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Evidence of frost-cracking inferred from acoustic emissions in a high-alpine rock-wall

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Abstract

Ice formation within rock is known to be an important driver of near-surface frost weathering as well as of rock damage at the depth of several meters, which may play a crucial role for the slow preconditioning of rock fall in steep permafrost areas. This letter reports results from an experiment where acoustic emission monitoring was used to investigate rock damage in a high-alpine rockwall induced by natural thermal cycling and freezing/thawing. The analysis of the large catalog of events obtained shows (i) robust power-law distributions in the time and energy domains, a footprint of rock micro-fracturing activity induced by stresses arising from thermal variations and associated freezing/thawing of rock; (ii) an increase in AE activity under sub-zero rocktemperatures, suggesting the importance of freezing-induced stresses. AE activity further increases in locations of the rock-wall that are prone to receiving melt water. These results suggest that the framework of further modeling studies (theoretical and numerical) should include damage, elastic interaction and poro-mechanics in order to describe freezing-related stresses.

Keywords: frost-cracking, acoustic emission, mechanical weathering, scaling properties

1 1. Introduction

The formation of ice within rock is likely to be an important driver of nearsurface frost weathering (Hallet et al., 1991) and rock damage at the depth of

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several meters (Murton et al., 2006), and in steep terrain, this process may be crucial for the slow preconditioning of rock fall from warming permafrost areas (Gruber and Haeberli, 2007). However, the transfer of corresponding theoretical insight and laboratory evidence to natural conditions characterized by strong spatial and temporal heterogeneity of the rock properties (e.g. fracture state, 8 water content, thermal and hydraulic conductivity) and thermal conditions is nontrivial. To examine rock fracture in natural conditions, we performed a 10 pilot experiment, monitoring acoustic emissions (AE) in a high-altitude rock-11 face during a four days period. In such conditions, the mechanical loading 12 of rock results from the combination of a constant gravity load and fluctuating 13 loads related to (i) thermal stresses, arising from the gradient of the temperature 14 field, (ii) pressure variations in rock pores and cracks, due to water or to ice 15 formation and (iii) short-term external loading such as earthquakes. While large 16 thermal stresses can only occur close to the rock surface, ice formation in pores, 17 cracks and fractures can potentially generate large stresses at greater depths, as 18 suggested by theoretical and lab studies. Reporting a preliminary analysis of the 19 microseismic activity monitored at a high alpine ridge, Amitrano et al. (2010) 20 recently stressed the importance of ice formation in fractures as they observed 21 micro-seismic activity corresponding with particular trends of the temperature 22 that could enhance ice formation. But the lack of details in the spatial and 23 temporal distribution of the seismic events precluded the full understanding of 24 the relationship between temperature evolution and related ice formation at 25 small spatial scale and the triggering of seismic events. 26

The mechanical loading of rocks involves local inelastic processes that produce elastic wave propagation so called acoustic emission (AE) at small scales and micro-seismicity (MS) at larger scales. Beside the common physical origin of the elastic wave emission, essentially induced by the propagation or shearing of cracks, these two terms denote differences in the frequency content of the recorded signals corresponding to sources of different size (see Hardy (2003) for a full presentation). MS relates to the range 1-10³ Hz whereas AE relates to the range 10^4 - 10^6 Hz. The corresponding source size is 1-10³ m for MS and ³⁵ 10⁻³-10⁻¹ m for AE. The material attenuation, that increases with frequency,
³⁶ precludes the detection of AE after approximatively 1 m of wave propagation,
³⁷ whereas MS can be detected at larger distances (up to km).

