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Inverse scattering at high energies for the multidimensional Newton equation in a long range potential

Alexandre Jollivet

June 15, 2013

Abstract

We define scattering data for the Newton equation in a potential $V \in C^2(\mathbb{R}^n, \mathbb{R})$, $n \geq 2$, that decays at infinity like $r^{-\alpha}$ for some $\alpha \in (0, 1]$. We provide estimates on the scattering solutions and scattering data and we prove, in particular, that the scattering data at high energies uniquely determine the short range part of the potential up to the knowledge of the long range tail of the potential. The Born approximation at fixed energy of the scattering data is also considered. We then change the definition of the scattering data to study inverse scattering in other asymptotic regimes. These results were obtained by developing the inverse scattering approach of [Novikov, 1999].

1 Introduction

Consider the multidimensional Newton equation in an external static force F deriving from a scalar potential V :

$$\ddot{x}(t) = F(x(t)) = -\nabla V(x(t)), \quad (1.1)$$

where $x(t) \in \mathbb{R}^n$, $\dot{x}(t) = \frac{dx}{dt}(t)$, $n \geq 2$.

When $n = 3$ then equation (1.1) is the equation of motion of a nonrelativistic particle of mass $m = 1$ and charge $e = 1$ in an external and static electric (or gravitational) field described by V (see [8]) where x denotes the position of the particle, \dot{x} denotes its velocity and \ddot{x} denotes its acceleration and t denotes the time.

We also assume throughout this paper that V satisfies the following conditions

$$F = F^l + F^s, \quad (1.2)$$

where $F^l := -\nabla V^l$, $F^s := -\nabla V^s$ and $(V^l, V^s) \in (C^2(\mathbb{R}^n, \mathbb{R}))^2$, and where V^l satisfies the following long range assumptions

$$|\partial_x^j V^l(x)| \leq \beta_{|j|}^l (1 + |x|)^{-(\alpha + |j|)}, \quad (1.3)$$

and V^s satisfies the following short range assumptions

$$|\partial_x^j V^s(x)| \leq \beta_{|j|+1}^s (1 + |x|)^{-(\alpha + 1 + |j|)}, \quad (1.4)$$

for $x \in \mathbb{R}^n$ and $|j| \leq 2$ and for some $\alpha \in (0, 1]$ (here j is the multiindex $j \in (\mathbb{N} \cup \{0\})^n$, $|j| = \sum_{m=1}^n j_m$, and β_m^l and $\beta_{m'}^s$ are positive real constants for $m = 0, 1, 2$ and for $m' = 1, 2, 3$). Note that the assumption $0 < \alpha \leq 1$ includes the decay rate of a Coulombian potential at infinity. Indeed for a Coulombian potential $V(x) = \frac{1}{|x|}$, estimates (1.3) are satisfied uniformly for $|x| > \varepsilon$ and $\alpha = 1$ for any $\varepsilon > 0$. Although our potentials (V^l, V^s) are assumed to be C^2 on the entire space \mathbb{R}^n , our present work may provide interesting results even in presence of singularities for the potentials (V^l, V^s) .

For equation (1.1) the energy

$$E = \frac{1}{2} |\dot{x}(t)|^2 + V(x(t)) \quad (1.5)$$

is an integral of motion.

For $\sigma > 0$ we denote by $\mathcal{B}(0, \sigma)$ the Euclidean open ball of center 0 and radius σ , $\mathcal{B}(0, \sigma) = \{y \in \mathbb{R}^n \mid |y| < \sigma\}$, and we denote by $\overline{\mathcal{B}(0, \sigma)} = \{y \in \mathbb{R}^n \mid |y| \leq \sigma\}$ its closure. We set $\mu := \sqrt{\frac{2^5 n \max(\beta_1^l, \beta_2^l)}{\alpha}}$. Under conditions (1.3) the following is valid (see Lemma 2.1 given in the next Section): for any $v \in \mathbb{R}^n \setminus \mathcal{B}(0, \mu)$, there exists a unique solution $z_{\pm}(v, \cdot)$ of the equation

$$\ddot{z}(t) = F^l(z(t)), \quad t \in \mathbb{R}, \quad (1.6)$$

so that

$$\dot{z}_{\pm}(v, t) - v = o(1), \quad \text{as } t \rightarrow \pm\infty, \quad z_{\pm}(v, 0) = 0,$$

and

$$|z_{\pm}(v, t) - tv| \leq \frac{2^{\frac{5}{2}} n^{\frac{1}{2}} \beta_1^l}{\alpha |v|} |t| \quad \text{for } t \in \mathbb{R}.$$

When $F^l \equiv 0$ then β_1^l , β_2^l and μ can be arbitrary close to 0, and we have that $z_{\pm}(v, t) = tv$ for $(t, v) \in \mathbb{R} \times \mathbb{R}^n$, $v \neq 0$.

Then under conditions (1.3) and (1.4), the following is valid: for any $(v_-, x_-) \in \mathbb{R}^n \setminus \mathcal{B}(0, \mu) \times \mathbb{R}^n$, the equation (1.1) has a unique solution $x \in C^2(\mathbb{R}, \mathbb{R}^n)$ such that

$$x(t) = z_-(v_-, t) + x_- + y_-(t), \quad (1.7)$$

where $|\dot{y}_-(t)| + |y_-(t)| \rightarrow 0$, as $t \rightarrow -\infty$; in addition for almost any $(v_-, x_-) \in \mathbb{R}^n \setminus \mathcal{B}(0, \mu) \times \mathbb{R}^n$,

$$x(t) = z_+(v_+, t) + x_+ + y_+(t), \quad (1.8)$$

for a unique $(v_+, x_+) \in \mathbb{R}^n \times \mathbb{R}^n$, where $|v_+| = |v_-| \geq \mu$ by conservation of the energy (1.5), and where $v_+ =: a(v_-, x_-)$, $x_+ =: b(v_-, x_-)$, and $|\dot{y}_+(t)| + |y_+(t)| \rightarrow 0$, as $t \rightarrow +\infty$. A solution x of (1.1) that satisfies (1.7) and (1.8) for some (v_-, x_-) , $v_- \neq 0$, is called a scattering solution.

We call the map $S : (\mathbb{R}^n \setminus \mathcal{B}(0, \mu)) \times \mathbb{R}^n \rightarrow (\mathbb{R}^n \setminus \mathcal{B}(0, \mu)) \times \mathbb{R}^n$ given by the formulas

$$v_+ = a(v_-, x_-), \quad x_+ = b(v_-, x_-), \quad (1.9)$$

the scattering map for the equation (1.1). In addition, $a(v_-, x_-)$, $b(v_-, x_-)$ are called the scattering data for the equation (1.1), and we define

$$a_{sc}(v_-, x_-) = a(v_-, x_-) - v_-, \quad b_{sc}(v_-, x_-) = b(v_-, x_-) - x_-. \quad (1.10)$$

Our definition of the scattering map is derived from constructions given in [6, 1]. We refer the reader to [6, 1] and references therein for the forward classical scattering theory.

By $\mathcal{D}(S)$ we denote the set of definition of S . Under the conditions (1.3) and (1.4) the map $S : \mathcal{D}(S) \rightarrow (\mathbb{R}^n \setminus \mathcal{B}(0, \mu)) \times \mathbb{R}^n$ is continuous, and $\text{Mes}((\mathbb{R}^n \setminus \mathcal{B}(0, \mu)) \times \mathbb{R}^n \setminus \mathcal{D}(S)) = 0$ for the Lebesgue measure on $\mathbb{R}^n \times \mathbb{R}^n$. In addition the map S is uniquely determined by its restriction to $\mathcal{M}(S) = \mathcal{D}(S) \cap \mathcal{M}$ and by F^l , where $\mathcal{M} = \{(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n \mid v_- \neq 0, v_- \cdot x_- = 0\}$. (Indeed if $x(t)$ is a solution of (1.1) then $x(t + t_0)$ is also a solution of (1.1) for any $t_0 \in \mathbb{R}$.)

One can imagine the following experimental setting that allows to measure the scattering data without knowing the potential V inside a (a priori bounded) region of interest. First find a potential V^l that generates the same long range effects as V does. Then compute the solutions $z_{\pm}(v, \cdot)$ of equation (1.6). Then for a fixed $(v_-, x_-) \in (\mathbb{R}^n \setminus \mathcal{B}(0, \mu)) \times \mathbb{R}^n$ send a particle far away from the region of interest with a trajectory asymptotic to $x_- + z_-(v_-, \cdot)$ at large and negative times. When the particle escapes any bounded region of the space at finite time, then detect the particle and find $S(v_-, x_-) = (v_+, x_+)$ so that the trajectory of the particle is asymptotic to $x_+ + z_+(v_+, \cdot)$ at large and positive times far away from the bounded region of interest.

In this paper we consider the following inverse scattering problem for equation (1.1):

$$\text{Given } S \text{ and given the long range tail } F^l \text{ of the force } F, \text{ find } F^s. \quad (1.11)$$

The main results of the present work consist in estimates and asymptotics for the scattering data (a_{sc}, b_{sc}) and scattering solutions for the equation (1.1)

and in application of these asymptotics and estimates to the inverse scattering problem (1.11) at high energies. Our main results include, in particular, Theorem 1.1 given below that provides the high energies asymptotics of the scattering data and the Born approximation of the scattering data at fixed energy.

Consider

$$T\mathbb{S}^{n-1} := \{(\theta, x) \in \mathbb{S}^{n-1} \times \mathbb{R}^n \mid \theta \cdot x = 0\},$$

and for any $m \in \mathbb{N}$ consider the x-ray transform P defined by

$$Pf(\theta, x) := \int_{-\infty}^{+\infty} f(t\theta + x)dt$$

for any function $f \in C(\mathbb{R}^n, \mathbb{R}^m)$ so that $|f(x)| = O(|x|^{-\tilde{\beta}})$ as $|x| \rightarrow +\infty$ for some $\tilde{\beta} > 1$. For $(\sigma, \tilde{\beta}, r, \tilde{\alpha}) \in (0, +\infty)^2 \times (0, 1) \times (0, 1]$, let $s_0 = s_0(\sigma, r, \tilde{\beta}, \tilde{\alpha})$ be defined as the root of the equation

$$1 = \frac{4\tilde{\beta}n(\sigma + 1)}{\tilde{\alpha}r(\frac{s_0}{2^{\frac{3}{2}}} - r)(1 - r)^{\tilde{\alpha}+2}} \left(1 + \frac{1}{\frac{s_0}{2^{\frac{3}{2}}} - r}\right)^2, \quad s_0 > 2^{\frac{3}{2}}r. \quad (1.12)$$

Then we have the following results.

Theorem 1.1. *Let $(\theta, x) \in T\mathbb{S}^{n-1}$. Under conditions (1.3) and (1.4) the following limits are valid*

$$\lim_{s \rightarrow +\infty} sa_{sc}(s\theta, x) = PF^l(\theta, x) + PF^s(\theta, x), \quad (1.13)$$

$$\lim_{s \rightarrow +\infty} s^2\theta \cdot \left(b_{sc}(s\theta, x) - W(s\theta, x)\right) = -PV^s(\theta, x), \quad (1.14)$$

where

$$\begin{aligned} W(v, x) &:= \int_{-\infty}^0 \int_{-\infty}^{\sigma} (F^l(z_-(v, \tau) + x) - F^l(z_-(v, \tau)))d\tau d\sigma \\ &\quad - \int_0^{+\infty} \int_{\sigma}^{+\infty} (F^l(z_+(a(v, x), \tau) + x) - F^l(z_+(a(v, x), \tau)))d\sigma d\tau, \end{aligned} \quad (1.15)$$

for $(v, x) \in \mathcal{D}(S)$.

In addition,

$$\left|a_{sc}(s\theta, x) - \int_{-\infty}^{+\infty} F(\tau s\theta + x)d\tau\right| \leq \frac{4n^2(3|x| + 5)\beta^2 \left(1 + \frac{1}{2^{-\frac{3}{2}}s - r}\right)^2}{\alpha^2(1 - r)^{2\alpha+3}(2^{-\frac{3}{2}}s - r)^2}, \quad (1.16)$$

$$\begin{aligned}
& \left| b_{sc}(s\theta, x) - W(s\theta, x) - \int_{-\infty}^0 \int_{-\infty}^{\sigma} F^s(\tau s\theta + x) d\tau d\sigma \right. \\
& \left. + \int_0^{+\infty} \int_{\sigma}^{+\infty} F^s(\tau s\theta + x) d\tau d\sigma \right| \leq \frac{4n^2(3|x| + 5)\beta^2 \left(1 + \frac{1}{2^{-\frac{3}{2}}s-r}\right)^2}{\alpha^2(1-r)^{2\alpha+2}(2^{-\frac{3}{2}}s-r)^3}, \quad (1.17)
\end{aligned}$$

for $(r, (\theta, x)) \in (0, 1) \times T\mathbb{S}^{n-1}$ and for $s > s_0(|x|, r, \beta, \alpha)$, where $\beta = \max(\beta_1^l, \beta_2^l, \beta_2^s, \beta_3^s)$.

Note that the vector W defined by (1.15) is known from the scattering data and from F^l . Then from (1.13) (resp. (1.14)) and inversion formulas for the X-ray transform for $n \geq 2$ (see [12, 4, 9, 10]) it follows that F^s can be reconstructed from a_{sc} (resp. b_{sc}).

Note that (1.16) and (1.17) also give the asymptotics of a_{sc} , b_{sc} , when the parameters α , n , s , θ and x are fixed and β decreases to 0 (where $\beta = \max(\beta_1^l, \beta_2^l, \beta_2^s, \beta_3^s)$). In that regime the leading term of $sa_{sc}(s\theta, x)$ and $s^2\theta \cdot (b_{sc}(s\theta, x) - W(s\theta, x))$ for $(\theta, x) \in T\mathbb{S}^{n-1}$ and for $s > s_0(|x|, r, \beta, \alpha)$ is given by the right hand sides of (1.13) and (1.14) respectively. Therefore Theorem 1.1 gives the Born approximation for the scattering data at fixed energy when the potential is sufficiently weak, and it proves that F^s can be reconstructed from the Born approximation of the scattering map at fixed energy.

Theorem 1.1 is a generalization of [10, formulas (4.8a), (4.8b), (4.9a) and (4.9b)] where inverse scattering for the classical multidimensional Newton equation was studied in the short range case ($F^l \equiv 0$). We develop Novikov's framework [10] to obtain our results. Note that results [10, formulas (4.8b) and (4.9b)] also provide the approximation of the scattering data $(a_{sc}(v_-, x_-), b_{sc}(v_-, x_-))$ for the short range case ($F^l \equiv 0$) when the parameters α , n , v_- and β are fixed and $|x_-| \rightarrow +\infty$. Such an asymptotic regime is not covered by Theorem 1.1. Therefore we shall modify in Section 3 the definition of the scattering map to study these modified scattering data in the following three asymptotic regimes: at high energies, Born approximation at fixed energy, and when the parameters α , n , v_- and β are fixed and $|x_-| \rightarrow +\infty$.

For inverse scattering at fixed energy for the multidimensional Newton equation, see for example [7] and references therein.

For the inverse scattering problem in quantum mechanics for the Schrödinger equation, see for example [3], [2], [11] and references given in [11].

Our paper is organized as follows. In Section 2 we transform the differential equation (1.1) with initial conditions (1.7) in an integral equation which takes the form $y_- = A(y_-)$. Then we study the operator A on a suitable space (Lemma 2.2) and we give estimates for the deflection $y_-(t)$ in (1.7) and for the scattering data $a_{sc}(v_-, x_-), b_{sc}(v_-, x_-)$ (Theorem 2.4). We prove

Theorem 1.1. Note that we work with small angle scattering compared to the dynamics generated by F^l through the “free” solutions $z_-(v_-, t)$: In particular, the angle between the vectors $\dot{x}(t) = \dot{z}_-(v_-, t) + \dot{y}_-(t)$ and $\dot{z}_-(v_-, t)$ goes to zero when the parameters β , α , n , $v_-/|v_-|$, x_- are fixed and $|v_-|$ increases. We also provide similar results when one replaces the “free” solutions $z_-(v, \cdot)$ by some other functions “ $z_{-,N}(v, \cdot)$ ” that may be easier to compute in practise (Formulas (2.47) and (2.48)). In Section 3 we change the definition of the scattering map so that one can obtain for the modified scattering data $(\tilde{a}_{sc}(v_-, x_-), \tilde{b}_{sc}(v_-, x_-))$ their approximation at high energies, or their Born approximation at fixed energy, or their approximation when the parameters α , n , v_- and β are fixed and $|x_-| \rightarrow +\infty$ (Theorems 3.3, 3.4 and formulas (3.40) and (3.41)). Sections 4, 5, 6 and 7 are devoted to proofs of our Theorems and Lemmas.

2 Scattering solutions

2.1 Integral equation

First we need the following Lemma 2.1 that generalizes the statements given in the Introduction on the existence of peculiar solutions z_\pm of the equation (1.6).

Lemma 2.1. *Assume conditions (1.3). Let $(v, x, w, h) \in (\mathbb{R}^n)^4$ so that $v \cdot x = 0$ and*

$$|w - v| \leq \frac{|v|}{4\sqrt{2}} \text{ and } |h| < 1 + \frac{|x|}{\sqrt{2}}, \quad (2.1)$$

and assume

$$\frac{2^5 n \max(\beta_1^l, \beta_2^l)}{\alpha |v|^2 (1 + \frac{|x|}{\sqrt{2}} - |h|)^\alpha} \leq 1. \quad (2.2)$$

Then there exists a unique solution $z_\pm(w, x + h, \cdot)$ of the equation (1.6) so that

$$\dot{z}_\pm(w, x + h, t) - w = o(1), \text{ as } t \rightarrow \pm\infty, \quad z_\pm(w, x + h, 0) = x + h, \quad (2.3)$$

and

$$|z_\pm(w, x + h, t) - x - h - tw| \leq \frac{2^{\frac{5}{2}} n^{\frac{1}{2}} \beta_1^l}{\alpha |v| (1 + \frac{|x|}{\sqrt{2}} - |h|)^\alpha} |t|, \quad (2.4)$$

for $t \in \mathbb{R}$.

