



Traveling fronts guided by the environment for reaction-diffusion equations

Henri Berestycki, Guillemette Chapuisat

► To cite this version:

Henri Berestycki, Guillemette Chapuisat. Traveling fronts guided by the environment for reaction-diffusion equations. *Networks and Heterogeneous Media*, 2013, 8 (1), pp.79-114. <hal-00712553v2>

HAL Id: hal-00712553

<https://hal.science/hal-00712553v2>

Submitted on 9 Mar 2016

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



HAL Authorization

Traveling fronts guided by the environment for reaction-diffusion equations

HENRI BERESTYCKI

CAMS, UMR 8557, EHESS
190-198 avenue de France, 75244 Paris Cedex 13, France

GUILLEMETTE CHAPUISAT

LATP, UMR 7353, Aix-Marseille Université
39 rue F. Joliot-Curie, 13453 Marseille Cedex 13, France
CAMS, UMR 8557, EHESS
190-198 avenue de France, 75244 Paris Cedex 13, France

Dedicated to Hiroshi Matano on the occasion of his *Kanreki*.

Abstract

This paper deals with the existence of traveling fronts for the reaction-diffusion equation:

$$\frac{\partial u}{\partial t} - \Delta u = h(u, y) \quad t \in \mathbb{R}, x = (x_1, y) \in \mathbb{R}^N.$$

We first consider the case $h(u, y) = f(u) - \alpha g(y)u$ where f is of KPP or bistable type and $\lim_{|y| \rightarrow +\infty} g(y) = +\infty$. This equation comes from a model in population dynamics in which there is spatial spreading as well as phenotypic mutation of a quantitative phenotypic trait that has a locally preferred value. The goal is to understand spreading and invasions in this heterogeneous context. We prove the existence of threshold value α_0 and of a nonzero asymptotic profile (a stationary limiting solution) $V(y)$ if and only if $\alpha < \alpha_0$. When this condition is met, we prove the existence of a traveling front. This allows us to completely identify the behavior of the solution of the parabolic problem in the KPP case.

We also study here the case where $h(y, u) = f(u)$ for $|y| \leq L_1$ and $h(y, u) \approx -\alpha u$ for $|y| > L_2 \geq L_1$. This equation provides a general framework for a model of cortical spreading depressions in the brain. We prove the existence of traveling front if L_1 is large enough and the non-existence if L_2 is too small.

1 Introduction

This paper deals with the existence of bounded traveling fronts for the reaction-diffusion equation

$$\frac{\partial u}{\partial t} - \Delta u = h(y, u) \quad t \in \mathbb{R}, x = (x_1, y) \in \mathbb{R}^N. \quad (1)$$

The function h will be of three different forms in this paper. The first two concern non-linear terms $h(y, u) = f(u) - \alpha g(y)u$ where $f : \mathbb{R} \rightarrow \mathbb{R}$ is C^1 , and is either of positive type, or of bistable type and $g : \mathbb{R}^{N-1} \rightarrow \mathbb{R}_+$ is C^0 , $g(0) = 0$ and $g \xrightarrow{|y| \rightarrow +\infty} +\infty$. The existence of traveling front depends on the value of $\alpha > 0$. The third case we consider here is when $h(y, u) = f(u)$ for $|y| \leq L_1$ and $h(y, u) \leq -mu$ for $|y| \geq L_2$ where $0 < L_1 \leq L_2 < \infty$ are given parameters and f is of bistable form and $h(y, u) + mu \rightarrow 0$ for $|y| \rightarrow +\infty$. We study the existence of traveling fronts depending on the value of L_1 and L_2 .

The problems we study in this paper bear some similarities with the question of traveling fronts in cylinders of [14]. However there are important differences that have to do with the fact that the cross section in [14] was bounded and only the Neumann condition was considered there. Whereas here, the problem is posed in the whole space and the solution vanish at infinity in directions orthogonal to the direction of propagation. We follow the same general scheme as in [14] and in particular make use of the sliding method. But some new ideas are also required. In particular, first, we treat directly the KPP case without the approximation of the KPP non-linearity by a combustion non-linearity as in [14]. Then in the approach of Berestycki - Nirenberg [14] to traveling fronts in cylinders for the bistable case, a useful result of H. Matano [33] was involved in the proof. Here, we rely on stability ideas but also use energy minimization properties to bound the speed of the solution in the finite domain approximation. In particular, we do not use the precise exponential behavior that was used in [14]. Actually the developments of this method that we present in this paper can be used to somewhat simplify parts of [14]. They can also be applied to traveling fronts in cylinder with Robin or Dirichlet boundary conditions.¹

Equation (1) in the first case comes from a model in population dynamics [22] that we briefly describe now. Let $u(t, x, v)$ represent the density of individuals at time t and position x that possess some given quantitative phenotypic trait represented by a continuous variable $v \in \mathbb{R}$. For example, the latter could be the size of wings or the height of an individual. We assume that individuals follow a brownian motion (i.e. they diffuse) in space with a constant diffusion coefficient ν , reproduce identically and disappear with a growth rate $k(x, v)$ that depends on the position x and on the trait v . Furthermore, they also reproduce with mutation that is represented by a kernel $K(x, v, w)$ and disappear due to competition with a constant $L > 0$. Thus, one is led to the following equation for u :

$$\begin{aligned} \partial_t u(t, x, v) - \nu \Delta_x u(t, x, v) &= k(x, v) u(t, x, v) \\ &+ \int_w K(x, v, w) u(t, x, w) dw - u(t, x, v) \int_w L u(t, x, w) dw. \end{aligned} \quad (2)$$

We assume moreover that there exists a most adapted trait $\phi = \phi(x)$ that may depends on the location x . The farther the trait of an individual is from the most adapted trait, the larger the probability of dying and not reproducing. Thus the growth rate can be written for example as $k(x, v) = a - b|v - \phi(x)|^2$ with a and $b > 0$.

¹The construction of traveling fronts for Neumann and Dirichlet conditions in cylinders given by [48] appears to be incomplete. Indeed, the continuity of the function ϕ on page 515 is not established so that using Dini's Theorem to derive Lemma 3.2 there is not justified.

Non-local reaction-diffusion equations of this type raise some new difficulties from a mathematical standpoint as shown in [13]. There, behaviors that are quite different from those in local equations are brought to light. After this paper was completed, we learned that in [1] the existence of traveling front was also derived for equation

$$\partial_t u(t, x, v) - \nu \Delta_{x,v} u(t, x, v) = k(x, v) u(t, x, v) - u(t, x, v) \int_w L(v, w) u(t, x, w) dw.$$

This work follows in part the methods of [13]. As in [13], the nonzero limiting stationary state is not prescribed. In a forthcoming numerical study [6], we study the full equation (2) and we discuss the monotonicity of fronts depending on the value of a and b .

In this paper, we introduce a simplified version of this model that emphasizes propagation guided by the environment. First, we assume that mutations are due to a diffusion process represented by a Brownian motion in the space of trait v . Furthermore, we assume that ϕ is linear. Then a rotation in the variables (x_1, y) allows one to reduce the problem to the case where the most adapted trait is $y = 0$. Therefore we assume $\phi(x) = 0$ and (2) can be rewritten as

$$\partial_t u(t, x, v) - \nu \Delta_{x,v} u(t, x, v) = (a - b|v|^2) u(t, x, v) - u(t, x, v) \int_w L u(t, x, w) dw. \quad (3)$$

Lastly we assume that competition is only between individuals sharing the same trait which leads us to equation

$$\partial_t u - \nu \Delta_{x,v} u = (a - Lu)u - b|v|^2 u. \quad (4)$$

Equation (1) is a generalization of this equation. In [22], the authors observe numerically a generalized transition front spreading along the graph of ϕ for equation (2) (see [9, 10, 11, 34] or [44] for the definition of generalized transition fronts). Here we want to prove theoretically (i) that there exists such a front for equation (1) at least for some values of the parameter $\alpha > 0$ and (ii) that extinction occurs if α is too large. The latter condition can be interpreted as saying that the “area” of adapted traits is too thin compared to the diffusion. To remain consistent with the biological motivation, we only consider here non-negative and bounded solutions of (1).

Other types of models related to this one have been proposed in the literature. For example, the model developed by Kirkpatrick and Barton in 1997 [31] also studies the evolution of a population and of its mean trait. The main difference is that they have a system in u and v where u represents the population and v the mean trait is described by a specific equation. This model has been further explored [25, 30]. It is worth noting that these models use the same type of non-linearity for the adaptation to the environment and model the mutation with the Laplace operator as well rather than integral operators.

This type of reaction-diffusion process in heterogeneous media also arises in many contexts in medicine. An important class of such models was treated in [17, 39]. They deal with the propagation of a *cortical spreading depression* (CSD) in the human brain. These CSD’s are transient depolarizations of the brain that

slowly propagate in the cortex of several animal species after a stroke, a head injury, or seizures [45]. They also are suspected of being responsible for the aura in *migraines with aura*. CSD's are the subject of intensive research in biology since experiments blocking them during strokes in rodents have produced very promising results [21, 36]. These observations however have not been confirmed in humans and the existence of CSD's in the human brain is still a matter of debate [35, 26, 2, 46]. Since very few experiments and measurements on human brain are available be it for obvious ethical or technical reasons, mathematical models of a CSD is helpful in understanding their existence and conditions for their propagation. In such a problem, the morphology of the brain and thus the geometry of the domain where CSD's propagate, is believed to play an important role.

The brain is composed of gray matter where neuron's soma are and of the white matter where only axons are to be found. The rodent brain (on which many of the biological experiments are done) is rather smooth and composed almost entirely of gray matter. On the opposite, the human brain is very tortuous. The gray matter is a thin layer at the periphery of the brain with much thickness variations and convolutions, the rest of the brain being composed of white matter. According to mathematical models of CSDs [19, 45, 43, 47], the depolarization amplitude follows a reaction-diffusion process of bistable type in the gray matter of the brain while it diffuses and is absorbed in the white matter of the brain. The modeling of CSD hence leads one to the study of equations of the following type:

$$\frac{\partial u}{\partial t} - \Delta u = f(u)\mathbf{1}_{|y|<L} - \alpha u\mathbf{1}_{|y|\geq L} \quad t \in \mathbb{R}, x = (x_1, y) \in \mathbb{R}^N. \quad (5)$$

Here, f is of bistable type and $|y| = L$ corresponds to the transition from gray matter to white matter. This equation is of type (1) and we also study it here in sections 7 and 8 where we extend earlier works on the subject. In [17], this equation was studied to prove that the thinness of the human gray matter (L small) may prevent the creation or the propagation of CSDs on large distances. It was proved by studying the energy in a traveling referential of the solution of (5) with a specific initial condition. The special case of (5) for $N = 2$ was described more completely in [20]. In [39], a numerical study shows that the convolutions of the brain have also a strong influence on the propagation of CSD. In [18], the effect of rapid variations of thickness of the gray matter was studied.

Lastly, let us note that the same kind of equation arises in the modeling of tumor cords but with a slightly more complicate KPP non-linearity. We plan to investigate this model in our forthcoming work [7].

As already mentioned, the study of propagation of fronts and spreading properties in heterogeneous media is of intense current interest. For instance, the existence of fronts propagating in non-homogeneous geometries with obstacles has been established in Berestycki, Hamel and Matano [8]. Definitions of generalized waves have been given by Berestycki and Hamel in [10] and [11] where they are called generalized transition waves. Somewhat different approaches to generalizing the notions of traveling fronts have been proposed by H. Matano [34] and W. Shen [44]. The existence of fronts for non-homogeneous equations are established in [37] and [50].

Let us first introduce some notations before stating the main results.

Notation. We note $x = (x_1, y) \in \mathbb{R}^N$ where $x_1 \in \mathbb{R}$ and $y \in \mathbb{R}^{N-1}$. Hence x is the space variable in \mathbb{R}^N , x_1 is its first coordinate and y is the vector of \mathbb{R}^{N-1} composed of all the other coordinates of x . As usual $B_R = B(0, R)$ denotes a ball of radius R centered at 0, but here it will always mean the ball in \mathbb{R}^{N-1} .

First we are interested in solutions of

$$\begin{cases} \frac{\partial u}{\partial t} - \Delta u = f(u) - \alpha g(y)u, & x = (x_1, y) \in \mathbb{R}^N, t \in \mathbb{R} \\ u \geq 0, & u \text{ bounded,} \end{cases} \quad (6)$$

with $\alpha > 0$. We will assume that $f : \mathbb{R} \rightarrow \mathbb{R}$ is C^1 and satisfies either one of the following conditions:

$$f(0) = f(1) = 0, f > 0 \text{ on } (0, 1) \text{ and } f'(0) > 0,$$

or

$$\begin{aligned} &\text{there exists } \theta \in]0, 1[\text{ such that } f(0) = f(\theta) = f(1) = 0, \\ &f < 0 \text{ on } (0, \theta) \text{ and } f > 0 \text{ on } (\theta, 1) \text{ with } f'(0) < 0, f'(1) < 0, \\ &\text{and } \int_0^1 f(s) ds > 0. \end{aligned}$$

The first case will be referred to as the positive case and the second one will be called bistable case. Furthermore, if f is in the positive case and if

$$s \mapsto \frac{f(s)}{s} \text{ is decreasing on } (0, 1]$$

we will say that f is of Fisher-KPP type. Since we are only interested in solutions of (6) in $[0, 1]$, we will further assume that $f(s) \leq 0$ for $s \geq 1$. Moreover we assume

$$g : \mathbb{R}^{N-1} \rightarrow \mathbb{R}_+ \text{ is continuous, } g(0) = 0, g > 0 \text{ on } \mathbb{R}^{N-1} \setminus \{0\} \quad (7)$$

(except in section 3.2 where g can vanish) and

$$\lim_{|y| \rightarrow +\infty} g(y) = +\infty. \quad (8)$$

Taking $g(y) = |y|^2$ and $f(s) = as(1-s)$ yields the particular case of equation (4).

This paper is concerned with the long term behavior of (6) and with the existence of curved traveling fronts, i.e. solutions $u(t, x) = U(x_1 - ct, y)$ with $c \in \mathbb{R}$ a constant and $U : \mathbb{R}^N \rightarrow \mathbb{R}$ such that the limits $\lim_{s \rightarrow \pm\infty} U(s, \cdot)$ exist uniformly and are not equal. Regarding these fronts, our main results are the following.

Theorem 1.1. *If f is of Fisher-KPP type, there exists $\alpha_0 > 0$ such that:*

- *For $\alpha \geq \alpha_0$, there exists no traveling front solution of (6),*
- *For $\alpha < \alpha_0$ there exists a threshold $c_* > 0$ such that there exists a traveling front of speed c of equation (6) if and only if $c \geq c_*$.*

This existence theorem gives us information on the behavior of the solution of the parabolic problem. In this paper we prove the following theorem:

Theorem 1.2. *If f is of Fisher-KPP type, for $u_0 \in L^\infty$, there exists a unique solution $u(t, x)$ of*

$$\begin{cases} \partial_t u - \Delta u = f(u) - \alpha g(y)u & \text{on } (0, +\infty) \times \mathbb{R}^N, \\ u(0, x) = u_0(x) & \text{on } \mathbb{R}^N. \end{cases}$$

- If $\alpha \geq \alpha_0$, it verifies $u(t, x) \xrightarrow{t \rightarrow +\infty} 0$ uniformly with respect to $x \in \mathbb{R}^N$.
- If $\alpha < \alpha_0$ and $u_0 \in C^0(\mathbb{R}^N)$ is compactly supported with $u_0 < V$ where $V = V(y)$ is the unique positive asymptotic profile (stationary solution), then

$$\begin{aligned} & \text{for any } c > c^* \quad \lim_{t \rightarrow +\infty} \sup_{|x_1| \geq ct} u(t, x) = 0, \\ & \text{for any } c \text{ with } 0 \leq c < c^* \quad \lim_{t \rightarrow +\infty} \sup_{|x_1| < ct} |u(t, x) - V(y)| = 0. \end{aligned}$$

This means that there is a threshold value α_0 such that for $\alpha \geq \alpha_0$, there is *extinction*. On the contrary, when $\alpha \leq \alpha_0$, there is spreading and the state $V(y)$ invades the whole space. The asymptotic speed of spreading is then c^* . The property of asymptotic spreading is in the same spirit of the theorem of asymptotic speed of spreading in cylinders established by Mallordy and Roquejoffre in [32].

Theorem 1.2 has interesting consequences for the dynamics of the phenotypic diversity in a population. Several studies have tried to understand population migrations through phenotypic diversity [24, 27, 28, 29, 41, 49]. Our invasion result states that for large times, one expect to see the state $V(y)$ at any location (and not the migration process) and it holds whatever the initial distribution of the population is. Note furthermore that the profile $V(y)$ is unique. Hence whatever the initial structure of the population is, the phenotypic diversity at large times is completely determined by the profile of the function g .

In the slightly more general case of a positive non-linearity, we will prove the following existence theorem.

Theorem 1.3. *If f is of positive type, there exists $\alpha_0 > 0$ such that for $\alpha < \alpha_0$ there exists a traveling front of equation (6).*

Regarding the case of bistable f we have the following result:

Theorem 1.4. *If f is of bistable type, there exist $\alpha^* \geq \alpha_* > 0$ such that*

- For $\alpha \geq \alpha^*$, there exists no traveling front solution of (6),
- For $\alpha < \alpha_*$, under condition 43 of Section 7, there exists a traveling front u of speed $c > 0$ solution of (6).

