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Influence of the supersaturation on Si diffusion and growth of Si nanoparticles in silicon-rich silica

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[Influence of the supersaturation on Si diffusion and growth of Si](http://dx.doi.org/10.1063/1.4792218) [nanoparticles in silicon-rich silica](http://dx.doi.org/10.1063/1.4792218)

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 $SiO_x/SiO₂$ multilayers have been prepared using magnetron sputtering and annealed in order to induce the growth of Si nanoparticles in Si-rich sublayers. This sample has undergone several successive annealing treatments and has been analyzed using a laser-assisted tomographic atom probe. This allows the phase separation between Si and $SiO₂$ and the growth process to be studied at the atomic scale as a function of annealing temperature. Si diffusion coefficient is estimated from the accurate measurement of matrix composition and Si particle size. We demonstrate that the diffusion coefficient in SiO_x is supersaturation dependent, leading to a decrease in silicon particle growth kinetics during annealing. In addition, we use our measurements to predict the critical thickness for efficient $SiO₂$ diffusion barriers. \odot 2013 American Institute of Physics. [\http://dx.doi.org/10.1063/1.4792218]

I. INTRODUCTION

The discovery of photoluminescence (PL) of porous sili-con by Canham^{[1](#page-8-0)} has demonstrated that silicon nanostructures are very promising materials for photonic applications. Thus, silicon nanoparticles (Si-nps) embedded in a dielectric matrix like $SiO₂$ have attracted major interest during the past few years. These materials exhibit quantum confinement effect in Si-nps and are fully compatible with silicon nanotechnology. Typical application fields concern the 3rd gener-ation of solar cells,^{[2](#page-8-0)} waveguide amplifiers^{[3](#page-8-0)} (if samples are doped with rare-earth atoms), and new memory devices based on carrier storage. 4 One of the major challenges for such applications is to control accurately structural characteristics of Si-nps, such as size distribution and density. For instance, tunable photoluminescent systems have been realized by accurately controlling the size distribution of Sinps.⁵ In such systems, optical properties also strongly depend on the nature of the interface between the Si-nps and the silica matrix. 6 In addition, in the case of electroluminescent devices or nanomemories, the quality of the dielectric matrix is of prime interest. A suboxide matrix can lead to undesirable charge trapping effects and, thereby, lead to the deterio-ration of electroluminescent or charge storage properties.^{[7–9](#page-8-0)} All these characteristics are determined during the formation and growth of Si-nps. Usually, Si-nps are obtained by using silicon-rich silicon oxide (SiO_X) prepared by deposition techniques (PECVD, 10 magnetron sputtering^{[2](#page-8-0)}) or silicon implan-tation in silica.^{[11](#page-8-0)} SiO_X is then annealed in order to induce the phase separation between Si and $SiO₂$, as predicted by the bi-nary Si-O phase diagram.^{[12](#page-8-0)} Controlling the size of Si-nps is commonly achieved by preparing SiO_X/SiO_2 multilayers $(MLs).¹³$ $(MLs).¹³$ $(MLs).¹³$ During phase separation, SiO₂ sublayers act as diffusion barriers and limit the growth of Si-nps to the thickness of SiO_X sublayers. Hence, it is essential to measure the diffu-

II. EXPERIMENTAL

A set of SiO_X/SiO_2 multilayer was deposited on [100] oriented silicon wafer using the reactive magnetron sputtering technique. $SiO₂$ sublayers were deposited using a pure silica target under Ar plasma while SiO_X ones were fabria)Electronic mail: etienne.talbot@univ-rouen.fr. cated by sputtering the silica target under 50% Ar $+50\%$ H₂

