# Brownian motion in the quadrant with oblique repulsion from the sides 

Dominique Lépingle

## To cite this version:

Dominique Lépingle. Brownian motion in the quadrant with oblique repulsion from the sides. 2013. hal-00725894v2

## HAL Id: hal-00725894 <br> https://hal.science/hal-00725894v2

Preprint submitted on 13 Feb 2013

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# BROWNIAN MOTION IN THE QUADRANT WITH OBLIQUE REPULSION FROM THE SIDES 

Dominique Lépingle ${ }^{1}$


#### Abstract

We consider the problem of strong existence and uniqueness of a Brownian motion forced to stay in the nonnegative quadrant by an electrostatic repulsion from the sides that works obliquely. When the direction of repulsion is normal, the question has previously been solved with the help of convex analysis. To construct the solution, we start from the normal case and then we use as main tool a comparison lemma. The results are reminiscent of the study of a Brownian motion with oblique reflection in a wedge. Actually, the same skew symmetry condition is involved when looking for a stationary distribution in product form. The terms of the product are now gamma distributions in place of exponential ones. An associated purely deterministic problem is also considered.


## 1. Introduction

In the late seventies the study of heavy traffic limits in open multi-station queueing networks has put the question of existence and properties of the Brownian motion obliquely reflected on the sides of a wedge and more generally on the faces of a polyhedron. In the following decade there was an extensive literary output on that topic, among which we mention the works of Harrison, Reiman, Williams and collaborators ([7],[8],[16],,[18],,[9],[4], to cite a few of them). But there is another way than normal or oblique reflection to prevent Brownian motion from overstepping a linear barrier. We may add as drift term the gradient of a concave function that explodes in the neighborhood of the faces of the polyhedron. To be more specific, let $\mathbf{n}_{\mathbf{1}}, \ldots, \mathbf{n}_{\mathbf{k}}$ be unit vectors in $\mathbb{R}^{d}$ and $b_{1}, \ldots, b_{k}$ be real numbers. The state space $S$ is defined by

$$
S:=\left\{\mathbf{x} \in \mathbb{R}^{d}: \mathbf{n}_{r} . \mathbf{x} \geq b_{r}, r=1, \ldots, k\right\}
$$

Let $\phi_{1}, \ldots, \phi_{k}$ be $k$ convex $C^{1}$ functions on $(0, \infty)$ with $\phi_{r}(0+)=+\infty$ for any $r=1, \ldots, k$. The potential function $\Phi$ on $S$ is defined by

$$
\Phi(\mathbf{x}):=\sum_{r=1}^{k} \phi_{r}\left(\mathbf{n}_{r} . \mathbf{x}-b_{r}\right) .
$$

From the general existence and uniqueness theorem on multivalued stochastic differential systems established in [1] and [2], completed with identification of the drift term as in Lemma 3.4 in [3], we know there exists a unique strong solution living in $S$ to the equation

$$
\begin{align*}
d \mathbf{X}_{t} & =d \mathbf{B}_{t}-\nabla \Phi\left(\mathbf{X}_{t}\right) d t \\
\mathbf{X}_{0} & \in S \tag{1}
\end{align*}
$$

where $\mathbf{B}$ is a Brownian motion in $\mathbb{R}^{d}$. As was proved in Proposition 4.1 of [13], the hypotheses $\phi_{r}(0+)=+\infty$ for any $r=1, \ldots, k$ entail that there is no additional boundary process of

[^0]local time type in the r.h.s. of (1) since the repulsion forces are sufficiently strong. As
$$
-\nabla \Phi(\mathbf{x})=-\sum_{r=1}^{k} \phi_{r}^{\prime}\left(\mathbf{n}_{r} . \mathbf{x}-b_{r}\right) \mathbf{n}_{r}
$$
the repulsion from the faces $F_{r}=\left\{\mathbf{x} \in S: \mathbf{n}_{r} \cdot \mathbf{x}=b_{r}\right\}$ points in a normal direction into the interior of $S$. We now introduce vectors $\mathbf{q}_{1}, \ldots, \mathbf{q}_{k}$ with $\mathbf{q}_{r} . \mathbf{n}_{r}=0$ for $r=1, \ldots, k$ and consider a new drift function
\[

$$
\begin{equation*}
-\sum_{r=1}^{k} \phi_{r}^{\prime}\left(\mathbf{n}_{r} . \mathbf{x}-b_{r}\right)\left(\mathbf{n}_{r}+\mathbf{q}_{r}\right) \tag{2}
\end{equation*}
$$

\]

This is a singular drift and convex analysis cannot be used as in [1] or [2] to get strong existence and uniqueness of the associated stochastic differential system.

In this paper we will concentrate on the nonnegative quadrant in $\mathbb{R}^{2}$ as state space and a drift term that derives from electrostatic repulsive forces. We obtain strong existence and uniqueness for a large set of parameters. Our results are reminiscent of the thorough study in [16] of the oblique reflection in a wedge. However a key tool in this work was an appropriate harmonic function that made a weak approach possible and fruitful. Thus full results were obtained, while our strong approach merely provides a partial answer to the crucial question of hitting the corner.

For theoretical as well as practical reasons, a great deal of interest was taken in the question of existence and computation of the invariant measure of the Brownian motion with a constant drift vector and oblique reflection ([8],[18],[9]). Under the assumption that the directions of reflection satisfy a skew symmetry condition, it was proved that the invariant measure has exponential product form density. Motivated by the so-called Atlas model of equity markets presented in [6], some authors ([15],[10],[12]) have recently studied Brownian motions on the line with rank dependent local characteristics. This model is strongly related to Brownian motion reflected in polyhedral domains. The invariant probability density has an explicit exponential product form when the volatility is constant [15] and a sum of products of exponentials form when the volatility coefficients depend on the rank ([12],[11]). This last kind of density was previously obtained in [5] for a Brownian motion in a wedge with oblique reflection.

A neighboring way has been recently explored in [14]. Here the process is a Brownian motion with a drift term that is continuous and depends obliquely, via a regular potential function, on the position of the process relative to a polyhedral domain. Under the same skew symmetry condition as in [9], the invariant density has an explicit product form again. In Section 5, we consider a Brownian motion with a constant drift living in the nonnegative quadrant with oblique electrostatic repulsion from the sides. Under the skew symmetry condition, there is still an invariant measure in product form. Now the terms of the product are two gamma distributions with explicit parameters.

In the last section we consider the same model of oblique electrostatic repulsion, this time without Brownian term. We are interested in the question of existence and uniqueness of this deterministic differential system when the starting point is the corner. There is a trivial solution and the proofs of uniqueness of the previous stochasting setting are still in force in this simpler case. Taking advantage of the explicit form of this solution, we can obtain weaker conditions for uniqueness.

## 2. The Setting

The general state space is the nonnegative quadrant $S=\mathbb{R}_{+} \times \mathbb{R}_{+}$. The corner $\mathbf{0}=(0,0)$ will play a crucial role and in some cases it will be necessary to restrict the state space to the punctured nonnegative quadrant $S^{\mathbf{0}}=S \backslash\{\mathbf{0}\}$.