Lastly, the model for CSD's leads one to equations of the type

$$\partial_t u - \Delta u = h(y, u) \quad x = (x_1, y) \in \mathbb{R}^N. \quad (9)$$

where $h(y, u)$ verifies

$$\begin{aligned} h(y, u) &= f(u) \text{ for } |y| \leq L_1 \\ h(y, u) &\leq -mu \text{ for } |y| \geq L_2 \\ h(y, u) + mu &\xrightarrow{|y| \rightarrow +\infty} 0 \quad \text{uniformly for } u \in \mathbb{R}^+ \end{aligned}$$

where $0 < L_1 \leq L_2 < \infty$ and $m > 0$ are given parameters and f is of bistable form.

In this paper we prove the following Theorem.

Theorem 1.5. *There exist critical radii $0 < L_* \leq L^* < \infty$ with the following properties:*

- *For $L_2 < L_*$, there is no traveling front solution of (9).*
- *For $L_1 > L^*$ (independently of L_2), assuming that there is a unique stable asymptotic profile of (53), there exists a traveling front of speed $c > 0$ solution of (9).*

The assumption on the uniqueness of the asymptotic profile is proved to be true for the case $N - 1 = 2$, $L_1 = L_2$ and $h(y, s) = -ms$ for $|y| \geq L_2$. This is done in [20] by phase plane method. For want of a uniqueness result for the profile equation in more general cases,

This theorem completes the study in [17] on the existence of CSD in the human brain. Indeed in [17] the transition from gray to white matter was instantaneous when biologically there is a smooth transition from gray to white matter. This Theorem confirms the intuition that CSD's can be found in part of the human brain where the gray matter is sufficiently thick but they can not propagate over large distances due to a thin gray matter in many parts of the human brain.

The paper is organized as follows. In section 2 we state some preliminary results that will be used in the sequel. Section 3 is dedicated to the study of the existence and uniqueness of non-zero asymptotic profiles for a traveling front solution of (6). In section 5 we study the large time behavior. There we prove extinction if $\alpha \geq \alpha_0$ and convergence towards the front of minimal speed if $\alpha < \alpha_0$. Section 6 extends existence of traveling front results to the case of a positive non-linearity. Then, section 7 is devoted to the study of the asymptotic profiles in the bistable case and section 8 to the existence of traveling front for $\alpha < \alpha_*$ in the bistable case. Lastly, in section 9 we describe the precise problem arising in the modeling of CSD's and state our main result in this framework.

2 Preliminary results

In our proofs, we will need several times the exponential decay of the asymptotic profile which can be easily proved from the following theorem established in [16].

Theorem 2.1. *Let $v \in H_{loc}^2(\mathbb{R}^N)$ be a positive function. Assume that there exists $\gamma > 0$ and $C > 0$ such that*

$$\forall x \in \mathbb{R}^N, \quad v(x) \leq Ce^{\sqrt{\gamma}|x|} \quad \text{and} \quad \liminf_{|x| \rightarrow \infty} \frac{\Delta v(x)}{v(x)} > \gamma.$$

Then, $\lim_{|x| \rightarrow \infty} v(x)e^{\sqrt{\gamma}|x|} = 0$.

This result is established in [16], lemma 2.2. In the context of equation (1), we thus have the following corollary.

Corollary 1. *Let u be a non-negative and bounded solution of*

$$\Delta v + f(v) - \alpha g v = 0 \quad \text{on } \mathbb{R}^{N-1}.$$

Then, for any $\gamma > 0$ there exists $C > 0$ such that

$$0 \leq v(y) \leq C e^{-\gamma|y|} \quad \text{and} \quad |\nabla v(y)| \leq C e^{-\gamma|y|}.$$

Proof. The estimate on v comes directly from Theorem 2.1 and the estimate on $|\nabla v|$ derives from standard global L^p estimates. \square

3 The case of a Fisher-KPP non-linearity. Asymptotic profiles.

In this section, we are interested in the asymptotic profiles of a traveling front solution of (6) as $x_1 \rightarrow \pm\infty$. Hence, we are looking for solutions of the following equation

$$\begin{cases} \Delta V + f(V) - \alpha g(y)V = 0, & y \in \mathbb{R}^{N-1}, \\ V \geq 0, & V \text{ bounded.} \end{cases} \quad (10)$$

We assume that $f : \mathbb{R} \rightarrow \mathbb{R}$ is C^1 ,

$$f(0) = f(1) = 0, \quad f > 0 \text{ on } (0, 1) \quad (11)$$

and

$$s \in (0, 1] \mapsto \frac{f(s)}{s} \text{ is decreasing.} \quad (12)$$

Since the constant function 0 is always a solution, the problem is to know when there exist non-zero solutions. As we will see here, the existence of such a positive asymptotic profile is characterized by the sign of the principal eigenvalue of the linearized operator around 0. We now make this notion precise.

3.1 Principal eigenvalue of the linearized operator

To start with, let us define the natural weighted space

$$\mathcal{H} = \{v \in H^1(\mathbb{R}^{N-1}), \sqrt{g}u \in L^2(\mathbb{R}^{N-1})\}$$

and its associated norm. For $v \in \mathcal{H}$, we set $\|v\|_{\mathcal{H}} = (\|v\|_{H^1}^2 + \|\sqrt{g}v\|_{L^2}^2)^{\frac{1}{2}}$. The linearized operator about 0 is $L\varphi = -\Delta\varphi + (\alpha g(y) - f'(0))\varphi$ for $\varphi \in \mathcal{H}$. We are interested in the eigenvalues of L . Even though the problem is set on all of \mathbb{R}^{N-1} , the term in $\alpha g(y)$ yields compactness of the injection $\mathcal{H} \hookrightarrow L^2(\mathbb{R}^{N-1})$. Hence the existence of a principal eigenvalue is obtained as usual.

Theorem 3.1. *Let us define*

$$R_{\alpha}(\varphi) = \frac{\int |\nabla \varphi|^2 + (\alpha g - f'(0))\varphi^2}{\int \varphi^2}.$$

The operator L has a smallest eigenvalue

$$\lambda_{\alpha} = \inf_{\varphi \in \mathcal{H} \setminus \{0\}} R_{\alpha}(\varphi). \quad (13)$$

Moreover there exists a unique positive eigenfunction associated with λ_α of L^2 -norm equal to 1, called φ_α in the following. The eigenspace associated with λ_α is spanned by φ_α .

The proof is classical due to the compactness of $\mathcal{H} \hookrightarrow L^2(R^{N-1})$. We refer for example to [23].

Remark 1. If $g(y) = |y|^2$, the problem can be rescaled and we obtain the harmonic oscillator for which principal eigenvalue and eigenfunction are well known [42]. In that case, $\lambda_\alpha = (N-1)\sqrt{\alpha} - f'(0)$ and $\varphi_\alpha = \left(\frac{\sqrt{\alpha}}{\pi}\right)^{\frac{1}{N-1}} e^{-\frac{\sqrt{\alpha}}{2}|y|^2}$.

Since the existence of a positive solution of (10) will depend on the sign of the principal eigenvalue, the following proposition describes the behavior of λ_α as a function of α .

Proposition 1. *The function $\alpha \mapsto \lambda_\alpha$ is continuous, increasing and concave for $\alpha \in (0, +\infty)$. Moreover $\lim_{\alpha \rightarrow 0} \lambda_\alpha = -f'(0)$ and for α large enough $\lambda_\alpha > 0$.*

Proof. Let us fix $\alpha > 0$ and $\eta > 0$. Equation (13) shows that

$$\lambda_{\alpha+\eta} \leq \int |\nabla \varphi_\alpha|^2 + ((\alpha + \eta)g - f'(0))\varphi_\alpha^2 = \lambda_\alpha + \eta \int g(y)\varphi_\alpha^2.$$

Similarly, we obtain $\lambda_\alpha \leq \lambda_{\alpha+\eta} - \eta \int g\varphi_{\alpha+\eta}^2$. From this we derive:

$$0 < \eta \int g\varphi_{\alpha+\eta}^2 \leq \lambda_{\alpha+\eta} - \lambda_\alpha \leq \eta \int g\varphi_\alpha^2.$$

This and similar computation for $\lambda_\alpha - \lambda_{\alpha-\eta}$ yields that $\alpha \mapsto \lambda_\alpha$ is increasing and locally Lipschitz on $(0, +\infty)$.

Concavity is classical. It suffices to observe that for each fixed φ ,

$$\alpha \mapsto R_\alpha(\varphi) = \frac{\int |\nabla \varphi|^2 + (\alpha g - f'(0))\varphi^2}{\int \varphi^2}$$

is an affine function of α and that $\lambda_\alpha = \inf_{\varphi \in \mathcal{H} \setminus \{0\}} R_\alpha(\varphi)$.

In order to prove that $\lambda_\alpha \xrightarrow{\alpha \rightarrow 0} -f'(0)$, for any $\varepsilon > 0$ choose a function ψ_ε of compact support with $\|\psi_\varepsilon\|_{L^2} = 1$ and $\int |\nabla \psi_\varepsilon|^2 < \varepsilon$. Let $\text{supp } \psi_\varepsilon \subset B_{R_\varepsilon}$. From (13) we get

$$-f'(0) \leq \lambda_\alpha \leq \varepsilon + \alpha \max_{B_{R_\varepsilon}} g - f'(0)$$

So for any $\alpha < \frac{\varepsilon}{\max_{B_{R_\varepsilon}} g}$,

$$-f'(0) \leq \lambda_\alpha \leq -f'(0) + 2\varepsilon.$$

Now we claim that $\lambda_\alpha > 0$ for large enough α . Argue by contradiction and assume that $\lambda_\alpha \leq 0$ for all $\alpha \in (0, +\infty)$. Since

$$0 \geq \lambda_\alpha = \int |\nabla \varphi_\alpha|^2 + \alpha \int g\varphi_\alpha^2 - f'(0),$$

we get

$$\int g\varphi_\alpha^2 \leq \frac{1}{\alpha}f'(0)$$

and $\varphi_\alpha \rightarrow 0$ in $L^2(\mathbb{R}^{N-1} \setminus B_R)$ for all $R > 0$. Furthermore, φ_α is bounded in \mathcal{H} and up to extraction we can assume that φ_α converges strongly in $L^2(\mathbb{R}^{N-1})$, thus φ_α converges to 0 in L^2 but this is impossible since $\int \varphi_\alpha^2 = 1$ for all $\alpha > 0$. \square

Corollary 2. *There exists $\alpha_0 > 0$ such that $\lambda_\alpha < 0$ for $\alpha < \alpha_0$, $\lambda_{\alpha_0} = 0$ and $\lambda_\alpha > 0$ for $\alpha > \alpha_0$.*

3.2 If g vanishes on B_{r_0}

The main part of the proof still holds if g vanishes on B_{r_0} but the result is slightly modified.

In this section, we assume that there exists $r_0 > 0$ such that (7) is substituted by the following assumption

$$g : \mathbb{R}^{N-1} \rightarrow \mathbb{R}_+ \in C^0, \quad g \equiv 0 \text{ on } B_{r_0} \text{ and } g > 0 \text{ on } \mathbb{R}^{N-1} \setminus B_{r_0}. \quad (14)$$

We define λ_Δ the principal eigenvalue of the Laplacian on B_{r_0} with Dirichlet boundary conditions, i.e.

$$\begin{cases} -\Delta\phi_0 = \lambda_\Delta\phi_0 & \text{on } B_{r_0}, \\ \phi_0 = 0 & \text{on } \partial B_{r_0}. \end{cases}$$

In this case, the principal eigenvalue of the linearized operator about 0 is well defined and Proposition 1 becomes

Proposition 2. *The function $\alpha \mapsto \lambda_\alpha$ is continuous, increasing and concave for $\alpha \in (0, +\infty)$, and $\lim_{\alpha \rightarrow 0} \lambda_\alpha = -f'(0)$. Now there are two cases:*

- i) *If $f'(0) < \lambda_\Delta$, then for α large enough $\lambda_\alpha > 0$.*
- ii) *If $f'(0) \geq \lambda_\Delta$, then $\lambda_\alpha \leq 0$ for all $\alpha > 0$.*

Proof. The proof of the first part of the proposition is exactly the same as in Proposition 1. We just have to prove i) and ii).

i) We assume that $f'(0) < \lambda_\Delta$ and argue by contradiction assuming that $\lambda_\alpha \leq 0$ for all $\alpha \in (0, +\infty)$. As in the proof of proposition 1, we have

$$\int g\varphi_\alpha^2 \leq \frac{1}{\alpha}f'(0)$$

and this yields $\varphi_\alpha \rightarrow 0$ in $L^2(\mathbb{R}^{N-1} \setminus B_R)$ for $\alpha \rightarrow +\infty$ but now for all $R > r_0$ only.

As before φ_α is bounded in \mathcal{H} and up to extraction, we have $\lambda_\alpha \rightarrow \lambda \leq 0$, weak convergence in \mathcal{H} and strong convergence in L^2 of φ_α to ϕ . The limit ϕ verifies $\int \phi^2 = 1$, $\phi \equiv 0$ for $|y| > r_0$ and

$$-\Delta\phi - f'(0)\phi = \lambda\phi$$

Thus $\phi \in H_0^1(B_{r_0})$ must coincide with ϕ_0 in B_{r_0} and $\lambda + f'(0) = \lambda_\Delta$ leading to $f'(0) \geq \lambda_\Delta$ since $\lambda \leq 0$. This is a contradiction.

ii) By taking $\varphi = \phi_0$ in the Rayleigh quotient (13), where ϕ_0 is the principal eigenvalue of the above problem in B_{r_0} with Dirichlet boundary conditions, we see that $\lambda_\alpha \leq \lambda_\Delta - f'(0) \leq 0$ for all $\alpha > 0$. \square

In the following, we will not state the results specifically for this case (14) and will rather assume (7). However, the proofs and results developed here carry over to this case with the obvious modifications.

3.3 Existence of non-zero asymptotic profile

Theorem 3.2. *For $\alpha \geq \alpha_0$, there is no solution of (10), where α_0 is defined in corollary 2. For $\alpha < \alpha_0$, there exists a unique positive solution of (10).*

Proof. Let us fix $\alpha \geq \alpha_0$. Then $\lambda_\alpha \geq 0$. Assume by contradiction that there exists a solution V of (10). Then the strong maximum principle shows that $V > 0$.

Since φ_α is an eigenfunction of the linearized operator L and V is solution of (10), we have

$$\begin{aligned} \int (\Delta V + f(V) - \alpha g V) \varphi_\alpha &= 0 \\ &= \int (\Delta \varphi_\alpha + (f'(0) - \alpha g) \varphi_\alpha + \lambda_\alpha \varphi_\alpha) V \end{aligned}$$

Now from corollary 1, V and ∇V are rapidly decreasing for $|y| \rightarrow \infty$ and so we can apply Stokes formula $\int \Delta V \varphi_\alpha = \int V \Delta \varphi_\alpha$. It yields $\int (f(V) - f'(0)V) \varphi_\alpha = \lambda_\alpha \int \varphi_\alpha V$ but $f(V) - f'(0)V < 0$ since f is of Fisher-KPP type and $\lambda_\alpha \geq 0$ thus a contradiction is obtained.

We now turn to the case $\alpha < \alpha_0$. For $\alpha < \alpha_0$, the eigenvalue λ_α is negative. Setting $\underline{V} = \varepsilon \varphi_\alpha$ with $\varepsilon > 0$, we get

$$\Delta \underline{V} + f(\underline{V}) - \alpha g \underline{V} = -\lambda_\alpha \varepsilon \varphi_\alpha + f(\varepsilon \varphi_\alpha) - f'(0) \varepsilon \varphi_\alpha \geq 0$$

if $\varepsilon > 0$ is chosen small enough. Hence \underline{V} is a sub-solution of (10). The constant function 1 is a super-solution and $\underline{V} \leq 1$ if ε is small enough. Therefore by the sub- and super-solution method, there exists a solution V such that $0 < \underline{V} \leq V \leq 1$.

Now consider V and W two non-zero solutions of (10). We argue by contradiction and assume that $V \not\equiv W$. Then for example $\Omega = \{y \in \mathbb{R}^{N-1}, V(y) < W(y)\}$ is not empty. Introduce a cutoff function $\beta \in C^\infty(\mathbb{R})$ with $\beta = 0$ on $(-\infty, 1/2]$, $\beta = 1$ on $[1, +\infty)$ and $0 < \beta' < 4$ on $(1/2, 1)$ and for all $\varepsilon > 0$, let us set $\beta_\varepsilon(s) = \beta(\frac{s}{\varepsilon})$. Using equation (10), we have

$$\begin{aligned} \int (-V \Delta W + \Delta V W) \beta_\varepsilon(W - V) &= \int (V f(W) - f(V) W) \beta_\varepsilon(W - V) \\ &\xrightarrow{\varepsilon \rightarrow 0} \int_\Omega (V f(W) - f(V) W) \end{aligned}$$

by Lebesgue's dominated convergence theorem. Owing to corollary 1, V , ∇V , W and ∇W have exponential decay and thus Stokes formula can be applied and we obtain

$$\begin{aligned} \int (-V \Delta W + \Delta V W) \beta_\varepsilon(W - V) &= \int \beta'_\varepsilon(W - V) \nabla(W - V) \cdot (V \nabla W - W \nabla V) \\ &= \underbrace{\int \beta'_\varepsilon(W - V) V |\nabla(W - V)|^2}_{=I_1} - \underbrace{\int \beta'_\varepsilon(W - V) (W - V) \nabla(W - V) \cdot \nabla V}_{=I_2}. \end{aligned}$$

In the term I_2 the integrand satisfies

$$|\beta'_\varepsilon(W - V)(W - V)\nabla(W - V) \cdot \nabla W| \leq 4|\nabla(W - V)| \cdot |\nabla W|$$

Therefore by Lebesgue's Theorem of dominated convergence, we infer that $I_2 \rightarrow 0$. Next the term I_1 satisfies $I_1 \geq 0$. Consequently, we may write:

$$0 \geq \int_{\Omega} (Vf(W) - Wf(V)) = \int_{\Omega} \left(\frac{f(W)}{W} - \frac{f(V)}{V} \right) VW$$

which is a contradiction in view of (12) as $W > V$ in Ω . Hence $V = W$ and the non-zero solution is unique. \square

The last point concerns the stability of the asymptotic profiles for $\alpha < \alpha_0$. Let us start by studying the energy of V . For $w \in \mathcal{H}$, we define the energy

$$J_\alpha(w) = \int_{\mathbb{R}^{N-1}} |\nabla w(y)|^2 + \frac{\alpha}{2} g(y) w^2(y) - F(w(y)) dy \quad (15)$$

where $F(u) = \int_0^u f(t)dt$.

