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# Time-Frequency Signal Processing for a Self-Mixing Laser Sensor for Vibration Measurement

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**Abstract**— In this paper, a novel time-frequency signal processing approach is presented for a Self-Mixing (SM) interferometric Laser Diode (LD) sensor that enables measurement of harmonic and arbitrarily shaped vibrations. The proposed time-frequency technique ameliorates the measurement precision of a previously published time-domain based displacement retrieval technique called the Phase Unwrapping Method (PUM). By incorporating a frequency-domain analysis to the PUM, we not only improve the measurement precision but also recover information about target movement harmonics that can be used for modal analysis applications. In addition, the time-frequency processing has been found to be robust in case of variations in optical feedback coupling factor. The time-frequency technique has thus provided a precision of approx. 15nm rms (while that of PUM is approx. 40nm rms) for micrometric harmonic and arbitrarily shaped vibrations by using a SM sensor based on a LD emitting at 785nm.

## I. INTRODUCTION

Self-mixing (SM) or optical feedback interferometry technique has been widely used for metrological sensing applications during the last two decades as it results in a simple, compact, self-aligned, and low-cost sensor [1].

SM effect occurs in a laser when a part of the beam backscattered by a target is coupled back into the laser cavity and causes interference with the emitted beam, thus modifying the spectral properties of the laser. The variations in the optical output power of the laser diode  $P(t)$  caused by this optical feedback can be written as [2]:

$$P(t) = P_0 [1 + m \cdot \cos(x_F(t))] \quad (1)$$

where  $P_0$  is the emitted optical power under free-running conditions,  $m$  is the modulation index and  $x_F(t)$  is the laser output phase in the presence of feedback, given by:

$$x_F(t) = 2\pi \frac{D(t)}{\lambda_F(t)/2} \quad (2)$$

where  $D(t)$  is the target displacement.

The emission wavelength subject to feedback  $\lambda_F(t)$  is given by the phase equation [2]:

$$x_0(t) - x_F(t) - C \sin[x_F(t) + \arctan(\alpha)] = 0 \quad (3)$$

where  $\alpha$  is the linewidth enhancement factor and  $x_0(t)$  is the laser output phase in the absence of feedback, given by:

$$x_0(t) = 2\pi \frac{D(t)}{\lambda_0/2} \quad (4)$$

where  $\lambda_0$  is the emission wavelength under free running conditions. The feedback coupling factor  $C$  is given by

$$C = \frac{\tau_D}{\tau_L} \gamma \sqrt{1 + \alpha^2} \kappa_{ext} \quad (4)$$

where  $\tau_L$  and  $\tau_D$  are the round trip times in the internal and external cavities respectively,  $\gamma$  is the coupling efficiency and  $\kappa_{ext}$  depends linearly on the surface reflectivity of the target.

The  $C$  parameter plays an important role in the SM interferometry as variations in  $C$  cause changes in the so called SM operating regimes [1], where each regime would require specific signal processing for the sake of displacement extraction [2]. For the present work, care has been taken to maintain the SM signal in the moderate regime with  $C$  ranging from 1 to 4 due to the fact that such a regime requires simplified SM fringe detection processing and is not affected by the fringe-loss issue [3].

A typical SM sensing set-up and a SM signal  $P(t)$  corresponding to a movement of target  $D(t)$  are presented in Fig. 1. The built-in photodiode available in the laser diode package can be used to measure the variations in the optical output power of the laser diode  $P(t)$  caused by the optical feedback. Each fringe in the SM signal represents a target displacement of  $\lambda/2$ . So, by using simple fringe counting, a sensor would ensure  $\lambda/2$  precision by detecting all fringes [4]. However, various algorithms exist that allow even better measurement precision [5-7].

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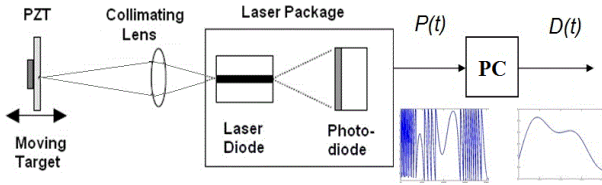


Figure 1. Block diagram of a typical Self-Mixing (SM) sensing set-up requiring only a laser diode package with its built-in photodiode and a collimating lens. A piezoelectric transducer (PZT) has been used as target. Variations in the optical output power of the laser diode  $P(t)$  are processed to retrieve the target motion  $D(t)$ .

