

On the invariant measures for the Ostrovsky equation Mohamad Darwich

▶ To cite this version:

Mohamad Darwich. On the invariant measures for the Ostrovsky equation. 2013. hal-00798781v2

HAL Id: hal-00798781 https://hal.science/hal-00798781v2

Preprint submitted on 11 Mar 2013

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

ON THE INVARIANT MEASURES FOR THE OSTROVSKY EQUATION.

DARWICH MOHAMAD.

ABSTRACT. In this paper, we construct invariant measures for the Ostrovsky equation associated with conservation laws. On the other hand, we prove the local well- posedness of the initial value problem for the periodic Ostrovsky equation with initial data in $H^s(\mathbb{T})$ for $s > -\frac{1}{2}$.

1. Introduction

In this paper, we construct an invariant measure for a dynamical system defined by the Ostrovsky equation (Ost)

$$\begin{cases} \partial_t u - u_{xxx} + \partial_x^{-1} u + u u_x = 0, \\ u(0, x) = u_0(x). \end{cases}$$
 (1.1)

associated to the conservation of the energie. The operator ∂_x^{-1} in the equation denotes a certain antiderivative with respect to the variable x defined for 0-mean value periodic function the Fourier transform by $\widehat{(\partial_x^{-1} f)} = \frac{\widehat{f}(\xi)}{i\xi}$.

Invariant measure play an important role in the theory of dynamical systems (DS). It is well known that the whole ergodic theory is based on this concept. On the other hand, they are necessary in various physical considerations.

Note that, one the well-known applications of invariant measures in the theory of dynamical is the Poincaré recurrence theorem: every flow which preserves a finite measure has the returning property modulo a set of measure zero.

Recently several papers [1], [10], [11]) have been published on invariant measures for dynamical system generated by nonlinear partial differential equations.

In [12] an infinite series of invariant measure associated with a higher conservation laws are constructed for the one-dimensional Korteweg de Vries (KdV) equation:

$$u_t + uu_x + u_{xxx} = 0,$$

by Zhidkov. In particular, invariant measure associated to the conservation of the energie are constructed for this equation.

Equation 1.1 is a perturbation of the Korteweg de Vries (KdV) equation with a nonlocal term and was deducted by Ostrovskii [9] as a model for weakly nonlinear long waves, in a rotating frame of reference, to describe the propagation of surface waves in the ocean.

We will construct invariant measures associated to the conservation of the Hamiltonian:

$$H(u(t)) = \frac{1}{2} \int (u_x)^2 + \frac{1}{2} \int (\partial_x^{-1} u)^2 - \frac{1}{6} \int u^3.$$

The paper is organized as follows. In Section 2 the basic notation is introduced and the basic results are formulated. In Section 3 the invariant measure which corresponds to the conservation of the Hamiltonian is constructed.

In Section 4 we will prove the local well-posedness for our equation in H^s , $s > -\frac{1}{2}$.

2. Notations and main results

We will use C to denote various time independent constants, usually depending only upon s. In case a constant depends upon other quantities, we will try to make it explicit. We use $A \lesssim B$ to denote an estimate of the form $A \leq CB$. similarly, we will write $A \sim B$ to mean $A \lesssim B$ and $B \lesssim A$. We writre $\langle \cdot \rangle := (1+|\cdot|^2)^{1/2} \sim 1+|\cdot|$. The notation a^+ denotes $a+\epsilon$ for an arbitrarily small ϵ . Similarly a- denotes $a-\epsilon$. Let

$$L_0^2 = \{ u \in L^2; \int_{\mathbb{T}} u dx = 0 \}.$$

On the circle, the Fourier transform is defined as

$$\hat{f}(n) = \frac{1}{2\pi} \int_{\mathbb{T}} f(x) \exp(-inx) dx.$$

We introduce the zero mean-value Sobolev spaces H^s defined by :

$$H_0^s =: \{ u \in \mathcal{S}'(\mathbb{T}); ||u||_{H^s} < +\infty \text{ and } \int_{\mathbb{T}} u dx = 0 \},$$
 (2.1)

where,

$$||u||_{H_0^s} = (2\pi)^{\frac{1}{2}} ||\langle . \rangle^s \hat{u}||_{l_x^2},$$
 (2.2)

and $X^{s,\frac{1}{2}}$ by

$$\{u \in \mathcal{S}'(\mathbb{T}); ||u||_{X^{s,\frac{1}{2}}} := ||\langle n \rangle^s \langle \tau + n^3 - \frac{1}{n} \rangle \hat{u}||_{l_n^2 L_\tau^2} < \infty \}.$$

Let

$$Y^{s} =: \{ u \in \mathcal{S}'(\mathbb{T}); ||u||_{Y^{s}} < +\infty \},$$

where

$$||u||_{Y^s} = ||u||_{X^{s,\frac{1}{2}}} + ||\langle n \rangle^s \hat{u}(n,\tau)||_{l_n^2 L_\tau^1}.$$

We will briefly remind the general construction of a Gaussian measure on a Hilbert space. Let X be a Hilbert space, and $\{e_k\}$ be the orthonormal basis in X which consists of eigenvectors of some operator $S = S^* > 0$ with corresponding eigenvalues $0 < \lambda_1 \le \lambda_2 \le \lambda_3 \le \lambda_k \le ...$ We call a set $M \subset X$ a cylindrical set iff:

$$M = \{x \in X; [(x, e_1), (x, e_2), ...(x, e_r)] \in F\}$$

for some Borel $F \subset \mathbb{R}^r$, and some integer r. We define the measure w as follows:

$$w(M) = (2\pi)^{-\frac{r}{2}} \prod_{j=1}^{r} \lambda_j^{\frac{1}{2}} \int_F e^{-\frac{1}{2} \sum_{j=1}^{r} \lambda_j y_j^2} dy.$$
 (2.3)

One can easily verify that the class \mathbb{A} of all cylindrical sets is an algebra on which the function w is additive. The function w is called the centered Gaussian measure on X with the correlation operator S^{-1} .

Definition 2.1. The measure w is called a countably additive measure on an algebra \mathbb{A} if $\lim_{n\to+\infty}(A_n)=0$ for any $A_n\in\mathbb{A}(n=1,2,3...)$ for which $A_1\supset A_2\supset A_3\supset\ldots \supset A_n\supset\ldots$ and $\bigcap_{n=1}^{\infty}A_n=\phi$

Now we give the following Lemma:

Lemma 2.1. The measure w is countably additive on the algebra \mathbb{A} iff S^{-1} is an operator of trace class, i.e iff $\sum_{k=1}^{+\infty} \lambda_k^{-1} < +\infty$.