sion coefficient in SiO_X and $SiO₂$ accurately in order to anticipate growth kinetics in Si-np based devices. Various studies have focused on the diffusion mechanism of Si in $SiO₂¹⁴$ $SiO₂¹⁴$ $SiO₂¹⁴$ by means of various techniques such as SIMS,¹⁵ TEM,^{[16](#page-8-0)} and RBS.^{[17](#page-8-0)} Unfortunately, these studies provide a wide range of diffusion coefficient values, and it appears that Si diffusion mechanism in $SiO₂$ is still unclear. In this work, we have studied the structural and chemical properties of annealed $SiO_X/SiO₂$ multilayers using atom probe tomography (APT). To our knowledge, APT appears to be the only technique that is able to provide a full description of these materials. Indeed, APT overcomes most of the limitations encountered by conventional techniques by providing a 3D chemical map of the sample at the atomic scale. For instance, APT is able to evidence all Si-nps regardless of their structure (crystalline or amorphous) unlike HRTEM. In addition, it provides a 3D reconstruction of the sample, which is, in the case of multilayered structures, much more convenient than planar projections provided by microscopy techniques. Furthermore, it permits to measure directly the composition of the surrounding oxide, which is, usually, only extrapolated when using EFTEM. Finally, it is a very sensitive technique, which is able to identify the smallest particles (smaller than 1 nm diameter). The purpose of this study is to investigate the influence of annealing treatments on the growth of Si-nps and to determine the diffusion coefficient of Si in SiO_X and $SiO₂$.

mixture plasma. During this process, $H₂$ reduces oxygen and a silicon-rich layer (containing approximately 50 at. %) is deposited. The ML was grown at 650° C with a power density of 1.3 W·cm⁻². The sputtering time permits to tune SiO_X and $SiO₂$ sublayer thicknesses, which have been measured by HRTEM to be 4 nm. More details about the deposition process have been given by Ternon et al. in a previous pa-per.^{[13](#page-8-0)} This ML is then annealed at 900 °C for 1 h under N_2 flow, which is efficient enough to initiate phase separation in such systems. This state will be referred as the reference state (a) thereafter. In order to study the phase separation and the Si-np growth, additional annealing treatments have been applied to the same sample (1 h at $900\,^{\circ}\text{C} + 1$ h at $1000\,^{\circ}\text{C}$ (b), 1 h at $900\degree C + 1$ h at $1100\degree C$ (c) and 1 h at $900\degree C + 2$ h at 1100° C (d)).

After each annealing treatment, the MLs have been analyzed by laser assisted wide angle tomographic atom probe (LAWATAP, CAMECA version) using UV (343 nm) femtosecond laser pulses (50 nJ, 350 fs). APT principle relies on field evaporation of surface atoms of a specimen and their chemical identification by time-of-flight mass spectrome-try.^{[18](#page-8-0)–[20](#page-8-0)} Before an analysis, the specimen is prepared in the form of a sharp tip with a radius of curvature smaller than 50 nm. Specimen preparation has been detailed by Thompson *et al.* in a previous paper.^{[21](#page-8-0)} The tip is placed under high vacuum ($\approx 10^{-8}$ Pa), at low temperature (80 K), and submitted to a high positive voltage ($V_0 \approx 3 - 15$ kV). Under these conditions, an intense electric field (inversely proportional to the radius) is created at the tip apex (several $V \cdot nm^{-1}$). Surface atoms are ionized and evaporated by means of laser pulses. Then, ions are collected on a position sensitive detector. The time of flight of an ion (measured between the laser pulse and the impact on the detector) permits to identify its chemical nature by time of flight mass spectrometry. APT analysis has already been demonstrated to provide relevant and accurate results for similar materials. $22,33$

III. RESULTS AND DISCUSSION

A. Reference sample

The 3D reconstruction obtained on the reference sample is presented in Figure 1. Aforementioned, this sample was annealed for 1 h at 900 °C. Figures $1(a)$ and $1(b)$ represent silicon and oxygen atom position maps, respectively. In such APT reconstruction, each red dot corresponds to a single silicon atom while each green dot corresponds to an oxygen atom. Silicon supersaturation and oxygen depletion allow the identification of SiO_X and $SiO₂$ sublayers in the stacking sequence of the ML and the composition of each sublayer can be measured by counting Si and O atoms. $SiO₂$ sublayers composition has been measured to be 34.3 ± 0.3 at. % of Si, which corresponds to the composition of a pure silica layer $(X_{\text{Si}} = 33.3 \text{ at. } \%)$. SiO_X sublayers contain 51.0±0.3 at. % of Si.