Let $\left(B_{t}, C_{t}\right)$ be a Brownian motion in the plane starting from $\mathbf{0}$, adapted to a filtration $\mathcal{F}=\left(\mathcal{F}_{t}\right)$ with usual conditions. Let $\alpha, \beta, \gamma, \delta$ be four real constants with $\alpha>0, \delta>0$. We say that an $\mathcal{F}$-adapted continuous process $(X, Y)$ with values in $S$ is a Brownian motion with electrostatic oblique repulsion from the sides if for any $t \geq 0$

$$
\begin{align*}
& X_{t}=X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}}+\beta \int_{0}^{t} \frac{d s}{Y_{s}} \geq 0 \\
& Y_{t}=Y_{0}+C_{t}+\gamma \int_{0}^{t} \frac{d s}{X_{s}}+\delta \int_{0}^{t} \frac{d s}{Y_{s}} \geq 0 \tag{3}
\end{align*}
$$

where $X_{0}$ and $Y_{0}$ are non-negative $\mathcal{F}_{0}$-measurable random variables. Each coordinate $X_{t}$ or $Y_{t}$ may vanish, so to make sense we must have a.s. for any $t \geq 0$

$$
\begin{array}{ll}
\int_{0}^{t} \mathbf{1}_{\left\{X_{s}=0\right\}} d s=0 & \int_{0}^{t} \mathbf{1}_{\left\{Y_{s}=0\right\}} d s=0 \\
\int_{0}^{t} \mathbf{1}_{\left\{X_{s}>0\right\}} \frac{d s}{X_{s}}<\infty & \int_{0}^{t} \mathbf{1}_{\left\{Y_{s}>0\right\}} \frac{d s}{Y_{s}}<\infty .
\end{array}
$$

The drift in (3) is of type (2) with $d=k=2$ and

$$
\begin{array}{ll}
\phi_{1}(x)=-\alpha \log (x) & \phi_{2}(y)=-\delta \log (y) \\
b_{1}=0 & b_{2}=0 \\
\mathbf{n}_{1}=(1,0) & \mathbf{n}_{2}=(0,1) \\
\mathbf{q}_{1}=\left(0, \frac{\gamma}{\alpha}\right) & \mathbf{q}_{2}=\left(\frac{\beta}{\delta}, 0\right)
\end{array}
$$

The case with $\beta=\gamma=0$ is a particular case of (1). In the sequel, we will note $(U, V)$ the solution of the system

$$
\begin{align*}
U_{t} & =X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{U_{s}} \geq 0 \\
V_{t} & =Y_{0}+C_{t}+\delta \int_{0}^{t} \frac{d s}{V_{s}} \geq 0 \tag{4}
\end{align*}
$$

The processes $U$ and $V$ are independent Bessel processes (if $X_{0}$ and $Y_{0}$ are independent variables). Actually, $U$ is a Bessel process with index $\alpha-\frac{1}{2}$, and the point 0 is intanstaneously reflecting for $U$ if $\alpha<\frac{1}{2}$ and polar if $\alpha \geq \frac{1}{2}$. Moreover, $U^{2}+V^{2}$ is the square of a Bessel process with index $\alpha+\delta$, and so the corner $\mathbf{0}$ is polar for $(U, V)$ in any case.

Comparison between $X$ and $U, Y$ and $V$ will play a key role in the construction of the solution $(X, Y)$ and the study of its behavior near the sides of the quadrant. The following simple lemma will be of constant use.

Lemma 1. For $T>0, \alpha>0$, let $x_{1}$ and $x_{2}$ be nonnegative continuous solutions on $[0, T]$ of the equations

$$
\begin{aligned}
& x_{1}(t)=v_{1}(t)+\alpha \int_{0}^{t} \frac{d s}{x_{1}(s)} \\
& x_{2}(t)=v_{2}(t)+\alpha \int_{0}^{t} \frac{d s}{x_{2}(s)}
\end{aligned}
$$

where $v_{1}, v_{2}$ are continuous functions such that $0 \leq v_{1}(0) \leq v_{2}(0)$, and $v_{2}-v_{1}$ is nondecreasing. Then $x_{1}(t) \leq x_{2}(t)$ on $[0, T]$.

Proof. Assume there exists $t \in(0, T]$ such that $x_{2}(t)<x_{1}(t)$. Set

$$
\tau:=\max \left\{s \leq t: x_{1}(s) \leq x_{2}(s)\right\}
$$

Then,

$$
\begin{aligned}
x_{2}(t)-x_{1}(t) & =x_{2}(\tau)-x_{1}(\tau)+\left(v_{2}(t)-v_{1}(t)\right)-\left(v_{2}(\tau)-v_{1}(\tau)\right)+\alpha \int_{\tau}^{t}\left(\frac{1}{x_{2}(s)}-\frac{1}{x_{1}(s)}\right) d s \\
& \geq 0
\end{aligned}
$$

a contradiction.

We will also need the following consequence of the results in [1] or [2] on multivalued stochastic differential systems, completed with the method used in [13] to check the lack of additional boundary process.
Proposition 2. Let $\alpha>0, \delta \geq 0, \sigma=\left(\sigma_{j}^{i} ; i, j=1,2\right)$ a $2 \times 2$-matrix, $\left(B_{1}, B_{2}\right)$ a Brownian motion in the plane, $b_{1}$ and $b_{2}$ two Lipschitz functions on $\mathbb{R}^{2}, Z_{0}^{1}$ and $Z_{0}^{2}$ two $\mathcal{F}_{0}$-measurable nonnegative random variables. There exists a unique solution $\left(Z^{1}, Z^{2}\right)$ to the system

$$
\begin{align*}
& Z_{t}^{1}=Z_{0}^{1}+\sigma_{1}^{1} B_{t}^{1}+\sigma_{2}^{1} B_{t}^{2}+\alpha \int_{0}^{t} \frac{d s}{Z_{s}^{s}}+\int_{0}^{t} b_{1}\left(Z_{s}^{1}, Z_{s}^{2}\right) d s \\
& Z_{t}^{2}=Z_{0}^{2}+\sigma_{1}^{2} B_{t}^{1}+\sigma_{2}^{2} B_{t}^{2}+\delta \int_{0}^{t} \frac{d s}{Z_{s}^{2}}+\int_{0}^{t} b_{2}\left(Z_{s}^{1}, Z_{s}^{2}\right) d s \tag{5}
\end{align*}
$$

with the conditions $Z_{t}^{1} \geq 0$ if $\delta=0$ and $Z_{t}^{1} \geq 0, Z_{t}^{2} \geq 0$ if $\delta>0$.
It is worth noticing that the solutions to (3) enjoy the Brownian scaling property. It means that if $(X, Y)$ is a solution to (3) starting from $\left(X_{0}, Y_{0}\right)$ with driving Brownian motion $\left(B_{t}, C_{t}\right)$, then for any $c>0$ the process ( $X_{t}^{\prime}:=c^{-1} X_{c^{2} t}, Y_{t}^{\prime}:=c^{-1} Y_{c^{2} t} ; t \geq 0$ ) is a solution to (3) starting from $\left(c^{-1} X_{0}, c^{-1} Y_{0}\right)$ with driving Brownian motion $\left(c^{-1} B_{c^{2}} t, c^{-1} C_{c^{2} t}\right)$.

## 3. Avoiding the corner

We shall see in the next section that existence and uniqueness of the solution to (3) are easily obtained as soon as the solution process keeps away from the corner. Thus the question of attaining the corner in finite time is of great interest.

Theorem 3. Let $(X, Y)$ be a solution to (3) in the interval $[0, \tau] \cap[0, \infty)$ where $\tau$ is a $\mathcal{F}$-stopping time. We set

$$
\tau^{0}:=\inf \left\{t \in(0, \tau] \cap(0, \infty):\left(X_{t}, Y_{t}\right)=\mathbf{0}\right\}
$$

with the usual convention $\inf \emptyset=\infty$. Then $\mathbb{P}\left(\tau^{\mathbf{0}}<\infty\right)=0$ if one of the following conditions is satisfied:
(1) $C_{1}: \beta \geq 0$ and $\gamma \geq 0$
(2) $C_{2 a}: \alpha \geq \frac{1}{2}$ and $\beta \geq 0$
(3) $C_{2 b}: \delta \geq \frac{1}{2}$ and $\gamma \geq 0$
(4) $C_{3}$ : There exist $\lambda>0$ and $\mu>0$ such that

- $\lambda \alpha+\mu \gamma \geq 0$
- $\lambda \beta+\mu \delta \geq 0$
- $\lambda(\lambda \alpha+\mu \gamma)+\mu(\lambda \beta+\mu \delta)-\frac{1}{2}\left(\lambda^{2}+\mu^{2}\right) \geq-2 \sqrt{\lambda \mu(\lambda \beta+\mu \delta)(\lambda \alpha+\mu \gamma)}$.

