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Xavier Lamy, Petru Mironescu

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Existence of critical points with semi-stiff boundary conditions for singular perturbation problems in simply connected planar domains

Preliminary version

Xavier Lamy ^{*} Petru Mironescu [†]

February 4, 2013

Abstract

Let Ω be a smooth bounded simply connected domain in \mathbb{R}^2 . We investigate the existence of critical points of the energy $E_\varepsilon(u) = 1/2 \int_\Omega |\nabla u|^2 + 1/(4\varepsilon^2) \int_\Omega (1 - |u|^2)^2$, where the complex map u has modulus one and prescribed degree d on the boundary. Under suitable nondegeneracy assumptions on Ω , we prove existence of critical points for small ε . In particular, we prove existence of critical points of prescribed degree one in domains close to a disc.

1 Introduction

Let $\Omega \subset \mathbb{R}^2$ be a smooth bounded simply connected domain. Let a map u belong to the space

$$\mathcal{E} := \{u \in H^1(\Omega, \mathbb{C}); |\operatorname{tr} u| = 1\},$$

where $\operatorname{tr} u$ denotes the trace of u on the boundary $\partial\Omega$. Then the trace $\operatorname{tr} u$ of u on $\partial\Omega$ belongs to the space $H^{1/2}(\partial\Omega; \mathbb{S}^1)$, and therefore we can define its

^{*}Université de Lyon, CNRS UMR 5208, Université Lyon 1, Institut Camille Jordan, 43 blvd. du 11 novembre 1918, F-69622 Villeurbanne cedex, France. Email address: xlamy@math.univ-lyon1.fr

[†]Université de Lyon, CNRS UMR 5208, Université Lyon 1, Institut Camille Jordan, 43 blvd. du 11 novembre 1918, F-69622 Villeurbanne cedex, France. Email address: mironescu@math.univ-lyon1.fr

winding number or degree, which we denote by $\deg(u, \partial\Omega)$ (see [10, Appendix]; see also [5, Section 2] for more details). This allows us to define the class

$$\mathcal{E}_d = \{u \in H^1(\Omega; \mathbb{C}); |\operatorname{tr} u| = 1, \deg(u, \partial\Omega) = d\}.$$

In this paper we study the existence of critical points of E_ε in \mathcal{E}_d , i.e., of critical points with prescribed degree d . The prescribed degree boundary condition is an intermediate model between the Dirichlet and the Neumann boundary conditions. The asymptotic of minimizers of the Ginzburg-Landau energy E_ε with Dirichlet boundary condition was first studied by Bethuel, Brezis and Hélein in their classical work [8]. In particular, it was shown in [8] that minimizers u_ε have zeros “well-inside” Ω , and that these zeros approach the singularities (vortices) of the limit u_* of the u_ε ’s as $\varepsilon \rightarrow 0$. In contrast, the only minimizers of E_ε with no boundary condition are constants. The same holds even for stable critical points of E_ε with Neumann boundary conditions [18]. The analysis of the prescribed degrees boundary condition (in domains which may be multiply connected) leads to a richer global picture [2], [14], [4], [3], [6], [12], [5]. More specifically, in multiply connected domains minimizers of E_ε may exist [14], [4] or not [3]. However, in such domains critical points of E_ε always exist [6], [12]. In simply connected domains, minimizers never exist [4]. More involved is the study of the existence of critical points in simply connected domains; this is our purpose. Typical methods in absence of absolute minimizers consist in constructing local minimizers, or constructing critical points by minimax methods. The first approach proved to be successful in multiply connected domains [6], but the arguments there do not adapt to our case. The latter approach led in [5] to the proof of the existence of critical points in simply connected domains for *large* ε ; but again we were unable to implement it in our case.

Our paper is devoted to the existence of critical points for *small* ε and thus complements [5]. Our approach is based on singular perturbations techniques. We explain this approach in the special case where the prescribed degree is $d = 1$. We first recall the main result in [8]. Consider the minimization of E_ε with Dirichlet boundary condition:

$$\min\{E_\varepsilon(u); \operatorname{tr} u = g \text{ on } \partial\Omega\}.$$

Here, $g : \partial\Omega \rightarrow \mathbb{S}^1$ is smooth, and we assume that $\deg(g, \partial\Omega) = 1$. Then there exists some $a \in \Omega$ such that, possibly up to a subsequence, minimizers u_ε satisfy $u_\varepsilon \rightarrow u_*$, with

$$u_*(z) = u_{*,a,g}(z) = \frac{z-a}{|z-a|} e^{iH}, \text{ with } H = H_{a,g} \text{ harmonic.} \quad (1.1)$$

In (1.1), the function H is uniquely determined (mod 2π) by the condition

$$u_* = g \text{ on } \partial\Omega. \quad (1.2)$$

The point a is not arbitrary: it has to be a critical point (actually, a point of minimum) of the “renormalized energy” $W(\cdot, g)$ associated with g .

In order to explain our main results in the case of prescribed degree boundary condition, we perform a handwaving analysis of our problem when $d = 1$. Assume that u_ε is a critical point of E_ε in \mathcal{E}_1 . Then u_ε has to vanish at some point a_ε , and up to a subsequence we have either

(i) $a_\varepsilon \rightarrow a \in \Omega$

or

(ii) $a_\varepsilon \rightarrow a \in \partial\Omega$.

Assume that (i) holds. Assume further, for the purpose of our discussion, that a_ε is the only zero of u_ε . Then the analysis in [8] suggests that the limit u_* of the u_ε 's should be again of the form $u_*(z) = \frac{z-a}{|z-a|} e^{i\psi}$. Formally, the fact that u_ε is a critical point of E_ε leads, as in [8], to the conclusion that the limiting point a is a critical point of a suitable renormalized energy $\widehat{W}(\cdot)$. Some basic properties of the energy \widehat{W} are studied in [15]; we will come back to this in Section 2. Of interest to us is the fact that \widehat{W} is smooth and does have critical points.

Let a be a critical point of \widehat{W} , and let u_* be as in (1.1)-(1.2). We plan to construct critical points u_ε of E_ε in \mathcal{E}_1 such that $u_\varepsilon \rightarrow u_*$ as $\varepsilon \rightarrow 0$. Our approach is inspired by the one of Pacard and Rivière [16]. In [16], critical points of E_ε with Dirichlet boundary condition g are constructed under a nondegeneracy assumption for the corresponding renormalized energy $W(\cdot, g)$. We encounter a similar situation in our problem: we are able to construct critical points of E_ε under some nondegeneracy assumptions that we explain below.

To start with, we will see in Section 2 that we may associate with each point $a \in \Omega$ a natural boundary datum g^a , solution of the minimization problem

$$\min\{W(a, g); g : \partial\Omega \rightarrow \mathbb{S}^1, \deg(g, \partial\Omega) = 1\}.$$

It turns out that, if a is a critical point of \widehat{W} , then a is also a critical point of $W(\cdot, g^a)$ (Section 2). Since \widehat{W} has a global maximum (this is easily seen from the explicit formula (3.9) of \widehat{W}), \widehat{W} has critical points, and thus there exists some $a \in \Omega$ critical point of $W(\cdot, g^a)$. Our first nondegeneracy assumption is

(ND1) there exists some $a \in \Omega$ nondegenerate critical point of $W(\cdot, g^a)$.

Assuming that (ND1) holds, set $g_0 := g^a$. Then we may prove that, for each g “close” to g_0 in a suitable sense, $W(\cdot, g)$ has a critical point $a(g)$ close to a (Section 5). Thus, to such $g \in C^{1,\alpha}(\partial\Omega; \mathbb{S}^1)$ we may associate the function

$$T_*(g) \in \dot{C}^\alpha(\partial\Omega; \mathbb{R}), \quad T_*(g) := u_* \wedge \frac{\partial u_*}{\partial \nu},$$

where $u_* = u_{*,a(g),g}$ is given by (1.1)-(1.2). One may prove that the map $g \mapsto T_*(g)$ is C^1 near g_0 , and that its differential L at g_0 is a Fredholm operator of index one (Section 10). Our second nondegeneracy assumption is

(ND2) L is onto.

We may now state our first result.

1.1 Theorem. *Assume that (ND1) and (ND2) hold. Then, for small ε , E_ε has critical points u_ε with prescribed degree one.*

The conditions (ND1) and (ND2) seem to be “generic”.¹ However, it is not clear whether the assumptions (ND1) and (ND2) are ever satisfied. Therefore, our next task is to exhibit nondegeneracy situations.

Loose Theorem. Assume that Ω is “close” to a disc. Then (ND1) and (ND2) hold. In particular, for small ε , E_ε has critical points of prescribed degree 1.

The above theorem applies to the unit disc \mathbb{D} . However, no sophisticated argument is needed for a disc. Indeed, when $\Omega = \mathbb{D}$ it is possible to construct explicit hedgehog type critical points of E_ε by minimizing E_ε in the class of the maps of the form $f(|z|)\frac{z}{|z|}$.

Our paper contains the proof of the two above theorems, as well as generalizations to higher degrees d . The plan of the paper is the following. In Section 2, we recall the definition and the main properties of the renormalized energies corresponding to either Dirichlet or prescribed degree boundary condition, and establish few new properties. In Sections 3 and 4, we derive new useful formulas for the renormalized energies. In Section 5, we prove that nondegeneracy of critical points of $W(\cdot, g)$ is stable with respect to small perturbations of g . Section 6 is devoted to the proof of a variant of the Pacard-Rivière [16] construction of critical points with Dirichlet condition; this is a key step in our proof. We prove Theorem 1.1 (for arbitrary degrees d) in Section 8. The proof relies on a Leray-Schauder degree argument, and the corresponding key estimate is obtained in Section 7. In Section 9, we prove that the couple of conditions (ND1)-(ND2) is stable with respect to small perturbations of the domain. This and the fact that $\Omega = \mathbb{D}$ satisfies (ND1)-(ND2) (Section 10) implies the Loose Theorem.

Notation

1. Points in \mathbb{R}^2 are denoted z in the Sections 3 and 4 relying on complex analysis techniques, and x or y elsewhere.

¹Critical points of smooth functionals are “generically” nondegenerate, and Fredholm operators of index one are “generically” onto.

2. $\mathbb{D}(z, r)$, $\overline{\mathbb{D}}(z, r)$ and $C(z, r)$ denote respectively the open disc, the closed disc and the circle of center z and radius r . We let $\mathbb{D} = \mathbb{D}(0, 1)$ denote the unit disc and set $\mathbb{D}_r = \mathbb{D}(0, r)$. \mathbb{S}^1 is the unit circle.
3. \wedge stands for the vector product of complex numbers or vectors. Examples: $(u_1 + \imath u_2) \wedge (v_1 + \imath v_2) = u_1 v_2 - u_2 v_1$, $(u_1 + \imath u_2) \wedge (\nabla v_1 + \imath \nabla v_2) = u_1 \nabla v_2 - u_2 \nabla v_1$, $(\nabla u_1 + \imath \nabla u_2) \wedge (\nabla v_1 + \imath \nabla v_2) = \nabla u_1 \cdot \nabla v_2 - \nabla u_2 \cdot \nabla v_1$.
4. If A is a set and k an integer, then we let
$$A_*^k = \{a = (a_1, \dots, a_k) \in A^k; a_j \neq a_l, \forall j \neq l\}.$$
5. When $k = 1$, we identify a collection $a = (a_1)$ with (the point or number) a_1 .
6. Additional indices emphasize the dependence of objects on variables. E.g.: $\psi_a = \psi_{a,g}$ recalls that ψ depends not only on a , but also on g .

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2 Renormalized energies and canonical maps

In the first part of this section, we follow [8] and [15].

We fix $k \in \mathbb{N}$ and a collection $\bar{d} = (d_1, \dots, d_k) \in \mathbb{Z}^k$, and we let $d := d_1 + \dots + d_k$. The bounded domain $\Omega \subset \mathbb{R}^2$ is assumed to be simply connected and $C^{1,\alpha}$.

We consider a collection of mutually distinct points in Ω , $a = (a_1, \dots, a_k) \in \Omega_*^k$ (the prescribed singularities), and also a boundary datum $g \in H^{1/2}(\partial\Omega; \mathbb{S}^1)$, of degree d .

For small $\rho > 0$, we define the open set $\Omega_\rho = \Omega \setminus \bigcup_{j=1}^k \bar{D}(a_j, \rho)$, and the classes of functions

$$\mathcal{F}_{\rho,g} = \{v \in H^1(\Omega_\rho; \mathbb{S}^1); \text{tr } v = g, \deg(v, C(a_j, \rho)) = d_j\}, \quad (2.1)$$

$$\widehat{\mathcal{F}}_\rho = \{v \in H^1(\Omega_\rho; \mathbb{S}^1); \deg(v, \partial\Omega) = d, \deg(v, C(a_j, \rho)) = d_j\}. \quad (2.2)$$

The functions in these classes have prescribed winding number d_j around each a_j , and prescribed boundary condition g (respectively prescribed degree d) on $\partial\Omega$. Of course, although we do not make this dependence explicit, the above classes depend not only on ρ and g , but also on a .

In [8] and [15], minimization of the Dirichlet energy $1/2 \int |\nabla v|^2$ over these spaces is studied, and the following asymptotic expansions are obtained as $\rho \rightarrow 0$:

$$\inf \left\{ \frac{1}{2} \int_{\Omega_\rho} |\nabla v|^2; v \in \mathcal{F}_{\rho,g} \right\} = \pi \left(\sum_{j=1}^k d_j^2 \right) \log \frac{1}{\rho} + W(a, g) + O(\rho), \quad (2.3)$$

$$\inf \left\{ \frac{1}{2} \int_{\Omega_\rho} |\nabla v|^2; v \in \widehat{\mathcal{F}}_\rho \right\} = \pi \left(\sum_{j=1}^k d_j^2 \right) \log \frac{1}{\rho} + \widehat{W}(a) + O(\rho). \quad (2.4)$$

In the above expressions, $W(a, g)$ and $\widehat{W}(a)$ are the so-called renormalized energies. These quantities depend not only on a and g , but also on \bar{d} and Ω .

Explicit formulae for the above renormalized energies can be found in [8] and [15], and involve the functions $\Phi_{a,g}$ and $\widehat{\Phi}_a$ defined as follows. $\Phi_{a,g}$ is the unique solution of

$$\begin{cases} \Delta \Phi_{a,g} = 2\pi \sum_{j=1}^k d_j \delta_{a_j} & \text{in } \Omega \\ \frac{\partial \Phi_{a,g}}{\partial \nu} = g \wedge \frac{\partial g}{\partial \tau} & \text{on } \partial\Omega, \\ \int_{\partial\Omega} \Phi_{a,g} = 0 \end{cases} \quad (2.5)$$

while $\widehat{\Phi}_a$ is the unique solution of

$$\begin{cases} \Delta \widehat{\Phi}_a = 2\pi \sum_{j=1}^k d_j \delta_{a_j} & \text{in } \Omega \\ \widehat{\Phi}_a = 0 & \text{on } \partial\Omega. \end{cases} \quad (2.6)$$

For further use, let us note that, if $\alpha \in \mathbb{S}^1$, then $\Phi_{a,g} = \Phi_{a,\alpha g}$. Therefore, we may naturally define $\Phi_{a,g}$ when g is an equivalence class in $H^{1/2}(\partial\Omega; \mathbb{S}^1)/\mathbb{S}^1$.

We also define the regular parts $R_{a,g}$ and \widehat{R}_a of $\Phi_{a,g}$ and $\widehat{\Phi}_a$ as follows:

$$R_{a,g}(x) = \Phi_{a,g}(x) - \sum_{j=1}^k d_j \log|x - a_j|, \quad \forall x \in \Omega, \quad (2.7)$$

respectively

$$\widehat{R}_a(x) = \widehat{\Phi}_a(x) - \sum_{j=1}^k d_j \log|x - a_j|, \quad \forall x \in \Omega. \quad (2.8)$$

The expressions of W and \widehat{W} are

$$W(a, g) = -\pi \sum_{j \neq l} d_j d_l \log|a_j - a_l| + \frac{1}{2} \int_{\partial\Omega} \Phi_{a,g} \left(g \wedge \frac{\partial g}{\partial \tau} \right) - \pi \sum_{j=1}^k d_j R_{a,g}(a_j), \quad (2.9)$$

respectively

$$\widehat{W}(a) = -\pi \sum_{j \neq l} d_j d_l \log|a_j - a_l| - \pi \sum_{j=1}^k d_j \widehat{R}_a(a_j). \quad (2.10)$$

The next result was proved in [15].

