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¹ Saturated hydraulic conductivity determined by on ² ground mono-offset Ground-Penetrating Radar inside ³ a single ring infiltrometer

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Abstract. In this study we show how to use GPR data acquired along the infiltration of water inside a single ring infiltrometer to inverse the sat- urated hydraulic conductivity. We used Hydrus-1D to simulate the water in- filtration. We generated water content profiles at each time step of infiltra- tion, based on a particular value of the saturated hydraulic conductivity, know- ing the other van Genuchten parameters. Water content profiles were con- verted to dielectric permittivity profiles using the Complex Refractive In-¹¹ dex Method relation. We then used the GprMax suite of programs to gen- erate radargrams and to follow the wetting front using arrival time of elec- tromagnetic waves recorded by a Ground-Penetrating Radar (GPR). The- oretically, the 1D time convolution between reflectivity and GPR signal at any infiltration time step is related to the peak of the reflected amplitude recorded in the corresponding trace in the radargram. We used this relation- ship to invert the saturated hydraulic conductivity for constant and falling head infiltrations. We present our method on synthetic examples and on two experiments carried out on sand soil. We further discuss on the uncertain- ties on the retrieved saturated hydraulic conductivity computed by our al-gorithm from the van Genuchten parameters.

1. Introduction

²² Soil hydraulic properties, represented by the soil water retention $\theta(h)$ and hydraulic 23 conductivity $K(h)$ functions, dictate water flow in the vadose zone, as well as partitioning ²⁴ between infiltration and runoff. Their evaluation has important implications for mod-²⁵ eling available water resources and for flood forecasting. It is also crucial in evaluating ²⁶ the dynamics of chemical pollutants in soil and in assessing the potential of groundwater ²⁷ pollution.

²⁸ Soil hydraulic functions can be described by several mathematical expression [Kosugi 29 et al., 2002], among them the van Genuchten function [van Genuchten, 1980]. The deter-³⁰ mination of the parameters defining the van Genuchten soil water retention function [van ³¹ Genuchten, 1980 is usually done using laboratory experiments, such as the water hanging $_{32}$ column [Dane and Hopmans, 2002].

³³ The hydraulic conductivity function can be estimated either in the laboratory, or in ³⁴ situ using infiltration tests. Among the large number of existing infiltration tests [*Angulo*-³⁵ Jaramillo et al., 2000, the single [*Muntz et al.*, 1905] or double ring infiltrometers [*Boivin*] ³⁶ et al., 1987] provide the field saturated hydraulic conductivity by applying a positive 37 water pressure on the soil surface, while the disk infiltrometer [*Perroux and White*, 1988; ³⁸ Clothier and White, 1981 allows to reconstruct the hydraulic conductivity curve, by ³⁹ applying different water pressures smaller than or equal to zero. For infiltration tests, the ⁴⁰ volume of infiltrated water versus time is fitted to infer the soil hydraulic conductivity at or ⁴¹ close to saturation. These tests are time-consuming and difficult to apply to landscape-⁴² scale forecasting of infiltration. Furthermore, their analysis involve various simplifying

X - 4 LÉGER, SAINTENOY AND COQUET: INVERTING SATURATED HYDRAULIC CONDUCTIVITY

 assumptions, partly due to the ignorance of the shape of the infiltration bulb. This lack of knowledge on the form of the infiltration bulb has to be filled to get accurate informations 45 on the soil water retention $\theta(h)$ function and consequently on hydraulic conductivity $K(h)$ function. This can be done by water content sensing.

⁴⁷ Vereecken *[Vereecken et al., 2008]* and Evett and Parkin *[Evett and Parkin, 2005]* give ⁴⁸ a state of the art on the different techniques available for soil moisture measurements. ⁴⁹ Among the large panel presented, geophysical methods take an important part, mainly ⁵⁰ because they are contact free and/or easy to use. The most commonly used hydro- $51 \text{ geophysical methods are electrical resistivity measurements}$ [*Goyal et al.*, 2006; *Zhou et al.*, \approx 2001] and electromagnetic methods [*Sheets and Hendrickx*, 1995; Akbar et al., 2005]. This ⁵³ paper focuses on the use of Ground-Penetrating Radar (GPR) as a tool for monitoring water infiltration in soil.