Proof. Assume $\tau=\infty$ a.s. for simplicity of notation.
Condition $C_{1}$. From Lemma 1 we get $X_{t} \geq U_{t}, Y_{t} \geq V_{t}$, where ( $U, V$ ) is the solution to (4), and we know that $\mathbf{0}$ is polar for $(U, V)$.
Condition $C_{2 a}$ (resp. $C_{2 b}$ ). From Lemma 1 we get $X_{t} \geq U_{t}$ (resp. $Y_{t} \geq V_{t}$ ) and in this case 0 is polar for $U$ (resp. V), so $U_{t}>0\left(\right.$ resp. $\left.V_{t}>0\right)$ for $t>0$.
Condition $C_{3}$. For $\epsilon>0$ let

$$
\begin{aligned}
& \sigma^{\epsilon}=\mathbf{1}_{\left\{\left(X_{0}, Y_{0}\right)=\mathbf{0}\right\}} \inf \left\{t>0: X_{t}+Y_{t} \geq \epsilon\right\} \\
& \boldsymbol{\tau}^{\mathbf{0}, \epsilon}=\inf \left\{t>\sigma^{\epsilon}:\left(X_{t}, Y_{t}\right)=\mathbf{0}\right\} .
\end{aligned}
$$

As $\epsilon \downarrow 0, \sigma^{\epsilon} \downarrow 0$ and $\tau^{0, \epsilon} \downarrow \tau^{0}$. We set $S_{t}=\lambda X_{t}+\mu Y_{t}$ for $t \geq 0, \lambda>0$ and $\mu>0$. From Ito formula we get for $t \in\left[\sigma^{\epsilon}, \tau^{\mathbf{0}, \epsilon}\right)$
$\log S_{t}$
$=\log S_{\sigma^{\epsilon}}+\int_{\sigma^{\epsilon}}^{t} \frac{\lambda d B_{s}+\mu d C_{s}}{S_{s}}+(\lambda \alpha+\mu \gamma) \int_{\sigma^{\epsilon}}^{t} \frac{d s}{X_{s} S_{s}}+(\lambda \beta+\mu \delta) \int_{\sigma^{\epsilon}}^{t} \frac{d s}{Y_{s} S_{s}}-\frac{1}{2}\left(\lambda^{2}+\mu^{2}\right) \int_{\sigma^{\epsilon}}^{t} \frac{d s}{S_{s}^{2}}$
$=\log S_{\sigma^{\epsilon}}+M_{t}+\int_{\sigma^{\epsilon}}^{t} \frac{P\left(X_{s}, Y_{s}\right)}{X_{s} Y_{s} S_{s}^{2}} d s$
where $M$ is a continuous local martingale and $P(x, y)$ is the second degree homogeneous polynomial

$$
P(x, y)=\lambda(\lambda \beta+\mu \delta) x^{2}+\mu(\lambda \alpha+\mu \gamma) y^{2}+\left(\lambda(\lambda \alpha+\mu \gamma)+\mu(\lambda \beta+\mu \delta)-\frac{1}{2}\left(\lambda^{2}+\mu^{2}\right)\right) x y
$$

Condition $C_{3}$ is exactly the condition for $P$ being nonnegative in $S$. Therefore

$$
0 \leq \int_{\sigma^{\epsilon}}^{t} \frac{P\left(X_{s}, Y_{s}\right)}{X_{s} Y_{s} S_{s}^{2}} d s<\infty
$$

and so

$$
0 \leq \int_{\sigma^{\epsilon}}^{\tau^{\mathbf{0}, \epsilon}} \frac{P\left(X_{s}, Y_{s}\right)}{X_{s} Y_{s} S_{s}^{2}} d s \leq \infty
$$

As $t \rightarrow \tau^{\mathbf{0}, \epsilon}$, the continuous local martingale $M$ either converges to a finite limit or oscillates between $+\infty$ and $-\infty$. It cannot converge to $-\infty$ and thus $S_{\tau^{\mathbf{0}, \epsilon}}>0$ on $\left\{\tau^{\mathbf{0}, \epsilon}<\infty\right\}$, proving that $\mathbb{P}\left(\tau^{\mathbf{0}, \epsilon}<\infty\right)=0$ and finally $\mathbb{P}\left(\tau^{\mathbf{0}}<\infty\right)=0$.

Example. When $\alpha=\delta$ and $|\beta|=|\gamma|$, condition $C_{3}$ is satisfied (with $\lambda=\mu$ ) if

$$
\begin{align*}
& \text { - } \quad \beta^{2} \leq \alpha-\frac{1}{4} \quad \text { when } \quad \beta=-\gamma \\
& \text { - } \quad-\beta \leq \alpha-\frac{1}{4} \quad \text { when } \quad \beta=\gamma<0 \text {. } \tag{6}
\end{align*}
$$

We may also be interested in hitting a single side. Then we set

$$
\begin{align*}
\tau_{X}^{0} & :=\inf \left\{t>0: X_{t}=0\right\}  \tag{7}\\
\tau_{Y}^{0} & :=\inf \left\{t>0: Y_{t}=0\right\}
\end{align*}
$$

We already know that $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=0$ if $\alpha \geq \frac{1}{2}$ and $\beta \geq 0$. Conversely we can prove that $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=1$ if $\alpha<\frac{1}{2}$ and $\beta \leq 0$. If we know that the corner is not hit and $\alpha \geq \frac{1}{2}$, we can get rid of the nonnegativity assumption on $\beta$.
Proposition 4. Assume $\mathbb{P}\left(\tau^{0}<\infty\right)=0$. If $\alpha \geq \frac{1}{2}$, then $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=0$.
Proof. For $\eta>0$ let

$$
\begin{aligned}
\theta_{X}^{\eta} & =\mathbf{1}_{\left\{X_{0}=0\right\}} \inf \left\{t>0: X_{t} \geq \eta\right\} \\
\tau_{X}^{0, \eta} & =\inf \left\{t>\theta_{X}^{\eta}: X_{t}=0\right\}
\end{aligned}
$$

As $\eta \downarrow 0, \theta_{X}^{\eta} \downarrow 0$ and $\tau_{X}^{0, \eta} \downarrow \tau_{X}^{0}$. For $t \in\left[\theta_{X}^{\eta}, \tau_{X}^{0, \eta}\right)$,

$$
\begin{equation*}
\log X_{t}=\log X_{\theta_{X}^{\eta}}+\int_{\theta_{X}^{\eta}}^{t} \frac{d B_{s}}{X_{s}}+\left(\alpha-\frac{1}{2}\right) \int_{\theta_{X}^{\eta}}^{t} \frac{d s}{X_{s}^{2}}+\beta \int_{\theta_{X}^{\eta}}^{t} \frac{d s}{X_{s} Y_{s}} \tag{8}
\end{equation*}
$$

Since $\mathbb{P}\left(\tau^{\mathbf{0}}<\infty\right)=0$, we have $Y_{\tau_{X}^{0, \eta}}>0$ on $\left\{\tau_{X}^{0, \eta}<\infty\right\}$. On this set,

$$
\int_{\theta_{X}^{\eta}}^{\tau_{X}^{0, \eta}} \frac{d s}{X_{s}}<\infty, \quad \int_{\theta_{X}^{\eta}}^{\tau_{X}^{0, \eta}} \frac{d s}{Y_{s}}<\infty
$$

and $X_{s}>0$ on $\left[\theta_{X}^{\eta}, \tau_{X}^{0, \eta}\right)$, which proves that

$$
\beta \int_{\theta_{X}^{\eta}}^{\tau_{X}^{0, \eta}} \frac{d s}{X_{s} Y_{s}}>-\infty
$$

As $t \rightarrow \tau_{X}^{0, \eta}$, the local martingale in the r.h.s. of (8) cannot converge to $-\infty$. This entails that $\mathbb{P}\left(\tau_{X}^{0, \eta}<\infty\right)=0$ and therefore $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=0$.

We may use again the method in Theorem 3 and Proposition 4 to learn more about hitting the sides of the quadrant.

Proposition 5. Assume $\alpha \geq \frac{1}{2}$ and $\delta \geq \frac{1}{2}$. Then $\mathbb{P}\left(\tau_{X}^{0}<\infty\right)=\mathbb{P}\left(\tau_{Y}^{0}<\infty\right)=0$ if one of the following conditions is satisfied:
(1) $\beta>0$
(2) $\gamma>0$
(3) $\beta \gamma \leq\left(\alpha-\frac{1}{2}\right)\left(\delta-\frac{1}{2}\right)$.