Measuring AE or MS activity therefore provides a powerful technique to 38 monitor the evolution damage at different scales. Due to their wide frequency 39 range, the simultaneous recording of AE and MS currently is technically not 40 possible. The AE has been extensively used as a tool at the laboratory rock 41 sample scale [e.g. Lockner (1993)] whereas MS has been mostly used at larger 42 scales in order to study seismicity and rockburst in mines, tunnels or quarries 43 (Hardy, 2003). In all these cases, AE/MS are considered to be an indicator of 44 inelastic behavior that can be related to damage increase or to shearing of exist-45 ing fractures (Cox and Meredith, 1993; Lockner, 1993). Several recent studies 46 report MS monitoring of slope instability (Amitrano et al., 2010; Gaffet et al., 47 2010). The originality of our study is to apply high frequency AE monitoring, 48 a technique traditionally used in laboratory experiments, to investigate rock 49 damage during freezing, in field conditions. The main advantage of using high 50 frequency monitoring is the sensitivity to emissions of relatively small energies. 51 This allows us to obtain a large catalog of events within a short monitoring 52 period, which is crucial to perform statistical analyses. At such high frequen-53 cies, acoustic signals are attenuated within about a meter, which determines the 54 spatial scale of our study. This is an advantage as most of the acoustic activity 55 related to freezing can be expected to occur within a meter from the surface. 56 This technique finally offers a high temporal resolution, as event rates up to 10^3 57 per second can be detected. 58

Fracturing dynamics during mechanical loading, usually displays scaling properties in the domains of size, space and time (Alava et al., 2006; Sethna et al., 2001). In the domain of size (magnitude) for example, the seismic events induced by damage processes display a power-law (PL) distribution, $N(s) \sim s^{-b}$, where s is an estimate of the event size (e.g. the maximum amplitude of the AE signal or its energy), N(s) is the probability distribution function (PDF) and b is a constant. This distribution is equivalent to the well-known Gutenberg-

Richter relationship observed for earthquakes (Gutenberg and Richter, 1954). 66 Scaling properties in space and time of the events have also been reported, char-67 acterizing their spatial and temporal clustering. The emergence of these scaling 68 properties is considered to be a universal feature of the damage dynamics in 69 heterogeneous media (Alava et al., 2006) as it is observed in a very robust man-70 ner for various loading conditions, various materials and scales ranging from 71 the micrometer (microcracks) to thousands of kilometers (the Earth's-crust or 72 the sea ice cover). In this letter we report an original in-situ experiment of 73 AE monitoring in high altitude thermal conditions. We show that AE activity 74 resulting from natural thermal cycling and induced freezing/thawing shows sim-75 ilar scaling properties, suggesting that the local stress fluctuations encountered 76 are high enough to induce micro-fracturing. 77

78 2. Measurement site and instrumental setting

The measurement site is a south-facing cliff of granitic gneiss (Wegmann and Keusen, 1998) that is situated at an elevation of 3500 m a.s.l. in the Swiss Alps. The local mean annual air temperature is about -7.3° C (1961–1990), whereas mean annual rock temperatures near the surface are between -2 and -3° C in this south face (Hasler et al., 2011). The site is directly next to the highaltitude research station Jungfraujoch and thus can be measured with standard AE equipment housed inside a heated building.

A six-channel high-frequency acquisition board (Mistras, Euro Physical Acous-86 tic) with 16-bit resolution and 10 V maximum amplitude was used. The AE 87 piezo-electrical sensors (EPA R6I) had an operating frequency range of 10-88 150 kHz with a peak sensitivity at 60 kHz. They included a pre-amplifier of 89 40 dB and were connected with 20 meters coaxial cables. The system was con-90 tinuously sampling with a frequency of 1 MHz and basic signal characteristics 91 such as time, amplitude, energy, duration, spectral content were calculated in 92 real time for events over 35 dB amplitude. The maximum noise amplitude has 93 been measured to be near 1 mV (30 dB), indicating that recorded events were 94

not induced by noise fluctuations. Full waveforms of 2 ms including a pre-trigger 95 of 0.4 ms were recorded for events over 40 dB (4 mV). Using an ultrasonic cou-96 pling gel (Sofranel), sensors were pressed on a steel plate with rubber bands. 97 We verified in the lab that the coupling gel behaves similarly for various tem-98 perature and do not generate AE when freezing (Weber et al., 2012). The steel 99 plates were screwed onto extension bolts (10 mm diameter) anchored about 5 cm 100 deep in the rock (Figure 1A). Each sensor installation was protected from water 101 with a plastic sleeve (Figure 1B). We verified that the crumpling of the plastic 102 cover due to wind or other factors do not cause AE. This has been tested during 103 the installation by moving by hand the plastic cover and no AE were generated. 104 Moreover, the windy periods we observed during the experiment have not been 105 recognized to generate AE. 106

Figure 2 shows an example of AE trace recorded at channel 1, with the trigger level, the maximal and minimal amplitude and the method for calculating the signal energy.

