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#### GROUND STATES FOR A STATIONARY MEAN-FIELD MODEL FOR A NUCLEON

#### MARIA J. ESTEBAN<sup>1</sup> AND SIMONA ROTA NODARI<sup>2,3</sup>

ABSTRACT. In this paper we consider a variational problem related to a model for a nucleon interacting with the  $\omega$  and  $\sigma$  mesons in the atomic nucleus. The model is relativistic, and we study it in a nuclear physics nonrelativistic limit, which is of a very different nature than the nonrelativistic limit in the atomic physics. Ground states are shown to exist for a large class of values for the parameters of the problem, which are determined by the values of some physical constants.

#### 1. INTRODUCTION

This article is concerned with the existence of minimizers for the energy functional

$$\mathcal{E}(\varphi) = \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{(1 - |\varphi|^2)_+} \, dx - \frac{a}{2} \int_{\mathbb{R}^3} |\varphi|^4 \, dx \tag{1.1}$$

under the  $L^2$ -normalization constraint

$$\int_{\mathbb{R}^3} |\varphi|^2 \, dx = 1. \tag{1.2}$$

More precisely, for a large class of values for the parameter a, we show the existence of solutions of the following minimization problem

$$I = \inf\left\{\mathcal{E}(\varphi); \ \varphi \in X, \int_{\mathbb{R}^3} |\varphi|^2 \, dx = 1\right\}, \tag{1.3}$$

where

$$X = \left\{ \varphi \in L^2(\mathbb{R}^3, \mathbb{C}^2) \; ; \; \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{(1 - |\varphi|^2)_+} \, dx < +\infty \right\} \; . \tag{1.4}$$

We remind that  $\boldsymbol{\sigma}$  denotes the vector of Pauli matrices  $(\sigma_1, \sigma_2, \sigma_3)$ ,

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

The Euler-Lagrange equation of the energy functional  $\mathcal{E}$  under the  $L^2$ -normalization constraint is given by the second order equation

$$-\boldsymbol{\sigma} \cdot \nabla \left(\frac{\boldsymbol{\sigma} \cdot \nabla \varphi}{(1-|\varphi|^2)_+}\right) + \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{(1-|\varphi|^2)_+^2} \varphi - a|\varphi|^2 \varphi + b\varphi = 0, \qquad (1.5)$$

where b is the Lagrange multiplier associated with the  $L^2$ -constraint (1.2). Hence a solution of the minimization problem (1.3) is a solution of the equation (1.5).

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Moreover, Lemma 2.1 below proves that any  $\varphi \in X$  satisfies  $|\varphi|^2 \leq 1$  a.e. in  $\mathbb{R}^3$ . So, a minimizer for (1.3) is actually a solution of

$$-\boldsymbol{\sigma} \cdot \nabla \left(\frac{\boldsymbol{\sigma} \cdot \nabla \varphi}{1 - |\varphi|^2}\right) + \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{(1 - |\varphi|^2)^2} \varphi - a|\varphi|^2 \varphi + b\varphi = 0.$$
(1.6)

Solutions of (1.6) which are minimizers for I are called ground states.

The equation (1.6) is a equivalent to the system

$$\begin{cases} i\boldsymbol{\sigma}\cdot\nabla\chi + |\chi|^2\varphi - a|\varphi|^2\varphi + b\varphi = 0, \\ -i\boldsymbol{\sigma}\cdot\nabla\varphi + (1-|\varphi|^2)\chi = 0. \end{cases}$$
(1.7)

As we formally derived in a previous paper ([1]), this system is the nuclear physics nonrelativistic limit of the  $\sigma$ - $\omega$  relativistic mean-field model ([9, 10]) in the case of a single nucleon.

In [1], we proved the existence of square integrable solutions of (1.7) in the particular form

$$\begin{pmatrix} \varphi(x)\\ \chi(x) \end{pmatrix} = \begin{pmatrix} g(r) \begin{pmatrix} 1\\ 0 \end{pmatrix}\\ if(r) \begin{pmatrix} \cos \vartheta\\ \sin \vartheta e^{i\phi} \end{pmatrix} \end{pmatrix},$$
(1.8)

where f and g are real valued radial functions. This ansatz corresponds to particles with minimal angular momentum, that is, j = 1/2 (for instance, see [8]). In this model, the equations for f and g read as follows:

$$\begin{cases} f' + \frac{2}{r}f = g(f^2 - ag^2 + b), \\ g' = f(1 - g^2), \end{cases}$$
(1.9)

where we assumed f(0) = 0 in order to avoid solutions with singularities at the origin, and we showed that given a, b > 0 such that a - 2b > 0, there exists at least one nontrivial solution of (1.9) such that

$$(f(r), g(r)) \longrightarrow (0, 0) \quad \text{as} \quad r \longrightarrow +\infty.$$
 (1.10)

In this paper, we prove the existence of solutions of the above nuclear physics nonrelativistic limit of the  $\sigma$ - $\omega$  relativistic mean-field model without considering any particular ansatz for the nucleon's wave function.

Note that (1.6) is the Euler-Lagrange equation of the energy functional

$$\mathcal{F}(\varphi) = \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{1 - |\varphi|^2} \, dx - \frac{a}{2} \int_{\mathbb{R}^3} |\varphi|^4 \, dx \tag{1.11}$$

under the  $L^2$  normalization constraint. In the Appendix, we prove that the energy functional  $\mathcal{F}$  is not bounded from below. So, trying to find solutions of (1.6) which minimize the energy  $\mathcal{F}$  is hopeless and the definition of ground states for (1.6) based on this functional is not clear.

In our previous work ([1]), we showed that for all the solutions of (1.9) which are square integrable,  $g^2(r) < 1$  in  $[0, +\infty)$ . Hence, according to this result, we conjecture that a solution of (1.6) has to satisfy  $|\varphi|^2 \leq 1$  a.e. in  $\mathbb{R}^3$ . As we prove in the Appendix, this assumption is also justified when we consider the intermediate model

$$\varphi = \left(\begin{array}{c} u\\0\end{array}\right) \tag{1.12}$$

with  $u : \mathbb{R}^3 \to \mathbb{R}$  and a > b. Moreover, in the physical literature finite nuclei are described via functions  $\varphi$  such that, in the right units,  $|\varphi|^2 \leq 1$  and  $|\varphi|$  is rather flat near the center of the nucleus, and is equal to 0 outside it, see [5, 2].

Note that if  $|\varphi|^2 \leq 1$  a.e. in  $\mathbb{R}^3$ , then  $\mathcal{F}(\varphi) = \mathcal{E}(\varphi)$ , and the ground states of (1.6) can be defined without further specification as the minimizers of  $\mathcal{E}$ .

