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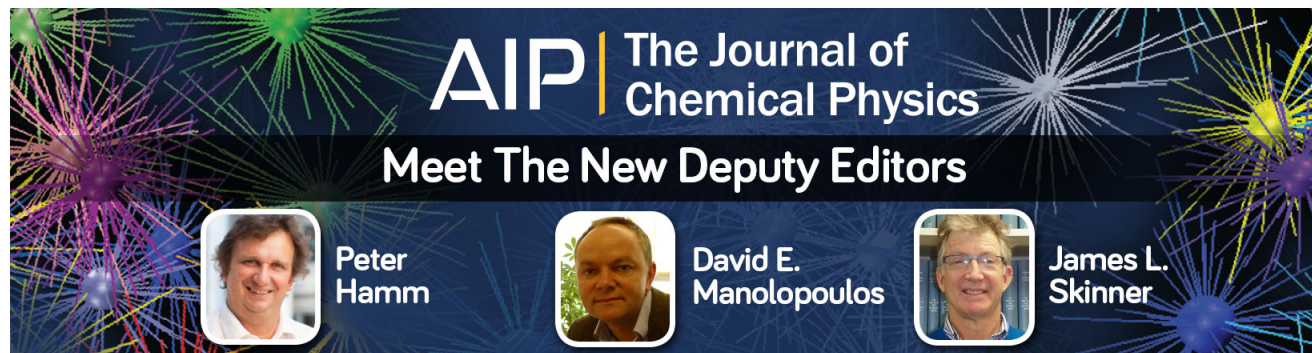
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The electronic structure of the triiodide ion from relativistic correlated calculations: A comparison of different methodologies

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The triiodide ion I_3^- exhibits a complex photodissociation behavior, the dynamics of which are not yet fully understood. As a first step toward determining the full potential energy surfaces of this species for subsequent simulations of its dissociation processes, we investigate the performance of different electronic structure methods [time-dependent density functional theory, complete active space perturbation theory to second order (CASPT2), Fock-space coupled cluster and multireference configuration interaction] in describing the ground and excited states of the triiodide ion along the symmetrical dissociation path. All methods apart from CASPT2 include scalar relativity and spin-orbit coupling in the orbital optimization, providing useful benchmark data for the more common two-step approaches in which spin-orbit coupling is introduced in the configuration interaction. Time-dependent density functional theory with the statistical averaging of model orbital potential functional is off the mark for this system. Another choice of functional may improve performance with respect to vertical excitation energies and spectroscopic constants, but all functionals are likely to face instability problems away from the equilibrium region. The Fock-space coupled cluster method was shown to perform clearly best in regions not too far from equilibrium but is plagued by convergence problems toward the dissociation limit due to intruder states. CASPT2 shows good performance at significantly lower computational cost, but is quite sensitive to symmetry breaking. We furthermore observe spikes in the CASPT2 potential curves away from equilibrium, signaling intruder state problems that we were unable to curb through the use of level shifts. Multireference configuration interaction is, in principle, a viable option, but its computational cost in the present case prohibits use other than for benchmarking purposes.   2010 American Institute of Physics. [doi:10.1063/1.3474571]

I. INTRODUCTION

Recent years have seen extraordinary advances in experimental techniques to probe chemical processes such as reaction dynamics in very short time frames. A wide range of techniques based on pump-probe schemes,¹ where the species under consideration are set in a nonstationary state by one light source (pump) and monitored by another (probe), providing information regarding the dynamical behavior of the system. A particularly interesting field, where such fast techniques are very helpful, is the study of the dynamics of stable negative ions.² By means of photodetachment or photodissociation these ions can be used to provide access to unstable neutral species that are difficult to study directly.

Some ionic species, however, are important in their own

right, apart from being used as precursors to other systems. A very well-known example is the triiodide ion (I_3^-). From a chemist's point of view,³⁻⁶ this relatively simple system is very interesting as it is (a) a structural analog of a transition state in S_N2 reactions and (b) an example of hyperconjugation, with a three-center four-electron bond. The widely accepted bonding picture of I_3^- is that of a (σ, π) system arising from the combinations of 5p orbitals on the three iodine atoms. In ΛS -coupling the ground state configuration is accordingly $\sigma_u^2 \pi_u^4 \pi_g^4 \pi_u^{*4} \sigma_g^2$, where the LUMO is the antibonding σ_u^* . This bonding picture is modified by spin-orbit coupling (SOC) (cf. Fig. 1), in particular for the MOs of *gerade* symmetry that are nonbonding and retain the atomic SO-mixing. In *ungerade* symmetry spin-orbit coupling is quenched due to bond formation and the orbitals are therefore essentially spin-pure. We have chosen to retain the ΛS notation of all orbitals for easier reference to other work, but note that the MOs denoted $\sigma_{1/2g}$ and $\pi_{1/2g}$ in Fig. 1 corre-

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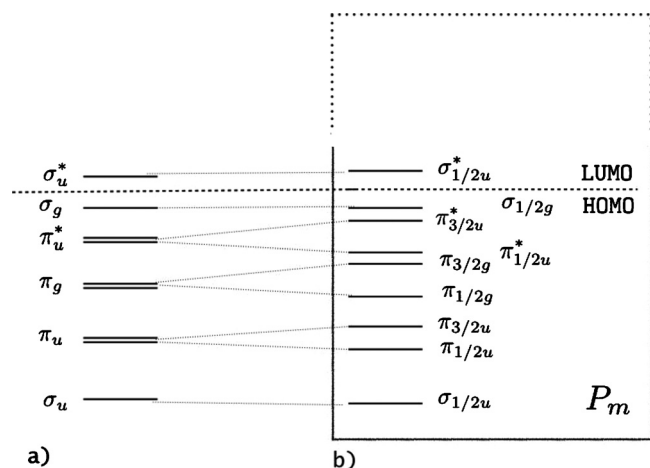


FIG. 1. Diagram of valence molecular orbitals of the triiodide ion, arising predominantly from the 5p manifold of the three iodine atoms: (a) spin-orbit free case and (b) spin-orbit case with the model P_m space employed in the IHFSCC calculations indicated.

spond to linear combinations of $p_{1/2}$ and $p_{3/2}$ on the terminal atoms, respectively.

While early experimental information on the electronic structure of I_3^- was obtained from UV spectroscopy,⁷ more detailed information has been gained from photodissociation studies in more recent works. These studies have originally been performed in solution.^{8–13} Questions arose as to whether it would be possible for the system to display a bent configuration at equilibrium, instead of the linear one proposed initially (the latter being the most accepted configuration). By contrast, experiments on the system in the gas-phase by Neumark and co-workers^{14–16} and by Nagata and co-workers^{17,18} strongly suggest that the ground state for this species is linear and centrosymmetric. Photofragment yield (PFY) spectra^{15,17} reveal two photoabsorption bands peaking at 3.43–3.45 and 4.25–4.28 eV. They correspond to the C and D bands observed in solution,⁷ and the peak values can be assigned as the vertical excitation energies from the ground to the excited $^3\Pi_u(0_u^+)$ and $^1\Sigma_u(0_u^+)$ states. Time of flight photofragment mass spectrometry by Nakanishi *et al.*¹⁸ reveal a 40:60 ratio between three- and two-body dissociations both in the C and D bands. The three-body dissociation is dominated by charge-asymmetric dissociation that is with the negative charge on one of the terminal iodine atoms. The two-body dissociation in the C and D bands is dominated by the (I^-+I_2) and $(I+I_2^-)$ channels, respectively. Although the same dissociation channels appear in the fast beam photofragment translation spectroscopy of Neumark and co-workers,¹⁶ there are some discrepancies, notably in the calculated branching fractions for the two- and three-body dissociations as well as the $(I^-+I_2)/(I+I_2^-)$ ratio, which call for further investigation.