Proof. For $\epsilon>0$ let

$$
\begin{aligned}
\rho^{\epsilon} & =\mathbf{1}_{\left\{X_{0} Y_{0}=0\right\}} \inf \left\{t>0: X_{t} Y_{t} \geq \epsilon\right\} \\
\rho^{0, \epsilon} & =\inf \left\{t>\rho^{\epsilon}: X_{t} Y_{t}=0\right\}
\end{aligned}
$$

For $\lambda>0$ and $\mu>0$ we set

$$
R_{t}=\lambda \log X_{t}+\mu \log Y_{t}
$$

From Ito formula we get for $t \in\left[\rho^{\epsilon}, \rho^{0, \epsilon}\right)$

$$
\begin{aligned}
R_{t} & =R_{\rho^{\epsilon}}+\int_{\rho^{\epsilon}}^{t}\left(\frac{\lambda}{X_{s}} d B_{s}+\frac{\mu}{Y_{s}} d C_{s}\right)+\int_{\rho^{\epsilon}}^{t}\left[\frac{\lambda\left(\alpha-\frac{1}{2}\right)}{X_{s}^{2}}+\frac{\mu\left(\delta-\frac{1}{2}\right)}{Y_{s}^{2}}+\frac{(\lambda \beta+\mu \gamma)}{X_{s} Y_{s}}\right] d s \\
& =R_{\rho^{\epsilon}}+N_{t}+\int_{\rho^{\epsilon}}^{t} \frac{Q\left(X_{s}, Y_{s}\right)}{X_{s}^{2} Y_{s}^{2}} d s
\end{aligned}
$$

where $N$ is a continuous local martingale and $Q(x, y)$ is the second degreee homogeneous polynomial

$$
Q(x, y)=\mu\left(\delta-\frac{1}{2}\right) x^{2}+\lambda\left(\alpha-\frac{1}{2}\right) y^{2}+(\lambda \beta+\mu \gamma) x y
$$

If we obtain that $Q$ is nonnegative on $S$, then the proof will terminate as in Theorem 3 under condition $C_{3}$. If $\beta>0$ or $\gamma>0$, we easily find $\lambda>0$ and $\mu>0$ such that $\lambda \beta+\mu \gamma \geq 0$, and then $Q$ is nonnegative on $\mathbb{R}^{2}$. If now $\beta \leq 0, \gamma \leq 0$ and $\beta \gamma \leq\left(\alpha-\frac{1}{2}\right)\left(\delta-\frac{1}{2}\right)$, we may take $\lambda=-\gamma, \mu=-\beta$ and we check that $Q(x, y)$ remains nonnegative on $\mathbb{R}^{2}$ as well.

## 4. Existence and uniqueness

We now proceed to the question of existence and uniqueness of a global solution to (3). We consider separately the three cases: $\beta \geq 0$ and $\gamma \geq 0$, then $\beta>0$ and $\gamma<0$, then $\beta \leq 0$ and $\gamma<0$.
4.1. Case $\beta \geq 0$ and $\gamma \geq 0$. This is exactly condition $C_{1}$.

Theorem 6. Assume $\beta \geq 0$ and $\gamma \geq 0$.
(1) There is a unique solution to (3) in $S^{\mathbf{0}}$.
(2) There is a solution to (3) in $S$ starting from $\mathbf{0}$.
(3) If $\alpha \delta \geq \beta \gamma$, there is a unique solution to (3) in $S$.

Proof. 1. Let $a>0, \epsilon>0$ and define for $(x, z) \in \mathbb{R}_{+} \times \mathbb{R}$

$$
\psi_{\epsilon}(x, z):=\frac{1}{\max (\gamma x+z, \alpha \epsilon)}
$$

This is a Lipschitz function. From Proposition 2 we know that the system

$$
\begin{align*}
X_{t}^{\epsilon} & =X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}^{\epsilon}}+\alpha \beta \int_{0}^{t} \psi_{\epsilon}\left(X_{s}^{\epsilon}, Z_{s}^{\epsilon}\right) d s \geq 0  \tag{9}\\
Z_{t}^{\epsilon} & =-\gamma X_{0}+\alpha\left(Y_{0}+\mathbf{1}_{\left\{Y_{0}=0\right\}} a\right)-\gamma B_{t}+\alpha C_{t}+\alpha(\alpha \delta-\beta \gamma) \int_{0}^{t} \psi_{\epsilon}\left(X_{s}^{\epsilon}, Z_{s}^{\epsilon}\right) d s
\end{align*}
$$

has a unique solution. Let

$$
\tau_{Y}^{\epsilon}:=\inf \left\{t>0: \gamma X_{t}^{\epsilon}+Z_{t}^{\epsilon}<\alpha \epsilon\right\}
$$

If $0<\eta<\epsilon<a$ we deduce from the uniqueness that $\left(X^{\epsilon}, Z^{\epsilon}\right)$ and $\left(X^{\eta}, Z^{\eta}\right)$ are identical on $\left[0, \tau_{Y}^{\epsilon}\right]$. Patching together we can set

$$
\begin{aligned}
X_{t} & :=\lim _{\epsilon \rightarrow 0} X_{t}^{\epsilon} \\
Y_{t} & :=\lim _{\epsilon \rightarrow 0} \frac{1}{\alpha}\left(\gamma X_{t}^{\epsilon}+Z_{t}^{\epsilon}\right)
\end{aligned}
$$

on $\left\{Y_{0}>0\right\} \times\left[0, \tau_{Y}^{0}\right)$, where

$$
\tau_{Y}^{0}:=\lim _{\epsilon \rightarrow 0} \tau_{Y}^{\epsilon}
$$

On this set, $(X, Y)$ is the unique solution to (3). As we noted in the proof of Theorem 3 with condition $C_{1}$, we have $X_{t} \geq U_{t}$ and $Y_{t} \geq V_{t}$. Therefore, on $\left\{Y_{0}>0\right\} \cap\left\{\tau_{Y}^{0}<\infty\right\}$,

$$
\int_{0}^{\tau_{Y}^{0}} \frac{d s}{X_{s}} \leq \int_{0}^{\tau_{Y}^{0}} \frac{d s}{U_{s}}<\infty \quad \text { and } \quad \int_{0}^{\tau_{Y}^{0}} \frac{d s}{Y_{s}} \leq \int_{0}^{\tau_{Y}^{0}} \frac{d s}{V_{s}}<\infty
$$

and we can define

$$
\begin{align*}
& X_{\tau_{Y}^{0}}:=\lim _{t \rightarrow \tau_{Y}^{0}} X_{t}=X_{0}+B_{\tau_{Y}^{0}}+\alpha \int_{0}^{\tau_{Y}^{0}} \frac{d s}{X_{s}}+\beta \int_{0}^{\tau_{Y}^{0}} \frac{d s}{Y_{s}} \\
& Y_{\tau_{Y}^{0}}:=\lim _{t \rightarrow \tau_{Y}^{0}} Y_{t}=Y_{0}+C_{\tau_{Y}^{0}}+\gamma \int_{0}^{\tau_{Y}^{0}} \frac{d s}{X_{s}}+\delta \int_{0}^{\tau_{Y}^{0}} \frac{d s}{Y_{s}} . \tag{10}
\end{align*}
$$

We have $Y_{\tau_{Y}^{0}}=0$ and as $\mathbf{0}$ is polar for $(U, V)$, then $X_{\tau_{Y}^{0}}>0$. In exactly the same way we can construct a solution on $\left\{Y_{0}>0\right\}$ in the interval $\left[T_{1}, T_{2}\right]$, where $T_{1}=\tau_{Y}^{0}, T_{2}=\inf \left\{t>T_{1}\right.$ : $\left.X_{t}=0\right\}$. Iterating, we get a solution on $\left\{Y_{0}>0\right\} \times\left[0, \lim _{n \rightarrow \infty} T_{n}\right.$ ) where

$$
\begin{aligned}
& T_{2 p}:=\inf \left\{t>T_{2 p-1}: X_{t}=0\right\} \\
& T_{2 p+1}:=\inf \left\{t>T_{2 p}: Y_{t}=0\right\} .
\end{aligned}
$$