⁵⁵ For few decades GPR has been known as an accurate method to highlight water vari-⁵⁶ ation in soils [Huisman et al., 2003; Annan, 2005]. Different techniques are available in ₅₇ the literature for monitoring water content in soils using GPR. Tomography imaging be- $\frac{1}{58}$ tween boreholes during an infiltration has been done by Binley [Binley et al., 2001] and 59 Kowalsky $[Kowalsky \;et \; al., 2005]$ among others. Many advances were done during the last ⁶⁰ years on Off-Ground GPR using full waveform inversion, for instance to invert soil hy- $_{61}$ draulic properties (Lambot [Lambot et al., 2006, 2009] and Jadoon [Jadoon et al., 2012]). ϵ_2 Grote [Grote et al., 2002] and Lunt [Lunt et al., 2005] used two-way travel time variations ⁶³ from a reflector at a known depth to monitor water content variation with time. Finally, μ multi-offset GPR survey techniques, i.e. CMP¹ or WARR², were carried out during infil-

⁶⁵ tration processes in the works of Greaves [*Greaves et al.*, 1996] or Mangel [*Mangel et al.*, ⁶⁶ 2012].

 σ The work presented here is based on mono-offset monitoring of infiltration with on-⁶⁸ ground surface GPR as related by Haarder [*Haarder et al.*, 2011], Moysey [*Moysey*, 2010], $\omega_{\rm e}$ Lai [Lai et al., 2012], Dagenbach [Dagenbach et al., 2013] and Saintenoy [Saintenoy et al., $\frac{1}{70}$ 2008]. Haarder [Haarder et al., 2011] used a constant offset on-ground GPR coupled with π dye tracing to exhibit preferential flows. They found that a GPR was able to map deep η infiltration comparing to dye tracer, but they did not manage to resolve the infiltration η_3 patterns (by-pass flow, fingering...). Moysey [*Moysey*, 2010] studied the infiltration inside 74 a sand box from the surface with on-ground GPR. He used the reflection from the wetting ⁷⁵ front as well as from the ground wave and the bottom of the box, to monitor the water ⁷⁶ content. He also modelled his experiment and estimated the van Genuchten parameters π using semblance analysis. As Léger [Léger et al., 2013], he found that the most poorly ⁷⁸ constrained parameter was n. Lai [Lai et al., 2012] used a joint time frequency analysis ⁷⁹ coupled with grayscale imaging to measure infiltration and drainage in controlled con-⁸⁰ ditions in laboratory. They were able to follow the peak frequency of the GPR wavelet ⁸¹ associated with the wetting front using time frequency analysis and then determined the ⁸² rate of water infiltration in unsaturated zone. Saintenoy [Saintenoy et al., 2008] mon-⁸³ itored the wetting bulb during an infiltration from a Porchet infiltrometer. They were ⁸⁴ able to identify the dimension of the bulb with time and good agreement was found with ⁸⁵ modelling.

⁸⁶ On the continuity of those studies, we present a method for monitoring the wetting ⁸⁷ front during infiltration using on-ground GPR with fixed offset inside a ring infiltrometer.

 The objectives of this paper were i) to check if the proposed method is accurate enough to monitor wetting front during infiltration with different boundary conditions, ii) to invert ⁹⁰ saturated hydraulic conductivity using the model of Mualem-van Genuchten [*Mualem*, 1976; van Genuchten, 1980] , and iii) to analyze the uncertainties using a simplified MC uncertainty analysis. The method has been tested on synthetic examples and on two field data sets.

2. Background

2.1. Unsaturated Flow Equation

⁹⁴ In this study we consider one-dimensionnal vertical water flow in a soil, described by ⁹⁵ the one-dimensional Richard's equation [Richards, 1931]. Its expression in term of water ⁹⁶ content is

$$
\frac{\partial \theta}{\partial t} = \frac{\partial K(\theta)}{\partial z} + \frac{\partial}{\partial z} \left[D(\theta) \frac{\partial \theta}{\partial z} \right],\tag{1}
$$

⁹⁷ where $K(\theta)$ is the hydraulic conductivity as a function of water content, and $D(\theta)$ is water diffusivity (Childs and Georges-Collis [Childs and Collis-George, 1950]), expressed ⁹⁹ in terms of water content as $D(\theta) = K(\theta) \frac{\partial h}{\partial \theta}$.