The site shows some heterogeneities in terms of the spatial distribution of 110 fractures (coexistence of compact and fractured zones), microtopography (con-111 vex and concave zones) and hydraulic conditions (dry and wet zones). Sensor 112 locations (Figure 1C), referred to as AE1–AE6, have been chosen to investigate 113 different configurations of fractured or compact rock and wet or dry conditions. 114 AE1 is installed in compact and homogeneous rock that is barely fractured in 115 a radius of about one meter. The rock surface bulges out slightly in this area. 116 It has been uncovered from snow one hour before installation, with some snow 117 remaining about 30 cm below the sensor. AE2 has compact, lightly fractured 118 rock that is slightly concave outward and subject to melt water flow from above. 119 AE3 is located in fractured, concave but rather dry rock. AE4 is in fractured 120 rock that is convex and apparently dry in surface. But a gully containing snow 121 is located 0.5 m above it that could supply water in the vicinity. AE5 is installed 122 in compact rock in an overhang underneath a gully and could receive melt water 123 through fractures. AE6 is installed in compact rock in a gully receiving melt 124 water from above. AE5 and AE6 are located near a fracture zone in the deepest 125

part of the gully; the rock mass to their left is slightly overhanging and may
be prone to movement. AE1-AE4 were installed on 6 April, AE5 and AE6 on
7 April 2010. All sensors were uninstalled on 10 April 2010. The Permasense
measurement site (Hasler et al., 2008; Beutel et al., 2009) at the same location
provides rock temperature data, measured close to sensor AE3, at depths of 10,
35, 60 and 85 cm (Hasler et al., 2011).

¹³² 3. Results

The system has been operated continuously for four days for AE1–AE4 and 133 three days for AE5–AE6. Air temperatures fluctuated between -2 and -10° C 134 during this time. There were virtually no clouds and radiative diurnal cycles 135 caused near-surface rock temperatures to rise to 10° C during the day and cool 136 to -5° C during nights. Diurnal thawing penetrated about 20 cm deep into the 137 rock wall which remained continuously frozen at greater depths. Due to snowfall 138 in the days before, small snow patches in concave portions of the rock wall locally 139 provided melt water flow during the day. Several thousand events were recorded 140 at each sensor, allowing a robust statistical analysis (Table 1). Because of the 141 large sensor spacing, attenuation precluded the detection of individual events 142 on several sensors. Given the frequency range at which the sensors operate, 143 the detected events can be expected to have their source within the meter scale 144 around the sensor. Since event source localization is impossible, detected events 145 are thus considered to be close to the receiving sensor. 146

Figure 3 reports the time evolution of AE activity and temperature of the 147 rock at 10, 35 and 60 cm depth as a function of the hour of the day. It high-148 lights the connection between the daily fluctuations of AE activity and that 149 of rock temperature during the monitoring period. For the entire duration of 150 the experiment, the temperature at 35 cm depth remains below 0° C whereas it 151 remains under 0°C at 60 cm depth. The temperature at 10 cm depth exhibits 152 freeze-thaw cycles with temperature ranging from -5° C to 10° C. At this depth, 153 the refreezing can by seen by the plateau characterizing the temperature near 154

¹⁵⁵ zero in the evening (i.e. at $h \simeq 20$ on days 1–2) when temperature decreases ¹⁵⁶ more slowly due to latent heat of water. This time period corresponds to a ¹⁵⁷ small peak of the AE activity visible on AE1, AE3, AE5 and AE6. We recall ¹⁵⁸ that our analysis is based on a single point measurement of rock temperature. ¹⁵⁹ The spatial variability of temperature at different locations could thus partly ¹⁶⁰ explain the shift seen in AE activity observed for the different sensors.