## II. PHASE UNWRAPPING METHOD

Let us start with a brief presentation of the Phase Unwrapping Method (PUM) [5] which is summarized in the schematic block diagram of Fig. 2. It can be split up into two principal steps.

The first one leads to the rough estimation of the phase  $x_F(t)$ . After an Automatic Gain Control (AGC) of  $P(t)$  to get  $P(t)/P_0$  ranging over a  $\pm 1$  interval, an arc-cosine function is used to get  $x_F(t) \bmod \pi$  (eq. 1). The derivative of the  $\arccos(P(t))$  is then compared with a threshold value to ascertain the presence of a transition (or a fringe) through the transition detector. An integrator then adds or subtracts  $2\pi$  to  $x_F(t)$ , depending on the sign of the transition associated with a SM fringe. Hence, the feedback phase  $x_F$  is thus roughly retrieved (The first step of PUM is then denoted as Rough PUM or RPUM for all subsequent discussion).

The second step of PUM deals with the joint estimation of parameters,  $C$  and  $\theta = x_F(t=0) + \arctan(\alpha)$ , by applying an iterative bi-dimensional minimization routine on the phase equation (3). This then leads to an approximation of free-running laser phase  $x_0$  proportional to target motion  $D(t)$ .

It is due to these two principal steps that the PUM is able to provide a target movement measurement precision of

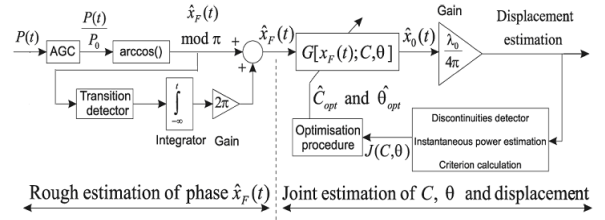


Figure 2. Block diagram of the Phase Unwrapping Method (PUM) indicating its two principal steps [5].

roughly  $\lambda/16$ . The experimental results that would be presented at the end of the paper will attest to such a precision of the PUM.

Let us now move on to our proposed time-frequency based processing of the self-mixing signal that reutilizes the building blocks provided by the PUM. The resultant algorithm is named Modified Phase Unwrapping Method (MPUM) and is detailed in the next section.

## III. MODIFIED PHASE UNWRAPPING METHOD

A closer look into the working of the two principal steps of PUM indicates that the first step (corresponding to the RPUM) is relatively easy to implement and results in a staircase shaped feedback phase signal as a result of  $2\pi$  addition/subtraction. The second step of PUM, however, is more complicated to implement due to its dependence on the use of an iterative bi-dimensional minimization routine. It is finally this step that allows recovering the free-running laser phase  $x_0$  from stair-case shaped feedback phase  $x_F$ .

So, keeping in mind the nature of these two principal steps of the PUM, it was decided to approach the second step of the PUM in a different manner that may allow us to avoid any iterative minimization routines while achieving at least equivalent or even better target movement measurement precision.

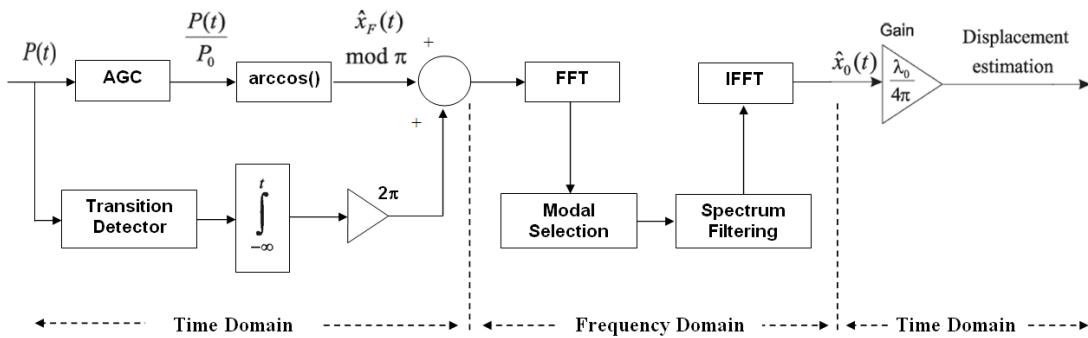


Figure 3. Block diagram of Time-Frequency signal processing of a SM interferometric sensor signal.