Since each sublayer has been clearly identified, it is possible to focus our investigation on SiO_X sublayers in which the phase separation and the particle growth occur. Si-nps are evidenced using a cluster identification algorithm. The principle of this algorithm is to place a sphere over each

FIG. 1. 3D reconstruction of SiO_X / SiO_2 MLs annealed at 900 °C during 1 h. Arrows indicate SiO_X sublayers; (volume: $22 \times 22 \times 55$ nm³). (a) Silicon mapping; (b) oxygen mapping; (c) single SiO_X sublayer. Silicon atoms belonging to Si-nps are evidenced with bold dots (volume: $20 \times 20 \times 4$ nm³); (d) suboxide matrix in this layer where Si-nps position is highlight by dashed-circle.

atom of the analyzed volume, then, the composition is measured in this sphere and compared to a threshold. If the composition is higher than this threshold, then the selected atom belongs to a cluster. If the composition is lower than this threshold, then the selected atom belongs to the matrix. The radius of the sphere (1 nm) and the concentration threshold (55% at. of silicon) have been tuned in order to identify clearly every Si-np in analyzed volumes without missing atoms from the Si-nps. During the cluster identification algorithm, silicon and oxygen atoms from the matrix are inevitably selected as belonging to Si-nps because of local magnification effects, which can lead to trajectory aberration in such material. Corrections are applied to the measured composition and to the estimated number of atoms inside each cluster, as explained in Ref. [22.](#page-8-0) After this data treatment, it is possible to differentiate a silicon atom which belongs to a nanoparticle from an atom which belongs to the surrounding silicon oxide. From this moment, two different phases are defined inside SiO_X sublayers: the Si-nps and the surrounding matrix. Figure $1(c)$ shows the chemical mapping of a single SiO_X sublayer in which only Si-nps are represented. In this configuration, chemical and structural properties of Si-nps, such as composition, size, and density, can be measured. The size of a single nanoparticle is determined by counting the number of Si atoms constituting the Si-np. Assuming a spherical shape the particle diameter is calculated using the following relation:

$$
d = 2 \times \sqrt[3]{\frac{3}{4\pi} n_{\text{Si}} V_{\text{Si}}},
$$
 (1)

where d is the diameter of the particle (in nm), n_{Si} is the number of Si atoms in the particle, and V_{Si} is the atomic volume of a Si atom $(in \ nm^3)$. In this sample, the mean diameter of Si-nps is estimated to be 2.9 nm, which is smaller than the thickness of SiO_X sublayers (4 nm). Figure [1\(d\)](#page-3-0) shows the chemical map of matrix atoms (Si and O) of a single SiO_X layer. The matrix contains 41.9 ± 0.3 at. % of Si distributed homogeneously. This concentration is significantly higher than in pure silica, which indicates that phase separation is not complete for this annealing treatment. Both measurements of mean radius and matrix composition indicate that further growth is achievable in SiO_X sublayers of the reference sample. In order to analyse the phase separation process, reference sample (a) $(1 h at 900 °C)$ was cut in three parts, each part submitted to an additional annealing treatment: $+ 1 h$ at 1000 °C (sample (b)); $+ 1 h$ at 1100 °C (sample (c)); $+ 2h$ at 1100 °C (sample (d)). These results are described in Sec. III B.

B. Additional annealing treatments

Results obtained by APT after each annealing treatment (reference state and three additional annealing treatments) are represented in Figure 2. For each analysis, the atomic mapping of a single SiO_X sublayer is presented. In order to compare each nanostructure, identical volumes are represented and undergo the same data treatment (cluster identification algorithm, calculation of Si-np sizes, and measurement of matrix composition). It is important to note that Si-nps are exclusively observed into Si-rich sublayer for each annealing treatment. This confirms that 4 nm-SiO_2 acts as a strong diffusion barrier for silicon atoms. SiO_X sublayer images show that the average Si-np size increases with annealing time and temperature.

