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DYNAMICS OF A FREE BOUNDARY PROBLEM WITH CURVATURE MODELING ELECTROSTATIC MEMS

JOACHIM ESCHER, PHILIPPE LAURENÇOT, AND CHRISTOPH WALKER

ABSTRACT. The dynamics of a free boundary problem for electrostatically actuated microelectromechanical systems (MEMS) is investigated. The model couples the electric potential to the deformation of the membrane, the deflection of the membrane being caused by application of a voltage difference across the device. More precisely, the electrostatic potential is a harmonic function in the angular domain that is partly bounded by the deformable membrane. The gradient trace of the electric potential on this free boundary part acts as a source term in the governing equation for the membrane deformation. The main feature of the model considered herein is that, unlike most previous research, the small deformation assumption is dropped, and curvature for the deformation of the membrane is taken into account what leads to a quasilinear parabolic equation. The free boundary problem is shown to be well-posed, locally in time for arbitrary voltage values and globally in time for small voltage values. Furthermore, existence of asymptotically stable steady-state configurations is proved in case of small voltage values as well as non-existence of steady-states if the applied voltage difference is large. Finally, convergence of solutions of the free boundary problem to the solutions of the well-established small aspect ratio model is shown.

1. INTRODUCTION

The focus of this paper is the analysis of a model describing the dynamics of an electrostatically actuated microelectromechanical system (MEMS) when the deformation of the devices are not assumed to be small. More precisely, consider an elastic plate held at potential V and suspended above a fixed ground plate held at zero potential. The potential difference between the two plates generates a Coulomb force and causes a deformation of the membrane, thereby converting electrostatic energy into mechanical energy, a feature used in the design of several MEMS-based devices such as micropumps or microswitches [26]. An ubiquitous phenomenon observed in such devices is the so called “pull-in” instability: a threshold value of the applied voltage V above which the elastic response of the membrane cannot balance the Coulomb force and the deformable membrane smashes into the fixed plate. Since this effect might either be useful or, in contrast, could damage the device, its understanding is of utmost practical importance and several mathematical models have been set up for its investigation [7, 21, 26].

In the following subsection we give a brief description of an idealized device as depicted in Figure 1, where the state of the device is characterized by the electrostatic potential in the region between the two plates and the deformation of the membrane which is not assumed to be small from the outset, cf. [5].

1.1. The Model. To derive the model for electrostatic MEMS with curvature we proceed similarly to [5, 7, 22]. We consider a rectangular thin elastic membrane that is suspended above a rigid plate. The $(\hat{x}, \hat{y}, \hat{z})$ -coordinate system is chosen such that the ground plate of dimension $[-L, L] \times [0, l]$ in (\hat{x}, \hat{y}) -direction is located at $\hat{z} = -H$, while the undeflected membrane

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with the same dimension $[-L, L] \times [0, l]$ in (\hat{x}, \hat{y}) -direction is located at $\hat{z} = 0$. The membrane is held fixed along the edges in \hat{y} -direction while the edges in \hat{x} -direction are free. Assuming homogeneity in \hat{y} -direction, the membrane may thus be considered as an elastic strip and the \hat{y} -direction is omitted in the sequel. The mechanical deflection of the membrane is caused by a voltage difference that is applied across the device. The membrane is held at potential V while the rigid plate is grounded. We denote the deflection of the membrane at position \hat{x} and time \hat{t} by $\hat{u} = \hat{u}(\hat{t}, \hat{x}) > -H$ and the electrostatic potential at position (\hat{x}, \hat{z}) and time \hat{t} by $\hat{\psi} = \hat{\psi}(\hat{t}, \hat{x}, \hat{z})$. We do not indicate the time variable \hat{t} for the time being. The electrostatic potential $\hat{\psi}$ is harmonic, i.e.

$$\Delta \hat{\psi} = 0 \quad \text{in } \hat{\Omega}(\hat{u}) \quad (1.1)$$

and satisfies the boundary conditions

$$\hat{\psi}(\hat{x}, -H) = 0, \quad \hat{\psi}(\hat{x}, \hat{u}(\hat{x})) = V, \quad \hat{x} \in (-L, L), \quad (1.2)$$

where

$$\hat{\Omega}(\hat{u}) := \{(\hat{x}, \hat{z}); -L < \hat{x} < L, -H < \hat{z} < \hat{u}(\hat{x})\}$$

is the region between the ground plate and the membrane. The total energy of the system constitutes of the electric potential and the elastic energy and reads

$$E(\hat{u}) = E_e(\hat{u}) + E_s(\hat{u}).$$

The *electrostatic energy* E_e in dependence of the deflection \hat{u} is given by

$$E_e(\hat{u}) = -\frac{\epsilon_0}{2} \int_{-L}^L \int_{-H}^{\hat{u}(\hat{x})} |\nabla \hat{\psi}(\hat{x}, \hat{z})|^2 \, d\hat{z} \, d\hat{x}$$

with ϵ_0 being the permittivity of free space while the *elastic energy* E_s only retains the contribution due to stretching (in particular, bending is neglected) and is proportional to the tension T and to the change of surface area of the membrane, i.e.

$$E_s(\hat{u}) = T \int_{-L}^L \left(\sqrt{1 + (\partial_{\hat{x}} \hat{u}(\hat{x}))^2} - 1 \right) \, d\hat{x}.$$

Introducing the dimensionless variables

$$x = \frac{\hat{x}}{L}, \quad z = \frac{\hat{z}}{H}, \quad u = \frac{\hat{u}}{H}, \quad \psi = \frac{\hat{\psi}}{V}$$

and denoting the aspect ratio of the device by $\varepsilon = H/L$, we may write the total energy in these variables in the form

$$\begin{aligned} E(u) &= TL \int_{-1}^1 \left(\sqrt{1 + \varepsilon^2 (\partial_x u(x))^2} - 1 \right) \, dx \\ &\quad - \frac{\epsilon_0 V^2}{2\varepsilon} \int_{\Omega(u)} \left(\varepsilon^2 |\partial_x \psi(x, z)|^2 + |\partial_z \psi(x, z)|^2 \right) \, d(x, z), \end{aligned} \quad (1.3)$$

with

$$\Omega(u) := \{(x, z) \in (-1, 1) \times (-1, \infty) : -1 < z < u(x)\},$$

so that, formally, the corresponding Euler-Lagrange equations are

$$0 = \varepsilon^2 \partial_x \left(\frac{\partial_x u}{\sqrt{1 + \varepsilon^2 (\partial_x u)^2}} \right) - \varepsilon^2 \lambda \left(\varepsilon^2 |\partial_x \psi(x, u(x))|^2 + |\partial_z \psi(x, u(x))|^2 \right) \quad (1.4)$$

for $x \in (-1, 1)$, where we have set

$$\lambda = \frac{\epsilon_0 V^2}{TL\varepsilon^3}.$$

We now take again time into account and derive the dynamics of the dimensionless deflection $u = u(\hat{t}, x)$ by means of Newton's second law. Letting ρ and δ denote the mass density per unit volume of the membrane and the membrane thickness, respectively, the sum over all forces equals $\rho\delta\partial_{\hat{t}}^2 u$. The elastic and electrostatic forces, given by the right hand side of equation (1.4), are combined with a damping force of the form $-a\partial_{\hat{t}} u$ being linearly proportional to the velocity. This yields

$$\rho\delta\partial_{\hat{t}}^2 u + a\partial_{\hat{t}} u = \varepsilon^2 \partial_x \left(\frac{\partial_x u}{\sqrt{1 + \varepsilon^2 (\partial_x u)^2}} \right) - \varepsilon^2 \lambda \left(\varepsilon^2 |\partial_x \psi(x, u(x))|^2 + |\partial_z \psi(x, u(x))|^2 \right).$$

Finally, scaling time based on the strength of damping according to $t = \hat{t}\varepsilon^2/a$ and setting $\gamma := \frac{\sqrt{\rho\delta\varepsilon}}{a}$, we derive for the dimensionless deflection the evolution problem

$$\gamma^2 \partial_t^2 u + \partial_t u = \partial_x \left(\frac{\partial_x u}{\sqrt{1 + \varepsilon^2 (\partial_x u)^2}} \right) - \lambda \left(\varepsilon^2 |\partial_x \psi(x, u(x))|^2 + |\partial_z \psi(x, u(x))|^2 \right) \quad (1.5)$$

for $t > 0$ and $x \in I := (-1, 1)$. Instead of considering this hyperbolic equation, however, we assume in this paper that viscous or damping forces dominate over inertial forces, i.e. we assume that $\gamma \ll 1$ and thus neglect the second order time derivative term in (1.5). The membrane displacement $u = u(t, x) \in (-1, \infty)$ then evolves according to

$$\partial_t u - \partial_x \left(\frac{\partial_x u}{\sqrt{1 + \varepsilon^2 (\partial_x u)^2}} \right) = -\lambda \left(\varepsilon^2 |\partial_x \psi(t, x, u(x))|^2 + |\partial_z \psi(t, x, u(x))|^2 \right), \quad (1.6)$$

for $t > 0$ and $x \in I$ with clamped boundary conditions

$$u(t, \pm 1) = 0, \quad t > 0, \quad (1.7)$$

and initial condition

$$u(0, x) = u^0(x), \quad x \in I. \quad (1.8)$$

In dimensionless variables, equations (1.1)-(1.2) read

$$\varepsilon^2 \partial_x^2 \psi + \partial_z^2 \psi = 0, \quad (x, z) \in \Omega(u(t)), \quad t > 0, \quad (1.9)$$

subject to the boundary conditions (extended continuously to the lateral boundary)

$$\psi(t, x, z) = \frac{1+z}{1+u(t, x)}, \quad (x, z) \in \partial\Omega(u(t)), \quad t > 0. \quad (1.10)$$

In the following we shall focus our attention on (1.6)-(1.10), its situation being depicted in Figure 1.

1.2. Simplified Models. Besides assuming that damping forces dominate over inertial forces and thus reducing equation (1.5) to (1.6), other simplifications of the model above have been considered as well in the literature. For instance, restricting attention to small deformations of the membrane yields a linearized stretching term $\partial_x^2 u$ in (1.6). The corresponding semilinear evolution problem with (1.6) being replaced by

$$\partial_t u - \partial_x^2 u = -\lambda \left(\varepsilon^2 |\partial_x \psi(t, x, u(x))|^2 + |\partial_z \psi(t, x, u(x))|^2 \right), \quad x \in I, \quad t > 0, \quad (1.11)$$

is investigated in [6]. It is shown therein that the problem (1.7)-(1.11) is well-posed locally in time. Moreover, solutions exist globally for small voltage values λ while global existence is shown not to hold for high voltage values. It is also proven that, for small voltage values, there is an asymptotically stable steady-state solution. Finally, as the parameter ε approaches zero, the solutions are shown to converge toward the solutions of the so-called small aspect ratio model,

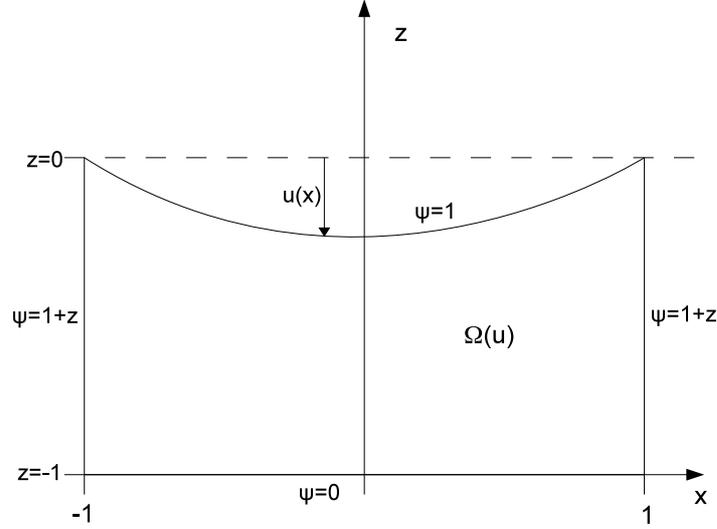


FIGURE 1. Idealized electrostatic MEMS device.

see (1.13) below. Indeed, letting $\varepsilon = 0$ (and applying a potential V with $V^2 \sim \varepsilon^3$ as suggested in [5]), one can solve (1.9)-(1.10) explicitly for the potential $\psi = \psi_0$, that is,

$$\psi_0(t, x, z) = \frac{1+z}{1+u_0(t, x)}, \quad (t, x, z) \in [0, \infty) \times I \times (-1, 0), \quad (1.12)$$

and the displacement $u = u_0$ satisfies

$$\begin{aligned} \partial_t u_0 - \partial_x^2 u_0 &= -\frac{\lambda}{(1+u_0)^2}, & x \in I, \quad t \in (0, \infty), \\ u_0(t, \pm 1) &= 0, & t \in (0, \infty), \\ u_0(0, x) &= u^0(x), & x \in I. \end{aligned} \quad (1.13)$$

In the limit $\varepsilon \rightarrow 0$, the free boundary problem is thus reduced to the singular semilinear heat equation (1.13) which has been studied thoroughly in recent years, see [7] for a survey as well as e.g. [8, 9, 13, 14, 15, 16, 18, 22, 25]. It is noteworthy to remark that the picture regarding pull-in voltage for the small aspect ratio model (1.13) is rather complete.

