

# **Semiclassical measure for the solution of the Helmholtz equation with an unbounded source**

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### **To cite this version:**

Julien Royer. Semiclassical measure for the solution of the Helmholtz equation with an unbounded source. Asymptotic Analysis, 2015, 91. hal-00784668

## **HAL Id: hal-00784668 <https://hal.science/hal-00784668v1>**

Submitted on 4 Feb 2013

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#### SEMICLASSICAL MEASURE FOR THE SOLUTION OF THE HELMHOLTZ EQUATION WITH AN UNBOUNDED SOURCE

#### JULIEN ROYER

ABSTRACT. We study the high frequency limit for the dissipative Helmholtz equation when the source term concentrates on a submanifold of  $\mathbb{R}^n$ . We prove that the solution has a unique semi-classical measure, which is precisely described in terms of the classical properties of the problem. This result is already known when the micro-support of the source is bounded, we now consider the general case.

#### 1. STATEMENT OF THE RESULT

We consider on  $\mathbb{R}^n$  the Helmholtz equation

<span id="page-1-0"></span>
$$
(H_h - E_h)u_h = S_h, \text{ where } H_h = -h^2 \Delta + V_1(x) - ihV_2(x). \tag{1.1}
$$

Here  $V_1$  and  $V_2$  are smooth and real-valued potentials which go to 0 at infinity. Thus for any  $h \in ]0, 1]$  the operator  $H_h$  is a non-symmetric (unless  $V_2 = 0$ ) Schrödinger operator with domain  $H^2(\mathbb{R}^n)$ . The energy  $E_h$  will be chosen in such a way that for  $h > 0$  small enough,  $\delta > \frac{1}{2}$ and  $S_h \in L^{2,\delta}(\mathbb{R}^n)$  the equation [\(1.1\)](#page-1-0) has a unique outgoing solution  $u_h \in L^{2,-\delta}(\mathbb{R}^n)$ . Here we denote by  $L^{2,\delta}(\mathbb{R}^n)$  the weighted space  $L^2(\langle x \rangle^{2\delta} dx)$ , where  $\langle x \rangle = (1 + |x|^2)^{\frac{1}{2}}$ .

The source term  $S_h$  we consider is a profile which concentrates on a submanifold  $\Gamma$  of dimension  $d \in [0, n-1]$  in  $\mathbb{R}^n$ , endowed with the Lebesgue measure  $\sigma_{\Gamma}$ . Given an amplitude  $A \in C^{\infty}(\Gamma)$ and S in the Schwartz space  $\mathcal{S}(\mathbb{R}^n)$  we set, for  $h \in ]0,1]$  and  $x \in \mathbb{R}^n$ :

<span id="page-1-1"></span>
$$
S_h(x) = h^{\frac{1-n-d}{2}} \int_{\Gamma} A(z) S\left(\frac{x-z}{h}\right) d\sigma_{\Gamma}(z)
$$
\n(1.2)

(this definition will make sense with the assumptions on  $\Gamma$  and A given below). Our purpose is to study the semiclassical measures for the family of corresponding solutions  $(u_h)$  when the submanifold  $\Gamma$  is allowed to be unbounded.

This work comes after a number of contributions which deal with more and more general situations. The first paper about the subject is [\[BCKP02\]](#page-15-0), where  $\Gamma = \{0\}$  (see also [\[Cas05\]](#page-15-1)). The result was generalized in [\[CPR02\]](#page-15-2) to the case where  $\Gamma$  is an affine subspace of  $\mathbb{R}^n$ , under the assumption that the refraction index is constant  $(V_1 = 0)$ . This restriction was overcome in [\[WZ06\]](#page-15-3). In [\[Fou06\]](#page-15-4) the source term concentrates on two points and in [\[Fou06\]](#page-15-4) the refraction index is discontinuous along an hyperplane of  $\mathbb{R}^n$ . All these papers study the semiclassical measure of the solution using its Wigner transform.

The approach in [\[Bon09\]](#page-15-5) is different. The semiclassical measures are defined with pseudodifferential calculus (see  $(1.14)$ ) and the resolvent is replaced by the integral over positive times of the propagator (as in  $[Cas05]$ ). We used this point of view in  $[Roy10]$  to deal with the case of a non-constant absorption index ( $V_2 \neq 0$ ,  $V_2 \geq 0$ ) and a general bounded submanifold Γ. We also considered in  $[{\rm Roy11}]$  an absorption index  $V_2$  which can take non-positive values. The purpose

<sup>2010</sup> Mathematics Subject Classification. 35J10, 47A55, 47B44, 47G30.

Key words and phrases. Non-selfadjoint operators, resolvent estimates, Helmholtz equation, semiclassical measures.

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of this paper is now to allow a general unbounded submanifold Γ.

Let us now state more precisely the assumptions. The potentials  $V_1$  and  $V_2$  are respectively of long and of short range: there exist constants  $\rho > 0$  and  $c_{\alpha}$  for  $\alpha \in \mathbb{N}^n$  such that

<span id="page-2-2"></span>
$$
\forall \alpha \in \mathbb{N}^n, \forall x \in \mathbb{R}^n, \quad |\partial^{\alpha} V_1(x)| \leq c_{\alpha} \langle x \rangle^{-\rho - |\alpha|} \quad \text{and} \quad |\partial^{\alpha} V_2(x)| \leq c_{\alpha} \langle x \rangle^{-1-\rho - |\alpha|}. \tag{1.3}
$$

Then we introduce the hamiltonian flow  $\phi^t$  corresponding to the classical symbol  $p : (x, \xi) \mapsto$  $|\xi|^2 + V_1(x)$  on the phase space  $\mathbb{R}^{2n}$ . For all  $w \in \mathbb{R}^{2n}$ ,  $t \mapsto \phi^t(w) = (X(t, w), \Xi(t, w)) \in \mathbb{R}^{2n}$  is the solution of the system

$$
\begin{cases}\n\partial_t X(t, w) = 2\Xi(t, w), \\
\partial_t \Xi(t, w) = -\nabla V_1(X(t, w)), \\
\phi^0(w) = w.\n\end{cases}
$$
\n(1.4)

For  $I \subset \mathbb{R}$  we set

$$
\Omega_b(I) = \left\{ w \in p^{-1}(I) : \sup_{t \in \mathbb{R}} |X(t, w)| < \infty \right\}
$$

We also denote by

$$
H_p q = \{p, q\} = \nabla_{\xi} p \cdot \nabla_x q - \nabla_x p \cdot \nabla_{\xi} q
$$
 the Poisson bracket of p with a symbol  $q \in C^{\infty}(\mathbb{R}^{2n})$ .

We now consider an energy  $E_0 > 0$  such that

<span id="page-2-0"></span>
$$
\forall w \in \Omega_b(\{E_0\}), \exists T > 0, \quad \int_0^T V_2(X(t, w)) dt > 0.
$$
 (1.5)

Let  $\delta > \frac{1}{2}$ . We know (see [\[Roy11\]](#page-15-7)) that under Assumption [\(1.5\)](#page-2-0) there exist an open neighborhood J of  $E_0$ ,  $h_0 > 0$  and  $c \ge 0$  such that for

$$
z \in \mathbb{C}_{J,+} = \{ z \in \mathbb{C} : \text{Re } z \in J, \text{Im } z > 0 \}
$$

and  $h \in ]0, h_0]$  the operator  $(H_h - z)$  has a bounded inverse on  $L^2(\mathbb{R}^n)$  and

$$
\left\| \left\langle x \right\rangle^{-\delta} \left( H_h - z \right)^{-1} \left\langle x \right\rangle^{-\delta} \right\|_{\mathcal{L}\left(L^2(\mathbb{R}^n)\right)} \leqslant \frac{c}{h}.
$$

Here  $\mathcal{L}(L^2(\mathbb{R}^n))$  denotes the space of bounded operators on  $L^2(\mathbb{R}^n)$ . Moreover for any  $\lambda \in J$  the limit

$$
(H_h - (\lambda + i0))^{-1} = \lim_{\beta \to 0^+} (H_h - (\lambda + i\beta))^{-1}
$$

exists in  $\mathcal{L}(L^{2,\delta}(\mathbb{R}^n), L^{2,-\delta}(\mathbb{R}^n)).$ 

Now let us be more explicit about the source term  $S_h$  we consider. We recall that  $\Gamma$  is a submanifold of dimension  $d \in [0, n-1]$  in  $\mathbb{R}^n$ , endowed with the Riemannian structure given by the restriction of the usual structure on  $\mathbb{R}^n$ , and the corresponding Lebesgue measure  $\sigma_{\Gamma}$ . We assume that there exist  $R_1 > 0$  and  $\sigma_1 \in ]0,1[$  such that

<span id="page-2-1"></span>
$$
N\Gamma \cap \mathcal{Z}_{-}(R_1, 0, -\sigma_1) = \emptyset, \qquad (1.6)
$$

.

where  $N\Gamma = \{(z,\xi) \in \Gamma \times \mathbb{R}^n : \xi \perp T_z\Gamma\}$  is the normal bundle of  $\Gamma$  and  $\mathcal{Z}_-(R_1,0,\sigma_1)$  is an incoming region: for  $R \geqslant 0$ ,  $\nu \geqslant 0$  and  $\sigma \in [-1, 1]$  we set

$$
\mathcal{Z}_{\pm}(R,\nu,\sigma) = \left\{ (x,\xi) \in \mathbb{R}^{2n} : |x| \geqslant R, |\xi| \geqslant \nu \text{ and } \langle x,\xi \rangle \geqslant \sigma |x| |\xi| \right\}
$$

Note that Assumption [\(1.6\)](#page-2-1) is satisfied for any bounded submanifold of  $\mathbb{R}^n$ . When  $d = 0$ , it actually implies that  $\Gamma$  is bounded, but this is not the case in higher dimension: this assumption holds for instance for any affine subspace of dimension  $d \in [1, n-1]$  in  $\mathbb{R}^n$ . Now that  $\sigma_1$  is fixed, we can choose a smaller neighborhood  $J$  of  $E_0$  and assume that

<span id="page-2-3"></span>
$$
J \subset [E_1, E_2]
$$
, were  $E_1 > 0$  and  $\left(\frac{1 + \sigma_1}{2}\right)^2 E_2 < E_1$ . (1.7)

We assume that

<span id="page-2-4"></span>
$$
\forall z \in \Gamma, \quad V_1(z) > E_0,\tag{1.8}
$$

and we define

$$
N_E\Gamma = N\Gamma \cap p^{-1}(\{E_0\}).
$$

 $N_E\Gamma$  is a submanifold of dimension  $(n-1)$  in  $\mathbb{R}^{2n}$ , endowed with the Riemannian structure defined as follows: for  $(z,\xi) \in N_E \Gamma$  and  $(Z,\Xi),(\tilde{Z},\tilde{\Xi}) \in T_{(z,\xi)}N_E \Gamma \subset \mathbb{R}^{2n}$  we set

$$
g_{(z,\xi)}\big((Z,\Xi),(\tilde{Z},\tilde{\Xi})\big)=\left\langle Z,\tilde{Z}\right\rangle_{\mathbb{R}^n}+\left\langle \Xi_{\perp},\tilde{\Xi}_{\perp}\right\rangle_{\mathbb{R}^n},
$$

where  $\Xi_{\perp}$ ,  $\tilde{\Xi}_{\perp}$  are the orthogonal projections of  $\Xi$ ,  $\tilde{\Xi} \in \mathbb{R}^n$  on  $(T_z\Gamma \oplus \mathbb{R}\xi)^{\perp}$  (see the discussion in [\[Roy10\]](#page-15-6)). We denote by  $\sigma_{N_{E}\Gamma}$  the canonical measure on  $N_{E}\Gamma$  given by g, and assume that