2.2. Hydraulic Properties Functions

 Several mathematical functions exist to model the hydraulic properties of porous me- dia [Kosugi et al., 2002]. We chose the van Genuchten model [van Genuchten, 1980] ¹⁰² with the relation of Mualem [*Mualem*, 1976], giving the following expression for the water retention curve:

$$
\theta(h) = \theta_r + (\theta_s - \theta_r)(1 + (\alpha h)^n)^{\frac{1}{n} - 1},\tag{2}
$$

104 where θ_s is the saturated water content, θ_r , the residual water content, and α and n , two ¹⁰⁵ fitting parameters which are respectively linked to the matric potential and the slope of ¹⁰⁶ the water retention curve at the inflexion point. The hydraulic conductivity function is 107 described by

$$
K(\theta) = K_s \Theta^{\lambda} \left[1 - \left[1 - \Theta^{\frac{n}{n-1}} \right]^{\frac{n}{n-1}} \right]^2,
$$
\n⁽³⁾

¹⁰⁸ with K_s the saturated hydraulic conductivity, $\Theta = \frac{\theta - \theta_r}{\theta_s - \theta_r}$ the effective saturation and λ a 109 factor that accounts for pore tortuosity. The λ parameter has an influence on the shape ¹¹⁰ of the hydraulic conductivity function. However in this study we concentrated on the μ ¹¹¹ inversion of only one parameter, the saturated hydraulic conductivity. We fixed λ equal $_{112}$ to 0.5 as reported in [*Mualem*, 1976].

2.3. Petrophysical Relationships

 Several empirical and conceptual relationships exist to convert soil dielectric permit- tivity to volumetric water content. Using the fact that the experiments presented here have been made in a quarry of Fontainebleau sand, considered as pure silica, we used the ¹¹⁶ CRIM relation *[Birchak et al.,* 1974; *Roth et al.,* 1990], which relates the relative dielectric 117 permittivity of bulk media, ε_b , to the volumetric summation of each components of it. Thus for a tri-phasic medium comprising water, air and silicium, we obtain

$$
\sqrt{\varepsilon_b} = \theta \sqrt{\varepsilon_w} + (1 - \phi) \sqrt{\varepsilon_s} + (\phi - \theta), \tag{4}
$$

119 where $\varepsilon_w = 80.1$, $\varepsilon_s = 2.5$ are respectively the relative dielectric permittivity of water 120 and silica, ϕ the porosity and θ the volumetric water content.

2.4. Dielectric Permittivity Versus Electromagnetic Wave Velocity

 Surface GPR consists in a transmitting antenna, being a dipole, positioned on the surface, that emits short pulses of spherical electromagnetic (EM) wave in response to an excitation current source, and a receiving antenna, also located at the surface, which converts the incoming EM fields to an electrical signal source to be treated. Following the ¹²⁵ works of Annan [*Annan*, 1999], the velocity of electromagnetic waves is

$$
v = \frac{c}{\sqrt{\varepsilon'\mu_r \frac{1+\sqrt{1+tan^2\delta}}{2}}},\tag{5}
$$

 $_{126}$ where δ is the loss factor as a function of the dielectric permittivity, frequency and electrical conductivity, ε' is the real part of the relative dielectric permittivity, μ_r the relative ¹²⁸ magnetic permeability and c is the velocity of EM waves in air equal to 0.3 m/ns . Con-¹²⁹ sidering the case of non magnetic soil with low conductivity, in the range of 10 MHz ¹³⁰ to 1 GHz, the real part dominates the imaginary part of the dielectric permittivity and $_{131}$ neglecting Debye [*Debye*, 1929] effect, equation (5) reduces to:

$$
v = \frac{c}{\sqrt{\varepsilon'}}.\tag{6}
$$

¹³² We used this equation to compute the travelling time of an EM wave through a layer of ¹³³ soil of known thickness with a given dielectric permittivity.

2.5. Electromagnetic Modelling

 Numerous techniques are available for simulating GPR data, e.g. ray-based methods (e.g. Cai and McMechan [Cai and McMechan, 1995] or Sethian and Popovici [Sethian and Popovici, 1999]), time-domain finite-difference full-waveform methods (e.g. Kunz and Luebbers [Kunz and Luebbers, 1996] or Kowalsky [Kowalsky et al., 2001]), or finite differences time domain (FDTD) (e.g. Irving and Knight [Irving and Knight, 2006]). We

 used the GprMax 2D codes of Giannopoulos [Giannopoulos, 2005], which uses FDTD modelling to solve the maxwell equations in 2 dimensions.