2.1 Proposition. *We have*

$$\widehat{W}(a) = \inf \{W(a, g); g \in \mathcal{E}_d\}, \quad (2.11)$$

and the infimum is attained in (2.11).

Recall that $\mathcal{E}_d := \{g \in H^{1/2}(\partial\Omega; \mathbb{S}^1); \deg(g, \partial\Omega) = d\}$.

We present here an alternative proof of Proposition 2.1, in the course of which we exhibit a formula of the form

$$W(a, g) = \widehat{W}(a) + \text{non negative terms},$$

which will be useful in the sequel.

Proof of Proposition 2.1. We identify a map $\psi \in H^{1/2}(\partial\Omega; \mathbb{R})$ with its harmonic extension to Ω , still denoted ψ . Given $\psi \in H^{1/2}(\partial\Omega; \mathbb{R})$, we define its (normalized) harmonic conjugate $\psi^* \in H^{1/2}(\partial\Omega; \mathbb{R})$ as follows. The harmonic extension of ψ^* (still denoted ψ^*) is the unique solution of

$$\begin{cases} \psi + i\psi^* \text{ is holomorphic in } \Omega, \\ \int_{\partial\Omega} \psi^* = 0. \end{cases} \quad (2.12)$$

Note that the Cauchy-Riemann equations imply

$$\frac{\partial \psi^*}{\partial \nu} = -\frac{\partial \psi}{\partial \tau} \quad \text{and} \quad \frac{\partial \psi^*}{\partial \tau} = \frac{\partial \psi}{\partial \nu}, \quad (2.13)$$

at least when ψ is smooth. When ψ is merely $H^{1/2}$, the distributions $\frac{\partial\psi}{\partial\nu}, \frac{\partial\psi^*}{\partial\nu} \in H^{-1/2}$ are respectively defined as the trace on $\partial\Omega$ of the normal derivatives of ψ and ψ^* , and the equalities in (2.13) are to be understood as equalities of distributions in $H^{-1/2}$.

We consider the space

$$H^{1/2}(\partial\Omega; \mathbb{R})/\mathbb{R} \simeq \dot{H}^{1/2}(\partial\Omega; \mathbb{R}) := \left\{ \psi \in H^{1/2}(\partial\Omega; \mathbb{R}); \int_{\partial\Omega} \psi = 0 \right\}, \quad (2.14)$$

which is endowed with the natural norm

$$|\psi|_{H^{1/2}}^2 = \int_{\Omega} |\nabla\psi|^2 = \int_{\Omega} |\nabla\psi^*|^2 = \int_{\partial\Omega} \frac{\partial\psi^*}{\partial\nu} \psi^* = - \int_{\partial\Omega} \frac{\partial\psi}{\partial\tau} \psi^*. \quad (2.15)$$

If ψ not smooth, then the two last integrals are to be understood as $H^{-1/2} - H^{1/2}$ duality brackets.

Given $a \in \Omega_*^k$, we define the canonical boundary datum associated with a as the unique element $g = g^a \in H^{1/2}(\partial\Omega; \mathbb{S}^1)/\mathbb{S}^1$ such that $\deg(g, \partial\Omega) = d$ and

$$g^a \wedge \frac{\partial g^a}{\partial\tau} = \frac{\partial \widehat{\Phi}_a}{\partial\nu}. \quad (2.16)$$

Our first observation is that g^a is well-defined and smooth. (It would be more accurate to assert that every map in the equivalence class defining g^a is smooth.) Indeed, existence of a smooth $g : \partial\Omega \rightarrow \mathbb{S}^1$ satisfying $g \wedge \frac{\partial g}{\partial\tau} = h$ (with given h) is equivalent to h smooth and

$$\int_{\partial\Omega} h = 2\pi d. \quad (2.17)$$

In addition, g (if it exists) is unique modulo \mathbb{S}^1 . In our case, we have $h = \frac{\partial \widehat{\Phi}_a}{\partial\nu}$, which is smooth (since $\widehat{\Phi}_a$ is smooth near $\partial\Omega$). In addition, using the equation (2.6), we see that (2.17) holds. If we compare the definition of g^a to the one of $\Phi_{a,g}$, we see that the canonical datum g^a is the unique g (modulo multiplication by a constant in \mathbb{S}^1) such that

$$\widehat{\Phi}_a = \Phi_{a,g}. \quad (2.18)$$

Given $g \in H^{1/2}(\partial\Omega; \mathbb{S}^1)$ with $\deg(g, \partial\Omega) = d$, we have $\deg(g/g^a, \partial\Omega) = 0$. Therefore, we may find $\psi = \psi_{a,g} \in H^{1/2}(\partial\Omega; \mathbb{R})$, unique modulo a constant, such that [9]

$$g = g^a e^{i\psi} = g^a e^{i\psi_{a,g}}. \quad (2.19)$$

Thus we have

$$\begin{cases} \Delta [\Phi_{a,g} - \widehat{\Phi}_a] = 0 & \text{in } \Omega \\ \frac{\partial}{\partial\nu} [\Phi_{a,g} - \widehat{\Phi}_a] = \frac{\partial\psi}{\partial\tau} & \text{on } \partial\Omega. \\ \int_{\partial\Omega} (\Phi_{a,g} - \widehat{\Phi}_a) = 0 \end{cases} \quad (2.20)$$

Combining the above with the definition of the harmonic conjugate, we find that

$$\Phi_{a,g} = \widehat{\Phi}_a - \psi^* = \widehat{\Phi}_a - \psi_{a,g}^*. \quad (2.21)$$

Plugging (2.21) into the expression of $W(a, g)$ given by formula (2.9), we find

$$\begin{aligned} W(a, g) &= -\pi \sum_{j \neq l} d_j d_l \log |a_j - a_l| + \frac{1}{2} \int_{\partial\Omega} (\widehat{\Phi}_a - \psi^*) \left(g^a \wedge \frac{\partial g^a}{\partial \tau} + \frac{\partial \psi}{\partial \tau} \right) \\ &\quad - \pi \sum_{j=1}^k d_j (\widehat{R}_a(a_j) - \psi^*(a_j)) \\ &= \widehat{W}(a) - \frac{1}{2} \int_{\partial\Omega} \psi^* \left(g^a \wedge \frac{\partial g^a}{\partial \tau} \right) - \frac{1}{2} \int_{\partial\Omega} \psi^* \frac{\partial \psi}{\partial \tau} + \pi \sum_{j=1}^k d_j \psi^*(a_j). \end{aligned} \quad (2.22)$$

In the last equality we used the fact that $\widehat{\Phi}_a = 0$ on $\partial\Omega$. Furthermore, using the definition of g^a and the fact that ψ^* is harmonic, we obtain

$$\begin{aligned} \int_{\partial\Omega} \psi^* \left(g^a \wedge \frac{\partial g^a}{\partial \tau} \right) &= \int_{\partial\Omega} \psi^* \frac{\partial \widehat{\Phi}_a}{\partial \nu} \\ &= \int_{\partial\Omega} \frac{\partial \psi^*}{\partial \nu} \widehat{\Phi}_a + \int_{\Omega} \psi^* \Delta \widehat{\Phi}_a = 2\pi \sum_{j=1}^k d_j \psi^*(a_j). \end{aligned} \quad (2.23)$$

Using (2.15), (2.22) and (2.23), we finally obtain

$$W(a, g) = \widehat{W}(a) + \frac{1}{2} |\psi_{a,g}|_{H^{1/2}}^2. \quad (2.24)$$

In particular, we recover the conclusion of Proposition 2.1 in the following stronger form: the minimum of $W(a, \cdot)$ is attained (exactly) when $g = g^a$ (modulo \mathbb{S}^1). \square

2.2 Remark. The canonical boundary datum g^a will play a crucial role in our subsequent analysis. We emphasize here the fact that g^a is the (unique modulo \mathbb{S}^1) solution of

$$g^a \wedge \frac{\partial g^a}{\partial \tau} = \frac{\partial \widehat{\Phi}_a}{\partial \nu} \quad \text{on } \partial\Omega. \quad (2.25)$$

The limit (as $\rho \rightarrow 0$) of the variational problem (2.3) is also connected to the so-called canonical harmonic map $u_{*,a,g}$ associated to prescribed singularities $a \in \Omega_*^k$ and to the Dirichlet condition $g \in H^{1/2}(\partial\Omega; \mathbb{S}^1)$. In fact, in [8, Chapter I] it is proved that the unique solution $u_{\rho,g}$ of the minimization problem $\inf \{ \int |\nabla u|^2; u \in \mathcal{F}_{\rho,g} \}$ tends to $u_{*,a,g}$, in $C_{loc}^k(\Omega \setminus \{a_j\})$ as $\rho \rightarrow 0$.²

²Actually, in [8, Chapter I] the map g is supposed smooth, but the argument adapts to a general $g \in H^{1/2}(\partial\Omega; \mathbb{S}^1)$.

The canonical harmonic map is defined by the formula

$$\begin{cases} u = u_{*,a,g} = e^{iH} \prod_{j=1}^k \left(\frac{z-a_j}{|z-a_j|} \right)^{d_j} & \text{in } \Omega \\ \Delta H = 0 & \text{in } \Omega \\ u = g & \text{on } \partial\Omega \end{cases} \quad (2.26)$$

The fact that $\deg(g, \partial\Omega) = d = \sum d_j$ guarantees that $H = H_g$ is well defined. Indeed, there exists $\psi \in H^{1/2}(\partial\Omega; \mathbb{R})$ such that

$$g \prod_{j=1}^k \left(\frac{z-a_j}{|z-a_j|} \right)^{-d_j} = e^{i\psi},$$

and then we can simply let H be the harmonic extension of ψ . On the other hand, we note that H is uniquely defined up to a multiple of 2π .

Equivalently, u in (2.26) is characterized by [8, Chapter I]

$$\begin{cases} |u| = 1 \\ u \wedge \frac{\partial u}{\partial x_1} = -\frac{\partial \Phi_{a,g}}{\partial x_2} \\ u \wedge \frac{\partial u}{\partial x_2} = \frac{\partial \Phi_{a,g}}{\partial x_1} \\ u = g \text{ on } \partial\Omega \end{cases} \quad (2.27)$$

In particular, we have

$$u_{*,a,g} \wedge \frac{\partial u_{*,a,g}}{\partial \nu} = -\frac{\partial \Phi_{a,g}}{\partial \tau} \quad \text{on } \partial\Omega \quad (2.28)$$

and

$$\int_{\partial\Omega} u_{*,a,g} \wedge \frac{\partial u_{*,a,g}}{\partial \nu} = 0. \quad (2.29)$$

2.3 Remark. For the minimization problem (2.4), the situation is similar. As established in [15], the solution v_ρ to $\inf \left\{ \int |\nabla v|^2; v \in \widehat{\mathcal{F}}_\rho \right\}$ converges (in an appropriate sense) as $\rho \rightarrow 0$, to $v_{*,a} := u_{*,a,g^a}$. Since g^a is defined modulo \mathbb{S}^1 , $v_{*,a}$ is also defined modulo \mathbb{S}^1 . Therefore, in this context the convergence actually means that subsequences of (v_ρ) converge to representatives (modulo \mathbb{S}^1) of $v_{*,a}$.

We end this section with the definition of the following quantity, which will play a very important role in what follows. For $a \in \Omega_*^k$ and $g \in H^{1/2}(\partial\Omega; \mathbb{S}^1)$, we set

$$N(a, g) := u_{*,a,g} \wedge \frac{\partial u_{*,a,g}}{\partial \nu} = -\frac{\partial \Phi_{a,g}}{\partial \tau} \in H^{-1/2}(\partial\Omega; \mathbb{R}). \quad (2.30)$$

3 Transport of formulas onto the unit disc

Let $f : \mathbb{D} \rightarrow \Omega$ be a conformal representation. The assumption $\Omega \in C^{1,\alpha}$ ensures that f and its inverse $\varphi := f^{-1} : \Omega \rightarrow \mathbb{D}$ are $C^{1,\alpha}$ up to the boundary.

The goal of this section is to understand how the objects defined in Section 2 are transported by φ and f .

We will stress the dependence on the domain by using superscripts (e.g. $W = W^\Omega$). For $\alpha \in \mathbb{D}_*^k$, the notation $a = f(\alpha)$ stands for $a := (f(\alpha_1), \dots, f(\alpha_k)) \in \Omega_*^k$.

First of all, for $\alpha \in \Omega_*^k$, we have

$$\Phi_{a,g}^\Omega = \Phi_{\varphi(\alpha),g \circ f}^\mathbb{D} \circ \varphi + C, \quad (3.1)$$

where $C = C(\alpha, g, f) = -\int_{\mathbb{S}^1} \Phi_{\varphi(\alpha),g \circ f}^\mathbb{D} |f'|$. Indeed, (3.1) is justified as follows. By a direct calculation, both sides of (3.1) satisfy the same Poisson equation, with the same Neumann boundary condition. The constant C comes from the normalization condition $\int_{\partial\Omega} \Phi_{a,g} = 0$. The same argument applies to show that

$$\widehat{\Phi}_a^\Omega = \widehat{\Phi}_{\varphi(\alpha)}^\mathbb{D} \circ \varphi. \quad (3.2)$$

Here there is no renormalization constant since $\widehat{\Phi}_a$ satisfies a Dirichlet boundary condition.

Normal and tangential derivatives transform in the following way. If $v : \mathbb{D} \rightarrow \mathbb{C}$, then

$$\frac{\partial}{\partial \tau} [v \circ \varphi](z) = |\varphi'(z)| \frac{\partial v}{\partial \tau}(\varphi(z)), \quad z \in \partial\Omega, \quad (3.3)$$

$$\frac{\partial}{\partial \nu} [v \circ \varphi](z) = |\varphi'(z)| \frac{\partial v}{\partial \nu}(\varphi(z)), \quad z \in \partial\Omega. \quad (3.4)$$

Using (3.2), (3.3), (3.4) together with formula (2.16) characterizing g^a , we find, for $\alpha \in \Omega_*^k$,

$$g^a \circ f = g^{\varphi(\alpha)}. \quad (3.5)$$

On the other hand, we claim that

$$u_{*,a,g}^\Omega = u_{*,\varphi(\alpha),g \circ f}^\mathbb{D} \circ \varphi. \quad (3.6)$$

Indeed, this follows from the observation that the two sides of (3.6) agree on $\partial\Omega$, combined with (2.26) and with the fact, when H is harmonic in \mathbb{D} , we may write

$$\frac{\varphi(z) - \varphi(a)}{|\varphi(z) - \varphi(a)|} e^{iH \circ \varphi(z)} = \frac{z - a}{|z - a|} e^{iK(z)}, \quad \text{with } K \text{ harmonic in } \Omega.$$

As a consequence of (3.6) and (3.4), we obtain, recalling the definition (2.30) of N ,

$$N^\Omega(a, g) = |\varphi'| N^\mathbb{D}(\varphi(a), g \circ f) \circ \varphi. \quad (3.7)$$

The formulas of the renormalized energies \widehat{W} and W transport in a more complicated way.