<span id="page-3-3"></span>
$$
\sigma_{N_E \Gamma} \left( \left\{ (z,\xi) \in N_E \Gamma : \exists t > 0, \phi^t(z,\xi) \in N_E \Gamma \right\} \right) = 0. \tag{1.9}
$$

<span id="page-3-4"></span>We now introduce the amplitude  $A \in C^{\infty}(\Gamma)$ . We assume that there exist  $\delta > \frac{1}{2}$  and  $c \geq 0$ such that

$$
\int_{\Gamma} \langle z \rangle^{\delta} \left( |A(z)| + \|d_z A\| + |A(z)| \left\| \Pi_z \right\| \right) d\sigma_{\Gamma}(z) < +\infty. \tag{1.10a}
$$

Moreover for all  $r \in ]0,1]$  and  $x \in \mathbb{R}^n$  we have

$$
\int_{B(x,r)\cap\Gamma} \langle z \rangle^{\delta} \left( |A(z)| + \|d_z A\| + |A(z)| \|\Pi_z\| \right) d\sigma_{\Gamma}(z) \leqslant cr^d. \tag{1.10b}
$$

Here  $B(x, r)$  is the ball of radius r and centered at  $x, d_zA : T_z\Gamma \to \mathbb{R}$  is the differential of A at point z and II is the second fundamental form of the submanifold Γ. For any  $z \in \Gamma$ ,  $\prod_z$  is a bilinear form from  $T_z\Gamma$  to  $N_z\Gamma$  (see Appendix [A\)](#page-14-0), and

$$
\|{\rm II}_{z}\|=\sup_{\|X\|_{T_{z}\Gamma}= \|Y\|_{T_{z}\Gamma}=1}\|{\rm II}_{z}(X,Y)\|_{N_{z}\Gamma}\,.
$$

Note that all these estimates hold when  $A \in C_0^{\infty}(\Gamma)$ . Here A is allowed to have a non-compact support, but it still has to stay away from the bundary of  $\Gamma$ :

<span id="page-3-5"></span>
$$
\forall \theta \in C_0^{\infty}(\mathbb{R}^n), \quad z \mapsto A(z)\theta(z) \in C_0^{\infty}(\Gamma). \tag{1.11}
$$

Then it remains to consider  $S \in \mathcal{S}(\mathbb{R}^n)$  and define the source term  $S_h$  by [\(1.2\)](#page-1-1).

Let  $(E_h)_{h\in[0,h_0]}$  be a family of energy in  $\mathbb{C}_{J,+}\cup J$  such that

<span id="page-3-1"></span>
$$
E_h = E_0 + h\tilde{E} + \underset{h \to 0}{\circ} (h) \tag{1.12}
$$

for some  $\tilde{E} \in \overline{\mathbb{C}_+}.$  Since for all  $h \in ]0, h_0]$  the source term  $S_h$  given by  $(1.2)$  belongs to  $L^{2,\delta}(\mathbb{R}^n)$ (see Proposition [2.2\)](#page-6-0) we can define

<span id="page-3-2"></span>
$$
u_h = (H_h - (E_h + i0))^{-1} S_h \in L^{2,-\delta}(\mathbb{R}^n).
$$
\n(1.13)

Here  $(H_h - (E_h + i0))^{-1}$  stands for  $(H_h - E_h)^{-1}$  when  $E_h \in \mathbb{C}_{J,+}$ .

Our purpose is to study the semiclassical measures for this family  $(u_h)_{h\in]0,h_0]}$ . In other words, non-negative measures  $\mu$  on the phase space  $\mathbb{R}^{2n} \simeq T^* \mathbb{R}^n$  such that

<span id="page-3-0"></span>
$$
\forall q \in C_0^{\infty}(\mathbb{R}^{2n}), \quad \left\langle \text{Op}_{h_m}^w(q)u_{h_m}, u_{h_m} \right\rangle \xrightarrow[m \to +\infty]{} \int_{\mathbb{R}^{2n}} q \, d\mu,
$$
\n(1.14)

for some sequence  $(h_m)_{m \in \mathbb{N}} \in ]0, h_0]^{\mathbb{N}}$  such that  $h_m \to 0$  (see [Gér91]). Here  $\text{Op}_h^w(q)$  denotes the Weyl quantization of the symbol  $q$ :

$$
\mathrm{Op}_{h}^{w}(q)u(x) = \frac{1}{(2\pi h)^{n}} \int_{\mathbb{R}^{n}} \int_{\mathbb{R}^{n}} e^{\frac{i}{h}\langle x-y,\xi\rangle} q\left(\frac{x+y}{2},\xi\right) u(y) \, dy \, d\xi.
$$

We will also use the standard quantization:

$$
\mathrm{Op}_h(q)u(x) = \frac{1}{(2\pi h)^n} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} e^{\frac{i}{h}\langle x-y,\xi\rangle} q(x,\xi)u(y) \, dy \, d\xi.
$$

We denote by  $C_b^{\infty}(\mathbb{R}^{2n})$  the set of smooth symbols whose derivatives are bounded. For  $\delta \in \mathbb{R}$ , we also denote by  $\mathcal{S}(\langle x \rangle^{\delta})$  the set of symbols  $a \in C^{\infty}(\mathbb{R}^{2n})$  such that

$$
\forall \alpha, \beta \in \mathbb{N}^n, \exists c_{\alpha, \beta} \geqslant 0, \forall (x, \xi) \in \mathbb{R}^{2n}, \quad \left| \partial_x^{\alpha} \partial_{\xi}^{\beta} a(x, \xi) \right| \leqslant c_{\alpha, \beta} \left\langle x \right\rangle^{\delta},
$$

and by  $\mathcal{S}_{\delta}(\mathbb{R}^{2n})$  the set of symbols  $a \in C^{\infty}(\mathbb{R}^{2n})$  such that

$$
\forall \alpha, \beta \in \mathbb{N}^n, \exists c_{\alpha, \beta} \geqslant 0, \forall (x, \xi) \in \mathbb{R}^{2n}, \quad \left| \partial_x^{\alpha} \partial_{\xi}^{\beta} a(x, \xi) \right| \leqslant c_{\alpha, \beta} \left\langle x \right\rangle^{\delta - |\alpha|}.
$$

We can similarly define the sets of symbols  $\mathcal{S}(\langle \xi \rangle^{\delta})$  for  $\delta \in \mathbb{R}$ . We refer to [\[Rob87,](#page-15-9) [Zwo12\]](#page-15-10) for more details about semiclassical analysis.

For  $(z, \xi) \in N_E \Gamma$  we set

$$
\kappa(z,\xi) = \pi (2\pi)^{d-n} |A(z)|^2 |\xi|^{-1} |\hat{S}(\xi)|^2,
$$

where  $\hat{S}$  is the Fourier transform of S. The theorem we want to prove is the following:

<span id="page-4-0"></span>**Theorem 1.1.** Let  $S_h$  and  $E_h$  be given by [\(1.2\)](#page-1-1) and [\(1.12\)](#page-3-1), and  $u_h$  defined by [\(1.13\)](#page-3-2). Assume *that the assumptions* [\(1.3\)](#page-2-2)*,* [\(1.5\)](#page-2-0)*,* [\(1.6\)](#page-2-1)*,* [\(1.7\)](#page-2-3)*,* [\(1.8\)](#page-2-4)*,* [\(1.9\)](#page-3-3)*,* [\(1.10\)](#page-3-4) *and* [\(1.11\)](#page-3-5) *hold.*

(*i*) Then there exists a non-negative Radon measure  $\mu$  on  $\mathbb{R}^{2n}$  such that for all  $q \in C_0^{\infty}(\mathbb{R}^{2n})$ *we have*

<span id="page-4-2"></span>
$$
\langle \operatorname{Op}_h^w(q)u_h, u_h \rangle \xrightarrow[h \to 0]{} \int_{\mathbb{R}^{2n}} q \, d\mu. \tag{1.15}
$$

- *(ii) This measure is characterized by the following three properties:*
	- *a.*  $\mu$  *is supported in*  $p^{-1}(\{E_0\})$ *.*
	- *b.* For all  $\sigma \in ]\sigma_1, 1[$  there exists  $R \geq 0$  such that  $\mu$  is zero in the incoming region  $\mathcal{Z}_-(R, 0, -\sigma)$ .
	- *c.* µ *satisfies the following Liouville equation:*

$$
(H_p + 2\operatorname{Im} \tilde{E} + 2V_2)\mu = \kappa \sigma_{N_E\Gamma}.
$$

*This means that for all*  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  *we have* 

$$
\int_{\mathbb{R}^{2n}} (-H_p + 2 \operatorname{Im} \tilde{E} + 2V_2) q d\mu = \int_{N_E \Gamma} q(z, \xi) \kappa(z, \xi) d\sigma_{N_E \Gamma}(z, \xi).
$$

*(iii)* These three properties imply that for any  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  we have

<span id="page-4-1"></span>
$$
\int_{\mathbb{R}^{2n}} q \, d\mu = \int_{\mathbb{R}_+} \int_{N_E \Gamma} \kappa(z, \xi) q(\phi^t(z, \xi)) e^{-2t \operatorname{Im} \tilde{E} - 2 \int_0^t V_2(X(s, z, \xi)) ds} \, d\sigma_{N_E \Gamma}(z, \xi) \, dt. \tag{1.16}
$$

This result is known when A is compactly supported on  $\Gamma$  (see [\[Roy10,](#page-15-6) [Roy11\]](#page-15-7)). The idea for the proof is to write the resolvent as the integral of the propagator over positive times, and to approximate  $u_h$  by a partial solution which only takes into account finite times:

$$
u_h^T = \frac{i}{h} \int_0^\infty \chi_T(t) e^{-\frac{it}{h}(H_h - E_h)} S_h dt.
$$

Here  $\chi_T(t) = \chi(t-T)$ , where  $\chi \in C^{\infty}(\mathbb{R}, [0,1])$  is equal to 1 in a neighborhood of  $]-\infty,0]$  and supported in  $]-\infty, \tau_0]$  for some well-chosen  $\tau_0 \in ]0,1]$ . Note that for all  $h \in ]0,1]$  the semi-group  $t \mapsto e^{-\frac{it}{h}H_h}$  is well-defined for all  $t \geq 0$ . However this is not a contractions semi-group since  $V_2$ is not assumed to be non-negative (see for instance Corollary 3.6 in [\[EN00\]](#page-15-11))

The idea will be the same to deal with the case of an amplitude A whose support is not bounded. Let  $\Theta \in C_0^{\infty}(\mathbb{R}^n, [0, 1])$  be equal to 1 on  $B(0, 1)$ . For any  $R > 0$  we set  $A_R : z \in \Gamma \mapsto$  $A(z)\Theta(z/R)$  and

$$
S_h^R(x) = h^{\frac{1-n-d}{2}} \int_{z \in \Gamma} A_R(z) S\left(\frac{x-z}{h}\right) d\sigma_{\Gamma}(z).
$$
 (1.17)

Given any  $R > 0$ , the proof of [\[Roy10,](#page-15-6) [Roy11\]](#page-15-7) applies for the source term  $S_h^R$ . Since the choice of  $\chi$  mentionned above depends on the support of  $A_R$ , we denote it by  $\chi_{0,R}$ . Moreover  $\chi_{0,R}$  can be chosen non-increasing. Then for any  $T \geq 0$  we set  $\chi_{T,R}: t \mapsto \chi_{0,R}(t-T)$  and for any  $h \in ]0, h_0]$ :

$$
u_h^{T,R} = \frac{i}{h} \int_0^\infty \chi_{T,R}(t) e^{-\frac{it}{h}(H_h - E_h)} S_h^R dt.
$$

The key point is to prove that in some suitable sense  $u_h^{T,R}$  is a good approximation of  $u_h$  for large T and R, and  $h > 0$  small enough.

