



Some remarks on the L^2 -critical nonlinear Schrödinger Equation with a nonlinear damping.

Mohamad Darwich

► To cite this version:

Mohamad Darwich. Some remarks on the L^2 -critical nonlinear Schrödinger Equation with a nonlinear damping.. 2012. hal-00712285v2

HAL Id: hal-00712285

<https://hal.science/hal-00712285v2>

Preprint submitted on 1 Jul 2012 (v2), last revised 15 Jan 2013 (v4)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

SOME REMARKS ON THE L^2 -CRITICAL NONLINEAR SCHRÖDINGER EQUATION WITH A NONLINEAR DAMPING.

DARWICH MOHAMAD

ABSTRACT. We consider the Cauchy problem for the L^2 -critical nonlinear Schrödinger equation with a nonlinear damping. According to the power of the damping term, we prove the global existence or the existence of finite time blowup dynamics with the log-log blow-up speed for $\|\nabla u(t)\|_{L^2}$.

1. INTRODUCTION

In this paper, we study the blowup and the global existence of solutions for the focusing NLS equation with a nonlinear damping:

$$\begin{cases} iu_t + \Delta u + |u|^{\frac{4}{d}}u + ia|u|^p u = 0, (t, x) \in [0, \infty[\times \mathbb{R}^d, d = 1, 2, 3, 4. \\ u(0) = u_0 \in H^1(\mathbb{R}^d) \end{cases} \quad (1.1)$$

with initial data $u(0) = u_0 \in H^1(\mathbb{R}^d)$ where $a > 0$ is the coefficient of friction and $p \in N^*$. Note that if we replace $+|u|^{\frac{4}{d}}u$ by $-|u|^{\frac{4}{d}}u$, (1.1) becomes the defocusing NLS equation.

Equation (1.1) arises in various areas of nonlinear optics, plasma physics and fluid mechanics. Fibich [7] noted that in the nonlinear optics context, the origin of nonlinear damping is multiphoton absorption. For example, in the case of solids the number p corresponds to the number of photons it takes to make a transition from the valence band to the conduction band. Similar behavior can occur with free atoms, in this case p corresponds to the number of photons needed to make a transition from the ground state to some excited state or to the continuum.

The Cauchy problem for (1.1) was studied by Kato [11] and Cazenave [3] and it is known that if $p < \frac{4}{d-2}$, then the problem is locally well-posed in $H^1(\mathbb{R}^d)$: For any $u_0 \in H^1(\mathbb{R}^d)$, there exist $T \in (0, \infty]$ and a unique solution $u(t)$ of (1.1) with $u(0) = u_0$ such that $u \in C([0, T]; H^1(\mathbb{R}^d))$. Moreover, T is the maximal existence time of the solution $u(t)$ in the sense that if $T < \infty$ then $\lim_{t \rightarrow T} \|u(t)\|_{H^1(\mathbb{R}^d)} = \infty$.

Let us notice that for $a = 0$ (1.1) becomes the L^2 -critical nonlinear Schrödinger equation:

$$\begin{cases} iu_t + \Delta u + |u|^{\frac{4}{d}}u = 0 \\ u(0) = u_0 \in H^1(\mathbb{R}^d) \end{cases} \quad (1.2)$$

Key words and phrases. Damped Nonlinear Schrödinger Equation, Blow-up, Global existence.

For $u_0 \in H^1$, a sharp criterion for global existence for (1.2) has been exhibited by Weinstein [24]:

For $\|u_0\|_{L^2} < \|Q\|_{L^2}$, the solution of (1.2) is global in H^1 . This follows from the conservation of the energy and the L^2 norm and the sharp Gagliardo-Nirenberg inequality:

$$\forall u \in H^1, E(u) \geq \frac{1}{2} \left(\int |\nabla u|^2 \right) \left(1 - \left(\frac{\int |u|^2}{\int |Q|^2} \right)^{\frac{2}{d}} \right).$$

On the other hand, there exists explicit solutions with $\|u_0\|_{L^2} = \|Q\|_{L^2}$ that blow up in finite time.

In the series of papers [14, 22], Merle and Raphael have studied the blowup for (1.2) with $\|Q\|_{L^2} < \|u_0\|_{L^2} < \|Q\|_{L^2} + \delta$, δ small and have proven the existence of the blowup regime corresponding to the log-log law:

$$\|u(t)\|_{H^1(\mathbb{R}^d)} \sim \left(\frac{\log |\log(T-t)|}{T-t} \right)^{\frac{1}{2}}. \quad (1.3)$$

Darwich [6] has proved in case of the linear damping ($p = 0$), the global existence in H^1 for $\|u_0\|_{L^2} \leq \|Q\|_{L^2}$, and has showed that the log-log regime is stable by perturbations and the solutions blows up in finite time with the log-log law.

