-(\lambda)} \frac{2n\pi}{\sqrt{\lambda - \lambda_0^-}}.$$

We argue as in (5.17) and write for $\lambda \in \Gamma_n^-$ with $n > L$

$$\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\tilde{\sigma}_\ell^{N,k} - \lambda}{\sigma_\ell^{-,k} - \lambda} = \exp \left(\sum_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \log \left(1 + \frac{\tilde{\sigma}_\ell^{N,k} - \sigma_\ell^{-,k}}{\sigma_\ell^{-,k} - \lambda} \right) \right).$$

Note that by (5.11), $|\sigma_\ell^{N,k} - \tau_\ell^-| \leq \gamma_\ell^-/2 + C\frac{M^2}{N}$. Taking into account that $|\sigma_\ell^{-,k} - \tau_\ell^-| \leq \gamma_\ell^-/2$ it then follows that

$$|\tilde{\sigma}_\ell^{N,k} - \sigma_\ell^{-,k}| \leq \gamma_\ell^- + C\frac{M^2}{N} = O\left(\frac{1}{\ell^2} + \frac{M^2}{N}\right).$$

As $(\sigma_\ell^{-,k} - \lambda)^{-1} = O\left(\frac{1}{L^2}\right)$ for $\ell \leq L/2$ and $n > L$, one concludes from (5.1)

$$\sum_{\substack{1 \leq \ell \leq L/2 \\ \ell \neq k}} \left| \frac{\tilde{\sigma}_\ell^{N,k} - \sigma_\ell^{-,k}}{\sigma_\ell^{-,k} - \lambda} \right| = O\left(\frac{1}{L^2} \left(\sum_{\ell \geq 1} \frac{1}{\ell^2} + \frac{M^2}{N} L \right)\right) = O\left(\frac{1}{L^2}\right).$$

Furthermore

$$\sum_{\substack{L/2 < \ell \leq L \\ \ell \neq k}} \left| \frac{\tilde{\sigma}_\ell^{N,k} - \sigma_\ell^{-,k}}{\sigma_\ell^{-,k} - \lambda} \right| = O\left(\frac{\log L}{L^3}\right).$$

Altogether we thus have proved that for $n > L$

$$\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\tilde{\sigma}_\ell^{N,k} - \lambda}{\sigma_\ell^{-,k} - \lambda} - 1 = O\left(\frac{1}{L^2}\right).$$

Actually, these asymptotic estimates hold not only for $\lambda \in \Gamma_n^-$, but also for any λ in the interior of Γ_n^- , uniformly in $1 \leq k \leq L$ and $n > L$. Furthermore, one verifies in a straightforward way that $\zeta_n^{-,k}(\lambda) = O(1)$ for $\lambda \in \Gamma_n^-$ and in the interior of Γ_n^- . To prove the claimed estimate note that $\int_{\Gamma_n^-} \zeta_n^{-,k}(\lambda) \frac{\sigma_n^{-,k} - \lambda}{w_n^-(\lambda)} d\lambda = 0$ by the definition of $(\sigma_\ell^{-,k})_{\ell \neq k}$, implying that

$$\mathfrak{F}_n^k(\tilde{s}^{N,k}) = \int_{\Gamma_n^-} \zeta_n^{-,k}(\lambda) \left(\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\tilde{\sigma}_\ell^{N,k} - \lambda}{\sigma_\ell^{-,k} - \lambda} - 1 \right) \frac{\sigma_n^{-,k} - \lambda}{w_n^-(\lambda)} d\lambda.$$

In the case $\gamma_n^- = 0$, one has $\frac{\sigma_n^{-,k} - \lambda}{w_n^-(\lambda)} = 1$ and as the integrand is analytic in λ the latter integral vanishes by Cauchy's theorem. In the case $\gamma_n^- > 0$ we deform the contour Γ_n^- to the interval $[\lambda_{2n-1}^-, \lambda_{2n}^-]$ to conclude by the mean value theorem that

$$\mathfrak{F}_n^k(\tilde{s}^{N,k}) = O\left(\frac{1}{L^2} \gamma_n^-\right) = O\left(\frac{1}{L^4} n^2 \gamma_n^-\right).$$

The claimed estimate then follows from the assumption that $\alpha, \beta \in C_0^2(\mathbb{T})$. Going through the arguments of the proof one verifies the claimed uniformity statement. \square

Proof of Theorem 5.1. Combining (5.12) and (5.18) yields

$$\|\mathfrak{F}^k(\tilde{s}^{N,k})\|_{\ell^2} = O\left(\frac{1}{L^{5/2}}\right)$$

uniformly for $1 \leq k \leq L$. As for any $1 \leq k \leq L$, $\mathfrak{F}^k(\tilde{s}^{-,k}) = 0$ (cf (5.10)), $\tilde{s}^{N,k} \in V^k$ (Lemma 5.3), \mathfrak{F}^k is 1-1 and together with its inverse uniformly Lipschitz on V^k (Lemma 5.2) we then get the claimed estimate

$$\|\tilde{\sigma}^{N,k} - \sigma^{-,k}\|_{\ell^2} = \|\tilde{s}^{N,k} - s^{-,k}\|_{\ell^2} = O\left(\frac{1}{L^{5/2}}\right).$$

Going through the arguments of the proof one verifies the claimed uniformity statement. \square

6 Leading order asymptotics

In this section we compute the leading order asymptotics of the Toda frequencies ω_n^N . In particular we prove the asymptotics (1.5) in the bulk. In addition we identify the principal contributions in the formulas (4.9) to the asymptotics of ω_n^N at the two edges.

Remark 6.1 *We recall that for the remainder of the paper we set $M = [F(N)]$ and $L = [F(M)]$ where $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$ satisfies (F) with $\eta \leq 1/3$. Then (5.1) holds and the asymptotics of Theorem 1.2 and Theorem 3.2 at the right and left edges read as follows: for any $1 \leq n \leq L$*

$$8N^2 I_n^N = I_n^- + O\left(\left(\frac{M^2}{N} + \gamma_n^-\right) \frac{L}{M^{1/2}}\right) \quad 8N^2 I_{N-n}^N = I_n^+ + O\left(\left(\frac{M^2}{N} + \gamma_n^+\right) \frac{L}{M^{1/2}}\right)$$

$$J_n^N - J_n^-, \quad J_{N-n}^N - J_n^+ = O\left(\frac{L}{M^{1/2}}\right)$$

Proposition 6.2 *Uniformly for $M < n < N - M$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\omega_n^N = 2\left(\sin \frac{n\pi}{N}\right) \left(1 + O\left(\frac{\log M}{M^2}\right)\right).$$

Proof: By (4.7) there exists $\mu_* \in [\lambda_{2n-1}^N, \lambda_{2n}^N]$ so that

$$\omega_n^N = \sqrt[+]{(\lambda_{2N-1}^N - \mu_*)(\mu_* - \lambda_0^N)} \prod_{\substack{k \neq n \\ 1 \leq k < N}} \frac{\sqrt[+]{(\lambda_{2k}^N - \mu_*)(\lambda_{2k-1}^N - \mu_*)}}{|\sigma_k^{N,n} - \mu_*|}$$

By Theorem 2.1, $\mu_* = -2 \cos \frac{n\pi}{N} + O\left(\frac{1}{N^2 M}\right)$, $\lambda_0^N = -2 + O(N^{-2})$, and $\lambda_{2N-1}^N = 2 + O(N^{-2})$. Hence

$$(\lambda_{2N-1}^N - \mu_*)(\mu_* - \lambda_0^N) = 4 \sin^2 \frac{n\pi}{N} + O\left(\frac{1}{N^2}\right).$$

As $\sin \frac{n\pi}{N} \geq \sin \frac{M\pi}{N} = O\left(\frac{M}{N}\right)$ for $M < n < N - M$,

$$+\sqrt{(\lambda_{2N-1}^N - \mu_*)(\mu_* - \lambda_0^N)} = (2 \sin \frac{n\pi}{N})(1 + O(\frac{1}{M^2})). \quad (6.1)$$

Next we write

$$\prod_{\substack{k \neq n \\ 0 < k < N}} \frac{\sqrt{(\lambda_{2k}^N - \mu_*)(\lambda_{2k-1}^N - \mu_*)}}{|\sigma_k^{N,n} - \mu_*|} = \left(\prod_{\substack{k \neq n \\ 0 < k < N}} \frac{|\lambda_{2k}^N - \mu_*|}{|\sigma_k^{N,n} - \mu_*|} \prod_{\substack{k \neq n \\ 0 < k < N}} \frac{|\lambda_{2k-1}^N - \mu_*|}{|\sigma_k^{N,n} - \mu_*|} \right)^{1/2}.$$

The two products are estimated in the same way, so we concentrate on the first one. Note that $0 < \frac{\lambda_{2k}^N - \mu_*}{\sigma_k^{N,n} - \mu_*} = 1 + \frac{\lambda_{2k}^N - \sigma_k^{N,n}}{\sigma_k^{N,n} - \mu_*}$ for any $k \neq n$. It will turn out that $\frac{\lambda_{2k}^N - \sigma_k^{N,n}}{\sigma_k^{N,n} - \mu_*}$ is very small so that we can write

$$\prod_{\substack{k \neq n \\ 0 < k < N}} \frac{\lambda_{2k}^N - \mu_*}{\sigma_k^{N,n} - \mu_*} = \exp \left(\sum_{\substack{k \neq n \\ 0 < k < N}} \log \left(1 + \frac{\lambda_{2k}^N - \sigma_k^{N,n}}{\sigma_k^{N,n} - \mu_*} \right) \right) \quad (6.2)$$

and get the estimate

$$\exp \left(\sum_{\substack{k \neq n \\ 0 < k < N}} \log \left(1 + \frac{\lambda_{2k}^N - \sigma_k^{N,n}}{\sigma_k^{N,n} - \mu_*} \right) \right) - 1 = O \left(\sum_{\substack{k \neq n \\ 0 < k < N}} \frac{|\lambda_{2k}^N - \sigma_k^{N,n}|}{|\sigma_k^{N,n} - \mu_*|} \right).$$

The latter sum is split up and as $\lambda_{2k-1}^N \leq \sigma_k^{N,n} \leq \lambda_{2k}^N$ can be estimated as follows

$$\sum_{\substack{k \neq n \\ 0 < k < N}} \frac{|\lambda_{2k}^N - \sigma_k^{N,n}|}{|\sigma_k^{N,n} - \mu_*|} \leq \sum_{0 < k < n} \frac{\gamma_k^N}{\lambda_{2n-1}^N - \lambda_{2k}^N} + \sum_{n < k < N} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N}.$$

Taking into account that $\eta \leq 1/3$ and hence $\frac{M^3}{N} = o(1)$, the right hand side of the latter inequality is $O(\frac{\log M}{M^2})$ by Proposition A.1 (ii). Altogether we thus have proved that

$$\prod_{\substack{k \neq n \\ 0 < k < N}} \frac{\sqrt{(\lambda_{2k}^N - \mu_*)(\lambda_{2k-1}^N - \mu_*)}}{|\sigma_k^{N,n} - \mu_*|} = 1 + O\left(\frac{\log M}{M^2}\right).$$

Combined with (6.1), we then obtain the stated estimate. Going through the arguments of the proof, one verifies the claimed uniformity statement. \square

Proposition 6.3 *Uniformly for $1 \leq n \leq M$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$ one has (i) $\omega_n^N = O\left(\frac{n}{N}\right)$ and (ii) $\omega_{N-n}^N = O\left(\frac{n}{N}\right)$.*

Proof: The estimates (i) and (ii) are proved in the same way and so we concentrate on (i) only. We use again formula (4.7). As by Theorem 2.1

$$\lambda_{2N-1}^N - \mu_* \leq 4 - \frac{\lambda_0^+ + \lambda_{2n-1}^-}{4N^2} + O\left(\frac{M^2}{N^3}\right), \quad \mu_* - \lambda_0^N \leq \frac{\lambda_{2n}^- - \lambda_0^-}{4N^2} + O\left(\frac{M^2}{N^3}\right)$$

one gets

$$\sqrt[+]{(\lambda_{2N-1}^N - \mu_*)(\mu_* - \lambda_0^N)} = O\left(\frac{\sqrt{\lambda_{2n}^- - \lambda_0^-}}{N}\right) = O\left(\frac{n}{N}\right). \quad (6.3)$$

Next we need to estimate the product

$$\prod_{k \neq n} \frac{\sqrt[+]{(\lambda_{2k}^N - \mu_*)(\lambda_{2k-1}^N - \mu_*)}}{|\sigma_k^{N,n} - \mu_*|} = \left(\prod_{\substack{k \neq n \\ 1 \leq k < N}} \frac{|\lambda_{2k}^N - \mu_*|}{|\sigma_k^{N,n} - \mu_*|} \right)^{1/2} \left(\prod_{\substack{k \neq n \\ 1 \leq k < N}} \frac{|\lambda_{2k-1}^N - \mu_*|}{|\sigma_k^{N,n} - \mu_*|} \right)^{1/2}.$$

The two products in the latter expression are estimated in the same way and thus we concentrate again on the first one. Argue as in the proof of Proposition 6.2 and use that by Proposition A.1(i)

$$\sum_{\substack{k \neq n \\ 0 < k < N}} \frac{|\lambda_{2k}^N - \sigma_k^{N,n}|}{|\sigma_k^{N,n} - \mu_*|} \leq \sum_{0 < k < n} \frac{\gamma_k^N}{\lambda_{2n-1}^N - \lambda_{2k}^N} + \sum_{n < k < N} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O(1)$$

to conclude that $\left(\prod_{\substack{k \neq n \\ 1 \leq k < N}} \frac{|\lambda_{2k}^N - \mu_*|}{|\sigma_k^{N,n} - \mu_*|} \right)^{1/2} = O(1)$. Substituting the obtained estimates into formula (4.7) yields (i). Going through the arguments of the proof one verifies the claimed uniformity. \square

We now use the rough estimates of the Toda frequencies ω_n^N obtained above to identify the principal contributions in the formulas (4.9),

$$iN\omega_n^N = \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{(\mu - \frac{1}{N}\mathbf{p}_N)\dot{\Delta}_N(\mu)d\mu}{\sqrt[+]{\Delta_N(\mu)^2 - 4}} - \sum_{k \in \mathcal{J}_N} I_k^N \omega_k^N \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\varphi_k^N(\mu)d\mu}{\sqrt[+]{\chi_N(\mu)}},$$

to their asymptotics at the two edges. Recall that $\mathbf{p}_N = \sum_{i=1}^N b_i^N$ and χ_N is the characteristic polynomial, defined by (4.1).

