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# The influence of boundary conditions on the contact problem in a 3D Navier-Stokes Flow 

David Gerard-Varet; Matthieu Hillairet ${ }^{\dagger}$ and Chao Wang ${ }^{\ddagger}$

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#### Abstract

We consider the free fall of a sphere above a wall in a viscous incompressible fluid. We investigate the influence of boundary conditions on the finite-time occurrence of contact between the sphere and the wall. We prove that slip boundary conditions enable to circumvent the "no-collision" paradox associated with no-slip boundary conditions. We also examine the case of mixed boundary conditions.


## 1 Introduction

Understanding the close-range dynamics of particles inside a fluid is crucial to several fields of application, such as rheology, sedimentation or slurry erosion. A related problem is to compute the drag induced by nearby solid bodies: it is important to lubrication theory, or to the control of microswimmers, among others. In the case of a viscous incompressible fluid, a standard way to model the fluid-solid interaction is through the coupling of the Navier-Stokes equations for the fluid part and the conservation of momentum for the bodies. However, such model reveals unrealistic features: in particular, it can overvalue the drag force induced on bodies that are close to collision. Consider for instance the free fall of a rigid sphere over a wall, and assume standard no-slip conditions at the solid surface of the sphere and the wall. Then, the model predicts that no collision is possible between the sphere and the wall, no matter their relative density and the viscosity of the fluid.This is a "no-collision paradox": it has been recognized at first for reduced models (Stokes approximation, see [5, 6]), and more recently for the full constant-density incompressible Navier Stokes system [13, 19, 21]. These results show the need for a more accurate model of the fluid behavior in the "small-distance" regime. Several refinements were proposed in the past [ $1,8,27$ ], see the introduction of [26] for a review. In this paper, we investigate the influence of boundary conditions at the fluid-solid interfaces.

[^0]Boundary conditions are constitutive equations that can be derived theoretically or experimentally. A popular general type of boundary conditions was proposed by Navier in the early time of analytical fluid dynamics. This Navier law states that, on the one hand, the normal component of the velocity is continuous at the fluid/solid interface (impermeability condition) and that, on the other hand, the amount of slip in the tangential part of the velocity is proportional to the tangential part of the normal stress exerted by the fluid on the boundary. The proportionality coefficient is called "slip length". Assuming the slip length vanishes (i.e., assuming no-slip boundary conditions) suffices to make theoretical and experimental results to coincide in numerous cases. However, developments such as the "no-collision paradox" lead to question this assumption [23]. We refer the reader to [2, 28] for a further description of experimental and theoretical validations of the Navier boundary conditions with slip. We propose here new insights on the contact-problem, by considering the effect of Navier boundary conditions. We shall study this effect in the context of the full Navier-Stokes system, which is, up to our knowledge, new.

We consider a simplified case where only one homogeneous solid is moving in the fluid. We denote $S(t) \subseteq \mathbb{R}^{3}, F(t) \subseteq \mathbb{R}^{3}$ the solid and the fluid domains at time $t$, and $\Omega \triangleq \overline{S(t)} \cup F(t)$ the total domain. We assume that the fluid is governed by the constant-density incompressible Navier-Stokes equations :

$$
\begin{align*}
\rho_{F}\left(\partial_{t} u_{F}+u_{F} \cdot \nabla u_{F}\right)-\mu_{F} \Delta u_{F} & =-\nabla p_{F}-\rho_{F} g e_{3}, & & t>0,  \tag{1.1}\\
\operatorname{div} u_{F} & =0, & & t>0, \tag{1.2}
\end{align*}
$$

where $u_{F}$ and $p_{F}$ stand respectively for the velocity and internal pressure, $\rho_{F}$ the density, $\mu_{F}$ the viscosity in the fluid. We emphasize that we consider here the influence of gravity whose direction is the third vector of the basis $\left(e_{1}, e_{2}, e_{3}\right)$, i.e., the "vertical" unit vector. The solid motion is governed by Newton laws. Denoting $x_{S}(t) \in \mathbb{R}^{3}$ the position of the center of mass of the solid at time $t, U_{S}(t) \in \mathbb{R}^{3}$ its velocity and $\omega_{S}(t) \in \mathbb{R}^{3}$ the angular velocity, the Newton laws read

$$
\begin{align*}
m_{S} \frac{d}{d t} U_{S}(t) & =-\int_{\partial S(t)}\left(2 \mu_{F} D\left(u_{F}\right)-p_{F} I\right) n d \sigma-m_{S} g e_{3}  \tag{1.3}\\
\frac{d}{d t}\left(J_{S} \omega_{S}(t)\right) & =-\int_{\partial S(t)}\left(x-x_{S}(t)\right) \times\left(2 \mu_{F} D\left(u_{F}\right)-p_{F} I\right) n d \sigma \tag{1.4}
\end{align*}
$$

Here $\rho_{S}$ and $m_{S} \triangleq \rho_{S}|S(0)|$ are the density and the mass of the solid respectively. Symbol $n$ stands for the unit normal vector to $\partial S(t)$ pointing outside the fluid domain. The $3 \times 3$ inertia matrix $J_{S}(t)$ is defined by

$$
J_{S}(t) \triangleq \rho_{S} \int_{S(t)}\left(\left|x-x_{S}(t)\right|^{2} I_{d}-\left(x-x_{S}(t)\right) \otimes\left(x-x_{S}(t)\right)\right) d x
$$

We emphasize that the solid is homogeneous, so that no angular momentum is induced by gravity in (1.4). Given $\left(U_{S}, \omega_{S}, x_{S}\right)$, the solid velocity-field $u_{S}$ at each point $x \in S(t)$ reads

$$
\begin{equation*}
u_{S}(t, x) \triangleq U_{S}(t)+\omega_{S} \times\left(x-x_{S}(t)\right) \tag{1.5}
\end{equation*}
$$

In this paper, we complete the system with Navier boundary conditions:

$$
\begin{align*}
\left.\left(u_{F}-u_{S}\right) \cdot n\right|_{\partial S(t)} & =0, & \left(u_{F}-u_{S}\right) \times\left. n\right|_{\partial S(t)} & =-2 \beta_{S}\left(D\left(u_{F}\right) n\right) \times n,  \tag{1.6}\\
\left.u_{F} \cdot n\right|_{\partial \Omega} & =0, & u_{F} \times\left. n\right|_{\partial \Omega} & =-2 \beta_{\Omega}\left(D\left(u_{F}\right) n\right) \times n, \tag{1.7}
\end{align*}
$$

where $\beta_{S}, \beta_{\Omega} \geq 0$ are the slip lengths at the fluid-solid and fluid-container interface respectively. We remind that assuming that $\beta_{S}$ or $\beta_{\Omega}$ vanishes means assuming no-slip boundary conditions at the corresponding interface.

System (1.1)-(1.7) is completed with initial conditions :

$$
\begin{align*}
& S(0)=S_{0}, \quad u_{F}(0, x)=u_{F, 0}(x), \quad x \in \Omega \backslash \bar{S}_{0}  \tag{1.8}\\
& U(0)=U_{0}, \quad \omega(0)=\omega_{0} . \tag{1.9}
\end{align*}
$$

Equation (1.9) may also be written :

$$
\begin{equation*}
u_{S}(0, x)=u_{S, 0}(x), \quad x \in S_{0} . \tag{1.9}
\end{equation*}
$$

where $u_{S, 0}$ is computed with respect to $U_{0}, \omega_{0}$ and $S_{0}$ via (1.5). When $S_{0} \subset \Omega$ and $u_{F, 0}$ has finite energy, namely $u_{F, 0} \in L^{2}\left(\Omega \backslash \bar{S}_{0}\right)$, standard computations show that a reasonable solution to (1.1)-(1.9) should satisfy the energy estimate :

$$
\begin{align*}
& \frac{1}{2}\left[\int_{F(t)} \rho_{F}\left|u_{F}(t)\right|^{2}+\int_{S(t)} \rho_{S}\left|u_{S}(t)\right|^{2}\right] \\
+ & \mu_{F}\left(2 \int_{0}^{t} \int_{F(t)}\left|D\left(u_{F}(t)\right)\right|^{2}+\frac{1}{\beta_{\Omega}} \int_{0}^{t} \int_{\partial \Omega}\left|u_{F} \times n\right|^{2} \mathrm{~d} \sigma+\frac{1}{\beta_{S}} \int_{0}^{t} \int_{\partial S(t)}\left|\left(u_{F}-u_{S}\right) \times n\right|^{2} \mathrm{~d} \sigma\right) \\
& \leq-\int_{0}^{t} \int_{F(t)} \rho_{F} g e_{3} \cdot u_{F}-\int_{0}^{t} \int_{S(t)} \rho_{S} g e_{3} \cdot u_{S}+\frac{1}{2} \int_{F_{0}} \rho_{F}\left|u_{F, 0}\right|^{2}+\frac{1}{2} \int_{S_{0}} \rho_{S}\left|u_{S}(0)\right|^{2} . \tag{1.10}
\end{align*}
$$

where $F_{0} \triangleq \Omega \backslash \bar{S}_{0}$. When $\beta_{S}=0$ or $\beta_{\Omega}=0$, the corresponding boundary term vanishes.
Roughly, with regards to the classical theory of the Navier-Stokes equation in a fixed domain, one can expect two kinds of solutions for (1.1)-(1.9):

- Strong solutions, locally in time: they should be unique, with possible blow-up prior to collision in 3D.
- Weak solutions: they should exist at least up to collision, and be possibly non-unique.

In the no-slip case ( $\beta_{S}=\beta_{\Omega}=0$ ), both kinds of solutions have been built in various contexts: 2D/3D problems, bounded/unbounded container, smooth/singular shape of the solid bodies. Existence of strong solutions locally in time (prior to collision) is now well-known, see $[9$, $12,13,17,32,33]$. The existence of weak solutions up to collision has been studied in $[4,7,10,18,24,25,30,34]$. Uniqueness of weak solutions up to collision in 2D is established in the recent paper [16]. Finally, extension of weak solutions after contact is given in [29] for 2D problems and [11] for 3D problems.

In the case of slip conditions, when $\beta_{S}>0$ and $\beta_{\Omega}>0$, much less in known. It is partly because genuine new mathematical difficulties arise: for instance, at the level of weak solutions, discontinuity of the tangential velocity at the fluid/solid interface forbids global $H^{1}$ bound on the extended velocity-field $u$ as defined by:

$$
\begin{equation*}
u=u_{S} \text { in } S(t), \quad u=u_{F} \text { in } F(t) . \tag{1.11}
\end{equation*}
$$

Such global bound is a key point in the treatment of the Dirichlet case. Hence, to investigate the contact problem for system (1.1)-(1.9), we shall rely on the recent paper [15], which shows the existence of weak solutions up to collision. Both the definition of weak solutions and the existence result given in [15] will be recalled in the next section. Let us also mention the work [31], related to a single body without container.

Considering a special class of weak solutions, we will show that slip conditions allow collision in finite time. Hence, taking into account slight slip at the solid boundaries allows to clear the no-collision paradox. More precisely, we shall consider configurations obeying the following assumptions:

A1. The solid $S(t) \sim S_{0}$ is a ball of radius 1 .
A2. $\Omega$ is a smooth convex domain of $\left\{x_{3}>0\right\}$, flat near:

$$
\partial \Omega \supset\left\{x_{3}=0,|x|<2 \delta\right\} \text { for some } \delta \in\left(0, \frac{1}{4}\right) .
$$

A3. $S(t)$ is axisymmetric with respect to $\left\{x_{1}=x_{2}=0\right\}$.
A4. The only possible contact between the solid and the container is at $x=0$. More precisely,

$$
\operatorname{dist}\left(S(t), \partial \Omega \backslash\left\{x_{3}=0,|x|<\delta\right\}\right) \geq d_{\delta}>0
$$

Parameters $\delta$ and $d_{\delta}$ will be involved in the proof. Let us comment on these assumptions. Their only goal is to ensure that the geometry of the potential contact zone does not change with time. Namely, it is an aperture between a sphere and a plane, only varying with the distance $h(t)$ between the south pole of $S(t)$ and $x=0$. The whole point is then to determine if $h(t)$ can vanish in finite time or not. This simplifies greatly our computations. However, as will be clear from our proof, our result is genuinely local in space time: it is independent from the global characteristics of the solid/container geometry.

Note that A1-A2 are just assumptions on $S_{0}$ and $\Omega$. If $S_{0} \subset \subset \Omega$, and if the initial data ( $u_{F, 0}, u_{S, 0}$ ) has finite energy, article [15] ensures the existence of a weak solution up to contact between $S(t)$ and $\partial \Omega$ (see next section). Assumption A3 states that the axisymmetry of the solid is preserved with time. Actually, it will be automatically satisfied if $S_{0}, \Omega$ and the initial velocities are axisymmetric. More precisely, in this case, one can adapt the construction of weak solutions from [15] by including the symmetry constraint, and obtain directly an axisymmetric weak solution. Finally, assumption A4 guarantees that the sphere remains away from the top and lateral boundaries. Again, this can be ensured by a proper choice of the container and the initial data. As regards lateral boundaries, it is enough by assumption A3 that they are at distance $>1$ to the vertical axis $\left\{x_{1}=x_{2}=0\right\}$. Also, there are several ways to avoid merging of the solid body and the top boundary. One can for instance assume that the solid is heavier than the fluid $\left(\rho_{S}>\rho_{F}\right)$, and that the initial kinetic energy of the fluid and the solid is small enough. As the total energy of weak solutions is non-increasing, this means that the increase of the potential energy through time can only be small: hence, the solid will not rise too much, and so will remain away from the top boundary. Again, we insist that our method of proof could extend to more general settings.

We can now state our main results:

Theorem 1.1 (Slip Case). Assume $\beta_{S}, \beta_{\Omega}>0$ and $\rho_{S}>\rho_{F}$. For any weak solution ( $S, u$ ) (see Definition 2.1) satisfying A1-A4, the solid body $S$ touches $\partial \Omega$ in finite time.

Theorem 1.2 (Mixed Case). Let $\beta_{S}$ or $\beta_{\Omega}$ vanish. Any weak solution ( $S, u$ ) (see Definition 2.3) satisfying A1-A4 is global. In particular, the solid body $S$ never touches $\partial \Omega$.

We recall that the full no-slip case (when $\beta_{S}=\beta_{\Omega}=0$ ) was already treated in [21]. Theorems 1.1 and 1.2 show that one possible way to obtain realistic contact, and to circumvent the no-collision paradox, is to introduce slip in the boundary conditions on all solid boundaries of the fluid domain.

We conclude with a few hints on the proof of Theorem 1.1 (ideas for Theorem 1.2 are similar). The proof elaborates on the paper [14], devoted to a simplified linear system: one takes there $\Omega=\mathbf{R}_{+}^{3}$, while the fluid outside the sphere $S(t)$ is governed by a steady Stokes flow (with $\mu_{F}=1$ for simplicity):

$$
-\Delta u_{F}=-\nabla p_{F}-\rho_{F} g e_{3}, \quad \nabla \cdot u_{F}=0, \quad x \in F(t)
$$

In this simplified setting, the fluid and solid domains $F(t)=F_{h(t)}$ and $S(t)=S_{h}(t)$ are characterized by the distance $h(t)$ beween the sphere and the plane wall. Then, it is easily seen that $u_{F}(t)=h^{\prime}(t) w_{h(t)}$ where $w_{h}$ is the solution of a normalized Stokes equation, set in the domain $F_{h}$ "frozen" at distance $h$ :

$$
\left\{\begin{align*}
-\Delta w_{h}+\nabla p_{h}=0, & \nabla \cdot w_{h}=0, \quad x \in F_{h},\left.\quad\left(w_{h}-e_{3}\right) \cdot n\right|_{\partial S_{h}}=0,\left.\quad w_{h} \cdot n\right|_{\partial \Omega}=0  \tag{1.12}\\
\left(w_{h}-e_{3}\right) \times\left. n\right|_{\partial S_{h}} & =-2 \beta_{S}\left(D\left(w_{h}\right) n\right) \times n \\
w_{h} \times\left. n\right|_{\partial \Omega} & =-2 \beta_{\Omega}\left(D\left(w_{h}\right) n\right) \times n
\end{align*}\right.
$$

We refer to [14] for all details. Eventually, the dynamics reduces to an ODE of the type

$$
h^{\prime \prime}(t)=-h^{\prime}(t) \mathcal{D}(h(t))+\frac{\rho_{S}-\rho_{F}}{\rho_{S}} g
$$

where the drag term $\mathcal{D}(h)$ is the energy of $w_{h}$ :

$$
\mathcal{D}(h)=\mathcal{E}_{h}\left(w_{h}\right), \quad \mathcal{E}_{h}(u):=\int_{F_{h}}|\nabla u|^{2}+\left(\frac{1}{\beta_{S}}+1\right) \int_{\partial S_{h}}\left|\left(u-e_{3}\right) \times n\right|^{2}+\frac{1}{\beta_{\Omega}} \int_{\partial \Omega}|u \times n|^{2} .
$$

For this reduced model, the contact problem resumes to the computation of the drag $\mathcal{D}_{h}$ in the limit $h \rightarrow 0$. The difficulty is that there is no simple formula for the solution $w_{h}$ of the Stokes system in $F_{h}$.