Now we present some definitions related to invariant measure :

Definition 2.2. Let M be a complete separable metric space and let a function $h : \mathbb{R} \times M \longmapsto M$ for any fixed t be a homeomorphism of the space M into itself satisfying the properties:

- (1) $h(0,x) = x \text{ for any } x \in M$.
- (2) $h(t, h(\tau, x)) = h(t + \tau, x)$ for any $t, \tau \in \mathbb{R}$ and $x \in M$.

Then, we call the function h a dynamical system with the space M. If μ is a Borel measure defined on the phase space M and $\mu(\Omega) = \mu(h(\Omega, t))$ for an arbitrary Borel set $\Omega \subset M$ and for all $t \in \mathbb{R}$, then it is called an invariant measure for the dynamical system h.

Let us now state our results:

Theorem 2.1. Let s > -1/2, and $\phi \in H_0^s$. Then there exists a time $T = T(||\phi||_{H_0^s}) > 0$ and a unique solution u of (1.1) in $C([0,T],H_0^s) \cap Y^s$ and the map $\phi \longmapsto u$ is C^{∞} from H_0^s to $C([0,T],H_0^s)$.

Theorem 2.2. Let $\phi \in L_0^2$, then the Problem 1.1 is global well-posedness in L^2 and the Borel measure μ on L^2 defined for any Borel set $\Omega \subset L^2$ by the rule

$$\mu(\Omega) = \int_{\Omega} e^{-g(u)} dw(u)$$

where w is the centered Gaussian measure corresponding to the correlation operator $S^{-1} = (-\Delta + \Delta^{-1})^{-1}$, and $g(u) = \frac{1}{3} \int u^3 dx$ the nonlinear term of the Hamiltonian is an invariant measure for (1.1).

3. Invariance of Gibbs Measure

In this section, we construct an invariant measure to Equation 1.1 with respect to the conservation of the Hamiltonian. Let us first present result on invariant measures for systems of autonomous ordinary differential equations. Consider the following system of ordinary differential equations:

$$\dot{x} = b(x),\tag{3.1}$$

where $x(t): \mathbb{R} \longrightarrow \mathbb{R}^n$ is an unknown vector-function and $b(x): \mathbb{R}^n \longrightarrow \mathbb{R}^n$ is a continuously differentiable map. Let h(t,x) be the corresponding function ("dynamical system") from $\mathbb{R} \times \mathbb{R}^n$ into \mathbb{R}^n transforming any $t \in \mathbb{R}$ and $x_0 \in \mathbb{R}^n$ into the solution x(t), taken at the moment of time t, of the above system supplied with the initial data $x(0) = x_0$.

Theorem 3.1. Let P(x) be a continuously differentiable function from \mathbb{R}^n into \mathbb{R} . For the Borel measure

$$\nu(\Omega) = \int_{\Omega} P(x) dx$$

to be invariant for the function h(t,x) in the sense that $\nu(h(t,\Omega)) = \nu(\Omega)$ for any bounded domain Ω and for any t, it is sufficient and necessary that

$$\sum_{i=1}^{n} \frac{\partial}{\partial x_i} (P(x)b_i(x)) = 0,$$

for all $x \in \mathbb{R}^n$.

We shall construct an invariant measure for (1.1). Let A > 0, the space $L^2(0, A)$ be real equipped with the scalar product:

$$(u,v)_{L^2(0,A)} = \int_0^A u\overline{v}dx.$$

and $J = \frac{\partial}{\partial x}Q$ where the operator Q maps $v^* \in L^2$ into $v \in L^2$ such that $v^*(g) = (v,g)_{L^2(0,A)}$. Finally, let $S = -\Delta + \Delta^{-1}$. We set $H(u) = \frac{1}{2}(\int (u_x)^2 - \int (\partial_x^{-1}u)^2) + \frac{1}{3}\int u^3 = \frac{1}{2}(Su,u) + g(u)$. Note that System 1.1 takes the form:

$$\begin{cases} \frac{\partial u}{\partial t}(t) = J \frac{\delta}{\delta u} H(u(t)), t \in \mathbb{R} \\ u(t_0) = \phi \in H^s, \end{cases}$$
 (3.2)

Let $e_{2k-1}(x) = \frac{\sqrt{2}}{\sqrt{A}}\sin(\frac{2\pi nx}{A})$, $e_{2k} = \frac{\sqrt{2}}{\sqrt{A}}\cos(\frac{2\pi nx}{A})$ where k = 1, 2, 3... Then $(e_k)_{k=1,2,...}$ is an orthonormal basis of the space $L_0^2(0,A)$ consisting of eigenfunctions of the operator Δ with the corresponding eigenvalues $0 < \lambda_1 = \lambda_2 < ... < \lambda_{2k-1} = \lambda_{2k} < ...$ Let P_m be the orthogonal projector in L_0^2 onto the subspace $L_m = \text{span}\{e_1,...,e_{2m}\}$ and P_m^{\perp} be the orthogonal projector in $L_0^2(0,A)$ onto the orthogonal complement L_m^{\perp} to the subspace L_m . Let also $v_i = -\lambda_i + \lambda_i^{-1}$, then v_i are eigenvalues of S.

Consider the following problem:

$$\begin{cases} \partial_t u^m - u_{xxx}^m + \partial_x^{-1} u^m + P_m(u^m u_x^m) = 0, \\ u^m(0, x) = P_m u_0(x). \end{cases}$$
(3.3)

The existence of u is global in L^2 in time (see later) and the solution of (3.3) converges to u in $C([0,T],L^2)$ for any fixed T, more precisely we have the following lemma:

Lemma 3.1. (1) The solution u_m of (3.3) converges in $C([0,T], L^2)$ to the solution u of (1.1).