The main result of our paper is the following

**Theorem 1.1.** If I < 0 there exists a minimizer of (1.3). Moreover, I < 0 if and only if  $a > a_0$  where  $a_0$  is a strictly positive constant. In particular,  $10.96 \approx \frac{2}{S^2} < a_0 < 48.06$ , where S the best constant in the Sobolev embedding of  $H^1(\mathbb{R}^3)$  into  $L^6(\mathbb{R}^3)$ .

## **Remark 1.** The upper estimate for $a_0$ is obtained by using a particular test function and is probably not optimal.

The proof of the above theorem is an application of the concentration-compactness principle ([3, 4]) with some new ingredients. The main new difficulty is due to the presence of the term  $\int_{\mathbb{R}^3} \frac{|\sigma \cdot \nabla \varphi|^2}{(1-|\varphi|^2)_+} dx$  in the energy functional. As we will see below, to rule out the dichotomy case in the concentration-compactness lemma we have to choose *ad-hoc* cut-off functions allowing us to deal with possible singularities of the integrand. This is also necessary in order to show the localization properties of  $\int_{\mathbb{R}^3} \frac{|\sigma \cdot \nabla \varphi|^2}{(1-|\varphi|^2)_+} dx$ .

In the next section, we will establish a concentration-compactness lemma in X and then apply it to prove our main result. The Appendix contains some auxiliary results about various properties of the model problem that we consider here.

#### 2. Proof of Theorem 1.1

To prove this theorem, we are going to apply a concentration-compactness lemma that we state below. The reader may refer to [3] and [4] for more details on this kind of approach. The particular shape of the energy functional, where the kinetic energy term is multiplied by a function which could present singularities as  $|\varphi|$  gets close to 1 creates some complications in the use of concentration-compactness, that we deal with by using very particular cut-off functions.

Let us introduce

$$I_{\nu} = \inf \left\{ \mathcal{E}(\varphi) \; ; \; \varphi \in X, \int_{\mathbb{R}^3} |\varphi|^2 \, dx = \nu \right\}$$
(2.1)

where  $\nu > 0$  and  $I_1 = I$ , and we make a few preliminary observations.

**Lemma 2.1** ([6]). Let 
$$\varphi \in X$$
. Then,  $\varphi \in H^1(\mathbb{R}^3, \mathbb{C}^2)$  and  $|\varphi|^2 \leq 1$  a.e. in  $\mathbb{R}^3$ .

*Proof.* First, by a straightforward calculation, we obtain

$$\int_{\mathbb{R}^3} |\nabla \varphi|^2 \, dx = \int_{\mathbb{R}^3} |\boldsymbol{\sigma} \cdot \nabla \varphi|^2 \, dx \le \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{(1-|\varphi|^2)_+} \, dx < +\infty.$$

Hence,  $\varphi \in H^1(\mathbb{R}^3, \mathbb{C}^2)$ . Next, let  $n \in \mathbb{C}^2$  such that |n| = 1. Note that for  $\varphi \in X$ ,  $1_{\operatorname{Re}\{n \cdot \varphi\} \ge 1}(\boldsymbol{\sigma} \cdot \nabla \varphi) = 0$ , a.e. in  $\mathbb{R}^3$ . Define the functions  $f = (\operatorname{Re}\{n \cdot \varphi\} - 1)_+$ 

and  $\psi = fn$ . (Note that for 2 complex vectors  $A, B \in \mathbb{C}^2$ ,  $A \cdot B$  denotes the scalar product  $\sum_{i=1}^{2} \overline{A_i} B$ , where  $\overline{z}$  stands for the complex conjugate of any complex number z).

We have  $f \in H^1(\mathbb{R}^3, \mathbb{R})$  and  $\psi \in H^1(\mathbb{R}^3, \mathbb{C}^2)$ . Moreover, for k = 1, 2, 3,

.

 $\partial_k \psi = \partial_k f n$  and  $\partial_k f = \operatorname{Re} \{ n \cdot \partial_k \varphi \} \mathbf{1}_{\operatorname{Re} \{ n \cdot \varphi \} \ge 1} = n \cdot \partial_k \psi.$ 

Hence, we obtain

$$\begin{split} \int_{\mathbb{R}^3} |\nabla f|^2 \, dx &= \int_{\mathbb{R}^3} |\nabla \psi|^2 \, dx = \int_{\mathbb{R}^3} \sum_{k=1}^3 \operatorname{Re} \left\{ \operatorname{Re} \left\{ n \cdot \partial_k \varphi \right\} n \cdot \partial_k \psi \right\} \, dx \\ &= \int_{\mathbb{R}^3} \sum_{k=1}^3 \operatorname{Re} \left\{ n \cdot \partial_k \varphi \right\} \operatorname{Re} \left\{ n \cdot \partial_k \psi \right\} \, dx = \int_{\mathbb{R}^3} \sum_{k=1}^3 \operatorname{Re} \left\{ \partial_k f \, n \cdot \partial_k \varphi \right\} \, dx \\ &= \int_{\mathbb{R}^3} \operatorname{Re} \left\{ \nabla \psi \cdot \nabla \varphi \right\} \, dx = \int_{\mathbb{R}^3} \operatorname{Re} \left\{ (\boldsymbol{\sigma} \cdot \nabla \psi) \cdot (\boldsymbol{\sigma} \cdot \nabla \varphi) \right\} \, dx \\ &= \int_{\mathbb{R}^3} \operatorname{Re} \left\{ (\boldsymbol{\sigma} \cdot \nabla \psi) \cdot 1_{\operatorname{Re}\{n \cdot \varphi\} \ge 1} (\boldsymbol{\sigma} \cdot \nabla \varphi) \right\} \, dx = 0 \end{split}$$

As a consequence, f = 0 a.e. in  $\mathbb{R}^3$  that means  $\operatorname{Re}\{n \cdot \varphi\} \leq 1$  a.e. for all  $n \in \mathbb{C}^2$  such that |n| = 1. This clearly implies that  $|\varphi| \leq 1$  a.e. in  $\mathbb{R}^3$ .

In what follows, we say that a sequence  $\{\varphi_n\}_n$  is X-bounded if there exists a positive constant C independent of n such that

$$\|\varphi_n\|_{L^2}^2 + \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_n|^2)_+} \, dx \le C \,. \tag{2.2}$$

**Lemma 2.2.** Let  $\{\varphi_n\}_n$  be a minimizing sequence of (2.1), then  $\{\varphi_n\}_n$  is X-bounded, bounded in  $H^1(\mathbb{R}^3)$  and  $I_{\nu} > -\infty$ .

