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# A mixing effect induced by sources concentrated in a soft junction and the gradient concentration phenomenon 

Anne-Laure Bessoud ${ }^{*}$, Pongpol Juntharee ${ }^{\dagger \ddagger}$, Christian Licht ${ }^{\dagger}$, and Gérard Michaille ${ }^{\ddagger}$


#### Abstract

We show that the variational limit of a $\varepsilon$-soft and thin junction problem $\left(\mathcal{P}_{\varepsilon}\right)$ with sources concentrated in the junction gives rise to a surface energy mixing the internal energy and sources. The surface energy functional possesses an integral representation with respect to the Gradient YoungConcentration measures generated by sequences $\left(\bar{u}_{\varepsilon}\right)_{\varepsilon>0}$ of minimizers of $\left(\mathcal{P}_{\varepsilon}\right)$.


AMS subject classifications: 35B40, 49Q20, 28A33
Keywords: $\Gamma$-convergence, Gradient Young measures, concentration measures, minimization problems.

[^0]
## 1 Introduction

This paper concerns a soft thin junction subjected to concentrated sources. More precisely, let $\Omega$ be a domain in $\mathbb{R}^{N}$ and let $B_{\varepsilon}:=\Sigma \times\left(-\frac{\varepsilon}{2}, \frac{\varepsilon}{2}\right) \subset \Omega, \Sigma \subset \mathbb{R}^{N-1}$, be the layer occupied by the soft thin junction (cf Figure 1). We consider the minimization problem

$$
\min _{u \in W_{\Gamma_{0}}^{1,2}(\Omega)}\left\{\int_{\Omega \backslash B_{\varepsilon}} f(\nabla u) d x+\varepsilon \int_{B_{\varepsilon}} g(\nabla u) d x-\left\langle\mathcal{S}^{\varepsilon}, u\right\rangle_{\varepsilon}\right\}
$$

where $W_{\Gamma_{0}}^{1,2}(\Omega)$ denotes the space of Sobolev functions with null trace on a part $\Gamma_{0}$ of the boundary of $\Omega$, and the linear form $\left\langle\mathcal{S}^{\varepsilon}, .\right\rangle_{\varepsilon}$ represents the work of the source (or the loading). Let $B:=\Sigma \times$ $\left(-\frac{1}{2}, \frac{1}{2}\right)$. A suitably rescaled $\mathcal{S}_{\varepsilon}$ of $\mathcal{S}^{\varepsilon}$ is assumed to strongly converge to some $\mathcal{S}$ in the dual of the space $V(B):=\left\{u \in L^{2}(\Omega): \frac{\partial u}{\partial x_{N}} \in L^{2}(\Omega)\right\}$ when $\varepsilon$ tends to zero. A general example of such sources which are measures on $B_{\varepsilon}$ is given in Section 4 of the paper. Sources of the form $c \frac{1}{L(\varepsilon)} \mathbb{1}_{B_{\varepsilon}}$ where $c$ is any constant and $L(\varepsilon) \sim \varepsilon$, is a trivial example of measures satisfying this condition with $\mathcal{S}=\mathbb{1}_{B}$. Note that in this paper the source (or the loading) $\mathcal{S}_{\varepsilon}$ is a non $L^{2}$-continuous perturbation of the energy functional $\int_{\Omega \backslash B_{\varepsilon}} f(\nabla u) d x+\varepsilon \int_{B_{\varepsilon}} g(\nabla u) d x$.

Among the physical motivations of $\left(\mathcal{P}_{\varepsilon}\right)$ one may mention various applications to heat conduction or electrostatic problems involving sources concentrated in the layer $B_{\varepsilon}$ with conductivity or permittivity of order the size of $B_{\varepsilon}$. One may also think of membrane problems with an exterior loading concentrated in $B_{\varepsilon}$ occupied by a material with stiffness of order the small size of $B_{\varepsilon}$. Such a problem with a source concentrated in the junction was considered in [3] in a one dimensional case in order to highlight and illustrate a gradient concentration phenomenon, but the authors were not able to express the variational limit problem.

This paper illustrates the same gradient concentration phenomenon with a complete description of the limit problem in the sense of $\Gamma$-convergence (Theorem 3.3). When the size $\varepsilon$ of the layer goes to zero, fields $u_{\varepsilon}$ of bounded energy develop a discontinuity through $\Sigma$. More precisely, at the variational limit, the internal energy functional of the junction $\varepsilon \int_{B_{\varepsilon}} g(\nabla u) d x$ and the work of the loading $\left\langle\mathcal{S}^{\varepsilon}, u\right\rangle_{\varepsilon}$ are combined into a functional of the type

$$
\bar{H}(u)=\int_{\Sigma} \bar{h}\left(\hat{x}, u^{+}-u^{-}, \frac{u^{+}+u^{-}}{2}\right) d \hat{x}
$$

and the limit problem reads as

$$
\begin{equation*}
\min _{u \in W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma)}\left\{\int_{\Omega} f(\nabla u) d x+\bar{H}(u)\right\} \tag{P}
\end{equation*}
$$

where $u^{ \pm}$denote the traces on $\Sigma$. When regarding the various studies devoted to the asymptotic modeling of junction problems (see $[2,9,7,10]$ and references therein) the main novelty is that the density $\bar{h}$ depends also of the mean $\frac{u^{+}+u^{-}}{2}$. Furthermore we show that the sequence of minimizers of $\left(\mathcal{P}_{\varepsilon}\right)$ (which converges to a minimizer $\bar{u}$ of the limit problem $(\mathcal{P})$ ) generates a gradient Young-concentration measure $\bar{\mu}$ in the sense defined in [3]. Then we can give an integral representation of the internal part of $\bar{H}$ with respect to the measure $\bar{\mu}$ (Theorem 5.5) so that it can be localized in $\Sigma \times\{ \pm 1\}$. Finally this provides new bounds on the measure $\bar{\mu}$ (Corollary 5.6).

The paper is organized as follows: in Section 2 we fix notation and give a detailed description of the problem $\left(\mathcal{P}_{\varepsilon}\right)$. Section 3 is devoted to the asymptotic analysis of $\left(\mathcal{P}_{\varepsilon}\right)$ in the sense of the $\Gamma$-convergence of the functional energy extended to $L^{2}(\Omega)$ equipped with its strong topology. In Section 4 we describe a large class of suitable sources $\mathcal{S}_{\varepsilon}$. Finally Section 5 is concerned with the analysis of the gradient concentration phenomenon generated by sequences of minimizers of $\left(\mathcal{P}_{\varepsilon}\right)$. We stress the fact that one could treat the problem in $L^{p}(\Omega), 1<p<+\infty$ in the same way without additional difficulties.

## 2 Description of the minimization problem

Let $\varepsilon>0$ be a small parameter intended to go to zero, more precisely taking values in a countable subset of $\left(0, \varepsilon_{0}\right.$ ] whose 0 is the only cluster point. The reference configuration of the assembly of the two adherents and the adhesive is a cylinder $\Omega:=\Sigma \times(-r, r)$ (with $r>\varepsilon$ ), where $\Sigma$ is a bounded domain in $\mathbb{R}^{N-1}, N \geq 2$, with Lipschitz boundary. For $x \in \mathbb{R}^{N}$ we sometimes write $x=\left(\hat{x}, x_{N}\right)$ where $\hat{x} \in \mathbb{R}^{N-1}$. In all the paper, $C$ denotes a non negative constant which does not depend on $\varepsilon$ and may vary from line to line. We do not relabel the various considered subsequences and the symbols $\rightarrow$ and $\rightharpoonup$ denote various strong convergences and weak convergences respectively. We define the following sets:


Figure 1: physical domain
. $B_{\varepsilon}:=\Sigma \times\left(-\frac{\varepsilon}{2}, \frac{\varepsilon}{2}\right)$;
. $B:=\Sigma \times\left(-\frac{1}{2}, \frac{1}{2}\right)$;
. $\Omega_{\varepsilon}=\Omega \backslash \overline{B_{\varepsilon}} ;$
. $\Gamma_{0}$ is a subset of the boundary $\partial \Omega$ of $\Omega$ such that $\operatorname{dist}\left(\bar{\Gamma}_{0}, \overline{\partial B_{\varepsilon} \cap \partial \Omega}\right)>0$ for all $\varepsilon<\varepsilon_{0}$;
. we write $\Omega_{\varepsilon}^{-}, \Omega_{\varepsilon}^{+}, \Omega^{-}, \Omega^{+}, B_{\varepsilon}^{+}$and $B_{\varepsilon}^{-}$for the sets $\Omega_{\varepsilon} \cap\left[x_{N}<0\right]$ and $\Omega_{\varepsilon} \cap\left[x_{N}>0\right], \Omega \cap\left[x_{N}<0\right]$, $\Omega \cap\left[x_{N}>0\right]$ and $B_{\varepsilon} \cap\left[x_{N}>0\right], B_{\varepsilon} \cap\left[x_{N}<0\right]$ respectively.

We will be concerned with the following spaces:

- $W_{\Gamma_{0}}^{1,2}\left(\Omega_{\varepsilon}\right):=\left\{u \in W^{1,2}\left(\Omega_{\varepsilon}\right): u=0\right.$ on $\left.\Gamma_{0}\right\} ;$
. $W_{\Gamma_{0}}^{1,2}(\Omega):=\left\{u \in W^{1,2}(\Omega): u=0\right.$ on $\left.\Gamma_{0}\right\} ;$
. $W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma):=\left\{u \in W^{1,2}(\Omega \backslash \Sigma): u=0\right.$ on $\left.\Gamma_{0}\right\}$, and for every $z \in W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma), z^{ \pm}$will stand for the traces of $z$ on $\Sigma$ considered as a Sobolev function on $\Omega^{+}$and $\Omega^{-}$respectively.