Theorem 3.3. *For $\alpha < \alpha_0$, the unique positive solution of (10) V is stable in the energy sense, i.e. V is the global minimum of J_α and, furthermore $J_\alpha(V) < 0 = J_\alpha(0)$.*

Proof. Owing to the maximum principle, solutions of (10) are between 0 and 1. Hence we can modify f on $] -\infty, 0[$ such that it becomes odd and as a consequence, F can be considered as even. Since λ_α the principal eigenvalue of the linearized operator about the zero solution is negative for $\alpha < \alpha_0$, 0 cannot be the global minimum of J_α . Now J_α admits a global minimum that will be called \tilde{V} for the argument. One can prove that $|\tilde{V}|$ is also a global minimum of J_α and hence $|\tilde{V}|$ is a positive solution of (10). By uniqueness, $|\tilde{V}| = V$ and thus V is a global minimum of J_α . Since 0 is not a global minimum, necessarily $J_\alpha(V) < 0 = J_\alpha(0)$. \square

We now conclude with the linearized stability of V .

Theorem 3.4. *For $\alpha < \alpha_0$, consider the linearized operator about V and denote $\lambda_1[V]$ the principal eigenvalue of this operator. Then $\lambda_1[V] > 0$.*

Proof. Denote by ψ a positive eigenfunction associated with $\lambda_1[V]$ and assume by contradiction that $\lambda_1[V] \leq 0$. If $\lambda_1[V] < 0$, it is easy to see that for $\varepsilon > 0$ small enough $V + \varepsilon\psi < 1$ is a sub-solution of (10). From there, it would follow that there exists a solution of (10) between $V + \varepsilon\psi$ and 1 but this contradicts the uniqueness of V .

Now if $\lambda_1[V] = 0$, letting ψ be as above, we get

$$-\Delta\psi + \alpha g(y)\psi - f'(V)\psi = 0. \quad (16)$$

From the equation and since V is unique for every given $0 < \alpha < \alpha_0$, it is clear that V is differentiable with respect to α and that $w := \frac{\partial V}{\partial \alpha}$ satisfies:

$$-\Delta w + \alpha g(y)w - f'(V)w = -g(y)V. \quad (17)$$

We know that $w \leq 0$ and from (17) which shows that $w \not\equiv 0$, we actually see from the maximum principle that $w < 0$ in \mathbb{R}^{N-1} . It is also easily seen that w has exponential decay at infinity. From (16) and (17), it then follows that $\int_{\mathbb{R}^{N-1}} gVw = 0$ which is a contradiction. Hence $\lambda_1[V] > 0$. \square

4 Traveling fronts for a Fisher-KPP non-linearity

This section is devoted to the definition of a speed c^* for which a traveling front of equation (6) exists for $\alpha \in (0, \alpha_0)$. The threshold of existence of the non-zero asymptotic profile is called α_0 as in the previous section. For $0 < \alpha < \alpha_0$, V denotes the unique non-zero asymptotic profile. As shown in the previous section, the energy of the non-zero profile $J_\alpha(V)$ is negative.

A curved traveling front of speed c is a function $u(x_1 - ct, y)$ solution of equation (6) and connecting the non-zero asymptotic state V to 0. Thus we are looking for a solution of

$$\begin{cases} -\Delta u - c\partial_1 u + \alpha g(y)u = f(u), & x = (x_1, y) \in \mathbb{R}^N \\ u(x_1, \cdot) \xrightarrow{x_1 \rightarrow -\infty} V, & u(x_1, \cdot) \xrightarrow{x_1 \rightarrow +\infty} 0 \text{ uniformly in } y \in \mathbb{R}^{N-1}, \\ u \geq 0, & u \text{ bounded} \end{cases} \quad (18)$$

where $c \in \mathbb{R}$ is also an unknown of the problem.

The construction of c^* in Theorem 1.1 uses the sliding method following ideas of [14]. Note however that there are important differences with [14]. In that paper, the Fisher-KPP case is derived by first solving the “combustion non-linearity” and then approach the Fisher-KPP non-linearity as a limiting case of truncated functions. Contrary to [14] here, we derive directly the existence of a solution of the Fisher-KPP case. Actually the method we present here can be applied to somewhat simplify the proof of [14] in the Fisher-KPP case for cylinder with Neumann conditions.

4.1 Problem on a domain bounded in x_1 .

Let us fix $a > 1$ and $c \in \mathbb{R}$ for this subsection and consider the following problem:

$$\begin{cases} -\Delta u - c\partial_1 u + \alpha g(y)u = f(u), & x = (x_1, y) \in (-a, a) \times \mathbb{R}^{N-1} \\ u(-a, \cdot) = V, & u(a, \cdot) = 0, \\ u \geq 0, & u \text{ bounded.} \end{cases} \quad (19)$$

The aim of this subsection is to prove the following theorem:

Theorem 4.1. *There exists a unique solution of (19), denoted u_a^c in the following. This solution decreases in the x_1 -direction, i.e. $\partial_1 u_a^c < 0$. Thus $0 < u_a^c < V$ on $(-a, a) \times \mathbb{R}^{N-1}$. Moreover $c \mapsto u_a^c$ is decreasing and continuous from \mathbb{R} to $L^\infty([-a, a] \times \mathbb{R}^{N-1})$.*

To prove this theorem, we require the following two propositions.

Proposition 3. *Let u be a solution of (19). Then $u(x_1, y) \leq V(y)$ for $(x_1, y) \in [-a, a] \times \mathbb{R}^{N-1}$.*

Proof. Let $M \geq 1$ be such that $u \leq M$ and consider ψ_R defined on B_R the largest solution of

$$\begin{cases} -\Delta_y \psi_R + \alpha g(y)\psi_R = f(\psi_R) & \text{for } y \in B_R, \\ \psi_R = M & \text{for } y \in \partial B_R, \quad 0 \leq \psi_R \leq M. \end{cases} \quad (20)$$

Here we think of f as having been extended by 0 outside $[0, 1]$. Since $f(s) \leq 0$ for all $s \geq 1$, we observe that:

- by the strong maximum principle, $0 < \psi_R < M$ on B_R .
- since $V \leq 1 \leq M$ and V is a sub-solution of (20), through monotone iterations we have $V \leq \psi_R$.
- if $R' > R$, $\psi_{R'}$ is once again a sub-solution of (20) on B_R and thus $\psi_{R'} \leq \psi_R$ on B_R .
- therefore ψ_R tends to a function when $R \rightarrow +\infty$ and through local elliptic estimates, this function is a non-zero solution ($\geq V$) of the asymptotic problem (10). By uniqueness, we obtain $\psi_R \xrightarrow{R \rightarrow +\infty} V$

Now we consider the problem

$$\begin{cases} -\Delta w - c\partial_1 w + \alpha g(y)w = f(w) & \text{for } x \in (-a, a) \times B_R, \\ w = M & \text{for } x \in (-a, a) \times \partial B_R, \quad w = \psi_R \text{ for } x_1 = \pm a, y \in B_R. \end{cases} \quad (21)$$

The solution u of (19) is a sub-solution of (21) and the constant function M is a super-solution. Using monotone iterations starting from the super-solution M , we build the same sequence as previously (for problem (20)) since by induction the solutions do not depend on $x_1 \in (-a, a)$. Hence the sequence converges toward ψ_R and we have $u \leq \psi_R \leq M$. Now letting $R \rightarrow +\infty$ yields $u \leq V$. \square

Proposition 4. *Let R be such that $g(y) > \frac{K}{\alpha}$ for $y \notin \overline{B_R}$ where K is the Lipschitz norm of f on $[0, 1]$. We set $\Omega = I \times (\mathbb{R}^{N-1} \setminus \overline{B_R})$ where I is an open bounded interval of \mathbb{R} .*

Suppose u and $v \in C^2(\Omega) \cap C^0(\overline{\Omega})$ are solutions of

$$-\Delta w - c\partial_1 w + \alpha g(y)w = f(w) \quad \text{on } \Omega \quad (22)$$

and $u \leq v$ on $\partial\Omega$. Then $u \leq v$ on Ω .

Proof. By contradiction, suppose this is not true. Due to corollary 1 and proposition 3, $u(x_1, y)$ and $v(x_1, y)$ converge uniformly to 0 for $|y| \rightarrow +\infty$. Consequently, there exist $(x_0, y_0) \in \Omega$ such that

$$0 > \min_{\overline{\Omega}}(v - u) = (v - u)(x_0, y_0).$$

Since $(x_0, y_0) \in \Omega$, we have $\partial_1(v - u)(x_0, y_0) = 0$ and $\Delta(v - u)(x_0, y_0) \geq 0$, and subtracting the equation (22) with u from the one with v , we obtain

$$\alpha g(|y_0|)(v - u)(x_0, y_0) \geq f(v(x_0, y_0)) - f(u(x_0, y_0)) \geq -K|(v - u)(x_0, y_0)|$$

which is impossible since $\alpha g(|y_0|) > K$ and $(v - u)(x_0, y_0) < 0$. \square

Let us now turn to the proof of Theorem 4.1 using sliding method.

First $\bar{u}(x, y) = V(y)$ is a super-solution, 0 is a sub-solution and $0 \leq \bar{u}$, so by monotone iterations, there exists a solution u of (19).

Lemma 4.2. *Assume u and v are two solutions of (19). Then*

$$v(x_1 + h, y) \leq u(x_1, y) \text{ for all } h \in [0, 2a) \text{ and all } (x_1, y) \in [-a, a - h] \times \mathbb{R}^{N-1}.$$

Proof of the lemma. By proposition 3, we have $0 \leq u \leq V$ (resp. $0 \leq v \leq V$) and using the strong maximum principle, we obtain $0 < u < V$ (resp. $0 < v < V$) on $(-a, a) \times \mathbb{R}^{N-1}$.

For $h \in [0, 2a)$, let $I_h = (-a, a - h)$ and for $(x_1, y) \in \overline{I_h} \times \mathbb{R}^{N-1}$, set $v_h(x_1, y) = v(x_1 + h, y)$.

Let us fix $R > 0$ such that $g(y) > \frac{K}{\alpha}$ for $y \notin \overline{B_R}$. By compactness and continuity of u and v , there exists $\varepsilon > 0$ such that $v_h \leq u$ on $\overline{I_h} \times \overline{B_R}$ for any h such that $2a - \varepsilon \leq h < 2a$. Proposition 4 shows that $v_h \leq u$ on $\overline{I_h} \times \mathbb{R}^{N-1}$ for any $h \geq 2a - \varepsilon$. This enables us to define

$$h^* = \inf\{h \geq 0, v_h \leq u \text{ on } \overline{I_h} \times \mathbb{R}^{N-1}\}.$$

Let us prove that $h^* = 0$ and argue by contradiction that $h^* > 0$. By continuity, $v_{h^*} \leq u$ on $\overline{I_{h^*}} \times \mathbb{R}^{N-1}$.

Suppose that $\min_{I_{h^*} \times B_R} u - v_{h^*} > 0$. This would imply that for $h^* - h > 0$ small, $\min_{I_h \times B_R} u - v_h > 0$ and by Proposition 4, $v_h \leq u$ on $\overline{I_h} \times \mathbb{R}^{N-1}$ in contradiction with the definition of h^* . Therefore $\min_{I_{h^*} \times B_R} u - v_{h^*} = 0$. This implies the existence of $(x_1^*, y^*) \in I_{h^*} \times \overline{B_R}$ such that $v_{h^*}(x_1^*, y^*) = u(x_1^*, y^*)$ (note that $u - v_{h^*} > 0$ for $x_1 = -a$ or $a - h^*$). Writing in the usual way that $u - v_{h^*}$ is solution of a linear elliptic equation in $I_{h^*} \times \mathbb{R}^{N-1}$ and $u - v_{h^*} \geq 0$ with $u - v_{h^*}$ vanishing at the point (x_1^*, y^*) , the strong maximum principle implies that $u - v_{h^*} \equiv 0$ which is impossible. \square

Applying the preceding lemma with $h = 0$ yields the uniqueness of the solution of (19). Taking $u = v = u_a^c$, one sees that u_a^c is monotone decreasing. Thus $\partial_1 u_a^c \leq 0$ and deriving equation (19) and applying once more the maximum principle gives $\partial_1 u_a^c < 0$.

It remains to study the behavior of u_a^c with respect to c . The continuity is deduced from the uniqueness of the solution and a priori estimates in the standard way. Now let $c_1 < c_2$ and denote by u_1 (resp. by u_2) the solution of (19) with $c = c_1$ (resp. $c = c_2$). Since $\partial_1 u_1 < 0$,

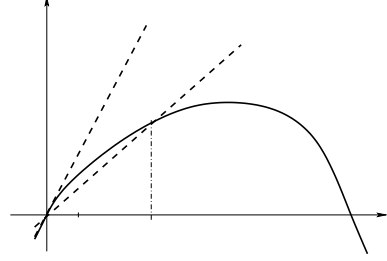
$$\Delta u_1 + c_2 \partial_1 u_1 + f(u_1) - \alpha g(y) u_1 = (c_2 - c_1) \partial_1 u_1 < 0$$

and $u_1 > 0$ is a super-solution of equation (19) with $c = c_2$. By uniqueness of the solution, necessarily $u_2 \leq u_1$. Once more the strong maximum principle implies $u_2 < u_1$.

4.2 Convergence to a solution on \mathbb{R}^N

Now that the equation is solved on a domain bounded in the x_1 -direction, the idea is to increase a up to infinity so that the domain tends to \mathbb{R}^N . However if c is chosen arbitrarily, the function u_a^c may converge toward the constant 0 or to V when a tends to infinity. Hence we adopt a normalization method as in [14]. The following theorem will define the value of the speed c depending on a to avoid those situations. We recall that since $\alpha < \alpha_0$, λ_α the principal eigenvalue of the linearized operator about the solution 0 is negative.

Theorem 4.3. *Let us fix $\varepsilon > 0$. Let $\delta > 0$ be such that $\delta < -\lambda_\alpha \leq f'(0)$. Let $\eta > 0$ be such that $\forall s \in [0, \eta] f(s) \geq (f'(0) - \delta)s$. We fix $\theta \in (0, \frac{\eta}{2})$. Then there exists $A(\varepsilon) > 0$ such that for all $a \geq A(\varepsilon)$, there exists a unique speed $c_a \in (0, 2\sqrt{-\lambda_\alpha} + \varepsilon)$ with $u_a^{c_a}(0, 0) = \theta$.*



Proof. By continuity and monotonicity, it suffices to prove:

- i) $u_a^0(0, 0) > \theta$,
- ii) $u_a^c(0, 0) < \theta$ for $c = 2\sqrt{-\lambda_\alpha} + \varepsilon$ and a large enough.

i) Assume $c = 0$. Let φ_α be the positive eigenfunction of the linearized operator L associated with the first eigenvalue $\lambda_\alpha < 0$ and with the normalization $\|\varphi_\alpha\|_\infty = 1$. Let us introduce $v(x_1, y) = h(x_1)\varphi_\alpha(y)$ for $(x_1, y) \in [-a, a] \times \mathbb{R}^N$ where $h(x_1) = \eta \frac{a-x_1}{2a}$. Then $0 < v \leq \eta$ on $[-a, a] \times \mathbb{R}^N$, which yields

$$\begin{aligned} -\Delta v + \alpha g(y)v - f(v) &\leq -\Delta v + \alpha g(y)v - (f'(0) - \delta)v \\ &= (\lambda_\alpha + \delta)v \leq 0. \end{aligned}$$

Moreover by construction of V (cf section 3), $v(-a, y) = \eta\varphi_\alpha(y) < V(y)$ if η is small enough. Then $v(a, y) = 0$ and $v \leq 1$. Hence, v is a sub-solution of (19) for $c = 0$. Thus $v \leq u_a^0$ and therefore, $u_a^0(0, 0) \geq v(0, 0) = \frac{\eta}{2} > \theta$.

ii) Let us construct an explicit super-solution for $c = 2\sqrt{-\lambda_\alpha} + \varepsilon$. We recall from section 3.1 that $\lambda_\beta < \lambda_\alpha$ for $\beta < \alpha$ and that $\lim_{\beta \rightarrow \alpha} \lambda_\beta = \lambda_\alpha$. Thus there exists $\beta \in (0, \alpha)$ such that $2\sqrt{-\lambda_\beta} \leq 2\sqrt{-\lambda_\alpha} + \varepsilon$. As before, let ψ_β denote the positive eigenfunction of the linearized operator L associated with the first eigenvalue $\lambda_\beta < 0$ with the normalization $\psi_\beta(0) = 1$. Choose R such that for all $r \geq R$, $(\alpha - \beta)g(r) + \lambda_\beta > 0$ and $\alpha g(r) > f'(0)$, and choose $k > 0$ such that $k\psi_\beta \geq V$ on \overline{B}_R . The constant k only depends on β hence on ε .