Si-np size distribution is calculated for the four annealing treatments and is reported in Figure $3(a)$ along with Sinp mean diameter. In the case of the reference state, most of FIG. 2. Plan view of silicon nanoparticles in SiO_X sublayers (volume: $16 \times 16 \times 4 \text{ nm}^3$) for the different annealing treatments. Red dots correspond to silicon atoms and green ones correspond to oxygen atoms. Si atoms belonging to the Si-nps are presented with bold dots. (a) 1 h at 900° C; (b) 1 h at 900° C + 1 h at 1000 °C; (c) 1 h at 900 °C + 1 h at 1100 °C; (d) 1 h at $900\degree C + 2h$ at 1100 °C (dashed line indicates some particle coalescence).

the nanoparticles (55.2%) have a diameter, which range from 3 nm to 4 nm; however, many nanoparticles (39.6%) are smaller than 3 nm. APT is capable of detecting Si-nps smaller than 1 nm. Additional annealing treatment induces further Si-np growth, as evidenced by a size distribution

FIG. 3. (a) Si-np size distributions measured for reference state and three additional heat treatments; (b) mean diameter for each annealing treatment.

shifting towards higher diameter values. Si-nps smaller than 1 nm were not observed after one additional hour of annealing at 1000 °C or higher. The number of nanoparticles inside the 1–2 nm range continuously decreases from 15% in the case of annealing treatment (a), to 3% in the case of annealing treatment (d). Concerning the last annealing treatment (d) the diameter of some nanoparticles appears to be larger than the thickness of the SiO_X sublayer. These nanoparticles can be located within analyzed volumes and correspond to coalescing nanoparticles as shown previously in Figure $2(d)$. Si-np mean diameter continuously increases with annealing time from 2.9 nm up to 3.4 nm, as shown in Figure $3(b)$. This value is still lower than SiO_X sublayer thickness (4 nm), indicating that more Si-np could still grow. However, coalescence may lead to some changes in the morphology towards more elongated rather than spherical nanoparticles. In that case, APT would not be able to distinguish coagulated particles from single elongated particles.

Measuring the density of nanoparticle by counting the number of nanoparticles in each sublayer leads to a high uncertainty. This is due to the variability of the number of particles inside analyzed volumes, which are very small: ± 1 particle in $20 \times 20 \times 4$ nm³ typically. In order to avoid such an uncertainty, we have calculated the density of Si-nps (N_{Si-nps}) in SiO_X sublayers using

$$
N_{\rm Si-nps} = \frac{\Delta N_{\rm Si}^{\rm ex}}{n_{\rm Si}^{\rm particles}},\tag{2}
$$

where $n_{Si}^{\text{particles}}$ is the mean number of Si atoms belonging to a nanoparticle (deduced from the size distribution). This value is compared to the total amount of excess silicon atoms precipitated in the form of Si-nps $(\Delta N_{\rm Si}^{\rm ex}$ in at/cm³). Calculated Si-np density is plotted as a function of annealing treatment in Figure $4(a)$. For the reference sample (a), the density of nanoparticles is estimated to be $11 \pm 1 \times 10^{18}$ Si-np/cm³. The expected value assuming a complete phase separation (pure silica matrix and 4 nm Si-nps) is 9.9×10^{18} Si-np/cm³. The first additional annealing treatment increases the density of nanoparticles up to $22 \pm 1 \times 10^{18}$ Si-np/cm³. Further annealing induces a lower density $17\pm1\times10^{18}$ Si-np/cm³ and $15\pm1\times10^{18}$ Si-np/cm³ for (c) and (d), respectively. Along with the evolution of the density, Si concentration in the matrix is also represented as a function of annealing treatment in Figure 4(b). Initial and final silicon concentrations have been represented for each annealing treatment in order to estimate how much silicon excess has precipitated during annealing. In the case of the first annealing treatment (a), silicon concentration decreases from 51 at. % to 42 at. %. In the case of additional annealing treatments (b) and (c), silicon concentration starts from 42 at. % and reaches its equilibrium value (pure silica 33.3 at. %). Finally, it is important to note that for annealing treatment (d), the matrix initial composition is already the equilibrium value of silica: 33.3 at. %.

