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- Saturated hydraulic conductivity determined by on
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- 4 **Abstract.** In this study we show how to use GPR data acquired along
- 5 the infiltration of water inside a single ring infiltrometer to inverse the sat-
- 6 urated hydraulic conductivity. We used Hydrus-1D to simulate the water in-
- <sup>7</sup> filtration. We generated water content profiles at each time step of infiltra-
- 8 tion, based on a particular value of the saturated hydraulic conductivity, know-
- 9 ing the other van Genuchten parameters. Water content profiles were con-
- verted to dielectric permittivity profiles using the Complex Refractive In-
- dex Method relation. We then used the GprMax suite of programs to gen-
- erate radargrams and to follow the wetting front using arrival time of elec-
- tromagnetic waves recorded by a Ground-Penetrating Radar (GPR). The-
- oretically, the 1D time convolution between reflectivity and GPR signal at
- any infiltration time step is related to the peak of the reflected amplitude
- recorded in the corresponding trace in the radargram. We used this relation-
- ship to invert the saturated hydraulic conductivity for constant and falling
- 18 head infiltrations. We present our method on synthetic examples and on two
- experiments carried out on sand soil. We further discuss on the uncertain-
- ties on the retrieved saturated hydraulic conductivity computed by our al-
- 21 gorithm from the van Genuchten parameters.

#### 1. Introduction

Soil hydraulic properties, represented by the soil water retention  $\theta(h)$  and hydraulic 22 conductivity K(h) functions, dictate water flow in the vadose zone, as well as partitioning 23 between infiltration and runoff. Their evaluation has important implications for modeling available water resources and for flood forecasting. It is also crucial in evaluating the dynamics of chemical pollutants in soil and in assessing the potential of groundwater pollution. Soil hydraulic functions can be described by several mathematical expression [Kosuqi et al., 2002, among them the van Genuchten function [van Genuchten, 1980]. The determination of the parameters defining the van Genuchten soil water retention function vanGenuchten, 1980 is usually done using laboratory experiments, such as the water hanging column [Dane and Hopmans, 2002]. The hydraulic conductivity function can be estimated either in the laboratory, or in situ using infiltration tests. Among the large number of existing infiltration tests [Angulo-Jaramillo et al., 2000, the single [Muntz et al., 1905] or double ring infiltrometers [Boivin et al., 1987 provide the field saturated hydraulic conductivity by applying a positive water pressure on the soil surface, while the disk infiltrometer [Perroux and White, 1988; Clothier and White, 1981 allows to reconstruct the hydraulic conductivity curve, by applying different water pressures smaller than or equal to zero. For infiltration tests, the volume of infiltrated water versus time is fitted to infer the soil hydraulic conductivity at or close to saturation. These tests are time-consuming and difficult to apply to landscapescale forecasting of infiltration. Furthermore, their analysis involve various simplifying

- assumptions, partly due to the ignorance of the shape of the infiltration bulb. This lack of
- 44 knowledge on the form of the infiltration bulb has to be filled to get accurate informations
- on the soil water retention  $\theta(h)$  function and consequently on hydraulic conductivity K(h)
- 46 function. This can be done by water content sensing.
- Vereecken [Vereecken et al., 2008] and Evett and Parkin [Evett and Parkin, 2005] give
- a state of the art on the different techniques available for soil moisture measurements.
- <sup>49</sup> Among the large panel presented, geophysical methods take an important part, mainly
- because they are contact free and/or easy to use. The most commonly used hydro-
- geophysical methods are electrical resistivity measurements [Goyal et al., 2006; Zhou et al.,
- <sup>52</sup> 2001] and electromagnetic methods [Sheets and Hendrickx, 1995; Akbar et al., 2005]. This
- paper focuses on the use of Ground-Penetrating Radar (GPR) as a tool for monitoring
- water infiltration in soil.
- For few decades GPR has been known as an accurate method to highlight water vari-
- ation in soils [Huisman et al., 2003; Annan, 2005]. Different techniques are available in
- 57 the literature for monitoring water content in soils using GPR. Tomography imaging be-
- tween boreholes during an infiltration has been done by Binley [Binley et al., 2001] and
- 59 Kowalsky [Kowalsky et al., 2005] among others. Many advances were done during the last
- 60 years on Off-Ground GPR using full waveform inversion, for instance to invert soil hy-
- draulic properties (Lambot [Lambot et al., 2006, 2009] and Jadoon [Jadoon et al., 2012]).
- Grote [Grote et al., 2002] and Lunt [Lunt et al., 2005] used two-way travel time variations
- from a reflector at a known depth to monitor water content variation with time. Finally,
- 64 multi-offset GPR survey techniques, i.e. CMP<sup>1</sup> or WARR<sup>2</sup>, were carried out during infil-

tration processes in the works of Greaves [Greaves et al., 1996] or Mangel [Mangel et al., 2012].