Numerical observations suggest that this finite time blowup phenomena persists in the case of the nonlinear damping for $p < \frac{4}{d}$ (see Fibich [7] and [20]). Passot and Sulem [20] have proved that the solutions are global in $H^1(\mathbb{R}^2)$ for $p > 2$. Therefore, it is interesting now to study for which values of p , we have either the blow-up or the global existence of solutions, it is our aim in this paper and we will show that:

If $p \geq \frac{4}{d}$, then the solution of (1.1) is global in H^1 .

The global existence for small data in H^1 , for $p \leq 2$.

The existence of finite time blowup solutions.

More precisely, we have the following theorem:

Theorem 1.1. *Let u_0 in $H^1(\mathbb{R}^d)$ with $d = 1, 2, 3, 4$:*

- (1) *if $\frac{4}{d-2} > p \geq \frac{4}{d}$, then the solution of (1.1) is global in H^1 .*
- (2) *For $p = 1$ and $d = 1, 2, 3$ or $p = 2$ and $d = 2$, the solution is global in H^1 , for $u_0 \in H^1$ such that $\|u_0\|_{L^2} < \alpha$ with $\alpha > 0$ small enough.*
- (3) *let $p < \frac{4}{d}$, then there exists $\delta_0 > 0$ such that $\forall a > 0$ and $\forall \delta \in]0, \delta_0[$, there exists $u_0 \in H^1$ with $\|u_0\|_{L^2} = \|Q\|_{L^2} + \delta$, such that the solution of (1.1) blows up in finite time in the log-log regime.*

In the case $p = \frac{4}{d}$, we have the following theorem:

Theorem 1.2. *Let $p = \frac{4}{d}$, then the initial-value problem (1.1) is globally well posed in $H^s(\mathbb{R}^d)$, $s \geq 0$.*

Remark 1.1. *Theorem 1.2 can be given an answer for an open problem for Sparber and Antonelli in their paper [1] and for Fibich and Klein in their paper [8], at least for the L^2 -critical case.*

Remark 1.2. *Theorem 1.2 and part (1) and (2) of Theorem 1.1 still true in the defocusing case.*

In this paper we will prove the existence of log-log explosive solutions in the case $p < \frac{4}{d}$. We establish the global existence of the solutions as soon as $\frac{4}{d} \leq p < \frac{4}{d-2}$, where $\frac{4}{d-2}$ is the H^1 -critical exponent.

Acknowledgments. I would like to thank my advisor prof Luc Molinet for his rigorous attention to this paper, and Dr. Christof Sparber for his remarks.

2. GLOBAL EXISTENCE.

In this section, we prove assertion (1) and (2) of Theorem 1.1 and Theorem 1.2. To prove part (1), we will prove that the H^1 -norm of u is bounded for any time. To prove part (2), we use generalised Gagliardo-Nirenberg inequality to show that the energy is non increasing. Finally to prove Theorem 1.2, we establish an a priori estimate on the critical Strichartz norm.

Theorem 2.1. *Let $p \leq \frac{4}{d-2}$ (for $d > 2$), then the initial-value problem (1.1) is locally well posed in $H^1(\mathbb{R}^d)$ (If $p < \frac{4}{d-2}$ the minimal time of the existence depends on $\|u_0\|_{H^1}$).*

Proof: See [3] page 93 Theorem 4.4.1.

To prove the following proposition, we will proceed in the same way as in the section 3.1 in [20].

Proposition 2.1. *Let u be a solution of (1.1) and $\frac{4}{d-2} \geq p > \frac{4}{d}$ then*

$$\|\nabla u(t)\|_{L^2(\mathbb{R}^d)} \leq \|\nabla u(0)\|_{L^2(\mathbb{R}^d)} e^{a^{(\frac{-4t}{pd-4})}}.$$

Proof: Multiply Eq. (1.1) by $\Delta \bar{u}$, integrate and take the imaginary part, this gives

$$\frac{1}{2} \frac{d}{dt} \int |\nabla u|^2 dx + a \int |u|^p |\nabla u|^2 + a \Re \int u \nabla |u|^p \nabla \bar{u} dx = -\frac{4}{d} \Im \int u \nabla \bar{u} \Re(u \nabla u) |u|^{\frac{4}{d}-2}. \quad (2.1)$$

In the l.h.s, a simple calculation shows that the third term rewrites in the form $\frac{p}{4} \int |u|^{p-2} (\nabla |u|^2)^2$. Equation (2.1) becomes:

$$\frac{1}{2} \frac{d}{dt} \int |\nabla u|^2 dx + a \int |u|^p |\nabla u|^2 + a \frac{p}{4} \int |u|^{p-2} (\nabla |u|^2)^2 \leq \frac{2}{d} \int |u|^{\frac{4}{d}} |\nabla u|^2. \quad (2.2)$$

To estimate the r.h.s of (2.2), we rewrite it as ($p > \frac{4}{d}$)

$$\int |u|^{\frac{4}{d}} |\nabla u|^2 = \int |u|^{\frac{4}{d}} |\nabla u|^{\frac{8}{pd}} |\nabla u|^{2-\frac{8}{pd}}.$$

Now by Hölder inequality we obtain that

$$\int |u|^{\frac{4}{d}} |\nabla u|^2 \leq \left(\int |u|^p |\nabla u|^2 \right)^{\frac{4}{pd}} \left(\int |\nabla u|^2 \right)^{1-\frac{4}{pd}}.$$

Then inequality (2.2) takes the form:

$$\frac{d}{dt} w(t) + 2av(t) \leq \frac{4}{d} v(t)^{\frac{4}{pd}} w(t)^{1-\frac{4}{pd}}.$$

where $w(t) = \int |\nabla u|^2$ and $v(t) = \int |u|^p |\nabla u|^2$.