Proposition 6.4 *Uniformly for $1 \leq n \leq M$, $n \leq M_0 \leq M$, and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\omega_n^N = \frac{2\pi n}{N} - \frac{\pi}{4N^3} \sum_{j=1}^n \lambda_{2j-2}^- - (P1) - (P2) - (P3) + \text{Error} \quad (6.4)$$

where

$$(P1) := \frac{1}{4N^3} \sum_{j=1}^n \int_{\nu_{2j-2}^N}^{\nu_{2j-1}^N} \arccos\left(\frac{(-1)^{N+j}}{2} \Delta_N\left(-2 + \frac{\lambda}{4N^2}\right)\right) d\lambda$$

$$(P2) := \frac{1}{N} \sum_{\substack{1 \leq k \leq M_0 \\ I_k^N \neq 0}} I_k^N \omega_k^N \int_{\mathcal{I}_n^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu, \quad (P3) := \frac{1}{N} \sum_{M_0 < k \leq M} I_k^N \omega_k^N \int_{\mathcal{I}_n^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu$$

with $\nu_j^N = 4N^2(2 + \lambda_j^N)$, $\mathcal{I}_n^N := [\lambda_0^N, \lambda_{2n-1}^N] \setminus (\cup_{j=1}^{n-1} [\lambda_{2j-1}^N, \lambda_{2j}^N])$, and

$$\text{Error} = O\left(\frac{1}{N^3} \left(\frac{nM^2}{N} + \frac{n}{M^3}\right)\right) \quad (6.5)$$

Furthermore,

$$(P3) = O\left(\frac{\log(n+1)}{N^3} \left(\frac{1}{M_0^5} + \frac{M^4 \log M}{N^2}\right)\right). \quad (6.6)$$

Remark 6.5 *The integer M_0 in the statement of Proposition 6.4 is a free parameter for which different choices will be made in the course of the proof of Theorem 1.1: in the proof of Proposition 7.1, Proposition 6.4 will be applied for $1 \leq n \leq M$ with $M_0 = M$ whereas in the proof of Theorem 7.4 we apply it for $1 \leq n \leq L$ with $M_0 = L$.*

Proof. For any $1 \leq n \leq M$, write the formula (4.9) for the n th frequency, recalled above, as a sum $\omega_n^N = (T1) + (T2) - (T3) - (T4)$ where

$$(T1) := -\frac{2}{N} \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\dot{\Delta}_N(\mu) d\mu}{i^c \sqrt{\Delta_N^2(\mu) - 4}} - \frac{\mathfrak{p}_N}{N^2} \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\dot{\Delta}_N(\mu) d\mu}{i^c \sqrt{\Delta_N^2(\mu) - 4}} \quad (6.7)$$

$$(T2) := \frac{1}{N} \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{(\mu + 2) \dot{\Delta}_N(\mu) d\mu}{i^c \sqrt{\Delta_N^2(\mu) - 4}} \quad (6.8)$$

$$(T3) := \frac{1}{N} \sum_{\substack{1 \leq k \leq M_0 \\ I_k^N \neq 0}} I_k^N \omega_k^N \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\varphi_k^N(\mu) d\mu}{i \sqrt[c]{\chi_N(\mu)}}. \quad (6.9)$$

$$(T4) := \frac{1}{N} \sum_{\substack{M_0 < k < N \\ I_k^N \neq 0}} I_k^N \omega_k^N \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\varphi_k^N(\mu) d\mu}{i \sqrt[c]{\chi_N(\mu)}} \quad (6.10)$$

The various terms are treated individually. We begin with the term (T1). Notice that $|\Delta_N(\mu)| < 2 \forall \mu \in (\lambda_{2j-2}^N, \lambda_{2j-1}^N)$ and

$$\Delta_N(\lambda_{2j-1}^N) = 2(-1)^{N+j}, \quad \Delta_N(\lambda_{2j-2}^N) = 2(-1)^{N+j+1}$$

and denote by $\arccos(x)$ the principal branch of the inverse of cosine, defined on $[-1, 1]$ with $\arccos(-1) = \pi$ and $\arccos(1) = 0$. Then for $s \in \{\pm 1\}$

$$\frac{d}{d\mu} \arccos\left(\frac{(-1)^s \Delta_N(\mu)}{2}\right) = -\frac{\dot{\Delta}_N(\mu) (-1)^s}{\sqrt[+]{4 - \Delta_N(\mu)^2}} \quad \forall \mu \in (\lambda_{2j-2}^N, \lambda_{2j-1}^N).$$

As by the definition of the c -root, $\sqrt[+]{4 - \Delta_N(\mu)^2} = (-1)^{N+1+j} i \sqrt[c]{\Delta_N(\mu)^2 - 4}$ one concludes that $\frac{\dot{\Delta}_N(\mu)}{i \sqrt[c]{\Delta_N(\mu)^2 - 4}} = \frac{d}{d\mu} \arccos\left(\frac{(-1)^{N+j} \Delta_N(\mu)}{2}\right)$ and hence

$$\int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\dot{\Delta}_N(\mu)}{i \sqrt[c]{\Delta_N(\mu)^2 - 4}} d\mu = -\pi.$$

Finally, as $\mathfrak{p}_N = O(N^{-3})$ by [2, Proposition 8.1],

$$(T1) = \frac{2\pi n}{N} + O\left(\frac{n}{N^5}\right). \quad (6.11)$$

To estimate the term (T2) integrate by parts to get for (T2)

$$\begin{aligned} & \sum_{j=1}^n \frac{\mu + 2}{N} \arccos\left(\frac{(-1)^{N+j} \Delta_N(\mu)}{2}\right) \Big|_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} - \frac{1}{N} \sum_{j=1}^n \int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \arccos\left(\frac{(-1)^{N+j} \Delta_N(\mu)}{2}\right) d\mu \\ &= -\frac{\pi}{4N^3} \sum_{j=1}^n \nu_{2j-2}^N - \frac{1}{4N^3} \sum_{j=1}^n \int_{\nu_{2j-2}^N}^{\nu_{2j-1}^N} \arccos\left(\frac{(-1)^{N+j}}{2} \Delta_N\left(-2 + \frac{\lambda}{4N^2}\right)\right) d\lambda \end{aligned}$$

where we made the change of variable of integration $\mu = -2 + \lambda/4N^2$. As by Theorem 2.1 $\nu_j^N = \lambda_j^- + O(M^2/N)$

$$(T2) = -\frac{\pi}{4N^3} \sum_{j=1}^n \lambda_{2j-2}^N + O\left(\frac{nM^2}{N^4}\right) - (P1). \quad (6.12)$$

Next observe that $(T3) = (P2)$. Concerning $(T4)$ note that the summation indices k, j in the definition of $(T4)$ satisfy $n < k$ and $j \leq n$, implying that the improper Riemann integral $\int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\varphi_k^N(\mu) d\mu}{i \sqrt{\chi_N(\mu)}}$ exists even if $I_k^N = 0$. Hence splitting up $(T4)$ we get $(T4) = (P3) + (E1) + (E2) + (E3)$ where

$$(E1) := \sum_{M < k \leq \frac{N}{2}} \dots \quad (E2) := \sum_{\frac{N}{2} < k < N-M} \dots \quad (E3) := \sum_{N-M \leq k < N} \dots$$

In fact, $\int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\varphi_k^N(\mu)}{\sqrt{\chi_N(\mu)}} d\mu = O\left(\frac{1}{\tau_k^N - \tau_j^N} \frac{1}{N}\right)$ for such k, j by Lemma B.1 yielding

$$(T4) = O\left(\sum_{n < k < N} I_k^N \omega_k^N \sum_{j=1}^n \frac{1}{\tau_k^N - \tau_j^N} \frac{1}{N^2}\right).$$

Concerning the estimate (6.6), note that for $1 \leq j \leq n$ and $n < k \leq M$, $\frac{1}{\tau_k^N - \tau_j^N} = O\left(\frac{N^2}{k^2 - j^2}\right)$ (cf Theorem 2.1), $I_k^N = O\left(\frac{1}{kN^2}((\gamma_k^-)^2 + \frac{M^4}{N^2})\right)$ (Proposition 3.1), and $\omega_k^N = O\left(\frac{k}{N}\right)$ (Proposition 6.3). As $((k^2 \gamma_k^-)^2)_{k \geq 1}$ is summable and $\sum_{j=1}^n \frac{1}{k^2 - j^2} = O\left(\frac{\log(n+1)}{k}\right)$ (Lemma C.7) one has for any $n \leq M_0 \leq M$,

$$\begin{aligned} (P3) &= O\left(\sum_{\substack{M_0 < k \leq M \\ 1 \leq j \leq n}} \frac{1}{N^3} \left(\frac{(k^2 \gamma_k^-)^2}{k^4} \frac{1}{k^2 - j^2} + \frac{M^4}{N^2} \frac{1}{k^2 - j^2}\right)\right) \\ &= O\left(\frac{\log(n+1)}{N^3} \left(\frac{1}{M_0^5} + \frac{M^4 \log M}{N^2}\right)\right). \end{aligned}$$

To estimate $(E1)$ note that for $M < k \leq N/2$, $\omega_k^N = O\left(\sin \frac{k\pi}{N}\right)$ (Proposition 6.2), $I_k^N = O\left(\frac{1}{M^2 N^2} \frac{1}{k}\right)$ (Proposition 3.1), and $\frac{1}{\tau_k^N - \tau_j^N} = O\left(\frac{1}{\cos \frac{k\pi}{N} - \cos \frac{M\pi}{N}}\right)$ (Theorem 2.1, Lemma A.2). By Lemma C.5 one then concludes that

$$(E1) = \sum_{M < k \leq \frac{N}{2}} \sum_{1 \leq j \leq n} O\left(\frac{1}{M^2 N^2} \frac{1}{k} \frac{\sin \frac{k\pi}{N}}{\cos \frac{k\pi}{N} - \cos \frac{M\pi}{N}} \frac{1}{N^2}\right) = O\left(\frac{n}{M^3 N^3}\right). \quad (6.13)$$

To estimate (E2), observe that $\omega_k^N = O(\sin \frac{k\pi}{N})$ (Proposition 6.3), $I_k^N = O(\frac{1}{M^2 N^2} \frac{1}{N-k})$ (Proposition 3.1), and $\frac{1}{\tau_k^N - \tau_j^N} = O(1)$ (Theorem 2.1, cf Lemma A.2)) and hence

$$(E2) = O\left(\sum_{\frac{N}{2} \leq k < N-M} \sum_{1 \leq j \leq n} \frac{1}{M^2 N^2} \frac{\sin \frac{k\pi}{N}}{N-k} \frac{1}{N^2}\right) = O\left(\frac{n}{M^2 N^4}\right) \quad (6.14)$$

where we used that $\sum_{\frac{N}{2} \leq k < N-M} \frac{1}{N-k} \sin \frac{k\pi}{N} = \sum_{M < \ell \leq \frac{N}{2}} \frac{1}{\ell} \sin \frac{(N-\ell)\pi}{N}$ is bounded as $\sin \frac{(N-\ell)\pi}{N} = \sin \frac{\ell\pi}{N} \leq \frac{\ell\pi}{N}$. Finally we consider (E3). Changing the index of summation, $\ell = N - k$ and using that $\omega_{N-\ell}^N = O(\frac{\ell}{N})$ (Proposition 6.3), $I_{N-\ell}^N = O(\frac{1}{\ell N^2} ((\gamma_\ell^+)^2 + \frac{M^4}{N^2}))$ (Proposition 3.1), and $\frac{1}{\tau_{N-\ell}^N - \tau_j^N} = O(1)$ (Theorem 2.1, cf Lemma A.2))

$$(E3) = O\left(\sum_{1 \leq \ell \leq M} \sum_{1 \leq j \leq n} I_{N-\ell}^N \omega_{N-\ell}^N \frac{1}{N^2}\right) = O\left(\frac{n}{N^5}\right). \quad (6.15)$$

Combining the estimates obtained for (E1), (E2), and (E3) yields the claimed bound (6.5). Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

7 End of proof of Theorem 1.1

In this section we complete the proof of Theorem 1.1 using the results obtained so far together with auxiliary estimates obtained in Appendix D. First we prove the asymptotics (1.6) for the frequencies near the edges.

Proposition 7.1 *Uniformly for $1 \leq n \leq M$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$ one has $\omega_n^N, \omega_{N-n}^N = \frac{2\pi n}{N} + O(\frac{n^3}{N^3})$.*

Proof. The asymptotics of ω_n^N and ω_{N-n}^N are shown in a similar way so we concentrate on ω_n^N . Starting point is formula (6.4) with the choice $M_0 = M$,

$$\omega_n^N - \frac{2\pi n}{N} = -\frac{\pi}{4N^3} \sum_{j=1}^n \lambda_{2j-2}^- - (P1) - (P2) + Error.$$

Note that $\sum_{j=1}^n \lambda_{2j-2}^- = O(n^3)$ (asymptotics of $(\lambda_j^-)_{j \geq 1}$) and $(P1) = O(\frac{n^2}{N^3})$ (Theorem 2.1) where $(P1) = \frac{1}{4N^3} \sum_{j=1}^n \int_{\nu_{2j-2}^N}^{\nu_{2j-1}^N} \arccos\left(\frac{(-1)^{N+j}}{2} \Delta_N(-2 + \frac{\lambda}{4N^2})\right) d\lambda$.

Concerning $(P2) = \frac{1}{N} \sum_{I_k^N \neq 0, 1 \leq k \leq M} I_k^N \omega_k^N \int_{\mathcal{I}_n^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu$ note that $N^2 I_k^N = O\left(\frac{1}{k}((\gamma_k^-)^2 + \frac{M^4}{N^2})\right)$ (Proposition 3.1), $\sum_{k=1}^{\infty} k^4 (\gamma_k^-)^2 < \infty$ ($q_- \in C_0^2(\mathbb{T})$), $\frac{I_k^N}{\gamma_k^-} = O(1)$ (Remark 3.3, Theorem 3.2), and $\omega_k^N = O\left(\frac{k}{N}\right)$ (Proposition 6.4). It then follows from Corollary D.2, D.4, D.6, D.8, D.10, and Corollary D.12 that $(P2) = O\left(\frac{1}{N^3}\right)$. As $Error = O\left(\frac{1}{N^3}\left(\frac{nM^2}{N} + \frac{n}{M^3}\right)\right)$ (Proposition 6.4) and thus $Error = O\left(\frac{1}{N^3}\right)$, the stated asymptotics follow. Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

To prove the asymptotics (1.3) for the frequencies at the left edge. the starting point is formula (6.4) with the choice $M_0 = L$,

$$\omega_n^N - \frac{2\pi n}{N} = -\frac{\pi}{4N^3} \sum_{j=1}^n \lambda_{2j-2}^- - (P1) - (P2) - (P3) + Error. \quad (7.1)$$

First consider $(P1) = \frac{1}{4N^3} \sum_{j=1}^n \int_{\nu_{2j-2}^N}^{\nu_{2j-1}^N} \arccos\left(\frac{(-1)^{N+j}}{2} \Delta_N\left(-2 + \frac{\lambda}{4N^2}\right)\right) d\lambda$.