The work presented here is based on mono-offset monitoring of infiltration with on-67 ground surface GPR as related by Haarder [Haarder et al., 2011], Moysey [Moysey, 2010], Lai [Lai et al., 2012], Dagenbach [Dagenbach et al., 2013] and Saintenoy [Saintenoy et al., 2008. Haarder [Haarder et al., 2011] used a constant offset on-ground GPR coupled with dye tracing to exhibit preferential flows. They found that a GPR was able to map deep 71 infiltration comparing to dye tracer, but they did not manage to resolve the infiltration patterns (by-pass flow, fingering...). Moysey [Moysey, 2010] studied the infiltration inside 73 a sand box from the surface with on-ground GPR. He used the reflection from the wetting front as well as from the ground wave and the bottom of the box, to monitor the water content. He also modelled his experiment and estimated the van Genuchten parameters using semblance analysis. As Léger [Léger et al., 2013], he found that the most poorly constrained parameter was n. Lai [Lai et al., 2012] used a joint time frequency analysis coupled with grayscale imaging to measure infiltration and drainage in controlled conditions in laboratory. They were able to follow the peak frequency of the GPR wavelet associated with the wetting front using time frequency analysis and then determined the 81 rate of water infiltration in unsaturated zone. Saintenoy [Saintenoy et al., 2008] monitored the wetting bulb during an infiltration from a Porchet infiltrometer. They were able to identify the dimension of the bulb with time and good agreement was found with modelling.

On the continuity of those studies, we present a method for monitoring the wetting front during infiltration using on-ground GPR with fixed offset inside a ring infiltrometer.

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- The objectives of this paper were i) to check if the proposed method is accurate enough to
- monitor wetting front during infiltration with different boundary conditions, ii) to invert
- saturated hydraulic conductivity using the model of Mualem-van Genuchten [Mualem,
- 1976; van Genuchten, 1980, and iii) to analyze the uncertainties using a simplified MC
- uncertainty analysis. The method has been tested on synthetic examples and on two field
- 93 data sets.

# 2. Background

# 2.1. Unsaturated Flow Equation

- In this study we consider one-dimensionnal vertical water flow in a soil, described by
- the one-dimensional Richard's equation [Richards, 1931]. Its expression in term of water
- 96 content is

$$\frac{\partial \theta}{\partial t} = \frac{\partial K(\theta)}{\partial z} + \frac{\partial}{\partial z} \left[ D(\theta) \frac{\partial \theta}{\partial z} \right],\tag{1}$$

- where  $K(\theta)$  is the hydraulic conductivity as a function of water content, and  $D(\theta)$  is
- water diffusivity (Childs and Georges-Collis [Childs and Collis-George, 1950]), expressed
- in terms of water content as  $D(\theta) = K(\theta) \frac{\partial h}{\partial \theta}$ .

#### 2.2. Hydraulic Properties Functions

Several mathematical functions exist to model the hydraulic properties of porous media [Kosugi et al., 2002]. We chose the van Genuchten model [van Genuchten, 1980]
with the relation of Mualem [Mualem, 1976], giving the following expression for the water
retention curve:

$$\theta(h) = \theta_r + (\theta_s - \theta_r)(1 + (\alpha h)^n)^{\frac{1}{n} - 1},\tag{2}$$

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where  $\theta_s$  is the saturated water content,  $\theta_r$ , the residual water content, and  $\alpha$  and n, two fitting parameters which are respectively linked to the matric potential and the slope of the water retention curve at the inflexion point. The hydraulic conductivity function is described by

$$K(\theta) = K_s \Theta^{\lambda} \left[ 1 - \left[ 1 - \Theta^{\frac{n}{n-1}} \right]^{\frac{n}{n-1}} \right]^2, \tag{3}$$

with  $K_s$  the saturated hydraulic conductivity,  $\Theta = \frac{\theta - \theta_r}{\theta_s - \theta_r}$  the effective saturation and  $\lambda$  a factor that accounts for pore tortuosity. The  $\lambda$  parameter has an influence on the shape of the hydraulic conductivity function. However in this study we concentrated on the inversion of only one parameter, the saturated hydraulic conductivity. We fixed  $\lambda$  equal to 0.5 as reported in [Mualem, 1976].