Lemma 4.4. *Then $k\psi_\beta \geq V$ on \mathbb{R}^{N-1} .*

Proof of the lemma. We follow a similar proof to that of lemma 4:

If the lemma does not stand, since $k\psi_\beta - V$ tends to 0 at ∞ , there exists $y_0 \in \mathbb{R}^{N-1} \setminus \overline{B}_R$ such that $(k\psi_\beta - V)(y_0) = \min_{\mathbb{R}^{N-1}}(k\psi_\beta - V) < 0$. At this point, $\Delta(k\psi_\beta - V) \geq 0$ but

$$\begin{aligned} -\Delta(k\psi_\beta - V) + \alpha g(y_0)(k\psi_\beta - V) - f'(0)(k\psi_\beta - V) &= \\ ((\alpha - \beta)g(|y_0|) + \lambda_\beta)k\psi_\beta + f'(0)V - f(V) &\geq ((\alpha - \beta)g(|y_0|) + \lambda_\beta)k\psi_\beta. \end{aligned}$$

By the choice of R , we get $\Delta(k\psi_\beta - V)(y_0) < 0$ which yields a contradiction. \square

Let us now build the super-solution when $c = 2\sqrt{-\lambda_\alpha} + \varepsilon$. We set $w(x_1, y) = z(x_1)k\psi_\beta(y)$ where z is the solution of

$$\begin{cases} z'' + cz' - \lambda_\beta z = 0 & \text{on } (-a, a), \\ z(-a) = 1, \quad z(a) = 0. \end{cases}$$

Then w verifies

$$\begin{cases} -\Delta w - c\partial_1 w + \alpha g(y)w = (\alpha - \beta)g(y)w + f'(0)w \geq f(w) \\ w(-a, \cdot) = k\psi_\beta \geq V, \quad w(a, \cdot) = 0. \end{cases}$$

so it is indeed a super-solution of (19). Moreover

$$z(x) = \frac{e^{\rho_+(x-a)} - e^{\rho_-(x-a)}}{e^{-\rho_+2a} - e^{-\rho_-2a}} \geq 0$$

where $\rho_- < \rho_+ < 0$ are the roots of $\rho^2 + c\rho - \lambda_\beta = 0$, i.e. $\rho_\pm = \frac{-c \pm \sqrt{c^2 + 4\lambda_\beta}}{2}$ (note that $c^2 + 4\lambda_\beta \geq 0$). Hence

$$\begin{aligned} 0 < u_a^c(0, 0) &< w(0, 0) = \frac{e^{-\rho_+a} - e^{-\rho_-a}}{e^{-\rho_+2a} - e^{-\rho_-2a}} k\psi_\beta(0) \\ &= \frac{1}{e^{-\rho_+a} + e^{-\rho_-a}} k\psi_\beta(0) \leq e^{-\frac{c}{2}a} k \leq e^{-a\sqrt{-\lambda_\alpha}} k \end{aligned}$$

Thus if a is large enough to have $e^{-a\sqrt{-\lambda_\alpha}} < \frac{\theta}{k}$, then for $c = 2\sqrt{-\lambda_\alpha} + \varepsilon$, we get $u_a^c(0, 0) < \theta$. \square

With the bounds on the speed c_a it is now possible to pass to the limit as a tends to infinity.

Proposition 5. *There exists a sequence $(a_j)_{j \in \mathbb{N}}$ such that $a_j \rightarrow +\infty$, $c_{a_j} \rightarrow c^* \in [0, 2\sqrt{-\lambda_\alpha}]$ and $u_{a_j}^{c_{a_j}} \rightarrow u$ in $C_{loc}^2(\mathbb{R}^N)$. The limit u is solution of*

$$\begin{cases} -\Delta u - c^*\partial_1 u + \alpha g(y)u = f(u) & \text{on } \mathbb{R}^N \\ 0 \leq u \leq V, \quad u(0, 0) = \theta, \quad \partial_1 u \leq 0. \end{cases} \quad (23)$$

Then u is necessarily a traveling front solution of (18) with $c = c^*$.

Proof. First for $j \in \mathbb{N}$, fix $\varepsilon_j = \frac{1}{j+1}$ and $a_j \geq A(\varepsilon_j)$ with $a_j \rightarrow +\infty$. Since $c_{a_j} \in (0, 2\sqrt{-\lambda_\alpha} + \varepsilon_j]$ is bounded, $u_{a_j}^{c_{a_j}}$ is uniformly bounded in $C^{2,\gamma}$ for any $\gamma \in (0, 1)$. Hence up to an extraction of a subsequence, there exist $c^* \in [0, 2\sqrt{-\lambda_\alpha}]$ and $u \in C_{loc}^2$ such that $c_a \rightarrow c^*$ and $u_a^{c_a} \rightarrow u$. Clearly the function u is a solution of (23). Owing to the normalization $u(0, 0) = \theta$, u is not a constant, moreover by the maximum principle $0 < u < V$ and $\partial_x u < 0$. Since u is decreasing in x , $u_\pm = \lim_{x \rightarrow \pm\infty} u(x, \cdot)$ are solutions of (10) and $u_-(0) > \theta > 0$ and $0 \leq u_+(0) < \theta$. This implies that $u_- = V$ and $u_+ \equiv 0$. Thus u is indeed a traveling front solution of (18). \square

4.3 Existence of traveling front for $c \geq 2\sqrt{-\lambda_\alpha}$.

In this section we still assume $0 < \alpha < \alpha_0$ and we will prove the following theorem.

Theorem 4.5. *There exists a traveling front of speed c of equation (18) if and only if $c \geq 2\sqrt{-\lambda_\alpha}$.*

We start with the Proposition

Proposition 6. *For $c < 2\sqrt{-\lambda_\alpha}$ there exists no traveling front solution of (18). Thus $c^* = 2\sqrt{-\lambda_\alpha}$ (where c^* is the traveling speed constructed in the previous section).*

Proof. We argue by contradiction and assume that there exists a traveling front u of speed $c < 2\sqrt{-\lambda_\alpha}$ of (18). We are going to construct a small positive sub-solution with compact support. To this end, we can find $\delta \in (0, f'(0))$ such that $c^2 + 4(\lambda_\alpha + 2\delta) < 0$ and $\eta > 0$ such that for all $s \in [0, \eta]$, $f(s) \geq (f'(0) - \delta)s$.

Since the linearized operator $L = -\Delta + \alpha g(y) - f'(0)$ is self adjoint, the principal eigenvalue λ_α is the limit of the Dirichlet principal eigenvalue in B_R when $R \rightarrow \infty$ (see [15] for more details):

$$\begin{cases} -\Delta \psi^R + \alpha g(y) \psi^R - f'(0) \psi^R = \lambda_\alpha^R \psi^R, & y \in B_R, \\ \psi^R > 0 & \text{on } B_R, \quad \psi^R = 0 & \text{on } \partial B_R. \end{cases} \quad (24)$$

Precisely $\lambda_\alpha^R > \lambda_\alpha$ and $\lambda_\alpha^R \xrightarrow{R \rightarrow \infty} \lambda_\alpha$. In the following, ψ^R denotes the positive eigenfunction with $\|\psi^R\|_\infty = 1$ and let us fix R sufficiently large so that $\lambda_\alpha < \lambda_\alpha^R < \lambda_\alpha + \delta$.

Let $\tilde{\sigma} = \sigma + i\frac{\pi}{2L}$, $L > 0$, be an imaginary root of $X^2 + cX - \lambda_\alpha^R - \delta = 0$ which is possible since $c^2 + 4(\lambda_\alpha^R + \delta) < c^2 + 4(\lambda_\alpha + 2\delta) < 0$. Finally let us fix $\varepsilon > 0$ small enough such that $\varepsilon e^{\sigma x_1} < \eta$ and $\varepsilon e^{\sigma x_1} \psi^R(y) < u(x_1, y)$ for $x \in [-L, L]$ and $y \in \overline{B_R}$. We set

$$w(x_1, y) = \begin{cases} \varepsilon e^{\sigma x_1} \cos(\frac{\pi}{2L} x_1) \psi^R(y) & \text{if } -L < x_1 < L, y \in B_R, \\ 0 & \text{otherwise.} \end{cases} \quad (25)$$

Then w verifies

$$-\Delta w - c\partial_1 w + \alpha g(y)w = (f'(0) - \delta)w \leq f(w)$$

since $0 \leq w \leq \eta$. Moreover $w \leq u$ and $w > 0$ on $(-L, L) \times B_R$. Thus w is a generalized sub-solution with compact support [12].

Let us now derive a contradiction with the existence of a traveling front u . Translate u to the left by defining $u_\tau(x_1, y) = u(x_1 + \tau, y)$ for $\tau > 0$. Since $u(x_1, \cdot) \xrightarrow{x_1 \rightarrow +\infty} 0$, there exists $\tau^* \geq 0$ such that $u_{\tau^*} \geq w$ but $u_{\tau^*}(x_1^*, y^*) = w(x_1^*, y^*)$. Since $u_{\tau^*} > 0$, $x_1^* \in (-L, L)$ and $y^* \in B_R$ (an interior point of the support of w). Now since w is a sub-solution, the strong maximum principle yields $u_{\tau^*} \equiv w$ on $[-L, L] \times \overline{B_R}$, but this is impossible on the boundary. \square

We have already proved that for $c < 2\sqrt{-\lambda_\alpha}$, there exists no traveling front of speed c solution of (18) and that for $c = c^* = 2\sqrt{-\lambda_\alpha}$ there exists a traveling front of speed c . Let us prove that for any $c > c^*$ there exists at least a traveling front to conclude with Theorem 4.5. The proof goes as usual. We consider the following problem

$$\begin{cases} -\Delta u - c\partial_1 u + \alpha g(y)u = f(u), & x = (x_1, y) \in (-a, a) \times \mathbb{R}^{N-1} \\ u(-a, \cdot) = u^*(-a + r, \cdot), & u(a, \cdot) = u^*(a + r, \cdot) \end{cases} \quad (26)$$

where u^* is the traveling front of speed c^* . The function $u^*(\cdot + r, \cdot)$ is a strict super-solution of (26) (since $c > c^*$) when 0 is a strict sub-solution and $0 < u^*(\cdot + r, \cdot)$.

Hence as in theorem 4.1, it can be proved that there exists a unique solution v_a^r of (26) and moreover $\partial_x w_a^r < 0$ and

$$\forall (x_1, y) \in [-a, a] \times \mathbb{R}^{N-1} \quad V(y) > u^*(-a + r, y) \geq v_a^r(x_1, y) \geq u^*(a + r, y) > 0.$$

By uniqueness, w_a^r depends continuously on $r \in \mathbb{R}$, so $w_a^r \xrightarrow{r \rightarrow +\infty} 0$ and $w_a^r \xrightarrow{r \rightarrow -\infty} V$ uniformly on $[-a, a] \times \mathbb{R}^{N-1}$. Let us denote $u_a = v_a^r$ where r is chosen in order that $v_a^r(0, 0) = \theta$ (see previous section for definition of θ). Once again taking any sequence $a_n \rightarrow +\infty$, up to an extraction $u_{a_n} \rightarrow u$ in C_{loc}^2 and u is a traveling front of speed c solution of (18).

5 The case of a Fisher-KPP non-linearity. Asymptotic speed of spreading.

This section is concerned with the asymptotic behavior of the solutions of the parabolic problem

$$\begin{cases} \partial_t u - \Delta u = f(u) - \alpha g(y)u & \text{on } \mathbb{R} \times \mathbb{R}^N \\ u(0, x) = u_0(x) & \text{on } \mathbb{R}^N \end{cases} \quad (27)$$

where f is Fisher-KPP and u_0 is an initial condition at least bounded.

5.1 Extinction for $\alpha \geq \alpha_0$

Let us fix $\alpha \geq \alpha_0$. We recall that there is no positive asymptotic profile of (10).

Theorem 5.1. *For $u_0 \in L^\infty$, there exists a unique solution $u(t, x)$ of (27) and it verifies $u(t, x) \xrightarrow{t \rightarrow +\infty} 0$ uniformly for $x \in \mathbb{R}^N$.*

This section is devoted to the proof of this theorem.

Let us fix $S = \max(1, \|u_0\|_\infty)$. Then the constant functions 0 and S are respectively sub- and super-solutions of (27). Thus there exists $u(t, x)$ a solution of (27) such that $0 \leq u \leq S$. By the parabolic maximum principle, this solution is unique.

Let us define w the solution of (27) with the initial condition $w(0, x) = S$. Since the problem and the initial condition do not depend on x_1 , neither does w thus we will write $w(t, y)$. By the maximum principle, $0 \leq u \leq w \leq S$ and since S is a super-solution, $\partial_t w \leq 0$. Thus $w(t, y) \xrightarrow{t \rightarrow +\infty} W(y)$ and

$$0 \leq \limsup_{t \rightarrow +\infty} u \leq W \leq S.$$

Now by parabolic local estimates, W is necessarily solution of

$$-\Delta_y W = f(W) - \alpha g(y)W$$

and thus is a nonnegative asymptotic profile. Since $\alpha \geq \alpha_0$, $W \equiv 0$. So $u(t, x)$ tends to 0 for $t \rightarrow +\infty$ uniformly in \mathbb{R}^N .

5.2 Spreading for $\alpha < \alpha_0$

In this section we assume $\alpha < \alpha_0$. So there exists a critical speed c^* of existence of traveling front for (18). We assume that $u_0 \in C_0^0(\mathbb{R}^N)$, i.e. u_0 is continuous and compactly supported, and that $u_0 < V$ where V is the positive asymptotic profile solution of (10). We will prove the spreading of the solution of (27) but we first need the following theorem.

Theorem 5.2. *The unique solution of*

$$\begin{cases} -\Delta z - c\partial_1 z + \alpha g(y)z = f(z) & (x_1, y) \in \mathbb{R}^N, \\ 0 < z(x_1, y) \leq V(y) & (x_1, y) \in \mathbb{R}^N. \end{cases} \quad (28)$$

with $c < c^*$ is $z(x, y) \equiv V(y)$.

Proof. Let us consider the generalized sub-solution with compact support $w(x_1, y)$ defined in (25). This is possible since $c < c^*$. Up to a decrease of $\varepsilon > 0$, we can assume that $w \leq z$ on \mathbb{R}^N . Now by applying the sliding method to w^τ where $w^\tau(x_1, y) = w(x_1 + \tau, y)$ and z , one can prove that $w^\tau \leq z$ for all $\tau \in \mathbb{R}$. We can thus define

$$\forall y \in \mathbb{R}^{N-1} \quad \underline{z}(y) = \inf_{x_1 \in \mathbb{R}} z(x_1, y) \geq 0$$

and state that $z \not\equiv 0$. Now \underline{z} is a super-solution of (10) since $\underline{z} = \inf_{h \in \mathbb{R}} z(\cdot + h, \cdot)$ and an infimum of solutions is a super-solution.

Finally as in section 3, we can build a positive sub-solution of (10) smaller than \underline{z} and thus by monotone iteration we have a solution of (10) between these sub- and super-solution. By uniqueness of the positive solution, we obtain $V \leq z$. And due to condition in (28), we have $z \equiv V$. \square

Let us now turn to the precise study of the spreading of the solution of (27)

Theorem 5.3. *For $u_0 \in C_0^0(\mathbb{R}^N)$ with $u_0 < V$, there exists a unique solution u of (27) and*

$$\text{for any } c > c^* \quad \lim_{t \rightarrow +\infty} \sup_{|x_1| \geq ct} u(t, x) = 0, \quad (29)$$

$$\text{for any } c \text{ with } 0 \leq c < c^* \quad \lim_{t \rightarrow +\infty} \sup_{|x_1| < ct} |u(t, x) - V(y)| = 0. \quad (30)$$

Proof. Fix $c > c^*$. Let U denote a traveling front of speed c^* . Since $U(x_1, \cdot) \rightarrow V$ for $x_1 \rightarrow -\infty$ locally uniformly, there exists $L \in \mathbb{R}$ such that $U(x_1 - L, y) > u_0(x_1, y)$ for all $(x_1, y) \in \mathbb{R}^N$. Now considering $v(t, x) = U(x_1 - L - c^*t, y)$ and applying the comparison principle, we have $u(t, x) \leq v(t, x)$ for all $t \geq 0$ and $x \in \mathbb{R}^N$. Thus since U is decreasing in x_1

$$\sup_{|x_1| > ct} u(t, x) \leq \sup_{|x_1| > ct} U(x_1 - L - c^*t, y) = \sup_{y \in \mathbb{R}^{N-1}} U((c - c^*)t - L, y).$$

Since $c > c^*$ and $U(x_1, y) \xrightarrow{x_1 \rightarrow +\infty} 0$ uniformly in $y \in \mathbb{R}^{N-1}$. We see that $\sup_{x_1 \geq ct} u(t, x) \rightarrow 0$ as $t \rightarrow +\infty$. Since $u(t, -x_1, y)$ satisfies the same equation (27), this shows that $\sup_{x_1 \leq -ct} u(t, x) \rightarrow 0$ as well as $t \rightarrow +\infty$. Thus (29) is proved.

Assume now $c < c^*$. Let us first prove the following weaken version of (30):

Lemma 5.4. For any $c \in \mathbb{R}$ with $|c| \leq c^*$,

$$\forall (x_1, y) \in \mathbb{R}^N \quad \lim_{t \rightarrow +\infty} |u(t, x_1 - ct, y) - V(y)| = 0. \quad (31)$$

Proof of lemma 5.4. Let us assume that $c \geq 0$, the proof being similar for $c \leq 0$.