3. Materials and Methods

3.1. Experimental Set-up

¹⁴¹ We studied infiltration of a 5-cm thick water layer inside of a single ring infiltrometer $_{142}$ in a sandy soil. The scheme of the apparatus is presented in Figure 1. The single ring infiltrometer was a 1-mm thick aluminum cylinder with a 60-cm diameter, approximately 20-cm high, buried in the soil to a depth of 10 cm. GPR antennae (namely the transmitter $_{145}$ T and the receiver R) were set up at a variable distance from the edge of the cylinder, $_{146}$ noted X, in Figure 1. In all our field experiments, we used a Mala RAMAC system with antennae centered on 1600 MHz, shielded at the top. The inner part of the cylinder was covered with a plastic waterproof sheet. This allowed us to fill the cylinder with water and create an initial 5-cm thick water layer, while preventing infiltration into the sand before starting data acquisition. The beginning of the acquisition was launched by pulling away the plastic sheet to trigger water infiltration. The GPR system was set to acquire a trace every 10 s. With this apparatus, we performed two types of infiltration: i) a falling head infiltration consisting of pulling away the plastic sheet and leaving water to infiltrate into the sand freely with no additional refill, and ii) a constant head infiltration, when water was continuously added to the ring to maintain a 5-cm thick water layer during the infiltration experiment. In the following examples, we will show that GPR data acquired every 10 s during the infiltration experiment can be used to estimate the ¹⁵⁸ saturated soil hydraulic conductivity, K_s . In all GPR data presented below, we subtracted the average trace and applied an Automatic Gain Control (AGC) to the data in order

160 to make them clearer. The van Genuchten parameters, α , n , θ_r , θ_i of the sand have ¹⁶¹ been determined in laboratory by several classical hanging water column experiments. $\frac{162}{162}$ We assumed arbitrarily a 5 % uncertainty for all the measured parameters. The sand was ¹⁶³ considered homogeneous. Its initial water content, θ_i , and porosity, ϕ , of the soil were ¹⁶⁴ determined using gravimetric measurements on field samples.

3.2. Modelling

 Infiltration experiments were simulated by solving Richards equation (Eq. (1)) using Hydrus-1D. The soil profile was 50 cm deep, assumed to be homogeneous, and divided into ¹⁶⁷ 1001 layers. We used either an atmospheric boundary condition (BC) with no rain and no evaporation at the soil surface, for the falling head infiltration, or a constant pressure head of 5 cm to the top node, for the constant head infiltration, and for both case free drainage BC at the bottom. To simulate the 5-cm layer of water, the initial condition was set to a 5 cm pressure head in the top node. We simulated the first 10 minutes of $_{172}$ the experiment with a time step of 10 s, i.e., with 60 water content snapshots. Using the CRIM relation (Eq. 4), each water content snapshot was converted to permittivity profiles (made of 1001 points), considering a three-phase media: sand (considered as pure silica), water, and air. Each one of these permittivity profiles were the input for the GprMax2D program [Giannopoulos, 2005]. GprMax2D gave simulated GPR monitoring $_{177}$ of the infiltration process. We then picked the maximum amplitude of the signal to get the Two Way Travel (TWT) time of the wetting front reflection.

3.3. Inversion Algorithm

179 3.3.1. Convolution Algorithm

 Our inversion algorithm was based on the comparison between the arrival times of the wetting front reflection observed in the radargrams acquired during the water infiltra- tion experiment and the arrival times of these reflections computed from the theoretical water content profiles modeled by Hydrus-1D. If a suitable relationship between water content and dielectric permittivity is known, water content profiles, obtained by the reso-¹⁸⁵ lution of the Richards [*Richards*, 1931] equation (done by Hydrus1D in our case), can be transformed to a 2D series of reflection coefficients:

$$
R_{i,t} = \frac{\sqrt{\varepsilon_{i+1,t}} - \sqrt{\varepsilon_{i,t}}}{\sqrt{\varepsilon_{i+1,t}} + \sqrt{\varepsilon_{i,t}}},\tag{7}
$$