A general approach to this problem was recently proposed in [14]. It works for no-slip conditions, slip ones, and for more general geometries. The starting idea is to use a variational characterization of the drag: one can identify $\mathcal{D}(h)$ as the minimizer of an energy $\mathcal{E}_{h}$, over an admissible set of fields $\mathcal{A}_{h}$ :

$$
\mathcal{D}(h)=\inf _{u \in \mathcal{A}_{h}} \mathcal{E}_{h}(u)=\mathcal{E}_{h}\left(w_{h}\right)
$$

Then, instead of computing the true minimizer, the point is to find an appropriate relaxed minimization problem, for which the minimizer can be easily computed, and such that the corresponding minimum is close to the exact one asymptotically in $h$. The choice of such
relaxed problem is developed in various situations in [14]. Roughly, in the case of slip conditions, one obtains that $c|\ln (h)| \leq D(h) \leq C|\ln (h)|$, in agreement with older computations of Hocking [23]. Hence, the drag is weak enough to allow for collision, within this linear Stokes approximation. See also [20] for numerical insight.

In the present paper, we manage to use ideas of the linear study, in the context of the nonlinear Navier-Stokes flow. At a formal level, the first idea is to multiply the momentum equation by $w(t, x)=w_{h(t)}(x), w_{h}$ solution of (1.12) (where this time $\Omega$ is the container). After formal manipulations, one obtains

$$
\begin{equation*}
h^{\prime \prime}(t)=-h^{\prime} \mathcal{D}(h)+\frac{\rho_{S}-\rho_{F}}{\rho_{S}} g+\text { remainder } \tag{1.13}
\end{equation*}
$$

where the remainder term involves the test function $w$ and the convective derivative of $u$. However, the Stokes solution $w_{h}$ is not known explicitly. Hence, the main point in our proof is to use another test function, based on the relaxed minimizer introduced in [14]. Then, one must show that for such test function, the remainder term is indeed small, that is controlled by the gravity when the solid is heavier than the fluid. Let us stress that several difficulties arise in this process. For instance, the relaxed minimizer introduced in the linear study is not accurate enough, so that we must improve it. Moreover, due to the slip conditions, the natural quantity is the symmetric gradient of the weak solution, not the full gradient. This makes the derivation of the various bounds more difficult than in the no-slip case.

The outline of the paper is as follows. In the next section, we introduce a suitable notion of weak solutions for system (1.1)-(1.9), in the slip case and in the mixed case. We recall the existence theorems up to collision available for such weak solutions. in the slip case and in the mixed case. The two last sections are devoted to the proofs of the two theorems.

## 2 On the definition of weak solutions

We introduce here a definition of weak solutions for (1.1)-(1.8) in two cases: slip boundary conditions at all boundaries, mixed slip/no-slip boundary conditions. We mostly follow [15]. Given $\Omega$ a Lipschitz domain, we set:

$$
\begin{array}{rlrl}
\mathcal{D}_{\sigma}(\Omega) & \triangleq\left\{\varphi \in C_{c}^{\infty}(\Omega), \quad \operatorname{div} \varphi=0\right\}, & & \mathcal{D}_{\sigma}(\bar{\Omega}) \triangleq\left\{\left.\varphi\right|_{\Omega}, \varphi \in \mathcal{D}_{\sigma}\left(\mathbb{R}^{3}\right)\right\}, \\
L_{\sigma}^{2}(\Omega) & \triangleq \text { the closure of } \mathcal{D}_{\sigma}(\Omega) \text { in } L^{2}(\Omega), & H_{\sigma}^{1}(\Omega) \triangleq H^{1}(\Omega) \cap L_{\sigma}^{2}(\Omega) \\
H_{\sigma}^{1}(\bar{\Omega}) & \triangleq \text { the closure of } \mathcal{D}_{\sigma}(\bar{\Omega}) \text { in } H^{1}(\Omega), & & \\
\mathcal{R} & \triangleq\left\{\varphi_{s}, \quad \varphi_{s}=V+\omega \times x, \quad\right. \text { for some } & \left.V \in \mathbb{R}^{3}, \omega \in \mathbb{R}^{3}\right\},
\end{array}
$$

### 2.1 The case of slip boundary conditions: $\beta_{S}>0$ and $\beta_{\Omega}>0$.

We consider here that the slip lengths $\beta_{S}$ and $\beta_{\Omega}$ do not vanish i.e., that the fluid slips on both boundaries. The corresponding Cauchy problem is studied in [15], where the following definition of weak solution is given:

Definition 2.1. Let $\Omega$ and $S_{0} \Subset \Omega$ be two Lipschitz bounded domains of $\mathbb{R}^{3}$. Let $u_{F, 0} \in$ $L_{\sigma}^{2}(\Omega), u_{S, 0} \in \mathcal{R}$ such that $u_{F, 0} \cdot n=u_{S, 0} \cdot n$ on $\partial S_{0}$. Let $T>0$, or $T=\infty$.

A weak solution of (1.1)-(1.9) on $[0, T)$ is a pair $(S, u)$ satisfying

1) $S(t) \Subset \Omega$ is a bounded domain of $\mathbb{R}^{3}$ for all $t \in[0, T)$ such that

$$
\chi_{S}(t, x) \triangleq 1_{S(t)}(x) \in L^{\infty}((0, T) \times \Omega)
$$

2) $u$ belongs to the space

$$
\mathcal{S}_{T} \triangleq\left\{u \in L^{\infty}\left(0, T ; L_{\sigma}^{2}(\Omega)\right), \text { there exists } u_{F} \in L_{l o c}^{2}\left([0, T) ; H_{\sigma}^{1}(\Omega)\right), u_{S} \in L^{\infty}(0, T ; \mathcal{R})\right.
$$

$$
\text { such that } \left.u(t, \cdot)=u_{F}(t, \cdot) \text { on } F(t), \quad u(t, \cdot)=u_{S}(t, \cdot) \text { on } S(t), \text { for all } t \in[0, T)\right\}
$$

where $F(t) \triangleq \Omega \backslash \overline{S(t)}$.
3) For all $\varphi$ belonging to the space

$$
\mathcal{T}_{T}^{s} \triangleq\left\{\varphi \in C\left([0, T] ; L_{\sigma}^{2}(\Omega)\right), \text { there exists } \varphi_{F} \in \mathcal{D}\left([0, T) ; \mathcal{D}_{\sigma}(\bar{\Omega})\right), \varphi_{S} \in \mathcal{D}([0, T) ; \mathcal{R})\right.
$$

$$
\text { such that } \left.\varphi(t, \cdot)=\varphi_{F}(t, \cdot) \text { on } F(t), \quad \varphi(t, \cdot)=\varphi_{S}(t, \cdot) \text { on } S(t), \text { for all } t \in[0, T]\right\}
$$

there holds

$$
\begin{align*}
& -\int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \cdot \partial_{t} \varphi_{F}-\int_{0}^{T} \int_{S(t)} \rho_{S} u_{S} \cdot \partial_{t} \varphi_{S}-\int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \otimes u_{F}: \nabla \varphi_{F} \\
& +\frac{\mu_{F}}{\beta_{\Omega}} \int_{0}^{T} \int_{\partial \Omega}\left(u_{F} \times n\right) \cdot\left(\varphi_{F} \times n\right)+\frac{\mu_{F}}{\beta_{S}} \int_{0}^{T} \int_{\partial S(t)}\left(\left(u_{F}-u_{S}\right) \times n\right) \cdot\left(\left(\varphi_{F}-\varphi_{S}\right) \times n\right) \\
& +\int_{0}^{T} \int_{F(t)} 2 \mu_{F} D\left(u_{F}\right): D\left(\varphi_{F}\right)=-\int_{0}^{T} \int_{F(t)} \rho_{F} g e_{3} \varphi_{F}-\int_{0}^{T} \int_{S(t)} \rho_{S} g e_{3} \varphi_{S} \\
& +\int_{F_{0}} \rho_{F} u_{F, 0} \varphi_{F}(0, \cdot)+\int_{S_{0}} \rho_{S} u_{S, 0} \varphi_{S}(0, \cdot) \tag{2.14}
\end{align*}
$$

4) For all $\psi \in \mathcal{D}([0, T) ; \mathcal{D}(\bar{\Omega}))$, we have

$$
-\int_{0}^{T} \int_{S(t)} \partial_{t} \psi-\int_{0}^{T} \int_{S(t)} u_{S} \cdot \nabla \psi=\int_{S_{0}} \psi(0, \cdot)
$$

5) We have the following energy inequality for a.a. $t \in(0, T)$ :

$$
\begin{align*}
& \frac{1}{2}\left[\int_{F(t)} \rho_{F}\left|u_{F}(t)\right|^{2}+\int_{S(t)} \rho_{S}\left|u_{S}(t)\right|^{2}\right] \\
+ & \mu_{F}\left(2 \int_{0}^{t} \int_{F(t)}\left|D\left(u_{F}(t)\right)\right|^{2}+\frac{1}{\beta_{\Omega}} \int_{0}^{t} \int_{\partial \Omega}\left|u_{F} \times n\right|^{2} \mathrm{~d} \sigma+\frac{1}{\beta_{S}} \int_{0}^{t} \int_{\partial S(t)}\left|\left(u_{F}-u_{S}\right) \times n\right|^{2} \mathrm{~d} \sigma\right) \\
& \leq-\int_{0}^{t} \int_{F(t)} \rho_{F} g e_{3} \cdot u_{F}-\int_{0}^{t} \int_{S(t)} \rho_{S} g e_{3} \cdot u_{S}+\frac{1}{2} \int_{F(0)} \rho_{F}\left|u_{F, 0}\right|^{2}+\frac{1}{2} \int_{S(0)} \rho_{S}\left|u_{S}(0)\right|^{2} . \tag{2.15}
\end{align*}
$$

We refer to [15] for a detailed discussion of this definition. Note that it is restricted to solutions up to collision, through the condition $S(t) \Subset \Omega$. We recall that item 3) is the identity one gets by multiplying (1.1) by a divergence-free test-function $\varphi_{F}$ and then combining with
(1.3)-(1.4). Item 4) is the weak form of the transport equation satisfied by $\chi_{S}$, the indicator function of $S(t)$ as defined in 1). Finally, 2) summarizes the regularity that is implied by the energy estimate (2.15). Note that $u_{F}$ and $u_{S}$ share the same normal component, so that the $L^{\infty} L^{2}$ bound on $u_{F}$ and $u_{S}$ yields a global $L^{\infty} L_{\sigma}^{2}(\Omega)$ on the extended velocity field $u$ (see (1.11) for the definition of $u$ ). Moreover, the $L^{\infty} L^{2}$ bound on $u_{F}$ and $L^{2} L^{2}$ bound on $D\left(u_{F}\right)$ imply an $L_{l o c}^{2} H^{1}$ bound, through Korn inequality (that is valid as long as the solid is away from the boundary of the container).

From [15], there is a weak solution for the system up to a contact between $S(t)$ and $\partial \Omega$ :
Theorem 2.2. Let $\Omega$ and $S_{0} \in \Omega$ two $C^{1,1}$ bounded domain of $\mathbb{R}^{3}$. Given $\left(\beta_{S}, \beta_{\Omega}\right) \in(0, \infty)^{2}$, and initial data $u_{F, 0} \in L_{\sigma}^{2}(\Omega), u_{S, 0} \in \mathcal{R}$ satisfying:

$$
u_{F, 0} \cdot n=u_{S, 0} \cdot n, \quad \text { on } \quad \partial S_{0},
$$

there exists $T \in \mathbf{R}_{+}^{*} \cup\{\infty\}$ and a weak solution of (1.1)-(1.9) on $[0, T)$. Moreover, such weak solution exists up to collision, that is either we can take $T=\infty$, or we can take $T \in \mathbf{R}_{+}^{*}$ in such a way that

$$
S(t) \Subset \Omega, \quad \text { for all } \quad t \in[0, T), \quad \text { and } \quad \lim _{t \rightarrow T^{-}} \operatorname{dist}(\partial S(t), \partial \Omega)=0 .
$$

The aim of Theorem 1.1 is to prove that $T<\infty$ under assumption A1-A4.

### 2.2 The mixed case: $\beta_{S}=0$ and $\beta_{\Omega}>0$.

We proceed with the case where slip is imposed on one of the boundaries only. For simplicity, we consider that no-slip boundary conditions are imposed on $\partial S(t)$ i.e. $\beta_{S}=0$ and $\beta_{\Omega}>0$. To our knowledge, this particular system has not been treated explicitly in previous studies. Nevertheless, as slip is allowed on the exterior boundary of $F(t)$ only, it is straightforward to adapt the method of [10] or [25], for instance, to tackle a Cauchy theory for weak solutions before contact. In particular, the extension trick (1.11) yields a sufficiently smooth velocityfield in this case so that we can reduce the full system to a global weak formulation on this extended velocity field. This reads as follows:

Definition 2.3. Let $\Omega$ and $S_{0} \subset \Omega$ be two Lipschitz bounded domains of $\mathbb{R}^{3}$. Let $u_{F, 0} \in$ $L_{\sigma}^{2}(\Omega), u_{S, 0} \in \mathcal{R}$ such that $u_{F, 0} \cdot n=u_{S, 0} \cdot n$ on $\partial S_{0}$. A weak solution of (1.1)-(1.9) is a pair ( $S, u$ ) satisfying

1) $S(t) \Subset \Omega$ is a bounded domain of $\mathbb{R}^{3}$ for all $t \in[0, T)$ such that

$$
\chi_{S}(t, x) \triangleq 1_{S(t)}(x) \in L^{\infty}((0, T) \times \Omega) .
$$

2) $u \in L^{\infty}\left(0, T ; L_{\sigma}^{2}(\Omega)\right) \cap L^{2}\left(0, T ; H_{\sigma}^{1}(\Omega)\right)$ satisfies

$$
u(t, \cdot)_{\mid S(t)} \in \mathcal{R} \quad \text { for a.a. } t \in(0, T) .
$$

3) For all $\varphi$ belonging to

$$
\mathcal{T}_{T}^{m} \triangleq\left\{\varphi \in \mathcal{D}([0, T) \times \bar{\Omega}), \quad \varphi\left(t,\left.\cdot\right|_{\mid S(t)} \in \mathcal{R} \text { for all } t \in[0, T)\right\}\right.
$$

there holds

$$
\begin{align*}
-\int_{0}^{T} \int_{\Omega} \rho\left(u \partial_{t} \varphi\right. & +u \otimes u: D(\varphi))+\int_{0}^{T} \int_{\Omega} 2 \mu_{F} D(u): D(\varphi) \\
& +\frac{1}{\beta_{\Omega}} \int_{0}^{T} \int_{\partial \Omega}(u \times n) \cdot(\varphi \times n)=-\int_{0}^{T} \int_{\Omega} \rho g e_{3} \varphi-\int_{\Omega} \rho_{0} u_{0} \varphi(0, \cdot) \tag{2.16}
\end{align*}
$$

where

$$
\begin{array}{rlc}
\rho(t, x) & =\rho_{S} 1_{S(t)}(x)+\rho_{F} 1_{F(t)}(x), \quad \text { for }(t, x) \in(0, T) \times \Omega, \\
\rho_{0}(x) & =\rho_{S} 1_{S_{0}}(x)+\rho_{F} 1_{F_{0}}(x), \quad \text { for } x \in \Omega, \\
u_{0}(x) & =u_{S, 0} 1_{S_{0}}(x)+u_{F, 0} 1_{F_{0}}(x), \quad \text { for } x \in \Omega,
\end{array}
$$

4) For all $\psi \in \mathcal{D}([0, T) ; \mathcal{D}(\bar{\Omega}))$, we have

$$
-\int_{0}^{T} \int_{S(t)} \partial_{t} \psi-\int_{0}^{T} \int_{S(t)} u \cdot \nabla \psi=\int_{S_{0}} \psi(0, \cdot)
$$

5) We have the following energy inequality for a.a. $t \in(0, T)$ :

$$
\begin{align*}
\frac{1}{2} \int_{\Omega} \rho|u(t)|^{2}+2 \mu_{F} \int_{0}^{t} \int_{\Omega}|D(u)|^{2}+\frac{1}{\beta_{\Omega}} \int_{0}^{t} & \int_{\partial \Omega}|u \times n|^{2} \mathrm{~d} \sigma \\
& \leq-\int_{0}^{t} \int_{F(t)} \rho g e_{3} u+\frac{1}{2} \int_{\Omega} \rho(0, \cdot)\left|u_{0}\right|^{2} \tag{2.17}
\end{align*}
$$

Following the method of proof of [10], one obtains:
Theorem 2.4. Let $\Omega$ and $S_{0} \Subset \Omega$ two $C^{1,1}$ bounded domain of $\mathbb{R}^{3}$. Given $\beta_{\Omega} \in(0, \infty)$ and initial data $u_{F, 0} \in L_{\sigma}^{2}(\Omega), u_{S, 0} \in \mathcal{R}$ satisfying

$$
u_{F, 0} \cdot n=u_{S, 0} \cdot n, \quad \text { on } \quad \partial S_{0}
$$

the same conclusion as in Theorem 2.2 holds.