On $\left\{Y_{0}>0\right\} \cap\left\{\lim _{n \rightarrow \infty} T_{n}<\infty\right\}$ we set $X_{\lim _{n \rightarrow \infty} T_{n}}:=\lim _{p \rightarrow \infty} X_{T_{2 p}}=0$ and $Y_{\lim _{n \rightarrow \infty} T_{n}}:=$ $\lim _{p \rightarrow \infty} Y_{T_{2 p+1}}=0$. The polarity of $\mathbf{0}$ entails this is not possible in finite time and thus $\lim _{n \rightarrow \infty} T_{n}=\infty$. So we have obtained a unique global solution on $\left\{Y_{0}>0\right\}$. In the same way we obtain a unique global solution on $\left\{X_{0}>0\right\}$ and as $\mathbb{P}\left(\left(X_{0}, Y_{0}\right)=\mathbf{0}\right)=0$ the proof is complete.
2. Assume now $X_{0}=Y_{0}=0$. Let $\left(y_{n}\right)_{n \geq 1}$ be a sequence of real numbers (strictly) decreasing to 0 . From the above paragraph it follows there exists for any $n \geq 1$ a unique solution ( $X^{n}, Y^{n}$ ) with values in $S^{0}$ to the system

$$
\begin{aligned}
X_{t}^{n} & =B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}^{n}}+\beta \int_{0}^{t} \frac{d s}{Y_{s}^{n}} \\
Y_{t}^{n} & =y_{n}+C_{t}+\gamma \int_{0}^{t} \frac{d s}{X_{s}^{n}}+\delta \int_{0}^{t} \frac{d s}{Y_{s}^{n}} .
\end{aligned}
$$

Let

$$
\tau:=\inf \left\{t>0: X_{t}^{n+1}<X_{t}^{n}\right\} .
$$

Using Lemma 1 we obtain $Y_{t}^{n+1} \leq Y_{t}^{n}$ on $[0, \tau]$. We note that $\left(X_{\tau}^{n}, Y_{\tau}^{n}\right) \in S^{0}$ on $\{\tau<\infty\}$. On $\left\{Y_{\tau}^{n+1}=Y_{\tau}^{n}\right\} \cap\{\tau<\infty\}$, since $X_{\tau}^{n+1}=X_{\tau}^{n}$ and the solution starting at time $\tau$ is unique, it follows that $X_{t}^{n+1}=X_{t}^{n}$ and $Y_{t}^{n+1}=Y_{t}^{n}$ on $[\tau, \infty)$. On $\left\{Y_{\tau}^{n+1}<Y_{\tau}^{n}\right\} \cap\{\tau<\infty\}$, the continuity of solutions at time $\tau$ entails there exists $\rho>0$ such that $Y_{t}^{n+1} \leq Y_{t}^{n}$ on $[\tau, \tau+\rho]$. A second application of Lemma 1 proves that $X_{t}^{n+1} \geq X_{t}^{n}$ on $[\tau, \tau+\rho]$, a contradiction to the definition of $\tau$. Therefore $\mathbb{P}(\tau=\infty)=1$. It follows that $X_{t}^{n+1} \geq X_{t}^{n}$ and $Y_{t}^{n+1} \leq Y_{t}^{n}$ for any $t \in[0, \infty)$, and we may define

$$
X_{t}:=\lim _{n \rightarrow \infty} \uparrow X_{t}^{n} \quad Y_{t}:=\lim _{n \rightarrow \infty} \downarrow Y_{t}^{n}
$$

As $Y_{t}^{n} \geq V_{t}$ where $(U, V)$ is the solution to (4) with $X_{0}=Y_{0}=0$, we have

$$
\begin{aligned}
X_{t} & =B_{t}+\alpha \lim _{n \rightarrow \infty} \int_{0}^{t} \frac{d s}{X_{n}^{n}}+\beta \lim _{n \rightarrow \infty} \int_{0}^{t} \frac{d s}{Y_{s}^{n}} \\
& =B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}}+\beta \int_{0}^{t} \frac{d s}{Y_{s}} \\
& <\infty
\end{aligned}
$$

and also

$$
\begin{aligned}
Y_{t} & =\lim _{n \rightarrow \infty} y_{n}+C_{t}+\gamma \lim _{n \rightarrow \infty} \int_{0}^{t} \frac{d s}{X_{s}^{n}}+\delta \lim _{n \rightarrow \infty} \int_{0}^{t} \frac{d s}{Y_{s}^{n}} \\
& =C_{t}+\gamma \int_{0}^{t} \frac{d s}{X_{s}}+\delta \int_{0}^{t} \frac{d s}{Y_{s}} \\
& <\infty
\end{aligned}
$$

3. Assume finally $\alpha \delta-\beta \gamma \geq 0$. As the conclusion holds true if $\beta=\gamma=0$, we may also assume $\beta>0$. Let $(X, Y)$ be the solution to (3) with $X_{0}=Y_{0}=0$ obtained in the previous paragraph and let ( $X^{\prime}, Y^{\prime}$ ) be another solution. Considering ( $X^{n}, Y^{n}$ ) again and replacing ( $X^{n+1}, Y^{n+1}$ ) with ( $X^{\prime}, Y^{\prime}$ ), the previous proof works and we finally obtain $X_{t}^{\prime} \geq X_{t}$ and $Y_{t}^{\prime} \leq Y_{t}$. Then,

$$
\begin{align*}
& \left(\delta\left(X_{t}-X_{t}^{\prime}\right)-\beta\left(Y_{t}-Y_{t}^{\prime}\right)\right)^{2} \\
= & 2 \int_{0}^{t}\left(\delta\left(X_{s}-X_{s}^{\prime}\right)-\beta\left(Y_{s}-Y_{s}^{\prime}\right)\right)(\alpha \delta-\beta \gamma)\left(\frac{1}{X_{s}}-\frac{1}{X_{s}^{\prime}}\right) d s  \tag{11}\\
\leq & 0
\end{align*}
$$

and thus $X_{t}^{\prime}=X_{t}$ and $Y_{t}^{\prime}=Y_{t}$, proving uniqueness. Replacing $(\delta, \beta)$ with $(\gamma, \alpha)$ in equation (11) we obtain the same conclusion if $\gamma>0$.
4.2. Case $\beta>0$ and $\gamma<0$.

Theorem 7. Assume $\beta>0, \gamma<0$ and one of the conditions $C_{2 a}$ or $C_{3}$ is satisfied. Then, there exists a unique solution to (3) in $S^{0}$.

Proof. The proof is similar to the proof of 1 in Theorem 6. The only change is that now $Y_{t} \leq V_{t}$. Therefore, on $\left\{Y_{0}>0\right\} \cap\left\{\tau_{Y}^{0}<\infty\right\}$,

$$
\delta \int_{0}^{\tau_{Y}^{0}} \frac{d s}{Y_{s}} \leq V_{\tau_{Y}^{0}}-Y_{0}-C_{\tau_{Y}^{0}}-\gamma \int_{0}^{\tau_{Y}^{0}} \frac{d s}{U_{s}}<\infty
$$

and we can define $X_{\tau_{Y}^{0}}$ and $Y_{\tau_{Y}^{0}}$ as previously done.
4.3. Case $\beta \leq 0$ and $\gamma<0$. In this case we can give a full answer to the question of existence and uniqueness. Our condition of existence is exactly the condition found in [17] for the reflected Brownian in a wedge being a semimartingale, i.e. there is a convex combination of the directions of reflection that points into the wedge from the corner.

Theorem 8. Assume $\beta \leq 0$ and $\gamma<0$.
(1) If $\alpha \delta>\beta \gamma$, there exists a unique solution to (3) in $S$.
(2) If $\alpha \delta \leq \beta \gamma$, there does not exist any solution.