We say that a function $h: \mathbb{R}^{N} \longrightarrow \mathbb{R} \cup\{+\infty\}$ satisfies a growth condition of order 2 if there exist $\alpha$ and $\beta$ in $\mathbb{R}^{+}$such that

$$
\alpha|\xi|^{2} \leq h(\xi) \leq \beta\left(1+|\xi|^{2}\right) \text { for all } \xi \in \mathbb{R}^{N}
$$

We consider two convex functions $f, g: \mathbb{R}^{N} \longrightarrow \mathbb{R}$ satisfying a growth condition of order 2 , and we assume that there exists a positively 2 -homogeneous function $g^{\infty, 2}$ satisfying

$$
\begin{equation*}
\left|g(\xi)-g^{\infty, 2}(\xi)\right| \leq \beta\left(1+|\xi|^{2-\delta}\right) \text { for all } \xi \in \mathbb{R}^{N} \tag{2.1}
\end{equation*}
$$

for some $\delta, 0<\delta<2$. Note that $g^{\infty, 2}$ is the positively 2 -homogeneous recession function of $g$, i.e.,

$$
g^{\infty, 2}(\xi)=\lim _{t \rightarrow+\infty} \frac{g(t \xi)}{t^{2}}
$$

is convex and satisfies the same growth condition of order 2 . We define the space

$$
V\left(B_{\varepsilon}\right):=\left\{u \in L^{2}\left(B_{\varepsilon}\right): \frac{\partial u}{\partial x_{N}} \in L^{2}\left(B_{\varepsilon}\right)\right\}
$$

equipped with the norm

$$
\|u\|_{V\left(B_{\varepsilon}\right)}:=\left(\int_{B_{\varepsilon}}|u|^{2} d x+\int_{B_{\varepsilon}}\left|\frac{\partial u}{\partial x_{N}}\right|^{2} d x\right)^{\frac{1}{2}}
$$

and we denote the duality bracket between the topological dual space $V^{\prime}\left(B_{\varepsilon}\right)$ and $V\left(B_{\varepsilon}\right)$ by $\langle,\rangle_{\varepsilon}$. The considered total energy functional $F_{\varepsilon}: L^{2}(\Omega) \longrightarrow \mathbb{R} \cup\{+\infty\}$ is defined by

$$
F_{\varepsilon}(u)=\left\{\begin{array}{l}
\int_{\Omega_{\varepsilon}} f(\nabla u) d x+\varepsilon \int_{B_{\varepsilon}} g(\nabla u) d x-\left\langle\mathcal{S}^{\varepsilon}, u\right\rangle_{\varepsilon} \text { if } u \in W_{\Gamma_{0}}^{1,2}(\Omega) \\
+\infty \text { otherwise }
\end{array}\right.
$$

where $\mathcal{S}^{\varepsilon}$ is given in $V^{\prime}\left(B_{\varepsilon}\right)$. Our aim is to describe the asymptotic behavior of the minimization problem

$$
\left(\mathcal{P}_{\varepsilon}\right) \quad \min \left\{F_{\varepsilon}(u): u \in L^{2}(\Omega)\right\}
$$

namely, the limit of $\min \left\{F_{\varepsilon}(u): u \in L^{2}(\Omega)\right\}$ together with the limit of the minimizer $\bar{u}_{\varepsilon}$, and to identify the limit problem in the framework of $\Gamma$-convergence.

Let us consider the space $V(B):=\left\{u \in L^{2}(B): \frac{\partial u}{\partial x_{N}} \in L^{2}(B)\right\}$ equipped with the norm

$$
\|u\|_{V(B)}:=\left(\int_{B}|u|^{2} d x+\int_{B}\left|\frac{\partial u}{\partial x_{N}}\right|^{2} d x\right)^{\frac{1}{2}}
$$

and denote the duality bracket between $V^{\prime}(B)$ and $V(B)$ by $\langle.,$.$\rangle . The linear continuous operator$

$$
\tau_{\varepsilon}: V\left(B_{\varepsilon}\right) \longrightarrow V(B)
$$

is defined for every $x=\left(\hat{x}, x_{N}\right) \in B$ by $\tau_{\varepsilon} u\left(\hat{x}, x_{N}\right):=u\left(\hat{x}, \varepsilon x_{N}\right)$ and we denote its transposed operator by ${ }^{T} \tau_{\varepsilon}$ :

$$
\left\langle^{T} \tau_{\varepsilon} \theta, u\right\rangle_{\varepsilon}=\left\langle\theta, \tau_{\varepsilon} u\right\rangle, \forall(\theta, u) \in V^{\prime}(B) \times V\left(B_{\varepsilon}\right)
$$

We make the following assumption on the source $\mathcal{S}^{\varepsilon}$ : there exists $\mathcal{S}$ in $V^{\prime}(B)$ such that

$$
S_{\varepsilon}:=\left({ }^{T} \tau_{\varepsilon}\right)^{-1} \mathcal{S}^{\varepsilon} \text { strongly converges to } \mathcal{S} \text { in } V^{\prime}(B)
$$

Then, in order to identify the $\Gamma$-limit of the functional $F_{\varepsilon}$, it will be more convenient to write the functional $F_{\varepsilon}$ as

$$
F_{\varepsilon}(u)=\left\{\begin{array}{l}
\int_{\Omega_{\varepsilon}} f(\nabla u) d x+\varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \tau_{\varepsilon} u, \frac{1}{\varepsilon} \frac{\partial \tau_{\varepsilon} u}{\partial x_{N}}\right) d x-\left\langle\mathcal{S}_{\varepsilon}, \tau_{\varepsilon} u\right\rangle \text { if } u \in W_{\Gamma_{0}}^{1,2}(\Omega) \\
+\infty \text { otherwise. }
\end{array}\right.
$$

## 3 The variational asymptotic model

Let $H: V(B) \longrightarrow \mathbb{R}$ be the functional defined by

$$
\begin{align*}
H(\theta) & :=\int_{B} g^{\infty, 2}\left(\hat{0}, \frac{\partial \theta}{\partial x_{N}}\right) d x-\langle\mathcal{S}, \theta\rangle \\
& =H_{\text {in }}(\theta)-\langle\mathcal{S}, \theta\rangle \tag{3.1}
\end{align*}
$$

We refer the fiunctional $H_{\text {in }}$ as the internal part of $H$. We claim that, when $L^{2}(\Omega)$ is equipped with its strong topology, the functional $F_{\varepsilon} \Gamma$-converges to the functional $F_{0}: L^{2}(\Omega) \longrightarrow \mathbb{R} \cup\{+\infty\}$ given by

$$
F_{0}(u)=\left\{\begin{array}{l}
\int_{\Omega} f(\nabla u) d x+\inf _{\theta \in X(u)} H(\theta) \text { if } u \in W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma) \\
+\infty \text { otherwise },
\end{array}\right.
$$

where $X(u):=\left\{\theta \in V(B): \theta\left(., \pm \frac{1}{2}\right)=u^{ \pm}\right\}$.
Before addressing the variational convergence process, we begin by establishing some compactness properties for sequences with bounded energy. Let us introduce the $\varepsilon$-translate operator $T_{\varepsilon}$ from $W^{1,2}(\Omega)$ into $W^{1,2}(\Omega \backslash \Sigma)$. For any function $w \in W^{1,2}(\Omega), \tilde{w}$ stands for its extension by reflexion on $\Sigma \times(-2 r,-r) \cup$ $(r, 2 r)$ and we define the $\varepsilon$-translate $T_{\varepsilon} w$ of $w$ by

$$
T_{\varepsilon} w\left(\hat{x}, x_{N}\right)= \begin{cases}\widetilde{w}\left(\hat{x}, x_{N}+\frac{\varepsilon}{2}\right) & \text { if } x \in \Omega^{+} \\ \widetilde{w}\left(\hat{x}, x_{N}-\frac{\varepsilon}{2}\right) & \text { if } x \in \Omega^{-}\end{cases}
$$

Lemma 3.1 (compactness). Let $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ be a sequence in $L^{2}(\Omega)$ such that $\sup _{\varepsilon>0} F_{\varepsilon}\left(u_{\varepsilon}\right)<+\infty$. Then

$$
\begin{equation*}
\int_{B_{\varepsilon}}\left|u_{\varepsilon}\right|^{2} d x \leq C \varepsilon\left(\int_{\Omega_{\varepsilon}}\left|\nabla u_{\varepsilon}\right|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{N}}\right|^{2} d x\right) \tag{i}
\end{equation*}
$$

(ii)

$$
\begin{equation*}
\sup _{\varepsilon>0}\left(\int_{\Omega_{\varepsilon}}\left|\nabla u_{\varepsilon}\right|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{N}}\right|^{2} d x\right)<+\infty \tag{3.3}
\end{equation*}
$$

(iii) there exist $u \in W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma)$ and a subsequence of $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ such that $u_{\varepsilon} \rightarrow u$ in $L^{2}(\Omega)$ and $u_{\varepsilon} \rightharpoonup u$ in $W_{\Gamma_{0}}^{1,2}\left(\Omega_{\eta}\right)$ for all $\eta>0$;
(iv) there exist $\theta \in V(B)$ and a subsequence such that $\tau_{\varepsilon} u_{\varepsilon} \rightharpoonup \theta$ in $V(B)$, i.e.

$$
\begin{aligned}
& \tau_{\varepsilon} u_{\varepsilon} \rightharpoonup \theta \text { in } L^{2}(B) \\
& \frac{\partial \tau_{\varepsilon} u_{\varepsilon}}{\partial x_{N}} \rightharpoonup \frac{\partial \theta}{\partial x_{N}} \text { in } L^{2}(B)
\end{aligned}
$$

moreover, $\varepsilon \widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon} \rightharpoonup 0$ in $L^{2}\left(B, \mathbb{R}^{N-1}\right) ;$
(v) $\theta\left(., \pm \frac{1}{2}\right)=u^{ \pm}$.

Proof. Proof of (i). Without loss of generality, we may assume that the $N-1$-dimensional Hausdorff measure of the intersection of $\Gamma_{0}$ with $\left[x_{N}>0\right]$ is positive so that (i) is a mere consequence of the following Poincaré-like inequality:

$$
\begin{equation*}
\exists C>0, \int_{B_{\varepsilon}}|\varphi|^{2} d x \leq C \varepsilon\left(\int_{\Omega_{\varepsilon}^{+}}|\nabla \varphi|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial \varphi}{\partial x_{N}}\right|^{2} d x\right) \forall \varphi \in W_{\Gamma_{0}}^{1,2}(\Omega) . \tag{3.4}
\end{equation*}
$$

Indeed, because

$$
\varphi\left(\hat{x}, x_{N}\right)=T_{\varepsilon} \varphi(\hat{x}, 0)+\int_{\frac{\varepsilon}{2}}^{x_{N}} \frac{\partial}{\partial x_{N}} \varphi(\hat{x}, t) d t \quad \forall x \in B_{\varepsilon}
$$

for all smooth function $\varphi \in W_{\Gamma_{0}}^{1,2}(\Omega)$, we get

$$
\left|\varphi\left(\hat{x}, x_{N}\right)\right|^{2} \leq 2\left(\left|T_{\varepsilon} \varphi(\hat{x}, 0)\right|^{2}+\varepsilon \int_{-\frac{\varepsilon}{2}}^{\frac{\varepsilon}{2}}\left|\frac{\partial}{\partial x_{N}} \varphi(\hat{x}, t)\right|^{2} d t\right) .
$$

Hence, integrating on $B_{\varepsilon}$ and using trace inequality and Poincaré inequality in $\Omega^{+}$give the desired inequality (3.4) for smooth $\varphi$, thus for all $\varphi$ in $W_{\Gamma_{0}}^{1,2}(\Omega)$ by a density argument.