(2) For any $\epsilon > 0$, and T > 0 there exists $\delta > 0$ such that

$$Max_{t \in [t_0 - T, t_0 + T]} \|u_m(., t) - v_m(., t)\|_{L^2} < \epsilon,$$

for any two solutions u_m and v_m of the problem (3.3), satisfying the condition

$$||u_m(.,t_0)-v_m(.,t_0)||_{L^2}<\delta.$$

Proof.: By the Duhamel formula, $u - u^m$ satisfies

$$u(t) - u^{m}(t) = e^{-itS}(u_0 - P_m u_0) - \frac{1}{2} \int_0^t e^{-i(t-t')S}(\partial_x(u^2(t')) - P_m(\partial_x((u^m)^2(t'))))dt'.$$

We can white that $R(t) := \partial_x(u^2(t')) - P_m(\partial_x((u^m)^2(t'))) = \partial_x(u^2 - u^2)$ $(P_{\frac{m}{2}}u)^2) + P_m\partial_x\left((P_{\frac{m}{2}}u)^2 - u^2\right) + P_m\partial_x(u^2 - (u_m)^2)$. Now, using the linear and bilinear estimates proved in section 4, we obtain that

$$||u - u_m||_{Y^s} \lesssim ||u_0 - P_m u_0||_{H^s} + T^{\gamma} ||u - u_m||_{Y^s} ||u + u_m||_{Y^s} + ||u - P_{\frac{m}{2}} u||_{Y^s},$$
(3.4)

then $u_m \longrightarrow u$ in Y^s , but $Y^s \hookrightarrow L_t^{\infty} L_x^2$, this gives the uniform convergence in L^2 .

The proof of part (2) is similar to part (1).

By $h_m(u_0,t)$ we denote the function mapping any $u_0 \in L^2$ and $t \in$ \mathbb{R} into $u_m(.,t+t_0)$ where $u_m(.,t)$ is the solution of the problem (3.3). It is clear that the function h_m is a dynamical system with the phase space $X^m = \text{span}\{e_1, ... e_m\}$. In addition, the direct verification shows that $\frac{d}{dt}||u_m(.,t)||^2_{L^2}=0$ and $\int u_m dx=0.$. For each m=1,2... let us consider in the space X^m the centered Gaussian measure w_m with the correlation operator S^{-1} . Since $S = S^*$ in X^m , the measure w_m is well-defined in X^m . Also, since $g(u) = \frac{1}{3} \int u^3$ is a continuous functional in X^m , the following Borel measures

$$\mu_m(\Omega) = \int_{\Omega} e^{-g(u)} dw_m(u).$$

(where Ω is an arbitrary Borel set in L^2) are well defined.

Definition 3.1. A set Π of measures defined on the Borel sets of a topologogical space is called tight if, for each $\epsilon > 0$, there exist a compact set K such that

$$\mu(K) > 1 - \epsilon$$

For all $\mu \in \Pi$.

We will use the following theorem:

Theorem 3.2. (Prokhorov) A tight set, Π , of measures on the Borel sets of a metric topological space, X, is relatively compact in the sense that for each sequence, $P_1, P_2, ...$ in Π there exists a subsequence that converges to a probability measure P, not necessarily in Π , in the sense that

$$\int gdP_{n_j} \longrightarrow \int gdP$$

for all bounded continuous integrands. Conversely, if the metric space is separable and complete, then each relatively compact set is tight.

To prove Theorem 2.2, we will prove the following Lemma:

Lemma 3.2. μ_m is an invariant measure for the dynamical system h_m with the phase space X^m .

Proof: Let us rewrite the system (3.3) for the coefficients a_k , where $u^m(t) = \sum_{k=1}^{k=2m} a_k(t)e_k$. Let $h(a) = H(\sum_{k=1}^{k=2m} a_ke_k)$ and J is a skew-symmetric matrix, $(J_m)_{2k-1,2k} = -\frac{2\pi k}{A} = -(J_m)_{2k,2k-1} (k=1,2,... m)$ then the problem take the form

 $\begin{cases} a'(t) = J_m \nabla_a h(a(t)), \\ a_k(t_0) = (u_0, e_k), k = 1, 2, \dots 2m \end{cases}$ (3.5)

Using Theorem 3.1, we can easily verify that the Borel measure:

$$\mu'_m(A) = (2\pi)^{-\frac{2m+1}{2}} \prod_{i=1}^{2m} v_i^{\frac{1}{2}} \int_A e^{-\frac{1}{2} \sum_{j=1}^{2m} v_j a_j^2 - g(\sum_{j=1}^{2m} a_j e_j(x))} da,$$

(with $v_j = -\lambda_j + \lambda_j^{-1}$ the eigenvalues of S) is invariant for the problem (3.5). Also, we introduce the measures

$$w_m(A) = (2\pi)^{-\frac{2m+1}{2}} \prod_{j=1}^{2m} v_j^{\frac{1}{2}} \int_A e^{-\frac{1}{2} \sum_{j=1}^{2m} v_j a_j^2} da.$$

Let $\Omega_m \subset X^m$ and $\Omega_m = \{u \in L^2, u = \sum_{j=1}^{2m} a_j e_j, a \in A\}$ where $A \subset \mathbb{R}^{2m}$ is a

Borel set. We set $\mu_m(\Omega_m) = \mu'_m(A)$. Since the measure μ'_m is invariant for (3.5), the measure μ_m is invariant for the problem (3.3).

Although the measure is defined on X^m , we can define it on the Borel sigmaalgebra of L^2 by the rule: $\mu_m(\Omega) = \mu_m(\Omega \cap X^m)$. Since the set $\Omega \cap X^m$ is open as a set in X^m for any open set $\Omega \subset L^2$, this procedure is correct.

Lemma 3.3. $(w_m)_m$ weakly converges to w in L^2 .

Proof: S^{-1} is an operator of trace since the trace $Tr(S^{-1}) = \sum_{k} v_k^{-1} = \sum_{k} v_k^{-1} = \sum_{k} v_k^{-1}$

$$\sum_{k} \frac{1}{\frac{1}{\frac{4\pi^{2}k^{2}}{A^{2}}} + \frac{4\pi^{2}k^{2}}{A^{2}}} < +\infty.$$
 Thus we can find a continuous positive func-

tion d(x) defined on $(0,\infty)$ with the property $\lim_{x\to+\infty}d(x)=+\infty$ such that

$$\sum_{k} v_k^{-1} d(\lambda_k) < +\infty$$
. We define the operator $T = d(S)$, the operator defined

by $T(e_k) = d(v_k)e_k$ and let $B = S^{-1}T$. According to the definition of d(x), $Tr(B) < +\infty$. Let R > 0 and $B_R = \{u \in L^2, T^{\frac{1}{2}}u \in L^2 \text{and} ||T^{\frac{1}{2}}u|| \leq R\}$, it is clear that the closure of B_R is compact for any R > 0. Combined the following inequality (see [4] for the proof)

$$w_n(\overline{B_R}^C) = w_n(\{u; (Tu, u)_{L^2} > R\}) \le \frac{Tr(B)}{R^2}.$$

with the Prokhorov theorem, this ensure that (w_n) is weakly compact on L^2 .

In view of the definition $w_n(M) \to w(M)$ for any cylindrical set $M \subset L^2$. (because $w_n(M) = w(M)$ for all sufficiently large n). Hence, since the

extension of a measure from an algebra to a minimal sigma-algebra is unique, we have proved that the sequence w_n converges to w weakly in L^2 and Lemma 3.3 is proved.