Proof. 1. Assume first $\alpha \delta>\beta \gamma$.
a) Existence. Let ( $h_{n}, n \geq 1$ ) be a (strictly) increasing sequence of bounded positive nonincreasing Lipschitz functions converging to $1 / x$ on $(0, \infty)$ and to $+\infty$ on $(-\infty, 0]$. For instance we can take

$$
\begin{array}{rlrl}
h_{n}(x) & =\left(1-\frac{1}{n}\right) \frac{1}{x} & & \text { on } \quad \\
& {\left[\frac{1}{n}, \infty\right)} \\
& =n-1 & & \text { on } \\
\left(-\infty, \frac{1}{n}\right] .
\end{array}
$$

We consider for each $n \geq 1$ the system

$$
\begin{align*}
X_{t}^{n} & =X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}^{n}}+\beta \int_{0}^{t} h_{n}\left(Y_{s}^{n}\right) d s \\
Y_{t}^{n} & =Y_{0}+C_{t}+\gamma \int_{0}^{t} h_{n}\left(X_{s}^{n}\right) d s+\delta \int_{0}^{t} \frac{d s}{Y_{s}^{n}} . \tag{12}
\end{align*}
$$

From Proposition 2 it follows there exists a unique solution to this system. We set

$$
\tau:=\inf \left\{s>0: X_{s}^{n+1}>X_{s}^{n}\right\} .
$$

We have $h_{n+1}\left(X_{t}^{n+1}\right) \geq h_{n}\left(X_{t}^{n}\right)$ on $[0, \tau]$. A first application of Lemma 1 shows that $Y_{t}^{n+1} \leq$ $Y_{t}^{n}$ on $[0, \tau]$. Since $h_{n+1}\left(Y_{\tau}^{n+1}\right)>h_{n}\left(Y_{\tau}^{n}\right)$ on $\{\tau<\infty\}$, we deduce from the continuity of solutions that there exists $\rho>0$ such that $h_{n+1}\left(Y_{t}^{n+1}\right) \geq h_{n}\left(Y_{t}^{n}\right)$ on $[\tau, \tau+\rho]$. A second application of Lemma 1 shows that $X_{t}^{n+1} \leq X_{t}^{n}$ on $[\tau, \tau+\rho]$, a contradiction to the definition of $\tau$. Thus $\mathbb{P}(\tau=\infty)=1$ proving that on the whole $[0, \infty)$ we have $X_{t}^{n+1} \leq X_{t}^{n}$ and $Y_{t}^{n+1} \leq Y_{t}^{n}$. Then we can set for any $t \in[0, \infty)$

$$
X_{t}:=\lim _{n \rightarrow \infty} X_{t}^{n} \quad \text { and } \quad Y_{t}:=\lim _{n \rightarrow \infty} Y_{t}^{n} .
$$

If $\alpha \delta>\beta \gamma$, there is a convex combination of the directions of repulsion that points into the positive quadrant, i.e. there exist $\lambda>0$ and $\mu>0$ such that $\lambda \alpha+\mu \gamma>0$ and $\mu \delta+\lambda \beta>0$. For $n \geq 1$ and $t \geq 0$,

$$
\begin{align*}
\lambda U_{t}+\mu V_{t} & \geq \lambda X_{t}^{n}+\mu Y_{t}^{n} \\
& \geq \lambda X_{0}+\mu Y_{0}+\lambda B_{t}+\mu C_{t}+(\lambda \alpha+\mu \gamma) \int_{0}^{t} \frac{d s}{X_{s}^{n}}+(\mu \delta+\lambda \beta) \int_{0}^{t} \frac{d s}{Y_{s}^{n}} \tag{13}
\end{align*}
$$

Letting $n \rightarrow \infty$ in (13) we obtain

$$
\int_{0}^{t} \frac{d s}{X_{s}}<\infty \quad \text { and } \quad \int_{0}^{t} \frac{d s}{Y_{s}}<\infty
$$

Then we may let $n$ go to $\infty$ in (12) proving that $(X, Y)$ is a solution to (3).
b) Uniqueness. Let $\left(X^{\prime}, Y^{\prime}\right)$ be another solution to (3). Replacing ( $X^{n+1}, Y^{n+1}$ ) with $\left(X^{\prime}, Y^{\prime}\right)$ we follow the above proof to obtain for $t \in[0, \infty)$ and $n \geq 1$

$$
X_{t}^{\prime} \leq X_{t}^{n} \quad \text { and } \quad Y_{t}^{\prime} \leq Y_{t}^{n}
$$

Letting $n \rightarrow \infty$ we conclude

$$
X_{t}^{\prime} \leq X_{t} \quad \text { and } \quad Y_{t}^{\prime} \leq Y_{t}
$$

With the same $\lambda>0$ and $\mu>0$ as above,

$$
\begin{aligned}
& \left(\lambda\left(X_{t}-X_{t}^{\prime}\right)+\mu\left(Y_{t}-Y_{t}^{\prime}\right)\right)^{2} \\
= & 2 \int_{0}^{t}\left(\lambda\left(X_{s}-X_{s}^{\prime}\right)+\mu\left(Y_{s}-Y_{s}^{\prime}\right)\right)\left[(\lambda \alpha+\mu \gamma)\left(\frac{1}{X_{s}}-\frac{1}{X_{s}^{\prime}}\right)+(\mu \delta+\lambda \beta)\left(\frac{1}{Y_{s}}-\frac{1}{Y_{s}^{\prime}}\right)\right] d s \\
\leq & 0
\end{aligned}
$$

and therefore $X_{t}^{\prime}=X_{t}, Y_{t}^{\prime}=Y_{t}$.
2. If $\alpha \delta \leq \beta \gamma$ there exist $\lambda>0$ and $\mu>0$ such that $\lambda \alpha+\mu \gamma \leq 0$ and $\mu \delta+\lambda \beta \leq 0$. Thus, if $(X, Y)$ is a solution to (3),

$$
0 \leq \lambda X_{t}+\mu Y_{t} \leq \lambda X_{0}+\mu Y_{0}+\lambda B_{t}+\mu C_{t}
$$

This is not possible since the paths of the Brownian motion $\left(\lambda^{2}+\mu^{2}\right)^{-1 / 2}\left(\lambda B_{t}+\mu C_{t}\right)$ are not bounded below. So there is no global solution.

In the following pictures, we display the $x$-repulsion direction vector $r_{x}=(\alpha, \gamma)$ and the $y$-repulsion direction vector $r_{y}=(\beta, \delta)$ in three illustrative instances.

$\beta>0, \gamma>0, \alpha \delta>\beta \gamma$

$\beta<0, \gamma>0, \alpha \beta+\gamma \delta=0$

$\beta<0, \gamma<0, \alpha \delta>\beta \gamma$

## 5. Product form stationary distribution

We introduce an additional constant drift $(-\mu,-\nu)$ in the nonnegative quadrant and consider the system

$$
\begin{align*}
& X_{t}=X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{X_{s}}+\beta \int_{0}^{t} \frac{d s}{Y_{s}}-\mu t  \tag{14}\\
& Y_{t}=Y_{0}+C_{t}+\gamma \int_{0}^{t} \frac{d s}{X_{s}}+\delta \int_{0}^{t} \frac{d s}{Y_{s}}-\nu t
\end{align*}
$$

with the conditions $X_{t} \geq 0, Y_{t} \geq 0$. If now

$$
\begin{align*}
U_{t} & =X_{0}+B_{t}+\alpha \int_{0}^{t} \frac{d s}{U_{s}}-\mu t \\
V_{t} & =Y_{0}+C_{t}+\delta \int_{0}^{t} \frac{d s}{V_{s}}-\nu t \tag{15}
\end{align*}
$$

with $U_{t} \geq 0, V_{t} \geq 0$, we can check that $\mathbf{0}$ is still polar for $(U, V)$ (as well, 0 is polar for $U$ if $\alpha \geq \frac{1}{2}$ and for $V$ if $\delta \geq \frac{1}{2}$ ). Therefore the results of the previous sections are still valid for the solution to (14). We are now looking for conditions on the set of parameters in order to obtain a stationary distribution for the Markov process $(X, Y)$ in the form of a product of two gamma distributions.
Theorem 9. Assume there exists a unique solution to (14) in $S^{\mathbf{0}}$ or in $S$. This process has an invariant distribution in the form $\Gamma(a, c) \otimes \Gamma(b, d)$ if and only if

- $\alpha \beta+\gamma \delta=0 \quad$ (skew symmetry)
- $a=2 \alpha+1, \quad b=2 \delta+1$
- $c=2 \delta \frac{\mu \alpha+\nu \gamma}{\alpha \delta-\beta \gamma}, \quad d=2 \alpha \frac{\mu \beta+\nu \delta}{\alpha \delta-\beta \gamma}$
- $\mu \alpha+\nu \gamma>0, \quad \mu \beta+\nu \delta>0$.