Let  $R > 0$ . For  $h \in ]0, h_0]$  we set

$$
\tilde{u}_h^R = (H_h - (E_h + i0))^{-1} S_h^R \in L^{2, -\delta}(\mathbb{R}^n).
$$

Since  $A_R$  is compactly supported on  $\Gamma$ , we know that Theorem [1.1](#page-4-0) holds for  $\tilde{u}_h^R$ . In particular there exists a non-negative Radon measure  $\tilde{\mu}_R$  on  $\mathbb{R}^{2n}$  such that

$$
\forall q \in C_0^{\infty}(\mathbb{R}^{2n}), \quad \left\langle \text{Op}_{h}^w(q)\tilde{u}_h^R, \tilde{u}_h^R \right\rangle \xrightarrow[h \to 0]{} \int_{\mathbb{R}^{2n}} q \, d\tilde{\mu}_R.
$$

Moreover, according to [\(1.16\)](#page-4-1),  $\tilde{\mu}_R$  is supported on the classical trajectories coming from  $N_E^R \Gamma =$  $\{(z,\xi)\in N_E\Gamma : z\in \text{supp }A_R\}$ . Let K be a compact subset of  $\mathbb{R}^{2n}$ . Assumption  $(1.6)$  ensures that for  $R_K > 0$  large enough and  $R \ge R_K$ , the trajectories coming form  $N_E^R \Gamma \setminus N_E^{R_K} \Gamma$  do not meet K (see Proposition [3.1\)](#page-10-0), and hence  $\tilde{\mu}_R = \tilde{\mu}_{R_K}$  on K. This is the idea we are going to use to prove existence of the semiclassical measure  $\mu$ . And as expected,  $\mu$  will coincide with  $\tilde{\mu}_{R_K}$  on K.

The plan of this paper is the following. In section [2](#page-5-0) we give some estimates for the source term  $S_h$ , and in Section [3](#page-10-1) we show that  $u_h^{T,R}$  is a good approximation of  $u_h$  in order to prove Theorem [1.1.](#page-4-0) In Appendix [A](#page-14-0) we recall some basic facts about differential geometry, and in particular the second fundamental form which appears when integrating by parts on Γ.

#### 2. Estimates of the source term

<span id="page-5-0"></span>In this section we prove that  $S_h$  and  $S_h^R$  for  $R > 0$  are (uniformly) of size  $O(\sqrt{h})$  in  $L^{2,\delta}(\mathbb{R}^n)$ , where  $\delta > \frac{1}{2}$  is given by [\(1.10\)](#page-3-4). Then we use Assumption [\(1.6\)](#page-2-1) to prove that if  $\omega_{-} \in \mathcal{S}_0(\mathbb{R}^{2n})$ is supported in  $\mathcal{Z}_-(R, 0, -\sigma)$  for some  $R > R_1$  and  $\sigma \in ]\sigma_1, 1[$ , then  $\text{Op}_h^w(\omega_-)S_h = O(h^{\frac{3}{2}})$  in  $L^{2,\delta}(\mathbb{R}^n)$ . Since we even have an estimate of size  $O(h^{\infty})$  when  $\omega_{-}$  is compactly supported, this proves in particular that  $S_h$  is microlocally supported outside an incoming region.

<span id="page-5-4"></span>**Lemma 2.1.** *Consider*  $B \in C^{\infty}(\Gamma)$ *, a family*  $(f_h^z)_{z \in \Gamma, h \in ]0,1]}$  *of functions in*  $L^{2,2+\delta+d/2}(\mathbb{R}^n)$ *, and assume that for some*  $C_1 \geq 0$  *we have* 

<span id="page-5-2"></span>
$$
\forall h \in ]0,1], \quad \int_{\Gamma} \left\langle z \right\rangle^{\delta} |B(z)| \left( 1 + \|f_h^z\|_{L^{2,2+\delta+d/2}(\mathbb{R}^n)}^2 \right) d\sigma_{\Gamma}(z) < C_1 \tag{2.1}
$$

*and*

<span id="page-5-1"></span>
$$
\forall x \in \mathbb{R}^n, \forall r \in ]0,1], \quad \int_{B(x,r)\cap\Gamma} \langle z \rangle^{\delta} |B(z)| \, d\sigma_{\Gamma}(z) \leq C_1 \, r^d. \tag{2.2}
$$

*For*  $h \in ]0,1]$  *we consider* 

$$
\tilde{S}_h: x \mapsto h^{\frac{1-n-d}{2}} \int_{\Gamma} B(z) f_h^z \left( \frac{x-z}{h} \right) d\sigma_{\Gamma}(z).
$$

*Then*  $\tilde{S}_h(x)$  *is well-defined for all*  $h \in ]0,1]$  *and*  $x \in \mathbb{R}^n$ *, and there exists a constant*  $C \geq 0$  *which only depends on* C<sup>1</sup> *and such that*

$$
\forall h \in ]0,1], \quad \left\|\tilde{S}_h\right\|_{L^{2,\delta}(\mathbb{R}^n)} \leqslant C\sqrt{h}.
$$

The idea of the proof is the same as for a compactly supported amplitude but we now have to be careful with the decay at infinity for functions on Γ.

*Proof.* We first remark that Assumption [\(2.2\)](#page-5-1) holds for all  $r > 0$  since for  $x \in \mathbb{R}^n$  and  $r \geq 1$ Assumption [\(2.1\)](#page-5-2) gives

<span id="page-5-3"></span>
$$
\int_{B(x,r)\cap\Gamma} \langle z \rangle^{\delta} |B(z)| \, d\sigma_{\Gamma}(z) \leqslant C_1 \leqslant r^d C_1. \tag{2.3}
$$

Let  $h \in ]0,1]$  and  $x \in \mathbb{R}^n$ . According to Cauchy-Schwarz inequality and  $(2.3)$  we can write

$$
\left(\int_{\Gamma} |B(z)| \left| f_h^z \left( \frac{x-z}{h} \right) \right| d\sigma_{\Gamma}(z) \right)^2 = \left( \sum_{m \in \mathbb{N}} \int_{mh \leq |x-z| < (m+1)h} |B(z)| \left| f_h^z \left( \frac{x-z}{h} \right) \right| d\sigma_{\Gamma}(z) \right)^2
$$
  

$$
\leq c \sum_{m \in \mathbb{N}} \langle m \rangle^2 \left( \int_{mh \leq |x-z| < (m+1)h} |B(z)| \left| f_h^z \left( \frac{x-z}{h} \right) \right| d\sigma_{\Gamma}(z) \right)^2
$$
  

$$
\leq c C_1 h^d \sum_{m \in \mathbb{N}} \langle m \rangle^{2+d} \int_{mh \leq |x-z| < (m+1)h} \langle z \rangle^{-\delta} |B(z)| \left| f_h^z \left( \frac{x-z}{h} \right) \right|^2 d\sigma_{\Gamma}(z),
$$

where  $c \geqslant 0$  stands for different universal constants. Now using  $(2.1)$  we obtain  $||S_h||^2_{L^{2,\delta}(\mathbb{R}^n)}$ 

$$
\leqslant c C_1 h^{1-n} \int_{\mathbb{R}^n} \langle x \rangle^{2\delta} \sum_{m \in \mathbb{N}} \langle m \rangle^{2+d} \int_{mh \leqslant |x-z| < (m+1)h} \langle z \rangle^{-\delta} |B(z)| \left| f_h^z \left( \frac{x-z}{h} \right) \right|^2 d\sigma_{\Gamma}(z) dx
$$
  
\n
$$
\leqslant c C_1 h \sum_{m \in \mathbb{N}} \langle m \rangle^{2+d} \int_{\Gamma} \int_{m \leqslant |y| < m+1} \langle z + hy \rangle^{2\delta} \langle z \rangle^{-\delta} |B(z)| |f_h^z(y)|^2 dy d\sigma_{\Gamma}(z)
$$
  
\n
$$
\leqslant c C_1 h \sum_{m \in \mathbb{N}} \langle m \rangle^{-2} \int_{\Gamma} \langle z \rangle^{\delta} |B(z)| \int_{m \leqslant |y| < m+1} \langle y \rangle^{4+d+2\delta} |f_h^z(y)|^2 dy d\sigma_{\Gamma}(z)
$$
  
\n
$$
\leqslant c C_1^2 h.
$$

 $\Box$ 

Applied with  $f_h^z = S$  and  $B = A$  or  $A_R$  for  $R > 0$  this proposition gives:

<span id="page-6-0"></span>**Proposition 2.2.** *There exists a constant*  $C \geq 0$  *such that* 

$$
\forall h \in ]0,1], \forall R > 0, \quad ||S_h||_{L^{2,\delta}(\mathbb{R}^n)} + ||S_h^R||_{L^{2,\delta}(\mathbb{R}^n)} \leqslant C\sqrt{h}.
$$

<span id="page-6-1"></span>For  $z \in \Gamma$  and  $\xi \in \mathbb{R}^n$  we denote by  $\xi_z^T$  the orthogonal projection of  $\xi$  on  $T_z \Gamma$ , and  $\xi_z^N = \xi - \xi_z^T$ . **Proposition 2.3.** Let  $q \in C_b^{\infty}(\mathbb{R}^{2n})$  and assume that for some  $\varepsilon > 0$  we have

 $\forall (x, \xi) \in \text{supp } q, \forall z \in \text{supp } A, \quad (|x - z| \geqslant \varepsilon \quad or \quad |\xi_z^T| \geqslant \varepsilon).$ 

*Then we have*

$$
\|\text{Op}_h(q)S_h\|_{L^{2,\delta}(\mathbb{R}^n)} = \bigcap_{h\to 0} \left(h^{\frac{3}{2}}\right).
$$

*Proof.* • We have

$$
\mathrm{Op}_h(q)S_h(x) = \frac{h^{\frac{1-n-d}{2}}}{(2\pi h)^n} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \int_{\Gamma} e^{\frac{i}{h}\langle x-y,\xi\rangle} q(x,\xi)A(z)S\left(\frac{y-z}{h}\right) d\sigma_{\Gamma}(z) dy d\xi.
$$

