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# Initial trace of solutions of general semilinear heat equation 

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#### Abstract

We study the initial trace problem for positive solutions of semilinear heat equations with strong absorption. We amphasize the particular case where $u$ solves $\partial u-\Delta u+t^{\alpha} u^{q}=0$, with $q>1$ and $\alpha>-1$.


## 1 Introduction

In this paper we study the initial trace problem for positive solutions of

$$
\begin{equation*}
\partial_{t} u-\Delta u+g(x, t, u)=0 \text { dans } Q_{T}^{\Omega}:=\Omega \times(0, T) \tag{1.1}
\end{equation*}
$$

where $\Omega$ is a smooth domain in $\mathbb{R}^{N}$ containing $O$ and $g \in C\left(\Omega \times \mathbb{R}_{+} \times \mathbb{R}\right), g(x, t,$.$) is$ nondecreasing $\forall(x, t) \in \Omega \times \mathbb{R}$ and $g(x, t, r) \geq 0$ if $r \geq 0$. We give general conditions for the existence of an initial trace which belongs to the class of positive Borel measures.
Theorem A Assume $g$ satisfies the above conditions and that equation (1.1) possesses a barrier at any $z \in \Omega$ Then if $\left.u \in C^{1}\left(Q_{T}^{\Omega}\right) \cap C(\Omega \times[0, T] \backslash\{(O, 0))\}\right)$ is a positive solution of (1.1) which vanishes on $\Omega \backslash\{O\}$ at $t=0$ it admits an initial trace which is a outer regular positive Borel measure.

The initial trace of positive solutions of (1.1) exists in the following sense: there exists a relatively closed set $\mathcal{S} \subset \Omega$ and a positive Radon measure $\mu$ on $\mathcal{R}:=\Omega \backslash \mathcal{S}$ with the following properties:
(i) for any $x_{0} \in \mathcal{S}$ and any $\epsilon>0$

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{B_{\epsilon}\left(x_{0}\right) \cap \Omega} u(x, t) d x=\infty \tag{1.2}
\end{equation*}
$$

(ii) for any $\zeta \in C_{c}(\mathcal{R})$

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\Omega} u(x, t) \zeta(x) d x=\int_{\Omega} \zeta d \mu \tag{1.3}
\end{equation*}
$$

The couple $(\mathcal{S}, \mu)$ is unique and characterizes a unique positive outer regular Borel measure $\nu$ on $\Omega$. The barrier property is always verified if $g$ satisfies the Keller-Osserman condition, that is $g(x, t, r) \geq h(r) \geq 0$ where $h$ is nondecreasing and

$$
\begin{equation*}
\int_{a}^{\infty} \frac{d s}{\sqrt{H(s)}} \quad \text { where } H(s)=\int_{0}^{s} h(t) d t \tag{1.4}
\end{equation*}
$$

It is also satisfied if $g(r)=t^{\alpha}|u|^{q-1} u$, with $q>1$ and $\alpha>-1$. This notion has been introduced by Marcus and Véron [10] in the study of the boundary trace of positive solutions of

$$
\begin{equation*}
-\Delta u+g(x, u)=0 \text { dans } \Omega \tag{1.5}
\end{equation*}
$$

In the second part we concentrate on the particular case of equation

$$
\begin{equation*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q}=0 \text { dans } Q_{T}^{\Omega} \tag{1.6}
\end{equation*}
$$

where $T>0, \alpha>-1$ and $q>1$. Among the main tools are the solutions with an isolated singularity at $(0, O)$, which are solutions $u$ of (1.6) in $Q_{T}^{\Omega}$, which belong to $C^{2,1}\left(Q_{T}^{\Omega}\right) \cap$ $C(\Omega \times[0, T) \backslash\{(O, 0)\})$ and satisfy

$$
\begin{equation*}
u(x, 0)=0 \text { in } \Omega \backslash\{O\} \tag{1.7}
\end{equation*}
$$

When $\alpha=0$, Brezis and Friedman prove in [2] that if $B_{2 R} \subset \Omega$, the following estimate holds:

$$
\begin{equation*}
u(x, t) \leq \frac{C(N, q, R)}{\left(|x|^{2}+t\right)^{\frac{1}{q-1}}} \quad \forall(x, t) \in B_{R} \backslash\{0\} \times[0, T] . \tag{1.8}
\end{equation*}
$$

They proved that if $1<q<q_{N}:=1+\frac{2}{N}$ and $k>0$ there exist singular solutions with initial data $u(., 0)=k \delta_{O}$, unique if $u$ vanishes on $\partial \Omega \times[0, T]$. In this range of exponents, Brezis, Peletier, Terman obtain in [1] the existence and uniqueness of a very singular solution of (1.6), always with $\alpha=0$, that is a positive solution in $Q_{\infty}$ under the form

$$
v_{0}(x, t)=t^{-1 /(q-1)} V_{0}\left(\frac{x}{\sqrt{t}}\right),
$$

where $V_{0}>0$ is $C^{2}$ and satisfies

$$
\begin{align*}
-\Delta V_{0}-\frac{1}{2} \eta \cdot \nabla V_{0}-\frac{1}{q-1} V_{0}+V_{0}^{q} & =0 \quad \text { in } \mathbb{R}^{N}  \tag{1.9}\\
\lim _{|\eta| \rightarrow \infty}|\eta|^{\frac{2}{q-1}} V_{0}(\eta) & =0 .
\end{align*}
$$

Actually, Kamin and Peletier proved that $v_{0}$ is the limit of the solutions $u_{k}$ of (1.6) in $Q_{\infty}$ which satisfy $u(., 0)=k \delta_{O}$. The very singular singular solution play a fundamental role in Marcus and Véron's description [7] of the initial trace of positive solutions of (1.6) with $\alpha=0$. In [9], Marcus et Véron study this equation when $\alpha \geq 0$ and $1<q<q_{\alpha, N}=1+\frac{2(1+\alpha)}{N}$. They prove the existence of a self-similar solution of (1.6) under the form

$$
v_{\alpha}(x, t)=t^{-\frac{1+\alpha}{q-1}} V_{\alpha}\left(\frac{x}{\sqrt{t}}\right)
$$

which satisfies

$$
\lim _{t \rightarrow 0} v_{\alpha}(x, t)=0 \quad \forall x \neq 0
$$

and

$$
\lim _{t \rightarrow 0} \int_{B_{\epsilon}} v_{\alpha}(x, t) d x=\infty \quad \forall \epsilon>0
$$

The function $V_{\alpha}$ verifies

$$
\begin{equation*}
-\Delta V_{\alpha}-\frac{1}{2} \eta \cdot \nabla V_{\alpha}-\frac{1+\alpha}{q-1} V_{\alpha}+V_{\alpha}^{q}=0 \text { in } \mathbb{R}^{N} \tag{1.10}
\end{equation*}
$$

Furthermore

$$
V_{\alpha}(\eta)=C|\eta|^{\frac{2(1+\alpha)}{q-1}-N} e^{\frac{-|\eta|^{2}}{4}}(1+o(1)) \text { as }|\eta| \rightarrow \infty
$$

If $1<q<q_{\alpha, N}$, they show that for every $k>0$ there exists a unique solution $u_{k}$ of (1.6) with initial data $k \delta_{O}$. Furthermore $\lim _{k \rightarrow \infty} u_{k}=v_{\alpha}$. Actually the limitation $\alpha \geq 0$ can be relaxed to $\alpha>-1$ has we will see it later on.

In this article we extend Brezis-Friedman removability result to equation (1.6) and prove that if $q \geq q_{\alpha, N}$, any solution in $Q_{T}$ which satisfies (1.7) is identically 0 . The starting point of our study is the following extension of estimate (1.8) valid for any $\alpha>-1$ and $q>1$.

$$
\begin{equation*}
u(x, t) \leq \frac{C(N, q, \alpha, R)}{\left(|x|^{2}+t\right)^{\frac{1+\alpha}{q-1}}} \quad \forall(x, t) \in B_{R} \backslash\{0\} \times[0, T] \tag{1.11}
\end{equation*}
$$

We prove that the following classification results holds.
Theorem B Assume $1<q<q_{\alpha, N}$ and $\left.u \in C^{1}\left(Q_{T}^{\Omega}\right) \cap C(\Omega \times[0, T] \backslash\{(O, 0))\}\right)$ is a positive solution of (1.6) which vanishes on $\Omega \backslash\{O\}$ at $t=0$. Then
(i) either there exists $k \geq 0$ such that $u(., 0)=k \delta_{O}$ and

$$
\begin{equation*}
u(x, t) \sim k E(x, t) \quad \text { as }(x, t) \rightarrow(O, 0) \tag{1.12}
\end{equation*}
$$

where $E(x, t)=(4 \pi t)^{-\frac{N}{2}} e^{-\frac{|x|^{2}}{4 t}}$,
(ii) or

$$
\begin{equation*}
u(x, t) \sim v_{\alpha}(x, t) \quad \text { as }(x, t) \rightarrow(O, 0) . \tag{1.13}
\end{equation*}
$$

In the supercritical case we obtain a removability statement.
Theorem C Assume $q \geq q_{\alpha, N}$ and $u$ is as in Theorem $A$, then $u$ can be extended by continuity as a function in $C(\Omega \times[0, T])$.

We prove that equation (1.6) admits a barrier at any $z \in \Omega$. More precisely we construct a positive solution $w_{B_{R}}$ of (1.6) in $Q_{\infty}^{B_{R}}$ which tends to 0 locally uniformly in $B_{R}$ when $t \rightarrow 0$ and which blows-up uniformly on $\partial B_{R} \times[\tau, \infty)$, for any $\tau>0$. Using sharp asymptotics of the function $V_{\alpha}$, we have the following result which extends Theorem A.
Theorem D Assume $1<q<q_{\alpha, N}$ and $\Omega \subset \mathbb{R}^{N}$ is open with a $C^{2}$ compact boundary, eventually empty. Then for any couple $(\mathcal{S}, \mu)$ where $\mathcal{S} \subset \Omega$ is relatively closed and $\mu \in$ $\mathfrak{M}_{+}(\Omega \backslash \mathcal{S})$, there exists a maximal and a minimal $u \in C^{1}\left(Q_{T}^{\Omega}\right) \cap C^{1}(\bar{\Omega} \times(0, T])$ which both satisfy (i) and (ii) and vanish on $\partial \Omega \times(0, T]$. If $\inf \left\{\left|z-z^{\prime}\right|: z \in \mathcal{S}, z^{\prime} \in \Omega^{c}\right\}>0$ and $\mu$ is bounded in a neighborhood of $\partial \Omega$, then $u$ is unique.

## 2 Initial trace

In this section $\Omega \subset \mathbb{R}^{N}$ is an open set, $Q_{T}^{\Omega}=\Omega \times(0, T), \partial_{\ell} Q_{T}^{\Omega}=\Omega \times\{0\} \cup \partial \Omega \times[0, T)$ and $g \in C\left(\Omega \times \mathbb{R}_{+} \times \mathbb{R}\right)$. If $u$ is defined in $\Omega \times \mathbb{R}_{+}$, we denote by $g \circ u$ the function $(x, t) \mapsto g(x, t, u(x, t))$. We say that $g$ belongs to $\mathcal{H}\left(\right.$ resp $\left.\mathcal{H}_{0}\right)$ if

$$
\begin{align*}
& g(x, t, r) \geq 0 \quad \forall(x, t, r) \in \Omega \times \mathbb{R}_{+} \times \mathbb{R}_{+} \\
& (\text {resp. } g \in \mathcal{H} \text { and } r \mapsto g(x, t, r) \text { is nondecreasing }) . \tag{2.1}
\end{align*}
$$

We denote by $\mathfrak{M}(\Omega)$ the set of Radon measures in $\Omega$, and by $\mathfrak{M}^{b}(\Omega)$ (resp. $\mathfrak{M}^{b, \rho}(\Omega)$ ) the subset of Radon measures such that

$$
\int_{\Omega} d|\mu|<\infty \quad\left(\text { resp. } \int_{\Omega} \rho d|\mu|<\infty\right)
$$

where $\rho(x):=\operatorname{dist}(x, \partial \Omega)$. Their positive cones are respectively $\mathfrak{M}_{+}(\Omega), \mathfrak{M}_{+}^{b}(\Omega)$ and $\mathfrak{M}_{+}^{b, \rho}(\Omega)$.

Definition 2.1 Let $\mathcal{S}$ be a relatively closed subset of $\Omega$ and $\mu$ a Radon measure on $\mathcal{R}:=$ $\Omega \backslash \mathcal{S}$. We say that a nonnegative function $u \in C\left(Q_{T}^{\Omega}\right)$ admits the couple $(\mathcal{S}, \mu)$ for initial trace if

$$
\lim _{t \rightarrow 0} \int_{\mathcal{R}} u(x, t) \zeta(x) d x=\int_{\mathcal{R}} \zeta d \mu \quad \forall \zeta \in C_{c}(\mathcal{R})
$$

and

$$
\lim _{t \rightarrow 0} \int_{U} u(x, t) d x=\infty \quad \forall U \subset \Omega, U \text { open, } U \cap \mathcal{S} \neq \emptyset
$$

The set $\mathcal{S}$ is the set of singular initial points of $u$ and its complement $\mathcal{R}$ the set of regular initial points. We write $\operatorname{tr}_{\Omega}(u)=(\mathcal{S}, \mu)$.

Let $\tilde{\mu}$ be the extension of $\mu$ as a locally bounded Borel measure. To the couple $(\mathcal{S}, \tilde{\mu})$ we can associate a unique Borel measure $\nu$ defined by

$$
\nu(E)= \begin{cases}\infty & \forall E \subset \Omega: E \text { Borel, } E \cap \mathcal{S} \neq \emptyset  \tag{2.2}\\ \tilde{\mu}(E) & \forall E \subset \Omega: E \text { Borel, } E \subset \mathcal{R}\end{cases}
$$

Lemma 2.1 Assume $\Omega$ is a bounded open set with a $C^{2}$ boundary, $T>0, g \in \mathcal{H}$, and let $u \in C(\bar{\Omega} \times(0, T])$ be a positive solution of

$$
\begin{equation*}
\partial_{t} u-\Delta u+g \circ u=0 \quad \text { in } Q_{T}^{\Omega} . \tag{2.3}
\end{equation*}
$$

If $g \circ u \in L_{\rho}^{1}\left(Q_{T}^{\Omega}\right)$, then $u \in L^{\infty}\left(0, T, L_{\rho}^{1}(\Omega)\right)$ and there exists $\mu \in \mathfrak{M}_{+}^{\rho}(\Omega)$ such that

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\Omega} u(x, t) \zeta(x) d x=\int_{\Omega} \zeta d \mu \quad \forall \zeta \in C_{c}(\Omega) \tag{2.4}
\end{equation*}
$$

Proof. If $\phi_{1}>0$ is the first eigenfunction of $-\Delta$ in $W_{0}^{1,2}(\Omega)$ and $\lambda_{1}$ is the corresponding eigenvalue, we have

$$
\frac{d}{d t} \int_{\Omega} u \phi_{1} d x+\lambda_{1} \int_{\Omega} u \phi_{1} d x+\int_{\Omega} g \circ u \phi_{1} d x+\int_{\partial \Omega} u \frac{\partial \phi_{1}}{\partial \nu} d S=0
$$

where $\nu$ is the normal vector. Set $X=\int_{\Omega} u \phi_{1} d x$, then by Hopf Lemma,

$$
X^{\prime}+\lambda_{1} X+\int_{\Omega} g \circ u \phi_{1} d x \geq 0
$$

which yields to

$$
\frac{d}{d t}\left(e^{\lambda_{1} t} X\right)+e^{\lambda_{1} t} \int_{\Omega} g \circ u \phi_{1} d x \geq 0
$$

For $s \in(t, T)$

$$
\frac{d}{d t}\left(e^{\lambda_{1} t} X-\int_{t}^{T} e^{\lambda_{1} s} \int_{\Omega} g \circ u \phi_{1} d x d s\right) \geq 0
$$

which means that the mapping

$$
t \rightarrow e^{\lambda_{1} t} X-\int_{t}^{T} e^{\lambda_{1} s} \int_{\Omega} g \circ u \phi_{1} d x d s
$$

is nondecreasing. Therefore

$$
e^{\lambda_{1} t} X-\int_{t}^{T} e^{\lambda_{1} s} \int_{\Omega} g \circ u \phi_{1} d x d s \leq e^{\lambda_{1} T} X
$$

and finally

$$
X \leq e^{\lambda_{1}(T-t)} X+e^{-\lambda_{1} t} \int_{t}^{T} e^{\lambda_{1} s} \int_{\Omega} g \circ u \phi_{1} d x d s
$$

Since $\rho^{-1} \phi_{1}$ is positively bounded from above and from below, $u \in L^{\infty}\left(0, T, L_{\rho}^{1}(\Omega)\right)$ and there exists a sequence $\left\{t_{n}\right\}$ decreasing to 0 and a measure $\mu \in \mathfrak{M}_{+}^{\rho}(\Omega)$ such that

$$
\lim _{t_{n} \rightarrow 0} \int_{\Omega} u\left(x, t_{n}\right) \zeta d x=\int_{\Omega} \zeta d \mu \quad \forall \zeta \in C_{c}(\Omega)
$$

If $\zeta \in C_{c}^{2}(\Omega)$ there holds

$$
\int_{\Omega} u\left(x, t_{n}\right) \zeta d x=\int_{t_{n}}^{T} \int_{\Omega}(g \circ u \zeta-u \Delta \zeta) d x d t+\int_{\Omega} u(x, T) \zeta d x
$$

thus

$$
\int_{\Omega} \zeta d \mu=\iint_{Q_{T}^{\Omega}}(g \circ u \zeta-u \Delta \zeta) d x d t+\int_{\Omega} u(x, T) \zeta d x .
$$

This implies that $\mu$ is uniquely determined and $u(., t)$ converges to $\mu$ in the weak sense of measures.

Corollary 2.1 Assume $\Omega$ is an open domain, $g \in \mathcal{H}$ and $u \in C^{2}\left(Q_{T}^{\Omega}\right)$ is a positive solution of (2.4). Suppose that for any $z \in \Omega$ there exists an open neighborhood $U \subset \Omega$ such that

$$
\int_{0}^{T} \int_{U} g \circ u d x d t<\infty
$$

Then $u(x, t) \in L^{\infty}\left(0, T, L_{l o c}^{1}(\Omega)\right)$ and there exists a positive Radon measure $\mu$ on $U$ such that

$$
\lim _{t \rightarrow 0} \int_{\mathcal{R}} u(x, t) f(x) d x=\int_{\mathcal{R}} d \mu \forall f \in C_{c}(\mathcal{R})
$$

Proof. We apply the previous lemma in replacing $U$ by a ball $B_{\epsilon}(z)$ and conclude by a partition of unity.

The following class of nonlinearity has been introduced by Marcus and Véron [10] in order to study the boundary trace of solutions of elliptic equations

Definition 2.2 A function $g \in \mathcal{H}$ is a coercive nonlinearity in $Q_{T}^{\Omega}$ if, for every subdomain $\Omega^{\prime}$ of $\Omega$ and every $\epsilon \in(0, T)$, the set of positive solutions of (1.1) in $Q_{\epsilon, T}^{\Omega^{\prime}}:=\Omega^{\prime} \times(\epsilon, T)$ is uniformly bounded in compact subsets of $Q_{\epsilon, T}^{\Omega^{\prime}}$.

Definition 2.3 Let $z \in \Omega$. We say that equation (1.1) possesses a strong barrier at $z$ if there exists a number $r_{0} \in(0, \rho(z))$ such that, for every $r \in\left(0, r_{0}\right)$, there exists a positive super solution $w=w_{r, z}$ of (1.1) in $B_{r}(z) \times(0, T)$ such that

$$
\begin{equation*}
w \in C\left(B_{r}(z) \times[0, T)\right), \lim _{|x-z| \rightarrow r} w(x, t)=\infty \text { locally uniformly if } t \in(0, T) \tag{2.5}
\end{equation*}
$$

Lemma 2.1 Assume $g \in \mathcal{H}$ is a coercive nonlinearity in $Q_{T}^{\Omega}$, then the set of solutions of (1.1) in $Q_{T}^{\Omega}$ is uniformly bounded from above in every compact subset of $Q_{T}^{\Omega}$. Furthermore, if $g \in \mathcal{H}_{0}, A \subset \Omega$ is open and (1.1) possesses a strong barrier at every point of $z \in A$, then the set of solutions $u$ of (1.1) such that $u \in C(A \times[0, T))$ and $u(x, 0)=0$ on $A$ is uniformly bounded from above in every compact subset of $A \times[0, T)$.