## 3 Proof of collision with slip conditions

This section is devoted to the proof of Theorem 1.1. We assume $\beta_{S}$ and $\beta_{\Omega}$ are fixed nonnegative slip lengths, and $\rho_{S}>\rho_{F}$. Let $(S, u)$ be a weak solution over $(0, T)$, as given by Theorem 2.2. We consider a solution up to its "maximal time of existence": that is $T=\infty$ if there is no collision, or $T<\infty$ the time of collision. We assume that this weak solution satisfies A1-A4 for $t \in[0, T)$. We denote by $h=h(t)$ the distance between the south pole of $S(t)$ and $x=0$. By assumption A4, when $h(t)$ is small enough, it is equal to the distance between $S(t)$ and the boundary $\partial \Omega$. Our goal is to prove that $h(t)$ goes to zero in finite time. In other words, $T<\infty$, that is collision occurs in finite time.

To build up our proof, we first need some general considerations on the connection between the weak formulation (2.14) and the distance $h$, for an arbitrary weak solution satisfying A4. We then design a special test-function in order to prove contact in finite time.

### 3.1 Some features of the weak formulation (2.14)

Let $(S, u)$ a weak solution defined on $(0, T)$ and $\varphi \in \mathcal{T}_{T}^{s}$ of the form

$$
\begin{equation*}
\varphi=\zeta(t) \varphi_{h(t)} \text { with } \quad \varphi_{S}=\zeta(t) e_{3} \tag{3.18}
\end{equation*}
$$

for some functions $\zeta \in \mathcal{D}([0, T))$ and $h \mapsto \varphi_{h}$ to be made precise. Integrating by parts, we get that for any $q \in \mathcal{D}([0, T) \times \bar{\Omega})$, for a.a. $t \in(0, T)$

$$
\begin{aligned}
& 2 \int_{F(t)} D\left(\varphi_{F}\right): D\left(u_{F}\right) \\
= & \int_{F(t)}\left(2 D\left(\varphi_{F}\right)-q I\right): \nabla u_{F} \\
= & -\int_{F(t)}\left(\Delta \varphi_{F}-\nabla q\right) \cdot u_{F}+\int_{\partial F(t)}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot u_{F} d \sigma \\
= & -\int_{F(t)}\left(\Delta \varphi_{F}-\nabla q\right) \cdot u_{F}+\int_{\partial \Omega}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot u_{F} d \sigma+\int_{\partial S(t)}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot u_{F} d \sigma \\
= & -\int_{F(t)}\left(\Delta \varphi_{F}-\nabla q\right) \cdot u_{F}+\int_{\partial \Omega}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot u_{F} d \sigma \\
& +\int_{\partial S(t)}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot\left(u_{F}-u_{S}\right) d \sigma+\int_{\partial S(t)}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot u_{S} d \sigma,
\end{aligned}
$$

Chosing also $q$ of the form $q=\zeta(t) q_{h(t)}$, the last term can be calculated by
$\int_{\partial S(t)}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot u_{S} d \sigma=h^{\prime}(t) \zeta(t) \int_{\partial S(t)}\left(2 D\left(\varphi_{h(t)}\right)-q_{h(t)} I\right) n \cdot e_{3} d \sigma \triangleq \zeta(t) h^{\prime}(t) n(h(t))$.
We emphasize that the domain $S(t)$ is completely fixed by the value of $h(t)$ (we shall denote $S_{h(t)}$ in what follows) so that the last integral does depend on $h(t)$ only. Combining all the equations with (2.14), we find

$$
\begin{align*}
& \mu_{F} \int_{0}^{T} \zeta(t) h^{\prime}(t) n(h(t)) \mathrm{d} t \\
= & \int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \cdot \partial_{t} \varphi_{F}+\int_{0}^{T} \int_{S(t)} \rho_{S} u_{S} \cdot \partial_{t} \varphi_{S}+\int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \otimes u_{F}: \nabla \varphi_{F} \\
- & \int_{0}^{T} \int_{F(t)} \rho_{F} g e_{3} \varphi_{F}-\int_{0}^{T} \int_{S(t)} \rho_{S} g e_{3} \varphi_{S}+\int_{F_{0}} \rho_{F} u_{F, 0} \varphi_{F}(0, \cdot)+\int_{S_{0}} \rho_{S} u_{S, 0} \varphi_{S}(0, \cdot) \\
- & \frac{\mu_{F}}{\beta_{\Omega}} \int_{0}^{T} \int_{P}\left(u_{F} \times n\right) \cdot\left(\varphi_{F} \times n\right)-\frac{\mu_{F}}{\beta_{S}} \int_{0}^{T} \int_{\partial S(t)}\left(\left(u_{F}-u_{S}\right) \times n\right) \cdot\left(\left(\varphi_{F}-\varphi_{S}\right) \times n\right) \\
+ & \mu_{F}\left(\int_{0}^{T} \int_{F(t)}\left(\Delta \varphi_{F}-\nabla q\right) \cdot u_{F}-\int_{0}^{T} \int_{\partial \Omega}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot u_{F} d \sigma\right. \\
- & \left.\int_{0}^{T} \int_{\partial S(t)}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot\left(u_{F}-u_{S}\right) d \sigma\right) \tag{3.1}
\end{align*}
$$

This identity enables to describe the time evolution of $h$ through the function $n$ that we have to choose now. To compute $n(h)$, we shall take advantage of the following identity:

$$
\begin{align*}
n(h) & =\int_{\partial S_{h}}\left(2 D\left(\varphi_{h}\right)-q_{h} I\right) n \cdot \varphi_{h} d \sigma+\int_{\partial S_{h}}\left(D\left(\varphi_{h}\right)-q_{h} I\right) n \cdot\left(e_{3}-\varphi_{h}\right) d \sigma \\
& =\int_{F_{h}}\left(\Delta \varphi_{h}-\nabla q_{h}\right) \cdot \varphi_{h} d x+2 \int_{F_{h}} D\left(\varphi_{h}\right): D\left(\varphi_{h}\right) d x \\
& -\int_{\partial \Omega}\left(2 D\left(\varphi_{h}\right)-q_{h} I\right) n \cdot \varphi_{h} d \sigma+\int_{\partial S_{h}}\left(D\left(\varphi_{h}\right)-q_{h} I\right) n \cdot\left(e_{3}-\varphi_{h}\right) d \sigma . \tag{3.20}
\end{align*}
$$

where $F_{h} \triangleq \Omega \backslash \bar{S}_{h}$. This identity holds true for all $h$ such that $S_{h} \subset \subset \Omega$.

### 3.2 Construction of test-functions

We construct now the test function $\varphi_{h}$ and the pressure $q_{h}$ associated with the solid particle $S_{h}$ frozen at distance $h$. It will be defined for $h \in\left(0, h_{M}\right)$, with $h_{M} \triangleq \sup _{t \in[0, T)} h(t)$. When $h \rightarrow 0$, a cusp arises in $F_{h}=\Omega \backslash \overline{S_{h}}$. This cusp is contained in the domain :

$$
\Omega_{h, r}:=\left\{\left(x_{1}, x_{2}, x_{3}\right) \in F_{h}, x_{3}<\frac{1}{2}+h \text { and }\left|\left(x_{1}, x_{2}\right)\right|<r\right\}
$$

for arbitrary $r \in\left(0, \frac{1}{2}\right)$. In particular, in all forthcoming estimates, we shall pay special attention to the region $\Omega_{h, 2 \delta}$, where $\delta$ was introduced in assumption A2.

To compute test-functions and associated pressures, we introduce cylindrical coordinates $(r, \theta, z)$ associated with cartesian coordinates $\left(x_{1}, x_{2}, x_{3}\right)$ in such a way that $x_{3}=z$. The local basis reads $\left(e_{r}, e_{\theta}, e_{z}\right)$. First, we introduce an approximate solution of the steady Stokes equation in the domain $\Omega_{h, 2 \delta}$. It will satisfy locally the slip boundary conditions (1.6)-(1.7), with $u_{S}$ replaced by $e_{z}$ :

$$
\widetilde{\varphi}_{h}=\frac{1}{2}\left(-\partial_{z} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right) x_{1},-\partial_{z} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right) x_{2}, \frac{1}{r} \partial_{r}\left(\Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right) r^{2}\right)\right)
$$

where

$$
\begin{equation*}
\gamma_{s}(r)=1-\sqrt{1-r^{2}}, \quad \Phi(r, t)=P_{1}(r) t+P_{2}(r) t^{2}+P_{3} t^{3} . \tag{3.21}
\end{equation*}
$$

Here

$$
\begin{align*}
P_{1}(r) & =\frac{6\left(2+\alpha_{S}\right)}{12+4\left(\alpha_{S}+\alpha_{P}\right)+\alpha_{S} \alpha_{P}}  \tag{3.22}\\
P_{2}(r) & =\frac{3\left(2+\alpha_{S}\right) \alpha_{P}}{12+4\left(\alpha_{S}+\alpha_{P}\right)+\alpha_{S} \alpha_{P}}  \tag{3.23}\\
P_{3}(r) & =-\frac{2\left(\alpha_{S}+\alpha_{S} \alpha_{P}+\alpha_{P}\right)}{12+4\left(\alpha_{S}+\alpha_{P}\right)+\alpha_{S} \alpha_{P}} \tag{3.24}
\end{align*}
$$

where

$$
\begin{equation*}
\alpha_{P}=\frac{h+\gamma_{s}(r)}{\beta_{\Omega}}, \quad \alpha_{S}=\left(\frac{1}{\beta_{S}}+2\right)\left(h+\gamma_{s}(r)\right) . \tag{3.25}
\end{equation*}
$$

This approximation is inspired by the computations in [14]. We emphasize that the value of $\alpha_{S}$ does not match exactly the one in [14]. This modification is required so that further boundary estimates hold true (see (3.39)). We stress that notation $\partial_{z}^{k} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right.$ ), resp. $\partial_{r}^{k} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)$, stands for $\partial_{z}^{k} \Psi(r, z)$, resp. $\partial_{r}^{k} \Psi(r, z), \quad \Psi(r, z) \triangleq \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)$. With this structure, $\tilde{\varphi}_{h}$ is divergence-free. Moreover, the function $\Phi$ is constructed so that it satisfies the following boundary conditions:

- for $z=0$

$$
\begin{align*}
& \Phi(r, 0)=0,  \tag{3.26}\\
& \partial_{z z} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)-\frac{1}{\beta_{\Omega}} \partial_{z} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)=0 \tag{3.27}
\end{align*}
$$

- for $z=h+\gamma_{S}(r)$

$$
\begin{align*}
& \Phi(r, 1)=1  \tag{3.28}\\
& \partial_{z z} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)+\left(\frac{1}{\beta_{S}}+2\right) \partial_{z} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)=0 \tag{3.29}
\end{align*}
$$

In the local cylindrical basis, the approximate solution $\widetilde{\varphi}_{h}$ reads:

$$
\begin{equation*}
\widetilde{\varphi}_{h}=\frac{1}{2}\left[-\partial_{z} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right) r e_{r}+\frac{1}{r} \partial_{r}\left(\Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right) r^{2}\right) e_{z}\right] . \tag{3.30}
\end{equation*}
$$

In particular, if we remark that $n$ satisfies $n=-r e_{r}+(1+h-z) e_{z}$ on $\partial S_{h} \cap \Omega_{h, 1 / 2}$ and $n=e_{3}$ on $\partial \Omega \cap \Omega_{h, 1 / 2}$, we obtain (see the introduction of Appendix A for further details) :

$$
\left\{\begin{align*}
\left.\left(\widetilde{\varphi}_{h}-e_{3}\right) \cdot n\right|_{\partial S_{h} \cap \partial \Omega_{h, 2 \delta}} & =0,  \tag{3.31}\\
\left.\widetilde{\varphi}_{h} \cdot n\right|_{\partial \Omega \cap \partial \Omega_{h, 2 \delta}} & =0, \\
\widetilde{\varphi}_{h} \times\left. n\right|_{\partial \Omega \cap \partial \Omega_{h, 2 \delta}} & =-2 \beta_{\Omega}\left(D\left(\widetilde{\varphi}_{h}\right) n\right) \times\left. n\right|_{\partial \Omega \cap \partial \Omega_{h, 2 \delta}}
\end{align*}\right.
$$

Thus, up to its extension to the whole of $\Omega$, the function $\widetilde{\varphi}_{h}$ is a good candidate for applying the computations of the previous section. Such extension $\varphi_{h}$ should satisfy $\left.\varphi_{h}\right|_{S_{h}}=e_{z}$. It remains to define it into the whole of $F_{h}$.

To this end, we introduce a smooth function $\phi=\phi(x)$ with compact support satisfying

$$
\begin{array}{ll}
\phi=1 & \text { on a } \frac{d_{\delta}}{2} \text {-neighborhood of the unit ball } B(0,1), \\
\phi=0 & \text { outside a } d_{\delta} \text {-neighborhood of } B(0,1) .
\end{array}
$$

We recall that $d_{\delta}$ was introduced in A4. We also define a cut-off function $\chi: \mathbb{R}^{3} \rightarrow[0,1]$ s.t.

$$
\begin{equation*}
\chi=1 \text { in }(-\delta, \delta)^{3}, \quad \chi=0 \text { outside }(-2 \delta, 2 \delta)^{3} . \tag{3.32}
\end{equation*}
$$

With these cut-off functions, we set (in cartesian coordinates):

$$
\begin{align*}
& \varphi_{h}=\frac{1}{2}\left(\begin{array}{c}
-x_{1} \partial_{z}\left[(1-\chi(x)) \phi\left(x-(1+h) e_{3}\right)+\chi(x) \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)\right] \\
-x_{2} \partial_{z}\left[(1-\chi(x)) \phi\left(x-(1+h) e_{3}\right)+\chi(x) \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)\right] \\
\frac{1}{r} \partial_{r}\left[r^{2}\left((1-\chi(x)) \phi\left(x-(1+h) e_{3}\right)+\chi(x) \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)\right)\right]
\end{array}\right) \text { in } F_{h}, \\
& \varphi_{h}=e_{3} \text { in } S_{h} . \tag{3.33}
\end{align*}
$$

By the definition of $\varphi_{h}$, we get

Lemma 3.1. The map $(h, x) \mapsto \varphi_{h}(x)$ is smooth on all domains:

$$
\left\{(h, x), h \in\left(0, h_{M}\right), x \in F_{h}\right\},\left\{(h, x) \in\left(0, h_{M}\right), x \in S_{h}\right\},\left\{(h, x), h \in\left[0, h_{M}\right), x \in F_{h} \backslash \Omega_{h, \delta}\right\} .
$$

Furthermore, there holds:

$$
\begin{equation*}
\varphi_{h} \cdot n=0 \quad \text { on } \partial \Omega \tag{3.34}
\end{equation*}
$$

and

$$
\begin{equation*}
\left.\varphi_{h}\right|_{F_{h}} \cdot n=e_{3} \cdot n \quad \text { on } \partial S_{h} . \tag{3.35}
\end{equation*}
$$

This comes from direct calculations, that we leave to the reader for brevity. We now state some refined estimates, to be used in the proof of Theorem 1.1. The whole point is to have an accurate bound on the blow-up of the test function as $h$ goes to zero.