Lemma 3.4. $\liminf_m \mu_m(\Omega) \ge \mu(\Omega)$ for any open set $\Omega \subset L^2$. $\limsup_m \mu_m(K) \le \mu(K)$ for any closed bounded set $K \subset L^2$.

Proof: Let $\Omega \subset L^2$ be open and let $B_R = \{u \in L^2, ||u||_L^2 < R\}$ for some R > 0.

Consider $\phi(u): 0 < \phi(u) < 1$ with the support belonging to $\Omega_R = \Omega \cap B_R$ such that

$$\int_X \phi(u)e^{-g(u)}dw(u) \geqslant \mu(\Omega_R) - \epsilon.$$

Then,

$$\lim_{m} \inf \mu_{m}(\Omega_{R}) = \lim_{m} \inf \int_{\Omega_{R}} e^{-g(u)} dw_{m}(u) \ge \lim_{m} \inf \int \phi(u) e^{-g(u)} dw_{m}(u)$$

$$= \int \phi(u) e^{-g(u)} dw(u) \ge \mu(\Omega_{R}) - \epsilon.$$

Therefore, due to the arbitrariness of $\epsilon > 0$ one has:

$$\liminf_{m} \mu_m(\Omega) \ge \limsup_{m} \mu_m(\Omega_R) \ge \mu(\Omega_R).$$

Taking $R \longrightarrow +\infty$ in this inequality, we obtain the first statement the lemma.

Let K be a closed bounded set. Fix $\epsilon > 0$. We take a continuous function $\phi \in [0,1]$ such that $\phi(u) = 1$ for any $u \in K$, $\phi(u) = 0$ if $\operatorname{dist}(u,K) > \epsilon$ and $\int \phi(u)e^{-g(u)}w(du) < \mu(K) + \epsilon$. Then

$$\limsup_{m} \mu_{m}(K) \leq \limsup_{m} \int \phi(u)e^{-g(u)}dw_{m}(u)$$
$$= \int \phi(u)e^{-g(u)}dw(u) \leq \mu(K) + \epsilon,$$

and due to the arbitrariness of $\epsilon > 0$, Lemma 3.4 is proved.

Lemma 3.5. Let $\Omega \subset L^2$ an open set and $t \in \mathbb{R}$. Then $\mu(\Omega) = \mu(h(\Omega, t))$.

Proof: Let $\Omega_1 = h(\Omega, t)$. Fix an arbitrary $t \in \mathbb{R}$, then Ω_1 is open too. First, let us suppose that $\mu(\Omega) < \infty$.

Fix an arbitrary $\epsilon > 0$, by Prokhorov Theorem there exists a compact set $K \subset \Omega$ such that $\mu(\Omega \setminus K) < \epsilon$, note that $K_1 = h(K, t)$ is a compact set, too, and $K_1 \subset \Omega_1$.

For any $A \subset L^2$, let ∂A be the boundary of the set A and let

$$\beta = \min\{dist(K, \partial\Omega); dist(K_1, \partial\Omega_1)\}\$$

(where $\operatorname{dist}(A, B) = \inf_{x \in A, y \in B} \|x - y\|_{L^2}$). Then, $\beta > 0$. According to Lemma 3.1, for any $z \in K$, there exists $\delta > 0$ such that for any $x, y \in B_{\delta}(z)$ one has $\|h_n(x,t) - h_n(y,t)\|_{L^2} < \frac{\beta}{3}$. Lets $\Omega^{\alpha} = \{q \in \Omega_1; \operatorname{dist}(q, \partial \Omega_1) \geq \alpha\}$ and $B_{\delta_1}(z_1), ... B_{\delta_l}(z_l)$ be a finite covering of the compact set K by these balls and let $B = \bigcup_{i=1}^{l} B_{\delta_i}(z_i)$.

Since $h_n(z_i,t) \longrightarrow h(z_i,t)(n \longrightarrow +\infty)$ for any i we obtain that $\operatorname{dist}(h_n(z,t),K_1) < \infty$ $\frac{\beta}{3}$, $\forall z \in B$ and large n. Thus, $h_n(B,t)$ belongs to a closed bounded subset of $\Omega^{\frac{p}{2}}$ for all sufficiently large n.

Further, we get by the invariance of μ_n and Lemma 3.4

$$\mu(\Omega) \le \mu(B) + \epsilon \le \liminf \mu_n(B) + \epsilon \le \liminf \mu_n(h_n(B,t)) + \epsilon \le \mu(\Omega_1) + \epsilon$$

because
$$\mu_n(B) = \mu_n(B \cap X_n) = \mu_n(h_n(B \cap X_n, t))$$
, and $h_n(B \cap X_n, t) \subset$

 $h_n(B,t)$). Hence, due to the arbitariness of $\epsilon > 0$, we have $\mu(\Omega) \leq \mu(\Omega_1)$.

By analogy $\mu(\Omega) \geqslant \mu(\Omega_1)$. Thus $\mu(\Omega) = \mu(\Omega_1)$.

Now if Ω is open and $\mu(\Omega) = +\infty$, then we take the sequence

$$\Omega^k = \Omega \cap \{ u \in L^2; ||u||_{L^2} + ||h(u, t)|| < k \}$$

and set $\Omega_1^k = h(\Omega^k, t)$. Then $\Omega = \cup \Omega^k$ and $\mu(\Omega^k) = \mu(\Omega_1^k) < \infty$. Taking $k \longrightarrow +\infty$, we obtain the statement of the lemma.

4. Well-posedness in $X^{s,\frac{1}{2}}$

In this section, we prove a global wellposedness result for the Ostrovsky equation by following the idea of Kenig, Ponce, and Vega in [8].