Using Young's inequality $ab \leq \epsilon a^q + C\epsilon^{-\frac{1}{q-1}} b^{q'}$, $\frac{1}{q} + \frac{1}{q'} = 1$, with $q = \frac{pd}{4}$ and $\epsilon = \frac{ad}{2}$ we obtain :

$$\frac{d}{dt}w(t) \leq a^{-\frac{1}{\frac{pd}{4}-1}} w(t).$$

This ensures that:

$$w(t) \leq w(0)e^{a\left(-\frac{4t}{pd-4}\right)}. \square$$

This show that, the H^1 -norm of u is bounded for any time and gives directly the proof of part one of Theorem 1.1.

Now we will prove the global existence for small data, for this we will use the following lemma, for the proof see [9]:

Lemma 2.1. *Let q, r be any real numbers satisfying $1 \leq q, r \leq \infty$, and let j, m be any integers satisfying $0 \leq j < m$. If u is any functions in $C_0^m(R^n)$, then*

$$|D^j u|_{L^s} \leq C |D^m u|_r^a |u|_q^{1-a}$$

where

$$\frac{1}{s} = \frac{j}{n} + a\left(\frac{1}{r} - \frac{m}{n}\right) + (1-a)\frac{1}{q},$$

for all a in the interval

$$\frac{j}{m} \leq a \leq 1,$$

where C is a constant depending only on n, m, j, q, r and a .

Now we have the following one:

Lemma 2.2. *Let $p \leq 2$ and $v \in C_0^\infty(R^n)$ then:*

$$\int |v|^{\frac{4}{d}+2+p} \leq C \left(\int |\nabla(|v|^{\frac{p+2}{2}})|^2 \right) \times \left(\int |v|^2 \right)^{\frac{2}{d}}.$$

where $c > 0$ depending only on d and p .

Proof: Take $s = \frac{\frac{4}{d}+2+p}{1+\frac{p}{2}}$, $q = \frac{2}{1+\frac{p}{2}}$, $r = 2$, $j = 0$ and $m = 1$, then by Lemma 2.1 we obtain that:

$$|u|_{\frac{\frac{4}{d}+2+p}{1+\frac{p}{2}}}_L \leq C |\nabla u|_{L^2}^{\frac{4+2p}{\frac{8}{d}+4+2p}} |u|_{\frac{\frac{2}{1+\frac{p}{2}}}}^{\frac{\frac{8}{d}}{\frac{8}{d}+4+2p}}.$$

Now take $u = |v|^{1+\frac{p}{2}}$, we obtain our lemma.

Now we can prove the following proposition:

Proposition 2.2. *let $p \leq 2$. There exists $0 < \alpha = \alpha(p, d) < \|Q\|_{L^2}$, such that for any $u_0 \in H^1$ with $\|u_0\|_{L^2} < \alpha$, it holds*

$$\frac{d}{dt}E(u(t)) \leq 0, \quad \forall t > 0.$$

Proof: It is a easy to prove that :

$$\frac{d}{dt}E(u(t)) = a\left(\int |u|^{\frac{4}{d}+p+2} - C_p \int |\nabla(|u|^{\frac{p+2}{2}})|^2\right),$$

then by Lemma 2.2 we obtain that:

$$\frac{d}{dt}E(u(t)) \leq a\left(\int |\nabla(|u|^{\frac{p+2}{2}})|^2\right)(C\left(\int |u|^2\right)^{\frac{2}{d}} - C_p)$$

Choosing $\alpha^{\frac{2}{d}} < \frac{C_p}{C}$, and using (3.2) we get the result.

Now we will prove part (2) of Theorem 1.1: We have by the sharp Gagliardo-Nirenberg inequality :

$$\forall u \in H^1, E(u) \geq \frac{1}{2}\left(\int |\nabla u|^2\right)\left(1 - \left(\frac{\int |u|^2}{\int |Q|^2}\right)^{\frac{2}{d}}\right).$$

Note that $E(u) \leq E(u_0)$, then by this inequality we can found a bound for the H^1 -norm and we obtain the global existence.

2.1. Critical case ($p = \frac{4}{d}$). Now we will treat the critical case. First let us prove that, if the solution blows up in finite time T , then $\|u\|_{L^{\frac{4}{d}+2}([0,T];L^{\frac{4}{d}+2}(R^d))} = +\infty$.

