

A new method for the in vivo identification of mechanical properties in arteries from cine MRI images: theoretical framework and validation.

Alexandre Franquet, Stéphane Avril, Rodolphe Le Riche, Pierre Badel, Fabien Schneider, Zhi-Yong Li, Chrisitan Boissier, Jean-Pierre Favre

► To cite this version:

Alexandre Franquet, Stéphane Avril, Rodolphe Le Riche, Pierre Badel, Fabien Schneider, et al.. A new method for the in vivo identification of mechanical properties in arteries from cine MRI images: theoretical framework and validation.. IEEE Transactions on Medical Imaging, 2013, 32 (8), pp.1448 - 1461. 10.1109/TMI.2013.2257828 . hal-00805124

HAL Id: hal-00805124 https://hal.science/hal-00805124

Submitted on 27 Mar 2013

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

A new method for the *in vivo* identification of mechanical properties in arteries from cine MRI images: theoretical framework and validation.

Alexandre Franquet, Stéphane Avril, Rodolphe Le Riche, Pierre Badel, Fa bien C Schneider, Zhi Yong Li, Christian Boissier and Jean Pierre Favre

5

q

10

19

 A. Franquet, S. Avril and P. Badel are with the CIS-EMSE, CNRS UMR 5146, Ecole Nationale Supérieure des Mines, Saint-Etienne F-42023, France. e-mail: avril@emse.fr

2. R. Le Riche is with CNRS UMR 6158 and the Ecole nationale Supérieure des Mines, Saint-Etienne F-42023

113. Z.Y. Li is with the School of Biological Science and Medical Engineering, Southeast University, Nanjing12210096, China

4. F.C. Schneider, C. Boissier and J.P. Favre are with the Radiology Department, Vascular Department, and
 Cardiovascular surgery Department, respectively, University Hospital of Saint Etienne France and GRT PAR UR4672 Jean Monnet University, Saint Etienne F-42000, France.

This work has been submitted to the IEEE for possible publication. Copy right may be transferred without notice, after which this version may no
 longer be accessible

Abstract

Quantifying the stiffness properties of soft tissues is essential for the di-20 agnosis of many cardiovascular diseases such as atherosclerosis. In these 21 pathologies it is widely agreed that the arterial wall stiffness is an indicator 22 of vulnerability. The present paper focuses on the carotid artery and pro-23 poses a new inversion methodology for deriving the stiffness properties of 24 the wall from cine-MRI (Magnetic Resonance Imaging) data. We address 25 this problem by setting-up a cost function defined as the distance between 26 the modeled pixel signals and the measured ones. Minimizing this cost 27

function yields the unknown stiffness properties of both the arterial wall and the surrounding tissues. The sensitivity of the identified properties to various sources of uncertainty is studied. Validation of the method is performed on a rubber phantom. The elastic modulus identified using the developed methodology lies within a mean error of 9.6 %. It is then applied to two young healthy subjects as a proof of practical feasibility, with identified values of 625 kPa and 587 kPa for one of the carotid of each subject.

Keywords Finite elements, mechanical properties, elasticity, inverse method,
 identification, artery, *in vivo* analysis.

37 **1** Introduction

Quantifying the stiffness properties of soft tissues is essential for the diagnosis of many diseases. This is especially important for cardiovascular pathologies where it has been shown that the stiffness of the arterial wall is an indicator of the stroke risk for diseases such as atherosclerosis [1]. This could be used to improve the diagnosis in clinical practice.

The mechanical properties of arteries have been a research topic of major attention for thirty years ([2, 3, 4, 5, 6, 7, 8, 9, 10]). Many authors have characterized the mechanical properties of excised arteries using either tensile tests, indentation tests or inflation tests. However the results reported in the literature are not very consistent. The large scattering observed can be attributed to the differences in experimental conditions and conservation processes, to the anisotropic and behavioral non linearities ([5, 11, 12]), or to the natural interindividual variability.

Therefore, the identification of the mechanical properties in vivo provides 51 another essential point of view. In vivo identification techniques are more con-52 straining because they are usually non-intrusive. An overview of the existing 53 methods and results of identification is presented in Table 1 which will be thor-54 oughly discussed in the "Discussion". In clinical applications, for assessing ar-55 terial stiffness, it is still commonplace to use the Moens-Korteweg theoretical 56 equation relating the Pulse Wave Velocity and the incremental modulus in a 57 round cylinder [13]. Recent developments based on the propagation of shear 58 waves in soft tissues have enabled clinicians to improve the evaluation of the arterial stiffness [14]. Although these types of methods are well suited for rapid 60

medical examinations, they suffer, among other things, from a poor precision
 because of the strong assumptions made about the mechanical behavior of ar teries.

In order to improve precision, medical imaging techniques can be employed. The basic method consists in tracking the diameter changes in the images and relating these to blood pressure variations ([15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 14]). The advantage is that a non-linear behavior can possibly be identified [26] however the precision remains questionable due to the assumptions about the geometry of the artery.

Imaging techniques providing access to kinematic fields (displacements,
strains) offer an attractive flexibility and enable the identification of more complex mechanical behaviors with a limited number of mechanical assumptions.
However, they require the development of appropriate inverse methods [27].
The most natural one is the Finite Element Model Updating (FEMU) method
which has been applied several times to data acquired *in vivo*.

The FEMU method consists in finding the parameters of a Finite Element (FE) model that minimize the gap between numerical and experimental kine-77 matic fields (displacements, strains). This method has been successfully ap-78 plied to identify the heterogeneous elastic properties of a pig artery with an atherosclerotic plaque ([28, 29]). The elastic properties and the contours of an 80 atherosclerotic plaque in the human coronary artery were also determined in a 81 similar way using a Parametric FEMU method (P-FEMU) [30]. Using the same idea, other authors [31] have designed a specific algorithm where the contours 83 are recovered by a watershed method and the heterogeneous elastic properties 84 are found at the same time. 85

All these applications of the FEMU method are based on experimental strain fields obtained by IntraVascular UltraSounds (IVUS). IVUS is an ultrasonic imaging technique where the ultrasound probe is inserted inside the arterial conduit using a catheter. Specific algorithms have been developed for deriving radial strains from the Radio Frequency (RF) signals with a suitable spatial resolution (pixel size ~0.1 mm). But the IVUS technique is very invasive due to the use of catheters thereby limiting its applications.

Other non-invasive imaging techniques are available for mapping the strain fields but their spatial resolution is inferior. For instance, strain maps can be measured *in vivo* in the vicinity of the carotid artery using a specific Mag⁹⁶ netic Resonance Imaging (MRI) sequence: the DENSE (Displacement ENcod-

⁹⁷ ing with Stimulated Echoes) sequence [32]. The resulting pixel size is over
⁹⁸ 0.5 mm and moreover requires highly skilled operators.

Despite the great potential of MRI for these types of problems, the pixel
 size is usually not fine enough to track the strain fields of the arterial wall, lim iting the use of MRI for the purpose of accurately identifying the mechanical
 behavior of arterial walls.

In this paper we address this issue by setting-up a model relating the tem-103 poral variations of the pixel MR signal to the strains and deformations of the 10 tissue. Using this approach, we are able to combine the advantages of a full 105 field based identification method with a non invasive imaging device. A novel 106 inverse method is developed called Magnitude-Based FEMU (MB-FEMU). The 107 cost function is defined as the distance between the modeled pixel signals and 108 the measured ones. Minimizing this cost function yields the unknown stiffness 109 properties of the arterial wall and possibly of the surrounding tissues. This 110 type of method has been developed in the context of elastography ([33, 34]) for 111 assessing hard nodules in breast cancers. However the issue of the low spatial 112 resolution of MRI compared with the size of the specimen may prevent from 113 using this method on arteries. In this study we used the partial volume effect 114 that occurs in MRI, which corresponds to the partial filling of the pixels by 115 different materials. 116

After presenting the theoretical and numerical foundations of the method, the method is validated on a rubber phantom and then applied to two subjects as a proof of practical feasibility.