¹⁸⁷ where $\sqrt{\varepsilon_{i,t}}$ and $\sqrt{\varepsilon_{i+1,t}}$ are the relative dielectric permittivity at the infiltration time t ¹⁸⁸ for two successive model cells centered at depth z_i and z_{i+1} . The effective depth where the reflection coefficient is calculated is $z_R = \frac{z_i + z_{i+1}}{2}$ ¹⁸⁹ reflection coefficient is calculated is $z_R = \frac{z_i + z_{i+1}}{2}$. Knowing the dielectric permittivity of ¹⁹⁰ each layer of the profile, the electromagnetic wave velocity (Eq. 6) and travel time can ¹⁹¹ be computed. The travel time is used to interpolate reflection coefficients to a constant ¹⁹² sampling interval. We used this depth to time conversion to compute a Ricker signal ¹⁹³ in this time interval. The center frequency of the Ricker was set to 1000 MHz, central ¹⁹⁴ frequency of the GPR signal recorded on the field. We derived it twice with respect to ¹⁹⁵ time to simulate the transformation made by the emitter and the receiver in real antennae. ¹⁹⁶ We then performed the convolution between this pseudo-GPR signal and the reflectivity ¹⁹⁷ to obtain

$$
O(t) = R(t) * \frac{\partial^2}{\partial t^2} I(t),
$$
\n(8)

¹⁹⁸ where $O(t)$ is the output signal, $R(t)$ is the reflectivity and $I(t)$ is the input source of the ¹⁹⁹ antenna.

X - 12 LEGER, SAINTENOY AND COQUET: INVERTING SATURATED HYDRAULIC CONDUCTIVITY ´

²⁰⁰ Some remarks have to be made about the comparison between 1D-temporal convolution and real electromagnetic signal. First of all, our inversion algorithm is based on the assumption that soil can be represented as a stack of homogeneous layers. The assumption of horizontal interfaces forces the reflection coefficient (equation (7)) to be expressed as a normal incidence case. Secondly, we considered that the 2-D plane waves computed by FDTD algorithm (modelling) and 3-D plane waves (experiments) could be treated as a 1- D temporal convolution. Third we neglect relaxation effects occurring when propagating an electromagnetic wave in water saturated sand.

3.3.2. Inversion Procedure

 We used the TWT time obtained from the radargram (modelled or experimental) as data to be fitted to derive the saturated hydraulic conductivity, assuming the other 4 van Genuchten parameters and initial water content were known. Using Hydrus-1D, we generated 60 water content snapshots using the saturated hydraulic conductivity in the ²¹³ range from 0.01 to 1 cm/min, with a step of 0.001 cm/min. For each value of K_s , we calculated the TWT time using our convolution algorithm and we computed the Root ²¹⁵ Mean Square Error (RMSE) between these times and the data as an objective function, ²¹⁶ to be computed as function of saturated hydraulic conductivity. The K_s value which corresponds to the minimum of the objective function was used as inverted value.

4. Falling Head Infiltration Experiment

4.1. Numerical Example

218 4.1.1. Forward modelling

 The set of hydrodynamical parameters used for this numerical example is presented in Table 1. The permittivity profiles, resulting from water content conversions from Hydrus-

²²¹ 1D to permittivity and which were used as input of GprMax2D program $[Giannopoulos,$ 2005] are presented in Figure 2-a. The simulated GPR monitoring of the infiltration process is shown in Figure 2-b. The horizontal axis is the number of traces simulated with GprMax2D, two traces being separated by 10 seconds, as permittivity profiles are. The vertical axis is the TWT time of the EM wave coming back to the receiver.

 On the profile presented in Figure 2-b, we denote one particular reflection, labeled A. Its arrival time is increasing as the wetting front moves deeper. This reflection is interpreted 228 as coming from the wetting front. The reflections labeled A' and A'' are primary and secondary multiples of reflection A. The reflection labeled B is the wave traveling in air ²³⁰ directly between the two antennae. After the 40^{th} trace, the 5-cm layer of water has been infiltrated, and drainage is starting. As a consequence, the permittivity of the upper part of the medium decreases and the velocity increases (Eq. 6). The TWT time of reflection ²³³ A increases more slowly, creating a change of slope in the reflection time curve (Fig. 2-b). In Figure 2-c, we display two curves: the TWT time of the maximum peak of reflection A (obtained from Figure 2-b) and the TWT time calculated by the convolution Algorithm. The result of the convolution algorithm is in good agreement with the GprMax2D 237 modelling.

4.1.2. Inverse Modeling

 We used the TWT time obtained from the radargram of Figure 2-b as data to be fitted to derive the saturated hydraulic conductivity, assuming the other 4 van Genuchten parameters and initial water content were known (see Table 1). The RMSE was minimized ²⁴² for $K_s = 0.121$ cm/min, which has to be compared with the value set as input, i.e.,

²⁴³ $K_s = 0.120$ cm/min. This result confirms the ability of our algorithm to invert saturated ²⁴⁴ hydraulic conductivity.