Key structural and chemical parameters of the samples presented in this section are summarized in Table I. For each annealing treatment, it gives initial and final mean radius,

FIG. 4. (a) Density of particles in SiO_X sublayers for each annealing treatment. (b) Initial and final matrix supersaturation for each annealing treatment. Dashed lines represent expected density value (a) and matrix composition (b) in a theoretical layer in which all Si excess has precipitated in the form of precipitates with a diameter of 4 nm. Solid lines connecting the experimental points are drawn to guide the eye only.

initial and final matrix concentration, as well as initial and final density of nanoparticles.

C. Apparent diffusion coefficient

As highlighted by Tsoukalas et al.,^{[14](#page-8-0)} Si diffusion coefficient (D_{Si}) measurements appear to be highly dependent on experimental approach and especially on silicon concentration in silicon oxide. For instance, D_{Si} appears to be higher in silicon-implanted silica^{[14](#page-8-0)} or grown silicon-rich silica^{[16](#page-8-0)}

TABLE I. Structural and chemical characteristics of analyzed multilayers. For each annealing treatment, initial and final radius of Si-nps (calculated from Eq. [\(1\)](#page-3-0)), initial and final matrix composition, and initial and final number density of nanoparticles are given.

| Annealing treatment | R_{1} (nm) | R_{2} (nm) | Initial matrix Final matrix composition $X_{\rm SiO_X}$ $(at, \%)$ | composition $X_{\rm SiOx}$ (at, %) | Initial density $(10^{18}$ $Si-np/cm^3$) | Final density |
|------------------------|-----------------|-----------------|-----------------------------------------------------------------------------|------------------------------------------|----------------------------------------------------|------------------|
| (a) | 0 | 1.47 | 51% | 42% | Ω | 11 |
| (b) | 1.47 | 1.49 | 42% | 33.3% | 11 | 22 |
| (c) | 1.47 | 1.62 | 42% | 33.3% | 11 | 17 |
| (d) | 1.62 | 1.69 | 33.3% | 33.3% | 17 | 15 |

than in stoichiometric silica.^{[15,17](#page-8-0)} The growth of Si-nps is related to an apparent diffusion coefficient. Hence, we used the data collected in the previous section in order to estimate this coefficient. Previous publications propose several mechanisms for the formation of Si-nps: diffusion controlled growth,^{[16,24](#page-8-0)} Ostwald ripening, 25 25 25 and more complex processes such as oxygen out-diffusion from precipitates. 26 In the present study, both diffusion controlled growth and Ostwald ripening have been considered, depending on the supersaturation and the growing sequence: growth followed by coarsening. Drawing inspiration from Nesbit, 16 a classical growth approach is used to relate the apparent silicon diffusion coefficient D_{Si} to Si-np mean radius. This model assumes that Sinp growth is related to the consumption of Si excess atoms from the surrounding matrix. Thus, it can only be used in supersaturated silica. D_{Si} can be deduced from a simple application of the classical growth equation

$$
D_{\text{Si}} = \frac{R_2^2 - R_1^2}{2(t_2 - t_1)} \frac{X_{\text{Si}} - X_{\text{SiO}_2}}{X_{\text{SiO}_X} - X_{\text{SiO}_2}},\tag{3}
$$

where $(t_2 - t_1)$ represents the annealing time, R_1 and R_2 are Si-np mean radius before and after annealing, respectively. X_{Si} , X_{SiO_2} , and X_{SiO_X} are Si concentration in Si-nps (X_{Si}) $= 1.00$), Si concentration in pure SiO₂ ($X_{SiO₂} = 0.33$), and Si concentration in SiO_X ($X_{SiO_x} = 1/(1 + x)$), respectively.