Let $v(t, x_1, y) = u(t, x_1 - ct, y)$. Then v satisfies the equation

$$\partial_t v - \Delta v - c\partial_1 v + \alpha g(y)v = f(v) \quad (32)$$

with the initial datum $v(0, x_1, y) = u_0(x, y) \geq 0$ and $\not\equiv 0$. Hence by the parabolic maximum principle, for all $(x_1, y) \in \mathbb{R}^N$ $v(1, x_1, y) > 0$. Now since $c < c^*$, in (25), we constructed $w(x_1, y) \geq 0$ a stationary non-zero sub-solution of (32) with compact support and w could be chosen arbitrary small. Hence we can assume $w \leq v(1, \cdot, \cdot)$. So if \tilde{w} is the solution of

$$\begin{cases} \partial_t \tilde{w} - \Delta \tilde{w} - c\partial_1 \tilde{w} + \alpha g(y)\tilde{w} = f(\tilde{w}) & t > 0, (x_1, y) \in \mathbb{R}^N \\ \tilde{w}(0, x_1, y) = w(x_1, y) & (x_1, y) \in \mathbb{R}^N \end{cases}$$

then by comparison principle, $\forall t \geq 1 \forall (x_1, y) \in \mathbb{R}^N$ $v(t, x_1, y) \geq \tilde{w}(t-1, x_1, y)$. Now since w is a sub-solution, \tilde{w} is increasing with respect to t and $0 \leq \tilde{w}(t, x_1, y) \leq V(y)$. Therefore, by standard elliptic estimates, $\tilde{w}(t, x_1, y) \xrightarrow{t \rightarrow +\infty} z(x, y)$ and z is a solution of (28). By theorem 5.2, we have $z \equiv V$ and this complete the proof of the lemma since by the comparison principle $\tilde{w}(t-1, x_1, y) \leq v(t, x_1, y) \leq V(y)$ thus

$$\lim_{t \rightarrow +\infty} v(t, x_1, y) = V(y)$$

which yields (31). \square

Let us now prove (30), that is the uniform convergence to V in the expanding slab $\{x_1 \leq ct\}$. We will only prove it for $0 \leq x_1 \leq ct$. Indeed using as before $u(t, -x_1, y)$, the general result follows from the convergence in the set $\{0 \leq x_1 \leq ct\}$.

Let c with $0 < c < c^*$ be fixed and let $\varepsilon > 0$ be given (arbitrarily small). For $R > 0$ sufficiently large, we know that the principal eigenvalue λ_α^R of the problem (24) above is such that $\lambda_\alpha^R < 0$. Denote by $\psi^R > 0$ the corresponding eigenfunction of (24). Under these conditions we know that there exists a unique solution $V^R(y) > 0$ of the profile equation in B_R with Dirichlet condition:

$$\begin{cases} -\Delta V^R + \alpha g(y)V^R = f(V^R) & \text{in } B_R \\ V^R = 0 & \text{on } \partial B_R, \quad V^R > 0 \text{ in } B_R. \end{cases} \quad (33)$$

(Compare e.g. [5]). Moreover, it is straightforward to show that V^R is increasing with R and that $\lim_{R \rightarrow +\infty} V^R(y) = V(y)$.

Let us choose $R > 0$ sufficiently large so that for all $y \notin B_R$ $V(y) < \varepsilon$ and for all $y \in \overline{B_R}$ $0 < V(y) - V^R(y) < \varepsilon$. The proof of the uniform convergence to V for $c < c^*$ will rest on the following Proposition.

Proposition 7. Let c be such that $0 < c < c^*$. Then, with R chosen as above, there exists a solution $v_c(x_1, y)$ defined for $x_1 \in \mathbb{R}^-$, $y \in \overline{B_R}$ of equation

$$-\Delta v - c\partial_1 v + \alpha g(y)v = f(v) \quad x_1 \leq 0, y \in \overline{B_R} \quad (34)$$

satisfying the following properties:

$$\begin{cases} v_c > 0 \text{ and } \partial_1 v_c < 0 \text{ in } \mathbb{R}_*^- \times B_R, \\ v_c(0, y) = 0 \text{ for } y \in \overline{B_R}, \\ v_c(x_1, y) = 0 \text{ for } y \in \partial B_R, \ x_1 \leq 0, \\ v_c(-\infty, y) = V^R(y) \text{ for } y \in \overline{B_R}. \end{cases}$$

Postponing the proof of this proposition, let us complete the proof of Theorem 5.3. Extending v_c by 0 for $x_1 \geq 0$ turns v_c into a (generalized) sub-solution of equation (34) in the cylinder $\mathbb{R} \times B_R$ (see [12]). Therefore $v_c(x_1 - c(t - t_0), y)$ is a sub-solution of the equation (27) in this cylinder for all $t_0 \geq 0$ and all $c \in (0, c^*)$.

By Lemma 5.4 (applied here in the case $c = 0$), we can fix $t_0 > 0$ sufficiently large such that for $t \geq t_0$ we have

$$u(t, 0, y) \geq V(y) - \frac{\delta}{2} \quad \text{for all } y \in \overline{B_R}$$

where $\delta = \min_{\overline{B_R}} (V - V^R) > 0$. Therefore,

$$u(t, 0, y) > V^R(y) \quad \text{for all } t \geq t_0 \text{ and all } y \in \overline{B_R}.$$

We fix $\tilde{c} \in (c, c^*)$ and we consider $v(t, x_1, y) = v_{\tilde{c}}(x_1 - \tilde{c}(t - t_0), y)$. In the region $D = (0, +\infty) \times B_R$, u is a solution and v a sub-solution of equation (27) and for any time $t \geq t_0$

$$u(t, x_1, y) \geq v(t, x_1, y) \quad \text{for } (x_1, y) \in \partial D.$$

Moreover, $u(t_0, x_1, y) \geq v(t_0, x_1, y) = 0$ in D . The comparison principle then yields

$$u(t, x_1, y) \geq v(t, x_1, y) \quad \text{in } D.$$

Therefore

$$\begin{aligned} \limsup_{t \rightarrow +\infty} \sup_{\substack{0 \leq x_1 \leq ct \\ y \in B_R}} (V(y) - u(t, x)) &\leq \limsup_{t \rightarrow +\infty} \sup_{\substack{0 \leq x_1 \leq ct \\ y \in B_R}} (V(y) - v_{\tilde{c}}(x_1 - \tilde{c}(t - t_0), y)) \\ &\leq \limsup_{t \rightarrow +\infty} \sup_{y \in B_R} (v_{\tilde{c}}((c - \tilde{c})t + \tilde{c}t_0, y)) \\ &\leq \sup_{y \in B_R} (V(y) - V^R(y)) < \varepsilon. \end{aligned}$$

Outside of B_R we already know that $0 < u < V < \varepsilon$ for any $t \geq 0$, $x_1 \in \mathbb{R}$ and $|y| \geq R$. Therefore

$$\limsup_{t \rightarrow +\infty} \sup_{0 \leq x_1 \leq ct} (V(y) - u(t, x)) \leq \varepsilon$$

Since this is true for all $\varepsilon > 0$ (and for $-ct \leq x_1 \leq 0$), we have thereby established (30).

It now remains to prove Proposition 7 which we carry now. As in (25), we construct a sub-solution of the equation (34) with compact support, namely:

$$w(x_1, y) = \begin{cases} \varepsilon e^{\sigma x_1} \cos(\frac{\pi}{2L} x_1 + \frac{\pi}{2}) \psi^R(y) & \text{if } -2L < x_1 < 0, \ y \in B_R, \\ 0 & \text{otherwise.} \end{cases}$$

In comparison with (25), there is a translation in x_1 such that the support of w now lies in $\mathbb{R}_- \times B_R$.

For any $b < 0$, let z_b be the solution of

$$\begin{cases} -\Delta z_b - c\partial_1 z_b + \alpha g(y)z_b = f(z_b) & \text{in } (b, 0) \times B_R, \\ z_b(b, y) = V^R(y), \quad z_b(0, y) = 0 & \text{for } y \in \overline{B_R} \\ z_b(x_1, y) = 0 & \text{for } x_1 \in (b, 0), |y| = R. \end{cases}$$

Since V^R is a super-solution and 0 a sub-solution, there exists a solution of this problem. By the sliding method of [4], we know that this solution is unique and satisfies $\partial_1 z_b < 0$ in $(b, 0) \times B_R$.

Next, for $b < -L$, since w is a sub-solution, we also know that

$$\forall b \leq -L \quad \forall (x_1, y) \in (b, 0) \times \overline{B} \quad z_b(x_1, y) > w(x_1, y).$$

This allows us to pass to the limit when $b \rightarrow -\infty$. Clearly $z_b(x_1, y) \xrightarrow{b \rightarrow -\infty} v_c(x_1, y)$. By the lower bound, $v_c(x_1, y) > w(x_1, y)$ which shows that $v_c(x_1, y) > 0$ in $\mathbb{R}_-^* \times B_R$. Since $\partial_1 v_c \leq 0$ and $v_c \not\equiv 0$, we also know that $\partial_1 v_c < 0$ in $\mathbb{R}_-^* \times B_R$. Now since $\lim_{x_1 \rightarrow -\infty} v_c(x_1, y)$ must be a positive solution of (33). Hence by uniqueness we get $v_c(-\infty, y) = V^R(y)$. This completes the proof of Proposition 7 and therefore of Theorem 5.3 \square

6 The case of a positive non-linearity.

In this section, we prove Theorem 1.3 about the existence of traveling fronts in the positive case. We use the notations of the preceding sections, in particular λ_α still denotes the principal eigenvalue of the linearized operator around 0 and φ_α an associated eigenfunction. We are interested in a traveling front solution of (6) when f is only assume to be of the positive type, that is $f : \mathbb{R} \rightarrow \mathbb{R}$ is C^1 with

$$f(0) = f(1) = 0, \quad f > 0 \text{ on } (0, 1) \text{ and } f'(0) > 0. \quad (35)$$

6.1 Asymptotic profiles

The linearized operator around 0 is exactly the same as in the Fisher-KPP case thus there exists $\alpha_0 > 0$ such that $\lambda_\alpha < 0$ for $\alpha < \alpha_0$ and $\lambda_\alpha \geq 0$ for $\alpha \geq \alpha_0$. In the same way as before, we can prove the existence result:

Proposition 8. *For $\alpha < \alpha_0$, there exists $V(y)$ a maximal positive asymptotic profile solution of (10).*

However, in this more general case, we have no information on the uniqueness of the positive asymptotic profile nor on the non-existence of profiles for $\alpha \geq \alpha_0$. Actually, this will depend on the non-linearity f .

Proof. As in Theorem 3.2, $\varepsilon\varphi_\alpha$ is a subsolution for $\varepsilon > 0$ small enough and 1 is a supersolution. Using monotone iterations, we can construct a maximal positive asymptotic profile. \square

Since the positive asymptotic profile may not be unique, we will need the following lemma before turning to the construction of traveling fronts.

Lemma 6.1. *For any $\alpha < \alpha_0$, there exists $\theta_\alpha > 0$ such that any positive asymptotic profile $W(y)$ solution of (10) satisfies*

$$W(0) \geq 2\theta_\alpha. \quad (36)$$

Proof. By contradiction, assume that there exist $W_n > 0$ solution of (10) with $W_n(0) \xrightarrow{n \rightarrow \infty} 0$. Then $\psi_n = \frac{W_n}{\|W_n\|_\infty}$ is a solution of

$$-\Delta\psi_n + \alpha g(y)\psi_n = \frac{\|W_n\|_\infty \psi_n}{\|W_n\|_\infty}.$$

Since ψ_n is bounded, up to extraction of a subsequence, we can let n tend to ∞ to obtain $\psi_n \rightarrow \psi_\infty \geq 0$ with $\|\psi_\infty\|_\infty = 1$ and

$$-\Delta\psi_\infty + \alpha g(y)\psi_\infty = f'(0)\psi_\infty.$$

In the previous limit, we made use of the compactness argument derived from the fact that $\psi_n(y) \rightarrow 0$ as $|y| \rightarrow \infty$ uniformly in n as is obtained from Theorem 2.1. Thus, ψ_∞ is a principal eigenfunction associated with the eigenvalue 0 which contradicts the fact that $\lambda_\alpha < 0$. \square

6.2 Existence of traveling fronts

In this section, we will use the same method as in section 4 to construct a traveling front solution of (10). However due to the possible non-uniqueness of the positive asymptotic profile, the result will be somewhat weaker in that the limiting profile is not prescribed. More precisely, we will prove the existence of $c \in \mathbb{R}$ and u solutions of

$$\begin{cases} -\Delta u - c\partial_1 u + \alpha g(y)u = f(u), & x = (x_1, y) \in \mathbb{R}^N \\ u(x_1, \cdot) \xrightarrow{x_1 \rightarrow +\infty} 0 \text{ uniformly in } y \in \mathbb{R}^{N-1}, \\ u > 0, \quad u \text{ bounded and } \partial_1 u < 0. \end{cases} \quad (37)$$

The construction of the solution follows the same line as before. We start by solving the problem on a domain bounded in x_1 , precisely we study solution of (19) where V is the maximal positive asymptotic profile. Since V is maximal, Theorem 4.1 still holds true. The only difficulty is to translate Theorem 4.3 to the case of a positive non-linearity. Having this aim in mind, we introduce two notations:

$$m = \sup_{(0,1]} \frac{f(s)}{s} \text{ and } \mu_\alpha = \lambda_\alpha + f'(0) - m \leq \lambda_\alpha < 0.$$

In the Fisher-KPP case, we observe that $m = f'(0)$ and thus $\mu_\alpha = \lambda_\alpha$. We will prove the following result.

Theorem 6.2. *Let $\varepsilon > 0$ be fixed such that $\varepsilon\varphi_\alpha < V$ (see previous section). Let $\delta > 0$ be such that $\delta < -\lambda_\alpha < f'(0)$ and let $\eta > 0$ be such that $\eta < \varepsilon$ and $\forall s \in [0, \eta] \ f(s) \geq (f'(0) - \delta)s$. We fix $\theta \in (0, \frac{\eta}{2})$ such that $\theta < \theta_\alpha$ (see Lemma 6.1 for definition of θ_α). Then there exists $A_\varepsilon > 0$ such that for all $a \geq A_\varepsilon$ there exists a unique speed $ca \in (0, 2\sqrt{-\mu_\alpha} + \varepsilon)$ such that $u_a^{c_a}$ the solution of (19) satisfies $u_a^{c_a}(0, 0) = \theta$.*

Proof. The only difference with the proof of Theorem 4.3 is in the upper bound of c_a (section ii) in the proof of Theorem 4.3). It goes as before but we need to replace λ_α by μ_α .

So let us construct an explicit super-solution of (19) for $c = 2\sqrt{-\mu_\alpha} + \varepsilon$. As before, we can fix $\beta < \alpha$ such that $2\sqrt{\mu_\beta} \leq 2\sqrt{\mu_\alpha} + \varepsilon$ and we consider ψ_β the positive eigenfunction of the linearized operator around 0 with the normalization $\|\psi_\beta\|_\infty = 1$. We fix $R > 0$ such that for all $r \geq R$

$$(\alpha - \beta)g(r) + \mu_\beta > 0 \text{ and } \alpha g(r) > m.$$

Let us then fix $k > 0$ such that $k\psi_\beta \geq V$ on \bar{B}_R . Then $k\psi_\beta \geq V$ on \mathbb{R}^{N-1} . Indeed we argue by contradiction and assume that $\min_{\mathbb{R}^{N-1}}(k\psi_\beta - V) = (k\psi_\beta - V)(y_0) < 0$ but at this point $y_0 \in \mathbb{R}^{N-1} \setminus \bar{B}_R$, we have

$$\begin{aligned} -\Delta(k\psi_\beta - V)(y_0) + (\alpha g(y_0) - m)(k\psi_\beta - V)(y_0) = \\ \mu_\beta k\psi_\beta(y_0) + (\alpha - \beta)g(y_0)k\psi_\beta(y_0) + mV(y_0) - f(V(y_0)) > 0 \end{aligned}$$

and this yields a contradiction.

Then we prove as in Theorem 4.3 that $w(x_1, y) = z(x_1)k\psi_\beta(y)$ is a super-solution if z is a solution of

$$\begin{cases} z'' + cz' - \mu_\beta z = 0 & \text{on } (-a, a), \\ z(-a) = 1, \quad z(a) = 0 \end{cases}$$

and as before we obtain that

$$0 < u_a^c(0, 0) < e^{-\frac{\varepsilon}{2}a}k \leq e^{-a\sqrt{-\mu_\alpha}}k$$

and the upper bound of c_a for large a is thus proved. \square

Then the convergence of $u_a^{c_a}$ to a solution of (37) when a tends to $+\infty$ is exactly the same except that the non-uniqueness of the positive asymptotic profile prevents us from determining the precise limit of $u(x_1, \cdot)$ for $x_1 \rightarrow -\infty$.

We leave it as an open problem to know whether there always is a traveling front connecting the maximum profile $V(y)$ to the 0 solution.

7 The case of a bistable non-linearity. Asymptotic profiles.

In this section we consider again equation (1) but in the bistable framework. That is, we assume that f is a C^1 function that satisfies the following assumptions for some $\theta \in (0, 1)$:

$$f(0) = f(\theta) = f(1) = 0, \quad f(s) < 0 \text{ for } s \in (0, \theta) \text{ and } f(s) > 0 \text{ in } (\theta, 1), \quad (38)$$

$$f'(0) > 0, \quad f'(1) > 0. \quad (39)$$

We also assume that

$$\int_0^1 f(s)ds > 0. \quad (40)$$

We are concerned here with the existence of traveling front solutions of (1), that is, (c, u) solution of (18). First we require some preliminary results on the equation (10) in the bistable case.