Lemma 7.2 *Uniformly for $1 \leq n \leq L$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$-\pi \sum_{j=1}^n \lambda_{2j-2}^- - 4N^3(P1) = \sum_{1 \leq j \leq n} \int_{\lambda_{2j-2}^-}^{\lambda_{2j-1}^-} \frac{\lambda \dot{\Delta}_-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\frac{n^2 L}{M^{1/2}}\right). \quad (7.2)$$

Proof. Integrating by parts, one gets for any $1 \leq j \leq n$

$$\int_{\lambda_{2j-2}^-}^{\lambda_{2j-1}^-} \frac{\lambda \dot{\Delta}_-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda = -\pi \lambda_{2j-2}^- - \int_{\lambda_{2j-2}^-}^{\lambda_{2j-1}^-} \arccos\left(\frac{(-1)^j}{2} \Delta_-(\lambda)\right) d\lambda.$$

To estimate the difference of $(Q_j^N) = \int_{\nu_{2j-2}^N}^{\nu_{2j-1}^N} \arccos\left(\frac{(-1)^{N+j}}{2} \Delta_N\left(-2 + \frac{\lambda}{4N^2}\right)\right) d\lambda$

with the corresponding KdV integral, $(Q_j^-) = \int_{\lambda_{2j-2}^-}^{\lambda_{2j-1}^-} \arccos\left(\frac{(-1)^j}{2} \Delta_-(\lambda)\right) d\lambda$,

we use that \arccos is $\frac{1}{2}$ -Hölder continuous. Indeed, making in the integral (Q_j^N) the change of variable $[0, \delta_j^-] \rightarrow [\lambda_{2j-1}^-, \lambda_{2j-2}^-]$, $x \mapsto \lambda(x) = \nu_{2j-2}^N + h_j^N x$ with $\delta_j^- = \lambda_{2j-1}^- - \lambda_{2j-2}^-$ and $h_j^N = \frac{\nu_{2j-1}^N - \nu_{2j-2}^N}{\lambda_{2j-1}^- - \lambda_{2j-2}^-}$, the difference $(Q_j^N) - (Q_j^-)$ takes the form

$$h_j^N \int_0^{\delta_j^-} f(x) dx + (h_j^N - 1) \int_0^{\delta_j^-} \arccos\left(\frac{(-1)^j}{2} \Delta_-(\lambda_{2j-2}^- + x)\right) dx \quad (7.3)$$

where $f(x) := \arccos\left(\frac{(-1)^j}{2}\Delta_N(\lambda(x))\right) - \arccos\left(\frac{(-1)^j}{2}\Delta_-(\lambda_{2j-2}^- + x)\right)$ satisfies by Theorem 2.2

$$f(x) = \arccos\left(\frac{(-1)^j}{2}\Delta_-(\lambda(x)) + O\left(\frac{L^2}{M}\right)\right) - \arccos\left(\frac{(-1)^j}{2}\Delta_-(\lambda_{2j-2}^- + x)\right).$$

Let us first consider the second term on the right hand side of (7.3). As the length δ_j^- of the j 'th band satisfies $1/\delta_j^- = O(1/j)$, one concludes from Theorem 2.1 that $h_j^N = 1 + O\left(\frac{M^2}{jN}\right)$. Hence

$$(h_j^N - 1) \int_0^{\delta_j^-} \arccos\left(\frac{(-1)^j}{2}\Delta_-(\lambda_{2j-2}^- + x)\right) dx = O\left(\frac{M^2}{N}\right).$$

Towards the first term on the right hand side of (7.3), note that by the Lipschitz continuity of Δ_- , $\Delta_-(\lambda(x)) - \Delta_-(\lambda_{2j-2}^- + x) = O(\lambda(x) - \lambda_{2j-2}^- - x)$. As $\lambda(x) - (\lambda_{2j-2}^- + x) = \nu_{2j-2}^N - \lambda_{2j-2}^- + (h_j^N - 1)x$ it follows that

$$\Delta_-(\lambda(x)) - \Delta_-(\lambda_{2j-2}^- + x) = O\left(j\frac{M^2}{N}\right),$$

and hence, by the $\frac{1}{2}$ -Hölder continuity of \arccos , $f(x) = O\left(\frac{L}{M^{1/2}}\right)$ (use that by (5.1), $j\frac{M^2}{N} \leq \frac{L^2}{M}$), implying that

$$h_j^N \int_0^{\delta_j^-} f(x) dx = O\left(j\frac{L}{M^{1/2}}\right).$$

Combining the obtained estimates, the stated asymptotics follow. Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

Next we consider the term $(P2) = \frac{1}{N} \sum_{\substack{1 \leq k \leq L \\ I_k^N \neq 0}} I_k^N \omega_k^N \int_{\mathcal{I}_n^N} \frac{\varphi_k^N(\mu)}{i \sqrt{\chi_N(\mu)}} d\mu$ in (7.1).

In Appendix D we show that the analogous expression for KdV plays an important role in describing the asymptotics of $(P2)$. The asymptotics proved in the lemmas of Appendix D show that $\gamma_k^N \int_{\mathcal{I}_n^N} \frac{\varphi_k^N(\mu)}{i \sqrt{\chi_N(\mu)}} d\mu$ stays bounded if I_k^N approaches 0. By Theorem 3.2, $(P2)$ can thus be written alternatively as $(P2) = \frac{1}{N} \sum_{1 \leq k \leq L} J_k^N \omega_k^N \cdot \gamma_k^N \int_{\mathcal{I}_n^N} \frac{\varphi_k^N(\mu)}{i \sqrt{\chi_N(\mu)}} d\mu$.

Lemma 7.3 *Uniformly for $1 \leq n \leq L$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$(P2) = \frac{1}{8N^3} \sum_{\substack{1 \leq k \leq L \\ I_k^- \neq 0}} I_k^- \int_{\mathcal{I}_n^-} \frac{\psi_k^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\frac{1}{N^3} \left(\frac{L}{M^{1/2}} + \frac{1}{L^{5/2}}\right)\right). \quad (7.4)$$

Proof. Using the terminology of Appendix D, decompose (P2) into three parts,

$$(P2) = \frac{1}{N} \sum_{\substack{1 \leq k \leq L \\ I_k^N \neq 0, k \neq n}} I_k^N \omega_k^N \int_{\mathcal{I}_n^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu \\ + \frac{1}{N} I_n^N \omega_n^N \int_{\mathcal{I}_n^N \setminus C_n^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu + \frac{1}{4N^3} J_n^N \omega_n^N \tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu$$

where we used that $I_n^N = J_n^N \tilde{\gamma}_n^N$ and $\tilde{\gamma}_n^N = 4N^2 \gamma_n^N$. The three parts are treated separately, but in a similar way, using the asymptotics of I_k^N (Appendix 3), ω_k^N (Proposition 7.1), and the integrals (Section D). Each of them yields $O\left(\frac{1}{N^3} \left(\frac{L}{M^{1/2}} + \frac{1}{L^{5/2}}\right)\right)$ as an error term. Let us compute the asymptotics for the third part in more detail. Write it as

$$J_n^N \omega_n^N \tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = J_n^- \gamma_n^- \int_{C_n^-} \frac{\psi_n^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + (T1) + (T2) + (T3)$$

where $(T1) := (J_n^N - J_n^-) \omega_n^N \tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu) d\mu}{i^c \sqrt{\chi_N(\mu)}}$, $(T2) := J_n^- (\omega_n^N - \frac{2\pi n}{N}) \tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu) d\mu}{i^c \sqrt{\chi_N(\mu)}}$,

$$(T3) := J_n^- \left(\frac{2\pi n}{N} \tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu - \gamma_n^- \int_{C_n^-} \frac{\psi_n^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda \right).$$

For any $1 \leq k \leq L$, one has $J_k^N - J_k^- = O\left(\frac{L}{M^{1/2}}\right)$ (Theorem 3.2, Remark 6.1), $\omega_n^N = \frac{2\pi n}{N} + O\left(\frac{n^3}{N^3}\right)$ (Proposition 7.1) and by Lemma D.9,

$$\frac{2\pi n}{N} \tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = \gamma_n^- \int_{C_n^-} \frac{\psi_n^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\frac{L}{M^{1/2}} + \frac{1}{L^{5/2}}\right).$$

Furthermore, $J_n^- = O(n^{-5})$ ([16]), $\omega_n^N = O\left(\frac{n}{N}\right)$ (Proposition 7.1), and finally $\tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = O\left(\frac{N}{n}\right)$ (Corollary D.10). Combining all these estimates

and using that $J_n^- \gamma_n^- = I_n^-/2$ one then concludes that

$$J_n^N \omega_n^N \tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = \frac{1}{2} I_n^- \int_{C_n^-} \frac{\psi_n^-(\lambda)}{i^c \sqrt{\Delta_-(\lambda)^2 - 4}} d\lambda + O\left(\frac{L}{M^{1/2}} + \frac{1}{L^{5/2}}\right).$$

As already mentioned above, the first second part in the decomposition of (P2) are estimated in a similar way. Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

Combining the results obtained so far we get the following asymptotics for the frequencies at the left edge.

Theorem 7.4 *Uniformly for $1 \leq n \leq L$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$,*

$$\omega_n^N = \frac{2\pi n}{N} - \frac{1}{24} \frac{1}{(2N)^3} \omega_n^- + O\left(\frac{1}{N^3} \left(\frac{n^2 L}{M^{1/2}} + \frac{1}{L^{5/2}}\right)\right) \quad (7.5)$$

Remark 7.5 *Correspondingly, the asymptotics of ω_{N-n}^N are given by $\omega_{N-n}^N = \frac{2\pi n}{N} - \frac{1}{24} \frac{1}{(2N)^3} \omega_n^+ + O\left(\frac{1}{N^3} \left(\frac{n^2 L}{M^{1/2}} + \frac{1}{L^{5/2}}\right)\right)$.*

Proof. By (7.1), Lemma 7.2 and Lemma 7.3, one has for any $1 \leq n \leq L$,

$$\begin{aligned} \omega_n^N &= \frac{2\pi n}{N} + \frac{1}{4N^3} \sum_{1 \leq j \leq n} \int_{\lambda_{2j-2}^-}^{\lambda_{2j-1}^-} \frac{\lambda \dot{\Delta}_-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\frac{1}{4N^3} \frac{n^2 L}{M^{1/2}}\right) \\ &- \frac{1}{8N^3} \sum_{\substack{1 \leq k \leq L \\ I_k^- \neq 0}} I_k^- \int_{\mathcal{I}_n^-} \frac{\psi_k^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\frac{1}{N^3} \left(\frac{L}{M^{1/2}} + \frac{1}{L^{5/2}}\right)\right) - (P3) + Error. \end{aligned}$$

By the formula for the KdV frequencies of Proposition 4.6, we thus get

$$\begin{aligned} \omega_n^N &= \frac{2\pi n}{N} - \frac{1}{8N^3} \frac{1}{24} \omega_n^- + \frac{1}{8N^3} \frac{1}{24} \sum_{\substack{k > L \\ I_k^- \neq 0}} I_k^- \int_{\mathcal{I}_n^-} \frac{\psi_k^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda \\ &+ O\left(\frac{1}{N^3} \left(\frac{n^2 L}{M^{1/2}} + \frac{1}{L^{5/2}}\right)\right) - (P3) + Error. \end{aligned}$$

By Proposition 6.4, $Error = O\left(\frac{1}{N^3} \left(\frac{nM^2}{N} + \frac{n}{M^3}\right)\right)$. As by (5.1), $\frac{nM^2}{N} \leq \frac{n^2 L}{M^{1/2}}$ and $\frac{n}{M^3} \leq \frac{1}{L^{5/2}}$, one concludes that $Error = O\left(\frac{1}{N^3} \left(\frac{n^2 L}{M^{1/2}} + \frac{1}{L^{5/2}}\right)\right)$. Taking into

account in addition that $M_0 = L$, Proposition 6.4 also implies that $(P_3) = O\left(\frac{1}{N^3}\left(\frac{n^2L}{M^{1/2}} + \frac{1}{L^{5/2}}\right)\right)$. Arguing as in Appendix B one sees that for any $k > L$, $\int_{\mathcal{I}_n^-} \frac{\psi_k^-(\lambda)}{i^\epsilon \sqrt{\Delta_-^2(\lambda)-4}} d\lambda = O\left(\frac{\log(n+1)}{k^2}\right)$. As $\alpha, \beta \in C_0^2(\mathbb{T})$, one has $\sum_{k=1}^\infty k^5 I_k^- < \infty$ ([17]), implying that $\sum_{\substack{k>L \\ I_k^- \neq 0}} I_k^- \int_{\mathcal{I}_n^-} \frac{\psi_k^-(\lambda)}{i^\epsilon \sqrt{\Delta_-^2(\lambda)-4}} d\lambda = O\left(\frac{\log L}{L^5}\right)$. Hence the stated asymptotics are proved. Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

Proof of Theorem 1.1 By Theorem 7.4 and the definitions of \mathcal{H}_{KdV}^N and ω_n^- , the claimed asymptotics (1.3) of the frequencies at the left edge hold. The asymptotics (1.4) of the frequencies at the right edge are proved using the symmetry of Toda chains, discussed in Appendix E. Indeed, let

$$\tilde{\beta}(x) := -\beta(-x); \quad \tilde{\alpha}(x) := \alpha(-x); \quad \tilde{q}_-(x) = -2\tilde{\alpha}(2x) + \tilde{\beta}(2x)$$

and note that $\tilde{q}_-(-x) = q_+(x)$. Furthermore, one verifies in a straightforward way that the periodic eigenvalues of $-\partial_x^2 + \tilde{q}_-(-x)$ coincide with the ones of $-\partial_x^2 + \tilde{q}_-(x)$. Hence q_+ and \tilde{q}_- have the same actions and the same KdV frequencies. In view of Appendix E one has $\tilde{b}_n^N = -b_{N-n}^N = \frac{1}{4N^2}\tilde{\beta}\left(\frac{n}{N}\right)$ and

$$\tilde{a}_n^N = a_{N-n-1}^N = 1 + \frac{1}{4N^2}\tilde{\alpha}\left(\frac{n+1}{N}\right) = 1 + \frac{1}{4N^2}\tilde{\alpha}\left(\frac{n}{N}\right) + O\left(\frac{1}{N^3}\right).$$

By (E.1), the claimed asymptotics (1.4) then follow from (1.3), applied to $\tilde{\beta}$ and $\tilde{\alpha}$. Alternatively, one can derive the asymptotics (1.4) from (1.3) using Corollary E.2. The asymptotics (1.6) of the frequencies near the left and right edges are established in Proposition 7.1. Finally the asymptotics (1.5) of the frequencies in the bulk are given in 6.2. \square

A Spectral estimates

In this appendix we apply Theorem 2.1 to prove formulas involving, the eigenvalues of the Jacobi matrices $Q_N^{\alpha,\beta}$, used at various places of the paper.

Proposition A.1 *Under the same assumptions as in Theorem 2.1 the following estimates hold uniformly in $1 \leq n \leq N/2$ and on bounded subsets of functions $\alpha, \beta \in C_0^2(\mathbb{T})$:*

(i) If $1 \leq n \leq M$, then

$$\sum_{0 < k < n} \frac{\gamma_k^N}{\lambda_{2n-1}^N - \lambda_{2k}^N}, \quad \sum_{n < k \leq M} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{1}{n}\right), \quad \sum_{M < k \leq N/2} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{\log M}{M^2}\right),$$

$$\sum_{N/2 < k < N-M} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{1}{NM}\right), \quad \sum_{N-M \leq k < N} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{1}{N^2} + \frac{M^3}{N^3}\right).$$

(ii) If $M < n \leq N/2$, then

$$\sum_{0 < k \leq M} \frac{\gamma_k^N}{\lambda_{2n-1}^N - \lambda_{2k}^N} = O\left(\frac{1}{M^2} + \frac{M \log M}{N}\right)$$

$$\sum_{M < k < n} \frac{\gamma_k^N}{\lambda_{2n-1}^N - \lambda_{2k}^N}, \quad \sum_{n < k \leq \frac{N}{2}} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{\log M}{M^2}\right),$$

$$\sum_{\frac{N}{2} < k < N-M} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{\log M}{M^2}\right), \quad \sum_{N-M \leq k < N} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{1}{N^2} + \frac{M^3}{N^3}\right).$$

Analogous results hold for $N/2 < n < N$.

To prove Proposition A.1 we first establish the following auxiliary result.

Lemma A.2 *Under the same assumptions as in Theorem 2.1 there exist $N_0 \geq 3$ and $C \geq 1$ so that*

$$\lambda_{2n-1}^N - \lambda_{2k}^N \geq \frac{1}{C} \frac{n^2 - k^2}{N^2} \quad \forall 0 \leq k < n \leq \frac{N}{2}$$

$$\lambda_{2k-1}^N - \lambda_{2n}^N \geq \frac{1}{C} \frac{k^2 - n^2}{N^2} \quad \forall 0 \leq n < k \leq \frac{N}{2}$$

where C and N_0 can be chosen uniformly for k, n and on bounded subsets of $\alpha, \beta \in C_0^2(\mathbb{T})$. Similar estimates hold for $\frac{N}{2} \leq k < n \leq N$ respectively $\frac{N}{2} \leq n < k \leq N$.