In the case of a very low supersaturation, LSW theory (or Ostwald ripening) $2^{7,28}$ predicts that Si-np growth is related to the coarsening of nanoparticles.^{[25](#page-8-0)} In this model, silicon atoms coming from dissolving small nanoparticles diffuse towards bigger nanoparticles through the unsaturated matrix. In that process, the diffusion coefficient D_{Si} can be extracted as

$$
D_{\text{Si}} = \frac{9R_g T}{8\gamma X_{\text{SiO}_2} V_m} \frac{(R_2^3 - R_1^3)(X_{\text{Si}} - X_{\text{SiO}_2})}{(t_2 - t_1)},\tag{4}
$$

where R_g stands for the universal gas constant ($R_g = 8.31$) $J \cdot K^{-1} \cdot mol^{-1}$, T stands for the annealing temperature, γ stands for the interfacial energy, and V_m stands for the volume of 1 mole of Si $(V_m = 12 \text{ cm}^3)$. γ , which corresponds to the interfacial energy of a Si-np into $SiO₂$, has been estimated to be about $\gamma \approx 1 \text{ J} \cdot \text{m}^{-2}$.^{[29](#page-8-0)}

Thanks to the nanostructure characterization and the composition measurements carried out in the previous section, the growth mechanism of Si-nps is identified for each annealing treatment. Since high silicon composition is detected before annealing (a), (b), and (c), we considered a diffusion controlled growth mechanism in these cases. Thus, for (a) , (b) , and (c) , Eq. (3) has been used in order to calculate D_{Si} .

The density increases from sample (a) to samples (b) and (c). This indicates that nucleation still takes place in these samples. An overlap between nucleation and growth should be taken into account, as explained in Ref. [30.](#page-8-0) However, once a nucleus overcomes its critical radius, it is relevant to consider that its growth is diffusion controlled. The critical radius r^* of a nucleus is given by

TABLE II. Calculated value of silicon diffusion coefficient for the different annealing treatment. For each annealing treatment, the temperature and the duration of the annealing, the initial and final radius of Si-nps, and the initial matrix composition are given.

| Treatment | Temperature $t_2 - t_1$ R_1 $(^\circ C)$ | | R ₂ | Initial matrix composition (h) (in nm) (in nm) X_{SiO_X} (at. %) $(cm^2 \cdot s^{-1})$ | $D_{\rm Si}$ |
|-----------|-----------------------------------------------|----------|----------------|------------------------------------------------------------------------------------------------|-----------------------|
| (a) | 900 | θ | 1.47 | 51% | 1.2×10^{-17} |
| (b) | 1000 | 1.47 | 1.49 | 42% | 5.6×10^{-19} |
| (c) | 1100 | 1.47 | 1.62 | 42% | 4.6×10^{-18} |
| (d) | 1100 | 1.62 | 1.69 | 33.3% | 3.8×10^{-18} |

$$
r^* = \frac{-2\gamma}{\Delta G_V},\tag{5}
$$

where ΔG_V is the driving force per volume unit and γ is the interfacial energy. As predicted by the Si-O phase diagram, 12 Si solubility limit in $SiO₂$ is extremely low $(SiO₂$ is a stoichiometric compound). The driving force for this transformation is close to infinity and the critical radius of a nuclei is close to zero. Hence, it is possible to relate particle growth to a diffusion controlled process, as mentioned by Nesbit.^{[16](#page-8-0)}

In the case of the last annealing treatment (d), initial supersaturation is extremely low (as mentioned previously, initial silicon concentration for this annealing treatment is 33.3%, which correspond to pure silica). Thus, in that case, LSW theory and Eq. (4) have been used to calculate D_{Si} . All the parameters required to calculate D_{Si} values are reported in Table II along with calculated diffusion coefficients.

Apparent silicon diffusion coefficients are represented as a function of inverse temperature and compared to previous results obtained by Tsoukalas et $al.$ ^{[14](#page-8-0)} Nesbit,^{[16](#page-8-0)} and Jaoul et al ^{[17](#page-8-0)} in Figure 5. Differences between all these studies can be explained by the variation of silicon supersaturation in studied systems. In the system studied by Nesbit, silicon diffusion coefficient has been estimated using samples containing 55 at. % and 58 at. % of silicon. This

FIG. 5. Si diffusivity as a function of temperature in SiO_X . Existing results referred in bibliography are illustrated with straight lines. Dots represent results obtained in the present study.