Proof. Let

$$
\rho(x, y)=x^{a-1} e^{-c x} y^{b-1} e^{-d y} \quad \text { for } x \geq 0, y \geq 0
$$

The infinitesimal generator of the diffusion (14) is given by

$$
L=\frac{1}{2}\left(\frac{\partial^{2}}{\partial x^{2}}+\frac{\partial^{2}}{\partial y^{2}}\right)+\left(\frac{\alpha}{x}+\frac{\beta}{y}-\mu\right) \frac{\partial}{\partial x}+\left(\frac{\gamma}{x}+\frac{\delta}{y}-\nu\right) \frac{\partial}{\partial y} .
$$

By a density argument, to prove that $\rho$ is an invariant density, it is enough to check that

$$
\int_{0}^{\infty} \int_{0}^{\infty} L f(x, y) \rho(x, y) d x d y=0
$$

for any $f(x, y)=g(x) h(y)$ with $g, h \in C_{c}^{2}((0, \infty))$ (compactly supported twice continuously differentiable functions on $(0, \infty))$. Integrating by parts, we get

$$
\int_{0}^{\infty} \int_{0}^{\infty} L(g h)(x, y) \rho(x, y) d x d y=\int_{0}^{\infty} \int_{0}^{\infty} g(x) h(y) J(x, y) d x d y
$$

where

$$
J(x, y)=\rho(x, y)\left[A+B x^{-1}+C x^{-2}+D y^{-1}+E y^{-2}+F x^{-1} y^{-1}\right]
$$

with

$$
\begin{aligned}
A & =\frac{1}{2} c^{2}+\frac{1}{2} d^{2}-\mu c-\nu d \\
B & =-(a-1) c+\mu(a-1)+\alpha c+\gamma d \\
C & =\frac{1}{2}(a-1)(a-2)-\alpha(a-2) \\
D & =-(b-1) d+\nu(b-1)+\beta c+\delta d \\
E & =\frac{1}{2}(b-1)(b-2)-\delta(b-2) \\
F & =\beta(a-1)+\gamma(b-1)
\end{aligned}
$$

Letting $A=B=C=D=E=F=0$ we obtain the specified values for $a, b, c, d$ and the skew symmetry condition $\alpha \beta+\gamma \delta=0$, which means that $r_{x}$ and $r_{y}$ are orthogonal. The last condition in the statement of the theorem is written out so that the invariant density $\rho$ is
integrable on $S$. It is satisfied if $(\mu, \nu)$ points into the interior of the quadrant designed by $r_{x}$ and $r_{y}$.

Remark. With the same proof, we may check that under the skew symmetry condition, when $\mu=\nu=0$, the function $\rho(x, y)=x^{2 \alpha} y^{2 \delta}$ is a non-integrable invariant density that does not depend on the obliqueness parameter $\beta$.

## 6. A SINGULAR DIFFERENTIAL SYSTEM

In this last section we give up the stochastic setting and consider the integral system

$$
\begin{align*}
x(t) & =\alpha \int_{0}^{t} \frac{d s}{x(s)}+\beta \int_{0}^{t} \frac{d s}{y(s)} \\
y(t) & =\gamma \int_{0}^{t} \frac{d s}{x(s)}+\delta \int_{0}^{t} \frac{d s}{y(s)} \tag{16}
\end{align*}
$$

where $x$ and $y$ are continuous functions from $[0, \infty)$ to $[0, \infty)$ with the conditions

$$
\begin{array}{ll}
\int_{0}^{t} \mathbf{1}_{\{x(s)=0\}} d s=0 & \int_{0}^{t} \mathbf{1}_{\{y(s)=0\}} d s=0 \\
\int_{0}^{t} \mathbf{1}_{\{x(s)>0\}} \frac{d s}{x(s)}<\infty & \int_{0}^{t} \mathbf{1}_{\{y(s)>0\}} \frac{d s}{y(s)}<\infty
\end{array}
$$

for any $t \geq 0$. Here $\alpha, \beta, \gamma, \delta$ are again four real constants with $\alpha>0, \delta>0$ and the starting point is the corner of the nonnegative quadrant.

An explicit solution is easily obtained.
Theorem 10. If $\max \{\beta, \gamma\} \geq 0$ or if $\beta \gamma<\alpha \delta$, there is a solution to (16) given by

$$
\begin{aligned}
& x(t)=c \sqrt{t} \\
& y(t)=d \sqrt{t}
\end{aligned}
$$

where

$$
\begin{aligned}
& c=\left(2 \alpha+\frac{\beta}{\delta}\left(\beta-\gamma+\sqrt{(\beta-\gamma)^{2}+4 \alpha \delta}\right)\right)^{1 / 2} \\
& d=\left(2 \delta+\frac{\gamma}{\alpha}\left(\gamma-\beta+\sqrt{(\beta-\gamma)^{2}+4 \alpha \delta}\right)\right)^{1 / 2} .
\end{aligned}
$$

There is no solution if $\beta<0, \gamma<0$ and $\alpha \delta \leq \beta \gamma$.
Proof. Writing down $x(t)=c \sqrt{t}$ and $y(t)=d \sqrt{t}$, we have to solve equations

$$
\begin{aligned}
& \frac{c}{2}=\frac{\alpha}{c}+\frac{\beta}{d} \\
& \frac{d}{2}=\frac{\gamma}{c}+\frac{\delta}{d}
\end{aligned}
$$

Simple computations lead to the requested values. We have to check that

$$
C=2 \alpha+\frac{\beta}{\delta}\left(\beta-\gamma+\sqrt{(\beta-\gamma)^{2}+4 \alpha \delta}\right)
$$

and

$$
D=2 \delta+\frac{\gamma}{\alpha}\left(\gamma-\beta+\sqrt{(\beta-\gamma)^{2}+4 \alpha \delta}\right)
$$

are positive. If $\beta \geq 0, C$ is clearly positive. This is also true if $\beta<0$ and $\beta \gamma<\alpha \delta$ since $C$ may be written

$$
C=\frac{4 \alpha(\alpha \delta-\beta \gamma)}{2 \alpha \delta-\beta \gamma+\beta^{2}-\beta \sqrt{4(\alpha \delta-\beta \gamma)+(\beta+\gamma)^{2}}} .
$$

The proof for $D$ is similar.
Assume now $\beta<0, \gamma<0, \alpha \delta \leq \beta \gamma$. As was noted in Section 4, there exist $\lambda>0$ and $\mu>0$ such that $\lambda \alpha+\mu \gamma \leq 0$ and $\mu \delta+\lambda \beta \leq 0$. If $(x, y)$ is a solution to (16), then

$$
0 \leq \lambda x(t)+\mu y(t) \leq 0
$$

and therefore $x(t)=y(t)=0$ for any $t \geq 0$, which is impossible.
We now take up the question of uniqueness.

Theorem 11. Let $\varepsilon=\min \left\{\beta\left(\beta-\gamma+\sqrt{(\beta-\gamma)^{2}+4 \alpha \delta}\right), \gamma\left(\gamma-\beta+\sqrt{(\beta-\gamma)^{2}+4 \alpha \delta}\right)\right\}$. The solution to (16) is unique

$$
\begin{array}{ll}
\text { a) when } & \beta \geq 0, \gamma \geq 0 \text {, if } \beta \gamma<2 \alpha \delta+\frac{\varepsilon}{2} \\
\text { b) when } & \beta \leq 0, \gamma \leq 0, \\
\text { if } \beta \gamma<\alpha \delta .
\end{array}
$$

Proof. We may use Proposition 2 with $\sigma_{j}^{i}=0$ for $i, j=1,2$ and consequently the proofs of Theorem 6 and Theorem 8 are still valid.
a) Case $\beta \geq 0, \gamma \geq 0$.