Compared to the wealth of experimental data on the triiodide species, relatively few theoretical studies have been performed. Kosloff and co-workers,¹⁹ at about the time of the gas-phase experiments by Neumark,^{14,15} presented an extensive study of the potential energy curves of the I_3^- , I_2 , and I_2^- species, which are thought to be involved in the photodissociation dynamics of I_3^- . Their results, obtained with a combi-

nation of multireference configuration interaction (MRCI) calculations and a diatomics-in-molecule (DIM) treatment where SOC effects were included, were the most accurate published at the time. However, the authors themselves considered the basis set used in the DIM treatment relatively small allowing for quantitative improvements on this early study. More recently, as a complement to their experimental work, Nakanishi *et al.*¹⁸ performed comprehensive spin-orbit configuration interaction (SOC-ICI) calculations which explored potential surfaces not only along the symmetric stretch of I_3^- but also along the asymmetric stretch as well. Given the limited treatment of electron correlation and the still rather modest basis set size, also this theoretical investigation still leaves room for improvement.

Accordingly, the theoretical studies performed so far, although very helpful in understanding some aspects of the experimental data (such as the kind of states involved in the two experimentally observed absorption bands), are probably not sufficiently accurate for a direct comparison with experiment in terms of quantities associated with the dynamical behavior of the system, such as branching ratios. To enable this kind of analysis, it should be possible to calculate the whole of the potential energy surfaces accurately. This has motivated us to investigate the use of different methodologies for describing the electronic structure of the triiodide species to find the best candidate for calculation of a faithful representation of the potential energy surfaces that can be used in studies of the dynamics of the system.

Even though I_3^- itself is a closed-shell species and can undergo two-body dissociation into two closed-shell species as well, it is important to account for SOC to achieve an accurate description of the full potential energy surfaces. Other possible channels involve dissociation into radicals, and furthermore SOC can have a significant impact on transition probabilities and intersystem crossings in the excitation/de-excitation processes. Apart from including SOC, theoretical methods should be able to describe the full potential energy surface (PES) and account for dynamical electron correlation at a reasonable cost, given the amount of calculations involved in determining a full PES.

As not all methods will fit this description we have chosen to compare expensive but accurate multireference coupled cluster (CC) and configuration interactions (CI) at selected points along the symmetric stretch of the molecule to methods that will allow full coverage of the ground and excited state potential energy surfaces. The wave function-based methods employed were (a) the intermediate Hamiltonian Fock-space coupled cluster method of Eliav and co-workers,^{20–22} (b) the relativistic multireference CI method of Fleig and co-workers,^{23,24} and (c) the spin-orbit complete active space (CAS) perturbation theory to second order (PT2) method (SO-CASPT2) of Roos and co-workers.^{25,26} Apart from these, we decided to also explore the very economical time-dependent DFT (TDDFT) method²⁷ to provide a simple orbital picture for the excitation processes. These TDDFT calculations are of course not expected to yield better results than any of the methods above due to the well-known shortcomings of the current-day functionals (as, for

instance, in describing charge-transfer excitations^{28,29} or that, within the adiabatic approximation, TDDFT can only describe single excitations^{30–32}).

The paper is organized as follows: In Sec. II we present the details of the calculations performed with the different methodologies; in Sec. III we compare the performance of the different methods in calculating excitation energies of I_3^- , both at selected bond lengths and at the equilibrium geometries for the symmetric configuration. We also take the opportunity to discuss results for the triiodide radical (I_3), a species that was experimentally observed in photoionization studies involving I_3^- in the gas-phase,^{33,34} and that was also investigated theoretically by Kosloff and co-workers. Finally, in Sec. IV we assess the relative strengths and weaknesses of each method and provide concluding remarks.

II. COMPUTATIONAL DETAILS

All calculations described were performed along the symmetric stretch coordinate, with I–I bond lengths r_{I-I} in the range of [2.60:6.00] Å. In order to compare the different methodologies we have chosen two geometries in which to calculate the vertical excitation energies, apart from those obtained at the equilibrium structures for the different methods. These geometries are $r=2.84$ and 2.93 Å, corresponding to bond lengths in the vicinity of the equilibrium geometries for I_3 (from the MRCI calculations of Kosloff and co-workers¹⁹) and I_3^- (equilibrium geometry reported from solid state studies³⁵), respectively.

A. Intermediate Hamiltonian Fock-space coupled cluster

Intermediate Hamiltonian Fock-space coupled cluster^{20–22} (IHFSCC) calculations were performed with a development version of the DIRAC (Ref. 36) program. For describing the spectrum of I_3^- with the IHFSCC method we have taken the anion as a starting point, and proceeded from the ground state through the $(1h, 0p)$ and $(0h, 1p)$ sectors in order to arrive at the $(1h, 1p)$ sector and, therefore, at the excitation energies,

$$I_3^-(0h, 0p) \rightarrow \{I_3^-(1h, 0p), I_3^-(0h, 1p)\} \rightarrow I_3^{-(*)}(1h, 1p). \quad (2.1)$$

For reasons of computational efficiency, the exact two-component Hamiltonian (X2C) scheme of Iliaš and Saue³⁷ was used. Two-electron spin same-orbit (SSO) and spin-other orbit (SOO) contributions were included via atomic mean-field integrals obtained with the AMFI^{38,39} code.

The triple zeta basis sets of Dyall^{40,41} were employed in all calculations. The starting large component 28s21p15d SCF set was kept uncontracted and augmented by 1s1p1d1f diffuse functions. It thus contains the 2s2p2d primitives recommended for polarization and valence correlation as well as the 2s2p2d primitives recommended for core-valence correlation. By further augmentation two different sets were defined: a valence correlation set (“aVTZ”), where one correlating f-type function was added to the augmented SCF set,

yielding a (29s22p16d2f) basis set, and a core-valence correlating set (“aCVTZ”), which is a superset of aVTZ with additional 2f1g set of functions added.

In combination with these two sets different correlation spaces were employed: the first (Q_1) is used together with basis set aVTZ and includes the orbitals with orbital energies (ϵ) between -1 and 4 hartree. This means that in the occupied orbital space the σ, π bonding system and three other σ -type orbitals arising from the 5s5p orbitals of iodine are present. The second correlation space (Q_2) is used together with basis set aCVTZ and includes the orbitals with orbital energies between -3 and 12 hartree. This corresponds to enlarging Q_1 by including the occupied 4d electrons of iodine apart from more virtual orbitals. The combinations aVTZ/ Q_1 and aCVTZ/ Q_2 will be referred to as IHFSCC(a) and IHFSCC(b), respectively.

A crucial ingredient of IHFSCC calculations to prevent convergence problems due to intruder states is the definition of the model (P_m) and intermediate (P_I) spaces that comprise the active space $P=P_m+P_I$.^{21,22} After testing different spaces at the fixed geometries mentioned above, we found that convergence problems were generally avoided when the P_m space contained 8 occupied ($1\sigma_g, 1\pi_g, 2\pi_u$, and $1\sigma_u$ orbitals) and 11 virtual orbitals ($2\sigma_g, 1\pi_g, 3\sigma_u$, and $2\pi_u$ orbitals) with a full P space containing 11 occupied (5 of g and 6 of u parities) and 22 virtual (12 of g and 10 of u parities) orbitals (cf. Fig. 1). For bond lengths larger than 3.58 Å, however, it was not possible to obtain convergence for the $(1h, 1p)$ sector for this partition.