4.2. Field experiment

245 4.2.1. Experimental Data and its Analysis

²⁴⁶ The experiment took place in a quarry of Fontainebleau sand in Cernay-La-Ville (Yve-²⁴⁷ lines, France). The middle of the antennae was positioned 11 cm away from the cylinder ²⁴⁸ wall $(X = 11$ cm in Fig. 1). The 5-cm water layer was fully infiltrated after about 10 ²⁴⁹ minutes, although in certain areas of the soil surface this time has been slightly shorter. ²⁵⁰ The sand parameters measured by the hanging water column are given in Table 1 and ²⁵¹ initial volumetric water content is $\theta_i = 0.09 \pm 0.01$ cm³/cm³. The recorded GPR data ²⁵² are shown in Fig. 3. In this profile, we denote three particular reflections. The one ²⁵³ interpreted as coming from the infiltration front, labeled A, is visible during the first 30 ²⁵⁴ minutes of the acquisition, with an arrival time varying from 2 ns down to 9 ns. The other ²⁵⁵ reflections come from the cylinder and are interpreted in [*Léger and Saintenoy*, 2012]. We ²⁵⁶ determined the arrival time of the A reflection peak and inverted the saturated hydraulic ²⁵⁷ conductivity using the same algorithm as for the synthetic case. We obtained the mini-²⁵⁸ mum of the objective function for $K_s = 0.120$ cm/min. In parallel, we also carried out ²⁵⁹ disk infiltrometer experiments, using the multi-potential method [*Ankeny et al.*, 1991; ₂₆₀ Reynolds and Elrick, 1991. We obtained a value of the saturated hydraulic conductivity ²⁶¹ of $K_{Disk} = 0.108 \pm 0.01$ cm/min.

²⁶² 4.2.2. Uncertainty Analysis

²⁶³ We attempted to evaluate the uncertainty in the saturated hydraulic conductivity re-²⁶⁴ trieved from GPR data fitting by using a modified Monte Carlo method. We qualified this

²⁶⁵ method as "modified Monte Carlo" in the sense that it is nor the Tarantola method \lceil Taran-²⁶⁶ tola, 1987] and neither the adaptive method proposed by the Guide to the expression of ²⁶⁷ uncertainty in measurement [BIPM et al., 2011] published by the Joint Committee for ²⁶⁸ Guides in Metrology (JCGM). We consider five major uncertainty sources, four from the ²⁶⁹ van Genuchten parameters, α , n , θ_r , θ_s and one from the initial water content θ_i . We do as-₂₇₀ sume that all uncertainties can by described by gaussian distribution probability function ₂₇₁ centered on the value found by several water hanging column experiments with a stan- $_{272}$ dard deviation of 5 % of this value. With this definition we obtained the following set of ²⁷³ a priori density function for experimental case: $\mathcal{N}_{\alpha}(\alpha^{\mu} = 0.023 \text{ cm}^{-1}, \alpha^{\sigma} = 0.001 \text{ cm}^{-1}),$ $N_n(n^\mu~=~6.7,~~n^\sigma~=~0.3),~\, \mathcal{N}_{\theta_r}(\theta_r^\mu~=~0.062~~cm^3/cm^3,~~\theta_r^\sigma~=~0.001cm^3/cm^3),~\, \mathcal{N}_{\theta_s}(\theta_s^\mu~=~0.001cm^3)$ \int_{275} 0.39 cm^3/cm^3 , $\theta_s^{\sigma} = 0.01$ cm^3/cm^3), and $\mathcal{N}_{\theta_i}(\theta_i^{\mu} = 0.09$ cm^3/cm^3 , $\theta_i^{\sigma} = 0.01$ cm^3/cm^3), ²⁷⁶ where the N stands for the gaussian/normal probability density function and the μ and 277 σ represent the mean and standard deviation. We generate multiple sets of parameters ²⁷⁸ by sampling each gaussian distribution, $\{\alpha^i, n^i, \theta^i_r, \theta^i_s, \theta^i_i\}$, where the subscript "i" is ₂₇₉ the iteration number. For each set the value of K_s minimising the objective function ²⁸⁰ was computed by our inversion procedure presented above. We generated enough sets ²⁸¹ of parameters such as the histogram of K_s values look like a gaussian function with a ²⁸² stabilized standard deviation. We used this standard deviation as uncertainty on K_s .

