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Deformations of third order Peregrine breather solutions of the NLS equation with four parameters.

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Abstract

In this paper, we give new solutions of the focusing NLS equation as a quotient of two determinants. This formulation gives in the case of the order 3, new deformations of the Peregrine breather with four parameters.

This gives a very efficient procedure to construct families of quasi-rational solutions of the NLS equation and to describe the apparition of multi rogue waves. With this method, we construct the analytical expressions of deformations of the Peregrine breather of order $N = 3$ depending on 4 real parameters and plot different types of rogue waves.

1 Introduction

The first results concerning the nonlinear Schrödinger equation (NLS) date from the Seventies. Precisely, in 1972 Zakharov and Shabat solved it using the the inverse scattering method [21, 22]. The first quasi-rational solutions of NLS equation were constructed in 1983 by Peregrine [20]. In 1986 Eleon-ski, Akhmediev and Kulagin obtained the two-phase almost periodic solution to the NLS equation and obtained the first higher order analogue of the Peregrine breather[3]. Other families of higher order were constructed in a series of articles by Akhmediev et al. [1, 2] using Darboux transformations. It has been shown in [8] in 2010, that rational solutions of the NLS equation

can be written as a quotient of two Wronskians. Recently, it has been constructed in [11] a new representation of the solutions of the NLS equation in terms of a ratio of two Wronskians determinants of even order $2N$ composed of elementary functions; the related solutions of NLS are called of order N . Quasi-rational solutions of the NLS equation were obtained by the passage to the limit when some parameter tended to 0.

These results can be compared with those obtain recently by Akhmediev et al. in [5] with Darboux dressing method and numerical approach.

Recently, an other representation of the solutions of the focusing NLS equation, as a ratio of two determinants has been given in [14] using generalized Darboux transform.

A new approach has been done in [19] which gives a determinant representation of solutions of the focusing NLS equation, obtained from Hirota bilinear method, derived by reduction of the Gram determinant representation for Davey-Stewartson system.

A little later, in 2012 one obtained a representation in terms of determinants which does not involve limits [13]. But, these first two formulations given in [11, 13] did depend in fact only on two parameters; this remark was first made by V.B. Matveev. Then, in the middle of the year 2012, one determined multi-parametric families of quasi rational solutions of NLS in terms of determinants of order N (determinants of order $2N$) dependent on $2N-2$ real parameters. They are similar to those previously explicitly found by V.B Matveev and P. Dubard by the method given in [2], for the first time in the case $N = 3$.

With this method, we obtain news deformations at order 3 with 4 real parameters. With this representation, one finds at the same time the circular forms well-known, but also the triangular forms put recently in obviousness by Ohta and Yang [19], and also by Akhmediev et al. [17].

The following orders will be the object of other publications.

2 Wronskian representation of solutions of NLS equation and quasi-rational limit

2.1 Solutions of NLS equation in terms of Wronskian determinant

We consider the focusing NLS equation

$$iv_t + v_{xx} + 2|v|^2v = 0. \quad (1)$$

From the works [11, 10], the solution of the NLS equation can be written in the form

$$v(x, t) = \frac{\det(I + A_3(x, t))}{\det(I + A_1(x, t))} \exp(2it - i\varphi). \quad (2)$$

In (2), the matrix $A_r = (a_{\nu\mu})_{1 \leq \nu, \mu \leq 2N}$ ($r = 3, 1$) is defined by

$$a_{\nu\mu} = (-1)^{\epsilon_\nu} \prod_{\lambda \neq \mu} \left| \frac{\gamma_\lambda + \gamma_\nu}{\gamma_\lambda - \gamma_\mu} \right| \exp(i\kappa_\nu x - 2\delta_\nu t + x_{r,\nu} + e_\nu). \quad (3)$$

$\kappa_\nu, \delta_\nu, \gamma_\nu$ are functions of the parameters $\lambda_\nu, \nu = 1, \dots, 2N$ satisfying the relations

$$0 < \lambda_j < 1, \quad \lambda_{N+j} = -\lambda_j, \quad 1 \leq j \leq N. \quad (4)$$

They are given by the following equations,

$$\kappa_\nu = 2\sqrt{1 - \lambda_\nu^2}, \quad \delta_\nu = \kappa_\nu \lambda_\nu, \quad \gamma_\nu = \sqrt{\frac{1 - \lambda_\nu}{1 + \lambda_\nu}}, \quad (5)$$

and

$$\kappa_{N+j} = \kappa_j, \quad \delta_{N+j} = -\delta_j, \quad \gamma_{N+j} = 1/\gamma_j, \quad j = 1 \dots N. \quad (6)$$

The terms $x_{r,\nu}$ ($r = 3, 1$) are defined by

$$x_{r,\nu} = (r - 1) \ln \frac{\gamma_\nu - i}{\gamma_\nu + i}, \quad 1 \leq j \leq 2N. \quad (7)$$

The parameters e_ν are defined by

$$e_j = ia_j - b_j, \quad e_{N+j} = ia_j + b_j, \quad 1 \leq j \leq N, \quad (8)$$

where a_j and b_j , for $1 \leq j \leq N$ are arbitrary real numbers. The terms ϵ_ν are defined by :

$$\begin{aligned}\epsilon_\nu &= 0, & 1 \leq \nu \leq N \\ \epsilon_\nu &= 1, & N + 1 \leq \nu \leq 2N.\end{aligned}$$

We use the following notations :

$$\Theta_{r,\nu} = \kappa_\nu x/2 + i\delta_\nu t - ix_{r,\nu}/2 + \gamma_\nu y - ie_\nu, \quad 1 \leq \nu \leq 2N.$$

We consider the functions

$$\begin{aligned}\phi_{r,\nu}(y) &= \sin \Theta_{r,\nu}, & 1 \leq \nu \leq N, \\ \phi_{r,\nu}(y) &= \cos \Theta_{r,\nu}, & N + 1 \leq \nu \leq 2N.\end{aligned}\tag{9}$$

$W_r(y) = W(\phi_{r,1}, \dots, \phi_{r,2N})$ is the Wronskian

$$W_r(y) = \det[(\partial_y^{\mu-1} \phi_{r,\nu})_{\nu, \mu \in [1, \dots, 2N]}].\tag{10}$$

Then we get the following link between Fredholm and Wronskian determinants [12]

Theorem 2.1

$$\det(I + A_r) = k_r(0) \times W_r(\phi_{r,1}, \dots, \phi_{r,2N})(0),\tag{11}$$

where

$$k_r(y) = \frac{2^{2N} \exp(i \sum_{\nu=1}^{2N} \Theta_{r,\nu})}{\prod_{\nu=2}^{2N} \prod_{\mu=1}^{\nu-1} (\gamma_\nu - \gamma_\mu)}.$$

In (11), the matrix A_r is defined by (3).

It can be deduced the following result :

Theorem 2.2 *The function v defined by*

$$v(x, t) = \frac{W_3(0)}{W_1(0)} \exp(2it - i\varphi).\tag{12}$$

is solution of the NLS equation (1)

$$iv_t + v_{xx} + 2|v|^2v = 0.$$

2.2 Quasi-rational solutions of NLS equation

In the following, we take the limit when the parameters $\lambda_j \rightarrow 1$ for $1 \leq j \leq N$ and $\lambda_j \rightarrow -1$ for $N + 1 \leq j \leq 2N$.