120 2 Methodology

The approach is based on the minimization of a cost function which measures 121 the distance between a template image and a registered image. The template 122 image is the experimental image of a deformed artery provided by the magni-123 tude of a cine Phase-Contrast MR sequence. The registered image is obtained 124 from the experimental image of the undeformed artery and a displacement 125 field provided by a FE analysis. After introducing the theoretical framework, 126 the implementation of the method is described and then the experimental set 127 up used to validate the method is presented. 128

Team	Year	Artery	Imaging system	Number of subjects	Methodology/index	Value (kPa)
Riley et al.	1992	Common carotid	B-mode echo	419	Variation of diameter	701 ± 324
Laurent et al.	1994	Radial artery	RF US signal	22	Variation of diameter / E _{inc}	2680 ± 1810
Brands et al.	1999	Common carotid	RF US signal	1	Variation of diameter / E _{inc}	480
Aggoun et al.	2000	Common carotid	B-mode echo	21	Variation of diameter / E _{inc}	187 ± 67
Bussy et al.	2000	Common carotid	RF US signal	40	Variation of diameter / E _{inc}	710 ± 290
Boutouyrie et al.	2001	Common carotid	RF US signal	15	Variation of diameter / E _{inc}	431 ± 243
Selzer et al.	2001	Common carotid	B-mode echo	24	Variation of diameter	630 ± 230
Kanai et al.	2003	Common carotid	RF US signal	2	Variation of diameter.	960
Boutouyrie et al.	2004	Common carotid	RF US signal	16	Variation of diameter / E _{inc}	230 ± 110
Hasegawa and Kanai	2004	Common carotid	RF US signal	1	Variation of diameter	1400
Baldewsing et al.	2005	Coronary artery	IVUS	1	Strain based FEMU	188
Taviani et al.	2008	Common carotid	MRI	3	Lumen's area based FEMU	250
Avril et al.	2009	Common carotid	MRI	1	Pulse wave Analysis	99
Couade et al.	2010	Common carotid	RF US signal	1	Shear waves propagation	402

Table 1: Literature summary of methods, imaging systems and Young's moduli of arteries identified *in vivo*.

129 2.1 Theoretical framework

130 2.1.1 Kinematics principles

¹³¹ In continuum mechanics a body can be viewed as a set of elements called ma-

terial points that occupy regions in space. The reference and current configura-

¹³³ tions of a body refer respectively to the undeformed or initial configuration of

 $_{^{134}}$ the body, denoted $\mathcal{C}_0,$ and to the deformed or current configuration, denoted

 \mathcal{C}_t , after the application of forces and moments. The position of a material point

¹³⁶ in the initial configuration C_0 is written <u>X</u> while the position of a material point

¹³⁷ in the current configuration C_t is written <u>x</u>. <u>x</u> is a function of the coordinates of

the material points and of time such as $\underline{x} = \Phi_{\theta}(\underline{X}) = \underline{X} + \underline{U}(\underline{X})$. Similarly it is possible to write $\underline{X} = \Phi_{\theta}^{-1}(\underline{x}) = \underline{x} - \underline{U}(\underline{X})$ where the transformation Φ_{θ} is a function of the material properties θ and of U the displacement. It is possible to introduce the function u which will be used in this paper so that \underline{X} is a function of \underline{x} only:

$$\underline{X} = \underline{x} - \underline{U}(\underline{X})
= \underline{x} - \underline{U}\left(\Phi_{\theta}^{-1}(\underline{x})\right)
= \underline{x} - \underline{u}(\underline{x})$$
(1)

143 2.1.2 Signal magnitude in MRI

In medical imaging, contrast results from materials which are different in nature. A cine-MRI sequence allows to measure the structure's response at different times of the cardiac cycle. We assumed that the magnitude at position \underline{x} of the signal in the deformed configuration C_t is identical to the magnitude of the signal at position $\Phi_{\theta}^{-1}(\underline{x})$ in the initial configuration C_0 . This is justified by the fact that \underline{x} and $\Phi_{\theta}^{-1}(\underline{x})$ are the positions of the same material particle in two different configurations. This means that:

$$I(\underline{x}) = I_0\left(\Phi_{\theta}^{-1}(\underline{x})\right)$$
(2)

where *I* is the magnitude in the current configuration C_t and I_0 is the magnitude in the initial configuration C_0 . An implementation of Φ_{θ} will be proposed later in the present article based on FE displacement fields.

154 2.1.3 Identification of material properties

Equation (2) is verified if the transformation Φ_{θ} corresponds to the real experimental transformation. Part of the inverse method is based on the evaluation of a cost function which calculates the distance between a measured quantity and an estimation of this quantity which is a function of unknown parameters. In our case the transformation Φ_{θ} of the structure is unknown. It is affected by the material parameters which can therefore be identified by minimizing the following distance or cost function J_2 :

$$J_{2}(\theta) = \frac{1}{2} \iint_{\Omega} \left(I(\underline{x}) - I_{0}\left(\Phi_{\theta}^{-1}(\underline{x})\right) \right)^{2} d\underline{x}$$
(3)

¹⁶² The domain Ω corresponds to the imaged body. I_0 and I are experimental im-¹⁶³ ages measured by the imaging system in the configurations C_0 and C_t , respec-¹⁶⁴ tively. An important feature of our method is that the distance J_2 is directly ¹⁶⁵ expressed in terms of image magnitudes, as opposed to displacements, strains, ¹⁶⁶ stresses, or energies.

¹⁶⁷ Minimizing Equation (3) on θ gives the material properties for which the ¹⁶⁸ image associated to the transformation Φ_{θ} is the closest to the measured image.

¹⁶⁹ 2.2 Implementation

The proposed MB-FEMU methodology is an instance of the Finite Element 170 Model Updating method. This method can be described by the association 171 of: (i) experimental data, (ii) a numerical model based on finite elements and 172 (iii) a cost function which measures the distance between (i) and (ii). MRI was 173 chosen for providing the experimental data. The numerical model consists 174 in computing the image of the deformed geometry from displacement fields 175 which were estimated by a FE computation. The cost function is an elabora-176 tion of Equation (3). These three aspects are now further detailed. Later in the 177 article, the method is applied in vitro to a rubber phantom and in vivo to two 178 healthy subjects. 17

180 2.2.1 Experimental data

A Phase Contrast sequence was used to provide cine-MRI data [35]. A bipolar 181 magnetic gradient is applied to the imaged body during the excitation phase 182 which means that a first gradient is applied along one direction followed by 183 the application of a second gradient in the opposite direction during the same 184 time frame. The moving protons and the stationary protons are respectively 185 dephased and rephased by the second gradient. The phase information of this 186 sequence is often used to obtain the flow velocity in the out-of-plane direction. 187 In our case however, only the magnitude of the signal is considered since it is 18 mainly correlated to the T_1 intrinsic property of the materials. 189

190 2.2.2 Numerical model

¹⁹¹ Finite element modeling was chosen to model the structure (phantom and ar-¹⁹² teries). 2D plane strain was assumed inasmuch as out-of-plane displacements were small compared to in-plane displacements and MR slice thickness. Quasi
static conditions were used as heart beats have a low frequency (approximately
1 Hz). In the loading interval, only the tangent behavior of both the phantom and the arteries was considered: linear elasticity in the finite deformations
framework was assumed for all materials.

¹⁹⁸ 2.2.3 Cost function description

The FEMU process iteratively adjusts the mechanical parameters of a model based on its mechanical response as obtained by solving FE models until convergence is reached. The iterations are created by the minimization of a distance with respect to the parameters:

$$min_{\theta}J_{2}(\theta)$$
 (4)

In practice the spatial resolution of MRI depends on the acquisition time, which should be small enough to satisfy the constraints of a clinical implementation, and on the number of \underline{k} sets which can be measured. Equation (3) is then a discrete sum with N terms (= number of pixels). In order to save computation time we assumed that the pixels were undeformable and that the transformation of a pixel is a translation only. As a consequence, the inverse transformation Φ_{θ}^{-1} can be simplified for a pixel centered at coordinates ξ_k as:

$$\Phi_{\theta}^{-1}\left(\underline{\xi}_{\underline{k}}\right) \approx \underline{\xi}_{\underline{k}} - \underline{u}_{\theta}\left(\underline{\xi}_{\underline{k}}\right)$$
(5)

where u_{θ} is the displacement in the current configuration C_t (see notation introduced in Equation (1)) which depends on the mechanical properties θ .

Fig. 1 describes the calculation of the signal magnitude provided by the pixel *k* which, following the above assumptions, is a discrete sum:

$$I_0\left(\Phi_{\theta}^{-1}\left(\underline{\xi_k}\right)\right) \approx I_0\left(\underline{\xi_k} - \underline{u_{\theta}}\left(\underline{\xi_k}\right)\right) = \sum_{j=1}^4 \alpha_{kj}\left(\theta\right) \cdot I_0\left(\underline{\xi_{kj}}\right) \tag{6}$$

As illustrated in Fig. 1, k1, ..., k4 are the pixels which, in the current configuration, intersect the initial image $I_0(\underline{\xi}_{kj})$. $\alpha_{k1}, ..., \alpha_{k4}$ are the normalized intersection areas.

The whole image can be reconstructed in this way. However, in order to provide more sensitivity to the method, a restrained set of pixels called "control



Figure 1: Illustration of the principle applied to generate the current image from the reference image and the computed displacements $u_{\theta}\left(\frac{\xi_k}{k}\right)$ in order to recover the coefficients α_{kj} for each pixel. In this example, the magnitude of the pixel in the middle of the deformed image is calculated using the coefficients $\{\alpha_{k1}, \alpha_{k2}, \alpha_{k3}, \alpha_{k4}\} = \{50\%, 20\%, 10\%, 20\%\}$ (see Equation 6).

pixels" was selected. Control pixels were located in areas where a large change
in signal magnitude is likely to be observed i.e. at the lumen contours. A binary
mask is defined to select the control pixels (see Fig. 2a) so that N is less than
the total number of pixels. The cost function is normalized by the number of
pixels considered.

