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Convergence to steady states for solutions of a reaction-diffusion equation with mass conservation

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Abstract. We consider a reaction-diffusion equation with mass conservation, which was originally proposed by Rubinstein and Sternberg as a model for phase separation in a binary mixture. We study the large time behavior of the solution and show that it converges to a stationary solution as t tends to infinity. We also evaluate the rate of this convergence and compute precisely the limit stationary solution in one dimension.

Keywords: Gradient flow, Lojasiewicz inequality, Stabilization of solutions, Mass-conserved Allen-Cahn equation, Infinite dimensional dynamical systems.

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1 Introduction

We consider a reaction-diffusion equation with mass conservation

$$(P) \begin{cases} u_t = \Delta u + f(u) - \int_{\Omega} f(u) & \text{in } \Omega \times \mathbb{R}^+, \\ \partial_{\nu} u = 0 & \text{on } \partial\Omega \times \mathbb{R}^+, \\ u(x,0) = g_0(x) & x \in \Omega, \end{cases}$$

where $\Omega \subset \mathbb{R}^N (N \ge 1)$ is a smooth bounded domain, ∂_{ν} is the outer normal derivative to $\partial \Omega$ and

$$\int_{\Omega} f(u) := \frac{1}{|\Omega|} \int_{\Omega} f(u(x)) \, dx.$$

Problem (P) was introduced by Rubinstein and Sternberg [24] as a model for phase separation in a binary mixture. Our goal is to study of the large time behavior of the solutions of Problem (P). More precisely, we show that any solution of (P) converges to a steady state. We also evaluate the rate of this convergence and compute the limit steady state in one dimension. The main tool is a Lojasiewicz inequality that was first proposed by Lojasiewicz himself [18], [20]. He showed that any bounded solution to gradient systems in \mathbb{R}^n , (which is an ODE system), converges to a stationary point. This idea was subsequently developed for infinite-dimensional gradient systems by L. Simon, who showed a version of this inequality and applied it to prove the stabilization in the Allen-Cahn equation see [26].

This model is **mass preserving**, namely

$$\int_{\Omega} u(x,t) \, dx = \int_{\Omega} g_0(x) \, dx \quad \text{for all } t > 0,$$

and it possesses a free energy functional which coincides with the usual Allen-Cahn functional

$$\mathcal{E}(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 \, dx - \int_{\Omega} F(u) \, dx,$$

where $F(u) := \int_0^u f(s) ds$.

In [24], the authors consider the model in which f is bistable type, typical example $f(s) = s - s^3$. In this paper, we assume that the function f is of the following form

$$f(s) = \sum_{i=1}^{n} a_i s^i \quad \text{where } n \ge 3 \text{ is an odd number, } a_n < 0.$$
(1)

Note that there exists a constant $c_1 > 0$ satisfying

$$f'(s) \le c_1 \text{ for all } s \in \mathbb{R}.$$
 (2)

Definition 1.1. Two constants $s_1 < s_2$ are said to satisfy Property (C) with respect to f, if

(C):
$$f(s_2) < f(s) < f(s_1)$$
 for all $s \in (s_1, s_2)$.

We have two following theorems.

Theorem 1.2. Let $g_0 \in L^{\infty}(\Omega)$. Problem (P) possesses a unique solution $u \in C([0,\infty); L^2(\Omega)))$ which satisfies for every T > 0

$$u \in L^{\infty}(Q_T) \cap L^2(0,T; H^1(\Omega)) \text{ and } u_t \in L^2(0,T; (H^1(\Omega))'),$$

where $Q_T := \Omega \times (0,T)$. Moreover $u \in C^{\infty}(\overline{\Omega} \times (0,+\infty))$ and

$$\{u(t), t \geq 1\}$$
 is relatively compact in $C^m(\Omega)$ for all $m \in \mathbb{N}$.

Theorem 1.3. Let $g_0 \in L^{\infty}(\Omega)$ and let u be the unique solution of Problem (P). Then

$$\lim_{t\to\infty}\|u(t)-\varphi\|_{C^m(\overline{\Omega})}=0 \text{ as } t\to\infty, \text{ for all } m\in\mathbb{N},$$

where φ is a smooth solution of the stationary problem

$$(S) \begin{cases} \Delta \varphi + f(\varphi) - \int_{\Omega} f(\varphi) &= 0 \text{ in } \Omega, \\ \partial_{\nu} \varphi &= 0 \text{ on } \partial \Omega. \end{cases}$$

Moreover

$$\int_{\Omega} \varphi = \int_{\Omega} g_0.$$

The proofs of Theorem 1.2 and Theorem 1.3 are based on following auxiliary problem.

Problem (**P**): In order to prove the main theorems, we first study Problem (P) when f is replaced by a smooth bounded function \overline{f} . More precisely, we consider the following problem

$$(\hat{P}) \begin{cases} \bar{u}_t = \Delta \bar{u} + \bar{f}(\bar{u}) - \int_{\Omega} \bar{f}(\bar{u}) & \text{in } \Omega \times \mathbb{R}^+, \\ \partial_{\nu} \bar{u} = 0 & \text{on } \partial\Omega \times \mathbb{R}^+, \\ \bar{u}(x,0) = g_0(x) & x \in \Omega, \end{cases}$$

and \bar{f} is supposed to be smooth and satisfy

 $|\bar{F}(s)|, |\bar{f}(s)|, |\bar{f}(s)s|, |\bar{f}'(s)| < c_2 \text{ for all } s \in \mathbb{R} \text{ and a constant } c_2 > 0.$ (3)

Here, $\bar{F}(s) = \int_0^s \bar{f}(\tau) d\tau$ is a primitive of \bar{f} . The corresponding Lyapunov functional is given by

$$E(\bar{u}) := \frac{1}{2} \int_{\Omega} |\nabla \bar{u}|^2 \, dx - \int_{\Omega} \bar{F}(\bar{u}) dx.$$

The choice of \overline{f} will be given in the proof of Theorem 1.2 (section 2). We shall show with such \overline{f} that the solution of Problem (\overline{P}) coincides with the solution of Problem (P) and deduce the result for Problem (P).

The paper is organized as follows: in section 2 we prove Theorem 1.2. In section 3, first we give a version of Lojasiewicz inequality for the energy functional E (Theorem 3.11), then apply it to prove Theorem 1.3 and to establish rate of this convergence. Section 4 is devoted to compute the limit stationary solution in one dimension.

2 The existence and uniqueness of solution

We first prove the uniqueness of the solutions of the Problem (P) and (\bar{P}) .

Lemma 2.1. (i) For any $g_0 \in L^{\infty}(\Omega)$, Problem (P) possesses at most one solution u such that

$$u \in L^{\infty}(Q_T) \cap L^2(0, T, H^1(\Omega)) \text{ and } u_t \in L^2(0, T; H^1(\Omega)').$$

Moreover,

$$\int_{\Omega} u(x,t)dx = \int_{\Omega} g_0(x) \, dx \text{ for all } t > 0.$$
(4)

(ii) For any $g_0 \in L^2(\Omega)$, Problem (\bar{P}) possesses at most one solution \bar{u} such that

$$\bar{u} \in L^{\infty}(0,T;L^{2}(\Omega)) \cap L^{2}(0,T,H^{1}(\Omega)) \text{ and } \bar{u}_{t} \in L^{2}(0,T;H^{1}(\Omega)').$$

Moreover,

$$\int_{\Omega} \bar{u}(x,t)dx = \int_{\Omega} g_0(x)dx \text{ for all } t > 0.$$
(5)

Proof. (ii) The identity (5) follows from

$$\frac{d}{dt} \int_{\Omega} \bar{u}(x,t) \, dx = 0$$

Let $g'_0 \in L^2(\Omega)$ be another initial condition and let \bar{u}' be a corresponding solution. We first set $w := \bar{u} - \bar{u}'$, multiply the difference of the equations for \bar{u} and \bar{u}' by w, then integrate over Ω to obtain

It follows from (3) that

$$\frac{1}{2}\frac{d}{dt}\int_{\Omega}w^{2}(t) + \int_{\Omega}|\nabla w(t)|^{2} \leq c_{2}\int_{\Omega}w^{2} + \frac{c_{2}}{|\Omega|}\left(\int_{\Omega}|w|\right)^{2}$$
$$\leq c_{3}\int_{\Omega}w^{2}.$$

Using Gronwall's lemma, we have

$$\int_{\Omega} |\bar{u}(t) - \bar{u}'(t)|^2 \le \exp(2c_3 t) \int_{\Omega} |g_0 - g_0'|^2, \tag{6}$$

which implies the uniqueness.

(i) The proof of part (i) is similar to the one of part (ii). Let $\tilde{g}_0 \in L^{\infty}(\Omega)$ such that $\int_{\Omega} \tilde{g}_0 = \int_{\Omega} g_0$ and let \tilde{u} be a solution corresponding with the initial condition \tilde{g}_0 . Note that

$$\int_{\Omega} u(t) - \int_{\Omega} \tilde{u}(t) = \int_{\Omega} g_0 - \int_{\Omega} \tilde{g}_0 = 0.$$
(7)

We use (2) and the method developed in the proof of part (ii) to obtain

$$\int_{\Omega} |u(t) - \tilde{u}(t)|^2 \le \exp(2c_1 t) \int_{\Omega} |g_0 - \tilde{g}_0|^2.$$
(8)

The uniqueness then follows from this estimate.

Lemma 2.2. Let $g_0 \in L^2(\Omega)$. Then for any T > 0 arbitrary

(i) Problem (\hat{P}) possesses a unique solution \bar{u} which satisfies

$$\bar{u} \in L^{\infty}(0,T;L^{2}(\Omega)) \cap L^{2}(0,T;H^{1}(\Omega)); \ \bar{u}_{t} \in L^{2}(0,T;(H^{1}(\Omega))').$$
 (9)

Moreover, $\bar{u} \in C([0,\infty); L^2(\Omega)).$

(ii) If, in addition, $g_0 \in H^1(\Omega)$ then

$$\bar{u} \in L^{\infty}(0,T; H^1(\Omega)) \cap L^2(0,T; H^2(\Omega)) \text{ and } \bar{u}_t \in L^2(Q_T).$$
 (10)

So that also $\bar{u} \in C([0,\infty); H^1(\Omega))$.

The mapping

$$T(t): g_0 \longmapsto \bar{u}(t)$$

is Lipschitz continuous on $L^2(\Omega)$ for all t > 0 and $T(t)_{t \ge 0}$ is a semigroup on $L^2(\Omega)$.

Remark 2.3. Consider the operator $-\Delta$ with homogeneous Neumann boundary conditions. Denote by

$$0 = \lambda_1 < \lambda_2 \le \dots \le \lambda_i \le \dots$$

the eigenvalues of the Laplacian operator, and $w_i, i = 1, \ldots$ the corresponding unit eigenfunctions. Remark that $w_1 = \frac{1}{|\Omega|^{1/2}}, \int_{\Omega} w_i = 0$ for $i \ge 2$ and that the w_i 's are smooth functions up to boundary. They constitute an orthonormal basis of $L^2(\Omega)$ and also an orthogonal basis of $H^1(\Omega)$.