We consider the parameter λ_j written in the form

$$\lambda_j = 1 - 2j^2\epsilon^2, \quad 1 \leq j \leq N. \quad (13)$$

When ϵ goes to 0, we realize limited expansions at order p , for $1 \leq j \leq N$, of the terms

$$\begin{aligned} \kappa_j &= 4j\epsilon(1 - \epsilon^2j^2)^{1/2}, \quad \delta_j = 4j\epsilon(1 - 2\epsilon^2j^2)(1 - \epsilon^2j^2)^{1/2}, \\ \gamma_j &= j\epsilon(1 - \epsilon^2j^2)^{-1/2}, \quad x_{r,j} = (r - 1) \ln \frac{1+i\epsilon j(1-\epsilon^2j^2)^{-1/2}}{1-i\epsilon j(1-\epsilon^2j^2)^{-1/2}}, \\ \kappa_{N+j} &= 4j\epsilon(1 - \epsilon^2j^2)^{1/2}, \quad \delta_{N+j} = -4j\epsilon(1 - 2\epsilon^2j^2)(1 - \epsilon^2j^2)^{1/2}, \\ \gamma_{N+j} &= 1/(j\epsilon)(1 - \epsilon^2j^2)^{1/2}, \quad x_{r,N+j} = (r - 1) \ln \frac{1-i\epsilon j(1-\epsilon^2j^2)^{-1/2}}{1+i\epsilon j(1-\epsilon^2j^2)^{-1/2}}. \end{aligned}$$

Then we get quasi-rational solutions of the NLS equation given by :

Theorem 2.3 *With the parameters λ_j defined by (13), a_j and b_j chosen as in (??), for $1 \leq j \leq N$, the function v defined by*

$$v(x, t) = \exp(2it - i\varphi) \lim_{\epsilon \rightarrow 0} \frac{W_3(0)}{W_1(0)}, \quad (14)$$

is a quasi-rational solution of the NLS equation (1)

$$iv_t + v_{xx} + 2|v|^2v = 0.$$

3 Expression of solutions of NLS equation in terms of a ratio of two determinants

We construct here solutions of the NLS equation which is expressed as a quotient of two determinants.

For this we need the following notations :

$$\begin{aligned} A_\nu &= \kappa_\nu x/2 + i\delta_\nu t - ix_{3,\nu}/2 - ie_\nu/2, \\ B_\nu &= \kappa_\nu x/2 + i\delta_\nu t - ix_{1,\nu}/2 - ie_\nu/2, \end{aligned}$$

for $1 \leq \nu \leq 2N$, with κ_ν , δ_ν , $x_{r,\nu}$ defined in (5), (6) and (7).

The parameters e_ν are defined by (8).

With particular special choices of the parameters a_j and b_j , for $1 \leq N$, we get new deformations depending on four parameters. Below we use the following functions :

$$\begin{aligned} f_{4j+1,k} &= \gamma_k^{4j-1} \sin A_k, & f_{4j+2,k} &= \gamma_k^{4j} \cos A_k, \\ f_{4j+3,k} &= -\gamma_k^{4j+1} \sin A_k, & f_{4j+4,k} &= -\gamma_k^{4j+2} \cos A_k, \end{aligned} \quad (15)$$

for $1 \leq k \leq N$, and

$$\begin{aligned} f_{4j+1,k} &= \gamma_k^{2N-4j-2} \cos A_k, & f_{4j+2,k} &= -\gamma_k^{2N-4j-3} \sin A_k, \\ f_{4j+3,k} &= -\gamma_k^{2N-4j-4} \cos A_k, & f_{4j+4,k} &= \gamma_k^{2N-4j-5} \sin A_k, \end{aligned} \quad (16)$$

for $N+1 \leq k \leq 2N$.

We define the functions $g_{j,k}$ for $1 \leq j \leq 2N$, $1 \leq k \leq 2N$ in the same way, we replace only the term A_k by B_k .

$$\begin{aligned} g_{4j+1,k} &= \gamma_k^{4j-1} \sin B_k, & g_{4j+2,k} &= \gamma_k^{4j} \cos B_k, \\ g_{4j+3,k} &= -\gamma_k^{4j+1} \sin B_k, & g_{4j+4,k} &= -\gamma_k^{4j+2} \cos B_k, \end{aligned} \quad (17)$$

for $1 \leq k \leq N$, and

$$\begin{aligned} g_{4j+1,k} &= \gamma_k^{2N-4j-2} \cos B_k, & g_{4j+2,k} &= -\gamma_k^{2N-4j-3} \sin B_k, \\ g_{4j+3,k} &= -\gamma_k^{2N-4j-4} \cos B_k, & g_{4j+4,k} &= \gamma_k^{2N-4j-5} \sin B_k, \end{aligned} \quad (18)$$

for $N+1 \leq k \leq 2N$.

Then we get the following result :

Theorem 3.1 *The function v defined by*

$$v(x, t) = \exp(2it - i\varphi) \times \frac{\det((n_{jk})_{j,k \in [1, 2N]})}{\det((d_{jk})_{j,k \in [1, 2N]})} \quad (19)$$

is a quasi-rational solution of the NLS equation (1)

$$iv_t + v_{xx} + 2|v|^2v = 0,$$

where

$$n_{j1} = f_{j,1}(x, t, 0), \quad 1 \leq j \leq 2N \quad n_{jk} = \frac{\partial^{2k-2} f_{j,1}}{\partial \epsilon^{2k-2}}(x, t, 0), \quad 2 \leq k \leq N, \quad 1 \leq j \leq 2N$$

$$n_{jN+1} = f_{j,N+1}(x, t, 0), 1 \leq j \leq 2N \quad n_{jN+k} = \frac{\partial^{2k-2} f_{j,N+1}}{\partial \epsilon^{2k-2}}(x, t, 0), 2 \leq k \leq N, 1 \leq j \leq 2N$$

$$d_{j1} = g_{j,1}(x, t, 0), 1 \leq j \leq 2N \quad d_{jk} = \frac{\partial^{2k-2} g_{j,1}}{\partial \epsilon^{2k-2}}(x, t, 0), 2 \leq k \leq N, 1 \leq j \leq 2N$$

$$d_{jN+1} = g_{j,N+1}(x, t, 0), 1 \leq j \leq 2N \quad d_{jN+k} = \frac{\partial^{2k-2} g_{j,N+1}}{\partial \epsilon^{2k-2}}(x, t, 0), 2 \leq k \leq N, 1 \leq j \leq 2N$$

The functions f and g are defined in (15), (16), (17), (18).

We don't give here the proof of this result in order to not weight down the text. We postpone the redaction of the proof to a next publication.

The solutions of the NLS equation can also be written in the form :

$$v(x, t) = \exp(2it - i\varphi) \times Q(x, t)$$

where $Q(x, t)$ is defined by :

$$Q(x, t) := \frac{\begin{vmatrix} f_{1,1}[0] & \dots & f_{1,1}[N-1] & f_{1,N+1}[0] & \dots & f_{1,N+1}[N-1] \\ f_{2,1}[0] & \dots & f_{2,1}[N-1] & f_{2,N+1}[0] & \dots & f_{2,N+1}[N-1] \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ f_{2N,1}[0] & \dots & f_{2N,1}[N-1] & f_{2N,N+1}[0] & \dots & f_{2N,N+1}[N-1] \end{vmatrix}}{\begin{vmatrix} g_{1,1}[0] & \dots & g_{1,1}[N-1] & g_{1,N+1}[0] & \dots & g_{1,N+1}[N-1] \\ g_{2,1}[0] & \dots & g_{2,1}[N-1] & g_{2,N+1}[0] & \dots & g_{2,N+1}[N-1] \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ g_{2N,1}[0] & \dots & g_{2N,1}[N-1] & g_{2N,N+1}[0] & \dots & g_{2N,N+1}[N-1] \end{vmatrix}} \quad (20)$$

4 Quasi-rational solutions of order 3 with four parameters

We have already constructed in [11] solutions for the cases $N = 1, 2, 3$, and in [13] with two parameters.