A proof of Lemma 2.1 is given in Section 4.

For the rest of this Section we set

$$\mu := \sqrt{\frac{2^5 n \max(\beta_1^l, \beta_2^l)}{\alpha}}, \quad (2.5)$$

$$z_{\pm}(v, t) = z_{\pm}(v, 0, t) \text{ for } t \in \mathbb{R}, \text{ when } \mu \leq |v|, \quad (2.6)$$

$$\beta_2 := \max(\beta_2^l, \beta_2^s). \quad (2.7)$$

Let $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $v_- \cdot x_- = 0$ and $|v_-| \geq \mu$. Then the function y_- in (1.7) satisfies the integral equation $y_- = A(y_-)$ where

$$A(f)(t) = \int_{-\infty}^t \int_{-\infty}^{\sigma} \left(F(z_-(v_-, \tau) + x_- + f(\tau)) - F^l(z_-(v_-, \tau)) \right) d\tau d\sigma \quad (2.8)$$

for $t \in \mathbb{R}$ and for $f \in C(\mathbb{R}, \mathbb{R}^n)$, $\sup_{(-\infty, 0]} |f| < \infty$. Under conditions (1.3) and (1.4) we have $A(f) \in C^2(\mathbb{R}, \mathbb{R}^n)$ for $f \in C(\mathbb{R}, \mathbb{R}^n)$ so that $\sup_{(-\infty, 0]} |f| < \infty$.

For $r > 0$ we introduce the following complete metric space M_r defined by

$$M_r = \{f \in C(\mathbb{R}, \mathbb{R}^n) \mid \sup_{(-\infty, 0]} |f| + \sup_{t \in [0, +\infty)} \left(\frac{|f(t)|}{1 + |t|} \right) \leq r\}, \quad (2.9)$$

and endowed with the norm $\|\cdot\|$ where $\|f\| = \sup_{(-\infty, 0]} |f| + \sup_{t \in [0, +\infty)} \left(\frac{|f(t)|}{1 + |t|} \right)$. Then we have the following estimate and contraction estimate for the map A restricted to M_r .

Lemma 2.2. *Let $(v_-, x_-) \in (\mathbb{R}^n \setminus \mathcal{B}(0, \mu)) \times \mathbb{R}^n$, $v_- \cdot x_- = 0$, and let $r > 0$, $r < \max(\frac{|v_-|}{2^{\frac{3}{2}}}, 1 + \frac{|x_-|}{\sqrt{2}})$. Then the following estimates are valid:*

$$\begin{aligned} \|A(f)\| &\leq \rho(n, \alpha, \beta_2, |x_-|, |v_-|, r) \\ &:= \frac{\beta_2(n(3|x_-| + 2r) + 2\sqrt{n})}{(\frac{|v_-|}{2^{\frac{3}{2}}} - r)(1 - r)^{\alpha}} \left(\frac{2}{\alpha(\frac{|v_-|}{2^{\frac{3}{2}}} - r)} + \frac{1}{(\alpha + 1)(1 - r)} \right), \end{aligned} \quad (2.10)$$

and

$$\|A(f_1) - A(f_2)\| \leq \lambda(n, \alpha, \beta_2, \beta_3^s, |x_-|, |v_-|, r) \|f_1 - f_2\|, \quad (2.11)$$

$$\begin{aligned} \lambda(n, \alpha, \beta_2, \beta_3^s, |x_-|, |v_-|, r) &:= \frac{2n}{\alpha(\frac{|v_-|}{2^{\frac{3}{2}}} - r)(1 - r + \frac{|x_-|}{\sqrt{2}})^{\alpha}} \left(\beta_2 + \frac{\beta_3^s}{1 - r + \frac{|x_-|}{\sqrt{2}}} + \frac{\beta_3^s}{\frac{|v_-|}{2^{\frac{3}{2}}} - r} \right) \\ &\quad \times \left(\frac{1}{1 - r + \frac{|x_-|}{\sqrt{2}}} + \frac{1}{\frac{|v_-|}{2^{\frac{3}{2}}} - r} \right), \end{aligned}$$

for $(f, f_1, f_2) \in M_r^3$.

A proof of Lemma 2.2 is given in Section 5.

We also need the following result.

Lemma 2.3. *Let $(v_-, x_-) \in (\mathbb{R}^n \setminus \mathcal{B}(0, \mu)) \times \mathbb{R}^n$, $v_- \cdot x_- = 0$, and let $r > 0$, $r < \max(\frac{|v_-|}{2^{\frac{3}{2}}}, 1 + \frac{|x_-|}{\sqrt{2}})$. When $y_- \in M_r$ is a fixed point for the map A then $z_-(v_-, \cdot) + x_- + y_-(\cdot)$ is a scattering solution for equation (1.1) and*

$$z_-(v_-, t) + x_- + y_-(t) = z_+(a(v_-, x_-), t) + b(v_-, x_-) + y_+(t), \quad (2.12)$$

for $t \geq 0$, where

$$a(v_-, x_-) := v_- + \int_{-\infty}^{+\infty} F(z_-(v_-, \tau) + x_- + y_-(\tau)) d\tau, \quad (2.13)$$

$$b(v_-, x_-) := x_- + l(v_-, x_-, y_-) + l_1(v_-, x_-) + l_2(v_-, x_-, y_-), \quad (2.14)$$

$$y_+(t) := \int_t^{+\infty} \int_{\sigma}^{+\infty} (F(z_-(v_-, \tau) + x_- + y_-(\tau)) - F^l(z_+(a(v_-, x_-), \tau))) d\tau d\sigma, \quad (2.15)$$

for $t \geq 0$, and where

$$\begin{aligned} l(v_-, x_-, y_-) &:= \int_{-\infty}^0 \int_{-\infty}^{\sigma} (F(z_-(v_-, \tau) + x_- + y_-(\tau)) - F^l(z_-(v_-, \tau))) d\tau d\sigma \\ &\quad - \int_0^{+\infty} \int_{\sigma}^{+\infty} F^s(z_-(v_-, \tau) + x_- + y_-(\tau)) d\tau d\sigma, \end{aligned} \quad (2.16)$$

$$l_1(v_-, x_-) := - \int_0^{+\infty} \int_{\sigma}^{+\infty} (F^l(z_+(a(v_-, x_-), \tau) + x_-) - F^l(z_+(a(v_-, x_-), \tau))) d\tau d\sigma, \quad (2.17)$$

$$l_2(v_-, x_-, y_-) := - \int_0^{+\infty} \int_{\sigma}^{+\infty} (F^l(z_-(v_-, \tau) + x_- + y_-(\tau)) - F^l(z_+(a(v_-, x_-), \tau) + x_-)) d\tau d\sigma, \quad (2.18)$$

for $t \geq 0$.

Lemma 2.3 is proved in Section 4. Note that l_1 is known from the scattering data and the knowledge of F^l .

2.2 Estimates on the scattering solutions

In this Section our main results consist in estimates and asymptotics for the scattering data (a_{sc}, b_{sc}) and scattering solutions for the equation (1.1).

Theorem 2.4. *Under the assumptions of Lemma 2.3 the following estimates are valid*

$$|\dot{y}_-(t)| \leq \frac{\beta_2(n(|x_-| + r) + \sqrt{n})}{(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\left(1 - r + |t|\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+1}}, \quad (2.19)$$

$$|y_-(t)| \leq \frac{\beta_2(n(|x_-| + r) + \sqrt{n})}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 - r + |t|\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^\alpha}, \quad (2.20)$$

for $t \leq 0$. In addition

$$|a_{sc}(v_-, x_-)| \leq \frac{2n^{\frac{1}{2}}}{\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} \left(\frac{\beta_1^l}{\alpha} + \frac{\beta_2}{(\alpha + 1)\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)} \right). \quad (2.21)$$

$$|l(v_-, x_-, y_-)| \leq \frac{\beta_2 n^{\frac{1}{2}}(n^{\frac{1}{2}}(|x_-| + r) + 2)}{\alpha(\alpha + 1)(1 - r)^\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2}, \quad (2.22)$$

and

$$\begin{aligned} & \left| a_{sc}(v_-, x_-) - \int_{-\infty}^{+\infty} F(z_-(v_-, \tau) + x_-) d\tau \right| \\ & \leq \frac{4 \max(\beta_2, \beta_3^s)^2 n^{\frac{3}{2}} (n^{\frac{1}{2}}(3|x_-| + 2r) + 2)}{\alpha^2 \left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 (1 - r)^{2\alpha+3}} \left(\frac{1}{\frac{|v_-|}{2\sqrt{2}} - r} + 1 \right)^2, \end{aligned} \quad (2.23)$$

$$\begin{aligned} |l(v_-, x_-, y_-) - l(v_-, x_-, 0)| & \leq \frac{4 \max(\beta_2, \beta_3^s)^2 n^{\frac{3}{2}} (n^{\frac{1}{2}}(3|x_-| + 2r) + 2)}{\alpha^2(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^3 (1 - r)^{2\alpha+2}} \\ & \quad \times \left(\frac{1}{\frac{|v_-|}{2\sqrt{2}} - r} + 1 \right)^2. \end{aligned} \quad (2.24)$$

In addition when

$$\frac{8n \max(\beta_1^l, \beta_2)}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 (1 - r)^{\alpha+1}} \leq 1, \quad (2.25)$$

then

$$|l_1(v_-, x_-)| \leq \frac{8\beta_2 n |x_-|}{\alpha(\alpha+1)|v_-|^2}, \quad (2.26)$$

$$|l_2(v_-, x_-, y_-)| \leq \frac{2n^{\frac{3}{2}}\beta_2^2(n^{\frac{1}{2}}(2|x_-|+r)+3)}{\alpha^2(\alpha+1)^2(1-r)^{2\alpha}(\frac{|v_-|}{2^{\frac{3}{2}}}-r)^4}, \quad (2.27)$$

$$\begin{aligned} |y_+(t)| &\leq \frac{2n^{\frac{1}{2}}\beta_2}{\alpha(\alpha+1)(\frac{|v_-|}{2^{\frac{3}{2}}}-r)^2(1-r+\frac{|x_-|}{\sqrt{2}}+t(\frac{|v_-|}{2^{\frac{3}{2}}}-r))^\alpha} \\ &\times \left(1 + \frac{2n\beta_2(n^{\frac{1}{2}}(2|x_-|+r)+3)}{\alpha(\alpha+1)(\frac{|v_-|}{2^{\frac{3}{2}}}-r)^2(1-r)^\alpha}\right), \end{aligned} \quad (2.28)$$

for $t \geq 0$.

A proof of Theorem 2.4 is given in Section 6. We now prove Theorem 1.1 combining Theorem 2.4, Lemma 2.2 and estimate (2.4).

2.3 Proof of Theorem 1.1

Let $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $v_- \cdot x_- = 0$ and $|v_-| \geq \mu$. We first prove estimates (2.32) and (2.34) given below. We use the following estimate (2.29)

$$\begin{aligned} |x_- + \eta v_- \tau + (1-\eta)z_-(v_-, \tau)| &\geq |x_- + \tau v_-| - |z_-(v_-, \tau) - \tau v_-| \\ &\geq \frac{|x_-|}{\sqrt{2}} + \left(\frac{|v_-|}{\sqrt{2}} - \frac{2^{\frac{5}{2}}n^{\frac{1}{2}}\beta_1^l}{\alpha|v_-|}\right)|\tau| \geq \frac{|x_-|}{\sqrt{2}} + |\tau|\frac{|v_-|}{2^{\frac{3}{2}}}, \end{aligned} \quad (2.29)$$

for $\eta \in (0, 1)$ and $\tau \in \mathbb{R}$ (we used (2.4) and (2.2)). Then from (1.3), (1.4), (2.29) and (2.4) it follows that

$$\begin{aligned} |F^s(z_-(v_-, \tau) + x_-) - F^s(\tau v_- + x_-)| &\leq \sup_{\eta \in (0,1)} \frac{n\beta_3^s |z_-(v_-, \tau) - \tau v_-|}{(1 + |x_- + \eta v_- + (1-\eta)z_-(v_-, \tau)|)^{\alpha+3}} \\ &\leq \frac{2^{\frac{5}{2}}n^{\frac{3}{2}}\beta_3^s\beta_1^l|\tau|}{\alpha|v_-|(1 + \frac{|x_-|}{\sqrt{2}} + |\tau|\frac{|v_-|}{2^{\frac{3}{2}}})^{\alpha+3}}, \end{aligned} \quad (2.30)$$

for $\tau \in \mathbb{R}$. Similarly

$$|F^l(z_-(v_-, \tau) + x_-) - F^l(\tau v_- + x_-)| \leq \frac{2^{\frac{5}{2}}n^{\frac{3}{2}}|\tau|\beta_2\beta_1^l}{\alpha|v_-|(1 + \frac{|x_-|}{\sqrt{2}} + |\tau|\frac{|v_-|}{2^{\frac{3}{2}}})^{\alpha+2}}, \quad (2.31)$$

for $\tau \in \mathbb{R}$. Then using (2.30) and (2.31) we have

$$\left| \int_{-\infty}^{+\infty} F(z_-(v_-, \tau) + x_-) d\tau - \int_{-\infty}^{+\infty} F(\tau v_- + x_-) d\tau \right|$$

$$\leq \frac{2^{\frac{7}{2}} n^{\frac{3}{2}} \max(\beta_2, \beta_3^s) \beta_1^l}{\alpha |v_-|} \int_{-\infty}^{+\infty} \frac{|\tau| d\tau}{(1 + \frac{|x_-|}{\sqrt{2}} + \frac{|v_-|}{2^{\frac{3}{2}}} |\tau|)^{\alpha+2}} \leq \frac{2^{\frac{15}{2}} n^{\frac{3}{2}} \max(\beta_2, \beta_3^s) \beta_1^l}{\alpha^2 |v_-|^3 (1 + \frac{|x_-|}{\sqrt{2}})^{\alpha}}. \quad (2.32)$$

Set

$$\Delta_1(v_-, x_-) = \int_{-\infty}^0 \int_{-\infty}^{\sigma} F^s(z_-(v_-, \tau) + x_-) d\tau d\sigma - \int_0^{+\infty} \int_{\sigma}^{+\infty} F^s(z_-(v_-, \tau) + x_-) d\tau d\sigma. \quad (2.33)$$

Then using (2.30) we have

$$\begin{aligned} & |\Delta_1(v_-, x_-) - \int_{-\infty}^0 \int_{-\infty}^{\sigma} F^s(\tau v_- + x_-) d\tau d\sigma + \int_0^{+\infty} \int_{\sigma}^{+\infty} F^s(\tau v_- + x_-) d\tau d\sigma| \\ & \leq \frac{2^{\frac{7}{2}} n^{\frac{3}{2}} \beta_1^l \beta_3^s}{\alpha |v_-|} \int_{-\infty}^0 \int_{-\infty}^{\sigma} \frac{|\tau| d\tau d\sigma}{(1 + \frac{|x_-|}{\sqrt{2}} + \frac{|v_-|}{2\sqrt{2}} |\tau|)^{\alpha+3}} \leq \frac{2^{\frac{13}{2}} n^{\frac{3}{2}} \beta_1^l \beta_3^s}{\alpha^2 (\alpha + 1) |v_-|^3 (1 + \frac{|x_-|}{\sqrt{2}})^{\alpha}}. \end{aligned} \quad (2.34)$$

Let $r > 0$, $r < \max(\frac{|v_-|}{2^{\frac{3}{2}}}, 1)$. Note that

$$\max\left(\frac{\rho}{r}, \lambda, \frac{8n \max(\beta_1^l, \beta_2)}{\alpha (\frac{|v_-|}{2^{\frac{3}{2}}} - r)^2 (1 - r)^{\alpha+1}}\right) \leq \frac{4\beta n(|x_-| + 1)}{\alpha r (\frac{|v_-|}{2^{\frac{3}{2}}} - r) (1 - r)^{\alpha+2}} \left(1 + \frac{1}{\frac{|v_-|}{2^{\frac{3}{2}}} - r}\right)^2, \quad (2.35)$$

where ρ and λ are defined by (2.10) and (2.11) respectively. Assume that $|v_-| > s_0(|x_-|, r, \beta, \alpha)$ where s_0 is the root of the equation (1.12). Then from (1.12) and Lemma 2.2 it follows that A has a unique fixed point in M_r denoted by y_- . Then adding (2.23) and (2.32) we obtain (1.16). Note also that

$$\begin{aligned} l(v_-, x_-, 0) &= \int_{-\infty}^0 \int_{-\infty}^{\sigma} (F^l(z_-(v_-, \tau) + x_-) - F^l(z_-(v_-, \tau))) d\tau d\sigma \\ &+ \Delta_1(v_-, x_-). \end{aligned}$$

Hence adding (2.34), (2.27) and (2.24) we obtain (1.17). Theorem 1.1 is proved. \square

2.4 Motivations for changing the definition of the scattering map

For a solution x at a nonzero energy for equation (1.1) we say that it is a scattering solution when there exists $\varepsilon > 0$ so that $1 + |x(t)| \geq \varepsilon(1 + |t|)$ for $t \in \mathbb{R}$ (see [1]). In the Introduction and in the previous subsections we choose to parametrize the scattering solutions of equation (1.1) by the solutions $z_{\pm}(v, \cdot)$

of the equation (1.6) (see the asymptotic behaviors (1.7) and (1.8)), and then to formulate the inverse scattering problem (1.11) using this parametrization. To compute the "free" solutions $z_{\pm}(v, \cdot)$ one has to integrate equation (1.6). For some cases solving (1.6) leads to simple exact formulas (see [8, Section 15] when F^l is a Coulombian force). In general one may choose to approximate the solutions $z_{\pm}(v, \cdot)$ by the functions $z_{\pm, N+1}(v, \cdot)$ defined below. In general the functions $z_{\pm, N+1}(v, \cdot)$ are easier to compute, and in this Subsection we use these approximations to obtain an other formulation of the inverse scattering problem and to mention results similar to Theorem 1.1 and to those given in the previous subsections.

