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Substituent effect in unsymmetrical lutetium bisphthalocyanines: a DFT analysis†

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DFT calculations have been carried out on a series of unsymmetrical bisphthalocyanine lutetium complexes in which one of the ligands is substituted by 8 or 16 chlorine atoms. It is shown that their unpaired electron is predominantly localized on the non-substituted ligand. An orbital explanation is provided to rationalize this effect. A good agreement is found between the computed and experimental ionization potentials and electroaffinities and the redox potentials of a closely related series, as well as between their TDDFT-computed and experimental UV-visible transitions which are analysed.

Introduction

The lanthanide bisphthalocyanines are sandwich complexes having two organic redox centers, the macrocyclic units, linked together by the metal ion, Ln(III), which apparently does not play a leading role as most properties are associated with the π -electron orbitals: the complex is equivalent to Ln^{3+} linked to two ligands, one phthalocyanine ring, formally Pc^{2-} , and one oxidized macrocycle, formally $\text{Pc}^{\bullet-}$.¹ The close proximity of the two phthalocyanines, 2.69 Å in the lutetium derivative,² allows the delocalization of the unpaired electron over both the macrocyclic rings *i.e.*, two equivalent $\text{Pc}^{1.5-}$ ligands, a fact which is now well accepted but has been debated,^{2,3} mainly because the two phthalocyanines have different shapes in the crystalline forms, one being domed and the other one flatter. This could be caused by packing in the solid phase or by differences in the environment of each of the two phthalocyanines. It also makes sense that, in an isotropic environment, in vacuum for example, the electron must be equally shared by the two π -systems of the symmetrical $\text{Lu}(\text{Pc})_2$ molecule. However, it is reasonably expected that asymmetry in the conjugated π -macrocycles of the sandwich complex would influence the electron density. Numerous molecules coupling two phthalocyanines differently substituted have been synthesized.^{4,5} From the electronic absorption spectra of phthalocyaninato-naphthalocyaninato Lu^{III} complexes, it has been concluded that the electron is delocalized and that the hole density is higher on the naphthalocyanine.^{3a,6} Near IR and Raman spectroscopies of heteroleptic complexes allowed some authors to conclude that the orbital levels keep

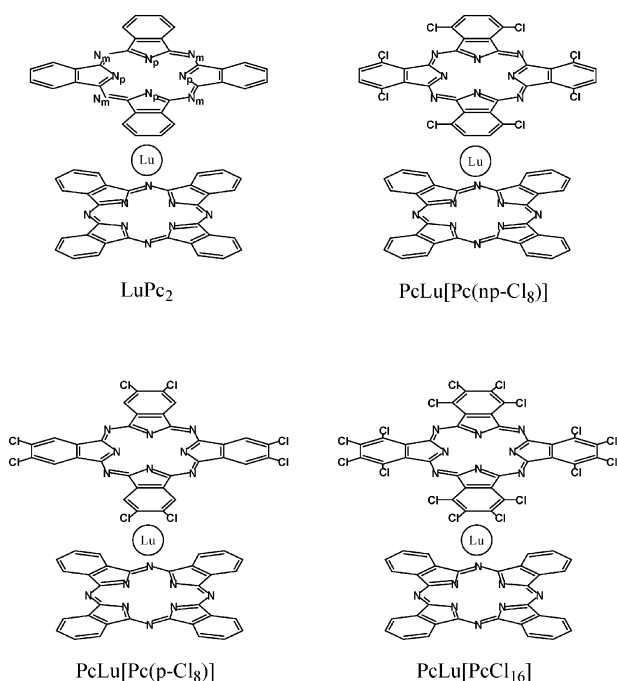
the porphyrin and phthalocyanine characters, a spectro-electrochemical study indicating that oxidation involves the porphyrin ring, reduction occurring on the phthalocyanine.^{5,7–10} However, from IR and Raman spectroscopies of numerous unsymmetrically substituted bisphthalocyanines, studied under their neutral as well as their reduced and oxidized forms, we have been unable to find any genuine marker of the redox state of the phthalocyanine units. This is why experimental works, as well as a theoretical study, have been undertaken to improve knowledge about the charge partitioning in heteroleptic bisphthalocyaninato lutetium complexes.

Unsymmetrical lutetium bisphthalocyanines, bearing four electron donor groups, *tert*-butyl substituents, on one ring and four, eight or sixteen acceptors, Cl, on the second unit, have been synthesized. They show very unusual properties, evidenced by spontaneous dimerization, in solutions at ambient temperature, and also by their electrochemical properties.¹¹ It appeared interesting to know whether this could be related to the influence of asymmetry on the charge and spin densities over the two macrocycles. Numerous theoretical calculations have already been performed for phthalocyanines and phthalocyanine complexes of which selected references are provided herein.^{12–28} However, to our knowledge, no density functional theory (DFT) calculations have been performed on bisphthalocyanine lutetium complexes in different oxidation states and no complete MO rationalization of the electronic structure in unsymmetrical complexes has been undertaken so far. In the following, we analyze by means of DFT calculations the bonding and the localization of the unpaired electron in LuPc_2 and the perturbation caused by the presence of chlorine substituents on one of the Pc ligands. The investigated compounds, namely LuPc_2 , $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]$ and $\text{PcLu}[\text{PcCl}_{16}]$, are shown in Scheme 1. Calculations have been carried out on the neutral, cationic and anionic forms of each bisphthalocyanine. The UV-visible spectra of all these complexes have been also computed at the TDDFT level and compared to the experimental ones. Details of the calculations are given below.