4.1 Analytical expressions of the solutions of NLS equation with Four parameters

Here, we give the expression v of the solution of NLS equation with four parameters; it is defined by

$$v_3(x, t, a, b) = \frac{n(x, t)}{d(x, t)} \exp(2it) = \left(1 - 4 \frac{G_3(2x, 4t) + iH_3(2x, 4t)}{Q_3(2x, 4t)}\right) e^{2it}$$

with

$$G_3(X, T) = \sum_{k=0}^{12} g_k(T)X^k, \quad H_3(X, T) = \sum_{k=0}^{12} h_k(T)X^k, \quad Q_3(X, T) = \sum_{k=0}^{12} q_k(T)X^k.$$

$$\begin{aligned} \mathbf{g}_{12} &= 0, & \mathbf{g}_{11} &= 0, & \mathbf{g}_{10} &= 6, & \mathbf{g}_9 &= 0, & \mathbf{g}_8 &= 90T^2 + 90, & \mathbf{g}_7 &= 0, \\ \mathbf{g}_6 &= 300T^4 - 360T^2 - 960Tb_1 + 1260, & \mathbf{g}_5 &= -1440T^2a_1 - 1440a_1 + 18a_2, \\ \mathbf{g}_4 &= 420T^6 - 900T^4 - 2400T^3b_1 + 2700T^2 + 3600a_1^2 - 1200b_1^2 + (-2490b_1 \\ &- 30b_2)T - 2700, & \mathbf{g}_3 &= -4800T^4a_1 - 19200Ta_1b_1 + (-28800a_1 + 60a_2)T^2 \\ &+ 14400a_1 + 60a_2, & \mathbf{g}_2 &= 270T^8 + 2520T^6 + 40500T^4 + (57060b_1 - 180b_2)T^3 \\ &+ (-7200a_1^2 + 21600b_1^2 - 81000)T^2 - 7200a_1^2 - 7200b_1^2 + (-58140b_1 - 180b_2)T \\ &- 4050, & \mathbf{g}_1 &= -3360T^6a_1 + 19200T^3a_1b_1 + (7200a_1 - 150a_2)T^4 + 9600a_1^3 + 9600a_1b_1^2 \\ &+ (-21600a_1 - 540a_2)T^2 + (19920a_1b_1 + 240a_1b_2 - 240a_2b_1)T + 21600a_1 + 90a_2, \\ \mathbf{g}_0 &= 66T^{10} + 2970T^8 + 1440T^7b_1 + 13140T^6 + (17886b_1 + 42b_2)T^5 + (8400a_1^2 \\ &+ 3600b_1^2 - 45900)T^4 + (-44340b_1 + 420b_2)T^3 + (7200a_1^2 + 240a_1a_2 \\ &- 30480b_1^2 + 240b_1b_2 - 12150)T^2 + 18000a_1^2 + 240a_1a_2 + a_2^2 + 10489b_1^2 \\ &+ 166b_1b_2 + b_2^2 + (-9600a_1^2b_1 - 9600b_1^3 + 7470b_1 + 90b_2)T + 4050 \end{aligned}$$

$$\begin{aligned} \mathbf{h}_{12} &= 0, & \mathbf{h}_{11} &= 0, & \mathbf{h}_{10} &= 6T, & \mathbf{h}_9 &= 0, & \mathbf{h}_8 &= 30T^3 - 90T - 60b_1, & \mathbf{h}_7 &= 0, \\ \mathbf{h}_6 &= 60T^5 - 840T^3 - 480T^2b_1 - 900T - 305b_1 + 5b_2, & \mathbf{h}_5 &= -480T^3a_1 + 960a_1b_1 \\ &+ (1440a_1 + 18a_2)T, & \mathbf{h}_4 &= 60T^7 - 1260T^5 - 600T^4b_1 - 2700T^3 + (-1245b_1 - 15b_2)T^2 \\ &+ (3600a_1^2 - 1200b_1^2 - 8100)T - 555b_1 + 15b_2, & \mathbf{h}_3 &= -960T^5a_1 - 9600T^2a_1b_1 \\ &+ (-9600a_1 + 20a_2)T^3 + 3320a_1b_1 + 40a_1b_2 - 40a_2b_1 + (14400a_1 - 60a_2)T, \\ \mathbf{h}_2 &= 30T^9 - 360T^7 + 10260T^5 + (21465b_1 - 45b_2)T^4 + (-2400a_1^2 + 7200b_1^2 \\ &- 37800)T^3 + 4800a_1^2b_1 + 4800b_1^3 + (-14130b_1 + 90b_2)T^2 + (7200a_1^2 + 7200b_1^2 \\ &+ 28350)T + 22005b_1 + 135b_2, & \mathbf{h}_1 &= -480T^7a_1 + 4800T^4a_1b_1 + (10080a_1 - 30a_2)T^5 \\ &+ (21600a_1 - 60a_2)T^3 + (9960a_1b_1 + 120a_1b_2 - 120a_2b_1)T^2 + 4440a_1b_1 - 120a_1b_2 \\ &+ 120a_2b_1 + (9600a_1^3 + 9600a_1b_1^2 + 64800a_1 + 450a_2)T, & \mathbf{h}_0 &= 6T^{11} + 150T^9 + 180T^8b_1 \\ &- 5220T^7 + (101b_1 + 7b_2)T^6 + (1680a_1^2 + 720b_1^2 - 57780)T^5 + (-63975b_1 + 75b_2)T^4 \\ &+ (-12000a_1^2 + 80a_1a_2 - 24560b_1^2 + 80b_1b_2 - 14850)T^3 - 7760a_1^2b_1 + 80a_1^2b_2 \\ &- 160a_1a_2b_1 - 1840b_1^3 - 80b_1^2b_2 + (-4800a_1^2b_1 - 4800b_1^3 - 41085b_1 - 495b_2)T^2 \\ &+ (-25200a_1^2 - 240a_1a_2 + a_2^2 - 14951b_1^2 - 314b_1b_2 + b_2^2 + 28350)T + 11835b_1 + 45b_2 \end{aligned}$$

$$\begin{aligned}
\mathbf{q}_{12} &= 1, \quad \mathbf{q}_{11} = 0, \quad \mathbf{q}_{10} = 6T^2 + 6, \quad \mathbf{q}_9 = 40a_1, \quad \mathbf{q}_8 = 15T^4 - 90T^2 - 120Tb_1 \\
&+ 135, \quad \mathbf{q}_7 = -2a_2, \quad \mathbf{q}_6 = 20T^6 - 180T^4 - 320T^3b_1 + 540T^2 + 240a_1^2 + 560b_1^2 \\
&+ (350b_1 + 10b_2)T + 2340, \quad \mathbf{q}_5 = -240T^4a_1 + 1920Ta_1b_1 + (1440a_1 + 18a_2)T^2 \\
&- 2160a_1 + 18a_2, \quad \mathbf{q}_4 = 15T^8 + 60T^6 - 240T^5b_1 - 1350T^4 + (-5630b_1 - 10b_2)T^3 \\
&+ (3600a_1^2 - 1200b_1^2 + 13500)T^2 + 3600a_1^2 - 40a_1a_2 + 280b_1^2 - 40b_1b_2 \\
&+ (3330b_1 - 90b_2)T + 3375, \quad \mathbf{q}_3 = -320T^6a_1 - 6400T^3a_1b_1 + (-14400a_1 \\
&+ 10a_2)T^4 - 3200a_1^3 - 3200a_1b_1^2 + (-43200a_1 - 60a_2)T^2 + (-31760a_1b_1 \\
&+ 80a_1b_2 - 80a_2b_1)T + 14400a_1 + 90a_2, \quad \mathbf{q}_2 = 6T^{10} + 270T^8 + 13500T^6 \\
&+ (11466b_1 - 18b_2)T^5 + (-1200a_1^2 + 3600b_1^2 + 78300)T^4 + (114660b_1 \\
&- 180b_2)T^3 + (7200a_1^2 + 64800b_1^2 - 36450)T^2 - 10800a_1^2 + a_2^2 - 9431b_1^2 \\
&- 74b_1b_2 + b_2^2 + (9600a_1^2b_1 + 9600b_1^3 - 58950b_1 - 450b_2)T + 12150, \\
\mathbf{q}_1 &= -120T^8a_1 + 1920T^5a_1b_1 + (-480a_1 - 10a_2)T^6 + (10800a_1 - 270a_2)T^4 \\
&+ (45040a_1b_1 + 80a_1b_2 - 80a_2b_1)T^3 + 9600a_1^3 + 160a_1^2a_2 - 2240a_1b_1^2 \\
&+ 320a_1b_1b_2 - 160a_2b_1^2 + (9600a_1^3 + 9600a_1b_1^2 - 108000a_1 - 990a_2)T^2 \\
&+ (-26640a_1b_1 + 720a_1b_2 - 720a_2b_1)T - 27000a_1 - 90a_2, \quad \mathbf{q}_0 = T^{12} + 126T^{10} \\
&+ 40T^9b_1 + 3735T^8 + (2086b_1 + 2b_2)T^7 + (560a_1^2 + 240b_1^2 + 15300)T^6 \\
&+ (-5214b_1 + 102b_2)T^5 + (3600a_1^2 + 40a_1a_2 - 12280b_1^2 + 40b_1b_2 + 143775)T^4 \\
&+ 6400a_1^4 + 12800a_1^2b_1^2 + 6400b_1^4 + (-3200a_1^2b_1 - 3200b_1^3 + 179730b_1 \\
&- 90b_2)T^3 + (32400a_1^2 - 240a_1a_2 + a_2^2 + 100249b_1^2 - 314b_1b_2 + b_2^2 \\
&+ 93150)T^2 + 39600a_1^2 + 360a_1a_2 + a_2^2 + 27649b_1^2 + 286b_1b_2 + b_2^2 + (22880a_1^2b_1 \\
&+ 160a_1^2b_2 - 320a_1a_2b_1 + 34720b_1^3 - 160b_1^2b_2 + 96750b_1 + 450b_2)T + 2025
\end{aligned}$$