Proof of Lemma A.2. Recall from [16, Proposition B.11] that the periodic eigenvalues of Hill operators are compact functions of the potential on $L^2(\mathbb{T})$. It then follows from the Counting Lemma that for any $k \geq 0$, the k 'th band length $\lambda_{2k+1}^- - \lambda_{2k}^-$ can be estimated from below uniformly on bounded subsets

of functions $\alpha, \beta \in C_0^2(\mathbb{T})$. Furthermore, $\lambda_{2\ell}^-, \lambda_{2\ell-1}^- = \frac{\ell^2 \pi^2}{N^2} + O(\frac{1}{\ell^2})$ as $\ell \rightarrow \infty$ where the error term is uniform on bounded subsets of functions $\alpha, \beta \in C_0^2(\mathbb{T})$ (cf e.g. [17]). After these preparations we can prove the claimed estimates. For the following we choose $N \geq N_0$ with $N_0 \geq 3$ sufficiently large. Let us first consider the case where $0 < n \leq M$. If $0 \leq k < n$, it follows from the above lower bound of the spectral bands and Theorem 2.1 that $\lambda_{2n-1}^N - \lambda_{2k}^N \geq \frac{1}{C} \frac{n^2 - k^2}{N^2}$ for some $C \geq 1$. The case where $n < k \leq M$ can be treated in the same way. If $M < k \leq \frac{N}{2}$, write

$$\lambda_{2k-1}^N - \lambda_{2n}^N = (\lambda_{2k-1}^N + 2 \cos \frac{M\pi}{N}) + (-2 \cos \frac{M\pi}{N} - \lambda_{2M}^N) + (\lambda_{2M}^N - \lambda_{2n}^N)$$

and estimate each expression in a bracket separately. We have already seen that $\lambda_{2M}^N - \lambda_{2n}^N \geq \frac{1}{C} \frac{M^2 - n^2}{N^2}$. As by Theorem 2.1, $\lambda_{2k-1}^N = -2 \cos \frac{k\pi}{N} + O(\frac{1}{N^2M})$ it follows from (C.1) that

$$\lambda_{2k-1}^N + 2 \cos \frac{M\pi}{N} = (2 \cos \frac{M\pi}{N} - 2 \cos \frac{k\pi}{N})(1 + O(\frac{1}{M^2})).$$

Hence by (C.2) there exists $C \geq 1$ so that $\lambda_{2k-1}^N + 2 \cos \frac{M\pi}{N} \geq \frac{1}{C} \frac{k^2 - M^2}{N^2}$. Finally, by Theorem 2.1 and the asymptotics of the eigenvalues λ_ℓ^- reviewed above, $\lambda_{2M}^N = -2 + \frac{M^2 \pi^2}{N^2} + O(\frac{1}{N^2 M^2})$ whereas by a straightforward Taylor expansion, $2 \cos \frac{M\pi}{N} = 2 - \frac{M^2 \pi^2}{N^2} + O(\frac{M^4}{N^4})$. Hence we get $-2 \cos \frac{M\pi}{N} - \lambda_{2M}^N = O(\frac{1}{N^2}(\frac{1}{M^2} + \frac{M^4}{N^2}))$. In view of the assumption $\eta < 1/2$ in (F), we then conclude that for $C \geq 1$ sufficiently large, $\lambda_{2k-1}^N - \lambda_{2n}^N \geq \frac{1}{C} \frac{k^2 - n^2}{N^2}$. The case where $M < n \leq \frac{N}{2}$ is treated in a similar fashion. Going through the arguments of the proof one verifies that $C \geq 1$ and $N_0 \geq 3$ can be chosen uniformly for k, n and for bounded subsets of functions $\alpha, \beta \in C_0^2(\mathbb{T})$. The estimates for $\frac{N}{2} \leq k < n \leq N$ respectively $\frac{N}{2} \leq n < k \leq N$ are proven in the same way. \square

Proof of Proposition A.1. (i) By Lemma A.2 and Theorem 2.1

$$\sum_{0 < k < n} \frac{\gamma_k^N}{\lambda_{2n-1}^N - \lambda_{2k}^N} = O\left(\sum_{0 < k < n} \frac{\gamma_k^-}{n^2 - k^2} + \frac{M^2}{N} \sum_{0 < k < n} \frac{1}{n^2 - k^2}\right) = O\left(\frac{1}{n}\right)$$

Similarly, $\sum_{n < k \leq M} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O(\frac{1}{n})$ and in view of Lemma C.7 one gets

$$\sum_{M < k \leq N/2} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} \leq O\left(\frac{1}{M} \sum_{M < k \leq N/2} \frac{1}{k^2 - n^2}\right) = O\left(\frac{\log M}{M^2}\right).$$

As for any $\frac{N}{2} < k < N - M$ and N sufficiently large, $\lambda_{2k-1}^N - \lambda_{2n}^N \geq \frac{1}{C}$ and $\gamma_k^N = O(\frac{1}{N^2M})$ one concludes

$$\sum_{N/2 < k < N-M} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{1}{NM}\right)$$

Finally, by similar arguments and taking into account that $\gamma_k^N = O(\frac{\gamma_{N-k}^+}{N^2} + \frac{M^2}{N^3})$ for $N - M \leq k < N$ one gets

$$\sum_{N-M \leq k < N} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{1}{N^2} \sum_{1 \leq \ell \leq M} \gamma_\ell^+ + \frac{M^3}{N^3}\right) = O\left(\frac{1}{N^2} + \frac{M^3}{N^3}\right).$$

(ii) The claimed estimates are proved in a similar way as the ones of item (i). We only mention that to prove $\sum_{0 < k \leq M} \frac{\gamma_k^N}{\lambda_{2n-1}^N - \lambda_{2k}^N} = O(\frac{1}{M^2} + \frac{M \log M}{N})$ we split the sum $\sum_{0 < k \leq M} = \sum_{0 < k \leq \frac{M}{2}} + \sum_{\frac{M}{2} < k \leq M}$ and use that $\sum_{\ell > 0} (\ell^2 \gamma_\ell^-)^2 < \infty$ and for the estimate $\sum_{\frac{N}{2} < k < N-M} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O(\frac{\log M}{M^2})$ we use that $\lambda_{2k-1}^N - \lambda_{2n}^N = -2 \cos \frac{k\pi}{N} + 2 \cos \frac{n\pi}{N} + O(\frac{1}{N^2M})$, $\cos \frac{n\pi}{N} \geq 0$ and $-\cos \frac{k\pi}{N} = \cos \frac{(N-k)\pi}{N}$ to conclude that by Lemma C.4

$$\sum_{\frac{N}{2} < k < N-M} \frac{\gamma_k^N}{\lambda_{2k-1}^N - \lambda_{2n}^N} = O\left(\frac{1}{N^2M} \sum_{M < k < \frac{N}{2}} \frac{1}{\cos \frac{k\pi}{N}}\right) = O\left(\frac{\log M}{M^2}\right).$$

Going through the arguments of the proofs of item (i) and (ii) one verifies that the estimates hold uniformly for $1 \leq n \leq \frac{N}{2}$ and on bounded subsets of functions $\alpha, \beta \in C_0^2(\mathbb{T})$. The estimates for $\frac{N}{2} < n \leq N$ are proven in the same way. \square

B Estimates of products

We prove estimates on products needed at various places of the paper. Throughout this appendix, $F : \mathbb{N} \rightarrow \mathbb{R}_{\geq 1}$ denotes a function satisfying (F) with $\eta \leq 1/3$, $M = [F(N)]$, $L = [F(M)]$, and $w_k^N(\mu)$, $w_k^\pm(\lambda)$ the standard roots introduced in Section 5.

In Section 5 and Section 6 we need the following lemmas.

Lemma B.1 For any $1 \leq j \leq M$ and $0 < k < N$, $k \neq j, j-1$

$$\int_{\lambda_{2j-2}^N}^{\lambda_{2j-1}^N} \frac{\varphi_k^N(\mu)}{\sqrt[3]{\chi_N(\mu)}} d\mu = O\left(\frac{1}{\tau_k^N - \tau_j^N} \frac{1}{N}\right). \quad (\text{B.1})$$

The estimate holds uniformly in j, k and uniformly on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

Proof. Let us first consider the case $j = 1$. Then by assumption, $k \geq 2$. As $\lambda_{2\ell-1}^N \leq \sigma_\ell^{N,k} \leq \lambda_{2\ell}^N$ one sees that for $\lambda_0^N \leq \mu \leq \lambda_1^N$, $\left| \frac{\varphi_k^N(\mu)}{\sqrt[3]{\chi_N(\mu)}} \right|$ is bounded by

$$\frac{1}{\sqrt[3]{(\lambda_1^N - \mu)(\mu - \lambda_0^N)}} \frac{\sqrt[3]{\lambda_2^N - \lambda_0^N}}{\sqrt[3]{\lambda_{2N-1}^N - \lambda_1^N}} \cdot \frac{1}{\lambda_{2k-1}^N - \lambda_1^N} \cdot \prod_{\substack{1 < \ell < N \\ \ell \neq k}} \left(\frac{\lambda_{2\ell}^N - \mu}{\lambda_{2\ell-1}^N - \mu} \right)^{1/2}.$$

By Lemma C.6 there exists $\lambda_0^N < \rho \equiv \rho_1^{N,k} < \lambda_1^N$ so that

$$\int_{\lambda_0^N}^{\lambda_1^N} \left| \frac{\varphi_k^N(\mu)}{\sqrt[3]{\chi_N(\mu)}} \right| d\mu \leq \pi \frac{\sqrt[3]{\lambda_2^N - \lambda_0^N}}{\sqrt[3]{\lambda_{2N-1}^N - \lambda_1^N}} \cdot \frac{1}{\lambda_{2k-1}^N - \lambda_1^N} \cdot \prod_{\substack{1 < \ell < N \\ \ell \neq k}} \left(\frac{\lambda_{2\ell}^N - \rho}{\lambda_{2\ell-1}^N - \rho} \right)^{1/2}.$$

As $\frac{\lambda_{2\ell}^N - \rho}{\lambda_{2\ell-1}^N - \rho} = 1 + \frac{\gamma_\ell^N}{\lambda_{2\ell-1}^N - \rho} \leq \exp\left(\frac{\gamma_\ell^N}{\lambda_{2\ell-1}^N - \rho}\right) \leq \exp\left(\frac{\gamma_\ell^N}{\lambda_{2\ell-1}^N - \lambda_2^N}\right)$ one has

$$\prod_{\substack{\ell \neq k \\ \ell > 1}} \left(\frac{\lambda_{2\ell}^N - \rho}{\lambda_{2\ell-1}^N - \rho} \right)^{1/2} \leq \exp\left(\frac{1}{2} \sum_{\substack{\ell \neq k \\ \ell > 1}} \frac{\gamma_\ell^N}{\lambda_{2\ell-1}^N - \lambda_2^N}\right) = O(1)$$

where for the latter estimate we used Proposition A.1. Furthermore, by Theorem 2.1, $\lambda_{2N-1}^N - \lambda_1^N = 4 + O\left(\frac{1}{N^2}\right)$, $\sqrt[3]{\lambda_2^N - \lambda_0^N} = O\left(\frac{1}{N}\right)$, and $(\lambda_{2k-1}^N - \lambda_1^N)^{-1} = O\left((\tau_k^N - \tau_1^N)^{-1}\right)$. This proves the claim for $j = 1$. The case where $2 \leq j \leq M$ is treated in a similar way. Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

In Section 5 we introduced the rectangles Γ_n^- with top and bottom side given by $[\lambda_{2n-1}^- - 2\rho, \lambda_{2n}^- + 2\rho] \pm i\rho$ where $\rho > 0$ satisfies $\lambda_{2\ell}^- + 3\rho < \lambda_{2\ell+1}^- - 3\rho \quad \forall \ell \geq 0$. Furthermore, we chose $N_0 \geq 3$ in such a way that the error term in Theorem 2.1 is smaller than ρ for any $N \geq N_0$ (cf Lemma 5.3). Then

$$\lambda_{2\ell-1}^- - \rho < \nu_{2\ell-1}^N \leq \nu_{2\ell}^N < \lambda_{2\ell}^- + \rho \quad \forall 1 \leq \ell \leq M. \quad (\text{B.2})$$

Lemma B.2 For any $\mu \equiv \mu(\lambda) = -2 + \frac{\lambda}{4N^2}$ with $\lambda \in \Gamma_n^-$ and $N \geq N_0$,

$$\frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} = 1 + O\left(\frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N}\right)$$

uniformly in $1 \leq n, k \leq L$, $L < \ell < N$, and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

Proof. Given $\mu \equiv \mu(\lambda)$ with $\lambda \in \Gamma_n^-$, write $\frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} - 1 = \frac{\sigma_\ell^{N,k} - \mu - w_\ell^N(\mu)}{w_\ell^N(\mu)}$. We estimate denominator and numerator separately. In view of the definition of Γ_n^- and the lower bound $\lambda_{2n}^N - (-2 + \frac{\lambda_{2n}^-}{4N^2}) \geq -\frac{\rho}{4N^2}$ from (B.2) it follows that

$$|w_\ell^N(\mu)| \geq \lambda_{2\ell-1}^N - (-2 + \frac{\lambda_{2n}^- + \rho}{4N^2}) \geq \lambda_{2\ell-1}^N - \lambda_{2n}^N - \frac{2\rho}{4N^2}.$$

Hence $\frac{1}{w_\ell^N(\mu)} = O(\frac{1}{\tau_\ell^N - \tau_n^N})$. Concerning the numerator, it follows from the definition of the standard root that

$$w_\ell^N(\mu) = (\tau_\ell^N - \mu) \sqrt[+]{1 - \left(\frac{\gamma_\ell^N/2}{\tau_\ell^N - \mu}\right)^2} = (\tau_\ell^N - \mu) + (\tau_\ell^N - \mu) \left(\sqrt[+]{1 - \left(\frac{\gamma_\ell^N/2}{\tau_\ell^N - \mu}\right)^2} - 1\right).$$

Note that by Theorem 2.1, $\gamma_\ell^N = O(\frac{M^2}{N^3} + \frac{1}{MN^2})$ and by the definition of ρ , $|\tau_\ell^N - \mu| \geq \frac{\rho}{N^2}$. Thus by choosing N_0 larger if necessary one can assume that for any $N \geq N_0$

$$\left|\frac{\gamma_\ell^N/2}{\tau_\ell^N - \mu}\right| \leq \frac{1}{2} \quad \forall L < \ell < N, \quad \forall \mu \in \Gamma_n^- \text{ with } 1 \leq n \leq L$$

and hence

$$|\tau_\ell^N - \mu| \left|1 - \sqrt[+]{1 - \left(\frac{\gamma_\ell^N/2}{\tau_\ell^N - \mu}\right)^2}\right| \leq \frac{|\gamma_\ell^N|}{4}.$$

Finally, as $|\sigma_\ell^{N,k} - \tau_\ell^N| \leq |\gamma_\ell^N/2|$, one has $\sigma_\ell^{N,k} - \mu - w_\ell^N(\mu) = O(\gamma_\ell^N)$. Altogether we thus proved that

$$\frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} = 1 + O\left(\frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N}\right).$$

Going through the arguments of the proof one verifies the claimed uniformity statement. \square

Lemma B.2 is used to get an estimate for

$$Q_k^{N,L}(\mu) = \frac{2}{\sqrt[4]{\lambda_{2N-1}^N - \mu}} \cdot \prod_{L < \ell < N} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)}$$

introduced in (5.14). With the notation introduced above we have

Lemma B.3 *For any $\mu \equiv \mu(\lambda) = -2 + \frac{\lambda}{4N^2}$ with $\lambda \in \Gamma_n^-$, $1 \leq n, k \leq L$, and $N \geq N_0$,*

$$Q_k^{N,L}(\mu) = 1 + O\left(\frac{1}{L^3}\right) \quad (\text{B.3})$$

uniformly in $1 \leq n, k \leq L$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.