7.1 Existence of asymptotic profiles in the bistable case

Consider equation

$$\begin{cases} \Delta u + f(u) - \alpha g(y)u = 0, & y \in \mathbb{R}^{N-1}, \\ u \geq 0, & u \text{ bounded}, \end{cases} \quad (41)$$

under the same assumption (7) and (8) as above for the function g .

The existence of solutions depends on α and is obtained in the following theorem.

Theorem 7.1. *Let f and g satisfy the above assumptions. There exists a threshold value $\alpha^* \in (0, \infty)$, such that:*

- i) *For any $\alpha \in (\alpha^*, +\infty)$, (41) does not have any positive (non-zero) solution.*
- ii) *For any $\alpha \in (0, \alpha^*]$, (41) admits a maximal positive solution $V(y)$.*
- iii) *For any $\alpha \in (0, \alpha^*)$, (41) admits a second positive solution $W(y)$ with $0 < W(y) < V(y)$.*

The rest of this section is devoted to the proof of this Theorem.

This Theorem follows from the observation that for $\alpha > 0$ any positive solution $u(y)$ of (41) satisfies $u(y) \rightarrow 0$ as $|y| \rightarrow \infty$. This is obtained from Corollary 1.

Next, by the maximum principle, any solution of (41) satisfies $0 \leq u \leq 1$ (we think of $f(s)$ as having been extended by 0 outside $[0, 1]$).

Now $\bar{u} \equiv 1$ is a super-solution of problem (41). Any solution of (41) for α is a sub-solution of (41) for any parameter $\beta \leq \alpha$. Therefore, if there exists a positive bounded solution of (41) for α , there also exists a positive solution for any $0 < \beta \leq \alpha$.

Next, we claim that for small enough $\alpha > 0$, (41) admits a positive solution. Indeed, consider the functional defined on \mathcal{H} :

$$J(w) = J_\alpha(w) = \int_{\mathbb{R}^{N-1}} \left(\frac{1}{2} |\nabla w|^2 + \frac{\alpha}{2} g(y) w^2 - F(w) \right) dy$$

where $F(z) = \int_0^z f(s) ds$. Recall that f is extended by 0 outside $[0, 1]$, thus F is bounded. Since $g(r) \rightarrow \infty$ as $r \rightarrow \infty$, it is straightforward to show that there exists a minimizer v of $J(w)$: $J(v) = \min\{J(w), w \in \mathcal{H}\}$. Furthermore, we know that $v \geq 0$ and v is a solution of (41) (see Theorem 3.3 for details).

Let us show that for $\alpha > 0$ small enough $J(v) < 0$. To this end, let ζ_R be defined by

$$\zeta_R(y) = \begin{cases} 1 & \text{if } |y| \leq R \\ R + 1 - |y| & \text{if } R \leq |y| \leq R + 1 \\ 0 & \text{if } |y| \geq R + 1 \end{cases}$$

Then $\zeta_R \in \mathcal{H}$ and

$$J_0(\zeta_R) = \int_{\mathbb{R}^{N-1}} \frac{|\nabla \zeta_R|^2}{2} - F(\zeta_R) \leq -F(1)|B_R| + C|B_{R+1} \setminus B_R|$$

where $|A|$ denotes the volume of A and C is a constant. Since $-F(1) < 0$ by (40), we see that by choosing R large enough, $J_0(\zeta_R) < 0$. Then for such an R fixed,

we see that $J_\alpha(\zeta_R) < 0$ provided $\alpha > 0$ is small enough. This guarantees that $J_\alpha(v) < 0$.

It follows that $v \not\equiv 0$. By the maximum principle, we then have $0 < v < 1$. This shows that for small $\alpha > 0$, (41) admits a positive solution.

Next, we show that if α is large enough (41) does not admit any positive solution. This can be seen by multiplying the equation by u and integrating to yield:

$$\int |\nabla u|^2 + \alpha \int g(y)u^2 = \int f(u)u \leq m \int u^2 \quad (42)$$

where $m = \sup_{s>0} \frac{f(s)}{s} > 0$. We conclude with the following lemma:

Lemma 7.2. *Under the assumption (8) $g(r) \xrightarrow{r \rightarrow \infty} \infty$, for any $\varepsilon > 0$, there exists a constant $K(\varepsilon) > 0$ such that for all $u \in \mathcal{H}$ one has:*

$$\int_{\mathbb{R}^{N-1}} u^2 \leq \varepsilon \int_{\mathbb{R}^{N-1}} |\nabla u|^2 + K(\varepsilon) \int_{\mathbb{R}^{N-1}} g(y)u^2.$$

Indeed choosing in the lemma $\varepsilon = \frac{1}{2m}$, we get from (42)

$$\frac{1}{2} \int |\nabla u|^2 + \left(\alpha - mK\left(\frac{1}{2m}\right) \right) \int g(y)u^2 \leq 0$$

This shows that for $\alpha \geq mK(\frac{1}{2m})$, the only solution of (41) is $u \equiv 0$.

Proof of Lemma 7.2. Let $\delta = \delta(\varepsilon) > 0$ be chosen such that the principal eigenvalue of $-\Delta$ in $H_0^1(B_{2\delta})$ is larger than $\frac{4}{\varepsilon}$. Let χ be a smooth cutoff function such that $\chi(r) = 1$ if $0 \leq r \leq \delta$, $\chi(r) = 0$ if $r \geq 2\delta$ and $0 \leq \chi \leq 1$. Consider $u_1 = \chi u$ and $u_2 = (1 - \chi)u$ so that $u = u_1 + u_2$. Using $(a + b)^2 \leq 2(a^2 + b^2)$, since $u_1 \in H_0^1(B_{2\delta})$ by Poincaré's inequality, we have

$$\int_{\mathbb{R}^{N-1}} u_1^2 = \int_{B_{2\delta}} u_1^2 \leq \frac{\varepsilon}{4} \int_{B_{2\delta}} |\nabla u_1|^2 \leq \frac{\varepsilon}{2} \left(\int_{B_{2\delta}} |\nabla u|^2 \chi^2 + \int_{B_{2\delta} \setminus B_\delta} u^2 |\nabla \chi|^2 \right).$$

So that

$$\int_{\mathbb{R}^{N-1}} u_1^2 \leq \frac{\varepsilon}{2} \int_{B_{2\delta}} |\nabla u|^2 + \varepsilon k_1(\varepsilon) \int_{|y| \geq \delta} u^2$$

where $k_1(\varepsilon) \geq |\nabla \chi|^2$.

Next $\int_{\mathbb{R}^{N-1}} u_2^2 \leq \int_{|y| \geq \delta} u^2$. Therefore

$$\int_{\mathbb{R}^{N-1}} u^2 \leq 2 \left(\int u_1^2 + \int u_2^2 \right) \leq \varepsilon \int |\nabla u|^2 + K(\varepsilon) \int g(y)u^2$$

where $K(\varepsilon) = 2 \frac{\varepsilon k_1(\varepsilon) + 1}{g(\delta)}$. The lemma is thus proved. \square

The next step is to prove that the set of $\alpha > 0$ such that (41) has a solution is a closed set. Let $\alpha_j \rightarrow \alpha^*$ be a sequence such that (41) admits a solution u_j such that $0 < u_j < 1$ for all j . Note that by the maximum principle, $\theta < \max u_j < 1$. The sequence (u_j) is bounded by 1 and by standard elliptic estimates is locally

compact. Therefore, one can extract a subsequence u_j such that $u_j \rightarrow u^*$ uniformly on compact sets in the C^2 -norm. Therefore u^* is a solution of (41) for the value $\alpha = \alpha^*$. We know that $u^* \geq 0$, but since $\max u_j > \theta$, we see that $\max u^* \geq \theta$. Indeed by section 2, $u_j(y) \rightarrow 0$ as $|y| \rightarrow \infty$ uniformly with respect to j . Therefore $u^* > 0$ and (41) also has a positive solution for α^* . This shows that the set of α such that (41) has a positive solution is an interval $(0, \alpha^*]$ with $0 < \alpha^* < \infty$.

Considering the evolution equation

$$\begin{cases} \partial_t z - \Delta z + \alpha g(y)z = f(z), & t > 0, y \in \mathbb{R}^{N-1}, \\ z(0, y) = 1, \end{cases}$$

we see that $t \mapsto z(t, y) \geq 0$ is decreasing and therefore has a limit. This limit is necessarily the maximal positive solution $V = V_\alpha$ for the α for which (41) has a positive solution, that is $\alpha \in (0, \alpha^*]$, or is 0 in the opposite case, that is when $\alpha > \alpha^*$.

The existence of a second solution when $0 < \alpha < \alpha^*$ is inspired from a work of P. Rabinowitz [40]. In a slightly different formulation, the existence of pairs of solutions is established in [40] by a topological degree argument for bistable type nonlinearities and another type of parameter dependence. The use of the topological degree involves compact operators and the results of [40] are set in the framework of bounded domains. A similar construction can be carried here owing to the condition (8) $g(r) \rightarrow +\infty$ as $r \rightarrow +\infty$. Indeed, under this condition, the injection $\mathcal{H} \hookrightarrow L^2(\mathbb{R}^{N-1})$ is compact. This allows one to construct a compact operator and to carry the argument of [40] to the present framework.

Since we will not use the second solution, we will leave out the details of the proof of the existence of a second solution.

7.2 Stable asymptotic profiles

As we have seen, a solution of (41) is obtained by the minimization of $J = J_\alpha$ defined above. The proof of the existence of the previous solution for $\alpha > 0$ small yields the following result.

Proposition 9. *There exists $0 < \alpha_* \leq \alpha^*$ such that for all $\alpha \in (0, \alpha_*)$ there exists a minimum $v_\alpha > 0$ of J_α and such that*

$$J_\alpha(v_\alpha) = \min_{H^1(\mathbb{R}^{N-1})} J_\alpha < 0.$$

In the following, we require the notion of stable solution.

Definition 7.3. Let v be a solution of (41). Eigenvalues of the linearized problem about v are defined as the eigenvalues λ of

$$-\Delta \varphi + \alpha g(y)\varphi - f'(v)\varphi = \lambda \varphi \quad \text{in } \mathbb{R}^{N-1}.$$

The principle eigenvalue is uniquely determined by the existence of a corresponding eigenfunction φ with $\varphi > 0$. We say that v is (weakly) stable if the principal eigenvalue $\lambda = \lambda_1[v]$ of the linearized problem satisfies $\lambda_1[v] \geq 0$.

It is well known that the maximal solution $V(y)$ given by Theorem 7.1 when $0 < \alpha \leq \alpha^*$ is weakly stable. Likewise, the minimum solution of the energy of the Proposition 9 above, when $0 < \alpha < \alpha_*$, is a weakly stable solution.

In the following we consider the case $0 < \alpha < \alpha_*$ and we make the following assumption.

There exists a unique positive stable solution of (10). (43)

This condition implies that the minimizer solution v_α : $J_\alpha(v_\alpha) = \min\{J_\alpha(v), v \in \mathcal{H}\}$ coincides with the maximum solution V .

We leave it as an open problem to give sufficient conditions for the uniqueness of the stable solution. Uniqueness results have been given for analogous problems but with $\alpha = 0$, which would rather correspond to the minimal solution in our framework [38]. Likewise it would be interesting to give sufficient conditions that ensure that $\alpha_* = \alpha^*$. Condition (43) has several implications that we can state.

Proposition 10. *For $\alpha \in (0, \alpha_*]$, under condition (43), there does not exist a pair of distinct ordered functions (v_1, v_2) with $0 < v_1 \leq v_2 < V$, v_1 is a sub-solution and v_2 is a non-maximal solution. That is, if $0 < v_1 \leq v_2 < V$ are respectively sub-solution and solution of (41), then $v_1 \equiv v_2$.*

Proof. The proof follows the observation in [14]. However, it requires new elements in view of the unbounded domain. If $v_1 < v_2$, let φ_2 be a principal eigenfunction of the linearized problem corresponding to $\lambda_1[v_2]$. Since 0 and V are the only stable solutions, $\lambda_1[v_2] < 0$. We claim that for $\varepsilon > 0$ sufficiently small, $\bar{v} = v_1 - \varepsilon\varphi_2$ is a super-solution of (41). Indeed

$$\begin{aligned} -\Delta\bar{v} + \alpha g(y)\bar{v} - f(\bar{v}) &= f(v_2) - f(\bar{v}) - f'(v_2)\varepsilon\varphi_2 - \lambda_1[v_2]\varepsilon\varphi_2 \\ &= \left(\frac{f(v_2) - f(v_2 - \varepsilon\varphi_2)}{\varepsilon\varphi_2} - f'(v_2) - \lambda_1[v_2] \right) \varepsilon\varphi_2. \end{aligned}$$

The right hand side is positive if $\varepsilon > 0$ is sufficiently small.

Next, given $R > 0$, we can choose $\varepsilon > 0$ small enough so that $v_1 < v_2 - \varepsilon\varphi_2$ in $\overline{B_R}$. We choose R so that $v_1(y) \leq \delta$ for all $|y| \geq R$ and f is decreasing on $[0, \delta]$. We claim that then $v_1 \leq v_2 - \varepsilon\varphi_2$ in $\mathbb{R}^{N-1} \setminus \overline{B_R}$. Argue by contradiction. In this were not the case, then, since v_1, v_2 and φ_2 converge to 0 at infinity, there exists $y, |y| > R$ such that

$$\min_{\mathbb{R}^{N-1}} \{v_2 - \varepsilon\varphi_2 - v_1\} = v_2(y) - \varepsilon\varphi_2(y) - v_1(y) < 0$$

This implies that $0 < \bar{v}(y) < v_1(y) \leq \delta$. Denote L the operator $L = -\Delta + \alpha g(y)$. Since $0 \leq L(\bar{v} - v_1) - (f(\bar{v}) - f(v_1))$ and $f(\bar{v}(y)) - f(v_1(y)) > 0$, at the point y we get $L(\bar{v} - v_1)(y) > 0$. Therefore, we have reached a contradiction. This shows that $v_1 \leq v_2 - \varepsilon\varphi_2$. Now we have a super-solution \bar{v} above a sub-solution v_1 . This implies that there exists a stable solution v such that $v_1 \leq v \leq v_2 - \varepsilon\varphi_2 < V$. This however is in contradiction with condition (43). \square

From this property, we derive the following useful consequence.

Proposition 11. *Let $\alpha \in (0, \alpha_*)$ and let W be the maximal solution of equation (41) with the value α_* of the parameter. Then, any other solution v of (41) with parameter α that is not the maximal solution cannot be above W .*

This immediately follows from the previous proposition as W is a sub-solution of the equation for the value $\alpha < \alpha_*$ and $W < V$.

A consequence of this proposition is

Proposition 12. *For $\alpha \in (0, \alpha_*]$ and under condition (43), the maximal solution V is isolated in L^∞ topology. Therefore, there exists $\theta_1 > \theta$ such that if v is a solution of (41) with $v(0) \geq \theta_1$ then $v \equiv V$.*

As we have done before, we can prove that if v is a solution such that $v \geq W$ in $\overline{B_R}$, then $v \geq W < V$ in \mathbb{R}^{N-1} . Then any solution $v \neq V$ is such that there exists $y \in \overline{B_R}$ such that $v(y) \leq W(y)$ and therefore $\|v - V\|_{L^\infty} \geq \min_{\overline{B_R}} V - W = \delta > 0$.

Now, if there exist a sequence v_n of solutions of (41) such that $v_n(0) \rightarrow V(0)$ then by elliptic estimates $v_n \rightarrow W$ a positive solution of (41) and $W(0) = V(0)$ so $W \equiv V$ by the maximum principle which contradicts the fact that V is isolated.

8 Traveling fronts for a bistable non-linearity

In this section we assume that f is of bistable type and satisfies (38)-(40). In addition, we assume that $0 < \alpha < \alpha_*$ and that condition (43) is fulfilled. Therefore, there exists a unique non-zero stable solution $V(y) = V_\alpha(y)$ of the profile equation (41). Therefore $V > 0$, $J_\alpha(V) = \min\{J(w), w \in \mathcal{H}\}$, $\lambda_1[V] \geq 0$ and V is isolated in the L^∞ topology. Furthermore V is the maximal solution. Any other non-zero solution w satisfies $0 < w < V$ in \mathbb{R}^{N-1} and $\lambda_1[w] < 0$ where $\lambda_1[w]$ is the principal eigenvalue of the linearized problem defined in definition 7.3.

In this section, we prove the existence of a traveling front solution of (1) representing an invasion of 0 by the state V at positive speed. Such a solution is given as a pair (c, u) of

$$\begin{cases} -\Delta u - c\partial_1 u + \alpha g(y)u = f(u) & \text{in } \mathbb{R}^N \\ u(-\infty, y) = V(y), \quad u(+\infty, y) = 0 \end{cases} \quad (44)$$

with $c < 0$ and $u : \mathbb{R}^N \rightarrow (0, 1)$.