4.2 Plots of the solutions of NLS equation with four parameters

Conversely to the study with two parameters given in preceding works [10, 11, 13], we get other type of symmetries in the plots in the (x, t) plane, in particular we obtain beside already known circular shapes, triangular configurations. We give some examples of this fact in the following.

4.2.1 Peregrine breather of order 3

If we choose $\tilde{a}_1 = \tilde{b}_1 = \tilde{a}_2 = \tilde{b}_2 = 0$, we obtain the classical Peregrine breather

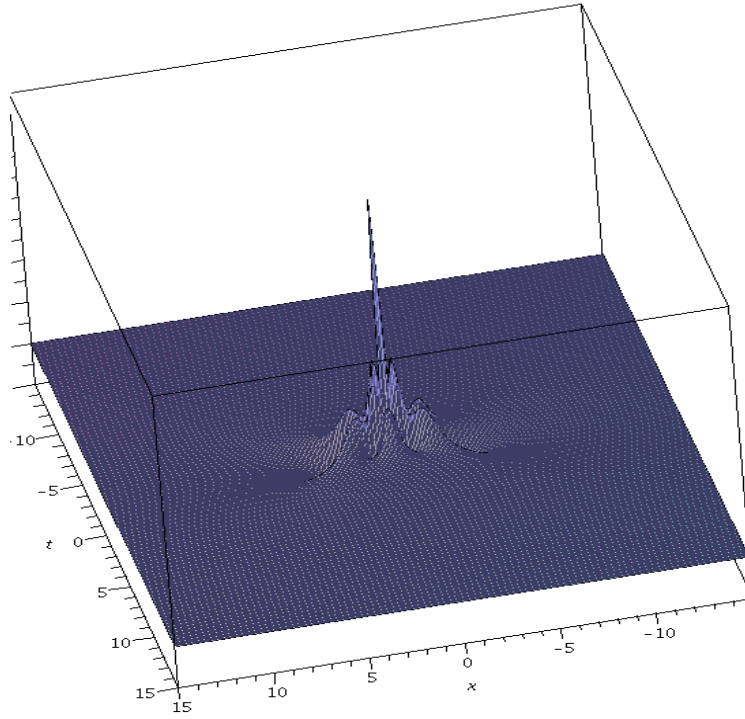


Figure 1: Solution of NLS, $N=3$, $\tilde{a}_1 = \tilde{b}_1 = \tilde{a}_2 = \tilde{b}_2 = 0$.

With other choices of parameters, we obtain all types of configurations : triangular with 6 peaks, circular with 6 peak, different cases with 1 until 6 peaks.

4.2.2 Variation of one parameter

If we choose $\tilde{a}_1 = 10^9$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 0$, we obtain :

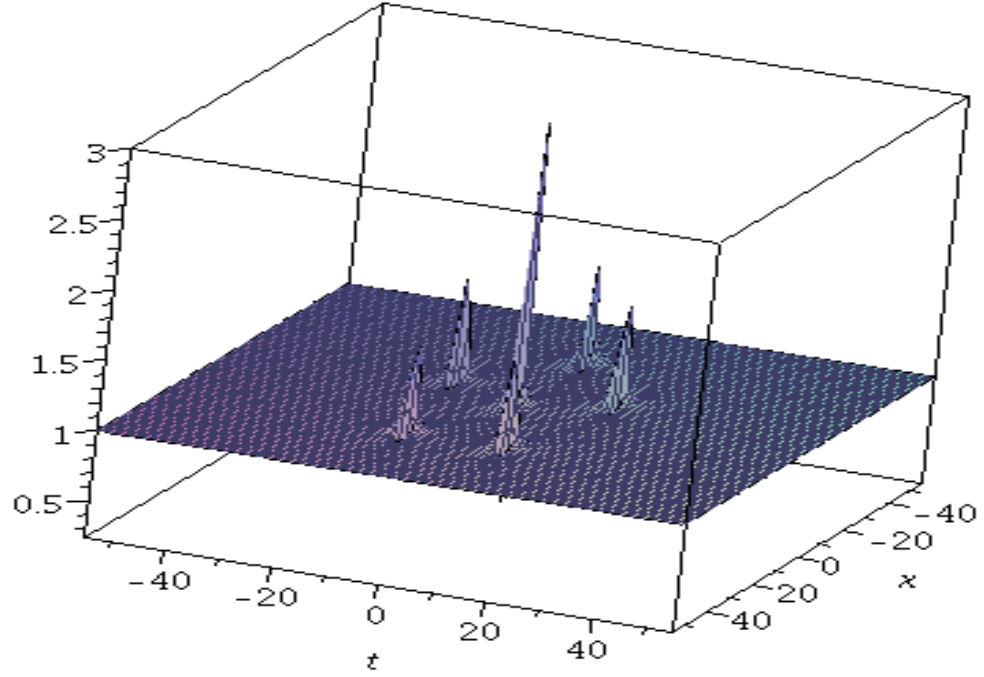


Figure 2: Solution of NLS, $N=3$, $\tilde{a}_1 = 10^9$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 0$.

If we choose $\tilde{a}_1 = 0$, $\tilde{b}_1 = 10^6$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 0$, we obtain :