B. MRCI

MRCI calculations have been carried out with the relativistic double group CI program LUCIAREL,^{23,24} which recently has been extended^{42,43} for parallel computer applications and to allow for the computation of molecular properties.⁴⁴ In all of the calculations reported here the aCVTZ basis set and the X2C Hamiltonian including two-electron SSO and SOO corrections provided by the AMFI^{38,39} code have been employed.

The molecular spinors for the CI calculations have been obtained by an average-of-configuration Hartree–Fock calculation, where the open shells were defined as the 8 occupied Kramers pairs as in the CC application above, and in addition the antibonding σ_u orbital, corresponding to an active subspace with 16 electrons in 9 Kramers pairs. This type of Hartree–Fock wave function comprises a good starting point for relativistic MRCI studies since it provides a balanced description of ground and electronically excited states. The concept of general active spaces has been employed for constructing the CI expansion. In the present case, all Slater determinants with zero, one, and two particles in the external space (truncated at 3 hartrees) were included and all possible active space distributions were allowed for the remaining electrons. These active space distributions were defined by a CAS space corresponding to the above space used in the average-of-configuration Hartree–Fock calculation [16 electrons in 9 Kramers pairs, or 15 electrons in 9 Kramers pairs in the case of neutral I_3 , “MRCI(a)”], and an additional space

including the iodine 5s orbitals where a restriction of up to two holes was imposed [“MRCI(b)”]. The resulting relativistic CI wave function describes the correlation of 16 (15 in the case of neutral I_3) or 22 electrons and consists of roughly 77×10^6 Slater determinants in the latter case.

C. CASPT2

Calculations with the CASPT2 method were carried out with the MOLCAS code⁴⁵ (version 7.0), within the CASSCF/CASPT2/SO-RASSI approach. In this approach, scalar relativistic effects are included in the CASSCF (Ref. 25) and CASPT2 (Ref. 26) calculations via the Douglas–Kroll–Hess⁴⁶ Hamiltonian, and in a subsequent calculation the CASPT2 spin-free states are used by the RASSI program to set up a spin-orbit coupling Hamiltonian.⁴⁷ In this Hamiltonian the one- and two-electron spin-orbit integrals are calculated in a mean-field fashion via the AMFI^{38,39} code. The basis set used in these calculations was the ANO-RCC (Ref. 48) set of TZP quality (7s6p4d2f1g). All calculations were carried out in C_{2h} symmetry.

The CASSCF active space used consisted of 16 electrons in the nine orbitals arising from the 5p orbitals of the three iodine atoms. We have kept the core orbitals (i.e., up to and including the 3d orbitals) frozen in the CASSCF (and subsequent CASPT2) calculations. The CASSCF states that entered the multistate CASPT2 calculations have been obtained from state-averaged calculations over nine roots for $^1\Sigma$, six roots for $^1\Pi$, eight roots for $^3\Sigma$, and six roots for $^3\Pi$, respectively. In the CASPT2 calculations an IPEA shift of 0.25 a.u. was used.⁴⁹

D. TDDFT

TDDFT calculations²⁷ were carried out with the ADF (Ref. 50) code. Scalar relativistic and spin-orbit effects were included via the zeroth-order regular approximation.⁵¹ In the TDDFT calculations noncollinear spin magnetization was used. We are also within the adiabatic local density approximation (LDA) approximation to TDDFT, therefore disregarding any time-dependence of the exchange-correlation kernel, while using only the derivatives of the LDA functional for computing the exchange-correlation contributions to the excitation energies.

We used the statistical averaging of model orbital potentials (SAOPs),⁵² in combination with the TZ2P basis set.⁵³ It should be noted that in the ADF implementation SAOP is used to construct the Kohn–Sham potential during the SCF cycle, whereas the ground state energy is calculated with the PW91 (Ref. 54) energy expression. This makes the ground state properties close to those obtained with PW91, the differences arising from the modification of the potential. More details can be found in the ADF documentation (for instance, Ref. 55). Given the dependence of TDDFT on the functional used, we shall use the label TDDFT/SAOP.

III. RESULTS AND DISCUSSION

In this section we discuss the results for the triiodide ion species obtained with the different methods. We begin by investigating the performance of the different methods in de-

TABLE I. Ground state spectroscopic constants (r_e in Å and ω_e in cm^{-1}) of I_3^- calculated with the DFT, MRCI(a), CASPT2, and IHFSCC(a,b) methods.

Method	r_e	ω_e
DFT	3.007	102
MRCI(a)	2.982	108
CASPT2	2.888	119
IHFSCC(a)	2.971	112
IHFSCC(b)	2.946	114
MRCI (Ref. 19, spin-free)	2.930	114
DIM+SO (Ref. 19)	2.966	95
Expt. (Ref. 35)	2.93	112 ± 1

termining ground state spectroscopic constants, before addressing the electronic spectrum. Given the exploratory nature of this paper, we will restrict ourselves to selected structures along the symmetric stretch coordinate.

When comparing methodologies, it should be kept in mind that each method, with the exception of TDDFT/SAOP, can be tuned through the selection of active electrons and orbitals. As reference we take IHFSCC(b), which we consider the most accurate methodology employed in this work.

We will focus mainly on vertical excitations calculated at selected geometries, but will also address “adiabatic” excitations for this particular cut of the surface, as a way to gain insight on how the different methods represent the overall shape of the surfaces. Finally, we will take a more detailed look the 0_u states, in particular, the absorbing 0_u^+ states, comparing the energetics and the excitation picture, in terms of the respective molecular orbitals for the different methods, before discussing the triiodide radical.

A. Ground state spectroscopic constants for I_3^-

The ground state spectroscopic constants of I_3^- obtained here are shown in Table I. We take the experimental bond length³⁵ in the solid state ($r_e=2.93$ Å) as a measure since to the best of our knowledge, there are no experimental bond length determinations for the triiodide ion in gas-phase.

With respect to this reference value CASPT2 shows a slight (about 0.04 Å) underestimation, whereas IHFSCC(a) shows a similar overestimation. Better agreement, with a deviation of less than 0.02 Å, is obtained with IHFSCC(b), where the 4d shell is included in the occupied space and the virtual space is truncated at a higher energy. The effect of electron correlation can be studied in a more systematic manner at the CASPT2 level. Our default scheme is a CAS consisting of the 5p manifold and in addition correlating 4s4p4d5s at the PT2 level, giving a bond length of 2.888 Å. Freezing 4s4p, which corresponds to the default CASPT2 correlation scheme in MOLCAS 7.2, increases the bond length to 2.914 Å. Freezing 4d gives an even more significant bond length extension to 2.966 Å, whereas in addition freezing 5s shortens the bond slightly to 2.958 Å. These results suggest that the inclusion of the 4s4p shells in the occupied space of the IHFSCC calculations will bring the bond length in even closer agreement with experiment and thus implies that environmental effects on the bond length are small. We were unable to optimize the I_3^- bond distance at the MRCI(b) level

TABLE II. Comparison of vertical excitation energies T_v (in eV) obtained with the different methods [TDDFT/SAOP, MRCI(a,b), CASPT2, and IHFSC-C(a,b)] for I_3^- at $r_1=r_2=2.93$ Å. States of the same symmetry as those for the optically active excited states are shown in boldface. Statistical measures of the error compared to IHFSCC(b) are also shown (see text for discussion).