We follow the construction of a solution given above. Namely, let $a \geq 1$ and in the slab $\Sigma_a = (-a, a) \times \mathbb{R}^{N-1}$, consider the problem

$$\begin{cases} -\Delta u - c\partial_1 u + \alpha g(y)u = f(u) & \text{in } \Sigma_a, \\ u(-a, y) = V(y), \quad u(+a, y) = 0. \end{cases} \quad (45)$$

We recall that for any $c \in \mathbb{R}$, for a fixed, there exists a unique solution $u = u^c$ of (45). Furthermore, $0 < u < V$ and $\partial_1 u < 0$ in Σ_a . The mapping $c \mapsto u^c$ is decreasing.

Up to here, the procedure is the same as before. From this point on however, we need to modify the above argument since we used the fact that f was positive.

Our first task is to prove the following

Proposition 13. *There exists a unique (c_a, u_a) such that u_a is a solution of (45) for speed c_a and u_a satisfies the normalization condition*

$$\max_{y \in \mathbb{R}^{N-1}} u_a(0, y) = \theta. \quad (46)$$

Let us first prove the existence of c_a . The uniqueness is clear.

The parameter c_a is bounded independently of $a \geq 1$. Moreover,

$$\liminf_{a \rightarrow +\infty} c_a \geq 0.$$

Proof. The bound from above is obtained simply by comparison with the one dimensional problem. Indeed, consider the ODE problem for $z = z(x_1)$:

$$\begin{cases} -z'' - \gamma z' = f(z) & \text{in } (-a, a) \\ z(-a) = 1, \quad z(+a) = 0, \quad z(0) = \theta \end{cases} \quad (47)$$

It is known that there exists a unique value γ^a for which (47) has a (unique) solution z . Furthermore, $\lim_{a \rightarrow +\infty} \gamma^a = \gamma^*$ where γ^* is the unique speed of traveling fronts for the 1D equation

$$\begin{cases} -z'' - \gamma^* z' = f(z) & \text{in } (-a, a) \\ z(-\infty) = 1, \quad z(+\infty) = 0 \end{cases}$$

Comparing (47) with (45), we see that for each $c = \gamma^a$, the solution z of (47) is a super-solution of (45), thus $z > u^{\gamma^a}$ and for all $y \in \mathbb{R}^{N-1}$, $u^{\gamma^a}(0, y) < z(0) = \theta$. Since $c \mapsto u^c$ is decreasing, we see that

$$\max_{y \in \mathbb{R}^{N-1}} u^c(0, y) < \theta \text{ for all } c \geq \gamma_a.$$

Assume now that $\max_{y \in \mathbb{R}^{N-1}} u^c(0, y) < \theta$ for all $c \in \mathbb{R}$. Passing to the limit for $c \rightarrow -\infty$, u^c converges toward a positive solution v of (41) with $\max v < \theta$. By the maximum principle, it is impossible thus there exists a unique $c_a \in (-\infty, \gamma^a)$ such that (46) is fulfilled.

Since $\gamma^a \rightarrow \gamma^* < \infty$ as $a \rightarrow +\infty$ and $a \mapsto \gamma^a$ is a continuous function, this shows that $\sup_{a \geq 1} c_a < \infty$.

Since $a \mapsto c_a$ is continuous, in order to complete the proof of the Proposition, it suffices to show that $\liminf_{a \rightarrow \infty} c_a \geq 0$. For this, we argue by contradiction and assume that for a sequence $a_j \rightarrow +\infty$ there holds $c_{a_j} < 0$. For the sake of simplicity, we write a instead of a_j . Since $c \mapsto u^c$ is decreasing, from this we infer that along this subsequence, the solution $v = v^a$ of

$$\begin{cases} -\Delta v + \alpha g(y)v = f(v) & \text{in } \Sigma_a \\ v(-a, y) = V(y), \quad v(+a, y) = 0 \end{cases}$$

satisfies $\max_{y \in \mathbb{R}^{N-1}} v(0, y) \leq \theta$.

Due to Proposition 12, there exist $\theta_1 > \theta$ such that if an asymptotic profile v solution of (41) verifies $v(0) \geq \theta_1$ then $v \equiv V$.

There is a point $b = b_j$, $-a < b < 0$ such that $v_a(b, 0) = \theta_1$. We now translate the solution to center it on $x_1 = b$. That is, we let $\tilde{v}_a(x_1, y) = v_a(x_1 + b, y)$ defined for $x_1 \in (-a - b, a - b)$ and $y \in \mathbb{R}^{N-1}$. The interval $(-a - b, a - b)$ either converges (along a subsequence) to $(-\infty, +\infty)$ or to some $(-d, +\infty)$ with $0 \leq d < \infty$. In both cases, by standard elliptic estimates, one can strike out a subsequence of \tilde{v}_a , denoted again \tilde{v}_a , such that \tilde{v}_a converges locally to some function w where w satisfies:

$$\begin{cases} -\Delta w + \alpha g(y)w = f(w) & \text{in } (-d, +\infty) \times \mathbb{R}^{N-1} \\ \partial_1 w \leq 0, \quad w(0, 0) = \theta_1. \end{cases} \quad (48)$$

In case the interval is converging to $(-d, +\infty)$, in addition we know that $w(-d, y) = V(y)$. If the interval converges to \mathbb{R} , then $\lim_{x \rightarrow -\infty} w(x_1, y)$ exists and is some

function $W(y)$ which is then a solution of the profile equation (41). But since $w(0,0) = \theta_1$, we know that $W(0) \geq \theta_1$. By the definition of θ_1 , this implies that $W \equiv V$. Therefore, denoting $d = \infty$ in case $(-a-b, a-b) \rightarrow \mathbb{R}$, in both cases, we get

$$\forall y \in \mathbb{R}^{N-1} \quad w(-d, y) = V(y)$$

where now $0 \leq d \leq +\infty$. We also know that $w(+\infty, y) = \psi(y)$ exists with $0 \leq \psi < V$.

Multiply (48) by $\partial_1 w$ and integrate over $(-d, +\infty) \times \mathbb{R}^{N-1}$ to get

$$\int_{\{x_1=-d\}} \frac{1}{2} (\partial_1 w)^2 + J(\psi) - J(V) = 0$$

where $\partial_1 w = 0$ if $d = \infty$. In all cases, we get

$$J(V) \geq J(\psi)$$

Since V minimize J_α , we obtain $V \equiv \psi$ and $w(x_1, y) = V(y)$ for all $x_1 \in (-d, +\infty)$ but this contradicts the renormalization $w(0,0) = \theta_1$.

We have thus reached a contradiction. This shows that for large a , $c_a \geq 0$, which completes the proof of the Proposition. \square

Let us now turn to the proof of the existence of traveling front solutions of (44). Since c_a and u_a are bounded, by standard elliptic estimates, we can strike out a sequence $a = a_j \rightarrow \infty$ (we continue to denote subsequences by a) such that $c_a \rightarrow c \geq 0$ and $u_a \rightarrow u$. We know that (c, u) satisfies the equation

$$-\Delta u + c \partial_1 u + \alpha g(y)u = f(u) \quad \text{in } \mathbb{R}^N$$

with $\partial_1 u \leq 0$ and $\max_{\mathbb{R}^{N-1}} u(0, \cdot) = \theta$. It remains to identify the limits as $x_1 \rightarrow \pm\infty$. These $\lim_{x_1 \rightarrow \pm\infty} u(x_1, y) = u_\pm(y)$ exist and are solutions of the asymptotic profile equation (41). Now since $0 \leq u_+(y) = \lim_{x_1 \rightarrow +\infty} u(x_1, y) \leq \theta$ and all positive solutions w of (41) satisfy $\max w > \theta$, we have $u_+ \equiv 0$.

We claim that $u_-(y) = \lim_{x_1 \rightarrow -\infty} u(x_1, y)$ coincides with $V(y)$. Clearly, $0 < u_- \leq V$. Argue by contradiction that $u_- \not\equiv V$, implying $u_- < V$. By assumption, u_- is an unstable solution of (41) in the sense that $\lambda_1[u_-] < 0$. Let us construct a super-solution of the stationary equation, that is a w with

$$-\Delta w + \alpha g(y)w \geq f(w)$$

such that w is a compact perturbation of u_- and as close as we wish to u_- .

Consider the linearized equation about u_- :

$$-\Delta \psi - f'(u_-(y))\psi + \alpha g(y)\psi = \lambda_1[u_-]\psi$$

with $\lambda_1[u_-] < 0$. We know that $\lambda_1[u_-]$ is the limit of the Dirichlet principal eigenvalue in a ball when the radius goes to infinity (This follows from the Rayleigh quotient minimization). Therefore, $R > 0$ can be chosen sufficiently large so that the principal eigenvalue μ and associated eigenfunction ψ of

$$\begin{cases} -\Delta \psi + \alpha g(y)\psi - f'(u_-)\psi = \mu \psi & \text{in } B_R \\ \psi = 0 & \text{on } \partial B_R, \quad \psi > 0 \quad \text{in } B_R \end{cases}$$

satisfy $\mu < 0$.

Consider the function $\zeta(x_1, y) = \cos(\omega x_1)\psi(y)$ defined for $x_1 \in (-L, L)$ with $L = \frac{\pi}{2\omega}$ and $|y| < R$. We note $D = (-L, L) \times B_R$. This function is positive and satisfies:

$$\begin{cases} -\Delta\zeta + \alpha g(y)\zeta - f'(u_-)\zeta = (\mu + \omega^2)\zeta & \text{in } D \\ \zeta = 0 & \text{on } \partial D \end{cases}$$

Choose L large enough so that $\mu + \omega^2 < 0$. Then let $w(x_1, y) = u_-(y) - \varepsilon\zeta(x_1, y)$ with $\varepsilon > 0$ and $(x_1, y) \in D$. This function satisfies

$$-\Delta w + \alpha g(y)w - f(w) = \left(-(\mu + \omega^2) + \frac{f(u_-) - f(u_- - \varepsilon\zeta)}{\varepsilon\zeta} - f'(u_-) \right) \varepsilon\zeta.$$

Since $\mu + \omega^2 < 0$, we can choose ε sufficiently small so that

$$-\Delta w + \alpha g(y)w - f(w) \geq 0 \quad \text{in } D \text{ and } w > 0.$$

Furthermore, because $\varepsilon\zeta = 0$ on ∂D and $\varepsilon\zeta > 0$ in D , that is $w < u_-$ in D , if we extend w by choosing $w(x_1, y) = u_-(y)$ for all $(x_1, y) \notin D$, we have constructed a (generalized) super-solution of the problem (see e.g. [12]).

Let us now derive a contradiction. We consider two cases.

• Cases (i): Suppose $c > 0$. Then $U(t, x_1, y) = u(x_1 - ct, y)$ is a solution of the evolution equation

$$\partial_t U - \Delta U + \alpha g(y)U = f(U) \quad t \in \mathbb{R}, (x_1, y) \in \mathbb{R}^N.$$

Now $U \xrightarrow{t \rightarrow -\infty} 0$ locally uniformly in (x_1, y) . Furthermore, for all times $U(t, x_1, y) \leq u_-(y)$. Since w is a compact perturbation of u_- for a time t_0 sufficiently negative, we get

$$\forall (x_1, y) \in \mathbb{R}^N \quad U(t_0, x_1, y) \leq w(x_1, t).$$

Now when $U(t, x_1, y) \xrightarrow{t \rightarrow +\infty} u_-(y)$ locally uniformly and we get a contradiction since $U(t, x_1, y) \leq w(x_1, y) < u_-(y)$ for all $(x_1, y) \in D$.

• Case (ii): The case that remains to be studied is $c = 0$ (since we already have $c \geq 0$). Then $u(x_1, y)$ is a stationary solution of the same equation that w is a super-solution of. Since $u(-\infty, y) = u_-(y)$ and $u(+\infty, y) = 0$, and since $w = u_-$ outside a compact set, after a translation, we can assume that $u_h = u(x_1 + h, y) \leq w(x_1, y)$ (for large enough h). Define

$$h^* = \inf\{h \in \mathbb{R}, u(x_1 + h, y) \leq w(x_1, y) \text{ in } \mathbb{R}^N\}.$$

Clearly, $h^* > -\infty$ (for $w < u_-$ at some points). Then $w(x_1, y) \geq u(x_1 + h^*, y) = u_{h^*}$ and $\min(w - u_{h^*}) = 0$ is necessarily achieved at a point of \bar{D} . Since $w(x_1, y) = u_-(y) > u(x_1 + h, y)$ for all h if $(x_1, y) \notin D$, we see that the maximum is achieved at an interior point of D . Writing $w - u_{h^*} \geq 0$ as a super-solution of a linear elliptic equation in D , we derive a contradiction with the strong maximum principle.

Therefore in all cases, the solution u satisfies the limiting condition:

$$u(-\infty, y) = V(y), \quad u(+\infty, y) = 0.$$

Therefore (c, u) is a solution of the traveling front equation (44).

9 The model of cortical spreading depression

We consider here more general versions of the model (5) described in the Introduction. The problems studied in this paper have the following general form

$$\partial_t u - \Delta u = h(y, u) \quad x = (x_1, y) \in \mathbb{R}^N. \quad (49)$$

In the modeling context $N = 2$ and 3 are the cases of interest. As indicated in the Introduction, this equation also describes cortical spreading depressions (CSD). There the wave propagates in a medium composed of two different components, the gray and white matters of the brain, with a narrow transition area separating them.

Thus we consider in this section functions $h(y, u)$ of the following type:

$$h(y, u) = f(u) \text{ for } |y| \leq L_1 \quad (50)$$

$$h(y, u) \leq -mu \text{ for } |y| \geq L_2 \quad (51)$$

$$h(y, u) + mu \xrightarrow{|y| \rightarrow +\infty} 0 \quad \text{uniformly for } u \in \mathbb{R}^+ \quad (52)$$

where $0 < L_1 \leq L_2 < \infty$ and $K \geq m > 0$ are given parameters and f is of bistable form. That is we assume that f verifies conditions (38)-(40) of section 7. Note that in particular, we assume

$$\int_0^1 f(s) ds > 0.$$

We also assume that $y \mapsto h(y, s)$ is continuous and that $s \mapsto h(y, s)$ is Lipschitz continuous for all $s \in [0, 1]$ (and $|y| \neq L_1$ in case $L_1 = L_2$). Lastly we assume that

$$\forall s \in [0, 1] \quad \forall y \in \mathbb{R}^{N-1} \quad h(y, s) \leq \max\{f(s), -ms\}.$$

9.1 The asymptotic profile equation

We start as usual with the profile equation

$$\begin{cases} -\Delta V = h(y, V) & y \in \mathbb{R}^{N-1}, \\ V \geq 0, & V \text{ bounded.} \end{cases} \quad (53)$$

We recall that $\lambda_1[V]$ is the principal eigenvalue of the linearized equation about V . This eigenvalue can be defined as

$$\lambda_1[V] = \inf_{\varphi \in H^1(\mathbb{R}^{N-1})} \frac{\int |\nabla \varphi|^2 - \partial_2 h(y, V) \varphi^2}{\int \varphi^2}.$$

Associated with (53) is the energy functional:

$$J(w) = \int_{\mathbb{R}^{N-1}} \left(\frac{1}{2} |\nabla w|^2 - H(y, w) \right) dy$$

where $H(y, z) = \int_0^z h(y, s) ds$. Note that owing to condition (51), $J(w)$ is well defined for all $w \in H^1(\mathbb{R}^{N-1})$.

Theorem 9.1. *There exist critical radii $0 < L_* \leq L^* < \infty$ with the following properties:*

- i) For $L_2 < L_*$, there is no solution other than 0 to the asymptotic profile equation (53).
- ii) For $L_1 > L^*$ (independently of L_2), there exists a maximum solution V of (53) and this solution is stable in the sense that $\lambda_1[V] \geq 0$.
- iii) For all $L_1 > L^*$, the minimum of the energy functional is achieved at some non-zero function $V_J(y)$, i.e.

$$J(V_J) = \min_{w \in H^1(\mathbb{R}^{N-1})} J(w) < 0.$$

Proof. i) Since the equation implies that $-\Delta u + mu \leq 0$ for all $|y| \geq L_2$ in \mathbb{R}^{N-1} , and $u > 0$ is bounded, by Theorem 2.1 we know that u and $|\nabla u|$ have exponential decay as $|y| \rightarrow +\infty$. Then we get

$$\min(1, m)\|u\|_{H^1(\mathbb{R}^{N-1})} \leq \int_{\mathbb{R}^{N-1}} |\nabla u|^2 + mu^2 \leq \int_{B_{L_2}} (f(u) + mu)u \leq K \int_{B_{L_2}} u^2.$$

We know, by Sobolev embedding and Hölder inequality, that

$$\int_{B_{L_2}} u^2 \leq \eta(L_2)\|u\|_{H^1(\mathbb{R}^{N-1})}$$

where $\eta(L_2) \rightarrow 0$ as $L_2 \rightarrow 0$. Therefore for L_2 small enough, these inequalities yield $u \equiv 0$.

ii) Next, since 1 is a super-solution of the equation in \mathbb{R}^{N-1} , there exists a maximum solution of equation (53) that we denote V . By what we have just seen, $V \equiv 0$ for L_2 sufficiently small. Let us now show that $V > 0$ for L_1 sufficiently large.

Let us consider the energy restricted to the ball of radius $R \leq L_1$

$$J_R(w) = \int_{B_R} \left(\frac{1}{2} |\nabla w|^2 - F(w) \right) dy$$

where $F(z) = \int_0^1 f(s)ds$. We know (see the proof of Theorem 7.1) that for R sufficiently large there exists a minimum w_R of

$$J_R(w_R) = \inf_{w \in H_0^1(B_R)} J_R(w) < 0.$$

Then $w_R > 0$ in B_R and w_R is solution of $-\Delta w_R = f(w_R)$ in B_R and $w_R = 0$ on ∂B_R . Extending w_R by 0 outside the ball B_R , we get a global (generalized) sub-solution. The solution V is such that $V \geq w_R$ (since V is the maximum solution). This implies that $V \not\equiv 0$ and therefore $V > 0$ in \mathbb{R}^{N-1} for $L_1 \geq R$.

iii) Now for $L_1 \geq R$, clearly

$$\inf_{w \in H^1(\mathbb{R}^{N-1})} J(w) \leq J_R(w_R) < 0.$$

Let us now show that the infimum is achieved.