*Proof.* Indeed, since  $\{\varphi_n\}_n$  is a minimizing sequence, there exists a constant C such that

$$C \ge \mathcal{E}(\varphi_n) \ge \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_n|^2)_+} \, dx - \frac{a}{2}\nu \ge \int_{\mathbb{R}^3} |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2 \, dx - \frac{a}{2}\nu$$
$$= \int_{\mathbb{R}^3} |\nabla \varphi_n|^2 \, dx - \frac{a}{2}\nu \ge -\frac{a}{2}\nu.$$

As a conclusion,  $\|\varphi_n\|_{H^1}$  is bounded independently of n and  $I_{\nu}$  is bounded from below.

**Lemma 2.3.** For all  $\nu \in (0,1)$ ,  $I_{\nu} \leq 0$ . Moreover, the strict inequality I < 0 is equivalent to the strict concentration-compactness inequalities

$$I < I_{\nu} + I_{1-\nu}$$
,  $\forall \nu \in (0,1)$ . (2.3)

*Proof.* Indeed, let  $\varphi \in \mathcal{D}(\mathbb{R}^3)$  such that  $\int_{\mathbb{R}^3} |\varphi|^2 = \nu$  and  $\int_{\mathbb{R}^3} \frac{|\sigma \cdot \nabla \varphi|^2}{(1-|\varphi|^2)_+} dx < +\infty$ , and let  $\varphi_{\gamma}(x) = \gamma^{-3/2} \varphi(\gamma^{-1}x)$  for  $\gamma > 1$ . Then

$$I_{\nu} \leq \mathcal{E}(\varphi_{\gamma}) = \frac{1}{\gamma^2} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{\left(1 - \frac{1}{\gamma^3} |\varphi|^2\right)_+} \, dx - \frac{1}{\gamma^3} \frac{a}{2} \int_{\mathbb{R}^3} |\varphi|^4 \, dx \,,$$

and letting  $\gamma \to +\infty$ , we prove  $I_{\nu} \leq 0$ .

By a scaling argument, we obtain

$$I_{\vartheta\nu} \leq \inf\left\{\vartheta^{1/3} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{(1-|\varphi|^2)_+} \, dx - \frac{\vartheta \, a}{2} \int_{\mathbb{R}^3} |\varphi|^4 \, dx |\varphi \in X, \int_{\mathbb{R}^3} |\varphi|^2 \, dx = \nu\right\},$$

and, if  $I_{\nu} < 0$ , we may restrict the infimum  $I_{\nu}$  to elements  $\varphi$  satisfying

$$K(\varphi) = \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{(1 - |\varphi|^2)_+} \, dx \ge \delta > 0 \,,$$

for some  $\delta > 0$ . Indeed, if there is a minimizing sequence  $\{\varphi_n\}_n$  of  $I_{\nu}$  such that  $K(\varphi_n) \xrightarrow[n]{} 0$ , then, by Sobolev embeddings,  $\varphi_n \xrightarrow[n]{} 0$  in  $L^p(\mathbb{R}^3)$  for  $2 and <math>I_{\nu} \ge 0$ . As a conclusion, if  $I_{\nu} < 0$ , then, for all  $\vartheta > 1$  and for all  $\nu > 0$ ,

$$I_{\vartheta\nu} < \vartheta \inf\left\{ \mathcal{E}(\varphi) | \varphi \in X, K(\varphi) > 0, \int_{\mathbb{R}^3} |\varphi|^2 \, dx = \nu \right\} = \vartheta I_\nu \,. \tag{2.4}$$

Hence, a straightforward argument (see lemma II.1 of [3]) proves that (2.3) is equivalent to I < 0.

In order to prove Theorem 1.1 we need to analyse the possible behaviour of minimizing sequences for I. This is done in the following lemma.

**Lemma 2.4.** Let  $\{\varphi_n\}_n$  be a X-bounded sequence such that  $\int_{\mathbb{R}^3} |\varphi_n|^2 dx = 1$  for all  $n \ge 0$ . Then there exists a subsequence that we still denote by  $\{\varphi_n\}_n$  such that one of the following properties holds:

(1) Compactness up to a translation: there exists a sequence  $\{y_n\}_n \subset \mathbb{R}^3$  such that, for every  $\varepsilon > 0$ , there exists  $0 < R < \infty$  with

$$\int_{B(y_n,R)} |\varphi_n|^2 \, dx \ge 1 - \varepsilon;$$

(2) Vanishing: for all  $0 < R < \infty$ 

$$\sup_{y \in \mathbb{R}^3} \int_{B(y,R)} |\varphi_n|^2 \, dx \xrightarrow[n]{} 0;$$

(3) Dichotomy: there exist  $\alpha \in (0,1)$  and  $n_0 \geq 0$  such that there exist two X-bounded sequences,  $\{\varphi_1^n\}_{n\geq n_0}$  and  $\{\varphi_2^n\}_{n\geq n_0}$ , satisfying the following properties:

$$\|\varphi_n - (\varphi_1^n + \varphi_2^n)\|_{L^p} \xrightarrow[n]{} 0, \text{ for } 2 \le p < 6,$$
 (2.5)

and

$$\int_{\mathbb{R}^3} |\varphi_1^n|^2 \, dx \xrightarrow[n]{} \alpha \ and \ \int_{\mathbb{R}^3} |\varphi_2^n|^2 \, dx \xrightarrow[n]{} 1 - \alpha, \tag{2.6}$$

dist(supp 
$$\varphi_1^n$$
, supp  $\varphi_2^n$ )  $\xrightarrow{n} +\infty$ . (2.7)

Moreover, in this case we have that

$$\liminf_{n \to +\infty} \mathcal{E}(\varphi_n) - \mathcal{E}(\varphi_1^n) - \mathcal{E}(\varphi_2^n) \ge 0, \qquad (2.8)$$

which implies  $I \ge I_{\alpha} + I_{1-\alpha}$ .