²²⁴ Finally, the cost function is written as:

$$J_{2}(\theta) = \frac{1}{2 \cdot N} \sum_{k=1}^{N} \left(I(\underline{\xi}_{k}) - I_{0} \left(\Phi_{\theta}^{-1} \left(\underline{\xi}_{k} \right) \right) \right)^{2}$$

$$= \frac{1}{2 \cdot N} \sum_{k=1}^{N} \left(I(\underline{\xi}_{k}) - \sum_{j=1}^{4} \alpha_{kj} \left(\theta \right) \cdot I_{0} \left(\underline{\xi}_{kj} \right) \right)^{2}$$
(7)

The minimization of Equation (7) is an iterative process where different 225 combinations of material parameters θ were tested in a FE model. The updated 226 vector of parameters ($\theta \leftarrow \theta + \Delta \theta$) was determined by a Levenberg-Marquardt 227 algorithm with bounds handling [36]. It requires gradients of the computed in-228 tensity $I_0\left(\Phi_{\theta}^{-1}\left(\xi_k\right)\right)$ with respect to the parameters θ which were estimated 22 by backward finite differences (or forward differences at lower bounds). The 230 identification process stopped when one of the following termination criteria 231 was satisfied: $J_2(\theta) \le \varepsilon_a = 10^{-7}$ (accuracy on J_2 was reached) or $\|\Delta\theta\| \le \varepsilon_b =$ 232 10^{-3} (the step size was too small, no more improvement was expected.). We 233



Figure 2: Example of images processing during the phantom study. A mask is defined to select the control pixels.

- ²³⁴ previously showed the importance of setting these two criteria ([37]).
- ²³⁵ The method is summarized in the flow chart of Fig. 3.

236 2.3 Validation of the MB-FEMU method

237 2.3.1 Experimental subjects

Phantom description An arterial rubber phantom was used to perform *in vitro* experimentation. The phantom was a 50-cm-long tube with a 7 mm inner diameter and 10 mm outer diameter made of PVA gel (1 freeze-thaw-cycle), surrounded by gadolinium enriched water mimicking the surrounding tissue (see Fig. 4). A peristaltic pump generated a pulsating flow mimicking physio-logical pressure while the wall motion was imaged.

Patients description Two healthy human subjects were selected for the experiment. They were 25-year-old and 24-year-old healthy subjects (185 cm/85 kg
and 183 cm/80 kg, respectively). Their common carotid arteries were imaged
using MRI.



Figure 3: Flow chart of the MB-FEMU methodology to identify mechanical and other parameters. From the cine Phase Contrast MRI images, initial (diastole) and deformed (systole) images were chosen. A FE model is created from the initial image. Then for a given set of initial parameters θ_0 a FE model is computed to provide a displacement field $U_{\theta}(\underline{X})$. This field is applied to the initial image $I_0\left(\underline{\zeta}_k\right)$ according to the method described in the text in order to generate a registered image, $I_0\left(\Phi_{\theta}^{-1}\left(\underline{\zeta}_k\right)\right)$. The cost function J_2 is the distance between this registered image and the experimental deformed image, calculated for a given region of interest (binary mask: control pixels). A Levenberg-Marquardt with bounds handling is used to minimize the cost function J_2 with respect to the parameters θ . The procedure stops if the desired accuracy on J_2 is reached (ε_a) or if the step size is too small (ε_b) to expect significant further improvements of the cost function J_2 .

248 2.3.2 Reference data

249 Measurement technique A 3T Siemens Verio system MRI scanner was used

²⁵⁰ for the measurements. A 2D spin-echo FLASH sequence was used to record a



(a) The 500 mm long PVA phantom was surrounded by water. A maximum pressure of 70 mmHg was applied to the tube's lumen.



(b) Experimental set-up (Diagram not to scale).



single 3 mm thick slice with a matrix size of 256×256 giving in-plane dimen-251 sions 0.27 mm \times 0.27 mm for the *in vitro* experiments and 0.586 mm \times 0.586 mm 252 for the in vivo study. The image acquisition was synchronized either with the 253 signal command for the phantom or with the physiological signal for the sub-25 jects. Regarding the physiological signal, infrared spectroscopy was used to 255 measure the blood flow in one of the subject's finger. We obtained $N_t = 50$ and 25 $N_t = 43$ snapshots for the phantom and for the subjects respectively, evenly 257 distributed throughout the cardiac cycle. It corresponds to a mean sampling 258 frequency of 50 Hz and 40 Hz, respectively. The magnitude of the signal was 259 digitized with a 12 bits resolution (integer between 0 and 4 095). 260

Choice of experimental images The raw magnitude image was a 3D array of size $256 \times 256 \times N_t$, denoted $I_n^{raw} \left(\underline{\xi}_l\right)$ for l varying from 0 to 65 535, and nvarying from 0 to $N_t - 1$. The images were filtered along the time axis using a Gaussian filter with a kernel size $\sigma = 1.5$ to produce $I_n \left(\underline{\xi}_k\right)$.

²⁶⁵ A region of interest of 63 × 63 pixels was selected around the vessel. The ²⁶⁶ first image $I_0\left(\underline{\xi}_k\right)$ with k varying from 0 to 3 968 was chosen as the initial ²⁶⁷ image. The deformed image was selected through the images $I_n\left(\underline{\xi}_k\right)$ for $1 \le$ ²⁶⁸ $n \le N_t - 1$ according to the following definition:

$$\max\left(\overline{\mid I_n\left(\underline{\xi}_{\underline{k}}\right) - I_0\left(\underline{\xi}_{\underline{k}}\right)\mid}\right) \tag{8}$$

An example of experimental images is shown in Fig. 5.



269

(a) Raw global image. The yellow square is a region of interest of size 63×63 pixels.



(b) Region of interest around the common carotid artery. The jugular vein touches the carotid artery. The vertebral artery can be observed at the bottom left.



ing filter. The kernel size (standard deviation) is 1.5.

Figure 5: Example of a magnitude image obtained by PC MRI. The raw image is filtered after having chosen a region of interest.

Choice of control pixels A binary mask was defined before the identification
procedure to discriminate the control pixels (see Fig. 2): (i) pixels which contain
the moving contours were selected automatically. (ii) Additional control pixels
were added by dilating the previous binary mask according to:

$$M_{(ii)} = M_{(i)} \oplus S \tag{9}$$

where $M_{(ii)}$ is the final binary mask used; $M_{(i)}$ is the binary mask at step (i); and *S* is a square structuring element of size 3×3 with $S_{ij} = 1$. The choice of a

²⁷⁶ limited set of pixels allows to improve the sensitivity of the cost function.

277 2.3.3 Description of the FE models

The 2D FE geometry was either derived from the known geometry of the phantom or segmented from PC images. The contours of the artery were extracted in order to separate the arterial wall from the surrounding tissue and then separately estimate the elastic moduli of each. Quadratic triangular elements were used for the mesh generation (we previously showed that this type of elements provides accurate displacement fields for this type of problems [37]). A uniform pressure was applied to the phantom or artery inner walls. The implicit solver of Abaqus© [38] was used.

Phantom A 2D cross section of the phantom was modeled. It was made up
 of 1 100 finite elements surrounded by approximately 5 000 similar elements
 mimicking the fluid.

The boundary conditions were enforced in the experiment. A maximal pressure of 9.33 kPa (70 mmHg) was applied on the phantom's inner wall. The free edges of the surrounding box were fixed in the FE model (see Fig. 4).

Linear elastic material properties were used for the phantom. The Poisson's ratio was set to 0.49 while the elastic modulus has to be recovered. The external water had no mechanical influence because the external pressure applied on the phantom was constant during the test. It was nonetheless incorporated into the model in order to provide the necessary external displacement fields. Linear elasticity was considered with a Poisson's ratio of 0 and an elastic modulus of 10^{-6} kPa.

Subjects The FE geometry of each subject (see Fig. 6) was derived from the 299 PC MR images. The contours of the neck, trachea, vertebral column and jugu-300 lar vein were segmented manually from the global undeformed image (see 30 Fig. 5a). The inner wall of the artery was recovered automatically using an al-302 gorithm based on watershed and Fourier polynomial descriptors: (i) gradients 303 of the mean magnitude along the cardiac cycle were calculated using finite 304 differences; (ii) watershed sources were chosen in the artery's lumen, in the 305 vein's lumen and in the surrounding tissue; (iii) "contour pixels" were then deduced; (iv) the final contour was the approximation of these pixels' centers 307 with a first degree Fourier descriptor. The thickness of the artery was mea-308 sured by echography: the IMT (Intima-Media Thickness) was $e_A = 0.45$ mm 309





artery

(b) Subject B (left commor carotid).