Remark B.4 *Similarly one shows that (B.3) holds for $\mu \leq (\lambda_{2L}^N + \lambda_{2L+1}^N)/2$ uniformly in $1 \leq k \leq L$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.*

Proof. By Theorem 2.1, $\frac{2}{\sqrt[4]{\lambda_{2N-1}^N - \mu}} = 1 + O\left(\frac{L^2}{N^2}\right)$ and by Lemma B.2

$$\prod_{L < \ell < N} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} = \prod_{L < \ell < N} \left(1 + O\left(\frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N}\right)\right) = \exp\left(\sum_{L < \ell < N} \log\left(1 + O\left(\frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N}\right)\right)\right)$$

yielding the estimate

$$\prod_{L < \ell < N} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} = 1 + O\left(\sum_{L < \ell < N} \frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N}\right).$$

Hence it suffices to show that $\sum_{L < \ell < N} \frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N} = O\left(\frac{1}{L^3}\right)$. The latter sum is split up into four parts $\sum_{L < \ell \leq M} + \sum_{M < \ell \leq \frac{N}{2}} + \sum_{\frac{N}{2} < \ell < N-M} + \sum_{N-M \leq \ell < N}$. We argue as in the proof of Proposition A.1. In particular there exists $C \geq 1$, independent of n , so that

$$\frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N} \leq C \frac{\gamma_\ell^- + \frac{M^2}{N}}{\ell^2 - L^2}, \quad L < \ell \leq M, \quad \frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N} \leq \frac{C}{N^2} \left(\gamma_{N-\ell}^+ + \frac{M^2}{N}\right), \quad N-M \leq \ell < N.$$

By Lemma C.7 and as $(\ell^2 \gamma_\ell^-)_{\ell \geq 1} \in \ell^2$ due to $\alpha, \beta \in C_0^2(\mathbb{T})$ one then has

$$\begin{aligned} \sum_{L < \ell \leq M} \frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N} &\leq C \left(\sum_{L < \ell \leq M} \frac{\ell^2 \gamma_\ell^-}{L^3(\ell - L)} + \sum_{L < \ell \leq M} \frac{M^2}{N} \frac{1}{L} \frac{1}{\ell - L} \right) \\ &= O\left(\frac{1}{L^3} + \frac{M^2}{N} \frac{\log L}{L}\right) = O\left(\frac{1}{L^3}\right) \end{aligned}$$

and as by Theorem 2.1, $\frac{1}{\tau_\ell^N - \tau_n^N} = O(1)$ for any $N - M \leq \ell < N$ and $\eta \leq 1/3$,

$$\sum_{N-M \leq \ell < N} \frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N} = O\left(\frac{1}{N^2}\right).$$

Finally, for $M < \ell \leq \frac{N}{2}$ one has $\gamma_\ell^N = O(N^{-2}M^{-1})$ and

$$\tau_\ell^N - \tau_n^N \geq \tau_\ell^N - \tau_M^N = -2 \cos \frac{\ell\pi}{N} + 2 \cos \frac{M\pi}{N} + O(N^{-2}M^{-1}).$$

Using $-\cos \frac{\ell\pi}{N} + \cos \frac{M\pi}{N} \geq \pi \frac{\ell^2 - M^2}{N^2}$ (cf (C.2)) and Lemma C.6 one sees that

$$\sum_{M < \ell \leq \frac{N}{2}} \frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N} = O\left(\frac{1}{N^2 M} \sum_{\ell > M} \frac{N^2}{\ell^2 - M^2}\right) = O\left(\frac{\log M}{M^2}\right).$$

Note that $\frac{\log M}{M^2} = O\left(\frac{1}{L^3}\right)$. Finally, as $(\tau_\ell^N - \tau_n^N)^{-1} = O(1)$ for $\frac{N}{2} < \ell < N - M$, one has $\sum_{\frac{N}{2} < \ell < N - M} \frac{\gamma_\ell^N}{\tau_\ell^N - \tau_n^N} = O\left(\frac{1}{NM}\right)$. Going through the arguments of the proof one verifies the claimed uniformity statement. \square

In a straightforward way one estimates $Q_k^{-,L}(\lambda) = \prod_{\ell > L} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^{-}(\lambda)}$, defined by (5.15) and obtains the following result.

Lemma B.5 *Uniformly for $1 \leq n, k \leq L$, $\lambda \in \Gamma_n^-$, and uniformly on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$Q_k^{-,L}(\lambda) = 1 + O\left(\sum_{\ell > L} \frac{\gamma_\ell^-}{\ell^2 - L^2}\right) = 1 + O\left(\frac{1}{L^3}\right). \quad (\text{B.4})$$

Remark B.6 *Similarly, one shows that (B.4) holds for $\lambda \leq (\lambda_{2L}^- + \lambda_{2L+1}^-)/2$ uniformly in $1 \leq k \leq L$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.*

Using the lemmas above we now estimate $\mathfrak{G}_k^{N,L}(\lambda)$, introduced in (5.16),

$$\mathfrak{G}_k^{N,L}(\lambda) = \frac{Q_k^{-,L}(\lambda)}{Q_k^{N,L}(\mu)} \cdot \frac{2N \sqrt[3]{\mu - \lambda_0^N}}{\sqrt[3]{\lambda - \lambda_0^-}} \cdot \frac{4N^2 w_k^N(\mu)}{w_k^-(\lambda)}$$

where $\mu \equiv \mu(\lambda) = -2 + \frac{\lambda}{4N^2}$.

Lemma B.7 *Uniformly for $1 \leq n, k \leq L$ with $n \neq k$, $\lambda \in \Gamma_n^-$, and uniformly on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$, one has*

$$\mathfrak{G}_k^{N,L}(\lambda) = 1 + O\left(\frac{1}{L^3}\right). \quad (\text{B.5})$$

Remark B.8 *Similarly, one shows that (B.5) holds for $\lambda \in [\lambda_1^-, \lambda_{2k-2}^-] \cup [\lambda_{2k+1}^-, \lambda_{2L}^-]$ uniformly in $1 \leq k \leq L$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$.*

Proof. By Theorem 2.1 one has, with $\nu_j^N = 4N^2(\lambda_j^N + 2)$,

$$\frac{4N^2 w_k^N(\mu)}{w_k^-(\lambda)} = \sqrt[+]{\frac{(\nu_{2k}^N - \lambda)(\nu_{2k-1}^N - \lambda)}{(\lambda_{2k}^- - \lambda)(\lambda_{2k-1}^- - \lambda)}} = 1 + O\left(\frac{M^2}{N} \frac{1}{|k^2 - n^2|}\right)$$

and

$$\frac{2N \sqrt[+]{\mu - \lambda_0^N}}{\sqrt[+]{\lambda - \lambda_0^-}} = \sqrt[+]{1 + \frac{\lambda_0^- - \nu_0^N}{\lambda - \lambda_0^-}} = 1 + O\left(\frac{M^2}{N} \frac{1}{n^2}\right).$$

The claimed estimate then follows from Lemma B.3, Lemma B.5, and (5.1). Going through the arguments of the proof one verifies the uniformity statement. \square

C Auxiliary lemmas

For the convenience of the reader, we collect in this appendix elementary lemmas which are used at various places of the paper. The first two lemmas concern uniform bounds for the Hölder continuity of some special functions.

Lemma C.1 (i) $0 \leq \sqrt{x} - \sqrt{y} \leq \sqrt{x-y} \quad \forall x \geq y \geq 0$;
(ii) $0 \leq \log x - \log y \leq x - y \quad \forall x \geq y \geq 1$;
(iii) $0 \leq \log(x + \sqrt{x^2 - 1}) - \log(y + \sqrt{y^2 - 1}) \leq 2\sqrt{x+y}\sqrt{x-y} \quad \forall x \geq y \geq 1$.

Remark C.2 *Note that for $x \geq 1$, the principal branch of $\arccos(x)$ is given by $\operatorname{arcosh}(x) = \log(x + \sqrt{x^2 - 1})$.*

Proof: (i) The claimed estimate is obtained from combining the following inequalities, valid for any $x \geq y \geq 0$,

$$x - y = (\sqrt{x} - \sqrt{y})(\sqrt{x} + \sqrt{y}), \quad x - y \leq \sqrt{x-y}\sqrt{x+y}, \quad \sqrt{x+y} \leq \sqrt{x} + \sqrt{y}.$$

(ii) For any $x \geq y \geq 1$ one has

$$0 \leq \log x - \log y = \int_0^1 \partial_t \log(y + t(x - y)) dt \leq x - y.$$

(iii) By applying first (ii) and then (i) one sees that for any $x \geq y \geq 1$,

$$\log(x + \sqrt{x^2 - 1}) - \log(y + \sqrt{y^2 - 1}) \leq (x - y) + \sqrt{x^2 - y^2}.$$

As $\sqrt{x^2 - y^2} = \sqrt{x - y}\sqrt{x + y}$, the claimed estimate follows. \square

Lemma C.3 For any $a, b > 0$ and $x, y \geq 0$,

$$\left| \frac{a}{\sqrt{a+x}} - \frac{b}{\sqrt{b+y}} \right| \leq |a - b|^{1/2} + |x - y|^{1/2}.$$

By a limiting argument, the inequality continues to hold for $a, b \geq 0$.

Proof: For any $a, b > 0$ and $x, y \geq 0$ one has

$$\begin{aligned} \left| \frac{a}{\sqrt{a+x}} - \frac{b}{\sqrt{b+y}} \right|^2 &\leq \left| \left(\frac{a}{\sqrt{a+x}} - \frac{b}{\sqrt{b+y}} \right) \left(\frac{a}{\sqrt{a+x}} + \frac{b}{\sqrt{b+y}} \right) \right|^2 \\ &= \left| \frac{a^2}{a+x} - \frac{b^2}{b+y} \right|^2 = \left| \frac{a^2b - ab^2 + a^2y - b^2x}{ab + xb + ay + xy} \right|^2 \end{aligned}$$

As $a^2b - ab^2 + a^2y - b^2x = (ab + xb + ay)(a - b) + ab(y - x)$ one gets $\left| \frac{a}{\sqrt{a+x}} - \frac{b}{\sqrt{b+y}} \right|^2 \leq |a - b| + |y - x|$ or

$$\left| \frac{a}{\sqrt{a+x}} - \frac{b}{\sqrt{b+y}} \right| \leq |a - b|^{1/2} + |y - x|^{1/2}.$$

\square

Lemma C.4 For any $M < n, k < N - M, k \neq n$

$$\left| \cos \frac{n\pi}{N} - \cos \frac{k\pi}{N} \right| \geq \frac{M}{N} \cdot \frac{\pi}{N} \quad (\text{C.1})$$

whereas for any $0 < \ell \leq \frac{N}{2}$ and $0 \leq j < \ell$

$$\cos \frac{j\pi}{N} - \cos \frac{\ell\pi}{N} \geq \frac{(\ell^2 - j^2)\pi}{N^2}. \quad (\text{C.2})$$

Furthermore, for any $M < n < N - M$ one has

$$\sum_{\substack{k \neq n \\ M < k < N - M}} \frac{1}{\left| \cos \frac{n\pi}{N} - \cos \frac{k\pi}{N} \right|} = O\left(N^2 \frac{\log M}{M}\right) \quad (\text{C.3})$$

Proof. The difference $\cos \frac{n\pi}{N} - \cos \frac{k\pi}{N}$ is bounded in absolute value from below by $\min_{\pm} |\cos \frac{n\pi}{N} - \cos \frac{(n\pm 1)\pi}{N}|$. Writing $|\cos x - \cos y| = |\int_x^y \sin t dt|$ and using that

$$\sin x \geq \sin \frac{M\pi}{N} \geq \frac{2}{\pi} \cdot \frac{M\pi}{N} \quad \forall x \in [\frac{M\pi}{N}, \pi - \frac{M\pi}{N}]$$

(C.1) follows. Next, for any $0 < \ell \leq \frac{N}{2}$ and $0 \leq j < \ell$ one has

$$\cos \frac{j\pi}{N} - \cos \frac{\ell\pi}{N} = \int_{\frac{j\pi}{N}}^{\frac{\ell\pi}{N}} \sin x dx \geq \frac{x^2}{\pi} \Big|_{\frac{j\pi}{N}}^{\frac{\ell\pi}{N}} = \frac{(\ell^2 - j^2)\pi}{N^2}$$

which proves (C.2). Concerning (C.3) let us concentrate on the case where $M < n \leq \frac{N}{2}$. (The case where $\frac{N}{2} \leq n < N - M$ is treated in the same way.) Then one sees in a straightforward way that $\sum_{M < k < N - M, k \neq n} \cdots \leq 3 \sum_{M < k \leq N/2, k \neq n} \cdots$. Hence (C.3) follows from (C.2) and Lemma C.7. \square

Lemma C.5 *The following estimate holds*

$$\sum_{M < k \leq \frac{N}{2}} \frac{\sin \frac{k\pi}{N}}{\cos \frac{M\pi}{N} - \cos \frac{k\pi}{N}} = O(N).$$

Proof. Note that

$$\sum_{M < k \leq \frac{N}{2}} \frac{\sin \frac{k\pi}{N}}{\cos \frac{M\pi}{N} - \cos \frac{k\pi}{N}} \leq \frac{\sin \frac{(M+1)\pi}{N}}{\cos \frac{M\pi}{N} - \cos \frac{(M+1)\pi}{N}} + \int_{\frac{(M+1)\pi}{N}}^{\frac{\pi}{2}} \frac{\sin x}{\cos \frac{M\pi}{N} - \cos x} dx$$

and

$$\begin{aligned} \int_{\frac{(M+1)\pi}{N}}^{\pi/2} \frac{\sin x}{\cos \frac{M\pi}{N} - \cos x} dx &= \log \left(\cos \frac{M\pi}{N} - \cos x \right) \Big|_{\frac{(M+1)\pi}{N}}^{\pi/2} \\ &\leq -\log \left(\cos \frac{M\pi}{N} - \cos \frac{(M+1)\pi}{N} \right). \end{aligned}$$

By (C.2) one has $\cos \frac{M\pi}{N} - \cos \frac{(M+1)\pi}{N} \geq \frac{(2M+1)\pi}{N^2}$ and hence

$$-\log \left(\cos \frac{M\pi}{N} - \cos \frac{(M+1)\pi}{N} \right) \leq 2 \log N.$$

Furthermore, together with the estimate $0 \leq \sin \frac{(M+1)\pi}{N} \leq \frac{(M+1)\pi}{N}$ one gets

$$\frac{\sin((M+1)\pi/N)}{\cos \frac{M\pi}{N} - \cos \frac{(M+1)\pi}{N}} \leq N.$$

Combining the two latter estimates the claim follows. \square

Lemma C.6 (Mean value theorem in integral form) Let f be continuous and $g \geq 0$ integrable on the interval $[a, b]$ with $a < b$. Then there exists $a < y < b$ so that $\int_a^b f(x)g(x)dx = f(y) \int_a^b g(x)dx$.

Lemma C.7 For any $k \in \mathbb{Z}_{\geq 1}$ and any $1 \leq n < k$, $\sum_{0 \leq j \leq n} \frac{1}{k^2 - j^2} = O\left(\frac{\log(n+1)}{k}\right)$. Similarly, for any $n > k$, $\sum_{j \geq n} \frac{1}{j^2 - k^2} = O\left(\frac{\log(k+1)}{k}\right)$.