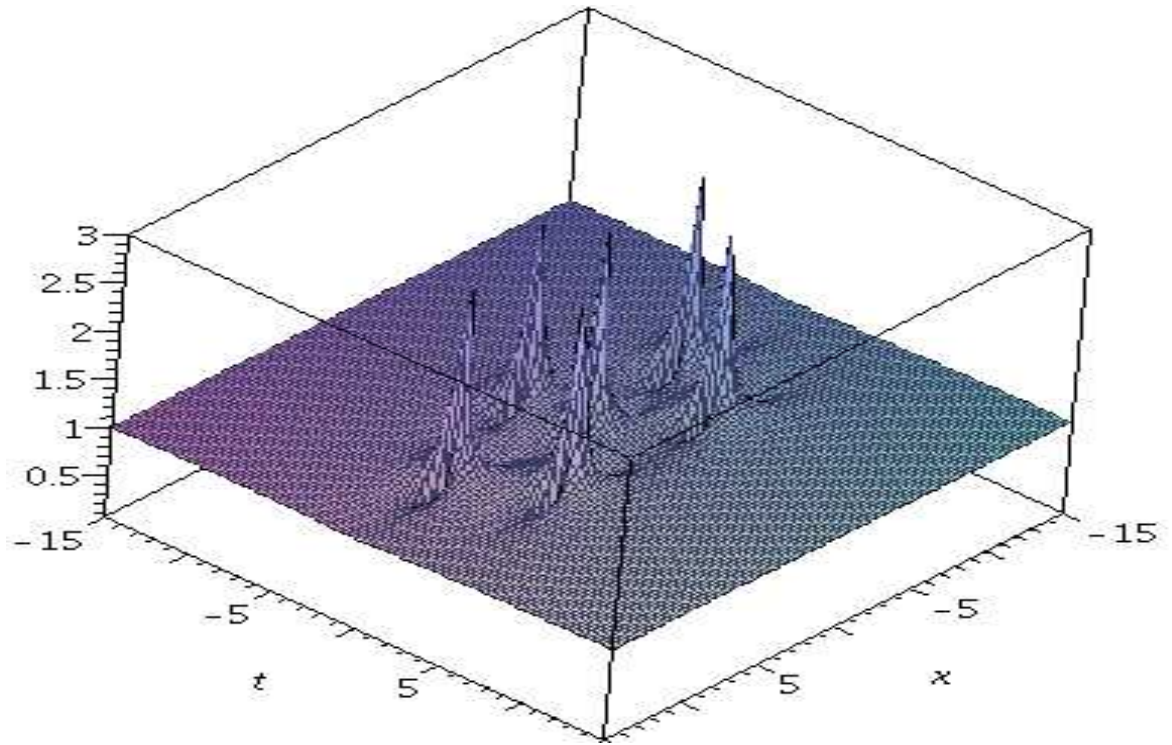


Figure 3: Solution of NLS, $N=3$, $\tilde{a}_1 = 0$, $\tilde{b}_1 = 10^6$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 0$.

If we choose $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 10^4$, $\tilde{b}_2 = 0$, we obtain :

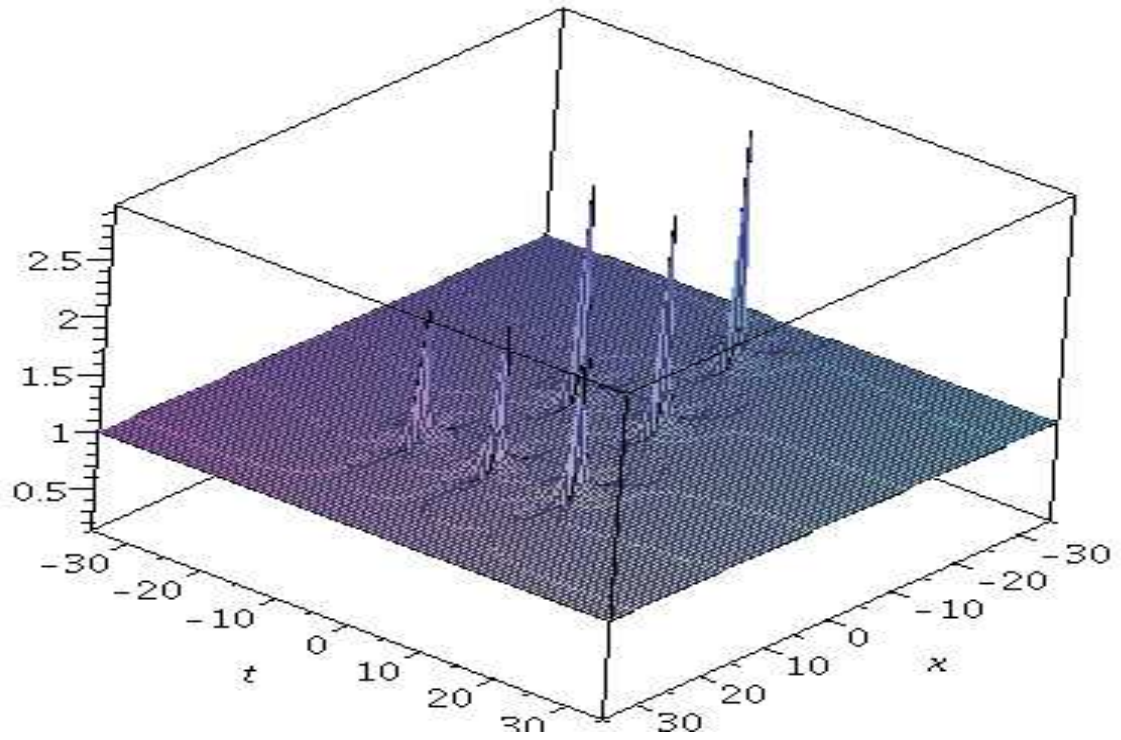


Figure 4: Solution of NLS, $N=3$, $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 10^4$, $\tilde{b}_2 = 0$, .

If we choose $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 10^5$, we obtain :

We obtain circular shapes in the case $a_2 = b_2 = 0$, and triangular configurations for $a_1 = b_1 = 0$. In the following we present the apparition of different configurations with 1 until 6 peaks.

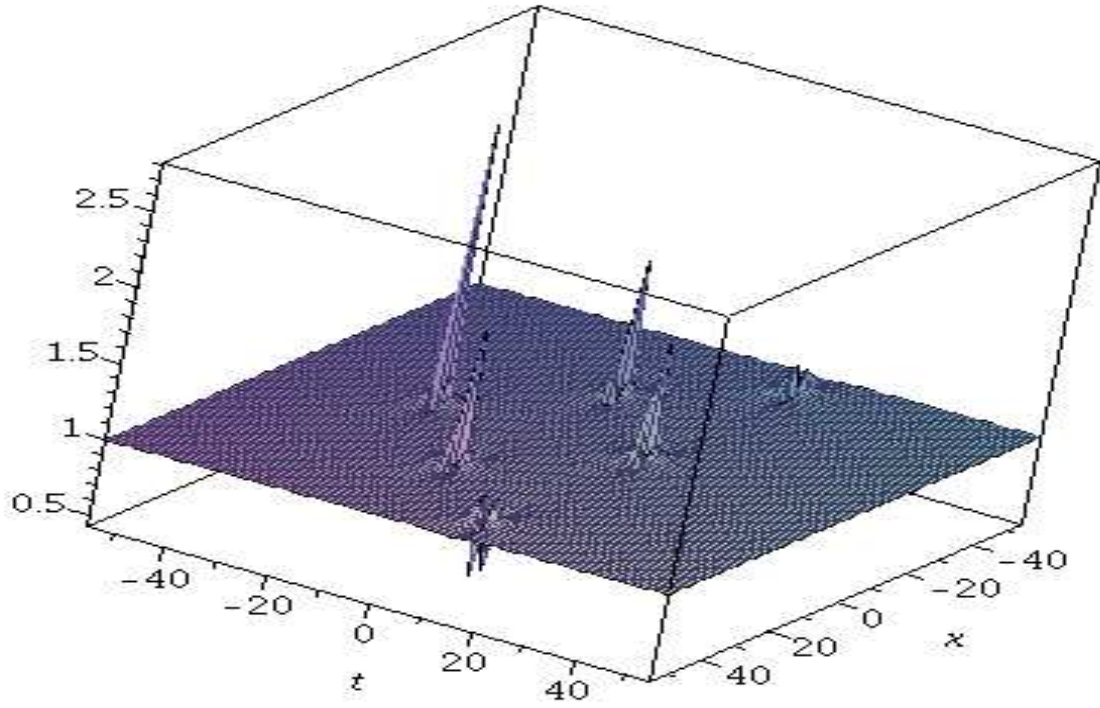


Figure 5: Solution of NLS, $N=3$, $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 10^5$.