(d) Zoom on the carotid artery

Figure 6: FE meshes were made of triangular quadratic plane strain elements. The arteries were composed of approximately 1 000 elements while the surrounding tissues were composed of approximately 7 000 elements. A uniform pressure was applied to the inner lumen of the carotid artery. Vertical symmetry was assumed. Vertebral column was fixed (see bold black lines).

- and $e_B = 0.43$ mm for each subject respectively. The outer contour of the artery
- was deduced from the inner contour and from the IMT.

The artery was represented using approximately 1 000 elements and the surrounding tissue with about 7 000 elements. Both the artery and the surrounding tissue were assumed linear elastic in the range of loading. Their Poisson's ratio was set to 0.49 as arteries were quasi incompressible. The elastic modulus of the artery was unknown and had to be identified while the elastic modulus of the surrounding tissue was set to 30 kPa. The influence of this
modulus on the identified elastic modulus of the artery is studied herein (see
Section "Results").

Regarding the boundary conditions, sagittal symmetry conditions were im-320 posed for each model (horizontal symmetry). The vertebral column was as-321 sumed undeformable and the corresponding contour was fixed. Other con-322 tours were left free. The mean measured pulse pressure (differential pres-323 sure between diastolic and systolic states), measured before and after the exam 32 with a digital sphygmomanometer, was applied to the artery's lumen ($\Delta P_A =$ 325 11.33 \pm 0.13 kPa with $P_A^{diastole} =$ 72 mmHg and $P_A^{systole} =$ 157 mmHg, and 326 $\Delta P_B = 6.67 \pm 0.13$ kPa with $P_B^{diastole} = 84$ mmHg and $P_B^{systole} = 134$ mmHg). 327 As a first approximation, we assumed that the blood pressure in the carotid 328 was equal to the blood pressure in the brachial artery. The values of these 329 pulse pressures were discussed below. The reference state was set to the dias-330 tolic state. 331

332 2.3.4 Definition of the parametric study

A parametric study was carried out to evaluate the sensitivity of the described methodology on its internal parameters. The uniqueness of the solution was verified by changing the initial elastic modulus. The influences of (i) the choice of control pixels, (ii) potential errors on the estimation of physical parameters (mechanical properties of the surrounded tissues or measured pressure) and (iii) potential errors on the geometry of the artery were studied. The reference configuration is described in Table 2.

340 **3 Results**

- ³⁴¹ In this section we focus on the final identified elastic moduli and on the values
- $_{_{342}}$ of the cost function at the optima. Results are given as mean \pm standard deviation.
- A full identification procedure requires approximately 5 min (8 Levenberg-
- ³⁴⁴ Marquardt iterations).

	Phantom	Subject A	Subject B
Initial elastic modulus of the vessel	300 kPa	500 kPa	500 kPa
Choice of control pixels	automatic ($N = 704$)	automatic $(N^A = 130)$	automatic $(N^B = 120)$
Smoothing filter	temporal Gaussian with $\sigma = 1.5$	temporal Gaussian with $\sigma = 1.5$	temporal Gaussian with $\sigma = 1.5$
Pressure applied	9.33 kPa	11.33 kPa	6.666 kPa
Elastic modulus of the surrounding tis- sue	0 kPa	30 kPa	30 kPa
Thickness of the vessel considered	1.5 mm	0.45 mm	0.43 mm

Table 2: Reference test configuration. Tests were conducted by modifying one parameter at a time.

³⁴⁵ 3.1 Validation of the MB-FEMU method with the phantom

346 3.1.1 Validation of the value of the identified properties

A rubber phantom was used to validate the method. Its elastic properties in the range of loading were determined by an uni-axial tensile test (elastic modulus E = 250 kPa and Poisson's ratio $\nu = 0.49$). 32 tests combining different choices of control pixels, different initial elastic moduli and different smoothing parameters were conducted. The details of these results are presented herein. The mean identified elastic modulus is E = 274 kPa ± 19.8 which corresponds to an average error $\varepsilon = 9.6$ %.

354 3.1.2 Uniqueness of solution

Our identification is based on solving the minimization problem (see Equation 355 4) with the Levenberg-Marquardt algorithm which is efficient for least squares 356 formulations but which may get trapped at local minima of I_2 . In order to 357 study the uniqueness of the solution to the minimization problem, five differ-358 ent initial values of the elastic modulus between 150 and 350 kPa have been 359 input to the identification algorithm. Results are $E_{mean} = 270$ kPa ± 1.09 and 360 $J_2^{mean} = 3.34 \cdot 10^4 \pm 2.64$. Note that the standard deviations are very low which 36 means that the identification is marginally affected by a change of the initial 362 value of the elastic modulus. 363

364 3.1.3 Effect of the choice of control pixels

The choice of the pixels considered in the cost function may have an influence 365 on the identified elastic modulus. The control pixels can be either selected 366 automatically using the procedure described previously or chosen manually 367 by defining a binary mask which contains the vessel. In this section, 15 dif-368 ferent binary masks were tested. For each one of them, a different number 369 of control pixels was assigned (from 704 to 1320 control pixels). Results are 370 $E_{mean} = 280 \text{ kPa} \pm 4.93$ which means that the variations of the binary mask 371 don't affect the identified modulus. Note nevertheless that the best identified 372 elastic modulus occurs when the control pixels are chosen automatically (704 373 control pixels; E = 271 kPa). The value of the cost function at optimum tends 374 to slightly decreases with the increase of the number of control pixels (results 375 are $J_2^{mean} = 2.67 \cdot 10^4 \pm 0.23 \cdot 10^4$). However there is no correlation between 376 this decrease of values and a change on the identified moduli. 377

378 3.1.4 Effect of the smoothing filter

We also studied the effect of filtering the data before the identification proce-379 dure (see Fig. 7a and 7b). Two types of filters were tested: temporal filters 38 (Gaussian filter as described previously with various kernel sizes), and spatial 381 filters (Wiener and median with different kernel sizes). The overall results ac-382 counting for both filters are $E_{mean} = 268$ kPa ± 31.09 and $J_2^{mean} = 3.27 \cdot 10^4 \pm$ 383 $0.88 \cdot 10^4$. The temporal filter tends to increase the estimated elastic modulus 384 (Fig. 7a). The elastic modulus ranges from E = 265 kPa to E = 341 kPa while 385 the value of the cost function decreases linearly with the increase of kernel size 386 σ (Fig. 7b). Median spatial filtering yields moduli close to the tensile test ref-387 erence value (kernels 3×3 and 5×5 : E = 256 kPa and E = 254 kPa). Wiener 388 spatial filtering underestimates the elastic modulus (kernels 3×3 and 5×5 : 380 E = 231 kPa and E = 217 kPa).

391 3.1.5 Effect of measurement errors in the pressure

Errors of plus or minus 15 % were considered for the nominal pressure *P* in the phantom (see Fig. 8a). Results show that the value of the cost function slightly decreases with the increase of pressure. Nevertheless, the standard deviation is very low which makes the cost function almost unaffected ($J_2^{mean} =$



(a) Identified elastic modulus for the phantom.



(b) Values of the cost function at the optima for the phantom.



(c) Identified elastic modulus for subject A.

Figure 7: Effect of the raw images filtering. The filter was applied to both the initial and the deformed images. Final identified elastic moduli (squares, left axis) and the values of the cost function at the optima (diamonds, right axis) are shown for the phantom and the two subjects.

³⁹⁶ $3.34 \cdot 10^4 \pm 3.74$). The identified elastic modulus ranges from $E_{min} = 231$ kPa ³⁹⁷ to $E_{max} = 307$ kPa which corresponds to errors of -7.6% to +22.8%. There ³⁹⁸ is a clear linear relationship between the pressure and the identified modulus



(d) Values of the cost function at the optima for subject A.



(e) Identified elastic modulus for subject B.



(f) Values of the cost function at the optima for subject B.

Figure 7: (continued with last three figures)

(r = 0.99).



Figure 8: Effect of the applied pressure in the FE model for the phantom and both subjects. Squares and diamonds correspond to the identified elastic moduli (left axis) and to the values of the cost function at the optima (right axis), respectively.

400 3.2 Healthy subjects study

⁴⁰¹ For the subjects, the same numerical experiments were conducted for the vari-

ations of the identification parameters and for the effect of the estimation of
 pressure on the identified modulus.