Proof of Lemma 2.2

We look for an approximate solution of the form

$$\bar{u}_m(t) = \sum_{i=1}^m g_{mi}(t) w_i$$

satisfying

$$\int_{\Omega} \bar{u}_{mt} w_j + \int_{\Omega} \nabla \bar{u}_m \nabla w_j = \int_{\Omega} \bar{f}(\bar{u}_m) w_j - \oint_{\Omega} \bar{f}(\bar{u}_m) \int_{\Omega} w_j$$
(11)

for $j = 1, \ldots, m$ and

$$\bar{u}_m(0) = \bar{u}_{m0} := \sum_{i=1}^m g_{mi}^0 w_i \to g_0 \text{ in } L^2(\Omega) \text{ as } m \to \infty.$$
 (12)

Note that $\int_{\Omega} w_j = 0$ for $2 \le j \le m$, we deduce that the equations (11) form a nonlinear differential system for the functions g_{m1}, \ldots, g_{mm} :

$$g'_{m1}(t) = 0, (13)$$

$$g'_{mj} - \lambda_j g_{mj} = \int_{\Omega} \bar{f} \Big(\sum_{i=1}^m g_{mi}(t) w_i \Big) w_j \quad \text{for } 2 \le j \le m.$$
(14)

The condition (12) is equivalent to the *m* scalar initial conditions

$$g_{m1}(0) = g_{m1}^0, (15)$$

$$g_{mj}(0) = g_{mj}^0 \text{ for } 2 \le j \le m.$$
 (16)

It follows from (13) and (15) that

$$g_{m1}(t) = g_{m1}^0 \text{ for all } t \ge 0.$$
 (17)

Substituting (17) in (14), we obtain a nonlinear differential system of m-1 variables. Then the nonlinear differential system (14) with the initial condition (16) has a maximal solution defined on some interval $(0, T_m)$. In fact $T_m = \infty$ because of the following *a priori* estimates.

A priori estimates for the proof of Lemma 2.2(i): First, by (12) and (17), we have

$$\int_{\Omega} \bar{u}_m(t) \, dx = \int_{\Omega} \bar{u}_{m0}(x) \, dx \to \int_{\Omega} g_0(x) \, dx \text{ as } m \to \infty \tag{18}$$

for all $0 \le t < T_m$. Consequently, there exists a constant c_4 independent of m such that

$$\left| \int_{\Omega} \bar{u}_m(t) \, dx \right| \le c_4 \text{ for all } 0 \le t < T_m.$$
(19)

Multiplying (11) by g_{mj} and summing on $j = 1, \ldots, m$ we obtain

$$\frac{1}{2}\frac{d}{dt}\int_{\Omega}|\bar{u}_m|^2 + \int_{\Omega}|\nabla\bar{u}_m|^2 = \int_{\Omega}\bar{f}(\bar{u}_m)\bar{u}_m - \int_{\Omega}\bar{f}(\bar{u}_m)\int_{\Omega}\bar{u}_m.$$

We use (3) and (19) to deduce that

$$\frac{1}{2}\frac{d}{dt}\int_{\Omega}|\bar{u}_m|^2 + \int_{\Omega}|\nabla\bar{u}_m|^2 \le c_2(|\Omega| + c_4).$$

Therefore we integrate this inequality from 0 to t with t arbitrary, and then take t = T to deduce that

$$\|\bar{u}_m\|_{L^{\infty}(0,T;L^2(\Omega))}, \|\bar{u}_m\|_{L^2(0,T;H^1(\Omega))} \le K_1(\|g_0\|_{L^2(\Omega))} + T),$$
(20)

where K_1 is a constant independent of m.

Now we will give an estimate for \bar{u}_{mt} . Fix any $\eta \in H^1(\Omega)$, with $\|\eta\|_{H^1(\Omega)} \leq 1$, and write $\eta = \eta_1 + \eta_2$ where $\eta_1 \in \operatorname{span}\{w_i\}_{i=1}^m$ and $\int_{\Omega} \eta_2 w_i = 0$ for all $i = 1, \ldots, m$. Note that $\eta_2 \in H^1(\Omega)$ and that

$$\int_{\Omega} \nabla \eta_2 \nabla w_i = -\lambda_i \int_{\Omega} \eta_2 w_i = 0 \text{ for all } 1 \le i \le m$$

so that η_2 is also orthogonal to span $\{w_i\}_{i=1}^m$ with respect to the scalar product in $H^1(\Omega)$. In particular, η_2 is orthogonal to η_1 in $H^1(\Omega)$, hence

$$\|\eta_1\|_{H^1(\Omega)} \le \|\eta\|_{H^1(\Omega)} \le 1.$$

We deduce from (11) that

$$\int_{\Omega} \bar{u}_{mt} \eta = \int_{\Omega} \bar{u}_{mt} \eta_1 = -\int_{\Omega} \nabla \bar{u}_m \nabla \eta_1 + \int_{\Omega} \bar{f}(\bar{u}_m) \eta_1 - \oint_{\Omega} \bar{f}(\bar{u}_m) \int_{\Omega} \eta_1.$$

Therefore,

$$\left| \int_{\Omega} \bar{u}_{mt} \eta \right| \leq \left(\sum_{i=1}^{N} \| \frac{\partial \bar{u}_m}{\partial x_i} \|_{L^2(\Omega)} \right) \left(\sum_{i=1}^{N} \| \frac{\partial \eta_1}{\partial x_i} \|_{L^2(\Omega)} \right) + 2c_2 \| \eta_1 \|_{L^1(\Omega)}$$
$$\leq \| \nabla \bar{u}_m \|_{L^2(\Omega)} \| \nabla \eta_1 \|_{L^2(\Omega)} + c_5 \| \eta_1 \|_{L^2(\Omega)}$$
$$\leq \| \nabla \bar{u}_m \|_{L^2(\Omega)} + c_5,$$

where c_5 is a constant independent of m. Since $\eta \in H^1(\Omega)$ is an arbitrary function such that $\|\eta\|_{H^1(\Omega)} \leq 1$, it follows that

$$\|\bar{u}_{mt}\|_{(H^1(\Omega))'} \le \|\nabla \bar{u}_m\|_{L^2(\Omega)} + c_5.$$

This together with (20) implies that there exists a positive constant K_2 independent of m such that

$$\|\bar{u}_{mt}\|_{L^2(0,T;(H^1(\Omega))')} \le K_2.$$
(21)

The estimates (20) and (21) will be necessary for the proof (i).

A priori estimates for the proof of Lemma 2.2(ii): We suppose that $\bar{u}_{0m} = \sum_{i=1}^{m} g_{mi}^{0} w_i \to g_0$ in $H^1(\Omega)$. First, note that

$$\int_{\Omega} \Delta \bar{u}_m \, dx = 0, \text{ and } \int_{\Omega} \bar{u}_{mt} \, dx = 0$$

Multiplying Equation (11) by $-\lambda_j g_{mj}$, summing on $j = 1, \ldots, m$ yields

$$\frac{1}{2}\frac{d}{dt}\int_{\Omega}|\nabla \bar{u}_{m}|^{2} + \int_{\Omega}|\Delta \bar{u}_{m}|^{2} = -\int_{\Omega}\bar{f}(\bar{u}_{m})\Delta \bar{u}_{m} + \int_{\Omega}\bar{f}(\bar{u}_{m})\int_{\Omega}\Delta \bar{u}_{m}$$
$$= \int_{\Omega}\bar{f}'(\bar{u}_{m})|\nabla \bar{u}_{m}|^{2}$$
$$\leq c_{2}\int_{\Omega}|\nabla \bar{u}_{m}|^{2}.$$

Integrating this inequality from 0 to T (T > 0 is arbitrary), we deduce that there exists a constant K_3 independent of m such that

$$\|\bar{u}_m\|_{L^{\infty}(0,T;H^1(\Omega))}, \|\bar{u}_m\|_{L^2(0,T;H^2(\Omega))} \le K_3.$$
(22)

Finally, we multiply (11) by $g'_{mj}(t)$ and sum on j = 1, ..., m to obtain

$$\int_{\Omega} \bar{u}_{mt}^2 + \int_{\Omega} \nabla \bar{u}_m \nabla \bar{u}_{mt} = \int_{\Omega} \bar{f}(\bar{u}_m) \bar{u}_{mt} - \int_{\Omega} \bar{f}(\bar{u}_m) \int_{\Omega} \bar{u}_{mt} = \frac{d}{dt} \int_{\Omega} \bar{F}(\bar{u}_m).$$

So that

$$\int_0^T \int_\Omega \bar{u}_{mt}^2 + \frac{1}{2} \int_\Omega |\nabla \bar{u}_m(T)|^2 = \int_\Omega \bar{F}(\bar{u}_m(T)) - \int_\Omega \bar{F}(\bar{u}_m(0)) + \frac{1}{2} \int_\Omega |\nabla \bar{u}_m(0)|^2,$$

which implies that

$$\|\bar{u}_{mt}\|_{L^2(Q_T)} \le K_4,\tag{23}$$

where K_4 is independent of m.

Passing to the limit: It follows from (20) and (21) that there exists a subsequence of \bar{u}_m , still denoted by \bar{u}_m , such that

 $\bar{u}_m \rightarrow \bar{u}$ weakly in $L^2(0, T; H^1(\Omega))$, weak-star in $L^{\infty}(0, T; L^2(\Omega))$, $\bar{u}_m \rightarrow \bar{u}$ strongly in $L^2(0, T; L^2(\Omega))$, $\bar{u}_{mt} \rightarrow \bar{u}_t$ weakly in $L^2(0, T; (H^1(\Omega))')$.

Moreover, since \bar{f} is Lipschitz continuous on \mathbb{R} ,

 $\bar{f}(\bar{u}_m) \to \bar{f}(\bar{u})$ strongly in $L^2(0,T;L^2(\Omega))$.

Passing to the limit in (11) we deduce that

$$\bar{u}_t = \Delta \bar{u} + \bar{f}(\bar{u}) - \int_{\Omega} \bar{f}(\bar{u}) \text{ in } L^2(0,T;(H^1(\Omega))').$$

Since $\bar{u} \in L^2(0,T; H^1(\Omega))$ and since $\bar{u}_t \in L^2(0,T; (H^1(\Omega))')$, one can shows that $\bar{u} \in C([0,T]; L^2(\Omega))$. Moreover $\bar{u}(0,x) = g_0(x)$ by classical arguments.

If $g_0 \in H^1(\Omega)$, it follows from (22) and (23) that \bar{u} satisfies (10). By [10, Theorem 4, p. 288] it follows $\bar{u} \in C([0, \infty); H^1(\Omega))$.

The uniqueness is proved in Lemma 2.1. The Lipstchitz continuity of T(t) on $L^2(\Omega)$ follows from (6). \Box

The following corollary follows directly from Lemma 2.2.