Assume without loss of generality that $\alpha \notin \{\frac{1}{m} \mid m \in \mathbb{N}, m > 0\}$ and set $N = \lfloor \alpha^{-1} \rfloor$ the integer part of α^{-1} . Then let $(x, v, w, h) \in (\mathbb{R}^n)^4$ so that $v \cdot x = 0$ and (2.1) and (2.2) are satisfied. We define by induction (see also [6])

$$z_{\pm, 0}(w, x + h, t) = x + h + tw, \quad (2.36)$$

$$z_{-, m+1}(w, x + h, t) = x + h + tw + \int_0^t \int_{-\infty}^{\sigma} F^l(z_{-, m}(w, x + h, \tau)) d\tau, \quad (2.37)$$

$$z_{+, m+1}(w, x + h, t) = x + h + tw - \int_0^t \int_{\sigma}^{+\infty} F^l(z_{+, m}(w, x + h, \tau)) d\tau, \quad (2.38)$$

for $t \in \mathbb{R}$ and for $m = 0 \dots N$. Then one can prove the following estimates by induction (see the proof of Lemma 2.1 and see also [6])

$$|z_{\pm, m}(w, x + h, t) - wt - x - h| \leq \frac{2^{\frac{5}{2}} n^{\frac{1}{2}} \beta_1^l}{\alpha |v| (1 + \frac{|x|}{\sqrt{2}} - |h|)^{\alpha}} |t|, \quad t \in \mathbb{R}, \quad m = 1 \dots N+1, \quad (2.39)$$

$$\begin{aligned} |z_{\pm, m+1}(w, x + h, t) - z_{\pm, m}(w, x + h, t)| &\leq \frac{2^{3(m+1)} n^{m+\frac{1}{2}} (\beta_2^l)^m \beta_1^l}{\alpha^{m+1} |v|^{2m+2} \prod_{j=1}^{m+1} (1 - j\alpha) j} \\ &\times \left(\left(1 + \frac{|x|}{\sqrt{2}} - |h| + |t| \frac{|v|}{2^{\frac{3}{2}}} \right)^{1-(m+1)\alpha} - \left(1 + \frac{|x|}{\sqrt{2}} - |h| \right)^{1-(m+1)\alpha} \right), \end{aligned} \quad (2.40)$$

for $m = 0 \dots N-1$ and for $\pm t \geq 0$,

$$|z_{\pm, N+1}(w, x + h, t) - z_{\pm, N}(w, x + h, t)| \leq \frac{2^{3(N+1)} n^{N+\frac{1}{2}} (\beta_2^l)^N \beta_1^l}{\alpha^{N+1} |v|^{2N+2} \prod_{j=1}^{N+1} j |1 - j\alpha| \left(1 + \frac{|x|}{\sqrt{2}} - |h| \right)^{(N+1)\alpha-1}}, \quad (2.41)$$

for $\pm t \geq 0$, and

$$|z_{\pm, m+1}(w, x + h, t) - z_{\pm, m}(w, x + h, t)| \leq \frac{2^{4m+\frac{5}{2}} n^{m+\frac{1}{2}} (\beta_2^l)^m \beta_1^l}{\alpha^{m+1} |v|^{2m+1} \left(1 + \frac{|x|}{\sqrt{2}} - |h| \right)^{(m+1)\alpha}} |t|, \quad (2.42)$$

for $t \in \mathbb{R}$ and for $m = 1 \dots N$. We set $z_{\pm, m}(v, \cdot) := z_{\pm, m}(v, 0, \cdot)$ for $m = 1 \dots N + 1$ and $|v| \geq \mu$. For $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $|v_-| \geq \mu$, there exists a unique solution $x(t)$ of equation (1.1) so that

$$x(t) = x_- + z_{-, N+1}(v_-, t) + y_-(t), \quad t \in \mathbb{R} \text{ and } \lim_{t \rightarrow -\infty} (|y_-(t)| + |\dot{y}_-(t)|) = 0. \quad (2.43)$$

In addition when the solution x in (2.43) is a scattering solution then there exists a unique $(v_+, x_+) \in \mathbb{R}^n \times \mathbb{R}^n$, $|v_+| = |v_-|$ so that

$$x(t) = x_+ + z_{+, N+1}(v_+, t) + y_+(t), \quad t \in \mathbb{R} \text{ and } \lim_{t \rightarrow +\infty} (|y_+(t)| + |\dot{y}_+(t)|) = 0. \quad (2.44)$$

In that case we define the scattering data $(a_N(v_-, x_-), b_N(v_-, x_-)) := (v_+, x_+)$, and we consider the inverse scattering problem

$$\text{Given } (a_N, b_N) \text{ and given the long range tail } F^l \text{ of the force } F, \text{ find } F^s. \quad (2.45)$$

The function y_- in (2.43) satisfies the following integral equation $y_- = A_N(y_-)$ where

$$A_N(f)(t) = \int_{-\infty}^t \int_{-\infty}^{\sigma} \left(F(z_{-, N+1}(v_-, \tau) + x_- + f(\tau)) - F^l(z_{-, N}(v_-, \tau)) \right) d\tau d\sigma \quad (2.46)$$

for $t \in \mathbb{R}$ and for $f \in C(\mathbb{R}, \mathbb{R}^n)$, $\sup_{(-\infty, 0]} |f| < \infty$. Then with appropriate changes in the proof of Lemma 2.2 we can study the operator A_N restricted to M_r and we can obtain estimate and contraction estimate similar to (2.10) and (2.11). We also obtain the analog of Lemma 2.3 by appropriate change in its proof, and the decomposition (2.12) remains valid by replacing a , b and $z_-(v_-, \tau) + x_- + y_-(\tau)$ by a_N , b_N and $z_{-, N+1}(v_-, \tau) + x_- + y_-(\tau)$ in (2.12)–(2.18), and by replacing z_+ and $F^l(z_-(v_-, \tau))$ by $z_{+, N+1}$ and $F^l(z_{-, N}(v_-, \tau))$ in (2.12) and (2.16), and by replacing z_+ by $z_{+, N}$ in (2.15), (2.17) and (2.18). An analog of Theorem 2.4 can be proved for the scattering solutions and scattering data (a_N, b_N) . Set $a_{sc, N}(v_-, x_-) := a_N(v_-, x_-) - v_-$ and $b_{sc, N}(v_-, x_-) := b_N(v_-, x_-) - x_-$. Finally the following high energies limits are valid. Let $(\theta, x) \in T\mathbb{S}^{n-1}$, then

$$\lim_{s \rightarrow +\infty} s a_{sc, N}(s\theta, x) = PF^l(\theta, x) + PF^s(\theta, x) \quad (2.47)$$

$$\lim_{s \rightarrow +\infty} s^2 \theta \cdot (b_{sc, N}(s\theta, x) - W_N(s\theta, x)) = -PV^s(\theta, x), \quad (2.48)$$

where

$$W_N(v, x) := \int_{-\infty}^0 \int_{-\infty}^{\sigma} (F^l(z_{-,N}(v, \tau) + x) - F^l(z_{-,N}(v, \tau))) d\tau d\sigma \quad (2.49)$$

$$- \int_0^{+\infty} \int_{\sigma}^{+\infty} (F^l(z_{+,N}(a_N(v, x), \tau) + x) - F^l(z_{+,N}(a_N(v, x), \tau))) d\sigma d\tau,$$

for $(v, x) \in \mathbb{R}^n \times \mathbb{R}^n$, $v \cdot x = 0$ and $|v| > C$ for some constant C . The vector W_N defined by (2.49) is known from the scattering data and from F^l . For the Problem (2.45), from (2.47) (resp. (2.48)) and inversion formulas for the X-ray transform for $n \geq 2$ (see [12, 4, 9, 10]) it follows that F^s can be reconstructed from $a_{sc,N}$ (resp. $b_{sc,N}$).

The limits (2.47) and (2.48) follow from estimates similar to (1.16) and (1.17) that also give the Born approximation of $a_{sc,N}$, $b_{sc,N}$ at fixed energy. However these similar estimates also do not provide the asymptotics of the scattering data $(a_{sc,N}, b_{sc,N})$ when the parameters α , n , v_- and β are fixed and $|x_-| \rightarrow +\infty$. Motivated by this disadvantage, in the next section we modify the definition of the scattering map given in the Introduction so that one can obtain a result on this asymptotic regime.

3 A modified scattering map

3.1 Changing the parametrization of the scattering solutions

We set

$$\mu(\sigma) := \sqrt{\frac{2^5 n \max(\beta_1^l, \beta_2^l)}{\alpha(1 + \frac{\sigma}{\sqrt{2}})^\alpha}}, \text{ for } \sigma \geq 0. \quad (3.1)$$

Under conditions (1.3) and (1.4), the following is valid: for any $(v_-, x_-) \in \mathbb{R}^n \setminus \{0\} \times \mathbb{R}^n$ so that $|v_-| \geq \mu(|x_-|)$ and $v_- \cdot x_- = 0$, then the equation (1.1) has a unique solution $x \in C^2(\mathbb{R}, \mathbb{R}^n)$ such that

$$x(t) = z_-(v_-, x_-, t) + y_-(t), \quad (3.2)$$

where $|\dot{y}_-(t)| + |y_-(t)| \rightarrow 0$, as $t \rightarrow -\infty$, and where $z_-(v_-, x_-, \cdot)$ is defined in Lemma 2.1 (for " $(w, x, v, h) = (v_-, x_-, v_-, 0)$ ").

In addition the function y_- in (3.2) satisfies the integral equation $y_- = \mathcal{A}(y_-)$ where

$$\mathcal{A}(f)(t) = \int_{-\infty}^t \int_{-\infty}^{\sigma} (F(z_-(v_-, x_-, \tau) + f(\tau)) - F^l(z_-(v_-, x_-, \tau))) d\tau d\sigma \quad (3.3)$$

for $t \in \mathbb{R}$ and for $f \in C(\mathbb{R}, \mathbb{R}^n)$, $\sup_{(-\infty, 0]} |f| < \infty$. Under conditions (1.3) and (1.4) we have $\mathcal{A}(f) \in C^2(\mathbb{R}, \mathbb{R}^n)$ for $f \in C(\mathbb{R}, \mathbb{R}^n)$ so that $\sup_{(-\infty, 0]} |f| < \infty$. We study the map \mathcal{A} defined by (3.3) on the metric space M_r defined by (2.9). Set

$$\tilde{k}(v_-, x_-, f) = v_- + \int_{-\infty}^{+\infty} F(z_-(v_-, x_-, \tau) + f(\tau)) d\tau, \quad (3.4)$$

for $f \in M_r$. For the rest of the section we set $\beta_2 = \max(\beta_2^l, \beta_2^s)$.

The following Lemma 3.1 is the analog of Lemma 2.2.

Lemma 3.1. *Let $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $v_- \cdot x_- = 0$, $|v_-| \geq \mu(|x_-|)$, and let $r > 0$, $r < \max(\frac{|v_-|}{2^{\frac{3}{2}}}, 1 + \frac{|x_-|}{\sqrt{2}})$. Then the following estimates are valid:*

$$\begin{aligned} \|\mathcal{A}(f)\| &\leq \tilde{\rho}(n, \alpha, \beta_2, |x_-|, |v_-|, r) \\ &= \frac{2\beta_2 n^{\frac{1}{2}} (n^{\frac{1}{2}} r + 1)}{(\frac{|v_-|}{2^{\frac{3}{2}}} - r)(1 - r + \frac{|x_-|}{\sqrt{2}})^{\alpha}} \left(\frac{1}{(\alpha + 1)(1 - r + \frac{|x_-|}{\sqrt{2}})} + \frac{2}{\alpha(\frac{|v_-|}{2^{\frac{3}{2}}} - r)} \right), \end{aligned} \quad (3.5)$$

$$\|\mathcal{A}(f_1) - \mathcal{A}(f_2)\| \leq \lambda(n, \alpha, \beta_2, \beta_3^s, |x_-|, |v_-|, r) \|f_1 - f_2\|, \quad (3.6)$$

and

$$|\tilde{k}(v_-, x_-, f) - v_-| \leq \frac{2n^{\frac{1}{2}}}{(\frac{|v_-|}{2\sqrt{2}} - r)(1 + \frac{|x_-|}{\sqrt{2}} - r)^{\alpha}} \left(\frac{\beta_1^l}{\alpha} + \frac{\beta_2}{(\alpha + 1)(1 + \frac{|x_-|}{\sqrt{2}} - r)} \right), \quad (3.7)$$

for $(f, f_1, f_2) \in M_r^3$, where λ is defined in (2.11).

Proof of Lemma 3.1 is given in Section 5.

Let $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $v_- \cdot x_- = 0$, and let $r \in (0, \max(\frac{1}{2} + \frac{|x_-|}{2^{\frac{3}{2}}}, 2^{-\frac{3}{2}}|v_-|))$.

Assume that

$$\frac{20n \max(\beta_1^l, \beta_2)}{\alpha(\frac{|v_-|}{2\sqrt{2}} - r)^2 (\frac{1}{2} + \frac{|x_-|}{2^{\frac{3}{2}}} - r)^{\alpha}} \leq 1. \quad (3.8)$$

Then using Lemma 2.1 and (3.7) one can consider the free solution $z_+(\tilde{k}(v_-, x_-, f), x_-, \cdot)$ (for " $(w, v, x, h) = (\tilde{k}(v_-, x_-, f), v_-, x_-, 0)$ ") for $f \in M_r$. In addition, with appropriate changes in the proof of Lemma 2.3 one can prove that when $y_- \in M_r$ is a fixed point of the operator \mathcal{A} then $z_-(v_-, x_-, \cdot) + y_-(\cdot)$ is a scattering solution of (1.1) (in the sense given in Section 2.4), and one can prove that the following decomposition is valid

$$z_-(v_-, x_-, t) + y_-(t) = z_+(\tilde{a}(v_-, x_-), x_- + h, t) + (\mathcal{G}_{v_-, x_-}(h) - h) + H(v_-, x_-, y_-, h)(t), \quad (3.9)$$

for $t \geq 0$, where

$$\tilde{a}(v_-, x_-) = v_- + \int_{-\infty}^{+\infty} F(z_-(v_-, x_-, \tau) + y_-(\tau)) d\tau. \quad (3.10)$$

$$\begin{aligned} \mathcal{G}_{v_-, x_-}(h) &= \tilde{l}(v_-, x_-, y_-) - \int_0^{+\infty} \int_{\sigma}^{+\infty} (F^l(z_-(v_-, x_-, \tau) + y_-(\tau)) \\ &\quad - F^l(z_+(\tilde{a}(v_-, x_-), x_- + h, \tau))) d\tau d\sigma, \end{aligned} \quad (3.11)$$

$$\begin{aligned} \tilde{l}(v_-, x_-, y_-) &= \int_{-\infty}^0 \int_{-\infty}^{\sigma} (F(z_-(v_-, x_-, \tau) + y_-(\tau)) - F^l(z_-(v_-, x_-, \tau))) d\tau d\sigma \\ &\quad - \int_0^{+\infty} \int_{\sigma}^{+\infty} F^s(z_-(v_-, x_-, \tau) + y_-(\tau)) d\tau d\sigma \end{aligned} \quad (3.12)$$

and

$$\begin{aligned} H(v_-, x_-, y_-, h)(t) &= \int_t^{+\infty} \int_{\sigma}^{+\infty} (F(z_-(v_-, x_-, \tau) + y_-(\tau)) \\ &\quad - F^l(z_+(\tilde{a}(v_-, x_-), x_- + h, \tau))) d\tau d\sigma, \end{aligned} \quad (3.13)$$

for $t \geq 0$ and for $|h| \leq \frac{1}{2} + \frac{|x_-|}{2\sqrt{2}}$. We need the following Lemma.

Lemma 3.2. *Let $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $v_- \cdot x_- = 0$ and let $r > 0$, $r < \frac{1}{2} + \frac{|x_-|}{2^{\frac{5}{2}}}$. Under conditions (1.3), (1.4) and (3.8), when $y_- \in M_r$ is a fixed point of the operator \mathcal{A} then the following estimates are valid:*

$$|\mathcal{G}_{v_-, x_-}(h)| \leq \frac{\beta_2(6(nr + \sqrt{n}) + n(1 + \frac{|x_-|}{\sqrt{2}}))}{2\alpha(\alpha + 1)(\frac{|v_-|}{2\sqrt{2}} - r)^2(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}} - r)^{\alpha}} \quad (3.14)$$

$$\leq \frac{1}{4} + \frac{|x_-|}{10\sqrt{2}}, \quad (3.15)$$

$$|\mathcal{G}_{v_-, x_-}(h) - \mathcal{G}_{v_-, x_-}(h')| \leq \frac{16n\beta_2^l|h - h'|}{\alpha(\alpha + 1)|v_-|^2(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}})^{\alpha}} \leq \frac{|h - h'|}{10}, \quad (3.16)$$

for $(h, h') \in \mathbb{R}^n \times \mathbb{R}^n$, $|h'| \leq |h| \leq \frac{1}{4} + \frac{|x_-|}{2^{\frac{5}{2}}}$.

Lemma 3.2 is proved in Section 7.