Proposition 2.3. *Let u be the unique maximal solution of (1.1) in $[0, T^*)$; if $T^* < \infty$, then $\|u\|_{L^\sigma([0,T],L^\sigma)} = \infty$ where $\sigma = \frac{4}{d} + 2$.*

For this we will need the following proposition:

Proposition 2.4. *There exists $\delta > 0$ with the following property. If $u_0 \in L^2(R^d)$ and $T \in (0, \infty]$ are such that $\|S(\cdot)u_0\|_{L^\sigma([0,T],L^\sigma)} < \delta$, there exists a unique solution $u \in C([0, T], L^2(R^d)) \cap L^\sigma([0, T], L^\sigma(R^n))$ of (1.1). In addition, $u \in L^q([0, T], L^r(R^d))$ for every admissible pair (q, r) ; for $t \in [0, T]$. Finally, u depends continuously in $C([0, T], L^2(R^d)) \cap L^\sigma([0, T], L^\sigma(R^d))$ on $u_0 \in L^2(R^n)$. If $u_0 \in H^1(R^d)$, then $u \in C([0, T], H^1(R^d))$.*

See [4] for the proof.

We need also the following lemma (see [4]):

Lemma 2.3. *Let $T \in (0, \infty]$, let $\sigma = \frac{4}{d} + 2$, and let (q, r) be an admissible pair. Then, whenever $u \in L^\sigma([0, T], L^\sigma(R^d))$, it follows that $F(u) = -i \int_0^t S(t-s)(|u|^{\frac{4}{d}}u + ia|u|^p u)ds \in C([0, T], H^{-1}(R^d)) \cap L^q(0, T, L^r(R^d))$. Furthermore, there exists K , independent of T , such that*

$$\|Fv - Fu\|_{L^q([0,T],L^r)} < K(\|u\|_{L^\sigma([0,T],L^\sigma(R^d))}^{\frac{4}{d}} + \|v\|_{L^\sigma([0,T],L^\sigma(R^d))}^{\frac{4}{d}})\|u-v\|_{L^\sigma([0,T],L^\sigma(R^d))} \quad (2.3)$$

for every $u, v \in L^\sigma([0, T], L^\sigma(R^d))$.

Proof of proof 2.3:

Let $u_0 \in L^2(R^d)$. Observe that $\|S(\cdot)u_0\|_{L^\sigma(0,T,L^\sigma)} \rightarrow 0$, as $T \rightarrow 0$. Thus for sufficiently small T , the hypotheses of Proposition 2.4 are satisfied. Applying iteratively this proposition, we can construct the maximal solution

$u \in C([0, T^*), L^2(\mathbb{R}^d)) \cap L^\sigma([0, T^*), L^\sigma(\mathbb{R}^d))$ of (1.1). We proceed by contradiction, assuming that $T^* < \infty$, and $\|u\|_{L^\sigma([0, T], L^\sigma)} < \infty$. Let $t \in [0, T^*)$. For every $s \in [0, T^* - t]$, we have

$$S(s)u(t) = u(t + s) - F(u(t + \cdot))(s).$$

From (2.3), we obtain

$$\|S(\cdot)u(t)\|_{L^\sigma([0, T^* - t], L^\sigma(\mathbb{R}^d))} \leq \|u\|_{L^\sigma([t, T^*], L^\sigma)} + K(\|u\|_{L^\sigma([t, T^*], L^\sigma)})^{\frac{4}{d}+1}$$

Therefore, for t fixed close enough to T^* , it follows that

$$\|S(\cdot)u(t)\|_{L^\sigma([0, T^* - t], L^\sigma(\mathbb{R}^d))} \leq \delta.$$

Applying Proposition 2.4, we find that u can be extended after T^* , which contradicts the maximality.

Corollary 2.1. *For $p = \frac{4}{d}$, the solution of (1.1) is global.*

Proof Multiply equation (1.1) by $\bar{u}$, and take the imaginary part to obtain:

$$\frac{d}{dt}\|u(t)\|_{L^2}^2 + 2a\|u\|_{L^{\frac{4}{d}+2}}^{\frac{4}{d}+2} = 0.$$

Hence $\forall t \in \mathbb{R}_+$

$$\|u\|_{L^{\frac{4}{d}+2}[0, t], L^{\frac{4}{d}+2}(\mathbb{R}^d)} \leq \frac{1}{2a}\|u_0\|_{L^2}^2.$$

Theorem 1.2 follows then directly from Proposition 2.3.

3. BLOW UP SOLUTION.

Now we will prove the existence of the explosive solutions in the regime log-log.