Proof. Let $K$ be a compact subset of $Q_{T}^{\Omega}$ and let $\Omega^{\prime}$ be a smooth, bounded domain of $\Omega$ and $\epsilon>0$ such that $K \subset Q_{\epsilon, T}^{\Omega^{\prime}}$ Let $U=U_{Q_{\infty}^{\Omega^{\prime}}}$ be the minimal large solution of (1.1) in $Q_{T}^{\Omega^{\prime}}$, i.e. the limit, when $k \rightarrow \infty$, of solutions with Cauchy-Dirichlet data $k$ on $\partial_{\ell} Q_{\epsilon, T}^{\Omega^{\prime}}:=$ $\Omega^{\prime} \times\{\epsilon\} \cup \partial \Omega^{\prime} \times[\epsilon, T)$. By the maximum principle, if $u \in C\left(Q_{T}^{\Omega}\right)$ is a solution of (1.1), then $u \leq U$ in $\Omega^{\prime}$.
For the second statement, let $K$ be a compact subset of $A$. For any $z \in K$ there exists $r_{z}>0$ such that for any $r \in\left(0, r_{z}\right)$ there exists a positive super solution of (1.1) in $Q_{T}^{B_{r}(z)}$ which satisfies (2.5). Since $K$ is compact, there exist $z_{1}, \ldots, z_{p}$ such that $K \subset \cup_{j=1}^{p} B_{r_{z_{j}} / 2}\left(z_{j}\right)$. For any $j \in\{1, \ldots, p\}$ we denote by $w_{j}$ the super solution in $Q_{T}^{B_{2 r_{z} / 3}\left(z_{j}\right)}$. By comparison principle, there holds

$$
\begin{equation*}
u(x, t) \leq \sup \left\{w_{j}(x, t):(x, t) \in B_{r_{z_{j}} / 2}\left(z_{j}\right) \times(0, T)\right\}:=M_{j} \tag{2.6}
\end{equation*}
$$

for $(x, t) \in Q_{T}^{B_{r_{z_{j}} / 2}\left(z_{j}\right)}$. Therefore $u \leq M=\max _{j=1, \ldots, p} M_{j}$ in $K \times(0, T)$.

Lemma 2.2 Let $g \in \mathcal{H}$ and $u \in C^{2}\left(Q_{T}^{\Omega}\right)$ be a positive solution of (1.1) and suppose $z \in \Omega$ is such that

$$
\begin{equation*}
\int_{0}^{T} \int_{B_{\epsilon}(z) \cap \Omega} g \circ u d x d t=\infty \quad \forall \epsilon>0 \tag{2.7}
\end{equation*}
$$

Suppose that at least one of the following sets of conditions holds:
(i) There exists an open neighborhood $U^{\prime} \subset \Omega$ of $z$ such that $u \in L^{1}\left(U^{\prime} \times(0, T)\right)$.
(ii) The following hold:

1- $g \in \mathcal{H}_{0}$,
2- (1.1) possesses a strong barrier at $z$.
Then,

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{B_{\epsilon}(z) \cap \Omega} u(x, t) d x=\infty \quad \forall \epsilon>0 \tag{2.8}
\end{equation*}
$$

Proof. Assume that $\Omega$ is bounded. First consider the case when condition (i). holds. Let $\phi \in C^{2,1}\left(U^{\prime} \times[0, T)\right)$ with compact support in $U^{\prime} \times[0, T)$ and such that $\phi(x, 0)=1$ in a neighborhood of $z$. Then

$$
\begin{equation*}
\int_{t}^{T} \int_{U^{\prime}}\left(u\left(-\phi_{t}-\Delta \phi\right)+g \circ u \phi\right) d x d t=\int_{U^{\prime}} u(t) \phi d x-\int_{U^{\prime}} u(T) \phi d x \tag{2.9}
\end{equation*}
$$

By assumption $\int_{t}^{T} \int_{U^{\prime}} u\left(\phi_{t}+\Delta \phi\right) d x d t$ is bounded. We let $t$ tend to 0 , the result follows from (2.7).

Next we assume that condition (ii) holds, $u \notin L^{1}\left(U^{\prime} \times(0, T)\right)$ for any neighborhood $U^{\prime}$ of $z$ and that the conclusion is not valid. Thus there exist $r^{*}>0$, such that $\bar{B}_{r^{*}}(z) \subset U_{z}$ and a sequence $\left\{t_{n}\right\}$ decreasing to 0 such that

$$
\int_{B_{r^{*}}(z)} u\left(x, t_{n}\right) d x \leq M
$$

for some $M>0$. Furthermore $g$ is coercive in $B_{r^{*}}(z) \times(0, T)$. Let $\left\{h_{n, k}\right\} \subset C^{\infty}\left(Q_{T}^{\Omega}\right)$ an increasing sequence with respect to $k$ and $n$ of nonnegative functions such that $h_{n, k}=0$ on $B_{r^{*}}(z) \times\{0\}, 0 \leq h_{n, k} \leq k$ and $h_{n, k}=k$ on $\left(t_{n}, T\right) \times \partial B_{r^{*}}(z)$. Let $w_{h_{n, k}}$ be the solution of (1.1) in $B_{r^{*}}(z) \times(0, T)$ such that $w_{h_{n, k}}=h_{n, k}$ on $\partial_{\ell} Q_{T}^{B_{r^{*}}(z)}$. By the maximum principle and condition (ii)-1, the sequence $\left\{w_{h_{n, k}}\right\}$ is monotone increasing with respect to $k$ and $n$. Condition (ii)-3 implies that, for every $r<r^{*}$ and $\beta<T$, the sequence is bounded in $\bar{B}_{r}(z) \times[0, \beta]$, and since $u$ is locally bounded in $Q_{T}^{\Omega}$ there exists $k=k(n)$ such that $k \geq u$ on $\left(t_{n}, T\right) \times \partial B_{r}(z)$ and $k(n) \rightarrow \infty$ when $n \rightarrow \infty$. Then $w=\lim _{n \rightarrow \infty} w_{h_{n, k}}$ is a solution of (1.1) which blows up on $\partial B_{r^{*}}(z) \times(0, T)$ and vanishes on $B_{r^{*}}(z) \times\{0\}$. Let $v_{n}$ be the solution of the heat equation in $B_{r^{*}}(z) \times\left(t_{n}, T\right)$ such that $v_{n}\left(., t_{n}\right)=u\left(., t_{n}\right)$ in $B_{r^{*}}(z)$ and $v_{n}=0$ on $\partial B_{r^{*}}(z) \times\left(t_{n}, T\right)$. Then $w_{h_{n, k(n)}}+v_{n}$ is a super solution of (1.1) in $B_{r^{*}}(z) \times\left(t_{n}, T\right)$ which dominates $u$ on $\partial_{\ell} Q_{t_{n}, T}^{B_{r^{*}}(z)}$. By the maximum principle,

$$
u \leq w_{h_{n, k(n)}}+v_{n} \quad \text { in } Q_{t_{n}, T}^{B_{r^{*}}(z)}
$$

And we have in particular

$$
\int_{B_{r}(z)} u(x, t) d x \leq \int_{B_{r}(z)}\left(w_{h_{n, k(n)}}+v_{n}\right)(x, t) d x \leq M+\int_{B_{r}(z)} w(x, t) d x \quad \forall t \in\left(t_{n}, T\right)
$$

Since it holds for any $n$, it implies $u \in L^{1}\left(Q_{T}^{B_{r}(z)}\right)$, which leads to a contradiction. Example 1. If $g(x, t, r)=h(r)$ where $h$ is continuous, nondecreasing and $h(0) \geq 0$, then

$$
\begin{equation*}
u_{t}-\Delta u+h(u)=0 \tag{2.10}
\end{equation*}
$$

possesses a strong barrier at any $z \in \Omega$ if and only if $h$ satisfies the Keller-Osserman condition, that is there exists some $a \geq 0$ such that

$$
\begin{equation*}
\int_{a}^{\infty} \frac{d s}{\sqrt{H(s)}}<\infty \text { where } H(s)=\int_{0}^{s} h(\tau) d \tau \tag{2.11}
\end{equation*}
$$

The super solution can be chosen to be the maximal solution $\phi_{r}$ of the elliptic equation

$$
\begin{equation*}
-\Delta \phi+h(\phi)=0 \quad \text { in } B_{r}(z) \tag{2.12}
\end{equation*}
$$

If we assume moreover that $h$ is super-additive, i.e. $h(a+b) \geq h(a)+h(b)$ for all $a, b \geq 0$, then there holds

$$
\begin{equation*}
\int_{a}^{\infty} \frac{d s}{h(s)}<\infty \tag{2.13}
\end{equation*}
$$

and any solution $u$ of (2.10) is dominated in $Q_{T}^{B_{r}(z)}$ by $\phi_{r}(x)+\psi(t)$ where $\psi$ is defined by inversion from

$$
\int_{\psi(t)}^{\infty} \frac{d s}{h(s)}=t \quad \forall t>0
$$

Example 2. If $g(x, t, r)=a(x) b(t) h(r)$ where $a \in C(\Omega), b \in C((0, T)), a, b>0$, then $g$ is a coercive nonlinearity if $h$ is super-additive and satisfies the Keller-Osserman condition. This is not sufficient for the existence of a barrier as it is shown in [9] with $h(r)=r^{q}(q>1)$ $a \equiv 1$ and $b(t)=e^{-\frac{1}{t}}$.

Proposition 2.2 Let $g \in \mathcal{H}_{0}$ such that at any $z \in \Omega$ there exists a strong barrier. We assume also

$$
\begin{equation*}
g(x, t, a)+g(x, t, b) \leq g(x, t, a+b) \quad \forall(x, t, a, b) \in Q_{T}^{\Omega} \times \mathbb{R}_{+} \times \mathbb{R}_{+} \tag{2.14}
\end{equation*}
$$

Let $\left\{u_{n}\right\}$ be a sequence of positive solutions of (1.1) which converges to $u$ locally uniformly in $Q_{T}^{\Omega}$. Denote by $\operatorname{tr}_{\Omega}\left(u_{n}\right)=\left(\mathcal{S}_{n}, \mu_{n}\right)$ and $\operatorname{tr}_{\Omega}(u)=(\mathcal{S}, \mu)$ their respective initial trace. If $\mathcal{A} \subset \cap_{n} \mathcal{R}_{n}$ is open and if $\mu_{n}\left(\mathcal{A}_{n}\right)$ remains bounded independently of $n \in \mathbb{N}$, then $\mathcal{A} \subset \mathcal{R}:=$ $\Omega \backslash \mathcal{S}$.
Proof. Let $z \in A$ and $\tilde{r} \in(0, \rho(z))$ such that for any $r \in(0, \tilde{r}]$ there exists a positive super solution $w_{r, z}$ satisfying (2.5) and $\bar{B}_{\tilde{r}}(z) \subset \mathcal{A}$. For any $n \in \mathbb{N}$ and $\tau \in(0, T)$, we denote by $u_{\tau, \chi_{B_{\tilde{r}}(z)}} \mu_{n}$ the solution of

$$
\begin{align*}
\partial_{t} u-\Delta u+g \circ u & =0 & & \text { in } B_{\tilde{r}}(z) \times(\tau, T) \\
u(., \tau) & =\chi_{B_{\tilde{r}}(z)} u_{n}(., \tau) & & \text { in } B_{\tilde{r}}(z)  \tag{2.15}\\
u & =0 & & \text { in } \partial B_{\tilde{r}}(z) \times(\tau, T)
\end{align*}
$$

By the maximum principle $u_{\tau, \chi_{B_{\tilde{r}}(z)}} \mu_{n} \leq u_{n}$ in $B_{\tilde{r}}(z) \times(\tau, T)$, and $g \circ u_{\tau, \chi_{B_{\tilde{r}}(z)}} \mu_{n} \leq g \circ u_{n}$ Furthermore, if $\zeta \in C^{1,1 ; 1}\left(Q_{T}^{\bar{B}_{\tilde{r}}(z)}\right)$ vanishes on $\partial B_{\tilde{r}}(z) \times[0, T)$ and for $t=T$, there holds

$$
\begin{equation*}
\iint_{B_{\tilde{r}}(z) \times(\tau, T)}\left(-u_{\tau, \chi_{B_{\tilde{r}}(z)}} \mu_{n}\left(\partial_{t} \zeta+\Delta \zeta\right)+\zeta g \circ u_{\tau, \chi_{B_{\tilde{r}}(z)} \mu_{n}}\right) d x d t=\int_{B_{\tilde{r}}(z)} u_{n}(x, \tau) \zeta(x, \tau) d x \tag{2.16}
\end{equation*}
$$

Since $u_{\tau, \chi_{B_{\tilde{r}}(z)}} \mu_{n}$ and $g \circ u_{\tau, \chi_{B_{\tilde{r}}(z)}} \mu_{n}$ are bounded independently of $\tau$, standard regularity theory for parabolic equations implies that they converge a.e. in $B_{\tilde{r}}(z) \times(0, T)$ when $\tau \rightarrow 0$ to $u_{\chi_{B_{\tilde{r}}(z)} \mu_{n}}$ and $g \circ u_{\chi_{B_{\tilde{r}}(z)} \mu_{n}}$. Furthermore

$$
\lim _{\tau \rightarrow 0} \int_{B_{\tilde{r}}(z)} u_{n}(x, \tau) \zeta(x, \tau) d x=\int_{B_{\tilde{r}}(z)} \zeta(x, 0) d \mu_{n}(x) .
$$

Using the dominated convergence theorem, it follows from (2.15) that

$$
\begin{equation*}
\iint_{Q_{T}^{B_{\tilde{r}}(z)}}\left(-u_{\chi_{B_{\tilde{r}}(z)} \mu_{n}}\left(\partial_{t} \zeta+\Delta \zeta\right)+\zeta g \circ u_{\chi_{B_{\tilde{r}}(z)} \mu_{n}}\right) d x d t=\int_{B_{\tilde{r}}(z)} \zeta(x, 0) d \mu_{n}(x) \tag{2.17}
\end{equation*}
$$

and $u_{{X_{\tilde{r}}(z)} \mu_{n}}$ is the (unique) solution of

$$
\begin{align*}
\partial_{t} u-\Delta u+g \circ u & =0 & & \text { in } Q_{T}^{B_{\tilde{T}}(z)} \\
u(., 0) & =\chi_{B_{\tilde{r}}(z)} \mu_{n} & & \text { in } B_{\tilde{r}}(z) \\
u & =0 & & \text { in } \partial B_{\tilde{r}}(z) \times(0, T) . \tag{2.18}
\end{align*}
$$

Furthermore, if $\eta$ is the solution of the backward problem

$$
\begin{align*}
\partial_{t} \eta+\Delta \eta & =-1 & & \text { in } Q_{T}^{B_{\tilde{r}}(z)} \\
\eta(., T) & =0 & & \text { in } B_{\tilde{r}}(z)  \tag{2.19}\\
\eta & =0 & & \text { in } \partial B_{\tilde{r}}(z) \times(0, T),
\end{align*}
$$

there holds

$$
\begin{equation*}
\iint_{Q_{T}^{B_{\tilde{r}}(z)}}\left(u_{\left.\chi_{B_{\tilde{r}}(z)} \mu_{n}+\eta g \circ u_{\chi_{B_{\tilde{r}}(z)} \mu_{n}}\right) d x d t=\int_{B_{\tilde{r}}(z)} \eta(x, 0) d \mu_{n}(x) \leq M, ~, ~, ~}\right. \tag{2.20}
\end{equation*}
$$

for some $M>0$ independent of $n$. Next we set $Z_{\tau, n}:=u_{\tau \chi_{B_{\tilde{r}}(z)}} \mu_{n}+w_{\tilde{r}, z}$. It is a super solution of (2.10) in $(\tau, T) \times B_{\tilde{r}}(z)$ which is infinite on $\partial B_{\tilde{r}}(z) \times[\tau, T)$ and dominates $u_{n}$ in $B_{\tilde{r}}(z)$ at $t=\tau$. Thus $Z_{\tau, n} \geq u_{n}$ in $(\tau, T) \times B_{\tilde{r}}(z)$. Letting $\tau \rightarrow 0$ we finally obtain

$$
\begin{equation*}
u_{\chi_{B_{\tilde{r}}(z)} \mu_{n}}(x, t) \leq u_{n}(x, t) \leq u_{\chi_{B_{\tilde{r}}(z)} \mu_{n}}(x, t)+w_{\tilde{r}, z}(x, t) \quad \forall(x, t) \in Q_{T}^{B_{\tilde{r}}(z)} \tag{2.21}
\end{equation*}
$$

For any $r<\tilde{r}$ and $T^{\prime}<T$, there exists $\delta, \sigma>0$ such that $\eta(x, t) \geq \delta$ and $w_{\tilde{r}, z}(x, t) \leq \sigma$ for all $(x, t) \in Q_{T^{\prime}}^{B_{r}(z)}$. It follows from (2.20), (2.21) and Fatou's lemma that $u$ and $g \circ u$ are integrable in $L^{1}\left(Q_{T^{\prime}}^{B_{r}(z)}\right)$. By Lemma 2.1, $B_{r^{\prime}}(z) \subset \mathcal{R}$. Since it holds for any $z \in \mathcal{A}$, the result is proved.

## 3 Construction of a barrier

In the next results we construct the barrier function
Lemma 3.1 Assume $\alpha>-1$ and $q>1$, then there exists a unique positive function $W_{\alpha} \in$ $C^{2}([0, \infty))$ satisfying

$$
\begin{align*}
W^{\prime \prime}+\frac{r}{2} W^{\prime}+\frac{1+\alpha}{q-1} W-W^{q} & =0 \quad \text { in }(0, \infty) \\
\lim _{r \rightarrow 0} W(r) & =\infty  \tag{3.1}\\
\lim _{r \rightarrow \infty} r^{\frac{2}{q-1}} W(r) & =0
\end{align*}
$$

Furthermore $W_{\alpha}$ is decreasing and

$$
\begin{equation*}
W_{\alpha}(r)=C r^{\frac{2(1+\alpha)}{q-1}-1} e^{-\frac{r^{2}}{4}}(1+\circ(1)) \quad \text { as } r \rightarrow \infty \tag{3.2}
\end{equation*}
$$

Proof. Consider the functional

$$
\begin{equation*}
J(\phi):=\frac{1}{2} \int_{0}^{\infty}\left(\phi^{\prime 2}-\frac{1+\alpha}{q-1} \phi^{2}+\frac{2}{q+1}|\phi|^{q+1}\right) e^{\frac{r^{2}}{4}} d r \tag{3.3}
\end{equation*}
$$

defined over the convex set

$$
H_{k}:=\left\{\phi \in W_{2}^{1}\left(0, \infty ; e^{\frac{r^{2}}{4}} d r\right) \cap L^{q+1}\left(0, \infty ; e^{\frac{r^{2}}{4}} d r\right): \phi(0)=k\right\}
$$

Note that if $\phi \in H_{k}$,

$$
\begin{aligned}
e^{\frac{r^{2}}{4}} \phi^{2}(r) & =\int_{r}^{\infty}\left(e^{\frac{s^{2}}{4}} \phi^{2}(s)\right)^{\prime} d s \\
& =2 \int_{r}^{\infty} e^{\frac{s^{2}}{4}} \phi \phi^{\prime}(s) d s+\frac{1}{2} \int_{r}^{\infty} s e^{\frac{s^{2}}{4}} \phi^{2}(s) d s
\end{aligned}
$$

In this set $J$ admits a positive minimizer $w_{k}$ which is the unique solution of

$$
\begin{align*}
w^{\prime \prime}+\frac{r}{2} w+\frac{1+\alpha}{q-1} w^{\prime}-w^{q} & =0 \quad \text { in }(0, \infty)  \tag{3.4}\\
w(0) & =k
\end{align*}
$$

Furthermore, $w_{k}=\lim _{n \rightarrow \infty} w_{k, n}$ where $w_{k, n}$ is the unique positive solution of

$$
\begin{align*}
w^{\prime \prime}+\frac{r}{2} w^{\prime}+\frac{1+\alpha}{q-1} w-w^{q} & =0 \quad \text { in }(0, n) \\
w(0) & =k  \tag{3.5}\\
w(n) & =0
\end{align*}
$$

and, by the maximum principle, $(k, n) \mapsto w_{k, n}$ is increasing. If we consider the linear equation

$$
\begin{equation*}
z^{\prime \prime}+\frac{r}{2} z^{\prime}+\frac{1+\alpha}{q-1} z=0 \quad \text { in }(0, \infty) \tag{3.6}
\end{equation*}
$$

it admits two linearly independent positive solutions $z_{1}$ and $z_{2}$ with the following asymptotic behaviour as $r \rightarrow \infty$

$$
\begin{equation*}
z_{1}(r)=r^{-\frac{2(1+\alpha)}{q-1}}(1+\circ(1)) \tag{3.7}
\end{equation*}
$$

and

$$
\begin{equation*}
z_{2}(r)=r^{\frac{2(1+\alpha)}{q-1}-1} e^{-\frac{r^{2}}{4}}(1+\circ(1)) \tag{3.8}
\end{equation*}
$$

(see [8, Appendix]. Since any solution of 3.4, and 3.5 as well, satisfies an a priori estimate of Keller-Osserman type (see [15])

$$
\begin{equation*}
w(r) \leq C r^{-\frac{2}{q-1}} \quad \text { for } 0<r<1 \tag{3.9}
\end{equation*}
$$

there holds

$$
w_{k, n} \leq C z_{2}(r) \quad \text { for } k \geq r \geq 1
$$

Letting $n$ and $k$ go to infinity successively, it follows that $W_{\alpha}=\lim k, n \rightarrow \infty w_{k, n}$ exists. It is a positive solution of problem (3.1) and it satisfies

$$
W_{\alpha}(r) \leq C\left(r^{-\frac{2}{q-1}}+z_{2}(r)\right) \quad \text { for } r>0
$$

The singular behaviour at $r=0$ is standard (see e.g. [15]) and yields to

$$
\begin{equation*}
W_{\alpha}(r)=\frac{2(q+1)}{(q-1)^{2}} r^{-\frac{2}{q-1}}(1+\circ(1)) \quad \text { as } r \rightarrow 0 \tag{3.10}
\end{equation*}
$$

Thus uniqueness follows by the maximum principle and estimates (3.3) and (3.12) is obtained via standard linearization, using the upper estimate at infinity.