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† Electronic supplementary information (ESI) available: Atomic Cartesian coordinates of the optimized compounds (Table S1) and distances between the $(\text{N}_p)_4$ plane and the planes formed by the various ligand type atoms (Table S2). See DOI: 10.1039/b810131k



Result and discussion

The free phthalocyanine ligands

In order to derive the electronic structure of the above mentioned bisphthalocyanine complexes from that of their free ligands, considered here as hypothetical building block units for the whole complexes, we have carried out calculations at the same level of theory on the neutral and dianionic forms of the isolated ligand fragments, *i.e.*, $[\text{Pc}]^{0/2-}$, $[\text{Pc}(\text{p-Cl}_8)]^{0/2-}$, $[\text{Pc}(\text{np-Cl}_8)]^{0/2-}$ and $[\text{PcCl}_{16}]^{2-}$. D_{4h} symmetry and singlet states were firstly assumed. In terms of level ordering, energy gaps and orbital localization, our results match, at least qualitatively, those of previous investigations on $[\text{Pc}]^{0/2-}$ obtained at various levels of theory.^{13b,16a,23} To our knowledge, no optimized structures of free $[\text{Pc}]^{0/2-}$ have been published so far. Nevertheless, our metrical data (see ESI†) are consistent with the most recent results on metal-complexes MPc and free H_2Pc species.^{19c,20a,b,d,e,24c,28} The MO diagrams of the computed dianionic forms are shown in Fig. 1.

They exhibit two nearly degenerate non-bonding HOMO's of a_{1u} and b_{1g} symmetry lying above an a_{2u} MO, their relative ordering depending on the nature of the substituent (H or Cl) at C_γ . These orbitals are plotted in Fig. 2 in the case of Pc^{2-} . They are very similar in the case of the other substituted phthalocyanines, in which they exhibit only small chlorine participation.

It is noteworthy that the π -type a_{1u} MO has no participation by symmetry of the nitrogen atoms and exhibits a pseudo-8-fold symmetry on the macrocycle. On the other hand, the σ -type b_{1g} MO can be viewed as the out-of-phase combination of the four ligand nitrogen lone-pairs. This orbital is expected to be strongly involved in the phthalocyanine complexation by a metal. The π -type a_{2u} HOMO-2 has also significant localization on the four basic nitrogen atoms. Thus, the b_{1g} orbital (and to a

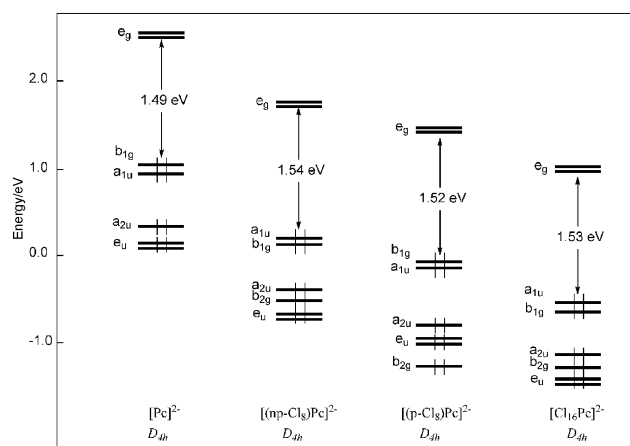


Fig. 1 One-electron energy diagrams of the dianionic forms of the free phthalocyanine ligands. Only the highest occupied and lowest unoccupied levels are considered.

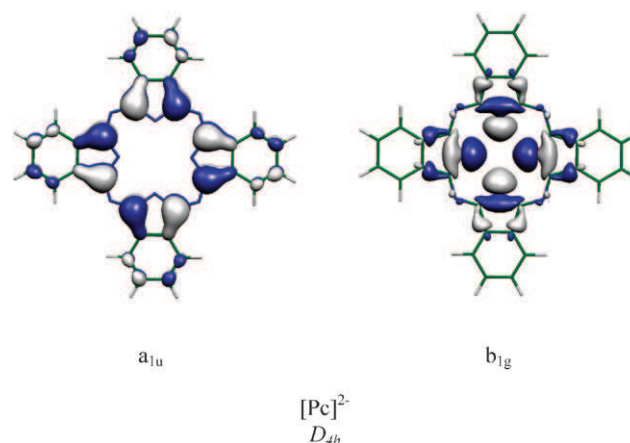


Fig. 2 Plots of the two highest occupied MO's of Pc^{2-} .

lesser extend the a_{2u} MO) is expected to be stabilized upon metal complexation, whereas the a_{1u} orbital will remain unperturbed.

The doubly oxidized phthalocyaninato ligands, with two electrons less are the neutral phthalocyanines Pc^0 which exhibit, in the assumed D_{4h} symmetry, a very small HOMO(b_{1g})/LUMO(a_{1u}) gap, with no significant change in the ordering of the other levels when compared to their dianions. It turns out that this small HOMO/LUMO gap induces second-order Jahn–Teller instability for the singlet state of the neutral D_{4h} species which are all found to be more stable in a distorted D_{2h} structure by ~ 0.5 eV. On the other hand, the dianionic species are more stable in the D_{4h} symmetry.

An important feature arising from a look at Fig. 1 is the variation of energy of the frontier orbital levels with respect to the number of chlorine substituents. The larger the number of chlorines, the lower the energy of the phthalocyanine frontier orbitals. This trend is independent of the phthalocyanine oxidation state. For example, the energy of the a_{1u} level varies in the order $\text{Pc} > \text{Pc}(\text{p-Cl}_8) \approx \text{Pc}(\text{np-Cl}_8) > \text{PcCl}_{16}$ in both the neutral and dianionic species. This chlorine dependence is the consequence of the halogen electro-attracting inductive effect which tends to stabilize the phthalocyanine frontier orbital. Thus, it largely overcomes the halogen π -donating

effect which tends to destabilize the same levels. The difference between the frontier orbital energies of the $\text{Pc}(\text{p-Cl}_8)$ and $\text{Pc}(\text{np-Cl}_8)$ isomers is due to the closeness to the N_p atoms (see Scheme 1) of the chlorine substituents in the latter which renders stronger the inductive effect.