The AE activity appears to be significantly larger when the temperature 161 is below zero. The largest AE activity takes place during colder (night-time) 162 periods, when the near-surface of the rock refreezes. This is observed for all 163 the sensors with the notable exception of AE5. Sensors AE1, AE3 and AE5 164 show activity peaks significantly smaller than AE2, AE4 and AE6. AE1 and 165 AE3 are located in dry areas whereas AE2, and AE6 are located in wet areas 166 with melting water coming from upslope. AE4 was located in an apparently 167 dry area but the presence of snow 0.5 m above could provide melting water that 168 was not visible on the surface. In contrast to the other sites, any meltwater 169 reaching AE5 would have to percolate through fractures. No water comes from 170 the surface due to its situation under an overhang. These observations express 171 the large spatial variability of the near-surface AE activity. They also suggest 172 that the availability of water has a strong control on AE activity. In order 173 to verify the dependence between AE activity and temperature we calculate 174 the distribution of events as a function of the temperature (Fig. 4 left). The 175 relationship between negative temperatures and AE activity appears clearly for 176 all the sensors except AE5. Between 70% and 95% of the events are recorded 177 during negative temperature periods, whereas for AE5, this proportion is only of 178 37% (see Tab. 1 for details). One may note also that a slight AE peak is visible 179 around temperature zero. The difference in the amplitude of the AE activity 180 between dry and wet areas is confirmed. In order to take into account the fact 181 that much more time is spent in each temperature increment below zero than 182 above, we normalized the event number by the time spent into each temperature 183 interval (Fig. 4 right). the temperature dependence appears stronger for sensors 184 AE2 AE4 and AE6, and weaker for sensors AE1 and AE3. The particularity of 185

¹⁸⁶ the sensor AE5 appears stronger.

We analyzed the scaling properties of the AE focusing on the domains of 187 energy and time and described other striking behavior of the measured data. 188 For each sensor we calculated the distribution of event energy for the entire 189 duration of the experiments. The PDF's of event energy (Fig. 5) show a power-190 law (PL) behavior spanning several orders of magnitude with an exponent b =191 1.55 ± 0.05 . For E smaller than 100, the pdf shows a clear departure from PL 192 trend characterized by a decrease of the slope toward the smallest events. This 193 effect is commonly related to the completeness of the catalog (see e.g. Wiemer 194 and Wyss, 2000). Under this value, due to the signal attenuation, the smallest 195 events are detected only on the vicinity of the sensor and so the sampling of these 196 events is incomplete. On the contrary, above this value there is no statistical 197 sampling bias on the PL trend can be estimated with confidence. 198

We verified that the completeness and the PL trend remain unchanged when selecting the events recorded during positive or negative temperature periods. We also verified the effect of selecting only events of energy larger than the completeness. The trends observed on Figures 3 and 4 remain the same, in particular with respect to the temperature dependence of AE activity. The only notable effect is the reduction of the amplitude of AE peaks.

As the PL distribution is characteristic of the rupture processes (Alava et al., 205 2006), the later observation is the first direct evidence published showing that 207 natural thermal cycling and associated freezing/thawing induce near-surface 208 damage in a rock wall.

In order to investigate scaling properties of the temporal distribution of AE events we use the correlation integral (Grassberger and Procaccia, 1983),

$$C(\Delta t) = 2\mathcal{N}(\Delta t)/(N(N-1)) \tag{1}$$

where N is the total number of damage events, $\mathcal{N}(\Delta t)$ is the number of pairs of events separated by a time smaller than Δt . This integral expresses how the events are distributed in time. If the correlation integral exhibits a PL $C(\Delta t) \sim$ Δt^{D_2} , the population can be considered as a fractal set, i.e. characterized by a scaling invariance in time domain, with a correlation dimension D_2 . A value of D_2 smaller than 1 indicates a time clustering, i.e. the probability of observing an event is larger when a previous one has occurred within a short time. In other words, an event is more likely to occur within a short separating time from the previous one. Departure from power-law trend or slope changes indicate the existence of characteristic time scale limiting the extent of scale invariance.

