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# Initial value problems for diffusion equations with singular potential 

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#### Abstract

Let $V$ be a nonnegative locally bounded function defined in $Q_{\infty}:=\mathbb{R}^{n} \times(0, \infty)$. We study under what conditions on $V$ and on a Radon measure $\mu$ in $\mathbb{R}^{d}$ does it exist a function which satisfies $\partial_{t} u-\Delta u+V u=0$ in $Q_{\infty}$ and $u(., 0)=\mu$. We prove the existence of a subcritical case in which any measure is admissible and a supercritical case where capacitary conditions are needed. We obtain a general representation theorem of positive solutions when $t V(x, t)$ is bounded and we prove the existence of an initial trace in the class of outer regular Borel measures.


## Contents

1 Introduction ..... 2
2 The subcritical case ..... 6
3 The supercritical case ..... 12
3.1 Capacities ..... 12
3.2 The singular set of $V$ ..... 13
3.3 Properties of positive solutions and representation formula ..... 15
3.4 Reduced measures ..... 21
4 Initial trace ..... 23
4.1 The direct method ..... 23
4.2 The sweeping method ..... 27

[^0]
## 5 Appendix: the case of a bounded domain

5.1 The subcritical case . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 28
5.2 The supercritical case . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . . 30

## 1 Introduction

In this article we study the initial value problem for the heat equation

$$
\begin{align*}
\partial_{t} u-\Delta u+V(x, t) u & =0 & & \text { in } Q_{T}:=\mathbb{R}^{n} \times(0, T)  \tag{1.1}\\
u(., 0) & =\mu & & \text { in } \mathbb{R}^{n},
\end{align*}
$$

where $V \in L_{\text {loc }}^{\infty}\left(Q_{T}\right)$ is a nonnegative function and $\mu$ a Radon measure in $\mathbb{R}^{n}$. By a (weak) solution of (1.1) we mean a function $u \in L_{l o c}^{1}\left(\bar{Q}_{T}\right)$ such that $V u \in L_{l o c}^{1}\left(\bar{Q}_{T}\right)$, satisfying

$$
\begin{equation*}
-\iint_{Q_{T}}\left(\partial_{t} \phi+\Delta \phi\right) u d x d t+\iint_{Q_{T}} V u \phi d x d t=\int_{\Omega} \zeta d \mu \tag{1.2}
\end{equation*}
$$

for every function $\zeta \in C_{c}^{1,1 ; 1}\left(\bar{Q}_{T}\right)$ which vanishes for $t=T$. Besides the singularity of the potential at $t=0$, there are two main difficulties which appear for constructing weak solutions : the growth of the measure at infinity and the concentration of the measure near some points in $\mathbb{R}^{n}$.

We denote by $H(x, t)=\left(\frac{1}{4 \pi t}\right)^{\frac{n}{2}} e^{-\frac{|x|^{2}}{4 t}}$ the Gaussian kernel in $\mathbb{R}^{n}$ and by $\mathbb{H}[\mu]$ the corresponding heat potential of a measure $\mu \in \mathfrak{M}\left(\mathbb{R}^{n}\right)$. Thus

$$
\begin{equation*}
\mathbb{H}[\mu](x, t)=\left(\frac{1}{4 \pi t}\right)^{\frac{n}{2}} \int e^{-\frac{|x-y|^{2}}{4 t}} d \mu(y), \tag{1.3}
\end{equation*}
$$

whenever this expression have a meaning: for example it is straightforward that if $\mu \in \mathfrak{M}\left(\mathbb{R}^{n}\right)$ satisfies

$$
\begin{equation*}
\|\mu\|_{\mathfrak{M}_{T}}:=\int_{\mathbb{R}^{n}} e^{-\frac{|y|^{2}}{4 T}} d|\mu|(y)<\infty \tag{1.4}
\end{equation*}
$$

then (1.3) has a meaning as long as $t<T$, and let be $\mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$ the set of Radon measures in $\mathbb{R}^{n}$ satisfying (1.4). If $G \subset \mathbb{R}^{n}$, let $Q_{T}^{G}$ be the cylinder $G \times(0, T), B_{R}(x)$ the ball of center $x$ and radius $R$ and $B_{R}=B_{R}(0)$. We prove
Theorem A Let the measure $\mu$ verifies

$$
\begin{equation*}
\iint_{Q_{T}^{B_{R}}} \mathbb{H}[|\mu|](x, t) V(x, t) d x d t \leq M_{R} \quad \forall R>0 \tag{1.5}
\end{equation*}
$$

Then (1.1) admits a solution in $Q_{T}$.
A measure which satisfies (1.5) is called an admissible measure and a measure for which there exists a solution to problem (1.1) is called a good measure. Notice that even when $V=0$, uniqueness without any restriction on $u$ is not true, however the next uniqueness result holds:
Theorem B Let $u$ be a weak solution of (1.1) with $\mu=0$. If $u$ satisfies

$$
\begin{equation*}
\iint_{Q_{T}}(1+V(x, t)) e^{-\lambda|x|^{2}}|u(x, t)| d x d t<\infty \tag{1.6}
\end{equation*}
$$

for some $\lambda>0$, then $u=0$.
We denote by $\mathcal{E}_{\nu}\left(Q_{T}\right)$ the set of functions $u \in L_{l o c}^{1}\left(Q_{T}\right)$ for which (1.6) holds for some $\lambda>0$. The general result we prove is the following.
Theorem C Let $\mu \in \mathfrak{M}\left(\mathbb{R}^{n}\right)$ be an admissible measure satisfying (1.4). Then there exists a unique solution $u_{\mu} \in \mathcal{E}_{\nu}\left(Q_{T}\right)$ to problem (1.1). Furthermore

$$
\begin{equation*}
\iint_{Q_{T}}\left(\frac{n}{2 T}+V\right)|u| e^{-\frac{|x|^{2}}{4(T-t)}} d x d t \leq \int_{\mathbb{R}^{n}} e^{-\frac{|y|^{2}}{4 T}} d|\mu|(y) \tag{1.7}
\end{equation*}
$$

We consider first the subcritical case, which means that any positive measure satisfying (1.4) is a good measure and we prove that such is the case if for any $R>0$ there exist $m_{R}>0$ such that

$$
\begin{equation*}
\iint_{Q_{T}^{B_{R}}} H(x-y, t) V(x, t) d x d t \leq m_{R} e^{-\frac{|y|^{2}}{4 T}} . \tag{1.8}
\end{equation*}
$$

Moreover we prove a stability result among the measures satisfying (1.4): if $V$ verifies for all $R>0$

$$
\begin{equation*}
\sup _{y \in \mathbb{R}^{n}} e^{\frac{|y|^{2}}{4 T}} \iint_{E} H(x-y, t) V(x, t) d x d t \rightarrow 0 \quad \text { when }|E| \rightarrow 0, E \text { Borel subset of } Q_{T}^{B_{R}} \tag{1.9}
\end{equation*}
$$

then if $\left\{\mu_{k}\right\}$ is a sequence of Radon measures bounded in $\mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$ which converges in the weak sense of measures to $\mu$, then $\left\{\left(u_{\mu_{k}}, V u_{\mu_{k}}\right)\right\}$ converges to $\left(u_{\mu}, V u_{\mu}\right)$ in $L_{l o c}^{1}\left(\bar{Q}_{T}\right)$.

In the supercritical case, that is when not any measure in $\mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$ is a good measure, we develop a capacitary framework in order to characterize the good measures. We denote by $\mathfrak{M}^{V}\left(\mathbb{R}^{n}\right)$ the set of Radon measures such that $V \mathbb{H}[\mu] \in L^{1}\left(Q_{T}\right)$ and $\|\mu\|_{\mathfrak{M}^{V}}:=\|V \mathbb{H}[\mu]\|_{L^{1}}$. If $E \subset Q_{T}$ is a Borel set, we set

$$
\begin{equation*}
C_{V}(E)=\sup \left\{\mu(E): \mu \in \mathfrak{M}_{+}^{V}\left(\mathbb{R}^{n}\right), \mu\left(E^{c}\right)=0,\|\mu\|_{\mathfrak{M}^{V}} \leq 1\right\} \tag{1.10}
\end{equation*}
$$

This defines a capacity. If

$$
\begin{equation*}
C_{V}^{*}(E)=\inf \left\{\|f\|_{L^{\infty}}: \check{H}[f](y) \geq 1 \quad \forall y \in E\right\} \tag{1.11}
\end{equation*}
$$

where

$$
\begin{equation*}
\check{H}[f](y)=\iint_{Q_{T}} H(x-y, t) V(x, t) f(x, t) d x d t=\int_{0}^{T} \mathbb{H}[V f](y, t) d t \quad \forall y \in \mathbb{R}^{n}, \tag{1.12}
\end{equation*}
$$

then $C_{V}^{*}(E)=C_{V}(E)$ for any compact set. Denote by $Z_{V}$ the singular set of $V$, that is the largest set with zero $C_{V}$ capacity. Then

$$
\begin{equation*}
Z_{V}=\left\{x \in \mathbb{R}^{n}: \iint_{Q_{T}} H(x-y, t) V(y, t) d x d t=\infty\right\} \tag{1.13}
\end{equation*}
$$

and the following result characterizes the good measures.
Theorem D If $\mu$ is an admissible measure then $\mu\left(Z_{V}\right)=0$. If $\mu \in \mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$ satisfies $\mu\left(Z_{V}\right)=0$, then it is a good measure. Furthermore $\mu$ is a positive good measure if and only if there exists
an increasing sequence of positive admissible measures $\left\{\mu_{k}\right\}$ which converges to $\mu$ in the weak * topology.

Since many important applications deal with the nonlinear equation

$$
\begin{equation*}
\partial_{t} u-\Delta u+|u|^{q-1} u=0 \quad \text { in } Q_{\infty}:=\mathbb{R}^{n} \times(0, \infty) \tag{1.14}
\end{equation*}
$$

where $q>1$ and due to the fact that any solution defined in $Q_{\infty}$ satisfies

$$
\begin{equation*}
|u(x, t)|^{q-1} \leq \frac{1}{t(q-1)} \quad \forall(x, t) \in Q_{\infty}, \tag{1.15}
\end{equation*}
$$

we shall concentrate on potentials $V$ which satisfy

$$
\begin{equation*}
0 \leq V(x, t) \leq \frac{C_{1}}{t} \quad \forall(x, t) \in Q_{T} \tag{1.16}
\end{equation*}
$$

for some $C_{1}>0$. For such potentials we prove the existence of a representation theorem for positive solutions of

$$
\begin{equation*}
\partial_{t} u-\Delta u+V(x, t) u=0 \quad \text { in } Q_{T} \tag{1.17}
\end{equation*}
$$

If $u$ is a positive solution of (1.1) in $Q_{T}$ with $\mu \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right)$, it is the increasing limit of the solutions $u=u_{R}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+V(x, t) u & =0 & & \text { in } Q_{T}^{B_{R}} \\
u & =0 & & \text { in } \partial B_{R} \times(0, T)  \tag{1.18}\\
u(., 0) & =\chi_{B_{R}} \mu & & \text { in } B_{R},
\end{align*}
$$

when $R \rightarrow \infty$, thus there exists a positive fonction $H_{V} \in C\left(\mathbb{R}^{n} \times \mathbb{R}^{n} \times(0, T)\right)$ such that

$$
\begin{equation*}
u(x, t)=\int_{\mathbb{R}^{N}} H_{V}(x, y, t) d \mu(y) \tag{1.19}
\end{equation*}
$$

Furthermore we show how to construct $H_{V}$ from $V$ and we prove the following formula

$$
\begin{equation*}
H_{V}(x, y, t)=e^{\psi(x, t)} \Gamma(x, y, t) \tag{1.20}
\end{equation*}
$$

where

$$
\begin{equation*}
\psi(x, t)=\int_{t}^{T} \int_{\mathbb{R}^{n}}\left(\frac{1}{4 \pi(s-t)}\right)^{\frac{n}{2}} e^{-\frac{|x-y|^{2}}{4(s-t)}} V(y, s) d y d s \tag{1.21}
\end{equation*}
$$

and $\Gamma$ satisfies the following estimate

$$
\begin{equation*}
c_{1} t^{-\frac{n}{2}} e^{-\gamma_{1} \frac{|x-y|^{2}}{t}} \leq \Gamma(x, y, t) \leq c_{2} t^{-\frac{n}{2}} e^{-\gamma_{2} \frac{|x-y|^{2}}{t}} \tag{1.22}
\end{equation*}
$$

where $A_{i}, c_{i}$ depends on $T, d$ and $V$. Conversely, we first prove the following representation result Theorem E Assume $V$ satisfies (1.16). If $u$ is a positive solution of (1.1) in $Q_{T}$, there exists a positive Radon measure $\mu$ in $\mathbb{R}^{n}$ such that (1.19) holds.

If $\mu \in \mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$ is positive, we can define for any $k>0$ the solution $u_{k}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+V_{k}(x, t) u & =0 & & \text { in } Q_{T}  \tag{1.23}\\
u(., 0) & =\mu & & \text { in } \mathbb{R}^{n}
\end{align*}
$$

where $V_{k}(x, t)=\min \{k, V(x, t)\}$, and

$$
\begin{equation*}
u_{k}(x, t)=\int_{\mathbb{R}^{N}} H_{V_{k}}(x, y, t) d \mu(y) \tag{1.24}
\end{equation*}
$$

Moreover $\left\{H_{V_{k}}\right\}$ and $\left\{v_{k}\right\}$ decrease respectively to $H_{V}$ and $u^{*}$ there holds

$$
\begin{equation*}
u^{*}(x, t)=\int_{\mathbb{R}^{N}} H_{V}(x, y, t) d \mu(y) \tag{1.25}
\end{equation*}
$$

However $u^{*}$ is not a solution of (1.1), but of a relaxed problem where $\mu$ is replaced by a smaller measure $\mu^{*}$ called the reduced measure associated to $\mu$. If we define the zero set of $V$ by

$$
\begin{equation*}
\mathcal{S i n g}_{V}:=\left\{y \in \mathbb{R}^{N}: H_{V}(x, y, t)=0\right\}, \tag{1.26}
\end{equation*}
$$

we prove
Theorem $\mathbf{F}$ Assume $V$ satisfies (1.15) and $\mu \in \mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$. Then
(i) $\operatorname{supp}\left(\mu-\mu^{*}\right) \subset \operatorname{Sing}_{V}$.
(ii) If $\mu\left(\operatorname{Sing}_{V}\right)=0$, then $\mu^{*}=0$.
(iii) $\operatorname{Sing}_{V} \subset Z_{V}$.