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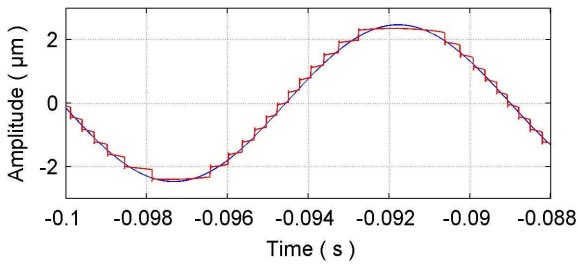


Figure 4. RPUM (red) and MPUM (blue) for an experimental SM signal with  $C=3.85$ .

The block diagram of the ensuing MPUM that is based on a time domain-frequency domain signal processing of SM interferometric sensor signal is thus presented in Fig. 3.

As already mentioned, it was decided to keep the first step of PUM. This results in recovering the roughly retrieved feedback phase  $x_F$ . Then, the second step is replaced with a frequency domain analysis of the roughly retrieved feedback phase  $x_F$  (see Fig. 3). Thus, by using FFT (Fast Fourier Transform), the frequency components making up the target movement are analyzed. By using a threshold, those frequency components that make up the main mode of the target movement (or modes in case of non-harmonic movement of the target) can be selected. Then, the unwanted noise and harmonics contained in the frequency spectrum of  $x_F$  are filtered and finally, using the Inverse FFT (IFFT) allows us to retrieve  $x_0$ .

Fig. 4 shows RPUM and MPUM based signals for an experimental SM signal whose  $C$  value is estimated to be 3.85. It can be seen that the time domain-frequency domain signal processing has allowed us to recover the target movement.

Let us now present the experimental set-up used to acquire SM signals subsequently used for processing and comparison.

#### IV. EXPERIMENTAL SET-UP

The experimental set-up used to validate the working of

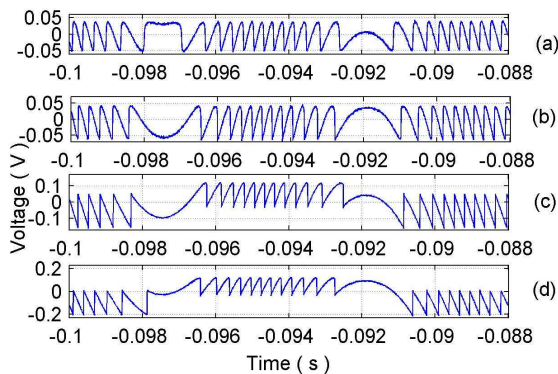


Figure 6. Experimental SM signals with different  $C$  values corresponding to a target vibration of  $5\mu\text{m}$  at  $90\text{Hz}$ .

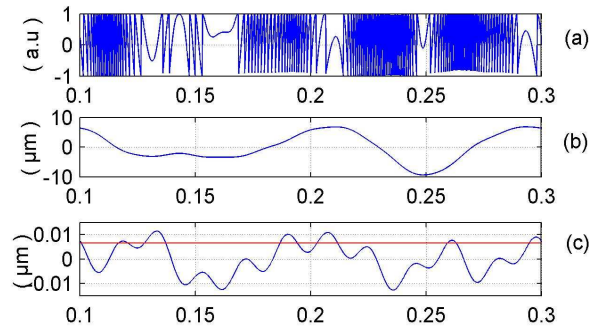


Figure 5. Simulated arbitrary movement: (a) SM signal, (b) retrieved movement, and (c) error (blue) and RMS error (red).

the time domain-frequency domain signal processing of SM interferometric laser sensor signals is shown in Fig. 1. The SM sensor is based on a Hitachi® HL7851 laser diode package emitting at  $\lambda=785\text{ nm}$  with an output power of  $50\text{ mW}$ . A commercial PZT (piezoelectric transducer) actuator from Physik Instrumente (P753.2CD) was used as target. This device has a built-in capacitive feedback sensor with  $2\text{nm}$  resolution that was subsequently used as a reference sensor to quantify the error of our processing method.

#### V. RESULTS

Various simulations and experimental tests were done to validate the proposed signal processing as explained below.

##### A. Simulations

Firstly, different harmonic vibrations having a varying feedback coupling factor  $C$  value were simulated using [8] and then processed using the presented signal processing methods. It was thus observed that the MPUM provided better precision results as compared with the PUM for a variation in  $C$  in the range of 1 to 4.