# 2.3. Petrophysical Relationships

Several empirical and conceptual relationships exist to convert soil dielectric permittivity to volumetric water content. Using the fact that the experiments presented here
have been made in a quarry of Fontainebleau sand, considered as pure silica, we used the
CRIM relation [Birchak et al., 1974; Roth et al., 1990], which relates the relative dielectric
permittivity of bulk media,  $\varepsilon_b$ , to the volumetric summation of each components of it.
Thus for a tri-phasic medium comprising water, air and silicium, we obtain

$$\sqrt{\varepsilon_b} = \theta \sqrt{\varepsilon_w} + (1 - \phi)\sqrt{\varepsilon_s} + (\phi - \theta), \tag{4}$$

where  $\varepsilon_w = 80.1$ ,  $\varepsilon_s = 2.5$  are respectively the relative dielectric permittivity of water and silica,  $\phi$  the porosity and  $\theta$  the volumetric water content.

# 2.4. Dielectric Permittivity Versus Electromagnetic Wave Velocity

Surface GPR consists in a transmitting antenna, being a dipole, positioned on the surface, that emits short pulses of spherical electromagnetic (EM) wave in response to an excitation current source, and a receiving antenna, also located at the surface, which converts the incoming EM fields to an electrical signal source to be treated. Following the works of Annan [Annan, 1999], the velocity of electromagnetic waves is

$$v = \frac{c}{\sqrt{\varepsilon' \mu_r \frac{1 + \sqrt{1 + tan^2 \delta}}{2}}},\tag{5}$$

where  $\delta$  is the loss factor as a function of the dielectric permittivity, frequency and electrical conductivity,  $\varepsilon'$  is the real part of the relative dielectric permittivity,  $\mu_r$  the relative magnetic permeability and c is the velocity of EM waves in air equal to 0.3 m/ns. Considering the case of non magnetic soil with low conductivity, in the range of 10 MHz to 1 GHz, the real part dominates the imaginary part of the dielectric permittivity and neglecting Debye [Debye, 1929] effect, equation (5) reduces to:

$$v = \frac{c}{\sqrt{\varepsilon'}}. ag{6}$$

We used this equation to compute the travelling time of an EM wave through a layer of soil of known thickness with a given dielectric permittivity.

# 2.5. Electromagnetic Modelling

Numerous techniques are available for simulating GPR data, e.g. ray-based methods

(e.g. Cai and McMechan [Cai and McMechan, 1995] or Sethian and Popovici [Sethian

and Popovici, 1999]), time-domain finite-difference full-waveform methods (e.g. Kunz

and Luebbers [Kunz and Luebbers, 1996] or Kowalsky [Kowalsky et al., 2001]), or finite

differences time domain (FDTD) (e.g. Irving and Knight [Irving and Knight, 2006]). We

used the GprMax 2D codes of Giannopoulos [Giannopoulos, 2005], which uses FDTD modelling to solve the maxwell equations in 2 dimensions.

# 3. Materials and Methods

#### 3.1. Experimental Set-up

We studied infiltration of a 5-cm thick water layer inside of a single ring infiltrometer 141 in a sandy soil. The scheme of the apparatus is presented in Figure 1. The single ring 142 infiltrometer was a 1-mm thick aluminum cylinder with a 60-cm diameter, approximately 143 20-cm high, buried in the soil to a depth of 10 cm. GPR antennae (namely the transmitter 144 T and the receiver R) were set up at a variable distance from the edge of the cylinder, noted X, in Figure 1. In all our field experiments, we used a Mala RAMAC system with antennae centered on 1600 MHz, shielded at the top. The inner part of the cylinder was 147 covered with a plastic waterproof sheet. This allowed us to fill the cylinder with water and create an initial 5-cm thick water layer, while preventing infiltration into the sand before starting data acquisition. The beginning of the acquisition was launched by pulling away the plastic sheet to trigger water infiltration. The GPR system was set to acquire a trace every 10 s. With this apparatus, we performed two types of infiltration: i) a 152 falling head infiltration consisting of pulling away the plastic sheet and leaving water to 153 infiltrate into the sand freely with no additional refill, and ii) a constant head infiltration, 154 when water was continuously added to the ring to maintain a 5-cm thick water layer 155 during the infiltration experiment. In the following examples, we will show that GPR 156 data acquired every 10 s during the infiltration experiment can be used to estimate the 157 saturated soil hydraulic conductivity,  $K_s$ . In all GPR data presented below, we subtracted 158 the average trace and applied an Automatic Gain Control (AGC) to the data in order 159

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to make them clearer. The van Genuchten parameters,  $\alpha$ , n,  $\theta_r$ ,  $\theta_i$  of the sand have been determined in laboratory by several classical hanging water column experiments. We assumed arbitrarily a 5 % uncertainty for all the measured parameters. The sand was considered homogeneous. Its initial water content,  $\theta_i$ , and porosity,  $\phi$ , of the soil were determined using gravimetric measurements on field samples.