Figure 6 shows the correlation integrals obtained for each sensor. At small 221 time scales, a first PL trend is identified over two to three orders of magnitude 222 in time, for all sensors except AE5, with a correlation dimension D_2 between 223 0.75 and 0.9. The value $D_2 < 1$ expresses the temporal clustering of AE events, 224 i.e. the detected events are strongly correlated over this temporal scale range. 225 The extent of the PL trend towards small time scales is limited by the duration 226 of the recording (1 ms) so no event can be detected with a lower separating 227 time. The upper limit of PL extends to about 0.5s for AE1 and AE4, 1s for 228 AE2 and AE6, and $5 \,\mathrm{s}$ for sensor AE3. This mean that the events are strongly 229 correlated in time within this time scale. So this value can be interpreted as 230 the duration of correlated events series that we may consider as sequences of 231 cascading events, i.e. the temporal correlation length, i.e. the duration for 232 which a pair of events is more correlated in time than in a random serie. For 233 larger separating time the correlation integral displays a rather flat shape until 234 recovering a secondary PL trend. In this region, AE4 and AE5 are characterized 235 by a slope $D_2 \simeq 0.9$ indicating a slightly clustered temporal distribution. The 236 duration of the correlated sequences is about 2-3 hours. For AE3 the secondary 237 power-law trend appears of poor quality, with a low D_2 value corresponding 238 to highly clustered events. On the other hand, a clear secondary power law 239 trend of exponent $D_2 \simeq 0.75$ is observed for AE1, AE2 and AE6, a temporal 240 correlation dimension similar to the one observed at smaller time scales. This 241 power law trend extends up to about $\Delta t = 10 - 12$ hours, indicating scaling 242 properties spreading over the duration of the freezing period. 243

244 **4.** Discussion and conclusions

The first striking aspect of the measurements presented is the relationship 245 between negative temperatures and AE activity, suggesting that damage is re-246 lated to freezing-induced stresses. As ice is a better wave transmitter than liquid 247 water one may ask if this observation could be an artefact induced by the lack 248 of detection instead of a lack of AE at temperatures $> 0^{\circ}$. In order to verify 249 this, we compare the two sets of events distinguishing positive and negative 250 temperature periods. A first evidence is that the completeness is the same for 251 the two sets (see section 3). This indicates that there is no significant changes 252 in the event detection. As the attenuation is known to affect the frequency 253 content (Hardy, 2003), we analyse the evolution of the mean frequency of the 254 events. Surprisingly, we observed a slight increase of the mean frequency for 255 warmer periods (Fig. 7), particularly for AE2, AE5 and AE6. This could be 256 interpreted as an effect of the thermal dilation of cracks. The rock material in 257 between cracks expands inducing a closure of cracks and then a reduction of 258 the attenuation. The fact that variations of the frequency are limited could be 259 explained by the resonant nature of the sensors we used. The cracks closure 260 can also induce AE, but this has been shown to be limited compared to the 261 crack shearing process (e.g. Moradian et al., 2010). We also examine the rela-262 tionship between the signal amplitude and duration, considering that a higher 263 attenuation should correspond to a shorter duration for a given amplitude. We 264 observed no (or slight) variations regarding the temperature. Consequently, the 265 increase of AE activity in sub-zero temperatures is unlikely do be an artefact in 266 the AE detection. 267

We now discuss the different mechanisms that could induce damage through thermal cycling. A possible mechanism for explaining the relationship between AE and temperature is the differential thermal dilation inducing thermal stressing. In the present case, the amplitude of thermal gradients (in space and time) is limited compared to the one needed for inducing damage within intact rock, as observed at the laboratory or estimated analytically considering

thermo-mechanical coupling (e.g. Fredrich and Wong, 1986; Wai et al., 1982). 274 In the case of already damaged or fractured material, it has been shown that 275 daily thermal cycling can induce shearing along existing fractures (Gunzburger 276 et al., 2005) even in absence of freeze-thaw cycles. The most favorable peri-277 ods are the ones corresponding to high thermal gradient in time $(\delta T/\delta t)$ and 278 in space $(\delta T/\delta z, z)$ representing the depth). Referring to our experiment, this 279 mechanism could induce AE events when the absolute value of the spatial tem-280 perature gradient is high, independently of positive or negative temperature. 281 We estimated the temperature gradient in depth as the difference between tem-282 perature measured at 10 and 35 cm depth, divided by their separating distance 283 $\delta T/\delta z = (T_{35cm} - T_{10cm})/0.25m$. Figure 8 shows that most of the AE activity 284 occurred when $\delta T/\delta z > 20^{\circ}C/m$, when the rock 10 cm below the surface is 285 much cooler than 35 cm below. 286

Thermo-elastic stresses arise when non-uniform (spatial) gradients in temperature develop in elastic materials, which tend to result from rapidly varying temperatures. Note that, somewhat non-intuitively, thermo-elastic stresses would no be expected in unconfined elastic media if the temperature gradient is uniform even if temperatures vary considerably in space. So this mechanism can be evoked only when the temporal gradient is also high.