The bisphthalocyanine lutetium complexes

We start the analysis by looking at the $(\text{LuPc}_2)^{+/0/-1}$ series. Although the $\text{Ln}(\text{III})$ bisphthalocyanine complexes usually crystallize in the staggered *pseudo-D*_{4d} conformations,^{2,29} geometry optimizations were made starting from the *D*_{4h} (eclipsed) and *D*_{4d} (staggered) conformations, but assuming lower *C*_{4v} or *C*₂ symmetry constraint, thus allowing distortion away from the more symmetrical conformations. In any case, the lowest energy was found to correspond to the staggered conformation (by ~ 0.6 eV) with very little or no significant distortion away from the ideal *D*_{4d} symmetry. It should be noted that optimization of the less stable eclipsed conformation leads to unsymmetrical phthalocyanines in the case of $(\text{LuPc}_2)^{+/0/-1}$. The results reported here correspond to the more stable staggered conformation optimized assuming *C*_{4v} symmetry constraint. Relevant calculated data are given in Table 1.

The optimized geometry of the neutral form is shown in Fig. 3, some relevant metrical data are given in Table 1 and the MO diagrams of the $(\text{LuPc}_2)^{+/0/-1}$ series are sketched in Fig. 4. It should be noted that the optimized metric data of LuPc_2 are in a good agreement with the available experimental ones, with slightly longer optimized Lu–N_p distances (2.41 Å vs. 2.38 Å).² The slight longitudinal elongation of the molecule with the

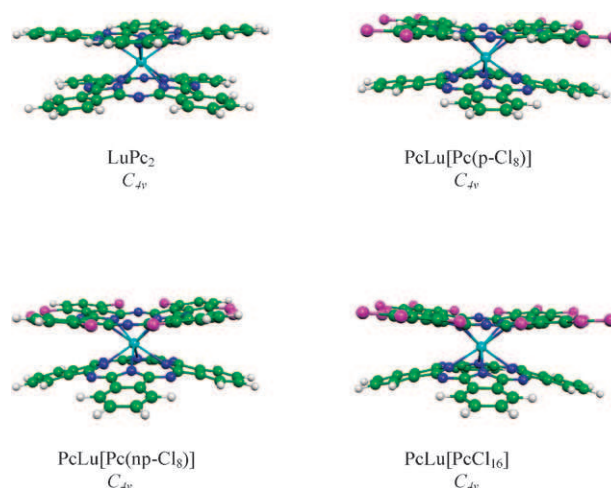


Fig. 3 Optimized structures of LuPc_2 , $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]$ and $\text{PcLu}[\text{PcCl}_{16}]$.

increasing number of electrons occupying the $2a_2$ level is consistent with the π – π antibonding character of this orbital. This effect has been proved experimentally on a related terbium series.³² On the other hand, the Pc doming is quite insensitive to the oxidation state of the complex (Table 1), as exemplified by the average intra-ligand distances between the $(\text{N}_\text{p})_4$ and the $(\text{C}_\text{p})_4$ planes which all lie between 0.55 Å and 0.65 Å (see Table S2† for full metrical data describing Pc doming).

The singlet ground states of the cationic and anionic forms are associated with large HOMO/LUMO gaps. They differ from each other by the occupation of the $2a_2$ level (also labeled

Table 1 Major computed data for the $(\text{LuPc}_2)^{+/0/-1}$, $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]^{+/0/-1}$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]^{+/0/-1}$ and $\text{PcLu}[\text{PcCl}_{16}]^{+/0/-1}$ series (see computational details). Values in parentheses are experimental distances taken from ref. 2

	$[\text{LuPc}_2]^q$ <i>C</i> _{4v}			$\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]^q$ <i>C</i> _{4v}			$\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]^q$ <i>C</i> _{4v}			$\text{PcLu}[\text{PcCl}_{16}]^q$ <i>C</i> _{4v}		
	<i>q</i> = +1	= 0	= −1	<i>q</i> = +1	= 0	= −1	<i>q</i> = +1	= 0	= −1	<i>q</i> = +1	= 0	= −1
HOMO–LUMO Gap/eV	0.71	0.88	0.92	0.70	0.85	0.83	0.49	0.92	0.90	0.48	0.87	0.76
First ionisation energy/eV		5.74			6.12			5.98			6.21	
Electronic affinity/eV		3.38			3.74			3.64			3.93	
μ/Debye	0.00	0.05	0.14	2.99	2.98	1.99	2.26	2.74	1.71	3.98	5.35	3.61
Lu–N _p distances/Å ^a	2.408	2.415 (2.372)	2.419 (2.392)	2.397	2.415	2.439	2.407	2.420	2.445	2.414	2.425	2.458
	2.398	2.418 (2.387)	2.416 (2.374)	2.390	2.400	2.407	2.396	2.405	2.410	2.397	2.403	2.399
(Np) ₄ ···(Np) ₄ interplane distance/Å	2.730	2.781 (2.690)	2.770 (2.701)	2.716	2.755	2.798	2.701	2.731	2.764	2.719	2.736	2.770
Pc ^{Cl} doming angle/° ^b	172	173	172	173	174	174	171	171	171	173	172	172
Pc doming angle/° ^b	172	173	172	172	173	173	167	168	166	169	167	167
Mulliken net charges												
Lu	1.56	1.55	1.57	1.56	1.56	1.56	1.60	1.60	1.61	1.61	1.63	1.63
Pc ^{Cl}	−0.28	−0.77	−1.28	−0.39	−0.90	−1.35	−0.39	−0.92	−1.36	−0.48	−1.02	−1.42
Pc	−0.28	−0.78	−1.29	−0.18	−0.66	−1.21	−0.21	−0.69	−1.25	−0.13	−0.61	−1.21
% 2a ₂ Pc ^{Cl}	50	50	50	46	41	31	46	41	21	43	34	16.5
% 2a ₂ Pc	50	50	50	54	59	69	54	59	74	57	66	83.5
Spin density												
Lu		0			0			0				
Pc ^{Cl}		0.50			0.40			0.39			0.32	
Pc		0.50			0.60			0.61			0.68	

^a Bold values correspond to the substituted Pc^{Cl} ligand. ^b The doming angle is defined as the angle between the (Np)₄ plane and each of the planar C₈N pyrrolic units.