Corollary 2.4. For any $g_0 \in L^2(\Omega)$ and $\delta > 0$, we have

$$\bar{u} \in C([\delta, \infty); H^1(\Omega)).$$

We now prove more regularity properties of the solution. For this, we recall some technical lemmas which will be used in what follows. Denote by $Q_a^b = \Omega \times (a, b)$ for $a \in \mathbb{R}$ and $b \in \mathbb{R} \cup \{+\infty\}$.

Lemma 2.5. Let $u_0 \in L^2(\Omega), g \in L^p(Q_T)$ for some $p \in (1, \infty)$ and let u be the solution of the time evolution problem

$$\begin{cases} u_t - \Delta u = f & \text{in } Q_T, \\ \partial_{\nu} u = 0 & \text{on } \partial \Omega \times (0, T), \\ u(x, 0) = u_0(x) & x \in \Omega. \end{cases}$$

Then

(i) for each $0 < \delta < 1$, there exists a positive constant C such that

$$\|u\|_{W_p^{2,1}(Q_{\delta}^T)} \le C(\|u_0\|_{L^2(\Omega)} + \|f\|_{L^p(Q_T)}).$$

(ii) If we suppose that $u_0 \in W^{2,p}(\Omega)$ then we obtain the estimate

 $\|u\|_{W_p^{2,1}(Q_T)} \le C(\|u_0\|_{W^{2,p}(\Omega)} + \|f\|_{L^p(Q_T)}).$

Lemma 2.6. One has the following embedding

$$W_p^{2,1}(Q_T) \subset C^{\lambda,\lambda/2}(\bar{Q}_T) \text{ with } \lambda = 2 - \frac{N+2}{p} \text{ if } p > \frac{N+2}{2} \text{ and } p \neq N+2.$$

Lemma 2.5 and Lemma 2.6 follow from [17, chapter 4, section 3 and chapter 2, section 3] which are stated in [4, p. 206].

Proposition 2.7. Let $\bar{u} \in C^{2,1}(\overline{\Omega} \times (0,T]) \cap C(\overline{\Omega} \times [0,T])$ be the solution of Problem (\bar{P}) with the initial condition g_0 . Assume that $s_1 < g_0 < s_2$ then

$$s_1 < \bar{u}(x,t) < s_2$$

for $x \in \overline{\Omega}$ and $0 < t \leq T$.

Proof. For the purpose of contradiction, we suppose that there exists a first time $0 < t_0 \leq T$ such that $\bar{u}(x_0, t_0) = s_1$ or $\bar{u}(x_0, t_0) = s_2$ for some $x_0 \in \overline{\Omega}$. Without loss of generality, assume that $\bar{u}(x_0, t_0) = s_2$. By the continuity of \bar{u} and the definition of t_0 , we have

$$s_1 \leq \overline{u}(x, t_0) \leq s_2$$
 for all $x \in \Omega$, and $\overline{u}(x, t) < s_2$ for all $x \in \Omega$ and $0 \leq t < t_0$
(24)

Since $\partial_{\nu}\bar{u} = 0$, we deduce from Hopf's maximum principle that $x_0 \in \Omega$. Therefore the function $\bar{u}(\cdot, t_0)$ attains its maximum at $x_0 \in \Omega$, which implies that $\Delta \bar{u}(x_0, t_0) \leq 0$. By (24), we have

$$\bar{u}_t(x_0, t_0) = \lim_{\Delta t \to 0^+} \frac{\bar{u}(x_0, t_0 - \Delta t) - \bar{u}(x_0, t_0)}{-\Delta t_0} \ge 0,$$

which we substitute in Problem (\bar{P}) to obtain $\oint_{\Omega} (\bar{f}(s_2) - \bar{f}(\bar{u}(x, t_0))) dx \ge 0$. Since $s_1 \le \bar{u}(x, t_0) \le s_2$ for all $x \in \Omega$, it follows that $\bar{f}(s_2) \le \bar{f}(\bar{u}(x, t_0))$ for all $x \in \Omega$ so that $\bar{f}(s_2) = \bar{f}(\bar{u}(x, t_0))$ and hence $\bar{u}(x, t_0) = s_2$ for all x in Ω . But the inequality $\bar{u}(x, \frac{t_0}{2}) < s_2 = \bar{u}(x, t_0)$ for all $x \in \Omega$ implies that

$$\int_{\Omega} \bar{u}(x, \frac{t_0}{2}) \, dx < \int_{\Omega} \bar{u}(x, t_0) \, dx,$$

which contradicts the mass preserving property.

Lemma 2.8. Let $0 < \delta < \frac{1}{2}$, $\alpha \in (0,1)$ be arbitrary. Let s_1, s_2 be two constants which satisfy Property (\mathcal{C}) with respect to \overline{f} . We assume that $s_1 < g_0 < s_2$ then

(i)
$$s_1 \leq \overline{u}(x,t) \leq s_2$$
 for all $x \in \overline{\Omega}, t > 0$,
(ii) $\|\overline{u}\|_{C^{2m+1+\alpha,m+\frac{1+\alpha}{2}}(Q_{\delta}^{\infty})} \leq C(m,\delta,s_1,s_2,\Omega)$ for any $m \in \mathbb{N}$.

Proof. (i) The main idea of the proof is to approximate the initial function g_0 by a sequence of smooth functions and to first obtain uniform a priori estimates for the corresponding solutions of Problem (\bar{P}) . To begin with, we choose a sequence $g_{0n} \in C^{\infty}(\bar{\Omega})$ such that $s_1 < g_{0n} < s_2$ converges to g_0 in $L^2(\Omega)$ as $n \to \infty$ (cf. [1, Lemma 2.18, p. 29]). We denote by \bar{u}_n the corresponding solutions of Problem (\bar{P}) . Since $|\bar{f}| \leq c_2$, we deduce from Lemma 2.5(ii) that $\bar{u}_n \in W_p^{2,1}(Q_T)$ for all p > 1, which by the Sobolev embedding in Lemma 2.6 implies that $\bar{u}_n \in C^{1+\alpha,(1+\alpha)/2}(\bar{Q}_T)$ for all $\alpha \in (0, 1)$. Applying a standard bootstrap argument (cf. [17, Theorem 10.1, p, 351]), we deduce that $\bar{u}_n \in C^{3+\alpha,1+\frac{1+\alpha}{2}}(\bar{Q}_T)$, which by Proposition 2.7 is such that $s_1 < \bar{u}_n < s_2$.

In view of Lemma 2.5(i) we also deduce that $\bar{u} \in W_p^{2,1}(Q_{\delta}^T)$ for all $\delta \in (0, \frac{1}{2})$ and p > 1, which in turn implies that $\bar{u} \in C(\overline{\Omega} \times (0, T])$. Since by (6),

$$\int_{\Omega} |\bar{u}(t) - \bar{u}_n(t)|^2 \le \exp(2c_3 t) \int_{\Omega} |g_0 - g_{0n}|^2$$

it follows that

$$\bar{u}_n(t) \to \bar{u}(t)$$
 in $L^2(\Omega)$ for all $t > 0$.

Thus $s_1 \leq \overline{u} \leq s_2$ in $\overline{\Omega} \times (0, T]$. Finally, since T is arbitrary,

$$s_1 \leq \overline{u}(x,t) \leq s_2$$
 for all $x \in \overline{\Omega}, t > 0$.

(ii) Reasoning in the same way as for the function \bar{u}_n on the domains Q_k^{k+1} and $Q_{k+1/2}^{k+3/2}$, we deduce that

$$\|\bar{u}\|_{C^{1+\alpha,(1+\alpha)/2}(Q_{k+\delta}^{k+1})} \le C \left(\|\bar{u}(k)\|_{L^2(\Omega)} + \left\|\bar{f}(\bar{u}) - \int_{\Omega} \bar{f}(\bar{u})\right\|_{L^p(Q_k^{k+1})} \right) \le C(\delta, s_1, s_2, \Omega)$$

and a similar one on the domain $Q_{k+1/2}^{k+3/2}$. Therefore, we have

 $\|\bar{u}\|_{C^{1+\alpha,(1+\alpha)/2}(Q^{\infty}_{\delta})} \le C(\delta, s_1, s_2, \Omega).$

By standard bootstrap arguments [17, Theorem 10.1, p. 351], Lemma 2.8(ii) holds. $\hfill \Box$

Corollary 2.9. Let s_1, s_2 be two constants which satisfy Property (C) with respect to \overline{f} . Assume that $s_1 < g_0 < s_2$ then $\{\overline{u}(t) : t \ge 1\}$ is relatively compact in $C^m(\overline{\Omega})$ for all $m \in \mathbb{N}$.

In particular, $\{\bar{u}(t): t \geq 1\}$ is relatively compact in $H^1(\Omega)$.

Proof of Theorem 1.2

Construction function \overline{f} : first, choose two constants s_1, s_2 (such that $|s_1|, |s_2|$ large enough) which satisfy Property (\mathcal{C}) and satisfy

 $s_1 < g_0(x) < s_2$ for almost $x \in \Omega$.

It is easy to see that such s_1, s_2 exist. Now, let $\overline{f} \in C_c^{\infty}(\mathbb{R})$ be given by

$$\bar{f}(s) = \begin{cases} 0 & \text{if } s \le s_1 - 2, \\ f(s) & \text{if } s \in [s_1 - 1, s_2 + 1], \\ 0 & \text{if } s \ge s_2 + 2. \end{cases}$$

Obviously, this function satisfy (3).

Note that s_1, s_2 satisfy Property (\mathcal{C}) with respect to \bar{f} . According to Lemma 2.7, we have $\bar{u}(t) \in [s_1, s_2]$. On the other hand, $\bar{f}(s) = f(s)$ for all $s \in [s_1, s_2]$, so that $\bar{f}(\bar{u}) = f(u)$. It follows that \bar{u} coincides with the unique solution of Problem (P). Now, the assertions of the theorem follow from Lemma 2.2, Lemma 2.8(ii) and Corollary 2.9. \Box **Remark 2.10.** We can extend all results in this section to f being a smooth function in \mathbb{R} . We summarize two important results.

- (i) Theorem 1.2 holds provided that f is smooth.
- (ii) A comparison result: let f be a smooth function and let $s_1 < s_2$ be two constants which satisfy Property (C) with respect to f. Assume that $s_1 < g_0 < s_2$ then $s_1 \le u(x,t) \le s_2$.

3 Large time behavior

In this section, we first give some modifications for Problem (P), which bright out a slight view for this problem. Lojasiewicz inequality is proven in section 3.2, then we apply the inequality to prove the convergence of the solution to a steady state. A result of velocity of the convergence is also established in this section.

3.1 Modification setting

In the following, by mass conservation property, we may without loss of generality assume that u satisfies the condition

$$\int_{\Omega} u(x,t) \, dx = 0 \quad \text{for all } t \ge 0.$$

In fact, it suffices to replace the solution u by $u-m_0$ with $m_0 := \frac{1}{|\Omega|} \int_{\Omega} g_0(x) dx$ to note that $u - m_0$ satisfies Problem (P) with f replaced by $f_1(s) := f(s + m_0)$. The question, whether u converges as $t \to \infty$ is not affected by this normalization.