Let us point out that [6] is apparently the first mathematical analysis of the parabolic free boundary problem (1.7)-(1.11) while the corresponding elliptic (i.e. steady-state) free boundary problem is investigated in [20]. Moreover, we shall emphasize that the inclusion of non-small deformations is a feature of great physical relevance and, even though the results presented herein are reminiscent of the ones in [6], the quasilinear structure of (1.6) is by no means a trivial mathematical extension of (1.11).

1.3. Main Results. To state our findings on (1.6)-(1.10), we introduce the spaces

$$W_{q,D}^{2\alpha}(I) := \begin{cases} \{u \in W_q^{2\alpha}(I); u(\pm 1) = 0\} & \text{for } 2\alpha \in \left(\frac{1}{q}, 2\right], \\ W_q^{2\alpha}(I) & \text{for } 0 \leq 2\alpha < \frac{1}{q}. \end{cases} \quad (1.14)$$

We shall prove the following result regarding local and global existence of solutions:

Theorem 1.1 (Well-Posedness). *Let $q \in (2, \infty)$, $\varepsilon > 0$, and consider an initial value $u^0 \in W_{q,D}^2(I)$ such that $u^0(x) > -1$ for $x \in I$. Then, the following are true:*

- (i) *For each voltage value $\lambda > 0$, there is a unique maximal solution (u, ψ) to (1.6)-(1.10) on the maximal interval of existence $[0, T_m^\varepsilon)$ in the sense that*

$$u \in C^1([0, T_m^\varepsilon), L_q(I)) \cap C([0, T_m^\varepsilon), W_{q,D}^2(I))$$

satisfies (1.6)-(1.8) together with

$$u(t, x) > -1, \quad (t, x) \in [0, T_m^\varepsilon) \times I,$$

and $\psi(t) \in W_2^2(\Omega(u(t)))$ solves (1.9)-(1.10) on $\Omega(u(t))$ for each $t \in [0, T_m^\varepsilon)$.

- (ii) *If for each $\tau > 0$ there is $\kappa(\tau) \in (0, 1)$ such that $u(t) \geq -1 + \kappa(\tau)$ and $\|u(t)\|_{W_q^2(I)} \leq \kappa(\tau)^{-1}$ for $t \in [0, T_m^\varepsilon) \cap [0, \tau]$, then the solution exists globally, that is, $T_m^\varepsilon = \infty$.*
- (iii) *If $u^0(x) \leq 0$ for $x \in I$, then $u(t, x) \leq 0$ for $(t, x) \in [0, T_m^\varepsilon) \times I$. If $u^0 = u^0(x)$ is even with respect to $x \in I$, then, for all $t \in [0, T_m^\varepsilon)$, $u = u(t, x)$ and $\psi = \psi(t, x, z)$ are even with respect to $x \in I$ as well.*
- (iv) *Given $\kappa \in (0, 1)$, there are $\lambda_*(\kappa) > 0$ and $r(\kappa) > 0$ such that $T_m^\varepsilon = \infty$ with $u(t, x) \geq -1 + \kappa$ for $(t, x) \in [0, \infty) \times I$ provided that $\lambda \in (0, \lambda_*(\kappa))$ and $\|u^0\|_{W_q^2(I)} \leq r(\kappa)$. In that case, u enjoys the following additional regularity properties:*

$$u \in BUC^\rho([0, \infty), W_{q,D}^{2-\rho}(I)) \cap L_\infty([0, \infty), W_{q,D}^2(I))$$

for some $\rho > 0$ small.

Note that part (iv) of Theorem 1.1 provides uniform estimates on u in the $W_q^2(I)$ -norm and ensures that u never touches down on -1, not even in infinite time. In contrast to the semilinear case considered in [6], the global existence result for the quasilinear equation (1.6) requires initially a small deformation, see also Remark 3.3 below. The proof of Theorem 1.1 is the content of Section 3. It is based on interpreting (1.6) as a abstract quasilinear Cauchy problem which allows us to employ the powerful theory of evolution operators developed in [4]. Let us emphasize at this point that the regularity properties of the right-hand side of equation (1.6) established in [6] are not sufficient to handle the quasilinear character of the curvature operator and we consequently have to derive Lipschitz properties of the right-hand side of (1.6) in weaker topologies than in [6]. This is the purpose of Section 2.

Regarding existence and asymptotic stability of steady-state solutions to (1.6)-(1.10) we have a similar result as in [20, Thm. 1] and [6, Thm.1.3].

Theorem 1.2 (Asymptotically Stable Steady-State Solutions). *Let $q \in (2, \infty)$ and $\varepsilon > 0$.*

- (i) *Let $\kappa \in (0, 1)$. There are $\delta = \delta(\kappa) > 0$ and an analytic function*

$$[\lambda \mapsto U_\lambda] : [0, \delta) \rightarrow W_{q,D}^2(I)$$

such that $(U_\lambda, \Psi_\lambda)$ is the unique steady-state to (1.6)-(1.10) satisfying $\|U_\lambda\|_{W_{q,D}^2(I)} \leq 1/\kappa$ with $U_\lambda \geq -1 + \kappa$ and $\Psi_\lambda \in W_2^2(\Omega(U_\lambda))$ when $\lambda \in (0, \delta)$. The steady-state possesses the additional

regularity

$$\begin{aligned} U_\lambda &\in C^{2+\alpha}([-1, 1]), \\ \Psi_\lambda &\in W_2^2(\Omega(U_\lambda)) \cap C(\overline{\Omega(U_\lambda)}) \cap C^{2+\alpha}(\Omega(U_\lambda) \cup \Gamma(U_\lambda)), \end{aligned} \quad (1.15)$$

where $\alpha \in [0, 1)$ is arbitrary and $\Gamma(U_\lambda)$ denotes the boundary of $\Omega(U_\lambda)$ without corners. Moreover, U_λ is negative, convex, and even with $U_0 = 0$ and $\Psi_\lambda = \Psi_\lambda(x, z)$ is even with respect to $x \in I$.

- (ii) Let $\lambda \in (0, \delta)$. There are $\omega_0, m, R > 0$ such that for each initial value $u^0 \in W_{q,D}^2(I)$ with $\|u^0 - U_\lambda\|_{W_{q,D}^2} < m$, (1.6)-(1.10) has a unique global solution (u, ψ) with

$$u \in C^1([0, \infty), L_q(I)) \cap C([0, \infty), W_{q,D}^2(I)), \quad \psi(t) \in W_2^2(\Omega(u(t))), \quad t \geq 0,$$

and

$$u(t, x) > -1, \quad (t, x) \in [0, \infty) \times I.$$

Moreover,

$$\|u(t) - U_\lambda\|_{W_{q,D}^2(I)} + \|\partial_t u(t)\|_{L_q(I)} \leq Re^{-\omega_0 t} \|u^0 - U_\lambda\|_{W_{q,D}^2(I)}, \quad t \geq 0. \quad (1.16)$$

Part (ii) of Theorem 1.2 shows local exponential stability of the steady-states derived in part (i). We also point out that the potential ψ converges exponentially to Ψ_λ in the W_2^2 -norm as $t \rightarrow \infty$, see Remark 4.1 for a precise statement. The proof of Theorem 1.2 is given in Section 4 and relies on the Implicit Function Theorem for part (i) and the Principle of Linearized Stability for part (ii).

Clearly, Theorem 1.2 is just a local result with respect to λ values. However, we next show that there is an upper threshold for λ above which no steady-state solution exists. This is expected on physical grounds and is related to the ‘‘pull-in’’ instability already mentioned in the introduction.

Theorem 1.3 (Non-Existence of Steady-State Solutions). *Let $q \in (2, \infty)$ and $\varepsilon > 0$. There is $\bar{\lambda}(\varepsilon) > 0$ such that, if $\lambda \geq \bar{\lambda}(\varepsilon)$, then there is no steady-state solution (u, ψ) to (1.6)-(1.10) such that $u \in W_{q,D}^2(I)$, $\psi \in W_2^2(\Omega(u))$, and $u(x) > -1$ for $x \in I$. In addition, $\bar{\lambda}(\varepsilon) \rightarrow 2$ as $\varepsilon \rightarrow 0$.*

Similar results have already been obtained for related models, including the small aspect ratio model [5, 7] and for the stationary free boundary problem corresponding to (1.7)-(1.11), see [20]. The proof of Theorem 1.3 relies on a lower bound on $\partial_z \psi(x, u(x))$ established in the latter paper and is given in Section 5.

The final issue we address is the connection between the free boundary problem (1.6)-(1.10) and the small aspect ratio limit (1.13). More precisely, we show the following convergence result:

Theorem 1.4 (Small Aspect Ratio Limit). *Let $\lambda > 0$, $q \in (2, \infty)$, and let $u^0 \in W_{q,D}^2(I)$ with $-1 < u^0(x) \leq 0$ for $x \in I$. For $\varepsilon > 0$ we denote the unique solution to (1.6)-(1.10) on the maximal interval of existence $[0, T_m^\varepsilon)$ by $(u_\varepsilon, \psi_\varepsilon)$. There are $\tau > 0$, $\varepsilon_0 > 0$, and $\kappa_0 \in (0, 1)$ depending only on q and u^0 such that $T_m^\varepsilon \geq \tau$ with $u_\varepsilon(t) \geq -1 + \kappa_0$ and $\|u_\varepsilon(t)\|_{W_q^2(I)} \leq \kappa_0^{-1}$ for all $(t, \varepsilon) \in [0, \tau] \times (0, \varepsilon_0)$. Moreover, as $\varepsilon \rightarrow 0$,*

$$u_\varepsilon \longrightarrow u_0 \quad \text{in} \quad C^{1-\theta}([0, \tau], W_q^{2\theta}(I)), \quad 0 < \theta < 1,$$

and

$$\psi_\varepsilon(t) \mathbf{1}_{\Omega(u_\varepsilon(t))} \longrightarrow \psi_0(t) \mathbf{1}_{\Omega(u_0(t))} \quad \text{in} \quad L_2(I \times (-1, 0)), \quad t \in [0, \tau], \quad (1.17)$$

where

$$u_0 \in C^1([0, \tau], L_q(I)) \cap C([0, \tau], W_{q,D}^2(I))$$

is the unique solution to the small aspect ratio equation (1.13) and ψ_0 is the potential given in (1.12). Furthermore, there is $\Lambda(u^0) > 0$ such that the results above hold true for each $\tau > 0$ provided that $\lambda \in (0, \Lambda(u^0))$.

The proof is given in Section 6. A similar result has been established in [20, Thm. 2] for the stationary problem and in [6, Thm.1.4] for the semilinear parabolic version (1.11). As in the latter paper, the crucial step is to derive the ε -independent lower bound $\tau > 0$ on T_m^ε , which is not guaranteed by the analysis leading to Theorem 1.1. The proof of Theorem 1.4 uses several properties of (1.9)-(1.10) with respect to the ε -dependence shown in [6].

2. ON THE ELLIPTIC EQUATION (1.9)-(1.10)

We shall first derive properties of solutions to the elliptic equation (1.9)-(1.10) in dependence of a given (free) boundary. To do so, we transform the free boundary problem (1.9)-(1.10) to the fixed rectangle $\Omega := I \times (0, 1)$. More precisely, let $q \in (2, \infty)$ be fixed and consider an arbitrary function $v \in W_{q,D}^2(I)$ taking values in $(-1, \infty)$. We then define a diffeomorphism $T_v : \overline{\Omega(v)} \rightarrow \overline{\Omega}$ by setting

$$T_v(x, z) := \left(x, \frac{1+z}{1+v(x)} \right), \quad (x, z) \in \overline{\Omega(v)}, \quad (2.1)$$

with $\Omega(v) := \{(x, z) \in I \times (-1, \infty) ; 1 < z < v(x)\}$. Clearly, its inverse is

$$T_v^{-1}(x, \eta) = (x, (1+v(x))\eta - 1), \quad (x, \eta) \in \overline{\Omega}, \quad (2.2)$$

and the Laplace operator from (1.9) is transformed to the v -dependent differential operator

$$\begin{aligned} \mathcal{L}_v w := & \varepsilon^2 \partial_x^2 w - 2\varepsilon^2 \eta \frac{\partial_x v(x)}{1+v(x)} \partial_x \partial_\eta w + \frac{1 + \varepsilon^2 \eta^2 (\partial_x v(x))^2}{(1+v(x))^2} \partial_\eta^2 w \\ & + \varepsilon^2 \eta \left[2 \left(\frac{\partial_x v(x)}{1+v(x)} \right)^2 - \frac{\partial_x^2 v(x)}{1+v(x)} \right] \partial_\eta w. \end{aligned}$$

An alternative formulation of the boundary value problem (1.9)-(1.10) is then

$$(\mathcal{L}_{u(t)} \phi)(t, x, \eta) = 0, \quad (x, \eta) \in \Omega, \quad t > 0, \quad (2.3)$$

$$\phi(t, x, \eta) = \eta, \quad (x, \eta) \in \partial\Omega, \quad t > 0, \quad (2.4)$$

for $\phi = \psi \circ T_{u(t)}^{-1}$. With this notation, the quasilinear evolution equation (1.6) for u becomes

$$\partial_t u - \partial_x \left(\frac{\partial_x u}{\sqrt{1 + \varepsilon^2 (\partial_x u)^2}} \right) = -\lambda \left[\frac{1 + \varepsilon^2 (\partial_x u)^2}{(1+u)^2} \right] |\partial_\eta \phi(\cdot, 1)|^2, \quad x \in I, \quad t > 0, \quad (2.5)$$

where we have used $\partial_x \phi(t, x, 1) = 0$ for $x \in I$ and $t > 0$ due to $\phi(t, x, 1) = 1$ by (2.4). The investigation of the dynamics of (2.5) involves the properties of its nonlinear right hand side as well as the properties of the quasilinear curvature term. We shall see that these two features of (2.5) are somewhat opposite as the treatment of the former requires a functional analytic setting in $W_q^2(I)$ to handle the second order terms of $\mathcal{L}_{u(t)}$ in (2.3), while a slightly weaker setting has to be chosen to guarantee Hölder continuity of u with respect to time which is required in quasilinear evolution equations (see Remark 3.3 for further details). To account for these features of (2.5) we have to refine the Lipschitz property of the right-hand side of (2.5) derived in [6] as stated in (2.8) below.