Proof of (ii). From the coercivity conditions satisfied by $f$ and $g$, estimate (3.2), and the strong convergence of $\mathcal{S}_{\varepsilon}$ in $V^{\prime}(B)$, one has

$$
\begin{aligned}
\alpha\left(\int_{\Omega_{\varepsilon}}\left|\nabla u_{\varepsilon}\right|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{N}}\right|^{2} d x\right) & \leq C+\left|\left\langle\mathcal{S}^{\varepsilon}, u_{\varepsilon}\right\rangle_{\varepsilon}\right| \\
& =C+\left|\left\langle\mathcal{S}_{\varepsilon}, \tau_{\varepsilon} u_{\varepsilon}\right\rangle\right| \\
& \leq C+\left\|\mathcal{S}_{\varepsilon}\right\|_{V^{\prime}(B)}\left\|\tau_{\varepsilon} u_{\varepsilon}\right\|_{V(B)} \\
& =C+\left\|\mathcal{S}_{\varepsilon}\right\|_{V^{\prime}(B)}\left(\frac{1}{\varepsilon} \int_{B_{\varepsilon}}\left|u_{\varepsilon}\right|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{N}}\right|^{2} d x\right)^{\frac{1}{2}} \\
& \leq C+C\left(\int_{\Omega_{\varepsilon}}\left|\nabla u_{\varepsilon}\right|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{N}}\right|^{2} d x\right)^{1 / 2} .
\end{aligned}
$$

Then, setting $X_{\varepsilon}:=\left(\int_{\Omega_{\varepsilon}}\left|\nabla u_{\varepsilon}\right|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{N}}\right|^{2} d x\right)^{1 / 2},(3.3)$ follows from the estimate $\alpha X_{\varepsilon}^{2} \leq$ $C+C X_{\varepsilon}$.

Proof of (iii).
Step 1. We claim that there exist $z \in W^{1,2}(\Omega \backslash \Sigma)$ and a subsequence of $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ such that $T_{\varepsilon} u_{\varepsilon} \rightharpoonup z$ in $W^{1,2}(\Omega \backslash \Sigma)$ and strongly in $L^{2}(\Omega \backslash \Sigma)$. Clearly,

$$
\begin{equation*}
T_{\varepsilon} u_{\varepsilon} \in W^{1,2}(\Omega \backslash \Sigma) \text { and } \frac{\partial}{\partial x_{i}} T_{\varepsilon} u_{\varepsilon}=T_{\varepsilon} \frac{\partial}{\partial x_{i}} u_{\varepsilon} \text { for all } \varepsilon>0 \tag{3.5}
\end{equation*}
$$

Combining the Poincaré inequality, (3.3) and (3.5), we deduce

$$
\sup _{\varepsilon>0}\left\|T_{\varepsilon} u_{\varepsilon}\right\|_{W^{1,2}(\Omega \backslash \Sigma)}^{2} \leq C \sup _{\varepsilon>0}\left(\int_{\Omega_{\varepsilon}}\left|\nabla u_{\varepsilon}(x)\right|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{N}}(x)\right|^{2} d x\right)<+\infty .
$$

Therefore, $\left(T_{\varepsilon} u_{\varepsilon}\right)_{\varepsilon>0}$ is bounded in $W^{1,2}(\Omega \backslash \Sigma)$ and the claim follows immediately.
Step 2. We establish that there exists $u$ in $L^{2}(\Omega)$ such that we can extract from the previous subsequence $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ a subsequence strongly converging to $u$ in $L^{2}(\Omega)$. We can write

$$
\int_{\Omega_{\varepsilon}}\left|u_{\varepsilon}(x)\right|^{2} d x=\int_{\Omega^{+} \cup \Omega^{-}}\left|T_{\varepsilon} u_{\varepsilon}(x)\right|^{2} d x-\int_{\Sigma \times\left(\left(r-\frac{\varepsilon}{2}, r\right) \cup\left(-r,-r+\frac{\varepsilon}{2}\right)\right)}\left|T_{\varepsilon} u_{\varepsilon}(x)\right|^{2} d x,
$$

so that

$$
\begin{equation*}
\left\|u_{\varepsilon}\right\|_{L^{2}(\Omega)}^{2}=\int_{\Omega^{+} \cup \Omega^{-}}\left|T_{\varepsilon} u_{\varepsilon}(x)\right|^{2} d x+\int_{B_{\varepsilon}}\left|u_{\varepsilon}(x)\right|^{2} d x-\int_{\Sigma \times\left(\left(r-\frac{\varepsilon}{2}, r\right) \cup\left(-r,-r+\frac{\varepsilon}{2}\right)\right)}\left|T_{\varepsilon} u_{\varepsilon}(x)\right|^{2} d x \tag{3.6}
\end{equation*}
$$

From step 1 and (3.2), we deduce that $\sup _{\varepsilon>0}\left\|u_{\varepsilon}\right\|_{L^{2}(\Omega)}<+\infty$. Thus there exist $u \in L^{2}(\Omega)$ and a not relabelled subsequence such that $u_{\varepsilon} \rightharpoonup u$ in $L^{2}(\Omega)$. Let us prove that $u=z$. Since $u_{\varepsilon} \rightharpoonup u$ in $L^{2}(\Omega)$ and $T_{\varepsilon} u_{\varepsilon} \rightharpoonup z$ in $W^{1,2}(\Omega \backslash \Sigma)$, we have for any $\varphi \in C_{c}^{\infty}(\Omega)$,

$$
\begin{aligned}
\int_{\Omega} u(x) \varphi(x) d x & =\lim _{\varepsilon \rightarrow 0} \int_{\Omega} u_{\varepsilon}(x) \varphi\left(\widehat{x}, x_{N}-\frac{\varepsilon}{2}\right) d x \\
& =\lim _{\varepsilon \rightarrow 0} \int_{\Omega} T_{\varepsilon} u_{\varepsilon}(x) \varphi(x) d x \\
& =\int_{\Omega} z(x) \varphi(x) d x .
\end{aligned}
$$

Thus $u=z$ almost everywhere in $\Omega$ and we deduce that $u \in W^{1,2}(\Omega \backslash \Sigma)$.
Moreover, from (3.2) we have that $\int_{B_{\varepsilon}}\left|u_{\varepsilon}(x)\right|^{2} d x \rightarrow 0$ as $\varepsilon \rightarrow 0$. On the other hand, since $T_{\varepsilon} u_{\varepsilon} \rightarrow z$
in $L^{2}(\Omega)$, we infer $\int_{\Omega^{+} \cup \Omega^{-}}\left|T_{\varepsilon} u_{\varepsilon}(x)\right|^{2} d x \rightarrow \int_{\Omega}|z(x)|^{2} d x$ and $\int_{\Sigma \times\left(\left(r-\frac{\varepsilon}{2}, r\right) \cup\left(-r,-r+\frac{\varepsilon}{2}\right)\right)}\left|T_{\varepsilon} u_{\varepsilon}(x)\right|^{2} d x \rightarrow 0$. Then we deduce that $\left\|u_{\varepsilon}\right\|_{L^{2}(\Omega)} \rightarrow\|z\|_{L^{2}(\Omega)}=\|u\|_{L^{2}(\Omega)}$ and thus that $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ strongly converges to $u$ in $L^{2}(\Omega)$.

Step 3. It remains to establish that for any $\eta>0$, there exists a subsequence of $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ such that $u_{\varepsilon\left\lfloor\Omega_{\eta}\right.} \rightharpoonup u_{\left\lfloor\Omega_{\eta}\right.}$ in $W_{\Gamma_{0}}^{1,2}\left(\Omega_{\eta}\right)$. It will immediately result that $u \in W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma)$.
Let $\eta>0$. Clearly, there exists $0<\varepsilon_{1}<\eta$ such that $\Omega_{\eta} \subseteq \Omega_{\varepsilon}$ for all $\varepsilon \leq \varepsilon_{1}$. By the Poincaré inequality we have

$$
\sup _{\varepsilon>0}\left\|u_{\varepsilon}\right\|_{W^{1,2}\left(\Omega_{\eta}\right)}^{2} \leq C \sup _{\varepsilon>0}\left(\int_{\Omega_{\varepsilon}}\left|\nabla u_{\varepsilon}(x)\right|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{N}}(x)\right|^{2} d x\right)<+\infty
$$

Thus, $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ is bounded in $W_{\Gamma_{0}}^{1,2}\left(\Omega_{\eta}\right)$, and there exist $w \in W_{\Gamma_{0}}^{1,2}\left(\Omega_{\eta}\right)$ and a not relabelled subsequence of $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ satisfying $u_{\varepsilon} \rightarrow w$ in $L^{2}\left(\Omega_{\eta}\right)$ and $u_{\varepsilon} \rightharpoonup w$ in $W_{\Gamma_{0}}^{1,2}\left(\Omega_{\eta}\right)$. It is easily seen that in fact $w=u_{\left\lfloor\Omega_{\eta}\right.}$.

Proof of (iv). The weak convergence of $\tau_{\varepsilon} u_{\varepsilon}$ to some $\theta$ in $V(B)$ follows from (3.2) and (3.3). Indeed

$$
\begin{aligned}
\sup _{\varepsilon>0}\left\|\tau_{\varepsilon} u_{\varepsilon}\right\|_{V(B)} & =\sup _{\varepsilon>0}\left(\frac{1}{\varepsilon} \int_{B_{\varepsilon}}\left|u_{\varepsilon}\right|^{2} d x+\varepsilon \int_{B_{\varepsilon}}\left|\frac{\partial u_{\varepsilon}}{\partial x_{N}}\right|^{2} d x\right)^{\frac{1}{2}} \\
& \leq C \sup _{\varepsilon>0} X_{\varepsilon}<+\infty
\end{aligned}
$$

Now we deduce that $\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon} \rightharpoonup \widehat{\nabla} \theta$ in the distributional sense so that $\varepsilon \widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon} \rightharpoonup 0$ in the distributional sense. On the other hand, from the coercivity of $g, \varepsilon \widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}$ weakly converges to some $L^{2}\left(B, \mathbb{R}^{N-1}\right)$ function. Hence, $\varepsilon \widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon} \rightharpoonup 0$ in $L^{2}\left(B, \mathbb{R}^{N-1}\right)$.

Proof of $(v)$. Note that $\theta\left(., \pm \frac{1}{2}\right)$ is well defined. Indeed, one has

$$
V(B) \subset W^{1,2}\left(\left(-\frac{1}{2}, \frac{1}{2}\right), L^{2}(\Sigma)\right) \subset \mathcal{C}\left(\left[-\frac{1}{2}, \frac{1}{2}\right], L^{2}(\Sigma)\right)
$$

Clearly, $\tau_{\varepsilon} u_{\varepsilon}\left(\hat{x}, \pm \frac{1}{2}\right)=\left(T_{\varepsilon} u_{\varepsilon}\right)^{ \pm}(\hat{x})$ (in the sense of traces on $\Sigma$ of $W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma)$-functions) so that $\tau_{\varepsilon} u_{\varepsilon}\left(\hat{x}, \pm \frac{1}{2}\right) \rightarrow u^{ \pm}$in $L^{2}(\Sigma)$. On the other hand, since

$$
\tau_{\varepsilon} u_{\varepsilon}\left(\hat{x}, x_{N}\right)=\tau_{\varepsilon} u_{\varepsilon}\left(\hat{x}, \pm \frac{1}{2}\right)+\int_{ \pm \frac{1}{2}}^{x_{N}} \frac{\partial \tau_{\varepsilon} u_{\varepsilon}}{\partial x_{N}}(\hat{x}, s) d s
$$

for a.e. $x$ in $B$, we infer that for all $\varphi \in \mathcal{C}_{c}(\Sigma)$,

$$
\begin{equation*}
\int_{-\frac{1}{2}}^{\frac{1}{2}} \int_{\Sigma} \tau_{\varepsilon} u_{\varepsilon}\left(\hat{x}, x_{N}\right) \varphi(\hat{x}) d x=\int_{\Sigma}\left(T_{\varepsilon} u_{\varepsilon}\right)^{ \pm}(\hat{x}) \varphi(\hat{x}) d \hat{x}+\int_{-\frac{1}{2}}^{\frac{1}{2}} \int_{\Sigma} \int_{ \pm \frac{1}{2}}^{x_{N}} \frac{\partial \tau_{\varepsilon} u_{\varepsilon}}{\partial x_{N}}(\hat{x}, s) \varphi(\hat{x}) d s d x \tag{3.7}
\end{equation*}
$$

By passing to the limit in (3.7), we obtain

$$
\int_{-\frac{1}{2}}^{\frac{1}{2}} \int_{\Sigma} \theta\left(\hat{x}, x_{N}\right) \varphi(\hat{x}) d x=\int_{\Sigma} u^{ \pm}(\hat{x}) \varphi(\hat{x}) d \hat{x}+\int_{-\frac{1}{2}}^{\frac{1}{2}} \int_{\Sigma} \int_{ \pm \frac{1}{2}}^{x_{N}} \frac{\partial \theta}{\partial x_{N}}(\hat{x}, s) \varphi(\hat{x}) d s d x
$$

from which we deduce

$$
\int_{\Sigma} u^{ \pm}(\hat{x}) \varphi(\hat{x}) d \hat{x}=\int_{\Sigma} \theta\left(\hat{x}, \pm \frac{1}{2}\right) \varphi(\hat{x}) d \hat{x}
$$

Thus $\theta\left(., \pm \frac{1}{2}\right)=u^{ \pm}$almost everywhere in $\Sigma$.