State	TDDFT/SAOP		MRCI(a)		MRCI(b)		CASPT2		IHFSCC(a)		IHFSCC(b)	
	Ω	T_v	Ω	T_v	Ω	T_v	Ω	T_v	Ω	T_v	Ω	T_v
1	2_g	1.92	2_g	2.32	2_g	2.30	2_g	2.24	2_g	2.10	2_g	2.05
2	1_g	2.04	1_g	2.44	0_u^-	2.40	1_g	2.32	1_g	2.23	1_g	2.18
3	0_u^-	2.42	0_u^-	2.47	1_u	2.41	0_u^-	2.47	0_u^-	2.26	0_u^-	2.20
4	1_u	2.43	1_u	2.48	1_g	2.41	1_u	2.47	1_u	2.27	1_u	2.20
5	0_g^-	2.50	0_g^-	2.94	0_g^-	2.91	0_g^-	2.76	0_g^-	2.68	0_g^-	2.64
6	0_g^+	2.56	0_g^+	2.98	0_g^+	2.95	0_g^+	2.82	0_g^+	2.73	0_g^+	2.69
7	1_g	2.70	1_g	3.13	1_g	3.08	1_g	2.85	1_g	2.90	1_g	2.86
8	2_u	2.61	2_u	3.25	2_u	3.21	2_u	3.10	2_u	3.22	2_u	3.17
9	1_u	2.72	1_u	3.30	1_u	3.25	1_u	3.11	1_u	3.30	1_u	3.24
10	0_u^+	3.14	0_u^+	3.71	0_u^+	3.66	0_u^+	3.52	0_u^+	3.52	0_u^+	3.51
11	2_g	3.50	0_u^-	4.04	0_u^-	3.98	0_u^-	3.79	0_u^-	3.95	2_g	3.88
12	0_u^-	3.42	1_u	4.09	1_u	4.02	1_u	3.80	2_g	3.96	0_u^-	3.91
13	1_g	3.63	2_g	4.19	2_g	4.16	2_g	3.98	1_u	4.03	1_g	4.00
14	1_u	3.56	1_g	4.29	1_g	4.25	1_g	4.06	1_g	4.07	1_u	4.00
15	0_g^-	4.10	0_g^-	4.81	0_u^+	4.75	0_u^+	4.51	0_u^+	4.33	0_u^+	4.33
16	0_g^+	4.11	0_g^+	4.82	0_g^-	4.77	0_g^-	4.51	0_g^-	4.54	0_g^-	4.48
17	1_g	4.34	0_u^+	4.83	0_g^+	4.78	0_g^+	4.53	0_g^+	4.54	0_g^+	4.48
18	0_u^+	4.46	1_g	4.96	1_g	4.91	1_g	4.60	2_u	4.7	1_g	4.68
$\bar{\Delta}$		-0.24		0.25		0.20		0.05		0.05		
Δ_{std}		0.24		0.11		0.11		0.15		0.02		
$\bar{\Delta}_{\text{abs}}$		0.31		0.25		0.21		0.12		0.05		
Δ_{max}		0.56		0.50		0.41		0.27		0.08		

due to computational constraints. Based on the trends observed with CASPT2, one would expect that MRCI(b), which correlates the $5s5p$, would *overestimate* the bond length further compared to MRCI(a). This is surprising since Vala *et al.*¹⁹ with the same active space and basis set quality reproduced the experimental bond length in their spin-orbit free MRCI calculations. One possible source of error in the present MRCI calculations is the truncation of the virtual space, but from previous experience we deem it unlikely that this truncation affects the bond length by 0.05 Å. We instead believe that the results of Vala *et al.*¹⁹ are fortuitous since our studies clearly show that correlation of the $4s4p4d$ manifold has a significant effect on the spectroscopic constants. Indeed, if we subtract the 0.07 Å gained by correlating $4s4p4d5s$ at the CASPT2 level from our MRCI(a) result we come quite close to the experimental value.

It is interesting to note that DFT, where in principle all orbitals are correlated, overestimates r_e by about 0.07 Å. Clearly, the SAOP potential was not derived with the aim of providing accurate ground state structures. On the other hand, other generalized gradient approximation (GGA) functionals do not perform better, e.g., PW91 (2.996 Å), PBE (Ref. 56) (2.996 Å), M06-L (Ref. 57) (3.065 Å), and BP86 (3.007 Å). The latter value contrasts significantly with the bond length of 3.14 Å reported by Landrum *et al.*,⁵ but they employed a TZP basis, thus not including the diffuse functions of the TZ2P basis as we do. Hybrid functionals do perform better, e.g., PBE0 (Ref. 58) (2.946 Å) and M06-2X (Ref. 59) (2.925 Å), whereas LDA overbinds, giving (2.913 Å). We note in passing that the effect of spin-orbit

coupling on the bond lengths is on the order of 0.02 Å for the cited functionals, again an indication that the good agreement of the spin-orbit free MRCI calculation of Vala *et al.*¹⁹ with the experimental bond length is likely to be fortuitous.

A similar picture is seen for the harmonic frequencies. Discrepancies between MRCI, CASPT2, IHFSCC, and both experiment⁶⁰ and the spin-free MRCI calculations of Kosloff are of the order of a few cm^{-1} , whereas DFT underestimates the frequency by 10 cm^{-1} . The spin-orbit numbers of Kosloff, on the other hand, show rather large discrepancies (17 cm^{-1}), which may be due to artifacts arising from their diatomics-in-molecule treatment.

B. Benchmark calculations of the electronic spectra of I_3^-

In the comparison of calculated spectra, we consider all methods both at the experimental bond length ($r=2.93$ Å) and at their own optimal ground state equilibrium geometries (shown in Table I). The corresponding spectra can be found in Tables II and III, respectively.

As already mentioned, we discuss the performance of the other methods relative to IHFSCC(b). This method is accurate for the first 18 excited states, which go up to about 4.5–5.0 eV, but becomes less trustworthy for higher energies due to the increasing importance of double excitations in these states. Double excitations are readily captured by MRCI and CASPT2 but neither by TDDFT/SAOP nor by the IHFSCC calculations on the $(1h, 1p)$ sector of Fock-space.

TABLE III. Vertical (T_v) excitation energies (in eV) calculated with the TDDFT/SAOP, MRCI(a), CASPT2, and IHFSCC(a,b) methods for I_3^- calculated at the optimum bond length for each method. States of the same symmetry as those for the optically active excited states are shown in boldface. Statistical measures of the error compared to IHFSCC(b) are also shown (see text for discussion).

State	TDDFT/SAOP		MRCI(a)		CASPT2		IHFSCC(a)		IHFSCC(b)	
	Ω	T_v	Ω	T_v	Ω	T_v	Ω	T_v	Ω	T_v
1	2_g	1.67	2_g	2.15	2_g	2.38	2_g	1.98	2_g	2.00
2	1_g	1.78	1_g	2.27	1_g	2.45	1_g	2.10	1_g	2.13
3	0_u^-	2.13	0_u^-	2.28	1_u^-	2.63	0_u^-	2.11	0_u^-	2.13
4	1_u	2.14	1_u	2.28	0_u^-	2.63	1_u	2.12	1_u	2.14
5	0_g^-	2.24	0_g^-	2.77	0_g^-	2.90	0_g^-	2.55	0_g^-	2.58
6	2_u	2.26	0_g^+	2.81	0_g^+	2.95	0_g^+	2.60	0_g^+	2.64
7	0_g^+	2.31	1_g	2.94	1_g	3.00	1_g	2.77	1_g	2.81
8	1_u	2.38	2_u	3.00	2_u	3.31	2_u	3.04	2_u	3.10
9	1_g	2.44	1_u	3.05	1_u	3.32	1_u	3.11	1_u	3.17
10	0_u^+	2.82	0_u^+	3.47	0_u^+	3.72	0_u^+	3.44	0_u^+	3.36
11	2_g	3.07	0_u^-	3.79	0_u^-	3.99	2_g	3.73	2_g	3.79
12	0_u^-	3.09	1_u	3.84	1_u	4.00	0_u^-	3.77	0_u^-	3.84
13	1_g	3.20	2_g	3.89	2_g	4.23	1_g	3.84	1_g	3.90
14	1_u	3.22	1_g	3.98	1_g	4.31	1_u	3.85	1_u	3.93
15	0_g^-	3.66	0_g^-	4.52	0_u^+	4.67	0_u^+	4.17	0_u^+	4.27
16	0_g^+	3.68	0_g^+	4.52	0_g^-	4.75	0_g^-	4.31	0_g^-	4.38
17	1_g	3.91	0_u^+	4.61	0_g^+	4.78	0_g^+	4.32	0_g^+	4.39
18	0_u^+	4.29	1_g	4.65	1_g	4.85	1_g	4.50	1_g	4.59
$\bar{\Delta}$		-0.49		0.09		0.32		-0.05		
Δ_{std}		0.29		0.12		0.12		0.04		
$\bar{\Delta}_{\text{abs}}$		0.49		0.13		0.32		0.06		
Δ_{max}		0.84		0.34		0.50		0.10		