3.1 Lemma. *Let $\alpha \in \mathbb{D}_*^k$, $a := f(\alpha)$ and $g : \partial\Omega \rightarrow \mathbb{S}^1$. Then*

$$W^\Omega(a, g) = W^\mathbb{D}(\alpha, g \circ f) + \pi \sum_j d_j^2 \log |f'(\alpha_j)|, \quad (3.8)$$

$$\widehat{W}^\Omega(a) = \widehat{W}^\mathbb{D}(\alpha) + \pi \sum_j d_j^2 \log |f'(\alpha_j)|. \quad (3.9)$$

Proof. Using definition of $R_{a, g}$ (2.7), together with (3.1), we compute, for $z \in \mathbb{D}$,

$$\begin{aligned} R_{a, g}^\Omega(f(z)) &= \Phi_{\alpha, g \circ f}^\mathbb{D}(z) - \sum_{l=1}^k d_l \log |f(z) - f(\alpha_l)| + C \\ &= R_{\alpha, g \circ f}^\mathbb{D}(z) - \sum_{l=1}^k d_l \log \left| \frac{f(z) - f(\alpha_l)}{z - \alpha_l} \right| + C. \end{aligned} \quad (3.10)$$

The above is well-defined when $z \neq \alpha_j$, and extends by continuity at $z = \alpha_j$. In particular,

$$R_{a, g}^\Omega(f(\alpha_j)) = R_{\alpha, g \circ f}^\mathbb{D}(\alpha_j) - \sum_{l \neq j} d_l \log \left| \frac{f(\alpha_j) - f(\alpha_l)}{\alpha_j - \alpha_l} \right| - d_j \log |f'(\alpha_j)| + C. \quad (3.11)$$

Finally, we plug (3.1) and (3.11) into formula (2.9) expressing W in terms of $\Phi_{a, g}$ and $R_{a, g}$. We obtain, using also the fact that $\deg(g, \partial\Omega) = d = \sum d_j$,

$$\begin{aligned} W^\Omega(a, g) &= -\pi \sum_{j \neq l} d_j d_l \log |\alpha_j - \alpha_l| + \frac{1}{2} \int_{\partial\Omega} \Phi_{\alpha, g \circ f}^\mathbb{D} \circ \varphi \left(g \wedge \frac{\partial g}{\partial \tau} \right) \\ &\quad + \frac{1}{2} C \int_{\partial\Omega} g \wedge \frac{\partial g}{\partial \tau} - \pi \sum_j d_j C - \pi \sum_{j=1}^k d_j R_{\alpha, g \circ f}^\mathbb{D}(\alpha_j) + \pi \sum_{j=1}^k d_j^2 \log |f'(\alpha_j)| \\ &= -\pi \sum_{j \neq l} d_j d_l \log |\alpha_j - \alpha_l| + \frac{1}{2} \int_{\partial\mathbb{D}} \Phi_{\alpha, g \circ f}^\mathbb{D}(g \circ f) \wedge \frac{\partial}{\partial \tau}(g \circ f) \\ &\quad - \pi \sum_{j=1}^k d_j R_{\alpha, g \circ f}^\mathbb{D}(\alpha_j) + \pi \sum_{j=1}^k d_j^2 \log |f'(\alpha_j)| \\ &= W^\mathbb{D}(\alpha, g \circ f) + \pi \sum_j d_j^2 \log |f'(\alpha_j)|. \end{aligned}$$

Formula (3.9) can be proved following the same lines (the calculations are even simpler than for (3.8)). Alternatively, we can obtain (3.9) via the relation $\widehat{W}(a) = W(a, g^a)$. \square

4 Explicit formulas in the unit disc

In this section we derive explicit formulas for $\widehat{W}^{\mathbb{D}}$, $W^{\mathbb{D}}$ and $N^{\mathbb{D}}$.

We start by recalling the explicit formulas for $\widehat{\Phi}_\alpha^{\mathbb{D}}$ and $\widehat{W}^{\mathbb{D}}$ [15]: for $\alpha \in \mathbb{D}_*^k$, we have

$$\widehat{\Phi}_\alpha^{\mathbb{D}}(z) = \sum_{j=1}^k d_j (\log|z - \alpha_j| - \log|1 - \bar{\alpha}_j z|), \quad \forall z \in \mathbb{D}, \quad (4.1)$$

$$\begin{aligned} \widehat{W}^{\mathbb{D}}(\alpha) &= -\pi \sum_{j \neq l} d_j d_l \log|\alpha_j - \alpha_l| + \pi \sum_{j \neq l} d_j d_l \log|1 - \bar{\alpha}_j \alpha_l| \\ &\quad + \pi \sum_j d_j^2 \log(1 - |\alpha_j|^2). \end{aligned} \quad (4.2)$$

The formulas for W and N are more involved.

4.1 Lemma. *Let $\alpha^0 \in \mathbb{D}_*^k$ be fixed, and $g^0 := g^{\alpha^0} : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ be an associated canonical boundary map. (Recall that g^0 is defined up to a multiplicative constant.) Then it holds:*

(i) For $\alpha \in \mathbb{D}_*^k$ and for $\psi \in H^{1/2}(\mathbb{S}^1; \mathbb{R})$,

$$W^{\mathbb{D}}(\alpha, g^0 e^{i\psi}) = \widehat{W}^{\mathbb{D}}(\alpha) + \frac{1}{2} \int_{\mathbb{D}} \left| \nabla (\psi_{\alpha, g^0}^* + \psi^*) \right|^2, \quad (4.3)$$

and, for $z \in \mathbb{D}$,

$$\nabla \psi_{\alpha, g^0}^*(z) = 2 \sum_{j=1}^k d_j \left(\frac{\alpha_j (1 - \bar{\alpha}_j z)}{|1 - \bar{\alpha}_j z|^2} - \frac{\alpha_j^0 (1 - \bar{\alpha}_j^0 z)}{|1 - \bar{\alpha}_j^0 z|^2} \right) \in \mathbb{C} \simeq \mathbb{R}^2. \quad (4.4)$$

(ii) For $\alpha \in \mathbb{D}_*^k$ and for $\psi \in H^{1/2}(\mathbb{S}^1; \mathbb{R})$,

$$N^{\mathbb{D}}(\alpha, g^0 e^{i\psi}) = \frac{\partial \psi^*}{\partial \tau} + 2 \sum_j d_j \frac{\alpha_j^0 \wedge z}{|z - \alpha_j^0|^2} - 2 \sum_j d_j \frac{\alpha_j \wedge z}{|z - \alpha_j|^2}. \quad (4.5)$$

Proof of (i). Since we will always work in the unit disc, we drop the superscript \mathbb{D} .

We know from (2.24) that for $g \in H^{1/2}(\mathbb{S}^1; \mathbb{S}^1)$,

$$W(\alpha, g) = \widehat{W}(\alpha) + \frac{1}{2} \|\psi_{\alpha, g}\|_{H^{1/2}}^2, \quad (4.6)$$

where $\psi_{\alpha, g}$ is defined (modulo a constant) in (2.19) by

$$g = g^\alpha e^{i\psi_{\alpha, g}}. \quad (4.7)$$

Taking $g = g^0 e^{i\psi}$, and using $g^0 = g^\alpha e^{i\psi_{\alpha, g^0}}$, we find

$$g = g^\alpha e^{i\psi_{\alpha, g^0}} e^{i\psi} = g^\alpha e^{i(\psi_{\alpha, g^0} + \psi)}, \quad (4.8)$$

so that it holds

$$\psi_{\alpha, g} = \psi_{\alpha, g^0} + \psi. \quad (4.9)$$

This leads to

$$W(\alpha, g^0 e^{i\psi}) = \widehat{W}(\alpha) + \frac{1}{2} |\psi_{\alpha, g^0} + \psi|_{H^{1/2}}^2 = \widehat{W}(\alpha) + \frac{1}{2} \int_{\mathbb{D}} |\nabla(\psi_{\alpha, g^0}^* + \psi^*)|^2, \quad (4.10)$$

i.e., (4.3) holds. In order to complete the proof of (i), it remains to compute $\nabla \psi_{\alpha, g^0}^*$.

Recall that ψ_{α, g^0}^* is characterized by

$$\begin{cases} \Delta \psi_{\alpha, g^0}^* = 0 & \text{in } \mathbb{D} \\ \frac{\partial}{\partial \nu} \psi_{\alpha, g^0}^* = -\frac{\partial}{\partial \tau} \psi_{\alpha, g^0} & \text{on } \mathbb{S}^1. \\ \int_{\mathbb{S}^1} \psi_{\alpha, g^0}^* = 0 \end{cases} \quad (4.11)$$

Since $e^{i\psi_{\alpha, g^0}} = g^0/g^\alpha$, we have

$$\frac{\partial \psi_{\alpha, g^0}}{\partial \tau} = g^0 \wedge \frac{\partial g^0}{\partial \tau} - g^\alpha \wedge \frac{\partial g^\alpha}{\partial \tau}. \quad (4.12)$$

By definition of g^α and $g^0 = g^{\alpha^0}$, and using (4.1), we obtain

$$\begin{aligned} g^\alpha \wedge \frac{\partial g^\alpha}{\partial \tau} &= \frac{\partial \widehat{\Phi}_\alpha}{\partial \nu} = \sum_j d_j \frac{\partial}{\partial \nu} [\log |z - \alpha_j| - \log |1 - \overline{\alpha_j} z|], \\ g^0 \wedge \frac{\partial g^0}{\partial \tau} &= \frac{\partial \widehat{\Phi}_{\alpha^0}}{\partial \nu} = \sum_j d_j \frac{\partial}{\partial \nu} [\log |z - \alpha_j^0| - \log |1 - \overline{\alpha_j^0} z|]. \end{aligned} \quad (4.13)$$

We also note the identity

$$1 = \frac{\partial}{\partial \nu} [\log |1 - \overline{\alpha} z| + \log |z - \alpha|], \quad \forall \alpha \in \mathbb{D}. \quad (4.14)$$

Combining (4.12)-(4.14), we obtain

$$\frac{\partial \psi_{\alpha, g^0}^*}{\partial \nu} = -\frac{\partial \psi_{\alpha, g^0}}{\partial \tau} = \frac{\partial}{\partial \nu} \left[2 \sum_j d_j (\log |1 - \overline{\alpha_j^0} z| - \log |1 - \overline{\alpha_j} z|) \right]. \quad (4.15)$$

Therefore, there exists a constant $c(\alpha) \in \mathbb{R}$ such that

$$\psi_{\alpha, g^0}^*(z) = 2 \sum_j d_j (\log |1 - \overline{\alpha_j^0} z| - \log |1 - \overline{\alpha_j} z|) + c(\alpha), \quad \forall z \in \mathbb{D}. \quad (4.16)$$

Indeed, the right-hand side of (4.16) satisfies (4.11), and so does ψ_{α, g^0}^* . The constant $c(\alpha)$ is determined by the normalization condition $\int \psi_{\alpha, g^0}^* = 0$. From (4.16) we immediately obtain (4.4). \square

Proof of (ii). In view of formula (2.26), we have

$$N(\alpha, g^0 e^{\iota\psi}) = \frac{\partial H}{\partial \nu} + \frac{\partial}{\partial \nu} \left[\sum_j d_j \theta(z - \alpha_j) \right] = \frac{\partial H^*}{\partial \tau} - \sum_j d_j \frac{\partial}{\partial \tau} [\log|z - \alpha_j|], \quad (4.17)$$

where H^* is the harmonic conjugate of H , characterized (up to a constant) by

$$\begin{cases} \Delta H^* = 0 & \text{in } \mathbb{D} \\ \frac{\partial H^*}{\partial \nu} = -\frac{\partial H}{\partial \tau} & \text{on } \mathbb{S}^1. \end{cases} \quad (4.18)$$

On the boundary \mathbb{S}^1 , we have

$$e^{\iota H} = \prod_j \left(\frac{z - \alpha_j}{|z - \alpha_j|} \right)^{-d_j} \quad g = \prod_j \left(\frac{z - \alpha_j}{|z - \alpha_j|} \right)^{-d_j} g^0 e^{\iota\psi}, \quad (4.19)$$

so that

$$\begin{aligned} \frac{\partial H}{\partial \tau} &= \frac{\partial \psi}{\partial \tau} + g^0 \wedge \frac{\partial g^0}{\partial \tau} - \sum_j d_j \frac{\partial}{\partial \tau} [\theta(z - \alpha_j)] \\ &= -\frac{\partial \psi^*}{\partial \nu} + \frac{\partial \widehat{\Phi}_{\alpha^0}}{\partial \nu} - \sum_j d_j \frac{\partial}{\partial \nu} [\log|z - \alpha_j|] \\ &= -\frac{\partial \psi^*}{\partial \nu} + \sum_j d_j \frac{\partial}{\partial \nu} [\log|z - \alpha_j^0| - \log|1 - \bar{\alpha}_j^0 z|] - \sum_j d_j \frac{\partial}{\partial \nu} [\log|z - \alpha_j|.] \end{aligned} \quad (4.20)$$

Here we have used the definition of $g^0 = g^{\alpha^0}$ and the explicit formula (4.1) for $\widehat{\Phi}_\alpha$. Using (4.14), we obtain

$$\frac{\partial H}{\partial \tau} = -\frac{\partial}{\partial \nu} \left[\psi^* + \sum_j d_j (2 \log|1 - \bar{\alpha}_j^0 z| - \log|1 - \bar{\alpha}_j z|) \right]. \quad (4.21)$$

We deduce that there exists a constant $c = c(\psi, \alpha)$ such that

$$H^* = \psi^* + \sum_j d_j (2 \log|1 - \bar{\alpha}_j^0 z| - \log|1 - \bar{\alpha}_j z|) + c. \quad (4.22)$$

From (4.22) and (4.17) we obtain

$$N(\alpha, g^0 e^{\iota\psi}) = \frac{\partial \psi^*}{\partial \tau} + \sum_j d_j \frac{\partial}{\partial \tau} [2 \log|1 - \bar{\alpha}_j^0 z| - \log|1 - \bar{\alpha}_j z| - \log|z - \alpha_j|]. \quad (4.23)$$

Using the fact that for every $\alpha \in \mathbb{D}$ we have

$$\frac{\partial}{\partial \tau} [\log|z - \alpha|] = \frac{\partial}{\partial \tau} [\log|1 - \bar{\alpha} z|] = \frac{\alpha \wedge z}{|z - \alpha|^2}, \quad \forall z \in \mathbb{S}^1, \quad (4.24)$$

we finally obtain

$$N(\alpha, g^0 e^{\iota\psi}) = \frac{\partial \psi^*}{\partial \tau} + 2 \sum_j d_j \frac{\alpha_j^0 \wedge z}{|z - \alpha_j^0|^2} - 2 \sum_j d_j \frac{\alpha_j \wedge z}{|z - \alpha_j|^2}, \quad (4.25)$$

as claimed. \square

5 Nondegeneracy of W is stable

In this section we show that, if $a^0 \in (\Omega_0)_*^k$ is a nondegenerate critical point of $W^{\Omega_0}(\cdot, g_0)$, with $g_0 : \partial\Omega_0 \rightarrow \mathbb{S}^1$, then for Ω “close to” Ω_0 , and for $g : \partial\Omega \rightarrow \mathbb{S}^1$ “close to” g_0 , there exists a unique nondegenerate critical point a of $W^\Omega(\cdot, g)$ “close to” a^0 . Unlike the analysis we perform in subsequent sections, smoothness (of the domain or of the boundary datum) is not crucial here. In order to emphasize this fact, we first state and prove a result concerning rough boundary datum (Proposition 5.1). We next present a “smoother” variant of the stability result (Proposition 5.3).

The notion of closeness will be expressed in terms of conformal representations. Let us first introduce some definitions. Let X be the space

$$X := \{f \in C^1(\overline{\mathbb{D}}; \mathbb{C}); f \text{ is holomorphic in } \mathbb{D}\}, \quad (5.1)$$

which is a Banach space with the $\|\cdot\|_{C^1}$ norm. In X we will consider the open set

$$V := \{f \in X; f \text{ is bijective and } f^{-1} \in X\}. \quad (5.2)$$

Every $f \in V$ induces a conformal representation $f : \mathbb{D} \rightarrow \Omega := f(\mathbb{D})$, which is C^1 up to the boundary. In what follows, we denote by f^{-1} both the inverse of $f : \overline{\mathbb{D}} \rightarrow \overline{\Omega}$ and of $f|_{\mathbb{S}^1} : \mathbb{S}^1 \rightarrow \partial\Omega$.

Similar considerations apply to the space

$$X_\alpha := \{f \in C^{1,\alpha}(\overline{\mathbb{D}}; \mathbb{C}); f \text{ is holomorphic in } \mathbb{D}\}, \quad (5.3)$$

and to the open set

$$V_\alpha := \{f \in X; f \text{ is bijective and } f^{-1} \in X_\alpha\}. \quad (5.4)$$

Here, $0 < \alpha < 1$.