Proof of Lemma 2.4. Let  $\{\varphi_n\}_n$  be a X-bounded sequence such that  $\int_{\mathbb{R}^3} |\varphi_n|^2 dx = \nu$  for all  $n \ge 0$ . We remind that X-bounded means that there exists C > 0 such that

$$\|\varphi_n\|_{L^2}^2 + \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_n|^2)_+} \, dx \le C.$$

Moreover, thanks to Lemma 2.1, if  $\{\varphi_n\}_n$  is a X-bounded sequence then  $\{\varphi_n\}_n$  is bounded in  $L^{\infty}$  (by the constant 1) and in  $H^1(\mathbb{R}^3)$ . Then, along the lines of [3], we introduce the so-called Lévy concentration functions

$$Q_n(R) = \sup_{y \in \mathbb{R}^3} \int_{|x-y| < R} |\varphi_n|^2 \, dx,$$
(2.9)

$$K_n(R) = \sup_{y \in \mathbb{R}^3} \int_{|x-y| < R} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1-|\varphi_n|^2)_+} \, dx \tag{2.10}$$

for R > 0. Note that  $Q_n$  and  $K_n$  are continuous non-decreasing functions on  $[0, +\infty)$ , such that for all  $n \ge 0$  and for all R > 0

$$Q_n(R) + K_n(R) \le C$$

since  $\{\varphi_n\}_n$  is X-bounded. Then, up to a subsequence, we have for all R > 0

$$Q_n(R) \xrightarrow[n]{} Q(R), \tag{2.11}$$

$$K_n(R) \xrightarrow[n]{} K(R),$$
 (2.12)

where Q and K are nonnegative, non-decreasing functions. Clearly, we have that

$$\alpha = \lim_{R \to +\infty} Q(R) \in [0, 1],$$

and we denote  $l = \lim_{R \to +\infty} K(R)$ .

If  $\alpha = 0$ , then the situation (2) of the lemma arises as a direct consequence of Definition (2.9). If  $\alpha = 1$ , then (1) follows, see [3] for details. Assume that  $\alpha \in (0, 1)$ , we have to show that (3) holds.

First of all, consider  $\varepsilon > 0$ , small, and  $R_{\varepsilon} > 0$  such that  $Q(R_{\varepsilon}) = \alpha - \varepsilon$  and  $K(R_{\varepsilon}) \leq l - \varepsilon$ . Then, for *n* large enough,

$$Q_n(R_{\varepsilon}) - Q(R_{\varepsilon}) < 1/n, \quad K_n(R_{\varepsilon}) - K(R_{\varepsilon}) < 1/n$$

and by definition of the Lévy functions  $Q_n$ , extracting subsequences if necessary, there exists  $y_n \in \mathbb{R}^3$  such that

$$\left| \int_{|x-y_n| < R_{\varepsilon}} |\varphi_n|^2 dx - Q_n(R_{\varepsilon}) \right| \le \frac{1}{n},$$
$$\left| \int_{|x-y_n| < R_{\varepsilon}} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1-|\varphi_n|^2)_+} dx - K_n(R_{\varepsilon}) \right| \le \frac{1}{n}$$

Next define  $R_n > R_{\varepsilon}$  such that

$$\int_{R_{\varepsilon} < |x-y_n| < R_n} |\varphi_n|^2 \, dx = \frac{3}{n} + \varepsilon \, .$$

Necessarily,  $R_n \to +\infty$  as  $n \to +\infty$ . Indeed, if  $R_n \leq M$  for some M > 0, then  $Q(M) > \alpha$ , which is impossible. We then deduce that for n large enough,

$$\int_{\frac{R_n}{8} \le |x - y_n| \le R_n} |\varphi_n|^2 \, dx \le \frac{3}{n} + \varepsilon$$

Let  $\xi$ ,  $\zeta$  be cut-off functions:  $\xi, \zeta \in \mathcal{D}(\mathbb{R}^3)$  such that

$$\xi(x) = \begin{cases} 1 & |x| \le 1\\ 1 - \exp\left(1 - \frac{1}{1 - \exp\left(1 - \frac{1}{2 - |x|}\right)}\right) & 1 < |x| < 2\\ 0 & |x| \ge 2\\ \end{cases}$$
$$\zeta(x) = \begin{cases} 0 & |x| \le 1\\ \exp\left(1 - \frac{1}{1 - \exp\left(1 - \frac{1}{2 - |x|}\right)}\right) & 1 < |x| < 2\\ 1 & |x| \ge 2 \end{cases}$$

and let  $\xi_{\mu}$ ,  $\zeta_{\mu}$  denote  $\xi\left(\frac{\cdot}{\mu}\right)$ ,  $\zeta\left(\frac{\cdot}{\mu}\right)$ . We define

$$\varphi_1^n(\cdot) = \xi_{\frac{R_n}{8}}(\cdot - y_n)\varphi_n(\cdot) = \xi_{\frac{R_n}{8},y_n}(\cdot)\varphi_n(\cdot)$$
(2.13)

$$\varphi_2^n(\cdot) = \zeta_{\frac{R_n}{2}}(\cdot - y_n)\varphi_n(\cdot) = \zeta_{\frac{R_n}{2},y_n}(\cdot)\varphi_n(\cdot)$$
(2.14)

with  $R_n \to +\infty$ . (2.7) follows easily from these definitions. Furthermore, (2.5) and (2.6) are obtained in the following way:

$$\lim_{n \to +\infty} \int_{\mathbb{R}^3} |\varphi_n - (\varphi_1^n + \varphi_2^n)|^2 dx = \lim_{n \to +\infty} \int_{\frac{R_n}{8} \le |x - y_n| \le R_n} |(1 - \xi_{\frac{R_n}{8}} - \zeta_{\frac{R_n}{2}})\varphi_n|^2 dx$$
$$\leq \lim_{n \to +\infty} \int_{\frac{R_n}{8} \le |x - y_n| \le R_n} |\varphi_n|^2 dx \le \varepsilon,$$

Now by taking a sequence of  $\varepsilon$  tending to 0, and by taking a diagonal sequence of the functions  $\varphi_n$ , and calling it by the same name, we find

$$\int_{\frac{R_n}{8} \le |x - y_n| \le R_n} |\varphi_n|^2 \, dx \xrightarrow[n]{} 0,$$

and, since  $\{\varphi_1^n\}_n$  and  $\{\varphi_2^n\}_n$  are bounded in  $H^1(\mathbb{R}^3)$ , we also obtain

$$\lim_{n \to +\infty} \|\varphi_n - (\varphi_1^n + \varphi_2^n)\|_{L^p} \xrightarrow[n]{} 0,$$

for  $2 \leq p < 6$ . Next, we have to prove that  $\{\varphi_1^n\}_{n \geq n_0}$  and  $\{\varphi_2^n\}_{n \geq n_0}$  are X-bounded. To this purpose, we show that

$$\lim_{n \to +\infty} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_1^n|^2}{(1 - |\varphi_1^n|^2)_+} \, dx - \int_{\mathbb{R}^3} \frac{\xi_{\frac{R_n}{8}, y_n}^2 |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_1^n|^2)_+} \, dx = 0 \tag{2.15}$$

and

$$\lim_{n \to +\infty} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_2^n|^2}{(1 - |\varphi_2^n|^2)_+} \, dx - \int_{\mathbb{R}^3} \frac{\zeta_{\frac{R_n}{2}, y_n}^2 |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_2^n|^2)_+} \, dx = 0 \tag{2.16}$$