In the statistical analysis of the errors for each method we therefore include only the first 18 states that are dominated by single excitations. In this analysis that is found at the bottom of Tables II and III, we report the mean error $\bar{\Delta}$ and its standard deviation Δ_{std} , as well as the mean absolute error $\bar{\Delta}_{\text{abs}}$ and the maximum absolute error Δ_{max} relative to the IHFSCC(b) results.

1. Vertical excitations

Inspecting Table II we immediately see some general trends: CASPT2 and MRCI tend to overestimate the IHFSCC excitation energies, whereas TDDFT/SAOP shows the lowest excitation energies among all methods considered. One can also see that the degree of electron correlation introduced by the different methods affects distinct regions of the spectrum differently. For the lowest ten excited states, the ordering of states is consistent with most methods whereas for higher states, where the density of states is higher, small variations in the correlation treatment result in significant reorderings. The IHFSCC(a) scheme yields only small deviations relative to the larger IHFSCC(b) calculation. This is quantified by the small mean errors (signed and absolute) and a small standard deviation (0.02 eV), suggesting that core-valence correlation does not play a prominent role in describing the transitions to low-lying excited states. Such effects can be studied with more ease at the CASPT2 level. We find that freezing the 4s4p shells at the PT2 level has no effect on the vertical excitation energies, whereas freezing the 4d shell brings about an upward shift of about 0.12 eV, somewhat larger than what is observed at the IHFSCC level.

Freezing in addition the 5s shell leads to a mean upward shift of around 0.21 eV compared to our default CASPT2 scheme. With these observations in mind one can readily understand why CASPT2 reproduces better the IHFSCC(b) reference than the MRCI calculations since the latter have much more limited active spaces. TDDFT/SAOP underestimates ($\bar{\Delta}=0.24$ eV) the excitation energies, but perhaps more worrisome is that Δ_{std} is rather large (0.24 eV), indicating large nonsystematic errors.

While it is beyond the scope of this paper to present a detailed analysis of the difference between IHFSCC and other methods, it is nevertheless instructive to consider comparisons of excitation energies calculated by IHFSCC and the linear response coupled cluster (LRCC) method done by various authors.^{61–65} From those studies it becomes clear that LRCC and IHFSCC may yield significantly different excitation energies. Some evidence from recent studies for different molecular systems performed by some of us⁶⁶ as well as by other authors^{67–70} seems to point to a systematic upward shift in LRCC excitation energies compared to IHFSCC ones, similar to what is observed here when comparing MRCI and IHFSCC. This could be consistent with the different parametrization used in describing the excitations from the ground to the excited states (linear in MRCI and LRCC, and exponential for IHFSCC). It is less straightforward to rationalize the discrepancies between CASPT2 and IHFSCC as these methods describe the ground state with a rather different wave function. Based on previous experience^{66,71–73} we expect CASPT2 results for low-lying

excited states to deviate up to a few tenths of an eV from IHFSCC, but is more difficult to predict either an increase or a decrease of the excitation energies.

In Table III we present vertical excitation energies at the ground state equilibrium geometry optimized for each method. With a spread of 0.12 Å between the optimized bond lengths, the discrepancies with respect to the reference IHFSCC(b) data increase. The effect is accentuated by the fact that the vertical excitation energies probe the repulsive wall of the excited states, as will become clear in Sec. III B 2. We first note that the vertical excitation energies obtained with the IHFSCC(b) scheme to the optically active states—3.36 and 4.27 eV—are close to the experimentally observed peaks in the PFY spectra [3.43 and 4.25 eV (Ref. 18)]. The agreement is also quite acceptable with the IHFSCC(a) scheme, which again indicates that core-valence correlation (correlating deeper than the 5s shell) can possibly be disregarded when constructing the full potential energy surfaces. The corresponding CASPT2 results—3.72 and 4.67 eV—are close to the SOCI values reported by Nakanishi *et al.*¹⁸ (3.79 and 4.70 eV). MRCI(a) is reasonably close to the lower excitation energy, but severely overestimates the second one, thus providing a peak separation energy of 1.14 eV which is far from the experimentally observed 0.82 eV. In fact, it is only the IHFSCC method that is able to reproduce the experimental peak positions and separation; all other methods explored in this paper overestimate the splitting significantly with TDDFT/SAOP being simply off the mark for these excitations.

The statistical analysis at these selected geometries indicates that both MRCI and CASPT2 are able to provide a balanced treatment of the ground and excited state surfaces. A problem with MRCI(a), currently the only feasible CI approach to cover the whole surface, is the fact that the optically active states (indicated in boldface in both Tables II and III) that are well described with IHFSCC(b) exhibit relatively large errors when treated with MRCI(a). An important remark concerning CASPT2 is that our conclusions are based on calculations carried out in C_{2h} symmetry with the rotation axis aligned with the molecular axis. Since we are presently only exploring the symmetric stretch one would rather be inclined to use the highest possible point group symmetry available for a centrosymmetric configuration, which is D_{2h} with the MOLCAS code. However, the introduction of a vertical mirror plane places π_x and π_y orbitals in different irreducible representations and introduces a symmetry breaking to which CASPT2 is extremely sensitive. Vertical excitation energies at bond length 2.93 Å calculated in D_{2h} and C_{2v} symmetries give upper and lower bounds, respectively, for the corresponding C_{2h} values within a span approaching 0.1 eV. The choice of symmetry can therefore seriously affect conclusions regarding the performance of the method. On the other hand, restricting symmetry to C_{2h} means that the important distinction between 0^+ and 0^- states is lost. The use of supersymmetry is hardly an option in the scan of the potential surfaces, so we have in the present work rather relied on scripts that analyze the CASSCF orbitals to extract this symmetry information. Having to work at re-

duced symmetry also implies that orbitals are optimized and thus averaged for a larger number of states.

We also note that the two other experimentally observed⁷ bands—the A band (at 2.19 eV in CH_2Cl_2 and 2.16 eV in Et_4NI_3 crystal), assigned to a ${}^3\Pi_u$, and the B band (at 2.82 eV in CH_2Cl_2 and 2.71 eV in Et_4NI_3 crystal), assigned to a ${}^3\Sigma_u^+$ state—can indeed be correlated with the first and second 1_u states calculated for the vertical excitations with IHFSCC(b) (2.13 and 3.17 eV), IHFSCC(a) (2.11 and 3.11 eV), MRCI(a) (2.28 eV and 3.05 eV), and TDDFT/SAOP (2.14 and 2.38 eV), while agreement with SO-CASPT2 is less good (2.63 and 3.32 eV). While the agreement for the maxima of band A is rather good, for band B it is much less so. This could be due to the fact that this band is rather weak and only visible as a shoulder on the C band in the spectra reported by Gabes and Stufkens,⁷ making the exact position dependent on the fitting procedure employed to establish the band maximum.