Proof. Note that for any $n > k$, $\sum_{j \geq n} \frac{1}{j^2 - k^2} \leq \frac{1}{2k} + \int_n^\infty \frac{1}{x^2 - k^2} dx$. Using that $\frac{1}{x^2 - k^2} = \left(\frac{1}{x-k} - \frac{1}{x+k}\right) \frac{1}{2k} = \frac{1}{2k} \frac{d}{dx} \log \frac{x-k}{x+k} = \frac{1}{2k} \frac{d}{dx} \log\left(1 - \frac{2k}{x+k}\right)$ one gets

$$\begin{aligned} \int_n^N \frac{1}{x^2 - k^2} dx &= \frac{1}{2k} \log\left(1 - \frac{2k}{N+k}\right) - \frac{1}{2k} \log\left(1 - \frac{2k}{n+k}\right) \\ &= \frac{1}{2k} \log\left(1 - \frac{2k}{N+k}\right) + \frac{1}{2k} \log\left(1 + \frac{2k}{n-k}\right). \end{aligned}$$

Hence $\lim_{N \rightarrow \infty} \int_n^N \frac{1}{x^2 - k^2} dx = \frac{1}{2k} \log\left(1 + \frac{2k}{n-k}\right) \leq \frac{1}{2k} \log(1 + 2k)$. The sum $\sum_{0 \leq j \leq n} \frac{1}{k^2 - j^2}$ can be estimated in a similar fashion. \square

D Auxiliary estimates

In this appendix we prove auxiliary estimates needed in Section 7 to derive the claimed asymptotics of ω_n^N for $1 \leq n \leq L$. In view of (6.5), (6.6), and the assumptions about F , formula (6.4) reads in the case $1 \leq n \leq L$ and $M_0 = L$ as $\omega_n^N = \frac{2\pi n}{N} - \frac{\pi}{4N^3} \sum_{j=1}^n \lambda_{2j-2}^- - (P1) - (P2) + Error_L$ where

$$(P2) = \frac{1}{N} \sum_{\substack{1 \leq k \leq L \\ I_k^N \neq 0}} I_k^N \omega_k^N \int_{\mathcal{I}_n^N} \frac{\varphi_k^N(\mu)}{i \sqrt{\chi_N(\mu)}} d\mu, \quad (D.1)$$

$$Error_L = O\left(\frac{1}{N^3} \left(\frac{\log L}{L^5} + \frac{nM^2}{N} + \frac{n}{M^3}\right)\right).$$

We will need estimates for the integral in (D.1), relating it to a corresponding one of KdV. The integral in (D.1) is a sum of convergent improper Riemann integrals on intervals at the endpoints of which the integrand might be singular. More specifically, the improper Riemann integrals $\int_{\lambda_{2k-2}^N}^{\lambda_{2k-1}^N} \frac{\varphi_k^N(\mu)}{i \sqrt{\chi_N(\mu)}} d\mu$

and $\int_{\lambda_{2k}^N}^{\lambda_{2k+1}^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu$ can be shown to be of the order $O(\frac{N}{k} \log \gamma_k^N)$ and to diverge if $I_k^N = 0$. On the other hand, $i^c \sqrt{\chi_N(\mu)}$ is real valued on both intervals $[\lambda_{2k-2}^N, \lambda_{2k-1}^N]$ and $[\lambda_{2k}^N, \lambda_{2k+1}^N]$ and has opposite signs which will lead to cancelations. Hence a careful analysis is needed to obtain estimates which are uniform on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$. We split the domain \mathcal{I}_n^N of the integral in (D.1) as follows. Fix $\rho > 0$ in such a way that $\lambda_{2j}^- + 3\rho < \lambda_{2j+1}^- - 3\rho$ for any $j \geq 0$ (cf (B.2)), set $\rho^N = \frac{\rho}{4N^2}$ and introduce

$$\begin{aligned} A_j^N &:= [\lambda_{2j-1}^N - \rho^N, \lambda_{2j-1}^N] \cup [\lambda_{2j}^N, \lambda_{2j}^N + \rho^N], \quad 0 < j < n \\ B_j^N &:= [\lambda_{2j}^N + \rho^N, \lambda_{2j+1}^N - \rho^N], \quad 0 < j < n \\ A_0^N &:= [\lambda_0^N, \lambda_1^N - \rho^N], \quad C_n^N := [\lambda_{2n-1}^N - \rho^N, \lambda_{2n-1}^N] \end{aligned}$$

so that $\mathcal{I}_n^N = (\cup_{j=0}^{n-1} A_j^N) \cup (\cup_{j=1}^{n-1} B_j^N) \cup C_n^N$. Correspondingly define

$$\begin{aligned} A_j^- &:= [\lambda_{2j-1}^- - \rho, \lambda_{2j-1}^-] \cup [\lambda_{2j}^-, \lambda_{2j}^- + \rho], \quad , \quad 0 < j < n \\ B_j^- &:= [\lambda_{2j}^- + \rho, \lambda_{2j+1}^- - \rho], \quad 0 < j < n \\ A_0^- &:= [\lambda_0^-, \lambda_1^- - \rho], \quad C_n^- := [\lambda_{2n-1}^- - \rho, \lambda_{2n-1}^-] \end{aligned}$$

and set $\mathcal{I}_n^- = (\cup_{j=0}^{n-1} A_j^-) \cup (\cup_{j=1}^{n-1} B_j^-) \cup C_n^-$. In the lemmas below we will derive the asymptotics of the integrals over the intervals $(A_j^N)_{0 \leq j < n}$, $(B_j^N)_{0 < j < n}$, and C_n^N in terms of corresponding integrals for KdV over the intervals $(A_j^-)_{0 \leq j < n}$, $(B_j^-)_{0 < j < n}$, and C_n^- . We begin with the analysis of the integral over the interval $A_0^M = [\lambda_0^N, \lambda_1^N - \rho^N]$.

Lemma D.1 *Uniformly for $1 \leq k \leq L$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\int_{A_0^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = \frac{N}{2\pi k} \int_{A_0^-} \frac{\psi_k^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\frac{N}{k^2} \frac{1}{L^{5/2}}\right). \quad (\text{D.2})$$

Proof. For $\mu \in A_0^N$, and $1 \leq k \leq L$ one has by (5.2), (5.3), (5.14), (5.15)

$$\frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} = \frac{1}{2} \left(\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} \right) \frac{1}{\sqrt{\mu - \lambda_0^N}} \frac{1}{w_k^N(\mu)} Q_k^{N,L}(\mu).$$

Here we used that $\varphi_k^N(\mu) = (-1)^{N-2} \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} (\sigma_\ell^{N,k} - \mu)$ and that by (5.4) $(-1)^N i \sqrt[5]{\chi_N(\mu)} > 0$ for any $\mu \in (\lambda_0^N, \lambda_1^N)$ to get the correct sign. Setting $\mu \equiv \mu(\lambda) = -2 + \frac{\lambda}{4N^2}$ and writing $\frac{1}{w_k^N(\mu)} \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)}$ as

$$\frac{4N^2}{w_k^-(\lambda)} \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} \cdot \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{4N^2(\sigma_\ell^{N,k} - \mu)}{\sigma_\ell^{-,k} - \lambda} \cdot \prod_{1 \leq \ell \leq L} \frac{w_\ell^-(\lambda)}{4N^2 w_\ell^N(\mu)}$$

one argues as in the proof of (5.17) to conclude that for $\lambda \leq \lambda_1^- - \rho/2$, $\prod_{1 \leq \ell \leq L} \frac{w_\ell^-(\lambda)}{4N^2 w_\ell^N(\mu)} = 1 + O(\frac{M^2}{N})$ and, taking into account also Theorem 5.1, $\prod_{\substack{1 \leq \ell \leq L \\ \ell \neq k}} \frac{4N^2(\sigma_\ell^{N,k} - \mu)}{\sigma_\ell^{-,k} - \lambda} = 1 + O(\frac{1}{L^{5/2}})$. By Lemma B.3 (cf Remark B.4) and Lemma B.5 (cf Remark B.6), $Q_k^{N,L}(\mu) = 1 + O(\frac{1}{L^3})$ respectively $Q_k^{-,L}(\lambda)^{-1} = 1 + O(\frac{1}{L^3})$. By (5.1) one then gets uniformly for $\lambda \leq \lambda_1^- - \rho/2$,

$$\frac{\varphi_k^N(\mu)}{i \sqrt[5]{\chi_N(\mu)}} = \frac{1}{w_k^-(\lambda)} \left(\prod_{\substack{0 < \ell < \infty \\ \ell \neq k}} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} \right) \frac{2N^2}{\sqrt[5]{\mu - \lambda_0^N}} \left(1 + O\left(\frac{1}{L^{5/2}}\right) \right).$$

Note that uniformly for $\lambda \in A_0^-$, $\prod_{\substack{0 < \ell < \infty \\ \ell \neq k}} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} = O(1)$ and, by the asymptotics of the periodic eigenvalues of H_- , $\tau_k^- = (\lambda_{2k-1}^- + \lambda_{2k}^-)/2 = 4k^2\pi^2 + O(1)$, whence $w_k^-(\lambda)^{-1} = O(k^{-2})$. Thus for $\lambda \leq \lambda_1^- - \rho/2$,

$$f_0^k(\lambda) := \frac{1}{w_k^-(\lambda)} \prod_{\substack{0 < \ell < \infty \\ \ell \neq k}} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} = O(k^{-2}). \quad (\text{D.3})$$

To prove the claimed asymptotics of the integral $\int_{A_0^N} \frac{\varphi_k^N(\mu)}{i \sqrt[5]{\chi_N(\mu)}} d\mu$ introduce the change of variable $[\lambda_0^-, \lambda_1^- - \rho] \rightarrow [\lambda_0^N, \lambda_1^N - \rho^N]$, given by

$$x \mapsto \mu = \lambda_0^N + \frac{h_0^N}{4N^2}(x - \lambda_0^-), \quad h_0^N := 4N^2 \frac{\lambda_1^N - \rho^N - \lambda_0^N}{\lambda_1^- - \rho - \lambda_0^-} > 0. \quad (\text{D.4})$$

Then $h_0^N = 1 + O(\frac{M^2}{N})$ (cf Theorem 2.1) and for any $x \in [\lambda_0^-, \lambda_1^- - \rho]$,

$$\lambda(x) := (\mu(x) + 2)4N^2 = x + O\left(\frac{M^2}{N}\right), \quad \frac{d\mu}{\sqrt[5]{\mu - \lambda_0^N}} = \frac{\sqrt[5]{h_0^N}}{2N} \frac{dx}{\sqrt[5]{x - \lambda_0^-}}.$$

Using (D.3) and (5.1) we get

$$\int_{A_0^N} \frac{\varphi_k^N(\mu)}{i \sqrt[5]{\chi_N(\mu)}} d\mu = N \int_{\lambda_0^-}^{\lambda_1^- - \rho} f_0^k(\lambda(x)) \frac{\sqrt[5]{h_0}}{\sqrt[5]{x - \lambda_0^-}} dx + O\left(\frac{N}{k^2} \frac{1}{L^{5/2}}\right). \quad (\text{D.5})$$

As by (5.7) and (5.8), $\frac{1}{2\pi k} \frac{\psi_k^-(\lambda)}{i \sqrt[5]{\Delta_-^2(\lambda) - 4}} = \frac{1}{\sqrt[5]{\lambda - \lambda_0^-}} f_0^k(\lambda)$ and as $\lambda(x) = h_0^N x + c_N$ with $c_N = O(\frac{M^2}{N})$ the stated asymptotics then follow by using once more (D.3). Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

For the proof of Theorem 1.1 we need bounds for $\int_{A_0^N} \frac{\varphi_k^N(\mu)}{i \sqrt[5]{\chi_N(\mu)}} d\mu$. In particular, in the proof of Proposition 7.1 we need a bound for these integrals also for k in the range $L < k \leq M$. By similar arguments as in the proof of Lemma D.1 one gets the following result.

Corollary D.2 *Uniformly for $1 \leq k \leq M$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\int_{A_0^N} \frac{\varphi_k^N(\mu)}{i \sqrt[5]{\chi_N(\mu)}} d\mu = O\left(\frac{N}{k^2}\right). \quad (\text{D.6})$$

Proof. Going through the arguments of the proof of Lemma D.1 one sees that $\prod_{1 \leq \ell \leq M} \frac{w_\ell^-(\lambda)}{4N^2 w_\ell^N(\mu)} = 1 + O(\frac{M^2}{N})$ and $\prod_{\substack{1 \leq \ell \leq M \\ \ell \neq k}} \frac{4N^2(\sigma_\ell^{N,k} - \mu)}{\sigma_\ell^{-,k} - \lambda} = O(1)$ (Theorem 2.1, $\sigma_\ell^N \in [\lambda_{2\ell-1}^N, \lambda_{2\ell}^N]$). As $Q_k^{N,M}(\mu)$, $Q_k^{-,M}(\lambda)^{-1} = O(1)$ the claimed estimate then follows from (D.3). \square

Let us now analyze the integrals over A_j^N in the case $0 < j < n$, $j \neq k$.

Lemma D.3 *Uniformly for $1 \leq n, k \leq L$, $0 < j < n$ with $j \neq k$, and on bounded sets of functions α, β in $C_0^2(\mathbb{T})$, $\int_{A_j^N} \frac{\varphi_k^N(\mu)}{i \sqrt[5]{\chi_N(\mu)}} d\mu$ has the asymptotics*

$$\frac{N}{2\pi k} \int_{A_j^-} \frac{\psi_k^-(\lambda)}{i \sqrt[5]{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\left(\frac{M}{N^{1/2}} + \frac{1}{L^{5/2}}\right) \frac{N}{j(k^2 - j^2)}\right). \quad (\text{D.7})$$

Proof. We present a proof which can be easily adapted to the case where $j = k$. Similarly as in the proof of Lemma D.1, write for $\mu \in A_j^N$

$$\frac{\varphi_k^N(\mu)}{i \sqrt[5]{\chi_N(\mu)}} = \frac{1}{2 \sqrt[5]{\mu - \lambda_0^N}} \left(\frac{\sigma_j^{N,k} - \mu}{w_k^N(\mu)} \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq j,k}} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} \right) Q_k^{N,L}(\mu) \frac{1}{w_j^N(\mu)}.$$

Set $\mu \equiv \mu(\lambda) = -2 + \frac{\lambda}{4N^2}$ and write $\frac{\sigma_j^{N,k} - \mu}{w_k^N(\mu)} \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq j,k}} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)}$ as

$$\frac{4N^2(\sigma_j^{N,k} - \mu)}{w_k^-(\lambda)} \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq j,k}} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} \cdot \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq j,k}} \frac{4N^2(\sigma_\ell^{N,k} - \mu)}{\sigma_\ell^{-,k} - \lambda} \cdot \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq j}} \frac{w_\ell^-(\lambda)}{4N^2 w_\ell^N(\mu)}.$$

Arguing as in the proof of Lemma D.1 and taking into account that by Theorem 2.1, $4N^2(\mu - \lambda_0^N) = (\lambda - \lambda_0^-)(1 + O(\frac{1}{j^2} \frac{M^2}{N}))$, and hence $\frac{1}{2 \sqrt[3]{\mu - \lambda_0^N}} = \frac{N}{\sqrt[3]{\lambda - \lambda_0^-}}(1 + O(\frac{1}{j^2} \frac{M^2}{N}))$, one sees that for $\lambda \in \tilde{A}_j^- := [\lambda_{2j-1}^- - 2\rho, \lambda_{2j}^- + 2\rho]$

$$\frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} = f_j^{N,k}(\lambda) \frac{N}{w_j^N(\mu)} (1 + g_j^{N,k}(\lambda)),$$

where

$$f_j^{N,k}(\lambda) = \frac{1}{\sqrt[3]{\lambda - \lambda_0^-}} \frac{4N^2(\sigma_j^{N,k} - \mu)}{w_k^-(\lambda)} \prod_{\substack{0 < \ell < \infty \\ \ell \neq j,k}} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} \quad (\text{D.8})$$

and $g_j^{N,k}$ is defined by the equation above and satisfies $g_j^{N,k}(\lambda) = O(\frac{1}{L^{5/2}})$. Actually, $f_j^{N,k}$ and $g_j^{N,k}$ are defined and analytic on the rectangle $R_j \subset \mathbb{C}$ with top and bottom side given by $\tilde{A}_j^- \pm i2\rho$. Note that for $\lambda \in R_j$, one has $4N^2(\sigma_j^{N,k} - \mu) = O(1)$ (Theorem 2.1, $\sigma_j^{N,k} \in [\lambda_{2j-1}^N, \lambda_{2j}^N]$), $1/\sqrt[3]{\lambda - \lambda_0^-} = O(\frac{1}{j})$, $1/w_k^-(\lambda) = O(\frac{1}{k^2 - j^2})$, and $\prod_{\substack{0 < \ell < \infty \\ \ell \neq j,k}} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} = O(1)$. Hence for $\lambda \in R_j$

$$f_j^{N,k}(\lambda) = O\left(\frac{1}{j(k^2 - j^2)}\right). \quad (\text{D.9})$$