Defining for $\kappa \in (0, 1)$ the open subset

$$S_q(\kappa) := \left\{ u \in W_{q,D}^2(I) ; \|u\|_{W_{q,D}^2(I)} < 1/\kappa \quad \text{and} \quad -1 + \kappa < u(x) \text{ for } x \in I \right\}$$

of $W_{q,D}^2(I)$ (defined in (1.14)) with closure

$$\overline{S}_q(\kappa) = \left\{ u \in W_{q,D}^2(I); \|u\|_{W_{q,D}^2(I)} \leq 1/\kappa \quad \text{and} \quad -1 + \kappa \leq u(x) \text{ for } x \in I \right\},$$

the crucial properties of the nonlinear right-hand side of (2.5) are collected in the following proposition:

Proposition 2.1. *Let $q \in (2, \infty)$, $\kappa \in (0, 1)$, and $\varepsilon > 0$. For each $v \in \overline{S}_q(\kappa)$ there is a unique solution $\phi_v \in W_2^2(\Omega)$ to*

$$(\mathcal{L}_v \phi_v)(x, \eta) = 0, \quad (x, \eta) \in \Omega, \quad (2.6)$$

$$\phi_v(x, \eta) = \eta, \quad (x, \eta) \in \partial\Omega. \quad (2.7)$$

If \tilde{v} is defined by $\tilde{v}(x) := v(-x)$ for $x \in I$, then $\phi_{\tilde{v}}(x, \eta) = \phi_v(-x, \eta)$ for $(x, \eta) \in \Omega$. Moreover, for $2\sigma \in [0, 1/2)$, the mapping

$$g_\varepsilon : S_q(\kappa) \longrightarrow W_{2,D}^{2\sigma}(I), \quad v \longmapsto \frac{1 + \varepsilon^2(\partial_x v)^2}{(1 + v)^2} |\partial_\eta \phi_v(\cdot, 1)|^2$$

is analytic and bounded with $g_\varepsilon(0) = 1$. Finally, if $\zeta \in [0, 1/2)$ and $v \in [0, (1 - 2\zeta)/2)$, then there exists a constant $c_1(\kappa, \varepsilon) > 0$ such that

$$\|g_\varepsilon(v) - g_\varepsilon(w)\|_{W_{2,D}^v(I)} \leq c_1(\kappa, \varepsilon) \|v - w\|_{W_{q,D}^{2-\zeta}(I)}, \quad v, w \in \overline{S}_q(\kappa). \quad (2.8)$$

According to [6, Prop. 2.1] we actually only have to prove (2.8). Notice that this global Lipschitz property is in the weaker topology of $W_{q,D}^{2-\zeta}(I)$ instead of $W_{q,D}^2(I)$ and improves [6, Prop. 2.1] where it was established for $\zeta = 0$. The property (2.8) will be a consequence of a sequence of lemmas. For the remainder of this section we fix $\varepsilon > 0$, $\kappa \in (0, 1)$, and $q \in (2, \infty)$.

In the following, if $\alpha > 1/2$ we let $W_{2,D}^\alpha(\Omega)$ denote the subspace of elements in $W_2^\alpha(\Omega)$ whose boundary trace is zero, and if $0 \leq \alpha < 1/2$ we set $W_{2,D}^\alpha(\Omega) := W_2^\alpha(\Omega)$. We equip $W_{2,D}^1(\Omega)$ with the norm

$$\|\Phi\|_{W_{2,D}^1(\Omega)} := \left(\|\partial_x \Phi\|_{L_2(\Omega)}^2 + \|\partial_\eta \Phi\|_{L_2(\Omega)}^2 \right)^{1/2},$$

and introduce the notation

$$W_{2,D}^{-\theta}(\Omega) := (W_{2,D}^\theta(\Omega))', \quad 0 \leq \theta \leq 1.$$

Lemma 2.2. *For each $v \in \overline{S}_q(\kappa)$ and $F \in W_{2,D}^{-1}(\Omega)$ there is a unique solution $\Phi \in W_{2,D}^1(\Omega)$ to the boundary value problem*

$$-\mathcal{L}_v \Phi = F \quad \text{in } \Omega, \quad (2.9)$$

$$\Phi = 0 \quad \text{on } \partial\Omega, \quad (2.10)$$

and there is a constant $c_2(\kappa, \varepsilon) > 0$ such that

$$\|\Phi\|_{W_{2,D}^1(\Omega)} \leq c_2(\kappa, \varepsilon) \|F\|_{W_{2,D}^{-1}(\Omega)}. \quad (2.11)$$

Furthermore, if $F \in L_2(\Omega)$, then $\Phi \in W_{2,D}^2(\Omega)$ and

$$\|\Phi\|_{W_{2,D}^2(\Omega)} \leq c_2(\kappa, \varepsilon) \|F\|_{L_2(\Omega)}. \quad (2.12)$$

Proof. According to [12, Def. 1.3.2.3, Eq. (1,3,2,3)], we may write any $F \in W_{2,D}^{-1}(\Omega)$ in the form $F = f_0 + \partial_x f_1 + \partial_\eta f_2$ with $(f_0, f_1, f_2) \in L_2(\Omega)^3$. Consequently, [10, Thm. 8.3] ensures that the

boundary value problem (2.9)-(2.10) has a unique solution $\Phi \in W_{2,D}^1(\Omega)$. Furthermore, taking Φ as a test function in the weak formulation of (2.9)-(2.10) gives

$$\begin{aligned} \langle F, \Phi \rangle &= \int_{\Omega} \left[\varepsilon^2 (\partial_x \Phi)^2 - 2 \varepsilon^2 \eta \frac{\partial_x v}{1+v} \partial_x \Phi \partial_{\eta} \Phi - \varepsilon^2 \eta \partial_x \left(\frac{\partial_x v}{1+v} \right) \Phi \partial_{\eta} \Phi - \varepsilon^2 \frac{\partial_x v}{1+v} \Phi \partial_x \Phi \right] d(x, \eta) \\ &\quad + \int_{\Omega} \left[\frac{1 + \varepsilon^2 \eta^2 (\partial_x v)^2}{(1+v)^2} (\partial_{\eta} \Phi)^2 + 2 \varepsilon^2 \eta \left(\frac{\partial_x v}{1+v} \right)^2 \Phi \partial_{\eta} \Phi \right] d(x, \eta) \\ &\quad + \int_{\Omega} \varepsilon^2 \eta \left[\partial_x \left(\frac{\partial_x v}{1+v} \right) - \left(\frac{\partial_x v}{1+v} \right)^2 \right] \Phi \partial_{\eta} \Phi d(x, \eta) \\ &= \int_{\Omega} \left[\varepsilon^2 (\partial_x \Phi)^2 - 2 \varepsilon^2 \eta \frac{\partial_x v}{1+v} \partial_x \Phi \partial_{\eta} \Phi + \frac{1 + \varepsilon^2 \eta^2 (\partial_x v)^2}{(1+v)^2} (\partial_{\eta} \Phi)^2 \right] d(x, \eta) \\ &\quad - \int_{\Omega} \varepsilon^2 \frac{\partial_x v}{1+v} \left[\partial_x \Phi - \eta \frac{\partial_x v}{1+v} \partial_{\eta} \Phi \right] \Phi d(x, \eta) \end{aligned}$$

and thus

$$\begin{aligned} &\int_{\Omega} \left[\varepsilon^2 \left(\partial_x \Phi - \eta \frac{\partial_x v}{1+v} \partial_{\eta} \Phi \right)^2 + \left(\frac{\partial_{\eta} \Phi}{1+v} \right)^2 \right] d(x, \eta) \\ &\leq \varepsilon^2 \left\| \frac{\partial_x v}{1+v} \right\|_{L^{\infty}(I)} \left\| \partial_x \Phi - \eta \frac{\partial_x v}{1+v} \partial_{\eta} \Phi \right\|_{L_2(\Omega)} \|\Phi\|_{L_2(\Omega)} + \|F\|_{W_{2,D}^{-1}(\Omega)} \|\Phi\|_{W_{2,D}^1(\Omega)}. \end{aligned} \quad (2.13)$$

Note then that by definition of $\bar{S}_q(\kappa)$ and Sobolev's embedding, there is $c_3 > 0$ such that

$$1 + v(x) \geq \kappa, \quad x \in I, \quad \|v\|_{C^1([-1,1])} \leq \frac{c_3}{\kappa} \quad (2.14)$$

for all $v \in \bar{S}_q(\kappa)$. Also, if $\zeta = (\zeta_1, \zeta_2) \in \mathbb{R}^2$, Young's inequality ensures that, for $(x, \eta) \in \Omega$,

$$\begin{aligned} \varepsilon^2 \zeta_1^2 &\leq 2 \varepsilon^2 \left(\zeta_1 - \eta \frac{\partial_x v(x)}{1+v(x)} \zeta_2 \right)^2 + 2 \varepsilon^2 \eta^2 \left(\frac{\partial_x v(x)}{1+v(x)} \right)^2 \zeta_2^2 \\ &\leq 2 \left(1 + 2 \varepsilon^2 \|\partial_x v\|_{\infty}^2 \right) \left(\varepsilon^2 \left(\zeta_1 - \eta \frac{\partial_x v(x)}{1+v(x)} \zeta_2 \right)^2 + \frac{1}{2} \left(\frac{\zeta_2}{1+v(x)} \right)^2 \right). \end{aligned}$$

Therefore, introducing

$$v(\kappa, \varepsilon) := \frac{1}{2} \min \left\{ \frac{\varepsilon^2 \kappa^2}{\kappa^2 + 2 \varepsilon^2 c_3^2}, \frac{\kappa^2}{(\kappa + c_3)^2} \right\},$$

we infer from (2.14) that

$$v(\kappa, \varepsilon) (\zeta_1^2 + \zeta_2^2) \leq \varepsilon^2 \left(\zeta_1 - \eta \frac{\partial_x v(x)}{1+v(x)} \zeta_2 \right)^2 + \left(\frac{\zeta_2}{1+v(x)} \right)^2. \quad (2.15)$$

Consequently, (2.13), (2.14), and (2.15) give

$$\begin{aligned} &\varepsilon^2 \left\| \partial_x \Phi - \eta \frac{\partial_x v}{1+v} \partial_{\eta} \Phi \right\|_{L_2(\Omega)}^2 + \left\| \frac{\partial_{\eta} \Phi}{1+v} \right\|_{L_2(\Omega)}^2 \\ &\leq \varepsilon^2 \frac{c_3}{\kappa^2} \left\| \partial_x \Phi - \eta \frac{\partial_x v}{1+v} \partial_{\eta} \Phi \right\|_{L_2(\Omega)} \|\Phi\|_{L_2(\Omega)} \\ &\quad + \frac{\|F\|_{W_{2,D}^{-1}(\Omega)}}{\sqrt{v(\kappa, \varepsilon)}} \left(\varepsilon^2 \left\| \partial_x \Phi - \eta \frac{\partial_x v}{1+v} \partial_{\eta} \Phi \right\|_{L_2(\Omega)}^2 + \left\| \frac{\partial_{\eta} \Phi}{1+v} \right\|_{L_2(\Omega)}^2 \right)^{1/2}, \end{aligned}$$

whence, using again (2.15),

$$\|\Phi\|_{W_{2,D}^1(\Omega)} \leq \frac{\varepsilon}{\sqrt{\nu(\kappa, \varepsilon)}} \frac{c_3}{\kappa^2} \|\Phi\|_{L_2(\Omega)} + \frac{\|F\|_{W_{2,D}^{-1}(\Omega)}}{\nu(\kappa, \varepsilon)}. \quad (2.16)$$

We now proceed as in [10, Lem. 9.17] and argue by contradiction to show (2.11) (see also [6, Lem. 6.2] for a similar argument in a slightly different functional setting). The last statement of Lemma 2.2 is proved in [6, Lem. 6.2]. \square

Now, introducing

$$f_v(x, \eta) := \mathcal{L}_v \eta = \varepsilon^2 \eta \left[2 \left(\frac{\partial_x v(x)}{1+v(x)} \right)^2 - \frac{\partial_x^2 v(x)}{1+v(x)} \right], \quad (x, \eta) \in \Omega, \quad (2.17)$$

for $v \in \overline{S}_q(\kappa)$ given, we readily deduce that $f_v \in L_2(\Omega)$ with

$$\|f_v\|_{L_2(\Omega)} \leq c_4(\kappa, \varepsilon). \quad (2.18)$$

Consequently, Lemma 2.2 provides a unique solution $\Phi_v \in W_{2,D}^2(\Omega)$ to

$$\begin{aligned} -(\mathcal{L}_v \Phi_v)(x, \eta) &= f_v, \quad (x, \eta) \in \Omega, \\ \Phi_v(x, \eta) &= 0, \quad (x, \eta) \in \partial\Omega. \end{aligned}$$