Figure 8 shows the temperature gradient in time at 10 cm depth. The 293 maximum absolute value of the temporal gradient occurs at approximatively 13 294 and 18 hours. These two periods corresponds to AE activity peaks visible on 295 Figures 3 and 8, suggesting that AE could be related to fast contraction/dilation 296 of the rock. As the temperature in depth is stable, the temperature gradient is 297 highly correlated with the shallow temperature (Fig. 8). The highest absolute 298 value of the gradient corresponds to the minimum of the shallow temperature. 299 Consequently it is difficult to distinguish between the impacts of the temperature 300 and of the gradient of temperature on the AE activity. 301

AE activity was shown to be more intense at locations with melt-water, where large bursts of events where recorded during night-time refreezing. This suggests that water freezing plays an important role in the mechanical load-

ing. A possible origin for the damage is the volume expansion of freezing water 305 contained in fractures, that is often evoked for explaining cryo-fracturing. The 306 volume change provokes considerable ice pressure in the fractures that can prop-307 agate and release the induced stresses. This could explain the AE activity peak 308 observed when temperature decreases near zero (Fig. 4). Note that the tempo-309 ral gradient is very small at this time precluding the effect of fast contraction 310 evoked before. When occurring within in a porous medium, the phase change 311 of water spreads over a temperature range that depends on the pore dimension 312 (Coussy, 2005) and not only for $T \approx 0^{\circ}C$. This could explain why increased 313 AE is sustained across the entire observed range of negative temperatures and 314 not restricted to temperatures near $0^{\circ}C$. 315

Another possible mechanism related to ice formation is the cryo-suction or 316 ice segregation (Hallet et al., 1991; Coussy, 2005). Ice in pores or cracks is 317 surrounded by an unfrozen water film due to disjoining (intermolecular) forces. 318 Disjoining forces between the ice and the rock can then cause ice-filled cracks to 319 widen as water is drawn in from the surrounding medium by a free-energy gra-320 dient. This phenomenon operates at temperatures below zero in a temperature 321 range depending mostly on the pore size distribution, permeability and fracture 322 mechanical properties of the rock (Hallet et al., 1991). This could explain why 323 AE activity increases for temperatures several degrees below $0^{\circ}C$. Moreover, 324 one may keep in mind that the surface was probably much cooler than at a 325 depth of 10 cm during the freezing period. 326

The former mechanism is related to an increase of the ice pressure within 327 fractures. One may reasonably ask if the AE could be induced by the cracking 328 of ice itself instead of the embedding rock. Ice growth is supposed to induce 329 compression stress in the ice and, due to the reaction of the embedding material, 330 tensile stress at the cracks tips. The ice, as rock do, behaves according the 331 Coulomb failure criterion (Weiss and Schulson, 2009). So its strength is larger 332 in compression than in tension. The symmetrical case applies for the rock 333 around the cracks: its strength is smaller in tension than in compression, more 334 over the crack tips acts as a stress concentrator. This suggests that the rupture 335

is more likely to occur in the rock rather than in the ice. To test this,one may
realize waves velocity measurement during periods without ice, to verify that
the damage induced in the rock is increasing with time.

Beside the results common to all sensors we also observed differences be-339 tween sensors. The amplitude of the AE activity seems to depend both on the 340 fracture state and the water availability. The relationship between AE activity 341 and temperature displays also some variations from sensor to sensor. This could 342 be attributed to the fact that we use a single temperature measurement point 343 that we considered as representative of the whole monitoring area. Spatial vari-344 ability probably exists (cf Gubler et al., 2011) that we were not able to take into 345 account in more detail here. The sensor AE5 behaves differently than others 346 in the sense that the AE crisis occurred during the day whereas the night time 347 was relatively quiet. In contrast to all others, this sensor was sensitive to AE 348 activity associated with fractures supplied with meltwater from above. The AE 349 activity could then be related to water pressure increase in the fractures that 350 has been shown to be able to produce seismicity even with limited water table 351 fluctuations (Guglielmi et al., 2008). Here we based our discussion on qualita-352 tive considerations of water availability. For a more quantitative discussion, it 353 is clear that further measurements are necessary, in particular for assessing the 354 amount of liquid water available for the formation of ice, and the water pres-355 sure. This could help us to define the respective roles played by thermal dilation, 356 ice formation and cryosuction in the generation of AE. Such measures should 357 be spatially distributed for better understand the spatial heterogeneity of AE 358 activity. 359