4.2.3 Apparition of 1 until 6 peaks

If we choose $\tilde{a}_1 = 0$, $\tilde{b}_1 = 10^7$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 10^7$, we obtain :

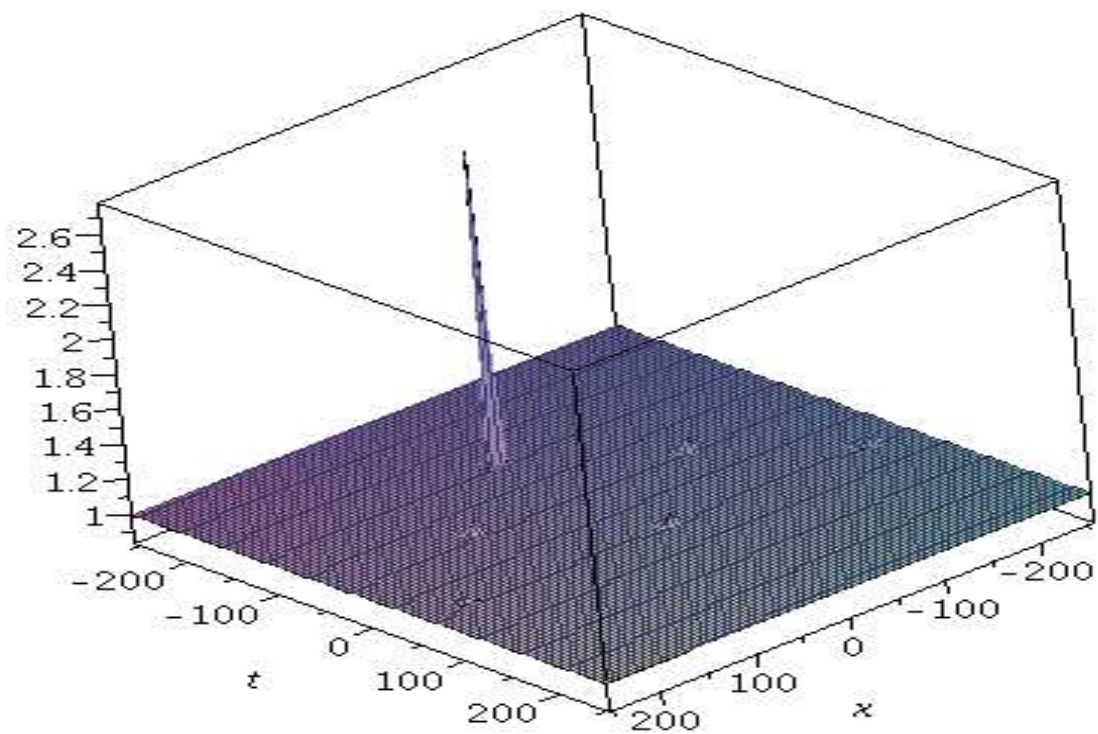


Figure 6: Solution of NLS, $N=3$, $\tilde{a}_1 = 0$, $\tilde{b}_1 = 10^7$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 10^7$.

If we choose $\tilde{a}_1 = 10^5$, $\tilde{b}_1 = 10^5$, $\tilde{a}_2 = 10^5$, $\tilde{b}_2 = 0$, we obtain :

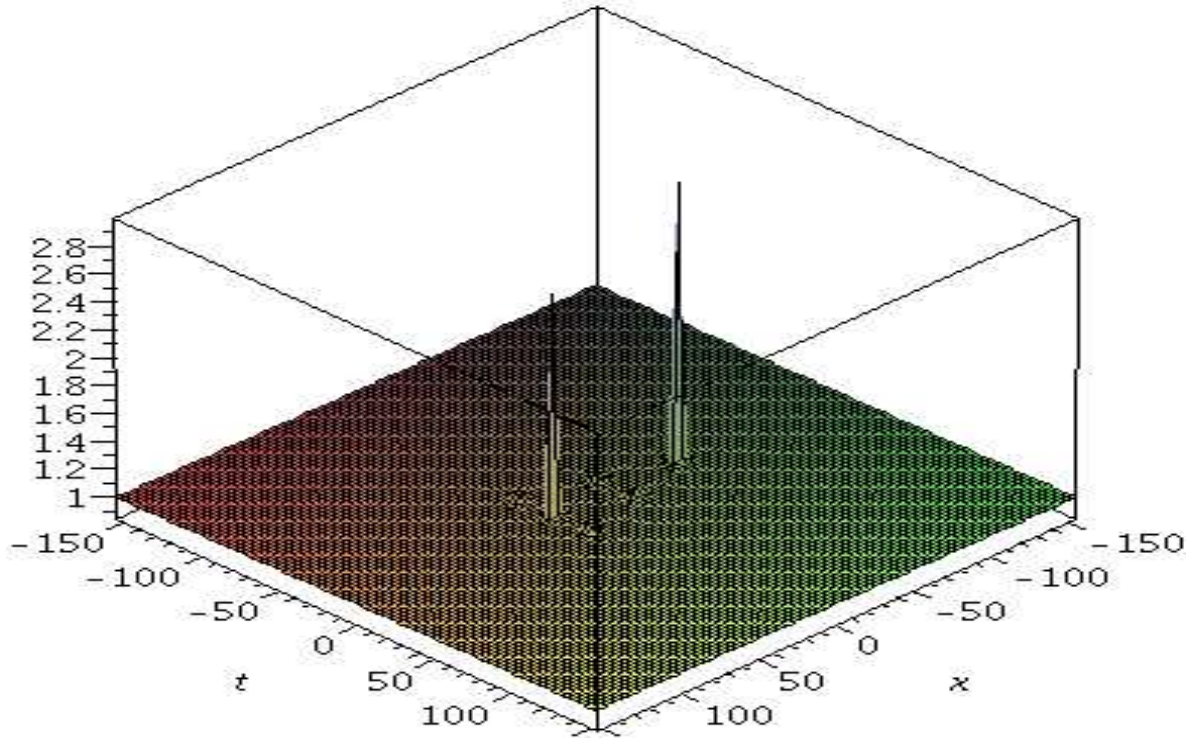


Figure 7: Solution of NLS, $N=3$, $\tilde{a}_1 = 10^5$, $\tilde{b}_1 = 10^5$, $\tilde{a}_2 = 10^5$, $\tilde{b}_2 = 0$.

If we choose $\tilde{a}_1 = 0$, $\tilde{b}_1 = 10^6$, $\tilde{a}_2 = 10^6$, $\tilde{b}_2 = 0$, we obtain :

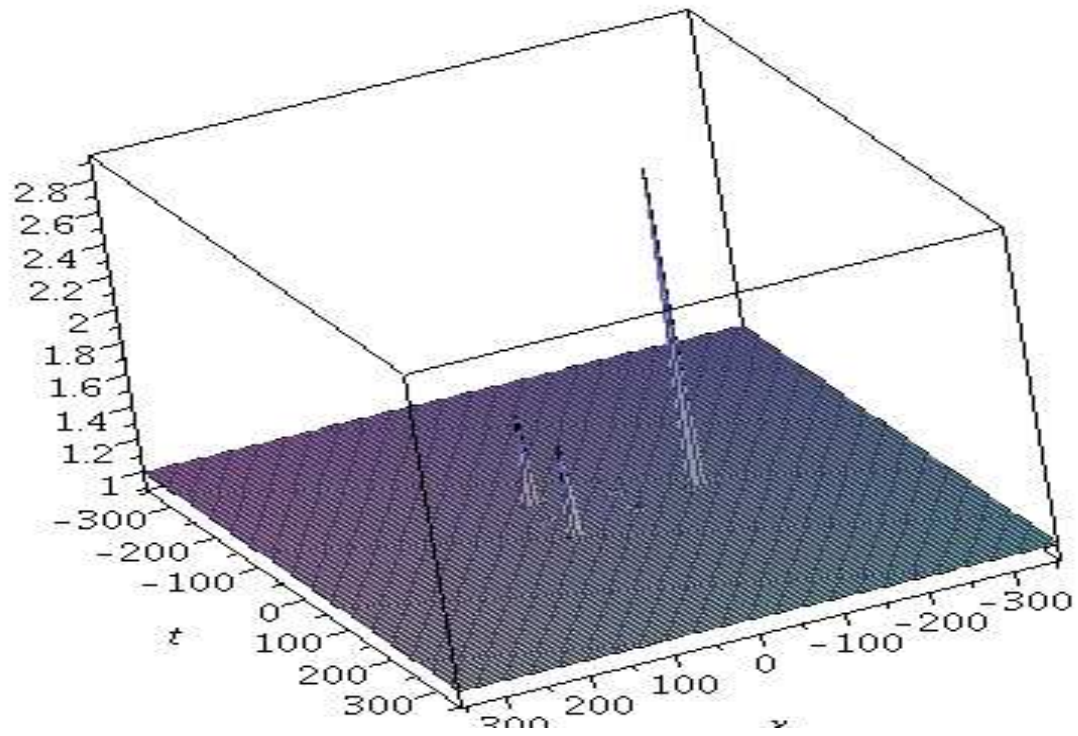


Figure 8: Solution of NLS, $N=3$, $\tilde{a}_1 = 0$, $\tilde{b}_1 = 10^6$, $\tilde{a}_2 = 10^6$, $\tilde{b}_2 = 0$.