As in the proof of Theorem 6.3, there exist two extremal solutions $\left(x_{1}, y_{1}\right)$ and $\left(x_{2}, y_{2}\right)$ such that for any solution $(x, y)$ and any $t \geq 0$,

$$
\begin{aligned}
& x_{2}(t) \leq x(t) \leq x_{1}(t) \\
& y_{1}(t) \leq y(t) \leq y_{2}(t) .
\end{aligned}
$$

Uniqueness for equation (16) follows as well if $\beta \gamma \leq \alpha \delta$. We can improve the result by using the explicit form of the known solution and Gronwall lemma. Assume $\beta \gamma>\alpha \delta$. As in the proof of Theorem 6.2, we consider the unique solution $\left(x^{n}, y^{n}\right)$ to the system

$$
\begin{aligned}
& x^{n}(t)=\xi_{n}+\alpha \int_{0}^{t} \frac{d s}{x^{n}(s)}+\beta \int_{0}^{t} \frac{d s}{y^{n}(s)} \geq 0 \\
& y^{n}(t)=\gamma \int_{0}^{t} \frac{d s}{x^{n}(s)}+\delta \int_{0}^{t} \frac{d s}{y^{n}(s)} \quad \geq 0
\end{aligned}
$$

where $\left(\xi_{n}\right)$ is a sequence of decreasing numbers tending to 0 . Remarking that for any $t \geq 0$ we have $c \sqrt{t} \leq x^{n}(t)$ and $d \sqrt{t} \geq y^{n}(t)$, we set

$$
v_{n}(t)=\left(\delta\left(x^{n}(t)-c \sqrt{t}\right)-\beta\left(y^{n}(t)-d \sqrt{t}\right)\right)^{2} .
$$

We develop

$$
\begin{aligned}
v_{n}(t) & =\delta^{2} \xi_{n}^{2}+2 \int_{0}^{t} \sqrt{v_{n}(s)}(\beta \gamma-\alpha \delta)\left(\frac{1}{c \sqrt{s}}-\frac{1}{x^{n}(s)}\right) d s \\
& =\delta^{2} \xi_{n}^{2}+2 \frac{\beta \gamma-\alpha \delta}{c} \int_{0}^{t} \sqrt{v_{n}(s)} \frac{x^{n}(s)-c \sqrt{s}}{\sqrt{s} x^{n}(s)} d s \\
& \leq \delta^{2} \xi_{n}^{2}+2 \frac{\beta \gamma-\alpha \delta}{c \delta} \int_{0}^{t} \frac{v_{n}(s)}{\sqrt{5} x^{n}(s)} d s .
\end{aligned}
$$

From $x^{n}(t) \geq \xi_{n}$ and $x^{n}(t) \geq c \sqrt{t}$ it follows that for $t \geq \xi_{n}^{2} / c^{2}$

$$
v_{n}(t) \leq \delta^{2} \xi_{n}^{2}+2 \frac{\beta \gamma-\alpha \delta}{c \delta}\left(\int_{0}^{\xi_{n}^{2} / c^{2}} \frac{v_{n}(s)}{\xi_{n} \sqrt{s}} d s+\int_{\xi_{n}^{2} / c^{2}}^{t} \frac{v_{n}(s)}{c s} d s\right) .
$$

From Gronwall lemma,

$$
\begin{aligned}
v_{n}(t) & \leq \delta^{2} \xi_{n}^{2} \exp \left\{2 \frac{\beta \gamma-\alpha \delta}{c \delta}\left(\frac{2 \sqrt{\xi_{n}^{2} / c^{2}}}{\xi_{n}}+\frac{1}{c} \log t-\frac{1}{c} \log \frac{\xi_{n}^{2}}{c^{2}}\right)\right\} \\
& =C(t) \xi_{n}^{2\left(1-2 \frac{\beta \gamma-\alpha \delta}{c^{2} \delta}\right)}
\end{aligned}
$$

where $C(t)$ does not depend on $\xi_{n}$. Therefore $v_{n}(t) \rightarrow 0$ when $\xi_{n} \rightarrow 0$ if

$$
\beta \gamma<2 \alpha \delta+\frac{\beta}{2}\left(\beta-\gamma+\sqrt{(\beta-\gamma)^{2}+4 \alpha \delta}\right) .
$$

We conclude that

$$
x_{1}(t)=c \sqrt{t}, \quad y_{1}(t)=d \sqrt{t} .
$$

In a similar way,

$$
x_{2}(t)=c \sqrt{t}, \quad y_{2}(t)=d \sqrt{t}
$$

under the condition

$$
\beta \gamma<2 \alpha \delta+\frac{\gamma}{2}\left(\gamma-\beta+\sqrt{(\beta-\gamma)^{2}+4 \alpha \delta}\right) .
$$

Uniqueness follows if both conditions are satisfied.
b) Case $\beta \leq 0, \gamma<0$.

The proof is the same as for Theorem 8.1.

Remark. The question of uniqueness for all $\beta>0, \gamma>0$ remains open. Let us notice that for $\alpha=\delta=0$ uniqueness fails because the solutions are given by

$$
\begin{aligned}
& x(t)=C t^{\frac{\beta}{\beta+\gamma}} \\
& y(t)=\frac{\beta+\gamma}{C} t^{\frac{\gamma}{\beta+\gamma}}
\end{aligned}
$$

and depend on the positive parameter $C$.

Acknowledgement. The author is grateful to N. Demni for drawing his attention to reference [14], which motivated this reseach.

## References

[1] Cépa E. Equations différentielles stochastiques multivoques. Sém Probab. XXIX, Lecture Notes in Math. 1613,86-107, Springer 1995.
[2] Cépa E. Problème de Skorohod multivoque. Ann. Probab. 26,500-532, 1998.
[3] Cépa E., Lépingle D. Diffusing particles with electrostatic repulsion. Probab. Theory Related Fields 107,429-449, 1997.
[4] Dai J.G., Williams R.J. Existence and uniqueness of semimartingale reflecting Brownian motions in convex polyhedra. Theory Probab. Appl. 40,1-40, 1996.
[5] Dieker A.B., Moriarty J. Reflected Brownian motion in a wedge: sum-of-exponential stationary densities. Elec. Comm. Probab. 15,0-16, 2009.
[6] Fernholz E.R. Stochastic portfolio theory. Applications of Mathematics 48, Springer Verlag, New York, 2002.
[7] Harrison J.M. The diffusion approximation for tandem queues in heavy traffic. Adv. in Appl. Probab. 10,886-905, 1978.
[8] Harrison J.M., Reiman M.I. Reflected Brownian motion in an orthant. Annals Probab. 9,302-308, 1981.
[9] Harrison J.M., Williams R.J. Brownian models of open queueing networks with homogeneous customer populations. Stochastics 22,77-115, 1987.
[10] Ichiba T., Karatzas I. On collisions of Brownian particles. Ann. Appl. Probab. 20,951-977, 2010.
[11] Ichiba T., Karatzas I., Prokaj V. Diffusions with rank-based characteristics and values in the nonnegative quadrant. arXiv:1202.0036v2.
[12] Ichiba T., Papathanakos V., Banner A., Karatzas I., Fernholz R. Hybrid Atlas models. Ann. Appl. Probab. 21,609-644, 2011.
[13] Lépingle D. Boundary behavior of a constrained Brownian motion between reflecting-repelling walls. Probab. Math. Stat. 30,273-287, 2010.
[14] O'Connell N., Ortmann J. Product-form invariant measures for Brownian motion with drift satisfying a skew-symmetry type condition. arXiv:1201.5586v2.
[15] Pal S., Pitman J. One-dimensional Brownian particle systems with rank-dependent drifts. Ann. Appl. Probab. 18,2179-2207, 2008.
[16] Varadhan S.R.S., Williams R.J. Brownian motion in a wedge with oblique reflection. Comm. Pure Appl. 38,405-443, 1985.
[17] Williams R.J. Reflected Brownian motion in a wedge: semimartingale property. Z. Wahr. verw. Gebiete 69,161-176, 1985.
[18] Williams R.J. Reflected Brownian motion with skew symmetric data in a polyhedral domain. Probab. Theory Related Fields 75,459-485, 1987.


[^0]:    Key words and phrases. Singular stochastic differential equation; electrostatic repulsion; Bessel process; product form stationary distribution.
    AMS classification: 60H10; 34A12.
    ${ }^{1}$ Université d'Orléans, MAPMO-FDP, F-45067 Orléans (dominique.lepingle@univ-orleans.fr)