404 3.2.1 Identification parameter variations

Uniqueness of solution The initial elastic modulus of the artery was changed six times to study the uniqueness of the solution in the range 400 to 900 kPa. Results are $E^A_{mean} = 625$ kPa ± 0.27 and $E^B_{mean} = 586$ kPa ± 3.25 . The value of the cost function at the optimum is unaffected: $J^{mean}_{2B} = 166.38 \pm 5.10^{-4}$ and $J^{mean}_{2B} = 187.48 \pm 0.10$. As well as the results on the phantom, the identified elastic modulus is independent of the initial elastic modulus.

Effect of the choice of control pixels The choice of different sets of control pixels doesn't affect the identified elastic modulus (results are $E_{mean}^{A} =$ 630 kPa ± 13.50 and $E_{mean}^{B} =$ 589 kPa ± 9.66 for 15 sets of control pixels which contain from 101/89 to 220/167 pixels for A/B, respectively). In each case the identified elastic moduli are similar while the values of the cost function decrease with the increase of the number of control pixels. This observation is similar to what we found with the phantom.

Effect of the smoothing filter Results are summarized in Fig. 7c, 7d, 7e and 418 7f. The average results are $E_{mean}^A = 602$ kPa \pm 31.45 and $E_{mean}^B = 598$ kPa \pm 419 33.38. The identified moduli range from 555 kPa to 664 kPa for subject A and 420 from 547 kPa to 670 kPa for subject B. The minimum and maximum values 421 were obtained in both cases for the 3×3 median filter and for the 2.5 Gaussian 422 filter, respectively. The value of the cost function at optimum decreases with 423 the augmentation of the "strength" of the filtering although the comparison 424 between the temporal and the spatial filters is not easy. However the filtering 425 of images always reduces the value of the cost function at optimum compared 426 to using the raw images. 427

428 3.2.2 Effect of the estimation of physical parameters

Effect of measurement errors in the pressure Pulse blood pressure (differ-420 ence between maximal and minimal pressures) was measured with a digital 430 sphygmomanometer ($\Delta P_A = 11.33$ kPa and $\Delta P_B = 6.67$ kPa). Similarly to the 431 phantom validation, a variation of plus or minus 15 % were considered for the 432 nominal pressure ΔP applied in the FE model. Results are detailed in Fig. 8b 433 and 8c. Identified elastic moduli range from $E_{min}^A = 517$ kPa to $E_{max}^A = 734$ kPa 434 for subject A and from $E_{min}^B = 487$ kPa to $E_{max}^B = 691$ kPa for subject B. The 43 relationship between the pressure and the identified modulus is linear which 436 confirms the results from the phantom. Nevertheless and contrary to the phan-437 tom study, the cost function value at optimum tends to increase linearly with a 438 higher pressure in both subjects. 439

Effect of the estimation of the mechanical properties of the surrounding tis-440 sue The artery was assumed to be surrounded by a soft elastic media. Its 441 elastic properties were supposed to be known. The choice of its elastic prop-442 erties may have an influence on the identification of the artery's stiffness. For 443 instance, in Fig. 9, we show how the choice of the elastic modulus of the sur-444 rounding tissue affects the identification of the elastic modulus of the artery. 445 As the surrounding tissue stiffness is fixed to increasingly higher values, the 446 identified elastic modulus of the artery decreases. The values of the cost func-447 tion exhibits a second order polynomial shape with a minimum reached for 448 $E_{surrounding}^A \approx 30$ kPa and $E_{surrounding}^B \approx 150$ kPa. This shows that the cost func-449 tion is sensitive to the stiffness of the surrounding tissue. The modification of 450 the elasticity of the surrounding tissue corresponds somehow to a modification 451 of the pressure applied at the outer boundary of the artery. Previous results on 452 the effect of the pressure shows a linear relationship between the identified 453 moduli or the value of the cost function at optimum and the inner pressure 454 applied. Here the correlation between the elastic modulus of the surrounding 455 tissue and the identified artery's modulus is linear in both cases ($r^A = 0.98$ 456 and $r^B = 0.94$). For the cost function however, the relationship with the elastic 457 properties of the surrounding tissue seems to be a quadratic form. 458



Figure 9: Influence of the elastic modulus of the surrounding tissue on the identified elastic modulus of the artery (squares, left axis) and on the value of the cost function value at the optima (diamonds, right axis).

459 3.2.3 Effects of errors in the geometry

Errors in the artery's position The FE model geometry was defined from the undeformed image. Potential errors in the positioning of the artery were considered (see Fig. 10). No tendency appears with the X-axis or Y-axis offsets neither for the identified moduli nor for the values of the cost function at the optima.

Effect of measurement errors in the artery's diameter The inner lumen of the artery is determined automatically (see section "Methodology"). This section aims at studying the effect of a wrong estimation of the inner diameter due to the algorithm described previously. Errors of plus or minus 10 % of the nominal diameter were considered (see Fig. 11). For subject B, the diameter of the artery could not be increased further because it was already in contact



(a) Error of the position of the artery along the X-axis for subject A.



(b) Error of the position of the artery along the Y-axis for subject A.

Figure 10: Study of potential errors in the positioning of the artery in the FE model. Squares represent the final identified elastic moduli (left axis) and diamonds the values of the cost function at the optima (right axis).

with the jugular vein. The resulting elastic moduli identified lie within the following ranges respectively for subjects A and B : $E_{range}^{A} = [553 \text{ kPa}, 659 \text{ kPa}]$ and $E_{range}^{B} = [487 \text{ kPa}, 587 \text{ kPa}]$. The identified moduli increase with diameter. The value of the cost function at the optimum decreases when the diameter is reduced.

Effect of measurement errors in the artery's thickness The thickness of the artery was measured with ultrasounds. Errors of plus or minus 10 % in the thickness of the artery were considered (see Fig. 12). The elastic moduli ranges resulting from the identification process are respectively $E_{range}^{A} = [566 \text{ kPa}, 696 \text{ kPa}]$ and $E_{range}^{B} = [533 \text{ kPa}, 660 \text{ kPa}]$. The correlation between the identified modulus and the thickness of the artery is linear, with a decrease of the identified



(c) Error of the position of the artery along the X-axis for subject B.



(d) Error of the position of the artery along the Y-axis for subject B.



modulus when the thickness increases. The same effect can be observed for the
value of the cost function at optimum, although the range of variation is very
low.

485 4 Discussion

The discussion is organized in two parts. The first part focuses on the robustness of the MB-FEMU method whereas the second part discusses the elastic
properties of arteries obtained.

489 4.1 Robustness of the MB-FEMU method

⁴⁹⁰ In the three cases studied in the present work (the phantom and the two sub-

₄₉₁ jects) the robustness of the MB-FEMU method was tested with respect to dif-



Figure 11: Effect of the measurement error of the artery's diameter. Squares represents the elastic moduli (left axis) and diamonds the values of the cost function at the optima (right axis).

ferent sources of uncertainty and artifacts. The normalized input parameters were noted X_i . Assuming that the relationship between the identified elastic modulus and parameter X_i is linear, it is possible to rank the influence of each parameter by calculating the following normalized regression coefficient:

$$a_{i} = \left| \frac{Cov\left(X_{i}, E\right)}{Var\left(X_{i}\right)} \right| \tag{10}$$

where *E* is the identified elastic modulus and X_i is the normalized parameter *i* ($X_i \in [0,1]$). The higher the coefficient a_i , the higher the influence of the parameter on the identified variable. Results are shown in Table 3.

The sources of uncertainty and artifacts can be sorted into three groups according to their relative influence:

The most important parameter is the pressure applied on the vein wall
 (i) by the surrounding soft tissue on the outer surface, (ii) and by the



Figure 12: Effect of the measurement error of the artery's thickness. External diameter remains constant in all cases and only the internal diameter was modified. The identified elastic moduli (left axis) and the values of the cost function at the optima (right axis) are represented as squares and diamonds, respectively.

blood circulation on the inner surface. It is particularly obvious for sub-503 ject B (see Fig. 9b). The estimation of the stiffness of the surrounding 504 tissue and of the pressure applied is crucial to accurately identify the 505 elastic modulus of the artery. It can be observed in Fig. 9a and 9b that 506 the value of the cost function evolves when the stiffness of the surround-507 ing tissue is increased in the model. A minimum is reached for the values 508 $E_{surrounding}^A \approx 30$ kPa and $E_{surrounding}^B \approx 150$ kPa. This result suggests that 500 the identification process can be used to determine the elastic moduli of 510 both the artery and the surrounding tissue simultaneously. In this case 511 the uniqueness of the solution would not be guaranteed and the compu-512 tation time would increase. The asymmetric displacement of the artery 513 between the configurations (see Fig. 13) suggests the presence of a stiffer 514

tissue on the right hand side of the artery (tendons for instance).