In the sequel we set

$$
\begin{equation*}
w_{\alpha}(s, t)=t^{-\frac{1+\alpha}{q-1}} W_{\alpha}\left(\frac{s}{\sqrt{t}}\right) \quad \forall s>0, t>0 \tag{3.11}
\end{equation*}
$$

Proposition 3.2 Assume $\alpha>-1$ and $q>1$. Then for any $R>0$, there exists $C=$ $C(q, \alpha, R)>0$ such that any solution $u$ of (1.6) in $Q_{\infty}^{B_{R}}$ which vanishes on $B_{R} \times\{0\}$ satisfies

$$
\begin{equation*}
u(x, t) \leq 2 N t^{-\frac{1+\alpha}{q-1}} W_{\alpha}\left(\frac{R-|x|}{\sqrt{t}}\right) \quad \forall(x, t) \in Q_{\infty}^{B_{R}} \tag{3.12}
\end{equation*}
$$

Proof. For $m>0$, set $S_{m}=\left\{x=\left(x_{1, \ldots, x_{N}}:\left|x_{j}\right|<m\right\}\right.$. For $R^{\prime}<R$

$$
\begin{equation*}
\tilde{w}_{R^{\prime}}(x, t)=t^{-\frac{1+\alpha}{q-1}} \sum_{j=1}^{N}\left(W_{\alpha}\left(\frac{R^{\prime}-x_{j}}{\sqrt{t}}\right)+W_{\alpha}\left(\frac{R^{\prime}+x_{j}}{\sqrt{t}}\right)\right) \quad \forall(x, t) \in Q_{\infty}^{S_{R^{\prime}}} \tag{3.13}
\end{equation*}
$$

Then $\tilde{w}_{R^{\prime}}$ is a super solution of (1.6) in $Q_{\infty}^{S_{R^{\prime}}}$ which is infinite on $\partial S_{R^{\prime}} \times(0, \infty)$, thus $u \leq \tilde{w}_{R^{\prime}}$. Letting $R^{\prime} \rightarrow R$ yields to $u \leq \tilde{w}_{R}$ in $Q_{\infty}^{S_{R}}$. Since the equation is invariant by rotation, for any $x \in B_{R}$, there is a rotation $\mathcal{R}$ such that $\mathcal{R}(x)$ has only a positive $x_{1}$-coordinate. Thus

$$
\begin{align*}
u(x, t) & \leq \tilde{w}_{R}(|x|, t) \\
& \leq t^{-\frac{1+\alpha}{q-1}}\left(W_{\alpha}\left(\frac{R-|x|}{\sqrt{t}}\right)+(2 N-1) W_{\alpha}\left(\frac{R}{\sqrt{t}}\right)\right)  \tag{3.14}\\
& \leq 2 N t^{-\frac{1+\alpha}{q-1}} W_{\alpha}\left(\frac{R-|x|}{\sqrt{t}}\right),
\end{align*}
$$

which is (3.12).
Proposition 3.3 Assume $\alpha>-1, q>1$ and $R>0$. Then there exists a unique positive solution $w_{B_{R}}$ of (1.6) in $Q_{\infty}^{B_{R}}$, continuous in $B_{R} \times[0, \infty)$, which vanishes on $B_{R} \times\{0\}$ and satisfies $\lim _{|x| \rightarrow R} w_{B_{R}}(x, t)=\infty$, locally uniformly in $(0, \infty)$. In particular

$$
\begin{equation*}
t^{-\frac{1+\alpha}{q-1}} W_{\alpha}\left(\frac{R-|x|}{\sqrt{t}}\right) \leq w_{B_{R}}(x, t) \leq 2 N t^{-\frac{1+\alpha}{q-1}} W_{\alpha}\left(\frac{R-|x|}{\sqrt{t}}\right) \quad \forall(x, t) \in Q_{\infty}^{B_{R}} \tag{3.15}
\end{equation*}
$$

Proof. For $k>0$, let $w_{B_{R}}^{k}$ be the solution of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{\infty}^{B_{R}} \\
u & =k & & \text { in } \partial_{\ell} Q_{\infty}^{B_{R}}  \tag{3.16}\\
u(., 0) & =0 & & \text { in } B_{R} .
\end{align*}
$$

By (3.12), $w_{B_{R}}^{k}(x, t) \leq 2 N t^{-\frac{1+\alpha}{q-1}} W_{\alpha}\left(\frac{R-|x|}{\sqrt{t}}\right)$. There there exists $w_{B_{R}}=\lim _{k \rightarrow \infty} w_{B_{R}}^{k}$ and $w_{B_{R}}$ is a solution of (1.6) in $Q_{\infty}^{B_{R}}$ which vanishes on $B_{R} \times\{0\}$ and is infinite on $\partial B_{R} \times(0, \infty)$. Consider the similarity transformation $T_{m}$ which leaves equation (1.6) invariant

$$
T_{m}[u](x, t)=m^{\frac{1+\alpha}{q-1}} u(\sqrt{m} x, m t) \quad \forall m>0
$$

then $T_{m}\left[w_{B_{R}}^{k}\right]=w_{B_{\frac{R}{m}}^{m}}^{\frac{1+\alpha}{q-1}} k$ which implies

$$
T_{m}\left[w_{B_{R}}\right]=w_{B} \frac{R}{\sqrt{m}} \quad \forall m>0
$$

If $u \in C\left(B_{R} \times[0, \infty)\right)$ is any positive solution of problem

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{\infty}^{B_{R}} \\
\lim _{|x| \rightarrow R} u(x, t) & =\infty & & \text { locally uniformly on }(0, \infty)  \tag{3.17}\\
u(., 0) & =0 & & \text { in } B_{R},
\end{align*}
$$

then for any $m>1$ and $\epsilon>0$, there exists $\tau_{\epsilon}>0$ such that $u(x, t) \leq \epsilon$ in $B_{\frac{R}{\sqrt{m}}}$ for $0 \leq t \leq \tau_{\epsilon}$. Therefore, for any $T>0$,

$$
u \leq \epsilon+w_{B} \frac{R}{\sqrt{m}} \quad \text { on } \partial_{\ell} Q_{T}^{B} \frac{R}{\sqrt{m}} \cup \overline{B_{\frac{R}{\sqrt{m}}}} \times\{0\}
$$

Since $\epsilon+w_{B_{\frac{R}{\sqrt{m}}}^{\sqrt{m}}}$ is a super solution $u \leq \epsilon+w_{\frac{R}{R}^{\sqrt{m}}}$ in $Q_{T}^{B \frac{R}{\sqrt{m}}}$. Letting $\epsilon \rightarrow 0, m \rightarrow 1$ and $T \rightarrow \infty$ yields to $u \leq w_{B_{R}}$. In the same way, with $0<m<1$, we obtain $u \geq w_{B_{R}}$.

The next result is an immediate consequence of Proposition 3.2.
Proposition 3.4 Assume $\alpha>-1, q>1$ and $K \subset \Omega$ is compact. Let $u$ be any solution $u$ of (1.6) in $Q_{\infty}^{\Omega}$ which vanishes on $\Omega \backslash\{K\} \times\{0\}$ and on $\partial \Omega \times[0, \infty)$, then

$$
\begin{equation*}
u(x, t) \leq 2 N t^{-\frac{1+\alpha}{q-1}} W_{\alpha}\left(\frac{\operatorname{dist}(x, K)}{\sqrt{t}}\right) \tag{3.18}
\end{equation*}
$$

## 4 Upper estimates

We start with the following upper estimate already obtained by Shishkov and Véron [12] in the case $\alpha \geq 0$.
Proposition 4.1 Let $q>1$ and $\alpha>-1$. If $u$ is a solution of (1.6) vanishing on $\partial \Omega \times[0, T)$, there holds

$$
\begin{equation*}
u(x, t) \leq c_{\alpha} t^{-\frac{\alpha+1}{q-1}} \quad \text { for all }(x, t) \in Q_{T}^{\Omega} \tag{4.1}
\end{equation*}
$$

with $c_{\alpha}=\left(\frac{\alpha+1}{q-1}\right)^{\frac{1}{q-1}}$.
Proof. Let $\phi(t)=c_{\alpha} t^{-\frac{\alpha+1}{q-1}}$ be the maximal solution of

$$
\begin{aligned}
\phi^{\prime}+t^{\alpha} \phi^{q} & =0 \\
\phi(0) & =\infty .
\end{aligned}
$$

with $c_{\alpha}=\left(\frac{\alpha+1}{q-1}\right)^{\frac{1}{q-1}}$.
Case $\alpha \geq 0$. For $\tau>0$, we denote by $\Phi_{1, \tau}$ the solution of

$$
\begin{align*}
-\Delta \Phi_{1, \tau}+\tau^{\alpha} \Phi_{1, \tau}^{q} & =0 \quad \text { in } B_{1} \\
\lim _{|x| \rightarrow 1} \Phi_{1, \tau}(x) & =\infty \tag{4.2}
\end{align*}
$$

and for $R>0$

$$
\Phi_{R, \tau}(x)=R^{\frac{-2}{q-1}} \Phi_{1, \tau}\left(\frac{x}{R}\right)
$$

Note that $\Phi_{R, \tau}(x)$ is the solution of the problem (4.2) in the ball $B_{R}$. The function $\Phi_{R, \tau}$ tends to 0 uniformly on every compact set of $\mathbb{R}^{N}$ when $R \rightarrow \infty$. Set

$$
\tilde{v}(x, t)=\phi(t-\tau)+\Phi_{R, \tau}(x),
$$

then $\tilde{v}$ is a supersolution of (1.6) in $B_{R} \times[\tau, T)$ which is infinite on $\partial B_{R} \times[\tau, T) \cup B_{R} \times\{0\}$. Then $u(x, t) \leq \tilde{v}(x, t)$. Letting $R \rightarrow \infty$ and $\tau \rightarrow 0$, we obtain

$$
u(x, t) \leq c_{\alpha} t^{-\frac{\alpha+1}{q-1}} \quad \text { for all }(x, t) \in Q_{T}
$$

Case $-1<\alpha<0$. Let $\tau>0$ and $\phi_{\tau}(t)=c_{\alpha}\left(t^{\alpha+1}-\tau^{\alpha+1}\right)^{-\frac{1}{q-1}}$ be the solution of

$$
\begin{aligned}
\phi_{\tau}^{\prime}+t^{\alpha} \phi_{\tau}^{q} & =0 \quad \text { on }(\tau, \infty) \\
\phi_{\tau}(\tau) & =\infty
\end{aligned}
$$

If $\Phi_{1, T}$ is the solution of (4.2) with $\tau=T$, we set

$$
\Phi_{R, T}(x)=R^{\frac{-2}{q-1}} \Phi_{1, T}\left(\frac{x}{R}\right)
$$

Clearly $\Phi_{R, T}$ tends to 0 uniformly on every compact of $\mathbb{R}^{N}$ when $R \rightarrow \infty$. Set

$$
\hat{v}(x, t)=\phi_{\tau}(t)+\Phi_{R, T}(x)
$$

$\hat{v}$ is a supersolution of (1.6) in $B_{R} \times(\tau, T)$, thus $u(x, t) \leq \hat{v}(x, t)$, as in the first case. Letting $R \rightarrow \infty$ and $\tau \rightarrow 0$, we obtain the desired estimate.

Combining Proposition 3.3 and Proposition 4.1 we obtain,

Corollary 4.2 Assume $q>1, \alpha>-1$ and $K \subset \Omega$ is compact. If $u \in C^{2}\left(Q_{T}^{\Omega}\right) \cap C\left(\bar{Q}_{T}^{\Omega} \backslash\right.$ $K \times\{0\})$ is a solution of (1.6) which vanishes on $\partial \Omega \times[0, T) \cup\{(\Omega \backslash K) \times\{0\}\}$, there holds

$$
\begin{equation*}
u(x, t) \leq \min \left\{2 N W_{\alpha}\left(\frac{\operatorname{dist}(x, K)}{\sqrt{t}}\right), c_{\alpha}\right\} t^{-\frac{\alpha+1}{q-1}} \quad \text { for all }(x, t) \in Q_{T}^{\Omega} \tag{4.3}
\end{equation*}
$$

In the particular case where $K=\{O\}$, (4.3) yields to

$$
\begin{equation*}
u(x, t) \leq \min \left\{2 N W_{\alpha}\left(\frac{|x|}{\sqrt{t}}\right), c_{\alpha}\right\} t^{-\frac{\alpha+1}{q-1}} \leq \frac{c_{1}}{\left(|x|^{2}+t\right)^{\frac{1+\alpha}{q-1}}} \quad \text { for all }(x, t) \in Q_{T}^{\Omega} \tag{4.4}
\end{equation*}
$$

for some $c_{1}=c_{1}(\alpha, q)>0$.
Remark. If $\Omega$ is replaced by $\mathbb{R}^{N}$, the previous estimates (4.1), (4.3) and (4.4) remains valid. Furthermore, $K$ needs only to be closed.

## 5 Isolated singularities

In this section we present the results of classifications of isolated singularities of positive solutions of (1.6), always in the range $q>1$ and $\alpha>-1$. Since the proofs are somewhat similar to the ones of classification of boundary singularities of positive solutions of

$$
\begin{equation*}
\partial_{t} u-\Delta u+u^{q}=0 \quad \text { in } Q_{T}^{\Omega} \tag{5.1}
\end{equation*}
$$

obtained by Marcus and Véron in [8], we only indicate their main ideas.
If we look for solution of (1.6) in $Q_{T}:=Q_{T}^{\mathbb{R}^{N}}$ under the form

$$
u(x, t)=t^{\gamma} V\left(\frac{x}{\sqrt{t}}\right)
$$

it is immediate that $\gamma=-\frac{1+\alpha}{q-1}$ and $V$ is a solution of

$$
\begin{equation*}
-\Delta V-\frac{\eta}{2} \cdot \nabla V-\frac{1+\alpha}{q-1} V+V^{q}=0 \quad \text { in } \mathbb{R}^{N} \tag{5.2}
\end{equation*}
$$

It is proved by Escobedo and Kavian [4] that if $\frac{1+\alpha}{q-1}>\frac{N}{2}$, or equivalently if

$$
\begin{equation*}
1<q<q_{c, \alpha}:=1+\frac{2(1+\alpha)}{q-1} \tag{5.3}
\end{equation*}
$$

there exists a positive solution of (5.2) which is the minimizer of the functional

$$
\begin{equation*}
\omega \mapsto J(\omega):=\frac{1}{2} \int_{\mathbb{R}^{N}}\left(|\nabla \omega|^{2}-\frac{1+\alpha}{q-1} \omega^{2}+\frac{2}{q+1}|\omega|^{q+1}\right) e^{\frac{|\eta|^{2}}{4}} d \eta . \tag{5.4}
\end{equation*}
$$

This minimizing solution $V_{\alpha}$ is unique and radial. Furthermore, by adapting the results of [8, Th 2.1], it satisfies

$$
\begin{equation*}
V_{\alpha}(\eta)=C|\eta|^{2 \frac{1+\alpha}{q-1}-N} e^{-\frac{|\eta|^{2}}{4}}(1+\circ(1)) \quad \text { as } \quad|\eta| \rightarrow \infty \tag{5.5}
\end{equation*}
$$

The function

$$
\begin{equation*}
v_{\alpha}(x, t)=t^{-\frac{1+\alpha}{q-1}} V_{\alpha}\left(\frac{x}{\sqrt{t}}\right) \tag{5.6}
\end{equation*}
$$

is a positive solution of (5.2) in $Q_{\infty}$, continuous in $\bar{Q}_{T} \backslash\{(O, 0)\}$, vanishes on $\mathbb{R}^{N} \times\{0\} \backslash$ $\{(O, 0)\}$ and satisfies

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{B_{\epsilon}} v_{\alpha}(x, t) d x=\infty \quad \forall \epsilon>0 \tag{5.7}
\end{equation*}
$$

It is called the very singular solution of (1.6).
When $\mathbb{R}^{N}$ is replaced by a a proper open domain $\Omega$ with a compact $C^{2}$ boundary there exists no self-similar solution to (1.6). For any $k>0$ and $a \in \Omega$ there exists a unique solution $u:=u_{k \delta_{a}}$ to the initial value problem

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{\infty}^{\Omega} \\
u & =0 & & \text { on } \partial \Omega \times[0, \infty)  \tag{5.8}\\
u(., 0) & =k \delta_{a} & & \text { in } \Omega .
\end{align*}
$$

(see e.g. [9]). The function $u$ belongs to $L^{q}\left(Q_{T}^{\Omega} ; t^{\alpha} d x d t\right) \cap L^{1}\left(Q_{T}^{\Omega}\right), T>0$ arbitrary, and satisfies

$$
\begin{equation*}
\iint_{Q_{T}^{\Omega}}\left(-u\left(\partial_{t} \zeta+\Delta \zeta\right)+t^{\alpha} u^{q} \zeta\right) d x d t=k \zeta(a) \tag{5.9}
\end{equation*}
$$

for all $\zeta \in C^{1,1 ; 1}\left(\bar{Q}_{T}^{\Omega}\right)$ which vanishes on $\partial \Omega \times[0, T]$ and on $\Omega \times\{T\}$. It is unique in the class of of weak solutions, i.e. the functions belonging to $L^{q}\left(Q_{T}^{\Omega} ; t^{\alpha} d x d t\right) \cap L^{1}\left(Q_{T}^{\Omega}\right)$ and satisfying the above identity. When $k \rightarrow \infty, u_{k \delta_{a}} \uparrow u_{\infty \delta_{a}}$, where $u_{\infty \delta_{a}}:=u_{\infty, a}$ is a solution of (1.6) in $Q_{\infty}^{\Omega}$ which vanishes on $\partial \Omega \times[0, \infty)$ and on $\bar{\Omega} \times\{0\} \backslash\{(a, 0)\}$ and satisfies (5.7). Finally

$$
\begin{equation*}
u_{k \delta_{a}} \sim k E(x-a, t) \quad \text { when }(x, t) \rightarrow(a, 0) \tag{5.10}
\end{equation*}
$$

and

$$
\begin{equation*}
u_{\infty, a} \sim v_{\alpha}(x-a, t) \quad \text { when }(x, t) \rightarrow(a, 0) . \tag{5.11}
\end{equation*}
$$

The following classification of isolated singularities holds
Theorem 5.1 Assume $\alpha>-1,1<q<q_{c, \alpha}$ and $a \in \Omega$. If $u \in \bar{Q}_{T}^{\Omega} \backslash\{(a, 0)\}$ is a positive solution of (1.6) which vanishes on $\Omega \times\{0\} \backslash\{(a, 0)\}$, then
(i) either there exists $k \geq 0$ such that

$$
\begin{equation*}
u(x, t) \sim k E(x-a, t) \quad \text { when }(x, t) \rightarrow(a, 0) \tag{5.12}
\end{equation*}
$$

and $u$ is a solution of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{T}^{\Omega} \\
u(., 0) & =k \delta_{a} & & \text { in } \Omega, \tag{5.13}
\end{align*}
$$

(ii) or

$$
\begin{equation*}
u(x, t) \sim v_{\alpha}(x-a, t) \quad \text { when }(x, t) \rightarrow(a, 0) \tag{5.14}
\end{equation*}
$$

and $u$ satisfies (5.7).