Next, we fix two constants s_1, s_2 (such that $|s_1|, |s_2|$ large enough) which satisfy Property (C) and satisfy

$$s_1 < g_0(x) < s_2$$
 for almost $x \in \Omega$.

Function \bar{f} is given as in the proof of Theorem 1.2. More precisely, $\bar{f} \in C_c^{\infty}(\mathbb{R})$ and

$$\bar{f}(s) = \begin{cases} 0 & \text{if } s \le s_1 - 2, \\ f(s) & \text{if } s \in [s_1 - 1, s_2 + 1], \\ 0 & \text{if } s \ge s_2 + 2. \end{cases}$$

By using the fact that the solution of Problem (\bar{P}) coincides with the unique solution of Problem (P), we shall study the large time behavior for Problem (\bar{P}) and deduce the result for Problem (P).

3.2 A version of Lojasiewicz inequality

We define the spaces

 $H = \{ u \in L^2(\Omega), \int_{\Omega} u(x) dx = 0 \}, \text{ equipped with the norm } \|\cdot\|_H := \|\cdot\|_{L^2(\Omega)},$ $V = \left\{ u \in H^1(\Omega) : \int_{\Omega} u(x) dx = 0 \right\}, \text{ equipped with the norm } \|\cdot\|_V := \|\cdot\|_{H^1(\Omega)}.$ Let V^* be the dual space of V and . We identify H with its dual to obtain:

 $V \hookrightarrow H \hookrightarrow V^*,$

where the embeddings $V \hookrightarrow H$, $H \hookrightarrow V^*$ are continuous, dense and compact (see e.g. [15, p. 677]). We use $\langle \cdot, \cdot \rangle$ to denote the scalar product in H and the scalar product for the duality V^*, V . We denote by $\mathcal{L}(X, Y)$ the space of bounded linear operators from a Banach space X to a second Banach space Y, and we write $\mathcal{L}(X) := \mathcal{L}(X, X)$.

Throughout the following, we denote by $C \ge 0$ a generic constant which may vary from line to line. We start with the following result.

Lemma 3.1. Let $1 \leq p < \infty$ and let g be a continuous, bounded function on \mathbb{R} . We define an operator $B : H^1(\Omega) \to L^p(\Omega)$ by $B(u)(x) := \overline{f}(u(x))$. Then B is continuous from $H^1(\Omega)$ to $L^p(\Omega)$.

Proof. It follows from [16, Lemma 16.1, p. 60] that B is continuous from $L^2(\Omega)$ to $L^p(\Omega)$. The result holds because of the continuous embedding $H^1(\Omega) \hookrightarrow L^2(\Omega)$.

Lemma 3.2. The functional E is twice continuously Fréchet differentiable on V. We denote by E', L be the first and second derivative of E, respectively. Then

(i) The first derivative

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$$E': V \longrightarrow V^* \text{ is given by}$$

$$E'(u), h\rangle_{V^*, V} = \int_{\Omega} \nabla u \nabla h - \int_{\Omega} \bar{f}(u)h \quad \text{for all } u, h \in V. \quad (25)$$

(ii) The second derivative

$$L: V \longrightarrow \mathcal{L}(V, V^*) \text{ is given by}$$
$$\langle L(u)h, k \rangle_{V^*, V} = \int_{\Omega} \nabla h \nabla k - \int_{\Omega} \bar{f}'(u)hk \quad \text{for all } u, h, k \in V.$$
(26)

Proof. We write E as the difference of E_1 and E_2 , where

$$E_1(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx \text{ and } E_2(u) = \int_{\Omega} \bar{F}(u) dx.$$
 (27)

Obviously, E_1 is twice continuously Fréchet differentiable. Its derivatives are easily identified in the formula (25) and (26). We now prove the differentiability of E_2 .

By Taylor's formula, there exists $\theta(x) \in (0, 1)$ such that

$$\overline{F}(u+h) - \overline{F}(u) = \overline{f}(u+\theta h)h$$
 for all $u, h \in V$.

Il follows that

$$\left| E_2(u+h) - E_2(u) - \int_{\Omega} \bar{f}(u)h \, dx \right|$$

 $\leq \int_{\Omega} |\bar{f}(u+\theta h) - \bar{f}(u)| \, |h| \, dx \leq C \|\bar{f}(u+\theta h) - \bar{f}(u)\|_{L^2(\Omega)} \|h\|_V.$

Note that $u + \theta h$ tends to u in $H^1(\Omega)$ as $h \to 0$ in V; it follows from Lemma 3.1 that $\|\bar{f}(u+\theta h) - \bar{f}(u)\|_{L^2(\Omega)}$ tends to 0 as $h \to 0$ in V. Thus

$$\left| E_2(u+h) - E_2(u) - \int_{\Omega} \bar{f}(u)h \, dx \right| = o(\|h\|_V) \text{ as } h \to 0.$$

This implies that the first derivative E'_2 exists and

$$\langle E'_2(u),h\rangle_{V^*,V} = \int_{\Omega} \bar{f}(u)h\,dx.$$

The Fréchet differentiability of E'_2 is shown in a similar way. Choose p > 2 such that V is continuously embedded in $L^p(\Omega)$. Let T be a linear mapping from V to V^* given by

$$\langle T h, k \rangle_{V^*, V} = \int_{\Omega} \bar{f}'(u) h k \, dx.$$

We will use below a generalized Holder inequality based on the identity

$$\frac{1}{p} + \frac{1}{p} + \frac{p-2}{p} = 1.$$

For every $u, h, k \in V$, there exist $\eta(x) \in (0, 1)$ such that

$$\begin{aligned} \left| \langle E_{2}'(u+h) - E_{2}'(u) - Th, k \rangle_{V^{*}, V} \right| \\ &\leq \int_{\Omega} \left| \bar{f}'(u+\eta h) - \bar{f}'(u) \right| |h| \, |k| \, dx \\ &\leq \| \bar{f}'(u+\eta h) - \bar{f}'(u) \|_{L^{p/(p-2)}(\Omega)} \|h\|_{L^{p}(\Omega)} \|k\|_{L^{p}(\Omega)} \\ &\leq C \| \bar{f}'(u+\eta h) - \bar{f}'(u) \|_{L^{p/(p-2)}(\Omega)} \|h\|_{V} \|k\|_{V}, \end{aligned}$$
(28)

It follows from (28) that

$$||E'_{2}(u+h) - E'_{2}(u) - Th||_{V^{*}} \le C ||\bar{f}'(u+\eta h) - \bar{f}'(u)||_{L^{p/(p-2)}(\Omega)} ||h||_{V}.$$

Since $p/(p-2) < +\infty$ and since \bar{f}' is bounded, $\|\bar{f}'(u+\eta h) - \bar{f}'(u)\|_{L^{p/(p-2)}(\Omega)}$ tends to 0 as $h \to 0$. Thus

$$||E'_2(u+h) - E'_2(u) - Th||_{V^*} = o(||h||_V)$$

which implies that

$$\langle E_2''(u)h,k\rangle_{V^*,V} = \int_{\Omega} \bar{f}'(u)h\,k$$
 for all $u,h,k \in V$.

On the other hand,

$$\begin{aligned} |\langle (E_2''(u) - E_2''(v))h, k \rangle_{V^*, V}| &\leq \int_{\Omega} |\bar{f}'(u) - \bar{f}'(v)| |h| |k| \, dx \\ &\leq C \|\bar{f}'(u) - \bar{f}'(v)\|_{L^{p/(p-2)}(\Omega)} \|h\|_V \|k\|_V, \end{aligned}$$

so that

$$\|E_2''(u) - E_2''(v)\|_{\mathcal{L}(V,V^*)} \le C \|\bar{f}'(u) - \bar{f}'(v)\|_{L^{p/p-2}(\Omega)}$$

This estimate implies the continuity of E_2'' .

We define a continuous bilinear form from $V\times V\to \mathbb{R}$ by

$$a(u,v) = \int_{\Omega} \nabla u \nabla v \, dx.$$

Lemma 3.3. There exists an isomorphism A from V onto V^* such that

$$a(u,v) = \langle Au, v \rangle_{V^*,V} \text{ for all } u, v \in V.$$
(29)

Proof. For each $u \in V$ we define the functional T_u on V by $T_u(v) = a(u, v)$. By continuity of the bilinear form a, $|T_u(v)| \leq C ||u||_V ||v||_V$. Thus, $T_u \in V^*$ and

$$||T_u||_{V^*} \le C ||u||_V. \tag{30}$$

We define the operator $A: V \to V^*$ by $A(u) = T_u$. It follows from (30) that A is continuous.

Next, we prove that A is injective. Let $u \in V$ such that Au = 0. Then a(u, v) = 0 for all $v \in V$. In particular, a(u, u) = 0. Recall that with Poincaré-Wirtinger inequality for all $w \in V$,

$$\int_{\Omega} w^2 dx = \int_{\Omega} (w - f_{\Omega} w)^2 dx \le C(\Omega) \int_{\Omega} |\nabla w|^2 dx,$$
(31)

we deduce that $\int_{\Omega} u^2 dx = 0$, hence u = 0.

We now claim that A is surjective. First, note that a is coercive because of Poincaré-Wirtinger inequality (31). It follows from Lax-Milgram theorem that for all $T \in V^*$ there exists a unique element $u \in V$ such that

$$a(u, v) = \langle T, v \rangle_{V^*, V}$$
 for all $v \in V$.

This follows $\langle T, v \rangle_{V^*, V} = \langle T_u, v \rangle_{V^*, V}$, hence $T = T_u = Au$ so that A is surjective. By a consequence of the open mapping theorem [3, Corollary 2.7, p. 35] theorem, we conclude that A is an isomorphism from V onto V^* . \Box

Corollary 3.4. The first and second derivatives of E can be represented in V^* as:

$$E'(u) = Au - \bar{f}(u) + \oint_{\Omega} \bar{f}(u), \qquad (32)$$

$$L(u)h = Ah - \bar{f}'(u)h + \int_{\Omega} \bar{f}'(u)h, \qquad (33)$$

for all $u, h \in V$.

Proof. Since \bar{f} is bounded, $\bar{f}(u) - \int_{\Omega} \bar{f}(u) \in H \hookrightarrow V^*$. Therefore,

$$Au - \bar{f}(u) + \int_{\Omega} \bar{f}(u) \in V^*.$$

Since

$$\int_{\Omega} \left(\oint_{\Omega} \bar{f}(u) \right) h = \oint_{\Omega} \bar{f}(u) \int_{\Omega} h = 0 \text{ for all } h \in V,$$

it follows that

$$\langle Au - \bar{f}(u) + \int_{\Omega} \bar{f}(u), h \rangle_{V^*, V} = \int_{\Omega} \nabla u \nabla h - \int_{\Omega} \bar{f}(u) h.$$

This togother with (25) implies that

$$E'(u) = Au - \bar{f}(u) + \int_{\Omega} \bar{f}(u)$$

Identity (33) may be proved in a similar way.