²⁸³ We did not consider the uncertainties on radargram picking, because we evaluated it ²⁸⁴ has a very weak influence comparing to the other uncertainties considered.

²⁸⁵ Using our analysis, we found in the case of falling head infiltration that K_s was equal ²⁸⁶ to 0.12 ± 0.01 cm/min. This narrow range of possible values is in agreement with disk ²⁸⁷ infiltrometer value, and clearly shows the accuracy of our method.

5. Constant Head Infiltration Experiment

5.1. Numerical Example

²⁸⁸ 5.1.1. Forward Modelling

²⁸⁹ In this second case, a water layer of 5 cm above the ground was kept constant during ²⁹⁰ the entire experiment. Similarly as above, using the same van Genuchten parameters ²⁹¹ as in the first synthetic example (Table 1), we modeled infiltration of water inside a $_{292}$ ring infiltrometer by applying a constant pressure head of 5 cm to the top node during 10 ²⁹³ minutes. The permittivity profiles are presented in Fig. 4-a, with each curve plotted every $_{294}$ 10 s as in the previous case. Fig. 4-b shows the radargram simulated with GprMax2D. ²⁹⁵ As can be seen, the reflection labeled A describing the position of the infiltration front, is ²⁹⁶ returning at increasing times, because infiltration is being constantly fed by the constant ₂₉₇ ponding depth, contrary to the previous falling head case. In Fig. 4-c, we computed ²⁹⁸ the TWT time of the wetting front using the convolution algorithm and picking the A ²⁹⁹ reflection from the radargram in Fig. 4-b.

300 5.1.2. Inverse Modelling

³⁰¹ We inverted for the saturated hydraulic conductivity by minimizing the differences ³⁰² between the arrival times of the wetting front reflection obtained by the convolution ³⁰³ algorithm and the arrival times picked from the simulated radargram in Fig. 4-b. The ³⁰⁴ objective function was minimized for $K_s = 0.119$ cm/min, to be compared with the value 305 used for simulating the data: $K_s = 0.120 \text{ cm/min}$.

5.2. Field Experiment

³⁰⁶ The experiment took place in the same quarry of Fontainebleau sand as the previous 307 experiment. The middle of the antennae was positioned in the middle of the ring ($X = 30$)

³⁰⁸ cm in Fig. 1). The GPR data are shown in Fig. 5 and were recorded during 80 minutes ³⁰⁹ (only a part of the radargram is presented). We used the van Genuchten parameters ³¹⁰ determined in the laboratory using the hanging column experiments (Table 1) and we 311 measured on sand core samples an initial volumetric water content of $\theta_i = 0.07 \pm 0.02$.

 $_{312}$ In the profile presented in Fig. 5, the arrival time of reflection A ranges from 0 at the $_{313}$ beginning of the experiment to about 6 ns after 10 min. We picked the arrival time of the A ³¹⁴ reflection peak and computed the objective function using the same procedure as described 315 before. We obtained the minimum of the objective function for $K_s = 0.089$ cm/min. ³¹⁶ Again, this value has to be compared with the one obtained by the disk infiltrometer 317 experiment, $K_{Disk} = 0.108 \pm 0.01$ cm/min. Using the same procedure as presented in ³¹⁸ the earlier field example, we found a range of possible values for the saturated hydraulic 319 conductivity, $K_s = 0.089 \pm 0.005$ cm/min. Despite the fact that we are not in the same ³²⁰ range as the disk infiltrometer method the discrepancy is very small and allows us to 321 conclude on the good accuracy of our method.

6. Discussion

³²² The results presented above indicate clearly that a commercial surface GPR can be used ³²³ as a tool for monitoring the wetting front. Although the use of surface-based GPR data $_{324}$ to estimate the parameters of unsaturated flow models is not new [Moysey, 2010], our ³²⁵ method gives accurate values of the saturated hydraulic conductivity with uncertainties ³²⁶ comparable or smaller than those obtained with disk infiltrometer measurements. A dis-³²⁷ tinct advantage of our approach is the simplicity of the algorithm and its rapidity to con-³²⁸ verge, which is very encouraging for more complicated models (stack of non-homogeneous 329 layers).