If $q \geq q_{c, \alpha}$ there holds the following result which extends Brezis and Freidman's classical one.

Theorem 5.2 Assume $\alpha>-1, q \geq q_{c, \alpha}$ and $a \in \Omega$. If $u \in \bar{Q}_{T}^{\Omega} \backslash\{(a, 0)\}$ is a positive solution of (1.6) which vanishes on $\Omega \times\{0\} \backslash\{(a, 0)\}$, then it can be extended as a continuous function $\tilde{u}$ which vanishes on $\Omega \times\{0\}$.

Proof. Up to modifying a few parameters the proof is similar to Brezis-Friedman's construction. The first step is to prove that $u \in L^{q}\left(Q_{T}^{B_{R}}\right)$ for some $R>0$. This is done by using (4.4) and appropriate test functions. Then the same choice of test functions yields to the fact that $u$ satisfies

$$
\lim _{t \rightarrow 0} \int_{\Omega} u(x, t) \zeta(x) d x=0 \quad \forall \zeta \in C_{0}^{\infty}(\Omega)
$$

Finally the extension of $u$ by zero at $t=0$ satisfies the equation in $\Omega \times[0, T)$.
Remark. We recall that $E(x, t)=(4 \pi t)^{-\frac{N}{2}} e^{-\frac{|x|^{2}}{4 t}}$. Then if $1 \leq r<q_{c, \alpha}$, there holds

$$
\begin{equation*}
\iint_{Q_{T}} E^{r}(x, t) t^{\alpha} d x d t<\infty \tag{5.15}
\end{equation*}
$$

while if $r \geq q_{c, \alpha}$

$$
\begin{equation*}
\iint_{Q_{T}} E^{r}(x, t) t^{\alpha} d x d t=\infty \tag{5.16}
\end{equation*}
$$

## 6 The trace theorem

In all this section we assume that $\Omega \subset \mathbb{R}^{N}$ is an open domain with a compact $C^{2}$ boundary, $\alpha>-1$ and $1<q<q_{c, \alpha}$, and let $u \in C(\bar{\Omega} \times(0, T])$ be a positive solution of (1.6) in $Q_{T}^{\Omega}$ which vanishes on $\partial \Omega \times(0, T]$. By Section $2 u$ possesses an initial trace $\operatorname{tr}_{\Omega}(u)=(\mathcal{S}, \mu)$ where $\mathcal{S}$ is a relatively closed subset of $\Omega$ and $\mu$ is a Radon measure on $\mathcal{R}:=\Omega \backslash \mathcal{S}$. To this couple we can associate a unique outer regular Borel measure $\nu$ defined by

$$
\nu(E)= \begin{cases}\mu(E) & \text { if } E \subset \mathcal{R}  \tag{6.1}\\ \infty & \text { if } E \cap \mathcal{S} \neq \emptyset\end{cases}
$$

for any Borel subset $E$ of $\Omega$. Conversely, to any outer Borel measure $\nu$ on $\Omega$ we can associate the regular set $\mathcal{R} \subset \Omega$ which is the set of points $y \in \Omega$ which possess an open neighborhood $\mathcal{O}_{y}$ such that $\nu\left(\mathcal{O}_{y}\right)<\infty$. Clearly $\mathcal{R}$ is open and the restriction of $\nu$ to $\mathcal{R}$ is a positive Radon measure. The set $\mathcal{S}=\Omega \backslash \mathcal{R}$ is relatively closed and it is the singular part of $\nu$. It has the property that $\nu(E)=\infty$ for any Borel set $E$ such that $E \cap \mathcal{S} \neq \emptyset$. We shall denote by $\mathfrak{B}^{\text {reg }}(\Omega)$ the set of outer regular Borel measures in $\Omega$ and by $\mathfrak{B}_{c}^{\text {reg }}(\Omega)$ the subset of $\mathfrak{B}^{\text {reg }}(\Omega)$ for which $\mathcal{S}$ is a compact subset of $\Omega$. Thus $u$ is a solution of the following problem,

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{T}^{\Omega} \\
u \geq 0, u & =0 & & \text { on } \partial \Omega \times[0, T)  \tag{6.2}\\
\operatorname{tr}_{\Omega}(u) & =\nu \approx(\mathcal{S}, \mu) \in \mathfrak{B}^{r e g}(\Omega) . & &
\end{align*}
$$

Definition 6.1 We denote by $\mathcal{U}_{\mathcal{S}, \mu}(\Omega)$ the set of solutions of problem (6.2).

The first step in the characterization of the singular set is the following lower estimate.
Theorem 6.2 Let $u \in \mathcal{U}_{\mathcal{S}, \mu}(\Omega)$ and $a \in \mathcal{S}$, then

$$
\begin{equation*}
u(x, t) \geq u_{\infty, a}(x, t) \tag{6.3}
\end{equation*}
$$

Furthermore, if $\mathcal{S}$ has a nonempty interior $\mathcal{A}$, there holds

$$
\begin{equation*}
\lim _{t \rightarrow 0} t^{\frac{1+\alpha}{q-1}} u(x, t)=c_{\alpha} \tag{6.4}
\end{equation*}
$$

uniformly on compact subsets of $\mathcal{A}$.
We first give a proof of (6.53) in the case where either $-1<\alpha \leq 0$, or $\alpha>0$ and $1<q<q_{c, 0}$.

Proposition 6.3 Assume either $-1<\alpha \leq 0$ and $1<q<q_{c, \alpha}$, or $\alpha>0$ and $1<q<q_{c, 0}$, then inequality (6.53) holds.
Proof. For any $\epsilon>0$, there holds

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{B_{\epsilon}(a)} u(x, t) d x=\infty \tag{6.5}
\end{equation*}
$$

If $k>0$ is fixed, and $\left\{t_{n}\right\}$ is a sequence decreasing to 0 . There exists $t_{n_{1}}$ such that

$$
\int_{B_{2-1}(a)} u\left(x, t_{n_{1}}\right) d x>k
$$

and there exists $m=m_{1}(k)$ such that

$$
\int_{B_{2-1}(a)} m_{1}(k) \wedge u\left(x, t_{n_{1}}\right) d x=k,
$$

where $A \wedge B=\inf \{A, B\}$. Assume we have constructed $t_{n_{j}}<t_{n_{j-1}}$ and $m_{j}(k)>0$ such that

$$
\int_{B_{2^{-j}}(a)} m_{j}(k) \wedge u\left(x, t_{n_{j}}\right) d x=k
$$

Since (6.5) holds with $\epsilon=2^{-j-1}$, there exists $t_{n_{j+1}}<t_{n_{j}}$ such that

$$
\int_{B_{2^{-j-1}}(a)} u\left(x, t_{n_{j}}\right) d x>k .
$$

and thus $m_{j+1}(k)>0$ such that

$$
\int_{B_{2-j-1}(a)} m_{j+1}(k) \wedge u\left(x, t_{n_{j+1}}\right) d x=k .
$$

Next we denote by $u_{j}$ the solution of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } \Omega \times\left(t_{n_{j}}, T\right) \\
u & =0 & & \text { on } \partial \Omega \times\left[t_{n_{j}}, T\right)  \tag{6.6}\\
u\left(., t_{n_{j}}\right) & =m_{j+1}(k) \wedge u\left(., t_{n_{j+1}}\right) & & \text { in } \Omega .
\end{align*}
$$

By the maximum principle $u_{j} \leq u$ in $\Omega \times\left[t_{n_{j}}, T\right)$, or equivalently $u_{j}\left(x, t+t_{n_{j}}\right) \leq u\left(x, t+t_{n_{j}}\right)$ in $Q_{T-t_{n_{j}}}^{\Omega}$. Clearly

$$
u_{j, 0}:=u\left(., t_{n_{j}}\right) \rightharpoonup k \delta_{a} \quad \text { as } j \rightarrow \infty
$$

in the weak sense of measures. In order to prove that $u_{j}$ converges to $u_{k \delta_{a}}$ we notice that $u_{j}\left(x, t+t_{n_{j}}\right) \leq E * u_{j, 0}(x, t)$ in $Q_{\infty}^{\Omega}$. If $-1<\alpha \leq 0$ and $r \in\left(1, q_{c, \alpha}\right)$

$$
\begin{aligned}
\iint_{Q_{T}^{\Omega}} u_{j}^{r}\left(x, t+t_{n_{j}}\right)\left(t+t_{n_{j}}\right)^{\alpha} d x d t & \leq \iint_{Q_{T}}\left(E * u_{j, 0}\right)^{r}(x, t)\left(t+t_{n_{j}}\right)^{\alpha} d x d t \\
& \leq\left\|u_{j, 0}\right\|_{L^{1}}^{r} \iint_{Q_{T}} E^{r}(x, t) t^{\alpha} d x d t
\end{aligned}
$$

using Young's inequality and (5.15). If $\alpha \geq 0$ and $r \in\left(1, q_{c, 0}\right)$

$$
\begin{aligned}
\iint_{Q_{T}^{\Omega}} u_{j}^{r}\left(x, t+t_{n_{j}}\right)\left(t+t_{n_{j}}\right)^{\alpha} d x d t & \leq 2^{\alpha} T^{\alpha} \iint_{Q_{T}}\left(E * u_{j, 0}\right)^{r}(x, t) d x d t \\
& \leq 2^{\alpha} T^{\alpha}\left\|u_{j, 0}\right\|_{L^{1}}^{r} \iint_{Q_{T}} E^{r}(x, t) d x d t
\end{aligned}
$$

Furthermore, for $s \in\left(1, q_{c, 0}\right)$

$$
\begin{aligned}
\iint_{Q_{T}^{\Omega}} u_{j}^{s}\left(x, t+t_{n_{j}}\right) d x d t & \leq \iint_{Q_{T}}\left(E * u_{j, 0}\right)^{s}(x, t) d x d t \\
& \leq\left\|v_{j, 0}\right\|_{L^{1}}^{s} \iint_{Q_{T}} E^{s}(x, t) d x d t
\end{aligned}
$$

Thus the set of functions $\left\{u_{j}^{q}\left(., .+t_{n_{j}}\right)\left(.+t_{n_{j}}\right)^{\alpha}\right\}$ and $\left\{u_{j}\left(., .+t_{n_{j}}\right)\right\}$ are uniformly integrable in $L^{1}\left(Q_{T}^{\Omega}\right)$. Since $u_{j}$ satisfies

$$
\begin{equation*}
\iint_{Q_{T}^{\Omega}}\left(-u_{j}\left(x, t+t_{n_{j}}\right)\left(\partial_{t} \zeta+\Delta \zeta\right)+\left(t+t_{n_{j}}\right)^{\alpha} u_{j}^{q}\left(x, t+t_{n_{j}}\right) \zeta\right) d x d t=\int_{\Omega} u_{j, 0} \zeta d x \tag{6.7}
\end{equation*}
$$

for any test function $\zeta \in C^{1,1 ; 1}\left(\bar{Q}_{T}^{\Omega}\right)$ vanishing on $\partial \Omega \times[0, T] \cup \Omega \times\{T\}$ and converges a.e. in $Q_{T}^{\Omega}$ to some $u^{*}$, it follows by Vitali's theorem that

$$
\begin{equation*}
\iint_{Q_{T}^{\Omega}}\left(-v\left(\partial_{t} \zeta+\Delta \zeta\right)+t^{\alpha} v^{q} \zeta\right) d x d t=k \zeta(a) \tag{6.8}
\end{equation*}
$$

thus $u^{*}=u_{k \delta_{a}}$ by uniqueness, which implies the claim since $u \geq u^{*}$.
When $\alpha>0$ and $q_{c, 0} \leq q<q_{c, \alpha}$, this argument cannot work since the sequence $u_{j, 0}$ could concentrate too quickly with respect to $t$ to a Dirac mass and isolated singularities are removable for solutions of

$$
\partial_{t} u-\Delta u+c u^{q}=0 .
$$

We develop below a proof valid for any $\alpha>-1$ and $1<q<q_{c, \alpha}$, which is based upon the parabolic Harnack inequality and shows that such a fast concentration never occurs.

Lemma 6.4 Assume $\alpha>-1$ and $1<q<q_{c, \alpha} . \operatorname{Let}\left\{\left(x_{n}, t_{n}\right)\right\} \subset Q_{T}^{\Omega}$ be a sequence converging to $(a, 0)$ and $\ell>0$. Put $V_{n}=B_{\ell \sqrt{t_{n}}}\left(x_{n}\right)$ and suppose that there exist nonnegative functions $h_{n} \in L^{\infty}\left(\mathbb{R}^{N}\right)$ with support in $V_{n}$ such that $0 \leq h_{n} \leq c_{1} t_{n}^{-\frac{N}{2}}$ and

$$
\begin{equation*}
h_{n} \rightharpoonup k \delta_{0} . \tag{6.9}
\end{equation*}
$$

Then the solutions $u_{n}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } \Omega \times\left(t_{n}, T\right) \\
u & =0 & & \text { on } \partial \Omega \times\left[t_{n}, T\right)  \tag{6.10}\\
u\left(., t_{n}\right) & =h_{n} & & \text { in } \Omega,
\end{align*}
$$

satisfy $u_{n} \rightarrow u_{k \delta_{a}}$ when $n \rightarrow \infty$.
Proof. The estimate $h_{n} \leq c t_{n}^{-\frac{N}{2}} \chi_{V_{n}}$ can be written under the form

$$
h_{n}(x) \leq c_{2} E\left(x-x_{n}, t_{n}\right) \chi_{V_{n}}
$$

where $c_{2}=(4 \pi)^{\frac{N}{2}} e^{\frac{\ell}{4}} c_{1}$. By the maximum principle

$$
u_{n}(x, t) \leq c_{2} E\left(x-x_{n}, t\right) \quad \text { in } \Omega \times\left(t_{n}, \infty\right) .
$$

By (5.15), (5.16), the sequences $\left\{E^{q}\left(.-x_{n},.\right) t^{\alpha}\right\}$ and $\left\{E\left(.-x_{n},.\right)\right\}$ are uniformly integrable in $Q_{T}^{\Omega}$, therefore, if we extend $u_{n}$ by zero in $Q_{t_{n}}^{\Omega}$ and denote by $\tilde{u}_{n}$ the extended function defined in $Q_{T}^{\Omega}$, we infer that the sequences $\left\{t^{\alpha} \tilde{u}_{n}^{q}\right\}$ and $\left\{\tilde{u}_{n}\right\}$ are uniformly integrable in $Q_{T}^{\Omega}$. Using standard regularity estimates there exists a function $u^{*}$ defined in $Q_{T}^{\Omega}$ and a subsequence $u_{n_{j}}$ such that $\tilde{u}_{n_{j}} \rightarrow u^{*}$ locally uniformly in $Q_{T}^{\Omega}$. It follows by uniform integrability and Vitali's convergence theorem that

$$
\tilde{u}_{n_{j}} \rightarrow u^{*} \quad \text { in } L^{q}\left(Q_{T}^{\Omega} ; t^{\alpha} d x d t\right) \cap L^{1}\left(Q_{T}^{\Omega}\right) .
$$

Let $\zeta \in C^{1,1 ; 1}\left(\bar{Q}_{T}^{\Omega}\right)$ vanishing on $\partial \Omega \times[0, T] \cup \Omega \times\{T\}$, then

$$
\iint_{Q_{T}^{\Omega}}\left(-\tilde{u}_{n_{j}}\left(\partial_{t} \zeta+\Delta \zeta\right)+t^{\alpha} \tilde{u}_{n_{j}}^{q} \zeta\right) d x d t=\int_{\Omega} \zeta\left(., t_{n_{j}}\right) h_{n} d x
$$

Using the previous convergence results and the assumption (6.9), we derive

$$
\iint_{Q_{T}^{\Omega}}\left(-u^{*}\left(\partial_{t} \zeta+\Delta \zeta\right)+t^{\alpha} u^{* q} \zeta\right) d x d t=k \zeta(O) .
$$

Thus $u^{*}=u_{k \delta_{0}}$ and $\tilde{u}_{n} \rightarrow u_{k \delta_{0}}$ locally uniformly in $Q_{T}^{\Omega}$.
Lemma 6.5 Let $u$ be a positive solution of (1.6) in $Q_{T}^{\Omega}$ vanishing on $\partial \Omega \times[0, T)$. Then for any $\Omega^{\prime} \subset \bar{\Omega} \subset \Omega$ there exists a constant $C=C\left(N, q, \alpha, \Omega^{\prime}\right)>0$ such that

$$
\begin{equation*}
u(y, s) \leq u(x, t) e^{C\left(\frac{|x-y|^{2}}{t-s}+\frac{t}{s}+1\right)} \quad \forall(x, t),(y, s) \in Q_{T}^{\Omega^{\prime}}, s<t \tag{6.11}
\end{equation*}
$$

Proof. By (4.1), $V(x, t):=t^{\alpha} u^{q-1} \leq c_{\alpha}^{q-1} t^{-1}$. If we write (1.6) under the form

$$
\begin{equation*}
\partial_{t} u-\Delta u+V(x, t) u=0 \quad \text { in } Q_{T}^{\Omega} \tag{6.12}
\end{equation*}
$$

it follows (6.11) from parabolic Harnack inequality (see e.g. [5, Lemma 3.16] although the result is much older).