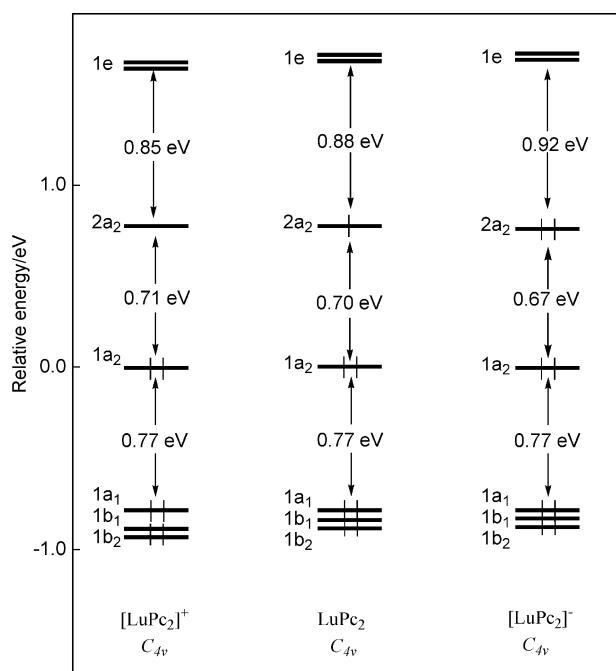


Fig. 4 MO diagrams of the $(\text{LuPc}_2)^{+/0/-1}$ models. The energy of the $1a_2$ level has been arbitrarily set to zero in the three diagrams.

a_2 in the D_{4d} symmetry). This orbital, which is the SOMO of the neutral species, is shown in Fig. 5. It has no metal participation by symmetry and is equally localized on both Pc ligands. Consistently, the spin density on each Pc ligand is equal to 0.5. Any attempt to dissymmetrize the SOMO by rendering non-equivalent the Pc ligands in the geometry optimization process lead to the same symmetrically delocalized mixed valence.

Although this result is fully consistent with the recent conclusion of a near-infrared intervalence band analysis on LuPc_2 ,³⁰ one may argue that the considered level of theory tends to systematically favor delocalized distributions, due to the self-interaction errors contained in the DFT formalism. Since the aim of this paper is not to deeply analyze the

mixed-valent nature of LuPc_2 , but mainly investigate its unsymmetrical relatives, we did not carry out calculations on LuPc_2 at a higher level of theory. Nevertheless, taken as a whole, the DFT results, the experimental X-ray structure² and the near-infrared data³⁰ suggest that LuPc_2 has a delocalized or nearly delocalized mixed valence.

As said above, the $2a_2$ SOMO of LuPc_2 has no metal (including 4f) character. As one can see in Fig. 5, it is the out-of-phase combination on the π -type a_{1u} HOMO/SOMO of $\text{Pc}^{-2/-1}$ (see Fig. 1 and 2). The corresponding in-phase combination (b_1 in D_{4d} symmetry) is the next lower $1a_2$ MO, which is occupied in the three $(\text{LuPc}_2)^{+/0/-1}$ species. Thus, the two highest occupied orbitals of LuPc_2 do not depend on the metal–ligand interaction since they are only combinations of the a_{1u} Pc frontier orbitals (see left side of Scheme 2) with no metal participation. The energy splitting between these two combinations is not negligible (~ 0.7 eV), resulting from significant overlap between the two Pc a_{1u} orbitals. This splitting is about the same in the eclipsed conformation due to the near 8-fold symmetry (Fig. 2) of the Pc a_{1u} orbitals. Therefore, the occupation number of the a_2 system in the $(\text{LuPc}_2)^{+/0/-1}$ series has no significant effect on the rotational conformation preference, which, as well as the macrocycle doming, has been shown to be essentially driven by the necessity of maximizing metal–ligand bonding together with minimizing the steric hindrance between the ligands.³³ On the other hand, because of the π – π antibonding nature of the $2a_2$ MO, its occupation number has some effect on the distance between the two Pc ligands. This result as well as the energy diagrams of Fig. 4 agree with recently published DFT calculations on $(\text{YPc}_2)^{\pm}$.^{18b}

Unlike the $(\text{LuPc}_2)^{+/0/-1}$ series, the chlorinated relatives $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]^{+/0/-1}$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]^{+/0/-1}$ and $\text{PcLu}[\text{PcCl}_{16}]^{+/0/-1}$ have non-equivalent ligands. They were also found to be more stable in the staggered conformation, assuming C_{4v} symmetry constraint. Relevant computed data are given in Table 1. The optimized structures and the MO diagrams corresponding to the neutral species are shown in Fig. 3 and Fig. 6, respectively. Although unsymmetrical, the substituted complexes have quite similar structural characteristics to their non-substituted relatives. The $\text{Pc}(\text{np-Cl}_8)$ and PcCl_{16}

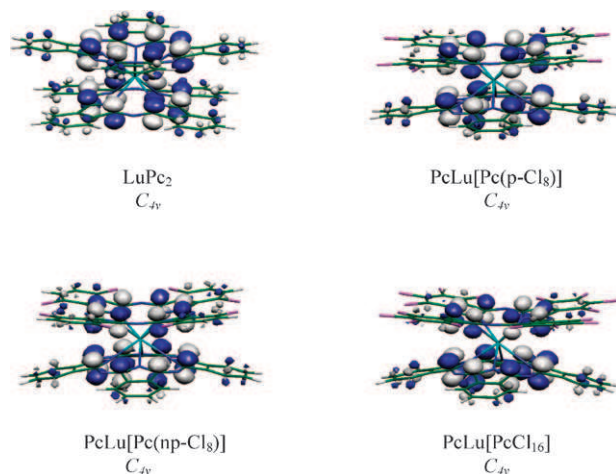
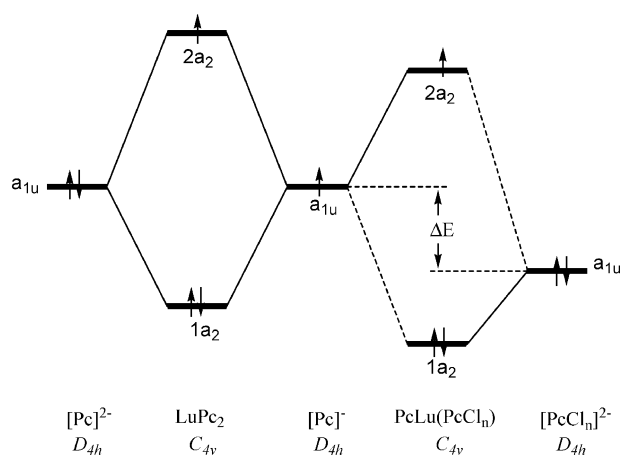


Fig. 5 Plots of the SOMO's of LuPc_2 , $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]$ and $\text{PcLu}[\text{PcCl}_{16}]$.