We look for a solution of (1.1) such that for t close enough to blowup time, we shall have the following decomposition:

$$u(t, x) = \frac{1}{\lambda^{\frac{d}{2}}(t)}(Q_{b(t)} + \epsilon)(t, \frac{x - x(t)}{\lambda(t)})e^{i\gamma(t)}, \quad (3.1)$$

for some geometrical parameters $(b(t), \lambda(t), x(t), \gamma(t)) \in (0, \infty) \times (0, \infty) \times \mathbb{R}^d \times \mathbb{R}$, here $\lambda(t) \sim \frac{1}{\|\nabla u(t)\|_{L^2}}$, and the profiles Q_b are suitable deformations of Q related to some extra degeneracy of the problem.

Now we define the following quantities:

$$L^2\text{-norm} : \|u(t, x)\|_{L^2} = \int |u(t, x)|^2 dx.$$

$$\text{Energy} : E(u(t, x)) = \frac{1}{2}\|\nabla u\|_{L^2}^2 - \frac{d}{4+2d}\|u\|_{L^{\frac{4}{d}+2}}^{\frac{4}{d}+2}.$$

$$\text{Kinetic momentum} : P(u(t)) = \text{Im}(\int \nabla u \bar{u}(t, x)).$$

Remark 3.1. *It is easy to prove that if u is a solution of (1.1) then:*

$$\frac{d}{dt}\|u(t)\|_{L^2} = -2a \int |u|^{p+2}, t \in [0, T), \quad (3.2)$$

$$\frac{d}{dt}E(u(t)) = -a(\|u^{\frac{p}{2}}\nabla u\|_{L^2}^2 - \|u\|_{L^{\frac{4}{\frac{4}{d}+2+p}}}^{\frac{4}{d}+2+p}) \quad (3.3)$$

and

$$\frac{d}{dt}P(u(t)) = -2a \operatorname{Im} \int \bar{u}|u|^p \nabla u, t \in [0, T]. \quad (3.4)$$

Now we have the following theorem:

Theorem 3.1. *There exist a set of initial data Ω in H^1 , such that for $0 < a < a_0$ with a_0 small enough, the corresponding solution $u(t)$ of (1.1) blows up in finite time in the log-log regime.*

The set of initial data Ω is open in H^1 , using the continuity with regard to the initial data and the parameters, we can prove the following corollary:

Corollary 3.1. *Let $u_0 \in H^1$ be an initial data such that the corresponding solution $u(t)$ of (1.2) blows up in the loglog regime. There exist $\beta_0 > 0$ and $a_0 > 0$ such that if $v_0 = u_0 + h_0$, $\|h_0\|_{H^1} \leq \beta_0$ and $a \leq a_0$, the solution $v(t)$ for (1.1) with the initial data v_0 blowup in finite time.*

Proof of Corollary 3.1: Let $S(t)$ be the propagator for the linear equation:

$$i\partial_t u + \Delta u = 0, \quad (t, x) \in [0, \infty[\times \mathbb{R}^d.$$

The Cauchy problem for (1.1) with $u(0) = u_0 \in H^1(\mathbb{R}^d)$ is equivalent to the integral equation:

$$u(t) = S(t)u_0 + i \int_0^t S(t-s)(|u(s)|^{\frac{4}{d}}u(s) + ia|u(s)|^p u)ds.$$

We know from Lemma 3.2: there exist $T(\|u_0\|_{H^1(\mathbb{R}^d)}) > 0$ such that:

$$\forall 0 \leq a \leq 1, \|u\|_{L^\infty([0,T];H^1)} \leq 2\|u_0\|_{H^1(\mathbb{R}^d)}.$$

Let u a solution for (1.1) and v solution for (1.2) we have:

$$u-v = S(t)(u_0-v_0) + i \int_0^t S(t-s)(|u(s)|^{\frac{4}{d}}u(s) - |v(s)|^{\frac{4}{d}}v(s))ds + ia \int_0^t S(t-t')u(t')|u(t')|^p dt'.$$

By Strichartz we obtain (see the proof of Lemma 3.2):

$$\begin{aligned} \|u-v\|_{L^\infty([0,T];H^1)} &\leq \|u_0-v_0\|_{H^1(\mathbb{R}^d)} \\ &+ CT^\gamma (\|u\|_{L^\infty([0,T];H^1)}^{\frac{4}{d}} + \|v\|_{L^\infty([0,T];H^1)}^{\frac{4}{d}}) \|u-v\|_{L^\infty([0,T];H^1)} \\ &+ CaT \|u\|_{L^\infty([0,T];H^1)}^{p+1}. \end{aligned}$$

Thus for $T_1 = \min(T(\|u_0\|_{H^1(\mathbb{R}^d)}), T(\|v_0\|_{H^1(\mathbb{R}^d)}))$ we obtain $\forall 0 \leq t \leq T_1$:

$$\begin{aligned} \|u-v\|_{L^\infty([0,T];H^1)} &\leq \|u_0-v_0\|_{H^1(\mathbb{R}^d)} \\ &+ CT^\gamma (\|u_0\|_{H^1}^{\frac{4}{d}} + \|v_0\|_{H^1}^{\frac{4}{d}}) \|u-v\|_{L^\infty([0,T];H^1)} \\ &+ CaT \|u\|_{L^\infty([0,T];H^1)}^{p+1}. \end{aligned}$$