# 3.2. Modelling

Infiltration experiments were simulated by solving Richards equation (Eq. (1)) using Hydrus-1D. The soil profile was 50 cm deep, assumed to be homogeneous, and divided into 1001 layers. We used either an atmospheric boundary condition (BC) with no rain and 167 no evaporation at the soil surface, for the falling head infiltration, or a constant pressure 168 head of 5 cm to the top node, for the constant head infiltration, and for both case free 169 drainage BC at the bottom. To simulate the 5-cm layer of water, the initial condition 170 was set to a 5 cm pressure head in the top node. We simulated the first 10 minutes of 171 the experiment with a time step of 10 s, i.e., with 60 water content snapshots. Using 172 the CRIM relation (Eq. 4), each water content snapshot was converted to permittivity 173 profiles (made of 1001 points), considering a three-phase media: sand (considered as pure 174 silica), water, and air. Each one of these permittivity profiles were the input for the 175 GprMax2D program [Giannopoulos, 2005]. GprMax2D gave simulated GPR monitoring 176 of the infiltration process. We then picked the maximum amplitude of the signal to get 177 the Two Way Travel (TWT) time of the wetting front reflection.

#### 3.3. Inversion Algorithm

#### 3.3.1. Convolution Algorithm

Our inversion algorithm was based on the comparison between the arrival times of the
wetting front reflection observed in the radargrams acquired during the water infiltration experiment and the arrival times of these reflections computed from the theoretical
water content profiles modeled by Hydrus-1D. If a suitable relationship between water
content and dielectric permittivity is known, water content profiles, obtained by the resolution of the Richards [Richards, 1931] equation (done by Hydrus1D in our case), can be
transformed to a 2D series of reflection coefficients:

$$R_{i,t} = \frac{\sqrt{\varepsilon_{i+1,t}} - \sqrt{\varepsilon_{i,t}}}{\sqrt{\varepsilon_{i+1,t}} + \sqrt{\varepsilon_{i,t}}},\tag{7}$$

where  $\sqrt{\varepsilon_{i,t}}$  and  $\sqrt{\varepsilon_{i+1,t}}$  are the relative dielectric permittivity at the infiltration time t 187 for two successive model cells centered at depth  $z_i$  and  $z_{i+1}$ . The effective depth where the reflection coefficient is calculated is  $z_R = \frac{z_i + z_{i+1}}{2}$ . Knowing the dielectric permittivity of 189 each layer of the profile, the electromagnetic wave velocity (Eq. 6) and travel time can 190 be computed. The travel time is used to interpolate reflection coefficients to a constant 191 sampling interval. We used this depth to time conversion to compute a Ricker signal in this time interval. The center frequency of the Ricker was set to 1000 MHz, central frequency of the GPR signal recorded on the field. We derived it twice with respect to time to simulate the transformation made by the emitter and the receiver in real antennae. We then performed the convolution between this pseudo-GPR signal and the reflectivity to obtain 197

$$O(t) = R(t) * \frac{\partial^2}{\partial t^2} I(t), \tag{8}$$

where O(t) is the output signal, R(t) is the reflectivity and I(t) is the input source of the antenna.

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Some remarks have to be made about the comparison between 1D-temporal convolution 200 and real electromagnetic signal. First of all, our inversion algorithm is based on the 201 assumption that soil can be represented as a stack of homogeneous layers. The assumption 202 of horizontal interfaces forces the reflection coefficient (equation (7)) to be expressed as 203 a normal incidence case. Secondly, we considered that the 2-D plane waves computed by 204 FDTD algorithm (modelling) and 3-D plane waves (experiments) could be treated as a 1-205 D temporal convolution. Third we neglect relaxation effects occurring when propagating 206 an electromagnetic wave in water saturated sand. 207

#### 3.3.2. Inversion Procedure

We used the TWT time obtained from the radargram (modelled or experimental) as 209 data to be fitted to derive the saturated hydraulic conductivity, assuming the other 4 210 van Genuchten parameters and initial water content were known. Using Hydrus-1D, we 211 generated 60 water content snapshots using the saturated hydraulic conductivity in the 212 range from 0.01 to 1 cm/min, with a step of 0.001 cm/min. For each value of  $K_s$ , we 213 calculated the TWT time using our convolution algorithm and we computed the Root 214 Mean Square Error (RMSE) between these times and the data as an objective function, 215 to be computed as function of saturated hydraulic conductivity. The  $K_s$  value which 216 corresponds to the minimum of the objective function was used as inverted value. 217