If $G \subset \mathbb{R}^{N}$ is a bounded open subset, we denote by $\phi_{G}$ is the first eigenfunction of $-\Delta$ in $W_{0}^{1,2}(G)$ normalized by $\sup \phi_{G}=1$ with corresponding eigenvalue $\lambda_{G}$.
Lemma 6.6 Let $G \subset \mathbb{R}^{N}$ be a bounded open subset with a smooth boundary $\omega \in C^{1}\left(Q_{T}^{G}\right)$, $\omega \geq 0$, such that

$$
\begin{equation*}
\int_{0}^{T}\|\omega(., t)\|_{L^{\infty}(G)} d t<\infty \tag{6.13}
\end{equation*}
$$

If $v \in C^{2 ; 1}(\bar{G} \times(0, T])$ is a positive solution of

$$
\begin{equation*}
\partial_{t} v-\Delta v+\omega v=0 \quad \text { in } Q_{T}^{G} \tag{6.14}
\end{equation*}
$$

then $v \phi_{G} \in L^{1}\left(Q_{T}^{G}\right), \omega v \phi_{G}^{3} \in L^{1}\left(Q_{T}^{G}\right)$ and there exists $\mu \in \mathfrak{M}_{+}(G)$ such that

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{G} v(x, t) \zeta(x) d x=\int_{G} \zeta d \mu \quad \forall \zeta \in C_{c}(\bar{G}) \tag{6.15}
\end{equation*}
$$

Proof. Set $\gamma(t)=\|\omega(., t)\|_{L^{\infty}(G)}$. Denote $X(t)=\int_{G} v(x, t) \phi_{G}(x) d x$. Then, from (6.14),

$$
X^{\prime}+\lambda_{G} X+\gamma(t) X \geq 0
$$

This implies that the function

$$
t \mapsto e^{t \lambda_{G}+\int_{0}^{t} \gamma(s) d s} X(t)
$$

(which exists thanks to (6.13)) is nondecreasing. Therefore there exists $X(0)=\lim _{t \rightarrow 0} X(t)$ and $v \phi_{G} \in L^{1}\left(Q_{T}^{G}\right)$. Furthermore, if we set $Y(t)=\int_{G} v(x, t) \phi_{G}^{3}(x) d x$

$$
\frac{d}{d t}\left(e^{3 t \lambda_{G}} X(t)\right)+e^{3 t \lambda_{1}} \int_{\Omega}\left(\omega \phi_{G}^{3}-6 \phi_{G}\left|\nabla \phi_{G}\right|^{2}\right) v d x=0
$$

Since $\nabla \phi_{G}$ is bounded and $v \phi_{G} \in L^{1}\left(Q_{T}^{G}\right)$, it implies

$$
\left.e^{3 t \lambda_{G}} Y(T)+\int_{0}^{T} e^{3 t \lambda_{G}} \int_{\Omega} \omega(x, t) v(x, t) \phi_{G}^{2} d x d t=Y(0)+6 \int_{0}^{T} e^{3 t \lambda_{G}} \int_{\Omega} \phi_{G}\left|\nabla \phi_{G}\right|^{2}\right) v d x d t
$$

which implies that $\omega v \phi_{G}^{3} \in L^{1}\left(Q_{T}^{G}\right)$. The argument given in the proof of Lemma 2.1 shows that $v$ admits an initial trace which belongs to $\mathfrak{M}_{+}(G)$.
Proof of Theorem 6.2. We define the parabolic distance in $\mathbb{R}^{N} \times \mathbb{R}$ by

$$
d_{P}((x, t),(y, s))=\sqrt{|x-y|^{2}+|t-s|}
$$

Step 1. We first prove that if $u$ satisfies

$$
\begin{equation*}
\limsup _{\left.d_{P}((x, t))(a, 0)\right) \rightarrow 0} t^{\frac{N}{2}} u(x, t)<\infty \tag{6.16}
\end{equation*}
$$

then $a \in \mathcal{R}(u)$. If (6.16) holds there exists $\epsilon, c>0$ such that

$$
u(x, t) \leq c t^{-\frac{N}{2}} \quad \forall(x, t) \text { s.t. } \sqrt{|x-a|^{2}+t} \leq \epsilon
$$

If we set $\omega(x, t)=t^{\alpha} u^{q-1}(x, t)$, then

$$
\omega(x, t) \leq c^{q-1} t^{-\left(\frac{N(q-1)}{2}-\alpha\right)} \quad \forall(x, t) \in B_{\frac{\epsilon}{\sqrt{2}}}(a) \times\left(0, \frac{\epsilon^{2}}{2}\right] .
$$

Since $\left.q<q_{c, \alpha}, \frac{N(q-1)}{2}\right)-\alpha<1$, thus the assumptions of Lemma 6.6 are fulfilled and there exists a positive Radon measure $\mu$ in $B_{\frac{\epsilon}{\sqrt{2}}}(a)$ such that

$$
\int_{\Omega} u(x, t) \zeta(x) d x \rightarrow \int_{\Omega} \zeta d \mu \text { when } t \rightarrow 0, \quad \forall \zeta \in C_{c}^{\infty}\left(B_{\frac{\epsilon}{\sqrt{2}}}(a)\right) .
$$

Furthermore $t^{\alpha} u^{q} \in L^{1}\left(B_{\frac{\epsilon}{\sqrt{3}}}(a) \times(0, T)\right)$, which is the claim.
Step 2. Since $a \in \mathcal{S}(u)$, there holds

$$
\begin{equation*}
\limsup _{\left.d_{P}((x, t))(a, 0)\right) \rightarrow 0} t^{\frac{N}{2}} u(x, t)=\infty . \tag{6.17}
\end{equation*}
$$

Then there exists a sequence $\left\{\left(x_{n}, s_{n}\right)\right\}$ converging to $(a, 0)$ such that

$$
\begin{equation*}
u\left(x_{n}, t_{n}\right) \geq n t_{n}^{-\frac{N}{2}} \tag{6.18}
\end{equation*}
$$

We apply Lemma 6.5 with $s=s_{n}, t=2 s_{n}:=t_{n}, y=x_{n}\left|x-x_{n}\right| \leq \sqrt{s_{n}}$. Then

$$
u\left(x, t_{n}\right) \geq C n t_{n}^{-\frac{N}{2}} \quad \forall x \in V_{n}:=B_{\frac{\sqrt{t_{n}}}{2}}\left(x_{n}\right)
$$

This implies

$$
\int_{V_{n}} u\left(x, t_{n}\right) d x \geq C_{N} n
$$

For $k<n$ fixed, we denote by $v:=v_{n, k}$ the solution of

$$
\begin{align*}
\partial_{t} v-\Delta v+t^{\alpha} v^{q} & =0 & & \text { in } \Omega \times\left(t_{n}, \infty\right) \\
v & =0 & & \text { in } \partial \Omega \times\left(t_{n}, \infty\right) \\
v\left(., t_{n}\right) & =C k t_{n}^{-\frac{N}{2}} \chi_{V_{n}} & & \text { in } \Omega .
\end{align*}
$$

By the maximum principle $u \geq v_{n, k}$ in $\Omega \times\left(t_{n}, \infty\right)$. Furthermore

$$
\int_{V_{n}} v\left(., t_{n}\right) d x=C_{N} k
$$

Thus $v\left(., t_{n}\right) \rightharpoonup C_{N} k \delta_{a}$ in the weak sense of measures. It follows from Lemma 6.4 that $v_{n, k} \rightarrow u_{C_{N} k \delta_{a}}$ locally uniformly in $Q_{T}^{\Omega}$. Therefore $u \geq u_{C_{N} k \delta_{a}}$ in $Q_{T}^{\Omega}$. Since $k$ is arbitrary, we obtain (6.53).
Step 3. Formula (6.4) holds. Denote by $S_{m}(a)=\left\{x \in \mathbb{R}^{N}:\left|x_{j}\right|<m\right\}$. If $\bar{S}_{R}(a) \subset \mathcal{S}$, the function

$$
(x, t) \mapsto c_{\alpha}(t-\tau)^{-\frac{1+\alpha}{q-1}}+w_{B_{R}}(x-a, t-\tau)
$$

is a super solution of $(1.6)$ in $S_{R}(a) \times(\tau, \infty)$ which is infinite on $S_{R}(a) \times\{\tau\} \cup \partial S_{R}(a) \times[\tau, \infty)$ by Proposition 3.3, while $u$ is finite, thus it dominates $u$ in $S_{R}(a) \times(\tau, \infty)$. Letting $\tau \rightarrow 0$ yields to

$$
\begin{equation*}
u(x, t) \leq c_{\alpha} t^{-\frac{1+\alpha}{q-1}}+w_{B_{R}}(x-a, t) \quad \forall(x, t) \in Q_{T}^{S_{R}(a)} \tag{6.20}
\end{equation*}
$$

Conversely, the function

$$
(x, t) \mapsto u(x, t-\tau)+w_{B_{R}}(x-a, t-\tau)
$$

is a super solution in $S_{R}(a) \times(\tau, \infty)$ which dominates $c_{\alpha} t^{-\frac{1+\alpha}{q-1}}$ on $S_{R}(a) \times\{\tau\} \cup \partial S_{R}(a) \times$ $[\tau, \infty)$, thus as above, we obtain (6.20). Since

$$
\lim _{t \rightarrow 0} w_{B_{R}}(x-a, t)=0
$$

uniformly on $B_{R^{\prime}}(a)$ for any $R^{\prime}<R$, we derive (6.4).
Proposition 6.7 For any relatively closed $\mathcal{S} \in \Omega$, the set $\mathcal{U}_{\mathcal{S}, 0}(\Omega)$ is not empty and it admits a minimal element $\underline{u}_{\mathcal{S}, 0}$ and a maximal element $\bar{u}_{\mathcal{S}, 0}$.

Proof. Step 1: Existence of a maximal solution. The maximal solution is constructed by thickening $\Omega$ and $\mathcal{S}$ in defining for $0<\sigma$

$$
\Omega_{\sigma}=\left\{x \in \mathbb{R}^{N}: \operatorname{dist}(x, \Omega)<\sigma\right\}, \quad \mathcal{S}_{\sigma}=\left\{x \in \mathbb{R}^{N}: \operatorname{dist}(x, \overline{\mathcal{S}}) \leq \sigma\right\}
$$

If $z \in \partial \Omega$, we denote by $\mathbf{n}_{z}$ the outward unit normal vector to $\Omega$ at $z$. Since $\partial \Omega$ is compact and $C^{2}$, there exists $\sigma_{0}>0$ such that for any $(z, \sigma) \in \partial \Omega \times\left[0 \sigma_{0}\right]$, the mapping $\Pi:(z, \sigma) \mapsto$ $z+\sigma \mathbf{n}_{z}$ is a $C^{2}$ diffeomorphism from $\partial \Omega \times\left[0, \sigma_{0}\right]$ to $\Theta_{\sigma_{0}}^{\prime}:=\bar{\Omega}_{\sigma_{0}} \backslash \Omega$. The mapping $\Pi$ defines the flow coordinates near $\partial \Omega$.

If $0<\delta<\sigma$, there exists a unique solution $u=u_{n, \sigma, \delta}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{\infty}^{\Omega_{\sigma}} \\
u & =0 & & \text { in } \partial \Omega_{\sigma} \times(0, \infty)  \tag{6.21}\\
u(., 0) & =n \chi_{s_{\delta}} & & \text { in } \Omega_{\sigma}
\end{align*}
$$

Notice that $\mathcal{S}_{\delta}$ is closed in $\Omega_{\sigma}$ and $\inf \left\{\left|z-z^{\prime}\right|: z \in \mathcal{S}_{\delta}, z^{\prime} \in \Omega_{\delta}^{c}\right\}=\delta-\sigma$. Existence is standard as well as uniqueness in the case where $\Omega$ is bounded. If $\Omega^{c}$ is bounded the proof goes as in the uniqueness proof in Proposition 6.10. When $n \rightarrow \infty,\left\{u_{n, \sigma, \delta}\right\} \uparrow u_{\sigma, \delta}$ which is a solution of (1.6) in $Q_{\infty}^{\Omega_{\sigma}}$. Since $u_{n, \sigma, \delta}$ satisfies (4.1), for any $r, \tau>0$ and any $a \in$ $\partial \Omega_{\sigma}, u_{n, \sigma, \delta}$ remains uniformly continuous with respect to $n$ in $\bar{Q}_{\infty}^{\Omega_{\sigma}} \cap\left(\bar{B}_{r}(a) \times\left[2^{-2} \tau, 2 \tau\right]\right)$. Consequently $u_{\sigma, \delta}(x, t)=0$ on $\partial \Omega_{\sigma} \times(0, \infty) \cap\left(\bar{B}_{r}(a) \times\left[2^{-2} \tau, 2 \tau\right]\right)$. Therefore $u_{\sigma, \delta}$ vanishes
on $\partial \Omega_{\sigma} \times(0, \infty)$. If $a \in \Omega_{\sigma}$ and $\operatorname{dist}\left(a, \mathcal{S}_{\delta}\right)=r>0, u_{n, \sigma, \delta}(x, t) \leq w_{B_{r}}(x-a, t)$ in $Q_{B_{r}(a)}^{\infty}$. This implies that $u_{n, \sigma, \delta}$ remains uniformly continuous with respect to $n$ in $\left.\bar{B}_{r^{\prime}}(a) \times[0, T)\right]$ for any $0<r^{\prime}<r$ and $T>0$. Since $u_{n, \sigma, \delta}(x, t) \rightarrow 0$ in $\bar{B}_{r^{\prime}}(a), u_{\sigma, \delta}$ inherits the same property. Consequently $u_{\sigma, \delta}$ has initial trace $\left(\mathcal{S}_{\delta}, 0\right)$ in $\Omega_{\sigma}$. By the maximum principle the mapping $(n, \delta) \mapsto u_{n, \sigma, \delta}$ is increasing with respect to $n$ and decreasing with respect to $\delta$. Furthermore, if $0<\delta^{\prime}<\sigma^{\prime}<\sigma$ and $0<\delta<\sigma$, there holds $u_{n, \sigma^{\prime}, \delta^{\prime}}<u_{n, \sigma, \delta}$ in $Q_{\infty}^{\Omega \sigma^{\prime}}$, which implies $u_{\sigma^{\prime}, \delta^{\prime}}<u_{\sigma, \delta}$. We define

$$
\begin{equation*}
\bar{u}_{\mathcal{S}, 0}=\lim _{\sigma \rightarrow 0} \lim _{\delta \rightarrow 0} \lim _{n \rightarrow \infty} u_{n, \sigma, \delta} . \tag{6.22}
\end{equation*}
$$

Then $\bar{u}_{\mathcal{S}, 0}$ is a solution of (1.6) in $Q_{\infty}^{\Omega}$. Since $\lim _{t \rightarrow 0} u_{\sigma, \delta}(x, t)=0$ uniformly on any compact subset $K \subset \mathcal{R}=\Omega \backslash \mathcal{S}, \bar{u}_{\mathcal{S}, 0}$ has initial trace 0 on $\mathcal{R}$. If $a \in \mathcal{S}$, we denote by $\tilde{u}_{\infty, a}$ the function defined in $Q_{\infty}^{\Omega_{\sigma}}$ by

$$
\tilde{u}_{\infty, a}(x, t)= \begin{cases}u_{\infty, a}(x, t) & \text { if }(x, t) \in Q_{\infty}^{\Omega} \\ 0 & \text { if }(x, t) \in Q_{\infty}^{\Omega_{\sigma}} \backslash Q_{\infty}^{\Omega}\end{cases}
$$

Then $\tilde{u}_{\infty, a}$ is a subsolution of of (1.6) in $Q_{\infty}^{\Omega_{\sigma}}$ which is smaller than $u_{\infty, a, \sigma}$ which is the limit, when $k \rightarrow \infty$ of the solution $u_{k \delta_{a}, \sigma}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{\infty}^{\Omega_{\sigma}} \\
u & =0 & & \text { in } \partial \Omega_{\sigma} \times(0, \infty)  \tag{6.23}\\
u(., 0) & =k \delta_{a} & & \text { in } \Omega_{\sigma} .
\end{align*}
$$

There holds, by Theorem 6.2,

$$
u_{\sigma, \delta}(x, t) \geq u_{\infty, a, \sigma} \geq \tilde{u}_{\infty, a}(x, t) \quad \forall(x, t) \in Q_{\infty}^{\Omega_{\sigma}}
$$

Letting successively $\delta \rightarrow 0$ and $\sigma \rightarrow 0$ yields to $\bar{u}_{\mathcal{S}, 0} \geq \tilde{u}_{\infty, a}=u_{\infty, a}$ in $Q_{\infty}^{\Omega_{\sigma}}$. Therefore any $a \in \mathcal{S}$ is a singular initial point of $\bar{u}_{\mathcal{S}, 0}$. Since $\mathcal{S} \cup \mathcal{R}=\Omega$, it follows that $\operatorname{tr}_{\Omega}(u)=(\mathcal{S}, 0)$. Since $u_{\sigma, \delta}$ satisfies (4.1) and $\partial \Omega_{\sigma}$ has bounded curvature, independent of $\sigma$, there holds classicaly

$$
\begin{equation*}
\left|\nabla u_{\sigma, \delta}(x, t)\right| \leq c t^{-\frac{q+\alpha}{q-1}} \quad \forall(x, t) \in \partial \Omega_{\sigma} \times(0, \infty) \tag{6.24}
\end{equation*}
$$

If $z \in \partial \Omega$ and then by the mean value theorem there exists $\theta \in(0,1)$ such that,

$$
0 \leq u_{\sigma, \delta}(z, t)=u_{\sigma, \delta}\left(z+\sigma \mathbf{n}_{z}, t\right)-\sigma \nabla u_{\sigma, \delta}\left(z+\theta \sigma \mathbf{n}_{z}, t\right) \cdot \mathbf{n}_{z} \leq c \sigma t^{-\frac{q+\alpha}{q-1}}
$$

This implies that $\bar{u}_{\mathcal{S}, 0}$ vanishes on $\partial \Omega_{\sigma} \times(0, \infty)$.
Let $u$ be any positive solution of (1.6) in $Q_{\infty}^{\Omega}$, vanishing on $\partial \Omega \times(0, \infty)$, with initial trace $(\mathcal{S}, 0)$. For $0<\delta<\sigma$ fixed and for $R, \epsilon>0$, there exists $\tau_{\epsilon}>0$ such that, for any $\tau \in\left(0, \tau_{\epsilon}\right]$,

$$
u(x, \tau) \leq \epsilon \quad \forall x \in \bar{B}_{R} \cap \Omega \backslash \mathcal{S}_{\sigma}
$$

This is due to the fact that $u(x, \tau) \rightarrow 0$ when $\tau \rightarrow 0$, uniformly on compact subset of $\bar{B}_{R} \cap \mathcal{R}$. Assume that $\Omega$ is unbounded (the case where $\Omega$ is bounded is simpler since it does not require to introduce the barrier $w_{B_{R}}$ ) and let $R>0$ large enough so $\Omega^{c} \subset B_{R}$. By (4.1)-(4.3) there exists $0<\tau_{1} \leq \tau_{0}$ such that for any $\tau \in\left(0, \tau_{1}\right]$,

$$
u(x, \tau) \leq w_{B_{R}}(x, \tau)+u_{\sigma, \delta}(x, \tau) \quad \forall x \in \bar{B}_{R} \cap \Omega \cap \mathcal{S}_{\sigma}
$$

Furthermore $u(x, t)<w_{B_{R}}(x, t)$ for all $t>0$ and $x \in \partial B_{R} \cap \Omega$. Since $\epsilon+w_{B_{R}}+u_{\sigma, \delta}$ is a supersolution for (1.6) in $B_{R} \cap \Omega \times(0, \infty)$, it follows that

$$
u(x, t) \leq \epsilon+w_{B_{R}}(x, t)+u_{\sigma, \delta}(x, t) \quad \forall(x, t) \in B_{R} \cap \Omega \times(0, \infty)
$$

Letting successively $\delta \rightarrow 0, \sigma \rightarrow 0, R \rightarrow \infty$ (here we use the fact that $w_{B_{R}}(x, t) \rightarrow 0$ when $R \rightarrow \infty$ by Proposition 3.4) and $\epsilon \rightarrow 0$ yields to $u \leq \bar{u}_{\mathcal{S}, 0}$.
Step 2: Existence of a minimal solution. The set $\mathcal{U}_{\mathcal{S}, 0}(\Omega)$ is not empty since it contains $\bar{u}_{\mathcal{S}, 0}$ and we may define

$$
\begin{equation*}
\tilde{u}_{\mathcal{S}, 0}=\sup \left\{u_{\infty, a}: a \in \mathcal{S}\right\} \tag{6.25}
\end{equation*}
$$

and

$$
\begin{equation*}
\hat{u}_{\mathcal{S}, 0}=\inf \left\{u: u \in \mathcal{U}_{\mathcal{S}, 0}(\Omega)\right\} . \tag{6.26}
\end{equation*}
$$

The functions $\tilde{u}_{\mathcal{S}, 0}$ and $\hat{u}_{\mathcal{S}, 0}$ are respectively positive sub and super solutions of (1.6) in $Q_{\infty}^{\Omega}$. They are bounded from above by $\bar{u}_{\mathcal{S}, 0}$ and from below by $u_{\infty, a}$ for any $a \in \mathcal{S}$. Since $u_{\infty, a} \leq u$ for any $a \in \mathcal{S}$ and $u \in \mathcal{U}_{\mathcal{S}, 0}(\Omega)$, it follows that $\tilde{u}_{\mathcal{S}, 0} \leq \hat{u}_{\mathcal{S}, 0}$. Therefore there exists a solution $\underline{u}_{\mathcal{S}, 0}$ of (1.6) in $Q_{\infty}^{\Omega}$ which satisfies and

$$
\begin{equation*}
\tilde{u}_{\mathcal{S}, 0} \leq \underline{u}_{\mathcal{S}, 0} \leq \hat{u}_{\mathcal{S}, 0} . \tag{6.27}
\end{equation*}
$$