If we choose $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 10^5$, $\tilde{b}_2 = 0$, we obtain :

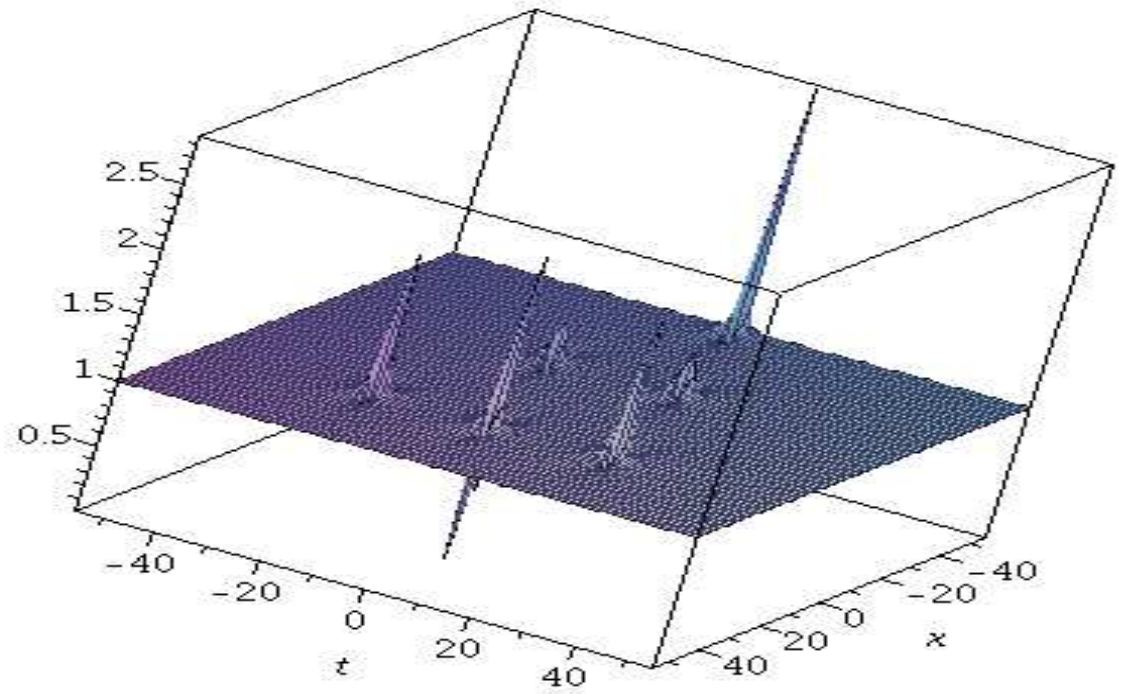


Figure 9: Solution of NLS, $N=3$, $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 10^5$, $\tilde{b}_2 = 0$.

If we choose $\tilde{a}_1 = 10^4$, $\tilde{b}_1 = 10^4$, $\tilde{a}_2 = 10^4$, $\tilde{b}_2 = 10^4$, we obtain :

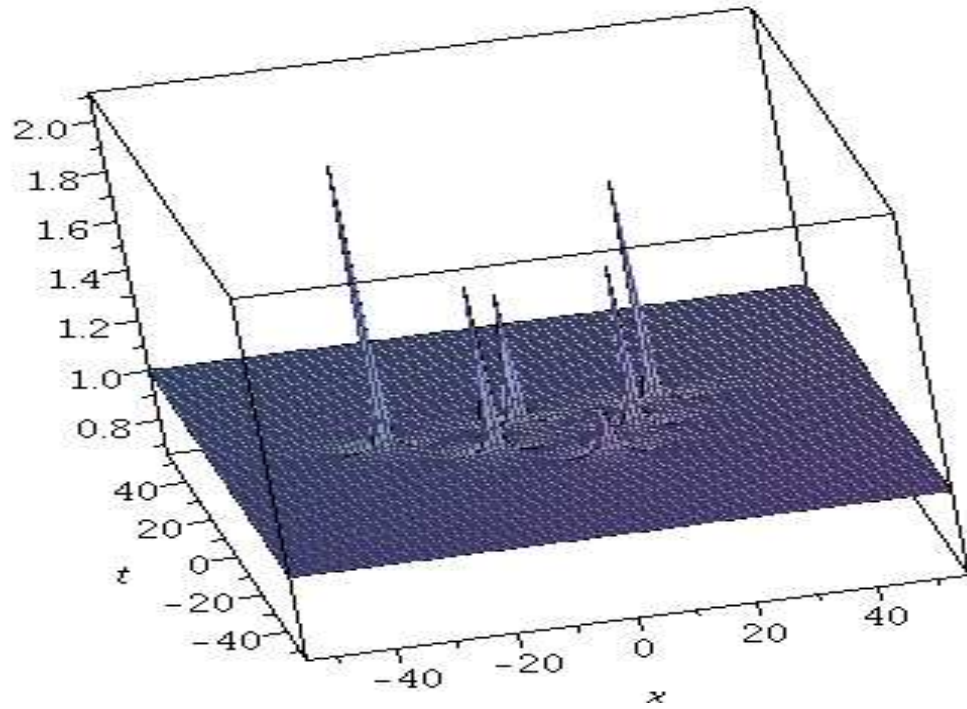


Figure 10: Solution of NLS, $N=3$, $\tilde{a}_1 = 10^4$, $\tilde{b}_1 = 10^4$, $\tilde{a}_2 = 10^4$, $\tilde{b}_2 = 10^4$.

If we choose $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 10^4$, $\tilde{b}_2 = 0$, we obtain :

4.2.4 Circular and triangular configurations

In general we obtain generically circular shapes in the case $a_2 = b_2 = 0$, and triangular configurations in the case $a_1 = b_1 = 0$. We present here some examples.

If we choose $\tilde{a}_1 = 0$, $\tilde{b}_1 = 10^7$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 0$, we obtain :