Clearly, defining

$$\phi_v(x, \eta) := \Phi_v(x, \eta) + \eta, \quad (x, \eta) \in \overline{\Omega}, \quad (2.19)$$

gives then the unique solution $\phi_v \in W_2^2(\Omega)$ to (2.6)-(2.7). To prove a Lipschitz dependence of ϕ_v on $v \in \overline{S}_q(\kappa)$, we introduce a bounded linear operator

$$\mathcal{A}(v) \in \mathcal{L}(W_{2,D}^1(\Omega), W_{2,D}^{-1}(\Omega)) \cap \mathcal{L}(W_{2,D}^2(\Omega), L_2(\Omega))$$

by setting

$$\mathcal{A}(v)\Phi := -\mathcal{L}_v \Phi, \quad \Phi \in W_{2,D}^1(\Omega).$$

Note that $\mathcal{A}(v)$ is invertible according to Lemma 2.2 and that $\Phi_v = \mathcal{A}(v)^{-1} f_v$. For the inverse $\mathcal{A}(v)^{-1}$ we have:

Lemma 2.3. *Given $\theta \in [0, 1] \setminus \{1/2\}$, there is a constant $c_5(\kappa, \varepsilon) > 0$ such that*

$$\|\mathcal{A}(v)^{-1}\|_{\mathcal{L}(W_{2,D}^{\theta-1}(\Omega), W_{2,D}^{\theta+1}(\Omega))} \leq c_5(\kappa, \varepsilon), \quad v \in \overline{S}_q(\kappa).$$

Proof. Due to Lemma 2.2, $\mathcal{A}(v)^{-1}$ belongs for each $v \in \overline{S}_q(\kappa)$ to both $\mathcal{L}(W_{2,D}^{-1}(\Omega), W_{2,D}^1(\Omega))$ and $\mathcal{L}(L_2(\Omega), W_{2,D}^2(\Omega))$ with

$$\|\mathcal{A}(v)^{-1}\|_{\mathcal{L}(W_{2,D}^{-1}(\Omega), W_{2,D}^1(\Omega))} + \|\mathcal{A}(v)^{-1}\|_{\mathcal{L}(L_2(\Omega), W_{2,D}^2(\Omega))} \leq 2c_2(\kappa, \varepsilon).$$

Hence, using complex interpolation, we derive for $\theta \in [0, 1]$,

$$\|\mathcal{A}(v)^{-1}\|_{\mathcal{L}([W_{2,D}^{-1}(\Omega), L_2(\Omega)]_\theta, [W_{2,D}^1(\Omega), W_{2,D}^2(\Omega)]_\theta)} \leq 2c_2(\kappa, \varepsilon),$$

and it then remains to characterize the interpolation spaces for $\theta \in [0, 1] \setminus \{1/2\}$. For this we first invoke [28, Thm. 1.1.11] to obtain that

$$[W_{2,D}^{-1}(\Omega), L_2(\Omega)]_\theta = [(W_{2,D}^1(\Omega))', (L_2(\Omega))']_\theta = [W_{2,D}^1(\Omega), L_2(\Omega)]'_\theta$$

with equivalent norms and hence, using [12, Cor. 1.4.4.5] and [19, Prop. 2.11 & Prop. 3.3] to characterize the last interpolation space,

$$[W_{2,D}^{-1}(\Omega), L_2(\Omega)]_\theta = (W_{2,D}^{1-\theta}(\Omega))' = W_{2,D}^{\theta-1}(\Omega).$$

Finally, since the embedding

$$[W_{2,D}^1(\Omega), W_{2,D}^2(\Omega)]_\theta \hookrightarrow W_{2,D}^{\theta+1}(\Omega)$$

is obviously continuous by [28, Thm. 4.3.2/2] and definition of interpolation, the assertion follows. \square

Next, we show that $\mathcal{A}(v)$ is Lipschitz continuous with respect to v in a suitable topology. More precisely:

Lemma 2.4. *Given $\xi \in [0, (q-1)/q]$ and $\alpha \in (\xi, 1)$, there exists $c_6(\kappa, \varepsilon) > 0$ such that*

$$\|\mathcal{A}(v) - \mathcal{A}(w)\|_{\mathcal{L}(W_{2,D}^2(\Omega), W_{2,D}^{-\alpha}(\Omega))} \leq c_6(\kappa, \varepsilon) \|v - w\|_{W_q^{2-\xi}(I)}, \quad v, w \in \bar{S}_q(\kappa).$$

Proof. Consider $v, w \in \bar{S}_q(\kappa)$ and $\Phi \in W_{2,D}^2(\Omega)$. Then $\mathcal{A}(v)\Phi \in L_2(\Omega) \hookrightarrow W_{2,D}^{-\alpha}(\Omega)$ and so, for any $\varphi \in W_{2,D}^\alpha(\Omega)$,

$$\begin{aligned} \int_{\Omega} [(\mathcal{A}(v) - \mathcal{A}(w)) \Phi] \varphi \, d(x, \eta) &= 2\varepsilon^2 \int_{\Omega} \eta \left(\frac{\partial_x v}{1+v} - \frac{\partial_x w}{1+w} \right) \partial_x \partial_\eta \Phi \varphi \, d(x, \eta) \\ &\quad - \int_{\Omega} \left(\frac{1 + \varepsilon^2 \eta^2 (\partial_x v)^2}{(1+v)^2} - \frac{1 + \varepsilon^2 \eta^2 (\partial_x w)^2}{(1+w)^2} \right) \partial_\eta^2 \Phi \varphi \, d(x, \eta) \\ &\quad - 2\varepsilon^2 \int_{\Omega} \eta \left[\left(\frac{\partial_x v}{1+v} \right)^2 - \left(\frac{\partial_x w}{1+w} \right)^2 \right] \partial_\eta \Phi \varphi \, d(x, \eta) \\ &\quad + \varepsilon^2 \int_{\Omega} \eta \left(\frac{\partial_x^2 v}{1+v} - \frac{\partial_x^2 w}{1+w} \right) \partial_\eta \Phi \varphi \, d(x, \eta) \\ &=: J_1 + J_2 + J_3 + J_4. \end{aligned}$$

Since $\xi \in [0, (q-1)/q]$ it follows from the continuous embedding $W_q^{2-\xi}(I) \hookrightarrow W_\infty^1(I)$ that

$$\begin{aligned} |J_1| &\leq 2\varepsilon^2 \left\| \frac{\partial_x v}{1+v} - \frac{\partial_x w}{1+w} \right\|_{L_\infty(I)} \|\partial_x \partial_\eta \Phi\|_{L_2(\Omega)} \|\varphi\|_{L_2(\Omega)} \\ &\leq \varepsilon^2 c(\kappa) \|v - w\|_{W_q^{2-\xi}(I)} \|\Phi\|_{W_{2,D}^2(\Omega)} \|\varphi\|_{W_{2,D}^\alpha(\Omega)}. \end{aligned}$$

Similarly, we have

$$|J_2| + |J_3| \leq (1 + \varepsilon^2) c(\kappa) \|v - w\|_{W_q^{2-\xi}(I)} \|\Phi\|_{W_{2,D}^2(\Omega)} \|\varphi\|_{W_{2,D}^\alpha(\Omega)}.$$

Finally, we infer from the embedding $W_2^1(\Omega) \hookrightarrow L_{2q/(q-2)}(\Omega)$ and the fact that $(W_{q'}^\xi(I))' = W_q^{-\xi}(I)$ since $\xi \in [0, (q-1)/q]$, see, e.g., [3, Eq. (5.14)], that

$$\begin{aligned} |J_4| &\leq \varepsilon^2 \left| \int_{\Omega} \partial_x^2 (v - w) \frac{\eta \partial_\eta \Phi \varphi}{1+v} \, d(x, \eta) \right| \\ &\quad + \varepsilon^2 \|\partial_x^2 w\|_{L_q(I)} \left\| \frac{1}{1+v} - \frac{1}{1+w} \right\|_{L_\infty(I)} \|\partial_\eta \Phi\|_{L_{2q/(q-2)}(\Omega)} \|\varphi\|_{L_2(\Omega)} \\ &\leq \varepsilon^2 \|\partial_x^2 (v - w)\|_{W_q^{-\xi}(I)} \left\| \frac{1}{1+v} \int_0^1 \eta \partial_\eta \Phi(\cdot, \eta) \varphi(\cdot, \eta) \, d\eta \right\|_{W_{q'}^\xi(I)} \\ &\quad + \varepsilon^2 c(\kappa) \|v - w\|_{W_q^{2-\xi}(I)} \|\partial_\eta \Phi\|_{W_{2,D}^1(\Omega)} \|\varphi\|_{W_{2,D}^\alpha(\Omega)}. \end{aligned}$$

Now, choosing $\beta \in [\xi, \alpha)$, pointwise multiplication is continuous as mappings

$$W_q^2(I) \cdot W_2^\beta(I) \hookrightarrow W_{q'}^\xi(I), \quad C^2(\bar{\Omega}) \cdot W_2^1(\Omega) \cdot W_2^\alpha(\Omega) \hookrightarrow W_2^\beta(\Omega)$$

according to Theorem 7.1, while an interpolation argument shows

$$\left\| \int_0^1 \zeta(\cdot, \eta) \, d\eta \right\|_{W_2^\gamma(I)} \leq c \|\zeta\|_{W_2^\gamma(\Omega)}, \quad \zeta \in W_2^\gamma(\Omega), \quad (2.20)$$

for all $\gamma \in [0, 1]$, since $\zeta \mapsto \int_0^1 \zeta(\cdot, \eta) \, d\eta$ belongs to both $\mathcal{L}(L_2(\Omega), L_2(I))$ and $\mathcal{L}(W_2^1(\Omega), W_2^1(I))$. Therefore, we deduce that

$$\begin{aligned} \left\| \frac{1}{1+v} \int_0^1 \eta \partial_\eta \Phi(\cdot, \eta) \varphi(\cdot, \eta) \, d\eta \right\|_{W_q^\xi(I)} &\leq \left\| \frac{1}{1+v} \right\|_{W_q^2(I)} \left\| \int_0^1 \eta \partial_\eta \Phi(\cdot, \eta) \varphi(\cdot, \eta) \, d\eta \right\|_{W_2^\beta(I)} \\ &\leq c(\kappa) \|\eta \partial_\eta \Phi \varphi\|_{W_{2,D}^\beta(\Omega)} \\ &\leq c(\kappa) \|\partial_\eta \Phi\|_{W_{2,D}^1(\Omega)} \|\varphi\|_{W_{2,D}^\alpha(\Omega)}, \end{aligned}$$

and we arrive at

$$|J_4| \leq \varepsilon^2 c(\kappa) \|v - w\|_{W_q^{2-\xi}(I)} \|\Phi\|_{W_{2,D}^2(\Omega)} \|\varphi\|_{W_{2,D}^\alpha(\Omega)}.$$

Consequently, gathering the above estimates on $|J_i|$, we obtain

$$\|(\mathcal{A}(v) - \mathcal{A}(w))\Phi\|_{W_{2,D}^{-\alpha}(\Omega)} \leq (1 + \varepsilon^2) c(\kappa) \|v - w\|_{W_q^{2-\xi}(I)} \|\Phi\|_{W_{2,D}^2(\Omega)},$$

whence the claim. \square

Next, we prove that f_v from (2.17) depends Lipschitz continuously on $v \in \overline{S}_q(\kappa)$.

Lemma 2.5. *Given $\xi \in [0, (q-1)/q]$ and $\alpha \in (\xi, 1)$, there exists $c_7(\kappa) > 0$ such that*

$$\|f_v - f_w\|_{W_{2,D}^{-\alpha}(\Omega)} \leq c_7(\kappa) \varepsilon^2 \|v - w\|_{W_q^{2-\xi}(I)}, \quad v, w \in \overline{S}_q(\kappa).$$

Proof. Consider $v, w \in \overline{S}_q(\kappa)$ and recall that $f_v \in L_2(\Omega) \hookrightarrow W_{2,D}^{-\alpha}(\Omega)$ by (2.18). Then, it follows from the embedding $W_2^{2-\xi}(I) \hookrightarrow W_\infty^1(I)$ and (2.14) that, for $\varphi \in W_{2,D}^\alpha(\Omega)$,

$$\begin{aligned} &\left| \int_\Omega (f_v - f_w) \varphi \, d(x, \eta) \right| \\ &\leq 2\varepsilon^2 \int_\Omega \left| \left(\frac{\partial_x v}{1+v} \right)^2 - \left(\frac{\partial_x w}{1+w} \right)^2 \right| |\varphi| \, d(x, \eta) + \varepsilon^2 \int_\Omega |\partial_x^2 v| \left| \frac{1}{1+v} - \frac{1}{1+w} \right| |\varphi| \, d(x, \eta) \\ &\quad + \varepsilon^2 \left| \int_\Omega (\partial_x^2 v - \partial_x^2 w) \frac{\varphi}{1+w} \, d(x, \eta) \right| \\ &\leq \varepsilon^2 c(\kappa) \|v - w\|_{W_\infty^1(I)} \|\varphi\|_{L_2(\Omega)} + \varepsilon^2 c(\kappa) \|\partial_x^2 v\|_{L_2(I)} \|v - w\|_{L_\infty(I)} \|\varphi\|_{L_2(\Omega)} \\ &\quad + \varepsilon^2 \|\partial_x^2(v - w)\|_{W_q^{-\xi}(I)} \left\| \frac{1}{1+w} \int_0^1 \varphi(\cdot, \eta) \, d\eta \right\|_{W_q^\xi(I)} \\ &\leq \varepsilon^2 c(\kappa) \|v - w\|_{W_q^{2-\xi}(I)} \left(\|\varphi\|_{W_{2,D}^\alpha(\Omega)} + \left\| \frac{1}{1+w} \right\|_{W_q^2(I)} \left\| \int_0^1 \varphi(\cdot, \eta) \, d\eta \right\|_{W_2^\alpha(I)} \right), \end{aligned}$$

where we have used for the last inequality continuity of pointwise multiplication

$$W_q^2(I) \cdot W_2^\alpha(I) \hookrightarrow W_q^\xi(I)$$

as in the proof of Lemma 2.4. Taking $\gamma = \alpha$ in (2.20) and using (2.14), we end up with

$$\left| \int_\Omega (f_v - f_w) \varphi \, d(x, \eta) \right| \leq \varepsilon^2 c(\kappa) \|v - w\|_{W_q^{2-\xi}(I)} \|\varphi\|_{W_{2,D}^\alpha(\Omega)},$$

which yields the assertion. \square

Combining the previous lemmas we now readily obtain the Lipschitz continuity of ϕ_v with respect to $v \in \bar{S}_q(\kappa)$.