We recall that this is defined in the sense of an oscillatory integral. After a finite number of partial integrations with the operator  $\frac{1+ih\xi\cdot\nabla_y}{1+|\xi|^2}$  we can assume that  $q \in \mathcal{S}(\langle\xi\rangle^{-(n+1)})$ . Then we can use Fubini's Theorem and make the change of variables  $y = z + hv$  for any fixed z. Let  $\chi_1 \in C_0^{\infty}(\mathbb{R}^n, [0,1])$  be supported in  $B(0,\varepsilon)$  and equal to 1 on a neighborhood of 0. We set  $\chi_2 = 1 - \chi_1$ . For  $x \in \mathbb{R}^n$  and  $h \in ]0,1]$  we can write

$$
Op_h(q)S_h(x) = I_1(x,h) + I_2(x,h),
$$

where for  $j \in \{1, 2\}$ :

$$
I_j(x,h) = \frac{h^{\frac{1-n-d}{2}}}{(2\pi)^n} \int_{\Gamma} A(z) \chi_j(x-z) \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} e^{\frac{i}{h} \langle x-z,\xi \rangle} e^{-i \langle v,\xi \rangle} q(x,\xi) S(v) dv d\xi d\sigma_{\Gamma}(z).
$$

• Let  $N \in \mathbb{N}$  and

$$
B(x, z, h) = \frac{\chi_2(x - z)}{(2\pi)^n} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} e^{\frac{i}{h}\langle x - z, \xi \rangle} S(v) L^N \left( e^{-i\langle v, \xi \rangle} q(x, \xi) \right) dv d\xi,
$$

where for  $f \in C^{\infty}(\mathbb{R}^n \times \Gamma \times \mathbb{R}^n \times ]0,1]$  we have set

$$
Lv: (x, z, \xi, h) \mapsto -ih \operatorname{div}_{\xi} \left( \frac{(x-z)f(x, z, \xi, h)}{|x-z|^2} \right).
$$

Then there exists a constant c such that

$$
\forall h \in ]0,1], \forall x \in \mathbb{R}^n, \forall z \in \Gamma, \quad |B(x,z,h)| \leqslant ch^N \left\langle x-z\right\rangle^{-N},
$$

and hence

$$
\begin{split} |I_2(x,h)|^2 &= h^{1-n-d} \left| \int_{\Gamma} A(z)B(x,z,h) \, d\sigma_{\Gamma}(z) \right|^2 \\ &\le ch^{2N+1-n-d} \left( \int_{\Gamma} \langle z \rangle^{\delta} |A(z)| \, d\sigma_{\Gamma}(z) \right) \left( \int_{\Gamma} \langle z \rangle^{-\delta} |A(z)| \langle x-z \rangle^{-2N} \, d\sigma(z) \right), \end{split}
$$

where c depends neither on  $x \in \mathbb{R}^n$  nor on  $h \in ]0,1]$ . We obtain

$$
||I_2(h)||_{L^{2,\delta}(\mathbb{R}^n)}^2 \leq c h^{2N+1-n-d} \int_{\Gamma} \int_{\mathbb{R}^n} \langle z \rangle^{-\delta} |A(z)| \langle x \rangle^{2\delta} \langle x-z \rangle^{-2N} dx d\sigma(z)
$$
  

$$
\leq c h^{2N+1-n-d} \int_{\Gamma} \int_{\mathbb{R}^n} \langle z \rangle^{\delta} |A(z)| \langle x-z \rangle^{2\delta} \langle x-z \rangle^{-2N} dx d\sigma(z)
$$

and finally, if  $N$  was chosen large enough:

$$
||I_2(h)||_{L^{2,\delta}(\mathbb{R}^n)}^2 = O_{h\to 0}(h^3).
$$

• We now turn to  $I_1$ . Let  $x, \xi \in \mathbb{R}^n$ . The function  $z \mapsto \xi_z^T \in T_z\Gamma$  defines a vector field on  $\Gamma$ , which we denote by  $\xi^T$ , and the norm of  $\xi^T$  in  $T_z\Gamma$  is the same as in  $\mathbb{R}^n$ . By assumption, if  $A(z)q(x,\xi)\chi_1(x-z)\neq 0$  then  $\left|\xi_z^T\right|\geqslant \varepsilon$ . And we remark that when  $\xi_z^T\neq 0$  we have

$$
e^{\frac{i}{h}\langle x-z,\xi\rangle} = i h \big|\xi_z^T\big|^{-2} \xi_z^T \cdot e^{\frac{i}{h}\langle x-z,\xi\rangle},
$$

where  $\xi_z^T \cdot f(z)$  is the derivative of  $f \in C^{\infty}(\Gamma)$  at point z and in the direction of  $\xi_z^T$ . For any  $z \in \text{supp } A \cap B(x, \varepsilon)$  (which is compact according to Assumption [\(1.11\)](#page-3-5)) there exists an open neighborhood  $v_z$  of z in  $\Gamma$  which is orientable, and we can find  $z_1, \ldots, z_K \in \text{supp } A \cap B(x, \varepsilon)$  such that supp  $A \cap \overline{B(x,\varepsilon)} \subset \bigcup_{k=1}^K \mathcal{V}_{z_k}$ . We consider  $\zeta_1,\ldots,\zeta_K \in C_0^\infty(\Gamma,[0,1])$  such that  $\sum_{k=1}^K \zeta_k = 1$ in a neighborhood of supp  $A \cap B(x, \varepsilon)$  in  $\Gamma$  and  $\zeta_k$  is supported in  $\mathcal{V}_k$  for all  $k \in [1, K]$ . Let  $k \in [1, K]$ . According to Green's Theorem [A.1](#page-15-12) we have

$$
\int_{\Gamma} \operatorname{div} \left( e^{\frac{i}{\hbar} \langle x - z, \xi \rangle} \chi_1(x - z) (A \zeta_k)(z) \left| \xi_z^T \right|^{-2} \xi_z^T \right) d\sigma_{\Gamma}(z) = 0
$$

and hence

$$
\int_{\Gamma} e^{\frac{i}{h}\langle x-z,\xi\rangle} (A\zeta_k)(z) \chi_1(x-z) d\sigma_{\Gamma}(z)
$$
\n
$$
= -ih \int_{\Gamma} \operatorname{div} \xi^T(z) |\xi_z^T|^{-2} e^{\frac{i}{h}\langle x-z,\xi\rangle} (A\zeta_k)(z) \chi_1(x-z) d\sigma_{\Gamma}(z)
$$
\n
$$
-ih \int_{\Gamma} e^{\frac{i}{h}\langle x-z,\xi\rangle} \xi_z^T \cdot ((A\zeta_k)(z) \chi_1(x-z)) |\xi_z^T|^{-2} d\sigma_{\Gamma}(z)
$$
\n
$$
+ih \int_{\Gamma} e^{\frac{i}{h}\langle x-z,\xi\rangle} (A\zeta_k)(z) \chi_1(x-z) |\xi_z^T|^{-4} \xi_z^T \cdot |\xi_z^T|^2 d\sigma_{\Gamma}(z).
$$

Taking the sum over  $k \in [1, K]$  gives

$$
I_1(h) = -ih(I_{1,1}(h) + I_{1,2}(h) + I_{1,3}(h))
$$

where, for instance,

$$
I_{1,1}(x,h)
$$
\n
$$
= \frac{h^{\frac{1-n-d}{2}}}{(2\pi)^n} \int_{\Gamma} A(z) \chi_1(x-z) \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \text{div}\,\xi^T(z) |\xi_z^T|^{-2} e^{\frac{i}{h}\langle x-z,\xi\rangle} e^{-i\langle v,\xi\rangle} q(x,\xi) S(v) \,dv \,d\xi \,d\sigma_{\Gamma}(z)
$$
\n
$$
= \frac{h^{\frac{1-n-d}{2}}}{(2\pi)^n} \int_{\Gamma} A(z) \langle ||\Pi_z|| \rangle \chi_1(x-z) \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} e^{\frac{i}{h}\langle x-z,\xi\rangle} e^{-i\langle v,\xi\rangle} q_{1,z}(x,\xi) S(v) \,dv \,d\xi \,d\sigma_{\Gamma}(z)
$$
\n
$$
= h^{\frac{1-n-d}{2}} \int_{\Gamma} A(z) \langle ||\Pi_z|| \rangle \chi_1(x-z) \big(\text{Op}_h(q_{1,z})S\big) \left(\frac{x-z}{h}\right) \,d\sigma_{\Gamma}(z).
$$

Here we have set

$$
q_{1,z}(x,\xi) = \langle ||\Pi_z|| \rangle^{-1} \operatorname{div} \xi^T(z) |\xi_z^T|^{-2} q(x,\xi).
$$

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• We recall that the Levi-Civita connection  $\nabla^{\Gamma}$  on the submanifold  $\Gamma$  at point z is given by the orthogonal projection on  $T_z\Gamma$  of the usual differential on  $\mathbb{R}^n$  (see Appendix [A\)](#page-14-0). Let  $z \in \Gamma$ . Let Y be a vector field on a neighborhood of z in  $\Gamma$  and let  $\overline{Y}$  be an extension of Y on a neighborhood of z in  $\mathbb{R}^n$ . Using [\(A.1\)](#page-14-1) we see that on a neighborhood of z in  $\Gamma$  we have

$$
\langle \nabla_Y \xi^T, Y \rangle_{T\Gamma} = Y \cdot \langle \xi^T, Y \rangle_{T\Gamma} - \langle \xi^T, \nabla_Y Y \rangle_{T\Gamma} = \overline{Y} \cdot \langle \xi, \overline{Y} \rangle_{\mathbb{R}^n} - \langle \xi^T, \nabla_Y Y \rangle_{T\Gamma}
$$
  
=  $\langle \xi, \nabla_Y Y + \Pi(Y, Y) \rangle_{\mathbb{R}^n} - \langle \xi^T, \nabla_Y Y \rangle_{T\Gamma}$   
=  $\langle \xi^N, \Pi(Y, Y) \rangle_{\mathbb{R}^n}$ .