The spatial variability we observed concerning AE activity, and consequently damage, has major consequences on the potential erosion induced by frostcracking. Considering the frost-cracking as homogeneous and deriving the potential of damage from calculation based on the air temperature only (e.g. Hales and Roering, 2007; Delunel et al., 2010), could lead to poor estimates of weathering. Another important point for the evolution of the morphology in mountain areas is that fractured zones appears more prone to frost-cracking. This enhances the localization of damage that tends to concentrate in already damaged
zones. As a consequence, the spatial variability of frost-cracking should increase
with time.

The second striking aspect is the PL distribution of the AE event energy, 370 which is clearly identified on all sensors with an exponent $b = 1.55 \pm 0.05$ and 371 verified over six orders of magnitude. This scaling relationship is the signature of 372 the rock micro-fracturing activity. Beside the claim of the possible universality 373 of the *b*-value, this parameter has been proposed to be dependent on various 374 parameters, in particular on the strength heterogeneity of the material, the 375 applied stress and the proximity of failure (see Amitrano, 2012, and references 376 therein for an extended discussion). 377

The *b*-value of 1.55 for the energy pdf (corresponding to 0.55 for the cdf), we 378 observed for our data, is in the range expected for rocks experiencing uniaxial 379 or triaxial compression stress state in laboratory experiences (Lockner, 1993) 380 and is very close to that observed prior to the peak load (Amitrano, 2003; 381 Lockner, 1993) and for creep of compression tests after the onset of tertiary 382 creep (Grgic and Amitrano, 2009) and for seismic forerunners recorded in a cliff 383 before its collapse (Amitrano et al., 2005) although the loading mode is very 384 different. This provides an indication that the stress induced by thermal cycling 385 and/or freezing/thawing of water in rock pores and cracks reaches values close 386 to the rock strength. The fact that the b-value is found to be similar for all 387 sensors indicates that strength heterogeneity and stress, which are the two main 388 factors influencing this parameter, are comparable in the different locations we 389 investigated. 390

The time distribution of AE also reveals power-law distribution, which is a supplementary indication of the complex behavior the frost induced damage and of the presence of strong interaction between damage events. The damage activity appears to be clustered at two different time scales, for $\Delta t < 5 - 10s$ and for $100s < \Delta t < 10h$. This could be interpreted as the effect of two loading mechanisms for which the interaction operates at different time scale. The first could correspond to cascading events related to the elastic redistribution of stress when damage occurs. The second could be related to the reloading induced by temperature changes and/or water migration that operates at longer time scales. In order to investigate these questions more precisely, further measurements estimating the source depth of AE events, the evolution of the liquid water content and of the temperature at depth are crucial. Moreover, the mechanisms through which freeze/thaw-induced stresses can occur also need to be better differentiated under in-situ conditions. differentiated under in-situ conditions.

The experiments of Hallet et al. (1991), revealing sustained microfracturing 405 activity throughout a 3-day period during which temperature and temperature 406 gradients were held constant in the sample, demonstrated the ability of ice 407 segregation to fracture rock. However, in the field, such conditions are never 408 achieved and it may be difficult to robustly distinguish the role of volumetric 409 expansion (as water turns into ice) from that of ice segregation. A number of 410 factors such as solutes, pressure, pore size and pore material can depress the 411 freezing point of water contained in rock (Krautblatter et al. 2010) down to 412 several degrees below 0°C. This explains why ice formation in rock, such as in 413 any porous media, is progressive (Coussy and Fen-Chong, 2005) occurring over 414 a whole range of sub-zero temperatures. Volumetric expansion-induced damage 415 could potentially occur over this whole range of temperatures. Using theoretical 416 arguments, Vlahou and Worster (2010) reported that volumetric expansion can 417 only develop significant pressures (~ 10 MPa) in water saturated confined (spher-418 ical) cavities larger than 1cm in diameter of very low permeability $(10^{-15} cm^2)$. 419 This prediction basically rules out the role of volumetric expansion on rock 420 fracture, since such conditions are very seldom (if never) achieved in nature. 421 Contrastingly, a different body of work from the cement and concrete research, 422 has shown theoretically that both crystallization and micro-cryosuction mech-423 anisms can induce pressures of several tens of MPa in a single water-saturated, 424 micrometer-size pore embedded in a porous material (Coussy and Fen-Chong, 425 2005). While the pore structure of rocks certainly does contain such small fea-426 tures (Fredrich et al, 1995), a partial saturation of pore space, which is often 427 achieved in natural conditions, might give a completely different picture. 428