The last section is devoted to the initial trace problem: to any positive solution $u$ of (1.1) we can associate an open subset $\mathcal{R}(u) \subset \mathbb{R}^{n}$ which is the set of points $y$ which possesses a neighborhood $U$ such that

$$
\begin{equation*}
\iint_{Q_{T}^{U}} V(x, t) u(x, t) d x d t<\infty . \tag{1.27}
\end{equation*}
$$

There exists a positive Radon measure $\mu_{u}$ on $\mathcal{R}(u)$ such that

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathbb{R}^{n}} u(x, t) \zeta(x) d x=\int_{\mathbb{R}^{n}} \zeta d \mu \quad \forall \zeta \in C_{c}(\mathcal{R}(u)) \tag{1.28}
\end{equation*}
$$

The set $\mathcal{S}(u)=\mathbb{R}^{n} \backslash \mathcal{R}(u)$ is the set of points $y$ such that for any open set $U$ containing $y$, there holds

$$
\begin{equation*}
\iint_{Q_{T}^{U}} V(x, t) u(x, t) d x d t=\infty \tag{1.29}
\end{equation*}
$$

If $V$ satisfies (1.17), $\mathcal{S}(u)$ it has the property that

$$
\begin{equation*}
\limsup \sin _{t \rightarrow 0} \int_{U} u(x, t) d x=\infty \tag{1.30}
\end{equation*}
$$

Furthermore, if is satisfies (1.9), then $\mathcal{S}(u)=\emptyset$.
An alternative construction of the initial trace based on the sweeping method is also developed.

## 2 The subcritical case

Let $Q_{T}=\mathbb{R}^{n} \times(0, T]$. In this section we consider the linear parabolic problem

$$
\begin{align*}
\partial_{t} u-\Delta u+V u & =0 & & \text { in } \quad Q_{T} \\
u(., 0) & =\mu & & \text { in } \mathbb{R}^{n} \times\{0\},
\end{align*}
$$

where $V \in L_{l o c}^{1}\left(Q_{T}\right)$ is nonnegative and $\mu$ is a Radon measure.
Definition 2.1 We say that $\mu \in \mathfrak{M}\left(\mathbb{R}^{n}\right)$ is a good measure if problem (2.1) has a weak solution $u$ i.e. there exists a function $u \in L_{l o c}^{1}\left(\bar{Q}_{T}\right)$, such that $V u \in L_{l o c}^{1}\left(\bar{Q}_{T}\right)$ which satisfies

$$
\begin{equation*}
-\iint_{Q_{T}} u\left(\partial_{t} \phi+\Delta \phi\right) d x d t+\iint_{Q_{T}} V u \phi d x d t=\int_{\mathbb{R}^{n}} \phi(x, 0) d \mu \quad \forall \phi \in X\left(Q_{T}\right), \tag{2.2}
\end{equation*}
$$

where $X\left(Q_{T}\right)$ is the space of test functions defined by

$$
X\left(Q_{T}\right)=\left\{\phi \in C_{c}\left(\bar{Q}_{T}\right), \partial_{t} \phi+\Delta \phi \in L_{l o c}^{\infty}\left(\bar{Q}_{\infty}\right), \phi(x, T)=0\right\}
$$

Definition 2.2 Let $H(x, t)$ be the heat kernel of heat equation in $\mathbb{R}^{n}$, we say that $\mu \in \mathfrak{M}\left(\mathbb{R}^{n}\right)$ is an admissible measure if

$$
\begin{equation*}
\|V \mathbb{H}[|\mu|] \mid\|_{L^{1}\left(Q_{T}^{B_{R}}\right)}=\iint_{Q_{T}^{B_{R}}}\left(\int_{\mathbb{R}^{n}} H(x-y, t) d|\mu(y)|\right) V(x, t) d x d t<M_{R, T} \tag{i}
\end{equation*}
$$

where $M_{R, T}$ is a positive constant.
Definition 2.3 $A$ function $u(x, t)$ will be said to belong to the class $\mathcal{E}_{\mathcal{V}}\left(Q_{T}\right)$ if there exists $\lambda>0$ such that

$$
\iint_{Q_{T}} e^{-\lambda|x|^{2}}|u(x, t)|(1+V(x, t)) d x d t<\infty .
$$

A measure in $\mathbb{R}^{n}$ belongs to the class $\mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$ if

$$
\|\mu\|_{\mathfrak{M}_{T}}:=\int_{\mathbb{R}^{n}} e^{-\frac{|x|^{2}}{4 T}} d|\mu|<\infty
$$

Lemma 2.4 There exists at most one weak solution of problem (2.1) in the class $\mathcal{E}_{\mathcal{V}}\left(Q_{T}\right)$.
Proof. Let $u_{1}$ and $u_{2}$ be two solutions in the class $\mathcal{E}_{\mathcal{V}}\left(Q_{T}\right)$ then $w=u_{1}-u_{2}$ is a solution with initial data 0 . Choose a standard mollifier $\rho: B(0,1) \mapsto[0,1]$ and define

$$
w_{j}(x, t)=j^{n} \int_{B_{\frac{1}{j}}(x)} \rho(j(x-y)) w(y, t) d y \equiv \int_{B_{\frac{1}{j}}(x)} \rho_{j}(x-y) w(y, t) d y
$$

Then $w_{j}(., t)$ is $C^{\infty}$ and from the equation satisfied by $w$, it holds

$$
\partial_{t} w_{j}-\Delta w_{j}+\int_{B_{\frac{1}{j}}(x)} V(y, t) \rho_{j}(x-y) w(y) d y=0,
$$

where $\partial_{t} w_{j}$ is taken in the weak sense.
First we consider the case $\lambda>0$ and $t \leq \min \left\{\frac{1}{16 \lambda}, T\right\}=T^{\prime}$.
Set $\phi(x, t)=\xi(x, t) \zeta(x)$, where $\xi(x, t)=e^{-\frac{|x|^{2}}{4\left(\frac{1}{8 \lambda}-t\right)}}$ and $\zeta \in C_{c}^{\infty}\left(\mathbb{R}^{n}\right)$. Given $\varepsilon>0$ we define

$$
g_{j}=\sqrt{w_{j}^{2}+\varepsilon}
$$

Because $\partial_{t}\left(g_{j} \phi\right)=\frac{\partial_{t} w_{j}}{\sqrt{w_{j}^{2}+\varepsilon}} \phi+g_{j} \partial_{t} \phi$,, by a straightforward calculation we have

$$
\begin{aligned}
\int_{\mathbb{R}^{n}}\left[g_{j} \phi(., s)\right]_{s=t}^{s=0} d x & =\iint_{Q_{t}} \frac{w_{j}}{\sqrt{w_{j}^{2}+\varepsilon}} \phi \Delta w_{j} d x d s \\
& -\iint_{Q_{t}} \frac{w_{j}(x, s)}{\sqrt{w_{j}^{2}(x, s)+\varepsilon}} \phi(x, s)\left(\int_{B_{\frac{1}{j}}(x)} V(y, t) \rho_{j}(x-y) w(y, s) d y\right) d x d s \\
& +\iint_{Q_{t}} g_{j} \phi_{s} d x d s \\
& =I_{1}+I_{2}+I_{3} .
\end{aligned}
$$

By integration by parts, we obtain

$$
\begin{aligned}
I_{1} & =-\iint_{Q_{t}} \frac{\left|\nabla w_{j}\right|^{2}}{\sqrt{w_{j}^{2}(x, s)+\varepsilon}} \phi d x d s+\iint_{Q_{t}} \frac{\left|\nabla w_{j}\right|^{2} w_{j}^{2}}{\left(w_{j}^{2}(x, s)+\varepsilon\right)^{\frac{3}{2}}} \phi d x d s-\iint_{Q_{t}} \frac{w_{j}}{\sqrt{w_{j}^{2}+\varepsilon}} \nabla w_{j} \cdot \nabla \phi d x d s \\
& \leq-\iint_{Q_{t}} \frac{w_{j}}{\sqrt{w_{j}^{2}+\varepsilon}} \nabla w_{j} \cdot \nabla \phi d x d s \\
& \leq-\iint_{Q_{t}} \nabla g_{j} \cdot \nabla \phi d x d s \\
& =-\iint_{Q_{t}} \zeta \nabla g_{j} \cdot \nabla \xi d x d s-\iint_{Q_{t}} \xi \nabla g_{j} \cdot \nabla \zeta d x d s \\
& =\iint_{Q_{t}} \zeta g_{j} \Delta \xi d x d s+\iint_{Q_{t}} g_{j} \nabla \zeta \cdot \nabla \xi d x d s .
\end{aligned}
$$

Since $t \leq T$, there holds $\xi\left|\nabla g_{j}\right| \in L^{1}\left(Q_{T^{\prime}}\right), \xi g_{j} \in L^{1}\left(Q_{T^{\prime}}\right),|\Delta \xi| g_{j} \in L^{1}\left(Q_{T^{\prime}}\right), \partial_{s} \xi g_{j} \in L^{1}\left(Q_{T^{\prime}}\right)$ and

$$
\iint_{Q_{t}} \frac{w_{j}(x, s)}{\sqrt{w_{j}^{2}(x, s)+\varepsilon}}\left(\int_{B_{\frac{1}{j}}(x)} V(y, t) \rho_{j}(x-y) w(y, s) d y\right) \xi d x d s<\infty
$$

The reason for which $\xi\left|\nabla g_{j}\right| \in L^{1}\left(Q_{T^{\prime}}\right)$ follows from the next inequality

$$
\begin{aligned}
\iint_{Q_{T^{\prime}}}\left|\nabla g_{j}\right| \xi d x d s & =\iint_{Q_{T^{\prime}}} \frac{\left|\nabla w_{j}\right|}{\sqrt{\varepsilon+w_{j}^{2}}} \xi d x d s \\
& \leq \iint_{Q_{T^{\prime}}} e^{-\frac{|x|^{2}}{4\left(\frac{1}{8 \lambda}-t\right)}}\left(\int_{B_{\frac{1}{j}}(x)}\left|\nabla \rho_{j}(x-y)\right| w(y, s) d y\right) d x d s .
\end{aligned}
$$

Since $\forall y \in B_{\frac{1}{j}}(x)$, we have $|x|^{2} \geq\left(|y|^{2}-\frac{1}{j}\right)^{2}=|y|^{2}+\frac{1}{j^{2}}-2 \frac{|y|}{j} \geq \frac{|y|^{2}}{2}-(C-1) \frac{1}{j^{2}}$, for some positive constant $C>0$ independent on $j, y$ and $x$. Thus we have, using the fact that $e^{-\lambda|y|^{2}} w \in L^{1}\left(Q_{T}\right)$,

$$
\iint_{Q_{T^{\prime}}}\left|\nabla g_{j}\right| \xi d x d s \leq C(j, \lambda) \iint_{Q_{T^{\prime}}} \int_{B_{\frac{1}{j}}(x)} e^{-\frac{|y|^{2}}{8\left(\frac{y}{8 \lambda}-t\right)}}\left|\nabla \rho_{j}(x-y)\right| w(y, s) d y d x d s<\infty
$$

Also
$\iint_{Q_{t}} \frac{w_{j}(x, s)}{\sqrt{w_{j}^{2}(x, s)+\varepsilon}} \xi\left(\int_{B_{\frac{1}{j}}(x)} V(y, t) \rho_{j}(x-y) w(y, s) d y\right) d x d s \rightarrow_{j \rightarrow \infty} \iint_{Q_{t}} \frac{w^{2}(x, s)}{\sqrt{w^{2}(x, s)+\varepsilon}} \xi V(y, t) d x d s$
and

$$
\int_{\mathbb{R}^{n}} \sqrt{w_{j}^{2}(x, s)+\varepsilon}\left(\xi_{s}+\Delta \xi\right) d x d s \rightarrow_{j \rightarrow \infty} \int_{\mathbb{R}^{n}} \sqrt{w^{2}(x, s)+\varepsilon}\left(\xi_{s}+\Delta \xi\right) d x d s
$$

We choose $\zeta_{R}=1$ in $B_{R}, 0 \leq \zeta_{R} \leq 1$ in $B_{R+1} \backslash B_{R}$ and 0 otherwise. Letting successively $j \rightarrow \infty$, $R \rightarrow \infty$ and finally $\varepsilon \rightarrow 0$, we derive

$$
\int_{\mathbb{R}^{n}}|w(x, t)| \xi(x, t) d x \leq \iint_{Q_{t}}|w|\left(\xi_{s}+\Delta \xi\right) d x d s-\iint_{Q_{t}} w(x, s) \xi V(y, t) d x d s
$$

Since

$$
\xi_{s}+\Delta \xi=-\frac{n}{2\left(\frac{1}{8 \lambda}-s\right)}
$$

and $V \geq 0$, we have $w(x, t)=0 \forall(x, t) \in Q_{T^{\prime}}$. If $T^{\prime}=T$ this complete the proof for $\lambda \geq 0$, otherwise the proof can be completed by a finite number of interations of the same argument on $\mathrm{R}^{n} \times\left(T^{\prime}, 2 T^{\prime}\right), \mathrm{R}^{n} \times\left(2 T^{\prime}, 3 T^{\prime}\right)$, etc. If $\lambda=0$ we set $\xi=1$ and the result follows by similar argument

Theorem 2.5 If $\mu \in \mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$ is an admissible measure, there exists a unique $u=u_{\mu} \in \mathcal{E}_{\mathcal{V}}\left(Q_{T}\right)$ solution of (2.1). Furthermore the following estimate holds

$$
\begin{equation*}
\frac{n}{2 T} \iint_{Q_{T}}|u| e^{-\frac{|x|^{2}}{4(T-t)}} d x d s+\iint_{Q_{T}}|u| V e^{-\frac{|x|^{2}}{4(T-t)}} d x d s \leq \int_{\mathbb{R}^{n}} e^{-\frac{|x|^{2}}{4 T}} d|\mu| \tag{2.3}
\end{equation*}
$$

Proof. First we assume that $\mu \geq 0$. Let $\mu_{R}=\chi_{B_{R}} \mu$. It is well known that the heat kernel $H^{B_{R}}(x, y, t)$ in $\Omega=B_{R}$ is increasing with respect to $R$ and $H^{B_{R}} \rightarrow H$, as $R \rightarrow \infty$ in $L^{1}\left(Q_{T}\right)$ for any $T>0$. Thus $\mu_{R}$ is an admissible measure in $B_{R}$ and by Proposition 5.4, there exists a unique weak solution $u_{R}$ of problem 5.2 on $\Omega=B_{R}$. By (ii) of Proposition 5.5 we have

$$
-\iint_{Q_{T}}\left|u_{R}\right|\left(\partial_{t} \phi+\Delta \phi\right) d x d t+\iint_{Q_{T}}\left|u_{R}\right| V \phi d x d t \leq \int_{B_{R}} \phi(x, 0) d\left|\mu_{R}\right|
$$

If we set $\phi_{\varepsilon}(x, t)=e^{-\frac{|x|^{2}}{4(T+\varepsilon-t)}} ; \varepsilon>0$, then

$$
\partial_{t} \phi+\Delta \phi=-\frac{n}{2(T+\varepsilon-t)} e^{-\frac{|x|^{2}}{4(T+\varepsilon-t)}}
$$

thus we have

$$
\iint_{Q_{T}}\left|u_{R}\right| \frac{n}{2(T+\varepsilon-t)} e^{-\frac{|x|^{2}}{4(T+\varepsilon-t)}} d x d t+\iint_{Q_{T}}\left|u_{R}\right| V e^{-\frac{|x|^{2}}{4(T+\varepsilon-t)}} d x d t \leq \int_{B_{R}} e^{\frac{-|x|^{2}}{4 T+4 \varepsilon}} d \mu_{R},
$$

which implies

$$
\frac{n}{2 T+\varepsilon} \int_{0}^{T} \int_{B_{R}}\left|u_{R}\right| e^{-\frac{|x|^{2}}{4(T+\varepsilon-t)}} d x d t+\int_{0}^{T} \int_{B_{R}}\left|u_{R}\right| V e^{-\frac{|x|^{2}}{4(T+\varepsilon-t)}} d x d t \leq \int_{\mathbb{R}^{n}} e^{\frac{-|x|^{2}}{4 T+4 \varepsilon}} d \mu_{R}
$$

Letting $\varepsilon \rightarrow 0$, we derive

$$
\frac{n}{2 T} \iint_{Q_{T}}\left|u_{R}\right| e^{-\frac{|x|^{2}}{4(T-t)}} d x d t+\iint_{Q_{T}}\left|u_{R}\right| V e^{-\frac{|x|^{2}}{4(T-t)}} d x d t \leq \int_{\mathbb{R}^{n}} e^{\frac{-|x|^{2}}{4 T}} d \mu_{R} \leq \int_{\mathbb{R}^{n}} e^{\frac{-|x|^{2}}{4 T}} d \mu
$$

Now by the maximum principle $\left\{u_{R}\right\}$ is increasing with respect to $R$ and converges to some function $u$. By the above inequality $u \in \mathcal{E}_{\mathcal{V}}\left(Q_{T}\right)$ satisfies the estimate (2.5) and $u$ is a weak solution of problem (2.1). By Lemma 2.4 it is unique. In the general case we write $\mu=\mu^{+}-\mu^{-}$and the result follows by the above arguments and Lemma 2.4. In the sequel we shall denote by $u_{\mu}$ this unique solution.