This implies that $\underline{u}_{\mathcal{S}, 0}$ has initial trace $(\mathcal{S}, 0)$, it vanishes on $\partial \Omega \times(0, \infty)$ and it is therefore the minimal element of $\mathcal{U}_{\mathcal{S}, 0}(\Omega)$.
Remark. If dist $\left(\mathcal{S}, \Omega^{c}\right)>0$, it is not needed to replace $\Omega$ by a larger set $\Omega_{\sigma}$ in order to construct the maximal solution. The construction of $\bar{u}_{\mathcal{S}, 0}$ can be done in replacing $u_{n, \sigma, \delta}$ by the solution $u=u_{n, \sigma}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{\infty}^{\Omega} \\
u & =0 & & \text { in } \partial \Omega \times(0, \infty)  \tag{6.28}\\
u(., 0) & =n \chi_{\mathcal{s}_{\delta}} & & \text { in } \Omega,
\end{align*}
$$

with $\delta<\delta_{0}:=\operatorname{dist}\left(\mathcal{S}, \Omega^{c}\right)$.
The next result is an extension of Proposition 2.2.
Proposition 6.8 Assume $\alpha>-1$ and $1<q<q_{c, \alpha}$. Let $\left\{u_{n}\right\}$ be a sequence of positive solutions of (1.6) which converges to $u$ locally uniformly in $Q_{T}^{\Omega}$, and denote by $\left(\mathcal{S}_{n}, \mu_{n}\right)$ and $(\mathcal{S}, \mu)$ the respective initial trace of $u_{n}$ and $u$. If $\mathcal{A}$ is an open subset of $\cap_{n} \mathcal{R}_{n}$ and $\mu_{n}(\mathcal{A})$ remains bounded independently of $n \in \mathbb{N}$ (where $\mathcal{R}_{n}=\Omega \backslash \mathcal{S}_{n}$ and $\mathcal{R}=\Omega \backslash \mathcal{S}$ ), then $\mathcal{A} \subset \mathcal{R}$ and $\chi_{\mathcal{A}} \mu_{n} \rightharpoonup \chi_{\mathcal{A}} \mu$ in the weak sense of measures. Conversely, if $\mathcal{A} \subset \mathcal{R}$, then for any compact $K \subset \mathcal{A}$, there exist $C_{K}>0$ and $n_{K} \in \mathbb{N}$ such that $\mu_{n}(K) \leq C_{K}$ for any $n \geq n_{K}$.
Proof. Clearly (2.14) holds. We keep the notations of the proof of Proposition 2.2 where the first statement has been proved in assuming $\bar{B}_{\tilde{r}}(z) \subset \mathcal{A}$. Since $\mu_{n}(\mathcal{A})$ remains bounded, there exists a subsequence $\left\{n_{j}\right\}$ and a positive measure $\mu^{\prime}$ on $\mathcal{A}$ such that $\mu_{n_{j}} \rightharpoonup \mu^{\prime}$ in the weak sense of measures in $\mathcal{A}$. Then $u_{\chi_{B_{\tilde{r}}(z)} \mu_{n_{j}}}$ converges locally uniformly in $Q_{\infty}^{B_{\tilde{r}}(z)}$ to the solution $u_{\chi_{B_{\tilde{r}}(z)} \mu^{\prime}}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{\infty}^{B_{\tilde{r}}(z)} \\
u & =0 & & \text { in } \partial B_{\tilde{r}}(z) \times(0, \infty)  \tag{6.29}\\
u(., 0) & =\chi_{B_{\tilde{r}}(z)} \mu^{\prime} & & \text { in } B_{\tilde{r}}(z) .
\end{align*}
$$

Since $q<q_{c, \alpha}$, the convergence of $u_{{X_{B_{\bar{r}}(z)}}^{\mu_{n_{j}}}}$ and $t^{\alpha} u_{\chi_{B_{\bar{r}}(z)}^{q} \mu_{n_{j}}}$ respectively to $u_{{X_{B_{\bar{r}}(z)}} \mu_{n_{j}}}$ and $t^{\alpha} u_{\chi_{B_{\bar{r}}(z)} \mu^{\prime}}^{q}$ holds in $L^{1}\left(Q_{T}^{B_{\bar{\tau}}(z)}\right)$ for any $T>0$. Relation (2.21) reads

$$
\begin{equation*}
u_{\chi_{B_{\bar{r}}(z)} \mu_{n}}(x, t) \leq u_{n}(x, t) \leq u_{\chi_{B_{\bar{r}}(z)} \mu_{n}}(x, t)+w_{B_{\bar{r}}(z)}(x, t) \quad \text { in } Q_{\infty}^{B_{\tilde{F}}(z)} . \tag{6.30}
\end{equation*}
$$

(see Proposition 3.3). Then $u_{n_{j}}$ and $t^{\alpha} u_{n_{j}}^{q}$ converge to $u$ and $t^{\alpha} u^{q}$ respectively, in $L^{1}\left(Q_{T}^{B_{r}(z)}\right)$ for any $r<\tilde{r}$. From (2.17), we derive

$$
\begin{equation*}
\iint_{Q_{\infty}^{B_{\infty}(z)}}\left(-u\left(\zeta_{t}+\Delta \zeta\right)+\zeta t^{\alpha} u^{q}\right) d x d t=\int_{B_{r}(z)} \zeta(x, 0) d \mu^{\prime}(x), \tag{6.31}
\end{equation*}
$$

for any $\zeta \in C_{c}^{1,1 ; 1}\left(\bar{Q}_{\infty}^{B_{r}(z)}\right)$ which vanishes for $t$ large enough. This implies that $\mu^{\prime}$ is the initial trace of $u$ in $B_{r}(z)$, i.e. $\chi_{B_{r}(z)} \mu^{\prime}=\chi_{B_{r}(z)} \mu$ and $\chi_{B_{r}(z)} \mu_{n} \rightharpoonup \chi_{B_{r}(z)} \mu$. Using a partition of unity, we conclude that $\chi_{\mathcal{A}} \mu_{n} \rightharpoonup \chi_{\mathcal{A}} \mu$.

Conversely, we assume that there exist a compact set $K \subset \mathcal{A}$ and a subsequence $\mu_{n_{j}}$ such that $\mu_{n_{j}}(K) \rightarrow \infty$. Thus, using the diagonal process, there exist $z \in K$ and another subsequence that we still denote $\mu_{n_{j}}$ such that

$$
\lim _{n_{j} \rightarrow \infty} \mu_{n_{j}}\left(B_{\epsilon}(z)\right)=\infty \quad \forall \epsilon>0 .
$$

Therefore, we can construct a subsequence $\left\{n_{j_{e}}\right\} \subset\left\{n_{j}\right\}$ such that

$$
\mu_{n_{j_{\ell}}}\left(B_{2^{-n_{j_{\ell}}}}(z)\right)=m_{n_{j_{\ell}}} \rightarrow \infty
$$

when $n_{j_{\ell}} \rightarrow \infty$. Since the solution $u_{\chi_{B_{2}-n_{j_{\ell}}}(z)} \mu_{n_{j_{\ell}}}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } Q_{\infty}^{B_{\tilde{r}}(z)} \\
u & =0 & & \text { in } \partial B_{\tilde{r}}(z) \times(0, \infty)  \tag{6.32}\\
u(., 0) & =\chi_{B_{2}-n_{j_{\ell}}(z)} \mu_{n_{j_{\ell}}} & & \text { in } B_{\tilde{r}}(z) .
\end{align*}
$$

converges to $u_{\infty \delta_{z}}^{B_{\bar{F}}(z)}$ which is the limit of the solution $u_{k \delta_{z}}$ of (6.32) with initial data $u(., 0)=$ $k \delta_{z}$, and is dominated by $u_{n_{j_{\ell}}}$ in $Q_{T}^{B_{r}(z)}$ we conclude that $u \geq u_{\infty \delta_{z}}^{B_{\bar{F}}(z)}$ in $Q_{T}^{B_{\bar{z}}(z)}$, which implies that $z \in \mathcal{S}$, contradiction.
Proposition 6.9 Assume $u_{1}$ and $u_{2}$ are two positive solutions of (1.6) in $Q_{\infty}^{\Omega}$ with initial trace $(\mathcal{S}, \mu)$. Then for any $a \in \mathcal{R}$ and $R>0$ such that $B_{R}(a) \subset \mathcal{R}$, there holds

$$
\begin{equation*}
\left|u_{1}(x, t)-u_{2}(x, t)\right| \leq w_{B_{R}}(x-a, t) \quad \forall(x, t) \in Q_{\infty}^{B_{R}(a)} \tag{6.33}
\end{equation*}
$$

In particular $\lim _{t \rightarrow 0}\left|u_{1}(x, t)-u_{2}(x, t)\right|=0$ uniformly on any compact subset of $\mathcal{R}$.
Proof. Since $u$ and $u^{\prime}$ are solution of (1.6) and $B_{R}(a) \in \mathcal{R}$, for any $i=1,2, R^{\prime}<R$ and $T>0$,

$$
\iint_{Q_{T}^{B_{R}^{\prime}}(a)} t^{\alpha} u_{i}^{q}(x, t) d x d t+\iint_{Q_{T}^{B_{R}^{\prime}(a)}} u_{i}(x, t) d x d t<\infty,
$$

furthermore

$$
\lim _{t \rightarrow 0} \int_{B_{R}(a)} u_{i}(x, t) \zeta(x) d x=\int_{B_{R}(a)} \zeta(x) d \mu(x) \quad \forall \zeta \in C_{c}\left(B_{R}(a)\right)
$$

This implies that $u_{i}$ has a Sobolev trace on $\partial_{\ell} Q_{T}^{B_{R}^{\prime}(a)}$ which belongs to $L^{1}$ and they are the limit, when $k \rightarrow \infty$ of the solutions $u_{i, k}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha}|u|^{q-1} u & =0 & & \text { in } Q_{\infty}^{B_{R}^{\prime}(a)} \\
u & =\min \left\{k,\left.u_{i}\right|_{Q_{\infty}^{B_{R}^{\prime}(a)}}\right\} & & \text { in } \partial_{\ell} Q_{\infty}^{B_{R}^{\prime}(a)}  \tag{6.34}\\
u(., 0) & =\mu & & \text { in } B_{R}^{\prime}(a) .
\end{align*}
$$

Since $u_{2, k}+w_{B_{R}^{\prime}}(.-a)$ is a supersolution

$$
u_{1, k} \leq u_{2, k}+w_{B_{R}^{\prime}}(.-a) \Longrightarrow\left|u_{1, k}-u_{2, k}\right| \leq w_{B_{R}^{\prime}}(.-a) \text { in } Q_{\infty}^{B_{R}^{\prime}(a)}
$$

Letting $k \rightarrow \infty, R^{\prime}$ to $R$, we derive (6.33). The second statement is a consequence of the fact that $\lim _{t \rightarrow 0} w_{B_{R}}(.-a)=0$, uniformly on $B_{R}^{\prime}$ by Proposition 3.3.
Remark. The previous estimate does not use the fact that $\partial \Omega$ is smooth and bounded. If the $u_{i}$ belong to $\mathcal{U}_{\mathcal{S}, \mu}(\Omega)$, estimate (6.33) can be improved since the $u_{i}$ vanish on $\partial_{\ell} Q_{\infty}^{B_{R}(a)}$, and we obtain,

$$
\begin{equation*}
\left|u_{1}(x, t)-u_{2}(x, t)\right| \leq \min \left\{w_{B_{R}}(x-a, t), c_{\alpha} t^{-\frac{1+\alpha}{q-1}}\right\} \quad \forall(x, t) \in Q_{\infty}^{B_{R}(a)} \tag{6.35}
\end{equation*}
$$

Proposition 6.10 Assume $\Omega \subseteq \mathbb{R}^{N}$ is either $\mathbb{R}^{N}$ or an open domain with a $C^{2}$ compact boundary, $\alpha>-1$ and $1<q<q_{c, \alpha}$. Then for any measure $\mu$ in $\Omega$ such that $\mu\left\lfloor\Omega_{R} \in\right.$ $\mathfrak{M}_{+}^{b, \rho}\left(\Omega_{R}\right)$ where $\Omega_{R}=\Omega \cap B_{R}$, there exists a unique solution $u_{\mu}$ to

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha}|u|^{q-1} u=0 & \text { in } Q_{\infty}^{\Omega} \\
u=0 & \text { in } \partial_{\ell} Q_{\infty}^{\Omega}  \tag{6.36}\\
u(., 0)=\mu & \text { in } \Omega,
\end{align*}
$$

and the mapping $\mu \mapsto u_{\mu}$ is increasing. Furthermore, if $\left\{\mu_{n}\right\}$ is a sequence of positive measures such that $\mu_{n}\left\lfloor\Omega_{R} \in \mathfrak{M}_{+}^{b, \rho}\left(\Omega_{R}\right)\right.$ which converges weakly to $\mu \Omega_{R} \in \mathfrak{M}_{+}^{b, \rho}\left(\Omega_{R}\right)$, then $\left\{u_{\mu_{n}}\right\} \rightarrow u_{\mu}$ locally uniformly in $Q_{\infty}^{\Omega}$.
Proof. We recall that $u$ is a solution of (6.36) if $u \in L_{l o c}^{1}\left(\overline{Q_{\infty}^{\Omega}}\right),|u|^{q} \in L_{l o c}^{1}\left(\overline{Q_{\infty}^{\Omega}} ; t^{\alpha} \rho d x d t\right)$ satisfies

$$
\begin{equation*}
\iint_{Q_{\infty}^{\Omega}}\left(-u\left(\partial_{t} \zeta+\Delta \zeta\right)+t^{\alpha}|u|^{q-1} u \zeta\right) d x d t=\int_{\Omega} \zeta(., 0) d \mu \tag{6.37}
\end{equation*}
$$

for any test function $\zeta \in C_{0}^{2,1 ; 1}\left(\overline{Q_{\infty}^{\Omega}}\right)$. When $\Omega$ is bounded, existence, uniqueness and stability are proved in [9]. Thus we assume that $\Omega$ is unbounded and we assume $R \geq R_{0}$ such that $\Omega^{c} \subset B_{R_{0}}$. There exists a unique solution $u_{R}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha}|u|^{q-1} u=0 & \text { in } Q_{\infty}^{\Omega_{R}} \\
u=0 & \text { in } \partial_{\ell} Q_{\infty}^{\Omega_{R}}  \tag{6.38}\\
u(., 0)=\mu\left\lfloor\Omega_{R}\right. & \text { in } \Omega_{R} .
\end{align*}
$$

The function $u_{R}$ is nonnegative, $R \mapsto u_{R}$ is increasing. For $R>R_{1}, u_{R}$ admits a Sobolev trace $f_{R_{1}}$ on $\partial B_{R_{1}} \times(0, T)$ which is an integrable function, and $u_{R}$ is the unique solution of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha}|u|^{q-1} u=0 & \text { in } Q_{\infty}^{\Omega_{R_{1}}} \\
u=0 & \text { in } \partial \Omega \times(0, \infty)  \tag{6.39}\\
u=f_{R_{1}} & \text { in } \partial B_{R_{1}} \times(0, \infty) \\
u(., 0)=\mu\left\lfloor\Omega_{R_{1}}\right. & \text { in } \Omega_{R_{1}} .
\end{align*}
$$

Furthermore, $u_{R} \bigsqcup_{Q_{\infty} \Omega_{R_{1}}}=\lim _{m \rightarrow \infty} v_{m}$, where $v_{m}$ is the unique solution of (6.36) where the boundary data on $\partial B_{R_{1}} \times(0, \infty)$ is replaced by $f_{R_{1}, m}=f_{R_{1}} \wedge m\left(m \in \mathbb{N}^{*}\right)$. Let $v_{R_{1}}$ be the unique solution of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha}|u|^{q-1} u=0 & \text { in } Q_{\infty}^{\Omega_{R_{1}}} \\
u=0 & \text { in } \partial \Omega \times(0, \infty)  \tag{6.40}\\
u=0 & \text { in } \partial B_{R_{1}} \times(0, \infty) \\
u(., 0)=\mu\left\lfloor\Omega_{R_{1}}\right. & \text { in } \Omega_{R_{1}},
\end{align*}
$$

If $w_{B_{R_{1}}}$ is the barrier function in $Q_{\infty}^{\Omega_{R_{1}}}$ which has been constructed in Proposition 3.3, $v_{R_{1}}+w_{B_{R_{1}}}$ is a supersolution for problem (6.39). Since it is larger than $v_{m}$ in $Q_{\infty}^{\Omega_{R_{1}}}$ for any $m>0$, there holds $u_{R} \leq v_{R_{1}}+w_{B_{R_{1}}}$, for any $R>R_{1}$. Then $u_{R} \uparrow u_{\mu}$ which is a solution of 1.6 in $Q_{\infty}^{\Omega_{R}}$. By Proposition 3.4, $w_{B_{R_{1}}}$ remains uniformly bounded in $Q_{\infty}^{\Omega_{R^{\prime}}}$ for any $R_{0}<R^{\prime}<R_{1}$. Therefore $u_{\mu}$ shares the same property. If $\zeta \in C_{c}^{1,1 ; 1}\left(Q_{\infty}^{\bar{\Omega}}\right)$ vanishes on $\left.Q_{\infty}^{\partial \Omega}\right)$ and for $|x|>R^{\prime}>R_{0}$, there holds for $R>R^{\prime}>R_{0}$ and $T>0$,

$$
\begin{equation*}
\iint_{Q_{\infty}^{\Omega}}\left(-u_{R}(x, t)\left(\partial_{t} \zeta+\Delta \zeta\right)+\zeta t^{\alpha} u_{R}^{q}\right) d x d t=\int_{\Omega} \zeta(x, 0) d \mu(x)-\int_{\Omega} \zeta(x, T) u_{R}(x, T) d x \tag{6.41}
\end{equation*}
$$

If we let $R \rightarrow \infty$ we deduce by the monotone convergence theorem that $u_{\mu}$ is a weak solution of (6.36). This proves existence.

For uniqueness, we consider $u_{\mu}$ and $u_{\mu}^{\prime}$ two solutions of (6.36). By the same argument as in the existence part, for any $R>0, u_{\mu}$ is smaller than the supersolution $u_{\mu}^{\prime}+w_{B_{R_{1}}}$ in $Q_{\infty}^{\Omega_{R}}$. Since $\lim _{R \rightarrow \infty} w_{B_{R}}=0$ by Proposition 3.4 we obtain $u_{\mu} \leq u_{\mu}^{\prime}$. Similarly $u_{\mu}^{\prime} \leq u_{\mu}$. Uniqueness implies the monotonicity of the mapping $\mu \mapsto u_{\mu}$.

For proving the stability, assume $\left\{\mu_{n}\left\lfloor\Omega_{R}\right\}\right.$ converges to $\mu \varrho_{\Omega_{R}}$ in the weak sense of measures in $\mathfrak{M}_{+}^{b, \rho}\left(\Omega_{R}\right)$ for any $R>R_{0}$. Then the sequence of solutions $v_{n, R}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha}|u|^{q-1} u=0 & \text { in } Q_{\infty}^{\Omega_{R}} \\
u=0 & \text { in } \partial \Omega \times(0, \infty)  \tag{6.42}\\
u=0 & \text { in } \partial B_{R} \times(0, \infty) \\
u(., 0)=\mu_{n}\left\lfloor\Omega_{R}\right. & \text { in } \Omega_{R},
\end{align*}
$$

converges to the solution $v_{R}$ of (6.40) with $R_{1}=R$. In particular $v_{n, R} \rightarrow v_{R}$ and $t^{\alpha} v_{n, R}^{q} \rightarrow$ $t^{\alpha} v_{R}^{q}$ in $L^{1}\left(Q_{T}^{\Omega_{R}}\right)$ and by standard regularity result the convergence of $v_{n, R}$ towards $v_{R}$ holds uniformly on $\bar{\Omega}_{R} \times[\epsilon, T]$ for any $0<\epsilon<T$. Furthermore $u_{\mu_{n}} \leq v_{n, R}+W_{B_{R}}$ in $Q_{\infty}^{\Omega_{R}}$. This
jointly with standard local regularity results for heat equation, implies that $\left\{u_{\mu_{n}}\right\}$ remains uniformly bounded and hence relatively compact for the topology of uniform convergence on any compact set of $\bar{\Omega} \times[\epsilon, \infty)$. Thus there exist a subsequence $\left\{u_{\mu_{n_{k}}}\right\}$ and a function $u^{*} \in$ $\left.C^{( } \bar{Q}_{\infty}^{\Omega}\right)$ such that $u_{\mu_{n_{k}}} \rightarrow u^{*}$ locally uniformly in $\bar{\Omega} \times(0, \infty)$. Since $t^{\alpha} u_{\mu_{n}}^{q} \leq t^{\alpha} v_{n, R}^{q}+t^{\alpha} W_{B_{R}}^{q}$, there also holds by the dominated convergence theorem $t^{\alpha} u_{\mu_{n_{k}}}^{q} \rightarrow t^{\alpha} u^{* q}$ in $L_{l o c}^{1}(\bar{\Omega} \times[0, \infty))$. Henceforth letting $n_{k} \rightarrow \infty$ in the expression

$$
\begin{align*}
\iint_{Q_{\infty}^{\Omega}}\left(-u_{\mu_{n_{k}}}(x, t)\left(\partial_{t} \zeta\right.\right. & \left.+\Delta \zeta)+\zeta t^{\alpha} u_{\mu_{n_{k}}}^{q}\right) d x d t  \tag{6.43}\\
& =\int_{\Omega} \zeta(x, 0) d \mu_{n_{k}}(x)-\int_{\Omega} \zeta(x, T) u_{n_{k}}(x, T) d x
\end{align*}
$$

where $\zeta \in C_{c}^{1,1 ; 1}\left(Q_{\infty}^{\bar{\Omega}}\right)$, we conclude that $u^{*}=u_{\mu}$ and that $u_{\mu_{n}} \rightarrow u_{\mu}$.
Proposition 6.11 Assume $F$ is a non-empty relatively closed subset of $\Omega, \mathcal{R}=\Omega \backslash F$ and $\mu \in \mathfrak{M}_{+}(\mathcal{R})$. If we set

$$
\begin{equation*}
\partial_{\mu} F=\left\{z \in F: \mu\left(\mathcal{R} \cap B_{r}(z)\right)=\infty, \forall r>0\right\} \tag{6.44}
\end{equation*}
$$

then $\partial_{\mu} F$ is relatively closed in $\Omega$. If $\mathcal{R}^{\mu}=\Omega \backslash \partial_{\mu} F$, it contains $\mathcal{R}$ and if $\mu^{*}$ is the measure defined in $\mathcal{R}^{\mu}$ by $\mu$ on $\mathcal{R}$ and 0 in $\mathcal{R}^{\mu} \cap \mathcal{R}^{c}$, then there exist a minimal positive solution $\underline{u}_{\mu^{*}}$ and a maximal solution $\bar{u}_{\mu^{*}}$ of (1.6) vanishing on $\partial \Omega \times(0, \infty)$ satisfying $\operatorname{tr}_{\Omega}(u)=$ $\left.\stackrel{u}{\mu}_{\mu} F, \mu^{*}\right)$. Furthermore $\underline{u}_{\mu^{*}}$ and $\bar{u}_{\mu^{*}}$ are respectively the minimal and the maximal element of $\mathcal{U}_{\partial_{\mu} F, \mu^{*}}(\Omega)$.