2. Explorations of the dissociative region

The discussion so far has dealt with geometries at and close to the ground state equilibrium structures. However, in order to be useful in modeling the dissociation process of I_3^- the potential surfaces of the excited states far from the equilibrium region have to be properly described as well. In order to probe the relative accuracy of the methods, we have chosen to investigate displacements along the symmetric stretch coordinate, along which photodissociation via the C and D bands generally initiates. This represents a one-dimensional cut through the full three-dimensional surface enabling us to define local minima (and harmonic frequencies) for the different excited states in this restricted geometry. One should thereby keep in mind, however, that such extrema do not necessarily correspond to the true spectroscopic constants for the corresponding states as we did not investigate the curvature of the surface in directions orthogonal to the symmetric stretch coordinate.

A further restriction in the current study is the presence of intruder states that made it impossible to converge the IHFSCC calculations at internuclear distances larger than 3.6 Å. This means that we can only compare with the IHFSCC(b) reference at relatively short distances. Similarly, we could not obtain TDDFT/SAOP results beyond about 3.9 Å due to triplet instabilities. In the CASPT2 calculations we observe downward spikes in the potential curves, visible in Fig. 2, appear around 3.5 Å, signaling a weak intruder state problem. These features appear already in the spin-orbit free CASPT2 calculations, whereas the CASSCF calculations produce smooth potential curves. We were unable to curb these features, applying real⁷⁴ and imaginary⁷⁵ level shifts of up to 0.20 E_h . On the other hand, outside the spikes the potential curves appear smooth and so a pragmatic approach to fitting the potential surfaces would be to remove points for which the weight of the reference CASSCF state drops below a selected value since this characterizes the appearance of the spikes. A more satisfying solution would be to increase the active space, but this easily brings the computational cost out of practical reach.

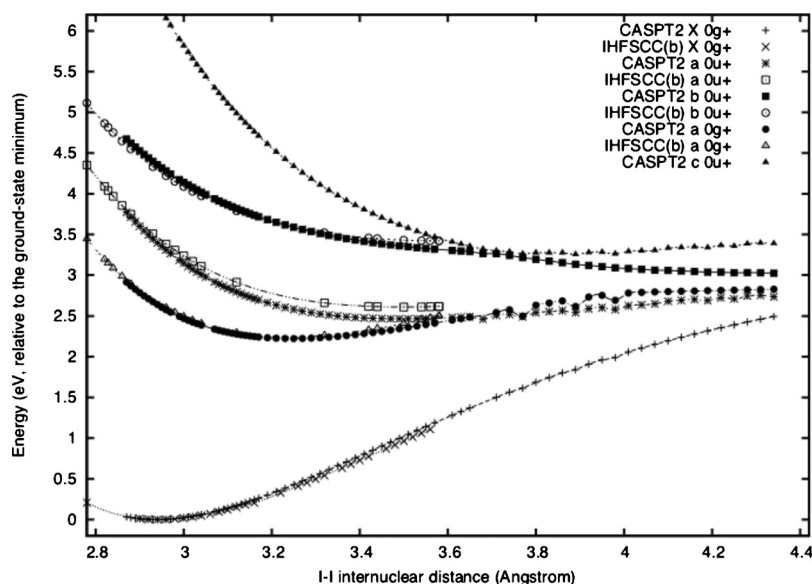


FIG. 2. Potential energy curves along the symmetric stretch coordinate for the ground state (GS, 0_g^+) and first excited 0_g^+ state, and lower 0_u^+ states (the first two corresponding to the optically active states) from CASPT2 and IHFSCC(b) calculations. All potentials were scaled so that the ground state minimum corresponds to the origin; the CASPT2 excited states were further shifted (a) downward by 0.31 eV and (b) toward larger internuclear distances by 0.05 Å, to compensate the systematic errors compared to IHFSCC(b).

Potential curves along the symmetric stretch for the lower two 0_g^+ and three 0_u^+ states, calculated at the CASPT2 and IHFSCC(b) level, are given in Fig. 2. The crossing of the first excited 0_g^+ and 0_u^+ states around 3.7 Å translate into a conical intersection in a 2-dimensional (2D) plot involving also the antisymmetric stretch and plays an important role in the photodissociation via the C band, as discussed by Nakanishi *et al.*¹⁸ These authors indicate a number of avoided crossings and conical intersections appearing in their SOCI calculations, such as an avoided crossing between the second and third 0_u^+ state, clearly visible around 3.7 Å in Fig. 2, which renders the second 0_u^+ state, associated with the D band, dissociative. Due to these features in the potential curves we have restricted the calculation of spectroscopic constants and adiabatic excitation energies to the ten lowest excited states of I_3^+ , presented in Table IV.

The lower excited states are dominated by single excita-

tions to the antibonding $\sigma_{1/2u}^*$ LUMO and bond lengths are therefore systematically considerably longer than observed for the ground state. This also implies, as already mentioned, that vertical excitations probe the repulsive wall of the excited states. We note, for instance, that the vertical and adiabatic excitation energies to the first excited 0_u^+ , calculated at the IHFSCC(b) level, are 3.36 and 2.61 eV, respectively, a difference of 0.75 eV, meaning that the photodissociation process starts off with a significant acceleration of the vibrational wave packet.

With respect to the spectroscopic constants given in Table IV, we observe excellent agreement between IHFSCC(a) and IHFSCC(b), again suggesting that correlation of the 4d shell is not crucial for a good description of the excited states potential surfaces. CASPT2 tends to consistently underestimate the bond lengths compared to IHFSCC(b) by about 0.07 Å, in line with the difference of 0.058 Å for the

TABLE IV. Spectroscopic constants (r_e in Å and ω_e in cm^{-1}) and adiabatic (T_e) excitation energies (in eV) for the ten lowest excited states of I_3^+ , calculated with the TDDFT/SAOP, MRCI(a), CASPT2, IHFSCC(a), and IHFSCC(b) methods. The lowest, optically active 0_u^+ state is shown in boldface. Statistical measures of the error compared to IHFSCC(b) are also shown (see text for discussion). Values in parenthesis have reduced accuracy due to triplet instabilities (TDDFT) or lack of convergence (IHFSCC) and have not been included in the statistical analysis.

State	TDDFT/SAOP				MRCI(a)				CASPT2				IHFSCC(a)				IHFSCC(b)			
	Ω	r_e	ω_e	T_e	Ω	r_e	ω_e	T_e	Ω	r_e	ω_e	T_e	Ω	r_e	ω_e	T_e	Ω	r_e	ω_e	T_e
1	2_g	3.443	55	1.19	0_u^-	3.501	56	1.53	2_g	3.188	80	1.97	0_u^-	3.439	60	1.45	0_u^-	3.417	60	1.45
2	1_g	3.462	52	1.29	1_u	3.503	56	1.53	1_u	3.318	71	1.98	1_u	3.436	61	1.46	1_u	3.414	61	1.46
3	0_u^-	(3.719)	(22)	(1.38)	2_g	3.383	61	1.64	0_u^-	3.315	72	1.99	2_g	3.270	81	1.59	2_g	3.245	83	1.59
4	1_u	(3.696)	(29)	(1.38)	1_g	3.421	56	1.71	1_g	3.208	76	2.02	1_g	3.282	78	1.70	1_g	3.255	80	1.71
5	2_u	3.616	46	1.41	2_u	3.644	47	1.89	1_u	3.427	59	2.33	0_g^-	3.383	52	2.05	0_g^-	3.358	53	2.07
6	1_u	3.669	39	1.52	1_u	3.657	47	1.89	2_u	3.417	60	2.35	1_g	1_g	(3.514)	(29)	(2.20)
7	0_g^+	3.388	63	1.68	0_g^-	3.478	48	2.16	0_g^-	3.222	73	2.44	2_u	3.475	63	2.18	2_u	3.446	64	2.22
8	1_g	(3.877)	(54)	(1.81)	1_g	3.562	43	2.22	1_g	3.260	68	2.47	0_g^+	3.264	82	2.22	0_g^+	3.240	84	2.24
9	0_g^-	3.554	38	1.88	0_g^+	3.388	62	2.28	0_g^+	3.198	78	2.53	1_u	3.492	62	2.21	1_u	3.462	63	2.26
10	0_u^+	3.513	57	2.10	0_u^+	3.731	34	2.36	0_u^+	3.445	53	2.78	0_u^+	0_u^+	3.498	50	2.61
$\bar{\Delta}$		0.163	-18	-0.52		0.152	-15	-0.07		-0.066	3	0.31		0.025	-1	-0.02				
Δ_{std}		0.106	20	0.18		0.051	8	0.19		0.036	9	0.16		0.003	1	0.02				
$\bar{\Delta}_{\text{abs}}$		0.163	23	0.52		0.152	15	0.14		0.066	7	0.31		0.025	1	0.02				
Δ_{max}		0.314	-46	-0.81		0.233	-24	-0.37		-0.136	20	0.54		0.030	-2	-0.05				