Lemma 2.6. *Given $\xi \in [0, (q-1)/q]$ and $\alpha \in (\xi, 1)$, there exists $c_8(\kappa, \varepsilon) > 0$ such that*

$$\|\phi_v - \phi_w\|_{W_{2,D}^{2-\alpha}(\Omega)} \leq c_8(\kappa, \varepsilon) \|v - w\|_{W_q^{2-\xi}(I)}, \quad v, w \in \bar{S}_q(\kappa).$$

Proof. Let $v, w \in \bar{S}_q(\kappa)$ and recall that $\phi_v = \Phi_v + \eta$ with $\Phi_v = \mathcal{A}(v)^{-1}f_v \in W_{2,D}^2(\Omega)$. Then obviously $\phi_v - \phi_w = \Phi_v - \Phi_w \in W_{2,D}^{2-\alpha}(\Omega)$ and so

$$\begin{aligned} \|\phi_v - \phi_w\|_{W_{2,D}^{2-\alpha}(\Omega)} &= \|\mathcal{A}(v)^{-1}(f_v - f_w) + (\mathcal{A}(v)^{-1} - \mathcal{A}(w)^{-1})f_w\|_{W_{2,D}^{2-\alpha}(\Omega)} \\ &\leq \|\mathcal{A}(v)^{-1}\|_{\mathcal{L}(W_{2,D}^{-\alpha}(\Omega), W_{2,D}^{2-\alpha}(\Omega))} \|f_v - f_w\|_{W_{2,D}^{-\alpha}(\Omega)} \\ &\quad + \|\mathcal{A}(v)^{-1}\|_{\mathcal{L}(W_{2,D}^{-\alpha}(\Omega), W_{2,D}^{2-\alpha}(\Omega))} \|\mathcal{A}(v) - \mathcal{A}(w)\|_{\mathcal{L}(W_{2,D}^2(\Omega), W_{2,D}^{-\alpha}(\Omega))} \\ &\quad \times \|\mathcal{A}(w)^{-1}\|_{\mathcal{L}(L_2(\Omega), W_{2,D}^2(\Omega))} \|f_w\|_{L_2(\Omega)} \end{aligned}$$

and the statement follows from (2.18), Lemma 2.3, Lemma 2.4, and Lemma 2.5. \square

Proof of Proposition 2.1. As previously noted, we are left to prove assertion (2.8). For this let $q \in (2, \infty)$, $\xi \in [0, 1/2)$, and $\nu \in [0, (1-2\xi)/2)$. By [6, Eq. (2.30)], given $2\sigma \in [0, 1/2)$ there is $c_9(\kappa, \varepsilon)$ such that

$$\|\partial_\eta \phi_v(\cdot, 1)\|_{W_2^{1/2}(I)} + \|\partial_\eta \phi_v(\cdot, 1)\|_{W_2^{2\sigma}(I)} \leq c_9(\kappa, \varepsilon), \quad v \in \bar{S}_q(\kappa). \quad (2.21)$$

Given $v, w \in \bar{S}_q(\kappa)$, it follows from the definition of g_ε that

$$\begin{aligned} &\|g_\varepsilon(v) - g_\varepsilon(w)\|_{W_2^\nu(I)} \\ &\leq \varepsilon^2 \left\| \frac{\partial_x v + \partial_x w}{(1+v)^2} (\partial_x v - \partial_x w) |\partial_\eta \phi_v(\cdot, 1)|^2 \right\|_{W_2^\nu(I)} \\ &\quad + \left\| \left(1 + \varepsilon^2 (\partial_x w)^2\right) \frac{2+v+w}{(1+v)^2(1+w)^2} (v-w) |\partial_\eta \phi_v(\cdot, 1)|^2 \right\|_{W_2^\nu(I)} \\ &\quad + \left\| \frac{1 + \varepsilon^2 (\partial_x w)^2}{(1+w)^2} (\partial_\eta \phi_v(\cdot, 1) + \partial_\eta \phi_w(\cdot, 1)) (\partial_\eta \phi_v(\cdot, 1) - \partial_\eta \phi_w(\cdot, 1)) \right\|_{W_2^\nu(I)} \\ &=: J_1 + J_2 + J_3. \end{aligned}$$

We now fix $2\sigma \in (\xi + \nu, 1/2)$ and $s \in [\nu, 1-\xi)$ with $s \geq 1/q$ so that pointwise multiplication

$$W_q^s(I) \cdot W_2^{2\sigma}(I) \hookrightarrow W_2^\nu(I) \quad (2.22)$$

is continuous by Theorem 7.1. Then,

$$J_1 \leq c \varepsilon^2 \left\| \frac{\partial_x v + \partial_x w}{(1+v)^2} (\partial_x v - \partial_x w) \right\|_{W_q^s(I)} \left\| |\partial_\eta \phi_v(\cdot, 1)|^2 \right\|_{W_2^{2\sigma}(I)}.$$

We infer from (2.21) and continuity of pointwise multiplication

$$W_q^2(I) \cdot W_q^1(I) \cdot W_q^{1-\xi}(I) \hookrightarrow W_q^s(I) \quad (2.23)$$

guaranteed by Theorem 7.1 that

$$J_1 \leq c(\kappa, \varepsilon) \left\| \frac{1}{(1+v)^2} \right\|_{W_q^2(I)} \|\partial_x v + \partial_x w\|_{W_q^1(I)} \|\partial_x v - \partial_x w\|_{W_q^{1-\xi}(I)},$$

hence, since both v and w belong to $\overline{S}_q(\kappa)$ and $W_q^2(I)$ is an algebra,

$$J_1 \leq c(\kappa, \varepsilon) \|v - w\|_{W_q^{2-\xi}(I)}. \quad (2.24)$$

Using again the continuity (2.22) and (2.23) of pointwise multiplication as well as (2.21), we obtain

$$J_2 \leq c(\kappa, \varepsilon) \left\| 1 + \varepsilon^2 (\partial_x w)^2 \right\|_{W_q^1(I)} \left\| \frac{2 + v + w}{(1 + v)^2 (1 + w)^2} \right\|_{W_q^2(I)} \|v - w\|_{W_q^{1-\xi}(I)}.$$

As $W_q^1(I)$ and $W_q^2(I)$ are algebras, we deduce from the properties of $\overline{S}_q(\kappa)$ that

$$J_2 \leq c(\kappa, \varepsilon) \|v - w\|_{W_q^{2-\xi}(I)}. \quad (2.25)$$

Finally, fix $\alpha \in (\xi, 2\sigma - \nu)$. Invoking once more Theorem 7.1 gives continuity of pointwise multiplication

$$W_q^1(I) \cdot W_2^{1/2}(I) \cdot W_2^{1/2-\alpha}(I) \hookrightarrow W_2^\nu(I), \quad (2.26)$$

and thus

$$J_3 \leq \left\| \frac{1 + \varepsilon^2 (\partial_x w)^2}{(1 + w)^2} \right\|_{W_q^1(I)} \|\partial_\eta \phi_v(\cdot, 1) + \partial_\eta \phi_w(\cdot, 1)\|_{W_2^{1/2}(I)} \|\partial_\eta \phi_v(\cdot, 1) - \partial_\eta \phi_w(\cdot, 1)\|_{W_2^{1/2-\alpha}(I)}.$$

Since $w \in \overline{S}_q(\kappa)$, it follows from (2.21), the properties of the trace operator [12, Thm. 1.5.1.1], and Lemma 2.6 that

$$J_3 \leq c(\kappa, \varepsilon) \|\phi_v - \phi_w\|_{W_{2,D}^{2-\alpha}(\Omega)} \leq c(\kappa, \varepsilon) \|v - w\|_{W_q^{2-\xi}(I)}. \quad (2.27)$$

Gathering (2.24), (2.25), and (2.27) completes the proof of Proposition 2.1. \square

3. WELL-POSEDNESS OF THE EVOLUTION PROBLEM: PROOF OF THEOREM 1.1

We next focus our attention on the quasilinear curvature part of equation (1.6). Let us first note that

$$\partial_x \left(\frac{\partial_x u}{\sqrt{1 + \varepsilon^2 (\partial_x u)^2}} \right) = \frac{1}{(1 + \varepsilon^2 (\partial_x u)^2)^{3/2}} \partial_x^2 u, \quad (3.1)$$

which motivates the definition of

$$A(w)v := -\frac{1}{(1 + (\partial_x w)^2)^{3/2}} \partial_x^2 v, \quad v \in W_{q,D}^2(I), \quad (3.2)$$

where w belongs, for $q \in (2, \infty)$, $\kappa \in (0, 1)$, and $\xi \in (0, (q-1)/q)$ given, to the set

$$Z_q(\kappa) := \{w \in W_q^{2-\xi}(I); \|w\|_{W_q^{2-\xi}(I)} \leq 1/\kappa\}.$$

Note that $W_q^{2-\xi}(I) \hookrightarrow C^1([-1, 1])$. Also observe that $-A(\varepsilon u)u$ coincides with (3.1). For $\omega > 0$ and $k \geq 1$, let $\mathcal{H}(W_{q,D}^2(I), L_q(I); k, \omega)$ be the set of all $\mathcal{A} \in \mathcal{L}(W_{q,D}^2(I), L_q(I))$ such that $\omega + \mathcal{A}$ is an isomorphism from $W_{q,D}^2(I)$ onto $L_q(I)$ and satisfies the resolvent estimates

$$\frac{1}{k} \leq \frac{\|(\mu + \mathcal{A})z\|_{L_q(I)}}{|\mu| \|z\|_{L_q(I)} + \|z\|_{W_{q,D}^2(I)}} \leq k, \quad \operatorname{Re} \mu \geq \omega, \quad z \in W_{q,D}^2(I) \setminus \{0\}.$$

Then $\mathcal{A} \in \mathcal{H}(W_{q,D}^2(I), L_q(I); k, \omega)$ implies that $\mathcal{A} \in \mathcal{H}(W_{q,D}^2(I), L_q(I))$, that is, $-\mathcal{A}$ generates an analytic semigroup on $L_q(I)$ with domain $W_{q,D}^2(I)$, see [4, I.Thm.1.2.2].

Lemma 3.1. *Let $q \in (2, \infty)$, $\kappa \in (0, 1)$, and $\xi \in (0, (q-1)/q)$. Then there are $k := k(\kappa) \geq 1$ and $\omega := \omega(\kappa) > 0$ such that, for each $w \in Z_q(\kappa)$,*

$$-2\omega + A(w) \in \mathcal{H}(W_{q,D}^2(I), L_q(I); k, \omega)$$

and $A(w)$ is resolvent positive. Moreover, there is a constant $\ell(\kappa) > 0$ such that

$$\|A(w_1) - A(w_2)\|_{\mathcal{L}(W_{q,D}^2(I), L_q(I))} \leq \ell(\kappa) \|w_1 - w_2\|_{W_q^{2-\xi}(I)}, \quad w_1, w_2 \in Z_q(\kappa). \quad (3.3)$$

Proof. Note that the continuous embedding $W_q^{2-\xi}(I) \hookrightarrow C^1([-1, 1])$ ensures the existence of $c(\kappa) > 1$ such that

$$1 \leq (1 + (\partial_x w)^2)^{3/2} < c(\kappa), \quad w \in Z_q(\kappa). \quad (3.4)$$