Now for $T_2 = \frac{1}{2} \min(\max^{-\frac{1}{\gamma}}(\|u_0\|_{H^1}^p, \|v_0\|_{H^1}^p), T_1)$:

$$\begin{aligned} \|u-v\|_{L^\infty([0,t];H^1)} &\leq \|u_0-v_0\|_{H^1(\mathbb{R}^d)} + \frac{1}{2} \|u-v\|_{L^\infty([0,t];H^1)} \\ &+ a, \end{aligned}$$

thus

$$\|u - v\|_{L^\infty([0,t];H^1)} \leq C\|u_0 - v_0\|_{H^1(\mathbb{R}^d)} + a. \quad \forall 0 < t < T_2$$

Thus the map $(a, \phi) \rightarrow u(\cdot, a, \phi)$ is continuous in $(0, u_0)$ from $\mathbb{R} \times H^1(\mathbb{R}^d)$ to $C([0, T_2], H^1(\mathbb{R}^d))$. Since T_2 only depends on $\|u_0\|_{H^1(\mathbb{R}^d)}$, this continuity extends to any interval $[0, T]$ in the maximal interval of existence of u .

Note that we will abbreviate our proof because it is very very close to the case of linear damping ($p = 0$ see Darwich [6]). Actually, as noticed in [21], we only need to prove that in the log-log regime the L^2 norm does not grow, and the growth of the energy (resp the momentum) is below $\frac{1}{\lambda^2}$ (resp $\frac{1}{\lambda}$). In this paper, we will prove that in the log-log regime, the growth of the energy and the momentum are bounded by:

$$E(u(t)) \lesssim \log(\lambda(t))\lambda(t)^{-\frac{pd}{2}}, \quad P(u(t)) \lesssim \log(\lambda(t))\lambda(t)^{1-\frac{pd}{4}}.$$

Let us recall that a function $u : [0, T] \mapsto H^1$ follows the log-log regime if the following uniform controls on the decomposition (3.1) hold on $[0, T]$:

- Control of $b(t)$

$$b(t) > 0, \quad b(t) < 10b(0). \quad (3.5)$$

- Control of λ :

$$\lambda(t) \leq e^{-e^{\frac{\pi}{100b(t)}}} \quad (3.6)$$

and the monotonicity of λ :

$$\lambda(t_2) \leq \frac{3}{2}\lambda(t_1), \quad \forall 0 \leq t_1 \leq t_2 \leq T. \quad (3.7)$$

Let $k_0 \leq k_+$ be integers and $T^+ \in [0, T]$ such that

$$\frac{1}{2^{k_0}} \leq \lambda(0) \leq \frac{1}{2^{k_0-1}}, \quad \frac{1}{2^{k_+}} \leq \lambda(T^+) \leq \frac{1}{2^{k_+-1}} \quad (3.8)$$

and for $k_0 \leq k \leq k_+$, let t_k be a time such that

$$\lambda(t_k) = \frac{1}{2^k}, \quad (3.9)$$

then we assume the control of the doubling time interval:

$$t_{k+1} - t_k \leq k\lambda^2(t_k). \quad (3.10)$$

- control of the excess of mass:

$$\int |\nabla \epsilon(t)|^2 + \int |\epsilon(t)|^2 e^{-|y|} \leq \Gamma_{b(t)}^{\frac{1}{4}}. \quad (3.11)$$

Remark 3.2. We will use a bootstrap argument, in fact the set of initial data Ω will be to initialize the bootstrap.

3.1. Control of the energy and the kinetic momentum in the log-log regime. We recall the Strichartz estimates. An ordered pair (q, r) is called admissible if $\frac{2}{q} + \frac{d}{r} = \frac{d}{2}$, $2 < q \leq \infty$. We define the Strichartz norm of functions $u : [0, T] \times \mathbb{R}^d \mapsto C$ by:

$$\|u\|_{S^0([0, T] \times \mathbb{R}^d)} = \sup_{(q, r) \text{ admissible}} \|u\|_{L_t^q L_x^r([0, T] \times \mathbb{R}^d)} \quad (3.12)$$

and

$$\|u\|_{S^1([0, T] \times \mathbb{R}^d)} = \sup_{(q, r) \text{ admissible}} \|\nabla u\|_{L_t^q L_x^r([0, T] \times \mathbb{R}^d)} \quad (3.13)$$

We will sometimes abbreviate $S^i([0, T] \times \mathbb{R}^2)$ with S_T^i or $S^i[0, T]$, $i = 1, 2$. Let us denote the Hölder dual exponent of q by q' so that $\frac{1}{q} + \frac{1}{q'} = 1$. The Strichartz estimates may be expressed as:

$$\|u\|_{S_T^0} \leq \|u_0\|_{L^2} + \|(i\partial_t + \Delta)u\|_{L_t^{q'} L_x^{r'}} \quad (3.14)$$

where (q, r) is any admissible pair. Now we will derive an estimate on the energy, to check that it remains small with respect to λ^{-2} :