Definition 2.6 $A$ potential $V$ is called subcritical in $Q_{T}$ if for any $R>0$ there exists $m_{R}>0$ such that

$$
\begin{equation*}
\iint_{Q_{T}^{B_{R}}} H(x-y, t) V(x, t) d x d t \leq m_{R} e^{-\frac{|y|^{2}}{4 T}} \quad \forall y \in \mathbb{R}^{n} . \tag{2.4}
\end{equation*}
$$

It is called strongly subcritical if moreover

$$
\begin{equation*}
e^{\frac{|y|^{2}}{4 T}} \iint_{E} H(x-y, t) V(x, t) d x d t \rightarrow 0 \quad \text { when }|E| \rightarrow 0, E \text { Borel subset of } Q_{T}^{B_{R}} \tag{2.5}
\end{equation*}
$$

uniformly with respect to $y \in \mathbb{R}^{n}$
Theorem 2.7 Assume $V$ is subcritical. Then any measure in $\mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$ is admissible. Furthermore, if $V$ is strongly subcritical and $\left\{\mu_{k}\right\}$ is a sequence of measures uniformly bounded in $\mathfrak{M}_{T}\left(\mathbb{R}^{N}\right)$ which converges weakly to $\mu$, then the corresponding solutions $\left\{u_{\mu_{k}}\right\}$ converge to $u_{\mu}$ in $L_{l o c}^{1}\left(\bar{Q}_{T}\right)$, and $\left\{V u_{\mu_{k}}\right\}$ converges to $V u_{\mu}$ in $L_{l o c}^{1}\left(\bar{Q}_{T}\right)$.

Proof. For the first statement we can assume $\mu \geq 0$ and there holds

$$
\begin{aligned}
\iint_{Q_{T}^{B_{R}}} H(x-y, t) d \mu(y) V(x, t) d x d t & =\int_{\mathbb{R}^{n}}\left(\iint_{Q_{T}^{B_{R}}} H(t, x-y) V(x, t) d x d t\right) d \mu(y) \\
& \leq m_{R} \int_{\mathbb{R}^{n}} e^{-\frac{|y|^{2}}{4 T}} d \mu(y) \\
& \leq m_{R}\|\mu\|_{\mathfrak{M}_{T}} .
\end{aligned}
$$

Thus $\mu$ is admissible. For the second statement, we assume first that $\mu_{k} \geq 0$. By lower semicontinuity $\mu \in \mathfrak{M}_{T}\left(\mathbb{R}^{N}\right)$ and $\|V \mathbb{H}[\mu]\|_{L^{1}\left(Q_{T}^{B_{R}}\right)} \leq M_{R, T}$ for any $k$. Since $0 \leq u_{\mu_{k}} \leq \mathbb{H}\left[\mu_{k}\right]$ and
$\mathbb{H}\left[\mu_{k}\right] \rightarrow \mathbb{H}[\mu]$ in $L_{l o c}^{1}\left(\bar{Q}_{T}\right)$, the sequence $\left\{u_{\mu_{k}}\right\}$ is uniformly integrable and thus relatively compact in $L_{l o c}^{1}\left(\bar{Q}_{T}\right)$. Furthermore $0 \leq V u_{\mu_{k}} \leq V \mathbb{H}\left[\mu_{k}\right]$. Let $E \subset Q_{T}^{B_{R}}$ be a Borel subset, then

$$
\begin{aligned}
\iint_{E} V \mathbb{H}\left[\mu_{k}\right] d x d t & =\int_{\mathbb{R}^{n}}\left(\iint_{E} V H(x-y, t) d x d t\right) d \mu_{k}(y) \\
& =\int_{\mathbb{R}^{n}}\left(e^{\frac{|y|^{2}}{4 T}} \iint_{E} V(x) H(x-y, t) d x d t\right) e^{-\frac{|y|^{2}}{4 T}} d \mu_{k}(y) \\
& \leq \epsilon(|E|)\left\|\mu_{k}\right\|_{\mathfrak{M}_{T}}
\end{aligned}
$$

where $\epsilon(r) \rightarrow 0$ as $r \rightarrow 0$. Thus $\left\{\left(u_{\mu_{k}}, V u_{\mu_{k}}\right)\right\}$ is locally compact in $L_{l o c}^{1}\left(\bar{Q}_{T}\right)$ and, using a diagonal sequence, there exist $u \in L_{l o c}^{1}\left(\bar{Q}_{T}\right)$ with $V u \in L_{l o c}^{1}\left(\bar{Q}_{T}\right)$ and a subsequence $\left\{k_{j}\right\}$ such that $\left\{\left(u_{\mu_{k_{j}}}, V u_{\mu_{k_{j}}}\right)\right\}$ converges to $\left(u_{\mu}, V u_{\mu}\right)$ a.e. and in $L_{l o c}^{1}\left(\bar{Q}_{T}\right)$. From the integral expression (2.2) satisfied by the $u_{\mu_{k}}, u$ is a weak solution of problem (2.1). Since the $u_{\mu_{k}}$ satisfy (2.3), the property holds for $u$, therefore $u=u_{\mu}$ is the unique solution of (2.1), which ends the proof.

As a variant of the above result which will be useful later on we have
Proposition 2.8 Assume $V$ satisfies

$$
\begin{equation*}
e^{\frac{|y|^{2}}{4 T}} \iint_{E} H(x-y, t) V(x, t+\tau) d x d t \rightarrow 0 \quad \text { when }|E| \rightarrow 0, E \text { Borel subset of } Q_{T}^{B_{R}} \tag{2.6}
\end{equation*}
$$

uniformly with respect to $y \in \mathbb{R}^{n}$ and $\tau \in\left[0, \tau_{0}\right]$. Let $\tau_{k}>0$ with $\tau_{k} \rightarrow 0$ and $\left\{\mu_{k}\right\}$ be a sequence uniformly bounded in $\mathfrak{M}_{T}\left(\mathbb{R}^{N}\right)$ which converges weakly to $\mu$. Then the solutions $\left\{u_{\tau_{k}, \mu_{k}}\right\}$ of

$$
\begin{align*}
\partial_{t} u-\Delta u+V u & =0 & & \text { on } \tag{2.7}
\end{align*} \quad \mathbb{R}^{n} \times\left(\tau_{k}, T\right)
$$

(extended by 0 on $\left(0, \tau_{k}\right)$ ) converge to $u_{\mu}$ in $L_{l o c}^{1}\left(Q_{T}\right)$, and $\left\{V u_{\mu_{k}}\right\}$ converges to $V u_{\mu}$ in $L_{l o c}^{1}\left(Q_{T}\right)$.
Condition (2.5) may be very difficult to verify and we give below a sufficient condition for it to hold.

Proposition 2.9 Assume $V$ satisfies

$$
\begin{equation*}
\lim _{\lambda \rightarrow 0} e^{\frac{|y|^{2}}{4 T}} \lambda^{-n} \int_{0}^{\lambda} \int_{B_{\lambda^{2}}(y)} V(x, t) d x d t=0 \tag{2.8}
\end{equation*}
$$

uniformly with respect to $y \in \mathbb{R}^{n}$, then $V$ is strongly subcritical.
Proof. Let $E \subset Q_{T}^{B_{R}}$ be a Borel set. For $\delta>0$, we define the weighted heat ball of amplitude $\delta e^{-\frac{|y|^{2}}{4 T}}$ by

$$
P_{\delta}=P_{\delta}(y, T)=\left\{(x, t) \in Q_{T}: H(x-y, t) \geq \delta e^{-\frac{|y|^{2}}{4 T}}\right\}
$$

By an straightforward computation, one sees that

$$
P_{\delta}(y, T) \subset B_{a_{n} \delta^{-\frac{1}{n}} e^{\frac{|y|^{2}}{4 n T}}}(y) \times\left[0, b_{n} \delta^{-\frac{2}{n}} e^{\frac{|y|^{2}}{2 n T}}\right]:=R_{\delta}(y, T)
$$

for some $a_{n}, b_{n}>0$. We write

$$
\iint_{E} H(x-y, t) V(x, t) d x d t=\iint_{E \cap P_{\delta}} H(x-y, t) V(x, t) d x d t+\iint_{E \cap P_{\delta}^{c}} H(x-y, t) V(x, t) d x d t
$$

Then

$$
\iint_{E \cap P_{\delta}^{c}} H(x-y, t) V(x, t) d x d t \leq \delta e^{-\frac{|y|^{2}}{4 T}} \iint_{E} V(x, t) d x d t
$$

and

$$
\begin{aligned}
& \iint_{E \cap P_{\delta}} H(x-y, t) V(x, t) d x d t \leq \int_{0}^{\delta} \int_{\left\{(x, t) \in Q_{T}^{B_{R}}: H(x-y, t)=\tau e^{\left.-\frac{|y|^{2}}{4 T}\right\}}\right.} V(x, t) d S_{\tau}(x, t) \tau d \tau \\
& \leq {\left[\tau \int_{0}^{\tau} \int_{\left\{(x, t) \in Q_{T}^{B_{R}}: H(x-y, t)=\sigma e^{\left.-\frac{|y|^{2}}{4 T}\right\}}\right.} V(x, t) d S_{\sigma}(x, t) d \sigma\right]_{\tau=0}^{\tau=\delta} } \\
&-\int_{0}^{\delta} \int_{0}^{\tau} \int_{\left\{(x, t) \in Q_{T}^{B_{R}}: H(x-y, t)=\sigma e^{\left.-\frac{|y|^{2}}{4 T}\right\}}\right.} V(x, t) d S_{\sigma}(x, t) d \sigma d \tau \\
& \leq \delta \int_{0}^{\delta} \int_{\left\{(x, t) \in Q_{T}^{B_{R}}: H(x-y, t)=\sigma e^{\left.-\frac{|y|^{2}}{4 T}\right\}}\right.} V(x, t) d S_{\sigma}(x, t) d \sigma .
\end{aligned}
$$

The first integration by parts is justified since $V \in L^{1}\left(Q_{T}^{B_{R}}\right)$. Notice that

$$
\delta \int_{0}^{\delta} \int_{\left\{(x, t) \in Q_{T}^{B_{R}}: H(x-y, t)=\sigma e^{\left.-\frac{|y|^{2}}{4 T}\right\}}\right.} V(x, t) d S_{\sigma}(x, t) d \sigma=\delta \iint_{Q_{T}^{B_{R} \cap P_{\delta}}} V(x, t) d x d t
$$

and

$$
\begin{aligned}
\delta \iint_{Q_{T}^{B_{R}} \cap P_{\delta}} V(x, t) d x d t & \leq \delta \iint_{Q_{T}^{B_{R} \cap R_{\delta}(y, T)}} V(x, t) d x d t \\
& \leq \beta r^{-n} \int_{0}^{\alpha r} \int_{B_{R} \cap B_{(\alpha r)^{2}}(y)} V(x, t) d x d t,
\end{aligned}
$$

for some $\alpha, \beta>0$ and if we have set $r=\delta^{-\frac{1}{n}}$. Notice also that $B_{R} \cap B_{(\alpha r)^{2}}(y)=\emptyset$ if $|y| \geq R+(\alpha r)^{2}$, or, equivalently, if $|y| \geq R+\alpha^{2} \delta^{-\frac{2}{n}}$.
(i) If $|y| \geq R+\alpha^{2}$, we fix $\delta$ such that $1<\delta$, then

$$
e^{\frac{|y|^{2}}{4 T}} \iint_{E} H(x-y, t) V(x, t) d x d t \leq \delta \iint_{E} V(x, t) d x d t
$$

which can be made smaller than $\epsilon$ provided $|E|$ is small enough.
(ii) If $|y|<R+\alpha^{2}$, then

$$
\begin{aligned}
e^{\frac{|y|^{2}}{4 T}} \iint_{E \cap P_{\delta}^{c}} H(x-y, t) V(x, t) d x d t & \leq e^{\frac{R^{2}+\alpha^{4}}{2 T}} \iint_{E \cap P_{\delta}^{c}} H(x-y, t) V(x, t) d x d t \\
& \leq \delta e^{\frac{R^{2}+\alpha^{4}}{2 T}} \iint_{E} V(x, t) d x d t .
\end{aligned}
$$

Given $\epsilon>0$, we fix $\delta=r^{-n}$ such that

$$
e^{\frac{R^{2}+\alpha^{4}}{2 T}} \iint_{E \cap P_{\delta}} H(x-y, t) V(x, t) d x d t \leq \beta e^{\frac{R^{2}+\alpha^{4}}{2 T}} r^{-n} \int_{0}^{\alpha r} \int_{B_{R} \cap B_{(\alpha r)^{2}(y)}} V(x, t) d x d t \leq \frac{\epsilon}{2}
$$

and then $\eta>0$ such that $|E| \leq \eta$ implies

$$
e^{\frac{|y|^{2}}{4 T}} \iint_{E \cap P_{\delta}^{c}} H(x-y, t) V(x, t) d x d t \leq \delta e^{\frac{R^{2}+\alpha^{4}}{2 T}} \iint_{E} V(x, t) d x d t \leq \frac{\epsilon}{2} .
$$

Therefore

$$
e^{\frac{|y|^{2}}{4 T}} \iint_{E} H(x-y, t) V(x, t) d x d t \leq \epsilon,
$$

which is (2.5).
Remark In Theorem 2.7 and Proposition 2.9, the assumption of uniformity with respect to $y \in \mathbb{R}^{n}$ in (2.5), (2.6) and (2.8) can be replaced by uniformity with respect to $y \in B_{R_{0}}$ if all the measures $\mu_{k}$ have their support in $B_{R_{0}}$. A extension of these assumptions, valid when the convergent measures $\mu_{k}$ have their support in a fixed compact set is to assume that $V$ is locally strongly subcritical, which means that (2.5) holds uniformly with respect to $y$ in a compact set. Similar extension holds for (2.8).