Our work space is Y^s , the completion of functions that are Schwarz in time and C^{∞} in space with norm:

$$||u||_{Y^s} = ||u||_{X^{s,\frac{1}{2}}} + ||\langle n \rangle^s \hat{u}(n,\tau)||_{l_n^2 L_{\tau}^1}$$

 Y^s is a slight modification of $X^{s,\frac{1}{2}}$ such that $||u||_{L^{\infty}_t H^s_x} \lesssim ||u||_{Y^s}$. We see that the nonlinear part of the Ostrovsky equation is $u\partial_x u$, and by Fourier transform we write it in frequency as

$$n\sum_{n_1\in\dot{\mathbb{Z}}}\int_{\tau_1\in\mathbb{R}}\hat{u}(n_1,\tau_1)\hat{u}(n-n_1,\tau-\tau_1)d\tau_1.$$

The resonance function is given by:

$$R(n, n_1) = \tau + m(n) - (\tau_1 + m(n_1) - (\tau - \tau_1 + m(n - n_1)) = 3nn_1(n - n_1) - \frac{1}{n} \left(1 - \frac{n^3}{nn_1(n - n_1)}\right)$$

where $m(n) = n^3 - \frac{1}{n}$. Now we have the following lower bound on the resonance function:

Lemma 4.1. If $|n||n_1||n-n_1| \neq 0$, and $\frac{1}{|n|} < 1$, then:

$$|R(n, n_1)| \gtrsim |n||n_1||n - n_1|,$$
 (4.1)

and

$$|n|^2 \le 2|nn_1(n-n_1)|. \tag{4.2}$$

Proof: (4.2) is obvious.

Now

$$R^{2}(n, n_{1}) = 9n^{2}n_{1}^{2}(n - n_{1})^{2} - 6n_{1}(n - n_{1}) + 6n^{2} + \frac{1}{n^{2}}\left(1 - \frac{n^{3}}{n(n_{1}(n - n_{1}))}\right)^{2}$$

$$= n^{2}n_{1}^{2}(n - n_{1})^{2} + 8n^{2}n_{1}^{2}(n - n_{1})^{2} - 6n_{1}(n - n_{1}) + 6n^{2} + \frac{1}{n^{2}}\left(1 - \frac{n^{3}}{n(n_{1}(n - n_{1}))}\right)^{2}$$

$$\geq n^{2}n_{1}^{2}(n - n_{1})^{2} + 8n^{2}n_{1}^{2}(n - n_{1})^{2} - 6n_{1}(n - n_{1})$$

$$= n^{2}n_{1}^{2}(n - n_{1})^{2} + |n_{1}(n - n_{1})| (8n^{2} |n_{1}(n - n_{1})| - 6)$$

Using (4.2) we obtain that:

$$R^2(n, n_1) \gtrsim n^2 n_1^2 (n - n_1)^2$$

By the same argument employed in [8], we state the following elemental estimates without proof.

Lemma 4.2. For any $\epsilon > 0$, $\alpha \in \mathbb{R}$ and $0 < \rho < 1$, we have:

$$\int_{\mathbb{R}} \frac{d\beta}{(1+|\beta|)(1+|\alpha-\beta|)} \lesssim \frac{\log(2+|\alpha|)}{(1+|\alpha|)}.$$

$$\int_{\mathbb{R}} \frac{d\beta}{(1+|\beta|)^{\rho}(1+|\alpha-\beta|)} \lesssim \frac{1+\log(1+|\alpha|)}{(1+|\alpha|)^{\rho}}.$$

$$\int_{\mathbb{R}} \frac{d\beta}{(1+|\beta|)^{1+\epsilon}(1+|\alpha-\beta|)^{1+\epsilon}} \lesssim \frac{1}{(1+|\alpha|)^{1+\epsilon}}.$$

Lemma 4.3. There exists c > 0 such that for any $\rho > \frac{2}{3}$ and any τ , $\tau_1 \in \mathbb{R}$, the following is true:

$$\sum_{n_1 \neq 0} \frac{\log(2 + |\tau + m(n_1) + m(n - n_1)|)}{(1 + |\tau + m(n_1) + m(n - n_1)|)} \leq C.$$

$$\sum_{n \neq 0} \frac{\log(2 + |\tau_1 + m(n_1) - m(n - n_1)|)}{(1 + |\tau_1 + m(n_1) - m(n - n_1)|)} \leq C.$$

$$\sum_{n \neq 0} \frac{\log(1 + |\tau_1 + m(n_1) - m(n - n_1)|)}{(1 + |\tau_1 + m(n_1) - m(n - n_1)|)^{\rho}} \leq C.$$

Proposition 4.1. Let $s \ge -\frac{1}{2}$, then for all f, g with compact support in time included in the subset $\{(t,x),t\in[-T,T]\}$, there exists $\theta > 0$ such that:

$$\|\partial_x(fg)\|_{X^{s,-\frac{1}{2}}} \lesssim T^{\theta} \|f\|_{X^{s,\frac{1}{2}}} \|g\|_{X^{s,\frac{1}{2}}}.$$

Remark 4.1. This proposition is false for $s < -\frac{1}{2}$. We can exhibit a counterexample to the bilinear estimate in the Prop (4.1) inspired by the similar argument in [8].

We now use the lower bound of the resonance function to recover the derivative on the non-linear term $u\partial_x u$.

Lemma 4.4. Let

$$F_s = \frac{\mid n \mid^{2s+2} \mid n_1(n-n_1) \mid^{-2s}}{\sigma(\tau, \tau_1, n, n_1)}$$

and

$$F_{s,r} = \frac{\mid n \mid^{2s+2} \mid n_1(n-n_1) \mid^{-2s}}{\sigma^{2(1-r)}(\tau, \tau_1, n, n_1)}$$

where $\sigma(\tau, \tau_1, n, n_1) = \max\{ \mid \tau + m(n) \mid, \mid \tau_1 + m(n_1) \mid, \mid \tau - \tau_1 + m(n - n_1) \mid \}$. Then, for $s \ge -\frac{1}{2}$, $0 < r < \frac{1}{4}$, we have

$$F_s \lesssim 1$$

and

$$F_{s,r} \lesssim \frac{1}{\mid n \mid^{2-4r}}.$$

Proof: This follows from Lemma 4.1.

According to [6] we have the following Lemma:

Lemma 4.5. For any $u \in X^{s,\frac{1}{2}}$ supported in [-T,T] and for any $0 < b < \frac{1}{2}$, it holds:

$$||u||_{X^{s,b}} \lesssim T^{(\frac{1}{2}-b)-}||u||_{X^{s,1/2-}} \lesssim T^{(\frac{1}{2}-b)-}||u||_{X^{s,1/2}}.$$
 (4.3)