We define

$$\mathcal{L}^{p}(\Omega) := \{ u \in L^{p}(\Omega), \int_{\Omega} u(x) \, dx = 0 \},\$$

equipped with the norm $\|\cdot\|_{\mathcal{L}^p(\Omega)} := \|\cdot\|_{L^p(\Omega)}$ and

$$X_p := \{ u \in W^{2,p}(\Omega), \partial_{\nu} u = 0, \int_{\Omega} u(x) dx = 0 \}.$$

Lemma 3.5. Let $p \geq 2$, then for any $g \in \mathcal{L}^p(\Omega)$, there exists a unique solution $u \in X_p$ of the equation

$$Au = g \text{ in } V^*. \tag{34}$$

Moreover, $A = -\Delta$ on X_p .

Proof. It follows from Lemma 3.3 that Equation (34) has a unique solution $u \in V$. We now claim that $u \in X_p$. Consider the elliptic problem

$$\begin{cases} -\Delta \tilde{u} = g & \text{in } \Omega, \\ \partial_{\nu} \tilde{u} = 0 & \text{on } \partial \Omega. \end{cases}$$

First, since $g \in H$, it follows from the Fredholm alternative that this problem possesses a unique solution $\tilde{u} \in V$. Next, since $g \in \mathcal{L}^p(\Omega)$, we deduce from [2] that $\tilde{u} \in W^{2,p}(\Omega)$ so that also $\tilde{u} \in X_p$. In fact, \tilde{u} satisfies Equation (34) since

$$\langle g, v \rangle_{V^*, V} = \langle -\Delta \tilde{u}, v \rangle_{V^*, V} = \int_{\Omega} \nabla \tilde{u} \nabla v \, dx = a(\tilde{u}, v) = \langle A \tilde{u}, v \rangle_{V^*, V}$$

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for all $v \in V$. By the uniqueness of the solution of Equation (34), $u = \tilde{u} \in X_p$. On the other hand, for all $w \in X_p, v \in V$

$$\langle -\Delta w, v \rangle_{V^*, V} = \int_{\Omega} \nabla w \nabla v \, dx = \langle Aw, v \rangle_{V^*, V},$$

so that $A = -\Delta$ on X_p .

Definition 3.6. We say that $\varphi \in V$ is a critical point of E if $E'(\varphi) = 0$.

Lemma 3.7. For every $\varphi \in V$, the following assertions are equivalent:

(i) φ is a critical point of E,

(ii) $\varphi \in X_2$ and φ satisfies the equations

$$-\Delta\varphi - \bar{f}(\varphi) + \int_{\Omega} \bar{f}(\varphi) = 0 \ in \ \Omega, \tag{35}$$

$$\partial_{\nu}\varphi = 0$$
 on $\partial\Omega$. (36)

Moreover, φ is $C^{\infty}(\overline{\Omega})$.

Proof. $(ii) \Rightarrow (i)$. It follows directly from Lemma 3.5 and the formula (32). $(i) \Rightarrow (ii)$. Assume that $\varphi \in V$ is a critical point of E. We deduce from (32) that

$$A(\varphi) = \bar{f}(\varphi) - \int_{\Omega} \bar{f}(\varphi) \quad \text{in } V^*.$$

Since $\bar{f}(\varphi) - \int_{\Omega} \bar{f}(\varphi) \in H$, then $A(\varphi) \in H$. It follows from Lemma 3.5 that $\varphi \in X_2$ and $A = -\Delta$. Therefore φ satisfies (35).

Finally, we deduce that $\varphi \in C^{\infty}(\overline{\Omega})$ from the boundedness of $\overline{f}(\varphi) - \int_{\Omega} \overline{f}(\varphi)$, Sobolev embedding theorem and a standard bootstrap argument.

Lemma 3.8. Let φ be a critical point of E. Then operator $L(\varphi)$ is Fredholm from V to V^* . Moreover,

- (i) ker $L(\varphi)$ is finite-dimensional and contained in $C^{\infty}(\overline{\Omega})$.
- (ii) $\langle u, v \rangle_{V,V^*} = 0$ for all $u \in \ker L(\varphi)$ and $v \in \operatorname{Rg} L(\varphi)$,

(iii) V^* is the topological direct sum of ker $L(\varphi) \subset V \hookrightarrow V^*$ and $\operatorname{Rg} L(\varphi)$,

(iv) if $g \in \mathcal{L}^p(\Omega) \cap \operatorname{Rg} L(\varphi)$ for $p \geq 2$ and $u \in V$ solves the equation

$$L(\varphi)u = g \text{ in } V$$

then $u \in X_p$. Consequently,

$$\operatorname{Rg}(L(\varphi)|_{X_p}) = \operatorname{Rg} L(\varphi) \cap \mathcal{L}^p(\Omega).$$

Proof. We first prove that the linear operator

$$\begin{split} T:V &\longrightarrow V^* \\ h &\longmapsto -\bar{f}'(\varphi)h + \int_\Omega \bar{f}'(\varphi)h \end{split}$$

is compact. Indeed, it follows from the compact embedding $H \hookrightarrow V^*$ and the following estimate

$$\begin{aligned} \|Th\|_{H} &\leq \|\bar{f}'(\varphi)h\|_{L^{2}(\Omega)} + \left\| \oint_{\Omega} \bar{f}'(\varphi)h \right\|_{L^{2}(\Omega)} \\ &\leq C(\|h\|_{L^{2}(\Omega)} + \|h\|_{L^{1}(\Omega)}) \\ &\leq C\|h\|_{V}. \end{aligned}$$

Recall that since A is an isomorphism from V onto V^* , it is also a Fredholm operator of index

 $\operatorname{ind} A := \dim \ker A - \operatorname{codim} \operatorname{Rg} A = 0.$

It follows that $L(\varphi) = A + T$, as a sum of a Fredholm operator and a compact operator, is also a Fredholm operator with the same index [3, p. 168]. Therefore,

 $\operatorname{Rg} L(\varphi)$ is closed in V^* and $\dim \ker L(\varphi) = \operatorname{codim} \operatorname{Rg} L(\varphi) < \infty$. (37)

(i) Using similar arguments as the proof in Lemma 3.7, we deduce that if $h \in \ker L(\varphi)$ then $h \in X_2$ and satisfies the equation:

$$-\Delta h - \bar{f}'(\varphi)h + \int_{\Omega} \bar{f}'(\varphi)h = 0 \text{ in } \Omega,$$

$$\partial_{\nu}h = 0 \quad \text{on } \partial\Omega.$$

Note that $\overline{f}'(\varphi) \in C^{\infty}(\overline{\Omega})$; we deduce that $h \in C^{\infty}(\overline{\Omega})$ from a Sobolev embedding theorem and a bootstrap argument.

(ii) We may identify the linear operator $L(\varphi)$ with a bilinear symmetric form on $V \times V$ (e.g see [29, Section 10.5.3 p. 82]). Thus, for every $u \in \ker L(\varphi), v = L(\varphi)w, w \in V$,

$$\langle u,v\rangle_{V,V^*} = \langle u,L(\varphi)w\rangle_{V,V^*} = \langle L(\varphi)u,w\rangle_{V^*,V} = 0,$$

which implies (ii).

(iii) Using part (ii), we deduce that for every $u \in \ker L(\varphi) \cap \operatorname{Rg} L(\varphi)$, $\langle u, u \rangle_{V^*, V} = 0$, hence u = 0. It follows that $\ker L(\varphi) \cap \operatorname{Rg} L(\varphi) = \{0\}$. On the other hand, dim $\ker L(\varphi) = \operatorname{codim} \operatorname{Rg} L(\varphi)$ so that V^* is the algebraic direct sum of $\ker L(\varphi)$ and $\operatorname{Rg} L(\varphi)$.

Since ker $L(\varphi)$ is finite-dimensional, it is closed in V^* . It follows from (37) that $\operatorname{Rg} L(\varphi)$ is closed in V^* , thus V^* is the topological direct sum of ker $L(\varphi)$ and $\operatorname{Rg} L(\varphi)$.

(iv) Since $g \in \operatorname{Rg} L(\varphi)$, there exists $u \in V$ satisfying

$$Au = \bar{f}'(\varphi)u - \int_{\Omega} \bar{f}'(\varphi)u + g.$$

We write

$$Au = \bar{f}'(\varphi)u - \oint_{\Omega} \bar{f}'(\varphi)u + g \in H,$$

thus $u \in X_2$ and $A = -\Delta$. We have

$$-\Delta u - \bar{f}'(\varphi)u = -\oint_{\Omega} \bar{f}'(\varphi)u + g \in L^p(\Omega),$$

note that $\overline{f}'(\varphi) \in C^{\infty}(\overline{\Omega})$ and use elliptic regularity theory to deduce that $u \in X_p$. From this we obtain (iv).

We equip the norm $\|\cdot\|_{X_p} := \|\cdot\|_{W^{2,p}(\Omega)}$ on X_p .

Lemma 3.9. Assume that $\varphi \in X_p, p > \max\{N, 2\}$ then we can consider E' as a mapping from X_p to $\mathcal{L}^p(\Omega)$ and it may be represented of the following form

$$E'(\varphi) = -\Delta \varphi - \bar{f}(\varphi) + \oint_{\Omega} \bar{f}(\varphi).$$

If, in addition $s_1 \leq \varphi \leq s_2$ then there exists a neighborhood $U(\varphi)$ of φ in X_p such that

$$E'\big|_{U(\varphi)} \mapsto \mathcal{L}^p(\Omega)$$

is analytic.

Remark 3.10. We consider here the standard definition of analyticity (see [28, Definition 8.8, p. 362]):

A mapping T from a Banach space X into a Banach space Y is called analytic on a neighborhood of $z \in X$ if it may be represented as

$$T(z+h) - T(z) = \sum_{k \ge 1} T_k(z)[h, \dots, h] \text{ in } Y,$$

for any $h \in X$, $||h||_X \le \varepsilon$, ε small enough,

where $T_k(z)$ is a symmetric k-linear form on X with values in Y and

$$\sum_{k \ge 1} \|T_k(z)\|_{\mathcal{L}_k(X,Y)} \|h\|^k < \infty.$$

Here, $\mathcal{L}_k(X, Y)$ is the space of bounded k-linear operators $X^k \to Y$.

Proof. If $\varphi \in X_p$, then we can see easily that

$$E'(\varphi) = -\Delta \varphi - \bar{f}(\varphi) + \int_{\Omega} \bar{f}(\varphi),$$

and obviously, $E'(\varphi) \in \mathcal{L}^p(\Omega)$.