Lemma 3.1. *Assuming that (3.5)-(3.11) hold, then the energy and kinetic momentum of the solution u to (1.1) are controlled on $[0, T]$ by:*

$$|E(u(t))| \leq C(\log(\lambda(t))\lambda(t)^{-\frac{pd}{4}}), \quad (3.15)$$

$$|P(u(t))| \leq C(\log(\lambda(t))\lambda(t)^{1-\frac{pd}{4}}). \quad (3.16)$$

To prove this lemma, we shall need the following one:

Lemma 3.2. *Let u be a solution of (1.1) emanating for u_0 in H^1 . Then $u \in C([0, \Delta T], H^1)$ where $\Delta T = \|u_0\|_{L^2}^{\frac{d-4}{d}} \|u_0\|_{H^1}^{-2}$, and we have the following control*

$$\|u\|_{S^0[t, t+\Delta T]} \leq 2\|u_0\|_{L^2}, \quad \|u\|_{S^1[t, t+\Delta T]} \leq 2\|u_0\|_{H^1(\mathbb{R}^d)}.$$

Proof of Lemma 3.1: According to (3.10) each interval $[t_k, t_{k+1}]$, can be divided into k intervals, $[\tau_k^j, \tau_k^{j+1}]$ such that the estimates of the previous lemma are true. From (3.3) and the Gagliardo-Nirenberg inequality, we obtain that:

$$\frac{d}{dt}E(u(t)) \leq \|u\|_{L^2}^{\frac{4}{d}+p-\frac{pd}{2}} \|\nabla u\|_{L^2}^{2+\frac{pd}{2}}$$

(3.2) Then

$$\int_{\tau_k^j}^{\tau_k^{j+1}} \frac{d}{dt}E(u(t))dt \leq C \int_{\tau_k^j}^{\tau_k^{j+1}} \|\nabla u(t)\|_{L^2}^{2+\frac{pd}{2}}$$

Then by Lemma 3.2, we obtain that:

$$\int_{\tau_k^j}^{\tau_k^{j+1}} \frac{d}{dt}E(u(t))dt \leq C(\tau_k^{j+1} - \tau_k^j)\lambda^{-2-\frac{pd}{2}}(\tau_k^j)$$

Note that $\tau_k^{j+1} - \tau_k^j \sim \lambda^2(\tau_k^j) \sim \lambda^2(t_k)$, then

$$\int_{\tau_k^j}^{\tau_k^{j+1}} \frac{d}{dt}E(u(t))dt \leq C\lambda^{-\frac{pd}{2}}(t_k)$$

Summing from $j = 1$ to $J_k \leq CK$, we obtain that:

$$\sum_{j=1}^{J_k} \int_{\tau_k^j}^{\tau_k^{j+1}} \frac{d}{dt} E(u(t)) dt \leq Ck\lambda^{-\frac{pd}{2}}(t_k)$$

Now taking $T^+ = T$ and summing from K_0 to K^+ , we obtain:

$$\int_0^{T^+} \frac{d}{dt} E(u(t)) dt \leq CK^+ \lambda^{-\frac{pd}{2}}(T^+) \lesssim C \log(\lambda(T)) \lambda^{-\frac{pd}{2}}(T).$$

Note that $\log(\lambda(T)) \lambda^{-\frac{pd}{2}}(T)$ is small with respect to $\frac{1}{\lambda^{\frac{p}{2}}}$ because $p < \frac{4}{d}$.

Now we prove **(3.16)**: From (3.4) we have :

$$|\frac{d}{dt} P(u(t))| \leq \int |\bar{u}| |\nabla u| |u|^p$$

By Gagliardo-Nirenberg inequality we have:

$$\|u\|_{L^{2p+2}}^{2p+2} \leq \|u\|_{L^2}^{2p+2-dp} \|\nabla u\|_{L^2}^{dp}$$

then

$$\frac{d}{dt} P(u(t)) \leq \left(\int |u|^{2(p+1)} \right)^{\frac{1}{2}} \|\nabla u\|_{L^2} \leq C \|\nabla u\|_{L^2}^{1+\frac{pd}{2}}.$$

Then:

$$\int_{\tau_k^j}^{\tau_k^{j+1}} \frac{d}{dt} P(u(t)) \leq C(\tau_k^{j+1} - \tau_k^j) \left\| \nabla u(\tau_k^j) \right\|_{L^2}^{1+\frac{pd}{2}} \leq C \|\nabla u(t_k)\|_{L^2}^{-1+\frac{pd}{2}}$$

Summing successively into j and k we obtain that:

$$\int_0^{T^+} \frac{d}{dt} P(u(t)) \lesssim \log(\lambda(T^+)) \lambda^{1-\frac{pd}{2}}(T^+).$$

Remark that this quantity is small with respect to $\frac{1}{\lambda}$ because $p < \frac{4}{d}$. $\square$