TABLE III. Diffusion coefficient and mean free path of a silicon atom in SiO_2 and SiO_X during 1 h as a function of annealing temperature.

| $T(^{\circ}C)$ | | SiO ₂ layers \approx 33.3 at. % of Si (from Tsoukalas <i>et al.</i>) | | $SiOx$ layers \approx 50.0 at. % of Si (from Nesbit) | | |
|--------------------------------------------------|------------------------|------------------------------------------------------------------------------------|------------------------|--------------------------------------------------------|------------------------|------------------------|
| | 900 | 1000 | 1100 | 900 | 1000 | 1100 |
| $D_{\rm Si}$ (cm ² /s ⁻¹) | 6.19×10^{-21} | 2.46×10^{-19} | 5.70×10^{-18} | 8.38×10^{-18} | 3.66×10^{-17} | 1.29×10^{-16} |
| $\langle d \rangle$ (nm) during 1 h | 0.07 | 0.42 | 2.03 | 2.46 | 5.14 | 9.64 |

diffusion coefficient seems to be the highest. Tsoukalas et al. studied $SiO₂$, which had been implanted with $Si³O$ ions at 50 keV to a dose of 2×10^{-15} cm². The implanted dose represents a much smaller silicon supersaturation (less than 0.1 at. % of excess silicon). The diffusion coefficient measured by Tsoukalas et al. is lower than the one estimated by Nesbit. As for Jaoul et al., they studied Si diffusion in pure quartz and measure the lowest diffusion coefficient. Thus, it seems that D_{Si} is higher in highly supersaturated silica. Our measurements of D_{Si} are coherent with this statement. In the case of annealing (a), where the SiO_X matrix initially contained 51 at. % of silicon, we measure the highest value for D_{Si} (approximately the same as the one measured by Nesbit). In the case of annealing (b), (c), and (d), where the SiO_X matrix initially contains less Si excess (respectively, 42 at. % and 33.3 at. % of silicon), we measure a lower value for D_{Si} (approximately the same as the one measured by Tsoukalas).

Our experimental results, obtained for different temperatures and for various silicon concentrations, clearly evidence the dependence of the apparent silicon diffusion coefficient with Si supersaturation in SiO_X . The highest value of diffusivity is obtained for the highest supersaturation (case (a), which has been annealed 1 h at $900\,^{\circ}\text{C}$). Diffusivity drastically decreases with supersaturation, as observed for annealing treatments (b), (c), and (d) (even if they have been annealed at higher temperatures). This dependence of silicon diffusion coefficient with the silicon supersaturation into SiO_X has been computationally predicted by Orlandini *et al.*^{[31](#page-8-0)} and is attributed to a change in diffusion mechanisms. For low silicon supersaturations, diffusion is oxygen-driven, and this process is extremely slow. But for high silicon supersaturations, diffusion becomes silicon-driven and is much faster. Here we present experimental evidence that confirms this prediction. Moreover, the low coefficient diffusion in $SiO₂$ is in good agreement with the fact that Si-nps are exclusively observed in Si-rich layers. This confirms $SiO₂$ diffusion barriers efficiency.

D. Silica diffusion barrier

We demonstrated that D_{Si} varies with SiO_X sublayer composition. Thus, in $SiO_X/SiO₂ MLs$, there are two different regions with two different diffusion coefficients. In SiO_x sublayers, which initially contain more than 50% at. of silicon, we demonstrated that D_{Si} can be compared to the diffusion coefficient measured by Nesbit in $1985¹⁶$ $1985¹⁶$ $1985¹⁶$

$$
D_{\rm Si}^{\rm Nesbit} = 1.2 \times 10^{-13} \exp\left(\frac{-1.9 \times 1.6 \times 10^{-19}}{k_{\rm B}T}\right). \quad (6)
$$

In $SiO₂$ sublayers though, D_{Si} can be compared to the diffusion coefficient measured by Tsoukalas et al .^{[14](#page-8-0)}

$$
D_{\rm Si}^{\rm Tsoukalas} = 1.378 \times 10^{-4} \exp\left(\frac{-4.74 \times 1.6 \times 10^{-19}}{k_{\rm B}T}\right). \tag{7}
$$