Since \bar{f} is a polynomial on $(s_1 - 1, s_2 + 1)$, using Taylor's expansion we have

$$\bar{f}(\varphi(x) + h(x)) - \bar{f}(\varphi(x)) = \sum_{k=1}^{n} \frac{\bar{f}^{(k)}(u(x))}{k!} h^k(x),$$

for all $h \in X_p$ such that $||h||_{C(\overline{\Omega})} \leq C ||h||_{X_p} < 1$. It follows that

$$E'(\varphi+h) - E'(\varphi) = -\Delta h + \sum_{k=1}^{n} \frac{\overline{f}^{(k)}(\varphi)}{k!} h^k - \sum_{k=1}^{n} \int_{\Omega} \frac{\overline{f}^{(k)}(\varphi)}{k!} h^k dx$$
$$= \sum_{k=1}^{n} T_k[h, \dots, h],$$

where

$$T_1[h] = -\Delta h + \bar{f}'(\varphi)h - \int_{\Omega} \bar{f}'(\varphi)h \quad \text{and}$$
$$T_k[h, \dots, h] = \frac{\bar{f}^{(k)}(\varphi)}{k!}h^k - \int_{\Omega} \frac{\bar{f}^{(k)}(\varphi)}{k!}h^k \quad \text{for all } 1 < k \le n$$

Note that for all $1 < k \le n$ and for all $h_1, \ldots, h_k \in X_p$

$$\left| T_k[h_1, \dots, h_k] \right| \leq \left| \frac{\bar{f}^{(k)}(\varphi)}{k!} h_1 \dots h_k \right| + \left| \int_{\Omega} \frac{\bar{f}^{(k)}(\varphi)}{k!} h_1 \dots h_k \right|$$
$$\leq C \prod_{i=1}^k \|h_i\|_{L^{\infty}(\Omega)} \leq C \prod_{i=1}^k \|h_i\|_{X_p},$$

so that

$$\left\| T_k[h_1, .., h_k] \right\|_{\mathcal{L}^p(\Omega)} \le C \prod_{i=1}^k \|h_i\|_{X_p}.$$

This estimate implies that $T_k \in \mathcal{L}_k(X_p, \mathcal{L}^p(\Omega))$ for all $1 < k \leq n$. It is easy to see that $T_1 \in \mathcal{L}(X_p, \mathcal{L}^p(\Omega))$. It follows that E' is analytic in a neighborhood of φ .

Theorem 3.11. (Lojasiewicz inequality). Let $\varphi \in V$ be a critical point of the functional E such that $s_1 \leq \varphi \leq s_2$. Then there exist constants $\theta \in (0, \frac{1}{2}]$ and $C, \sigma > 0$ such that

$$|E(u) - E(\varphi)|^{1-\theta} \le C ||E'(u)||_{V^*}$$

for all $||u - \varphi||_V \leq \sigma$.

In order to study Lojasiewicz inequality of a functional $E \in C^2(V)$ near a critical point φ , Chill gives the following conditions:

(C₁) The kernel $V_0 := \ker L(\varphi)$ is a complemented subspace of V, that means there exists a projection $P \in \mathcal{L}(V)$ such that $V_0 = \operatorname{Rg} P$. We denote by $P' \in \mathcal{L}(V^*)$ the adjoint projection of P.

 (C_2) There exists a Banach space W with the following properties

(i) W is continuously embedded in V^* and invariant under P' (i.e. $P'(W) \subset W$),

- (ii) $E' \in C^1(V, W)$,
- (iii) $\operatorname{Rg} L(\varphi) = \ker P' \cap W.$
- (C₃) There exists Banach spaces $X \subset V$ and $Y \subset W$ such that
 - (i) the spaces X and Y are invariant under projection P and P', respectively,
 - (ii) the restriction of the derivative E' in X is analytic in a neighborhood of φ with value in Y,
- (iii) ker $L(\varphi)$ is contained in X and finite dimensional,
- (iv) $\operatorname{Rg} L(\varphi)|_X = \ker P' \cap Y.$

It follow form [8, Corollary 3.11] that if the conditions (C₁), (C₂), (C₃) hold then there exist $\sigma > 0, \theta \in (0, \frac{1}{2}], C \ge 0$ such that

$$|E(u) - E(\varphi)|^{1-\theta} \le C ||E'(u)||_W \text{ for all } ||u - \varphi||_V \le \sigma.$$

We say that the functional E satisfies the Lojasiewicz inequality near φ and the constant θ will be called the Lojasiewicz exponent.

Our task is to verify the conditions (C_1) , (C_2) (C_3) . We need the following lemma

Lemma 3.12. The conditions $(\mathbf{C_1}), (\mathbf{C_2})$ hold with $W = V^*$ and P the projection onto ker $L(\varphi)$ along $\operatorname{Rg} L(\varphi) \cap V$. In this case, the adjoint projection P' is an extension of P to V^* .

Remark 3.13. Before giving the proof of Lemma 3.12, we recall a useful property in linear algebra. Let T be a linear mapping from a vector space to X into itself. Assume that $T^2 = T$, then

$$X = \ker T \oplus \operatorname{Rg} T.$$

Proof. Because of Lemma 3.8(iii), there exists a projection $Q \in \mathcal{L}(V^*)$ onto $\ker L(\varphi)$ along $\operatorname{Rg} L(\varphi)$. Since V and H are invariant under the projection Q, V and H can be represented as

$$V = \ker L(\varphi) \oplus (\operatorname{Rg} L(\varphi) \cap V) \text{ and } H = \ker L(\varphi) \oplus (\operatorname{Rg} L(\varphi) \cap H)$$
(38)

where the sums are algebraic direct sums. Moreover, the identities (38) are also topological direct sums since their components are closed in the corresponding spaces.

Let P, P_H be the restriction of Q to V, H, respectively. Then, $P \in \mathcal{L}(V), P_H \in \mathcal{L}(H)$. We shall prove that P' = Q.

By Lemma 3.8(ii), ker $L(\varphi)$ is orthogonal to $\operatorname{Rg} L(\varphi) \cap H$ with respect to the scalar product in H. Thus, P_H is an orthogonal projection in H. As a consequence P_H is an auto-adjoint operator of H.

Let $P' \in \mathcal{L}(V^*)$ be the adjoint projection of P. For all $u \in V, v \in H$,

$$\langle u, P'v \rangle_{V,V^*} = \langle Pu, v \rangle_{V,V^*} = \langle P_Hu, v \rangle_{H,H} = \langle u, P_Hv \rangle_{H,H} = \langle u, Qv \rangle_{V,V^*}$$

It follows that P' = Q on H. On the other hand, H is dense in V^* we deduce that P' = Q. Put $W = V^*$ then it is easy to see that the conditions (C_1) and (C_2) hold with the above projections P and P'.

Proof of Theorem 3.11

It remain to verify the condition (C₃). Let $p > \max\{N, 2\}$, choose $X = X_p$ and $Y = \mathcal{L}^p(\Omega)$.

- (i) It follows from Lemma 3.8(i) that ker $L(\varphi) \subset C^{\infty}(\overline{\Omega})$, thus P(V), P(H), P'(V), P'(H)are contained in ker $L(\varphi) \subset C^{\infty}(\overline{\Omega})$, which is a subspace of $X_p, \mathcal{L}^p(\Omega)$.
- (ii) Condition $C_3(ii)$ follows from Lemma 3.9(i).
- (iii) Condition $C_3(iii)$ follows from Lemma 3.8(i).
- (iv) Condition $C_3(iv)$ follows from Lemma 3.8(iv).

These prove theorem. \Box

3.3 Application of Lojasiewicz inequality

Lemma 3.14. Let \bar{u} be the solution of Problem (\bar{P}) corresponding with the initial condition g_0 . Then

(i) For all $0 < s \le t < \infty$,

$$E(\bar{u}(s)) = E(\bar{u}(t)) + \int_{s}^{t} \int_{\Omega} |\bar{u}_{t}|^{2} dx.$$
(39)

Moreover if $g_0 \in V$ we can take s = 0 in (39).

(ii) Further, $e := \lim_{t \to \infty} E(\bar{u}(t))$ exists.

Proof. (i) By Lemma 2.8, \bar{u} is a smooth function on $\overline{\Omega} \times (0, \infty)$ so that

$$\frac{d}{dt}E(\bar{u}(t)) = \int_{\Omega} \left(-\Delta \bar{u} - f(\bar{u})\right)\bar{u}_t$$
$$= \int_{\Omega} \left(-\Delta \bar{u} - f(\bar{u}) + \int_{\Omega} f(\bar{u})\right)\bar{u}_t$$
$$= -\int_{\Omega} \bar{u}_t^2(x,t) \, dx \le 0.$$

As a consequence, for all $0 < s \le t < \infty$

$$E(\bar{u}(s)) = E(\bar{u}(t)) + \int_s^t \int_\Omega |\bar{u}_t|^2 dx.$$

If $g_0 \in V$, using the continuity of E in V, noting that $\bar{u}_t \in L^2(\Omega \times (0, T))$ we deduce that (39) also holds for s = 0.

(ii) We recall that the function F is bounded on \mathbb{R} . Therefore the function $t \to E(\bar{u}(t))$, which is nonincreasing and bounded from below, converges to a limit as $t \to \infty$.

We denote by S(t) the semigroup on H that corresponding with problem (\bar{P}) and define the ω -limit set of g_0 by

 $\omega(g_0) := \{ \varphi \in V : \text{ there exists } \{t_n\} \text{ such that } S(t_n)v_0 \to \varphi \text{ in } V \text{ as } t_n \to \infty \}.$

We have the following result

- **Lemma 3.15.** (i) $\omega(g_0)$ is non-empty, compact of V. Furthermore, $\omega(g_0)$ is positive invariant under S(t), i.e., $S(t)(\omega(g_0)) \subset \omega(g_0)$.
- (ii) The functional E is constant on $\omega(g_0)$. If $\varphi \in \omega(g_0)$ then $s_1 \leq \varphi \leq s_2$ and φ is a critical points of E, i.e., $E'(\varphi) = 0$.
- (iii) $d(S(t)v_0, \omega(g_0)) \to 0 \text{ as } t \to \infty.$

Proof. (i) Since $\{\bar{u}(t), t \geq 1\}$ is compact in V, we can easily show that $\omega(g_0)$ is non-empty, compact of V. Next, note that if $\psi \in \omega(g_0)$, then there exists a sequence $t_n \to \infty$ such that $\psi = \lim_{n \to \infty} S(t_n)g_0$. For all $t \geq 0$, we have

$$S(t)\psi = \lim_{n \to \infty} S(t+t_n)g_0 \in \omega(g_0).$$

This shows that $\omega(g_0)$ is positive invariant.

(ii) First, we prove that E is constant on $\omega(g_0)$. Let $e = \lim_{t\to\infty} E(S(t)g_0)$ as in Lemma 3.14. For any $\varphi \in \omega(g_0)$, we have $\varphi = \lim_{n\to\infty} S(t_n)g_0$ for some sequence $t_n \to \infty$. Since E is continuous in V,

$$E(\varphi) = \lim_{n \to \infty} E(S(t_n)g_0) = e.$$

i.e., E is constant on $\omega(g_0)$.