Lemma 3.2. There exist constants $0<c<C$ so that, when $h<h_{M}$ :

$$
\begin{align*}
\left\|\varphi_{h}\right\|_{L^{2}\left(F_{h}\right)} & \leq C,  \tag{3.36}\\
\left\|\nabla \varphi_{h}\right\|_{L^{2}\left(F_{h}\right)}^{2} & \leq C|\ln h|,  \tag{3.3}\\
\left\|D\left(\varphi_{h}\right)\right\|_{L^{2}\left(F_{h}\right)}^{2} & \geq c|\ln (h)|, \tag{3.38}
\end{align*}
$$

and

$$
\begin{equation*}
\left\|2 \beta_{S}\left(D\left(\varphi_{h}\right) n\right) \times n+\left(\varphi_{h}-e_{3}\right) \times n\right\|_{L^{2}\left(\partial S_{h}\right)} \leq C, \tag{3.39}
\end{equation*}
$$

Furthermore, for $i=1,2,3$, denoting by exponent $i$ cartesian coordinates, there holds:

$$
\begin{align*}
\left\|\int_{0}^{h+\gamma_{s}(r)} \partial_{h} \varphi_{h}^{i} d z\right\|_{L^{2}\left(\partial \Omega \cap \partial \Omega_{h, \delta}\right)}^{2} & \leq C|\ln h|,  \tag{3.40}\\
\left\|\int_{z}^{h+\gamma_{s}(r)} \partial_{h} \varphi_{h}^{i} d s\right\|_{L^{2}\left(\Omega_{h, \delta}\right)} & \leq C,  \tag{3.41}\\
\left\|\left(h+\gamma_{s}(r)\right) \partial_{i} \int_{z}^{h+\gamma_{s}(r)} \partial_{h} \varphi_{h}^{i}\left(x_{1}, x_{2}, s\right) d s\right\|_{L^{2}\left(\Omega_{h, \delta}\right)} & \leq C, \tag{3.42}
\end{align*}
$$

and for $(i, j) \in\{1,2,3\}^{2}$ and $x \in \Omega_{h, \delta}$ :

$$
\begin{align*}
\left|\int_{z}^{h+\gamma_{S}(r)}\left(\partial_{i} \varphi_{h}^{j}\left(x_{1}, x_{2}, s\right)+\partial_{j} \varphi_{h}^{i}\left(x_{1}, x_{2}, s\right)\right) d s\right| & \leq C,  \tag{3.43}\\
\left|\left(h+\gamma_{S}(r)\right) \partial_{i} \int_{z}^{h+\gamma_{S}(r)}\left(\partial_{i} \varphi_{h}^{j}\left(x_{1}, x_{2}, s\right)+\partial_{j} \varphi_{h}^{i}\left(x_{1}, x_{2}, s\right)\right) d s\right| & \leq C . \tag{3.44}
\end{align*}
$$

The proof of this lemma is postponed to Appendix A.
When computing (3.19), we introduced a pressure $q$. The aim was to smoothen singularities in $\Delta \varphi_{h}$ which must arise when $h \rightarrow 0$. The construction of this pressure is the content of the next lemma:

Lemma 3.3. There exists a map $(h, x) \mapsto q_{h}$ which is smooth on the domains

$$
\left\{(h, x), h \in\left(0, h_{M}\right), x \in F_{h}\right\} \text { and }\left\{(h, x), h \in\left[0, h_{M}\right), x \in F_{h} \backslash \Omega_{h, \delta}\right\}
$$

and which satisfies :

- there exist constants $0<c<C<\infty$ such that :

$$
c|\ln h| \leq \int_{\partial S_{h}}\left(2 D\left(\varphi_{h}\right)-q_{h} I\right) n \cdot\left(e_{3}-\varphi_{h}\right) d \sigma-\int_{\partial \Omega}\left(2 D\left(\varphi_{h}\right)-q_{h} I\right) n \cdot \varphi_{h} d \sigma \leq C|\ln h|
$$

- there exists a positive constant $C$ such that, for any $h \in\left(0, h_{M}\right)$ and $v \in H^{1}\left(F_{h}\right)$ satisfying $v \cdot n=0$ on $\partial \Omega$, we have:

$$
\begin{equation*}
\left|\int_{F_{h}}\left(\Delta \varphi_{h}-\nabla q_{h}\right) v\right| \leq C\left(\|D(v)\|_{L^{2}\left(F_{h}\right)}+\|v\|_{L^{2}(\partial \Omega)}\right) \tag{3.45}
\end{equation*}
$$

The proof of this lemma is also postponed to Appendix A.

### 3.3 Conclusion of the proof of Theorem 1.1

We are now in a position to prove Theorem 1.1. Combining the energy inequality (2.15) with the classical identity

$$
-\int_{0}^{t} \int_{F(t)} \rho_{F} g e_{z} \cdot u_{F}-\int_{0}^{t} \int_{S(t)} \rho_{S} g e_{z} \cdot u_{S}=\frac{4\left(\rho_{F}-\rho_{S}\right) g \pi}{3}(h(t)-h(0))
$$

we obtain easily (remember that $\rho_{S}>\rho_{F}$ ):

$$
\begin{align*}
\sup _{(0, T)} \int_{\Omega}|u(t)|^{2}+\int_{0}^{T} & \int_{F(t)}\left|D\left(u_{F}(t)\right)\right|^{2} \\
& +\int_{0}^{T} \int_{\partial \Omega}\left|u_{F} \times n\right|^{2} \mathrm{~d} \sigma+\int_{0}^{T} \int_{\partial S(t)}\left|\left(u_{F}-u_{S}\right) \times n\right|^{2} \mathrm{~d} \sigma \leq C_{0} \tag{3.46}
\end{align*}
$$

This energy bound will be of constant use.
As mentioned before, we take a special test function in the variational formulation (2.14). Namely, we take $\varphi(t, x)=\zeta(t) \varphi_{h(t)}(x)$, with the field $\varphi_{h}$ built in paragraph 3.2. From Lemma 3.1, we have that $\varphi \in \mathcal{T}_{T}$ so that identity (3.19) holds true. The whole point is then to bound properly the r.h.s. of (3.19):

$$
\begin{aligned}
R H S & \triangleq \int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \cdot \partial_{t} \varphi_{F}+\int_{0}^{T} \int_{S(t)} \rho_{S} u_{S} \cdot \partial_{t} \varphi_{S}+\int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \otimes u_{F}: \nabla \varphi_{F} \\
& -\int_{0}^{T}\left(\int_{F(t)} \rho_{F} g e_{z} \varphi_{F}+\int_{S(t)} \rho_{S} g e_{z} \varphi_{S}\right)+\left(\int_{F_{0}} \rho_{F} u_{F, 0} \varphi_{F}(0, \cdot)+\int_{S_{0}} \rho_{S} u_{S, 0} \varphi_{S}(0, \cdot)\right) \\
& -\mu_{F} \int_{0}^{T} \int_{\partial \Omega}\left(u_{F} \times n\right) \cdot\left(\frac{1}{\beta_{\Omega}}\left(\varphi_{F} \times n\right)+\left(2 D\left(\varphi_{F}\right)-q I\right) n\right) d \sigma \\
& \left.-\mu_{F} \int_{0}^{T} \int_{\partial S(t)}\left(\left(u_{F}-u_{S}\right) \times n\right) \cdot\left(\frac{1}{\beta_{S}}\left(\varphi_{F}-\varphi_{S}\right) \times n\right)+\left(2 D\left(\varphi_{F}\right)-q I\right) n\right) d \sigma \\
& +\mu_{F} \int_{0}^{T} \int_{F(t)}\left(\Delta \varphi_{F}-\nabla q\right) \cdot u_{F} \\
& =I_{1}+I_{2}+I_{3}+I_{4}+I_{5}-\mu_{F}\left(I_{6}+I_{7}+I_{8}\right)
\end{aligned}
$$

where we have applied the continuity of normal traces on $\partial S(t)$ and $\partial \Omega$ in the boundary terms.

- We first deal with

$$
I_{3}=\int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \otimes u_{F}: \nabla \varphi_{F} \triangleq \int_{0}^{T} \rho_{F} \zeta(t) I_{3}(t)
$$

We introduce the functions $x \mapsto P_{z} \varphi_{h}^{i, j}(x)$ for $(i, j) \in\{1,2,3\}^{2}$ defined by:

$$
\begin{aligned}
P_{z} \varphi_{h}^{i, j}(x) & =\int_{z}^{h+\gamma_{S}(r)}\left(\partial_{i} \varphi_{h}^{j}+\partial_{j} \varphi_{h}^{i}\right)\left(x_{1}, x_{2}, s\right) d s, \quad \forall r<1, \quad z<h+\gamma_{s}(r) \\
P_{z} \varphi_{h}^{i, j}(x) & =\int_{z}^{h+1}\left(\partial_{i} \varphi_{h}^{j}+\partial_{j} \varphi_{h}^{i}\right)\left(x_{1}, x_{2}, s\right) d s \quad \text { otherwise. }
\end{aligned}
$$

Let us stress that $P_{z} \varphi_{h}^{i, j}$ is smooth in $F_{h}$, notably across $r=1$ : it follows from the cancellation of the symmetric gradient $\partial_{i} \varphi_{h}^{j}+\partial_{j} \varphi_{h}^{i}$ near $(r=1, z=1+h)$ that both expressions coincide there.

Moreover, by (3.43), there exists a constant $C$ independent of $t$ such that, for all $(i, j) \in$ $\{1,2,3\}^{2}$, there holds:

$$
\begin{array}{r}
\left|P_{z} \varphi_{h}^{i, j}\left(x_{1}, x_{2}, x_{3}\right)\right|+\left|\left(h+\gamma_{S}(r)\right) \partial_{i} P_{z} \varphi_{h}^{i, j}\left(x_{1}, x_{2}, x_{3}\right)\right| \leq C, \quad \forall x \in \Omega_{h(t), \delta} \\
\left|P_{z} \varphi_{h}^{i, j}\left(x_{1}, x_{2}, x_{3}\right)\right|+\left|\partial_{i} P_{z} \varphi_{h}^{i, j}\left(x_{1}, x_{2}, x_{3}\right)\right| \leq C, \quad \forall x \in F(t) \backslash \Omega_{h(t), \delta} . \tag{3.48}
\end{array}
$$

Then, by performing integration by parts, we obtain (summation on indices $i$ and $j$ is implicit):

$$
\begin{aligned}
I_{3}(t)= & -\int_{F(t)} u_{F}^{i} u_{F}^{j} \partial_{z}\left[P_{z} \varphi_{h}^{i, j}\right]=\int_{F(t)} \partial_{3}\left(u_{F}^{i} u_{F}^{j}\right) P_{z} \varphi_{h}^{i, j}-\int_{\partial \Omega} u_{F}^{i} u_{F}^{j} P_{z} \varphi_{h}^{i, j} n^{3} \mathrm{~d} \sigma \\
= & \int_{F(t)}\left(\left(\partial_{3} u_{F}^{i}+\partial_{i} u_{F}^{3}\right) u_{F}^{j}+\left(\partial_{3} u_{F}^{j}+\partial_{j} u_{F}^{3}\right) u_{F}^{i}\right) P_{z} \varphi_{h}^{i, j} \\
& -\int_{F(t)}\left(\partial_{i} u_{F}^{3} u_{F}^{j}+\partial_{j} u_{F}^{3} u_{F}^{i}\right) P_{z} \varphi_{h}^{i, j}-\int_{\partial \Omega}\left(u_{F}^{i} u_{F}^{j}\right) P_{z} \varphi_{h}^{i, j} n^{3} \mathrm{~d} \sigma \\
= & \int_{F(t)}\left(\left(\partial_{3} u_{F}^{i}+\partial_{i} u_{F}^{3}\right) u_{F}^{j}+\left(\partial_{3} u_{F}^{j}+\partial_{j} u_{F}^{3}\right) u_{F}^{i}\right) P_{z} \varphi_{h}^{i, j} \\
& +\int_{F(t)}\left(\left(\partial_{i} u_{F}^{j}+\partial_{j} u_{F}^{i}\right) u_{F}^{3} P_{z} \varphi_{h}^{i, j}+u_{F}^{3} u_{F}^{j} \partial_{i} P_{z} \varphi_{h}^{i, j}+u_{F}^{3} u_{F}^{i} \partial_{j} P_{z} \varphi_{h}^{i, j}\right) \\
& -\int_{\partial \Omega}\left(u_{F}^{i} u_{F}^{j} n^{3}+u_{F}^{j} u_{F}^{3} n^{i}+u_{F}^{i} u_{F}^{3} n^{j}\right) P_{z} \varphi_{h}^{i, j} \mathrm{~d} \sigma .
\end{aligned}
$$

In this last identity, we apply (3.47)-(3.48) to obtain :

$$
\left|\int_{\partial \Omega}\left(u_{F}^{i} u_{F}^{j} n^{3}+u_{F}^{j} u_{F}^{3} n^{i}+u_{F}^{i} u_{F}^{3} n^{j}\right) P_{z} \varphi_{h}^{i, j} \mathrm{~d} \sigma\right| \leq C\left\|u_{F}(t)\right\|_{L^{2}(\partial \Omega)}^{2}
$$

Concerning the other terms, we have for instance :

$$
\begin{aligned}
& \left|\int_{F(t)} u_{F}^{3} u_{F}^{j} \partial_{i} P_{z} \varphi_{h}^{i, j}\right| \leq\left|\int_{F(t) \cap \Omega_{h(t), \delta}} \frac{u_{F}^{3}}{h+\gamma_{S}(r)} u_{F}^{j}\left(h+\gamma_{S}(r)\right) \partial_{i} P_{z} \varphi_{h}^{i, j}\right| \\
& +\left|\int_{F(t) \backslash \Omega_{h(t), \delta}} u_{F}^{3} u_{F}^{j} \partial_{i} P_{z} \varphi_{h}^{i, j}\right| \\
& \leq\left\|\frac{u_{F}^{3}}{h+\gamma_{S}(r)}\right\|_{L^{2}\left(F(t) \cap \Omega_{h(t), \delta}\right)}\left\|u_{F}\right\|_{L^{2}(F(t))}\left\|\left(h+\gamma_{S}(r)\right) \partial_{i} P_{z} \varphi_{h}^{i, j}\right\|_{\left.L^{\infty}\left(F(t) \cap \Omega_{h(t), \delta}\right)\right)} \\
& \quad+\left\|u_{F}^{3}\right\|_{L^{2}\left(F(t) \cap\left\{\left|\left(x_{1}, x_{2}\right)\right|<\left(1+d_{\delta}\right)\right\}\right)}\left\|u_{F}\right\|_{L^{2}(F(t))}\left\|\partial_{i} P_{z} \varphi_{h}^{i, j}\right\|_{L^{\infty}\left(F(t) \backslash \Omega_{h(t), \delta \delta}\right)} .
\end{aligned}
$$

As $u_{F}^{3}(x)=0$ for $x \in\left\{\left|\left(x_{1}, x_{2}\right)\right|<1+d_{\delta}, x_{3}=0\right\}$, a straightforward Poincaré inequality yields that there exists a constant $C>0$ independent of time so that:

$$
\begin{aligned}
\left\|\frac{u_{F}^{3}}{h+\gamma_{S}(r)}\right\|_{L^{2}\left(F(t) \cap \Omega_{h(t), \delta}\right)} & \leq C\left\|\partial_{3} u_{F}^{3}\right\|_{L^{2}(F(t))} \\
\left\|u_{F}^{3}\right\|_{L^{2}\left(F(t) \cap\left\{\left|\left(x_{1}, x_{2}\right)\right|<\left(1+d_{\delta}\right)\right\}\right)} & \leq C\left(\left\|\partial_{3} u_{F}^{3}\right\|_{L^{2}(F(t))}+\left\|u_{F} \times n\right\|_{L^{2}(\partial \Omega)}\right) .
\end{aligned}
$$

Introducing (3.47)-(3.48) then yields :

$$
\left|I_{3}(t)\right| \leq C\left[\left\|D\left(u_{F}\right)\right\|_{L^{2}(F(t))}^{2}+\left\|D\left(u_{F}(t)\right)\right\|_{L^{2}(F(t))}\left\|u_{F}(t)\right\|_{L^{2}(F(t))}+\left\|u_{F}(t)\right\|_{L^{2}(\partial \Omega)}^{2}\right]
$$

so that (3.46) implies there exists a constant $C$ depending only on initial data for which

$$
\left|I_{3}\right| \leq C(1+\sqrt{T})\|\zeta\|_{L^{\infty}(0, T)} .
$$

- To compute $I_{1}$, we split:
$I_{1}=\int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \cdot \partial_{t} \varphi_{F}=\int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \zeta^{\prime}(t) \varphi_{h(t)}+\int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \zeta(t) h^{\prime}(t) \partial_{h} \varphi_{F} \triangleq I_{1}^{a}+I_{1}^{b}$.
Applying (3.36) yields :

$$
\left|I_{1}^{a}\right| \leq C\left\|\zeta^{\prime}\right\|_{L^{1}(0, T)}\|u\|_{L^{\infty}\left(0, T ; L^{2}(\Omega)\right)} \leq C_{0}\left\|\zeta^{\prime}\right\|_{L^{1}(0, T)} .
$$

Then, we have $I_{1}^{b}=\rho_{F} \int_{0}^{t} \zeta(t) h^{\prime}(t) I_{1}^{b}(t)$ where, for all $t \in(0, T)$ :

$$
I_{1}^{b}(t)=\sum_{i=1}^{3} \int_{F(t)} u_{F}^{i} \cdot \partial_{h} \varphi_{h(t)}^{i} .
$$

To bound this term, we define as previously $P_{z} \varphi_{h}^{i}$ which satisfies, for all $i=1,2,3$ :

$$
\begin{aligned}
& P_{z} \varphi_{h}^{i}(x)=\int_{z}^{h+\gamma_{S}(r)} \partial_{h} \varphi_{h}^{i}\left(x_{1}, x_{2}, s\right) d s, \quad \forall r<1, \quad z<h+\gamma_{s}(r), \\
& P_{z} \varphi_{h}^{i}(x)=\int_{z}^{h+1} \partial_{h} \varphi_{h}^{i}\left(x_{1}, x_{2}, s\right) d s \quad \text { otherwise. }
\end{aligned}
$$

Again, $P_{z} \varphi_{h}^{i}$ is continuous in $F_{h}$, notably near $r=1$, as $\partial_{h} \varphi_{h}$ vanishes near $(r=1, z=1+h)$.