Indeed, if (2.15) and (2.16) hold, we obtain that for all  $\varepsilon > 0$ , there exists  $n_0 \ge 0$  such that for all  $n \ge n_0$ , we have

$$\int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_1^n|^2}{(1 - |\varphi_1^n|^2)_+} \, dx \le \int_{\mathbb{R}^3} \frac{\xi_{\frac{R_n}{8}, y_n}^2 |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_1^n|^2)_+} \, dx + o(1)_{n \to +\infty} \\ \le \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_n|^2)_+} \, dx + o(1)_{n \to +\infty} \le C + o(1)_{n \to +\infty},$$

and

$$\int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_2^n|^2}{(1 - |\varphi_2^n|^2)_+} \, dx \le \int_{\mathbb{R}^3} \frac{\zeta_{\frac{R_n}{2}, y_n}^2 |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_2^n|^2)_+} \, dx + o(1)_{n \to +\infty} \\ \le \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_n|^2)_+} \, dx + o(1)_{n \to +\infty} \le C + o(1)_{n \to +\infty}.$$

To prove (2.15) we proceed as follows. We remark that

$$\int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_1^n|^2}{(1 - |\varphi_1^n|^2)_+} \, dx - \int_{\mathbb{R}^3} \frac{\xi_{\frac{R_n}{8}, y_n}^2 |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_1^n|^2)_+} \, dx = A_n + B_n$$

where

$$A_{n} := \int_{\mathbb{R}^{3}} \frac{|\boldsymbol{\sigma} \cdot (\nabla \xi_{\frac{R_{n}}{8}, y_{n}})\varphi_{n}|^{2}}{(1 - |\varphi_{1}^{n}|^{2})_{+}} dx = \int_{\frac{R_{n}}{8} \le |x - y_{n}| \le \frac{R_{n}}{4}} \frac{|\boldsymbol{\sigma} \cdot (\nabla \xi_{\frac{R_{n}}{8}, y_{n}})\varphi_{n}|^{2}}{(1 - |\varphi_{1}^{n}|^{2})_{+}} dx$$
$$\leq \int_{\frac{R_{n}}{8} \le |x - y_{n}| \le \frac{R_{n}}{4}} \frac{|\boldsymbol{\sigma} \cdot (\nabla \xi_{\frac{R_{n}}{8}, y_{n}})\varphi_{n}|^{2}}{1 - \xi_{\frac{R_{n}}{8}, y_{n}}^{2}} dx := C_{n}$$

and

$$|B_n| \le 2 (C_n)^{\frac{1}{2}} \left( \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_n|^2)_+} dx \right)^{\frac{1}{2}}$$

Let us now prove that  $C_n$  tends to 0 as n goes to  $+\infty$ . Using spherical coordinates, we obtain

$$C_{n} \leq \int_{\frac{R_{n}}{8}}^{\frac{R_{n}}{4}} \int_{0}^{\pi} \int_{0}^{2\pi} \frac{|(\sigma \cdot \boldsymbol{e}_{r}) \varphi_{n}(s, \theta, \phi)|^{2} \left(\xi'_{\frac{R_{n}}{8}}(s)\right)^{2}}{1 - \xi_{\frac{R_{n}}{8}}^{2}(s)} s^{2} \sin \theta \, ds \, d\theta \, d\phi$$

$$\leq \frac{64}{R_{n}^{2}} \int_{\frac{R_{n}}{8}}^{\frac{R_{n}}{4}} \int_{0}^{\pi} \int_{0}^{2\pi} \frac{|\varphi_{n}(s, \theta, \phi)|^{2} \left(\xi'\left(\frac{8}{R_{n}}s\right)\right)^{2}}{1 - \xi^{2}\left(\frac{8}{R_{n}}s\right)} s^{2} \sin \theta \, ds \, d\theta \, d\phi$$

$$\leq \frac{64}{R_{n}^{2}} \max_{1 \leq r \leq 2} \frac{(\xi'(r))^{2}}{1 - \xi^{2}(r)} \int_{0}^{+\infty} \int_{0}^{\pi} \int_{0}^{2\pi} |\varphi_{n}(s, \theta, \phi)|^{2} s^{2} \sin \theta \, ds \, d\theta \, d\phi = O\left(\frac{1}{R_{n}^{2}}\right)$$

$$(r'(\sigma))^{2}$$

since  $\max_{1 \le r \le 2} \frac{(\xi'(r))^2}{1-\xi^2(r)} \le C$ . Indeed, since  $\xi^2(r) = 1$  if and only if r = 1,  $\frac{(\xi'(r))^2}{1-\xi^2(r)}$  is a continuous function on (1,2). Moreover, by a straightforward calculation, we obtain  $\lim_{r \to 1^+} \frac{(\xi'(r))^2}{1-\xi^2(r)} = 0 = \lim_{r \to 2^-} \frac{(\xi'(r))^2}{1-\xi^2(r)}$ . Hence, we can conclude, that  $\frac{(\xi'(r))^2}{1-\xi^2(r)}$  is bounded in [1,2]. As a conclusion, since  $R_n \to +\infty$ , we obtain

$$\lim_{n \to +\infty} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_1^n|^2}{(1 - |\varphi_1^n|^2)_+} \, dx - \int_{\mathbb{R}^3} \frac{\xi_{\frac{R_n}{8}, y_n}^2 |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_1^n|^2)_+} \, dx = 0$$

With the same argument, we prove (2.16).

Finally, it remains to show that

$$\liminf_{n \to +\infty} \mathcal{E}(\varphi_n) - \mathcal{E}(\varphi_1^n) - \mathcal{E}(\varphi_2^n) \ge 0.$$

First of all, using the definitions (2.13) and (2.14), we obtain

$$\lim_{n \to +\infty} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_n|^2)_+} \, dx \ge \lim_{n \to +\infty} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_1^n|^2 - |\varphi_2^n|^2)_+} \, dx$$

Next, we remark that

$$\begin{split} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1-|\varphi_1^n|^2-|\varphi_2^n|^2)_+} \, dx &- \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_1^n|^2}{(1-|\varphi_1^n|^2)_+} \, dx - \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_2^n|^2}{(1-|\varphi_2^n|^2)_+} \, dx \\ &= \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1-|\varphi_1^n|^2-|\varphi_2^n|^2)_+} \, dx - \int_{\mathbb{R}^3} \frac{\xi_{\frac{R_n}{8},y_n}^2 |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1-|\varphi_1^n|^2)_+} \, dx \\ &- \int_{\mathbb{R}^3} \frac{\xi_{\frac{R_n}{2},y_n}^2 |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1-|\varphi_2^n|^2)_+} \, dx + o(1)_{n \to \infty} \\ &= \int_{\mathbb{R}^3} \frac{\left(1-\xi_{\frac{R_n}{8},y_n}^2 - \xi_{\frac{R_n}{2},y_n}^2\right) |\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1-|\varphi_1^n|^2-|\varphi_2^n|^2)_+} \, dx + o(1)_{n \to \infty} \\ &\geq o(1)_{n \to \infty}. \end{split}$$