515

2. Less important but not negligible are the sources of uncertainty that may 516 arise from the geometry definition. On one hand, we observe that the 517 error made on the artery position does not significantly significantly af-518 fect the identified elastic modulus. This constitutes one of the strength 519 of the MB-FEMU methodology. On the other hand, the diameter and the 520 thickness have a moderate but significant influence and thus, must be 521 characterized accurately. It is possible to define a parametric model as 522 in [30] to recover simultaneously both the elastic modulus and the ge-523 ometry of the artery (in this case the uniqueness of the solution would 524 not be guaranteed and the computation time would probably increase). 525 Note that the entire thickness of the artery is not taken into account here 526 (only the Intima-Media Thickness is considered) but that extra informa-527 tion on the total thickness could help reduce this uncertainty as the total 528 wall thickness can be measured by ultrasounds or MRI with a Turbo Spin 529 Echo flash sequence (Blackblood-FATSAT) [39]. 530

3. The least important sources of uncertainty results from the so-called iden-531 tification parameters: (i) position and number of control pixels, (ii) initial 532 elastic modulus. The induced errors can therefore be neglected as com-533 pared to those induced by the other sources. Identification results (very 534 close one to the other) obtained from different sets of initial parameters 535 are a good hint that the solution of the problem is very likely to be unique. 536 Only the kernel sizes of the filters must be chosen with care because an 537 increase of the kernel size decreases the apparent deformation of the im-538 ages between the initial and deformed configurations (see Fig. 7). 539

540 4.2 Identified elastic moduli

Arterial compliance has been widely studied over the last decades. Stiffness
values found in the literature are widespread for two main reasons: (i) subjects
doing the experiment are different (important inter-individual differences), (ii)
methods are different. The review of [40] on the calculation of arterial stiffness
indexes in clinical practice mentioned nine different stiffness indexes from the
"elastic modulus" to the "stiffness index". Recently [41] listed existing meth-



Figure 13: Difference between experimental deformed and initial images for subject B. A white crown appears around the artery. The asymmetry of displacement of the artery can be due to the high stiffness of tissues at the right hand side (tendons, etc.).

Table 3: Normalized regression coefficients a_i , measuring the influence of the different parameters on the identified elastic modulus of the artery, are calculated with Equation (10).

	Subject A	Subject B
E surrounding tissue	339.87	487.76
Pressure	192.64	181.41
Thickness	111.31	108.76
Diameter	90.64	76.96
Control pixels	43.61	26.40
Translation dX	11.85	44.48
Translation dY	8.40	10.18
Initial elastic modulus	0.39	2.84

ods which are used in clinical practice to assess arterial stiffness. The same
author defined in [16] the incremental modulus of an artery such as:

$$E_{inc} = \left(3\left(1 + \frac{R_d^2}{(R_d + IMT)^2 - R_d^2}\right)\right) \cdot \frac{R_d^2 \cdot \Delta P}{R_s^2 - R_d^2}$$
(11)

where R_s and R_d are respectively the systolic (highest pressure) and diastolic (lowest pressure) radii of the artery; *IMT* is the Intima-Media Thickness; ΔP is the pulse pressure between diastole and systole. This formula was determined for an isotropic, incompressible, semi closed thick cylinder with an inner pressure applied.

554 4.2.1 Comparison with ultrasounds measurements

The MB-FEMU method presented in this paper yields elastic moduli of 625 kPa 555 and 587 kPa when applied to two different young healthy subjects. In ad-55 dition an echography was conducted on these subjects in order to acquire 557 movies of their carotid arteries during three heart beats. Initial and deformed 55 images were extracted manually for both subjects (see Fig. 14). Pulse blood 559 pressures were measured during the exam with a digital sphygmomanome-560 ter: $\Delta P^A = 8.27$ kPa and $\Delta P^B = 6.40$ kPa. Arteries lumens were segmented 56 manually on the initial and deformed ultrasounds images and the correspond-562 ing incremental modulus was calculated with Equation (11). The implemen-563 tation of the MB-FEMU method was based on the assumption of plane strain 564 in a cross section of the artery which is different from the assumptions used 565 for deriving Equation (11). In order to make some equivalence between the methods, initial and deformed areas (and subsequently radii) were segmented 567 directly from the results of the FE computation at the optima for calculating 568 E_{inc} using Equation (11). The pressures used in the case of the FE model were 569 $\{\Delta P_A, \Delta P_B\} = \{11.33 \text{ kPa}, 6.666 \text{ kPa}\}.$ Results are $E_{inc}^{echo} = 374 \text{ kPa}; E_{inc}^{FE} =$ 570 496 kPa and $E_{inc}^{echo} = 434$ kPa; $E_{inc}^{FE} = 509$ kPa for the subjects A and B, respec-571 tively. 572

It can be noted that the results obtained with both imaging systems are 573 consistent with the values reported in the literature (see below) though small 57 differences can be noted. The differences between the two imaging systems 575 can be justified by three main reasons. First, the boundary conditions enforced 576 in each model are different because the ultrasound probe applied a pressure 577 on the neck during the exam. Secondly, the height on the neck at which the 578 ultrasound images were taken is not managed and is potentially different from the height of the MRI. Third, the blood pressure of the patients were sensibly 580 different during the MRI and the echography exams ($\Delta P_A^{echo} = 8.27$ kPa vs 581 $\Delta P_A^{MRI} = 11.33 \text{ kPa}; \Delta P_B^{echo} = 6.40 \text{ kPa vs } \Delta P_B^{MRI} = 6.666 \text{ kPa}$). This potentially could impact the behavior of the arterial wall which is non-linear. Dif-583 ferences in pulse pressures between MRI and echography as well as higher 584 pulse pressure in MRI were observed several times for some subjects. It can 585 be explained by some psychological stress induced by the noisy and confined 586 environment. 587

588 4.2.2 Comparison with the literature



(c) Initial image for subject B.

from movies taken during three heart beats.



Figure 14: Ultrasounds image (echography) of the carotid arteries of both subjects. Initial (diastole) and deformed (systole) images were extracted manually

Artery's stiffness A survey of different in vivo elastic moduli identified by 589 other authors is reported in Table 1. Different imaging systems were used: B-590 mode echography (see Fig. 14), Radio-Frequency UltraSounds signal, IntraVas-591 cular UltraSounds or MRI. The reported elastic moduli are in good agreement 592 with our results in most of the studies except for: (i) [16] ($E = 2680 \pm 1810$ kPa) 593 where the difference may be explained by the location of the artery (radial 59 artery at the wrist instead of carotid artery); (ii) [25] (E = 1400 kPa) where 595 the elastic modulus depends on the excitation frequency of the ultrasounds be-596 cause a viscoelastic behavior is considered; (iii) [42] (E = 188 kPa) where a dif-597 ferent artery was also studied (pathological coronary artery); (iv) [18] ($E_{inc} =$ 598

⁵⁹⁹ 187 ± 67 kPa) where the subjects were significantly younger (9-year-old); (v) ⁶⁰⁰ [13] (E = 99 kPa) where the authors suggest that the small elastic modulus ⁶⁰¹ found may be related to a bad estimation of the thickness.

MB-FEMU method vs other methods The methods used in the literature for 602 quantifying the artery's stiffness in vivo are generally based on the variation 603 of the artery's diameter during heart beats. In these cases authors generally 60 used theoretical formula by considering several mechanical and geometrical 605 assumptions: [16, 17, 18, 19, 20, 24] calculated the incremental modulus according to Equation (11); [15, 21] applied the Laplace law by considering a 2D plane 607 stress artery: [22, 25] derived an analytical formula of the circumferential elas-608 tic modulus using the assumptions that the artery is incompressible, isotropic and that the arterial pressure decreases linearly from the artery's inner diam-610 eter to adventitia). Other studies derived the artery's stiffness from velocity 611 Phase Contrast MR images and Moens-Korteweg equation which relates the 612 stiffness of a round artery to the pulse wave velocity and a round artery's stiff-613 ness [13] or studied the propagation of shear waves induced by an ultrasonic 614 device to deduce shear modulus (denoted μ) and the elastic modulus (denoted 615 *E*) according to $E = 3 \cdot u$ [14]. 616

Few studies measured the stiffness of the human carotid arteries by taking 617 into account the surrounding tissue and the environment of the artery (verte-618 bral column, tendons, etc.). [42] used a strain-based FEMU method to recover 619 both the contours and the elastic moduli of the constituents of an atheroscle-620 rotic plaque. This method is very invasive since the IVUS technique was em-621 ployed to obtain strain fields in the wall of the artery. [43] used another FEMU method with a cost function defined as a difference of areas but the areas were 623 segmented manually from Phase Contrast images which may be critical with 624 regard to the pixel size. 625

Our method and these alternative FEMU methods have in common the advantage that they are not based on strong mechanical assumptions. Indeed the arteries are known to behave highly non-linearly ([5, 44, 45, 12]). Only the tangent (or incremental) behavior is considered for the moment by linearizing the stress/strain relationship of arteries between diastole and systole. The identified elastic modulus is therefore a function of the diastolic and pulse pressures due to the intrinsic non-linear behavior of arteries. More complex constitutive equations could potentially be identified by including into the FE model anisotropy, viscoelastic effects, heterogeneity, etc. The limit is the question of the identifiability and uniqueness of the solution that should be investigated for each other constitutive equations. Compared to alternative FEMU methods reported in the literature, the MB-FEMU method has also the advantage of being compatible with non-invasive imaging techniques such as MRI. Indeed, clinical MRI was used and the impact on the patient is very limited (standard and routinely-used exam for a patient before carotid endarterectomy).