Now if  $Y_1, \ldots, Y_d$  are vector fields such that  $(Y_1(z), \ldots, Y_d(z))$  is an orthonormal basis of  $T_z\Gamma$  we obtain

$$
\operatorname{div} \xi^{T}(z) = \operatorname{Tr} \left( Y \mapsto \nabla_{Y} \xi^{T}(z) \right) = \sum_{j=1}^{d} \left\langle \xi_{z}^{N}, \operatorname{II}_{z}(Y_{j}(z), Y_{j}(z)) \right\rangle_{\mathbb{R}^{n}},
$$

and hence

$$
\left|\mathrm{div}\,\xi^T(z)\right|\leqslant d\left|\xi^N_z\right|\left|\left|\mathrm{II}_z\right|\right|.
$$

We can apply Lemma [2.1](#page-5-4) with  $B = A \langle ||\Pi_z|| \rangle$  and  $f_h^z(x) = \chi_1(hx) (\text{Op}_h^w(q_{1,z})S)(x)$ . Indeed, since q is assumed to be in  $\mathcal{S}(\langle \xi \rangle^{-1})$ ,  $q_{1,z}$  is in  $C_0^{\infty}(\mathbb{R}^{2n})$  uniformly in  $z \in \Gamma$ , and hence for all  $k \in \mathbb{N}$ the operator  $\mathrm{Op}_{h}^{w}(q_{1,z})$  belongs to  $\mathcal{L}(L^{2,k}(\mathbb{R}^{n}))$  and the norm is uniform in z (see [\[Wan88\]](#page-15-13)). This proves that

$$
||I_{1,1}(h)||_{L^{2,\delta}(\mathbb{R}^n)} = \bigcirc_{h \to 0} (\sqrt{h}).
$$

• We have

$$
\xi_z^T \cdot (A(z)\chi_1(x-z)) = \chi_1(x-z)d_z A(\xi_z^T) + A(z)\,\xi_z^T \cdot \chi_1(x-z).
$$

We set

$$
\begin{cases} B_{2,1}(z) = \|d_z A\| \\ q_{z,2,1}(x,\xi) = \chi_1(hx) d_z A(\xi_z^T) \|d_z A\|^{-1} q(x,\xi) \end{cases} \quad \text{and} \quad \begin{cases} B_{2,2}(z) = A(z) \\ q_{z,2,2}(x,\xi) = q(x,\xi) d(\chi_1)_{hx}(\xi_z^T) \end{cases}
$$

 $(d_z A(\xi_z^T) || d_z A||^{-1}$  can be replaced by  $|\xi|$  when  $|| d_z A || = 0$ ). As above, applying Lemma [2.1](#page-5-4) with  $B_{2,j}$  and  $f_{h,2,j}^z = \text{Op}_{h}^w(q_{z,2,j})S$  for  $j \in \{1,2\}$  gives then

$$
||I_{1,2}(h)||_{L^{2,\delta}(\mathbb{R}^n)} = \bigcirc_{h \to 0} (\sqrt{h}).
$$

We now deal with  $I_{1,3}(h)$ . For any vector field Y on  $\Gamma$  we have

$$
Y \cdot \langle \xi^T, \xi^T \rangle_{\text{TT}} = Y \cdot \langle \xi, \xi^T \rangle_{\mathbb{R}^n}
$$
  
=  $\langle \xi, \nabla_Y \xi^T + \text{II}(Y, \xi^T) \rangle = \langle \xi^T, \nabla_Y \xi^T \rangle_{\text{TT}} + \langle \xi^N, \text{II}(Y, \xi^T) \rangle_{\mathbb{R}^n}$   
=  $\frac{1}{2} Y \cdot \langle \xi^T, \xi^T \rangle_{\text{TT}} + \langle \xi^N, \text{II}(Y, \xi^T) \rangle_{\mathbb{R}^n}$ ,

and in particular:

$$
\xi^T \cdot \langle \xi^T, \xi^T \rangle_{T\Gamma} = 2 \langle \xi^N, \Pi \left( \xi^T, \xi^T \right) \rangle_{\mathbb{R}^n}.
$$

Thus we can estimate  $I_{1,3}(h)$  as  $I_{1,1}(h)$ , and this concludes the proof.

We now check that according to [\(1.6\)](#page-2-1) the assumptions of Proposition [2.3](#page-6-1) are satisfied for a symbol q supported in an incoming region:

<span id="page-8-0"></span>**Proposition 2.4.** *Let*  $\sigma_2 \in ]\sigma_1, 1]$ *,*  $R_2 > R_1$  *and*  $\nu_0 > 0$ *. Then there exists*  $\varepsilon > 0$  *such that for all*  $(x, \xi) \in \mathcal{Z}_-(R_2, \nu_0, -\sigma_2)$  *and*  $z \in \Gamma$  *we have* 

$$
|x - z| \geqslant \varepsilon \quad or \quad \left| \xi_z^T \right| \geqslant \varepsilon.
$$

*Proof.* Let  $(x, \xi) \in \mathcal{Z}_-(R_2, \nu_0, -\sigma_2), z \in \Gamma$ , and assume that

$$
|x-z| \leqslant \min\left(R_2 - R_1, \frac{R_1(\sigma_2 - \sigma_1)}{4}\right).
$$

In particular  $|z| \ge R_1$  and hence, according to [\(1.6\)](#page-2-1), we have

$$
|z| \left| \xi_z^T \right| \geqslant -\left\langle z, \xi_z^T \right\rangle = \left\langle z, \xi_z^N - \xi \right\rangle = \left\langle z, \xi_z^N \right\rangle - \left\langle x, \xi \right\rangle + \left\langle x - z, \xi \right\rangle
$$
  
\n
$$
\geqslant -\sigma_1 |z| \left| \xi_z^N \right| + \sigma_2 |x| \left| \xi \right| - \left| x - z \right| \left| \xi \right| \geqslant \left( -\sigma_1 + \sigma_2 \frac{|x|}{|z|} - \frac{|x - z|}{|z|} \right) |z| \left| \xi \right|
$$
  
\n
$$
\geqslant \left( \sigma_2 - \sigma_1 - (1 + \sigma_2) \frac{|x - z|}{R_1} \right) |z| \nu_0
$$
  
\n
$$
\geqslant |z| \frac{\nu_0 (\sigma_2 - \sigma_1)}{2}.
$$

Now we can estimate the solution  $u<sub>h</sub>$  in an incoming region. We recall the following result:

<span id="page-9-0"></span>**Proposition 2.5.** Let  $\tilde{R} > 0$ ,  $0 \leq \tilde{\nu} < \nu$  and  $-1 < \tilde{\sigma} < \sigma < 1$ . Then there exists  $R > \tilde{R}$  $such that for \omega \in S_0(\mathbb{R}^{2n})$  supported outside  $\mathcal{Z}_-(\tilde{R}, \tilde{\nu}, -\tilde{\sigma})$  and  $\omega_- \in S_0(\mathbb{R}^{2n})$  supported in  $\mathcal{Z}_-(R,\nu,-\sigma)$ *, we have* 

$$
\sup_{z \in \mathbb{C}_{J,+}} \left\| \langle x \rangle^{-\delta} \operatorname{Op}_h(\omega_-) (H_h - z)^{-1} \operatorname{Op}_h(\omega) \langle x \rangle^{-\delta} \right\|_{\mathcal{L}(L^2(\mathbb{R}^n))} = \bigcup_{h \to 0} (h^{\infty}).
$$

*Moreover the estimate remains true for the limit*  $(H_h - (\lambda + i0))^{-1}$ ,  $\lambda \in J$ *.* 

This theorem is proved in [\[RT89\]](#page-15-14) for the self-adjoint case and extended in [\[Roy10,](#page-15-6) [Roy11\]](#page-15-7) for our non-selfadjoint setting. Before giving an estimate of  $u<sub>h</sub>$  in an incoming region, we recall that it concentrates on the hypersurface of energy  $E_0$ . For  $h \in ]0, h_0], t \geq 0$  and  $z \in \mathbb{C}$  we set

$$
U_h(t,z) = e^{-\frac{it}{h}(H_h-z)}.
$$

<span id="page-9-1"></span>**Proposition 2.6.** Let  $q \in C_b^{\infty}(\mathbb{R}^{2n})$  be a symbol which vanishes on  $p^{-1}(I)$  for some neighborhood  $I \subset J$  *of*  $E_0$ *. Let*  $T \geq 0$ *. Then there exists*  $C \geq 0$  *such that for*  $h \in ]0,1]$ *,*  $z \in \mathbb{C}_{I,+}$  *and*  $\chi \in C_0^{\infty}(\mathbb{R}, [0, 1])$  *non-increasing and supported in*  $[0, T + 1]$  *we have* 

$$
\left\| \frac{i}{h} \int_0^\infty \chi(t) \mathcal{O} \mathcal{P}_h^w(q) U_h(t, z) dt \right\|_{\mathcal{L}(L^2(\mathbb{R}^n))} \leqslant C
$$

*and*

$$
\left\|{\rm Op}_h^w(q)(H_h-z)^{-1}\right\|_{\mathcal{L}(L^{2,\delta}(\mathbb{R}^n),L^{2,-\delta}(\mathbb{R}^n))}\leqslant C.
$$

*In particular if*  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  *is supported outside*  $p^{-1}(\lbrace E_0 \rbrace)$  *we have* 

$$
\langle \operatorname{Op}_h^w(q)u_h, u_h \rangle \xrightarrow[h \to 0]{} 0.
$$

This is Proposition 2.11 in [\[Roy10\]](#page-15-6). Note that the same holds if  $\text{Op}_{h}^{w}(q)$  is on the right of the propagator or the resolvent.

<span id="page-9-2"></span>**Proposition 2.7.** Let  $\sigma \in ]\sigma_1, 1[$ . Then there exists  $R \geq 0$  such that for  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  supported *in*  $\mathcal{Z}_-(R, 0, -\sigma)$  *we have* 

$$
\langle \operatorname{Op}_h^w(q)u_h, u_h \rangle \xrightarrow[h \to 0]{} 0.
$$

*Proof.* Let  $\sigma_2, \sigma_3$  be such that  $\sigma_1 < \sigma_2 < \sigma_3 < \sigma$ ,  $R_2 > R_1$  and  $\nu_0 \in ]0, (\text{inf } J)/3[$ . Let  $\omega_- \in$  $\mathcal{S}_0(\mathbb{R}^{2n})$  be supported in  $\mathcal{Z}_-(R_2,\nu_0,-\sigma_2)$  and equal to 1 on  $\mathcal{Z}_-(2R_2,2\nu_0,-\sigma_3)$ . Let  $R > 2R_2$  be chosen large enough and consider  $q_-, \tilde{q}_-\in C_0^{\infty}(\mathbb{R}^{2n})$  supported in  $\mathcal{Z}_-(R, 0, -\sigma)\cap p^{-1}(J)$  and such that  $\tilde{q}_- = 1$  on a neighborhood of supp  $q_-\$ . If R is large enough we have  $\mathcal{Z}_-(R, 0, -\sigma) \cap p^{-1}(J) \subset$  $\mathcal{Z}_-(R, 3\nu_0, -\sigma)$ , so according to Propositions [2.3](#page-6-1) and [2.5](#page-9-0) we have

$$
\|Op_h^w(q_-)(H_h - (E_h + i0))^{-1}S_h\|_{L^2(\mathbb{R}^n)}
$$
  
\$\leq \left\|Op\_h^w(q\_-)(H\_h - (E\_h + i0))^{-1}\right\|\_{\mathcal{L}(L^{2,\delta}(\mathbb{R}^n), L^2(\mathbb{R}^n))}\left\|Op\_h(\omega\_-)S\_h\right\|\_{L^{2,\delta}(\mathbb{R}^n)}  
+ \left\|Op\_h^w(q\_-)(H\_h - (E\_h + i0))^{-1}(1 - Op\_h(\omega\_-)) \langle x \rangle^{-\delta}\right\|\_{\mathcal{L}(L^2(\mathbb{R}^n))}\left\|S\_h\right\|\_{L^{2,\delta}(\mathbb{R}^n)}  
= O(\sqrt{h}).