Let $w \in Z_q(\kappa)$ be fixed and put $W := (1 + (\partial_x w)^2)^{3/2}$. Since e.g. [3, Rem.4.2(c)] ensures that $-A(w)$ generates an analytic semigroup on $L_q(I)$, the operator $\mu + A(w)$ is boundedly invertible for each $\mu \in \mathbb{C}$ with sufficiently large real part. Clearly, μ belongs to the spectrum of $-A(w)$ only if it is an eigenvalue, and if $\varphi \in W_{q,D}^2(I, \mathbb{C})$ is a corresponding eigenvector, then testing the equation

$$\mu\varphi - \frac{1}{W}\partial_x^2\varphi = 0$$

with $W\bar{\varphi}$ readily gives $\mu \in (-\infty, 0)$ and $\mu \leq -\mu_1/\|W\|_\infty$, where $\mu_1 := \pi^2/4$ denotes the principal eigenvalue of $-\partial_x^2$ subject to Dirichlet boundary conditions. Consequently, due to (3.4) the resolvent set of $-A(w)$ contains the half plane $[\operatorname{Re} \mu \geq -\mu_1/c(\kappa)]$. Put $\omega := \mu_1/(2^{1+q/2}c(\kappa)) > 0$. Then ω is independent of w and $-\omega + A(w)$ is an isomorphism from $W_{q,D}^2(I)$ onto $L_q(I)$. Moreover, if $F \in L_q(I, \mathbb{C})$ and $\operatorname{Re} \mu > 0$, then the equation

$$\mu u - 2\omega u - \frac{1}{W}\partial_x^2 u = F$$

has a unique solution $u \in W_{q,D}^2(I, \mathbb{C})$. Testing this equation by $W|u|^{q-2}\bar{u}$ and using (3.4) yield the resolvent estimate (see, e.g. the proof of [23, Prop.2.4.2] for details)

$$\|(\mu - 2\omega + A(w))^{-1}\|_{\mathcal{L}(L_q(I))} \leq \frac{c'(\kappa)}{|\mu|},$$

where $c'(\kappa) > 0$ is a constant independent of w . Since we clearly have by (3.4)

$$\|-2\omega + A(w)\|_{\mathcal{L}(W_{q,D}^2(I), L_q(I))} \leq c'(\kappa),$$

it follows from [4, I.Rem.1.2.1(a)] that

$$-2\omega + A(w) \in \mathcal{H}(W_{q,D}^2(I), L_q(I); k, \omega)$$

for some $k := k(\kappa) \geq 1$. That $A(w)$ is resolvent positive and Lipschitz property (3.3) are readily seen, the latter being a consequence of the continuous embedding $W_q^{2-\xi}(I) \hookrightarrow C^1([-1, 1])$. \square

If $w = w(t)$ is Hölder continuous in t , then $A(w)$ generates a parabolic evolution operator $U_{A(w)}(t, s)$ on $L_q(I)$ in the sense of [4, Sect. II]. More precisely:

Proposition 3.2. *Let $q \in (2, \infty)$, $\kappa \in (0, 1)$, and $\xi \in (0, (q-1)/q)$. Let $\omega(\kappa) > 0$ and $\ell(\kappa) > 0$ be as in Lemma 3.1 and define, for $\rho \in (0, 1)$ and $N, \tau > 0$ given,*

$$\mathcal{W}_\tau(\kappa) := \{w \in C([0, \tau], W_{q,D}^{2-\xi}(I)); \|w(t) - w(s)\|_{W_{q,D}^{2-\xi}(I)} \leq \frac{N}{\ell(\kappa)}|t - s|^\rho \text{ and } w(t) \in Z_q(\kappa) \text{ for } 0 \leq t, s \leq \tau\}.$$

Then, there is a constant $c_0(\rho) > 0$ being independent of N and τ such that the following is true: for each $w \in \mathcal{W}_\tau(\kappa)$, there exists a unique parabolic evolution operator $U_{A(w)}(t, s)$, $0 \leq s \leq t \leq \tau$, possessing $W_{q,D}^2(I)$ as a regularity subspace, and

$$\|U_{A(w)}(t, s)\|_{\mathcal{L}(W_{q,D}^{2\alpha}(I), W_{q,D}^{2\beta}(I))} \leq c_*(\kappa) (t-s)^{\alpha-\beta} e^{-\vartheta(t-s)}, \quad 0 \leq s < t \leq \tau,$$

for $0 \leq \alpha \leq \beta \leq 1$ with $2\alpha, 2\beta \neq 1/q$, where the constant $c_*(\kappa) \geq 1$ depends on N , α , and β but is independent of τ , and

$$-\vartheta := -\vartheta(\kappa, \rho, N) := c_0(\rho)N^{1/\rho} - \omega(\kappa). \quad (3.5)$$

Moreover, $U_{A(w)}(t, s) \in \mathcal{L}(L_q(I))$ is a positive operator for $0 \leq s \leq t \leq \tau$.

Proof. Notice that, for each $w \in \mathcal{W}_\tau(\kappa)$,

$$A(w) \in C^0([0, \tau], \mathcal{L}(W_{q,D}^2(I), L_q(I))), \quad -2\omega(\kappa) + A(w) \subset \mathcal{H}(W_{q,D}^2(I), L_q(I); k(\kappa), \omega(\kappa)) \quad (3.6)$$

with

$$\sup_{0 \leq s < t \leq \tau} \frac{\|A(w(t)) - A(w(s))\|_{\mathcal{L}(W_{q,D}^2(I), L_q(I))}}{|t-s|^\rho} \leq N \quad (3.7)$$

by Lemma 3.1. Hence, the assertion follows from [4, II. Thm. 5.1.1, Lem. 5.1.3, Thm. 6.4.2] and the interpolation results of [11, 27]. \square

We are now in a position to prove Theorem 1.1.

Proof of Theorem 1.1. Using the previously introduced notation, we may rewrite (2.5) subject to the boundary condition (1.7) and the initial condition (1.8) as an abstract λ -dependent quasilinear Cauchy problem of the form

$$\frac{d}{dt}u + A(\varepsilon u)u = -\lambda g_\varepsilon(u), \quad t > 0, \quad u(0) = u^0. \quad (3.8)$$

Existence of solutions to (3.8) is based on applying a fixed point argument.

Let $\lambda > 0$ and $q \in (2, \infty)$. For simplicity we restrict to the case $\varepsilon \in (0, 1]$. A similar proof works for $\varepsilon > 1$ by changing some of the constants occurring in the proof. Let us consider an initial value $u^0 \in W_{q,D}^2(I)$ with $u^0(x) > -1$ for $x \in I$. Clearly, there is $\kappa \in (0, 1/2)$ with

$$u^0 \in \overline{S}_q(2\kappa), \quad \|u^0\|_{W_{q,D}^{2-\xi}(I)} \leq \frac{1}{2\kappa}, \quad (3.9)$$

where we fix ξ and σ such that

$$0 < \xi < \frac{1}{q}, \quad 0 < \frac{1}{2} - \frac{1}{q} < 2\sigma < \frac{1}{2} - \xi.$$

Let $4\rho \in (0, \xi)$, let $c_0(\rho) > 0$ be as in Proposition 3.2 and choose then $N > 0$ with the property that $-\vartheta := c_0(\rho)N^{1/\rho} - \omega(\kappa) < 0$ in (3.5). Since

$$W_{2,D}^{2\sigma}(I) \hookrightarrow W_{q,D}^{2\sigma - \frac{1}{2} + \frac{1}{q}}(I) \hookrightarrow L_q(I), \quad (3.10)$$

it follows from Proposition 3.2 (with $(\alpha, \beta) = (1, 1)$ and $(\alpha, \beta) = (\sigma - (q-2)/4q, 1)$) that, for $w \in \mathcal{W}_\tau(\kappa)$ fixed,

$$\|U_{A(w)}(t, s)\|_{\mathcal{L}(W_{q,D}^2(I))} + (t-s)^{-\sigma+1+\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} \|U_{A(w)}(t, s)\|_{\mathcal{L}(W_{2,D}^{2\sigma}(I), W_{q,D}^2(I))} \leq c_*(\kappa) e^{-\vartheta(t-s)}, \quad (3.11)$$

for $0 \leq s \leq t \leq \tau$, where the constant $c_*(\kappa) \geq 1$ is independent of w and $\tau > 0$. Now, set

$$\mathcal{V}_\tau(\kappa) := \left\{ v \in \mathcal{W}_\tau(\kappa) ; \|v(t)\|_{W_{q,D}^2(I)} \leq \frac{1}{\kappa_0} \text{ and } v(t) \geq -1 + \kappa \text{ for } 0 \leq t, s \leq \tau \right\},$$

with $\kappa_0 := \kappa/c_*(\kappa) \leq \kappa$ and $\tau > 0$ to be chosen later and observe that, when endowed with the topology of $C([0, \tau], W_{q,D}^{2-\xi}(I))$, $\mathcal{V}_\tau(\kappa)$ is a complete metric space. In addition, since $\kappa_0 \leq \kappa$, we have $v(t) \in \overline{S}_q(\kappa)$ for all $t \in [0, \tau]$ and $v \in \mathcal{V}_\tau(\kappa)$. It is, moreover, worthwhile to point out that $\varepsilon v \in \mathcal{V}_\tau(\kappa)$ for $v \in \mathcal{V}_\tau(\kappa)$ since $\varepsilon \in (0, 1]$.¹ By Proposition 2.1 there is $c_1(\kappa, \varepsilon) > 0$ such that

$$\|g_\varepsilon(v) - g_\varepsilon(w)\|_{W_{2,D}^{2\sigma}(I)} \leq c_1(\kappa, \varepsilon) \|v - w\|_{W_{q,D}^{2-\xi}(I)}, \quad v, w \in \overline{S}_q(\kappa), \quad (3.12)$$

and

$$\|g_\varepsilon(v)\|_{W_{2,D}^{2\sigma}(I)} \leq c_1(\kappa, \varepsilon), \quad v \in \overline{S}_q(\kappa). \quad (3.13)$$

We then claim that

$$\Lambda(v)(t) := U_{A(\varepsilon v)}(t, 0) u^0 - \lambda \int_0^t U_{A(\varepsilon v)}(t, s) g_\varepsilon(v(s)) \, ds, \quad t \in [0, \tau], \quad v \in \mathcal{V}_\tau(\kappa),$$

defines a contraction from $\mathcal{V}_\tau(\kappa)$ into itself if either $\lambda > 0$ is arbitrary and $\tau = \tau(\kappa, \lambda) > 0$ is sufficiently small, or $\lambda > 0$ is sufficiently small and $\tau > 0$ is arbitrary. To see this let v, w be arbitrary elements of $\mathcal{V}_\tau(\kappa)$ and let $t \in [0, \tau]$. Since $U_{A(\varepsilon v)}(t, 0)$ is a positive operator and $u^0 \geq -1 + 2\kappa$, it follows from the embedding $W_q^2(I) \hookrightarrow L_\infty(I)$ with constant 2, (3.11), and (3.13) that

$$\begin{aligned} \Lambda(v)(t) &\geq -1 + 2\kappa - 2\lambda \int_0^t \left\| U_{A(\varepsilon v)}(t, s) g_\varepsilon(v(s)) \right\|_{W_{q,D}^2(I)} \, ds \\ &\geq -1 + 2\kappa - 2\lambda c_*(\kappa) \int_0^t e^{-\vartheta(t-s)} (t-s)^{\sigma-1-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} \|g_\varepsilon(v(s))\|_{W_{2,D}^{2\sigma}(I)} \, ds \\ &\geq -1 + 2\kappa - 2\lambda c_*(\kappa) c_1(\kappa, \varepsilon) \int_0^\tau e^{-\vartheta s} s^{\sigma-1-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} \, ds, \end{aligned} \quad (3.14)$$

while (3.11) and (3.13) ensure that

$$\begin{aligned} \|\Lambda(v)(t)\|_{W_{q,D}^2(I)} &\leq c_*(\kappa) \|u^0\|_{W_{q,D}^2(I)} \\ &\quad + \lambda c_*(\kappa) \int_0^t e^{-\vartheta(t-s)} (t-s)^{\sigma-1-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} \|g_\varepsilon(v(s))\|_{W_{2,D}^{2\sigma}(I)} \, ds \\ &\leq \frac{c_*(\kappa)}{2\kappa} + \lambda c_*(\kappa) c_1(\kappa, \varepsilon) \int_0^\tau e^{-\vartheta s} s^{\sigma-1-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} \, ds. \end{aligned} \quad (3.15)$$

Moreover, since $\xi > 0$ and owing to (3.6), (3.7), and [4, II. Thm. 5.2.1] (with the choice $\beta = 1 - \xi/2$, $\alpha = 1$, and $2\gamma = 2\sigma - 1/2 + 1/q$ of the parameters therein) there is a number $n_*(\kappa) > 0$ such that

$$\begin{aligned} &\|\Lambda(v)(t) - \Lambda(w)(t)\|_{W_{q,D}^{2-\xi}(I)} \\ &\leq n_*(\kappa) e^{-\vartheta t} \left\{ \lambda t^{\frac{\xi}{2}} \|g_\varepsilon(v) - g_\varepsilon(w)\|_{L_\infty((0,t), L_q(I))} \right. \\ &\quad \left. + t^{\frac{\xi}{2}} \|A(\varepsilon v) - A(\varepsilon w)\|_{C([0,\tau], \mathcal{L}(W_{q,D}^2(I), L_q(I)))} \left[\|u^0\|_{W_{q,D}^2(I)} + \lambda t^{\sigma-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} \|g_\varepsilon(v)\|_{L_\infty((0,t), W_{q,D}^{2\sigma-1/2+1/q}(I))} \right] \right\} \\ &\leq \lambda n_*(\kappa) t^{\frac{\xi}{2}} e^{-\vartheta t} c_1(\kappa, \varepsilon) \|v - w\|_{\mathcal{V}_\tau(\kappa)} \\ &\quad + n_*(\kappa) t^{\frac{\xi}{2}} e^{-\vartheta t} \ell(\kappa) \|v - w\|_{\mathcal{V}_\tau(\kappa)} \left[\|u^0\|_{W_{q,D}^2(I)} + \lambda t^{\sigma-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} c_1(\kappa, \varepsilon) \right], \end{aligned}$$