5.1 Proposition. *Let Ω_0 be a smooth bounded simply connected $C^{1,\alpha}$ domain and $f_0 : \mathbb{D} \rightarrow \Omega_0$ be a conformal representation. Assume that there exists $\alpha^0 \in \mathbb{D}_*^k$ and $g_0 \in H^{1/2}(\mathbb{S}^1; \mathbb{S}^1)$ such that $a^0 := f_0(\alpha^0)$ is a nondegenerate critical point of $W^{\Omega_0}(\cdot, g_0 \circ f_0^{-1})$.*

Then there exist a neighborhood \mathcal{V} of $(f_0, 0)$ in $V \times H^{1/2}(\mathbb{S}^1; \mathbb{R})$, a smooth map $\alpha : \mathcal{V} \rightarrow \mathbb{D}_^k$, and some $\delta > 0$, such that the following holds.*

Let $(f, \psi) \in \mathcal{V}$ and consider the domain $\Omega := f(\mathbb{D})$ together with the boundary datum $g := (g_0 e^{i\psi}) \circ f^{-1} \in H^{1/2}(\partial\Omega; \mathbb{S}^1)$. Then $W^\Omega(\cdot, g)$ admits a unique critical point $a \in \Omega_^k$ satisfying $|f^{-1}(a) - \alpha^0| < \delta$. This a is given by the map $a(f, \psi) = f(\alpha(f, \psi))$. Furthermore, a is a nondegenerate critical point of $W^\Omega(\cdot, g)$.*

Before proving Proposition 5.1 we state as a lemma the following smoothness result.

5.2 Lemma. *The map $\widetilde{W} : \mathbb{D}_*^k \times V \times H^{1/2}(\mathbb{S}^1; \mathbb{R}) \rightarrow \mathbb{R}$, defined by*

$$\widetilde{W}(\alpha, f, \psi) = W^{f(\mathbb{D})}(f(\alpha), (g_0 e^{i\psi}) \circ f^{-1}), \quad (5.5)$$

is smooth.

Similarly, the map $\widetilde{W}_\alpha : \mathbb{D}_^k \times V_\alpha \times C^{1,\alpha}(\mathbb{S}^1; \mathbb{R}) \rightarrow \mathbb{R}$, defined by*

$$\widetilde{W}_\alpha(\alpha, f, \psi) = W^{f(\mathbb{D})}(f(\alpha), (g_0 e^{i\psi}) \circ f^{-1}), \quad (5.6)$$

is smooth.

Proof of Lemma 5.2. The idea is to rely on the formulas derived in Sections 3 and 4 in order to obtain an explicit formula for \widetilde{W} , from which it will be clear that \widetilde{W} is smooth.

To start with, formula (3.8) gives

$$\widetilde{W}(\alpha, f, \psi) = W^{\mathbb{D}}(\alpha, g_0 e^{i\psi}) + \pi \sum_j d_j^2 \log |f'(\alpha_j)|. \quad (5.7)$$

Using the fact that for holomorphic f all derivatives can be estimated locally using only $\|f\|_\infty$, it can be easily shown that the maps

$$\mathbb{D}_*^k \times V \ni (\alpha, f) \mapsto \log |f'(\alpha_j)| \quad (5.8)$$

are smooth.

Therefore the second term in the right-hand side of (5.7) is smooth, and in order to complete the proof of Lemma 5.2 it suffices to prove that

$$\mathbb{D}_*^k \times H^{1/2}(\mathbb{S}^1; \mathbb{R}) \ni (\alpha, \psi) \mapsto W^{\mathbb{D}}(\alpha, g_0 e^{i\psi}) := P_{g_0}(\alpha, \psi) \quad (5.9)$$

is smooth. Clearly, if $g \in H^{1/2}(\mathbb{S}^1; \mathbb{S}^1)$ is such that $\deg(g, \mathbb{S}^1) = \deg(g_0, \mathbb{S}^1)$, then we may write $g = g_0 e^{i\psi_0}$ for some $\psi_0 \in H^{1/2}(\mathbb{S}^1; \mathbb{R})$, and then we have $P_g(\alpha, \psi) = P_{g_0}(\alpha, \psi + \psi_0)$. This implies that the smoothness of P_{g_0} does not depend on the choice of g_0 . Therefore, we may assume that $g_0 = g^{\alpha^0}$ for some $\alpha^0 \in \mathbb{D}_*^k$. This assumption allows us to use Lemma 4.1. Using (4.3), we obtain

$$\begin{aligned} W^{\mathbb{D}}(\alpha, g_0 e^{i\psi}) &= \widehat{W}^{\mathbb{D}}(\alpha) + \frac{1}{2} \int_{\mathbb{D}} |\nabla \psi_{\alpha, g_0}^*|^2 + \frac{1}{2} \int_{\mathbb{D}} |\nabla \psi^*|^2 \\ &\quad + \int_{\mathbb{D}} \nabla \psi_{\alpha, g_0}^* \cdot \nabla \psi^*. \end{aligned} \quad (5.10)$$

We examine the smoothness of the four terms on the right-hand side of (5.10). The first term depends only on α and is smooth thanks to formula (4.2). The second term depends only on α and is smooth thanks to formula (4.4). The third term depends only on ψ and is a continuous quadratic form, hence it is smooth. The fourth and last term depends linearly on ψ and is smooth thanks to formula (4.4) again.

Hence the map (5.9) is smooth, and the proof of the $H^{1/2}$ part of the lemma is complete.

The proof of the $C^{1,\alpha}$ part of the follows the same lines and is left to the reader. \square

Proof of Proposition 5.1. Let us first remark the following fact. Fix $f \in V$ and $\psi \in H^{1/2}(\partial\mathbb{D}; \mathbb{R})$ and consider the domain $\Omega = f(\mathbb{D})$ together with the boundary datum $g = (g_0 e^{i\psi}) \circ f^{-1}$. Then, for any $\alpha \in \mathbb{D}_*^k$, $f(\alpha)$ is a nondegenerate critical point of $W^\Omega(\cdot, g)$ if and only if α is a nondegenerate critical point of $\widetilde{W}(\cdot, f, \psi)$. This is a simple consequence of the fact that f induces a diffeomorphism from \mathbb{D}_*^k to Ω_*^k .

We consider the map $F : \mathbb{D}_*^k \times V \times H^{1/2}(\mathbb{S}^1; \mathbb{R}) \rightarrow \mathbb{R}^{2k}$,

$$F : (\alpha, f, \psi) \mapsto \nabla_\alpha \widetilde{W}(\alpha, f, \psi). \quad (5.11)$$

Lemma 5.2 ensures that F is smooth. Moreover, the assumption that α^0 is a nondegenerate critical point of $W^{\Omega_0}(\cdot, g_0 \circ f^{-1})$ ensures that α^0 is a nondegenerate critical point of $\widetilde{W}(\cdot, f_0, 0)$. Therefore $F(\alpha^0, f_0, 0) = 0$, and $D_\alpha F(\alpha^0, f_0, 0)$ is invertible.

This enables us to apply the implicit function theorem: there exist of a neighborhood \mathcal{V} of $(f_0, 0)$ in $V \times H^{1/2}(\mathbb{S}^1; \mathbb{R})$, a smooth map $\alpha : \mathcal{V} \rightarrow \mathbb{D}_*^k$, and $\delta > 0$, such that, for $(f, \psi) \in \mathcal{V}$ and $|\alpha - \alpha^0| < \delta$,

$$F(\alpha, f, \psi) = 0 \iff \alpha = \alpha(f, \psi). \quad (5.12)$$

We may also assume that $D_\alpha F(\alpha(f, \psi), f, \psi)$ is invertible, so that $\alpha(f, \psi)$ is a nondegenerate critical point of $\widetilde{W}(\cdot, f, \psi)$. This implies that $a := f(\alpha(f, \psi))$ is a nondegenerate critical point of $W^\Omega(\cdot, g)$, where $\Omega = f(\mathbb{D})$ and $g = (g_0 e^{i\psi}) \circ f^{-1}$. In view of (5.12), a is the unique critical point of $W^\Omega(\cdot, g)$ satisfying $|f^{-1}(a) - \alpha^0| < \delta$.

The proof of Proposition 5.1 is complete. \square

In what follows, we will use the following smoother version of Proposition 5.1.

5.3 Proposition. *Let Ω_0 be a smooth bounded simply connected $C^{1,\alpha}$ domain and $f_0 : \mathbb{D} \rightarrow \Omega_0$ be a conformal representation. Assume that there exists $\alpha^0 \in \mathbb{D}_*^k$ and $g_0 \in C^{1,\alpha}(\mathbb{S}^1; \mathbb{S}^1)$ such that $\alpha^0 := f_0(\alpha^0)$ is a nondegenerate critical point of $W^{\Omega_0}(\cdot, g_0 \circ f_0^{-1})$.*

Then there exist a neighborhood \mathcal{V} of $(f_0, 0)$ in $V_\alpha \times C^{1,\alpha}(\mathbb{S}^1; \mathbb{R})$, a smooth map $\alpha : \mathcal{V} \rightarrow \mathbb{D}_^k$, and some $\delta > 0$, such that the following holds.*

Let $(f, \psi) \in \mathcal{V}$ and consider the domain $\Omega := f(\mathbb{D})$ together with the boundary datum $g := (g_0 e^{i\psi}) \circ f^{-1} \in C^{1,\alpha}(\partial\Omega; \mathbb{S}^1)$. Then $W^\Omega(\cdot, g)$ admits a unique critical point $a \in \Omega_^k$ satisfying $|f^{-1}(a) - \alpha^0| < \delta$, given by the map $a(f, \psi) = f(\alpha(f, \psi))$. Furthermore, a is a nondegenerate critical point of $W^\Omega(\cdot, g)$.*

Here, V_α is given by (5.4). The proof of Proposition 5.3 is identical to the one of Proposition 5.1 and is left to the reader.

We will need later the following special case of Proposition 5.3, where Ω is fixed.

5.4 Corollary. *Let $a^0 \in \Omega_*^k$ be a nondegenerate critical point of $W(\cdot, g_0)$, for some $g_0 \in C^{1,\alpha}(\partial\Omega; \mathbb{S}^1)$. Then, for g in a small $C^{1,\alpha}$ -neighborhood \mathfrak{A} of g_0 , $W(\cdot, g)$ has, near a^0 , a unique nondegenerate critical point $a(g)$. In addition, the map $\psi \mapsto a(g_0 e^{i\psi})$, defined for ψ in a sufficiently small neighborhood of the origin in $C^{1,\alpha}(\partial\Omega; \mathbb{R})$, is smooth.*

We note here that Corollary 5.4 allows us to define a map

$$T_* = T_{*,a^0,g_0}^\Omega : \mathfrak{A} \rightarrow C^\alpha(\partial\Omega; \mathbb{R}), \quad T_*(g) := N^\Omega(a(g), g) = u_{*,a(g),g} \wedge \frac{\partial u_{*,a(g),g}}{\partial v}. \quad (5.13)$$

Since $W^\Omega(\cdot, g)$ does not depend on the class of g modulo \mathbb{S}^1 , neither do $a(g)$ and T_* . Moreover, in view of (2.28) and (2.29) we have

$$\int_{\partial\Omega} u_{*,a,g} \wedge \frac{\partial u_{*,a,g}}{\partial v} = \int_{\partial\Omega} \frac{\partial \Phi_{a,g}}{\partial \tau} = 0.$$

We find that the map T_* induces a map, still denoted T_* , from $\mathfrak{A}/\mathbb{S}^1$ into $\dot{C}^\alpha(\partial\Omega; \mathbb{R})$. Here, we define

$$\dot{C}^\alpha(\partial\Omega; \mathbb{R}) := \left\{ \psi \in C^\alpha(\partial\Omega; \mathbb{R}); \int_{\partial\Omega} \psi = 0 \right\}.$$

It is also convenient to consider, in a sufficiently small neighborhood \mathfrak{B} of the origin in $C^{1,\alpha}(\partial\Omega; \mathbb{R})$, the maps (both denoted U_*)

$$U_* = U_{*,a^0,g_0}^\Omega : \mathfrak{B} \rightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R}), \quad U_*(\psi) = T_*(g_0 e^{i\psi}) \quad (5.14)$$

and

$$U_* = U_{*,a^0,g_0}^\Omega : \mathfrak{B}/\mathbb{R} \rightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R}), \quad U_*(\psi) = T_*(g_0 e^{i\psi}). \quad (5.15)$$

The above U_* 's are smooth. Indeed, this is obtained by combining (3.7) with (4.5) and with the fact that $\psi \mapsto a(g_0 e^{i\psi})$ is smooth.

6 A uniform version of the Pacard-Rivière construction of critical points

We start by explaining how the results established in this section compare to the existent literature.

Let us first briefly recall the Bethuel-Brezis-Hélein analysis of critical points of the Ginzburg-Landau energy E_ε with prescribed Dirichlet boundary condition $g \in C^\infty(\partial\Omega; \mathbb{S}^1)$ [8, Chapter X]. Consider a fixed boundary condition $g \in C^\infty(\partial\Omega; \mathbb{S}^1)$, with $\deg(g, \partial\Omega) = d = \sum_{j=1}^k d_j$. Given $a \in \Omega_*^k$ of $W(\cdot, d_1, \dots, d_k, g)$, consider the canonical harmonic map given by (2.26). The fact that a is a

critical point of $W(\cdot, d_1, \dots, d_k, g)$ is equivalent to the fact that the harmonic function H_j , defined near a_j by

$$u_* = e^{iH_j} \left(\frac{z - a_j}{|z - a_j|} \right)^{d_j}, \quad (6.1)$$

satisfies $\nabla H_j(a_j) = 0$ [8, Chapter VII].

The main result in [8, Chapter X] asserts that, when Ω is starshaped critical points of E_ε converge, as $\varepsilon \rightarrow 0$ (up to subsequences and in appropriate function spaces), to a canonical harmonic map $u_* = u_{*,a,g}$ associated with a critical point $a \in \Omega_*^k$ of $W(\cdot, d_1, \dots, d_k, g)$.

Granted this result, one can address the converse: given a critical point $a \in \Omega_*^k$ of $W(\cdot, d_1, \dots, d_k, g)$, does there exist critical points u_ε of E_ε with prescribed boundary condition g , such that $u_\varepsilon \rightarrow u_{*,a,g}$ as $\varepsilon \rightarrow 0$? Here we will be interested in the answer provided by Pacard and Rivière [16].

6.1 Theorem ([16, Theorem 1.4]). *Assume that $g \in C^{2,\alpha}(\partial\Omega; \mathbb{S}^1)$ and $d_j \in \{\pm 1\}$. Let $a \in \Omega_*^k$ be a nondegenerate critical point of $W(\cdot, d_1, \dots, d_k, g)$.*

Then there exists $\varepsilon_0 > 0$ such that for every $\varepsilon \in (0, \varepsilon_0)$, there exists u_ε a critical point of E_ε with $u_\varepsilon = g$ on $\partial\Omega$, and

$$u_\varepsilon \rightarrow u_{*,a(g),g} \quad \text{as } \varepsilon \rightarrow 0 \quad (6.2)$$

in $C_{loc}^{2,\alpha}(\Omega \setminus \{a_1, \dots, a_k\})$.

The purpose of this section is to establish a variant of Theorem 6.1, in which g is assumed to be merely $C^{1,\alpha}$ and is not fixed anymore. In addition, we will obtain a uniform existence theorem, and uniform convergence rate. More specifically, we fix integers d_1, \dots, d_k . Since these integers do not depend on the boundary datum g we consider, we will omit the dependence of W with respect to d_1, \dots, d_k : we write $W(\cdot, g)$ instead of $W(\cdot, d_1, \dots, d_k, g)$. We consider $a_0 \in \Omega_*^k$ a nondegenerate critical point of the renormalized energy $W(\cdot, g_0)$ associated with $g_0 \in C^{1,\alpha}(\partial\Omega; \mathbb{S}^1)$. By Corollary 5.4, we know that, for g in a small $C^{1,\alpha}$ -neighborhood \mathfrak{A} of g_0 , $W(\cdot, g)$ has, near a_0 , a unique nondegenerate critical point $a(g)$.

In this section, we establish the following variant of Theorem 6.1.