Moreover, by the same arguments as above, one sees that $g_j^{N,k}(\lambda) = O(\frac{1}{L^{5/2}})$ is valid on R_j . To proceed further observe that, by an explicit computation,

$$\int_{A_j^N} \frac{d\mu}{w_j^N(\mu)} = 0,$$

so that the integral $\frac{1}{N} \int_{A_j^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu$ equals

$$\int_{A_j^N} \frac{f_j^{N,k}(\lambda(\mu))(1 + g_j^{N,k}(\lambda(\mu))) - f_j^{N,k}(\lambda(\tau_j^N))(1 + g_j^{N,k}(\lambda(\tau_j^N)))}{w_j^N(\mu)} d\mu.$$

Now the integrals over the intervals $[\lambda_{2j-1}^N - \rho^N, \lambda_{2j-1}^N]$ and $[\lambda_{2j}^N, \lambda_{2j}^N + \rho^N]$ of A_j^N can be analyzed separately. As the two cases are treated in the same way we concentrate on the integral (T) over the interval $[\lambda_{2j}^N, \lambda_{2j}^N + \rho^N]$ only. Write $(T) = (T1) - (T2) - (T3)$ where with $\lambda \equiv \lambda(\mu) = 4N^2(\mu + 2)$, $\tilde{\tau}_j^N = 4N^2(\tau_j^N + 2)$, and the identity $w_j^N(\mu) = -\frac{1}{2N} \sqrt[4]{\lambda - \nu_{2j-1}^N} \sqrt[4]{\mu - \lambda_{2j}^N}$ on $[\lambda_{2j}^N, \lambda_{2j}^N + \rho^N]$,

$$(T1) := \int_{\lambda_{2j}^N}^{\lambda_{2j}^N + \rho^N} \frac{f_j^{N,k}(\lambda) - f_j^{N,k}(\tilde{\tau}_j^N)}{w_j^N(\mu)} d\mu, \quad (D.10)$$

$$(T2) := \int_{\lambda_{2j}^N}^{\lambda_{2j}^N + \rho^N} f_j^{N,k}(\lambda) \frac{g_j^{N,k}(\lambda) - g_j^{N,k}(\tilde{\tau}_j^N)}{\sqrt[4]{\lambda - \nu_{2j-1}^N}} \frac{2Nd\mu}{\sqrt[4]{\mu - \lambda_{2j}^N}},$$

$$(T3) := g_j^{N,k}(\tilde{\tau}_j^N) \int_{\lambda_{2j}^N}^{\lambda_{2j}^N + \rho^N} \frac{f_j^{N,k}(\lambda) - f_j^{N,k}(\tilde{\tau}_j^N)}{\sqrt[4]{\lambda - \nu_{2j-1}^N}} \frac{2Nd\mu}{\sqrt[4]{\mu - \lambda_{2j}^N}}.$$

By Lemma C.6 there exists $\lambda_{2j}^N < \mu_* < \lambda_{2j-1}^N$ so that with $\lambda_* = \lambda(\mu_*)$ $(T2) = f_j^{N,k}(\lambda_*) \frac{g_j^{N,k}(\lambda_*) - g_j^{N,k}(\tilde{\tau}_j^N)}{\sqrt[4]{\lambda_* - \nu_{2j-1}^N}} \frac{\sqrt[4]{\rho}}{2}$. As $f_j^{N,k}(\lambda_*) = O(\frac{1}{j(k^2 - j^2)})$ (use (D.9)) and by Cauchy's theorem, $g_j^{N,k}(\lambda_*) - g_j^{N,k}(\tilde{\tau}_j^N) = O(L^{-5/2}(\lambda_* - \tilde{\tau}_j^N))$ (use $g_j^{N,k}(\lambda) = O(L^{-5/2})$) it then follows from $\nu_{2j-1}^N \leq \tilde{\tau}_j^N$ that

$$(T2) = O\left(\frac{1}{j(j^2 - k^2)} \frac{1}{L^{5/2}}\right).$$

The term $(T3)$ is studied in a similar way and admits the same bound. Similarly, write $\int_{A_j^-} \frac{\psi_k^-(\lambda)}{\sqrt{\Delta_-^2(\lambda) - 4}} d\lambda = 2\pi k \int_{A_j^-} \frac{f_j^{-,k}(\lambda)}{w_j^-(\lambda)} d\lambda$ where in view of (5.7), $f_j^{-,k}$ is given by (D.8) with $4N^2(\sigma_j^{N,k} - \mu)$ replaced by $(\sigma_j^{-,k} - \lambda)$. As $\int_{A_j^-} \frac{1}{w_j^-(\lambda)} d\lambda = 0$

$$\frac{1}{2\pi k} \int_{A_j^-} \frac{\psi_k^-(\lambda)}{\sqrt{\Delta_-^2(\lambda) - 4}} d\lambda = \int_{A_j^-} \frac{f_j^{-,k}(\lambda) - f_j^{-,k}(\tau_j^-)}{w_j^-(\lambda)} d\lambda,$$

allowing again to consider the integrals over the intervals $[\lambda_{2j-1}^- - \rho, \lambda_{2j-1}^-]$ and $[\lambda_{2j}^-, \lambda_{2j}^- + \rho]$ of A_j^- separately. It then remains to compare $(T1)$ with

$$(S1) := \int_{\lambda_{2j}^-}^{\lambda_{2j}^- + \rho} \frac{f_j^{-,k}(\lambda) - f_j^{-,k}(\tau_j^-)}{w_j^-(\lambda)} d\lambda. \quad (D.11)$$

Make the change of variables $[0, \rho] \rightarrow [\lambda_{2j}^N, \lambda_{2j}^N + \rho^N]$, $x \mapsto \mu(x) = \lambda_{2j}^N + 4N^2x$ in (T1) and $[0, \rho] \rightarrow [\lambda_{2j}^-, \lambda_{2j}^- + \rho]$, $x \mapsto \lambda = \lambda_{2j}^- + x$ in (S1) so that with $\tilde{\gamma}_j^N = 4N^2\gamma_j^N$, $\tilde{\tau}_j^N = 4N^2(\tau_j^N + 2)$, and the sign of the standard root

$$(T1) = - \int_0^\rho \frac{f_j^{N,k}(x + \nu_{2j}^N) - f_j^{N,k}(\tilde{\tau}_j^N)}{\sqrt[4]{x} \sqrt[4]{x + \tilde{\gamma}_j^N}} dx = - \int_0^\rho \frac{F_j^{N,k}(x) \cdot (x + \tilde{\gamma}_j^N/2)}{\sqrt[4]{x + \tilde{\gamma}_j^N}} \frac{dx}{\sqrt[4]{x}}$$

where

$$F_j^{N,k}(x) = \int_0^1 \partial_\lambda f_j^{N,k}(\tilde{\tau}_j^N + t(x + \tilde{\gamma}_j^N/2)) dt.$$

Similarly, with $F_j^{-,k}(x) = \int_0^1 \partial_\lambda f_j^{-,k}(\tau_j^- + t(x + \gamma_j^-/2)) dt$, one has

$$(S1) = - \int_0^\rho \frac{F_j^{-,k}(x) \cdot (x + \gamma_j^-/2)}{\sqrt[4]{x + \gamma_j^-}} \frac{dy}{\sqrt[4]{x}}. \quad (D.12)$$

Thus, by Lemma C.6, there exists $0 < z < \rho$ so that

$$\begin{aligned} ((S1) - (T1)) \frac{2}{\sqrt[4]{\rho}} &= F_j^{N,k}(z) \frac{z + \tilde{\gamma}_j^N/2}{\sqrt[4]{z + \tilde{\gamma}_j^N}} - F_j^{-,k}(z) \frac{z + \gamma_j^-/2}{\sqrt[4]{z + \gamma_j^-}} \\ &= \left(F_j^{N,k}(z) - F_j^{-,k}(z) \right) \frac{z + \tilde{\gamma}_j^N/2}{\sqrt[4]{z + \tilde{\gamma}_j^N}} + F_j^{-,k}(z) \left(\frac{z + \tilde{\gamma}_j^N/2}{\sqrt[4]{z + \tilde{\gamma}_j^N}} - \frac{z + \gamma_j^-/2}{\sqrt[4]{z + \gamma_j^-}} \right). \end{aligned}$$

As $F_j^{-,k} = O(\frac{1}{j(j^2 - k^2)})$ (by Cauchy's estimate and (D.9)), $\tilde{\gamma}_j^N - \gamma_j^- = O(\frac{M^2}{N})$ (by Theorem 2.1), Lemma C.3 implies that

$$F_j^{-,k}(z) \left(\frac{z + \tilde{\gamma}_j^N/2}{\sqrt[4]{z + \tilde{\gamma}_j^N}} - \frac{z + \gamma_j^-/2}{\sqrt[4]{z + \gamma_j^-}} \right) = O\left(\frac{1}{j(j^2 - k^2)} \frac{M}{N^{1/2}} \right).$$

Using that by (D.9) and Theorem 5.1

$$f_j^{N,k}(x + \lambda_{2j}^-) - f_j^{-,k}(x + \lambda_{2j}^-) = O\left(\frac{\tilde{\sigma}_j^{N,k} - \sigma_j^{-,k}}{j(k^2 - j^2)} \right) = O\left(\frac{1}{L^{5/2}} \frac{1}{j(k^2 - j^2)} \right),$$

and using that $f_j^{N,k}$ and $f_j^{-,k}$ are analytic, hence in particular Lipschitz, one concludes by Cauchy's estimate and by the boundedness of $\frac{x + \tilde{\gamma}_j^N/2}{\sqrt[4]{x + \tilde{\gamma}_j^N}}$ that

$$\left(F_j^{N,k}(z) - F_j^{-,k}(z) \right) \frac{z + \tilde{\gamma}_j^N/2}{\sqrt[4]{z + \tilde{\gamma}_j^N}} = O\left(\frac{1}{j(k^2 - j^2)} \frac{1}{L^{5/2}} \right).$$

By combining the estimates obtained the stated asymptotics follow. Going through the arguments of the proof one verifies that the claimed uniformity statement holds. \square

Corollary D.4 *Uniformly for $1 \leq n, k \leq M$, $0 < j < n$ with $j \neq k$, and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\int_{A_j^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = O\left(\frac{N}{j(k^2 - j^2)}\right). \quad (\text{D.13})$$

Proof. One verifies in a straightforward way that uniformly for any $1 \leq n, k \leq M$ and $0 < j < n$ with $j \neq k$, (D.9) holds and the function $g_j^{N,k}$, introduced in the proof of Lemma D.3, satisfies $1 + g_j^{N,k}(\lambda(\mu(x))) = O(1)$. Arguing as in the proofs of Corollary D.2 and Lemma D.3 yields the claimed estimate. \square

Next we analyze the integral $\int_{A_j^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu$ in the case where $j = k$.

Lemma D.5 *Uniformly for $1 \leq n \leq L$, $0 < k < n$, and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\int_{A_k^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = \frac{N}{2\pi k} \int_{A_k^-} \frac{\psi_k^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\left(\frac{M}{N^{1/2}} + \frac{1}{L^{5/2}}\right) \frac{N}{k}\right). \quad (\text{D.14})$$

Proof. Proceeding as in the proof of Lemma D.3 one gets

$$\frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} = f_k^{N,k}(\lambda) \frac{N}{w_k^N(\mu)} \left(1 + O\left(\frac{1}{L^{5/2}}\right)\right)$$

where now

$$f_k^{N,k}(\lambda) = \left(\prod_{\substack{0 < \ell < \infty \\ \ell \neq k}} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} \right) \frac{1}{\sqrt{\lambda - \lambda_0^-}} \quad (\text{D.15})$$

which is analytic in the rectangle R_k , introduced in the proof of Lemma D.3. Since $\lambda \sim 4\pi^2 k^2$ for $\lambda \in R_k$, it satisfies

$$f_k^{N,k}(\lambda) = O\left(\frac{1}{k}\right). \quad (\text{D.16})$$

Arguing as in the proof of Lemma D.3, the claimed statements follow. \square

Corollary D.6 *Uniformly for $1 \leq n \leq M$, $1 < k < n$, and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\int_{A_k^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = O\left(\frac{N}{k}\right). \quad (\text{D.17})$$

Proof. Arguing as in the proofs of Corollary D.4 and Lemma D.5, the claimed statement follows. \square

Next we analyse the integral $\int_{C_n^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu$ in the case where $k \neq n$.

Lemma D.7 *Uniformly for $1 \leq n \leq L$, $0 < k < n$, and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\int_{C_n^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = \frac{N}{2\pi k} \int_{C_n^-} \frac{\psi_k^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\frac{N}{n(k^2 - n^2)} \left(\frac{M}{N^{1/2}} + \frac{1}{L^{5/2}}\right)\right). \quad (\text{D.18})$$

Proof. We proceed similarly as in the proof of Lemma D.3, but have to take into account that possibly, $\lambda_{2n-1}^N = \lambda_{2n}^N$. Fortunately, in such a case $\sigma_n^{N,k} - \mu = \tau_n^N - \mu = w^N(\mu)$, implying that $\frac{\sigma_n^{N,k} - \mu}{w_n^N(\mu)} \equiv 1$. Therefore write $\frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}}$ as

$$\frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} = \frac{1}{2 \sqrt{\mu - \lambda_0^N}} \left(\frac{1}{w_k^N(\mu)} \prod_{\substack{1 \leq \ell \leq L \\ \ell \neq n, k}} \frac{\sigma_\ell^{N,k} - \mu}{w_\ell^N(\mu)} \right) Q_k^{N,L}(\mu) \frac{\sigma_n^{N,k} - \mu}{w_n^N(\mu)}.$$

Arguing as in the proof of Lemma D.3 one sees that for λ in the rectangle $R_n \subset \mathbb{C}$, as introduced there, and $\mu = -2 + \frac{\lambda}{4N^2}$

$$\frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} = N f_n^{-,k}(\lambda) (1 + g_n^{N,k}(\lambda)) \frac{4N^2(\sigma_n^{N,k} - \mu)}{w_n^N(\mu)},$$

where

$$f_n^{-,k}(\lambda) = \frac{1}{\sqrt{\lambda - \lambda_0^-}} \frac{1}{w_k^-(\lambda)} \prod_{\substack{0 \leq \ell < \infty \\ \ell \neq n, k}} \frac{\sigma_\ell^{-,k} - \lambda}{w_\ell^-(\lambda)} = O\left(\frac{1}{n(k^2 - n^2)}\right) \quad (\text{D.19})$$

and $g_n^{N,k}$ is defined by the equation above and satisfies $g_n^{N,k}(\lambda) = O\left(\frac{1}{L^{5/2}}\right)$. Making the change of variable $\mu(x) = \lambda_{2n-1}^N - \rho^N + \frac{1}{4N^2}x$, implying that