The results we have reported in this letter show the feasibility of studying 429 rock damage under natural conditions of thermal cycling and freezing using the 430 AE technique. AE activity was shown to significantly increase in sub-zero tem-431 peratures, especially in locations of the rock-wall that are prone to receiving 432 melt water, suggesting that freezing-induced stresses contribute to rock dam-433 age. The robust PL distribution of AE event in the domains of energy and 434 time distributions, a common observation in rupture processes dynamics, is an 435 indication that damage is acting in such conditions. These results finally sug-436 gest that the framework of further modeling studies (theoretical and numerical) 437 should include damage, elastic interaction and poro-mechanics in order to de-438 scribe freezing-related stresses. 439

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Sensor	AE1	AE2	AE3	AE4	AE5	AE6
Nb. AE events	1.2×10^4	1.5×10^4	$0.2 imes 10^4$	1.8×10^4	$0.3 imes 10^4$	1.6×10^4
R	0.80	0.86	0.70	0.94	0.37	0.75
Water	Snow	From	Dry	From	Through	From
supply	cover	above		above	fractures	above
Fractured	No	Slightly	Yes	Yes	Yes	No
Topography	Slightly	Slightly	Concave	Convex	Concave	Concave
	convex	concave				

Table 1: Number of AE events recorded and configuration at each sensor location. R is the proportion of events recorded during negative temperature periods.

Figure 1 : Sensor installation at rock surface (A) and weatherproofing (B) shown for AE4. Sensor locations (C) are referred to as AE1-AE6 in the text. Figure 2 : Example of an acoustic emission trace recorded on channel 1. The dotted lines indicate the trigger above which the trace is recorded. The signal energy, E, is calculated by summing the squared amplitude over the during of the signal. The maximal and minimal amplitude are shown for information.

Figure 3: AE event rates and rock temperature at 10, 35 and 60 cm depth as a function of the time. Time corresponds to the hour in the day.

Figure 4: AE event rates as a function of the temperature at 10 cm depth. 558 Left: Number of events detected in each temperature interval. The bins are 559 $0.5^{\circ}C$ width. Right: Event number detected in each temperature interval nor-560 malized by the time spent in each interval. Sensors 1 3 and 5 (Top), which 561 are in relatively dry sites, have an AE activity smaller than sensors 2 4 and 562 6. Except sensor 5, they all show a clear dependence on the temperature, with 563 a huge increase of the AE for negative temperatures and reduced activity for 564 positive temperatures. When normalizing by the time spent in each bin (right 565 draft) the temperature dependence appears stronger for sensors AE2 AE4 and 566 AE6, and weaker for sensors AE1 and AE3. The particularity of the sensor AE5 567 appears stronger. 568

Figure 5: AE event energy distribution for all channels for the complete experiment. The Power-law trend with b = 1.55 is shown as guideline. The completeness energy is also shown to correspond to 100 (Arbitrary unit).

Figure 6: Temporal correlation integral of AE events for all channels. The dashed lines provide guidelines corresponding to exponents $D_2 = 0.75$ and $D_2 =$ 0.9.

Figure 7 Mean frequency of the AE event for each sensors. Blue dots correspond to the individual values whereas red crosses give the moving average for 100 successive events and 90% of overlap.

Figure 8: Top: AE activity for all the sensors. Center: In depth gradient of the temperature. Bottom: rate of change temperature. All values are plotted as a function of hours.



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