Lemma 3.2. For every $u \in W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma), \inf _{\theta \in X(u)} H(\theta)>-\infty$ and there exists $\theta(u) \in X(u)$ such that $\inf _{\theta \in X(u)} H(\theta)=H(\theta(u))$.

Proof. The proof follows from standard arguments used in the direct method of the Calculus of Variation.

As a consequence of Lemma 3.2, in its domain $W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma)$, the functional $F_{0}$ may be written

$$
F_{0}(u)=\int_{\Omega} f(\nabla u) d x+H(\theta(u))
$$

Theorem 3.3 is the main result of this section.
Theorem 3.3. The sequence $\left(F_{\varepsilon}\right)_{\varepsilon>0} \Gamma$-converges to the functional $F_{0}$ when $L^{2}(\Omega)$ is equipped with its strong topology.

The proof results from the following two propositions.
Proposition 3.4. For every $u \in L^{2}(\Omega)$ and every $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ strongly converging to $u$ in $L^{2}(\Omega)$ one has

$$
F_{0}(u) \leq \liminf _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(u_{\varepsilon}\right)
$$

Proposition 3.5. For every $u \in L^{2}(\Omega)$ there exists $\left(v_{\varepsilon}\right)_{\varepsilon>0}$ strongly converging to $u$ in $L^{2}(\Omega)$ satisfying

$$
F_{0}(u) \geq \limsup _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(v_{\varepsilon}\right)
$$

Proof of Proposition 3.4. We may assume $\liminf _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(u_{\varepsilon}\right)<+\infty$. From Lemma $3.1 u \in W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma)$ and there exists $\theta \in X(u)$ such that $\tau_{\varepsilon} u_{\varepsilon} \rightharpoonup \theta$ in $V(B)$. Since $\mathcal{S}_{\varepsilon} \rightarrow \mathcal{S}$ in $V^{\prime}(B)$, one has

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0}\left\langle\mathcal{S}_{\varepsilon}, \tau_{\varepsilon} u_{\varepsilon}\right\rangle=\langle\mathcal{S}, \theta\rangle \tag{3.8}
\end{equation*}
$$

On the other hand, since from Lemma 3.1, $u_{\varepsilon} \rightharpoonup u$ in $W_{\Gamma_{0}}^{1,2}\left(\Omega_{\eta}\right)$ for all $\eta>0$, one has

$$
\begin{equation*}
\liminf _{\varepsilon \rightarrow 0} \int_{\Omega_{\varepsilon}} f\left(\nabla u_{\varepsilon}\right) d x \geq \int_{\Omega} f(\nabla u) d x \tag{3.9}
\end{equation*}
$$

Finally from (iv) of Lemma 3.1 and a standard lower semicontinuity argument

$$
\begin{align*}
& \liminf _{\varepsilon \rightarrow 0} \varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right) d x \\
\geq & \liminf _{\varepsilon \rightarrow 0}\left(\varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right) d x-\int_{B} g^{\infty, 2}\left(\varepsilon \widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right) d x\right) \\
& +\liminf _{\varepsilon \rightarrow 0} \int_{B} g^{\infty, 2}\left(\varepsilon \widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right) d x \\
\geq & \liminf _{\varepsilon \rightarrow 0}\left(\varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right) d x-\int_{B} g^{\infty, 2}\left(\varepsilon \widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right) d x\right) \\
& +\int_{B} g^{\infty, 2}\left(\hat{0}, \frac{\partial \theta}{\partial x_{N}}\right) d x \\
= & \int_{B} g^{\infty, 2}\left(\hat{0}, \frac{\partial \theta}{\partial x_{N}}\right) d x \tag{3.10}
\end{align*}
$$

provided that we establish

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0}\left(\varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right) d x-\int_{B} g^{\infty, 2}\left(\varepsilon \widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right) d x\right)=0 \tag{3.11}
\end{equation*}
$$

Since $g^{\infty, 2}$ is positively homogeneous of degree 2, and from (2.1), we have

$$
\begin{aligned}
& \int_{B}\left|\varepsilon^{2} g\left(\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right)-g^{\infty, 2}\left(\varepsilon \widehat{\nabla} \tau_{\varepsilon}^{-1} u_{\varepsilon}, \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right)\right| d x \\
= & \varepsilon^{2} \int_{B}\left|g\left(\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right)-g^{\infty, 2}\left(\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right)\right| d x \\
\leq & C \varepsilon^{2} \int_{B}\left[1+\left|\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}\right|^{2-\delta}+\left|\frac{1}{\varepsilon} \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right|^{2-\delta}\right] d x .
\end{aligned}
$$

Thus, by using Hölder's inequality (take $p=\frac{2}{2-\delta}, q=\frac{2}{\delta}$ ) we deduce

$$
\int_{B}\left|\varepsilon^{2} g\left(\widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial \tau_{\varepsilon} u_{\varepsilon}}{\partial x_{N}}\right)-g^{\infty, 2}\left(\varepsilon \widehat{\nabla} \tau_{\varepsilon} u_{\varepsilon}, \frac{\partial\left(\tau_{\varepsilon} u_{\varepsilon}\right)}{\partial x_{N}}\right)\right| d x \leq C \varepsilon^{\delta}
$$

which proves (3.11). The conclusion of Proposition 3.4 follows by collecting (3.8), (3.9) and (3.10).
Proof of Proposition 3.5. Let $u \in L^{2}(\Omega)$. We have to construct a sequence $\left(v_{\varepsilon}\right)_{\varepsilon>0}$ strongly converging to $u$ in $L^{2}(\Omega)$ such that $\limsup _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(v_{\varepsilon}\right) \leq F_{0}(u)$. If $F_{0}(u)=+\infty$, then $u \in L^{2}(\Omega) \backslash W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma)$, and clearly, for any sequence $\left(v_{\varepsilon}\right)_{\varepsilon>0}$ converging to $u$, $\limsup _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(v_{\varepsilon}\right) \leq F_{0}(u)$ is true. Now, for the harder part, we assume $F_{0}(u)<+\infty$. Then $u \in W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma)$ and

$$
F_{0}(u)=\int_{\Omega} f(\nabla u(x)) d x+\inf _{\theta \in X(u)} H(\theta)
$$

To complete the proof, from $\bar{\theta}:=\theta(u)$, i.e. $H(\bar{\theta})=\inf _{\theta \in X(u)} H(\theta)$, we construct a sequence $\left(v_{\varepsilon}\right)_{\varepsilon>0}$ strongly converging to $u$ in $L^{2}(\Omega)$ and satisfying

$$
F_{0}(u) \geq \limsup _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(v_{\varepsilon}\right)
$$

The proof is divided into four steps:
Step 1. Let us extend $u$ and $\bar{\theta}$ by 0 into $\left(\mathbb{R}^{N-1} \backslash \Sigma\right) \times(-r, r)$ and write these extensions $\widetilde{u}$ and $\widetilde{\bar{\theta}}$. For a sequence $\delta$ of positive numbers intended to go to 0 , consider a standard sequence of molifier $\left(\rho_{\delta}\right)_{\delta}$ and set

$$
\begin{aligned}
& u_{\delta}:=\rho_{\delta} * \widetilde{u} \text { defined by } \rho_{\delta} * \widetilde{u}\left(\widehat{x}, x_{N}\right)=\int_{\mathbb{R}^{N-1}} \rho_{\delta}(\widehat{x}-\widehat{y}) \widetilde{u}\left(\widehat{y}, x_{N}\right) d \widehat{y} \quad \text { for all }\left(\widehat{x}, x_{N}\right) \in \Omega \\
& \theta_{\delta}:=\rho_{\delta} * \tilde{\bar{\theta}} \text { defined by } \rho_{\delta} * \widetilde{\bar{\theta}}\left(\widehat{x}, x_{N}\right)=\int_{\mathbb{R}^{N-1}} \rho_{\delta}(\widehat{x}-\widehat{y}) \widetilde{\bar{\theta}}\left(\widehat{y}, x_{N}\right) d \widehat{y} \quad \text { for all }\left(\widehat{x}, x_{N}\right) \in \Omega
\end{aligned}
$$

Clearly,

$$
\left\{\begin{array}{l}
\theta_{\delta}\left(\widehat{x}, \pm \frac{1}{2}\right)=u_{\delta}(\widehat{x}, 0) \text { for all } \widehat{x} \in \Sigma  \tag{3.12}\\
u_{\delta} \in W^{1,2}(\Omega \backslash \Sigma), \theta_{\delta} \in W^{1,2}(B) \\
u_{\delta} \rightarrow u \text { in } W^{1,2}(\Omega \backslash \Sigma), \theta_{\delta} \rightarrow \bar{\theta} \text { in } V(B)
\end{array}\right.
$$

Next, for each $\delta>0$, we define the sequence $\left(v_{\delta, \varepsilon}\right)_{\varepsilon>0}$ as follows:

$$
v_{\delta, \varepsilon}\left(\widehat{x}, x_{N}\right)=\left\{\begin{array}{ll}
u_{\delta}\left(\widehat{x}, x_{N} \pm \frac{\varepsilon}{2}\right) & \text { on }  \tag{3.13}\\
\Omega_{\varepsilon}^{\mp} \\
\theta_{\delta}\left(\widehat{x}, \frac{x_{N}}{\varepsilon}\right) & \text { on }
\end{array} B_{\varepsilon}\right.
$$