#### 4. Falling Head Infiltration Experiment

# 4.1. Numerical Example

#### $_{218}$ 4.1.1. Forward modelling

The set of hydrodynamical parameters used for this numerical example is presented in
Table 1. The permittivity profiles, resulting from water content conversions from Hydrus-

<sup>221</sup> 1D to permittivity and which were used as input of GprMax2D program [Giannopoulos, 2005] are presented in Figure 2-a. The simulated GPR monitoring of the infiltration process is shown in Figure 2-b. The horizontal axis is the number of traces simulated with GprMax2D, two traces being separated by 10 seconds, as permittivity profiles are.

The vertical axis is the TWT time of the EM wave coming back to the receiver.

On the profile presented in Figure 2-b, we denote one particular reflection, labeled A. Its 226 arrival time is increasing as the wetting front moves deeper. This reflection is interpreted 227 as coming from the wetting front. The reflections labeled A' and A'' are primary and secondary multiples of reflection A. The reflection labeled B is the wave traveling in air 229 directly between the two antennae. After the  $40^{th}$  trace, the 5-cm layer of water has been 230 infiltrated, and drainage is starting. As a consequence, the permittivity of the upper part 231 of the medium decreases and the velocity increases (Eq. 6). The TWT time of reflection A increases more slowly, creating a change of slope in the reflection time curve (Fig. 2-b). 233 In Figure 2-c, we display two curves: the TWT time of the maximum peak of reflection A (obtained from Figure 2-b) and the TWT time calculated by the convolution Algorithm. 235 The result of the convolution algorithm is in good agreement with the GprMax2D 236 modelling. 237

#### $_{8}$ 4.1.2. Inverse Modeling

We used the TWT time obtained from the radargram of Figure 2-b as data to be fitted to derive the saturated hydraulic conductivity, assuming the other 4 van Genuchten parameters and initial water content were known (see Table 1). The RMSE was minimized for  $K_s = 0.121 \ cm/min$ , which has to be compared with the value set as input, i.e.,

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 $K_s = 0.120 \ cm/min$ . This result confirms the ability of our algorithm to invert saturated hydraulic conductivity.

# 4.2. Field experiment

# 4.2.1. Experimental Data and its Analysis

The experiment took place in a quarry of Fontainebleau sand in Cernay-La-Ville (Yve-246 lines, France). The middle of the antennae was positioned 11 cm away from the cylinder 247 wall (X = 11 cm in Fig. 1). The 5-cm water layer was fully infiltrated after about 10 248 minutes, although in certain areas of the soil surface this time has been slightly shorter. The sand parameters measured by the hanging water column are given in Table 1 and initial volumetric water content is  $\theta_i = 0.09 \pm 0.01 \ cm^3/cm^3$ . The recorded GPR data are shown in Fig. 3. In this profile, we denote three particular reflections. The one 252 interpreted as coming from the infiltration front, labeled A, is visible during the first 30 minutes of the acquisition, with an arrival time varying from 2 ns down to 9 ns. The other reflections come from the cylinder and are interpreted in [Léger and Saintenoy, 2012]. We determined the arrival time of the A reflection peak and inverted the saturated hydraulic conductivity using the same algorithm as for the synthetic case. We obtained the mini-257 mum of the objective function for  $K_s = 0.120 \ cm/min$ . In parallel, we also carried out 258 disk infiltrometer experiments, using the multi-potential method [Ankeny et al., 1991; 259 Reynolds and Elrick, 1991. We obtained a value of the saturated hydraulic conductivity 260 of  $K_{Disk} = 0.108 \pm 0.01 \ cm/min$ . 261

#### 4.2.2. Uncertainty Analysis

We attempted to evaluate the uncertainty in the saturated hydraulic conductivity retrieved from GPR data fitting by using a modified Monte Carlo method. We qualified this