The same applies to  $\tilde{q}_-$  and this finally gives

$$
|\langle \operatorname{Op}_h^w(q_-)u_h, u_h \rangle| \leq \| \operatorname{Op}_h^w(q_-)u_h \| \| \operatorname{Op}_h^w(\tilde{q}_-)u_h \| + \underset{h \to 0}{\vphantom{\big|}} C(h^{\infty}) \xrightarrow[h \to 0]{} 0.
$$

#### <span id="page-10-1"></span>10 JULIEN ROYER

3. Control of large times and of the source term far from the origin

As mentionned in introduction, we expect that the semiclassical measure of  $(u_h)_{h\in[0,h_0]}$  on some bounded subset of  $\mathbb{R}^{2n}$  does not depend on the values of the amplitude  $A(z)$  for large  $|z|$ . When restricting our attention to finite times, this is a consequence of Egorov's Theorem and the following proposition, proved in [\[Roy11\]](#page-15-7) (Proposition 2.1):

<span id="page-10-0"></span>**Proposition 3.1.** Let  $E_2 \geqslant E_1 > 0$ ,  $J \subset [E_1, E_2]$  and  $\sigma_3 \in [0, 1]$  such that  $\sigma_3^2 E_2 < E_1$ . Then *there exist*  $\mathcal{R} > 0$  *and*  $c_0 > 0$  *such that* 

$$
\forall t \geq 0, \forall (x, \xi) \in \mathcal{Z}_{\pm}(\mathcal{R}, 0, \mp \sigma_3) \cap p^{-1}(J), \quad |X(\pm t, x, \xi)| \geq c_0(t + |x|).
$$

For  $r > 0$  we set

$$
B_x(r) = \{(x, \xi) : |x| < r\} \subset \mathbb{R}^{2n}
$$

.

With Proposition [3.1](#page-10-0) we can check that on a bounded subset of  $\mathbb{R}^{2n}$  we can ignore the contribution of the source far from the origin:

<span id="page-10-2"></span>**Proposition 3.2.** Let  $r > 0$ . There exists  $R_0 > 0$  such that for  $T \ge 0$ ,  $R \ge R_0$ , Im  $z \ge 0$  and  $q \in \widehat{C}_0^{\infty}(\mathbb{R}^{2n}, [0,1])$  *supported in*  $B_x(r)$  *we have* 

$$
\left\| \frac{i}{h} \int_0^\infty \chi_{T,R}(t) \mathcal{O}_{p_h}^w(q) U_h(t,z) (S_h - S_h^R) dt \right\|_{L^2(\mathbb{R}^n)} = \bigcup_{h \to 0} (\sqrt{h}),
$$

*where the size of the rest depends on* T *but not on* z, q or  $R \ge R_0$ .

We recall that  $\chi_{T,R} \in C^{\infty}(\mathbb{R}, [0,1])$  is non-increasing, equal to 1 on  $]-\infty, T]$  and equal to 0 on  $[T+1, +\infty]$ .

*Proof.* Let R and  $c_0$  be given by Proposition [3.1](#page-10-0) applied with  $\sigma_3 = (1 + \sigma_1)/2$  (which is allowed according to Assumption [\(1.7\)](#page-2-3)). Let  $\mathcal{R} > \max(\mathcal{R}, 2r/c_0, R_1)$  be so large that

$$
\nu_0^2 := \inf J - \sup_{|x| \ge \tilde{\mathcal{R}}} |V_1(x)| > 0.
$$

Let  $\sigma_2 \in ]\sigma_1, \sigma_3[$ . We consider  $\omega \in C_b^{\infty}(\mathbb{R}^{2n})$  supported in  $\mathcal{Z}_+(\tilde{\mathcal{R}}, 0, -\sigma_3)$  and equal to 1 in a neighborhood of  $\mathcal{Z}_+(2\tilde{\mathcal{R}},0,-\sigma_2)$ . Given  $\theta \in C_0^{\infty}(\mathbb{R})$  supported in J and equal to 1 on a neighborhood of  $E_0$  we prove that

$$
\sum_{j=1}^3 \left\| \frac{i}{h} \int_0^\infty \chi_{T,R}(t) \mathcal{O}_{P_h}^w(q) U_h(t,z) B_j(h) (S_h - S_h^R) dt \right\|_{L^2(\mathbb{R}^n)} = \bigcup_{h \to 0} \left( \sqrt{h} \right),
$$

where

$$
B_1(h) = \text{Op}_h^w(1 - \theta \circ p), \quad B_2(h) = \text{Op}_h^w(\theta \circ p) \text{Op}_h^w(\omega), \quad B_3(h) = \text{Op}_h^w(\theta \circ p) \text{Op}_h^w(1 - \omega).
$$

The first term is estimated with Propositions [2.6](#page-9-1) and [2.2.](#page-6-0) The function  $\chi_{T,R}$  depends on R, but since it is non-increasing and always vanishes on  $[T+1, +\infty]$ , we can check that the estimate is actually uniform in  $R > 0$ . According to Proposition [3.1](#page-10-0) and Egorov's Theorem (see Theorem 7.2 in [\[Roy11\]](#page-15-7)), the second term is of size  $O(h^{\infty})$ , again uniformly in  $R \ge 0$ . For the third term we set  $R_0 = 3\tilde{\mathcal{R}}$  and choose  $R \ge R_0$ . The symbol  $(\theta \circ p)(1 - \omega)$  is supported in

$$
B_x(2\tilde{\mathcal{R}}) \cup \mathcal{Z}_-(2\tilde{R},0,-\sigma_2) \cap p^{-1}(J) \subset B_x(2\tilde{\mathcal{R}}) \cup \mathcal{Z}_-(2\tilde{R},\nu_0,-\sigma_2).
$$

Since  $A - A_R$  is supported outside  $B(0, 3\tilde{R})$  and according to Propositions [2.3](#page-6-1) and [2.4](#page-8-0) (applied with  $A - A_R$  instead of A) we obtain that  $\left\| \text{Op}_h(1 - \omega)(S_h - S_h^R) \right\|_{L^2(\mathbb{R}^n)} = O(h^{\frac{3}{2}})$  uniformly in  $R.$ 

We know that [\(1.15\)](#page-4-2) holds for some measure  $\mu_T^R$  when  $u_h$  is replaced by  $u_h^{T,R}$  (see Theorem 4.3 in [\[Roy10\]](#page-15-6)). We now check that  $u_h^{T,R}$  is actually a good approximation of  $u_h$  (in some sense) when  $h > 0$  is small and T, R are large enough:

<span id="page-10-3"></span>**Proposition 3.3.** Let  $r > 0$  and  $R_0 > 0$  given by Proposition [3.2.](#page-10-2) Let K be a compact subset *of*  $B_x(r) \cap p^{-1}(J)$  *and*  $\varepsilon > 0$ *. Then there exists*  $T_0 \geq 0$  *such that for*  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  *supported in*  $K, T \geq T_0$  *and*  $R \geq R_0$  *we have* 

$$
\limsup_{h\to 0} \left| \langle \operatorname{Op}_h^w(q) u_h, u_h \rangle - \langle \operatorname{Op}_h^w(q) u_h^{T,R}, u_h^{T,R} \rangle \right| \leq \varepsilon ||q||_{\infty}.
$$

This proposition relies on the following consequence of the non-selfadjoint version of Egorov's Theorem (see [\[Roy11,](#page-15-7) Prop.7.3]):

<span id="page-11-1"></span>**Proposition 3.4.** Let *J* be a neighborhood of E such that Assumption [\(1.5\)](#page-2-0) holds for all  $\lambda \in J$ . Let  $K_1$  and  $K_2$  be compact subsets of  $p^{-1}(J)$ *.* Let  $\varepsilon > 0$ *. Then there exists*  $T_0 \geq 0$  *such that for*  $T \geq T_0$  and  $q_1, q_2 \in C_0^{\infty}(\mathbb{R}^{2n})$  respectively supported in  $K_1$  and  $K_2$  we have

$$
\limsup_{h\to 0} \|\text{Op}_{h}^{w}(q_1)U_{h}(T)\text{Op}_{h}^{w}(q_2)\|_{\mathcal{L}(L^{2}(\mathbb{R}^n))} \leq \varepsilon \|q_1\|_{\infty} \|q_2\|_{\infty}.
$$

We also need the following result about the classical flow:

<span id="page-11-0"></span>**Proposition 3.5.** Let  $E_1, E_2 \in \mathbb{R}_+^*$  be such that  $E_1 \leqslant E_2$ , and  $\sigma \in [0,1]$ . If R is chosen large *enough then for any compact subset* K of  $p^{-1}([E_1, E_2])$  *there exists*  $T_0 \geq 0$  *such that* 

 $\forall w \in K, \forall t \geq T_0, \quad \phi^{\pm t}(w) \in B_x(\mathcal{R}) \cup \mathcal{Z}_{\pm}(\mathcal{R}, 0, \pm \sigma).$ 

This is slightly more general than Lemma 5.2 in [\[Roy10\]](#page-15-6). We recall the idea of the proof:

*Proof.* We consider  $\mathcal{R}_0$  such that

$$
\forall x \in \mathbb{R}^n, \quad |x| \geqslant \mathcal{R}_0 \quad \Longrightarrow \quad |V_1(x)| + |x| \, |\nabla V_1(x)| \leqslant \frac{E_1}{3} (1 - \sigma^2).
$$

Let  $\tau > 0$  be such that  $\int_0^{\tau} \frac{(1-\sigma^2)\sqrt{3(\mathcal{R}_0+4s)}}{\sqrt{3(\mathcal{R}_0+4s)}}$  $\frac{(1-\sigma^2)\sqrt{E_1}}{\sqrt{3}(\mathcal{D}_1+\Lambda_2\sqrt{E_2})}$  $\frac{(1-\sigma^2)\sqrt{E_1}}{3(\mathcal{R}_0+4s\sqrt{E_2})}$   $ds > \sigma$ . We set  $\mathcal{R} = \mathcal{R}_0 + 4\tau\sqrt{E_2}$  and  $\mathcal{U}_{\pm} = B_x(\mathcal{R}_0) \cup$  $\mathcal{Z}_{\pm}(\mathcal{R}, 0, \pm \sigma)$ . We first prove that for any  $w \in K$ , if  $\phi^{\pm t}(w) \in \mathcal{U}_{\pm}$  for some  $t_w \geq 0$ , then  $\phi^{\pm t}(w) \in B_x(\mathcal{R}) \cup \mathcal{Z}_{\pm}(\mathcal{R}, 0, \pm \sigma)$  for all  $t \geq t_w$ . Since  $\phi^{\pm (t-t_w)}$  maps  $\mathcal{Z}_{\pm}(\mathcal{R}, 0, \pm \sigma)$  into itself for all  $t \geq t_w$ , we can assume that  $|X(\pm t_w, w)| = \mathcal{R}_0$ ,  $|X(\pm t, w)| = \mathcal{R}$  and  $|X(\pm s, w)| \in [\mathcal{R}, \mathcal{R}_0]$  for  $s \in [t_w, t]$ . Assume by contradiction that  $\phi^{\pm s}(w) \notin \mathcal{Z}_{\pm}(\mathcal{R}, 0, \pm \sigma)$  when  $s \in [t_w, t]$ . Then we can check that

$$
\pm \frac{\partial}{\partial s} \frac{X(\pm s, w) \cdot \Xi(\pm s, w)}{|X(\pm s, w)| |\Xi(\pm s, w)|} \geq \frac{(1 - \sigma^2)\sqrt{E_1}}{\sqrt{3}(\mathcal{R}_0 + 4(s - t_0)\sqrt{E_2})},
$$

which gives a contradiction. Then it only remains to check that if  $T_0$  is chosen large enough, then for all  $w \in K$  we can find  $t_w \in [0, T_0]$  such that  $\phi^{\pm t_w}(w) \in \mathcal{U}_{\pm}$ . For this we use compactness of K and the fact that any trajectory has a limit point in  $\Omega_b([E_1, E_2]) \subset B_x(\mathcal{R}_0)$  or goes to infinity and meets  $\mathcal{Z}_b(\mathcal{R}_0) \to \pi$  when t is large enough infinity and meets  $\mathcal{Z}_{\pm}(\mathcal{R}, 0, \pm \sigma)$  when t is large enough.