6.2 Theorem. *Let $d_1, \dots, d_k \in \{-1, 1\}$. Let a_0 be a nondegenerate critical point of $W(\cdot, g_0)$. Then there exist $\delta > 0$ and $\varepsilon_0 > 0$ such that the following holds. For every $g \in C^{1,\alpha}(\partial\Omega; \mathbb{S}^1)$ satisfying $\|g - g_0\|_{C^{1,\alpha}} \leq \delta$, and for every $\varepsilon \in (0, \varepsilon_0)$, there exists $u_\varepsilon = u_{\varepsilon,g}$ a critical point of E_ε with prescribed boundary condition g , such that*

$$u_{\varepsilon,g} \rightarrow u_{*,a(g),g} \quad \text{as } \varepsilon \rightarrow 0 \quad (6.3)$$

in $C_{loc}^{2,\alpha}(\Omega \setminus \{a_1, \dots, a_k\})$.

As announced, the difference with Theorem 1.4 in [16] is that we merely assume that $g \in C^{1,\alpha}$; in addition, we prove that ε_0 can be chosen independent of g . Theorem 6.2 allows us to define a map $F_\varepsilon : g \mapsto u_{\varepsilon,g}$ for every $\varepsilon \in (0, \varepsilon_0)$.

Theorem 6.2 is obtained by following the proof of Theorem 6.1 in [16]. All we have to check (and we will do in what follows) is that the estimates in [16] are uniform in g ; we also have to modify some arguments relying on the regularity assumption $g \in C^{2,\alpha}$.

Proof of Theorem 6.2. For the convenience of the reader, we recall the main steps of the proof of Theorem 6.1 in [16], and examine the crucial points where the estimates depend on g , respectively where the regularity of g plays a role.

The general strategy in [16] is to construct an “approximate solution” \tilde{u}_ε of the Ginzburg Landau equation

$$\mathcal{N}_\varepsilon(u) = 0, \quad \text{where } \mathcal{N}_\varepsilon(u) := \Delta u + \frac{u}{\varepsilon^2}(1 - |u|^2), \quad (6.4)$$

using the fairly precise knowledge we have of the form of solutions for small ε . Then, using a fixed point argument, one can prove that some perturbation of \tilde{u}_ε is in fact an exact solution of (6.4). The main difficulty lies in finding the good functional setting that makes the linearized operator $L_\varepsilon = D\mathcal{N}_\varepsilon$ around \tilde{u}_ε invertible, uniformly with respect to ε . This is achieved in [16] in the frame of appropriate weighted Hölder spaces.

In [16] the proof of Theorem 6.1 is divided into five chapters: Chapters 3 through 7. In what follows, we detail the content of these chapters and explain how to adapt the arguments for the need of Theorem 6.2.

Chapters 3 and 4 in [16]

[16, Chapter 3] is devoted to the study of the radially symmetric solution $u(re^{i\theta}) = f(r)e^{i\theta}$ of the Ginzburg-Landau equation in \mathbb{C} satisfying $\lim_{r \rightarrow \infty} f(r) = 1$. In particular, [16, Chapter 3] characterizes the bounded solutions of the linearized equation about this radial solution. This characterization is used in [16, Chapter 4] in the study of the mapping properties of the linearization of the Ginzburg-Landau operator (at the radial solution) in the punctured unit disc $\mathbb{D} \setminus \{0\}$. In particular, it is shown that the linearized operator is invertible between appropriate weighted Hölder spaces.

These two chapters (3 and 4) are independent of the boundary condition g , so that they can be used with no changes in the proof of Theorem 6.2.

Chapter 5 in [16]

The next step, in [16, Chapter 5], consists in constructing and estimating the approximate solution \tilde{u}_ε . This approximate solution looks like $u_* = u_{*,g,a(g)}$ away from its zeros (which are close to the singularities of u_*), and like the

radial solution studied in [16, Chapter 3] near its zeros. Since \tilde{u}_ε is built upon u_* , the estimates satisfied by \tilde{u}_ε involve u_* , and thus g .

More specifically, in [16, Chapter 5], various quantities are estimated in terms of constants $c(u_*)$ depending on u_* and its derivatives. An inspection of the proofs there combined with (2.26) shows that these constants depend only on $a(g)$, on the harmonic function $H = H_g$ and on the derivatives of H_g .

We claim that the constants $c(u_*)$ can be chosen independent of g satisfying

$$\|g - g_0\|_{C^{1,\alpha}(\partial\Omega)} \leq \delta. \quad (6.5)$$

Here, δ is sufficiently small in order to have the conclusion of Corollary 5.4. Indeed, the key observation is that there exists a constant $C > 0$ independent of g such that

$$\|H\|_{C^{1,\alpha}(\bar{\Omega})} \leq C; \quad (6.6)$$

this follows from the fact that H is harmonic and $\|H\|_{C^{1,\alpha}(\partial\Omega)} \leq C$.

In particular, we have

$$\|H\|_{C^k(\omega)} \leq C(k, \omega) \quad \text{for } k \in \mathbb{N} \text{ and } \bar{\omega} \subset \Omega. \quad (6.7)$$

Estimate (6.7) implies that all the interior estimates in [16, Chapter 5] are satisfied uniformly in $g \in C^{1,\alpha}$ satisfying (6.5). This settles the case of estimates (5.8), (5.9), (5.33), (5.42) and (5.43) in [16, Chapter 5].

It remains to consider the global and boundary estimates (5.29), (5.32) and (5.41) in [16]. These estimates rely on bounds on the solution ξ of the problem

$$\begin{cases} \Delta \xi - |\nabla u_*|^2 \xi + \frac{1-\xi^2}{\varepsilon^2} \xi = 0 & \text{in } \Omega_{\delta/2} \\ \xi = S_\varepsilon + w_{j,r} & \text{on } \partial\mathbb{D}_{\delta/2}(a_j) \\ \xi = 1 & \text{on } \partial\Omega \end{cases}. \quad (6.8)$$

Here, $\Omega_\sigma := \Omega \setminus \bigcup_j \mathbb{D}_\sigma(a_j)$ (for sufficiently small $\sigma > 0$), and $\delta := \varepsilon^2$. The auxiliary function S_ε is independent of g and is defined in [16, Section 3.6]. Finally, $w_{j,r}$, defined in [16, (5.7)], depends only $a(g)$ and on the restriction of H_j to compacts of Ω ; therefore, the estimates involving $w_{j,r}$ are uniform in g .

In [16, Lemma 5.1] the following estimates (numbered as (5.29) in [16]) are shown to hold:

$$1 - c\varepsilon^2 \leq \xi \leq 1 \quad \text{in } \Omega_\sigma, \quad (6.9)$$

$$1 - c\varepsilon^2 r_j^{-2} \leq \xi \leq 1 \quad \text{in } \mathbb{D}_\sigma(a_j) \setminus \mathbb{D}_{\delta/2}(a_j), \quad (6.10)$$

$$|\nabla^k \xi| \leq c_k \varepsilon^2 r_j^{-2-k} \quad \text{in } \mathbb{D}_{2\sigma}(a_j) \setminus \mathbb{D}_\delta(a_j) \quad (k \geq 1). \quad (6.11)$$

Here $\sigma > 0$ is fixed, and $r_j = r_j(x)$ denotes the distance from x to a_j . Estimates (6.10) and (6.11) are interior estimates, and therefore they hold uniformly in

$g \in C^{1,\alpha}$ satisfying (6.5), as explained above. We claim that the same conclusion applies to (6.9). Indeed, an inspection of the proof in [16] shows that the constant c in (6.9) is controlled by $\sup_{\Omega_\sigma} |\nabla u_*|$. The latter quantity is uniformly bounded, thanks to (6.6), whence the conclusion. This settles the case of the estimate (5.29) in [16].

We next turn to the estimate (5.32) in [16, Lemma 5.2]. Under the assumption that $g \in C^{2,\alpha}$, this lemma asserts that

$$\sup_{\Omega_\sigma} |\nabla^k \xi| \leq c\varepsilon^{2-k}, \quad k = 1, 2. \quad (6.12)$$

In our case, we only assume $g \in C^{1,\alpha}$. The corresponding estimates are given by our next result.

6.3 Lemma. *Assume that (6.5) holds. Then we have*

$$\sup_{\Omega_\sigma} |\nabla \xi| \leq c\varepsilon \text{ and } |\nabla \xi|_{\alpha, \Omega_\sigma} \leq c\varepsilon^{1-\alpha}. \quad (6.13)$$

Here, $|\cdot|_{\alpha, \Omega_\sigma}$ denotes the C^α semi-norm in Ω_σ :

$$|u|_{\alpha, \Omega_\sigma} := \sup \left\{ \frac{|u(x) - u(y)|}{|x - y|^\alpha}; x, y \in \Omega_\sigma \right\}.$$

Proof. We apply Lemma 10.3 in the Appendix with $w = \xi - 1$ in $G := \Omega_{\sigma/2}$, and find that

$$\sup_{\Omega_\sigma} |\nabla \xi| \leq C \left(\|w\|_{L^\infty(\Omega_\sigma)}^{1/2} \|\Delta w\|_{L^\infty(\Omega_\sigma)}^{1/2} + \|w\|_{C^{1,\alpha}(\partial\Omega_\sigma)} \right), \quad (6.14)$$

$$|\nabla \xi|_{\alpha, \Omega_\sigma} \leq C \left(\|w\|_{L^\infty(\Omega_\sigma)}^{1/2-\alpha/2} \|\Delta w\|_{L^\infty(\Omega_\sigma)}^{1/2+\alpha/2} + \|w\|_{C^{1,\alpha}(\partial\Omega_\sigma)} \right). \quad (6.15)$$

The conclusion then follows by combining (6.14)-(6.15) with the equation (6.8) and with estimates (6.9) and (6.11). \square

Finally, we examine estimate (5.41) in the last section of [16, Chapter 5]; this is a global estimate for $\mathcal{N}_\varepsilon(\tilde{u}_\varepsilon)$. Recall here that \mathcal{N}_ε is the Ginzburg-Landau operator, and that \tilde{u}_ε is the approximate solution of (6.4) constructed in [16, Chapter 5]. The estimate [16, (5.41)] involves an interior estimate and a boundary estimate. As above, the interior estimate is settled with the help of (6.7). We now turn to the boundary estimate, which is the following:

$$\|\mathcal{N}_\varepsilon(\tilde{u}_\varepsilon)\|_{C^\alpha(\Omega_\sigma)} \leq c\varepsilon^{1-\alpha}. \quad (6.16)$$

The proof of (6.16) in [16] relies on the estimates (6.12) above (see [16, Proof of Lemma 5.2]). In our case, (6.12) need not hold, since we only assume that $g \in C^{1,\alpha}$. However, we still obtain (6.16) as follows. We note that

$$\mathcal{N}_\varepsilon(\tilde{u}_\varepsilon) = 2\nabla u_* \cdot \nabla \xi \quad \text{in } \Omega_\sigma \quad (6.17)$$

(this is formula (5.46) in [16]). By (6.17), we have

$$\|\mathcal{N}_\varepsilon(\tilde{u}_\varepsilon)\|_{C^\alpha(\Omega_\sigma)} \leq c \|\nabla u_*\|_{C^\alpha(\Omega_\sigma)} \|\nabla \xi\|_{C^\alpha(\Omega_\sigma)}. \quad (6.18)$$

We obtain (6.16) as a consequence of (6.6) and of Lemma 6.3.

As a conclusion of this inspection, we find that all the estimates in [16, Chapter 5] are uniform in g satisfying (6.5); the arguments there need only minor changes. The most relevant change is that [16, Lemma 5.2] has to be replaced by Lemma 6.3.

Chapter 6 in [16]

We now turn to [16, Chapter 6], which deals with the conjugate linearized operator $\tilde{\mathcal{L}}_\varepsilon$ around the approximate solution. The main result of this chapter is [16, Theorem 6.1], which states that $\tilde{\mathcal{L}}_\varepsilon$ is invertible for $\varepsilon \in (0, \varepsilon_0)$, with the norm of its inverse bounded independently of ε . In order to adapt this theorem to our situation, we need to check that this ε_0 , and the bound on $\tilde{\mathcal{L}}_\varepsilon^{-1}$, can be chosen independently of g satisfying (6.5).

The proof of [16, Theorem 6.1] is divided into three parts:

- (a) The “interior” problem, consisting in the study of the linearized operator $\tilde{\mathcal{L}}_\varepsilon$ near the zeros of \tilde{u}_ε [16, Section 6.2].
- (b) The “exterior” problem, requiring the study of the linearized operator $\tilde{\mathcal{L}}_\varepsilon$ away from the zeros of \tilde{u}_ε [16, Section 6.3].
- (c) The study of the Dirichlet-to-Neumann mappings [16, Section 6.4]. (These mappings are used later in order to “glue” the two first steps together.)

The interior and the exterior problem rely on the estimates obtained in [16, Chapter 5]. An inspection of the proofs shows that all the estimates obtained there are uniform in g , with one possible exception: the estimates in [16, Proposition 6.2]. Indeed, these estimates rely on [16, Lemma 5.2], and more specifically on (6.12) (which does not hold in our setting). However, a closer look to [16, Proof of Proposition 6.2] shows that the conclusion of [16, Proposition 6.2] still holds if we replace (6.12) by Lemma 6.3. In conclusion, the first two steps can be carried out with uniform estimates, provided (6.5) holds.

The third step (Dirichlet-to-Neumann mappings) requires more care. In [16, Section 6.4], the following two operators are defined, for fixed small $\zeta > 0$ and for sufficiently small ε :

$$DN_{int,\varepsilon}, DN_{ext,\varepsilon} : \prod_{j=1}^k C^{2,\alpha}(C(\zeta, a_j)) \longrightarrow \prod_{j=1}^k C^{1,\alpha}(C(\zeta, a_j)). \quad (6.19)$$

(These are the interior and exterior Dirichlet-to-Neumann mappings.) The crucial result in part (c) is [16, Proposition 6.5], which states the existence of

some ε_0 such that $DN_{int,\varepsilon} - DN_{ext,\varepsilon}$ is an isomorphism for $\varepsilon \in (0, \varepsilon_0)$. The proof of this fact goes as follows. First the convergence

$$DN_{int,\varepsilon} - DN_{ext,\varepsilon} \longrightarrow DN_{int,0} - DN_{ext,0} \quad \text{as } \varepsilon \rightarrow 0 \quad (6.20)$$

is shown to hold in operator norm. The proof of (6.20) relies on the interior estimate (6.7). Therefore, the convergence in (6.20) is uniform in g satisfying (6.5).

We now return to the proof in [16, Chapter 6]. Once (6.20) is established, it remains to prove that the limiting operator $DN_{int,0} - DN_{ext,0}$ is invertible. This is done in [16, Proposition 6.5]; this is where the nondegeneracy of a as a critical point of $W(\cdot, g)$ comes into the picture. In order to extend the conclusion of [16, Proposition 6.5] to our setting, and to obtain a uniform bound for the inverse of $DN_{int,\varepsilon} - DN_{ext,\varepsilon}$, it suffices to check that $DN_{int,0} - DN_{ext,0}$ depends continuously on g . Indeed, this will lead to a uniform bound for the inverse of $DN_{int,\varepsilon} - DN_{ext,\varepsilon}$ provided ε is sufficiently small, uniformly in g satisfying (6.5) (possibly with a smaller δ). For this purpose, we examine the formulas of $DN_{ext,0}$ and $DN_{int,0}$. The definition of $DN_{ext,0}$ is given in [16, Proposition 6.4], and it turns out that that $DN_{ext,0}$ does not depend on g . As for $DN_{int,0}$, it is a diagonal operator of the form

$$DN_{int,0}(\phi_1, \dots, \phi_k) = \left(DN_{int,0}^1(\phi_1), \dots, DN_{int,0}^k(\phi_k) \right), \quad (6.21)$$

with $DN_{int,0}^j : C^{2,\alpha}(C(\zeta, a_j)) \rightarrow C^{1,\alpha}(C(\zeta, a_j))$, $\forall j \in \llbracket 1, k \rrbracket$.

Furthermore, from [16, Proposition 6.3] we know that $DN_{int,0}^j$ further splits as

$$DN_{int,0}^j = T_1 \oplus T_2, \quad \text{with } \begin{cases} T_1 : \text{span}\{e^{\pm in\theta}\}_{n \geq 2} \rightarrow \text{span}\{e^{\pm in\theta}\}_{n \geq 2} \\ T_2 : \text{span}\{e^{\pm i\theta}\} \rightarrow \text{span}\{e^{\pm i\theta}\} \end{cases} . \quad (6.22)$$

Here, the operator T_1 does not depend on g . Therefore, we only need to check that T_2 depends continuously on g . As a linear operator on a two-dimensional space, T_2 is represented by a 2×2 matrix. It is clear from [16, Proposition 6.3] that the coefficients of this matrix are smooth functions of $\nabla^2 H(a_j)$. In turn, $\nabla^2 H(a_j)$ depends smoothly on g , by Corollary 5.4.