Let (w_n) be a minimizing sequence: $J(w_n) \rightarrow \inf J < 0$ for $n \rightarrow +\infty$. Note that J is bounded from below. Writing

$$\begin{aligned} J(w_n) &\geq \frac{1}{2} \int_{\mathbb{R}^{N-1} \setminus B_{L_2}} |\nabla w_n|^2 + m w_n^2 + \int_{B_{L_2}} \frac{1}{2} |\nabla w_n|^2 - H(y, w_n) \\ &\geq \frac{1}{2} \int_{\mathbb{R}^{N-1} \setminus B_{L_2}} |\nabla w_n|^2 + m w_n^2 + \int_{B_{L_2}} \frac{1}{2} |\nabla w_n|^2 - C + \varepsilon w_n^2 \\ &= -C|B_{L_2}| + \tilde{\varepsilon} \|w_n\|_{H^1(\mathbb{R}^{N-1})}^2 \end{aligned}$$

we can strike out a subsequence still denoted (w_n) such that $w_n \rightarrow w$ weakly in $H^1(\mathbb{R}^{N-1})$. Now using (52), for every $\varepsilon > 0$, there exists $R = R(\varepsilon) > 0$ such that $|F(y, s) + \frac{m}{2}s^2| < \varepsilon s^2$ for all $|y| \geq R$ and all $s \in \mathbb{R}^+$ (there is no loss in generality in assuming $w_n \geq 0$ as w_n^+ is also a minimizing sequence). Therefore

$$J(w_n) = \frac{1}{2} \int_{\mathbb{R}^{N-1}} |\nabla w_n|^2 + \frac{m}{2} \int_{|y| \geq R} w_n^2 + r(\varepsilon) - \int_{|y| \leq R} H(w_n, y)$$

where $|r(\varepsilon)| \leq C\varepsilon$ for some constant $C > 0$.

By compact injection of $H^1(\mathbb{R}^{N-1}) \hookrightarrow L^2(B_R)$, we can assume that $w_n \rightarrow w$ strongly in $L^2(B_{L_2})$. Then by standard arguments relying on Lebesgue's dominated convergence Theorem, we see that

$$\int_{B_R} H(y, w_n) \rightarrow \int_{B_R} H(y, w).$$

Owing to Sobolev embedding and Hölder inequality, it is straightforward to check that the quantity

$$\frac{1}{2} \int_{\mathbb{R}^{N-1}} |\nabla w_n|^2 + \frac{m}{2} \int_{|y| \geq R} w_n^2$$

defines the square of a norm equivalent to the usual $H^1(\mathbb{R}^{N-1})$ norm. Hence using the lower semi-continuity of the norm, we get

$$\lim_{n \rightarrow \infty} J(w_n) = \inf J \geq \frac{1}{2} \int_{\mathbb{R}^{N-1}} |\nabla w|^2 + \frac{m}{2} \int_{|y| \geq R} w^2 + r(\varepsilon) - \int_{B_R} H(y, w).$$

Now using again (52), we get:

$$J(w) = \int_{\mathbb{R}^{N-1}} |\nabla w|^2 - \int_{\mathbb{R}^{N-1}} H(y, w) \leq \inf_{H^1(\mathbb{R}^{N-1})} J + 2|r(\varepsilon)|.$$

Since $\varepsilon > 0$ is arbitrarily small we get $J(w) = \inf_{H^1(\mathbb{R}^{N-1})} J$. □

Remark 2. By using the method of [3], one can show that for L_1 large enough there exists a second solution in \mathbb{R}^{N-1} .

We will now make use of the condition that the stable solution of (53) is unique. In the paper of Chapuisat and Joly [20], it is argued by phase plane method, that for the case $N - 1 = 2$, $L_1 = L_2$ and $h(y, s) = -ms$ for $|y| \geq L_2$ that indeed this is the case. We note that it is an interesting open problem to derive such uniqueness results in more general situations or to complete the heuristic part of the argument of [20].

9.2 Traveling fronts for the CSD model

In this section, we prove Theorem 1.5. The proof is similar as in Section 8. There we used that $h(y, u) = f(u) - \alpha g(y)u$ with $g \rightarrow +\infty$. But actually, the same properties that were entailed one can derived for $h(y, u) \leq -mu$ for large $|y|$. We start by constructing a solution of

$$\begin{cases} \Delta u_a - c_a \partial_1 u_a = h(y, u_a) & \text{in } (-a, a) \times \mathbb{R}^{N-1} \\ u_a(-a, y) = V(y), \quad u_a(+a, y) = 0, \end{cases} \quad (54)$$

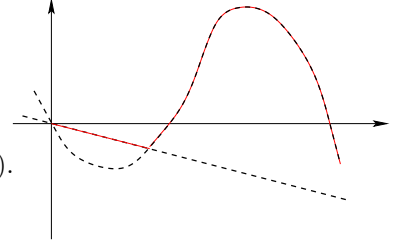
that verifies

$$\sup_{y \in \mathbb{R}^{N-1}} u_a(0, y) = \theta \quad (55)$$

for $a \geq 1$ and where θ is the unstable 0 of $f(u)$, that is $f(\theta) = 0$ and $0 < \theta < 1$. We recall that c_a is uniquely determined by the renormalization condition (55).

Let $\tilde{f}(u) = \max\{f(u), -m(u)\}$. Note that \tilde{f} itself is bistable:

$$\tilde{f}(0) = \tilde{f}(\theta) = \tilde{f}(1), \quad \tilde{f}(s) < 0 \text{ in } (0, \theta), \quad \tilde{f}(s) > 0 \text{ in } (\theta, 1).$$



We denote by z_a^c the solution of

$$\begin{cases} -z'' - cz' = \tilde{f}(z) \\ z(-a) = 1, \quad z(+a) = 0. \end{cases} \quad (56)$$

The function $z_a^{c_a}$ is a supersolution of (54) thus $u_a \leq z_a^{c_a}$. In view of (55) this implies that $z_a^{c_a}(0) \geq \theta$ and this implies that $c_a \leq \gamma_a$ where γ_a is the unique value of c such that the solution of (56) verifies $z_a^{\gamma_a}(0) = \theta$. This as before yields the upper bound for c_a .

The lower bound is achieved in the same manner as in the section 8 and the convergence for $a \rightarrow +\infty$ also.

Acknowledgments

This study was funded by the European Research Council under the European Union's Seventh Framework Programme (FP/2007-2013)/ERC Grant Agreement n.321186 - ReaDi - Reaction-Diffusion Equations, Propagation and Modelling. Henri Berestycki was also partially supported by an NSF FRG grant DMS-1065979. Part of this work was carried out while he was visiting the Department of Mathematics at the University of Chicago.

We wish to thank Laurent Desvillettes for bringing to our attention the population genetics model described in equations (2)-(3).

References

- [1] Matthieu Alfaro, Jérôme Coville, and Gaël Raoul. Travelling waves in a non-local reaction-diffusion equation as a model for a population structured by a space variable and a phenotypical trait. *arXiv:1211.3228*, 2013.
- [2] Tobias Back, Jochen G. Hirsch, Kristina Szabo, and Achim Gass. Failure to demonstrate peri-infarct depolarizations by repetitive mr diffusion imaging in acute human stroke. *Stroke*, 31(12):2901–2906, 2000.
- [3] H. Berestycki and P.-L. Lions. Une méthode locale pour l’existence de solutions positives de problèmes semi-linéaires elliptiques dans \mathbf{R}^N . *J. Analyse Math.*, 38:144–187, 1980.
- [4] H. Berestycki and L. Nirenberg. On the method of moving planes and the sliding method. *Bol. Soc. Brasil. Mat. (N.S.)*, 22(1):1–37, 1991.
- [5] Henri Berestycki. Le nombre de solutions de certains problèmes semi-linéaires elliptiques. *J. Funct. Anal.*, 40(1):1–29, 1981.
- [6] Henri Berestycki and Guillemette Chapuisat. A numerical study of a non local reaction-diffusion equation in population dynamics. In preparation., 2013.
- [7] Henri Berestycki and Guillemette Chapuisat. Propagation in a non homogeneous kpp equation arising in cancer modeling. In preparation, 2013.
- [8] Henri Berestycki, François Hamel, and Hiroshi Matano. Bistable traveling waves passing an obstacle. *Comm. Pures and Appl. Math.*, 62(6):729–788, 2009.
- [9] Henri Berestycki and François Hamel. Fronts and invasions in general domains. *C. R. Math. Acad. Sci. Paris*, 343(11-12):711–716, 2006.
- [10] Henri Berestycki and François Hamel. Generalized travelling waves for reaction-diffusion equations. In *Perspectives in nonlinear partial differential equations*, volume 446 of *Contemp. Math.*, pages 101–123. Amer. Math. Soc., Providence, RI, 2007.
- [11] Henri Berestycki and François Hamel. Generalized transition waves and their properties. *Comm. Pure Appl. Math.*, to appear, 2012.
- [12] Henri Berestycki and Pierre-Louis Lions. Some applications of the method of super and subsolutions. In *Bifurcation and nonlinear eigenvalue problems (Proc., Session, Univ. Paris XIII, Villetaneuse, 1978)*, volume 782 of *Lecture Notes in Math.*, pages 16–41. Springer, Berlin, 1980.
- [13] Henri Berestycki, Grégoire Nadin, Benoit Perthame, and Lenya Ryzhik. The non-local Fisher-KPP equation: travelling waves and steady states. *Nonlinearity*, 22(12):2813–2844, 2009.
- [14] Henri Berestycki and Louis Nirenberg. Travelling fronts in cylinders. *Ann. Inst. H. Poincaré Anal. Non Linéaire*, 9(5):497–572, 1992.

- [15] Henri Berestycki and Luca Rossi. On the principal eigenvalue of elliptic operators in \mathbb{R}^N and applications. *J. Eur. Math. Soc. (JEMS)*, 8(2):195–215, 2006.
- [16] Henri Berestycki and Luca Rossi. Reaction-diffusion equations for population dynamics with forced speed. I. The case of the whole space. *Discrete Contin. Dyn. Syst.*, 21(1):41–67, 2008.
- [17] Guillemette Chapuisat. Existence and nonexistence of curved front solution of a biological equation. *J. Differential Equations*, 236(1):237–279, 2007.
- [18] Guillemette Chapuisat and Emmanuel Grenier. Existence and nonexistence of traveling wave solutions for a bistable reaction-diffusion equation in an infinite cylinder whose diameter is suddenly increased. *Comm. Partial Differential Equations*, 30(10-12):1805–1816, 2005.
- [19] Guillemette Chapuisat, Emmanuel Grenier, Marie-Aimée Dronne, Marc Hommel, and Jean-Pierre Boissel. A global model of ischemic stroke with stress on spreading depression. *Progress in Biophysics and Molecular Biology*, 97(1):4–27, 2008.
- [20] Guillemette Chapuisat and Romain Joly. Asymptotic profiles for a traveling front solution of a biological equation. *Math. Models Methods Appl. Sci.*, 21(10):2155–2177, 2011.
- [21] Jacques De Keyser, Geert Sulter, and Paul G. Luiten. Clinical trials with neuroprotective drugs in acute ischaemic stroke: are we doing the right thing? *Trends Neurosci.*, 22(12):535–540, 1999.
- [22] Laurent Desvillettes, Régis Ferrières, and Céline Prevost. Infinite dimensional reaction-diffusion for population dynamics. Prépublication du CMLA No. 2003-04, 2003.
- [23] Lawrence C. Evans. *Partial differential equations*, volume 19 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, 1998.
- [24] Laurent Excoffier, Matthieu Foll, and Rémy J Petit. Genetic consequences of range expansions. *Annual Review of Ecology Evolution and Systematics*, 40:481–501, 2009.
- [25] Ido Filin, Robert D. Holt, and Michael Barfield. The relation of density regulation to habitat specialization, evolution of a species range, and the dynamics of biological invasions. *Am. Nat.*, 172(2):233–247, 2008.
- [26] Ali Gorji, Dieter Scheller, Heidrun Straub, Frank Tegtmeier, Rüdiger Köhling, Jörg Michael Höhling, Ingrid Tuxhorn, Alois Ebner, Peter Wolf, Hans Werner Panneck, Falk Oppel, and Erwin Josef Speckmann. Spreading depression in human neocortical slices. *Brain Res.*, 906(1-2):74–83, 2001.
- [27] Oskar Hallatschek, Pascal Hersen, Sharad Ramanathan, and David R. Nelson. Genetic drift at expanding frontiers promotes gene segregation. *Proceedings of the National Academy of Sciences*, 104(50):19926–19930, 2007.

- [28] Oskar Hallatschek and David R. Nelson. Gene surfing in expanding populations. *Theoretical Population Biology*, 73:158–170, 2008.
- [29] Oskar Hallatschek and David R. Nelson. Life at the front of an expanding population. *Evolution*, 64(1):193–206, 2010.
- [30] Robert D. Holt, Michael Barfield, Ido Filin, and Samantha Forde. Predation and the evolutionary dynamics of species ranges. *Am. Nat.*, 178(4):488–500, 2011.
- [31] Mark Kirkpatrick and Nicholas H. Barton. Evolution of a species’ range. *Am. Nat.*, 150(1):1–23, 1997.
- [32] Jean-François Mallordy and Jean-Michel Roquejoffre. A parabolic equation of the KPP type in higher dimensions. *SIAM J. Math. Anal.*, 26(1):1–20, 1995.
- [33] Hiroshi Matano. Asymptotic behavior and stability of solutions of semilinear diffusion equations. *Publ. Res. Inst. Math. Sci.*, 15(2):401–454, 1979.
- [34] Hiroshi Matano. Traveling waves in spatially inhomogeneous diffusive media – the non-periodic case. Preprint, 2012.
- [35] Avraham Mayevsky, Avi Doron, Tamar Manor, Sigal Meilin, Nili Zarchin, and George E. Ouaknine. Cortical spreading depression recorded from the human brain using a multiparametric monitoring system. *Brain Res.*, 740(1-2):268–274, 1996.
- [36] Maiken Nedergaard, Arthur J. Cooper, and Steven A. Goldman. Gap junctions are required for the propagation of spreading depression. *J. Neurobiol.*, 28(4):433–444, 1995.
- [37] James Nolen, Jean-Michel Roquejoffre, Lenya Ryzhik, and Andrej Zlatoš. Existence and non-existence of fisher-kpp transition fronts. *Arch. Ration. Mech. Anal.*, to appear, 2012.
- [38] L. A. Peletier and James Serrin. Uniqueness of nonnegative solutions of semilinear equations in \mathbf{R}^n . *J. Differential Equations*, 61(3):380–397, 1986.
- [39] Cristina Poggi, Ayman Moussa, Florence Hubert, and Guillemette Chapuisat. Numerical study of the stopping of aura during migraine. In *CEMRACS 2009: Mathematical modelling in medicine*, volume 30 of *ESAIM Proc.*, pages 44–52. EDP Sci., Les Ulis, 2010.
- [40] Paul H. Rabinowitz. Pairs of positive solutions of nonlinear elliptic partial differential equations. *Indiana Univ. Math. J.*, 23:173–186, 1973/74.
- [41] Lionel Roques, Jimmy Garnier, François Hamel, and Etienne K. Klein. Allee effect promotes diversity in traveling waves of colonization. *Proceedings of the National Academy of Sciences*, 109(23):8828–8833, 2012.
- [42] Laurent Schwartz. *Analyse hilbertienne*. Collection Méthodes. Hermann, Paris, 1979.
- [43] Bruno Shapiro. Osmotic forces and gap junctions in spreading depression: a computational model. *J. Comput. Neurosci.*, 10(1):99–120, 2001.

- [44] Wenxian Shen. Dynamical systems and traveling waves in almost periodic structures. *J. Differential Equations*, 169(2):493–548, 2001. Special issue in celebration of Jack K. Hale’s 70th birthday, Part 4 (Atlanta, GA/Lisbon, 1998).
- [45] George G. Somjen. *Ions in the Brain: Normal Function, Seizures, and Stroke*. Oxford University Press, New York, 2004.
- [46] Anthony J. Strong, Martin Fabricius, Martyn G. Boutelle, Stuart J. Hibbins, Sarah E. Hopwood, Robina Jones, Mark C. Parkin, and Martin Lauritzen. Spreading and synchronous depressions of cortical activity in acutely injured human brain. *Stroke*, 33(12):2738–2743, 2002.
- [47] Henry C. Tuckwell. Predictions and properties of a model of potassium and calcium ion movements during spreading cortical depression. *Int. J. Neurosci.*, 10(2-3):145–164, 1980.
- [48] José M. Vega. Travelling wavefronts of reaction-diffusion equations in cylindrical domains. *Comm. Partial Differential Equations*, 18(3-4):505–531, 1993.
- [49] Marcel O. Vlad, L. Luca Cavalli-Sforza, and John Ross. Enhanced (hydrodynamic) transport induced by population growth in reactiondiffusion systems with application to population genetics. *Proceedings of the National Academy of Sciences*, 101(28):10249–10253, 2004.
- [50] Andrej Zlatoš. Generalized traveling waves in disordered media: existence, uniqueness, and stability. Preprint, 2012.