TABLE V. Character of 0_u^+ states of I_3^- for TDDFT/SAOP, IHFSCC(b), CASPT2, and MRCI(a) at the respective ground state equilibrium geometries. These are given in terms of contributions from one-electron (for TDDFT, IHFSCC, and MRCI, where spin-orbit coupling is included at the SCF level) or many-electron states (for CASPT2). The oscillator strengths f associated to the transition from the ground state to these excited states are also shown (where available) along with the vertical excitation energy. The experimental absorption maxima are found at (a) 3.43 and 4.25 eV from gas-phase photofragment yield spectra (Ref. 18), (b) 3.41 and 4.22 eV in CH_2Cl_2 (solution) (Ref. 7), and (c) 3.38 and 4.28 eV in Et_4NI_3 (solid state) (Ref. 7).

Method	State	Excited state composition	T_v (eV)	f
TDDFT/SAOP	0_u^+	$51\% \sigma_{1/2g} \rightarrow \sigma_{1/2u}^* + 49\% \pi_{1/2g} \rightarrow \sigma_{1/2u}^*$	2.82	0.09
	0_u^+	$45\% \sigma_{1/2g} \rightarrow \sigma_{1/2u}^* + 50\% \pi_{1/2g} \rightarrow \sigma_{1/2u}^*$	4.29	1.54
MRCI	0_u^+	$46\% \sigma_{1/2g} \rightarrow \sigma_{1/2u}^* + 30\% \pi_{1/2g} \rightarrow \sigma_{1/2u}^* + 7\% (\pi_{1/2g}^1 \pi_{1/2u}^1 \rightarrow \sigma_{1/2u}^{*2})$	3.47	0.37
	0_u^+	$29\% \sigma_{1/2g} \rightarrow \sigma_{1/2u}^* + 47\% \pi_{1/2g} \rightarrow \sigma_{1/2u}^* + 6\% (\sigma_{1/2g}^1 \sigma_{1/2u}^1 \rightarrow \sigma_{1/2u}^{*2})$	4.61	1.83
IHFSCC	0_u^+	$40\% \sigma_{1/2g} \rightarrow \sigma_{1/2u}^* (\text{LUMO}) + 19\% \pi_{1/2g} \rightarrow \sigma_{1/2u}^* (\text{LUMO}) + 22\% \sigma_{1/2g} \rightarrow \sigma_{1/2u}^* + 15\% \pi_{1/2g} \rightarrow \sigma_{1/2u}^*$	3.36	n/a ^a
	0_u^+	$26\% \sigma_{1/2g} \rightarrow \sigma_{1/2u}^* (\text{LUMO}) + 36\% \pi_{1/2g} \rightarrow \sigma_{1/2u}^* (\text{LUMO}) + 8\% \sigma_{1/2g} \rightarrow \sigma_{1/2u}^* + 24\% \pi_{1/2g} \rightarrow \sigma_{1/2u}^*$	4.27	n/a ^a
CASPT2	0_u^+	$14\% {}^1\Sigma_u^+ (98\% \sigma_g \rightarrow \sigma_u^*) + 85\% {}^3\Pi_u (91\% \pi_g \rightarrow \sigma_u^*)$	3.72	0.45
	0_u^+	$84\% {}^1\Sigma_u^+ (98\% \sigma_g \rightarrow \sigma_u^*) + 15\% {}^3\Pi_u (91\% \pi_g \rightarrow \sigma_u^*)$	4.67	2.37

^aNot available.

ground state. On the other hand, MRCI(a) significantly overestimates the excited state bond lengths, approaching the errors of TDDFT/SAOP ($\bar{\Delta} \approx 0.16 \text{ \AA}$), although with lower standard deviation. The same tendency is observed for harmonic frequencies, where CASPT2 on the other hand shows rather good performance.

It is interesting then to note that MRCI(a) outperforms CASPT2 for adiabatic excitation energies, whereas TDDFT/SAOP severely underestimates them. CASPT2 has, on the other hand, a reasonably small and consistent standard deviation of around 0.16 eV in excitation energies. This systematic nature of the errors can therefore perhaps be exploited to further bring the results close to IHFSCC(b) by applying global shifts to the CASPT2 potential energy surfaces. Such an approximation is demonstrated in Fig. 2, where we show, for the optically active 0_u^+ states, the IHFSCC(b) cuts of the potential energy surface along corresponding cuts of the CASPT2 potential, after the CASPT2 bond lengths were uniformly shifted by 0.035 \AA and a shift of 0.15 eV was applied to the excited states. After these two corrections, we observe a very good agreement between the two methods, with rather small discrepancies for the second 0_u^+ state. Agreement for the first 0_u^+ state, however, seems to be somewhat poorer than for the second.

3. A closer look at the 0_u^+ states

The two strong, allowed transitions for I_3^- occur from the $\Omega=0_g^+$ ground state to two states with $\Omega=0_u^+$. Table V displays the dominant contributions to these transitions. If available we also show oscillator strengths for the corresponding transitions. Since the CASPT2 calculations were done in a two-step fashion we can also analyze the composition of these excited states in terms of spin-orbit free states. To a first approximation one can view these states as ${}^3\Pi_u$ and ${}^1\Sigma_u^+$, the former borrowing intensity from the latter through the spin-orbit coupling. This picture correlates very well with the DIM-SO results of Kosloff and co-workers.¹⁹ In an orbital

picture all methods give the same description, with different mixings of the two transitions from the occupied σ_g and π_g orbitals to the σ_u LUMO. In the IHFSCC(b) calculations two additional excitations appear that can be interpreted as providing orbital relaxation of the LUMO. Since the DFT calculations also include spin-orbit coupling at the SCF stage we see a similar picture as for IHFSCC and MRCI, except that relaxation effects do not play an important role (since the virtual orbitals see the same potential as the occupied ones).

Nakanishi *et al.*¹⁸ reported calculated squared transition moments (μ^2) 7.31 and 13.51 a_0^2 for the two lower 0_u^+ states, respectively, which translates into oscillator strengths of 0.68 and 1.56, thus giving somewhat more relative weight to the lower state than what we observe at the CASPT2 and IHFSCC(b) level and shown in Table V. Gabes and Stufkens⁷ reported an intensity ratio of around 1.6:1.0 between the upper and lower 0_u^+ states. Since CASPT2, MRCI, and TDDFT/SAOP all overestimate the peak separation between these states and thus underestimate the spin-orbit coupling between the underlying ${}^3\Pi_u$ and ${}^1\Sigma_u^+$ Λ S states, these methods are also unable to reproduce the experimental intensity ratio. IHFSCC does get the peak separation right, but unfortunately we do not have access to oscillator strengths for this methodology.