With (3.41) together with (3.42), there is a constant $C$ such that:

$$
\begin{gather*}
\left\|P_{z} \varphi_{h}^{i}\right\|_{L^{2}(F(t))}+\left\|\left(h+\gamma_{S}(r)\right) \partial_{i} P_{z} \varphi_{h}^{i}\right\|_{L^{2}\left(F(t) \cap \Omega_{h(t), 2 \delta}\right)}+\left\|\partial_{i} P_{z} \varphi_{h}^{i}\right\|_{L^{2}\left(F(t) \backslash \Omega_{h(t), 2 \delta)}\right.} \leq C,  \tag{3.49}\\
\left\|P_{z} \varphi_{h}^{i}\right\|_{L^{2}(\partial \Omega)} \leq C|\ln (h)|^{1 / 2} . \tag{3.50}
\end{gather*}
$$

From now on, we restrict implicitly the integral $I_{1}^{b}(t)$ to $\left|\left(x_{1}, x_{2}\right)\right|<\left(1+d_{\delta}\right)$ as $\partial_{h} \varphi_{h}$ vanishes elsewhere. Performing integration by parts (summation over index $i$ is implicit), we obtain :

$$
\begin{aligned}
I_{1}^{b}(t) & =-\int_{F(t)} u_{F}^{i} \partial_{z} P_{z} \varphi_{h}^{i} \\
& =\int_{F(t)} \partial_{3} u_{F}^{i} P_{z} \varphi_{h}^{i}-\int_{\partial \Omega} u_{F}^{i} P_{z} \varphi_{h}^{i} n^{3} \mathrm{~d} \sigma \\
& \leq \int_{F(t)}\left(\partial_{3} u_{F}^{i}+\partial_{i} u_{F}^{3}\right) P_{z} \varphi_{h}^{i}+\int_{F(t)} u_{F}^{3} \partial_{i} P_{z} \varphi_{h}^{i}-\int_{\partial \Omega}\left(u_{F}^{3} n^{i}+u_{F}^{i} n^{3}\right) P_{z} \varphi_{h}^{i}
\end{aligned}
$$

Relying on (3.49)-(3.50) and (3.46), we might then reproduce the arguments we used to compute $I_{3}$, resulting in:

$$
\left|I_{1}^{b}\right| \leq C\|\zeta\|_{L^{\infty}}\left(\int_{0}^{T}\left|h^{\prime}(t)\left\|\left.\ln h\right|^{\frac{1}{2}}\right\| u_{F} \times n \|_{L^{2}(\partial \Omega)}+\sqrt{T}\right) .\right.
$$

At this point, we remark that, given the computations of [14, Section 3.2], $|\ln (h(t))|^{1 / 2}$ is the minimum dissipation energy of functions $v \in H^{1}(F(t))$ satisfying $v \cdot n=e_{3} \cdot n$ on $\partial S(t)$ and $v \cdot n=0$ on $\partial \Omega$. Consequently, there exists a constant $C$ independent of time such that:

$$
|\dot{h}||\ln (h)|^{\frac{1}{2}} \leq C\left(\int_{F(t)}\left|D\left(u_{F}(t)\right)\right|^{2}+\int_{\partial \Omega}\left|u_{F} \times n\right|^{2} \mathrm{~d} \sigma+\int_{\partial S(t)}\left|\left(u_{F}-u_{S}\right) \times n\right|^{2} \mathrm{~d} \sigma\right)^{\frac{1}{2}}
$$

and (3.46) implies:

$$
\int_{0}^{T}\left|h^{\prime}(t)\left\|\left.\ln h\right|^{\frac{1}{2}}\right\| u_{F} \times n \|_{L^{2}(\partial \Omega)} \leq C .\right.
$$

- Next, we recall that $\int_{F(t)} \rho_{F} g e_{3} \varphi_{h(t)}=-\frac{4 \pi \rho_{F} g}{3}$, which implies:

$$
\begin{equation*}
I_{4}=-\int_{0}^{T} \int_{F(t)} \rho_{F} g e_{3} \varphi_{F}-\int_{0}^{T} \int_{S(t)} \rho_{S} g e_{3} \varphi_{S}=\frac{4 \pi\left(\rho_{F}-\rho_{S}\right) g}{3} \int_{0}^{T} \zeta(t) \mathrm{d} t . \tag{3.51}
\end{equation*}
$$

- Then, we apply (3.31) together with (3.46). It leads to

$$
\begin{aligned}
\left|I_{6}\right| & =\left|\int_{0}^{T} \zeta(t) \int_{\partial \Omega}\left(u_{F} \times n\right) \cdot\left(\frac{1}{\beta_{\Omega}}\left(\varphi_{h} \times n\right)+\left(2 D\left(\varphi_{h}\right)-q I\right) n\right) d \sigma\right| \\
& =\left|\int_{0}^{T} \zeta(t) \int_{\partial \Omega}\left(u_{F} \times n\right) \cdot\left(\frac{1}{\beta_{\Omega}}\left(\varphi_{h} \times n\right)+2 D\left(\varphi_{h}\right) n\right) d \sigma\right| \\
& \leq C \int_{0}^{T}|\zeta(t)|\left\|u_{F} \times n\right\|_{L^{2}(\partial \Omega)} \leq C^{\prime}\|\zeta\|_{L^{\infty}(0, T)} \sqrt{T},
\end{aligned}
$$

- By (3.39) together with (3.46):

$$
\begin{aligned}
\left|I_{7}\right| & \left.=\left\lvert\, \int_{0}^{T} \int_{\partial S(t)}\left(\left(u_{F}-u_{S}\right) \times n\right) \cdot\left(\frac{1}{\beta_{S}}\left(\varphi_{h}-e_{3}\right) \times n\right)+2 D\left(\varphi_{h}\right) n\right.\right) d \sigma \mid \\
& \left.\leq C \int_{0}^{T} \mid \zeta(t)\left\|\left(u_{F}-u_{S}\right) \times n\right\|_{L^{2}(\partial S(t))} \|\left(\varphi_{h}-e_{3}\right) \times n\right)+2 \beta_{S} D\left(\varphi_{h}\right) n \|_{L^{2}(\partial S(t))} \mathrm{d} t \\
& \leq C^{\prime}\|\zeta\|_{L^{\infty}(0, T)} \sqrt{T} .
\end{aligned}
$$

- As for $I_{8}$, applying (3.45), we have that

$$
\begin{aligned}
I_{8}=\int_{0}^{T} \int_{F(t)}\left(\Delta \varphi_{F}-\nabla q\right) u_{F} & \leq C \int_{0}^{T}|\zeta(t)|\left(\left\|D\left(u_{F}\right)\right\|_{L^{2}(F(t))}+\left\|u_{F}\right\|_{L^{2}(\partial \Omega)}\right) d t \\
& \leq C^{\prime}\|\zeta\|_{L^{\infty}(0, T)} \sqrt{T} .
\end{aligned}
$$

- Eventually, we get easily the following controls:

$$
\left|I_{2}\right| \leq C\left\|\zeta^{\prime}\right\|_{L^{1}(0, T)}, \quad I_{5} \leq C .
$$

Combining all these estimates to bound $R H S$, we end up with:

$$
\begin{align*}
& \mu_{F} \int_{0}^{T} h^{\prime}(t) \zeta(t) n(h(t)) \mathrm{d} t \\
& \quad=\frac{4 \pi\left(\rho_{F}-\rho_{S}\right) g}{3} \int_{0}^{T} \zeta(t) \mathrm{d} t+O\left(C_{0}\left(\|\zeta\|_{L^{\infty}(0, T)}+\left\|\zeta^{\prime}\right\|_{L^{1}(0, T)}\right)(1+\sqrt{T})\right) \tag{3.52}
\end{align*}
$$

We conclude by introducing a sequence $\zeta_{n}$ in $\mathcal{D}([0, T))$ such that

$$
\zeta_{n} \rightarrow 1_{[0, T)} \text { pointwise, } \quad \text { with } \quad\left\|\zeta_{n}\right\|_{L^{\infty}} \leq C, \quad\left\|\zeta_{n}^{\prime}\right\|_{L^{1}(0, T)} \leq C .
$$

A possible choice is $\zeta_{n}(t)=1_{[0, T)}(t)-\chi(n(T-t))$, for some smooth $\chi$ compactly supported with $\chi=1$ near 0 . We take $\zeta=\zeta_{n}$ in (3.52), and let $n$ go to infinity, to get

$$
\begin{equation*}
\mu_{F} \int_{h(0)}^{h(T)} n(s) \mathrm{d} s+\frac{4 \pi\left(\rho_{S}-\rho_{F}\right) g}{3} T \leq C_{0}(1+\sqrt{T}) . \tag{3.53}
\end{equation*}
$$

Introducing the bounds of Lemma 3.2 and Lemma 3.3 into (3.20), we show there exists a constant $C$ such that:

$$
n(h) \leq C|\ln h| \quad \forall h<h_{M},
$$

which implies that there exists a constant $C$ for which:

$$
\int_{h(0)}^{h} n(s) d s \geq-C, \quad \forall h<h_{M} .
$$

Hence, we obtain a contradiction in (3.53) if $T$ is large enough compared to $C$ and $C_{0}$ which are fixed by $h_{M}$ and the initial data only.

## 4 Proof of Theorem 1.2

The proof is very close to the one performed above for slip boundary conditions. Therefore, we only sketch the proof, insisting on its few specific features.

Computation of the drag. Let $(S, u)$ be a weak solution given by Theorem 2.4 on $[0, T)$. As before, we consider a solution up to its "maximal time of existence": that is $T=\infty$ if there is no collision, or $T<\infty$ the time of collision. We assume that $(S, u)$ satisfies A1-A4 over $[0, T)$. This time, we want to show that there is no collision, which means that $h$ does not go to zero in finite time (and thus $T=\infty$ ). We shall consider a test function $\varphi \in \mathcal{T}_{T}^{m}$ of the form

$$
\begin{equation*}
\varphi=\zeta(t) \varphi_{h(t)} \quad \text { with } \quad \varphi_{S}=\zeta(t) e_{z} \tag{4.54}
\end{equation*}
$$

associated with a pressure $q \in \mathcal{D}([0, T) \times \bar{\Omega})$, of the form

$$
q=\zeta(t) q_{h(t)}
$$

The values of $\varphi_{h}$ and $q_{h}$ are given below. As in the previous section, we shall use functions $\zeta=\zeta_{n}$ approximating $1_{(0, T)}$. Here and in what follows, we keep the convention that $\varphi_{F}=$ $\varphi 1_{F(t)}$ and $\varphi_{S}=\varphi 1_{S(t)}$. With computations similar to those of the previous section, we obtain:

$$
\begin{align*}
& \mu_{F} \int_{0}^{T} \zeta(t) h^{\prime}(t) n(h(t)) \mathrm{d} t \\
= & \int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \cdot \partial_{t} \varphi_{F}+\int_{0}^{T} \int_{F(t)} \rho_{F} u_{F} \otimes u_{F}: \nabla \varphi_{F} \\
- & \int_{0}^{T} \int_{F(t)} \rho_{F} g e_{z} \varphi_{F}-\int_{0}^{T} \int_{S(t)} \rho_{S} g e_{z} \varphi_{S}+\int_{F_{0}} \rho_{F} u_{F, 0} \varphi_{F}(0, \cdot)+\int_{S_{0}} \rho_{S} u_{S, 0} \varphi_{S}(0, \cdot) \\
- & \frac{\mu_{F}}{\beta_{\Omega}} \int_{0}^{T} \int_{P}\left(u_{F} \times n\right) \cdot\left(\varphi_{F} \times n\right) \\
+ & \mu_{F}\left(\int_{0}^{T} \int_{F(t)}\left(\Delta \varphi_{F}-\nabla q\right) \cdot u_{F}-\int_{0}^{T} \int_{\partial \Omega}\left(2 D\left(\varphi_{F}\right)-q I\right) n \cdot u_{F} d \sigma\right) \tag{4.55}
\end{align*}
$$

Again, we compute $n(h)$ applying the identity:

$$
\begin{align*}
n(h) & =\int_{F_{h}}\left(\Delta \varphi_{h}-\nabla q_{h}\right) \cdot \varphi_{h} d x+2 \int_{F_{h}} D\left(\varphi_{h}\right): D\left(\varphi_{h}\right) d x \\
& -\int_{\partial \Omega}\left(2 D\left(\varphi_{h}\right)-q_{h} I\right) n \cdot \varphi_{h} d \sigma \tag{4.56}
\end{align*}
$$

Construction of the test-function. For the function $(x, h) \mapsto \varphi_{h}(x)$, we keep the structure of (3.33) with a test-function $\widetilde{\varphi}_{h}$ still of the form (3.30). The only point that differs is the choice of the coefficients in the polynomial $\Phi$. We set now:

$$
P_{1}(r) \triangleq \frac{6}{4+\alpha_{P}} \quad P_{2}(r) \triangleq \frac{3 \alpha_{P}}{4+\alpha_{P}} \quad P_{3}(r) \triangleq-2 \frac{1+\alpha_{P}}{4+\alpha_{P}}
$$

with $\alpha_{P}$ given by (3.25). We remark that this is equivalent to taking the limit $\alpha_{S} \rightarrow \infty$ in our choice of test-function in the previous section. With this choice, the function $\Phi$ satisfies boundary conditions:

- for $z=0$

$$
\begin{align*}
& \Phi(r, 0)=0  \tag{4.57}\\
& \partial_{z z} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)-\frac{1}{\beta_{\Omega}} \partial_{z} \Phi\left(r, \frac{z}{h+\gamma_{s}(r)}\right)=0 \tag{4.58}
\end{align*}
$$

- for $z=h+\gamma_{S}(r)$

$$
\begin{align*}
& \Phi(r, 1)=1  \tag{4.59}\\
& \partial_{z} \Phi(r, 1)=0 . \tag{4.60}
\end{align*}
$$

Similarly to the previous case, we provide also the following technical lemmas :
Lemma 4.1. The map $(h, x) \mapsto \varphi_{h}(x)$ is smooth on all domains :

$$
\left\{(h, x), h \in\left(0, h_{M}\right), x \in F_{h}\right\},\left\{(h, x) \in\left(0, h_{M}\right), x \in S_{h}\right\},\left\{(h, x), h \in\left[0, h_{M}\right), x \in F_{h} \backslash \Omega_{h, \delta}\right\} .
$$

Furthermore, there holds:

$$
\begin{equation*}
\varphi_{h} \cdot n=0 \quad \text { on } \partial \Omega \tag{4.61}
\end{equation*}
$$

and

$$
\begin{equation*}
\left.\varphi_{h}\right|_{F_{h}}=e_{3} \quad \text { on } \partial S_{h} . \tag{4.62}
\end{equation*}
$$