Hence $DN_{int,0} - DN_{ext,0}$ depends continuously on g , as claimed.

This allows us to choose ε_0 independent of g satisfying (6.5) in [16, Proposition 6.5] and in [16, Theorem 6.1], and to obtain a uniform estimate for the inverse of $\widetilde{\mathcal{L}}_\varepsilon$.

Chapter 7 in [16]

Finally, in [16, Chapter 7] the results and estimates in [16, Chapters 3-6] are combined in order to prove Theorem 6.1. Our above analysis shows that these estimates are uniform, and therefore lead to the uniform version Theorem 6.2 of Theorem 6.1.

Conclusion

As a conclusion of our analysis, Theorem 6.2 holds. \square

For further use, we record two additional properties of the maps $u_{\varepsilon,g}$; these properties are immediate consequences of the construction in [16]. Let δ be as in Theorem 6.2. We consider the set

$$\mathfrak{A} := \{g \in C^{1,\alpha}(\partial\Omega; \mathbb{S}^1); \|g - g_0\|_{C^{1,\alpha}} < \delta\}. \quad (6.23)$$

6.4 Lemma. *Let $K \Subset \overline{\Omega} \setminus \{a_j\}$. Then we have $|u_{\varepsilon,g}| \rightarrow 1$ as $\varepsilon \rightarrow 0$, uniformly in K and in $g \in \mathfrak{A}$.*

Proof. This follows by an inspection of the construction in [16]. Formulas (5.36) and (5.37) in [16] ensure that, for small ε , the approximate solution \tilde{u}_ε satisfies $|\tilde{u}_\varepsilon| = |\xi|$ in K . The convergence then follows from the estimates on ξ , and from formula (7.1) in [16] connecting the approximate solution to the exact solution. \square

For the next result, it may be necessary to replace δ by a smaller value.

6.5 Lemma. *Let $g \in \mathfrak{A}$ and $\omega \in \mathbb{S}^1$. If $\omega g \in \mathfrak{A}$, then $u_{\varepsilon,\omega g} = \omega u_{\varepsilon,g}$.*

Proof. We have $W(\cdot, g) = W(\cdot, \omega g)$. Therefore, if a is a nondegenerate critical point of $W(\cdot, g)$, then a is also a nondegenerate critical point of $W(\cdot, \omega g)$. By Corollary 5.4, we find that $a(\omega g) = a(g)$. Using this equality, an inspection of the construction in [16] shows that

$$\tilde{u}_{\varepsilon,\omega g} = \omega \tilde{u}_{\varepsilon,g}. \quad (6.24)$$

Thanks to (6.24), we obtain that $\omega u_{\varepsilon,g}$ has all the properties satisfied by the solution $u_{\varepsilon,\omega g}$ constructed from $\tilde{u}_{\varepsilon,\omega g}$ via the inverse function theorem. Since the solution provided by the inverse function theorem is unique, we find that $u_{\varepsilon,\omega g} = \omega u_{\varepsilon,g}$, as claimed. \square

7 Convergence of the normal differentiation operators

In this section, we fix integers $d_1, \dots, d_k \in \{-1, 1\}$ as in Section 6. We assume that a^0 is a nondegenerate critical point of $W(\cdot, g_0)$. Let $g \in \mathfrak{A}$, where \mathfrak{A} is given by (6.23), and let $0 < \varepsilon < \varepsilon_0$. For such g and ε , we define $u_\varepsilon = u_{\varepsilon,g}$ as in Section 6. We also define $u_{*,g} := u_{*,a(g),g}$, where $a(g)$ is the unique critical point of $W(\cdot, g)$ close to a_0 (see Corollary 5.4). We consider the operators

$$T_\varepsilon, T_* : \mathfrak{A} \rightarrow C^\alpha(\partial\Omega; \mathbb{R}), \quad T_\varepsilon(g) := u_{\varepsilon,g} \wedge \frac{\partial u_{\varepsilon,g}}{\partial \nu} \quad \text{and} \quad T_*(g) := u_{*,g} \wedge \frac{\partial u_{*,g}}{\partial \nu}.$$

The main result of this section is the following

7.1 Proposition. *Let $0 < \beta < 1$. Then (possibly after replacing δ by a smaller number) we have*

$$\limsup_{\varepsilon \rightarrow 0} \sup_{g \in \mathfrak{A}} \|T_\varepsilon(g) - T_*(g)\|_{C^\beta(\partial\Omega)} = 0. \quad (7.1)$$

In particular, given $\mu > 0$ there exists some $\varepsilon_0 > 0$ such that, for $0 < \varepsilon < \varepsilon_0$, $T_\varepsilon - T_ : \mathfrak{A} \rightarrow C^\alpha(\partial\Omega; \mathbb{R})$ is compact and satisfies*

$$\|T_\varepsilon(g) - T_*(g)\|_{C^\alpha(\partial\Omega)} \leq \mu, \quad \forall \varepsilon < \varepsilon_0, \forall g \in \mathfrak{A}.$$

Proof. The last part of the proposition follows from the fact that the embedding $C^\beta(\partial\Omega; \mathbb{R}) \hookrightarrow C^\alpha(\partial\Omega; \mathbb{R})$ is compact when $\beta > \alpha$.

Whenever needed in the proof, we will replace δ by a smaller number. Let $a = a(g)$, $g \in \mathfrak{A}$, be such that $\nabla_a W(a, g) = 0$ and a is close to $a^0 = (a_1^0, \dots, a_k^0)$. Let $t > 0$ be a small number and set

$$\omega := \left\{ x \in \Omega; |x - a_j^0| > t, \forall j \in \llbracket 1, k \rrbracket \right\}.$$

We may assume that $|a(g) - a^0| < t/2$, $\forall g \in \mathfrak{A}$. In view of Theorem 6.2, we have $u_{\varepsilon, g} \rightarrow u_{*, g}$ in $C^{2, \beta}(K)$ as $\varepsilon \rightarrow 0$, for every $g \in \mathfrak{A}$ and for every K compact set such that $K \subset \bar{\omega} \setminus \partial\Omega$. In addition, by Lemma 6.4 we have $|u_{\varepsilon, g}| \rightarrow 1$ as $\varepsilon \rightarrow 0$ uniformly in $\bar{\omega}$ and in $g \in \mathfrak{A}$.

Let $\theta = \theta_g$ be the multi-valued argument of

$$z \mapsto \prod_{j=1}^k \frac{(z - a_j(g))^{d_j}}{|z - a_j(g)|^{d_j}}.$$

We note that $\nabla\theta_g$ is single-valued and that we have

$$\|\nabla\theta_g\|_{C^{1, \alpha}(\omega)} \leq C, \quad \forall g \in \mathfrak{A}. \quad (7.2)$$

For small ε (independent of g), we have $\deg(u_{\varepsilon, g}) = \deg(u_{*, g}) = d_j$ on $C(a_j^0, t)$, and thus we may write, locally in $\bar{\omega}$,

$$u_{\varepsilon, g} = \rho e^{i\varphi} = \rho_{\varepsilon, g} e^{i\varphi_{\varepsilon, g}} = \rho e^{i(\theta + \psi)} = \rho_{\varepsilon, g} e^{i(\theta_g + \psi_{\varepsilon, g})},$$

and similarly

$$u_{*, g} = e^{i(\theta + \psi_*)} = e^{i(\theta_g + \psi_{*, g})}.$$

We may choose $\psi_{*, g}$ in order to have

$$\|\psi_{*, g}\|_{C^{1, \alpha}(\omega)} \leq C, \quad \forall g \in \mathfrak{A}, \quad (7.3)$$

and we normalize $\psi_{\varepsilon, g}$ by the condition

$$\psi_{\varepsilon, g} = \psi_{*, g} \quad \text{on } \partial\Omega. \quad (7.4)$$

In terms of ρ , φ and ψ , the Ginzburg-Landau equation reads

$$\begin{cases} \operatorname{div}(\rho^2 \nabla \varphi) = \operatorname{div}(\rho^2(\theta + \psi)) = 0 \\ -\Delta \rho = \frac{1}{\varepsilon^2} \rho(1 - \rho^2) - \rho |\nabla \varphi|^2 \end{cases} .$$

Step 1. We have

$$\|\nabla \varphi_{\varepsilon, g}\|_{L^p(\omega)} \leq C_p, \quad \forall \varepsilon < \varepsilon_0, \forall g \in \mathfrak{A}, \forall 1 < p < \infty.$$

Indeed, we start by noting that we have $\|\nabla \theta_g\|_{L^p(\omega)} \leq C_p$; therefore, it suffices to prove that $\|\nabla \psi_{\varepsilon, g}\|_{L^p(\omega)} \leq C_p$. Using the equation $\operatorname{div}(\rho^2 \nabla \varphi) = 0$, we see that $\psi_{\varepsilon, g}$ satisfies

$$\Delta \psi_{\varepsilon, g} = \operatorname{div}((1 - \rho_{\varepsilon, g}^2) \nabla \theta_g + (1 - \rho_{\varepsilon, g}^2) \nabla \psi_{\varepsilon, g}) \quad \text{in } \omega. \quad (7.5)$$

We obtain

$$\begin{aligned} \|\nabla \psi_{\varepsilon, g}\|_{L^p(\omega)} &\leq C (\|\psi_{\varepsilon, g}\|_{W^{1-1/p, p}(\partial\omega)} + \|(1 - \rho_{\varepsilon, g}^2) \nabla \theta_g\|_{L^p(\omega)} + \|(1 - \rho_{\varepsilon, g}^2) \nabla \psi_{\varepsilon, g}\|_{L^p(\omega)}) \\ &\leq C_p + C \|1 - \rho_{\varepsilon, g}^2\|_{L^\infty(\omega)} \|\nabla \psi_{\varepsilon, g}\|_{L^p(\omega)}. \end{aligned}$$

Since $\rho_{\varepsilon, g} \rightarrow 1$ as $\varepsilon \rightarrow 0$ uniformly in $\bar{\omega}$ and in $g \in \mathfrak{A}$, the second term in the right-hand side of the above inequality can be absorbed in the left-hand side and we obtain the announced result.

Step 2. For $1 < p < \infty$ we have $\nabla \rho_{\varepsilon, g} \rightarrow 0$ in $L^p(\omega)$ as $\varepsilon \rightarrow 0$, uniformly in $g \in \mathfrak{A}$. This is obtained as follows. Let $\eta := \eta_{\varepsilon, g} := 1 - \rho_{\varepsilon, g} \in [0, 1]$, which satisfies

$$\begin{cases} -\Delta \eta + \frac{1}{\varepsilon^2} \rho(1 + \rho)\eta = \rho |\nabla \varphi|^2 & \text{in } \omega \\ \eta = 0 & \text{on } \partial\Omega \\ \eta \rightarrow 0 \text{ in } C^2 & \text{on } \partial\omega \setminus \partial\Omega \end{cases}, \quad (7.6)$$

the latter convergence being uniform in $g \in \mathfrak{A}$.

We may assume that $p \geq 2$. Multiplying (7.6) by η^{p-1} and using Step 1 and Hölder's inequality we find that, for small ε , we have

$$\begin{aligned} \frac{1}{4\varepsilon^2} \int_{\omega} \eta^p &\leq \frac{1}{\varepsilon^2} \int_{\omega} \rho(1 + \rho)\eta^p \\ &= \int_{\omega} \rho |\nabla \varphi|^2 \eta^{p-1} + \int_{\partial\omega \setminus \partial\Omega} \eta^{p-1} \frac{\partial \eta}{\partial \nu} - (p-1) \int_{\omega} \eta^{p-2} |\nabla \eta|^2 \\ &\leq \int_{\omega} \rho |\nabla \varphi|^2 \eta^{p-1} + \int_{\partial\omega \setminus \partial\Omega} \eta^{p-1} \frac{\partial \eta}{\partial \nu} \leq C \left(\int_{\omega} \eta^p \right)^{(p-1)/p} + C, \end{aligned}$$

and thus

$$\|\eta_{\varepsilon, g}\|_{L^p(\omega)} \leq C_p, \quad \forall \varepsilon < \varepsilon_0, \forall g \in \mathfrak{A}, \forall p < \infty. \quad (7.7)$$

Inserting (7.7) into (7.6), we find that $\Delta\eta$ is bounded in $L^p(\omega)$, $\forall p < \infty$. By standard elliptic estimates, we find that η (and thus ρ) is bounded in $W^{2,p}(\omega)$, $\forall p < \infty$. We conclude via the compact embedding $W^{2,p} \hookrightarrow W^{1,p}$ and the fact that, by Lemma 6.4, we have $\rho \rightarrow 1$ uniformly in $\bar{\omega}$.

Step 3. For every $\beta < 1$, we have $\psi_{\varepsilon,g} \rightarrow \psi_{*,g}$ in $C^{1,\beta}(\bar{\omega})$ as $\varepsilon \rightarrow 0$, uniformly in $g \in \mathfrak{A}$.

Indeed, $\psi_{\varepsilon,g} - \psi_{*,g}$ satisfies

$$\begin{cases} \Delta(\psi_{\varepsilon,g} - \psi_{*,g}) = -\frac{2}{\rho_{\varepsilon,g}} \nabla \rho_{\varepsilon,g} \cdot \nabla(\theta_g + \psi_{\varepsilon,g}) & \text{in } \omega \\ \psi_{\varepsilon,g} - \psi_{*,g} = 0 & \text{on } \partial\Omega \\ \psi_{\varepsilon,g} - \psi_{*,g} \rightarrow 0 \text{ in } C^2 & \text{on } \partial\omega \setminus \partial\Omega \end{cases}, \quad (7.8)$$

the latter convergence being uniform in g . By Steps 1 and 2, we have

$$\|\Delta(\psi_{\varepsilon,g} - \psi_{*,g})\|_{L^p(\omega)} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0, \text{ uniformly in } g.$$

Using (7.8), we find that $\psi_{\varepsilon,g} - \psi_{*,g} \rightarrow 0$ in $W^{2,p}(\omega)$. We conclude via the embedding $W^{2,p}(\omega) \hookrightarrow C^{1,\beta}(\bar{\omega})$, valid when $p > 2$ and $\beta = 1 - 2/p$.

Step 4. Conclusion.

We have

$$T_\varepsilon(g) = u_{\varepsilon,g} \wedge \frac{\partial u_{\varepsilon,g}}{\partial \nu} = \frac{\partial \varphi_{\varepsilon,g}}{\partial \nu} = \frac{\partial \theta_g}{\partial \nu} + \frac{\partial \psi_{\varepsilon,g}}{\partial \nu},$$

and similarly $T_*(g) = \frac{\partial \theta_g}{\partial \nu} + \frac{\partial \psi_{*,g}}{\partial \nu}$. Using Step 3, we find that

$$T_\varepsilon(g) - T_*(g) = \frac{\partial(\psi_{\varepsilon,g} - \psi_{*,g})}{\partial \nu} \rightarrow 0 \text{ in } C^\beta(\partial\Omega) \quad \text{as } \varepsilon \rightarrow 0, \text{ uniformly in } g \in \mathfrak{A}. \square$$

8 Existence of critical points in nondegenerate domains

Before stating the main result of this section, let us recall the definition (5.15) of the operator U_* in Section 5. Given a^0 a nondegenerate critical point of $W(\cdot, g)$, we first define, in a $C^{1,\alpha}$ neighborhood of g , the operator $T_* = T_{*,a^0,g}$. Then U_* is defined in a neighborhood \mathfrak{B} of the origin in $C^{1,\alpha}(\partial\Omega; \mathbb{R})$ by

$$U_*(\psi) = U_{*,a^0,g}(\psi) = T_*(ge^{i\psi}) = T_{*,a^0,g}(ge^{i\psi}).$$

We still denote by U_* the induced map $U_* : \mathfrak{B}/\mathbb{R} \rightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R})$, and recall that U_* is smooth.