Under the assumptions of Lemma 3.2 the map \mathcal{G}_{v_-, x_-} is a $\frac{1}{10}$ -contraction map from $\overline{\mathcal{B}}(0, \frac{1}{4} + \frac{|x_-|}{2^{\frac{5}{2}}})$ to $\overline{\mathcal{B}}(0, \frac{1}{4} + \frac{|x_-|}{2^{\frac{5}{2}}})$. We denote by $\tilde{b}_{sc}(v_-, x_-)$ its unique fixed point in $\overline{\mathcal{B}}(0, \frac{1}{4} + \frac{|x_-|}{2^{\frac{5}{2}}})$, and we set $\tilde{b}(v_-, x_-) := x_- + \tilde{b}_{sc}(v_-, x_-)$ and $\tilde{a}_{sc}(v_-, x_-) := \tilde{a}(v_-, x_-) - v_-$. The decomposition (3.9) becomes

$$z_-(v_-, x_-, t) + y_-(t) = z_+(\tilde{a}(v_-, x_-), \tilde{b}(v_-, x_-), t) + y_+(t), \quad (3.17)$$

$$y_+(t) = H(v_-, x_-, y_-, \tilde{b}_{sc}(v_-, x_-))(t), \quad (3.18)$$

for $t \geq 0$. The map $(\tilde{a}_{sc}, \tilde{b}_{sc})$ are our modified scattering data. The inverse scattering problem for equation (1.1) can now be formulated as follows

$$\text{Given } (\tilde{a}_{sc}, \tilde{b}_{sc}) \text{ and } F^l, \text{ find } F^s. \quad (3.19)$$

3.2 Estimates and asymptotics of the modified scattering data

Let $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $v_- \cdot x_- = 0$ and let $r > 0$, $r < \frac{1}{2} + \frac{|x_-|}{2^{\frac{3}{2}}}$ so that condition (3.8) is fulfilled. Then set

$$\tilde{W}(v_-, x_-) := \int_{-\infty}^0 F^l(z_-(v_-, x_-), \tau) d\tau + \int_0^{+\infty} F^l(z_+(\tilde{a}(v_-, x_-), x_-), \tau) d\tau. \quad (3.20)$$

Note that \tilde{W} is known from the modified scattering data and from F^l . We obtain the following analog of Theorem 2.4.

Theorem 3.3. *Under the assumptions of Lemma 3.2 and under conditions (1.3), (1.4) and (3.8), the following estimates are valid:*

$$|\dot{y}_-(t)| \leq \frac{\beta_2(nr + \sqrt{n})}{(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |t|\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+1}}, \quad (3.21)$$

$$|y_-(t)| \leq \frac{\beta_2(nr + \sqrt{n})}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |t|\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha}}, \quad (3.22)$$

for $t \leq 0$; and

$$|\tilde{a}_{sc}(v_-, x_-)| \leq \frac{6\sqrt{n} \max(\beta_1^l, \beta_2)}{\alpha\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^{\alpha}}, \quad (3.23)$$

$$|\tilde{b}_{sc}(v_-, x_-)| \leq \frac{4\beta_2(nr + \sqrt{n})}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}} - r\right)^{\alpha}}, \quad (3.24)$$

$$|y_+(t)| \leq \frac{2\beta_2\sqrt{n}}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2^{\frac{3}{2}}} - r\right)^2\left(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}} - r + t\left(\frac{|v_-|}{2^{\frac{3}{2}}} - r\right)\right)^{\alpha}}, \quad (3.25)$$

for $t \geq 0$. In addition

$$\begin{aligned} & |\tilde{a}_{sc}(v_-, x_-) - \tilde{W}(v_-, x_-) - \int_{-\infty}^{+\infty} F^s(z_-(v_-, x_-), \tau) d\tau| \\ & \leq \frac{4 \max(\beta_2, \beta_3^s)^2 n(nr + \sqrt{n})}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2^{\frac{3}{2}}} - r\right)^2(1 - r + \frac{|x_-|}{\sqrt{2}})^{2\alpha+1}} \left(3 + \frac{2}{\frac{|v_-|}{2^{\frac{3}{2}}} - r}\right)^2, \end{aligned} \quad (3.26)$$

$$|\tilde{b}_{sc}(v_-, x_-) - \tilde{l}(v_-, x_-, 0)| \leq \frac{10n(nr + \sqrt{n}) \max(\beta_2, \beta_3^s)^2 \left(3 + \frac{1}{\frac{|v_-|}{2\sqrt{2}} - r}\right)^2}{\alpha^2(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r\right)^3 \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^{2\alpha}}. \quad (3.27)$$

Theorem 3.3 is proved in Section 7.

For $(\sigma, \tilde{\beta}, r, \tilde{\alpha}) \in (0, +\infty)^3 \times (0, 1]$, $r < \frac{1}{2} + \frac{\sigma}{2^{\frac{3}{2}}}$, and let $\tilde{s}_0 = \tilde{s}_0(\sigma, r, \tilde{\beta}, \tilde{\alpha})$ be defined as the root of the equation

$$1 = \frac{12\tilde{\beta}n}{\tilde{\alpha}r\left(\frac{\tilde{s}_0}{2^{\frac{3}{2}}} - r\right)\left(\frac{1}{2} - r + \frac{\sigma}{2^{\frac{3}{2}}}\right)^{\tilde{\alpha}}} \left(1 + \frac{1}{\frac{\tilde{s}_0}{2^{\frac{3}{2}}} - r}\right)^2, \quad \tilde{s}_0 > 2^{\frac{3}{2}}r. \quad (3.28)$$

Then the high energies asymptotics of the modified scattering data $(\tilde{a}_{sc}, \tilde{b}_{sc})$ are given in the following Theorem 3.4.

Theorem 3.4. *Let $(\theta, x) \in T\mathbb{S}^{n-1}$. Under conditions (1.3) and (1.4) the following limits are valid:*

$$\lim_{s \rightarrow +\infty} s(\tilde{a}_{sc}(s\theta, x) - \tilde{W}(s\theta, x)) = PF^s(\theta, x), \quad (3.29)$$

$$\lim_{s \rightarrow +\infty} s^2\theta \cdot \tilde{b}_{sc}(s\theta, x) = -PV^s(\theta, x), \quad (3.30)$$

In addition,

$$\left| \tilde{a}_{sc}(s\theta, x) - \tilde{W}(s\theta, x) - \int_{-\infty}^{+\infty} F^s(\tau s\theta + x) d\tau \right| \leq \frac{12n^2\beta^2 \left(3 + \frac{2}{2^{-\frac{3}{2}}s - r}\right)^2}{\alpha(\alpha + 1)(2^{-\frac{3}{2}}s - r)^2 \left(1 - r + \frac{|x|}{\sqrt{2}}\right)^{2\alpha+1}}, \quad (3.31)$$

$$\begin{aligned} & \left| \tilde{b}_{sc}(s\theta, x) - \int_{-\infty}^0 \int_{-\infty}^{\sigma} F^s(\tau s\theta + x) d\tau d\sigma + \int_0^{+\infty} \int_{\sigma}^{+\infty} F^s(\tau s\theta + x) d\tau d\sigma \right| \\ & \leq \frac{24n^2\beta^2 \left(3 + \frac{1}{2^{-\frac{3}{2}}s - r}\right)^2}{\alpha^2(\alpha + 1)(2^{-\frac{3}{2}}s - r)^3 \left(1 - r + \frac{|x|}{\sqrt{2}}\right)^{2\alpha}}, \end{aligned} \quad (3.32)$$

for $r \in (0, \frac{1}{2} + \frac{|x|}{2^{\frac{3}{2}}})$ and for $s > \tilde{s}_0(|x|, r, \beta, \alpha)$, where $\beta = \max(\beta_1^l, \beta_2, \beta_3^s)$.

Formulas (3.29) and (3.30) prove that F^s can be reconstructed from the high energies asymptotics of the modified scattering data. For Problem (3.19) this implies that F^s can be reconstructed from \tilde{a}_{sc} and \tilde{b}_{sc} .

Estimates (3.31) and (3.32) also provide the Born approximation of the modified scattering data at fixed energy, i.e. the leading term of the asymptotics of \tilde{a}_{sc} , \tilde{b}_{sc} , when the parameters α , n , s , θ and x are fixed and β

decreases to 0. We obtain that F^s can be reconstructed from the Born approximation of the modified scattering data at fixed energy.

Estimates (3.31) and (3.32) also provide the first leading term in the asymptotics of the modified scattering data when the parameters α , n , s , θ and β are fixed and $|x|$ increases to $+\infty$.

Proof of Theorem 3.4. Let $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $v_- \cdot x_- = 0$ and $|v_-| \geq \mu(|x_-|)$. Similarly to estimate (2.29) we have

$$|\eta(x_- + \tau v_-) + (1 - \eta)z_-(v_-, x_-, \tau)| \geq \frac{|x_-|}{\sqrt{2}} + |\tau| \frac{|v_-|}{2^{\frac{3}{2}}}, \quad (3.33)$$

for $\eta \in (0, 1)$ and $\tau \in \mathbb{R}$. Using (3.33), (1.4) and (2.4), we obtain

$$\begin{aligned} & \left| \int_{-\infty}^{+\infty} (F^s(z_-(v_-, x_-, \tau)) - F^s(\tau v_- + x_-)) d\tau \right| \\ & \leq \frac{2^{\frac{5}{2}} n^{\frac{3}{2}} \beta_1^l \beta_3^s}{\alpha |v_-| (1 + \frac{|x_-|}{\sqrt{2}})^\alpha} \int_{-\infty}^{+\infty} \frac{|\tau|}{(1 + \frac{|x_-|}{\sqrt{2}} + |\tau| \frac{|v_-|}{2\sqrt{2}})^{\alpha+3}} d\tau \\ & \leq \frac{4n^{\frac{3}{2}} \beta_1^l \beta_3^s}{\alpha(\alpha+1) (\frac{|v_-|}{2\sqrt{2}})^3 (1 + \frac{|x_-|}{\sqrt{2}})^{2\alpha+1}}, \end{aligned} \quad (3.34)$$

and

$$\begin{aligned} & \int_{-\infty}^0 \int_{-\infty}^\sigma |F^s(z_-(v_-, x_-, \tau)) - F^s(\tau v_- + x_-)| d\tau d\sigma \\ & \leq \frac{2^{\frac{5}{2}} n^{\frac{3}{2}} \beta_1^l \beta_3^s}{\alpha |v_-| (1 + \frac{|x_-|}{\sqrt{2}})^\alpha} \int_{-\infty}^0 \int_{-\infty}^\sigma \frac{|\tau|}{(1 + \frac{|x_-|}{\sqrt{2}} + |\tau| \frac{|v_-|}{2\sqrt{2}})^{\alpha+3}} d\tau d\sigma \end{aligned} \quad (3.35)$$

The same estimate (3.35) holds for $\int_0^{+\infty} \int_\sigma^{+\infty} |F^s(z_-(v_-, x_-, \tau)) - F^s(\tau v_- + x_-)| d\tau d\sigma$. Therefore from the definition (3.12) it follows that

$$\begin{aligned} & |\tilde{l}(v_-, x_-, 0) - \int_{-\infty}^0 \int_{-\infty}^\sigma F^s(\tau v_- + x_-) d\tau + \int_0^{+\infty} \int_\sigma^{+\infty} F^s(\tau v_- + x_-) d\tau d\sigma| d\sigma \\ & \leq \frac{4n^{\frac{3}{2}} \beta_1^l \beta_3^s}{\alpha^2(\alpha+1) (\frac{|v_-|}{2\sqrt{2}})^4 (1 + \frac{|x_-|}{\sqrt{2}})^{2\alpha}}. \end{aligned} \quad (3.36)$$

Let $r > 0$, $r < \max(\frac{|v_-|}{2^{\frac{3}{2}}}, \frac{1}{2} + \frac{|x_-|}{\sqrt{2}})$. Note that

$$\max\left(\frac{\tilde{\rho}}{r}, \lambda, \frac{20n \max(\beta_1^l, \beta_2)}{\alpha (\frac{|v_-|}{2\sqrt{2}} - r)^2 (\frac{1}{2} + \frac{|x_-|}{2^{\frac{3}{2}}} - r)^\alpha}\right) \leq \frac{12\beta n (1 + \frac{1}{\frac{|v_-|}{2^{\frac{3}{2}}} - r})^2}{\alpha r (\frac{|v_-|}{2^{\frac{3}{2}}} - r) (\frac{1}{2} - r + \frac{|x_-|}{2^{\frac{3}{2}}})^\alpha}, \quad (3.37)$$

where $\tilde{\rho}$ and λ are defined in (3.5) and (2.11) respectively. Assume that $|v_-| > \tilde{s}_0(|x_-|, r, \beta, \alpha)$ where \tilde{s}_0 is the root of the equation (3.28). Then from (3.28) and Lemma 3.1 it follows that \mathcal{A} has a unique fixed point denoted by y_- in M_r . Then adding (3.26) and (3.34) we obtain (3.31). And adding (3.36) and (3.27) we obtain (3.32). Theorem 3.4 is proved. \square

3.3 Approximating the “free” solutions $z_{\pm}(v, x, \cdot)$

One may approximate the solutions $z_{\pm}(v, x, \cdot)$ by the functions $z_{\pm, N+1}(v, x, \cdot)$ defined in Section 2.4 that are easier to compute in general. Then we can repeat the study of the previous Subsections. For $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $v_- \cdot x_- = 0$, $|v_-| \geq \mu(|x_-|)$, there exists a unique solution $x(t)$ of equation (1.1) so that

$$x(t) = z_{-, N+1}(v_-, x_-, t) + y_-(t), \quad t \in \mathbb{R} \text{ and } \lim_{t \rightarrow -\infty} (|y_-(t)| + |\dot{y}_-(t)|) = 0. \quad (3.38)$$

The function y_- in (3.38) satisfies the following integral equation $y_- = \mathcal{A}_N(y_-)$ where

$$\mathcal{A}_N(f)(t) = \int_{-\infty}^t \int_{-\infty}^{\sigma} \left(F(z_{-, N+1}(v_-, x_-, \tau) + f(\tau)) - F^l(z_{-, N}(v_-, x_-, \tau)) \right) d\tau d\sigma \quad (3.39)$$

for $t \in \mathbb{R}$ and for $f \in C(\mathbb{R}, \mathbb{R}^n)$, $\sup_{(-\infty, 0]} |f| < \infty$. Then with appropriate changes in the proof of Lemma 3.1 we can study the operator \mathcal{A}_N restricted to M_r and we can obtain estimate and contraction estimate similar to (3.5) and (3.6). The decomposition (3.9) remains valid with the following changes: first we define \tilde{a}_N by the formula (3.10) where we replaced $z_-(v_-, x_-, \tau)$ by $z_{-, N+1}(v_-, x_-, \tau)$; then we replace \tilde{a} and $z_-(v_-, x_-, \tau) + y_-(\tau)$ by \tilde{a}_N and $z_{-, N+1}(v_-, x_-, \tau) + y_-(\tau)$ in (3.9), (3.11)–(3.13), and we replace z_+ and $F^l(z_-(v_-, x_-, \tau))$ by $z_{+, N+1}$ and $F^l(z_{-, N}(v_-, x_-, \tau))$ in (3.9) and (3.12), and we replace z_+ by $z_{+, N}$ in (3.11), (3.13). This defines a new map \mathcal{G}_{v_-, x_-} and the analog of Lemma 3.2 can be proved. Such a result then allows to define the scattering data \tilde{b}_N . Then an analog of the Theorem 3.3 can be proved for the scattering solutions and scattering data $(\tilde{a}_N, \tilde{b}_N)$. Set $\tilde{a}_{sc, N}(v_-, x_-) := \tilde{a}_N(v_-, x_-) - v_-$ and $\tilde{b}_{sc, N}(v_-, x_-) := \tilde{b}_N(v_-, x_-) - x_-$. Finally the following high energies limits are valid. Let $(\theta, x) \in T\mathbb{S}^{n-1}$, then

$$\lim_{s \rightarrow +\infty} s(\tilde{a}_{sc, N}(s\theta, x) - \tilde{W}_N(s\theta, x)) = PF^s(\theta, x), \quad (3.40)$$

$$\lim_{s \rightarrow +\infty} s^2 \theta \cdot \tilde{b}_{sc, N}(s\theta, x) = -PV^s(\theta, x), \quad (3.41)$$

where

$$\tilde{W}_N(v, x) := \int_{-\infty}^0 F^l(z_{-,N}(v, x, \tau)) d\tau + \int_0^{+\infty} F^l(z_{+,N}(\tilde{a}_N(v, x), x, \tau)) d\tau. \quad (3.42)$$

for $(v, x) \in \mathbb{R}^n \times \mathbb{R}^n$, $v \cdot x = 0$, $|v| > C$ for some constant C . The vector \tilde{W}_N defined by (3.42) is known from the scattering data and from F^l . From (3.40) (resp. (3.41)) it follows that F^s can be reconstructed from $\tilde{a}_{sc,N}$ (resp. $\tilde{b}_{sc,N}$).

The limits (3.40) and (3.41) follow from estimates similar to (3.31) and (3.32) that also give the Born approximation of $\tilde{a}_{sc,N}$, $\tilde{b}_{sc,N}$ at fixed energy, and the first leading term of the asymptotics of the scattering data $(\tilde{a}_{sc,N}, \tilde{b}_{sc,N})$ when the parameters α , n , v_- and β are fixed and $|x_-| \rightarrow +\infty$.