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method as "modified Monte Carlo" in the sense that it is nor the Tarantola method [Taran-265 tola, 1987] and neither the adaptive method proposed by the Guide to the expression of 266 uncertainty in measurement [BIPM et al., 2011] published by the Joint Committee for 267 Guides in Metrology (JCGM). We consider five major uncertainty sources, four from the 268 van Genuchten parameters,  $\alpha$ , n,  $\theta_r$ ,  $\theta_s$  and one from the initial water content  $\theta_i$ . We do as-269 sume that all uncertainties can by described by gaussian distribution probability function 270 centered on the value found by several water hanging column experiments with a stan-271 dard deviation of 5 % of this value. With this definition we obtained the following set of 272 a priori density function for experimental case:  $\mathcal{N}_{\alpha}(\alpha^{\mu} = 0.023 \ cm^{-1}, \ \alpha^{\sigma} = 0.001 \ cm^{-1}),$ 273  $\mathcal{N}_n(n^{\mu} = 6.7, n^{\sigma} = 0.3), \, \mathcal{N}_{\theta_r}(\theta_r^{\mu} = 0.062 \, cm^3/cm^3, \, \theta_r^{\sigma} = 0.001 cm^3/cm^3), \, \mathcal{N}_{\theta_s}(\theta_s^{\mu} = 0.001 cm^3/cm^3)$  $0.39 \ cm^3/cm^3, \ \theta_s^{\sigma} = 0.01 \ cm^3/cm^3), \ {\rm and} \ \mathcal{N}_{\theta_i}(\theta_i^{\mu} = 0.09 \ cm^3/cm^3, \ \theta_i^{\sigma} = 0.01 \ cm^3/cm^3),$ where the  $\mathcal{N}$  stands for the gaussian/normal probability density function and the  $\mu$  and  $\sigma$  represent the mean and standard deviation. We generate multiple sets of parameters 277 by sampling each gaussian distribution,  $\{\alpha^i, n^i, \theta^i_r, \theta^i_s, \theta^i_i\}$ , where the subscript "i" is the iteration number. For each set the value of  $K_s$  minimising the objective function 279 was computed by our inversion procedure presented above. We generated enough sets 280 of parameters such as the histogram of  $K_s$  values look like a gaussian function with a 281 stabilized standard deviation. We used this standard deviation as uncertainty on  $K_s$ . 282 We did not consider the uncertainties on radargram picking, because we evaluated it 283

Using our analysis, we found in the case of falling head infiltration that  $K_s$  was equal to  $0.12 \pm 0.01 \ cm/min$ . This narrow range of possible values is in agreement with disk infiltrometer value, and clearly shows the accuracy of our method.

has a very weak influence comparing to the other uncertainties considered.

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# 5. Constant Head Infiltration Experiment

# 5.1. Numerical Example

# 5.1.1. Forward Modelling

In this second case, a water layer of 5 cm above the ground was kept constant during 289 the entire experiment. Similarly as above, using the same van Genuchten parameters 290 as in the first synthetic example (Table 1), we modeled infiltration of water inside a 291 ring infiltrometer by applying a constant pressure head of 5 cm to the top node during 10 292 minutes. The permittivity profiles are presented in Fig. 4-a, with each curve plotted every 293 10 s as in the previous case. Fig. 4-b shows the radargram simulated with GprMax2D. 294 As can be seen, the reflection labeled A describing the position of the infiltration front, is 295 returning at increasing times, because infiltration is being constantly fed by the constant ponding depth, contrary to the previous falling head case. In Fig. 4-c, we computed 297 the TWT time of the wetting front using the convolution algorithm and picking the A reflection from the radargram in Fig. 4-b. 299

#### 5.1.2. Inverse Modelling

We inverted for the saturated hydraulic conductivity by minimizing the differences between the arrival times of the wetting front reflection obtained by the convolution algorithm and the arrival times picked from the simulated radargram in Fig. 4-b. The objective function was minimized for  $K_s = 0.119 \ cm/min$ , to be compared with the value used for simulating the data:  $K_s = 0.120 \ cm/min$ .

# 5.2. Field Experiment

The experiment took place in the same quarry of Fontainebleau sand as the previous experiment. The middle of the antennae was positioned in the middle of the ring (X = 30)

cm in Fig. 1). The GPR data are shown in Fig. 5 and were recorded during 80 minutes 308 (only a part of the radargram is presented). We used the van Genuchten parameters 309 determined in the laboratory using the hanging column experiments (Table 1) and we 310 measured on sand core samples an initial volumetric water content of  $\theta_i = 0.07 \pm 0.02$ . 311 In the profile presented in Fig. 5, the arrival time of reflection A ranges from 0 at the 312 beginning of the experiment to about 6 ns after 10 min. We picked the arrival time of the A 313 reflection peak and computed the objective function using the same procedure as described 314 before. We obtained the minimum of the objective function for  $K_s = 0.089 \ cm/min$ . 315 Again, this value has to be compared with the one obtained by the disk infiltrometer 316 experiment,  $K_{Disk} = 0.108 \pm 0.01$  cm/min. Using the same procedure as presented in 317 the earlier field example, we found a range of possible values for the saturated hydraulic 318 conductivity,  $K_s = 0.089 \pm 0.005 \ cm/min$ . Despite the fact that we are not in the same range as the disk infiltrometer method the discrepancy is very small and allows us to conclude on the good accuracy of our method.