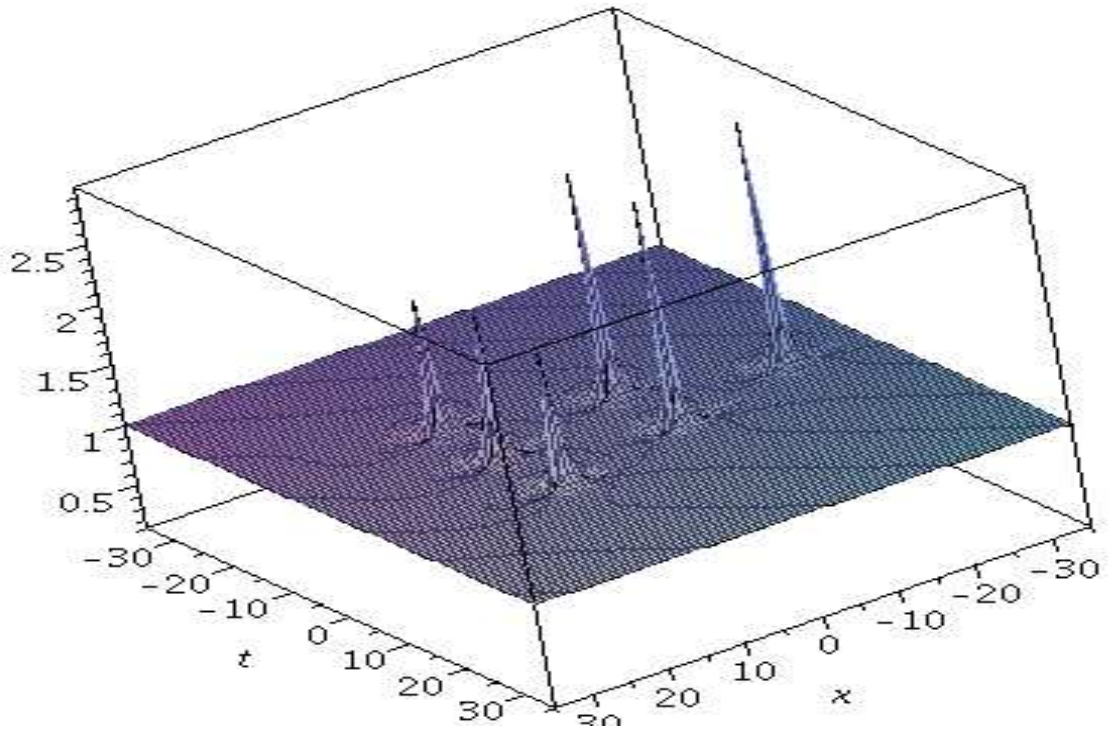


Figure 11: Solution of NLS, $N=3$, $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 10^4$, $\tilde{b}_2 = 0$.

If we choose $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 10^3$, we obtain :

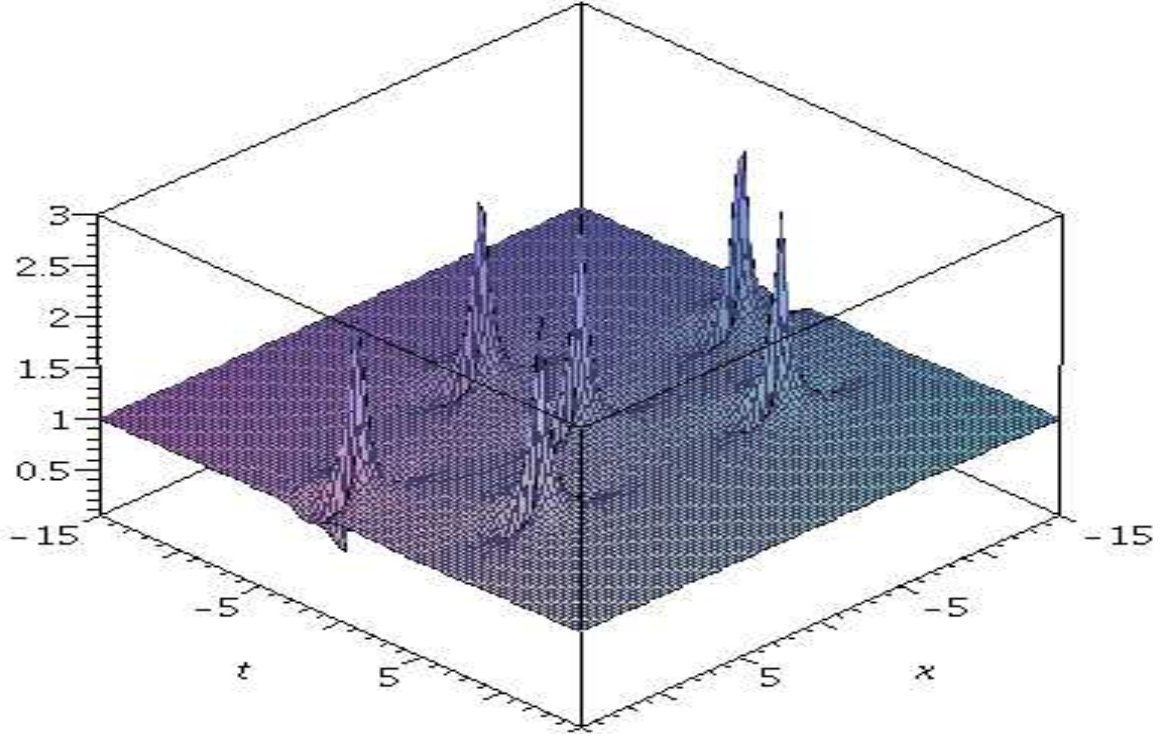


Figure 12: Solution of NLS, $N=3$, $\tilde{a}_1 = 0$, $\tilde{b}_1 = 10^7$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 0$.

5 Conclusion

The method described in the present paper provides a new tool to get explicitly solutions of the NLS equation.

The introduction of new parameters gives the appearance of new forms in conformity with those presented by Ohta and Yang [19], and also by Akhmediev et al. [17]. As it was already noted in previous studies with two parameters, one finds Peregrine breathers in the case where all the parameters are equal to 0.

We chose to present here the solutions of the NLS equation in the cases $N = 3$ only in order not to weigh down the text. We postpone the presentation of the higher orders in another publications.

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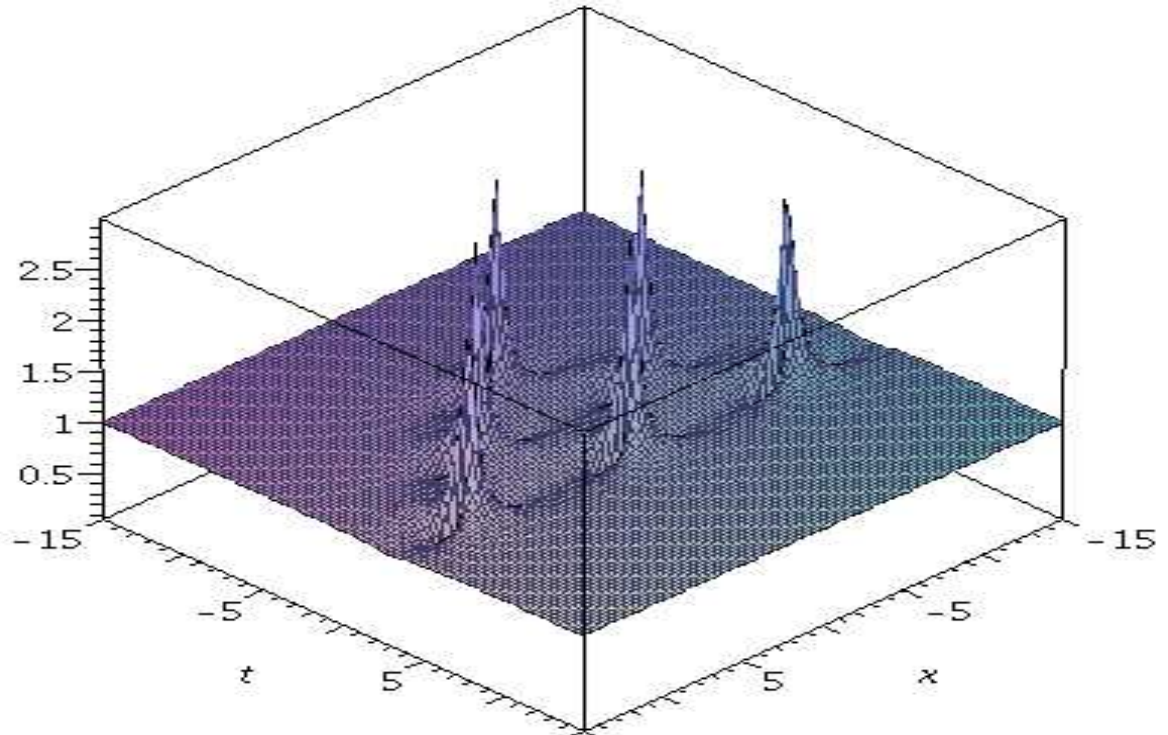


Figure 13: Solution of NLS, $N=3$, $\tilde{a}_1 = 0$, $\tilde{b}_1 = 0$, $\tilde{a}_2 = 0$, $\tilde{b}_2 = 10^3$.

universe of the nonlinear Schrödinger equation. I am very grateful to him for long fruitful discussions which we could have.

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