Scheme 2

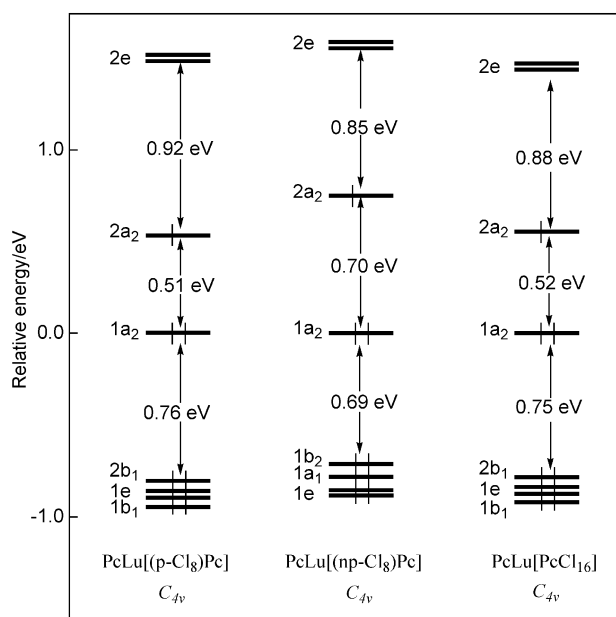


Fig. 6 MO diagrams of $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]$ and $\text{PcLu}[\text{PcCl}_{16}]$. The energy of the $1a_2$ level has been arbitrarily set to zero in the three diagrams.

complexed ligands exhibit a more pronounced doming than their Pc counterpart, whereas $\text{Pc}(\text{p-Cl}_8)$ does not. Clearly, the np position of the chlorine substituents (Scheme 1) on the macrocycle has a significant effect on its steric hindrance. Consistently, the $\text{Lu-N}_p(\text{Pc}^{\text{Cl}})$ distances are larger in the case of $\text{Pc}(\text{np-Cl}_8)$ and PcCl_{16} than in the case of $\text{Pc}(\text{p-Cl}_8)$. On the other hand, the $\text{Lu-N}_p(\text{Pc})$ distances of the three substituted complexes are almost equal and slightly shorter than in the non-substituted LuPc_2 species, whereas the doming of the unsubstituted Pc ligand is more pronounced than that of the substituted Pc^{Cl} one (Table 1). For example, in $\text{PcLu}[\text{PcCl}_{16}]$ the average distance between the $(\text{N}_p)_4$ and the $(\text{C}_7)_4$ planes in the Pc and PcCl_{16} ligands is 1.07 Å and 0.66 Å, respectively (see Table S2† for full metrical data associated with ligand doming).

As for LuPc_2 , the two highest occupied orbitals of $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]$ and $\text{PcLu}[\text{PcCl}_{16}]$ are the bonding and antibonding combinations of the phthalocyanine a_{1u} frontier orbitals. However, these two a_{1u} orbitals are no longer degenerate since they belong to different phthalocyanine ligands. As said above, the frontier orbitals of the chlorinated phthalocyanine lie at lower energy than those of the non-substituted Pc. The non degenerate a_{1u} interaction is sketched on the right side of Scheme 2. It follows that the bonding combination has a larger contribution of the lower a_{1u} orbital, *i.e.* PcCl_n , whereas the antibonding counterpart has a dominant Pc localization. As a consequence, the unpaired electron of the neutral species has a more important localization on the non-substituted Pc ligand. The mixing between the $a_{1u}(\text{Pc})$ and $a_{1u}(\text{PcCl}_n)$ orbitals depends on their energy difference ΔE . The smaller (larger) ΔE , the larger (smaller) the mixing. Thus, the lower the energy of $a_{1u}(\text{PcCl}_n)$, the smaller the mixing and the more Pc-polarized the unpaired electron of the neutral species. In other words, the density of

the unpaired electron on the non-substituted phthalocyanine ring increases with the number of Cl atoms on the chlorinated phthalocyanine. This trend is nicely reproduced by the computed data (Table 1) which show that the Pc participation in the $2a_2$ MO of the various computed models increases with the number of Cl atoms. This is also illustrated by the plots of the SOMO of the neutral species shown in Fig. 5. Consistently, the DFT-computed spin density on the Pc ligand increases when going from LuPc_2 (0.5) to $\text{PcLu}[\text{PcCl}_{16}]$ (0.68). It is close to 0.6 in both $\text{PcLu}[\text{PcCl}_8]$ isomers. It is noteworthy that in all the $\text{PcLu}[\text{PcCl}_n]$ complexes, the PcCl_n Mulliken charge is more negative than that of the Pc one, despite the fact that the unpaired electron is more localized on Pc. The explanation lies in the electron withdrawing effect of the chlorine atoms which act as global attractors on the whole electron density. As a matter of fact, this polarization effect is largely independent from the oxidation state of the $\text{PcLu}[\text{PcCl}_n]^{+/0/-}$ species, *i.e.*, from the occupation number of the $2a_2$ MO.