## 3 The supercritical case

### 3.1 Capacities

All the proofs in this subsection are similar to the ones of [12] and inspired by [6]; we omit them. We assume also that there exists a positive measure $\mu_{0}$ such that $\mathbb{H}\left[\mu_{0}\right] V \in L^{1}\left(Q_{T}\right)$.

Definition 3.1 If $\mu \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right)$ and $f$ is a nonnegative measurable function defined in $\Omega$ such that

$$
(t, x, y) \mapsto \mathbb{H}[\mu](y, t) V(x, t) f(x, t) \in L^{1}\left(Q_{T} \times \mathbb{R}^{n} ; d x d t \otimes d \mu\right)
$$

we set

$$
\mathcal{E}(f, \mu)=\int_{Q_{T}}\left(\int_{\mathbb{R}^{n}} H(x-y, t) d \mu(y)\right) V(x, t) f(x, t) d x d t
$$

If we put

$$
\breve{\mathbb{H}}[f](y)=\int_{Q_{T}} H(x-y, t) V(x, t) f(x, t) d x d t
$$

then by Fubini's Theorem, $\breve{H}[f](y)<\infty, \mu$-almost everywhere in $\mathbb{R}^{n}$ and

$$
\mathcal{E}(f, \mu)=\int_{\mathbb{R}^{n}}\left(\int_{Q_{T}} H(x-y, t) V(x, t) f(x, t) d x d t\right) d \mu(y)
$$

Proposition 3.2 Let $f$ be fixed. Then
(a) $y \mapsto \breve{H}[f](y)$ is lower semicontinuous in $\mathbb{R}^{n}$.
(b) $\mu \mapsto \mathcal{E}(f, \mu)$ is lower semicontinuous in $\mathfrak{M}_{+}\left(\mathbb{R}^{n}\right)$ in the weak* topology.

Definition 3.3 We denote by $\mathfrak{M}^{V}\left(\mathbb{R}^{n}\right)$ the set of all measures $\mu$ on $\mathbb{R}^{n}$ such that $V \mathbb{H}[|\mu|] \in$ $L^{1}\left(Q_{T}\right)$. If $\mu$ is such a measure, we set

$$
\|\mu\|_{\mathfrak{M}^{V}}=\int_{Q_{T}}\left(\int_{\mathbb{R}^{n}} H(x-y, t) d|\mu|(y)\right) V(x, t) d x d t=\|V \mathbb{H}[|\mu|]\|_{L^{1}\left(Q_{T}\right)}
$$

If $E \subset \mathbb{R}^{n}$ is a Borel set, we put

$$
\mathfrak{M}_{+}(E)=\left\{\mu \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right): \mu\left(E^{c}\right)=0\right\} \quad \text { and } \quad \mathfrak{M}_{+}^{V}(E)=\mathfrak{M}^{V}\left(\mathbb{R}^{n}\right) \cap \mathfrak{M}_{+}(E) .
$$

Definition 3.4 If $E \subset \mathbb{R}^{n}$ is any borel subset we define the set function $C_{V}$ by

$$
C_{V}(E):=\sup \left\{\mu(E): \mu \in \mathfrak{M}_{+}^{V}(E),\|\mu\|_{\mathfrak{M}^{v}} \leq 1\right\}
$$

this is equivalent to,

$$
C_{V}(E):=\sup \left\{\frac{\mu(E)}{\|\mu\|_{\mathfrak{M}^{V}}}: \mu \in \mathfrak{M}_{+}^{V}(E)\right\} .
$$

Proposition 3.5 The set function $C_{V}$ satisfies

$$
C_{V}(E) \leq \sup _{y \in E}\left(\int_{Q_{T}} H(x-y, t) V(x, t) d x d t\right)^{-1} \quad \forall E \subset \mathbb{R}^{n}, E \text { Borel. }
$$

Furthermore equality holds if $E$ is compact. Finally,

$$
C_{V}\left(E_{1} \cup E_{2}\right)=\sup \left\{C_{V}\left(E_{1}\right), C_{V}\left(E_{2}\right)\right\} \quad \forall E_{i} \subset \mathbb{R}^{n}, E_{i} \text { Borel. }
$$

Definition 3.6 For any Borel $E \subset \mathbb{R}^{n}$, we set

$$
C_{V}^{*}(E):=\inf \left\{\|f\|_{L^{\infty}}: \breve{\mathbb{H}}[f](y) \geq 1 \forall y \in E\right\}
$$

Proposition 3.7 For any compact set $E \subset \mathbb{R}^{n}$,

$$
C_{V}^{*}(E)=C_{V}(E)
$$

### 3.2 The singular set of $V$

In this section we assume that $V$ satisfies (1.16), although much weaker assumption could have been possible. We define the singular set of $V, Z_{V}$ by

$$
\begin{equation*}
Z_{V}=\left\{x \in \mathbb{R}^{n}: \iint_{Q_{T}} H(x-y, t) V(y, t) d y d t=\infty\right\} . \tag{3.1}
\end{equation*}
$$

Since the function $x \mapsto f(x)=\iint_{Q_{T}} H(x-y, t) V(y, t) d y d t$ is lower semicontinuous, it is a Borel function and $Z_{V}$ is a Borel set.

Lemma 3.8 If $x \in Z_{V}$ then for any $r>0$,

$$
\iint_{Q_{T}^{B_{r}(x)}} H(x-y, t) V(y, t) d y d t=\infty .
$$

Proof. We will prove it by contradiction, assuming that there exists $r>0$, such that

$$
\iint_{Q_{T}^{B r(x)}} H(x-y, t) V(y, t) d y \leq M
$$

Replacing $H$ by its value, we derive

$$
\begin{aligned}
\iint_{Q_{T}} H(x-y, t) V(y, t) d y d t & =\iint_{Q_{T}^{B_{r}(x)}} H(x-y, t) V(y, t) d y d t+\iint_{Q_{T}^{B_{r}^{C}(x)}} H(x-y, t) V(y, t) d y d t \\
& \leq M+C(n) \int_{0}^{T} t^{-\frac{n+2}{2}} e^{-\frac{r^{2}}{4 t}} d t<\infty .
\end{aligned}
$$

Which is clearly a contradiction.
Lemma 3.9 If $\mu$ is an admissible positive measure then $\mu\left(Z_{V}\right)=0$.
Proof. Let $K \subset Z_{V}$ be a compact set. In view of the above lemma there exists a $R>0$ such that $K \subset B_{R}$ and for each $x \in K$, we have

$$
\begin{equation*}
\iint_{Q_{T}^{B_{2 R}}} H(x-y, t) V(y) d y=\infty \tag{3.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\iint_{Q_{T}^{B_{2}^{c}}} H(x-y, t) V(y) d y<\infty . \tag{3.3}
\end{equation*}
$$

Now, $\mu_{K}=\chi_{K} \mu$ is an admissible measure and by Fubini theorem we have

$$
\begin{aligned}
\iint_{Q_{T}}\left(\int_{\mathbb{R}^{n}} H(x-y, t) d \mu_{K}(y)\right) V(x, t) d x d t= & \int_{K} \iint_{Q_{T}} H(x-y, t) V(x, t) d x d t d \mu(y) \\
= & \int_{K} \iint_{Q_{T}^{B_{2 R}}} H(x-y, t) V(x) d x d t d \mu(y) \\
& +\int_{K} \iint_{Q_{T}^{B_{2 R}^{c}}} H(x-y, t) V(x) d x d t d \mu y .
\end{aligned}
$$

By (3.3) the second integral above is finite and by (3.2)

$$
\iint_{Q_{T}^{B_{2 R}}} H(x-y, t) V(x) d x d t=\infty \quad \forall y \in K
$$

It follows that $\mu(K)=0$. This implies $\mu\left(Z_{V}\right)=0$ by regularity.
Theorem 3.10 If $\mu \in \mathfrak{M}_{T}\left(\mathbb{R}^{n}\right), \mu \geq 0$ such that $\mu\left(Z_{V}\right)=0$, then $\mu$ is a good measure.
Proof. We set $\mu_{R}=\chi_{B_{R}} \mu$. By Proposition 5.8, since $Z_{V}^{B_{R}} \subset Z_{V}, \mu_{R}$ is a good measure in $B_{R}$ with corresponding solution $u_{\mu}^{R}$. In view of Lemma 2.5, $u_{\mu}^{R}$ satisfies

$$
\iint_{Q_{T}^{B_{R}}}\left|u_{\mu}^{R}\right| \frac{n}{4(T-t)} e^{-\frac{|x|^{2}}{4(T-t)}} d x d t+\iint_{Q_{T}^{B_{R}}}\left|u_{\mu}^{R}\right| V e^{-\frac{|x|^{2}}{4(T-t)}} d x d t \leq \int_{B_{R}} e^{-\frac{|x|^{2}}{4 T}} d \mu .
$$

Also $\left\{u_{\mu}^{R}\right\}$ is an increasing function, thus converges to $u_{\mu}$. By the above estimate we have that $u_{\mu}$ belong to class $\mathcal{E}_{\mathcal{V}}\left(Q_{T}\right)$ and is a weak solution of (2.1).

Proposition 3.11 Let $\mu \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right)$. Then $\mu\left(Z_{V}\right)=0$ if and only if there exists an increasing sequence of positive admissible measures which converges to $\mu$ in the weak* topology.

Proof. The proof is similar as the one of [12, Th 3.11] and we present it for the sake of completeness. First, we assume that $\mu\left(Z_{V}\right)=0$. Then we define the set

$$
K_{N}=\left\{x \in \mathbb{R}^{n}: \int_{Q_{T}} H(x-y, t) V(y) d y d t \leq N\right\} .
$$

We note that $Z_{V} \cap K_{N}=\emptyset$. We set $\mu_{n}=\chi_{K_{N}} \mu$ then we have

$$
\int_{Q_{T}}\left(\int_{\mathbb{R}^{n}} H(x-y, t) d \mu_{n}(y)\right) V(x, t) d x d t \leq \mu\left(K_{N}\right)
$$

Thus $\mu_{n}$ is admissible, increasing with respect $n$. By the monotone theorem it follows that $\mu_{n} \rightarrow$ $\chi_{Z_{V}^{c}} \mu$. Since $\mu\left(Z_{V}\right)=0$ the result follows in this direction.
For the other direction. Let $\left\{\mu_{n}\right\}$ be an increasing sequence of positive admissible measure. Then by Lemma 3.9 we have that $\mu_{n}\left(Z_{V}\right)=0, \forall n \geq 1$. Since $\mu_{n} \leq \mu$, there exist an increasing functions $h_{n} \mu$-integrable such that $\mu_{n}=h_{n} \mu$. Since $0=\mu_{n}\left(Z_{V}\right) \rightarrow \mu\left(Z_{V}\right)$ the result follows.

### 3.3 Properties of positive solutions and representation formula

We first recall the construction of the kernel function for the operator $w \mapsto \partial_{t} w-\Delta w+V w$ in $Q_{T}$, always assuming that $V$ satisfies (1.16). For $\delta>0$ and $\mu \in \mathfrak{M}_{T}$, we denote by $w_{\delta}$ the solution of

$$
\begin{align*}
\partial_{t} w-\Delta w+V_{\delta} w & =0, & & \text { in } Q_{T} \\
w(., 0) & =\mu & & \text { in } \mathbb{R}^{n} . \tag{3.4}
\end{align*}
$$

where $V_{\delta}=V \chi_{Q_{\delta, T}}$ and $Q_{\delta, T}=(\delta, T) \times \mathbb{R}^{n}$. Then

$$
\begin{equation*}
w_{\delta}(x, t)=\int_{\mathbb{R}^{n}} H_{V_{\delta}}(x, y, t) d \mu(y) \tag{3.5}
\end{equation*}
$$

Lemma 3.12 The mapping $\delta \mapsto H_{V_{\delta}}(x, y, t)$ is increasing and converges to $H_{V} \in C\left(\mathbb{R}^{n} \times \mathbb{R}^{n} \times\right.$ $(0, T]))$ when $\delta \rightarrow 0$. Furthermore There exists a fonction $\left.H_{V} \in C\left(\mathbb{R}^{n} \times \mathbb{R}^{n} \times(0, T]\right)\right)$ such that for any $\mu \in \mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} w_{\delta}(x, t)=w(x, t)=\int_{\mathbb{R}^{n}} H_{V}(x, y, t) d \mu(y) \tag{3.6}
\end{equation*}
$$

Proof. Without loss of generality we can assume $\mu \geq 0$. By the maximum principle $\delta \mapsto H_{V_{\delta}}(x, y, t)$ is increasing and the result follows by the monotone convergence theorem.
If $\mathbb{R}^{n}$ is replaced by a smooth bounded domain $\Omega$, we can consider the problem

$$
\begin{align*}
\partial_{t} w-\Delta w+V_{\delta} w=0 & \text { in } Q_{T}^{\Omega} \\
w=0 & \text { in } \partial_{l} Q_{T}^{\Omega}:=\partial \Omega \times(0, T]  \tag{3.7}\\
w(., 0)=\mu & \text { in } \Omega .
\end{align*}
$$

where $V_{\delta}^{\prime}=V \chi_{Q_{\delta, T}^{\Omega}}$ and $Q_{\delta, T}^{\Omega}=(\delta, T) \times \Omega$. Then

$$
\begin{equation*}
w_{\delta}(x, t)=\int_{\Omega} H_{V_{\delta}}^{\Omega}(x, y, t) d \mu(y) \tag{3.8}
\end{equation*}
$$

The proof of the next result is straightforward.