4 Proof of Lemmas 2.1 and 2.3

Proof of Lemma 2.1. We prove the existence and uniqueness of the solution z_+ (similarly one can prove the existence and uniqueness of z_- or just use the relation “ $z_-(w, x + h, t) = z_+(-w, x + h, -t)$ ”). Set $C' = \frac{2^{\frac{5}{2}} n^{\frac{1}{2}} \beta_1^l}{\alpha |v| (1 + \frac{|x|}{\sqrt{2}} - |h|)^\alpha}$. Let \mathcal{V} be the complete metric space defined by

$$\mathcal{V} := \{g \in C(\mathbb{R}, \mathbb{R}^n) \mid |g(t)| \leq C'|t| \text{ for } t \in \mathbb{R}\},$$

endowed with the following norm $\|g\|_{\mathcal{V}} := \sup_{t \in \mathbb{R} \setminus \{0\}} \left| \frac{g(t)}{t} \right|$. We consider the integral equation

$$G_+ f(t) = - \int_0^t \int_{\sigma}^{+\infty} F^l(x + h + \tau w + f(\tau)) d\tau d\sigma, \quad t \in \mathbb{R}, \quad (4.1)$$

for $f \in \mathcal{V}$. First note that

$$\begin{aligned} |x + h + \tau w + f(\tau)| &\geq |x + \tau v| - |h| - |\tau| |v - w| - |f(\tau)| \\ &\geq \frac{|x|}{\sqrt{2}} - |h| + \left(\frac{3|v|}{4\sqrt{2}} - C' \right) |\tau| \geq \frac{|x|}{\sqrt{2}} - |h| + \frac{|v|}{2\sqrt{2}} |\tau|, \end{aligned} \quad (4.2)$$

for $\tau \in \mathbb{R}$ and $f \in \mathcal{V}$ (we used that $x \cdot v = 0$ and that $C' \leq \frac{|v|}{4\sqrt{2}}$). Using (1.3) we obtain that

$$|G_+ f(t)| \leq \sqrt{n} \beta_1^l \int_{-|t|}^0 \int_{\sigma}^{+\infty} \left(1 + \frac{|x|}{\sqrt{2}} - |h| + \frac{|v|}{2\sqrt{2}} |\tau| \right)^{-\alpha-1} d\tau d\sigma \leq C'|t|, \quad (4.3)$$

for $t \in \mathbb{R}$ and $f \in \mathcal{V}$. Now let $(f_1, f_2) \in \mathcal{V}^2$. Then using (1.3) we have

$$\begin{aligned} |F^l(x + h + \tau w + f_1(\tau)) - F^l(x + h + \tau w + f_2(\tau))| &\leq n\beta_2^l |f_1 - f_2|(\tau) \\ &\times \sup_{\varepsilon \in (0,1)} (1 + |x + h + \tau w + (\varepsilon f_1 + (1 - \varepsilon)f_2)(\tau)|)^{-\alpha-2}, \end{aligned}$$

for $\tau \in \mathbb{R}$. Hence using also (4.2) we have

$$\begin{aligned} |G_+ f_1(t) - G_+ f_2(t)| &\leq n\beta_2^l \|f_1 - f_2\|_{\mathcal{V}} \int_{-|t|}^0 \int_{\sigma}^{+\infty} \frac{|\tau| d\tau d\sigma}{\left(1 + \frac{|x|}{\sqrt{2}} - |h| + \frac{|v|}{2\sqrt{2}} |\tau|\right)^{\alpha+2}} \\ &\leq \frac{2\sqrt{2}n\beta_2^l \|f_1 - f_2\|_{\mathcal{V}}}{|v|} \int_{-|t|}^0 \int_{\sigma}^{+\infty} \frac{d\tau d\sigma}{\left(1 + \frac{|x|}{\sqrt{2}} - |h| + \frac{|v|}{2\sqrt{2}} |\tau|\right)^{\alpha+1}} \\ &\leq \frac{16n\beta_2^l |t| \|f_1 - f_2\|_{\mathcal{V}}}{\alpha |v|^2 \left(1 + \frac{|x|}{\sqrt{2}} - |h|\right)^{\alpha}} \leq 2^{-1} \|f_1 - f_2\|_{\mathcal{V}}, \end{aligned} \quad (4.4)$$

for $t \in \mathbb{R}$ (we used (2.2)). From (4.3) and (4.4) it follows that the operator G_+ is a contraction map from \mathcal{V} to \mathcal{V} . Set $z_+(w, x + h, t) = x + h + tw + f_{w, x+h}(t)$ for $t \in \mathbb{R}$, where $f_{w, x+h}$ denotes the unique fixed point of G_+ in \mathcal{V} . Then $z_+(w, x + h, \cdot)$ satisfies (1.6), (2.3) and (2.4). \square

Before proving Lemma 2.3 we recall the following standard result (see also [6, Lemma II.2]). For sake of consistency we provide a proof of Lemma 4.1 at the end of this Section.

Lemma 4.1. *Let $x(t)$ be a solution of equation (1.1) and let $z(t)$ be a solution of equation (1.6). Assume that there exists a vector $v \in \mathbb{R}^n$, $v \neq 0$, so that*

$$\lim_{t \rightarrow +\infty} \dot{z}(t) = \lim_{t \rightarrow +\infty} \dot{x}(t) = v. \quad (4.5)$$

Then

$$\sup_{(0, +\infty)} |x - z| < \infty. \quad (4.6)$$

Proof of Lemma 2.3. We need the following preliminary estimate (4.8). Using (2.9) we have for $\tau \in \mathbb{R}$ and $f \in M_r$,

$$|f(\tau)| \leq r|\tau| + r. \quad (4.7)$$

Hence

$$\begin{aligned} |z_-(v_-, \tau) + x_- + f(\tau)| &\geq |x_- + \tau v_-| - |z_-(v_-, \tau) - \tau v_-| - |f(\tau)| \\ &\geq \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{\sqrt{2}} - \frac{2^{\frac{5}{2}} n^{\frac{1}{2}} \beta_1^l}{|v_-|^{\alpha}} - r \right) \\ &\geq \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r \right). \end{aligned} \quad (4.8)$$

We used (2.4) (with " $(x, w, v, h) = (0, v_-, v_-, 0)$ "), the inequality $|x_- + \tau v_-| \geq \frac{|x_-|}{\sqrt{2}} + |\tau| \frac{|v_-|}{\sqrt{2}}$ ($x_- \cdot v_- = 0$) and (4.7) and we used the condition $|v_-| \geq \mu$. Hence the integral $\int_{-\infty}^{+\infty} F(z_-(v_-, \tau) + x_- + f(\tau)) d\tau$ is absolutely convergent for any $f \in M_r$. And when $y_- \in M_r$ is a fixed point for A then $z_-(v_-, \cdot) + x_- + y_-$ satisfies equation (1.1) (see (2.8)) and $\dot{z}_-(v_-, t) + x_- + \dot{y}_-(t) = v_- + \int_{-\infty}^t F(z_-(v_-, \tau) + x_- + \dot{y}_-(\tau)) d\tau \rightarrow a(v_-, x_-)$ as $t \rightarrow +\infty$, where $a(v_-, x_-)$ is defined in (2.13). Then from Lemma 4.1 it follows that $\sup_{t \in (0, +\infty)} |z_-(v_-, t) + x_- + y_-(t) - z_+(a(v_-, x_-), t)| < +\infty$. Using this latter estimate and (1.3) and (1.4), and using (4.8) we obtain that the integrals on the right hand sides of (2.15) and (2.18) are absolutely convergent. Then the decomposition (2.12) follows from the equality $A(y_-) = y_-$ and (2.8) and straightforward computations. \square

Proof of Lemma 4.1. We set $\delta(t) = x(t) - z(t)$ for $t \geq 0$. Property (4.5) shows that there exists $\varepsilon > 0$ so that

$$1 + |\eta x(t) + (1 - \eta)z(t)| \geq \varepsilon(1 + t), \text{ for } t \geq 0 \text{ and } \eta \in [0, 1]. \quad (4.9)$$

Then from equations (1.1) and (1.6) it follows that

$$\delta(t) = \delta(0) - \int_0^t \int_{\sigma}^{+\infty} F^s(x(\tau)) d\tau d\sigma - \int_0^t \int_{\sigma}^{+\infty} (F^l(x(\tau)) - F^l(z(\tau))) d\tau d\sigma, \quad (4.10)$$

for $t \geq 0$, where the integrals on the right hand side of (4.10) are absolutely convergent (see (4.9) and (1.3) and (1.4)). Note that

$$\int_0^t \int_{\sigma}^{+\infty} |F^s(x(\tau))| d\tau d\sigma \leq \sqrt{n} \beta_2 \varepsilon^{-\alpha-2} \int_0^t \int_{\sigma}^{+\infty} \frac{d\tau d\sigma}{(1 + \tau)^{\alpha+2}} \leq \frac{\sqrt{n} \beta_2}{\alpha(\alpha + 1) \varepsilon^{\alpha+2}}, \quad (4.11)$$

for $t \geq 0$. Hence using (4.10) we obtain

$$|\delta(t)| \leq C_0 + \int_0^t \int_{\sigma}^{+\infty} |F^l(x(\tau)) - F^l(z(\tau))| d\tau d\sigma, \text{ for } t \geq 0, \quad (4.12)$$

where $C_0 = |\delta(0)| + \frac{\sqrt{n} \beta_2}{\alpha(\alpha+1) \varepsilon^{\alpha+2}}$.

One may assume without loss of generality that $\alpha \neq \frac{1}{m}$ for any $m \in \mathbb{N}$. Otherwise replace α by some $\alpha' \in (0, \alpha)$ so that $\alpha' \neq \frac{1}{m'}$ for any $m' \in \mathbb{N}$. Then

$$\begin{aligned} \int_0^t \int_{\sigma}^{+\infty} (|F^l(x(\tau))| + |F^l(z(\tau))|) d\tau d\sigma &\leq 2\sqrt{n} \beta_1^l \varepsilon^{-\alpha-1} \int_0^t \int_{\sigma}^{+\infty} \frac{d\tau d\sigma}{(1 + \tau)^{\alpha+1}} \\ &\leq \frac{2\sqrt{n} \beta_1^l}{\alpha(1 - \alpha) \varepsilon^{\alpha+1}} (1 + t)^{1-\alpha}, \end{aligned} \quad (4.13)$$

for $t \geq 0$. Using also (4.12) we obtain that there exist positive constants C_1 and C'_1 so that $|\delta(t)| \leq C_1 + C'_1 t^{1-\alpha}$ for $t \geq 0$. Now using (4.12), the growth properties of F^l (1.3) and (4.9) we obtain

$$\begin{aligned} |\delta(t)| &\leq C_0 + \int_0^t \int_\sigma^{+\infty} \sup_{\eta \in (0,1)} \frac{n\beta_2 |x(\tau) - z(\tau)|}{(1 + |\eta x(\tau) + (1-\eta)z(\tau)|)^{\alpha+2}} d\tau d\sigma \\ &\leq C_0 + n\beta_2 \varepsilon^{-\alpha-2} \int_0^t \int_\sigma^{+\infty} \frac{|\delta(\tau)|}{(1+\tau)^{\alpha+2}} d\tau d\sigma, \end{aligned} \quad (4.14)$$

for $t \geq 0$. Then using (4.14) we prove by induction the following: For any $m = 1 \dots \lfloor \alpha^{-1} \rfloor$ there exist positive constants C_m and C'_m so that $|\delta(t)| \leq C_m + C'_m t^{1-m\alpha}$ for $t \geq 0$. Combining again this latter estimate for $m = \lfloor \alpha^{-1} \rfloor$ and the estimate (4.14) we obtain

$$\begin{aligned} |\delta(t)| &\leq C_0 + n\beta_2 \varepsilon^{-\alpha-2} \int_0^t \int_\sigma^{+\infty} \frac{C_m + C'_m \tau^{1-m\alpha}}{(1+\tau)^{\alpha+2}} d\tau d\sigma \\ &\leq C_0 + \frac{n\beta_2 C_m}{\alpha(\alpha+1)\varepsilon^{\alpha+2}} + \frac{C'_m n\beta_2}{(m+1)\alpha((m+1)\alpha-1)\varepsilon^{\alpha+2}}, \end{aligned}$$

for $t \geq 0$ ($(m+1)\alpha - 1 > 0$), which proves the lemma. \square

5 Proof of Lemmas 2.2 and 3.1

Proof of Lemma 2.2. We first prove (2.10). We need the following estimates for $A(f)(t)$ (5.5) and (5.11). Using (1.3), (1.4) and (4.8) we obtain

$$|F^s(z_-(v_-, \tau) + x_- + f(\tau))| \leq \frac{\beta_2 \sqrt{n}}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+2}}, \quad (5.1)$$

and

$$|F^l(z_-(v_-, \tau) + x_- + f(\tau)) - F^l(z_-(v_-, \tau))| \leq \frac{\beta_2 n(|x_-| + |f(\tau)|)}{\left(1 - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+2}}, \quad (5.2)$$

for $\tau \in \mathbb{R}$. In addition from (2.8) it follows that

$$\dot{A}(f)(t) = \int_{-\infty}^t \left(F(z_-(v_-, \tau) + x_- + f(\tau)) - F^l(z_-(v_-, \tau)) \right) d\tau, \quad \tau \in \mathbb{R}. \quad (5.3)$$

Therefore combining (5.1), (5.2) (with $|f(\tau)| \leq \sup_{(-\infty, 0)} |f|$ for $\tau \leq 0$) and (5.3) we obtain

$$\begin{aligned} |\dot{A}(f)(t)| &\leq \beta_2 \int_{-\infty}^t \frac{\sqrt{n} + n(|x_-| + |f(\tau)|)}{(1 - r + |\tau|(\frac{|v_-|}{2\sqrt{2}} - r))^{\alpha+2}} d\tau \\ &\leq \frac{\beta_2(n(|x_-| + \sup_{(-\infty, 0)} |f|) + \sqrt{n})}{(\alpha + 1)(\frac{|v_-|}{2\sqrt{2}} - r) \left(1 - r + |t|(\frac{|v_-|}{2\sqrt{2}} - r)\right)^{\alpha+1}}, \end{aligned} \quad (5.4)$$

for $t \leq 0$. Then integrating over $(-\infty, t)$ we obtain

$$|A(f)(t)| \leq \frac{\beta_2(n(|x_-| + \sup_{(-\infty, 0)} |f|) + \sqrt{n})}{\alpha(\alpha + 1)(\frac{|v_-|}{2\sqrt{2}} - r)^2 \left(1 - r + |t|(\frac{|v_-|}{2\sqrt{2}} - r)\right)^{\alpha}}, \quad (5.5)$$

for $t \leq 0$.

Now let $t \geq 0$. Using (5.3) we have

$$\begin{aligned} \dot{A}(f)(t) &= \dot{A}(f)(0) + \int_0^t (F^l(z_-(v_-), \tau) + x_- + f(\tau)) - F^l(z_-(v_-), \tau)) d\tau \\ &\quad + \int_0^t F^s(z_-(v_-), \tau) + x_- + f(\tau) d\tau. \end{aligned} \quad (5.6)$$

Hence from (5.4) (with " $t = 0$ "), (5.1) and (5.2) (with $|f(\tau)| \leq (1 + \tau) \sup_{s \in (0, +\infty)} \frac{|f(s)|}{1+s}$ for $\tau \geq 0$) it follows that

$$\begin{aligned} |\dot{A}(f)(t)| &\leq \frac{\beta_2(n(|x_-| + \sup_{(-\infty, 0)} |f|) + \sqrt{n})}{(\alpha + 1)(\frac{|v_-|}{2\sqrt{2}} - r)(1 - r)^{\alpha+1}} \\ &\quad + \beta_2 \int_0^t \frac{n^{\frac{1}{2}} + n|x_-| + n(1 + \tau) \sup_{s \in (0, +\infty)} \frac{|f(s)|}{1+s}}{\left(1 - r + |\tau|(\frac{|v_-|}{2\sqrt{2}} - r)\right)^{\alpha+2}} d\tau \\ &\leq \frac{\beta_2(2n|x_-| + n\|f\| + 2\sqrt{n})}{(\alpha + 1)(\frac{|v_-|}{2\sqrt{2}} - r)(1 - r)^{\alpha+1}} + \frac{n\beta_2 \sup_{s \in (0, +\infty)} \frac{|f(s)|}{1+s}}{\alpha(\frac{|v_-|}{2\sqrt{2}} - r)^2 (1 - r)^{\alpha}}. \end{aligned} \quad (5.7)$$

We also have

$$\begin{aligned} A(f)(t) &= A(f)(0) + t\dot{A}(f)(t) \\ &\quad - \int_0^t \int_{\sigma}^t \left(F^l(z_-(v_-), \tau) + x_- + f(\tau) - F^l(z_-(v_-), \tau)\right) d\tau d\sigma \\ &\quad - \int_0^t \int_{\sigma}^t F^s(z_-(v_-), \tau) + x_- + f(\tau) d\tau d\sigma. \end{aligned} \quad (5.8)$$

Using (5.1) and (5.2) we obtain

$$\left| \int_0^t \int_\sigma^t F^s(z_-(v_-, \tau) + x_- + f(\tau)) d\tau d\sigma \right| \leq \frac{\beta_2 \sqrt{n}}{\alpha(\alpha+1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha}, \quad (5.9)$$

$$\begin{aligned} & \left| \int_0^t \int_\sigma^t (F^l(z_-(v_-, \tau) + x_- + f(\tau)) - F^l(z_-(v_-, \tau))) d\tau d\sigma \right| \\ & \leq t \int_0^{+\infty} |F^l(z_-(v_-, \tau) + x_- + f(\tau)) - F^l(z_-(v_-, \tau))| d\tau \\ & \leq t \left(\frac{n\beta_2 \sup_{s \in (0, +\infty)} \frac{|f(s)|}{1+s}}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 (1-r)^\alpha} + \frac{n\beta_2 (|x_-| + \sup_{s \in (0, +\infty)} \frac{|f(s)|}{1+s})}{(\alpha+1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right) (1-r)^{\alpha+1}} \right). \quad (5.10) \end{aligned}$$

Combining (5.8), (5.5) (with "t = 0"), (5.7), (5.9) and (5.10) we obtain

$$\begin{aligned} |A(f)(t)| & \leq \frac{\beta_2(n(|x_-| + \sup_{(-\infty, 0)} |f|) + 2\sqrt{n})}{\alpha(\alpha+1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 (1-r)^\alpha} \\ & + t\beta_2 \left(\frac{(3n|x_-| + 2n\|f\| + 2\sqrt{n})}{(\alpha+1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right) (1-r)^{\alpha+1}} + \frac{2n \sup_{s \in (0, +\infty)} \frac{|f(s)|}{1+s}}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 (1-r)^\alpha} \right). \quad (5.11) \end{aligned}$$

Then (2.10) follows from (5.5) and (5.11) and the estimate $\|f\| \leq r$.