REFERENCES

- [1] P. Antonelli and C. Sparber. *Global well-posedness for cubic NLS with nonlinear damping*. Comm. Partial Differential Equations, 35 (2010) 48324845.
- [2] H. Berestycki and P.-L. Lions. *Nonlinear scalar field equations. II. Existence of infinitely many solutions*. Arch. Rational Mech. Anal., 82(1983):347–375.
- [3] T. Cazenave. *Semilinear Schrödinger equations*, volume 10 of *Courant Lecture Notes in Mathematics*. New York University Courant Institute of Mathematical Sciences, New York, 2003.
- [4] T. Cazenave and F. Weissler. *Some remarks on the nonlinear Schrödinger equation in the subcritical case*. New methods and results in nonlinear field equations (Bielefeld, 1987), 5969, Lecture Notes in Phys., 347, Springer, Berlin, 1989.
- [5] J. Colliander and P. Raphael. *Rough blowup solutions to the L^2 critical NLS*. Math. Ann., 345(2009):307–366.
- [6] M. Darwich. *Blowup for the Damped L^2 critical nonlinear Schrödinger equations*. Advances in Differential Equations. volume 17, Numbers 3-4 (2012), 337-367.
- [7] G. Fibich and F. Merle. *Self-focusing on bounded domains*. Phys. D, 155(2001):132–158.
- [8] G. Fibich and M. Klein. *Nonlinear-damping continuation of the nonlinear Schrödinger equation-a numerical study*. Physica D, 241 (2012), 519-527.
- [9] A. Friedman *Partial Differential Equations*.
- [10] T. Hmidi and S. Keraani. *Blowup theory for the critical nonlinear Schrödinger equations revisited*. Int. Math. Res. Not., 46(2005):2815–2828.

- [11] T. Kato. *On nonlinear Schrödinger equations* Ann. Inst. H. Poincaré Phys. Théor., 46(1987):113–129.
- [12] M.K Kwong. *Uniqueness of positive solutions of $\Delta u - u + u^p = 0$ in $\mathbf{R}^n$* . Arch. Rational Mech. Anal., 105(1989):243–266.
- [13] P.-L. Lions. *The concentration-compactness principle in the calculus of variations. The locally compact case. II*. Ann. Inst. H. Poincaré Anal. Non Linéaire, 1(1984):223–283.
- [14] F. Merle and P. Raphael. *Blow up dynamic and upper bound on the blow up rate for critical nonlinear Schrödinger equation*. In Journées “Équations aux Dérivées Partielles” (Forges-les-Eaux, 2002), pages Exp. No. XII, 5. Univ. Nantes, Nantes, 2002.
- [15] F. Merle and P. Raphael. *Sharp upper bound on the blow-up rate for the critical nonlinear Schrödinger equation*. Geom. Funct. Anal., 13(2003):591–642.
- [16] F. Merle and P. Raphael. *On universality of blow-up profile for L^2 critical nonlinear Schrödinger equation*. Invent. Math., 156(2004):565–672.
- [17] F. Merle and P. Raphael. *Profiles and quantization of the blow up mass for critical nonlinear Schrödinger equation*. Comm. Math. Phys., 253(2005):675–704.
- [18] F. Merle and P. Raphael. *On a sharp lower bound on the blow-up rate for the L^2 critical nonlinear Schrödinger equation*. J. Amer. Math. Soc., 19(2006):37–90 (electronic).
- [19] M. Ohta and G. Todorova. *Remarks on global existence and blowup for damped nonlinear Schrödinger equations*. Discrete Contin. Dyn. Syst., 23(2009):1313–1325.
- [20] T. Passota, C. Sulemb and P.L. Sulem. *Linear versus nonlinear dissipation for critical NLS equation*. Physica D, 203 (2005) 167184
- [21] F. Planchon and P. Raphaël. *Existence and stability of the log-log blow-up dynamics for the L^2 -critical nonlinear Schrödinger equation in a domain*. Ann. Henri Poincaré, 8(2007):1177–1219.
- [22] P. Raphael. *Stability of the log-log bound for blow up solutions to the critical nonlinear Schrödinger equation*. Math. Ann., 331(2005):577–609.
- [23] M. Tsutsumi. *Nonexistence of global solutions to the Cauchy problem for the damped nonlinear Schrödinger equations*. SIAM J. Math. Anal., 15(1984):357–366.
- [24] M.I. Weinstein. *Nonlinear Schrödinger equations and sharp interpolation estimates*. Comm. Math. Phys., 87(1982/83):567–576.

M. DARWICH: UNIVERSITÉ FRANÇOIS RABELAIS DE TOURS, LABORATOIRE DE MATHÉMATIQUES ET PHYSIQUE THÉORIQUE, UMR-CNRS 6083, PARC DE GRANDMONT, 37200 TOURS, FRANCE

E-mail address: Mohamad.Darwich@lmpt.univ-tours.fr