¹As already mentioned, the case $\varepsilon > 1$ can be handled by taking different values of κ and N in the definition of $\mathcal{V}_\tau(\kappa)$.

where we used (3.3), (3.10), (3.12), and (3.13) for the second inequality (recall that $\|\cdot\|_{\mathcal{V}_\tau(\kappa)} = \|\cdot\|_{C([0,\tau],W_q^{2-\xi}(I))}$). Thus, there is $c(\kappa) > 0$ such that

$$\begin{aligned} \|\Lambda(v)(t) - \Lambda(w)(t)\|_{W_{q,D}^{2-\xi}(I)} &\leq c(\kappa) \left\{ \left(\max_{0 \leq r \leq \tau} r^{\frac{\xi}{2}} e^{-\vartheta r} \right) \left[\|u^0\|_{W_{q,D}^2(I)} + \lambda \right] \right. \\ &\quad \left. + \lambda \left(\max_{0 \leq r \leq \tau} r^{\frac{\xi}{2} + \sigma - \frac{1}{2}(\frac{1}{2} - \frac{1}{q})} e^{-\vartheta r} \right) \right\} \|v - w\|_{\mathcal{V}_\tau(\kappa)}. \end{aligned} \quad (3.16)$$

We next observe that (3.6), (3.7), and [4, II. Thm. 5.3.1] ensure the existence of a number $m_*(\kappa) > 0$ such that, for $0 \leq s \leq t \leq \tau$, (since $\xi > 0$ and $u^0 \in W_{q,D}^2(I) \hookrightarrow W_{q,D}^{2-\xi+4\rho}(I)$)

$$\begin{aligned} \|\Lambda(v)(t) - \Lambda(v)(s)\|_{W_{q,D}^{2-\xi}(I)} &\leq m_*(\kappa) (t-s)^{2\rho} e^{-\vartheta t} \left(\|u^0\|_{W_{q,D}^{2-\xi+4\rho}(I)} + \lambda \|g_\varepsilon(v)\|_{L^\infty((0,t),W_{2,D}^{2\sigma}(I))} \right) \\ &\leq m_*(\kappa) \left(\max_{0 \leq r \leq \tau} r^\rho e^{-\vartheta r} \right) \left(\|u^0\|_{W_{q,D}^{2-\xi+4\rho}(I)} + \lambda c_1(\kappa, \varepsilon) \right) (t-s)^\rho \end{aligned} \quad (3.17)$$

by using (3.13). Similarly, from (3.9),

$$\begin{aligned} \|\Lambda(v)(t)\|_{W_{q,D}^{2-\xi}(I)} &\leq \|\Lambda(v)(t) - \Lambda(v)(0)\|_{W_{q,D}^{2-\xi}(I)} + \|u^0\|_{W_{q,D}^{2-\xi}(I)} \\ &\leq m_*(\kappa) \left(\max_{0 \leq r \leq \tau} r^{2\rho} e^{-\vartheta r} \right) \left(\|u^0\|_{W_{q,D}^{2-\xi+4\rho}(I)} + \lambda c_1(\kappa, \varepsilon) \right) + \frac{1}{2\kappa}. \end{aligned} \quad (3.18)$$

Gathering (3.14)-(3.18) we see that, for arbitrary $\lambda > 0$, we may choose $\tau := \tau(\kappa, \lambda) > 0$ sufficiently small such that the mapping $\Lambda : \mathcal{V}_\tau(\kappa) \rightarrow \mathcal{V}_\tau(\kappa)$ defines a contraction and thus has a unique fixed point u in $\mathcal{V}_\tau(\kappa)$.

Now observe that, owing to Lemma 3.1, (3.10), and (3.12) we have

$$(A(\varepsilon u), g_\varepsilon(u)) \in C^\rho([0, \tau], \mathcal{H}(W_{q,D}^2(I), L_q(I)) \times W_{q,D}^{2\sigma - \frac{1}{2} + \frac{1}{q}}(I))$$

where $2\sigma - \frac{1}{2} + \frac{1}{q} > 0$, and u is a mild solution to (3.8) on $[0, \tau]$ with $u^0 \in W_{q,D}^2(I)$. Thus,

$$u \in C^1([0, \tau], L_q(I)) \cap C([0, \tau], W_{q,D}^2(I))$$

is a strong solution to (3.8) by [1, Thm. 4.2] and [3, Thm. 10.1], which then clearly can be extended to a maximal solution on some maximal interval $[0, T_m^\varepsilon)$. This proves part (i) of Theorem 1.1. Since τ above depends only on κ and λ , we also obtain part (ii) of Theorem 1.1 while the proof of part (iii) is the same as in [6, Thm. 1.1 (iii)]. To prove part (iv) of Theorem 1.1 we note that, since $\vartheta > 0$, there are $\lambda_*(\kappa) > 0$ and $r(\kappa) > 0$ such that, according to (3.14)-(3.18), the mapping $\Lambda : \mathcal{V}_\tau(\kappa) \rightarrow \mathcal{V}_\tau(\kappa)$ defines a contraction for each $\tau > 0$ provided that $\lambda \in (0, \lambda_*(\kappa))$ and $\|u^0\|_{W_{q,D}^2(I)} \leq r(\kappa)$. Thus, in this case there is a unique fixed point u of Λ belonging to $\mathcal{V}_\tau(\kappa)$ for each $\tau > 0$. By definition of $\mathcal{V}_\tau(\kappa)$, this implies Theorem 1.1 (iv). \square

Remark 3.3. Compared to the semilinear case investigated in [6], where A is independent of v , the quasilinear problem (3.8) features a further and not negligible difficulty. Indeed, employing Banach's fixed point theorem in the proof of Theorem 1.1 requires two essential ingredients. First, to warrant the existence of a corresponding evolution operator $U_{A(\varepsilon v)}$, the operator $A(\varepsilon v(t)) \in \mathcal{L}(W_{q,D}^2(I), L_q(I))$ has to be a Hölder continuous function of time as shown in (3.17). And second, the mapping $\Lambda = \Lambda(v)$ has to depend Lipschitz continuously (with Lipschitz constant less than 1) on its argument v , see (3.16). Both of these properties can be guaranteed only in the topology of $W_{q,D}^{2-\xi}(I)$ for $v = v(t)$ with $\xi > 0$ (but not for $\xi = 0$), see [4, II. Thm. 5.2.1, Thm. 5.3.1]. This issue is the reason for refining [6, Proposition 2.1]

and deriving (2.8). Moreover, unlike the semilinear case the initial value u^0 plays an additional role in the quasilinear case when it comes to global existence as becomes apparent from (3.16) and (3.17). Finally, it is worthwhile to point out that solving the elliptic problem (2.6), (2.7) requires that $v \in W_{q,D}^2(I)$, pointwise with respect to time, rather than $v \in W_{q,D}^{2-\xi}(I)$.

4. ASYMPTOTICALLY STABLE STEADY-STATE SOLUTIONS: PROOF OF THEOREM 1.2

We now prove Theorem 1.2(i). For this let $q \in (2, \infty)$, $\varepsilon > 0$, and $\kappa \in (0, 1)$ be fixed and note that, due to (2.5) and (3.1), steady-state solutions to (1.6)-(1.10) are characterized by

$$\partial_x^2 u = \lambda \frac{(1 + \varepsilon^2 (\partial_x u)^2)^{5/2}}{(1 + u)^2} |\partial_\eta \phi_u(\cdot, 1)|^2 =: \lambda h_\varepsilon(u), \quad x \in I, \quad u(\pm 1) = 0, \quad (4.1)$$

the function ϕ_u being defined in (2.19). Thus we may proceed as in [6, Thm. 1.3 (i)]: First, it follows exactly as in [6, Prop. 2.1] that the mapping $h_\varepsilon : S_q(\kappa) \rightarrow L_q(I)$ is analytic. Thus, since $-A(0) \in \mathcal{L}(W_{q,D}^2(I), L_q(I))$ is invertible, the operator $A(\cdot)$ being defined in (3.2), we obtain that the mapping

$$F : \mathbb{R} \times S_q(\kappa) \rightarrow W_{q,D}^2(I), \quad (\lambda, v) \mapsto v + \lambda A(0)^{-1} h_\varepsilon(v)$$

is analytic with $F(0, 0) = 0$ and $D_v F(0, 0) = \text{id}_{W_{q,D}^2}$. Now, the Implicit Function Theorem ensures the existence of $\delta = \delta(\kappa) > 0$ and an analytic function

$$[\lambda \mapsto U_\lambda] : [0, \delta) \rightarrow W_{q,D}^2(I)$$

such that $F(\lambda, U_\lambda) = 0$ for $\lambda \in [0, \delta)$. Hence $(U_\lambda, \Psi_\lambda)$ is the unique steady-state to (1.6)-(1.10) satisfying $U_\lambda \in S_q(\kappa)$ and $\Psi_\lambda \in W_2^2(\Omega(U_\lambda))$ when $\lambda \in (0, \delta)$. To improve the regularity of $(U_\lambda, \Psi_\lambda)$ as stated in (1.15), we may argue as follows: Since (4.1) is basically the same equation as considered in [20, Thm. 1], a steady-state solution (u, ψ) to (1.6)-(1.10) possessing the regularity property (1.15) can be constructed by means of Schauder's fixed point theorem applied to (4.1) with $u \in W_\infty^2(I) \cap S_q(\kappa)$ for each λ sufficiently small. Hence, making $\delta > 0$ smaller, if necessary, uniqueness guarantees $(u, \psi) = (U_\lambda, \Psi_\lambda)$. This proves Theorem 1.2 (i).

To prove part (ii) of Theorem 1.2, we proceed similarly as in [6, Thm. 1.3 (ii)] and use the Principle of Linearized Stability. Let $\lambda \in (0, \delta)$ be given and write $v = u - U_\lambda$. Then, introducing $Q \in C^\infty(S_q(\kappa), L_q(I))$ with $Q(U_\lambda) = 0$ by $Q(u) := -A(\varepsilon u)u - \lambda g_\varepsilon(u)$, the linearization of (3.8) reads

$$\frac{d}{dt} v - D_u Q(U_\lambda) v = Q(v + U_\lambda) - D_u Q(U_\lambda) v =: G_\lambda(v)$$

with $G_\lambda \in C^\infty(O_\lambda, L_q(I))$ being defined on some open zero neighborhood O_λ in $W_{q,D}^2(I)$ such that $U_\lambda + O_\lambda \subset S_q(\kappa)$. In view of (3.1) we obtain

$$\frac{d}{dt} v + (A(\varepsilon U_\lambda) + B_\lambda) v = G_\lambda(v), \quad (4.2)$$

where

$$B_\lambda v := \lambda g_\varepsilon(U_\lambda) \frac{3\varepsilon^2 \partial_x U_\lambda}{1 + \varepsilon^2 (\partial_x U_\lambda)^2} \partial_x v + \lambda D_u g_\varepsilon(U_\lambda) v.$$

Since $U_\lambda \in S_q(\kappa)$, we have $A(\varepsilon U_\lambda) \in \mathcal{H}(W_{q,D}^2(I), L_q(I); k, \omega)$ with spectral bound less than $-\omega < 0$. Thus, since $\|B_\lambda\|_{\mathcal{L}(W_{q,D}^2(I), L_q(I))} \rightarrow 0$ as $\lambda \rightarrow 0$, it follows from [4, I. Cor. 1.4.3] that $-(A(\varepsilon U_\lambda) + B_\lambda)$ is the generator of an analytic semigroup on $L_q(I)$ and there is $\omega_1 > 0$ such that the complex half plane $[\text{Re } z \geq -\omega_1]$ belongs to the resolvent set of $-(A(\varepsilon U_\lambda) + B_\lambda)$ provided that λ is sufficiently small. Now we may apply [24, Thm. 9.1.2] and conclude statement (ii)

of Theorem 1.2 by making $\delta > 0$ smaller, if necessary.

Let us note that also the corresponding potential converges exponentially toward the steady-state as $t \rightarrow \infty$.

Remark 4.1. In [6] (see (2.29) in the proof of Prop.2.1 therein) the following analogue of Lemma 2.6

$$\|\phi_v - \phi_w\|_{W_2^2(\Omega)} \leq c(\kappa, \varepsilon) \|v - w\|_{W_q^2(I)}, \quad v, w \in \overline{S}_q(\kappa),$$

was shown to hold. Consequently, under the assumptions of Theorem 1.2 (ii) there is $R_1 > 0$ such that

$$\|\phi_{u(t)} - \phi_{U_\lambda}\|_{W_2^2(\Omega)} \leq R_1 e^{-\omega_0 t} \|u^0 - U_\lambda\|_{W_{q,D}^2(I)}, \quad t \geq 0,$$

which shows that also $\psi = \phi_{u(t)} \circ T_{u(t)}$ (with transformation $T_{u(t)}$ defined in (2.1)) converges exponentially to Ψ_λ as $t \rightarrow \infty$.