Proof of Proposition 4.1: Let

$$P_f^b(n,\tau) = |n|^s < \tau + m(n) >^b |\hat{f}(n,\tau)|,$$

then we have

$$||f||_{X^{b,s}} = (\sum_{n} \int_{\mathbb{R}} (P_f^b(n,\tau))^2 d\tau)^{\frac{1}{2}} = ||P_f^b(n,\tau)||_{l_n^2 L_\tau^2},$$

and

$$B(f,g)(n,\tau) = n^{s+1} < \tau + m(n) >^{-\frac{1}{2}} \sum_{n_1 \neq 0, n_1 \neq n} \int_{\mathbb{R}} \frac{(n_1(n-n_1))^{-s} P_f^{\frac{1}{2}-\gamma}(n_1, \tau_1) P_g^{\frac{1}{2}}(n-n_1, \tau-\tau_1) d\tau_1}{<\tau_1 + m(n_1) >^{\frac{1}{2}-\gamma} < \tau-\tau_1 + m(n-n_1) >^{\frac{1}{2}}}$$

$$(4.4)$$

Denote

$$F(n,\tau,n_1,\tau_1) = \frac{\mid n\mid^{s+1}\mid n_1(n-n_1)\mid^{-s}}{<\tau+m(n)>^{\frac{1}{2}}<\tau_1+m(n_1)>^{\frac{1}{2}}-\gamma}<\tau-\tau_1+m(n-n_1)>^{\frac{1}{2}}}.$$

Letting $E = \{(n, \tau, n_1, \tau_1) : |\tau - \tau_1 + m(n - n_1)| \le |\tau_1 + m(n_1)| \}$, then by symmetry, (4.4) is reduced to estimate

$$\left(\sum_{n\neq 0} \int_{\mathbb{R}} \left(\sum_{n_1\neq n, n_1\neq 0} \int_{\mathbb{R}} (1_E F)(n, \tau, n_1, \tau_1) P_f^{\frac{1}{2} - \gamma} (n - n_1, \tau - \tau_1) P_g^{\frac{1}{2}} (n_1, \tau_1) d\tau_1\right)^2 d\tau\right)^{\frac{1}{2}}.$$
(4.5)

We separate the two cases.

Case I: $|\tau_1 + m(n_1)| \le |\tau + m(n)|$

In this case, the set E is replaced by

$$E_I = \{(n, \tau, n_1, \tau_1) : |\tau - \tau_1 + m(n - n_1)| \le |\tau_1 + m(n_1)| \le |\tau + m(n)| \},$$

then by Cauchy-Schwarz inequality (4.5) is controlled by

$$\left\| \left(\sum_{n_1 \neq n, n_1 \neq 0} \int_{\mathbb{R}} (1_{E_I} F)^2(n, \tau, n_1, \tau_1) d\tau_1 \right)^{\frac{1}{2}} \right. \\ \times \left(\sum_{n_1 \neq n, n_1 \neq 0} \int_{\mathbb{R}} (P_f^{\frac{1}{2} - \gamma})^2(n - n_1, \tau - \tau_1) (P_g^{\frac{1}{2}})^2(n_1, \tau_1) d\tau_1 \right)^{\frac{1}{2}} \right\|_{l_n^2 L_\tau^2}.$$

$$(4.6)$$

Remark that

$$F^2 \approx F_s \frac{1}{\langle \tau_1 + m(n_1) \rangle^{1-2\gamma} \langle \tau - \tau_1 + m(n-n_1) \rangle}$$

with $F_s = \frac{|n|^{2s+2}|n_1(n-n_1)|^{-2s}}{\sigma(\tau,\tau_1,n,n_1)}$, then by Lemma 4.4, for $s \ge -\frac{1}{2}$, $(n,\tau,n_1,\tau_1) \in E_I$, we have

$$\sup_{n,\tau} \sum_{n_1} \int_{\mathbb{R}} (1_{E_I} F)^2(n,\tau,n_1,\tau_1) d\tau_1 \lesssim \sup_{n,\tau} \sum_{n_1} \int_{\mathbb{R}} \frac{d\tau_1}{<\tau_1 + m(n_1) >^{1-2\gamma} < \tau - \tau_1 + m(n-n_1) >}$$

we can easily see that

$$(4.6) \leq \sup_{n,\tau} \sum_{n_1} \int_{\mathbb{R}} \frac{d\tau_1}{\langle \tau_1 + m(n_1) \rangle^{1-2\gamma} \langle \tau - \tau_1 + m(n-n_1) \rangle} \|P_f^{\frac{1}{2}-\gamma}(n,\tau)\|_{l_n^2 L_\tau^2} \|P_g^{\frac{1}{2}}(n,\tau)\|_{l_n^2 L_\tau^2}$$

then by Lemma 4.2, 4.3(take $\alpha = \tau + m(n_1) + m(n - n_1)$ and $\beta = \tau_1 + m(n_1)$) and 4.5 we obtain that there exist $\theta > 0$ such that:

$$(4.5) \lesssim \|f\|_{X^{s,\frac{1}{2}-\gamma}} \|g\|_{X^{s,\frac{1}{2}}} \lesssim T^{\theta} \|f\|_{X^{s,\frac{1}{2}}} \|g\|_{X^{s,\frac{1}{2}}}.$$

Case II: $|\tau + m(n)| \le |\tau_1 + m(n_1)|$ Here the set E becomes:

$$E_{II} = \{ (n, \tau, n_1, \tau_1) : \mid \tau - \tau_1 + m(n - n_1) \mid \leq \mid \tau_1 + m(n_1) \mid, \mid \tau + m(n) \mid < \mid \tau_1 + m(n_1) \mid \}.$$

Then we will estimate

$$\|\sum_{n_1} \int_{\mathbb{R}} (1_{E_{II}} F)(n, \tau, n_1, \tau_1) P_f^{\frac{1}{2} - \gamma}(n - n_1, \tau - \tau_1) P_g^{\frac{1}{2}}(n_1, \tau_1) d\tau_1 \|_{l_n^2 L_\tau^2}$$
(4.7)

By duality, (4.7) equals to

$$\sup_{\|w\|_{l_n^2 L_\tau^2} = 1} \sum_{n, n_1} \int_{\mathbb{R}^2} w(n, \tau) (1_{E_{II}} F)(n, \tau, n_1, \tau_1) P_f^{\frac{1}{2} - \gamma}(n - n_1, \tau - \tau_1) P_g^{\frac{1}{2}}(n_1, \tau_1) d\tau_1 d\tau.$$

By Fubini's Theorem and Cauchy-Schwarz inequality, we could control (4.8) by

$$\sup_{\|w\|_{l_{x}^{2},L_{x}^{2}}=1} \left(\sum_{n_{1}} \int_{\mathbb{R}} \left[\sum_{n} \int_{\mathbb{R}} (1_{E_{II}} F)^{2}(n,\tau,n_{1},\tau_{1}) d\tau \right] \times$$
(4.9)

$$\Big[\sum_{n} \int_{\mathbb{R}} w^{2} (P_{f}^{\frac{1}{2}-\gamma})^{2} (n-n_{1},\tau-\tau_{1}) d\tau\Big] d\tau_{1}\Big)^{\frac{1}{2}} \|g\|_{X^{s,\frac{1}{2}}}.$$