It remains to prove (2.11). Estimate (2.11) will follow from (5.16) and (5.24) given below. Let $(f_1, f_2) \in M_r^2$. Using (1.3), (1.4) and (4.8) we obtain

$$\begin{aligned} |F^l(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^l(z_-(v_-, \tau) + x_- + f_2(\tau))| & \leq \frac{\beta_2 n |f_1 - f_2|(\tau)}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \right)^{\alpha+2}}, \\ & \quad (5.12) \\ |F^s(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^s(z_-(v_-, \tau) + x_- + f_2(\tau))| & \leq \frac{\beta_3^s n |f_1 - f_2|(\tau)}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \right)^{\alpha+3}}, \\ & \quad (5.13) \end{aligned}$$

for $\tau \in \mathbb{R}$. In addition from (5.3) it follows that

$$\begin{aligned} & \dot{A}(f_1)(t) - \dot{A}(f_2)(t) \\ & = \int_{-\infty}^t (F^l(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^l(z_-(v_-, \tau) + x_- + f_2(\tau))) d\tau \\ & + \int_{-\infty}^t (F^s(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^s(z_-(v_-, \tau) + x_- + f_2(\tau))) d\tau. \quad (5.14) \end{aligned}$$

Hence we integrate in the τ variable over the interval $(-\infty, t)$, $t \leq 0$, both sides of (5.12) and (5.13) where we use the inequality $|f_1 - f_2|(\tau) \leq \sup_{(-\infty, 0)} |f_1 - f_2|$, and we obtain

$$\begin{aligned} |\dot{A}(f_1)(t) - \dot{A}(f_2)(t)| &\leq \frac{n \sup_{(-\infty, 0]} |f_1 - f_2|}{\left(\frac{|v_-|}{2\sqrt{2}} - r\right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r + |t| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+1}} \left[\frac{\beta_2}{(\alpha+1)} \right. \\ &\quad \left. + \frac{\beta_3^s}{(\alpha+2) \left(1 + \frac{|x_-|}{\sqrt{2}} - r + |t| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)} \right], \end{aligned} \quad (5.15)$$

for $t \leq 0$. Then we integrate in the t variable both sides of (5.15) and we obtain

$$\begin{aligned} |A(f_1)(t) - A(f_2)(t)| &\leq \frac{n \sup_{(-\infty, 0]} |f_1 - f_2|}{(\alpha+1) \left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r + |t| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^\alpha} \left[\frac{\beta_2}{\alpha} \right. \\ &\quad \left. + \frac{\beta_3^s}{(\alpha+2) \left(1 + \frac{|x_-|}{\sqrt{2}} - r + |t| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)} \right], \end{aligned} \quad (5.16)$$

for $t \leq 0$.

From (5.6) it follows that

$$\begin{aligned} \dot{A}(f_1)(t) - \dot{A}(f_2)(t) &= \dot{A}(f_1)(0) - \dot{A}(f_2)(0) \\ &+ \int_0^t (F^l(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^l(z_-(v_-, \tau) + x_- + f_2(\tau))) d\tau \\ &+ \int_0^t (F^s(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^s(z_-(v_-, \tau) + x_- + f_2(\tau))) d\tau. \end{aligned} \quad (5.17)$$

We have

$$\begin{aligned} &\left| \int_0^t (F^s(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^s(z_-(v_-, \tau) + x_- + f_2(\tau))) d\tau \right| \\ &\leq n\beta_3^s \sup_{s \in [0, +\infty)} \frac{|(f_1 - f_2)(s)|}{1+s} \int_0^t \frac{1+\tau}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+3}} d\tau \\ &\leq \frac{n\beta_3^s \sup_{s \in [0, +\infty)} \frac{|(f_1 - f_2)(s)|}{1+s}}{\left(\frac{|v_-|}{2\sqrt{2}} - r\right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^{\alpha+1}} \left[\frac{1}{(\alpha+2) \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)} + \frac{1}{(\alpha+1) \left(\frac{|v_-|}{2\sqrt{2}} - r\right)} \right], \end{aligned} \quad (5.18)$$

and

$$\begin{aligned}
& \left| \int_0^t \left(F^l(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^l(z_-(v_-, \tau) + x_- + f_2(\tau)) \right) d\tau \right| \\
& \leq \frac{n\beta_2 \sup_{s \in [0, +\infty)} \frac{|(f_1 - f_2)(s)|}{1+s}}{\left(\frac{|v_-|}{2\sqrt{2}} - r \right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} \left[\frac{1}{(\alpha + 1) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)} + \frac{1}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r \right)} \right]. \tag{5.19}
\end{aligned}$$

Hence using also (5.15) (for “ $t = 0$ ”) we obtain

$$\begin{aligned}
|\dot{A}(f_1)(t) - \dot{A}(f_2)(t)| & \leq \frac{n\|f_1 - f_2\|}{\left(\frac{|v_-|}{2\sqrt{2}} - r \right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{\alpha+1}} \left[\frac{\beta_2}{(\alpha + 1)} \right. \\
& \quad \left. + \frac{\beta_3^s}{(\alpha + 2) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)} \right] \\
& \quad + \frac{n \sup_{s \in [0, +\infty)} \frac{|(f_1 - f_2)(s)|}{1+s}}{\left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} \left[\frac{\beta_3^s}{(\alpha + 1) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)} + \frac{\beta_2}{\alpha} \right]. \tag{5.20}
\end{aligned}$$

Next we consider

$$\begin{aligned}
A(f_1)(t) - A(f_2)(t) & = A(f_1)(0) - A(f_2)(0) + t(\dot{A}(f_1)(t) - \dot{A}(f_2)(t)) \tag{5.21} \\
& + \int_0^t \int_\sigma^t \left(F^l(z_-(v_-, \tau) + x_- + f_2(\tau)) - F^l(z_-(v_-, \tau) + x_- + f_1(\tau)) \right) d\tau d\sigma \\
& - \int_0^t \int_\sigma^t \left(F^s(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^s(z_-(v_-, \tau) + x_- + f_2(\tau)) \right) d\tau d\sigma
\end{aligned}$$

Using (5.12) we obtain

$$\begin{aligned}
& \int_0^t \int_\sigma^t \left| F^s(z_-(v_-, \tau) + x_- + f_1(\tau)) - F^s(z_-(v_-, \tau) + x_- + f_2(\tau)) \right| d\tau d\sigma \\
& \leq \frac{n\beta_3^s \sup_{s \in [0, +\infty)} \frac{|(f_1 - f_2)(s)|}{1+s}}{\alpha(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} \left[\frac{1}{\left(\frac{|v_-|}{2\sqrt{2}} - r \right)} + \frac{\alpha}{(\alpha + 2) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)} \right], \tag{5.22}
\end{aligned}$$

and

$$\begin{aligned}
& \int_0^t \int_\sigma^t \left| F^l(z_-(v_-, \tau) + x_- + f_2(\tau)) - F^l(z_-(v_-, \tau) + x_- + f_1(\tau)) \right| d\tau d\sigma \\
& \leq \frac{n\beta_2 \sup_{s \in [0, +\infty)} \frac{|(f_1 - f_2)(s)|}{1+s}}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} \left(\frac{1}{(\alpha + 1)} + t \right). \tag{5.23}
\end{aligned}$$

Therefore using also (5.16) (with " $t = 0$ "), (5.20) and (5.21), we have

$$\begin{aligned}
|A(f_1)(t) - A(f_2)(t)| &\leq \frac{n\|f_1 - f_2\|}{(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} \left[\frac{\beta_2}{\alpha} + \frac{\beta_3^s}{(\alpha + 2)\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)} \right] \\
&+ \frac{n\beta_3^s \sup_{s \in [0, +\infty)} \frac{|(f_1 - f_2)(s)|}{1+s}}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^3\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} \\
&+ t \frac{n\|f_1 - f_2\|}{\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^{\alpha+1}} \left[\frac{\beta_2}{(\alpha + 1)} + \frac{\beta_3^s}{(\alpha + 2)\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)} \right] \\
&+ t \frac{n \sup_{s \in [0, +\infty)} \frac{|(f_1 - f_2)(s)|}{1+s}}{\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} \left[\frac{\beta_3^s}{(\alpha + 1)\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)} + \frac{2\beta_2}{\alpha} \right]. \tag{5.24}
\end{aligned}$$

□

Proof of Lemma 3.1. We follow the proof of Lemma 2.2.

Estimate (4.7) for $\tau \in \mathbb{R}$, $f \in M_r$ still holds and we have

$$\begin{aligned}
|z_-(v_-, x_-, \tau) + f(\tau)| &\geq \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{\sqrt{2}} - \frac{2^{\frac{5}{2}} n^{\frac{1}{2}} \beta_1^l}{|v_-| \alpha \left(1 + \frac{|x_-|}{\sqrt{2}}\right)^\alpha} - r \right) \\
&\geq \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r \right). \tag{5.25}
\end{aligned}$$

We used (2.4) (with " $(x, w, h) = (x_-, v_-, 0)$ ") and the inequality $|v_-| \geq \mu(|x_-|)$. Then using (1.3), (1.4) and (5.25) we obtain

$$|F^s(z_-(v_-, x_-, \tau) + f(\tau))| \leq \frac{\beta_2 \sqrt{n}}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+2}}, \tag{5.26}$$

and

$$|F^l(z_-(v_-, x_-, \tau) + f(\tau)) - F^l(z_-(v_-, x_-, \tau))| \leq \frac{\beta_2 n |f(\tau)|}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+2}}, \tag{5.27}$$

for $\tau \in \mathbb{R}$. Then the proof of the following estimates (5.28), (5.29), (5.30), (5.31) is similar to the proof of the estimates (5.4), (5.5), (5.7) and (5.11) respectively, and we have

$$|\dot{\mathcal{A}}(f)(t)| \leq \frac{\beta_2 (n \sup_{(-\infty, 0)} |f| + \sqrt{n})}{(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r\right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r + |t| \left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+1}}, \tag{5.28}$$

for $t \leq 0$,

$$|\mathcal{A}(f)(t)| \leq \frac{\beta_2(n \sup_{(-\infty, 0)} |f| + \sqrt{n})}{\alpha(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r + |t| \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \right)^\alpha}, \quad (5.29)$$

for $t \leq 0$,

$$|\dot{\mathcal{A}}(f)(t)| \leq \frac{\beta_2(n \|f\| + 2\sqrt{n})}{(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{\alpha+1}} + \frac{n\beta_2 \sup_{s \in (0, +\infty)} \frac{|f(s)|}{1+s}}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha}, \quad (5.30)$$

for $t \geq 0$, and

$$\begin{aligned} |\mathcal{A}(f)(t)| &\leq \frac{\beta_2(n(\sup_{(-\infty, 0)} |f|) + 2\sqrt{n})}{\alpha(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} \\ &+ t\beta_2 \left(\frac{(2n\|f\| + 2\sqrt{n})}{(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{\alpha+1}} + \frac{2n \sup_{s \in (0, +\infty)} \frac{|f(s)|}{1+s}}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} \right), \end{aligned} \quad (5.31)$$

for $t \geq 0$. Estimate (3.5) follows from (5.29) and (5.31).

Let $(f_1, f_2) \in M_r^2$. Using (1.3), (1.4) and (5.25) we obtain

$$\begin{aligned} |F^l(z_-(v_-, x_-, \tau) + f_1(\tau)) - F^l(z_-(v_-, x_-, \tau) + f_2(\tau))| &\leq \frac{\beta_2 n |f_1 - f_2|(\tau)}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \right)^{\alpha+2}}, \\ &\quad (5.32) \\ |F^s(z_-(v_-, x_-, \tau) + f_1(\tau)) - F^s(z_-(v_-, x_-, \tau) + f_2(\tau))| &\leq \frac{\beta_3^s n |f_1 - f_2|(\tau)}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \right)^{\alpha+3}}, \\ &\quad (5.33) \end{aligned}$$

for $\tau \in \mathbb{R}$. Then similarly to the proof of (5.15) (resp. (5.16)) and (5.20) (resp. (5.24)) we prove the following: $\dot{\mathcal{A}}(f_1)(t) - \dot{\mathcal{A}}(f_2)(t)$ (resp. $\mathcal{A}(f_1)(t) - \mathcal{A}(f_2)(t)$) for $t \leq 0$ is bounded by the right-hand side of (5.15) (resp. (5.16)); $\dot{\mathcal{A}}(f_1)(t) - \dot{\mathcal{A}}(f_2)(t)$ (resp. $\mathcal{A}(f_1)(t) - \mathcal{A}(f_2)(t)$) for $t \geq 0$ is bounded by the right-hand side of (5.20) (resp. (5.24)). Then estimate (3.6) follows from these latter bounds on $\mathcal{A}(f_1)(t) - \mathcal{A}(f_2)(t)$, $t \in \mathbb{R}$.

It remains to prove (3.7). Let $f \in M_r$. Using (1.3) we obtain

$$|F^l(z_-(v_-, x_-, \tau) + f(\tau))| \leq \frac{\beta_1^l \sqrt{n}}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \right)^{\alpha+1}}, \quad (5.34)$$

for $\tau \in \mathbb{R}$. Then from (3.4), (5.26) and (5.34) it follows that

$$\begin{aligned} |\tilde{k}(v_-, x_-, f) - v_-| &\leq \int_{-\infty}^{+\infty} \frac{\sqrt{n}\beta_1^l d\tau}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau|\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+1}} \\ &\quad + \int_{-\infty}^{+\infty} \frac{\sqrt{n}\beta_2 d\tau}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau|\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+2}}, \end{aligned}$$

which proves (3.7). \square

6 Proof of Theorem 2.4

The estimate (2.19) (resp. (2.20)) follows from the assumption $y_- = A(y_-)$, and $\sup_{(-\infty, 0)} |y_-| \leq r$ and the estimate (5.4) (resp. (5.5)). Using (1.4) and (4.8) we obtain

$$|F^l(z_-(v_-, \tau) + x_- + y_-(\tau))| \leq \frac{\beta_1^l \sqrt{n}}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau|\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+1}}, \quad (6.1)$$

for $\tau \in \mathbb{R}$. Using (2.13), (5.1) and (6.1) we obtain

$$\begin{aligned} |a_{sc}(v_-, x_-)| &\leq \int_{-\infty}^{+\infty} \frac{\beta_1^l n^{\frac{1}{2}}}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + \left(\frac{|v_-|}{2\sqrt{2}} - r\right)|\tau|\right)^{\alpha+1}} d\tau \\ &\quad + \int_{-\infty}^{+\infty} \frac{\beta_2 n^{\frac{1}{2}}}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + \left(\frac{|v_-|}{2\sqrt{2}} - r\right)|\tau|\right)^{\alpha+2}} d\tau, \end{aligned}$$

which proves (2.21).

Then using (2.16), (5.1) and (5.2) we have

$$\begin{aligned} |l(v_-, x_-, y_-)| &\leq 2\beta_2 n^{\frac{1}{2}} \int_{-\infty}^0 \int_{-\infty}^{\sigma} \left(1 + \frac{|x_-|}{\sqrt{2}} - r + \left(\frac{|v_-|}{2\sqrt{2}} - r\right)|\tau|\right)^{-\alpha-2} d\tau d\sigma \\ &\quad + \beta_2 n \int_{-\infty}^0 \int_{-\infty}^{\sigma} \frac{|x_-| + \sup_{(-\infty, 0)} |y_-|}{\left(1 - r + \left(\frac{|v_-|}{2\sqrt{2}} - r\right)|\tau|\right)^{\alpha+2}} d\tau d\sigma \\ &\leq \frac{2\beta_2 n^{\frac{1}{2}}}{\alpha(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^{\alpha}} + \frac{\beta_2 n(|x_-| + \sup_{(-\infty, 0)} |y_-|)}{\alpha(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 (1-r)^{\alpha}} \\ &\leq \frac{\beta_2(n|x_-| + 2n^{\frac{1}{2}} + n \sup_{(-\infty, 0)} |y_-|)}{\alpha(\alpha+1)(1-r)^{\alpha}\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2}. \end{aligned}$$

This estimate with the estimate $\sup_{(-\infty, 0)} |y_-| \leq r$ proves (2.22). Now using (2.13) and (5.20) we obtain that

$$\begin{aligned}
& |a_{sc}(v_-, x_-) - \int_{-\infty}^{+\infty} F(z_-(v_-, \tau) + x_-)| = \lim_{t \rightarrow +\infty} |\dot{A}(y_-)(t) - \dot{A}(0)(t)| \\
& \leq \frac{n \|y_-\|}{\left(\frac{|v_-|}{2\sqrt{2}} - r\right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^{\alpha+1}} \left[\frac{\beta_2}{(\alpha+1)} + \frac{\beta_3^s}{(\alpha+2) \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)} \right] \\
& + \frac{n \sup_{s \in [0, +\infty)} \frac{|y_-(s)|}{1+s}}{\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} \left[\frac{\beta_3^s}{(\alpha+1) \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)} + \frac{\beta_2}{\alpha} \right]. \quad (6.2)
\end{aligned}$$

Then note that $\|y_-\| = \|A(y_-)\|$ is bounded by the right hand side of (2.10). Hence combining this latter bound on $\|y_-\|$ and (6.2) and the estimate $1 + \frac{|x_-|}{\sqrt{2}} - r \geq 1 - r$ we obtain (2.23).