Lemma 3.13 The mapping $\delta \mapsto H_{V_{\delta}}^{\Omega}(x, y, t)$ increases and converges to $\left.H_{V}^{\Omega} \in C(\Omega \times \Omega \times(0, T])\right)$ when $\delta \rightarrow 0$. Furthermore There exists a fonction $\left.H_{V}^{\Omega} \in C(\Omega \times \Omega \times(0, T])\right)$ such that for any $\mu \in \mathfrak{M}_{b}(\Omega)$

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} w_{\delta}(x, t)=w(x, t)=\int_{\Omega} H_{V}^{\Omega}(x, y, t) d \mu(y) \tag{3.9}
\end{equation*}
$$

Furthermore $H_{V}^{\Omega} \leq H_{V}^{\Omega^{\prime}} \leq H_{V}$ if $\Omega \subset \Omega^{\prime}$.
It is important to notice that the above results do not imply that $w$ is a weak solution of problem (1.1). This question will be considered later on with the notion of reduced measure.
Lemma 3.14 Assume $\mu \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right)$ is a good measure and let $u$ be a positive weak solution of problem (2.1). If $\Omega$ is a smooth bounded domain, then there exists a unique positive weak solution $v$ of problem

$$
\begin{align*}
\partial_{t} v-\Delta v+V v & =0, & & \text { in } Q_{T}^{\Omega}, \\
v & =0 & & \text { on } \partial_{l} Q_{T}^{\Omega}  \tag{3.10}\\
v(., 0) & =\chi_{\Omega} \mu & & \text { in } \Omega .
\end{align*}
$$

Furthermore

$$
\begin{equation*}
v(x, t)=\int_{\Omega} H_{V}^{\Omega}(x, y, t) d \mu(y) \tag{3.11}
\end{equation*}
$$

Proof. Let $\left\{t_{j}\right\}_{j=1}^{\infty}$ be a sequence decreasing to 0 , such that $t_{j}<T, \forall j \in \mathbb{N}$. We consider the following problem

$$
\begin{align*}
\partial_{t} v-\Delta v+V v & =0, & & \text { in } \quad \Omega \times\left(t_{j}, T\right] \\
v & =0 & & \text { on } \partial \Omega \times\left(t_{j}, T\right]  \tag{3.12}\\
v\left(., t_{j}\right) & =u\left(., t_{j}\right) & & \text { in } \Omega \times\left\{t_{j}\right\} .
\end{align*}
$$

Since $u, V u \in L^{1}\left(Q_{T}^{B_{R}}\right)$ for any $R>0, t \mapsto u(., t)$ is continuous with value in $L_{l o c}^{1}\left(\mathbb{R}^{n}\right)$, therefore $u\left(., t_{j}\right) \in L_{l o c}^{1}\left(\mathbb{R}^{n}\right)$ and there exists a unique solution $v_{j}$ to (3.12) (notice also that $V \in L^{\infty}\left(Q_{T}^{B_{R}}\right)$ ). By the maximum principle $0 \leq v_{j} \leq u$ and by standard parabolic estimates, we may assume that the sequence $v_{j}$ converges locally uniformly in $\Omega \times(0, T]$ to a function $v \leq u$. Also, if $\phi \in C^{1,1 ; 1}\left(\overline{Q_{T}^{\Omega}}\right)$ vanishes on $\partial_{l} Q_{T}^{\Omega}$ and satisfies $\phi(x, T)=0$, we have
$-\int_{t_{j}}^{T} \int_{\Omega} v_{j}\left(\partial_{t} \phi+\Delta \phi\right) d x d t+\int_{t_{j}}^{T} \int_{\Omega} V v_{j} \phi d x d t+\int_{\Omega} \phi\left(x, T-t_{j}\right) v_{j}\left(x, T-t_{j}\right) d x=\int_{\Omega} \phi(x, 0) u\left(x, t_{j}\right) d x$,
where in the above equality we have taken $\phi\left(., .-t_{j}\right)$ as test function. Since $\phi\left(., T-t_{j}\right) \rightarrow 0$ uniformly and $u\left(., t_{j}\right) \rightarrow \mu$ in the weak sense of measures, it follows by the dominated convergence theorem that

$$
-\iint_{Q_{T}^{\Omega}} v\left(\partial_{t} \phi+\Delta \phi\right) d x d t+\iint_{Q_{T}^{\Omega}} V v \phi d x d t=\int_{\Omega} \phi(y, 0) d \mu(y),
$$

thus $v$ is a weak solution of problem (3.10). Uniqueness follows as in Lemma 2.4. Finally, for $\delta>0$, we consider the solution $w_{\delta}$ of (3.7). Then it is expressed by (3.5). Furthermore

$$
-\iint_{Q_{T}^{\Omega}} w_{\delta}\left(\partial_{t} \phi+\Delta \phi\right) d x d t+\iint_{Q_{T}^{\Omega}} V_{\delta} w_{\delta} \phi d x d t=\int_{\Omega} \phi(x, 0) d \mu(x)
$$

The sequence $w_{\delta}$ is decreasing, with limit $w$. Since $w_{\delta} \geq v$, then $w \geq v$. If we assume $\phi \geq 0$, it follows from dominated convergence and Fatou's lemma that

$$
-\iint_{Q_{T}^{\Omega}} w\left(\partial_{t} \phi+\Delta \phi\right) d x d t+\iint_{Q_{T}^{\Omega}} V w \phi d x d t \leq \int_{\Omega} \phi(x, 0) d \mu(x)
$$

Thus $w$ is a subsolution for problem (3.10) for which we have comparison when existence. Finally $w=v$ and (3.11) holds.

Lemma 3.15 Assume $\mu \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right)$ is a good measure and let $u$ be a positive weak solution of problem (2.1). Then for any $(x, t) \in \mathbb{R}^{n} \times(0, T]$, we have

$$
\lim _{R \rightarrow \infty} u_{R}=u
$$

where $\left\{u_{R}\right\}$ is the increasing sequence of the weak solutions of the problem (3.10) with $\Omega=B_{R}$. Moreover, the convergence is uniform in any compact subset of $\mathbb{R}^{n} \times(0, T]$ and we have the representation formula

$$
u(x, t)=\int_{\mathbb{R}^{n}} H_{V}(x, y, t) d \mu(y)
$$

Proof. By the maximum principle, $u_{R} \leq u_{R^{\prime}} \leq u$ for any $0<R \leq R^{\prime}$. Thus $u_{R} \rightarrow w \leq u$. Also by standard parabolic estimates, this convergence is locally uniformly. Now by dominated convergence theorem, it follows that $w$ is a weak solution of problem (2.1) with initial data $\mu$. Now we set $\widetilde{w}=u-w \geq 0$. Since $\widetilde{w}$ satisfies in the weak sense

$$
\begin{aligned}
\widetilde{w}_{t}-\Delta \widetilde{w}+V \widetilde{w}=0 & \text { in } Q_{T} \\
\widetilde{w}(x, t) \geq 0 & \text { in } Q_{T} \\
\widetilde{w}(x, 0)=0 & \text { in } \mathbb{R}^{n}
\end{aligned}
$$

and $V \geq 0$, it clearly satisfies

$$
\begin{aligned}
\widetilde{w}_{t}-\Delta \widetilde{w} \leq 0 & \text { in } Q_{T} \\
\widetilde{w}(x, t) \geq 0 & \text { in } Q_{T} \\
\widetilde{w}(x, 0)=0 & \text { in } \mathbb{R}^{n}
\end{aligned}
$$

which implies $\widetilde{w}=0$. By the previous lemme $u_{R}$ admits the representation

$$
u^{R}(x, t)=\int_{B_{R}} H_{V}^{B_{R}}(x, y, t) d \mu(y)
$$

Since $\left\{H_{V}^{B_{R}}\right\}$ is an increasing sequence and $\lim _{R \rightarrow \infty} H_{V}^{B_{R}}=H_{V}$, we have using again Fatou's lemma as in the proof of Lemma 3.15

$$
u(x, t)=\lim _{R \rightarrow \infty} u^{R}(x, t)=\lim _{R \rightarrow \infty} \int_{B_{R}} H_{V}^{B_{R}}(x, y, t) d \mu(y)=\int_{\mathbb{R}^{n}} H_{V}(x, y, t) d \mu(y)
$$

Lemma 3.16 Harnack inequality Let $C_{1}>0$ and $V(x, t)$ be a potential satisfying (1.16) If $u$ is a positive solution of (1.17), then the Harnack inequality is valid:

$$
u(y, s) \leq u(x, t) \exp \left(C\left(n, C_{1}\right)\left(\frac{|x-y|^{2}}{t-s}+\frac{t}{s}+1\right)\right), \quad \forall(y, s),(x, t) \in Q_{T}, s<t
$$

Proof. We extend $V$ for $t \geq T$ by the value $C_{1} t^{-1}$. We consider the linear parabolic problem

$$
\begin{equation*}
\partial_{t} u-\Delta u+V u=0, \quad \text { in } \quad \mathbb{R}^{n} \times[1, \infty) \tag{3.13}
\end{equation*}
$$

It is well known that, under the assumption (1.16), every positive solution $u(x, t)$ of (3.13) satisfies the Harnack inequality

$$
u(y, s) \leq u(x, t) \exp \left(C\left(n, C_{1}\right)\left(\frac{|x-y|^{2}}{t-s}+\frac{t}{s}+1\right)\right), \quad \forall(x, t) \in \mathbb{R}^{n} \times[1, \infty)
$$

Set $\tilde{u}(x, t)=u\left(\frac{t}{\lambda^{2}} \frac{x}{\lambda}\right)$. Then $\tilde{u}$ satisfies

$$
\mathrm{u}_{t}-\Delta \mathrm{u}+\frac{1}{\lambda^{2}} V\left(\frac{t}{\lambda^{2}} \frac{x}{\lambda}\right) \tilde{u}=0, \quad \text { in } \mathbb{R}^{n} \times(0, \infty)
$$

We note here that $\frac{1}{\lambda^{2}} V\left(\frac{t}{\lambda^{2}} \frac{x}{\lambda}\right) \leq C_{1}, \forall t \geq \frac{1}{\lambda^{2}}$, thus $\tilde{u}$ satisfies the Harnack inequality

$$
\widetilde{u}(y, s) \leq \widetilde{u}(x, t) \exp \left(C\left(n, C_{1}\right)\left(\frac{|x-y|^{2}}{t-s}+\frac{t}{s}+1\right)\right), \quad \forall(x, t) \in \mathbb{R}^{n} \times\left[\frac{1}{\lambda^{2}}, \infty\right)
$$

By the last inequality and the definition of $\tilde{u}$ we derive the desired result.
Next, we set

$$
\begin{equation*}
\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right):=\left\{y \in \mathbb{R}^{n}: H_{V}(x, y, t)=0\right\} \tag{3.14}
\end{equation*}
$$

If $H_{V}(x, y, t)=0$ for some $(x, t) \in Q_{T}$, then $H_{V}\left(x^{\prime}, y, t^{\prime}\right)=0$ for any $\left(x^{\prime}, t^{\prime}\right) \in Q_{T}, t^{\prime}<t$ by Harnack inequality principle. We prove the Representation formula.

Theorem 3.17 Let $u$ be a positive solution of (1.17). Then there exists a measure $\mu \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right)$ such that

$$
u(x, t)=\int_{\mathbb{R}^{n}} H_{V}(x, y, t) d \mu(y)
$$

and $\mu$ is concentrated on $\left(\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)\right)^{c}$.
Proof. By Lemma 3.15 we have

$$
u(x, t)=\int_{\mathbb{R}^{n}} H_{V}(x, y, t-s) u(y, s) d y \quad \text { for any } s<t \leq T .
$$

We assume that $s \leq \frac{T}{2}$. By Harnack inequality on $x \mapsto H_{V}\left(x, y, \frac{T}{2}\right)$

$$
\int_{\mathbb{R}^{n}} H_{V}\left(0, y, \frac{T}{2}\right) u(y, s) d y \leq c(n) \int_{\mathbb{R}^{n}} H_{V}(0, y, T-s,) u(y, s) d y=c(n) u(0, T)
$$

For any Borel set $E$, we define the measure $\rho_{s}$ by

$$
\rho_{s}(E):=\int_{E} H_{V}\left(0, y, \frac{T}{2}\right) u(y, s) d y \leq \int_{\mathbb{R}^{n}} H_{V}\left(0, y, \frac{T}{2}\right) u(y, s) d y \leq c(n) u(T, 0)
$$

Thus there exists a decreasing sequence $\left\{s_{j}\right\}_{j=1}^{\infty}$ which converges to origin, such that the measure $\rho_{s_{j}}$ converges in the weak* topology to a positive Radon measure $\rho$. Also we have the estimate $\rho\left(\mathbb{R}^{n}\right) \leq C(n) u(0, T)$. Now choose $(x, t) \in Q_{T}$ and $j_{0}$ large enough such that $t>s_{j_{0}}$. Let $\varepsilon>0$, we set for any $j \geq j_{0}$,

$$
W_{j}(y)=\frac{H_{V}\left(x, y, t-s_{j}\right)}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon} .
$$

For any $R>0$ and $|y|>R$ we have

$$
W_{j}(y) \leq \frac{1}{\varepsilon} H_{V}\left(x, y, t-s_{j}\right) \leq \frac{1}{\varepsilon} H\left(x-y, t-s_{j}\right)<\frac{1}{\varepsilon} C\left(x, R, t-s_{j}\right),
$$

where $\lim _{R \rightarrow \infty} C\left(x, R, t-s_{j}\right)=0$. We have also

$$
\int_{|y| \geq R} W_{j}(y) d \rho_{j} \leq \frac{1}{\varepsilon} C\left(x, R, t-s_{j}\right) c(n) u(T, 0) .
$$

For any $|y|<R$, we have by standard parabolic estimates that $W_{j}(y) \rightarrow \frac{H_{V}(x, y, t)}{H_{V}\left(\frac{T}{2}, 0, y\right)+\varepsilon}$ when $j \rightarrow \infty$, uniformly with respect to $y$. Thus by the above estimates it follows

$$
\int_{\mathbb{R}^{n}} W_{j}(y) d \rho_{j} \rightarrow \int_{\mathbb{R}^{n}} \frac{H_{V}(x, y, t)}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon} d \rho
$$

For sufficiently large $j$ we have

$$
\begin{aligned}
\int_{\mathbb{R}^{n}} \frac{H_{V}\left(x, y, t-s_{j}\right)}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon} d \rho_{s_{j}} & =\int_{\mathbb{R}^{n}} \frac{H_{V}\left(x, y, t-s_{j}\right)}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon}\left(H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon-\varepsilon\right) u\left(y, s_{j}\right) d y \\
& =u(x, t)-\varepsilon \int_{\mathbb{R}^{n}} \frac{H_{V}\left(x, y, t-s_{j},\right)}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon} u\left(y, s_{j}\right) d y
\end{aligned}
$$

Note that this is a consequence of the identity

$$
\int_{\mathbb{R}^{n}} H_{V}\left(x, y, t-s_{j}\right) u\left(y, s_{j}\right) d y=u(x, t)
$$

Thus as before, we define $d \widetilde{\rho}_{j}=H_{V}\left(x, y, t-s_{j}\right) u\left(y, s_{j}\right) d y$ and thus there exists a subsequence, say $\left\{\widetilde{\rho}_{j}\right\}$ converges in the weak* topology to a positive Radon measure $\widetilde{\rho}$. Thus we have

$$
\begin{aligned}
\varepsilon \int_{\mathbb{R}^{n}} \frac{H_{V}\left(x, y, t-s_{j}\right)}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon} u\left(y, s_{j}\right) d y & =\varepsilon \int_{\mathbb{R}^{n}} \chi_{\left(\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)\right)^{c}} \frac{H_{V}\left(x, y, t-s_{j}\right)}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon} u\left(s_{j}, y\right) d y \\
& \rightarrow \varepsilon \int_{\mathbb{R}^{n}} \chi_{\left(\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)\right)^{c}} \frac{1}{H_{V}\left(\frac{T}{2}, 0, y\right)+\varepsilon} d \widetilde{\rho} .
\end{aligned}
$$

Combining the above relations, we derive

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} \frac{H_{V}(x, y, t)}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon} d \rho=u(x, t)-\varepsilon \int_{\mathbb{R}^{n}} \chi_{\left(\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)\right)^{c}} \frac{1}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon} d \widetilde{\rho} \tag{3.15}
\end{equation*}
$$

Now, we have

$$
\lim _{\varepsilon \rightarrow 0} \chi_{\left(\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)\right)^{c}} \frac{\varepsilon}{H_{V}\left(\frac{T}{2}, 0, y\right)+\varepsilon}=0
$$

and by Harnack inequality on the function $x \mapsto H_{V}(x, y, t)$

$$
\frac{H_{V}(x, y, t)}{H_{V}\left(0, y, \frac{T}{2}\right)+\varepsilon} \leq C(t, T)
$$

thus by dominated convergence theorem, we can let $\varepsilon$ tend to 0 in (3.15) and obtain

$$
\int_{\mathbb{R}^{n}} \frac{H_{V}(x, y, t)}{H_{V}\left(0, y, \frac{T}{2}\right)} d \rho=u(x, t) .
$$

And the result follows if we set

$$
d \mu=\chi_{\left(\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)\right)^{c}} \frac{1}{H_{V}\left(0, y, \frac{T}{2}\right)} d \rho .
$$

In the next result we give a construction of $H_{V}$, with some estimates and a different proof of the existence of an initial measure for positive solutions of (1.16).