Note that $S(t_n)v_0 \to \varphi$ in also $L^2(\Omega)$ so that we can extract a subsequence of $S(t_n)v_0$ which converges almost everywhere on Ω . On the other hand $s_1 \leq S(t_n)v_0 \leq s_2$ for all $n \geq 0$, therefore

$$s_1 \le \varphi \le s_2.$$

We now prove that φ is a critical of E. Since $\omega(g_0)$ is positive invariance, $E(S(t)\varphi) = E(\varphi)$ for all $t \ge 0$. It follows from Lemma 3.14 that

$$\int_0^t \int_\Omega |\varphi_t|^2 \, dx dt = 0 \text{ for all } t \ge 0,$$

so that $\varphi_t = 0$ for all $t \ge 0$. Hence, φ satisfies the equation

$$-\Delta \varphi - \bar{f}(\varphi) + \int_{\Omega} \bar{f}(\varphi) = 0 \text{ in } \Omega,$$
$$\partial_{\nu} \varphi = 0 \qquad \text{ on } \partial\Omega.$$

Therefore, φ is a critical point of the functional E by Lemma 3.7. (iii) Assume by contradiction that there exists a sequence $t_n \to \infty$ and $\varepsilon > 0$ such that $d(S(t_n)g_0, \omega(g_0)) \ge \varepsilon$. By compactness, there exists a subsequence $t_{n_k} \to \infty$ such that $S(t_{n_k})g_0 \to w \in \omega(g_0)$. Therefore, $d(S(t_{n_k})g_0, \omega(g_0)) = 0$ as $k \to \infty$, which is absurd.

Proof of Theorem 1.3

We will first apply Lojasiewicz inequality for Problem \overline{P} . Recall that since E is constant on $\omega(g_0)$, as in Lemma 3.15 we can write

$$e = E(v)$$
 for all $v \in \omega(g_0)$. (40)

It follows from Theorem 3.11 that E satisfies the Lojasiewicz inequality in the neighborhood of every $\varphi \in \omega(g_0)$; in other words, we have that for every $\varphi \in \omega(g_0)$ there exist constants $\theta \in (0, \frac{1}{2}], C \geq 0$ and $\delta > 0$ such that

$$|E(v) - E(\varphi)|^{1-\theta} \le C ||E'(v)||_{V^*} \text{ whenever } ||v - \varphi||_V \le \delta.$$
(41)

Since the functional E is continuous on V, we may choose δ small enough so that

$$|E(v) - E(\varphi)| < 1 \text{ whenever } ||v - \varphi||_V \le \delta.$$
(42)

It follows from the compactness of $\omega(g_0)$ in V that there exists a neighborhood \mathcal{U} of $\omega(g_0)$ composed of finitely many balls $B_j, j = 1, ..., J$, with center φ_j and radius δ_j . In each of the ball B_j , inequality (42) and the Lojasiewicz inequality (41) hold for some constants θ_j and C_j . We define $\bar{\theta} = \min \{\theta_j, j = 1, ..., J\}$ and $\bar{C} = \max \{C_j, j = 1, ..., J\}$ to deduce from (40), (41) and (42) that

$$|E(v) - e|^{1-\bar{\theta}} \le \bar{C} ||E'(v)||_{V^*} \text{ for } v \in \mathcal{U}.$$

Using Lemma 3.15(*iii*), there exists $t_0 \ge 0$ such that $\bar{u}(t) \in \mathcal{U}$ for all $t \ge t_0$. Hence, for every $t \ge t_0$, there holds

$$-\frac{d}{dt}|E(\bar{u}(t)) - e|^{\bar{\theta}} = \bar{\theta}|E(\bar{u}(t)) - e|^{\bar{\theta}-1} \left(-\frac{dE}{dt}(\bar{u}(t))\right)$$
$$\geq \frac{\bar{\theta}}{\bar{C}} \frac{\|\bar{u}_t\|_{L^2(\Omega)}^2}{\|E'(\bar{u}(t))\|_{V^*}}.$$
(43)

Note that for all $t \ge t_0, \bar{u}(t) \in C^{\infty}(\overline{\Omega})$, so that $E'(\bar{u}(t)) \in H$ and it can be written of the form

$$E'(\bar{u}(t)) = -\Delta \bar{u} - \bar{f}(\bar{u}) + \int_{\Omega} \bar{f}(\bar{u}) = -\bar{u}_t.$$

Applying continuous embedding $H \hookrightarrow V^*$, we have

$$\|E'(\bar{u}(t))\|_{V^*} \le \bar{\bar{C}} \|E'(\bar{u}(t))\|_{L^2(\Omega)} = \bar{\bar{C}} \|\bar{u}_t\|_{L^2(\Omega)} \text{ for all } t \ge t_0,$$
(44)

where \overline{C} is a positive constant. Combining (43) and (44) we obtain

$$-\frac{d}{dt}|E(\bar{u}(t)) - e|^{\bar{\theta}} \ge C_0 \|\bar{u}_t\|_{L^2(\Omega)}.$$

Here, $C_0 = \frac{\theta}{\bar{C}\bar{\bar{C}}}$. Thus

$$\|\bar{u}(t_1) - \bar{u}(t_2)\|_{L^2} \le \int_{t_1}^{t_2} \|\bar{u}_t\|_{L^2} \le \frac{1}{C_0} (|E(\bar{u}(t_1)) - e|^{\bar{\theta}} - |E(\bar{u}(t_2)) - e|^{\bar{\theta}})$$

for all $t_0 \leq t_1 \leq t_2$. Therefore, $\|\bar{u}(t_1) - \bar{u}(t_2)\|_{L^2(\Omega)}$ tends to zero as $t_1 \to \infty$ so that $\{\bar{u}(t)\}$ is a Cauchy sequence in H. Consequently, there exists $\varphi \in H$ such that $\lim_{t\to\infty} \bar{u}(t) = \varphi$ exists in H. It follows from Corollary 2.9(ii) that $\varphi \in C^{\infty}(\overline{\Omega})$ and that $\bar{u}(t) \to \varphi$ in $C^m(\overline{\Omega})$ for all $m \in \mathbb{N}$. We conclude that

$$\lim_{t \to \infty} \|u(t) - \varphi\|_{C^m(\overline{\Omega})} = 0,$$

and φ is a solution of the stationary problem (S). Moreover, by conserved mass property

$$\int_{\Omega} \varphi \, dx = \int_{\Omega} g_0 \, dx,$$

which completes the proof of Theorem 1.3. \Box

3.4 Rate of the convergence

In this section, we estimate the rate of the convergence of u(t) to φ . The proof is based on Lojasiewicz inequality. We consider two cases, when Lojasiewicz exponent $\theta = \frac{1}{2}$ and $\theta \in (0, \frac{1}{2})$.

3.4.1 When Lojasiewicz exponent $\theta \in (0, \frac{1}{2})$

We need the following lemma.

Lemma 3.16 (see [13], Lemma 3.3). Asume that for all $t \ge t_0$, some $\alpha > 0$ and a constant K > 0

$$\int_t^\infty \|u_t\|_H^2 \le K t^{-2\alpha - 1}$$

.

Then, we have

$$||u(t) - u(\tau)||_H \le \frac{\sqrt{K}}{1 - 2^{-\alpha}} t^{-\alpha} \quad \text{for all } \tau \ge t \ge t_0.$$

Consequently,

$$||u(t) - \varphi||_H \le \frac{\sqrt{K}}{1 - 2^{-\alpha}} t^{-\alpha} \quad \text{for all } t \ge t_0.$$

Proof.

Theorem 3.17. Assume that Theorem 3.11 holds for $\theta \in (0, \frac{1}{2})$, then for $\alpha := \frac{\theta}{1-2\theta} > 0$ and a constant K > 0, we have

$$\|u(t) - \varphi\|_H \le \frac{\sqrt{K}}{1 - 2^{-\alpha}} t^{-\alpha} \quad \text{for all } t > 0.$$

Proof. By the modification on section 3.1, we only need to prove this theorem for function \bar{u} . Since $\bar{u}(x,t)$ is smooth for all t > 0, we have

$$\frac{d}{dt}(E(\bar{u}) - E(\varphi)) = \langle E'(\bar{u}), \bar{u}_t \rangle = -\langle E'(\bar{u}), E'(\bar{u}) \rangle = -\|E'(\bar{u})\|_H^2.$$
(45)

Since $\bar{u}(t)$ tends to φ as $t \to \infty$. Therefore there exists $T_0 > 0$ such that for all $t \ge T_0$

 $\|\bar{u}(t) - \varphi\| \le \sigma \quad (\sigma \text{ in Theorem 3.11}).$

It follows that for all $t \ge T_0$

$$C \| E'(\bar{u}) \|_{V^*} \ge |E(\bar{u}) - E(\varphi)|^{1-\theta}.$$

Therefore, by using the continuous embedding $H \hookrightarrow V^*$, we obtain

$$C_1 \| E'(\bar{u}) \|_H \ge |E(\bar{u}) - E(\varphi)|^{1-\theta} = (E(\bar{u}) - E(\varphi))^{1-\theta}.$$
 (46)

Combining (47) and (48), we get

$$\frac{d}{dt}(E(\bar{u}) - E(\varphi)) \le -C_2(E(\bar{u}) - E(\varphi))^{2(1-\theta)} \text{ for all } t \ge T_0,$$

where $C_2 := 1/C_1^2$.

Note that
$$y(t) := \left((E(\bar{u}(T_0) - E(\varphi))^{2\theta - 1} + C_2(1 - 2\theta)(t - T_0))^{-1/(1 - 2\theta)} \right)^{-1/(1 - 2\theta)}$$

is the unique solution of the differential equation

$$\begin{cases} \frac{d}{dt}y(t) = -C_2 y^{2(1-\theta)} \text{ for } t \ge T_0, \\ y(T_0) = E(\bar{u}(T_0) - E(\varphi)). \end{cases}$$

We use a differential inequality in [22, Lemma 2.7, p.53] to deduce for all $t \ge T_0$ that

$$\begin{split} E(\bar{u}(t)) - E(\varphi) &\leq \left((E(\bar{u}(T_0) - E(\varphi))^{2\theta - 1} + C_2(1 - 2\theta)(t - T_0) \right)^{-1/(1 - 2\theta)} \\ &= \left((E(\bar{u}(T_0) - E(\varphi))^{2\theta - 1} - C_2(1 - 2\theta)T_0 + C_2(1 - 2\theta)t \right)^{-1/(1 - 2\theta)} \\ &= \left((E(\bar{u}(T_0) - E(\varphi))^{2\theta - 1} - C_2(1 - 2\theta)T_0 + C_2(1 - 2\theta)\frac{t}{2} + C_2(1 - 2\theta)\frac{t}{2} \right)^{-1/(1 - 2\theta)} \\ &\leq \left(C_2(1 - 2\theta)\frac{t}{2} \right)^{-1/(1 - 2\theta)} \text{ for all } t \geq \overline{T_0}, \text{ with some } \overline{T_0} > T_0 \text{ large enough.} \end{split}$$