Using (2.16), (5.12) and (5.13) we obtain

$$\begin{aligned}
& |l(v_-, x_-, y_-) - l(v_-, x_-, 0)| \\
& \leq \int_{-\infty}^0 \int_{-\infty}^\sigma |F^l(z_-(v_-, \tau) + x_- + y_-(\tau)) - F^l(z_-(v_-, \tau) + x_-)| d\tau d\sigma \\
& + \int_{-\infty}^0 \int_{-\infty}^\sigma |F^s(z_-(v_-, \tau) + x_- + y_-(\tau)) - F^s(z_-(v_-, \tau) + x_-)| d\tau d\sigma \\
& + \int_0^{+\infty} \int_\sigma^{+\infty} |F^s(z_-(v_-, \tau) + x_- + y_-(\tau)) - F^s(z_-(v_-, \tau) + x_-)| d\tau d\sigma \\
& \leq \frac{n}{(\alpha+1) \left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} \left(\frac{\beta_2 \sup_{s \in (-\infty, 0)} |y_-(s)|}{\alpha} + \frac{\beta_3^s \|y_-\|}{(\alpha+2) \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)} \right. \\
& \left. + \frac{\beta_3^s \sup_{s \in (-\infty, 0)} \frac{|y_-(s)|}{1+s}}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r\right)} \right). \quad (6.3)
\end{aligned}$$

Then $\|y_-\|$ is bounded by the right hand side of (2.10), and combining this latter bound on $\|y_-\|$ and (6.3) and the estimate $1 + \frac{|x_-|}{\sqrt{2}} - r \geq 1 - r$ we obtain (2.24).

It remains to prove (2.26), (2.27) and (2.28). From (2.21) it follows that

$$|a_{sc}(v_-, x_-)| \leq \frac{4n^{\frac{1}{2}} \max(\beta_1^l, \beta_2)}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r\right) (1-r)^{\alpha+1}}. \quad (6.4)$$

In addition using (2.4) (for $”(w, v, x, h) = (a(v_-, x_-), a(v_-, x_-), 0, 0)”$) and the identity $|a(v_-, x_-)| = |v_-|$ that follows from the conservation of energy,

we have

$$|z_+(a(v_-, x_-), t) - ta(v_-, x_-)| \leq \frac{2^{\frac{5}{2}} n^{\frac{1}{2}} \beta_1^l}{\alpha |v_-|} |t|, \quad t \in \mathbb{R}. \quad (6.5)$$

Then using (4.7) for " $f = y_-$ " and using (6.4), (6.5) and (2.4) (for " $(w, v, x, h) = (v_-, v_-, 0, 0)$ ") we obtain

$$\begin{aligned} & |x_- + \eta(z_-(v_-, t) + y_-(t)) + (1 - \eta)z_+(a(v_-, x_-), t)| \\ & \geq |x_- + tv_-| - |y_-(t)| - |z_-(v_-, t) - v_-t| - |z_+(a(v_-, x_-), t) - ta(v_-, x_-)| \\ & \quad - |t| |a_{sc}(v_-, x_-)| \\ & \geq \frac{|x_-|}{\sqrt{2}} - r + |t| \left(\frac{|v_-|}{\sqrt{2}} - \frac{2^{\frac{7}{2}} n^{\frac{1}{2}} \beta_1^l}{\alpha |v_-|} - \frac{4 \max(\beta_1^l, \beta_2) n^{\frac{1}{2}}}{\alpha (\frac{|v_-|}{2\sqrt{2}} - r)(1 - r)^{\alpha+1}} \right) \\ & \geq \frac{|x_-|}{\sqrt{2}} - r + |t| \left(\frac{|v_-|}{\sqrt{2}} - \frac{8 \max(\beta_1^l, \beta_2) n^{\frac{1}{2}}}{\alpha (\frac{|v_-|}{2\sqrt{2}} - r)(1 - r)^{\alpha+1}} \right) \geq \frac{|x_-|}{\sqrt{2}} - r + |t| \frac{|v_-|}{2\sqrt{2}}, \quad (6.6) \end{aligned}$$

for $\eta \in (0, 1)$ and $t \in \mathbb{R}$ (we used (2.25) and we used the estimate $|x_- + tv_-| \geq \frac{|x_-|}{\sqrt{2}} + |t| \frac{|v_-|}{\sqrt{2}}$ that follows from $x_- \cdot v_- = 0$). Similarly

$$|\eta x_- + z_+(a(v_-, x_-), t)| \geq \frac{|v_-|}{2\sqrt{2}} |t|, \quad \text{for } (\eta, t) \in (0, 1) \times \mathbb{R}. \quad (6.7)$$

From (2.17), (1.3) and (6.7) it follows that

$$\begin{aligned} |l_1(v_-, x_-)| & \leq \beta_2 n |x_-| \int_0^{+\infty} \int_{\sigma}^{+\infty} \sup_{\eta \in (0, 1)} (1 + |\eta x_- + z_+(a(v_-, x_-), \tau)|)^{-\alpha-2} d\tau d\sigma \\ & \leq \beta_2 n |x_-| \int_0^{+\infty} \int_{\sigma}^{+\infty} \left(1 + \frac{|v_-|}{2\sqrt{2}} |\tau|\right)^{-\alpha-2} d\tau d\sigma, \end{aligned}$$

which gives (2.26). Using (2.18), (1.3) and (6.6) we obtain

$$|l_2(v_-, x_-, y_-)| \leq n \beta_2 \int_0^{+\infty} \int_{\sigma}^{+\infty} \frac{|z_-(v_-, \tau) + y_-(\tau) - z_+(a(v_-, x_-), \tau)| d\tau d\sigma}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + |\tau| \frac{|v_-|}{2\sqrt{2}}\right)^{\alpha+2}}. \quad (6.8)$$

Then using (2.12) and (2.14) we have

$$\begin{aligned} & |z_-(v_-, \tau) + y_-(\tau) - z_+(a(v_-, x_-), \tau)| \\ & \leq |l(v_-, x_-, y_-)| + |l_1(v_-, x_-)| + |l_2(v_-, x_-, y_-)| + |y_+(\tau)|, \quad (6.9) \end{aligned}$$

for $\tau \in (0, +\infty)$. Combining (6.8) and (6.9) we obtain

$$\begin{aligned} & |l_2(v_-, x_-, y)| \leq \varepsilon(v_-, x_-, 0) \\ & \times (|l_2(v_-, x_-, y_-)| + |l(v_-, x_-, y_-)| + |l_1(v_-, x_-)| + \sup_{(0, +\infty)} |y_+|), \quad (6.10) \end{aligned}$$

where

$$\varepsilon(v_-, x_-, t) := \frac{n\beta_2}{\alpha(\alpha+1)\left(\frac{|v_-|}{2^{\frac{3}{2}}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r + t\left(\frac{|v_-|}{2^{\frac{3}{2}}} - r\right)\right)^\alpha}, \quad (6.11)$$

for $t \geq 0$. From (2.15) and (5.1) it follows that

$$\begin{aligned} |y_+(t)| &\leq n^{-\frac{1}{2}}\varepsilon(v_-, x_-, t) \\ &\quad + \int_t^{+\infty} \int_\sigma^{+\infty} |F^l(z_-(v_-, \tau) + x_- + y_-(\tau)) - F^l(z_+(a(v_-, x_-), \tau))| d\tau d\sigma, \end{aligned} \quad (6.12)$$

for $t \geq 0$. Then similarly to (6.10) we have

$$\begin{aligned} \sup_{(t, +\infty)} |y_+| &\leq \varepsilon(v_-, x_-, t)(n^{-\frac{1}{2}} + |l_2(v_-, x_-, y_-)| + |l(v_-, x_-, y_-)| \\ &\quad + |l_1(v_-, x_-)| + \sup_{(t, +\infty)} |y_+|), \end{aligned} \quad (6.13)$$

for $t \geq 0$. From (6.10) and (6.13) it follows that

$$(1 - \varepsilon(v_-, x_-, 0))|l_2(v_-, x_-, y_-)| \leq \varepsilon(v_-, x_-, 0)(|l(v_-, x_-, y_-)| + |l_1(v_-, x_-)| + \sup_{(0, +\infty)} |y_+|), \quad (6.14)$$

and

$$(1 - \varepsilon(v_-, x_-, t)) \sup_{(t, +\infty)} |y_+| \leq \varepsilon(v_-, x_-, t)(n^{-\frac{1}{2}} + |l_2(v_-, x_-, y_-)| + |l(v_-, x_-, y_-)| + |l_1(v_-, x_-)|), \quad (6.15)$$

for $t \geq 0$. Using (2.25) and (6.11) we have

$$\sup_{t \in (0, +\infty)} \varepsilon(v_-, x_-, t) = \varepsilon(v_-, x_-, 0) \leq 8^{-1}. \quad (6.16)$$

Then multiplying (6.14) by $(1 - \varepsilon(v_-, x_-, 0))$ and using (6.15) for $t = 0$ we obtain

$$\begin{aligned} (1 - 2\varepsilon(v_-, x_-, 0))|l_2(v_-, x_-, y_-)| &\leq \varepsilon(v_-, x_-, 0)^2 n^{-\frac{1}{2}} \\ &\quad + \varepsilon(v_-, x_-, 0)(|l(v_-, x_-, y_-)| + |l_1(v_-, x_-)|). \end{aligned} \quad (6.17)$$

Using (6.17) and (6.16) we have

$$|l_2(v_-, x_-, y_-)| \leq 2\varepsilon(v_-, x_-, 0)(n^{-\frac{1}{2}}\varepsilon(v_-, x_-, 0) + |l(v_-, x_-, y_-)| + |l_1(v_-, x_-)|). \quad (6.18)$$

Then (2.27) follows from (6.11) (for $t = 0$), (2.22) and (5.20). Using (6.16) and (6.15) we obtain

$$\sup_{(t, +\infty)} |y_+| \leq 2\varepsilon(v_-, x_-, t)(n^{-\frac{1}{2}} + |l_2(v_-, x_-, y_-)| + |l(v_-, x_-, y_-)| + |l_1(v_-, x_-)|), \quad (6.19)$$

for $t \geq 0$. Then (2.28) follows from (2.27) (combined with (2.25)), (2.22), (2.26) and (6.11). \square

7 Proof of Lemma 3.2 and Theorem 3.3

7.1 Preliminary lemmas

Lemma 7.1. *Let $(v, x, w, h) \in (\mathbb{R}^n)^4$ and $x \in \mathbb{R}^n$ so that $v \cdot x = 0$. Under the assumptions of Lemma 2.1 the following estimates are valid*

$$|z_+(w, x + h, t) - z_+(w, x + h', t)| \leq 2|h - h'|, \quad (7.1)$$

$$|\eta z_+(w, x + h', t) + (1 - \eta)z_+(w, x + h, t)| \geq \frac{|x|}{\sqrt{2}} - |h| + t \frac{|v|}{2\sqrt{2}}, \quad (7.2)$$

for $t \geq 0$, $\eta \in (0, 1)$ and $h' \in \mathbb{R}^n$, $|h'| \leq |h|$.

Proof of Lemma 7.1. First we prove (7.2). We estimate $\eta z_+(w, x + h', \tau) + (1 - \eta)z_+(w, x + h, t)$ for $\eta \in (0, 1)$ and $t \geq 0$, and for $(h, h') \in \mathcal{B}(0, 1 + \frac{|x|}{\sqrt{2}})^2$ so that $|h'| \leq |h|$. Using (2.4) we obtain

$$\begin{aligned} & |\eta z_+(w, x + h', t) + (1 - \eta)z_+(w, x + h, t)| \\ & \geq |x + tv| - |\eta h' + (1 - \eta)h| - t|v - w| - \eta|z_+(w, x + h', t) - x - h' - tw| \\ & \quad - (1 - \eta)|z_+(w, x + h, t) - x - h - tw| \\ & \geq \frac{|x|}{\sqrt{2}} - |h| + t \left(\frac{3|v|}{4\sqrt{2}} - (1 - \eta) \frac{2^{\frac{5}{2}} n^{\frac{1}{2}} \beta_1^l}{\alpha|v|(1 + \frac{|x|}{\sqrt{2}} - |h|)^\alpha} - \eta \frac{2^{\frac{5}{2}} n^{\frac{1}{2}} \beta_1^l}{\alpha|v|(1 + \frac{|x|}{\sqrt{2}} - |h'|)^\alpha} \right), \end{aligned} \quad (7.3)$$

for $t \geq 0$. Estimate (7.2) follows from (7.3) and assumption (2.2).

Now we set $\tilde{\delta}(t) := z_+(w, x + h', t) - z_+(w, x + h, t)$ for $t \geq 0$, where $|h'| \leq |h|$. From Lemma 4.1 it follows that $\sup_{(0, +\infty)} |\tilde{\delta}| < \infty$. Then from (1.6) with the boundary conditions (2.3) it follows that

$$\tilde{\delta}(t) = h' - h - \int_0^t \int_\sigma^{+\infty} (F^l(z_+(w, x + h', \tau)) - F^l(z_+(w, x + h, \tau))) d\tau d\sigma, \quad (7.4)$$

for $t \geq 0$, and using (1.3) and (7.2) we obtain

$$|\tilde{\delta}(t)| \leq |h' - h| + n\beta_2 \sup_{(0, +\infty)} |\tilde{\delta}| \int_0^t \int_\sigma^{+\infty} \left(1 + \frac{|x|}{\sqrt{2}} - |h| + \tau \frac{|v|}{2\sqrt{2}} \right)^{-\alpha-2} d\tau d\sigma, \quad (7.5)$$

for $t \geq 0$, which proves that

$$\left(1 - \frac{n\beta_2}{\alpha(\alpha+1) \left(\frac{|v|}{2\sqrt{2}} \right)^2 \left(1 + \frac{|x|}{\sqrt{2}} - |h| \right)^\alpha} \right) \sup_{(0, +\infty)} |\tilde{\delta}| \leq |h' - h|. \quad (7.6)$$

Using (2.2) we have $1 - \frac{8n\beta_2}{\alpha(\alpha+1)|v|^2 \left(1 + \frac{|x|}{\sqrt{2}} - |h| \right)^\alpha} \geq \frac{1}{2}$. Combining this latter estimate and (7.6) we obtain (7.1). \square

We also need the following lemma.

Lemma 7.2. *Let $(v_-, x_-) \in \mathbb{R}^n \times \mathbb{R}^n$, $v_- \cdot x_- = 0$ and let $r > 0$, $r < \frac{1}{2} + \frac{|x_-|}{2^{\frac{3}{2}}}$, and let $y_- \in M_r$. Under the assumptions of Lemma 3.2 the following estimates are valid*

$$|z_-(v_-, x_-, t) + y_-(t) - z_+(\tilde{a}(v_-, x_-), x_-, t)| \leq \frac{20\beta_2(nr + 3\sqrt{n})}{19\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha}, \quad (7.7)$$

$$|\eta(z_-(v_-, x_-, t) + y_-(t)) + (1 - \eta)z_+(\tilde{a}(v_-, x_-), x_-, t)| \geq \frac{|x_-|}{\sqrt{2}} - r + |t|\left(\frac{|v_-|}{2\sqrt{2}} - r\right), \quad (7.8)$$

for $t \geq 0$ and $\eta \in (0, 1)$.

Proof of Lemma 7.2. We set $\delta(t) := z_-(v_-, x_-, t) + y_-(t) - z_+(\tilde{a}(v_-, x_-), x_-, t)$ for $t \geq 0$. From Lemma 4.1 it follows that $\sup_{(0, +\infty)} |\delta| < \infty$. Using (1.1) and (1.6) we obtain

$$\begin{aligned} \delta(t) &= - \int_0^t \int_\sigma^{+\infty} (F^l(z_-(v_-, x_-, \tau) + y_-(\tau)) - F^l(z_+(\tilde{a}(v_-, x_-), x_-, \tau))) d\tau d\sigma \\ &\quad + \tilde{l}(v_-, x_-, y_-) + H_1(v_-, x_-, y_-)(t), \end{aligned} \quad (7.9)$$

for $t \geq 0$, where \tilde{l} is defined by (3.12) and where

$$H_1(v_-, x_-, y_-)(t) = \int_t^{+\infty} \int_\sigma^{+\infty} F^s(z_-(v_-, x_-, \tau) + y_-(\tau)) d\tau d\sigma. \quad (7.10)$$

Then from (3.12), (5.26), (5.27) and the estimate $|y_-(\tau)| \leq r$ for $\tau \leq 0$ we obtain

$$|\tilde{l}(v_-, x_-, y_-)| \leq \frac{\beta_2(nr + 2\sqrt{n})}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha}. \quad (7.11)$$

Using (7.10) and (5.26) we obtain

$$|H_1(v_-, x_-, y_-)(t)| \leq \frac{\beta_2\sqrt{n}}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha}, \quad (7.12)$$

for $t \geq 0$. From (3.7) and the inequality $r \leq \frac{1}{2} + \frac{|x_-|}{2^{\frac{3}{2}}}$ and from (3.10) it follows that

$$|\tilde{a}_{sc}(v_-, x_-)| \leq \frac{6\sqrt{n} \max(\beta_1^l, \beta_2)}{\alpha\left(\frac{|v_-|}{2^{\frac{3}{2}}} - r\right)(1 - r + \frac{|x_-|}{\sqrt{2}})^\alpha}. \quad (7.13)$$

Then from (3.8) it follows that $|a_{sc}(v_-, x_-)| \leq \frac{|v_-|}{4\sqrt{2}}$, and using (7.2) we obtain $|z_+(a(v_-, x_-), x_-, t)| \geq \frac{|x_-|}{\sqrt{2}} + t \frac{|v_-|}{2\sqrt{2}}$ for $t \geq 0$. And using (2.4) we obtain

$$\begin{aligned} & |\eta(z_-(v_-, x_-, t) + y_-(t)) + (1 - \eta)z_+(\tilde{a}(v_-, x_-), x_-, t)| \\ & \geq |x_- + tv_-| - (1 - \eta)(|z_+(\tilde{a}(v_-, x_-), x_-, t) - x_- - t\tilde{a}(v_-, x_-)| + t|\tilde{a}_{sc}(v_-, x_-)|) \\ & \quad - \eta|z_-(v_-, x_-, t) - x_- - v_-t| - \eta|y_-(t)| \\ & \geq \frac{|x_-|}{\sqrt{2}} - r + |t| \left(\frac{|v_-|}{\sqrt{2}} - \frac{6\sqrt{n}\beta}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} - \frac{2^{\frac{7}{2}} n^{\frac{1}{2}} \beta_1^l}{\alpha |v_-| \left(1 + \frac{|x_-|}{\sqrt{2}} \right)^\alpha} - r \right), \end{aligned} \quad (7.14)$$

for $t \geq 0$. Combining (7.14) and (3.8) we obtain (7.8).