Obviously $v_{\delta, \varepsilon}\left(\widehat{x}, x_{N}\right)$ belongs to $W^{1,2}(\Omega)$ and strongly converges to $u_{\delta}$ in $L^{2}(\Omega)$.
Step 2. We we claim that

$$
\begin{align*}
& \lim _{\varepsilon \rightarrow 0} \int_{\Omega_{\varepsilon}} f\left(\nabla v_{\delta, \varepsilon}\right)(x) d x=\int_{\Omega} f\left(\nabla u_{\delta}\right)(x) d x  \tag{3.14}\\
& \lim _{\varepsilon \rightarrow 0}\left(\varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \tau_{\varepsilon} v_{\delta, \varepsilon}, \frac{1}{\varepsilon} \frac{\partial \tau_{\varepsilon} v_{\delta, \varepsilon}}{\partial x_{N}}\right)(x) d x-\left\langle\mathcal{S}_{\varepsilon}, \tau_{\varepsilon} v_{\delta, \varepsilon}\right\rangle\right)=H\left(\theta_{\delta}\right) \tag{3.15}
\end{align*}
$$

Proof of (3.14): one has

$$
\begin{aligned}
\lim _{\varepsilon \rightarrow 0} \int_{\Omega_{\varepsilon}} f\left(\nabla v_{\delta, \varepsilon}\right)(x) d x & =\lim _{\varepsilon \rightarrow 0}\left(\int_{\Omega_{\varepsilon}^{+}} f\left(\nabla u_{\delta}\right)\left(\widehat{x}, x_{N}-\frac{\varepsilon}{2}\right) d x+\int_{\Omega_{\varepsilon}^{-}} f\left(\nabla u_{\delta}\left(\widehat{x}, x_{N}+\frac{\varepsilon}{2}\right)\right) d x\right) \\
& =\int_{\Omega^{+}} f\left(\nabla u_{\delta}\right)(x) d x+\int_{\Omega^{-}} f\left(\nabla u_{\delta}\right)(x) d x \\
& =\int_{\Omega^{-}} f\left(\nabla u_{\delta}\right)(x) d x .
\end{aligned}
$$

Proof of (3.15): Since $g^{\infty, 2}$ is positively homogeneous of degree 2 and $\mathcal{S}_{\varepsilon}$ strongly converges to $\mathcal{S}$ in $V^{\prime}(B)$, one has

$$
\lim _{\varepsilon \rightarrow 0}\left(\varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \theta_{\delta}, \frac{1}{\varepsilon} \frac{\partial \theta_{\delta}}{\partial x_{N}}\right)(x) d x-\left\langle\mathcal{S}_{\varepsilon}, \theta_{\delta}\right\rangle\right)=\int_{B} g^{\infty, 2}\left(\widehat{0}, \frac{\partial \theta_{\delta}}{\partial x_{N}}\right) d x-\left\langle\mathcal{S}, \theta_{\delta}\right\rangle=H\left(\theta_{\delta}\right)
$$

Step 3. We establish that $\lim _{\delta \rightarrow 0} \int_{\Omega} f\left(\nabla u_{\delta}\right) d x+H\left(\theta_{\delta}\right)=F_{0}(u)$. Since

$$
\int_{\Omega} f\left(\nabla u_{\delta}\right) d x+H\left(\theta_{\delta}\right)=\int_{\Omega} f\left(\nabla u_{\delta}\right) d x+\int_{B} g^{\infty, 2}\left(\widehat{0}, \frac{\partial \theta_{\delta}}{\partial x_{N}}\right) d x-\left\langle\mathcal{S}, \theta_{\delta}\right\rangle
$$

the result is a straightforward consequence of (3.12).
Step 4. By using a standard diagonalization argument, from step 2 and step 3, there exists a mapping $\varepsilon \mapsto \delta(\varepsilon)$ such that $v_{\delta(\varepsilon)} \longrightarrow u$ in $L^{2}(\Omega)$ and

$$
\lim _{\varepsilon \rightarrow 0}\left(\int_{\Omega} f\left(\nabla v_{\delta(\varepsilon)}\right)(x) d x+\varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \tau_{\varepsilon} v_{\delta(\varepsilon)}, \frac{1}{\varepsilon} \frac{\partial \tau_{\varepsilon} v_{\delta(\varepsilon)}}{\partial x_{N}}\right)(x) d x-\left\langle\mathcal{S}_{\varepsilon}, \tau_{\varepsilon} v_{\delta(\varepsilon)}\right\rangle\right)=F_{0}(u)
$$

The sequence $\left(v_{\varepsilon}\right)_{\varepsilon>0}$ where $v_{\varepsilon}:=v_{\delta(\varepsilon)}$ fullfils all the conditions except the boundary condition on $\Gamma_{0}$. From assumption $\operatorname{dist}\left(\bar{\Gamma}_{0}, \overline{\partial B_{\varepsilon} \cap \partial \Omega}\right)>0$, and by using a standard slicing method due to De Giorgi in a neighborhood of $\Gamma_{0}$ (see [4]), one can modify $v_{\varepsilon}$ in $\Omega_{\varepsilon}$ into a function $\tilde{v}_{\varepsilon}$ equal to $v_{\varepsilon}$ in $B_{\varepsilon}$, satisfying the boundary condition on $\Gamma_{0}$, and $\lim \sup _{\varepsilon \rightarrow 0} \int_{\Omega_{\varepsilon}} f\left(\nabla v_{\varepsilon}\right) d x=\limsup _{\varepsilon \rightarrow 0} \int_{\Omega_{\varepsilon}} f\left(\nabla \tilde{v}_{\varepsilon}\right) d x$. Still denoting by $v_{\varepsilon}$ this new function, we have $\lim _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(v_{\varepsilon}\right)=F_{0}(u)$ and the proof is complete.

Remark 3.6. In order to give an interpretation of the limit energy functional, it is worthwhile to write

$$
\begin{equation*}
\inf _{\theta \in X(u)} H(\theta)=\inf _{\theta \in V_{0}(B)}\left\{\int_{B} g^{\infty, 2}\left(\hat{0}, \frac{\partial \theta}{\partial x_{N}}(x)+[u](\hat{x})\right) d x-\langle\mathcal{S}, \theta\rangle\right\}-\langle\mathcal{S}, \tilde{u}\rangle \tag{3.16}
\end{equation*}
$$

where $[u]=u^{+}-u^{-}, V_{0}(B)=\left\{\theta \in V(B): \theta=0\right.$ on $\left.\Sigma \times\left\{ \pm \frac{1}{2}\right\}\right\}$ and $\tilde{u}(x)=x_{N}[u](\hat{x})+\frac{u^{+}(\hat{x})+u^{-}(\hat{x})}{2}$. Thererfore when the limit source $\mathcal{S}$ vanishes on $V(B)$, by using Jensen's inequality, $\inf _{\theta \in X(u)} H(\theta)$ reduces to

$$
\bar{H}(u)=\int_{\Sigma} g^{\infty, 2}(\hat{0},[u](\hat{x})) d \hat{x}
$$

which is nothing but the surface energy of the model obtained in [9]. When the limit source is not trivial, by using the Euler equation associated with (3.16), it is easily seen that $\bar{H}$ is a surface energy on $\Sigma$ of the form

$$
\bar{H}(u)=\int_{\Sigma} \bar{h}\left(\hat{x},[u](\hat{x}), \frac{u^{+}+u^{-}}{2}(\hat{x})\right) d \hat{x} .
$$

In this case we note that the energy density depends explicitly of the mean of the traces and that the surface energy $\bar{H}$ mixes the internal energy and the work of the loading.

## 4 Examples of measure sources $\mathcal{S}_{\varepsilon}$ concentrated in $B_{\varepsilon}$

The general form of elements of $V^{\prime}(B)$ is given for every $\theta$ in $V(B)$ by $\langle\mathcal{S}, \theta\rangle=\int_{B} s_{0} \theta d x+\int_{B} s_{1} \frac{\partial \theta}{\partial x_{N}} d x$ where $\left(s_{0}, s_{1}\right) \in L^{2}(B) \times L^{2}(B)$. The limit sources $\mathcal{S}$ considered in this section are generated by measures $\mathcal{S}^{\varepsilon}$ in $\mathbb{M}\left(B_{\varepsilon}\right)$ whose slicing structure $\mathcal{H}^{N-1}\left\lfloor\Sigma \otimes \mathcal{S}_{\hat{x}}^{\varepsilon}\right.$ is such that their slicing components $\mathcal{S}_{\hat{x}}^{\varepsilon}$ do not present a diffuse singular part in their Lebesgue-Nikodym decomposition in $\mathbb{M}\left(-\frac{\varepsilon}{2}, \frac{\varepsilon}{2}\right)$, i.e., are of the general form

$$
\mathcal{S}_{\hat{x}}^{\varepsilon}=\frac{1}{\varepsilon} a_{\varepsilon}(\hat{x}, \varepsilon t) d t+\sum_{n=-\infty}^{+\infty} b_{\varepsilon, n}(\hat{x}) \delta_{\varepsilon t_{n}(\hat{x})}
$$

where

$$
\left\{\begin{array}{l}
a_{\varepsilon} \in L^{2}(B), b_{\varepsilon, n} \in L^{2}(\Sigma) \\
t_{n}: \Sigma \longrightarrow\left(-\frac{1}{2}, \frac{1}{2}\right) \text { is a Borel measurable map. }
\end{array}\right.
$$

Roughly, such sources $\mathcal{S}^{\varepsilon}$ are sums of a function in $L^{2}(B)$ and a countable sum of surface sources, each of them being concentrated in the $N$-1-dimensional surface included in $B_{\varepsilon}$ whose graph is $\varepsilon t_{n}$. We make the following additional assumptions:
(H1) there exists $a \in L^{2}(B)$ such that $a_{\varepsilon} \rightarrow a$ in $L^{2}(B)$;
(H2) there exists $b_{n} \in L^{2}(\Sigma)$ such that $b_{\varepsilon, n} \rightarrow b_{n}$ in $L^{2}(\Sigma)$ when $\varepsilon \rightarrow 0$;
(H3) there exists $c_{n} \in \mathbb{R}^{+}$such that $\left\|b_{\varepsilon, n}\right\|_{L^{2}(\Sigma)} \leq c_{n}$ and $\sum_{n=-\infty}^{+\infty} c_{n}<+\infty$;
It is easy to check that the measure $\mathcal{S}_{\varepsilon}={ }^{T} \tau_{\varepsilon} \mathcal{S}^{\varepsilon}$ of $\mathbb{M}(B)$ is given by: $\mathcal{S}_{\varepsilon}=\mathcal{H}^{N-1}\left\lfloor\Sigma \otimes\left(\mathcal{S}_{\varepsilon}\right)_{\hat{x}}\right.$ where

$$
\left(\mathcal{S}_{\varepsilon}\right)_{\hat{x}}=a_{\varepsilon}(\hat{x}, t) d t+\sum_{n=-\infty}^{+\infty} b_{\varepsilon, n}(\hat{x}) \delta_{t_{n}(\hat{x})}
$$

Proposition 4.1. The measure $\mathcal{S}_{\varepsilon}$ strongly converges in $V^{\prime}(B)$ to the measure $\mathcal{S}$ defined for every $\theta \in V(B)$ by

$$
\langle\mathcal{S}, \theta\rangle=\int_{B} a(x) \theta(x) d x+\sum_{n=-\infty}^{+\infty} \int_{\Sigma} b_{n}(\hat{x}) \theta\left(\hat{x}, t_{n}(\hat{x})\right) d \hat{x}
$$