It follows that for all $t \ge \overline{T_0}$

$$\int_t^\infty \|\bar{u}_t(s)\|^2 \, ds \le K t^{-2\alpha - 1}$$

Here, $K := \left(\frac{C_2(1-2\theta)}{2}\right)^{-1/(1-2\theta)}$ and $\alpha := \frac{\theta}{1-2\theta} > 0$. Now, according to Lemma 3.16, we obtain

$$\|\bar{u}(t) - \varphi\|_H \le \frac{\sqrt{K}}{1 - 2^{-\alpha}} t^{-\alpha} \quad \text{for all } t \ge \overline{T_0}.$$

3.4.2 When Lojasiewicz exponent $\theta = \frac{1}{2}$

Lemma 3.18 (see [12], Lemma 2.2). Assume that there exists two constants $\gamma > 0$ and a > 0 such that for all $t \in [0, T]$,

$$\int_t^{+\infty} \|u_t\|_H^2 \, ds \le a \exp(-\gamma t).$$

Then for all $\tau \geq t \geq 0$, we have

$$\|u(t) - u(\tau)\|_H \le \sqrt{ab} \exp(-\frac{\gamma t}{2}).$$

Theorem 3.19. Assume that Theorem 3.11 holds for $\theta \in (0, \frac{1}{2})$, then for constants $K, \delta > 0$, we have

$$||u(t) - \varphi||_H \le K \exp(-\delta t) \quad for \ all \ t > 0.$$

Proof. By the modification on section 3.1, we only need to prove this theorem for function \bar{u} . Since $\bar{u}(x,t)$ is smooth for all t > 0, we have

$$\frac{d}{dt}(E(\bar{u}) - E(\varphi)) = \langle E'(\bar{u}), \bar{u}_t \rangle = -\langle E'(\bar{u}), E'(\bar{u}) \rangle = -\|E'(\bar{u})\|_H^2.$$
(47)

Since $\bar{u}(t)$ tends to φ as $t \to \infty$. Therefore there exists $T_0 > 0$ such that for all $t \ge T_0$

 $\|\bar{u}(t) - \varphi\| \le \sigma \quad (\sigma \text{ in Theorem 3.11}).$

It follows that for all $t \ge T_0$

$$C \| E'(\bar{u}) \|_{V^*} \ge |E(\bar{u}) - E(\varphi)|^{\frac{1}{2}}.$$

Therefore, by using the continuous embedding $H \hookrightarrow V^*$, we obtain

$$C_1 \|E'(\bar{u})\|_H \ge |E(\bar{u}) - E(\varphi)|^{\frac{1}{2}} = (E(\bar{u}) - E(\varphi))^{\frac{1}{2}}.$$
(48)

Combining (47) and (48), we get

$$\frac{d}{dt}(E(\bar{u}) - E(\varphi)) \le -C_2(E(\bar{u}) - E(\varphi)) \text{ for all } t \ge T_0.$$

where $C_2 := 1/C_1^2$.

Note that $y(t) := \left(E(\bar{u}(T_0) - E(\varphi)) \right) \exp(-C_2(t - T_0))$ is the unique solution of the differential equation

$$\begin{cases} \frac{d}{dt}y(t) = -C_2 y \text{ for } t \ge T_0, \\ y(T_0) = E(\bar{u}(T_0) - E(\varphi)). \end{cases}$$

We use a differential inequality in [22, Lemma 2.7, p.53] to deduce for all $t \ge T_0$ that

$$E(\bar{u}(t)) - E(\varphi) \le \left(E(\bar{u}(T_0) - E(\varphi)\right) \exp(-C_2(t - T_0))$$

In particular for all $t \geq T_0$,

$$\int_{t}^{\infty} \|\bar{u}_{t}(s)\|_{H}^{2} ds \leq \left(E(\bar{u}(T_{0}) - E(\varphi)) \right) \exp(-C_{2}(t - T_{0})).$$

Now, Using Lemma 3.18, we deduce the result of the theorem.

4 Stationary solution in one dimension

We consider Problem (P) in one dimension.

$$(P) \begin{cases} u_t = u_{xx} + f(u) - \int_{\Omega} f(u) & \text{in } \Omega \times \mathbb{R}^+, \\ u_x(x,t) = 0 & \text{on } \partial\Omega \times \mathbb{R}^+, \\ u(x,0) = g_0(x) & x \in \Omega. \end{cases}$$

Theorem 4.1. Set N = 1. Let a < b be such that satisfy Property (C) with respect to f. We also assume that $f'(s) \leq 0$ for all $s \in [a, b]$. Assume that $a < g_0(x) < b$ for almost $x \in \Omega$, then

$$\varphi = \int_{\Omega} g_0.$$

Proof. Firs, note that we have

$$a \le u(x,t) \le b$$
 for all $x \in \Omega, t \ge 0$.

It follows that $a \leq \varphi \leq b$. Set $w := \varphi_x$, differentiating the equations of φ in stationary problem (S) with respect to x, we obtain

$$(P_1) \begin{cases} w_{xx} + f'(\varphi)w = 0 & \text{in } \Omega, \\ w = 0 & \text{on } \partial\Omega. \end{cases}$$

Since $c(x) := f'(\varphi) \leq 0$, we apply maximum principle to deduce that $w \equiv 0$ is the unique solution of Problem (P_1) . Consequently, φ is constant. This together mass conservation property implies that

$$\varphi = \oint_{\Omega} g_0.$$

5 Appendix

Lemma 5.1. Let T be a linear mapping from V to V such that $T^2 = T$. then

(i)

$$V = \ker T \oplus \operatorname{Rg} T.$$

(ii) Let W be a subspace of V. Assume that W is invariant i.e. $T(W) \subset W$. Then

 $W = (\ker T \cap W) \oplus (\operatorname{Rg} T \cap W).$

Proof. (ii) First, note that we have

$$T_{|W}: W \longrightarrow W$$

satisfies $(T_{|W})^2 = T_{|W}$. It follows that

$$W = \ker \left(T_{|W} \right) \oplus \operatorname{Rg} \left(T_{|W} \right).$$

Since

$$\ker (T_{|W}) = \{ y \in W : \quad T(x) = 0 \} = \{ x \in W \text{ and } x \in \ker T \}$$
$$= \ker T \cap W,$$

and since

$$\operatorname{Rg}(T_{|W}) = T(W) \subset \operatorname{Rg} T \cap W,$$

then

$$W = \ker (T_{|W}) \oplus \operatorname{Rg} (T_{|W}) \subset (\ker T \cap W) \oplus (\operatorname{Rg} T \cap W).$$

But,

$$W \supset (\ker T \cap W) \oplus (\operatorname{Rg} T \cap W).$$

It follows that

$$W = (\ker T \cap W) \oplus (\operatorname{Rg} T \cap W).$$

References

- [1] R.A. Adams, *Sobolev spaces*, Academic Press, 1975.
- [2] S. Agmon, A. Douglis, L. Nirenberg, Estimates near the boundary for solutions of elliptic partial differential equations satisfying general boundary conditions. I Comm. Pure Appl. Math. 12, (1959), 623–727.
- [3] H. Brezis Functional analysis, Sobolev Spaces and Partial Differential Equations Spinger, (2010).
- [4] D. Brochet, D. Hilhorst and X. Chen, *Finite-dimensional exponential attractor for the phase field model*, Appl. Anal. 49 (1993), 197–212.
- [5] T. Cazenave, A. Haraux An Introduction to Semilinear Evolution Equations Oxford Science Publications, (1998)
- [6] R. Chill, The Lojasiewicz-Simon gradient inequality on Hilbert spaces, Proceedings of the 5th European-Maghrebian Workshop on Semigroup Theory, Evolution Equations, and Applications (M. A. Jendoubi, ed.), 2006, 25–36.
- [7] R. Chill, E. Fasangova, J. Pruss, Convergence to steady states of solutions of the Cahn-Hilliard and Caginal equations with boundary conditions, Math. Nachr. 279 (2006), 1448–1462.
- [8] R. Chill, On the Lojasiewicz-Simon gradient inequality, J. Funct. Anal. 201 (2003), 572-601.
- C.M. Elliott and H. Garcke, Existence results for diffusive surface motion laws, Adv. Math. Sci. Appl. 7 (1997), 467–490.
- [10] L.C. Evans, Partial differential equations, Graduate Studies in Mathematics, 19. American Mathematical Society, Providence, RI, 1998.
- [11] A. Haraux, Systèmes Dynamiques Dissipatifs et Applications, Recherches en Mathématiques Appliquées, Masson, Paris, 1991.
- [12] A. Haraux, M.A. Jendoubi, On the convergence of global and bounded solutions of some evolution equations, J. evol. equ. 7 (2007), 449–470.

- [13] A. Haraux, M.A. Jendoubi, O. Kavian, Rate of decay to equilibrium in some semilinear parabolic equations, J. evol. equ. 3 (2003), 463–484.
- [14] D. Henry, *Geometric Theory of Semilinear Parabolic Equations*, Springer, Berlin (1981).
- [15] S. Huang, P. Takáč Convergence in gradient-like systems which are asymptotically autonomous and analytic, Nonlinear Anal. Ser. A: Theory Methods 46 (2001) 675-698.
- [16] O. Kavian, Introduction à la Théorie des Points Critiques et Applications aux Problèmes Elliptiques, Spinger-Verlag, 1993
- [17] O.A. Ladyzenskaya, V.A. Solonnikov and N.N. Uraltseva, *Linear and quasi-linear equations of parabolic type*, Translations of Mathematical Monographs, vol 23, Providence R I. American Mathematical Society (1968).
- [18] S. Lojasiewicz, *Ensemble semi-analytique*. *I.H.E.S.*, Bures-sur-Yvette, 1965.
- [19] S. Lojasiewicz, Une propriété topologique des sous-ensembles analytiques reéls, Colloques Internationaux du C.N.R.S, 117, Les Equations aux Derivées Partielles (1963), 87–89.
- [20] S. Lojasiewicz, Sur la géométrie semi- et sous-analytique, Ann. Inst. Fourier (Grenoble), 43, (1993), 1575–1595.
- [21] M. Marion, Attractors for reaction-diffusion equations: existence and estimate of their dimension, Appl. Anal. 25 (1987), 10–147.
- [22] J.C. Robinson, Infinite-Dimensional Dynamical Systems, Cambridge University Press, Massachusetts, 2001.
- [23] H. Royden *Real analysis* Third edidion
- [24] J. Rubinstein and P. Sternberg, Nonlocal reaction-diffusion equations and nucleation, IMA J. of Appl. Math. 48 (1992), 249–264.
- [25] P. Rybka, K.-H. Hoffmann, Convergence of solutions to Cahn-Hilliard equation, Commun. PDE. 24, (1999), 1055–1077.

- [26] L. Simon, Asymptotics for a class of non-linear evolution equations, with applications to geometric problems, Ann. of Math. 118, (1983), 525–571.
- [27] R. Temam, Infinite-dimensional dynamical systems in mechanics and physics, Applied Mathematical Sciences, 68, Springer-Verlag, New York, 1988.
- [28] E. Zeidler, Nonlinear Functional Analysis and Its Applications Vol I, Spinger-Verlag
- [29] V. Zorich Mathematical analysis II, Spinger (2008).