From (7.9), (1.3) and (7.14) it follows that

$$|\delta(t)| \leq C(v_-, x_-, r) + n\beta_2 \sup_{(0, +\infty)} |\delta| \int_0^t \int_\sigma^{+\infty} \left(1 + \frac{|x_-|}{\sqrt{2}} - r + \tau \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \right)^{-\alpha-2} d\tau d\sigma, \quad (7.15)$$

for $t \geq 0$, where

$$C(v_-, x_-, r) := \frac{\beta_2(nr + 3\sqrt{n})}{\alpha(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} \geq |\tilde{l}(v_-, x_-, y_-)| + \sup_{(0, +\infty)} |H_1(v_-, x_-, y_-)| \quad (7.16)$$

(see (7.11) and (7.12)). From (7.15) it follows that

$$\left(1 - \frac{n\beta_2}{\alpha(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} \right) \sup_{(0, +\infty)} |\delta| \leq C(v_-, x_-, r). \quad (7.17)$$

Using (3.8) we have $1 - \frac{n\beta_2}{\alpha(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^\alpha} \geq \frac{19}{20}$, and then estimate (7.7) follows from (7.17) and (7.16). Hence Lemma 7.2 is proved. \square

7.2 Proof of Lemma 3.2

Let $(h, h') \in \overline{\mathcal{B}(0, \frac{1}{4} + \frac{|x_-|}{2\sqrt{2}})}$, $|h'| \leq |h|$. Using (7.1) and (7.8) we have

$$\begin{aligned} & |\eta(z_-(v_-, x_-, t) + y_-(t)) + (1 - \eta)z_+(\tilde{a}(v_-, x_-), x_- + h, t)| \\ & \geq |\eta(z_-(v_-, x_-, t) + y_-(t)) + (1 - \eta)z_+(\tilde{a}(v_-, x_-), x_-, t)| \\ & \quad - (1 - \eta)|z_+(\tilde{a}(v_-, x_-), x_-, t) - z_+(\tilde{a}(v_-, x_-), x_- + h, t)| \\ & \geq \frac{|x_-|}{\sqrt{2}} - r - 2|h| + |t| \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \geq \frac{|x_-|}{2\sqrt{2}} - \frac{1}{2} - r + |t| \left(\frac{|v_-|}{2\sqrt{2}} - r \right), \end{aligned} \quad (7.18)$$

for $t \geq 0$ and $\eta \in (0, 1)$. Using (7.1) and (7.7) we obtain

$$|z_-(v_-, x_-, t) + y_-(t) - z_+(\tilde{a}(v_-, x_-), x_- + h, t)| \leq \frac{20\beta_2(nr + 3\sqrt{n})}{19\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} + 2|h|, \quad (7.19)$$

Combining (7.11), (3.11), (1.3), (7.18) and (7.19) we have

$$\begin{aligned} |\mathcal{G}_{v_-, x_-}(h)| &\leq \frac{\beta_2(nr + 2\sqrt{n})}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} \\ &\quad + n\beta_2 \int_0^{+\infty} \int_\sigma^{+\infty} \frac{\frac{20\beta_2(nr + 3\sqrt{n})}{19\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} + 2|h|}{\left(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}} - r + \tau\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+2}} d\tau d\sigma \\ &\leq \frac{\beta_2(3(nr + \sqrt{n}) + 2|h|n)}{\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}} - r\right)^\alpha} \end{aligned} \quad (7.20)$$

(we used (3.8) which implies $\frac{60n\beta_2}{19\alpha(\alpha + 1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} \leq 1$). Then using

(7.20) and the estimate $|h| \leq \frac{1}{4} + \frac{|x_-|}{4\sqrt{2}}$ we obtain (3.14). Using again (3.8) and (3.14) and the estimate $r \leq \frac{1}{2} + \frac{|x_-|}{2\sqrt{2}}$ we obtain (3.15).

Then using (3.11), (1.3), (7.2) and (7.1) we have

$$\begin{aligned} |\mathcal{G}_{v_-, x_-}(h) - \mathcal{G}_{v_-, x_-}(h')| &\leq \int_0^{+\infty} \int_\sigma^{+\infty} |F^l(z_+(\tilde{a}(v_-, x_-), x_- + h', \tau)) \\ &\quad - F^l(z_+(\tilde{a}(v_-, x_-), x_- + h, \tau))| d\tau d\sigma \\ &\leq n\beta_2^l \int_0^{+\infty} \int_\sigma^{+\infty} \frac{2|h - h'| d\tau d\sigma}{\left(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}} + \tau\frac{|v_-|}{2\sqrt{2}}\right)^{\alpha+2}}, \end{aligned} \quad (7.21)$$

for $(h, h') \in \mathcal{B}(0, \frac{1}{4} + \frac{|x_-|}{4\sqrt{2}})^2$, which proves the first estimate in (3.16). The second estimate in (3.16) follows from (3.8). \square

7.3 Proof of Theorem 3.3

The estimate (3.21) (resp. (3.22)) follows from (5.28) (resp. (5.29)) and the equality $y_- = \mathcal{A}(y_-)$. For estimate (3.23) see also (7.13). We prove (3.24). Since $\tilde{b}_{sc}(v_-, x_-) = \mathcal{G}_{v_-, x_-}(\tilde{b}_{sc}(v_-, x_-))$ we use (7.20) with " $h = \tilde{b}_{sc}(v_-, x_-)$ "

and we obtain

$$\begin{aligned} & \left(1 - \frac{2n\beta_2}{\alpha(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}} - r\right)^\alpha} \right) |\tilde{b}_{sc}(v_-, x_-)| \\ & \leq \frac{3\beta_2(nr + \sqrt{n})}{\alpha(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}} - r\right)^\alpha}. \end{aligned} \quad (7.22)$$

Then use (3.8) to prove that $1 - \frac{2n\beta_2}{\alpha(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(\frac{1}{2} + \frac{|x_-|}{2\sqrt{2}} - r\right)^\alpha} \geq \frac{9}{10}$. This proves (3.24). We prove (3.25). From (1.3), (7.18) and (3.17) it follows that

$$\begin{aligned} & \int_t^{+\infty} \int_\sigma^{+\infty} |F^l(z_-(v_-, x_-, \tau) + y_-(\tau)) - F^l(z_+(\tilde{a}(v_-, x_-), \tilde{b}(v_-, x_-), \tau))| d\tau d\sigma \\ & \leq \beta_2 n \int_t^{+\infty} \int_\sigma^{+\infty} \frac{|y_+(\tau)| d\tau d\sigma}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + \tau\left(\frac{|v_-|}{2\sqrt{2}} - r\right)\right)^{\alpha+2}} \\ & \leq \frac{\beta_2 n \sup_{(t, +\infty)} |y_+|}{\alpha(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha}, \end{aligned} \quad (7.23)$$

for $t \geq 0$. Then using (3.18), (7.23) and (7.12) we obtain

$$|y_+(t)| \leq \frac{\beta_2(n^{\frac{1}{2}} + n \sup_{(t, +\infty)} |y_+|)}{\alpha(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha}, \quad (7.24)$$

for $t \geq 0$. Then from (3.8) it follows that $\frac{\beta_2 n}{\alpha(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha} \leq \frac{1}{20}$.

Hence from (7.24) it follows that

$$\sup_{(t, +\infty)} |y_+| \leq \frac{20\beta_2 n^{\frac{1}{2}}}{19\alpha(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^\alpha}, \quad (7.25)$$

which proves (3.25).

We prove (3.26). From the contraction estimate on $\dot{\mathcal{A}}(\cdot)(t)$ for $t \leq 0$ (see the proof of Lemma 3.1 and the right hand side of (5.15) for $(f_1, f_2, t) = (y_-, 0, 0)$) it follows that

$$\begin{aligned} & \left| \int_{-\infty}^0 (F(z_-(v_-, x_-, \tau) + y_-(\tau)) - F(z_-(v_-, x_-, \tau))) d\tau \right| \\ & \leq \frac{n \max(\beta_2, \beta_3^s) \sup_{(-\infty, 0]} |y_-|}{(\alpha+1)\left(\frac{|v_-|}{2\sqrt{2}} - r\right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)^{\alpha+1}} \left(1 + \frac{1}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r\right)}\right). \end{aligned} \quad (7.26)$$

Using (5.33) for $(f_1, f_2) = (y_-, 0)$ we also obtain

$$\begin{aligned}
& \left| \int_0^{+\infty} (F^s(z_-(v_-, x_-, \tau) + y_-(\tau)) - F^s(z_-(v_-, x_-, \tau))) d\tau \right| \\
& \leq \frac{n\beta_3^s \sup_{s \in [0, +\infty)} \frac{|y_-(s)|}{1+s}}{(\alpha+1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{\alpha+1}} \left[\frac{1}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)} \right. \\
& \quad \left. + \frac{1}{\left(\frac{|v_-|}{2\sqrt{2}} - r \right)} \right]. \tag{7.27}
\end{aligned}$$

Then using (1.3), (7.7) and (7.8) we have

$$\begin{aligned}
& \left| \int_0^{+\infty} (F^l(z_-(v_-, x_-, \tau) + y_-(\tau)) - F^l(z_+(\tilde{a}(v_-, x_-), x_-, \tau))) d\tau \right| \\
& \leq \frac{20n\beta_2^2(nr + 3\sqrt{n})}{19\alpha(\alpha+1)^2 \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^3 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{2\alpha+1}}. \tag{7.28}
\end{aligned}$$

Combining (3.10), (7.26), (7.27) and (7.28), we obtain

$$\begin{aligned}
& \left| \tilde{a}_{sc}(v_-, x_-) - \int_{-\infty}^0 F(z_-(v_-, x_-, \tau)) d\tau \right. \\
& \quad \left. - \int_0^{+\infty} F^s(z_-(v_-, x_-, \tau)) d\tau - \int_0^{+\infty} F^l(z_+(\tilde{a}(v_-, x_-), x_-, \tau)) d\tau \right| \\
& \leq \frac{n \max(\beta_2, \beta_3^s) \|y_-\|}{(\alpha+1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{\alpha+1}} \left(1 + \frac{1}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)} + \frac{1}{\left(\frac{|v_-|}{2\sqrt{2}} - r \right)} \right) \\
& \quad + \frac{20n\beta_2^2(nr + 3\sqrt{n})}{19\alpha(\alpha+1)^2 \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^3 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{2\alpha+1}}. \tag{7.29}
\end{aligned}$$

Then estimate (3.26) follows from (7.29), (3.5) and the equality $y_- = \mathcal{A}(y_-)$, and from the estimate $1 + \frac{|x_-|}{\sqrt{2}} - r \geq \frac{1}{2}$. We prove (3.27). From (3.12) it follows that

$$\begin{aligned}
& |\tilde{l}(v_-, x_-, y_-) - \tilde{l}(v_-, x_-, 0)| \\
& \leq \int_{-\infty}^0 \int_{-\infty}^{\sigma} |F(z_-(v_-, x_-, \tau) + y_-(\tau)) - F(z_-(v_-, x_-, \tau))| d\tau d\sigma \tag{7.30} \\
& \quad + \int_0^{+\infty} \int_{\sigma}^{+\infty} |F^s(z_-(v_-, x_-, \tau) + y_-(\tau)) - F^s(z_-(v_-, x_-, \tau))| d\tau d\sigma
\end{aligned}$$

From the contraction estimate on $\mathcal{A}(\cdot)(t)$ for $t \leq 0$ (see the proof of Lemma 3.1 and the right hand side of (5.16) for " $(f_1, f_2, t) = (y_-, 0, 0)$ ") we have

$$\begin{aligned} & \left| \int_{-\infty}^0 \int_{-\infty}^{\sigma} (F(z_-(v_-, x_-, \tau) + y_-(\tau)) - F(z_-(v_-, x_-, \tau))) d\tau d\sigma \right| \\ & \leq \frac{n \sup_{(-\infty, 0]} |y_-|}{(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{\alpha}} \left[\frac{\beta_2}{\alpha} + \frac{\beta_3^s}{(\alpha + 2) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)} \right], \quad (7.31) \end{aligned}$$

and using (5.33) for " $(f_1, f_2) = (y_-, 0)$ " we obtain

$$\begin{aligned} & \int_0^{+\infty} \int_{\sigma}^{+\infty} \left| F^s(z_-(v_-, x_-, \tau) + y_-(\tau)) - F^s(z_-(v_-, x_-, \tau)) \right| d\tau d\sigma \\ & \leq n\beta_3^s \sup_{s \in (0, +\infty)} \frac{|y_-(s)|}{1+s} \int_0^{+\infty} \int_{\sigma}^{+\infty} \frac{(1+\tau)d\tau d\sigma}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r + \tau \left(\frac{|v_-|}{2\sqrt{2}} - r \right) \right)^{\alpha+3}} \\ & \leq \frac{n\beta_3^s \sup_{s \in (0, +\infty)} \frac{|y_-(s)|}{1+s}}{(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{\alpha}} \left[\frac{1}{\alpha \left(\frac{|v_-|}{2\sqrt{2}} - r \right)} + \frac{1}{(\alpha + 2) \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)} \right]. \quad (7.32) \end{aligned}$$

Hence

$$\begin{aligned} & |\tilde{l}(v_-, x_-, y_-) - \tilde{l}(v_-, x_-, 0)| \leq \frac{n \max(\beta_2, \beta_3^s) \|y_-\|}{\alpha(\alpha + 1) \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^2 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{\alpha}} \\ & \times \left(1 + \frac{1}{\left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)} + \frac{1}{\left(\frac{|v_-|}{2\sqrt{2}} - r \right)} \right). \quad (7.33) \end{aligned}$$

And combining (7.23) and (7.25) we have

$$\begin{aligned} & \int_0^{+\infty} \int_{\sigma}^{+\infty} |F^l(z_-(v_-, x_-, \tau) + y_-(\tau)) - F^l(z_+(\tilde{a}(v_-, x_-), \tilde{b}(v_-, x_-), \tau))| d\tau d\sigma \\ & \leq \frac{20\beta_2^2 n^{\frac{3}{2}}}{19\alpha^2 (\alpha + 1)^2 \left(\frac{|v_-|}{2\sqrt{2}} - r \right)^4 \left(1 + \frac{|x_-|}{\sqrt{2}} - r \right)^{2\alpha}}. \quad (7.34) \end{aligned}$$

for $t \geq 0$. Then we use the estimates (7.33) and (3.5) and the equality $y_- = \mathcal{A}(y_-)$, and we use the estimate (7.34) and the estimate $1 + \frac{|x_-|}{\sqrt{2}} - r \geq \frac{1}{2}$ to obtain (3.27). \square

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References

- [1] J. DEREZINSKI, C. GÉRARD, *Scattering theory of classical and quantum N-particle systems*, Springer-Verlag, Berlin Heidelberg, 1997.
- [2] V. ENSS, R. WEDER, *Inverse potential scattering: a geometrical approach*. Mathematical Quantum Theory II: Schrödinger operators (Feldman, J., Froese, R. and Rosen, L., eds.), CRM Proc. Lecture Notes 8, Amer. Math. Soc., Providence, R.I., 1995.
- [3] L. D. FADDEEV, *Uniqueness of solution of the inverse scattering problem*, Vestnik. Leningrad. Univ. 11:(7), 126–130 (1956).
- [4] I. M. GEL'FAND, S. G. GINDIKIN, M. I. GRAEV, *Integral geometry in affine and projective spaces*, Itogi Nauki i Tekhniki, Sovr. Prob. Mat. 16, 53-226 (1980) (Russian).
- [5] S. HELGASON, *The Radon transform on Euclidean spaces, compact two-point homogeneous spaces and Grassmann manifolds*, Acta Math. 113, 153–180 (1965).
- [6] I. W. HERBST, *Classical scattering with long range forces*, Comm. Math. Phys. 35, 193–214 (1974).
- [7] A. JOLLIVET, *On inverse scattering at fixed energy for the multidimensional Newton equation in a non-compactly supported field*, to appear in J. Inverse Ill-Posed Probl.; (see also arXiv:1210.6552).
- [8] L. D. LANDAU, E. M. LIFSCHITZ, *Mechanics*, Pergamon Press Oxford, 1960.
- [9] F. NATTERER, *The Mathematics of Computerized Tomography*, Stuttgart: Teubner and Chichester: Wiley, 1986.
- [10] R. G. NOVIKOV, *Small angle scattering and X-ray transform in classical mechanics*, Ark. Mat. 37, 141–169 (1999).
- [11] R. G. NOVIKOV, *The $\overline{\partial}$ -approach to approximate inverse scattering at fixed energy in three dimensions*, IMRP Int. Math. Res. Pap. 6, 287–349 (2005).
- [12] J. RADON, *Über die Bestimmung von Funktionen durch ihre Integralwerte längs gewisser Mannigfaltigkeiten*. Ber. Verh. Sächs. Akad. Wiss. Leipzig, Math.-Nat. K1 69, 262-277 (1917).

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