Close to our method, the "hyper elastic warping" is a technique that was developed a few years ago ([46, 47, 48]). The objective is to perform image registration in order to quantify the kinematics of the transformation within a deforming body. Our approach is similar in the sense that we perform image registration as well, with the difference that our main objective is not to quantify the kinematics of the deformations but to identify the constitutive properties of the deforming body.

5 Conclusions

This study shows that identifying elastic moduli in common carotid arteries in 649 vivo from cine-MRI images is possible. An original method named the Mag-650 nitude Based Finite Element Model Updating (MB-FEMU) methodology has 65 been thoroughly described, validated on a rubber phantom and applied to two 652 healthy subjects. In our implementation MRI and Phase Contrast magnitude 653 images were used to provide the experimental data. The phase information of the Phase Contrast sequence could be used for improving the boundary con-655 ditions in the FE model as a perspective. We would like to bring the reader's 656 attention to the fact that the MB-FEMU can be easily extended to other imaging 657 techniques. The identified elastic moduli on healthy subjects are in good agree-658 ment with those found in the literature, especially if we consider the differences 659 in hypotheses made that we discussed in the article. The MB-FEMU method 660 offers several advantages compared to existing methods: it makes a maximum 661 use of experimental data since the FE results are used only as a perturbation of 662 the experimental acquisitions; it is possible to extract arterial stiffness from a 663 routine MRI exam on any patient; it can readily be extended to more complex 664 tissue constitutive behaviors. Heterogeneous geometries can be possibly iden-665

tified if an accurate FE model can be created, although the uniqueness of the
minimization problem will have to be verified. In the future this method will
be employed to study the modification of the mechanical behavior of arteries
with age and cardiovascular disease. Applications for the diagnosis of vulnerable atherosclerotic plaques in the carotid is probably the most promising
challenge for this method.

672 Acknowledgment

⁶⁷³ We are grateful to the two volunteers for their time and implication.

- 674 We would like to thank IEEE Transactions on Medical Imaging (please visit
- 675 http://ieeexplore.ieee.org/xpl/RecentIssue.jsp?punumber=42)

676 References

- [1] B. A. Kingwell, T. K. Wadell, T. L. Medley, J. D. Cameron, and A. M. Dart,
 "Large artery stiffness predicts ischemic threshold in patients with coronary artery disease," *Journal of the American College of Cardiology*, vol. 40,
 no. 4, pp. 773–779, 2002.
- [2] R. T. Lee, A. J. Grodzinsky, E. H. Frank, R. D. Kamm, and F. J. Schoen,
 "Structure-dependent dynamic mechanical behavior of fibrous caps from human atherosclerotic plaques," *Circulation*, vol. 83, no. 5, pp. 1764–1770,
 1991.
- [3] A. Duprey, K. Khanafer, M. Schlicht, S. Avril, D. Williams, and R. Berguer,
 "In vitro characterisation of physiological and maximum elastic modulus of ascending thoracic aortic aneurysms using uniaxial tensile testing,"
 European Journal of Vascular and Endovascular Surgery, vol. 39, no. 6, pp.
 700–707, 2010.
- [4] D. Beattie, C. Xu, R. Vito, S. Glagov, and M. C. Whang, "Mechanical anal ysis of heterogeneous, atherosclerotic human aorta," *Journal of Biomechan- ical Engineering*, vol. 120, no. 5, pp. 602–607, 1998.
- [5] Y. C. Fung, *Biomechanics : Mechanical properties of living tissues*, 2nd ed.
 New-York: Springer-Verlag, 1993.

695 696 697	[6]	G. A. Holzapfel, G. Sommer, and P. Regitnig, "Anisotropic Mechanical Properties of Tissue Components in Human Atherosclerotic Plaques," <i>Journal of Biomechanical Engineering</i> , vol. 126, no. 5, pp. 657–665, 2004.
698 699 700 701	[7]	C. L. Lendon, M. J. Davies, P. D. Richardson, and G. V. R. Born, "Testing of small connective tissue specimens for the determination of the mechanical behaviour of atherosclerotic plaques," <i>Journal of Biomechanical Engineering</i> , vol. 15, pp. 27–33, 1993.
702 703 704	[8]	H. M. Loree, B. J. Tobias, L. J. Gibson, R. D. Kamm, D. M. Small, and R. T. Lee, "Mechanical properties of model atherosclerotic lesion lipid pools," <i>Arteriosclerosis Thrombosis and Vascular Bioology</i> , vol. 14, pp. 230–234, 1994.
705 706 707 708	[9]	S. R. H. Barrett, M. P. F. Sutcliffe, S. Howarth, ZY. Li, and J. H. Gillard, "Experimental measurement of the mechanical properties of carotid atherothrombotic plaque fibrous cap," <i>Journal of Biomechanics</i> , vol. 42, pp. 1650–1655, 2009.
709 710	[10]	J. O. V. Delgadillo, "Mechanical properties of arterial wall," Ph.D. disser- tation, University of British Columbia, 2008.
711 712 713 714	[11]	G. A. Holzapfel, "Biomechanics of Soft Tissues with Application to Ar- terial Walls," in <i>Mathematical and Computational Modeling of Biological Sys-</i> <i>tems</i> , J. A. C. Martins and E. A. C. Borges Pires, Eds. Coimbra, Portugal: Centro Internacional de Matemática CIM, 2002, pp. 1–37.
715 716 717 718	[12]	E. Peña, V. Alastrué, A. Laborda, M. A. Martínez, and M. Doblaré, "A constitutive formulation of vascular tissue mechanics including viscoelasticity and softening behaviour," <i>Journal of Biomechanics</i> , vol. 43, no. 5, pp. 984–989, 2010.
719 720 721	[13]	S. Avril, J. M. Huntley, and R. Cusack, "In-vivo measurements of blood viscosity and wall stiffness in the carotid using PC-MRI," <i>European Journal of Computational Mechanics</i> , vol. 18, pp. 9–20, 2009.
722 723 724 725	[14]	M. Couade, M. Pernot, C. Prada, E. Messas, J. Emmerich, P. Bruneval, A. Criton, M. Fink, and M. Tanter, "Quantitative assessment of arterial wall biomechanical properties using shear wave imaging," <i>Ultrasound in Medicine & Biology</i> , vol. 36, no. 10, pp. 1662–1676, 2010.

726 727 728	[15]	W. A. Riley, R. W. Barnes, G. W. Evans, and G. L. Burke, "Ultrasonic Measurement of the Elastic Modulus of the Common Carotid Artery," <i>Stroke</i> , vol. 23, no. 7, pp. 952–956, 1992.
729 730 731 732	[16]	S. Laurent, X. Girerd, J. J. Mourad, P. Lacolley, L. Beck, P. Boutouyrie, J. P. Mignot, and M. Safar, "Elastic modulus of the radial artery wall material is not increased in patients with essential hypertension," <i>Arteriosclerosis and Thrombosis</i> , vol. 14, no. 7, pp. 1223–1231, 1994.
733 734 735 736	[17]	P. J. Brands, A. P. G. Hoeks, J. Willigers, C. Willekes, and R. S. Reneman, "An integrated system for the non-invasive assessment of vessel wall and hemodynamic properties of large arteries by means of ultrasound," <i>European Journal of Ultrasounds</i> , vol. 9, pp. 257–266, 1999.
737 738 739	[18]	Y. Aggoun, D. Sidi, B. I. Levy, S. Lyonnet, J. Kachaner, and D. Bonnet, "Me- chanical properties of the common carotid artery in Williams syndrome," <i>Heart</i> , vol. 84, pp. 290–293, 2000.
740 741 742	[19]	C. Bussy, P. Boutouyrie, P. Lacolley, P. Challande, and S. Laurent, "Intrinsic stiffness of the carotid arterial wall material in essential hypertensives," <i>Hypertension</i> , vol. 35, no. 5, pp. 1049–1054, 2000.
743 744 745 746	[20]	P. Boutouyrie, D. P. Germain, AI. Tropeano, B. Laloux, F. Carenzi, M. Zidi, X. Jeunemaitre, and S. Laurent, "Compressibility of the Carotid Artery in Patients With Pseudoxanthoma Elasticum," <i>Hypertension</i> , vol. 38, no. 5, pp. 1181–1184, 2001.
747 748 749 750	[21]	R. H. Selzer, W. J. Mack, P. L. Lee, H. Kwong-Fu, and H. N. Hodis, "Improved common carotid elasticity and intima-media thickness mea- surements from computer analysis of sequential ultrasound frames," <i>Atherosclerosis</i> , vol. 154, no. 1, pp. 185–193, 2001.
751 752 753	[22]	H. Kanai, H. Hasegawa, M. Ichiki, F. Tezuka, and Y. Koiwa, "Elastic- ity imaging of atheroma with transcutaneous ultrasound: preliminary study," <i>Circulation</i> , vol. 107, no. 24, pp. 3018–3021, 2003.
754 755	[23]	C. Stephanis, D. Mourmouras, and D. Tsagadopoulos, "On the elastic properties of arteries," <i>Journal of Biomechanics</i> , vol. 36, no. 11, pp. 1727–

756 1731, 2003.