Theorem 3.18 Assume $V$ satisfies (1.16), then

$$
\begin{equation*}
H_{V}(x, y, t)=e^{\psi(x, t)} \Gamma(x, y, t, 0) \tag{3.16}
\end{equation*}
$$

where

$$
\begin{equation*}
\psi(x, t)=\int_{t}^{T} \int_{\mathbb{R}^{n}} \frac{e^{-\frac{|x-y|^{2}}{4(s-t)}}}{4 \pi(t-s)} V(y, s) d y d s \tag{3.17}
\end{equation*}
$$

and

$$
\begin{equation*}
c_{1} \frac{e^{-\gamma_{1} \frac{|x-y|^{2}}{s-t}}}{(t-s)^{\frac{n}{2}}} \leq \Gamma(x, y, t, s) \leq c_{2} \frac{e^{-\gamma_{2} \frac{|x-y|^{2}}{s-t}}}{(t-s)^{\frac{n}{2}}} \tag{3.18}
\end{equation*}
$$

for some positive constants $c_{i}$ and $\gamma_{i}, i=1,2$.
Proof. Assuming that $u$ is a positive solution of (1.17), we set $u(x, t)=e^{\psi(x, t)} v(x, t)$. Then

$$
\begin{equation*}
\partial_{t} v-\Delta v-2 \nabla \psi \cdot \nabla v-|\nabla \psi|^{2} v+\left(\partial_{t} \psi-\Delta \psi+V\right) v=0 . \tag{3.19}
\end{equation*}
$$

We choose $\psi$ as the solution of problem

$$
\begin{align*}
-\partial_{t} \psi-\Delta \psi+V \psi=0 & \text { in } Q_{T} \\
\psi(., T)=0 & \text { in } \mathbb{R}^{n} . \tag{3.20}
\end{align*}
$$

Then $\psi$ is expressed by (3.17). Furthermore, by standard computations,
(i) $0 \leq \psi(x, t) \leq c \ln \frac{T}{t}$
(ii) $\quad|\nabla \psi(x, t)| \leq c_{1}(T)+c_{2}(T) \ln \frac{T}{t}$

The function $v$ satisfies

$$
\begin{equation*}
\partial_{t} v-\Delta v-2 \nabla \psi \cdot \nabla v-|\nabla \psi|^{2} v=0 . \tag{3.22}
\end{equation*}
$$

Then, by (3.21),

$$
\begin{array}{ll}
\text { (i) } & 0 \leq \int_{\mathbb{R}^{n}} \sup \left\{|\psi(x, s)|^{q}: x \in \mathbb{R}^{n}\right\} d s \leq M_{1} \\
\text { (ii) } & 0 \leq \int_{\mathbb{R}^{n}} \sup \left\{|\nabla \psi(x, s)|^{q}: x \in \mathbb{R}^{n}\right\} d s \leq M_{2} \tag{3.23}
\end{array}
$$

for any $1 \leq q<\infty$ for some $M_{i} \in \mathbb{R}_{+}$. This is the condition $H$ in [2] with $R_{0}=\infty$ and $p=\infty$. Therefore there exists a kernel function $\Gamma \in C\left(\mathbb{R}^{n} \times \mathbb{R}^{n} \times(0, T) \times(0, T)\right)$ which satisfies (3.18) and there exists also a positive Radon measure $\mu$ in $\mathbb{R}^{n}$ such that

$$
\begin{equation*}
v(x, t)=\int_{\mathbb{R}^{n}} \Gamma(x, y, t, 0) d \mu(y) . \tag{3.24}
\end{equation*}
$$

Finally $u$ verifies

$$
\begin{equation*}
u(x, t)=e^{\psi(x, t)} \int_{\mathbb{R}^{n}} \Gamma(x, y, t, 0) d \mu(y) . \tag{3.25}
\end{equation*}
$$

### 3.4 Reduced measures

In this section we assume that $V$ is nonnegative, but not necessarily satisfies (1.16), therefore we can construct $H_{V}[\mu]$ for $\mu \in \mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$. Furthermore, if $\mu$ is nonnegative we can consider the solution $u_{k}$ of the problem

$$
\begin{align*}
\partial_{t} u-\Delta u+V^{k} u=0, & \text { in } Q_{T} \\
u(., 0)=\mu & \text { in } \mathbb{R}^{n} \tag{3.26}
\end{align*}
$$

where $V^{k}=\min \{V, k\}$. Then there holds

$$
u_{k}(x, t)=\int_{\mathbb{R}^{n}} H_{V^{k}}(t, x, y) d \mu(y)=\mathbb{H}_{V^{k}}[\mu](x, t),
$$

and

$$
u_{k}+\int_{0}^{t} \int_{\mathbb{R}^{n}} H(t-s, x, y) V^{k} u_{k} d y d s=\mathbb{H}[\mu]
$$

Since $k \mapsto H_{V^{k}}$ is decreasing and converges to $H_{V}$, we derive

$$
\lim _{k \rightarrow \infty} u_{k}=u=\int_{\mathbb{R}^{n}} H_{V}(t, x, y) d \mu(y) .
$$

By Fatou's lemma

$$
\int_{0}^{t} \int_{\mathbb{R}^{n}} H(t-s, x, y) V u d y d s \leq \liminf _{k \rightarrow \infty} \int_{0}^{t} \int_{\mathbb{R}^{n}} H(t-s, x, y) V^{k} u_{k} d y d s
$$

It follows

$$
u(x, t)+\int_{0}^{t} \int_{\mathbb{R}^{n}} H(t-s, x, y) V u d y d s \leq \int_{\mathbb{R}^{n}} H_{V}(t, x, y) d \mu(y), \quad \forall(x, t) \in Q_{T}
$$

Now since $V u \in L_{l o c}^{1}\left(\bar{Q}_{T}\right)$ and

$$
\partial_{t} u-\Delta u+V u=0, \quad \text { in } \quad Q_{T}
$$

the function

$$
u(x, t)+\int_{0}^{t} \int_{\mathbb{R}^{n}} H(t-s, x, y) V u d y d s
$$

is nonnegative and satisfies the heat equation in $Q_{T}$. Therefore it admits an initial trace $\mu^{*} \in$ $\mathfrak{M}_{+}\left(\mathbb{R}^{n}\right)$ and actually $\mu^{*} \in \mathfrak{M}_{T}\left(\mathbb{R}^{n}\right)$. Furthermore, we have

$$
u(x, t)+\int_{0}^{t} \int_{\mathbb{R}^{n}} H(t-s, x, y) V u d y d s=\int_{\mathbb{R}^{n}} H(x-y, t) d \mu^{*}(y), \forall(x, t) \in Q_{T}
$$

or equivalently, $u$ is a positive weak solution of the problem

$$
\begin{aligned}
\partial_{t} u-\Delta u+V u & =0 & & \text { in } Q_{T} \\
u(., 0) & =\mu^{*} & & \text { in } \mathbb{R}^{n} .
\end{aligned}
$$

Note that $\mu^{*} \leq \mu$ and the mapping $\mu \mapsto \mu^{*}$ is nondecreasing.
Definition 3.19 The measure $\mu^{*}$ is the reduced measure associated to $\mu$
The proofs of the next results are similar to the ones of [12, Section 5].
Proposition 3.20 There holds $\mathbb{H}_{V}[\mu]=\mathbb{H}_{V}\left[\mu^{*}\right]$. Furthermore the reduced measure $\mu^{*}$ is the largest measure for which the following problem

$$
\begin{array}{rc}
\partial_{t} v-\Delta v+V v=0 & \text { in } Q_{T} \\
\lambda \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right), \lambda \leq \mu &  \tag{3.27}\\
v(., 0)=\lambda & \text { in } \mathbb{R}^{n}
\end{array}
$$

admits a solution.
Proposition 3.21 Let $W_{k}$ be an increasing sequence of nonnegative bounded measurable functions converging to $V$ a.e. in $Q_{T}$. Then the solution $u_{k}$ of

$$
\begin{aligned}
\partial_{t} v-\Delta v+W_{k} v & =0 & & \text { in } Q_{T} \\
v(., 0) & =\mu & & \text { in } \mathbb{R}^{n}
\end{aligned}
$$

converges to $u_{\mu^{*}}$.
We recall that $\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right):=\left\{y \in \mathbb{R}^{n}: H_{V}(x, y, t)=0\right\}$.
Proposition 3.22 Let $\mu$ be a nonnegative measure in $\mathcal{M}_{\mathcal{T}}\left(\mathbb{R}^{n}\right)$. Then
(i) $\left(\mu-\mu^{*}\right)\left(\left(\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)\right)^{c}\right)=0$
(ii) If $\mu\left(\left(\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)\right)^{c}\right)=0$, then $\mu^{*}=0$.
(iii) There always holds $\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right) \subset Z_{V}$.

## 4 Initial trace

### 4.1 The direct method

The initial trace that we developed in this section is an adaptation to the parabolic case of the notion of boundary trace for elliptic equations (see [9], [10], [12]). If $G \subset \overline{Q_{T}}$ is a relatively open set, we denote

$$
W(G)=\bigcap_{1 \leq p<\infty} W_{p}^{2,1}(G) \text { and } W_{l o c}(G)=\bigcap_{1 \leq p<\infty} W_{p l o c}^{2,1}(G)
$$

Since $V \in L_{l o c}^{\infty}\left(Q_{T}\right)$, any solution of (1.17) belongs to $W_{l o c}\left(Q_{T}\right)$.
Proposition 4.1 Let $u \in W_{l o c}\left(Q_{T}\right)$ be a positive solution (1.17). Assume that, for some $x \in \mathbb{R}^{n}$, there exists an open bounded neighborhood $U$ of $x$ such that

$$
\begin{equation*}
\iint_{Q_{T}^{U}} u(y, t) V(y, t) d x d t<\infty \tag{4.28}
\end{equation*}
$$

Then $u \in L^{1}(U \times(0, T))$ and there exists a unique positive Radon measure $\mu$ in $U$ such that

$$
\lim _{t \rightarrow 0} \int_{U} u(y, t) \phi(x) d x=\int_{U} \phi(x) d \mu, \quad \forall \phi \in C_{0}^{\infty}(U)
$$

Proof. Since $V u \in L^{1}(U \times(0, T))$ the following problem has a weak solution $v$ (see [9]).

$$
\begin{aligned}
\partial_{t} v-\Delta v & =V u, \quad \text { in } \quad U \times(0, T] \\
v(x, t) & =0 \quad \text { on } \partial U \times(0, T] \\
v(x, 0) & =0 \quad \text { in } U .
\end{aligned}
$$

Thus the function $w=u+v$ satisfies the heat equation. Thus there exists a unique Radon measure $\mu$ such that

$$
\lim _{t \rightarrow 0} \int_{U} w(y, t) \phi(x) d x=\int_{U} \phi(x) d \mu, \quad \forall \phi \in C_{0}^{\infty}(U)
$$

And the result follows since the initial data of $v$ is zero.
We set

$$
\begin{equation*}
\mathcal{R}(u)=\left\{y \in \mathbb{R}^{n}: \exists \text { bounded neighborhood } U \text { of } y, \iint_{Q_{T}^{U}} u(y, t) V(y, t) d x d t<\infty\right\} \tag{4.29}
\end{equation*}
$$

Then $\mathcal{R}(u)$ is open and there exists a unique positive Radon measure $\mu$ on $\mathcal{R}(u)$ such that

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathcal{R}} u(y, t) \phi(x) d x=\int_{\mathcal{R}} \phi(x) d \mu, \quad \forall \phi \in C_{0}^{\infty}(\mathcal{R}) \tag{4.30}
\end{equation*}
$$

Proposition 4.2 Let $u \in W_{l o c}\left(\mathbb{R}^{n} \times(0, T]\right)$ be a positive solution of (1.17). Assume that, for some $x \in \mathbb{R}^{n}$, there holds

$$
\begin{equation*}
\iint_{Q_{T}^{U}} u(y, t) V(y, t) d y d t=\infty \tag{4.31}
\end{equation*}
$$

for any bounded open neighborhood $U$ of $x$. Then

$$
\begin{equation*}
\limsup _{t \rightarrow 0} \int_{U} u(y, t) d y=\infty \tag{4.32}
\end{equation*}
$$