Similary to the previous case, we can show that:

$$\sup_{n_1, \tau_1} \sum_{n} \int_{\mathbb{R}} (1_{E_{II}} F)^2(n, \tau, n_1, \tau_1) d\tau \lesssim 1.$$

Finaly we obtain that

$$(4.9) \lesssim \|f\|_{X^{s,\frac{1}{2}-\gamma}} \|g\|_{X^{s,\frac{1}{2}}} \lesssim T^{\theta} \|f\|_{X^{s,\frac{1}{2}}} \|g\|_{X^{s,\frac{1}{2}}}.$$

Now we have the following proposition:

Proposition 4.2. Let $s \ge -\frac{1}{2}$ then for all f, g with compact support in time included in the subset $\{(t,x), t \in [-T,T]\}$, there exists $\theta > 0$ such that:

$$\left(\sum_{n \in \mathbb{Z}} |n|^{2s} \left[\int_{\mathbb{R}} \frac{|n\hat{f} * \hat{g}(n,\tau)|}{\langle \tau + m(n) \rangle} d\tau \right]^{2} \right)^{\frac{1}{2}} \lesssim T^{\theta} ||f||_{X^{s,\frac{1}{2}}} ||g||_{X^{s,\frac{1}{2}}}. \tag{4.10}$$

Proof: As in the proof of Prop 4.1, we consider (4.10) in the same two cases. It could be written as:

$$\left\| \int_{\mathbb{R}} \sum_{n_1} \int_{\mathbb{R}} (1_E F)(., \tau, n_1, \tau_1) P_f^{\frac{1}{2} - \gamma}(. - n_1, \tau - \tau_1) P_g^{\frac{1}{2}}(n_1, \tau_1) d\tau_1 d\tau \right\|_{l_n^2} \lesssim T^{\theta} \|f\|_{X^{s, \frac{1}{2}}} \|g\|_{X^{s, \frac{1}{2}}},$$

$$(4.11)$$

where

$$F(n,\tau,n_1,\tau_1) = \frac{\mid n\mid^{s+1}\mid n_1(n-n_1)\mid^{-s}}{<\tau + m(n)>^{\frac{1}{2}}<\tau_1 + m(n_1)>^{\frac{1}{2}}-\gamma}<\tau - \tau_1 + m(n-n_1)>^{\frac{1}{2}}}.$$

1)Case I: $|\tau_1 + m(n_1)| \le |\tau + m(n)|$. As before, the set E is replaced by

$$E_I = \{ | \tau - \tau_1 + m(n - n_1) | \le | \tau_1 + m(n_1) | \le | \tau + m(n) | \}.$$

By duality, we suffer to estimate

$$\sup_{\|w\|_{l_{p}^{2}}=1} \sum_{n,n_{1}} \int_{\mathbb{R}^{2}} w(n) (1_{E_{I}} F)(n,\tau,n_{1},\tau_{1}) P_{f}^{\frac{1}{2}-\gamma}(n-n_{1},\tau-\tau_{1}) P_{g}^{\frac{1}{2}}(n_{1},\tau_{1}) d\tau_{1} d\tau.$$

Now by Cauchy-Schwarz, we could control it by

$$\sup_{\|w\|_{l_n^2}=1} \Big(\sum_{n_1} \int_{\mathbb{R}} \Big[\sum_{n} \int_{\mathbb{R}} (1_{E_I} F)^2(n, \tau, n_1, \tau_1) d\tau \Big] \times \Big[\sum_{n} \int_{\mathbb{R}} w^2 (P_f^{\frac{1}{2} - \gamma})^2(n - n_1, \tau - \tau_1) d\tau \Big] d\tau_1 \Big)^{\frac{1}{2}} \|g\|_{X^{s, \frac{1}{2}}},$$

then it is sufficient to show that, for $s \geqslant -\frac{1}{2}$

$$D = \sup_{n_1} \sum_{n} \int_{\mathbb{R}} \int_{\mathbb{R}} (1_{E_I} F)^2(n, \tau, n_1, \tau_1) d\tau d\tau_1 \lesssim 1.$$

For some $0 < r < \frac{1}{4}$, D can be rewriten as:

$$D = \sup_{n_1} \sum_{n} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{1}{\langle \tau + m(n) \rangle^{2r}}, (1_{E_I} F_r)^2 (n, \tau, n_1, \tau_1) d\tau d\tau_1$$

where

$$F_r^2 = \frac{\mid n \mid^{2s+2} \mid n_1(n-n_1) \mid^{-2s}}{\langle \tau + m(n) \rangle^{2(1-r)}} \frac{1}{\langle \tau_1 + m(n_1) \rangle^{1-2\gamma} \langle \tau - \tau_1 + m(n-n_1) \rangle}.$$

Remark that

$$F_r^2 = F_{s,r} \frac{1}{\langle \tau_1 + m(n_1) \rangle^{1-2\gamma} \langle \tau - \tau_1 + m(n - n_1) \rangle},$$

then by Lemma 4.4, D could be controlled by

$$D \lesssim \sup_{n_1} \sum_{n} \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{1}{|n|^{2-4r}} \frac{d\tau_1 d\tau}{\langle \tau_1 + m(n_1) \rangle^{1-2\gamma+r} \langle \tau - \tau_1 + m(n-n_1) \rangle^{1+r}}$$

by Lemma 4.2 we have:

$$\int_{\mathbb{R}} \frac{d\tau_1}{<\tau_1 + m(n_1) >^{1-2\gamma+r} < \tau - \tau_1 + m(n-n_1) >^{1+r}} \lesssim \frac{1}{(1+\mid \tau + m(n_1) + m(n-n_1)\mid)^{1-2\gamma+r}}.$$

$$D \lesssim \sup_{n_1} \sum_{n} \frac{1}{|n|^{2-4r}} \int_{\mathbb{R}} \frac{d\tau}{(1+|\tau+m(n_1)+m(n-n_1)|)^{1-2\gamma+r}}.$$

Therefore, if $r < \frac{1}{4}$, we have $D \lesssim \sum_{n} \frac{1}{|n|^{2-4r}} < +\infty$. 2)Case II, $|\tau + m(n)| \le |\tau_1 + m(n_1)|$. Now we replace E with

$$E_{II} = \{ (n, \tau, n_1, \tau_1) : | \tau - \tau_1 + m(n - n_1) | \le | \tau_1 + m(n_1) |, | \tau + m(n) | < | \tau_1 + m(n_1) | \}.$$