757 758 759	[24]	P. Boutouyrie, D. P. Germain, JN. Fiessinger, B. Laloux, J. Perdu, and S. Laurent, "Increased Carotid Wall Stress in Vascular Ehlers-Danlos Syndrome," <i>Circulation</i> , vol. 109, no. 12, pp. 1530–1535, 2004.
760 761 762 763	[25]	H. Hasegawa and H. Kanai, "Measurement of Elastic Moduli of the Ar- terial Wall at Multiple Frequencies by Remote Actuation for Assessment of Viscoelasticity," <i>Japanese Journal of Applied Physics</i> , vol. 43, no. 5B, pp. 3197–3203, 2004.
764 765 766	[26]	I. Masson, "Contribution à la modélisation mécanique du comportement dynamique hyperélastrique et anisotrope de la paroi artérielle," Thèse de doctorat, Université Paris 12 Val de Marne, 2008.
767 768 769 770	[27]	S. Avril, M. Bonnet, AS. Bretelle, M. Grédiac, F. Hild, P. Ienny, F. La- tourte, D. Lemosse, S. Pagano, E. Pagnacco, and F. Pierron, "Overview of Identification Methods of Mechanical Parameters Based on Full-field Measurements," <i>Experimental Mechanics</i> , vol. 48, no. 4, pp. 381–402, 2008.
771 772 773 774 775	[28]	K. B. Chandran, J. H. Mun, K. K. Choi, J. S. Chen, A. Hamilton, A. Na- garaj, and D. D. McPherson, "A method for in-vivo analysis for regional arterial wall material property alterations with atherosclerosis: prelimi- nary results," <i>Medical Engineering & Physics</i> , vol. 25, no. 4, pp. 289–298, 2003.
776 777 778 779	[29]	A. J. Hamilton, H. Kim, A. Nagaraj, JH. H. Mun, L. L. Yan, S. I. Roth, D. D. McPherson, and K. B. Chandran, "Regional material property alter- ations in porcine femoral arteries with atheroma development," <i>Journal of Biomechanics</i> , vol. 38, no. 12, pp. 2354–2364, 2005.
780 781 782 783	[30]	R. A. Baldewsing, M. G. Danilouchkine, F. Mastik, J. A. Shaar, P. W. Serruys, and A. F. W. van der Steen, "An Inverse Method for Imaging the Local Elasticity of Atherosclerotic Coronary Plaques," <i>IEEE Transactions on Information Technology in Biomedicine</i> , vol. 12, no. 3, pp. 277–289, 2008.
784 785 786 787 788	[31]	S. Le Floc'h, J. Ohayon, P. Tracqui, G. Finet, A. M. Gharib, R. L. Maurice, G. Cloutier, and R. I. Pettigrew, "Vulnerable Atherosclerotic Plaque Elas- ticity Reconstruction Based on a Segmentation-driven Optimization Pro- cedure using Strain Measurements: Theoretical Framework," <i>IEEE Trans-</i> <i>actions on Medical Imaging</i> , vol. 28, no. 7, pp. 1126–1137, 2009.

789 790 791 792	[32]	A. P. Lin, E. Bennett, L. E. Wisk, M. Gharib, S. E. Fraser, H. Wen, and E. Benett, "Circumferential Strain in the Wall of the Common Carotid Artery: Comparing Displacement-Encoded and Cine MRI in Volunteers," <i>Magnetic resonance in Medicine</i> , vol. 60, no. 1, pp. 8–13, 2008.
793 794 795	[33]	M. I. Miga, "A new approach to elastography using mutual information and finite elements," <i>Physics in Medicine and Biology</i> , vol. 48, no. 4, pp. 467–480, 2003.
796 797 798	[34]	C. W. Washington and M. I. Miga, "Modality independent elastography (MIE): a new approach to elasticity imaging." <i>IEEE Transactions on Medical Imaging</i> , vol. 23, no. 9, pp. 1117–1128, 2004.
799 800	[35]	M. A. Bernstein, K. F. King, and X. J. Zhou, <i>Handbook of MRI pulse sequences</i> . Elsevier Inc., 2004.
801 802 803 804	[36]	F. Guyon and R. Le Riche, "Least Squares Parameter Estimation and the Levenberg-Marquardt Algorithm : Deterministic Analysis, Sensitivities and Numerical Experiments," Institut National des Sciences Appliquées, Tech. Rep., 2000.
805 806 807 808	[37]	A. Franquet, S. Avril, R. Le Riche, and P. Badel, "Identification of heterogeneous elastic properties in stenosed arteries: a numerical plane strain study," <i>Computer Methods in Biomechanics and Biomedical Engineering</i> , vol. 15, no. 1, pp. 49–58, 2012.
809 810	[38]	ABAQUS 6.8, <i>Abaqus Theory Manual</i> , Dassault Systèmes Simulia Corp. Providence (RI,USA), Ed., 2008.
811 812 813 814	[39]	M. R. Skilton, L. Boussel, F. Bonnet, S. Bernard, P. C. Douek, P. Moulin, and A. Serusclat, "Carotid intimamedia and adventitial thickening: Comparison of new and established ultrasound and magnetic resonance imaging techniques," <i>Atherosclerosis</i> , vol. 21, pp. 405–410, 2011.
815 816	[40]	I. S. Mackenzie, I. B. Wilkinson, and J. R. Cockcroft, "Assessment of Arterial Stiffness in Clinical Practice," <i>QJM</i> , vol. 95, no. 2, pp. 67–74, 2002.
817 818 819	[41]	S. Laurent, J. Cockcroft, L. Van Bortel, P. Boutouyrie, C. Giannattasio, D. Hayoz, B. Pannier, C. Vlachopoulos, I. Wilkinson, and H. Struijker-Boudier, "Expert Consensus Document on Arterial Stiffness: Method-

ological Issues and Clinical Applications," *European Heart Journal*, vol. 27,
 no. 21, pp. 2588–2605, 2006.

[42] R. A. Baldewsing, J. A. Schaar, F. Mastik, C. W. J. Oomens, and A. F. W.
van der Steen, "Assessment of vulnerable plaque composition by matching the deformation of a parametric plaque model to measured plaque
deformation," *IEEE Transactions on Medical Imaging*, vol. 24, no. 4, pp. 514–
528, 2005.

[43] V. Taviani, M. P. F. Sutcliffe, P. Wong, Z.-Y. Li, V. Young, M. J. Graves, and
J. H. Gillard, "In vivo non-invasive high resolution MR-based method for
the determination of the elastic modulus of arterial vessels," *Annual Inter- national Conference of the IEEE Engineering in Medicine and Biology Society.*,
vol. 2008, pp. 5569–5572, 2008.

- [44] G. A. Holzapfel and H. W. Weizsacker, "Biomechanical behavior of the
 arterial wall and its numerical characterization," *Computers in Biology and Medicine*, vol. 28, no. 4, pp. 377–392, 1998.
- [45] J. D. Humphrey, *Cardiovascular Solid Mechanics : Cells, Tissues and Organs*.
 New-York: Springer-Verlag, 2002.

[46] A. I. Veress, N. Phatak, and J. A. Weiss, "Deformable Image Registration with Hyperelastic Warping," in *Handbook of Biomedical Image Analysis*,
3rd ed., J. S. Suri, D. L. Wilson, and S. Laxminarayan, Eds. New-York:
Kluwer Academic/Plenum, 2004, pp. 487–534.

[47] J. A. Weiss and A. I. Veress, "Strain measurement using deformable image
registration," in *Mechanics of Biological Tissue*, G. A. Holzapfel and R. W.
Ogden, Eds., 2006, pp. 489–501.

[48] N. S. Phatak, S. A. Maas, A. I. Veress, N. A. Pack, E. V. R. Di Bella, and J. A.
Weiss, "Strain measurmeent in the left ventricule during systoo le with deformable image registration," *Medical Image Analysis*, vol. 13, no. 2, pp.

847 354–361, 2009.