C. Benchmark calculations on I_3^- : Electronic spectra and electron affinities

Given that the IHFSCC results for the $(1h, 0p)$ sector are generated as by-products of the excitation energies determination for I_3^- , we can also present spectroscopic constants and vertical and adiabatic excitation energies for this species, calculated with IHFSCC(b). We furthermore provide a comparison of this radical with IHFSCC, CASPT2, and MRCI at $r = 2.84 \text{ \AA}$, which is the geometry employed by Kosloff and co-workers¹⁹ for the same system. These results are also shown in Table VI.

TABLE VI. Comparison of excitation energies T_v (in eV) for different methods [CASPT2, IHFSCC(a,b)], and MRCI(a) for neutral I_3 at $r_1=r_2=2.84 \text{ \AA}$. Vertical (T_v) and adiabatic (T_e) excitation energies, as well as EAs are shown for IHFSCC(b).

State (Ω)	MRCI ^a	MRCI(a)	CASPT2	IHFSCC(a)	Expt.			IHFSCC(b)					
		$T_v(r_1=r_2=2.84 \text{ \AA})$			ω_e	T_e	EA ^b	R_e	ω_e	T_v ^c	T_v	T_e	EA ^d
X ($3/2_u$)	0.00	0.00	0.00	0.00	115 ± 5	0.00	4.15 ± 0.12	2.828	132	0.00	0.00	0.00	4.29
A ($1/2_g$)	0.28	0.33	0.42	0.27		0.28		2.884	113	0.24	0.25	0.23	
B ($1/2_u$)	0.61	0.65	0.55	0.64		0.62		2.837	129	0.65	0.65	0.65	
C ($3/2_g$)	0.78	0.78	0.77	0.88		0.68		2.951	115	0.86	0.88	0.78	
D ($1/2_g$)	1.62	1.63	1.55	1.68				2.933	118	1.68	1.70	1.62	

^aResults from Ref. 19 at $r_1=r_2=2.836 \text{ \AA}$.

^bReference 76; vertical detachment energy of I_3^- : 4.25 eV (Ref. 34).

^cCalculated at $r_1=r_2=2.84 \text{ \AA}$.

^dValue for the adiabatic EA. Vertical processes: $IP(I_3^-)=4.39 \text{ eV}$ and $EA(I_3)=4.20 \text{ eV}$.

The excitation energies calculated at $r=2.84 \text{ \AA}$ are in good agreement with those of Kosloff *et al.*, and, for the first two excited states, also with the experimental data. The $\Omega = 3/2_g$ third excited state has a much longer bond length than the ground state which may explain the 0.1–0.2 eV overestimation found relative to the experimental value when calculated as vertical excitation. The adiabatic results for IHFSCC(b) indeed show a decrease of 0.10 eV from the vertical to the adiabatic excitation energy.

In line with the discussion above, we see that the results obtained with the different methods are fairly consistent. We find a ground state with $\Omega=3/2_u$, followed by states of $\Omega = 1/2_g$, $1/2_u$, $3/2_g$, and $1/2_g$, respectively, in agreement with the experimental assignment.^{33,34}

The CASPT2 excitation energies are lower than those of IHFSCC(b) by about 0.1–0.2 eV with the exception of the first excited state, which is higher for CASPT2. In this case, however, MRCI and IHFSCC are much more alike than for I_3^- , with discrepancies generally smaller than 0.1 eV. This strong similarity could be due to the fact that for the ($1h, 0p$) sector used here the exponential parametrization of IHFSCC for the excited states is truncated at the linear term and therefore is essentially the same as in MRCI.^{62–65}

From the IHFSCC(b) calculations we furthermore obtain an adiabatic (vertical) electron affinity (EA) of 4.29 eV (4.20 eV) which compares well with the adiabatic EA of $4.15 \pm 0.12 \text{ eV}$, including zero-point vibrational corrections, obtained experimentally from a thermodynamic cycle,⁷⁶ as well as the vertical detachment energy of I_3^- of 4.25 eV reported by Choi *et al.*³⁴

IV. CONCLUSIONS

We have performed correlated electronic structure calculations including spin-orbit effects at high levels of theory on the triiodide ion (I_3^-) and the radical (I_3). The agreement between the different wave function-based methodologies employed is reasonable, as is their agreement with experimental results. Exploratory TDDFT calculations with the SAOP functional provide a qualitatively correct picture not too far from the equilibrium distance, but show unsystematic errors that prohibit use in quantitative description. The SAOP functional was chosen on the one hand because it has been constructed to correct the wrong asymptotic behavior of

pure exchange-correlation potentials, and on the other hand because it has shown good performance for excitation energies of molecules containing heavy elements.^{77–79} We believe that the principal reason for the relatively poor performance of TDDFT/SAOP is self-interaction errors.⁸⁰ This is indicated by the severe overestimation of the equilibrium bond lengths and the observation that the introduction of exact Hartree–Fock exchange through hybrid functionals gives spectroscopic constants in much better agreement with experiment. We have not undertaken a systematic study of the performance of other DFT functionals. Whereas there is a good reason to believe that some functionals will provide better vertical excitation energies at the equilibrium distance, we suspect that they will all encounter stability problems along the dissociation channels.

With respect to I_3^- , of the different methodologies evaluated, the intermediate Hamiltonian Fock-space coupled cluster is clearly the method that most accurately and consistently provides a picture which is both qualitatively and quantitatively correct for the excitation processes taking place in the initial steps in the photodissociation of the triiodide ion and the triiodide radical. Due to convergence problems this method is unfortunately not applicable to the complete potential energy surface, but we have been able to show that other wave function-based methods can reproduce the benchmark IHFSCC results rather well.

While it can be argued that for the triiodide species MRCI is slightly more accurate than SO-CASPT2, the latter has two interesting advantages: For one thing, it is computationally much more efficient than MRCI, and for another its errors seem to be very systematic for all electronic states considered. This systematic nature of errors is observed, in particular, for energies, allowing for global correction to be applied to the potential energy surfaces in order to bring them in agreement with IHFSCC. All is not well, however, with CASPT2. We observe a strong sensitivity of the method to symmetry breaking. A simulation of the photodissociation of I_3^- will require as a minimum the generation of the 2D potential surfaces of linear I_3^- . However, in order to avoid symmetry breaking such a scan can at best be carried out in C_2 symmetry, for which the distinction between 0^+ and 0^- states is lost and needs to be recovered by the use of scripts when analyzing the individual states. CASPT2 is furthermore susceptible to intruder states, and we have been unable

to remove the spikes observed in the potential curves toward dissociation by the use of level shifting techniques. One alternative is to explore a possibly more robust multireference perturbation theory such as *n*-electron valence second-order perturbation theory.⁸¹

MRCI clearly is another option, but presently too expensive to be able to generate a full set of potential energy surfaces for studies of the dynamics of the photodissociation process. The challenge of applying such methods to heavy-element systems is not only the mandatory treatment of spin-orbit coupling, but also the fact that heavy atoms are increasingly polarizable such that subvalence has to be correlated. In the present case we observe that the correlation of 4d is crucial for obtaining correct spectroscopic constants, less so for vertical excitation energies. Subvalence correlation can possibly be avoided by employing relativistic effective core potentials combined with core-polarization potentials.^{82,83} Finally, we would like to suggest that a successful combination of MCSCF and DFT (Ref. 84) could be the ideal tool for generating potential surfaces for dynamics.

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