Lemma 4.2. There exist constants $0<c<C$ so that, when $h<h_{M}$ :

$$
\begin{align*}
\left\|\varphi_{h}\right\|_{L^{2}\left(F_{h}\right)} & \leq C,  \tag{4.63}\\
\left\|\nabla \varphi_{h}\right\|_{L^{2}\left(F_{h}\right)}^{2} & \leq \frac{C}{h},  \tag{4.64}\\
\left\|D\left(\varphi_{h}\right)\right\|_{L^{2}\left(F_{h}\right)}^{2} & \geq \frac{c}{h}, \tag{4.65}
\end{align*}
$$

Furthermore, for $i=1,2,3$, denoting by exponent $i$ cartesian coordinates, there holds:

$$
\begin{align*}
\left\|\int_{0}^{h+\gamma_{s}(r)} \partial_{h} \varphi_{h}^{i} d z\right\|_{L^{2}\left(\partial \Omega \cap \partial \Omega_{h, \delta}\right)} & \leq C|\ln h|,  \tag{4.66}\\
\left\|\int_{z}^{h+\gamma_{s}(r)} \partial_{h} \varphi_{h}^{i} d s\right\|_{L^{2}\left(\Omega_{h, \delta}\right)} & \leq C,  \tag{4.67}\\
\left\|\left(h+\gamma_{s}(r)\right) \partial_{i} \int_{z}^{h+\gamma_{s}(r)} \partial_{h} \varphi_{F}^{i}\left(x_{1}, x_{2}, s\right) d s\right\|_{L^{2}\left(\Omega_{h, \delta}\right)} & \leq C, \tag{4.68}
\end{align*}
$$

and

- for $i \in\{1,2,3\}^{2}$, and $x \in \Omega_{h, \delta}$ :

$$
\begin{equation*}
\left|\partial_{i} \varphi_{h}^{3}+\partial_{3} \varphi_{h}^{i}\right| \leq \frac{C r}{\left(h+\gamma_{S}(r)\right)^{2}}, \tag{4.69}
\end{equation*}
$$

- for $(i, j) \in\{1,2\}^{2}$, and $x \in \Omega_{h, \delta}$ :

$$
\begin{align*}
\left|\int_{z}^{h+\gamma_{S}(r)}\left(\partial_{i} \varphi_{h}^{j}\left(x_{1}, x_{2}, s\right)+\partial_{j} \varphi_{h}^{i}\left(x_{1}, x_{2}, s\right)\right) d s\right| & \leq C  \tag{4.70}\\
\left|\left(h+\gamma_{S}(r)\right) \partial_{i} \int_{z}^{h+\gamma_{S}(r)}\left(\partial_{i} \varphi_{h}^{j}\left(x_{1}, x_{2}, s\right)+\partial_{j} \varphi_{h}^{i}\left(x_{1}, x_{2}, s\right)\right) d s\right| & \leq C . \tag{4.71}
\end{align*}
$$

Lemma 4.3. There exists a map $(h, x) \mapsto q_{h}$ which is smooth on the domain $\{(h, x), h \in$ $\left.\left(0, h_{M}\right), x \in F_{h}\right\}$ and which satisfies :

- there exists a constant positive constant $C$ such that :

$$
\left|\int_{\partial \Omega}\left(2 D\left(\varphi_{h}\right)-q_{h} I\right) n \cdot \varphi_{h} d \sigma\right| \leq C|\ln h|
$$

- there exists a positive constant $C$ such that, for any $h \in\left(0, h_{M}\right)$ and $v \in H^{1}(\Omega)$ satisfying $v_{S_{h}} \in \mathcal{R}$ and $v \cdot n=0$ on $\partial \Omega$, we have:

$$
\begin{equation*}
\left|\int_{F_{h}}\left(\Delta \varphi_{h}-\nabla q_{h}\right) v\right| \leq C\left(\|D(v)\|_{L^{2}\left(F_{h}\right)}+\|v\|_{L^{2}(\Omega)}\right) \tag{4.72}
\end{equation*}
$$

More details concerning these lemmas are provided in Appendix B.

Conclusion. Introducing the results of Lemma 4.2 and Lemma 4.3, we compute again the right-hand side of (4.55) with the help of energy estimates satisfied by $u$. We only point out two differences. First, there exists a universal constant $C$ so that

$$
\left\|u_{F}\right\|_{L^{2}(\partial S(t))}=\left\|u_{S}\right\|_{L^{2}(\partial S(t))} \leq C\left\|u_{S}\right\|_{L^{2}(S(t))} \leq C\|u\|_{L^{2}(\Omega)}
$$

by taking into account that $u_{S}$ is a rigid vector field. Second, one term requires a new treatment (we drop summation over index $i \in\{1,2,3\}$ for simplicity):

$$
\begin{aligned}
& \int_{F(t)} u_{F}^{i} u_{F}^{3}\left(\partial_{i} \varphi_{F}^{3}+\partial_{3} \varphi_{F}^{i}\right) \\
= & \int_{F(t) \cap \Omega_{h(t), 2 \delta}} \chi \frac{u_{F}^{3}}{\gamma_{s}(r)+h} u_{F}^{i}\left(\gamma_{s}(r)+h\right)\left(\partial_{i} \varphi_{F}^{3}+\partial_{3} \varphi_{F}^{i}\right)+O\left(\left\|u_{F}\right\|_{L^{2}}^{2}\right) \\
\leq & \left\|u_{F}^{i}\left(\gamma_{s}(r)+h\right)\left(\partial_{i} \varphi_{F}^{3}+\partial_{3} \varphi_{F}^{i}\right)\right\|_{L^{2}\left(F(t) \cap \Omega_{h(t), 2 \delta}\right)}\left\|\frac{u_{F}^{3}}{\gamma_{s}(r)+h}\right\|_{L^{2}\left(F(t) \cap \Omega_{h(t), 2 \delta}\right)}+O\left(\left\|u_{F}\right\|_{L^{2}}^{2}\right) \\
\leq & C\left\|D\left(u_{F}\right)\right\|_{L^{2}(F(t))}\left\|\chi u_{F}^{i}\left(\gamma_{s}(r)+h\right)\left(\partial_{i} \varphi_{F}^{3}+\partial_{3} \varphi_{F}^{i}\right)\right\|_{L^{2}\left(F(t) \cap \Omega_{h(t), 2 \delta}\right)}+O\left(\left\|u_{F}\right\|_{L^{2}(F(t))}^{2}\right) .
\end{aligned}
$$

where $\chi$ is given by (3.32). Because of (4.69), there holds

$$
\begin{aligned}
&\left\|\chi u_{F}^{i}\left(\gamma_{s}(r)+h\right)\left(\partial_{i} \varphi_{F}^{3}+\partial_{3} \varphi_{F}^{i}\right)\right\|_{L^{2}\left(F(t) \cap \Omega_{h, 2 \delta}\right)}^{2} \\
& \leq \int_{F(t)}\left|u_{F}^{i}\right|^{2} \cdot \frac{\chi}{\gamma_{s}(r)+h} \\
&= \int_{F(t)}\left(\partial_{z} u_{F}^{i}+\partial_{i} u_{F}^{3}\right) \chi u_{F}^{i} \frac{z-\gamma_{s}(r)-h}{\gamma_{s}(r)+h}-\int_{F(t)} \chi u_{F}^{i} \partial_{i} u_{F}^{3} \frac{z-\gamma_{s}(r)-h}{\gamma_{s}(r)+h} \\
& \quad+O\left(\left\|u_{F}\right\|_{L^{2}(F(t))}^{2}+\left\|u_{F}\right\|_{L^{2}(\partial \Omega)}^{2}\right) \\
&= \int_{F(t)} \frac{u_{F}^{3}}{\gamma_{s}(r)+h} u_{F}^{i} \chi\left(\gamma_{s}(r)+h\right) \partial_{i}\left(\frac{z-\gamma_{s}(r)-h}{\gamma_{s}(r)+h}\right) \\
&\left.\quad+O\left(\left\|u_{F}\right\|_{L^{2}(F(t))}^{2}+\left\|u_{F}\right\|_{L^{2}(\partial \Omega)}^{2}+\left\|D\left(u_{F}\right)\right\|_{L^{2}(F(t))}\right)\left\|u_{F}\right\|_{\left.L^{2}(F(t))\right)}\right) \\
&= O\left(\left\|u_{F}\right\|_{L^{2}(F(t))}^{2}+\left\|u_{F}\right\|_{L^{2}(\partial \Omega)}^{2}+\left\|D\left(u_{F}\right)\right\|_{L^{2}(F(t))}\left\|u_{F}\right\|_{\left.L^{2}(F(t))\right)}\right) .
\end{aligned}
$$

which implies that

$$
\left|\int_{F(t)} u_{F}^{i} u_{F}^{3}\left(\partial_{i} \varphi_{F}^{3}+\partial_{3} \varphi_{F}^{i}\right)\right|=O\left(\left\|D\left(u_{F}\right)\right\|_{L^{2}(F(t))}\left(\left\|D\left(u_{F}\right)\right\|+\left\|u_{F}\right\|_{L^{2}(\partial \Omega)}\right)+\left\|u_{F}\right\|_{L^{2}(F(t))}^{2}\right)
$$

Combining with estimates similar to the previous section for the other terms, we arrive at:

$$
\begin{equation*}
\left|\mu_{F} \int_{h(0)}^{h(T)} n(s) \mathrm{d} s\right| \leq C_{0}(1+T) . \tag{4.73}
\end{equation*}
$$

Introducing the bounds of Lemma 4.2 and Lemma 4.3 into (3.20), we show there exists a constant $C$ such that:

$$
n(h) \geq \frac{C}{h}, \quad \forall h \in\left(0, h_{M}\right),
$$

so that (4.73) yields finally :

$$
|\ln (h(T))| \leq C_{0}(1+T), \quad \forall T>0 .
$$

This ends the proof.

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## A Estimates on the test-function in the slip case

In this section, we collect estimates on the test-function of Section 3 as defined in (3.33). For further computations, we define,

$$
\Psi(r, z):=\Phi\left(r, \frac{z}{h+\gamma_{S}(r)}\right), \quad \forall h \in\left(0, h_{M}\right), \quad \forall(r, z) \in \Omega_{h, 1}
$$

In particular, we note that $\Psi\left(r, h+\gamma_{S}(r)\right)=1$ so that:

$$
\begin{aligned}
\frac{\mathrm{d}}{\mathrm{~d} r}\left[\frac{r^{2}}{2} \Psi\left(r, h+\gamma_{S}(r)\right)\right] & =r \\
& =\partial_{r}\left(\frac{r^{2}}{2} \Psi\right)\left(r, h+\gamma_{S}(r)\right)+\frac{r}{\sqrt{1-r^{2}}} \frac{r^{2}}{2} \partial_{z} \Psi\left(r, h+\gamma_{S}(r)\right)
\end{aligned}
$$

Thus, on $\partial S_{h} \cap \partial \Omega_{h, 1}$, there holds :

$$
\begin{aligned}
n \cdot \tilde{\varphi}_{h} & =\frac{r^{2}}{2} \partial_{z} \Psi\left(r, h+\gamma_{S}(r)\right)+\frac{\sqrt{1-r^{2}}}{r} \partial_{r}\left(\frac{r^{2}}{2} \Psi\right)\left(r, h+\gamma_{S}(r)\right) \\
& =\sqrt{1-r^{2}}=e_{3} \cdot n
\end{aligned}
$$

Before going to estimates on $\varphi_{h}$, we first obtain a precise decription of the size of $\Psi$ and its derivatives. Here and in what follows we use $a \lesssim b$ to denote that $a \leq C b$ for a constant $C$ independant of relevant parameters.

Proposition A.1. The following inequalities hold true, for all values of $h \in\left(0, h_{M}\right)$ and $(r, z) \in \Omega_{h, \delta}:$

$$
\begin{align*}
|\Psi(r, z)| & \lesssim 1,  \tag{A.74}\\
\left|\partial_{z} \Psi(r, z)\right| & \lesssim \frac{1}{h+\gamma_{S}(r)}, \quad\left|\partial_{r} \Psi(r, z)\right| \lesssim \frac{r}{h+\gamma_{S}(r)},  \tag{A.75}\\
\left|\partial_{r r} \Psi(r, z)\right| & \lesssim \frac{1}{h+\gamma_{S}(r)}, \quad\left|\partial_{z z} \Psi(r, z)\right| \lesssim \frac{1}{h+\gamma_{S}(r)},  \tag{A.76}\\
\left|\partial_{r z} \Psi(r, z)+\frac{r}{h+\gamma_{S}(r)} \partial_{z} \Psi(r, z)\right| & \lesssim \frac{r}{h+\gamma_{S}(r)},  \tag{A.77}\\
\left|\partial_{z z z} \Psi(r, z)\right|+\left|\partial_{r r z} \Psi(r, z)\right| & \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)^{2}},  \tag{A.78}\\
\left|\partial_{r z z} \Psi(r, z)\right| \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{2}}, & \left|\partial_{r r r} \Psi(r, z)\right| \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{2}}+\frac{1}{h+\gamma_{s}(r)} \tag{A.79}
\end{align*}
$$

and also

$$
\begin{align*}
\left|\partial_{h} \Psi(r, z)\right| & \lesssim \frac{1}{h+\gamma_{S}(r)},  \tag{A.80}\\
\left|\partial_{r h} \Psi(r, z)\right| & \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)}+\frac{r}{\left(h+\gamma_{S}(r)\right)^{2}},  \tag{A.81}\\
\left|\partial_{z z h} \Psi(r, z)\right|+\left|\partial_{r r h} \Psi(r, z)\right|+\left|\partial_{h z} \Psi(r, z)\right| & \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)^{2}},  \tag{A.82}\\
\left|\partial_{r z h} \Psi(r, z)\right| & \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)^{2}}+\frac{r}{\left(h+\gamma_{S}(r)\right)^{3}} . \tag{A.83}
\end{align*}
$$

Proof. At first, we note that

$$
\begin{aligned}
\left|\partial_{r} \alpha_{S}\right|+\left|\partial_{r} \alpha_{P}\right| & \lesssim r \\
\left|\partial_{r r r} \alpha_{S}\right|+\left|\partial_{r r} \alpha_{S}\right|+\left|\partial_{r r r} \alpha_{P}\right|+\left|\partial_{r r} \alpha_{P}\right| & \lesssim 1
\end{aligned}
$$

This yields, by (3.21)-(3.25), that we have:

$$
\left|\partial_{r} \Phi(r, t)\right|+\left|\partial_{r t t} \Phi(t, r)\right|+\left|\partial_{t r} \Phi(t, r)\right| \lesssim r, \quad \forall(r, t) \in(0, \delta) \times(0,1)
$$

and

$$
\left|\partial_{t} \Phi(r, t)\right|+\left|\partial_{r r r} \Phi(t, r)\right|+\left|\partial_{r r t} \Phi(t, r)\right|+\left|\partial_{r r} \Phi(t, r)\right| \lesssim 1 .
$$

We remark also that

$$
\left|\partial_{t t t} \Phi(t, r)\right|+\left|\partial_{t t} \Phi(t, r)\right| \lesssim\left(\alpha_{P}+\alpha_{S}\right) \lesssim\left(h+\gamma_{S}(r)\right)
$$

We introduce these bounds implicitly in the following computations.
For instance, we have:

$$
\begin{aligned}
\partial_{r} \Psi(r, z) & =\partial_{r} \Phi\left(r, \frac{z}{h+\gamma_{S}(r)}\right)-\frac{z \gamma_{S}^{\prime}(r)}{\left(h+\gamma_{S}(r)\right)^{2}} \partial_{t} \Phi\left(r, \frac{z}{\left(h+\gamma_{S}(r)\right.}\right) \\
& \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)}
\end{aligned}
$$

We obtain similarly (A.75)-(A.79) after tedious calculations. The only computation we make precise is the one concerning $\partial_{r z} \Psi$. Differentiating the last equality w.r.t. $z$, we get :

$$
\partial_{r z} \Psi(r, z)=\frac{1}{h+\gamma_{S}(r)} \partial_{r t} \Phi-\frac{\gamma_{S}^{\prime}(r)}{\left(h+\gamma_{S}(r)\right)^{2}} \partial_{t} \Phi-\frac{z \gamma_{S}^{\prime}(r)}{\left(h+\gamma_{S}(r)\right)^{3}} \partial_{t t} \Phi
$$

where :

$$
\partial_{z} \Psi(r, z)=\frac{1}{\left(h+\gamma_{S}(r)\right)} \partial_{t} \Phi
$$