$\lambda \equiv \lambda(\mu(x)) = \nu_{2n-1}^N - \rho + x$, and taking into account the sign of the standard root $w_n^N(\mu)$, the integral $\frac{1}{N} \int_{\lambda_{2n-1}^N - \rho^N}^{\lambda_{2n-1}^N} \frac{\varphi_k^N(\mu)}{i \varepsilon \sqrt{\chi_N(\mu)}} d\mu$ equals

$$\int_0^\rho f_n^{-,k}(\lambda) (1 + g_n^{N,k}(\lambda)) \frac{\tilde{\sigma}_n^{N,k} - \nu_{2n-1}^N + \rho - x}{\sqrt[4]{\tilde{\gamma}_n^N + \rho - x}} \frac{dx}{\sqrt{\rho - x}} \quad (\text{D.20})$$

where as usual, $\tilde{\sigma}_n^{N,k} = 4N^2(\sigma_n^{N,k} + 2)$, $\nu_{2n-1}^N = 4N^2(\lambda_{2n-1}^N + 2)$, and $\tilde{\gamma}_n^N = 4N^2\gamma_n^N$. As $0 \leq \tilde{\sigma}_n^{N,k} - \nu_{2n-1}^N \leq \tilde{\gamma}_n^N$ one has $\frac{\tilde{\sigma}_n^{N,k} - \nu_{2n-1}^N + \rho - x}{\sqrt[4]{\tilde{\gamma}_n^N + \rho - x}} = O(1)$ yielding

$$\frac{1}{N} \int_{\lambda_{2n-1}^N - \rho^N}^{\lambda_{2n-1}^N} \frac{\varphi_k^N(\mu)}{i \varepsilon \sqrt{\chi_N(\mu)}} d\mu = (T) + O\left(\frac{1}{L^{5/2}} \frac{N}{n(k^2 - n^2)}\right) \quad (\text{D.21})$$

$$(T) := \int_0^\rho f_n^{-,k}(x + \nu_{2n-1}^N - \rho) \frac{\tilde{\sigma}_n^{N,k} - \nu_{2n-1}^N + \rho - x}{\sqrt[4]{\tilde{\gamma}_n^N + \rho - x}} \frac{dx}{\sqrt{\rho - x}}$$

Arguing as in the proof of Lemma D.3 it remains to estimate the difference $(T) - (S)$ where

$$(S) := \int_0^\rho f_n^{-,k}(x + \lambda_{2n-1}^- - \rho) \frac{\sigma_n^{-,k} - \lambda_{2n-1}^- + \rho - x}{\sqrt[4]{\gamma_n^- + \rho - x}} \frac{dx}{\sqrt{\rho - x}}.$$

By Lemma C.6, it suffices to estimate, uniformly for $0 \leq x \leq \rho$,

$$f_n^{-,k}(x + \nu_{2n-1}^N - \rho) \frac{\tilde{\sigma}_n^{N,k} - \nu_{2n-1}^N + \rho - x}{\sqrt[4]{\tilde{\gamma}_n^N + \rho - x}} - f_n^{-,k}(x + \lambda_{2n-1}^- - \rho) \frac{\sigma_n^{-,k} - \lambda_{2n-1}^- + \rho - x}{\sqrt[4]{\gamma_n^- + \rho - x}}.$$

Note that by Cauchy's theorem, $f_n^{-,k}(x + \nu_{2n-1}^N - \rho) - f_n^{-,k}(x + \lambda_{2n-1}^- - \rho)$ is

$$O\left(\frac{1}{N} \frac{1}{n(k^2 - n^2)}\right) = O\left(\frac{M^2}{N} \frac{1}{n(k^2 - n^2)}\right)$$

whereas

$$(\tilde{\sigma}_n^{N,k} - \nu_{2n-1}^N) - (\sigma_n^{-,k} - \lambda_{2n-1}^-) = O\left(\frac{1}{L^{5/2}} + \frac{M^2}{N}\right) = O\left(\frac{1}{L^{5/2}}\right)$$

and, by Lemma C.3,

$$\left| \frac{\rho - x}{\sqrt[4]{\tilde{\gamma}_n^N + \rho - x}} - \frac{\rho - x}{\sqrt[4]{\gamma_n^- + \rho - x}} \right| \leq \sqrt[4]{|\tilde{\gamma}_n^N - \gamma_n^-|} = O\left(\frac{M}{N^{1/2}}\right).$$

Combining these estimates, the stated asymptotics follow. Going through the arguments of the proof the claimed uniformity statement follows. \square

Corollary D.8 *Uniformly for $1 \leq n \leq M$, $0 < k < n$, and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\int_{C_n^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = O\left(\frac{N}{n(k^2 - n^2)}\right). \quad (\text{D.22})$$

Proof. One verifies in a straightforward way that uniformly for any $1 \leq k, n \leq M$ with $k \neq n$, (D.19) holds and the function $g_n^{N,k}$, introduced in the proof of Lemma D.7, satisfies $1 + g_n^{N,k}(\lambda(\mu(x))) = O(1)$. Following the arguments of the proof of Lemma D.7, the claimed statement then follows. \square

Next we analyze the integral $\int_{C_n^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu$ in the case where $k = n$.

Lemma D.9 *Uniformly for $1 \leq n \leq L$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = \frac{N\gamma_n^-}{2\pi n} \int_{C_n^-} \frac{\psi_n^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\frac{N}{n} \left(\frac{M}{N^{1/2}} + \frac{1}{L^{5/2}}\right)\right). \quad (\text{D.23})$$

Proof. We proceed similarly as in the proof of Lemma D.7, but note that in the case $k = n$, the factor γ_n^N is a substitute for the missing factor $\sigma_n^{N,k} - \mu$. Using the same terminology as in the proof of Lemma D.7, we have

$$\frac{\tilde{\gamma}_n^N}{N} \int_{\lambda_{2n-1}^N - \rho^N}^{\lambda_{2n-1}^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} = \int_0^\rho f_n^{-,n}(\lambda) (1 + g_n^{N,k}(\lambda)) \frac{\tilde{\gamma}_n^N}{\sqrt{\tilde{\gamma}_n^N + \rho - x}} \frac{dx}{\sqrt{\rho - x}} \quad (\text{D.24})$$

where in the case at hand, $f_n^{-,n}$, defined on the rectangle R_n , satisfies

$$f_n^{-,n}(\lambda) = \frac{1}{\sqrt{\lambda - \lambda_0^-}} \prod_{\substack{0 < \ell < \infty \\ \ell \neq n}} \frac{\sigma_\ell^{-,n} - \lambda}{w_\ell^-(\lambda)} = O\left(\frac{1}{n}\right) \quad (\text{D.25})$$

and $g_n^{N,k}(\lambda)$ is again $O(\frac{1}{L^{5/2}})$. Following the arguments of the proof of Lemma D.7 yields the claimed statements. \square

Corollary D.10 *Uniformly for $1 \leq n \leq M$ and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\tilde{\gamma}_n^N \int_{C_n^N} \frac{\varphi_n^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = O\left(\frac{N}{n}\right). \quad (\text{D.26})$$

Proof. Arguing as in the proofs of Corollary D.8 and Lemma D.9 yields the claimed statement. \square

Finally we analyze the integrals over the intervals B_j^N with $1 \leq j \leq n-1$.

Lemma D.11 *Uniformly for $1 \leq n, k \leq L$, $0 < j < n$, and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\int_{B_j^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = \frac{N}{2\pi k} \int_{B_j^-} \frac{\psi_k^-(\lambda)}{i^c \sqrt{\Delta_-^2(\lambda) - 4}} d\lambda + O\left(\frac{N}{j \cdot d_{k,j}} \frac{1}{L^{5/2}}\right) \quad (\text{D.27})$$

where $d_{k,j} = 1 + \min(|j^2 - k^2|, |k^2 - (j+1)^2|)$.

Proof. As on the intervals B_j^N , the integrand $\frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}}$ is not singular, the integral is easier to handle. To get the claimed results, one argues as in the proof of Lemma D.1, taking into account that $\frac{1}{w_k^-(\lambda)} = O(\frac{1}{d_{k,j}})$. \square

Following the arguments of the proofs of Lemma D.11 and Corollary D.2 leads to the following

Corollary D.12 *Uniformly for $1 \leq n, k \leq M$, $0 < j < n$, and on bounded subsets of functions α, β in $C_0^2(\mathbb{T})$*

$$\int_{B_j^N} \frac{\varphi_k^N(\mu)}{i^c \sqrt{\chi_N(\mu)}} d\mu = O\left(\frac{N}{j \cdot d_{j,k}}\right). \quad (\text{D.28})$$

E Symmetry of the Toda chain

In this Appendix we discuss a symmetry of the Toda chain used to reduce the proof of the claimed asymptotics (1.4) of Theorem 1.1 of the frequencies ω_n^N at the right edge to the one of the asymptotics (1.3) of the frequencies at the left edge. It is more convenient to discuss the symmetry property for arbitrary Toda chains. Extend a vector $(b, a) = ((b_n)_{1 \leq n \leq N}, (a_n)_{1 \leq n \leq N})$ in $\mathbb{R}^N \times \mathbb{R}_{>0}^N$ and extend the vectors b and a periodically, $(b_n)_{n \in \mathbb{Z}}$, respectively $(a_n)_{n \in \mathbb{Z}}$. Define for any $n \in \mathbb{Z}$

$$\tilde{b}_n := -b_{N-n} \quad \tilde{a}_n := a_{N-1-n}$$

and consider the $2N \times 2N$ matrix $Q(\tilde{b}, \tilde{a})$. In this appendix, the superindex N is used to indicate the number of particles of the Toda chain considered and we use tilde for quantities such as eigenvalues, characteristic functions, discriminants, actions, ... when evaluated at (\tilde{b}, \tilde{a}) instead of (b, a) . So e.g. $\tilde{\lambda}_n^N = \lambda_n^N(\tilde{b}, \tilde{a})$ or $\tilde{\Delta}_N(\mu) = \Delta_N(\mu, (\tilde{b}, \tilde{a}))$. Then the following proposition holds.

Proposition E.1 *The following formulae hold:*

$$(i) \quad \tilde{\lambda}_n^N = -\lambda_{2N-1-n}^N \quad \forall 0 \leq n < 2N \quad \text{and hence} \quad \tilde{\chi}_N(\mu) = \chi_N(-\mu).$$

$$(ii) \quad \tilde{\Delta}_N(\mu) = (-1)^N \Delta_N(-\mu); \quad \partial_\mu \tilde{\Delta}_N(\mu) = -(-1)^N \dot{\Delta}_N(-\mu).$$

$$(iii) \quad \tilde{I}_n^N = I_{N-n}^N \quad \text{and} \quad \tilde{\varphi}_n^N(\mu) = (-1)^N \varphi_{N-n}^N(-\mu) \quad \forall 0 < n < N.$$

Combining these formulas one obtains

$$\tilde{\omega}_n^N = \omega_{N-n}^N \quad \forall 0 < n < N. \quad (\text{E.1})$$

Proof. To prove (i) note that an arbitrary vector $(F_n)_{1 \leq n \leq 2N} \in \mathbb{R}^{2N}$ is an eigenvector of $Q(b, a)$ with eigenvalue λ iff $((-1)^n F_{2N-n})_{1 \leq n \leq 2N}$ is an eigenvector of $Q(\tilde{b}, \tilde{a})$ with eigenvalue $-\lambda$. The identity for the characteristic polynomial $\tilde{\chi}_N(\mu)$ then follows from its product representation. Similarly, the first identity in (ii) follows from the product representation of $\Delta_N(\mu) - 2$, implying the claimed identity for the derivatives. Towards (iii) note that by the definition (1.10) of the c-root, one has $\sqrt[{}^c]{\Delta_N^2(-\mu) - 4} = (-1)^N \sqrt[{}^c]{\Delta_N^2(\mu) - 4}$. The stated identity for the actions then follows from formula (3.1) whereas the one for $\tilde{\varphi}_n^N$ follows from (1.12) and the fact that they are polynomials with leading term μ^{N-2} . Finally, the claimed formula for the frequencies then follows from (1.13). \square

By applying the identity (4.9) to $\tilde{\omega}_n^N$, Proposition E.1 leads to the following

Corollary E.2 *For any $0 < n < N$*

$$\begin{aligned} iN\omega_{N-n}^N &= \sum_{j=1}^n \int_{\lambda_{2N-1-(2j-1)}^N}^{\lambda_{2N-1-(2j-2)}^N} \frac{(\mu - \mathbf{p}_N/N) \dot{\Delta}_N(\mu)}{\sqrt[{}^c]{\Delta_N^2(\mu) - 4}} d\mu \\ &- \sum_{k \in \mathcal{J}_N} I_{N-k}^N \omega_{N-k}^N \sum_{j=1}^n \int_{\lambda_{2N-1-(2j-1)}^N}^{\lambda_{2N-1-(2j-2)}^N} \frac{\varphi_{N-k}^N(\mu)}{\sqrt[{}^c]{\chi_N(\mu)}} d\mu \end{aligned} \quad (\text{E.2})$$

References

- [1] D. Bambusi, T. Kappeler, T. Paul: *De Toda à KdV*, C.R. Acad. Sci Paris, Ser. I, 347, no. 17-18, 1025-1030, 2009
- [2] D. Bambusi, T. Kappeler, T. Paul: *From Toda to KdV*, preprint
- [3] D. Bambusi, A. Ponno: *On metastability in FPU*, Comm. Math. Phys. 264, 539-561, 2006
- [4] G. Berman, F. Izrailev, *The Fermi-Pasta-Ulam problem: 50 years of progress*, Chaos 15, 015104, 2005
- [5] A. Bloch, F. Golse, T. Paul, A. Uribe: *Dispersionless Toda and Toeplitz operators*, Duke Math. J. 117, 157-196, 2003
- [6] A. Carati, L. Galgani, A. Giorgilli, *The Fermi - Pasta - Ulam problem as a challenge for the foundations in physics*, Chaos, 15, 015105, 2005
- [7] C. Cercignani: *Solitons. Theory and application*, Riv. Nuovo Cim. 7, 429-469, 1977
- [8] E. Fermi, J. Pasta, S. Ulam, *Los Alamos report LA - 1940*, in E. Fermi, *Collected Papers*, 2, 977-988, 1955
- [9] H. Flaschka: *The Toda lattice I. Existence of integrals*, Phys. Rev. B 9, 1924-1925, 1974
- [10] H. Flaschka, D.W. McLaughlin: *Canonically conjugate variables for the Korteweg-de Vries equation and the Toda lattice with periodic boundary conditions*, Progress of Theoretical Physics, 55, 438-456, 1976
- [11] G. Gallavotti (Ed.), *Fermi-Pasta-Ulam problem: a status report*, Lect. Notes Phys. vol 728, Springer, 2010
- [12] A. Henrici, T. Kappeler: *Global action-angle variables for the periodic Toda lattice*, IMRN 2008; doi:10.1093/imrn/rnn031
- [13] A. Henrici, T. Kappeler: *Global Birkhoff coordinates for the periodic Toda lattice*, Nonlinearity 21, 2731-2758, 2008

- [14] A. Henrici, T. Kappeler: *Birkhoff normal form for the periodic Toda lattice*, Integrable Systems and Random Matrices, Contemp. Math., vol 458, 11 - 29, American Math. Soc., 2008
- [15] A. Henrici, T. Kappeler: *Nekhoroshev theorem for the periodic Toda lattice*, Chaos 19, 033120, 2009; doi:10.1063/1.3196783
- [16] T. Kappeler, J. Pöschel: *KdV & KAM*, Springer, 2003
- [17] T. Kappeler, B. Schaad, P. Topalov: *Qualitative features of periodic solutions of KdV*, arXiv:1110.0455, 2011
- [18] J. Laskar: *Frequency analysis for multi-dimensional systems. Global dynamics and diffusion*, Physica D, 67, 257281, 1993
- [19] J. Pöschel, E. Trubowitz: *Inverse spectral theory*, Academic Press, 1987
- [20] M. Toda: *Theory of nonlinear lattices*, 2nd edition, Springer Series on Solid-State Sciences 20, Springer, 1989
- [21] G. Schneider, C. Wayne: *Counter-propagating waves on fluid surfaces and the continuum limit of the Fermi-Pasta-Ulam model*, in International Conference on Differential Equations, Vol. 1.2 (Berlin 1999), 390-404, World Sci. Publ., 2000
- [22] N. Zabusky, M. Kruskal, *Interaction of solitons in a collisionless plasma and the recurrence of initial states*, Phys. Rev. Lett. 15, 240 - 245, 1965