# 6. Discussion

The results presented above indicate clearly that a commercial surface GPR can be used 322 as a tool for monitoring the wetting front. Although the use of surface-based GPR data 323 to estimate the parameters of unsaturated flow models is not new [Moysey, 2010], our 324 method gives accurate values of the saturated hydraulic conductivity with uncertainties 325 comparable or smaller than those obtained with disk infiltrometer measurements. A dis-326 tinct advantage of our approach is the simplicity of the algorithm and its rapidity to con-327 verge, which is very encouraging for more complicated models (stack of non-homogeneous 328 layers). 329

The discrepancy between saturated hydraulic conductivity determined by disk infiltrometry and that obtained with our GPR algorithm comes from different phenomena.
First of all, the van Genuchten parameters determined from the water hanging column
experiment are obtained with saturation coming from the bottom of the soil samples,
whereas in our case, the infiltration is a ponded one, thus coming from the top.

Despite the fact that we upgraded the single ring infiltrometer by the use of GPR to monitor the wetting front, we still suffer from the problem of entrapped-air, which causes reduction of saturated water content and hydraulic conductivity. This issue cannot be fixed with ponded infiltration. Disk infiltrometer measurement monitoring may cause less problems, working with negative matric potentials [Ankeny et al., 1991; Reynolds and Elrick, 1991].

During our modeling, we considered our soil as an homogeneous and isotropic one. Real soils exhibit heterogeneities, triggering preferential flows. Even in the case of our quarry of Fontainebleau sand, differences in packing and compaction could lead to creation of preferential flow paths.

One of the way to solve this issue could be to use a dual porosity model [Gerke and van Genuchten, 1993] and a Monte Carlo procedure to generate a high number of soil models with different parameters, as we did with the single porosity model in Hydrus-1D, and performed statistical analysis on the saturated hydraulic conductivity obtained.

An other source of error, already discussed above, comes from the assumption that a
350 3D infiltration monitored by 3D electromagnetic waves can be treated as a 1-D temporal
351 convolution. This limitation will be studied in future works, using Hydrus 2D/3D to
352 simulate 2D axisymmetrical infiltration and 2D infiltration.

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The results represent a promising step toward application of multi-parameters inversions. A first study in that direction was presented in Léger [Léger et al., 2013].

#### 7. Summary

This research investigated the use of on-ground surface GPR to monitor the wetting 355 front during infiltration inside a ring infiltrometer. We showed by modeling and exper-356 iments that a standard GPR device was able to monitor the displacement of the water 357 front in the soil. We tested in synthetic cases the ability of our algorithm to invert the 358 saturated hydraulic conductivity, knowing the other van Genuchten parameters and the 359 initial water content. Two infiltration experiments were performed, falling head infiltra-360 tion and constant head infiltration, in a quarry of Fontainebleau sand. The retrieved 361 saturated hydraulic conductivity was comparable to that obtained with disk infiltrometer 362 experiments. Uncertainty analysis accounting for all the van Genuchten parameters, was performed using a modified Monte Carlo method, and proved the robustness of our algorithm. Although results retrieved with GPR were in agreement with disk infiltrometry tests, we stress that further research is needed to improve our algorithm so as to determine the whole set of soil hydrodynamic parameters.

#### Notes

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- 1. Common MidPoints
  - 2. Wide-Angle Reflection- Refraction

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Falling Head Infiltration									
	$\theta_i$	$ \theta_r $	$\theta_s$	$\alpha$	n	$K_s$	Retrieved $K_s$		
				$(\mathrm{cm}^{-1})$		(cm/min)	(cm/min)		
Numerical	0.17	0.07	0.43	0.019	8.67	0.120	0.121		
Field	$0.09 \pm 0.01$	$0.062 \pm 0.003$	$0.39 \pm 0.01$	$0.023\ \pm0.001$	$ 6.7 \pm 0.3 $	$0.108 \pm 0.01^*$	$0.120 \pm 0.013$		
Constant Head Infiltration									
	$ \theta_i $	$  heta_r $	$\theta_s$	$\alpha$	n	$K_s$	Retrieved $K_s$		
				$(\mathrm{cm}^{-1})$		(cm/min)	(cm/min)		
Numerical	0.17	0.07	0.43	0.019	8.67	0.120	0.119		
Field	$0.07 \pm 0.02$	$0.062 \pm 0.003$	$0.39 \pm 0.01$	$0.023 \pm 0.001$	$6.7 \pm 0.3$	$0.108 \pm 0.01^*$	$0.089 \pm 0.009$		

**Table 1.** Hydrodynamic parameters for the numerical and field experiment. The \* indicates values of  $K_s$  measured from disk infiltrometer experiments.