Therefore, the functional $F_{\varepsilon} \Gamma$-converges to the functional $F_{0}: L^{2}(\Omega) \longrightarrow \mathbb{R} \cup\{+\infty\}$ given by

$$
F_{0}(u)=\left\{\begin{array}{l}
\int_{\Omega} f(\nabla u) d x+\inf _{\theta \in X(u)}\left\{\int_{B} g^{\infty, 2}\left(\hat{0}, \frac{\partial \theta}{\partial x_{N}}\right) d x-\int_{B} a \theta d x-\sum_{n=-\infty}^{+\infty} \int_{\Sigma} b_{n}(\hat{x}) \theta\left(\hat{x}, t_{n}(\hat{x})\right) d \hat{x}\right\} \\
+\infty \text { otherwise. }
\end{array}\right.
$$

Proof. The second assertion is a straightforward consequence of Theorem 3.3 provided that we establish the strong convergence of $\mathcal{S}_{\varepsilon}$ to $\mathcal{S}$ in $V^{\prime}(B)$. For every $\theta \in V(B)$ we have

$$
\left\langle\mathcal{S}_{\varepsilon}-\mathcal{S}, \theta\right\rangle=\int_{B}\left(a_{\varepsilon}-a\right) \theta d x+\int_{\Sigma} \sum_{n=-\infty}^{+\infty}\left(b_{\varepsilon, n}-b_{n}\right) \theta\left(\hat{x}, t_{n}(\hat{x})\right) d \hat{x}
$$

thus

$$
\begin{equation*}
\left|\left\langle\mathcal{S}_{\varepsilon}-\mathcal{S}, \theta\right\rangle\right| \leq\|\theta\|_{L^{2}(B)}\left\|a_{\varepsilon}-a\right\|_{L^{2}(B)}+\sum_{n=-\infty}^{+\infty}\left[\left\|b_{\varepsilon, n}-b_{n}\right\|_{L^{2}(\Sigma)}\left(\int_{\Sigma} \left\lvert\, \theta\left(\hat{x},\left.t_{n}(\hat{x})\right|^{2} d \hat{x}\right)^{\frac{1}{2}}\right.\right] .\right. \tag{4.1}
\end{equation*}
$$

But it is easy to establish that there exists a non negative constant $C$ such that

$$
\left(\int_{\Sigma} \left\lvert\, \theta\left(\hat{x},\left.t_{n}(\hat{x})\right|^{2} d \hat{x}\right)^{\frac{1}{2}} \leq C\|\theta\|_{V(B)}\right.\right.
$$

so that (4.1) yields

$$
\left\|\mathcal{S}_{\varepsilon}-\mathcal{S}\right\|_{V^{\prime}(B)} \leq\left\|a_{\varepsilon}-a\right\|_{L^{2}(B)}+C \sum_{n=-\infty}^{+\infty}\left\|b_{\varepsilon, n}-b_{n}\right\|_{L^{2}(\Sigma)}
$$

The conclusion follows from assumptions $(H 1),(H 2)$ and $(H 3)$.

## 5 The gradient concentration phenomenon

We first recall the notion of gradient Young-concentration measure introduced in [3]. Let us denote the unit sphere $\{-1,1\}$ of $\mathbb{R}$ by $\mathbb{S}^{0}$, and consider $\Sigma^{\prime} \subset \subset \Sigma, B_{\varepsilon}^{\prime}:=\Sigma^{\prime} \times\left(-\frac{\varepsilon}{2}, \frac{\varepsilon}{2}\right)$.
Definition 5.1. A pair $\left(v, \mu_{\Sigma^{\prime}}\right) \in L^{2}(\Omega) \times \mathbb{M}^{+}\left(\bar{\Omega} \times \mathbb{S}^{0}\right)$ is a gradient Young-concentration measure (localized on $\Sigma^{\prime}$ ) iff there exists a sequence $\left(v_{\varepsilon}\right)_{\varepsilon>0}$ in $W_{\Gamma_{0}}^{1,2}(\Omega)$ satisfying

$$
\left\{\begin{array}{l}
\sup _{\varepsilon>0} \int_{\Omega \backslash B_{\varepsilon}}\left|\nabla v_{\varepsilon}\right|^{2} d x<+\infty \\
v_{\varepsilon} \rightarrow v \text { in } L^{2}(\Omega), \\
\mu_{\varepsilon}:=\delta_{\frac{\partial v_{\varepsilon}}{\partial x_{N}} /\left|\frac{\partial v_{\varepsilon}}{\partial x_{N}}\right|(x)} \otimes \varepsilon \mathbb{1}_{B_{\varepsilon}^{\prime}}\left|\frac{\partial v_{\varepsilon}}{\partial x_{N}}\right|^{2} d x \stackrel{*}{\rightharpoonup} \mu_{\Sigma^{\prime}}
\end{array}\right.
$$

We say that the sequence $\left(v_{\varepsilon}\right)_{\varepsilon>0}$ generates the gradient Young-concentration measure $\left(v, \mu_{\Sigma^{\prime}}\right)$. We denote the set of gradient Young-concentration measures localized on $\Sigma^{\prime}$ by $\mathcal{Y C}\left(\Sigma^{\prime}\right)$.

Recall that the weak convergence $\stackrel{*}{\rightharpoonup}$ above is defined by

$$
\int_{\bar{\Omega}} \int_{\mathbb{S}^{0}} \theta(x) \varphi(\zeta) d \mu_{\varepsilon}=\int_{B_{\varepsilon}^{\prime}} \varepsilon \theta(x) \tilde{\varphi}\left(\frac{\partial v_{\varepsilon}}{\partial x_{N}}\right) d x \rightarrow \int_{\bar{\Omega}} \int_{\mathbb{S}^{0}} \theta(x) \varphi(\zeta) d \mu_{\Sigma^{\prime}}
$$

for all $\theta \in \mathcal{C}(\bar{\Omega})$ and all $\varphi \in \mathcal{C}\left(\mathbb{S}^{0}\right)$, where the 2-homogeneous extension $\tilde{\varphi}: \mathbb{R} \rightarrow \mathbb{R}$ of $\varphi \in \mathcal{C}\left(\mathbb{S}^{0}\right)$ is defined for all $\zeta \in \mathbb{R}^{m}$ by

$$
\tilde{\varphi}(\zeta)= \begin{cases}|\zeta|^{2} \varphi\left(\frac{\zeta}{|\zeta|}\right), & \text { if } \zeta \neq 0 \\ 0 & \text { otherwise }\end{cases}
$$

In [3], Theorem 3.1, the gradient Young-concentration measures was characterized as follows.
Theorem 5.2 (Characterization). A pair $\left(v, \mu_{\Sigma^{\prime}}=\mu_{x} \otimes \pi\right)$ belongs to $\mathcal{Y C}\left(\Sigma^{\prime}\right)$ if and only if $v \in W_{\Gamma_{0}}^{1,2}(\Omega \backslash$ $\Sigma), \pi$ is concentrated on $\bar{\Sigma}^{\prime}$ and, for every $\varphi \in \mathcal{C}\left(\mathbb{S}^{0}\right)$ such that $\varphi^{* *}>-\infty$,

$$
\begin{align*}
& \frac{d \pi}{d \mathcal{H}^{N-1}\left\lfloor\Sigma^{\prime}\right.}(x) \int_{\mathbb{S}^{0}} \varphi(\zeta) d \mu_{x} \geq \varphi^{* *}([v](x)) \quad \text { for } \mathcal{H}^{N-1} \text { a. e. } x \in \Sigma^{\prime}  \tag{5.1}\\
& \int_{\mathbb{S}^{0}} \varphi(\zeta) d \mu_{x} \geq 0 \quad \text { for } \pi_{s} \text { a. e. } x \in \bar{\Sigma}^{\prime}
\end{align*}
$$

where $\pi=\frac{d \pi}{d \mathcal{H}^{N-1}\left[\Sigma^{\prime}\right.} \mathcal{H}^{N-1}\left\lfloor\Sigma^{\prime}+\pi_{s}\right.$ is the Radon-Nikodym decomposition of $\pi$ with respect to the measure $\mathcal{H}^{N-1}\left\lfloor\Sigma^{\prime}\right.$.

Remark 5.3. Although from (3.3), $\delta_{\frac{\partial v_{\varepsilon}}{\partial x_{N}} /\left|\frac{\partial v_{\varepsilon}}{\partial x_{N}}\right|(x)} \otimes \varepsilon \mathbb{1}_{B_{\varepsilon}}\left|\frac{\partial v_{\varepsilon}}{\partial x_{N}}\right|^{2} d x$ possesses weak cluster points in the sense of the weak convergence $\stackrel{*}{\rightharpoonup}$ made precise above, for technical reason (proof of the sufficient conditions in Proposition 3.5 in [3]), it was not possible to state such a characterization for these cluster points because of possible concentration effects on the boundary of $\Sigma$. This is the reason why we deal with gradient Young-concentration measures localized on $\Sigma^{\prime} \subset \subset \Sigma$.

Taking into account that the 2-homogeneous extension $\tilde{\varphi}: \mathbb{R} \rightarrow \mathbb{R}$ of $\varphi \in \mathcal{C}\left(\mathbb{S}^{0}\right)$ satisfying $\varphi^{* *}>-\infty$ is of the form

$$
\varphi(\zeta)=\left\{\begin{array}{l}
c \zeta^{2} \text { if } \zeta \geq 0 \\
d \zeta^{2} \text { if } \zeta \leq 0
\end{array}\right.
$$

with $(c, d) \in \mathbb{R}^{+} \times \mathbb{R}^{+}$, the above characterization theorem can be reduced to the following (cf Corollary 3.6 in [3])