Proof. The set $\partial_{\mu} F$ is the blow-up set of the measure $\mu$. It is clearly a relatively closed subset of $\Omega$ included into $\overline{\mathcal{R}} \backslash \mathcal{R}$.
Step 1: Existence of a minimal solution. For $\delta>0$, we denote $\left(\partial_{\mu} F\right)_{\delta}:=\{x \in \Omega$ : $\left.\operatorname{dist}\left(x, \partial_{\mu} F\right) \leq \delta\right\}$ and $\mathcal{R}_{\delta}^{\mu}=\Omega \backslash\left(\partial_{\mu} F\right)_{\delta} \subset \mathcal{R}^{\mu}$. We define the Radon measure $\mu_{\delta}$ on $\Omega$ by

$$
\mu_{\delta}= \begin{cases}\mu & \text { on } \mathcal{R}_{\delta}^{\mu} \cap \mathcal{R} \\ 0 & \text { on } F \cup\left(\partial_{\mu} F\right)_{\delta}\end{cases}
$$

Then $\mu_{\delta}$ is a positive Radon measure in $\Omega$ and by Proposition 6.10 problem (6.36) with initial data $\mu_{\delta}$ admits a unique positive solution $u_{\mu_{\delta}}$. Furthermore the mapping $\delta \mapsto u_{\mu_{\delta}}$ is nonincreasing, and we set $\underline{u}_{\mu^{*}}=\lim _{\delta \rightarrow 0} u_{\mu_{\delta}}$. Then $\underline{u}_{\mu^{*}}$ is a positive solution of (1.6) in $Q_{\infty}^{\Omega}$ which vanishes on $\partial \Omega \times(0, \infty)$ and has initial trace $\left(\mathcal{S}^{\prime}, \mu^{\prime}\right)$. If $a \in \mathcal{R}^{\mu}$, there exists $R>0$ such that $\bar{B}_{R}(a) \subset \mathcal{R}^{\mu}$ and $\delta_{a}>0$ such that $\bar{B}_{R}(a) \subset \mathcal{R}_{\delta}^{\mu}$ for $0<\delta<\delta_{a}$, that we assume below. By the maximum principle there holds

$$
v_{\mu_{\delta}}(x, t) \leq u_{\delta}(x, t) \leq v_{\mu_{\delta}}(x, t)+w_{B_{R}}(x-a, t) \quad \text { in } Q_{\infty}^{B_{R}(a)}
$$

where $v_{\mu_{\delta}}$ is the solution of

$$
\begin{array}{cl}
\partial_{t} u-\Delta u+t^{\alpha} u^{q}=0 & \text { in } Q_{\infty}^{B_{R}(a)} \\
u \geq 0 & \text { in } Q_{\infty}^{B_{R}(a)} \\
u(x, t)=0 & \text { in } \partial B_{R}(a) \times(0, \infty) \\
u(., 0)=\mu_{\delta} \chi_{B_{R}(a)} & \text { in } B_{R}(a) .
\end{array}
$$

Letting $\delta \rightarrow 0$, then $\mu_{\delta} \chi_{B_{R}(a)} \uparrow \mu^{*} \chi_{B_{R}(a)}$, which yields to

$$
v_{\mu^{*}\left\lfloor_{B_{R}}(a)\right.}(x, t) \leq \underline{u}_{\mu^{*}}(x, t) \leq v_{\mu^{*}\left\lfloor_{B_{R}}(a)\right.}(x, t)+w_{B_{R}}(x-a, t) \quad \text { in } Q_{\infty}^{B_{R}(a)} .
$$

By a partition of unity, it implies that for any $\zeta \in C_{c}\left(\mathcal{R}^{\mu}\right)$, we have

$$
\lim _{t \rightarrow 0} \int_{\mathcal{R}^{\mu}} \underline{u}_{\mu^{*}}(x, t) \zeta(x) d x=\int_{\mathcal{R}^{\mu}} \zeta d \mu^{*}(x) .
$$

Therefore $\mathcal{R}^{\mu} \subset \mathcal{R}^{\prime}$ and $\mu^{\prime}\left\lfloor_{\mathcal{R}^{\mu}}=\mu^{*}\right.$. If $z \in \mathcal{R}^{\prime} \cap\left(\mathcal{R}^{\mu}\right)^{c}=\mathcal{R}^{\prime} \cap \partial_{\mu} F, \mu^{*}\left(B_{r}(z) \cap \mathcal{R}^{\mu}\right)=$ $\mu\left(B_{r}(z) \cap \mathcal{R}\right)=\infty$ for any $r>0$ while there exists $r_{0}>0$ such that $\mu^{\prime}\left(B_{r_{0}}(z)\right)<\infty$. By Proposition 6.8, for any $r^{\prime}<r_{0}$ there exists $C>0$ such that $\mu_{\delta}\left(\bar{B}_{r}(z)\right) \leq C$. By the monotone convergence theorem, it implies $\mu\left(\bar{B}_{r}(z) \cap \mathcal{R}\right) \leq C$, which contradicts the definition of $\partial_{\mu} F$. Thus $\mathcal{S}^{\prime}=\partial_{\mu} F$ and $\operatorname{tr}_{\Omega}\left(\underline{u}_{\mu^{*}}\right)=\left(\partial_{\mu} F, \mu^{*}\right)$.

Let us assume that $\Omega$ is unbounded, $R_{0}$ is such that $\Omega^{c} \subset B_{R_{0}}$ and $\Omega_{R}=\Omega \cap B_{R}$ for $R>R_{0}$. Let $\delta, \epsilon>0$, there exists $\tau_{\epsilon}$ such that

$$
\begin{equation*}
u_{\mu_{\delta}}(x, \tau) \leq \epsilon \quad \forall(x, \tau) \in \Omega_{R} \cap\left(\partial_{\mu} F\right)_{\frac{\delta}{2}} \times\left[0, \tau_{\epsilon}\right] \tag{6.45}
\end{equation*}
$$

Let $u \in \mathcal{U}_{\partial_{\mu} F, \mu^{*}}(\Omega)$. In order to compare $u_{\mu_{\delta}}$ and $v:=u+w_{B_{R}}+\epsilon$ in $\Omega_{R} \backslash\left(\partial_{\mu} F\right)_{\frac{\delta}{2}} \times\left(0, \tau_{\epsilon}\right]$ we see that

$$
u_{\mu_{\delta}}(., 0)=\chi_{\mathcal{R}_{R}^{\mu} \backslash\left(\partial_{\mu} F\right)_{\delta}} \mu^{*} \leq u(., 0)=\chi_{\mathcal{R}_{R}^{\mu} \backslash\left(\partial_{\mu} F\right)_{\frac{\delta}{2}}} \mu^{*}
$$

and both are bounded Radon measures. Since $u_{\mu_{\delta}}(t, 0) \leq v(x, t)$ in $\partial\left(\Omega_{R} \backslash\left(\partial_{\mu} F\right)_{\frac{\delta}{2}}\right) \times\left(0, \tau_{\epsilon}\right]$ and $v$ is a supersolution, it follows that $u_{\mu_{\delta}} \leq v$ in $\Omega_{R} \backslash\left(\partial_{\mu} F\right)_{\frac{\delta}{2}} \times\left(0, \tau_{\epsilon}\right]$. Using (6.45) we conclude that

$$
\begin{equation*}
u_{\mu_{\delta}}(x, \tau) \leq v(x, \tau) \quad \forall(x, \tau) \in \Omega_{R} \times\left[0, \tau_{\epsilon}\right] . \tag{6.46}
\end{equation*}
$$

Next, applying the comparison principle in $\Omega_{R} \times\left[\tau^{*}, \infty\right)$ between the solution $u_{\mu_{\delta}}$ and the supersolution $u+\epsilon+w_{B_{R}}$, we conclude that (6.46) holds in $\Omega_{R} \times[\tau, \infty)$ and thus in $Q_{\infty}^{B_{R}}$. Letting successively $R \rightarrow \infty, \epsilon \rightarrow 0$ and $\delta \rightarrow 0$, we conclude that $\underline{u}_{\mu^{*}} \leq u$, thus $u_{\mu^{*}}$ is the minimal element of $\mathcal{U}_{\partial_{\mu} F, \mu^{*}}(\Omega)$.
Step 2: Existence of a maximal solution. Let $\delta>0$ and $u \in \mathcal{U}_{\partial_{\mu} F, \mu^{*}}(\Omega)$. By Proposition 6.9, for any $R>0$ and $\epsilon>0$ there exists $\tau_{\epsilon}$ such that

$$
\begin{equation*}
u(x, t) \leq \underline{u}_{\mu^{*}}(x, t)+\epsilon \quad \forall(x, t) \in \Omega_{R} \backslash\left(\partial_{\mu} F\right)_{\delta} \times\left(0, \tau_{\epsilon}\right], \tag{6.47}
\end{equation*}
$$

and by (4.1), $u(x, t) \leq c_{\alpha} t^{-\frac{1+\alpha}{q-1}}$. Let $\tau \in\left(0, \tau_{\epsilon}\right]$ and $w_{\delta, \tau}$ be the solution of (6.36) in $Q_{\infty}^{\Omega_{R}}$ with initial data $\mu$ replaced by

$$
h_{\delta, \tau}= \begin{cases}\underline{u}_{\mu^{*}}(x, \tau) & \text { if } x \in \Omega_{R} \backslash\left(\partial_{\mu} F\right)_{\delta}  \tag{6.48}\\ c_{\alpha} \tau^{-\frac{1+\alpha}{q-1}} & \text { if } x \in \Omega_{R} \cap\left(\partial_{\mu} F\right)_{\delta}\end{cases}
$$

By (6.47), (6.48) and the maximum principle,

$$
\begin{equation*}
u(x, t+\tau) \leq w_{\delta, \tau}(x, t)+\epsilon+w_{B_{R}}(x, t+\tau) \quad \forall(x, t) \in Q_{\infty}^{\Omega_{R}} . \tag{6.49}
\end{equation*}
$$

Let $\bar{u}_{\left(\partial_{\mu} F\right)_{\delta}}$ be the maximal element of $\mathcal{U}_{\left(\partial_{\mu} F\right)_{\delta}, 0}(\Omega)$, which exists by Proposition 6.7. Then, by (4.3), for any $\delta^{\prime}>\delta$, there exists $\tau_{\delta^{\prime}} \in\left(0, \tau_{\epsilon}\right]$ such that for any $\tau \in\left(0, \tau_{\delta^{\prime}}\right]$

$$
\begin{equation*}
\max \left\{\underline{u}_{\mu^{*}}(., .+\tau), \bar{u}_{\left(\partial_{\mu} F\right)_{\delta}}(., .+\tau)\right\} \leq w_{\delta, \tau} \leq \underline{u}_{\mu^{*}}(., .+\tau)+\bar{u}_{\left(\partial_{\mu} F\right)_{\delta^{\prime}}}(., .+\tau) \quad \text { in } Q_{\infty}^{\Omega_{R}} \tag{6.50}
\end{equation*}
$$

Up to a sequence $\left\{\tau_{n}\right\}$ converging to $0,\left\{w_{\delta, \tau_{n}}\right\}$ converges, locally uniformly in $Q_{\infty}^{\Omega}$ to a solution $w_{\delta}$ of (1.6) in $Q_{\infty}^{\Omega}$ which satisfies

$$
\begin{equation*}
\max \left\{\underline{u}_{\mu^{*}}, \bar{u}_{\left(\partial_{\mu} F\right)_{\delta}}\right\} \leq w_{\delta} \leq \underline{u}_{\mu^{*}}+\bar{u}_{\left(\partial_{\mu} F\right)_{\delta^{\prime}}} \quad \text { in } Q_{\infty}^{\Omega_{R}} . \tag{6.51}
\end{equation*}
$$

We can replace $\delta^{\prime}$ by $\delta$ in this inequality, this proves that $w_{\delta}$ vanishes on $\partial_{\ell} Q_{\infty}^{\Omega_{R}}$ has initial trace $\left(\mu_{\delta},\left(\partial_{\mu} F\right)_{\delta}\right)$, therefore (6.49) becomes

$$
\begin{equation*}
u(x, t) \leq w_{\delta}(x, t)+\epsilon+w_{B_{R}}(x, t) \quad \forall(x, t) \in Q_{\infty}^{\Omega_{R}} \tag{6.52}
\end{equation*}
$$

Letting successively $R \rightarrow \infty$ and $\epsilon \rightarrow 0$ we deduce that $w_{\delta}$ is larger that any $u \in \mathcal{U}_{\partial_{\mu} F, \mu^{*}}(\Omega)$ in $Q_{\infty}^{\Omega_{R}}$. Since $h_{\delta, \tau}$ decreases with $\delta, w_{\delta}$ shares this property and the limit, denoted by $\bar{u}_{\mu^{*}}$ is a solution of (1.6) in $Q_{\infty}^{\Omega}$ which vanishes on $\partial_{\ell} Q_{\infty}^{\Omega_{R}}$ which is large than $u$, thus it is the maximal element of $\mathcal{U}_{\partial_{\mu} F, \mu^{*}}(\Omega)$.

Proposition 6.12 Under the assumptions of Proposition 6.11, we set $F_{\delta}:=\{x \in \Omega$ : $\operatorname{dist}(x, F) \leq \delta\}$ and $\mathcal{R}_{\delta}:=\Omega \backslash F_{\delta} \subset \mathcal{R}$. If we define the measure $\tilde{\mu}_{\delta}$ in $\Omega$ by

$$
\tilde{\mu}_{\delta}= \begin{cases}\mu & \text { on } \mathcal{R}_{\delta} \\ 0 & \text { on } F_{\delta}\end{cases}
$$

then $u_{\tilde{\mu}_{\delta}} \uparrow \underline{u}^{*}$ when $\delta \downarrow 0$.
Proof. There holds $\tilde{\mu}_{\delta} \leq \mu_{\delta}$ which implies $u_{\tilde{\mu}_{\delta}} \leq u_{\mu_{\delta}}$. When $\delta \rightarrow 0, u_{\tilde{\mu}_{\delta}} \uparrow \underline{u}^{*} \leq \underline{u}_{\mu^{*}}$, thus $\underline{u}^{*}$ is a positive solution of (1.6) in $Q_{\infty}^{\Omega}$ which vanishes on $\partial_{\ell} Q_{\infty}^{\Omega}$. Then $\operatorname{tr}_{\Omega}\left(\underline{u}^{*}\right)=\left(\mathcal{S}^{\prime \prime}, \mu^{\prime \prime}\right)$ and $\mathcal{S}^{\prime \prime} \subset \partial_{\mu} F$ and $\mu^{\prime \prime} \leq \mu^{*}$ on $\mathcal{R}^{\mu}$. Furthermore $\mu^{\prime \prime}=\mu=\mu^{*}$ on $\mathcal{R}$, as in the proof of Proposition 6.11. Since $\mu^{*}=0$ on $\mathcal{R}^{\mu} \backslash \mathcal{R}$ it follows that $\mu^{\prime \prime}=\mu^{*}$ on $\mathcal{R}^{\mu}$. Let $a \in \mathcal{R}^{\prime \prime} \cap \partial_{\mu} F$ and $R>0$ such that $\bar{B}_{R}(a) \subset \mathcal{R}^{\prime \prime}$. Then $\mu^{\prime \prime}\left(\bar{B}_{R}(a)\right)<\infty$. Therefore

$$
\mu^{\prime \prime}\left(\bar{B}_{R}(a) \cap \mathcal{R}\right)=\mu\left(\bar{B}_{R}(a) \cap \mathcal{R}\right)<\infty
$$

contradiction. Thus $\mathcal{S}^{\prime \prime}=\partial_{\mu} F$ and $\operatorname{tr}_{\Omega}\left(\underline{u}^{*}\right)=\left(\partial_{\mu} F, \mu^{*}\right)$. Since $\underline{u}^{*} \leq \underline{u}_{\mu^{*}}$ and $\underline{u}_{\mu^{*}}$ is minimal, it follows that $\underline{u}^{*}=\underline{u}_{\mu^{*}}$.

Proposition 6.13 Assume $v \in C(\bar{\Omega} \times(0, \infty))$ is a positive sub-solution (resp. supersolution) of (1.6) in $Q_{\infty}^{\Omega}$ which vanishes on $\partial \ell Q_{\infty}^{\Omega}$. Then there exists a minimal solution $\pi_{+}(v)$ larger than $v$ (resp. a maximal solution $\pi_{-}(v)$ smaller than $v$ and vanishing on $\left.\partial \ell Q_{\infty}^{\Omega}\right)$.

Proof. Assume $v$ is a subsolution. Let $\tau>0$ and let $u_{\tau}$ be the solution of

$$
\begin{align*}
\partial_{t} u-\Delta u+t^{\alpha} u^{q} & =0 & & \text { in } \Omega \times(\tau, \infty) \\
u \geq 0, u & =0 & & \text { in } \partial \Omega \times(\tau, \infty)  \tag{6.53}\\
u(., \tau) & =v(., \tau) & & \text { in } \Omega .
\end{align*}
$$

Existence and uniqueness follows from Proposition 6.10. Furthermore $u_{\tau} \geq v$ in $\Omega \times(\tau, \infty)$. This implies that for $0<\tau<\tau^{\prime}, u_{\tau} \geq u_{\tau^{\prime}}$. Since $u_{\tau}(x, t) \leq c_{\alpha}(t-\tau)^{-\frac{1+\alpha}{q-1}}$, there exists $\pi_{+}(v)=\lim _{\tau \rightarrow 0} u_{\tau}$, and $\pi_{+}(v)$ is a positive solution of (1.6) in $Q_{\infty}^{\Omega}$ and is larger than $v$. If $u$ is any positive solution of (1.6) in $Q_{\infty}^{\Omega}$, vanishing on $\partial \ell Q_{\infty}^{\Omega}$ and larger than $v$, for any $\tau>0$ it is larger than $u(., \tau)$, thus it is larger than $u_{\tau}$ on $\Omega \times(\tau, \infty)$. Therefore $u \geq \pi_{+}(v)$.

Assume now that $u$ is a supersolution. We define $u_{\tau}$ by (6.53). Then $u_{\tau} \leq v$ and $u_{\tau} \leq u_{\tau^{\prime}}$ for $0<\tau<\tau^{\prime}$. Then $\pi_{-}(v)=\lim _{\tau \rightarrow 0} u_{\tau}$, and $\pi_{-}(v)$ is a positive solution of (1.6) in $Q_{\infty}^{\Omega}$ and is smaller than $v$, and thus vanishing on $\partial \ell Q_{\infty}^{\Omega}$. Similarly as above $\pi_{-}(v)$ is larger than any positive solution smaller than $v$.