Thus,

$$
\begin{aligned}
\partial_{r z} \Psi(r, z)+\frac{r}{h+\gamma_{S}(r)} \partial_{z} \Psi(r, z) & =\frac{1}{\left(h+\gamma_{S}(r)\right)} \partial_{r t} \Phi+\frac{r-\gamma_{S}^{\prime}(r)}{\left(h+\gamma_{S}(r)\right)^{2}} \partial_{t} \Phi \\
& -\frac{z \gamma_{S}^{\prime}(r)}{\left(h+\gamma_{S}(r)\right)^{3}} \partial_{t t} \Phi
\end{aligned}
$$

where we dominate :

$$
\left|r-\gamma_{S}^{\prime}(r)\right| \lesssim r^{3}
$$

Combining with the previous estimates on the derivative of $\Phi$, we obtain the expected result.
Concerning derivatives w.r.t. $h$ the computations are similar. First, we note that :

$$
\left|\partial_{r}^{k} \partial_{h} \alpha_{S}(r)\right|+\left|\partial_{r}^{k} \partial_{h} \alpha_{P}(r)\right| \lesssim 1, \quad \forall k \in \mathbb{N}
$$

Correspondingly, we obtain that

$$
\left|\partial_{r}^{k} \partial_{t}^{l} \partial_{h} \Phi\right| \lesssim 1, \quad \forall(k, l) \in \mathbb{N}^{2}
$$

Then, we obtain, for instance :

$$
\begin{aligned}
\partial_{r h} \Psi(r, z) & =\partial_{r h} \Phi-\frac{z}{\left(h+\gamma_{S}(r)\right)^{2}} \partial_{r t} \Phi+2 \frac{z \gamma_{S}^{\prime}(r)}{\left(h+\gamma_{S}(r)\right)^{3}} \partial_{t} \Phi \\
& -\frac{z \gamma_{S}^{\prime}(r)}{\left(h+\gamma_{S}(r)\right)^{2}} \partial_{t h} \Phi+\frac{z^{2} \gamma_{S}^{\prime}(r)}{\left(h+\gamma_{S}(r)\right)^{4}} \partial_{t t} \Phi
\end{aligned}
$$

in which, we plug previous estimates on $\Phi$ and its derivative. This yields :

$$
\left|\partial_{r h} \Psi(r, z)\right| \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)}+\frac{r}{\left(h+\gamma_{S}(r)\right)^{2}}
$$

Other estimates are obtained similarly.
The rest of this appendix is devoted to the proofs of lemmas 3.2 and 3.3 . They rely on the following lemma whose proof can be found in [13]:

Lemma A.2. Given $(p, q) \in(0, \infty)^{2}$, when $h$ goes to 0 , the quantity

$$
\int_{0}^{\delta} \frac{|x|^{p} d x}{\left(h+|x|^{2}\right)^{q}}
$$

behaves like
(1) $h^{\frac{p+1}{2}-q}$, if $p+1<2 q$,
(2) $\ln h$, if $p+1=2 q$,
(3) $C$, if $p+1>2 q$,
with $C$ a constant depending on $(p, q)$.

## A. 1 Proof of Lemma 3.2

First, we recall that, by construction $\varphi_{h}$ is smooth outside $\Omega_{h, \delta}$. Thus we have get that

$$
\left\|\nabla^{k} \varphi_{h}\right\|_{L^{\infty}\left(F_{h} \backslash \Omega_{h, \delta}\right)} \lesssim 1, \quad \forall k \in \mathbb{N}
$$

Hence, we focus on $\Omega_{h, \delta}$ and replace $\varphi_{h}$ with $\widetilde{\varphi}_{h}$ in what remains. We keep notations of the previous lemma. In particular, we have by definition :

$$
\widetilde{\varphi}_{h}=\frac{1}{2}\left(-\partial_{z} \Psi(r, z) r e_{r}+\frac{1}{r} \partial_{r}\left[r^{2} \Psi(r, z)\right] e_{z}\right), \quad \text { in } \Omega_{h, \delta}
$$

Sobolev bounds. First, applying the bounds of the previous lemma, we get:

$$
\begin{aligned}
\left\|\widetilde{\varphi}_{h}\right\|_{L^{2}\left(\Omega_{h, \delta}\right)} & \lesssim\left\|r \partial_{z} \Psi\right\|_{L^{2}\left(\Omega_{h, \delta}\right)}+\left\|\frac{1}{r} \partial_{r}\left[r^{2} \Psi(r, z)\right]\right\|_{L^{2}\left(\Omega_{h, \delta}\right)} \\
& \lesssim\left[1+\int_{0}^{\delta} \frac{r^{3} \mathrm{~d} r}{\left(r^{2}+h\right)}\right]^{\frac{1}{2}} \lesssim 1 .
\end{aligned}
$$

Then, in cylindrical coordinates, we obtain:

$$
\begin{equation*}
\left|\nabla \widetilde{\varphi}_{h}\right| \lesssim\left|\partial_{z} \Psi\right|+\left|r \partial_{z z} \Psi\right|+\left|\partial_{r} \Psi\right|+\left|r \partial_{r r} \Psi\right|+\frac{r^{2}}{h+\gamma_{s}(r)}, \tag{A.84}
\end{equation*}
$$

where we used (A.77) to get rid of $\partial_{r z} \Psi$ terms. Consequently :

$$
\left\|\nabla \widetilde{\varphi}_{h}\right\|_{L^{2}\left(\Omega_{h, \delta}\right)}^{2} \lesssim 1+\int_{0}^{\delta} \frac{r \mathrm{~d} r}{\left(h+\gamma_{S}(r)\right)} \lesssim|\ln (h)| .
$$

On the other hand, we have

$$
\left\|\partial_{2} \varphi_{h}^{2}\right\|_{L^{2}\left(\Omega_{h, \delta}\right)}^{2}=\frac{1}{4}\left\|\left(1-\frac{r^{2}|\sin (\theta)|^{2}}{\left(h+\gamma_{S}(r)\right)^{2}}\right) \partial_{z} \Psi(r, z)\right\|_{L^{2}\left(\Omega_{h, \delta}\right)}^{2}+O(1)
$$

Let now $\varepsilon>0$ small enough. There holds:

$$
\begin{aligned}
\left\|\partial_{2} \varphi_{h}^{2}\right\|_{L^{2}\left(\Omega_{h, \delta}\right)}^{2} & \geq \frac{1}{8}\left\|\partial_{z} \Psi(r, z)\right\|_{L^{2}\left(\Omega_{h, \delta}\right) \cap\{|\theta|<\varepsilon\}}^{2}+O(1) \\
& \geq c \int_{0}^{1} \int_{0}^{\delta}\left|\partial_{t} \Phi(r, t)\right|^{2} \frac{r \mathrm{~d} r \mathrm{~d} t}{\left(h+\gamma_{S}(r)\right)}+O(1) .
\end{aligned}
$$

where $c=c_{\varepsilon}$ is a small positive constant. Finally, as $\partial_{t} \Phi$ is bounded from below, we get

$$
\left\|\partial_{2} \varphi_{h}^{2}\right\|_{L^{2}\left(\Omega_{h, \delta}\right)}^{2} \geq c^{\prime} \int_{0}^{1} \int_{0}^{\delta} \frac{r \mathrm{~d} r \mathrm{~d} t}{\left(h+\gamma_{S}(r)\right)}+O(1) \geq c|\ln h|
$$

Combining the above estimate yields :

$$
c|\ln h| \leq\left\|D\left(\varphi_{F}\right)\right\|_{L^{2}(F(t))}^{2} \leq\left\|\nabla \varphi_{F}\right\|_{L^{2}(F(t))}^{2} \leq C|\ln h| .
$$

This completes the proof of (3.36)-(3.38).
Boundary estimates. We proceed with estimate (3.39). By standard identities of differential geometry (see [3]), we have

$$
2 D\left(\widetilde{\varphi}_{h}\right) n \times n=\partial_{n} \widetilde{\varphi}_{h} \times n+\left(\widetilde{\varphi}_{h}-e_{3}\right) \times n .
$$

Thus, we only need to prove that

$$
\left\|\beta_{S} \partial_{n} \widetilde{\varphi}_{h} \times n+\left(\beta_{S}+1\right) \widetilde{\varphi}_{h} \times n\right\|_{L^{2}\left(\partial S_{h} \cap \partial \Omega_{h, 2 \delta}\right)} \leq C .
$$

In the aperture $\Omega_{h, 2 \delta}$ we recall that the normal to $\partial S_{h}$ directed toward $S_{h}$ satisfies $n=$ $-r e_{r}+((1+h)-z) e_{z}$. This yields:

$$
\begin{aligned}
2 \widetilde{\varphi}_{h} \times n & =\left(-\partial_{z} \Psi r e_{r}+\frac{1}{r} \partial_{r}\left[r^{2} \Psi\right] e_{z}\right) \times\left(-r e_{r}+((1+h)-z) e_{z}\right) \\
& =\left(r((1+h)-z) \partial_{z} \Psi-\partial_{r}\left[r^{2} \Psi\right]\right) e_{\theta}
\end{aligned}
$$

Note here that, on $\partial S_{h}$, we have $(1+h)-z-1 \lesssim r^{2}$ and that

$$
\partial_{r}\left[r^{2} \Psi\right] \lesssim r=O(1)
$$

Here and in the end of this proof $O(1)$ means bounded in $L^{2}\left(\partial S_{h} \cap \partial \Omega_{h, 2 \delta}\right)$. This yields:

$$
\widetilde{\varphi}_{h} \times n=\frac{r}{2} \partial_{z} \Psi e_{\theta}+O(1) .
$$

On the other hand, in the aperture $\Omega_{h, \delta}$ we have:

$$
\begin{aligned}
2 \partial_{n} \widetilde{\varphi}_{h} & =\left(r \partial_{r z}(\Psi r)-((1+h)-z) r \partial_{z z} \Psi\right) e_{r} \\
& +\left(-r \partial_{r}\left[\frac{1}{r} \partial_{r}\left[r^{2} \Psi\right]\right]+((1+h)-z) \partial_{z}\left[\frac{1}{r} \partial_{r}\left[r^{2} \Psi\right]\right]\right) e_{z}
\end{aligned}
$$

Introducing the estimates of Proposition A. 1 we get in particular:

$$
\begin{aligned}
\partial_{z}\left[\frac{1}{r} \partial_{r}\left[r^{2} \Psi\right]\right] & =2 \partial_{z} \Psi+r \partial_{r z} \Psi \\
& =\left(2-\frac{r^{2}}{h+\gamma_{S}(r)}\right) \partial_{z} \Psi+O(1)=\frac{2 h}{h+\gamma_{s}(r)} \partial_{z} \Psi+O(1)
\end{aligned}
$$

Similarly, we have:

$$
r \partial_{r z}(\Psi r)=r^{2} \partial_{r z} \Psi+r \partial_{z} \Psi=-r \partial_{z} \Psi+r\left(-\frac{r^{2}}{h+\gamma_{s}(r)}+2\right) \partial_{z} \Psi+O(1)=-r \partial_{z} \Psi+O(1)
$$

using that

$$
\begin{equation*}
r\left(-\frac{r^{2}}{h+\gamma_{s}(r)}+2\right) \partial_{z} \Psi=O\left(\frac{r h}{\left(h+\gamma_{s}(r)\right)^{2}}\right) \tag{A.85}
\end{equation*}
$$

is uniformly bounded in $L^{2}$ (use case 3 of Lemma A. 2 with $p=3, q=4$ ). Then,

$$
2 \partial_{n} \widetilde{\varphi}_{h}=-\left(r \partial_{z} \Psi+r \partial_{z z} \Psi\right) e_{r}+\frac{2 h}{h+\gamma_{s}(r)} \partial_{z} \Psi e_{z}+O(1)
$$

Recalling that $n=-r e_{r}+((1+h)-z) e_{z}$, and still applying (A.85), we obtain:

$$
\partial_{n} \widetilde{\varphi}_{h} \times n=\frac{r}{2}\left(\partial_{z z} \Psi+\partial_{z} \Psi\right) e_{\theta}+O(1)
$$

Finally, we obtain:

$$
\beta_{S} \partial_{n} \widetilde{\varphi}_{h} \times n+\left(\beta_{S}+1\right) \tilde{\varphi}_{h} \times n=\frac{r \beta_{S}}{2}\left(\partial_{z z} \Psi+\left(\frac{1}{\beta_{S}}+2\right) \partial_{z} \Psi\right)+O(1)
$$

By our choice of $\alpha_{S}$, it follows that (see (3.29))

$$
\partial_{z z} \Psi\left(r, h+\gamma_{S}(r)\right)+\left(\frac{1}{\beta_{S}}+2\right) \partial_{z} \Psi\left(r, h+\gamma_{S}(r)\right)=O(1)
$$

This yields (3.39).
$h$-derivatives. For $\partial_{h} \widetilde{\varphi}_{h}$, we note that, for $i=1,2,3$ we have:

$$
\left|\partial_{h} \tilde{\varphi}_{h}^{i}\right| \lesssim\left|r \partial_{z h} \Psi\right|+\left|\partial_{h} \Psi\right|+\left|r \partial_{r h} \Psi\right| \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{2}}+\frac{1}{\left(h+\gamma_{S}(r)\right)} .
$$

Consequently, as integrating with respect to the $z$-variable reduces to remove one power of ( $h+\gamma_{S}(r)$ ) in the denominator, we get

$$
\left\|\int_{z}^{h+\gamma_{s}(r)} \partial_{h} \widetilde{\varphi}_{F}^{i} d s\right\|_{L^{2}\left(F_{h} \cap \Omega_{h, \delta}\right)}^{2} \lesssim 1+\int_{0}^{\delta} \frac{r^{3} \mathrm{~d} r}{\left(h+\gamma_{S}(r)\right)} \lesssim 1 .
$$

and

$$
\left\|\int_{0}^{h+\gamma_{s}(r)} \partial_{h} \widetilde{\varphi}_{F}^{i} d z\right\|_{L^{2}\left(\partial \Omega \cap \partial \Omega_{h, \delta}\right)}^{2} \lesssim 1+\int_{0}^{\delta} \frac{r^{3} \mathrm{~d} r}{\left(h+\gamma_{S}(r)\right)^{2}} \lesssim C|\ln h| .
$$

Finally, we also compute: for $i=1,2$ :

$$
\begin{aligned}
\left|\partial_{i} \partial_{h} \widetilde{\varphi}_{h}\right| & \lesssim\left|\partial_{z h} \Psi\right|+\left|r \partial_{r z h} \Psi\right|+\left|\partial_{r h} \Psi\right|+\left|r \partial_{r r h} \Psi\right| \\
& \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)^{2}}
\end{aligned}
$$

so that for the same reason as above, we have:

$$
\left|\left(h+\gamma_{s}(r)\right) \partial_{i} \int_{z}^{h+\gamma_{s}(r)} \partial_{h} \widetilde{\varphi}_{h}^{i}\left(x_{1}, x_{2}, s\right) d s\right| \lesssim 1, \quad \text { for } i=1,2
$$

and

$$
\left\|\left(h+\gamma_{s}(r)\right) \partial_{i} \int_{z}^{h+\gamma_{s}(r)} \partial_{h} \widetilde{\varphi}_{F}^{i}\left(x_{1}, x_{2}, s\right) d s\right\|_{L^{2}\left(F_{h} \cap \Omega_{h, \delta}\right)} \lesssim 1, \quad \text { for } i=1,2,3 .
$$

This ends the proof of (4.66)-(4.68).
Pointwise bounds. We easily get that on $\Omega_{h, \delta}$

$$
\left|\nabla \varphi_{h}\right| \lesssim\left|r \partial_{r z} \Psi\right|+\left|\partial_{z} \Psi\right|+\left|r \partial_{r r} \Psi\right|+\left|\partial_{r} \Psi\right|+\left|r \partial_{z z} \Psi\right| \lesssim \frac{1}{h+\gamma_{S}(r)},
$$

from which we conclude:

$$
\left|\int_{z}^{h+\gamma_{S}(r)}\left(\partial_{i} \varphi_{h}^{j}\left(x_{1}, x_{2}, s\right)+\partial_{j} \varphi_{h}^{i}\left(x_{1}, x_{2}, s\right)\right) d s\right| \lesssim 1, \text { for }(i, j) \in\{1,2,3\}^{2} .
$$

Similarly, we get:
$\left|\partial_{1} \nabla \varphi_{h}\right|+\left|\partial_{2} \nabla \varphi_{h}\right| \lesssim\left|r \partial_{r r r} \Psi\right|+\left|\partial_{r r} \Psi\right|+\left|r \partial_{r r z} \Psi\right|+\left|\partial_{r z} \Psi\right|+\left|\partial_{z z} \Psi\right| \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{2}}+\frac{1}{h+\gamma_{s}}$. and we conclude:

$$
\left|\left(h+\gamma_{S}(r)\right) \partial_{i} \int_{z}^{h+\gamma_{S}(r)}\left(\partial_{i} \varphi_{F}^{j}\left(x_{1}, x_{2}, s\right)+\partial_{j} \varphi_{F}^{i}\left(x_{1}, x_{2}, s\right)\right) d s\right| \lesssim 1 .
$$

This ends the proof of our lemma.