Corollary 5.4. A measure $\left(v, \mu_{\Sigma^{\prime}}=\left(a(x) \delta_{1}+b(x) \delta_{-1}\right) \otimes \pi\right)$ belongs to $\mathcal{Y C}\left(\Sigma^{\prime}\right)$ if and only if $v \in$ $W_{\Gamma_{0}}^{1,2}(\Omega \backslash \Sigma), \pi$ is concentrated on $\bar{\Sigma}^{\prime}$ and

$$
\frac{d \pi}{d \mathcal{H}^{N-1}\left[\Sigma^{\prime}\right.}(x)(a(x) c+b(x) d) \geq \varphi([v](x)) \text { for } \mathcal{H}^{N-1}\left[\Sigma^{\prime} \text { a.e. } x \text { and for all }(c, d) \in \mathbb{R}^{+} \times \mathbb{R}^{+}\right.
$$

where $\varphi(\zeta)=\left\{\begin{array}{l}c \zeta^{2} \text { if } \zeta \geq 0 \\ d \zeta^{2} \text { if } \zeta \leq 0\end{array}\right.$.
As stated in [3] Remark 2.5, every sequence $\left(u_{\varepsilon}\right)_{\varepsilon>0}$ satisfying (3.3) generates a gradient Youngconcentration measure. Therefore every sequence $\left(\bar{u}_{\varepsilon}\right)_{\varepsilon>0}, \bar{u}_{\varepsilon} \in \operatorname{argmin} F_{\varepsilon}$, generates a measure $\bar{\mu}_{\Sigma^{\prime}} \in$ $\mathcal{Y C}\left(\Sigma^{\prime}\right)$. Let $\bar{u}$ be a strong limit of $\left(\bar{u}_{\varepsilon}\right)_{\varepsilon>0}$ in $L^{2}(\Omega)$, then, under the condition $g^{\infty, 2}\left(\hat{\xi}, \xi_{3}\right) \geq g^{\infty, 2}\left(\hat{0}, \xi_{3}\right)$, the next theorem states that the internal term $H_{\text {in }}(\bar{\theta})(\operatorname{cf}(3.1))$ where $\bar{\theta}$ is the solution of $\inf _{\theta \in X(\bar{u})} H(\theta)$, possesses an integral representation with respect to the Young-concentration measure $\bar{\mu}_{\Sigma^{\prime}}$. In some sense we localize $H_{\text {in }}$ on $S \times\{ \pm 1\}$. Moreover, by using Theorem 5.2 we will deduce some bounds on $\bar{\mu}_{\Sigma^{\prime}}$.
Theorem 5.5. Let $\bar{u}_{\varepsilon}$ be a minimizer of $\min \left\{F_{\varepsilon}(v): v \in L^{2}(\Omega)\right\}$ and, for every $\Sigma^{\prime} \subset \subset \Sigma$, ( $\left.\bar{u}, \bar{\mu}_{\Sigma^{\prime}}\right)$ be a gradient Young-concentration measure localized on $\Sigma^{\prime}$ generated by the sequence $\left(\bar{u}_{\varepsilon}\right)_{\varepsilon>0}$. Then the two following assertions hold:
i) $\bar{u}_{\varepsilon} \rightarrow \bar{u}$ in $L^{2}(\Omega), F_{\varepsilon}\left(\bar{u}_{\varepsilon}\right) \rightarrow F_{0}(\bar{u})=\min \left\{F_{0}(u): u \in L^{2}(\Omega)\right\}$;
ii) Let $\mathcal{F}$ be a countable familly of $\Sigma^{\prime} \subset \subset \Sigma$, then there exists $\bar{\mu} \in \mathbb{M}\left(\bar{\Omega} \times \mathbb{S}^{0}\right), \bar{\mu}=\bar{\mu}_{\hat{x}} \otimes \bar{\pi}$ with $\bar{\pi}$ concentrated on $\bar{\Sigma}$ such that for all $\Sigma^{\prime} \in \mathcal{F}, \bar{\mu}\left[\bar{\Sigma}^{\prime} \times \mathbb{S}^{0}=\bar{\mu}_{\Sigma^{\prime}}\right.$. Assume furthermore that $g^{\infty, 2}$ satisfies the condition

$$
\begin{equation*}
\forall \xi \in \mathbb{R}^{3}, g^{\infty, 2}\left(\hat{\xi}, \xi_{3}\right) \geq g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) . \tag{5.2}
\end{equation*}
$$

Then, every weak cluster point $\bar{\theta}$ of the sequence $\left(\tau_{\varepsilon} \bar{u}_{\varepsilon}\right)_{\varepsilon>0}$ in $V(B)$ satisfies $H(\bar{\theta})=\inf _{\theta \in X(\bar{u})} H(\theta)$ and

$$
\begin{align*}
& \int_{-\frac{1}{2}}^{\frac{1}{2}} g^{\infty, 2}\left(\hat{0}, \frac{\partial \bar{\theta}}{\partial x_{N}}\right)(\hat{x}, s) d s=\frac{d \bar{\pi}}{d \hat{x}}(\hat{x}) \int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}} \text { for a.e. } \hat{x} \text { in } \Sigma^{\prime} ;  \tag{5.3}\\
& H_{\text {in }}(\bar{\theta})=\int_{\Sigma}\left[\frac{d \bar{\pi}}{d \hat{x}}(\hat{x}) \int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}}\right] d \hat{x} .
\end{align*}
$$

Proof. According to the variational nature of the $\Gamma$-convergence, for a subsequence one has

$$
\begin{align*}
& \bar{u}_{\varepsilon} \rightarrow \bar{u} \text { in } L^{2}(\Omega) \\
& \lim _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(\bar{u}_{\varepsilon}\right)=F_{0}(\bar{u})=\min \left\{F_{0}(v): v \in L^{2}(\Omega)\right\} \\
& \quad=\int_{\Omega} f(\nabla \bar{u}) d x+\inf _{\theta \in X(\bar{u})} H(\theta) . \tag{5.4}
\end{align*}
$$

Fix $\Sigma^{\prime} \subset \subset \Sigma$. From (3.3), for the subsequence (possibly dependent on $\Sigma^{\prime}$ ) associated with the gradient Young-concentration measure ( $\bar{u}, \bar{\mu}_{\Sigma^{\prime}}$ ), there exist a subsequence and a measure $\bar{\mu}=\bar{\mu}_{\hat{x}} \otimes \bar{\pi}$ in $\mathbb{M}\left(\bar{\Omega} \times \mathbb{S}^{0}\right)$ with $\bar{\pi}$ concentrated in $\bar{\Sigma}$, such that

Thus, from (3.11) and (5.2) we infer

$$
\begin{align*}
\lim _{\varepsilon \rightarrow 0} \varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \tau_{\varepsilon} \bar{u}_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial \tau_{\varepsilon} \bar{u}_{\varepsilon}}{\partial x_{N}}\right) d x & =\lim _{\varepsilon \rightarrow 0} \int_{B} g^{\infty, 2}\left(\varepsilon \widehat{\nabla} \tau_{\varepsilon} \bar{u}_{\varepsilon}, \frac{\partial \tau_{\varepsilon} \bar{u}_{\varepsilon}}{\partial x_{N}}\right) d x \\
& \geq \lim _{\varepsilon \rightarrow 0} \int_{B_{\varepsilon}} g^{\infty, 2}\left(\hat{0}, \frac{\partial \bar{u}_{\varepsilon}}{\partial x_{N}}\right) d x \\
& =\int_{\bar{\Sigma}}\left(\int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}}\right) d \bar{\pi} . \tag{5.5}
\end{align*}
$$

Let $\bar{\theta}$ be the weak limit of $\tau_{\varepsilon} \bar{u}_{\varepsilon}$ in $V(B)$ for the considered subsequence. Then, from (5.5) and since

$$
\liminf _{\varepsilon \rightarrow 0} \int_{\Omega_{\varepsilon}} f\left(\nabla \bar{u}_{\varepsilon}\right) d x \geq \int_{\Omega} f(\nabla \bar{u}) d x \text { and } \lim _{\varepsilon \rightarrow 0}\left\langle\mathcal{S}_{\varepsilon}, \tau_{\varepsilon} \bar{u}_{\varepsilon}\right\rangle=\langle\mathcal{S}, \bar{\theta}\rangle,
$$

we infer

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(\bar{u}_{\varepsilon}\right) \geq \int_{\Omega} f(\nabla \bar{u}) d x+\int_{\bar{\Sigma}}\left(\int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}}\right) d \bar{\pi}-\langle\mathcal{S}, \bar{\theta}\rangle . \tag{5.6}
\end{equation*}
$$

Collecting (5.4) and (5.6) we obtain

$$
\int_{\Omega} f(\nabla \bar{u}) d x+\inf _{\theta \in X(\bar{u})} H(\theta) \geq \int_{\Omega} f(\nabla \bar{u}) d x+\int_{\bar{\Sigma}}\left(\int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}}\right) d \bar{\pi}-\langle\mathcal{S}, \bar{\theta}\rangle
$$

in particular

$$
\int_{\Omega} f(\nabla \bar{u}) d x+H(\bar{\theta}) \geq \int_{\Omega} f(\nabla \bar{u}) d x+\int_{\bar{\Sigma}}\left(\int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}}\right) d \bar{\pi}-\langle\mathcal{S}, \bar{\theta}\rangle
$$

thus

$$
\begin{align*}
\int_{B} g^{\infty, 2}\left(\hat{0}, \frac{\partial \bar{\theta}}{\partial x_{N}}\right) d x & \geq \int_{\bar{\Sigma}}\left(\int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}}\right) d \bar{\pi} \\
& \geq \int_{\bar{\Sigma}} \frac{d \bar{\pi}}{d \hat{x}}(\hat{x})\left(\int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}}\right) d \hat{x} \tag{5.7}
\end{align*}
$$

On the other hand, by a standard lower semicontinuity argument, for every $\varphi \in \mathcal{C}_{c}(\Sigma), \varphi \geq 0$,

$$
\begin{aligned}
\liminf _{\varepsilon \rightarrow 0} \int_{B} \varphi(\hat{x}) g^{\infty, 2}\left(\hat{0}, \frac{\partial \tau_{\varepsilon} \bar{u}_{\varepsilon}}{\partial x_{N}}\right) d x & =\int_{\Sigma} \varphi(\hat{x})\left(\int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}}\right) d \bar{\pi} \\
& \geq \int_{B} \varphi(\hat{x}) g^{\infty, 2}\left(\hat{0}, \frac{\partial \bar{\theta}}{\partial x_{N}}\right) d x
\end{aligned}
$$

so that

$$
\begin{equation*}
\int_{-\frac{1}{2}}^{\frac{1}{2}} g^{\infty, 2}\left(\hat{0}, \frac{\partial \bar{\theta}}{\partial x_{N}}\right)(\hat{x}, s) d s \leq \frac{d \bar{\pi}}{d \hat{x}}(\hat{x}) \int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}} \quad \text { for a.e. } \hat{x} \in \Sigma . \tag{5.8}
\end{equation*}
$$

Combining (5.7) and (5.8) we deduce

$$
\int_{-\frac{1}{2}}^{\frac{1}{2}} g^{\infty, 2}\left(\hat{0}, \frac{\partial \bar{\theta}}{\partial x_{N}}\right)(\hat{x}, s) d s=\frac{d \bar{\pi}}{d \hat{x}}(\hat{x}) \int_{\mathbb{S}^{0}} g^{\infty, 2}\left(\hat{0}, \xi_{3}\right) d \bar{\mu}_{\hat{x}} \quad \text { for a.e. } \hat{x} \in \Sigma .
$$