5. NON-EXISTENCE OF STEADY-STATE SOLUTIONS: PROOF OF THEOREM 1.3

Consider a steady-state solution (u, ψ) to (1.6)-(1.10) with $u \in W_{q,D}^2(I)$, $\psi \in W_2^2(\Omega(u))$, and $u(x) > -1$ for $x \in I$. Recall that, according to (1.6), (3.1), and the identity

$$\partial_x \psi(x, u(x)) = -\partial_x u(x) \partial_z \psi(x, u(x)),$$

the function u solves

$$\partial_x^2 u(x) = \lambda \left(1 + \varepsilon^2 (\partial_x u(x))^2\right)^{5/2} |\partial_z \psi(x, u(x))|^2, \quad x \in I,$$

with $u(\pm 1) = 0$, so that u is negative in I by the comparison principle and convex. Thanks to the latter property, [20, Lem. 4.1] ensures that

$$\partial_z \psi(x, u(x)) \geq 1, \quad x \in I.$$

Hence, u satisfies the following differential inequality

$$\partial_x^2 u(x) \geq \lambda \left(1 + \varepsilon^2 (\partial_x u(x))^2\right)^{5/2}, \quad x \in I.$$

Introducing

$$J(r) := \int_0^r \frac{ds}{(1+s^2)^{5/2}} = \frac{r(2r^2+3)}{3(r^2+1)^{3/2}}, \quad r \geq 0,$$

the function J maps $[0, \infty)$ onto $[0, 2/3)$ and the previous differential inequality for u reads

$$\partial_x J(\varepsilon \partial_x u) \geq \lambda \varepsilon, \quad x \in I. \quad (5.1)$$

Let x_m be a point of minimum of u . Since $u < 0$ in I and $x \mapsto u(-x)$ is also a steady-state, we may assume without loss of generality that $x_m \in (-1, 0)$. Integrating (5.1) over (x_m, x) for $x \in [0, 1]$ gives

$$J(\varepsilon \partial_x u(x)) \geq J(0) + \lambda \varepsilon (x - x_m) \geq \lambda \varepsilon x, \quad x \in [0, 1]. \quad (5.2)$$

Now, either $\lambda \varepsilon \geq 2/3$ and we deduce from (5.2) that $J(\varepsilon \partial_x u(1)) \geq \lambda \varepsilon \geq 2/3$ which contradicts the boundedness of $\partial_x u$. Or $\lambda \varepsilon < 2/3$ and, since J is concave, we infer from (5.2) after integration over $(0, 1)$ and Jensen's inequality that

$$J(-\varepsilon u(0)) = J\left(\int_0^1 \varepsilon \partial_x u(x) dx\right) \geq \int_0^1 J(\varepsilon \partial_x u(x)) dx \geq \frac{\lambda \varepsilon}{2}.$$

If $\lambda \geq 2J(\varepsilon)/\varepsilon$, the previous inequality and the monotonicity of J entail that $-u(0) \geq 1$ and a contradiction again. Now, defining $\bar{\lambda}(\varepsilon) := \min\{2J(\varepsilon), 2/3\}/\varepsilon$ and noticing that $\bar{\lambda}(\varepsilon) \rightarrow 2$ as $\varepsilon \rightarrow 0$, Theorem 1.3 follows.

6. SMALL ASPECT RATIO LIMIT: PROOF OF THEOREM 1.4

To prove Theorem 1.4 fix $\lambda > 0$, $q \in (2, \infty)$, and $u^0 \in W_{q,D}^2(I)$ such that $-1 < u^0 \leq 0$ in I . Clearly, there is $\kappa \in (0, 1)$ such that $u^0 \in S_q(\kappa)$. For $\varepsilon \in (0, 1)$ let $(u_\varepsilon, \psi_\varepsilon)$ be the unique solution to (1.6)-(1.10) which is defined on the maximal interval of existence $[0, T_m^\varepsilon)$. In the following, $(K_i)_{i \geq 1}$ denote positive constants depending only on q and κ , but not on $\varepsilon > 0$ sufficiently small.

Set $\kappa_1 := \kappa / (2c_*(\kappa)) < \kappa$, where $c_*(\kappa) \geq 1$ is the constant defined in (3.11). The continuity properties of u_ε ensure

$$\tau^\varepsilon := \sup \{t \in [0, T_m^\varepsilon) : u_\varepsilon(s) \in \overline{S}_q(\kappa_1) \text{ for all } s \in [0, t]\} > 0. \quad (6.1)$$

Owing to the continuous embedding of $W_q^2(I)$ in $W_\infty^1(I)$, there is a positive constant K_1 such that, for all $\varepsilon > 0$,

$$-1 + \kappa_1 \leq u_\varepsilon(t, x) \leq 0, \quad (t, x) \in [0, \tau^\varepsilon] \times [-1, 1], \quad (6.2)$$

$$\|u_\varepsilon(t)\|_{W_q^2(I)} + \|u_\varepsilon(t)\|_{W_\infty^1(I)} \leq K_1, \quad t \in [0, \tau^\varepsilon]. \quad (6.3)$$

As a consequence of (6.3) there is $\varepsilon_0 > 0$ depending only on q and κ such that

$$\varepsilon_0^2 \|\partial_x u_\varepsilon(t)\|_{L_\infty(I)}^2 \leq \frac{1}{2}, \quad (t, \varepsilon) \in [0, \tau^\varepsilon] \times (0, \varepsilon_0]. \quad (6.4)$$

For $\varepsilon \in (0, \varepsilon_0)$, we set

$$\phi_\varepsilon(t) := \phi_{u_\varepsilon(t)} = \psi_\varepsilon(t) \circ T_{u_\varepsilon(t)}^{-1}, \quad t \in [0, \tau^\varepsilon],$$

with $T_{u_\varepsilon(t)}^{-1}$ given by (2.2) and

$$\Phi_\varepsilon(t, x, \eta) := \phi_\varepsilon(t, x, \eta) - \eta, \quad (t, x, \eta) \in [0, \tau^\varepsilon] \times \overline{\Omega}.$$

We first recall uniform estimates on Φ_ε that have been established in [6, Lem. 5.1]:

Lemma 6.1. *There exists a positive constant K_2 such that, for $\varepsilon \in (0, \varepsilon_0)$ and $t \in [0, \tau^\varepsilon]$,*

$$\|\partial_x \Phi_\varepsilon(t)\|_{L_2(\Omega)} + \frac{1}{\varepsilon} \left(\|\Phi_\varepsilon(t)\|_{L_2(\Omega)} + \|\partial_\eta \Phi_\varepsilon(t)\|_{L_2(\Omega)} \right) \leq K_2, \quad (6.5)$$

$$\frac{1}{\varepsilon} \|\partial_x \partial_\eta \Phi_\varepsilon(t)\|_{L_2(\Omega)} + \frac{1}{\varepsilon^2} \|\partial_\eta^2 \Phi_\varepsilon(t)\|_{L_2(\Omega)} \leq K_2, \quad (6.6)$$

$$\frac{1}{\varepsilon} \|\partial_\eta \Phi_\varepsilon(t, \cdot, 1)\|_{W_2^{1/2}(I)} \leq K_2. \quad (6.7)$$

At this point let us mention that the assumption $u^0 \leq 0$ is used to obtain the previous lemma. We then deduce from Lemma 6.1 a positive lower bound on τ^ε .

Lemma 6.2. (i) *There is $\tau > 0$ depending only on q , λ , and κ such that $\tau^\varepsilon \geq \tau$ for all $\varepsilon \in (0, \varepsilon_0)$.*

(ii) *There is $\Lambda := \Lambda(\kappa) > 0$ such that $\tau^\varepsilon = T_m^\varepsilon = \infty$ for all $\varepsilon \in (0, \varepsilon_0)$ provided $\lambda \in (0, \Lambda)$.*

Proof. Owing to (6.7) we have $\|\partial_\eta \phi_\varepsilon(t, \cdot, 1)\|_{W_2^{1/2}(I)} \leq 1 + K_2 \varepsilon$ while (6.2) and (6.3) imply that

$$\left\| \frac{1 + \varepsilon^2 (\partial_x u_\varepsilon(t))^2}{(1 + u_\varepsilon(t))^2} \right\|_{W_q^1(I)} \leq K_3$$

for $t \in [0, \tau^\varepsilon]$. Consequently, from the continuity of pointwise multiplication

$$W_q^1(I) \cdot W_2^{1/2}(I) \cdot W_2^{1/2}(I) \hookrightarrow W_2^{2\sigma}(I)$$

for $2\sigma \in (0, 1/2)$ fixed (see Theorem 7.1 below), we conclude that there is $K_4 > 0$ such that

$$\|g_\varepsilon(u_\varepsilon(t))\|_{W_2^{2\sigma}(I)} \leq K_4, \quad t \in [0, \tau^\varepsilon]. \quad (6.8)$$

As in the proof of (3.14) and (3.15), we infer from (3.11), (6.8), the fact that $u^0 \in S_q(\kappa)$, and the Variation-of-Constant formula that, for $t \in [0, \tau^\varepsilon]$,

$$\begin{aligned} \|u_\varepsilon(t)\|_{W_{q,D}^2(I)} &\leq c_*(\kappa) \|u^0\|_{W_{q,D}^2(I)} + \lambda c_*(\kappa) \int_0^t e^{-\vartheta(t-s)} (t-s)^{\sigma-1-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} \|g_\varepsilon(u_\varepsilon(s))\|_{W_{2,D}^{2\sigma}(I)} ds \\ &\leq \frac{c_*(\kappa)}{\kappa} + \lambda c_*(\kappa) K_4 \int_0^t e^{-\vartheta s} s^{\sigma-1-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} ds, \end{aligned} \quad (6.9)$$

and, using in addition the embedding of $W_{q,D}^2(I)$ in $L_\infty(I)$ with constant 2,

$$\begin{aligned} u_\varepsilon(t) &\geq -1 + \kappa - 2\lambda \int_0^t \left\| U_{A(\varepsilon u_\varepsilon)}(t,s) g_\varepsilon(u_\varepsilon(s)) \right\|_{W_{q,D}^2(I)} ds \\ &\geq -1 + \kappa - 2\lambda c_*(\kappa) \int_0^t e^{-\vartheta(t-s)} (t-s)^{\sigma-1-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} \|g_\varepsilon(u_\varepsilon(s))\|_{W_{2,D}^{2\sigma}(I)} ds \\ &\geq -1 + \kappa - 2\lambda c_*(\kappa) K_4 \int_0^t e^{-\vartheta s} s^{\sigma-1-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} ds. \end{aligned} \quad (6.10)$$

Since the integral above converges to zero as $t \rightarrow 0$, there exists $\tau > 0$ which depends only on q and κ such that

$$\int_0^t e^{-\vartheta s} s^{\sigma-1-\frac{1}{2}(\frac{1}{2}-\frac{1}{q})} ds < \min \left\{ \frac{1}{\lambda \kappa K_4}, \frac{(2c_*(\kappa) - 1)\kappa}{4\lambda c_*(\kappa)^2 K_4} \right\} \quad \text{for all } t \in [0, \tau]. \quad (6.11)$$

Thanks to this choice, we readily deduce from (6.9) and (6.10) that $\|u_\varepsilon(t)\|_{W_{q,D}^2(I)} \leq 1/\kappa_1$ and $u_\varepsilon(t) \geq -1 + \kappa_1$ for all $t \in [0, \tau] \cap [0, \tau^\varepsilon]$. Therefore, $u_\varepsilon(t) \in \overline{S}_q(\kappa_1)$ for all $t \in [0, \tau] \cap [0, \tau^\varepsilon]$ and the definition of τ^ε implies that $\tau^\varepsilon \geq \tau$. Finally, it is obvious from (6.11) that there is $\Lambda(\kappa) > 0$ such that (6.11) holds for any $\tau > 0$ and $\lambda \in (0, \Lambda(\kappa))$. This implies that $\tau^\varepsilon \geq \tau$ for any $\tau > 0$ so that $\tau^\varepsilon = \infty$. \square

Based on these auxiliary Lemmas 6.1 and 6.2 we may proceed exactly as in the proof of [6, Thm.1.4] to conclude the statement of Theorem 1.4. In particular, an obvious consequence of Lemma 6.1 and Lemma 6.2 is that $\phi_\varepsilon \rightarrow \eta$ in suitable topologies as $\varepsilon \rightarrow 0$ and that $(u_\varepsilon)_{\varepsilon \in (0, \varepsilon_0)}$ is bounded in $C([0, \tau], W_{q,D}^2(I))$.

7. APPENDIX

In this appendix we recall a useful tool on pointwise multiplication of functions in Sobolev spaces which is used frequently in the previous sections.

Theorem 7.1. *Let Ω be an open and non-empty subset of \mathbb{R}^n with finite volume. Let $m \geq 2$ be an integer and let $p, p_j \in [1, \infty)$ and $s, s_j \in (0, \infty)$ for $1 \leq j \leq m$ be real numbers satisfying $s \leq \min\{s_j\}$ and*

$$s - \frac{n}{p} < \begin{cases} \sum_{s_j < \frac{n}{p_j}} (s_j - \frac{n}{p_j}) & \text{if } \min_{1 \leq j \leq m} \left\{ s_j - \frac{n}{p_j} \right\} < 0, \\ \min_{1 \leq j \leq m} \left\{ s_j - \frac{n}{p_j} \right\} & \text{otherwise.} \end{cases}$$

Then pointwise multiplication

$$\prod_{j=1}^m W_{p_j}^{s_j}(\Omega) \rightarrow W_p^s(\Omega)$$

is continuous.

Theorem 7.1 is a consequence of the more general result stated in [2, Thm. 4.1] (see also Remark 4.2 (d) therein). It follows by observing that the Sobolev spaces $W_p^s(\Omega)$ coincide with the Besov spaces $B_{p,p}^s(\Omega)$ provided $s \in (0, \infty) \setminus \mathbb{N}$ and $p \in [1, \infty)$ together with the fact that $W_p^{s_1}(\Omega) \hookrightarrow B_{p,p}^{s_2}(\Omega)$ if $s_1 > s_2 > 0$.

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