Proof. We will prove it by contradiction. We assume that there exists an open neighborhood of $x$ such that

$$
\int_{U} u(y, t) d y \leq M<\infty \quad \forall t \in(0, T)
$$

Then $\|u\|_{L^{1}\left(Q_{T}^{U}\right)} \leq M T$. Let $B_{r}(x) \subset \subset U$ for some $r>0$, and $\zeta \in C_{0}^{\infty}\left(B_{r}(x)\right)$, such that $\zeta=1$ in $B_{\frac{r}{2}}(x), \zeta=0$ in $B_{r}^{c}(x)$ and $0 \leq \zeta \leq 1$. Then since $u$ is a positive solution we have

$$
\begin{aligned}
& \int_{U} \partial_{t} u \zeta d x-\int_{U} u \Delta \zeta d x+\int_{U} V u \zeta d x=0 \Rightarrow \int_{B_{\frac{r}{2}}} V u d x \leq \int_{U} \partial_{t} u \zeta d x-\int_{U} u \Delta \zeta d x \Rightarrow \\
& \int_{U} \partial_{t} u \zeta d x-\int_{U} u \Delta \zeta d x+\int_{U} V u \zeta d x=0 \Rightarrow \int_{B_{\frac{r}{2}}} V u d x \leq-\int_{U} \partial_{t} u d x+M\|\Delta \zeta\|_{L^{\infty}}
\end{aligned}
$$

Integrating the above inequality on $(s, T)$, we get

$$
\begin{equation*}
\int_{s}^{T} \int_{B_{\frac{r}{2}}} V u d x d r \leq-\int_{U} u(x, T) d x+\int_{U} u(x, s) d x+\|u\|_{L^{1}\left(Q_{T}^{U}\right)}\|\Delta \zeta\|_{L^{\infty}} \tag{4.33}
\end{equation*}
$$

Letting $s \rightarrow 0$, we reach a contradiction.
Remark. It is not clear wether there holds

$$
\begin{equation*}
\liminf _{t \rightarrow 0} \int_{U} u(y, t) d y=\infty \tag{4.34}
\end{equation*}
$$

However, it follows from (4.33) that if $u \in L^{1}\left(Q_{T}^{U}\right)$, the above equality holds.
Definition 4.3 If $u$ is a positive solution of (1.17), we set $\mathcal{S}(u)=\mathbb{R}^{n} \backslash \mathcal{R}(u)$. The couple $(\mathcal{S}(u), \mu)$ is called the initial trace of $u$, denoted by $\operatorname{tr}_{\{t=0\}}(u)$. The sets $\mathcal{R}(u)$ and $\mathcal{S}(u)$ are respectively the regular and the singular sets of $\operatorname{tr}_{\{t=0\}}(u)$ and $\mu \in \mathfrak{M}_{+}(\mathcal{R}(u))$ is its regular part.
Example Take $V(x, t)=c t^{-1}, c>0$. If $u$ satisfies

$$
\begin{equation*}
\partial_{t} u-\Delta u+\frac{c}{t} u=0 \tag{4.35}
\end{equation*}
$$

then $v(x, t)=t^{c} u(x, t)$ satisfies the heat equation. Thus, if $u \geq 0$, there exists $\mu \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right)$ such that

$$
\begin{equation*}
u(x, t)=t^{-c} \mathbb{H}[\mu](x, t) . \tag{4.36}
\end{equation*}
$$

This is a representation formula. Notice that $V u(x, t)=c t^{-c-1} \mathbb{H}[\mu](x, t)$, therefore the regular set of $\operatorname{tr}_{\{t=0\}}(u)$ may be empty.

Proposition 4.4 Assume $V$ satisfies (1.16) and let $u \in W_{l o c}\left(Q_{T}\right)$ be a positive solution of (1.17) with initial trace $\left(\mathcal{S}(u), \mu_{u}\right)$. Then $u \geq u_{\mu_{u}}$.

Proof. We assume $\mathcal{S}(u) \neq \mathbb{R}^{n}$ otherwise the result is proved. Let $G$ and $E$ be open bounded domains such that $G \subset \subset E \subset \subset \mathcal{R}(u)$. Let $0<\delta=\inf \left\{|x-y|: x \in G, y \in E^{c}\right\}$. Choose $R>0$ such that $E \subset \subset B_{R}$. Let $\left\{t_{j}\right\}_{j=1}^{\infty}$ be a decreasing sequence converging to 0 . We denote by $u_{j}$ the weak solution of the problem

$$
\begin{aligned}
\partial_{t} v-\Delta v+V v & =0 & & \text { in } B_{R} \times\left(t_{j}, T\right] \\
v(x, t) & =0 & & \text { on } \partial B_{R} \times\left(t_{j}, T\right] \\
v\left(., t_{j}\right) & =\chi_{G} u\left(., t_{j}\right) & & \text { in } B_{R} \times\left\{t_{j}\right\},
\end{aligned}
$$

where $\chi$ is the characteristic function on $G$. Let $v_{j}^{R}$, be the solution

$$
\begin{aligned}
\partial_{t} v-\Delta v & =0 & & \text { in } \mathbb{R}^{n} \times\left(t_{j}, \infty\right] \\
v\left(., t_{j}\right) & =\chi_{G} u\left(., t_{j}\right) & & \text { in } \mathbb{R}^{n} \times\left\{t_{j}\right\} .
\end{aligned}
$$

Then by maximum principle we have $u_{j}^{R} \leq u$ and $u_{j}^{R} \leq v_{j}$ in $B_{R} \times\left(t_{j}, T\right]$, for any $j \in \mathbb{N}$. By standard parabolic estimates, we may assume that the sequence $u_{j}^{R}$ converges locally uniformly in $Q_{T}^{B_{R}}$ to a function $u^{R} \leq u$. Moreover, since $\chi_{G} \mu_{u}\left(., t_{j}\right) \rightharpoonup \chi_{G} \mu_{u}$ in the weak* topology, we derive from the representation formula that $v_{j} \rightarrow \mathbb{H}\left[\chi_{G} \mu_{u}\right]$. Furthermore $u^{R} \leq v$, which implies $\chi_{\left(t_{j}, T\right)} u_{j}^{R} \rightarrow u^{R}$ in $L^{1}\left(Q_{T}^{B_{R}}\right)$. There also holds

$$
\int_{t_{j}}^{T} \int_{B_{R}} u_{j}^{R} V d x d t=\int_{t_{j}}^{T} \int_{E} u_{j}^{R} V d x d t+\int_{t_{j}}^{T} \int_{B_{R} \backslash E} u_{j}^{R} V d x d t
$$

and, by the choice of $E$ and dominated convergence theorem,

$$
\int_{t_{j}}^{T} \int_{E} u_{j}^{R} V d x d t \leq \int_{0}^{T} \int_{E} u V d x d t<\infty \Rightarrow \lim _{j \rightarrow \infty} \int_{t_{j}}^{T} \int_{E} u_{j}^{R} V d x d t=\int_{0}^{T} \int_{E} u^{R} V d x d t
$$

Furthermore, for any $x \in B_{R} \backslash E$,

$$
v_{j}(x, t)=\left(\frac{1}{4 \pi\left(t-t_{j}\right)}\right)^{\frac{n}{2}} \int_{R^{n}} e^{-\frac{|x-y|^{2}}{4\left(t-t_{j}\right)}} \chi_{G} u\left(y, t_{j}\right) d y \leq\left(\frac{1}{4 \pi\left(t-t_{j}\right)}\right)^{\frac{n}{2}} e^{-\frac{\delta^{2}}{4\left(t-t_{j}\right)}} \int_{G} u\left(y, t_{j}\right) d y
$$

Next, since $V(x, t) \leq C t^{-1}$ and $u_{j}^{R} \leq v_{j}$, we obtain

$$
\begin{equation*}
\lim _{j \rightarrow \infty} \int_{t_{j}}^{T} \int_{B_{R} \backslash E} u_{j}^{R} V d x d t=\int_{0}^{T} \int_{B_{R} \backslash E} u^{R} V d x d t \tag{4.37}
\end{equation*}
$$

by using the previous estimate and the fact that $\chi_{G} \mu_{u}\left(x, t_{j}\right) \rightharpoonup \chi_{G} \mu_{u}$ in the weak* topology. It follows $\chi_{\left(t_{j}, T\right)} V u_{j}^{R} \rightarrow V u^{R}$ in $L^{1}\left(Q_{T}^{B_{R}}\right)$. There holds also $u_{G}^{R} \leq u$; by the maximum principle, the mapping $R \mapsto u_{G}^{R}$ is increasing and bounded from above by $u$. In view of Lemma 3.15,

$$
\lim _{R \rightarrow \infty} u_{G}^{R}=u_{G} \leq u
$$

and $u_{G}$ is a positive weak solution of

$$
\begin{aligned}
\partial_{t} v-\Delta v+V v & =0 & & \text { in } Q_{T} \\
v(., 0) & =\chi_{G} \mu_{u} & & \text { in } \mathbb{R}^{n} .
\end{aligned}
$$

Consider an increasing sequence $\left\{G_{i}\right\}_{i=1}^{\infty}$ of bounded open subsets, $G_{i} \subset \subset \mathcal{R}(u)$, with the property that $\bigcup_{i=1}^{\infty} G_{i}=\mathcal{R}(u)$. In view of Lemma 3.15 the sequence $\left\{u_{i}=u_{G_{i}}\right\}_{i=1}^{\infty}$ is increasing and converges to $\widetilde{u} \leq u$. Also we have

$$
u_{i}(x, t)+\int_{0}^{t} \int_{\mathbb{R}^{n}} H(t-s, x, y) V u_{i} d y d s=\int_{\mathbb{R}^{n}} H(x-y, t) d \mu_{i}, \forall(x, t) \in Q_{T}
$$

where $\mu_{i}=\chi_{G_{i}} \mu$. Now since $\mu_{i} \rightharpoonup \mu_{u}$, by the monotone convergence theorem we have

$$
\tilde{u}(x, t)+\int_{0}^{t} \int_{\mathbb{R}^{n}} H(t-s, x, y) V \tilde{u} d y d s=\int_{\mathbb{R}^{n}} H(x-y, t) d \mu_{u}, \forall(x, t) \in Q_{T},
$$

and $\tilde{u} \leq u$. this implies $\tilde{u}=u_{\mu_{u}}$, which ends the proof.
Remark. Assumption (1.17) is too strong and has only been used in (4.37). It could have been replaced by the following much weaker one: for any $R>0$ there exists a positive increasing function $\epsilon_{R}$ such that $\lim _{t \rightarrow 0} \epsilon(t)=0$ satisfying

$$
\begin{equation*}
V(x, t) \leq e^{t^{-1} \epsilon_{R}(t)} \quad \forall(x, t) \in Q_{T}^{B_{R}} \tag{4.38}
\end{equation*}
$$

We end this section with a result which shows that the stability of the initial value problem with respect to convergence the initial data in the weak* topology implies that the initial of positive solution has no singular part.

Theorem 4.5 Assume $V$ satisfies, for some $\tau_{0}>0$,

$$
\begin{equation*}
\lim _{|E| \rightarrow 0} \iint_{E} H(x-y, t) V(x, t+\tau) d x d t=0, \quad \text { E Borel subset of } Q_{T}^{B_{R}} \tag{4.39}
\end{equation*}
$$

for any $R>0$, uniformly with respect to $y$ is a compact set and $\tau \in\left[0, \tau_{0}\right]$. If $u$ is a positive solution of (1.17), then $\mathcal{R}(u)=\mathbb{R}^{n}$

Proof. We assume that $\mathcal{S}(u) \neq \emptyset$ and if $z \in \mathcal{S}(u)$ there holds

$$
\iint_{Q_{T}^{B_{r}(z)}} V u d x d t=\infty \quad \forall r>0 .
$$

In view of Proposition 4.1, there exist two sequences $\left\{r_{k}\right\}$ and $\left\{t_{j}\right\}$ decreasing to 0 such that

$$
\lim _{j \rightarrow \infty} \int_{B_{r_{k}}(z)} u\left(x, t_{j}\right) d x=\infty \quad \forall k \in \mathbb{N}
$$

For $k \in \mathbb{N}$ and $m>0$ fixed, there exists $j(k)$ such that

$$
\int_{B_{r_{k}}(z)} u\left(x, t_{j}\right) d x \geq m \quad \forall j \geq j(k) \in \mathbb{N}
$$

and there exists $\ell_{k}>0$ such that

$$
\int_{B_{r_{k}}(z)} \min \left\{u\left(x, t_{j(k)}\right), \ell_{k}\right\} d x=m
$$

Furthermore $j(k) \rightarrow \infty$ when $k \rightarrow \infty$. Let $R>\max \left\{r_{j}: j=1,2, \ldots\right\}$ and $u_{k}$ be the solution of

$$
\begin{aligned}
\partial_{t} v-\Delta v+V v & =0 & & \text { in } \mathbb{R}^{n} \times\left(t_{j(k)}, T\right] \\
v\left(., t_{j}\right) & =\chi_{B_{r_{k}}(z)} \min \left\{u\left(., t_{j(k)}\right), \ell_{k}\right\} & & \text { in } \mathbb{R}^{n} \times\left\{t_{j(k)}\right\},
\end{aligned}
$$

Then $\chi_{B_{r_{k}}(z)} \min \left\{u\left(., t_{j(k)}\right), \ell_{k}\right\} \rightarrow m \delta_{z}$ in the weak sense of measures. By Proposition 5.5 we obtain that $u \geq u_{k}$ on $B_{R}(z) \times\left(t_{j(k)}, T\right]$. Applying Proposition 2.8, and the remark here after, we conclude that $u_{k}\left(., .+t_{j(k)}\right) \rightarrow u_{m \delta_{z}}=m u_{\delta_{z}}$ in $L_{l o c}^{1}\left(\bar{Q}_{R}^{T}\right)$ This implies $u \geq m u_{\delta_{z}}$, and as $m$ is arbitrary, $u=\infty$, contradiction.

### 4.2 The sweeping method

In this subsection we adapt to equation (1.17) the sweeping method developed in [12] for constructing the boundary trace of solutions of stationnary Shrödinger equations. If $A \subset \mathbb{R}^{n}$ is a Borel set, we denote by

$$
\mathfrak{M}_{T+}(A)=\left\{\mu \in \mathfrak{M}_{+}\left(\mathbb{R}^{n}\right): \mu\left(A^{c}\right)=0, \int_{A} e^{-\frac{|x|^{2}}{4 T}} d \mu<\infty\right\}
$$

We recall that $\mu^{*}$ denotes the reduced measure associated to $\mu$.
Proposition 4.6 Let $u \in W_{\text {loc }}\left(Q_{T}\right)$ be a positive solution of (1.17) with singular set $\mathcal{S}(u) \nsubseteq \mathbb{R}^{n}$. If $\mu \in \mathfrak{M}_{T+}(\mathcal{S}(u))$, we set $v_{\mu}=\inf \left\{u, u_{\mu^{*}}\right\}$. Then

$$
\partial_{t} v_{\mu}-\Delta v_{\mu}+V v_{\mu} \geq 0 \quad \text { in } Q_{T}
$$

and $v_{\mu}$ admits a boundary trace $\gamma_{u}(\mu) \in \tilde{\mathfrak{M}}_{+}(\mathcal{S}(u))$. The mapping $\mu \mapsto \gamma_{u}(\mu)$ is nondecreasing and $\gamma_{u}(\mu) \leq \mu$.