Clearly, $\bar{\mu}\left\lfloor\bar{\Sigma}^{\prime} \times \mathbb{S}^{0}=\bar{\mu}_{\Sigma^{\prime}}\right.$. Now, by using a standard Cantor's diagonal process, the same equality holds for all $\Sigma^{\prime}$ of the countable familly $\mathcal{F}$. It remains to show that $H(\bar{\theta})=\inf _{\theta \in X(\bar{u})} H(\theta)$. It's enough to notice that

$$
\begin{aligned}
\lim _{\varepsilon \rightarrow 0} F_{\varepsilon}\left(\bar{u}_{\varepsilon}\right) & =\int_{\Omega} f(\nabla \bar{u}) d x+\inf _{\theta \in X(\bar{u})} H(\theta) \\
& \geq \liminf _{\varepsilon \rightarrow 0} \int_{\Omega_{\varepsilon}} f\left(\nabla \bar{u}_{\varepsilon}\right) d x+\liminf _{\varepsilon \rightarrow 0}\left(\varepsilon^{2} \int_{B} g\left(\widehat{\nabla} \tau_{\varepsilon} \bar{u}_{\varepsilon}, \frac{1}{\varepsilon} \frac{\partial \tau_{\varepsilon} \bar{u}_{\varepsilon}}{\partial x_{N}}\right) d x-\left\langle\mathcal{S}_{\varepsilon}, \tau_{\varepsilon} \bar{u}_{\varepsilon}\right\rangle\right) \\
& \geq \int_{\Omega} f(\nabla \bar{u}) d x+\int_{B} g^{\infty, 2}\left(\hat{0}, \frac{\partial \bar{\theta}}{\partial x_{N}}\right) d x-\langle\mathcal{S}, \bar{\theta}\rangle \\
& =\int_{\Omega} f(\nabla \bar{u}) d x+H(\bar{\theta})
\end{aligned}
$$

which completes the proof.
We define the following two constants associated with the function $g$ :

$$
c(g):=\min \left(\frac{g^{\infty, 2}(\widehat{0},-1)}{g^{\infty, 2}(\widehat{0}, 1)}, \frac{g^{\infty, 2}(\widehat{0}, 1)}{g^{\infty, 2}(\widehat{0},-1)}\right), C(g)=\frac{1}{c(g)}=\max \left(\frac{g^{\infty, 2}(\widehat{0},-1)}{g^{\infty, 2}(\widehat{0}, 1)}, \frac{g^{\infty, 2}(\widehat{0}, 1)}{g^{\infty, 2}(\widehat{0},-1)}\right)
$$

Recall that

$$
g^{\infty, 2}(\widehat{0}, \xi)=\left\{\begin{array}{cll}
g^{\infty, 2}(\widehat{0},-1)|\xi|^{2} & \text { if } & \xi \leq 0 \\
g^{\infty, 2}(\widehat{0}, 1)|\xi|^{2} & \text { if } & \xi>0
\end{array} .\right.
$$

Moreover, from the assumption on the function $g$, clearly,

$$
g^{\infty, 2}(\widehat{0}, 1)>0 \text { and } g^{\infty, 2}(\widehat{0},-1)>0
$$

We make precise the probability measure $\bar{\mu}_{\hat{x}}$ localized on $\Sigma^{\prime} \subset \subset \Sigma$ as follows:

$$
\bar{\mu}_{\hat{x}}:=p(\hat{x}) \delta_{1}+q(\hat{x}) \delta_{-1} \quad \text { with } \quad p(\hat{x})+q(\hat{x})=1 \text { a.e. } \hat{x} \in \Sigma^{\prime} .
$$

Corollary 5.6. Under the assumptions of Theorem 5.5, the three following estimates hold:
(i) for a.e. $\hat{x}$ in $\Sigma^{\prime}$

$$
\begin{equation*}
c(g) \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\hat{x}, s)\right|^{2} d s \leq \frac{d \bar{\pi}}{d x_{N}}(\hat{x}) \leq C(g) \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\hat{x}, s)\right|^{2} d s \tag{5.9}
\end{equation*}
$$

(ii) $\frac{c(g)|[\bar{u}](\hat{x})|^{2}}{\int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\hat{x}, s)\right|^{2} d s} \leq p(\hat{x}) \leq 1$ for a.e. $\hat{x}$ such that $[\bar{u}](\hat{x})>0$;

$$
\text { and } \frac{d \bar{\pi}}{d x_{N}}(\hat{x})=\int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\hat{x}, s)\right|^{2} d s \quad \text { when } g^{\infty, 2}(\widehat{0},-1)=g^{\infty, 2}(\widehat{0}, 1)
$$

(iii) $\frac{c(g)|[\bar{u}](\hat{x})|^{2}}{\int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\hat{x}, s)\right|^{2} d s} \leq q(\hat{x}) \leq 1$ for a.e. $\hat{x}$ such that $[\bar{u}](\hat{x})<0$.

Proof. Since $\bar{\mu}_{\hat{x}}=p(\widehat{x}) \delta_{1}+q(\widehat{x}) \delta_{-1}$, we have $\int_{\mathbb{S}_{0}} g^{\infty, 2}(\xi) d \bar{\mu}_{\hat{x}}=p(\widehat{x}) g^{\infty, 2}(\widehat{0}, 1)+q(\widehat{x}) g^{\infty, 2}(\widehat{0},-1)$ with $p(\widehat{x})+q(\widehat{x})=1$ a.e. $\widehat{x}$ in $\Sigma^{\prime}$ so that from (5.3), one has

$$
\begin{align*}
\int_{-\frac{1}{2}}^{\frac{1}{2}} g^{\infty, 2}\left(\widehat{0}, \frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right) d s & =\left(\int_{\mathbb{S}^{0}} g^{\infty, 2}(\xi) d \bar{\mu}_{\hat{x}}\right) \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \\
& =\frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x})\left\{p(\widehat{x}) g^{\infty, 2}(\widehat{0}, 1)+q(\widehat{x}) g^{\infty, 2}(\widehat{0},-1)\right\} \text { a.e. } \widehat{x} \in \Sigma^{\prime} \tag{5.10}
\end{align*}
$$

We are going to establish

$$
c(g) \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s \leq \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \leq C(g) \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s
$$

From (5.10) we deduce that

$$
\begin{align*}
\min \left\{g^{\infty, 2}(\widehat{0},-1), g^{\infty, 2}(\widehat{0}, 1)\right\} \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s & \leq \int_{-\frac{1}{2}}^{\frac{1}{2}} g^{\infty, 2}\left(\widehat{0}, \frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right) d s \\
& =\left(\int_{\mathbb{S}^{0}} g^{\infty, 2}(\xi) d \bar{\mu}_{x}\right) \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \\
& =\left\{p(\widehat{x}) g^{\infty, 2}(\widehat{0}, 1)+q(\widehat{x}) g^{\infty, 2}(\widehat{0},-1)\right\} \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \\
& \leq \max \left\{g^{\infty, 2}(\widehat{0},-1), g^{\infty, 2}(\widehat{0}, 1)\right\} \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \tag{5.11}
\end{align*}
$$

and

$$
\begin{aligned}
\min \left\{g^{\infty, 2}(\widehat{0},-1), g^{\infty, 2}(\widehat{0}, 1)\right\} \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) & =\min \left\{g^{\infty, 2}(\widehat{0},-1), g^{\infty, 2}(\widehat{0}, 1)\right\}\{p(\widehat{x})+q(\widehat{x})\} \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \\
& \leq\left\{p(\widehat{x}) g^{\infty, 2}(\widehat{0}, 1)+q(\widehat{x}) g^{\infty, 2}(\widehat{0},-1)\right\} \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \\
& =\left(\int_{\mathbb{S}_{0}} g^{\infty, 2}(\xi) d \bar{\mu}_{x}\right) \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \\
& =\int_{-\frac{1}{2}}^{\frac{1}{2}} g^{\infty, 2}\left(\widehat{0}, \frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right) d s \\
& \leq \max \left\{g^{\infty, 2}(\widehat{0},-1), g^{\infty, 2}(\widehat{0}, 1)\right\} \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s(5.12)
\end{aligned}
$$

Then, from (5.11) and (5.12) we have

$$
c(g) \int_{\frac{-1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s=\frac{\min \left\{g^{\infty, 2}(\widehat{0},-1), g^{\infty, 2}(\widehat{0}, 1)\right\}}{\max \left\{g^{\infty, 2}(\widehat{0},-1), g^{\infty, 2}(\widehat{0}, 1)\right\}} \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s \leq \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x})
$$

and

$$
\frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \leq \frac{\max \left\{g^{\infty, 2}(\widehat{0},-1), g^{\infty, 2}(\widehat{0}, 1)\right\}}{\min \left\{g^{\infty, 2}(\widehat{0},-1), g^{\infty, 2}(\widehat{0}, 1)\right\}} \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s=C(g) \int_{\frac{-1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s
$$

from which we deduce,

$$
c(g) \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s \leq \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x}) \leq C(g) \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s
$$

Let us prove (ii) and (iii). According to Theorem 5.2, for every $\varphi \in \mathcal{C}\left(\mathbb{S}^{0}\right)$ such that $\varphi^{* *}>-\infty$,

$$
\begin{equation*}
\frac{d \bar{\pi}}{d \mathcal{H}_{\left[\Sigma^{\prime}\right.}^{N-1}}(x) \int_{\mathbb{S}^{0}} \varphi(\zeta) d \bar{\mu}_{x} \geq \varphi^{* *}([v](x)) \quad \text { for } \mathcal{H}^{N-1} \text { a.e. } x \in \Sigma^{\prime} \tag{5.13}
\end{equation*}
$$

where $\bar{\pi}=\frac{d \bar{\pi}}{d \mathcal{H}_{\left[\Sigma^{\prime}\right.}^{N-1}} \mathcal{H}_{\left[\Sigma^{\prime}\right.}^{N-1}+\bar{\pi}_{s}$ is the Radon-Nikodym decomposition of $\bar{\pi}$ with respect to the measure $\mathcal{H}_{\left[\Sigma^{\prime}\right.}^{N-1}$. We assume that $[\bar{u}](\hat{x})>0$ and show that

$$
\frac{c(g)|[\bar{u}](\hat{x})|^{2}}{\int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\hat{x}, s)\right|^{2} d s} \leq p(\hat{x}) \leq 1
$$

Let

$$
\varphi(\xi)=\left\{\begin{array}{ccc}
\varphi(1)|\xi|^{2} & \text { if } & \xi \geq 0  \tag{5.14}\\
0 & \text { if } & \xi<0
\end{array}\right.
$$

Clearly, $\varphi^{* *}([\bar{u}](\widehat{x}))=\varphi([\bar{u}](\widehat{x}))=\varphi(1)|[u](\widehat{x})|^{2}$. From the inequality (5.13), it follows that

$$
\begin{aligned}
\varphi(1)|[\bar{u}](\widehat{x})|^{2} & =\varphi^{* *}([\bar{u}](\widehat{x})) \\
& \leq \frac{d \bar{\pi}}{d \widehat{x}}(\widehat{x})\left(\int_{\mathbb{S}^{0}} \varphi(\xi) d \bar{\mu}_{x}\right) \\
& \leq C(g)\left(\int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s\right)\left(\int_{\mathbb{S}_{0}} \varphi(\xi) d \bar{\mu}_{x}\right) \\
& =C(g) p(\widehat{x}) \varphi(1) \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\widehat{x}, s)\right|^{2} d s
\end{aligned}
$$

Then, we obtain

$$
\frac{|[\bar{u}](\hat{x})|^{2}}{C(g) \int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\hat{x}, s)\right|^{2} d s}=\frac{c(g)|[\bar{u}](\hat{x})|^{2}}{\int_{-\frac{1}{2}}^{\frac{1}{2}}\left|\frac{\partial \bar{\theta}}{\partial x_{N}}(\hat{x}, s)\right|^{2} d s} \leq p(\hat{x}) \leq 1
$$

The proof of (iii) is similar.

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