The bonding mode is very similar in all the computed species and involves only the 5d Lu orbitals. Indeed, the atomic 4f Mulliken population is unsurprisingly found to be always almost equal to 14 (no electron donation to the ligands), whereas the 6s and 6p populations are never significantly different from zero (no donation from the ligands). On the other hand, the formally vacant 5d orbitals of $\text{Lu}(\text{III})$ in $[\text{LuPc}_2]^+$ receive 1.41 electron from the ligands, indicating significant covalent interaction. This transfer decreases slightly to 1.38 and 1.37 in the more reduced forms LuPc_2 and $[\text{LuPc}_2]^-$, respectively. This is the consequence of the larger repulsion between the ligands in the reduced forms which in turn slightly weakens the metal–ligand covalent bonding interaction. Slightly lower 5d populations are computed in the chlorinated derivatives, the smaller ones corresponding to the $\text{PcLu}[\text{PcCl}_{16}]^{+/0/-}$ series (1.38, 1.36 and 1.34, respectively). To get a better insight into the metal–ligand bonding, we have carried out a decomposition of the energy corresponding to the interaction between the Lu^{3+} cation and the ligand system in a similar way as Ricciardi *et al.* for metal bisporphyrin complexes.³³ The calculations were carried out on the closed-shell $\text{PcLu}[\text{PcCl}_n]^{\pm}$ couples. The total bonding energy is decomposed in the usual way into the sum of the orbital interaction energy and the steric interaction energy, the latter being the sum of a Pauli repulsion term and of an electrostatic term. The corresponding computed values are given in Table 2. Considering first the four cationic species, one can see that the Pauli repulsion and the orbital interaction energies vary little across the series, as compared to the electrostatic interaction energy term which dominates the total interaction energy. The electron withdrawing character of the chlorine atoms tends to reduce the negative charge on the coordinated nitrogens, thus lowering the strength of the electrostatic interaction. Adding two electrons to the cationic species does not modify this effect that much since these two electrons cannot interact by symmetry with the metal valence orbitals. Thus, in the case of the anionic series, the total interaction energy is also dominated by its electrostatic component and a trend similar to that of the cationic series is observed. Finally, it should be noticed that the energy decomposition described above does not take into account

Table 2 Two-fragment energy decomposition of $\text{PcLu}[\text{PcCl}_n]^{\pm}$. Cation fragments: Lu^{+3} and $\{\text{Pc} + [\text{PcCl}_n]\}^{2-}$; anion fragments: Lu^{+3} and $\{\text{Pc} + [\text{PcCl}_n]\}^{4-}$

	$[\text{LuPc}_2]^q C_{4v}$		$\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]^q C_{4v}$		$\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]^q C_{4v}$		$\text{PcLu}[\text{PcCl}_{16}]^q C_{4v}$	
	$q = +1$	$q = -1$	$q = +1$	$q = -1$	$q = +1$	$q = -1$	$q = +1$	$q = -1$
Pauli repulsion energy/eV	9.20	8.66	9.35	8.55	9.12	8.41	9.01	8.43
Electrostatic Interaction energy/eV	-31.38	-45.65	-29.94	-43.47	-30.76	-44.65	-29.61	-43.24
Total steric interaction energy/eV	-22.18	-36.99	-20.58	-34.92	-21.63	-36.24	-20.60	-34.81
Orbital interaction energy/eV	-26.07	-25.98	-26.33	-26.07	-25.99	-25.64	-25.99	-25.69
Total bonding energy/eV	-48.25	-62.97	-46.92	-60.99	-47.62	-61.88	-46.59	-60.50

the repulsive interaction energy between the two phthalocyanine ligands, which is not negligible. For example, it was computed to be 2.9 eV and 9.5 eV in the case of $[\text{LuPc}_2]^+$ and $[\text{LuPc}_2]^-$.

A good way to test the quality of the DFT results described above is the correlation of the computed ionization potentials and electron-affinities of the neutral species with their experimental oxidation and reduction potentials, respectively. The experimental redox potentials are those of the *t*Bu derivatives $\text{Lu}[(\text{tBu}_4)\text{Pc}]_2$, $[(\text{tBu}_4)\text{Pc}]\text{Lu}[\text{Pc}(\text{p-Cl}_8)]$, and $[(\text{tBu}_4)\text{Pc}]\text{Lu}[\text{PcCl}_{16}]$. We have evaluated those of $[(\text{tBu}_4)\text{Pc}]\text{Lu}[\text{Pc}(\text{np-Cl}_8)]$ through calculations using Hammett coefficients.³⁴ The corresponding curves (Fig. 7) exhibit the expected qualitative linear correlations with very satisfying agreement, especially if one consider that the computed and experimental series are not exactly the same. This consistency between computed and experimental data brings full confidence in the whole theoretical results.

The opportunity of having determined the DFT wavefunctions of a whole family of lutetium bisphthalocyanines offered us the possibility of carrying out TDDFT calculations (see computational details) to calculate their major optical transitions and to rationalize their experimental spectra. We start with the description of the experimental UV-visible spectrum of LuPc_2 (recorded in CH_2Cl_2), shown in Fig. 8 together with the attributions of the major transitions as generally given in the literature.

The intense Soret band is associated with inter-phthalocyanine allowed transitions.^{12,15,30} The next band is known as the blue vibronic band (BV) and is due to various transitions which are forbidden in the idealized D_{4h} or D_{4d} symmetry.^{12,30} As for the Soret band, the *Q* band (and its vibronic component *Q*_{vib}) is associated with various allowed transitions.^{30a} The weak band around 900 nm results from a forbidden transition involving the SOMO. It is known as the red vibronic band (RV) and is, as well as the BV band, characteristic of the radical nature of LuPc_2 . Finally, a broad inter-valence band (IVB), not recorded on the spectrum of Fig. 8 but observed in the 1100–1600 nm range, is attributed to a HOMO–SOMO transition and sometimes described as associated with a charge transfer between the two macrocycles.^{15,12,30,31}