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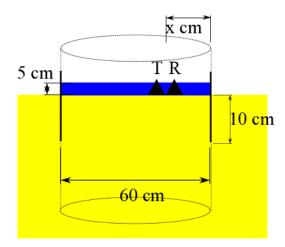


Figure 1. Experimental set up at its initial state.

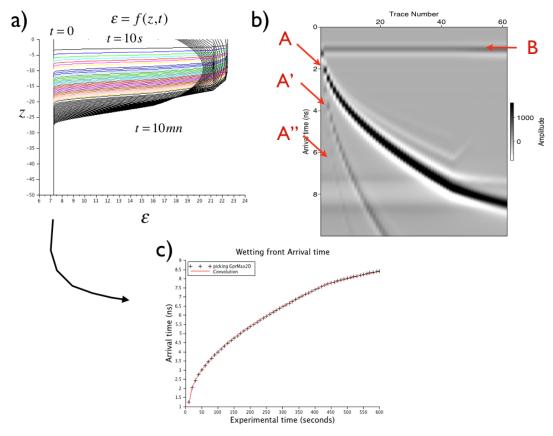
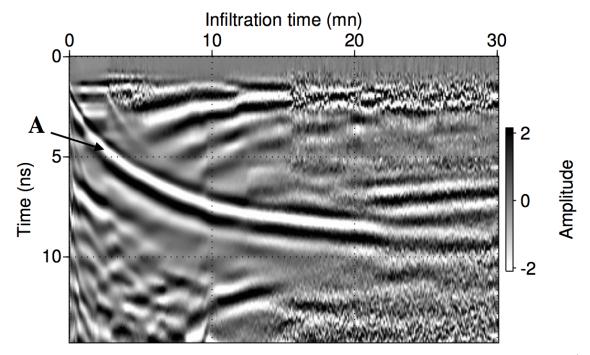


Figure 2. Falling head infiltration from a 5-cm thick water layer. a) Permittivity profiles: each curve is plotted every 10 s. b) Radargram simulated with GprMax2D; reflection A is coming from the wetting front, B is the direct wave, A' and A" are multiples of reflection A. c) TWT time computed by the convolution algorithm from the permittivity profiles (plain red line) and TWT time obtained by picking of A peak in fig b).



**Figure 3.** Experimental GPR data acquired during the falling head infiltration (using a 5-cm initial water layer). Reflection A is the reflection coming from the wetting front

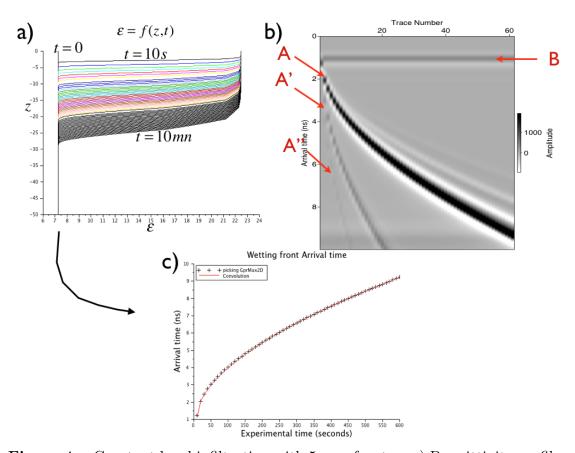
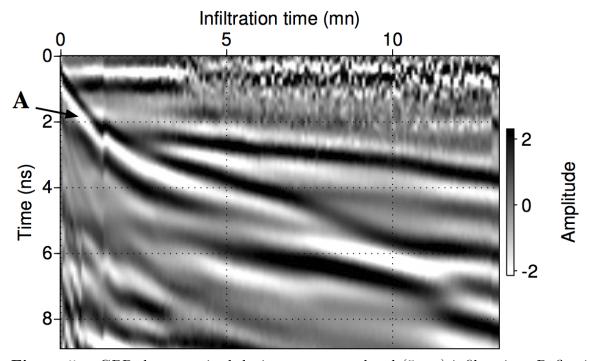


Figure 4. Constant head infiltration with 5 cm of water. a) Permittivity profiles, each curve is plotted every 10 s. b) Radargram simulated with GprMax2D, reflection A is the wetting front, B is the direct wave, A' and A" are multiples. c) Two Way Travel Time computed with our convolution algorithm from the simulated permittivity profiles.



**Figure 5.** GPR data acquired during a constant head (5 cm) infiltration. Reflection A is the reflection coming from the wetting front.