8.1 Theorem. Let $d_1, \dots, d_k \in \{-1, 1\}$ and set $d := d_1 + \dots + d_k$.

Let Ω be a bounded simply connected $C^{1,\alpha}$ domain satisfying the two following nondegeneracy conditions:

- (ND1) There exists $a^0 \in \Omega_*^k$ such that a^0 is a nondegenerate critical point of $W(\cdot, g^0) = W^\Omega(\cdot, d_1, \dots, d_k, g^0)$, with $g^0 = g^{a^0}$ the canonical boundary data associated with a^0 and d_1, \dots, d_k .
- (ND2) The corresponding operator $U_{*, a^0, g^0} : \mathfrak{B}/\mathbb{R} \rightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R})$ is a local diffeomorphism at the origin, i.e., the differential

$$DU_*(0) : C^{1,\alpha}(\partial\Omega; \mathbb{R})/\mathbb{R} \longrightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R})$$

is invertible.

Then there exists $\varepsilon_0 > 0$ such that, for $\varepsilon \in (0, \varepsilon_0)$, there exists $u_\varepsilon \in \mathcal{E}_d$ a critical point of E_ε with prescribed degree d .

8.2 Remark. It will be clear from the proof of Theorem 8.1 that the nondegeneracy condition (ND2) can actually be replaced by the following weaker condition:

(ND2') U_* is a local homeomorphism near the origin.

However in what follows it will be more convenient for us to consider the condition (ND2). The main reason for this is that (ND2) is stable under small perturbation of the domain, while it is not clear that (ND2') is stable.

8.3 Remark. We connect here the hypothesis (ND2) in Theorem 8.1 to the hypothesis (ND2) presented in the introduction. It is easy to see that $DU_*(0)$ is a Fredholm operator of index zero.³ Thus the above hypothesis (ND2) is equivalent to the fact that $DU_*(0)$ is onto. Is it not difficult to see (but will not be needed in what follows) that the surjectivity of $DU_*(0)$ is equivalent to the hypothesis (ND2) in the introduction, and that the index of the operator L that appears in the introduction is

$$\text{ind} L = \text{ind} DU_*(0) + 1 = 1.$$

Proof of Theorem 8.1. Since Ω satisfies (ND1), the results of Section 6 and 7 apply. We consider, as in Section 7, the operators

$$T_\varepsilon, T_* : \mathfrak{A} \rightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R})$$

and

$$U_* : \mathfrak{B}/\mathbb{R} \rightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R}),$$

³When Ω is the unit disc, this follows from the explicit formula (10.13). The general case is obtained by transport onto the unit disc, as in Section 3.

where $\mathfrak{A} = \{g \in C^{1,\alpha}(\Omega, \mathbb{S}^1); \|g - g^0\| < \delta\}$ and $\mathfrak{B} = \{\psi \in C^{1,\alpha}(\partial\Omega; \mathbb{R}); \|\psi\| < \delta\}$. Here, δ and ε are sufficiently small. We note that T_ε takes its values in $\dot{C}^\alpha(\partial\Omega; \mathbb{R})$. Indeed, $u = u_{\varepsilon, g}$ satisfies

$$\int_{\partial\Omega} u \wedge \frac{\partial u}{\partial \nu} = \int_{\Omega} \operatorname{div}(u \wedge \nabla u) = \int_{\Omega} u \wedge \Delta u = \int_{\Omega} \frac{|u|^2 - 1}{\varepsilon^2} u \wedge u = 0.$$

By Lemma 6.5, we may also consider the induced operators

$$U_\varepsilon : \mathfrak{B}/\mathbb{R} \rightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R}), U_\varepsilon(\psi) = T_\varepsilon(g^0 e^{i\psi}).$$

By definition of the canonical boundary datum, it holds

$$U_*(0) = u_{*, \alpha^0, g^0} \wedge \frac{\partial u_{*, \alpha^0, g^0}}{\partial \nu} = \frac{\partial \Phi_{\alpha^0, g^0}}{\partial \tau} = \frac{\partial \widehat{\Phi}_{\alpha^0}}{\partial \tau} = 0. \quad (8.1)$$

Thanks to (ND2), by considering a smaller δ if necessary, we may assume that U_* is a homeomorphism onto its image. By (8.1), there exists some $\eta > 0$ such that

$$U_*(\mathfrak{B}/\mathbb{R}) \supset \{\psi \in \dot{C}^\alpha(\partial\Omega; \mathbb{R}); \|\psi\|_{C^\alpha(\partial\Omega)} < \eta\} := B_\eta. \quad (8.2)$$

Recall the result of Proposition 7.1: for sufficiently small ε , $U_\varepsilon - U_*$ is compact and we have

$$\limsup_{\varepsilon \rightarrow 0} \{\|U_\varepsilon(\psi) - U_*(\psi)\|_{C^\alpha(\partial\Omega)}; \psi \in \mathfrak{B}\} = 0. \quad (8.3)$$

Using (8.2), (8.3) and standard properties of the Leray-Schauder degree, we find that

$$U_\varepsilon(\mathfrak{B}/\mathbb{R}) \supset \left\{ \psi \in \dot{C}^\alpha(\partial\Omega; \mathbb{R}); \|\psi\|_{C^\alpha(\partial\Omega)} < \frac{\eta}{2} \right\} = B_{\eta/2}, \quad (8.4)$$

for sufficiently small ε . Indeed, the argument goes as follows. We start from

$$U_\varepsilon(\mathfrak{B}/\mathbb{R}) = (\operatorname{Id} + (U_\varepsilon - U_*) \circ U_*^{-1})(U_*(\mathfrak{B}/\mathbb{R})) \supset (\operatorname{Id} + (U_\varepsilon - U_*) \circ U_*^{-1})(B_\eta). \quad (8.5)$$

Here, Id denotes the identity map in $\dot{C}^\alpha(\partial\Omega; \mathbb{R})$.

Let $L_\varepsilon := (U_\varepsilon - U_*) \circ U_*^{-1}$. Then $L_\varepsilon : B_\eta \rightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R})$ is compact and, by (8.3), there exists $\varepsilon_0 > 0$ such that

$$\sup \{\|L_\varepsilon(\psi)\|_{C^\alpha(\partial\Omega)}; \psi \in B_\eta\} < \eta/2, \quad \forall \varepsilon \in (0, \varepsilon_0). \quad (8.6)$$

We complete the proof of (8.4) by showing that $B_{\eta/2} \subset (\operatorname{Id} + L_\varepsilon)(B_\eta)$ for $\varepsilon \in (0, \varepsilon_0)$. For this purpose, we let $\psi_0 \in B_{\eta/2}$ and consider the compact operator $T : B_\eta \rightarrow \dot{C}^\alpha(\partial\Omega; \mathbb{R})$, $T(\psi) := L_\varepsilon(\psi) - \psi_0$. We claim that

$$(\operatorname{Id} + sT)(\psi) \neq 0, \quad \forall s \in [0, 1], \forall \psi \in \partial B_\eta. \quad (8.7)$$

Indeed, (8.7) is obtained by contradiction. Otherwise, using (8.6), we obtain, for some ψ such that $\|\psi\|_{C^\alpha(\partial\Omega)} = \eta$:

$$\eta/2 < \eta - s\|L_\varepsilon(\psi)\|_{C^\alpha(\partial\Omega)} \leq \|\psi + sL_\varepsilon(\psi)\|_{C^\alpha(\partial\Omega)} = \|s\psi_0\|_{C^\alpha(\partial\Omega)} < \eta/2.$$

By (8.7), the Leray-Schauder degree $\deg(\text{Id} + sT, B_\eta, 0)$ is well defined. By homotopy invariance, we find that

$$\deg(\text{Id} + sT, B_\eta, 0) = \deg(\text{Id}, B_\eta, 0) = 1.$$

As a consequence, the equation $(\text{Id} + T)(\psi) = 0$ admits at least a solution $\psi \in B_\eta$. This ψ satisfies $(\text{Id} + L_\varepsilon)(\psi) = \psi_0$. The proof of (8.4) is complete.

Let $\varepsilon \in (0, \varepsilon_0)$. Then, by (8.4), there exists some $\psi \in \mathfrak{B}$ such that $U_\varepsilon(\psi) = 0$. Let $g = g^0 e^{i\psi}$. Then $u_\varepsilon = u_{\varepsilon, g} \in \mathcal{E}_d$ is a solution of the Ginzburg-Landau equation, and it satisfies the semi-stiff boundary condition

$$u_\varepsilon \wedge \frac{\partial u_\varepsilon}{\partial \nu} = T_\varepsilon(g) = U_\varepsilon(\psi) = 0 \quad \text{on } \partial\Omega.$$

Therefore, u_ε is a critical point of E_ε with prescribed degree d . \square

9 Nondegeneracy of domains is stable

In this section we show that, if a domain Ω_0 satisfies the nondegeneracy conditions (ND1)-(ND2) required in Theorem 8.1, then a slightly perturbed domain $\Omega \approx \Omega_0$ still satisfies these nondegeneracy conditions.

9.1 Theorem. *Assume that Ω_0 satisfies (ND1)-(ND2). Fix a conformal representation $f_0 : \mathbb{D} \rightarrow \Omega_0$. There exists $\delta > 0$ such that, for every holomorphic map $f \in C^{1,\alpha}(\overline{\mathbb{D}})$ satisfying $\|f - f_0\|_{C^1} < \delta$, the domain $\Omega := f(\mathbb{D})$ satisfies (ND1)-(ND2).*

Proof. Let V_α be as in (5.4). We let $\tilde{g}^0 \in C^{1,\alpha}(\mathbb{S}^1; \mathbb{S}^1)/\mathbb{S}^1$ denote the canonical boundary datum associated with $\alpha^0 := f_0^{-1}(\alpha^0)$, so that $\tilde{g}^0 = g^0 \circ f_0$.

Since Ω_0 satisfies (ND1), we know from Proposition 5.3 that there exist: a neighborhood \mathcal{V}_1 of f_0 in V_α , a neighborhood \mathcal{V}_2 of the origin in $C^{1,\alpha}(\mathbb{S}^1; \mathbb{R})$, and a smooth map $\alpha : \mathcal{V}_1 \times \mathcal{V}_2 \rightarrow \mathbb{D}_*^k$, such that the following holds. For $f \in \mathcal{V}_1$ and $\psi \in \mathcal{V}_2$, let $\Omega = f(\mathbb{D})$ and $g = (\tilde{g}^0 e^{i\psi}) \circ f^{-1}$. Then $\alpha(f, \psi) := f(\alpha(f, \psi))$ is a nondegenerate critical point of $W^\Omega(\cdot, g)$.

By the above, we may define, as in (5.13), the smooth operator $U_{*,f} = U_{*,\alpha(f,0),\tilde{g}^0 \circ f^{-1}}$,

$$U_{*,f}(\psi) = N^\Omega(\alpha(f, \psi \circ f), (\tilde{g}^0 \circ f^{-1})e^{i\psi}) \quad \text{for small } \psi \in C^{1,\alpha}(\partial\Omega; \mathbb{R})/\mathbb{R}. \quad (9.1)$$

The spaces between which $U_{*,f}$ is defined vary with f . In order to deal with fixed spaces, we consider the linear isomorphisms

$$\Theta_f : \dot{C}^\alpha(\partial\Omega; \mathbb{R}) \rightarrow \dot{C}^\alpha(\mathbb{S}^1; \mathbb{R}), \psi \mapsto |f'| \psi \circ f, \quad (9.2)$$

$$\Xi_f : C^{1,\alpha}(\partial\Omega; \mathbb{R})/\mathbb{R} \rightarrow C^{1,\alpha}(\mathbb{S}^1; \mathbb{R})/\mathbb{R}, \psi \mapsto \psi \circ f, \quad (9.3)$$

and we let

$$U(f, \psi) = \Theta_f \circ U_{*,f} \circ \Xi_f^{-1}(\psi) \quad \text{for } (f, \psi) \in \mathcal{V}_1 \times (\mathcal{V}_2/\mathbb{R}), \quad (9.4)$$

so that $U_{*,f}$ is a local diffeomorphism if and only if $U(f, \cdot)$ is a local diffeomorphism.

Moreover, if we express N^Ω using (3.7), then we obtain

$$U(f, \psi) = N^\mathbb{D}(\alpha(f, \psi), \tilde{g}_0 e^{i\psi}). \quad (9.5)$$

By combining (9.5) with the explicit formula (4.5) for $N^\mathbb{D}$, we find that $U : \mathcal{V}_1 \times (\mathcal{V}_2/\mathbb{R}) \rightarrow \dot{C}^\alpha(\mathbb{S}^1; \mathbb{R})$ is smooth.

On the other hand, using the definition (2.16) of the canonical boundary datum, we have

$$u_{*,a,g^a} \wedge \frac{\partial u_{*,a,g^a}}{\partial \nu} = \frac{\partial \Phi_{a,g^a}}{\partial \tau} = \frac{\partial \widehat{\Phi}_a}{\partial \tau} = 0,$$

so that $U(f_0, 0) = 0$.

Moreover, since Ω_0 satisfies (ND2), $U(f_0, \cdot)$ is a local diffeomorphism near the origin, i.e., $D_\psi U(f_0, 0)$ is invertible. By the implicit function theorem, possibly after shrinking \mathcal{V}_1 , for every $f \in \mathcal{V}_1$ there exists $\psi(f) \in \mathcal{V}_2$ such that

$$U(f, \psi(f)) = 0. \quad (9.6)$$

In addition, the map $f \mapsto \psi(f)$ is smooth and we can assume that $D_\psi U(f, \psi(f))$ is invertible.

Let $f \in \mathcal{V}_1$ and set $\Omega := f(\mathbb{D})$. We claim that Ω satisfies (ND1)-(ND2). Assuming the claim proved for the moment, we complete the proof of Theorem 9.1 by taking any $\delta > 0$ such that

$$\{f \in X_\alpha; \|f - f_0\|_{C^{1,\alpha}} < \delta\} \subset \mathcal{V}_1.$$

We next turn to the proof of the claim. Let $g_\Omega := (\tilde{g}_0 e^{i\psi(f)}) \circ f^{-1} \in C^{1,\alpha}(\partial\Omega; \mathbb{S}^1)$, and $a_\Omega := \alpha(f, \psi(f)) \in \Omega_*^k$. By the definition (9.6) of $\psi(f)$ and the definition (9.4) of U , we obtain

$$U_{*,f}(\psi(f) \circ f^{-1}) = 0. \quad (9.7)$$

By combining (9.7) with the definition (9.1) of $U_{*,f}$, we find that

$$N^\Omega(a_\Omega, g_\Omega) = u_{*,a_\Omega, g_\Omega} \wedge \frac{\partial u_{*,a_\Omega, g_\Omega}}{\partial \nu} = \frac{\partial \Phi_{a_\Omega, g_\Omega}}{\partial \tau} = 0. \quad (9.8)$$

The normalization condition in (2.5) combined with (9.8) implies that $\Phi_{a_\Omega, g_\Omega} = 0$ on $\partial\Omega$, and thus

$$\Phi_{a_\Omega, g_\Omega} = \widehat{\Phi}_{a_\Omega}. \quad (9.9)$$

In turn, (9.9) implies that $g_\Omega = g^{a_\Omega}$ is the canonical boundary data associated with a_Ω . Since, by definition of the map $(f, \psi) \mapsto a(f, \psi)$, the configuration a_Ω is a nondegenerate critical point of $W(\cdot, g_\Omega)$, we find that the nondegeneracy condition (ND1) is satisfied by Ω .

On the other hand, since $D_\psi U(f, \psi(f))$ is invertible, $U_{*,f}$ is a local diffeomorphism near $\psi(f) \circ f^{-1}$, which means that U_{*,a_Ω, g_Ω} is a local diffeomorphism near the origin. We find that Ω satisfies (ND2).

The proof of Theorem 9.1 is complete. \square

10 The radial configuration is nondegenerate

In this section we let $d = 1$, $k = 1$, $d_1 = 1$, and prove that the unit disc \mathbb{D} satisfies (ND1)-(ND2). As a consequence, domains close to the unit disc satisfy the nondegeneracy conditions when $d = 1$, $k = 1$, $d_1 = 1$.

10.1 Proposition. *Assume $\Omega = \mathbb{D}$, $k = 1$, $d = 1$. Then $a = 0$ is a nondegenerate critical point of $W(\cdot, g^0)$, and $DU_{*,0, g^0}(0)$ is invertible.*

Proof. Step 1. 0 is a nondegenerate critical point of $W(\cdot, g^0)$.