The experimental λ_{max} associated with the above described bands are reported in Table 3 for $(\text{LuPc}_2)^{+/0/-1}$, together with the values corresponding to the major computed transitions (oscillator strength in parentheses) for the $(\text{LuPc}_2)^{+/0/-1}$, $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]^{+/0/-1}$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]^{+/0/-1}$ and $\text{PcLu}[\text{PcCl}_{16}]^{+/0/-1}$ series. The simulated spectra of the $(\text{LuPc}_2)^{+/0/-1}$ series are shown in Fig. 9. Because of the large

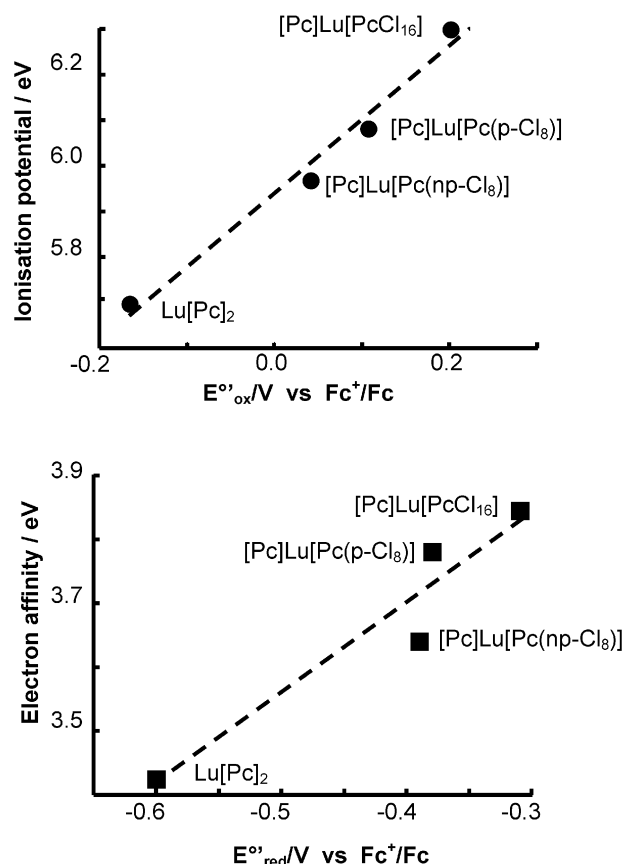


Fig. 7 Computed ionization potentials of the neutral species vs. electrochemical oxidation potentials (top) and computed electron affinities vs. electrochemical reduction potentials (bottom). Redox potentials are those of the $[(\text{tBu}_4)\text{Pc}]\text{Lu}[\text{Pc}(\text{np-Cl}_n)]$ derivatives. The redox potentials of $[(\text{tBu}_4)\text{Pc}]\text{Lu}[\text{Pc}(\text{np-Cl}_8)]$ have been evaluated using Hammett coefficients.

computing effort required, the transitions associated with the Soret band were not calculated.

There is a good agreement between the experimental and computed transitions of LuPc_2 . Calculations predict the $1a_2 \rightarrow 2a_2$ IVB transition to lie around 1300 nm. This value is close to the reported experimental one,^{30,31} suggesting that the Born–Oppenheimer approximation is still valid for this transition, thus supporting the fully delocalized nature of the unpaired electron. The RV band can be approximated to a $2a_2$ (SOMO) $\rightarrow 2e$ (LUMO + 1) ($a_2 \rightarrow e_3$ in D_{4d} symmetry) transition (78%), in agreement with an earlier indexation,^{30a} together with some $2a_2 \rightarrow 1e$ admixture (20%). Its small

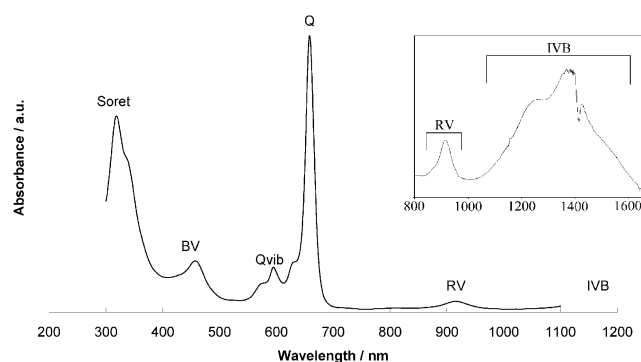


Fig. 8 Experimental UV-vis spectrum of LuPc_2 in CH_2Cl_2 .

oscillator strength comes from the strong pseudo- D_{4d} symmetry of the complex which renders the transition almost forbidden. However, vibronically allowed transitions are responsible for the observation of this band.³⁰ The major components of the Q band are found to be mainly associated with transitions from two occupied ligand levels of e symmetry to the $2a_2$ SOMO. Of course, the experimental Q_v band is not computed since it involves vibronic states. The BV band is mainly associated with $1a_1 \rightarrow 2a_2$ and $1b_1 \rightarrow 2a_2$ transitions. As for the RV band, associated vibronically allowed transitions are responsible for the observation of this band.^{30a} Although satisfying, the agreement between experiment and calculations in the case of the oxidized and reduced forms of LuPc_2 is not as good as in the case of the neutral species. This is probably due to the fact that the calculations do not take into account the solvent and counter-ion effects, which are likely not to be negligible in the absorption phenomenon, of these ionic species in solution. Nevertheless, the experimental λ_{max} values of $[\text{LuPc}_2]^+$ and $[\text{LuPc}_2]^-$ could be easily indexed to calculated transitions. They are associated with the same excitations as for the neutral form. In the particular case of $[\text{LuPc}_2]^+$, a weak absorption band (not reported in Table 3) appears in the RV