As a conclusion,

$$\lim_{n \to +\infty} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_n|^2}{(1 - |\varphi_n|^2)_+} \, dx \ge \lim_{n \to +\infty} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_1^n|^2}{(1 - |\varphi_1^n|^2)_+} \, dx$$
$$+ \lim_{n \to +\infty} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi_2^n|^2}{(1 - |\varphi_2^n|^2)_+} \, dx$$

and, using (2.5) and the localization properties of  $\varphi_1^n$  and  $\varphi_2^n$ , we have

$$I = \lim_{n \to +\infty} \mathcal{E}(\varphi_n) \ge \liminf_{n \to +\infty} \mathcal{E}(\varphi_1^n) + \liminf_{n \to +\infty} \mathcal{E}(\varphi_2^n) \ge I_{\alpha} + I_{1-\alpha}.$$

Proof of Theorem 1.1. Assume that I < 0. By Lemma 2.2, any minimizing sequence  $\{\varphi_n\}_n$  is X-bounded, and then we can use Lemma 2.4 to it. It is easy to rule out vanishing and dichotomy whenever I < 0.

Vanishing cannot occur. Indeed, If vanishing occurs, then, up to a subsequence,  $\forall R < +\infty$  we have

$$\lim_{n \to +\infty} \sup_{y \in \mathbb{R}^3} \int_{B(y,R)} |\varphi_n|^2 = 0.$$
(2.17)

This implies that  $\varphi_n$  converges strongly in  $L^p(\mathbb{R}^3)$  for  $2 and, as a consequence, <math>I \ge 0$ . Clearly, this contradicts I < 0. Moreover, if disheterny energy we have

Moreover, if dichotomy occurs, we have

$$I = \lim_{n \to +\infty} \mathcal{E}(\varphi_n) \ge \liminf_{n \to +\infty} \mathcal{E}(\varphi_1^n) + \liminf_{n \to +\infty} \mathcal{E}(\varphi_2^n) \ge I_{\alpha} + I_{1-\alpha}$$

which contradicts Lemma 2.3, since I < 0.

Hence, for n large enough, there exists  $\{y_n\}_n \in \mathbb{R}^3$  such that  $\forall \varepsilon > 0, \exists R < +\infty$ ,

$$\int_{B(y_n,R)} |\varphi_n|^2 \ge 1 - \varepsilon.$$

We denote by  $\tilde{\varphi}_n(\cdot) = \varphi_n(\cdot+y_n)$ . Since  $\{\tilde{\varphi}_n\}_n$  is bounded in  $H^1$ ,  $\{\tilde{\varphi}_n\}_n$  converges weakly in  $H^1$ , almost everywhere on  $\mathbb{R}^3$  and in  $L^p_{loc}$  for  $2 \leq p < 6$  to some  $\tilde{\varphi}$ . In particular, as a consequence of weak convergence in  $H^1$ ,  $\boldsymbol{\sigma} \cdot \nabla \tilde{\varphi}_n$  converges weakly to  $\boldsymbol{\sigma} \cdot \nabla \tilde{\varphi}$  in  $L^2$ . Moreover, thanks to the concentration-compactness argument,  $\{\tilde{\varphi}_n\}_n$  converges strongly in  $L^2$  and in  $L^p$  for  $2 \leq p < 6$ .

**Lemma 2.5.** Let  $\{f_n\}_n$  and  $\{g_n\}_n$  be two sequences of functions such that  $f_n : \mathbb{R}^3 \to \mathbb{R}_+, g_n : \mathbb{R}^3 \to \mathbb{C}^2$ ,  $f_n$  converges to f a.e.,  $g_n$  converges weakly to g in  $L^2$  and there exists a constant C, that does not depend on n, such that  $\int_{\mathbb{R}^3} f_n |g_n|^2 dx \leq C$ . Then

$$\int_{\mathbb{R}^3} f|g|^2 \, dx \le \liminf_{n \to +\infty} \int_{\mathbb{R}^3} f_n |g_n|^2 \, dx$$

*Proof.* Given a function  $h : \mathbb{R}^3 \to \mathbb{R}_+$ , let  $T_k$  be the function defined by

$$T_k(h)(x) = \begin{cases} h(x) & \text{if } h(x) \le k \\ k & \text{if } h(x) > k \end{cases}$$

for all  $k \in [0, \infty)$ . Hence, the following properties are satisfied for all  $k \in [0, \infty)$ :

$$T_k(f_n) \xrightarrow[n]{} T_k(f)$$
 a.e. in  $\mathbb{R}^3$ , (2.18)

$$T_k(f_n)|g|^2 \xrightarrow[n]{} T_k(f)|g|^2 \quad \text{in } L^1,$$

$$(2.19)$$

$$T_k(f_n)g \xrightarrow{}_n T_k(f)g \quad \text{in } L^2,$$
 (2.20)

$$||T_k(f_n)g||_{L^2} \to ||T_k(f)g||_{L^2},$$
 (2.21)

where to obtain (2.19) and (2.21), we use Lebesgue's dominated convergence theorem. Moreover, as a consequence of (2.20) and (2.21), we have

$$T_k(f_n)g \xrightarrow[n]{} T_k(f)g \quad \text{in } L^2.$$
 (2.22)

Next, we have

$$0 \leq \liminf_{n \to +\infty} \int_{\mathbb{R}^3} T_k(f_n) |g_n - g|^2 \, dx = \liminf_{n \to +\infty} \int_{\mathbb{R}^3} T_k(f_n) |g_n|^2 \, dx$$
$$+ \liminf_{n \to +\infty} \int_{\mathbb{R}^3} T_k(f_n) |g|^2 \, dx - \liminf_{n \to +\infty} \left( \int_{\mathbb{R}^3} T_k(f_n) \, \overline{g}_n \cdot g \, dx + \int_{\mathbb{R}^3} T_k(f_n) \, g_n \cdot \overline{g} \, dx \right)$$
$$= \liminf_{n \to +\infty} \int_{\mathbb{R}^3} T_k(f_n) |g_n|^2 \, dx + \int_{\mathbb{R}^3} T_k(f) |g|^2 \, dx - 2 \int_{\mathbb{R}^3} T_k(f) |g|^2 \, dx$$

thanks to (2.19), (2.22) and the fact that  $g_n$  converges weakly to g in  $L^2$ . As a consequence,

$$\int_{\mathbb{R}^3} T_k(f) |g|^2 \, dx \le \liminf_{n \to +\infty} \int_{\mathbb{R}^3} T_k(f_n) |g_n|^2 \, dx \tag{2.23}$$