Proof. It is classical that $v_{\mu}:=\inf \left\{u, u_{\mu^{*}}\right\}$ is a supersolution of (1.17) and $v_{\mu} \in \mathcal{E}_{\nu}\left(Q_{T}\right)$ as it holds with $u_{\mu^{*}}$. The function

$$
(x, t) \mapsto w(x, t)=\int_{0}^{t} \int_{\mathbb{R}^{n}} H(t-s, x, y) V(y, s) v_{\mu}(y, s) d y d s
$$

satisfies

$$
\begin{aligned}
\partial_{t} w-\Delta w-V w=0 & \text { in } Q_{T} \\
w(., 0)=0 & \text { in } \mathbb{R}^{n} \times\{0\} .
\end{aligned}
$$

Thus $v_{\mu}+w$ is a nonnegative supersolution of the heat equation in $Q_{T}$. It admits an initial trace in $\mathfrak{M}_{T+}(\mathcal{S}(u))$ that we denote by $\gamma_{u}(\mu)$. Clearly $\gamma_{u}(\mu) \leq \mu^{*} \leq \mu$ since $v_{\mu} \leq u_{\mu^{*}}$ and $\gamma_{u}(\mu)$ is nondecreasing with respect to $\mu$ as it is the case with $\mu \mapsto u_{\mu^{*}}$ is. Finally, since $v_{\mu}$ is a positive supersolution, it is larger that the solution of 2.1 where the initial data $\mu$ is replaced by $\gamma_{u}(\mu)$, that is $u_{\gamma_{u}(\mu)} \leq v_{\mu}$.

The proofs of the next four propositions are mere adaptations to the parabolic case of similar results dealing with elliptic equations and proved in [12]; we omit them.

Proposition 4.7 Let

$$
\nu_{S}(u):=\sup \left\{\gamma_{u}(\mu): \mu \in \mathfrak{M}_{T+}(\mathcal{S}(u))\right\} .
$$

Then $\nu_{S}(u)$ is a Borel measure on $\mathcal{S}(u)$.
Definition 4.8 The Borel measure $\nu(u)$ defined by

$$
\nu(u)(A):=\nu_{S}(u)(A \cap \mathcal{S}(u))+\mu_{u}(A \cap \mathcal{R}(u)), \quad \forall A \subset \mathbb{R}^{n}, A \text { Borel }
$$

is called the extended initial trace of $u$, denoted by $\operatorname{tr}_{\{t=0\}}^{e}(u)$.
Proposition 4.9 If $A \subset \mathcal{S}(u)$ is a Borel set, then

$$
\nu_{S}(A):=\sup \left\{\gamma_{u}(\mu)(A): \mu \in \mathfrak{M}_{T+}(A)\right\}
$$

Proposition 4.10 There always holds $\nu\left(\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)\right)=0$, where $\operatorname{Sing}_{V}\left(\mathbb{R}^{n}\right)$ is defined in (3.14).
Proposition 4.11 Assume $V$ satisfies condition (4.39). If $u$ is a positive solution of (1.17), then $\operatorname{tr}_{\{t=0\}}^{e}(u)=\mu_{u} \in \mathfrak{M}_{T+}\left(\mathbb{R}^{n}\right)$.

## 5 Appendix: the case of a bounded domain

### 5.1 The subcritical case

Let $\Omega$ be a bounded domain with a $C^{2}$ boundary. We denote by $\mathfrak{M}(\Omega)$ the space of Radon measures in $\Omega$, by $\mathfrak{M}_{+}(\Omega)$ its positive cone and by $\mathfrak{M}_{\rho}(\Omega)$ the space of Radon measures in $\Omega$ which satisfy

$$
\begin{equation*}
\int_{\Omega} \rho d|\mu|<\infty \tag{5.1}
\end{equation*}
$$

for some weight function $\rho: \Omega \mapsto \mathbb{R}_{+}$. As an important particular case $\rho(x)=d^{\alpha}(x)$, where $d(x)=\operatorname{dist}(x, \partial \Omega)$ and $\alpha \geq 0$. We consider the linear parabolic problem

$$
\begin{align*}
\partial_{t} u-\Delta u+V u & =0, & & \text { in } Q_{T}^{\Omega}=\Omega \times(0, T] \\
u & =0 & & \text { on } \partial_{l} Q_{T}^{\Omega}=\partial \Omega \times(0, T]  \tag{5.2}\\
u(., 0) & =\mu & & \text { in } \Omega .
\end{align*}
$$

Definition 5.1 We say that $\mu \in \mathfrak{M}_{d}(\Omega)$ is a good measure if the above problem has a weak solution $u$, i.e. there exists a function $u \in L^{1}\left(Q_{T}^{\Omega}\right)$, such that $V u \in L_{d}^{1}\left(Q_{T}^{\Omega}\right)$ which satisfies

$$
\begin{equation*}
-\int_{0}^{T} \int_{\Omega} u\left(\partial_{t} \phi+\Delta \phi\right) d x d t+\int_{0}^{T} \int_{\Omega} V u \phi d x d t=\int_{\Omega} \phi(x, 0) d \mu \tag{5.3}
\end{equation*}
$$

$\forall \phi \in C^{1,1 ; 1}\left(\overline{Q_{T}^{\Omega}}\right)$ which vanishes on $\partial_{l} Q_{T}^{\Omega}$ and satisfies $\phi(x, T)=0$.
Definition 5.2 Let $H^{\Omega}(x, y, t)$ be the heat kernel in $\Omega$. Then we say that $\mu \in \mathfrak{M}_{d}(\Omega)$ is a admissible measure if

$$
\left\|\mathbb{H}^{\Omega}[|\mu|]\right\|_{L^{1}\left(Q_{T}^{\Omega}\right)}=\int_{Q_{T}^{\Omega}}\left(\int_{\Omega} H^{\Omega}(x-y, t) d|\mu(y)|\right) V(x, t) \psi(x) d x d t<\infty .
$$

The next a proposition is direct consequence of [9, Lemma 2.4].
Proposition 5.3 Assume $\mu \in \mathfrak{M}_{d}(\Omega)$ and let $u$ be a weak solution of problem (5.2), then the following inequalities are valid
(i)

$$
\|u\|_{L^{1}\left(Q_{T}^{\Omega}\right)}+\|V u\|_{L_{\psi}^{1}\left(Q_{T}^{\Omega}\right)} \leq C(n, \Omega) \int_{\Omega} d d|\mu|
$$

(ii)

$$
-\int_{0}^{T} \int_{\Omega}|u|\left(\partial_{t} \phi+\Delta \phi\right) d x d t+\int_{0}^{T} \int_{\Omega}|u| V \phi d x d t \leq \int_{\Omega} \phi(x, 0) d|\mu|
$$

$\forall \phi \in C^{1,1 ; 1}\left(\overline{Q_{T}^{\Omega}}\right), \phi \geq 0$.
(iii)

$$
\lambda_{\Omega} \int_{0}^{T} \int_{\Omega}(x) u^{+} d x d t+\int_{0}^{T} \int_{\Omega} V u^{+} \psi d x d t \leq \int_{\Omega} \psi(x) d \mu^{+}
$$

where $\psi$ is the solution of

$$
\begin{align*}
-\Delta \psi & =1, & & \text { in } \Omega \\
\psi & =0 & & \text { on } \partial \Omega \tag{5.4}
\end{align*}
$$

Proof. For (ii), in [9, Lemma 2.4, p 1456], above from the relation (2.39), we can take $\tilde{\zeta}=\gamma(u) \zeta$ for some $0 \leq \zeta \in C^{1,1 ; 1}\left(\overline{Q_{T}^{\Omega}}\right)$, since $u=0$ on $\partial_{l} Q_{T}^{\Omega}$. For (iii) we consider (as in [9, Remark 2.5]) $\phi(x, t)=t \psi(x)$. The inequality holds by the same type of calculations as in [12].

Proposition 5.4 The problem (5.2) admits at most one solution. Furthermore, if $\mu$ is admissible, then there exists a unique solution; we denote it $u_{\mu}$.

Similarly as Theorem 2.7 and Proposition 2.7 , we have the following stability results
Proposition 5.5 (i) Assume that $V$ satisfies the stability condition

$$
\begin{equation*}
\lim _{|E| \rightarrow 0} \iint_{E} H^{\Omega}(x, y, t) V(y, t) d(x) d y d t=0, \quad \forall E \subset Q_{T}^{\Omega}, E \text { Borel. } \tag{5.5}
\end{equation*}
$$

uniformly with respect to $y \in \Omega$. If $\left\{\mu_{k}\right\}$ is a bounded sequence in $\mathfrak{M}_{d}(\Omega)$ converging to $\mu$ in the dual sense of $\mathfrak{M}_{d}(\Omega)$, then $\left(u_{\mu_{k}}, V u_{\mu_{k}}\right)$ converges to $\left(u_{\mu}, V u_{\mu}\right)$ in $L^{1}\left(Q_{T}^{\Omega}\right) \times L_{d}^{1}\left(Q_{T}^{\Omega}\right)$. (ii) Furthermore if

$$
\begin{equation*}
\lim _{|E| \rightarrow 0} \iint_{E} H^{\Omega}\left(x, y, t+\tau_{n}\right) V(y, t) d(x) d y d t=0, \quad \forall E \subset Q_{T}^{\Omega}, E \text { Borel. } \tag{5.6}
\end{equation*}
$$

uniformly with respect to $y \in \Omega$ and $\tau_{k} \in\left[0, \tau_{0}\right]$ converges to 0 and $\left\{\mu_{k}\right\}$ is in (i), then the solutions $u_{\tau_{k}, \mu_{k}}$ of the shifted problem

$$
\begin{array}{rlrl}
\partial_{t} u-\Delta u+V u & =0 & & \text { on } \quad \Omega \times\left(\tau_{k}, T\right) \\
u & =0 & & \text { on }  \tag{5.7}\\
u\left(., \tau_{k}\right) & =\mu_{k} \times\left(\tau_{k}, T\right) \\
& \text { on } & \Omega \times\left\{\tau_{k}\right\}
\end{array}
$$

(extended by 0 on $\left(0, \tau_{k}\right)$ ) converge to $u_{\mu}$ in $L_{d}^{1}\left(Q_{T}^{\Omega}\right)$, and $\left\{V u_{\mu_{k}}\right\}$ converges to $V u_{\mu}$ in $L_{d}^{1}\left(Q_{T}^{\Omega}\right)$.

Proof. We can easily see that the measure $\mu_{n}$ is admissible and uniqueness holds; furthermore any admissible measure is a good measure is a good measure as in Theorem 2.5, and

$$
\iint_{Q_{T}^{\Omega}} u_{\mu_{n}} d x d s+\iint_{Q_{T}^{\Omega}} u_{\mu_{n}} V \psi d x d s \leq C \int_{\Omega} d \mu_{n}<C .
$$

The remaining of the proof is similar to the one of Theorem 2.7.

### 5.2 The supercritical case

Lemma 5.6 Let $\left\{\mu_{n}\right\}_{n=1}^{\infty}$ be an increasing sequence of good measures converging to some measure $\mu$ in the weak* topology, then $\mu$ is good.

Proof. Let $u_{\mu_{n}}$ be the weak solution of (5.2) with initial data $\mu_{n}$. Then by Proposition 5.5 -(iii), $\left\{u_{\mu_{n}}\right\}$ is an increasing sequence. By 5.5 -(i) the sequence $\left\{u_{\mu_{n}}\right\}$ is bounded in $L^{1}\left(Q_{T}^{\Omega}\right)$. Thus $u_{\mu_{n}} \rightarrow u \in L^{1}\left(Q_{T}^{\Omega}\right)$. Also by (iii) of Proposition 5.5, we have that $V u_{\mu_{n}} \rightarrow V u$ in $L_{\psi}^{1}\left(Q_{T}^{\Omega}\right)$. Thus we can easily prove that $u$ is a weak solution of (5.2) with $\mu$ as initial data.

Let

$$
\begin{equation*}
Z_{V}^{\Omega}=\left\{x \in \Omega: \int_{Q_{T}^{\Omega}} H^{\Omega}(t, x, y) V(y) \psi(y) d y=\infty\right\} \tag{5.8}
\end{equation*}
$$

We note that, since $H^{\Omega}(t, x, y) \leq H(x-y, t)$ for any bounded $\Omega$ with smooth boundary, it holds $Z_{V}^{\Omega} \subset Z_{V}$. By the same arguments as in [12] we can prove the following results

Proposition 5.7 Let $\mu$ be an admissible positive measure. Then $\mu\left(Z_{V}^{\Omega}\right)=0$
Proposition 5.8 Let $\mu \in \mathfrak{M}_{d+}(\Omega)$ such that $\mu\left(Z_{V}^{\Omega}\right)=0$, then $\mu$ is good.
Proposition 5.9 Let $\mu \in \mathfrak{M}_{d+}(\Omega)$ be a good measure. Then the following assertions are equivalent:
(i) $\mu\left(Z_{V}^{\Omega}\right)=0$.
(ii) There exists an increasing sequence of admissible measures $\left\{\mu_{n}\right\}$ which converges to $\mu$ in the weak*-topology

## References

[1] Adams D. R. and Hedberg L. I., Function spaces and potential theory, Grundlehren Math. Wissen. 314, Springer (1996).
[2] Aronson D., Non-negative solutions of linear parabolic equations, Ann. Sc. Norm. Sup. Pisa Cl. Sc. 22, 607-694 (1968).
[3] Choquet G., Theory of capacities, Ann. Inst. Fourier 5, 131-295 (1953-54).
[4] Doob J. L., Classical potential theory and its probabilistic counterpart, Reprint of the 1984 edition, Classics in Mathematics, Springer-Verlag, Berlin (2001).
[5] Dynkin E. B., Diffusions, Superdiffusions and Partial Differential Equations, American Math. Soc., Providence, Rhode Island, Colloquium Publications 50 (2002).
[6] Fuglede B., Le théorème du minimax et la théorie fine du potentiel, Ann. Inst. Fourier 15, 65-87 (1965).
[7] Fuglede B., Application du théorème minimax l'étude de diverses capacités, C.R. Acad. Sci. Paris 266, 921-923 (1968).
[8] Kato T., Shrödinger operators with singular potentials, Israel J. Math. 13, 135-148 (1972).
[9] Marcus M. and Véron L., Initial trace of positive solutions of some nonlinear parabolic equations, Comm. Part. Diff. Equ. 24, 1445-1499 (1999).
[10] Marcus M. and Véron L., Initial trace of positive solutions to semilinear parabolic inequalities, Adv. Nonlinear Studies 2, 395-436 (2002)
[11] Richard Y. and Véron L., Isotropic singularities of solutions of nonlinear elliptic inequalities, Ann. Inst. H. Poincaré-Anal. Non Linéaire 6, 37-72 (1989).
[12] Véron L. and Yarur C., Boundary value problems with measures for elliptic equations with singular potentials, J. Funct. Anal 262, 733-772 (2012).


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