Theorem 6.14 Assume $\Omega \subseteq \mathbb{R}^{N}$ is either $\mathbb{R}^{N}$ or an open domain with a $C^{2}$ compact boundary, $\alpha>-1$ and $1<q<q_{c, \alpha}$. Then for any $\nu \approx(\mathcal{S}, \mu) \in \mathfrak{B}^{\text {reg }}(\Omega)$ there exist a maximal positive solution $\bar{u}_{\mathcal{S}, \mu}$ and a minimal positive solution $\underline{u}_{\mathcal{S}, \mu}$ of (1.6) in $Q_{\infty}^{\Omega}$ vanishing on $\partial \Omega \times(0, \infty)$ with initial trace $\nu$. Furthermore, if $\inf \left\{\left|z-z^{\prime}\right|: z \in \mathcal{S}, z^{\prime} \in \Omega^{c}\right\}>0$, then

$$
\begin{equation*}
\bar{u}_{\mathcal{S}, \mu}-\underline{u}_{\mathcal{S}, \mu} \leq \bar{u}_{\mathcal{S}, 0}-\underline{u}_{\mathcal{S}, 0} . \tag{6.54}
\end{equation*}
$$

Proof. Step 1: Construction of the maximal and minimal solutions. The functions $\underline{u}_{\mu^{*}}, \underline{u}_{\mathcal{S}, 0}$, $\bar{u}_{\mu^{*}}$ and $\bar{u}_{\mathcal{S}, 0}$ have been defined in Proposition 6.11 and Proposition 6.7. Since $\sup \left\{\underline{u}_{\mu^{*}}, \underline{u}_{\mathcal{S}, 0}\right\}$ is a subsolution of (1.6) which is smaller than the supersolution $\bar{u}_{\mu^{*}}+\bar{u}_{\mathcal{S}, 0}$ we set

$$
\begin{equation*}
\text { (i) } \underline{u}_{\mathcal{S}, \mu}=\pi_{+}\left(\sup \left\{\underline{u}_{\mu^{*}}, \underline{u}_{\mathcal{S}, 0}\right\}\right) \quad \text { and } \quad \text { (ii) } \quad \bar{u}_{\mathcal{S}, \mu}=\pi_{-}\left(\bar{u}_{\mu^{*}}+\bar{u}_{\mathcal{S}, 0}\right) . \tag{6.55}
\end{equation*}
$$

Then $\underline{u}_{\mathcal{S}, \mu}$ and $\bar{u}_{\mathcal{S}, \mu}$ are solutions which satisfy

$$
\begin{equation*}
\sup \left\{\underline{u}_{\mu^{*}}, \underline{u}_{\mathcal{S}, 0}\right\} \leq \underline{u}_{\mathcal{S}, \mu} \leq \bar{u}_{\mathcal{S}, \mu} \leq \bar{u}_{\mu^{*}}+\bar{u}_{\mathcal{S}, 0} \tag{6.56}
\end{equation*}
$$

Therefore $\underline{u}_{\mathcal{S}, \mu}$ and $\bar{u}_{\mathcal{S}, \mu}$ vanish on $\partial_{\ell} Q_{\infty}^{\Omega}$, they have initial trace $\mu$ on $\mathcal{R}$ and are larger than any $u_{\infty, a}$ for $a \in \mathcal{S}$ (notice that $\partial_{\mu} \mathcal{S} \subset \mathcal{S}$ ). This implies that $\underline{u}_{\mathcal{S}, \mu}$ and $\bar{u}_{\mathcal{S}, \mu}$ belong to $\mathcal{U}_{\mathcal{S}, \mu}(\Omega)$.

Let $u \in \mathcal{U}_{\mathcal{S}, \mu}(\Omega)$. Then for $\sigma>\delta>0$ and $\epsilon, R>0$, there exists $\tau_{1}>0$ such that

$$
u(x, t) \leq u_{\sigma, \delta}(x, t) \text { for }(x, t) \in\left(\mathcal{S}_{\frac{\delta}{2}} \cap \Omega_{R}\right) \times\left(0, \tau_{1}\right] .
$$

There exists $\tau_{2} \in\left(0, \tau_{1}\right]$ such that

$$
u(x, t) \leq \bar{u}_{\mu^{*}}(x, t)+\epsilon+w_{B_{R}}(x, t) \text { for }(x, t) \in\left(\Omega_{R} \backslash \mathcal{S}_{\frac{\delta}{2}}\right) \times\left(0, \tau_{2}\right]
$$

Therefore

$$
u(x, t) \leq u_{\sigma, \delta}(x, t)+\bar{u}_{\mu^{*}}(x, t)+\epsilon+w_{B_{R}}(x, t) \text { for }(x, t) \in \Omega_{R} \times\left(0, \tau_{2}\right] .
$$

Therefore

$$
\begin{equation*}
u \leq u_{\sigma, \delta}+u_{\mu^{*}}+\epsilon+w_{B_{R}} \quad \text { in } Q_{\infty}^{\Omega_{R}} \tag{6.57}
\end{equation*}
$$

Letting successively $R \rightarrow \infty, \epsilon \rightarrow 0, \delta \rightarrow 0$ and $\sigma \rightarrow 0$ we obtain $u \leq \bar{u}_{\mathcal{S}, 0}+\bar{u}_{\mu^{*}}$, and therefore $u \leq \pi_{+}\left(\bar{u}_{\mathcal{S}, 0}+\bar{u}_{\mu^{*}}\right)=\bar{u}_{\mathcal{S}, \mu}$. Next, we also have

$$
u \geq \tilde{u}_{\mathcal{S}, 0}:=\sup \left\{u_{\infty, a}: a \in \mathcal{S}\right\} \Longrightarrow u \geq \underline{u}_{\mathcal{S}, 0}=\pi_{+}\left(\tilde{u}_{\mathcal{S}, 0}\right)
$$

by (6.25). With the notations of Proposition 6.12 with $F=\mathcal{S}$, for any $R>0, \delta>0$ and $\epsilon>0$, there exists $\tau_{\epsilon}$ such that

$$
u_{\tilde{\mu}_{\delta}}(x, t) \leq u(x, t)+w_{B_{R}}(x, t)+\epsilon \quad \text { in } \Omega_{R} \times\left(0, \tau_{\epsilon}\right]
$$

because the support of $\tilde{\mu}_{\delta}$ is included in $\Omega \backslash \mathcal{S}_{\delta}$. Therefore this last inequality holds in $Q_{\infty}^{\Omega_{R}}$ and consequently

$$
\sup \left\{u_{\tilde{\mu}_{\delta}}, \underline{u}_{\mathcal{S}, 0}\right\} \leq u+\epsilon+w_{B_{R}} \quad \text { in } Q_{\infty}^{\Omega_{R}}
$$

and we can let $R \rightarrow \infty$ and $\epsilon \rightarrow 0$ to obtain $\sup \left\{u_{\tilde{\mu}_{\delta}}, \underline{u}_{\mathcal{S}, 0}\right\} \leq u$ in $Q_{\infty}^{\Omega}$. Letting $\delta \rightarrow 0$ and using Proposition 6.12 we get $\sup \left\{\underline{u}_{\mu^{*}}, \underline{u}_{\mathcal{S}, 0}\right\} \leq \underline{u}_{\mathcal{S}, \mu} \leq u$.
Step 2: Alternative construction. For $0<\delta<\sigma$ and $n \in \mathbb{N}_{*}$ we denote by $\bar{u}_{n, \sigma, \delta, \mu}$ the solution of

$$
\begin{array}{cl}
\partial_{t} u-\Delta u+t^{\alpha} u^{q}=0 & \text { in } Q_{\infty}^{\Omega_{\sigma}} \\
u \geq 0 & \text { in } Q_{\infty}^{\Omega_{\sigma}}  \tag{6.58}\\
u(x, t)=0 & \text { in } \partial_{\ell} Q_{\infty}^{\Omega_{\sigma}} \\
u(., 0)=\mu_{\delta}+n \chi_{\mathcal{S}_{\delta}} & \text { in } \Omega_{\sigma}
\end{array}
$$

We denote here $\Omega_{\sigma}=\left\{x \in \mathbb{R}^{N}: \operatorname{dist}(x, \Omega)<\sigma\right\}, \mathcal{S}_{\delta}=\left\{x \in \mathbb{R}^{N}: \operatorname{dist}(x, \mathcal{S})<\delta\right\}$ and $\mathcal{R}_{\delta}=\Omega \cap \mathcal{S}_{\delta}^{c}$ and $\mu_{\delta}=\chi_{\mathcal{R}_{\delta}} \mu$. The same arguments of monotonicity as in Proposition 6.7 and Proposition 6.11 show that

$$
\begin{equation*}
\lim _{\sigma \rightarrow 0} \lim _{\delta \rightarrow 0} \lim _{n \rightarrow \infty}=\bar{u}_{n, \sigma, \delta, \mu}=\bar{u}_{\mathcal{S}, \mu} \tag{6.59}
\end{equation*}
$$

If $\tau>0$ we denote by $\underline{u}_{\tau, \delta, \mu}$ the solution of

$$
\begin{array}{cl}
\partial_{t} u-\Delta u+t^{\alpha} u^{q}=0 & \text { in } Q_{\infty}^{\Omega} \\
u \geq 0 & \text { in } Q_{\infty}^{\Omega}  \tag{6.60}\\
u(x, t)=0 & \text { in } \partial_{\ell} Q_{\infty}^{\Omega} \\
u(., 0)=\mu_{\delta}+\chi_{\mathcal{s}_{\delta}} \underline{u}_{\mathcal{S}, 0}(., \tau) & \text { in } \Omega,
\end{array}
$$

Using estimate (6.53) and Proposition 6.9 is is easy to prove that $\underline{u}_{\tau, \delta, \mu} \leq u$ for any $u \in$ $\mathcal{U}_{\mathcal{S}, 0}(\Omega)$. Furthermore

$$
\max \left\{u_{\mu_{\delta}}, u_{\chi_{\mathcal{S}_{\delta}} \underline{u}_{\mathcal{S}, 0}(\cdot, \tau)}\right\} \leq \underline{u}_{\tau, \delta, \mu} \leq u_{\mu_{\delta}}+\underline{u}_{\mathcal{S}, 0}(., .+\tau),
$$

since have

$$
u_{\chi_{\mathcal{S}_{\delta}} \underline{u}_{\mathcal{S}, 0}(., \tau)} \leq \underline{u}_{\mathcal{S}, 0}(., \tau) \leq u_{\chi_{\mathcal{S}_{\delta}}} \underline{u}_{\mathcal{S}, 0}(,, \tau)+C\left(\frac{\delta^{2}}{\tau}\right)^{\frac{1+\alpha}{q-1}-\frac{1}{2}} e^{-\frac{\delta^{2}}{4 \tau}}
$$

by (5.5) with $N=1$. Set $c(\delta, \tau)=C\left(\frac{\delta^{2}}{\tau}\right)^{\frac{1+\alpha}{q-1}-\frac{1}{2}} e^{-\frac{\delta^{2}}{4 \tau}}$, then

$$
\max \left\{u_{\mu_{\delta}}, u_{\chi_{\mathcal{S}_{\delta}} \underline{u}_{\mathcal{S}, 0}(., \tau)}\right\} \geq \max \left\{u_{\mu_{\delta}}, \underline{u}_{\mathcal{S}, 0}(., \tau)-c(\delta, \tau)\right\} \geq \max \left\{u_{\mu_{\delta}}, \underline{u}_{\mathcal{S}, 0}(., \tau)\right\}-c(\delta, \tau)
$$

Therefore, if $\underline{u}_{\tau_{n}, \delta, \mu} \rightarrow \underline{u}_{0, \delta, \mu}$ locally uniformly in $Q_{\infty}^{\Omega}$, then $\underline{u}_{0, \delta, \mu}$ is a solution of (1.6) in $Q_{\infty}^{\Omega}$ which satisfies

$$
\begin{equation*}
\max \left\{u_{\mu_{\delta}}, \underline{u}_{\mathcal{S}, 0}\right\} \leq \underline{u}_{\delta, 0, \mu} \leq u_{\mu_{\delta}}+\underline{u}_{\mathcal{S}, 0} \tag{6.61}
\end{equation*}
$$

and is smaller than any $u \in \mathcal{U}_{\mathcal{S}, 0}(\Omega)$. There exists $\delta_{n} \rightarrow 0$ such that $\underline{u}_{0, \delta_{n}, \mu} \rightarrow \underline{u}_{0,0, \mu}$. Then

$$
\begin{equation*}
\max \left\{\underline{u}_{\mu^{*}}, \underline{u}_{\mathcal{S}, 0}\right\} \leq \underline{u}_{0,0, \mu} \leq \underline{u}_{\mu^{*}}+\underline{u}_{\mathcal{S}, 0}, \tag{6.62}
\end{equation*}
$$

and $\underline{u}_{0,0, \mu}$ is an element of $\mathcal{U}_{\mathcal{S}, 0}(\Omega)$ smaller than any $u \in \mathcal{U}_{\mathcal{S}, 0}(\Omega)$. Thus $\underline{u}_{0,0, \mu}=\underline{u}_{\mathcal{S}, \mu}$ and

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} \lim _{\tau \rightarrow 0} \underline{u}_{\tau, \delta, \mu}=\underline{u}_{\mathcal{S}, \mu} . \tag{6.63}
\end{equation*}
$$

Step 3: Proof of (6.54). We assume $\inf \left\{\left|z-z^{\prime}\right|: z \in \mathcal{S}, z^{\prime} \in \Omega^{c}\right\}=\delta_{0}>0$, so that we can take $\sigma=0$ in the construction of $\bar{u}_{\mathcal{S}, \mu}$. Put $\tau=\left(\frac{c_{\alpha}}{n}\right)^{\frac{q-1}{1+\alpha}}$ and

$$
\begin{equation*}
\bar{Z}_{n, \delta, \mu}=\bar{u}_{n, \delta, \mu}-\bar{u}_{n, \delta, 0}, \underline{Z}_{\tau, \delta, \mu}=\underline{u}_{\tau, \delta, \mu}-\underline{u}_{\tau, \delta, 0} \text { and } W_{n, \delta, \mu}=\bar{Z}_{n, \delta, \mu}-\underline{Z}_{\tau, \delta, \mu} . \tag{6.64}
\end{equation*}
$$

Then $w=W_{n, \delta, \mu}$ satisfies

$$
\begin{equation*}
\partial_{t} w-\Delta w+t^{\alpha}\left(\bar{u}_{n, \delta, \mu}^{q}-\bar{u}_{n, \delta, 0}^{q}-\underline{u}_{\tau, \delta, \mu}^{q}+\underline{u}_{\tau, \delta, 0}^{q}\right) \tag{6.65}
\end{equation*}
$$

in $Q_{\infty}^{\Omega}$ and we can write

$$
\begin{align*}
\bar{u}_{n, \delta, \mu}^{q}-\bar{u}_{n, \delta, 0}^{q}-\underline{u}_{\tau, \delta, \mu}^{q}+\underline{u}_{\tau, \delta, 0}^{q} & =\left(\bar{u}_{n, \delta, \mu}^{q}-\underline{u}_{\tau, \delta, \mu}^{q}\right)-\left(\bar{u}_{n, \delta, 0}^{q}-\underline{u}_{\tau, \delta, 0}^{q}\right)  \tag{6.66}\\
& =d_{\mu}\left(\bar{u}_{n, \delta, \mu}-\underline{u}_{\tau, \delta, \mu}\right)-d_{0}\left(\bar{u}_{n, \delta, 0}-\underline{u}_{\tau, \delta, 0}\right)
\end{align*}
$$

where

$$
d_{\mu}(x, t)= \begin{cases}\frac{\bar{u}_{n, \delta, \mu}^{q}-\underline{u}_{\tau, \delta, \mu}^{q}}{\bar{u}_{n, \delta, \mu}-\underline{u}_{\tau, \delta, \mu}} & \text { if } \bar{u}_{n, \delta, \mu} \neq \underline{u}_{\tau, \delta, \mu}  \tag{6.67}\\ 0 & \text { if } \bar{u}_{n, \delta, \mu}=\underline{u}_{\tau, \delta, \mu}\end{cases}
$$

and $d_{0}$ is defined accordingly. Since

$$
\bar{u}_{n, \delta, \mu} \geq \max \left\{\underline{u}_{\tau, \delta, \mu}, \bar{u}_{n, \delta, 0}\right\} \text { and } \underline{u}_{\tau, \delta, 0} \leq \min \left\{\underline{u}_{\tau, \delta, \mu}, \bar{u}_{n, \delta, 0}\right\},
$$

there holds $d_{\mu} \geq d_{0} \geq$ by convexity. Using the fact that $\bar{u}_{n, \delta, 0}-\underline{u}_{\tau, \delta, 0} \geq 0$ is infers that

$$
d_{\mu}\left(\bar{u}_{n, \delta, \mu}-\underline{u}_{\tau, \delta, \mu}\right)-d_{0}\left(\bar{u}_{n, \delta, 0}-\underline{u}_{\tau, \delta, 0}\right) \geq d_{\mu}\left(\bar{u}_{n, \delta, \mu}-\underline{u}_{\tau, \delta, \mu}-\bar{u}_{n, \delta, 0}+\underline{u}_{\tau, \delta, 0}\right)
$$

Finally (6.65) becomes

$$
\begin{equation*}
\partial_{t} w-\Delta w+t^{\alpha} d_{\mu} w \leq 0 \quad \text { in } Q_{\infty}^{\Omega} \tag{6.68}
\end{equation*}
$$

Furthermore, in the sense of measures,

$$
w(., 0)=\mu_{\delta}+n \chi_{\mathcal{S}_{\delta}}-n \chi_{\mathcal{S}_{\delta}}-\left(\mu_{\delta}+\chi_{\mathcal{S}_{\delta}} \underline{u}_{\mathcal{S}, 0}(., \tau)\right)+\chi_{\mathcal{S}_{\delta}} \underline{u}_{\mathcal{S}, 0}(., \tau)=0 .
$$

Because $w=0$ in $\partial_{\ell} Q_{\infty}^{\Omega}$ it follows $w \leq 0$ by the maximum principe. Therefore

$$
\begin{equation*}
\bar{Z}_{n, \delta, \mu} \leq \underline{Z}_{\tau, \delta, \mu} \Longrightarrow \bar{u}_{n, \delta, \mu}-\underline{u}_{\tau, \delta, \mu} \leq \bar{u}_{n, \delta, 0}-\underline{u}_{\tau, \delta, 0} . \tag{6.69}
\end{equation*}
$$

If we let successively $n \rightarrow \infty$ (and therefore $\tau \rightarrow 0$ ) and $\delta \rightarrow 0$, we obtain (6.54).

Remark. We do not know if (6.54 holds if we do not assume $\inf \left\{\left|z-z^{\prime}\right|: z \in \mathcal{S}, z^{\prime} \in \Omega^{c}\right\}=$ $\delta_{0}>0$. However, if for $\theta>0$ we set $\mathcal{S}_{\theta}=\mathcal{S} \cap\left\{x \in \Omega\right.$ : dist $\left.\left(x, \Omega^{c}\right) \geq \theta\right\}$, then we have

$$
\begin{equation*}
\bar{u}_{\mathcal{S}_{\theta}, \mu}-\underline{u}_{\mathcal{S}_{\theta}, \mu} \leq \bar{u}_{\mathcal{S}_{\theta}, 0}-\underline{u}_{\mathcal{S}_{\theta}, 0} . \tag{6.70}
\end{equation*}
$$

Furthermore all the four above functions increases when $\theta$ decreases to 0 . If we set

then we infer that

$$
\begin{equation*}
\bar{u}_{\underline{\mathcal{S}}, \mu}-\underline{u}_{\mathcal{S}, \mu} \leq \bar{u}_{\underline{\mathcal{S}}, 0}-\underline{u}_{\underline{\mathcal{S}}, 0} . \tag{6.72}
\end{equation*}
$$

Our final result is the following existence and uniqueness theorem the proof is similar to the one of [7, Th 3.5].

Theorem 6.15 Assume $\Omega \subseteq \mathbb{R}^{N}$ is either $\mathbb{R}^{N}$ or an open domain with a $C^{2}$ compact boundary, $\alpha>-1$ and $1<q<q_{c, \alpha}$. Then for any $\nu \approx(\mathcal{S}, \mu) \in \mathfrak{B}^{\text {reg }}(\Omega)$ such that $\rho \mu$ is bounded in any neighborhood of $\partial \Omega$ and $\inf \left\{\left|z-z^{\prime}\right|: z \in \mathcal{S}, z^{\prime} \in \Omega^{c}\right\}>0$, the set $\mathcal{U}_{\mathcal{S}, \mu}(\Omega)$ contains one and only one element.

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