We write

$$1+ |\tau + m(n)| = (1+ |\tau + m(n)|)^{r} (1+ |\tau + m(n)|)^{1-r}$$

where $\frac{1}{2} < r < 1$. As in case I, F_r denotes

$$\frac{\mid n\mid^{s+1}\mid n_1(n-n_1)\mid^{-s}}{<\tau+m(n)>^{(1-r)}} \frac{1}{<\tau_1+m(n_1)>^{\frac{1}{2}-\gamma}<\tau-\tau_1+m(n-n_1)>^{\frac{1}{2}}}.$$

It suffices to estimate

$$\left(\sum_{n} \left(\int_{\mathbb{R}} \sum_{n_{1}} \int_{\mathbb{R}} \frac{1}{\langle \tau + m(n) \rangle^{r}} (1_{E_{II}} F_{r})(n, \tau, n_{1}, \tau_{1}) P_{f}^{\frac{1}{2} - \gamma}(n - n_{1}, \tau - \tau_{1}) P_{g}^{\frac{1}{2}}(n_{1}, \tau_{1}) d\tau_{1} d\tau\right)^{2}\right)^{\frac{1}{2}}.$$

$$(4.12)$$

Applying the Cauchy-Schwarz inequality in τ we see that (4.12) is bounded

$$\left[\sum_{n} \left(\int_{\mathbb{R}} \frac{d\tau}{\langle \tau + m(n) \rangle^{2r}} \right) \right]$$

$$\times \left(\int_{\mathbb{R}} \left(\sum_{n_1} \int_{\mathbb{R}} (1_{E_{II}} F_r)(n, \tau, n_1, \tau_1) P_f^{\frac{1}{2} - \gamma}(n - n_1, \tau - \tau_1) P_g^{\frac{1}{2}}(n_1, \tau_1) d\tau_1 \right)^2 d\tau \right) \right]^{\frac{1}{2}}.$$

Since 2r > 1, then (4.12) is dominated by

$$\left[\sum_{n} \int_{\mathbb{R}} \left(\sum_{n_{1}} \int_{\mathbb{R}} (1_{E_{II}} F_{r})(n, \tau, n_{1}, \tau_{1}) P_{f}^{\frac{1}{2} - \gamma}(n - n_{1}, \tau - \tau_{1}) P_{g}^{\frac{1}{2}}(n_{1}, \tau_{1}) d\tau_{1}\right)^{2} d\tau\right]^{\frac{1}{2}},$$

then as the case II in the proof of Prop 4.1, we obtain the estimate, and this end the proof.

Now we return to the proof of **Theorem 2.1**: Let L defined by

$$L(u) = \psi(t)[S(t)\phi - \int_0^t S(t - t')\partial_x(\psi_T^2 u^2(t'))dt'], \tag{4.13}$$

where $t \in \mathbb{R}$, ψ indicates a time cutoff function :

$$\psi \in C_0^{\infty}(\mathbb{R}), \quad \sup \psi \subset [-2, 2], \quad \psi = 1 \text{ on } [-1, 1],$$
 (4.14)

$$\psi_T(.) = \psi(./T),$$

we will apply a fixed point argument to (4.13), using the following estimates:

Proposition 4.3. There exists a constant $C = C(\phi)$ such that:

$$||L(u)||_{Y^s} \le C||\phi||_{H^s} + CT^{\gamma}||u||_{Y^s}^2$$

and

$$||L(u) - L(v)||_{Y^s} \le CT^{\gamma}||u - v||_{Y^s}||u + v||_{Y^s}.$$

Proof: It follows from Propositions 4.1, 4.2 and classical linear estimates (see [3]).

Note that if we take $T = (4C^2||\phi||_{H^s})^{-1/\gamma}$ we deduce from Prop 4.3 that L is strictly contractive in the ball $B(0, \frac{1}{8C^2})$ in Y^s . This proves the existence of a unique solution u to (4.13) in Y^s .

4.1. Global existence in L^2 . Its easy to see that the L^2 -norm is conserved ($||u(t)||_{L^2} = ||u_0||_{L^2}$). Hence, if we take an initial data u_0 in L_0^2 , the solution u such that $u(0) = u_0$ can be extended for all positive times and the existence is global in $L_0^2(\mathbb{T})$.

References

- [1] A.A.Arsen'ev, On the invariant measures for the classical dynamical systems with infinite-dimensional phase space. Matem. Sbornik. 121, 297 (1983). (In Russian.)
- [2] J. Bourgain, Periodic nonlinear Schrödinger equation and invariant measures, Comm. Math. Phys. 166 (1994), 126.
- [3] J. Colliander, M. Keel, G. Staffilani, H. Takaoka, T. Tao, Sharp Global Well-Posedness for KdV and Modified KdV on R and T, J. Amer. Math. Soc. 16 (2003), no. 3, 705749.
- [4] Yu. L. Daletskii and S. V. Fomin, Measures and Differential Equations in Infinite-Dimensional Spaces, Nauka, Moscow, 1983.
- [5] M.X. Fernique, Intégrabilité des Vecteurs Gaussiens, Academie des Sciences, Paris, Comptes Rendus, 270, Séries A (1970), 1698-1699.
- [6] J. Ginibre, Y. Tsutsumi, G. Velo, On the Cauchy Problem for the Zakharov System. J. Funct. Anal., 133 (1995) 50-68.
- [7] H. Kuo, Gaussian Measures in Banach Spaces, Lec. Notes in Math. 463, Springer-Verlag, New York, 1975.
- [8] 15. C. Kenig, G. Ponce and L. Vega. A bilinear estimate with applications to the KdV equation. J. Amer. Math. Soc., 9(1996).
- [9] L.A. Ostrovskii, Nonlinear internal waves in a rotating ocean. Okeanologiya 18 (2) (1978) 181-191.
- [10] P.E.Zhidkov, On invariant measures infinite-dimensional for some dynamical systems. Ann. Inst. Henri Poincare 62(3)1995,267-287
- [11] P. E. Zhidkov, An invariant measure for a nonlinear wave Equation. Nonlinear Analysis: Theory, Methods & Applications Vol. 22, No. 3, pp. 319-325, 1994.
- [12] P.E. Zhidkov, Korteweg-de Vries and Nonlinear Schrödinger Equations: Qualitative theory. Springer, 2001.

M. Darwich: Université François rabelais de Tours, Laboratoire de Mathématiques et Physique Théorique, UMR-CNRS 7350, Parc de Grandmont, 37200 Tours, France

E-mail address: Mohamad.Darwich@univ-tours.fr