X - 18 LÉGER, SAINTENOY AND COQUET: INVERTING SATURATED HYDRAULIC CONDUCTIVITY

₃₃₀ The discrepancy between saturated hydraulic conductivity determined by disk infil-331 trometry and that obtained with our GPR algorithm comes from different phenomena. ³³² First of all, the van Genuchten parameters determined from the water hanging column ³³³ experiment are obtained with saturation coming from the bottom of the soil samples, ³³⁴ whereas in our case, the infiltration is a ponded one, thus coming from the top.

³³⁵ Despite the fact that we upgraded the single ring infiltrometer by the use of GPR to ³³⁶ monitor the wetting front, we still suffer from the problem of entrapped-air, which causes ³³⁷ reduction of saturated water content and hydraulic conductivity. This issue cannot be ³³⁸ fixed with ponded infiltration. Disk infiltrometer measurement monitoring may cause less 339 problems, working with negative matric potentials $[Ankeny$ et al., 1991; Reynolds and 340 Elrick, 1991].

³⁴¹ During our modeling, we considered our soil as an homogeneous and isotropic one. Real soils exhibit heterogeneities, triggering preferential flows. Even in the case of our quarry of Fontainebleau sand, differences in packing and compaction could lead to creation of preferential flow paths.

³⁴⁵ One of the way to solve this issue could be to use a dual porosity model [*Gerke and van* ³⁴⁶ Genuchten, 1993 and a Monte Carlo procedure to generate a high number of soil models ³⁴⁷ with different parameters, as we did with the single porosity model in Hydrus-1D, and ³⁴⁸ performed statistical analysis on the saturated hydraulic conductivity obtained.

³⁴⁹ An other source of error, already discussed above, comes from the assumption that a 3D infiltration monitored by 3D electromagnetic waves can be treated as a 1-D temporal convolution. This limitation will be studied in future works, using Hydrus 2D/3D to simulate 2D axisymmetrical infiltration and 2D infiltration.

³⁵³ The results represent a promising step toward application of multi-parameters inver- 354 sions. A first study in that direction was presented in Léger [Léger et al., 2013].

7. Summary

 This research investigated the use of on-ground surface GPR to monitor the wetting front during infiltration inside a ring infiltrometer. We showed by modeling and exper-³⁵⁷ iments that a standard GPR device was able to monitor the displacement of the water front in the soil. We tested in synthetic cases the ability of our algorithm to invert the saturated hydraulic conductivity, knowing the other van Genuchten parameters and the ³⁶⁰ initial water content. Two infiltration experiments were performed, falling head infiltra- tion and constant head infiltration, in a quarry of Fontainebleau sand. The retrieved ³⁶² saturated hydraulic conductivity was comparable to that obtained with disk infiltrometer experiments. Uncertainty analysis accounting for all the van Genuchten parameters, was performed using a modified Monte Carlo method, and proved the robustness of our al- gorithm. Although results retrieved with GPR were in agreement with disk infiltrometry tests, we stress that further research is needed to improve our algorithm so as to determine ³⁶⁷ the whole set of soil hydrodynamic parameters.

Notes

1. Common MidPoints 368

2. Wide-Angle Reflection- Refraction

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X - 20 LÉGER, SAINTENOY AND COQUET: INVERTING SATURATED HYDRAULIC CONDUCTIVITY

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Table 1. Hydrodynamic parameters for the numerical and field experiment. The $*$ indicates

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Figure 1. Experimental set up at its initial state.

Figure 2. Falling head infiltration from a 5-cm thick water layer. a) Permittivity profiles: each curve is plotted every 10 s. b) Radargram simulated with GprMax2D; reflection A is coming from the wetting front, B is the direct wave, A' and A" are multiples of reflection A. c) TWT time computed by the convolution algorithm from the permittivity profiles (plain red line) and TWT time obtained by picking of A peak in fig b).

Figure 3. Experimental GPR data acquired during the falling head infiltration (using a 5-cm initial water layer). Reflection A is the reflection coming from the wetting front

 $\frac{150}{150}$ $\frac{200}{250}$ $\frac{250}{300}$ $\frac{350}{350}$ $\frac{400}{450}$ $\frac{450}{450}$ 550 $\overline{600}$

Figure 4. Constant head infiltration with 5 cm of water. a) Permittivity profiles, each curve is plotted every 10 s. b) Radargram simulated with GprMax2D, reflection A is the wetting front, B is the direct wave, A' and A" are multiples. c) Two Way Travel Time computed with our convolution algorithm from the simulated permittivity profiles.

Figure 5. GPR data acquired during a constant head (5 cm) infiltration. Reflection A is the reflection coming from the wetting front.