Let  $\sigma_2 < \sigma_3 < \sigma_4 < \sigma_5 \in ]\sigma_1,1[$  and  $\nu_0 \in ]0,\sqrt{\inf J}/4[$ . Let R be given by Proposition [2.5](#page-9-0) applied with  $(R, \tilde{\nu}, \tilde{\sigma}) = (3R_1, 2\nu_0, \sigma_3)$  and  $(\nu, \sigma) = (3\nu_0, \sigma_4)$ . Choosing R larger if necessary, we can assume that  $|\xi| \geq 4\nu_0$  if  $p(x,\xi) \in J$  and  $|x| \geq \mathcal{R}$ . We can also assume that  $2\mathcal{R}$  satisfies the conclusion of Proposition [3.5](#page-11-0) applied with  $[E_1, E_2] \supset J$ .

**Lemma 3.6.** Let  $r \geq 4R$  and  $R_0 > 0$  given by Proposition [3.2.](#page-10-2) Let  $Q \in C_0^{\infty}(\mathbb{R}^{2n}, [0,1])$  be *supported in*  $B_x(r) \cap p^{-1}(J)$  *and equal to 1 in a neighborhood of*  $B_x(3\mathcal{R}) \cap p^{-1}(\overline{I})$  *for some open neighborhood* I of  $E_0$ . Let K be a compact subset of  $B_x(r) \cap p^{-1}(J)$  and  $\delta > 0$ . Then there exists  $T_0 \geq 0$  such that for  $T \geq T_0$ ,  $R \geq R_0$  and  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  supported in K we have

$$
\limsup_{h\to 0} \left\| \operatorname{Op}_h^w(q) \left( u_h - u_h^{T,R} - A_T^{\delta}(h) \operatorname{Op}_h^w(Q) u_h \right) \right\|_{L^2(\mathbb{R}^n)} \leq \delta \|q\|_{\infty},
$$

where  $A_T^{\delta}(h)$  *is a bounded operator such that* 

$$
\forall T \geqslant T_0, \quad \limsup_{h \to 0} \left\| A_T^{\delta}(h) \right\|_{\mathcal{L}(L^2(\mathbb{R}^n))} \leqslant \delta.
$$

For the proof we follows the same general idea as in [\[Roy10\]](#page-15-6):

*Proof.* • We consider  $\tilde{q} \in C_0^{\infty}(\mathbb{R}^{2n})$  supported in  $B_x(r) \cap p^{-1}(J)$  and equal to 1 on a neighborhood of K. Let  $\theta \in C_0^{\infty}(\mathbb{R}^n)$  be supported in  $B(0, 3\mathcal{R})$  and equal to 1 on  $B(0, 2\mathcal{R})$ . Let  $\omega_-\in\mathcal{S}_0(\mathbb{R}^{2n})$  be equal to 1 on  $\mathcal{Z}_-(2\mathcal{R}, 4\nu_0, -\sigma_5)$  and supported in  $\mathcal{Z}_-(\mathcal{R}, 3\nu_0, -\sigma_4)$ . Let  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  be supported in  $K, h \in ]0, h_0]$ ,  $T \geq 0$  and  $z \in \mathbb{C}_{I,+}$   $(h_0 > 0$  was fixed small enough in the introduction). Since z is not in the spectrum of  $H_h$  we can consider  $(H_h - z)^{-1}S_h \in H^2(\mathbb{R}^n)$  and write:

<span id="page-12-0"></span>
$$
Op_h^w(q)(H_h - z)^{-1}S_h - Op_h^w(q)U_h(T, z)(H_h - z)^{-1}S_h
$$
\n
$$
= \frac{i}{h} \int_0^\infty Op_h^w(q)\big(\chi_{T,R}(t) - \chi_{0,R}(t)U_h(T, z)\big)U_h(t, z)S_h dt
$$
\n
$$
= \frac{i}{h} \int_0^\infty Op_h^w(q)\big(\chi_{T,R}(t) - \chi_{0,R}(t)U_h(T, z)\big)U_h(t, z)S_h^R dt + O_h(\sqrt{h}),
$$
\n(3.1)

where the rest is estimated in  $L^2(\mathbb{R}^n)$  uniformly in  $R \ge R_0$  but not in T (see Proposition [3.2\)](#page-10-2).

• We have

$$
Op_h^w(q)U_h(T, z)(H_h - z)^{-1}S_h = Op_h^w(q)U_h(T, z)Op_h^w(Q)(H_h - z)^{-1}S_h
$$
  
+
$$
Op_h^w(q)U_h(T, z)Op_h^w(1 - Q)\theta(x)(H_h - z)^{-1}S_h
$$
  
+
$$
Op_h^w(q)U_h(T, z)Op_h^w(1 - Q)(1 - \theta(x))Op_h(\omega_{-})(H_h - z)^{-1}S_h
$$
  
+
$$
Op_h^w(q)U_h(T, z)Op_h^w(1 - Q)(1 - \theta(x))Op_h(1 - \omega_{-})(H_h - z)^{-1}S_h.
$$

The second term of the right-hand side is of size  $O(\sqrt{h})$  uniformly in  $z \in \mathbb{C}_{I,+}$  according to Proposition [2.6.](#page-9-1) Let  $\omega \in \mathcal{S}_0(\mathbb{R}^{2n})$  be supported in  $\mathcal{Z}_-(2R_1,\nu_0,-\sigma_2)$  and equal to 1 in  $\mathcal{Z}_-(3R_1, 2\nu_0, -\sigma_3)$ . According to Propositions [2.3,](#page-6-1) [2.4](#page-8-0) and [2.5](#page-9-0) we have

$$
\|Op_h^w(q)U_h(T,z)Op_h^w(1-Q)(1-\theta(x))Op_h(\omega_-)(H_h-z)^{-1}S_h\|
$$
  
\$\leq c\_{q,T} ||Op\_h(\omega\_-)(H\_h-z)^{-1}(1-Op\_h(\omega))S\_h||\_{L^{2,-\delta}(\mathbb{R}^n)} + \underset{h\to 0}{O}(\sqrt{h})\$  
= 
$$
\underset{h\to 0}{O}(\sqrt{h}),
$$

uniformly in  $z \in \mathbb{C}_{J,+} (\langle \mathcal{O}_{P_h}^w(q)U_h(T,z)\langle x\rangle^{\delta}) = O(1)$  uniformly in  $z \in \mathbb{C}_{J,+}$  but not in T. Finally the last term is of size  $O(h^{\infty})$  according to Egorov's Theorem and Proposition [3.5](#page-11-0) if  $T \geq T_0$ ,  $T_0$  being given by Proposition [3.5](#page-11-0) applied with  $\sigma_5$ . We consider  $\tilde{q}, \tilde{Q} \in C_0^{\infty}(\mathbb{R}^{2n}, [0,1])$ supported in  $p^{-1}(J)$  and equal to 1 respectively on a neighborhood of supp q and supp Q, and we set

$$
A_T^{\delta}(z, h) = \text{Op}_h^w(\tilde{q}) U_h(T, z) \text{Op}_h^w(\tilde{Q}).
$$

According to Proposition [3.4](#page-11-1) we have

$$
\limsup_{h \to 0} \sup_{z \in \mathbb{C}_{I,+}} ||A_T^{\delta}(z,h)||_{\mathcal{L}(L^2(\mathbb{R}^n))} \le \delta
$$

when  $T \geq T_0$ , if  $T_0$  was chosen large enough. We finally obtain

$$
\mathrm{Op}_{h}^{w}(q)U_{h}(T,z)(H_{h}-z)^{-1}S_{h}=\mathrm{Op}_{h}^{w}(q)A_{T}^{\delta}(z,h)\mathrm{Op}_{h}^{w}(Q)(H_{h}-z)^{-1}S_{h}+\underset{h\to 0}{O}(\sqrt{h})
$$

in  $L^2(\mathbb{R}^n)$  where the size of the rest is uniform in  $z \in \mathbb{C}_{I,+}$ .

**3.** For  $h \in ]0, h_0]$  and  $T \geq T_0$ , we can take the limit  $z \to E_h$  in [\(3.1\)](#page-12-0)  $(E_h \in \overline{\mathbb{C}_{I,+}}]$  if  $h_0$  is small enough). This gives in  $L^2(\mathbb{R}^n)$ :

$$
Op_h^w(q)u_h = Op_h^w(q)u_h^{T,R} - Op_h^w(q)U_h^E(T)u_h^{0,R} + Op_h^w(q)A_T^{\delta}(E_h, h)Op_h^w(Q)u_h + O_h(\sqrt{h}).
$$

Let  $q_1 \in C_0^{\infty}(\mathbb{R}^{2n}, [0,1])$  be equal to 1 on a neighborhood of  $\left\{ \phi^t(z,\xi), t \in [0,1], (z,\xi) \in N_E^{R_0} \Gamma \right\}$ . Using the results about the contribution of small times (see in particular Corollary 4.4 in  $[Roy10]$ ), we know that

$$
\limsup_{h \to 0} \left\| \text{Op}_{h}^{w}(q) U_{h}^{E}(T) u_{h}^{0,R_0} \right\|_{L^{2}(\mathbb{R}^n)} = \limsup_{h \to 0} \left\| \text{Op}_{h}^{w}(q) U_{h}^{E}(T) \text{Op}_{h}^{w}(q_{1}^{2}) u_{h}^{0,R_0} \right\|_{L^{2}(\mathbb{R}^n)} \leq C \limsup_{h \to 0} \left\| \text{Op}_{h}^{w}(q_{1}) U_{h}^{E}(T) \text{Op}_{h}^{w}(q_{1}) \right\|_{\mathcal{L}(L^{2}(\mathbb{R}^n))},
$$

for some constant C which does not depend on  $h, T$  or  $R$ . According to Proposition [3.4,](#page-11-1) this limit is less than  $\delta ||q||_{\infty}$  for any  $T \ge T_0$  if  $T_0$  was chosen large enough. Since for any  $T \ge T_0$  we have

$$
\phi^{-T}(\operatorname{supp} q) \cap \left\{ \phi^t(z,\xi), t \in [0,1], (z,\xi) \in N_E^R \Gamma \setminus N_E^{R_0} \Gamma \right\} = \emptyset
$$

Egorov's Theorem also gives

$$
\limsup_{h \to 0} \left\| \text{Op}_{h}^{w}(q) U_{h}^{E}(T) (u_{h}^{0,R} - u_{h}^{0,R_0}) \right\|_{L^{2}(\mathbb{R}^{n})} = 0,
$$

which concludes the proof.  $\Box$ 

Then Proposition [3.3](#page-10-3) is proved exactly as in [\[Roy10\]](#page-15-6) (see Proposition 5.5), and we can show existence of a semiclassical measure:

**Proposition 3.7.** *There exists a non-negative Radon measure*  $\mu$  *on*  $\mathbb{R}^{2n}$  *such that for*  $q \in \mathbb{R}$  $C_0^{\infty}(\mathbb{R}^{2n})$  *we have* 

$$
\int_{\mathbb{R}^{2n}} q \, d\mu = \lim_{T, R \to +\infty} \int_{\mathbb{R}^{2n}} q \, d\mu_T^R \tag{3.2}
$$