Indeed, by combining (2.25) and (4.1), we easily obtain that the canonical boundary datum $g^0 : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ corresponding to $a = 0$ is given by $g^0(z) = z$. From (4.10) we know that

$$W(a, g^0) = \widehat{W}(a) + \frac{1}{2} \int_{\mathbb{D}} \left| \nabla \psi_{a, g^0}^* \right|^2. \quad (10.1)$$

On the other hand, (4.16) leads to

$$\nabla \psi_{a, g^0}^*(x) = -2 \frac{a(\bar{a}x - 1)}{|1 - \bar{a}x|^2}, \quad \forall x \in \mathbb{D}, \quad (10.2)$$

and therefore

$$\frac{1}{2} \int_{\mathbb{D}} \left| \nabla \psi_{a, g^0}^* \right|^2 = 2|a|^2 \int_{\mathbb{D}} \frac{dx}{|1 - \bar{a}x|^2}. \quad (10.3)$$

Thanks to the $|a|^2$ factor, if we differentiate (10.3) with respect to a , and next let $a = 0$, we obtain

$$\nabla_a \left[\frac{1}{2} \int_{\mathbb{D}} \left| \nabla \psi_{a, g^0}^* \right|^2 \right] \Big|_{a=0} = 0. \quad (10.4)$$

If we differentiate twice (10.3) with respect to a , and next let $a = 0$, then we are left with only one non zero term (thanks to the $|a|^2$ factor again). More specifically, we obtain

$$\nabla_a^2 \left[\frac{1}{2} \int_{\mathbb{D}} |\nabla \psi_{a, g^0}^*|^2 \right] \Big|_{a=0} = 4 \int_{\mathbb{D}} \frac{dx}{|1 - \bar{a}x|^2} \Big|_{a=0} I_2 = 4 \int_{\mathbb{D}} dx I_2 = 4\pi I_2. \quad (10.5)$$

By combining (10.1) with (10.4) and (10.5), we find that

$$\nabla_a W(0, g^0) = \nabla \widehat{W}(0), \quad \nabla_a^2 W(0, g^0) = \nabla^2 \widehat{W}(0) + 4\pi I_2. \quad (10.6)$$

We next compute $\nabla \widehat{W}(0)$ and $\nabla^2 \widehat{W}(0)$. When $k = 1$ and $d = 1$, formula (4.2) reads

$$\widehat{W}(a) = \pi \log(1 - |a|^2), \quad \forall a \in \mathbb{D}. \quad (10.7)$$

Identifying the complex number a with a vector in \mathbb{R}^2 , the two first derivatives of \widehat{W} are respectively given by:

$$\nabla \widehat{W}(a) = \frac{2\pi}{|a|^2 - 1} a \in \mathbb{R}^2 \quad (10.8)$$

$$\nabla^2 \widehat{W}(a) = \frac{2\pi}{|a|^2 - 1} I_2 + \frac{4\pi}{(|a|^2 - 1)^2} a \otimes a \in M_2(\mathbb{R}). \quad (10.9)$$

In particular, we obtain $\nabla \widehat{W}(0) = 0$ and $\nabla^2 \widehat{W}(0) = -2\pi I_2$. Plugging this into (10.6) yields

$$\nabla_a W(0, g^0) = 0, \quad \nabla_a^2 W(0, g^0) = 2\pi I_2, \quad (10.10)$$

so that $a = 0$ is indeed a nondegenerate critical point of $W(\cdot, g^0)$.

Step 2. $DU_*(0)$ is invertible.

In our case, formula (4.5) becomes

$$N(a, g^0 e^{i\psi}) = \frac{\partial \psi^*}{\partial \tau} - 2 \frac{a \wedge z}{|z - a|^2}. \quad (10.11)$$

Therefore

$$U_*(\psi) = \frac{\partial \psi^*}{\partial \tau} - 2 \frac{a(\psi) \wedge z}{|z - a(\psi)|^2}, \quad (10.12)$$

where $\psi \mapsto a(\psi)$ is smooth, $a(0) = 0$ and $a(\psi)$ is a nondegenerate critical point of $W(\cdot, g^0 e^{i\psi})$.

Using (10.12) together with the fact that $a(0) = 0$, we obtain that

$$DU_*(0)\psi = \frac{\partial \psi^*}{\partial \tau} - 2(Da(0)\psi) \wedge z. \quad (10.13)$$

In (10.13), ψ is either a function in $C^{1,\alpha}(\mathbb{S}^1; \mathbb{R})$, or a class in $C^{1,\alpha}(\mathbb{S}^1; \mathbb{R})/\mathbb{R}$. Thus the linear operator $DU_*(0) : C^{1,\alpha}(\mathbb{S}^1; \mathbb{R})/\mathbb{R} \rightarrow \dot{C}^\alpha(\mathbb{S}^1; \mathbb{R})$ can be written $DU_*(0) = L - K$, where

$$L(\psi) := \frac{\partial \psi^*}{\partial \tau} \text{ and } K(\psi) := 2(Da(0)\psi) \wedge z, \quad \forall \psi \in C^{1,\alpha}(\mathbb{S}^1; \mathbb{R})/\mathbb{R}.$$

The operator L is an isomorphism, and K is compact since it has finite range. As a consequence, $DU_*(0)$ is Fredholm of index zero and, in order to complete Step 2, it suffices to prove that $DU_*(0)$ is injective. For this purpose, we compute $Da(0)$ using the implicit equation

$$F(a(\psi), \psi) := \nabla_a W(a(\psi), g^0 e^{i\psi}) = 0 \quad (10.14)$$

satisfied by a . By differentiating (10.14) with respect to ψ we obtain (via (10.10))

$$Da(0)\psi = -\nabla_a^2 W(0, g^0) D_\psi F(0, 0)\psi = -2\pi I_2 D_\psi F(0, 0)\psi. \quad (10.15)$$

Let us compute $D_\psi F(0, 0)$. Recalling (4.3), we find that

$$\begin{aligned} F(a, \psi) &= \nabla \widehat{W}(a) + \nabla_a \left[\frac{1}{2} \int_{\mathbb{D}} |\nabla \psi_{a, g^0}^* + \nabla \psi^*|^2 \right] \\ &= \nabla \widehat{W}(a) + \nabla_a \left[\frac{1}{2} \int_{\mathbb{D}} |\nabla \psi_{a, g^0}^*|^2 \right] + \nabla_a \left[\int_{\mathbb{D}} \nabla \psi_{a, g^0}^* \cdot \nabla \psi^* \right]. \end{aligned} \quad (10.16)$$

The two first terms do not depend on ψ , and the last term depends linearly on ψ . Hence we obtain

$$D_\psi F(a, 0)\psi = \nabla_a \left[\int_{\mathbb{D}} \nabla \psi_{a, g^0}^* \cdot \nabla \psi^* \right]. \quad (10.17)$$

Integrating by parts, using the explicit formula (4.16) for ψ_{a, g^0}^* , and the fact that

$$\int_{\mathbb{S}^1} \frac{\partial \psi^*}{\partial v} = \int_{\mathbb{S}^1} \frac{\partial \psi}{\partial \tau} = 0,$$

we find that

$$\int_{\mathbb{D}} \nabla \psi_{a, g^0}^* \cdot \nabla \psi^* = -2 \int_{\mathbb{S}^1} \log|1 - \bar{a}z| \frac{\partial \psi^*}{\partial v}. \quad (10.18)$$

If we first plug (10.18) into (10.17) and next let $a = 0$, then we obtain

$$D_\psi F(0, 0)\psi = 2 \int_{\mathbb{S}^1} z \frac{\partial \psi^*}{\partial v}, \quad (10.19)$$

and finally

$$Da(0)\psi = -4\pi \int_{\mathbb{S}^1} z \frac{\partial \psi^*}{\partial v}. \quad (10.20)$$

We are now in position to prove that $DU_*(0)$ is injective. Let $\psi \in \ker DU_*(0)$. Then we have

$$\frac{\partial \psi^*}{\partial \tau} = 2(D\alpha(0)\psi) \wedge z = \alpha \wedge z, \quad (10.21)$$

where

$$\alpha = -8\pi \int_{\mathbb{S}^1} z \frac{\partial \psi^*}{\partial v} \in \mathbb{C}. \quad (10.22)$$

Since ψ^* is harmonic and has zero average on \mathbb{S}^1 , we may write

$$\psi^*(re^{i\theta}) = \sum_{n \neq 0} a_n r^n e^{in\theta}. \quad (10.23)$$

Hence (10.21) yields

$$\frac{\bar{\alpha}}{2i} e^{i\theta} - \frac{\alpha}{2i} e^{-i\theta} = \frac{\partial \psi^*}{\partial \tau}(e^{i\theta}) = \sum_{n \neq 0} in a_n e^{in\theta}. \quad (10.24)$$

Identifying the Fourier coefficients, we obtain

$$a_n = 0 \text{ for } |n| > 1, \quad a_1 = -\frac{\bar{\alpha}}{2}, \quad a_{-1} = -\frac{\alpha}{2}, \quad (10.25)$$

so that (10.23) becomes

$$\psi^*(re^{i\theta}) = -\frac{\bar{\alpha}}{2} r e^{i\theta} - \frac{\alpha}{2} r e^{-i\theta}. \quad (10.26)$$

By (10.26), we have

$$\int_{\mathbb{S}^1} z \frac{\partial \psi^*}{\partial v} = -\frac{1}{2} \int_0^{2\pi} e^{i\theta} (\bar{\alpha} e^{i\theta} + \alpha e^{-i\theta}) d\theta = -\pi \alpha. \quad (10.27)$$

Plugging (10.27) into (10.22) we obtain $\alpha = 8\pi^2 \alpha$, so that $\alpha = 0$ and consequently $\psi^* = 0$. Therefore, we have $\psi = 0$ modulo \mathbb{R} , and thus $DU_*(0)$ is invertible. \square

10.2 Corollary. *If a domain Ω is sufficiently close to the unit disc, in the sense that there exists a conformal representation $f : \mathbb{D} \rightarrow \Omega$ such that $\|f - \text{Id}\|_{C^{1,\alpha}} < \delta$ for sufficiently small δ , then, for small ε , E_ε admits critical points with prescribed degree one.*

Appendix

The following is a $C^{1,\alpha}$ variant of [7, Lemmas A1, A2].

10.3 Lemma. *Let $G \subset \mathbb{R}^n$ be a bounded open set of class $C^{1,\alpha}$. Assume that*

$$\begin{cases} \Delta w = f & \text{in } G \\ w = \varphi & \text{on } \partial G \end{cases}. \quad (10.28)$$

Then

$$\sup_G |\nabla w| \leq C \left(\|f\|_{L^\infty}^{1/2} \left(\|w\|_{L^\infty}^{1/2} + \|\varphi\|_{L^\infty(\partial G)}^{1/2} \right) + \|\varphi\|_{C^{1,\alpha}(\partial G)} \right), \quad (10.29)$$

$$|\nabla w|_{0,\alpha,G} \leq C \left(\|f\|_{L^\infty}^{1/2+\alpha/2} \left(\|w\|_{L^\infty}^{1/2-\alpha/2} + \|\varphi\|_{L^\infty(\partial G)}^{1/2-\alpha/2} \right) + \|\varphi\|_{C^{1,\alpha}(\partial G)} \right), \quad (10.30)$$

for a constant C depending only on G . In addition, when $G = \Omega_\sigma$, where $\sigma_1 \leq \sigma \leq \sigma_2$ and σ_1, σ_2 are fixed small numbers, we may take C independent of σ .

Proof. We write $w = u + v$, where

$$\begin{cases} \Delta u = 0 & \text{in } G \\ u = \varphi & \text{on } \partial G \end{cases}, \quad (10.31)$$

and

$$\begin{cases} \Delta v = f & \text{in } G \\ v = 0 & \text{on } \partial G \end{cases}. \quad (10.32)$$

By standard elliptic estimates [13, Theorem 8.33] we have

$$\|u\|_{C^{1,\alpha}} \leq c \|\varphi\|_{C^{1,\alpha}}. \quad (10.33)$$

Therefore we only need to prove that v satisfies the estimates

$$\sup_G |\nabla v| \leq C \|f\|_{L^\infty}^{1/2} \|v\|_{L^\infty}^{1/2}, \quad (10.34)$$

$$|\nabla v|_{0,\alpha,G} \leq C \|f\|_{L^\infty}^{1/2+\alpha/2} \|v\|_{L^\infty}^{1/2-\alpha/2}. \quad (10.35)$$

Estimate (10.34) is proved in [7, Lemma A.2] by combining an interior estimate with a boundary estimate. Estimate (10.35) can be obtained following exactly the same lines. In order to see this, we detail for example the proof of the interior estimate corresponding to (10.35). Proceeding as in [7, Lemma A.1], we first show that

$$|\nabla v|_{0,\alpha,G_d} \leq C \left(\|f\|_{L^\infty}^{1/2+\alpha/2} \|v\|_{L^\infty}^{1/2-\alpha/2} + \frac{1}{d^{1+\alpha}} \|v\|_{L^\infty} \right), \quad (10.36)$$

where, for $d > 0$, we let $G_d := \{x \in G; \text{dist}(x, \partial G) > d\}$. In order to prove (10.36), we let $x_0 \in G_d$ and $\lambda \in (0, d]$, and define

$$v_\lambda(y) := v(x_0 + \lambda y), \quad y \in B_1(0). \quad (10.37)$$

Then the function v_λ satisfies the equation

$$\Delta v_\lambda = f_\lambda \text{ in } B_1(0), \quad \text{with } f_\lambda(y) := \lambda^2 f(x_0 + \lambda y). \quad (10.38)$$

Standard elliptic estimates [13, Theorem 8.33] yield

$$\begin{aligned} \lambda^{1+\alpha} |\nabla v|_{0,\alpha,B_{\lambda/2}(x_0)} &= |\nabla v_\lambda|_{0,\alpha,B_{1/2}(0)} \leq C(\|v_\lambda\|_{L^\infty} + \|f_\lambda\|_{L^\infty}) \\ &\leq C(\|v\|_{L^\infty} + \lambda^2 \|f\|_{L^\infty}). \end{aligned} \quad (10.39)$$

We next discuss the two following cases.

Case 1. $\frac{\|v\|_{L^\infty}}{\|f\|_{L^\infty}} \leq d^2$.

In this case, we apply (10.39) with $\lambda = (\|v\|_{L^\infty}/\|f\|_{L^\infty})^{1/2}$. We find that

$$|\nabla v|_{0,\alpha,B_{\lambda/2}(x_0)} \leq 2C\|v\|_{L^\infty}^{1/2-\alpha/2} \|f\|_{L^\infty}^{1/2+\alpha/2}, \quad (10.40)$$

so that (10.36) is satisfied.

Case 2. $\frac{\|v\|_{L^\infty}}{\|f\|_{L^\infty}} > d^2$.

In this case, we apply (10.39) with $\lambda = d$. We obtain

$$\begin{aligned} |\nabla v|_{0,\alpha,B_{\lambda/2}(x_0)} &\leq C(d^{-1-\alpha}\|v\|_{L^\infty} + d^{1-\alpha}\|f\|_{L^\infty}) \\ &\leq C\left(d^{-1-\alpha}\|v\|_{L^\infty} + \|v\|_{L^\infty}^{1/2-\alpha/2} \|f\|_{L^\infty}^{1/2+\alpha/2}\right), \end{aligned} \quad (10.41)$$

so that in both cases (10.36) is satisfied.

Once (10.36) is established, we easily obtain the interior estimate corresponding to (10.35). Indeed, standard elliptic estimates [13, Theorem 3.7] imply $\|v\|_{L^\infty} \leq C\|f\|_{L^\infty}$, so that from (10.36) we obtain

$$|\nabla v|_{0,\alpha,K} \leq C\|f\|_{L^\infty}^{1/2+\alpha/2} \|v\|_{L^\infty}^{1/2-\alpha/2}, \quad (10.42)$$

for every compact set $K \subset G$.

The proof of the boundary version of (10.42) is also a straightforward adaptation of the corresponding estimate established in [7, proof of Lemma A.2], and we omit it here. \square

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