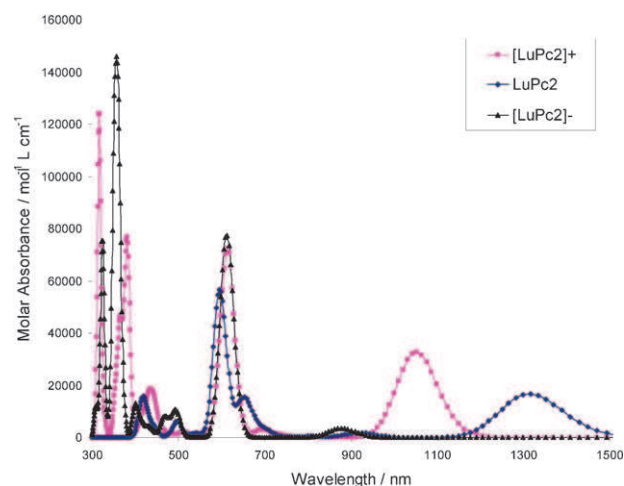


Fig. 9 Simulated absorption spectra of $(\text{LuPc}_2)^{+/0/-1}$. The Soret band was not calculated for the neutral species.

region of the experimental spectrum. It is assigned by the calculations to a forbidden $1b_2 \rightarrow 2a_2$ transition. The chlorine-substituted series exhibits similar features in their simulated spectra (Table 3 and Fig. 9).

Conclusion

This paper reports the first theoretical investigation on a series of unsymmetrical bisphthalocyanine lutetium complexes. The electronic structure of $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]$ and $\text{PcLu}[\text{PcCl}_{16}]$ indicates that, contrarily to the symmetrical LuPc_2 complex which exhibits an unpaired electron equally delocalized on both phthalocyanine ligands, the unsymmetrical ones have their unpaired electron predominantly localized on the non-substituted ligand. Although this fact can be rationalized on the basis of orbital interaction arguments, it is at first sight counter-intuitive since the chlorinated phthalocyanine is the most electron withdrawing ligand.

Table 3 Experimental λ_{max} values (in nm) recorded in dichloromethane for $(\text{LuPc}_2)^{+/0/-1}$, together with the corresponding computed transitions for $(\text{LuPc}_2)^{+/0/-1}$, $\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]^{+/0/-1}$, $\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]^{+/0/-1}$ and $\text{PcLu}[\text{PcCl}_{16}]^{+/0/-1}$. The first value (when available) is experimental; the bold value is computed and its oscillator strength $\times 10^{-5}$ is given in parenthesis

	Soret (experimental values only)	BV	Q	RV	IVB
LuPc_2					
Cationic	276, 312	476/ 434 (5)	625, 694/ 614 (32)	918/ 920 (1)	1051 (15)
Neutral	320	456/ 497 (2)	595, 658/ 596 (24), 652 (7)	—/ 878 (2)	1314 (8)
Anionic	275, 334		618, 698/ 612 (36)		
$\text{PcLu}[\text{Pc}(\text{p-Cl}_8)]$					
Cationic		598 (2), 591 (1)	653 (17)	928 (1)	1087 (14)
Neutral		486 (2), 473 (1)	619 (32), 690 (4)	826 (1)	1332 (7)
Anionic			624 (39)		
$\text{PcLu}[\text{Pc}(\text{np-Cl}_8)]$					
Cationic		553 (3), 573 (5)	624 (6), 682 (11)	717 (15)	1423 (11)
Neutral		502 (1), 521 (3)	607 (7), 656 (10)	599 (5)	1679 (11)
Anionic		474 (8)			
$\text{PcLu}[\text{PcCl}_{16}]$					
Cationic		562 (6), 585 (3),	799 (10), 687 (13)	737 (8), 732 (1)	1480 (9)
Neutral		513 (1), 518 (3), 525(2)	663 (12)		1701 (7)
Anionic		478 (3), 513 (2), 542 (5)	665 (39)		

Computational details

Density functional theory (DFT) calculations were carried out on the studied compounds using the Amsterdam Density Functional (ADF) program³⁵ developed by Baerends and coworkers.³⁶ Electron correlation was treated within the local density approximation (LDA) in the Vosko–Wilk–Nusair parametrization.³⁷ The non-local corrections of Becke and Perdew were added to the exchange and correlation energies, respectively.^{38,39} The numerical integration procedure applied for the calculations was developed by te Velde *et al.*^{36e} Relativistic corrections were added using the Zeroth Order Regular Approximation (ZORA) scalar Hamiltonian.⁴⁰ The atom electronic configurations were described by a triple- ζ Slater-type orbital (STO) basis set for H 1s, C 2s and 2p, N 2s and 2p and Cl 3s and 3p augmented with a 3d single- ζ polarization function for C, N and Cl atoms and with a 2p single- ζ polarization function for H. A triple- ζ STO basis set was used for Lu 4f, 5d and 6s, a double- ζ basis for Lu 5s and 5p augmented with a single- ζ 6p polarization function for Lu atom. A frozen-core approximation was used to treat the core shells up to 1s for C and N, 2p for Cl and 4d for Lu. Full geometry optimizations were carried out using the analytical gradient method implemented by Verluis and Ziegler.⁴¹ Spin-unrestricted calculations were performed for all the open-shell systems. The fragment interaction energy decomposition was made according to the method proposed by Ziegler and coworkers.⁴² In the case of the $\text{PcLu}[\text{PcCl}_n]^+$ cations, the considered fragments were Lu^{+3} and $\{\text{Pc} + [\text{PcCl}_n]\}^{2-}$, the electron configuration of the latter being $(1a_2)^2(2a_2)^0$ (see Scheme 2). In the case of the $\text{PcLu}[\text{PcCl}_n]^-$ anions, the considered fragments were Lu^{+3} and $\{\text{Pc} + [\text{PcCl}_n]\}^{4-}$, the electron configuration of the latter being $(1a_2)^2(2a_2)^2$.

The UV-visible transitions were calculated by means of time-dependent DFT (TDDFT) calculations,⁴³ at the same level of theory. Only spin-allowed transitions have been taken into account. Moreover, only transitions with non negligible oscillator strengths are reported and discussed.

Representation of the molecular structures were done using MOLEKEL4.1.⁴⁴ The UV/Visible spectra have been simulated from the computed TDDFT data and their oscillated strengths by using the SWizard program,⁴⁵ each transition being associated with a Gaussian function of half-height width equal to 1000 cm^{-1} .

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