From this equation, it is possible to estimate the mean free path $\langle d \rangle$ of a silicon atom through SiO_X or SiO₂ layers along the stacking direction

$$
\langle d \rangle = \sqrt{2 \times (Dt)}.\tag{8}
$$

 $\langle d \rangle$ is calculated for 1 h annealing treatment, in the case of SiO_X and $SiO₂$, and for different temperatures. They have been listed in Table III. These calculations permit to estimate the necessary thickness of $SiO₂$ barriers, depending on the annealing treatment which is going to be applied. For instance, considering an annealing treatment of 1h at $1000\degree$ C, 0.42 nm thick SiO₂ barriers will be necessary to prevent silicon to diffuse from one SiO_X layer to one other. Considering an annealing treatment of 1 h at 1100° C though, 2.03 nm thick barriers will be required.

In order to check the reliability of our model, we prepared two sets of SiO_X/SiO_2 multilayers. In both sets, SiO_X sublayer thickness and composition are the same (SiO_X) thickness is 4 nm and the layer contains \approx 50 at. %. of Si). The only changing parameter is $SiO₂$ barrier thickness. This thickness is equal to 3.0 nm in the first sample set and to 1.5 nm in the second one. Each set of multilayers has been annealed for 1 h at 1000° C and for 1 h at 1100° C. Figure [6](#page-8-0) represents 3D reconstructions obtained by APT of these four samples. In each volume, red polygons represent isoconcentration surfaces set to 55 at. % of Si and permit to evidence Si-rich areas. As expected, 3 nm thick SiO_2 barriers are sufficient to preserve the multilayered structure after 1 h at 1000 °C (Figure $6(a)$) and after 1 h at 1100 °C (Figure $6(b)$). This is consistent with our previous calculation. Concerning 1.5 nm thick $SiO₂$ barriers, the multilayered structure is preserved after 1 h at $1000\,^{\circ}$ C (Figure [6\(c\)](#page-8-0)). However, after 1 h at 1100 °C , the SiO_X/SiO₂ stacking sequence has disappeared (Figure [6\(d\)](#page-8-0)). For this annealing treatment, we predicted a mean free path of 2.03 nm. Since $\langle d \rangle$ is higher than SiO₂ barrier thickness, Si atoms are able to diffuse from one SiO_X layer to one other and the multilayered structure is lost.

IV. CONCLUSION

Silicon precipitation in supersaturated silica has been investigated by mean of APT. Several annealing treatments

FIG. 6. 3D reconstruction of two sets of sample with different $SiO₂$ layer thicknesses and submitted to two different annealing treatments. (a) 3 nm thick SiO_2 barriers, sample annealed for 1 h at 1000 °C. (b) 3 nm thick SiO_2 barriers, sample annealed for 1 h at 1100° C. (c) 1.5 nm thick SiO_2 barriers, sample annealed for 1 h at 1000° C. (d) 1.5 nm thick SiO_2 barriers, sample annealed for 1 h at 1100° C. In each volume, red polygons correspond to 55 at. % of Si isoconcentration surfaces.

(ranging from 1h at 900° C to 1h at 900° C + 2h at 1100 °C) have been applied to $4 \text{ nm-SiO}_X/4 \text{ nm-SiO}_2$ multilayered structures in order to induce Si precipitation. For each annealing treatment, changes in nanoparticle size, in nanoparticle density, and in matrix composition have been reported. By using diffusion controlled growth model and LSW theory, we have estimated an apparent silicon diffusion coefficients. These results lead to the conclusion that silicon supersaturation strongly influences the diffusion coefficient in SiO_X as predicted by Orlandini et al. These results can also explain the great variability in D_{Si} values reported in the literature and summed up by Tsoukalas et al. This phenomenon is already used in SiO_x/SiO_2 multilayers since SiO_2 is commonly used as a diffusion barrier. We were able to calculate D_{Si} into silica and to predict the critical thickness for these diffusion barriers depending on the annealing time and temperature. Nevertheless, since D_{Si} decreases along with the silicon excess during the phase separation, a supersaturation-dependent diffusion coefficient should be taken into account in order to predict accurately the kinetics of Si-nps growth in SiO_X systems.

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