Since

$$\liminf_{n \to +\infty} \int_{\mathbb{R}^3} T_k(f_n) |g_n|^2 \, dx \le \liminf_{n \to +\infty} \int_{\mathbb{R}^3} f_n |g_n|^2 \, dx \le C,$$

we can pass to the limit for k that goes to  $+\infty$  in (2.23) and we obtain

$$\int_{\mathbb{R}^3} f|g|^2 \, dx \le \liminf_{n \to +\infty} \int_{\mathbb{R}^3} f_n |g_n|^2 \, dx.$$

By applying Lemma (2.5) to  $f_n = \frac{1}{(1-|\tilde{\varphi}_n|^2)}_+$  and  $g_n = |\boldsymbol{\sigma} \cdot \nabla \tilde{\varphi}_n|$ , we obtain

$$\int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \tilde{\varphi}|^2}{(1-|\tilde{\varphi}|^2)_+} \, dx \le \liminf_{n \to +\infty} \int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \tilde{\varphi}_n|^2}{(1-|\tilde{\varphi}_n|^2)_+} \, dx$$

Hence,  $\tilde{\varphi} \in X$ ,  $\int_{\mathbb{R}^3} |\tilde{\varphi}|^2 dx = 1$ , and

$$\mathcal{C}(\tilde{\varphi}) \leq \liminf_{n \to +\infty} \mathcal{E}(\tilde{\varphi}_n) \leq \mathcal{E}(\tilde{\varphi}).$$

As a conclusion, the minimum of I is achieved by  $\tilde{\varphi}$ .

Finally, it remains to prove that there exists  $a_0 > 0$  such that for all  $a > a_0$  we have I < 0.

It is clear that I < 0 for a large enough. Since I is non-increasing with respect to a, we may denote by  $a_0$  the least positive constant such that I < 0 for  $a > a_0$ . We have to prove that  $a_0 > 0$  or in other words I = 0 for a small enough. Using Sobolev and Hölder inequalities, we find, for  $\varphi \in X$  such that  $\int_{\mathbb{R}^3} |\varphi|^2 dx = 1$ ,

$$\mathcal{E}(\varphi) \ge \frac{1}{S^2} \left( \int_{\mathbb{R}^3} |\varphi|^6 \, dx \right)^{1/3} - \frac{a}{2} \left( \int_{\mathbb{R}^3} |\varphi|^6 \, dx \right)^{1/3}$$

Hence, if  $a \leq \frac{2}{S^2}$ , I = 0. This implies  $a_0 > \frac{2}{S^2}$ . According to [7] the best constant for the Sobolev inequality

$$\|u\|_{L^q(\mathbb{R}^m)} \le C \|\nabla u\|_{L^p(\mathbb{R}^m)}$$

with  $1 and <math>q = \frac{mp}{(m-p)}$  is given by

$$C = \pi^{-1/2} m^{-1/p} \left(\frac{p-1}{m-p}\right)^{1-1/p} \left(\frac{\Gamma(1+m/2)\Gamma(m)}{\Gamma(m/p)\Gamma(1+m-m/p)}\right)^{1/m}.$$

In particular,

$$S = \frac{1}{\sqrt{3\pi}} \left(\frac{4}{\sqrt{\pi}}\right)^{1/3},$$

and

$$\frac{2}{S^2} = \frac{3\pi^{4/3}}{2^{1/3}} \approx 10.96.$$

To obtain an upper estimate for  $a_0$ , we consider the following test function

$$\bar{\varphi}(x) = \left(\begin{array}{c} \bar{f}_R(|x|) \\ 0 \end{array}\right)$$

where  $\bar{f}_R(|x|) = \bar{f}\left(\frac{|x|}{R}\right)$ ,

$$\bar{f}(|x|) = \begin{cases} \cos(|x|) & |x| \le \frac{\pi}{2} \\ 0 & |x| > \frac{\pi}{2} \end{cases}$$

and  $R \in (0,1)$  is such that  $\int |\bar{f}_R|^2 dx = 1$ . This implies

$$R = \left(\frac{2}{\pi}\right)^{2/3} \left(\frac{3}{\pi^2 - 6}\right)^{1/3}$$

Next, we denote by  $\bar{a}$  the positive constant such that  $\mathcal{E}(\varphi) = 0$ . By definition,

$$\bar{a} = \frac{2\int_{\mathbb{R}^3} \frac{|\boldsymbol{\sigma} \cdot \nabla \bar{\varphi}|^2}{(1-|\bar{\varphi}|^2)_+} \, dx}{\int |\bar{\varphi}|^4} = \frac{2\int_{\mathbb{R}^3} \frac{|\nabla \bar{f}_R|^2}{1-|\bar{f}_R|^2} \, dx}{\int |\bar{f}_R|^4} \,,$$

and, by a straightforward calculation, we obtain

$$\int_{\mathbb{R}^3} \frac{|\nabla \bar{f}_R|^2}{1 - |\bar{f}_R|^2} \, dx = \frac{\pi^4}{6} R = \frac{\pi^{10/3}}{3^{2/3} (2(\pi^2 - 6))^{1/3}},$$
$$\int |\bar{f}_R|^4 = \frac{\pi^2 (2\pi^2 - 15)}{32} R^3 = \frac{3(2\pi^2 - 15)}{8(\pi^2 - 6)}.$$

As a consequence,

$$\bar{a} = \frac{8\pi^{10/3} \left(\frac{2}{3}(\pi^2 - 6)\right)^{2/3}}{3(2\pi^2 - 15)} \approx 48.06$$

Since the energy functional  $\mathcal{E}$  is decreasing in a, if  $a > \bar{a}$  then  $I \leq \mathcal{E}(\bar{\varphi}) < 0$ . As a conclusion,  $a_0 \leq \bar{a} + \varepsilon$  for all  $\varepsilon > 0$ .