## A. 2 Proof of Lemma 3.3

Boundary integrals We focus on the domain $\Omega_{h, \delta}$ and we replace $\varphi_{h}$ with $\widetilde{\varphi}_{h}$ in the proof because everything is bounded elsewhere. First, it turns out that the first identity is independant of $q_{h}$. Indeed, applying boundary conditions (3.31) we obtain :

$$
\begin{aligned}
\int_{\partial \Omega_{h, \delta} \cap \partial \Omega}\left(2 D\left(\widetilde{\varphi}_{h}\right)-q_{h}\right) n \cdot \widetilde{\varphi}_{h} \mathrm{~d} \sigma & =\int_{\partial \Omega_{h, \delta} \cap \partial \Omega}\left[\left(2 D\left(\widetilde{\varphi}_{h}\right) n\right) \times n\right] \cdot\left(\widetilde{\varphi}_{h} \times n\right) \mathrm{d} \sigma \\
& =-\frac{1}{\beta_{\Omega}} \int_{\partial \Omega_{h, \delta} \cap \partial \Omega}\left|\widetilde{\varphi}_{h} \times n\right|^{2} \mathrm{~d} \sigma
\end{aligned}
$$

Similarly, introducing boundary conditions (3.31) and estimate (3.39), we compute:

$$
\begin{aligned}
\int_{\partial \Omega_{h, \delta} \cap \partial S_{h}}\left(2 D\left(\widetilde{\varphi}_{h}\right)\right. & \left.-q_{h}\right) n \cdot\left(e_{3}-\widetilde{\varphi}_{h}\right) \mathrm{d} \sigma=\int_{\partial \Omega_{h, \delta} \cap \partial S_{h}}\left[\left(2 D\left(\widetilde{\varphi}_{h}\right) n\right) \times n\right] \cdot\left(\left(e_{3}-\widetilde{\varphi}_{h}\right) \times n\right) \mathrm{d} \sigma \\
& =\frac{1}{\beta_{S}} \int_{\partial \Omega_{h, \delta} \cap \partial S_{h}}\left|\left(e_{3}-\widetilde{\varphi}_{h}\right) \times n\right|^{2} \mathrm{~d} \sigma+O\left(\left\|\left(e_{3}-\widetilde{\varphi}_{h}\right) \times n\right\|_{L^{2}\left(\partial S_{h} \cap \Omega_{h, \delta}\right)}\right)
\end{aligned}
$$

On the boundaries, there holds

$$
\widetilde{\varphi}_{h}=-\frac{P_{1}(r)}{\left(h+\gamma_{S}(r)\right)} r e_{r} \quad \text { in } \quad \partial \Omega_{h, \delta} \cap \partial \Omega, \quad\left|\widetilde{\varphi}_{h}\right| \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)} \quad \text { in } \quad \partial \Omega_{h, \delta} \cap \partial S_{h}
$$

which implies that

$$
|\ln (h)| \lesssim \int_{0}^{\delta} \frac{r^{3} d r}{\left(h+\gamma_{S}(r)\right)^{2}} \lesssim\left\|\widetilde{\varphi}_{h} \times n\right\|_{L^{2}\left(\partial \Omega_{h, \delta} \cap \partial \Omega\right)}^{2}=\left\|\widetilde{\varphi}_{h}\right\|_{L^{2}\left(\partial \Omega_{h, \delta} \cap \partial \Omega\right)}^{2}
$$

and which implies that
$|\ln h| \lesssim \int_{\partial \Omega_{h, \delta} \cap \partial S_{h}}\left(2 D\left(\widetilde{\varphi}_{h}\right)-q_{h}\right) n \cdot\left(e_{3}-\widetilde{\varphi}_{h}\right) \mathrm{d} \sigma-\int_{\partial \Omega \cap \partial \Omega_{h, \delta}}\left(D\left(\widetilde{\varphi}_{h}\right)-q_{h} I\right) n \cdot \varphi_{h} d \sigma \lesssim|\ln (h)|$.
Construction of $q_{h}$. Again, we focus on the domain $\Omega_{h, \delta}$ where $\varphi_{h}=\widetilde{\varphi}_{h}$ to compute $q_{h}$. The construction is extended to the whole $F_{h}$ by a standard truncation argument.

Given $h>0$, let define

$$
2 q_{h} \triangleq-\left(r \partial_{r z} \Psi+2 \partial_{z} \Psi+\int_{0}^{r} \partial_{z z z} \Psi(s, z) s d s\right)
$$

Obviously, the map $(h, x) \mapsto q_{h}(x)$ is $C^{1}$ on $\left\{h \in\left(0, h_{M}\right), x \in F_{h}\right\}$. and we have $\Delta \widetilde{\varphi}_{h}-\nabla q_{h}=$ $f_{h} e_{3}$ where

$$
\left|f_{h}\right| \lesssim\left|r \partial_{r r r} \Psi\right|+\left|\partial_{r r} \Psi\right|+\left|\frac{\partial_{r} \Psi}{r}\right|+\left|r \partial_{r z z} \Psi\right|+\left|\partial_{z z} \Psi\right| \lesssim \frac{1}{\left(h+\gamma_{s}(r)\right)}
$$

Then, for arbitrary $h \in\left(0, h_{M}\right)$ and $v \in H^{1}\left(F_{h}\right)$ we introduce $\chi$ the truncation function as defined in (3.32) and we split :

$$
\int_{F_{h}}\left(\Delta \varphi_{h}-\nabla q_{h}\right) v=\int_{F_{h}} \chi\left(\Delta \varphi_{h}-\nabla q_{h}\right) v+\int_{F_{h}}(1-\chi)\left(\Delta \varphi_{h}-\nabla q_{h}\right) v=I+I I
$$

Then, we have:

$$
\begin{aligned}
|I| & =\int_{F_{h} \cap \Omega_{h, 2 \delta}}\left|\chi f_{h} \| v_{3}\right| \\
& \leq\left\|\left(h+\gamma_{S}\right) f_{h}\right\|_{L^{2}\left(F_{h} \cap \Omega_{h, 2 \delta}\right)}\left\|\frac{v_{3}}{\gamma_{s}(r)+h}\right\|_{L^{2}\left(F_{h} \cap \Omega_{h, 2 \delta}\right)}
\end{aligned}
$$

where a standard Poincaré inequality yields:

$$
\left\|\frac{v_{3}}{\gamma_{s}(r)+h}\right\|_{L^{2}\left(F_{h} \cap \Omega_{h, 2 \delta}\right)} \lesssim\left\|\partial_{3} v_{3}\right\|_{L^{2}\left(F_{h} \cap \Omega_{h, \delta}\right)} \lesssim\|D(v)\|_{L^{2}\left(F_{h} \cap \Omega_{h, \delta}\right)}
$$

As for the second integral, we introduce $P_{z} R^{i}$ for $i=1,2,3$, which satisfy :

$$
\begin{align*}
& P_{z} R^{i}=\int_{z}^{h+\gamma_{s}(r)}(1-\chi)\left(\Delta \varphi_{h}^{i}-\partial_{i} q_{h}\right) \quad \text { for } r<1, z<r+\gamma_{s}(r)  \tag{A.86}\\
& P_{z} R^{i}=\int_{z}^{h+1}(1-\chi)\left(\Delta \varphi_{h}^{i}-\partial_{i} q_{h}\right) \quad \text { otherwise. } \tag{A.87}
\end{align*}
$$

Repeating the computations of $I_{1}(t)$ and $I_{2}(t)$ in Section 3.3, we obtain (the summation over index $i$ is implicit) :

$$
|I I|=\int_{F_{h}}\left(\partial_{3} v^{i}+\partial_{i} v^{3}\right) P_{z} R^{i}+\int_{F_{h}} v^{3} \partial_{i} P_{z} R^{i}-\int_{\partial \Omega}\left(v^{3} n^{i}+v^{i} n^{3}\right) P_{z} R^{i} \mathrm{~d} \sigma .
$$

As $\left(\Delta \varphi_{h}-\nabla q_{h}\right)$ is uniformly bounded outside $\Omega_{h, \delta}$, applying a Cauchy Schwarz inequality to dominate $K_{2}$ and combining with $K_{1}$ concludes the proof :

$$
\left|\int_{F_{h}}\left(\Delta \varphi_{h}-\nabla q_{h}\right) \cdot v\right| \lesssim\|D(v)\|_{L^{2}\left(F_{h}\right)}+\|v\|_{L^{2}(\partial \Omega)}
$$

## B Estimates on the test function in the mixed case

In this appendix, we collect estimates on the test function of Section 4. As previously, we define,

$$
\Psi(r, z):=\Phi\left(r, \frac{z}{h+\gamma_{S}(r)}\right), \quad \forall h \in\left(0, h_{M}\right), \quad \forall(r, z) \in \Omega_{h, 1}
$$

The only point which differs from the test-function in the slip case is the decription of this $\Psi$. Combining these estimates in the same manner as in the previous appendix, we obtain Lemma 4.2. As the computations are completely similar, all the proofs are left to the reader.

Proposition B.1. The following inequalities hold true, for all $h \in\left(0, h_{M}\right)$ and $(r, z) \in \Omega_{h, \delta}$ :

$$
\begin{array}{rlrl}
|\Psi(r, z)| & \lesssim 1, & \\
\left|\partial_{z} \Psi(r, z)\right| & \lesssim \frac{1}{h+\gamma_{S}(r)}, & & \left|\partial_{r} \Psi(r, z)\right| \lesssim \frac{r}{h+\gamma_{S}(r)}, \\
\left|\partial_{r r} \Psi(r, z)\right| & \lesssim \frac{1}{h+\gamma_{S}(r)}, & \left|\partial_{z z} \Psi(r, z)\right| \lesssim \frac{1}{\left(h+\gamma_{S}(r)^{2}\right.}, \\
\left|\partial_{r z} \Psi(r, z)\right| & \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{2}}, & \left|\partial_{z z z} \Psi(r, z)\right| \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)^{3}} \\
\left|\partial_{r r z} \Psi(r, z)\right| & \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)^{2}}, & \left|\partial_{r z z} \Psi(r, z)\right| \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{3}}, \\
\left|\partial_{r r r} \Psi(r, z)\right| & \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{2}}, & \left|\partial_{r r r z} \Psi(r, z)\right| \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{3}}, \tag{B.93}
\end{array}
$$

and also

$$
\begin{align*}
\left|\partial_{h} \Psi(r, z)\right| & \lesssim \frac{1}{h+\gamma_{S}(r)}, \quad\left|\partial_{r h} \Psi(r, z)\right| \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{2}}  \tag{B.95}\\
\left|\partial_{r r h} \Psi(r, z)\right| & \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)^{2}}, \quad\left|\partial_{h z} \Psi(r, z)\right| \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)^{2}}  \tag{B.96}\\
\left|\partial_{r z h} \Psi(r, z)\right| & \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{3}}+\frac{1}{\left(h+\gamma_{S}(r)\right)^{2}}  \tag{B.97}\\
\left|\partial_{z z h} \Psi(r, z)\right| & \leq \frac{1}{\left(h+\gamma_{S}(r)\right)^{3}} . \tag{B.98}
\end{align*}
$$

We also need a different way to construct the pressure $q_{h}$ in order to prove Lemma 4.3. Again, we construct $q_{h}$ only in the aperture domain $\Omega_{h, \delta}$, the extension to $F_{h}$ being done by a standard truncation. So, for $h>0$ we set :

$$
2 q_{h} \triangleq r \partial_{r z} \Psi+2 \partial_{z} \Psi-\int_{0}^{r} \partial_{z z z} \Psi s \mathrm{~d} s
$$

In this way, we obtain $\Delta \varphi_{h}-\nabla q_{h}=f^{r} e_{r}+f^{z} e_{z}$ where the $f_{i}$ 's are bounded outside $\Omega_{h, \delta}$ and, in $\Omega_{h, \delta}$, we have:

$$
\begin{align*}
\left|f^{z}\right| & \lesssim\left|r \partial_{r r r} \Psi\right|+\left|\partial_{r r} \Psi\right|+\left|\frac{\partial_{r} \Psi}{r}\right| \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)}  \tag{B.99}\\
\left|f^{r}\right| & \lesssim\left|r \partial_{r r z} \Psi\right|+\left|\partial_{r z} \Psi\right| \lesssim \frac{r}{\left(h+\gamma_{S}(r)\right)^{2}}  \tag{B.100}\\
\left|\partial_{r} f^{r}\right|+\left|\frac{f^{r}}{r}\right| & \lesssim\left|r \partial_{r r r z} \Psi\right|+\left|\partial_{r r z} \Psi\right|+\left|\frac{\partial_{r z} \Psi}{r}\right| \lesssim \frac{1}{\left(h+\gamma_{S}(r)\right)^{2}} . \tag{B.101}
\end{align*}
$$

Then, for any $v \in H^{1}(\Omega)$ such that $v \cdot n=0$ on $\partial \Omega$ and $v \in \mathcal{R}$ on $S_{h}$ we have :

$$
\int_{F_{h}}\left(\Delta \varphi_{h}-\nabla q_{h}\right) \cdot v=\int_{F_{h}} f^{z} v^{3}+\int_{F_{h}} f^{r} e_{r} \cdot v .
$$

First, we apply Poincaré inequality, as $v^{3}=0$ on $\partial \Omega \cap \partial \Omega_{h, \delta}$, together with (B.99). This yields:

$$
\begin{aligned}
\left|\int_{F_{h}} f^{z} v_{3}\right| & \lesssim\left|\int_{F_{h} \cap \Omega_{h, \delta}} f^{z}\left(h+\gamma_{S}(r)\right) \frac{v_{3}}{\left(h+\gamma_{S}(r)\right)}\right|+\|v\|_{L^{2}\left(F_{h}\right)} \\
& \lesssim\|D(v)\|_{L^{2}\left(F_{h}\right)}+\|v\|_{L^{2}\left(F_{h}\right)}
\end{aligned}
$$

For the second integral: we write $f^{r} e_{r}=f^{1} e_{1}+f^{2} e_{2}$, and split the integral, with $\chi$ as defined in (3.32) :

$$
\int_{F_{h}} f^{i} v_{i}=\int_{F_{h}} \chi f^{i} v^{i}+O\left(\|v\|_{L^{2}\left(F_{h}\right)}\right)
$$

as $f^{r}$ is uniformly bounded outside $\Omega_{h, \delta}$. Then, we introduce $P_{z} f^{i}$, for $i=1,2$, which satisfies:

$$
P_{z} f^{i}=\int_{0}^{z} \chi f^{i}
$$

Through integration by parts, we get (with the Einstein convention of repeated indices):

$$
\begin{aligned}
\int_{F_{h}} \chi f^{i} v^{i}= & \int_{F_{h} \cap \Omega_{h, 2 \delta}} P_{z} f^{i}\left(\partial_{3} v^{i}+\partial_{i} v^{3}\right)+\int_{F_{h} \cap \Omega_{h, 2 \delta}} \frac{v_{3}}{h+\gamma_{S}(r)}\left(h+\gamma_{S}(r)\right)\left(\partial_{i} P_{z} f^{i}\right) \\
& -\int_{\partial S_{h} \cap \Omega_{h, 2 \delta}} P_{z} f^{i}\left(v^{i} n^{3}+v^{3} n^{i}\right) \triangleq I+I I+I I I
\end{aligned}
$$

Thanks to the bounds (B.100)-(B.101), we get:

$$
|I|+|I I| \lesssim\|D(v)\|_{L^{2}\left(F_{h}\right)}
$$

As regards the last integral, we simply write, because $v_{\left.\right|_{S_{h}}} \in \mathcal{R}$,

$$
|I I I| \lesssim\left\|P_{z} f^{i}\right\|_{L^{1}\left(\partial S_{h}\right)}\|v\|_{L^{\infty}\left(\partial S_{h}\right)} \lesssim\|v\|_{L^{2}\left(S_{h}\right)} \lesssim\|v\|_{L^{2}(\Omega)}
$$

This ends the proof.


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