*and*

$$
\langle \operatorname{Op}_h^w(q)u_h, u_h \rangle \xrightarrow[h \to 0]{} \int_{\mathbb{R}^{2n}} q \, d\mu.
$$

*Proof.* The result is clear outside  $p^{-1}(\lbrace E_0 \rbrace)$ , so we focus on symbols supported in  $p^{-1}(J)$ . Let K be a compact subset of  $p^{-1}(J)$  and  $\varepsilon > 0$ . Let  $T_0$  and  $R_0$  given by Proposition [3.3.](#page-10-3) For  $T_1, T_2 \geq T_0, R_1, R_2 \geq R_0$  and  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  supported in K we have

$$
\left| \int_{\mathbb{R}^{2n}} q \, d\mu_{T_1}^{R_1} - \int_{\mathbb{R}^{2n}} q \, d\mu_{T_2}^{R_2} \right|
$$
  
=  $\lim_{h \to 0} \left| \left\langle \mathrm{Op}_h^w(q) u_h^{T_1, R_1}, u_h^{T_1, R_1} \right\rangle - \left\langle \mathrm{Op}_h^w(q) u_h^{T_2, R_2}, u_h^{T_2, R_2} \right\rangle \right| \leq 2\varepsilon \|q\|_{\infty}.$ 

This proves that  $(T, R) \mapsto \int q d\mu_T^R$  has a limit at infinity, which we denote by  $L(q)$ . The map  $q \mapsto \int q \, d\mu_T^R$  is a nonnegative linear form on  $C_0^{\infty}(\mathbb{R}^{2n})$  for all  $T, R > 0$ , and hence so is  $q \mapsto L(q)$ . Let  $T_0$  be as above for  $\varepsilon = 1$  and  $C_K$  be a constant such that for all  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  we have

$$
\left| \int_{\mathbb{R}^{2n}} q \, d\mu_{T_0}^{R_0} \right| \leqslant C_K \left\| q \right\|_{\infty}.
$$

Then we have

$$
|L(q)| \leq |L(q) - \int_{\mathbb{R}^{2n}} q \, d\mu_{T_0}^{R_0} + \left| \int_{\mathbb{R}^{2n}} q \, d\mu_{T_0}^{R_0} \right|
$$
  
\n
$$
\leq \lim_{T, R \to +\infty} \left| \int_{\mathbb{R}^{2n}} q \, d\mu_{T}^{R} - \int_{\mathbb{R}^{2n}} q \, d\mu_{T_0}^{R_0} \right| + C_K ||q||_{\infty}
$$
  
\n
$$
\leq (2 + C_K) ||q||_{\infty},
$$

which proves that the linear form L on  $C_0^{\infty}(\mathbb{R}^{2n})$  can be extended as a continuous linear form on the space of continuous and compactly supported functions on  $\mathbb{R}^{2n}$ . The first assertion is now a consequence of Riesz' Lemma. And the second can be proved as in [\[Roy10\]](#page-15-6), using the fact that  $\langle \text{Op}_h^w(q)u_h, u_h \rangle$  is close to  $\langle \text{Op}_h^w(q)u_h^{T,R}, u_h^{T,R} \rangle$  in the sense of Proposition [3.3,](#page-10-3) and  $\left\langle \text{Op}_{h}^{w}(q)u_{h}^{T,R}, u_{h}^{T,R}\right\rangle$  goes to  $\int q d\mu_{T}^{R}$  as h goes to 0.

It is now easy to prove the remark about the measures  $\tilde{\mu}_R$  mentionned in introduction:

<span id="page-13-0"></span>**Proposition 3.8.** Let  $r > 0$ ,  $R_0$  given by Proposition [3.2](#page-10-2) and  $R \ge R_0$ . Then the measures  $\mu$ *and*  $\tilde{\mu}_R$  *coincide on*  $B_x(r)$ *.* 

To prove this assertion, we only have to apply Proposition [3.3](#page-10-3) with  $u_h$  and  $\tilde{u}_h^R$ . Let  $q \in$  $C_0^{\infty}(\mathbb{R}^{2n})$  be supported in  $B_x(r) \cap p^{-1}(J)$ . Since for large T and small h the quantity  $\left\langle \operatorname{Op}_h^w(q) u_h^{T,R}, u_h^{T,R} \right\rangle$ is a good approximation both for  $\langle \text{Op}_h^w(q)u_h, u_h \rangle$  and  $\langle \text{Op}_h^w(q)\tilde{u}_h^R, \tilde{u}_h^R \rangle$ , these two quantities have the same limit as  $h$  goes to 0.

It only remains to prove the properties given in Theorem [1.1:](#page-4-0)

*Proof.* The first two properties are direct consequences of Propositions [2.6](#page-9-1) and [2.7.](#page-9-2) Let  $r > 0$ and  $R_0$  given by Proposition [3.2.](#page-10-2) Property (c) is already known for the compactly supported amplitude  $A_{R_0}$ , and hence for any  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  we have

$$
\int_{\mathbb{R}^{2n}} (-H_p + 2 \operatorname{Im} \tilde{E} + 2V_2) q d\tilde{\mu}_{R_0} = \pi (2\pi)^{d-n} \int_{N_E \Gamma} q(z,\xi) |A_{R_0}(z)|^2 |\xi|^{-1} |\hat{S}(\xi)|^2 d\sigma_{N_E \Gamma}(z,\xi).
$$

Now assume that q, and hence  $(H_p + 2 \operatorname{Im} \tilde{E} + 2V_2)q$ , are supported in  $B_x(r)$ . According to Proposition [3.8](#page-13-0) we have

$$
\int_{\mathbb{R}^{2n}} (-H_p + 2 \operatorname{Im} \tilde{E} + 2V_2) q d\tilde{\mu}_{R_0} = \int_{\mathbb{R}^{2n}} (-H_p + 2 \operatorname{Im} \tilde{E} + 2V_2) q d\mu.
$$

On the other hand, according to Proposition [3.1](#page-10-0) trajectories comming from the points of  $N_E\Gamma$  $N_E^{R_0}$ T never reach  $B_x(r)$ , so we also have

$$
\pi (2\pi)^{d-n} \int_{N_{E}\Gamma} q(z,\xi) (|A(z)|^2 - |A_{R_0}(z)|^2) |\xi|^{-1} |\hat{S}(\xi)|^2 d\sigma_{N_{E}\Gamma}(z,\xi) = 0.
$$

This proves that Property (c) holds when  $q \in C_0^{\infty}(\mathbb{R}^{2n})$  is supported in  $B_x(r)$ . Since this holds for any  $r > 0$ , the theorem is proved.

#### Appendix A. Short review about differential geometry

<span id="page-14-0"></span>We briefly recall in this appendix the basic results of differential geometry we have used. Detailed expositions can be found for instance in [\[dC92\]](#page-15-15) and [\[Spi99\]](#page-15-16).

Let M be a differential manifold. We denote by  $\mathcal{X}(M)$  the set of vector fields on M. An affine connection on  $M$  is a mapping

$$
\nabla: \left\{ \begin{array}{ccc} \mathcal{X}(M) \times \mathcal{X}(M) & \to & \mathcal{X}(M) \\ (X,Y) & \mapsto & \nabla_X Y \end{array} \right.
$$

which satisfies the following properties (for  $X, Y, Z \in \mathcal{X}(M)$  and  $f, g \in C^{\infty}(M)$ ):

(i)  $\nabla_{fX+gY}Z = f\nabla_{X}Z + g\nabla_{Y}Z,$ (ii)  $\nabla_X (Y + Z) = \nabla_X Y + \nabla_X Z,$ (iii)  $\nabla_X(fY) = f\nabla_X Y + (X \cdot f)Y$ .

The Levi-Civita connection on M is the unique connection  $\nabla$  on M which is

(i) symmetric:

$$
\forall X, Y \in \mathcal{X}(M), \quad \nabla_X Y - \nabla_Y X = [X, Y] = XY - YX,
$$

(ii) and compatible with the Riemannian metric:

<span id="page-14-1"></span>
$$
\forall X, Y, Z \in \mathcal{X}(M), \quad X \cdot \langle Y, Z \rangle = \langle \nabla_X Y, Z \rangle_M + \langle Y, \nabla_X Z \rangle_M. \tag{A.1}
$$

The Levi-Civita connection on  $\mathbb{R}^n$  endowed with the canonical metric is the usual differential. Now let  $\Gamma$  be a submanifold of  $\mathbb{R}^n$ , endowed by the Riemannian structure defined by restriction of the scalar product of  $\mathbb{R}^n$ . For  $X, Y \in \mathcal{X}(\Gamma)$ ,  $z \in \Gamma$ , and  $\overline{X}, \overline{Y} \in \mathcal{X}(\mathbb{R}^n)$  such that  $X = \overline{X}$  and  $Y = \overline{Y}$  in a neighborhood of z in  $\Gamma$  we set

$$
\nabla^{\Gamma}_X Y(z) = \left(\nabla^{\mathbb{R}^n}_{\overline{X}} \overline{Y}(z)\right)^T_z.
$$

This definition does not depend on the choice of  $\overline{X}$  or  $\overline{Y}$  and defines the Levi-Civita connexion on Γ.

Let  $X \in \mathcal{X}(\Gamma)$ . The divergence div  $X(z)$  at point  $z \in \Gamma$  is defined as the trace of the linear map  $Y \mapsto \nabla_Y^{\Gamma} X(z)$  on  $T_z \Gamma$ . If  $X \in \mathcal{X}(\Gamma)$  and  $f \in C^{\infty}(\Gamma)$  then we have

$$
\operatorname{div}(fX) = X \cdot f + f \operatorname{div} X. \tag{A.2}
$$

The main theorem we have used in Section [2](#page-5-0) is the following:

<span id="page-15-12"></span>**Theorem A.1** (Green's Theorem). *If* M *is an oriented Riemannian manifold and*  $X \in \mathcal{X}(M)$ *is compactly supported, then*

$$
\int_M \operatorname{div} X \, dV_M = 0,
$$

*where*  $dV_m$  *denotes the volume element on*  $M$ *.* 

We finally recall the basic properties of the second fundamental form on Γ. Given  $X, Y \in \mathcal{X}(\Gamma)$ and  $z \in \Gamma$  we set

$$
\mathrm{II}_z(X,Y) = \nabla_{\overline{X}}^{\mathbb{R}_n} \overline{Y}(z) - \nabla_X^{\Gamma} Y(z) = \left(\nabla_{\overline{X}}^{\mathbb{R}_n} \overline{Y}(z)\right)_z^N \in N_z \Gamma,
$$

where  $\overline{X}$  and  $\overline{Y}$  are extensions of X and Y on a neighborhood of z in  $\mathbb{R}^n$ . We can check that  $\text{II}_{z}(X, Y)$  is well-defined and actually only depends on  $X(z)$  and  $Y(z)$ . Moreover the bilinear form  $II_z$  is symmetric.

Acknowledgement. This work is partially supported by the French National Research Project NOSEVOL, No. ANR 2011 BS01019 01, entitled *Nonselfadjoint operators, semiclassical analysis and evolution equations*.

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