Then, arbitrarily shaped vibration signals were also tested. Such a kind of signal is presented in Fig. 5. It was again noted

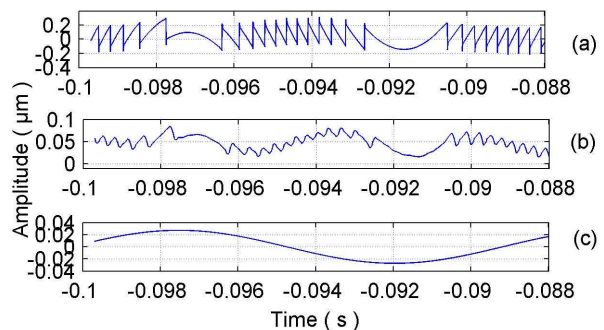


Figure 7. Residual difference between the reference commercial sensor and vibration retrieval using RPUM (a), PUM (b), and MPUM (c).

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TABLE I. RMS ERROR RESULTS FOR EXPERIMENTAL SM SIGNALS (SHOWN IN FIG. 6) FOR DIFFERENT RETRIEVAL METHODS.

SM signal	C	RPUM (nm)	PUM (nm)	MPUM (nm)
1	1,02	100,8	55,4	16,2
2	1,21	77,9	38,2	7,2
3	3,20	99,6	26,3	18,5
4	3,96	98,7	28,5	19,3

that the MPUM provided better precision results as compared with the PUM. The retrieved movement seen in Fig. 5 (b) presents an rms error of less than 10nm as seen in Fig. 5 (c).

### B. Experimental Results

Let us now look into some of the results obtained for experimental SM signals acquired using the set-up indicated in Fig. 1.

Firstly, different experimental SM signals corresponding to a PZT vibration of 5 $\mu$ m at 90Hz were acquired where each of these had a different *C* value (see Fig. 6). The variation in the *C* value was achieved either by varying the sensor-to-target distance or by changing the laser spot size on the PZT (by focusing or defocusing the lens). Thus, it was possible to obtain a variation in *C* in the range of 1 to 4 while maintaining the same PZT movement.

Then, these different SM signals were processed by using the afore-mentioned displacement retrieval methods i.e. PUM, RPUM and MPUM. The rms error results for all of these methods were measured by comparing the retrieved displacement with the reference commercial sensor signal. These error results can be seen in Table 1. It can thus be observed that our proposed MPUM based on a time-frequency signal processing provides better results than PUM. Note that the estimation of the *C* value, indicated in Table 1, was achieved by using the PUM.

Fig. 7 presents the difference between the reference commercial sensor and vibration retrieval using RPUM, PUM, and MPUM for one of the afore-mentioned experimental SM signals. It can then be seen that a frequency domain processing of the RPUM allows us to obtain a signal that is even better than that of the PUM.

Finally, an experimental SM signal for an arbitrary target movement and corresponding retrieval using the MPUM are presented in Fig. 8 which shows that the MPUM is equally capable of processing non-harmonic movements.

### VI. CONCLUSION

A novel time domain-frequency domain signal processing approach is presented for a Self-Mixing (SM) interferometric Laser Diode (LD) sensor which enables a better measurement of harmonic and arbitrarily shaped vibrations as compared to a previously published vibration retrieval technique, called the Phase Unwrapping Method (PUM).

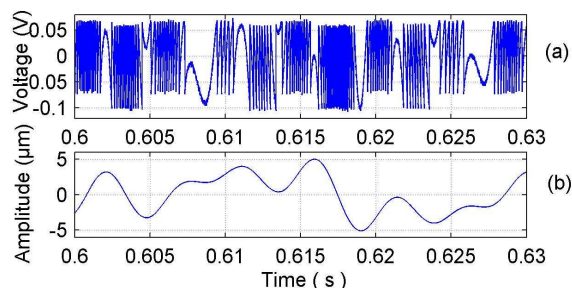


Figure 8. Experimental SM signal (a) for an arbitrary target movement and corresponding retrieval using MPUM (b).

The incorporation of a frequency-domain analysis to the PUM allows us to not only improve the measurement precision but also recover information about target movement harmonics that can be used for modal analysis applications. Furthermore, the time-frequency processing has been found to be robust in case of variations in the optical feedback coupling factor.

The time-frequency method has thus provided a precision of approx. 15nm rms (while that of PUM is approx. 40nm rms) for micrometric harmonic and arbitrarily shaped vibrations measured by using a self-mixing interferometric sensor based on a laser diode emitting at 785nm.

### VII. ACKNOWLEDGEMENT

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