#### Appendix A

**A.1.** We begin this section by proving that if  $(\varphi, \chi)$  a solution of (1.7) with  $\varphi \in H^1(\mathbb{R}^3)$  of the form (1.12), then  $|\varphi|^2 \leq 1$  a.e. in  $\mathbb{R}^3$ . As we saw before,  $\varphi$  is a solution of

$$-\boldsymbol{\sigma} \cdot \nabla \left(\frac{\boldsymbol{\sigma} \cdot \nabla \varphi}{1 - |\varphi|^2}\right) + \frac{|\boldsymbol{\sigma} \cdot \nabla \varphi|^2}{(1 - |\varphi|^2)^2} \varphi - a|\varphi|^2 \varphi + b\varphi = 0, \tag{A.1}$$

or equivalently,

$$\frac{\Delta\varphi}{|\varphi|^2-1} - \frac{|\boldsymbol{\sigma}\cdot\nabla\varphi|^2}{(|\varphi|^2-1)^2}\varphi - a|\varphi|^2\varphi + b\varphi = 0\,,$$

or still,

$$\Delta \varphi - \frac{|\nabla \varphi|^2}{|\varphi|^2 - 1} \varphi - (a|\varphi|^2 \varphi - b\varphi)(|\varphi|^2 - 1) = 0,$$

because for functions  $\varphi$  of the form (1.12),

$$|\sigma \cdot \nabla \varphi|^2 = |\nabla \varphi|^2$$
 and  $\sigma \cdot (\nabla \varphi \wedge \nabla \varphi) = 0$  a.e.

For any K > 1, we define the truncation function  $T_K(s)$  by  $T_K(s) = s$  if 1 < s < K, and  $T_K(s) = 0$  otherwise. Multiplying the above equation by  $\varphi T_K(|\varphi|^2) \in L^2(\mathbb{R}^3)$ , we obtain

$$-\int_{\mathbb{R}^{3}} |\nabla \varphi|^{2} T_{K}(|\varphi|^{2}) - \int_{\mathbb{R}^{3}} (\nabla \varphi \cdot \varphi) \nabla T_{K}(|\varphi|^{2}) - \int_{\mathbb{R}^{3}} \frac{|\nabla \varphi|^{2}}{|\varphi|^{2} - 1} |\varphi|^{2} T_{K}(|\varphi|^{2}) - \int_{\mathbb{R}^{3}} (a|\varphi|^{2} - b)(|\varphi|^{2} - 1)|\varphi|^{2} T_{K}(|\varphi|^{2}) = 0.$$
(A.2)

Moreover, for all K > 1,

$$\nabla T_K(|\varphi|^2) = \begin{cases} 2\varphi \cdot \nabla \varphi & 1 < |\varphi|^2 < K\\ 0 & |\varphi|^2 \le 1 \text{ or } |\varphi|^2 \ge K \end{cases}$$

Therefore, if a - b > 0 the l.h.s of (A.2) is negative and this implies that either  $|\varphi|^2 \leq 1$  or  $|\varphi|^2 \geq K$  a.e. As a conclusion, taking the limit  $K \to +\infty$ , if a - b > 0 then any solution  $\varphi$  of (A.1) of the form (1.12) satisfies  $|\varphi|^2 \leq 1$  a.e. in  $\mathbb{R}^3$ , and in the equation (A.1) we can replace the term  $(1 - |\varphi|^2)$  by  $(1 - |\varphi|^2)_+$  without changing its solution set. The same happens for solutions of the form (1.8).

**A.2.** Let us next prove that the functional  $\mathcal{F}$ , defined by (1.11), is not bounded from below. Consider the function  $\xi$  introduced in the proof of Lemma 2.4. Let us denote  $A := \int_{\mathbb{R}^3} |\xi(x)|^2 dx$ .

Then, let us define the radially symmetric function

$$f(r) = \begin{cases} e^{(r-\sqrt{\ln 2})^2}, & 0 \le r < \sqrt{\ln 2}, \\ \bar{\xi}(r+1-\sqrt{\ln 2}), & r \ge \sqrt{\ln 2}, \end{cases}$$

where  $\overline{\xi}(|x|) = \xi(x)$  for all x, and take  $a := \int_{\mathbb{R}^3} f(|x|)^2 dx$ . Note that  $\operatorname{supp}(f) \subset [0, 1 + \sqrt{\ln 2}]$  and  $\max_{0 \le r \le 1 + \sqrt{\ln 2}} \frac{(f'(r))^2}{1 - f^2(r)} \le C$ , for some constant C > 0.

Next, for all integers n > 0, define the rescaled functions  $\xi_n(x) := n^{3/2}\xi(nx)$ . This change of variables leaves invariant the  $L^2(\mathbb{R}^3)$  norm. Then for n large, consider the function

$$g_n(x) := \max_{\mathbb{T}^3} \{\xi_n, f\}$$

Note that the measure of the set  $\{x \in \mathbb{R}^3 : g_n = \xi_n\}$  tends to 0 as n goes to  $+\infty$ . This function satisfies  $\int_{\mathbb{R}^3} |g_n(x)|^2 dx = A + a + o(1)$ , as n goes to  $+\infty$ . In order to normalize it in the  $L^2$  norm, let us finally define the rescaled function  $g_n^R(x) := g_n\left(\frac{x}{R}\right), R > 0$  and choose  $R_n$  such that  $\int_{\mathbb{R}^3} |g_n^{R_n}(x)|^2 dx = 1$ . As n goes to  $+\infty, R_n \to \overline{R} := (A + a)^{-1/3} > 0$ . We compute now the energy  $\mathcal{F}$  of the vector function  $\varphi_n^{R_n}$  defined by

$$\varphi_n^{R_n}(x) = \left(\begin{array}{c} g_n^{R_n}(x) \\ 0 \end{array}\right) \,.$$

We find

$$\mathcal{F}(\varphi_n^{R_n}) = \int_{\xi_n^{R_n} \ge f^{R_n}} \frac{\left(\left(\xi_n^{R_n}\right)'(r)\right)^2}{1 - \left(\xi_n^{R_n}(r)\right)^2} \, dx - \frac{a \, n^3 \, R_n^3}{2} \int_{\xi_n^{R_n} \ge f^{R_n}} |\xi|^4 \, dx \\ + R_n \, \int_{\xi_n^{R_n} \le f^{R_n}} \frac{\left(f'\left(r\right)\right)^2}{1 - f^2\left(r\right)} \, dx - \frac{a \, R_n^3}{2} \, \int_{\xi_n^{R_n} \le f^{R_n}} f(x)^4 \, dx \\ \le -\frac{a \, n^3 \, R_n^3}{2} \int_{\mathbb{R}^3} |\xi|^4 \, dx + R_n \, \int_{\mathbb{R}^3} \frac{\left(f'\left(r\right)\right)^2}{1 - f^2\left(r\right)} \, dx - \frac{a \, R_n^3}{2} \, \int_{\mathbb{R}^3} f(x)^4 \, dx + o(n^3) \, ,$$

because whenever  $\xi_n^{R_n} \ge f^{R_n}$ ,  $(\xi_n^{R_n})^2 > 1$  and because the sequence  $\{R_n\}_n$  is bounded. This clearly shows that  $\mathcal{F}$  is unbounded from below.

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